

BFKL NLL phenomenology : forward jets and Mueller Navelet jets

Christophe Royon
IRFU-SPP, CEA Saclay

Low x 2008, 6-10 July 2008, Kolimbari, Crete

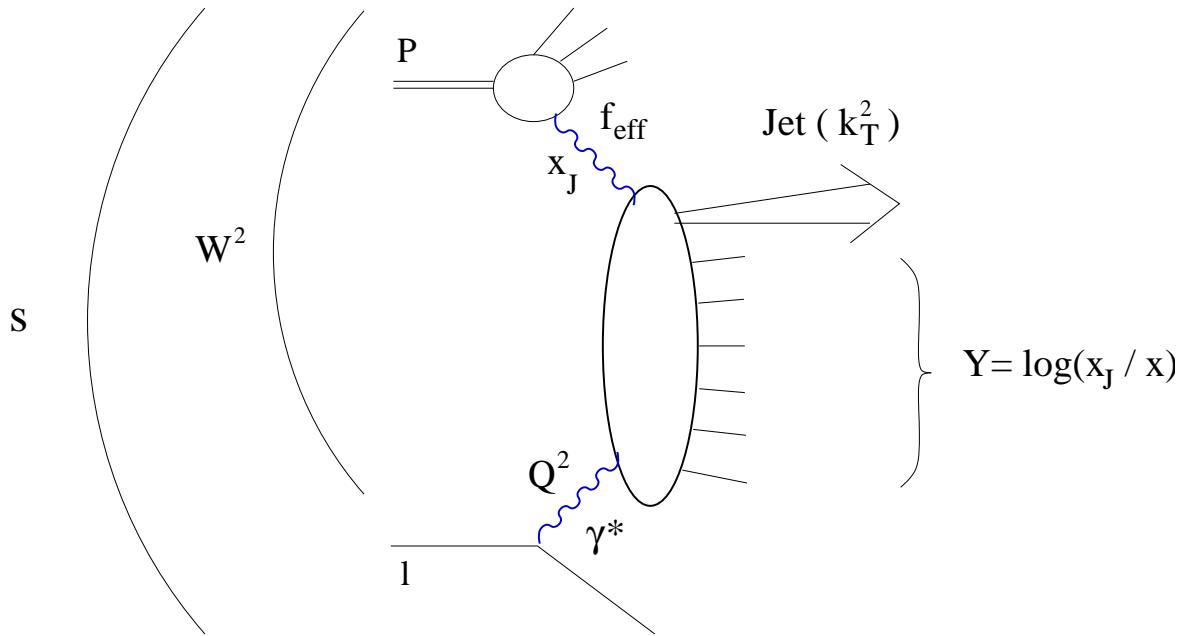
Contents:

- BFKL-NLL formalism
- Fit to H1 $d\sigma/dx$ data
- Prediction for the H1 triple differential cross section
- Prediction for Mueller Navelet jets at the Tevatron/LHC

Work done in collaboration with O. Kepka, C. Marquet, R. Peschanski

Phys. Lett. B 655 (2007) 236; Eur. Phys. J. C55 (2008) 259, arXiv:0704.3409

Forward jet measurement at HERA



- Typical kinematical domain where BFKL effects are supposed to appear with respect to DGLAP: $k_T^2 \sim Q^2$, and Q^2 not too large
- LO BFKL forward jet cross section: 2 parameters α_S , normalisation
- NLL BFKL cross section: one single parameter: normalisation (α_S running via RGE)

BFKL NLL and resummation schemes

- **NLO BFKL:** Corrections were found to be large with respect to LO, and lead to unphysical results
- **NLO BFKL kernels need resummation:** to remove additional spurious singularities in γ and $(1 - \gamma)$
- **NLO BFKL kernel:** (γ and ω associated to $\log Q^2$ and rapidity after Mellin transform)

$$\chi_{NLO}(\gamma, \omega) = \chi^{(0)}(\gamma, \omega) + \alpha(\chi_1(\gamma) - \chi_1^{(0)}(\gamma))$$

- $\chi_1(\gamma)$: calculated, NLO BFKL eigenvalues (Lipatov, Fadin, Camici, Ciafaloni)
- $\chi^{(0)}$ and $\chi_1(0)$: ambiguity of resummation at higher order than NLO, different ways to remove these singularities, not imposed by BFKL equation, Salam, Ciafaloni, Colferai
- Transformation of the energy scale: $\gamma \rightarrow \gamma - \omega/2$ (Salam) needed for F_2 but not for forward jet cross sections (the problem is symmetric contrary to F_2)
- **BFKL NLL full calculation available (no saddle point approximation):** resolution of implicit equation performed by numerical methods

BFKL NLL calculation

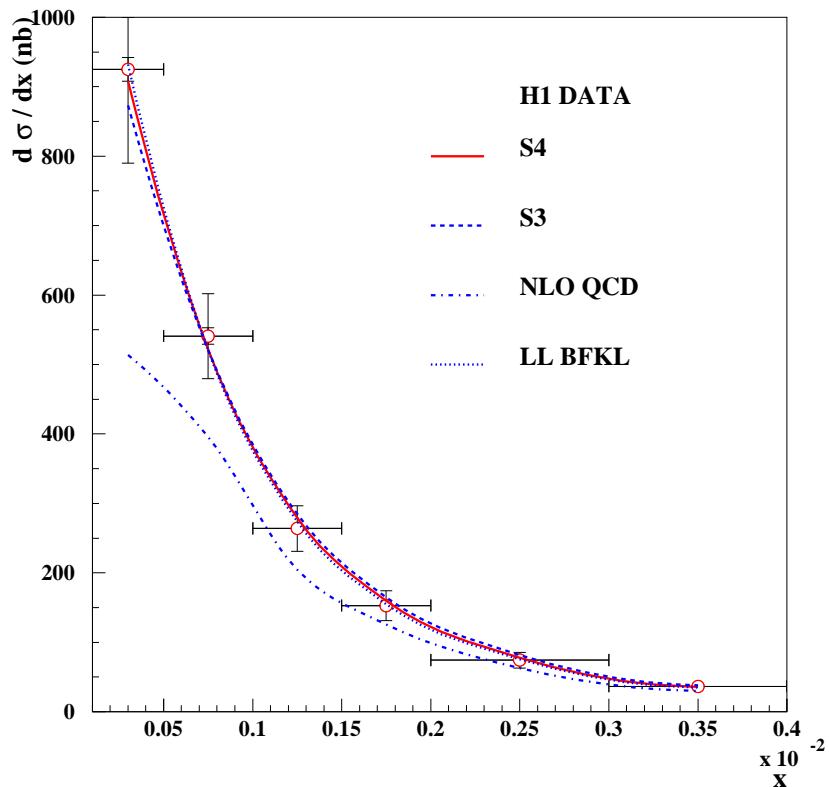
- Full BFKL NLL calculation available in S3 and S4 schemes for forward jet production (modulo the impact factors taken at LL)
- Equation:

