



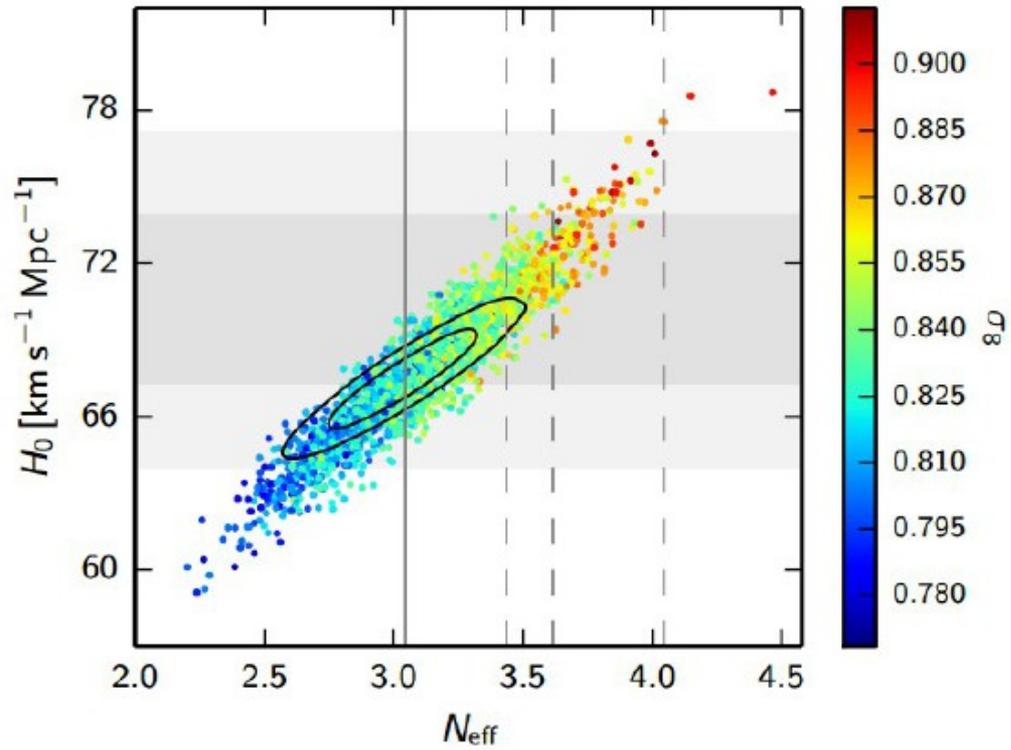
# Neutrino Cosmology (II)

Stanford, August 12th 2015

Alessandro Melchiorri

University of Rome “La Sapienza”

(Disclaimer: many thanks to Eleonora Di Valentino and Sergio Pastor for some of the slides)



$$N_{\text{eff}} = 3.13 \pm 0.32 \quad \textit{Planck TT+lowP};$$

$$N_{\text{eff}} = 3.15 \pm 0.23 \quad \textit{Planck TT+lowP+BAO};$$

$$N_{\text{eff}} = 2.99 \pm 0.20 \quad \textit{Planck TT, TE, EE+lowP};$$

$$N_{\text{eff}} = 3.04 \pm 0.18 \quad \textit{Planck TT, TE, EE+lowP+BAO}.$$

# Mechanisms for having $N_{\text{eff}} > 3.046$

Non-standard decoupling (Mangano et al. 2006)

$N_{\text{eff}} \sim 3.04 - 3.1$

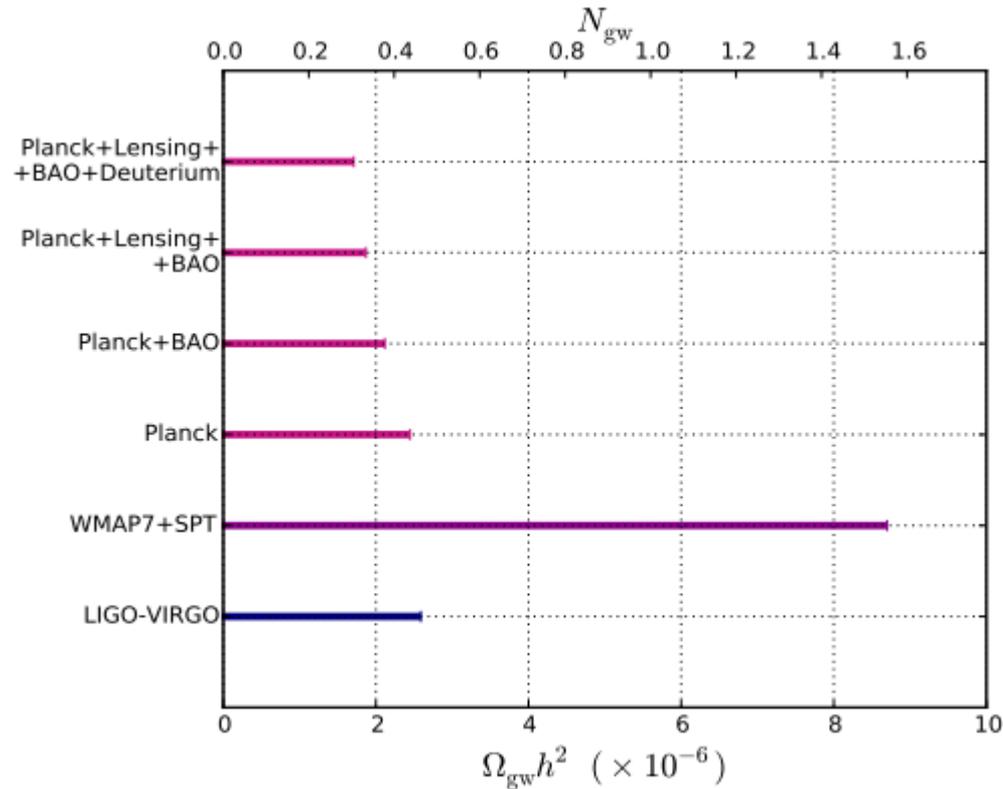
Thermal axions (see eg. Di Valentino et al. 2015)

Gravitational waves (Smith et al., 2006)

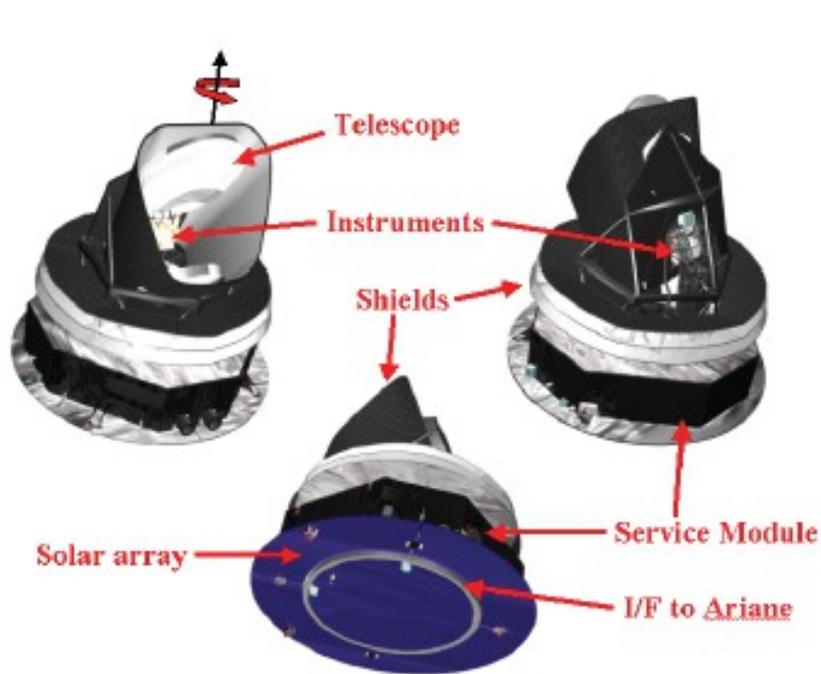
Extra dimensions

....

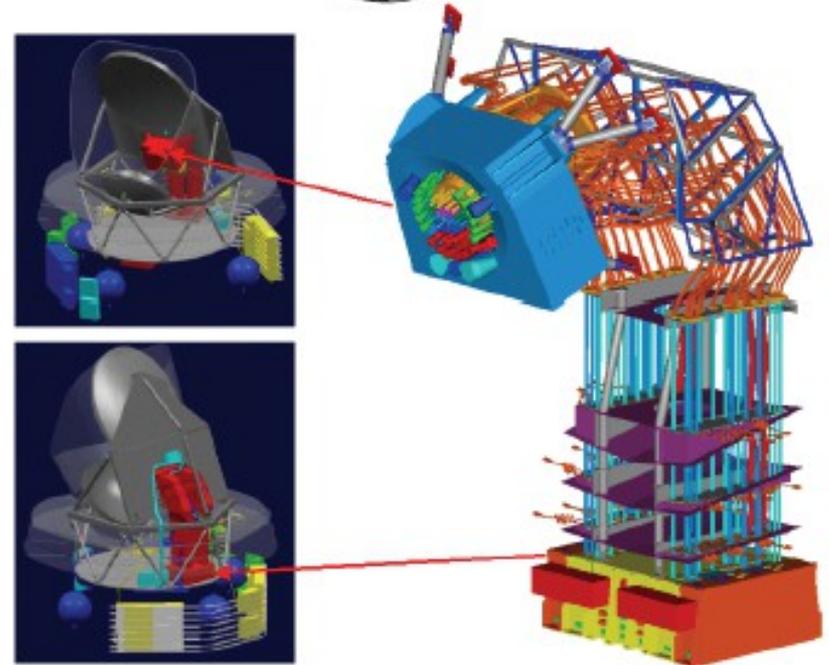
# New constraints from Planck on GW



# Planck in cartoons



Planck has two instruments, the Low Frequency Instrument (LFI) and the High Frequency Instrument (HFI) in a shared focal plane containing 74 channels and covering 8 degrees on the sky.



# Planck being assembled



# Planck in February 2009



Planck  
Satellite launch  
14/5/2009



## PLANCK'S LAST COMMAND



23 October 2013 ESA's Planck space telescope has been turned off after nearly 4.5 years soaking up the relic radiation from the Big Bang and studying the evolution of stars and galaxies throughout the Universe's history.

# Planck Collaboration 400+ names

## Planck 2015 results. I. Overview of products and scientific results

Planck Collaboration: R. Adam<sup>97</sup>, P. A. R. Ade<sup>114</sup>, N. Aghanim<sup>79</sup>, Y. Akrami<sup>83,134</sup>, M. I. R. Alves<sup>79</sup>, M. Arnaud<sup>95</sup>, F. Arroja<sup>87,101</sup>, J. Aumont<sup>79</sup>, C. Baccigalupi<sup>113</sup>, M. Ballardini<sup>66,68,40</sup>, A. J. Banday<sup>129,12</sup>, R. B. Barreiro<sup>86</sup>, J. G. Bartlett<sup>1,88</sup>, N. Bartolo<sup>41,87</sup>, S. Basak<sup>113</sup>, P. Battaglia<sup>124</sup>, E. Battaner<sup>132,133</sup>, R. Battye<sup>89</sup>, K. Benabed<sup>80,126</sup>, A. Benoit<sup>77</sup>, A. Benoit-Lévy<sup>30,80,126</sup>, J.-P. Bernard<sup>129,12</sup>, M. Bersanelli<sup>44,67</sup>, B. Bertinocourt<sup>79</sup>, P. Bielewicz<sup>129,12,113</sup>, A. Bonaldi<sup>89</sup>, L. Bonavera<sup>86</sup>, J. R. Bond<sup>11</sup>, J. Borrill<sup>17,119</sup>, F. R. Bouchet<sup>80,117</sup>, F. Boulanger<sup>79,110</sup>, M. Bucher<sup>1</sup>, C. Burigana<sup>66,42,68</sup>, R. C. Butler<sup>66</sup>, E. Calabrese<sup>122</sup>, J.-F. Cardoso<sup>96,1,80</sup>, P. Carvalho<sup>82,90</sup>, B. Casaponsa<sup>86</sup>, G. Castex<sup>1</sup>, A. Catalano<sup>97,93</sup>, A. Challinor<sup>82,90,15</sup>, A. Chamballu<sup>95,19,79</sup>, R.-R. Chary<sup>78</sup>, H. C. Chiang<sup>34,9</sup>, J. Chluba<sup>29,90</sup>, P. R. Christensen<sup>107,48</sup>, S. Church<sup>121</sup>, M. Clemens<sup>63</sup>, D. L. Clements<sup>75</sup>, S. Colombi<sup>80,136</sup>, L. P. L. Colombo<sup>28,88</sup>, C. Combet<sup>97</sup>, B. Comis<sup>97</sup>, D. Contreras<sup>27</sup>, F. Couchot<sup>91</sup>, A. Coullais<sup>93</sup>, B. P. Crill<sup>88,108</sup>, M. Cruz<sup>24</sup>, A. Curto<sup>8,86</sup>, F. Cuttaia<sup>66</sup>, L. Danese<sup>113</sup>, R. D. Davies<sup>89</sup>, R. J. Davis<sup>89</sup>, P. de Bernardis<sup>43</sup>, A. de Rosa<sup>66</sup>, G. de Zotti<sup>63,113</sup>, J. Delabrouille<sup>1</sup>, J.-M. Delouis<sup>80,126</sup>, F.-X. Désert<sup>72</sup>, E. Di Valentino<sup>43</sup>, C. Dickinson<sup>89</sup>, J. M. Diego<sup>86</sup>, K. Dolag<sup>131,102</sup>, H. Dole<sup>79,78</sup>, S. Donzelli<sup>67</sup>, O. Doré<sup>88,14</sup>, M. Douspis<sup>79</sup>, A. Ducout<sup>80,75</sup>, J. Dunkley<sup>122</sup>, X. Dupac<sup>52</sup>, G. Efstathiou<sup>42</sup>, P. R. M. Eisenhardt<sup>88</sup>, F. Elsner<sup>30,80,126</sup>, T. A. Enßlin<sup>102</sup>, H. K. Eriksen<sup>83</sup>, E. Falgarone<sup>93</sup>, Y. Fantaye<sup>83</sup>, M. Farhang<sup>11,111</sup>, S. Feeney<sup>75</sup>, J. Fergusson<sup>15</sup>, R. Fernandez-Cobos<sup>86</sup>, F. Feroz<sup>8</sup>, F. Finelli<sup>66,68</sup>, E. Florido<sup>132</sup>, O. Forni<sup>129,12</sup>, M. Frailis<sup>65</sup>, A. A. Fraisse<sup>34</sup>, C. Franceschet<sup>44</sup>, E. Franceschi<sup>66</sup>, A. Frejsel<sup>107</sup>, A. Frolov<sup>116</sup>, S. Galeotta<sup>65</sup>, S. Galli<sup>80</sup>, K. Ganga<sup>1</sup>, C. Gauthier<sup>1,101</sup>, R. T. Génova-Santos<sup>85</sup>, M. Gerbino<sup>43</sup>, T. Ghosh<sup>79</sup>, M. Giard<sup>129,12</sup>, Y. Giraud-Héraud<sup>1</sup>, E. Giusarma<sup>43</sup>, E. Gjerløw<sup>83</sup>, J. González-Nuevo<sup>86,113</sup>, K. M. Górski<sup>88,136</sup>, K. J. B. Grainge<sup>8,90</sup>, S. Gratton<sup>90,82</sup>, A. Gregorio<sup>45,65,71</sup>, A. Gruppuso<sup>66</sup>, J. E. Gudmundsson<sup>34</sup>, J. Hamann<sup>125,123</sup>, W. Handley<sup>80,8</sup>, F. K. Hansen<sup>83</sup>, D. Hanson<sup>104,88,11</sup>, D. L. Harrison<sup>82,90</sup>, A. Heavens<sup>75</sup>, G. Helou<sup>14</sup>, S. Henrot-Versillé<sup>91</sup>, C. Hernández-Monteagudo<sup>16,102</sup>, D. Herranz<sup>86</sup>, S. R. Hildebrandt<sup>88,14</sup>, E. Hivon<sup>80,126</sup>, M. Hobson<sup>8</sup>, W. A. Holmes<sup>80</sup>, A. Hornstrup<sup>20</sup>, W. Hovest<sup>102</sup>, Z. Huang<sup>11</sup>, K. M. Huffenberger<sup>32</sup>, G. Hurier<sup>79</sup>, S. Ilić<sup>129,12,79</sup>, A. H. Jaffe<sup>75</sup>, T. R. Jaffe<sup>129,12</sup>, T. 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Moss<sup>15</sup>, S. Mottet<sup>80</sup>, M. Müenchemeyer<sup>80</sup>, D. Munshi<sup>114</sup>, J. A. Murphy<sup>106</sup>, A. Narimani<sup>27</sup>, P. Naselsky<sup>107,48</sup>, A. Nastasi<sup>79</sup>, F. Nati<sup>34</sup>, P. Natoli<sup>42,5,66</sup>, M. Negrello<sup>63</sup>, C. B. Netterfield<sup>25</sup>, H. U. Nørgaard-Nielsen<sup>20</sup>, F. Noviello<sup>89</sup>, D. Novikov<sup>100</sup>, I. Novikov<sup>107,100</sup>, M. Olamaica<sup>8</sup>, N. Oppermann<sup>11</sup>, E. Orlando<sup>135</sup>, C. A. Oxborrow<sup>20</sup>, F. Paci<sup>113</sup>, L. Pagano<sup>43,69</sup>, F. Pajot<sup>79</sup>, R. Paladini<sup>76</sup>, S. Pandolfi<sup>22</sup>, D. Paoletti<sup>66,68</sup>, B. Partridge<sup>59</sup>, F. Pasian<sup>65</sup>, G. Patanchon<sup>1</sup>, T. J. Pearson<sup>14,76</sup>, M. Peel<sup>89</sup>, H. V. Peiris<sup>30</sup>, V.-M. Pelkonen<sup>78</sup>, O. Perdereau<sup>91</sup>, L. Perotto<sup>97</sup>, Y. C. Perrott<sup>8</sup>, F. Perrotta<sup>113</sup>, V. Pettorino<sup>38</sup>, F. Piacentini<sup>43</sup>, M. Piat<sup>1</sup>, E. Pierpaoli<sup>28</sup>, D. Pietrobon<sup>88</sup>, S. Plaszczyński<sup>91</sup>, D. Pogosyan<sup>35</sup>, E. Pointecouteau<sup>129,12</sup>, G. Polenta<sup>5,64</sup>, L. Popa<sup>81</sup>, G. W. Pratt<sup>95</sup>, G. Príncipe<sup>14,88</sup>, S. Prunet<sup>80,126</sup>, J.-L. Puget<sup>79</sup>, J. P. Rachen<sup>26,102</sup>, B. Racine<sup>1</sup>, W. T. Reach<sup>130</sup>, R. Rebolo<sup>85,18,49</sup>, M. Reinecke<sup>102</sup>, M. Remazeilles<sup>89,79,1</sup>, C. Renault<sup>97</sup>, A. Renzi<sup>47,70</sup>, I. Ristorcelli<sup>129,12</sup>, G. Rocha<sup>88,14</sup>, M. Roman<sup>1</sup>, E. Romelli<sup>45,65</sup>, C. Rosset<sup>1</sup>, M. Rossetti<sup>44,67</sup>, A. Rotti<sup>74</sup>, G. Roudier<sup>1,93,88</sup>, B. Rouillé d'Orfeuil<sup>91</sup>, M. Rowan-Robinson<sup>75</sup>, J. A. Rubiño-Martín<sup>85,49</sup>, B. Ruiz-Granados<sup>132</sup>, C. Rumsey<sup>8</sup>, B. Rusholme<sup>76</sup>, N. Saïd<sup>43</sup>, V. Salvatelli<sup>43,7</sup>, L. Salvati<sup>43</sup>, M. Sandri<sup>66</sup>, H. S. Sanghera<sup>82,90</sup>, D. Santos<sup>97</sup>, R. D. E. Saunders<sup>8,90</sup>, A. Sauvé<sup>129,12</sup>, M. Savelainen<sup>33,60</sup>, G. Savini<sup>109</sup>, B. M. Schaefer<sup>127</sup>, M. P. Schammel<sup>8</sup>, D. Scott<sup>27</sup>, M. D. Seiffert<sup>88,14</sup>, P. Serra<sup>79</sup>, E. P. S. Shellard<sup>15</sup>, T. W. Shimwell<sup>8</sup>, M. Shiraishi<sup>81,87</sup>, K. Smith<sup>34</sup>, T. Souradeep<sup>74</sup>, L. D. Spencer<sup>114</sup>, M. Spinelli<sup>91</sup>, S. A. Stanford<sup>37</sup>, D. Stern<sup>89</sup>, V. Stolyarov<sup>8,90,120</sup>, R. Stompor<sup>1</sup>, A. W. Strong<sup>103</sup>, R. Sudiwala<sup>114</sup>, R. Sunyaev<sup>102,118</sup>, P. Sutter<sup>80</sup>, D. Sutton<sup>82,90</sup>, A.-S. Suur-Uksi<sup>33,60</sup>, J.-F. Sygnet<sup>80</sup>, J. A. Tauber<sup>53</sup>, D. Tavagnacco<sup>65,45</sup>, L. Terenzi<sup>54,66</sup>, D. Texier<sup>51</sup>, L. Toffolatti<sup>23,86,66</sup>, M. Tomasi<sup>44,67</sup>, M. Tomikowski<sup>3</sup>, M. Tristram<sup>91</sup>, A. Troja<sup>44</sup>, T. Trombetti<sup>66</sup>, M. Tucci<sup>21</sup>, J. Tuovinen<sup>13</sup>, M. Türlér<sup>73</sup>, G. Umata<sup>62</sup>, L. Valenziano<sup>66</sup>, J. Valiviita<sup>33,60</sup>, B. Van Tent<sup>98</sup>, T. Vassallo<sup>65</sup>, M. Vidal<sup>89</sup>, M. Viel<sup>65,71</sup>, P. Vielva<sup>86</sup>, F. Villa<sup>66</sup>, L. A. Wade<sup>88</sup>, B. Walter<sup>89</sup>, B. D. Wandelt<sup>80,126,39</sup>, R. Watson<sup>89</sup>, I. K. Wehus<sup>88</sup>, N. Welikala<sup>122</sup>, J. Weller<sup>131</sup>, M. White<sup>36</sup>, S. D. M. White<sup>102</sup>, A. Wilkinson<sup>89</sup>, D. Yvon<sup>19</sup>, A. Zacchei<sup>65</sup>, J. P. Zibin<sup>27</sup>, and A. Zonca<sup>38</sup>

# Planck 2013 Core-Team

(a fraction of it)



# Planck 2015 Core-Team

(again, a fraction of it)



# PLANCK 2013: High Impact !



## Top Cited Articles during 2014

1. [1739](#) core citations in 2014

**Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC**

ATLAS Collaboration (Georges Aad (Freiburg U.) *et al.*). Jul 2012. 24 pp.

Published in *Phys.Lett. B716 (2012) 1-29*

CERN-PH-EP-2012-218

DOI: [10.1016/j.physletb.2012.08.020](#)

e-Print: [arXiv:1207.7214](#) [hep-ex] | [PDF](#)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)

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2. [1715](#) core citations in 2014

**Observation of a new boson at a mass of 125 GeV with the CMS experiment at the LHC**

CMS Collaboration (Serguei Chatrchyan (Yerevan Phys. Inst.) *et al.*). Jul 2012. 42 pp.

Published in *Phys.Lett. B716 (2012) 30-61*

CMS-HIG-12-028, CERN-PH-EP-2012-220

DOI: [10.1016/j.physletb.2012.08.021](#)

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3. [1664](#) core citations in 2014

**Planck 2013 results. XVI. Cosmological parameters**

Planck Collaboration (P.A.R. Ade (Cardiff U.) *et al.*). Mar 20, 2013. 67 pp.

Published in *Astron.Astrophys. 571 (2014) A16*

DOI: [10.1051/0004-6361/201321591](#)

e-Print: [arXiv:1303.5076](#) [astro-ph.CO] | [PDF](#)

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# Planck 2015 release (papers)

- I. Overview of products and results (*this paper*)
- II. Low Frequency Instrument data processing
  - III. LFI systematic uncertainties
  - IV. LFI beams and window functions
  - V. LFI calibration
  - VI. LFI maps
- VII. High Frequency Instrument data processing: Time-ordered information and beam processing
- VIII. High Frequency Instrument data processing: Calibration and maps
- IX. Diffuse component separation: CMB maps
- X. Diffuse component separation: Foreground maps
- XI. CMB power spectra, likelihood, and consistency of cosmological parameters
- XII. Simulations
- XIII. Cosmological parameters
- XIV. Dark energy and modified gravity
- XV. Gravitational lensing
- XVI. Isotropy and statistics of the CMB
- XVII. Primordial non-Gaussianity
- XVIII. Background geometry and topology of the Universe
- XIX. Constraints on primordial magnetic fields
- XX. Constraints on inflation
- XXI. The integrated Sachs-Wolfe effect
- XXII. A map of the thermal Sunyaev-Zeldovich effect
- XXIII. The thermal Sunyaev-Zeldovich effect–cosmic infrared background correlation
- XXIV. Cosmology from Sunyaev-Zeldovich cluster counts
- XXV. Diffuse, low-frequency Galactic foregrounds
- XXVI. The Second Planck Catalogue of Compact Sources
- XXVII. The Second Planck Catalogue of Sunyaev-Zeldovich Sources
- XXVIII. The Planck Catalogue of Galactic Cold Clumps

# Planck 2015: still very high impact !



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HEP

12 record trovati

La ricerca ha impiegato 25.19 secondi.

## 1. Planck 2015 results. XIII. Cosmological parameters

<sup>(535)</sup>Planck Collaboration (P.A.R. Ade (Cardiff U.) *et al.*), Feb 5, 2015.

e-Print: [arXiv:1502.01589](#) [[astro-ph.CO](#)] | [PDF](#)

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## 2. Planck 2015 results. XX. Constraints on inflation

<sup>(233)</sup>Planck Collaboration (P.A.R. Ade (Cardiff U.) *et al.*), Feb 7, 2015. 64 pp.

e-Print: [arXiv:1502.02114](#) [[astro-ph.CO](#)] | [PDF](#)

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## 3. Joint Analysis of BICEP2/Keck Array and Planck Data

<sup>(168)</sup>BICEP2 and Planck Collaborations (P. A. R. Ade (Cardiff U.) *et al.*), Feb 2, 2015. 17 pp.

Published in *Phys.Rev.Lett.* **114** (2015) 101301

DOI: [10.1103/PhysRevLett.114.101301](#)

e-Print: [arXiv:1502.00612](#) [[astro-ph.CO](#)] | [PDF](#)

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## 4. Precise determination of the mass of the Higgs boson and tests of compatibility of its couplings with the standard model predictions using proton collisions at 7 and 8 TeV

<sup>(162)</sup>CMS Collaboration (Vardan Khachatryan (Yerevan Phys. Inst.) *et al.*), Dec 30, 2014. 75 pp.

