

Neutrino Detectors



Robert Svoboda, SLAC Summer Institute, 2015



The age old question...

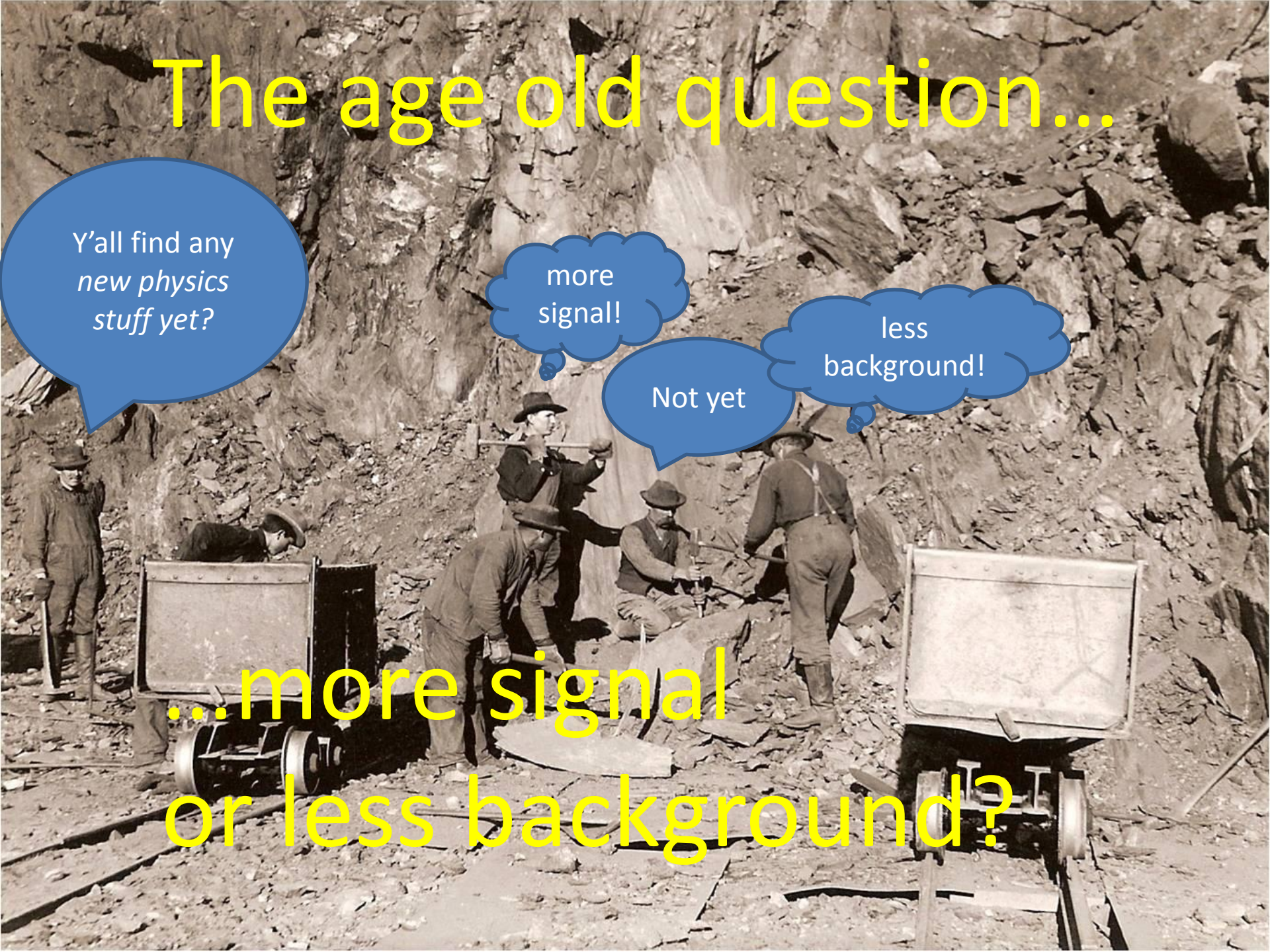
Y'all find any
new physics
stuff yet?

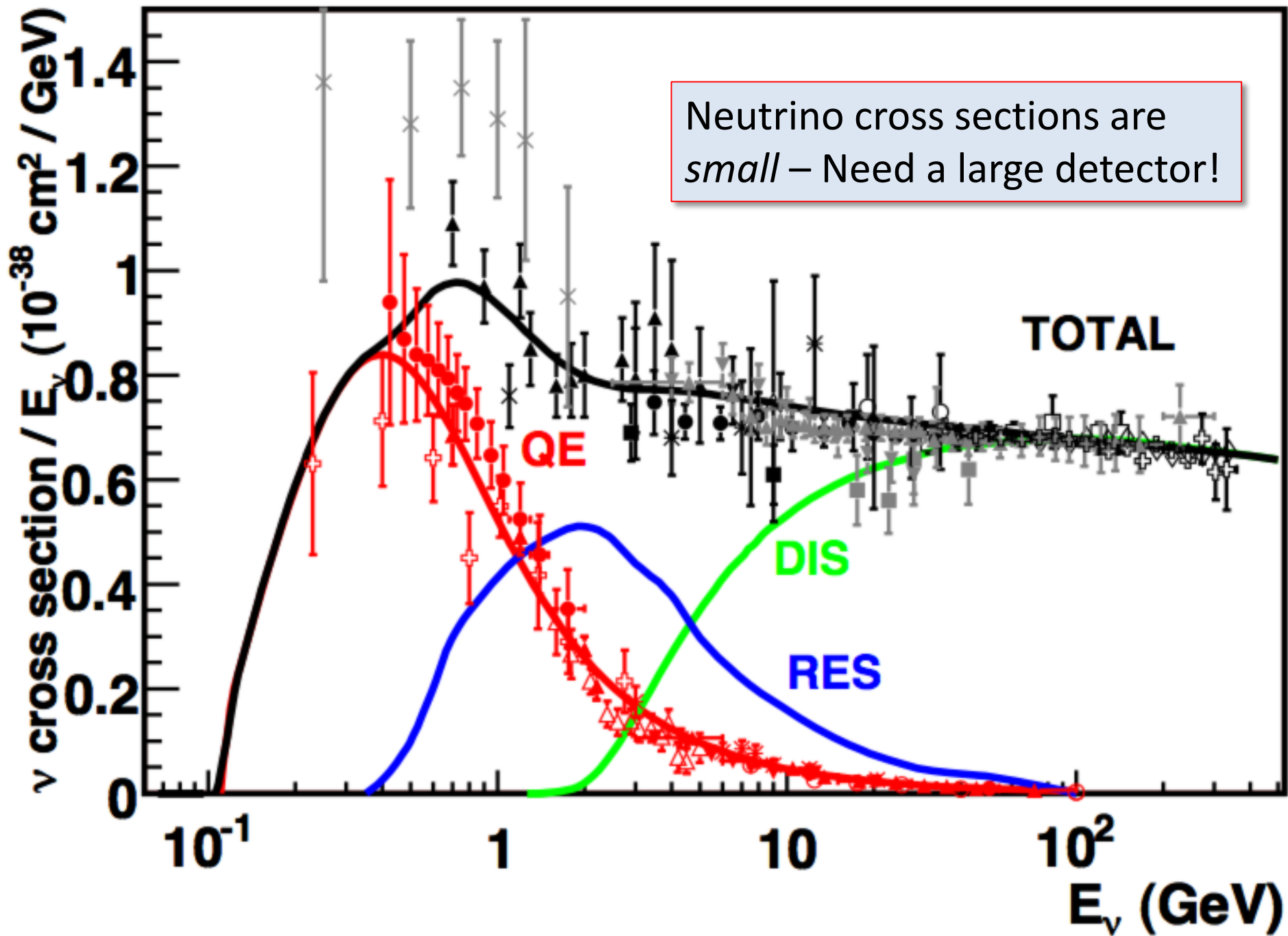
more
signal!

Not yet

less
background!

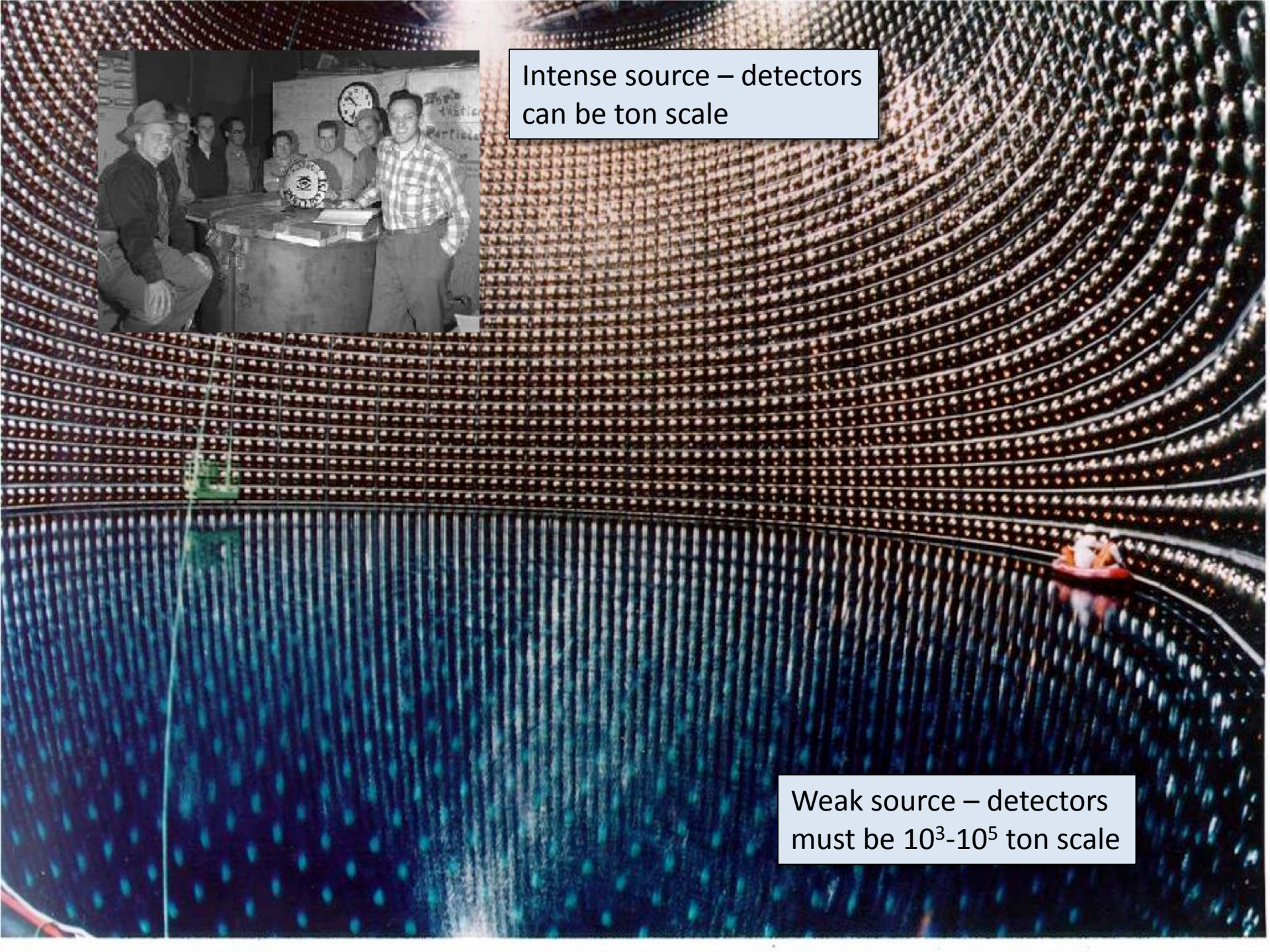
...more signal
or less background?





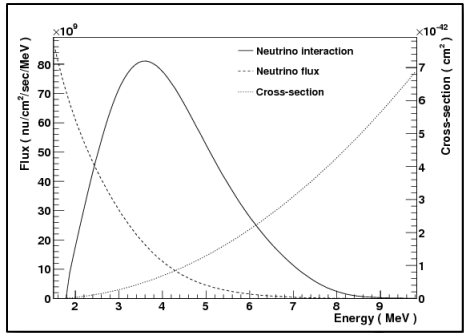
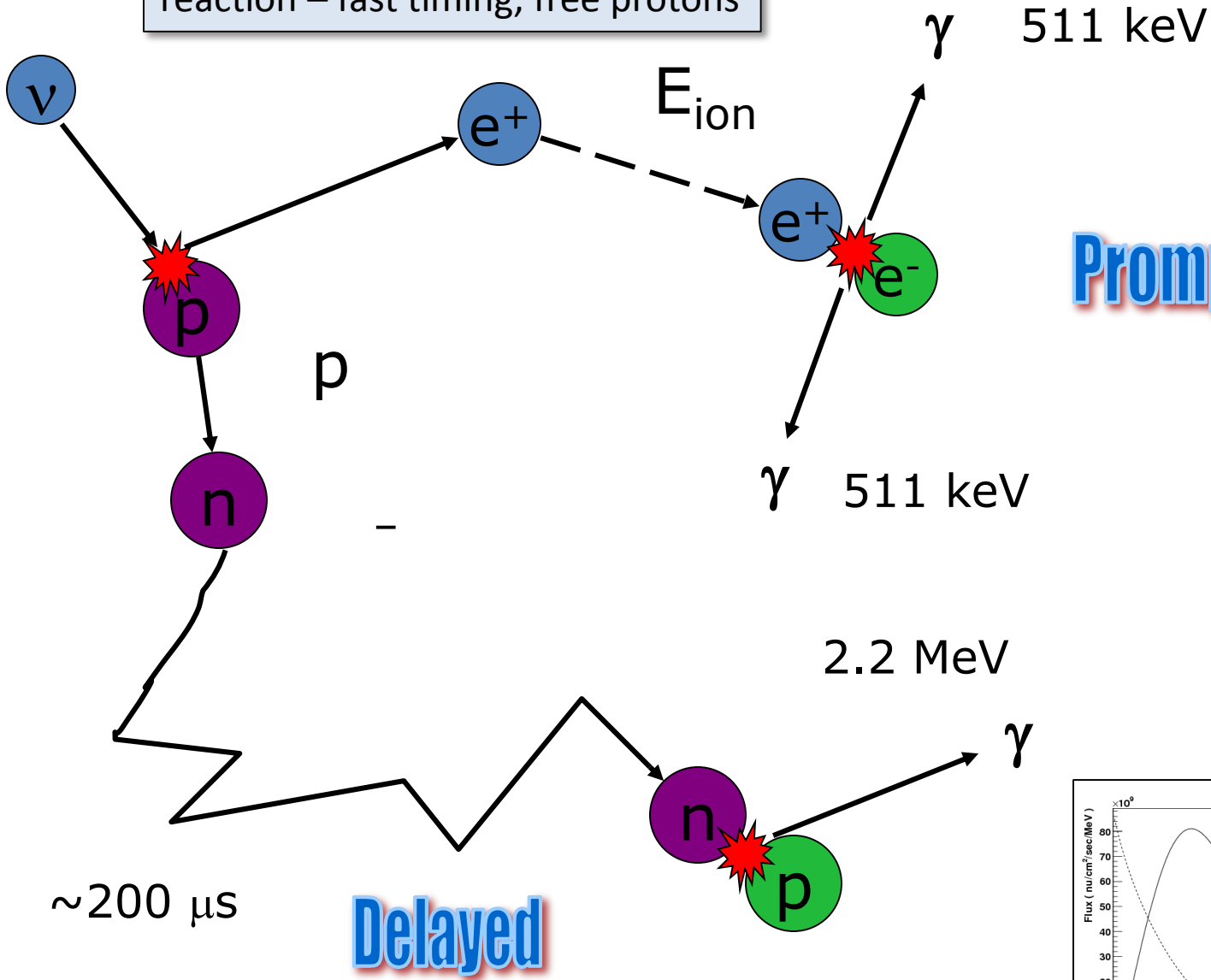


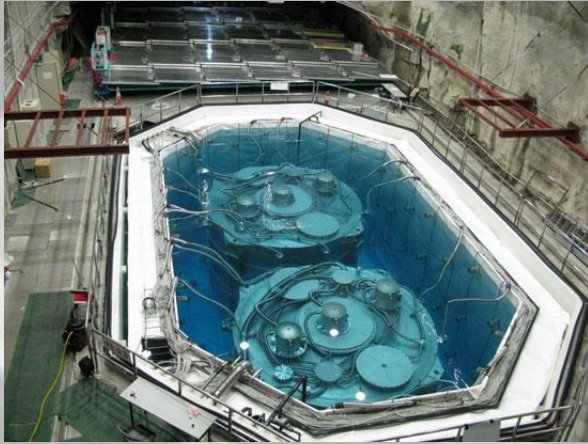
Intense source – detectors
can be ton scale



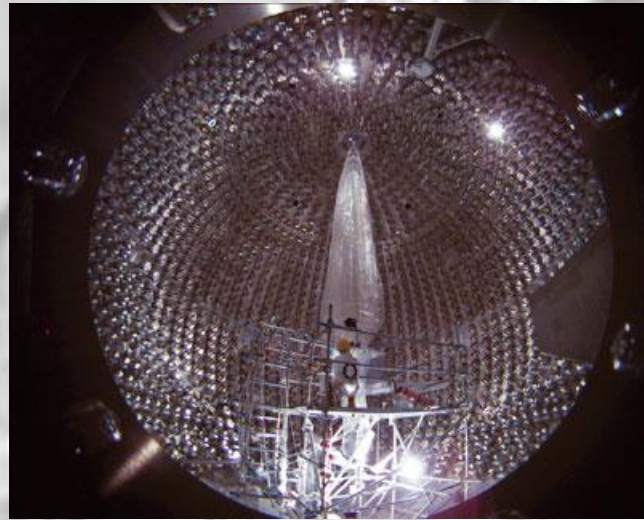
Weak source – detectors
must be 10^3 - 10^5 ton scale

Well-known Inverse Beta Decay reaction – fast timing, free protons

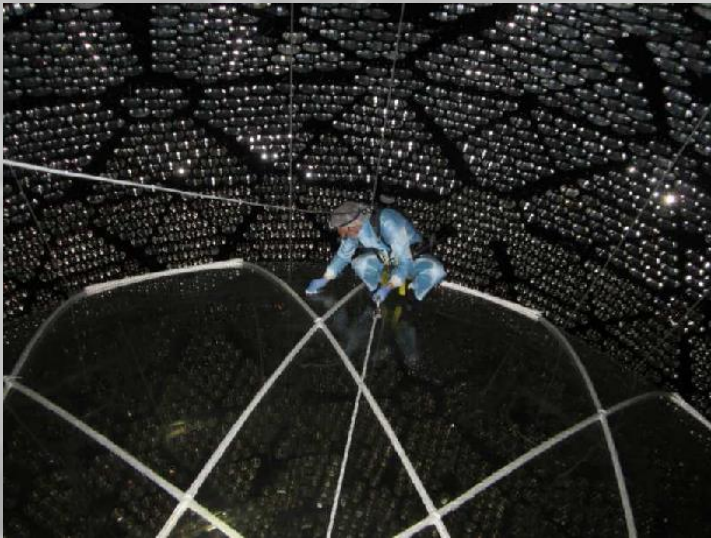




Daya Bay



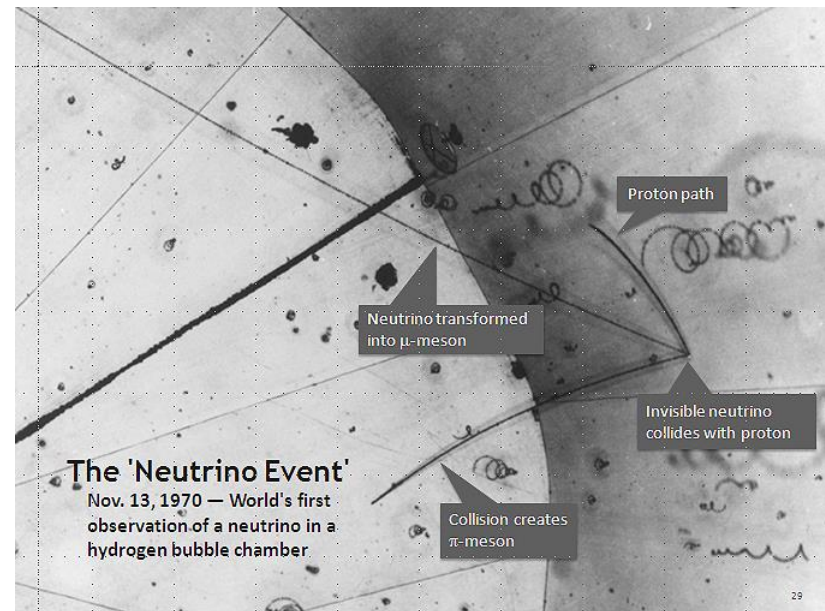
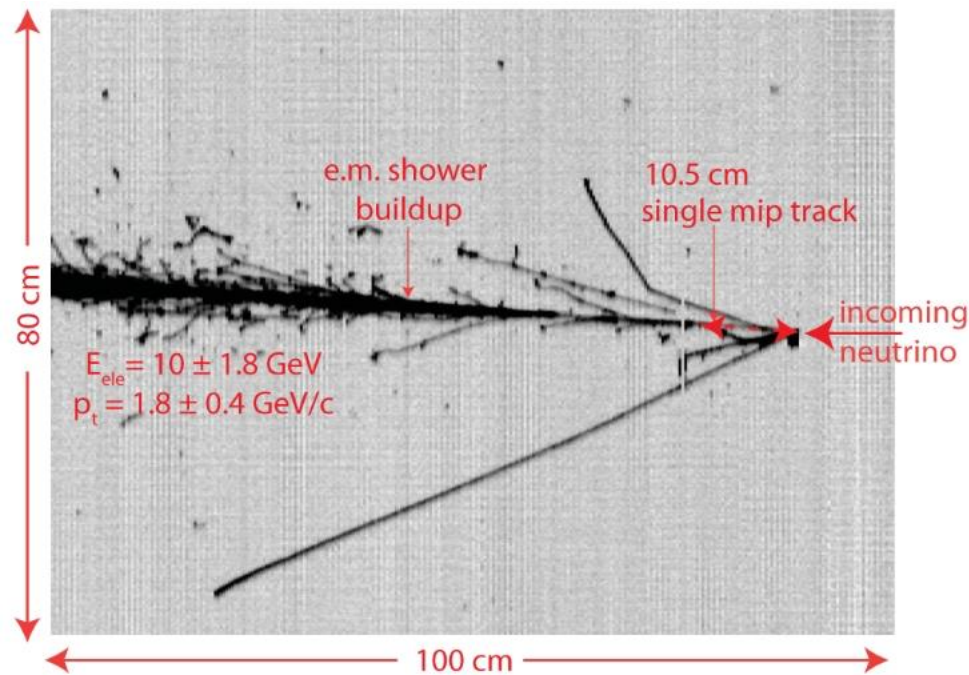
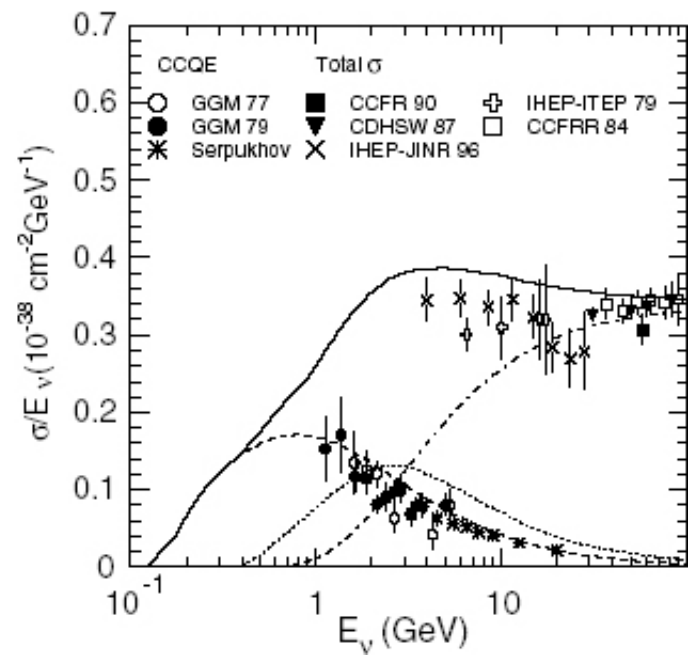
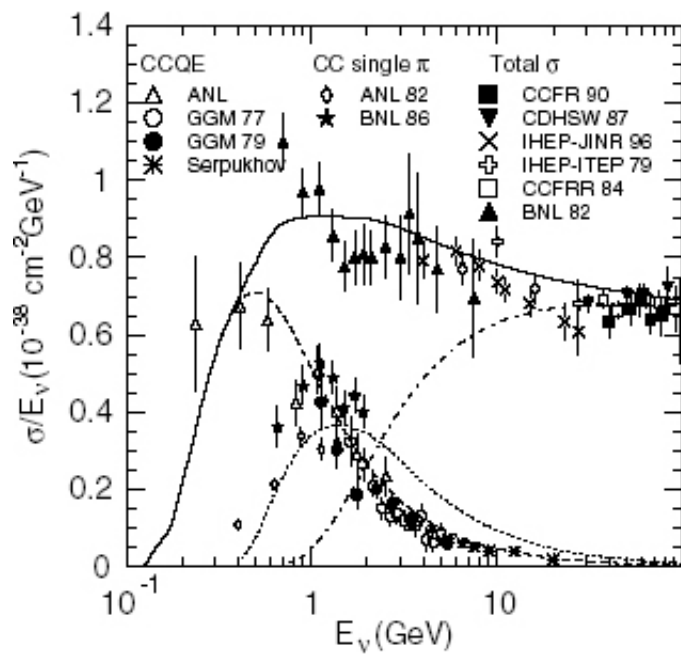
KamLAND



SNO+



Double Chooz





MINOS



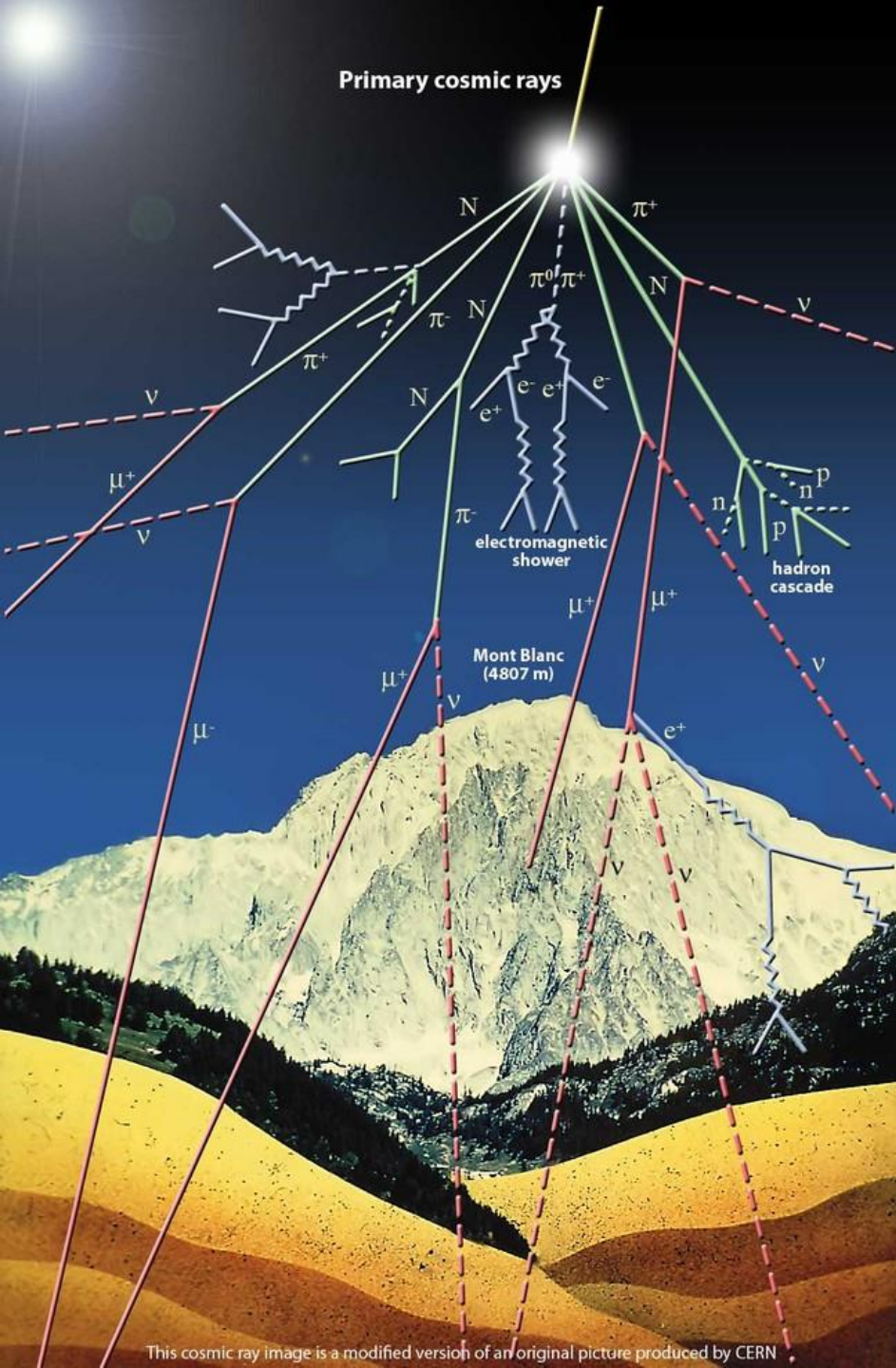
ICARUS



Super-K



Nova



- muons penetrate deep underground – to km of depth
- muon rate, fast neutrons can be a problem
- spallation products can live for seconds (e.g. ^{16}N) to days (e.g. ^{46}Sc)
- Most large detectors need to be located *underground*

How to detect neutrinos?

Y'all find any
new physics
stuff yet?

collect
Charge!

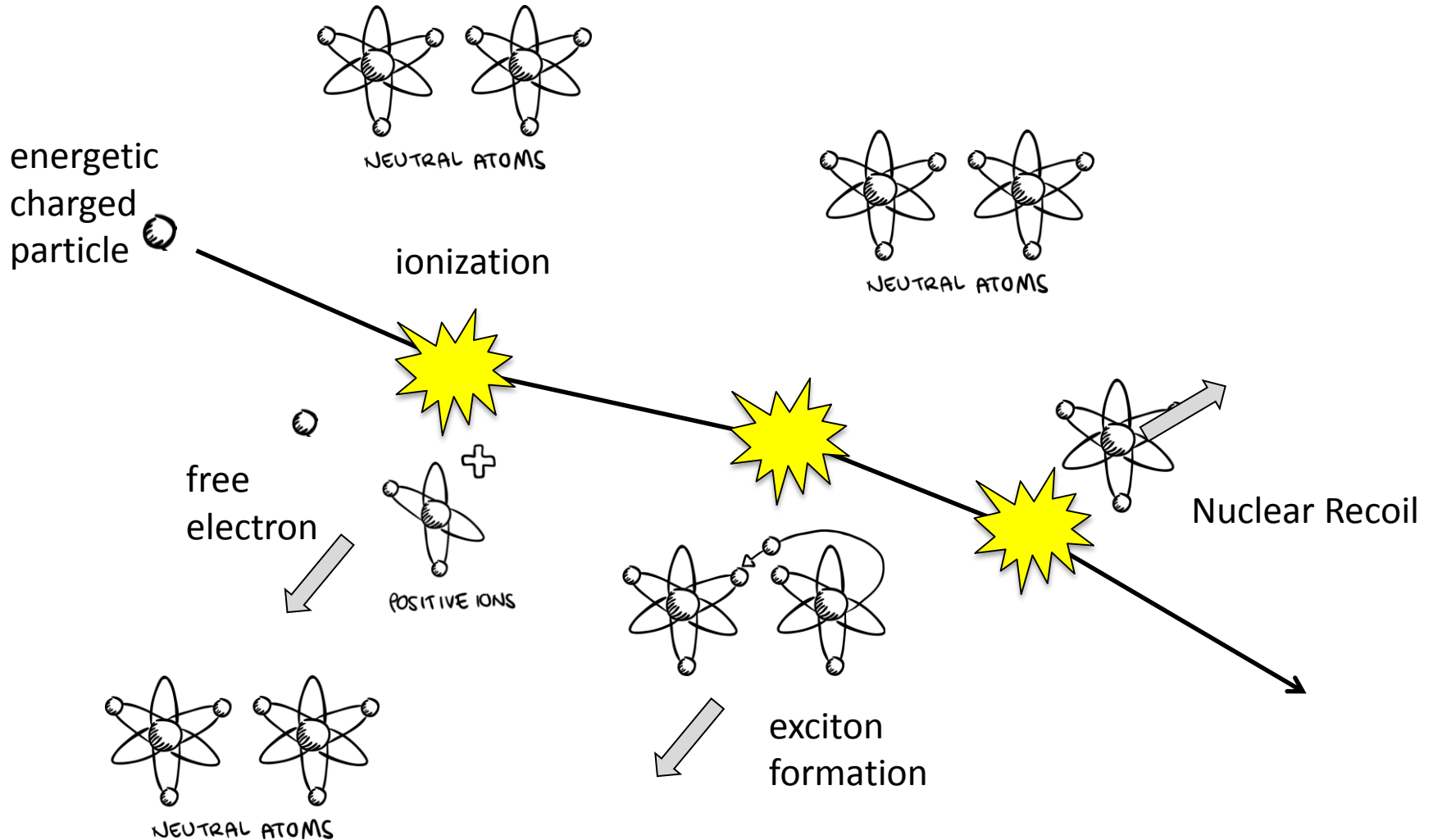
Not yet

collect
Photons!

...collect charge
or photons?



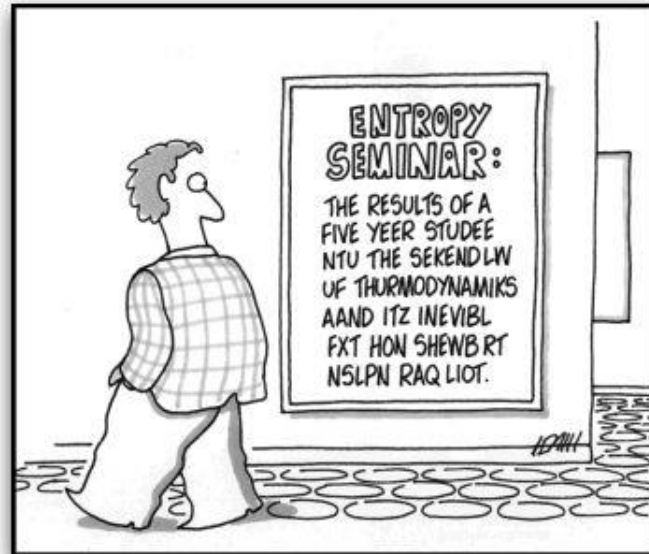
Energy Loss in Matter



Energy Balance

$$E_{\text{Total}} = E_{\text{ion}} + E_{\text{exciton}} + E_{\text{heat}}$$

All the energy eventually ends up as heat...



Neutrino Detectors

Detect Photons

- Water Cherenkov
- Organic Liquid Scintillator
- Plastic Scintillator
- Photon Detection

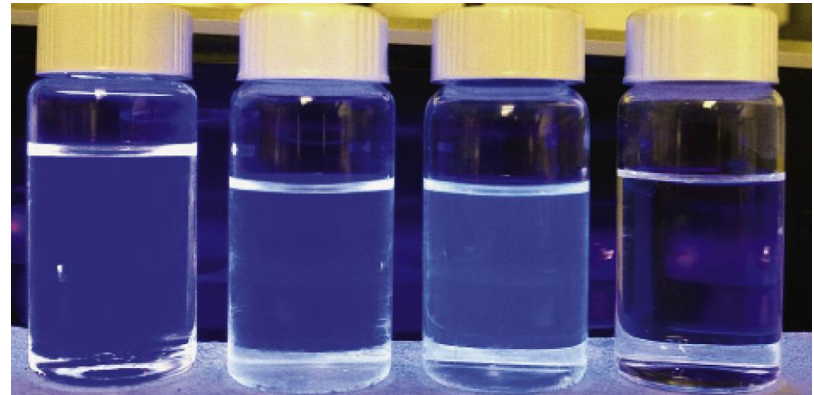
Detect Charge

- Time Projection Chamber (TPC)
- Charge Detection

The technique needs to match what you are trying to measure...

