

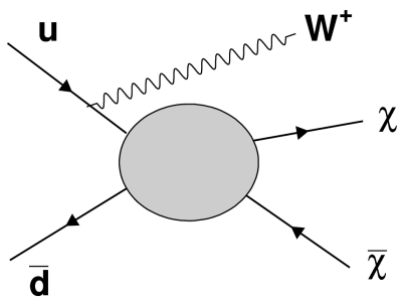
# Dark Matter Mono-Lepton Channels

## Simplified model in Madgraph

Kerstin Hoepfner, Viktor Kutzner, Klaas Padeken  
RWTH Aachen University, III. Phys. Inst. A

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# Mono-lepton Channels

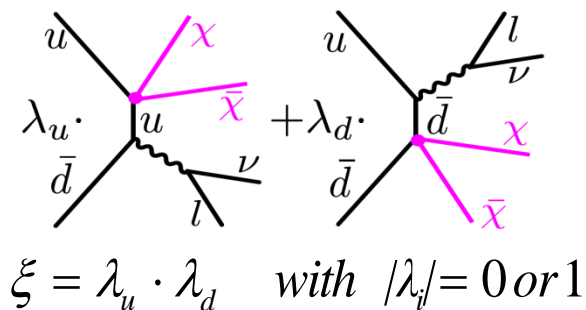


Channels used:

$W \rightarrow e\nu + \text{MET}$

$W \rightarrow \mu\nu + \text{MET}$

Signature: high  $p_T$  lepton + MET (dominated by  $\chi\chi$ )



$\sqrt{8}$  TeV: used EFT. Madgraph.

Vector (V) and axial-vector (AV) couplings

Emphasis on **interference**, limits for  $\xi = \pm 1, 0$

arXiv:1208.4361v2

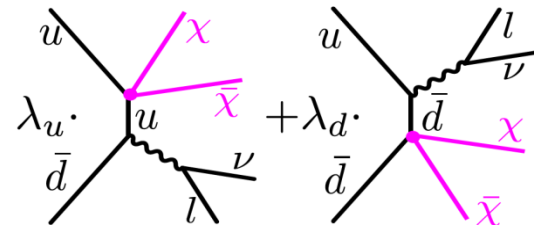
$\sqrt{13}$  and  $\sqrt{14}$  TeV: simplified model (DMV)

implemented in **Madgraph**

Ref: "Missing Energy Signatures of Dark Matter at the LHC" Patrick J. Fox, Roni Harnik, Joachim Kopp, Yuhsin Tsai (2011). arXiv:1109.4398

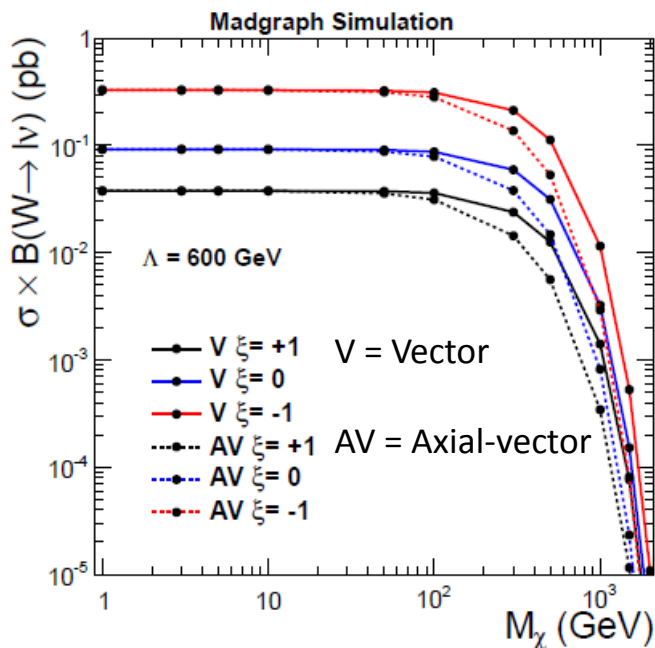
# Mono-Lepton (EFT)

Mono-lepton channel sensitive to possibly different coupling to u- and d-type quarks

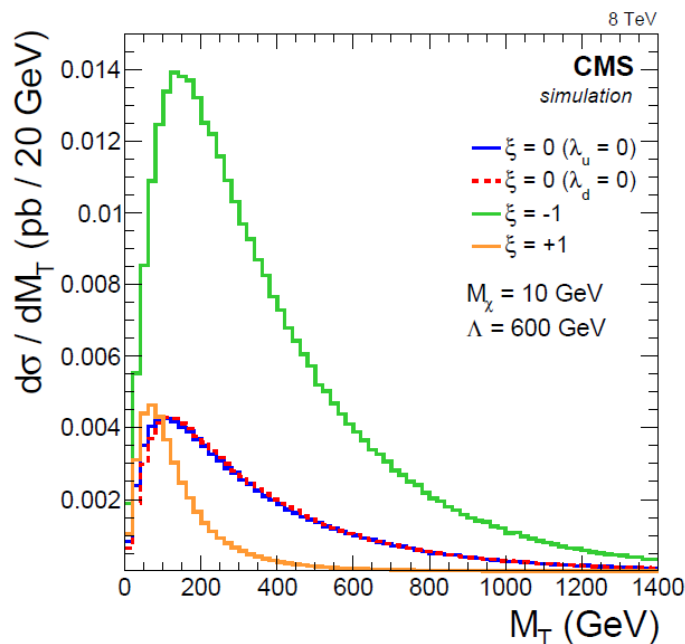


Parametrized by  $\xi = -1, 0, +1$   
(interference)

