

Time Resolution of Thin LGADs

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with

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- Results from the Nov 2014 Beam Test at CERN
- Improvement in hand: Sensor Capacitance
- Measurements with α Front TCT



Time Resolution and Slew Rate

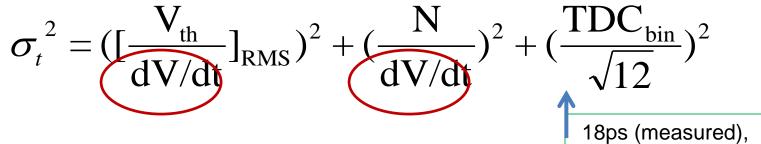
The time resolution σ_t depends on rise time τ_r ,

and τ_r depends on the collection time (i.e. the detector thickness).

The time resolution has 3 terms: **time walk** due to amplitude variation, **time jitter** due to noise, **binning** resolution:

$$\sigma_t^2 = ([\tau_r \frac{V_{th}}{S}]_{RMS})^2 + (\tau_r \frac{1}{S/N})^2 + (\frac{TDC_{bin}}{\sqrt{12}})^2$$

Introducing the slew-rate S/ $\tau_r = dV/dt$



(binning error is negligible at 50 ps sampling rate)

18ps (measured), vs 14ps (expected)

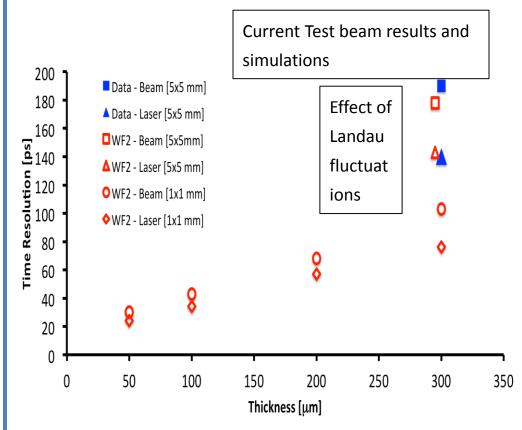
we find that for constant noise N, to minimize the time resolution, we need to maximize the slew-rate dV/dt of the signal (and threshold in the simple case)

Need both large and fast signals.



Time Resolution for thin LGAD

For the 300 μ m thick large LGAD pads (C \approx 10 pF), the time resolution measured in the beam tests (BT) at Frascati and CERN, is predicted by the **Weightfield (WF2)** simulation)



The time resolution is predicted to improve for smaller LGAD and optimized electronics. Reduced the thickness also improves the time resolution.

Expect for 50 μ m thick LGAD (C~2pF):

 $\sigma_t = 30 \text{ ps} \text{ (requires ASIC)}$

N. Cartiglia, F. Cenna et al. "Weightfield", 2014 IEEE NSS-MIC



Santander 2015

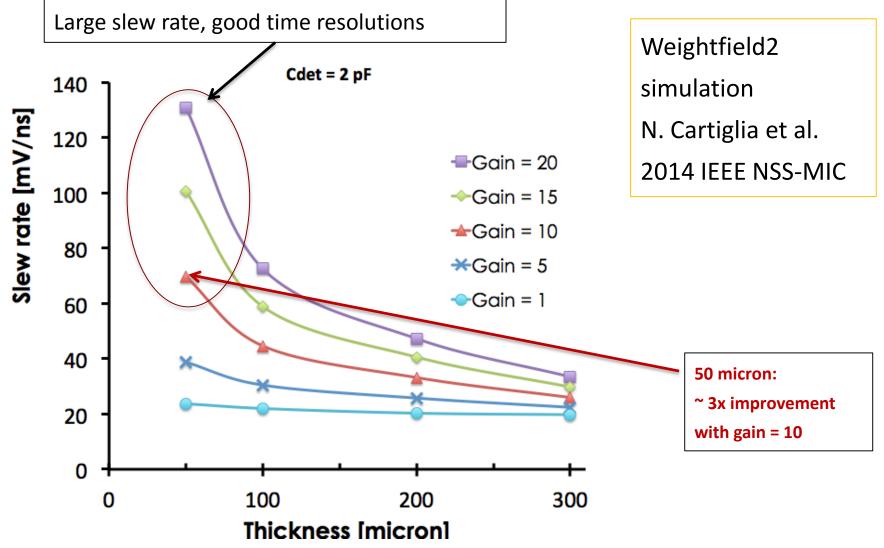
RD50

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Slew-rate as a Function of Sensor Thickness



