

Workshop on Top physics at the LC 2015

Higher order EW corrections to $t\bar{t}$ with GRACE

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Outline

1. Introduction

- ❖ Motivation

2. Method

- ❖ GRACE system

- ❖ Polarization

3. Results

- ❖ Cross section

- ❖ Angular distribution

- ❖ Forward and backward asymmetry

- ❖ The top quark decay

4. Summary

ILC e^+e^-

- ✿ The main physics goals in ILC are top physics, Higgs physics and Physics beyond the Standard Model.
- ✿ Precise measurement of top mass, width, top Yukawa coupling are crucial.

Motivation

- ✿ Loop calculation with the target <1% precision according to the data of experiment.
- ✿ Polarization: (e^- : 80%), (e^+ : 30%)
- ✿ Full electroweak loop correction for $t\bar{t}$ pair with polarization.

GRACE

- ✿ GRACE is a system to calculate tree and loop cross sections automatically with beam-polarization based on SM and MSSM.
- ✿ GRACE-Loop system can calculate:

SM processes:

$$e^+ e^- \rightarrow t\bar{t}\gamma$$

Loop Eur. Phys. J. C **73**, 2400 (2013)

$$e^+ e^- \rightarrow e^+ e^- \gamma$$

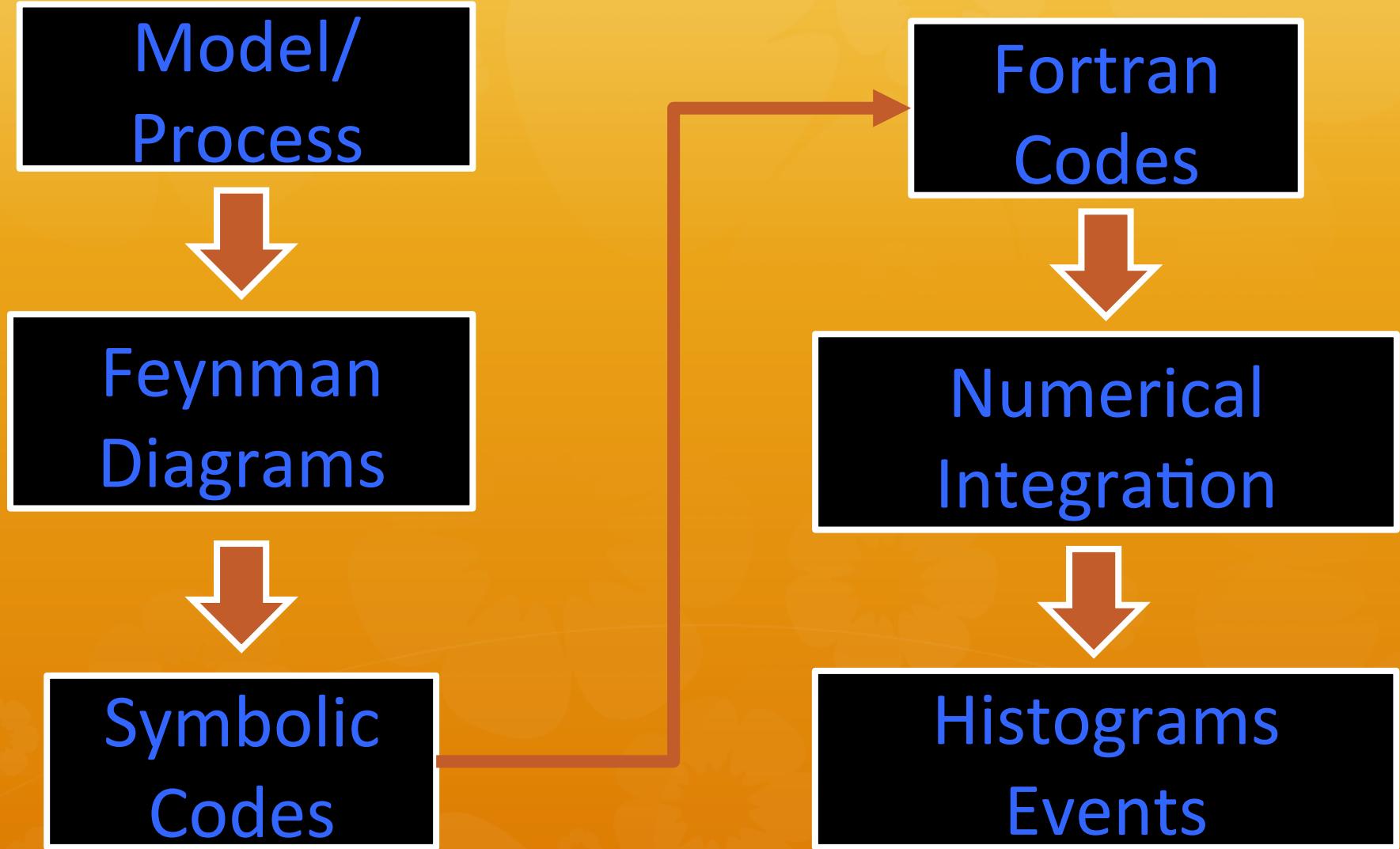
Loop Phys. Lett. **B740**, 192 (2014)

MSSM processes:

$$e^- e^\pm \rightarrow e^- e^\pm \tilde{\chi}_{1,2}^0 \tilde{\chi}_{1,2}^0$$

Tree 20910D, JPS September, 2014

What is GRACE?



- ✿ The renormalization has been carried out with the on-shell condition of the Kyoto scheme¹.

- ✿ The non-linear gauge fixing Lagrangian condition²:

$$\begin{aligned}\mathcal{L}_{GF} = & -\frac{1}{\xi_W} \left| (\partial_\mu - ie\tilde{\alpha}A_\mu - igc_W\tilde{\beta}Z_\mu)W^{\mu+} + \xi_W \frac{g}{2}(v + \tilde{\delta}H + i\tilde{\kappa}\chi_3)\chi^+ \right|^2 \\ & -\frac{1}{2\xi_Z}(\partial.Z + \xi_Z \frac{g}{2c_W}(v + \tilde{\epsilon}H)\chi_3)^2 - \frac{1}{2\xi_A}(\partial.A)^2.\end{aligned}$$

- ✿ The results must be independent of non-linear gauge parameters ($\tilde{\alpha}, \tilde{\beta}, \tilde{\delta}, \tilde{\kappa}, \tilde{\epsilon}$).

¹ K. Aoki et al, Suppl. Prog. Theor. Phys. **73 (1982) 1.**

² Phys, Rept. **430**, 117 (2006)

System checking

The total cross section is

$$\sigma_{full} = \sigma_{Tree} + \sigma_{Loop}(C_{UV}, \tilde{\alpha}, \tilde{\beta}, \tilde{\delta}, \tilde{\epsilon}, \tilde{\kappa}) \\ + \sigma_{Tree}\delta_{soft}(\lambda, E_\gamma < k_c) + \sigma_{Hard}(k_c)$$

Soft photon: $E_\gamma < k_c$

Hard photon: $E_\gamma \geq k_c$

Using the numerical check by GRACE, we confirmed:

- ✿ C_{UV} (Ultra violet coefficient) independence
- ✿ Photon mass (λ) independence
- ✿ Gauge invariance
- ✿ k_c independence

Polarization

Polarization: (e^- : 80%), (e^+ : 30%)

- ✿ e^-e^+ : non polarization
- ✿ e^-e^+ : left-right polarization
- ✿ e^-e^+ : right-left polarization

To include polarization effect, we multiply the projection operators into amplitude.

