t tbar + isolated photon production at NLO accuracy matched with PS

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in collaboration with

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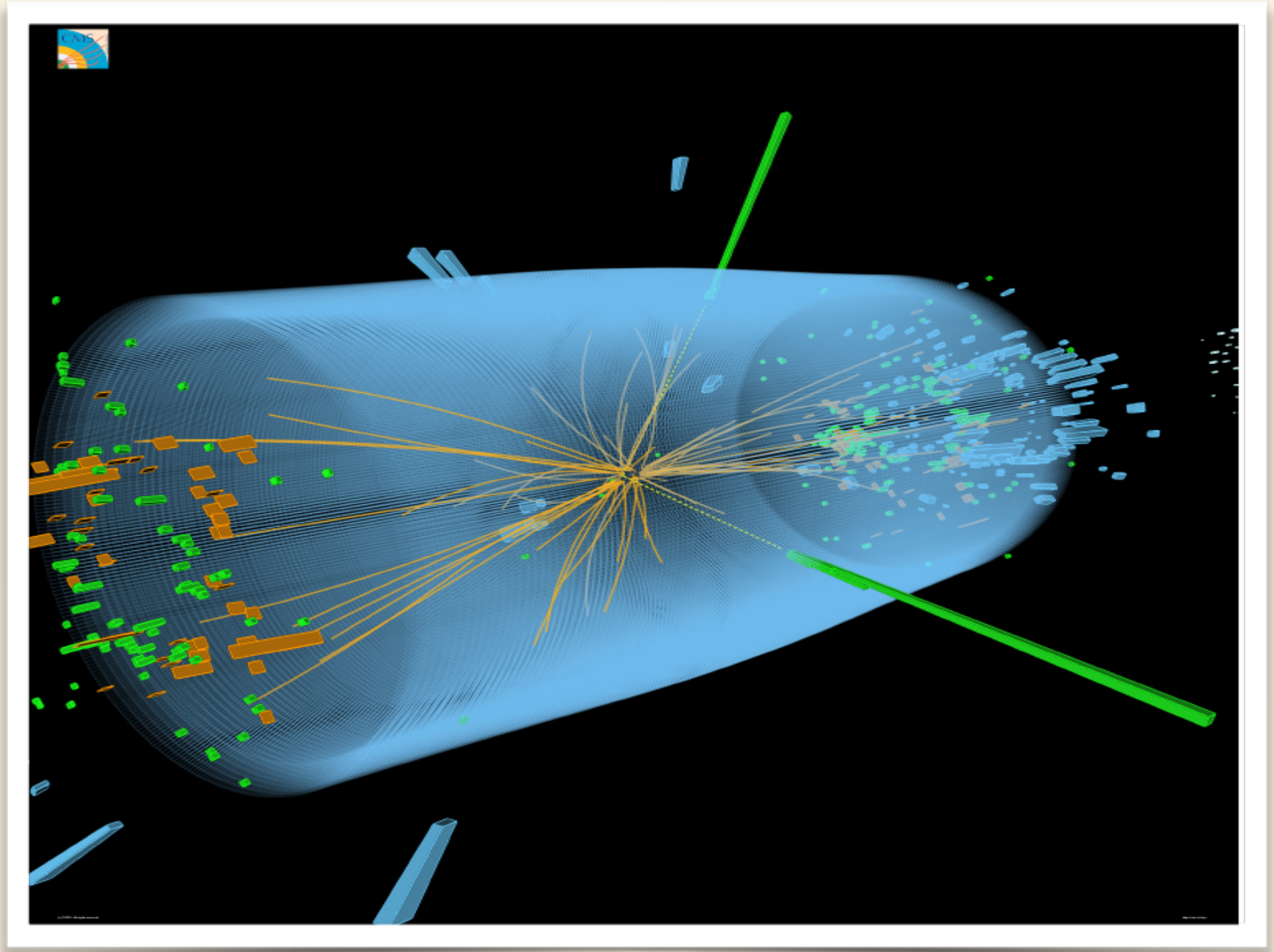
Higgs boson has been discovered



one of the two best channels

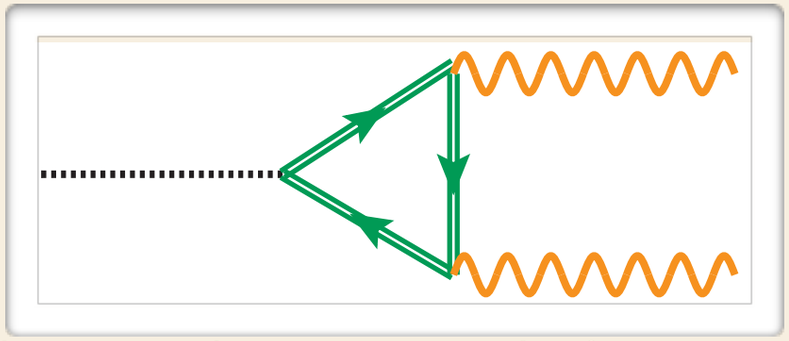
Higgs boson has been discovered

- m_H [GeV] = $125.09 \pm 0.21_{stat} \pm 0.11_{scale} \pm 0.02_{other} \pm 0.01_{theo}$ (CMS + ATLAS Run 1: $\gamma\gamma$ + 4 lepton)
- Γ_H [MeV] = $1.7^{+7.7}_{-1.8}$ (CMS), < 23 (95%, ATLAS)
- $\sigma/\sigma_{SM} = 1.00 \pm 0.13$ (ATLAS)
- All measured properties are consistent with SM expectations within experimental uncertainties
  - spin zero
  - parity +
  - couples to masses of W and Z (with c_V=1 within experimental uncertainty)
- Prime task at LHC: to measure its couplings (could signal BSM physics)
- y_t cannot be measured in $H \to tT$ decay ($m_t > m_H$)

How to measure y_t?

