



FROM HADRON COLLISIONS AND RAW DATA TO THE PHYSICS MEASUREMENTS

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**SPLIT SCHOOL OF
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Preface

GOAL OF THE LECTURES:

- **Make you familiar with:**
 - typical steps in the **data analysis chain** within HEP experiments,
 - important aspects of the **statistical data analysis**;
- **Introduce you to important aspects in the extraction of physics information.**

A FEW ASSUMPTIONS:

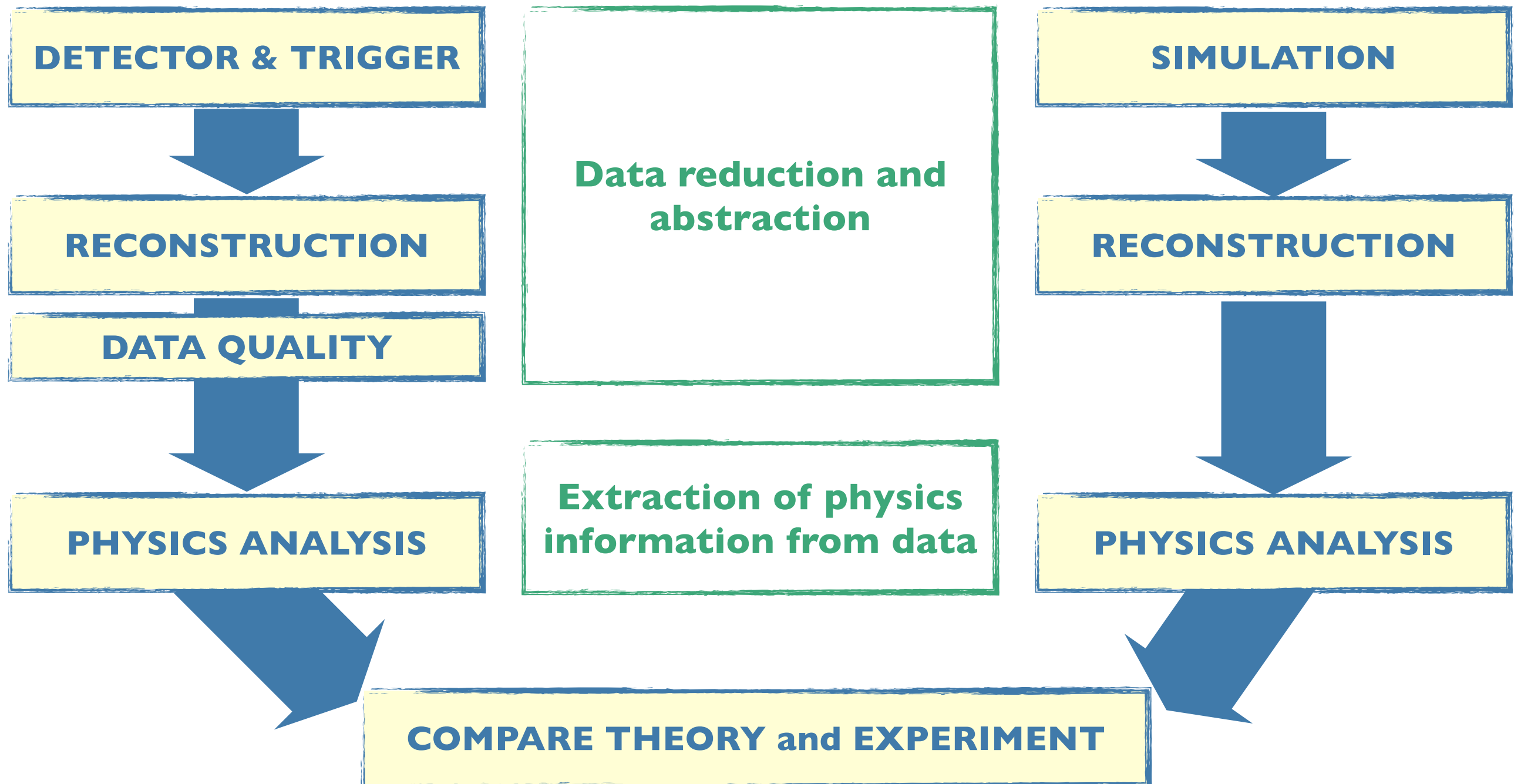
- You are familiar with basic concepts of the **LHC**, and **multi-purpose detectors** in HEP,
- You are familiar with basics of **instrumentation in HEP**,
- You have a limited experience with the physics analysis.

DISCLAIMER:

- **Examples in lectures will have a slight bias towards CMS detector and physics results.**

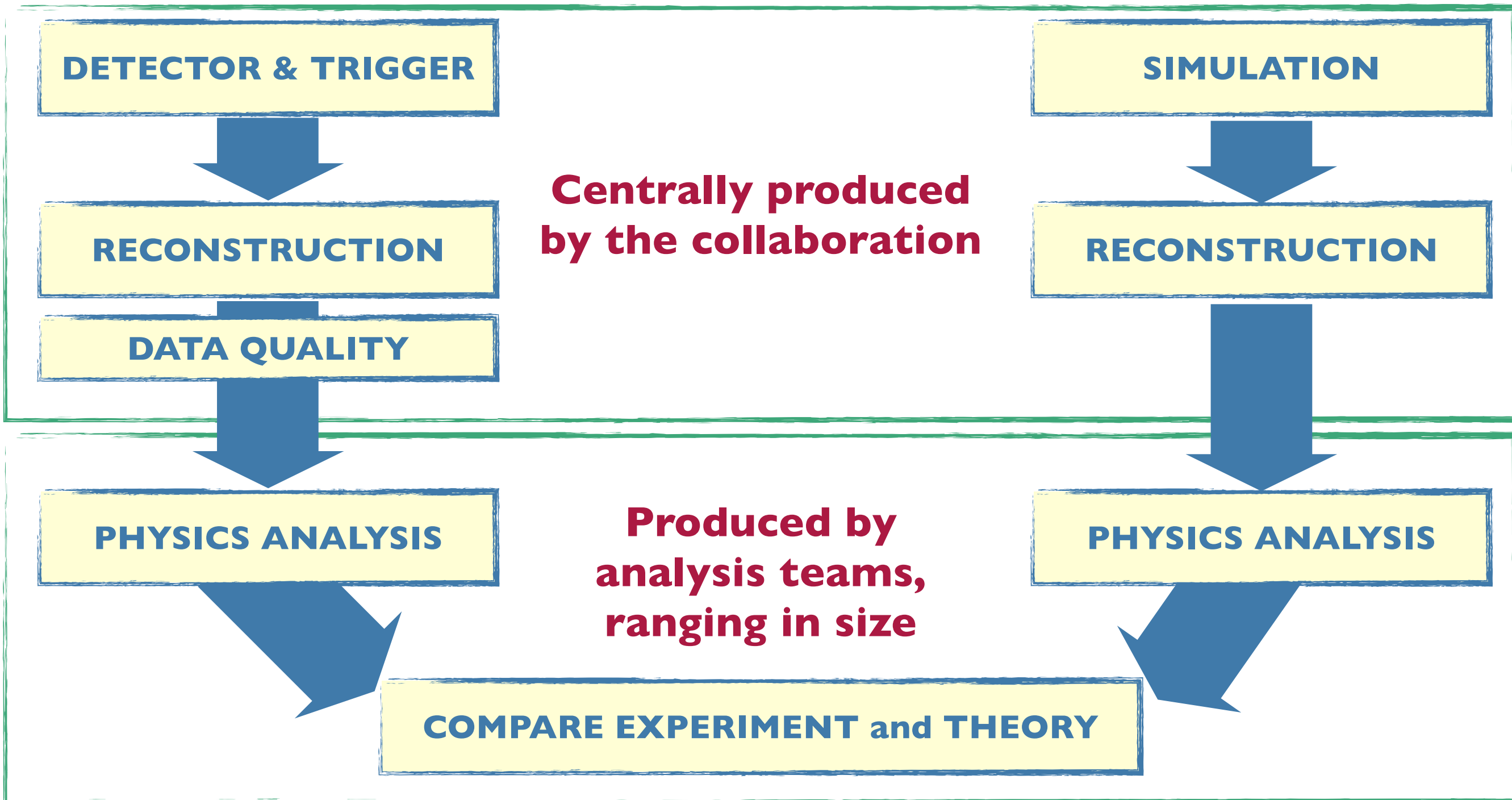
Overview

GENERAL ANALYSIS FLOW



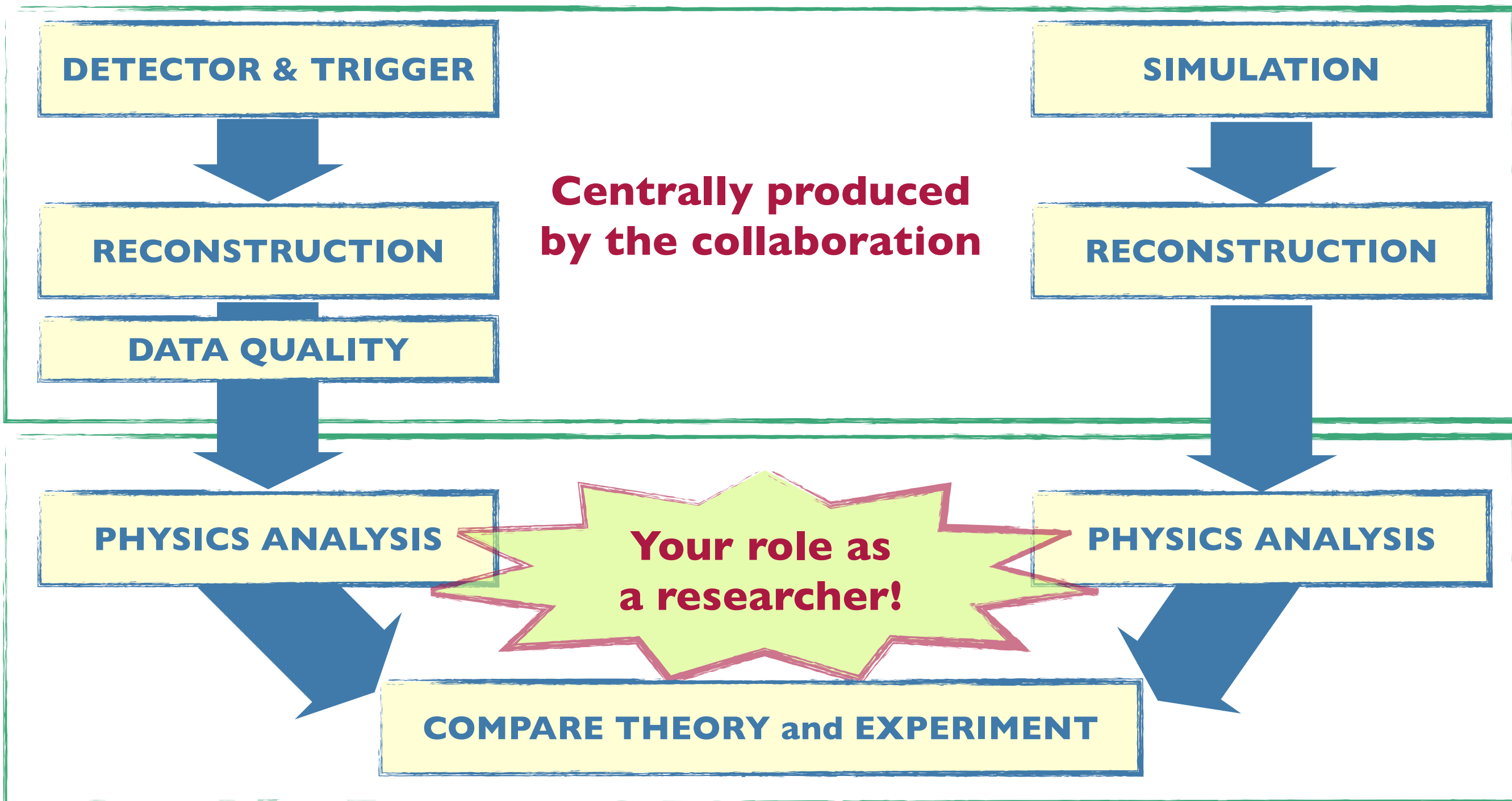
Overview

GENERAL ANALYSIS FLOW



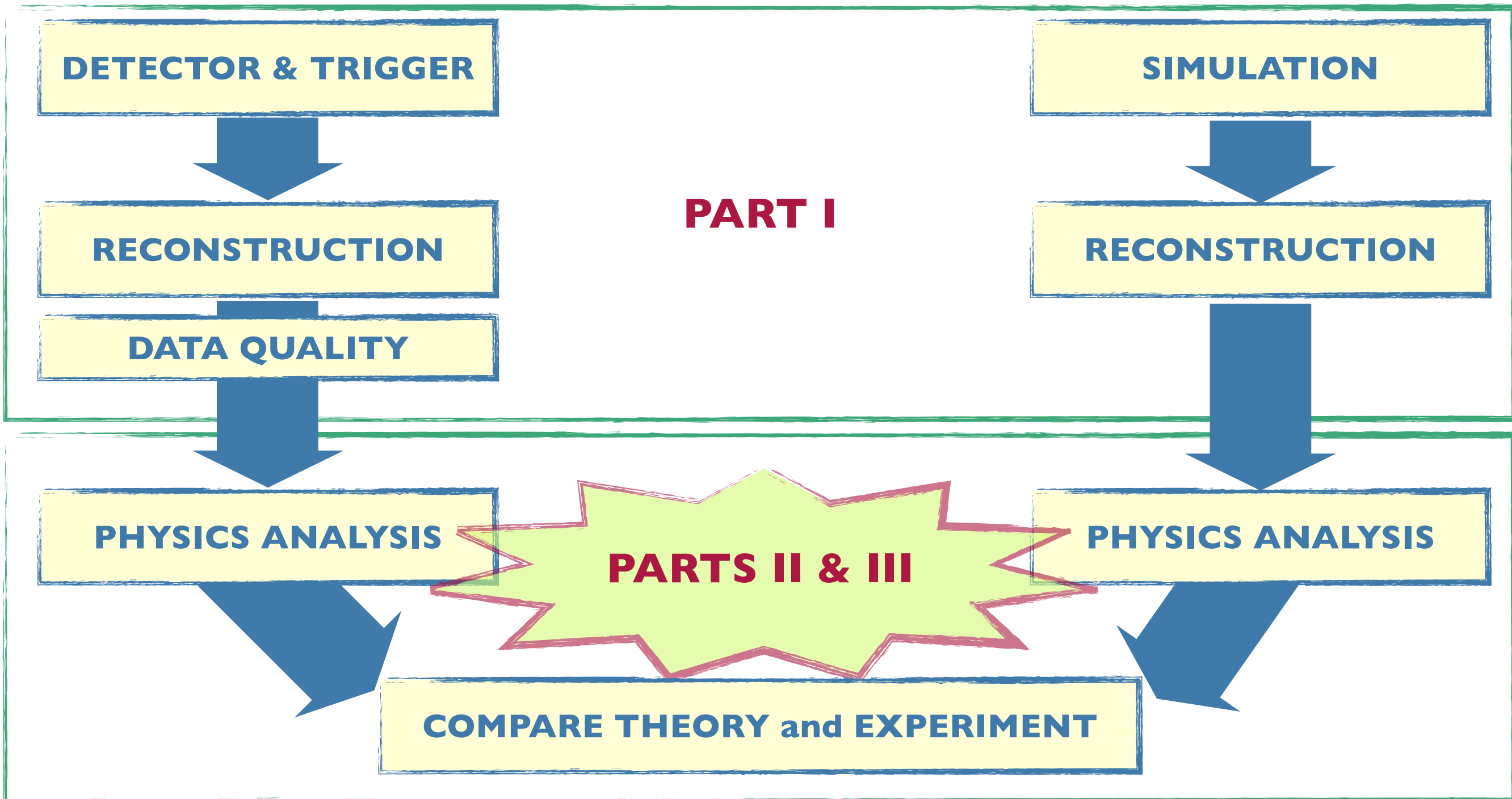
Overview

GENERAL ANALYSIS FLOW



Outline of the lecture

GENERAL ANALYSIS FLOW



Outline of the lecture

PART I: From hadron collisions to the data abstraction

- Key ingredients and steps in the event reconstruction
 - for detailed information on particle identification through interaction with matter see lecture by L.Dobrzynski

PART II: Data analysis and extraction of physics information

- Basic ingredients of the statistical / physics data analysis

PART III: Selected topics on the physics measurements

- Estimation of background processes using data
- Matrix Element Method for separation of physics processes
- Exploitation of interference effects in particle physics
- Fiducial cross section measurements



PART I

From hadron collisions to the data abstraction

Detector & Trigger

DETECTOR:

- Designed to allow for identification of particles through interaction with matter
 - Collect & digitize large amount of information (many channels, many sub-detectors)
 - Size of each event 1.5-2 MB (similar between ATLAS and CMS)



TRIGGER:

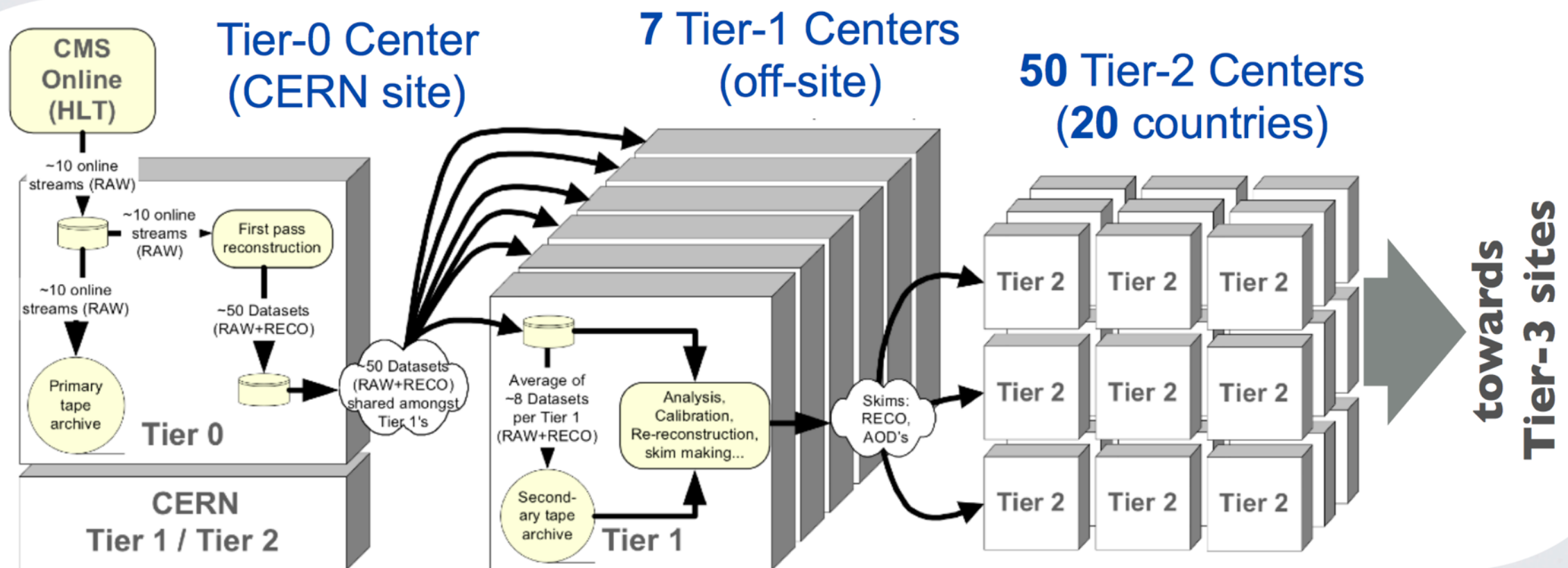
- Decide to readout and process the event, or to throw event away
 - Filtering necessary because of the high rate of collisions, and our interest in processes with widely different production rates (orders of magnitude)
 - Need to perform a fast event processing and selection - need for approximation

Data (re)processing & storage

- Need to allow physicists around the world to access and analyze data
 - **WLHC Grid**: Computing and storage resources distributed around the world
 - Enables: Calibration, re-reconstruction, skimming, simulations, storage, ...

The CMS computing model

Worldwide LHC Computing Grid Infrastructure

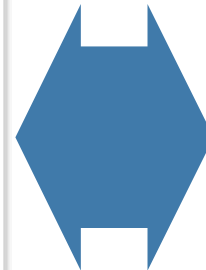


Theory, reality, experiment

Reality → Experiment

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0x01e84c10: 0x01e8 0x8848 0x01e8 0x83d8 0x6c73 0x6f72 0x7400 0x0000
0x01e84c20: 0x0000 0x0019 0x0000 0x0000 0x01e8 0x4d08 0x01e8 0x5b7c
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0x01e84d40: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
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0x01e84d80: 0x0000 0x0019 0x0000 0x0000 0x0000 0x0000 0x01e8 0x5b7c
0x01e84d90: 0x01e8 0x87c0 0x01e8 0x8718 0x7377 0x6974 0x6368 0x0000
    