Liquid Scintillator

- In use since 1950's
- Organic liquid with waveshifting fluor converts ionization to photons in 400-700 nm range
- Detectors with 0.2 MeV threshold up to ~ 1 kton have been built
- Fast, free protons, can be loaded with neutron capture agents like gadolinium



Disadvantage: Moderately expensive, combustible, semi-toxic, hazardous waste, light is isotropic – track reconstruction difficult

Plastic Scintillator

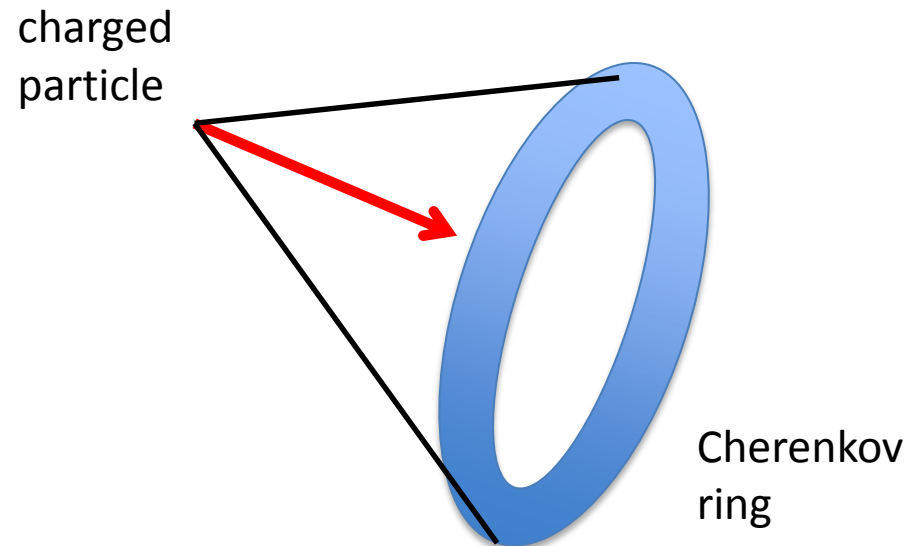
- In use since 1960's
- Can be shaped to fit into large detectors like MINOS or in specialized configurations like MINERVA
- Easier to handle than liquid scintillator
- Still fast with free protons



Disadvantage: Expensive, loading with neutron capture agents difficult, needs segmented readout

Water Cherenkov

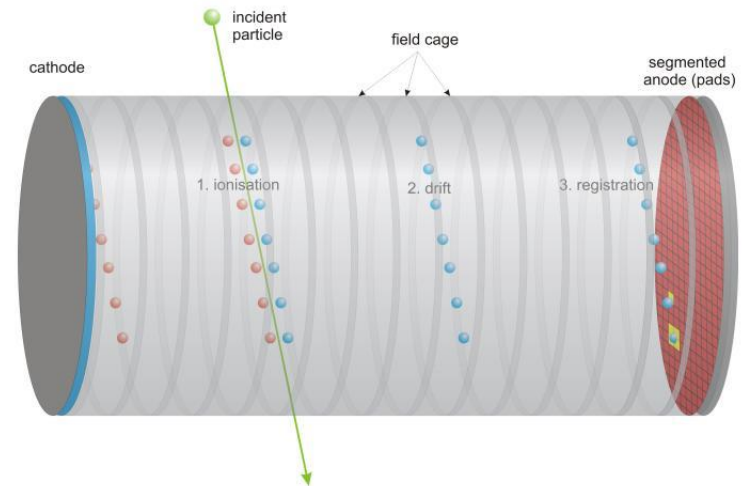
- Large detectors since 1980's. Particles moving faster than c/n give off UV light in a cone shape
- Can reconstruct tracks of particles and determine type from multiple scattering
- Water detectors up to 22 ktons have been built. Megatons in ice!
- Very inexpensive. Can load with neutron capture agent.



Disadvantage: low light yield (2-3% of scintillator), threshold >3 MeV, particles with velocity $v < c/n$ not visible

Liquid Argon TPC

- In development since 1980's. Can drift electrons over meters in liquid noble gases like argon
- Can image tracks of particles with excellent precision. Not subject to Cherenkov threshold
- Scintillate in far UV
- Argon can be obtained in kiloton quantities
- Detectors up to 300 tons have been built



Disadvantage: low light yield (2-3% of scintillator), threshold >3 MeV, particles with velocity $v < c/n$ not visible

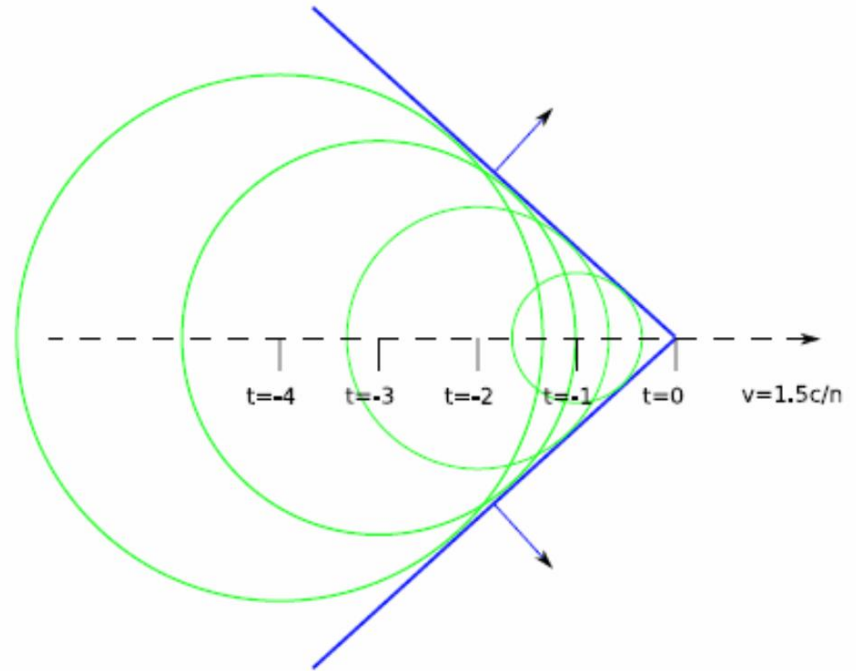
How Particles Produce Light

- charged particles can produce light in materials
 - Čerenkov radiation
 - scintillation in some materials due to excitation and ionization
 - transition radiation
 - neutral particles (e.g. γ and n) must interact first and the resulting electrons, protons, recoiling nuclei or other charged hadrons produce light
-

Čerenkov Light

- emitted by charged particles whose velocity exceeds c/n
- results in a coherent shock wavefront that's a cone of light
- half angle given by:
 - $\cos \theta = 1 / [\beta n(\lambda)]$
 - $\beta = v/c$
 - e.g. 41° in water for $\beta \approx 1$
- spectrum of photons emitted:

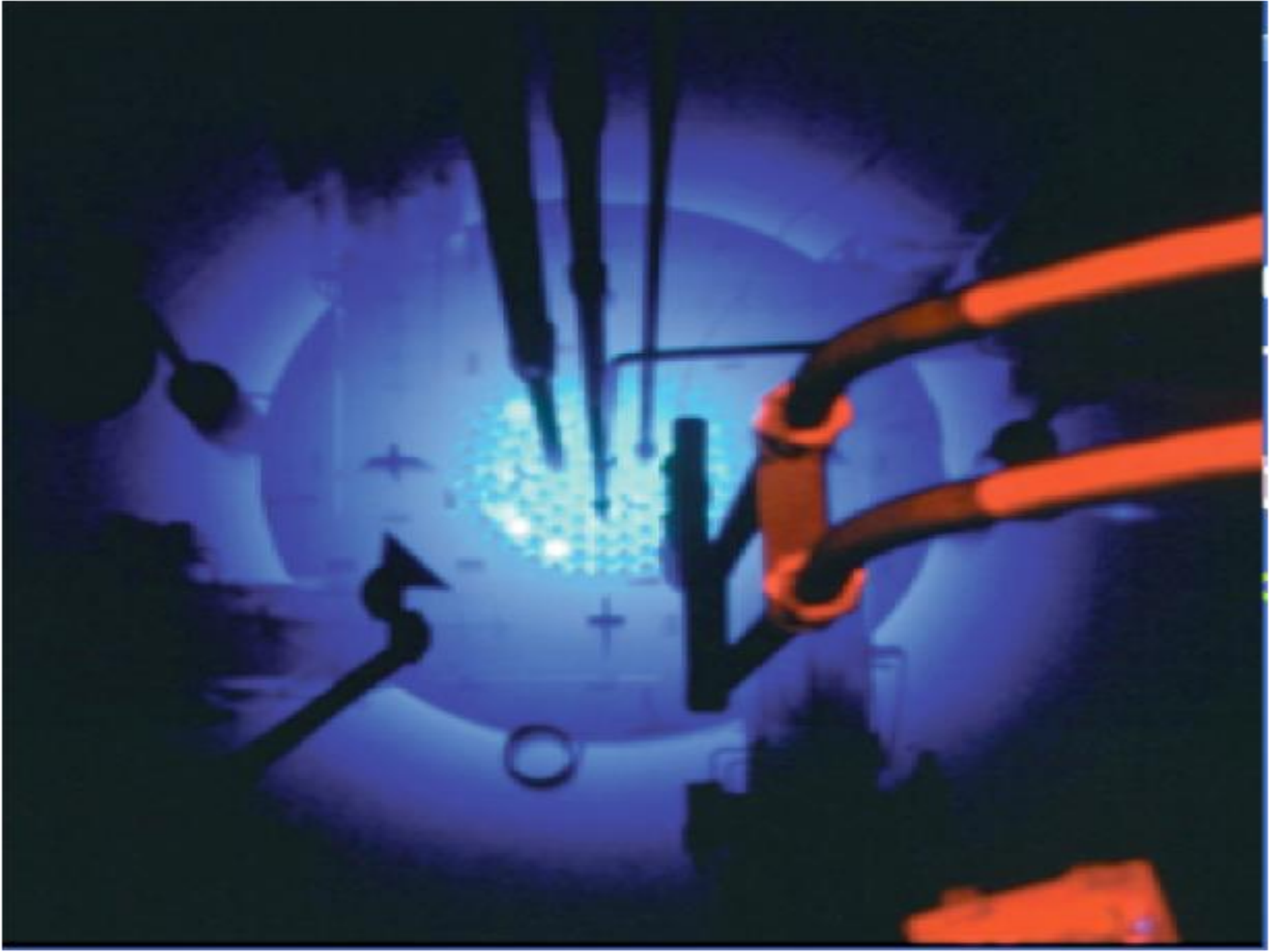
$$\frac{dN}{dx d\omega} = \frac{\alpha}{c} \left(1 - \frac{1}{n^2(\omega) \beta^2} \right)$$



from <http://wwwmagic.mppmu.mpg.de/publications/theses/JLopez.pdf>

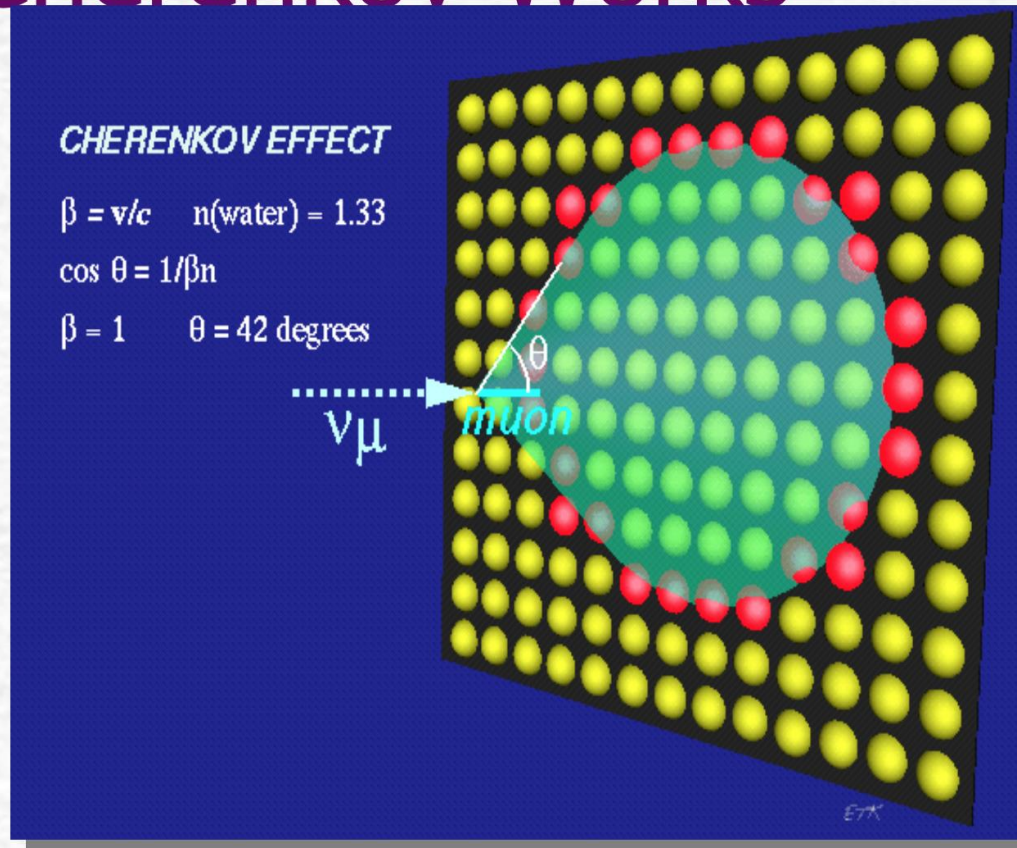
Characteristics of Čerenkov Light

- it's prompt
- basically flat frequency spectrum of the emitted light
 - that's what distinguished it from fluorescence which exhibits bands in the emitted spectrum
- which means it goes as $\frac{dN}{d\lambda} \propto \frac{1}{\lambda^2}$
- hence it's UV and blue light
 - cannot be harder UV because most materials start to absorb at shorter wavelengths and ϵ must be real



How Water Cherenkov Works

- ✓ Cheap target material
- ✓ Surface instrumentation
- ✓ Vertex from PMT timing
- ✓ Direction from ring edge
- ✓ Energy from pulse height, range and opening angle
- ✓ Particle ID from hit pattern and delayed muon decay signature
- ✓ Cherenkov threshold:
 - ✓ $\beta > 1/n \sim 0.75$

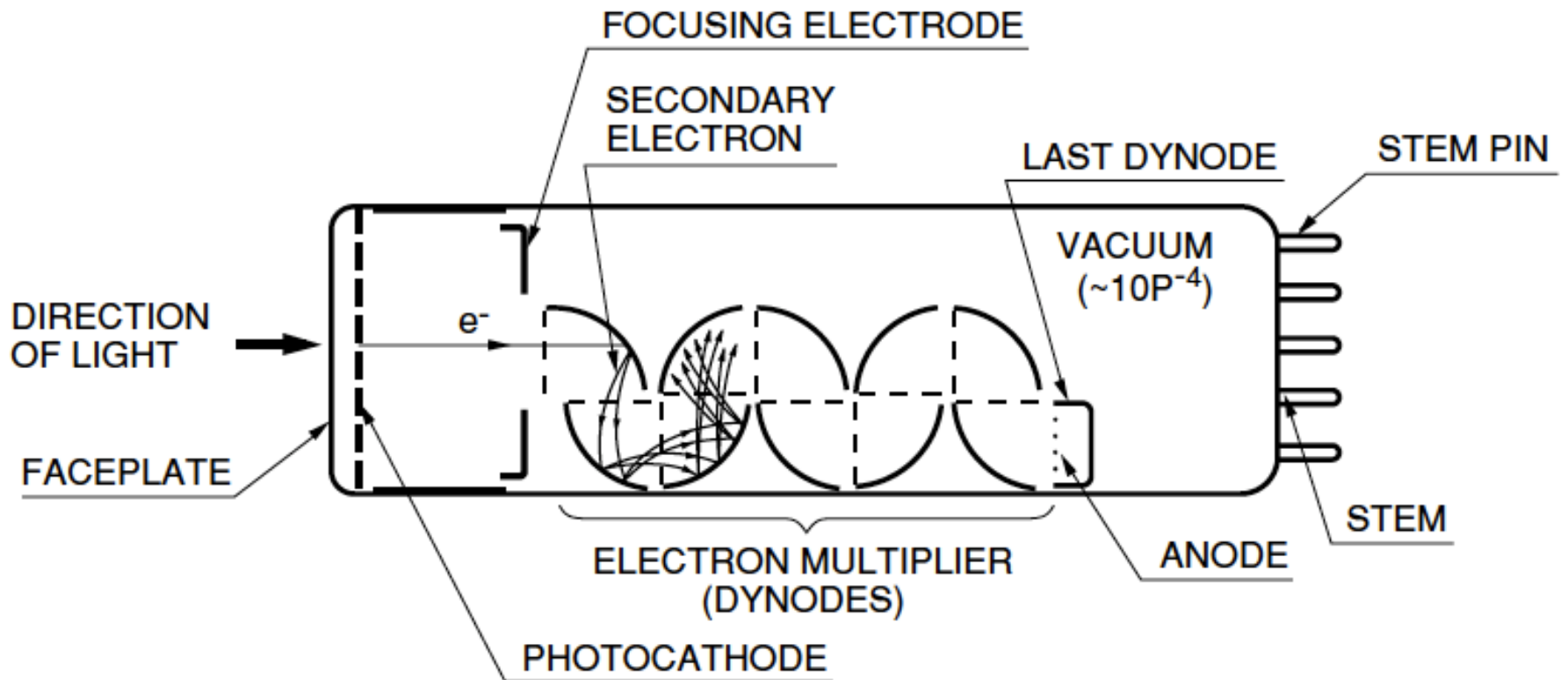


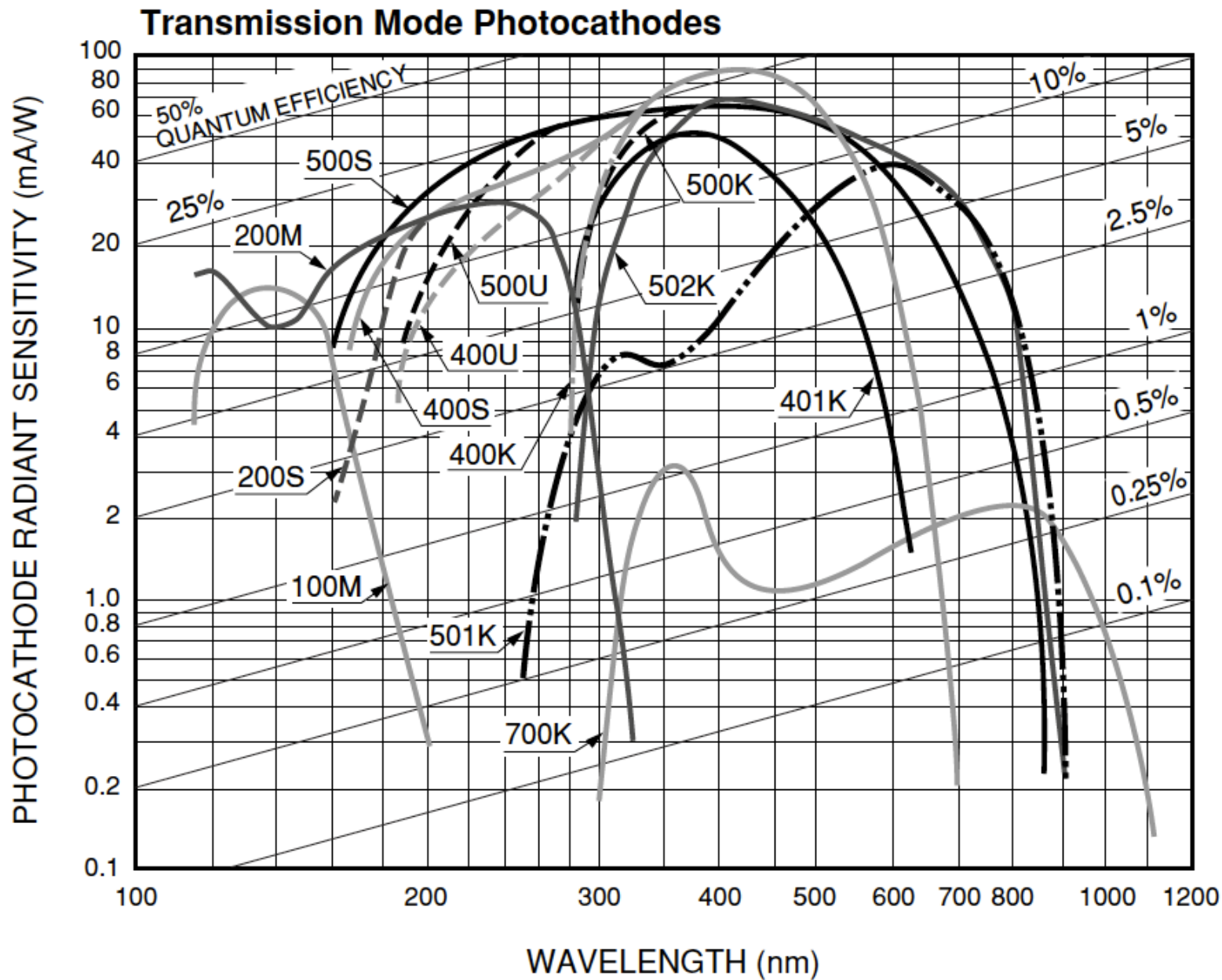
D.Casper

Water Cherenkov detectors require large, fast photon detectors in the UV/Visible spectral range. Particle reconstruction and ID possible with "CCD-like" ring image



Photomultiplier Tubes (PMT)





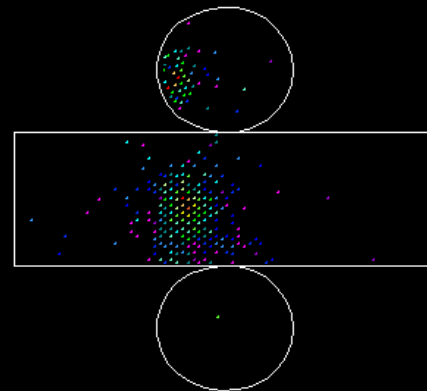
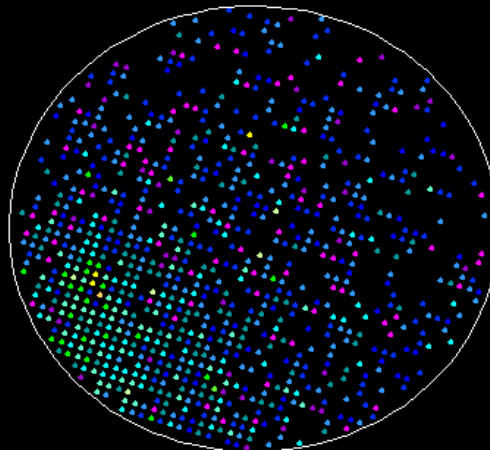
THBV3_0402Eb

Figure 4-2 (b): Typical spectral response characteristics of transmission mode photocathodes

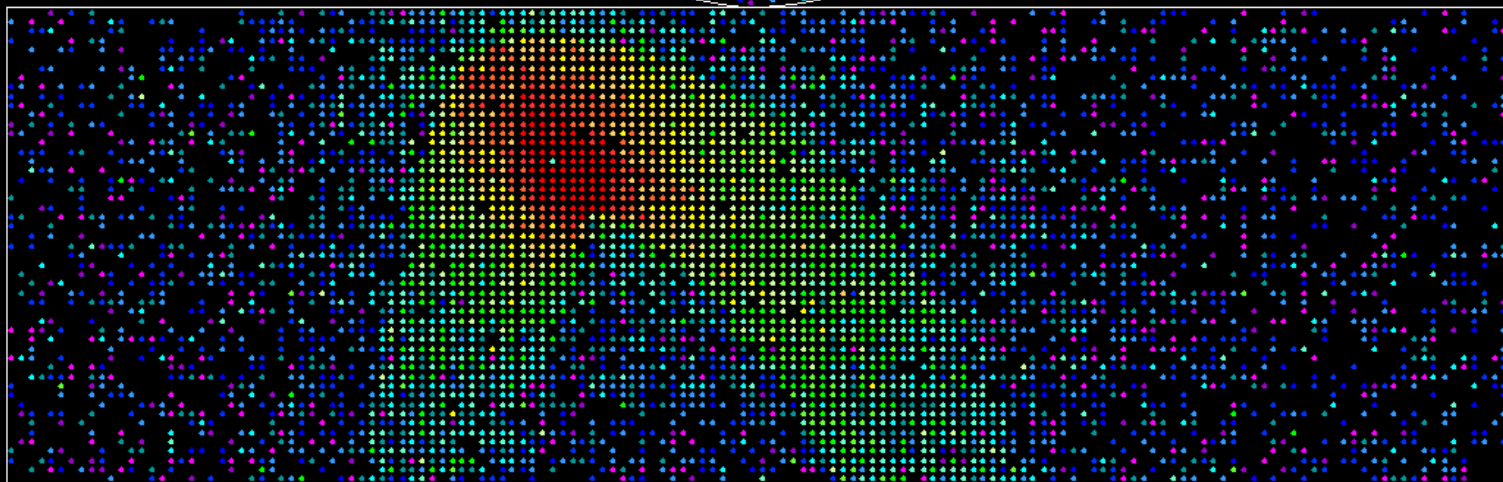
```

Display      :CHARGE INNER
Date         :Thu Apr 30 2015
Run          :73580 Normal
Event        :36607464
Event time   :13:12:00.324994
TRG Type(s) :LE HE SLE OD SHE
TotalPE ID/OD :74719.4 2169.6
NumHits ID/OD :6883 239
Time Diff    :19237.656250 us

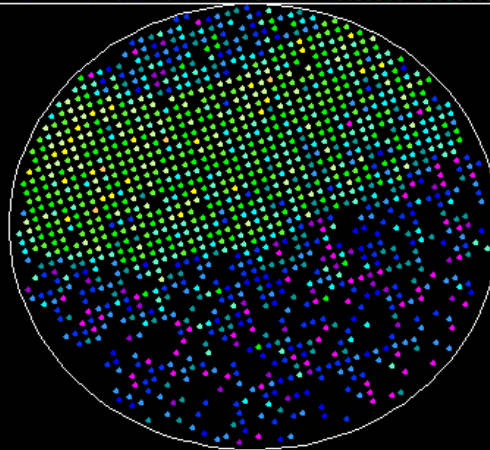
```



<- Rotate -> : 1 PI / 10



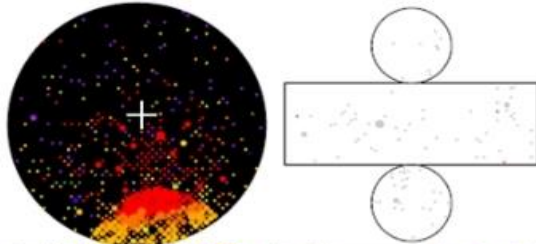
Time Window (ns):[-300.0, 1000.0]



0.0 0.5 1.4 1.9 2.6 3.7 5.1 7.0 9.7 13.4 18.5 25.6 35.4 49.0 67.8 179.4

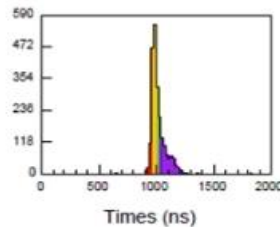
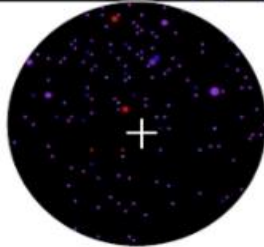
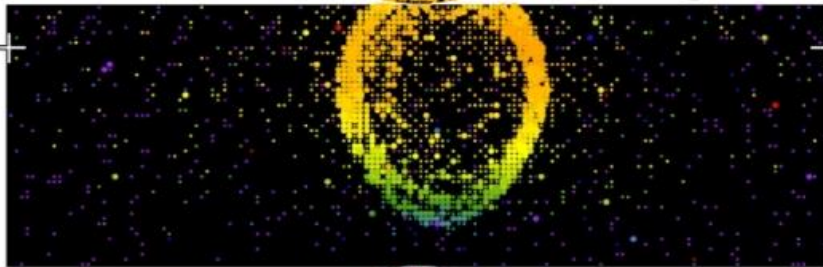
Super-Kamiokande

Run 3062 Event 475360
96-11-08:12:07.30
Inner: 2306 hits, 7763 pE
Outer: 5 hits, 4 pE (in-time)
Trigger ID: 0x03
D wall: 601.2 cm
PC mu-like, p = 1088.0 MeV/c



Time (ns)

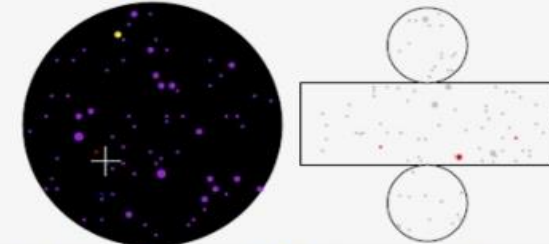
- * < 971
- * 971- 977
- * 977- 983
- * 983- 989
- * 989- 995
- * 995-1001
- * 1001-1007
- * 1007-1013
- * 1013-1019
- * 1019-1025
- * 1025-1031
- * 1031-1037
- * 1037-1043
- * 1043-1049
- * 1049-1055
- * >1055



Muon

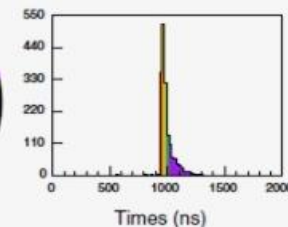
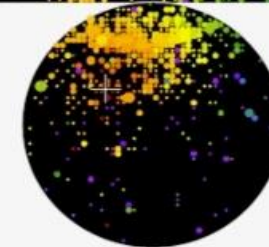
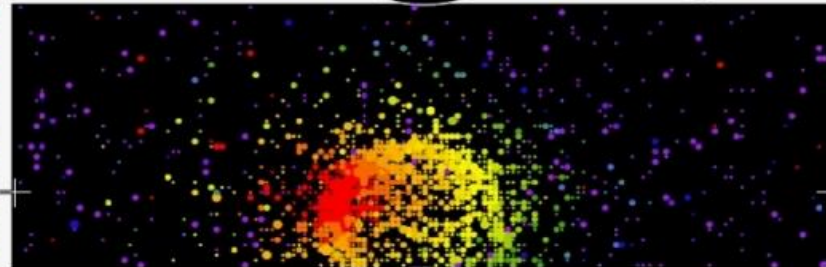
Super-Kamiokande

Run 3013 Event 149004
96-10-24:19:39:51
Inner: 1763 hits, 4003 pE
Outer: 3 hits, 5 pE (in-time)
Trigger ID: 0x03
D wall: 897.4 cm
PC e-like, p = 463.8 MeV/c



Time (ns)

- * < 950
- * 950- 956
- * 956- 962
- * 962- 968
- * 968- 974
- * 974- 979
- * 979- 985
- * 985- 991
- * 991- 997
- * 997- 1003
- * 1003-1008
- * 1008-1013
- * 1013-1018
- * 1018-1023
- * 1023-1028
- * >1028



Electron

Scattering makes electron tracks easily distinguishable from muon tracks with very high efficiency (>98%) in CCQE region.



What's New?

Super-Kamiokande

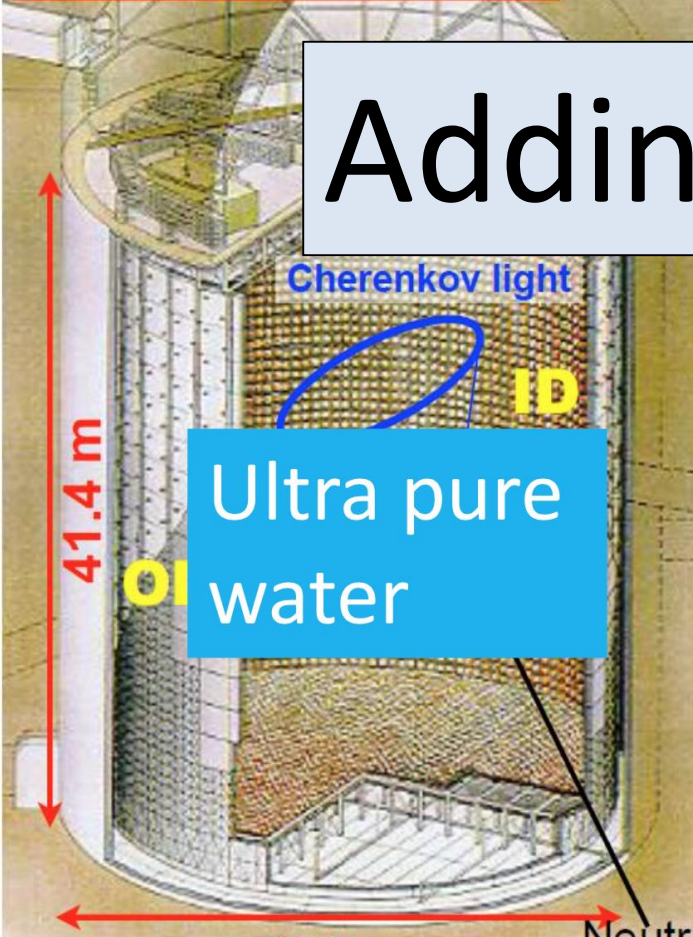
50000 tons of Water Cherenkov detector

Kamioka mine

~1km



Adding Gadolinium...



