



New CTEQ–Jefferson Lab (CJ15) analysis of parton distribution functions

Wally Melnitchouk



CTEQ-JLab (CJ) collaboration: <http://www.jlab.org/CJ>

Alberto Accardi, Jeff Owens, Nobuo Sato (theory)

Eric Christy, Thia Keppel, Simona Malace, Peter Monaghan (experiment)

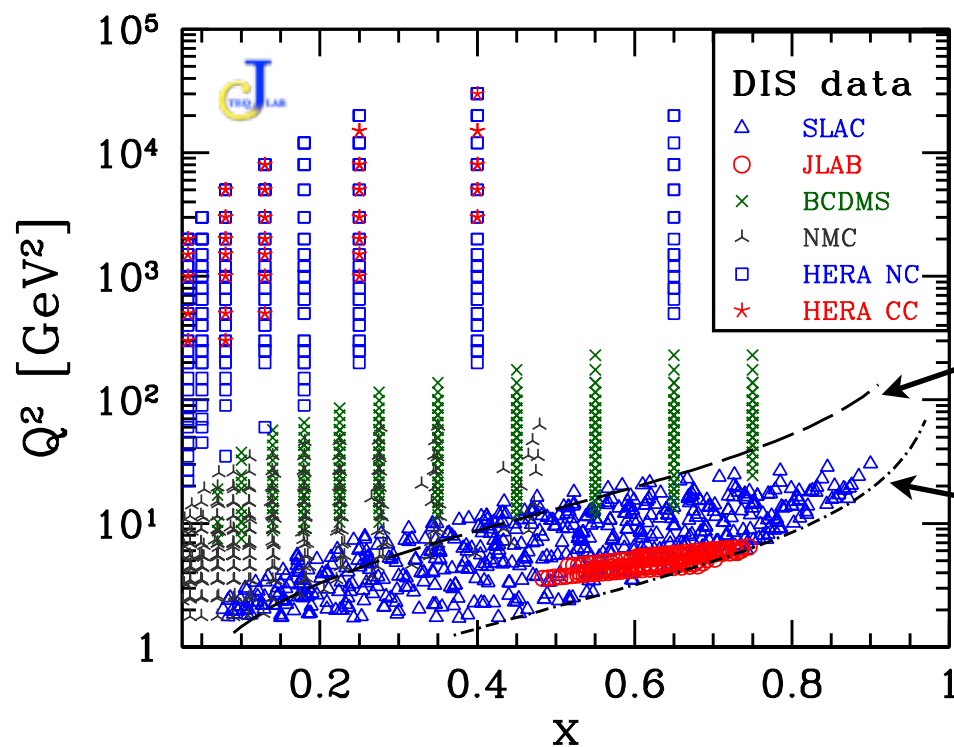
Outline

- CJ PDFs – motivations and goals
- New developments since CJ12
 - more complete treatment of nuclear corrections
 - impact of new lepton & W asymmetry data on d/u
 - inclusion of JLab data
 - analysis of $\bar{d} - \bar{u}$ at large x
- Future plans
 - inclusion of new (LHC & JLab) data
 - Monte Carlo based analysis

- Next-to-leading order (NLO) analysis of expanded set of proton and deuterium data (no heavy nuclei)

→ include high- x region ($x > 0.4$)

- High- x region requires use of data at lower W & Q^2



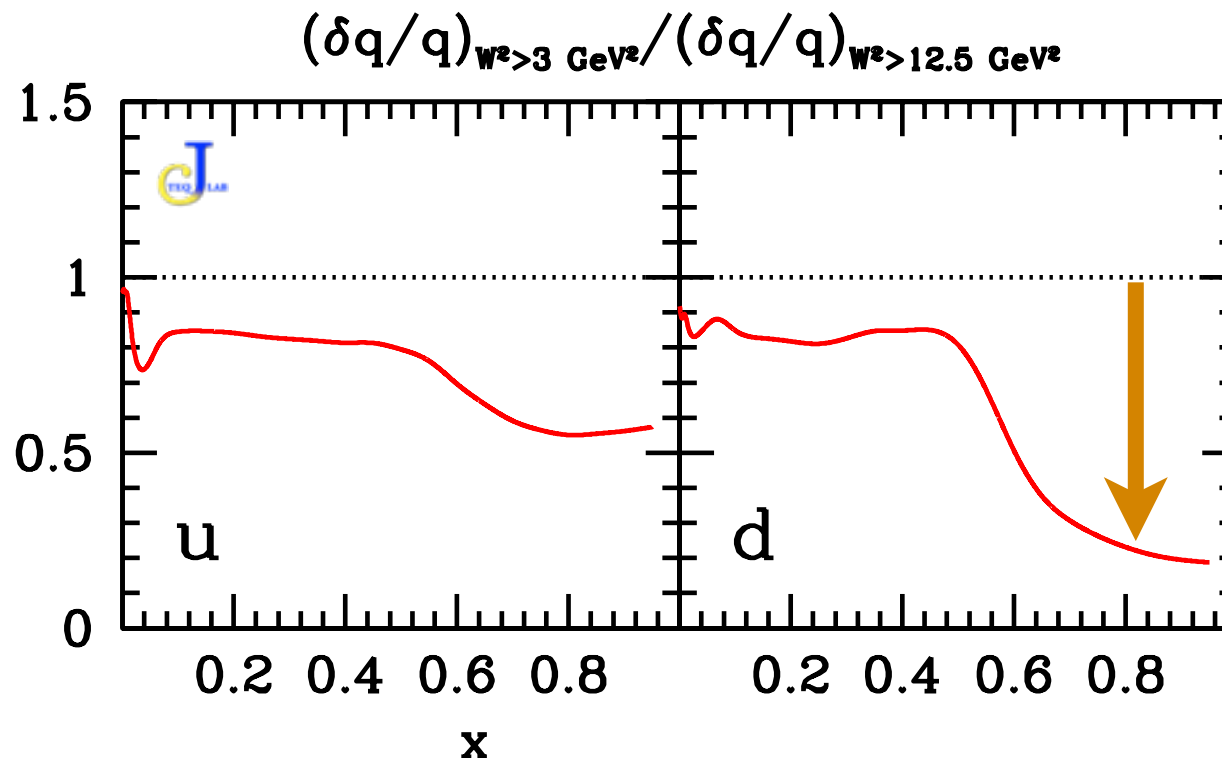
$$W^2 = M^2 + Q^2 \frac{(1-x)}{x}$$

strong cut:
 $Q^2 > 4 \text{ GeV}^2, W^2 > 12.25 \text{ GeV}^2$

weak cut:
 $Q^2 > m_c^2, W^2 > 3 \text{ GeV}^2$

→ factor 2 increase in # of DIS data points when relax strong cut (excludes most SLAC, all JLab data) → weak cut

- Next-to-leading order (NLO) analysis of expanded set of proton and deuterium data (no heavy nuclei)
 - include high- x region ($x > 0.4$)
- High- x region requires use of data at lower W & Q^2



→ significant error reduction at high x

- Next-to-leading order (NLO) analysis of expanded set of proton and deuterium data (no heavy nuclei)
 - include high- x region ($x > 0.4$)
- High- x region requires use of data at lower W & Q^2
- Analysis of high- x data requires careful treatment of subleading $1/Q^2$ corrections
 - target mass corrections, dynamical higher twists
- Correct for nuclear effects in deuteron (binding + off-shell)
 - binding + Fermi motion (well known), nucleon off-shell (less well known)
 - impact on d/u ratio in large- x region

■ CJ15 data sets and χ^2 values

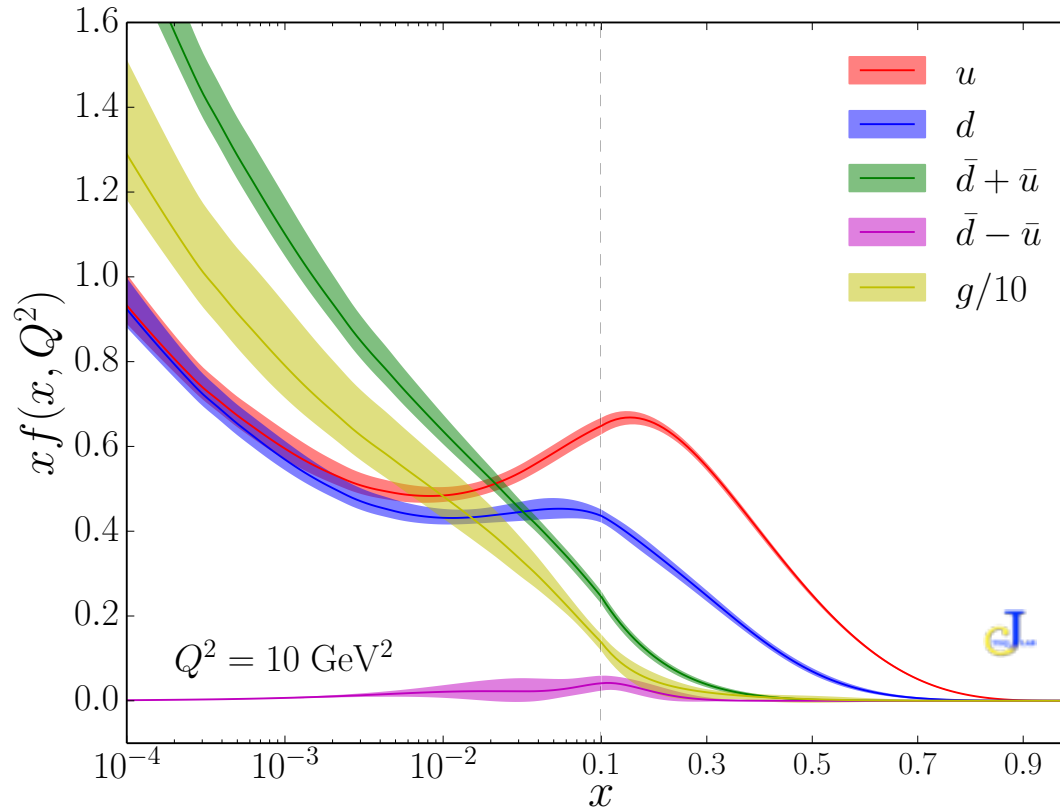
	experiment	# points	χ^2	
			NLO	LO
DIS F_2	BCDMS (p) [23]	351	437	432
	BCDMS (d) [23]	254	294	299
	NMC (p) [24]	275	407	414
	NMC (d/p) [25]	189	172	180
	SLAC (p) [26]	564	435	496
	SLAC (d) [26]	582	372	417
	JLab (p) [27]	136	166	164
	JLab (d) [27]	136	124	127
	JLab (n/d) [85]	191	217	224
DIS σ	HERA (NC e^-p) [28]	145	112	161
	HERA (NC e^+p) [28]	408	541	872
	HERA (CC e^-p) [28]	34	19	19
	HERA (CC e^+p) [28]	34	31	33
Drell-Yan	E605 (pCu) [45]	119	93	104
	E866 (pp) [29]	121	139	155
	E866 (pd) [29]	129	144	191
	E866 (pd/pp) [30]	12	9	9
W/charge asymmetry CDF (e) [31]		11	12	11
	DØ (μ) [32]	10	20	21
	DØ (e) [33]	13	27	56
	CDF (W) [34]	13	15	12
	DØ (W) [35]	14	16	47
Z rapidity	CDF (Z) [36]	28	27	79
	DØ (Z) [37]	28	16	23
jet	CDF (run 2) [39]	72	15	22
	DØ (run 2) [41]	110	21	46
γ +jet	DØ 1 [42]	16	6	20
	DØ 2 [42]	16	15	40
	DØ 3 [42]	12	25	35
	DØ 4 [42]	12	13	77
total		4035	3941	4786
total + norm			3950	4918
	χ^2/dof		0.98	1.22

