

OpenLoops: Supplementing Monte Carlos with Loop Matrix Elements

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Outline

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- Introduction
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- Usage

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- General considerations
- $W + 1, 2, 3$ jets

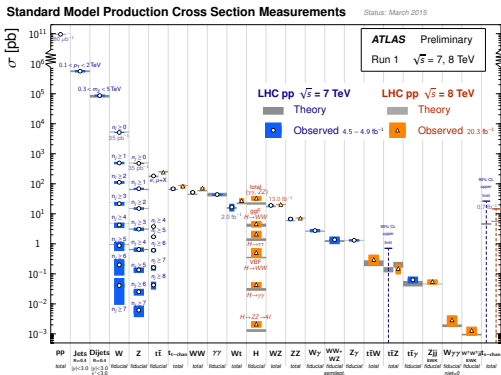
3 NNLO Applications: real-virtual corrections

- $pp \rightarrow W^+ W^-$

NLO automation

Many accurate measurements performed in LHC Run 1

- Consistent with Standard Model predictions.
- Run 2: Precision is likely to be the key to find BSM effects as small deviations from SM predictions.



- Many processes \rightarrow automated tools required.
- Feasibility 2 \rightarrow 4 NLO QCD well established.

Importance of NLO electroweak and NNLO QCD simulations becomes more and more apparent.

Shower Monte Carlos meet NLO

Before NLO matching: two different worlds

NLO calculations

- mostly fixed order
- high precision cross sections for inclusive observables
- until a few years ago: man-years of work per process

Shower Monte Carlos

- based on leading order matrix elements
- shapes and exclusive observables
- highly automated
- “close to experiments”

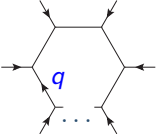
Ongoing: NLO QCD plus parton shower (and merging) is becoming the default accuracy for LHC simulations.

Upgrade existing tools to NLO, using loop matrix elements from a one-loop provider. Various more or less generic providers: BlackHat, GoSam, HELAC-NLO, MadLoop, NJet, OpenLoops, ...

Loop amplitudes

Diagrammatic representation of a loop amplitude:

factorised colour, **numerator** \mathcal{N} , **loop integral**, **momentum** q .



$$= \mathbf{c} \cdot \int d^d q \frac{\mathcal{N}(q)}{D_0 D_1 \dots D_{N-1}}$$

Must be reduced to scalar 1-, 2-, 3-, 4-point basis integrals:

$$= \mathbf{c} \cdot \sum a_i \text{---} \bigcirc \text{---} + \sum b_i \text{---} \bigcirc \text{---} + \sum c_i \text{---} \triangleleft \text{---} + \sum d_i \text{---} \square \text{---}$$

Reduction techniques/tools of our choice

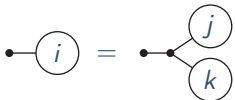
- OPP reduction: `CutTools` [Ossola, Papadopoulos, Pittau], scalar integrals from `OneLoop` [van Hameren]
- Tensor integral reduction: `Collier` [Denner, Dittmaier, Hofer]

Need to calculate the numerator function $\mathcal{N}(q)$.

From tree recursion to loop diagrams

Recursive construction of tree wave functions

Starting from external legs, connect wave functions w^α with vertices and propagators to recursively build “sub-trees”. Wave functions of sub-trees are 4-tuples of complex numbers (for the spinor/Lorentz index).



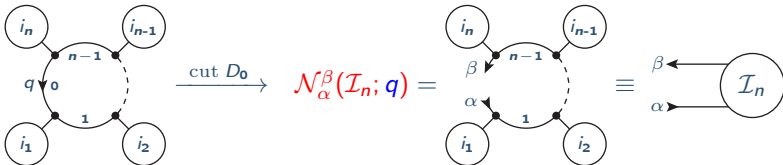
external lines are not depicted

$$w^\beta(i) = \frac{X_{\gamma\delta}^\beta}{p_i^2 - m_i^2} w^\gamma(j) w^\delta(k)$$

$X_{\gamma\delta}^\beta$ describes the interaction of i, j, k

Loop diagrams

A one-loop diagram is an ordered set of sub-trees $\mathcal{I}_n = \{i_1, \dots, i_n\}$, connected by loop propagators.