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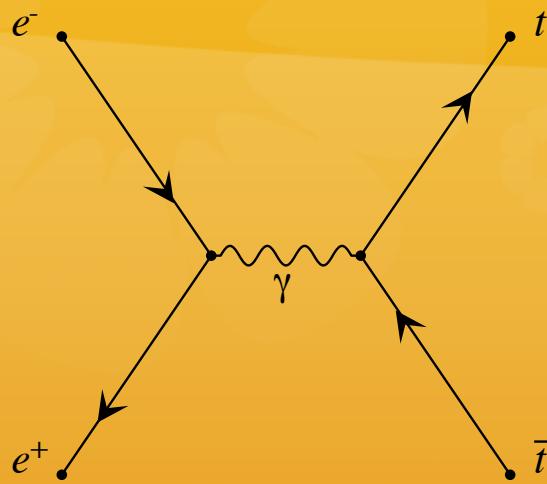


Program size becomes much larger!!!

Parameters

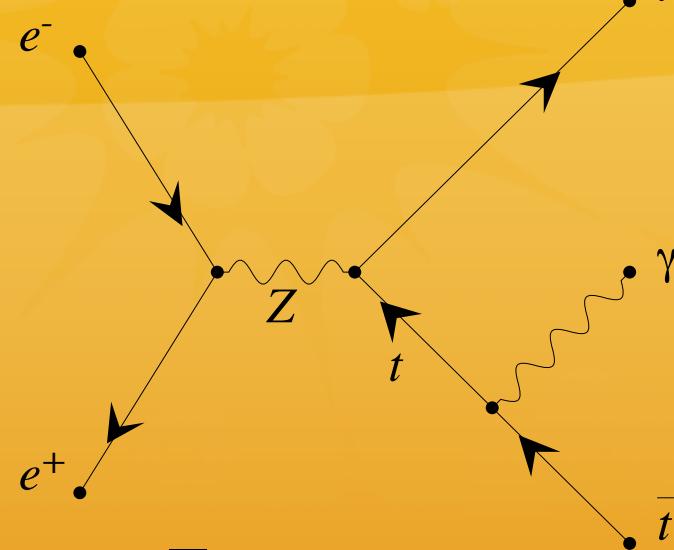
m_Z	91.187 GeV
m_W	80.22 GeV
m_H	126 GeV
m_t	173.5 GeV
α	1/137.0359895
$\sin^2 \theta_W$	$1 - \frac{m_W^2}{m_Z^2}$

$e^- e^+ \rightarrow t\bar{t}$ Tree
(4 diagrams)

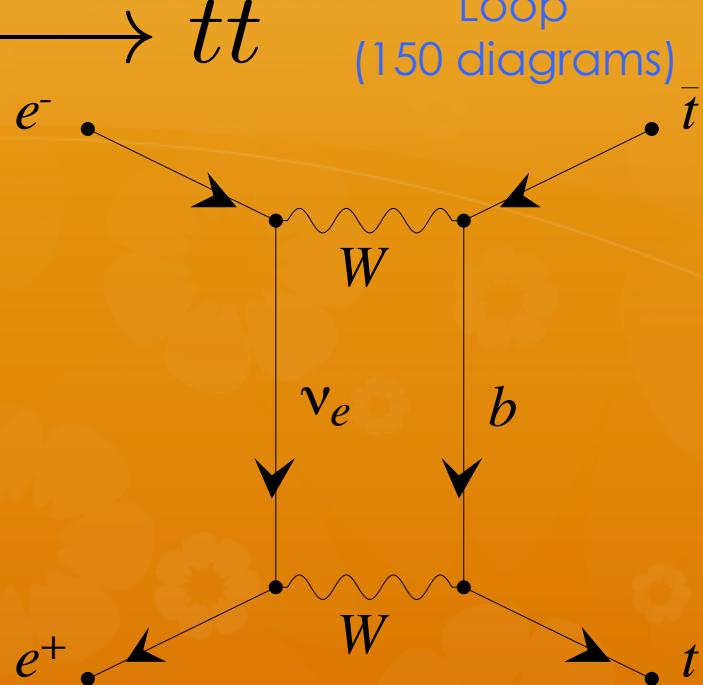
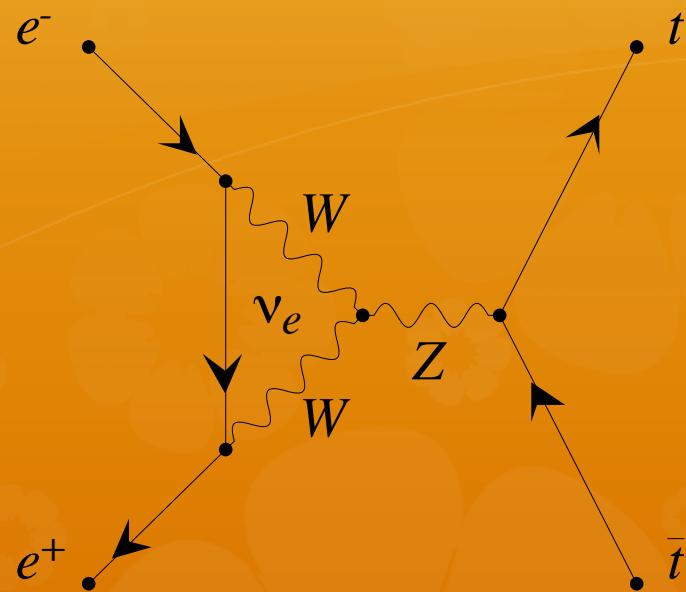


$e^- e^+ \rightarrow t\bar{t}$

Radiation
(16 diagrams)

