

# Jet Reconstruction at the LC

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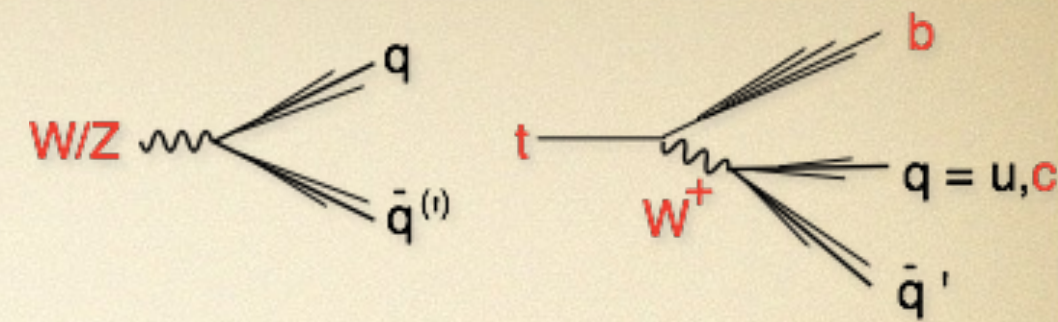
# outline

- brief introduction to jet clustering
- new challenges at future  $e^+e^-$  colliders
  - ▶ ongoing studies on beam jet removal
  - ▶ efforts towards new color-singlet clustering

Many thanks to I. Garcia for providing some material of nice study in Valencia group, see his talk @ CLIC Workshop 2015



# introduction to jet clustering



**motivation:** reconstruct parton level information of main hard process  
which is calculable using perturbative QCD

**requirement:** infrared safe — well behaved for the soft and collinear limit

## sequential type

- i) define inter-particle distance (jet resolution variable)  $y_{ij}$
- ii) find the **smallest**  $y_{ij}$  and combine particle i and j into one pseudo-particle
- iii) iterate the recombination process until  $y_{ij} > y_{\text{cut}}$  or  $\# \text{jet} = \# \text{fixed}$

## global type

- Cone algorithm (Sternan-Weinberg jet): maximize the energy inside one fixed cone
- Georgi algorithm (arxiv: 1408.1161 / 1408.3823): maximize Jet function

$$J_{\beta}(P_{\alpha}) \equiv E_{\alpha} - \beta \frac{P_{\alpha}^2}{E_{\alpha}} = E_{\alpha} [(1 - \beta) + \beta v_{\alpha}^2]$$



# introduction to jet clustering: sequential algorithms at e+e-

Algorithm		resolution scale	comment
JADE	J	$2E_i E_j (1 - \cos \theta_{ij})$	
DURHAM	D	$2 \min(E_i^2, E_j^2) (1 - \cos \theta_{ij})$	
DURHAM (LUCLUS $k_\perp$ ) (Also DURHAM/Lu)	DL	$\frac{2 \mathbf{p}_i ^2  \mathbf{p}_j ^2 (1 - \cos \theta_{ij})}{( \mathbf{p}_i  +  \mathbf{p}_j )^2}$	LUCLUS without reassignment and preclustering
LUCLUS	L	$\frac{2 \mathbf{p}_i ^2  \mathbf{p}_j ^2 (1 - \cos \theta_{ij})}{( \mathbf{p}_i  +  \mathbf{p}_j )^2}$	
GENEVA	G	$\frac{8}{9} E_{\text{vis}}^2 \frac{E_i E_j (1 - \cos \theta_{ij})}{(E_i + E_j)^2}$	
ANGULAR-ORDERED DURHAM	A	$2 \min(E_i^2, E_j^2) (1 - \cos \theta_{ij})$	CAMBRIDGE without soft-freezing
CAMBRIDGE	C	$2 \min(E_i^2, E_j^2) (1 - \cos \theta_{ij})$	
CAMBRIDGE (LUCLUS $k_\perp$ )	CL	$\frac{2 \mathbf{p}_i ^2  \mathbf{p}_j ^2 (1 - \cos \theta_{ij})}{( \mathbf{p}_i  +  \mathbf{p}_j )^2}$	
DICLUS mode 0	Di0	$\frac{(s_{ji} - (m_i + m_j)^2)(s_{ik} - (m_i + m_k)^2)}{s_{ijk}}$	3 $\rightarrow$ 2 clustering
DICLUS mode 1	Di1	$\frac{(s_{ji} - (m_i + m_j)^2)(s_{ik} - (m_i + m_k)^2)}{s_{ijk}}$	largest initial cluster retains its direction
DICLUS mode 2	Di2	$\frac{s_{ji} s_{ik}}{s_{ijk}}$	largest initial cluster retains its direction



# introduction to jet clustering: examples of improvement

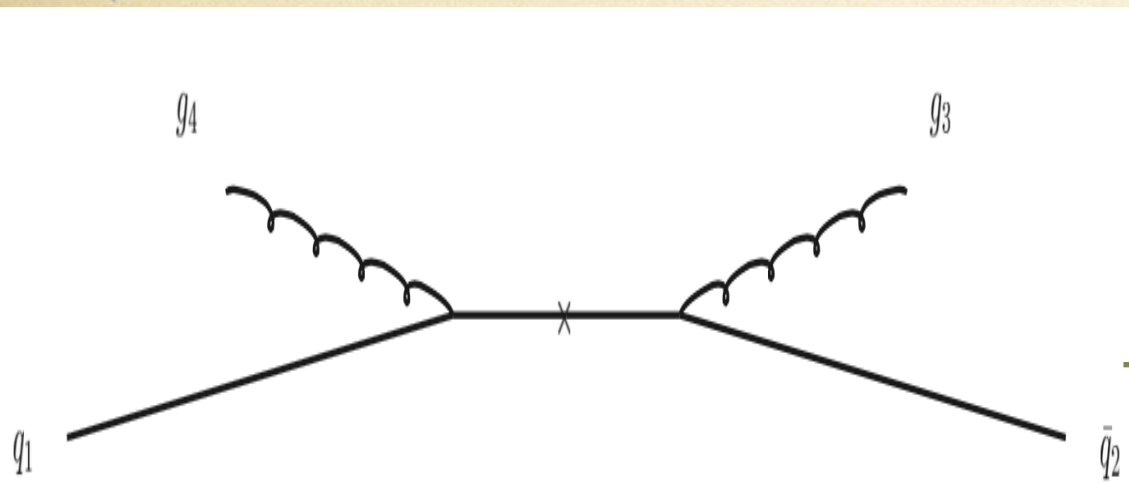


Figure 2: The seagull diagram with  $E_3, E_4 \ll E_1, E_2$  and  $\theta_{14}, \theta_{23} \ll \theta_{34}$ .

$$\text{JADE: } y_{ij} = \frac{2E_i E_j (1 - \cos\theta_{ij})}{E_{\text{vis}}^2}$$

remedy unnatural soft jets by  $g_3 g_4$

$$\text{Durham: } y_{ij} = \frac{2\min(E_i^2, E_j^2)(1 - \cos\theta_{ij})}{E_{\text{vis}}^2}$$

remedy soft large angle  $g_4$

**Angular-Ordered Durham:**

introduce angular measure  $v_{ij}$ , start pair with smallest angle,  $v_{ij} = 2(1 - \cos\theta_{ij})$

**Cambridge: “soft-freezing”**

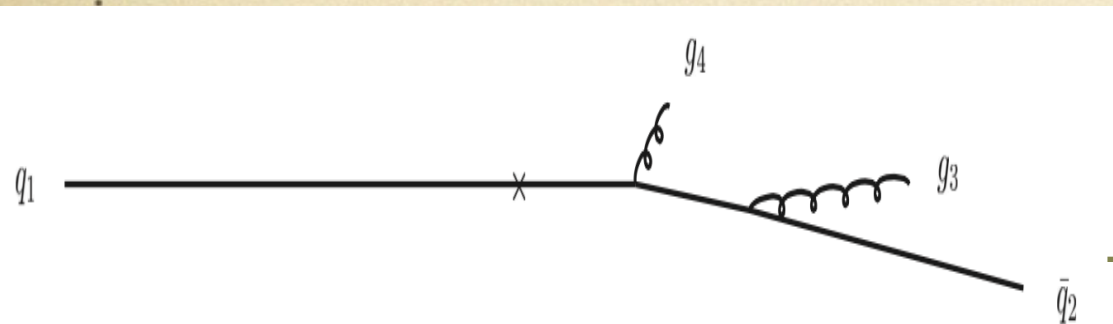


Figure 3: Parton branching with ‘unresolved’, soft, large-angle gluon emission  $g_4$ . Here, one has the following configuration:  $E_2 \gg E_3 \gg E_4$ , and  $\theta_{23} \ll \theta_{24} \approx \theta_{34}$ .

check complete reviews by

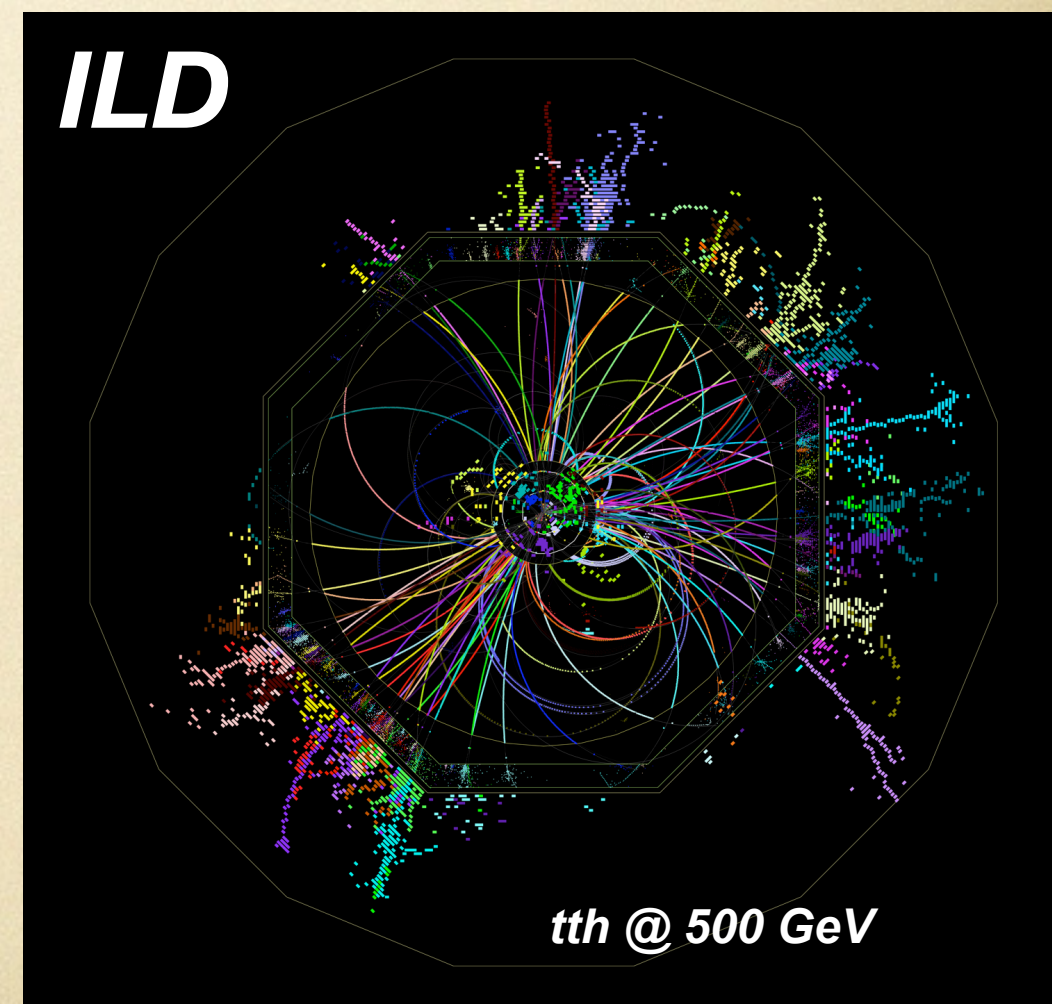
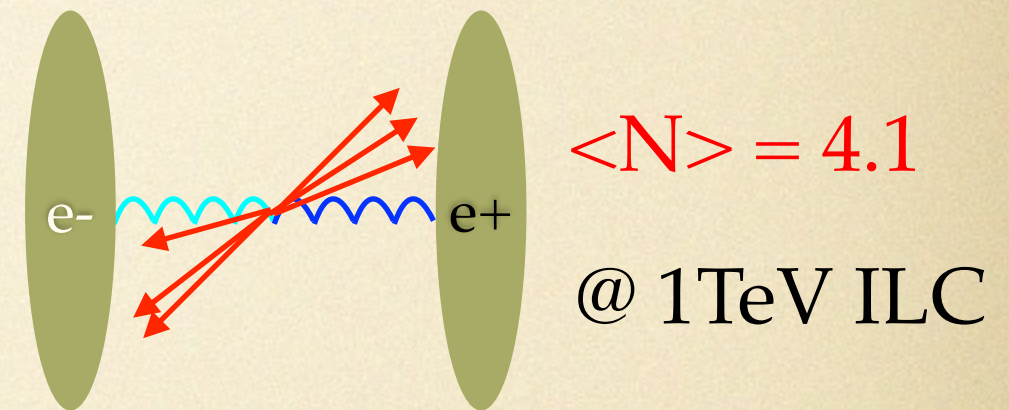
S. Morreti, et al., JHEP08 (1998) 001

