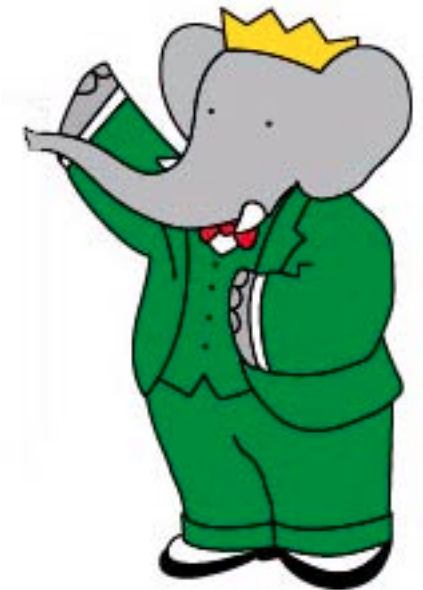
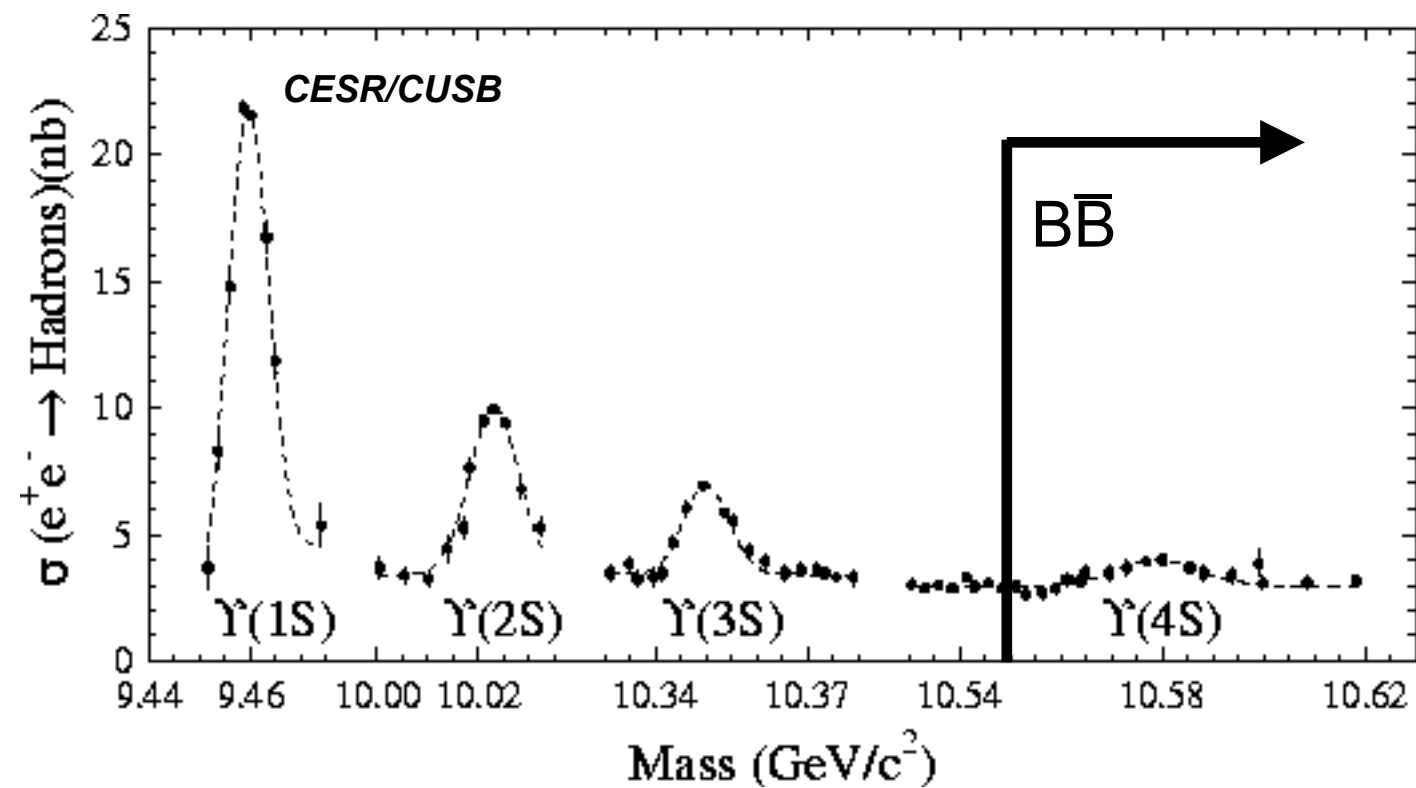


# Recent Results from BaBar: New Physics Searches

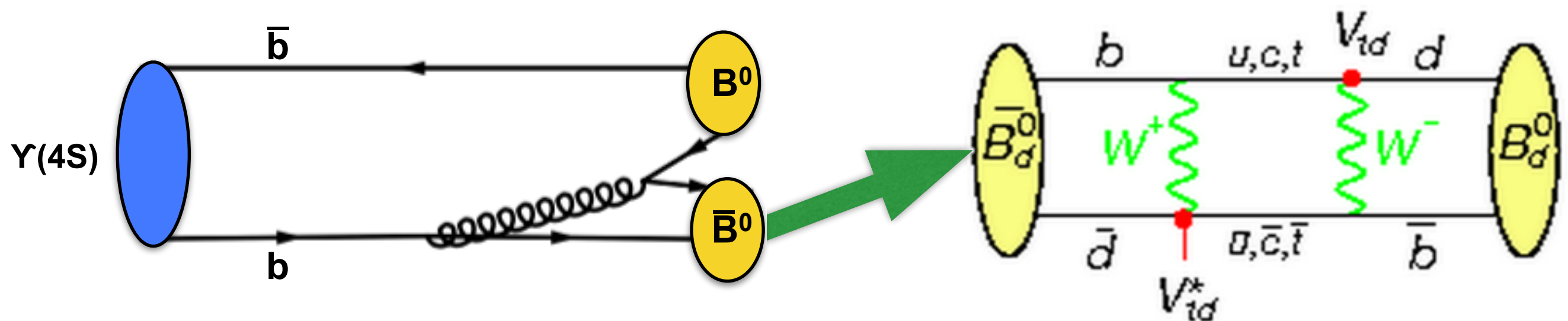
Matt Graham, SLAC  
2105 SLAC Summer Institute  
August 14 , 2015



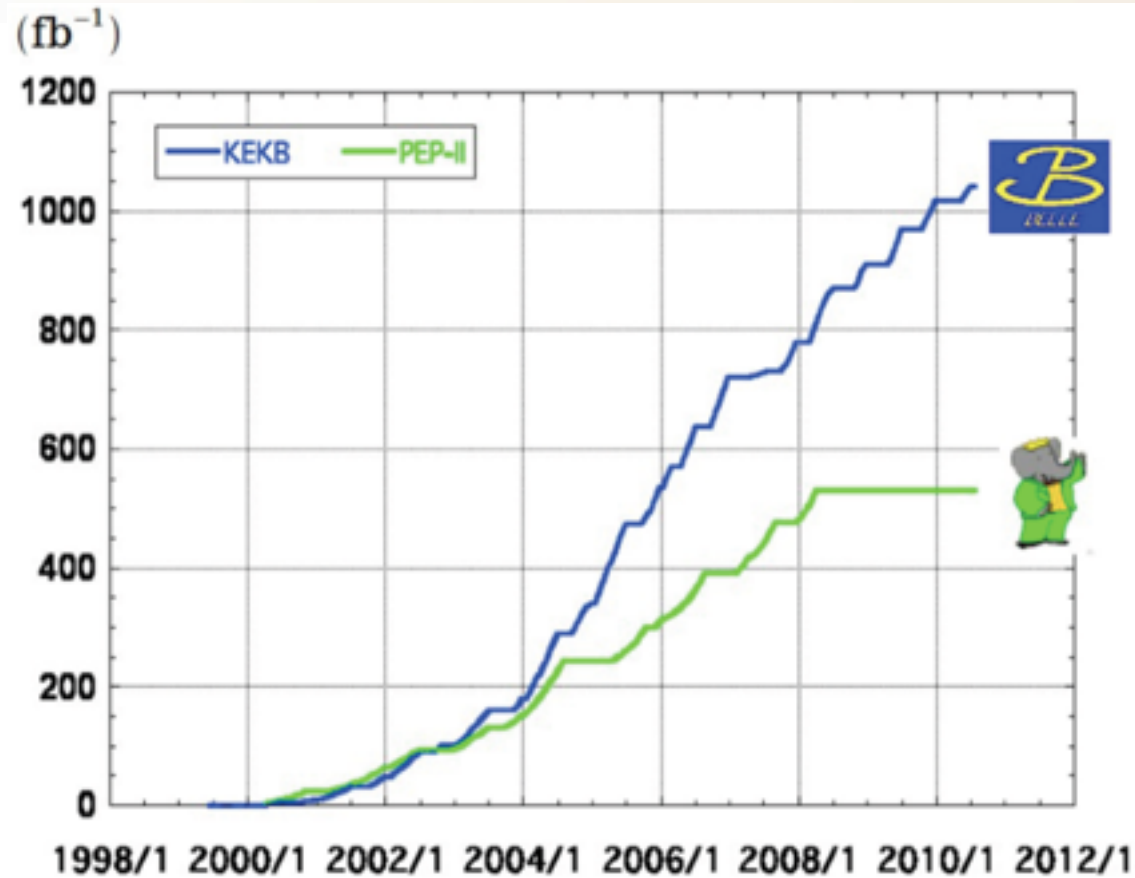
# What is a B-Factory?



The **main** thrust of a B-Factory is to produce lots and lots of coherent  $B^0\text{-}\bar{B}^0$  pairs in order to study CP-violation.

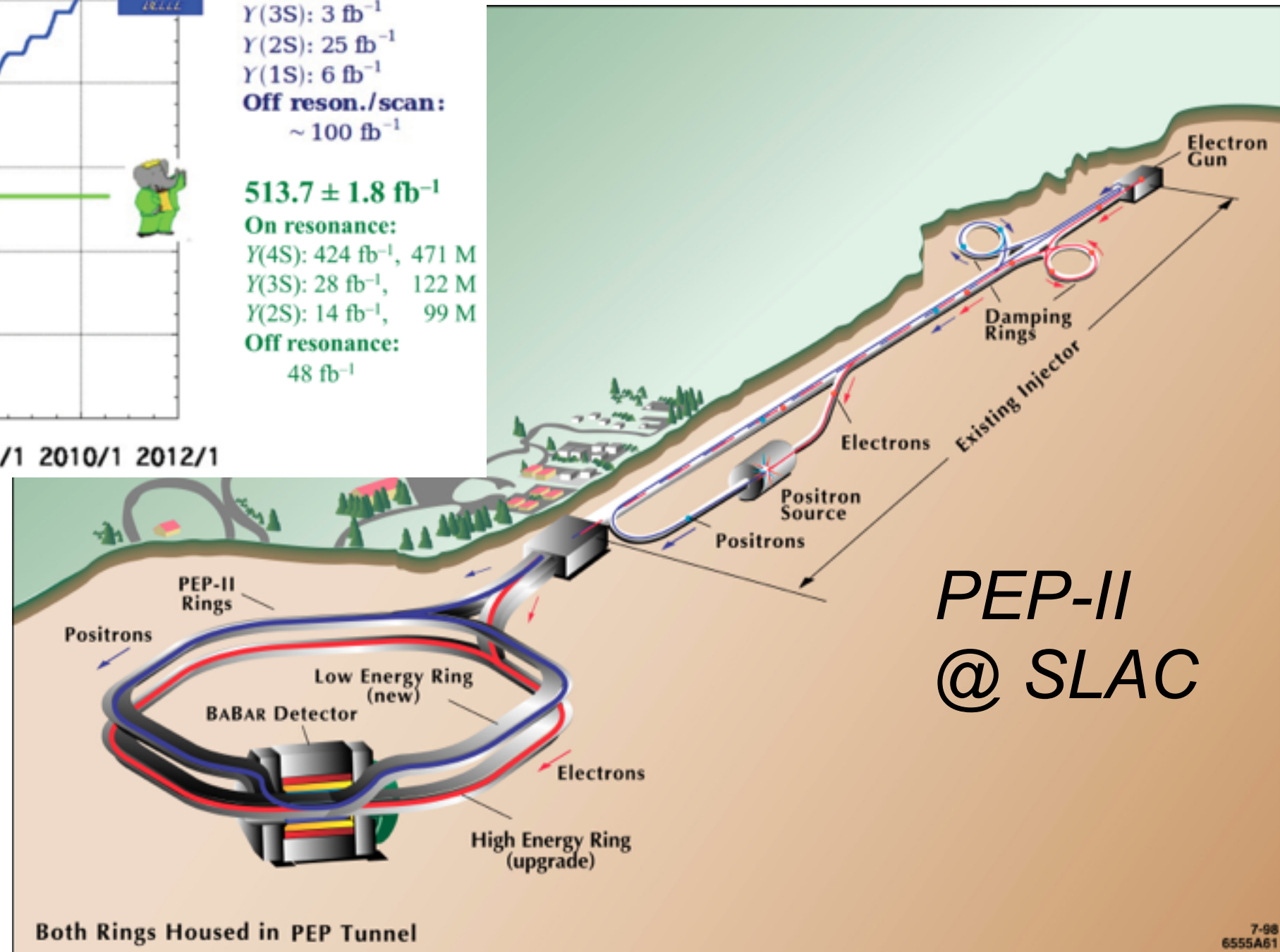


# Asymmetric B-Factories: PEP-II and KEKB



**> 1 ab<sup>-1</sup>**  
**On resonance:**  
 Y(5S): 121 fb<sup>-1</sup>  
 Y(4S): 711 fb<sup>-1</sup>  
 Y(3S): 3 fb<sup>-1</sup>  
 Y(2S): 25 fb<sup>-1</sup>  
 Y(1S): 6 fb<sup>-1</sup>  
**Off reson./scan:**  
 ~ 100 fb<sup>-1</sup>