Published in *Eur.Phys.J.* **C75** (2015) 5, 212

CMS-HIG-14-009, CERN-PH-EP-2014-288

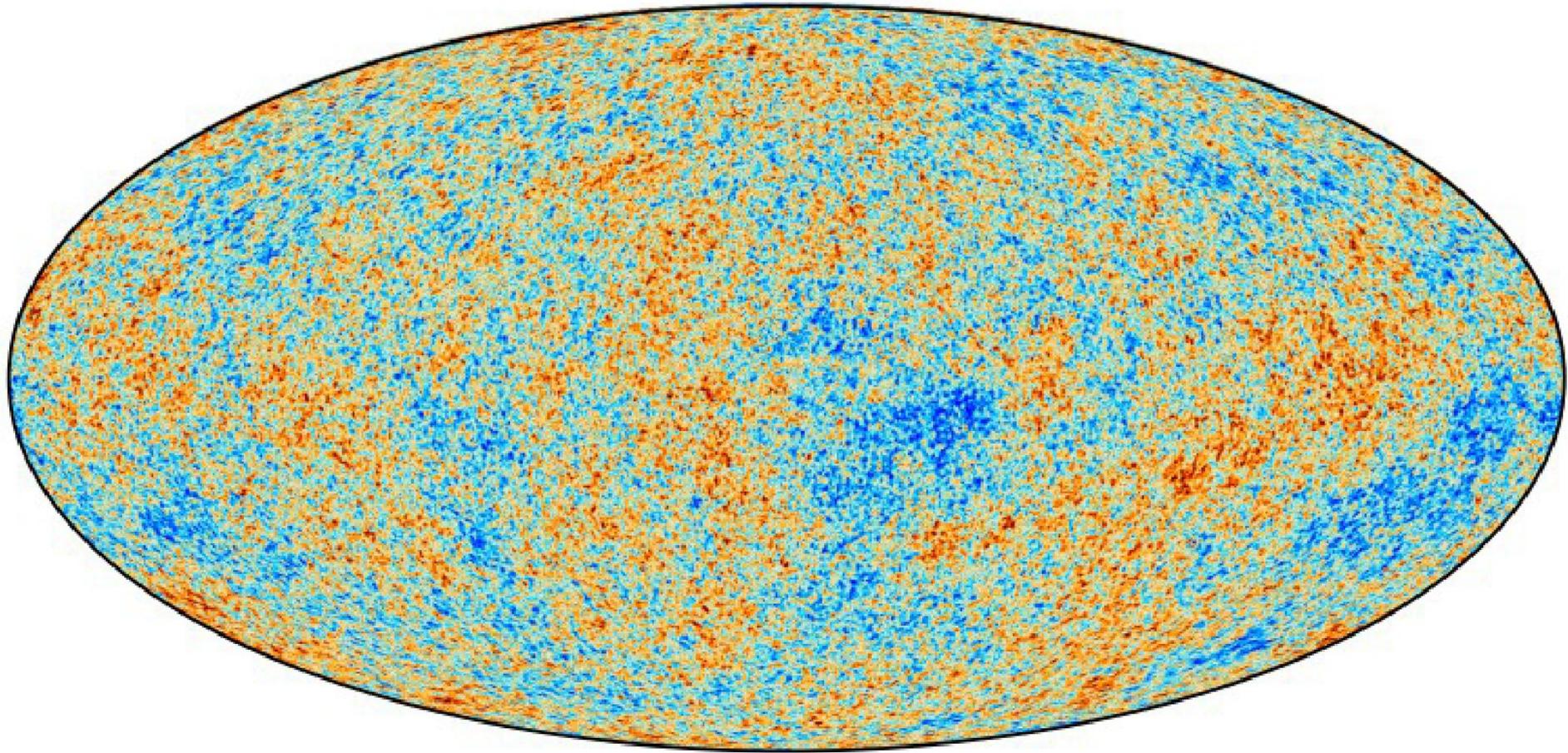
DOI: [10.1140/epjc/s10052-015-3351-7](#)

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# Planck 2015 CMB Temperature map

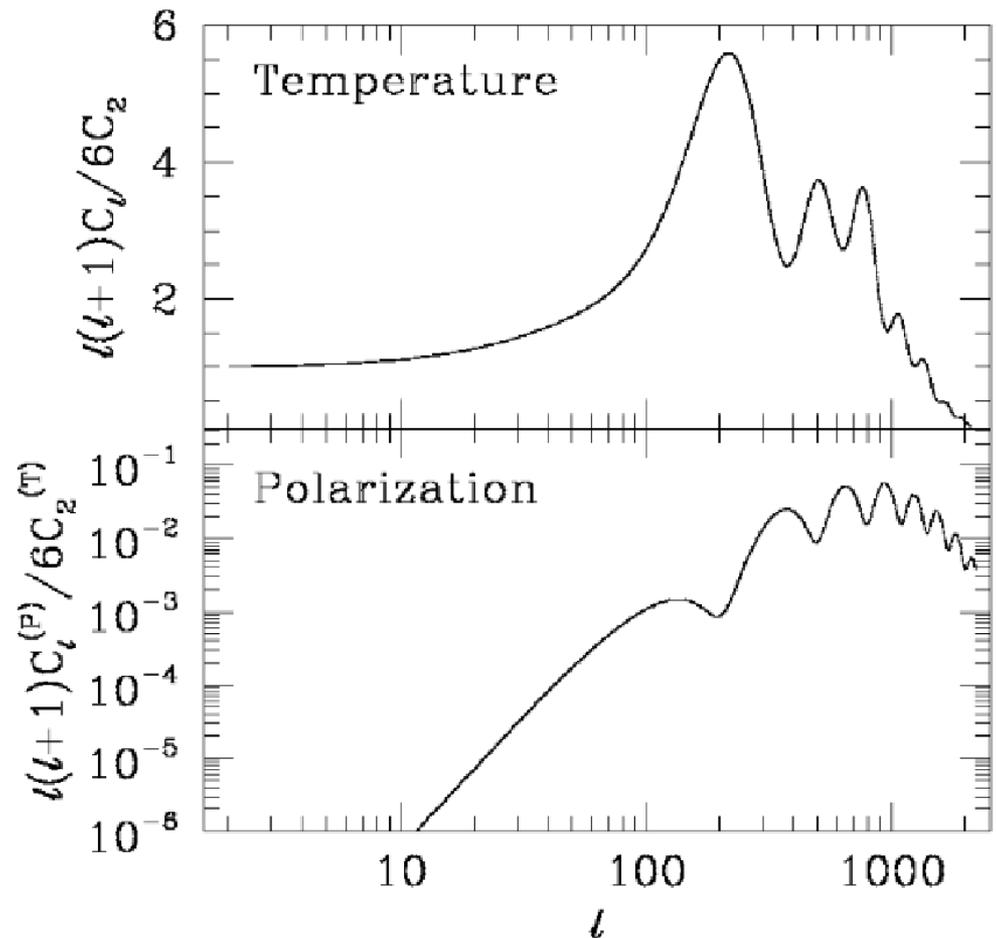


# The CMB Angular Power Spectrum

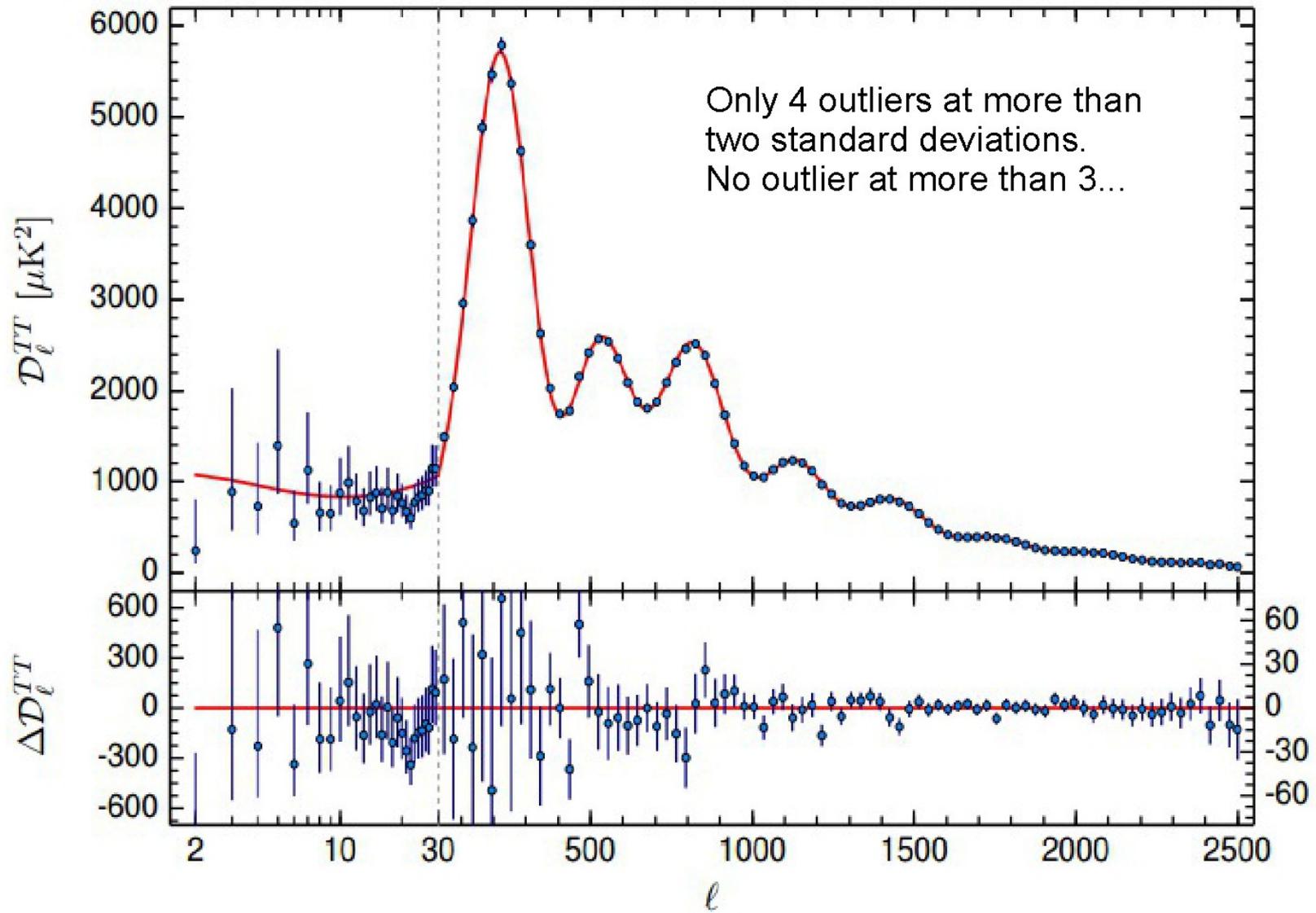
$$\left\langle \frac{\Delta T}{T}(\vec{\gamma}_1) \frac{\Delta T}{T}(\vec{\gamma}_2) \right\rangle = \frac{1}{2\pi} \sum_{\ell} (2\ell + 1) C_{\ell} P_{\ell}(\vec{\gamma}_1 \cdot \vec{\gamma}_2)$$

R.m.s. of  $\Delta T/T$  has  $l(l+1)C_l/2\pi$  power per decade in  $l$ :

$$\langle (\Delta T/T)^2 \rangle_{rms} = \sum_l \frac{(2l+1)}{4\pi} C_l \approx \int \frac{l(l+1)}{2\pi} C_l d \ln l$$



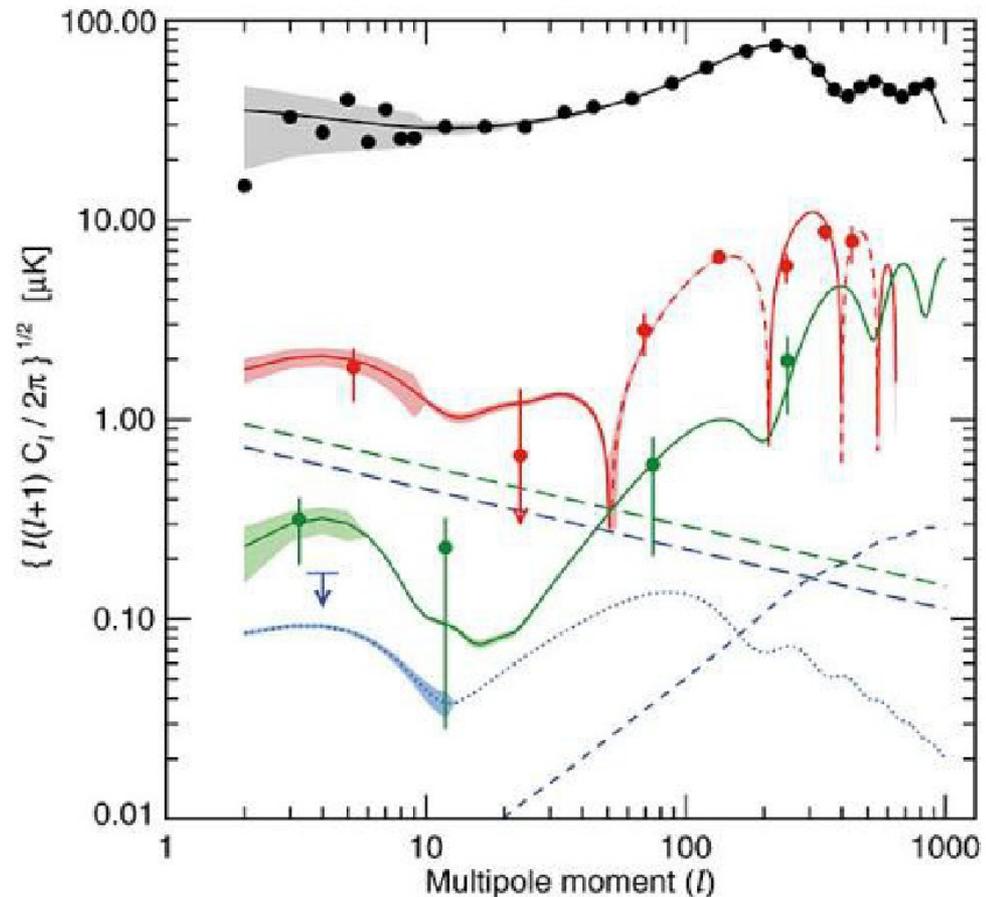
# Planck 2015 Temperature Angular Spectrum



# The CMB Angular Power Spectrum

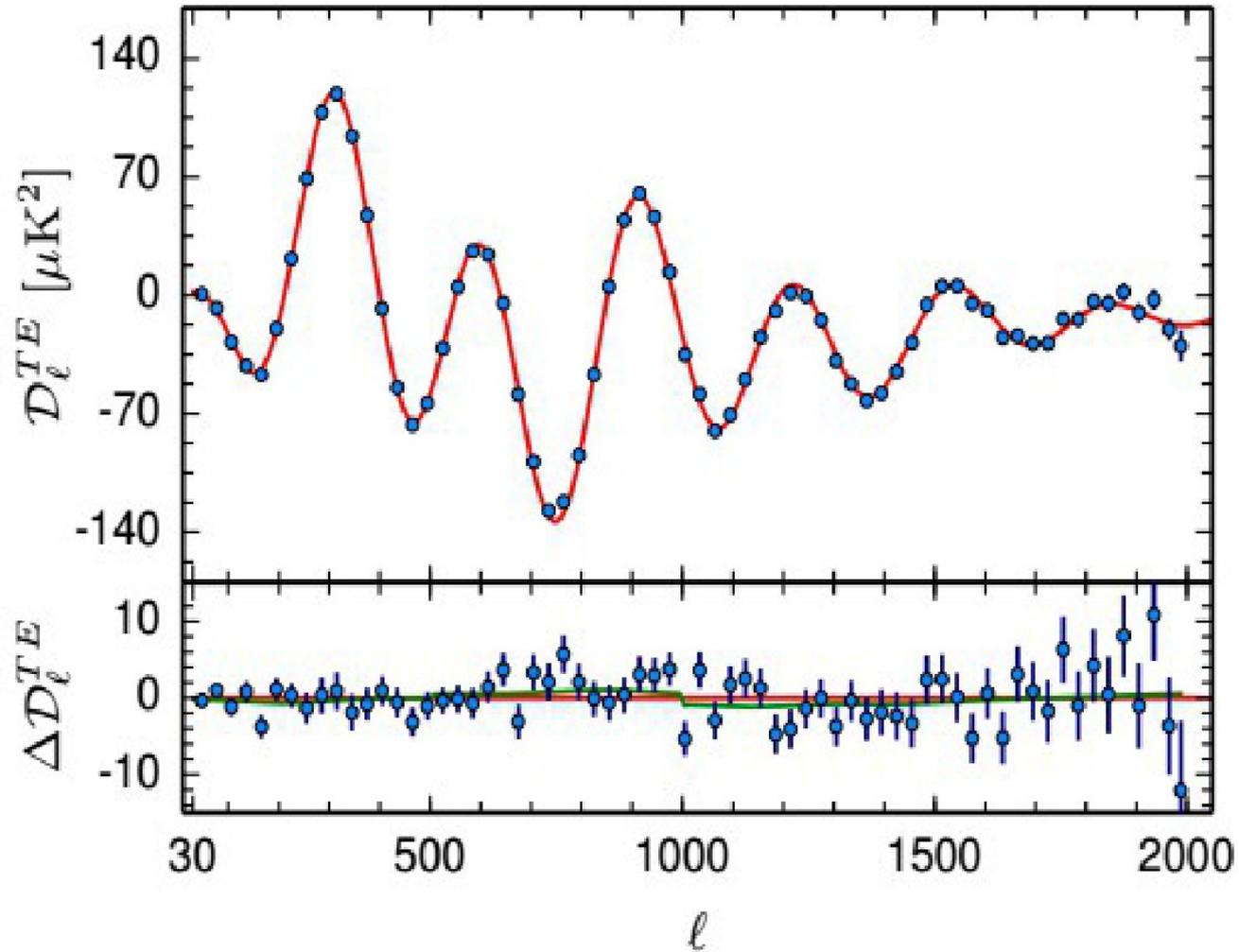
We can extract 4 independent angular spectra from the CMB:

- Temperature
- Cross Temperature Polarization
- Polarization type E (density fluctuations)
- Polarization type B (gravity waves)



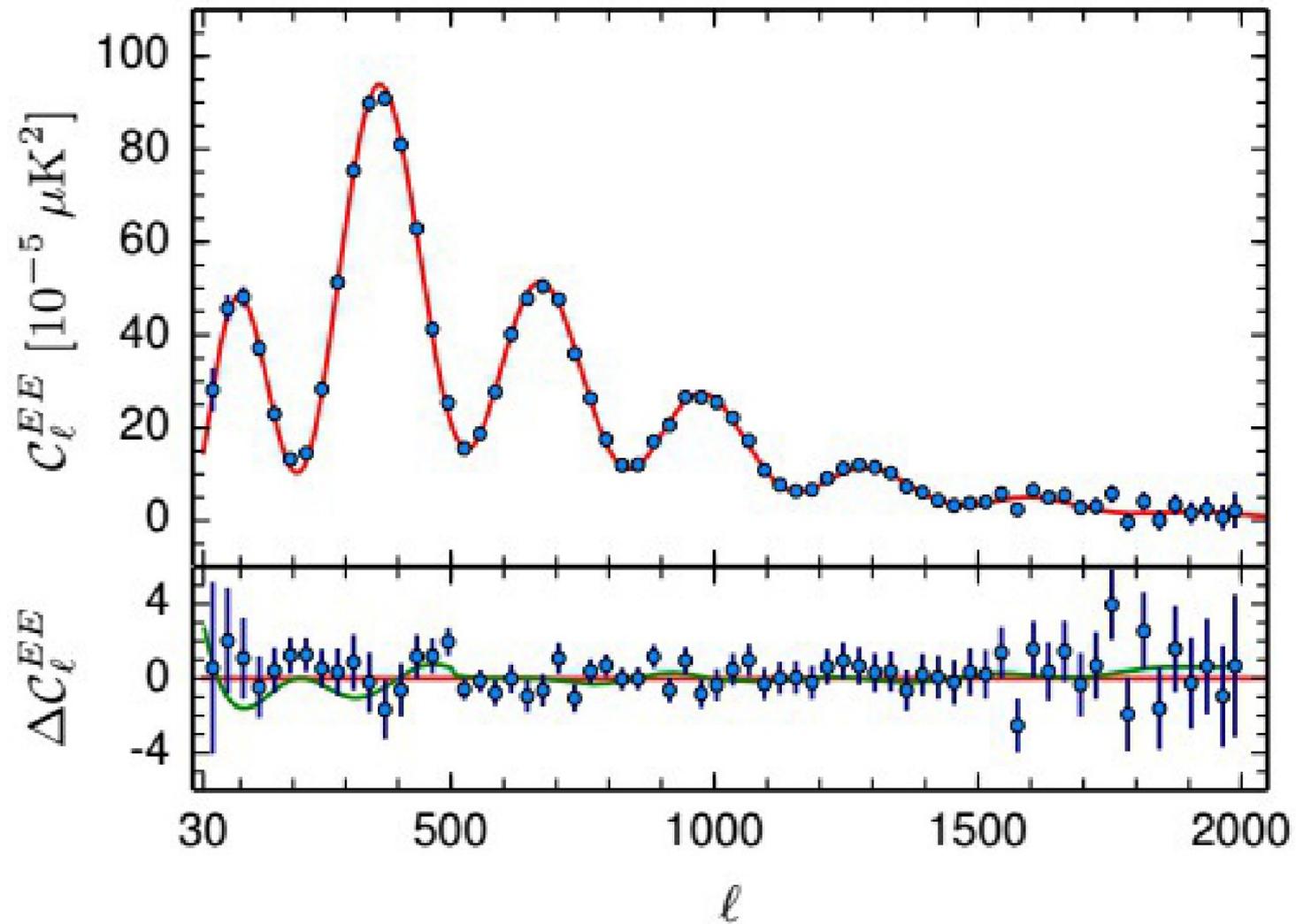
**Planck 2013 release was only temperature ps.**  
**Planck 2015 now includes polarization !**

# Planck 2015 TE angular spectrum



Red Line: Best fit to the TT angular spectrum

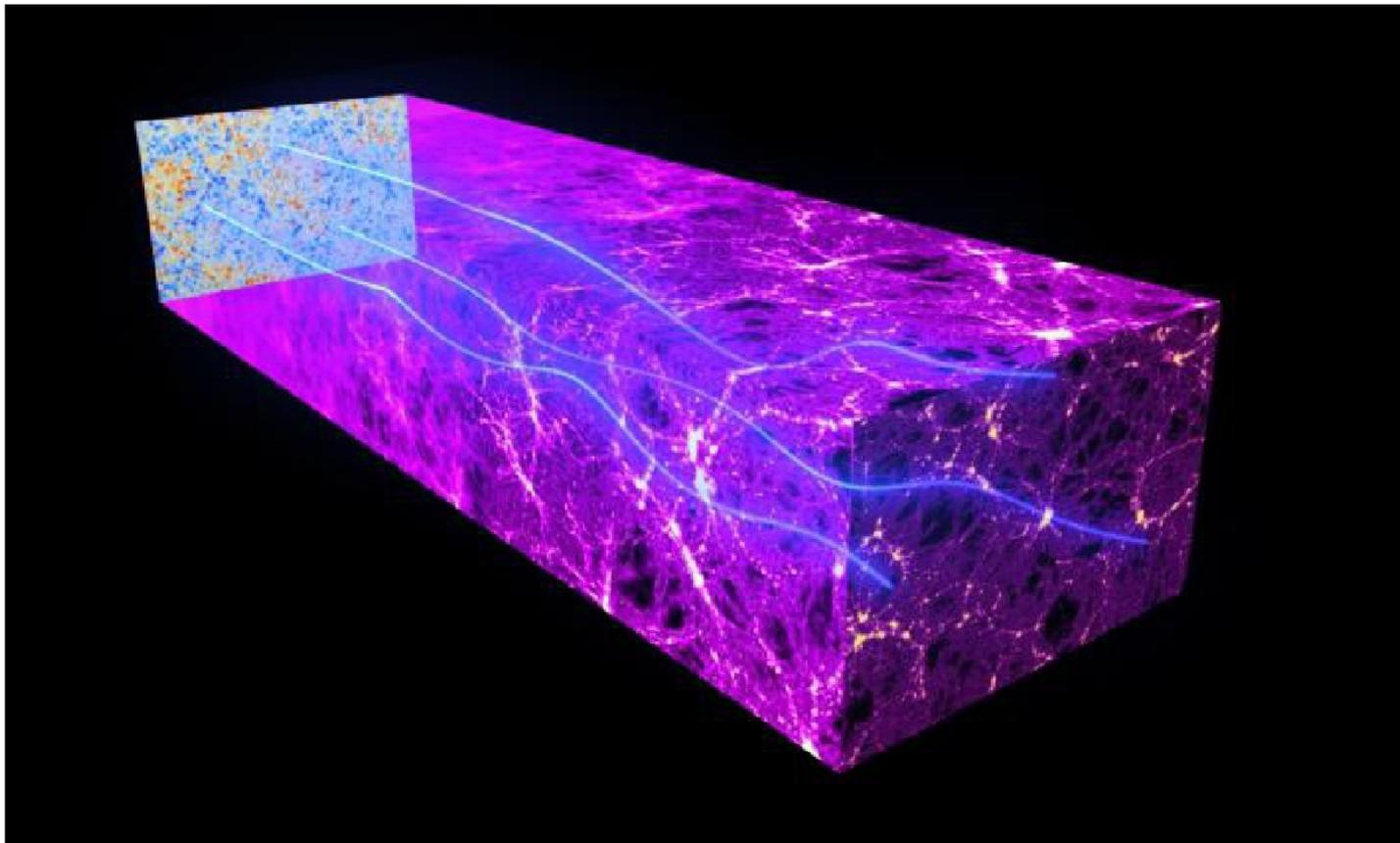
# Planck 2015 EE angular spectrum



Red Line: Best fit to the TT angular spectrum

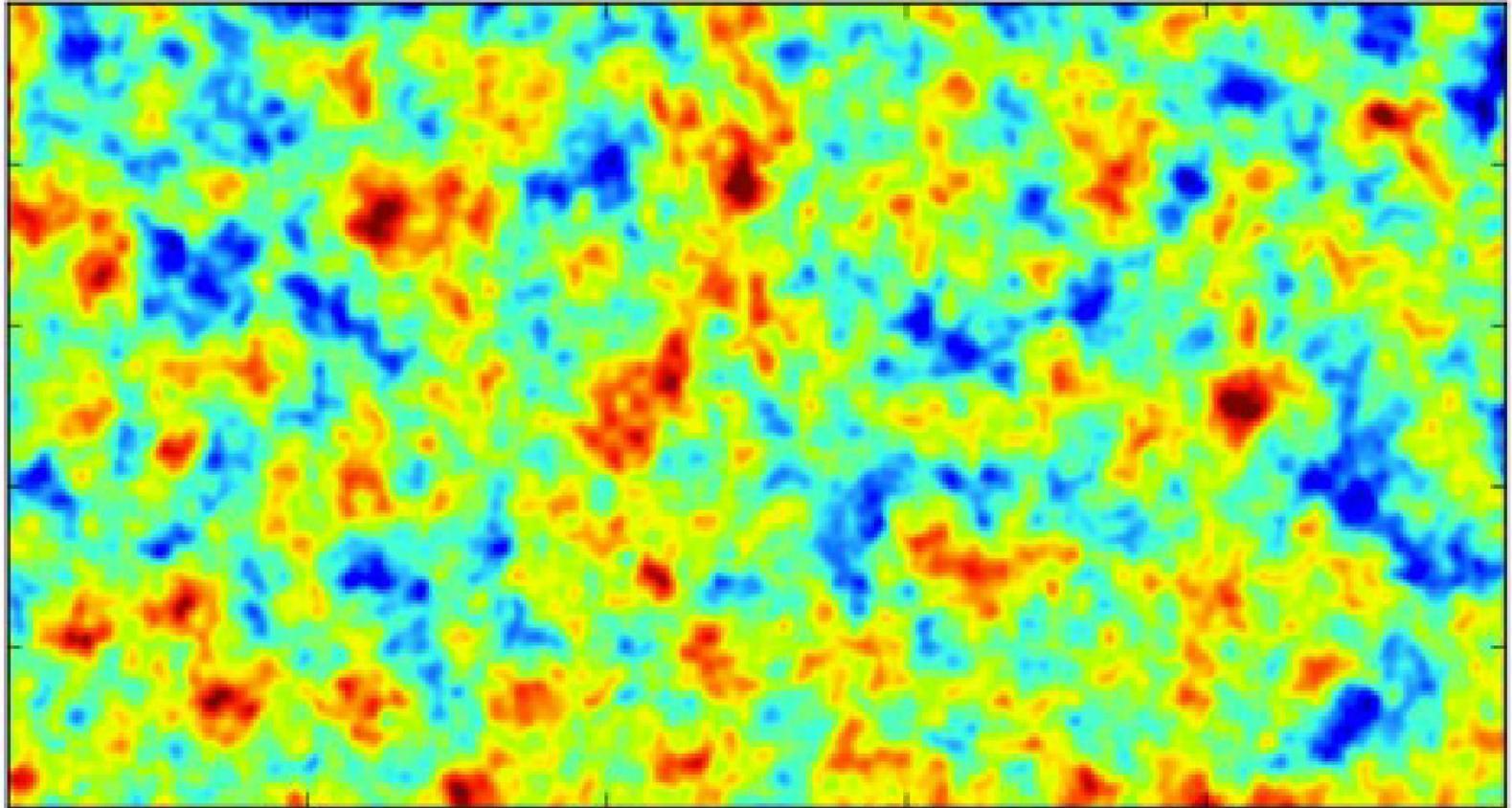
# Gravitational Lensing

The gravitational effects of intervening matter bend the path of CMB light on its way from the early universe to the Planck telescope. This “gravitational lensing” distorts our image of the CMB