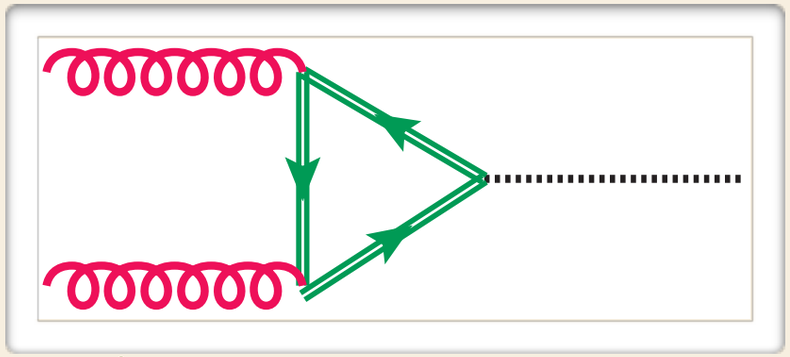
- $H \to \gamma\gamma$ is sensitive to y_t through t-quark loop,

  but rates are small and W loop also contributes



- $gg \to H$ is sensitive to y_t through t-quark loop

  if only SM model particles contribute (so far xsec is consistent with SM)

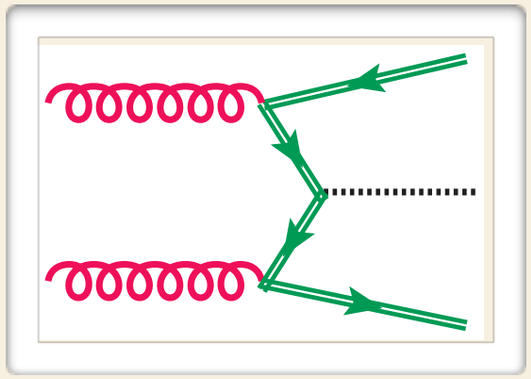


- $gg \to H$ is sensitive to BSM physics

  if y_t is measured separately

tTH hadroproduction

- y_t can be measured in pp $\rightarrow$ tTH through many decay channels (all very difficult):



  - hadrons with single lepton: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \mathrm{b}\bar{\mathrm{b}}$
  - hadrons with dilepton: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}\ell\nu,\ \mathrm{H} \rightarrow \mathrm{b}\bar{\mathrm{b}}$
  - hadrons with hadronic tau: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \tau_h^+\tau_h^-$
  - diphoton with lepton: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \gamma\gamma$
  - diphoton with hadrons: $\mathrm{t} \rightarrow \mathrm{b}\ jj,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \gamma\gamma$
  - same sign dilepton: $\mathrm{t} \rightarrow \mathrm{b}\ jj,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \ell\nu\ell[\nu]$
  - 3 leptons with di, trilepton: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}jj,\ \mathrm{H} \rightarrow \ell[\nu]\ell[\nu]$
  - 4 lepton with di, trilepton: $\mathrm{t} \rightarrow \mathrm{b}\ell\nu,\ \bar{\mathrm{t}} \rightarrow \bar{\mathrm{b}}\ell[\nu],\ \mathrm{H} \rightarrow \ell[\nu]\ell[\nu]$

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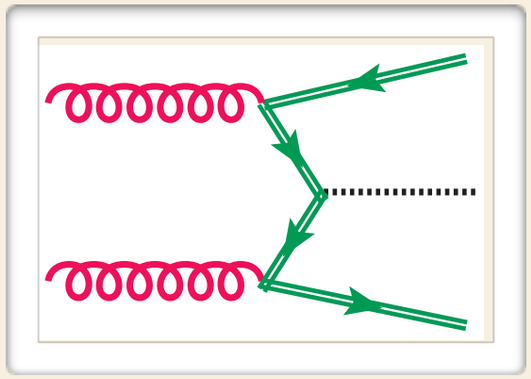

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the photons have to be isolated

Isolated photons in the final state

... are cumbersome theoretically

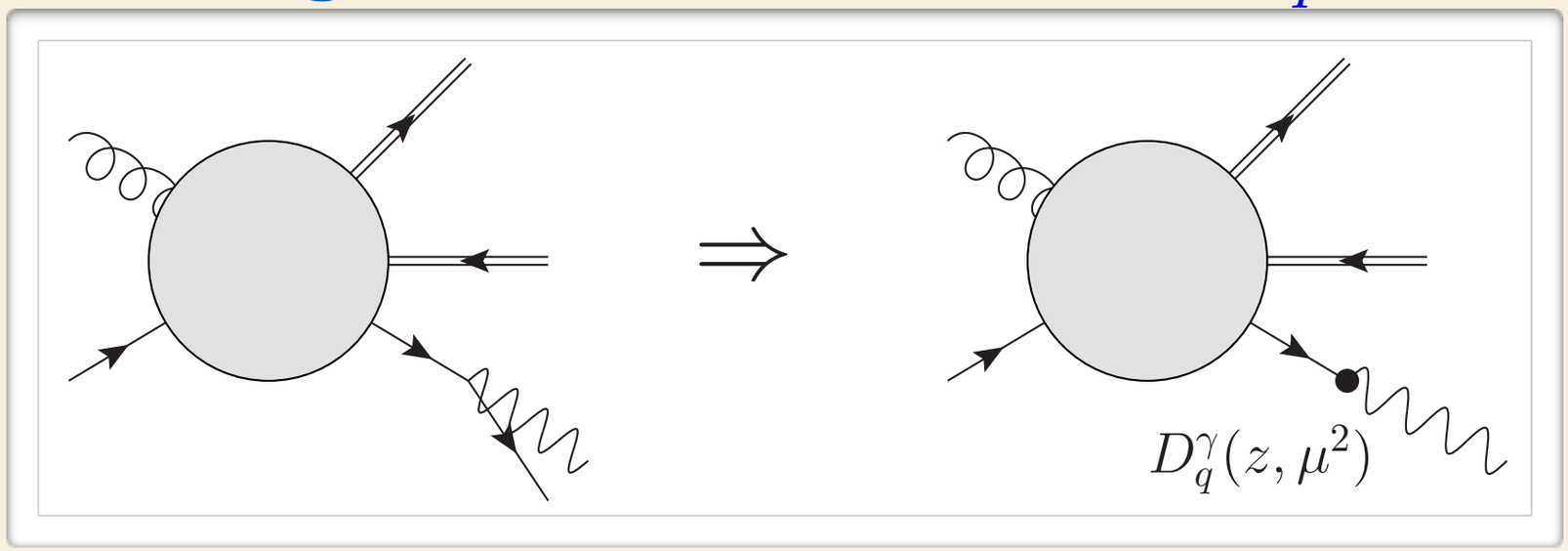
two contributions:

direct

fragmentation

$$\sigma = \sigma_d + \sigma_f$$

Origin of the fragmentation contribution is a light quark - photon singularity absorbed into the fragmentation function $D_q^\gamma(z, \mu^2)$



photon fragmentation is not known well

Fragmentation function

Fragmentation function

- cannot be computed in perturbation theory, only its evolution
- suppressed by isolation, hard to extract

Can we avoid it?

