

# EXCLUSIVE PRODUCTION OF DIJETS AT THE TEVATRON (AND AT THE LHC)

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# OUTLINE

- 1 INTEREST
- 2 INGREDIENTS
- 3 UNCERTAINTIES AND IR REGION
- 4 LHC AND HIGGS
- 5 CONCLUSIONS

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## EXCLUSIVE PRODUCTION



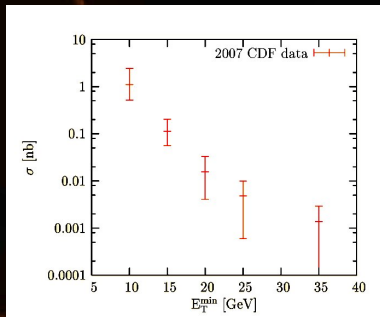
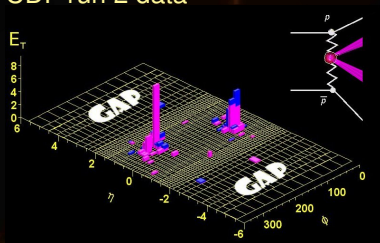
Embedding of a hard process into pomeron exchange →

- no underlying event
- little background if sharp resonance and measurement of the hadronic energy
- discovery tool for new physics decaying into hadrons

But is the rate sufficient?

# JET DATA

CDF run 2 data



- $\bar{p}$  measured,  
 $p$  reconstructed by  
Monte Carlo

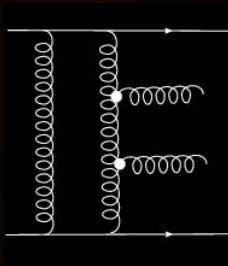
- $E_T$  from 10 to 35 GeV
- $\sigma = 1$  nb to 1 pb
- $M_{jj}$  up to 135 GeV  $\approx M_H$

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# BASIC INGREDIENTS

Collins-Berera

- Partonic singlet exchange:
  - $qq \rightarrow q + gg + q$
- Embed in hadrons:
  - $p\bar{p} \rightarrow p + gg + \bar{p}$
- Large vertex corrections
- Large screening corrections
- Make jets:
  - $p\bar{p} \rightarrow p + JJ + \bar{p}$

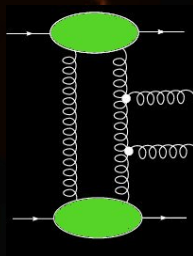


- fully calculable
- exact kinematics in  $\perp$  plane
- $\sigma = \infty$  (IR divergence)

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Cheng-Wu, Super-Gunion

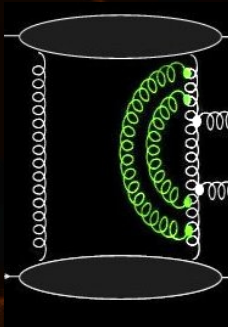
- regulates IR region
- changes the mass dependence
- $\sigma(E_T > 10 \text{ GeV}) \approx 600 \text{ nb}$



# BASIC INGREDIENTS

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*Diakonov-Dokshitzer-Trojan, Kaidalov-Khoze-Martin-Ryskin*



- very large correction
- double logs fully known
- $\sigma(E_T > 10 \text{ GeV}) \approx 25 \text{ nb}$

## BASIC INGREDIENTS

- Partonic singlet exchange:
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- under control if jet production at small distances
- otherwise depends on unitarisation scheme
- $\sigma(E_T > 10 \text{ GeV}) \approx 3 \text{ nb}$

## BASIC INGREDIENTS

- Partonic singlet exchange:

- $qq \rightarrow q + gg + q$

- Embed in hadrons:

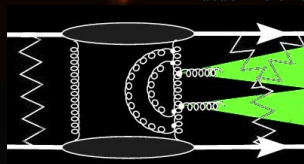
- $p\bar{p} \rightarrow p + gg + \bar{p}$

- Large vertex corrections

- Large screening corrections

- Make jets:

- $p\bar{p} \rightarrow p + JJ + \bar{p}$

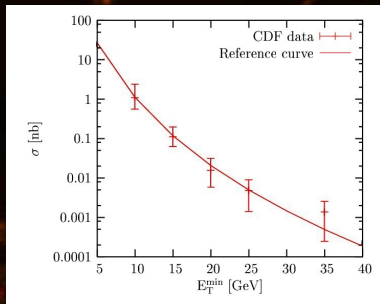


*Kaidalov-Khoze-Martin-Ryskin, Salam*

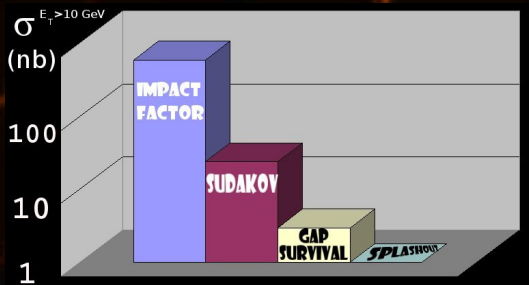
- several parametrisations
- $\sigma(E_T > 10 \text{ GeV}) \approx 1 \text{ nb}$

## POSSIBLE RESULT

	parameter	value
hard	scale of $\alpha_S$	$s_{gg}$
	$\Lambda_{QCD}^{(5)}$	200 MeV
Sudakov	scale of $\alpha_S$	loop momentum
	$\triangleleft$ ordering	yes
	terms	$\log^2 + \log + \text{constant}$
	lower scale	external off-shellness
	upper scale	$k_T^2/2$
impact factor	unintegrated	fitted to
	gluon density	$F_2$
gap survival	$\langle S^2 \rangle$	15%
splash-out	$E_T^{\text{jets}} / E_T^{\text{partons}}$	0.8



Huge correction factors:



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# UNCERTAINTIES

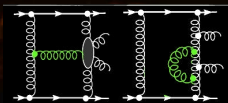
parameter	uncertainty (highest/lowest)
Sudakov	20
Impact factor	3
Gap survival	3
Splash-out	2
Total	~ 400

# EXAMPLE: SUDAKOV FORM FACTOR

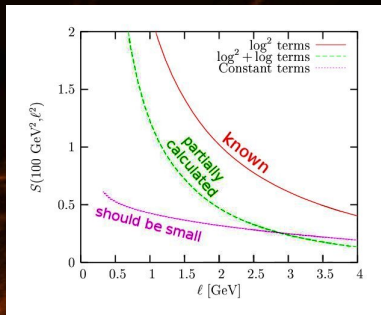
$$T(\mu^2, \ell^2) = \exp[-S(\mu^2, \ell^2)] , \quad S(\mu^2, \ell^2) = \int_{\ell^2}^{\mu^2} \frac{dq^2}{q^2} \frac{\alpha_s(q^2)}{2\pi} \int_0^{1-\Delta} dz [zP_{gg} + N_f P_{gq}]$$