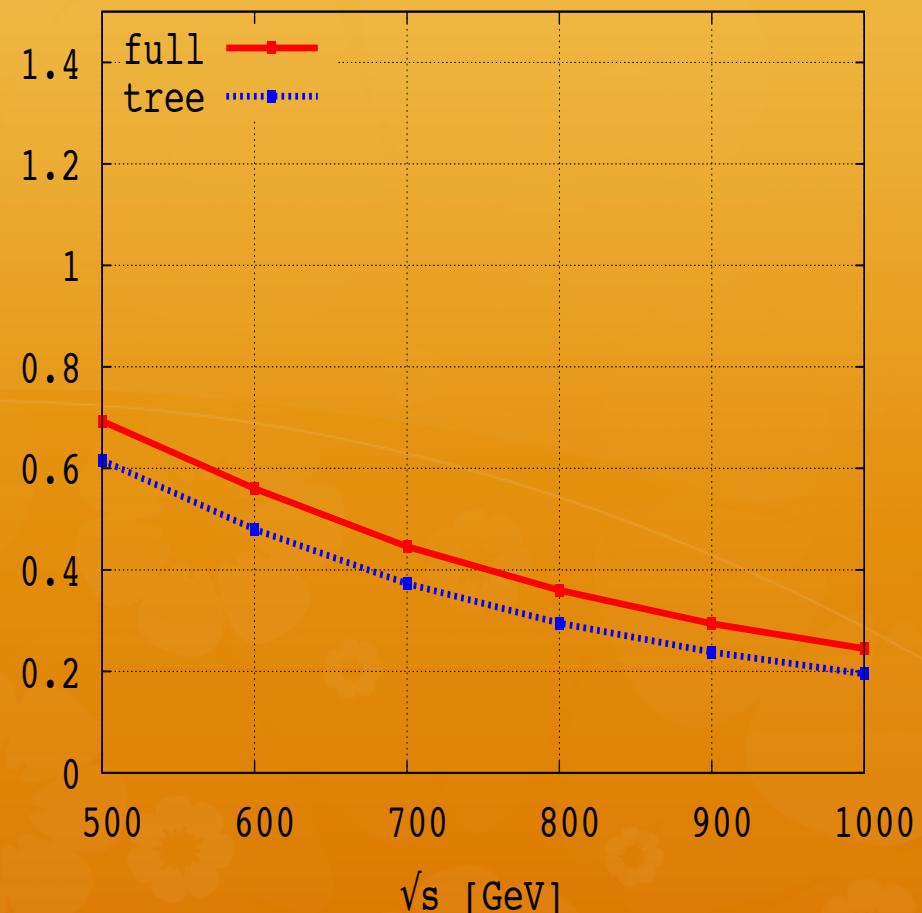
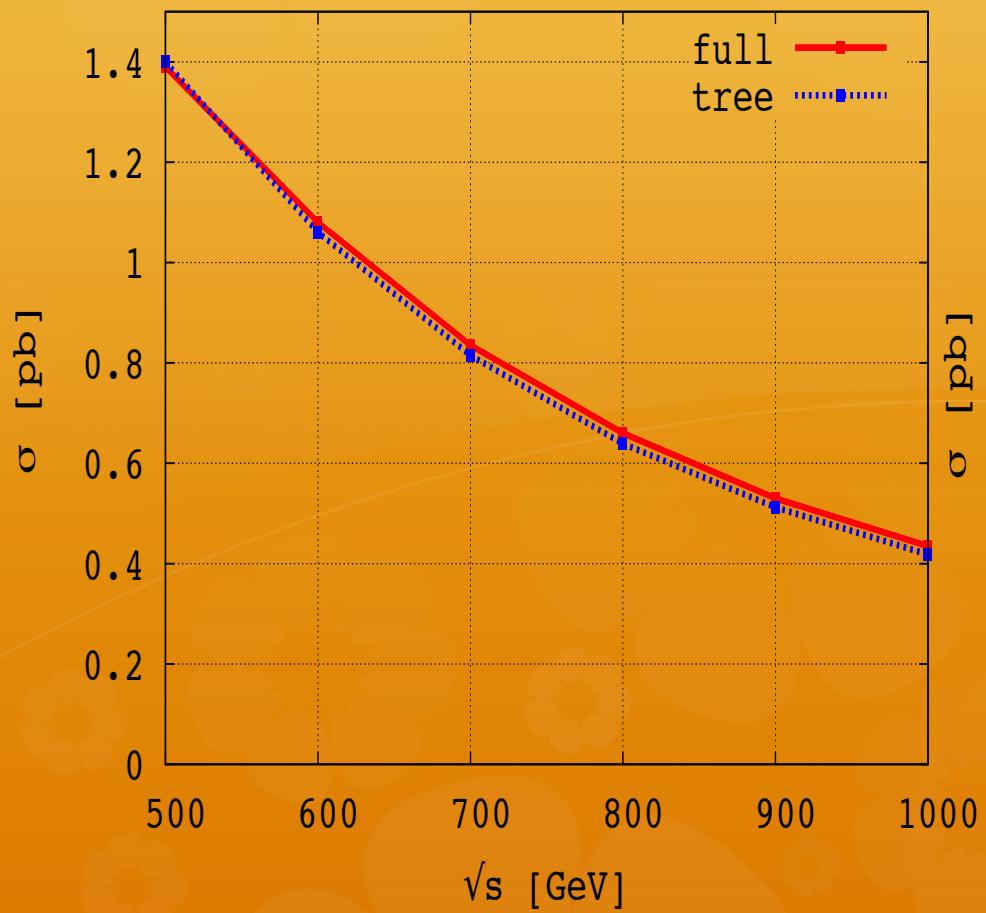


Loop
(150 diagrams)



Results

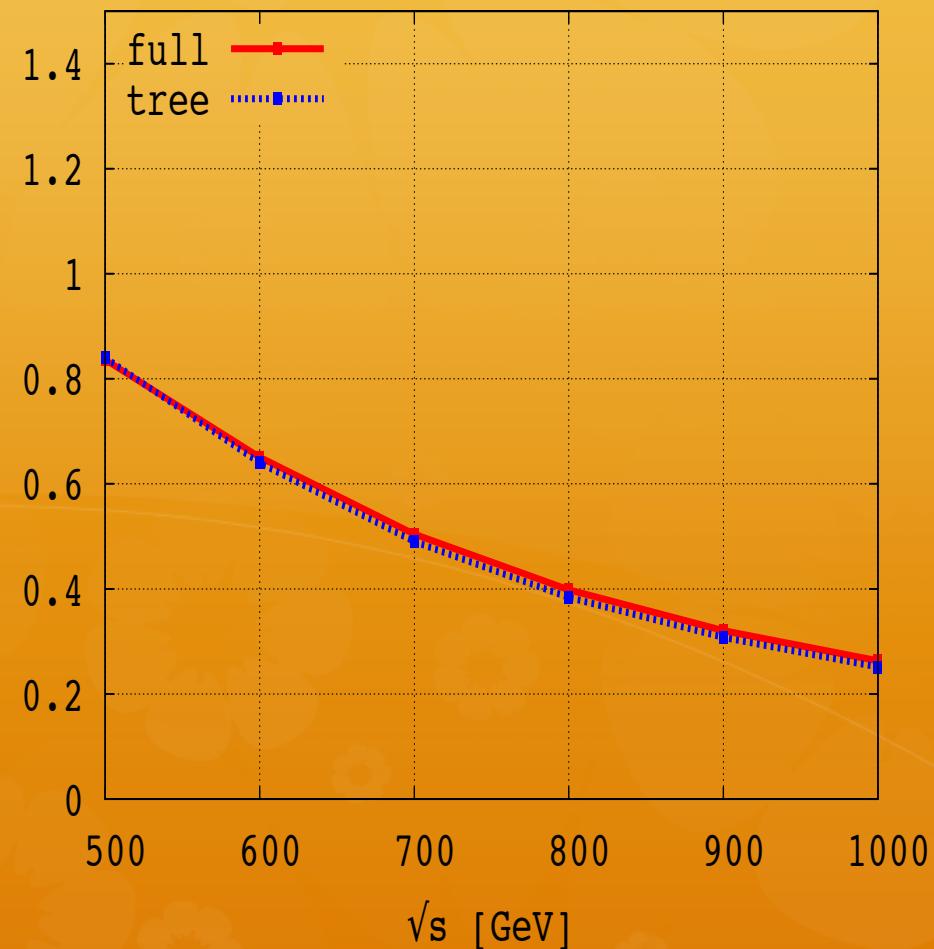
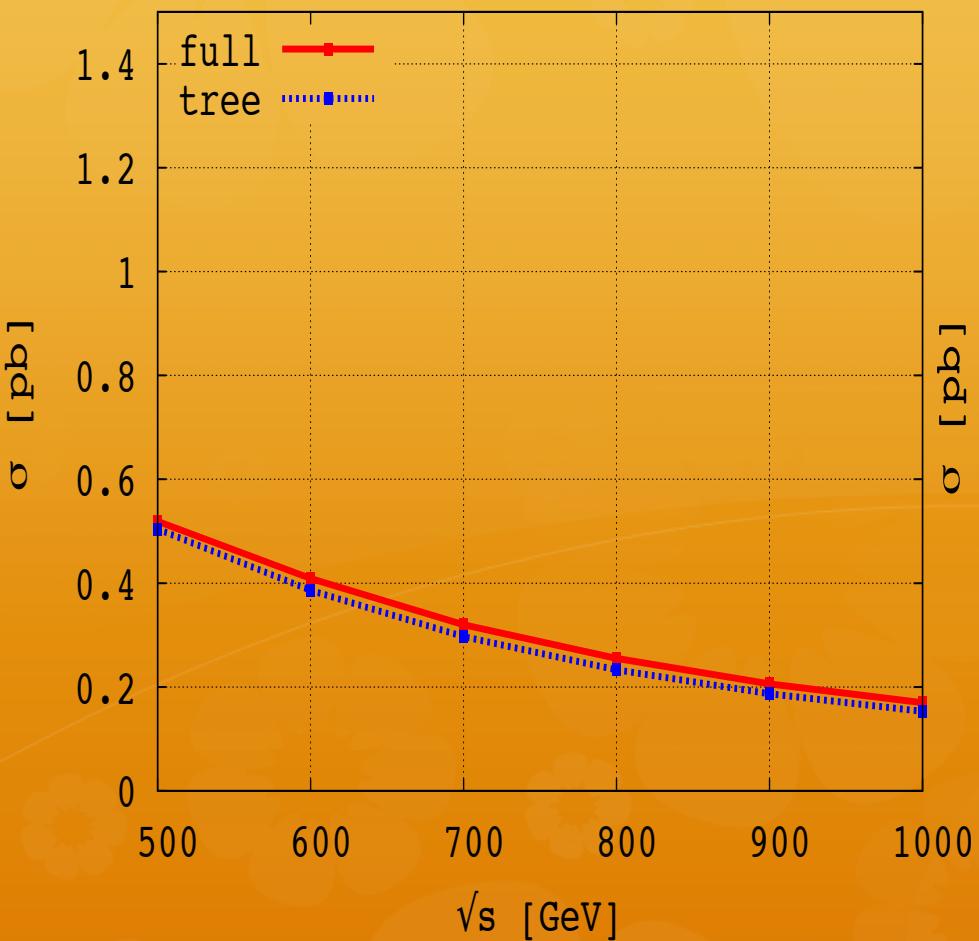
100% INITIAL POLARIZATION