```



Theory

$$\mathcal{L} = -\frac{1}{4} \mathbf{W}_{\mu\nu} \cdot \mathbf{W}^{\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu}$$

W^\pm, Z, γ kinetic energies and self-interactions

$$+\bar{L}\gamma^\mu (i\partial_\mu - g\frac{1}{2}\boldsymbol{\tau} \cdot \mathbf{W}_\mu - g'Y B_\mu) L$$

$$+\bar{R}\gamma^\mu (i\partial_\mu - g'Y B_\mu) R$$

lepton and quark kinetic energies and their interactions with W^\pm, Z, γ

$$+ \left| (i\partial_\mu - g\frac{1}{2}\boldsymbol{\tau} \cdot \mathbf{W}_\mu - g'Y B_\mu) \phi \right|^2 - V(\phi)$$

W^\pm, Z, γ and Higgs masses and couplings

$$-(G_1\bar{L}\phi R + G_2\bar{L}\phi_c R + h.c.)$$

lepton and quark masses and coupling to Higgs

L ... left-handed fermion (l or q) doublet

R ... right-handed fermion singlet

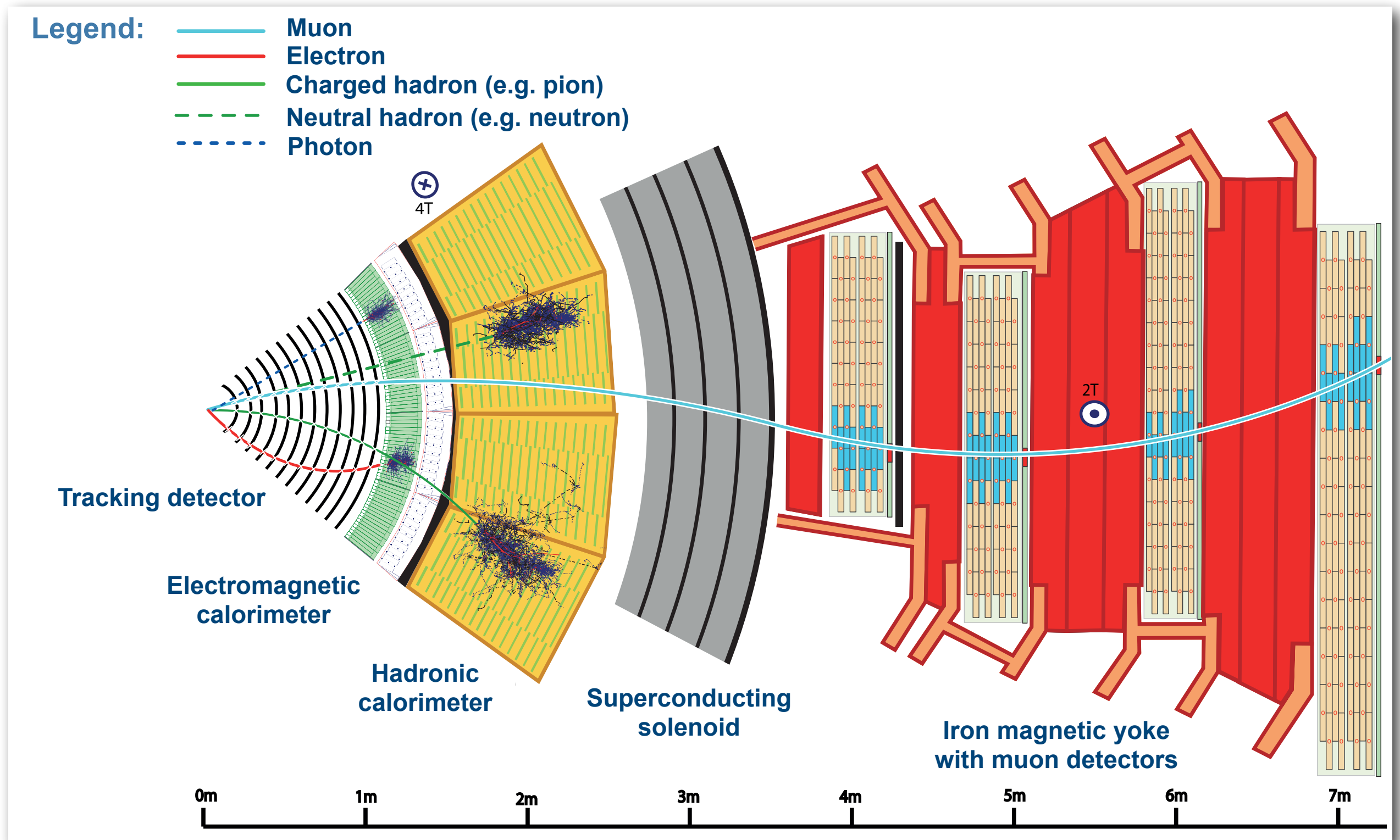
\mathcal{L} from QCD:

$$\mathcal{L} = \underbrace{\bar{q}(i\gamma^\mu\partial_\mu - m)q}_{E_{\text{kin}}(q)} - \underbrace{g(\bar{q}\gamma^\mu T_a q)G_\mu^a}_{\text{Interaction } q, g} - \underbrace{\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}}_{E_{\text{kin}}(g) \text{ includes self-interaction between gluons}}$$



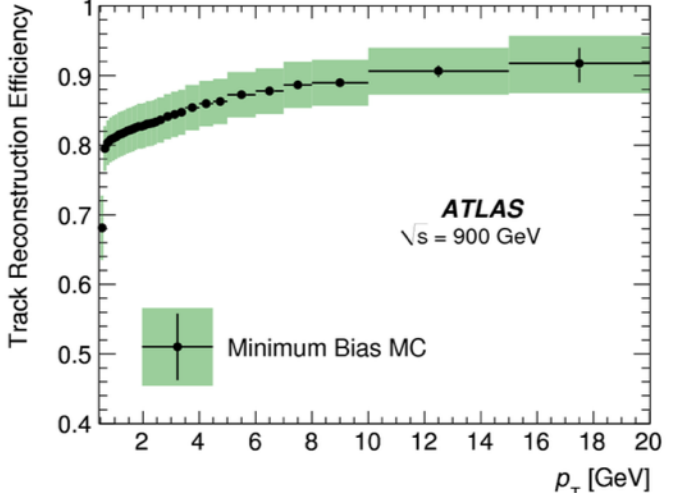
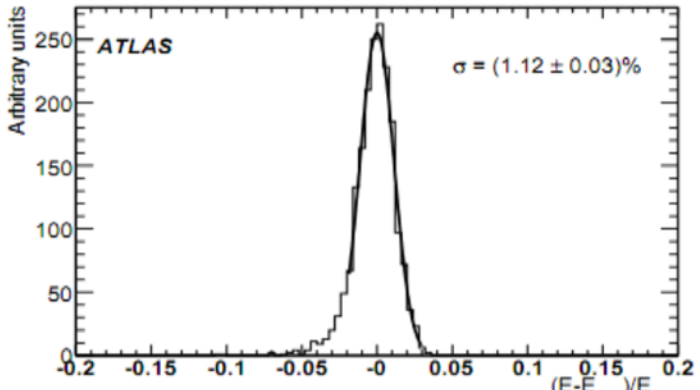
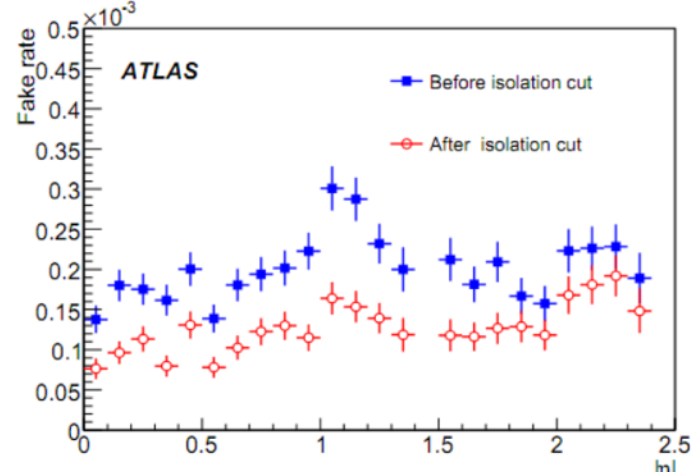
Make sense of these numbers through data abstraction based on physics

Reminder: Layered structure of HEP detector (CMS)



Event reconstruction (I)

IMPORTANT FIGURES OF MERIT:

Efficiency	how often do we reconstruct the object	tracking efficiency = (number of reconstructed tracks) / (number of true tracks)	 <p>ATLAS $\sqrt{s} = 900 \text{ GeV}$ Minimum Bias MC</p> <p>Track Reconstruction Efficiency vs p_T [GeV]. The plot shows a green shaded region representing the efficiency, which starts at approximately 0.8 for $p_T = 2 \text{ GeV}$ and increases to about 0.95 at $p_T = 20 \text{ GeV}$. A black line with error bars represents the Minimum Bias MC simulation.</p>	High
Resolution	how accurately do we reconstruct the quantity	energy resolution = (measured energy – true energy) / (true energy)	 <p>ATLAS $\sigma = (1.12 \pm 0.03)\%$</p> <p>Arbitrary units vs $(E - E_{\text{true}}) / E_{\text{true}}$. The plot shows a sharp peak centered at 0, indicating good energy resolution. The x-axis ranges from -0.2 to 0.2, and the y-axis ranges from 0 to 250.</p>	Good
Fake rate	how often we reconstruct a different object as the object we are interested in	a jet faking an electron, fake rate = (Number of jets reconstructed as an electron) / (Number of jets)	 <p>ATLAS</p> <p>Fake rate vs η. The plot shows two data series: 'Before isolation cut' (blue squares) and 'After isolation cut' (red circles). The y-axis is Fake rate $\times 10^{-3}$ (0 to 0.5), and the x-axis is η (0 to 2.5). The fake rate is significantly lower after the isolation cut.</p>	Low

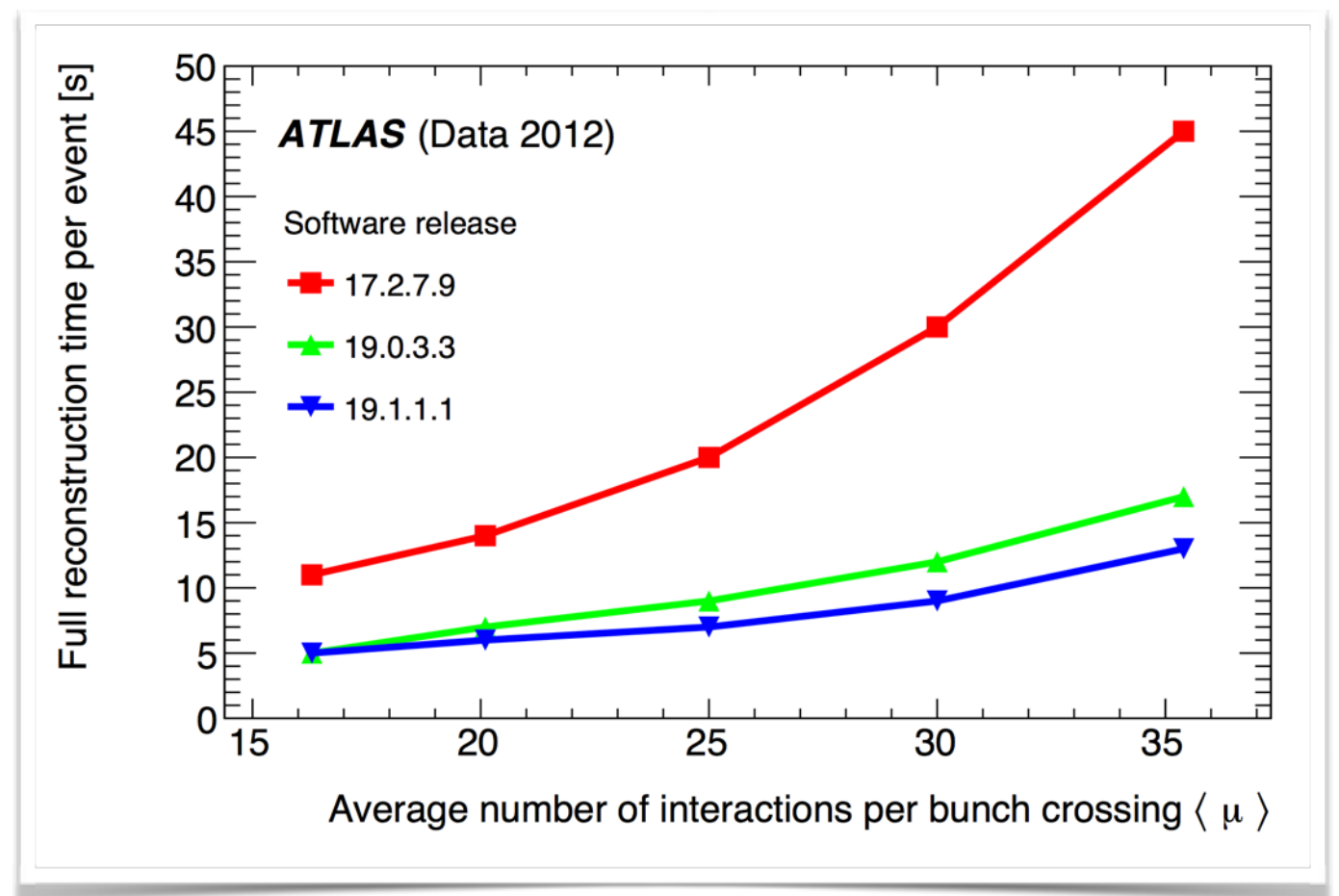
Event reconstruction (2)

GOAL:

- Data abstraction based on the particle physics principles

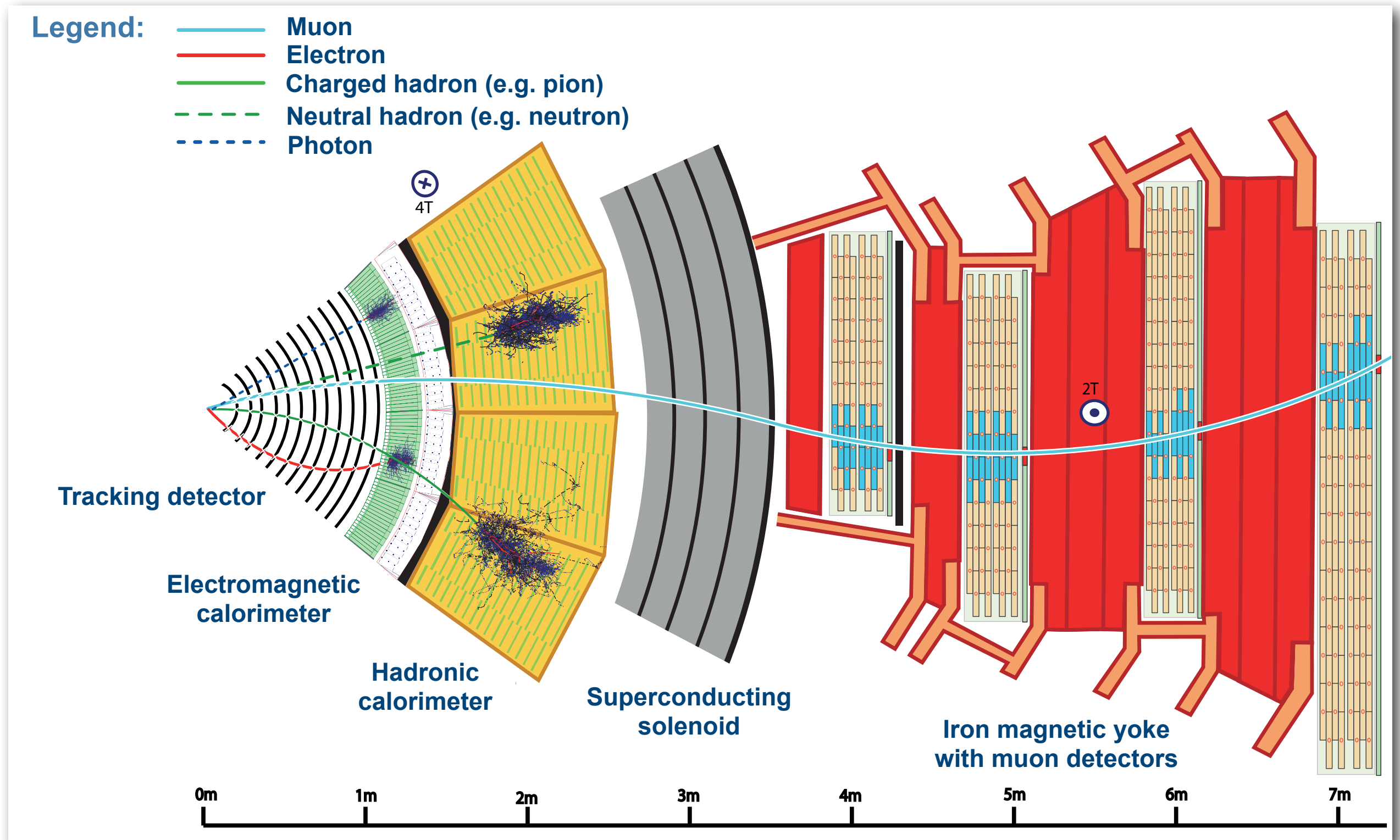
DESIRED PERFORMANCE:

- High efficiency + Good resolution + Low mis-reconstruction rate.
- Robust against detector problems and data-taking conditions:
 - Noise,
 - Dead regions of the detector,
 - Increased pile-up.
- Computing-friendly:
 - CPU time per event,
 - Memory use.



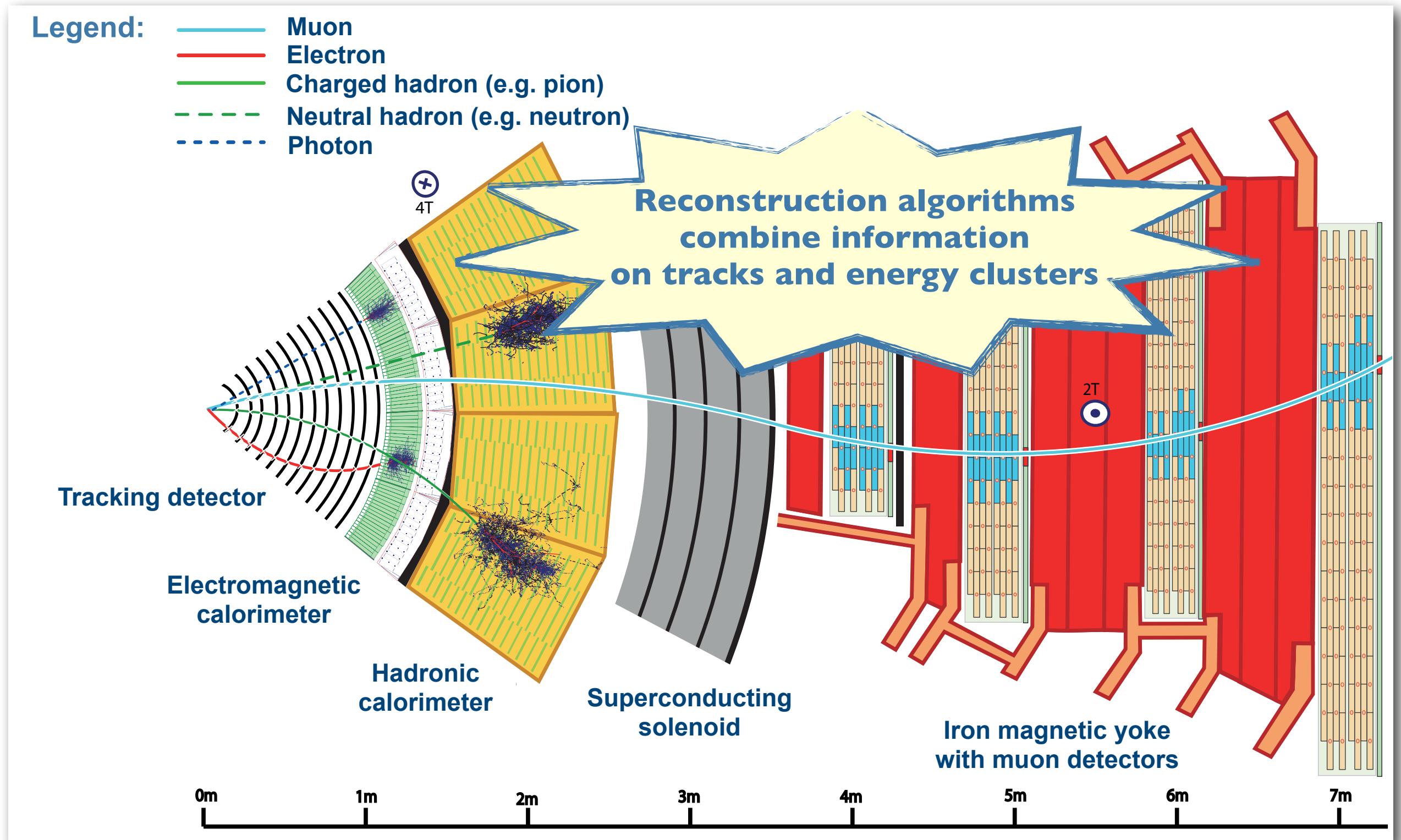
Individual particles

- Direct reconstruction of certain types of individual particles

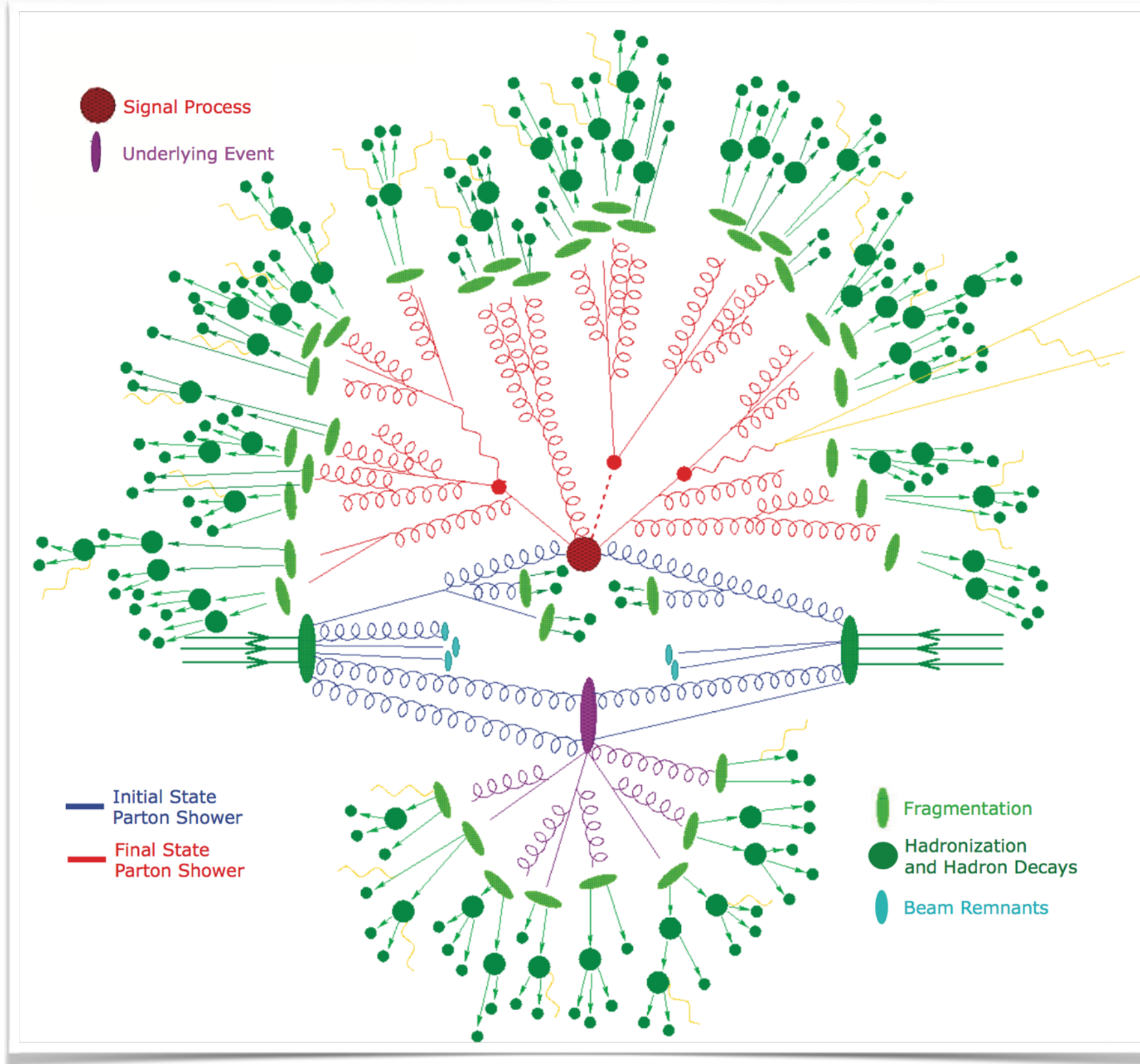


Individual particles

- Direct reconstruction of certain types of individual particles



Complex picture of proton-proton collisions



Jets

PRODUCTION:

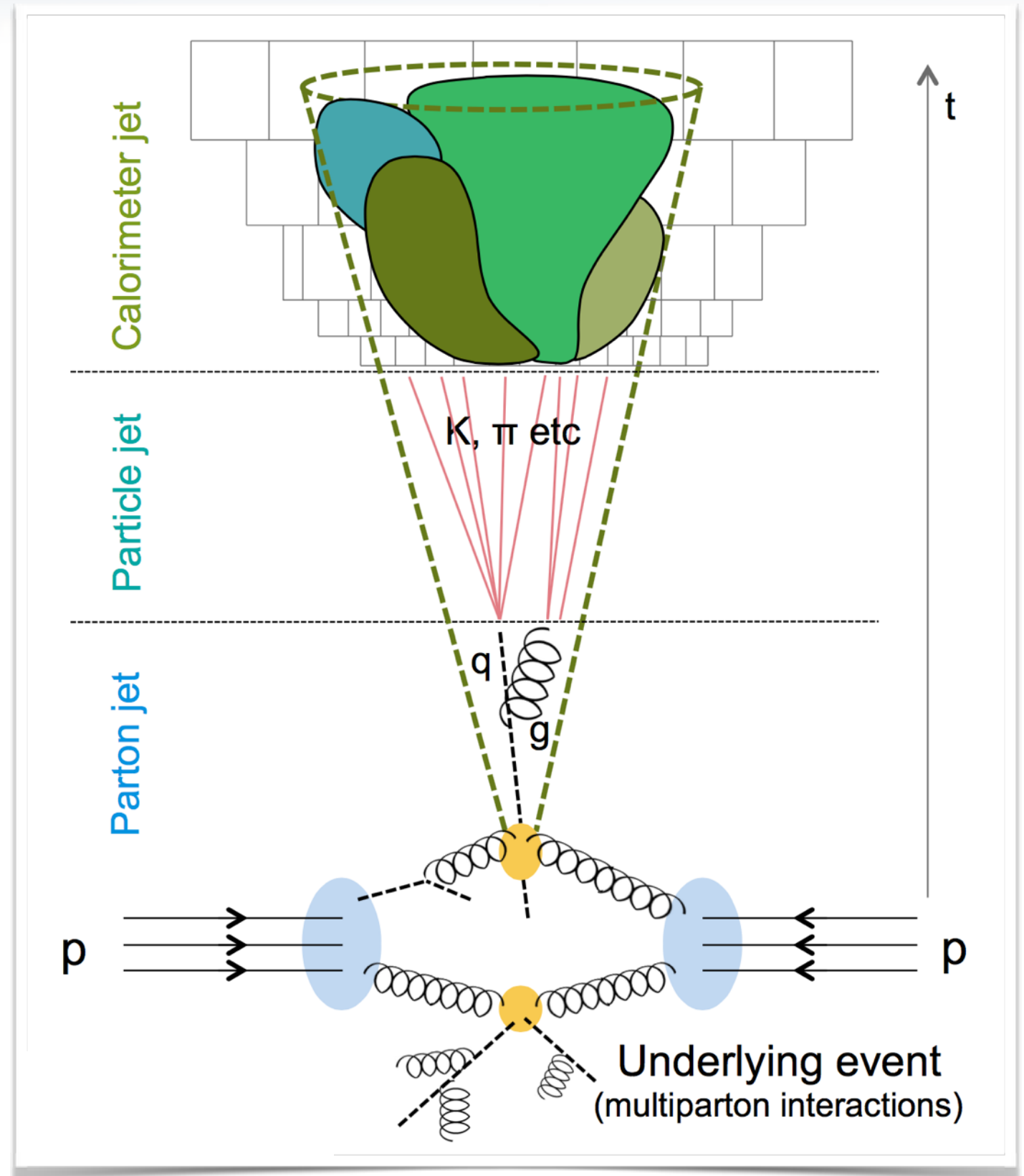
- by fragmentation of gluons and (light) quarks in QCD scattering

RECONSTRUCTION:

- Need to satisfy requirements:
 - theoretical requirements (infrared and collinear safety)
 - experimental requirements (detector & environment independent, easily implementable, etc.)
- Commonly used in ATLAS and CMS
 - 'anti-kt' algorithm (typical cone sizes: $R=0.4/0.5$)

CALIBRATION:

- Correct the energy and position measurement, and the resolution.
- Correct for instrumental & physics effects



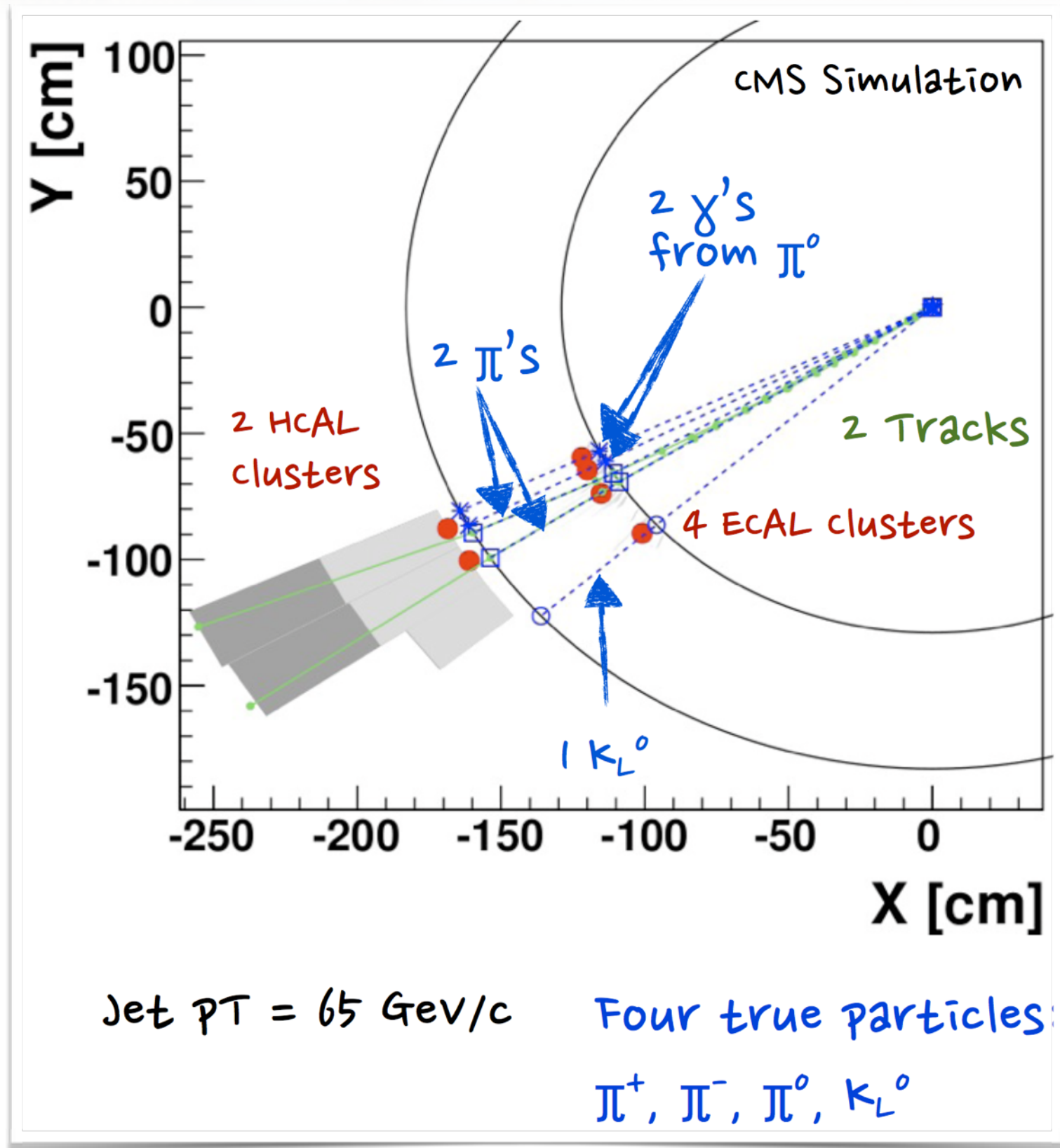
Particle flow approach

BASICS:

- “Flow of particles” through the detector.
- Reconstruct and identify all individual particles
- Combination of sub-detector info for measuring E , η , φ , ID.

IMPORTANT ASPECTS:

- High precision, and high efficiency tracking;
- Large magnetic field for good p_T resolution and charged-to-neutral particle separation;
- Highly granular calorimeter



PART I: From hadron collisions to the data abstraction





PART II

Data analysis and extraction of physics information

Physics/data analysis (I)

MAIN GOAL:

- **Learn about the nature** by extracting physics information from data

IMPORTANT ASPECTS:

- **Use particle physics principles** to explore interesting phenomena in data
- **Use statistics** for presentation and interpretation (explanation) of data
 - **Descriptive statistics:**
Describes the main features of a collection of data in quantitative terms
 - **Inductive statistics:** Makes inference about a random process from its observed behavior during a finite period of time

Physics/data analysis (2)

CONFIRMATORY AND EXPLORATORY DATA ANALYSIS:

- **Confirmatory analysis** - statistical hypothesis testing:
A method of making statistical decisions using experimental data, (**Frequentist** hypothesis testing & **Bayesian** inference).
- **Exploratory analysis** - uses data to suggest hypothesis to test:
Complements confirmatory data analysis.

QUANTITATIVE AND GRAPHICAL TECHNIQUES:

- **Quantitative techniques** - yield numeric / tabular output:
(point estimation, interval estimation, hypothesis testing, etc.),
- **Graphical techniques** - gain insight in data set, finding structures in data, checking assumptions on statistical models, communicate results in convincing way, (Includes: graphs, histograms, scatter plots, probability plots, residual plots, etc.).

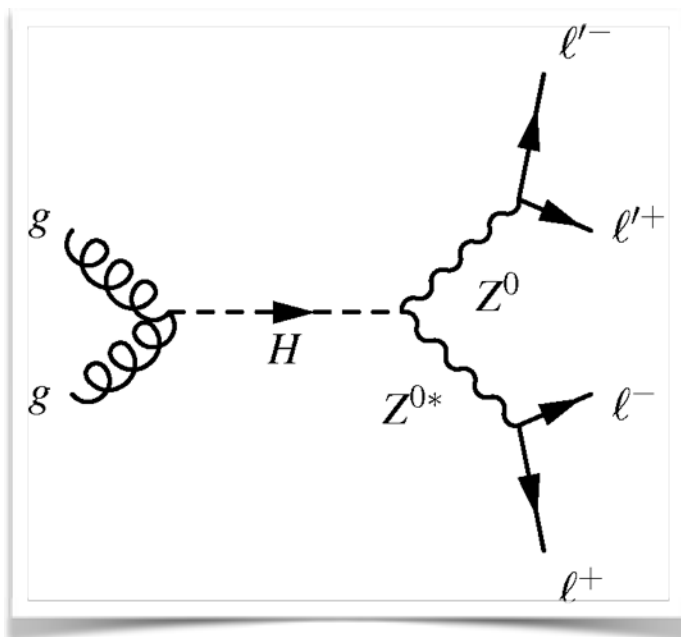
Signal vs. background processes

SIGNAL AND BACKGROUND:

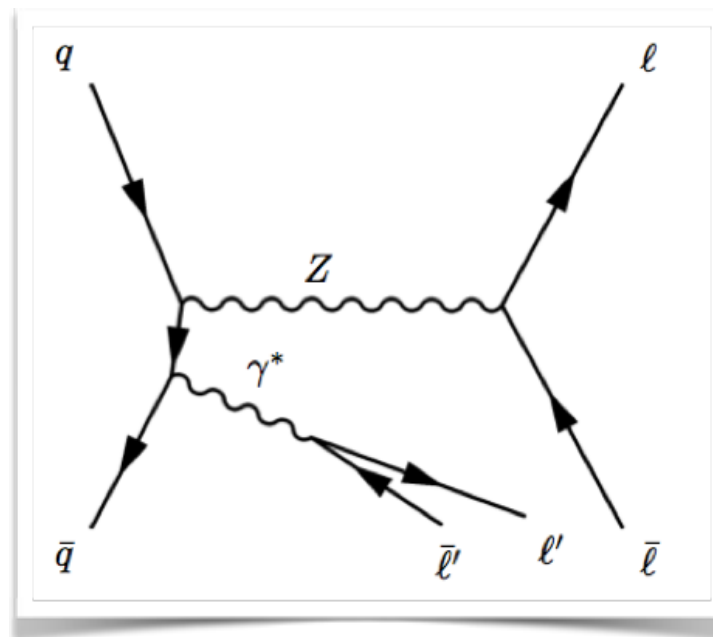
- **Signal:** an event coming from the physical process under study
- **Background:** any other event
 - “Dangerous”: events that after reconstruction have final state topology as signal
 - Irreducible** - intrinsically the same final state as the signal
 - Reducible / instrumental** - events with mis-reconstructed objects in final state (e.g. parts of jets reconstructed as electrons)

EXAMPLE PROCESS: $pp \rightarrow H \rightarrow ZZ \rightarrow 4\ell$

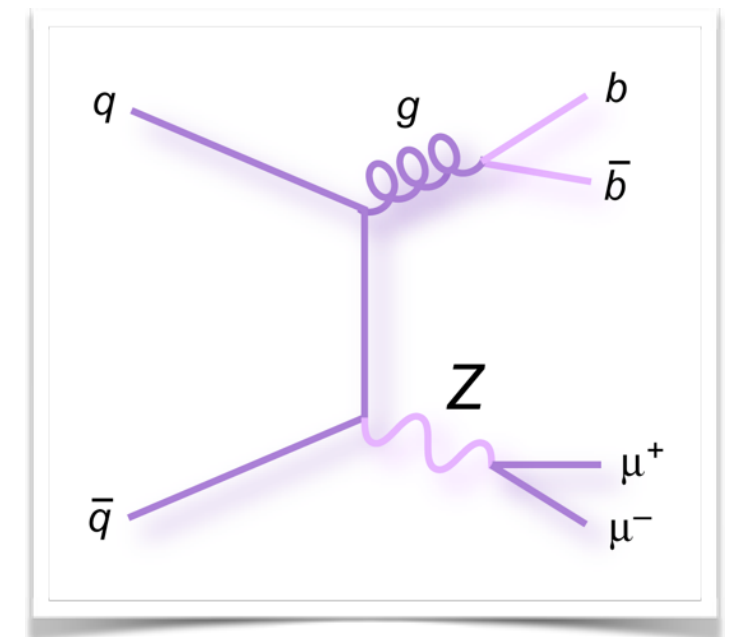
Signal



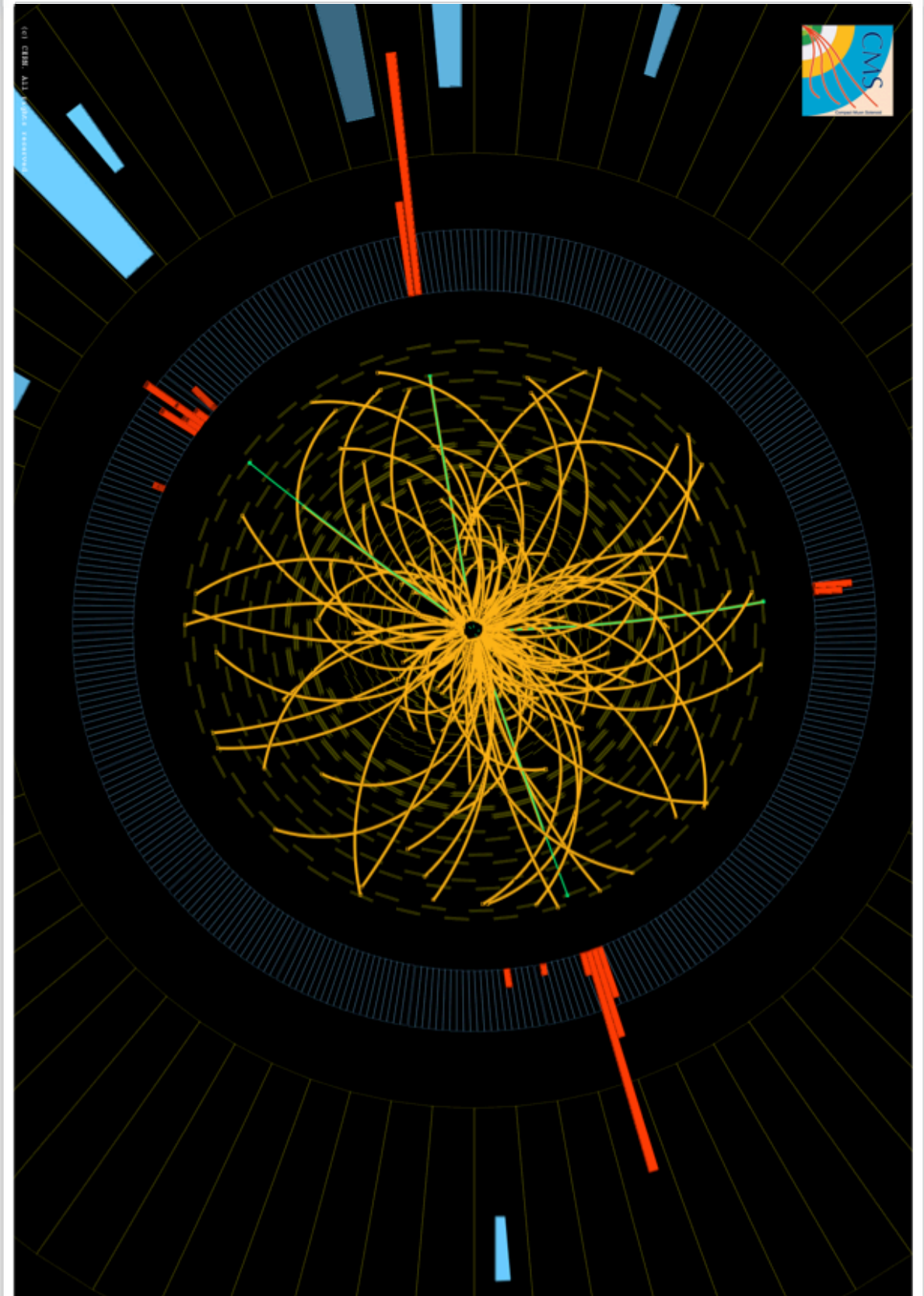
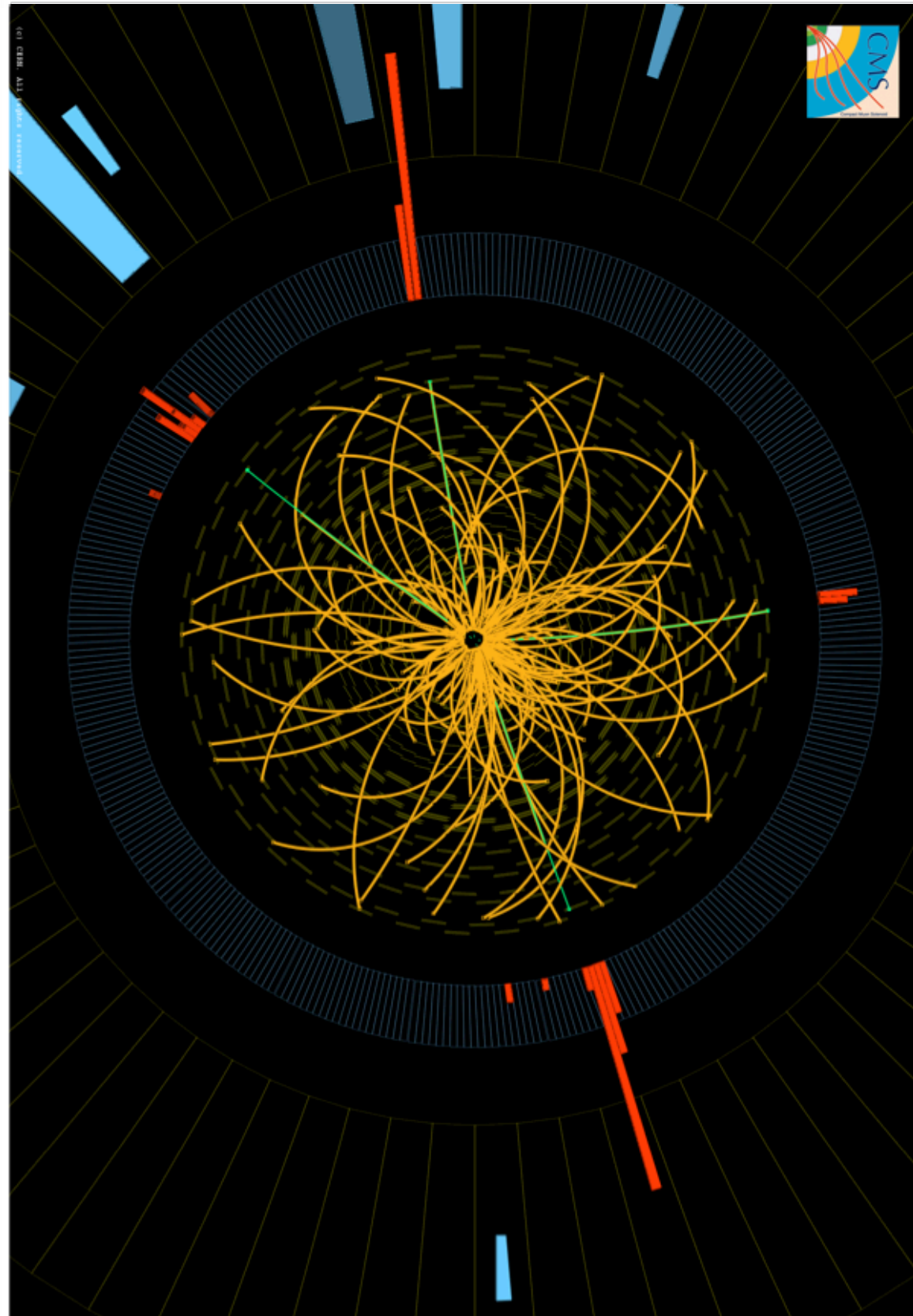
Irreducible background



Instrumental background



Signal? or background?



Signal vs. background processes (2)

ULTIMATE GOAL:

- Use physics principles to separate signal from background (as much as possible), **$H \rightarrow ZZ \rightarrow 4\ell$** example: require at least 4 reconstructed leptons with high p_T , etc.
- Typically obtained in a chain of several steps



IMPORTANT ASPECTS:

- **Nature & observations are non-deterministic:**
For a given event it is not possible to tell whether it's signal or background!
- Assign probabilities that the observed event comes from signal or background
 $p(\text{event}|\text{signal})$ and $p(\text{event}|\text{background})$