# Gravitational Lensing

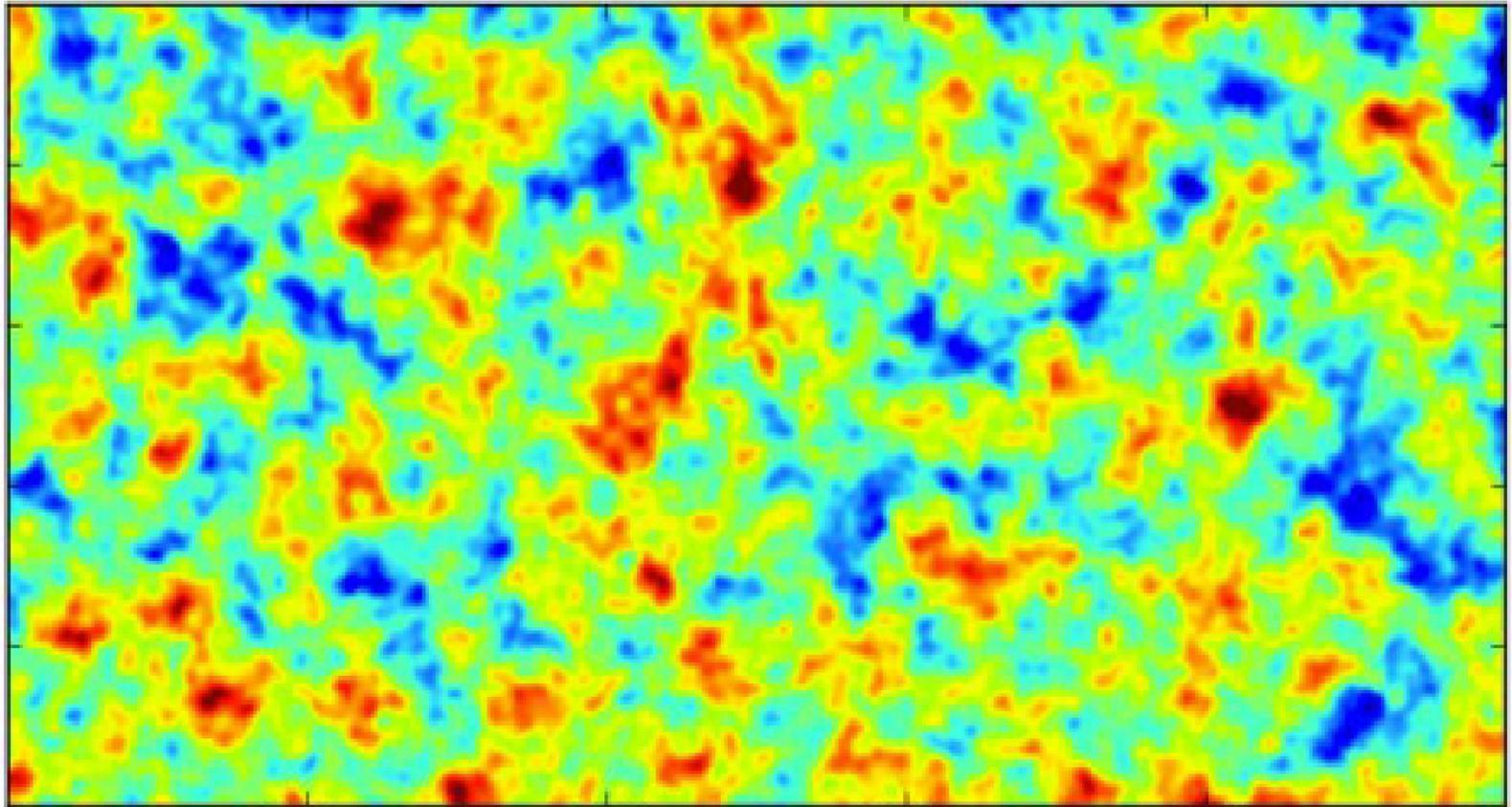
A simulated patch of CMB sky – **before lensing**



←  $10^\circ$  →

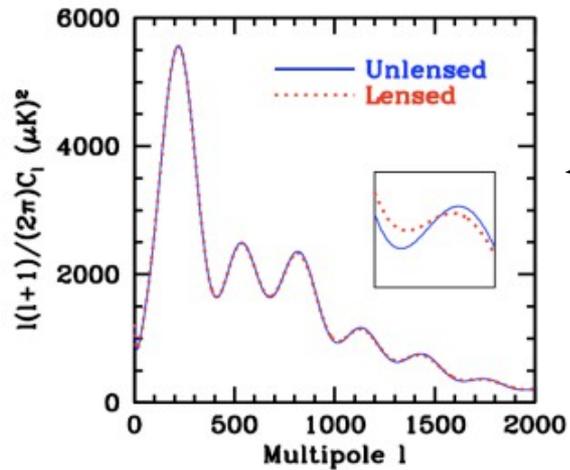
# Gravitational Lensing

A simulated patch of CMB sky – **after lensing**



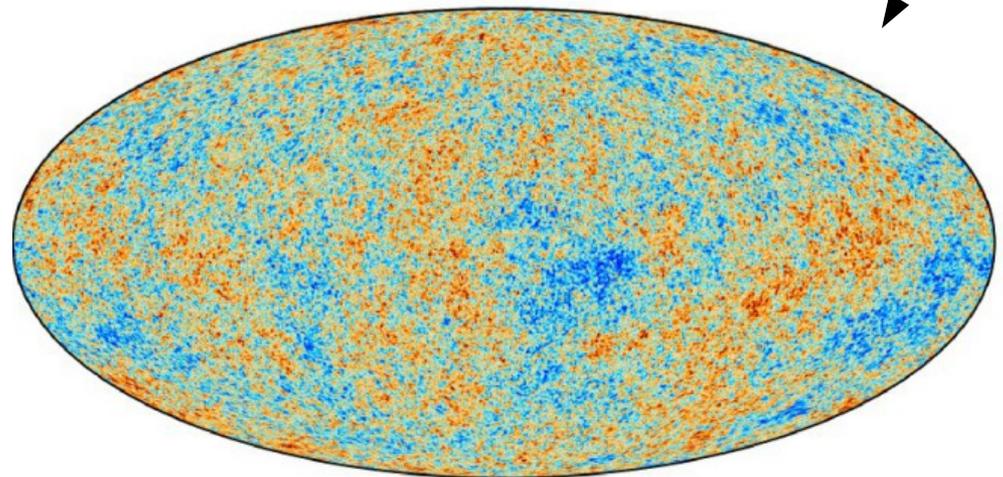
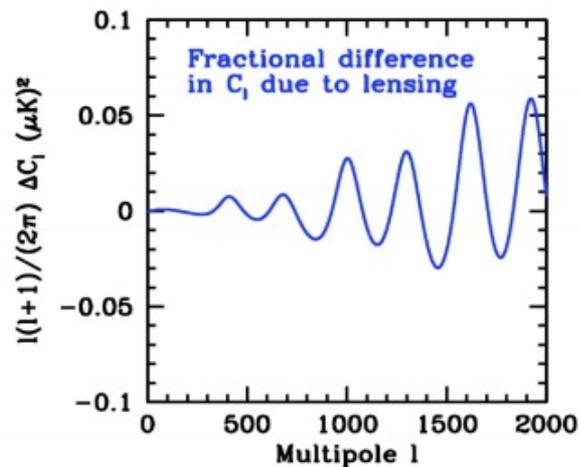
←  $10^\circ$  →

# Extracting Lensing Signal from the CMB

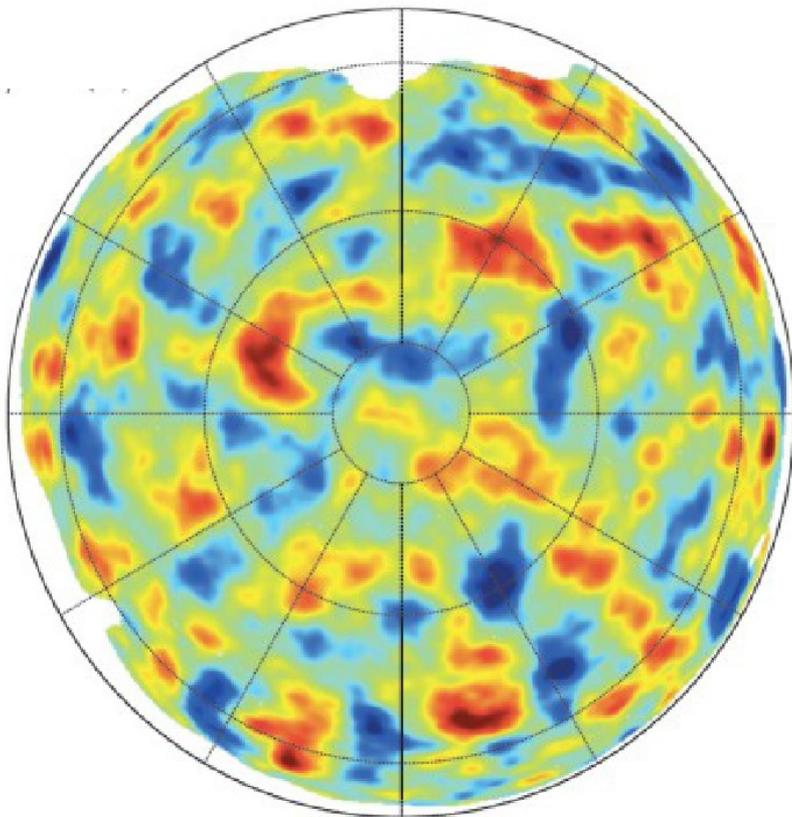


Extraction at Angular Spectrum Level  
Effects on CMB angular spectrum.  
More massive neutrinos  $\rightarrow$  less lensing.

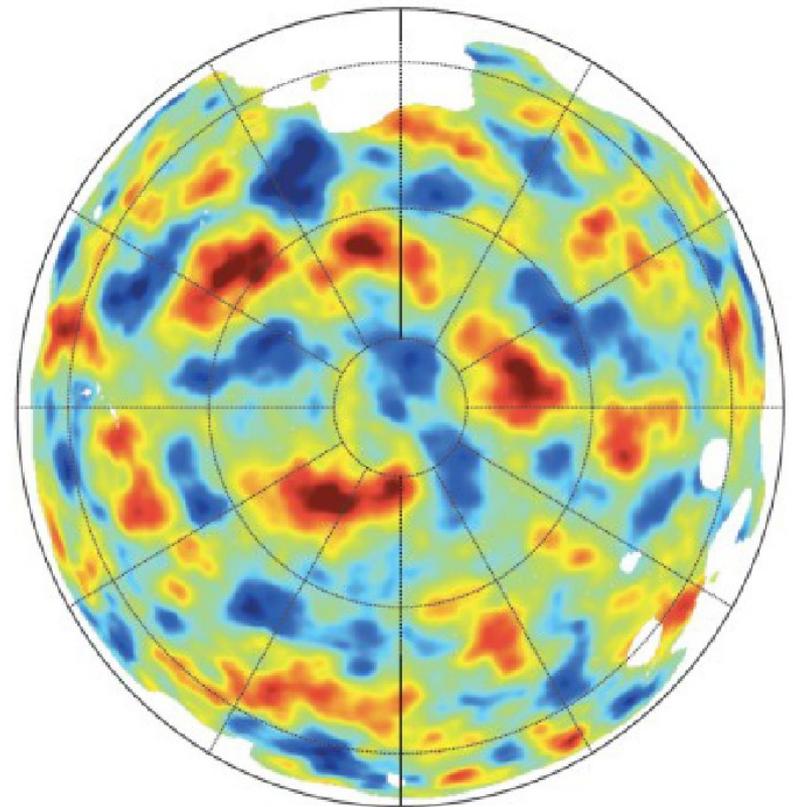
Extraction at map level.  
Lensing introduce a correlation between  $C_l$ 's.  
This creates a non-zero trispectrum (4-point c.f.).



# Planck dark matter distribution through CMB lensing

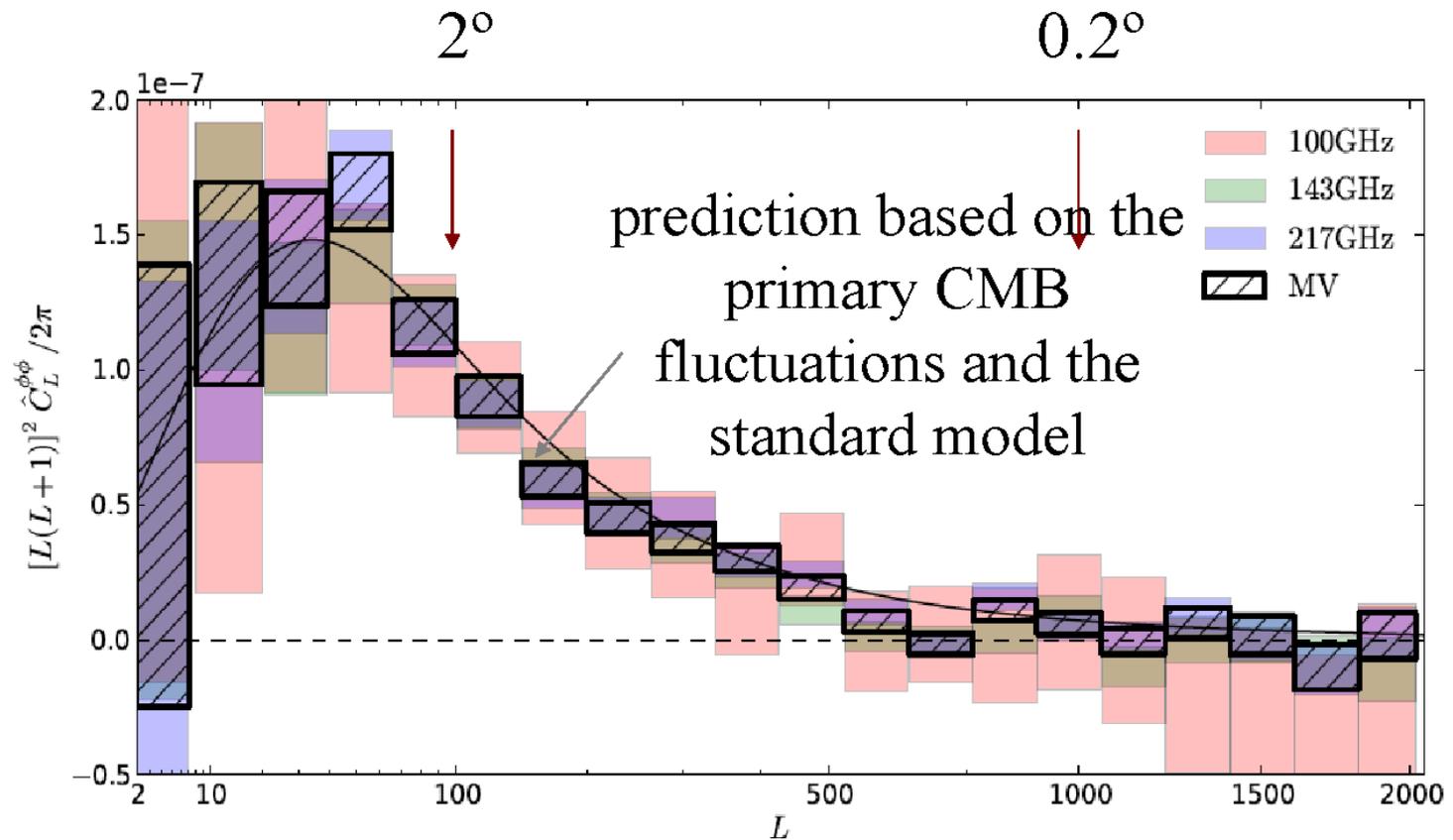


Galactic North



Galactic South

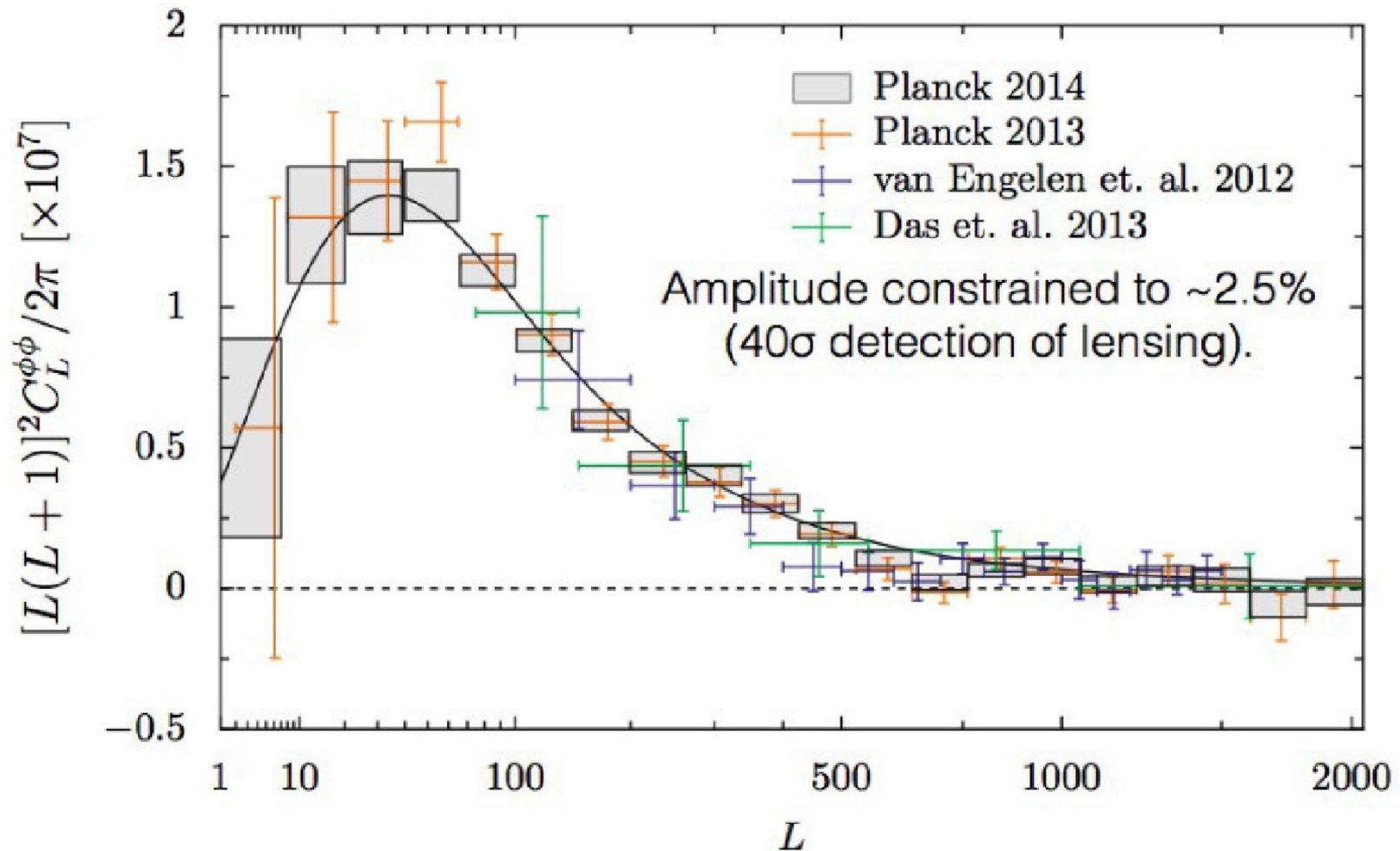
# PLANCK LENSING 2013 Power Spectrum Measured from the Trispectrum (4-point correlation)



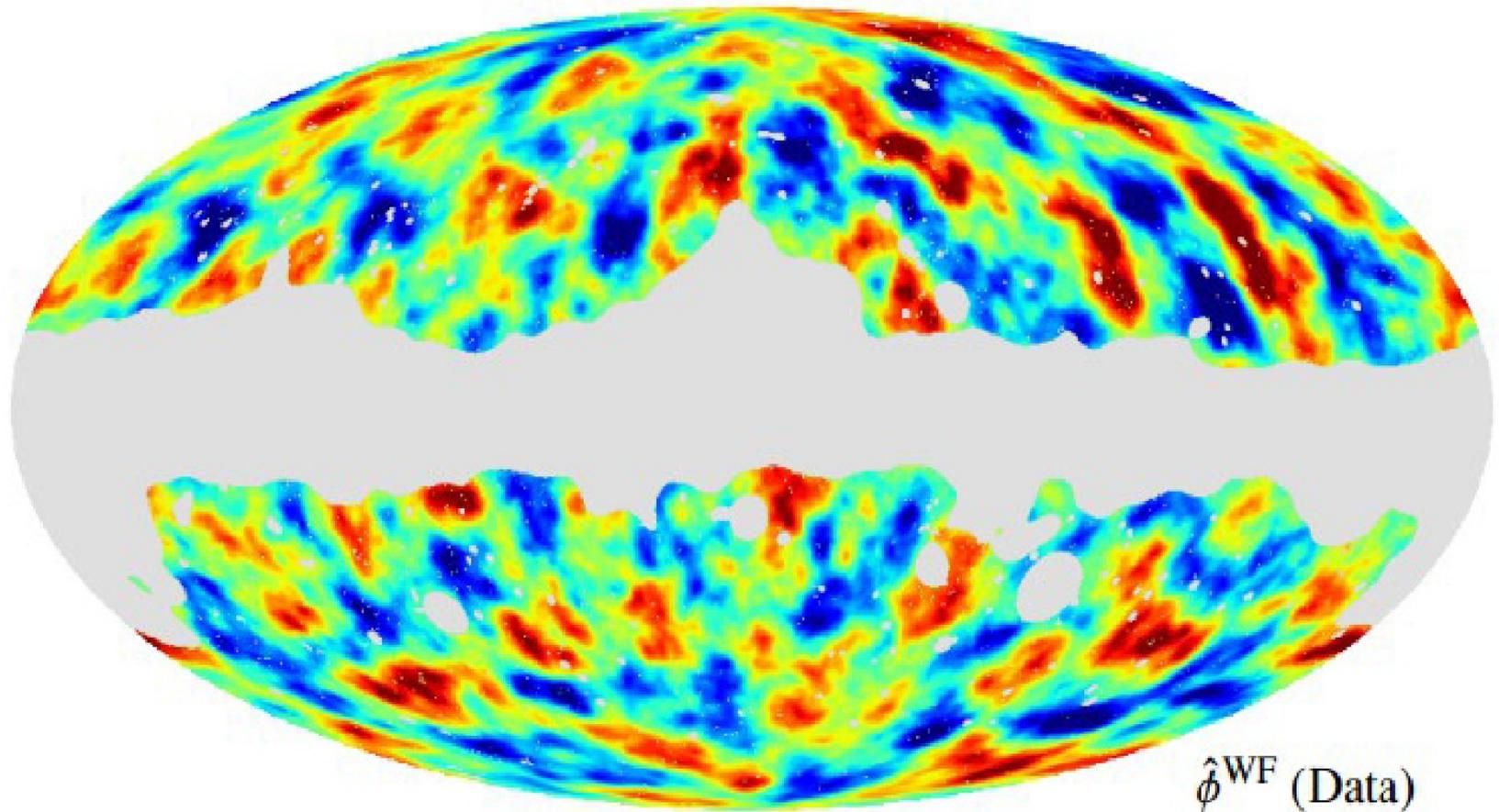
It is a 25 sigma effect!!

This spectrum helps in constraining parameters

# Planck 2015 Lensing Power spectrum



# Planck 2015 Lensing Map



# Constraints on Cosmological Parameters (assuming LCDM)

# The standard cosmological model

- Assumes General Relativity, Inflation, Adiabatic and Scalar Perturbations, flat universe.

- Friedmann-Robertson-Walker (or Friedmann-Lemaitre) metric. Hubble Constant (+1)

$$H_0 = 100 h \text{ km} / \text{s} / \text{Mpc}$$

- 3 Energy components: Baryons, Cold Dark Matter, Cosmological Constant (+3). Flat Universe (-1).

$$\omega_b = \Omega_b h^2 \quad \omega_{CDM} = \Omega_{CDM} h^2$$

- Initial conditions for perturbations given by Inflation: Adiabatic, nearly scale invariant initial power spectrum, only scalar perturbations. Two free parameters (+2): Amplitude and Spectral index.

$$P(k) \approx A_s \left( \frac{k}{k_0} \right)^{n_s}$$

Pivot scale is usually fixed to:

$$k_0 = 0.002 \text{ hMpc}^{-1}$$

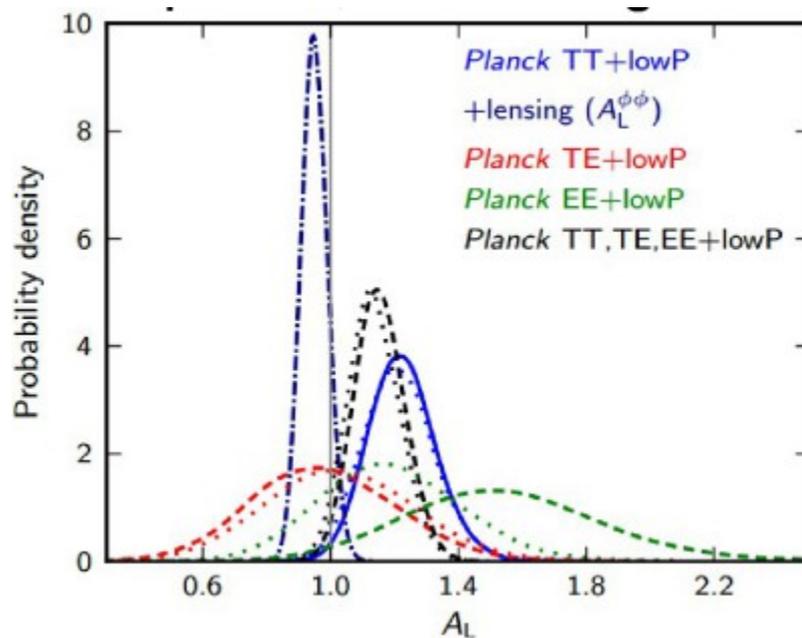
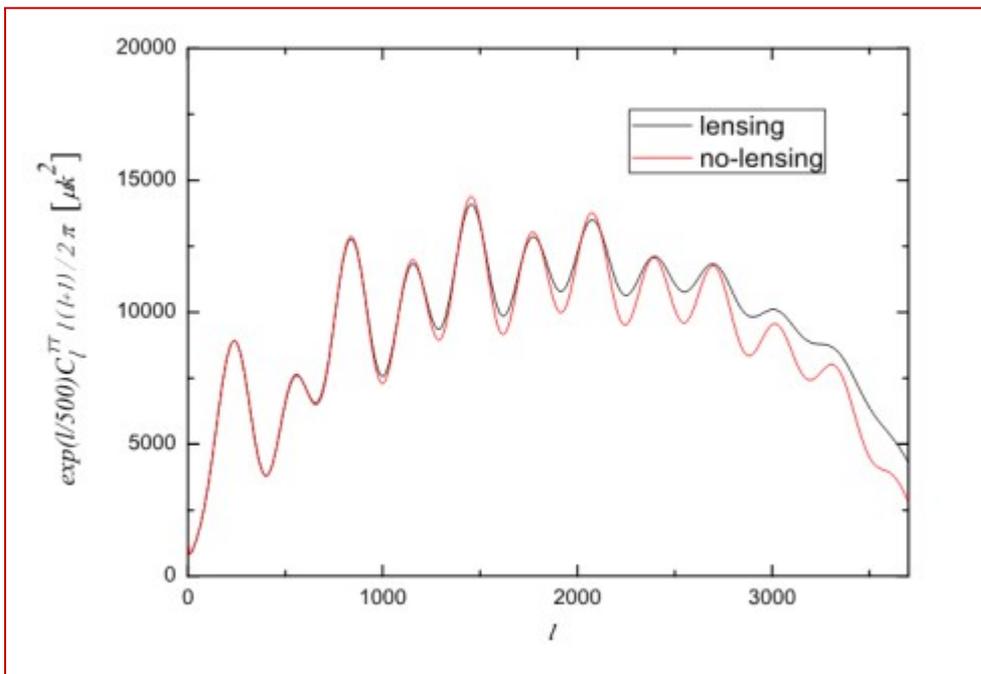
- Late universe reionization characterized with a single parameter(+1) : optical depth  $\tau$  or reionization redshift  $z_r$ .