$$\frac{d\sigma_{T,L}^{\gamma^* p \rightarrow JX}}{dx_J dk_T^2} = \frac{\alpha_s(k_T^2) \alpha_s(Q^2)}{k_T^2 Q^2} f_{eff}(x_J, k_T^2) \int \frac{d\gamma}{2i\pi} \left(\frac{Q^2}{k_T^2} \right)^\gamma \phi_{T,L}^\gamma(\gamma) e^{\bar{\alpha}(k_T Q) \chi_{eff}[\gamma, \bar{\alpha}(k_T Q)] Y}$$

- χ_{eff} computed using BFKL NLL formalism in the S3 and S4 schemes
- Implicit equation: $\chi_{eff}(\gamma, \alpha) = \chi_{NLL}(\gamma, \alpha, \chi_{eff}(\gamma, \alpha))$ solved numerically

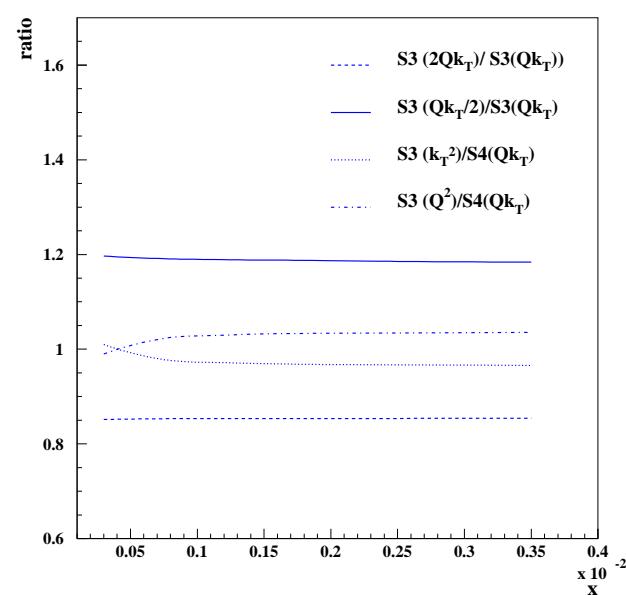
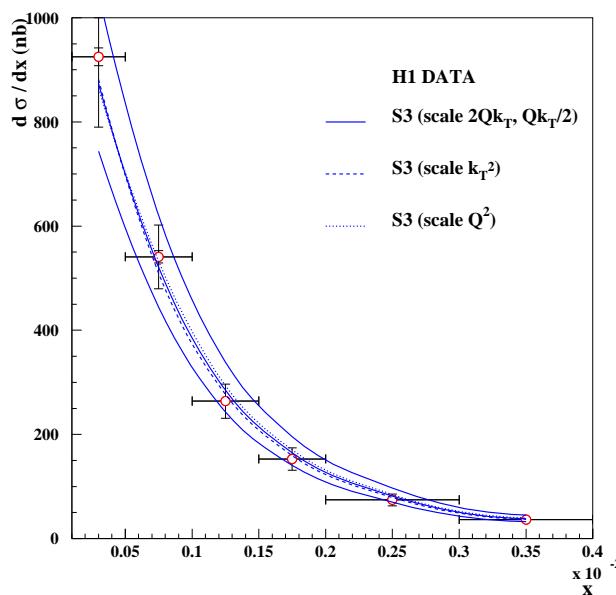
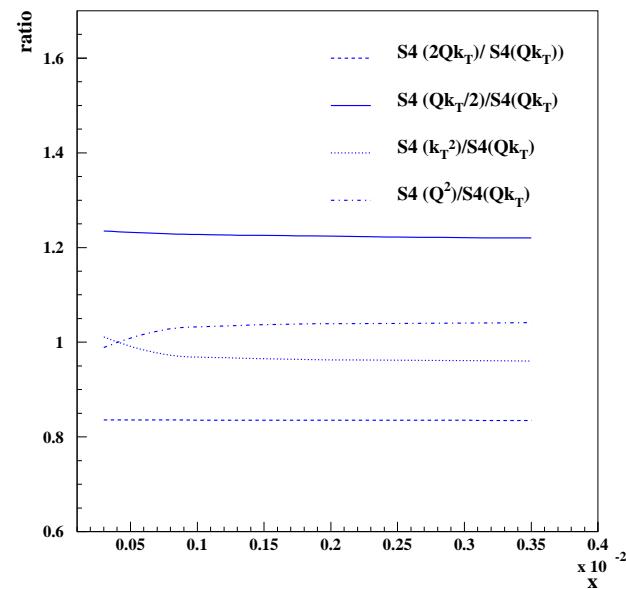
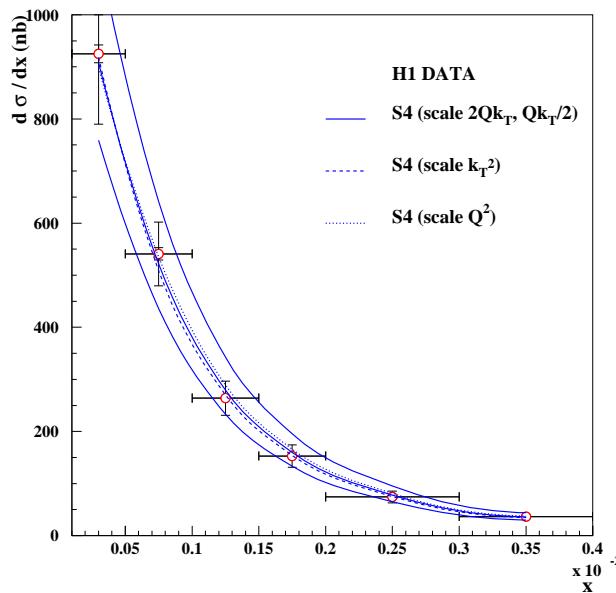
Fit results