Significant improvements in time resolution require thin and small detectors

Hartmut F.-W. Sadrozinski, Beam Test of LGADs, 10th Trento, Feb. 2015



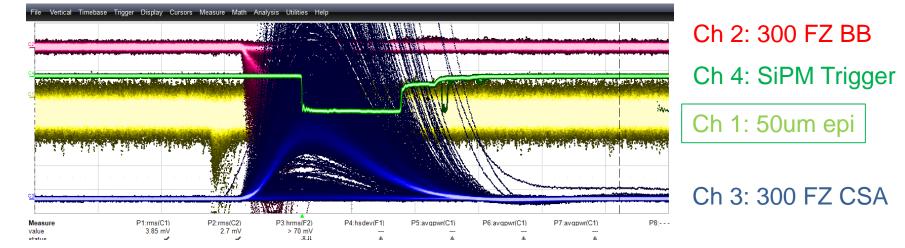
2014 November nights in the H6 Beam Line

2 300FZ LGAD (G = 10-15)

50um epi LGAD (G = 2-3).

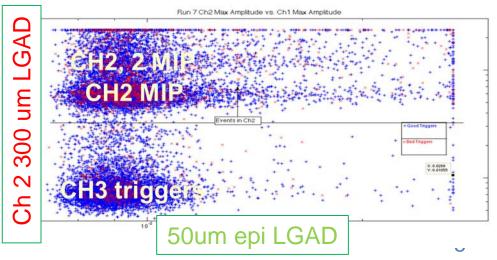
Trigger: fast SiPM

Bias voltage scans taken with digital scope at 10 ps sample speed.



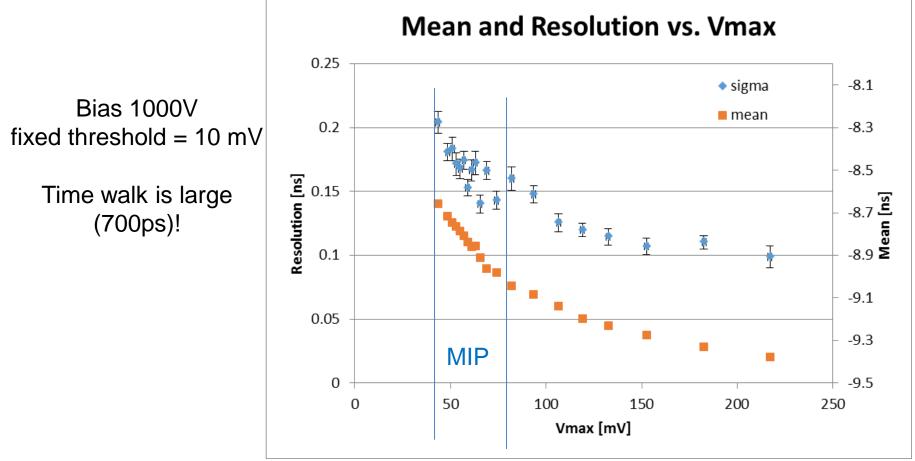
Trigger: (OR of 300um FZ sensors). AND.SiPM (SiPM resolution ~50 ps?).

Thanks to the AFP beam test group Joern Lange & Michael Rijssenbeek) for help and use of the trigger





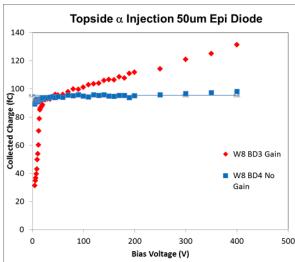
Analysis of Ch 2 (300 µm FZ, BB)





Gain in large (34pF) 50um epi Diode (Marta)

2015 Santander RD50 SCIPP, Sadrozinski, Hartmut F.-W.



Top-side α injection

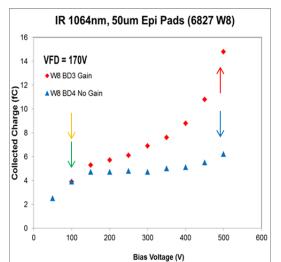
The epi structure prevents electron injection through back-side $\boldsymbol{\alpha}$ radiation.