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- yes: isolate photon from partons with cone radius decreasing as parton energy inside the cone decreases (smooth isolation of Frixione):

$$\sum_{i \in \mathrm{trks}} E_{\perp,i} \Theta\left(\delta_0 - \Delta R(\gamma, i)\right) \leq \epsilon_\gamma E_{\perp,\gamma} \left(\frac{1 - \cos \Delta R(\gamma, i)}{1 - \cos \delta_0} \right)^{n_\gamma}$$

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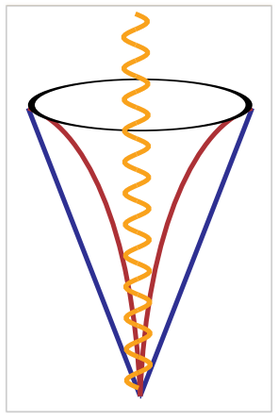

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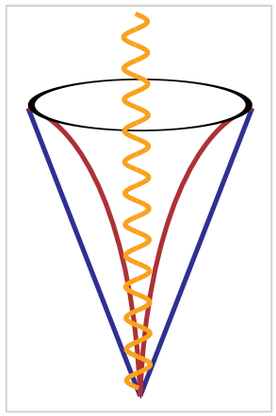

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- but experimenters prefer cone of fixed radius with reduced hadronic activity inside the cone:

$$E_{\perp,\text{had}} = \sum_{i\in\text{tracks}} E_{\perp,i}\Theta\left(R_\gamma - R(p_\gamma,p_i)\right) < E_{\perp,\text{had}}^{\max}$$

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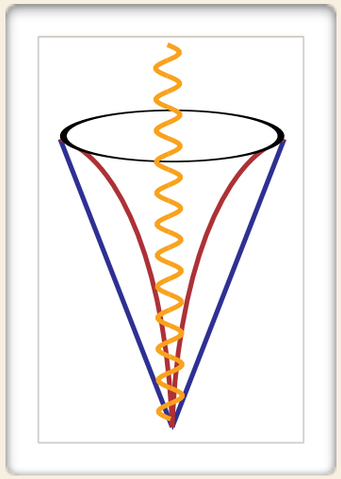

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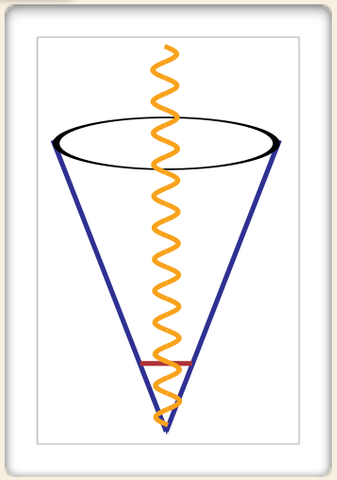

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Resolution:

Generate events with loose smooth photon generation isolation (i.e. small δ_0^{gen}) using the POWHEG method

[Kardos and Trócsányi, arXiv: 1406.2324]

PowHel framework

HELAC-NLO

[Bevilacqua et al, arXiv: 1110.1499]

POWHEG-BOX

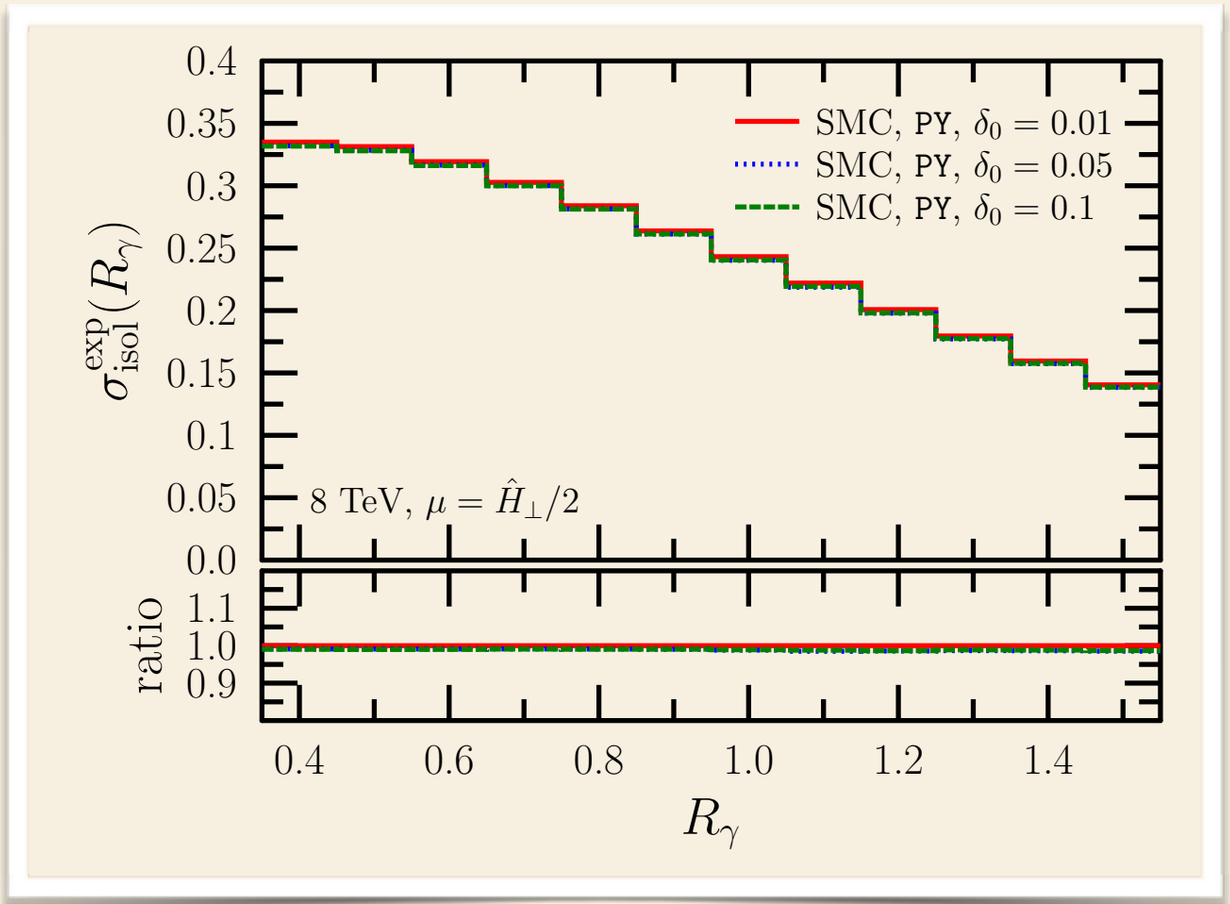
[Alioli, Nason, Oleari, Re, arXiv: 1002.2581]