Important examples in PART III:

- Estimation of reducible background from data,
- Separation of irreducible background from signal.

Physicists & statisticians : illustrative comparison

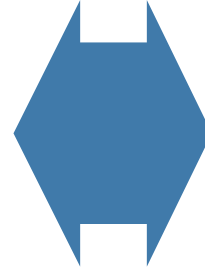
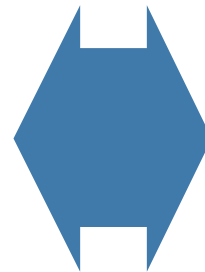
EXAMPLE: HISTOGRAM FITTING:

Physicists

1. Determining the “best fit” parameters of a curve

2. Determining the errors on the parameters

3. Judging the goodness of a fit



Statisticians

1. Point estimation

2. Confidence interval estimation

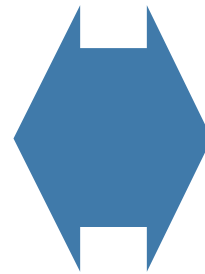
3. Goodness-of-fit (hypothesis) testing

Physicists & statisticians : illustrative comparison

EXAMPLE: HISTOGRAM FITTING:

Physicists

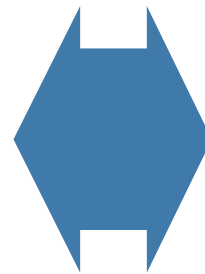
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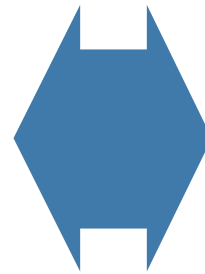
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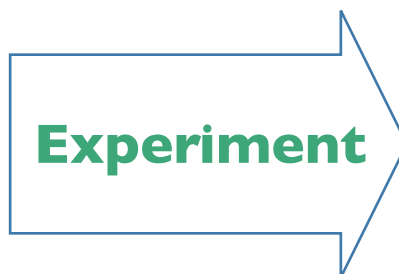
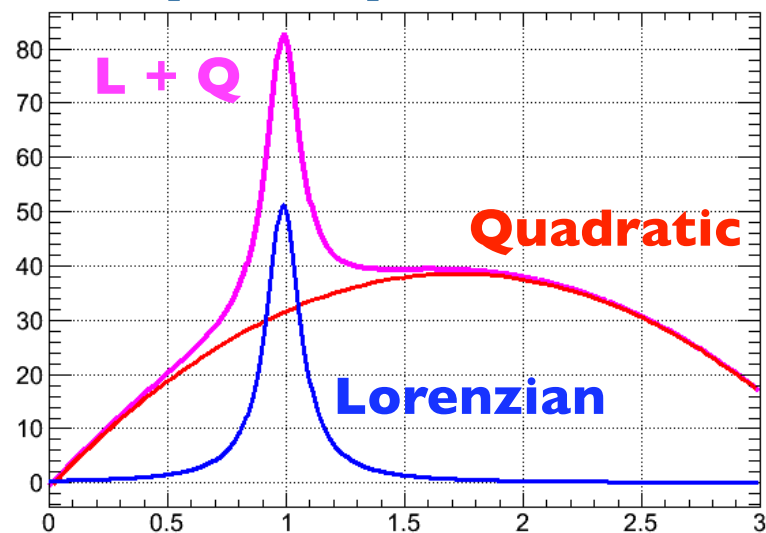


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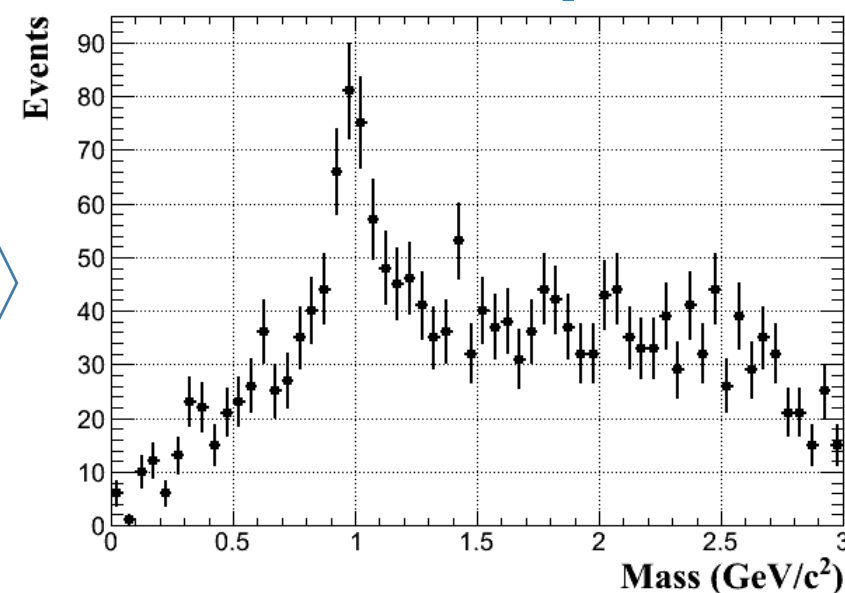
Parameter/point estimation

- **Physical phenomena:** described by a function that depends on p **unknown** parameters with **true values:** $\theta^{true} = (\theta_1^{true}, \theta_2^{true}, \dots, \theta_p^{true})$
- **Goal:** Estimate the true values of the parameters: θ_i^{true} (estimators: $\hat{\theta}_i$)
- **EXAMPLE:** Sampling the reality & estimating parameters

Physical phenomena



Data sample



$$\text{Lorentzian} = L(x; D, \Gamma, M) \sim \frac{D\Gamma}{(x^2 - M^2)^2 + 0.25\Gamma^2}$$

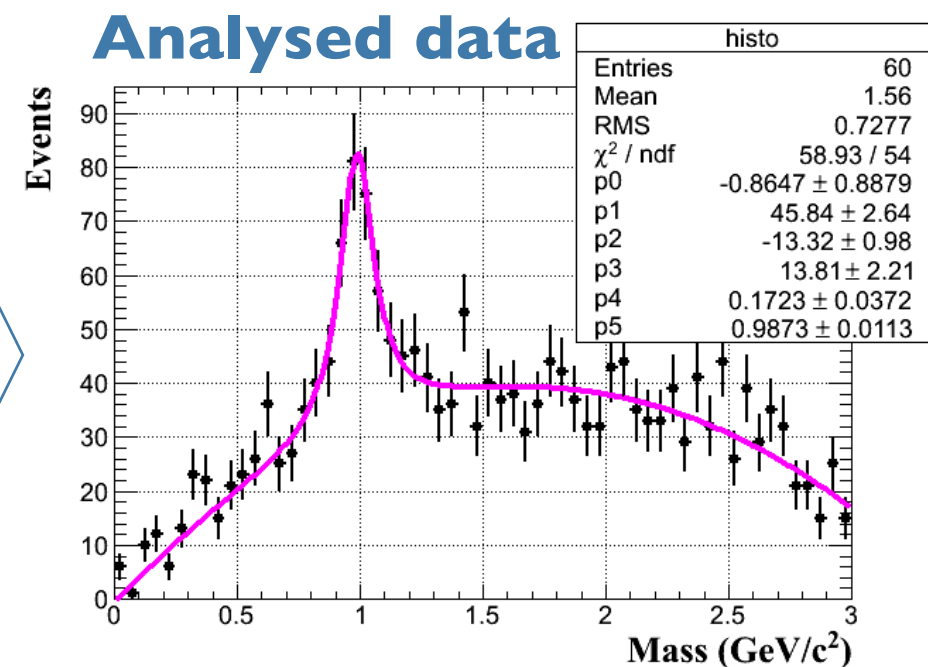
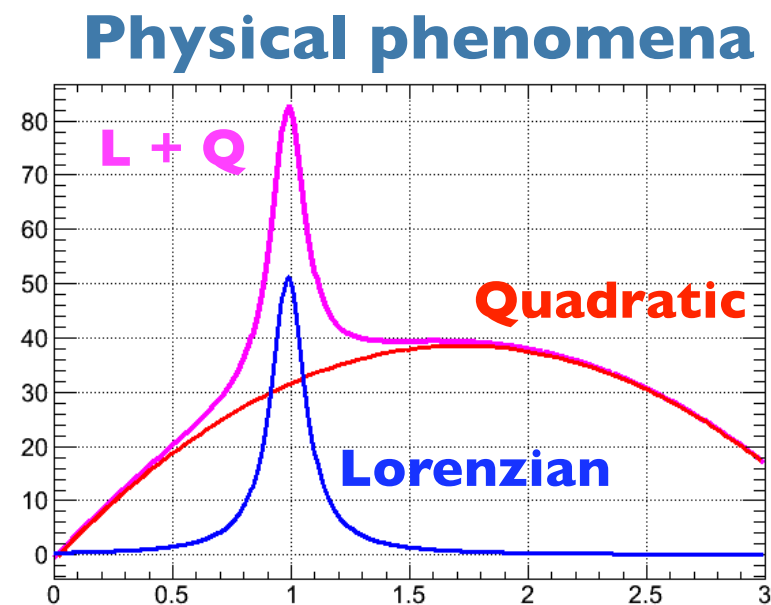
$$\text{Quadratic} = Q(x; A, B, C) \sim A + Bx + Cx^2$$

- Find the **function** that describes the measurements the best

$$F(x; D, \Gamma, M, A, B, C) = L(x; D, \Gamma, M) + Q(x; A, B, C) = F(x; \theta)$$

Parameter/point estimation

- **Physical phenomena:** described by a function that depends on p **unknown** parameters with **true values:** $\theta^{true} = (\theta_1^{true}, \theta_2^{true}, \dots, \theta_p^{true})$
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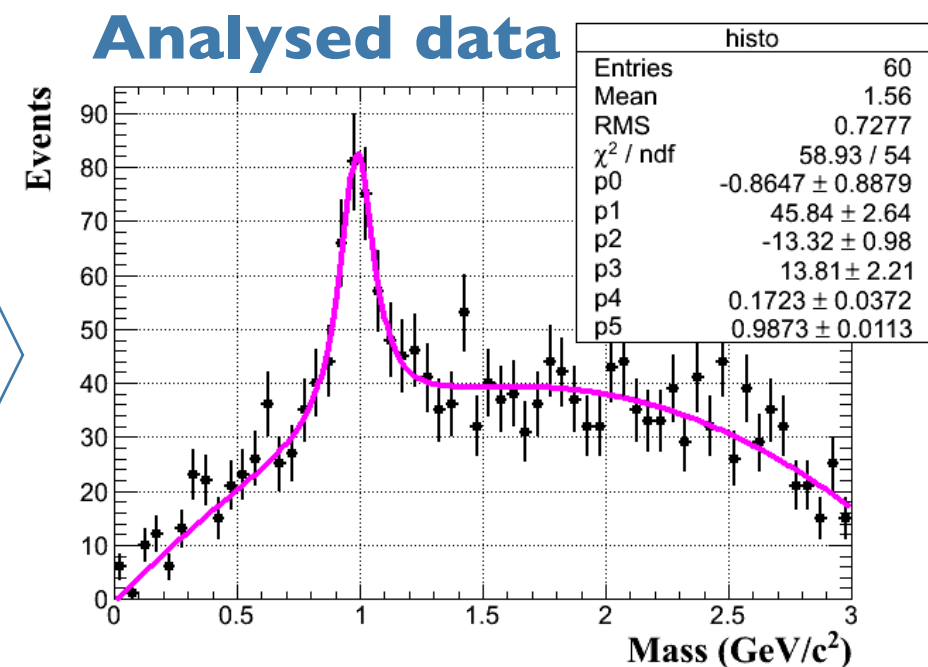
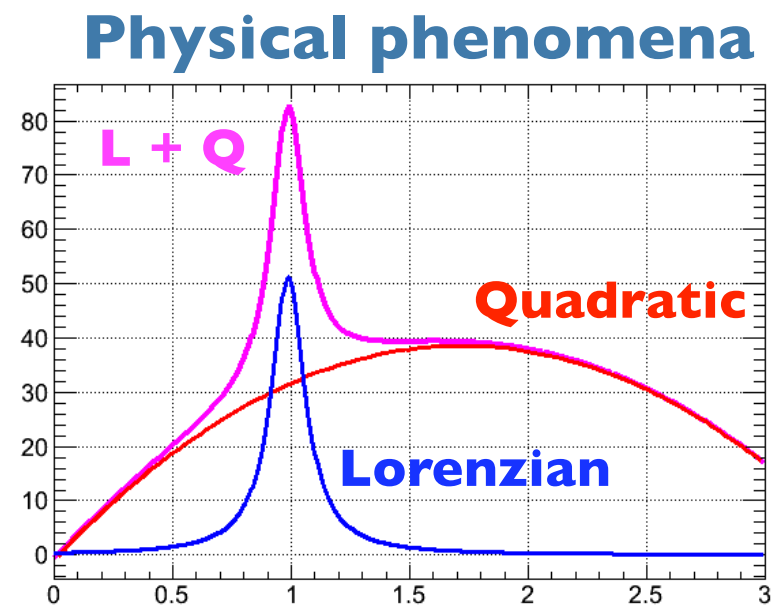


- Find the **function** that describes the measurements the best
- The parameters of that **function** are **estimators of the unknown parameters**

$$F(x; \hat{D}, \hat{\Gamma}, \hat{M}, \hat{A}, \hat{B}, \hat{C}) = L(x; \hat{D}, \hat{\Gamma}, \hat{M}) + Q(x; \hat{A}, \hat{B}, \hat{C}) = F(x; \hat{\theta})$$

Parameter/point estimation

- **Physical phenomena:** described by a function that depends on p **unknown** parameters with **true values:** $\theta^{true} = (\theta_1^{true}, \theta_2^{true}, \dots, \theta_p^{true})$
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- Find the **function** that describes the measurements the best
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$$F(x; \hat{D}, \hat{\Gamma}, \hat{M}, \hat{A}, \hat{B}, \hat{C}) = L(x; \hat{D}, \hat{\Gamma}, \hat{M}) + Q(x; \hat{A}, \hat{B}, \hat{C}) = F(x; \hat{\theta})$$

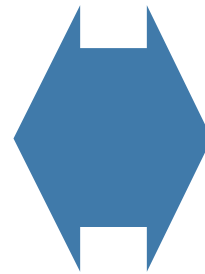
- **Methods:** **Maximum Likelihood, Least Squares, Method of Moments, etc.**

Physicists & statisticians : illustrative comparison

EXAMPLE: HISTOGRAM FITTING:

Physicists

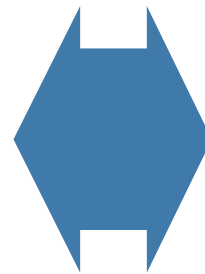
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Statisticians

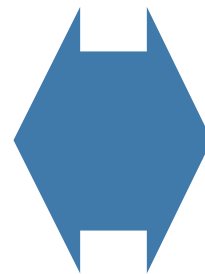
1. Point estimation

2. Determining the errors on the parameters



2. Confidence interval estimation

3. Judging the goodness of a fit



3. Goodness-of-fit (hypothesis) testing

Parameter error (confidence interval) estimation

Confidence intervals⁷

For a Gaussian estimator the result of an experiment is usually expressed by

- The parameter's estimated value, plus/minus an estimate of the **standard deviation**, $\hat{\theta} \pm \sigma_{\hat{\theta}}$

If the pdf is not Gaussian, or in the presence of physical boundaries

- One usually quotes instead an **interval**.

The quoted interval or limit should:

- Objectively communicate the result of the experiment,
- **Communicate incorporated prior beliefs** and relevant assumptions,
- Provide interval that covers the true value of the θ with specified probability,
- Make possible to draw conclusions about the parameter.

These goals are satisfied in **case of large data sample** by $\hat{\theta} \pm \sigma_{\hat{\theta}}$, and in the multi-parameter case by

- The parameter estimates and covariance matrix.

For **small data sample**, or in case of constrained variables, the Bayesian or the Neyman approach can be used.

⁷Adapted from [Particle Data Group](#).

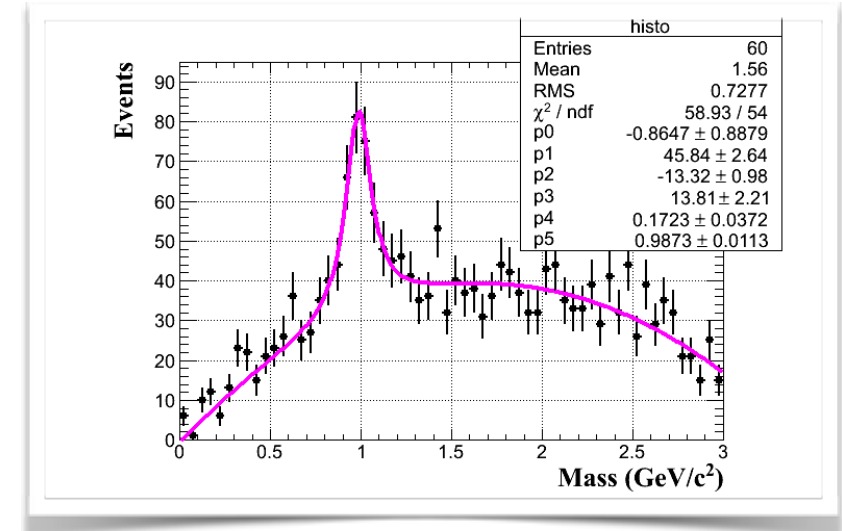
Errors on the Maximum Likelihood estimates

ML parameter errors:

Can be **extracted from the shape** of the likelihood function around its maximum

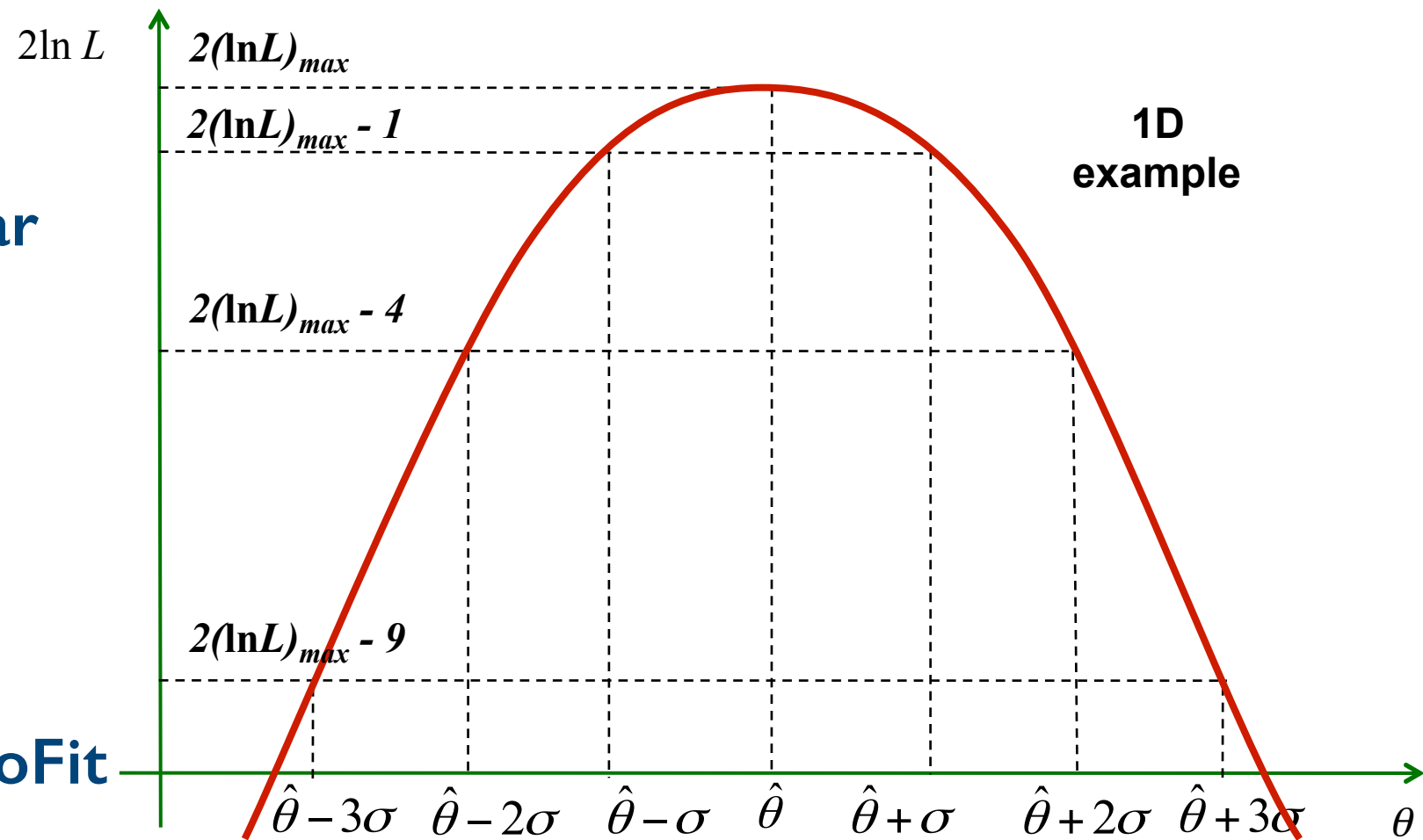
- In the large N limits 1D likelihood function is Gaussian, and the $\ln\mathcal{L}$ is paraboloid (χ^2 with 1 d.o.f.)
- 1D Interval $\pm 1\sigma$ are extracted as values for which $2\ln\mathcal{L}$ falls by 1 from its maximum value $(\ln\mathcal{L})_{max}$

Example



Important aspects:

- For finite samples or non-linear problems $\ln\mathcal{L}$ is asymmetric
- In case of P parameters, $\ln\mathcal{L}$ asymptotically behaves as χ^2 with P d.o.f.
- Can be computed using **MINUIT/MINOS** in **ROOT/RooFit**



Uncertainty in physics measurements

The sources of uncertainty in measurement⁹:

- **Incomplete definition** of the measurand; or its imperfect realization
- **Non-representative sampling**
- inadequate knowledge of the effects of environmental conditions; or imperfect measurements of these conditions
- **Personal bias** in reading instruments
- **Finite instrument resolution**
- Inexact values of measurement standards and reference materials
- **Inexact values of constants** and other parameters obtained from external sources and used in the data-reduction algorithm
- **Approximations and assumptions** incorporated in the measurement procedure
- **Variations of repeated observations** of the measurand under apparently identical conditions

⁹Adapted from the The International Organization for Standardization (ISO) Guide to the Expression of Uncertainty in Measurement.

Optimal presentation of search results

Optimal presentation of search results has some desired properties¹⁰:

- **Uncertainties due to systematic effects should be included in a clear and consistent way.**
 - Often it is useful to quote the statistical and systematical error separately, e.g. $\sigma = 45 \pm 4 \pm 1 \text{ mb}$.
- The result should summarize completely the experiment; so that no extra information should be required for further analysis.
- Results should be easily turned into probabilistic statements.
- Analysis should be transparent, and result should be stated in such a way that it cannot be misleading. The presentation of the result should not depend on the particular application.
- **If possible full pdf-distributions and even data sets can be attached into analysis results.**
- In **unified approach to data analysis**, the transitions between exclusion, observation, discovery, and measurement are kept as small as possible.

¹⁰Adapted from F. James, *Workshop on Confidence Limits*, CERN-2000-005, 2000.

Physicists & statisticians : illustrative comparison