Total: 1+3-1+2+1= 6 parameters.

Parameter	[1] <i>Planck</i> TT+lowP	[2] <i>Planck</i> TE+lowP	[3] <i>Planck</i> EE+lowP	[4] <i>Planck</i> TT,TE,EE+lowP	([1] - [4])/σ <sub>[1]</sub>
$\Omega_b h^2$ . . . . .	$0.02222 \pm 0.00023$	$0.02228 \pm 0.00025$	$0.0240 \pm 0.0013$	$0.02225 \pm 0.00016$	-0.1
$\Omega_c h^2$ . . . . .	$0.1197 \pm 0.0022$	$0.1187 \pm 0.0021$	$0.1150^{+0.0048}_{-0.0055}$	$0.1198 \pm 0.0015$	0.0
$100\theta_{MC}$ . . . . .	$1.04085 \pm 0.00047$	$1.04094 \pm 0.00051$	$1.03988 \pm 0.00094$	$1.04077 \pm 0.00032$	0.2
$\tau$ . . . . .	$0.078 \pm 0.019$	$0.053 \pm 0.019$	$0.059^{+0.022}_{-0.019}$	$0.079 \pm 0.017$	-0.1
$\ln(10^{10} A_s)$ . . . . .	$3.089 \pm 0.036$	$3.031 \pm 0.041$	$3.066^{+0.046}_{-0.041}$	$3.094 \pm 0.034$	-0.1
$n_s$ . . . . .	$0.9655 \pm 0.0062$	$0.965 \pm 0.012$	$0.973 \pm 0.016$	$0.9645 \pm 0.0049$	0.2
$H_0$ . . . . .	$67.31 \pm 0.96$	$67.73 \pm 0.92$	$70.2 \pm 3.0$	$67.27 \pm 0.66$	0.0
$\Omega_m$ . . . . .	$0.315 \pm 0.013$	$0.300 \pm 0.012$	$0.286^{+0.027}_{-0.038}$	$0.3156 \pm 0.0091$	0.0
$\sigma_8$ . . . . .	$0.829 \pm 0.014$	$0.802 \pm 0.018$	$0.796 \pm 0.024$	$0.831 \pm 0.013$	0.0
$10^9 A_s e^{-2\tau}$ . . . . .	$1.880 \pm 0.014$	$1.865 \pm 0.019$	$1.907 \pm 0.027$	$1.882 \pm 0.012$	-0.1

2015 Planck results are perfectly in agreement with the standard  $\Lambda$ CDM cosmological model.

# Lensing contribution in TT, TE angular spectra, still too high

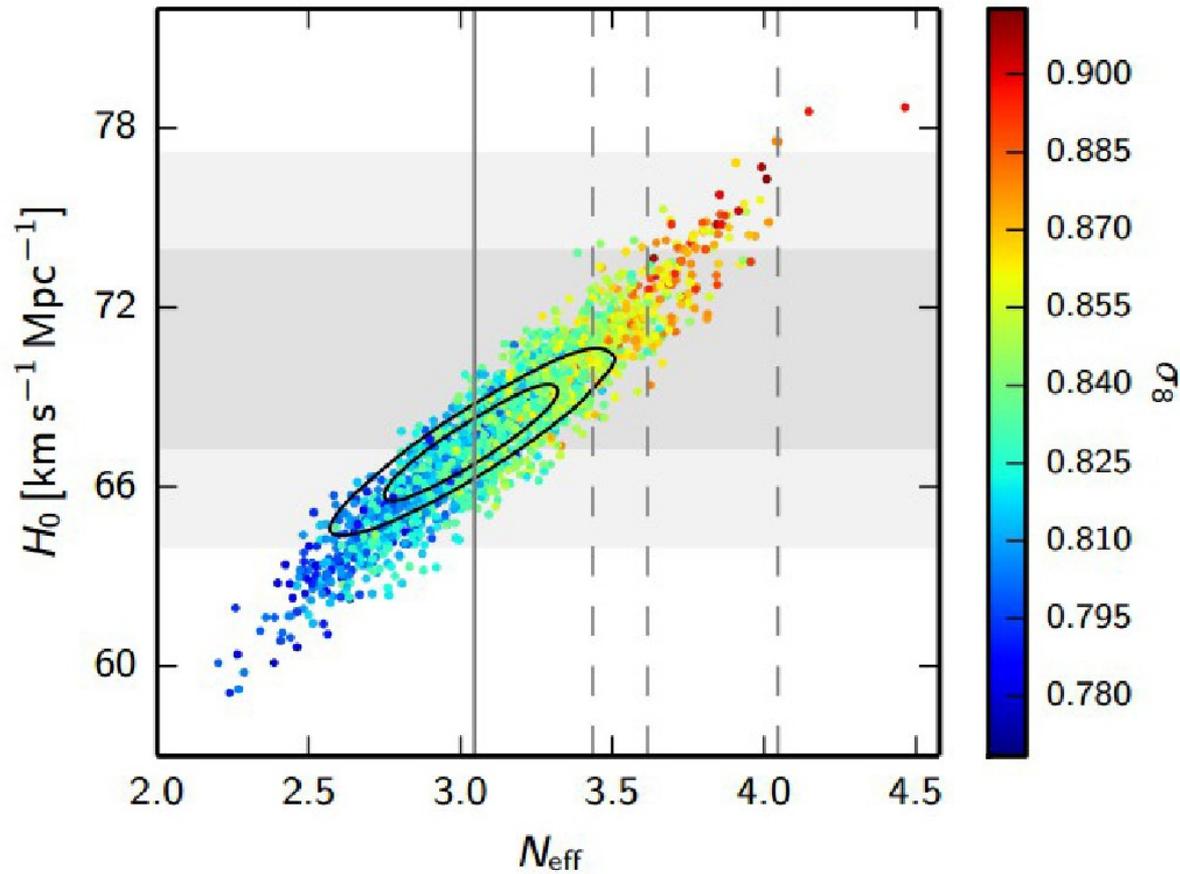


$$A_L = 1.22 \pm 0.10 \quad (68\%, \text{Planck TT+lowP})$$

$$A_L^{\phi\phi} = 0.95 \pm 0.04 \quad (68\%, \text{Planck TT+lowP+lensing}).$$

Is clearly a systematic in TT data since we don't see this in the TTTT lensing

Neutrino cosmology...



$$N_{\text{eff}} = 3.13 \pm 0.32 \quad \text{Planck TT+lowP};$$

$$N_{\text{eff}} = 3.15 \pm 0.23 \quad \text{Planck TT+lowP+BAO};$$

$$N_{\text{eff}} = 2.99 \pm 0.20 \quad \text{Planck TT, TE, EE+lowP};$$

$$N_{\text{eff}} = 3.04 \pm 0.18 \quad \text{Planck TT, TE, EE+lowP+BAO}.$$

# Neutrino Background Anisotropies

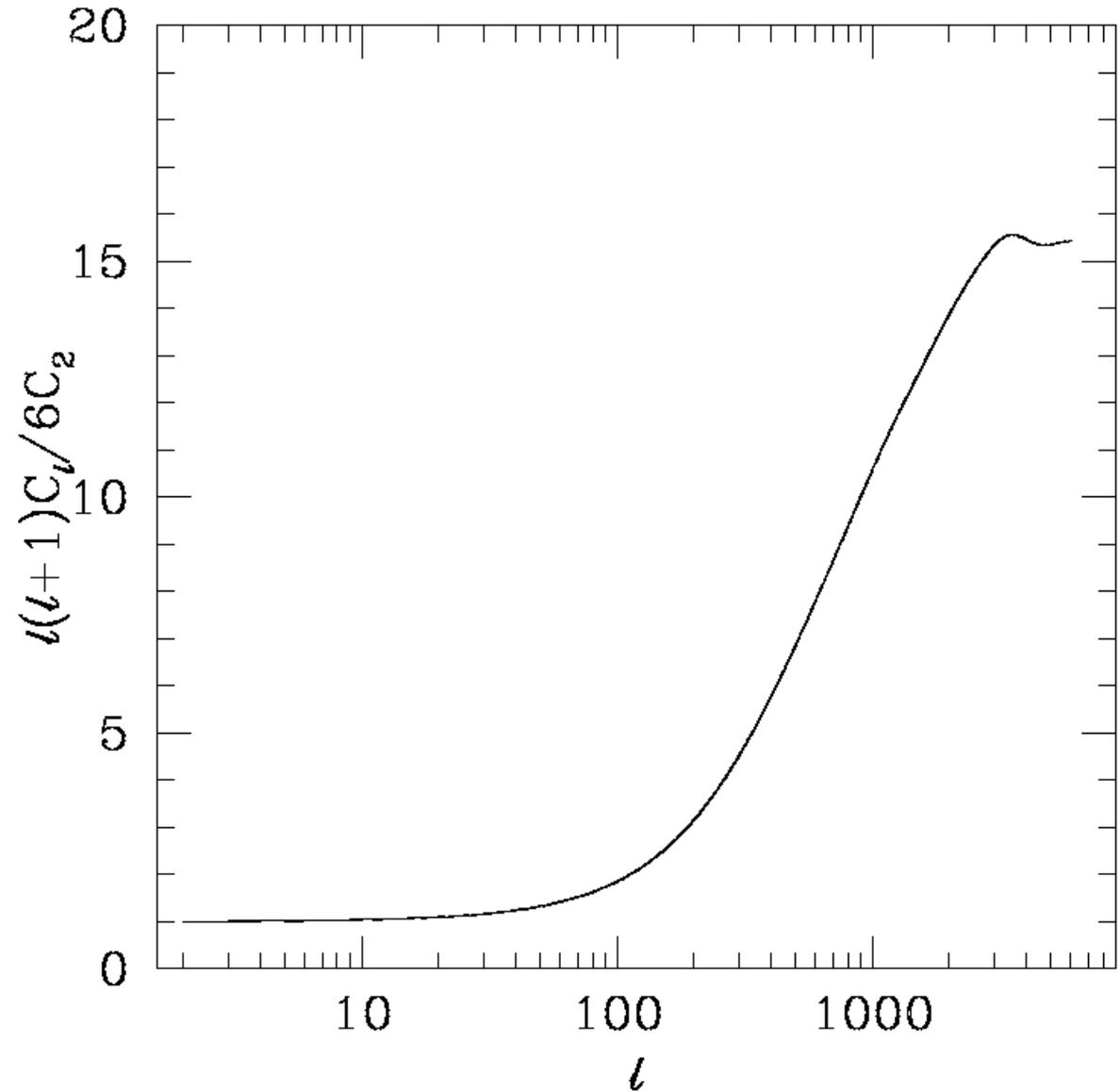
Massless neutrinos, like photons, have anisotropies which follow a Liouville differential equation:

$$\frac{\partial \mathcal{I}}{\partial t} + \frac{\gamma_i}{a} \frac{\partial \mathcal{I}}{\partial x^i} - 2\dot{h}_{jk} \gamma_j \gamma_k = 0$$

As in the case of photons, these anisotropies can be computed integrating a hierarchy of differential equations.

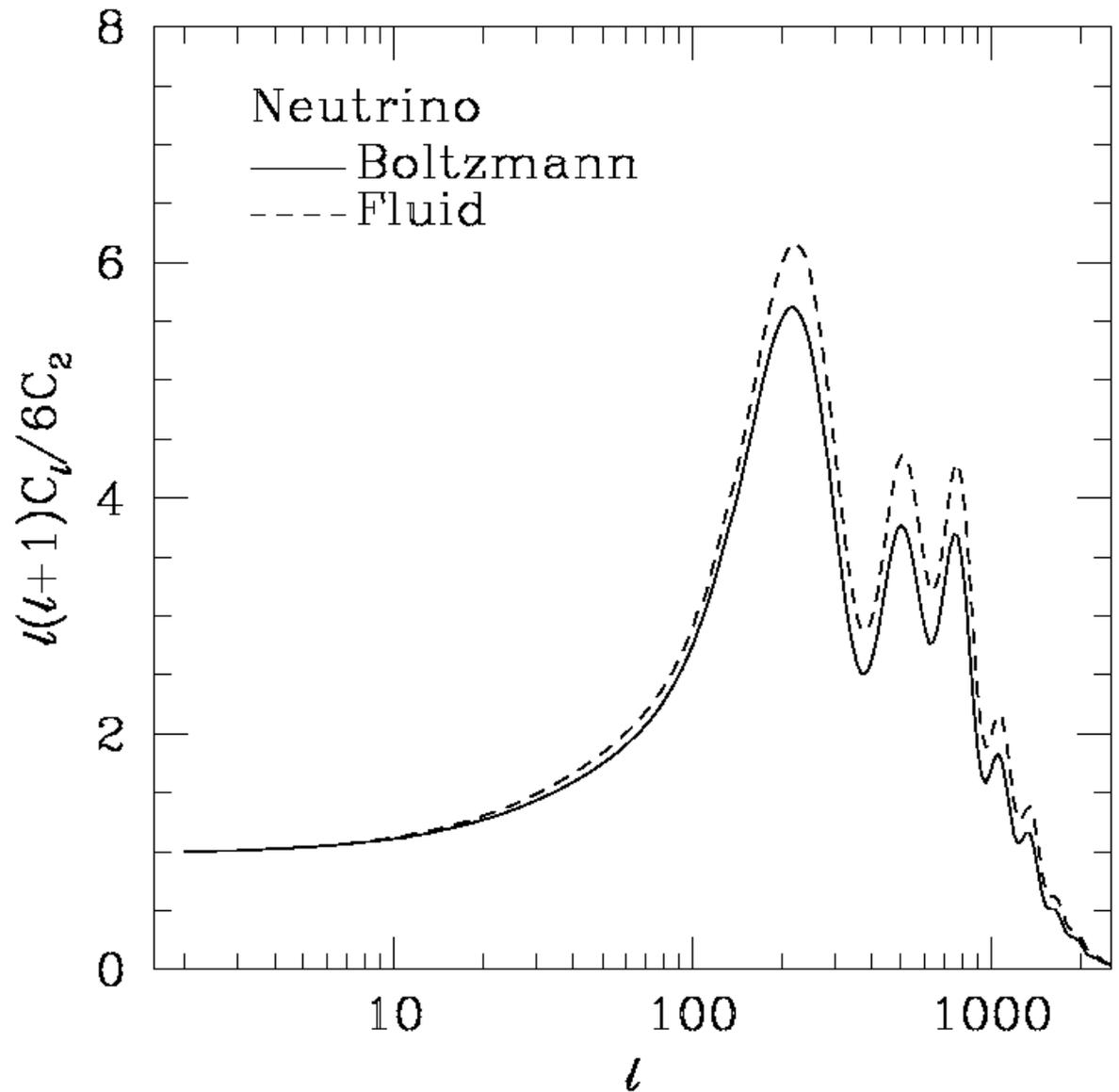
# Neutrino Background Anisotropies

Can we see them ?

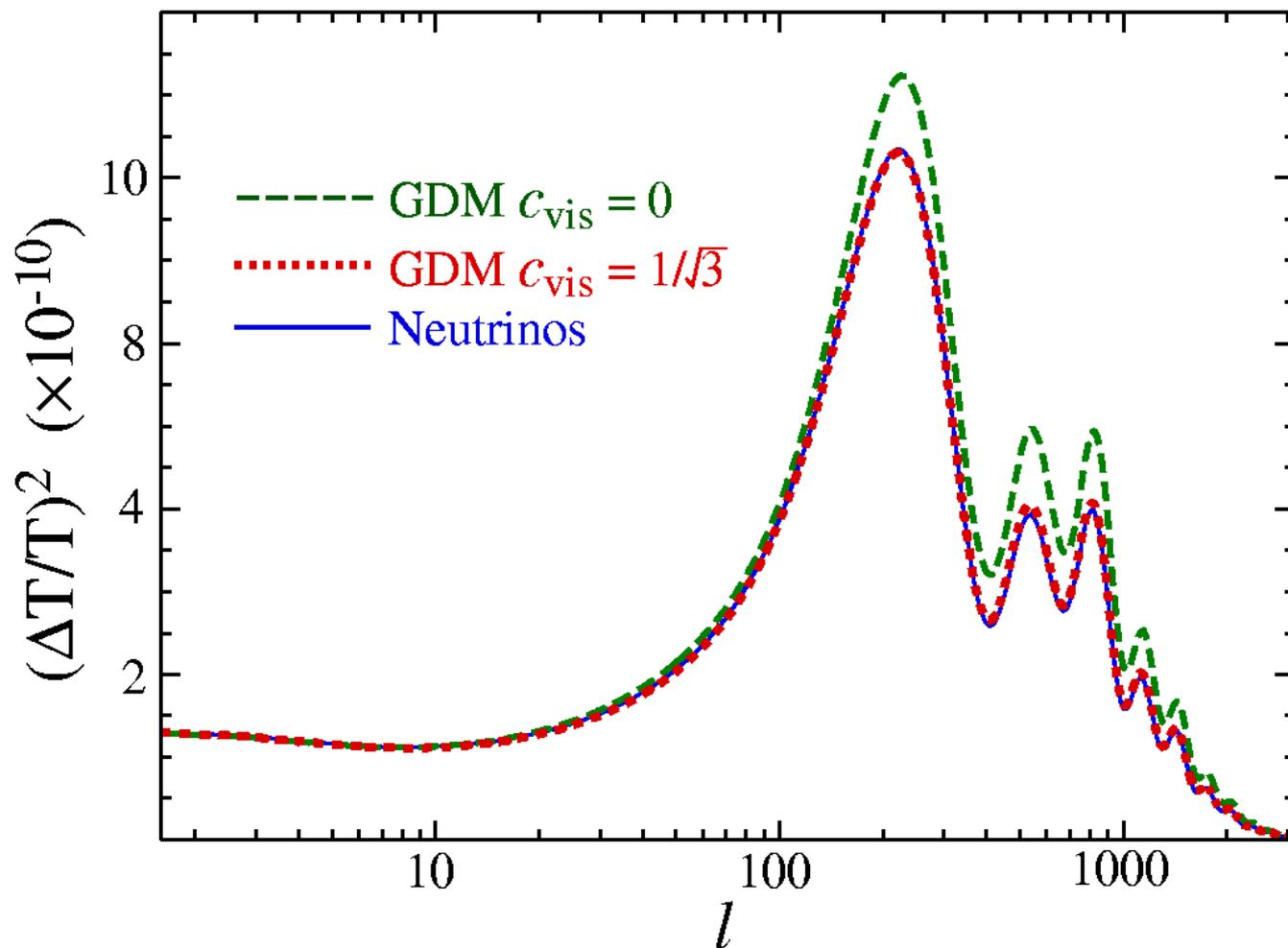


Hu et al., astro-ph/9505043

Not directly!  
But we can see the effects on the CMB angular spectrum!  
CMB photons see the NB anisotropies through gravity.

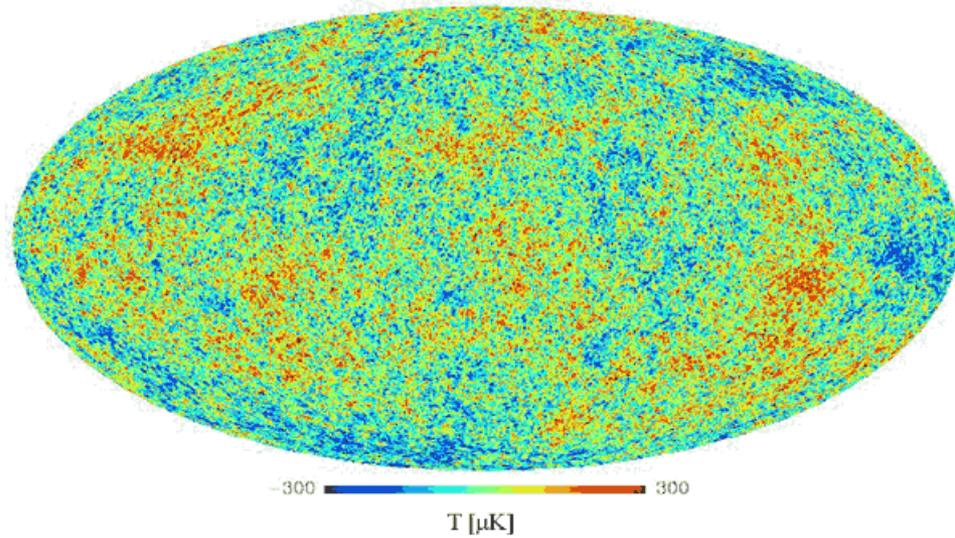


The Neutrino anisotropies can be parameterized through the “speed viscosity”  $c_{\text{vis}}$ , which controls the relationship between velocity/metric shear and anisotropic stress in the NB.



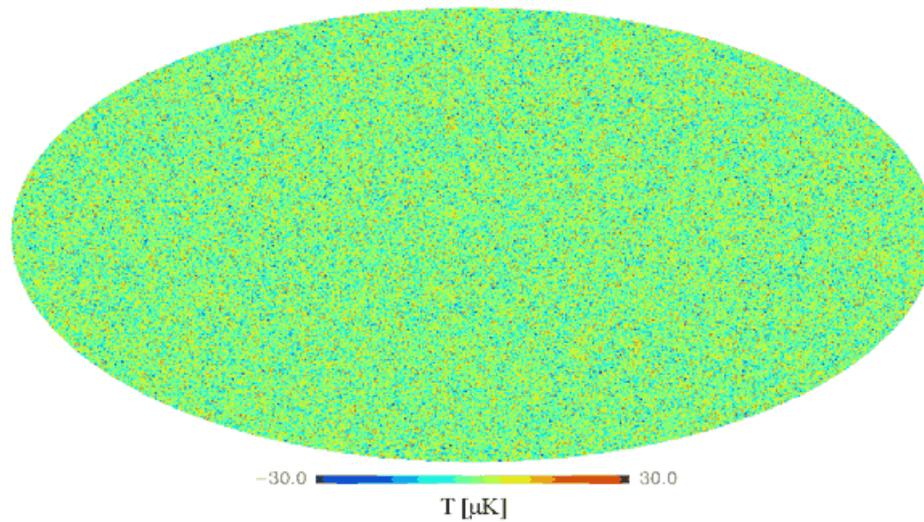
Hu, Eisenstein, Tegmark and White, 1999

### Standard Model



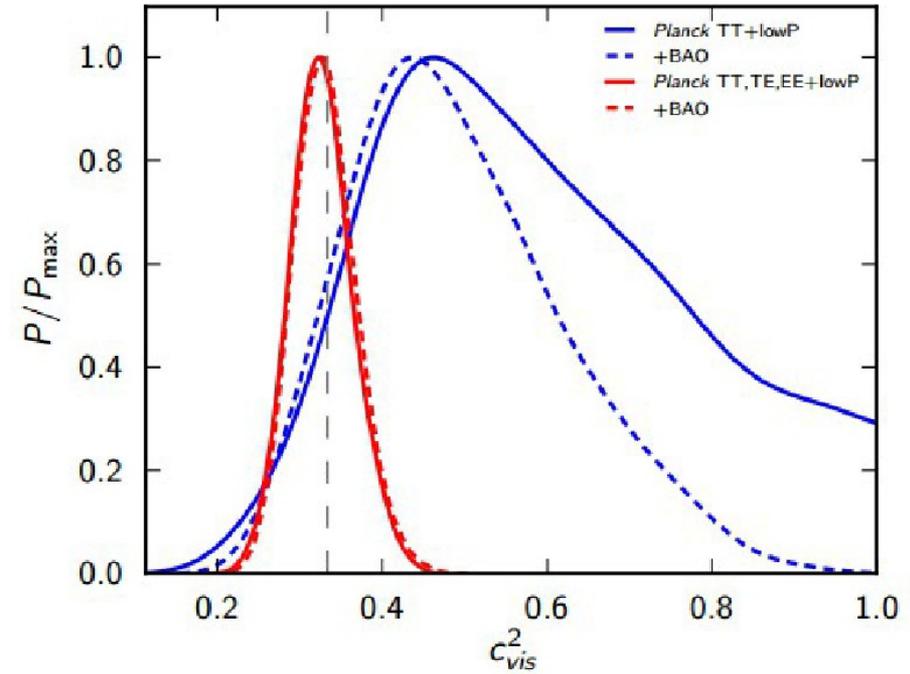
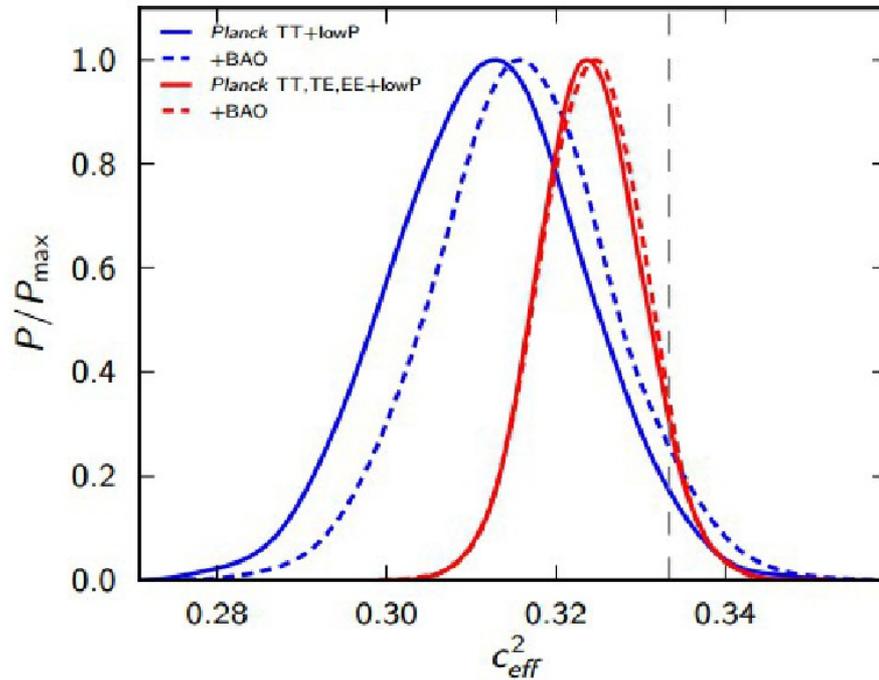
Simulation showing the distribution on the sky of temperature fluctuations in the Cosmic Microwave Background with neutrinos as in the Standard Model.

### Contribution from neutrino ripples



The net effect on the Microwave Background of the presence of neutrino ripples, interpreted as the signature of the existence of neutrino fluctuations as predicted in the Standard Model.

# Perturbations in the neutrino background



$$\left. \begin{aligned} c_{\text{eff}}^2 &= 0.3242 \pm 0.0059 \\ c_{\text{vis}}^2 &= 0.331 \pm 0.037 \end{aligned} \right\} \text{Planck TT,TE,EE+lowP+BAO.}$$

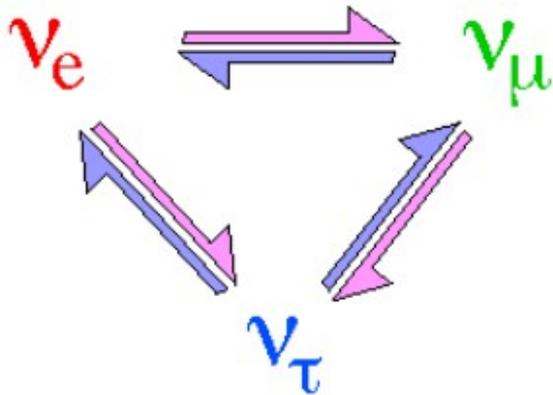
Planck collaboration 2015

$c_{\text{eff}}^2$  is the neutrino sound speed in its own reference frame and  $c_{\text{vis}}^2$  parameterizes the anisotropic stress. For standard non-interacting massless neutrinos  $c_{\text{eff}}^2 = c_{\text{vis}}^2 = 1/3$ . Any deviation from the expected values could provide a hint of non-standard physics in the neutrino sector. **A vanishing value of  $c_{\text{vis}}^2$** , which means that the CNB is a perfect fluid and might imply an interaction between neutrinos and other species, **is excluded at about  $9\sigma$**  when Planck polarization data are included!!!