A. Ali, et al. arxiv: 1012.2288



# new challenges at next high energy e+e- colliders

- beam jets from underlying events, mainly  $\gamma\gamma \rightarrow$  hadrons: not all particles reconstructed in one event are from interested hard processes
- more jet multiplicities: tt-6jet, ZHH-6jet, ttH-8jet, etc.
- higher available  $Q^2$ : deeper parton shower process
- possible  $y_{ij}$  inside one jet larger than  $y_{ij}$  between two jets  $\rightarrow$  mis-clustering due to jet overlap





# efforts to handle the beam jets — Valencia jet-clustering

- Durham algorithm at e+e- has been adapted to hadron collider where beam remnant is more relevant (quite long time ago): longitudinal invariant  $k_t$  algorithm (anti- $k_t$ ).
- recently, Valencia algorithm has been proposed to combine good features of lepton colliders (Durham like distance).

longitudinal inv.  $k_t$

$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \Delta R_{ij}^2 / R^2$$

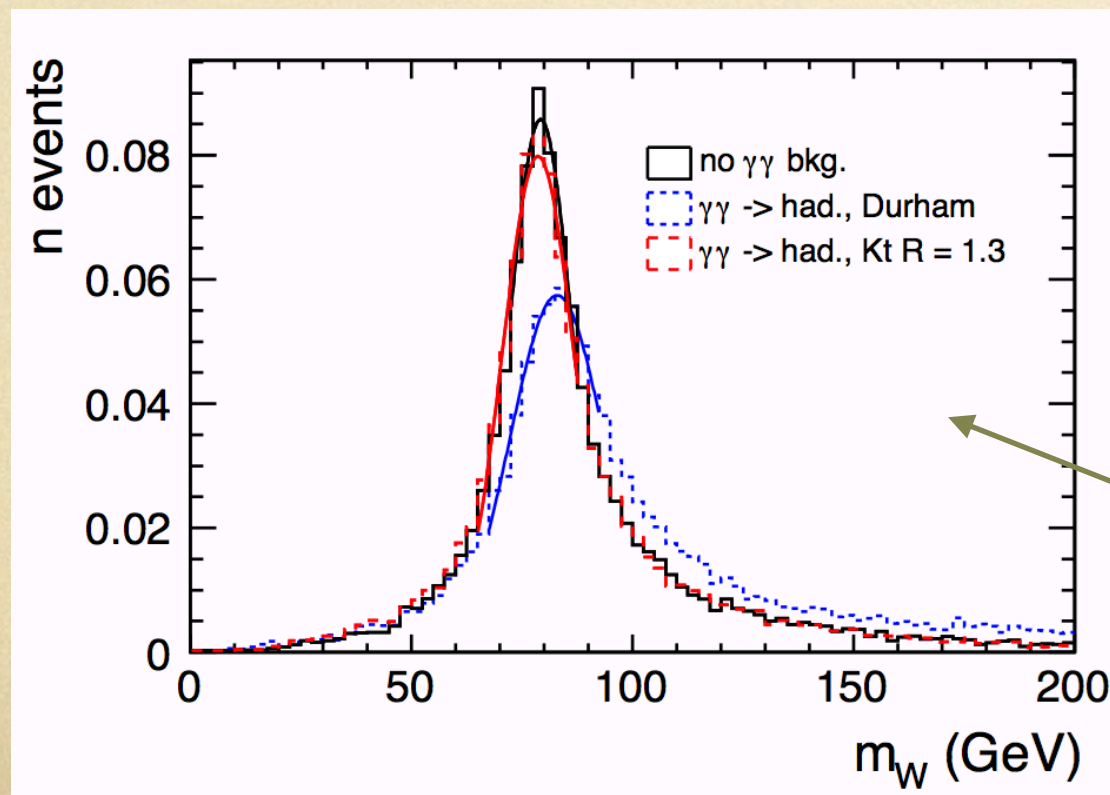
$$d_{iB} = p_{ti}^2$$

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

Valencia algorithm

$$d_{ij} = \min(E_i^{2\beta}, E_j^{2\beta})(1 - \cos \theta_{ij}) / R^2$$

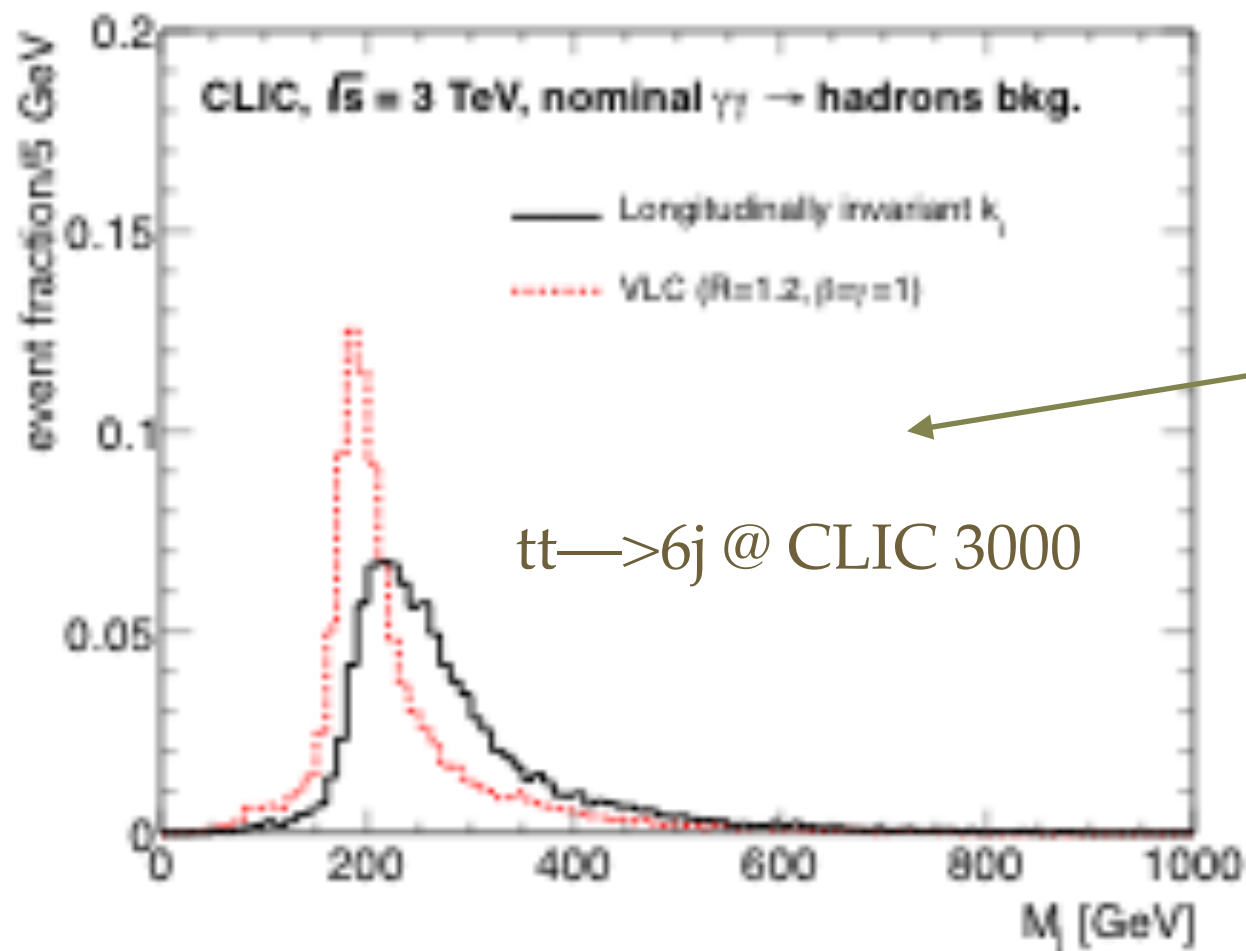
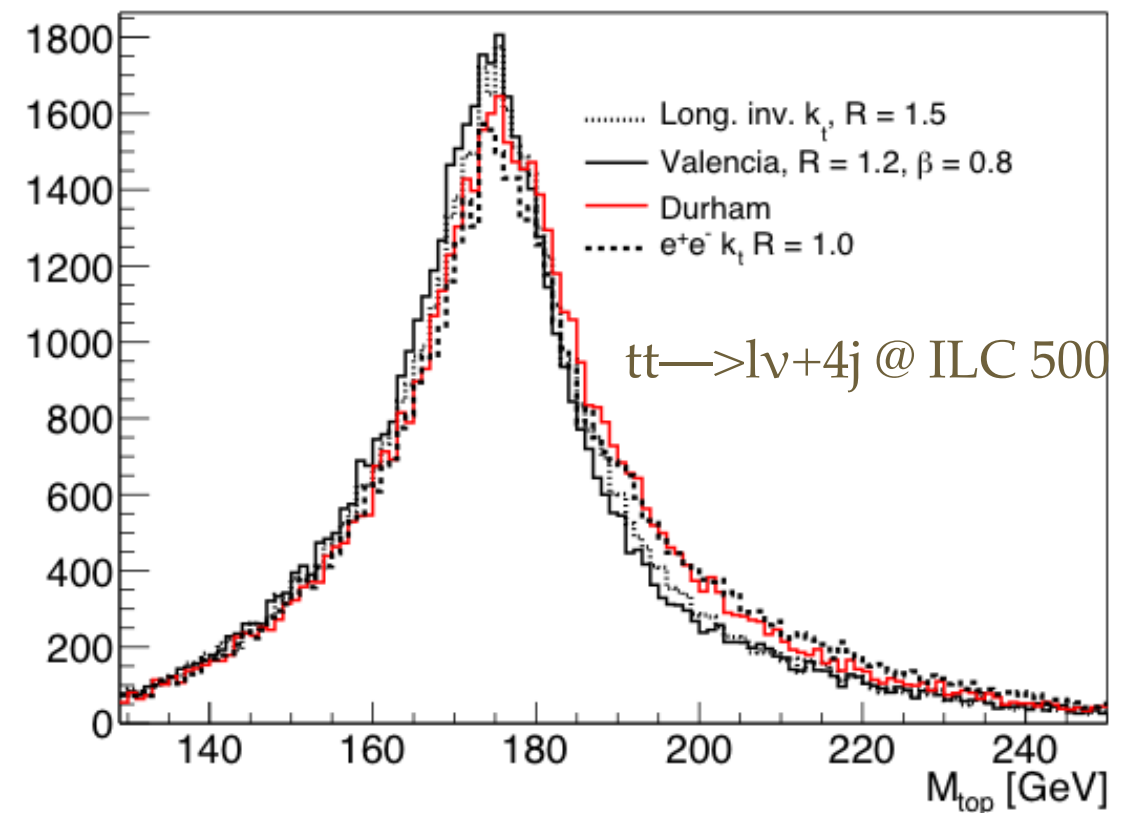
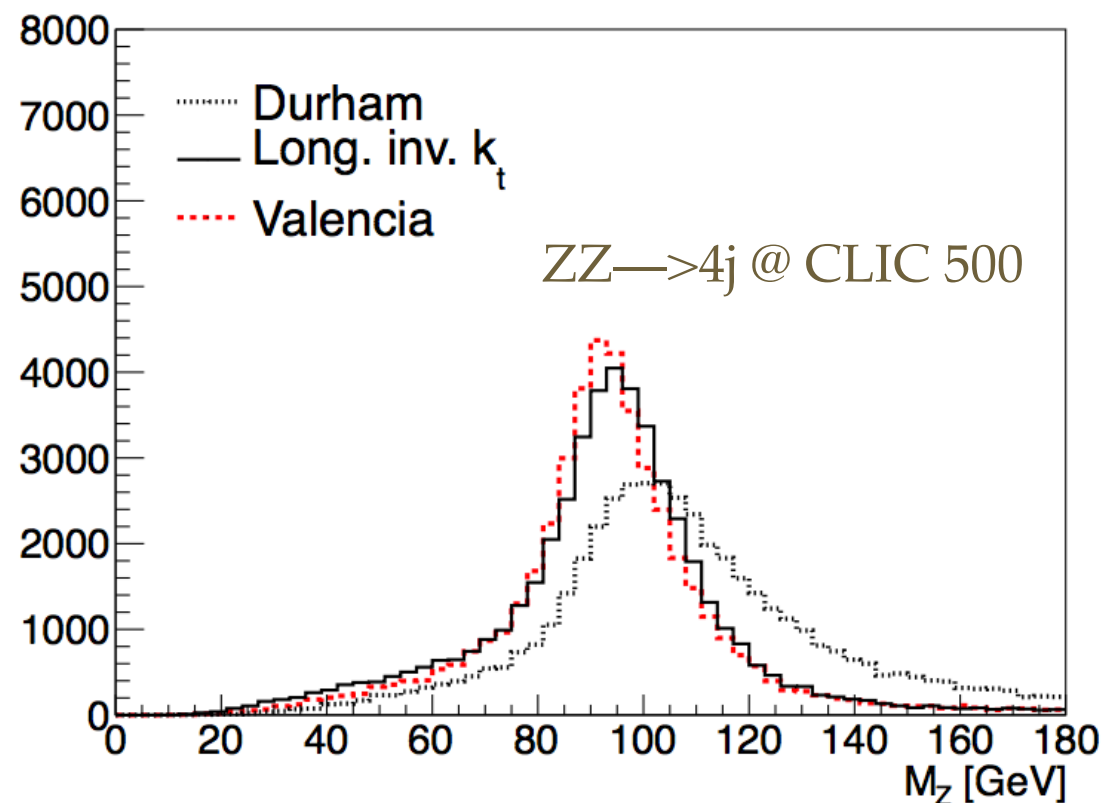
$$d_{iB} = p_{ti}^{2\beta} = E_i^{2\beta} \sin^2 \theta_i$$



benchmark using  $k_t$  for  
 $WW \rightarrow lvqq$  in ILD-DBD



# Valencia jet-clustering: more applications



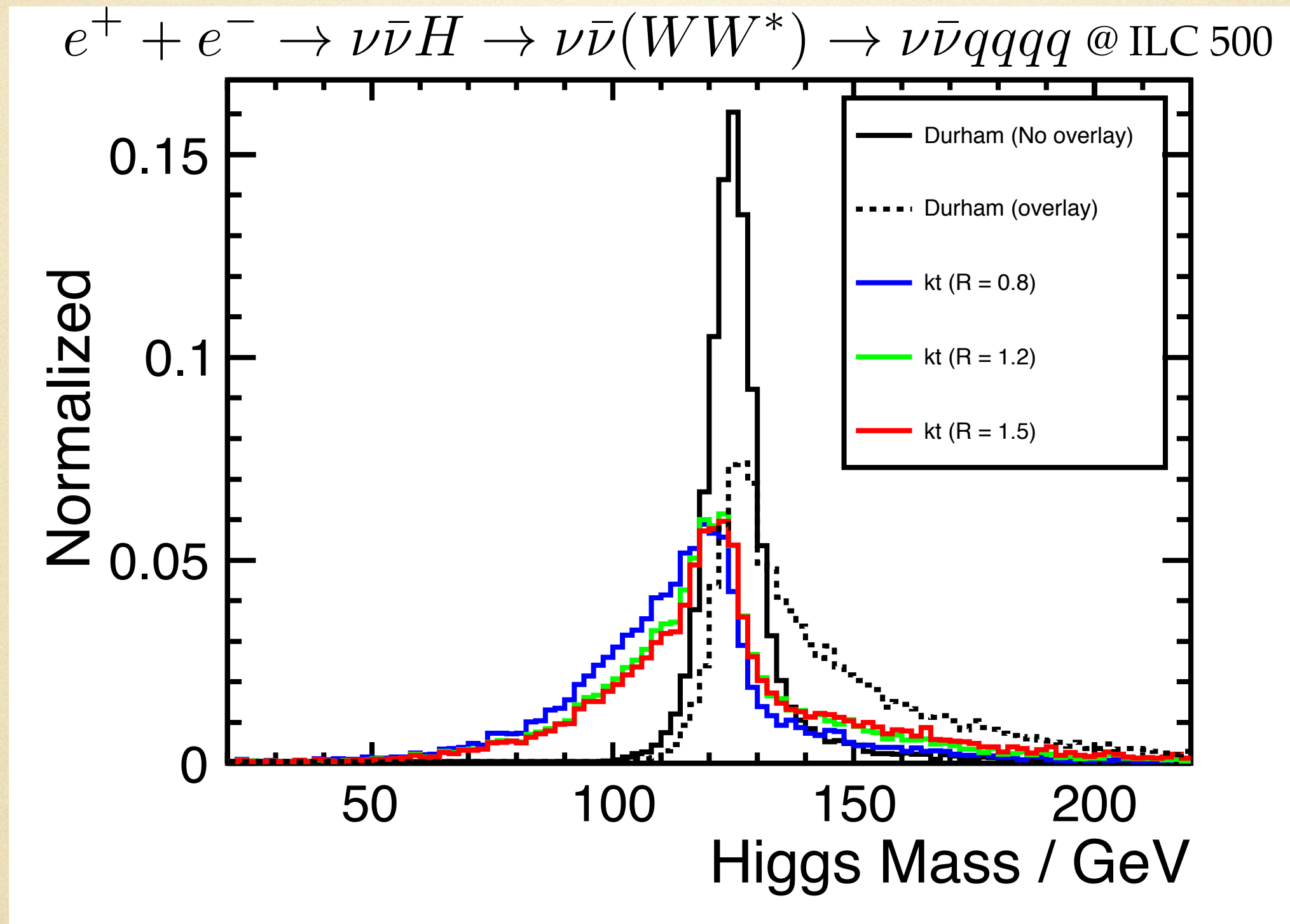
VLC is significantly better

I. Garcia @ CLIC Workshop 2015

M. Vos, et al., arxiv: 1404:4294



# impact of $\gamma\gamma \rightarrow \text{hadrons}$ overlay: one more example



longitudinal inv. kt is barely working in  
WW-fusion channel  $H \rightarrow WW^* \rightarrow 4j$

