

# Folded Supersymmetry

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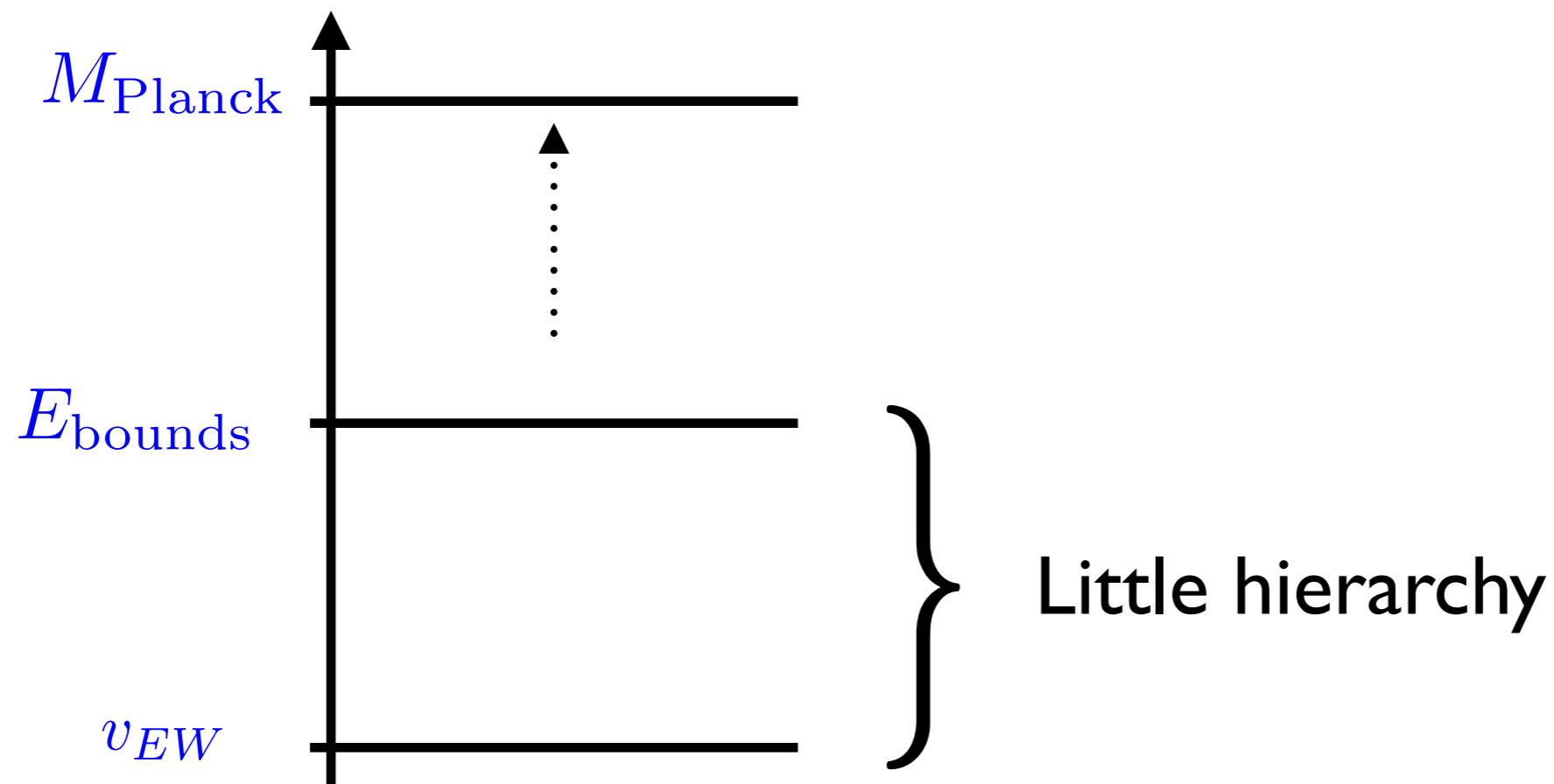
# Outline

- Can we save naturalness ?
- Folded Supersymmetry
- Some Phenomenology
- Outlook

# Naturalness and the LHC

Is the Electroweak scale natural ?

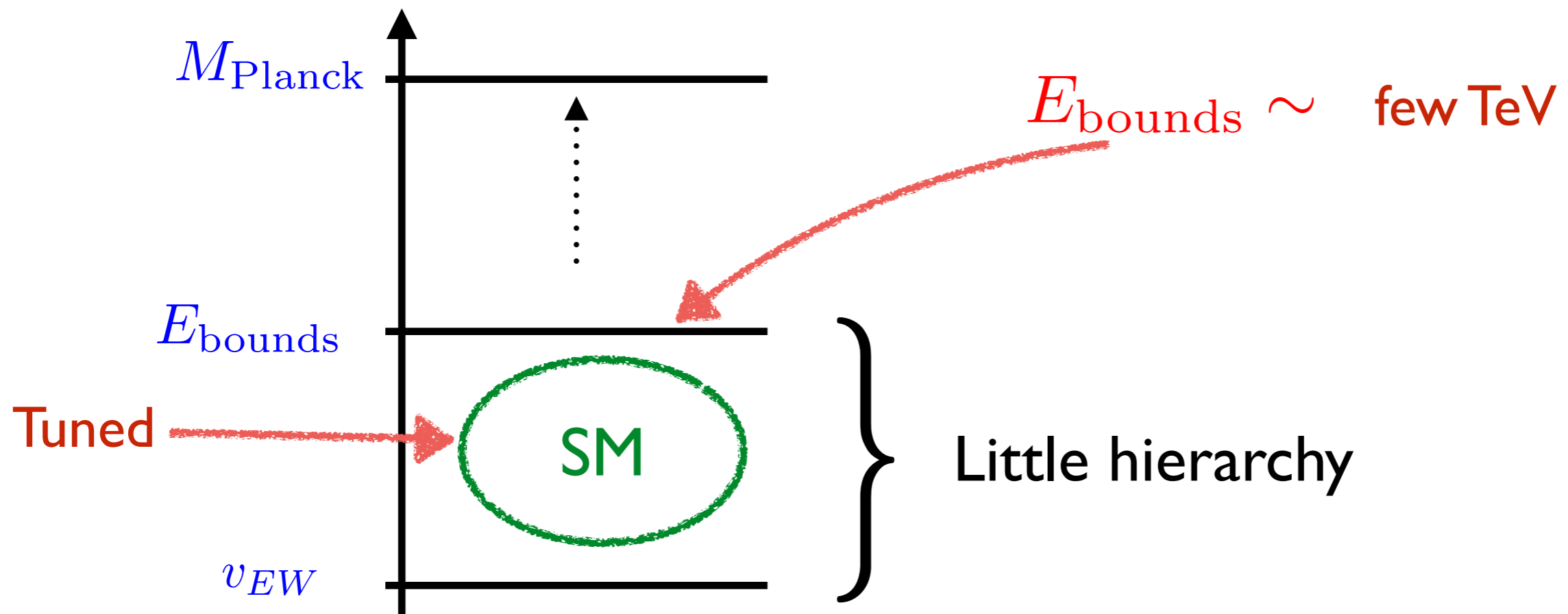
- The LHC found the Higgs
- Plus ... nothing else.