Using front-side α injection to compare LGAD and no-gain pads requires that the energy loss of the α 's in the active region are the same (or are known well).

The α signal in the no-gain diode is constant as a function of the bias voltage.

While the α signal in the LGAD reaches the same level only after sufficient depth of the p+ implant is depleted (we estimate ~ 7 um) at a bias of ~ 25V.

At higher bias, gain of about g = 1.4 is observed.



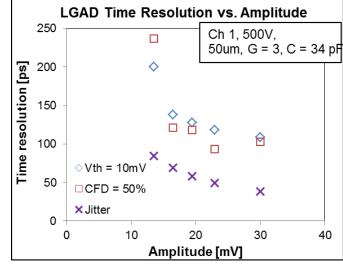
IR Laser Injection

IR (1064nm) laser injection from front shows characteristic LGAD voltage dependence of signal, while no-gain diode is ~ constant above full-depletion voltage VFD = 140 V. Gain of about 3 is observed.

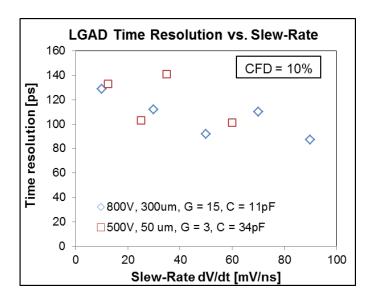
The observed difference of the gain factor for top-side α injection (1.4) wrt laser (g=3) is predicted by the "Weightfield2" program.



Preliminary Beam Test Time resolution



Expected strong dependence on amplitude



Using the slew-rate dV/dt yields unified description of both thin and thick sensor! (amplitudes of 50um epi and 300 FZ LGADs differ by factor 5!)

Surprisingly weak dependence on slew rate: trigger resolution??



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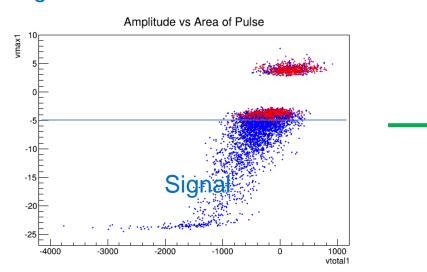
Time Stamp of Pulse Maximum

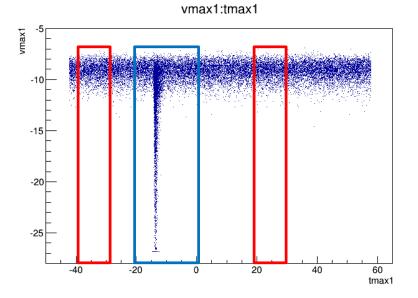
Introduce the time stamp tmax1 of the maximum pulse amplitude vmax1

Beam Data is in a very narrow band:

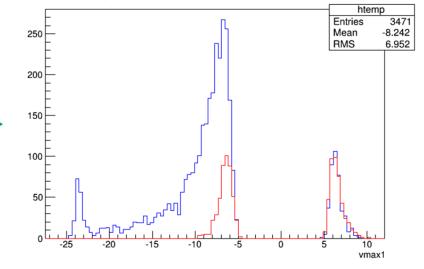
→ To get amplitude spectrum:
 Select signal window -20 < tmax1 < 0 ns
 Subtract noise from siide bands
 -40 < tmax1 < -30 ns + 20 < tmax1 < 30 ns

Introduce pulse area vtotal1 of the pulse. Pulse amplitude vs. Pulse area for signal window and side bands.