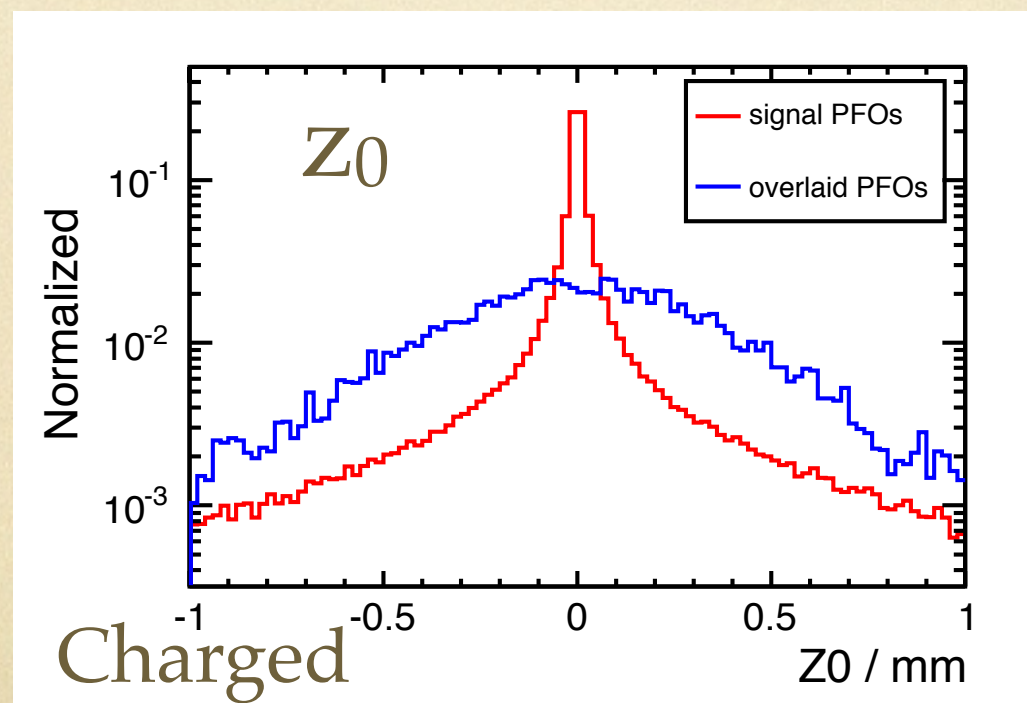
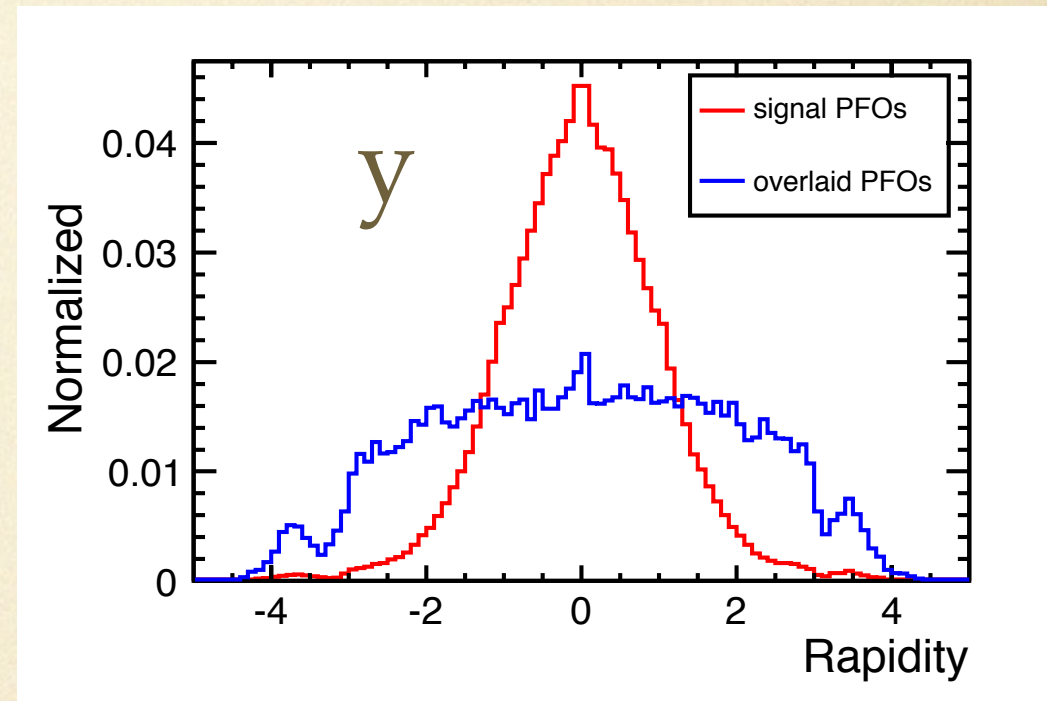
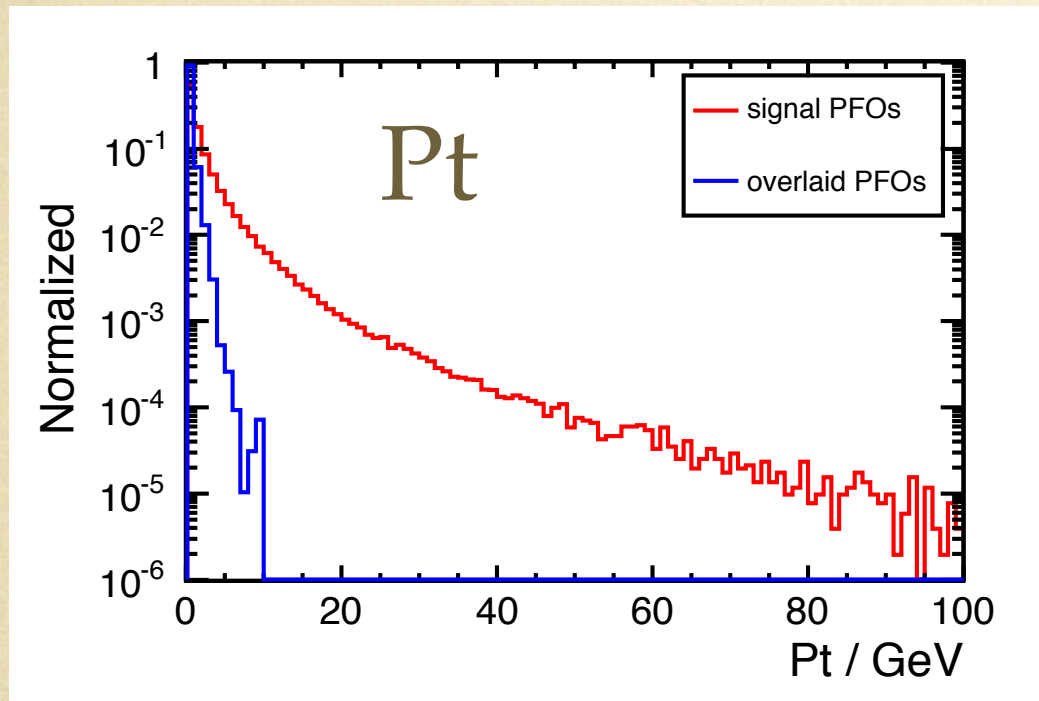


## efforts to handle the beam jets at LC — alternative approach

- the beam remnant at  $e^+e^-$  should be much simpler than those at hadron collider —> in principle we should be able to **reconstruct those remnant exclusively** without generic clustering algorithm.
- the overlay events, though in the same bunch, have a **shifted IP in z-direction** to the primary IP of signal events.
- the bunch **spread in z at ILC  $\sim 300\mu\text{m}$** ; track **impact parameter resolution  $\sim 10\mu\text{m}$** ; if the low momentum tracking and vertex finder work well, we would be able to **reconstructed the IP of those overlay events** when the shift is more than 10s of  $\mu\text{m}$ .
- as a first step approach, try to check how well we can separate signal and overly particle by particle