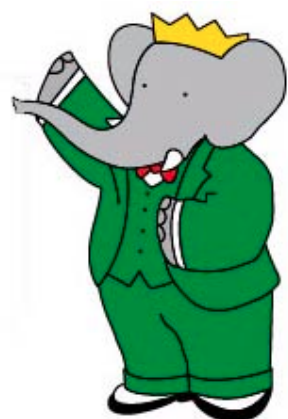
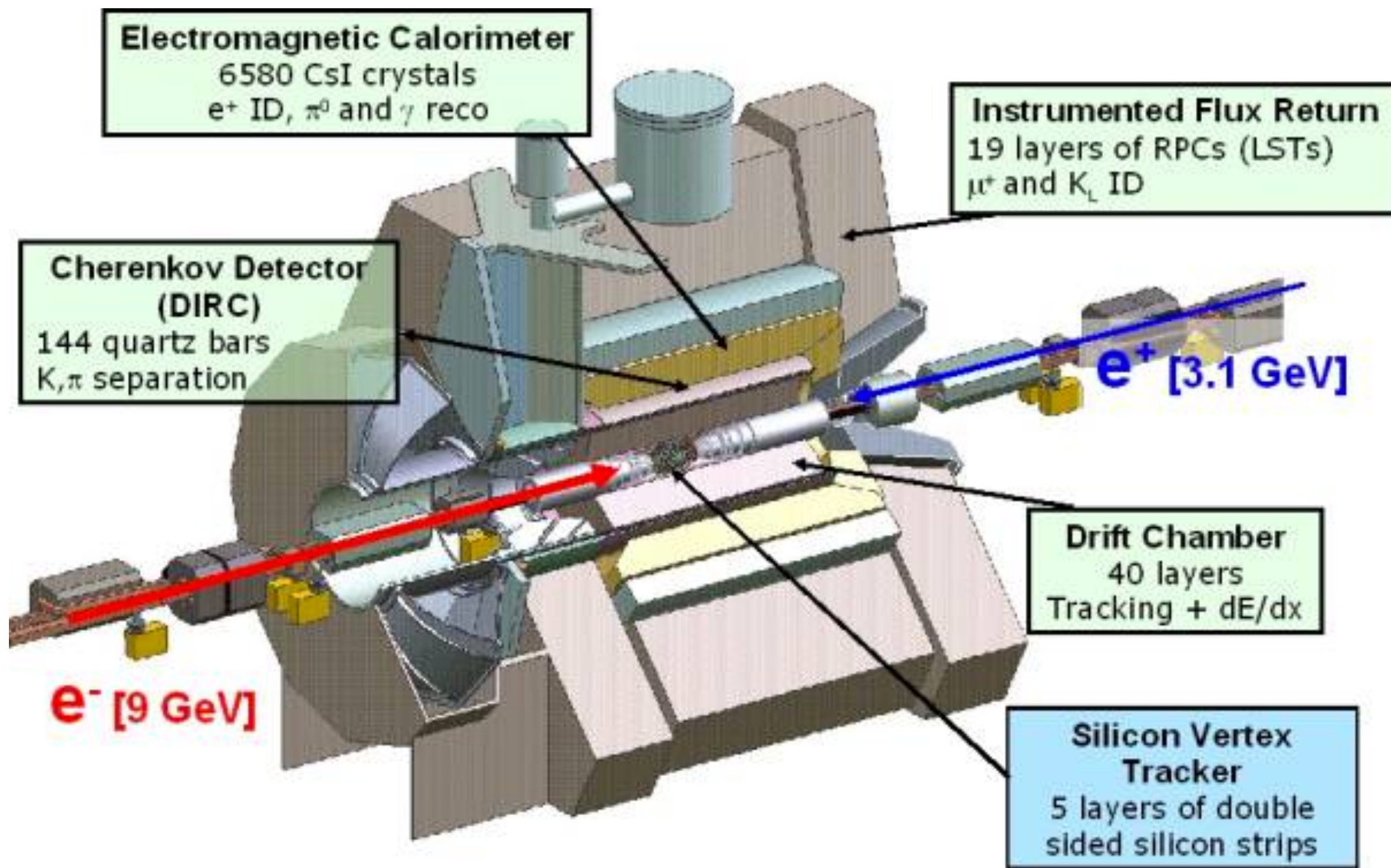
**513.7 ± 1.8 fb<sup>-1</sup>**  
**On resonance:**  
 Y(4S): 424 fb<sup>-1</sup>, 471 M  
 Y(3S): 28 fb<sup>-1</sup>, 122 M  
 Y(2S): 14 fb<sup>-1</sup>, 99 M  
**Off resonance:**  
 48 fb<sup>-1</sup>



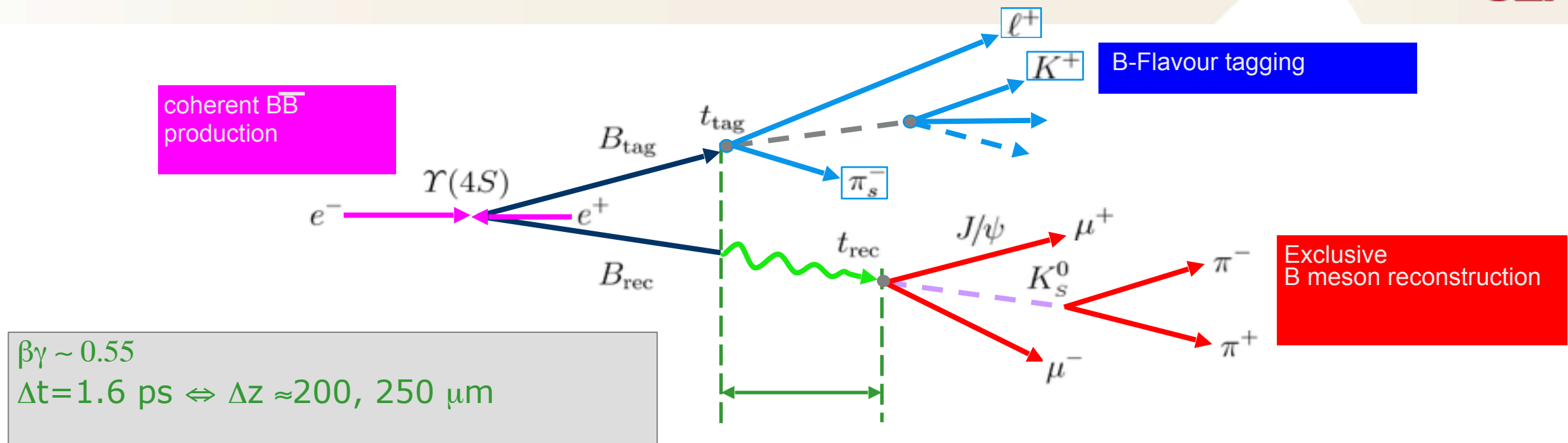
PEP-II  
 @ SLAC



# B-Factory Detectors (e.g. BaBar)



# Guiding the design...



## $e^+e^-$ B-Factories:

Typical error on  $\Delta t \sim 0.8 \text{ ps}$

The  $\Delta t$  resolution function obtained from high stat.  $B \rightarrow DX$

Events tagged using the charge of the leptons, kaons, pions

Effective tagging efficiency (including mistag-rate)  $\sim 30\%$

## LHCb:

Typical error on  $\Delta t \sim 0.05 \text{ ps}$

Effective tagging efficiency (including mistag-rate)  $\sim 5\%$

# Searches for New Physics in B-Factories

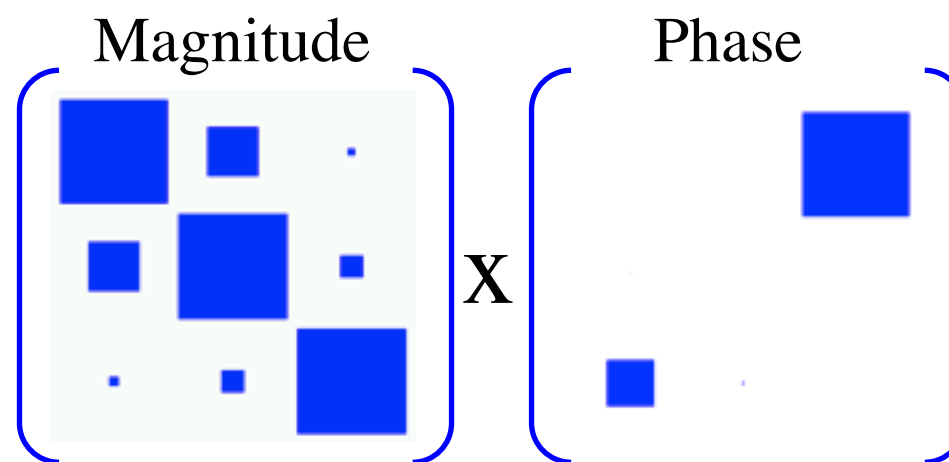
- Indirect searches:
  - This is what B-Factories are all about...

# Searches for New Physics in B-Factories

- Indirect searches:
  - This is what B-Factories are all about...

**The  
CKM  
Matrix**

$$V_{pq} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \propto$$



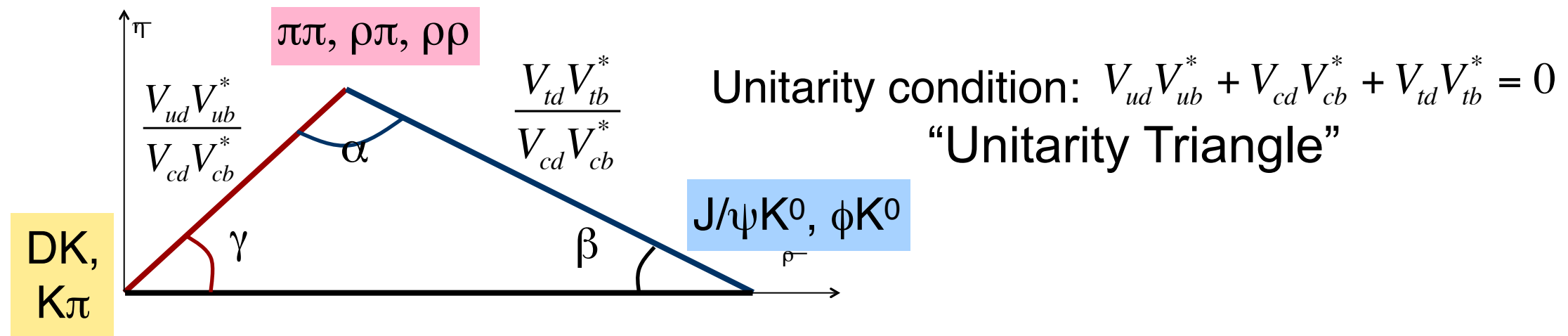
**CKM ↔ quarks**  
**PMNS ↔ leptons**

$$\approx \begin{pmatrix} 1 - \frac{\lambda^2}{2} & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix}$$

$$\lambda = \sin \theta_C = 0.22$$

# Searches for New Physics in B-Factories

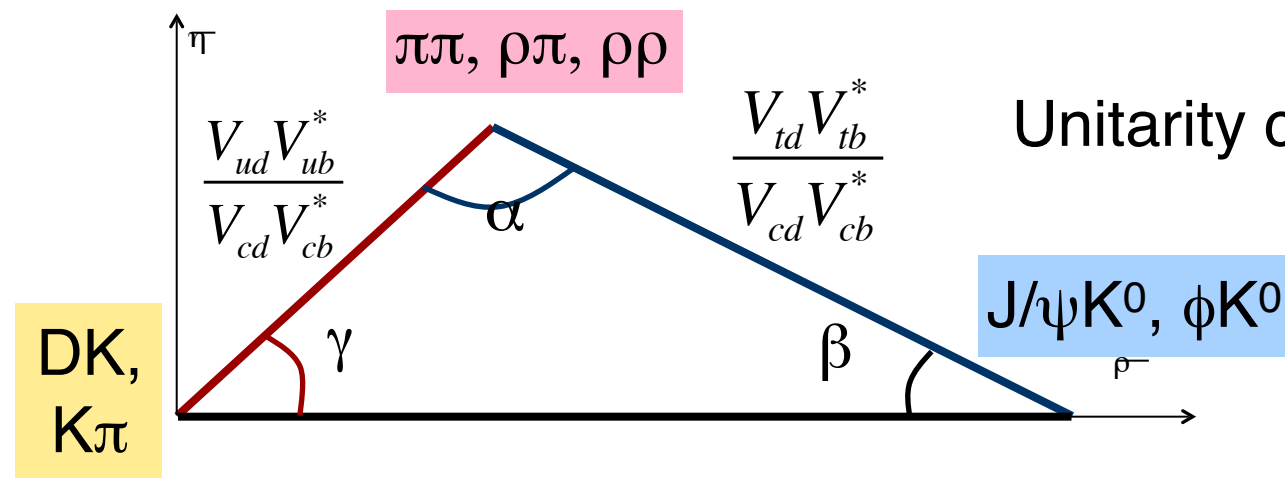
- Indirect searches:
  - This is what B-Factories are all about...





# Searches for New Physics in B-Factories

- Indirect searches:
  - This is what B-Factories are all about...

