

# Timelike Compton Scattering: polarization observables and experimental perspectives for JLab at 12 GeV

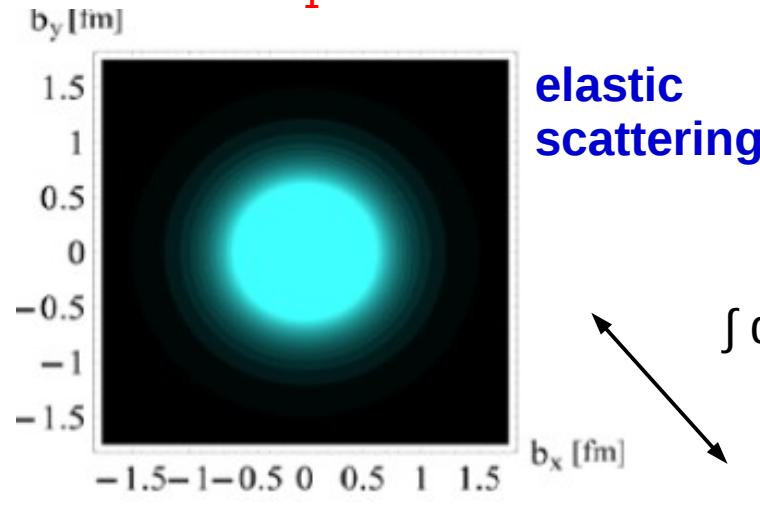
POETIC 2015 conference, Palaiseau, FRANCE

Marie Boér, IPN Orsay, FRANCE  
Sept. 11, 2015

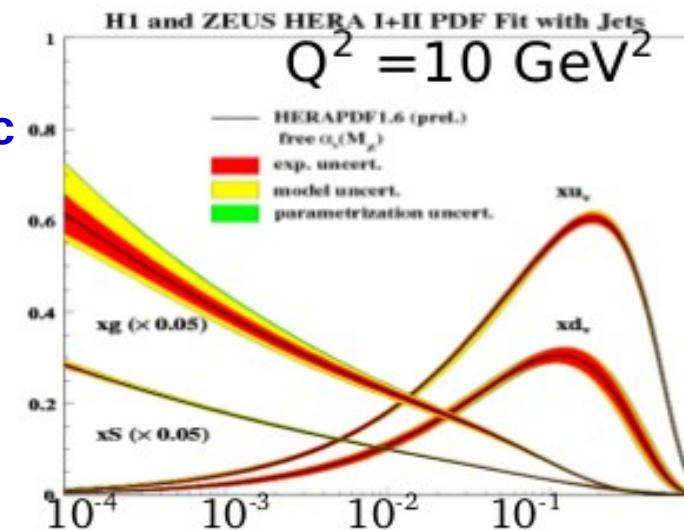
In collaboration with: M. Guidal, M. Vanderhaeghen ;  
A.&H. Mkrtchyan, V. Tadevosyan, P. Nadel-Turonski, J. Zhang, Z. Zhao (...)

# "3D" imaging : Generalized Parton Distributions (GPDs)

Form factors  $F_1(t) \rightarrow$  transverse charge densities



Parton Distributions  $q(x)$



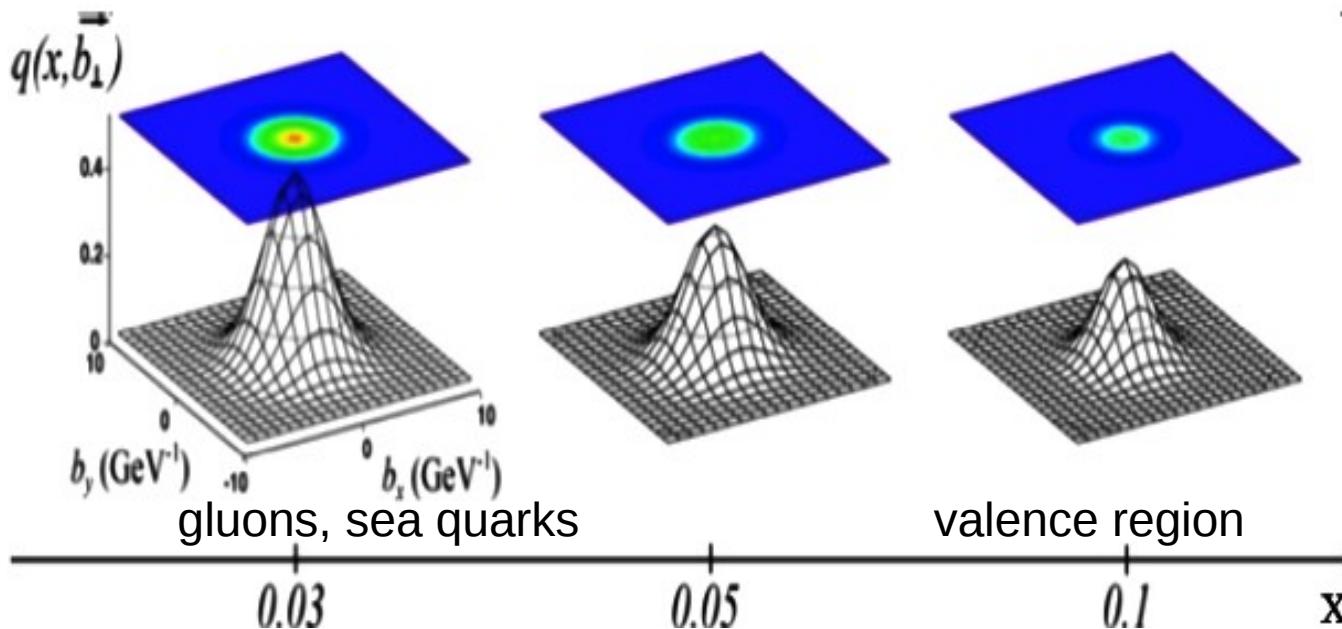
$$\int dx$$

$$t = 0$$

hard exclusive processes

Generalized Parton Distributions

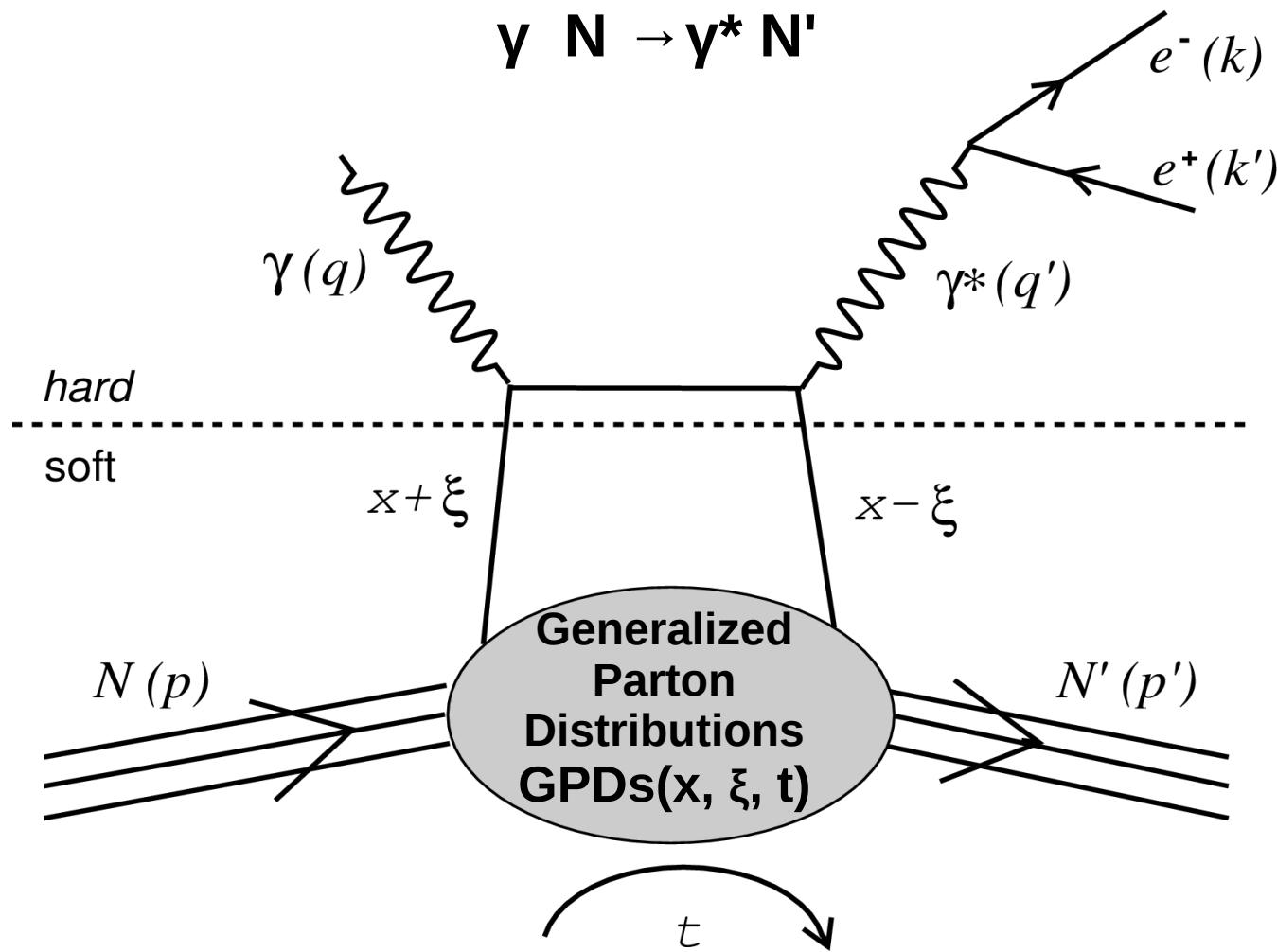
Correlation between longitudinal momentum fraction  $x$  and transverse charge densities  $b_{x,y}$