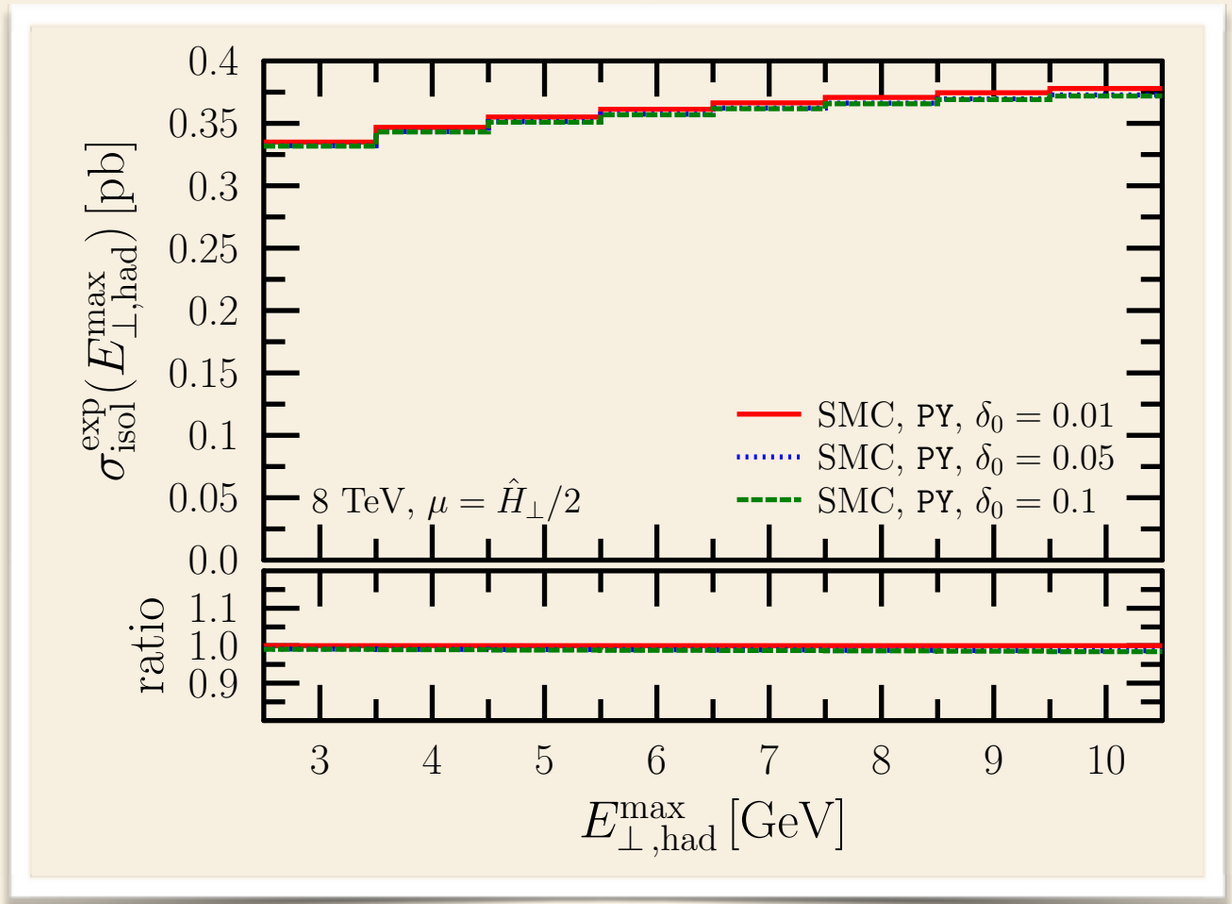
PowHel

RESULT of PowHel:

Les Houches file of Born and Born+1st radiation events (LHE) ready for processing with SMC followed by almost arbitrary experimental analysis