Madgraph simulation



Largest cross section for  $\xi = -1$   
For  $M_\chi < \sim 70$  GeV same cross section for V and AV coupling of fixed  $\xi$



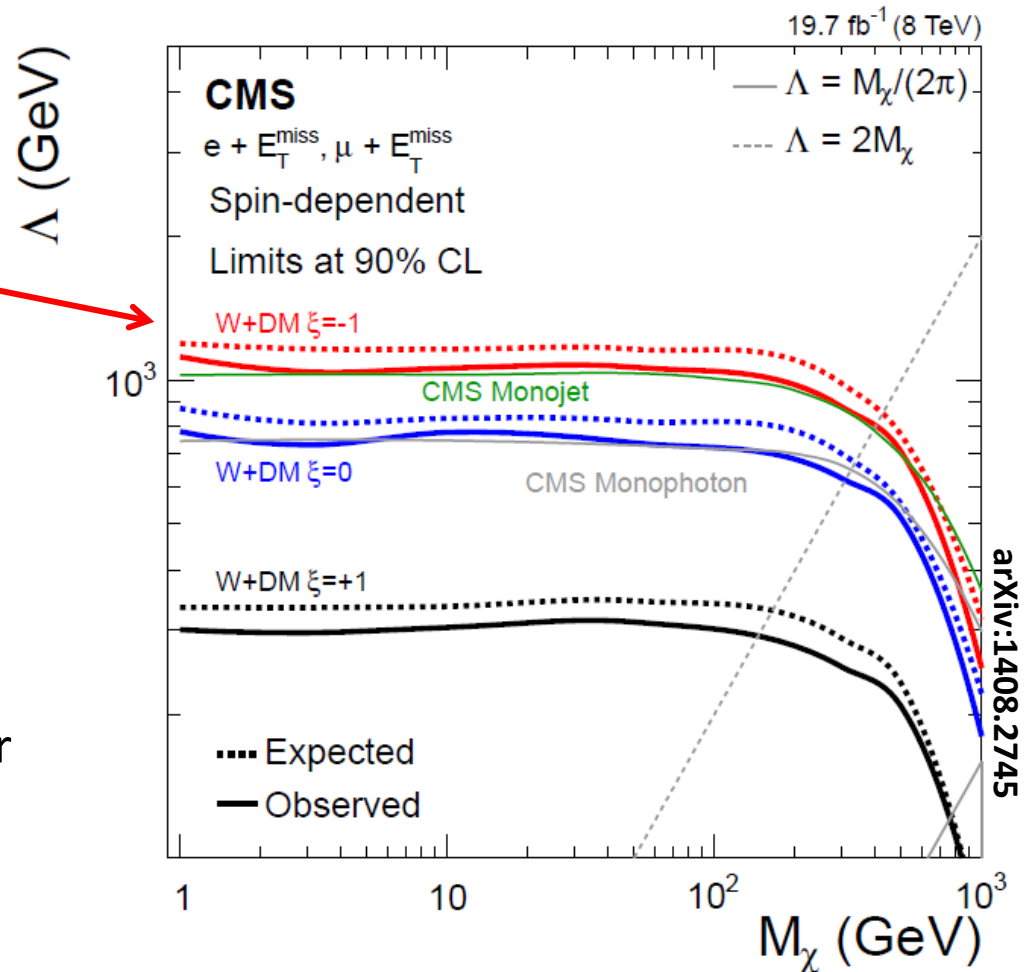
Interference type influences  $M_T$  shape  
→ impact on sensitivity

# Because of $\xi$ always three results for each coupling

For  $\xi = -1$  maximum cross section.  
Sensitivity comparable to monojet.

To reduce no. of parameters, concentrate on  $\xi = -1$  in the following, although  $\xi = +1$  has very different signal shape & sensitivity.

pp cross-section and lambda limits nearly independent on coupling (vector vs axial-vector) due to similar cross section.





# Turning to Simplified Models

## Switch to MG simplified model because it does include interference

(DMV) arXiv:1109.4398

Madgraph-version 2.1.2.

Model version 1.0 from P. Fox, J. Kopp et al.

**Mediator =  $Z'$**

**Four parameters in Madgraph:**

- DM mass  $M_\chi$
- Mediator mass  $M_{\text{Med}}$
- Mediator width  $\Gamma_{\text{med}}$
- Couplings =  $\sqrt{g_{\text{SM}} g_{\text{DM}}}$

DMV implemented in MG as **V-A**.