Amplitude Distribution of Good and Noise Events

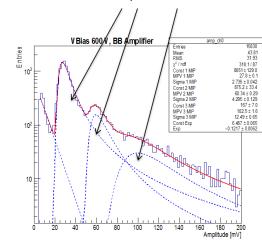


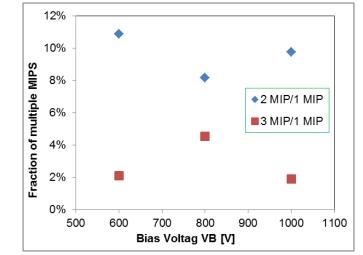


Signal Analysis in terms of MIPs

Extraction of Beam Composition from 300µm FZ sample

The 300um FZ data are well explained as a sum of 3 Landau distributions: 1, 2 and 3 MIPS

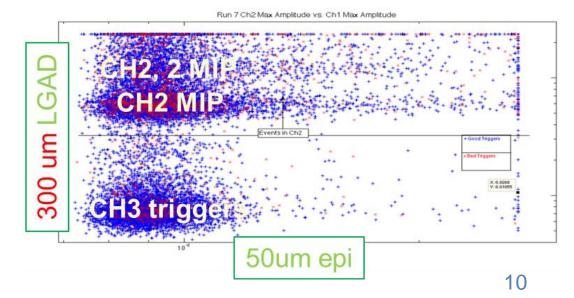




Beam Composition: 88% single MIP

Estimate > 50% of events in 50µm sample above threshold of 5 mV are single MIPs

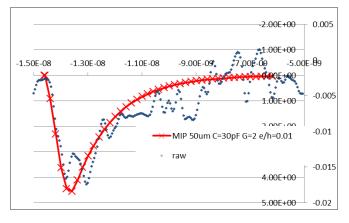
Correlation mainly with Ch 2

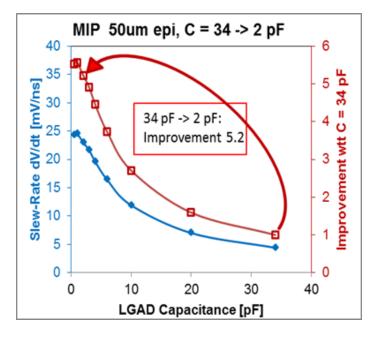


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Effect of Detector Capacitance (50µm epi)





