

LHCb results on CP violation in B decays

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CP violation

• The Standard Model predicted CP asymmetry is not sufficient to explain the baryon asymmetry of the Universe $V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$ \Rightarrow New Physics *CP* effects are expected 0.7 $\Delta m_d \& \Delta m_s$ ε_K 0.6 Precise measurements of heavy hadron 0.5 sin 2β decays \Rightarrow redundant determination sol. w/ cos $2\beta < 0$ excl. at CL > 0.95 0.4 of the CKM parameters α 0.3 0.2 Vur 0.1 In this talk Kief Run-I results 0.0 0.2 -0.4 -0.2 0.0 0.4 0.6 0.8 1.0 $\overline{\rho}$ ✓ Measurement of the B^0 mixing frequency ✓ Measurement of \mathcal{CP} in $B^0 \to J/\psi K_S^0$ decays ✓ Measurement of CP and polarization fractions in $B_s^0 \to J/\psi \overline{K}^{*0}$ decays Study of $B^- \to D^0 K^- \pi^+ \pi^-$ and $B^- \to D^0 \pi^- \pi^+ \pi^-$ decays and determination of the CKM angle γ Determination of the CKM parameter $|V_{ub}|$

Precision measurement of Δm_d

• In neutral meson systems, the flavour oscillation frequency is $\Delta m = m_H - m_L$





- Mixing asymmetry measured in: $B^0 \to D^{(*)-} \mu^+ \nu_\mu X$ decays

$$A(t) = \frac{N^{not \ osc}(t) - N^{osc}(t)}{N^{not \ osc}(t) + N^{osc}(t)} = \cos(\Delta m_d t)$$

Precision measurement of Δm_d

• In neutral meson systems, the flavour oscillation frequency is $\Delta m = m_H - m_L$





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Precision measurement of Δm_d

- Fit to the decay time distribution for unmixed and mixed events
- Mixing asymmetry projections in four flavour tagging categories



The Evergreens: $\sin(2\beta)$ using $B^0 \to J/\psi K_S^0$ decays

Mixing induced CP: in the interference of direct decay and decay after mixing



production asymmetries

PRL 115, 031601 (2015)



 Consistent with world average and similar precision to B factories

• 95690 $B_s^0 \to J/\psi [\to \mu^+ \mu^-] \phi [\to K^+ K^-]$

• time dependent angular analysis to disentangle CP even and CP odd



CP violation and polarization fractions in $B_s^0 \to J/\psi \overline{K}^{*0}$ decays



$$\phi_q^{\rm measured} = \phi_q^{SM} + \Delta \phi_{\rm Penguin} + \Delta \phi_{\rm NewPhysics} \begin{bmatrix} {\rm Nie} & {\rm Lie} & {\rm Lie} \\ {\rm Lie} & {\rm Lie} & {\rm Lie} \\ {\rm Lie} & {\rm Lie} & {\rm Lie} & {\rm Lie} & {\rm Lie} \\ {\rm Lie} & {\rm Li$$

Nierste et al. <u>arXiv:1503.00859</u>, Liu et al. <u>PRD 89, 094010 (2014)</u>



- $\Delta \phi_{\mathrm{Penguin}}$ and/or CP could be different for each polarization state
- Measurement of $\Delta\phi_{\rm Penguin}$ with decays where the penguin/tree ratio is not suppressed

$$A(B_s^0 \to (J/\psi \overline{K}^{*0})_i) = -\lambda \mathcal{A}_i (1 - a_i e^{i\theta_i} e^{i\gamma})$$
$$A(B_s^0 \to (J/\psi \phi)_i) = \left(1 - \frac{\lambda^2}{2}\right) \mathcal{A}'_i (1 - \frac{a'_i}{2} e^{i\theta'_i} e^{i\gamma})$$

 $i \in (0, \bot, \|, S)$

SU(3) flavor : $a_i = a'_i, \theta_i = \theta'_i$



CP violation and polarization fractions in $B_s^0 \rightarrow J/\psi \overline{K}^{*0}$ decays