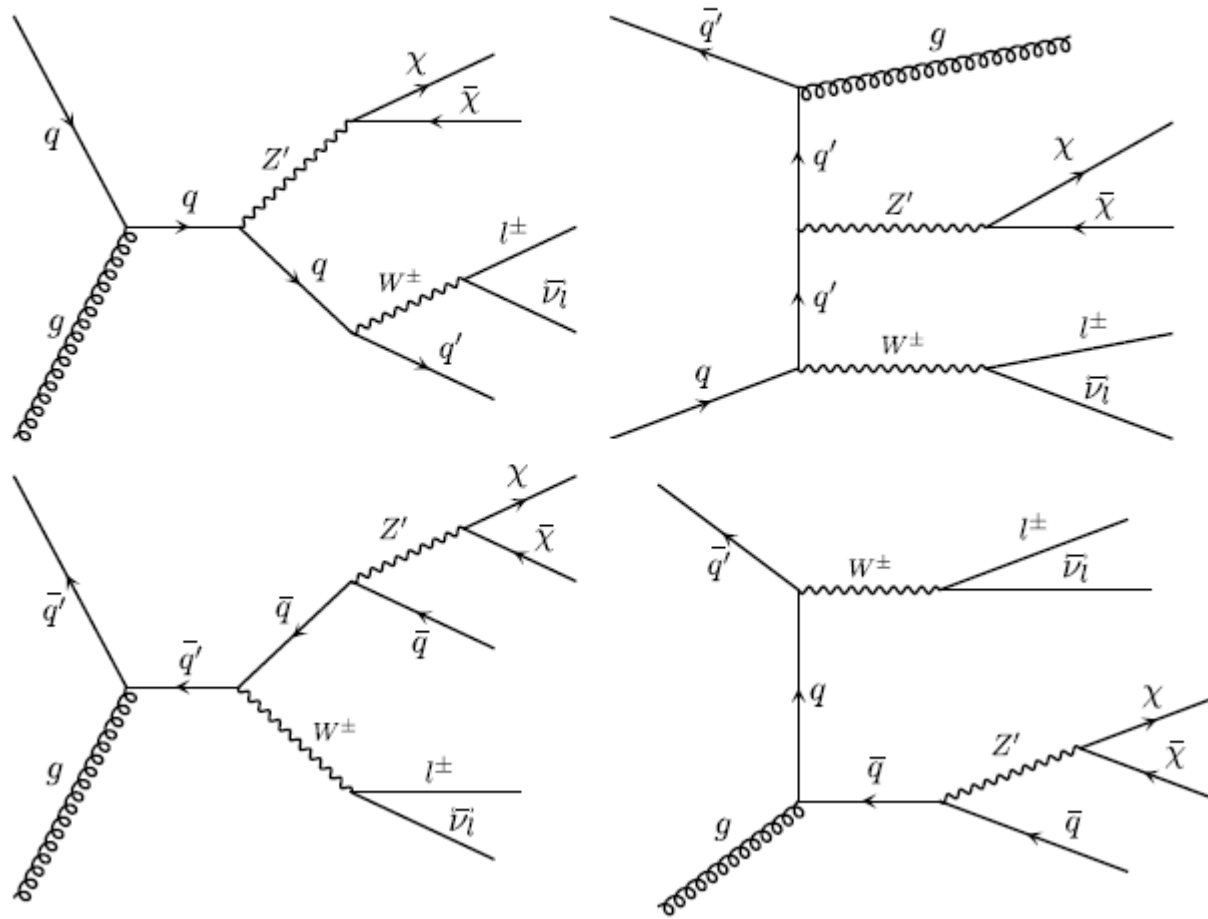
We modified it to either pure V or A. From EFT interpretation we know  $V \approx A$  cross sections very similar. Concentrate on axial-vector coupling where LHC competitive to direct detection for (not for V).

Up to now always used  $g_{\text{DM}} = g_{\text{SM}} = 1$

Discuss width later

# Feynman Graphs (Madgraph)

Feynman graph examples for process  $pp \rightarrow W\chi\bar{\chi}(j), W \rightarrow l\nu_l$ :



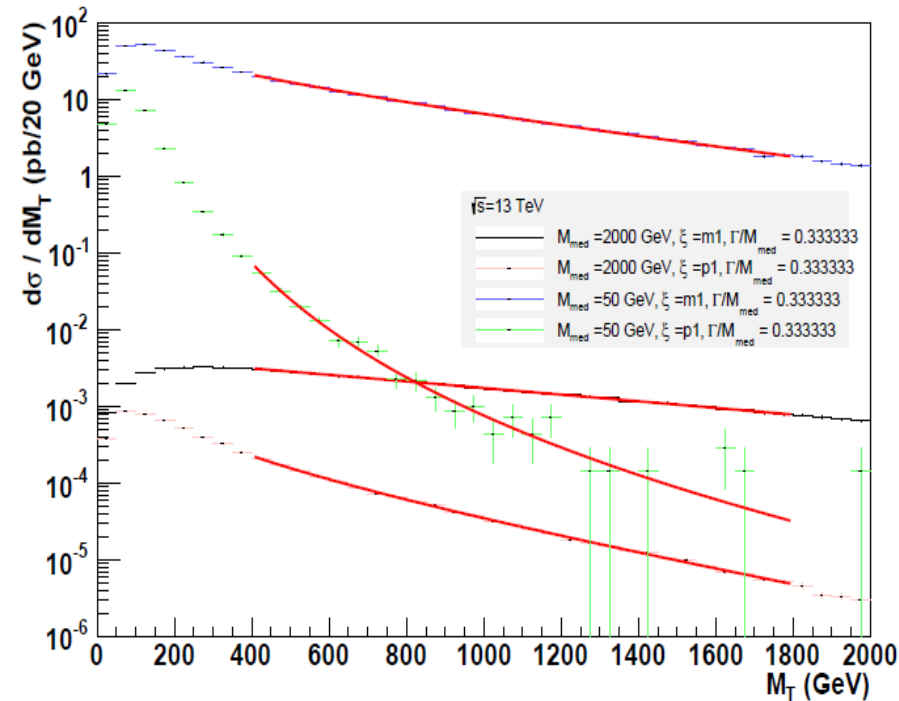
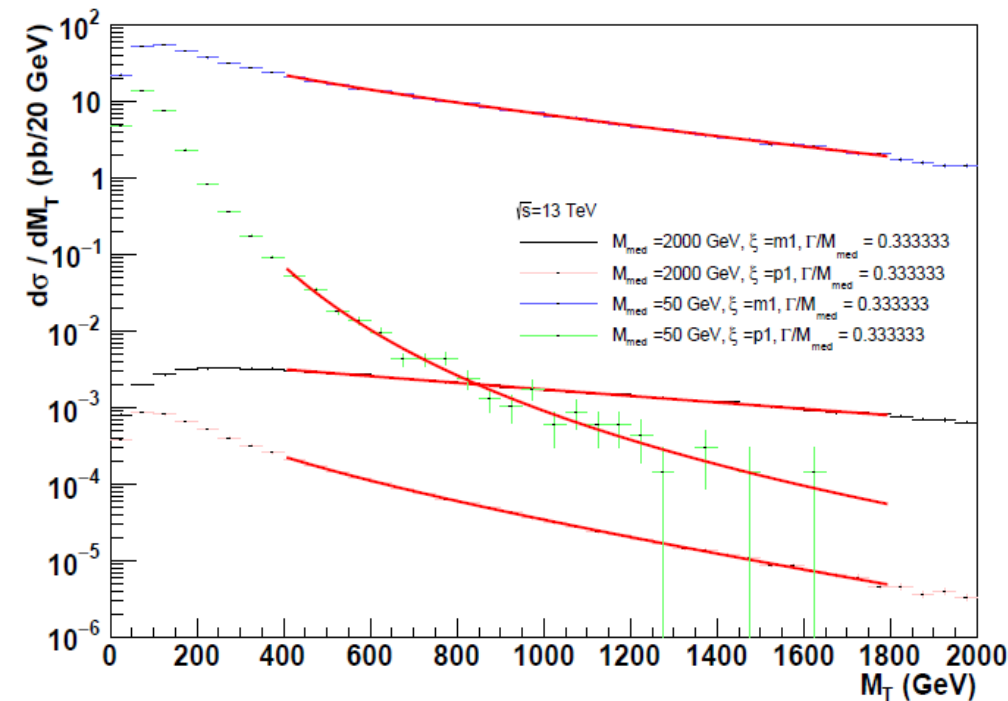
$$l = e, \mu$$