# Neutrino Masses

# We know that flavour neutrino oscillations exists

From present evidences of oscillations from experiments measuring atmospheric, solar, reactor and accelerator neutrinos

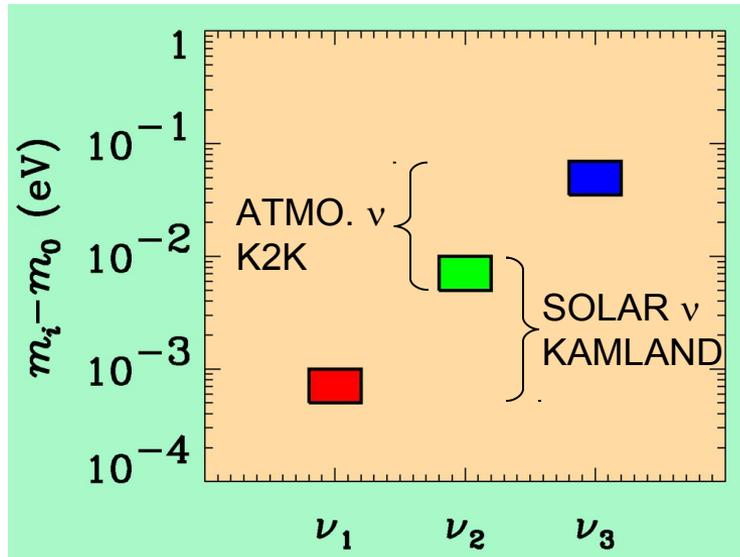


$$(e, \mu, \tau) \leftrightarrow (\nu_1, \nu_2, \nu_3)$$

$$\nu_{\alpha L} = \sum_{i=1}^3 U_{\alpha i} \nu_{iL}, \quad (\alpha = e, \mu, \tau)$$

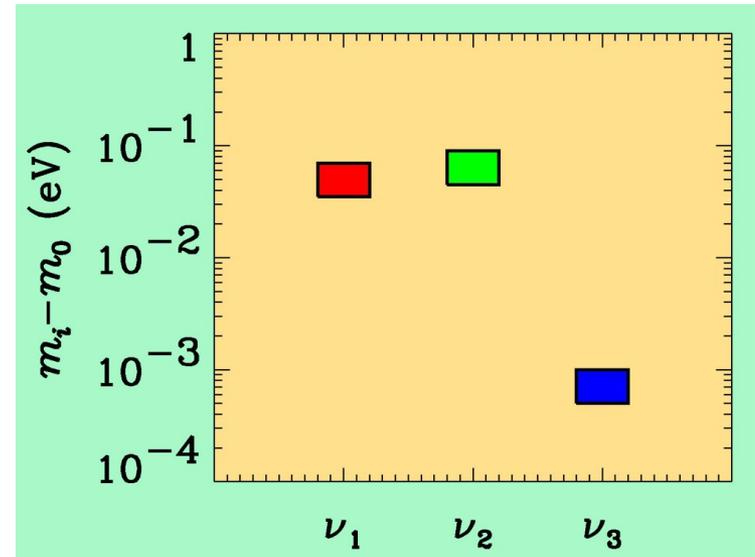
$$\begin{array}{c}
 \nu_e \\
 \nu_\mu \\
 \nu_\tau
 \end{array}
 \begin{bmatrix}
 & \nu_1 & \nu_2 & \nu_3 \\
 & c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\
 -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\
 s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13}
 \end{bmatrix}
 \times \text{diag}(e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1) .$$

If neutrino masses are hierarchical then oscillation experiments do not give information on the absolute value of neutrino masses



Normal hierarchy

$$m_3 > m_2 > m_1$$



Inverted hierarchy

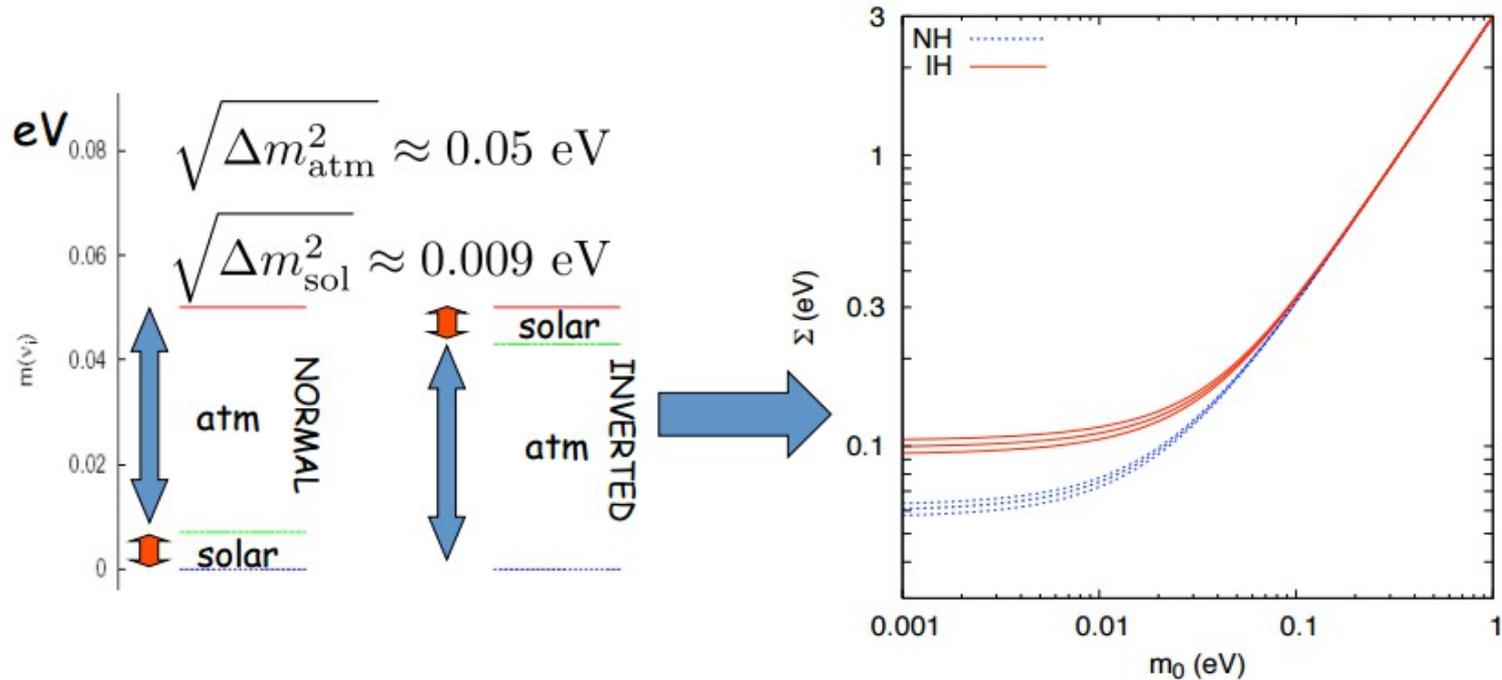
$$m_2 > m_1 > m_3$$

Moreover neutrino masses can also be degenerate

$$m_1, m_2, m_3 \gg \delta m_{\text{atmospheric}}$$

# Neutrino Absolute Mass Scale

Oscillation data do not fix the absolute mass scale.



$$0.06(0.1) \text{ eV} \lesssim \sum_i m_i \lesssim 6 \text{ eV}$$

A large range of values (about 2 orders of magnitude) is still allowed by current experiments !

# Energy density from neutrinos (after decoupling)

If they are relativistic:

$$\rho_\nu = \frac{7\pi^2}{120} N_{\text{eff}} \left(\frac{4}{11}\right)^{4/3} T_\gamma^4$$

$$\Omega_\nu h^2 = \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}} \Omega_\gamma h^2$$

$$\Omega_\nu h^2 = 0.58 \times 10^{-5} N_{\text{eff}} \\ (T_\gamma = 2.726 \text{ K})$$

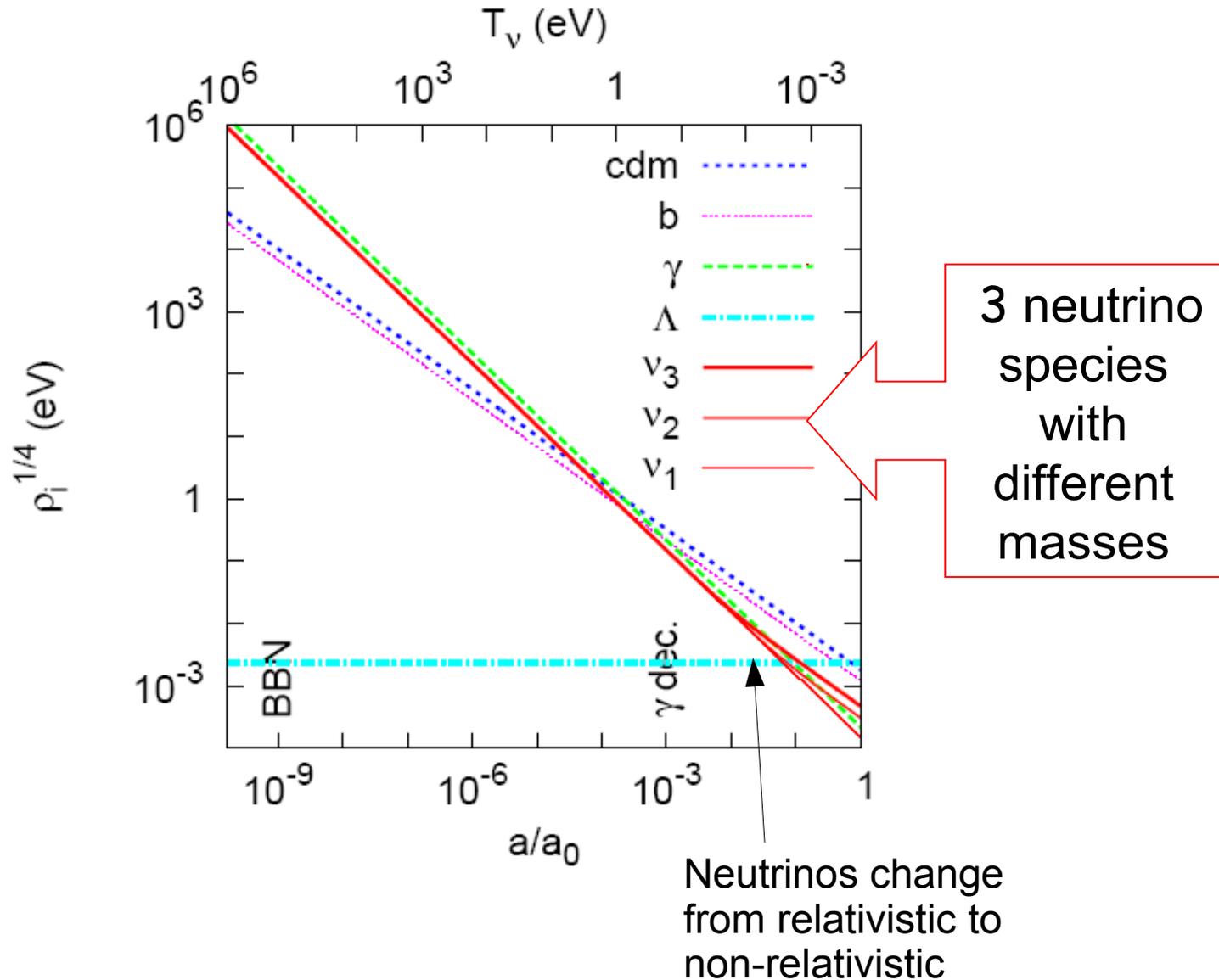
If they are not relativistic:

$$\rho_\nu = \sum n_i m_i \\ n_i = \frac{6\zeta(3)}{11\pi^2} T_\gamma^3$$

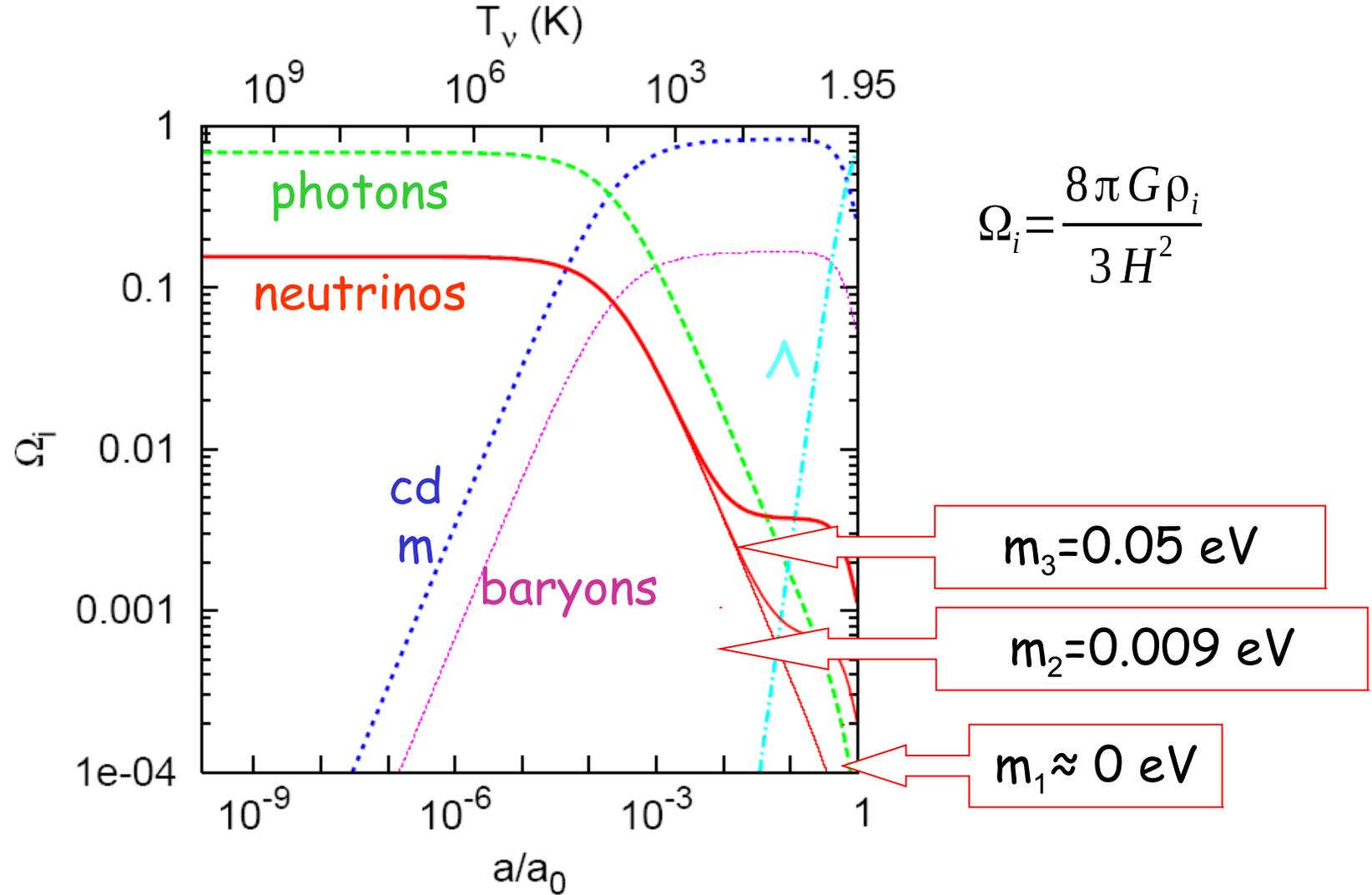


$$\Omega_\nu h^2 = \frac{\sum m_i}{93.2 \text{ eV}} \\ (T_\gamma = 2.726 \text{ K})$$

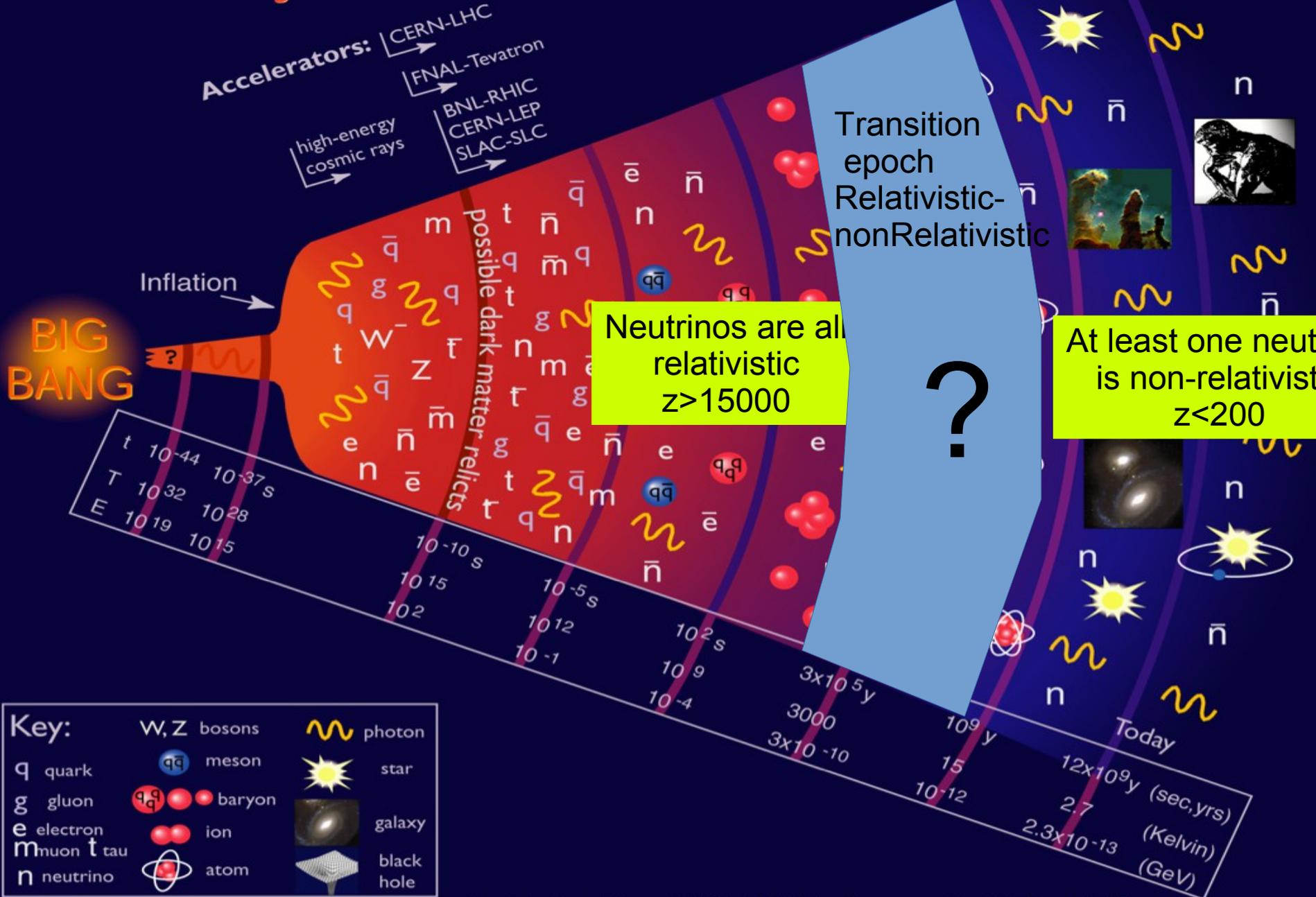
# Evolution of the background densities: 1 MeV $\rightarrow$ now

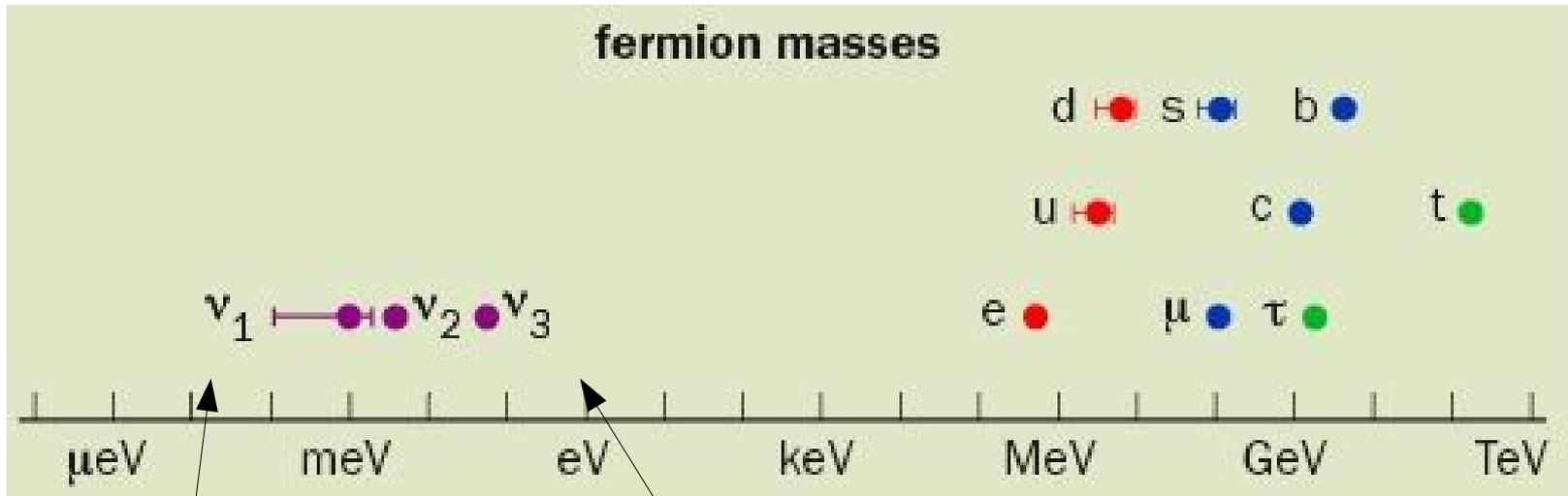


# Evolution of the background densities: 1 MeV $\rightarrow$ now



# History of the Universe





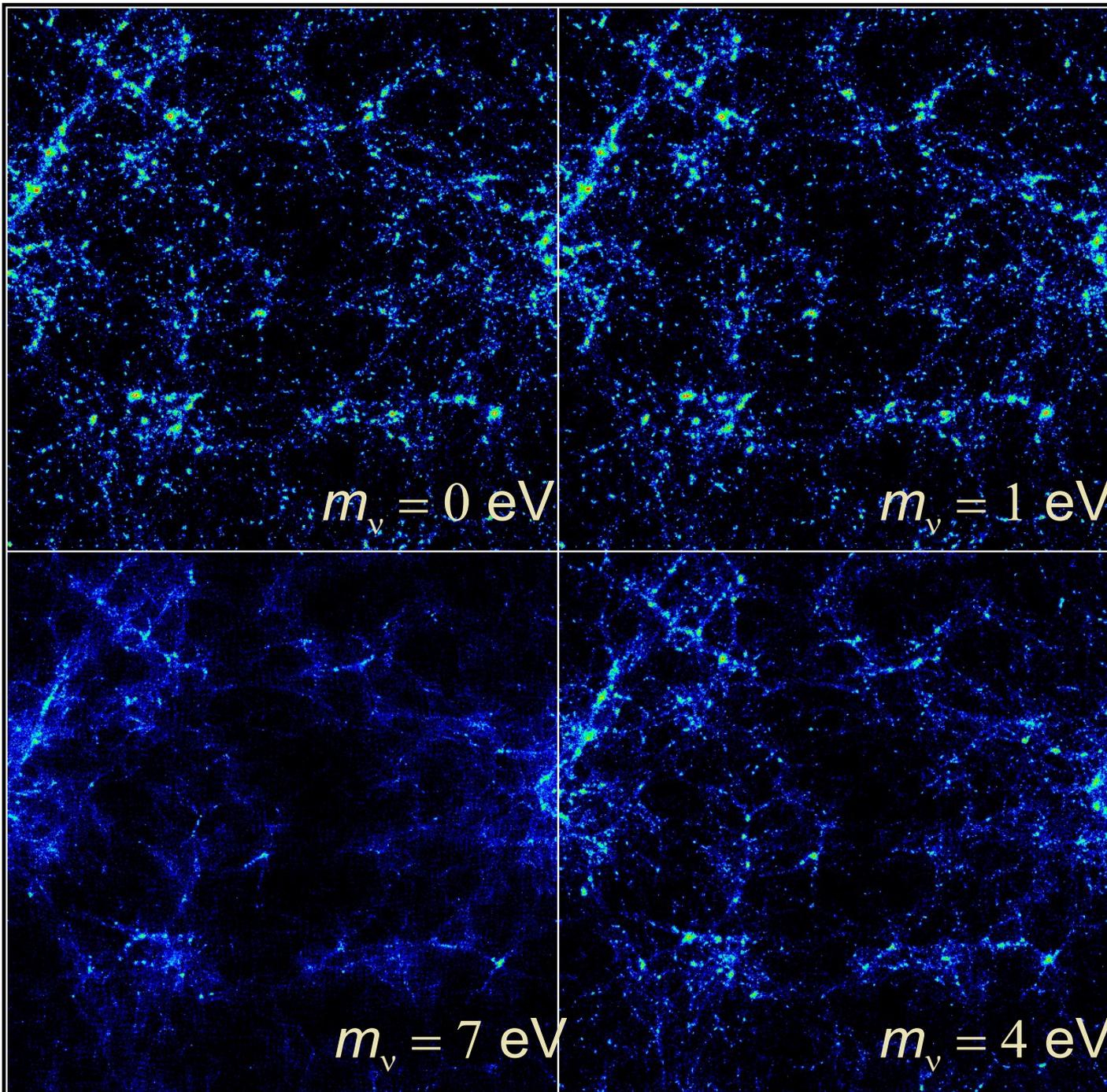
This mass range is very good for cosmologists !

# Massive Neutrinos affect large scale structure !

- ◆ We can relate the neutrino abundance in the universe to the total mass:

$$\Omega_{\nu} h^2 = \frac{\Sigma m_{\nu}}{93.6 eV}$$

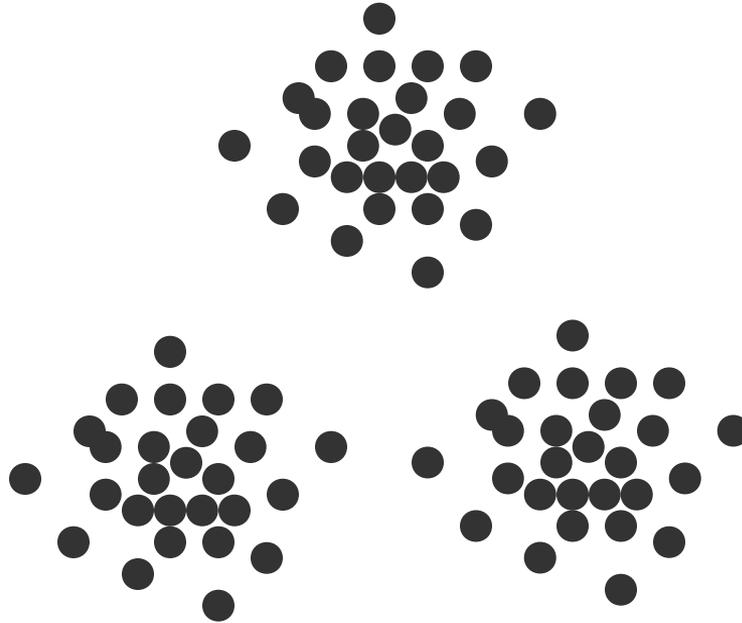
- ◆ Less clustering in universe with massive neutrinos.