# Naturalness and the LHC

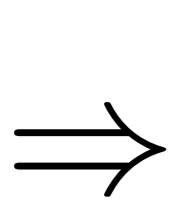
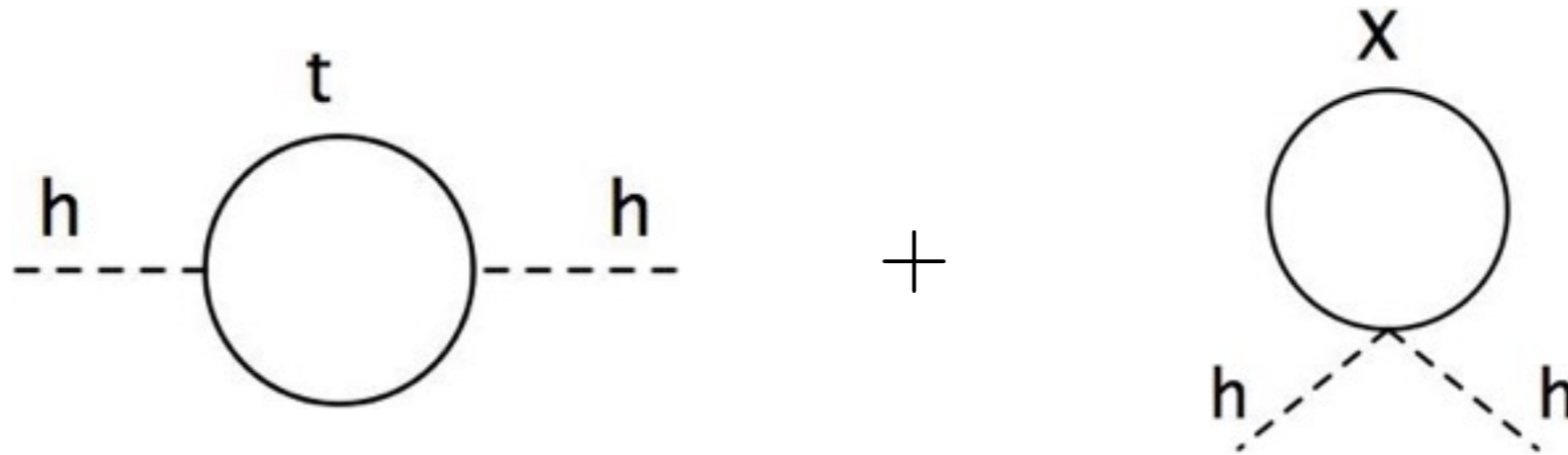
Is the Electroweak scale natural ?

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# Naturalness and the LHC

UV sensitivity of  $m_h^2$  dominated by top quark



- Top partners  $X$  carry color
- Easily produced at the LHC

# Colorless Top Partners

Last refuge of naturalness ?

Top partners need not carry color

If symmetry protecting  $m_h^2$  does not commute with  $SU(3)_c$

Exchanges  $SU(3)_c \rightarrow SU(3)'$

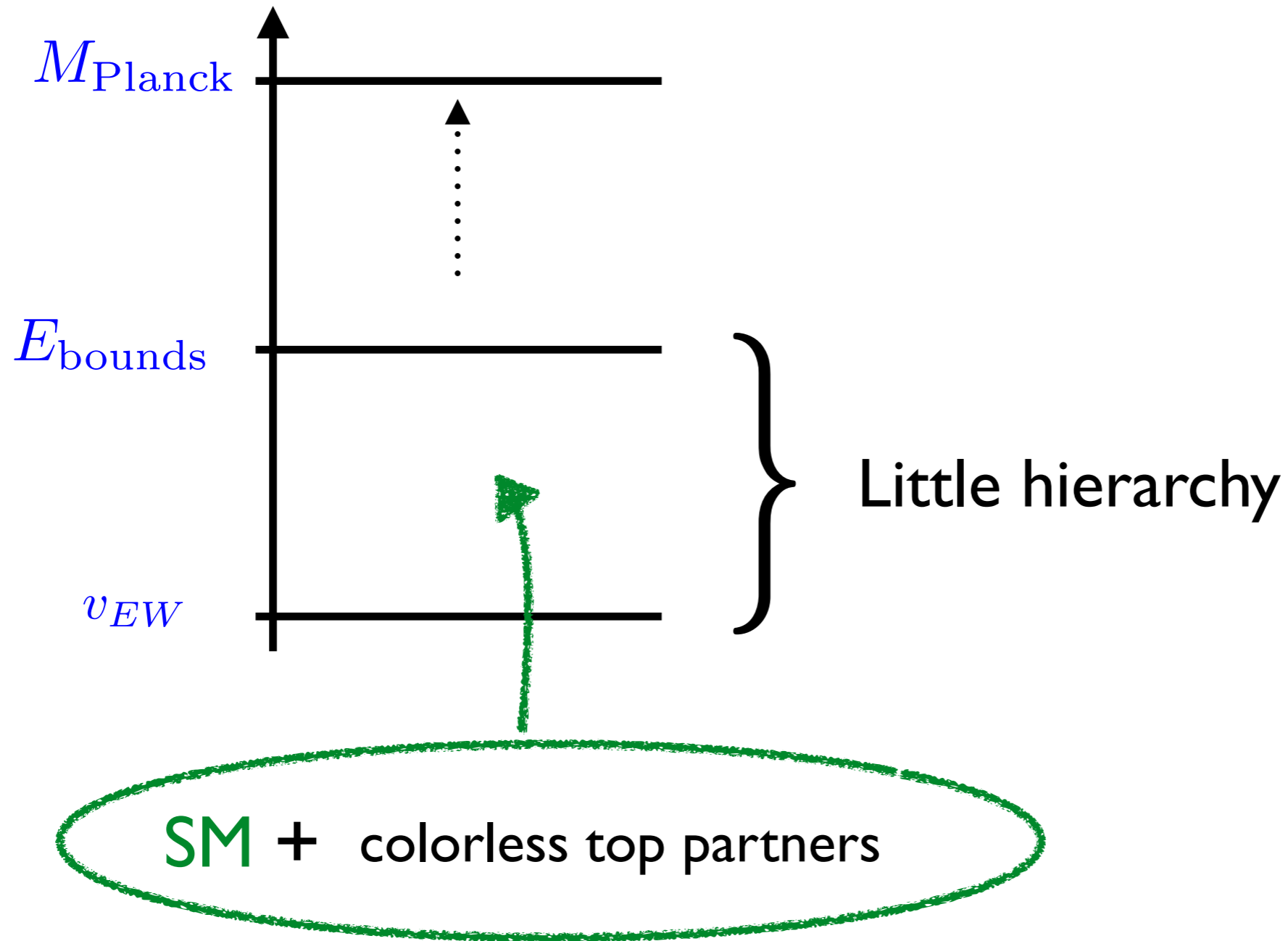
The X's are charged under  $SU(3)'$

Bounds on  $m_X$  not as stringent

Colorless X models are more natural

# Colorless Top Partners

General idea To solve the Little Hierarchy problem



# Colorless Top Partners

## Ingredients for neutral naturalness

- Symmetry protecting the Higgs: spontaneously broken global symmetry, SUSY, ...
- Extend the color gauge symmetry to have at least  $[SU(3)]^2$
- Either impose a discrete symmetry or orbifold  
In general, CTP theories can be obtained from orbifolding

*N.Craig, S.Knapen, P. Longhi, 1410.6806, 1411.7393*



# Folded Supersymmetry

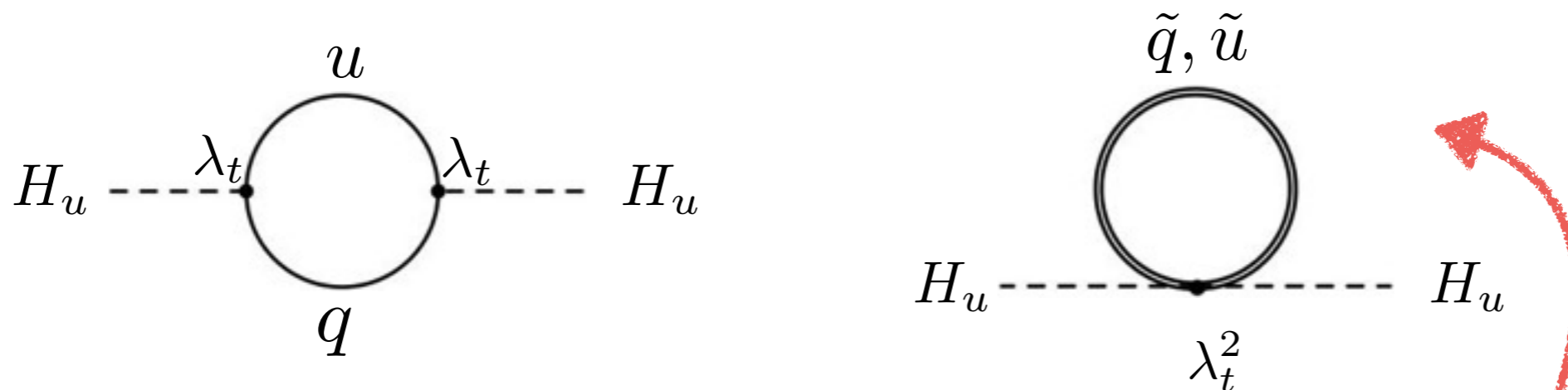
w/ Z.Chacko, H. Goh, R. Harnik , hep-ph/0609152

- Squarks need to be charged under  $SU(2)_L$
- Need not be charged under  $SU(3)_c$

# Folded Supersymmetry

w/ Z.Chacko, H. Goh, R. Harnik , hep-ph/0609152

- Squarks need to be charged under  $SU(2)_L$
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Can we have SUSY with colorless stops ?