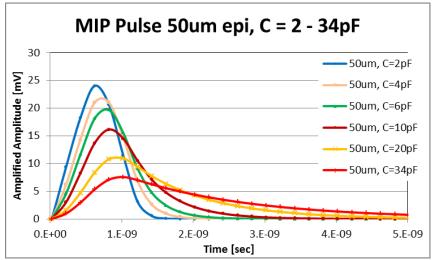


Lower capacitance will improve rise time

 Risetime vs. C

 34pF
 710 ps

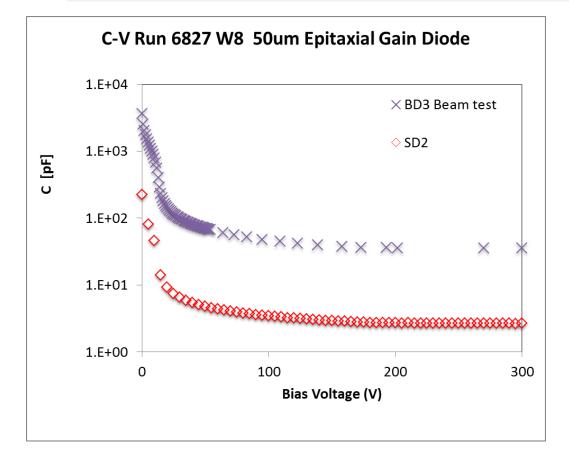
 2 pF
 390 ps



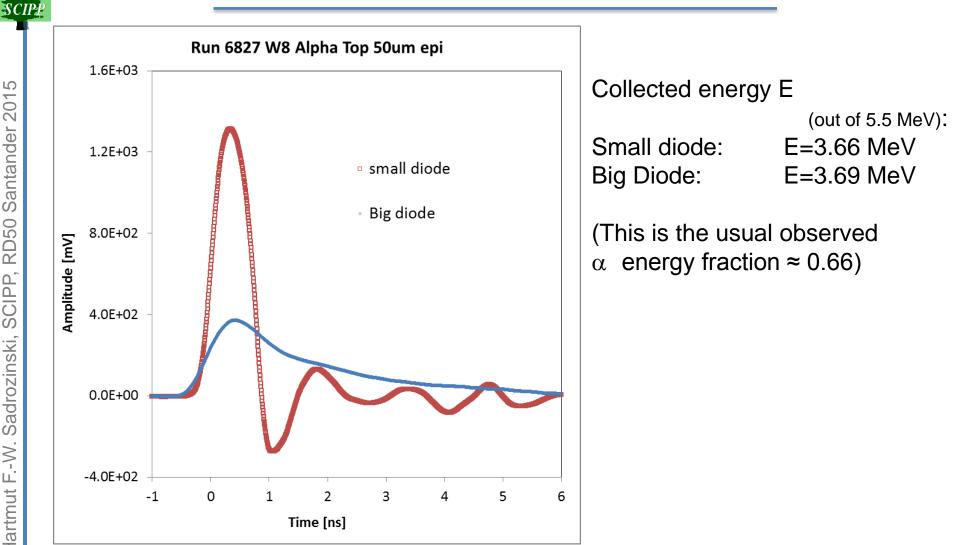
The LGAD diodes with C = 2.5 pF will increase the amplitude by a factor 3 and the slew rate by factor 5.2 wrt beam test diode.

This will improve the time resolution by a factor 5.2

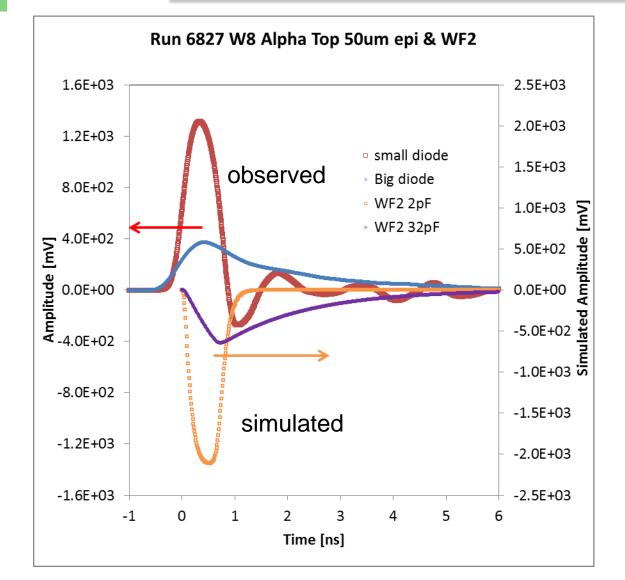




C = 2.5 & 34 pF 50 um epi Sensors: Front α TCT



Front α 50 um epi Sensors: Comparison with WF2



Excellent predictive power of Weightfield2 simulator

SCIP

Front α 50 um epi Sensors: Comparison with WF2

litude [mV

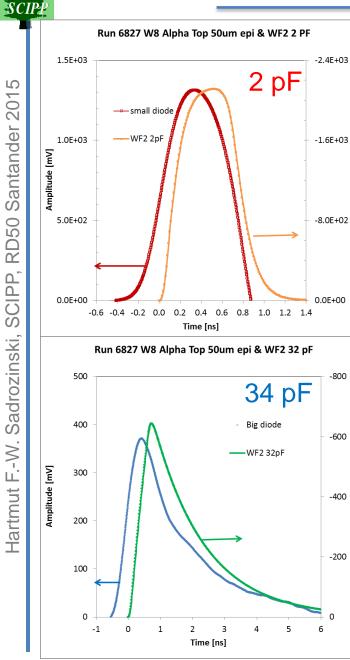
-800

-600

-400

-200

Simulated Amplitude [mV]



Measured Pulse shapes for the small diode (C = 2.5 pF) and big diode (C = 34 pF)compared to WF2 simulation.

Comparisons with Weightfield2 (WF2) are excellent:

(The WF2 simulation is scaled down in both plots by 1.5/2.4 = 0.63, close to the observed α energy fraction)



Thin LGAD promise excellent time resolution.

A 50 μ m epi sample in the Nov 2015 beam test was missing two needed ingredients:

large gain (gain = 3 instead of 10) small capacitance (C = 34 instead of 2.5 pF)

A comparison of the time resolution with a 300 μ m FZ LGAD shows the expected scaling with the slew rate dV/dt.

A small diode (C=2.5pF) is under preparation and we expect improvement of about 3 in S/N which will improve the trigger efficiency about 5 in dV/dt which will improve the time resolution