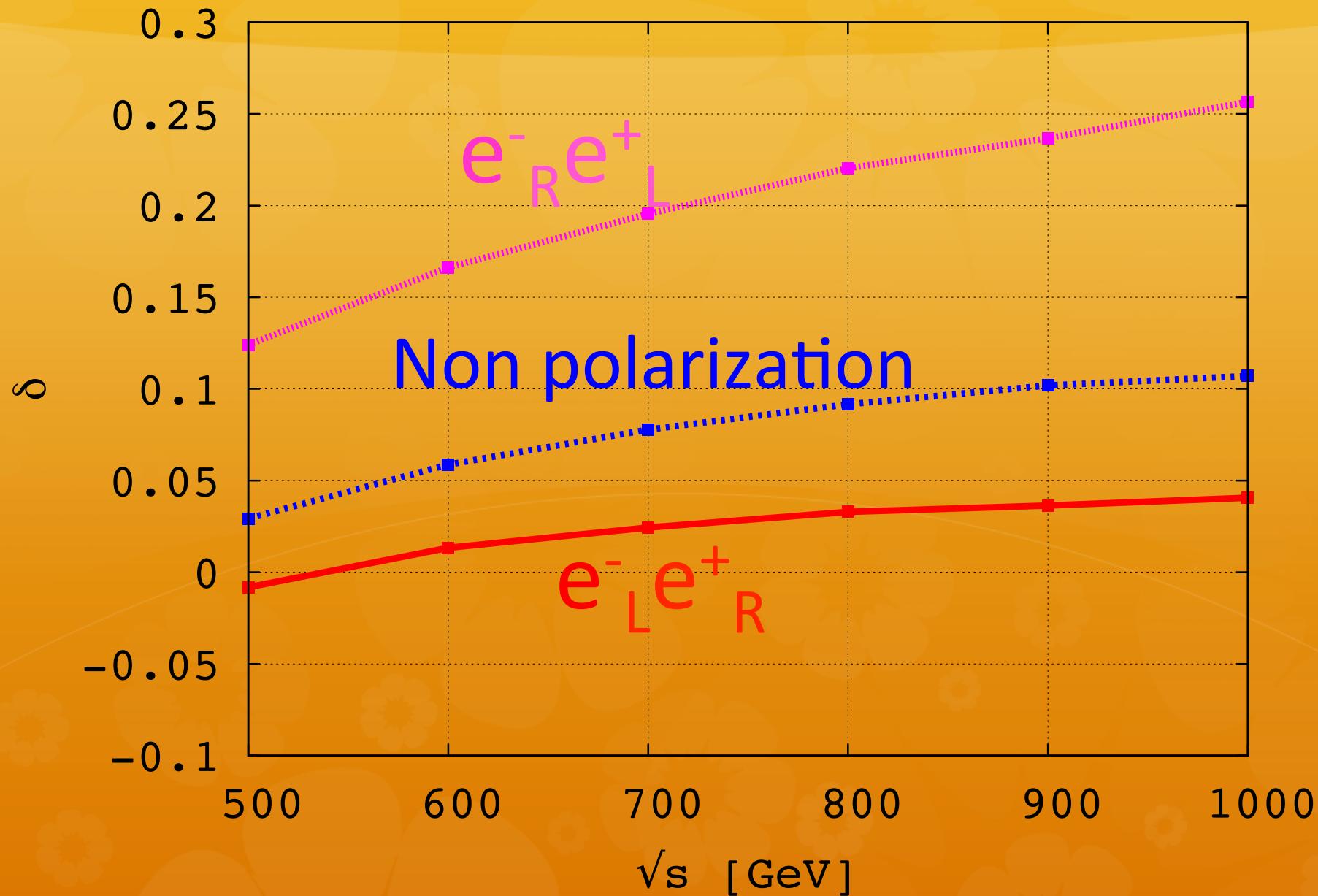
Results

INITIAL POLARIZATION

Non Polarization $e^- 80\%, e^+ 30\%$



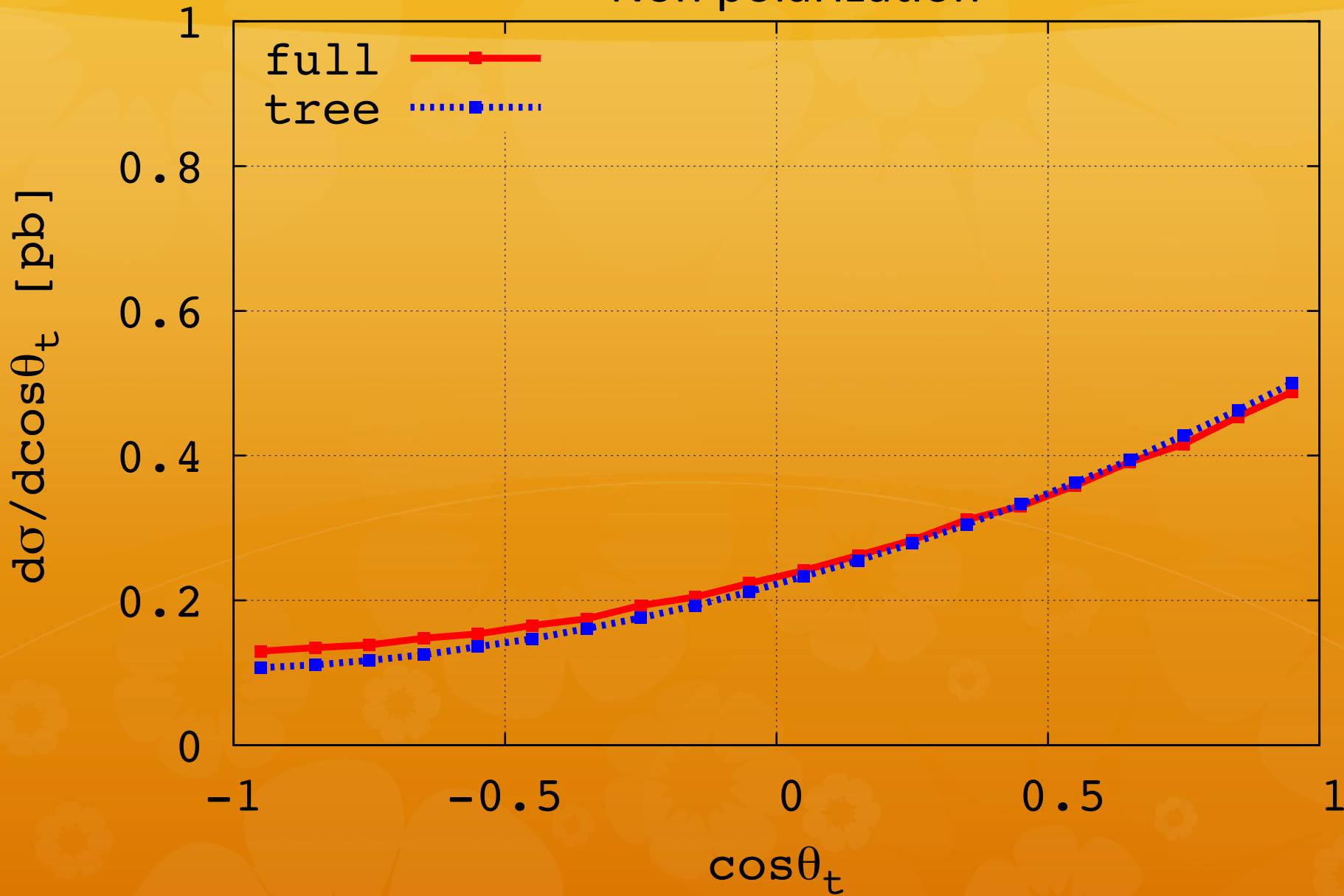
$$\delta = \frac{\sigma_{full} - \sigma_{tree}}{\sigma_{tree}}$$



$\sqrt{s} = 500 \text{ GeV}$

Angular distribution

Non polarization

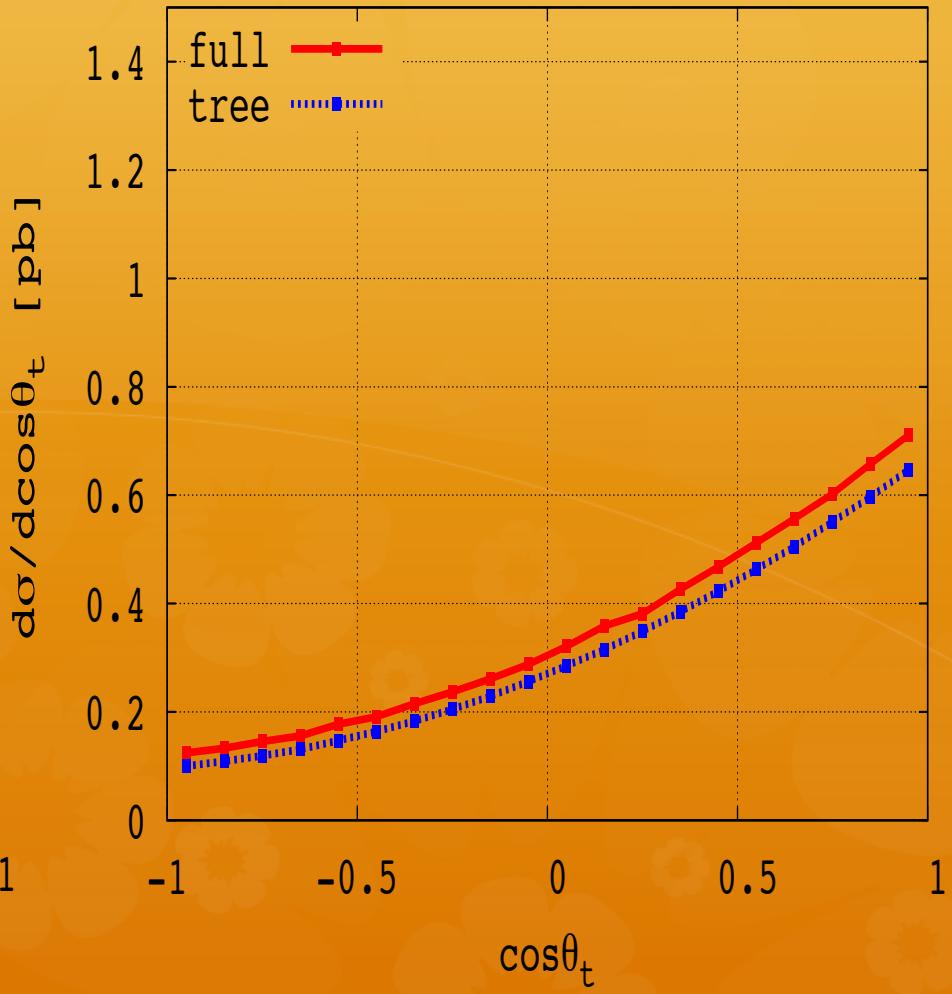
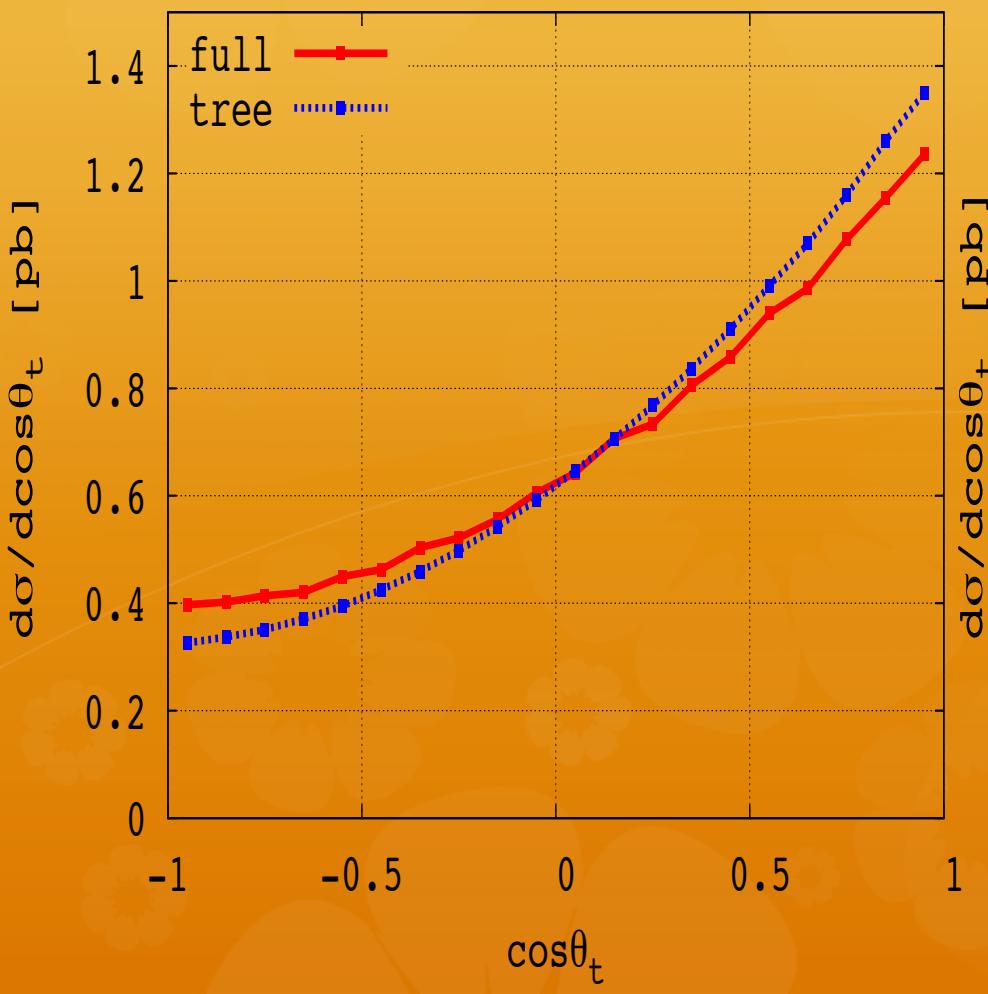