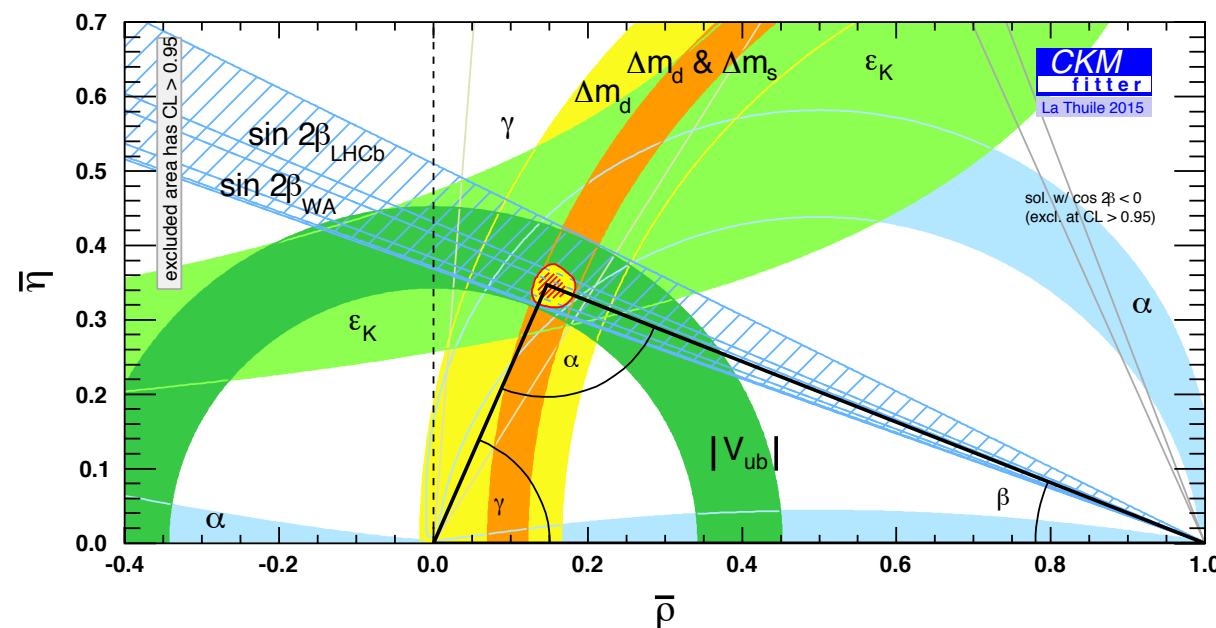


Unitarity condition:  $V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$   
 “Unitarity Triangle”

$$\bar{\rho} = 0.15 \pm 0.01$$

$$\bar{\eta} = 0.34 \pm 0.01$$

***This is a triumph of the HEP community!***



# Searches for New Physics in B-Factories

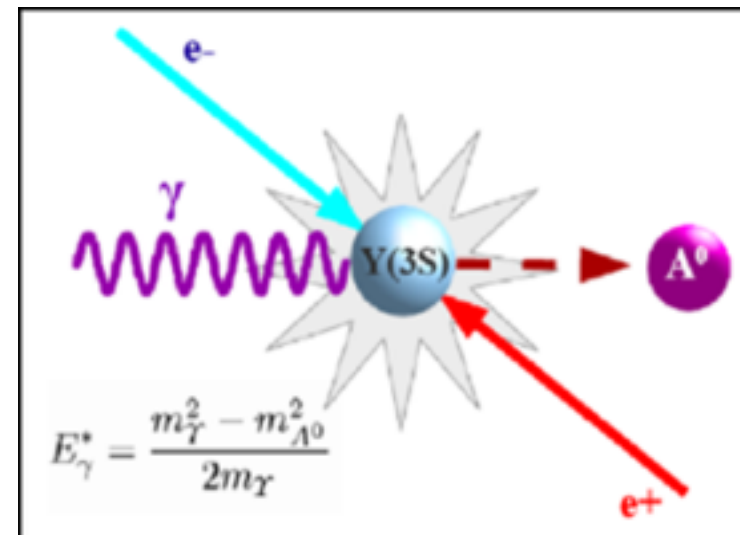
- Indirect searches:
  - This is what B-Factories are all about...
  - ...but not the only thing...
- Direct searches:
  - searches for light higgs ( $A^0$ )
  - searches for dark sectors
  - ...this is what the rest of the talk is about...

# Searches for the light CP-Odd Higgs: $A^0$

- The  $A^0$  is a CP-Odd Higgs particle that comes up in 2-Higgs Doublet Models (2HDM)...along with  $h^0$ ,  $H^0$ , and  $H^\pm$
- Its mass *could* be light ( $\sim$ MeV-GeV)  $\rightarrow$  search for it at B-Factories
- 2HDM can be seen as a generic extension of SM but also arise naturally in SUSY models (in particular the NMSSM)

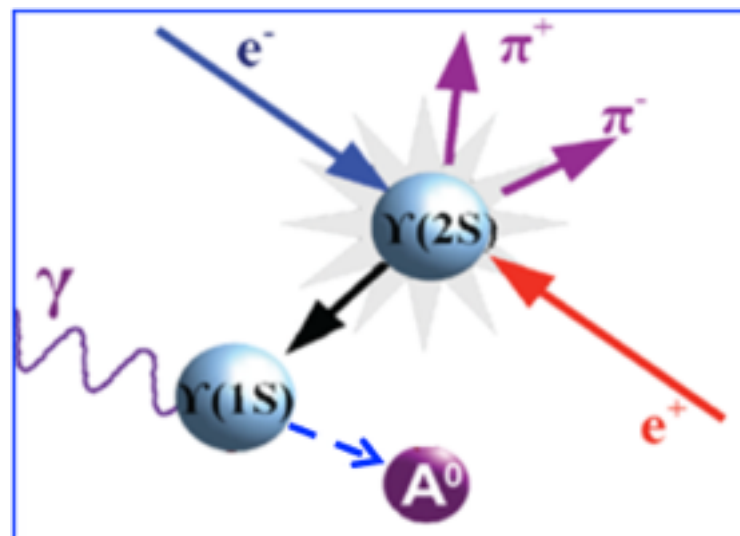
- B-Factories look for  $A^0$  in two ways:
  - directly from radiative decay of  $Y(nS)$ :

$$Y(nS) \rightarrow \gamma A^0 \rightarrow f \bar{f}$$



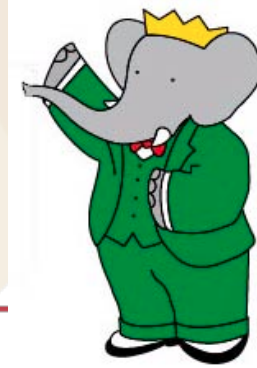
- from the cascade of  $Y(nS) \rightarrow Y(1S)$ :

$$Y(nS) \rightarrow Y(1S) \pi \pi \text{ (or } \gamma) \\ \downarrow \\ \gamma A^0 \rightarrow f \bar{f}$$



- ...looking for the cascade from  $Y(nS)$  greatly reduces the background (ensures you have a  $b\bar{b}$  state)

# Search for $A^0 \rightarrow c\bar{c}$

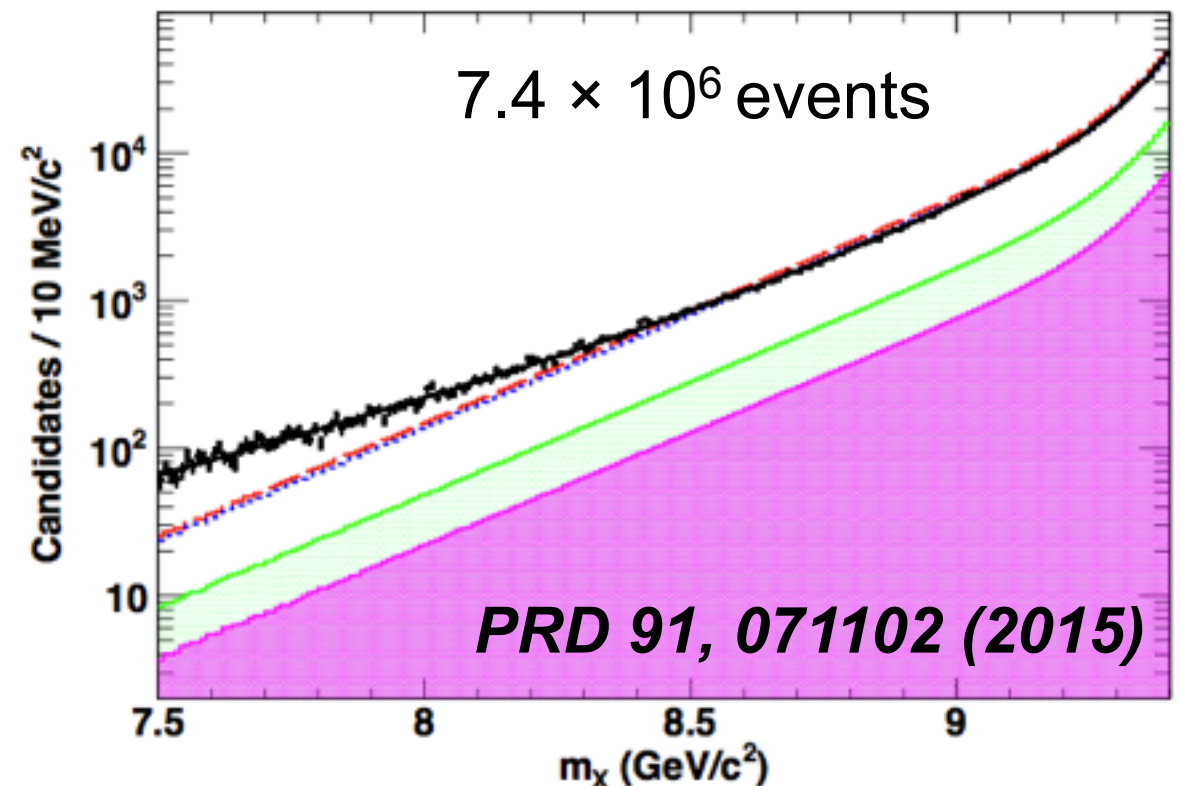
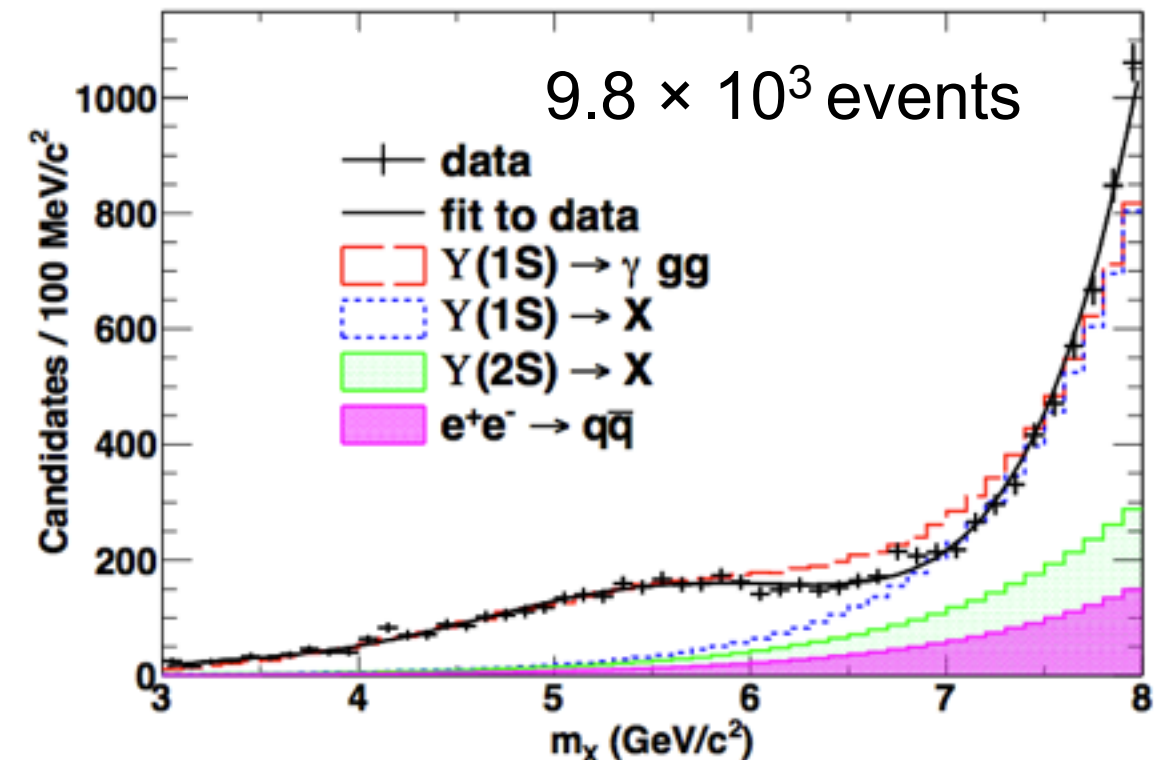


SLAC

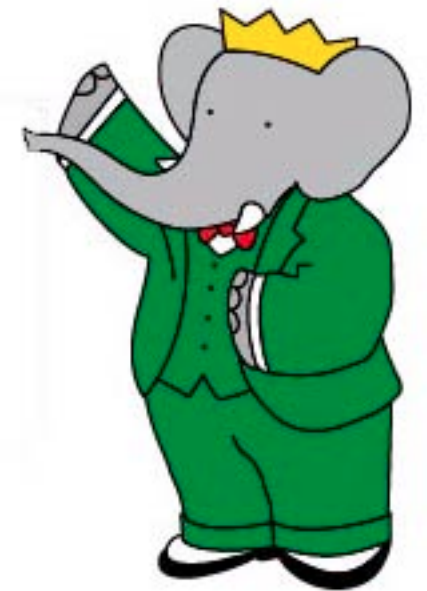
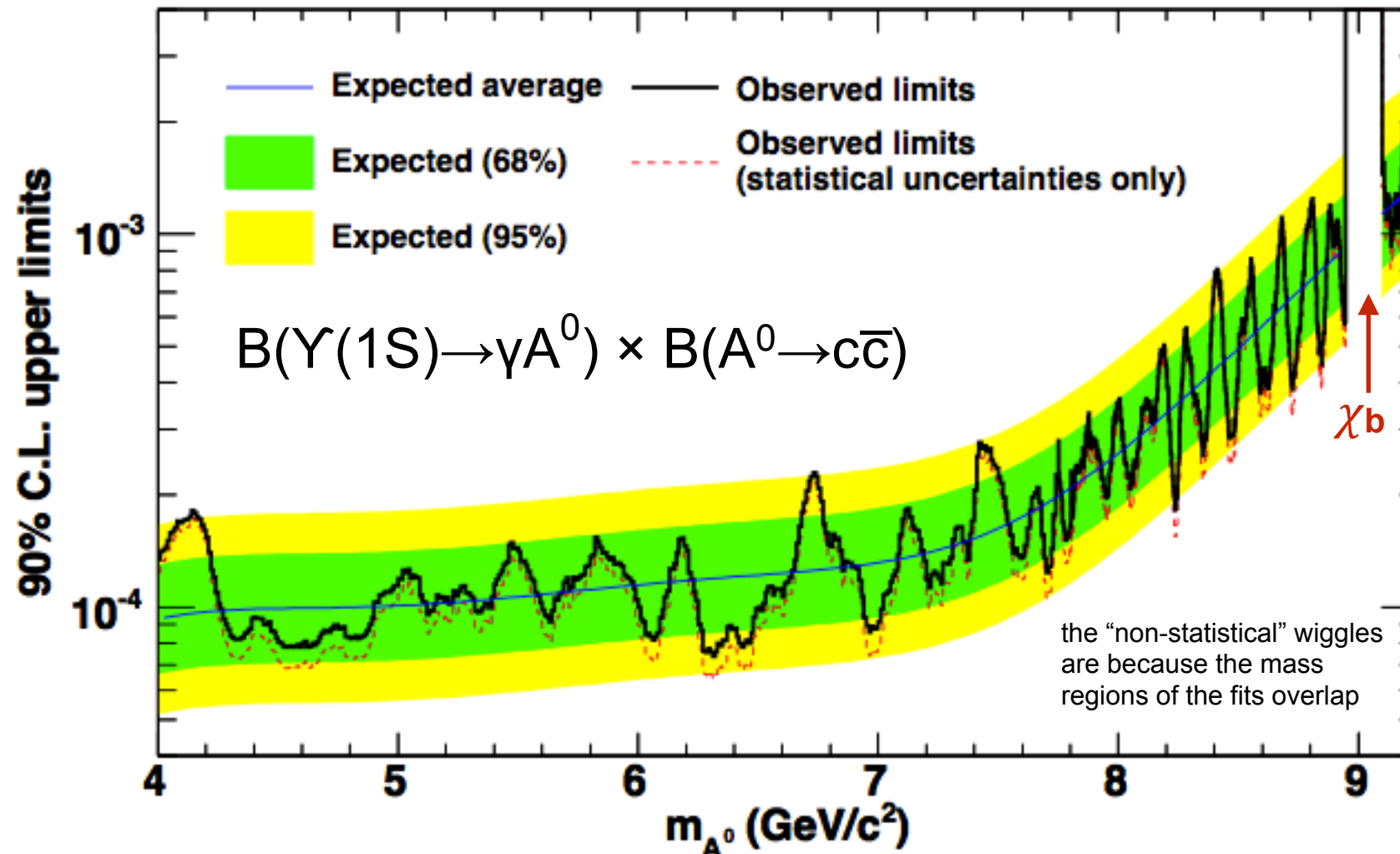
- 13.6 fb<sup>-1</sup> on  $\Upsilon(2S)$  resonance:
  - $(98.3 \pm 0.9) \times 10^6$   $\Upsilon(2S)$  mesons
  - $(17.5 \pm 0.3) \times 10^6$   $\pi^+\pi^- \Upsilon(1S)$  decays
- Reconstruct at least **one** D-meson in one of:
  - $D^0 \rightarrow K^-\pi^+, K^-\pi^+\pi^-\pi^+, K_s\pi^+\pi^-$
  - $D^+ \rightarrow K^-\pi^+\pi^+$
  - $D^{*+} \rightarrow \pi^+D^0 (\rightarrow K^-\pi^+\pi^0)$
  - ...fraction of possible final states included in efficiency (JETSET)...