• measurement of direct CP ($J/\psi K^+\pi^- \operatorname{vs} J/\psi K^-\pi^+$) decays are flavour specific

Candidates / 0.1

- time integrated
- polarization dependent
- Account for production and detection asymmetries







• Results $\begin{aligned} f_0 &= 0.497 \pm 0.025(\text{stat}) \pm 0.025(\text{syst}) \\ f_{\parallel} &= 0.179 \pm 0.027(\text{stat}) \pm 0.013(\text{syst}) \end{aligned}$ $A_0^{CP}(B_s^0 \to J/\psi \overline{K}^{*0}) &= -0.048 \pm 0.057(\text{stat}) \pm 0.020(\text{syst}) \\ A_{\parallel}^{CP}(B_s^0 \to J/\psi \overline{K}^{*0}) &= 0.171 \pm 0.152(\text{stat}) \pm 0.028(\text{syst}) \end{aligned}$ $A_{\perp}^{CP}(B_s^0 \to J/\psi \overline{K}^{*0}) &= -0.049 \pm 0.096(\text{stat}) \pm 0.025(\text{syst}) \end{aligned}$

$$\mathcal{B}(B_s^0 \to J/\psi \overline{K}^{*0}) = \left(4.13 \pm 0.16(\text{stat}) \pm 0.25(\text{syst}) \pm 0.24(\text{f}_d/\text{f}_s)\right) \times 10^{-5}$$

• Combined with $B^0 \to J/\psi \rho^0$: $\Delta \phi_{s,0}^{J/\psi\phi} = 0.000^{+0.009}_{-0.011} (\text{stat})^{+0.004}_{-0.009} (\text{syst}) \text{rad}$ $\Delta \phi_{s,\parallel}^{J/\psi\phi} = 0.001^{+0.010}_{-0.014} (\text{stat})^{+0.007}_{-0.008} (\text{syst}) \text{rad}$ $\Delta \phi_{s,\perp}^{J/\psi\phi} = 0.003^{+0.010}_{-0.014} (\text{stat})^{+0.007}_{-0.008} (\text{syst}) \text{rad}$

absolute shift smaller than 19 mrad current experimental precision:

 $\sigma(\phi_s) = \pm 0.035 \text{ rad}$

NO NEED TO WORRY YET

LHCB-PAPER-2015-034

• CKM angle $\gamma \equiv [-(V_{ud}V_{ub}^*)/(V_{cd}V_{cb}^*)]$

• Study the interference between $B^- \to D^0 X^-_{s,d}$ and $B^- \to \overline{D^0} X^-_{s,d}$ decays, selecting final states accessible to both D^0 and $\overline{D^0}$ (all tree level)



 r_B and δ_B are the amplitude ratio and strong phase difference between $B \to D^0 X$ and $B \to \overline{D^0} X$ $r_D e^{\delta_D}$ is the ratio between the Cabibbo Favoured and the Doubly Cabibbo Suppressed amplitudes Coherence factors κ appear with multibody states

CP observable:

$$R^{X^{\pm}} = \frac{\Gamma(B^{\pm} \to [K^{\mp}\pi^{\pm}]_D X^{\pm})}{\Gamma(B^{\pm} \to [K^{\pm}\pi^{\mp}]_D X^{\pm})} = \frac{r_B^2 + r_D^2 + 2\kappa r_B r_D \cos(\delta_B + \delta_D \pm \gamma)}{1 + r_B^2 r_D^2 + 2\kappa r_B r_D \cos(\delta_B - \delta_D \pm \gamma)}$$

$$LHCB-PAPER-2015-020$$

- CKM angle $\gamma \equiv \left[-(V_{ud}V_{ub}^*)/(V_{cd}V_{cb}^*)\right]$
- Study the interference between $B^- \to D^0 X^-_{s,d}$ and $B^- \to \overline{D^0} X^-_{s,d}$ decays, selecting final states accessible to both D^0 and $\overline{D^0}$ (all tree level)
- Different "methods":
- **ADS** selection of <u>quasi-flavor specific final states</u>, Cabibbo favoured (CF) and doubly Cabibbo suppressed (DCS) $D \to K^{\pm} \pi^{\mp}$