Same model could be used for hadronic W-channels

# Comparison unmodified V-A and modified axial-vector coupling

CMS Simulation - DMV model with V-A coupling

Modified DMV model with axial-vector coupling  
CMS Simulation

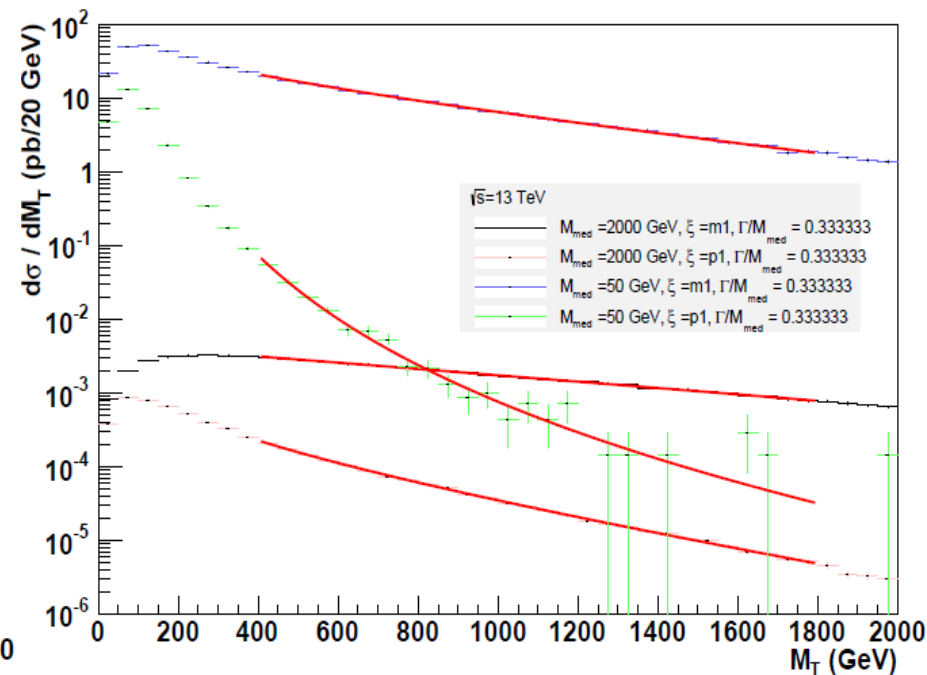
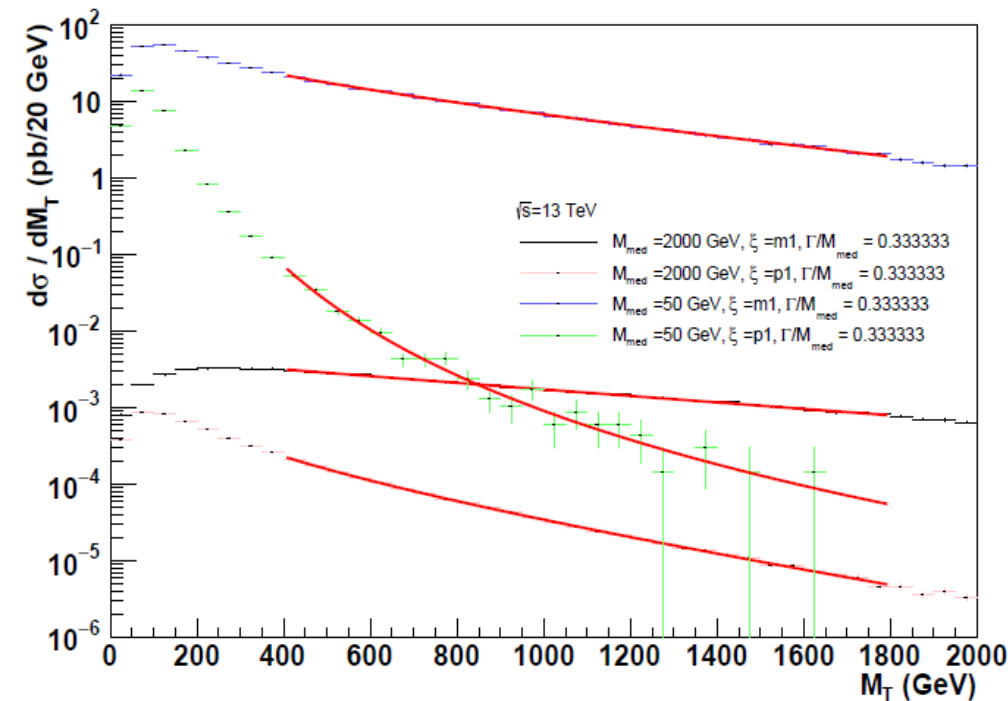