$$m_R^2 = M_{\Upsilon(2S)}^2 + m_{\pi\pi}^2 - 2M_{\Upsilon(2S)}E_{\pi\pi}$$

- Dataset split into high- and low-mass regions because of very different background shapes



# Search for $A^0 \rightarrow c\bar{c}$ upper limits



- Perform a scan in  $M_X$  on the plots on previous page..
  - step along in mass fitting the signal contribution to a Crystal Ball™ function and the background to a second-order polynomial (floating)
- No significant signal observed (either with or without accounting for the trials factor), so an upper limit is calculated

**PRD 91, 071102 (2015)**



# A comprehensive suite of searches...

|  |                          |
|--|--------------------------|
| $Y(3S) \rightarrow \gamma A^0, A^0 \rightarrow \text{invisible}$   | arXiv:0808.0017 [hep-ex] |
| $Y(3S) \rightarrow \gamma A^0, A^0 \rightarrow \tau^+\tau^-$       | PRL 103, 181801 (2009)   |
| $Y(2S, 3S) \rightarrow \gamma A^0, A^0 \rightarrow \mu^+\mu^-$     | PRL 103, 081801 (2009)   |
| $Y(2S, 3S) \rightarrow \gamma A^0, A^0 \rightarrow \text{hadrons}$ | PRL 107, 221801 (2011)   |
| $Y(1S) \rightarrow \gamma A^0, A^0 \rightarrow \text{invisible}$   | PRL 107, 021804 (2011)   |
| $Y(1S) \rightarrow \gamma A^0, A^0 \rightarrow \mu^+\mu^-$         | PRD 87, 031102 (2013)    |
| $Y(1S) \rightarrow \gamma A^0, A^0 \rightarrow \tau^+\tau^-$       | PRD 88, 031102 (2013)    |
| $Y(1S) \rightarrow \gamma A^0, A^0 \rightarrow gg, ss$             | PRD 88, 031701 (2013)    |
| $Y(1S) \rightarrow \gamma A^0, A^0 \rightarrow cc$                 | PRD 91, 071102 (2015)    |

# New interactions beyond the SM?



**vector**



$$\frac{1}{2} \epsilon F_{\mu\nu}^Y F'^{\mu\nu}$$

**dark photon**

**Higgs**



$$\epsilon_h |h|^2 |\phi|^2$$

**dark scalar**

**neutrino**



$$\epsilon_\nu (hL)\psi$$

**sterile  
neutrino**

**axion**



$$\frac{1}{f_a} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

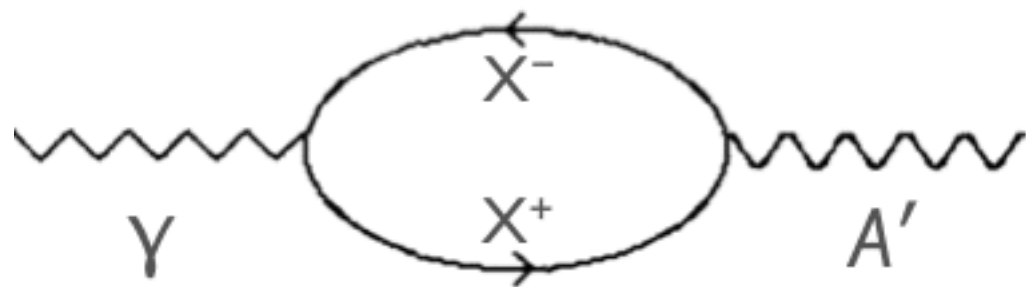
**ALPSs**

# Kinetic mixing from the vector portal

*an old idea: if there is an additional  $U(1)$  symmetry in nature, there will be mixing between the photon and the new gauge boson*

Holdom, Phys. Lett B 166, 1986

$$\mathcal{L}_{U(1)'} = -\frac{1}{4} V_{\mu\nu}^2 - \boxed{\frac{\epsilon}{2} V_{\mu\nu} F^{\mu\nu}} + |D_\mu \phi|^2 - V(\phi)$$

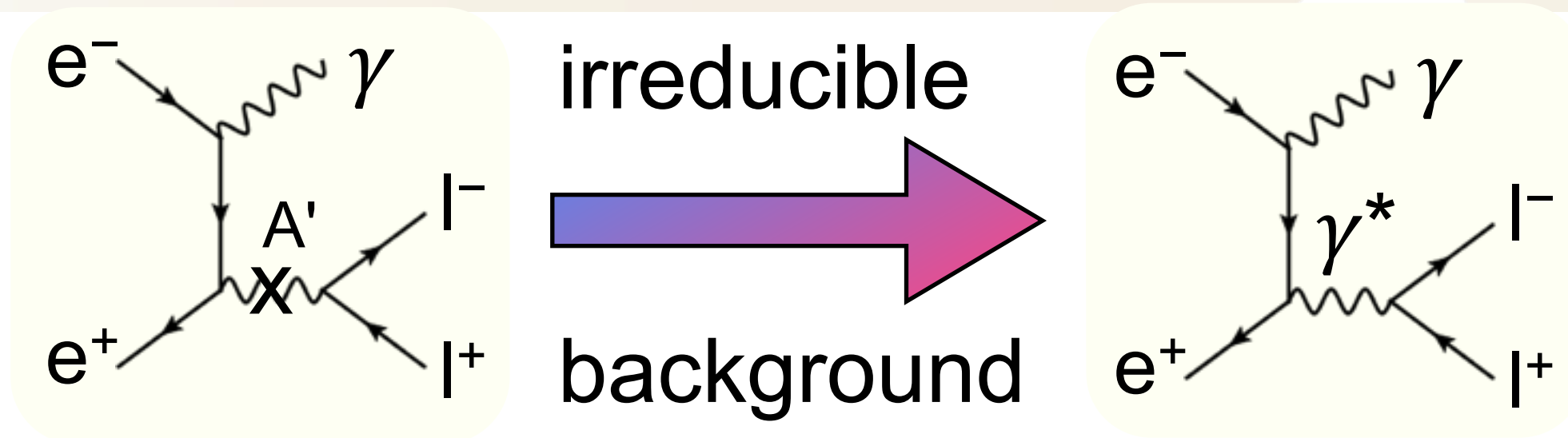


*Kinetic Mixing term*

NOTE:  $A^0$  (light higgs)  $\neq$   $A'$  (dark photon) !!!

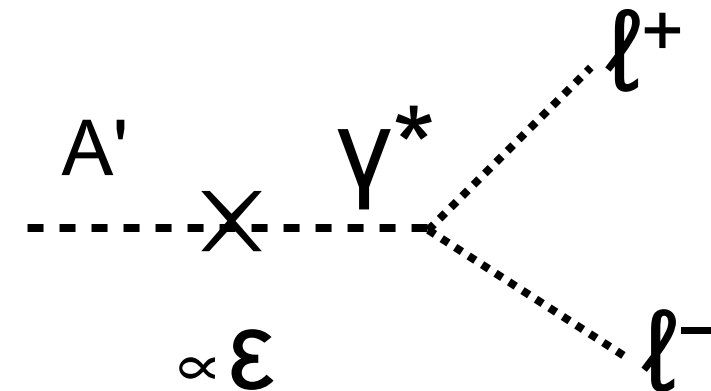
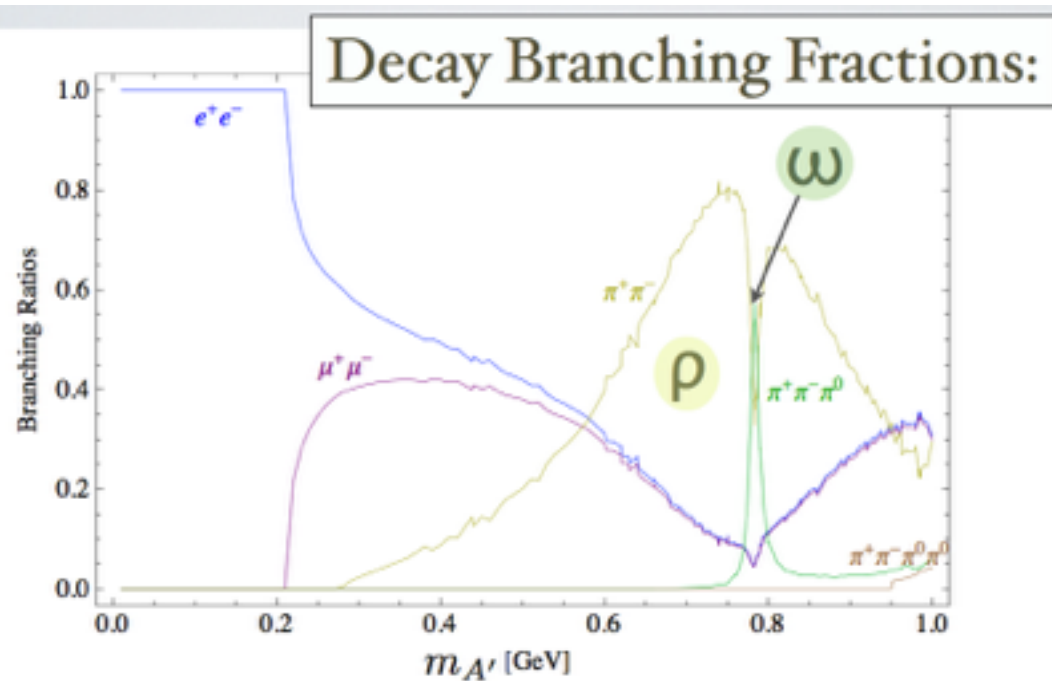
- extremely general conclusion...even arises from broken symmetries
- gives coupling of normal charged matter to the new “*dark photon*”  $q=\epsilon e$
- dark photon could get mass via a “dark higgs” mechanism (or elsewhere)

# Dark photon production at $e^+e^-$ colliders

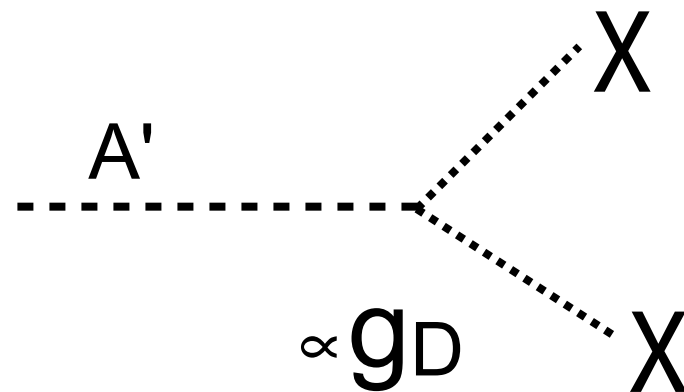


- The flavor factories (BaBar, Belle(-II), KLOE) are great places to look for dark photons:
  - high luminosity
    - not quite high enough to get to displaced-vertex regions
  - excellent acceptance & PID
    - also makes them good places to look for non-minimal dark sector ( $A'h$ , non-Abelian, etc)
  - good mass resolution  $\sim$  few MeV
- The primary *signature* is very similar to the  $A^0$  searches...look for a bump over a continuum spectrum...