Physical predictions do not depend on the generation isolation…





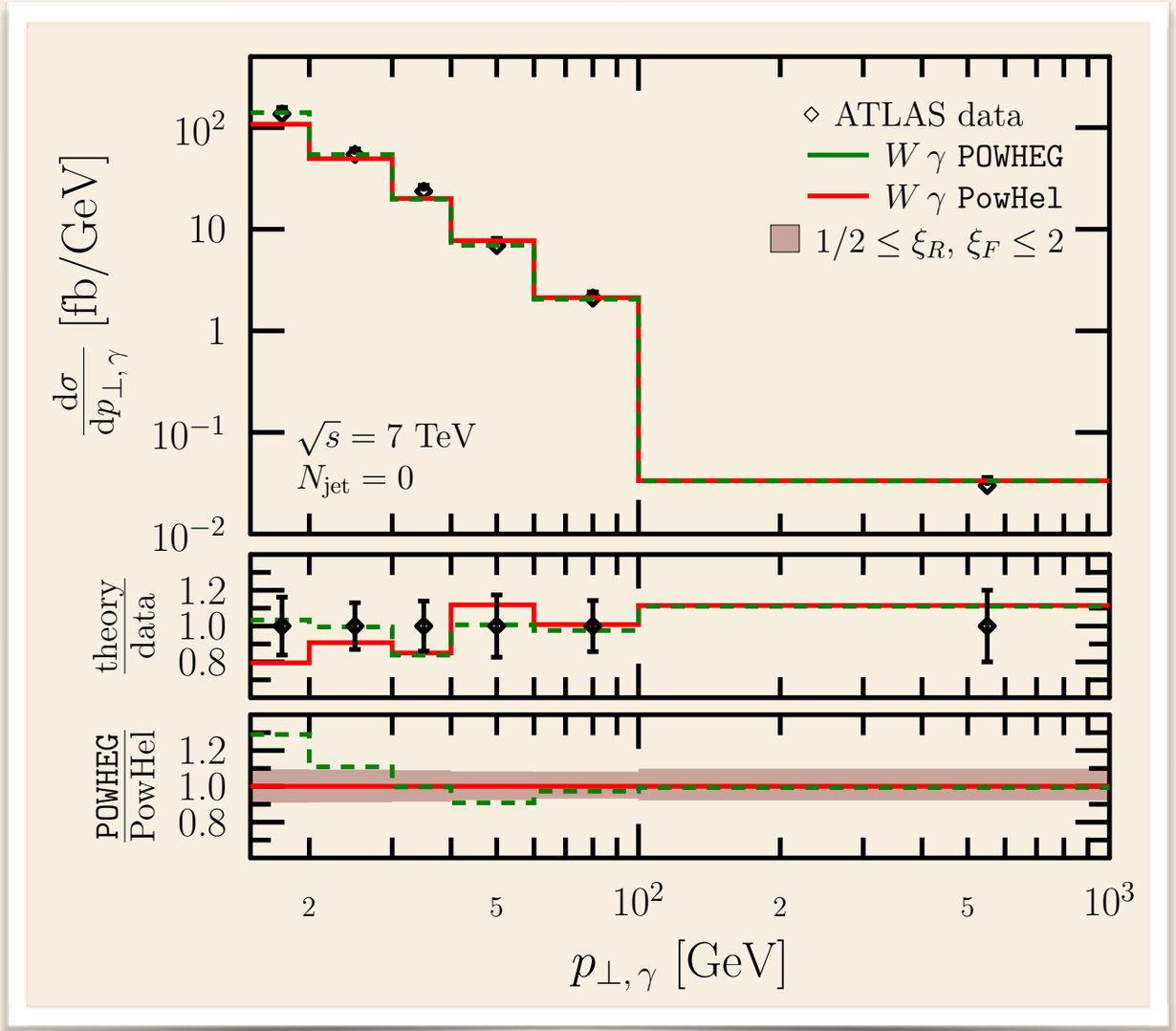
… if it is small enough

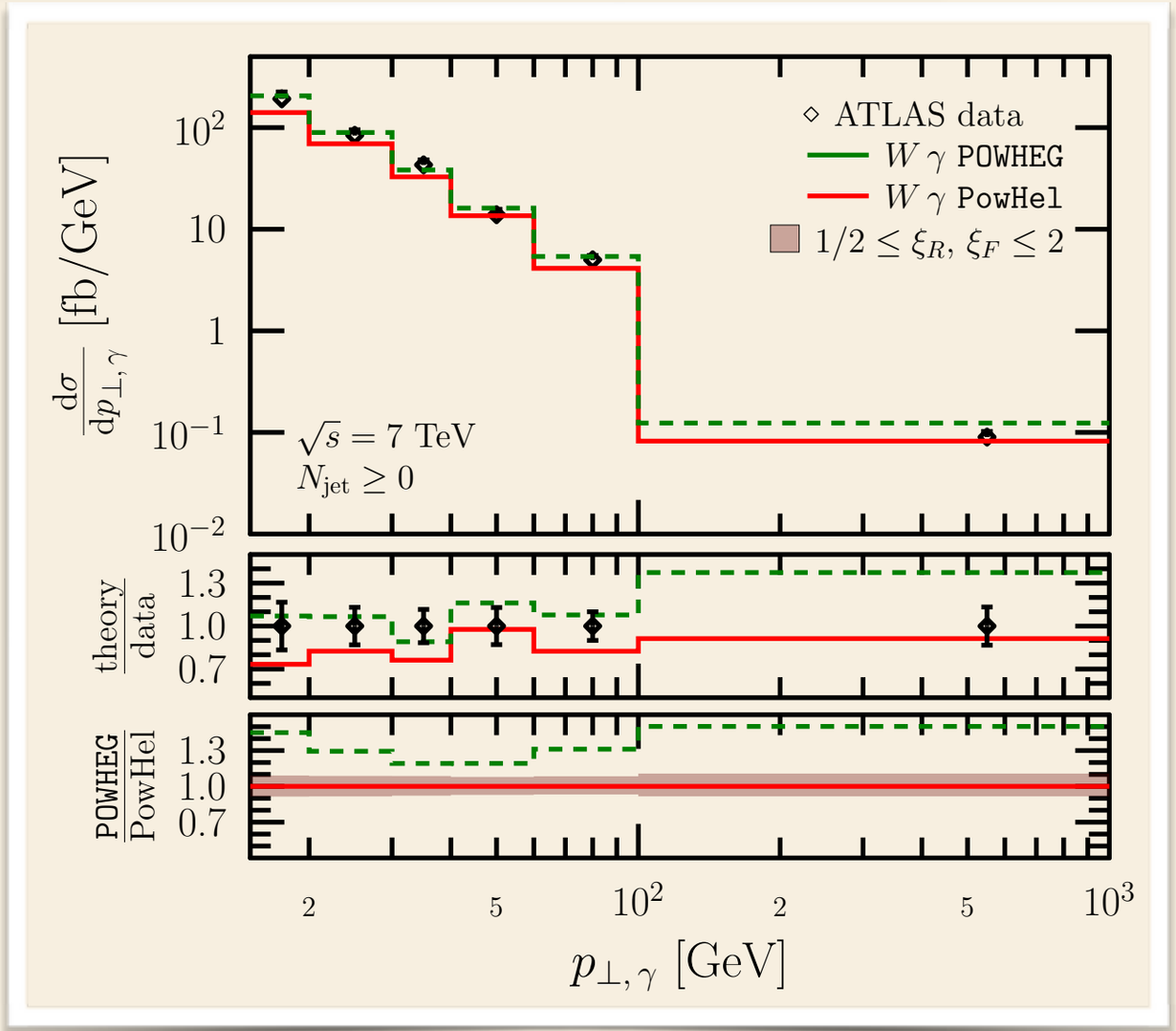
Expectation:

with sufficiently loose generation isolation
the fragmentation can be neglected
within the expected uncertainty
of matched NLO+PS predictions if the
photon is harder than the accompanying jets

[Kardos and Trócsányi, arXiv: 1406.2324]

Check method with Wγ production





good description of ATLAS data if

radiated photon is harder than accompanying jets

POWHEG (+MiNLO): [Barze et al: 1408.5766]