Nucleon tomography :  
 FT of GPD H ( $x, 0, |t| = \Delta_\perp^2$ )

$x$  : longitudinal momentum fraction  
 $t$  : momentum transfer squared

# Timelike Compton Scattering



$x$  : average longitudinal momentum fraction of the struck quark

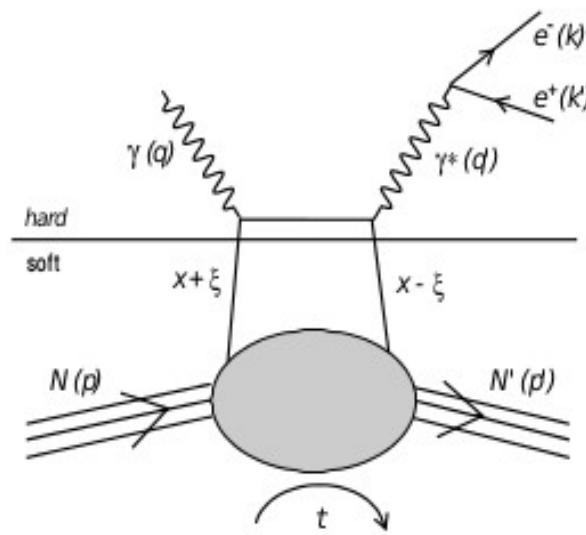
$\xi$  : longitudinal momentum transfer

$t$  : momentum transfer squared

$Q'^2 = +q'^2$  : invariant mass of the lepton pair

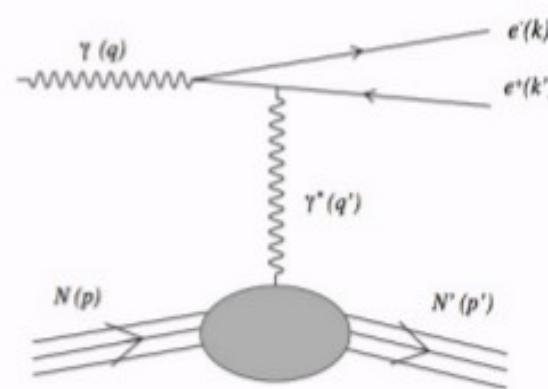
# Interference with Bethe-Heitler

$$\gamma N \rightarrow e^+ e^- N =$$



**Timelike Compton Scattering (TCS)  
sensitive to the nucleon GPDs**

**Bethe-Heitler (BH)  
sensitive to the nucleon Form Factors**



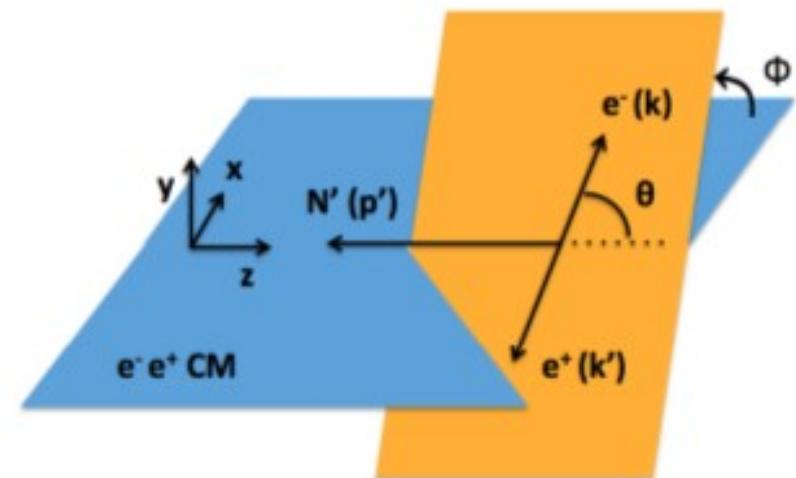
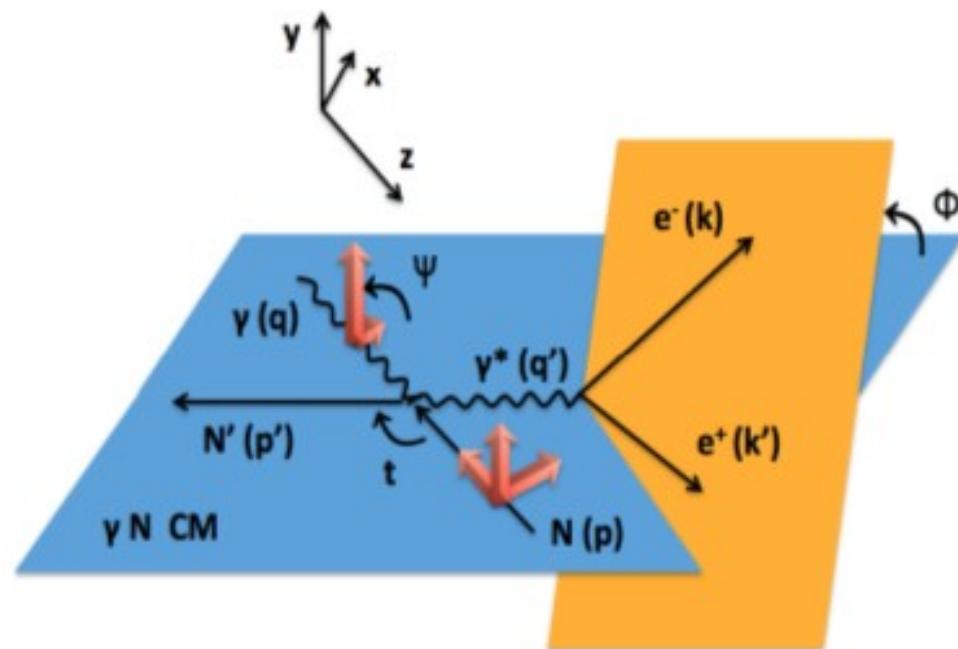
# TCS : angles and notations