# Folded Supersymmetry

## Ingredients

- Supersymmetry
- Extended gauge symmetry:

$$G \supset SU(3)_A \times SU(3)_B$$

- Break SUSY at a high scale to obtain at low energies

$$q_A, \tilde{q}_B$$

Accidental low energy SUSY cancels quadratic div. at one loop

# Large N Orbifold Correspondence

*S.Kachru, E. Silverstein, hep-th/9802183; M.Bershadsky, A.Johansen, hep-th/9803248;  
M.Schmaltz, hep-th/9805218*

Parent theory  $\xrightarrow{\text{Orbifold}}$  Daughter theory

$$(\text{C.F.})_{\text{Parent}} = (\text{C.F.})_{\text{Daughter}} + O\left(\frac{1}{N}\right)$$

# Bifold Protection

Global  $U(N)$        $\lambda S Q_i \bar{Q}_i$        $i = 1 \dots, N$

$M_S^2$  is quadratically divergent

- Supersymmetrize
- Duplicate index running in loop:  $i = 1, \dots, 2N$

# Bifold Protection

- Define

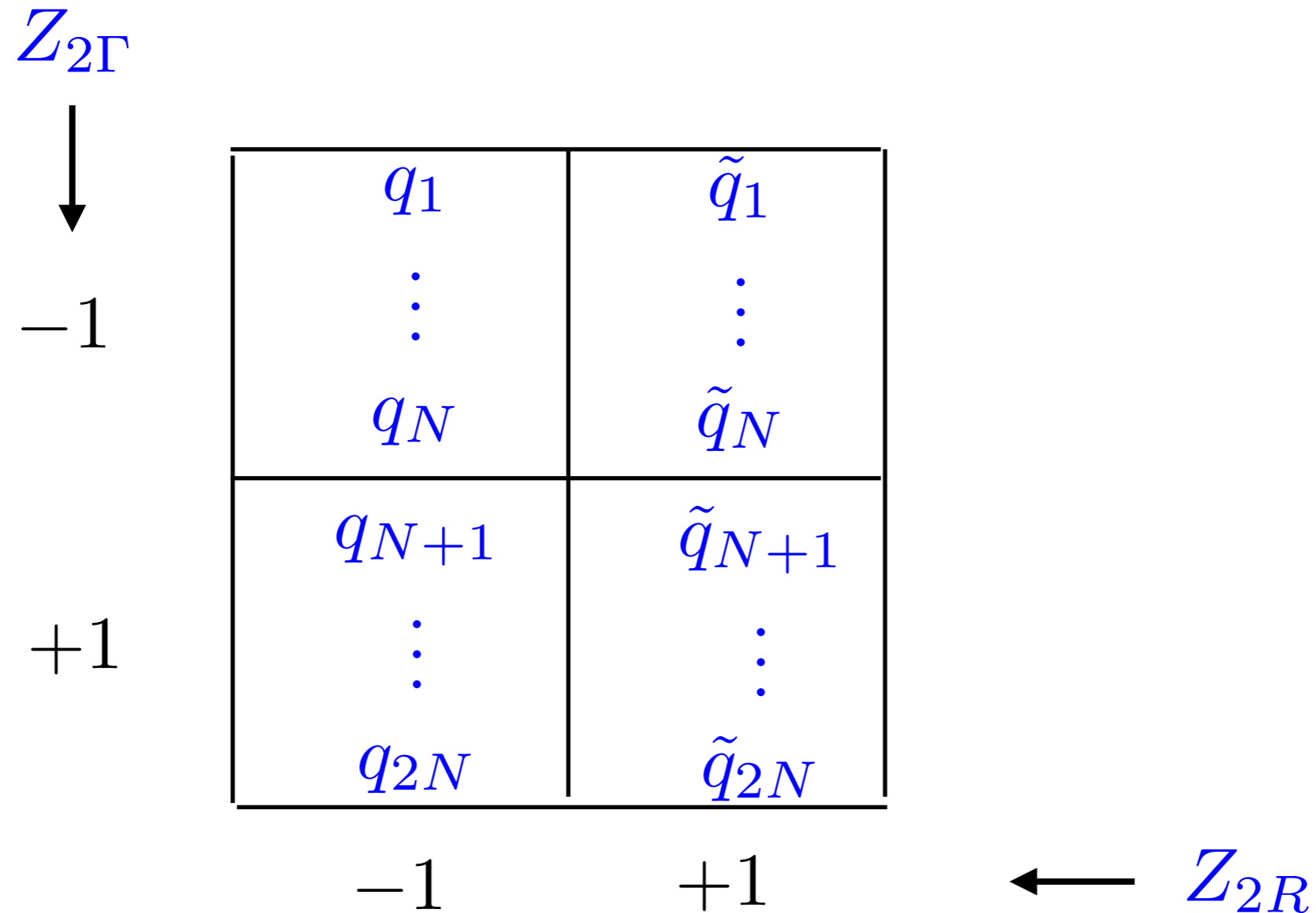
$$\Gamma = \begin{pmatrix} +1 & & & & & & \\ & \ddots & & & & & \\ & & +1 & & & & \\ & & & -1 & & & \\ & & & & \ddots & & \\ & & & & & & -1 \end{pmatrix}$$
$$\begin{matrix} 1 \\ \vdots \\ N \\ N+1 \\ \vdots \\ 2N \end{matrix}$$

- Theory is invariant under

$$Z_{2\Gamma} \left\{ \begin{array}{l} S \rightarrow S \\ Q_i \rightarrow -\Gamma Q_i \\ Q_i \rightarrow -\Gamma^* \bar{Q}_i \end{array} \right. \qquad Z_{2R} \left\{ \begin{array}{ll} \text{fermions} & \text{odd} \\ \text{bosons} & \text{even} \end{array} \right.$$