fragmentation is modeled by interleaved QCD+QED shower

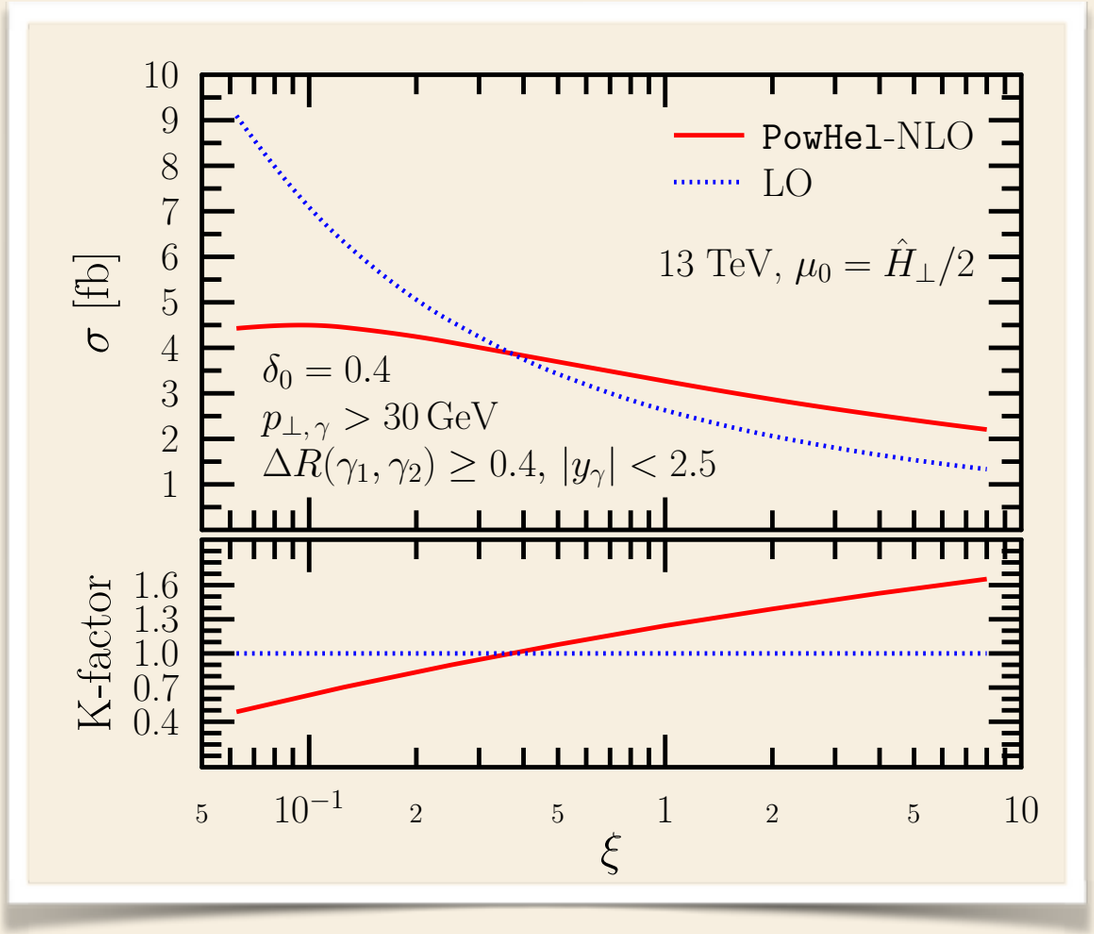
The method is general

$t\bar{t}\gamma\gamma$ hadroproduction at NLO
and matched to PS

[Kardos and Trócsányi, arXiv: 1408.0278]

Choice of scales

We use the dynamical scale $\mu_{dyn} = H_T/2$, where H_T is the scalar sum of transverse masses of final-state particles present in the underlying Born event



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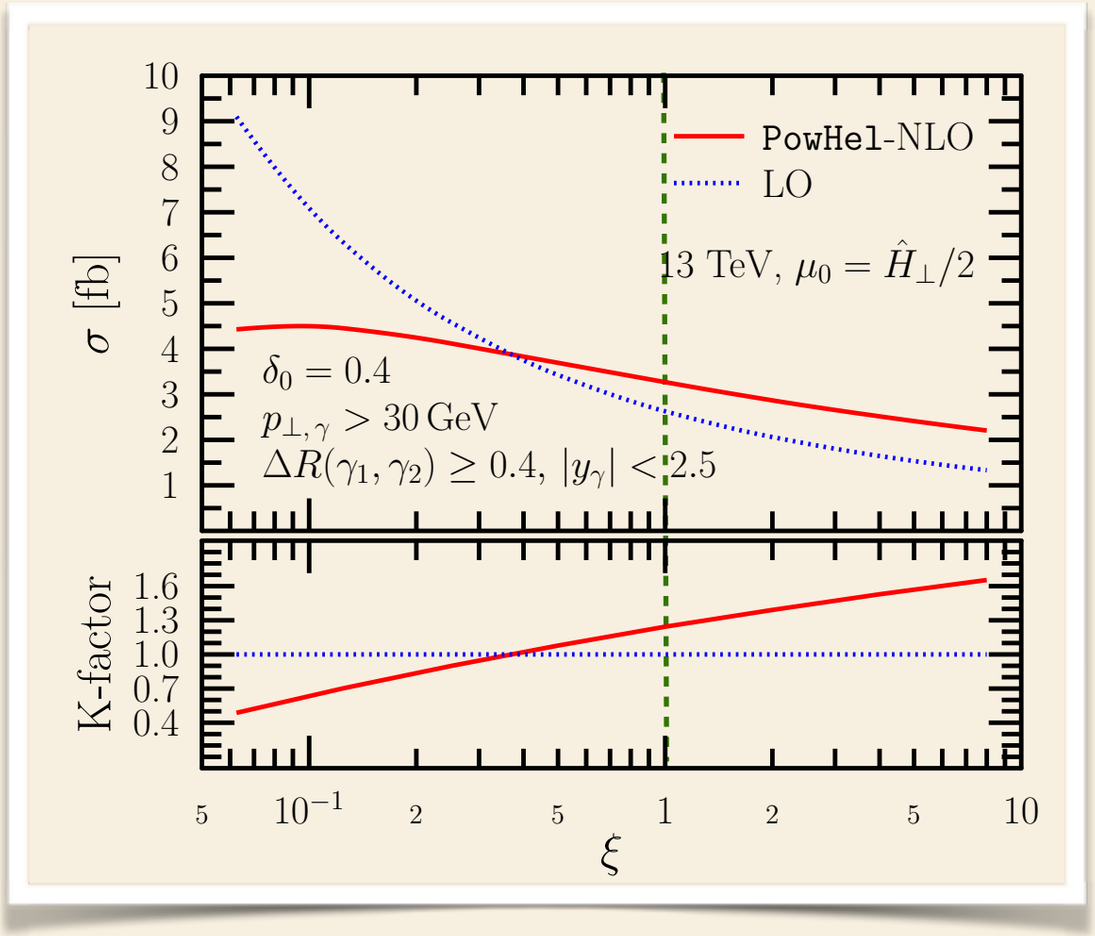
With this scale

✓ the K factor is moderate, implying good convergence

✓ σ_{LO} = 2.6 fb, σ_{NLO} = 3.3 fb
K = 1.24 (@ 13 TeV with cuts)



scale dependence:

LO: $^{+30\%}_{-27\%}$, NLO: $^{+14\%}_{-13\%}$, largest if $\mu_R = \mu_F = \mu_{dyn}$

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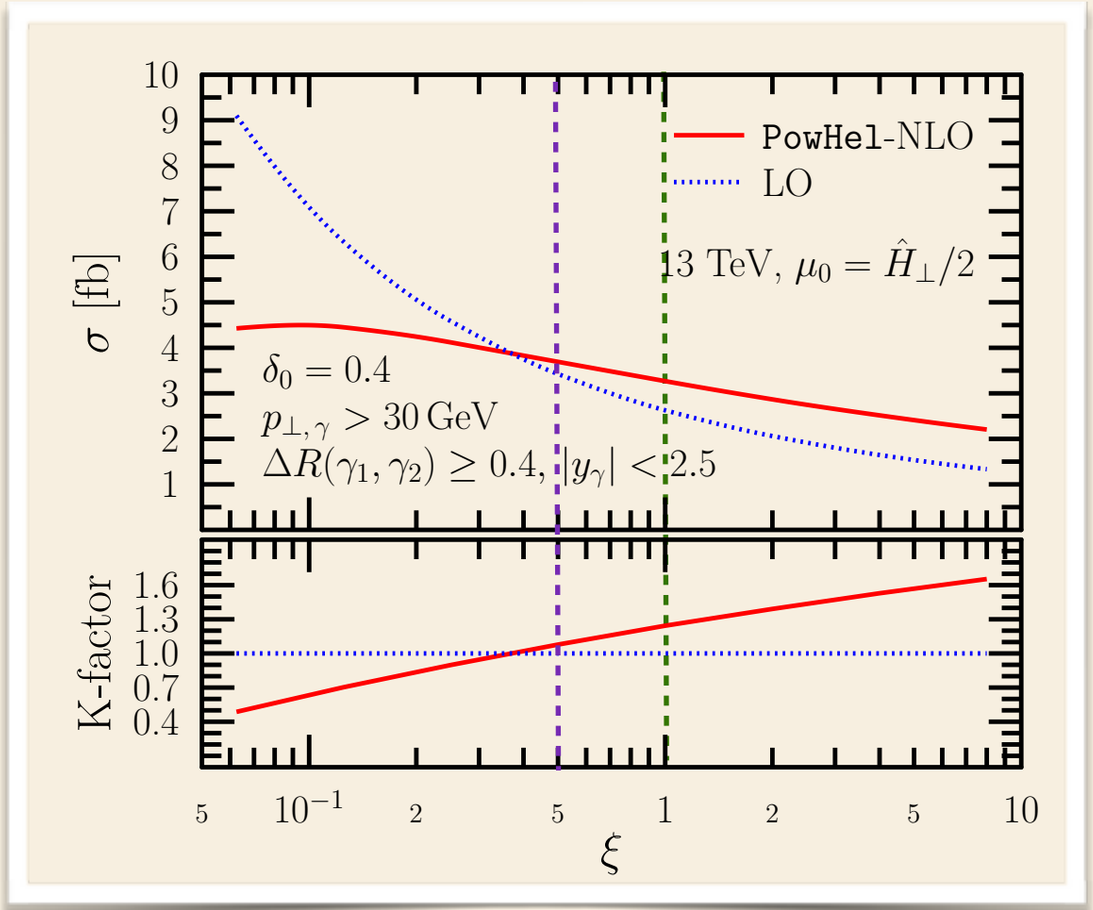