$\gamma N \rightarrow e^+ e^- N$

Fixed beam energy  
or  $\xi$

$$\frac{d\sigma}{dQ'^2 dt d\phi d(\cos\theta)}$$

- $\Psi$ : (reaction plane,  $\gamma$  spin)
- $\varphi$ : (hadronic plane,  $e^+ e^-$  pair)
- $\Theta$ : ( $\gamma^*$ ,  $e^-$ )



## Notations

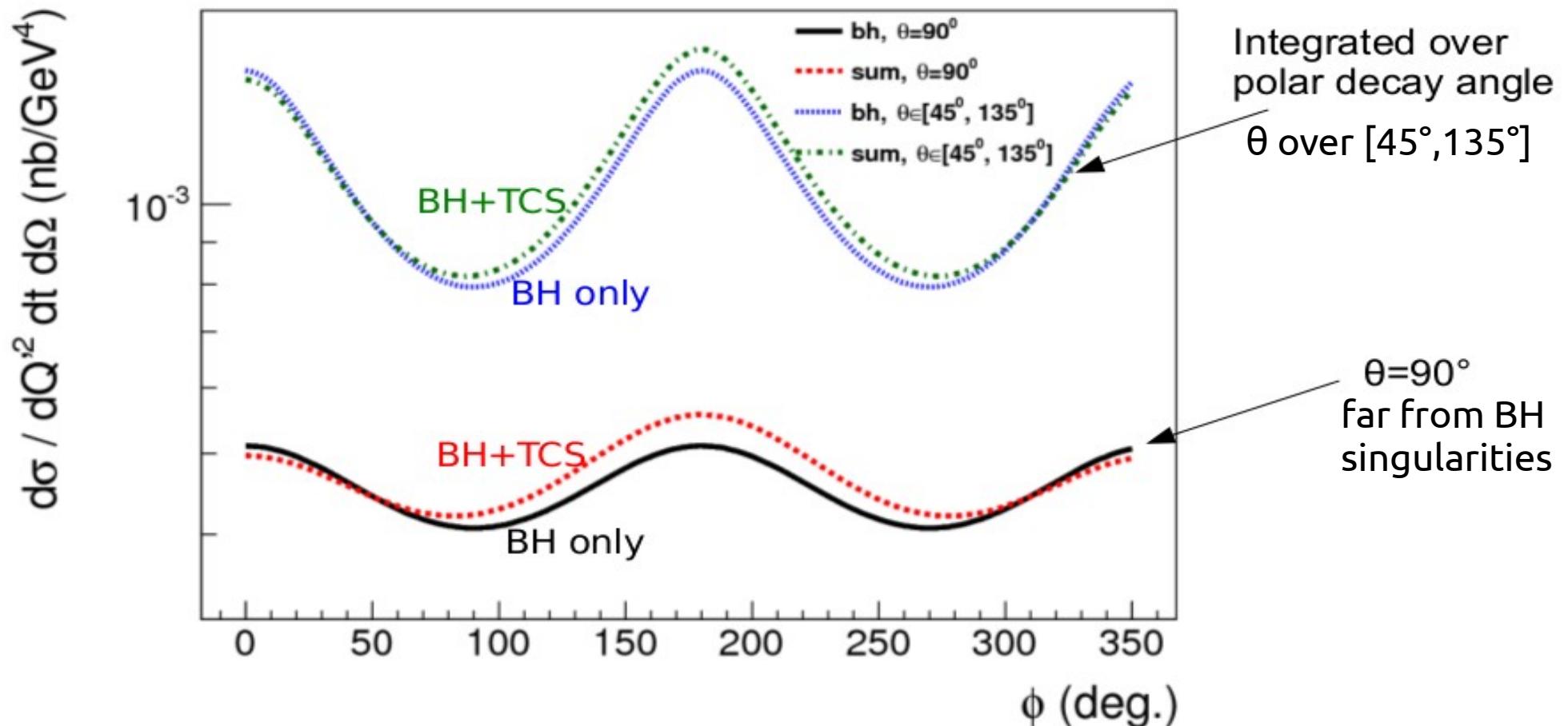
$A_{ij}$ : asymmetry

1st index: photon polarisation,  $\odot$  = circular,  $L$  = linear,  $U$  = unpolarized

2d index: nucleon polarisation,  $x$  (transverse, in plane),  $y$  (transverse),  $z$  (longitudinal)

# BH and TCS cross section angular dependencies

BH singularities :  $e^-$  in  $\gamma$  direction ( $\theta \rightarrow 0^\circ$ ) => singularity at  $\phi=180^\circ$   
 $e^+$  in  $\gamma$  direction ( $\theta \rightarrow 180^\circ$ ) => singularity at  $\phi=0^\circ$



BH is largely dominant, only few % from TCS

Integrated in all following figures

# BH and TCS kinematical dependencies

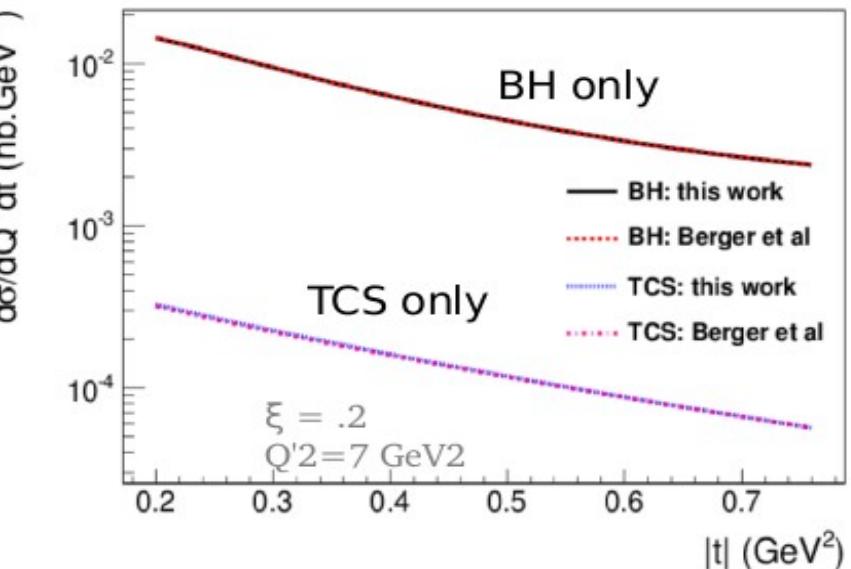
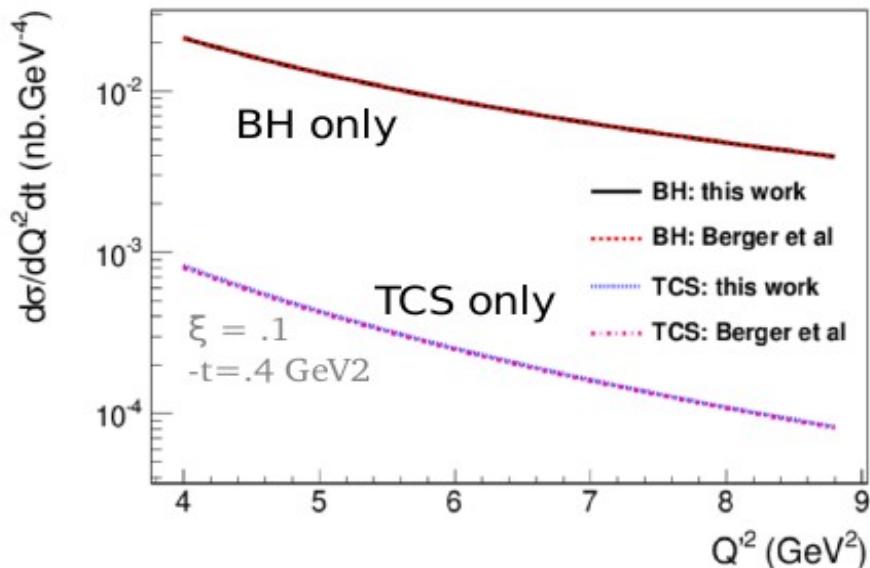
## cross sections vs $Q^2$ and vs $t$