→ Nearly no difference

# Comparison unmodified V-A and modified vector coupling

CMS Simulation - DMV model with V-A coupling

CMS Simulation DMV model with vector coupling



Similarities in xsec for V and AV xsec and  $M_T$  shape (as seen in EFT) still holds. Allows to reduce number of samples to be generated.



cross section values:

$m_{\text{med}}$ (GeV)	$\xi$	$\sigma \times \text{BR}$ (pb), V-A	$\sigma \times \text{BR}$ (pb), axial vector	$\sigma \times \text{BR}$ (pb), vector only
50	-1	524.5	$503.9 \pm 0.2934$	$503.9 \pm 0.2934$
500	-1	-	$5.306 \pm 0.00349$	$5.306 \pm 0.00349$
2000	-1	0.08155	$0.08153 \pm 4.974\text{e-}05$	$0.08153 \pm 4.974\text{e-}05$
50	+1	29.60	$28.92 \pm 0.01883$	$28.92 \pm 0.01883$
500	+1	-	$0.4943 \pm 0.0003043$	$0.4943 \pm 0.0003043$
2000	+1	0.005660	$0.005661 \pm 3.083\text{e-}06$	$0.005661 \pm 3.083\text{e-}06$

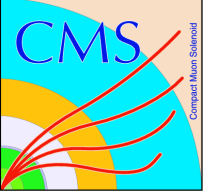
Same xs as well as  $m_T$  spectrum for supposedly pure vector and axial vector coupling after changes to the DMV model!

# Selected parameters for MG5 sample generation

- $\xi = -1, +1$        $\xi = -1$  with max xsec,  $\xi = +1$  with very different MT shape, hence impact on sensitivity.  $\xi = 0$  can be scaled from -1.
- $m_{DM} = 1 \text{ GeV}, 10 \text{ GeV}, 250 \text{ GeV}$
- $g_{DM} = g_{SM} = 1$       Keep SM and DM couplings at 1 for now.
- upper limit on  $\frac{\Gamma}{m_{med}} = \frac{1}{3}$
- AV coupling implemented in DMV model
- $m_{med} = 0.05, 0.5, 2 \text{ TeV}$  (selected from shape analysis parameter plots)

<sup>6</sup>  $\Gamma = M/8\pi$  corresponds to a mediator that can annihilate into only one quark flavor and helicity and has couplings  $g_\chi g_q = 1$ . Since in figure 7 we have assumed couplings to all quark helicities and flavors (collider production is dominated by coupling to up-quarks though), and since  $g_\chi g_q > 1$  in parts of the plot (see dashed contours),  $\Gamma = M/8\pi$  can be regarded as an approximate lower limit on the mediator width.

(rationale for lower limit, arxiv:1109.4398v2)



# Vary Mediator Width

Mediator width has impact on sensitivity. But little impact on  $M_T$  spectrum, no need to vary kinematic selections.

Width can be set in MG. Starting point: varied width between  $\Gamma = M_{\text{med}}/3$  and  $\Gamma = M_{\text{med}}/8\pi$

Alternatively, calculate width from axial couplings formulas:

$$\Gamma(Z' \rightarrow \bar{\chi}\chi)_{\text{vector}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{12\pi} \left(1 + \frac{2m_{\text{DM}}^2}{M_{\text{med}}^2}\right) \sqrt{1 - \frac{4m_{\text{DM}}^2}{M_{\text{med}}^2}}$$

$$\Gamma(Z' \rightarrow \bar{q}q)_{\text{vector}} = \frac{3g_q^2 M_{\text{med}}}{12\pi} \left(1 + \frac{2m_q^2}{M_{\text{med}}^2}\right) \sqrt{1 - \frac{4m_q^2}{M_{\text{med}}^2}}$$