OpenLoops recursion

Connect sub-trees along the loop to build the numerator $\mathcal{N} \equiv \mathcal{N}_\alpha^\alpha$

$$\mathcal{N}_\alpha^\beta(\mathcal{I}_n; q) = \mathcal{X}_{\gamma\delta}^\beta(q) \mathcal{N}_\alpha^\gamma(\mathcal{I}_{n-1}; q) w^\delta(i_n)$$

Naively, the numeric recursion only works by fixing the loop momentum to numeric values. Original idea behind OPP reduction: do this for multiple values of $q \rightarrow$ CPU expensive.

OpenLoops: regard \mathcal{N} as a polynomial in q and formulate a recursion for the polynomial coefficients.

$$\mathcal{N}_\alpha^\beta(\mathcal{I}_n; q) = \sum_{r=0}^n \mathcal{N}_{\mu_1 \dots \mu_r; \alpha}^\beta(\mathcal{I}_n) q^{\mu_1} \dots q^{\mu_r}, \quad \mathcal{X}_{\gamma\delta}^\beta = Y_{\gamma\delta}^\beta + q^\nu Z_{\nu; \gamma\delta}^\beta$$

$$\mathcal{N}_{\mu_1 \dots \mu_r; \alpha}^\beta(\mathcal{I}_n) = \left[Y_{\gamma\delta}^\beta \mathcal{N}_{\mu_1 \dots \mu_r; \alpha}^\gamma(\mathcal{I}_{n-1}) + Z_{\mu_1; \gamma\delta}^\beta \mathcal{N}_{\mu_2 \dots \mu_r; \alpha}^\gamma(\mathcal{I}_{n-1}) \right] w^\delta(i_n)$$

The “open-loops” coefficients $\mathcal{N}_{\mu_1 \dots \mu_r; \alpha}^\beta$ encode the **functional dependence** of the numerator on the loop momentum.

Remarks

- Inspired by a Dyson-Schwinger recursion for loop momentum polynomials [van Hameren '09], but diagrammatic.
- Numerical implementation requires only **universal building blocks**, derived from the Feynman rules of the theory.
- Naturally works with both, tensor integrals and OPP. OPP similarly efficient as tensor integrals thanks to
 - cheap multiple evaluations of the numerator for multiple values of the loop momentum.
 - interference with Born amplitude, colour and helicity summation before reduction to scalar integrals.
- By 1-2 orders of magnitudes faster than original OPP based algorithms. → **Performance breakthrough**.
- Rational contributions of type R_2 are restored by counterterm-like Feynman rules. [Draggiotis, Garzelli, Malamos, Papadopoulos, Pittau '09, '10; Shao, Zhang, Chao '11]

Performance

process	diags	size/MB	time/ms
$u\bar{u} \rightarrow t\bar{t}$	11	0.1	0.27(0.16)
$u\bar{u} \rightarrow W^+W^-$	12	0.1	0.14
$u\bar{d} \rightarrow W^+g$	11	0.1	0.24
$u\bar{u} \rightarrow Zg$	34		0.75
$gg \rightarrow t\bar{t}$	44	0.2	1.6(0.7)
$u\bar{u} \rightarrow t\bar{t}g$	114	0.4	4.8(2.4)
$u\bar{u} \rightarrow W^+W^-g$	198	0.4	3.4
$u\bar{d} \rightarrow W^+gg$	144	0.5	4.0
$u\bar{u} \rightarrow Zgg$	408		17
$gg \rightarrow t\bar{t}g$	585	1.2	40(14)
$u\bar{u} \rightarrow t\bar{t}gg$	1507	3.6	134(101)
$u\bar{u} \rightarrow W^+W^-gg$	2129	2.5	89
$u\bar{d} \rightarrow W^+ggg$	1935	4.2	120
$u\bar{u} \rightarrow Zggg$	5274		524
$gg \rightarrow t\bar{t}gg$	8739	16	1460(530)

i7-3770K (single thread),
 gfortran 4.8, dynamic
 (ifort static ~ 30% faster),
 tensor integral reduction
 with Collier.

Colour and helicity
 summed, heavy quark
 loops included.

W/Z production includes
 leptonic decays and non-
 resonant contributions.

$t\bar{t}$ production numbers
 in brackets are for
 massless decays.

CutTools provides similar performance for $2 \rightarrow 4$, but is slower for lower multiplicities.