← BONuS F_2^n / F_2^d

← DØ A_l

← DØ A_W

■ CJ15 PDFs



our “biased” fit form*

$$x f = a_0 x^{a_1} (1 - x)^{a_2} \times (1 + a_3 \sqrt{x} + a_4 x)$$

at $Q^2 = Q_0^2$

* except for d and \bar{d}/\bar{u}

- for strange assume $s = \bar{s} \propto (\bar{d} + \bar{u})$ (no neutrino data)
- charm computed perturbatively (s-ACOT scheme);
no compelling reason for intrinsic charm

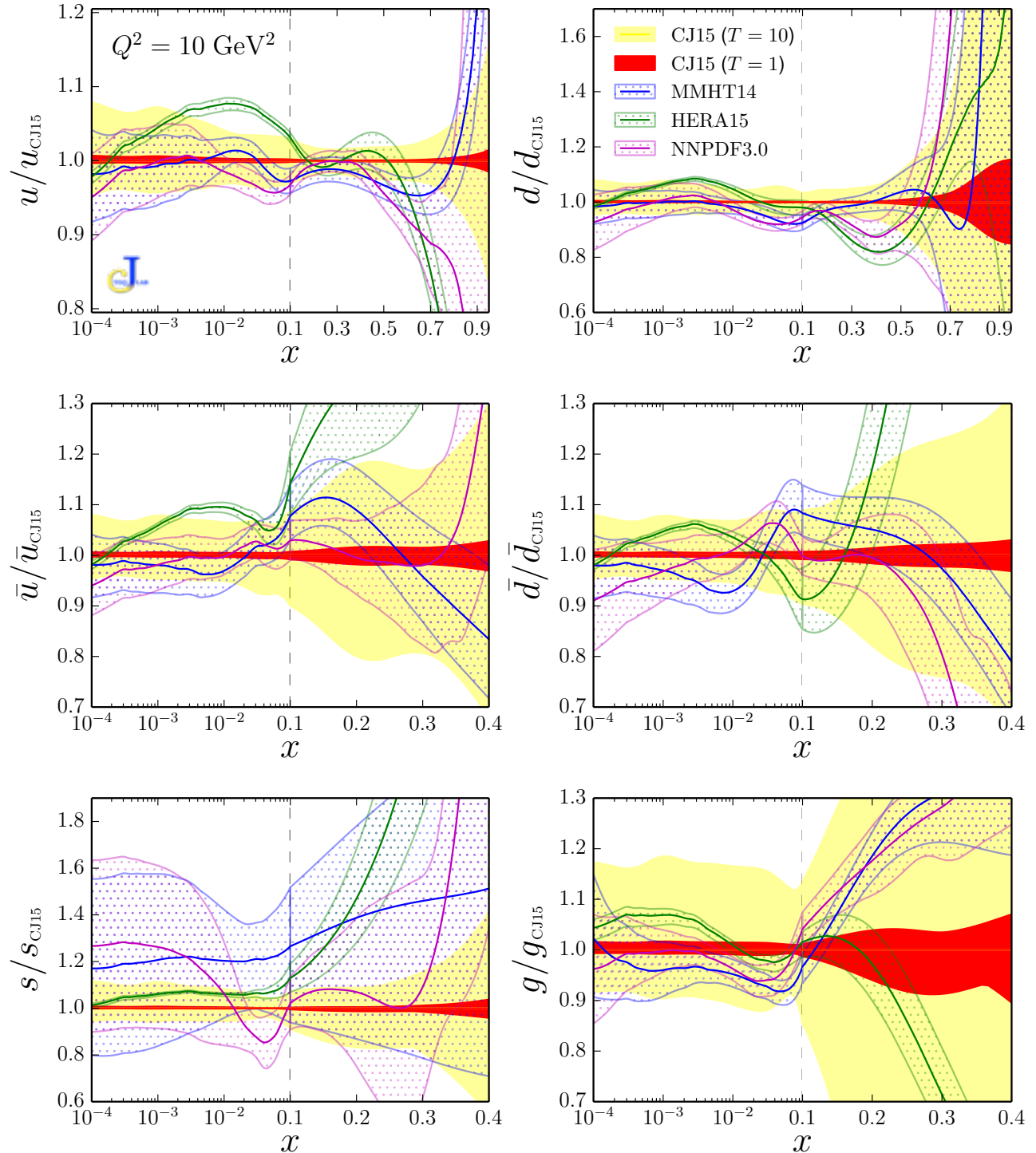
Jimenez-Delgado et al., PRL 114, 082002 (2015)

■ CJ15 vs. other PDFs

→ currently use $\Delta\chi^2 = 1$
but assume tolerance $T=10$ for some app's

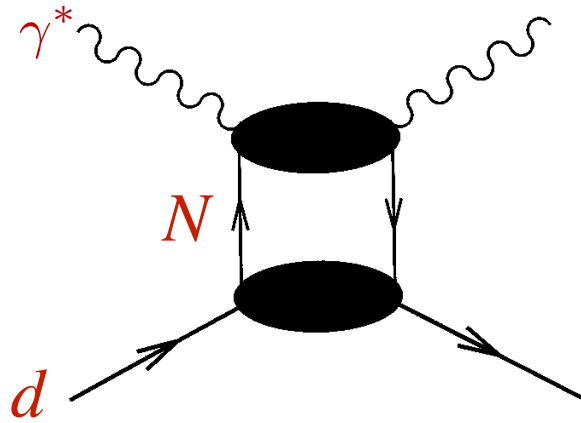
→ HERAPDF15 errors
comparable to $T=1$,
others generally
between $T=1$ & 10

→ larger uncertainty
for d PDF at high x
than for u



Nuclear corrections

- Nuclear structure function at $x \gg 0$ dominated by incoherent scattering from individual nucleons



$$q^d(x, Q^2) = \int \frac{dz}{z} dp^2 f_{N/d}(z, p^2) \tilde{q}^N(x/z, p^2, Q^2)$$

nucleon momentum
distribution in d
("smearing function")

PDF in bound
(off-shell) nucleon

$$\rightarrow z = \frac{p \cdot q}{p_d \cdot q} \approx 1 + \frac{p_0 + \gamma p_z}{M} \left[p_0 = M + \varepsilon, \quad \varepsilon = \varepsilon_d - \frac{\vec{p}^2}{2M} \right]$$

momentum fraction of d carried by N

\rightarrow at finite Q^2 , smearing function depends on $\gamma = \sqrt{1 + 4M^2 x^2 / Q^2}$

Nuclear corrections

- Expand off-shell nucleon PDF about on-shell ($p^2 = M^2$) limit

$$\tilde{q}^N(x, p^2) = q^N(x) \left[1 + \frac{(p^2 - M^2)}{M^2} \delta q^N(x) \right]$$

$$\delta q^N = \left. \frac{\partial \log \tilde{q}^N}{\partial \log p^2} \right|_{p^2=M^2}$$

- Deuteron PDF sum of on- and off-shell contributions

$$q^d = q^{d(\text{on})} + q^{d(\text{off})}, \text{ where}$$

$$q^{d(\text{on})}(x, Q^2) = \int \frac{dz}{z} f^{(\text{on})}(z) q^N(x/z, Q^2)$$

$$q^{d(\text{off})}(x, Q^2) = \int \frac{dz}{z} f^{(\text{off})}(z) \delta q^N(x/z, Q^2) q^N(x/z, Q^2)$$

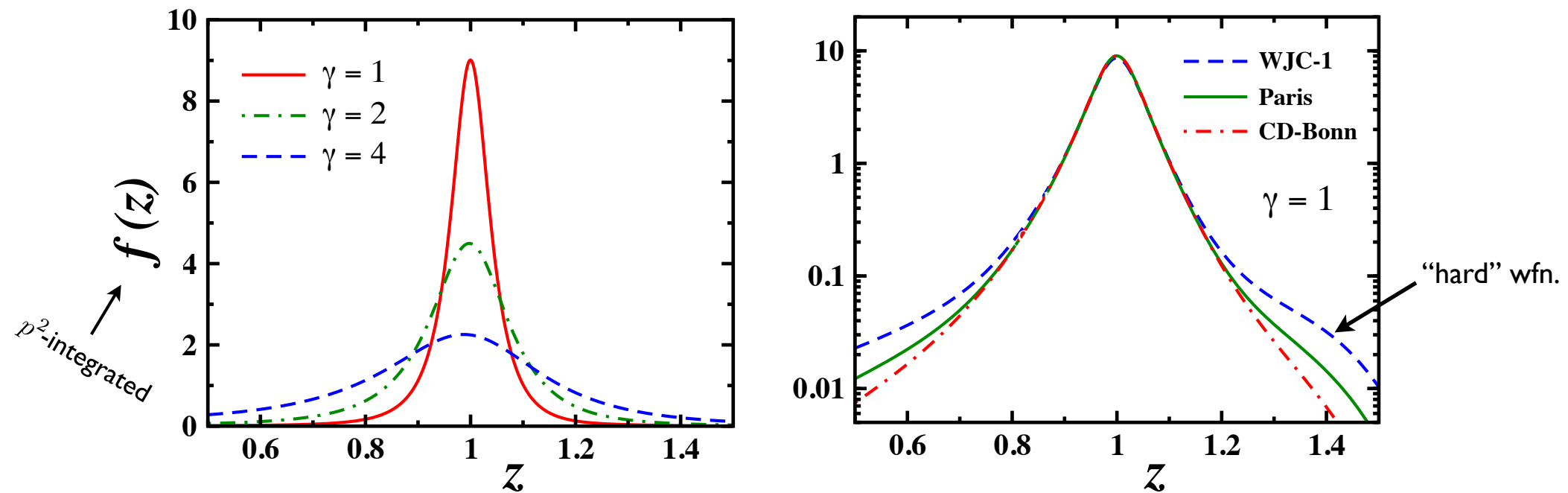
on-shell & off-shell smearing functions

$$\begin{aligned} &\rightarrow f^{(\text{on})}(z) = \int dp^2 f_{N/d}(z, p^2) \\ &\rightarrow f^{(\text{off})}(z) = \int dp^2 \frac{p^2 - M^2}{M^2} f_{N/d}(z, p^2) \end{aligned}$$

Nuclear corrections

- Smearing function in the deuteron computed in “weak binding approximation” – expand in powers of \vec{p}^2/M^2

$$f_{N/d}(z, p^2) = \frac{1}{(2\pi)^3} \frac{1}{\gamma^2} \left[1 + \frac{\gamma^2 - 1}{z^2} \left(1 + \frac{2\varepsilon}{M} + \frac{\vec{p}^2}{2M^2} (1 - 3\hat{p}_z^2) \right) \right] |\psi_d(p)|^2$$



- effectively more smearing for larger x and lower Q^2
- greater wave function dependence at large z (\rightarrow large x)

Nuclear corrections

CJ12

- Nucleon off-shell correction to quark PDF

→ off-shell covariant quark “spectator” (OCS) model

$$\tilde{q}^N(x, p^2) = \int d\hat{p}^2 \Phi_{q/N}(\hat{p}^2, \Lambda(p^2))$$

momentum distribution of quarks
with virtuality \hat{p}^2 in bound nucleon

→ scale parameter $\Lambda(p^2)$ suppresses large- p^2 contributions

→ off-shell “rescaling” parameter $\lambda = \frac{\partial \log \Lambda^2}{\partial p^2}$ fitted