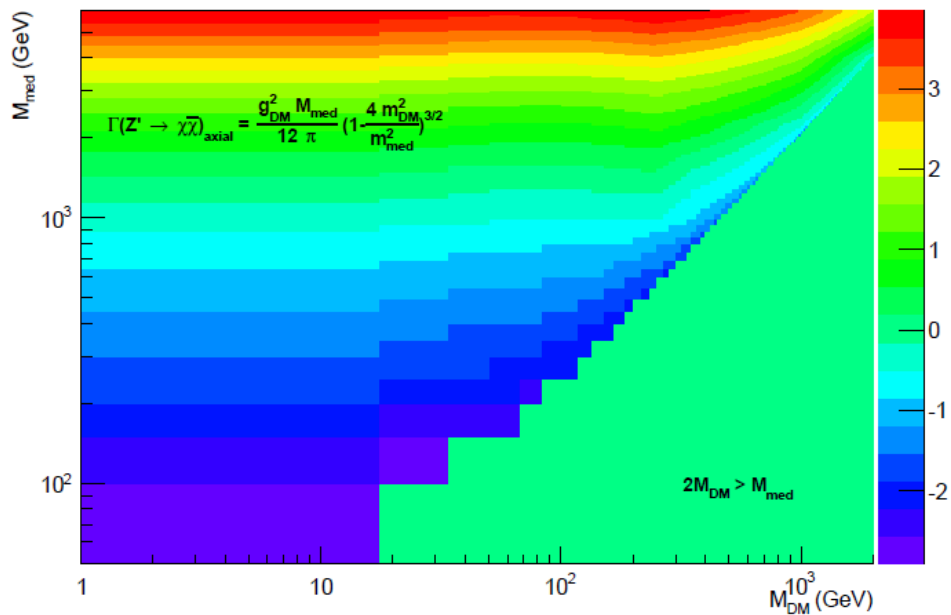
$$\Gamma(Z' \rightarrow \bar{\chi}\chi)_{\text{axial}} = \frac{g_{\text{DM}}^2 M_{\text{med}}}{12\pi} \left(1 - \frac{4m_{\text{DM}}^2}{M_{\text{med}}^2}\right)^{3/2}$$

$$\Gamma(Z' \rightarrow \bar{q}q)_{\text{axial}} = \frac{3g_q^2 M_{\text{med}}}{12\pi} \left(1 - \frac{4m_q^2}{M_{\text{med}}^2}\right)^{3/2}$$

For high  $M_{\text{med}}$  corresponds rather to  $\Gamma \approx M/12\pi$ . Resonance becomes very narrow and sensitivity increases.

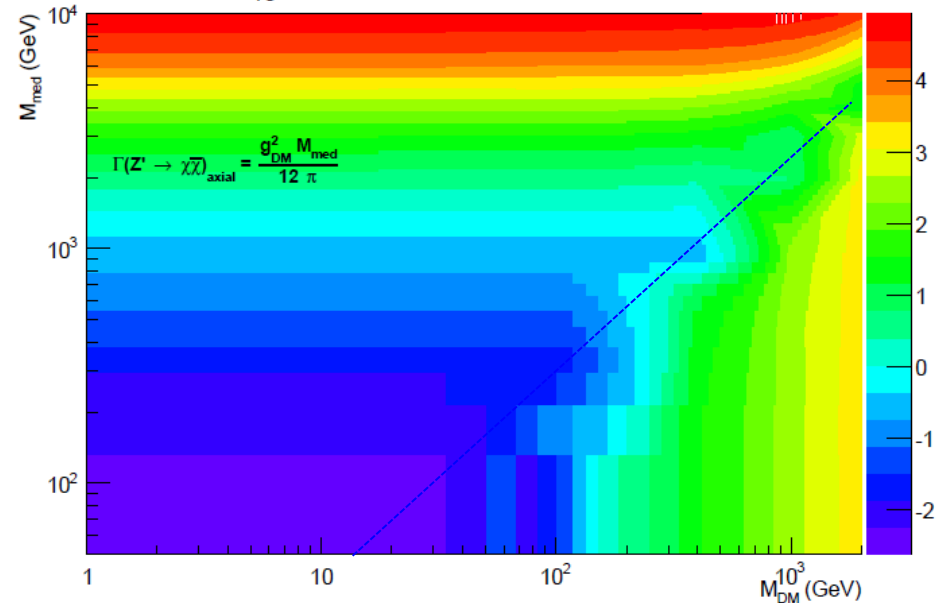
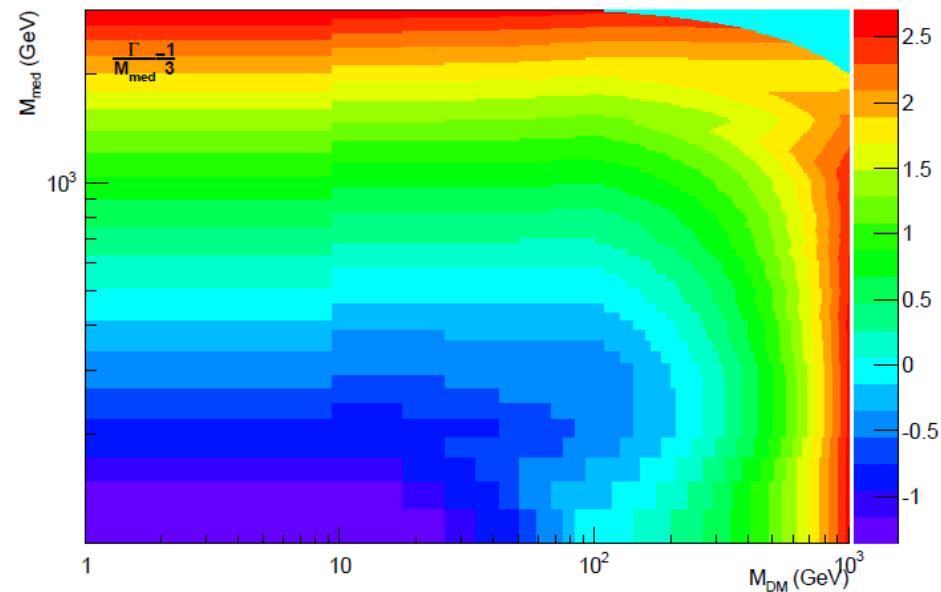
$\log_{10}(\sigma_{\text{EFT}}/\sigma_{\text{FT}}), \sqrt{s}=13 \text{ TeV}, \xi=-1, \text{AV}$