CP observable of interest: relative widths of DCS to CF, separated by charge

GLW selection of <u>CP eigenstates</u> $D \to K^+ K^-$, $D \to \pi^+ \pi^-$

CP observables of interest: charge-averaged yields ratios, charge asymmetries

- Counting analyses
- Ratios or double ratios reduce the systematic uncertainties

• Signal yields are obtained with a simultaneous unbinned extended maximum likelihood fit to the B candidate invariant mass spectra



• Fitted signal yields are corrected (for the events removed by the vetoes, charmless background, B production asymmetry, kaon and pion detection asymmetries)

Constraints on γ



• The precision is comparable or better than most previous measurements

LHCB-PAPER-2015-020

LHCb combination & previous measurements

- Measurements included:
- $B^+ \rightarrow Dh^+, D \rightarrow hh \text{ GLW/ADS}, 1 \text{ fb}^{-1}$ Phys. Lett. B712 (2012)203
- $B^+ \to Dh^+, D \to K\pi\pi\pi, \text{ADS}, 1 \text{ fb}^{-1}$ Phys. Lett. B723 (2013) 44

•
$$B^+ \to DK^+, D \to K^0_S hh$$
, model independent, GGSZ, 3 fb⁻¹

•
$$B^+ \to DK^+, D \to K^0_S K\pi$$
, GLS, 3 fb⁻¹ Phys. Lett. B733 (2014) 36

•
$$B^0 \rightarrow DK^{*0}$$
, $D \rightarrow hh$, GLW/ADS, 3 fb⁻¹_{Phys. Rev. D90 (2014) 112002}

•
$$B_s^0 \to D_s^{\mp} K^{\pm}$$
, time-dependent, 1 fb⁻¹ JHEP 11 (2014) 060

- Combination accounts for $D^0 \overline{D^0}$ mixing effect and supplementary informations from other experiments
- Taking the best fit value and 68% CL interval

$$\gamma = (73^{+9}_{-10})^{\circ}$$

LHCB-CONF-2014-004

Determination of $|V_{ub}|$

- |Vub| is complementary to γ and β in constraining the Unitarity Triangle
- Tension between the exclusive and inclusive measurements of |Vub|

• LHCb strategy: measure the ratio of branching fractions of the Λ_b^0 baryon into $p\mu^-\overline{\nu_\mu}$ and $\Lambda_c^+\mu^-\overline{\nu_\mu}$

$$\frac{|V_{ub}|^2}{|V_{cb}|^2} = \frac{\mathcal{B}(\Lambda_b^0 \to p\mu^- \overline{\nu_\mu})}{\mathcal{B}(\Lambda_b^0 \to \Lambda_c^+ \mu^- \overline{\nu_\mu})} R_{FF}$$

 $R_{FF} = 1.470 \pm 0.115 (stat) \pm 0.104 (syst)$ W. Detmold, C. Lehner and S. Meinel arXiv:1503.01421

Belle measurement arXiv:1312.7826

$$\frac{\mathcal{B}(\Lambda_b^0 \to p\mu^- \overline{\nu_{\mu}})_{q^2 > 15 GeV^2/c^4}}{\mathcal{B}(\Lambda_b^0 \to \Lambda_c^+ \mu^- \overline{\nu_{\mu}})_{q^2 > 7 GeV^2/c^4}} = \frac{N(\Lambda_b^0 \to p\mu^- \overline{\nu_{\mu}})}{N(\Lambda_b^0 \to \Lambda_c^+ (\to pK^- \pi^+)\mu^- \overline{\nu_{\mu}})} \times \frac{\epsilon(\Lambda_b^0 \to \Lambda_c^+ (\to pK^- \pi^+)\mu^- \overline{\nu_{\mu}})}{\epsilon(\Lambda_b^0 \to p\mu^- \overline{\nu_{\mu}})} \times \mathcal{B}(\Lambda_c^+ \to pK^- \pi^+)$$
world average
$$|V_{cb}| = (39.5 \pm 0.8) \times 10^{-3}$$