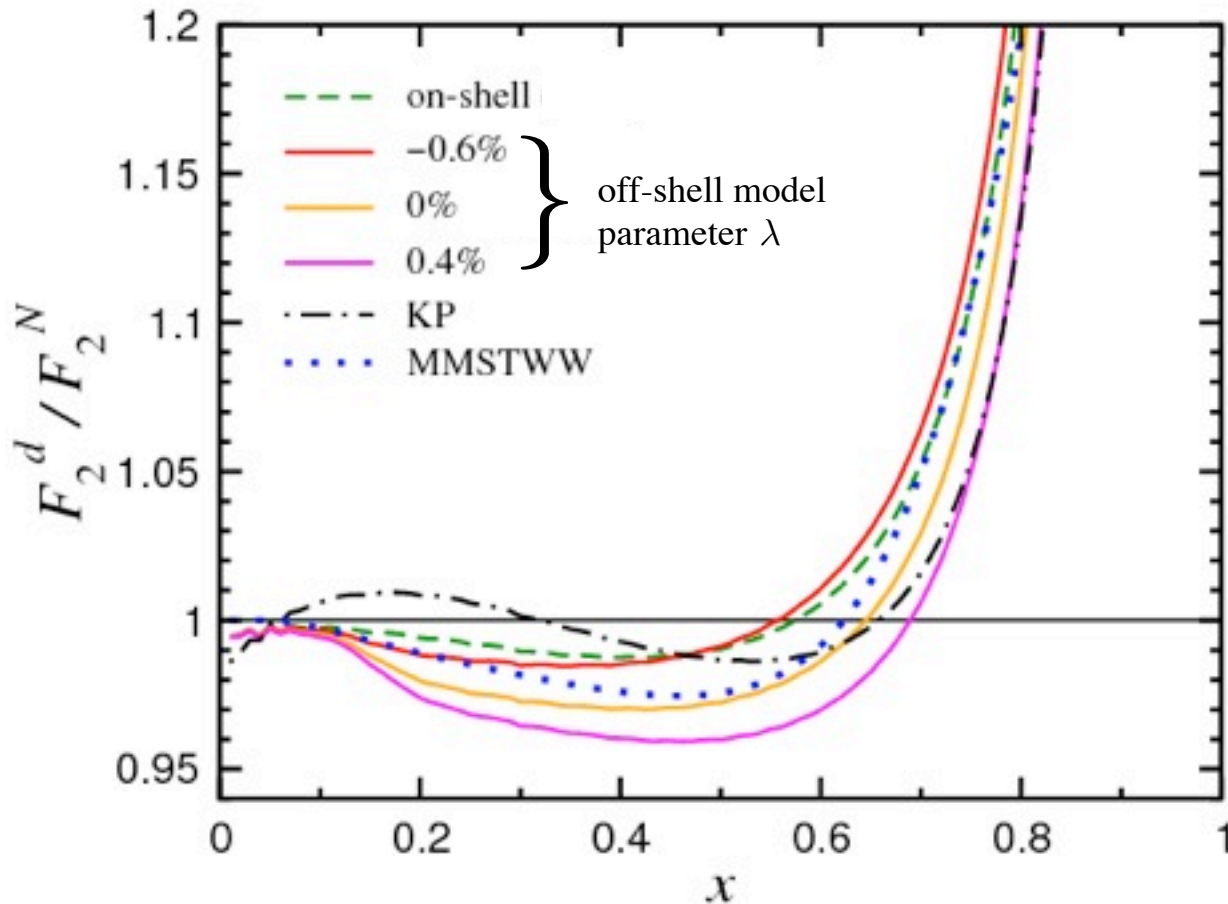
→ applied to valence, antiquark & gluon PDFs

Owens et al., PRD 87, 094012 (2013)

Ehlers et al., PRD 90, 014010 (2014)

Nuclear corrections

- Nucleon off-shell correction to quark PDF



- larger off-shell effects for larger λ , and for KP model
- enhancement (“antishadowing”) at $x \sim 0.2$ in KP model

Kulagin, Petti
NPA 765, 126 (2006)

Nuclear corrections

CJ15

- Nucleon off-shell correction to quark PDF

→ alternatively, parametrize δq^N phenomenologically

$$\delta q^N = C_N(x - x_0)(x - x_1)(1 + x - x_0)$$

Kulagin, Petti, NPA 765, 126 (2006)

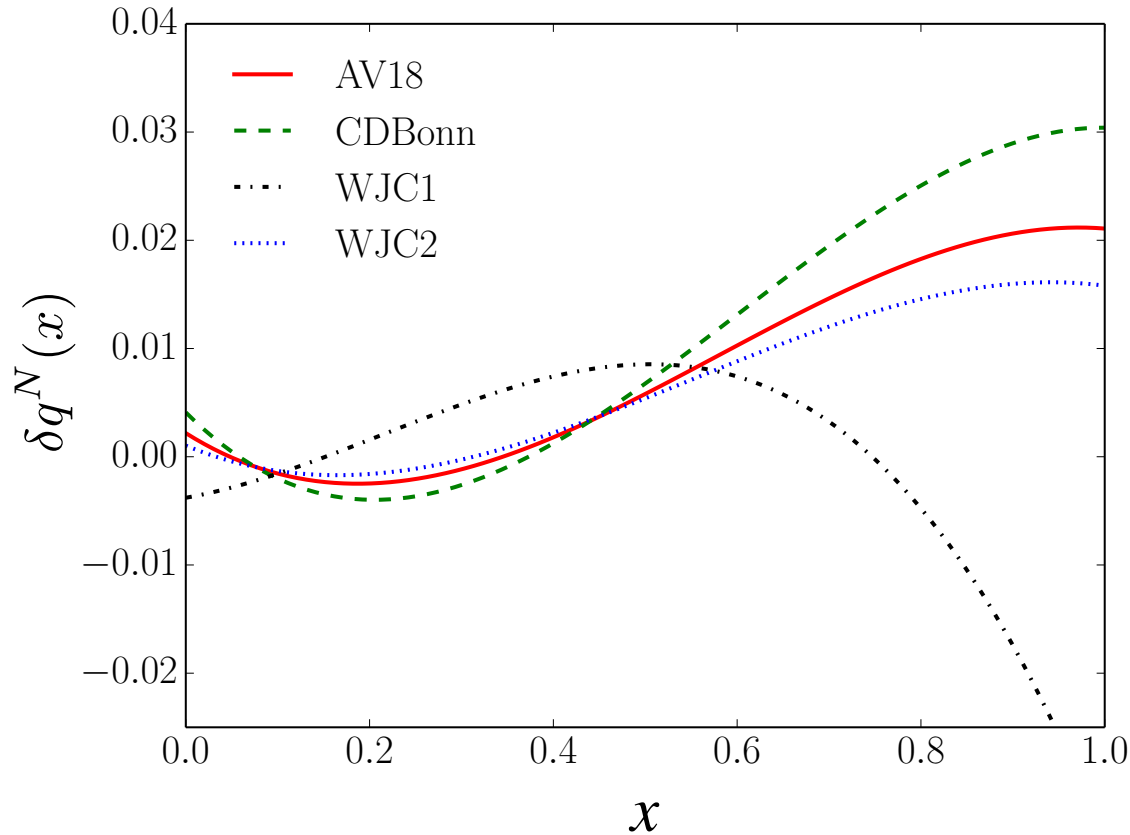
→ fit 2 of $\{C, x_0, x_1\}$ for given deuteron wave function;
fix third parameter from normalization condition

$$\int_0^1 dx \delta q^N(x) \left(q^N(x) - \bar{q}^N(x) \right) = 0$$

→ similar to Kulagin-Petti, but fitted only to *deuteron* data
(avoid uncontrolled extrapolations from large- A data)

Nuclear corrections

- Nucleon off-shell correction to quark PDF



deuteron wave functions from NN scattering

AV18: *Wiringa et al., PRC 51, 38 (1995)*

CD-Bonn: *Machleidt, PRC 63, 024001 (2001)*

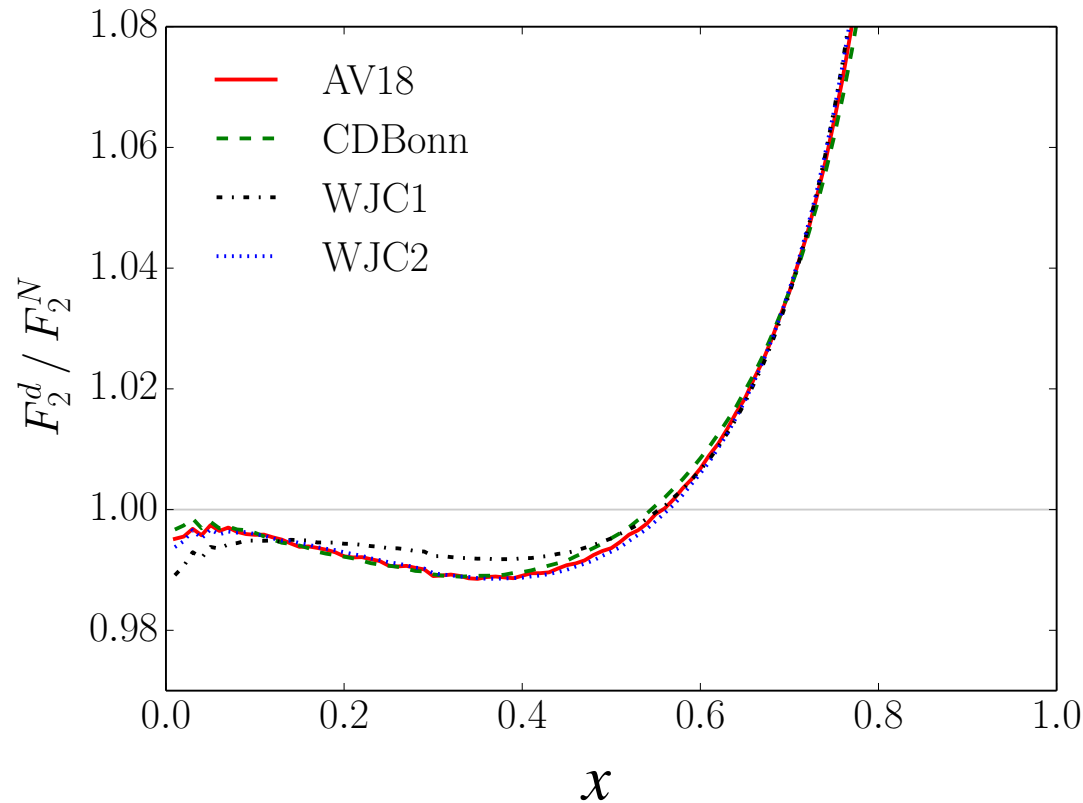
WJC: *Gross, Stadler, PRC 82, 034004 (2010)*

(long-distance part of wfns. is model independent, from chiral symmetry)

→ fitted off-shell corrections weakly dependent on wave function, except for WJC-1 (hardest momentum distribution – largest tail)