integrated over decay angles

$$\theta \in [45^\circ, 135^\circ]$$

$$\Phi \in [0^\circ, 360^\circ]$$

- BH is always 1 or 2 order of magnitude larger than TCS
- order of pb



Comparison to the pioneering theoretical work

Berger, Diehl, Pire, E.P.J. C23 (2002) 675

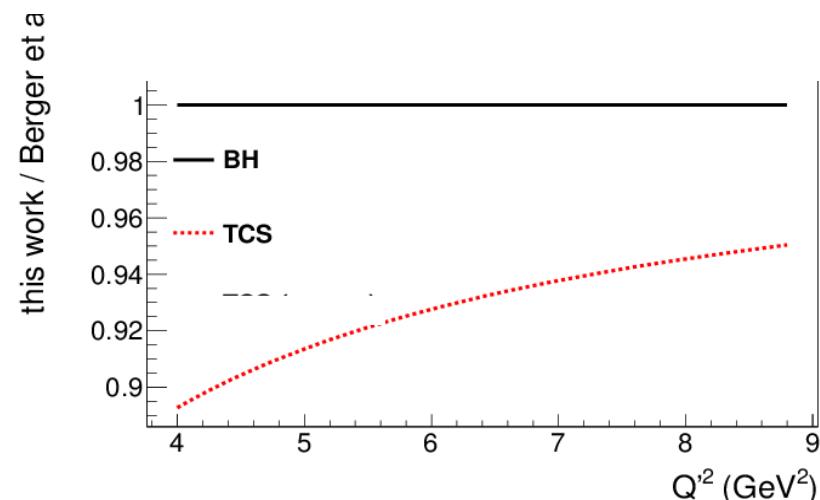


Bethe-Heitler:  $\approx$  equal

TCS: few % at low  $Q^2$

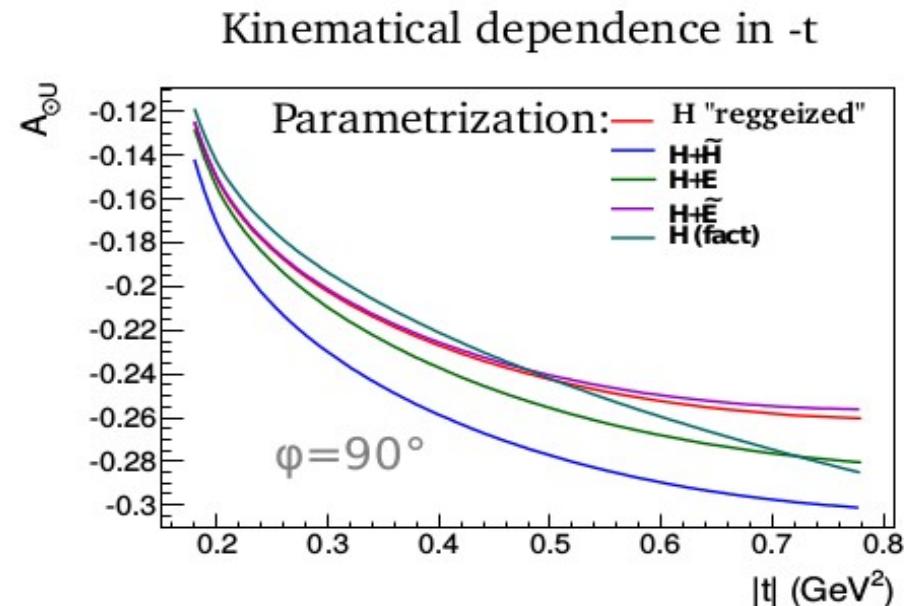
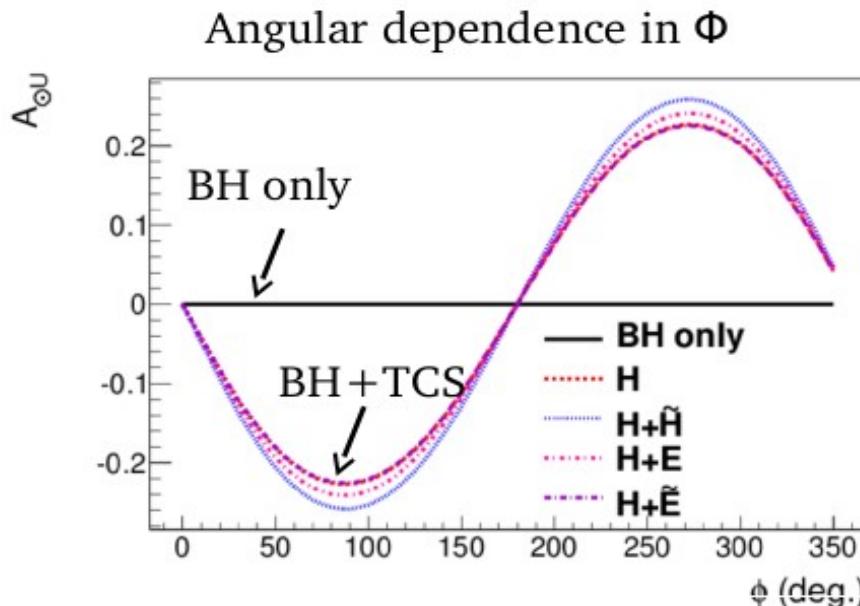
=> we waived some  $t/Q^2$  approximations

(higher twist corrections, gauge invariance)



# Beam Spin Asymmetries (BSA)

## Circularly Polarized



$A_{\odot U} \propto$  **imaginary part of amplitudes**  $\Rightarrow A_{\odot U} = 0$  for Bethe-Heitler

Asymmetry  $\approx 20\%$

This observable : mostly sensitive to  $H$  and  $\tilde{H}$

**$\approx 20\%$  asymmetry coming from interference  
BH x TCS and sensitive to GPDs**

Hall B CLAS12 proposal  
E12-12-01 PAC39 (2012)

Hall A SoLID run group proposal  
E12-12-006A PAC43 (2015)

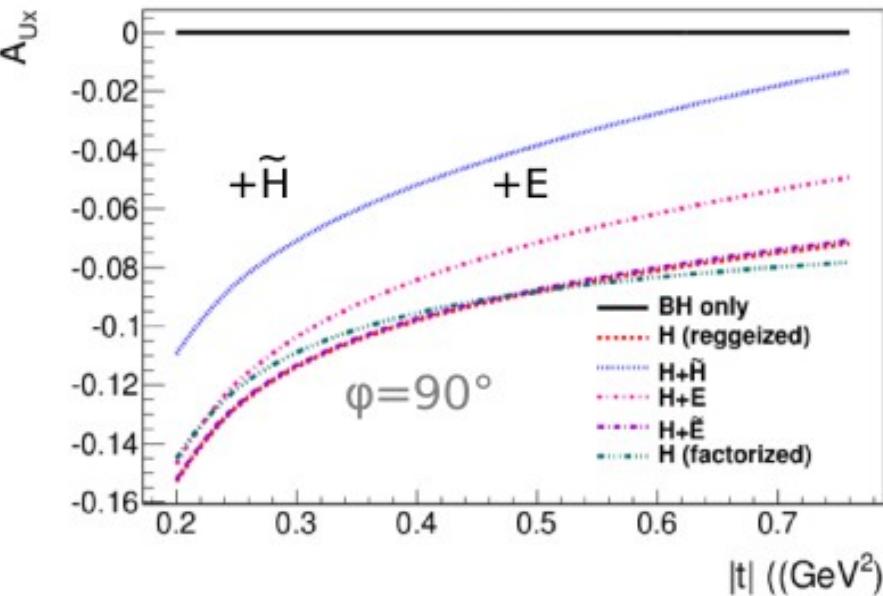
**Linearly Polarized** + sensitive to the real part of amplitudes  
- more difficult experimentaly

$$\xi=0.2, Q'^2=7 \text{ GeV}^2, -t=0.4 \text{ GeV}^2, \theta \in [45^\circ, 135^\circ]$$