- Fit of NLL BFKL calculation to the H1 $d\sigma/dx$ data: one single parameter, normalisation of cross section
- χ^2 for S3: 29.5 (1.15), S4: 10.0 (0.48)
- Good description of H1 data using BFKL LO and BFKL NLL formalism, DGLAP-NLO fails to describe the data
- BFKL higher corrections found to be small (We are in the BFKL-LO region, cut on $0.5 < k_T^2/Q^2 < 5$)



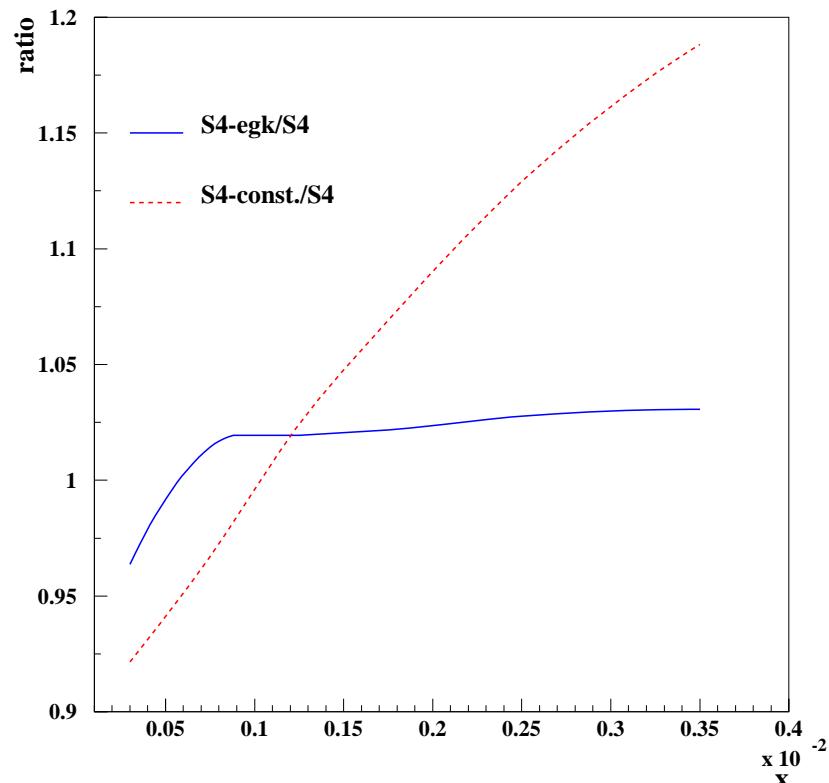
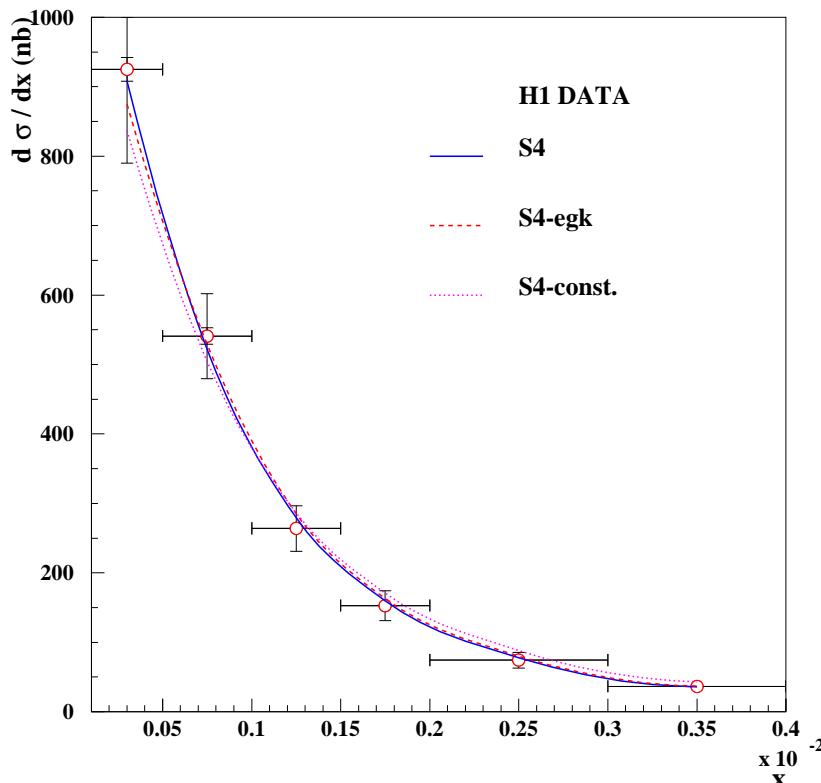
Scale variation - Resummation model variation

- **Scale dependence:** variation of the scale between $2Qk_T$, $Qk_T/2$, Q^2 , k_T^2 : $\sim 20\%$ difference
- **Resummation scheme dependence:** Use S3 and S4, S4 is slightly better