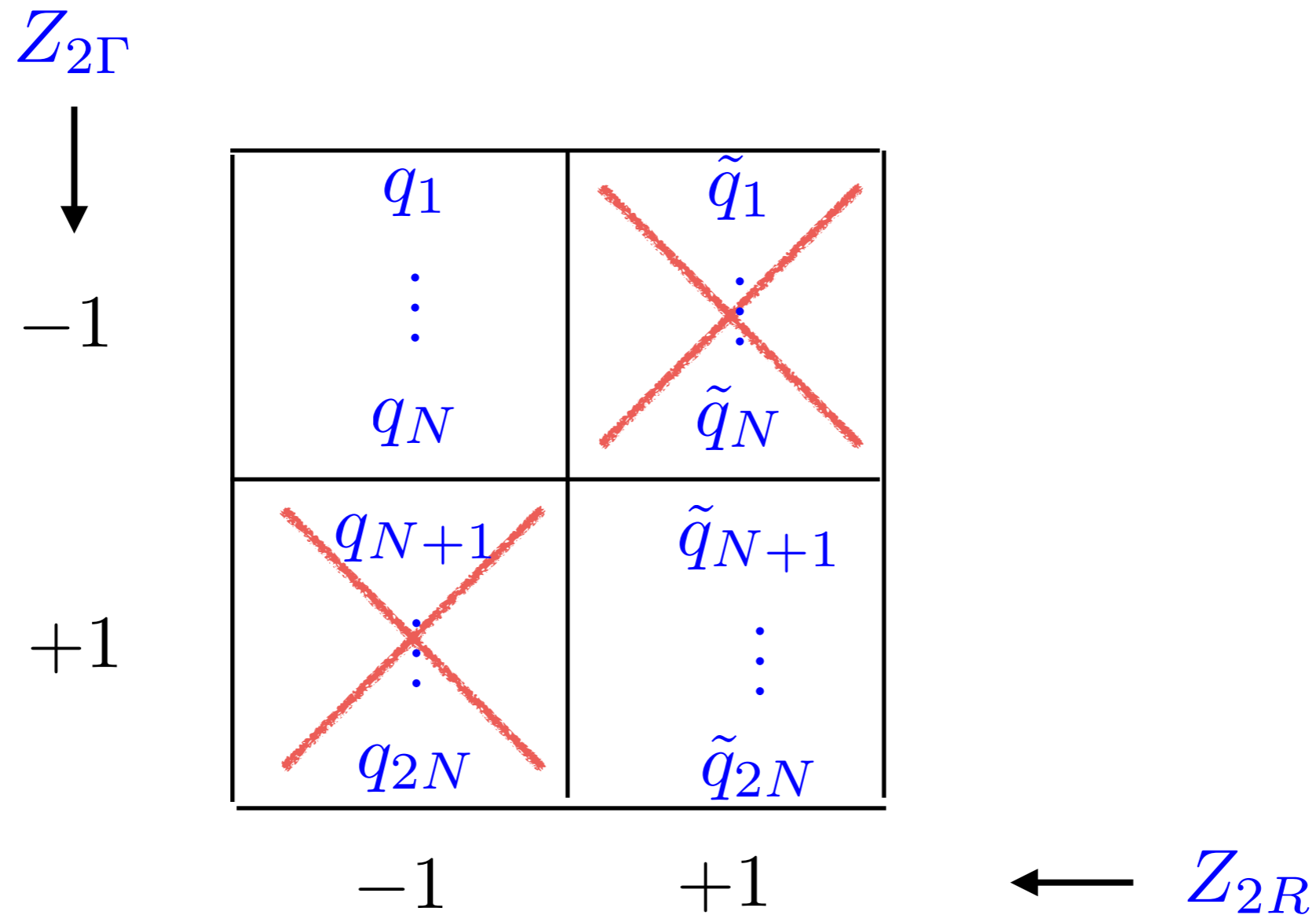
# Bifold Protection

Orbifold projection: Project out states odd under  $Z_{2\Gamma} \times Z_{2R}$



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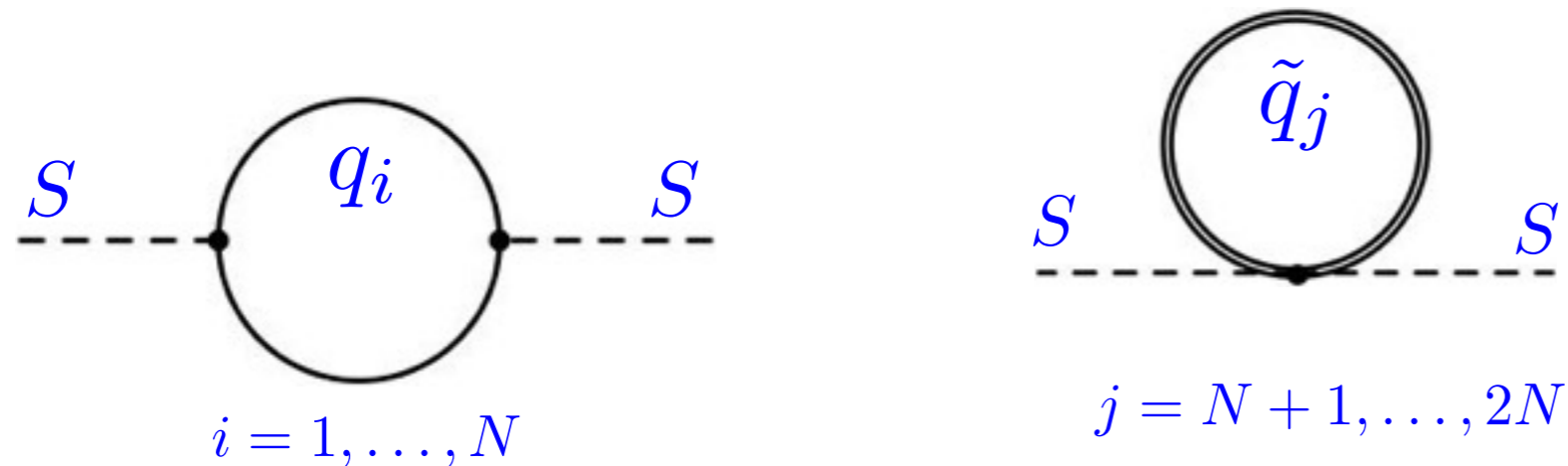


# Bifold Protection

Accidental SUSY: spectrum not supersymmetric

$$\begin{array}{ccc} q_1 & & \tilde{q}_{N+1} \\ \vdots & + & \vdots \\ q_N & & \tilde{q}_{2N} \end{array}$$

But still cancels one-loop quadratic divergence



# Folded SUSY Model

Extend color  $SU(3) \longrightarrow SU(3)_A \times SU(3)_B \times Z_2$

Orbifold so that:

$q_A, u_A + \tilde{q}_B, \tilde{u}_B$  remain in the spectrum

No gauginos

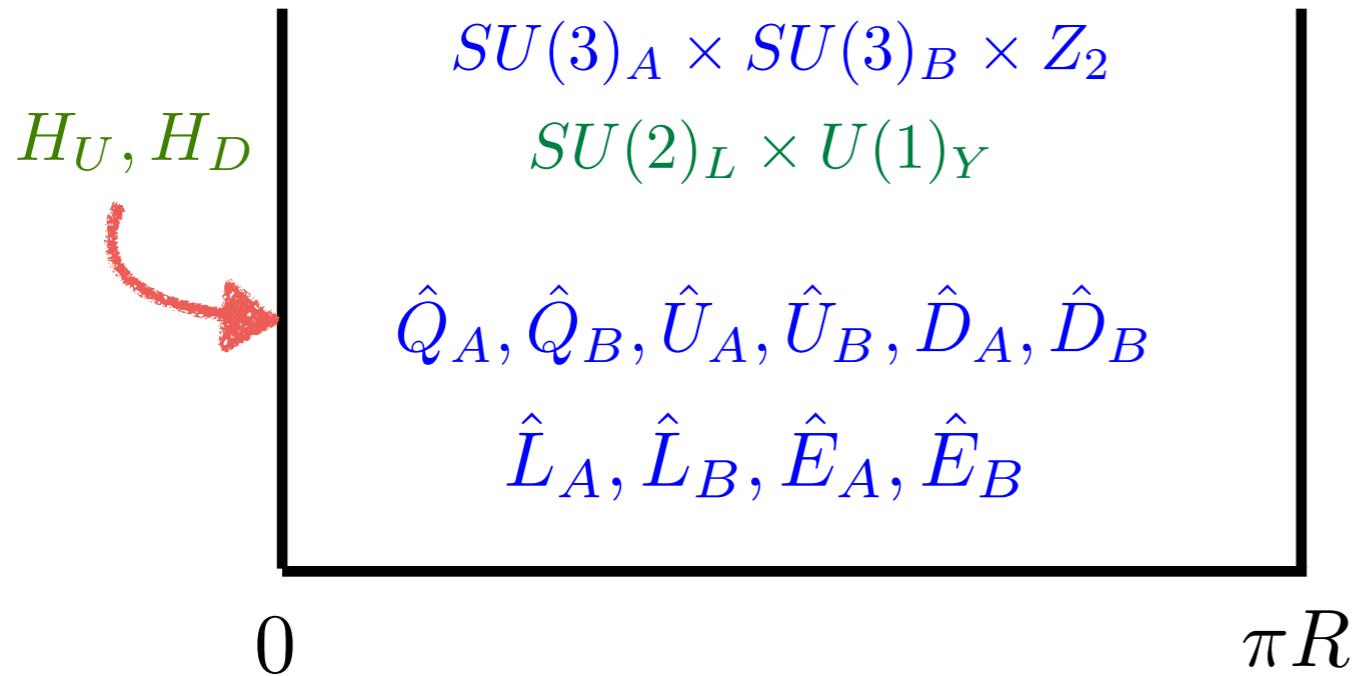
Yukawas obey

$$(\lambda_t h_u q_A u_A + \text{h.c.}) + \lambda_t^2 |\tilde{q}_B h_u|^2 + \lambda_t^2 |\tilde{u}_B|^2 |h_u|^2$$