Phase	Period	Livetime (days)	Fiducial vol. (kton)	# of PMTs	Energy thr.(MeV)
SK-I	1996.4 ~ 2001.7	1496	22.5	11146 (40%)	4.5
SK-II	2002.10 ~ 2005.10	791		5182 (20%)	6.5
SK-III	2006.7 ~ 2008.8	548	22.5 (>5.5MeV) 13.3 (<5.5MeV)	11129 (40%)	4.5
SK-IV	2008.9 ~	1669	22.5 (>5.5MeV) 13.3 (4.5<E<5.5) 8.8 (<4.5MeV)		3.5

total **4504** days

(coverage) (Kinetic energy)

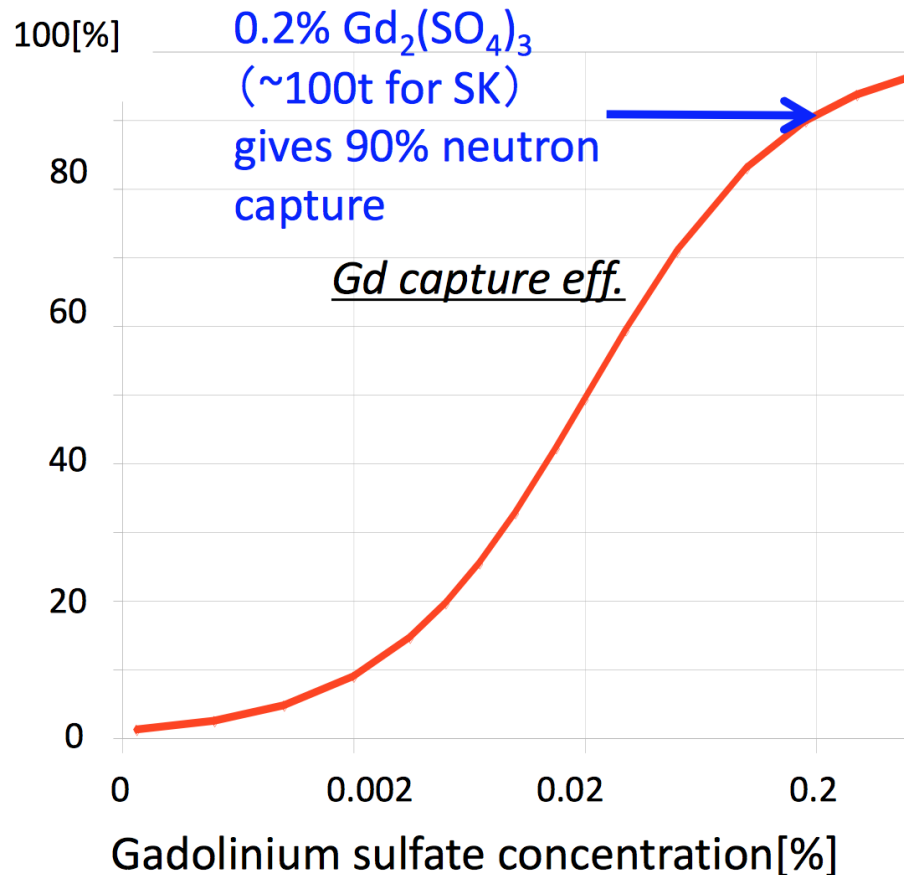
39.3 m

41.4 m

Super K-Gd

Beacom and Vagins PRL93,171101 (2004)

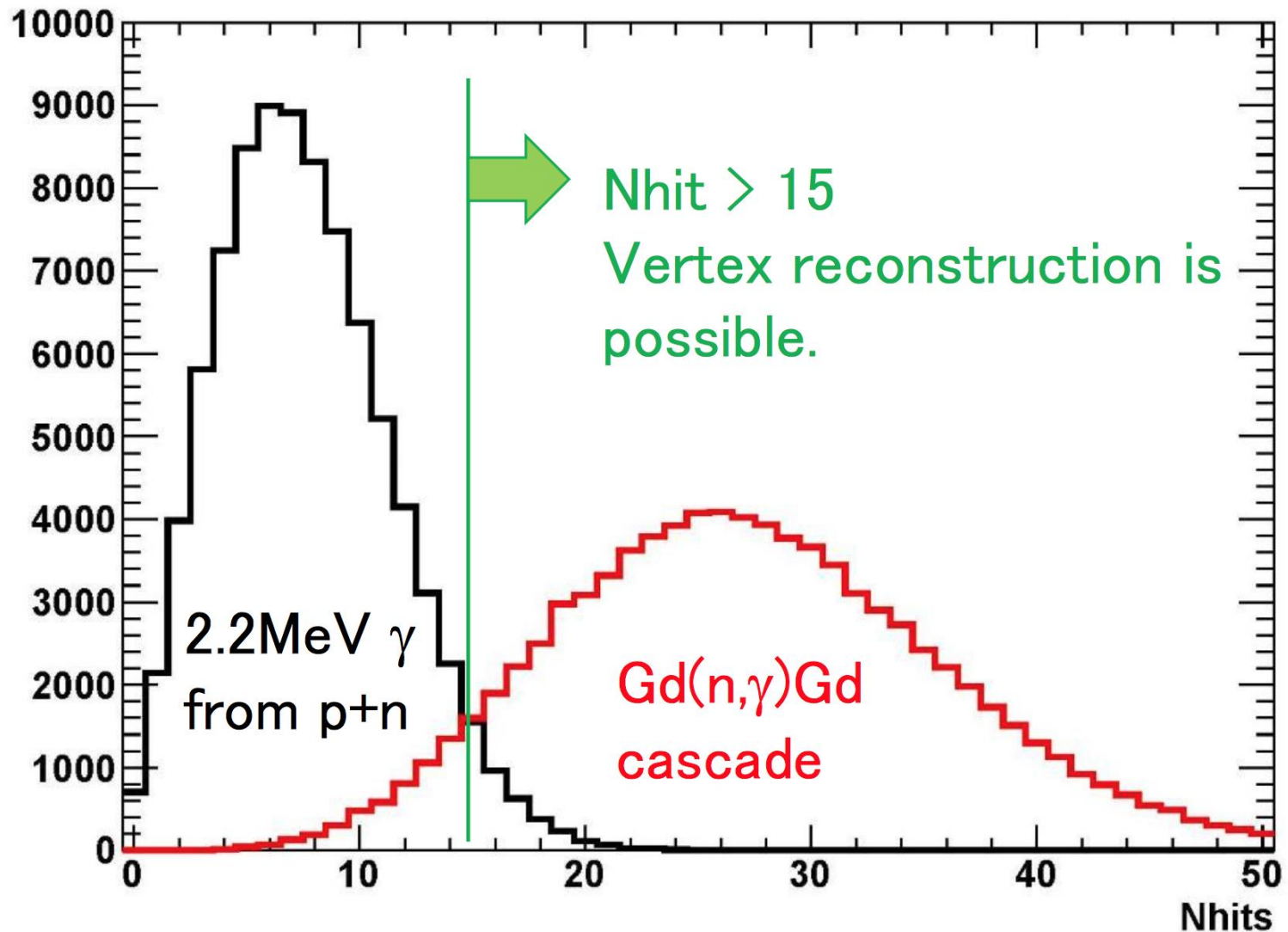
- Large cross section for thermal neutron (48.89kb)
- Neutron captured Gd emits 3-4 γ ray in total 8 MeV
- **We can tag $\bar{\nu}_e$ by using the delayed coincidence technique.**



Physics targets:

- (1) Supernova relic neutrino (SRN)
- (2) Improve pointing accuracy for galactic supernova
- (3) Precursor of nearby supernova by Si-burning neutrinos
- (4) Reduce proton decay background
- (5) Neutrino/anti-neutrino discrimination (Long-baseline and atm nu's)
- (6) Reactor neutrinos

Number of hit PMT (Nhit) distributions



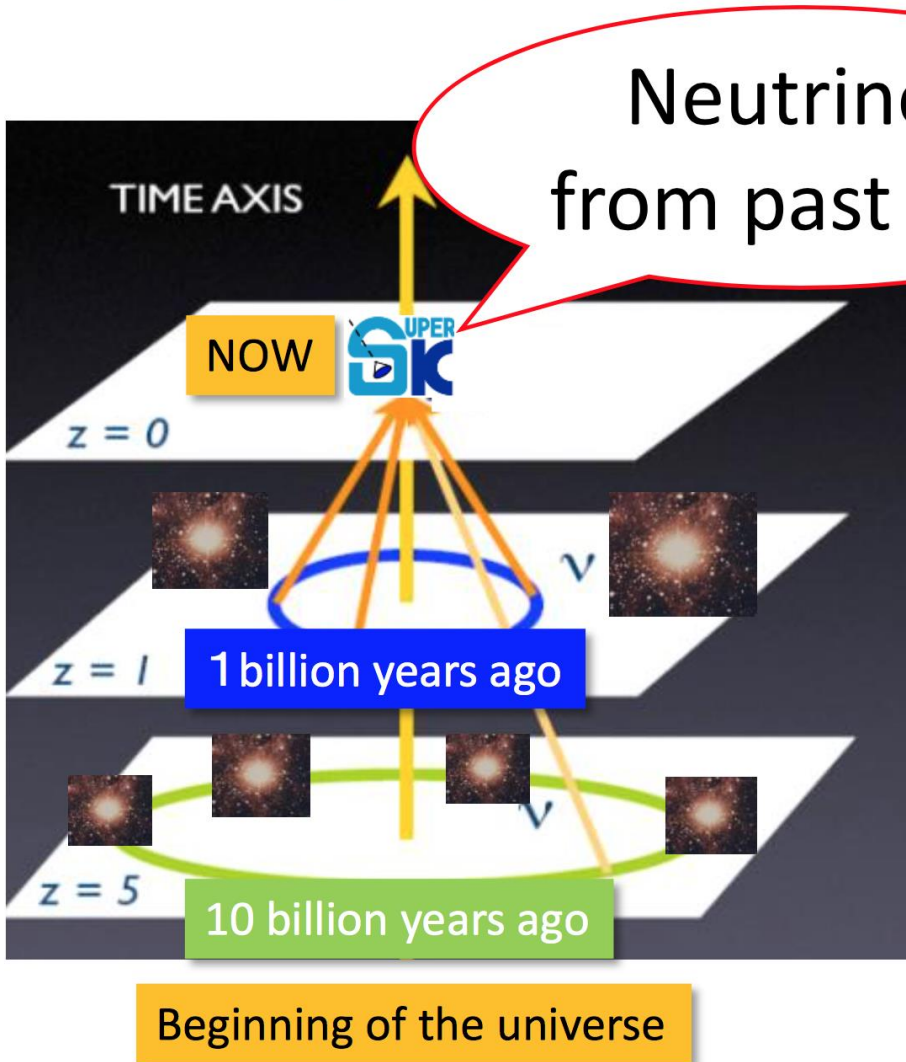
Efficiency and fake probability

2.2 MeV γ : Efficiency: 10 ~ 20%, fake probability: $\sim 10^{-2}$

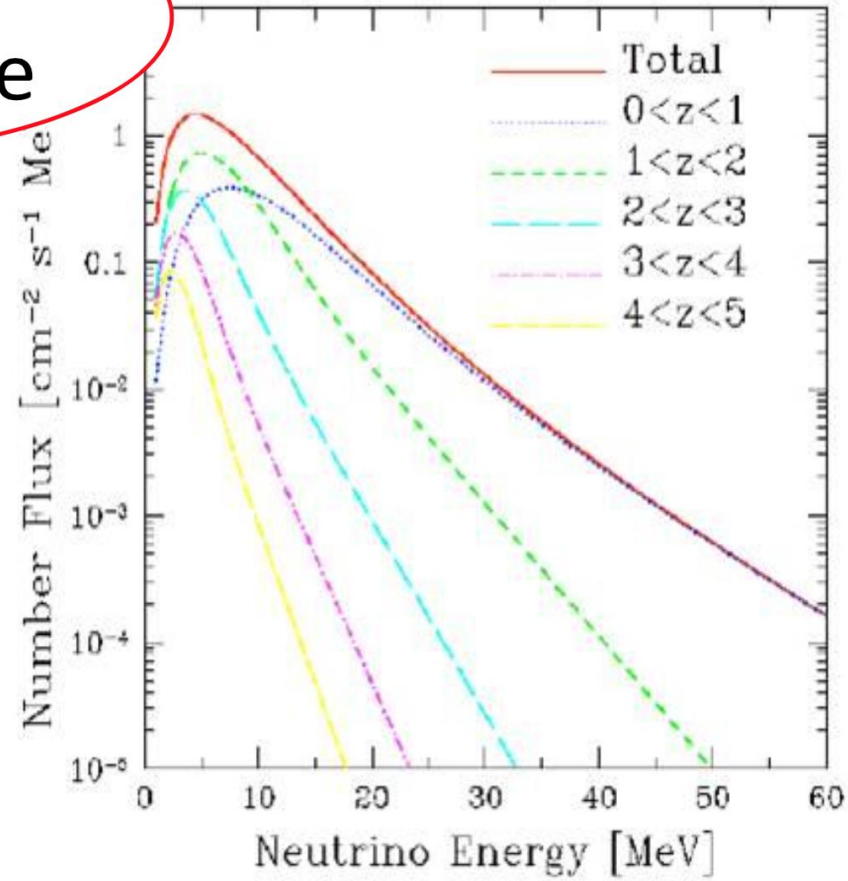
Gd(n, γ)Gd: Efficiency: >80%, fake probability: $< 10^{-4}$

Supernova Relic Neutrino (SRN)

10^{10} stellar/galaxy $\times 10^{10}$ galaxy $\times 0.3\%$ (become SNe) $\sim O(10^{17})$ SNe



Neutrinos
from past SNe

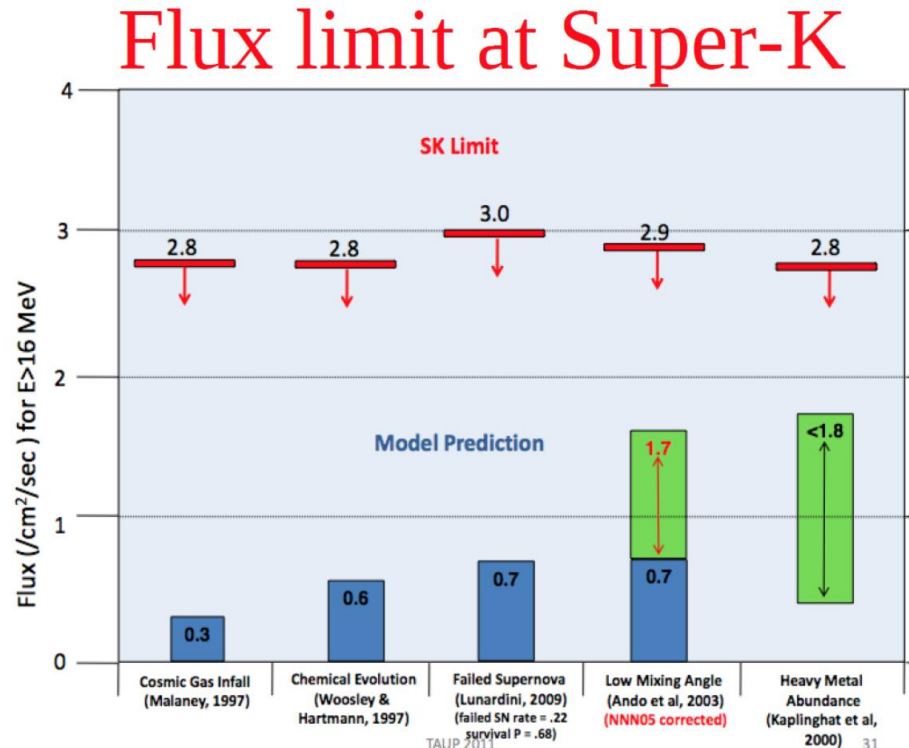


S.Ando, Astrophys.J. 607, 20(2004)

Theoretical flux prediction : $0.3 \sim 1.5 / \text{cm}^2/\text{s}$ (17.3MeV threshold)

Current limits from Super-K background dominated by "stealth" muon events and a threshold limited by muon spallation.

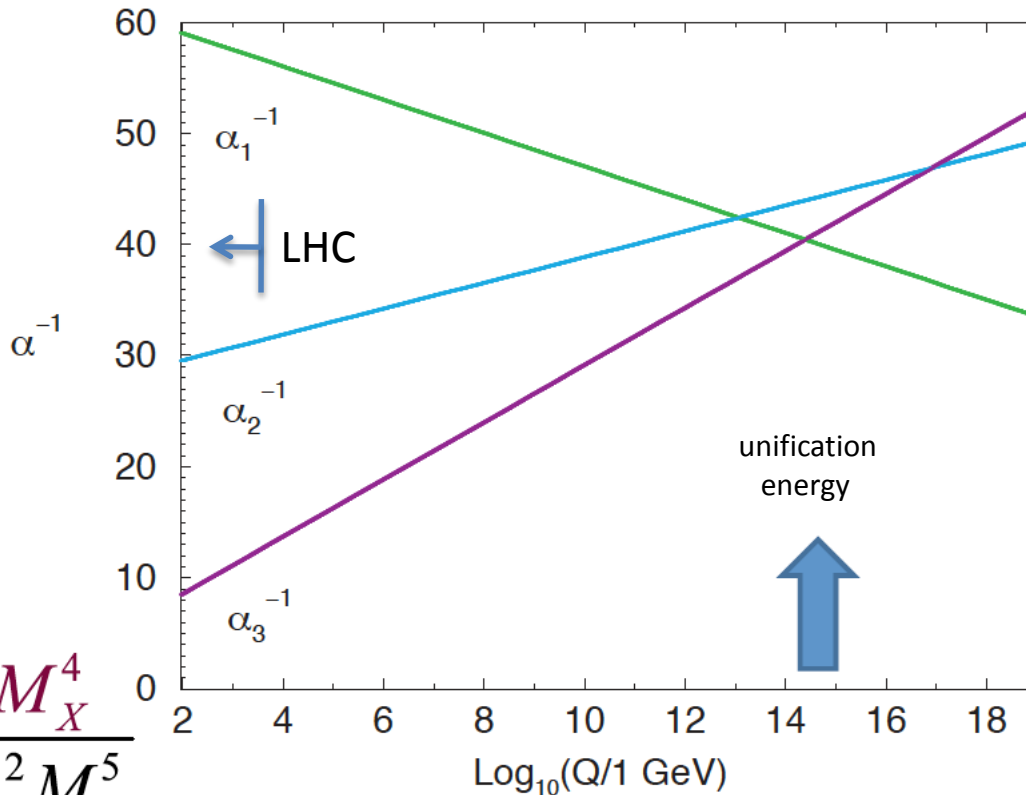
n-Tagging Inverse Beta Decay events from Diffuse SN should increase sensitivity to the levels implied by current models.



Phys. Rev. D 85, 052007 (2012)

Proton Decay: Is Normal Matter Stable?

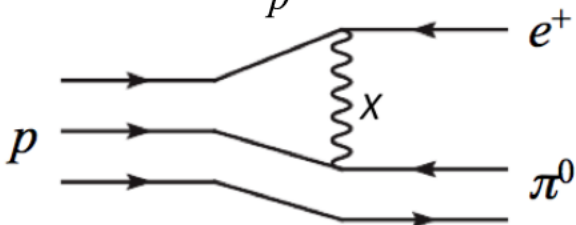
Gauge Coupling Unification



Three of the four forces of nature are thought to become similar in strength at very high energies – far above any conceivable accelerator

Simple unification theory ruled out by data – proton decay is an effective way to test such theories

$$\tau \approx \frac{M_X^4}{\alpha^2 M_p^5}$$

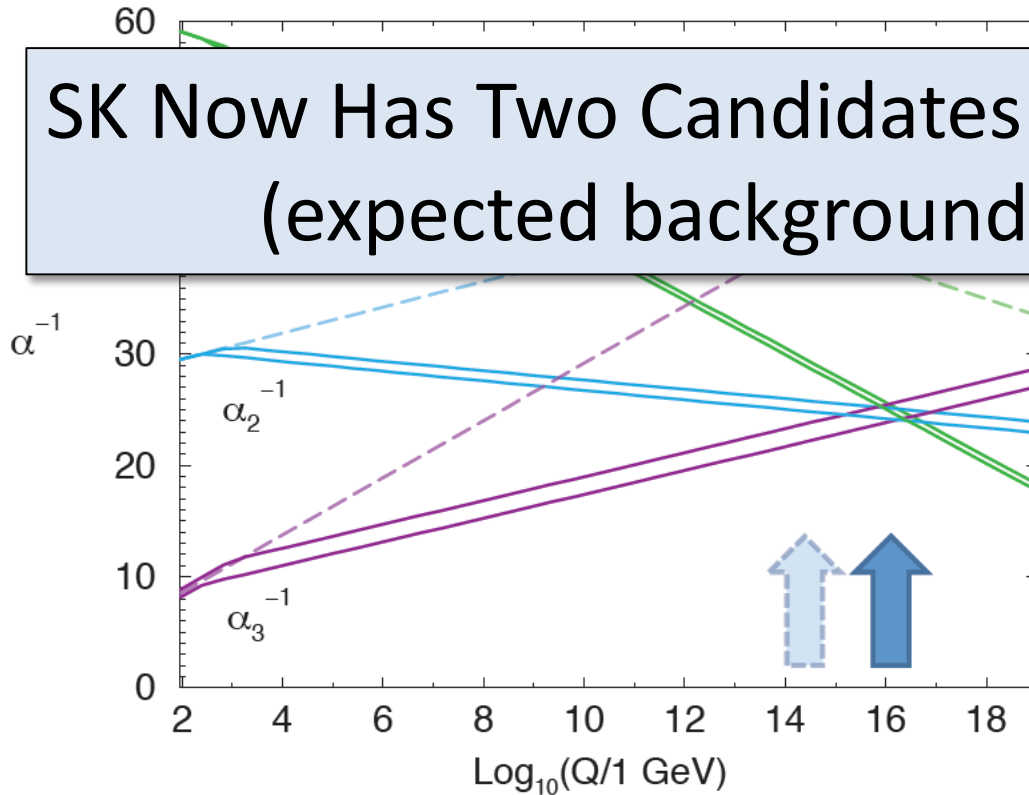


$$\tau(e^+ \pi^0) = 4.5 \times 10^{29 \pm 1.7} \text{ years (predicted)}$$

$$> 6.6 \times 10^{33} \text{ years (PRL 102 2009)}$$

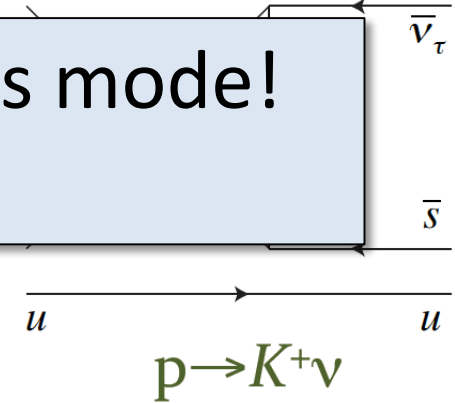
New theories (e.g. SUSY) can push up unification scale

SK Now Has Two Candidates in this mode!
(expected background 0.9)



Unification scale pushed up...

$$\tau(e^+ \pi^0) \approx 10^{35-38} \text{ years}$$



Example of a possible proton decay through supersymmetric particles.

Observation of virtual processes like proton decay is our **only known way** to access physics at these energies

Neutrons from Proton Decay in Water

- 2/10 of protons are free protons. **No neutrons.**
- 2/10 of protons are in $P_{1/2}$ shell. If they decay nucleus is already in the ground state. **No neutrons**
- 4/10 of protons are in $P_{3/2}$ shell. If they decay then a $P_{1/2}$ proton will drop down, giving a 6 MeV gamma. **No neutrons.** (Ejiri gives 94% B.R. for this)
- **~80% of proton decays should give neutrons only indirectly from FSI.** (such FSI usually makes them undetectable anyway) This is fairly model independent. Ejiri's more detailed estimate gives 81%
- Similar numbers for neutron decay.

Will Proton Decay Result in Neutrons?

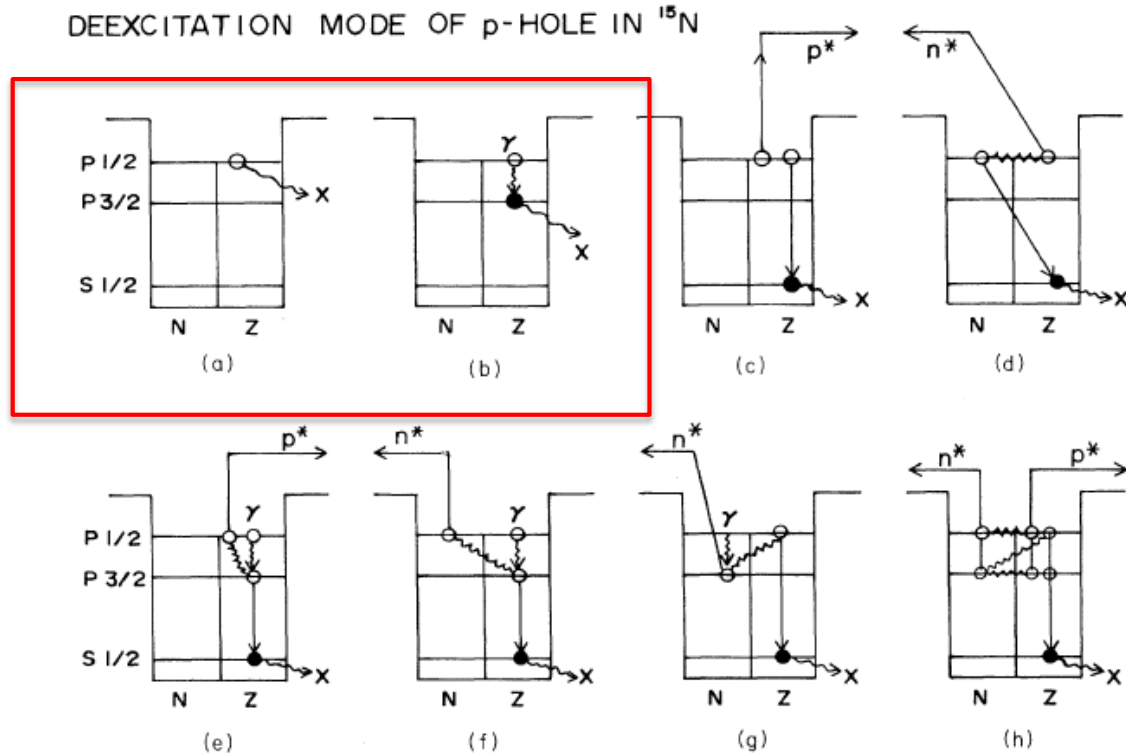


FIG. 1. Deexcitation scheme of a proton hole produced by proton decay ($p \rightarrow x$) in ^{16}O . N and Z stand for neutron and proton shells, respectively. p^* and n^* are protons and neutrons emitted into the continuum region.

Motivation

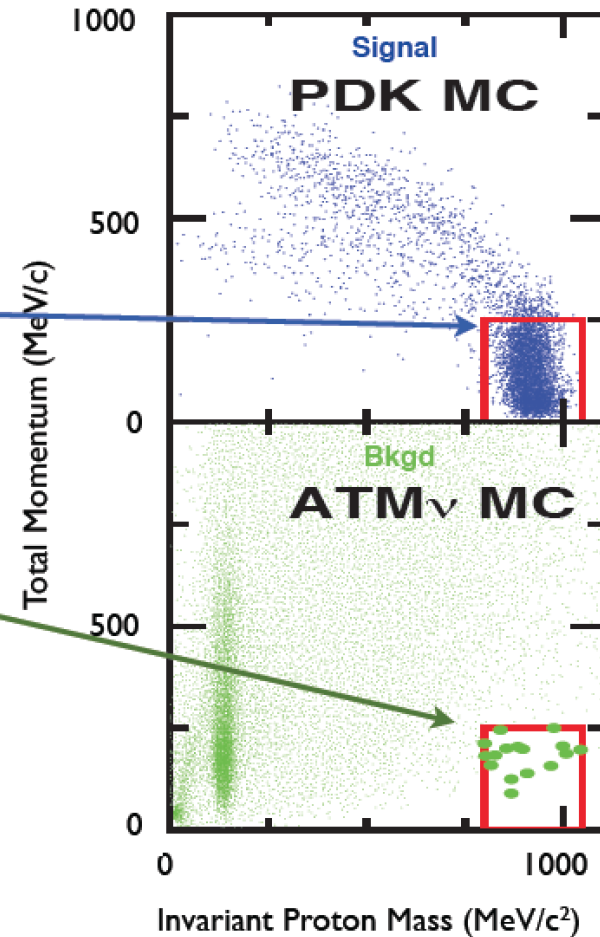
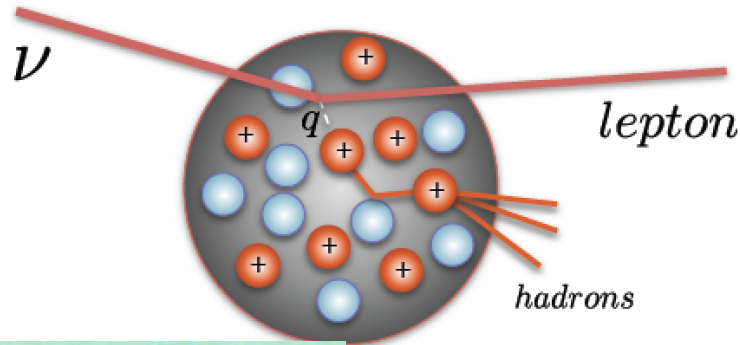
Backgrounds come almost exclusively from atmospheric neutrino interactions

Proton decay events are expected to only rarely produce neutrons in the final state.

High energy neutrino interactions typically produce neutrons in the final state

ANNIE Experiment

SK-IV



I. Anghel^{1,4}, G. Davies⁴, F. Di Lodovico¹¹, A. Elagin⁹, H. Frisch⁹, R. Hill⁹, G. Jocher⁵, T. Katori¹¹, J. Learned¹¹, R. Northrop⁹, C. Pilcher⁹, E. Ramberg³, M.C. Sanchez^{1,4}, M. Smy⁷, H. Sobel⁷, R. Svoboda⁶, S. Usman⁵, M. Vagins⁷, G. Varner¹⁰, R. Wagner¹, M. Wetstein⁹, L. Winslow⁸, and M. Yeh²

¹Argonne National Laboratory ²Brookhaven National Laboratory ³Fermi National Accelerator Laboratory ⁴Iowa State University
⁵National Geospatial-Intelligence Agency ⁶University of California at Davis ⁷University of California at Irvine
⁸University of California at Los Angeles ⁹University of Chicago ¹⁰University of Hawaii ¹¹Queen Mary University of London

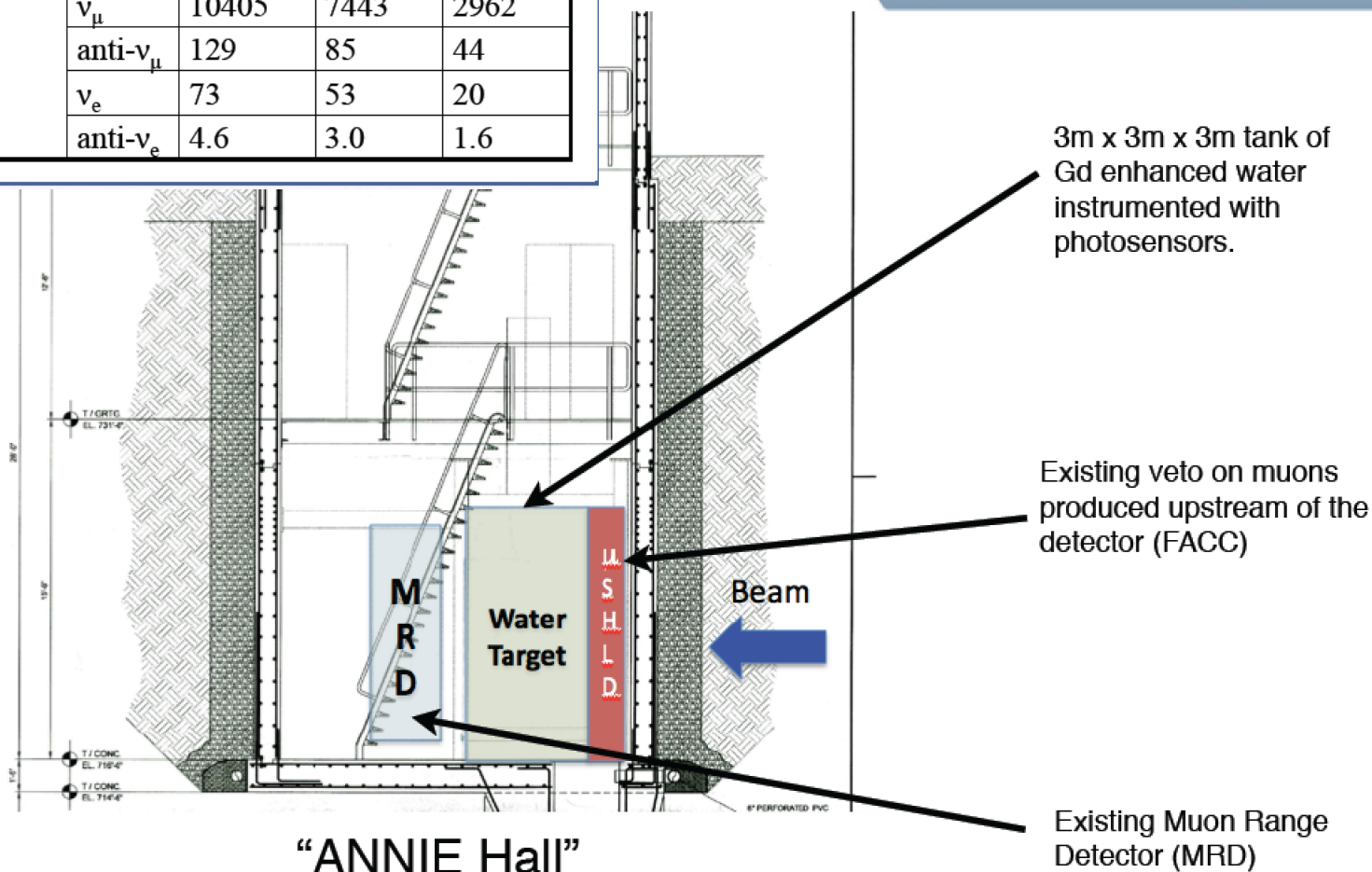


Rates Expected with 1×10^{20} POT exposure at SciBooNE pit

Djurcic

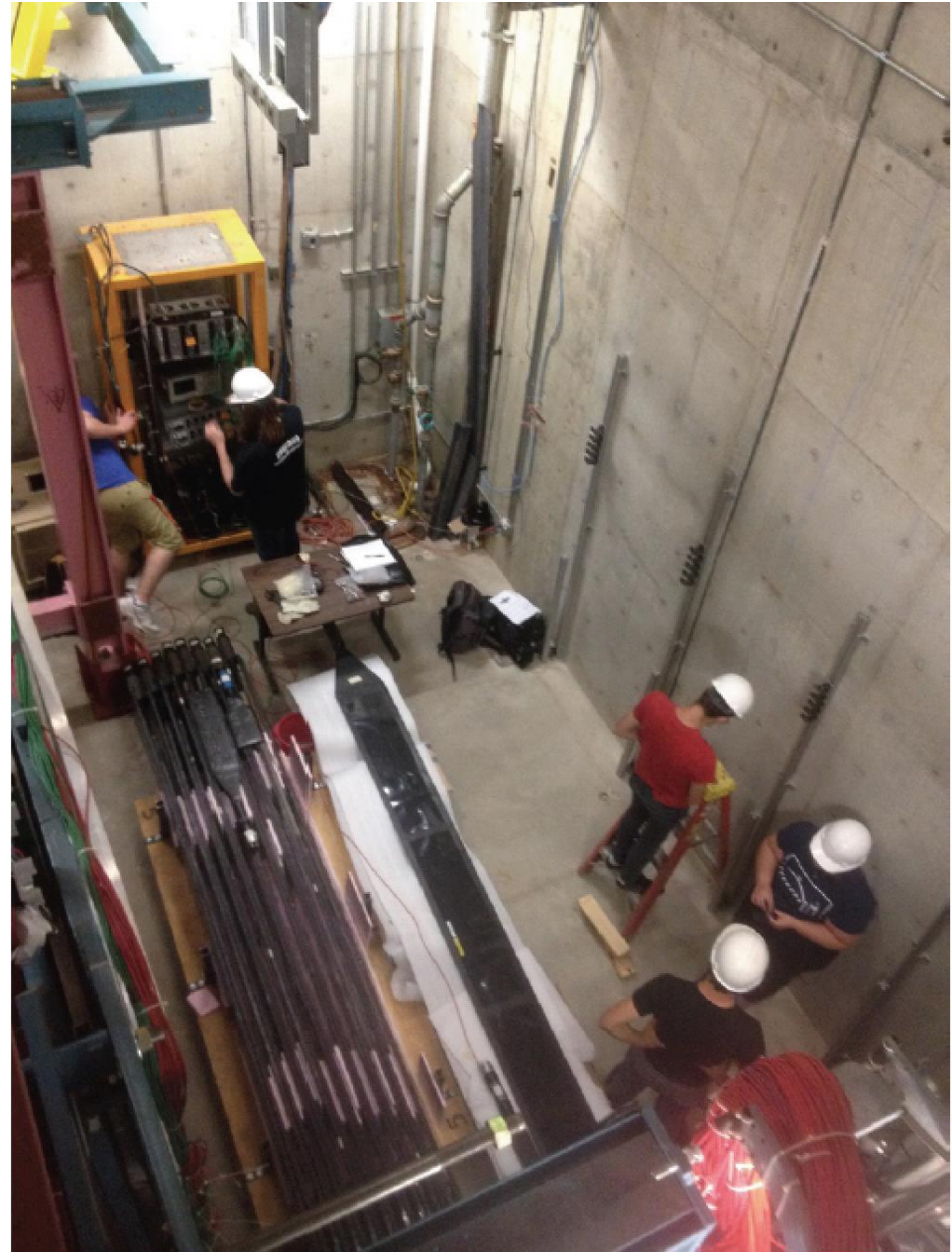


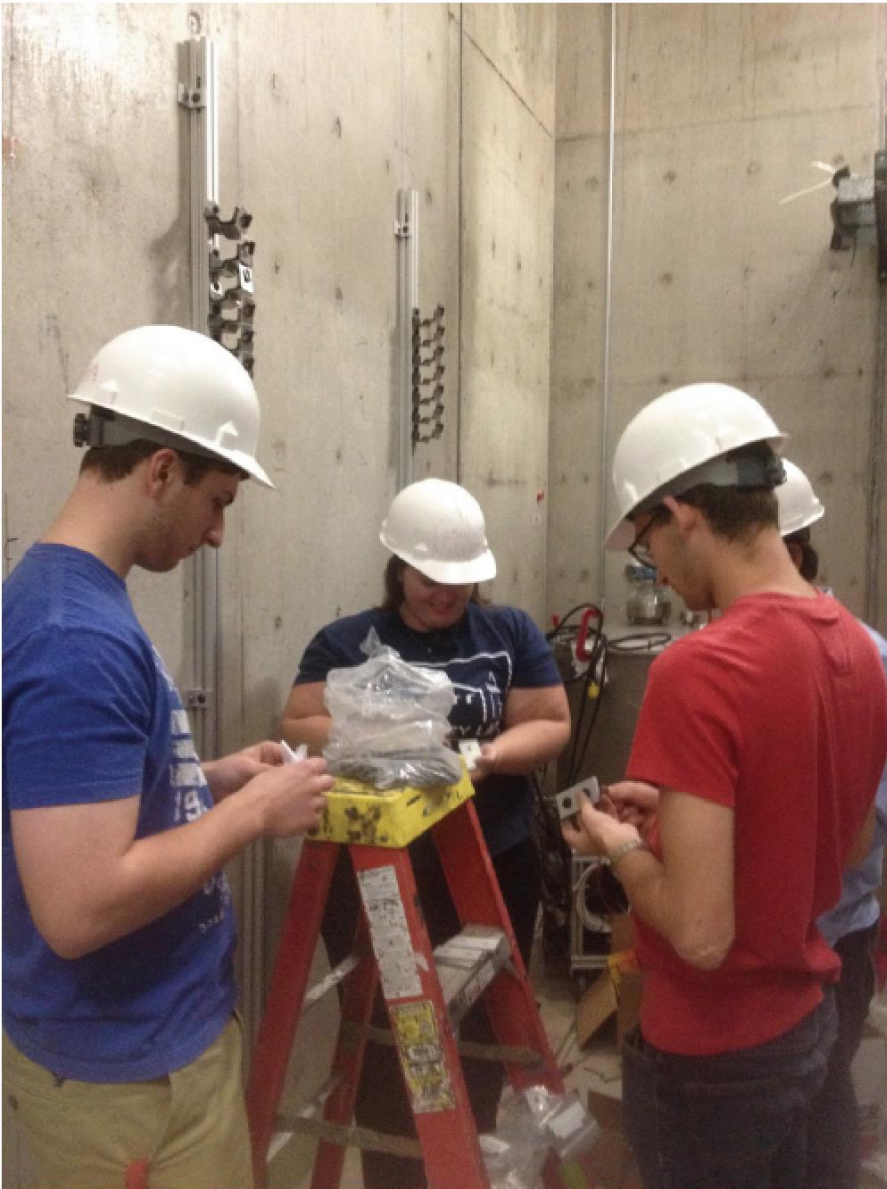
	Total Events [1/1ton/10 ²⁰ POT]	v-type	Total (per v-type)	Charged Current	Neutral Current
Booster Beam (v-mode, Target = CH ₂)	10419	ν_{μ}	10210	7265	2945
		anti- ν_{μ}	133	88	45
		ν_e	72	52	20
		anti- ν_e	4.4	3	1.4
Booster Beam (v-mode, Target = H ₂ O)	10612	ν_{μ}	10405	7443	2962
		anti- ν_{μ}	129	85	44
		ν_e	73	53	20
		anti- ν_e	4.6	3.0	1.6