Dependence on impact factor

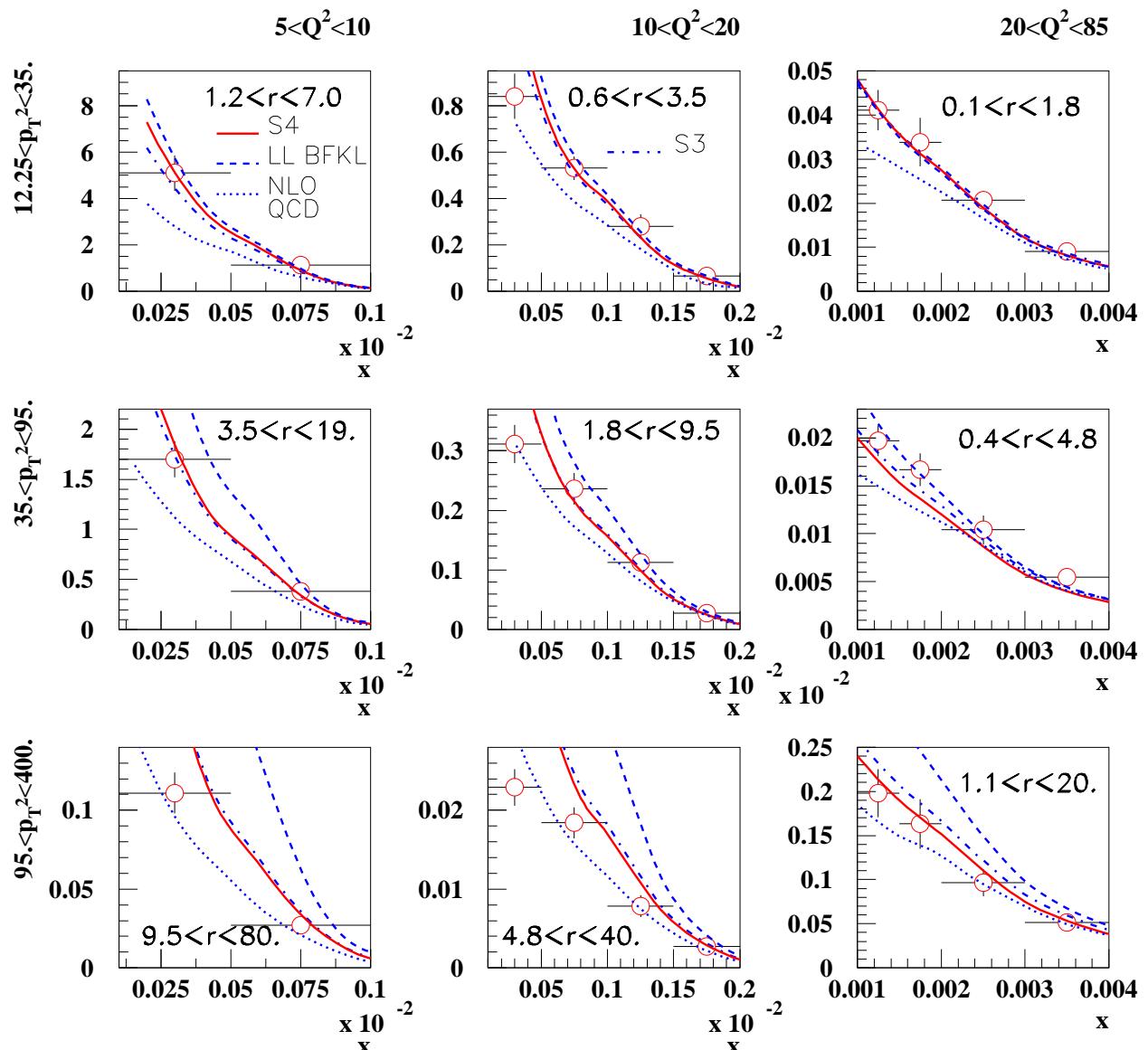
- Impact factor not yet fully known at NLL
- Variation of impact factor, 3 studies: h_T , $h_L(\gamma)$ at LO; h_T , $h_L(1/2)$ constant; implement the higher-order corrections in the impact factor due to exact gluon kinematics in the $\gamma^* \rightarrow q\bar{q}$ transition (see C.D. White, R. Peschanski, R.S. Thorne, Phys. Lett. B 639 (2006) 652)



Comparison with H1 triple differential data

- Triple differential cross section: Keep the normalisation from the fit to $d\sigma/dx$ and predict the triple differential cross section
- Good description over the full range