Installation

OpenLoops is available for download from hepforge

<http://openloops.hepforge.org>

Requirements: gfortran 4.6 or later, Python 2.x ($x \geq 4$)

The easiest way to install and keep your installation up to date is to pull a copy from the subversion repository (there's also a tarball available on the website)

svn checkout \

```
http://openloops.hepforge.org/svn/OpenLoops/branches/public \
OpenLoops cd OpenLoops
```

go to the website and look up the processes which are available for download and install the ones you need, e.g. for Z + up to three jets:

```
./openloops libinstall ppzj ppzjj ppzjjj
```

Or to install everything: `./openloops libinstall all.coll`

To update OpenLoops and installed processes: `./openloops update`

Use OpenLoops in your own programs

OpenLoops can be used via

- a Binoth Les Houches Accord (BLHA) interface
- its native interface in Fortran and C

Provides tree level matrix elements for Born and real corrections, including spin and colour correlations, 1-loop and (1-loop)².

For NLO simulations these must be supplemented by phase space integration and a subtraction procedure, possibly a parton shower, multi-jet merging, hadronisation, ...

$\mathcal{O}(100)$ processes available: almost the full range of $2 \rightarrow 4$ (plus decays) Standard Model processes with QCD corrections.

Physics models supported (processes will be added step-by-step)

- Full Standard Model with QCD and electroweak corrections
- Higgs Effective field theory
- Two Higgs Doublet Model Type I and II

Code example

```
program main
  use openloops
  implicit none
  integer :: id
  real(8) :: psp(0:3,5), m2l0, m2l1(0:2), acc, sqrt_s = 1000
  call set_parameter("order_ew", 1) ! coupling order
  call set_parameter("mass(23)", 91.2)
  id = register_process("1 -1 -> 23 2 -2", 11) ! dd~ -> Zuu~
  call start() ! initialise
  call set_parameter("alpha_s", 0.1)
  call set_parameter("mu", 100)
  call phase_space_point(id, sqrt_s, psp) ! phase-space pt. from Rambo
  call evaluate_loop(id, psp, m2l0, m2l1, acc) ! matrix element
  print *, "tree, loop ep^0,-1,-2, accuracy:", m2l0, m2l1, acc
  call finish()
end program main
```

Interfaces to Monte Carlo programs

■ Sherpa

Interface **included in the official release** 2.1.0 (or later).

Full-featured event generator, Sherpa+OpenLoops is well tested and used in several published simulations.

■ **Herwig++**: Official support in the next Herwig++ release.

■ **POWHEG**: Coming soon.

■ **MUNICH**: Parton-level Monte Carlo by **S. Kallweit** (to be published). Particularly used in NNLO calculations.

Applications

NLO: focus on applicability to experiments, “beyond fixed order”

- MEPS@NLO for $ll\nu\nu + 0, 1$ jets [Cascioli, Höche, Krauss, PM, Pozzorini, Siegert]
- (loop)² merging for $HH + 0, 1$ jets [PM, Papaefstathiou]
- NLO $W^+W^-b\bar{b}$ with $m_b > 0$ [Cascioli, Kallweit, PM, Pozzorini]
- MEPS@NLO $W^+W^-W^\pm + 0, 1$ jets [Höche, Krauss, Pozzorini, Schönherr, Thompson, Zapp]
- S-MC@NLO for $t\bar{t}b\bar{b}$ with $m_b > 0$ [Cascioli, PM, Moretti, Pozzorini, Siegert]
- MEPS@NLO for $t\bar{t} + 0, 1, 2$ j [Höche, Krauss, PM, Pozzorini, Schönherr, Siegert]

NNLO: for real-virtual corrections, partly for (1-loop)²

- $Z\gamma$ [Grazzini, Kallweit, Rathlev, Torre]
- ZZ [Cascioli, Gehrmann, Grazzini, Kallweit, PM, v. Manteuffel, Pozzorini, Rathlev, Tancredi, Weihs]
- W^+W^- [Gehrmann, Grazzini, Kallweit, PM, v. Manteuffel, Pozzorini, Rathlev, Tancredi]
- $t\bar{t}$ [Abelof, Gehrmann-De Ridder, PM, Pozzorini]

NLO-EW: $W + 1, 2, 3$ jets [Kallweit Lindert, PM, Pozzorini, Schönherr]

Electroweak Sudakov Effects

Naively, or more precisely at small momentum transfer $Q \lesssim M_W$, electroweak corrections are small compared to QCD corrections: $\alpha \sim \alpha_s^2$

But: Electroweak corrections can be strongly enhanced by $\log \frac{Q^2}{M_W^2}$.