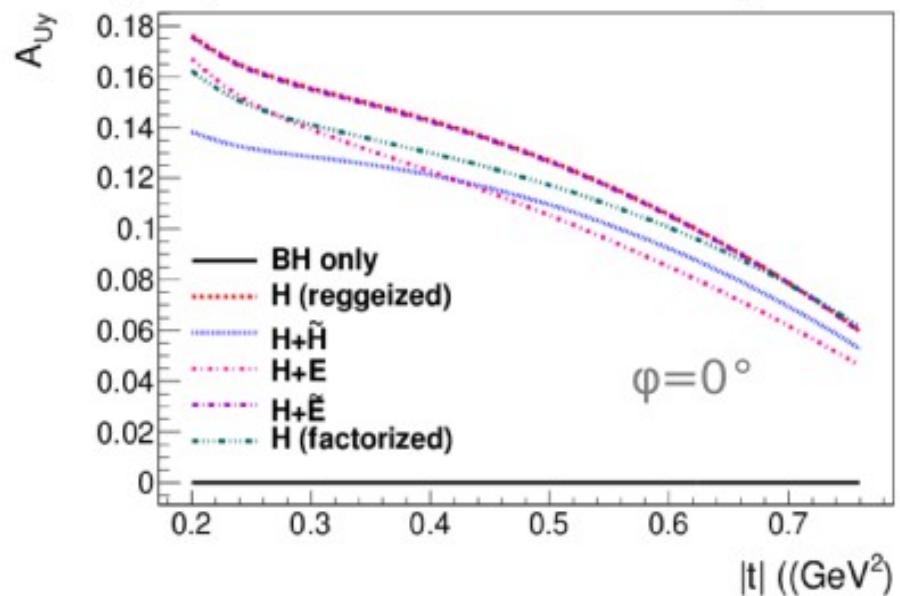
# Target Spin Asymmetries (TSA)

## Transversally polarized target asymmetries vs $|t|$

“in hadronic plane”



“perpendicular to hadronic plane”



- Im part of amplitudes  $\Rightarrow A_{U_i} [\text{BH}] = 0$
- Sensitive to  $H, \tilde{H}, E$  10% to 20% asymmetries

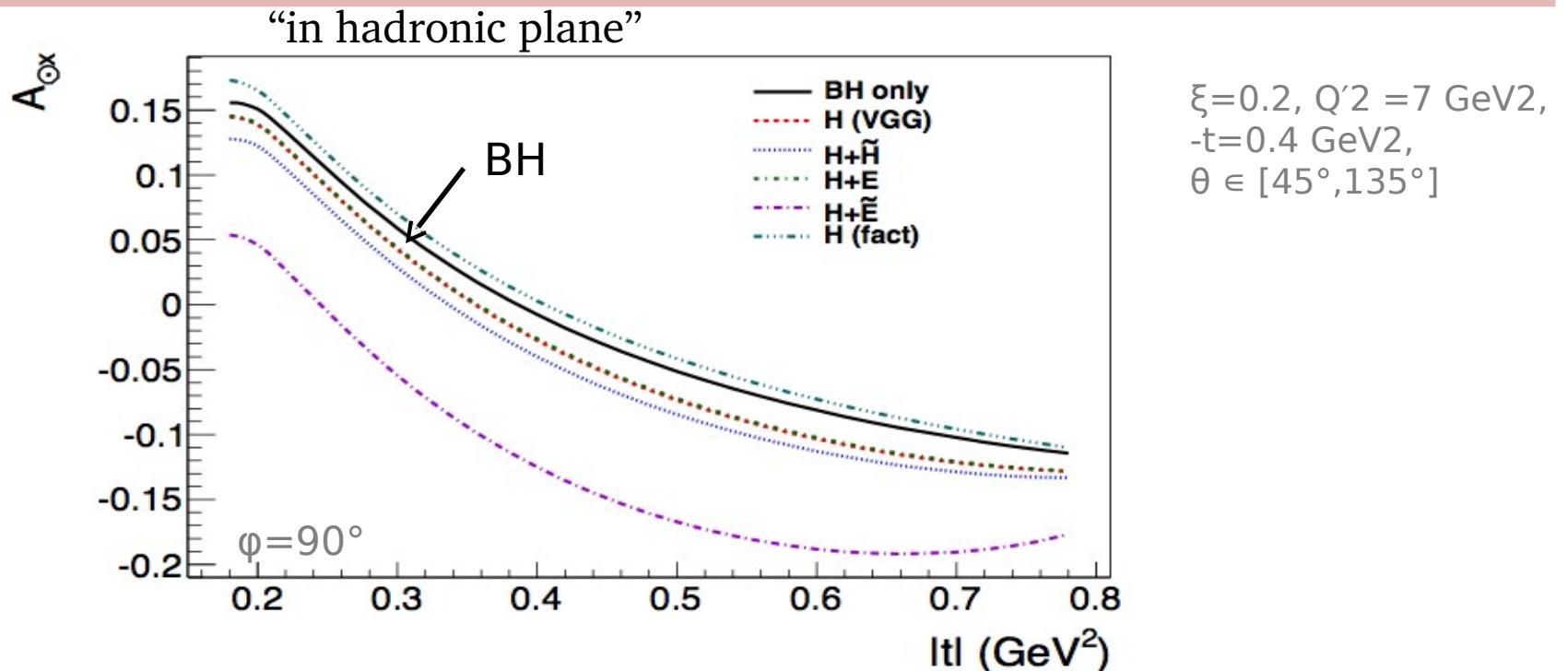
Hall C LOI 12-15-007  
PAC43 (2015)

→ Spin physics with GPD E...

$$\xi=0.2, Q'^2 = 7 \text{ GeV}^2, -t=0.4 \text{ GeV}^2, \theta \in [45^\circ, 135^\circ]$$

# Double Spin Asymmetries (BTSA)

Circularly polarized beam and transversally pol. target vs  $|t|$



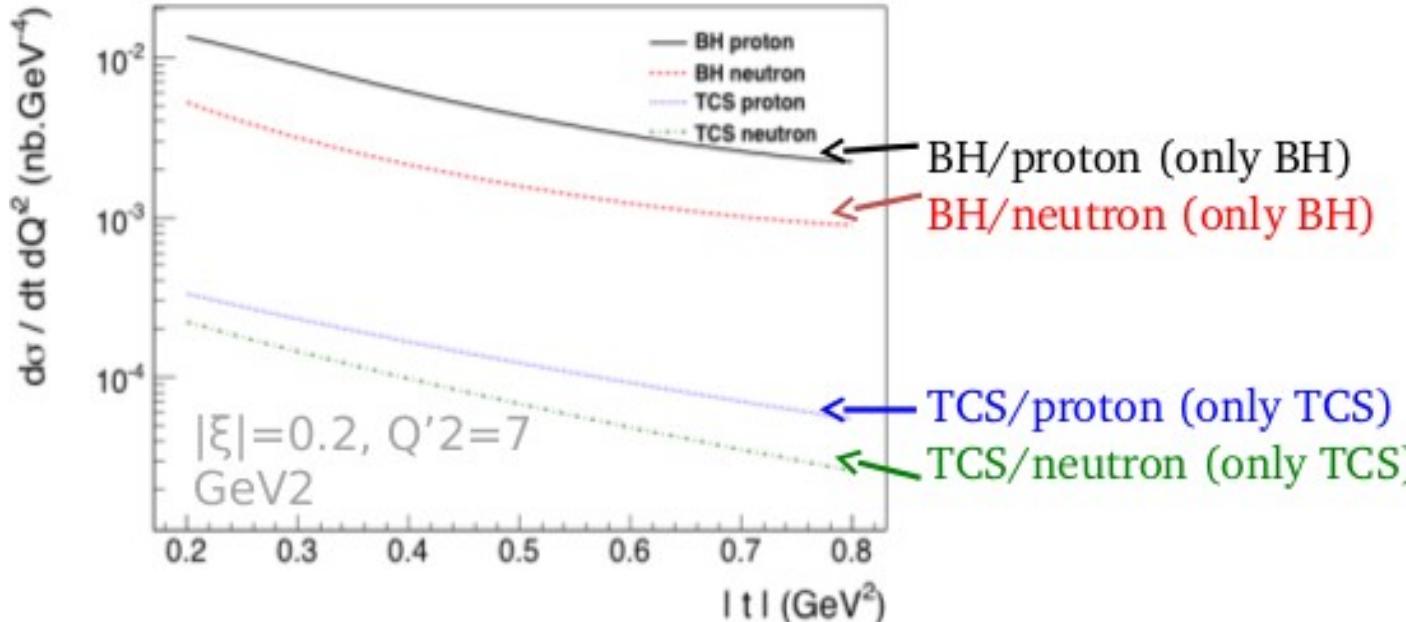
- Very sensitive to the GPDs parameterization
- Sensitive to the real part of amplitudes
- But
  - $A[BH] \neq 0$ , few % deviation from TCS signal
  - Bins in  $\phi$  and  $\theta$  preferable for signal
  - Experimental difficulties (stat...)