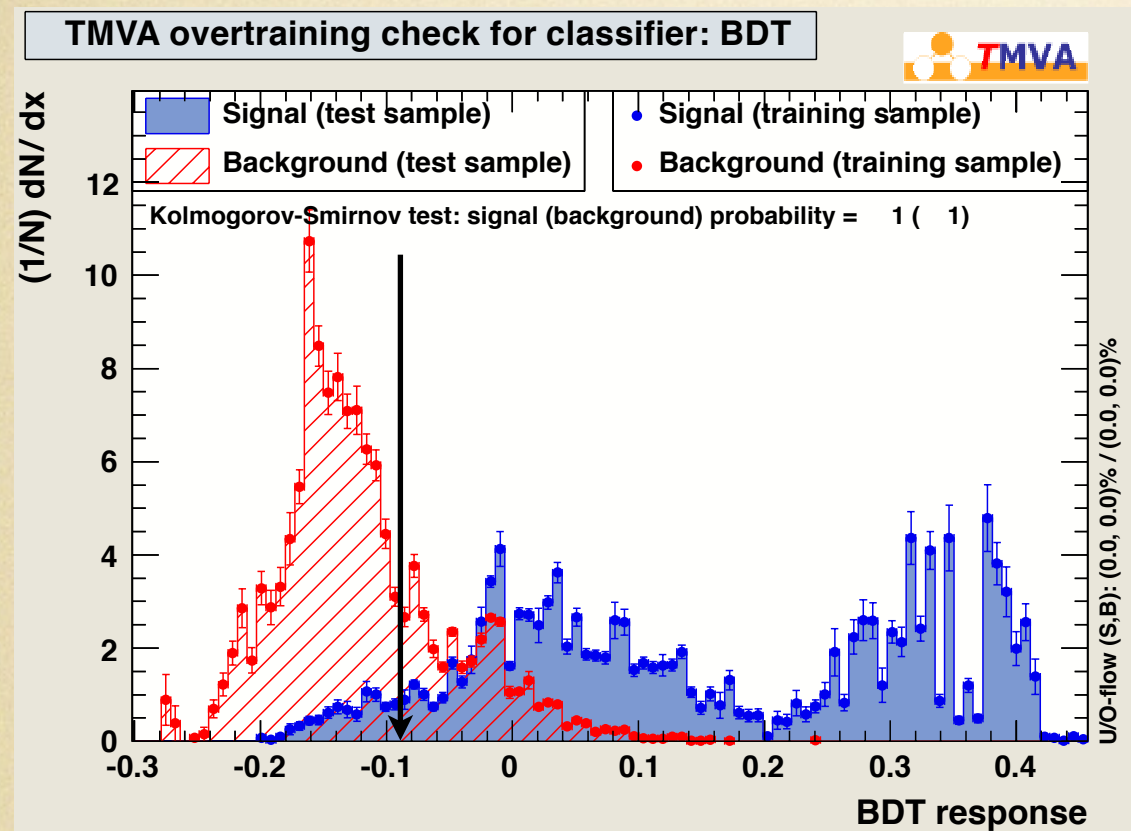


# characteristics of particles from overlay events



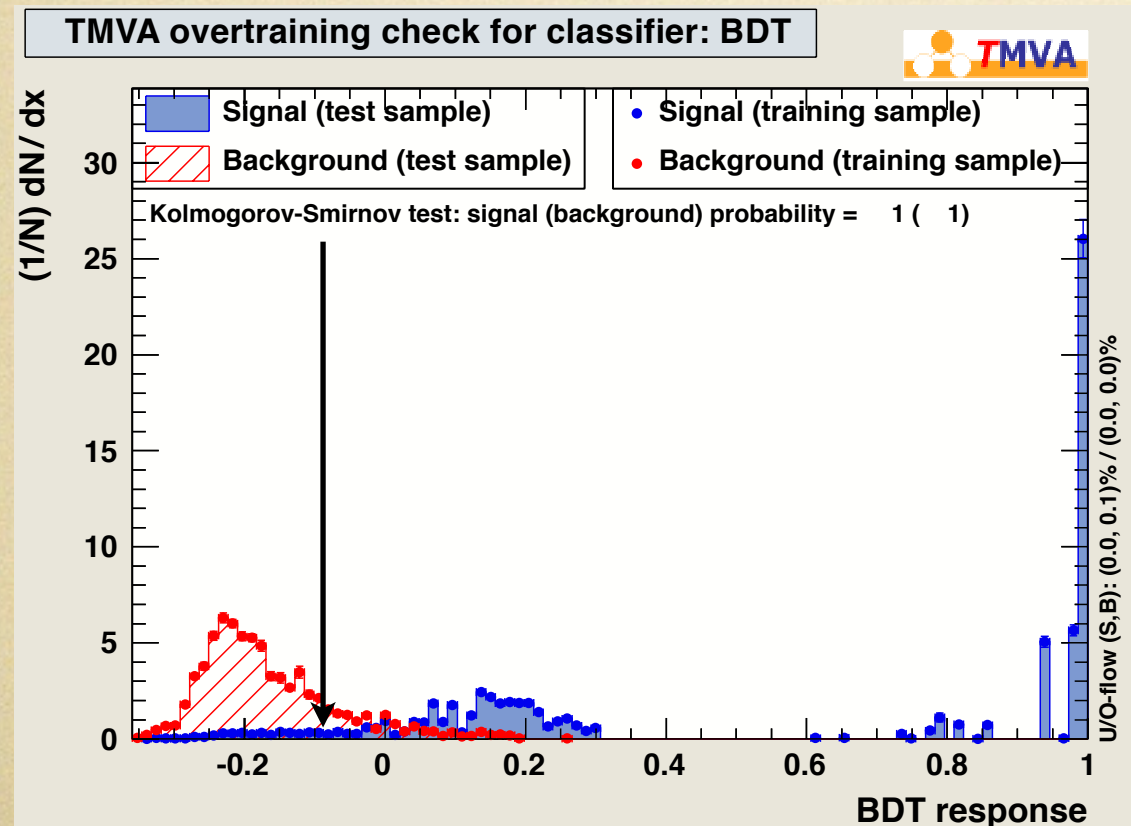


# MVA to separate signal particle and overlay particle



category 1: neural or large  $z_0$

input:  $P_t$ , Rapidity



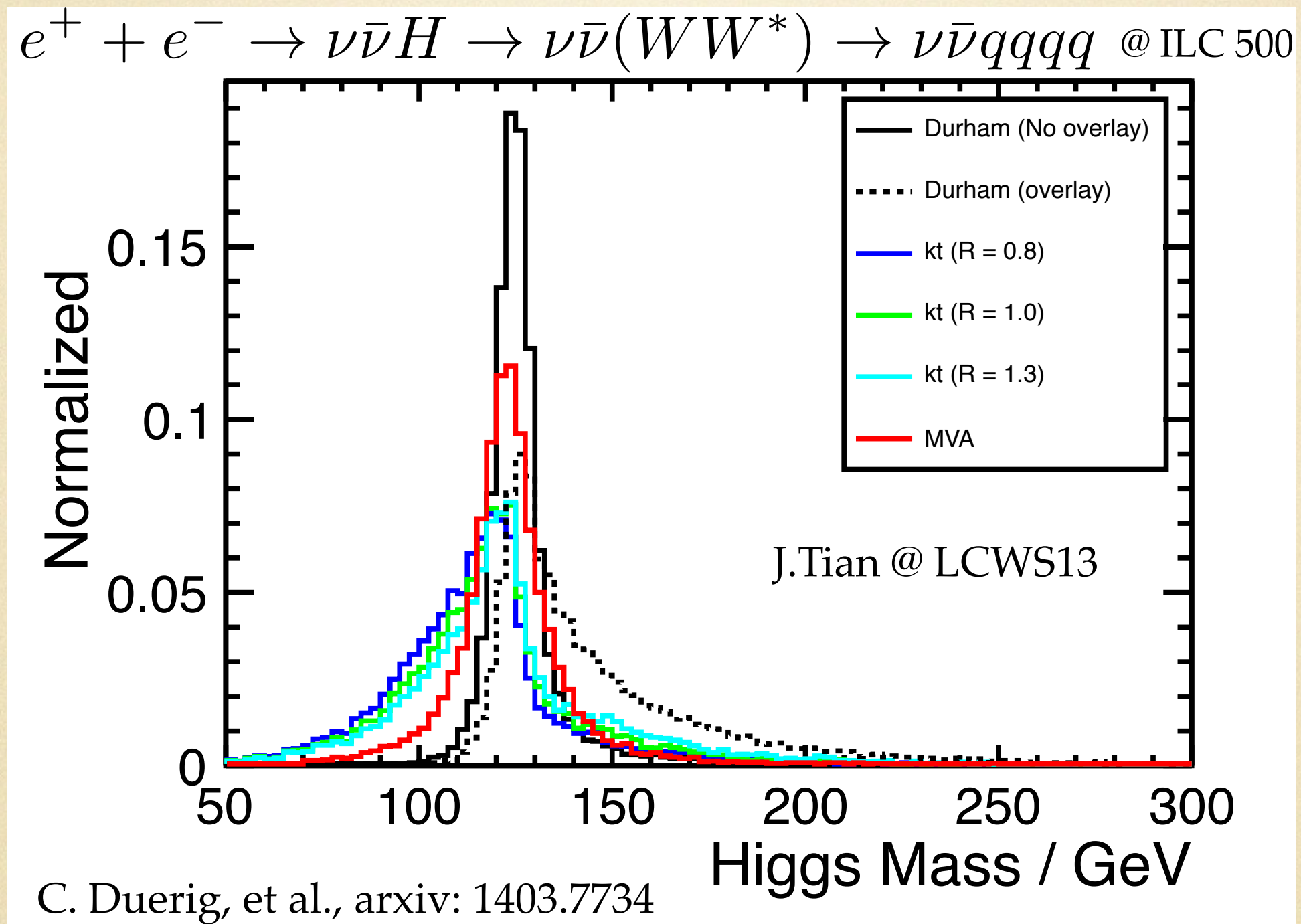
category 2: charged and small  $z_0$

input:  $P_t$ , Rapidity,  $z_0$

each PFO weighted by energy in both cases



# efforts to handle the beam jets at LC— MVA approach



as a first step, MVA seed particle tagging already gives better performance in  $H \rightarrow WW^*$  process; need generalize to clustering algorithm and add vertex seeds, and check whether it's better in other channels; ongoing



## towards to a generic color-singlet jet clustering

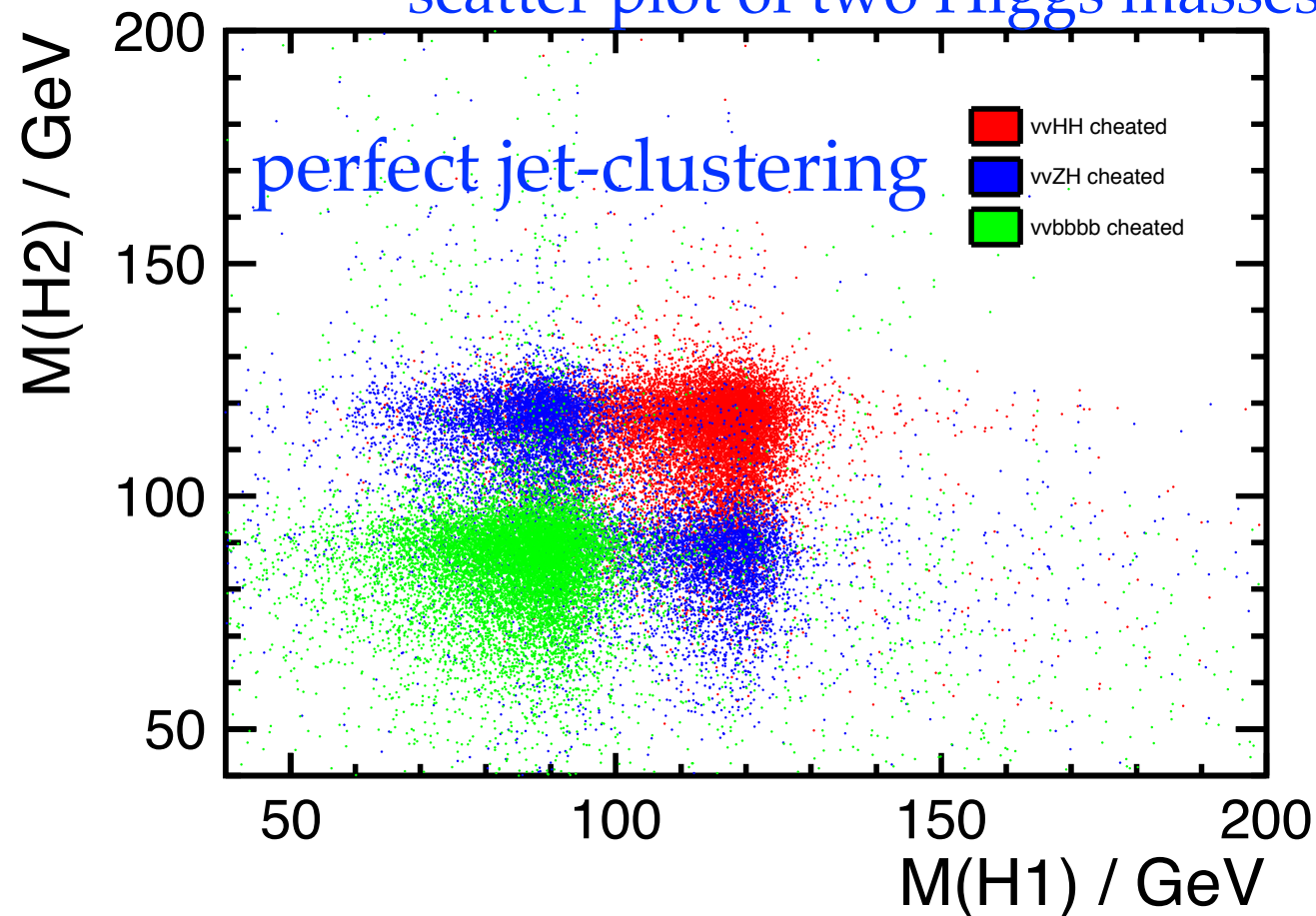
- first of all this idea is in a conceptual level, feel free to start your coffee break now,
- and I'm an experimentalist, feel free to ignore what's I'm going to propose towards a very ambitious new clustering algorithm.
- but, we do need help from you, in particular QCD experts.
- starting point: why do need a much better color-singlet clustering