Nuclear corrections

- Nuclear EMC ratio in deuteron

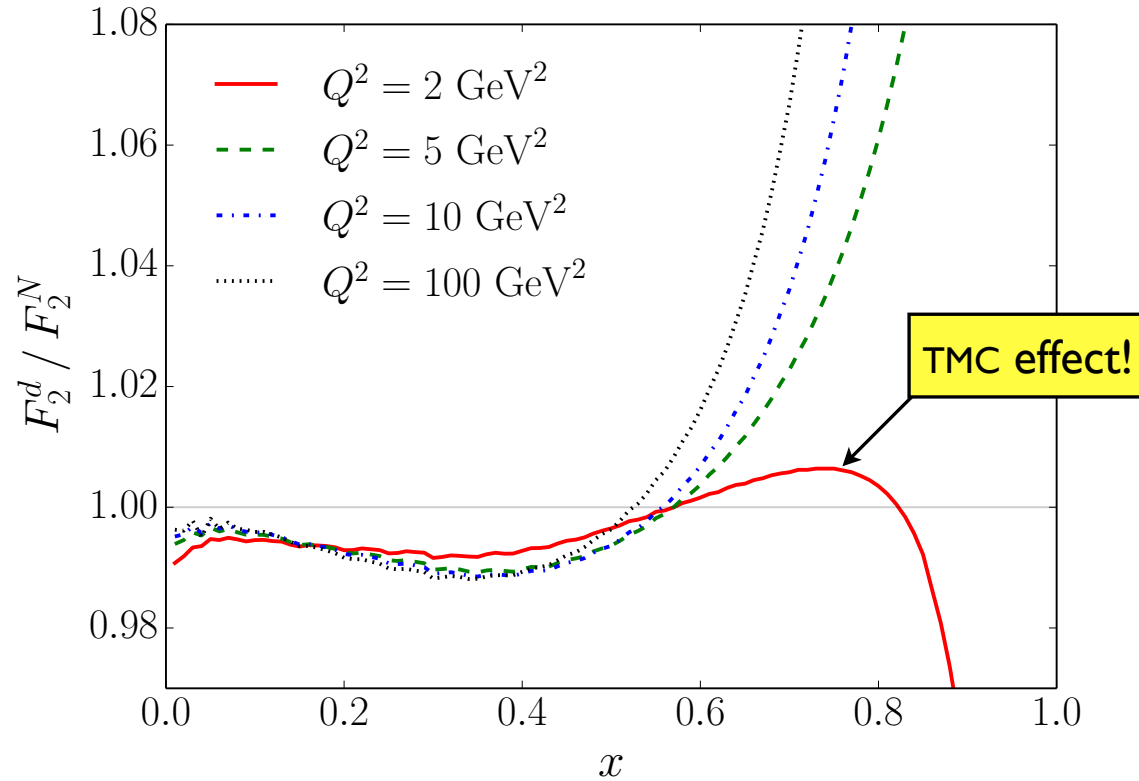


→ observables sensitive only to combined smearing *and* off-shell corrections

→ no evidence for antishadowing at $x \sim 0.1$

Nuclear corrections

- Nuclear EMC ratio in deuteron

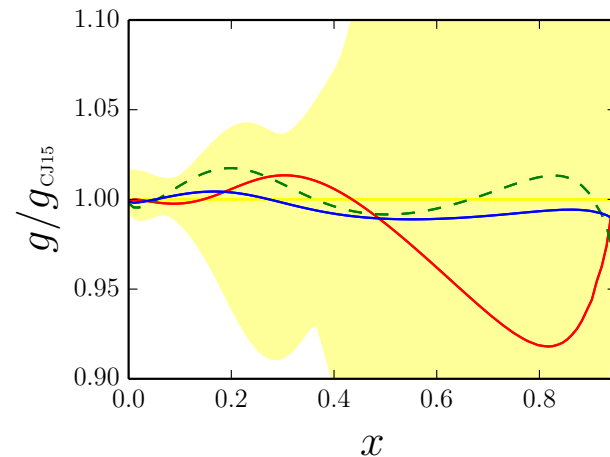
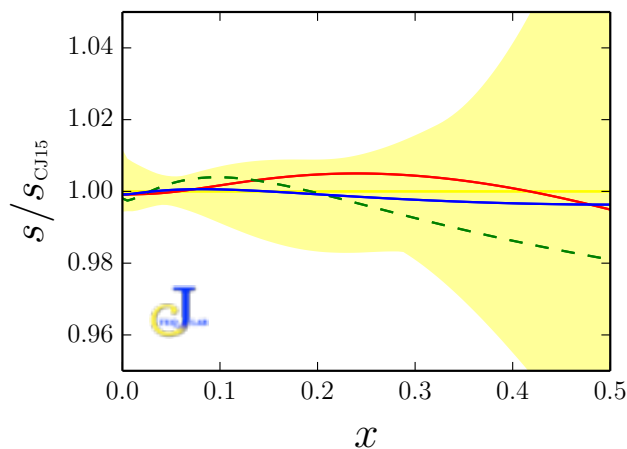
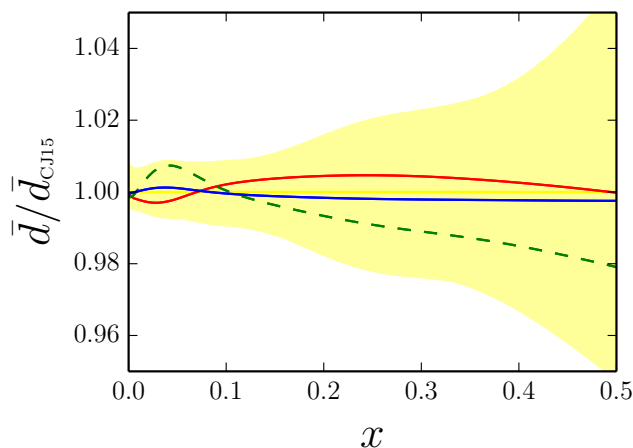
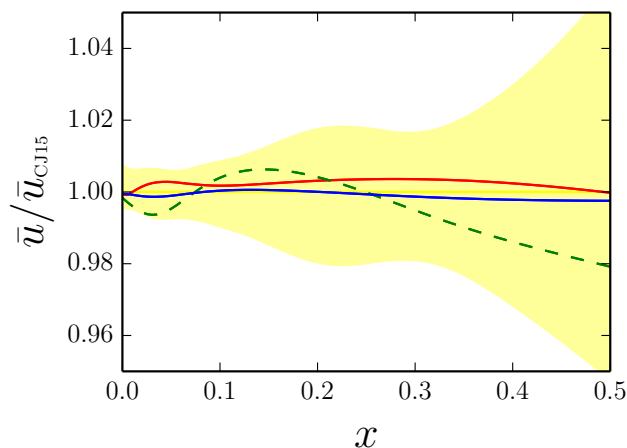
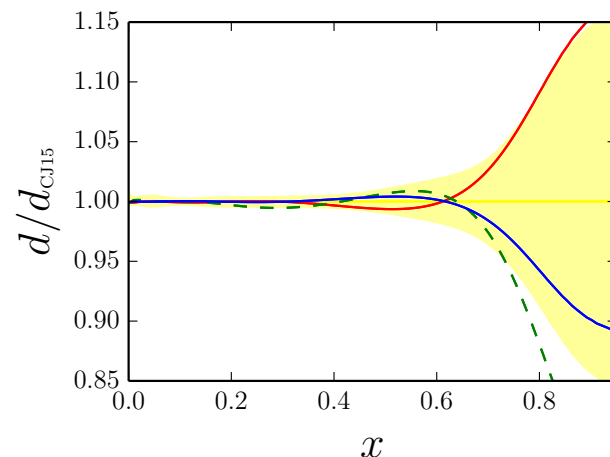
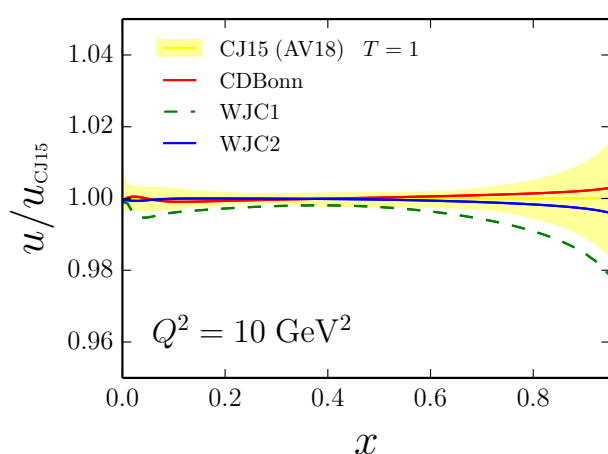


→ ratio has significant Q^2 dependence at low Q^2 from target mass effects

Nuclear corrections & CJ15 PDFs

→ results for PDFs using all wfns. other than WJC-1 are within 1σ

→ WJC-1 outside 1σ error for d at high x & u at low & high x

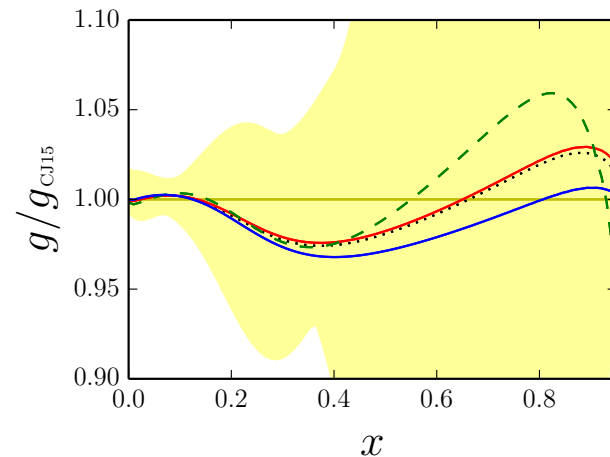
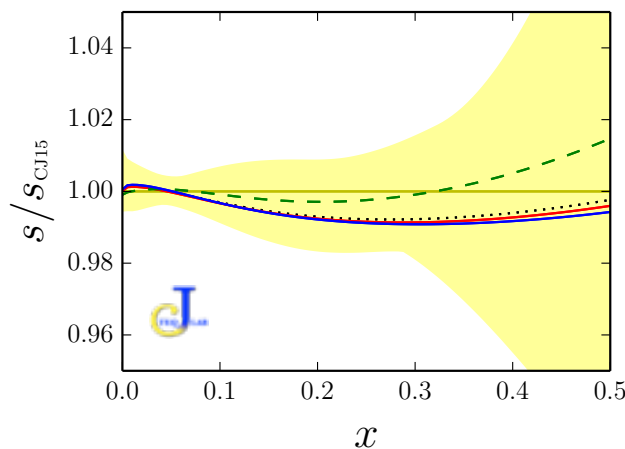
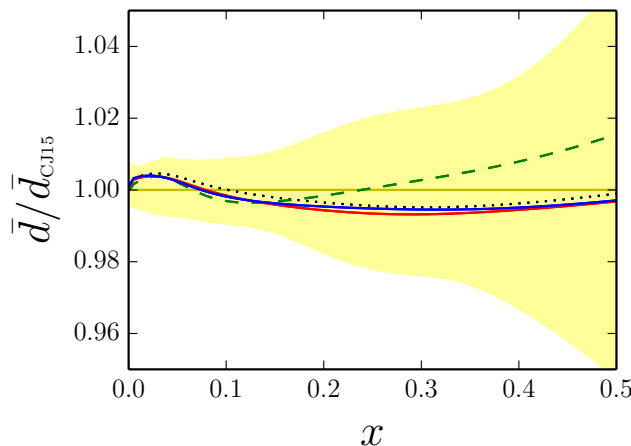
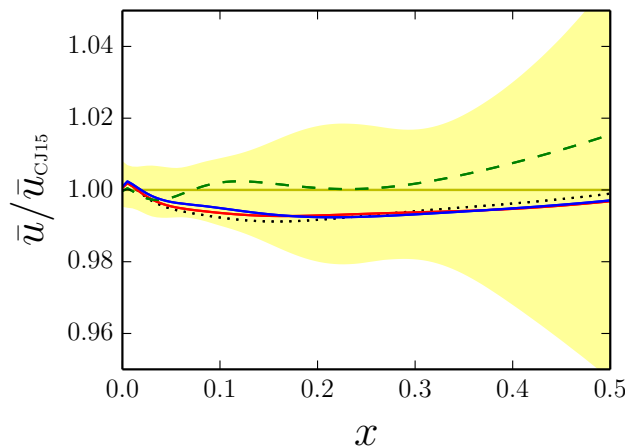
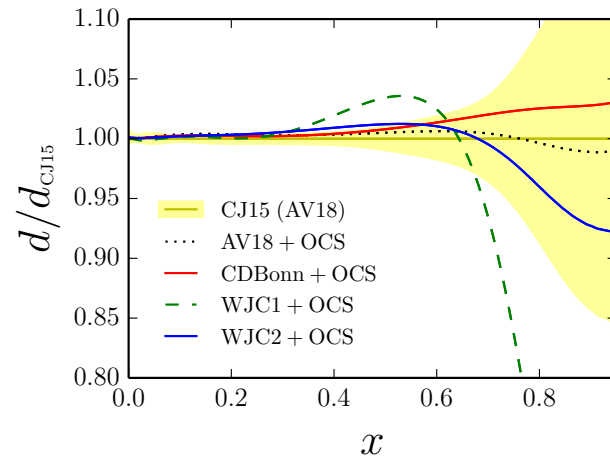
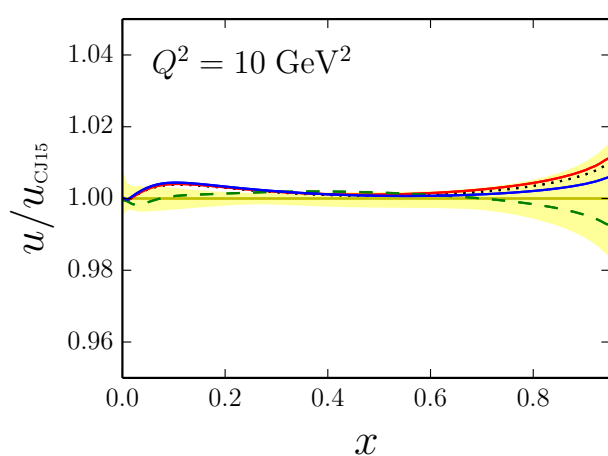