Ma '96

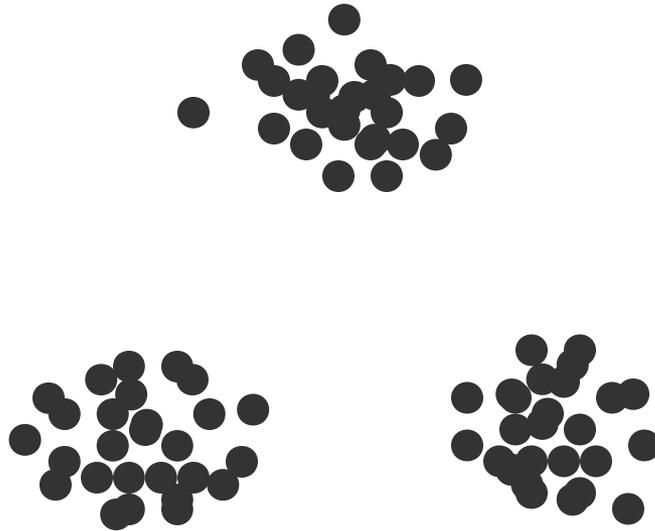
# Structure formation after equality

baryons and  
CDM  
experience  
gravitational  
clustering



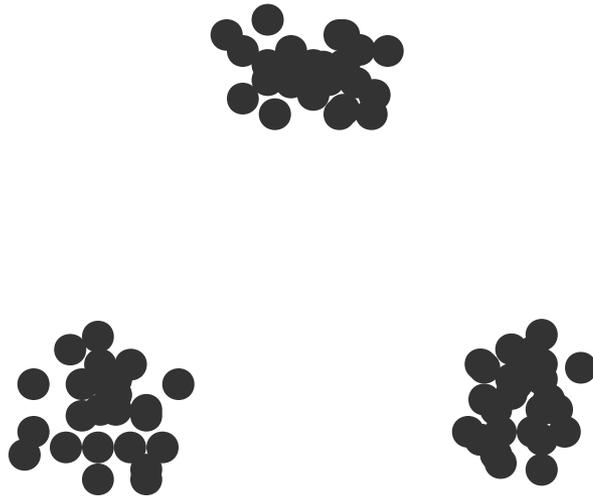
# Structure formation after equality

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experience  
gravitational  
clustering



# Structure formation after equality

baryons and  
CDM  
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gravitational  
clustering



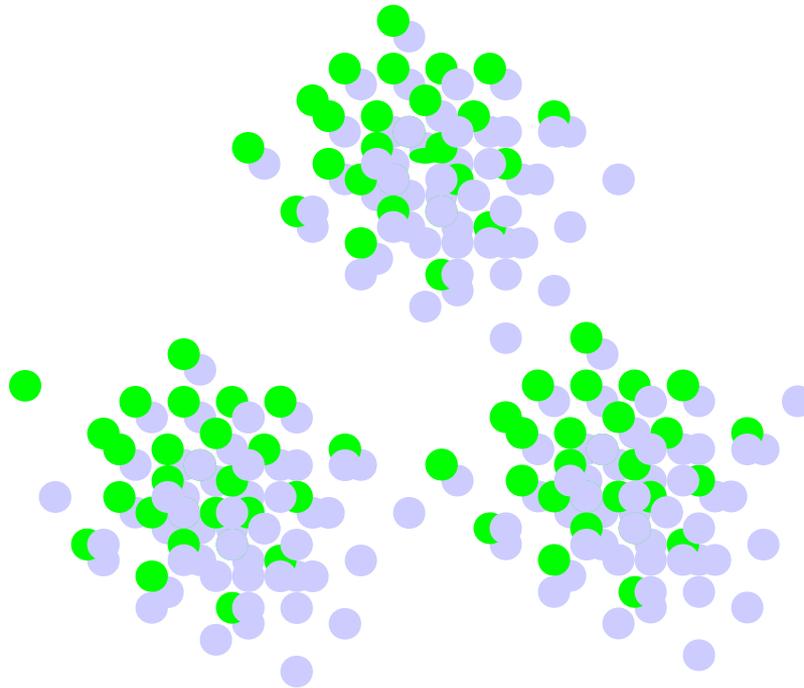
growth of  $\delta\rho/\rho(k,t)$  fixed by  
« gravity vs. expansion » balance

$$\delta\rho/\rho \propto a$$

# Structure formation after equality

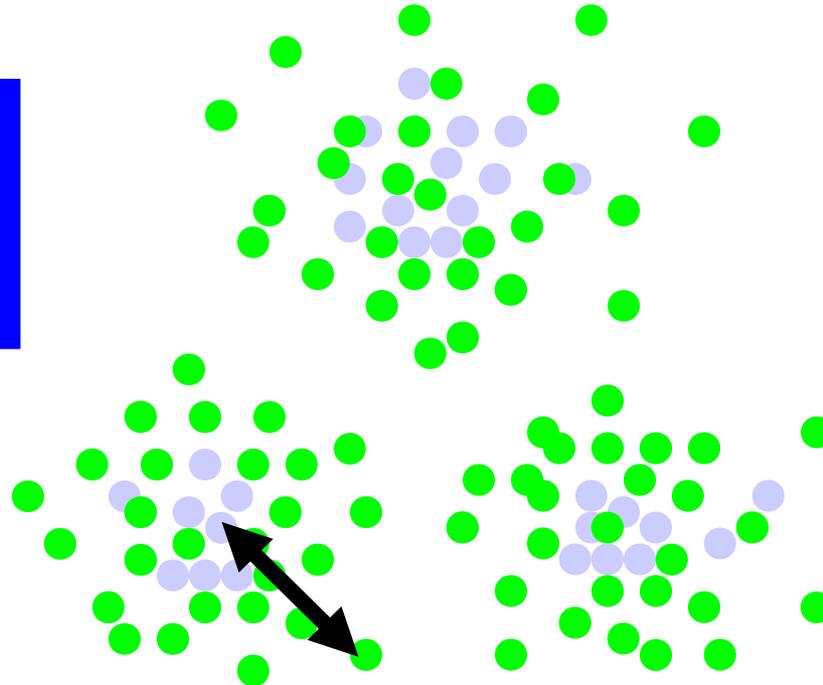
baryons and  
CDM  
experience  
gravitational  
clustering

neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle / m$



# Structure formation after equality

baryons and  
CDM  
experience  
gravitational  
clustering

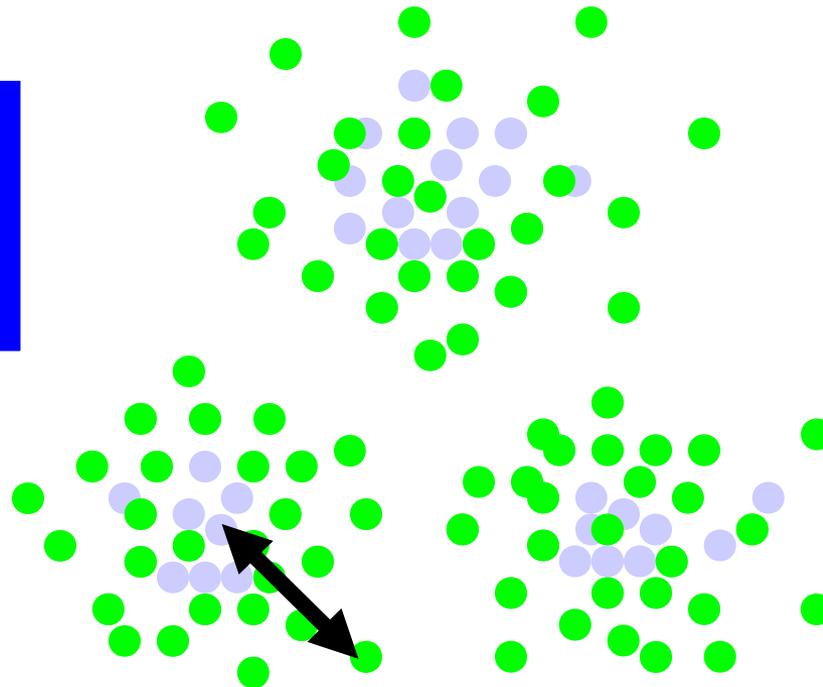


neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle / m$

neutrinos cannot cluster below a diffusion length

# Structure formation after equality

baryons and  
CDM  
experience  
gravitational  
clustering



neutrinos  
experience  
free-streaming  
with  
 $v = c$  or  $\langle p \rangle / m$

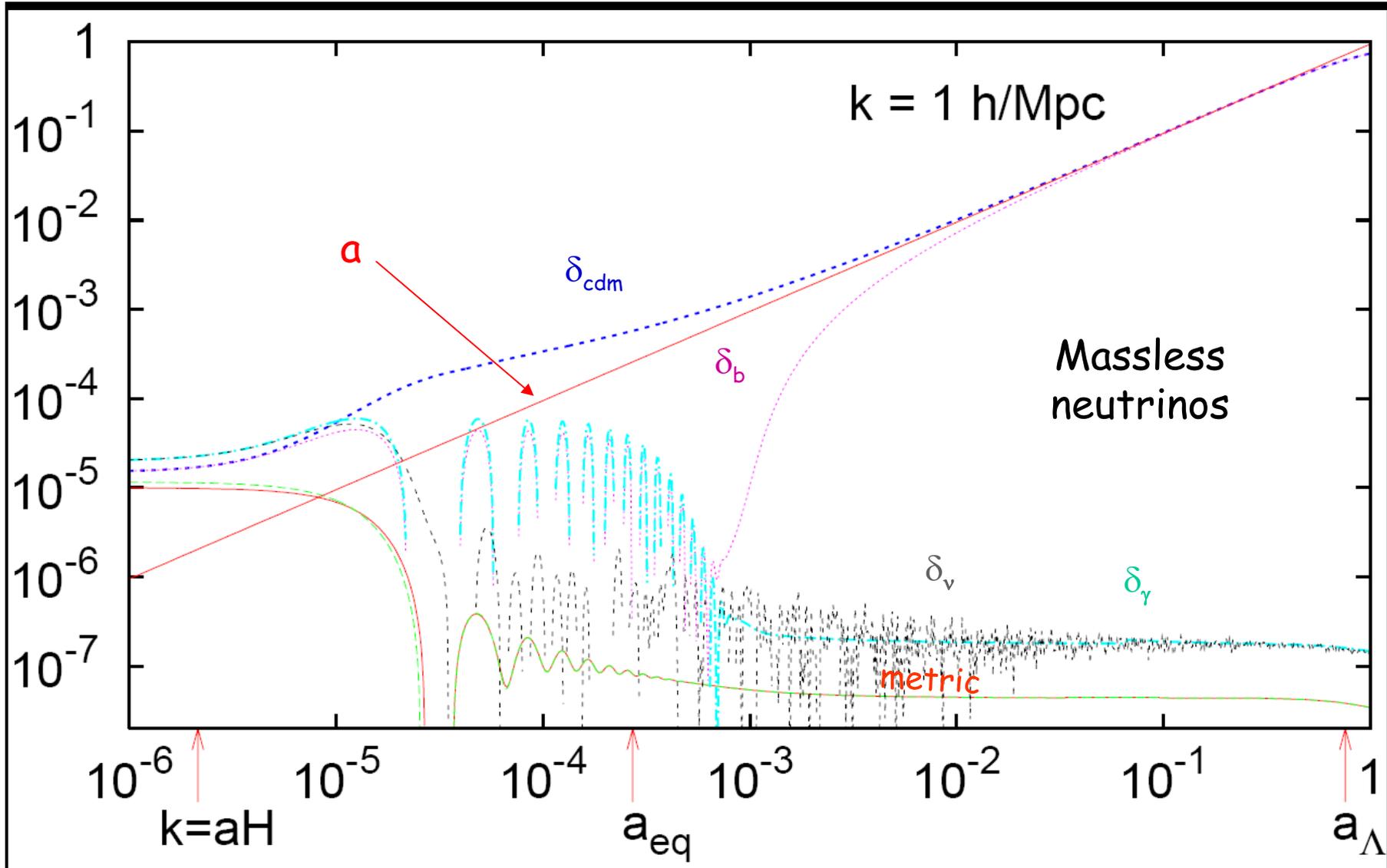
for  $(2\pi/k) < \lambda$  ,

free-streaming suppresses growth of structures during MD

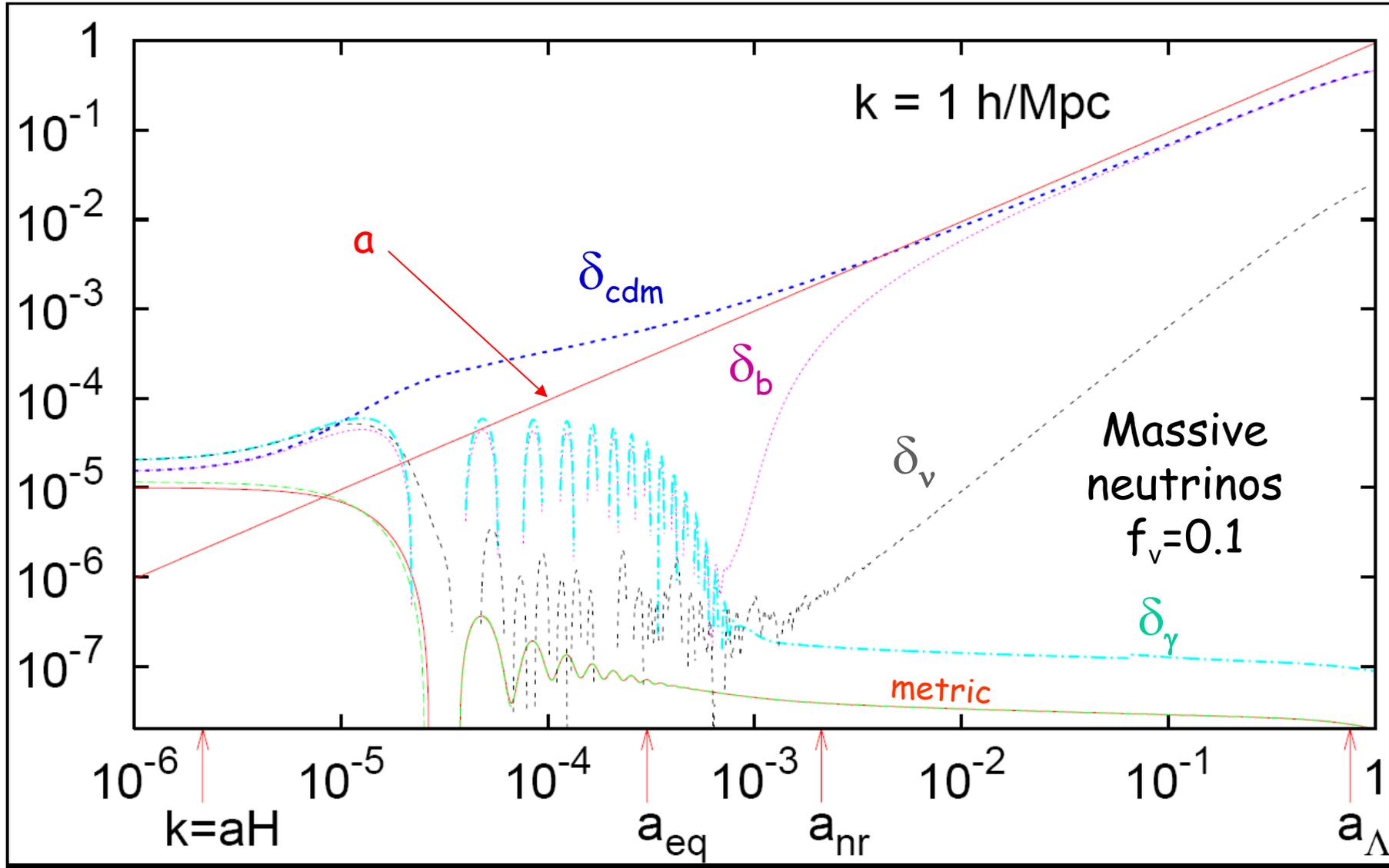
$$\delta\rho/\rho \propto a^{1-3/5} f_v$$

$$\text{with } f_v = \rho_v / \rho_m \approx (\Sigma m_\nu) / (15 \text{ eV})$$

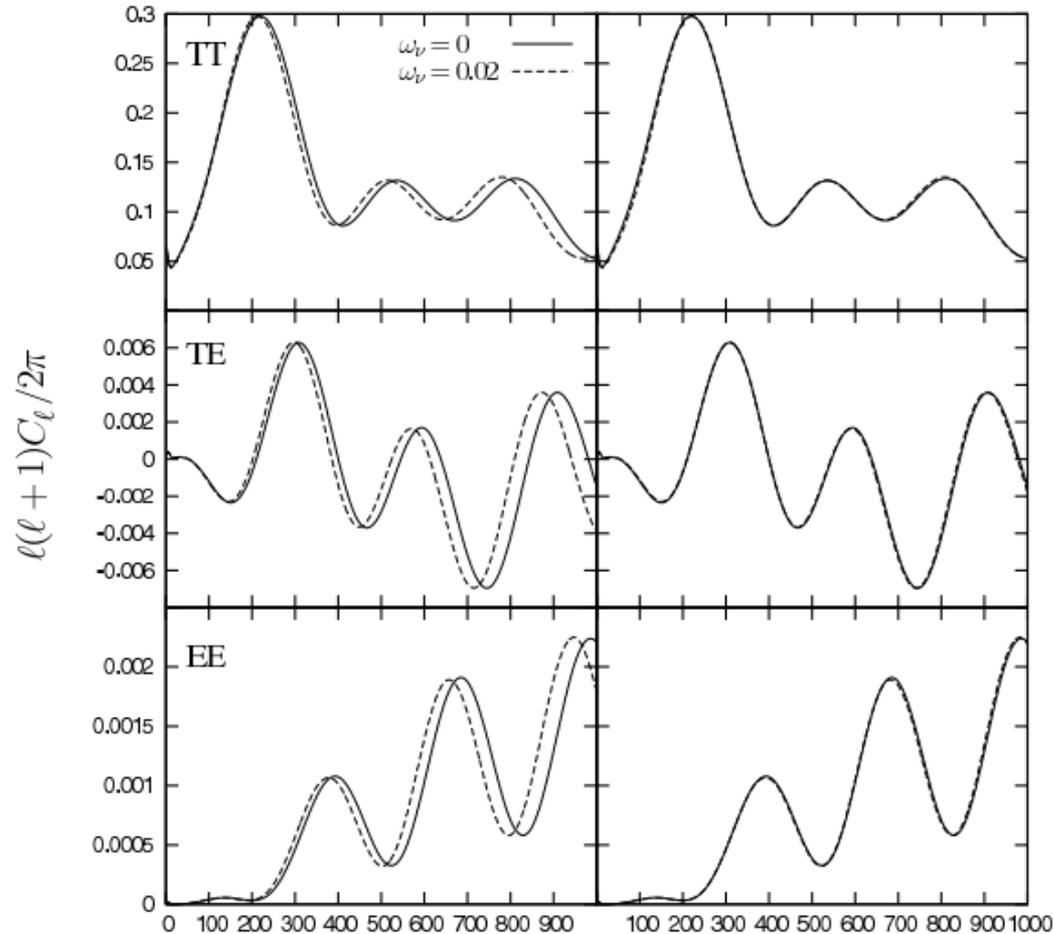
# Structure formation after equality



# Structure formation after equality



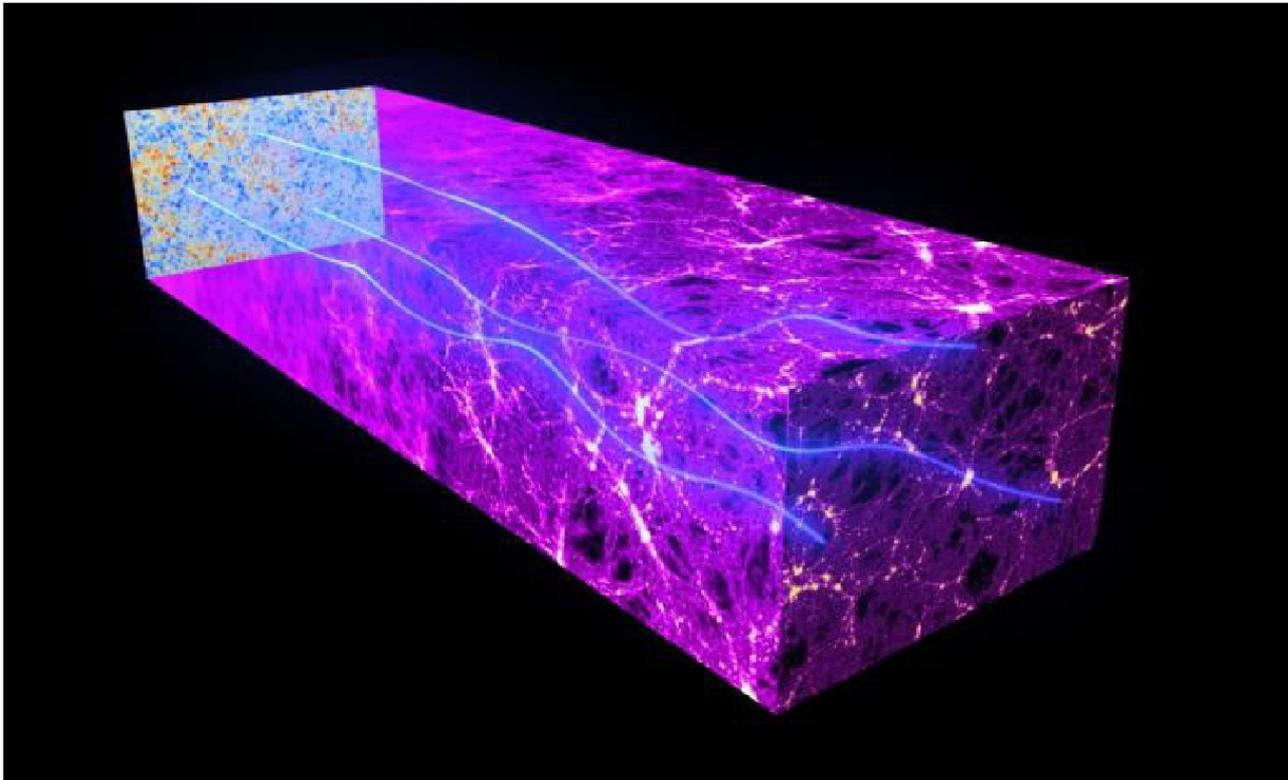
# CMB bounds on neutrino masses



If neutrinos have a mass larger than  $m \sim 0.3$  eV they are non-relativistic at recombination. At this point, increasing their mass increases the dark matter component at recombination. The change in the gravitational potential at early times with massive neutrinos leads to a slight changes in  $C_\ell$  below the first acoustic peak (see Ma & Bertshchinger 1998, Fukugita et al, 2006). In principle, CMB should constrain neutrino masses to  $m < 2$  eV but constraints are actually much better because of lensing.

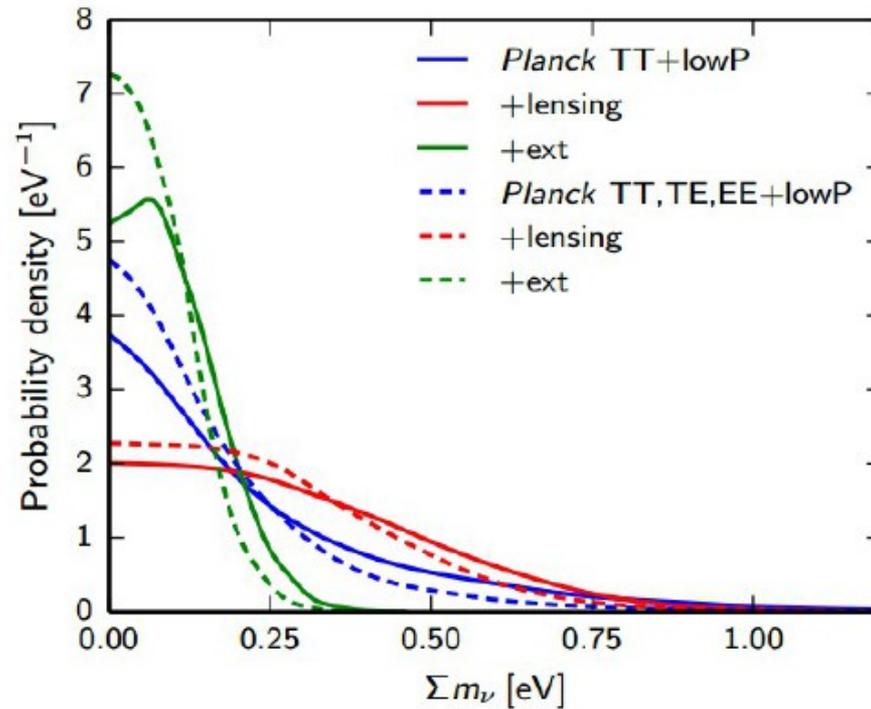
# Gravitational Lensing

The gravitational effects of intervening matter bend the path of CMB light on its way from the early universe to the Planck telescope. This “gravitational lensing” distorts our image of the CMB



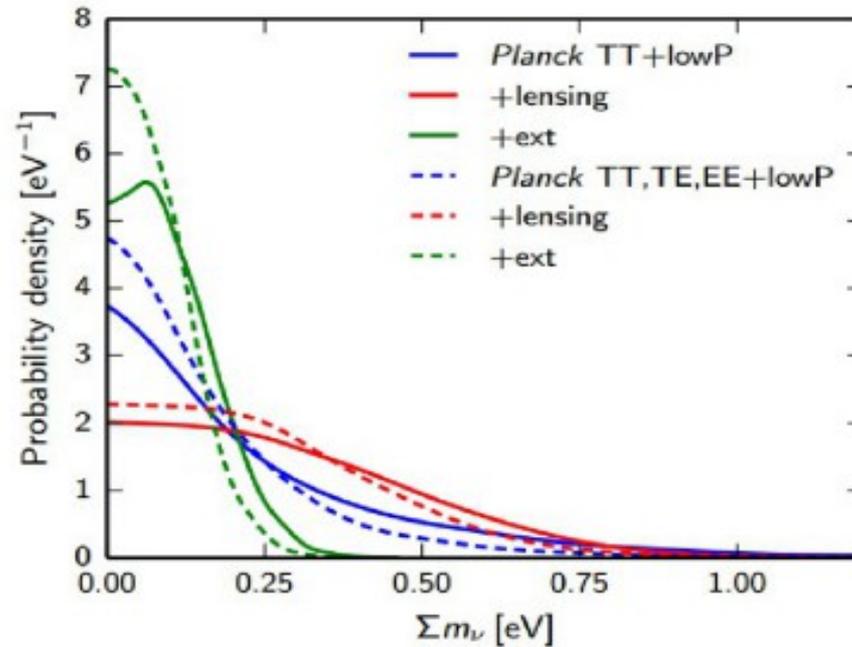
CMB anisotropy do not probe just recombination epoch !  
CMB lensing probes structure formation around  $z \sim 2$  (broadly)

# 3 standard Massive Neutrinos



$$\left. \begin{array}{l} \Sigma m_\nu < 0.23 \text{ eV} \\ \Omega_\nu h^2 < 0.0025 \end{array} \right\} 95\%, \text{ *Planck* TT+lowP+lensing+ext.}$$

# 3 standard Massive Neutrinos



$$\sum m_\nu < 0.72 \text{ eV} \quad \textit{Planck TT+lowP};$$

$$\sum m_\nu < 0.21 \text{ eV} \quad \textit{Planck TT+lowP+BAO};$$

$$\sum m_\nu < 0.49 \text{ eV} \quad \textit{Planck TT, TE, EE+lowP};$$

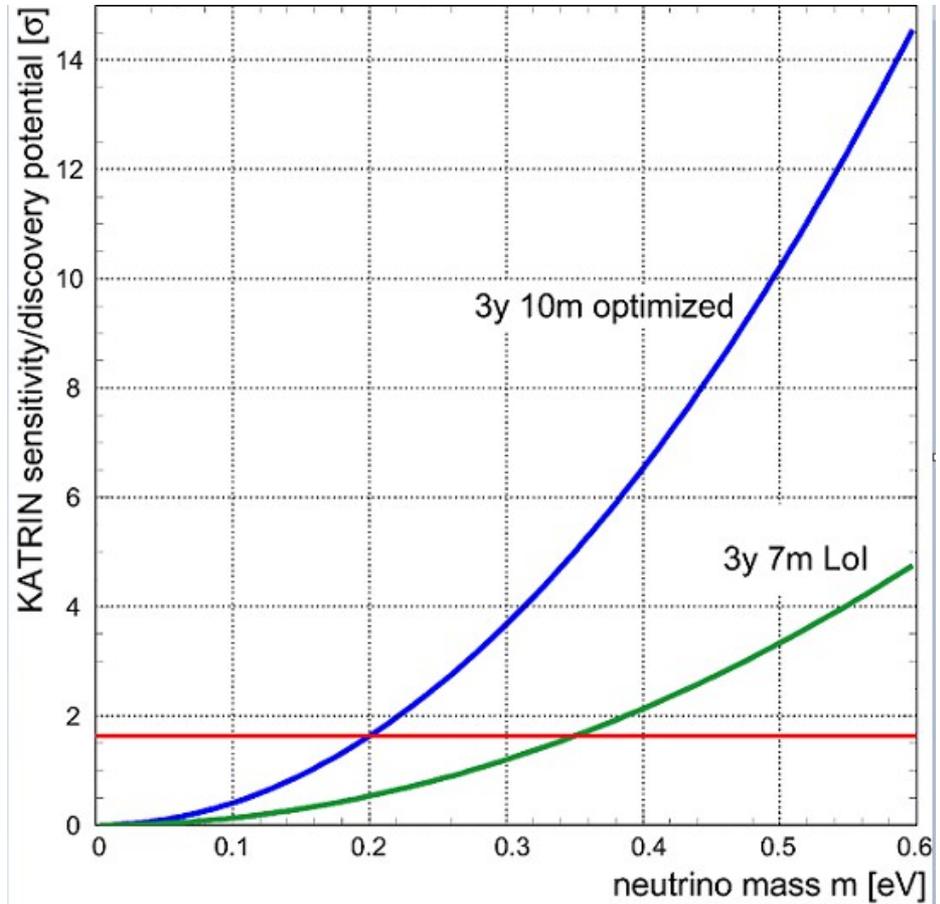
$$\sum m_\nu < 0.17 \text{ eV} \quad \textit{Planck TT, TE, EE+lowP+BAO}.$$

# Implications of a $<0.17$ eV constraint at 95 % c.l. from Planck+BAO

Parameter	Present	
	Normal	Inverted
$M_\nu, \xi$ free	[53 – 170]	[93 – 195]
$M_\nu, \xi \equiv 1$	[53 – 170]	[93 – 194]
$m_{\beta\beta}, \xi$ free	$< 36$	[16 – 59]
$m_{\beta\beta}, \xi \equiv 1$	$< 36$	[16 – 59]
$m_\beta, \xi$ free	[6 – 50]	[46 – 71]
$m_\beta, \xi \equiv 1$	[6 – 50]	[46 – 71]

Constraints in meV !!!!