$\sqrt{s} = 500 \text{ GeV}$

Angular distribution

$e^-_L e^+_R \rightarrow t\bar{t}$

$e^-_R e^+_L \rightarrow t\bar{t}$



Forward and Backward Asymmetry

$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B}$$

$$\sigma_F = \int_0^1 \frac{d\sigma}{d \cos \theta_t} d \cos \theta_t \quad \sigma_B = \int_{-1}^0 \frac{d\sigma}{d \cos \theta_t} d \cos \theta_t$$

$e^- e^+ \rightarrow t\bar{t}$	A _{FB} <i>Tree</i>	A _{FB} <i>Full</i>
$e^- e^+$	0.410	0.359
$e_L^- e_R^+$	0.385	0.317
$e_R^- e_L^+$	0.467	0.443
$e_L^- (80\%) e_R^+ (30\%)$	0.388	0.321

The correction of A_{FB} is sizable → we should take it into account at ILC.

Top decay $t \rightarrow b \mu^+ \nu_\mu$

Because the number of Feynman diagrams for 6 body final state $e^- e^+ \rightarrow b\bar{b} \mu^- \mu^+ \nu\bar{\nu}$ is too large.

Full EW direct calculating is almost impossible.



We use “narrow width approximation” method.

Due to the top width is so narrow compared with
the top mass

$\mathcal{O}(\alpha)$ correction

$$\Gamma_t^{Tree} = 1.429 \text{ GeV}$$

$$\Gamma_t^{Loop} = 1.524 \text{ GeV} \quad \leftarrow$$

Full-ELWK corrected
one calculated for a
three-body decay of a
top-quark.

$$m_t = 173.5 \text{ GeV}$$

Narrow width approximation method

Total cross section:

$$\begin{aligned}\sigma &= \frac{1}{flux} \int |\mathcal{M}|^2 d\Omega_N \\ &= \frac{1}{flux} \int \frac{\left| \sum \mathcal{M}_p u_\lambda(q) \bar{u}_\lambda(q) \mathcal{M}_d \right|^2}{(q^2 - m^2)^2 + m^2 \Gamma^2} \frac{dq^2}{2\pi} d\cos\theta_q d\varphi_q d\Omega_m d\Omega_{N-m} \\ &\simeq \frac{1}{flux} \sum_\lambda \int \left| \tilde{\mathcal{M}}_p^\lambda \right|^2 d\cos\theta_q d\varphi_q d\Omega_{N-m} \int \left| \tilde{\mathcal{M}}_d^\lambda \right|^2 d\Omega_m \\ &\quad \int_{-\infty}^{+\infty} \frac{1}{(q^2 - m^2)^2 + m^2 \Gamma^2} \frac{dq^2}{2\pi} \quad \text{On-shell condition}\end{aligned}$$

u_λ : spinor of decayed particle

q_μ : momentum (off-shell) of decayed particle

λ : spin of decayed particle

\mathcal{M}_p : product amplitude

\mathcal{M}_d : decay amplitude

$$q_0^2 = m^2$$

$$\tilde{\mathcal{M}}_p^\lambda = \mathcal{M}_p u_\lambda(q_0)$$

$$\tilde{\mathcal{M}}_d^\lambda = \mathcal{M}_d u_\lambda(q_0)$$

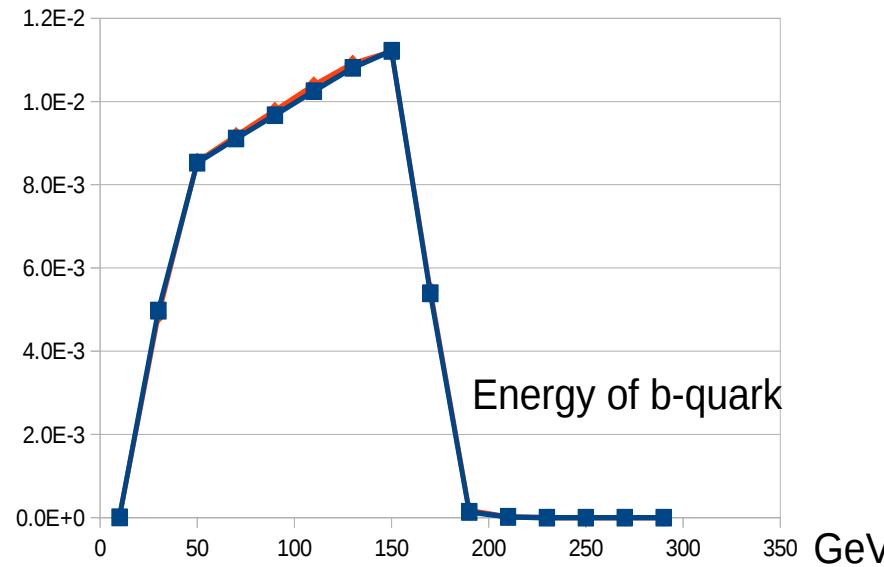
$$\begin{aligned}
I &= \int_{-\infty}^{+\infty} \frac{1}{(q^2 - m^2)^2 + m^2\Gamma^2} dq^2 \\
&= \frac{1}{m^2\Gamma^2} \int_{-\frac{\pi}{2}}^{+\frac{\pi}{2}} m\Gamma d\theta \\
&= \frac{\pi}{m\Gamma}
\end{aligned}$$

$$\begin{aligned}
\int \left| \tilde{\mathcal{M}}_d^\lambda \right|^2 d\Omega_m \int_{-\infty}^{+\infty} \frac{1}{(q^2 - m^2)^2 + m^2\Gamma^2} \frac{dq^2}{2\pi} &= \frac{1}{\Gamma} \frac{1}{2m} \int \left| \tilde{\mathcal{M}}_d^\lambda \right|^2 d\Omega_m \\
&\quad \downarrow \text{GRACE} \\
&= \frac{\Gamma_i^\lambda}{\Gamma} = Br
\end{aligned}$$