Nuclear corrections & CJ15 PDFs

→ results for PDFs using all wfns. other than WJC-1 are within 1σ

→ WJC-1 outside 1σ error for d at high x & u at low & high x

→ off-shell model dep. (cf. OCS) generally weak except for d PDF



Nuclear corrections & CJ15 PDFs

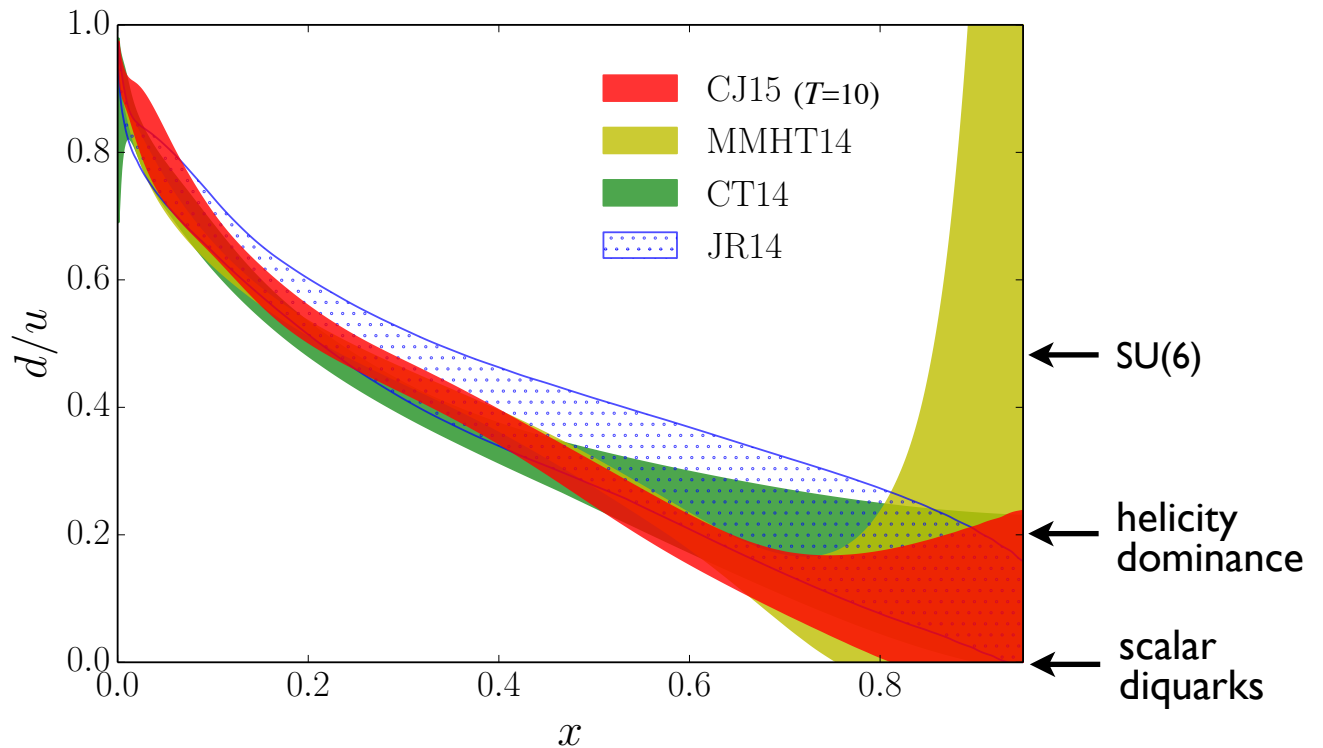
→ d/u ratio at high x
of interest for
nonperturbative
models of nucleon

→ more flexible
parametrization

$$d \rightarrow d + b x^c u$$

allows finite,
nonzero $x = 1$ limit

(standard PDF form gives
 0 or ∞ unless $a_2^d = a_2^u$)



MMHT14: fitted deuteron correction,
“standard” d parametrization

CT14: flexible d parametrization,
no nuclear corrections

JR14: similar deuteron correction,
no lepton/ W asymmetry data

Nuclear corrections & CJ15 PDFs

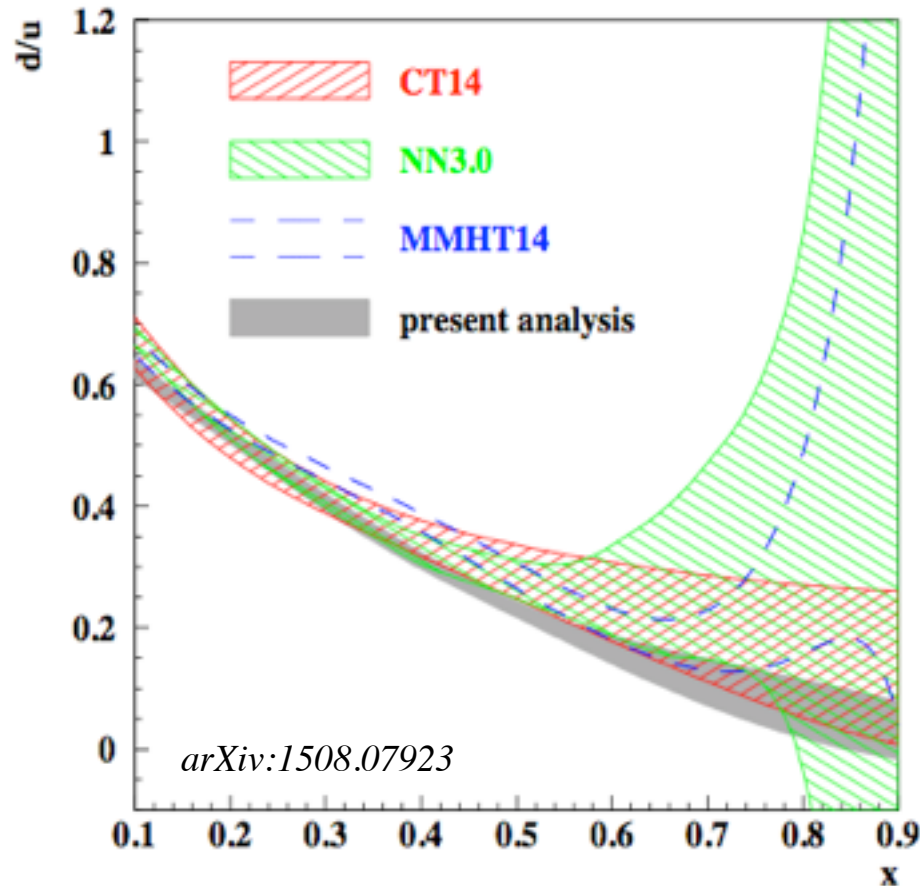
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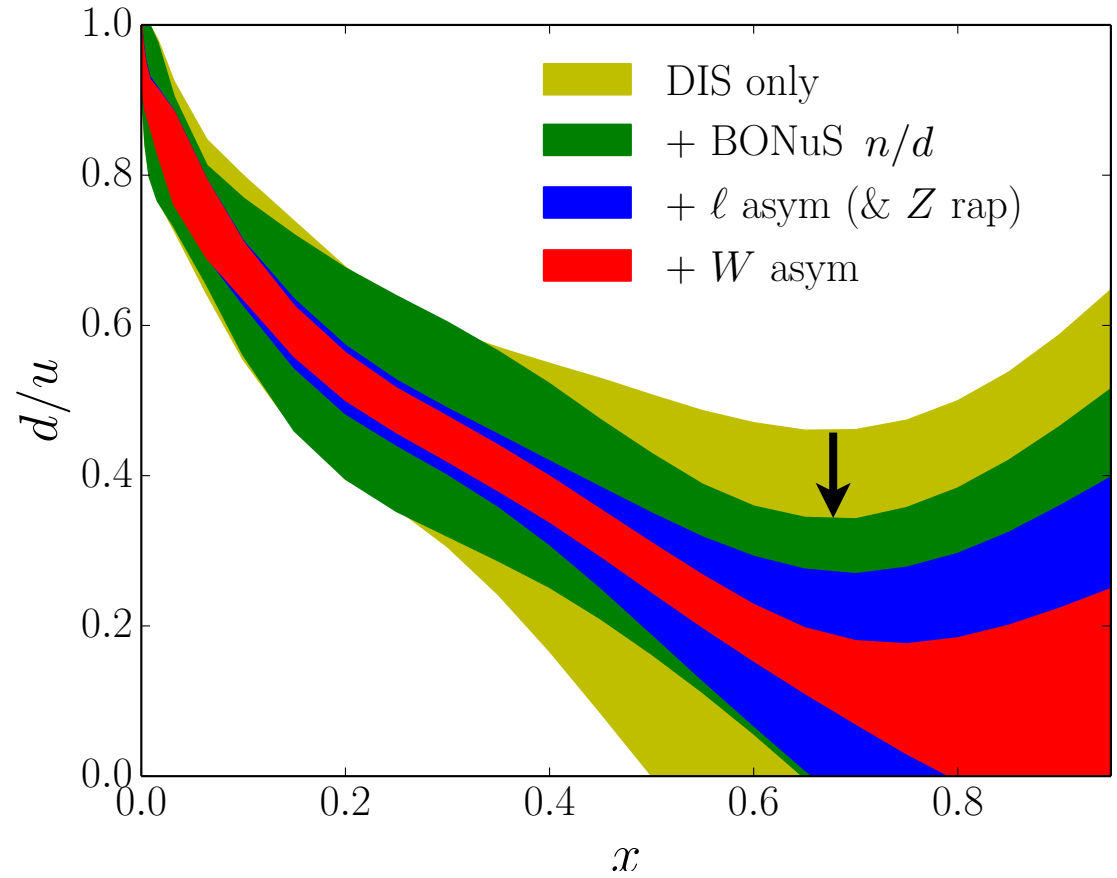
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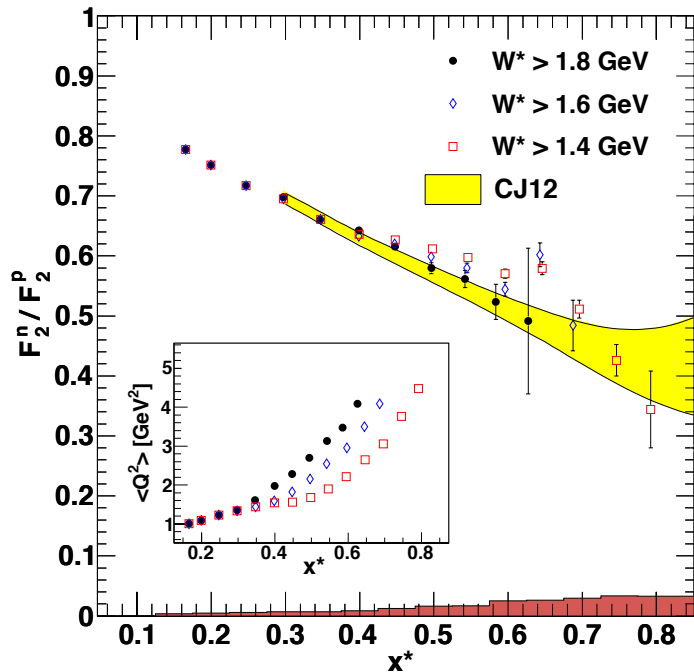
ABMP15: includes new lepton asymmetries,
K-P nuclear corrections