Large logarithms arise from weak boson exchange between on-shell legs.

Estimate for 1-loop leading log (LL) and next-to-leading log (NLL) at $Q = 1$ TeV:

$$\text{LL} \simeq -\frac{\alpha}{\pi s_W^2} \log^2 \frac{Q^2}{M_W^2} \simeq -26\%$$

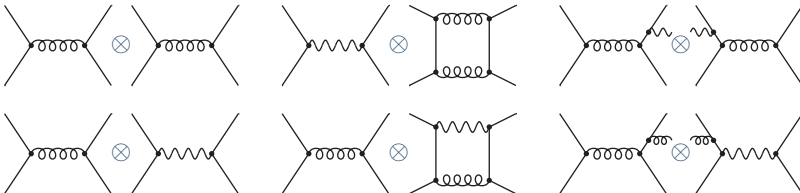
$$\text{NLL} \simeq +\frac{3\alpha}{\pi s_W^2} \log \frac{Q^2}{M_W^2} \simeq +16\%$$

- $\mathcal{O}(10\%)$ effect at 1 TeV.
- Could hide BSM effects in tails of distributions.
- Strongly process dependent.

Need electroweak NLO corrections

Power Counting

The notion of QCD and EW corrections is not well defined.



- EW correction to “QCD Born” or QCD correction to “EW Born”?
- For individual diagrams, the correction type can be defined via the origin of the UV poles (if any), but not for a full process.
- Real corrections involve photon and QCD parton emissions.
- For a consistent definition of the perturbative order and for infra-red finiteness:

Take into account complete $\mathcal{O}(\alpha^n \alpha_s^m)$ contributions.

NLO EW Automation

Things to consider

- Electroweak infra-red subtraction: photon emission.
Heavy boson emission: only partial cancellation between real and virtual. Usually best treated as separate processes.
- Jet/Photon definition: collinear $q \rightarrow q\gamma$ splitting.
- Matching to QED shower. Merging?

Performance (wrt. QCD corrections)

process	diagrams		size [MB]		time [ms/point]	
	QCD	EW	QCD	EW	QCD	EW
$d\bar{d} \rightarrow t\bar{t}$	11	33	0.1	0.2	0.27	0.69
$g\bar{g} \rightarrow t\bar{t}$	44	70	0.2	0.3	1.6	2.8
$d\bar{d} \rightarrow t\bar{t}g$	114	360	0.4	0.9	4.8	13
$g\bar{g} \rightarrow t\bar{t}g$	585	660	1.4	1.6	40	56
$d\bar{d} \rightarrow t\bar{t}u\bar{u}$	236	1274	0.8	2.8	12	48
$g\bar{g} \rightarrow t\bar{t}g\bar{g}$	8739	7614	18	16	1458	1557

CPU: i7-3770K, compiler: gfortran 4.8

EW corrections are moderately slower than QCD corrections.

NLO EW $W^+ + 1, 2, 3$ jets

W +jets: Testing ground for theory and tools in multi-jet cases.

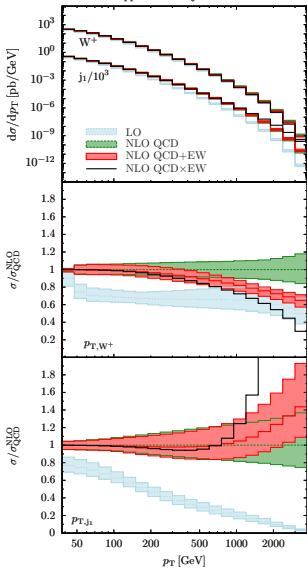
- NLO QCD available up to $W + 5j$ [Blackhat '09,'11,'13]
- $\mathcal{O}(< 10\%)$ QCD scale uncertainties.
- NLO EW $W + 1j$: $\mathcal{O}(-50\%)$ EW Sudakov effects at TeV scale
[Kühn, Kulesza, Pozzorini, Schulze '04-'07; Denner, Dittmaier, Kasprzik, Mück '09-'11]

Here: NLO EW $W^+ + 1, 2, 3j$

- On-shell external W bosons: $\Gamma_W = 0$
- Regulate s channel propagators of heavy particles with a small width $\Gamma_{\text{reg}} \rightarrow 0$.