σ = *Production* \cdot *Branching ratio*

Test of the Narrow Width Approximation

$d\sigma/dE(\text{pb}/\text{GeV})$

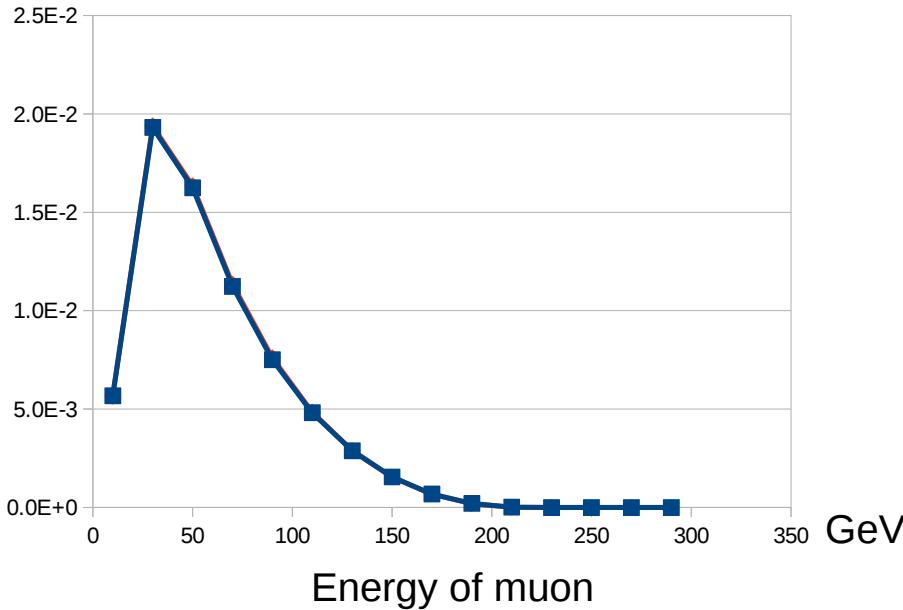


Energy of b-quark

$$e^- e^+ \rightarrow b\bar{b} \mu^- \mu^+ \nu\bar{\nu}$$

Blue : Exact 6-body
Orange : NWA

$d\sigma/dE(\text{pb}/\text{GeV})$



Energy of muon

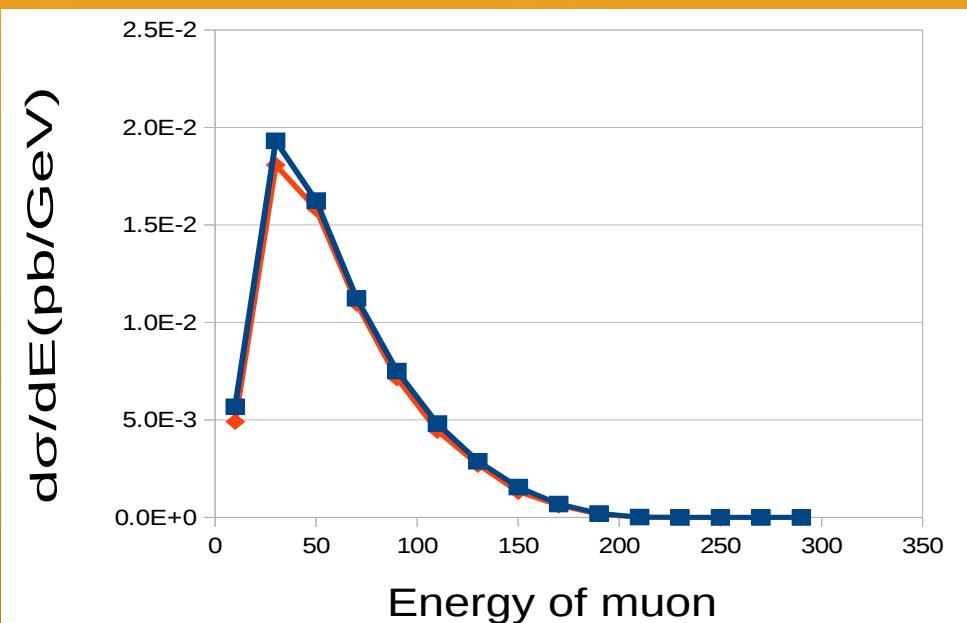
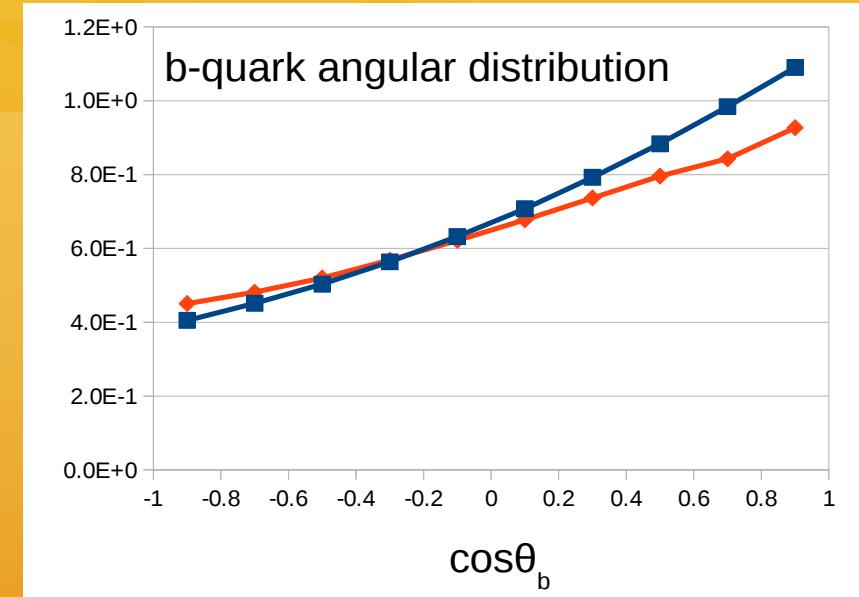
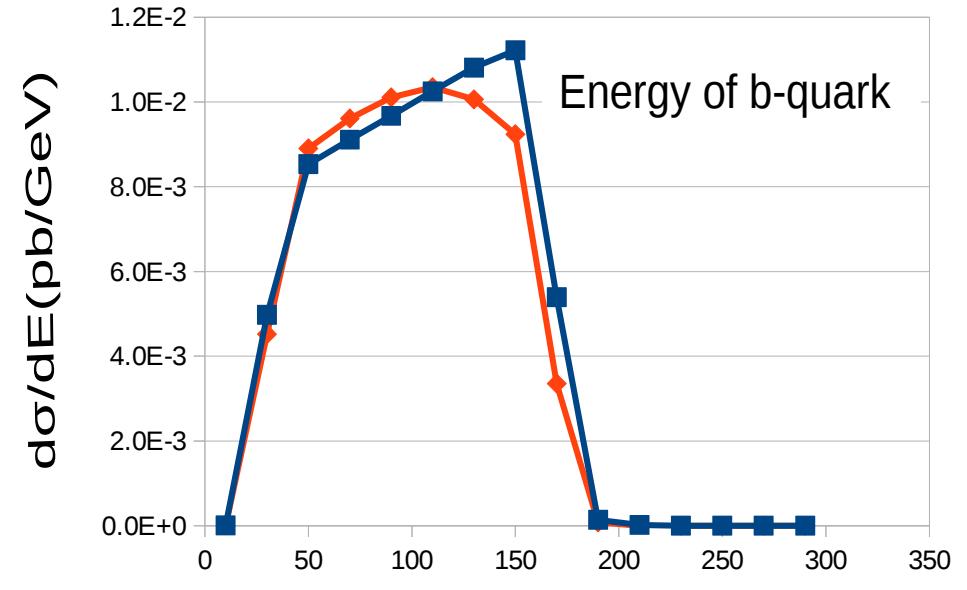
$\sqrt{s}=500\text{GeV}$
Electron : 100% Left
Positron : 100% Right