“ANNIE Hall”

(formerly the SciBooNE pit)





For documentation and animation of the wall construction, go to:

https://cdcv.s.fnl.gov/redmine/projects/annie_experiment/wiki/Veto_design



Liquid Scintillator Detectors



Scintillation

- the physics definition of scintillation:
 - the process by which ionization produced by charged particles excites a material and light is emitted by the de-excitation
 - one of the most common detection techniques in nuclear and particle physics
 - earliest use by Crookes in 1903
 - a ZnS-coated screen scintillates when struck by α particles
 - then Curran and Baker in 1944
 - coated a photomultiplier tube with ZnS producing the first scintillation counter that didn't require the human eye
 - since then many forms of scintillation detectors, large and small, have been developed
 - the scintillation process differs in different materials (e.g. inorganic crystals, organic liquids, noble gases and liquids, plastic scintillators)
 - we'll briefly examine each type...
-

Aside: Definitions

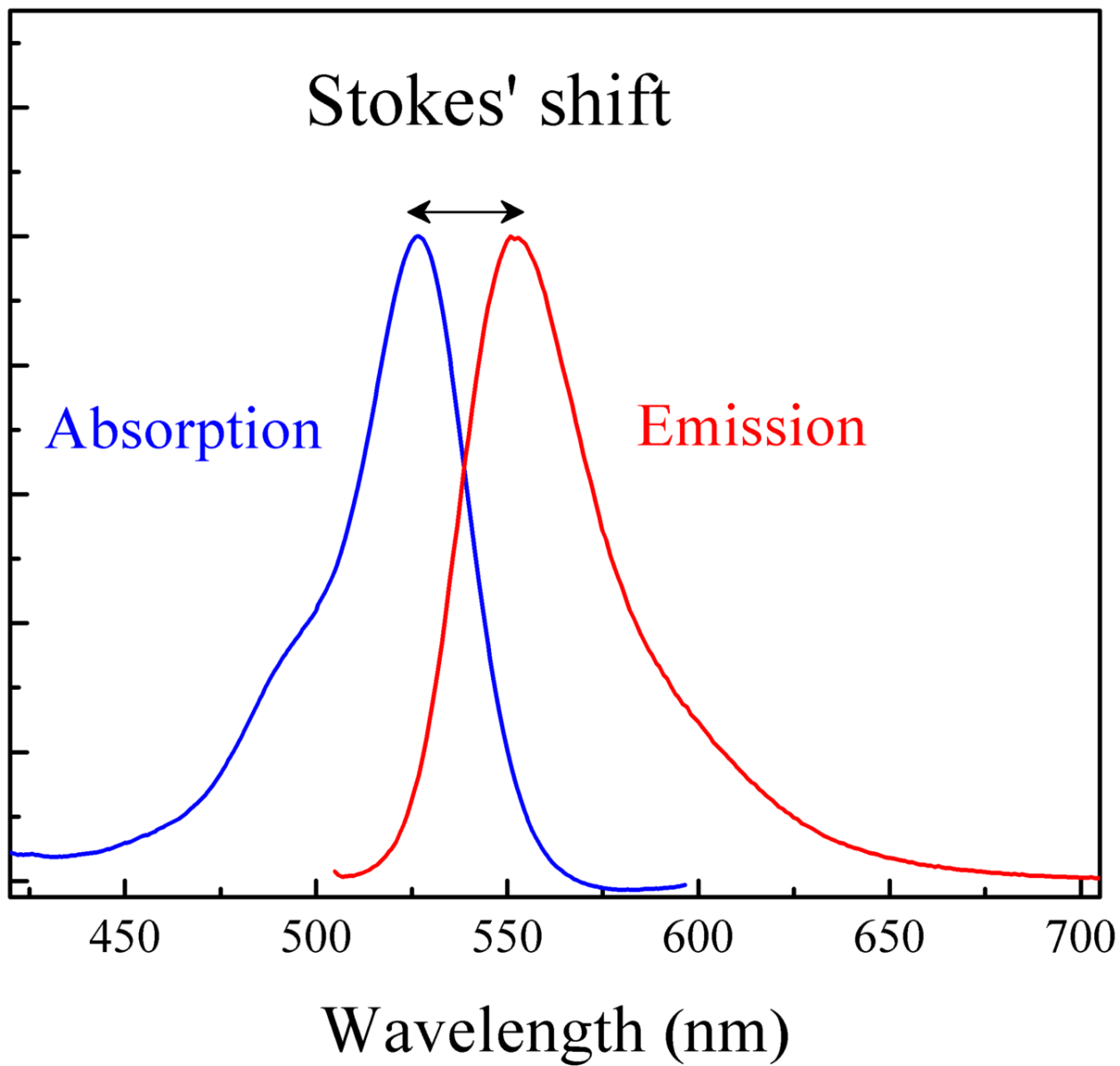
- when you excite a material (not thermally) and it subsequently gives off light, that is *luminescence*
 - how it's excited determines the type of luminescence (e.g. *photoluminescence*, *chemiluminescence*, *triboluminescence*)
 - *fluorescence* is photoluminescence or scintillation (i.e. excitation produced by ionizing radiation) that has a fast decay time (ns to μs)
 - *phosphorescence* is the same, only with a much slower decay time (ms to seconds)
-

Stokes Shift

an important, general concept to keep in mind for all scintillators

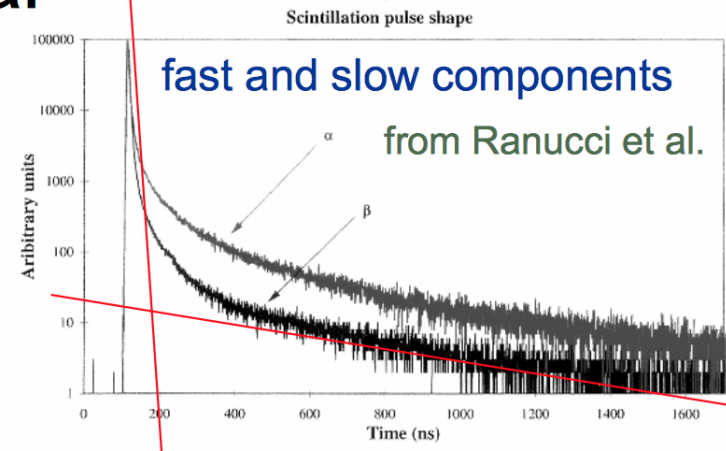
- emitted photons are at longer wavelengths (smaller energies) than the energy gap of the excitation
- the processes that produce this “Stokes shift” are different in different scintillating materials
- this allows the scintillation light to propagate through the material
 - emitted photons can't be self-absorbed by exciting the material again

from Wikipedia



Scintillator Characteristics of Interest

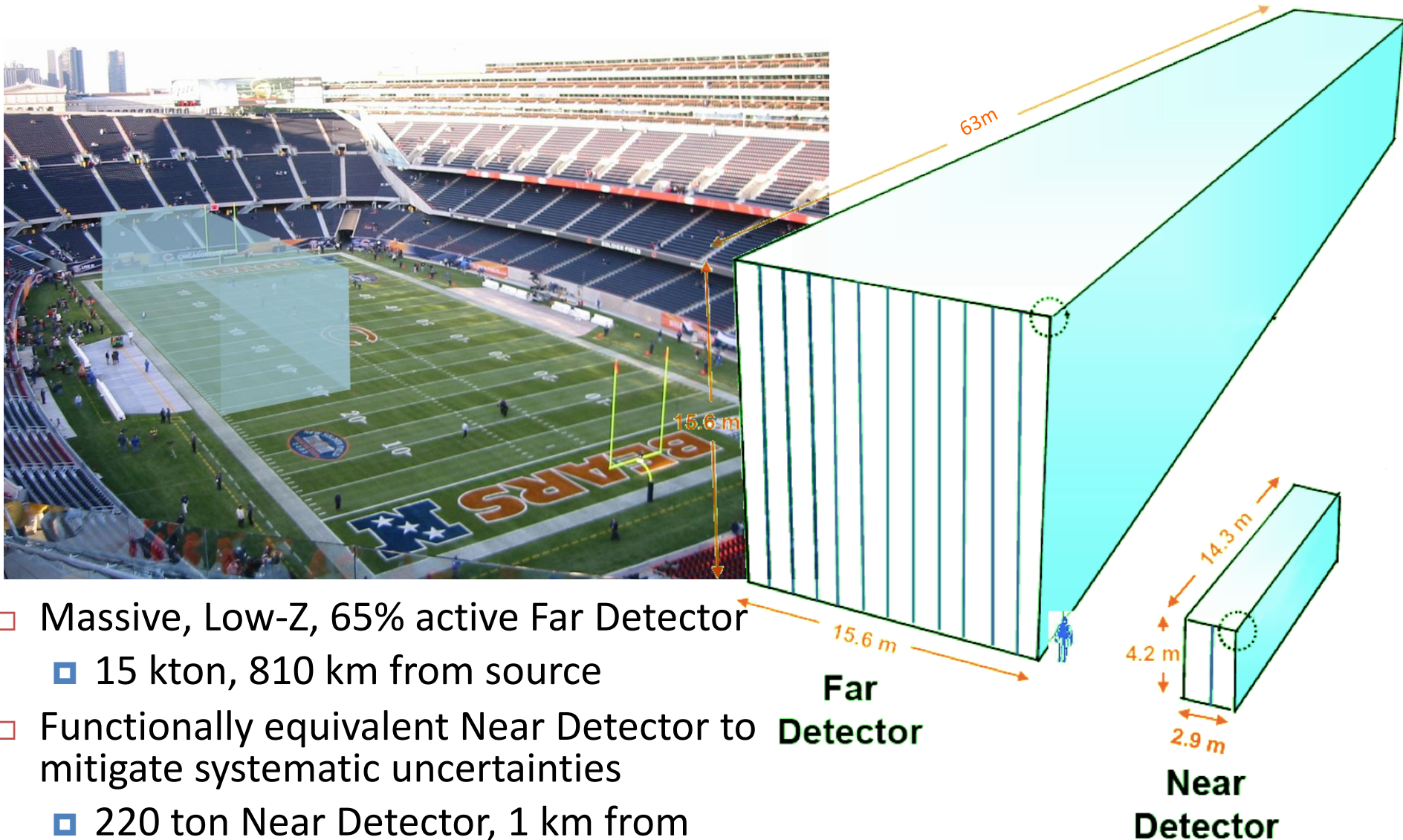
- **light yield:** high efficiency for converting ionization energy to light output [photons/MeV]
- **emission spectrum:** overlaps with spectral response of light detector (e.g. PMT)
- **decay time:** can have several time constants
- **density and Z:** determine response to γ , e^- and other electromagnetic processes



Organic Scintillators

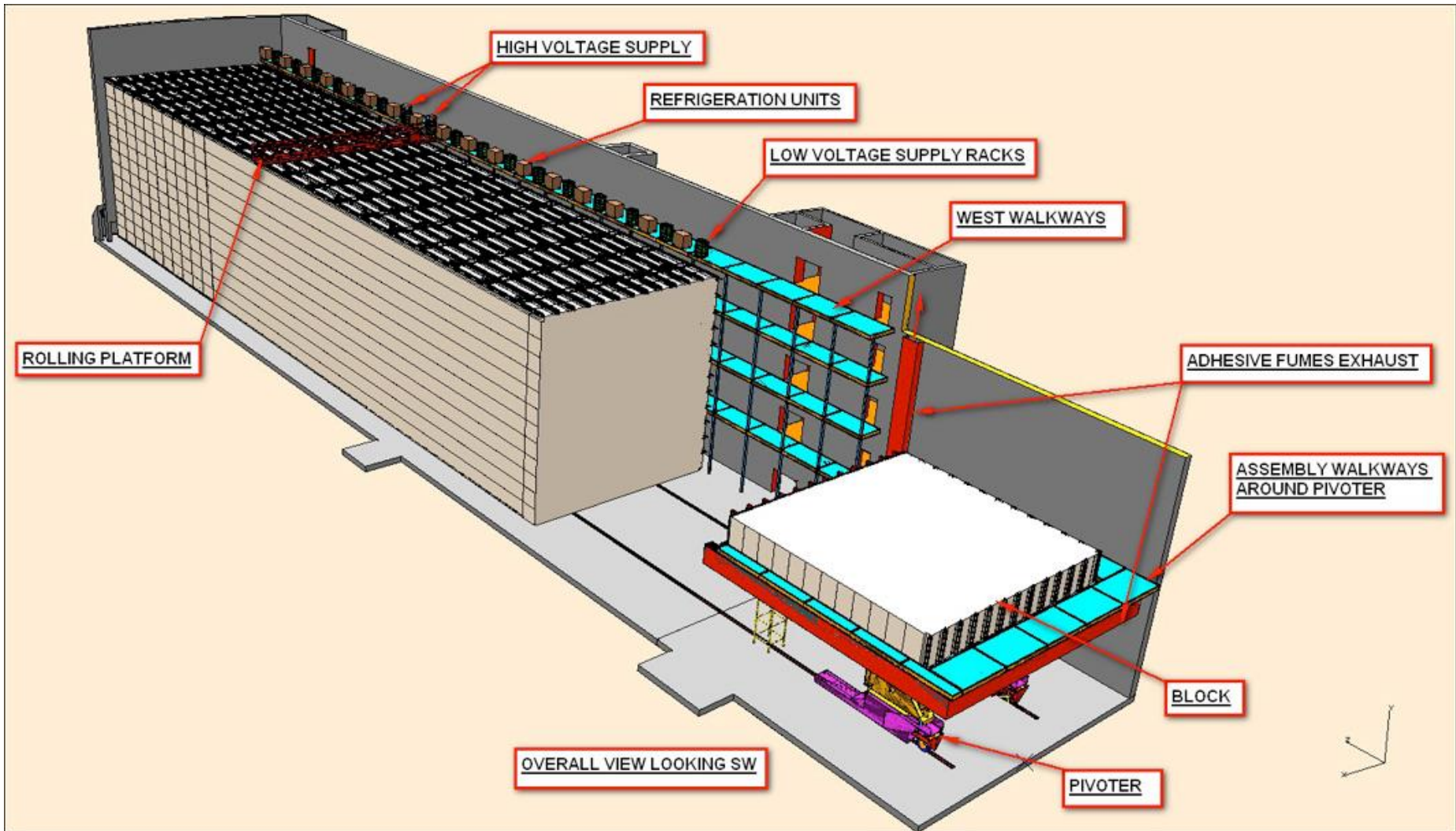
- the scintillation mechanism is determined by the chemistry and physics of the benzene ring
 - an organic scintillator will thus scintillate whether it's in a crystal form, is a liquid, a gas, or imbedded in a polymer
 - all organic scintillators in use employ aromatic molecules (i.e. have a benzene ring)
 - we'll examine the photophysics of aromatic molecules in detail...
-

The NOvA Detectors



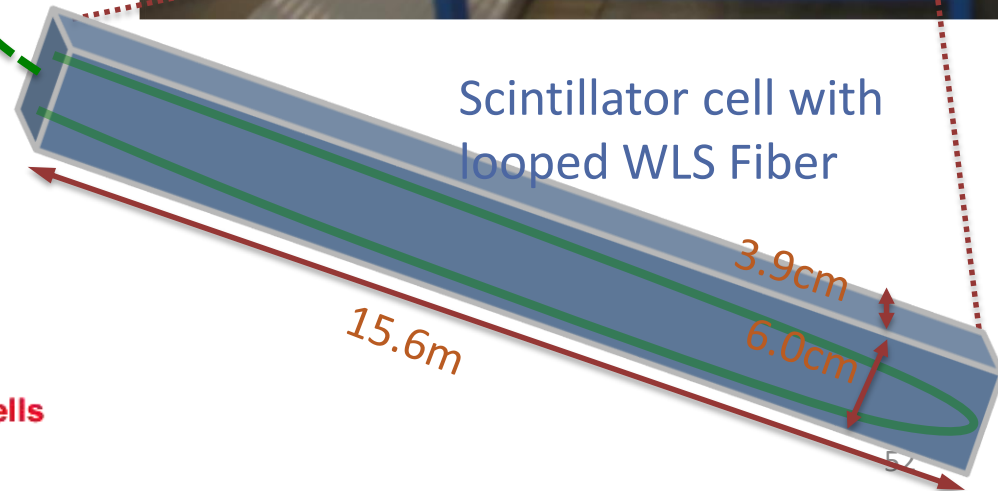
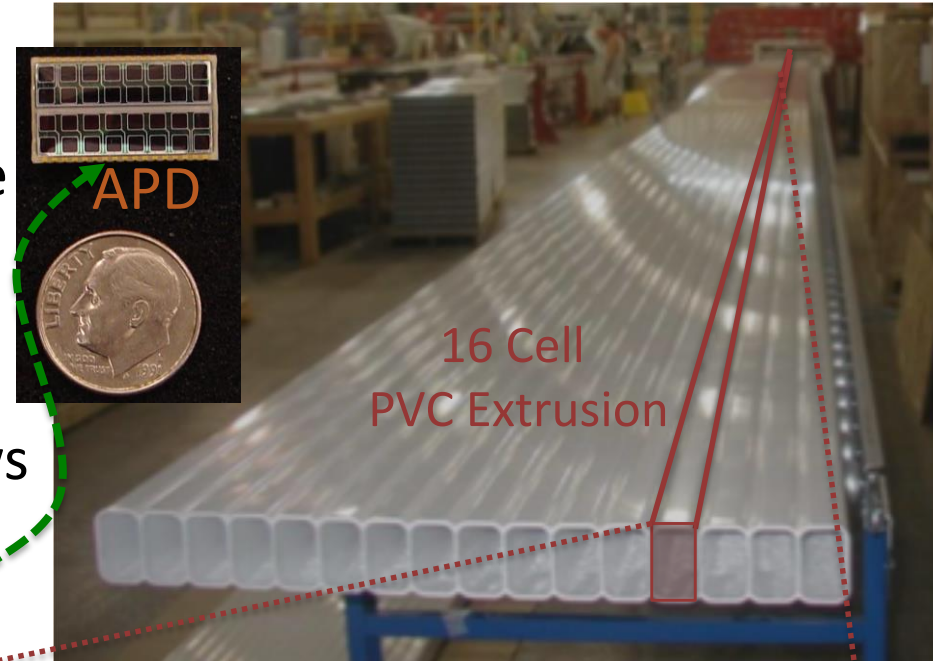
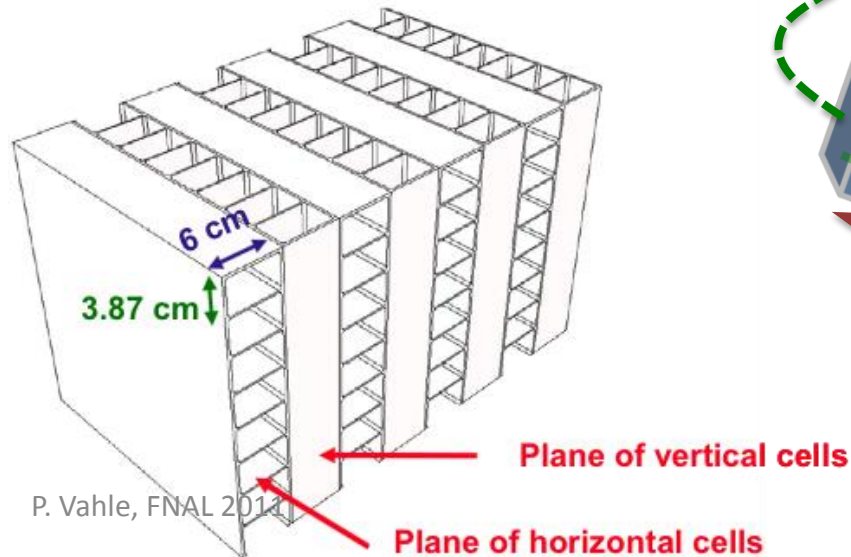
- Massive, Low-Z, 65% active Far Detector
 - 15 kton, 810 km from source
- Functionally equivalent Near Detector to mitigate systematic uncertainties
 - 220 ton Near Detector, 1 km from source

15,000 tons of LS



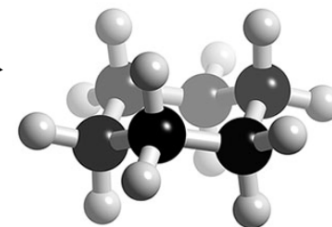
Detector Technology

- PVC extrusion + Liquid Scintillator
 - ▣ mineral oil + 5% pseudocumene
- Read out via WLS fiber to APD
 - ▣ muon crossing far end=38 PE
- Layered planes of orthogonal views
- $0.15 X_0$ per layer



Carbon Bonds

- carbon: $1s^2 2s^2 2p^2$ electronic ground state
- carbon bonds: $1s^2 2s^1 2p^3$ and the 2s and 2p orbitals are hybridized
 - sp^3 hybridization is tetrahedral (e.g. diamond, methane, cyclohexane)
 - not luminescent
 - sp^2 hybridization is planar
 - p_z orbital is unchanged
 - double-bonded carbon (e.g. ethylene, benzene)
 - is luminescent and the basis of organic scintillators
 - sp hybridization is linear
 - triple-bonded carbon (e.g. acetylene)
 - also luminescent



from 3Dchem.com

sp³ and sp² hybridization

P orbitals are shaped like figure eights and electrons in p orbitals are slightly farther away from the nucleus than electrons in s orbitals. The **sp³** orbital is 25% s character and 75% p character (1 2s orbital, + 3 2p (2p_x, 2p_y, 2p_z)). For sp² **hybridization** there is 33% s character and 66% p character.

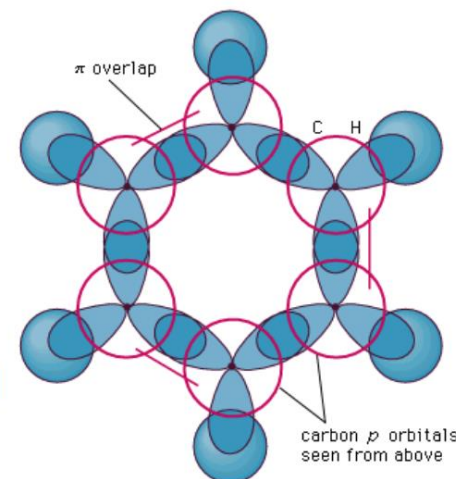
[Sp³ hybridization - Khan Academy](#)

[https://www.khanacademy.org/.../sp³-hybrid-orbital-jay-fin...](https://www.khanacademy.org/.../sp3-hybrid-orbital-jay-fin...) Khan Academy ▾

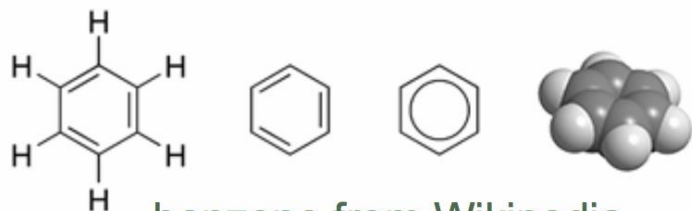
σ -Bonds and π -Bonds

from Encyclopedia Britannica web

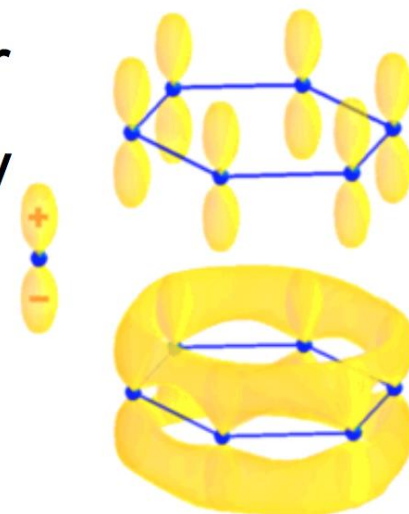
- σ -bonds are in the plane, bond angle 120° , from sp^2 hybridization
- π -orbitals are out of the plane
 - in the benzene ring (and other carbon double bonds) they overlap each other
 - result is the π -electrons are completely delocalized



©1996 Encyclopaedia Britannica, Inc.



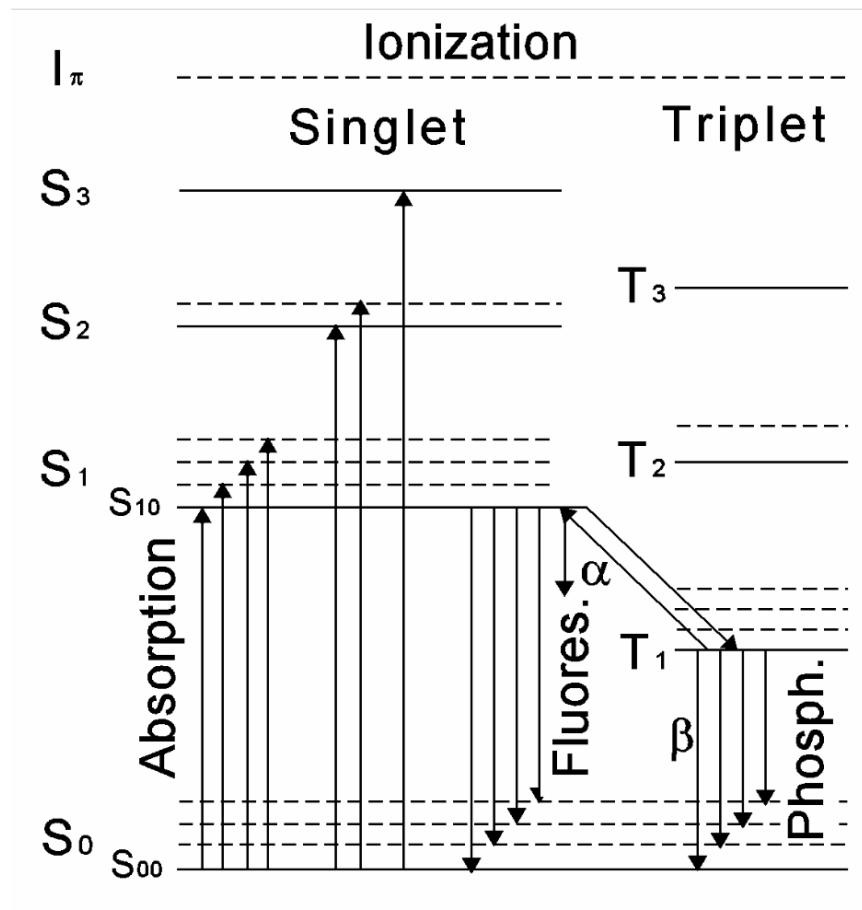
benzene from Wikipedia



from <http://www.monos.leidenuniv.nl/smo/index.html?basics/photophysics.htm>

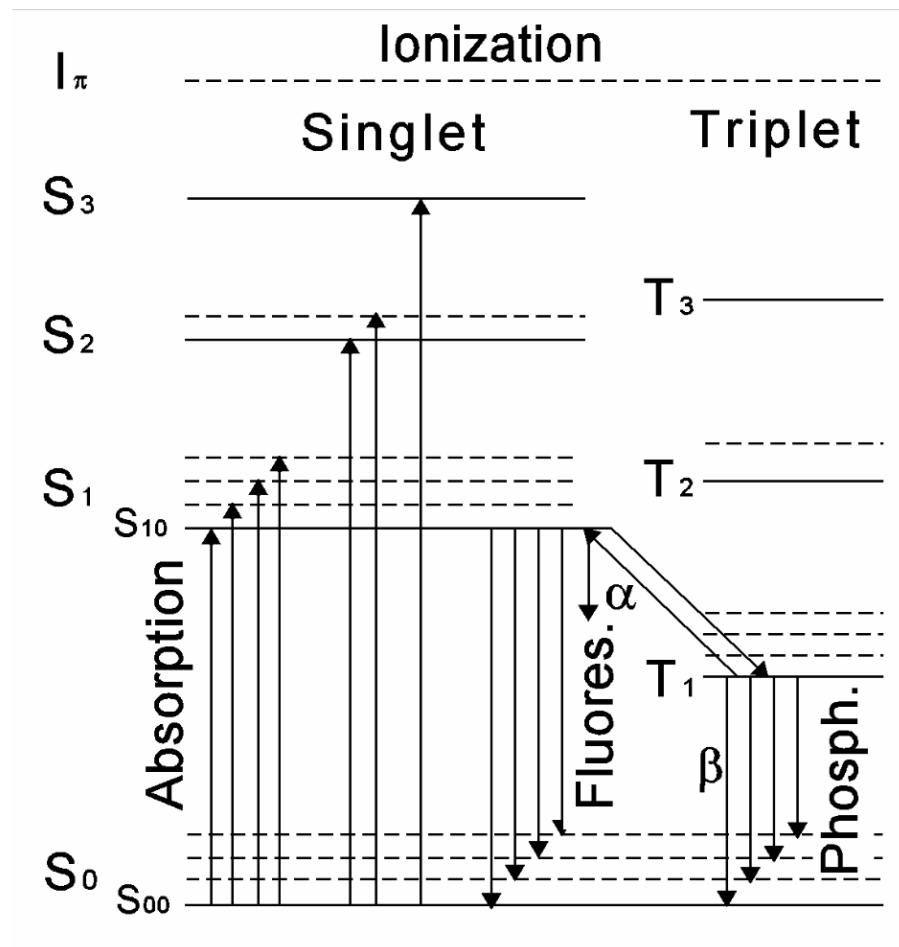
π -Electronic States

- theory of π -electron excited states uses a “perimeter free-electron model”, with electrons going around the ring in both directions, spin up and down
- absorption spectroscopy confirms the theory
- the π -electronic states diagram can be deduced from theory and spectroscopy
 - has vibrational sub-levels (e.g. S_{10} , S_{11}) with ~ 0.16 eV spacing
- after absorption of a photon or excitation by ionization, the molecule undergoes vibrational relaxation (or internal conversion or degradation) to S_{10}



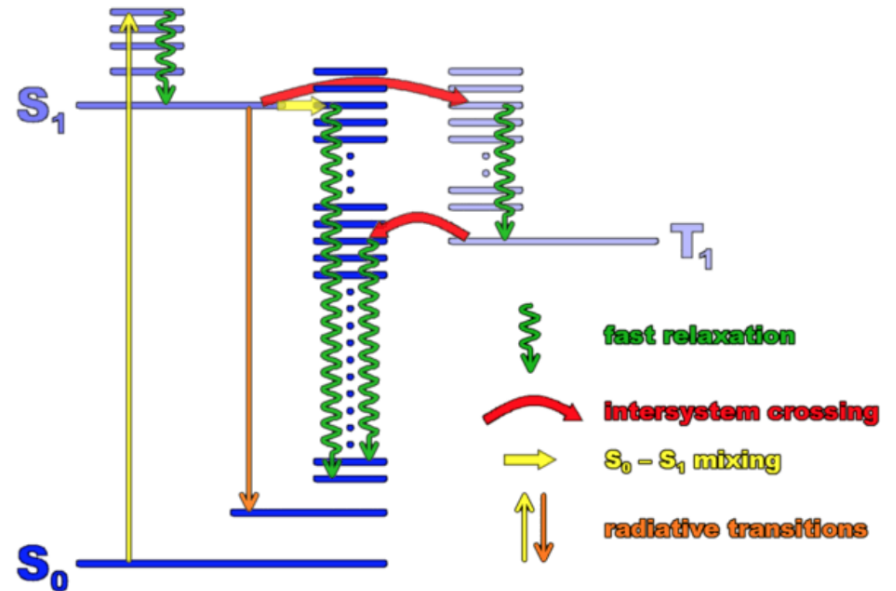
Scintillation in Organic Molecules

- the excited S_{10} state decays radiatively to vibrational sub-levels of the ground state
- the S_{10} lifetime is \sim ns
- thus the fluorescence emission spectrum is roughly a “mirror image” of the absorption spectrum (same spacing)
- emitted photons have less energy than S_{00} - S_{10} – that’s the important **Stokes shift**
- there is no S_2 - S_0 emission because internal conversion is efficient and fast \sim ps
- there are σ -electronic excited states too; but they are at higher energies above S_3