# Prospects for direct detection for KATRIN....



# Constraints on a massive sterile neutrino

We consider a particle that contributes to  $N_{\text{eff}}$  when is relativistic and behaves with a mass:

$$m_{\nu, \text{sterile}}^{\text{eff}} \equiv (94.1 \Omega_{\nu, \text{sterile}} h^2) \text{ eV}$$

when is not relativistic.

This is a very general parametrization of a sterile neutrino.

If the sterile neutrino is thermally distributed we have:

$$m_{\nu, \text{sterile}}^{\text{eff}} = (T_s/T_\nu)^3 m_{\text{sterile}}^{\text{thermal}} = (\Delta N_{\text{eff}})^{3/4} m_{\text{sterile}}^{\text{thermal}}$$

While for the Dodelson-Widrow case the relation is given by:

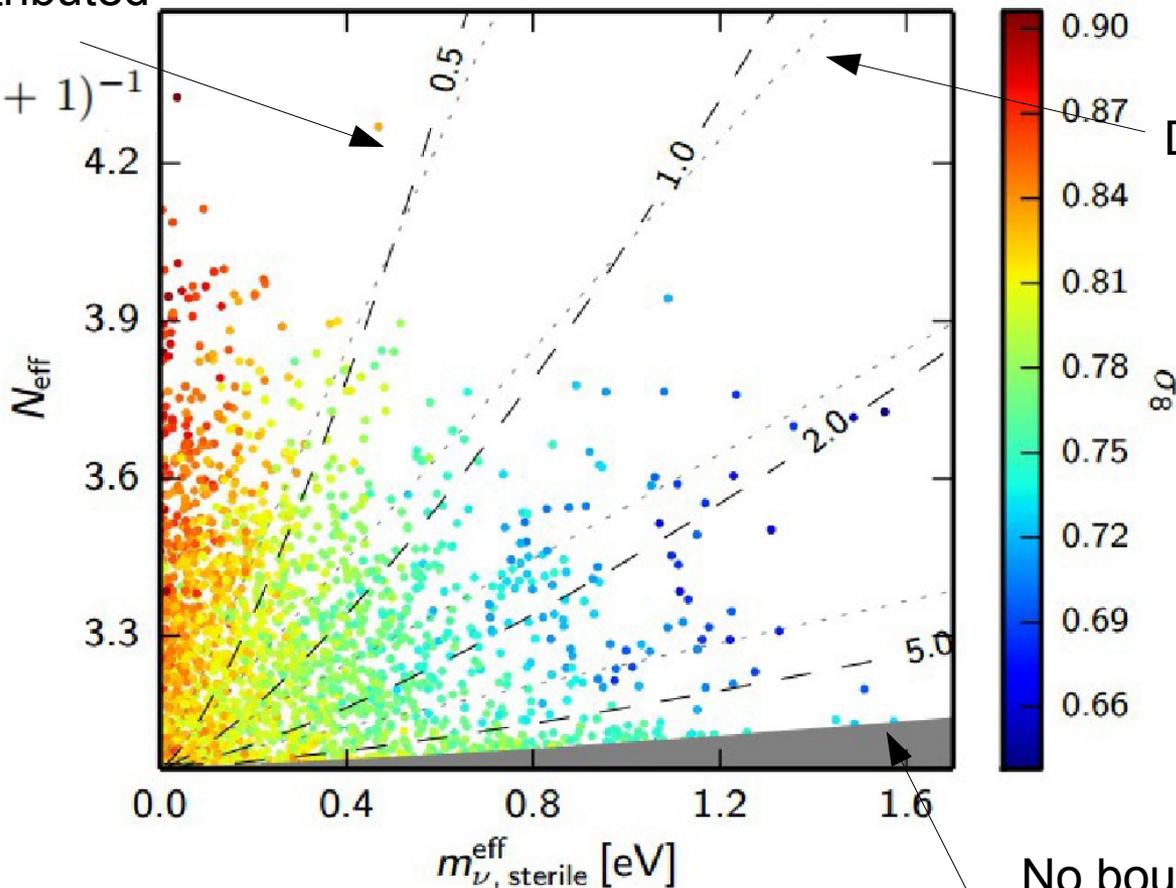
$$m_{\nu, \text{sterile}}^{\text{eff}} = \chi_s m_{\text{sterile}}^{\text{DW}}$$

$$\Delta N_{\text{eff}} = \chi_s$$

# Massive Sterile Neutrinos

Thermally distributed

$$f_s(E) = (e^{E/T_s} + 1)^{-1}$$



Dodelson-Widrow

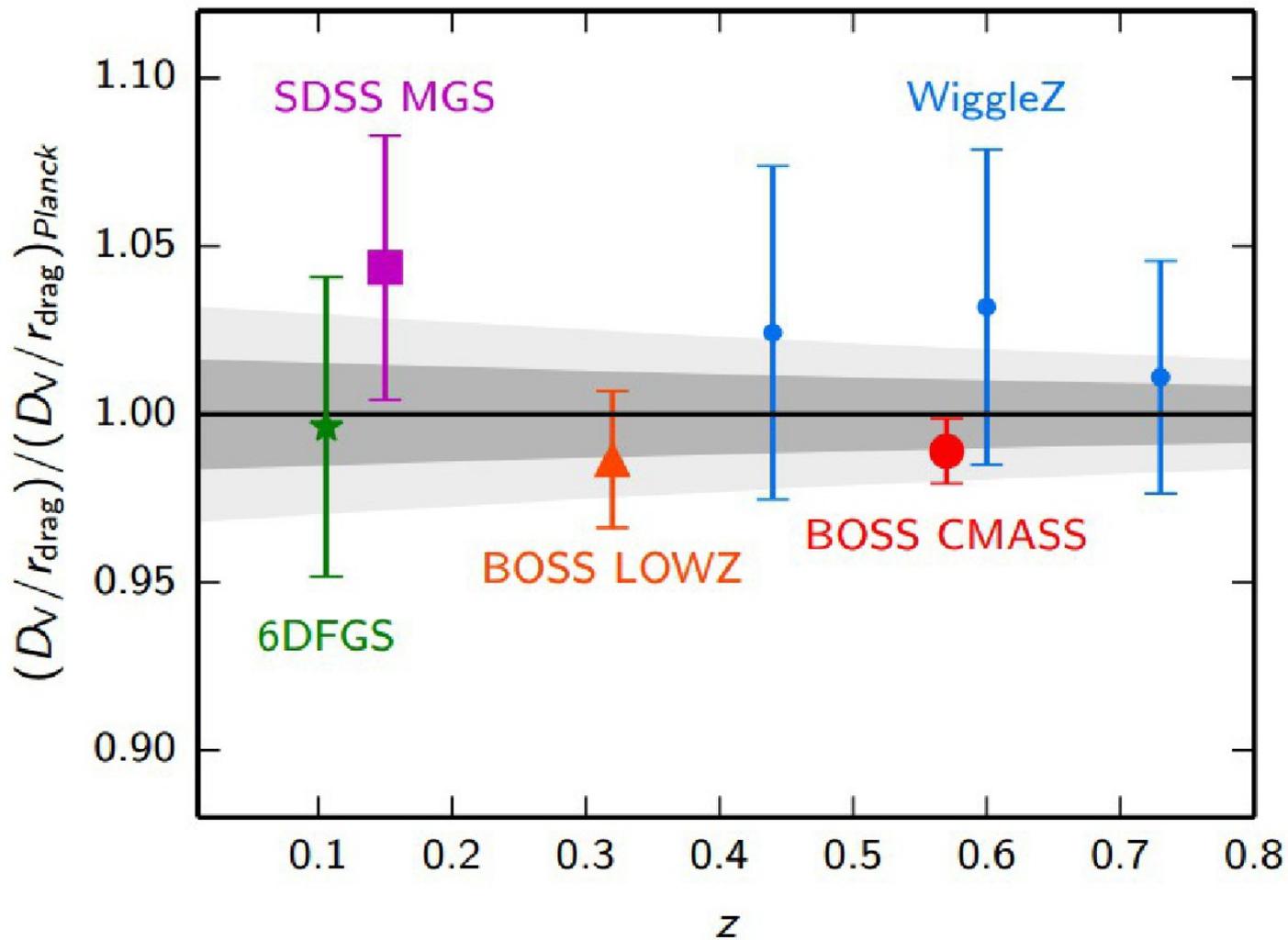
$$f_s(E) = \frac{\chi_s}{(e^{E/T_\nu} + 1)}$$

$\theta_{12}$

No bound !  
Sterile neutrino is cold  
dark matter (for the CMB)

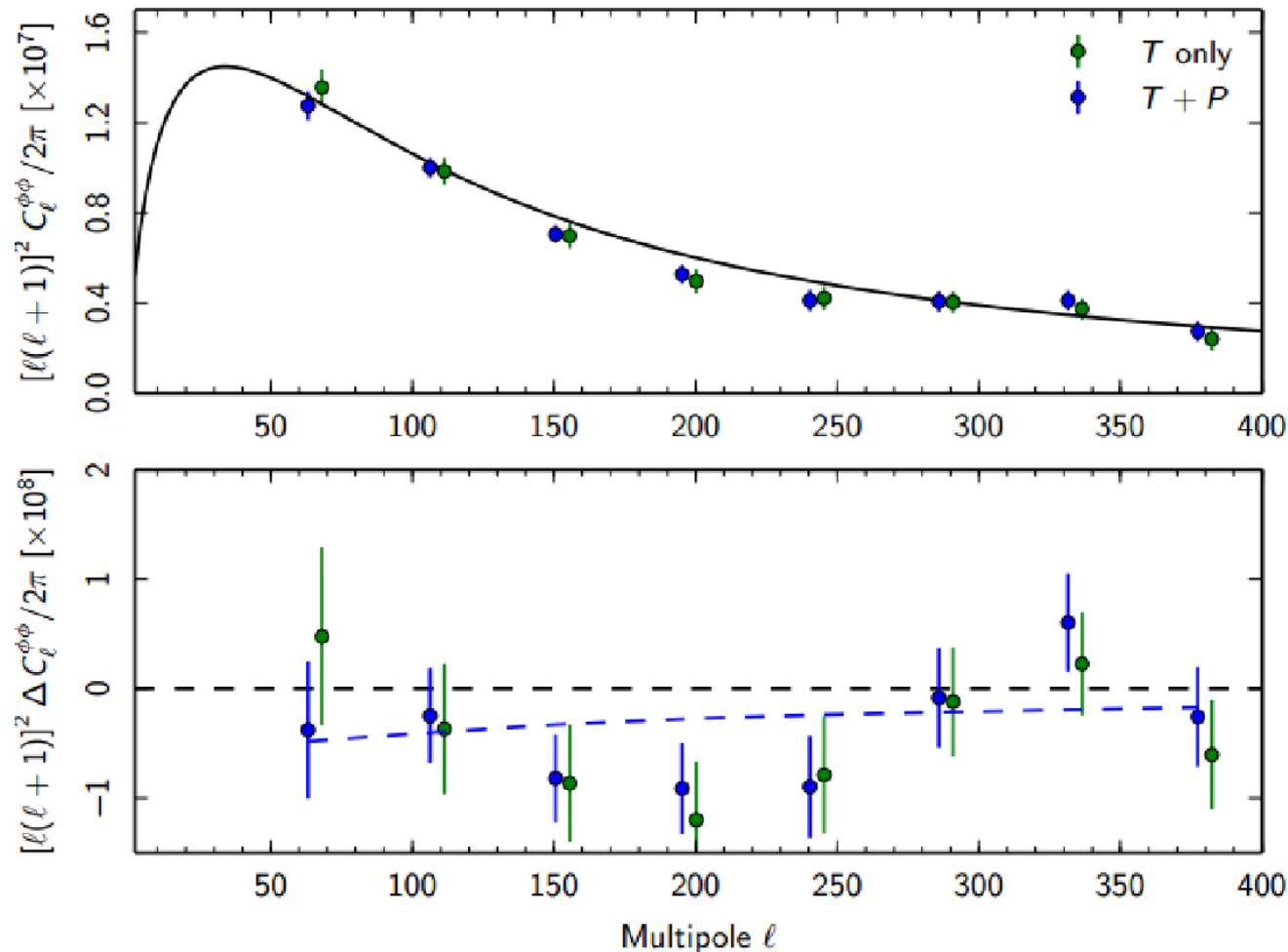
$$\left. \begin{array}{l} N_{\text{eff}} < 3.7 \\ m_{\nu, \text{sterile}}^{\text{eff}} < 0.38 \text{ eV} \end{array} \right\} 95\%, \text{ Planck TT+lowP+lensing+BAO.}$$

Consistency with other datasets...



Acoustic-scale distance ratio  $D_V(z)/r_{\text{drag}}$  in the base  $\Lambda$ CDM model divided by the mean distance ratio from Planck TT+lowP+lensing. The points with  $1\sigma$  errors are as follows: green star (6dFGS, Beutler et al. 2011); square (SDSS MGS, Ross et al. 2014); red triangle and large circle (BOSS “LOWZ” and CMASS surveys, Anderson et al. 2014); and small blue circles (WiggleZ, as analysed by Kazin et al. 2014). The grey bands show the 68% and 95% confidence ranges allowed by Planck TT+lowP+lensing.

# Small tension between Planck TT best fit (assuming LCDM) and CMB lensing dataset



Planck cosmology still in tension with:

- HST value of the Hubble constant
- Weak lensing data from CFHTlens survey
- SZ clusters counts
- Redshift space distortions

- **HST value of the Hubble constant:**

$$H_0 = (73.8 \pm 2.4) \text{ km s}^{-1}\text{Mpc}^{-1}$$

## Possible solution:

- In the paper we use the R11 Cepheid data reanalysed by Efstathiou (2014) using the revised geometric maser distance to NGC 4258 of Humphreys et al. (2013) as a distance anchor, finding a “conservative”  $H_0$  prior:

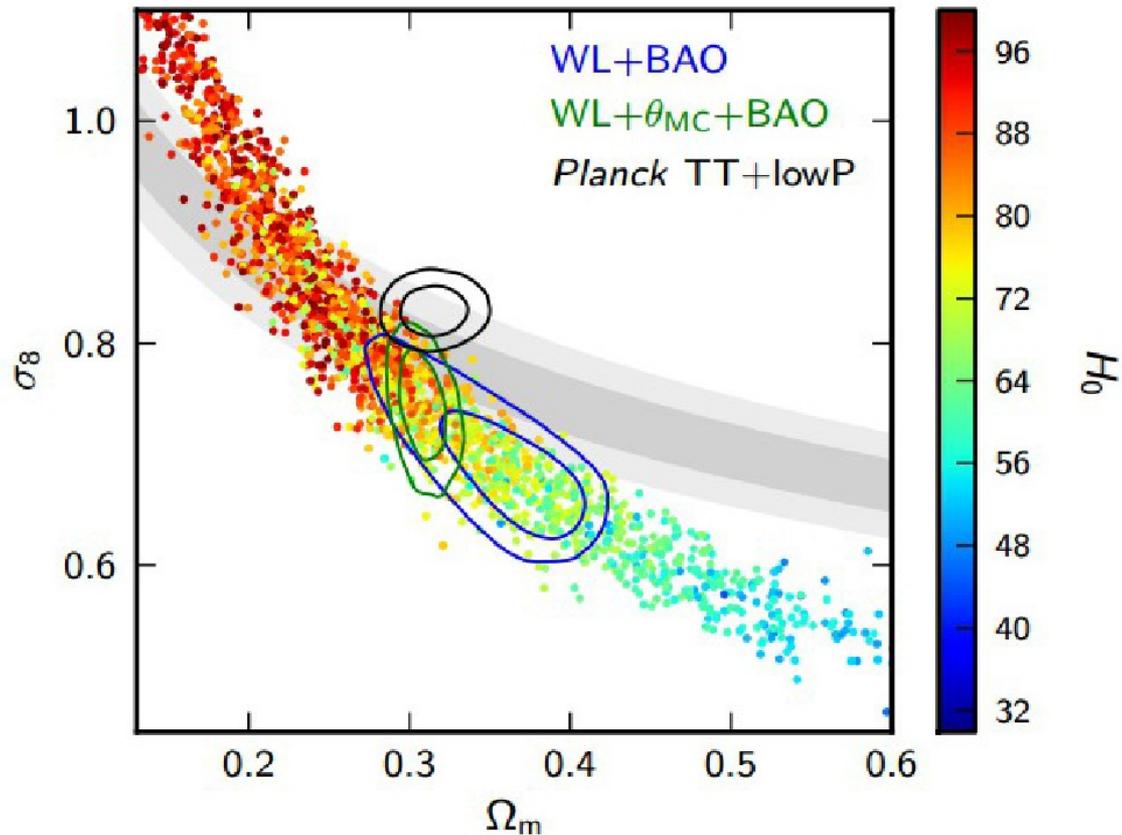
$$H_0 = (70.6 \pm 3.3) \text{ km s}^{-1}\text{Mpc}^{-1}$$

which is within  $1\sigma$  of the Planck TT estimate

$$H_0 = (69.7 \pm 2.1) \text{ km s}^{-1}\text{Mpc}^{-1}, \quad \text{WMAP9,}$$

$$H_0 = (68.0 \pm 0.7) \text{ km s}^{-1}\text{Mpc}^{-1}, \quad \text{WMAP9+BAO.}$$

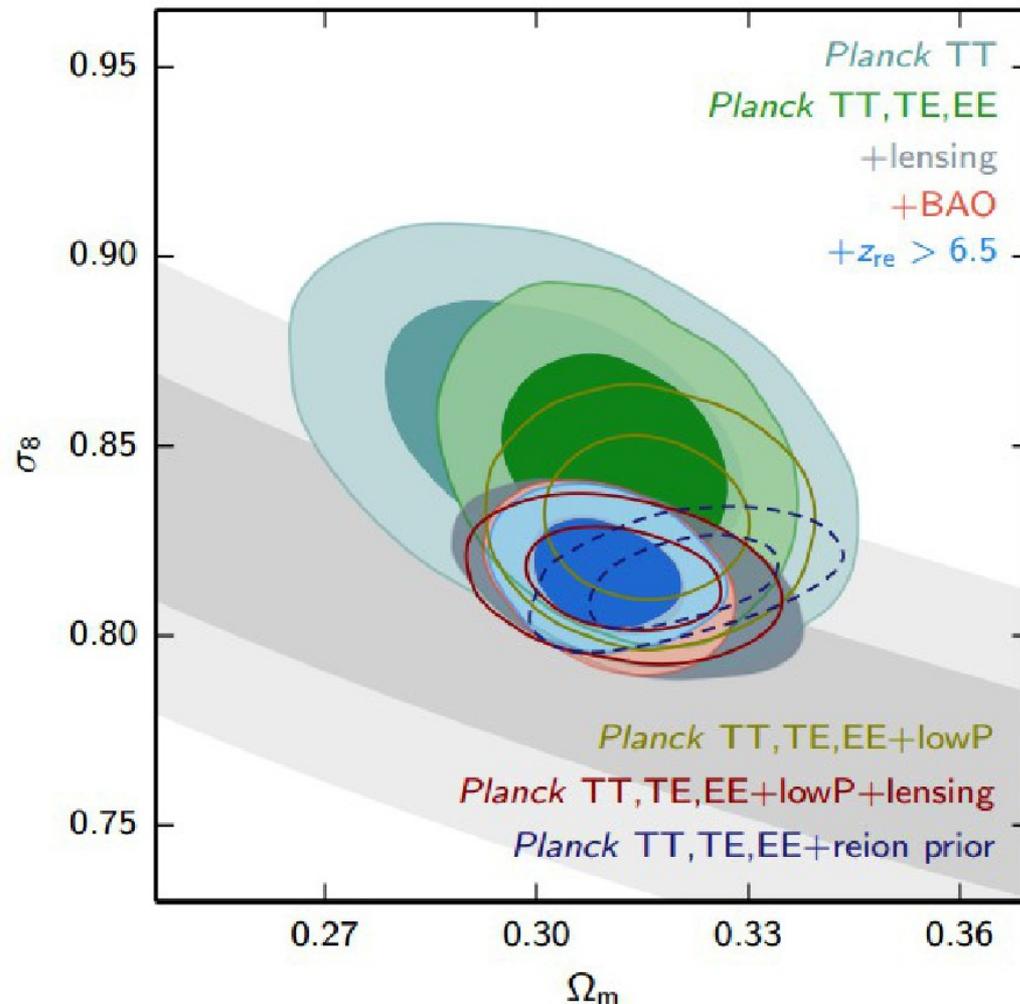
- Weak lensing data from CFHTlens survey.



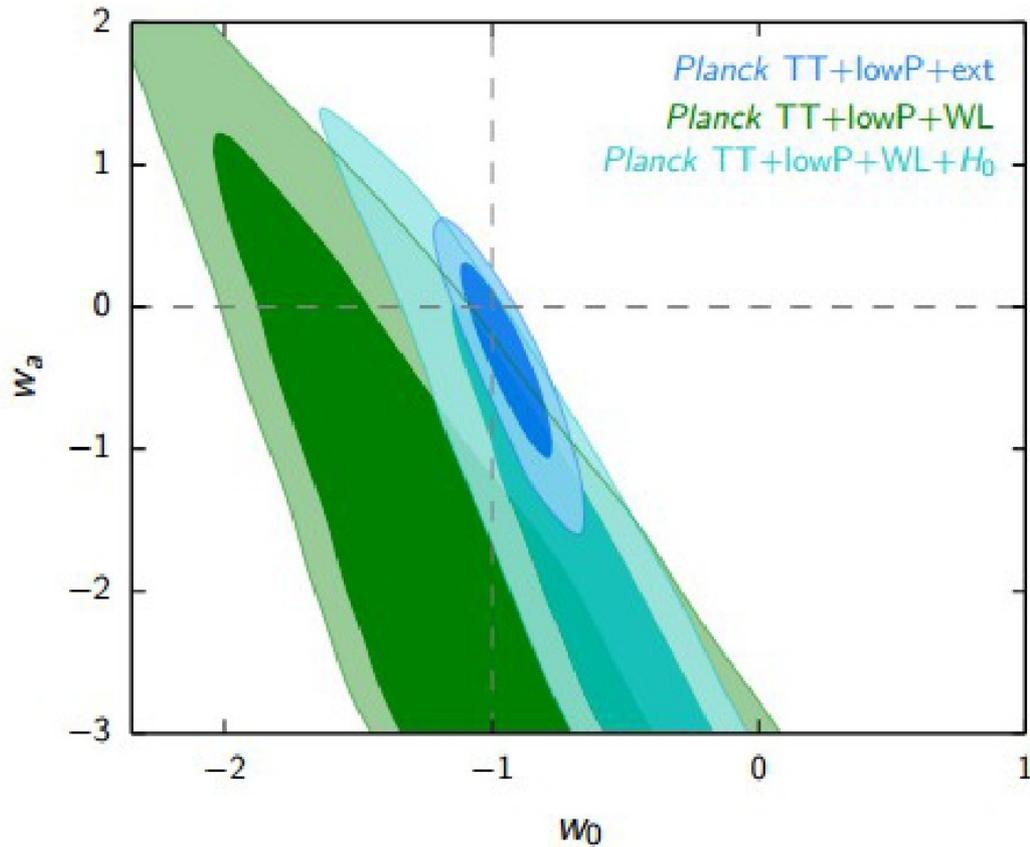
The grey band show the constraint from Planck CMB lensing.

# • SZ clusters counts

The Planck catalog consist of 439 clusters detected via their Sunyaev-Zeldovich (SZ) signal and prefers smaller value than Planck CMB for  $\sigma_8$ .