# Decays of the dark photon



$\Rightarrow$  **BUT**, if there is a state,  $\chi$ , with dark charge and is lighter than  $m(A')/2$

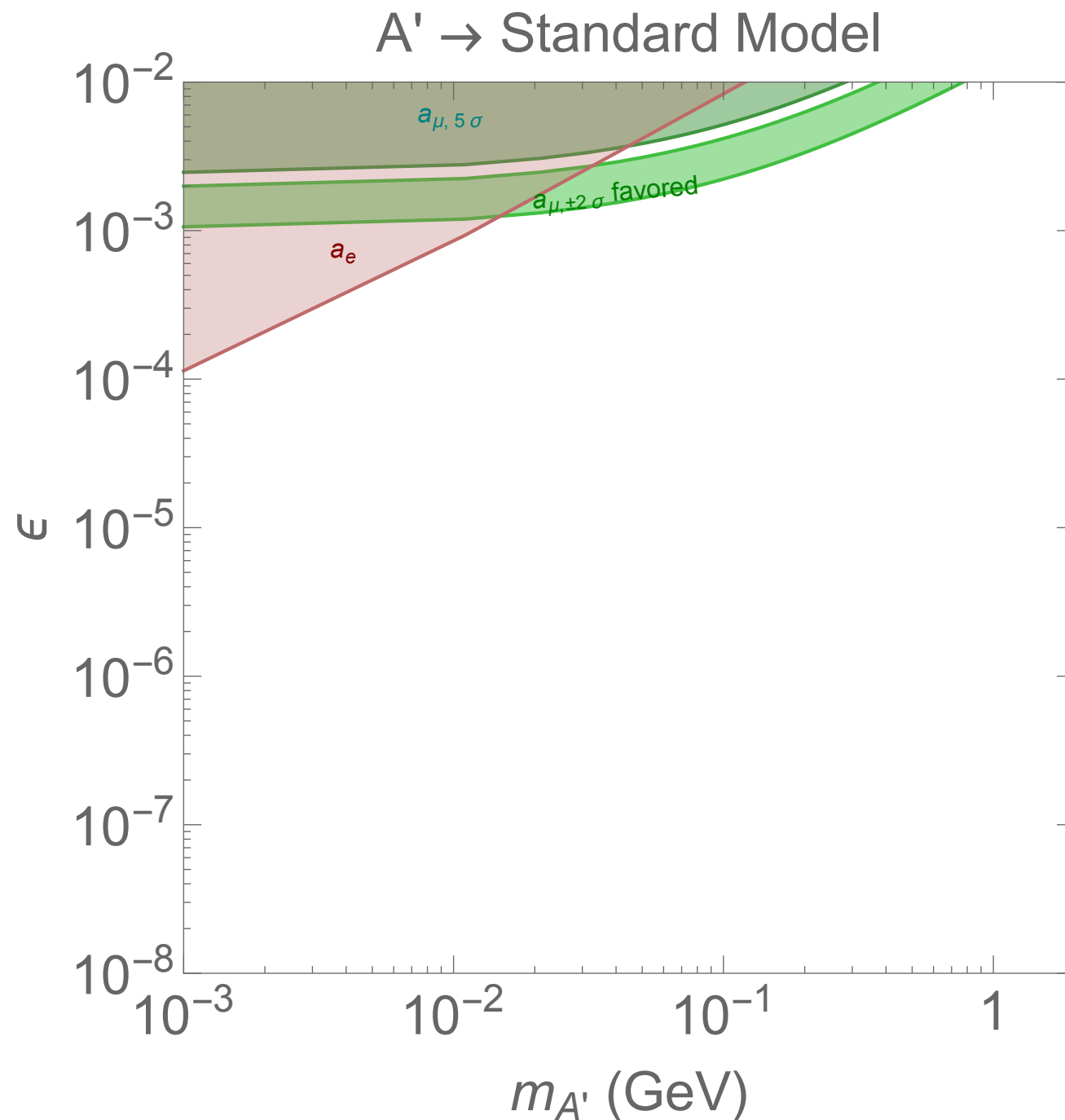


$\chi = \text{light, hidden sector dark matter}$

...unless  $g_D \ll \epsilon$ ,  $A'$  decays to dark sector will dominate



# Search phase-space

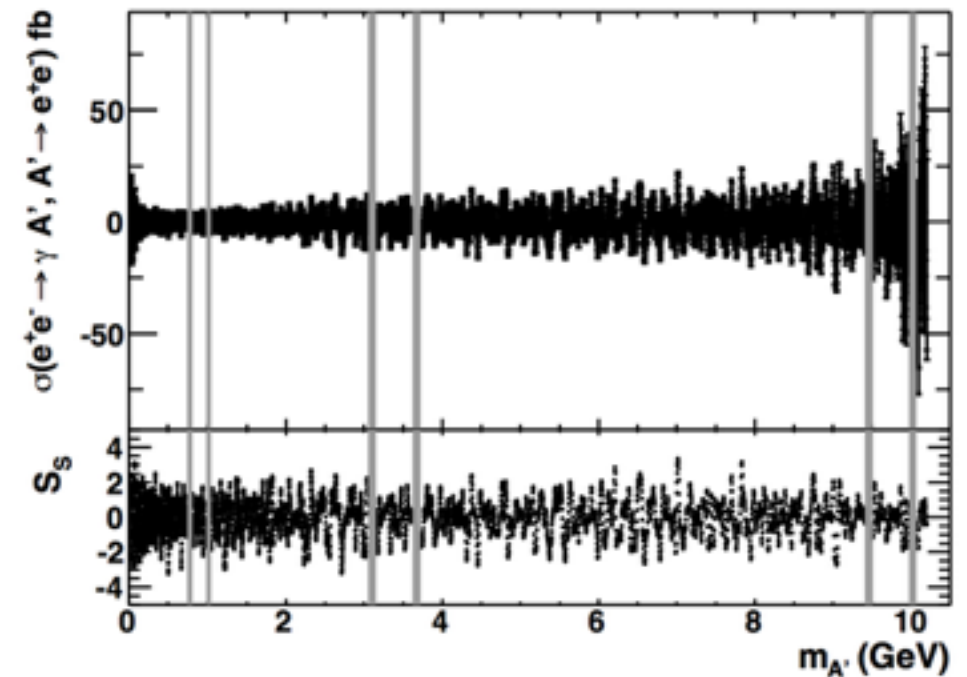


Start off with 2  
indirect constraints:  
g-2 of electron & muon

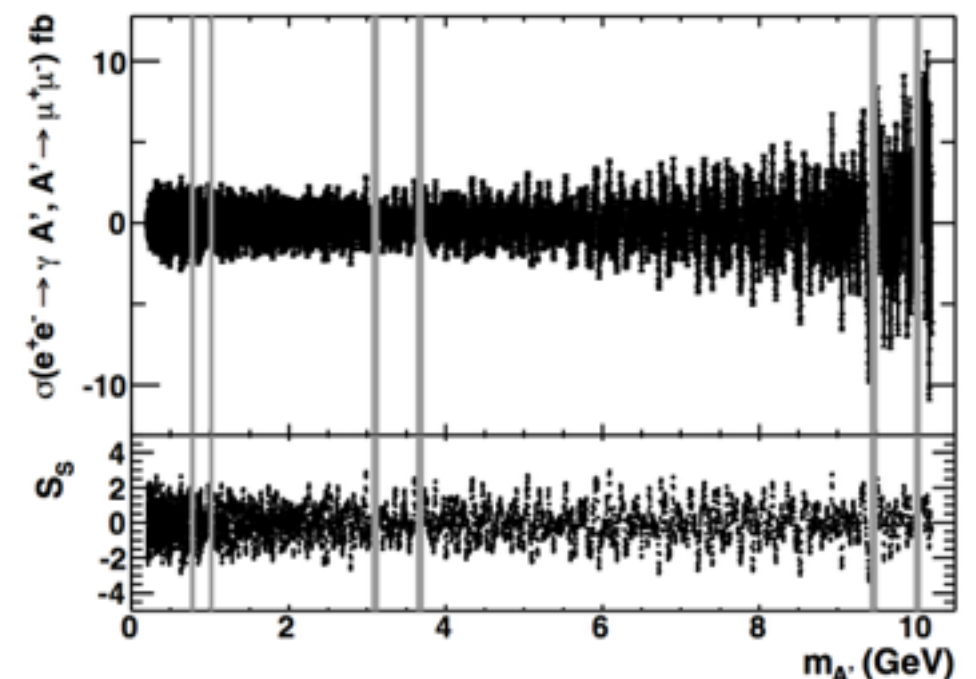
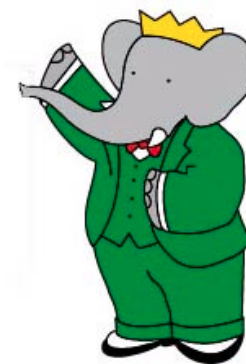
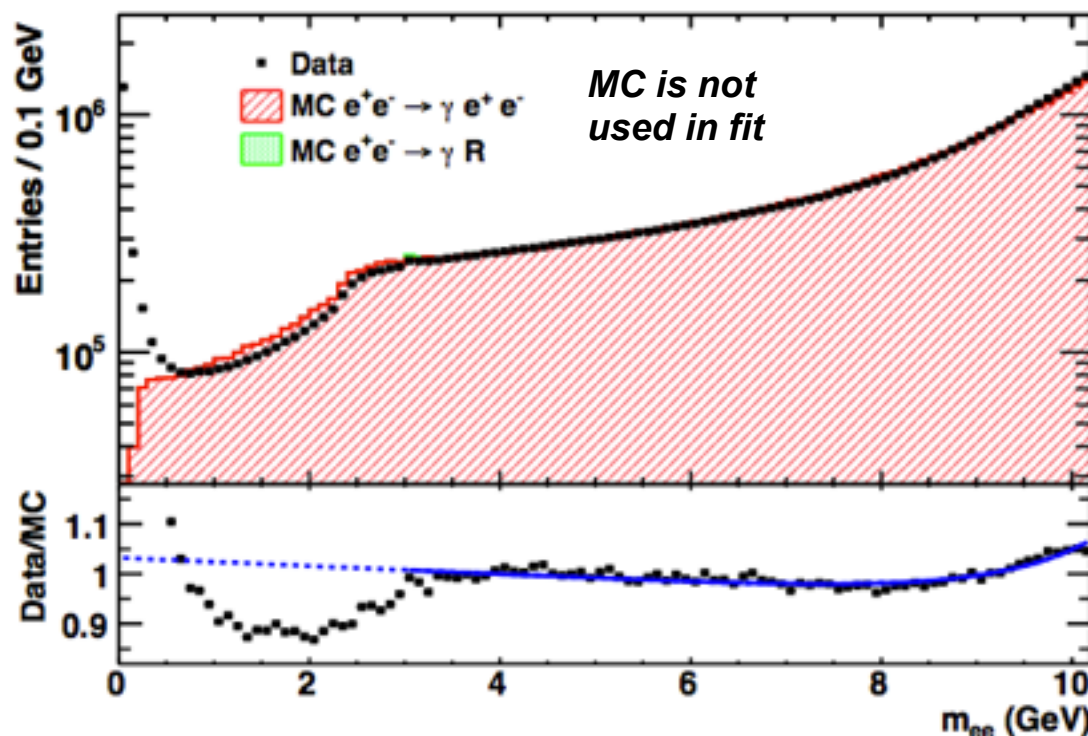
**But!** A  $\sim 3\sigma$  anomaly in  $g_{\mu}-2$   
...interesting territory!

# Search for dark photon in $e^+e^- \rightarrow \gamma \ell^+ \ell^-$

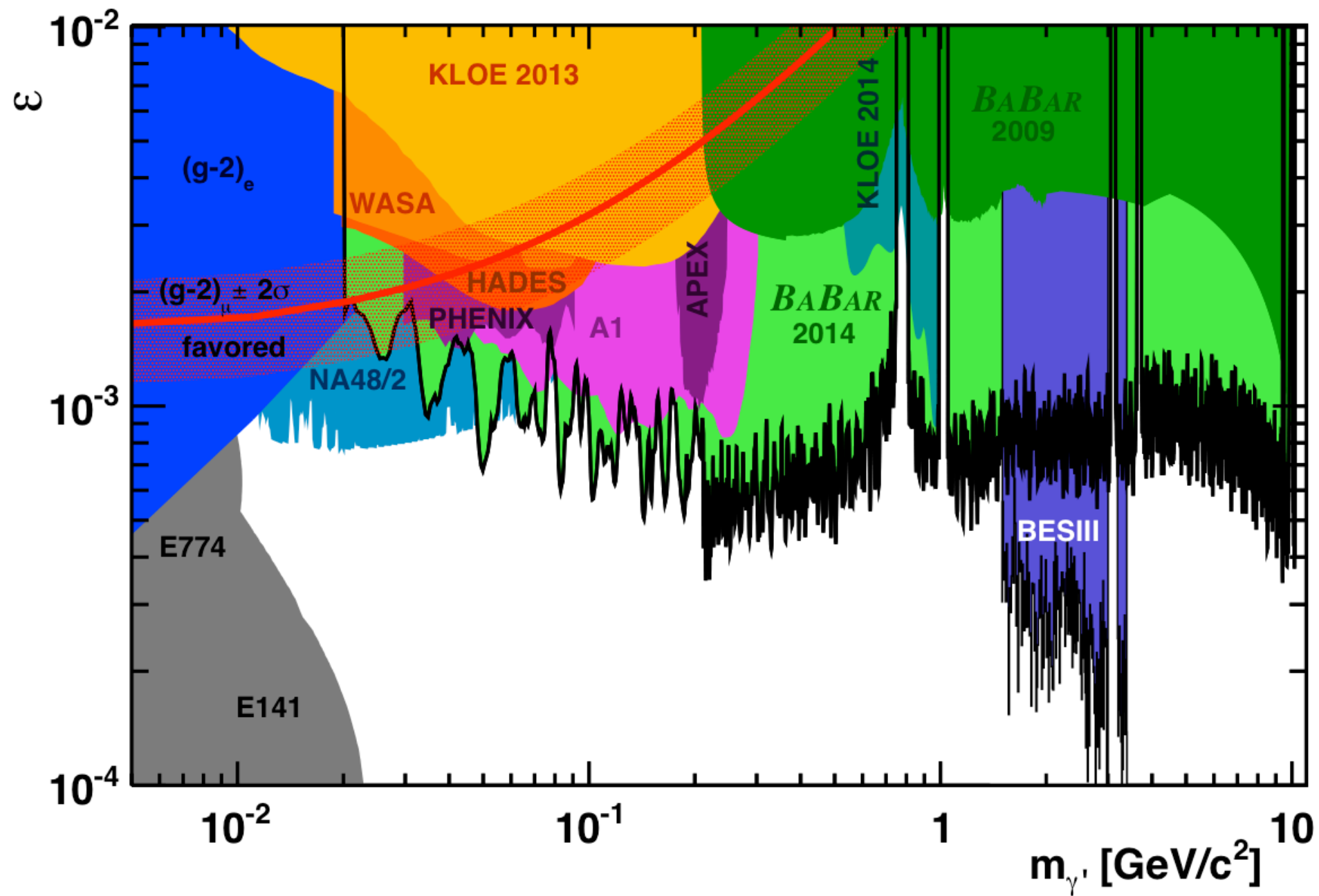
- this analysis uses the **FULL** BaBar dataset
  - 514 fb<sup>-1</sup> of data...Y(4S), Y(3S), Y(2S) & off resonance
- fully reconstruct the  $\gamma$  and both  $\ell$ 's, either  $e^+e^-$  or  $\mu^+\mu^-$ 
  - some care is taken to remove  $\gamma \rightarrow e^+e^-$  conversions at low  $e^+e^-$  mass
- scan in  $\ell^+\ell^-$  mass, fitting signal (non-parametric) + background (polynomial) in small mass steps