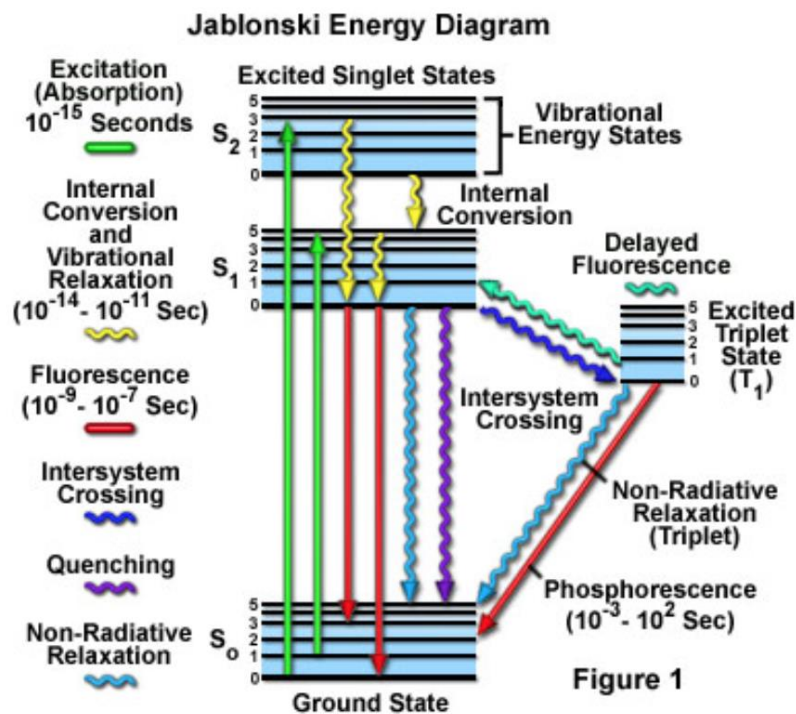
Competing Processes

- non-radiative decay
 - due to overlap of an S_0 vibrational state with S_1
- ratio of the radiative rate to the total decay rate (radiative plus non-radiative) is the *fluorescence quantum yield*
- quantum yields of 0.8 or greater are typical for “good” scintillators



Competing Processes cont'd

- an intersystem crossing can occur, populating the triplet state
 - S_0 - T_1 and T_1 - S_0 cannot occur directly due to angular momentum and parity selection rules
- T_1 has a long lifetime since decay to the S_0 is forbidden
 - this is phosphorescence
- delayed fluorescence from the triplet can also occur
 - $T_1 + T_1 \rightarrow S_1 + S_0 + \text{phonons}$
 - $\sim \mu\text{s}$

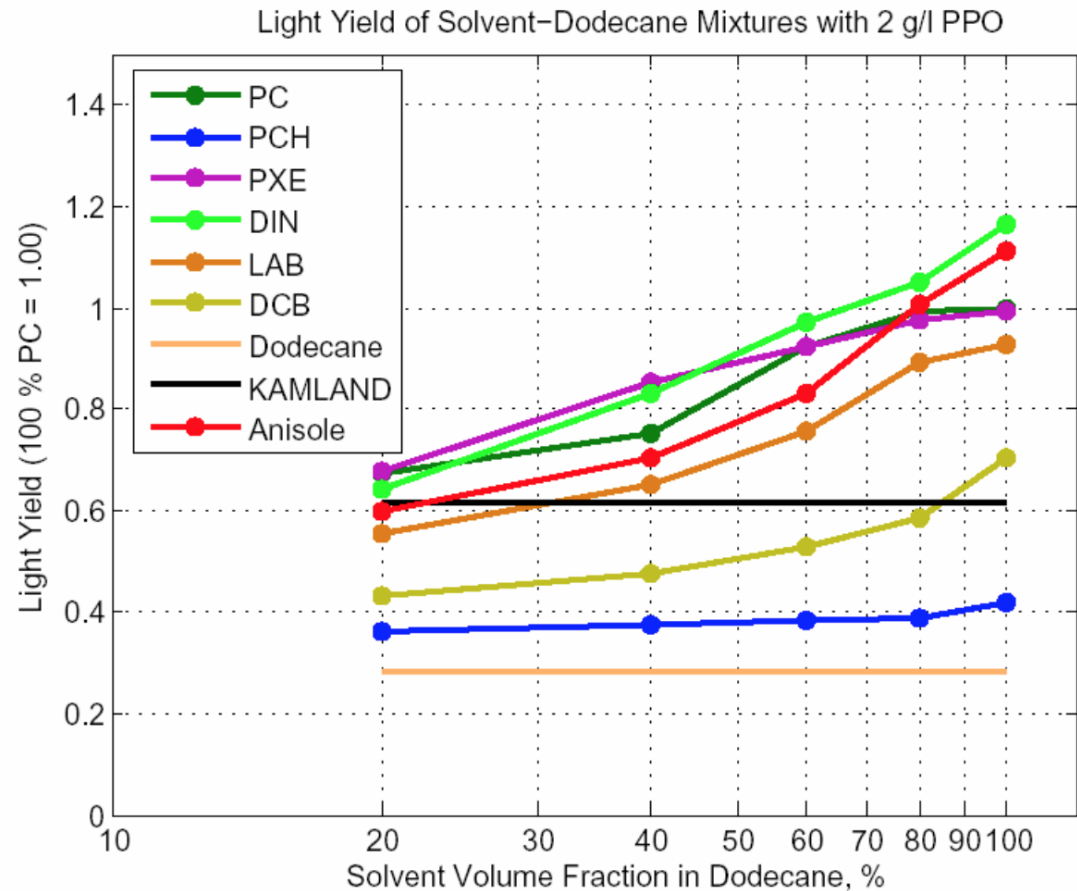


Organic Liquids

- **any aromatic organic solvent is a scintillator**
 - e.g. benzene, xylene, pseudocumene (1,2,4-trimethylbenzene)
- **to improve the performance, a scintillating solute is dissolved in the solvent – this is called the fluor**
 - e.g. PPO (2,5-diphenyloxazole), p-terphenyl, butyl-PBD
 - typically g/L quantities
- **this accomplishes the following:**
 - efficient non-radiative transfer of excitation energy from the solvent to the fluor
 - high fluorescence quantum yield of the fluor
 - emission of fluor at even longer wavelengths compared to the absorption of the solvent, further reducing self-absorption
 - typically fast decay times
- **if desired a secondary wavelength shifter (WLS) can be added**
 - e.g. bisMSB, POPOP
 - absorbs the light from the fluor and re-emits it at longer wavelengths that match the spectral response of the light detector
 - typically 10's mg/L quantities
- **the whole cocktail can be diluted in an optically-inert liquid**
 - e.g. mineral oil, dodecane
 - usually done for chemical reasons, though good for transparency too

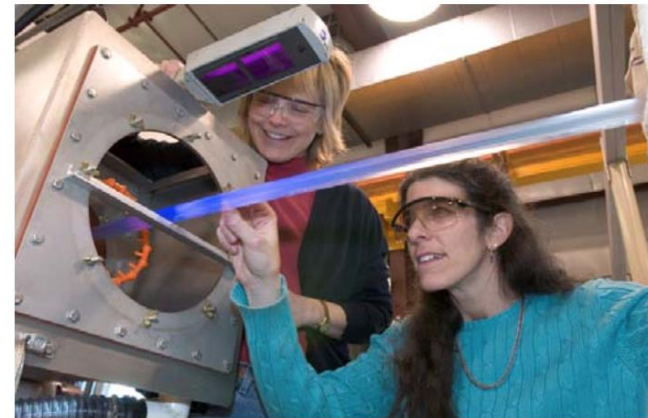
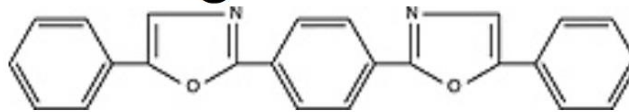
Diluting Liquid Scintillator

- diluting the cocktail in mineral oil or dodecane reduces the light output
- but note that 20% solvent concentration still exhibits ~70% of the light output
- it's not just the fraction of the excitation that was imparted to the solvent
- strongly suggests non-radiative energy transfer takes place between the dodecane and the solvent and/or fluor



Plastic Scintillator

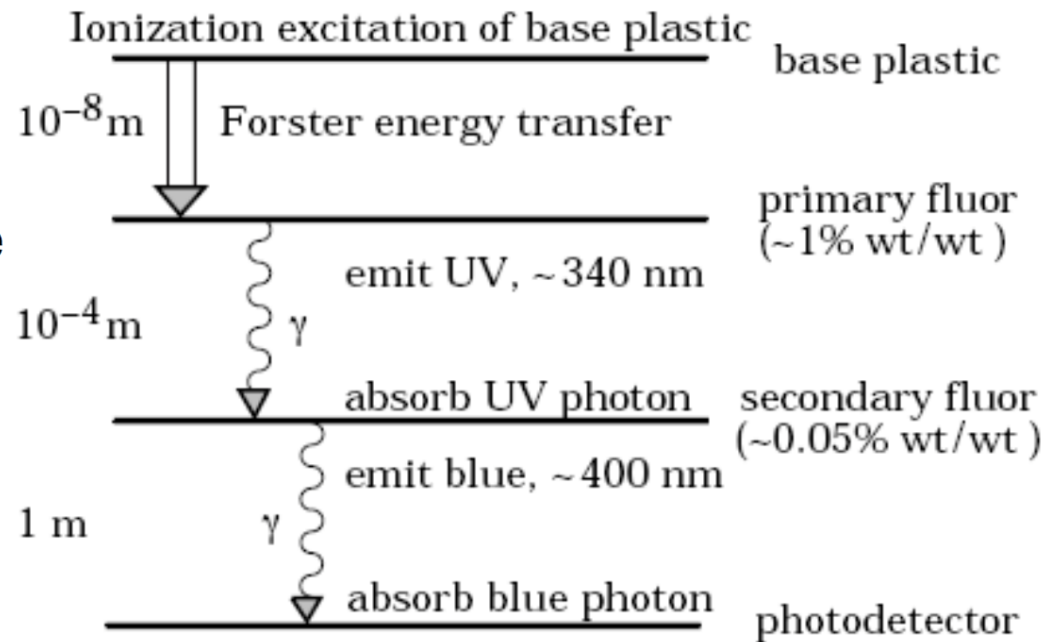
- very common in nuclear and particle physics
- polymer base is typically PVT (polyvinyltoluene) or polystyrene (PVT-based slightly higher light yield)
 - aliphatic plastics (e.g. acrylic) can be used also, but have half the light yield of aromatic plastics
- primary solute is dissolved in the plastic
 - e.g. PBD, p-terphenyl, PBO
- secondary wavelength shifter
 - e.g. POPOP
 - e.g. BBQ
- same scintillation mechanism



extruded plastic scintillator
for MINERvA from Wikipedia

Scintillation in a Plastic

- high concentrations of fluor shorten the decay time and increase the light yield
- at 1% fluor concentration, the average distance between a fluor molecule and an excited base is 10 nm, less than a wavelength of light
- Förster dipole-dipole resonance energy transfer efficiently excites the fluor with a faster decay time and high quantum yield



from Particle Data Group, Review of Particle Detectors

Plastic Scintillator: Advantages and Disadvantages

- reliable, robust, cheap, easy to fabricate into desired shapes
- can be used to detect neutrons by detecting proton recoils (as for a liquid scintillator)
- **drawbacks:**
 - subject to aging
 - degrades upon exposure to some chemicals (e.g. oils, solvents, fingerprints)
 - suffers radiation damage
 - not resistant to high temperatures
 - surface crazing affects light propagation because a plastic scintillator usually relies on total internal reflection at the surface



Care of Plastic Scintillator

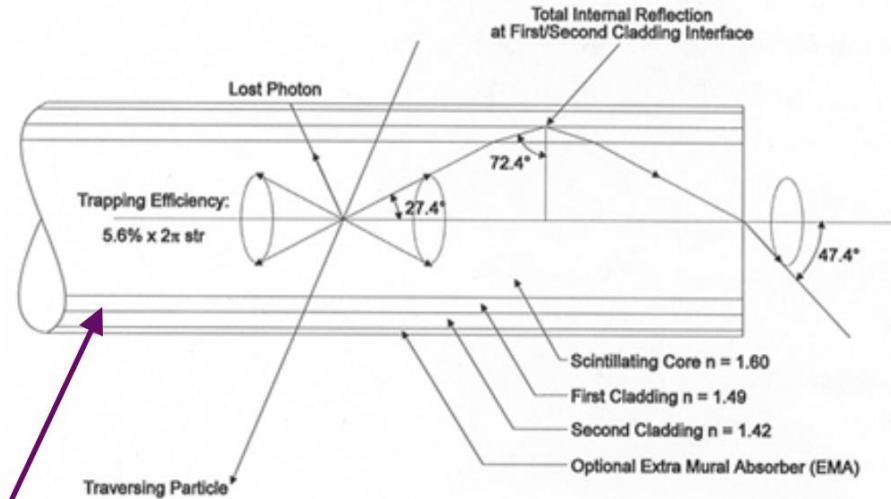
Jeff Wilkes (UW) cleaning up a plastic scintillator from NSF Public Affairs web page



Scintillating Fibres

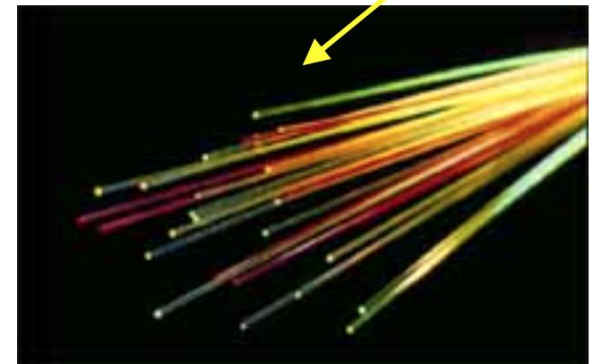
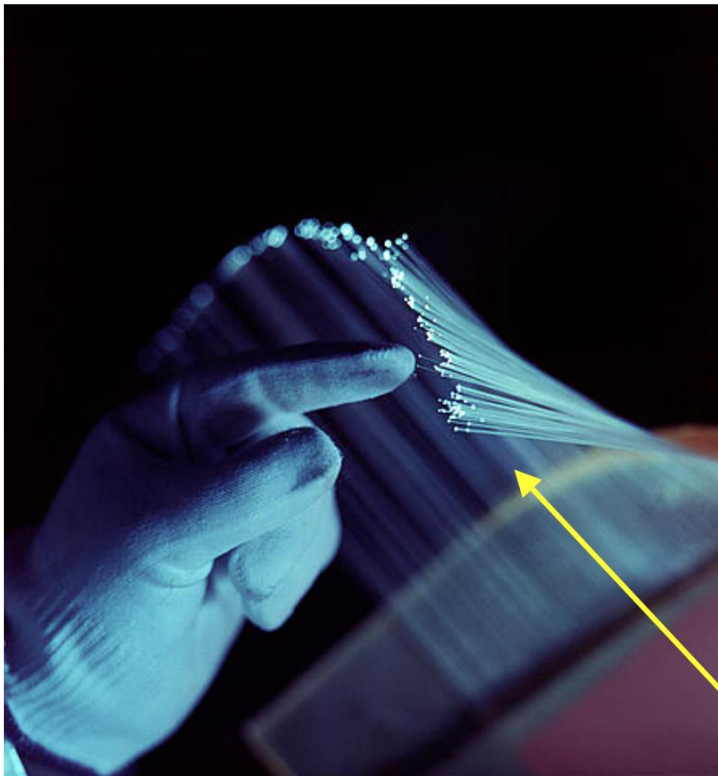
- the fibre core is
 - glass with activator
 - polystyrene with fluor

A Typical Round Multi Clad Scintillating Fiber



PMMA first cladding
fluorinated PMMA second cladding

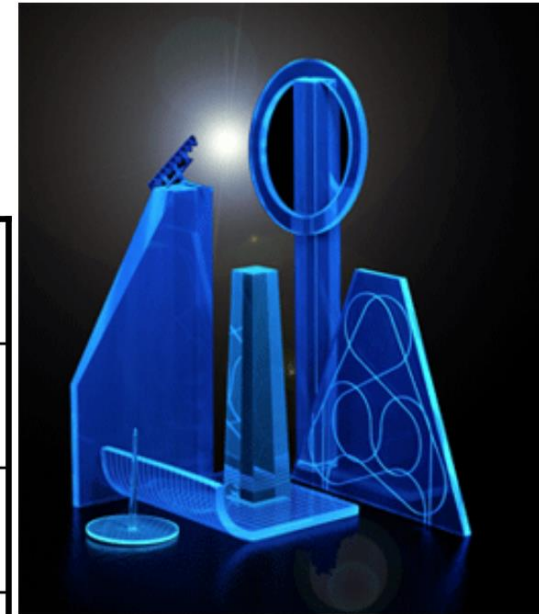
polystyrene core
fibres from Saint-
Gobain Crystals



glass fibers from Pacific Northwest National Lab

Comparing Plastic to Liquid

scintillator	light output	peak λ	decay constant	attenuation length	index of refraction	density [g/cm ³]
BC-400	65%	423 nm	2.4 ns	250 cm	1.58	1.032
BC-404	68%	408 nm	1.8 ns	160 cm	1.58	1.032
BC-416	38%	434 nm	3.3 ns	400 cm	1.58	1.032
BC-428	36%	480 nm	12.5 ns	150 cm	1.58	1.032
PC-based liquid	80%	425 nm	2.5 ns		1.505	0.877
30% PC diluted in mineral oil	60%	425 nm	3 ns	>500 cm	1.48	0.86
LAB-based	75%	425 nm	3.5 ns		1.47	0.87



from Saint-Gobain Crystals

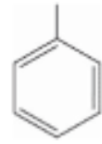
note: values for plastics are for commercial products while liquid scintillator cocktails are representative values that can be adjusted by selecting the fluor, the wavelength shifter and their concentrations

note: light yield is quoted as a fraction of anthracene, which is 17,400 photons/MeV

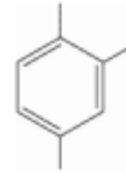


Intermission

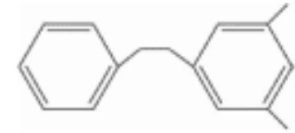
Solvents



Toluene



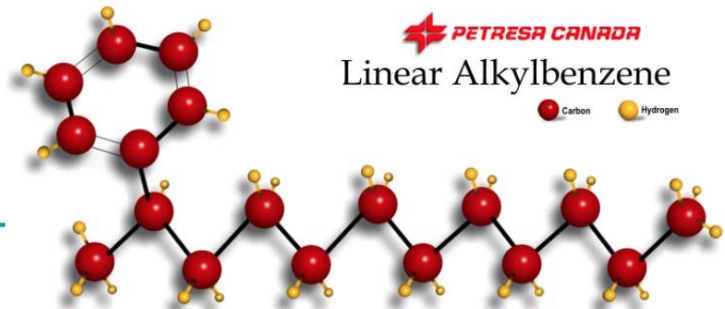
Pseudocumene



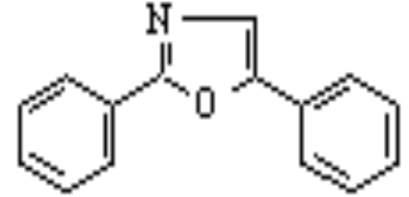
PXE (phenyl xylylene)

from http://nationaldiagnostics.com/article_info.php/articles_id/117

- pseudocumene is common
 - used in KamLAND, Borexino, Palo Verde, MACRO, many commercial liquid scintillator cocktails
- safer solvents are being used (especially in biology applications)
 - PXE has density 0.985
 - DIN is in many commercial cocktails, miscible with water
 - **however** absorption is higher so less suited for large detectors
 - LAB (linear alkylbenzene)
 - developed by SNO+ as a scintillator (solvent) for nuclear/particle physics experiments; we're the first experiment to propose using LAB
 - high flash point, low toxicity
 - high light yield (about high as pseudocumene-based liquids)
 - excellent transparency
 - compatible with plastics like acrylic
 - cheap



Binary Mixtures and the Fluor



PPO or 2,5-diphenyloxazole

- the solvent molecule x is excited, relaxes to S_{1x}
- either by radiative or non-radiative transfer, the S_{1y} in the fluor molecule y is excited
 - radiative transfer is absorption and re-emission by the fluor
 - the emission spectrum of the solvent chosen to overlap the absorption spectrum of the fluor
 - S_{1y} is thus slightly lower in energy than S_{1x}
- this has to occur preferentially over non-radiative de-excitation of the solvent molecules S_{1x}
- efficient non-radiative transfer depends on the concentration of fluor in the binary mixture
 - occurs by solvent-solvent dipole resonance interaction (**Förster energy transfer**) or by thermal diffusion of excited solvent molecules in a liquid
 - occurs by **Förster energy transfer** between the same or adjacent polymer chains in the plastic

Determining the Efficiency of Energy Transfer

LAB-PPO energy transfer efficiency

Method 1 (fluorescence emission method): measure the fluorescence excitation and emission of PPO-LAB solution; excitation at 318 nm (LAB excitation), integrate emission range from 340-550 nm (PPO emission)

Method 2 (decay time method): excitation at 250 nm, emission 360 nm, measure the decay times of PPO-LAB solution, compare to the decay time of pure PPO

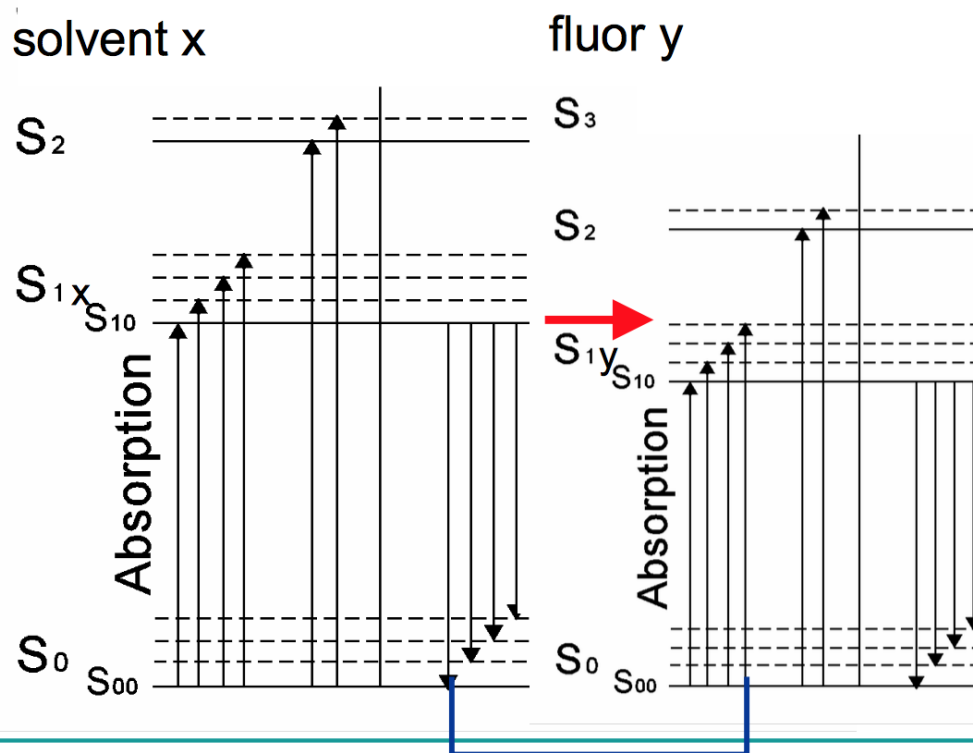
PPO [g/L]	Method 1	Method 2
32	n/a	90.7%
4	80%	80%
2	72%	73%
1	60%	61%
0.5	46%	55%
0.25	35%	47%
0.1	23%	n/a
0.02	5.6%	n/a

M.Chen

Energy Transfer: Solvent-Fluor

red: efficient non-radiative transfer is desired

aside: fluor-WLS energy transfer can be *radiative* because of the high quantum yield of the fluor and small WLS concentrations used to maximize transparency



blue: radiative transfer works; reduced by solvent fluorescent quantum yield

Additional Considerations

- quenching, **linearity** and pulse-shape discrimination (PSD) are related properties
 - PSD is the ability to distinguish between different particles exciting the scintillation by the shape of the light pulse
 - these properties can be considered, in general, for all types of scintillators
 - most of this discussion is directed at organic scintillators
-

Impurity Quenching

- presence of oxygen in an organic liquid at ppm levels can strongly quench fluorescence, especially in unitary scintillators
 - greatly reduces light yield
 - shortens lifetimes
 - oxygen and other impurities can compete with the fluor in energy transfer from the excited solvent
 - again reducing light yield
 - to combat oxygen quenching
 - increase the fluor concentration
 - deoxygenate the scintillator by bubbling with nitrogen or vacuum degassing the liquid
 - impurities in an inorganic crystal can quench the excitation, competing with activator centres
 - the obvious solution – grow pure crystals
-

Ionization Quenching

- ionization excites singlet and triplet π -electronic states
 - leading to fast fluorescence and slow delayed fluorescence from the triplet
 - results in multi-component scintillation decay time profile whereas UV excitation produces mainly single time constant fluorescence
 - high ionization density can quench the excited singlet π -electrons
 - the fast component is thus reduced for high dE/dx particles
 - picture it as overlapping excitations that interfere with each other
 - three important consequences:
 - non-linearity in energy response
 - heavy particles with higher dE/dx (e.g. α) produce less light for the same energy deposit, (by a factor of >10 for α in liquids)
 - the scintillation pulse shape (fast/slow components) is different for heavy particles, enabling pulse-shape discrimination
-

Ionization Quenching – Birks' Rule

- start with linearity at low ionization density

$$\frac{dL}{dx} = A \frac{dE}{dx}$$

- let the density of excited molecules be proportional to the ionization density

$$B \frac{dE}{dx}$$

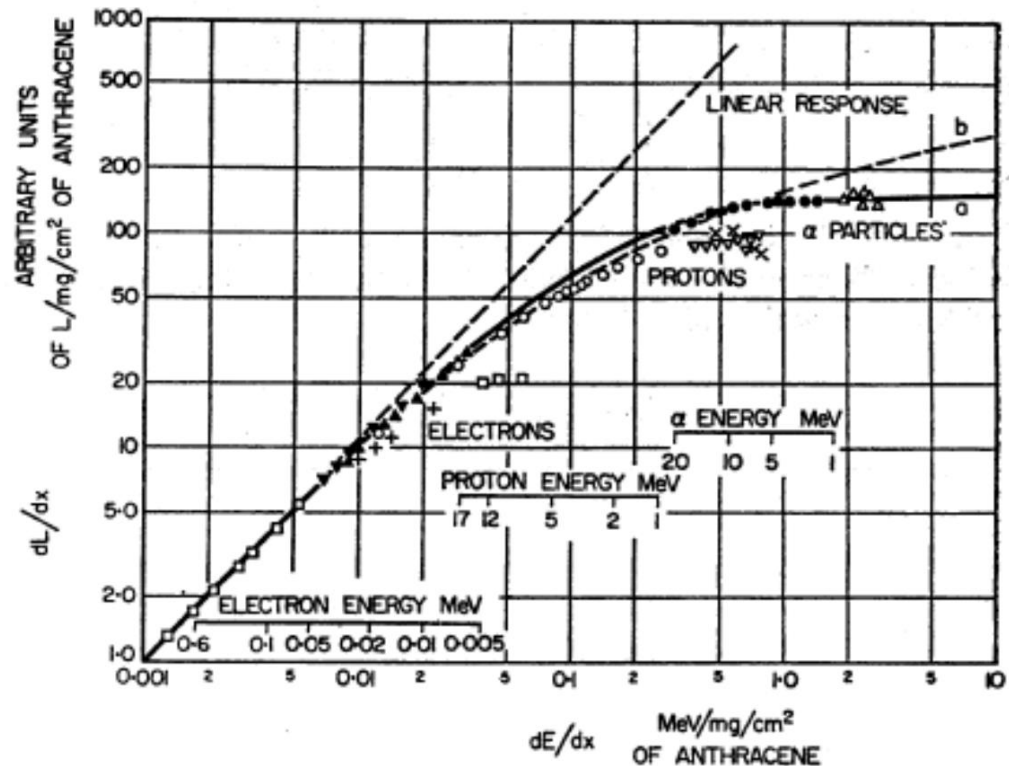
- let k be the fraction that is quenched

$$\frac{dL}{dx} = \frac{A \frac{dE}{dx}}{1 + kB \frac{dE}{dx}}$$

- for small dE/dx , approaches linearity

- for large dE/dx , approaches saturation

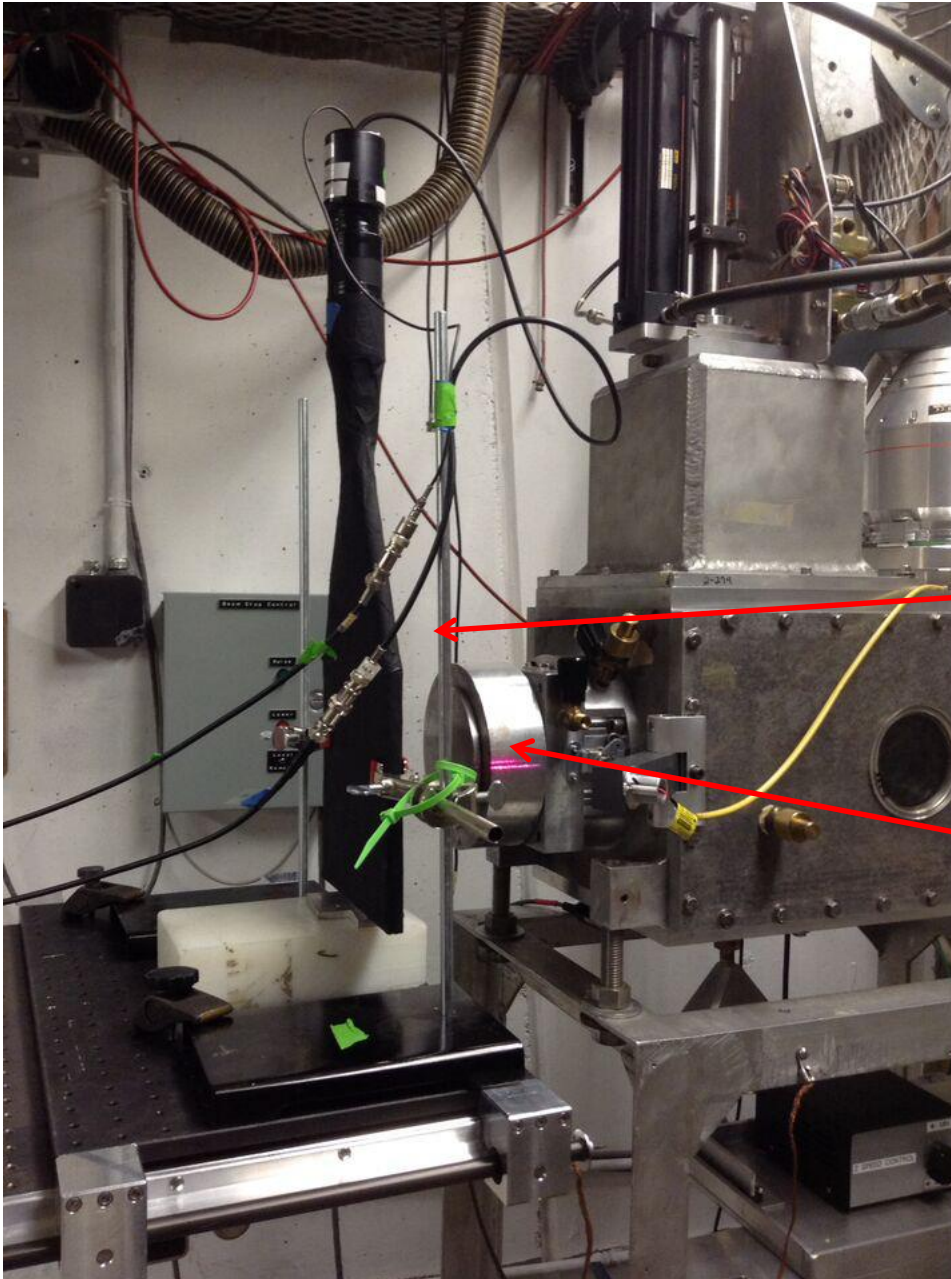
$$\frac{dL}{dx} = \frac{A}{kB}$$



from Birks

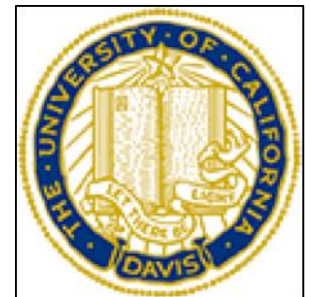
- for any given scintillator composition, must fit the quenching data to get kB
- A is the absolute scintillation efficiency