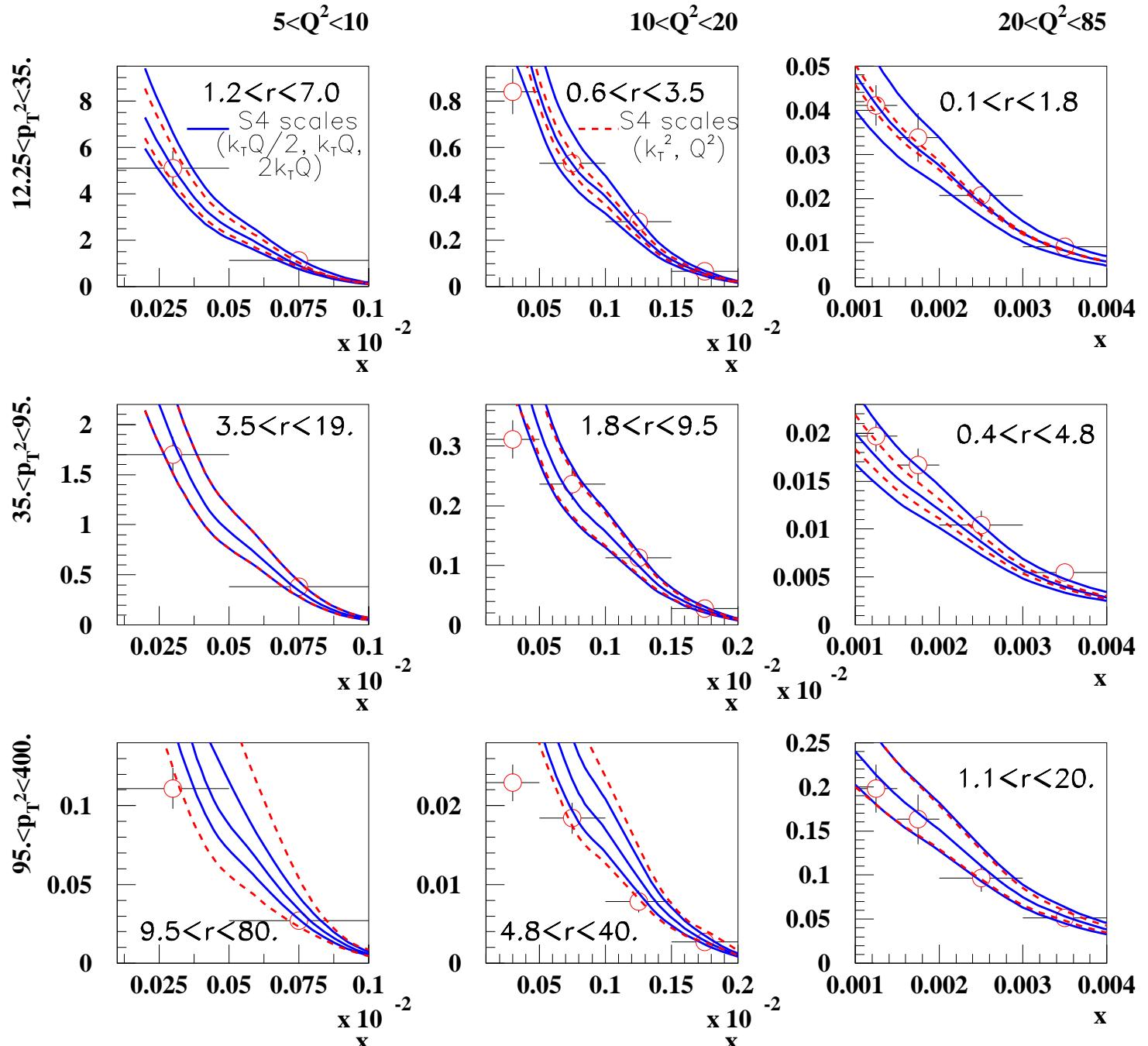
$d \sigma/dx dp_T^2 d Q^2 - H1 DATA$



Comparison with H1 triple differential data

Study of scale variation: 20% at low p_T^2 , > 70% at higher p_T^2
as for DGLAP

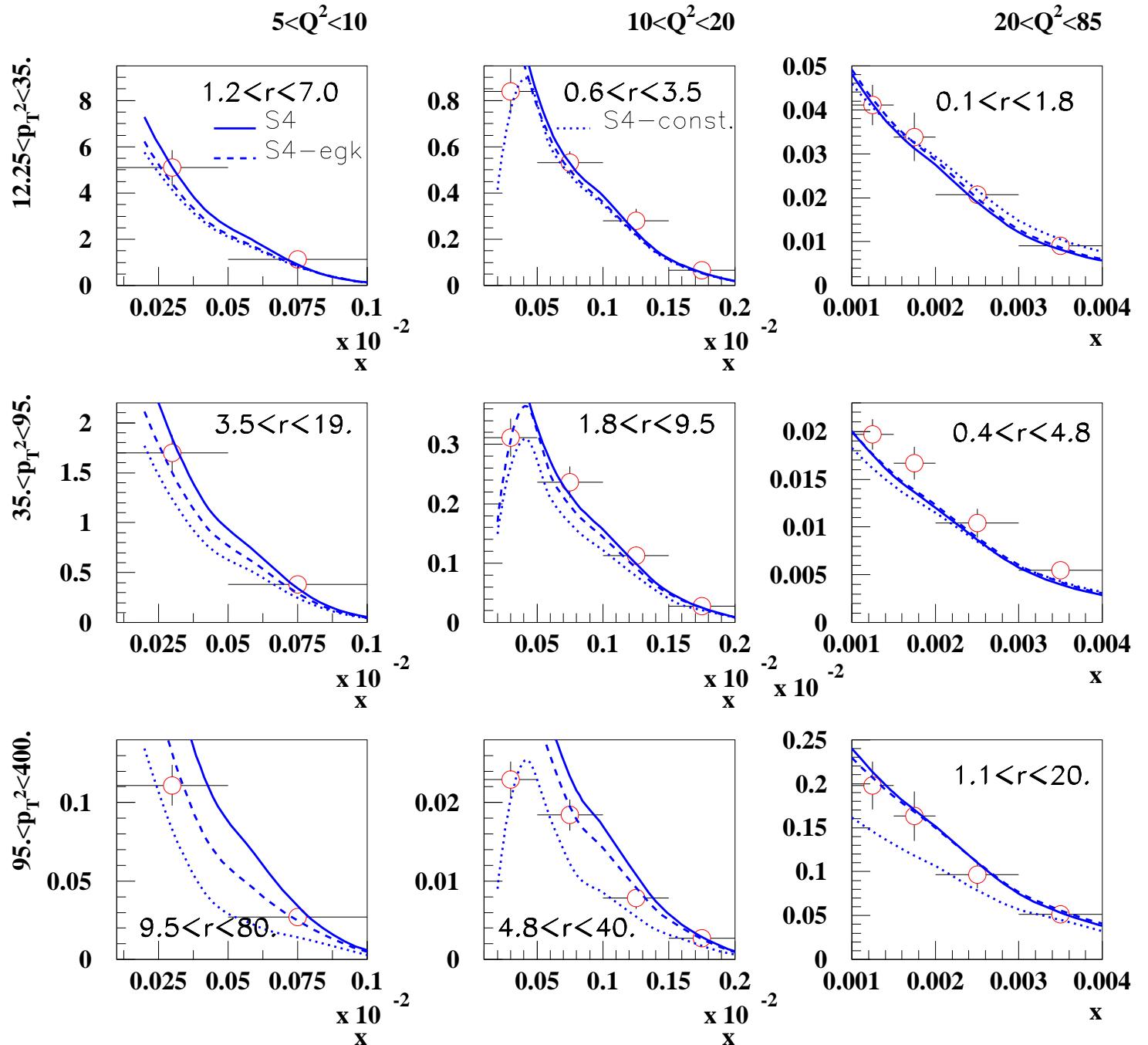
$d\sigma/dx dp_T^2 dQ^2 - \text{H1 DATA}$