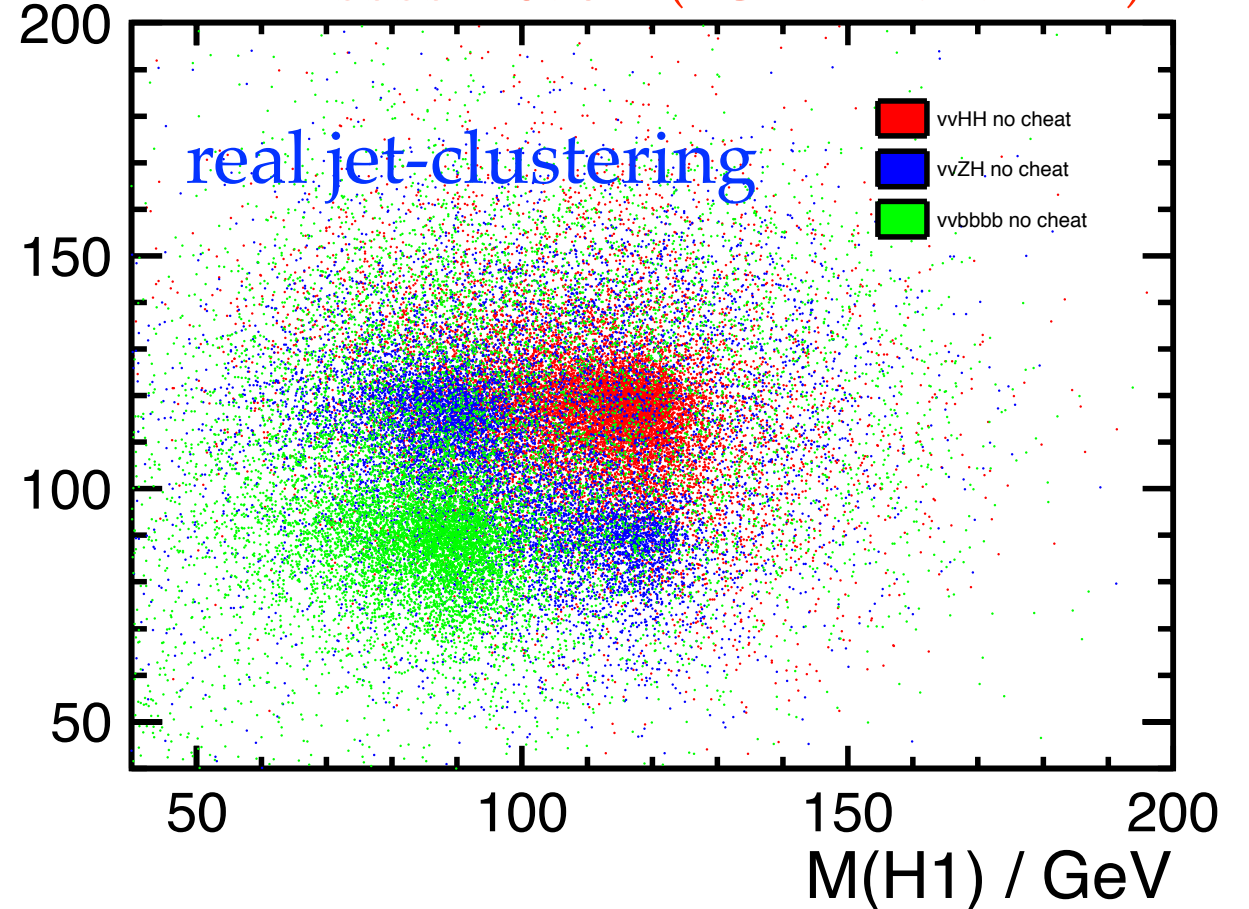
# impact of jet-clustering in Higgs self-coupling measurement

(without beam overlay now)

scatter plot of two Higgs masses



vVHH $\rightarrow$ vVbbbb mode: (BG: ZZH and ZZZ)

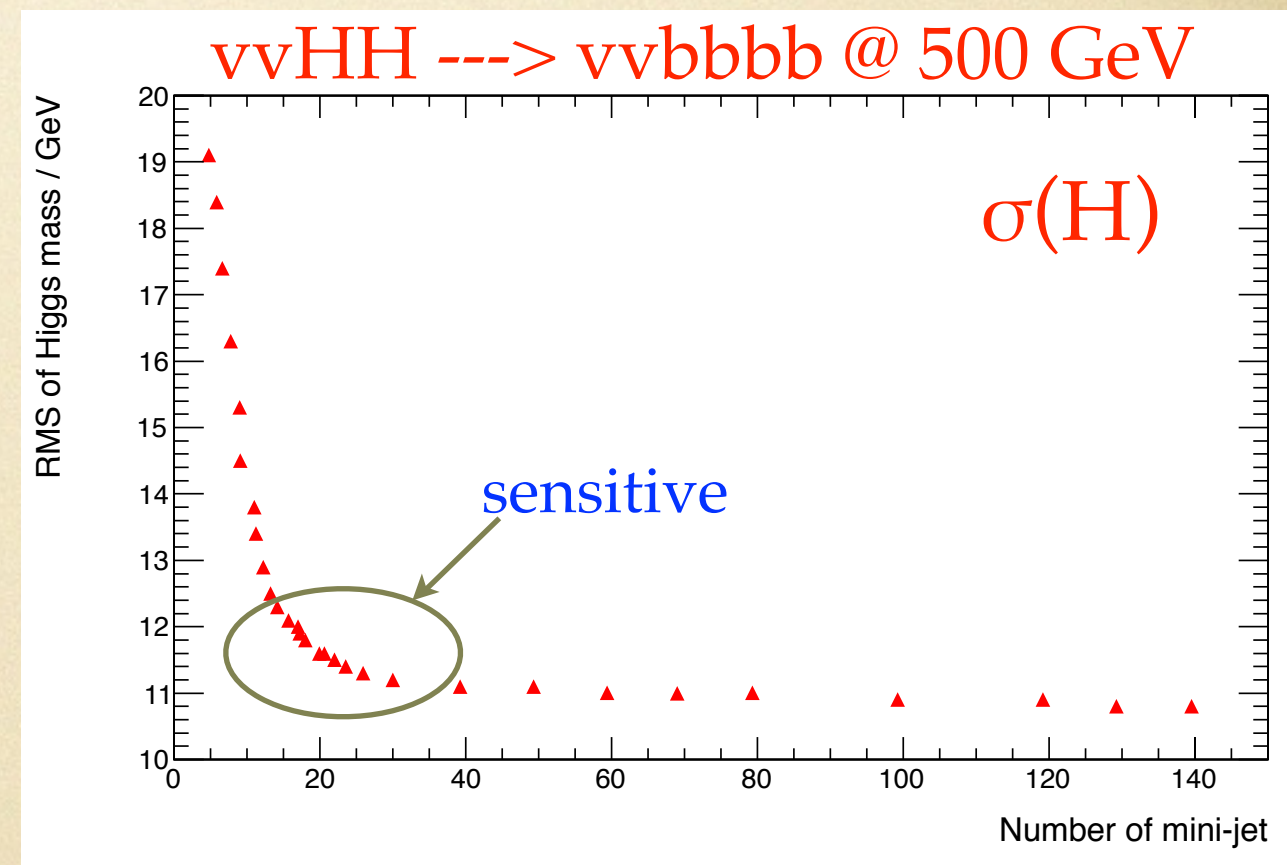
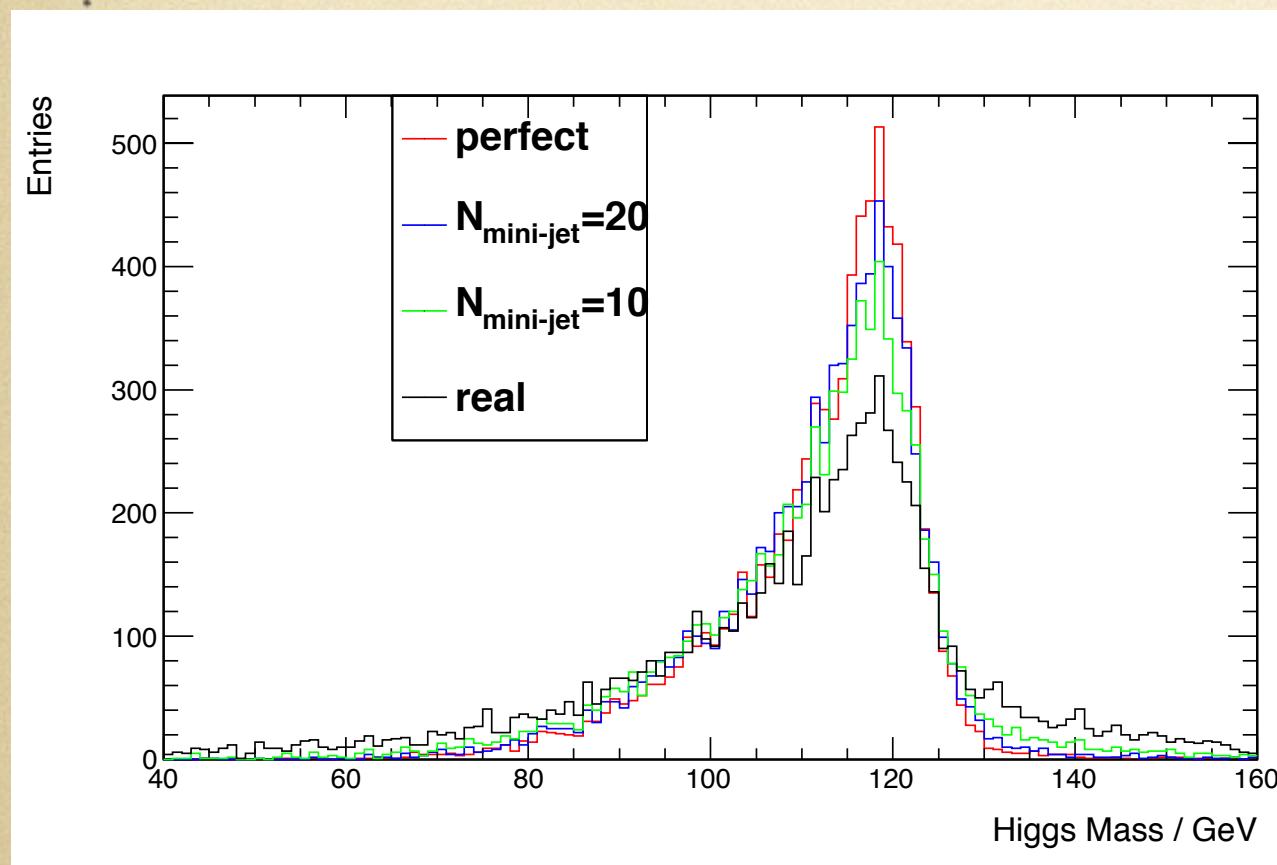


it has been studied if a color singlet jet clustering can be implemented for both signal and BG,  $\lambda_{HHH}$  measurement can be improved by 40%, which means 20%  $\delta\lambda_{HHH}/\lambda$  ( $5\sigma$ ) would already be possible at 500 GeV ILC with the H20 scenario.



# investigation so far: when does mis-clustering start?

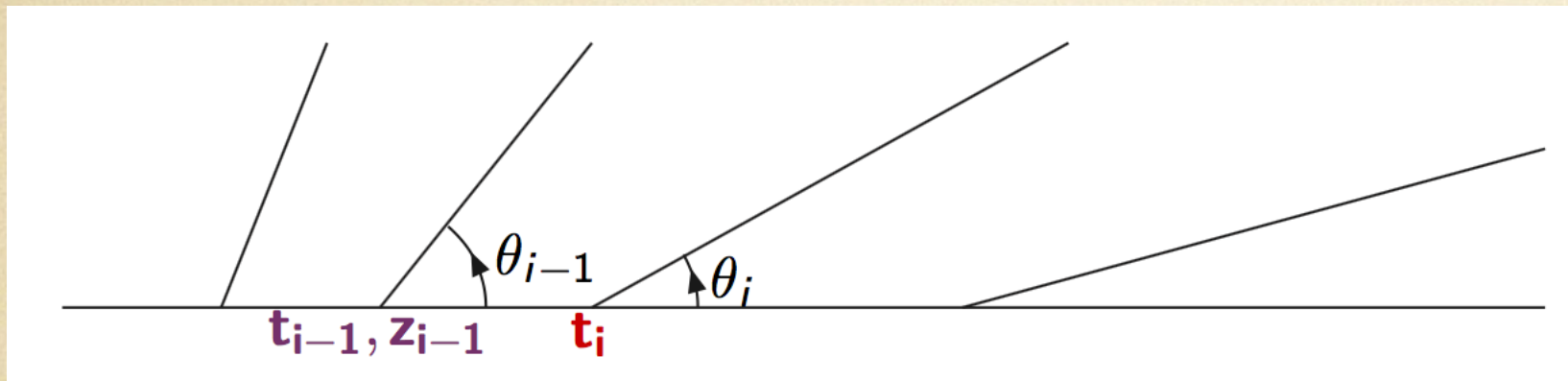
- i) mini-jet clustering (pre-clustering) with realistic Durham algorithm, stop at certain fixed number of jet
- ii) combine those mini-jets using MC truth information to two Higgs (color-singlet)
- iii) check the dependence of Higgs mass with the number of mini-jets