Measuring Birk's constant at CNL. June 2015



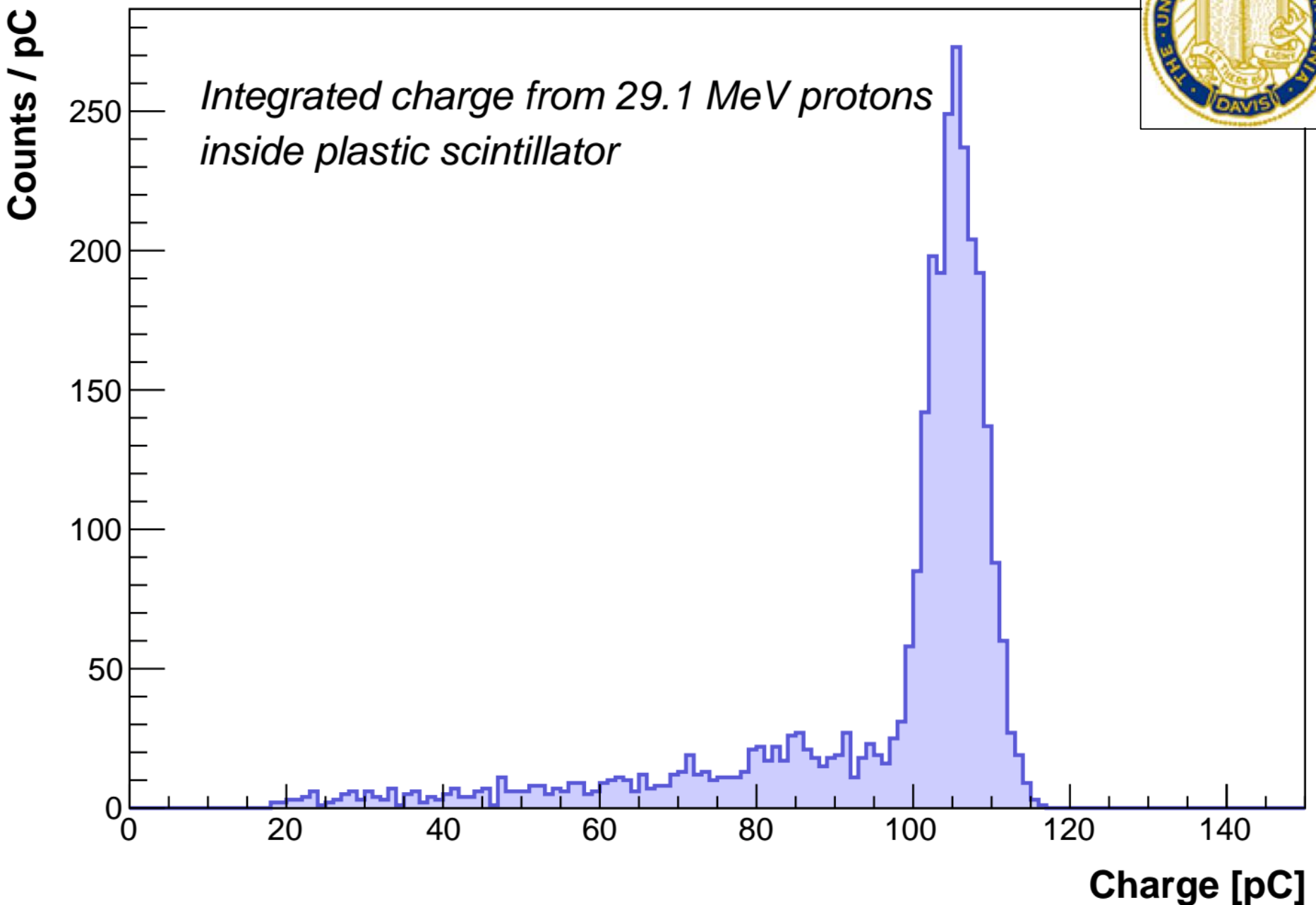
Plastic Scintillator

Proton beam pipe
window

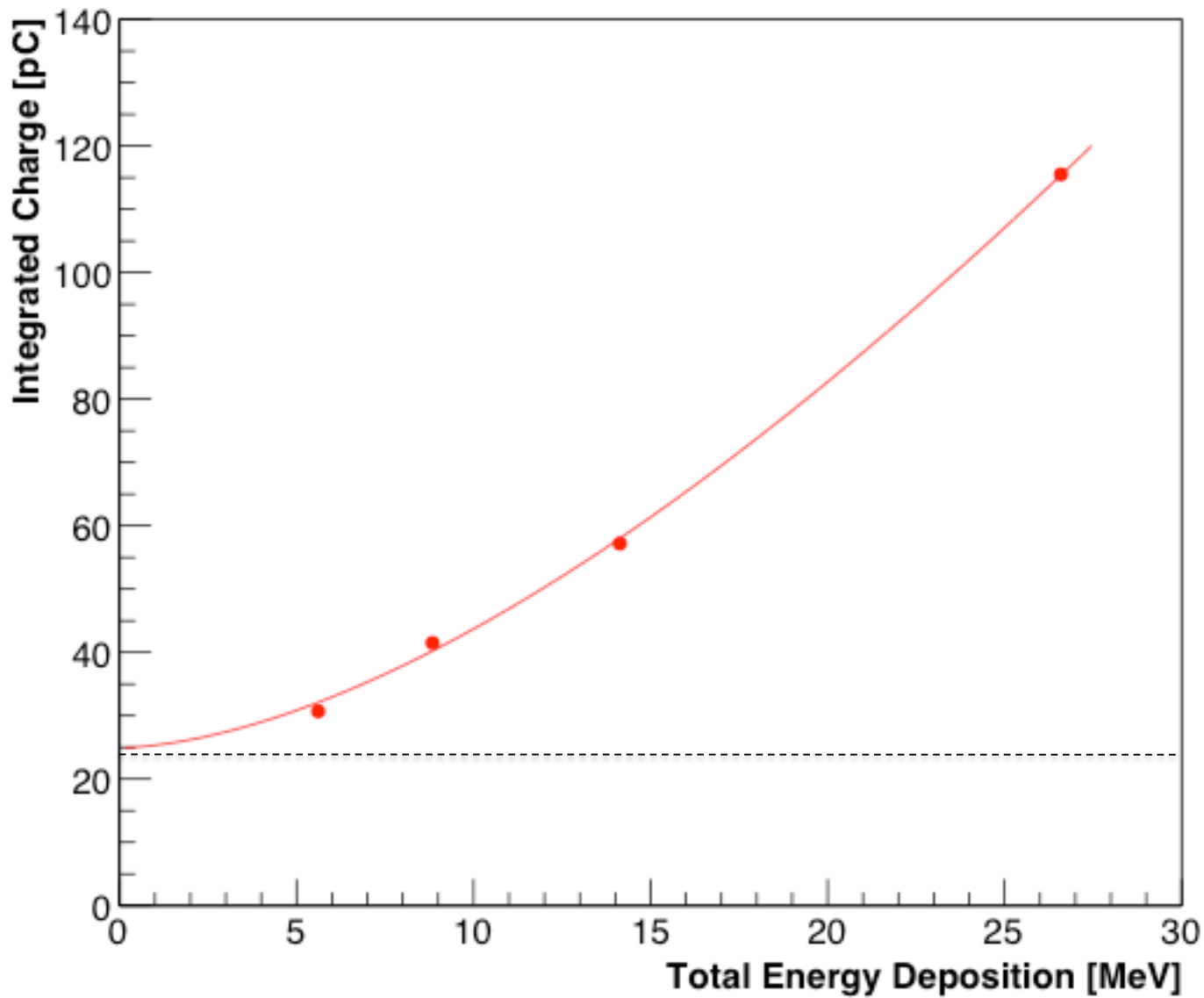




*Integrated charge from 29.1 MeV protons
inside plastic scintillator*

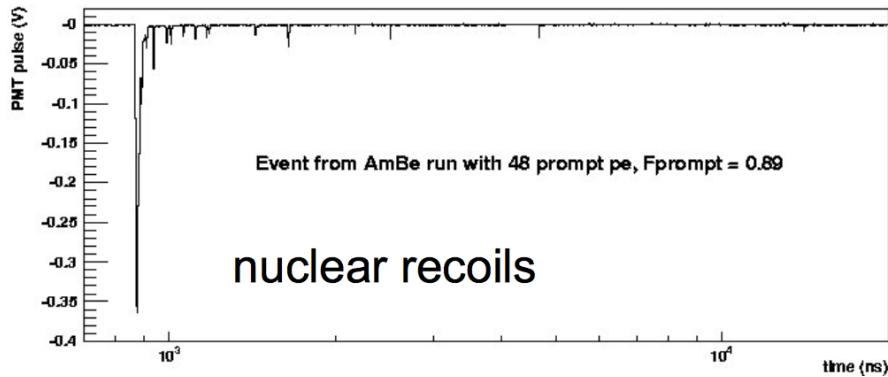
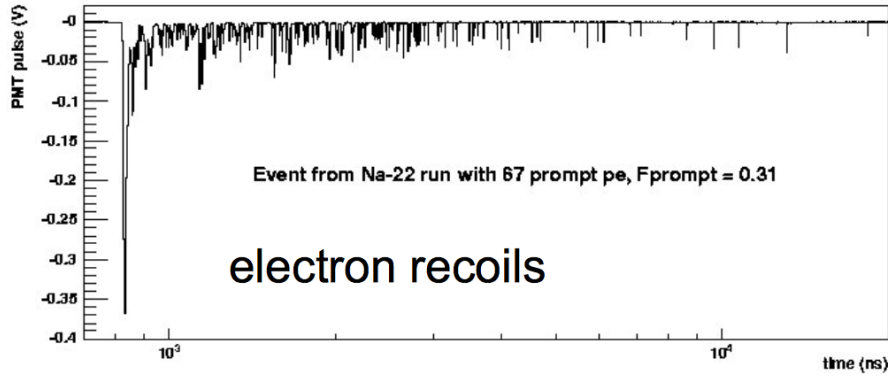


Plastic Scintillator Proton Quenching

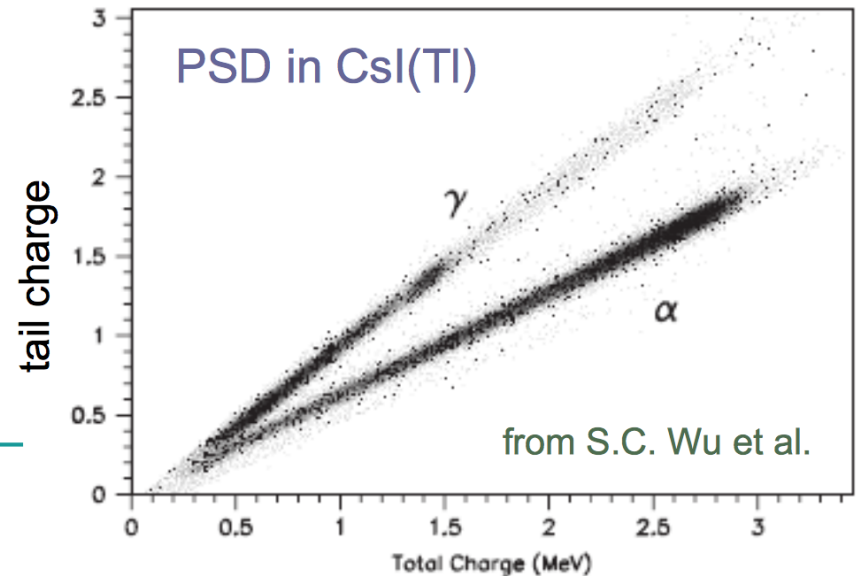
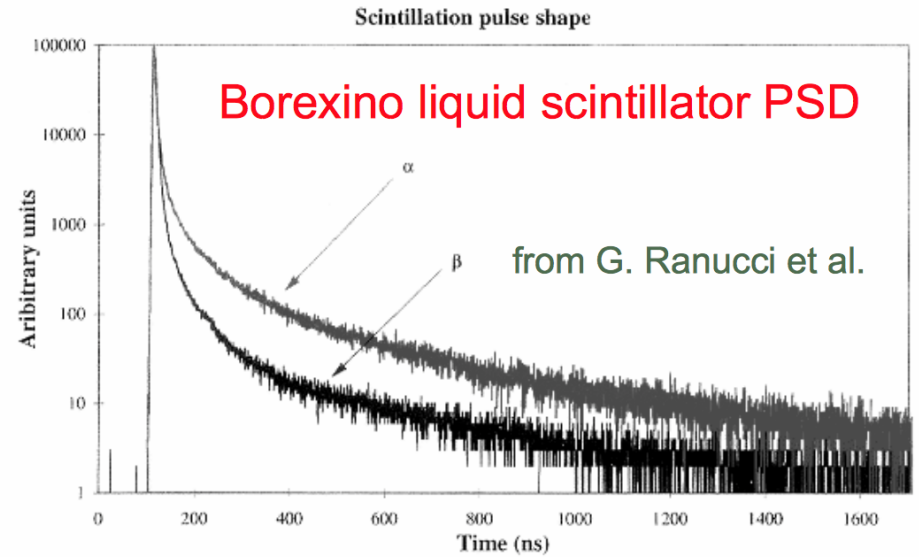


ADC
pedestal

PSD Scintillation Time Profiles



PSD = Pulse Shape Discrimination



Xenon Scintillation

Ionization in liquid:

- About one electron-ion pair /15 eV of deposited energy

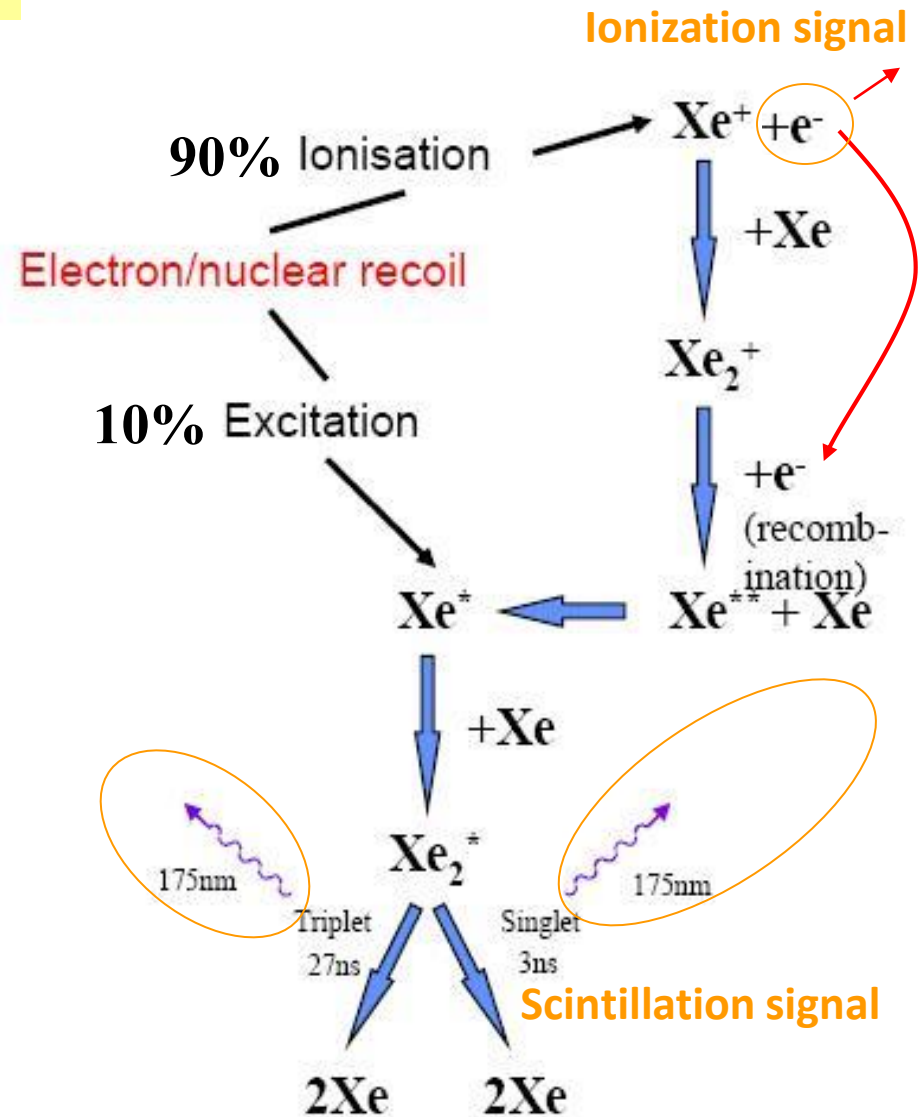
175 nm Scintillation:

- About one photon/20 eV of deposited energy

- Ionization/Scintillation ratio **changes with interaction type**

Ion pairs **more likely to recombine** in the dense tracks generated by nuclear interactions

- A **strong anti-correlation** between ionization and scintillation



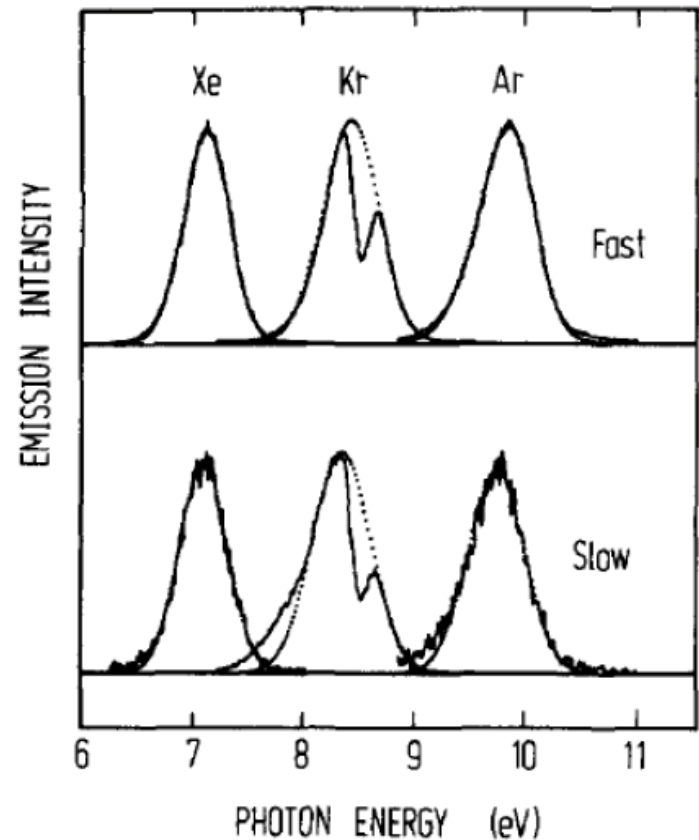
Scintillation in liquid noble elements

Xe is 175 nm

Ar is 128 nm

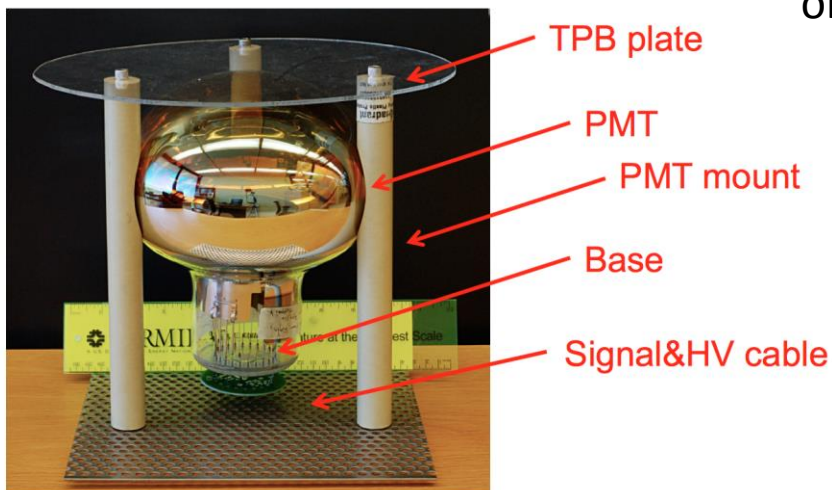
This means that light sensors need to be match to this wavelength OR one needs to shift the light to a longer wavelength.

E. Morikawa et al., J Chem Phys
vol 91 (1989) 1469

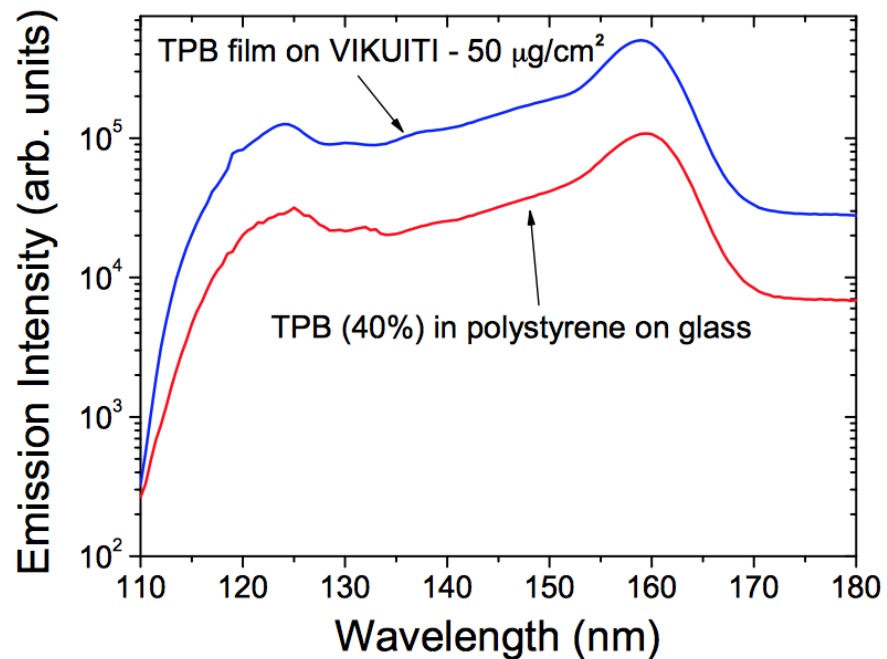
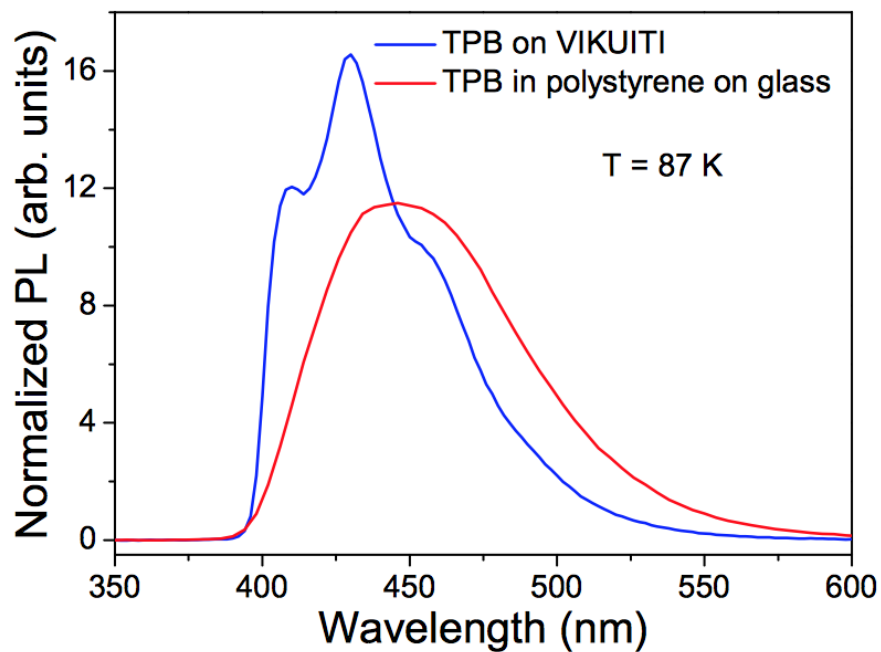
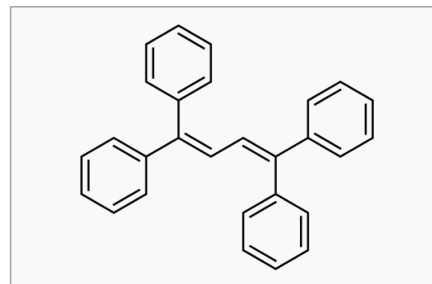


T.Katori

One needs to shift wavelength by coating a wavelength shifter (e.g. TPB) directly onto the PMT front face or on a plate of glass or acrylic in front of the PMT

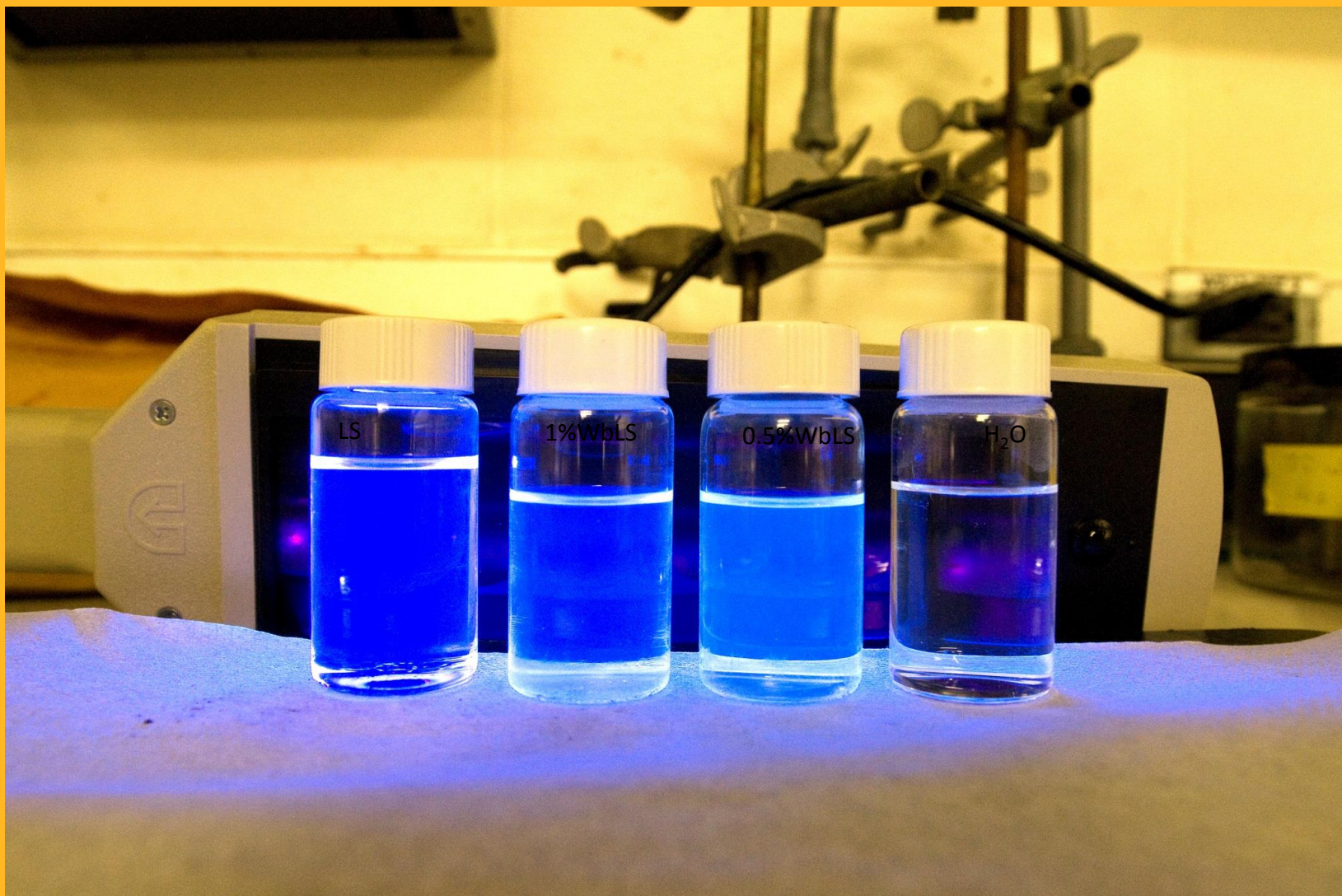


Tetraphenyl butadiene





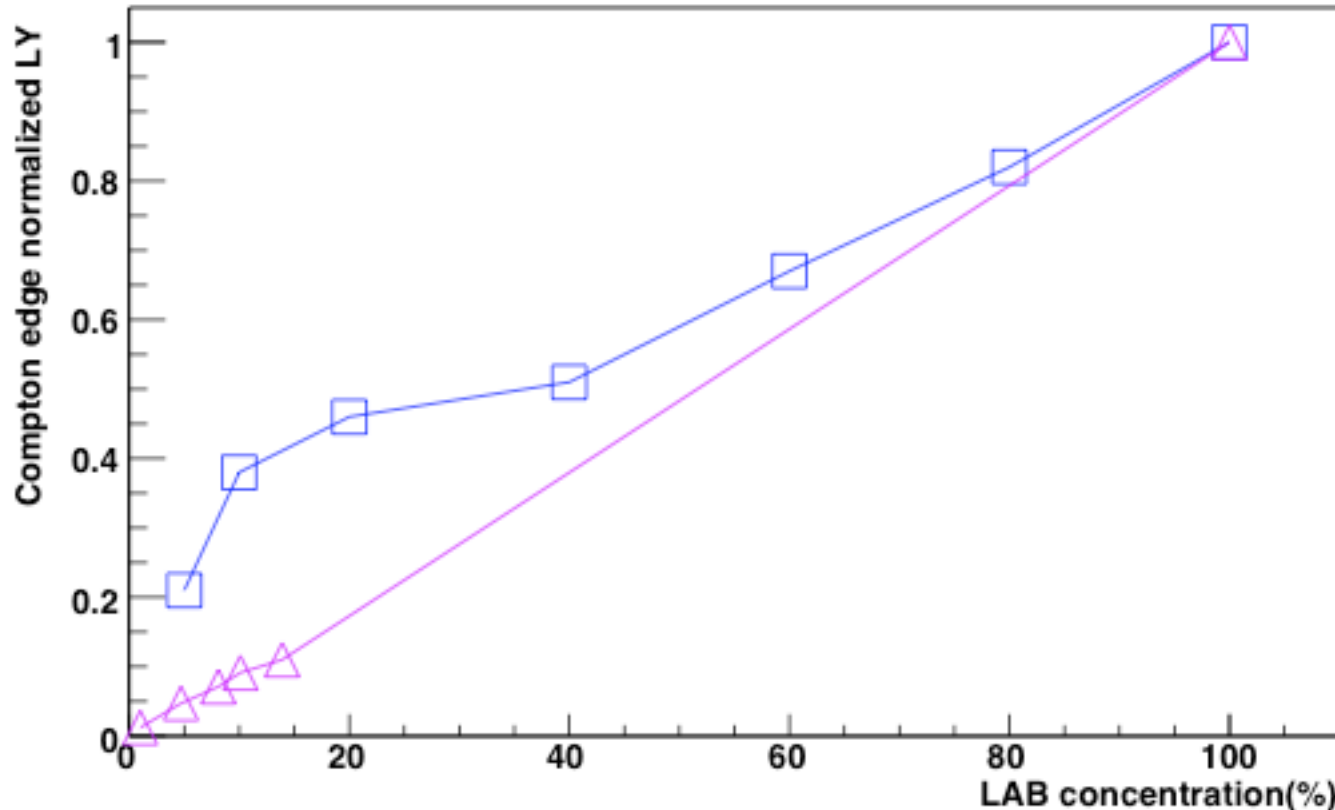
Water-base Liquid Scintillator (WbLS)



Different compositions of WbLS (1% ~ 100 optical photons/MeV), in comparison with pure LS and H₂O, under UV light ($\lambda=250\text{nm}$)

Dilution of WbLS in water allows for tuning light yield as desired to match the physics.

Compton edge normalized LY vs LAB concentration(%)



WbLS cocktail in water (violet) and cyclohexane (blue)

What can you do with this?

Advanced Scintillator Detector Concept (ASDC):

A Concept Paper on the Physics Potential of Water-Based Liquid Scintillator

J. R. Alonso,¹ N. Barros,² M. Bergevin,³ A. Bernstein,⁴ L. Bignell,⁵ E. Blucher,⁶ F. Calaprice,⁷
J. M. Conrad,¹ F. B. Descamps,⁸ M. V. Diwan,⁵ D. A. Dwyer,⁸ S. T. Dye,⁹ A. Elagin,⁶
P. Feng,¹⁰ C. Grant,³ S. Grullon,² S. Hans,⁵ D. E. Jaffe,⁵ S. H. Kettell,⁵ J. R. Klein,²
K. Lande,² J. G. Learned,¹¹ K. B. Luk,^{8,12} J. Maricic,¹¹ P. Marleau,¹⁰ A. Mastbaum,²
W. F. McDonough,¹³ L. Oberauer,¹⁴ G. D. Orebi Gann^a,^{8,12} R. Rosero,⁵ S. D. Rountree,¹⁵
M. C. Sanchez,¹⁶ M. H. Shaevitz,¹⁷ T. M. Shokair,¹⁸ M. B. Smy,¹⁹ M. Strait,⁶ R. Svoboda,³
N. Tolich,²⁰ M. R. Vagins,¹⁹ K. A. van Bibber,¹⁸ B. Viren,⁵ R. B. Vogelaar,¹⁵ M. J. Wetstein,⁶
L. Winslow,¹ B. Wonsak,²¹ E. T. Worcester,⁵ M. Wurm,²² M. Yeh,⁵ and C. Zhang⁵

¹Massachusetts Institute of Technology, Cambridge, MA 02139, USA

²Department of Physics and Astronomy, University of Pennsylvania, Philadelphia, PA 19104, USA

³Physics Department, University of California, Davis CA 95616, USA

⁴Lawrence Livermore National Laboratory, Livermore, CA 94550, USA

⁵Brookhaven National Laboratory, Upton, NY 11973, USA

⁶Enrico Fermi Institute, University of Chicago, Chicago, IL 60637, USA

⁷Department of Physics, Princeton University, NJ 08544, USA

⁸Lawrence Berkeley National Laboratory, Berkeley, CA 94720, USA

⁹Department of Natural Sciences, Hawaii Pacific University, Kaneohe, Hawaii 96744, USA

¹⁰Sandia National Laboratories, Livermore, CA 94550, USA

¹¹Department of Physics and Astronomy, University of Hawaii at Manoa, Honolulu, HI 96922 USA

¹²Department of Physics, University of California, Berkeley, CA 94720, USA

¹³Department of Geology, University of Maryland, College Park, MD 20742, USA

¹⁴TUM, Physik-Department, James-Frank-Str. 1, 85748 Garching, Germany

¹⁵Department of Physics, Virginia Polytechnic Institute and State University, Blacksburg, VA 24061, USA

¹⁶Department of Physics and Astronomy, Iowa State University, Ames, IA 50011, USA

¹⁷Department of Physics, Columbia University, New York, NY 10027, USA

¹⁸Department of Nuclear Engineering, University of California, Berkeley, CA 94720, USA

¹⁹Department of Physics and Astronomy, University of California, Irvine, CA 92697, USA

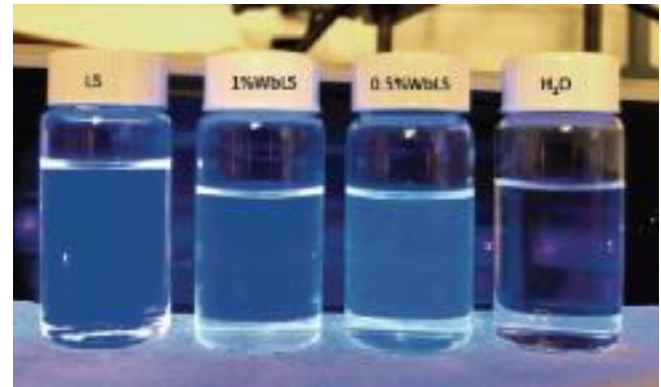
²⁰Center for Experimental Nuclear Physics and Astrophysics,

and Department of Physics, University of Washington, Seattle, WA 98195, USA

²¹Institute for Experimental Physics, University of Hamburg, Germany

²²Institute of Physics & EC PRISMA, Johannes Gutenberg-University Mainz, 55128 Mainz, Germany

Advanced Scintillator
Detector Concept
(ASDC) concept paper
posted on archive.



1% gives ~100 optical photons/MeV

4% WbLS gives approximately four
times the light yield of pure water

arXiv:1409.5864

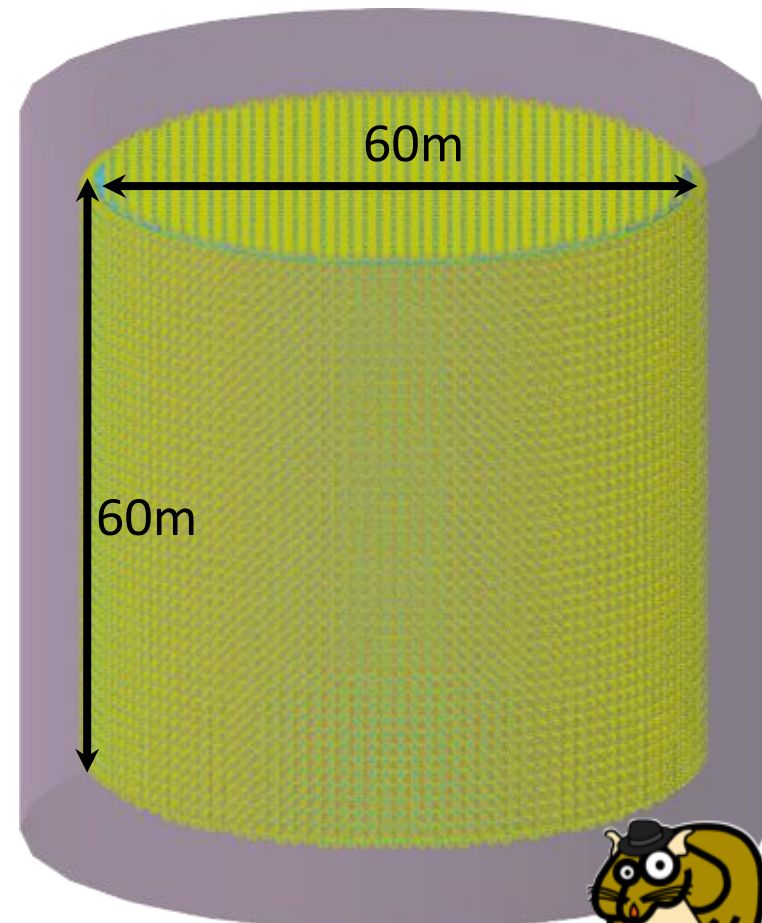
THEIA:

A realisation of the Advanced Scintillation Detector Concept (ASDC)

Concept paper - [arXiv:1409.5864](https://arxiv.org/abs/1409.5864)

- 50-100 kton WbLS target
- High coverage with ultra-fast, high efficiency photon sensors
- 4800 m.w.e. underground (Homestake).
- Is Kamioka a possibility?
- Comprehensive low-energy program: solar neutrinos, supernova, DSNB, proton decay, geo-neutrinos, DBD
- In the LBNF beam: long-baseline program complementary to proposed LAr detector

○ **Broad physics program!**

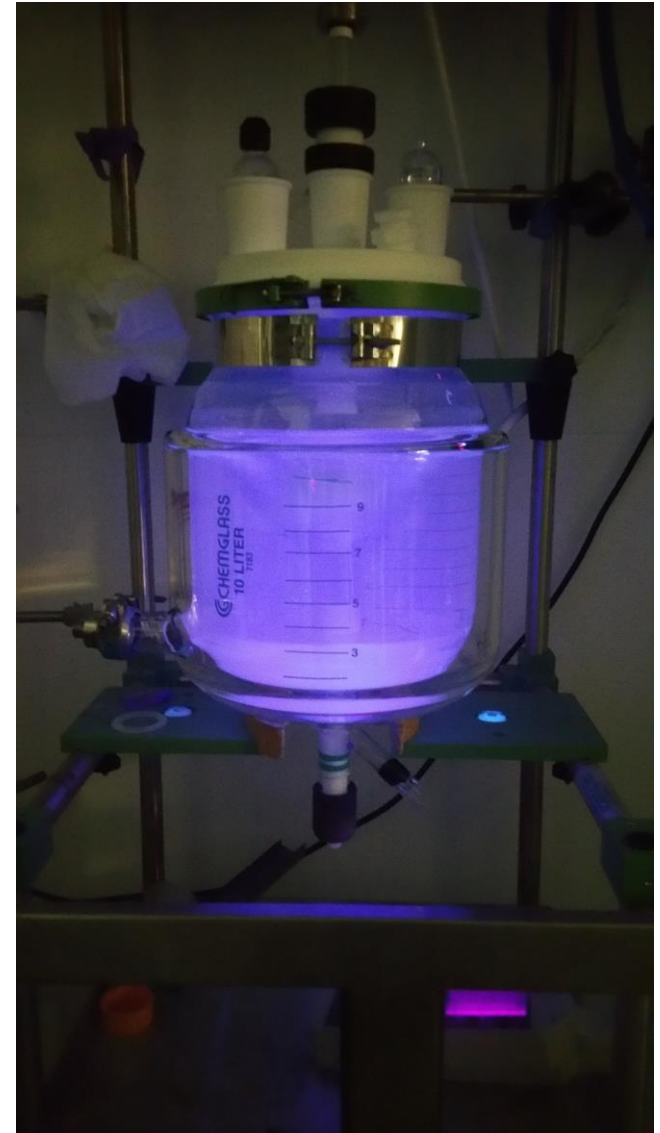


Detector image product of RAT-PAC

WbLS Development Status

- Light yield studies at BNL and soon at LBNL
- Stability and material compatibility studies at BNL (uncovered one problem so far – butyl rubber adhesive).
- Purification studies at UC Davis using NanoFiltration (NF) to separate organic components from water.
- Scaled up production: 10 liters produced in June with BNL 5 liter reactor. Used for NF and Material studies.
- 100 liter batch under production. Will be used for attenuation length studies (currently have on 1-meter arm). Will use UCI and/or LLNL facility.
- 1-ton BNL prototype approved and under construction.