Effect of data sets of d/u

→ new JLab (BONuS) data
reduces error at $x \sim 0.6$



BONuS: spectator proton tagging
in semi-inclusive ed DIS

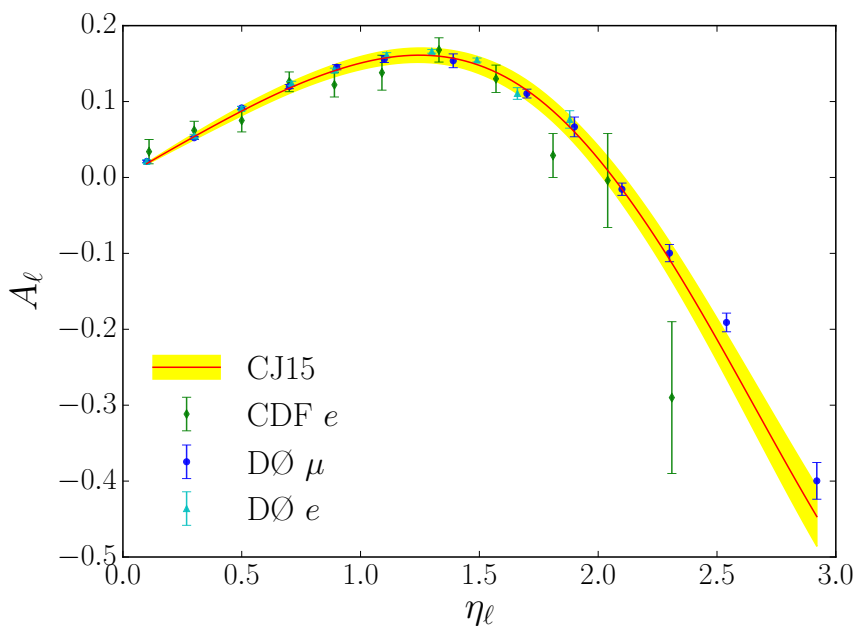
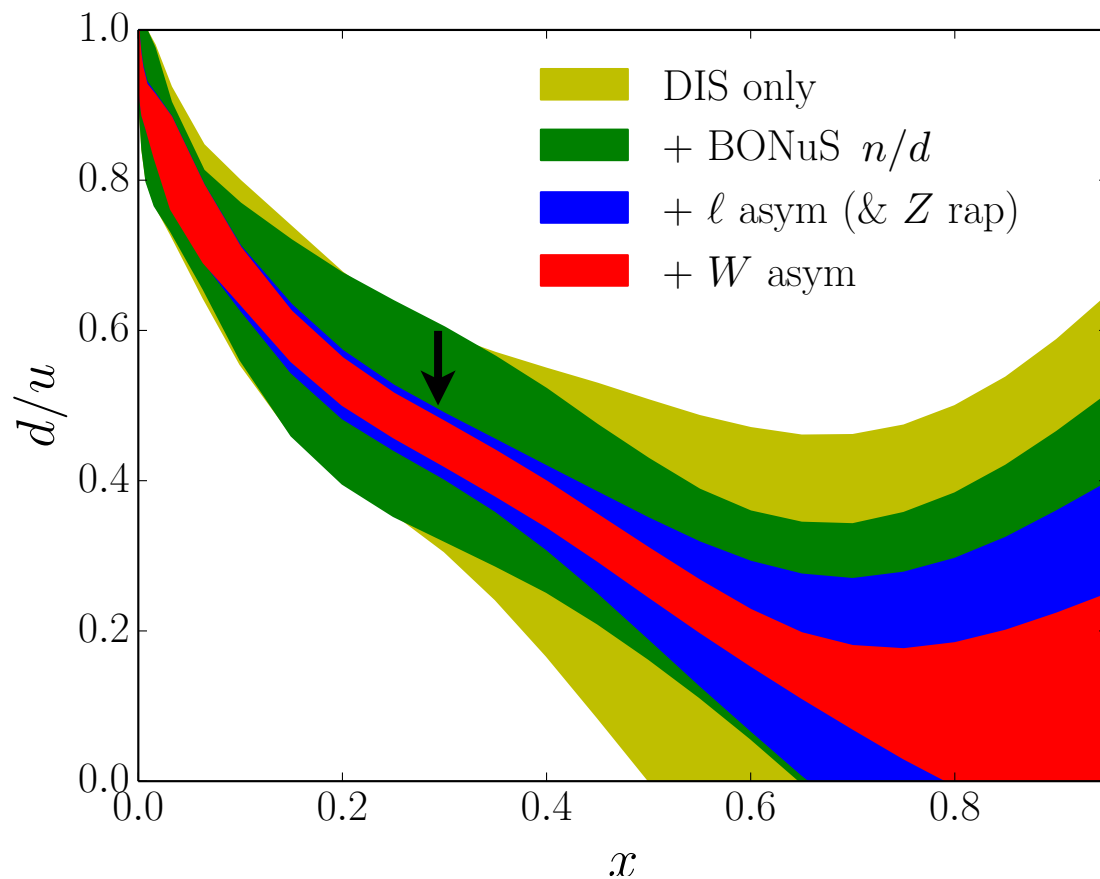


Baillie et al.
PRL 108, 142001 (2012)

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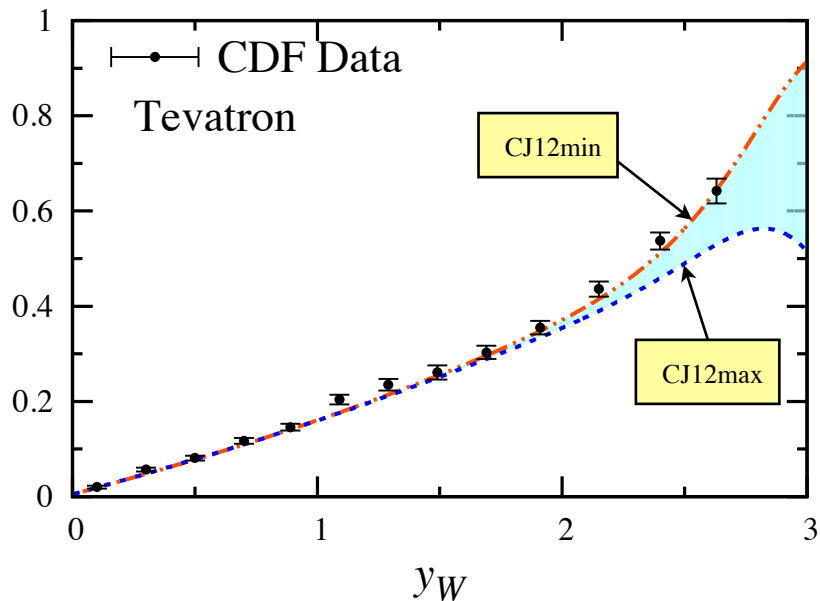
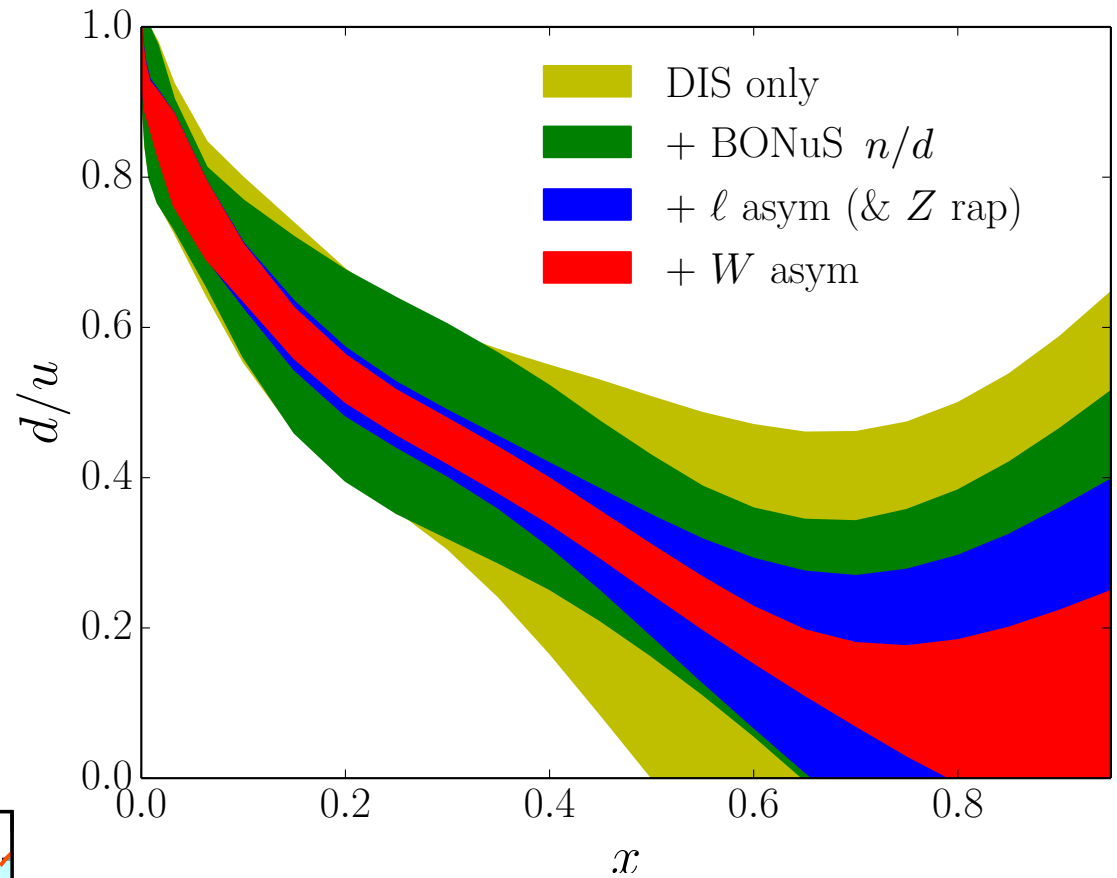
→ significant reduction
from new lepton
asymmetry data
(little effect from Z rap. data)