EXAMPLE: HISTOGRAM FITTING:

Physicists

1. Determining the “best fit” parameters of a curve

2. Determining the errors on the parameters

3. Judging the goodness of a fit

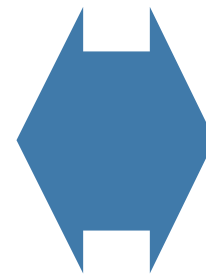
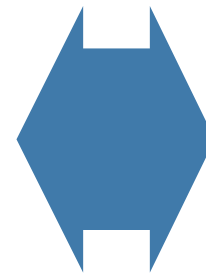
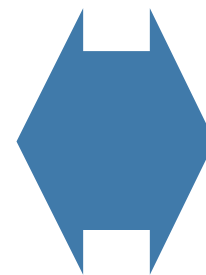
Statisticians

1. Point estimation

2. Confidence interval estimation

3. Goodness-of-fit (hypothesis) testing

discussed later, see references for more info



PART II: Data analysis and extraction of physics information





PART III

Selected topics on the physics measurements

PART III

Selected topics on the physics measurements

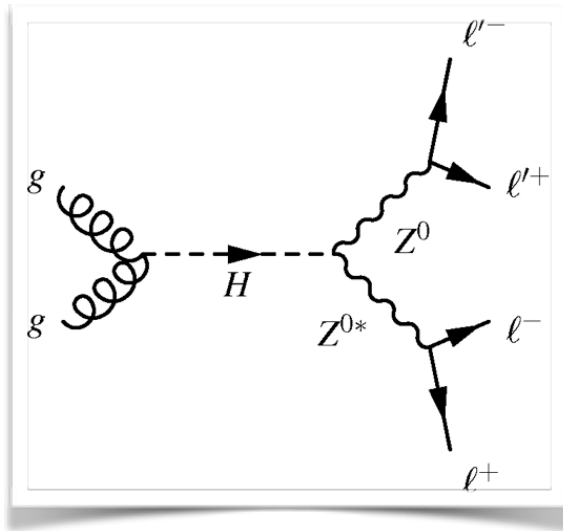
- Estimation of reducible background processes using data
- Matrix Element Method for discrimination between processes
- Exploitation of interference effects in particle physics
- Measurement of fiducial cross sections

Examples based on process: $pp \rightarrow H \rightarrow ZZ \rightarrow 4\ell$

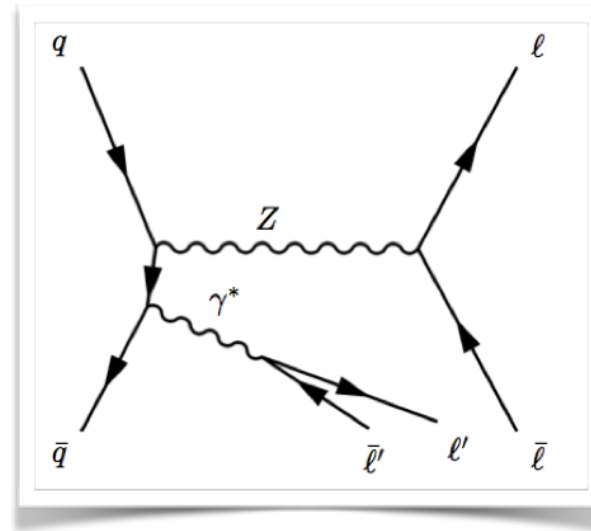
H → ZZ → 4ℓ

- Benefits from fully reconstructible decay with excellent mass resolution (1-2%)

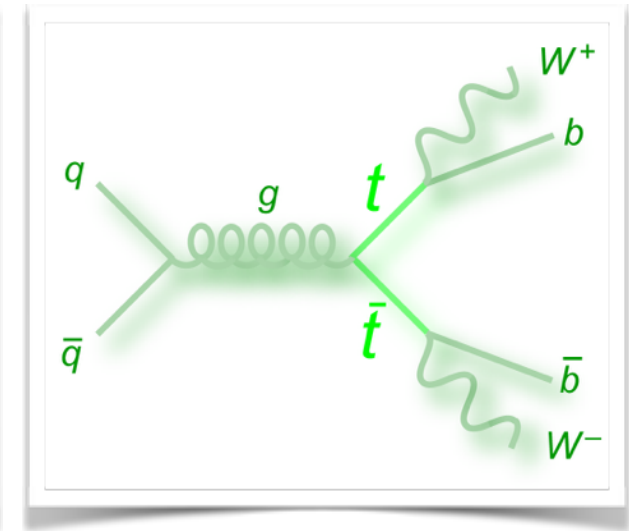
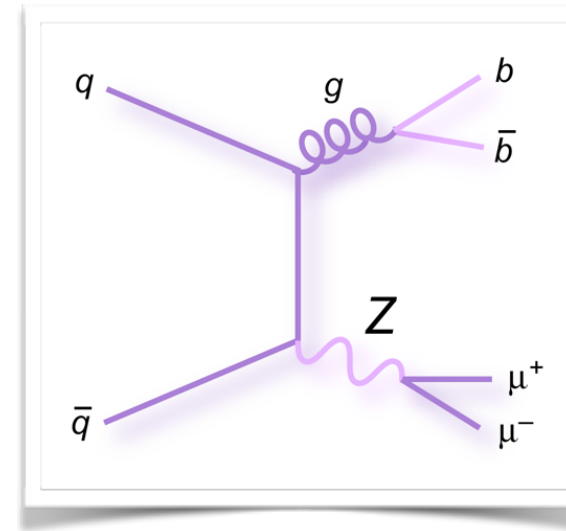
Signal process



Irreducible background

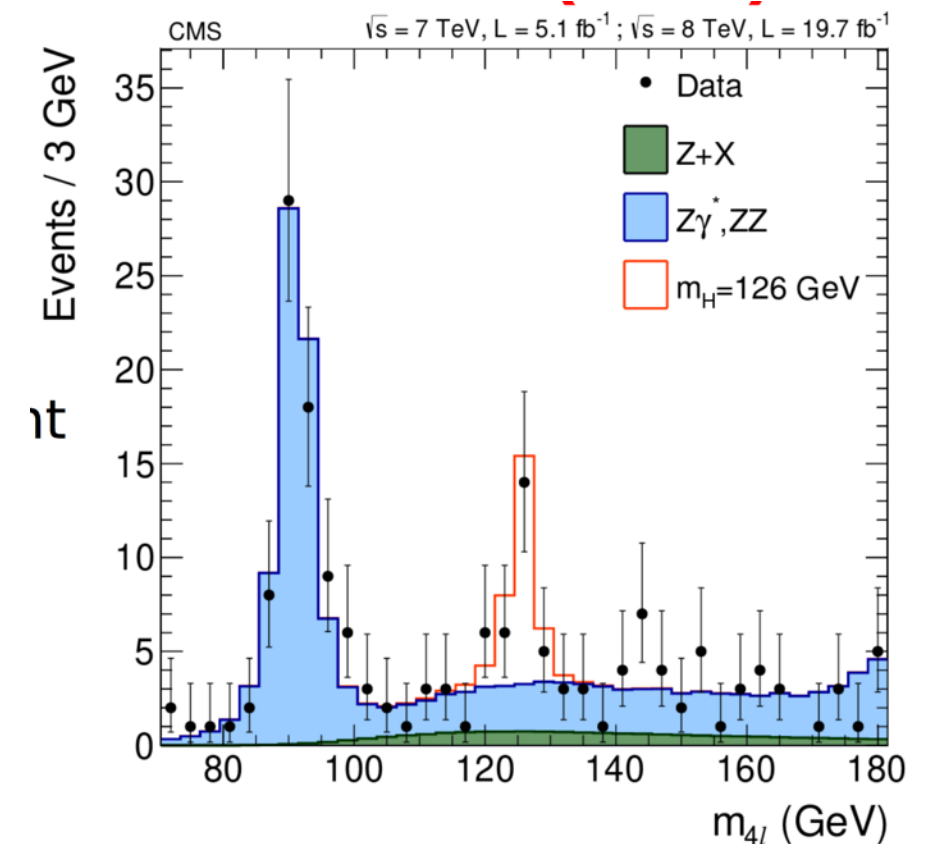


Instrumental backgrounds

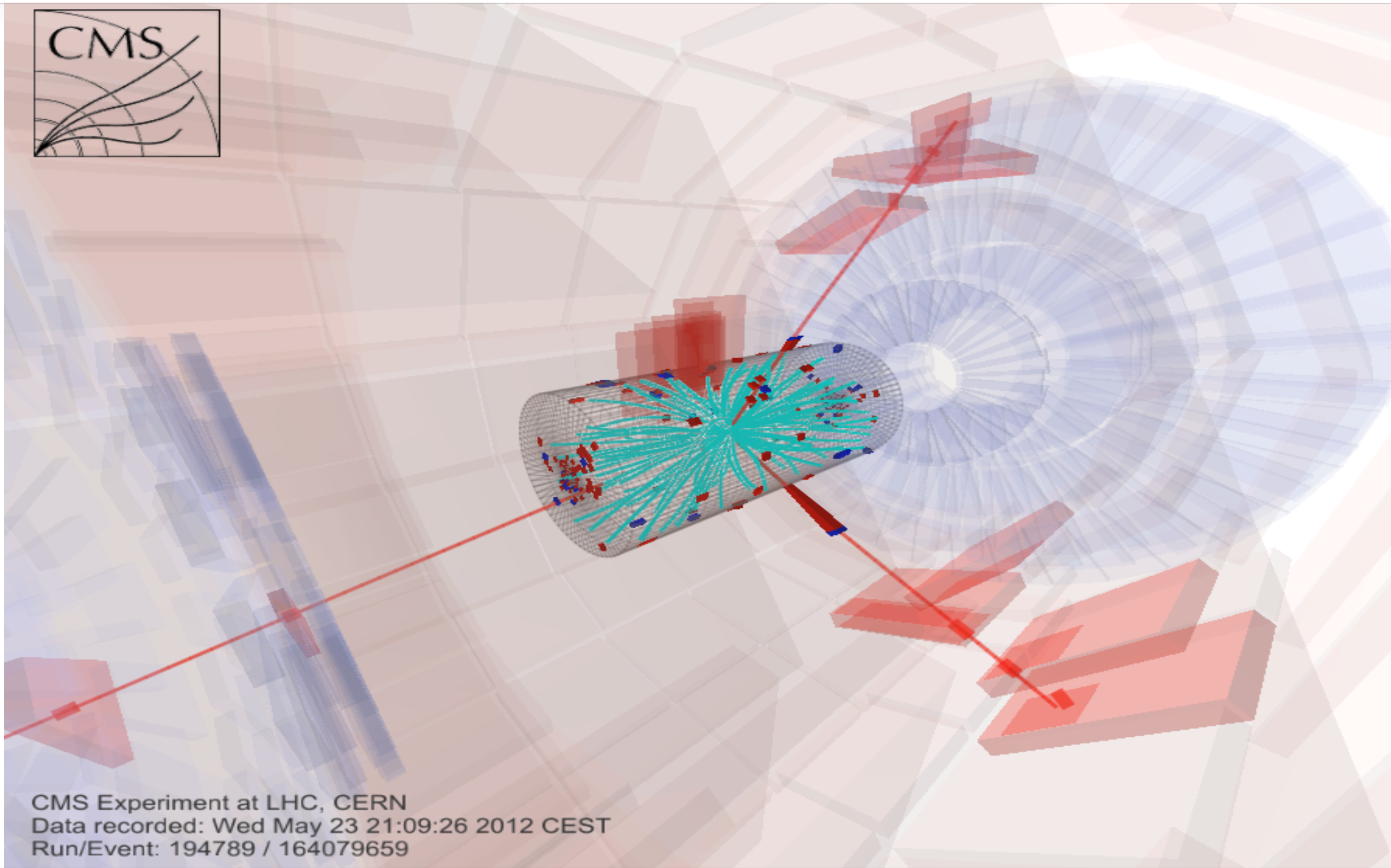


IMPORTANT CHARACTERISTICS:

- Very small branching fraction (~0.01%), but very clean signature
- About 15-20 reconstructed signal events expected after selection (8 TeV Run I)
- Very small background from irreducible $pp \rightarrow ZZ$ and reducible Z +jets.



Event display : $H \rightarrow ZZ \rightarrow 4\ell$



CMS Experiment at LHC, CERN
Data recorded: Wed May 23 21:09:26 2012 CEST
Run/Event: 194789 / 164079659



Estimation of reducible backgrounds using data

Estimation of reducible background

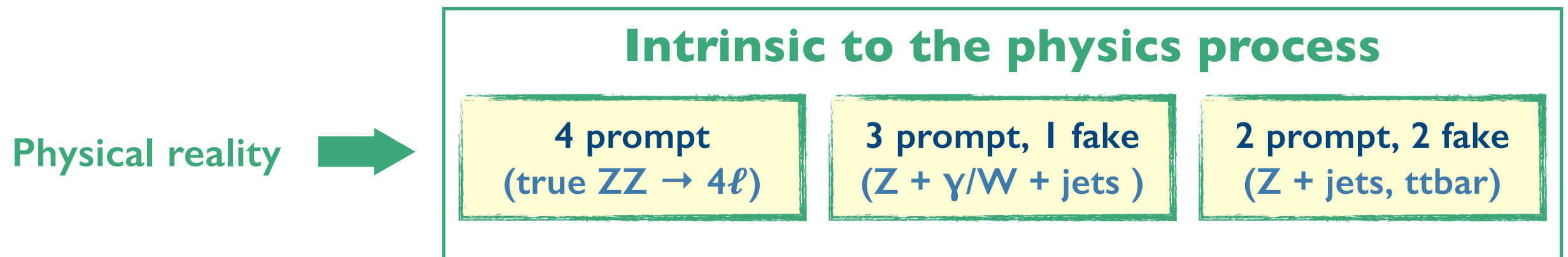
REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

- Events with non-prompt leptons (do not originate from the hard scattering)
 - misidentified / fake leptons (parts of jets, leptons from meson decays)

Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

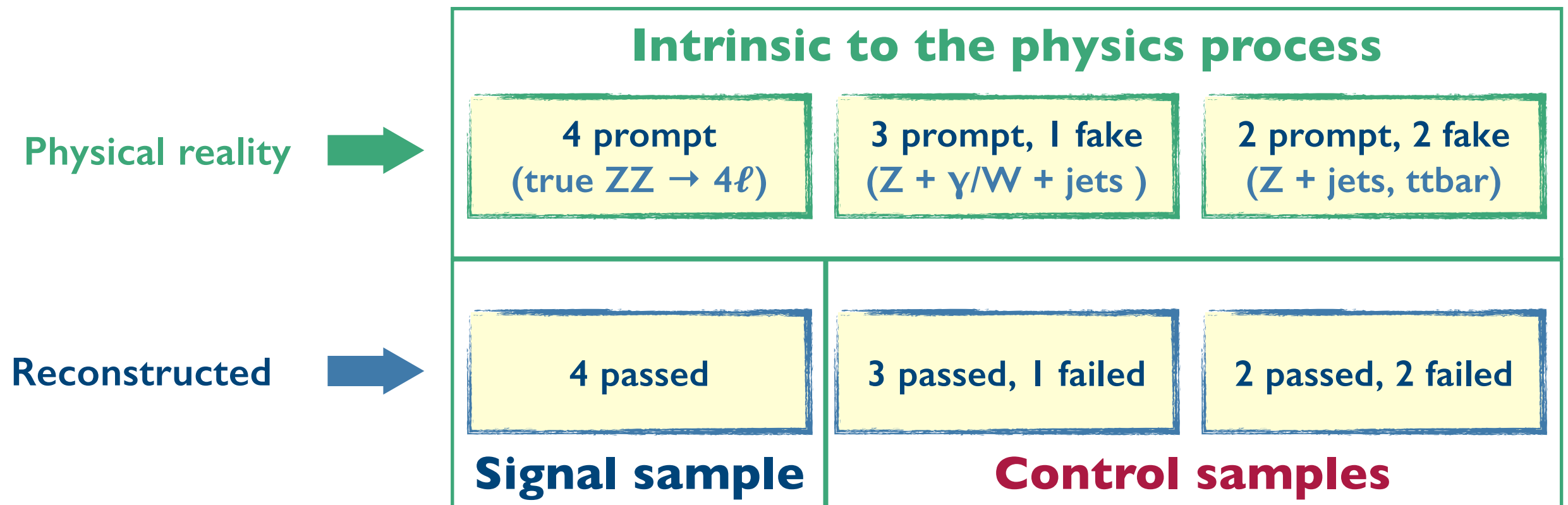
- Events with non-prompt leptons (do not originate from the hard scattering)
 - misidentified / fake leptons (parts of jets, leptons from meson decays)



Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

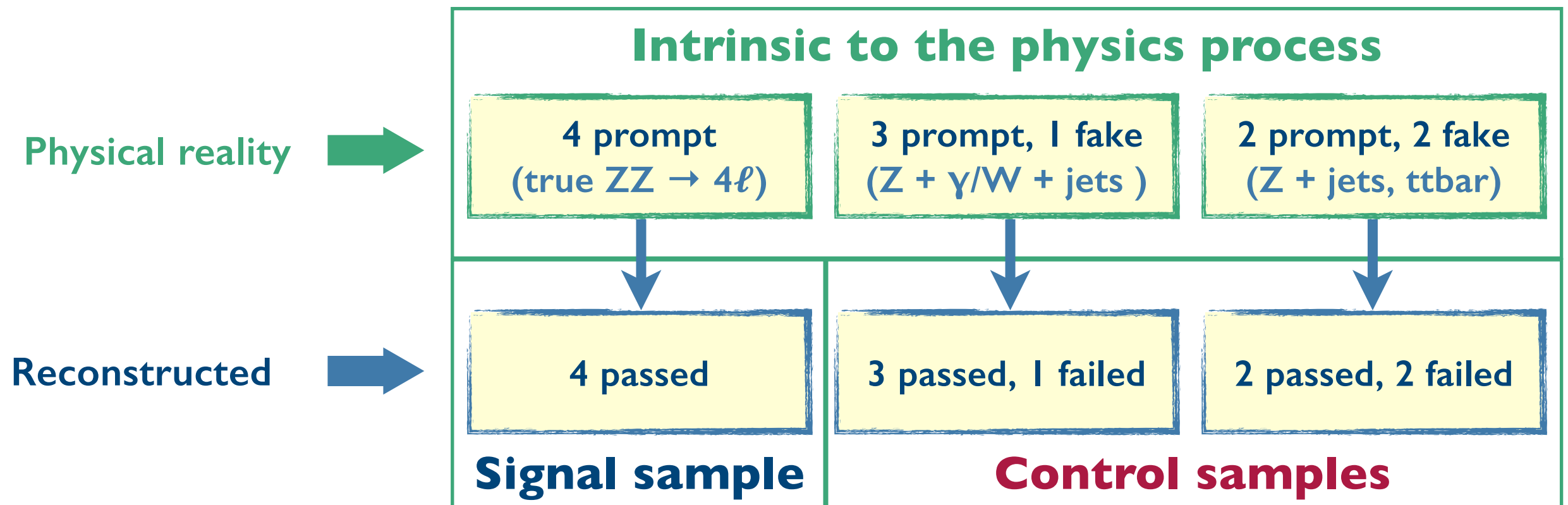
- Events with non-prompt leptons (do not originate from the hard scattering)
 - **misidentified / fake leptons** (parts of jets, leptons from meson decays)



Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

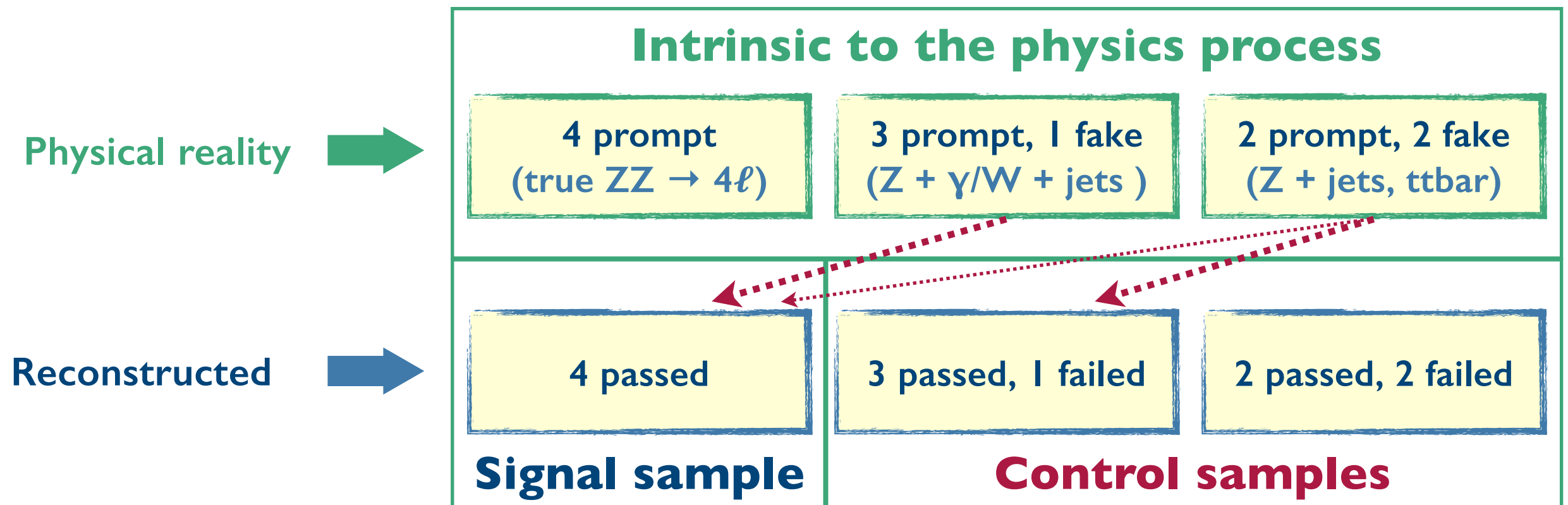
- Events with non-prompt leptons (do not originate from the hard scattering)
 - **misidentified / fake leptons** (parts of jets, leptons from meson decays)



Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

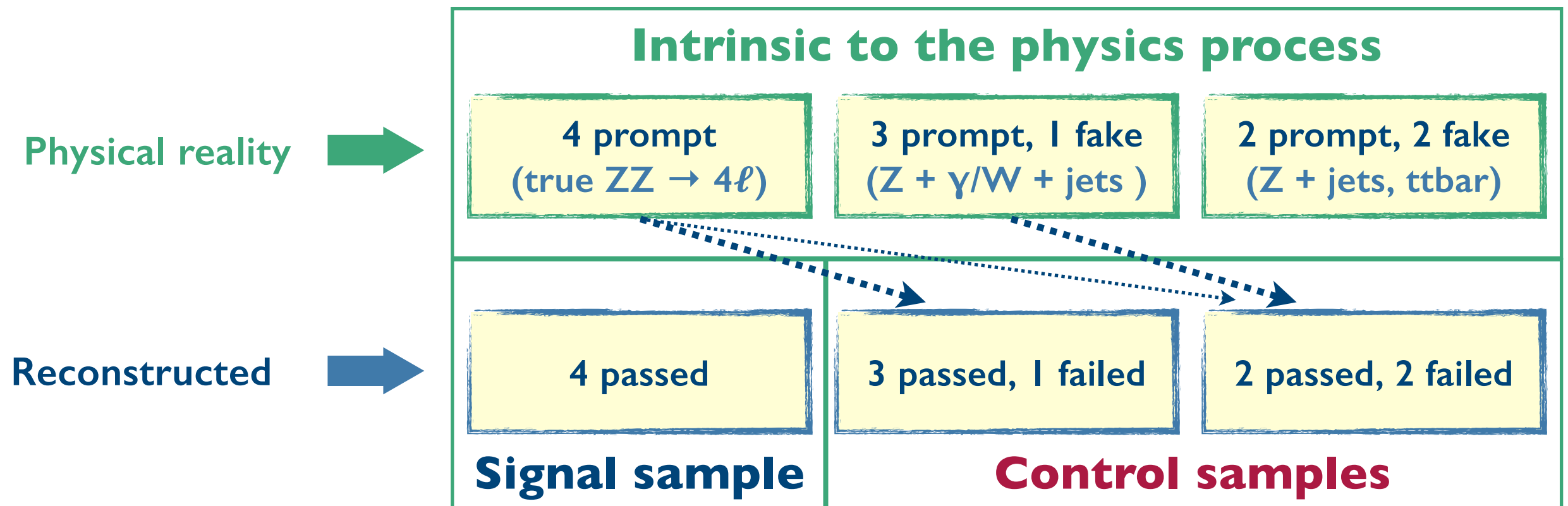
- Events with non-prompt leptons (do not originate from the hard scattering)
 - **misidentified / fake leptons** (parts of jets, leptons from meson decays)



Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

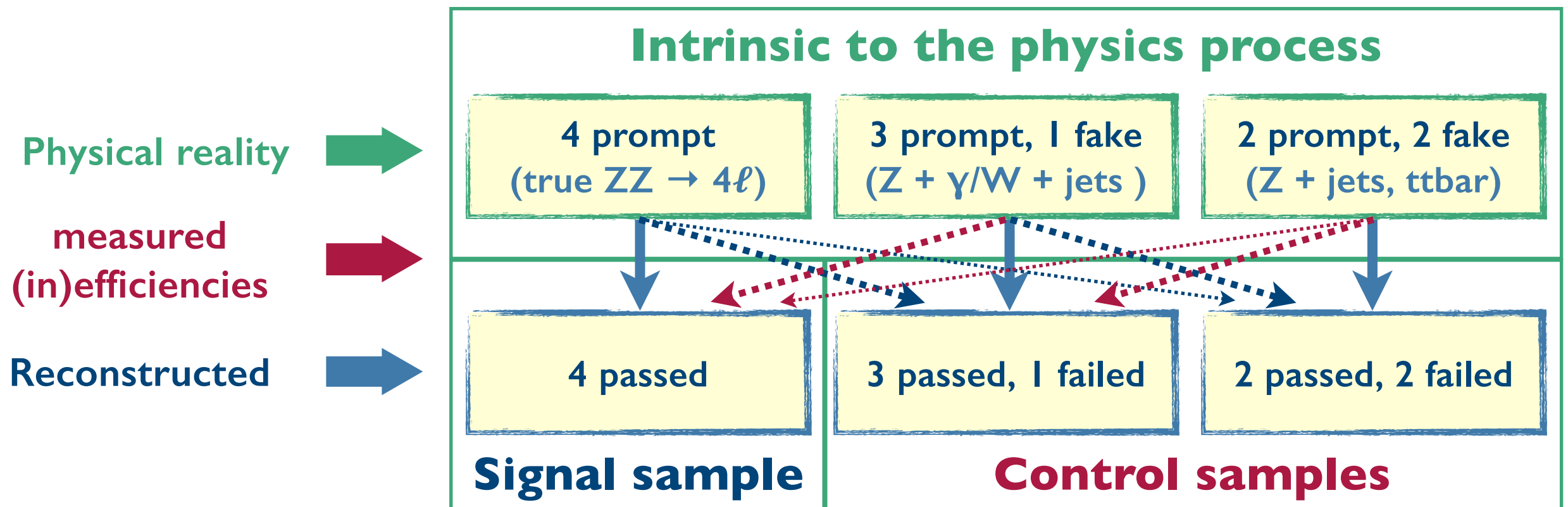
- Events with non-prompt leptons (do not originate from the hard scattering)
 - **misidentified / fake leptons** (parts of jets, leptons from meson decays)



Estimation of reducible background

REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

- Events with non-prompt leptons (do not originate from the hard scattering)
 - **misidentified / fake leptons** (parts of jets, leptons from meson decays)



APPROACH:

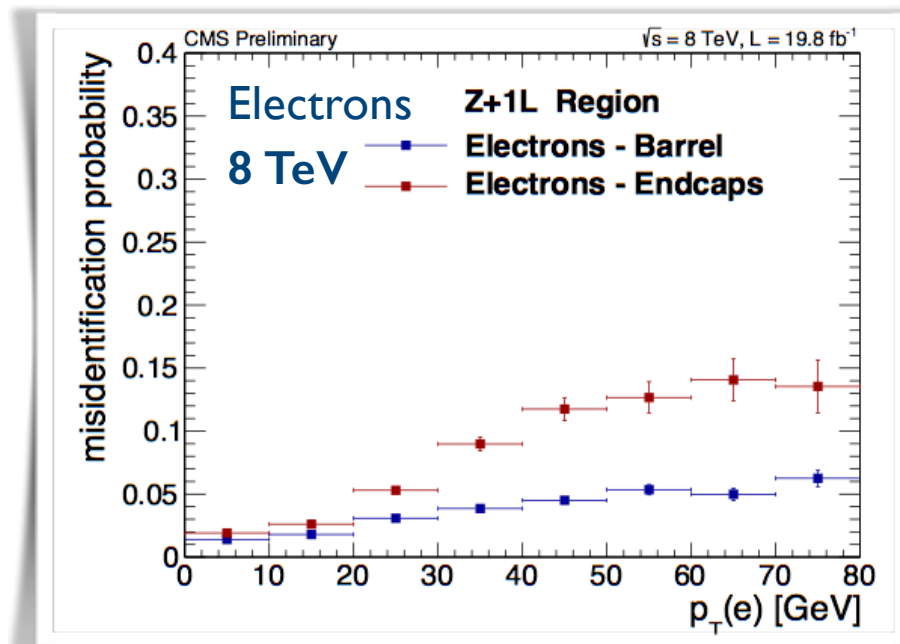
- Estimated from data control samples enriched with misidentified leptons
- Using the **probabilities to misidentify a lepton** ("fake ratios", measured in data)

Misidentification probabilities

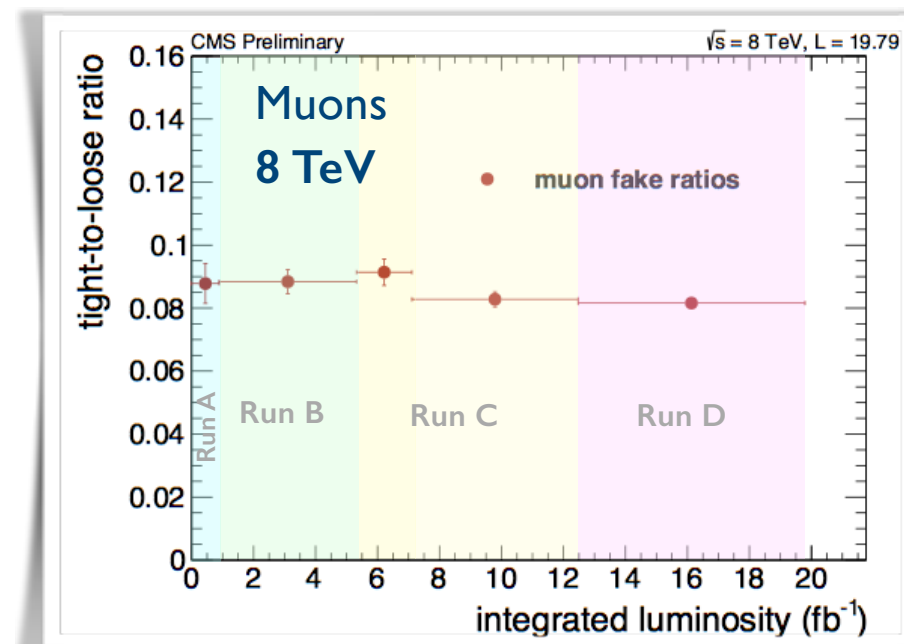
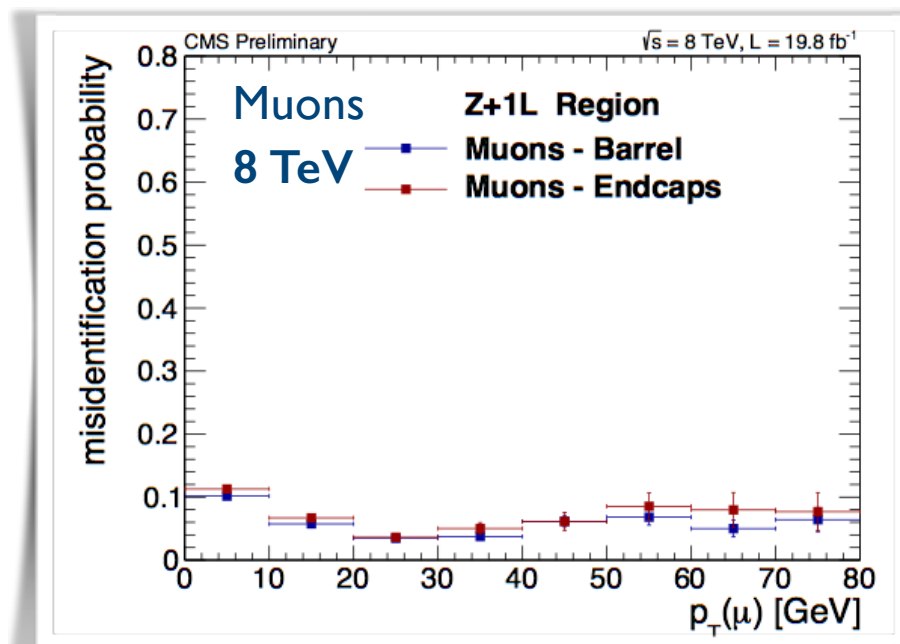
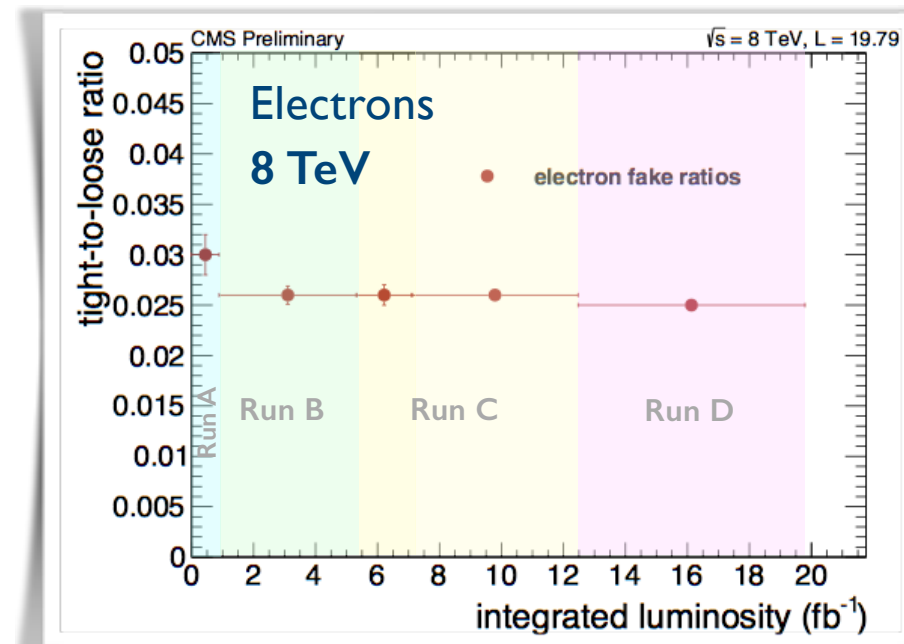
The fake lepton ratios (f_μ , f_e) are measured in sample **Z + 1 loose lepton (Z+1 ℓ)**

- Loose leptons - “analysis selection” leptons with relaxed ID and removed isolation requirements,
- Biases from prompt leptons (Z γ^* , ttbar) suppressed by $|m_{21} - m_Z| < 10$ GeV and miss. $E_T < 25$ GeV

p_T dependance



stability vs. data taking



Estimation of reducible background

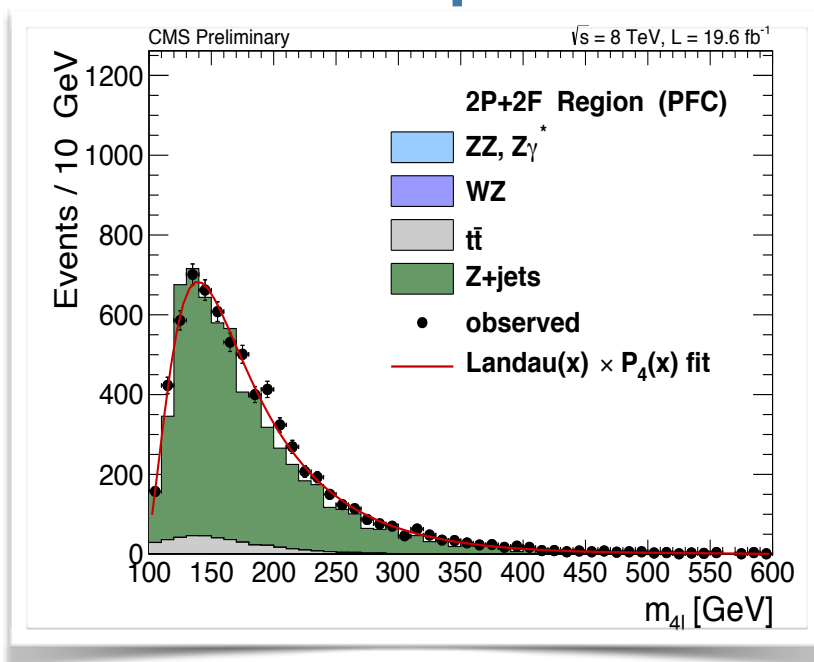
REDUCIBLE BACKGROUND FOR $H \rightarrow ZZ \rightarrow 4\ell$:

- Solve the system of linear equations (using measured “fake rates”):

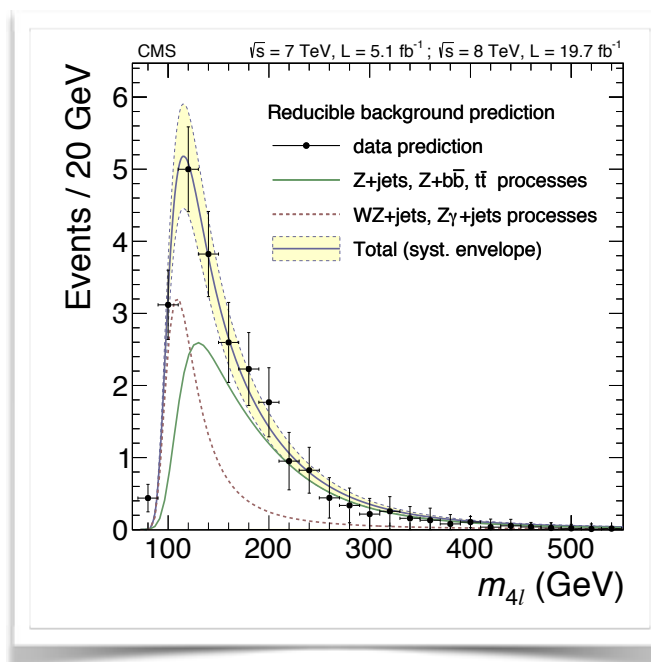
$$N_{\text{SR}}^{\text{bkg}} = \sum \frac{f_i}{1 - f_i} \left(N_{3\text{P}1\text{F}} - N_{3\text{P}1\text{F}}^{(\text{from } 2\text{P}2\text{F})} - N_{3\text{P}1\text{F}}^{(\text{from } ZZ)} \right) + \sum \frac{f_i}{1 - f_i} \frac{f_j}{1 - f_j} N_{2\text{P}2\text{F}} \quad (\text{symbolic})$$

- Examples of control regions and final estimates:

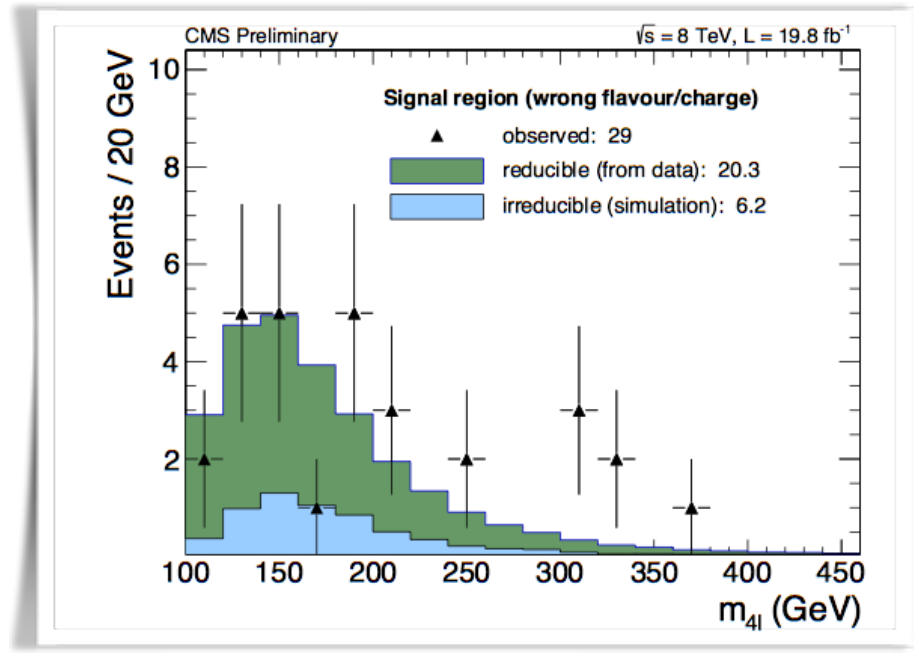
Control sample 2P+2F



Final estimate



Validation in data



- Sources of **systematic uncertainties**:

- Limited statistics in control regions,
- Different background composition in region $Z+1\ell$, and regions $2\text{P}+2\text{F}$ and $3\text{P}+1\text{F}$
- Validation in data** using events with “wrong” flavour/charge combination

Estimation of reducible backgrounds using data

To remember...

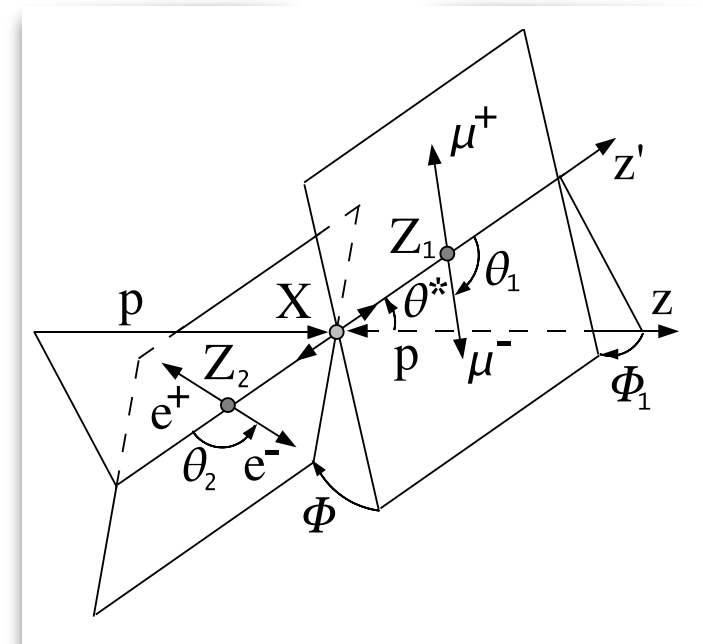
1. Reducible backgrounds can be estimated from data control samples enriched with misidentified objects by solving set of linear (recursive) equations,
2. Measurement of misidentification probabilities, and systematics need special care.



Matrix Element Method for processes discrimination

Matrix Element Method (MEM)

$$H \rightarrow ZZ \rightarrow 4l$$



Matrix element for process X - $|\text{ME}(X)|^2$:

- **Probability density** function for event of process X to occur in a given point of phase space
- **Inverted logic: Likelihood** that event observed in a given point of phase space originates from process X

Advantages:

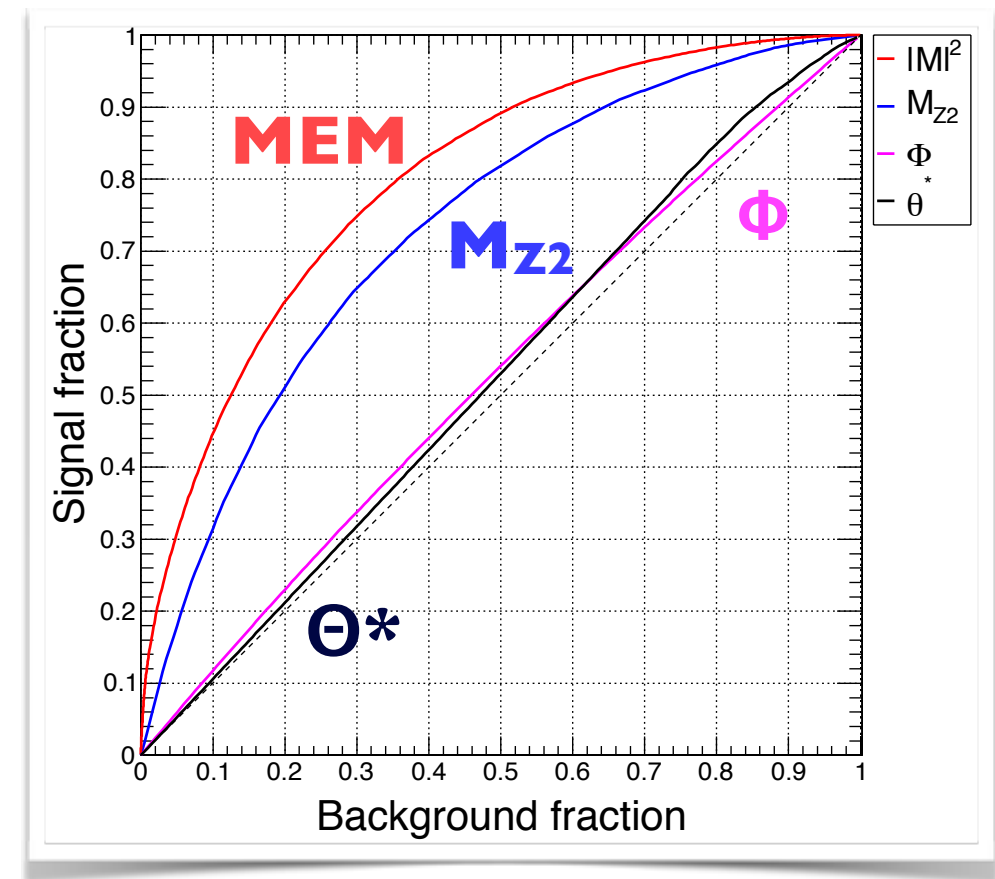
- Efficiently uses all available kinematic information
- Greater sensitivity than any other method

Example:

- ROC curves from MEM and 3 analyses using a single variable

Discrimination of processes A & B:

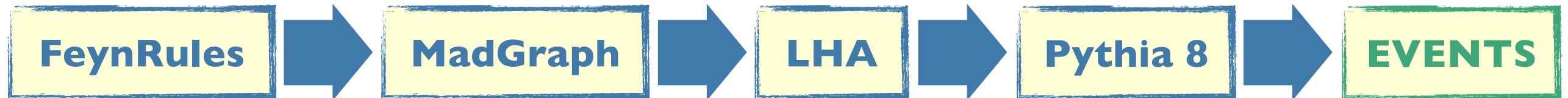
- Maximal discrimination between processes A & B from the ratio of $|\text{ME}(A)|^2$ and $|\text{ME}(B)|^2$



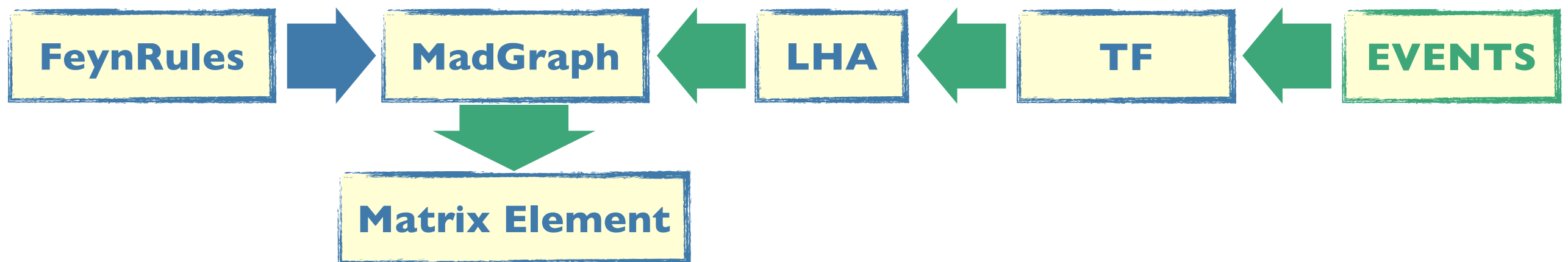
MEM & MC Simulation tools

MC Simulation tools - an example:

- From Lagrangians to events:



- From events to ... matrix elements:



- Automatic generation and calculation of matrix elements
- Many MC tools available on the market (MadGraph, MCFM, JHUGen, etc.)

Application:

- Separation of signal and background, and alternative signal hypotheses
- In case of final states with invisible or poorly reconstructed objects (neutrinos, jets) need to take into account detector effects (TF) or integrate out invisible degrees of freedom

H → ZZ → 4l: Kinematic Discriminants

- **Matrix Element Method:** Use ratio of LO matrix elements $|ME|^2$ to build discriminants
 - do not use system p_T and rapidity Y (NLO effects, PDFs)
 - use the assumption: $m_X = m_{4l}$

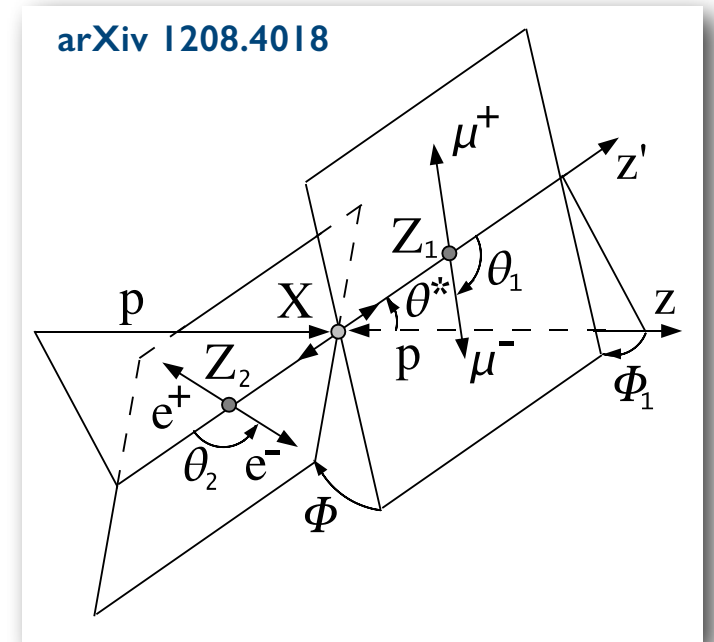
Basic ME-discriminator to separate SM Higgs from backgrounds:

$$KD(H;ZZ) = \frac{|ME_H(gg \rightarrow H \rightarrow 4l)|^2}{|ME_{ZZ}(q\bar{q} \rightarrow 4l)|^2}$$

Basic ME-discriminator to separate alternative J^P hypothesis from bkg.:

$$KD(J^{CP};ZZ) = \frac{|ME_{J^{CP}}(xx \rightarrow J^{CP} \rightarrow 4l)|^2}{|ME_{ZZ}(q\bar{q} \rightarrow 4l)|^2}$$

Use kinematics of the 4l system



H → ZZ → 4l: Kinematic Discriminants

- **Matrix Element Method:** Use ratio of LO matrix elements $|ME|^2$ to build discriminants
 - do not use system p_T and rapidity Y (NLO effects, PDFs)
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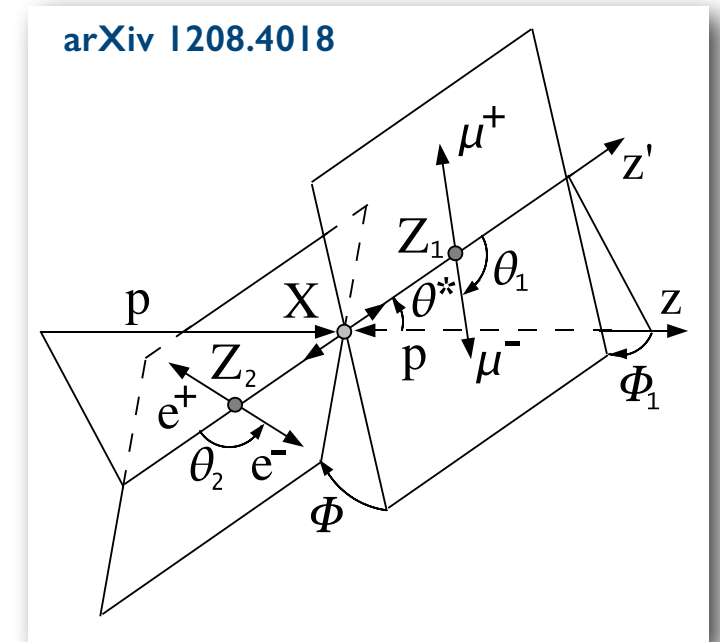
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$$KD(J^{CP};ZZ) = \frac{|ME_{J^{CP}}(xx \rightarrow J^{CP} \rightarrow 4\ell)|^2}{|ME_{ZZ}(q\bar{q} \rightarrow 4\ell)|^2}$$

Use kinematics of the 4l system



- **Extend discriminators to include the discriminating m_{4l} information:**

Extended discriminator to separate SM Higgs from backgrounds:

$$D(H;ZZ) = \frac{|ME_X(xx \rightarrow H \rightarrow 4\ell)|^2 \cdot pdf(m_{4\ell} | m_H)}{|ME_{ZZ}(q\bar{q} \rightarrow 4\ell)|^2 \cdot pdf(m_{4\ell} | ZZ)}$$

Extended discriminator to separate an alternative J^P hypothesis from backgrounds:

$$D(J^{CP};ZZ) = \frac{|ME_{J^{CP}}(xx \rightarrow J^{CP} \rightarrow 4\ell)|^2 \cdot pdf(m_{4\ell} | m_{J^{CP}})}{|ME_{ZZ}(q\bar{q} \rightarrow 4\ell)|^2 \cdot pdf(m_{4\ell} | ZZ)}$$

H → ZZ → 4l: Kinematic Discriminants

- **Matrix Element Method:** Use ratio of LO matrix elements $|ME|^2$ to build discriminants
 - do not use system p_T and rapidity Y (NLO effects, PDFs)
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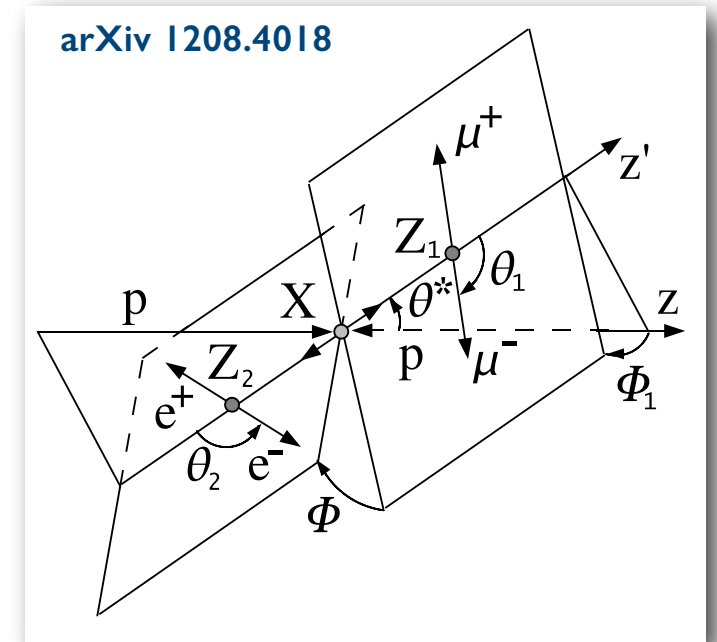
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$$KD(J^{CP};ZZ) = \frac{|ME_{J^{CP}}(xx \rightarrow J^{CP} \rightarrow 4l)|^2}{|ME_{ZZ}(q\bar{q} \rightarrow 4l)|^2}$$

Use kinematics of the 4l system



- **Extend discriminators to include the discriminating m_{4l} information:**

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$$D(H;ZZ) = \frac{|ME_X(xx \rightarrow H \rightarrow 4l)|^2 \cdot pdf(m_{4l} | m_H)}{|ME_{ZZ}(q\bar{q} \rightarrow 4l)|^2 \cdot pdf(m_{4l} | ZZ)}$$

change the variables

(without loss of information)

$$D(H;ZZ)$$

Extended discriminator to separate an alternative J^P hypothesis from backgrounds:

$$D(J^{CP};ZZ) = \frac{|ME_{J^{CP}}(xx \rightarrow J^{CP} \rightarrow 4l)|^2 \cdot pdf(m_{4l} | m_{J^{CP}})}{|ME_{ZZ}(q\bar{q} \rightarrow 4l)|^2 \cdot pdf(m_{4l} | ZZ)}$$

change the variables

(without loss of information)

$$D(J^{CP};H) = \frac{D(J^{CP};ZZ)}{D(H;ZZ)}$$

H → ZZ → 4l: Kinematic Discriminants

- **Matrix Element Method:** Use ratio of LO matrix elements $|ME|^2$ to build discriminants
 - do not use system p_T and rapidity Y (NLO effects, PDFs)
 - use the assumption: $m_X = m_{4l}$
 - Final discriminators D_{JCP} and D_{BKG} obtained by compressing $D(J^{CP};H)$ and $D(H;ZZ)$ between 0 and 1:

Discriminator D_{JP} to separate SM Higgs from an alternative J^P hypothesis

$$D_{JP} = \left[1 + const. \cdot \frac{|ME_{JP}(\vec{p}_i)|^2}{|ME_H(\vec{p}_i)|^2} \right]^{-1}$$

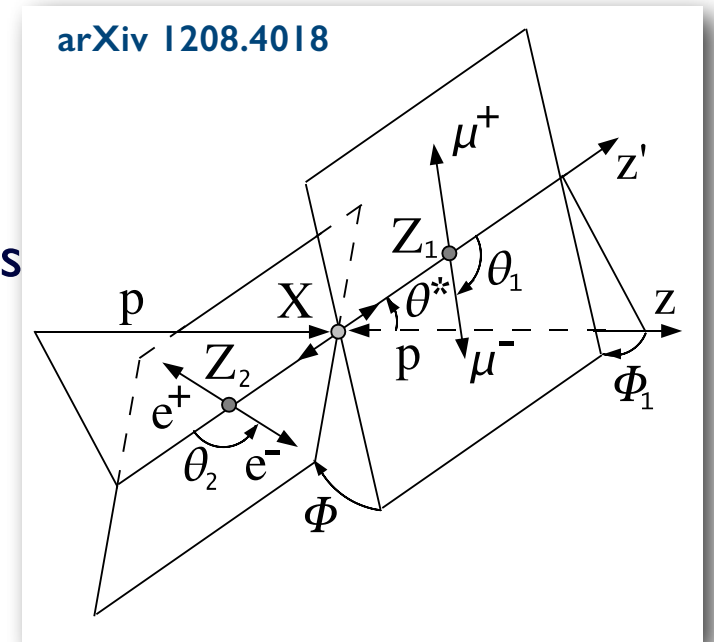
Discriminator D_{BKG} to separate signal(s) from backgrounds:

$$D_{BKG} = \left[1 + const. \cdot \frac{|ME_{ZZ}(\vec{p}_i)|^2 \cdot pdf(m_{4l}|ZZ)}{|ME_H(\vec{p}_i)|^2 \cdot pdf(m_{4l}|H)} \right]^{-1}$$

- LO MEs are computed using **JHUGen** (signal) and **MC FM** ($qq \rightarrow ZZ$) in **MELA** package
- Common subset of MEs validated with **MEKD** (**FeynRules** + **Madgraph**)

REFERENCES: arXiv 1210.0896 , arXiv 1001.3396 , arXiv 1108.2274, arXiv 1208.4018, arXiv 1211.1959

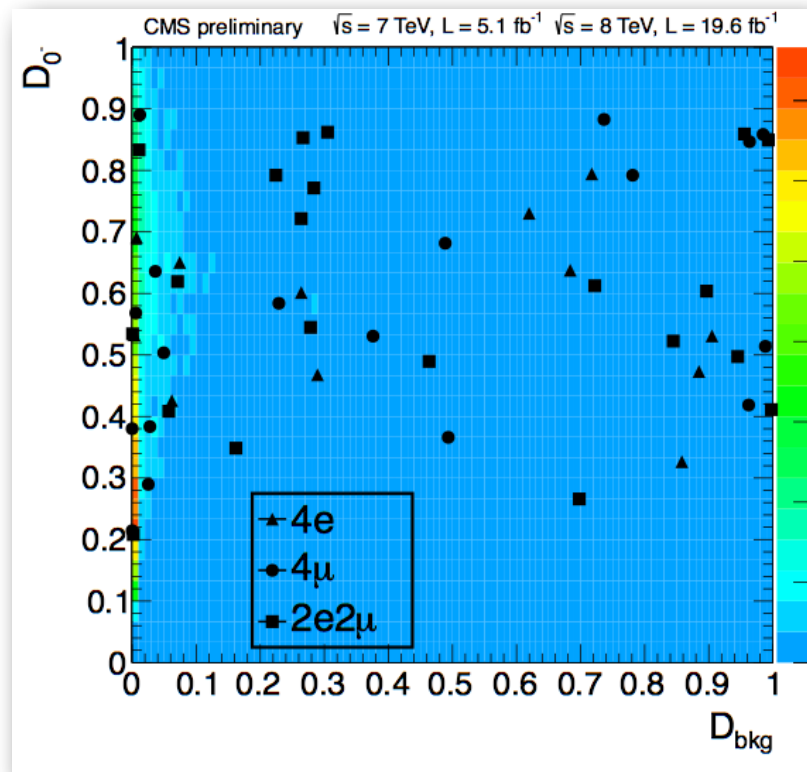
Use kinematics of the 4l system



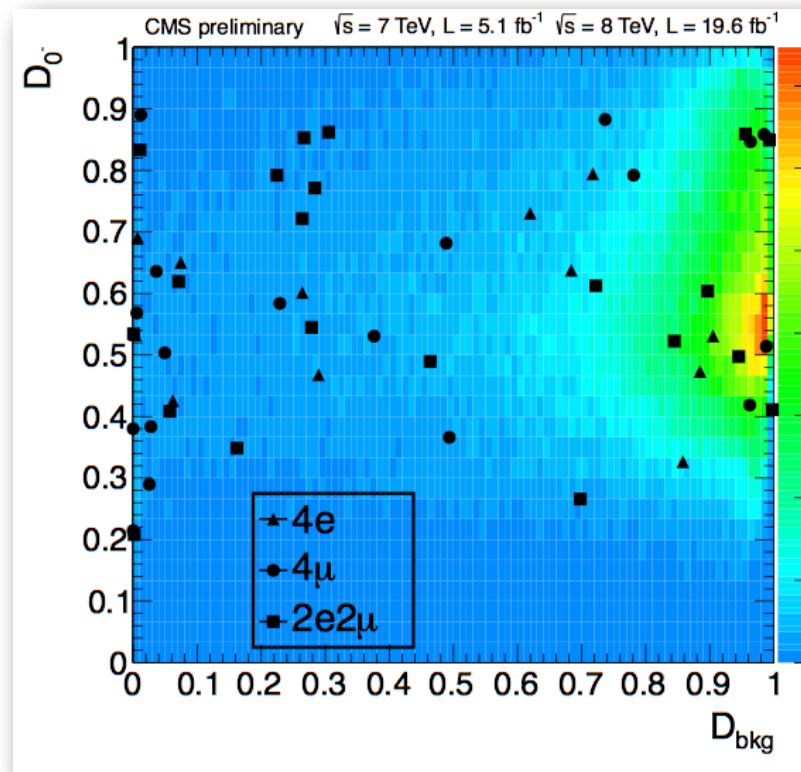
$H \rightarrow ZZ \rightarrow 4l$: Discrimination in $(D_{\text{BKG}}, D_{\text{JP}})$ plane

- Analysis performed for each alternative J^P hypotheses using 2D templates:
 - SM Higgs, alternative signal, qq/gg \rightarrow ZZ from simulation, Z+X from control region in data

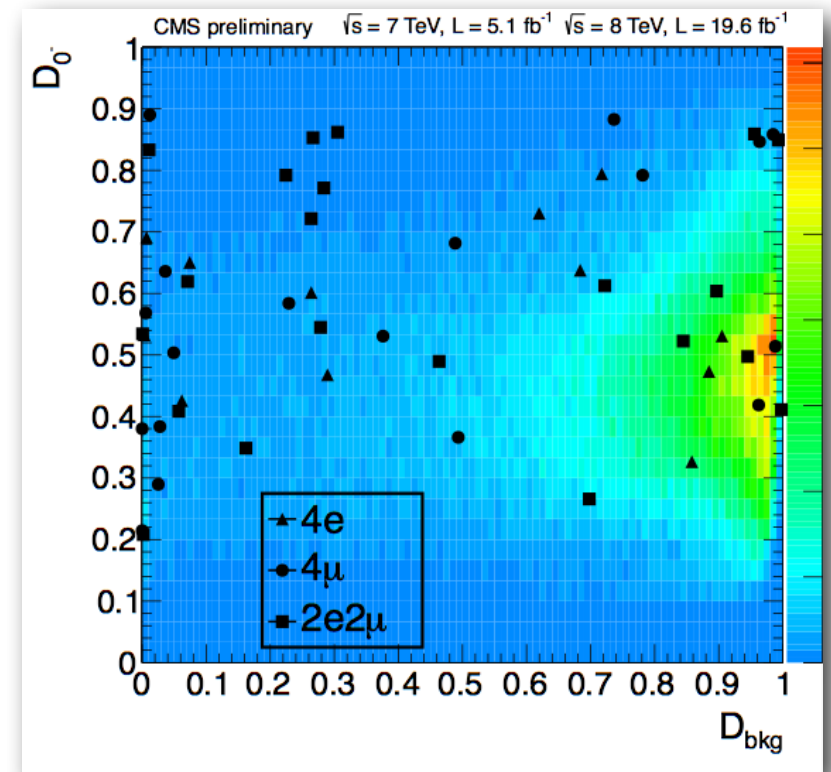
Background processes



SM Higgs boson



exotic Higgs-like boson

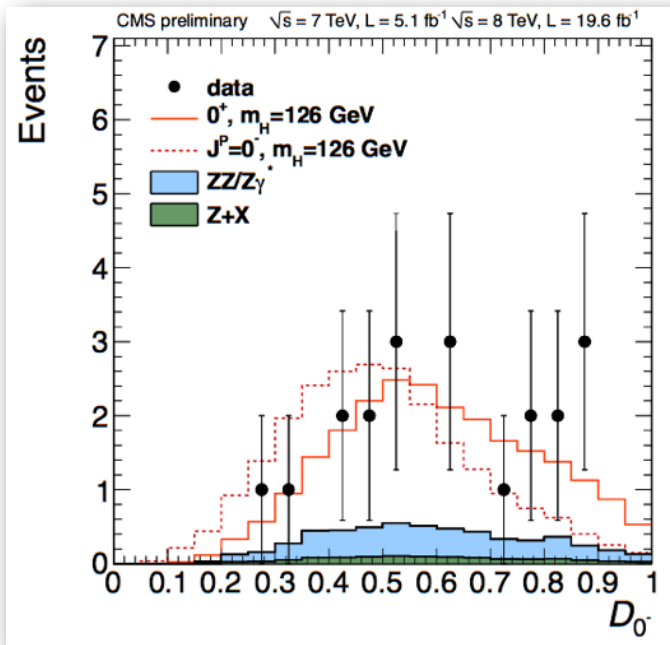


- Perform statistical analysis (hypothesis testing) by generating pseudo-observations (using $\mathcal{P}(D_{\text{JP}}, D_{\text{BKG}})$ distributions and “log likelihood ratio” as the test statistics)

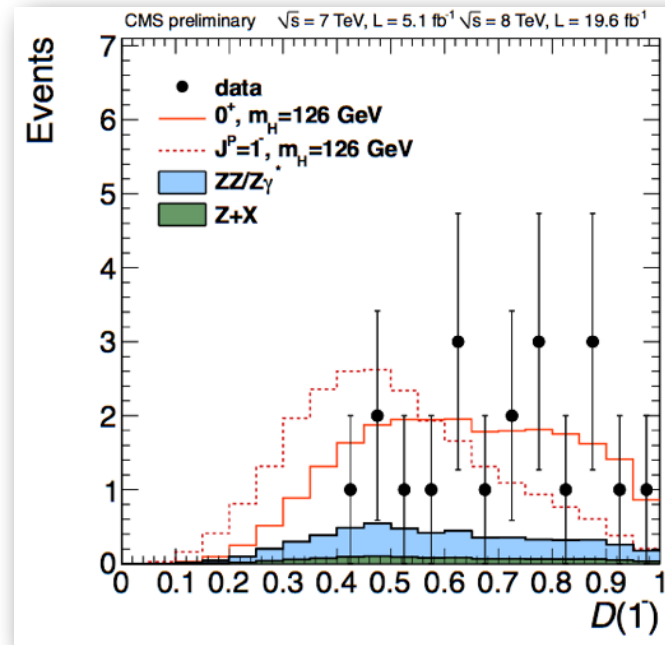
H → ZZ → 4l: Discriminator D_{J^P} ($D_{\text{BKG}} > 0.5$)

- $D_{\text{BKG}} > 0.5$ cut is just for illustration

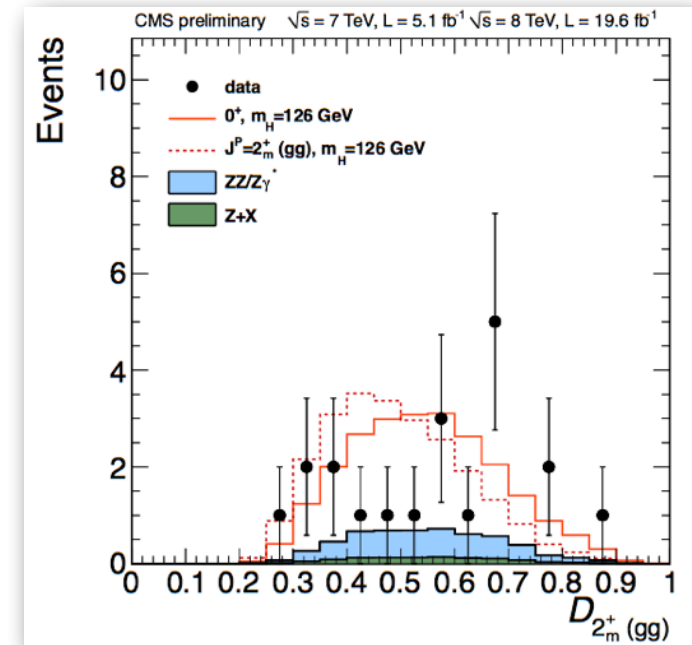
$gg \rightarrow 0^-$



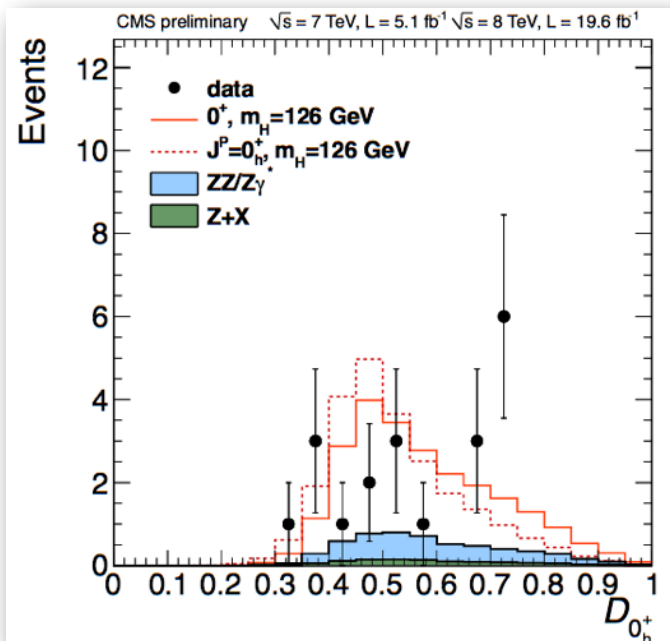
$qq \rightarrow 1^-$



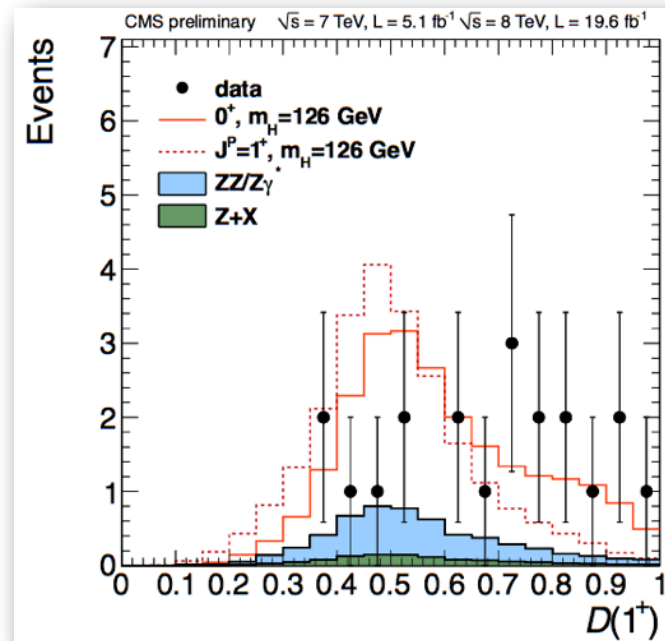
$gg \rightarrow 2_m^+$



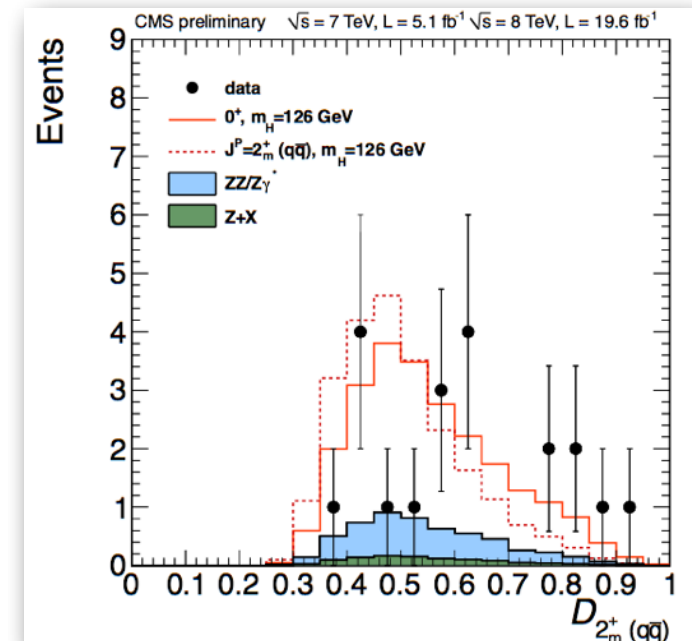
$gg \rightarrow 0_h^+$



$qq \rightarrow 1^+$

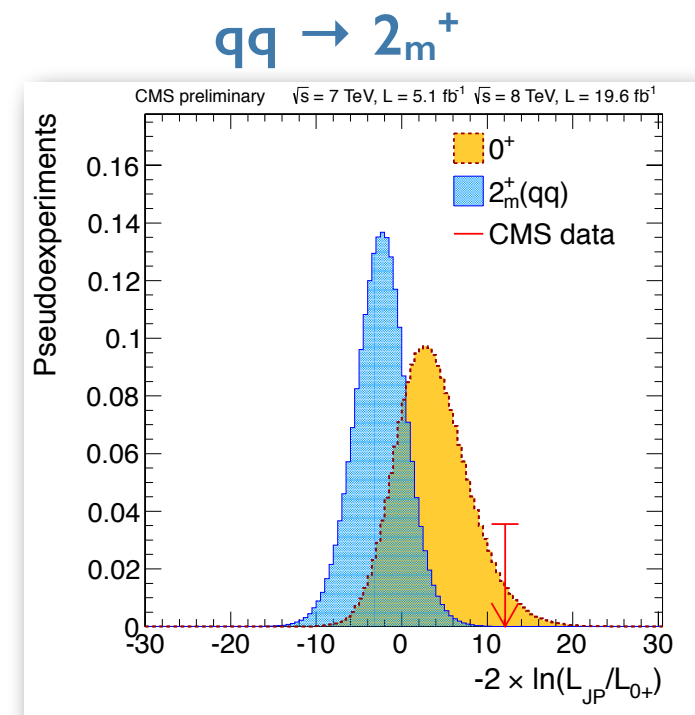
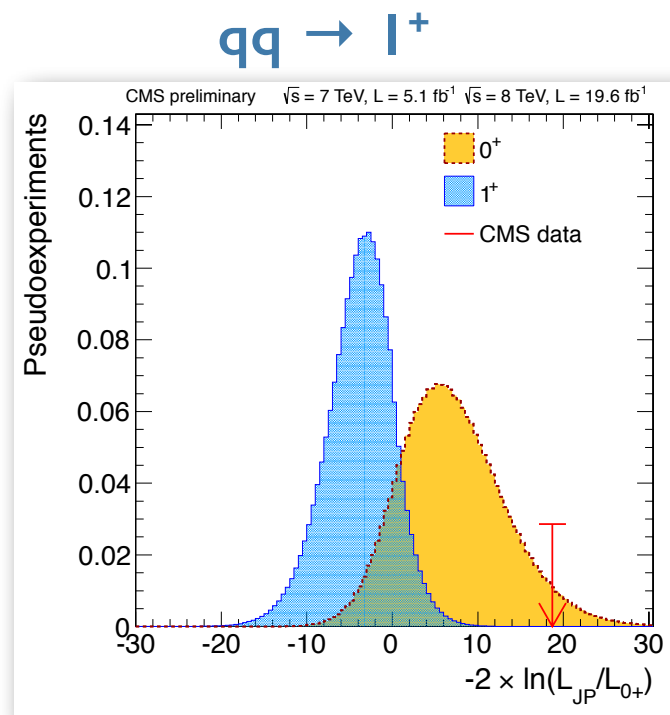
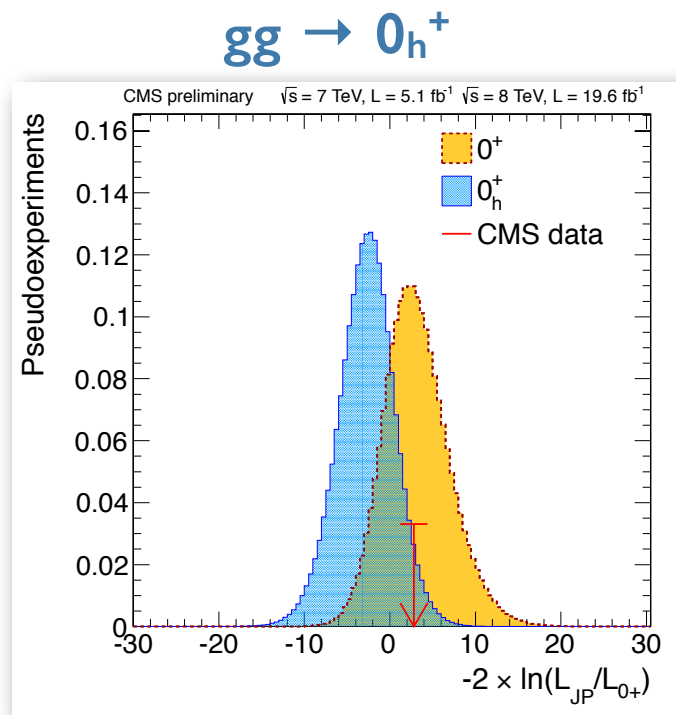
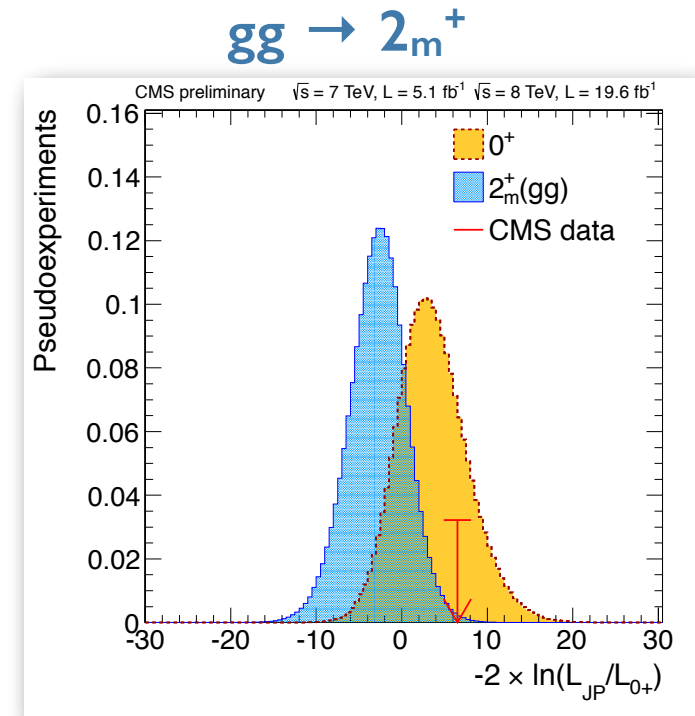
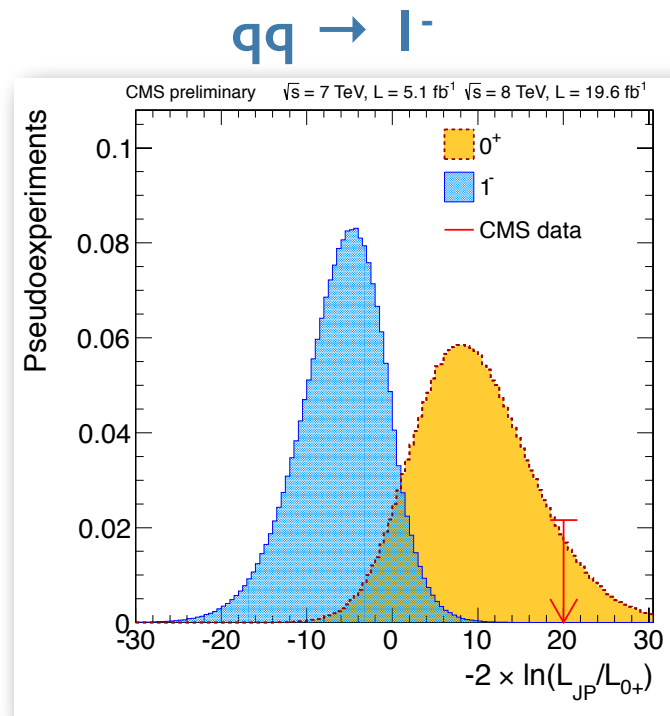
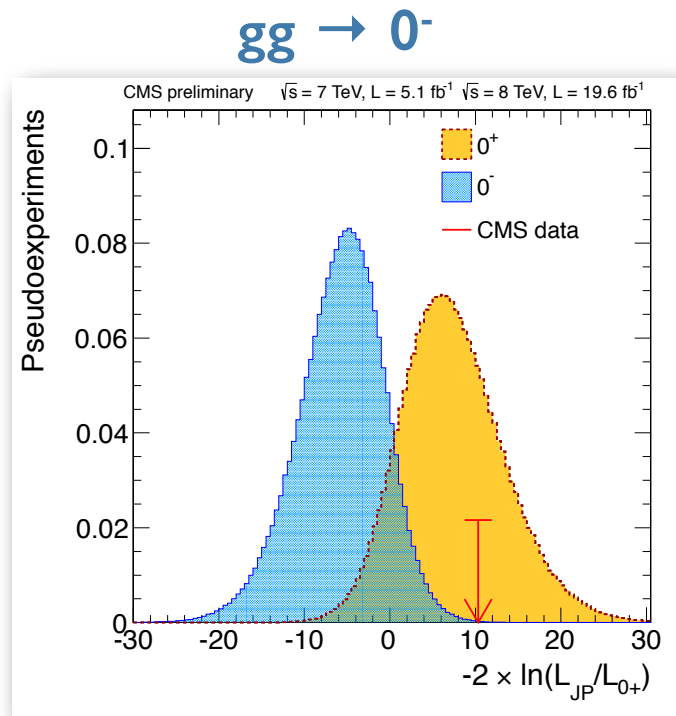


$qq \rightarrow 2_m^+$



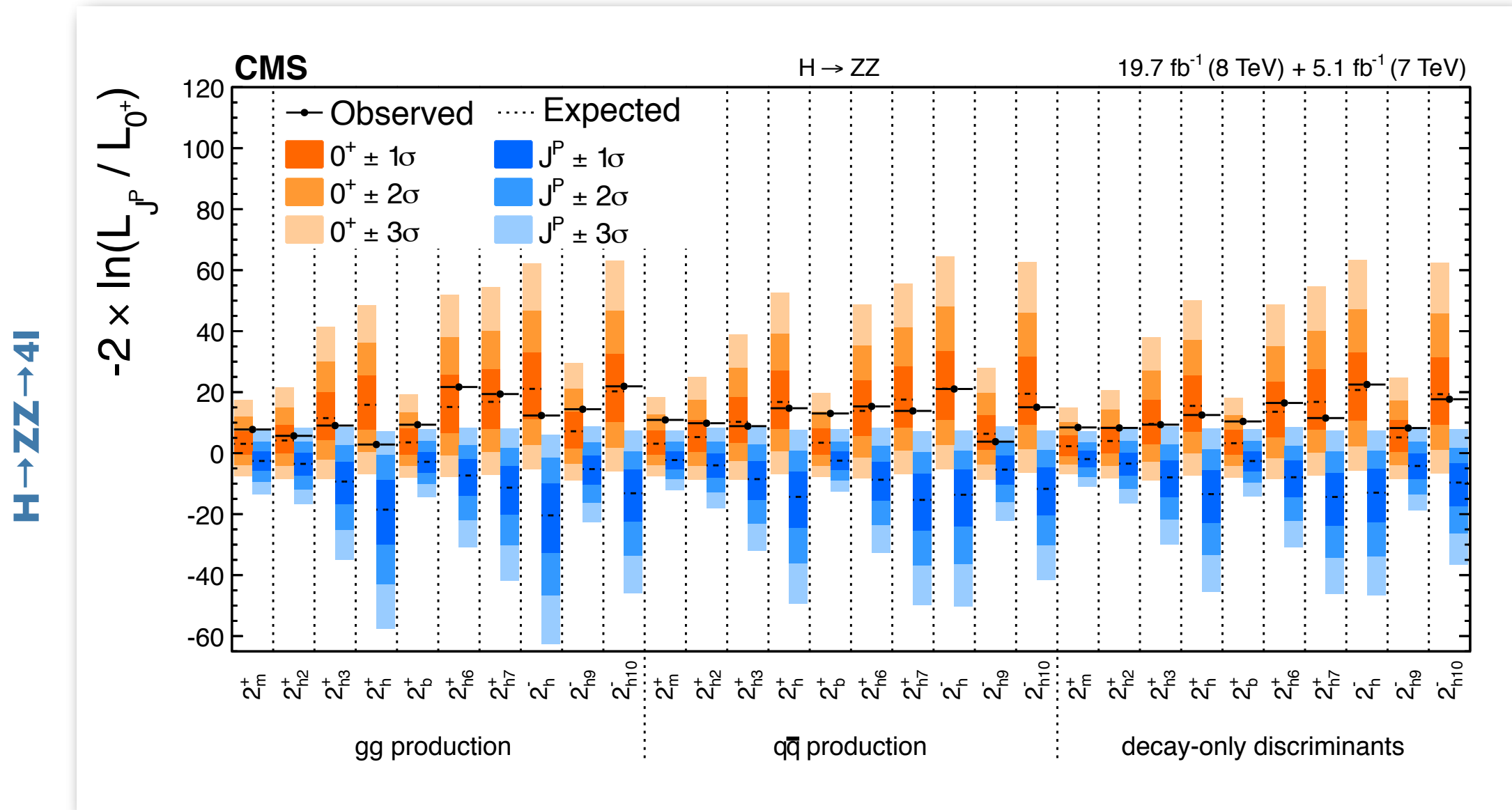
H → ZZ → 4l: J^{CP} hypotheses testing

- Test statistics for J^P hypotheses and the observed results (μ from data)



H → ZZ → 4l: Example of spin-two scenario exclusions

- **H → ZZ → 4l**: Test the pure state spin-two terms (qq production, gg production and using *production independent discriminants*):



**Excluded all pure state spin-two hypotheses
at 96.9% CL or better!**

Matrix Element Method for processes discrimination

To remember...

1. MEM offers very good discrimination of signal and background, and alternative signal hypotheses, by exploiting (all) available kinematic information,
2. Many tools and packages for automatic ME calculation publicly available,
3. Less efficient for final states with invisible or poorly reconstructed objects (jets, miss. E_T) - it needs special care.



Exploitation of interference effects in particle physics


Interference in particle physics

Interference of the process amplitudes:

- Amplitudes of the processes with same initial and final state **interfere**

$$\underbrace{|\text{ME}(A + B)|^2}_{\text{total process A+B}} = \underbrace{|\text{ME}(A)|^2}_{\text{pure process A}} + \underbrace{[\text{ME}(A)^*\text{ME}(B) + \text{ME}(B)^*\text{ME}(A)]}_{\text{interference term}} + \underbrace{|\text{ME}(B)|^2}_{\text{pure process B}}$$

Some examples:

- Interference associated with permutations of identical leptons in the $4e$ and 4μ final states
 - important improvement in separation power in $H \rightarrow ZZ \rightarrow 4l$ channel ([arxiv: 1210.0896v2](https://arxiv.org/abs/1210.0896v2))
- Interference effects between different $H \rightarrow ZZ$ amplitudes:
 - dramatic impact on differential distributions in case of scalar Higgs ([arxiv: 1310.1397](https://arxiv.org/abs/1310.1397))
 -  motivation to consider/exploit these interference effects (topic of this talk)
- Interference effects and the Higgs boson off-shell production

Example: Higgs effective couplings and interference

- Effective Lagrangian associated with ZZ-decays of a scalar particle X (on shell production)

$$\mathcal{L}_{\text{HZZ}} \ni \kappa_1 \frac{m_Z^2}{v} X Z_\mu Z^\mu + \frac{\kappa_2}{2v} X Z_{\mu\nu} Z^{\mu\nu} + \frac{\kappa_3}{2v} X Z_{\mu\nu} \tilde{Z}^{\mu\nu}$$

κ_1 - term (\mathcal{O}_1)
(SM Higgs-like)