BNL WbLS production





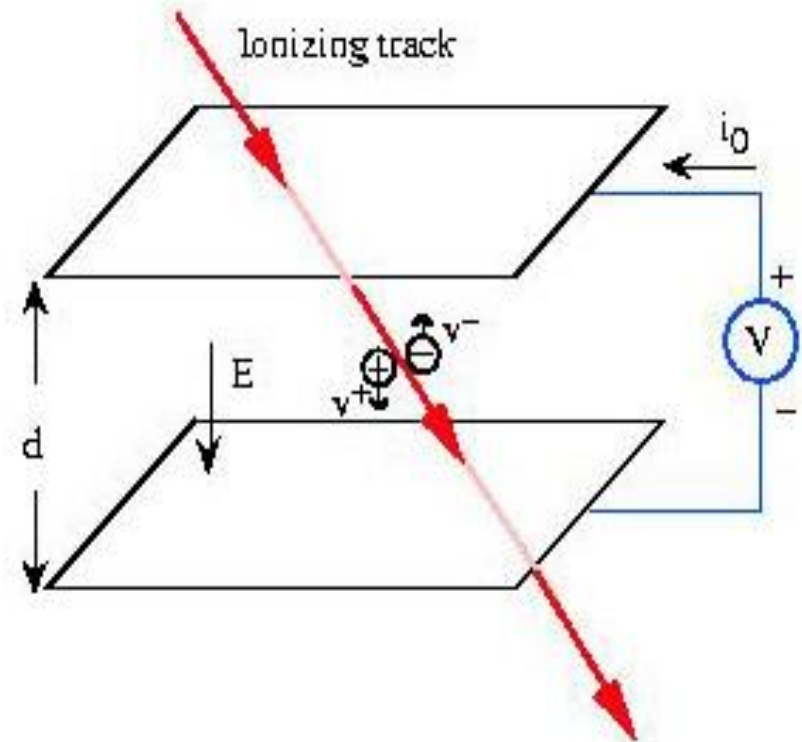
Liquid Argon Time Projection Chamber



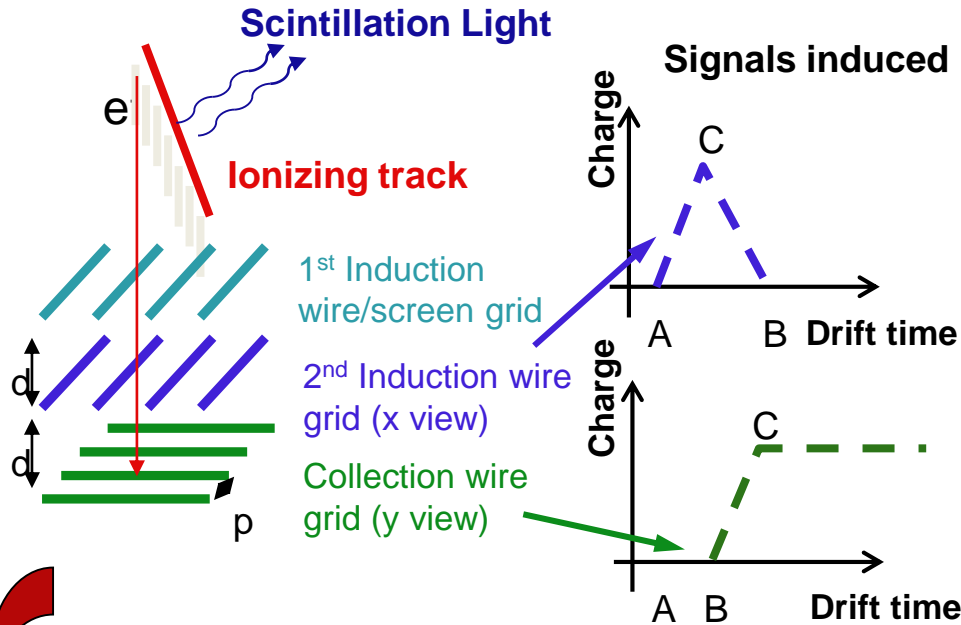
Single Phase LAr TPC

Operating principles

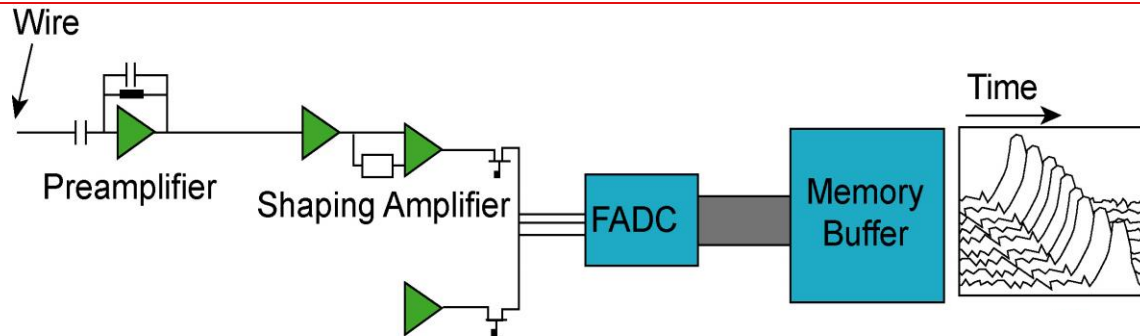
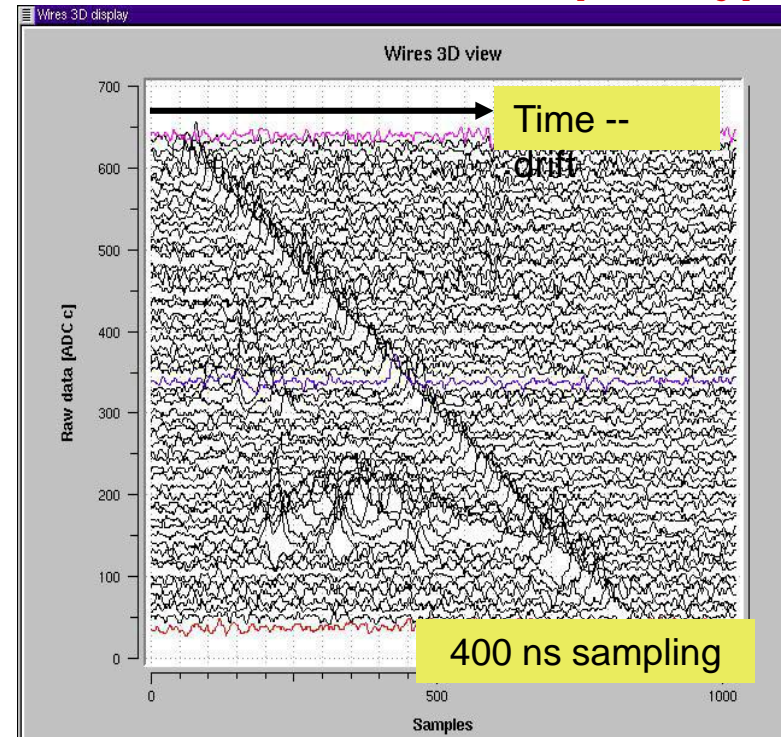
- **ionizing events** taking place in a volume of LAr (where a **uniform electric field** is applied) produce **electron-ion pairs**
- These charges **drift** along the field lines. The motion of the much **faster electrons** induces a **current on the anode**. The electrons can drift several metres if the LAr is high purified (electronegative impurities < 0.1 ppb O_2 equiv.)



Non-destructive multiple readout



Raw Data from a 10 m³ prototype

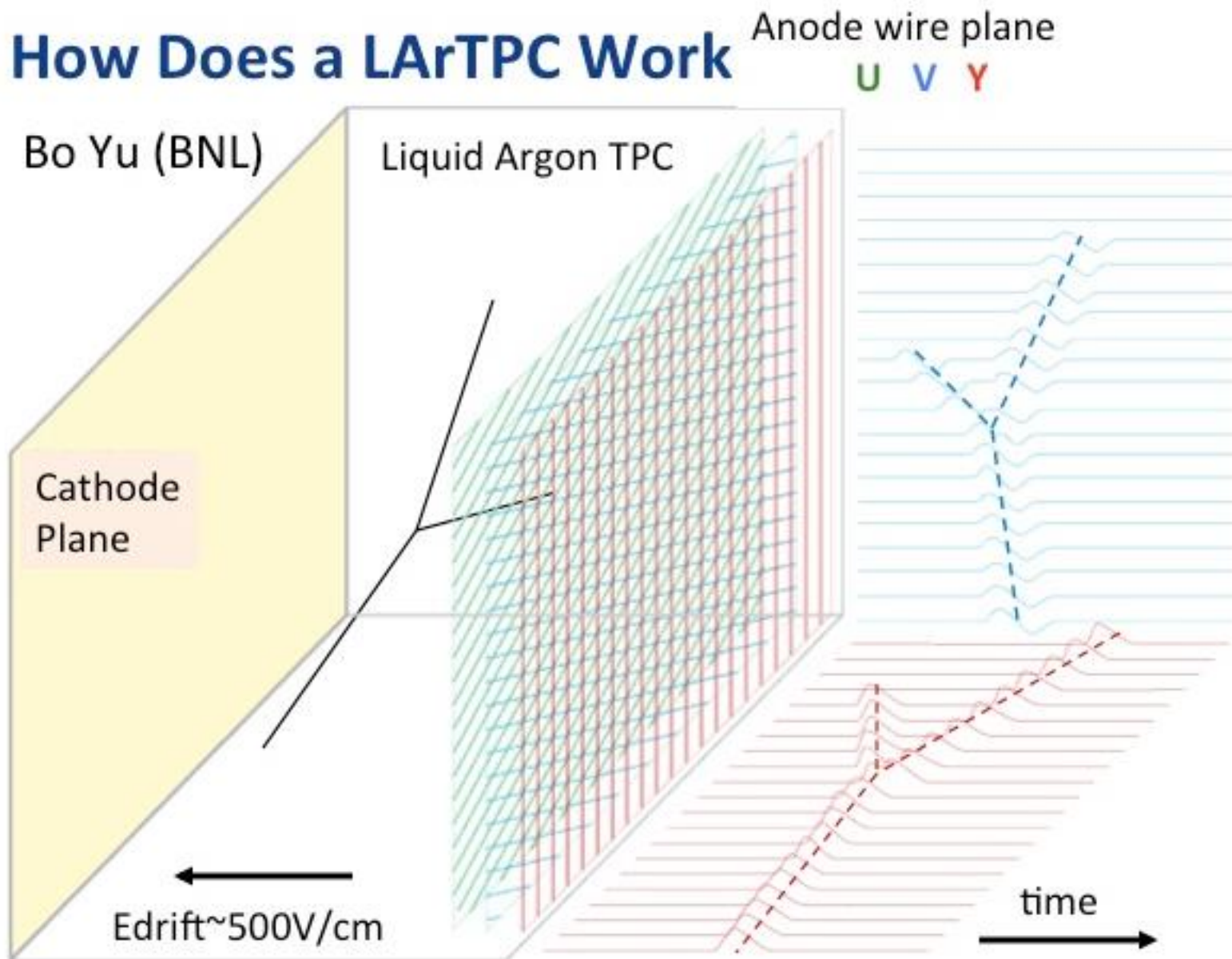


Continuous waveform recording

- **No charge multiplication** occurs in LAr
- LAr is also a **very good scintillator** scintillation light ($\lambda = 128 \text{ nm}$) provides a prompt signal to be used for triggering purposes and for absolute event time measurement

- High electron mobility ($\sim 500 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$)
- Possibility of extreme purification ($< 0.1 \text{ ppb O}_2$)
- \hookrightarrow Long electron life time ($> \text{ms}$) and drift paths ($> \text{m}$)
- High electron-ion pairs yield ($\sim 10000 \text{ e}^-$ for 2 mm of a m.i.p. track)
density: 1.4 g/cm^3 $dE/dx: 2 \text{ MeV/c}$
- Available in large quantities ($\text{GAr} \sim 0.9\% \text{ of air}$)

How Does a LArTPC Work



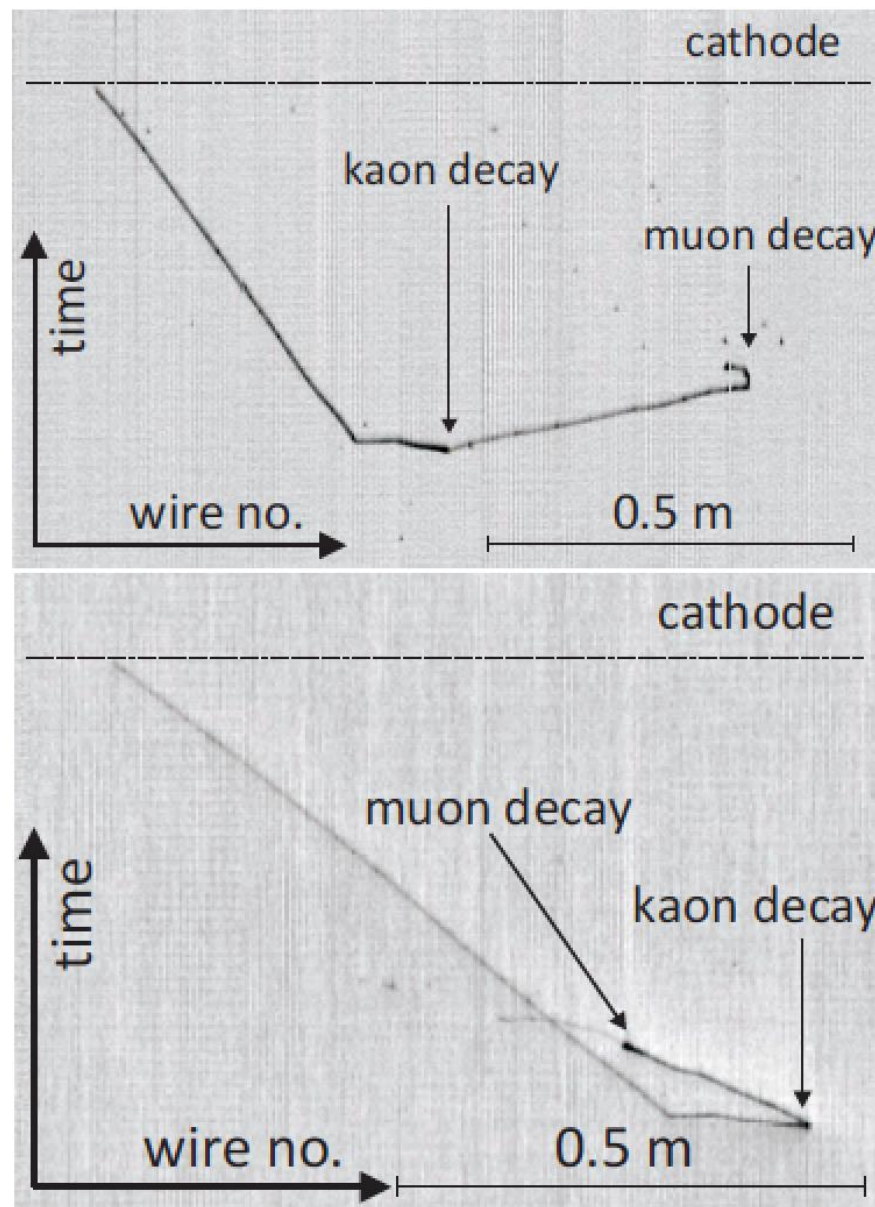
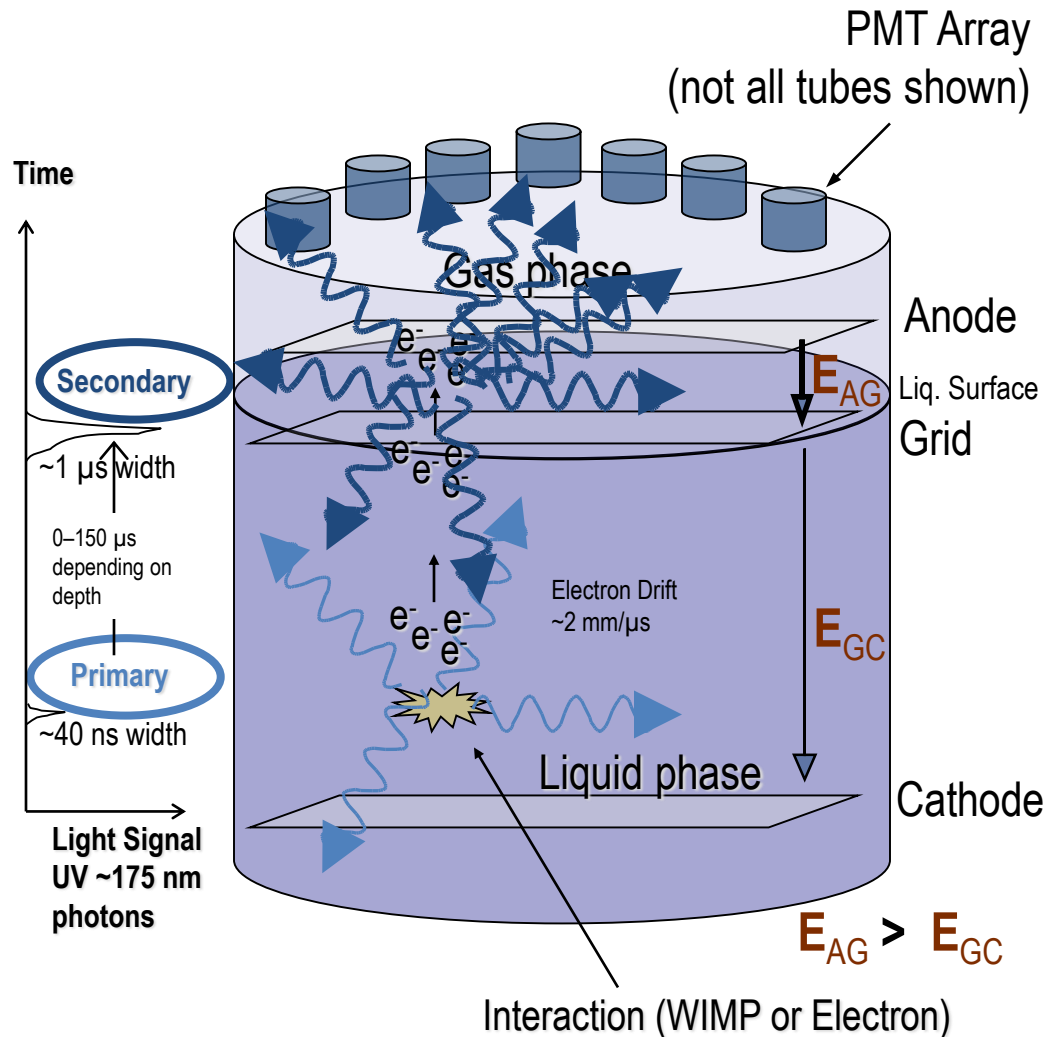
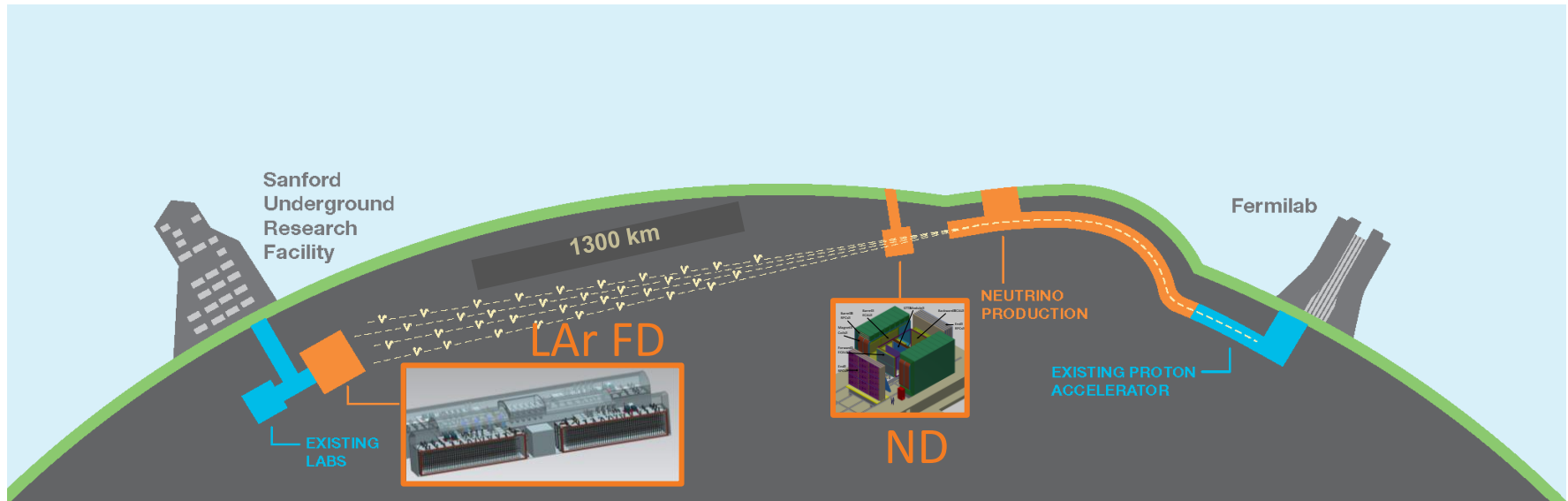


Figure 4.3: Event display for a decaying kaon candidate $K \rightarrow \mu\nu_\mu$ $\mu \rightarrow e\nu_e\nu_\mu$ in the ICARUS T600 detector observed in the CNGS data (K : 90 cm, 325 MeV; μ : 54 cm, 147 MeV; e : 13 cm, 27 MeV). The top figure shows the signal on the collection plane, and the bottom figure shows the signal on the second induction plane [82].

Dual Phase Technique

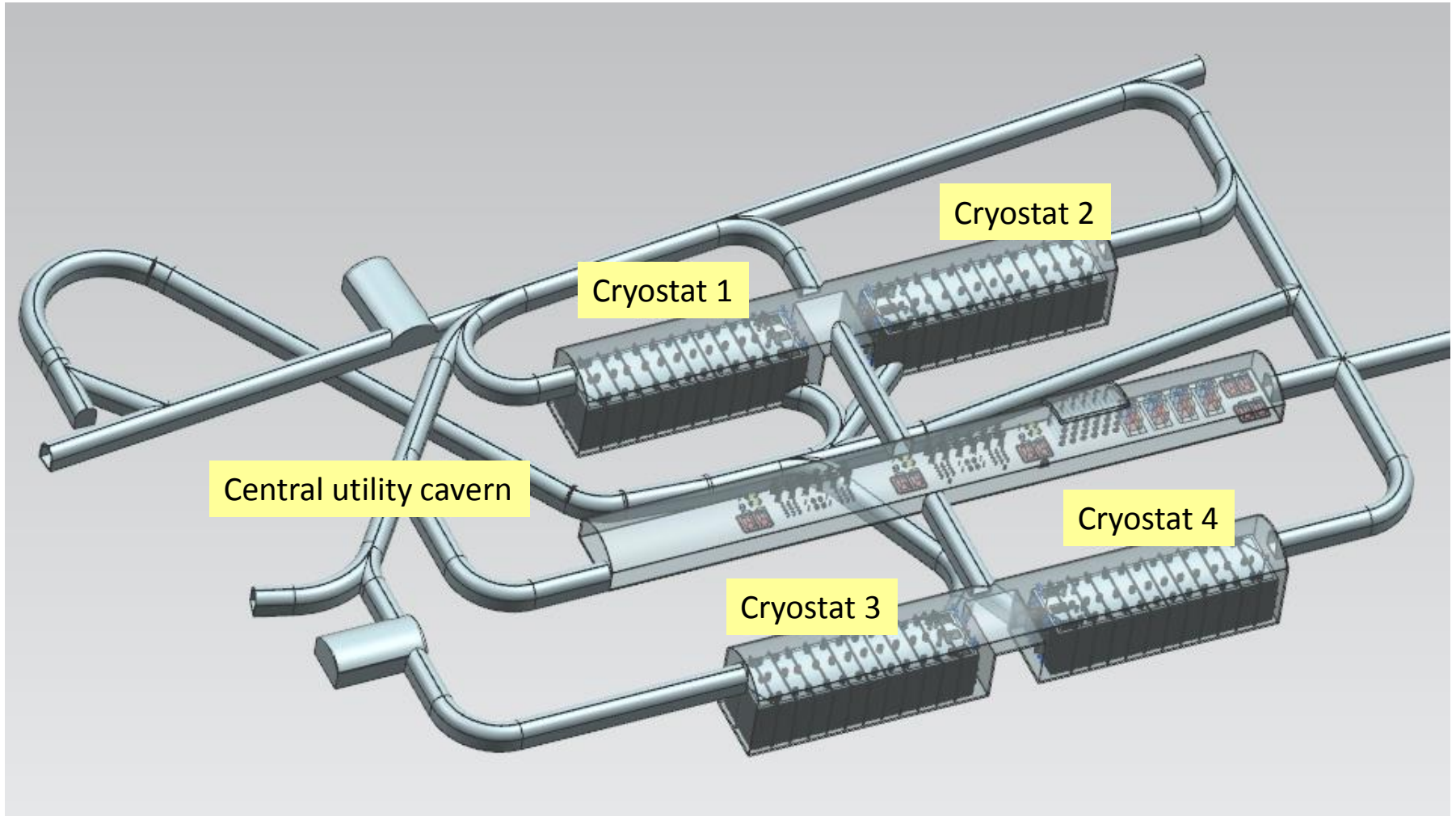


DUNE experiment overview



- 1.2 MW wide-band neutrino beam from FNAL, upgradeable to 2.4 MW
- Highly-capable near detector
- LAr 40-kton fiducial mass far detector
 - @ Sanford Underground Research Facility in SD
 - 1300 km baseline
 - 4850 ft (2300 mwe) depth
 - Four 10 kt modules, installation starting 2021

LBNF far detector facilities for DUNE

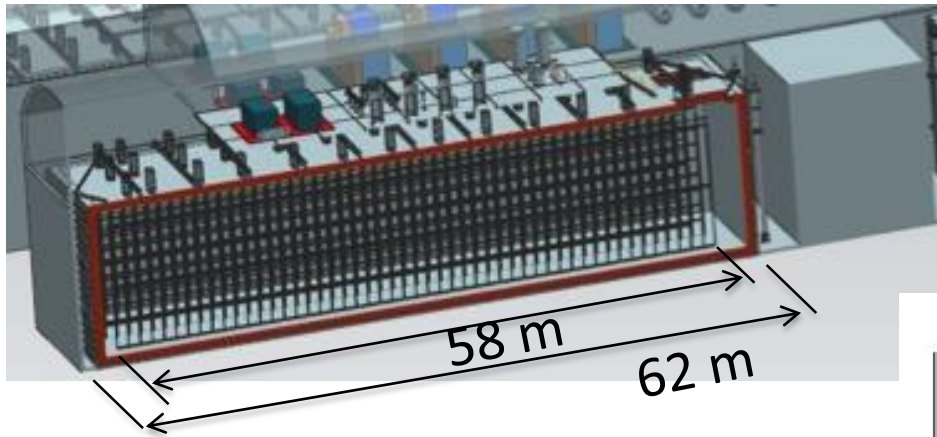


Cryostats: (CERN-FNAL design team) 17.1 kt LAr each

Free-standing steel-supported membrane cryostats

Central utility cavern: cryogenics support equipment

Single Phase Detector



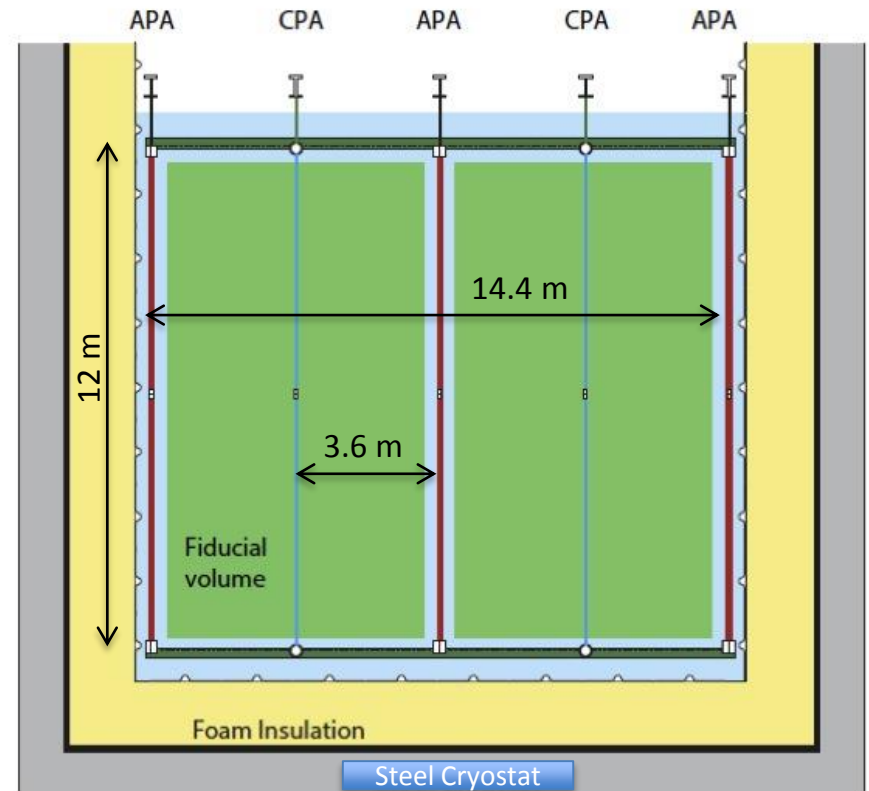
17.1-kt total, 13.8-kt active,
11.6-kt fiducial mass

3 Anode Plane Assemblies (APA)
w/ cold electronics

Cathode planes (CPA) at 180 kV
3.6 m max drift length

Photon detectors for fast
event timing (non-beam physics)

Reference design
(for first module):
horizontal-drift single-phase
time-projection chamber



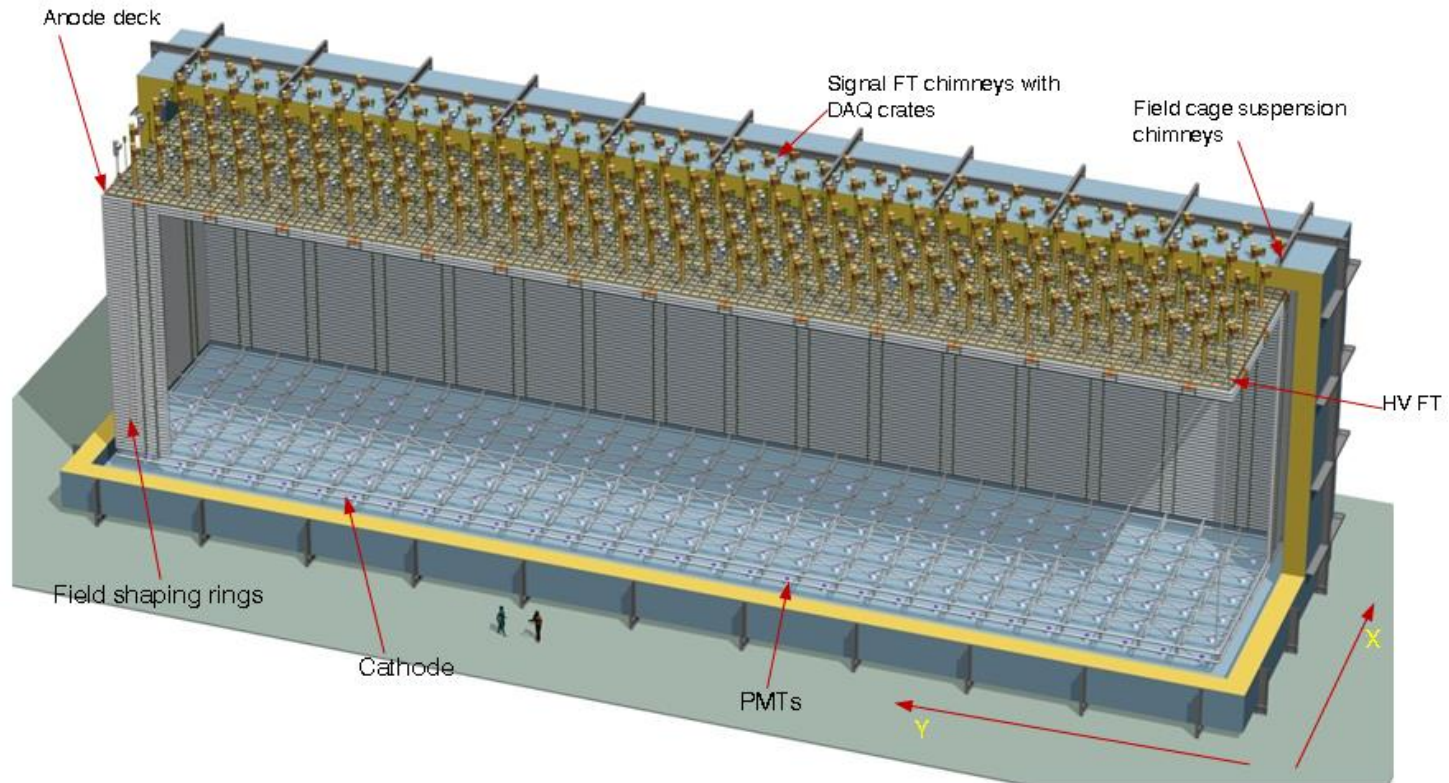
Dual Phase Detector

Dual-phase TPC is alternative design:

vertical drift w/ multiplication and readout at liquid-gas interface

Significant R&D by LBNO collaboration

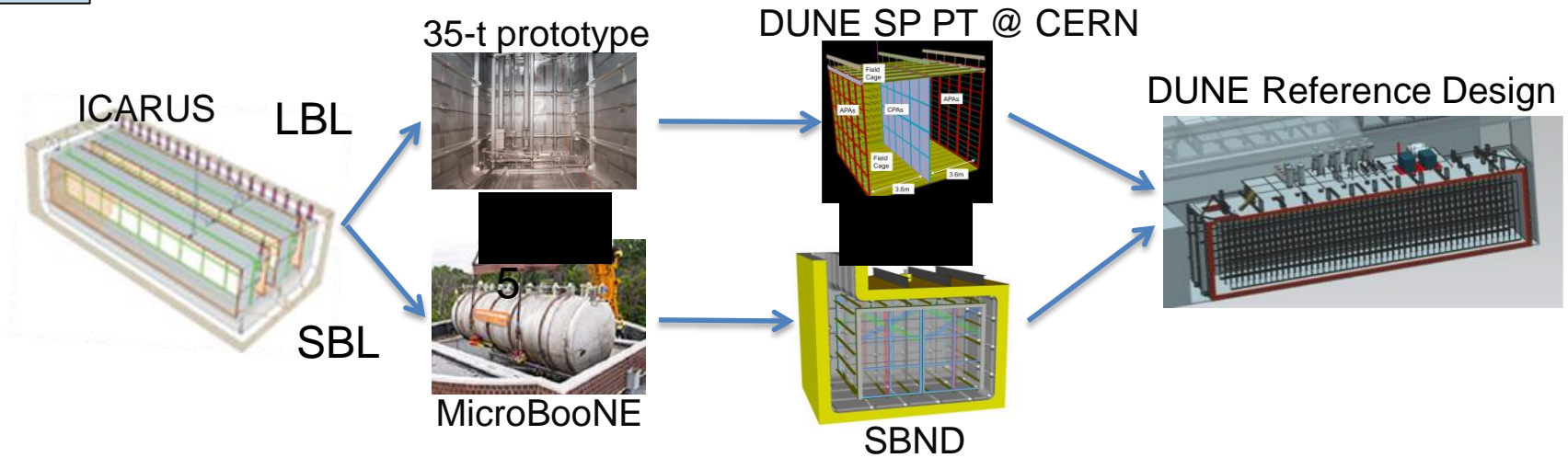
Could be implemented for module(s) 2-4



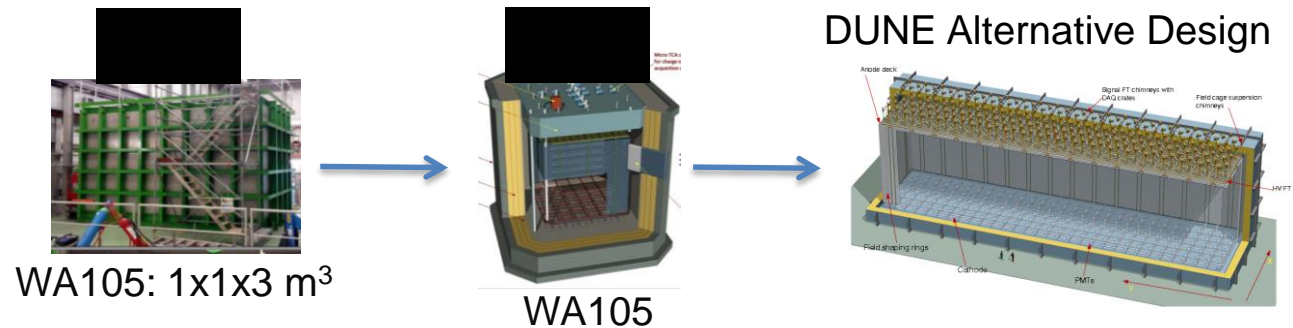
More from G.Feldman later in the school...