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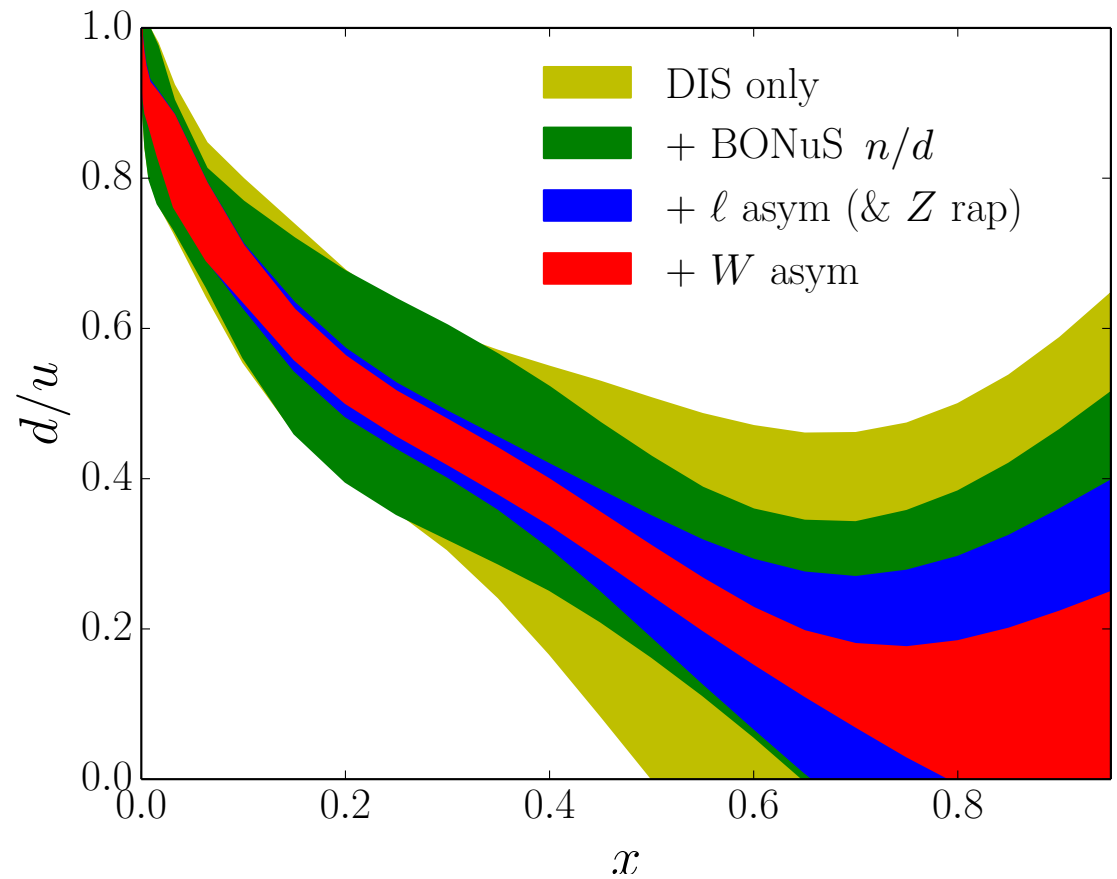
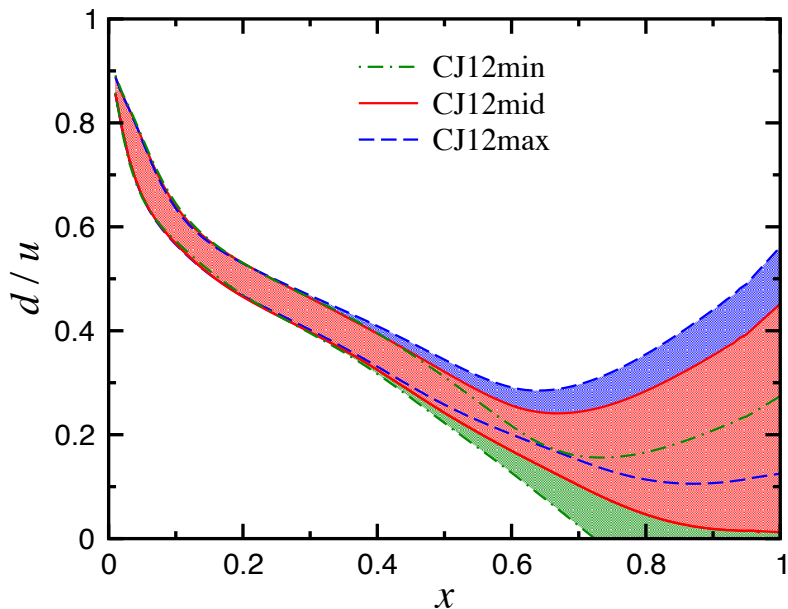
→ W asymmetry at large W rapidity
more sensitive to d/u at high x

→ earlier CDF data preferred smaller
("CJ12min") nuclear corrections

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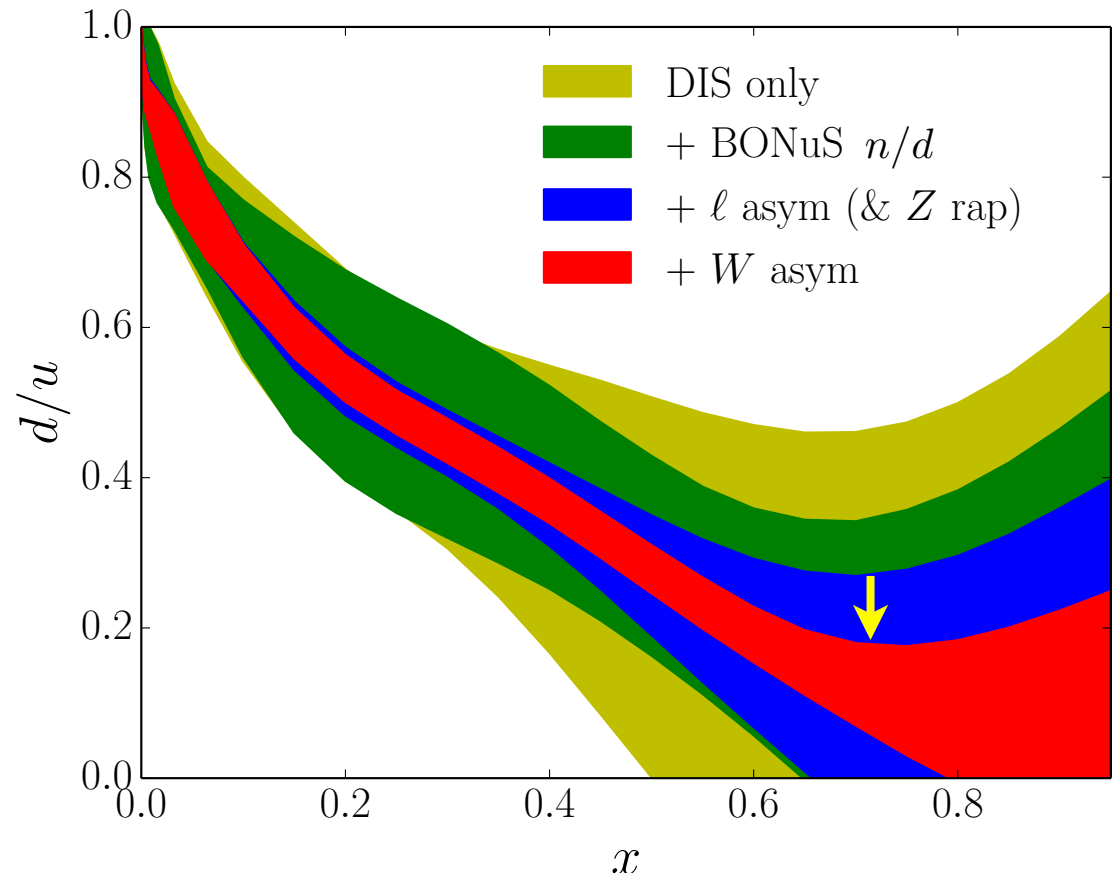
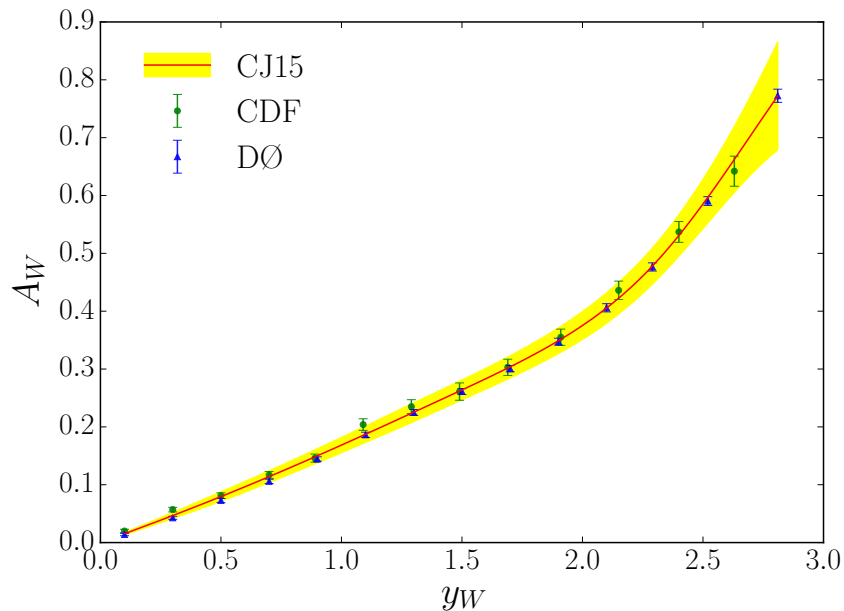
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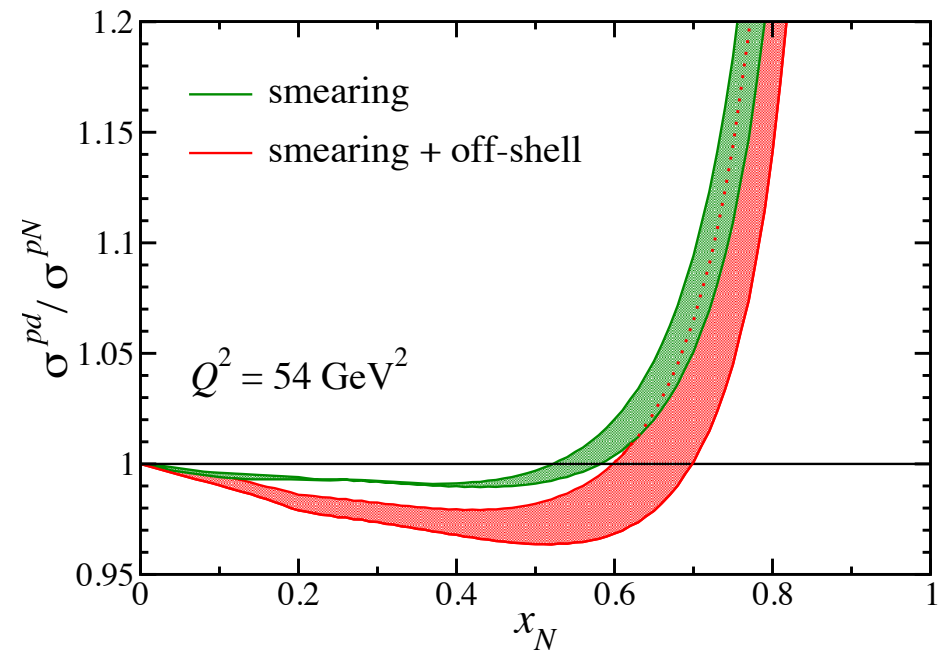
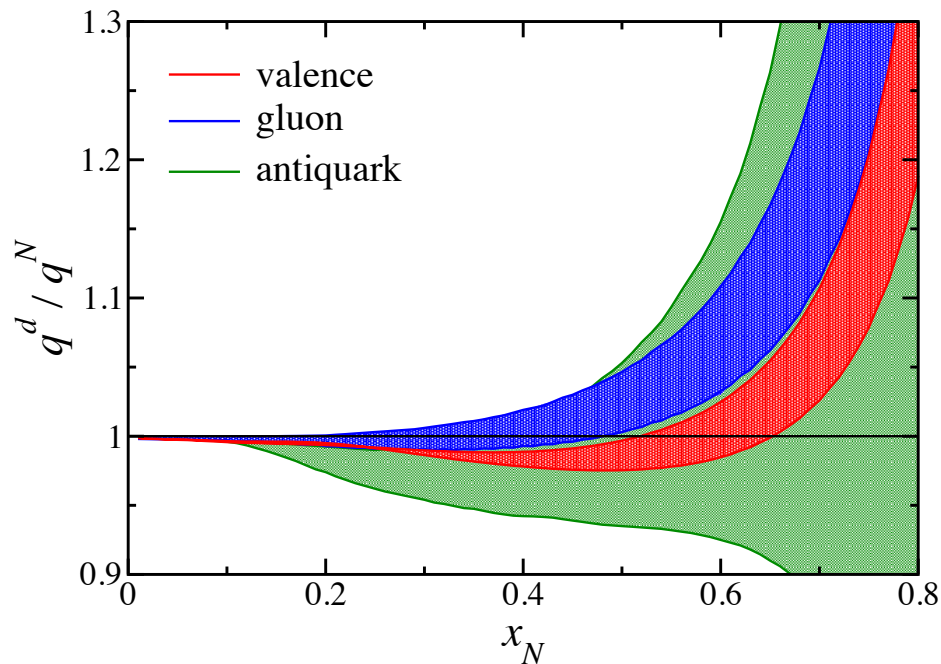