arXiv:1406.2980

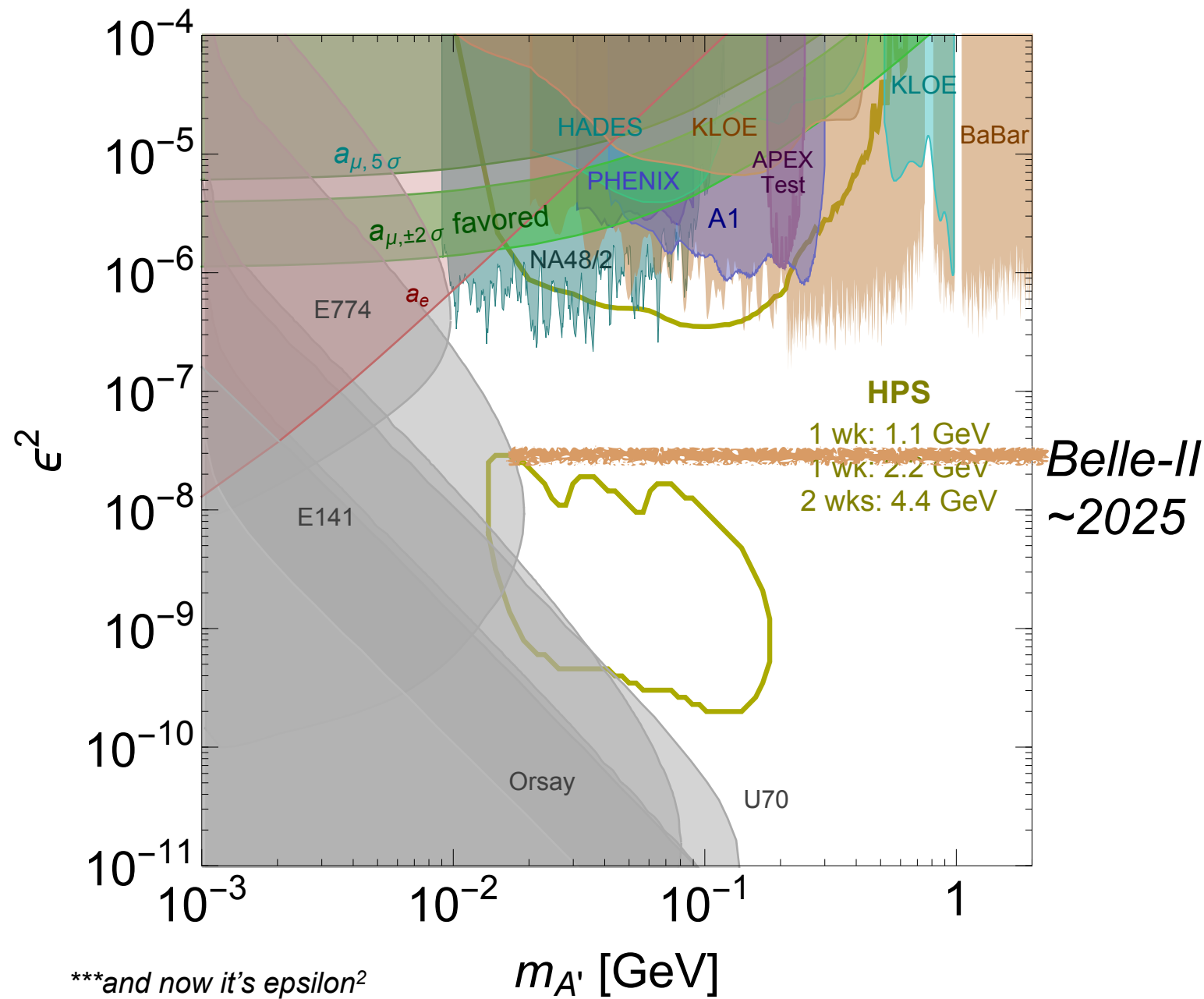


# Limits on dark photon coupling from BaBar



Plot courtesy of B. Echenard

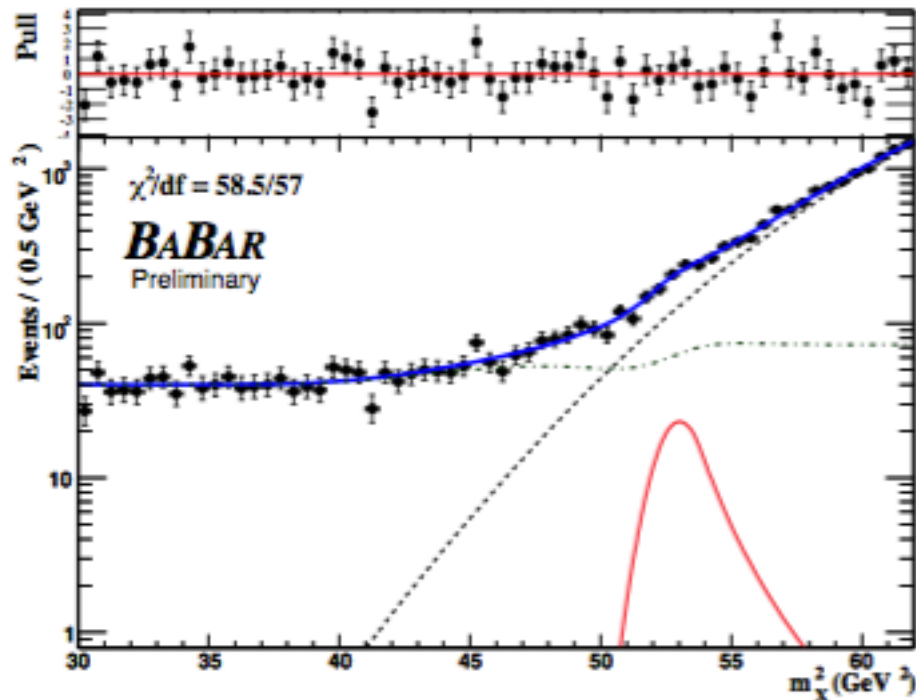
# Current dark photon limits (+ HPS)



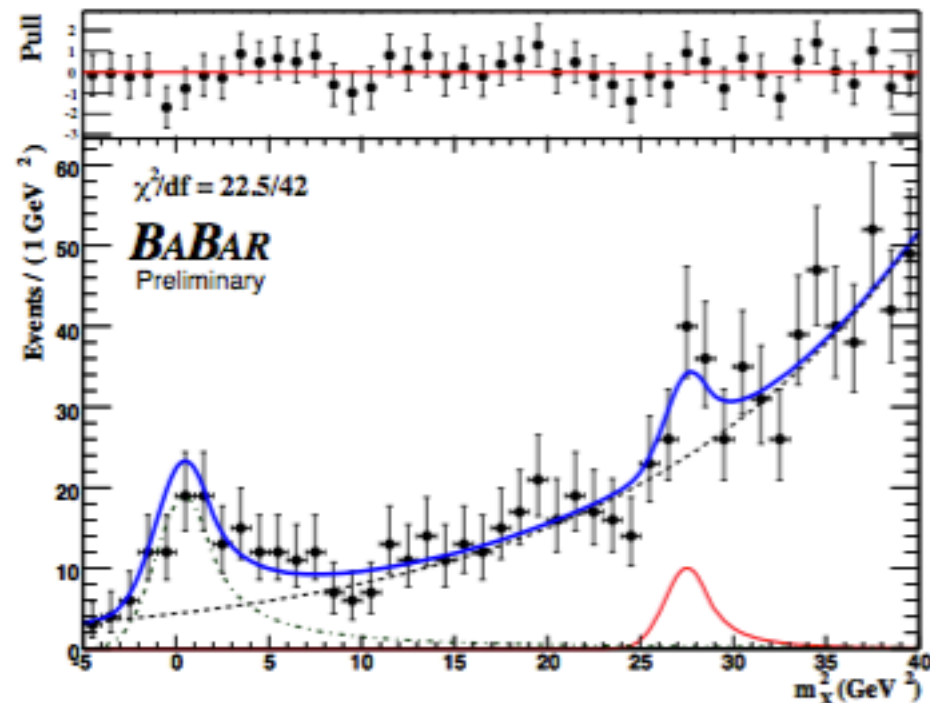
A number of other experiments have probed the g-2 region as well (with null results)...

...but still a large amount of phase space to probe and (near) future experiments will search a large fraction of it. Takashi Maruyama will talk about this in detail in a later talk.

# $A^0$ (and dark photon) $\rightarrow$ Invisible decays



[arxiv.org:0808.0017v1](http://arxiv.org:0808.0017v1)

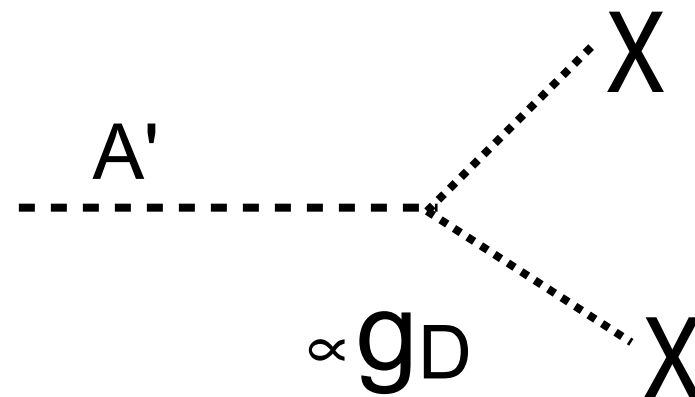
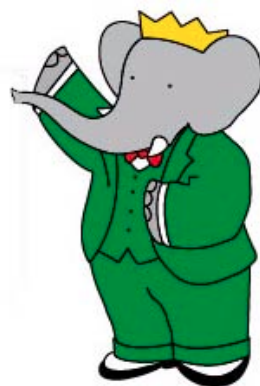


- Search for CP-Odd Higgs by looking for:  
 $Y(3S) \rightarrow \gamma A^0$   
↓  
nothing
- reconstruct the missing mass with:

$$E_\gamma^* = \frac{m_\gamma^2 - m_{A^0}^2}{2m_\gamma}$$

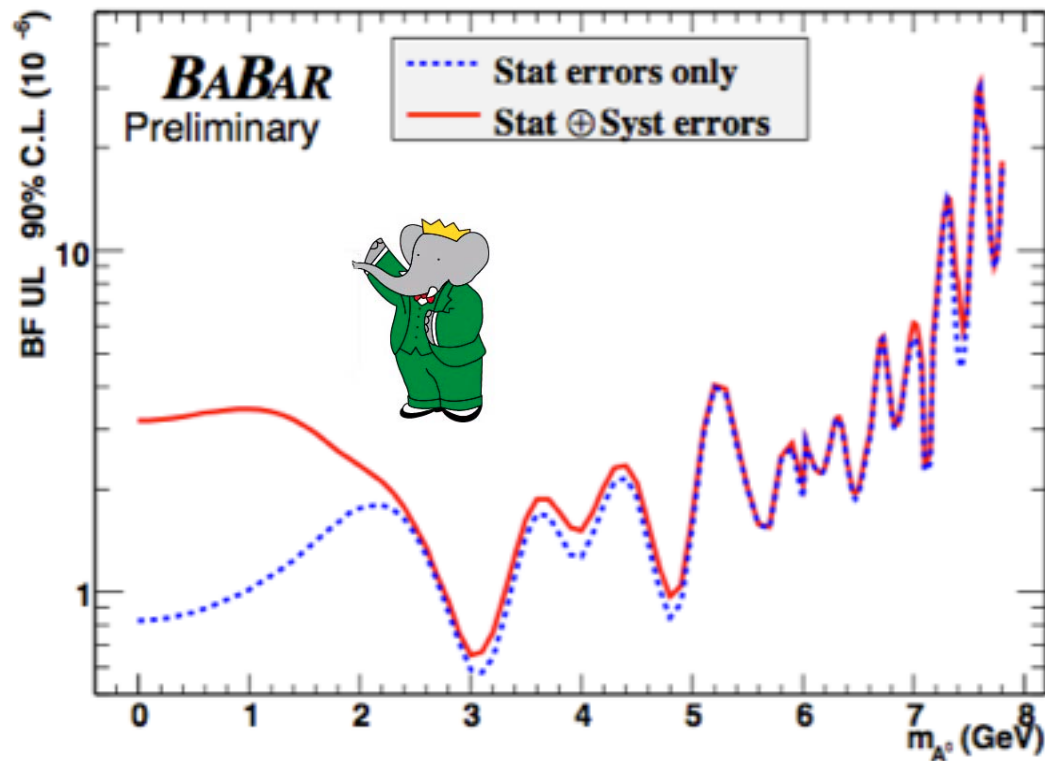
$$m_X^2 \equiv m_{\gamma(3S)}^2 - 2E_\gamma^* m_{\gamma(3S)}$$

- Scan in the missing mass, just like the others..
- This can also be used to look for invisible  $A'$  decays!