Comparison with H1 triple differential data

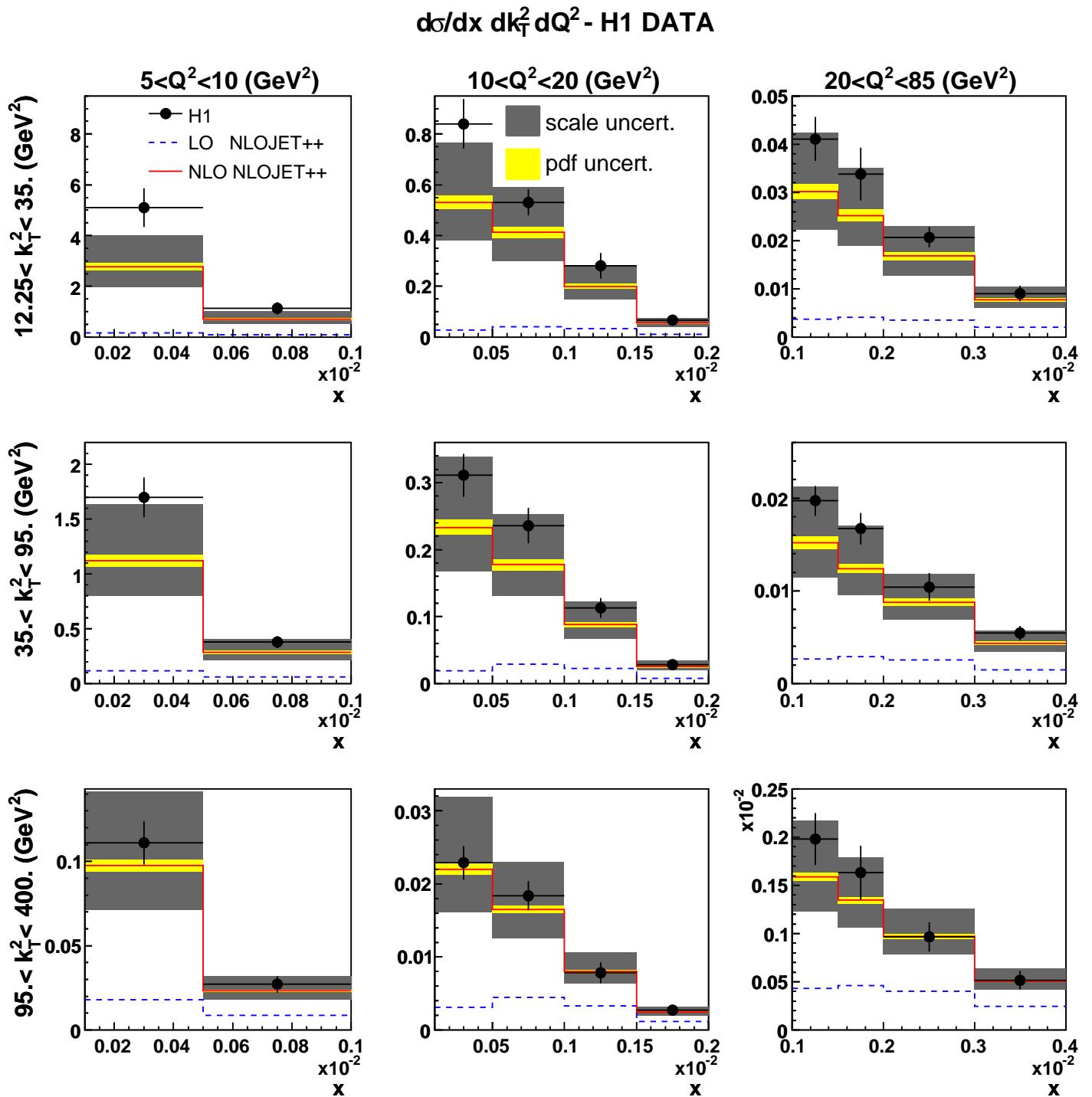
Study of dependence on impact factor

$d \sigma/dx dp_T^2 d Q^2 - H1 DATA$



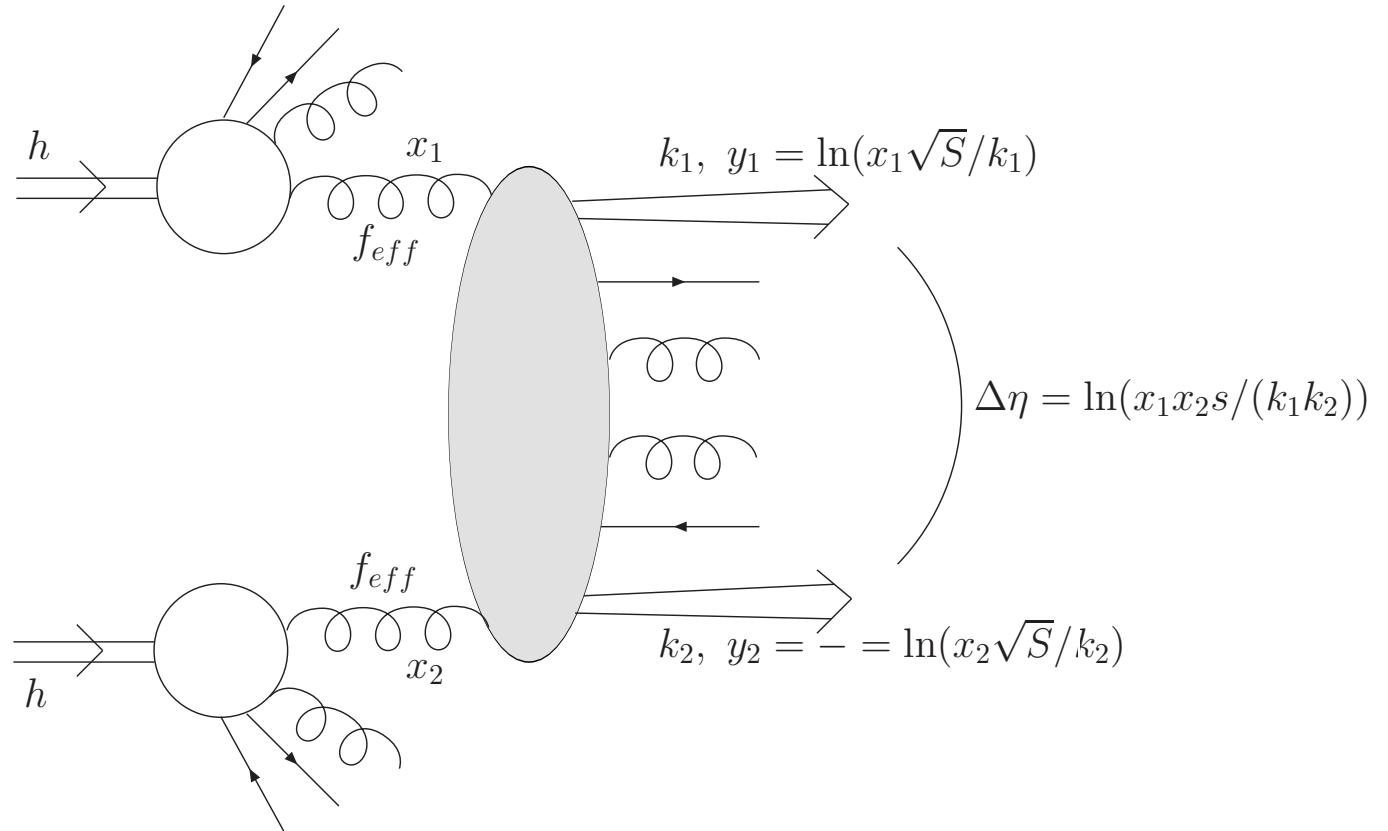
Comparison with H1 triple differential data