CERN neutrino platform + FNAL prototyping + experience from FNAL SBN program

Single
phase



Dual
phase



Liquid Argon TPC

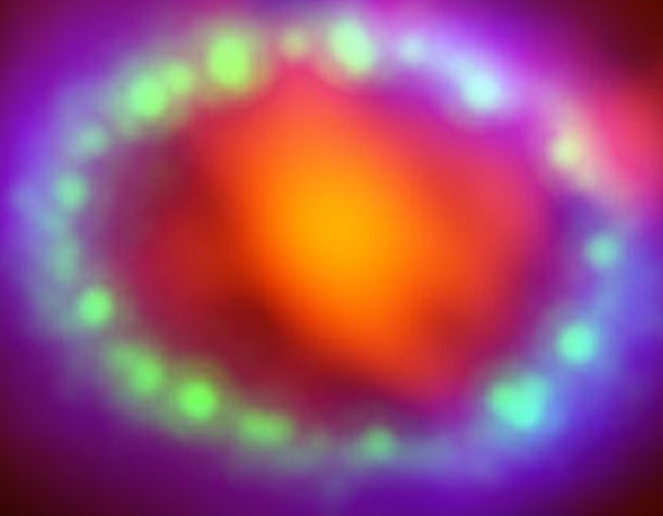
Advantages

- tracks recorded at sub 1mm level. Easy to distinguish multiple and/or complex events
- recoil hadrons visible even below Cherenkov threshold
- Argon is a scintillator
- Argon is relatively cheap (comparable to scintillator)
- **Significant ν_e cross section at low energy**

Disadvantages

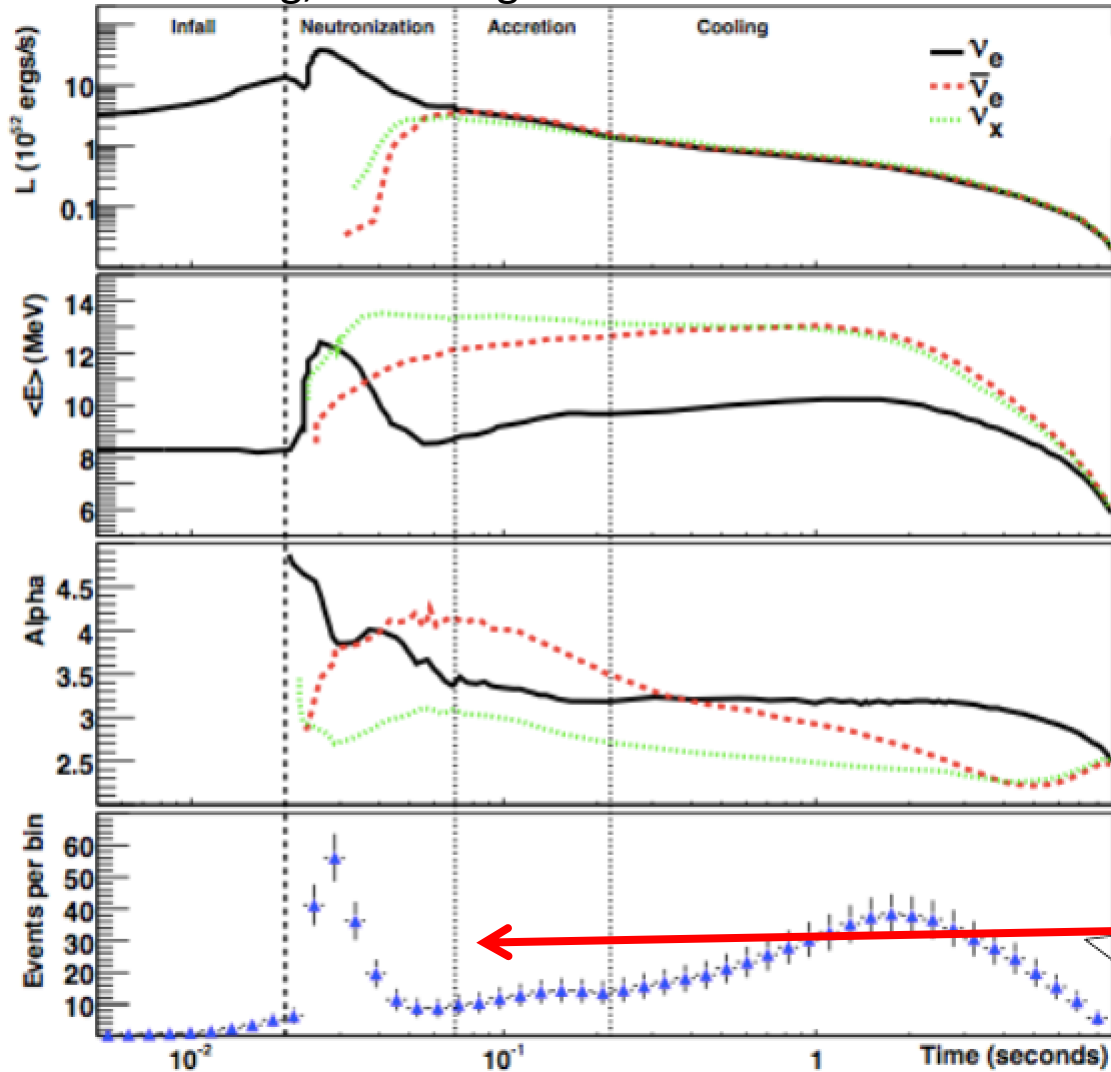
- Drift is slow compared to Cherenkov or scintillation
- Need scintillation light collection for a "t0"
- Complex nucleus compared to oxygen or carbon
- No free protons
- No neutron capture tagging

The Long(er) Baseline Neutrino Experiment



Detection of a SN neutronization burst requires ν_e sensitivity.

K.Scholberg, "Garching" model used



Burst is only 20 ms long and is essentially all ν_e

Mean energy of events is low, 10-12 MeV

IMB/Kamiokande detected higher energy cooling neutrinos, not neutrinos from the neutronization process

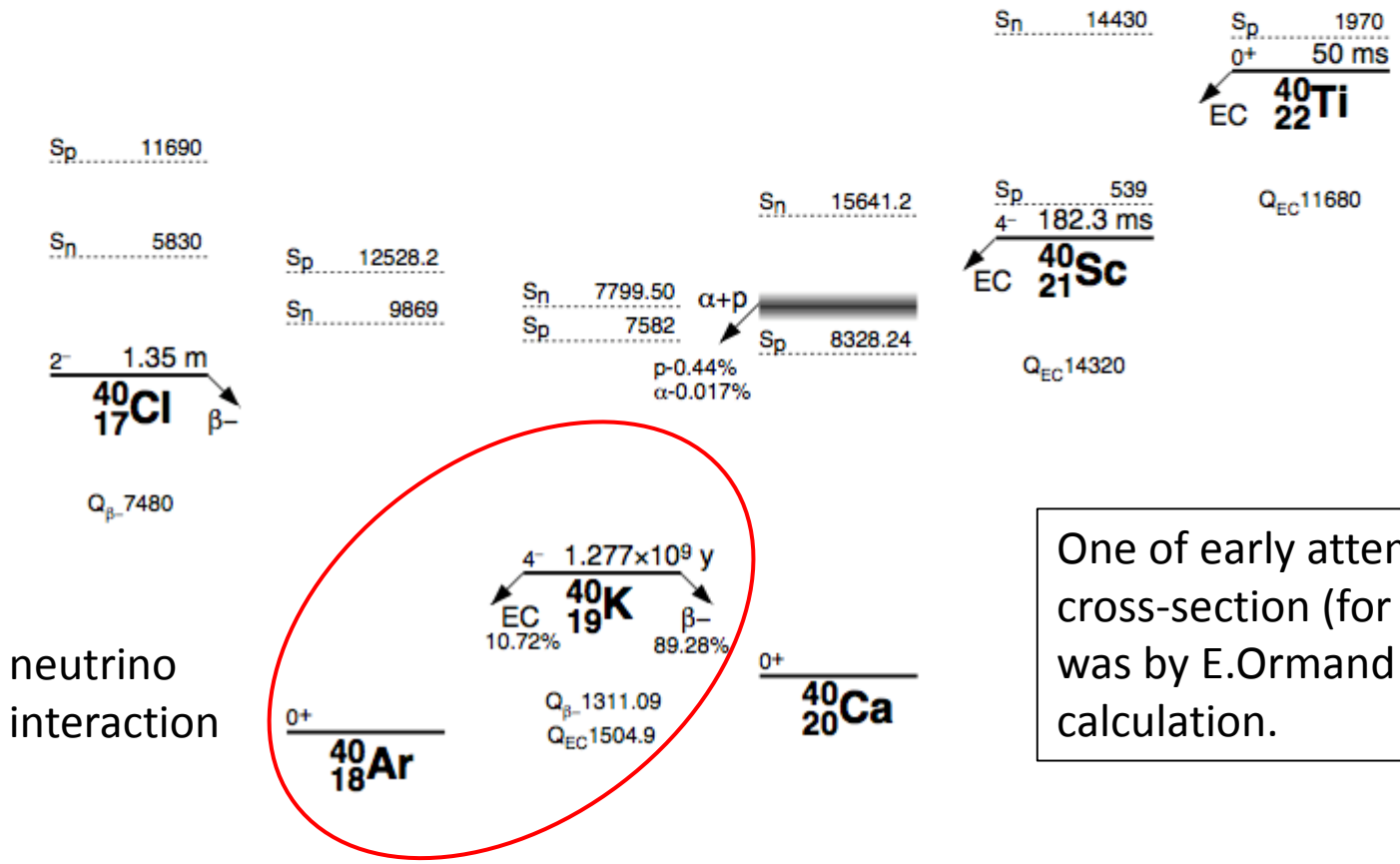
Potential for ν_e detection in liquid argon by LBNE

Do we really know how to detect a supernova
with a liquid argon TPC?

cross-sections?
triggering?
timing?



$\nu_e + {}^{40}\text{Ar} \rightarrow {}^{40}\text{K}_{\text{g.s.}} + e^-$ has a low threshold, but rarely occurs due to $0^+ \rightarrow 4^-$ (3rd forbidden) transition. Most all of the cross section is into excited states of ${}^{40}\text{K}$.



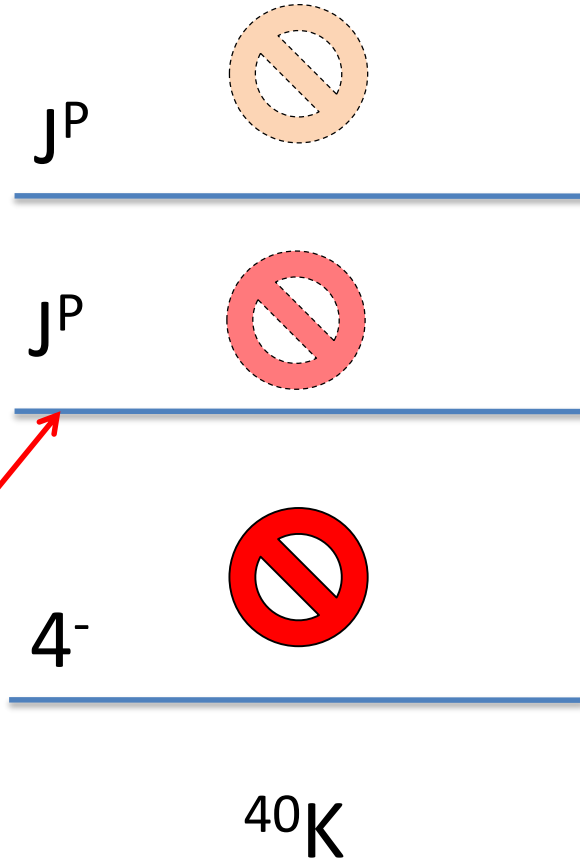
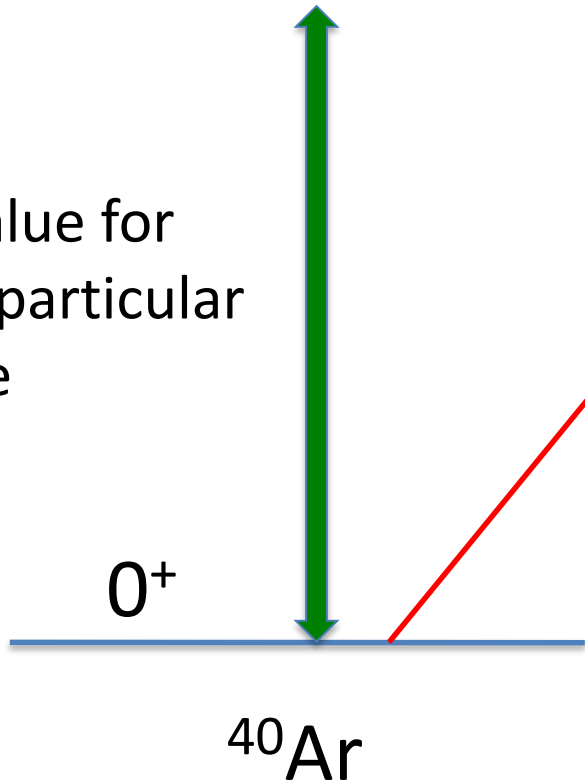
neutrino interaction

One of early attempts to calculate cross-section (for solar neutrinos) was by E.Ormand using shell model calculation.

~30 relevant excited states of various filling "strength".



Q-value for this particular state

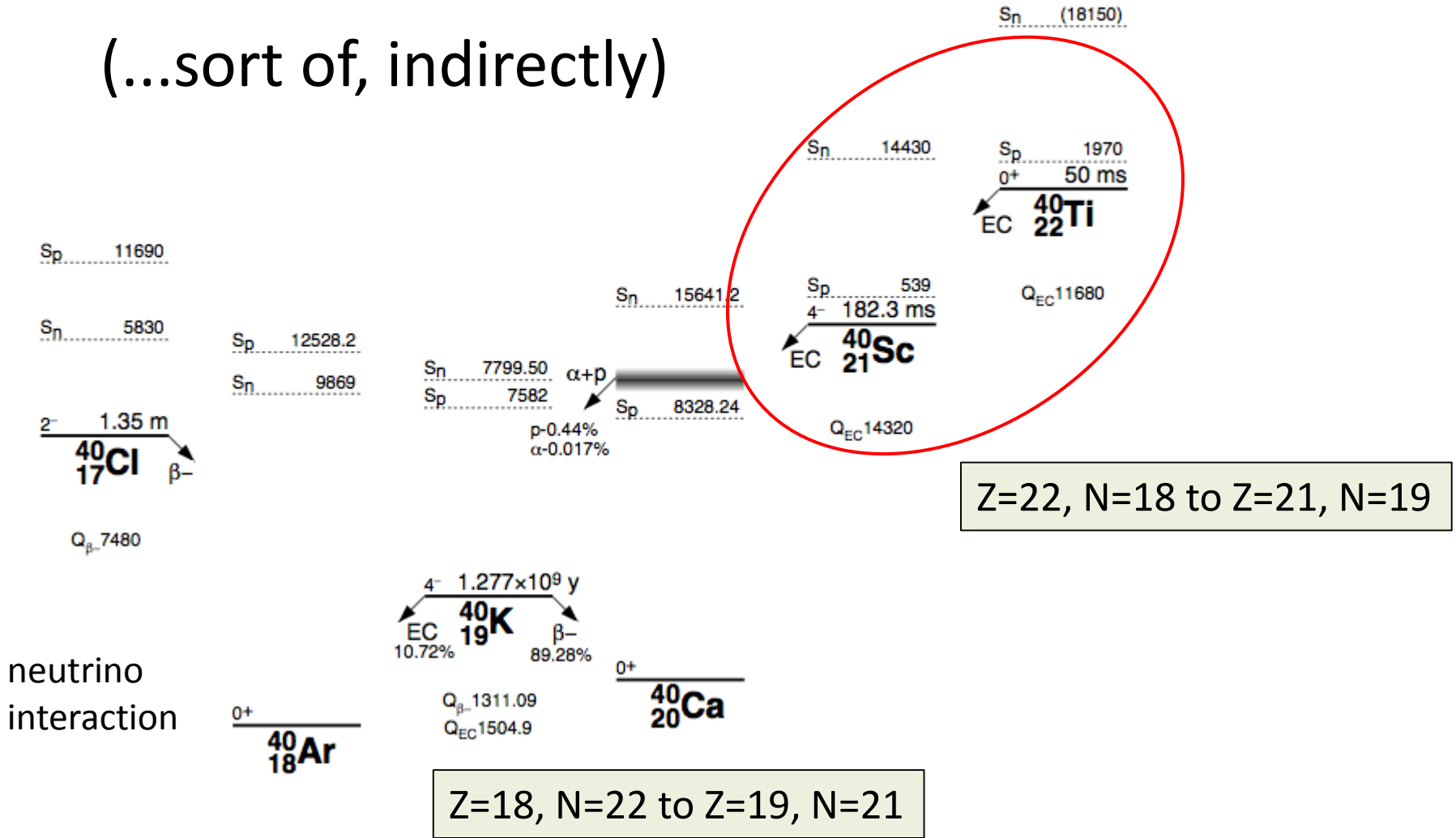


$$E_{e^-} = E_\nu - Q - K_{\text{recoil}}$$

Can we just measure the strengths?

A=40
NP A521, 1(1990)

(...sort of, indirectly)



Z=22, N=18 to Z=21, N=19

Z=18, N=22 to Z=19, N=21

neutrino
interaction

Neutrino absorption efficiency of an ^{40}Ar detector from the β decay of ^{40}Ti

M. Bhattacharya, A. García, and N. I. Kaloskamis*
University of Notre Dame, Notre Dame, Indiana 46556

E. G. Adelberger and H. E. Swanson
University of Washington, Seattle, Washington 98195

R. Anne, M. Lewitowicz, M. G. Saint-Laurent, and W. Trinder
GANIL, BP 5027, F-14021 Caen Cedex, France

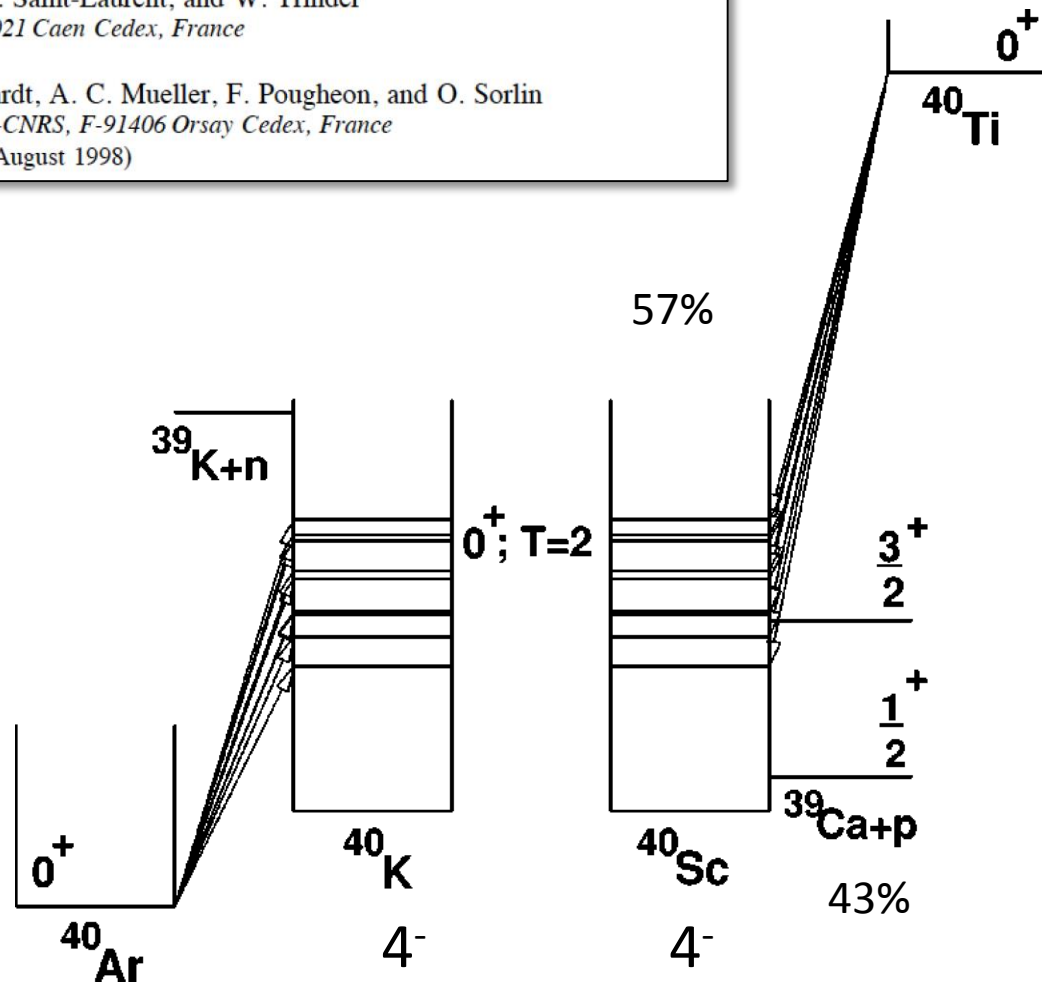
C. Donzaud, D. Guillemaud-Mueller, S. Leenhardt, A. C. Mueller, F. Pougheon, and O. Sorlin
Institut de Physique Nucléaire, IN2P3-CNRS, F-91406 Orsay Cedex, France

(Received 4 August 1998)

Make ^{40}Ti via heavy ions
 on a thin, heavy target (e.g.
 Cr on Ni). Embed ions in
 silicon detector.

Use TOF and dE/dx to separate
 ^{40}Ti out from ion "soup"

Observe beta decay to ^{40}Sc
 excited states, which decay via
 delayed proton emission



Transition Strengths and ν_e Cross Sections

TABLE VII. β decay of ^{40}Ti and the neutrino-capture cross section of ^{40}Ar . (Cross sections for ^8B and supernova ν 's are computed with a 5 MeV threshold on the total energy of the outgoing electron.) The isobaric correspondences suggested in this table are based solely on excitation energies and J^π assignments.

$E_x(\text{keV})(^{40}\text{Sc})$	$B(F)+B(GT)$	$E_x(\text{keV})(^{40}\text{K})$	$J^\pi(^{40}\text{K})^a$	$\sigma_B(10^{-43}\text{cm}^2)^b$	$\sigma_{\text{SN}}(10^{-43}\text{cm}^2)^c$
2281 ± 8	0.90 ± 0.04	2289.88 ± 0.03	1 ⁺	3.19 ± 0.14	38.04 ± 1.69
2752 ± 8	1.50 ± 0.06	2730.38 ± 0.04	1 ⁺ d	4.14 ± 0.14	59.86 ± 2.00
2937 ± 13	0.11 ± 0.02	2950.70 ± 0.50		0.27 ± 0.05	4.28 ± 0.78
3143 ± 20	0.06 ± 0.01	3109.75 ± 0.04	(1,2) ⁺	0.13 ± 0.02	2.29 ± 0.38
3334 ± 19	0.04 ± 0.01	3146.44 ± 0.09	1	0.09 ± 0.04	1.52 ± 0.76
3569 ± 56	0.01 ± 0.01	3293 ± 10		0.02 ± 0.02	0.37 ± 0.37
3652 ± 10	0.16 ± 0.02	3738.50 ± 0.05	1 ⁺	0.24 ± 0.02	5.62 ± 0.35
3786 ± 10	0.26 ± 0.03	3797.58 ± 0.06	1 ⁺	0.39 ± 0.03	9.06 ± 0.70
3861 ± 49	0.01 ± 0.01	3840.25 ± 0.05	(1,2) ⁺	0.02 ± 0.02	0.35 ± 0.35
4067 ± 24	0.05 ± 0.02	3898 ± 8		0.07 ± 0.03	1.72 ± 0.69
4111 ± 30	0.11 ± 0.03	3996 ± 10	U	0.14 ± 0.03	3.74 ± 0.68
4267 ± 10	0.29 ± 0.03	4352 ± 5	U	0.28 ± 0.03	9.38 ± 0.97
4364 ± 8	3.84 ± 0.17	4384.00 ± 0.30	0 ⁺ ;2	3.75 ± 0.17	124.07 ± 5.49
4522 ± 16	0.31 ± 0.05	4697 ± 10	U	0.24 ± 0.04	9.57 ± 1.54
4655 ± 12	0.38 ± 0.06	4761 ± 5	(1,2) ⁺	0.27 ± 0.04	11.64 ± 1.53
4825 ± 21	0.47 ± 0.08	4788.65 ± 0.17	1 ⁺	0.33 ± 0.06	14.34 ± 2.44
5017 ± 27	0.36 ± 0.09	4848 ± 10		0.24 ± 0.06	10.91 ± 2.73
5080 ± 35	0.23 ± 0.07	5027 ± 5		0.13 ± 0.05	6.80 ± 2.36
5223 ± 32	0.03 ± 0.03			0.02 ± 0.02	0.87 ± 0.87
5696 ± 23	0.11 ± 0.04			0.03 ± 0.01	2.97 ± 1.08
6006 ± 21	0.13 ± 0.05			0.03 ± 0.01	3.37 ± 1.30
Total	9.36 ± 0.26			14.02 ± 0.30	320.77 ± 8.27

- The transition strengths were calculated via

$$B_i(GT) + B_i(F) = \frac{K R_i}{t_{1/2} f(E_i)}$$

- Neutrino cross sections were then calculated via

$$\sigma(E_\nu) = \frac{G_F^2 \cos^2(\theta_{ud})}{\pi \hbar^4 c^3} \sum_i p_i W_i F(Z, W_i) [B_i(GT) + B_i(F)]$$

MARLEY

Model of ARgon Low Energy Yield



- Neutrino cross sections were then calculated via

$$\sigma(E_\nu) = \frac{G_F^2 \cos^2(\theta_{ud})}{\pi \hbar^4 c^3} \sum_i p_i W_i F(Z, W_i) [B_i(GT) + B_i(F)]$$

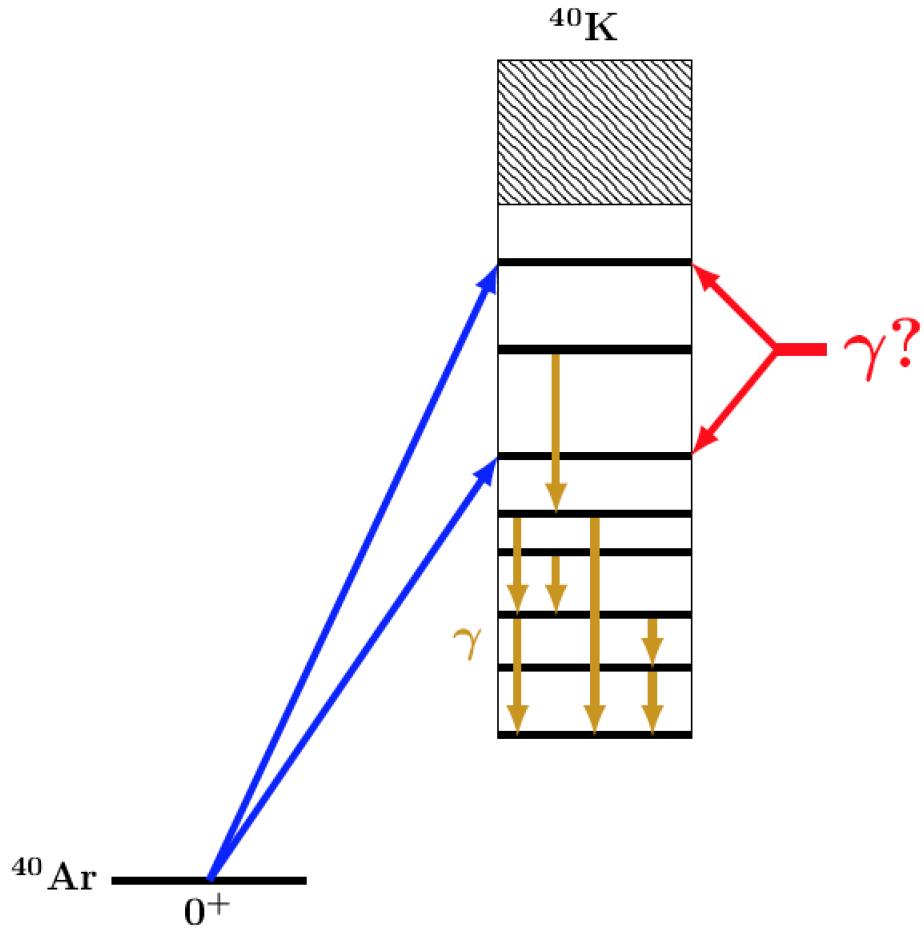
18 in

...based on 40-Ti measurements

What about gamma cascades?

without these neutrino energy cannot be reconstructed

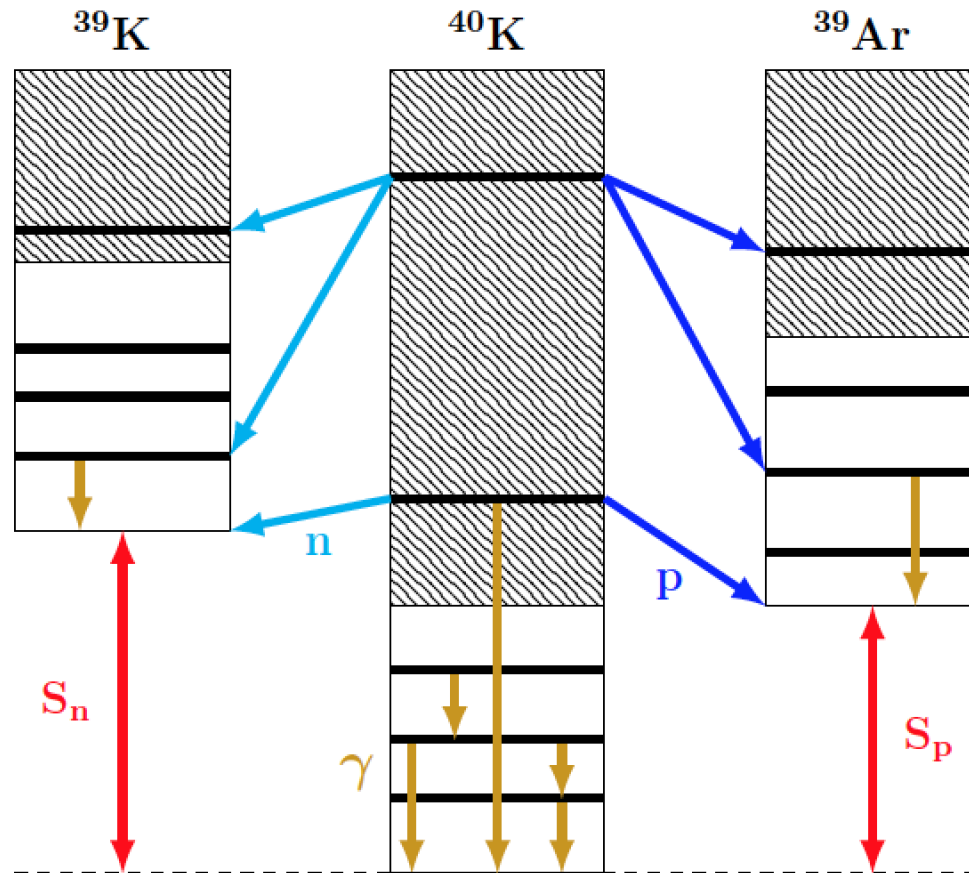
Predicting γ signal is difficult when there are no data



- Make educated guesses about level J^π s
- Use models (e.g., Weisskopf estimates) to approximate γ branching ratios
- Revert to ENSDF data whenever available



Unbound states are even more problematic



- Supplement any available data with nuclear structure and decay models
- Most previous work has ignored these details
- We are currently investigating the best approach

(Help!)



MARLEY

Maybe it would be better to try and measure the cross section directly. At least verify to the 10% level before spending >\$100M to go underground

We are thinking about this...



Summary

- Incredible rich field of neutrino detection
- 1 MeV-100 TeV
- Range of detector technologies to adapt to needed measurements
- ... a real need for R&D to expand the capabilities of existing technology and develop new ones
- CPAD meeting in Arlington Texas – **y'all come!**

NEW TECHNOLOGIES FOR DISCOVERY

Organized by the Coordinating Panel for Advanced Detectors of the
Division of Particles and Fields of the American Physical Society

October 5 - 7, 2015 - University of Texas at Arlington

Home

Committees

Working Groups

Scientific Program

Registration

List of Registrants

Accommodation

Travel Information

Social Program

Sponsors

Welcome to the CPAD Instrumentation Frontier Meeting at UTA

This 2.5 day workshop will provide an evaluation of the Detector R&D program being carried out in support of the High Energy Physics science mission. The goal of the workshop is to determine if the existing program meets the science needs of the five science drivers identified by P5 within the twenty year P5 vision. The program will place an emphasis on proposing new ideas for detection technologies, articulating grand challenges for the field and identifying instrumentation opportunities to enhance the program as an enabler of our science. In addition to plenary sessions and parallel session working groups there will also be an Instrumentation Slam. We guarantee new instrumentation ideas! The outcome of the workshop will be a formulation of the needs of the field to ensure a vibrant long term future.



The workshop is organized by the [Coordinating Panel for Advanced Detectors](#) (CPAD)* of the Division of Particles and Fields of the American Physical Society.

<http://www.uta.edu/physics/cpad2015/>