Total cross section



0.5% accuracy

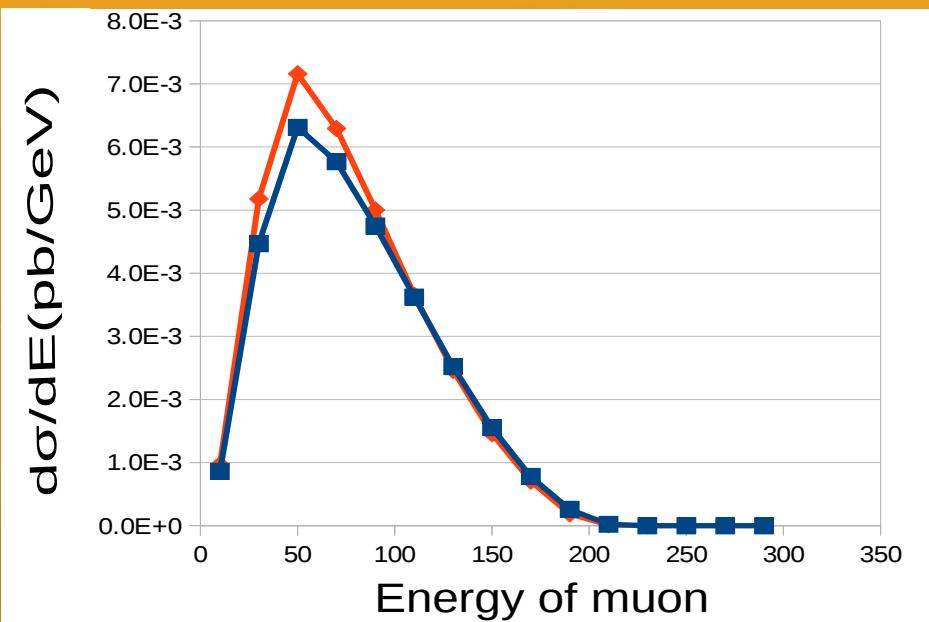
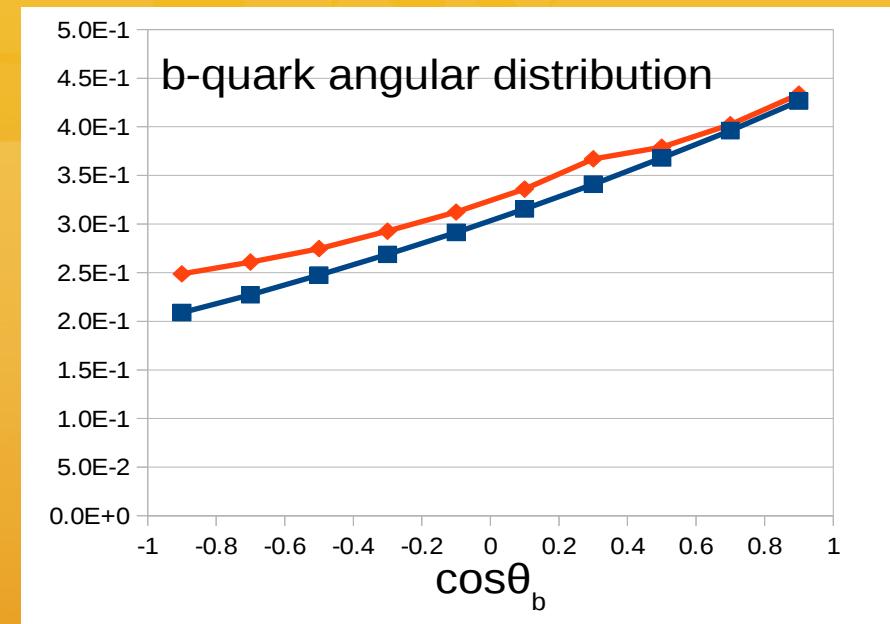
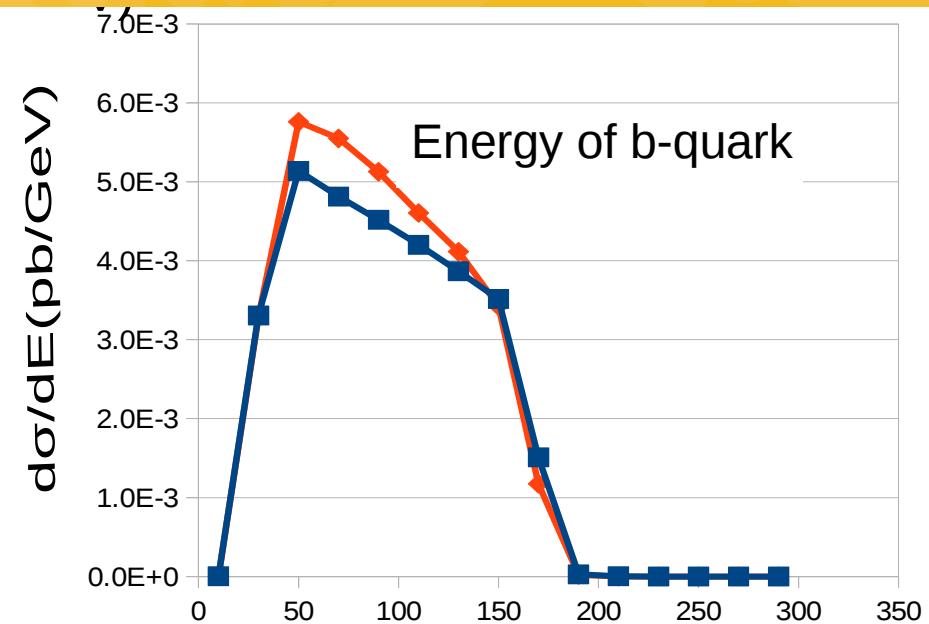
Full EW results for 6-body final process (L-R)



$\sqrt{s}=500\text{GeV}$
 Electron : 100% Left
 Positron : 100% Right

Blue : NWA Tree
 Orange : NWA NLO(ELWK)

Full EW results for 6-body final process (R-L)



$\left. \begin{array}{l} \sqrt{s}=500\text{GeV} \\ \text{Electron : 100\% Right} \\ \text{Positron : 100\% Left} \end{array} \right\}$

Blue : NWA Tree
 Orange : NWA NLO(ELWK)

Summary

- ✿ We are successful to calculate full electroweak one loop correction for $t\bar{t}$ pair with the effect of the polarization.
- ✿ Thanks to the polarization, $e^-_L e^+_R$ cross section is larger than $e^-_R e^+_L$ one and the loop correction of $e^-_L e^+_R$ is smaller than that of $e^-_R e^+_L$.
- ✿ Due to the sizable correction of forward-backward asymmetry, higher order correction with beam polarization should be considered at the ILC.
- ✿ Six body final state with polarization can be included into full electroweak loop correction by using GRACE with narrow width method.

**Thank you for your
attention!**