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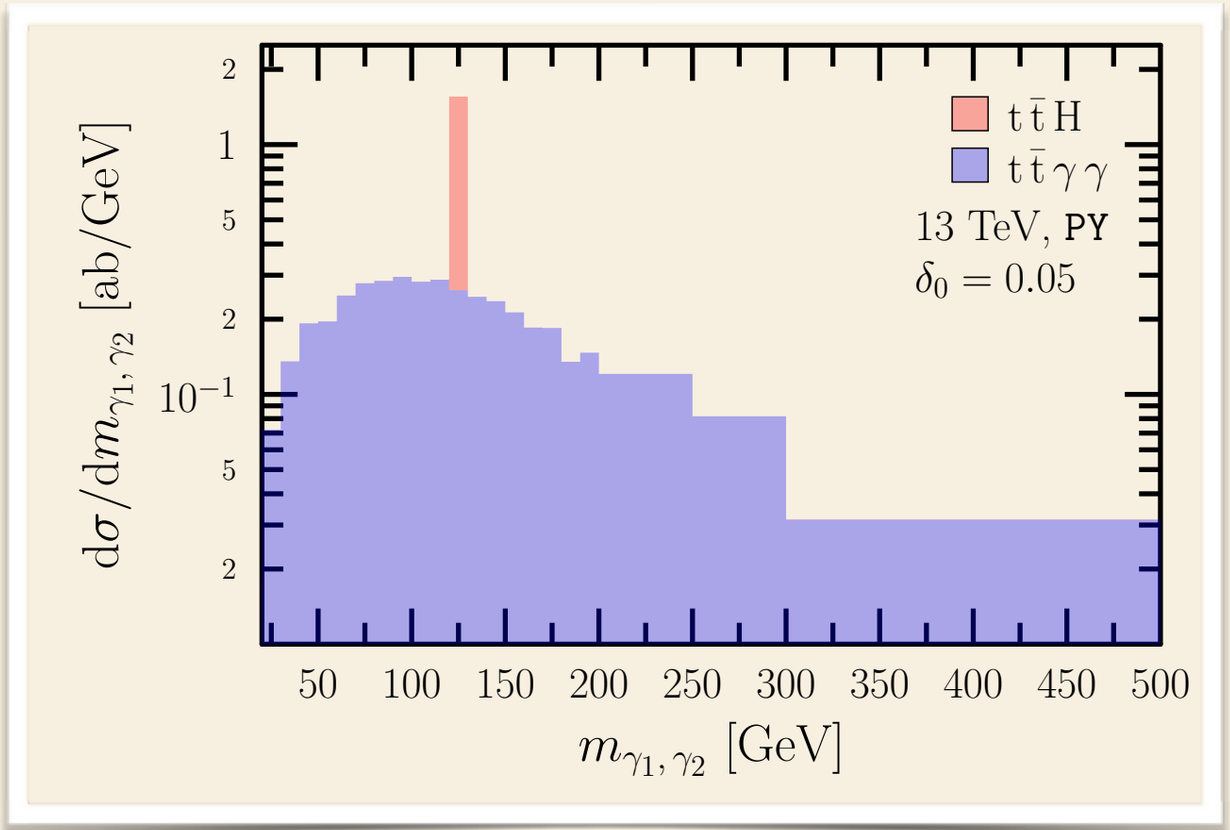
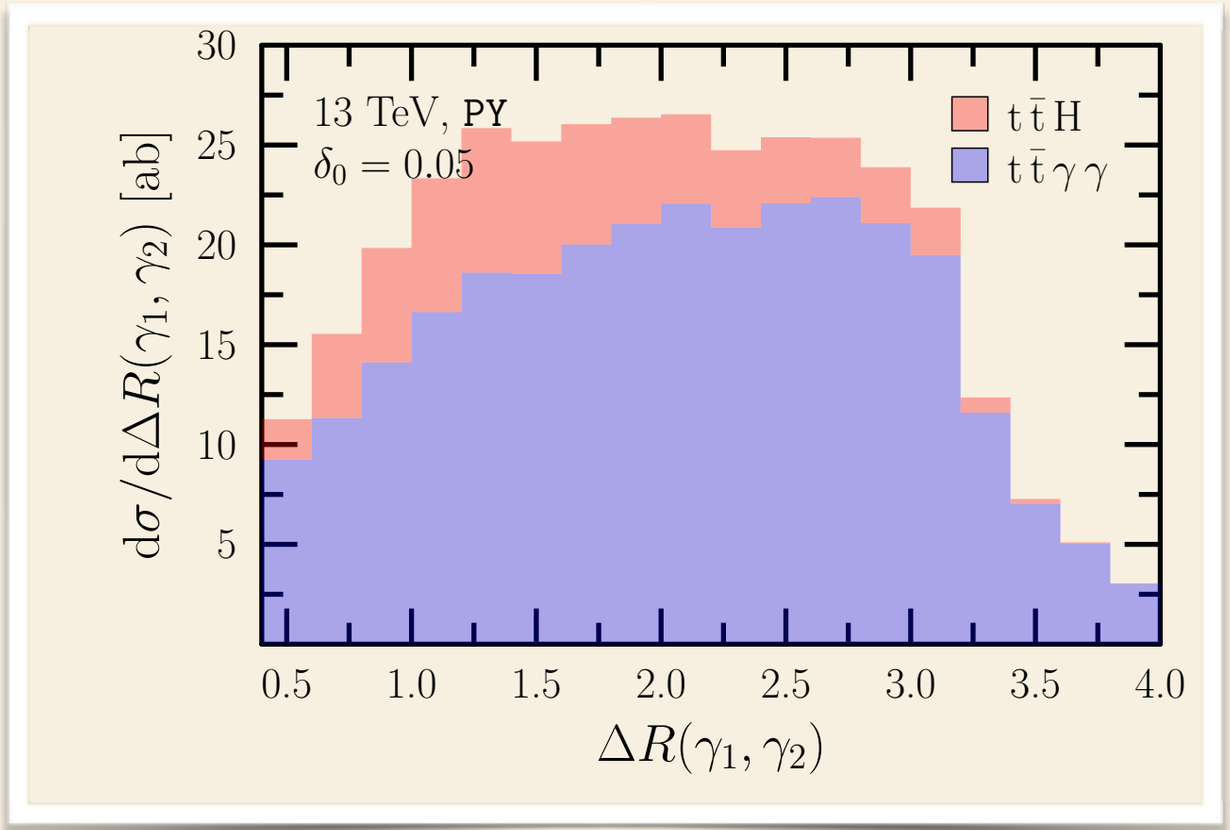
Combine with:

tTH hadroproduction at NLO and matched to PS

[Garzelli, Kardos, Papadopoulos and Trócsányi, arXiv: 1108.0387]

pp →tT (H→γγ) signal over pp →tT γγ background





obtained with HEPTopTagger, CT10NLO, PYTHIA-6.4.25

signal: $\mu_R = \mu_F = m_t + m_H/2 = (172.5 + 62.5)$ GeV

background: $\mu_R = \mu_F = H_T/2$

Conclusions

- POWHEG method offers a simple way for generating almost inclusive sample of pre-showered events with (loosely isolated) photons
- Predictions for hadroproduction of Wγ final states agree with ATLAS data if the photon is harder than the accompanying jet(s)
- Predictions presented for hadroproduction of tTγ and tTγγ final states at NLO accuracy (and matched with PS)
- Effects of PS are small
- Events are available on request or at http://grid.kfki.hu/twiki/bin/view/DbTheory/

Processes available in PowHel

✓ tT	[Kardos et al, arXiv:
✓ tT + Z	1111.0610,1111.1444,
✓ tT + W	1208.2665,
✓ tT + H/A	1108.0387,
✓ tT + j	1101.2672,
✓ WWbB	1405.5659,
✓ tT + bB	1303.6291, 1408.0266
✓ tT, W + γ	1406.2324
✓ tT + γ γ	1408.0278]

The end

Extra slides

HEPTopTagger

- Jet reconstruction using all the hadronic tracks with the C/A algorithm with R = 1.5 using FastJet
- Only those jets were considered for which $p_\perp > 200$ GeV and $|y| < 5$
- The t-quark candidate subjets were looked for in the jet-mass range of 150–200GeV
- We selected that particular subjet as a t-quark jet for which the jet mass was the closest to m_t

Selection cuts for tTH signal

- Two hard photons in central region with $p_{\perp,\gamma} > 30$ GeV and $|y_\gamma| < 2.5$.
- The photons are isolated from each other: $\Delta R(\gamma_1, \gamma_2) > 0.4$.
- Jet clustering on all the hadronic tracks with the anti-$k_\perp$ algorithm as implemented in FastJet with $R = 0.4$ and $p_{\perp,j} > 30$ GeV. Photons are isolated from these jets by $\Delta R(\gamma, j) > 0.4$ measured in the rapidity-azimuthal angle plane.
- Both of the hard photons are isolated from the top jet obtained by top tagging and from the three subjets of the top jet by $\Delta R(\gamma,j) > 0.4$ measured on the rapidity-azimuthal angle plane.
- One hard lepton in the central region with $p_\perp > 30$ GeV and $|y_l| < 2.5$, without distinction between leptons and antileptons.
- The lepton is isolated from both the jets and the photons with $\Delta R(l, i) > 0.4$, $i \in \{\gamma, j\}$ measured on the rapidity-azimuthal angle plane.
- Around photons the minimal hadronic activity in a cone of $R = 0.4$ is $E_{max} = 3$ GeV.