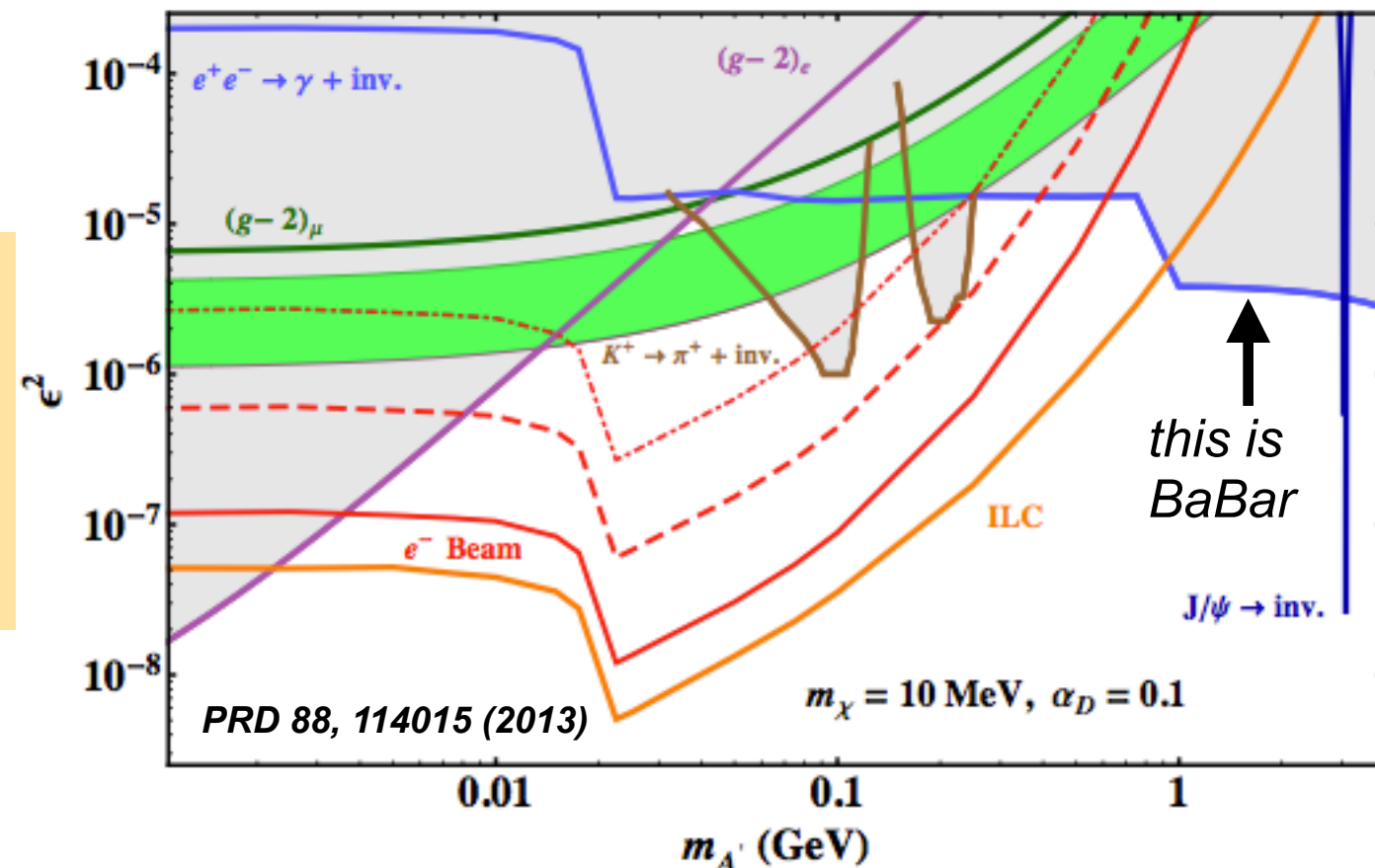


# Dark photon $\rightarrow$ invisible constraints



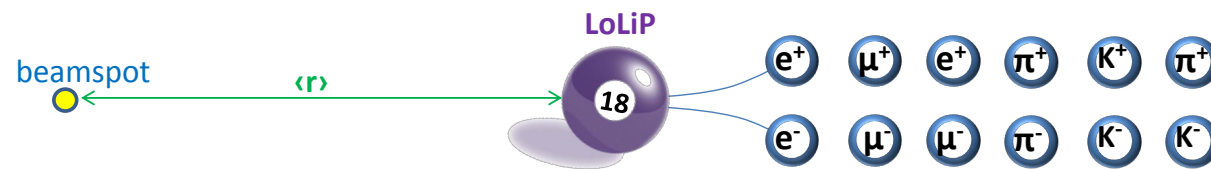
The constraint on  $A' \rightarrow \chi\chi$  is more complicated than the visible searches...depends on:  $\epsilon$ ,  $m(A')$ ,  $\alpha_D$ ,  $m(\chi)$   $\rightarrow$  look at projections

There are a few proposals being developed (in various stages of reality) for experiments designed to look for these sorts of decays and light dark matter in general...it's an exciting time!



# Search for displaced decays @ BaBar

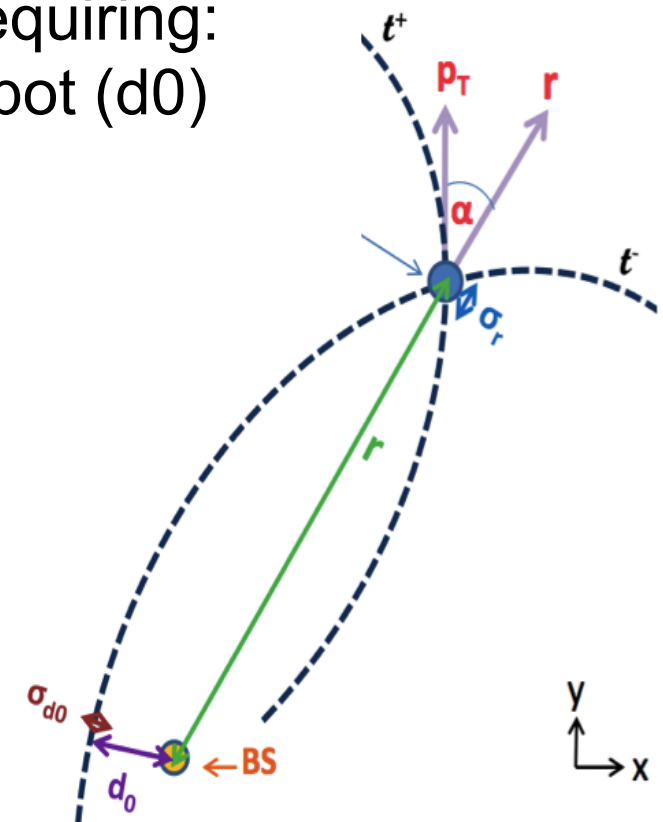
- Many BSM models include possible meta-stable states...i.e. particles that decay on the lab scale:  $\sim 100\mu$  - few m distances



- BaBar search vertexed  $f^+f^-$  tracks (where  $f = K, \pi, e, \text{ or } \mu$ ) requiring:
  - each track was inconsistent with originating from beam spot ( $d_0$ )
  - no tracker hits before the vertex
  - the vertex position wasn't near likely converter material (beampipe, support tube, drift chamber wall)

## Upper limits vs $m$ of

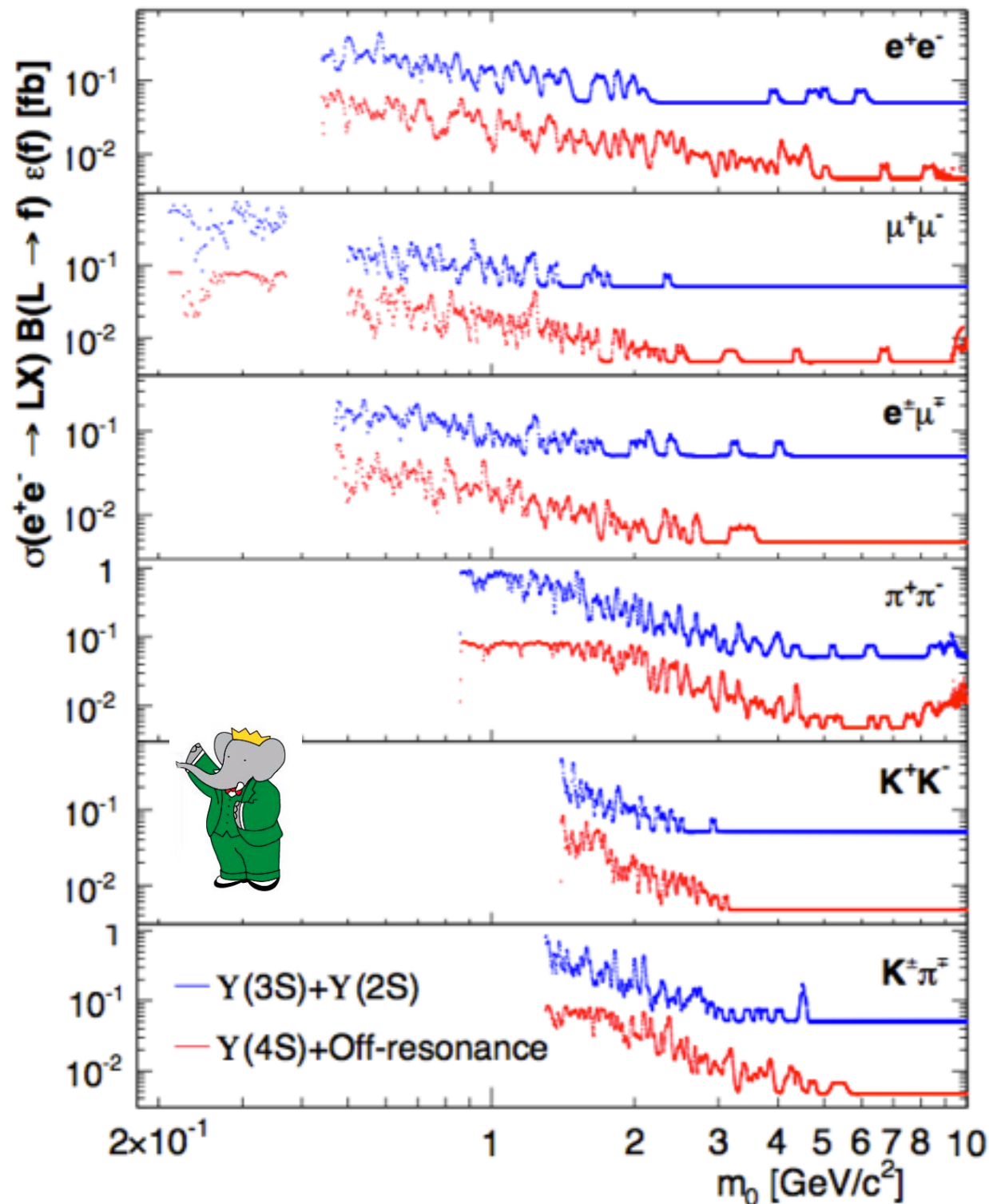
- > Model independent:  $\sigma(e^+e^- \rightarrow L) \times \text{BF}(L \rightarrow f) \times \epsilon(f)$   
at or near  $Y(4S)$ , at  $Y(2S, 3S)$   
giving tables of  $\epsilon(m, p_t, ct)$
- > Model dependent:  $\text{BF}(B \rightarrow X_s L) \times \text{BF}(L \rightarrow f)$



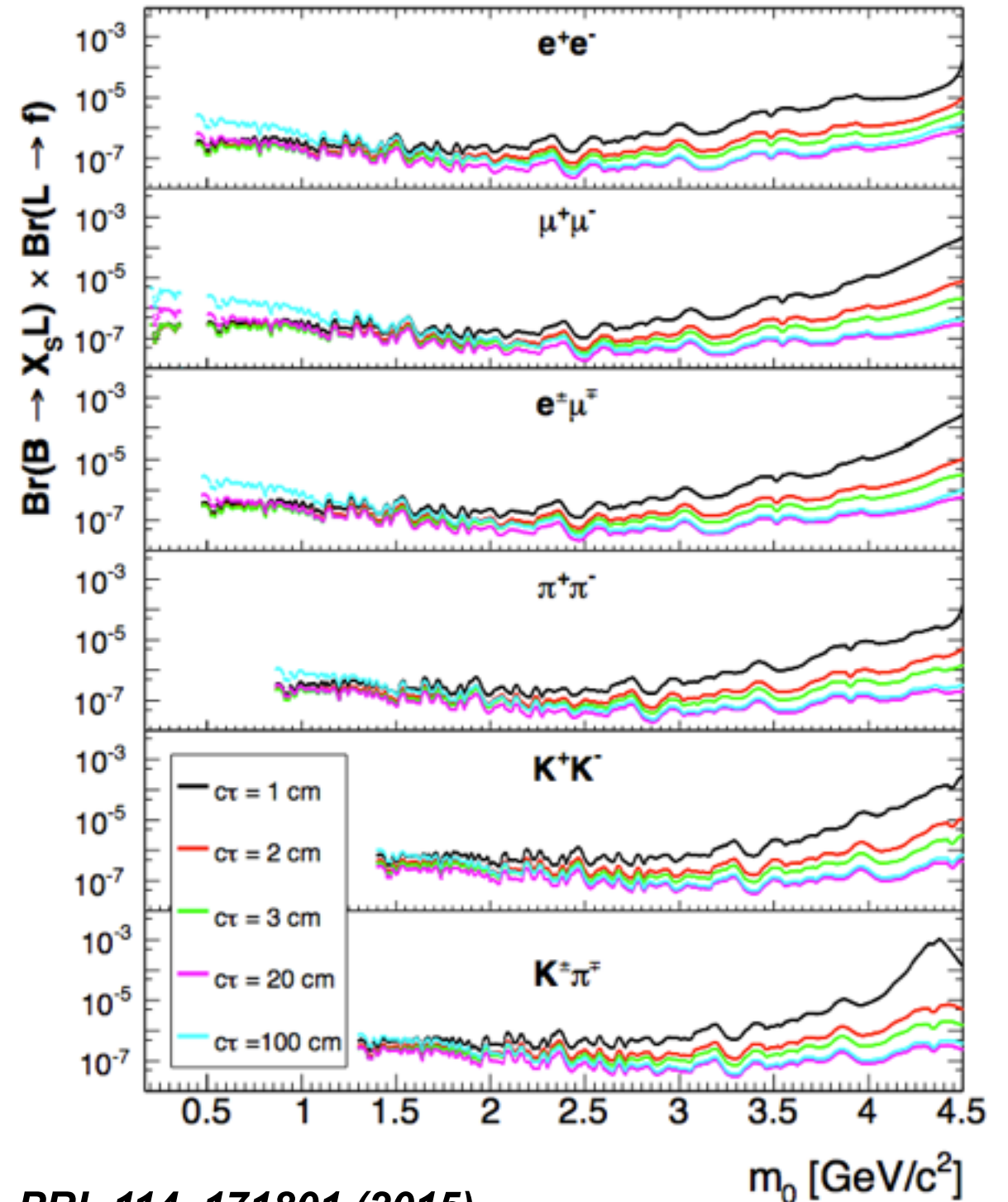
Supplement to PRL114, 171801 (2015)