κ_2 - term (\mathcal{O}_2)
(0_h^+ , loop-induced)

κ_3 - term (\mathcal{O}_3)
(0^- , loop-induced)

κ_2 and κ_3 terms - the lowest dim. effective operators with the given symmetry properties

- Magnitude of interference effects:

- Impact of interference on the total production rate

[arxiv:1304.4936 \[hep-ph\]](#)

$$\Gamma(X \rightarrow ZZ) = \Gamma_{SM} \sum_{i,j} \gamma_{ij} \kappa_i \kappa_j,$$

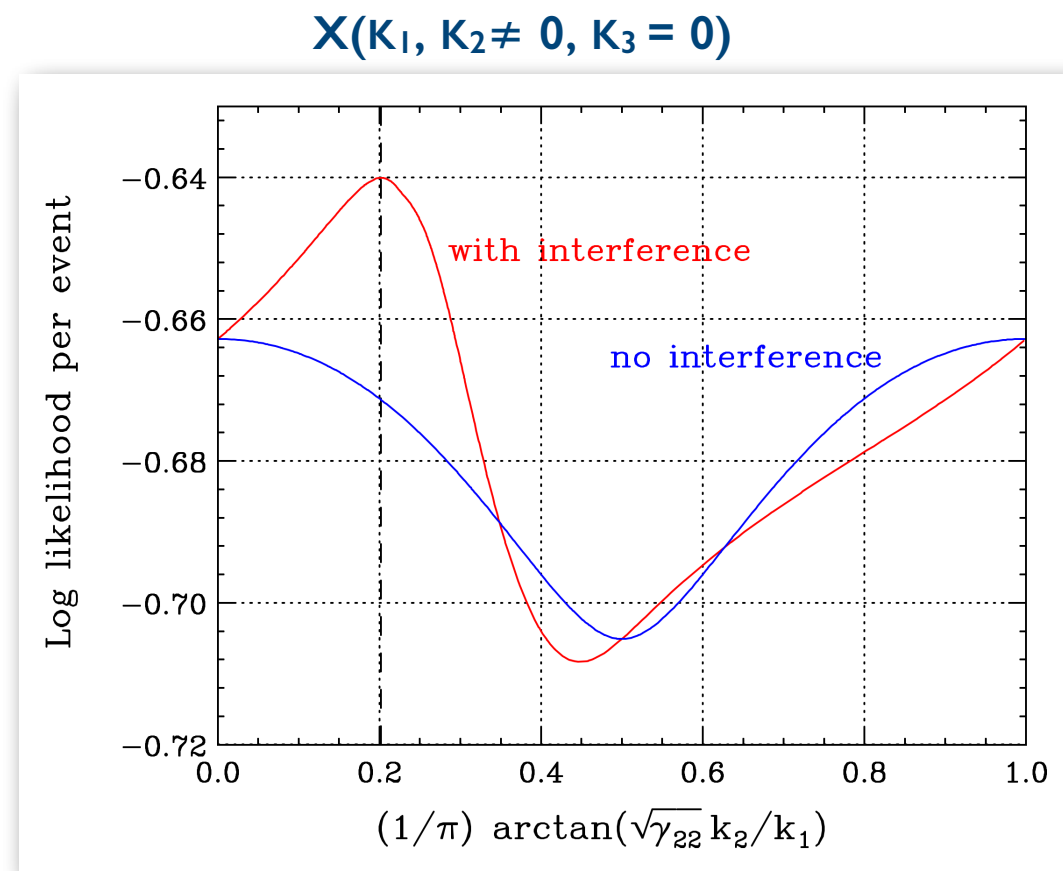
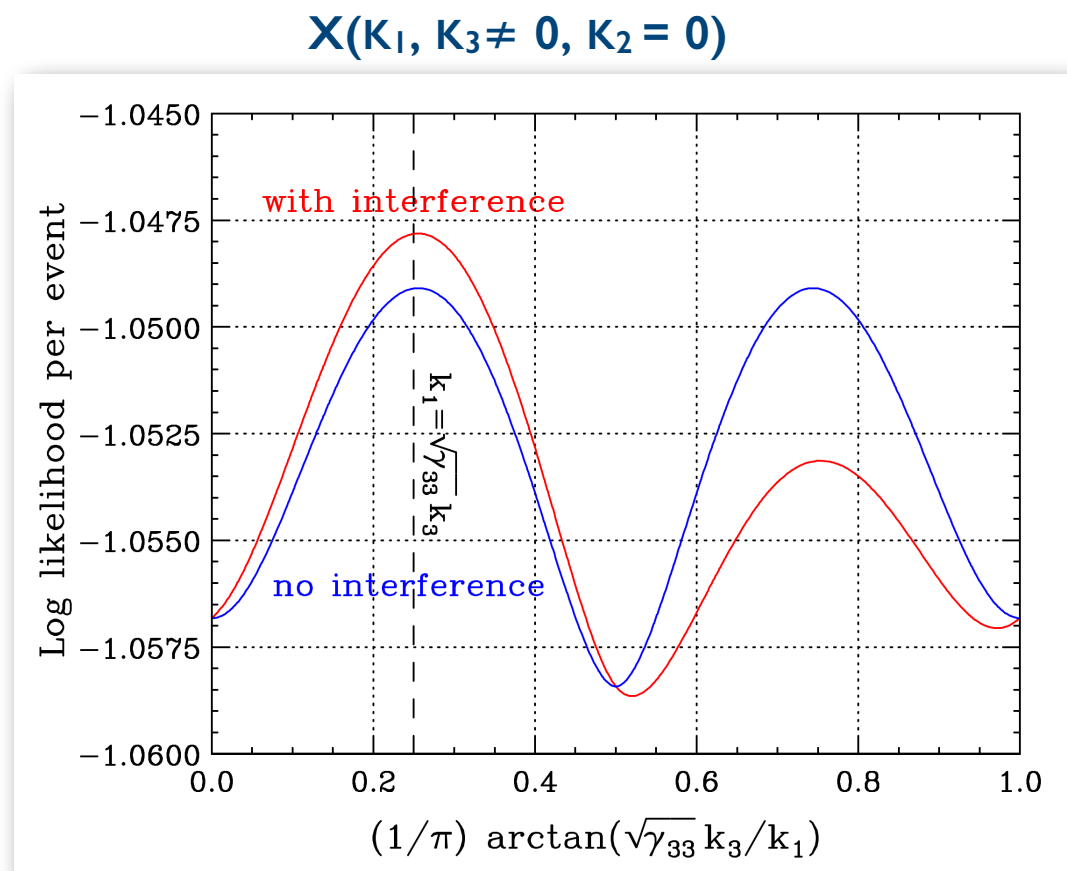
$$\gamma_{11} = 1, \quad \gamma_{22} = 0.090, \quad \gamma_{33} = 0.038, \quad \gamma_{12} = -0.250, \quad \gamma_{13} = \gamma_{23} = 0.$$

(for $2e2\mu$, before selection)

- Impact of interference on differential distributions (next slides)

Likelihood examples

- Per-event likelihoods built using the $|ME|^2$ for particular benchmark points
 - Demonstrate the potential to establish the presence or absence of interference, as well as to determine the relative sign of couplings



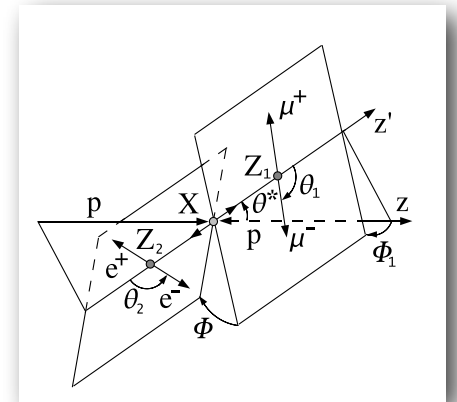
- The presence of interference breaks the $K_{2,3} \rightarrow -K_{2,3}$ symmetry and gives one sensitivity to the sign of the couplings.
 - case of interference between K_1 and the K_2 terms more straightforward to detect

Example: A prototypical analysis

- Basic idea:

- Construct observables/discriminants which take into account the interference effects and are “tuned” for each point of the parameter space
- Perform hypothesis tests for a discrete set of points in the parameter space.

Use kinematics of 4l system



- Kinematic discriminants based on Matrix Element Method:

- Discriminants $D(X;0^+)$** - computed for each point of the parameter space

$$D(X;0^+) = \frac{|\mathcal{M}(X)|^2}{|\mathcal{M}(0^+)|^2}$$



“ $0^+ \oplus 0^-$ ” example:

$$D(X;0^+) = \kappa_1^2 + \kappa_3^2 \frac{|\mathcal{M}(0^-)|^2}{|\mathcal{M}(0^+)|^2} + \kappa_1 \kappa_3 \frac{(\text{interference})}{|\mathcal{M}(0^+)|^2}$$

- $D(X;0^+)$ takes into account all aspects in which kinematics differ between the two hypotheses
- Discriminants such as $D(0^-;0^+)$ and $D(0_h^+;0^+)$ - inherently insensitive to interference effects:

$$D_{0^-} = D(0^-;0^+) = \frac{|\mathcal{M}(0^-)|^2}{|\mathcal{M}(0^+)|^2}, \quad D_{0_h^+} = D(0_h^+;0^+) = \frac{|\mathcal{M}(0_h^+)|^2}{|\mathcal{M}(0^+)|^2}$$

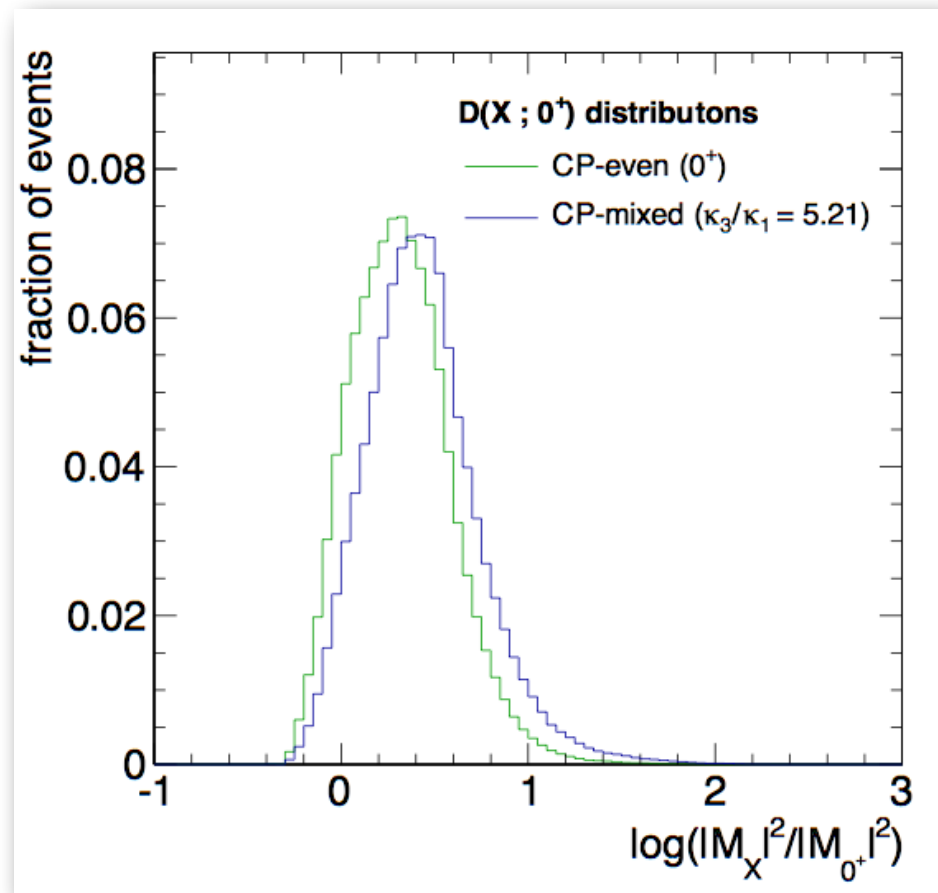
- Hypotheses separation extracted from test-statistics distributions of toy experiments

- Compute the exclusion regions for ratio of couplings (here κ_3/κ_1 and κ_2/κ_1 separately) for a given luminosity (up to 3000 fb^{-1} at 14TeV)
- Repeat same analysis using the $D(0^-;0^+)$ or $D(0_h^+;0^+)$ discriminants
 - to quantify the gain in sensitivity due to the interference effects

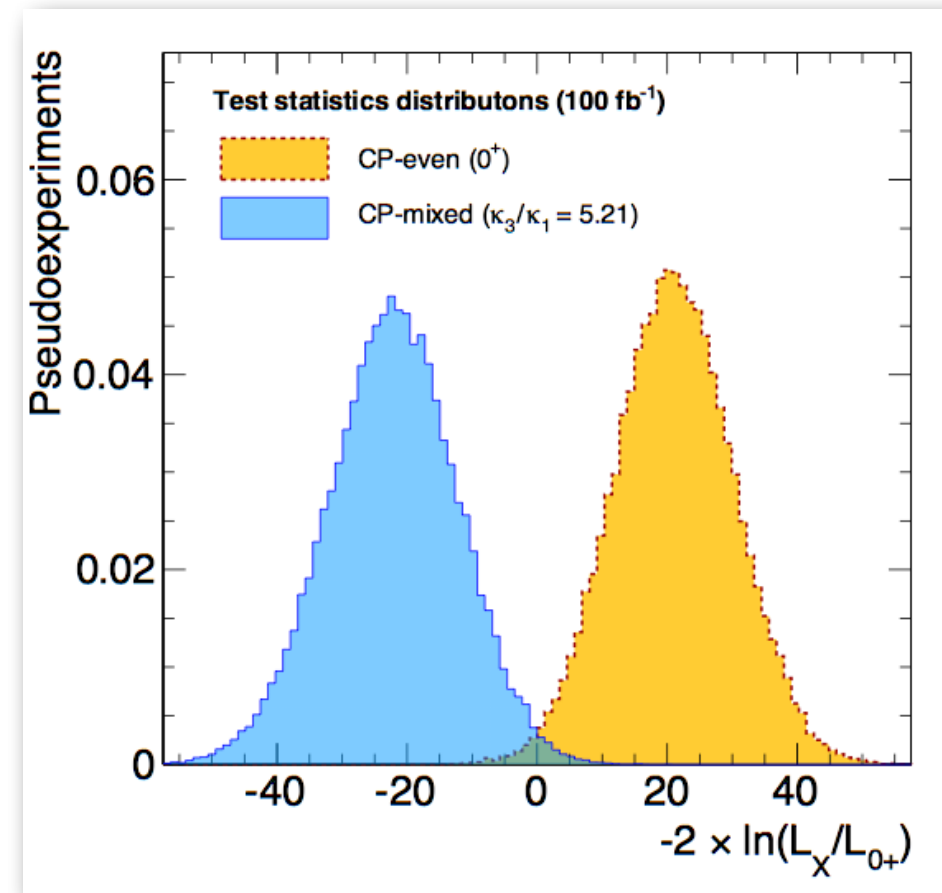
Example: Templates and test-statistics

- ID templates of $D(X;0^+)$ built for the two signal hypotheses and background events
 - all three final states together, 8 TeV samples only.

$D(X;0^+)$ for $K_3/K_1 = 5.21$



TS distributions, 100 fb^{-1} , $K_3/K_1 = 5.21$



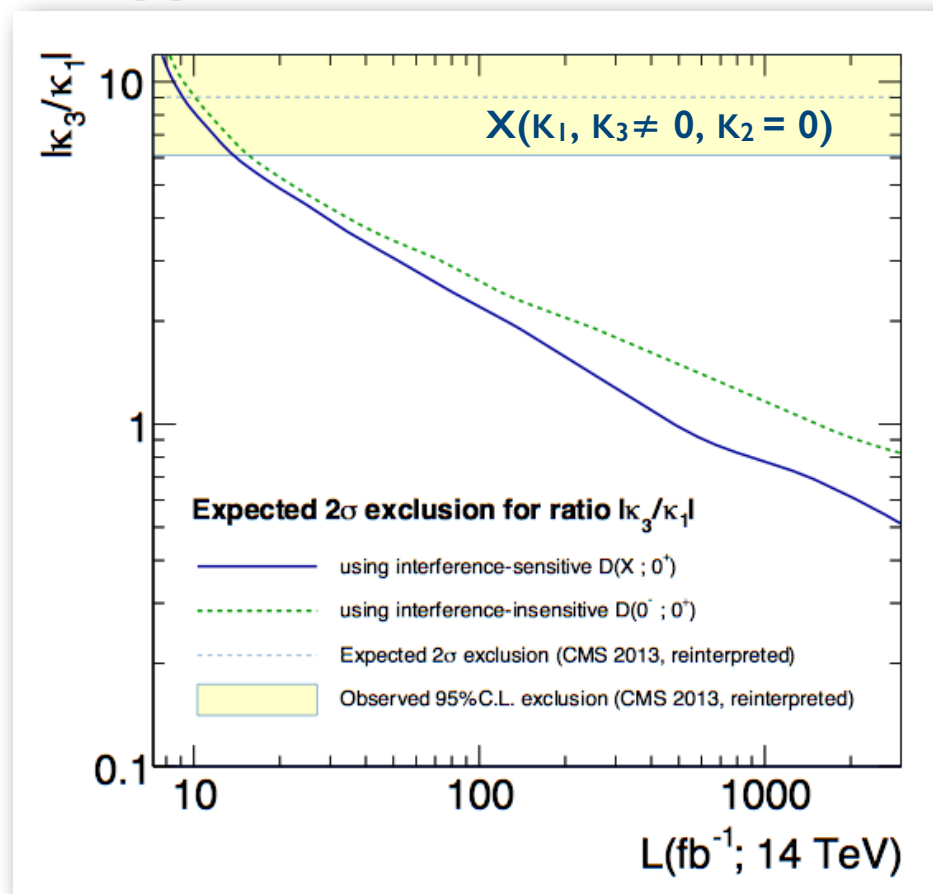
- Test-statistics (TS) for 0^+ and X hypotheses from tossing the toy experiments (50k),
 - used to compute the expected separation between the hypotheses for a given int. luminosity

Example: Results for the κ_3/κ_1 and κ_2/κ_1 ratios

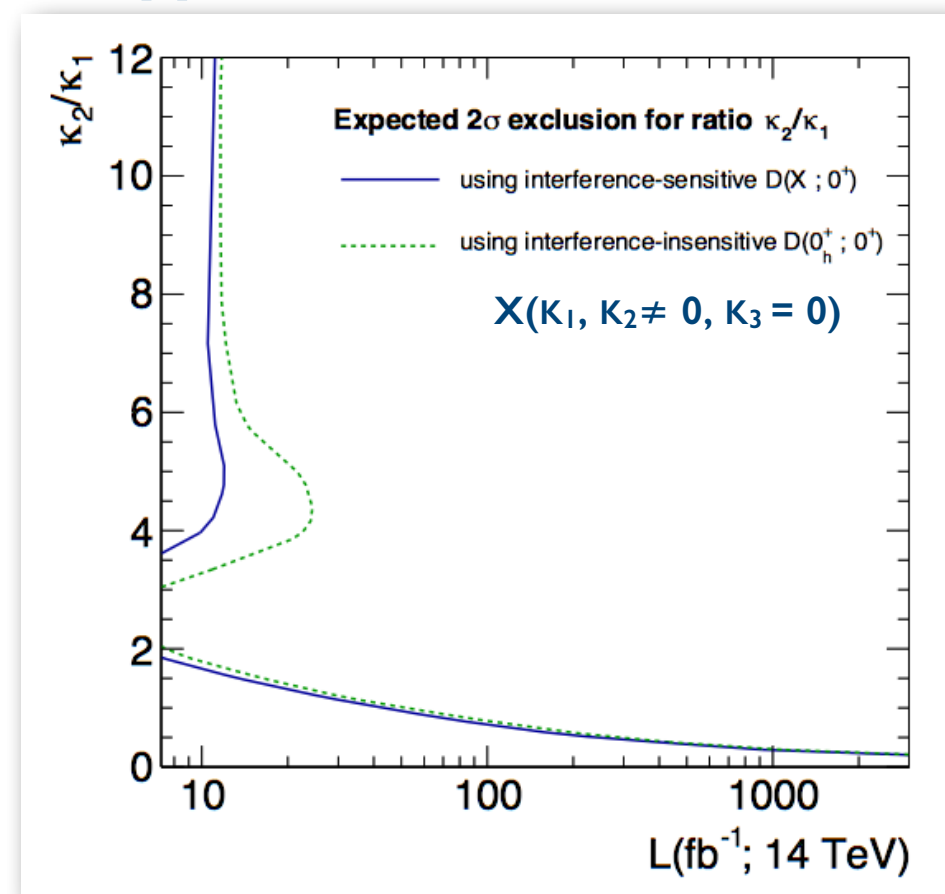
- Results in terms of integrated luminosity at 14TeV required to achieve the expected 2σ limit
 - results for interference-sensitive $D(X;0^+)$ and interference-blind $D(0^-;0^+) / D(0_h^+;0^+)$ discriminants

[arxiv: 1310.1397](https://arxiv.org/abs/1310.1397) [hep-ph]

upper limits on κ_3/κ_1 ratio



upper limits on κ_2/κ_1 ratio

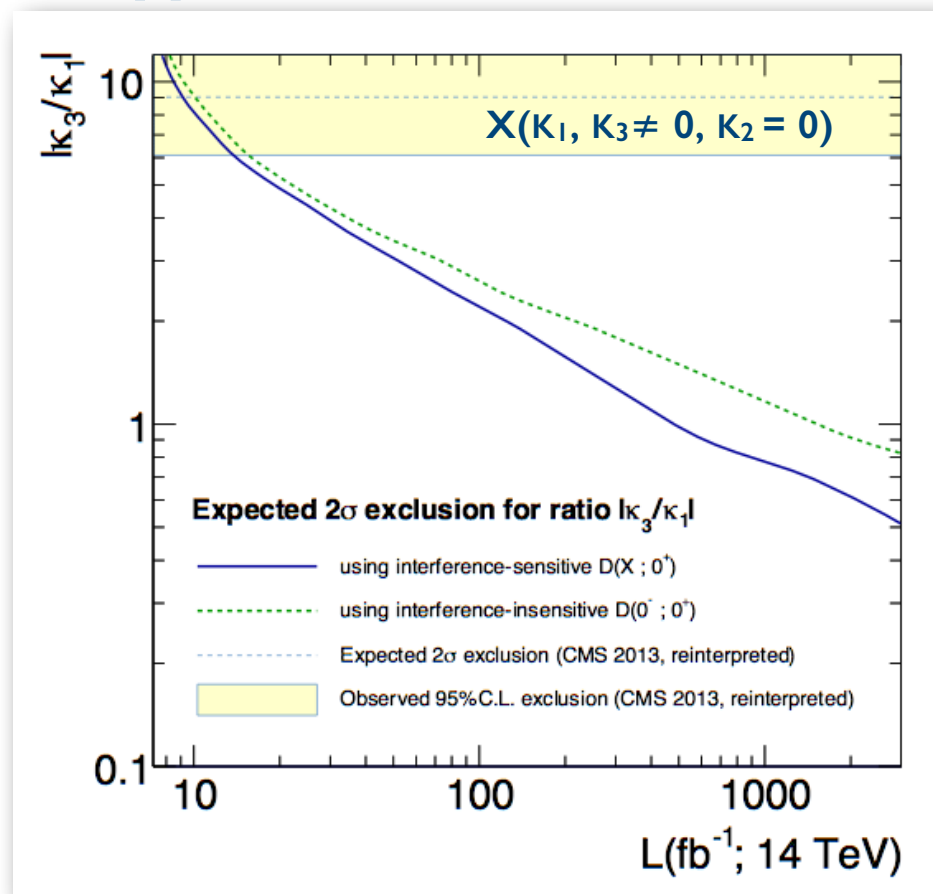


- Reduce the int. luminosity required to exclude the ratio $\kappa_3/\kappa_1 \sim 1$ by up to a factor of 4
 - Modest difference $\mathcal{O}(10\%)$ in sensitivities to κ_3/κ_1 at int. luminosity of $\sim 10 \text{ fb}^{-1}$ ($\sim 25 \text{ fb}^{-1}$ @ 8TeV).
- Substantial effects of interference in range $\kappa_2/\kappa_1 \approx 2 - 4$
 - allow us to exclude a range of $2 < \kappa_2/\kappa_1 < 4$ with the already existing data!

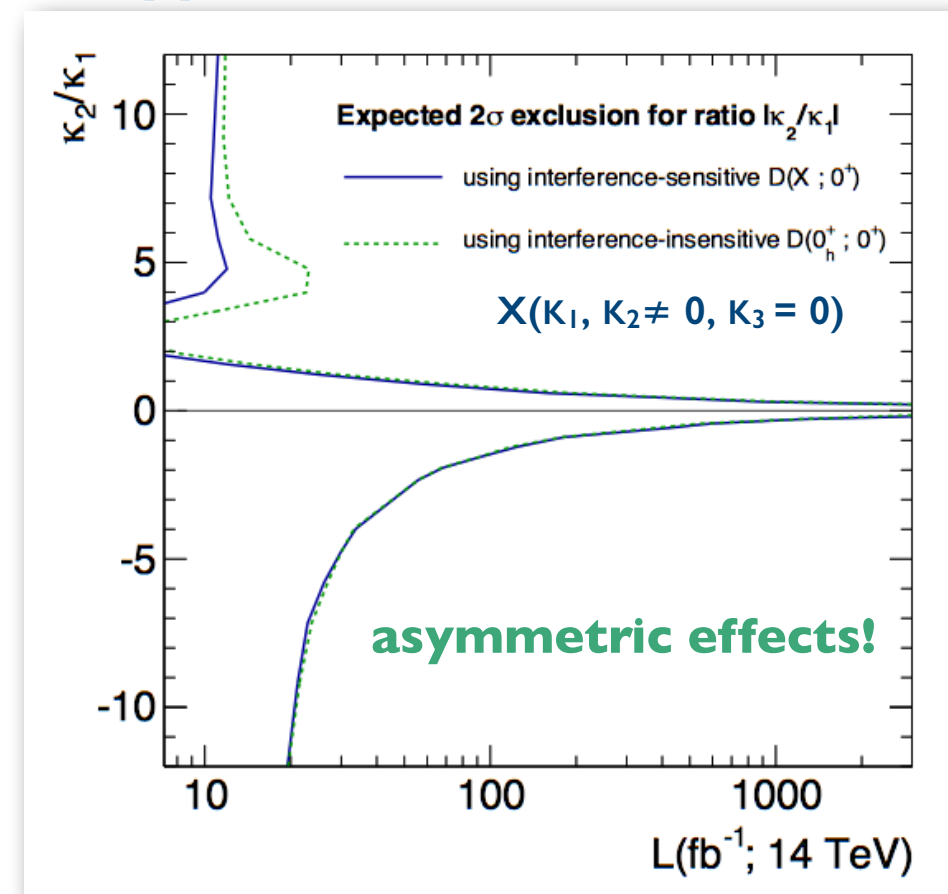
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upper limits on κ_3/κ_1 ratio



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Exploitation of interference effects in particle physics

To remember...

1. Interference effects can affect both inclusive and differential distributions of signal and background processes,
2. Exploiting interference effects can improve sensitivity to (certain) rare processes

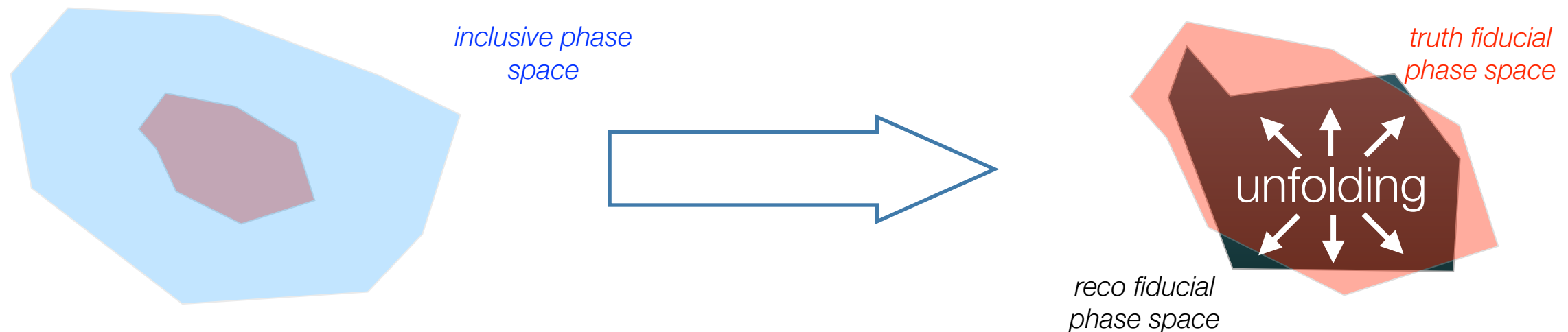


Measurement of fiducial cross sections

Some basics

MOTIVATION:

- Fiducial cross sections offer a **possibility to describe data in model independent way**
 - Maximise the applicability of LHC data to explore the **QCD effects** in the SM, and capture **BSM effects** in the Higgs boson physics.



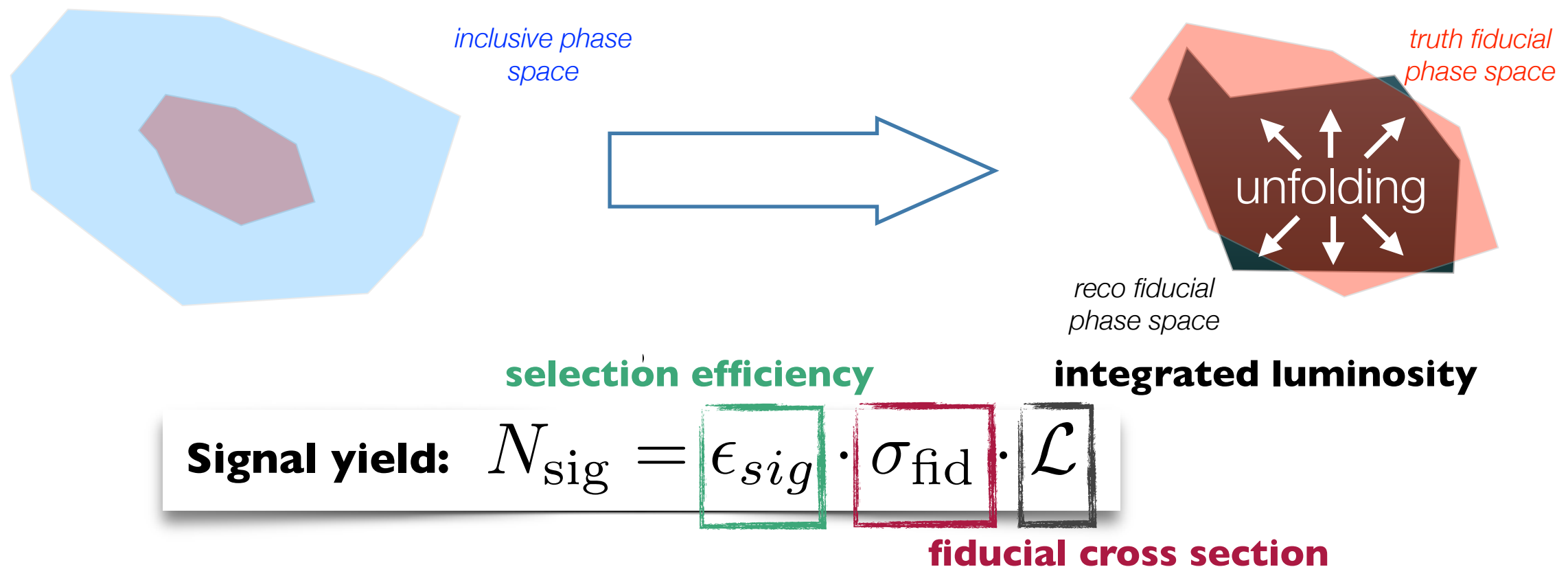
A FEW IMPORTANT ASPECTS:

- **Model independence** of the measurements
 - Factorise theory uncertainties from experimental ones (no extrapolation)
 - Need for the measurements to **survive the passage of time**

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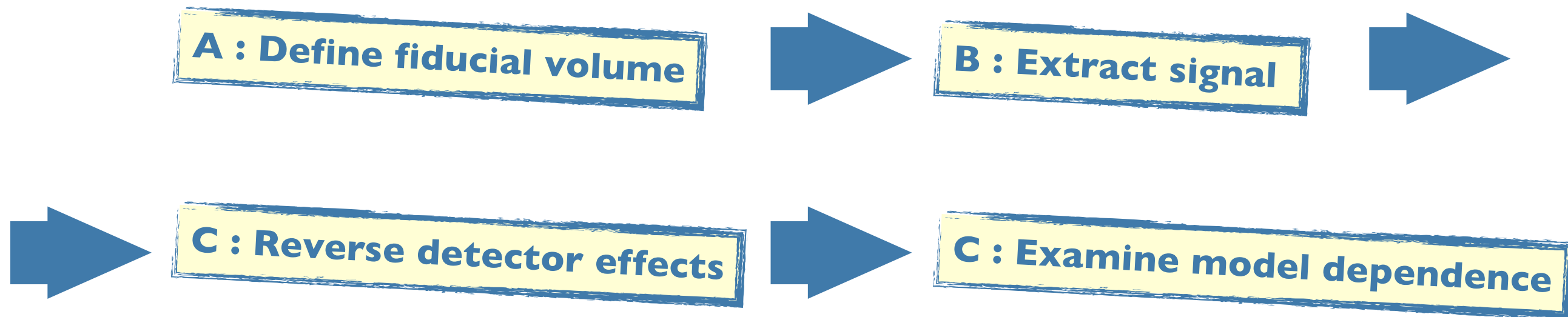
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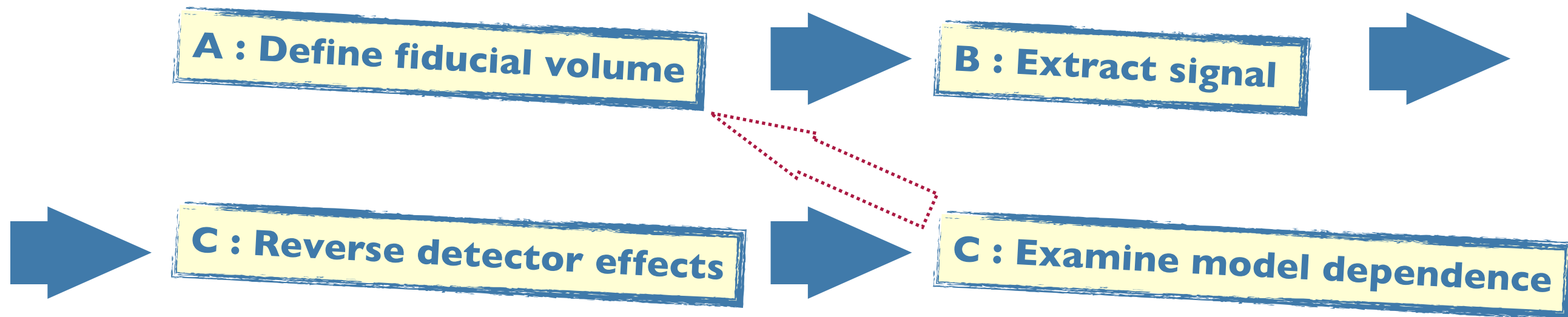
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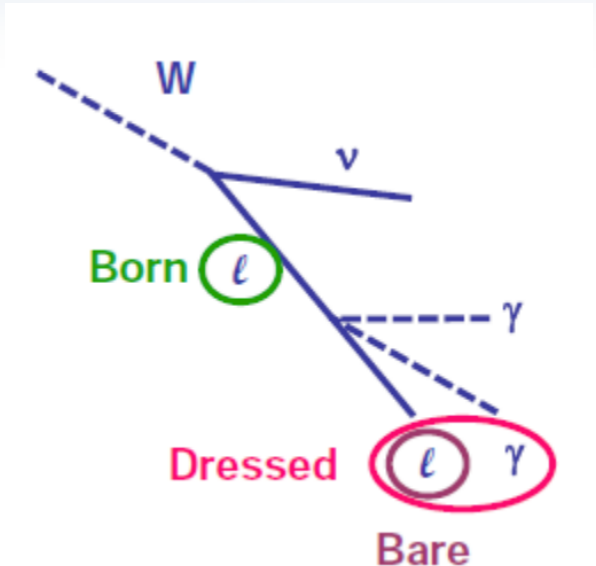
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Important ingredients

FIDUCIAL DEFINITIONS:

- Definition of the fiducial-level objects (leptons, photons, jets)
- **Isolation requirement** plays an important role
- Out-of-fiducial signal contributions need special care
- **NOTE:** Different kinematical cuts in ATLAS/CMS (optimised to exploit detector potential).



M(H) HYPOTHESIS:

- Use **best-fit value measured by experiment(s)** for comparisons with theory (either treat $m(H)$ as a free parameter and fit for it, or fix $m(H)$ to best-fit value).

MODEL DEPENDENCE:

- Build response matrix and repeat the unfolding procedure once per model
 - **SM studies:** vary production mode composition (e.g. within experimental constraints)
 - **BSM studies:** consider a predefined set of exotic models (with/without exp. constraints)

Unfolding

GOAL:

- Undo the effects of smearing due to detector resolution & efficiency
- Complex problem (and not very well defined), important for theory comparisons

POSSIBLE APPROACHES:

- First subtract background, then unfold with **inverted detector response** $[\epsilon_{ij}]$

Step I:
$$N_{\text{sig}}^{\text{f},i}(m_{4\ell}) = N_{\text{obs}}^{\text{f},i}(m_{4\ell}) - N_{\text{bkg}}^{\text{f},i}(m_{4\ell})$$

Step II:
$$\sigma_{\text{fid}}^j = [\epsilon_{i,j}]^{-1} \cdot \left[\frac{1}{\mathcal{L}} N_{\text{sig}}^i \right]$$

► Add systematics to cover for possible biases, cross-check the claimed coverage

- Fold detector response matrix $[\epsilon_{ij}]$ in the likelihood and **perform background subtraction and signal unfolding simultaneously.**

$$\begin{aligned} N_{\text{obs}}^{\text{f},i}(m_{4\ell}) &= N_{\text{sig}}^{\text{f},i}(m_{4\ell}) + N_{\text{bkg}}^{\text{f},i}(m_{4\ell}) \\ &= \epsilon_{i,j}^{\text{f}} \sigma_{\text{fid}}^{\text{f},j} \cdot \mathcal{L} \cdot \text{pdf}(m_{4\ell} | \text{H} \rightarrow 4\ell) + N_{\text{bkg}}^{\text{f},i} \cdot \text{pdf}(m_{4\ell} | \text{bkg}) \end{aligned}$$