Future  
need high luminosity  
+ resolution

BTSA with linearly polarized beam : sensitive to imaginary part of amplitudes

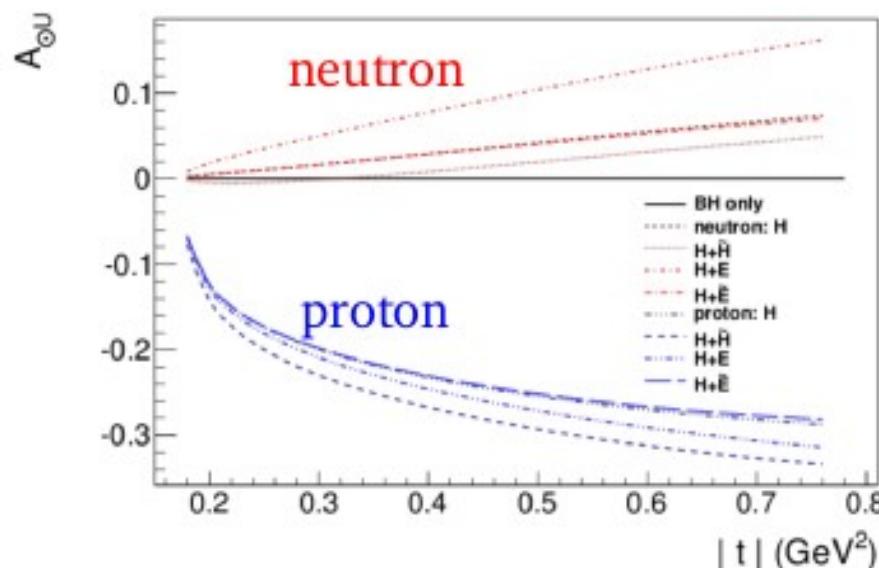
# Comparisons : TCS off the proton and off the neutron

unpolarized cross sections TCS and BH



Future

JLab Hall A  
SoLID  ${}^3\text{He}$  target ?



⇒ TCS off neutron is measurable  
but is more difficult experimentaly  
⇒ Asymmetries ≈ same  $\Phi$  and  $t$   
dependancies and same magnitudes

- flavor separation : u and d quarks
- GPD E (next slide) =>  $J_i$  sum rule

# Spin physics with TCS off the proton and off the neutron

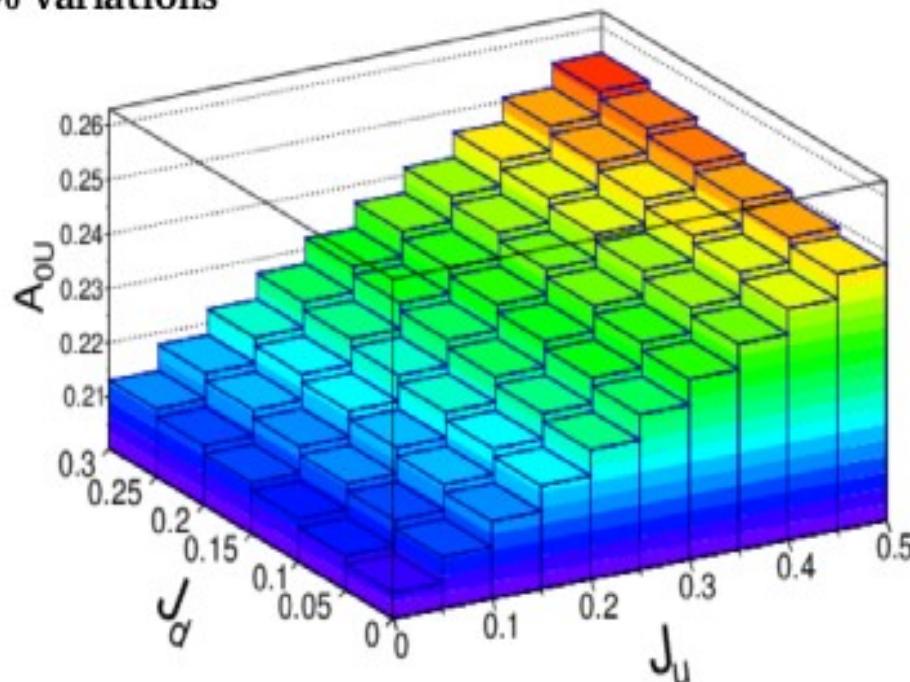
Sensitivity to GPD E in BSA:

BSA as a function of u and d quark angular momenta  $J_u$  and  $J_d$  (max amplitude)

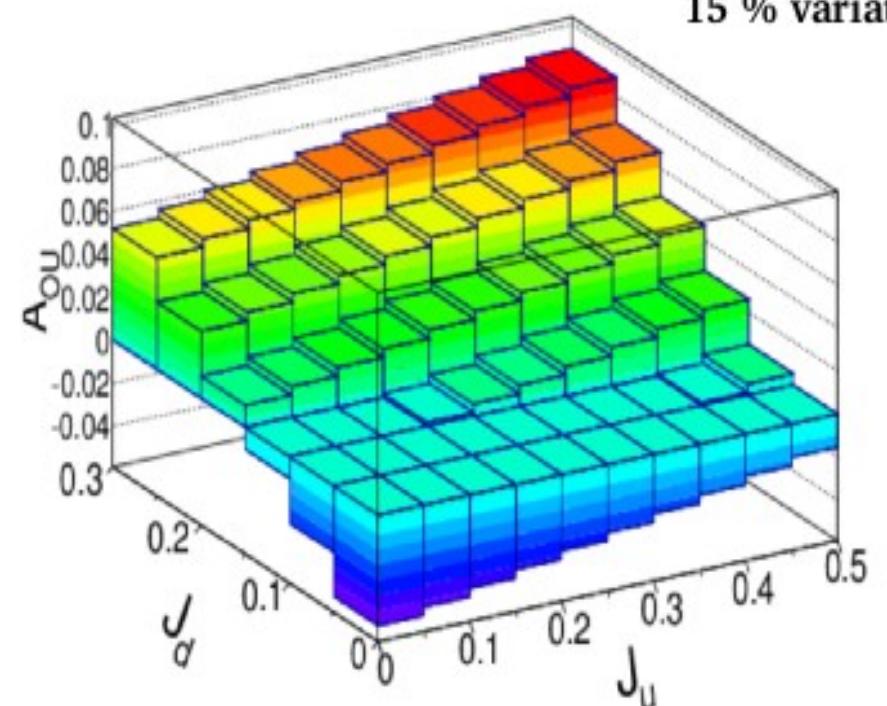
BSA(-t=.4 GeV<sup>2</sup>) vs  $J_u$  and  $J_d$  for **proton**:

for **neutron**:

5 % variations



15 % variations

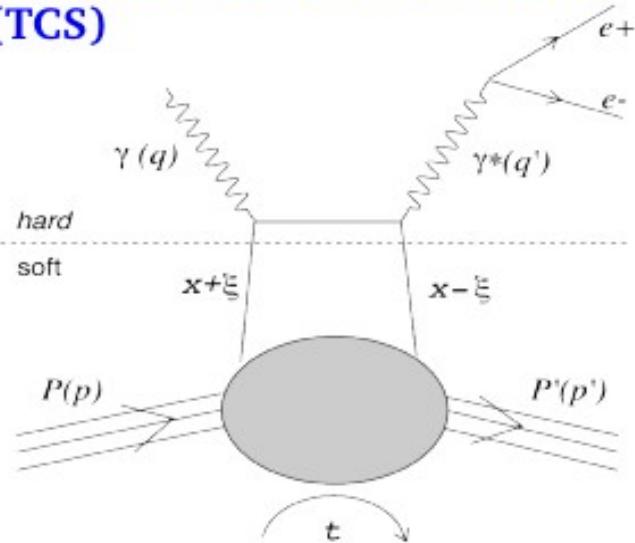


Stronger sensitivity to  $J_u$  and  $J_d$  + change of sign for BSA off the neutron

→ Ji sum rule... studies of angular momenta of quarks

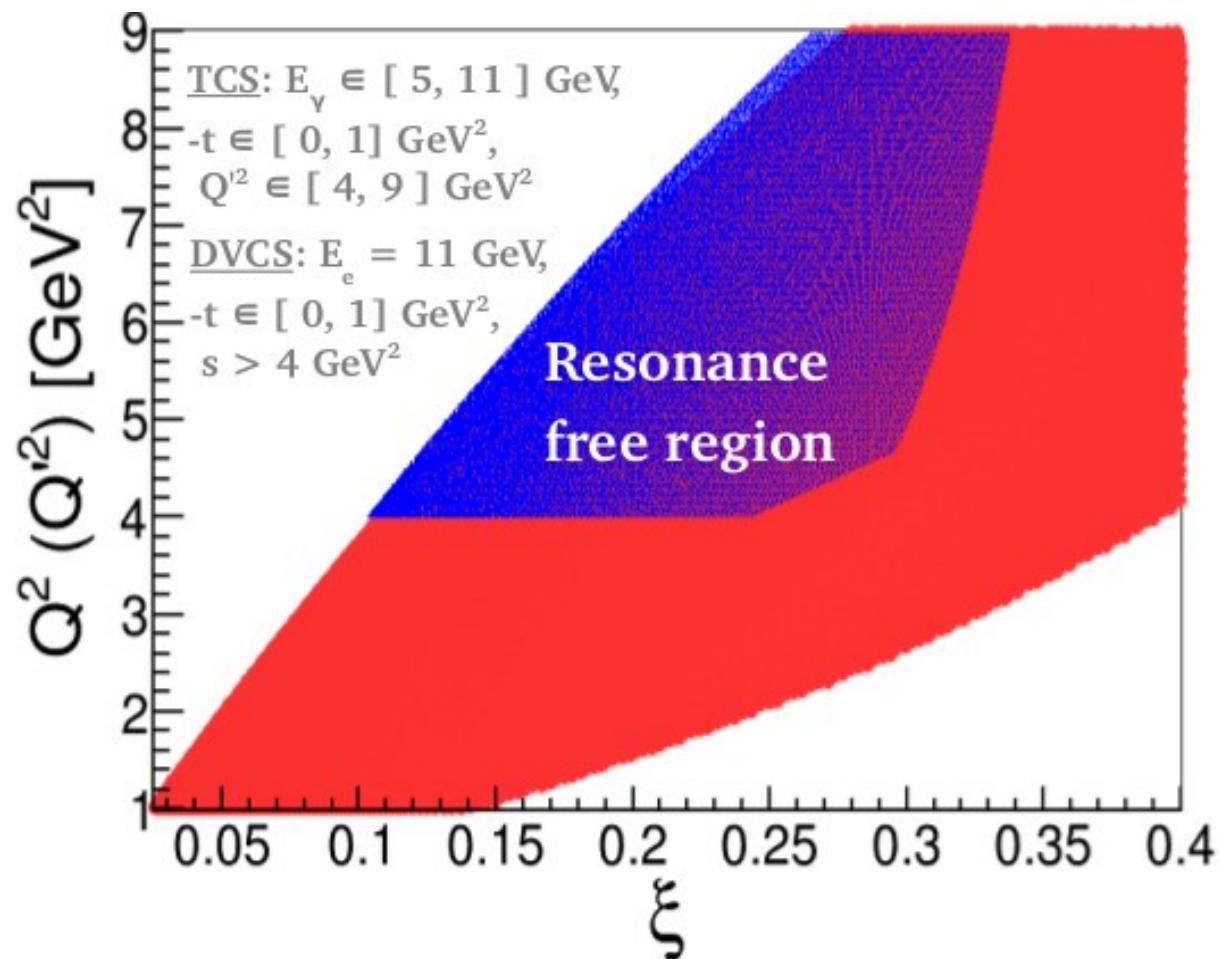
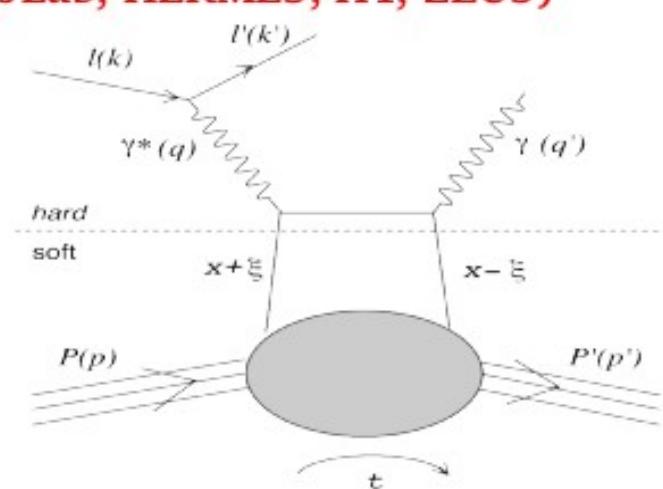
# Fits of DVCS and of TCS : phase space for JLab at 12 GeV

## Timelike Compton Scattering (TCS)



## Deeply Virtual Compton Scattering (DVCS)

Measurements already published  
(JLab, HERMES, H1, ZEUS)



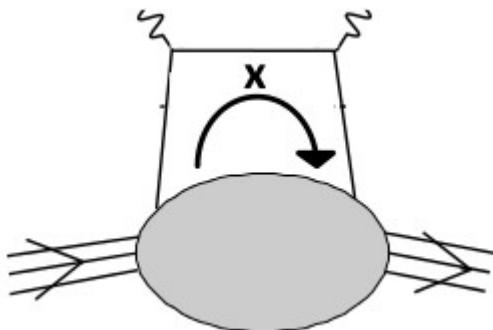
Interest of TCS and DVCS in parallel :

- Universality of GPDs
- Complementary observables
- Higher twist and higher order effects