$W^+ + j$ inclusive

$pp \rightarrow W^+ + 1j @ 13 \text{ TeV}$



Inclusive

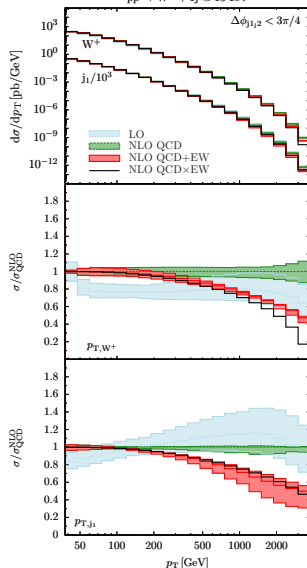
- -30% EW corr. wrt. NLO QCD at 2 TeV.
- Large QCD corrections, arising from (back-to-back) dijet configurations. → effectively LO.

Exclusive

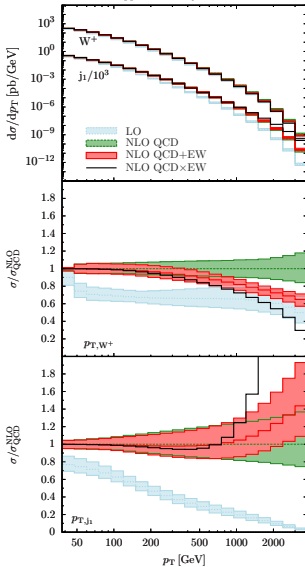
- -40% EW corr. wrt. NLO QCD at 2 TeV.
- Moderate QCD corrections.
- EW corrections exceed QCD uncertainties above $p_{T,W} > 300 \text{ GeV}$

$W^+ + j, \Delta\phi_{j_1,j_2} < 3\pi/4$

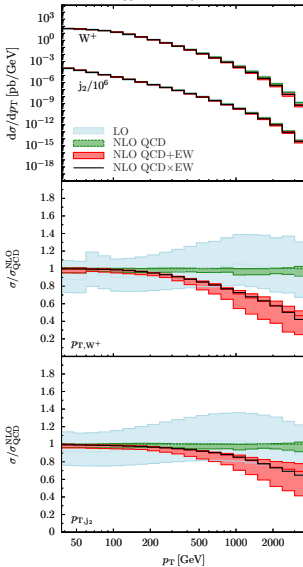
$pp \rightarrow W^+ + 1j @ 13 \text{ TeV}$



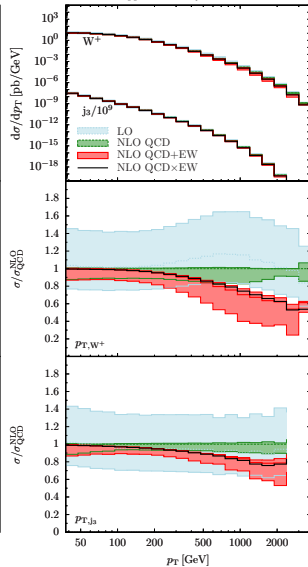
W+1j

pp → W⁺ + 1j @ 13 TeV

W+2j

pp → W⁺ + 2j @ 13 TeV

W+3j

pp → W⁺ + 3j @ 13 TeV

NNLO $pp \rightarrow W^+ W^-$ [Gehrmann, Grazzini, Kallweit, PM, v. Manteuffel, Pozzorini, Rathlev, Tancredi]

$W^+ W^-$ cross section: 20% excess (2σ) wrt. NLO+ gg theory.

- underestimated uncertainties?

Ingredients to on-shell $W^+ W^-$ production

- 2-loop master integrals [Gehrmann, Tancredi, v. Manteuffel, Weihs; Caola, Henn, Melnikov, Smirnov] and amplitudes [Gehrmann, Tancredi, v. Manteuffel]
- 1-loop amplitudes from OpenLoops
- NNLO q_T subtraction [Catani, Grazzini; Kallweit, Rathlev]

The real-virtual subtraction is particularly challenging, because one must integrate over phase space regions where the gluon is **soft or collinear** with a massless quark \rightarrow 1-loop matrix elements become **numerically unstable**.

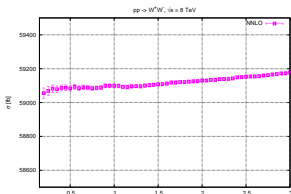
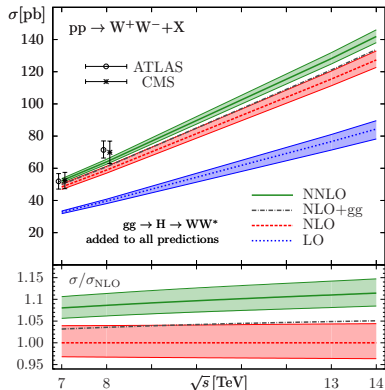
Matrix elements and subtraction terms are both divergent.