► Full correlation of relevant parameters

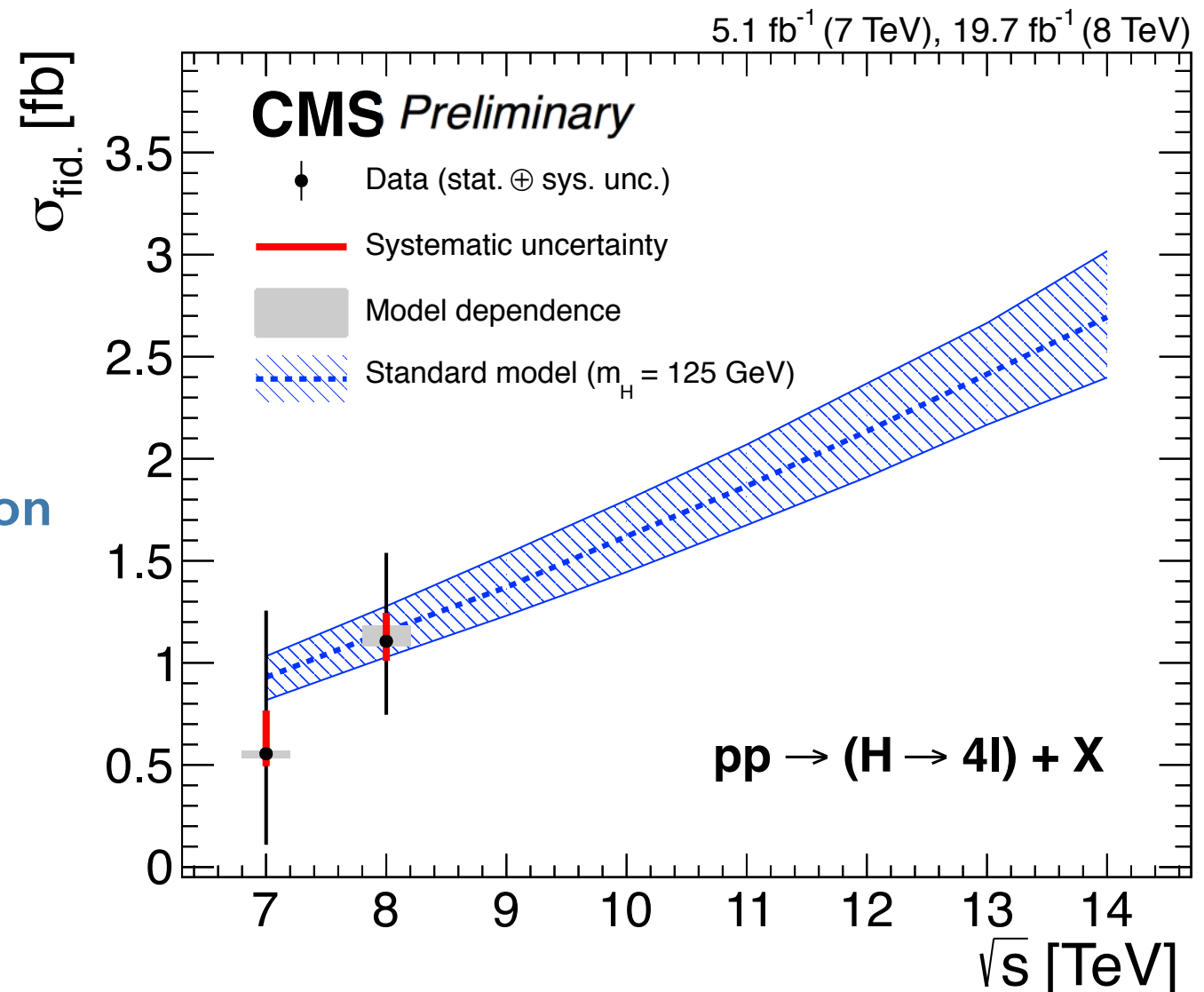
Inclusive cross sections: $H \rightarrow ZZ \rightarrow 4l$

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CMS-PAS-HIG-14-028

- Inclusive cross sections at 7 and 8 TeV, measured in $H \rightarrow 4l$ channel**

- Un-binned maximum likelihood fit to m_{4l}
- Model dependence estimated from range of SM and exotic Higgs models:
 - <1% using experimental constraints
 - ~7% without experimental constraints
- CMS: $Z \rightarrow 4l$ resonance used as validation**
- 7 TeV measurement statistically limited.



ATLAS @ 8 TeV:

$$\sigma_{\text{fid}} = 2.11 \pm_{-0.47}^{+0.53} (\text{stat.}) \pm 0.08 (\text{syst.}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 1.30 \pm 0.13 \text{ fb}$$

CMS @ 8 TeV:

$$\sigma_{\text{fid}} = 1.11 \pm_{-0.35}^{+0.41} (\text{stat.}) \pm_{-0.10}^{+0.14} (\text{syst.}) \pm_{-0.02}^{+0.08} (\text{model}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 1.15 \pm_{-0.13}^{+0.12} (\text{stat.}) \text{ fb}$$

CMS @ 7 TeV:

$$\sigma_{\text{fid}} = 0.56 \pm_{-0.44}^{+0.67} (\text{stat.}) \pm_{-0.06}^{+0.21} (\text{syst.}) \pm_{-0.02}^{+0.02} (\text{model}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 0.93 \pm_{-0.11}^{+0.10} (\text{stat.}) \text{ fb}$$

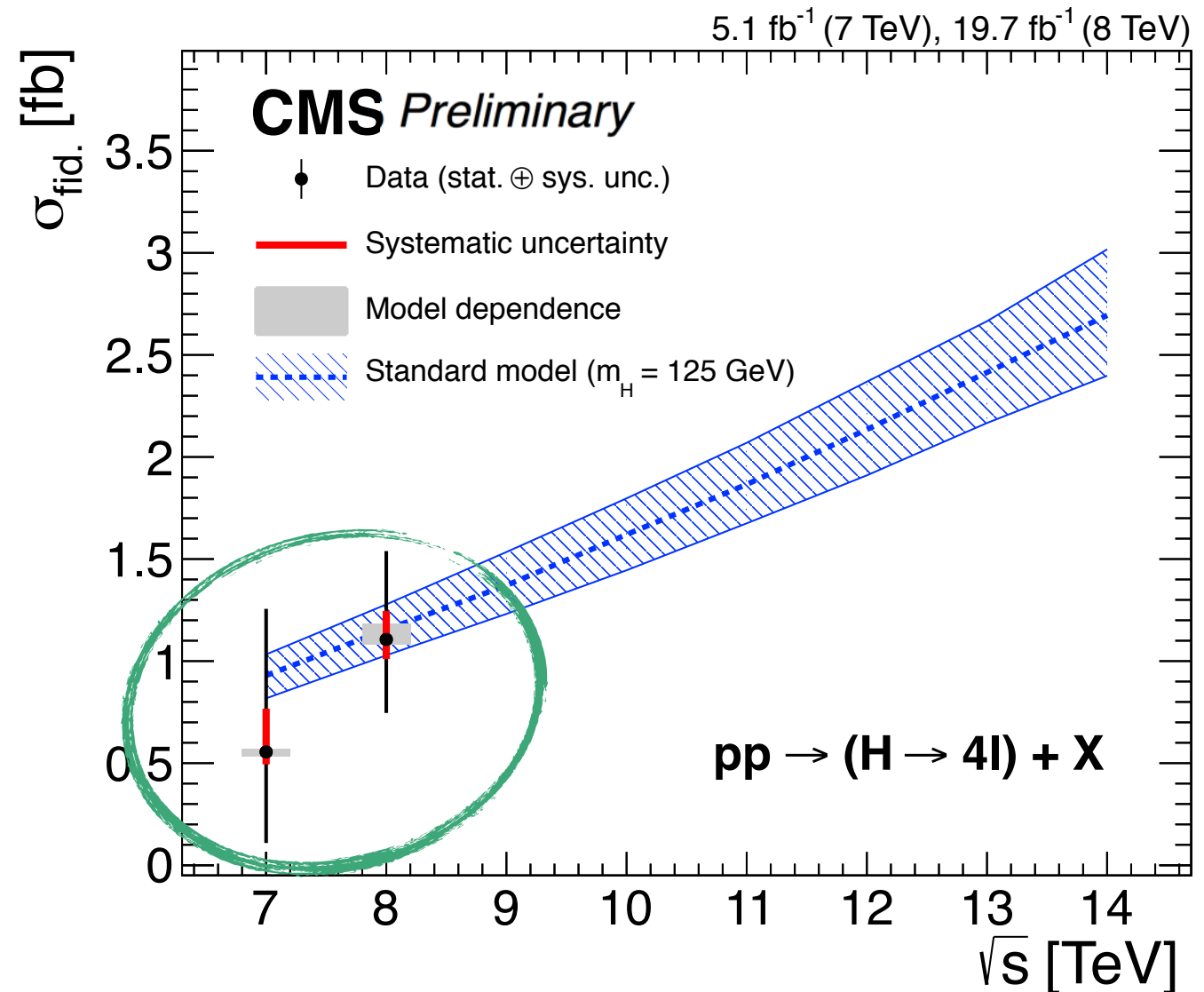
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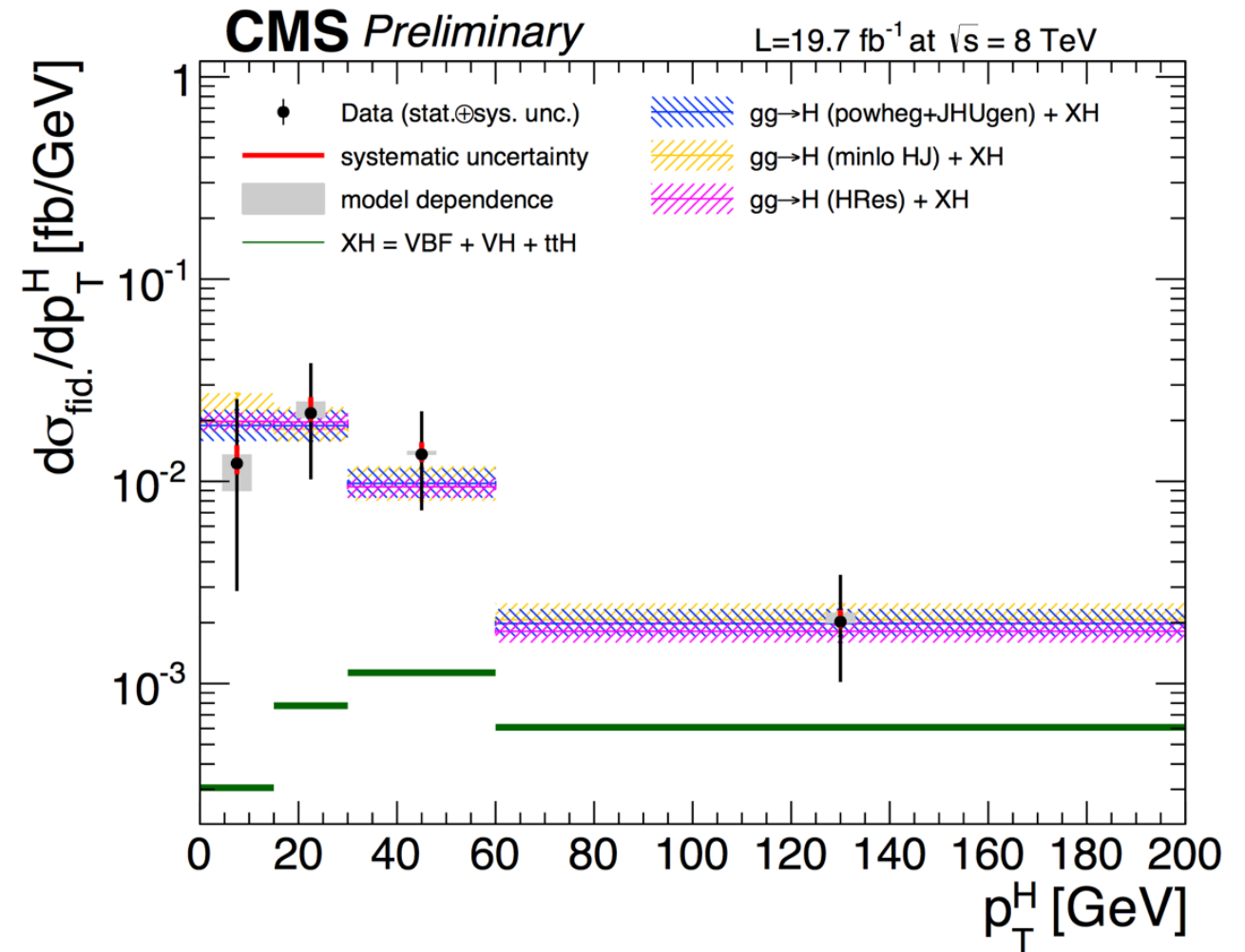
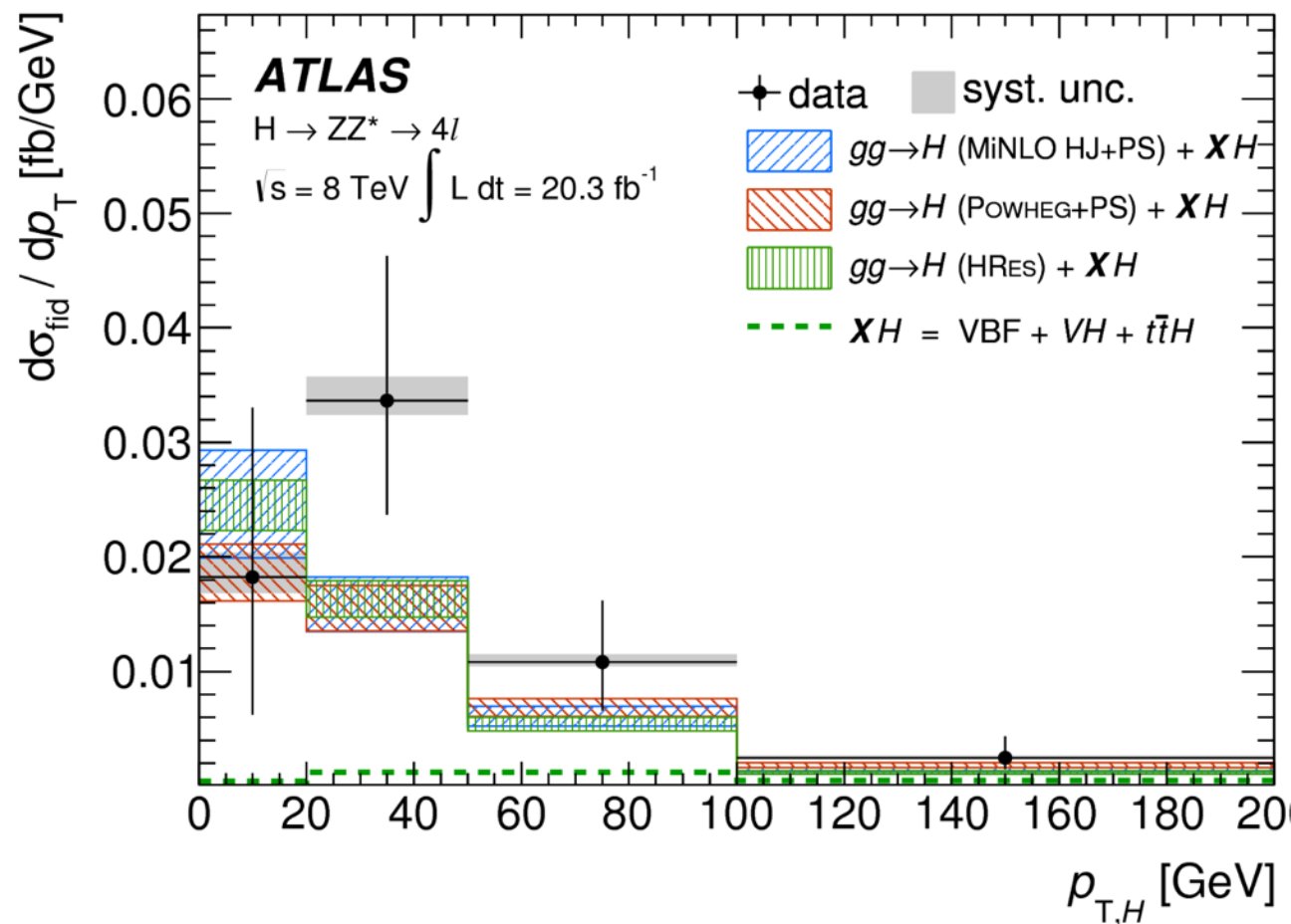
► **Compatible with theoretical estimates (slightly higher rates by ATLAS)**

Differential XS: $H \rightarrow ZZ \rightarrow 4l$

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CMS-PAS-HIG-14-028

- Measured as a function of Higgs candidate kinematic properties:
 - $p_T(\mathbf{H})$: sensitive to production mode, new physics in $gg \rightarrow H$ loop
 - $|\mathbf{y}(\mathbf{H})|$: sensitive to production mode, parton distribution functions (see backup).

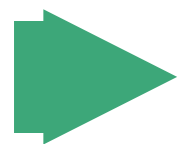
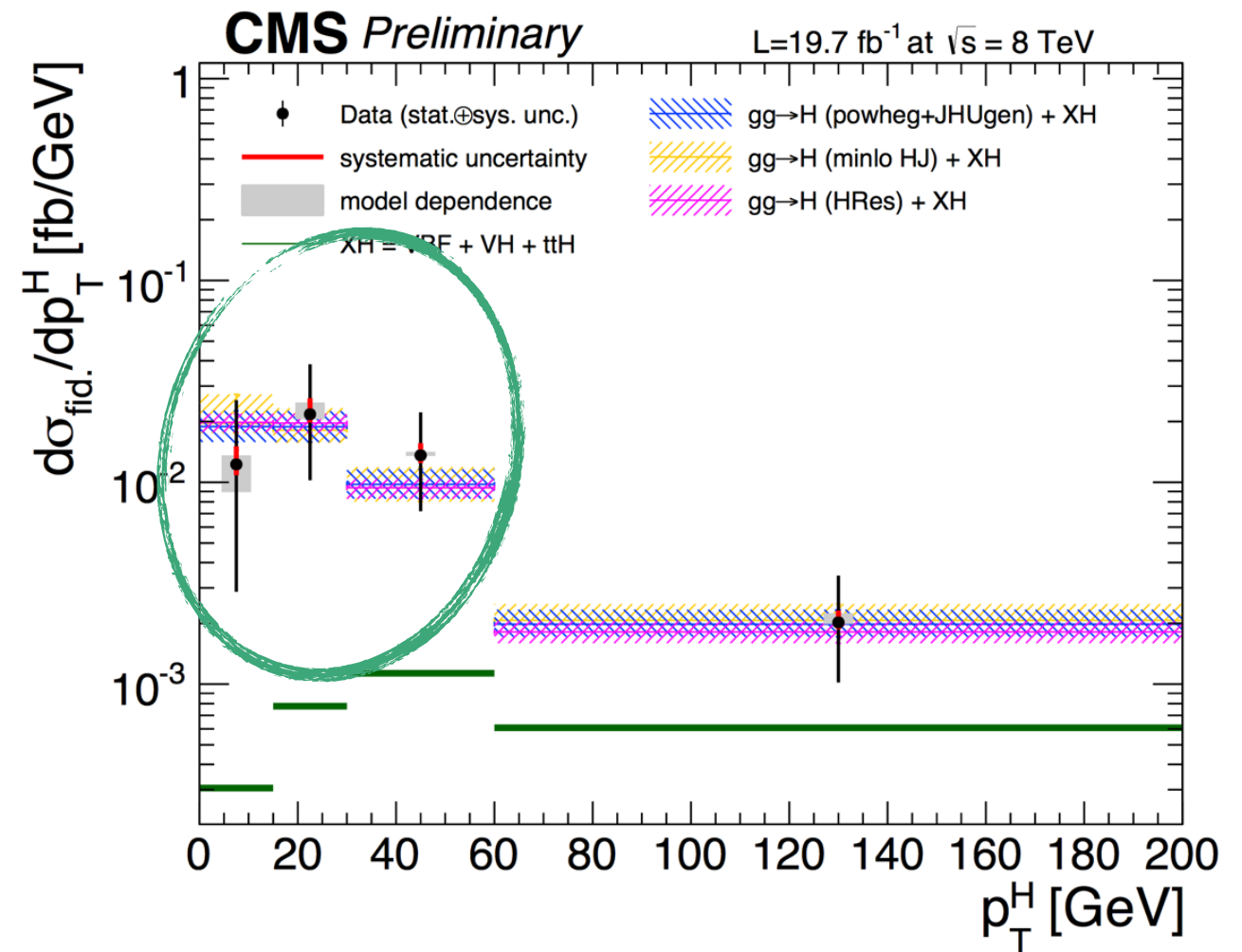
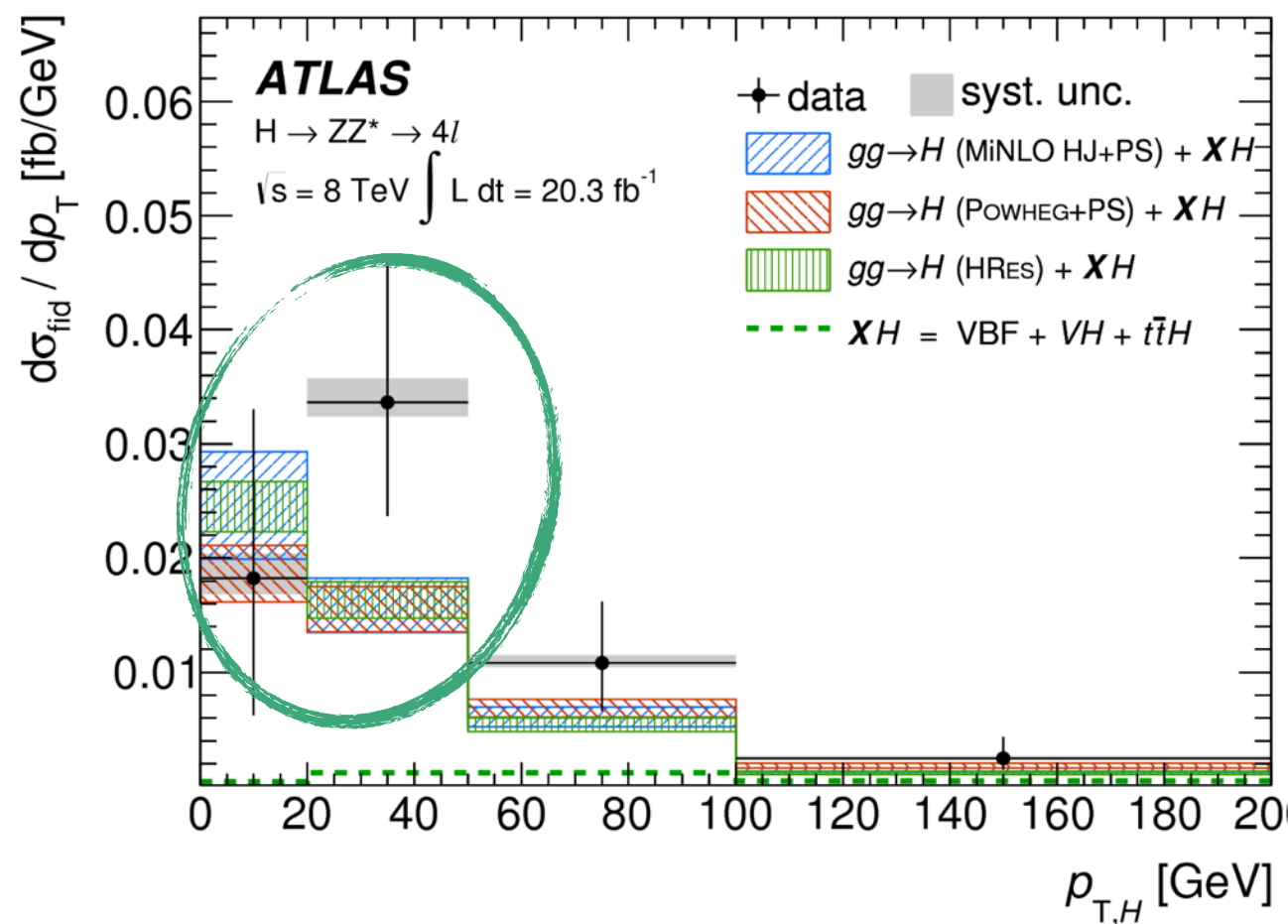


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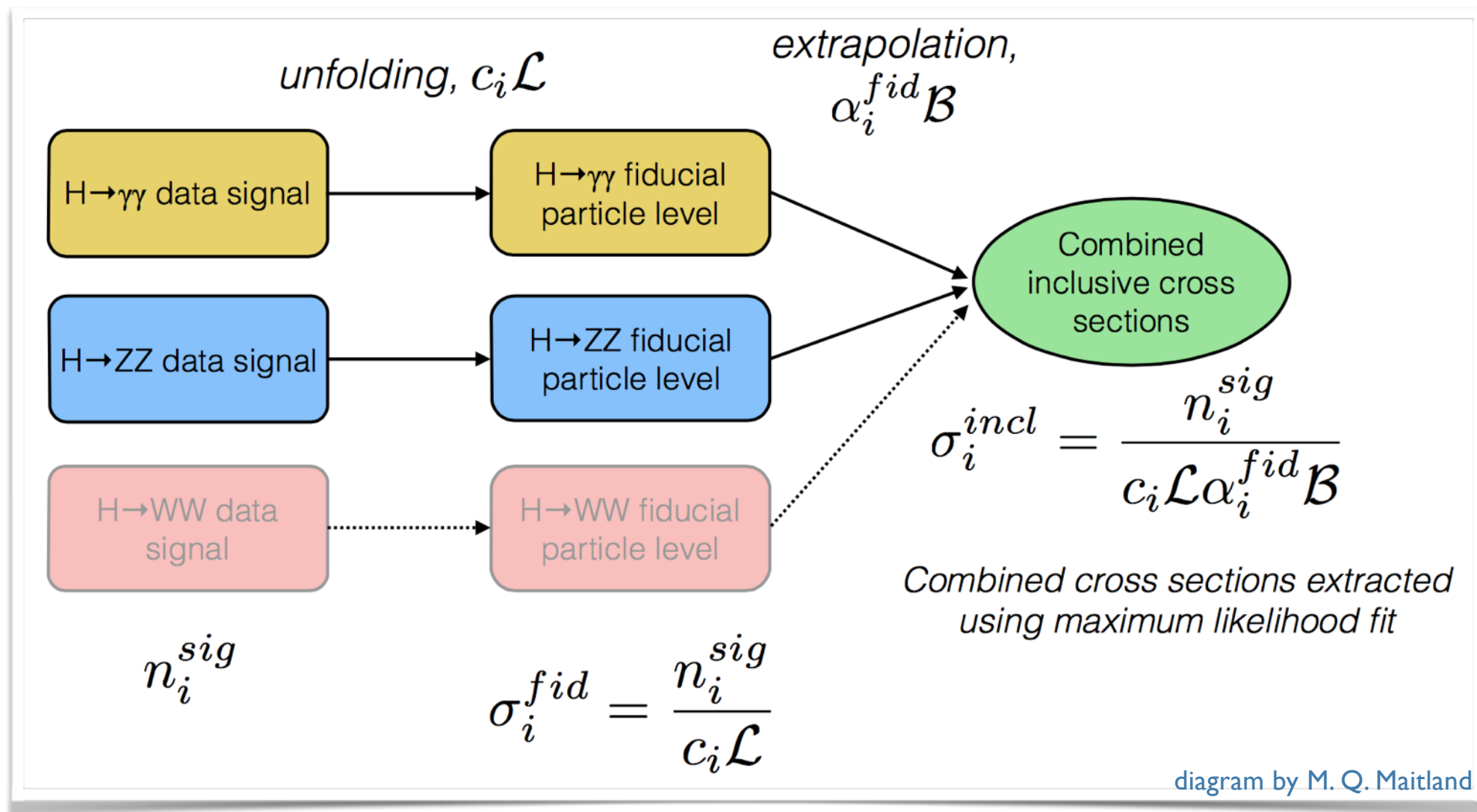


Somewhat harder p_T spectrum observed by ATLAS and CMS w.r.t theoretical estimates (p -value ~ 10 -15%)

Combination of measurements

Combination between decay channels ($H \rightarrow \gamma\gamma$, $H \rightarrow ZZ$, $H \rightarrow WW$, etc.):

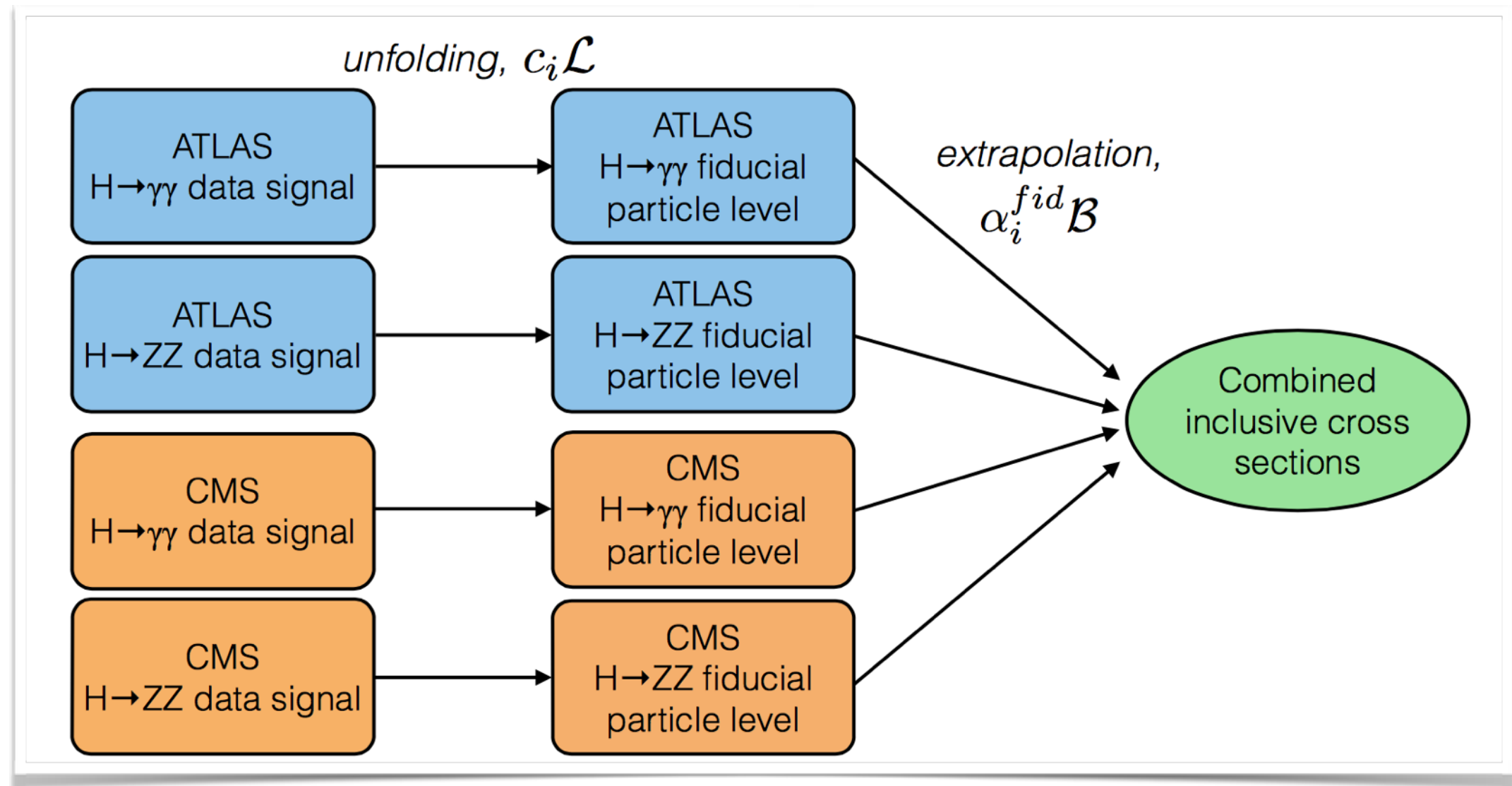
- Perform the fit to **inclusive XS in the full phase space** (inherent assumption of the same source of decays)
- **Statistical precision at the expense of model dependence due to extrapolation** (quote a total XS, check the compatibility between the measurements).



Combination of measurements

Combination between experiments:

- Potential to combine inclusive and differential cross sections (need harmonisation in fiducial objects, bin edges, unfolding, etc.)
 - Choose **common fiducial or inclusive phase space?**
- } still to be discussed





Measurement of fiducial cross sections

To remember...

1. Fiducial cross sections offer a possibility to describe data in a model independent way
2. Factorise theory uncertainties from experimental ones, and represent legacy results
3. Several approaches to the unfolding of detector effects. Requires a special care.

PART III: Selected topics on the physics measurements





THE END
Thank you for your attention!



ADDITIONAL MATERIAL

Reducible Background - contributions at a glance

- Definitions:

- Tight-to-loose ratio measured in **Z + 1 loose lepton events** \Rightarrow f
- observed events in the “**2 passed + 2 failed**” (**2P+2F**) region \Rightarrow N_{2P2F}
- observed events in the “**3 passed + 3 failed**” (**3P+1F**) region \Rightarrow N_{3P1F}
- ZZ** contribution in the “**1 passed + 1 failed**” (**3P+1F**) region \Rightarrow $N_{3P1F}^{(ZZ)}$
- Contributions from **ZZ** and **3P+1F** processes in **2P+2F** region are negligible

- The expected contributions from “**2 prompt + 2 fake**” processes (symbolic)

- in the **3P+1F** region \Rightarrow

$$N_{3P1F}^{(2P2F)} = \sum \left[\frac{f_i}{(1-f_i)} + \frac{f_j}{(1-f_j)} \right] N_{2P2F}$$

- in the **Signal region (SR)** \Rightarrow

$$N_{SR}^{(2P2F)} = \sum \frac{f_i}{(1-f_i)} \frac{f_j}{(1-f_j)} N_{2P2F}$$

(symbolic)

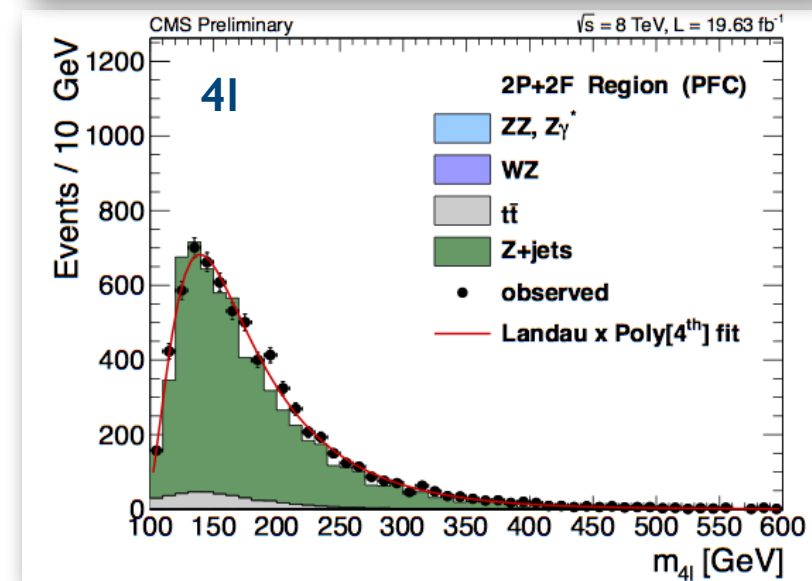
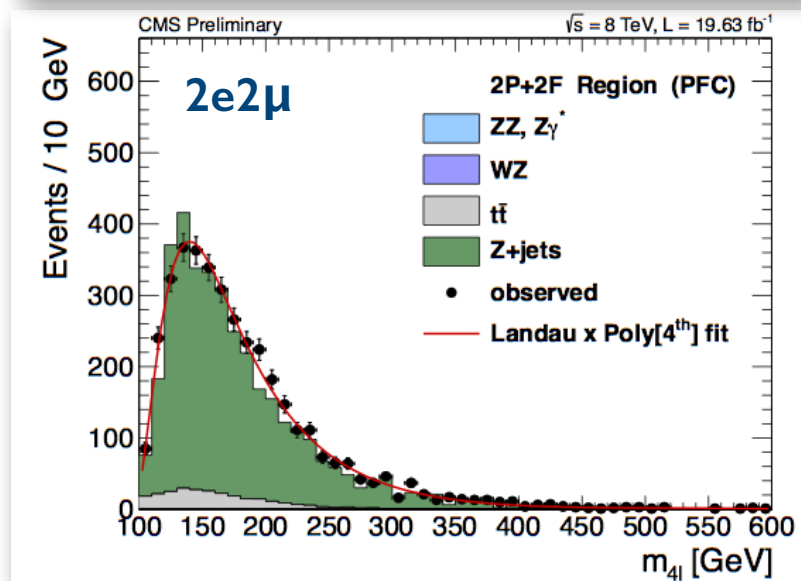
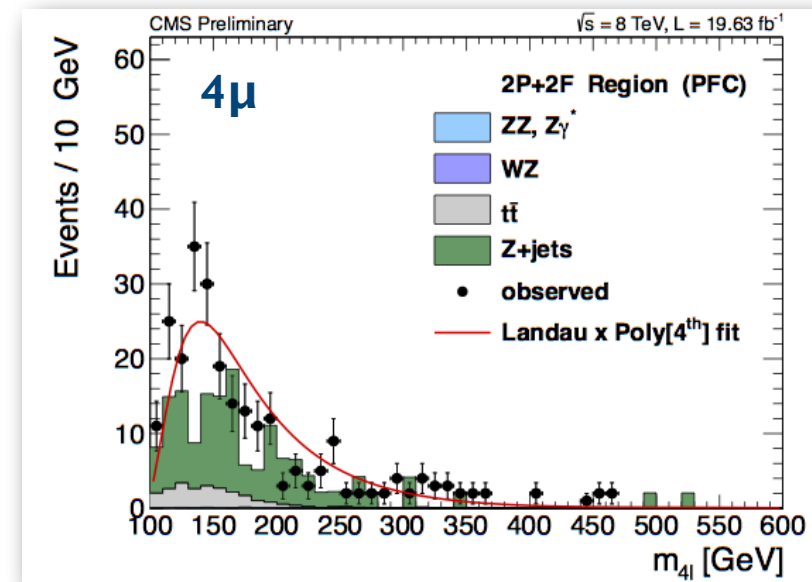
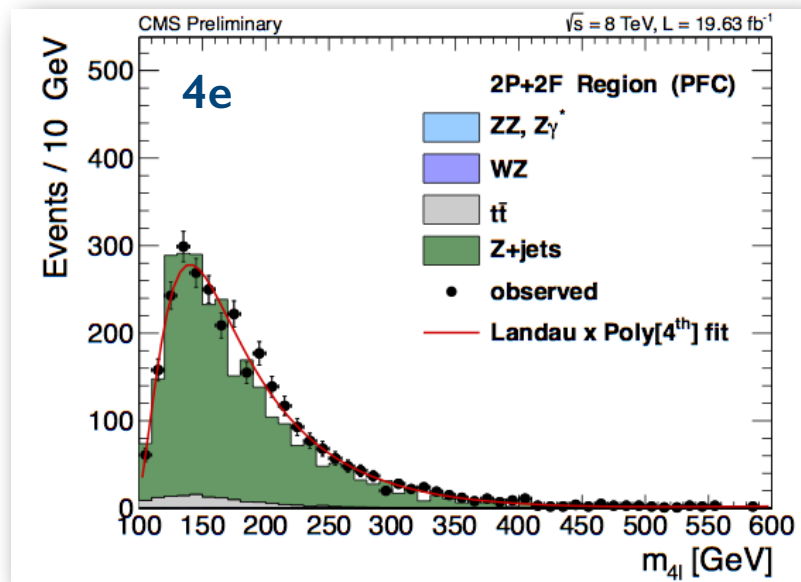
- Total expected contributions

- in the **Signal region** (symbolic)

$$N_{SR}^{bkg} = \sum \frac{f_i}{(1-f_i)} (N_{3P1F} - N_{3P1F}^{(2P2F)} - N_{3P1F}^{(ZZ)}) + \sum \frac{f_i}{(1-f_i)} \frac{f_j}{(1-f_j)} N_{2P2F}$$

Reducible Background - Control region 2P+2F

- Measure contribution from processes with 2 prompt + 2 fake leptons using
 - Control region with “2 passed + 2 failed” (2P+2F) leptons

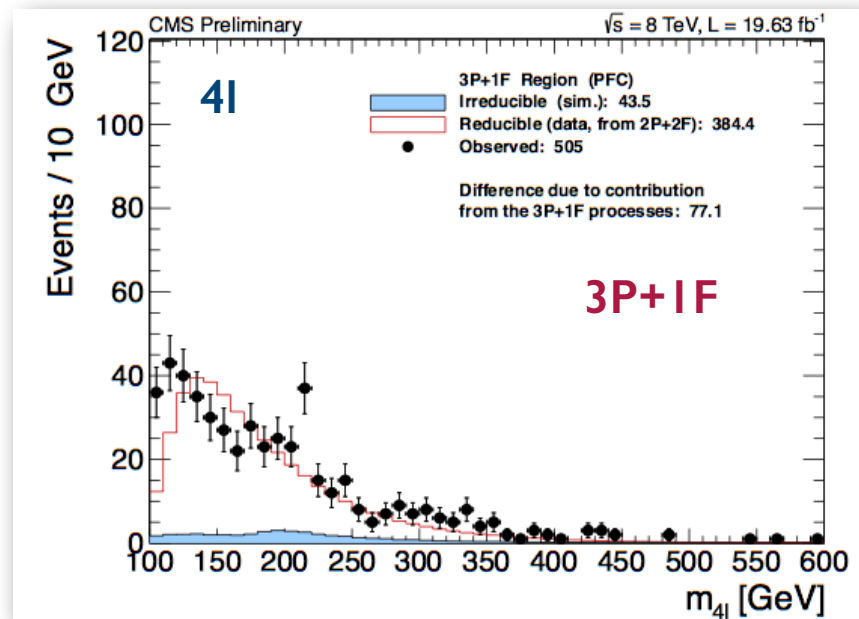


- 4e and 2e2μ channel - data are well described by MC
 - 4μ channel - data are not properly described by MC (as before)
- Reminder: MC predictions are not used in the analysis

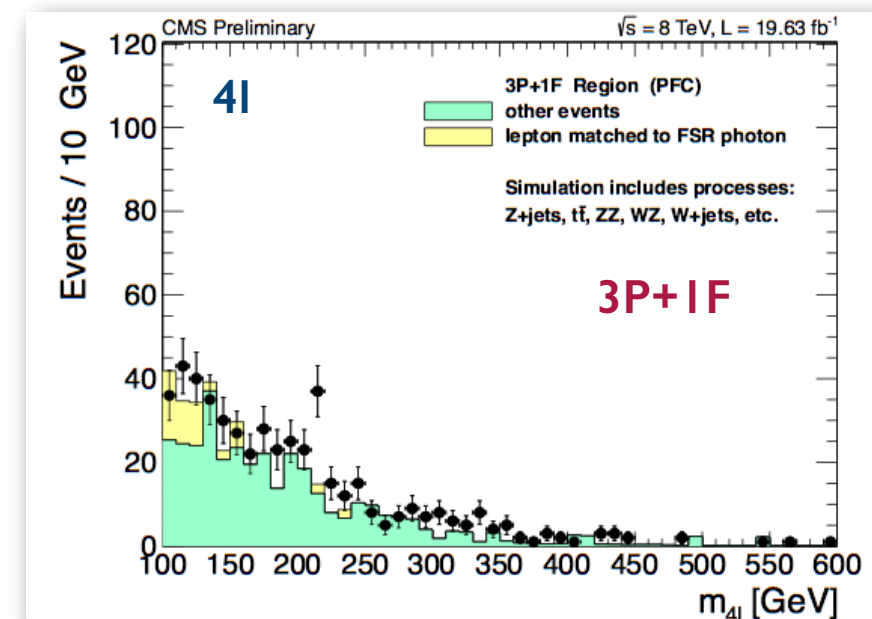
Reducible Background - Control region 3P+1F

- Measure contribution from processes with 3 prompt + 1 fake leptons ($Z\gamma, WZ, \text{etc.}$) using
 - Control region with “3 passed + 1 failed” (3P+1F) leptons
 - Estimated spill-over from 2P+2F region (**from data**) and Signal region (from simulation)

contributions from
2P+2F & Signal region



simulation
(MC-truth matching)



- 3P+1F events predominantly in the low m_{4l} region - $Z\gamma$ events (asymmetric conversions)
- Total expected contributions in the Signal region

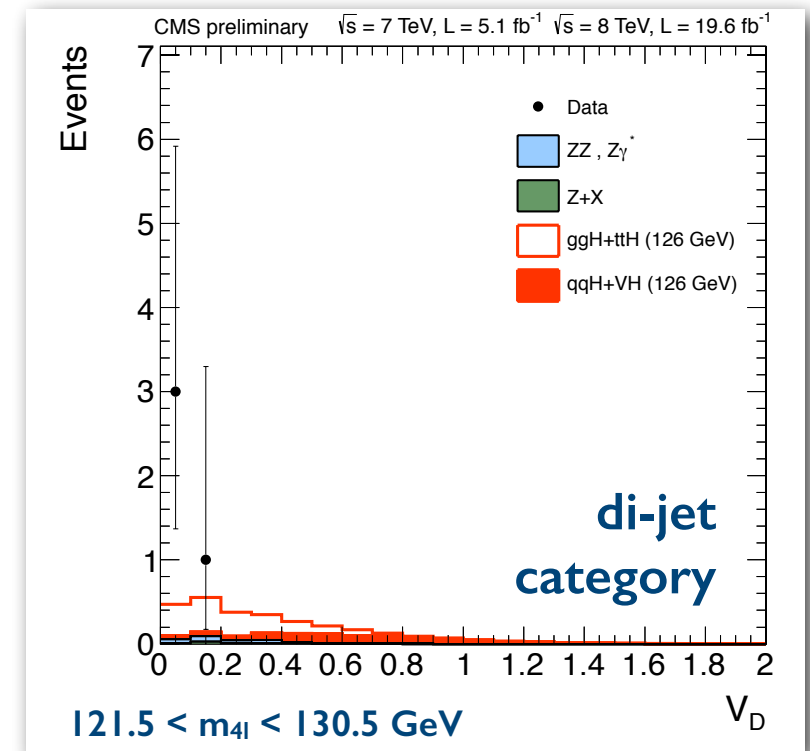
(symbolic)

$$N_{SR}^{bkg} = \sum \frac{f_i}{(1-f_i)} (N_{3P1F} - N_{3P1F}^{(2P2F)} - N_{3P1F}^{(ZZ)}) + \sum \frac{f_i}{(1-f_i)} \frac{f_j}{(1-f_j)} N_{2P2F}$$

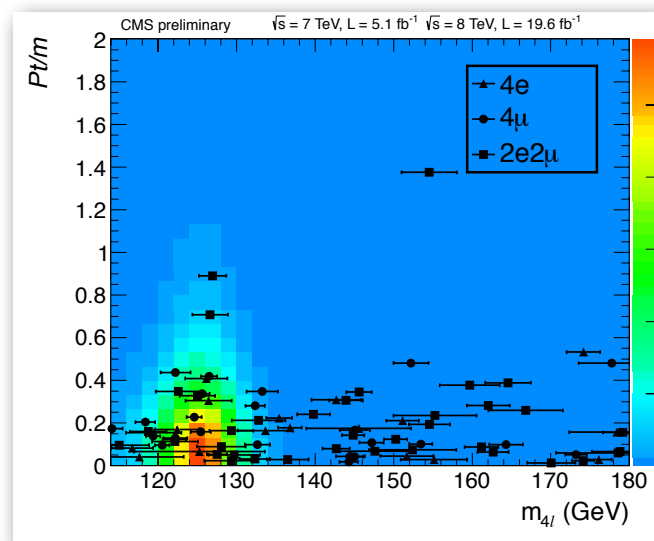
Event categories in the analysis

- The event sample is split into two categories:
 - Category I: Events with $N_{\text{JETS}} < 2$. (5% VBF)
 - Category II: Events with $N_{\text{JETS}} \geq 2$. (20% VBF)
- Discriminate production mechanisms (fermion- vs. vector-boson-induced):
 - Cat. I: using discriminant: p_{T}/m_{4l}
 - Cat. II: using linear discriminant: $V_{\text{D}} = \alpha \Delta\eta_{jj} + \beta m_{jj}$
- Analysis based on correlated 3D distributions:
 - Cat. I: $\mathcal{P}(m_{4l}) \times \mathcal{P}(\text{KD} | m_{4l}) \times \mathcal{P}(p_{\text{T}}/m_{4l} | m_{4l})$
 - Cat. II: $\mathcal{P}(m_{4l}) \times \mathcal{P}(\text{KD} | m_{4l}) \times \mathcal{P}(V_{\text{D}} | m_{4l})$

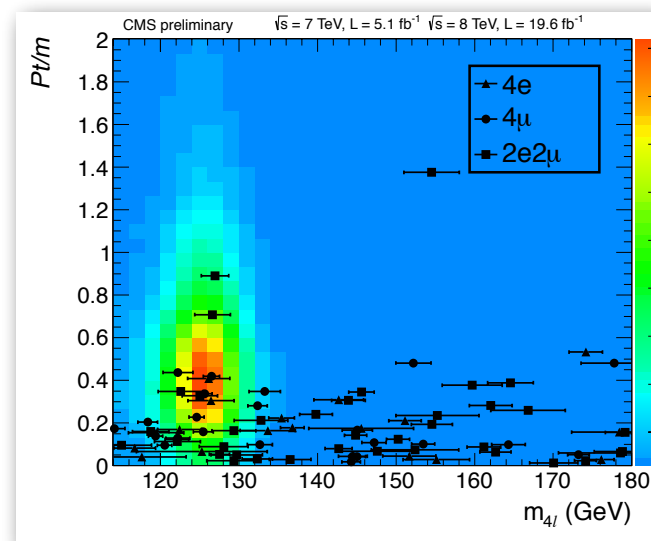
V_{D} distribution



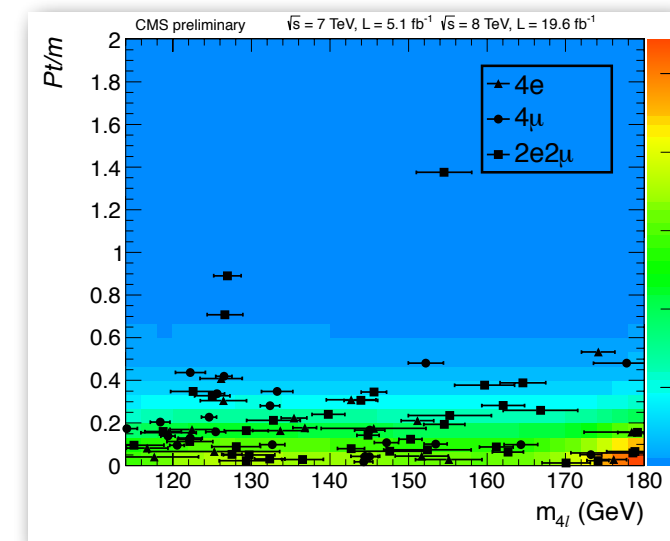
gg → H



VBF



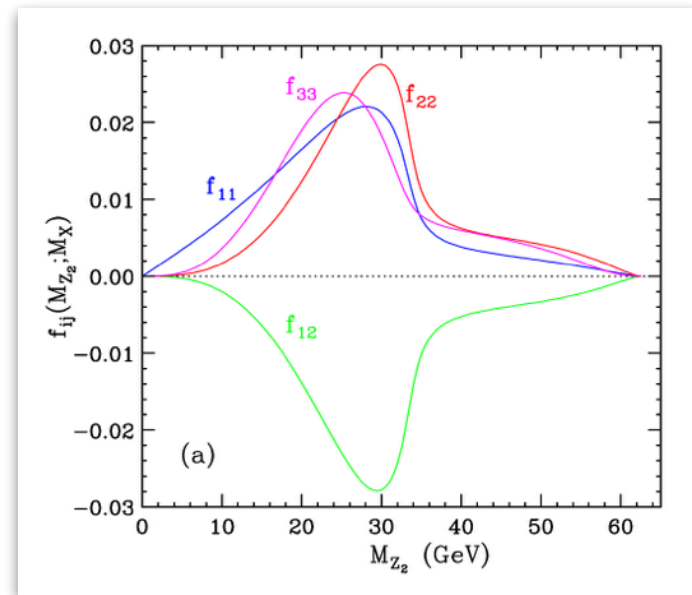
qq → ZZ



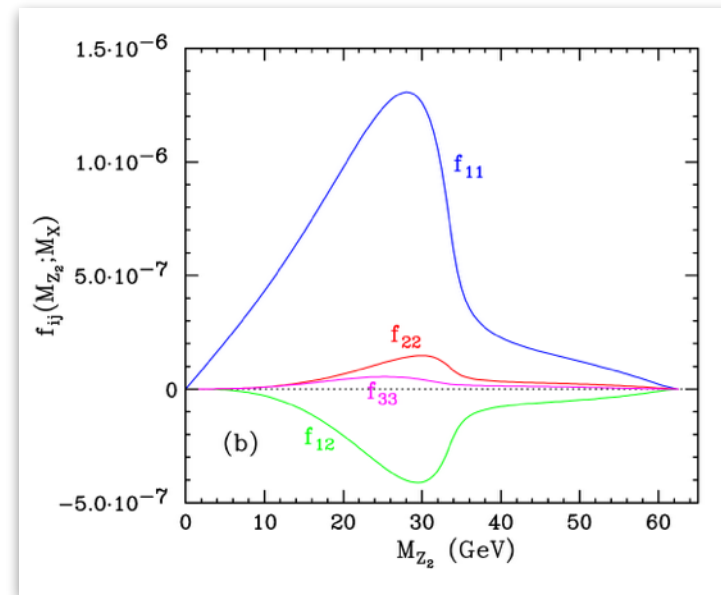
Interference effects - m_{Z_2} example

- Important interference effects between operators in presence of multiple non-zero couplings

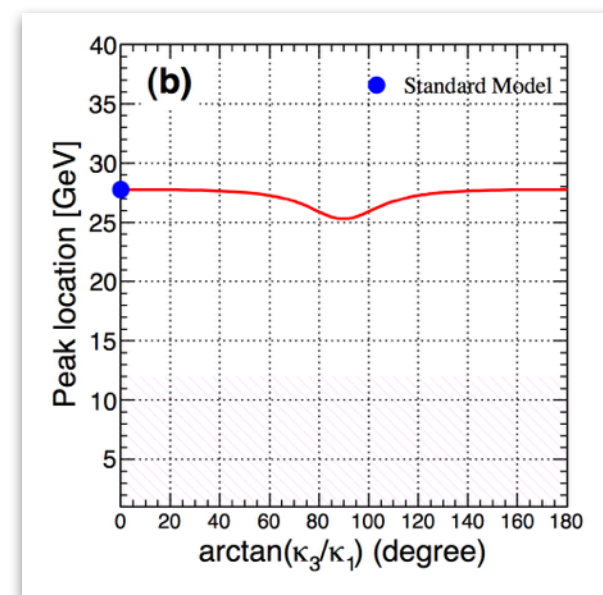
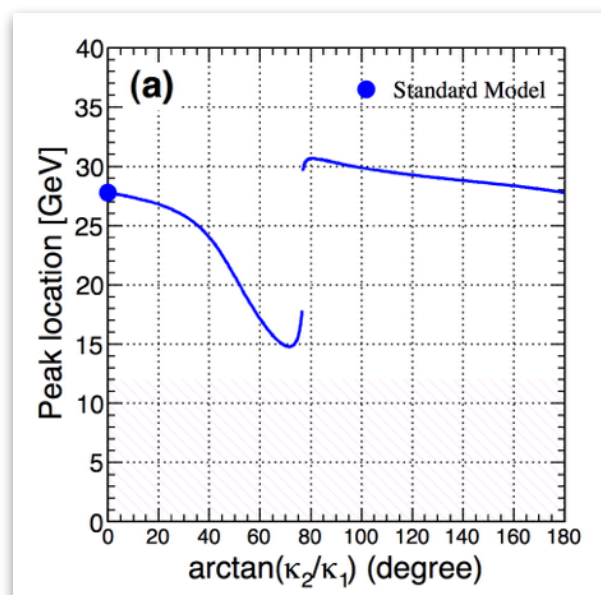
m_{Z_2} distribution (unit normalised)



m_{Z_2} distribution (normalised to yield)



- Peak of m_{Z_2} distribution displays “first order phase transition” from K_1 - K_2 interference, no such feature when considering K_1 and K_3



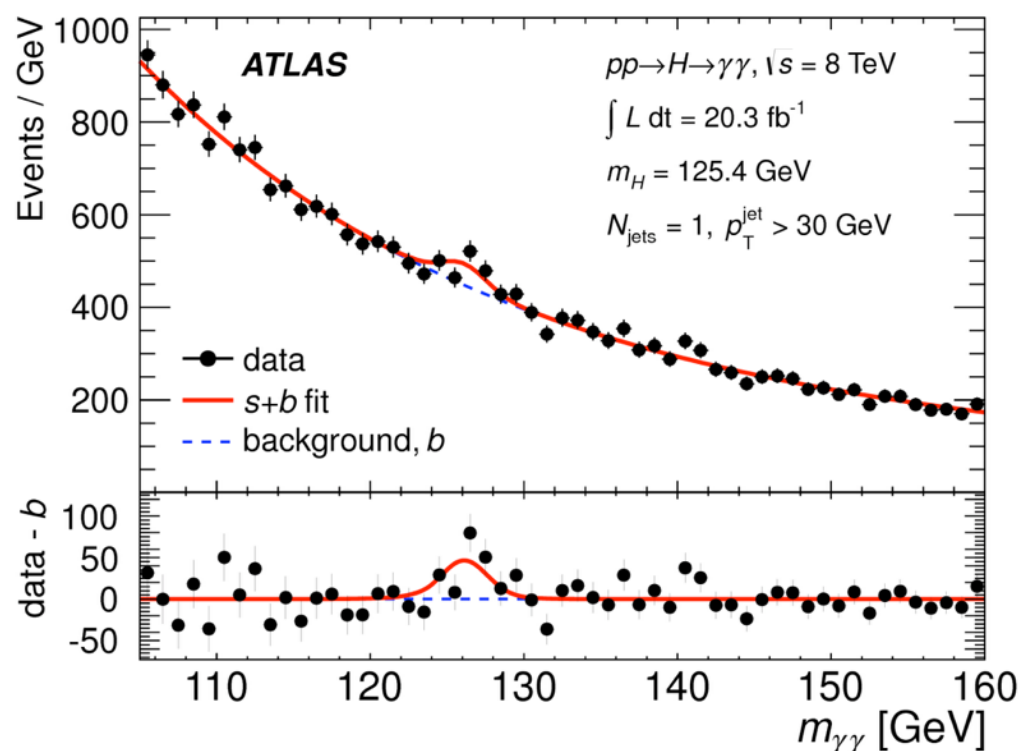