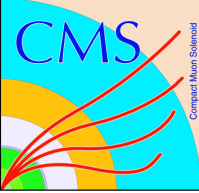
FT = DMV model



$\log_{10}(\sigma_{\text{EFT}}/\sigma_{\text{FT}}), \sqrt{s}=13 \text{ TeV}, \xi=-1, \text{AV}$

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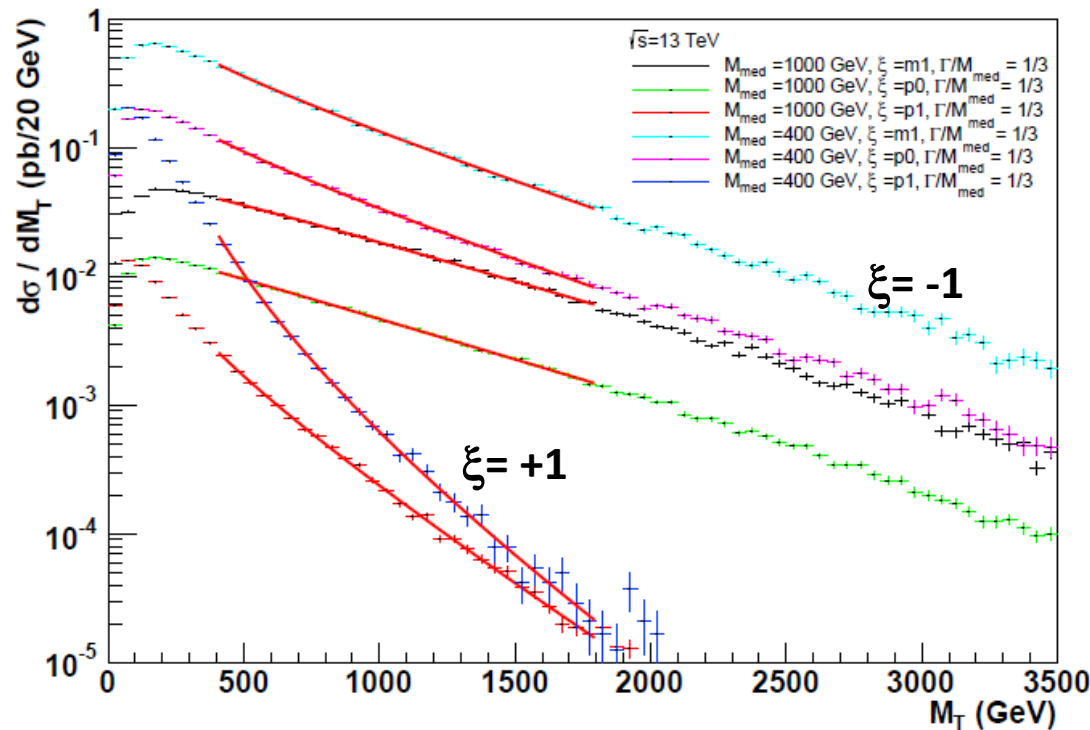




# Implications of Simplified Model on $M_T$ Spectrum

Studied impact of model parameters on transverse mass  $M_T$  spectrum for 13 TeV. Address question if kinematic selections ( $p_T/MET$ ,  $\Delta\phi$ ) have to be re-optimized?

CMS Simulation

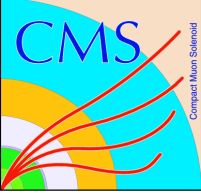


For  $\Gamma/M = 1/3$

See  $\Gamma/M=1/8\pi$  next page

## Conclusions:

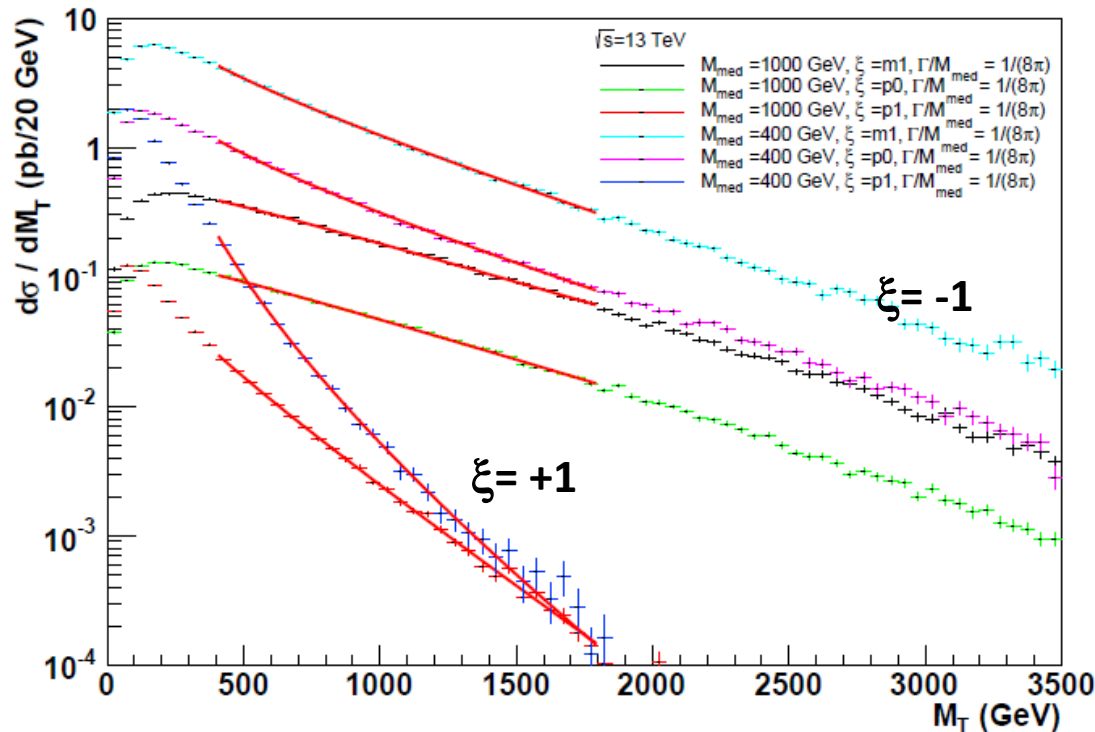
- Mediator width has little impact on  $M_T$  shape
- Mediator mass slightly changes  $M_T$  spectrum
- $M_T$  shape dependence on  $\xi$  as seen in EFT