$\longrightarrow$  Accidental SUSY still protects  $m_h^2$

# Folded SUSY UV Completion

Can be realized in 5D compactified on  $S_1/Z_2$



SUSY broken by BCs  
(Scherk-Schwarz)

BCs break  $Z_2$  at  $\pi R$

# Folded SUSY UV Completion

- $\mathcal{N} = 1$  in 5D  $\longrightarrow$   $\mathcal{N} = 2$  in 4D
- 5D Hypermultiplet  $\hat{Q} \sim (Q, Q^c)$  in terms of 4D superfields

$$Q = (\tilde{q}, q) \quad Q^c = (\tilde{q}^c, q^c)$$

or

$$Q' = (\tilde{q}^{*c}, q) \quad Q'^c = (-\tilde{q}^*, q^c)$$

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$$\left. \begin{array}{l}
 Q = (\tilde{q}, q) \quad Q^c = (\tilde{q}^c, q^c) \quad \longleftarrow \text{at } y = 0 \\
 \text{or} \\
 Q' = (\tilde{q}^{*c}, q) \quad Q'^c = (-\tilde{q}^*, q^c) \quad \longleftarrow \text{at } y = \pi R
 \end{array} \right\} \mathcal{N} = 1$$

related by  $SU(2)_R$  rotation parametrized by  $\alpha$

in this example  $\alpha = \frac{1}{2}$

(more in David Pinner's talk)

# Folded SUSY Spectrum

$$\hat{Q}_A \sim \left\{ \begin{array}{ll} (Q_A, Q_A^c) & +, - \text{ at } y = 0 \\ (Q'_A, Q'^c_A) & +, - \text{ at } y = \pi R \end{array} \right. \quad \text{fermion zero mode}$$

$$\hat{Q}_B \sim \left\{ \begin{array}{ll} (Q_B, Q_B^c) & +, - \text{ at } y = 0 \\ (Q'_B, Q'^c_B) & -, + \text{ at } y = \pi R \end{array} \right. \quad \text{scalar zero mode}$$

BCs break  $Z_2$  (A  $\longleftrightarrow$  B) at  $y = \pi R$

# Folded SUSY Spectrum

- Fermion zero modes from

$$\hat{Q}_A, \hat{U}_A, \hat{D}_A, \hat{L}_A, \hat{E}_A$$

- Scalar zero modes from

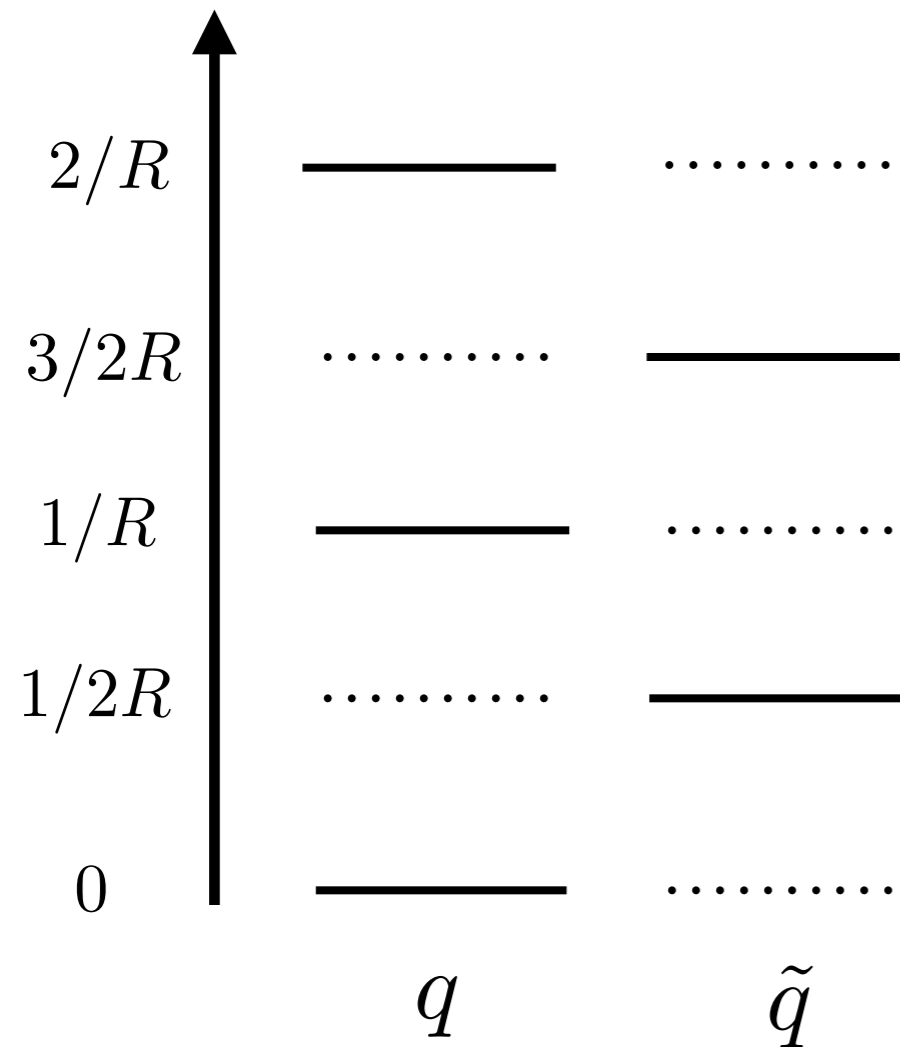
$$\hat{Q}_B, \hat{U}_B, \hat{D}_B, \hat{L}_B, \hat{E}_B$$

- Localize Higgses at  $y = 0$

$$\delta(y) \lambda_t \{Q_{3A} H_U U_{3A} + Q_{3B} H_U U_{3B}\}$$

→ generates desired Yukawas at low energies

# Folded SUSY Spectrum



Top contributions cancel in all  
KK levels at one loop

But weak gauge boson loops give (finite) contributions



# Folded SUSY Spectrum

Zero-mode Folded sfermions: *A. Delgado, A. Pomarol, M Quiros, hep-ph/9812489*

$$m_Q^2 = K \frac{1}{4\pi^4} \left( \frac{4}{3}g_3^2 + \frac{3}{4}g_2^2 + \frac{1}{36}g_1^2 \right) \frac{1}{R^2}$$

$$m_U^2 = K \frac{1}{4\pi^4} \left( \frac{4}{3}g_3^2 + \frac{4}{9}g_1^2 \right) \frac{1}{R^2}$$

$$m_D^2 = K \frac{1}{4\pi^4} \left( \frac{4}{3}g_3^2 + \frac{1}{9}g_1^2 \right) \frac{1}{R^2}$$

$$m_L^2 = K \frac{1}{4\pi^4} \left( \frac{3}{4}g_2^2 + \frac{1}{4}g_1^2 \right) \frac{1}{R^2}$$