Durham jet clustering starts to have major mis-clustering when remaining #mini-jet  $\sim 20$



# proposal of new clustering: reconstruct the parton shower likelihood



$$d\mathcal{P}_a = \sum_{b,c} \frac{\alpha_{abc}}{2\pi} P_{a \rightarrow bc}(z) dt dz .$$

$$\begin{aligned} P_{q \rightarrow qg}(z) &= C_F \frac{1+z^2}{1-z} , \\ P_{g \rightarrow gg}(z) &= N_C \frac{(1-z(1-z))^2}{z(1-z)} , \\ P_{g \rightarrow q\bar{q}}(z) &= T_R (z^2 + (1-z)^2) , \\ P_{q \rightarrow q\gamma}(z) &= e_q^2 \frac{1+z^2}{1-z} , \\ P_{\ell \rightarrow \ell\gamma}(z) &= e_\ell^2 \frac{1+z^2}{1-z} , \end{aligned}$$

- do mini-jet clustering until #mini-jet  $\sim 20$ , during which Durham algorithm works well  $\rightarrow$  which also helps retain infrared safe
- new algorithm comes in to combine the mini-jets: if these mini-jets are actually the relic of parton shower, we should aim for the reconstruction of parton shower history as the maximum information
- generic feature of parton shower can be used to help: angular ordering,  $\theta_{i-1} < \theta_i < \theta_{i+1}$ ,  $t_{i-1} < t_i < t_{i+1}$ .
- assign each branching with a probability  $P_{q \rightarrow qg}$ , etc.
- above is intraJet parton shower; we can use some feature of color correlation interJet, such as rapidity gap, etc. (not really sure).



# summary

- there have been many jet algorithms for  $e^+e^-$  colliders studied in the past decades, it's worth understanding the advantages in each of them, which could give us some useful hints.
- new challenges are expected at the next high energy  $e^+e^-$  collider, to handle beam jets and mis-clustering.
- improvements on removing beam jets by robust and generic Valencia clustering algorithm, and specific particle / vertex based removal method, have been observed. It would be interesting to take a look at the comparison and more applications. The vertex based method would benefit a lot from improved low momentum tracking and vertex finder.
- more general color-singlet jet clustering algorithm would help a lot for all Higgs self-coupling measurement and possibly most of others with multi-jet final states, a conceptual proposal to maximal use parton shower information is just started.

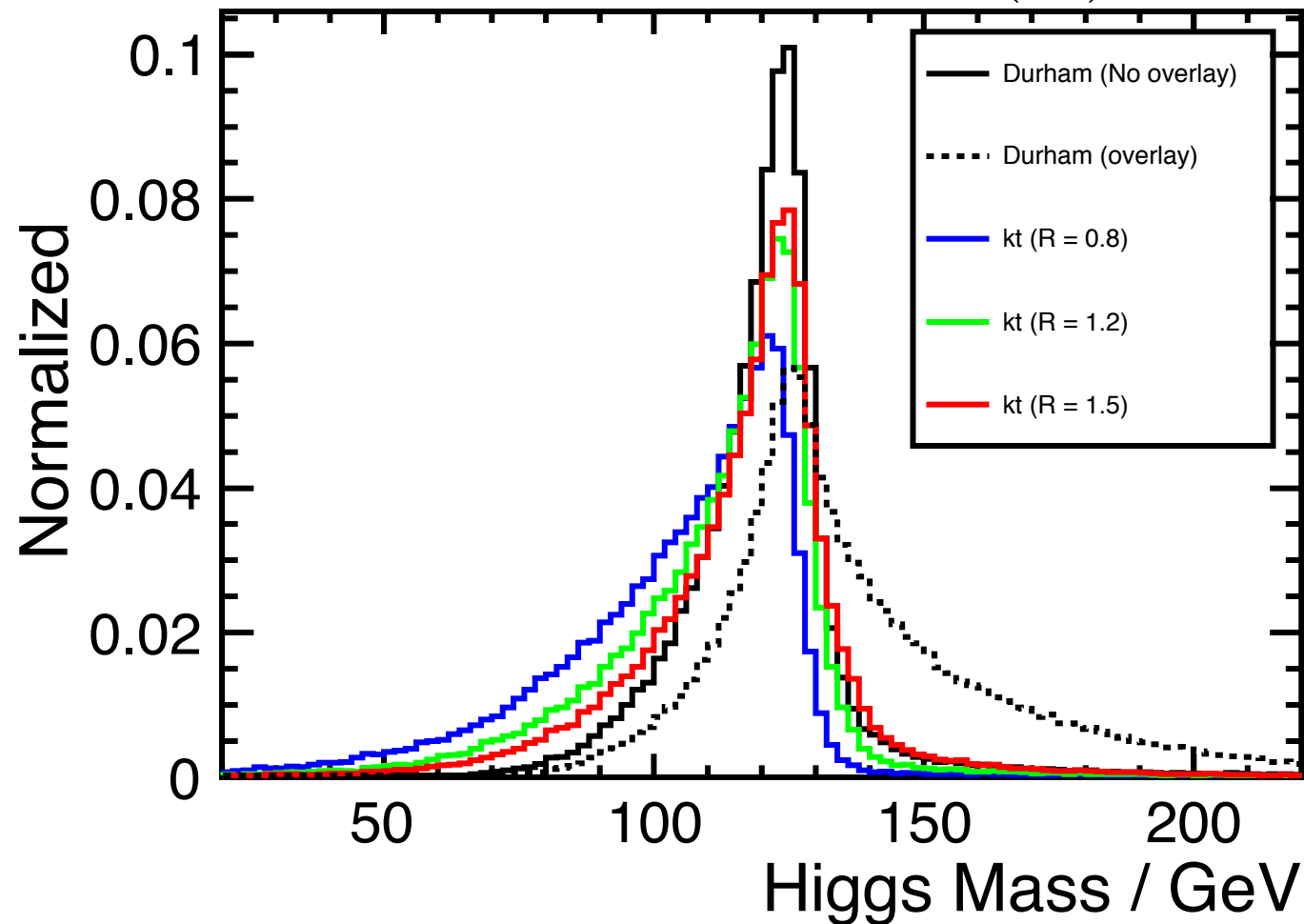


backup



# typical method: kt jet-clustering

$e^+ + e^- \rightarrow \nu\bar{\nu}H \rightarrow \nu\bar{\nu}(b\bar{b})$  @ 500 GeV



$$d_{ij} = \min(p_{ti}^2, p_{tj}^2) \Delta R_{ij}^2 / R^2$$

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

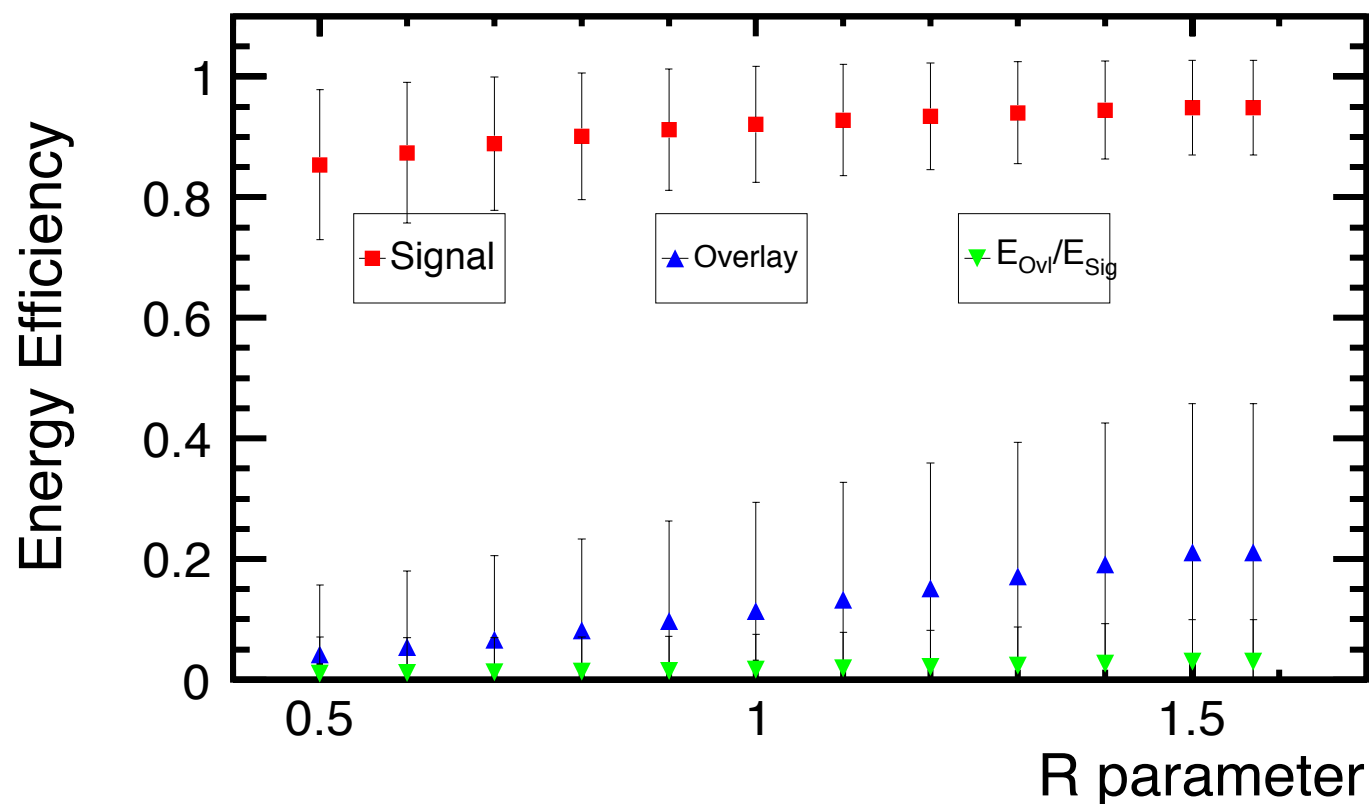
overlaid particles usually very forward  $\rightarrow$  large  $y \rightarrow$  far from physics jet