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CMS Simulation



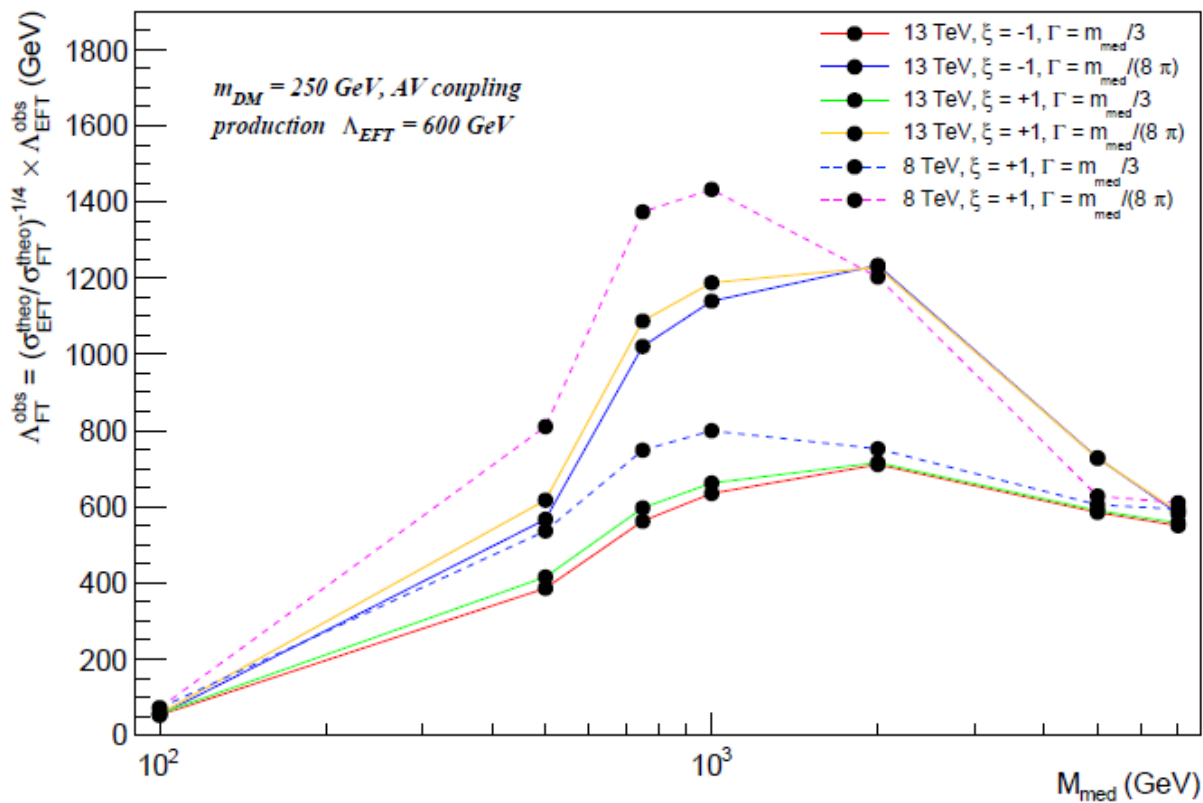
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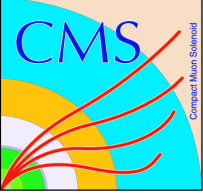
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# Mediator mass scan for $\sqrt{s} = 8, 13$ TeV

Varied width between  $\Gamma = M_{\text{med}}/3$  and  $\Gamma = M_{\text{med}}/8\pi$

$\Lambda$  mediator mass scan





# Plans and some Questions

Next round of data analysis in terms of EFT and DMV simplified model (MG implementation). Use MG as baseline to study interference?

Info of this model on twiki?

Vary the width or use only very narrow one?

Take  $m_{DM}$  and  $M_{med}$  mass points from ATLAS/CMS recommendation.

Need to simulate only  $\xi = -1$  and  $\xi = +1$  where  $M_T$  shape is different.  $\xi = 0$  can be scaled.

Also vary coupling strength (so far used  $g_{SM} = g_{DM} = 1$ ).

Proposed values:  $g_{SM}$  and  $g_{DM} = 0.3, 0.5, 1.0$  and  $1.45$

Test other models with mono-lepton final state discussed in this forum