how much we can gain in sensitivity if we optimise the analysis to exploit the interference effects?

Modelling of processes @ higher orders

- Three major improvements in the theoretical Higgs modelling in 2015:
 - **N³LO** cross section, **fully inclusive**
[Anastasiou, Duhr, Dulat, Herzog, Mistlberger (2015)]
 - **NNLO H+J** cross section, **fully exclusive**
[Boughezal, Caola, Melnikov, Petriello, Schulze (2015)], [Boughezal, Focke, Giele, Liu, Petriello (2015)][Chen, Gerhmann, Glover, Jaquier (in progress)]
 - **NNLO VBF** cross section, **fully exclusive**
[Cacciari, Dreyer, Karlberg, Salam, Zanderighi (2015)]
- Improvements in the theoretical modelling of VV' production:
 - **NNLO pp → VV' + X** cross section, **fully exclusive**
[Catani, Cieri, de Florian, Ferrera, Grazzini; Grazzini, Kallweit, Rathlev (Vgamma)] [Gehrmann, Grazzini, Kallweit, Maieroefer, v. Manteuffel, Pozzorini, Rathlev, Tancredi (WW)] [Cascioli, Gehrmann, Grazzini, Kallweit, Maieroefer, v. Manteuffel, Pozzorini, Rathlev, Tancredi, Weihs (ZZ)] [Grazzini, Kallweit, Rathlev (ZZ*, in progress)]

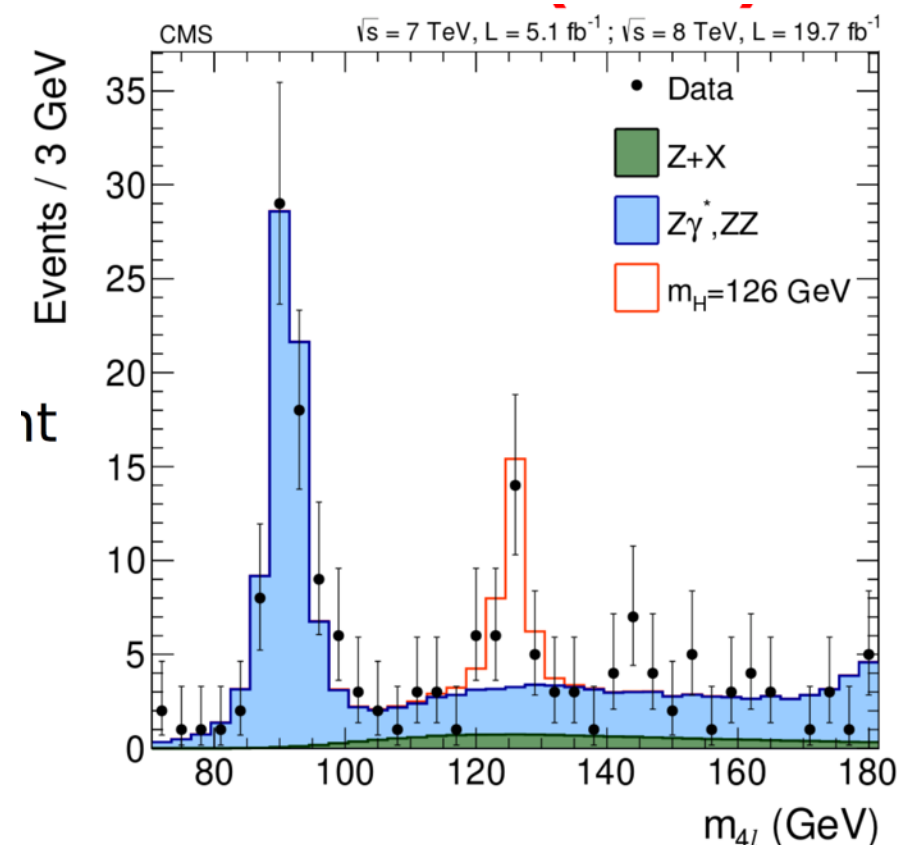
Overview of the channels

- **Fiducial XS measured in $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ \rightarrow 4l$ channels**
 - Fully reconstructible decays with excellent mass resolution ($H \rightarrow \gamma\gamma$: 1-2%, $H \rightarrow ZZ \rightarrow 4l$: 1-3%)



$H \rightarrow \gamma\gamma$ channel:

- Small branching fraction ($\sim 0.2\%$), but allows high selection efficiency ($\sim 40-45\%$)
- A few hundred reconstructed signal events expected after selection (8 TeV Run I)
- Large continuum background from QCD $\gamma\gamma$ and γ +jet.



$H \rightarrow ZZ \rightarrow 4l$ channel:

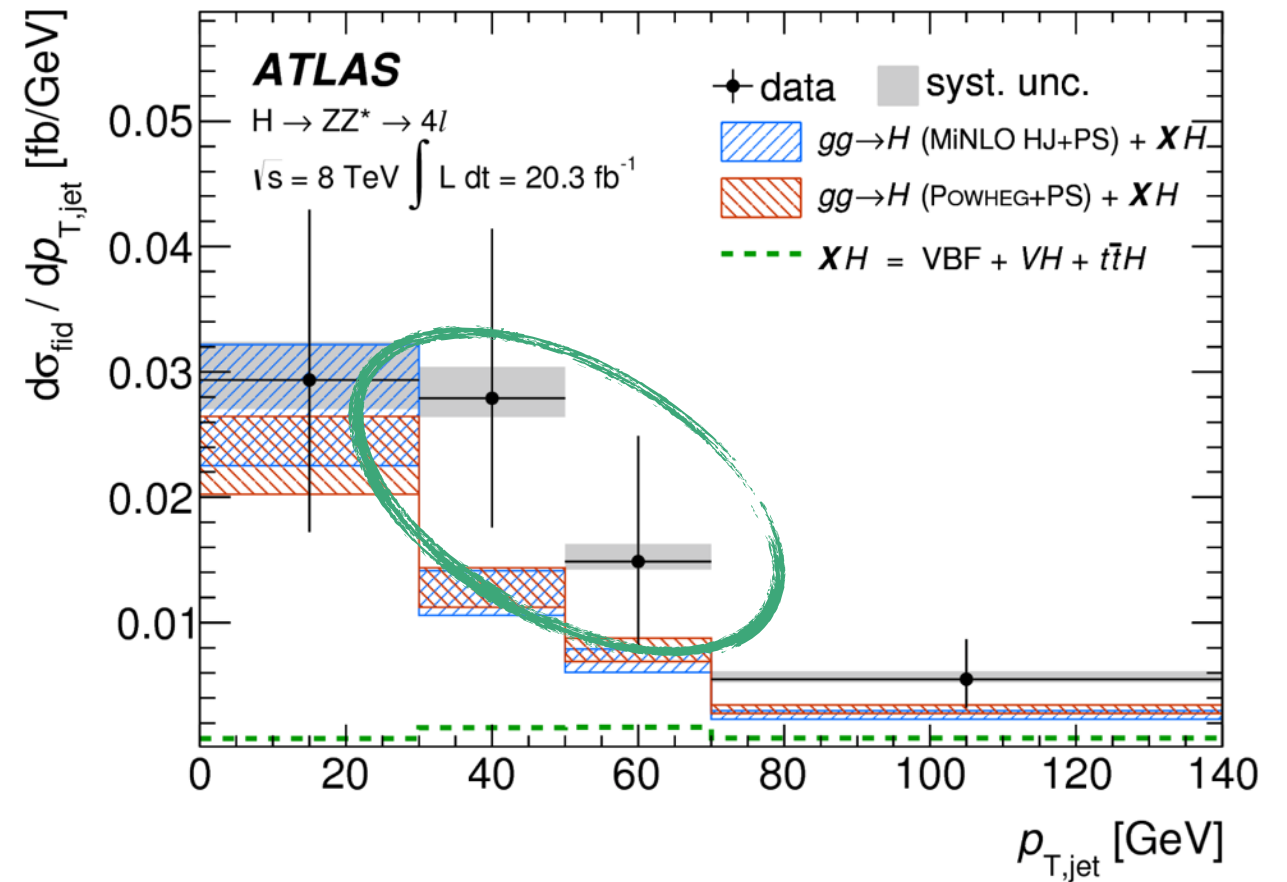
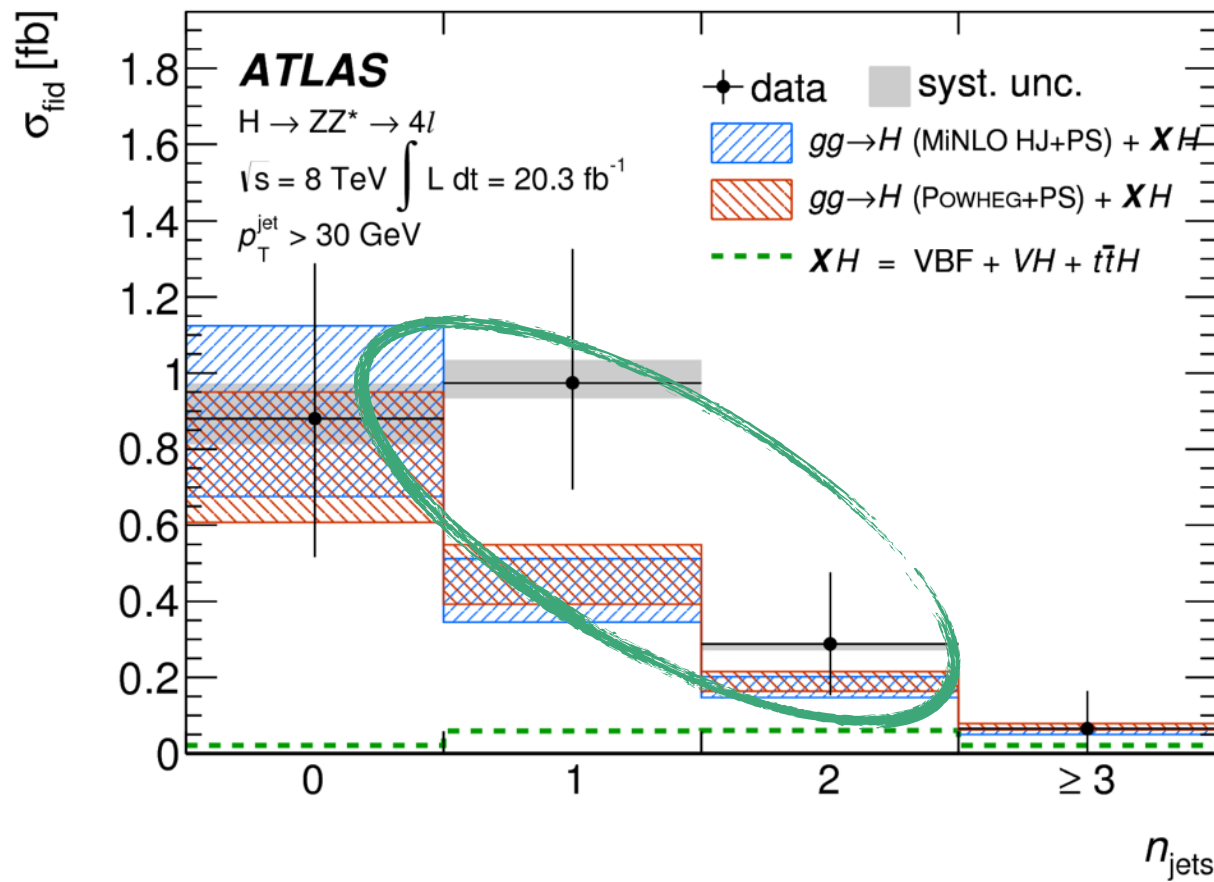
- Very small branching fraction ($\sim 0.01\%$), but very clean signature
- About 15-20 reconstructed signal events expected after selection (8 TeV Run I)
- Very small background from irreducible $pp \rightarrow ZZ$ and reducible Z +jets.

Differential XS: $H \rightarrow ZZ \rightarrow 4l$

PLB 738 (2014) 234

- Measured as a function of several jet-related observables:
 - $N(\text{jet}), p_{T}(j_1)$ (many other observables, see backup)
 - Sensitive to theoretical modelling of hard quark and gluon radiation, relative contributions of different production modes, BSM effects, etc.

ATLAS @ 8TeV:



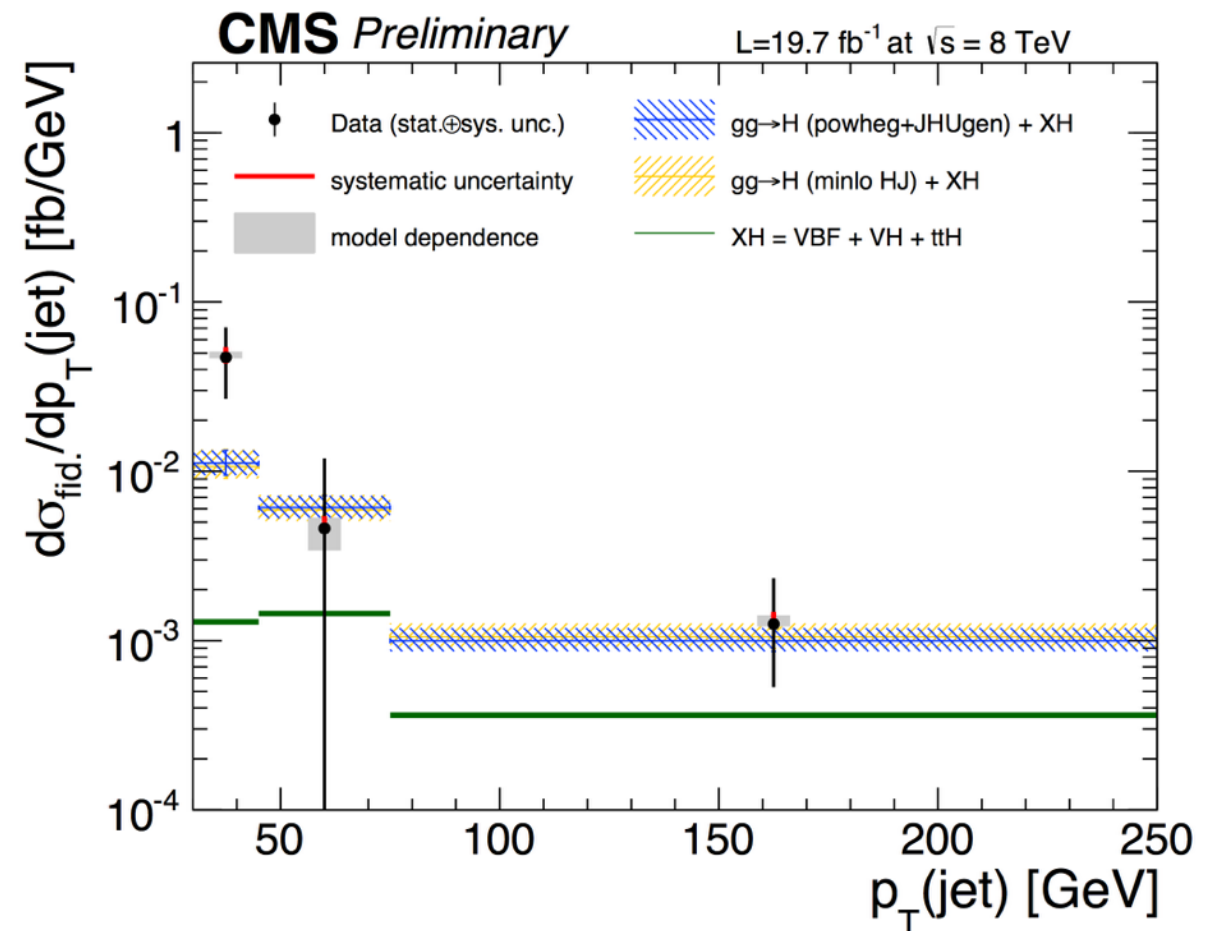
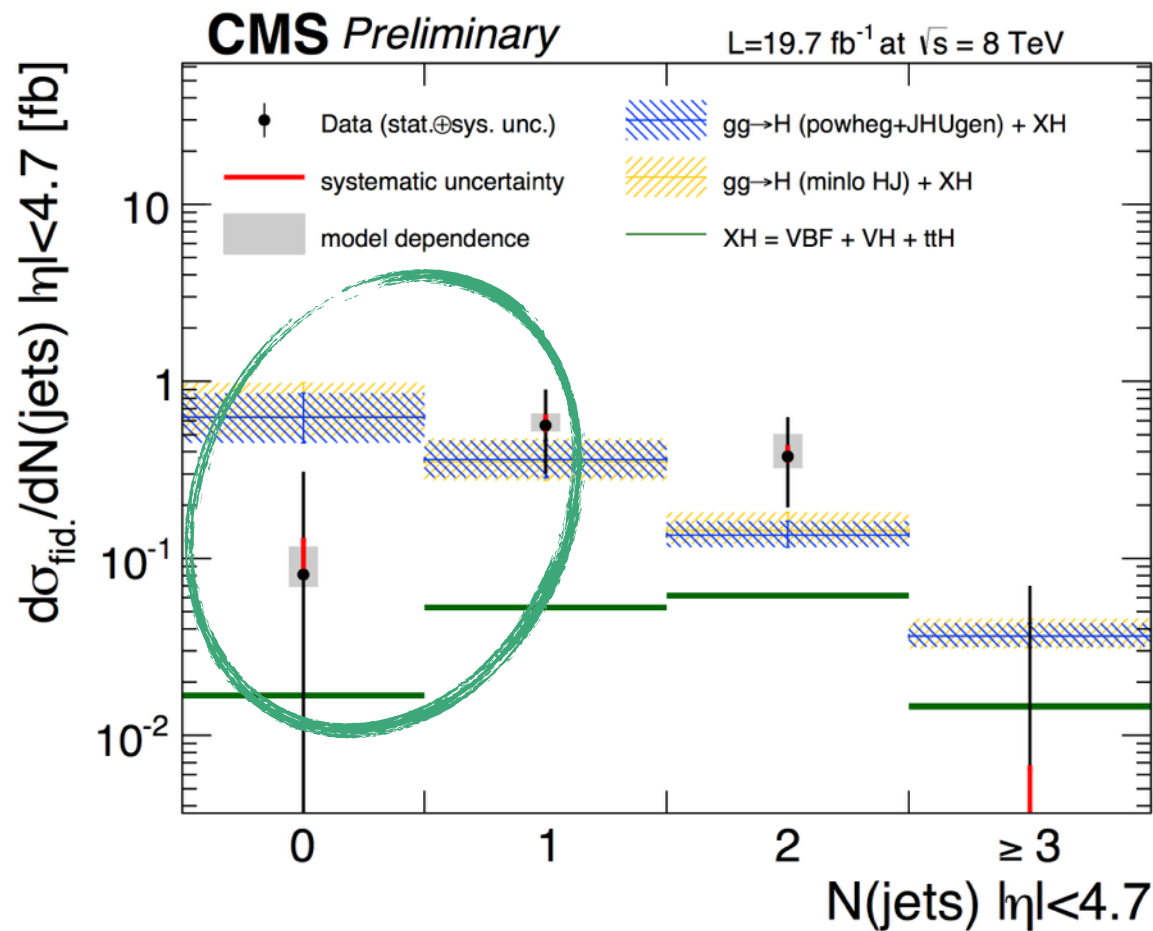
► Somewhat higher jet activity observed in data w.r.t theoretical estimates

Differential XS: $H \rightarrow ZZ \rightarrow 4l$

CMS-PAS-HIG-14-028

- Measured as a function of several jet-related observables:
 - $N(\text{jet}), p_T(j_1)$ (many other observables, see backup)
 - Sensitive to theoretical modelling of hard quark and gluon radiation, relative contributions of different production modes, BSM effects, etc.

CMS @ 8TeV:



Somewhat higher jet activity observed in data w.r.t theoretical estimates also seen by CMS: p -value $\sim 13\%$

Inclusive cross sections: $H \rightarrow \gamma\gamma$

JHEP 09 (2014) 112

CMS-PAS-HIG-14-016

- Inclusive cross sections at 8TeV, measured in $H \rightarrow \gamma\gamma$ channel
 - Un-binned maximum likelihood fit to $m_{\gamma\gamma}$ (primarily based on legacy analyses)
 - Compared to theoretical estimates with NNLO+NNLL QCD accuracy
 - ATLAS: compared also to other state-of-art predictions.

ATLAS @ 8TeV:

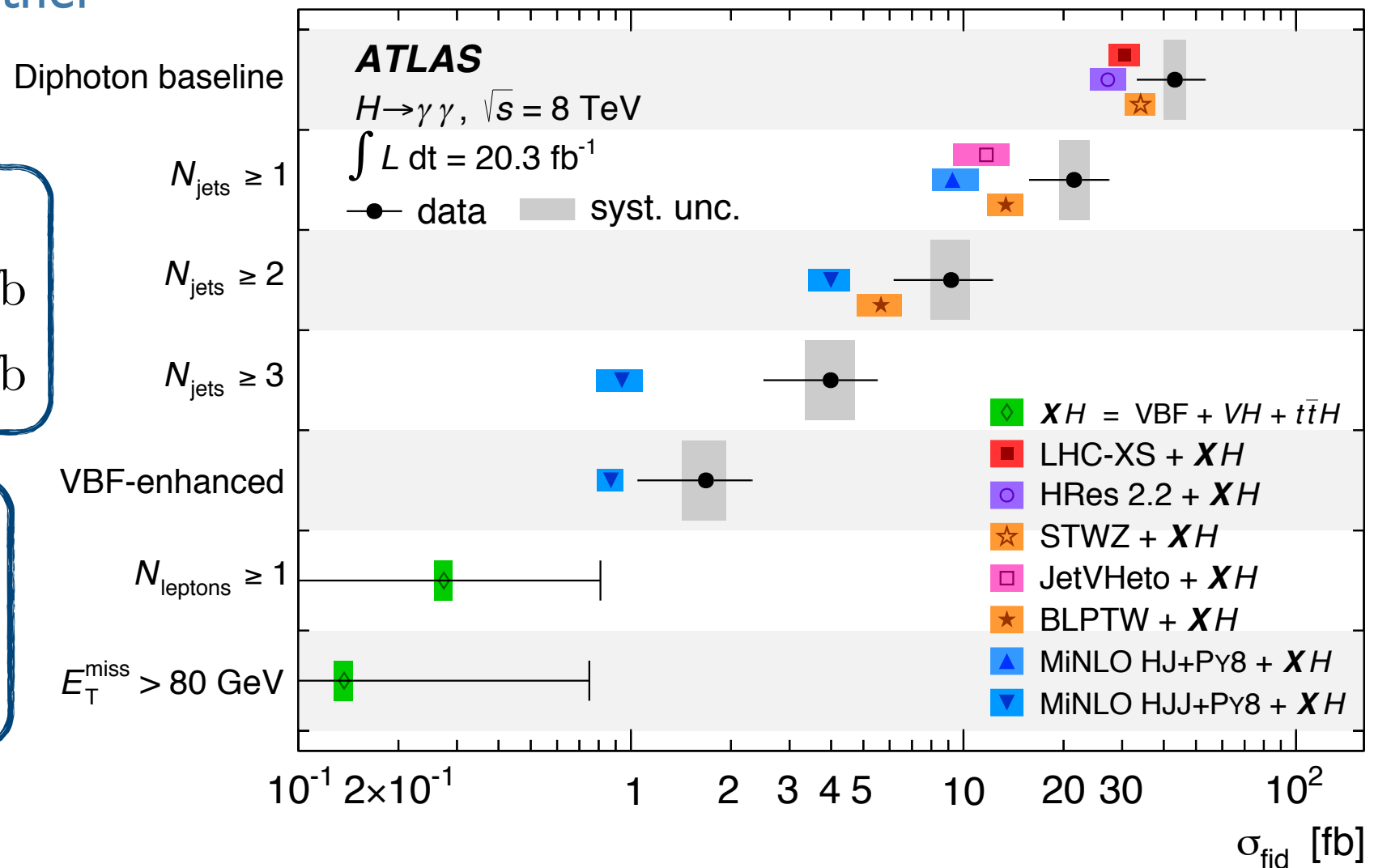
$$\sigma_{\text{fid}} = 43.2 \pm 9.4(\text{stat.})^{+3.2}_{-2.9}(\text{syst.}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 30.5 \pm 3.3 \text{ fb}$$

CMS @ 8TeV:

$$\sigma_{\text{fid}} = 32 \pm 10(\text{stat.}) \pm 3(\text{syst.}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 31^{+4}_{-3}(\text{stat.}) \text{ fb}$$



Inclusive cross sections: $H \rightarrow \gamma\gamma$

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CMS-PAS-HIG-14-016

- Inclusive cross sections at 8TeV, measured in $H \rightarrow \gamma\gamma$ channel
 - Un-binned maximum likelihood fit to $m_{\gamma\gamma}$ (primarily based on legacy analyses)
 - Compared to theoretical estimates with NNLO+NNLL QCD accuracy
 - ATLAS: compared also to other state-of-art predictions.

ATLAS @ 8TeV:

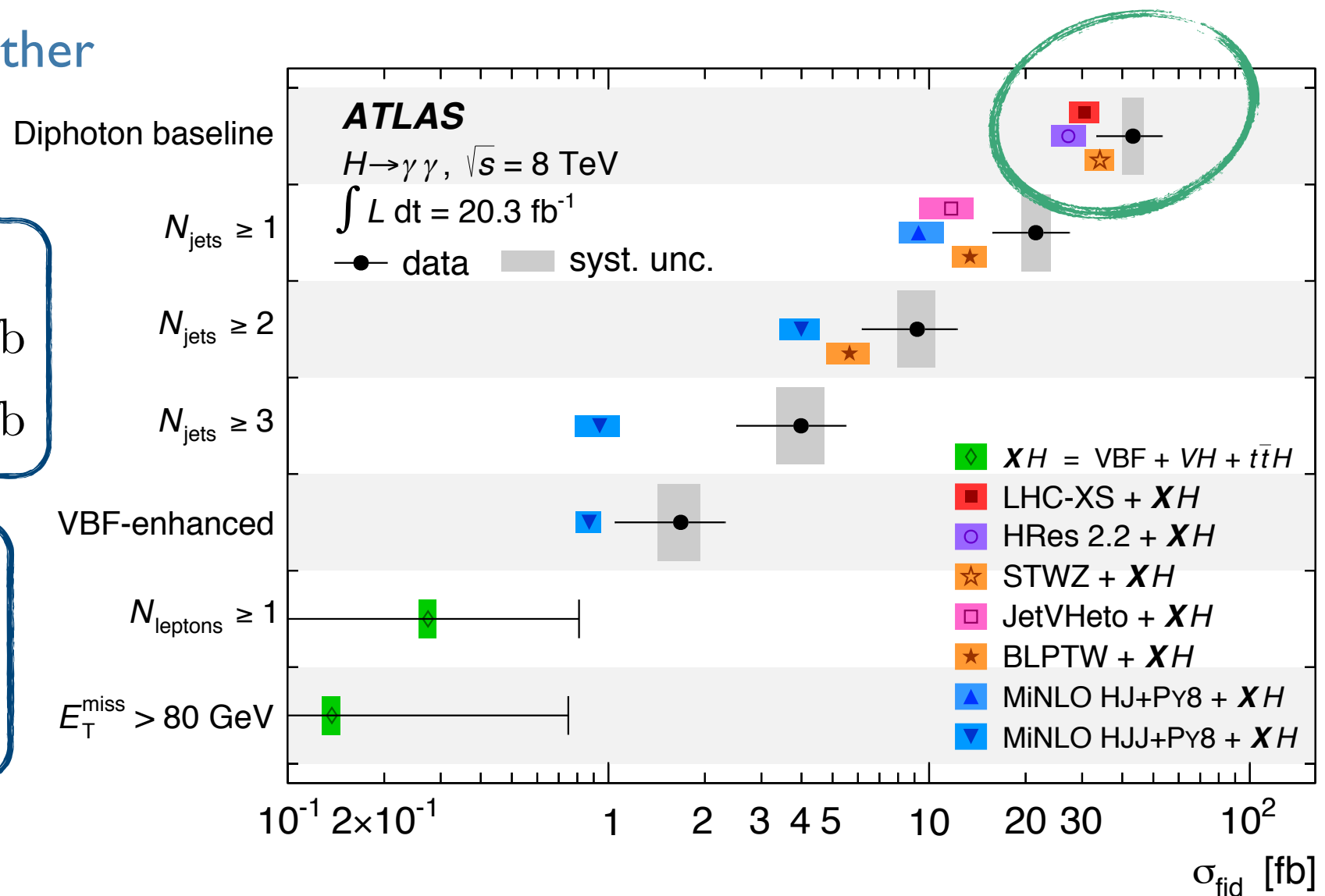
$$\sigma_{\text{fid}} = 43.2 \pm 9.4(\text{stat.})^{+3.2}_{-2.9}(\text{syst.}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 30.5 \pm 3.3 \text{ fb}$$

CMS @ 8TeV:

$$\sigma_{\text{fid}} = 32 \pm 10(\text{stat.}) \pm 3(\text{syst.}) \text{ fb}$$

$$\sigma_{\text{fid}}^{\text{SM}} = 31^{+4}_{-3}(\text{stat.}) \text{ fb}$$



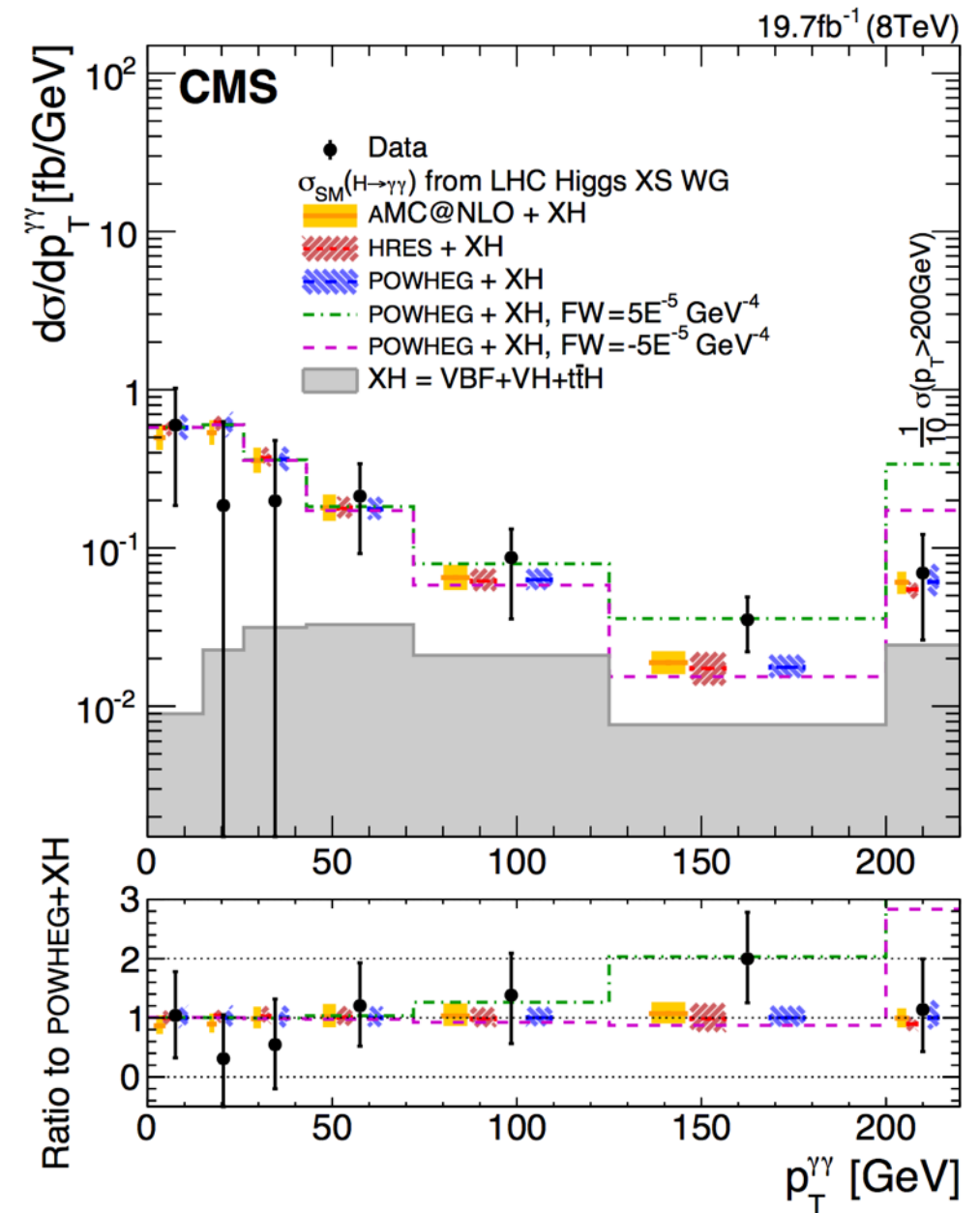
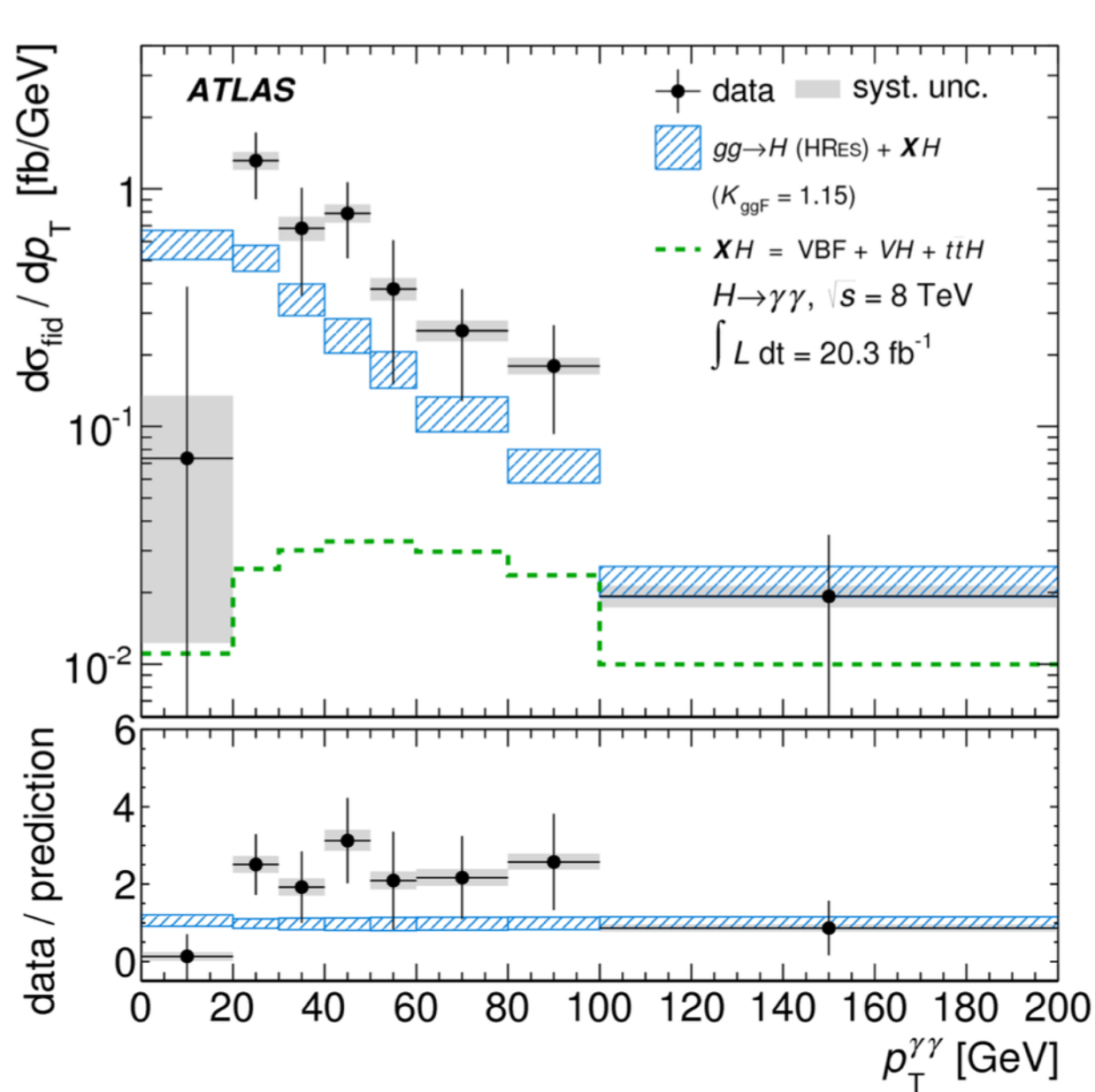
► **Compatible with theoretical estimates (slightly higher rates by ATLAS)**

Differential XS: $H \rightarrow \gamma\gamma$

JHEP 09 (2014) 112

CMS-PAS-HIG-14-016

- Measured as a function of Higgs candidate kinematic properties:
 - $p_T(\mathbf{H})$: sensitive to production mode, new physics in $gg \rightarrow H$ loop
 - $|\mathbf{y}(\mathbf{H})|$: sensitive to production mode, parton distribution functions (see backup).

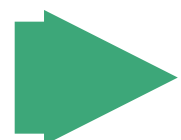
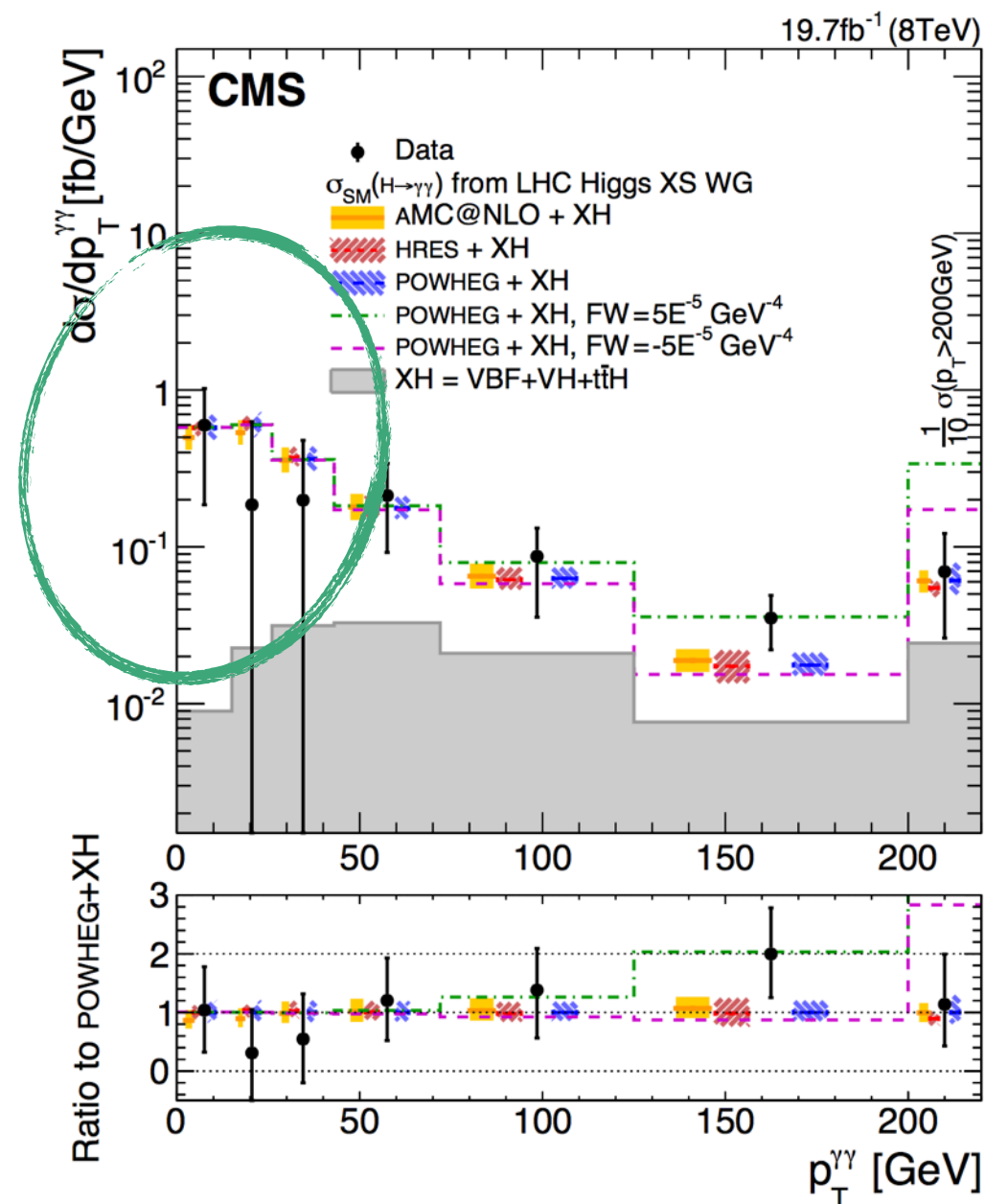
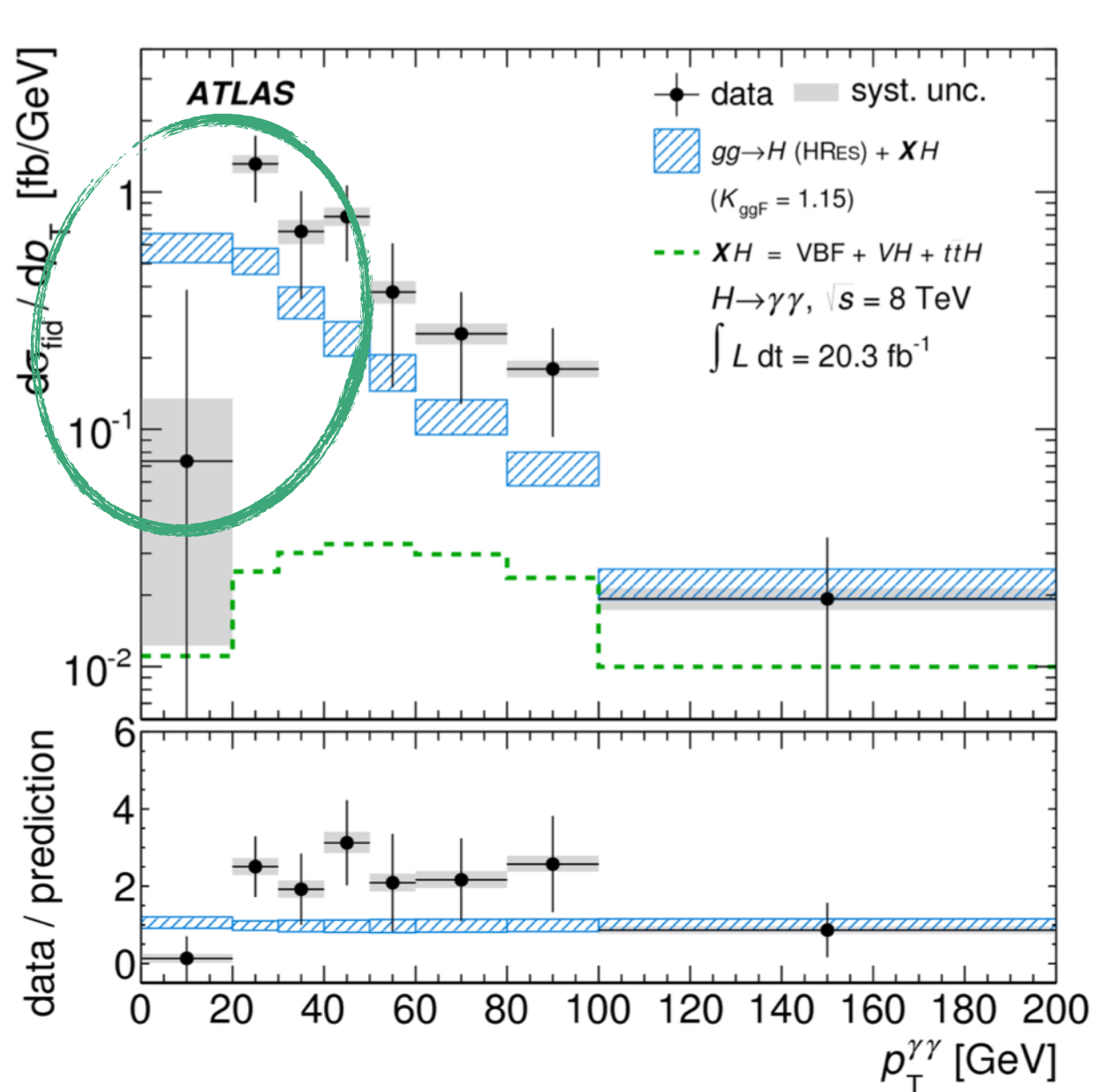


Differential XS: $H \rightarrow \gamma\gamma$

JHEP 09 (2014) 112

CMS-PAS-HIG-14-016

- Measured as a function of Higgs candidate kinematic properties:
 - $p_T(\mathbf{H})$: sensitive to production mode, new physics in $gg \rightarrow H$ loop
 - $|\mathbf{y}(\mathbf{H})|$: sensitive to production mode, parton distribution functions (see backup).



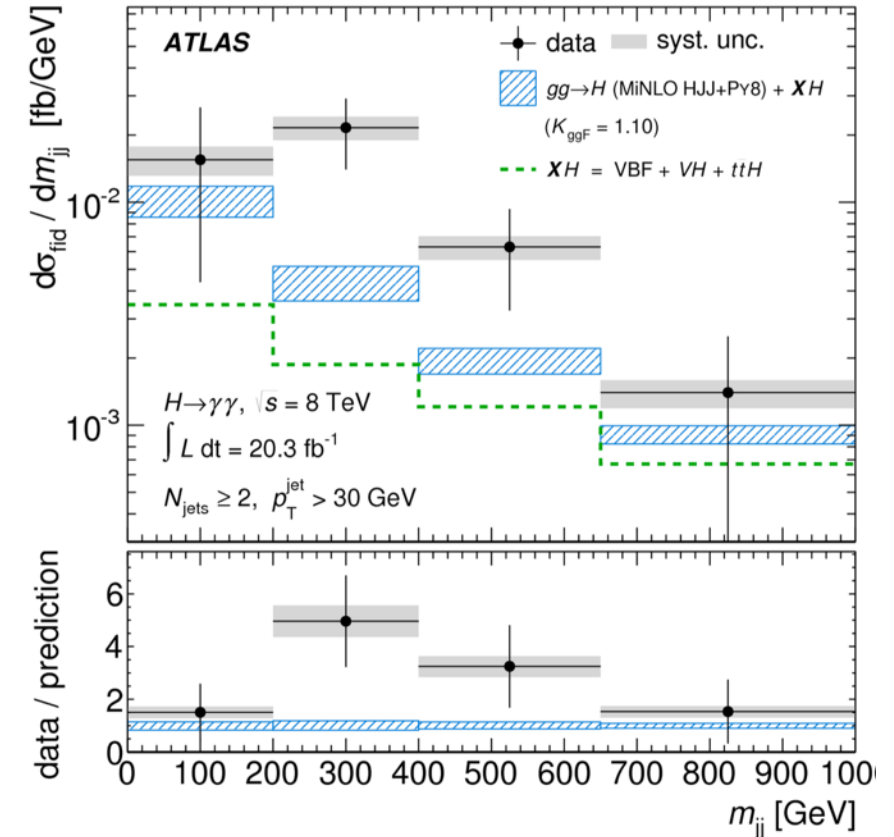
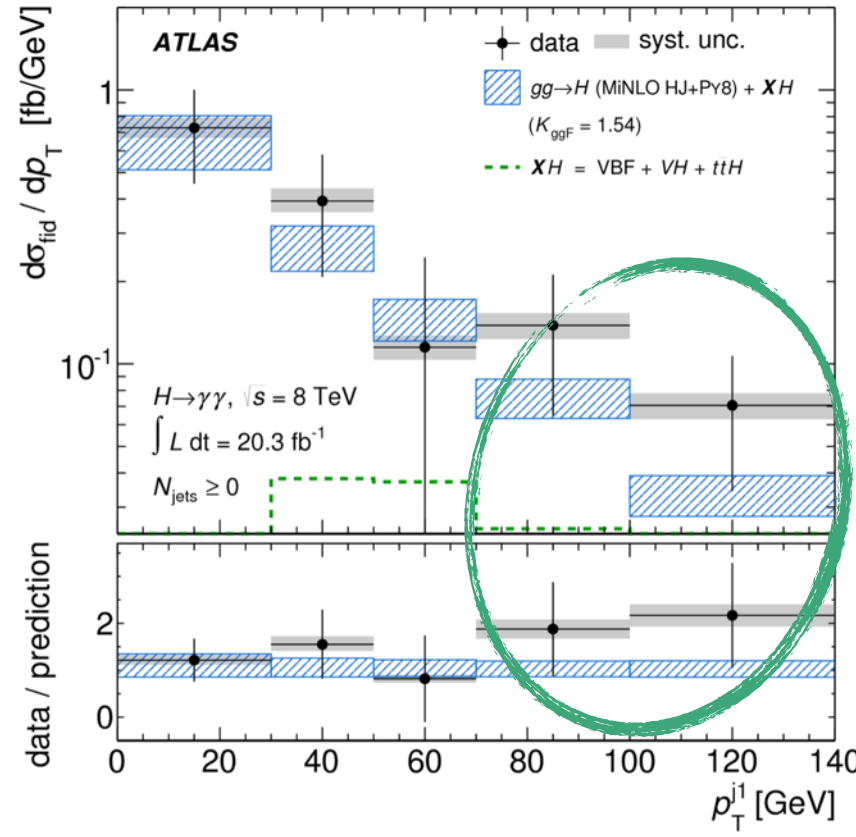
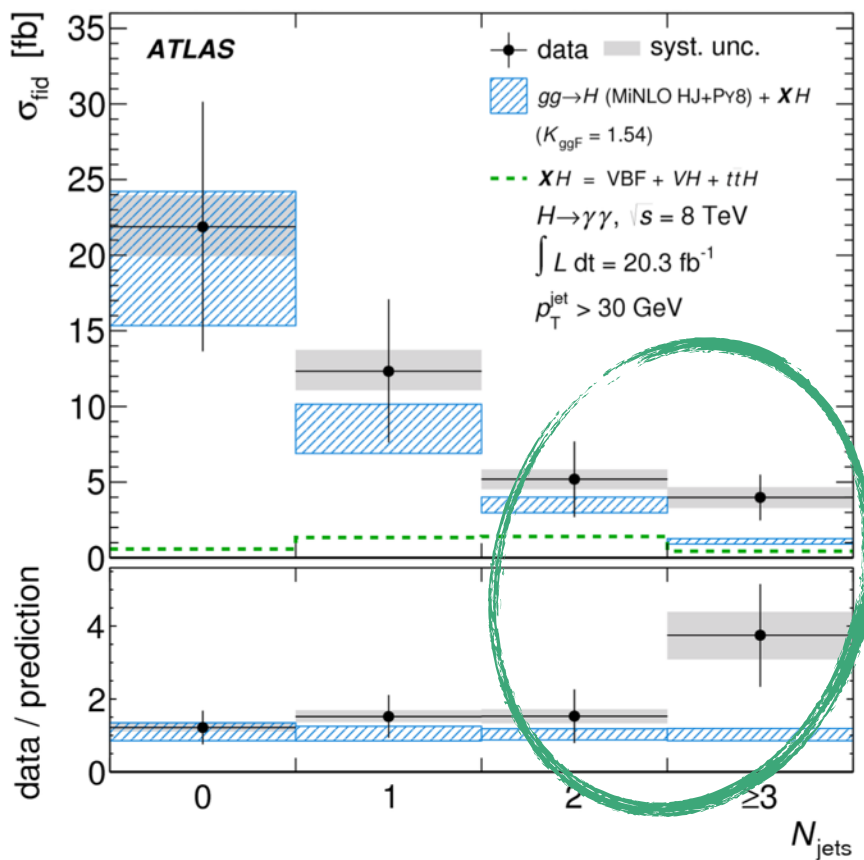
Some deficit in low- p_T range observed by ATLAS (not in CMS)

Differential XS: $H \rightarrow \gamma\gamma$

JHEP 09 (2014) 112

- Measured as a function of several jet-related observables:
 - $N(\text{jet})$, $p_T(j_1)$, $m(jj)$ (many other observables, see backup)
 - Sensitive to theoretical modelling of hard quark and gluon radiation, relative contributions of different production modes, BSM effects, etc.

ATLAS @ 8TeV:



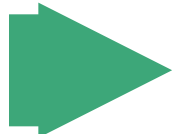
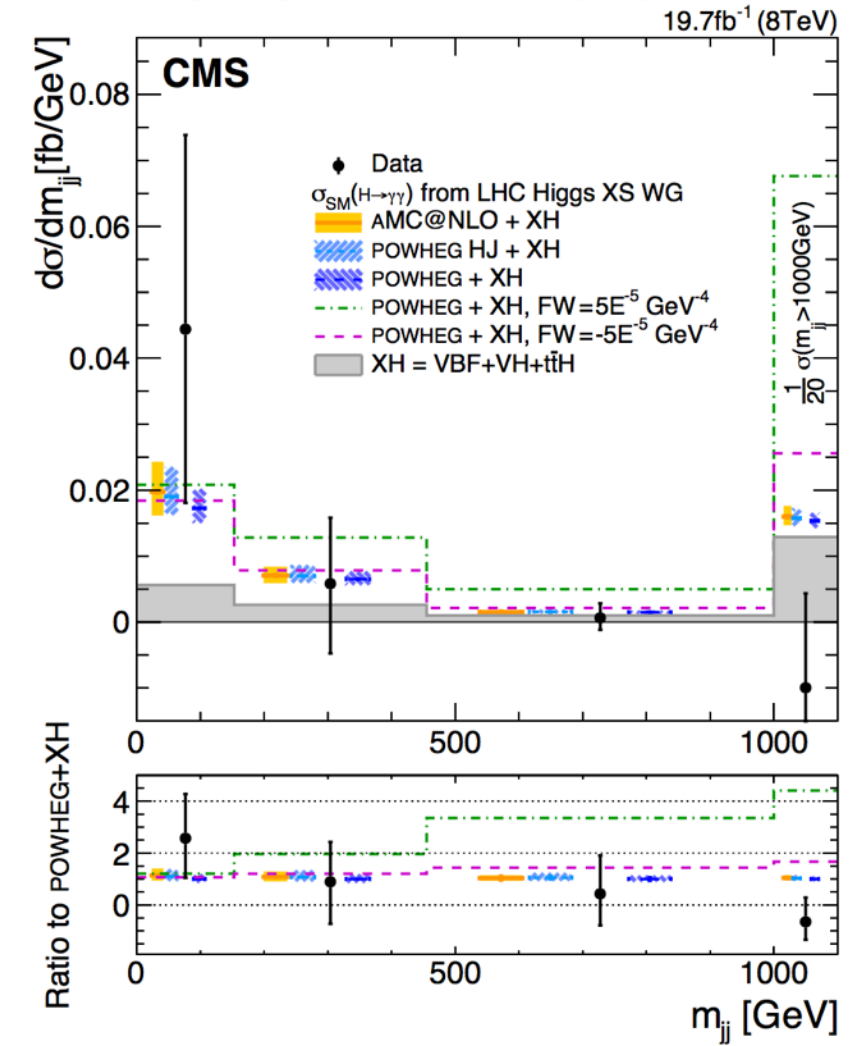
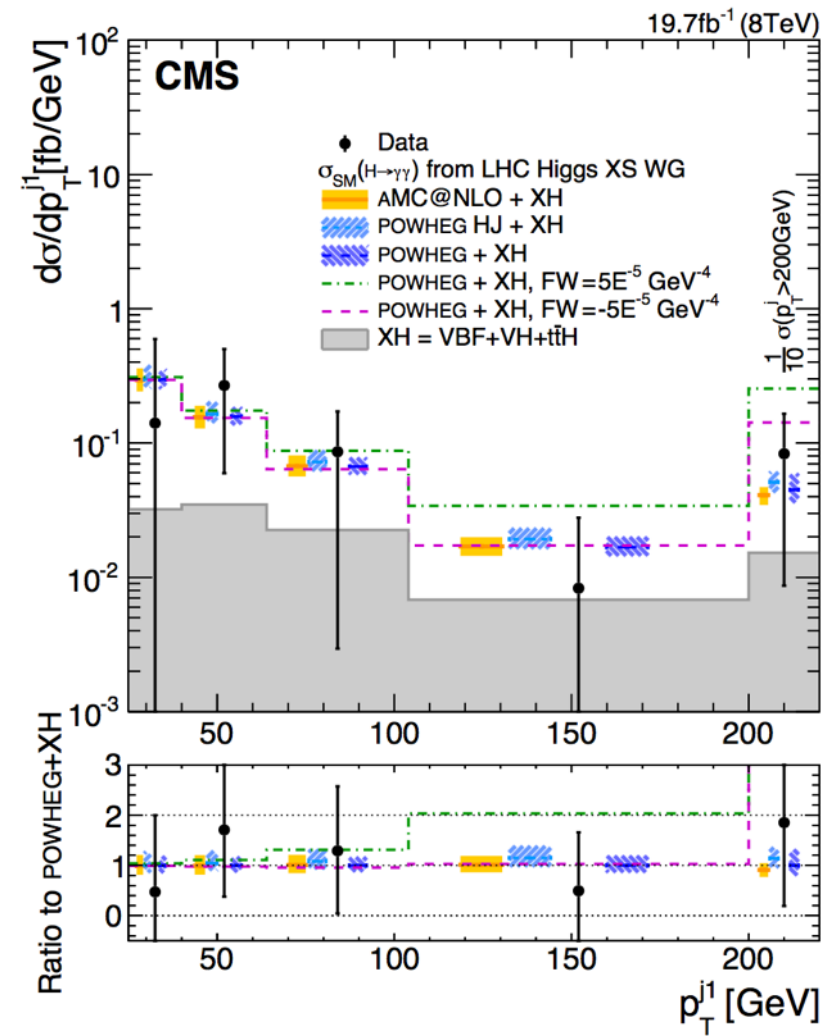
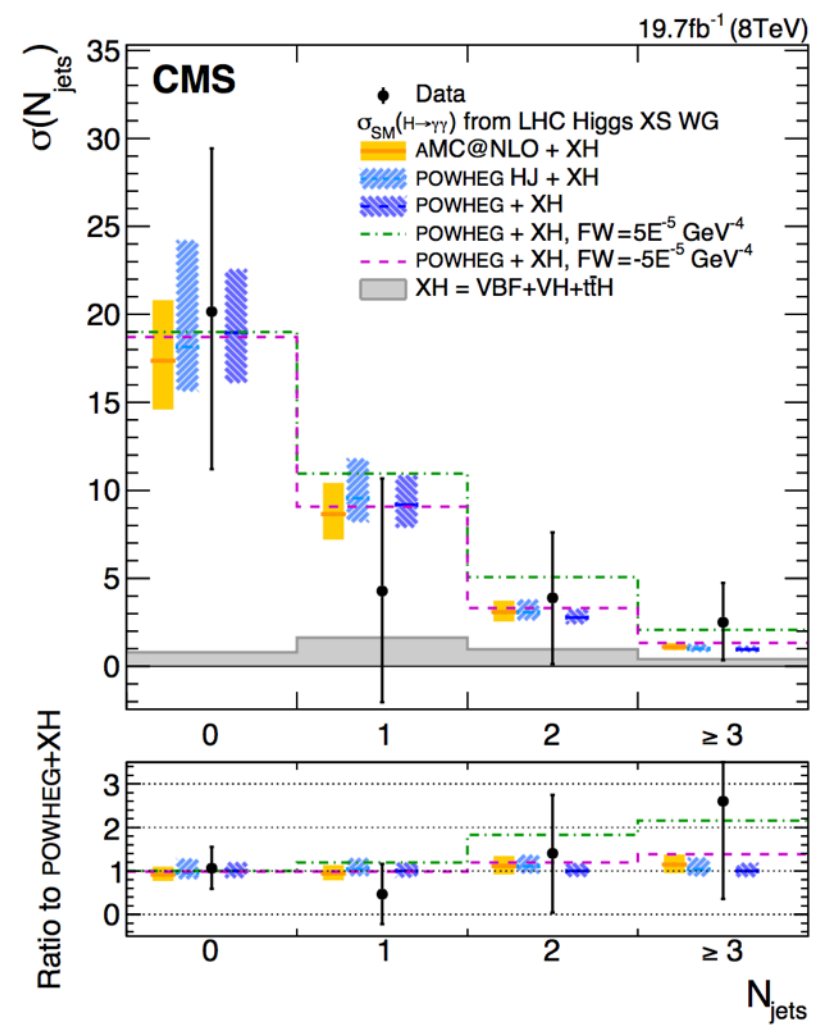
Slightly higher jet activity observed in data w.r.t theoretical estimates

Differential XS: $H \rightarrow \gamma\gamma$

CMS-PAS-HIG-14-016

- Measured as a function of several jet-related observables:
 - $N(\text{jet})$, $p_T(j_1)$, $m(jj)$ (many other observables, see backup)
 - Sensitive to theoretical modelling of hard quark and gluon radiation, relative contributions of different production modes, BSM effects, etc.

CMS @ 8TeV:

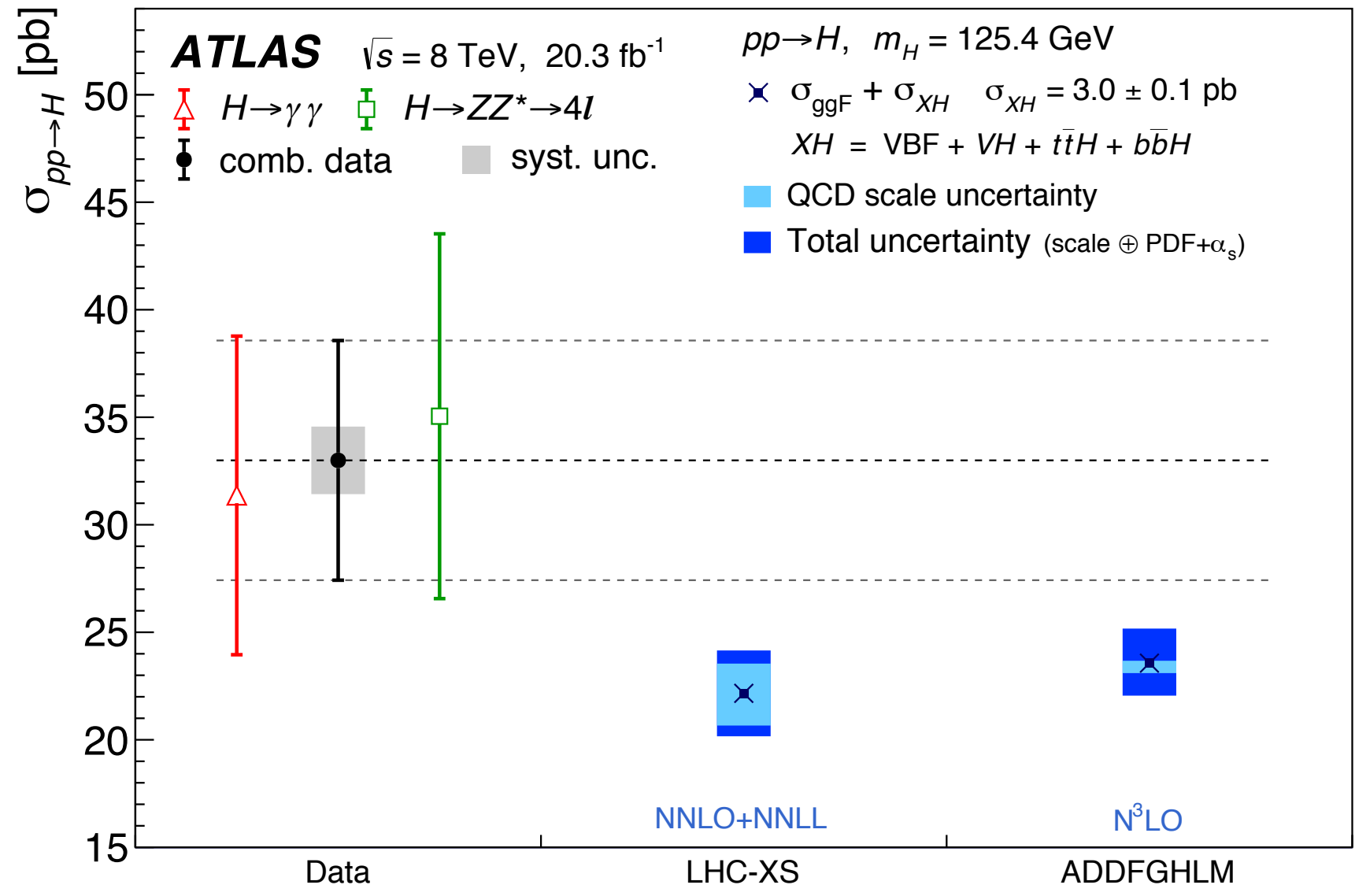


Good agreement with theoretical estimates in case of CMS

Combination: Inclusive XS ($H \rightarrow \gamma\gamma$, $H \rightarrow ZZ$)

PRL. 115 (2015) 091801

- Inclusive XS from the combination of two channels
 - Compared to theoretical estimates with **NNLO+NNLL** and **N³LO** accuracy in **QCD**.



p-values:
LHC-XS: 5.5 %
ADDFGHLM: 9%

H → 4l: 35.0 ± 8.4 (stat) ± 1.8 (sys) pb
H → yy: 31.4 ± 7.2 (stat) ± 1.6 (sys) pb
Combined: 33.0 ± 5.3 (stat) ± 1.6 (sys) pb

Total inclusive cross section measured by ATLAS is higher than state-of-art predictions