DGLAP study: large scale dependence



Mueller Navelet jets

Same kind of processes at the Tevatron and the LHC



- Same kind of processes at the Tevatron and the LHC:
Mueller Navelet jets
- Study the $\Delta\Phi$ between jets dependence of the cross section:

Mueller Navelet jets: $\Delta\Phi$ dependence

- Study the $\Delta\Phi$ dependence of the relative cross section
- Relevant variables:

$$\begin{aligned}\Delta\eta &= y_1 - y_2 \\ y &= (y_1 + y_2)/2 \\ Q &= \sqrt{k_1 k_2} \\ R &= k_2/k_1\end{aligned}$$

- Azimuthal correlation of dijets:

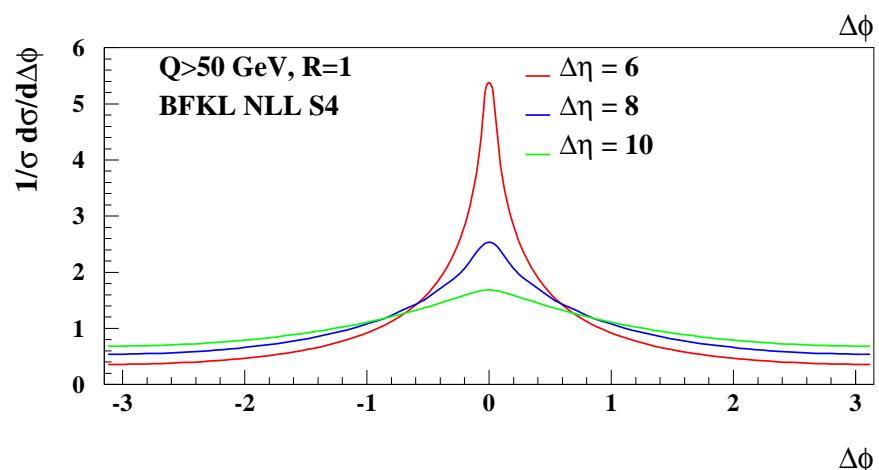
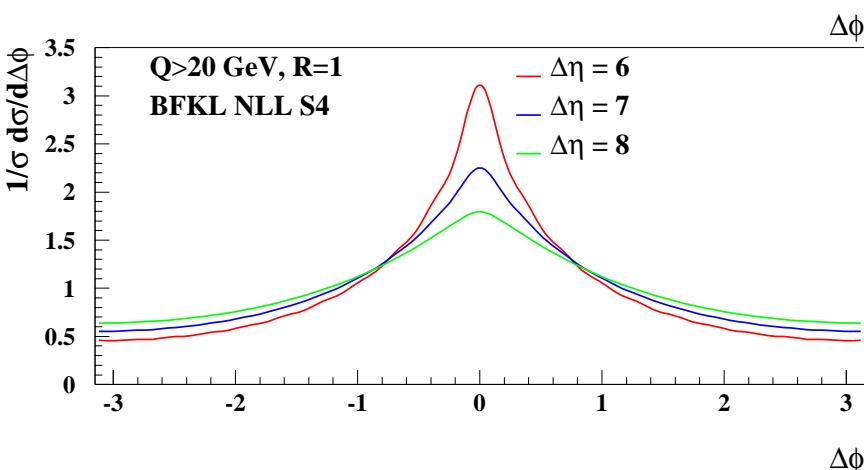
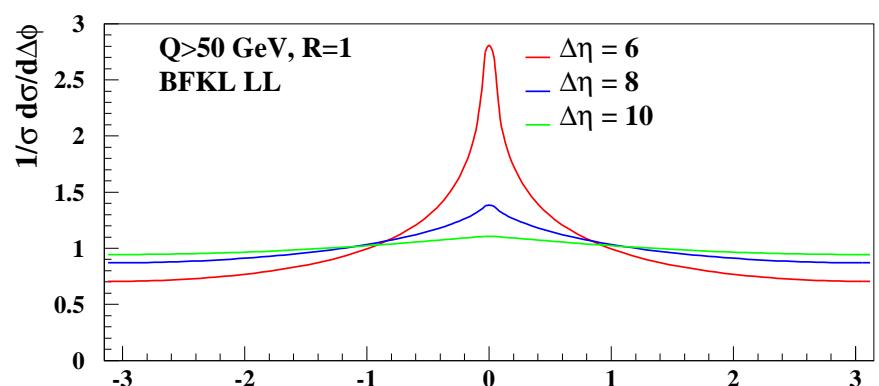
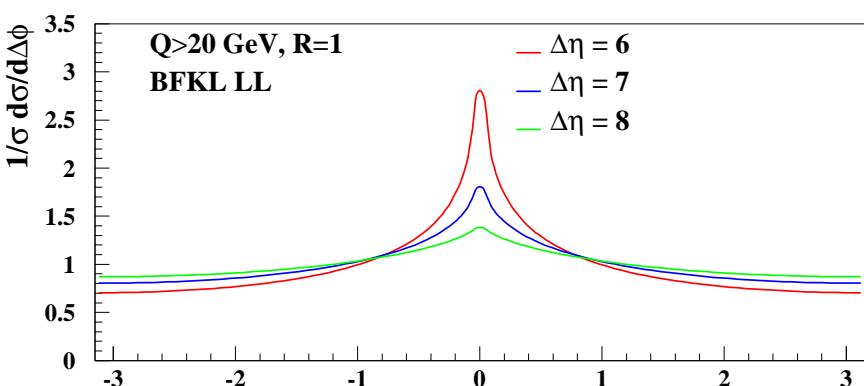
$$2\pi \frac{d\sigma}{d\Delta\eta dR d\Delta\Phi} \Bigg/ \frac{d\sigma}{d\Delta\eta dR} = 1 + \frac{2}{\sigma_0(\Delta\eta, R)} \sum_{p=1}^{\infty} \sigma_p(\Delta\eta, R) \cos(p\Delta\Phi)$$