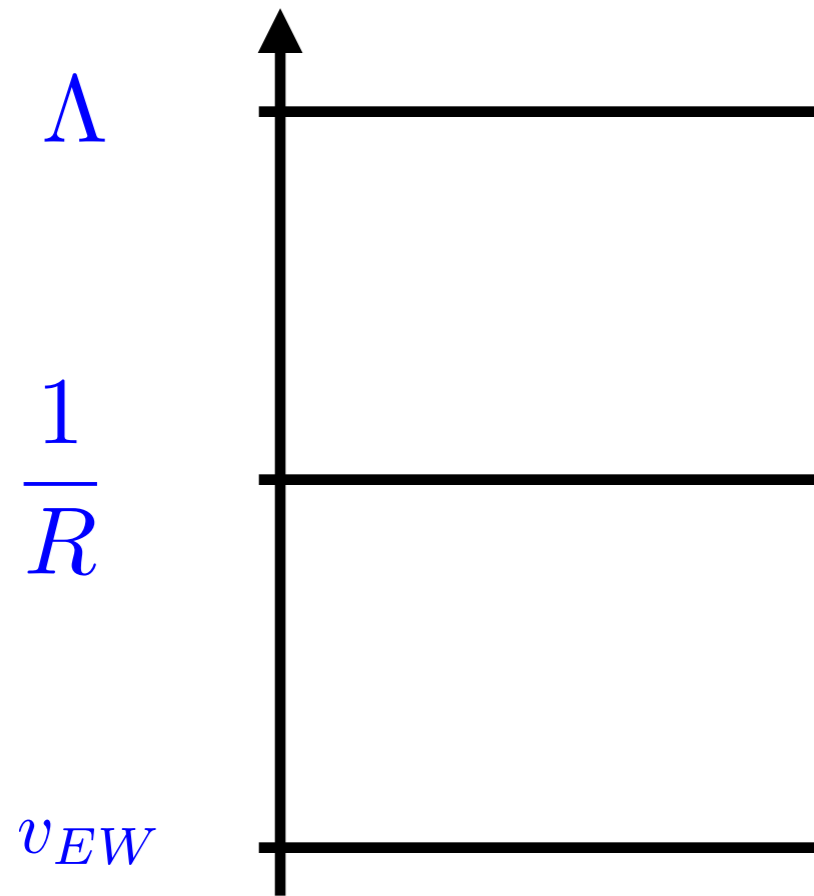
$$m_E^2 = K \frac{1}{4\pi^4} g_1^2 \frac{1}{R^2}$$

plus Yukawa contributions for 3rd generation

$$m_{Q_3}^2 = K \frac{\lambda_t^2}{8\pi^4} \frac{1}{R^2}$$

$$m_{U_3}^2 = K \frac{\lambda_t^2}{4\pi^4} \frac{1}{R^2}$$

# Folded SUSY



5D SUSY

Accidental SUSY

$$\Lambda R \lesssim 4$$

$$\frac{1}{R} \sim (5 - 7) \text{ TeV}$$

# Folded SUSY Model Building

Many important issues:

- Smallish  $\Lambda R$  forced by top Yukawa volume suppression  
→ Potentially dangerous brane kinetic terms at  $y = \pi R$

Requires additional discrete symmetry at  $y = \pi R$

- Electroweak symmetry breaking ?

$$(\delta m_H^2)_{\text{gauge}} \simeq \frac{K}{16\pi^4} (3g_2^2 + g_1^2) \frac{1}{R^2} \quad \text{from gauge loops}$$

$$(\delta m_H^2)_{\text{top}} \simeq -\frac{\lambda_t^2}{4\pi^2} M_{\tilde{t}}^2 \log \left( \frac{1}{RM_{\tilde{t}}} \right) \quad \text{from top sector 2 loops}$$

Does this work ? Probably not. See David Pinner's talk.

# Folded SUSY Model Building

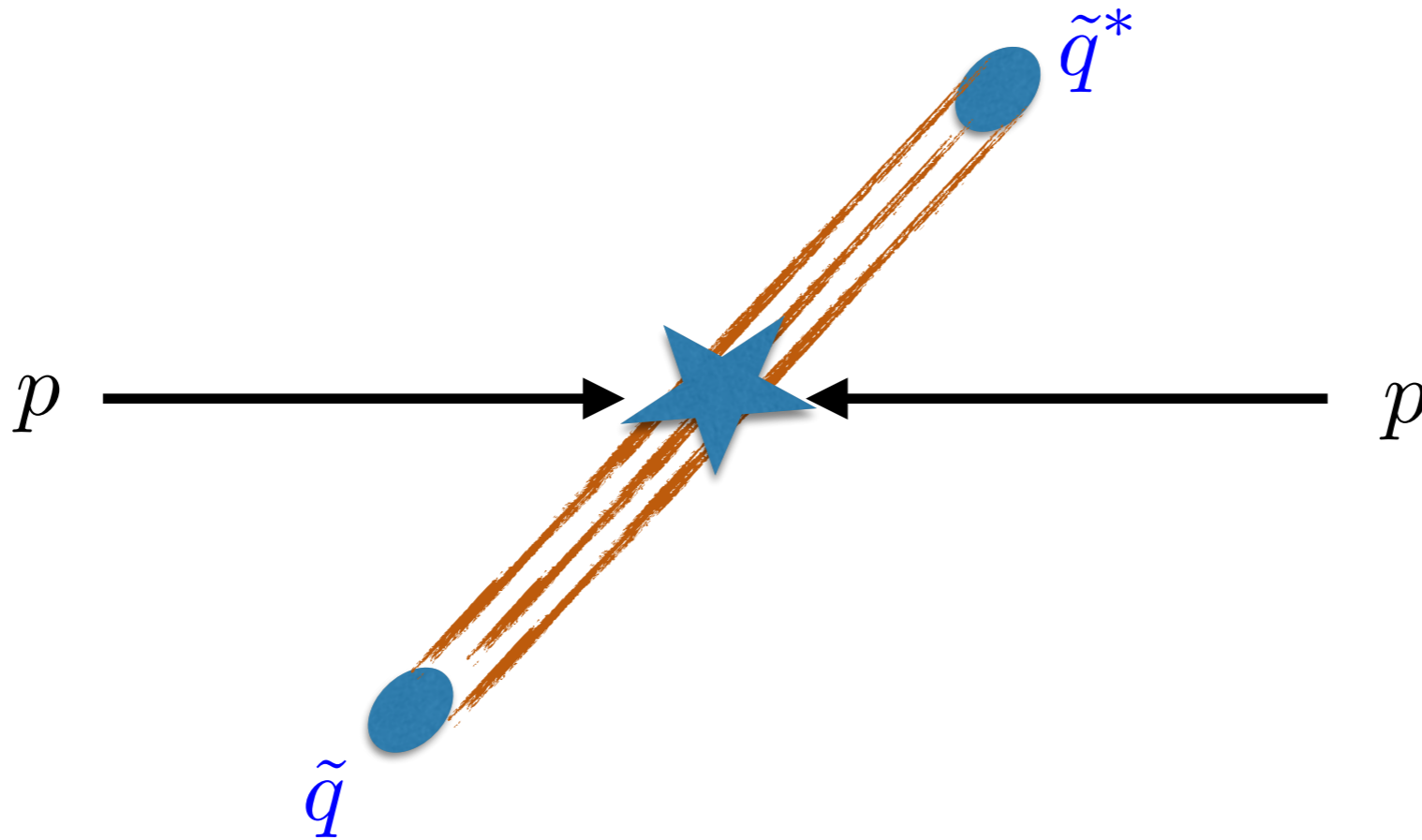
- Just as the MSSM, Folded MSSM needs help for  $m_h = 125$  GeV  
E.g.: non-decoupling D terms from a  $U(1)_X$