See talk of M. Guidal

# Accessing GPDs with Compton Form Factors

$\xi, t$  = measurable  
 $x$  = loop  
 $x \pm \xi$  = propagator

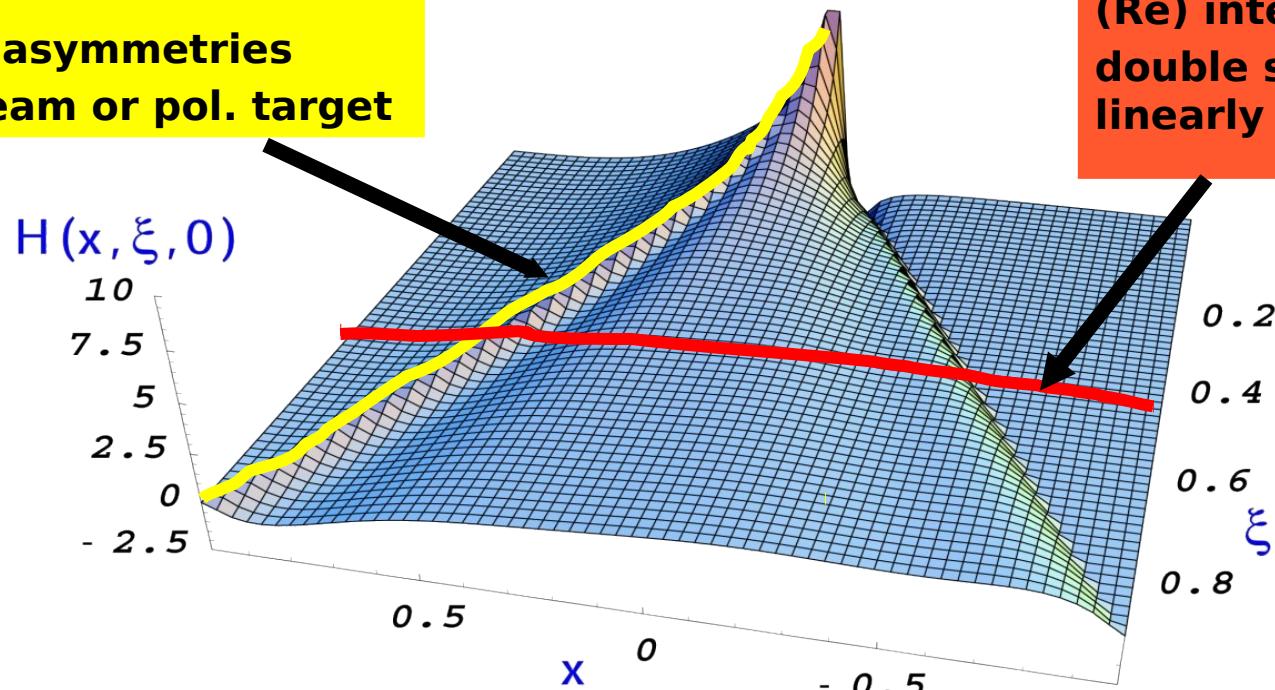


$$T^{DVCS} \sim \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi + i\epsilon} dx + \dots \sim \underbrace{P \int_{-1}^{+1} \frac{H(x, \xi, t)}{x \pm \xi} dx}_{\text{Re } (\mathcal{H})} - i\pi H(\pm \xi, \xi, t) + \dots \underbrace{i\pi H(\pm \xi, \xi, t)}_{\text{Im } (\mathcal{H})}$$

Compton Form Factor (CFF)

CFF are extracted from different observables,  
at different kinematical points:

(Im,  $x = \pm \xi$ )  
single spin asymmetries  
circ. pol. beam or pol. target



(Re) integrate GPD over x  
double spin asymmetries or  
linearly polarized beam

# Fits of GPDs with TCS

## Set of results (uncertainties)

Compton Form  
Factors (CFFs)

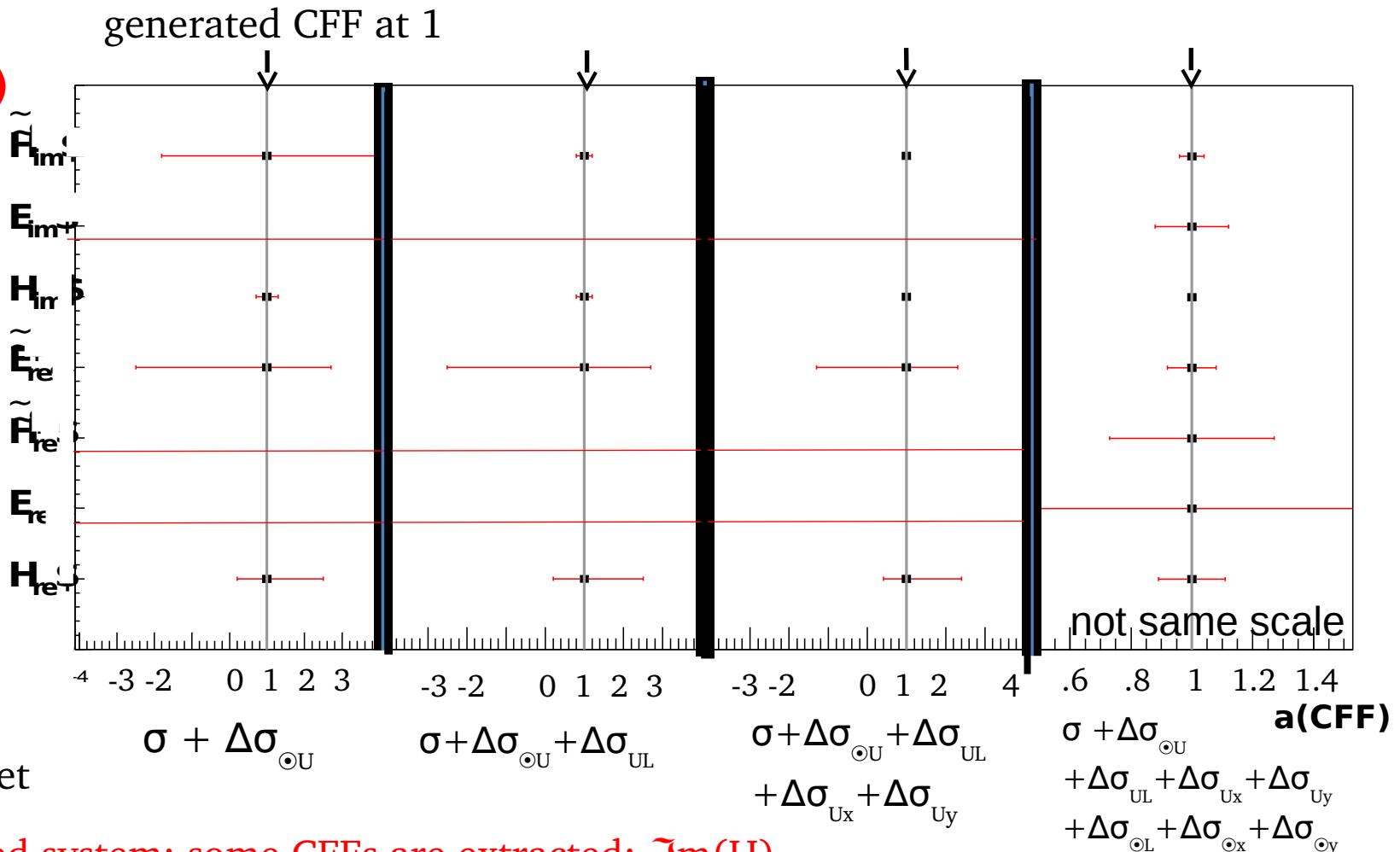
simulations;  
without smearing  
 $\delta\sigma = 5\%$ ,  $\delta\Delta\sigma = 2\%$

generated "CFF" = 1

$\xi=0.2$ ,  $Q^2 = 7 \text{ GeV}^2$ ,  
 $-t=0.4 \text{ GeV}^2$ ,  $\theta = 90^\circ$

Observables:  
With polarized  
beam and/or target

- underconstrained system: some CFFs are extracted:  $\Im(\mathcal{H})$
- 8 independant observables, 7 CFFs: all CFFs are extracted
- single spin asymmetries  $\propto \text{Im}T \implies \text{Im}(\text{CFFs})$  are extracted with smaller error bars
- compared to DVCS : more difficult with TCS, but complementary



CFFs can be extracted from TCS fits assuming 5% uncertainties on observables

# Summary

- Nucleon "3D" imaging with Generalized Parton Distributions
- TCS signal: through interference with BH, with beam and/or target spin asymmetries
- 2 accepted proposals for JLab at 12 GeV: CLAS12 (Hall B) and SoLID (Hall A), 1 LOI NPS (Hall C)
- Other experimental perspectives : linearly polarized beam, real photon beam, double spin asymmetries, TCS off the neutron..
- Fits : CFFs can be extracted, interests for DVCS+TCS combined fits...