where

$$\begin{aligned}\sigma_p &= \int_{E_T}^{\infty} \frac{dQ}{Q^3} \alpha_s(Q^2/R) \alpha_s(Q^2 R) \\ &\quad \left(\int_{y<}^{y>} dy x_1 f_{eff}(x_1, Q^2/R) x_2 f_{eff}(x_2, Q^2 R) \right) \\ &\quad \int_{1/2-\infty}^{1/2+\infty} \frac{d\gamma}{2i\pi} R^{-2\gamma} e^{\bar{\alpha}(Q^2)\chi_{eff}(p)\Delta\eta}\end{aligned}$$

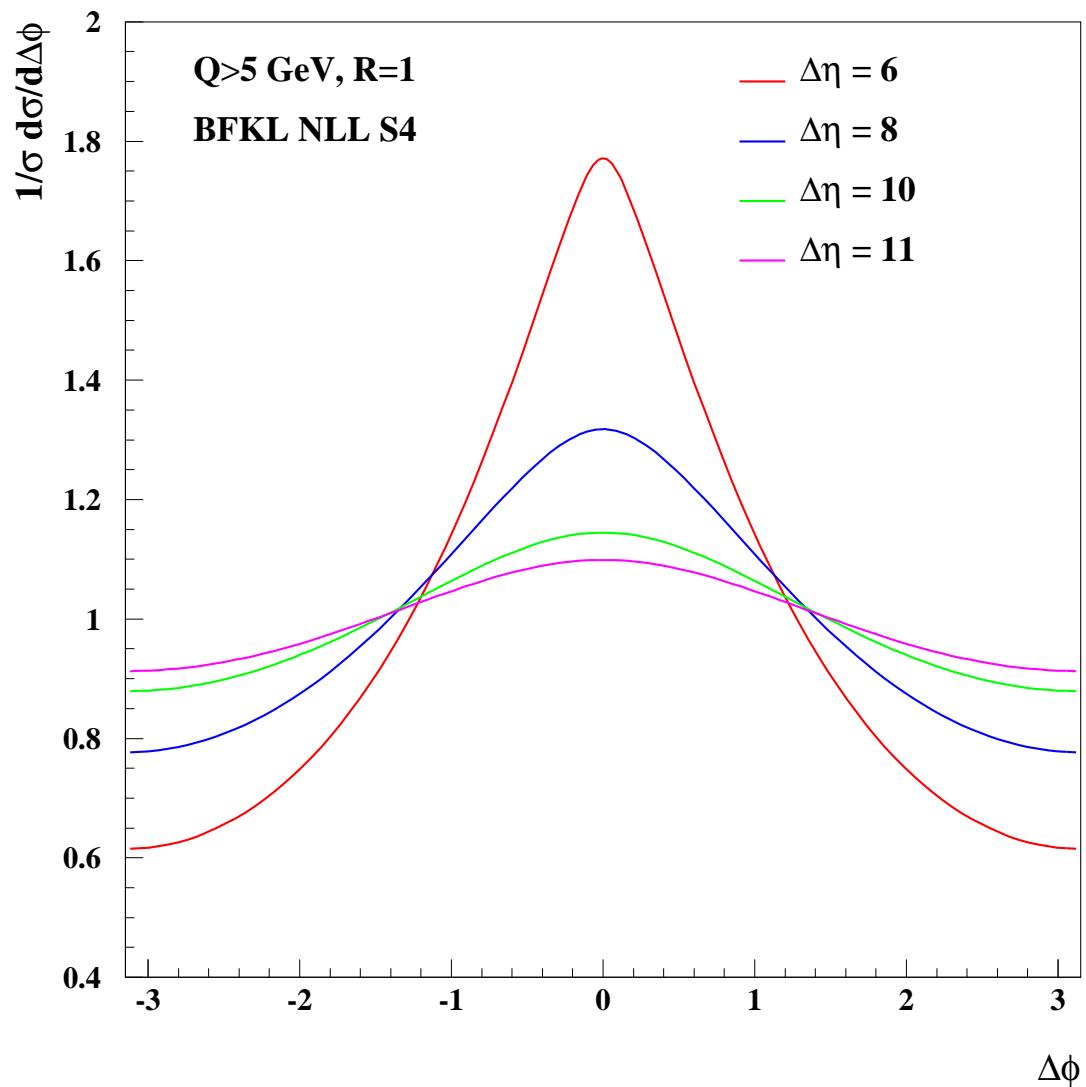
Mueller Navelet jets: $\Delta\Phi$ dependence

- $1/\sigma d\sigma/d\Delta\Phi$ spectrum for BFKL LL and BFKL NLL as a function of $\Delta\Phi$ for different values of $\Delta\eta$
- Measurement to be performed at the Tevatron/LHC



Mueller Navelet jets in CDF

Possibility to measure $\Delta\Phi$ distribution in CDF for large $\Delta\eta$ and low jet p_T ($p_T > 5$ GeV) using the CDF miniPLUG calorimeter



Conclusion

- DGLAP NLO fails to describe forward jet data
- BFKL NLL description of H1 and ZEUS forward jet data: very good description using full BFKL-NLL kernel and LO impact factors
- Study scale dependence and also dependence on assumption of impact factor: typically $\sim 20\%$ uncertainty, larger at high p_T
- Mueller Navelet jets: Full calculation available using S3 and S4 schemes
- Mueller Navelet jets $\Delta\Phi$ dependence: weak dependence even after NLL corrections, little sensitivity to chosen scale
- Mueller Navelet jets: Very nice measurement to be performed at the Tevatron/LHC, special use of CDF forward miniPLUG calorimeter which gives a good acceptance at large η and small p_T for jets