Trick: virtual corrections  
 $\sim 1$ -brehmstrahlung

- true for  $\log^2$
- true for *some* log



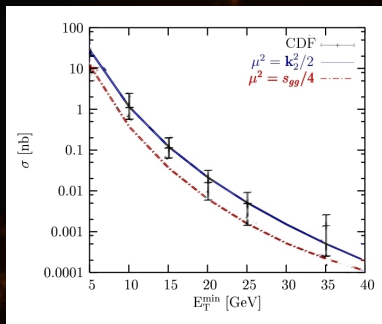
- not true for constant terms



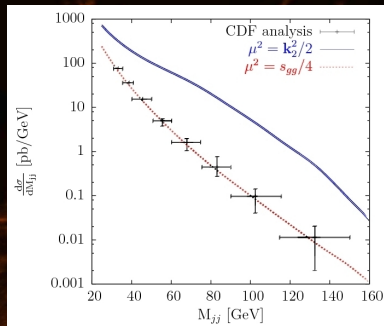


# CONSEQUENCES

Two curves fitting the  $E_T$  distribution



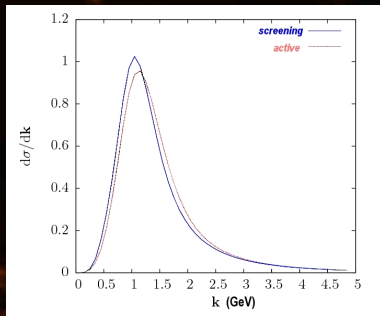
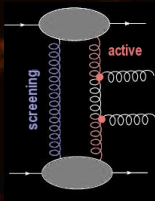
can produce **very** different mass distributions



The ExHuMe Monte-Carlo use to analyse the data takes  $\mu^2 = s_{gg}/2.62$ .

## GLUON MOMENTUM DISTRIBUTION

- One can write  $\sigma = \int dk \frac{d\sigma}{dk}$  with  $k$  the (non-measurable) momentum of one of the internal gluons.



- Only 30% of the cross section comes from the phase-space region with all off-shellnesses  $> 1$  GeV.

The calculation is tentative  
at best !

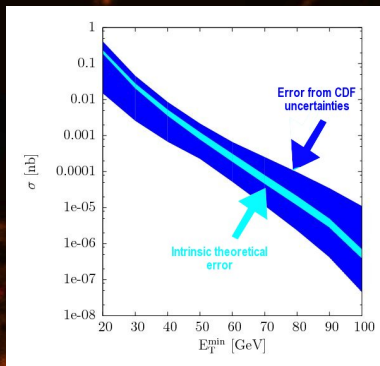
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# LHC DIJETS

One can predict the jet LHC cross section using the CDF result to calibrate it

- For typical cuts of FP420

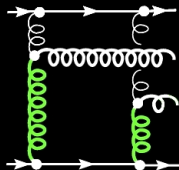
parameter	value
proton fractional momentum loss	0.002-0.02
jet rapidity	<1
mass of jet system	>50 GeV



## STANDARD CANDLE FOR HIGGS PRODUCTION?

Maybe not: there are diagrams in jet production which may be important, and which are not present in Higgs production:

Diagrams negligible at lowest order

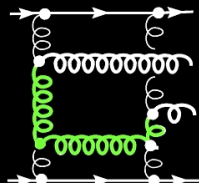


- extra propagator
- suppression by factor  $1/M_{jj}^2$

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Diagrams negligible at lowest order  
correspond to a specific routing of the hard momentum, without double logs

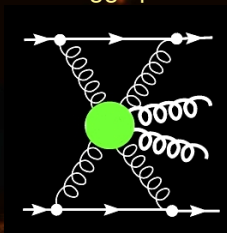


- the hard momentum must be routed outside of the hadron form factor

## STANDARD CANDLE FOR HIGGS PRODUCTION?

Maybe not: there are diagrams in jet production which may be important, and which are not present in Higgs production:

Diagrams negligible at lowest order correspond to a specific routing of the hard momentum, without double logs but further corrections to the localised hard momentum will contain double logs



- The localised hard vertex will have double logs

## STANDARD CANDLE FOR HIGGS PRODUCTION?

One knows the contribution to be negative, and the Sudakov corrections to be very different

$$\sim \frac{1}{M_{jj}^2} (1 + C_1 \alpha_S \log + C_2 \alpha_S^2 \log^3 + C_3 \alpha_S^3 \log^5 + \dots) \text{ vs.} \\ (1 - C \alpha_S \log^2 + \frac{C^2}{2} \alpha_S^2 \log^4 - \frac{C^3}{6} \alpha_S^3 \log^6 + \dots)$$

Tuning the theoretical uncertainties on the dijet data would only lead to a lower bound on the Higgs production cross section.



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# CONCLUSIONS

- Importance of CDF data to test theoretical ideas
- Calculation largely in the nonperturbative domain
- Large uncertainties (factor 20 up or down) due to Sudakov factor (and gap survival)
- Potential problems to relate the dijet production to Higgs boson production