

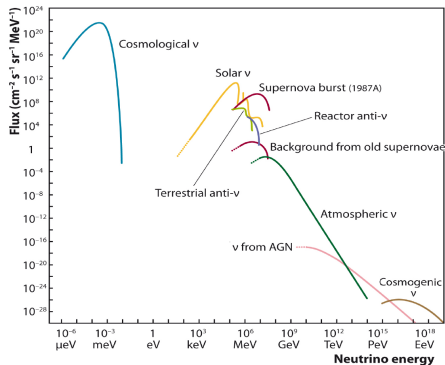


## Detection of the 2 Kelvin Cosmic $\nu$ Background

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43<sup>rd</sup> SLAC Summer Institute – The Universe of Neutrinos

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**Figure:** Flux density at earth's surface due to natural and reactor neutrinos (Eur. Phys. J. H **37**, 515 (2012), arXiv:1207.4952 [astro-ph.IM])

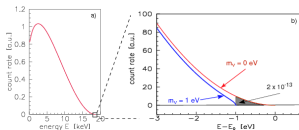
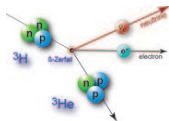
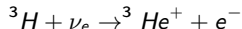
- The  $C\nu B$  is a degenerate Fermi sea of neutrinos that filled the universe [S. Weinberg, Phys. Rev. 128, 1457 (1962)]
- Decoupled from photons, baryons, etc. at  $t \sim 1s$  ( $T_{dec}^\nu \sim 1MeV$ )
- Current temperature of  $C\nu B$  is about  $2K$  ( $\sim 0.1meV$ )
- Current density (all flavors of  $\nu$  and  $\bar{\nu}$  combined)  $n \sim 300cm^{-3}$
- Flavor states decohered  $\implies$  currently in mass eigenstates

## Proposed experiments

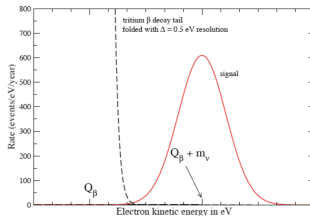
## $\beta$ endpoint experiments: Method

$\beta$  Decay: Used for measuring mass of neutrino (KATRIN, Project 8)

Neutrino Capture: Potentially a method to measure the C $\nu$ B (PTOLEMY)



- Non-zero mass of neutrino will distort the  $\beta$  decay tail



- Signal is a mono-energetic electron beyond  $\beta$  decay tail
- Process is exothermic, no threshold energy

## $\beta$ endpoint experiments: PTOLEMY

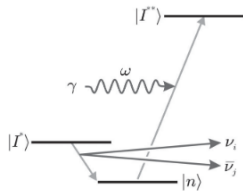
### Princeton Tritium Observatory for Light, Early-Universe, Massive-Neutrino Yield (PTOLEMY)

**Target:** Large active target  $\sim 1\text{M Ci}$  (100 grams of Tritium)

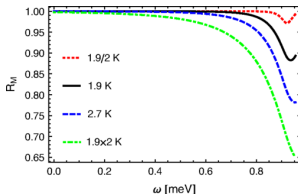
**Energy Resolution:** High resolution ( $\Delta E < m_\nu$ ) required for signal-background separation. Neutrino mass mixing parameters indicate mass eigenstates at least as massive as 0.05 eV while cosmological bounds  $< 0.3$  eV. Radio-Frequency techniques strive to reach sub-eV energy resolution.

**Background:** Below microHertz of background rate in signal region (PTOLEMY expected signal rate of approximately  $0.3 \mu\text{Hz}$  of neutrino capture events)

## Indirect detection by atomic de-excitation



Energy level diagram [Phys. Rev. D 75, 113007]



Spectral distortion caused by Pauli blocking of relic neutrinos assuming the lightest neutrino mass is 5meV,  $\epsilon_{eg} = 11\text{meV}$  and zero chemical potential.

[Phys. Rev. D 91, 063516]

- Radiative emission of neutrino pairs (RENP) can happen as atoms de-excite from a metastable state.
- Use laser to enhance emission probability.
- Compare LH and RH circular polarizations to reject non-parity-violating (non-weak) backgrounds (i.e. photon emission).
- Expected rate for  $10^{16}$  atoms  $\sim 1/\text{day}$ .
- RENP spectrum is modified in the presence of ambient relic neutrinos by a factor of  $(1 - f_i)(1 - \bar{f}_j)$ , where  $f_i$  ( $\bar{f}_j$ ) is the momentum distribution function of mass eigenstate  $\nu_i$  ( $\bar{\nu}_j$ ).
- Various caveats such as finding the right atom. No experiment has looked for this process.