# BaBar displaced vertex search results

**Model Independent**



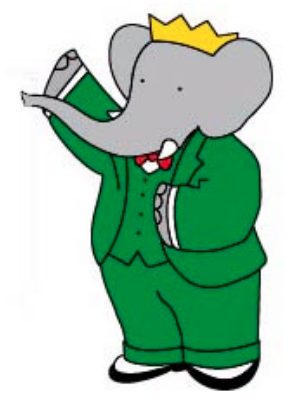
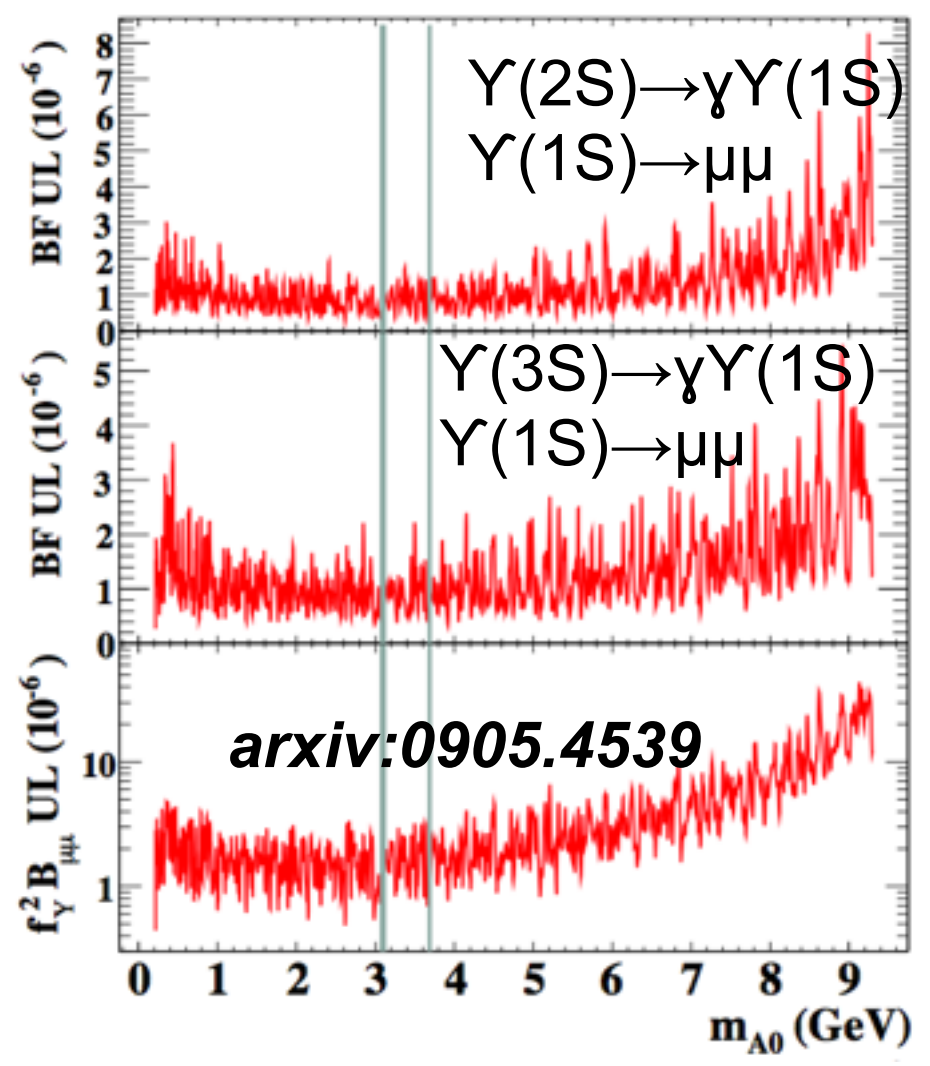
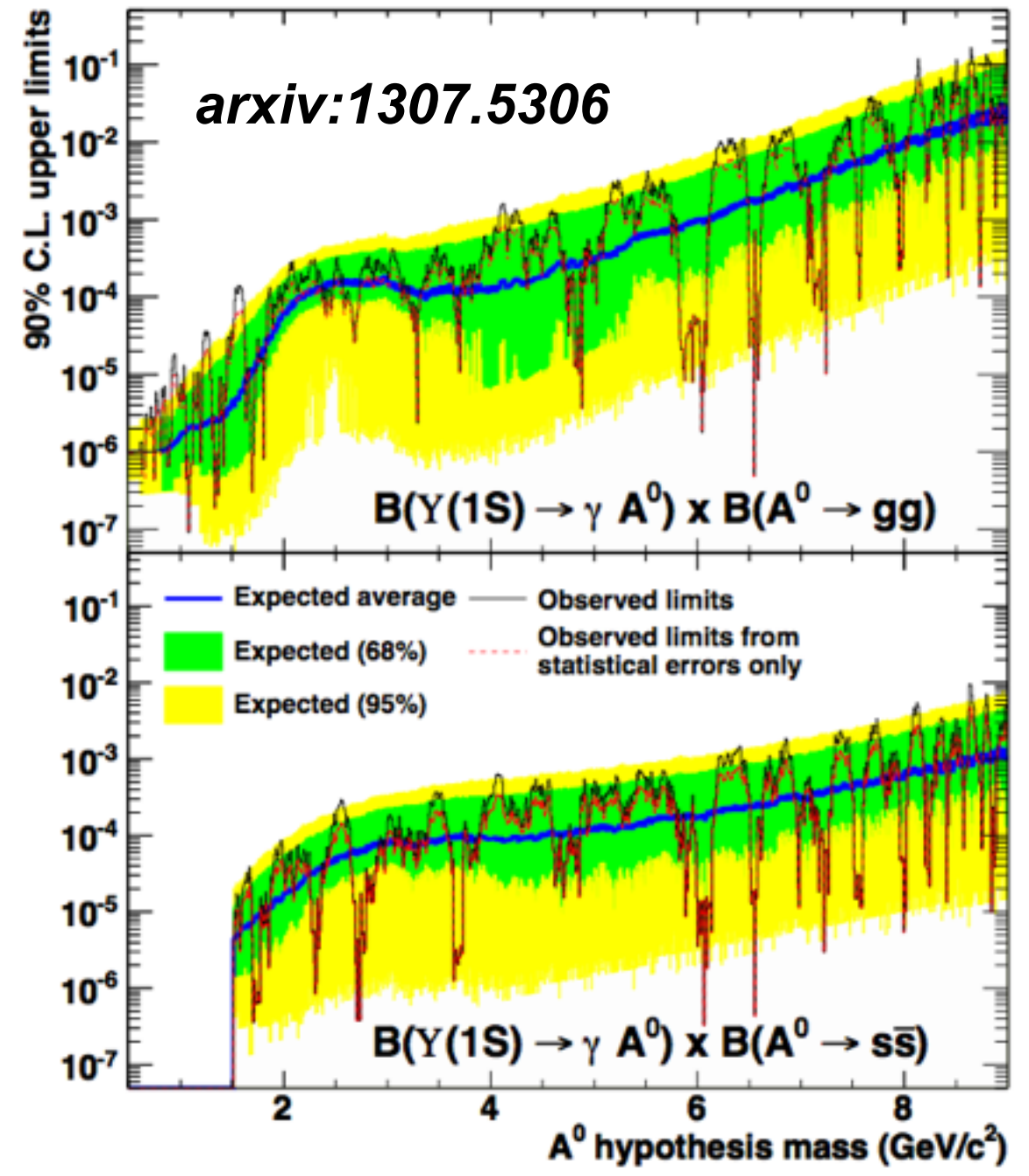
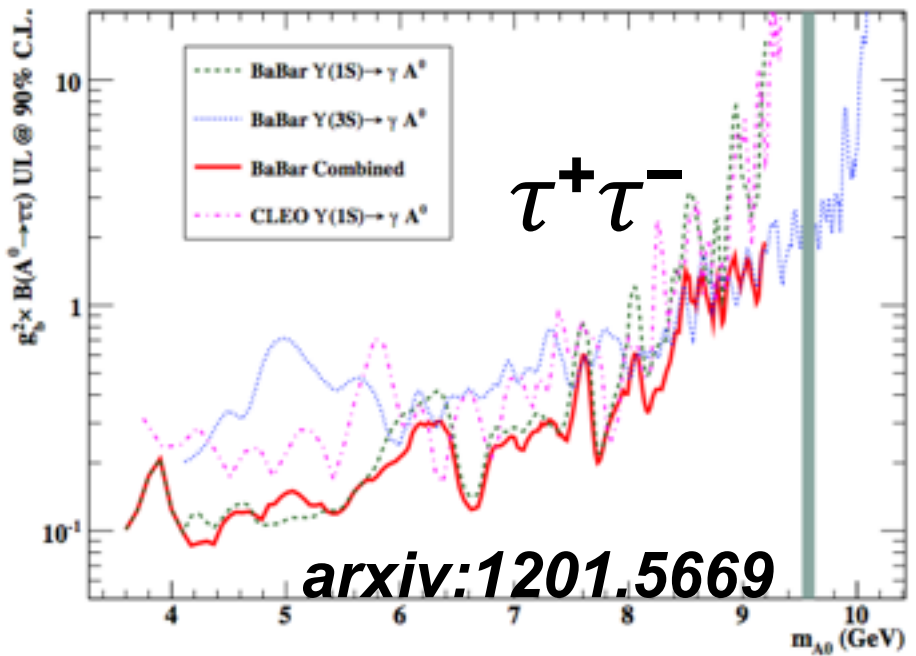
**Inflaton Model (with strangeness=-1)**



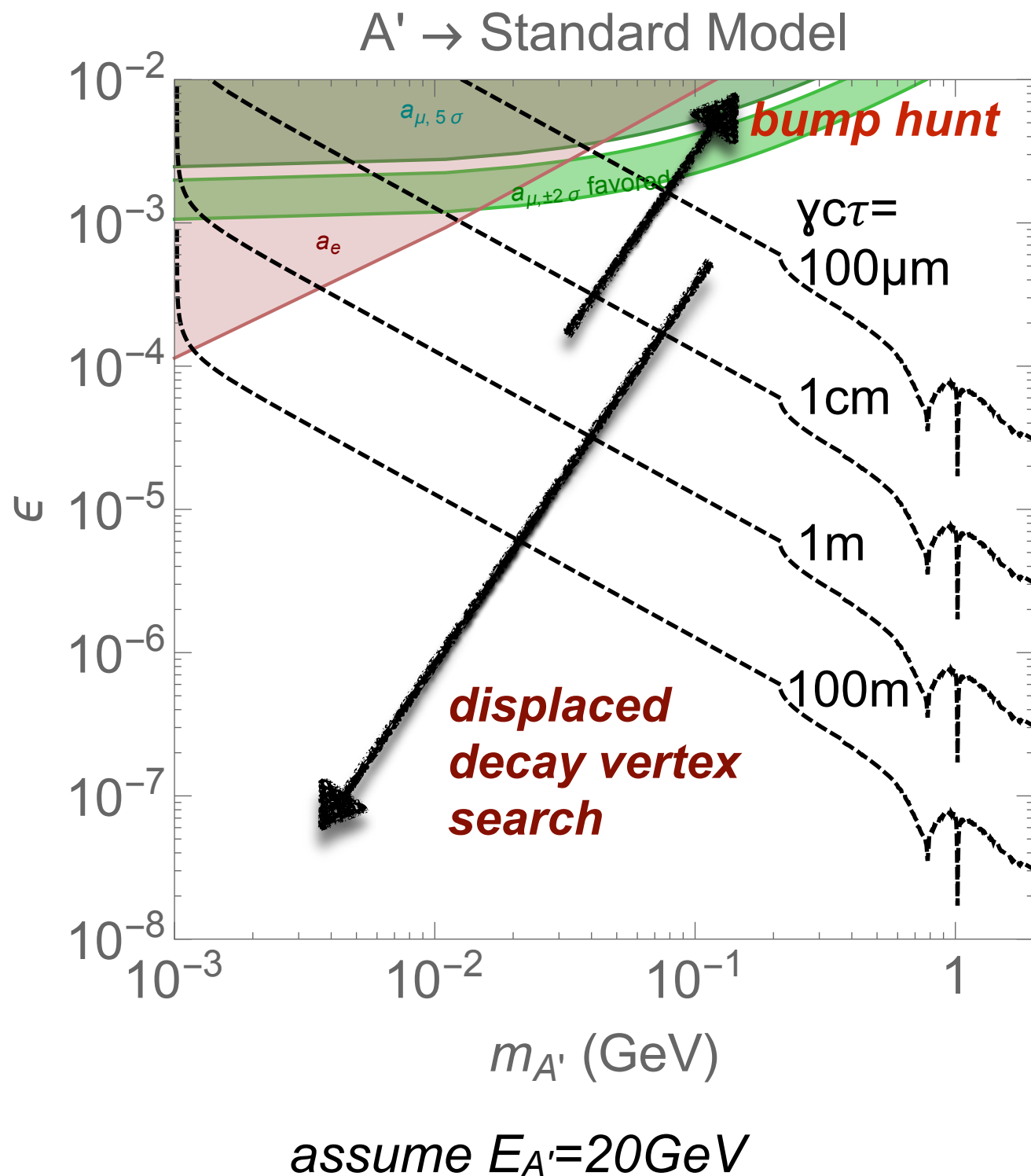
*PRL 114, 171801 (2015)*

- The B-Factories (not just BaBar/Belle but also LHCb) have been incredibly successful in probing the quark flavor sector...and testing its consistency with the SM
  - Belle-II will be carrying that to the next level starting in few years
- The combination of beam-intensity/precision detectors are also a great place to look for new physics
  - here I've talked about new “low-energy” physics — lights Higgs and dark sector particles — but there are other potential states as well
  - in addition, high-scale NP can have direct effects at these low energies... charged LFV in tau decays is an example where the B-Factories have an outsized impact
- The B-Factory searches (including Belle-II) for dark photons (both visible & invisible decays) from  $\sim 1-10$  GeV will be very hard to beat in the foreseeable future









thinking about direct searches...

if (  $\gamma c\tau < \sim 100 \mu\text{m}$  )  
bump-hunt;

if (  $\gamma c\tau > 1 \text{ m}$  )  
displaced decay;

*...some mushy middle where both handles are useful;*