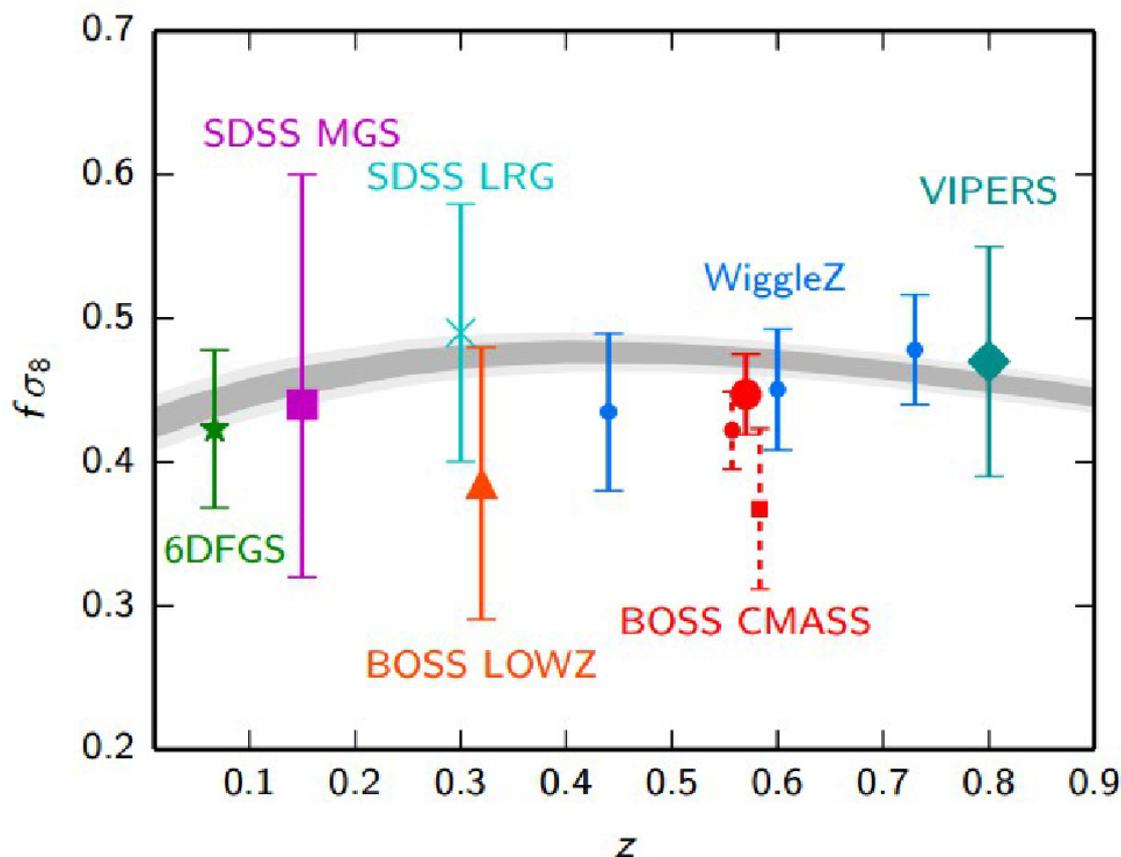


# Planck+external datasets



Ext: BAO+JLA+H0 (Efstathiou)

# • Redshift space distortions



Constraints on the growth rate of fluctuations from various redshift surveys in the base  $\Lambda$ CDM model. The large red circle is BOSS CMASS, as analysed by Samushia et al. 2014, while the points with dashed red error bars correspond to alternative analyses of BOSS CMASS from Beutler et al. (2014b, small circle) and Chuang et al. (2013, small square). The BOSS CMASS points are based on the same data set and are therefore not independent. The grey bands show the range allowed by Planck TT+lowP+lensing in the base  $\Lambda$ CDM model.

# Extending LCDM: going beyond 6 parameters

Di Valentino, AM, Silk, arXiv:1507.06646

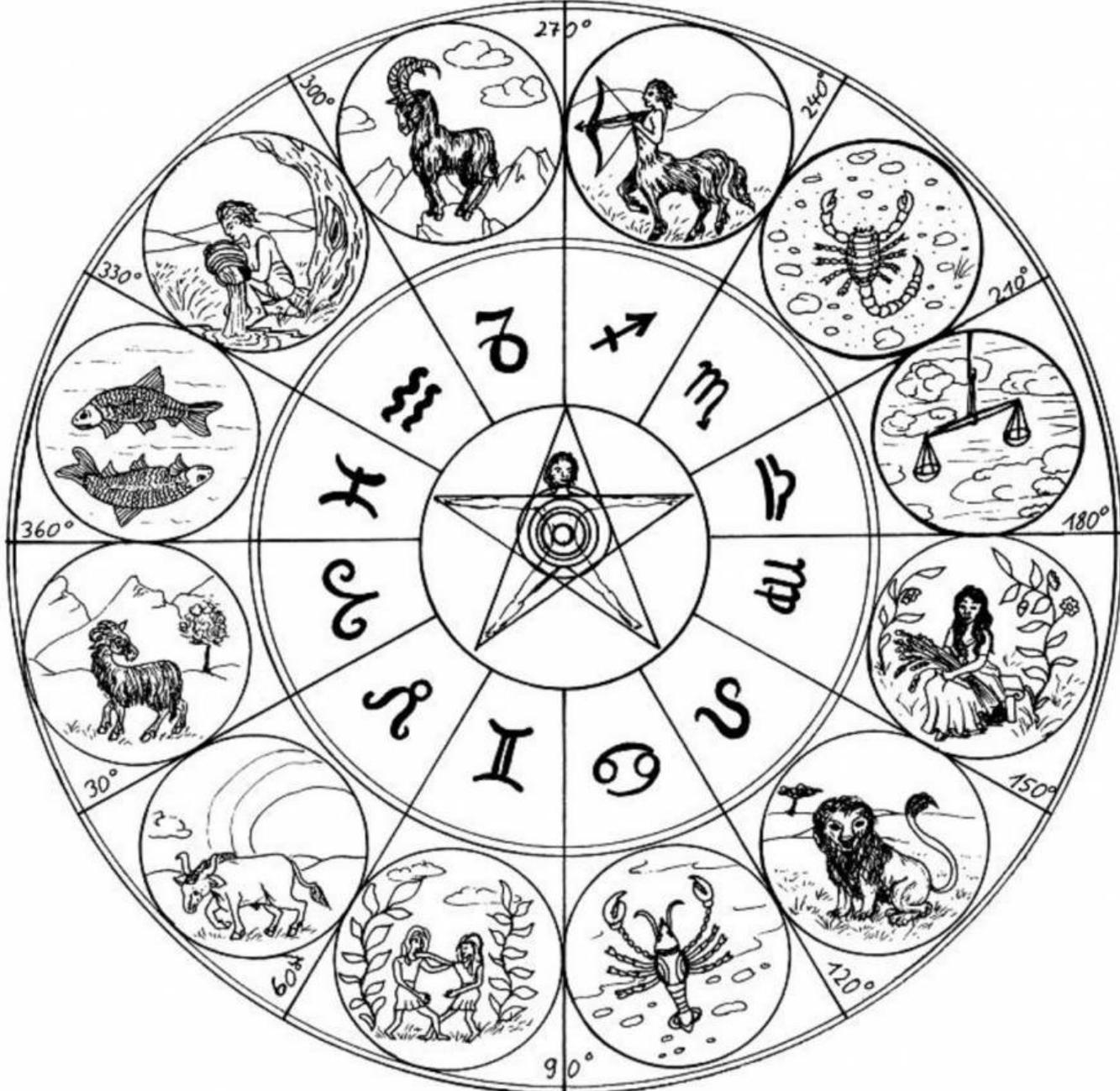
When we analyze current datasets we often limit ourselves to the 6 parameters of  $\Lambda$ CDM or to 1-2 parameter extensions to it.

Given the high quality of the new data we should extend the parameter space and analyze the data considering at the same time:

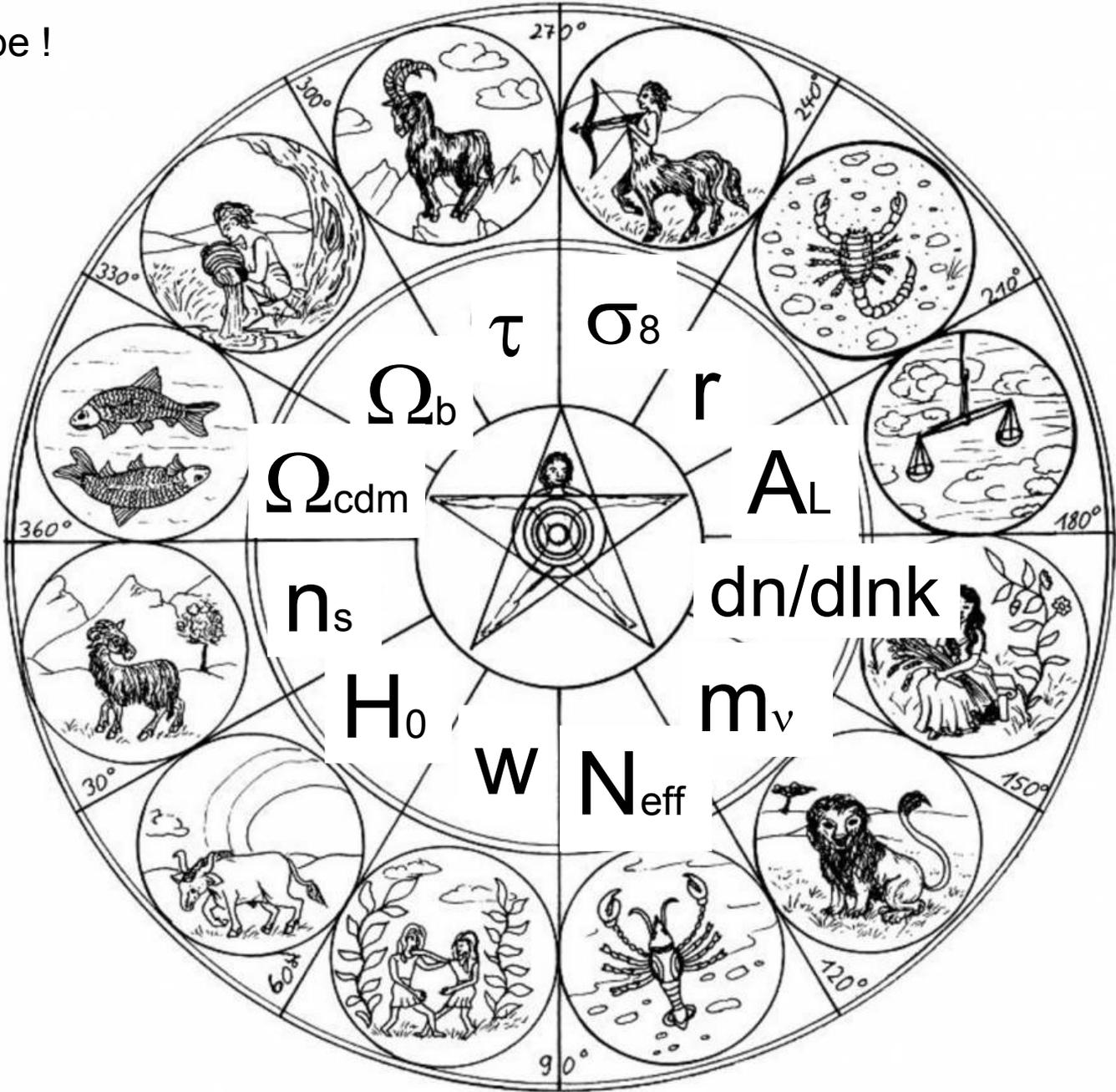
- Variation in the dark energy equation of state  $w$
- Neutrino mass.
- Neutrino effective number
- Gravitational waves (tensor component)
- Running of the primordial spectral index.
- Lensing amplitude in the CMB angular spectra.

i.e. we should, at least, vary 12 cosmological parameters instead of just 6 !

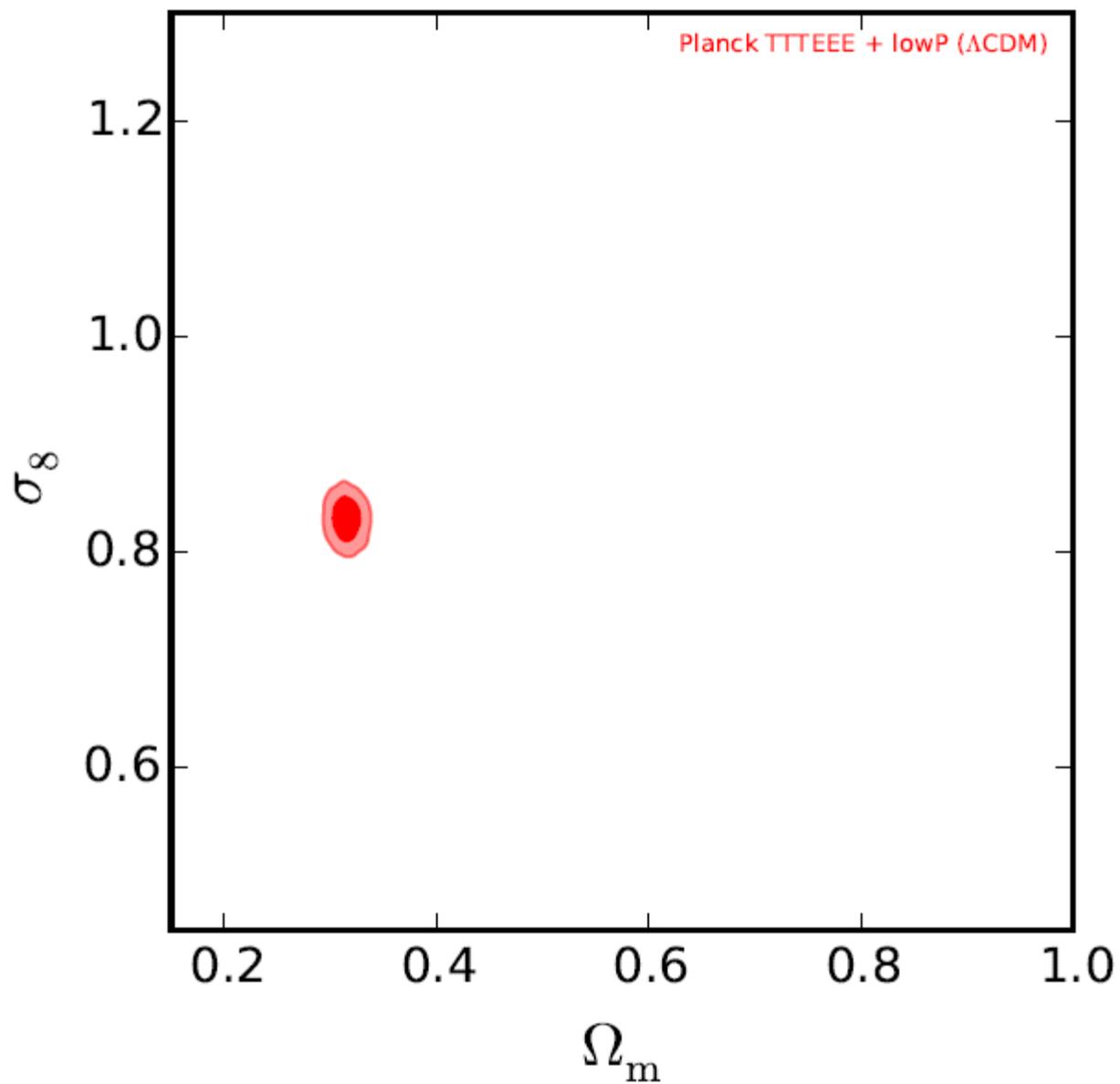
Horo-scope



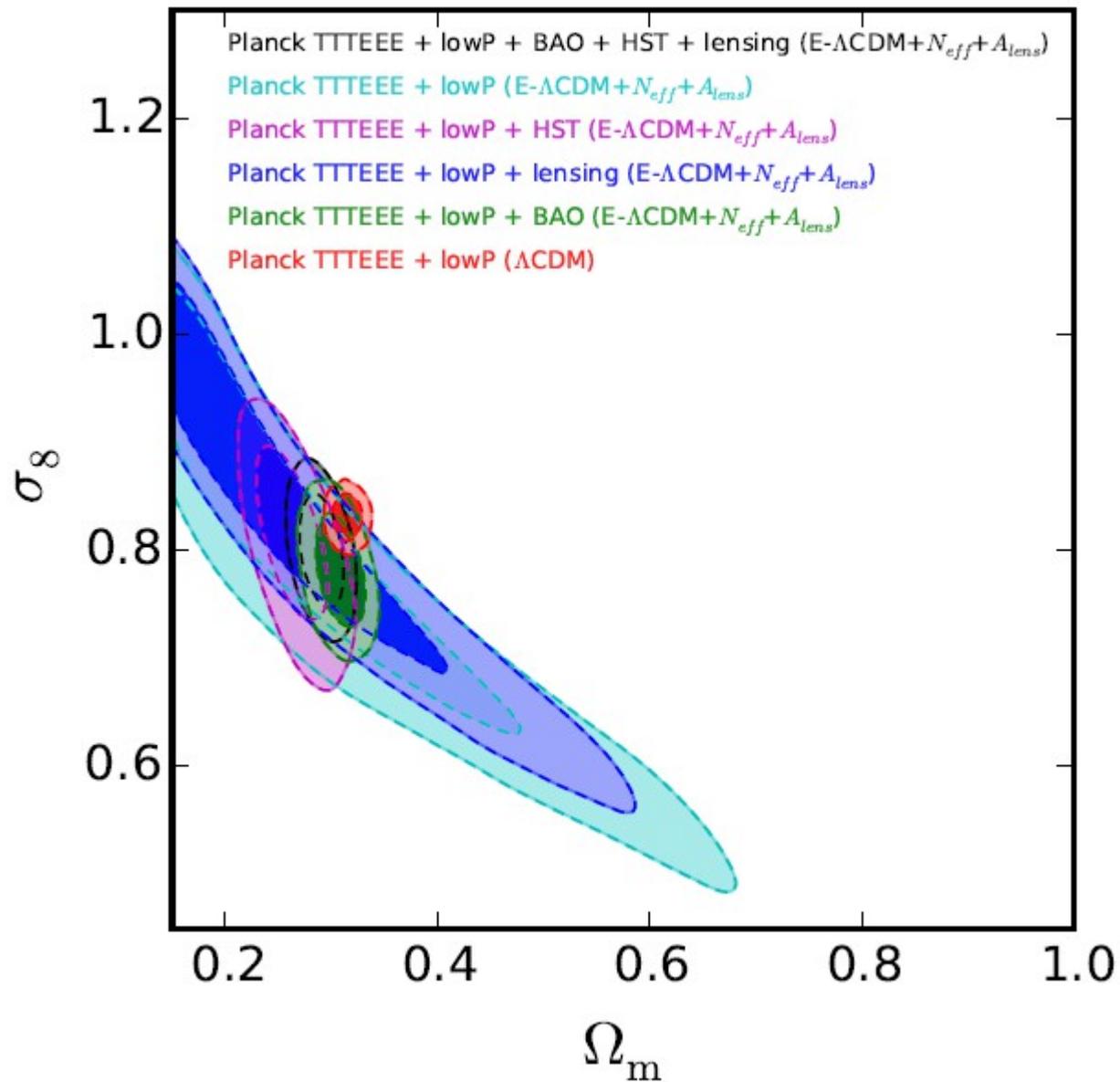
Cosmo-scope !



# Constraints from Planck assuming 6 parameters LCDM



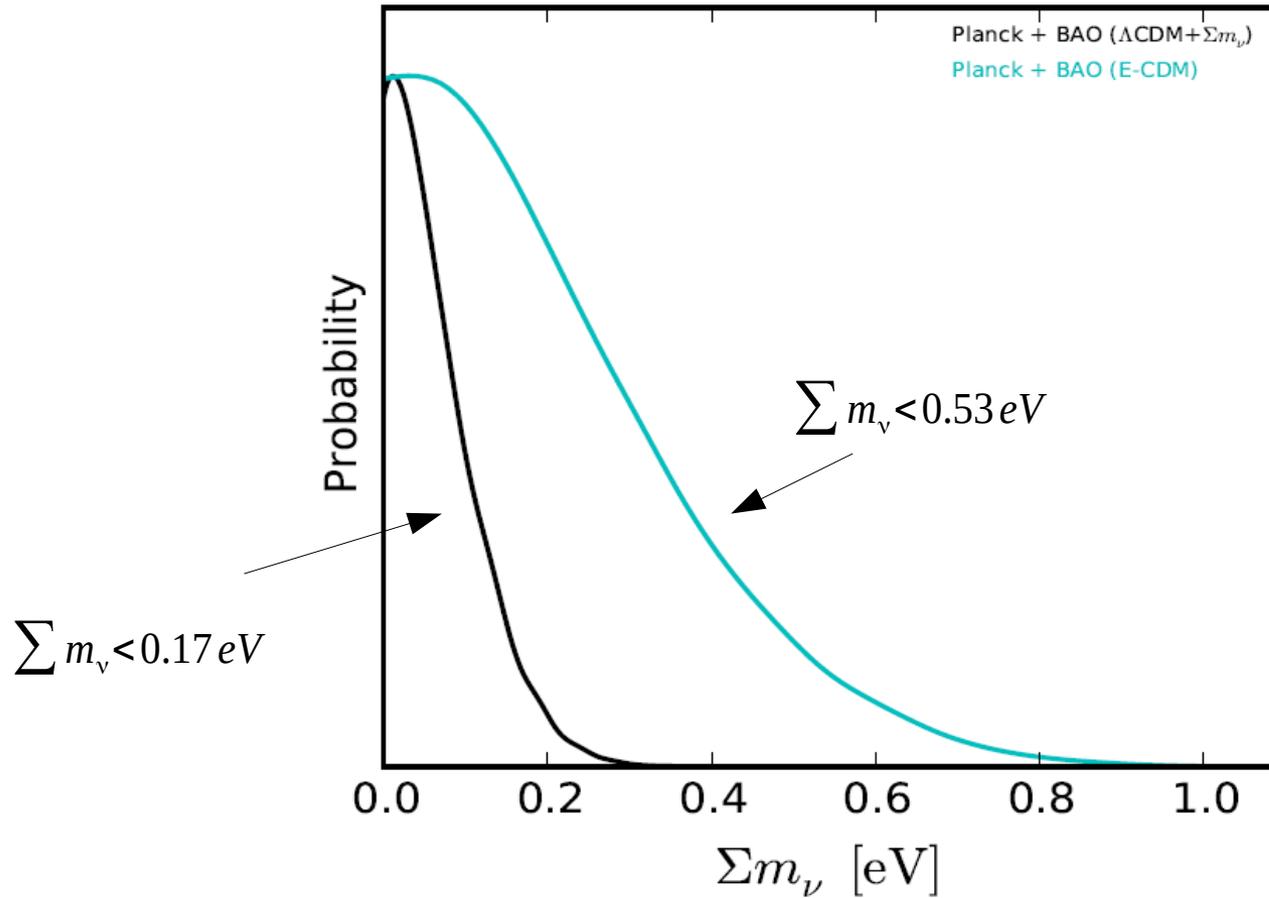
# Constraints from Planck assuming 12 parameters $\Lambda$ CDM



# Limits at 95 % c.l.

Model Dataset	$\Omega_b h^2$	$\Omega_c h^2$	$H_0$	$\tau$	$n_s$	$\sigma_8$	$\frac{dn_s}{d \ln k}$	$r$	$w$	$\Sigma m_\nu [eV]$	$N_{\text{eff}}$	$A_{\text{lens}}$
$\Lambda$ CDM Planck	$0.02226^{+0.00031}_{-0.00029}$	$0.1198^{+0.0028}_{-0.0028}$	$67.3^{+1.3}_{-1.3}$	$0.079^{+0.034}_{-0.035}$	$0.9646^{+0.0092}_{-0.0092}$	$0.831^{+0.026}_{-0.026}$	-	-	-	-	-	-
$\Lambda$ CDM Planck+ BAO	$0.02229^{+0.00028}_{-0.00027}$	$0.1193^{+0.0021}_{-0.0020}$	$67.52^{+0.93}_{-0.93}$	$0.082^{+0.031}_{-0.032}$	$0.9662^{+0.0078}_{-0.0079}$	$0.832^{+0.025}_{-0.025}$	-	-	-	-	-	-
$e$ CDM Planck	$0.02239^{+0.00060}_{-0.00056}$	$0.1186^{+0.0071}_{-0.0068}$	$> 51.2$	$0.058^{+0.040}_{-0.043}$	$0.967^{+0.025}_{-0.025}$	$0.81^{+0.24}_{-0.26}$	$-0.003^{+0.020}_{-0.019}$	$< 0.183$	$-1.32^{+0.98}_{-0.85}$	$< 0.959$	$3.08^{+0.57}_{-0.51}$	$1.21^{+0.27}_{-0.24}$
$e$ CDM Planck+BAO	$0.02251^{+0.00056}_{-0.00052}$	$0.1185^{+0.0069}_{-0.0069}$	$68.4^{+4.3}_{-4.1}$	$0.058^{+0.041}_{-0.043}$	$0.972^{+0.024}_{-0.024}$	$0.781^{+0.065}_{-0.063}$	$-0.004^{+0.018}_{-0.018}$	$< 0.187$	$-1.04^{+0.20}_{-0.21}$	$< 0.534$	$3.11^{+0.52}_{-0.48}$	$1.20^{+0.19}_{-0.19}$
$e$ CDM Planck+lensing	$0.02214^{+0.00053}_{-0.00052}$	$0.1176^{+0.0069}_{-0.0066}$	$> 54.5$	$0.058^{+0.040}_{-0.043}$	$0.959^{+0.024}_{-0.024}$	$0.85^{+0.21}_{-0.24}$	$-0.005^{+0.018}_{-0.018}$	$< 0.178$	$-1.45^{+0.96}_{-0.83}$	$< 0.661$	$2.93^{+0.51}_{-0.48}$	$1.04^{+0.16}_{-0.15}$
$e$ CDM Planck+HST	$0.02239^{+0.00059}_{-0.00057}$	$0.1187^{+0.0072}_{-0.0070}$	$74.4^{+5.1}_{-5.1}$	$0.057^{+0.040}_{-0.045}$	$0.966^{+0.025}_{-0.025}$	$0.81^{+0.10}_{-0.11}$	$-0.003^{+0.020}_{-0.019}$	$< 0.186$	$-1.32^{+0.29}_{-0.31}$	$< 0.957$	$3.09^{+0.58}_{-0.55}$	$1.18^{+0.19}_{-0.18}$
$e$ CDM Planck+JLA	$0.02242^{+0.00058}_{-0.00056}$	$0.1188^{+0.0071}_{-0.0067}$	$67.4^{+4.4}_{-4.2}$	$0.058^{+0.040}_{-0.043}$	$0.968^{+0.025}_{-0.025}$	$0.759^{+0.088}_{-0.089}$	$-0.004^{+0.020}_{-0.019}$	$< 0.183$	$-1.06^{+0.13}_{-0.14}$	$< 0.854$	$3.10^{+0.57}_{-0.54}$	$1.20^{+0.19}_{-0.17}$
$e$ CDM Planck+WL	$0.02251^{+0.00056}_{-0.00055}$	$0.1188^{+0.0073}_{-0.0069}$	$> 54.2$	$< 0.0835$	$0.972^{+0.024}_{-0.024}$	$0.82^{+0.22}_{-0.25}$	$0.000^{+0.020}_{-0.019}$	$< 0.197$	$-1.41^{+0.98}_{-0.79}$	$< 0.974$	$3.16^{+0.58}_{-0.56}$	$1.24^{+0.23}_{-0.22}$
$e$ CDM Planck+BAO-RSD	$0.02253^{+0.00052}_{-0.00050}$	$0.1184^{+0.0069}_{-0.0067}$	$68.6^{+4.2}_{-3.9}$	$0.056^{+0.038}_{-0.042}$	$0.972^{+0.023}_{-0.023}$	$0.774^{+0.055}_{-0.058}$	$-0.004^{+0.018}_{-0.018}$	$< 0.188$	$-1.05^{+0.17}_{-0.19}$	$< 0.626$	$3.12^{+0.51}_{-0.48}$	$1.22^{+0.18}_{-0.17}$
$e$ CDM Planck+BKP	$0.02237^{+0.00057}_{-0.00056}$	$0.1186^{+0.0072}_{-0.0069}$	$> 52.3$	$0.058^{+0.039}_{-0.044}$	$0.966^{+0.026}_{-0.026}$	$0.81^{+0.23}_{-0.25}$	$-0.003^{+0.019}_{-0.018}$	$< 0.101$	$-1.31^{+0.96}_{-0.89}$	$< 0.876$	$3.07^{+0.57}_{-0.55}$	$1.20^{+0.24}_{-0.22}$

Di Valentino, AM, Silk, arXiv:1507.06646



When moving from 6+1 to 12 parameters, the Planck+BAO 95% c.i. constraint on the sum of neutrino masses is relaxed by a factor  $\sim 3$ .

# Conclusions

- The 6-parameters cosmological standard model has been further confirmed By the Planck 2015 data release.
- We have strong implications for neutrino physics: neutrino effective number and total neutrino mass are now strongly constrained.
- no evidence for extra radiation or extra sterile neutrino.
- no evidence for neutrino mass. If constraints under LCDM are true, it will Be really difficult to detect a neutrino mass with planned beta and double beta Decay experiments.
- Cosmological constraints are model dependent ! Increasing the number of parameters could relax current neutrino constraints by a factor 3.

Future direct neutrino mass measurements will provide a wonderful test for LCDM cosmology !