# Our Ideas

## Superconductor experiments

- The binding energy of Cooper pairs in superconductors is  $\mathcal{O}(\text{meV})$ .
- Held near its transition temperature, depositions of this order of energy create sharp changes in the resistance of the semiconductor.
- Intriguingly, this is the same order of energy as the average neutrino in the  $C\nu B$ .
- However, even electrons are much more massive than neutrinos that only a tiny fraction of the neutrino energy is transferred in elastic scattering at  $C\nu B$  energies, so this idea does not apply well to the question.
- This idea is pursued in the context of light WIMP detection by Hochberg et al. (1504.07237).



## Proton-Stimulated Inverse $\beta$ Decay

What if we invert the energy distribution of a typical Inverse Beta Decay neutrino experiment?

The rate of inverse beta decay in the laboratory frame is

$$\Gamma \sim 10^{-27} \left( \frac{E_p}{\text{GeV}} \right)^2 N_{\text{particles}} \text{Hz}$$

For nominal beam parameters (e.g. the high-luminosity LHC) with  $E_p \sim 7$  TeV and  $\sim 3 \times 10^{15}$  particles per beam (and two beams),

$$\Gamma \sim 2 \times 10^{-4} \text{Hz}$$

which is  $\sim 24$  events per day along the entire 27 km in circumference.

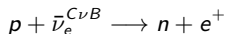
This comes with significant experimental challenges:

- For 100 events per year, we then only have to instrument  $\sim 300$  m of beam tube cumulatively
- Background primarily high energy neutrinos; need at least 100 keV resolution to resolve separate cosmic neutrino signal
- We need high ( $>90\%$ ) reconstruction efficiency for coincident neutron-positron pairs

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## However...

This experiment would be looking to detect proton Inverse  $\beta$  Decay



Which is allowed only if the center-of-mass energy is above the threshold

$$\sqrt{s} \geq M_n + m_e - M_p - m_\nu \simeq M_n + m_e - M_p \simeq 1.80 \text{ MeV}$$

For  $E_{\nu^{C\nu B}} \simeq m_\nu$  we have  $\sqrt{s} = \sqrt{2E_p m_\nu}$  so that a proton beam with energy  $E_p$  can only observe this process if

$$m_\nu \geq \frac{(M_n + m_e - M_p)^2}{2E_p} \quad \text{i.e.} \quad m_\nu [\text{eV}] \gtrsim \frac{1.63}{E_p [\text{TeV}]}$$

This means that LHC ( $E_p \sim 7 \text{ TeV}$ ) can only probe  $m_\nu$  regions of  $m_\nu \gtrsim 0.23$  eV; however, it's close. If next generation of accelerators were to have a beam energy that was higher by a factor of 5 or more, we would be guaranteed to be above threshold for at least one neutrino.

## Outtakes: Ideas that didn't quite pan out

**Detection of  $W^\pm$  peak from  $e^- + \nu_e^{C\nu B}$  scattering:** can be performed at  $e^-$  accelerator. Required beam energy is YeV i.e.  $10^{24}$  eV to achieve  $\sqrt{s} \sim 100$  GeV.  $\nu$  accelerator can look for  $Z^0$  peak with same energy issue, cleaner signature but need  $\nu$  beam.

**Annual modulation of deflection of torsion pendulum:** A Cavendish-type torsion balance that measures the "cosmic  $\nu$  wind" due to the solar system motion through the galaxy. We need 2-3 orders magnitude higher precision than today's technology.

**" $1\nu\beta\beta$ ":** Idea is to detect  $\frac{A}{Z}X + \nu_e^{C\nu B} \rightarrow \frac{A}{Z+2}Y + \bar{\nu}_e + 2e^-$  at  $0\nu\beta\beta$  experiments; hope is to exclude energetic  $\nu$  because  $\frac{A}{Z}X + \nu_e^{C\nu B} \rightarrow \frac{A}{Z+1}Z + e^-$  is above threshold.  
 $\text{Rate}_{1\nu\beta\beta} / \text{Rate}_{2\nu\beta\beta} \sim 10^{-21}$

**Ultrarelativistic water Čerenkov detector:** Idea is to detect  $\nu^{C\nu B}$  by being in a reference frame in which the products of their interactions emit Čerenkov light. Requires a detector to be moving really fast.

# Conclusions

## Conclusions

- There is increasing interest in a  $C\nu B$  experiment due to advances in low-background and high-resolution detector technology
  - For current technologies, cost and detector mass are the main current limitations
  - The first cosmic neutrino detector will have to be very different from our current neutrino detectors!
- Current proposals include ideas from particle physics, condensed matter and atomic physics
- Many of these ideas may need to wait for significant technological advances
  - If we could generate  $YeV$  particle beams, many of the experiments would be immediately feasible with reasonable detector masses
  - Very high energy resolution is a requirement for any cosmic neutrino detector, given the unavoidable background from cosmogenic neutrinos and other low-background processes