# Folded SUSY Signals

# Folded SUSY Signals at the LHC

Electroweak pair production of F-squarks

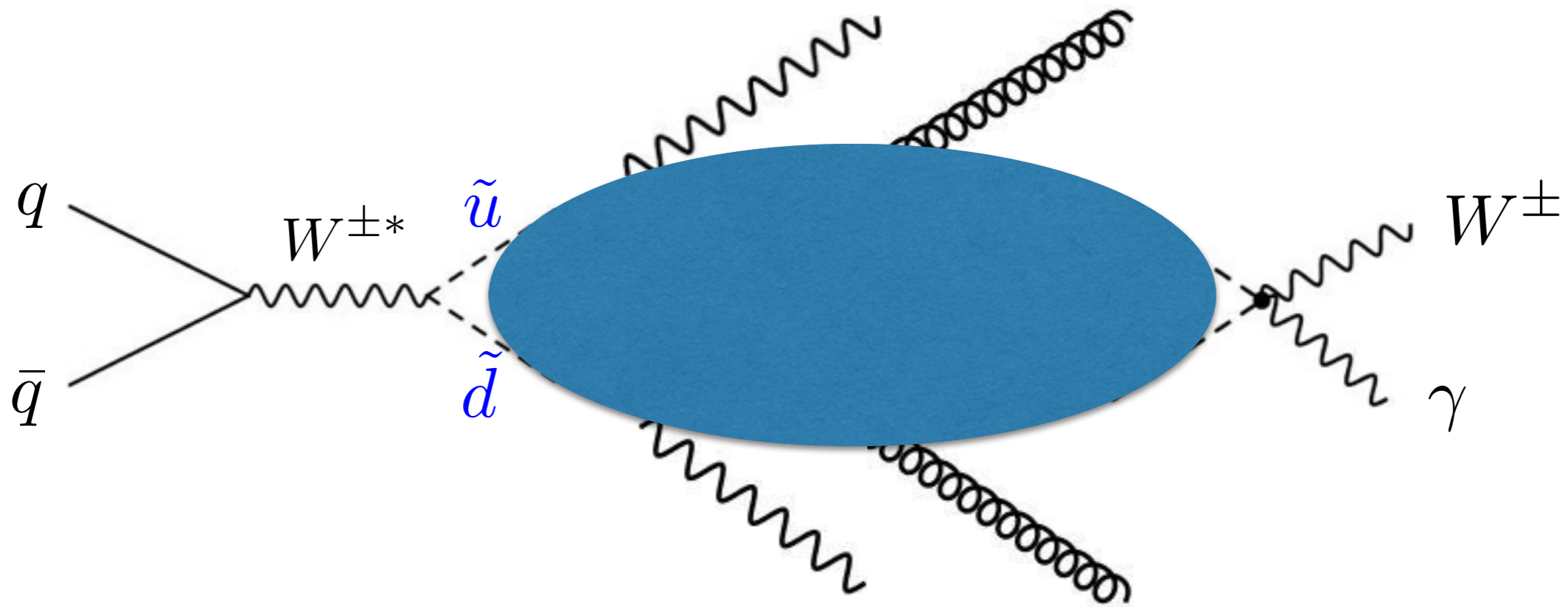
But  $m_T \gg \Lambda'_{QCD} \simeq \text{few GeV} \longrightarrow$  they do not hadronize