paras opt in kt jet clustering

Max No. of Jets = 2

$R = 1.5$

overlay is removed efficiently



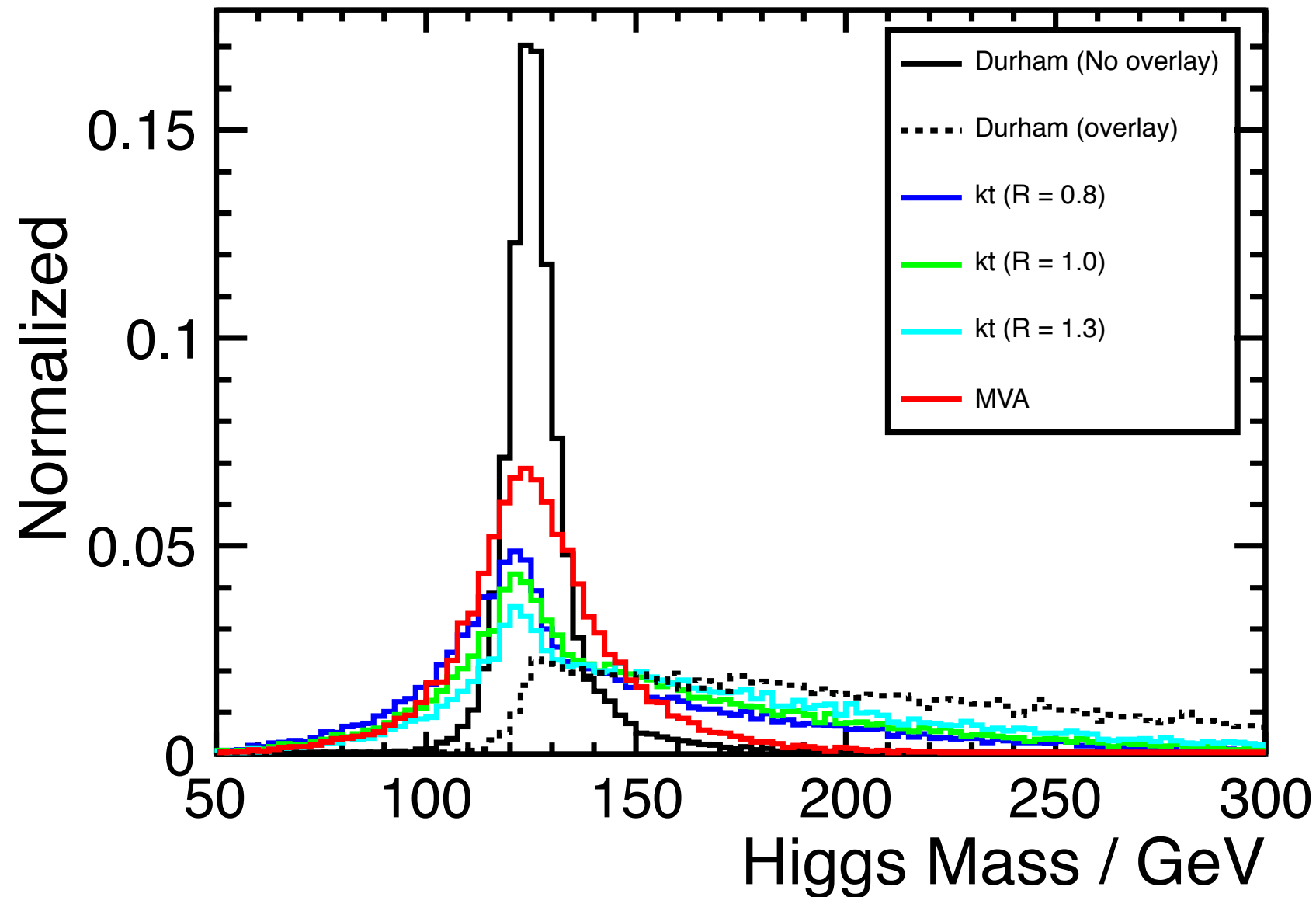
Eff(sig)  $\sim$  95%

Eff(ovl)  $\sim$  21%

purity  $\sim$  97%



$$e^+ + e^- \rightarrow \nu\bar{\nu}H \rightarrow \nu\bar{\nu}(WW^*) \rightarrow \nu\bar{\nu}qqqq \quad @ 1 \text{ TeV}$$



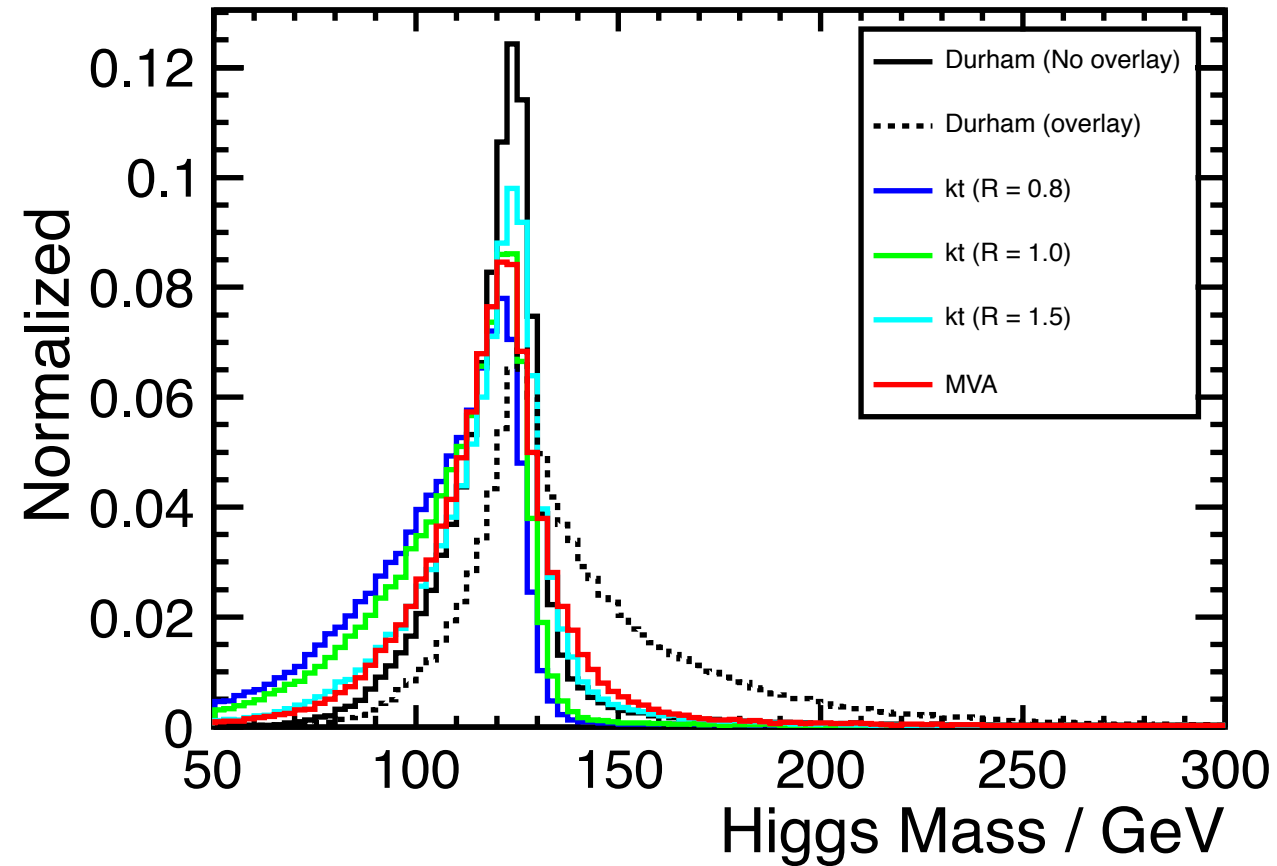
Eff(sig) ~ 89%  
 Eff(ovl) ~ 16%  
 purity ~ 95%

better than kt, but obviously not satisfactory

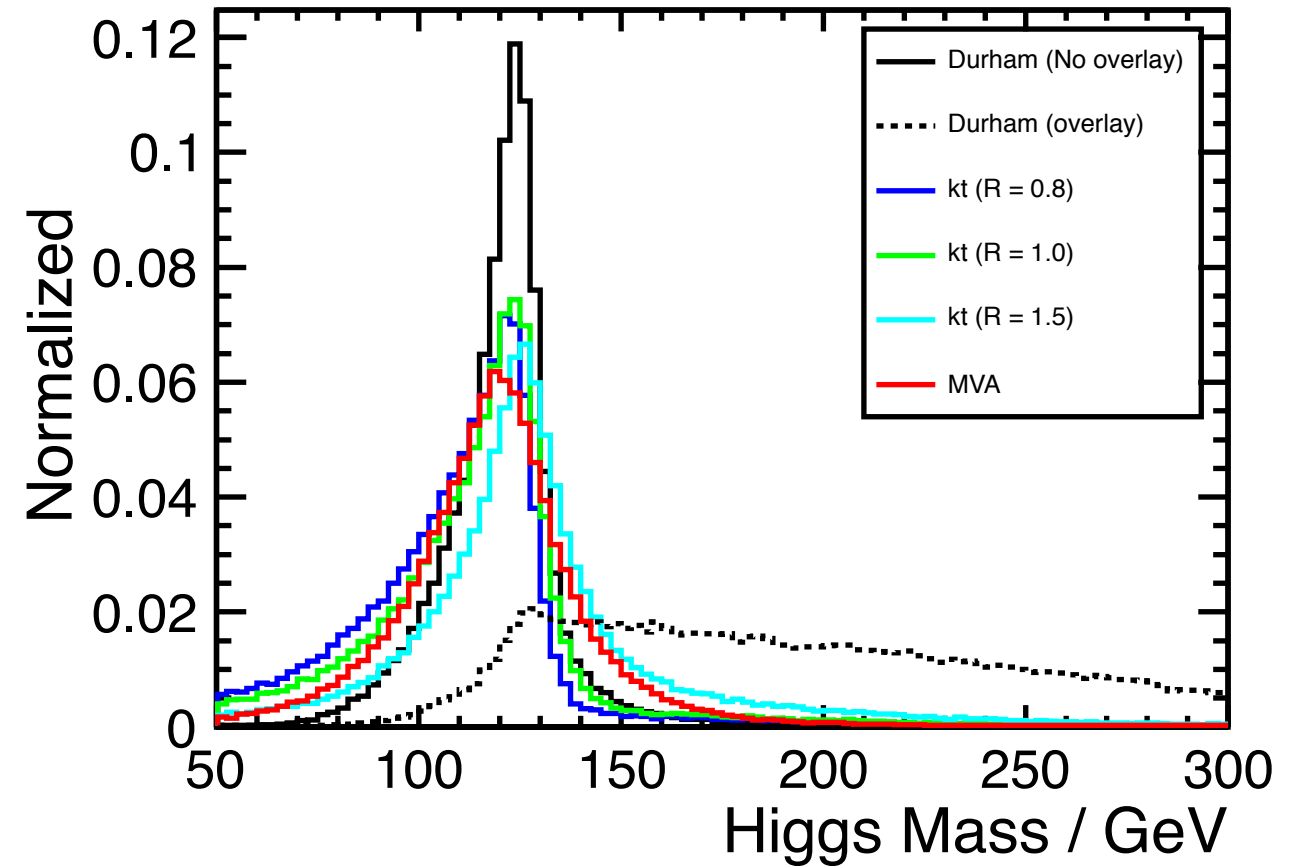


$$e^+ + e^- \rightarrow \nu\bar{\nu}H \rightarrow \nu\bar{\nu}(b\bar{b})$$

@ 500 GeV



@ 1 TeV



in this channel, MVA is not better than kt...