\rightarrow **High precision required** to have precision left after subtraction.

Works with the “vanilla” OpenLoops stability trigger with Collier and CutTools (could be tuned further).

NNLO $pp \rightarrow W^+W^-$: Results

- NNLO corrections between 9% (7 TeV) and 12% (14 TeV).
- Only 35% of the NNLO correction arises from $gg \rightarrow W^+W^-$.
- 3% uncertainty estimated from scale variations.
 No significant reduction wrt. NLO.
 More reliable, because at NNLO all channels are open.
- Eases tension with ATLAS/CMS 8 TeV data.



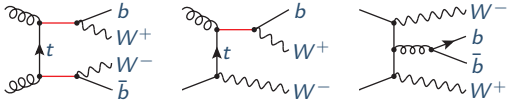
Dependence of the cross section on $q_T/m_{WW} \rightarrow 0$

- Stable at 1%.
- Increasing uncertainties from real-virtual and double real to similar amounts.

Subtraction of Top Quark Contributions

Huge top contamination due to

$pp \rightarrow Wt \rightarrow W^+ W^- b$ (30-60% at NLO for 7-14 TeV) and
 $pp \rightarrow t\bar{t} \rightarrow W^+ W^- b\bar{b}$ (400%-800% at NNLO).



“Top-free” definition of $W^+ W^-$

- Experimental: suppress top contribution with jet vetos and subtract remnant top background.
- 4-flavour scheme: massive b -quarks, excluding b -quark emission
- 5-flavour scheme: massless b -quark emissions needed to cancel collinear $g \rightarrow b\bar{b}$ singularity.

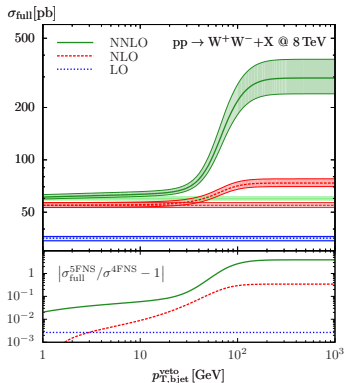
Subtract top contribution by exploiting $\xi_t = \Gamma_t / \Gamma_t^{phys}$ scaling.

Numerical fit $\sigma^{5F} = \xi_t^{-2} \sigma_{t\bar{t}}^{5F} + \xi_t^{-1} \sigma_{Wt}^{5F} + \sigma_{W^+W^-}^{5F}$

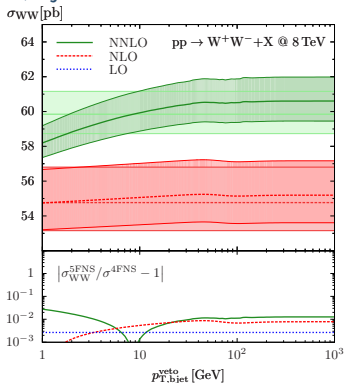
Assess theory ambiguities by comparing 4F and top subtracted 5F.

Top Subtraction and b-jet Veto

b-jet veto: cross section for $p_T < p_{T,b\text{-jet}}^{\text{veto}}$



Full 5F cross section vs. 4F
 Top contamination still 10%
 at $p_{T,b\text{-jet}}^{\text{veto}} = 10$ GeV.
 No plateau, due to singularity.



Top-free 5F cross section vs. 4F
 Stable top subtraction for
 $p_{T,b\text{-jet}}^{\text{veto}} = 10$ GeV.
 1% agreement between 4F and 5F.

Conclusions

OpenLoops makes using 1-loop matrix elements easy

- provides all matrix elements which are needed for NLO
- wide range of processes available, continuously extended
- very fast
- numerically stable, even for NNLO real-virtuals
- Standard Model, Higgs Effective Theory, soon 2HDM

Electroweak corrections

- NLO EW $2 \rightarrow 4$ at moderate higher CPU cost wrt. QCD
- Large corrections due to EW Sudakov logs in the TeV region
- $W + 1, 2, 3j$: multi-jet significantly different from $W + 1j$

<http://openloops.hepforge.org>