$\longrightarrow$  “squirks” have to come back for annihilation

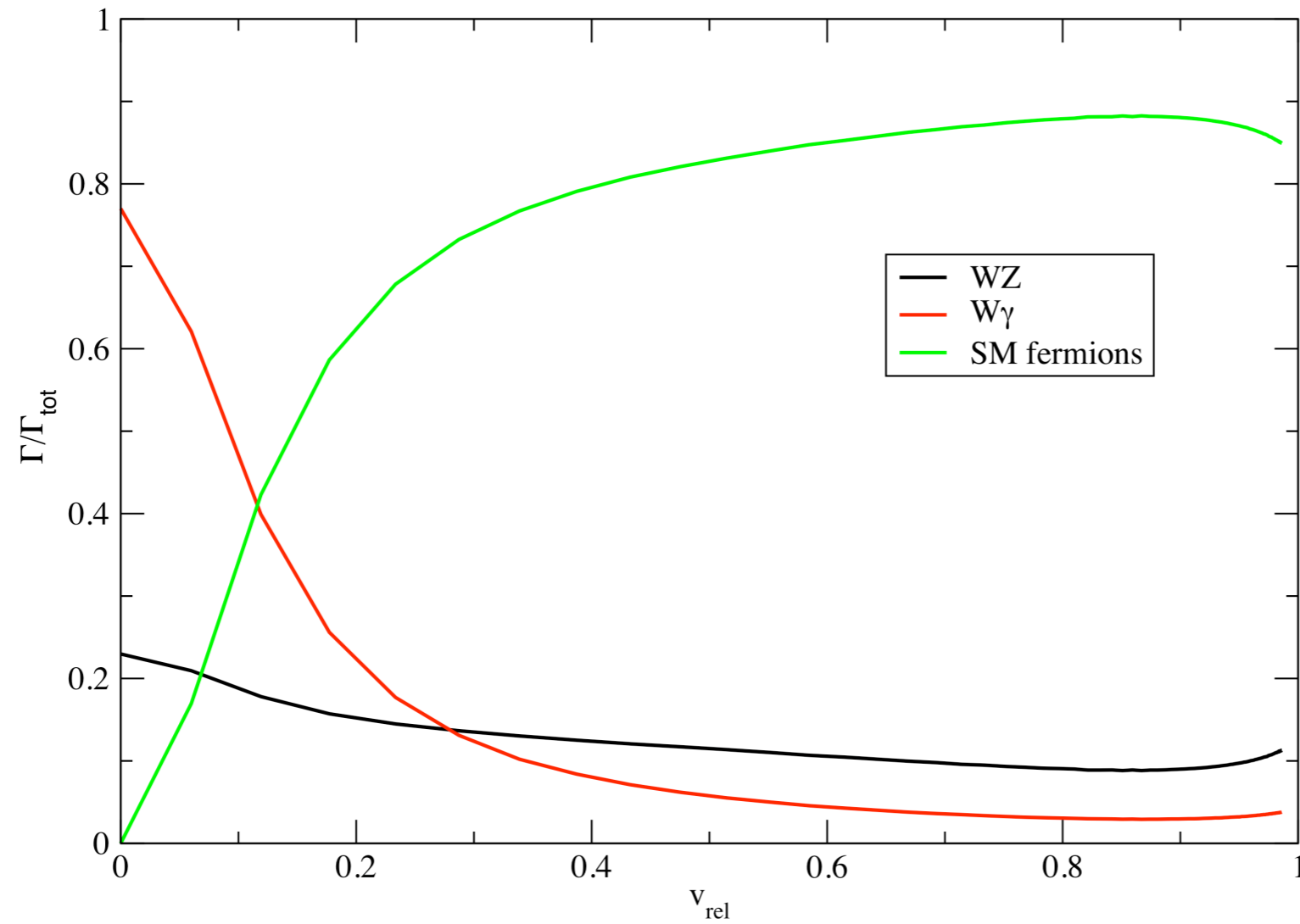
# Squirk Annihilation

w/Z.Chacko, H.Goh, R. Harnik, C. Krenke, 0805.4667



- Annihilation is prompt
- Onium is in s-wave before annihilation

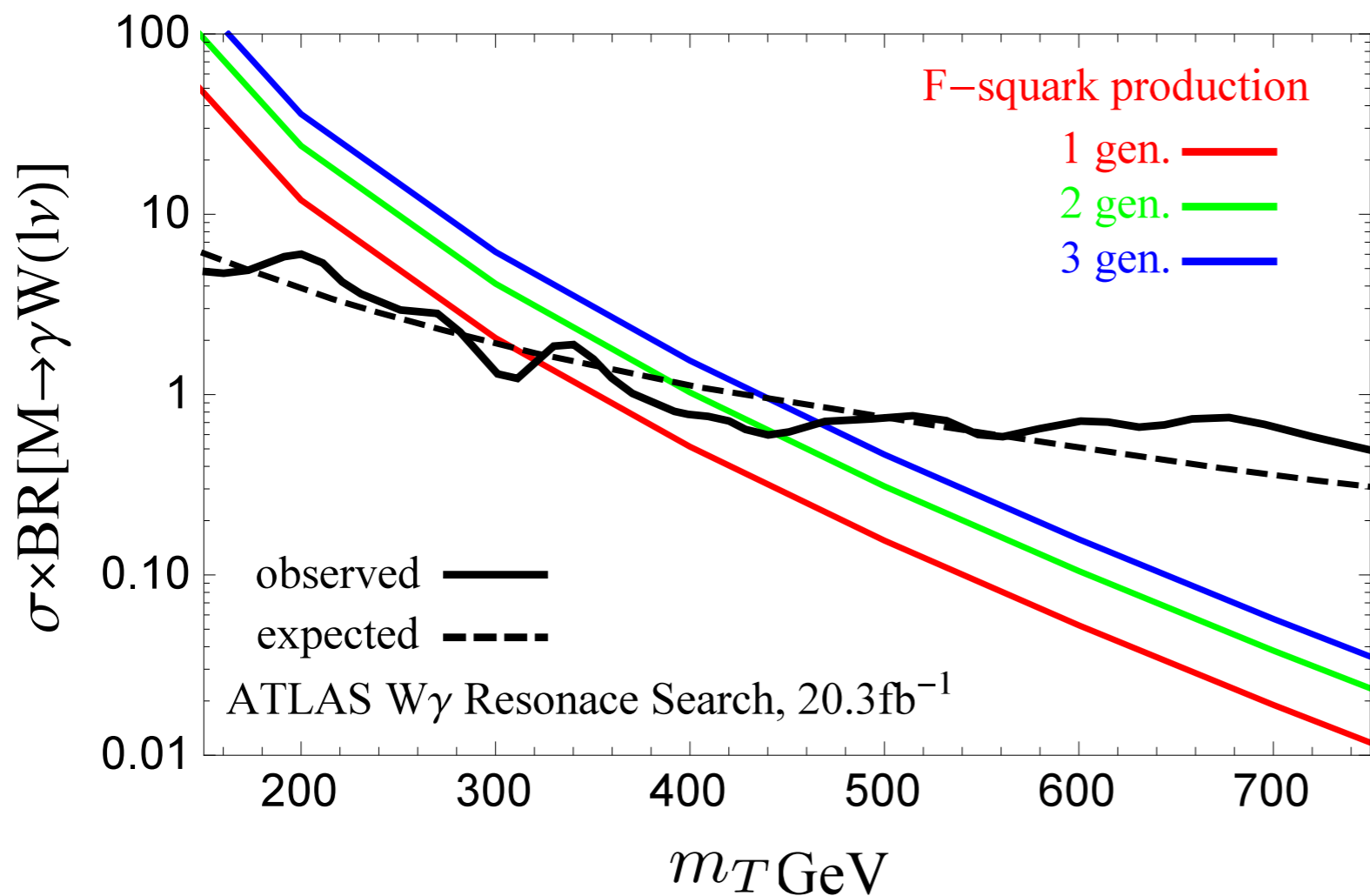
# Squirk Annihilation





# Bounds from the LHC

w/ Z.Chacko et al., 1411.3310



## Assumes

- No smearing
- No  $\beta$  decay

$m_T > \left\{ \begin{array}{l} 320 \\ 445 \\ 465 \end{array} \right\} \text{ GeV} \longrightarrow \text{Direct search better than Higgs couplings}$

# Folded Sleptons (in progress)

In the minimal model, lepton hypermultiplets

$$\begin{array}{cc} \hat{L}_A(1, 1, 2, -1/2) & \hat{L}_B(1, 1, 2, -1/2) \\ \hat{E}_A(1, 1, 1, 1) & \hat{E}_B(1, 1, 1, 1) \end{array}$$

Zero modes:                      **Leptons**                      **F-sleptons**

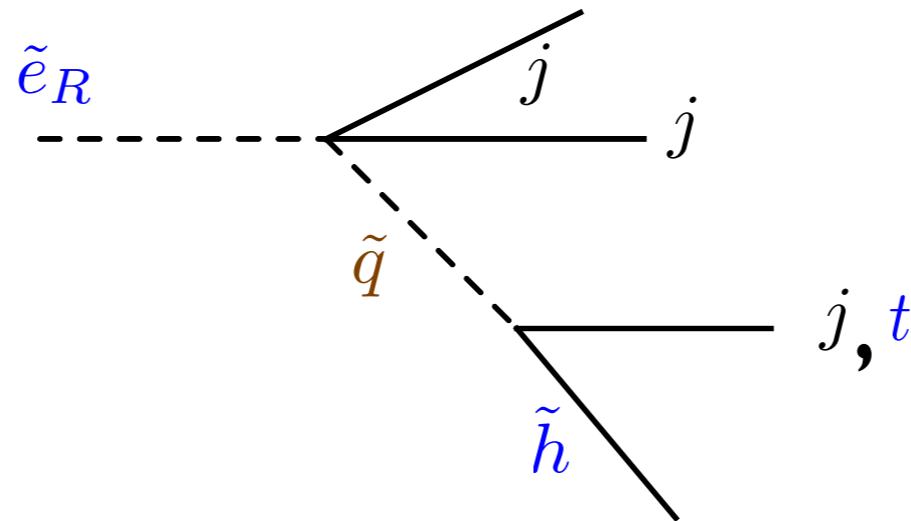
Lightest F-slepton is stable !

Need to add  $Z_2$  preserving HDOs

E.g. 
$$\delta(y) \int d\theta^2 \left( \frac{U_A U_A D_A E_B}{\Lambda} + \frac{U_B U_B D_B E_A}{\Lambda} \right)$$

# Folded Sleptons Decays

Leading to, e.g.:



- Long lifetime

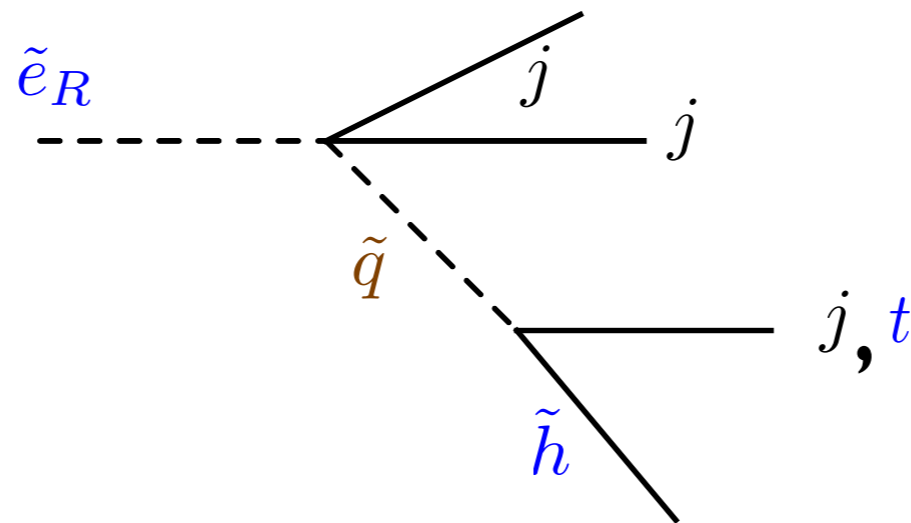
goes through detector below or close to top threshold

- Bounds from CMS (1305.0491):

$$M_{\tilde{e}_R} > 340 \text{ GeV}$$

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Leading to, e.g.:



- Long lifetime

goes through detector below or close to top threshold

- Bounds from CMS (1305.0491):

$$M_{\tilde{e}_R} > 340 \text{ GeV}$$

- Could be an important bound on the naturalness of the model

# Folded SUSY Prospects

- LHC searches for  $W \gamma$
- Model the associated radiation ( $\gamma'$ s and glueball)
- Possible signals for glueball decay to SM fermions with displaced vertices (*see David Curtain's talk*)
- Folded sleptons

# Conclusions/Outlook

- Folded SUSY: colorless stops
- Experimental signals:
  - ★ Quirky dynamics: details of highly excited states
  - ★ Glueball decay
  - ★ Folded sleptons
- Model building issues: UV completions, dark matter, ...
- For the future (100 TeV): explore the physics of the cut-off.