→ W asymmetry at large W rapidity
more sensitive to d/u at high x

→ new D0 data reduce uncertainties at
 $x \sim 0.6 - 0.7$, strongly favor models with
small (but nonzero) nuclear corrections

Light antiquark sea

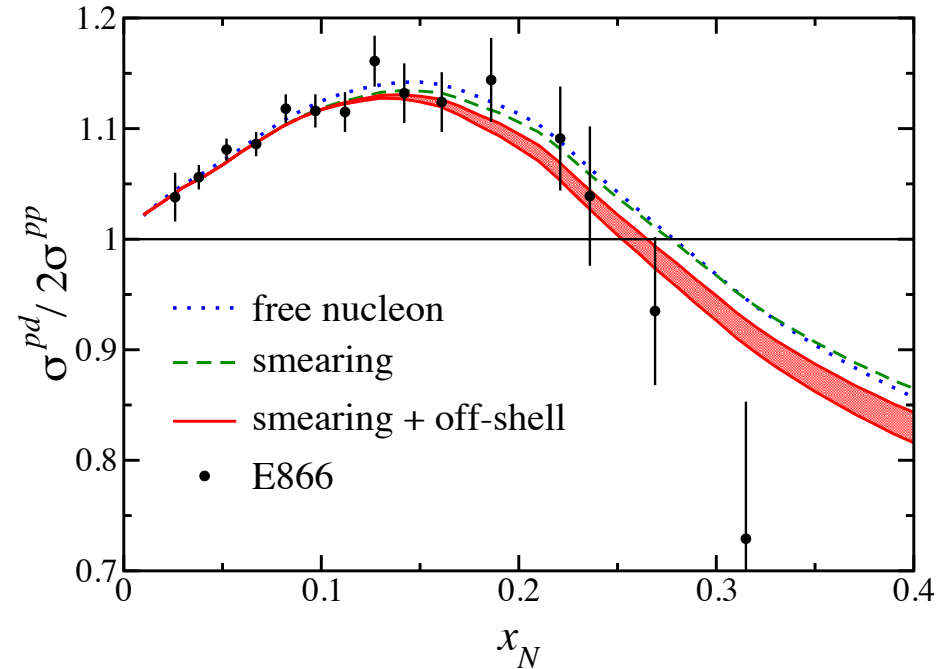
- Flavor asymmetry $\bar{d} - \bar{u}$ in proton sea constrained mostly by FNAL E866 $\sigma^{pd} / \sigma^{pp}$ Drell-Yan data
- Recently nuclear corrections to pd data have been computed in same framework (smearing + off-shell) as in DIS
 - requires nuclear modifications in antiquark and gluon PDFs



Ehlers et al., PRD 90, 014010 (2014)

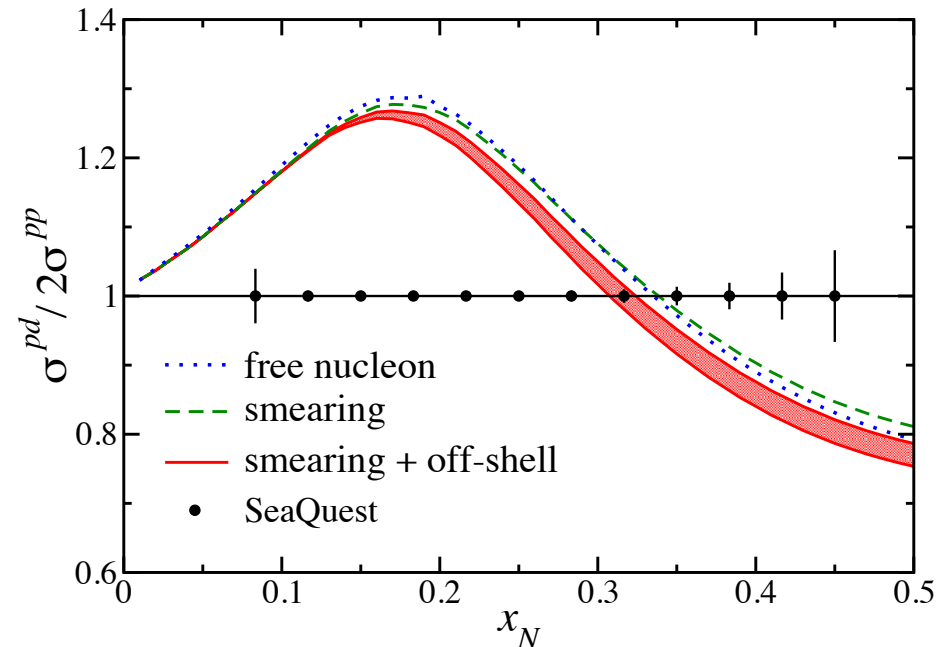
Light antiquark sea

- Modest effect at E866 kinematics, given large errors at high x



- More important effects expected at E906/SeaQuest kinematics at $x > 0.2$

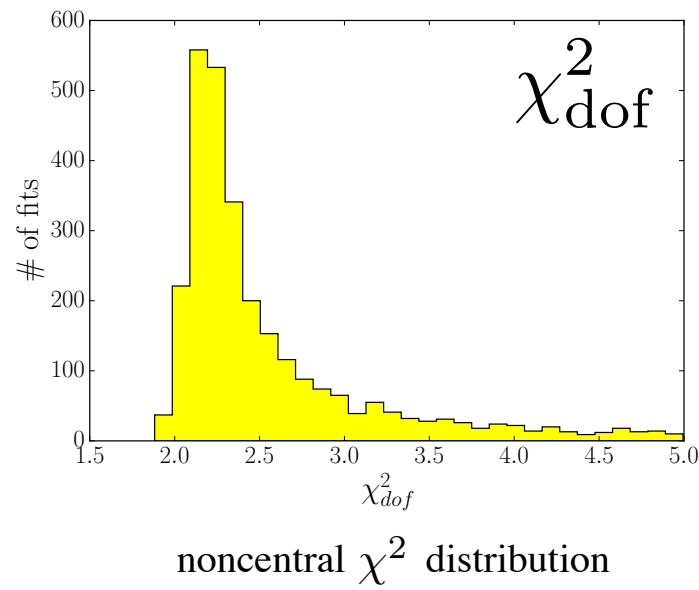
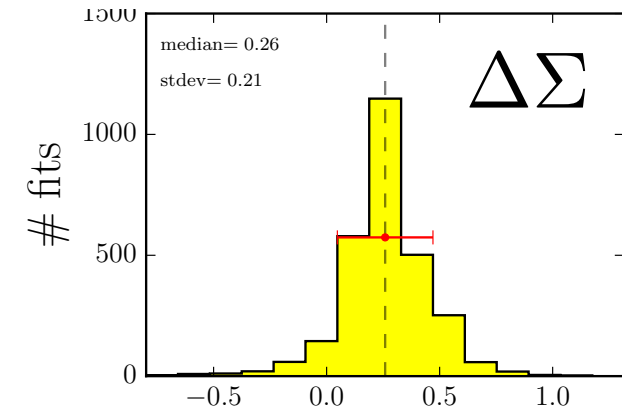
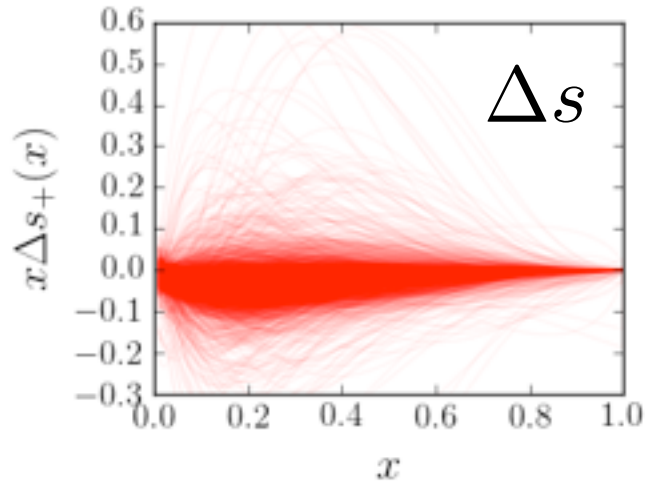
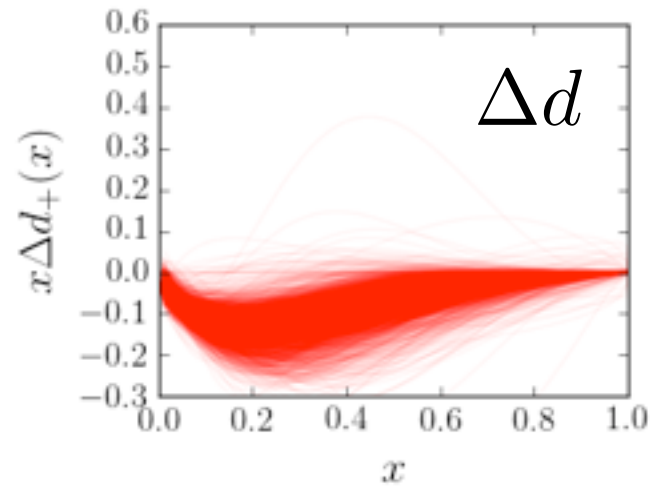
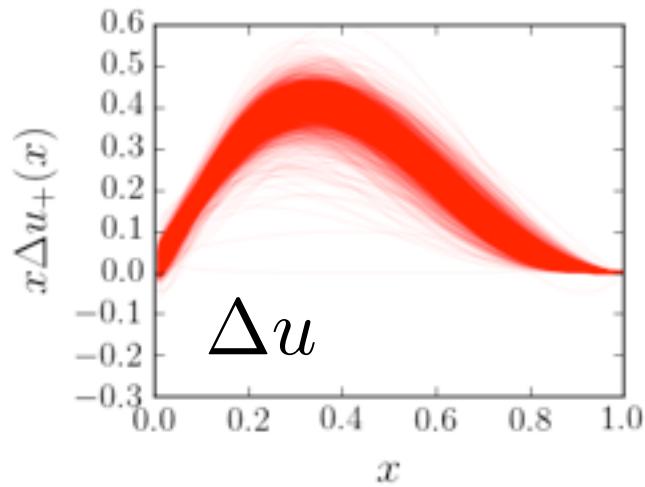
→ relevant for possible sign change of $\bar{d} - \bar{u}$ at high x ?



Summary & outlook

- **New CJ15 PDFs will be released soon** (~ autumn 2015)
 - constraints on large- x PDFs from new data on lepton & W -asymmetries; first JLab F_2^n data
 - reduced nuclear uncertainties on d quark *cf.* CJ12, with smaller d/u ratio for $x \rightarrow 1$ (smaller nuclear corrections)
 - treatment of nuclear corrections in deuteron extended to sea quarks and gluons
(important for pd Drell-Yan cross sections for \bar{q} , and for F_L)
- **Future direction: minimize input parameter bias through Monte Carlo analysis**
 - explore parameter space through MC sampling
 - no ambiguity in “tolerance criteria” (or assumptions about Gaussian errors)
 - cross validation (random data partition) & bootstrap (data resampling)

JAM (JLab Angular Momentum)
polarized PDF analysis:
Nobuo Sato et al. (2015)



- $\sim 10,000$ fits,
 $\sim 3,000$ data points
(~ 40 parameters)
- fast implementation
in Mellin space
- easily extendable to
unpolarized sector

Thank you