- Boosted decision tree removes backgrounds with additional charged tracks that could vertex with a pµ candidate.
- Efficiency from simulation, with many data-driven corrections

Determination of $|V_{ub}|$

• Corrected mass, $m_{corr} = \sqrt{m_{h\mu}^2 + p_{\perp}^2 + p_{\perp}}$ fits are used to extract signal and control sample yields, accounting for the per-event uncertainty



- can check the consistency of $|V_{ub}|/|V_{cb}|$ with eta
- does not support a right-handed coupling of significant magnitude

Nature Physics 10 (2015)

Summary and Conclusions

- Precise measurements using b-hadron decays multiply the complementary constraints to the CKM picture
- CP-conserving quantities:
 - world's best measurement of Δm_d
 - first measurement at Hadron Colliders of $|V_{ub}|/|V_{cb}|$

CP-violating quantities:

- γ determination with B→DKππ and B→Dπππ decays and γ combination
- Run-I measurements of $sin(2\beta)$ and ϕs
- Contributions from "Penguin pollution" are shown to be small!
- Stay tuned for Run-II!





Backup slides

Δm_d : B decay time and Resolutions

- The momentum of the B meson cannot be measured precisely due to the partial reconstruction of the decay.
- The B decay time is corrected using the factor:

$$k = p_{reco}/p_{true}$$

• The k-factors are also used to model the decay time resolution:



k factor

 $M_{PDG} \cdot k_{av}(M)$

Flavour tagging



- CKM angle $\gamma \equiv \left[-(V_{ud}V_{ub}^*)/(V_{cd}V_{cb}^*)\right]$
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- Different "methods":

GLW

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CP observable of interest:

$$R^{X^{\pm}} = \frac{\Gamma(B^{\pm} \to [K^{\mp}\pi^{\pm}]_D X^{\pm})}{\Gamma(B^{\pm} \to [K^{\pm}\pi^{\mp}]_D X^{\pm})} = \frac{r_B^2 + r_D^2 + 2\kappa r_B r_D \cos(\delta_B + \delta_D \pm \gamma)}{1 + r_B^2 r_D^2 + 2\kappa r_B r_D \cos(\delta_B - \delta_D \pm \gamma)}$$

selection of CP eigenstates $D \to K^+ K^-$, $D \to \pi^+ \pi^-$

CP observables of interest:

$$\begin{aligned} R_{CP+}^{h^+h^-} &= 2 \frac{\Gamma(B^- \to [h^+h^-]_D X_s^-) + \Gamma(B^+ \to [h^+h^-]_D X_s^+)}{\Gamma(B^- \to [K^-\pi^+]_D X_s^-) + \Gamma(B^+ \to [K^+\pi^-]_D X_s^+)} \qquad R_{CP+} = \frac{R_{s/d}^{n^+n^-}}{R_{s/d}^{K\pi}} \\ &= 1 + r_B^2 + 2kr_B \cos\delta_B \cos\gamma \\ \mathcal{A}_{X^\pm}^f &\equiv \frac{\Gamma(B^- \to f_D X^-) - \Gamma(B^+ \to \bar{f_D} X^+)}{\Gamma(B^- \to f_D X^-) + \Gamma(B^+ \to \bar{f_D} X^+)} = 2kr_B \sin\delta_B \sin\gamma/R_{CP+} \end{aligned}$$

Ratios or double ratios reduce the systematic uncertainties

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$|V_{ub}|$: significant right-handed coupling



 Hypothesis of a right-handed coupling (Bernlochner et al., <u>arXiv:1408.2516</u>, Crivellin, <u>arXiv.0907.2461</u>) is not supported by the combination after inclusion of the LHCb measurement