

# How to (path) integrate by differentiating

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# Overview

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- Applications to QFT
  - Expresses functional integrations and functional transforms in terms of functional differentiation.
  - Offers new perturbative approaches.



## New representations of integration:

$$\int_a^b f(x) dx = \lim_{\epsilon \rightarrow 0} f(\partial_\epsilon) \frac{e^{\epsilon b} - e^{\epsilon a}}{\epsilon}$$

$$\int_{-\infty}^{\infty} f(x) dx = \lim_{\epsilon \rightarrow 0^+} (f(\partial_\epsilon) + f(-\partial_\epsilon)) \frac{1}{\epsilon}$$

Compare with:

$$f'(x) = \lim_{\epsilon \rightarrow 0} (f(x + \epsilon) - f(x)) \frac{1}{\epsilon}$$

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## Examples for integration

Similarly, one quickly obtains, e.g.,

$$\int_{-\infty}^{\infty} \frac{\sin^5(x)}{x} dx = 3\pi/8$$

$$\int_{-\infty}^{\infty} \frac{\sin^2(x)}{x^2} dx = \pi$$

$$\int_{-\infty}^{\infty} \frac{(1 - \cos(tx))}{x^2} dx = \pi|t|$$

$$\int_{-\infty}^{\infty} x^2 \cos(x) e^{-x^2} dx = \sqrt{\pi} e^{-1/4} / 4$$

etc ...

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**Fourier transforming is even easier than integrating!**

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## How are they related?

The zero-frequency value of the Fourier transform is the integral (up to a prefactor of  $\sqrt{2\pi}$ ).

## Examples for Fourier

For example, for  $f(x) = \sin(x)/x$ , recall:

$$\begin{aligned}
 \int_{-\infty}^{\infty} \frac{\sin(x)}{x} dx &= 2\pi \lim_{x \rightarrow 0} \frac{1}{2i} \left( e^{\partial_x} - e^{-\partial_x} \right) \frac{1}{-i\partial_x} \delta(x) \\
 &= \pi \lim_{x \rightarrow 0} \left( e^{\partial_x} - e^{-\partial_x} \right) (\Theta(x) + c') \\
 &= \pi \lim_{x \rightarrow 0} (\Theta(x+1) - \Theta(x-1)) \\
 &= \pi
 \end{aligned}$$

By not taking the limit and by dividing by  $\sqrt{2\pi}$ , we obtain immediately:

$$\mathcal{F}[f](x) = \sqrt{\pi/2} (\Theta(x+1) - \Theta(x-1))$$

# Proof of the Fourier formula

**Why does this work?**

## Proof of the Fourier formula

The claim is:

$$\mathcal{F}[f](x) = \sqrt{2\pi} f(-i\partial_x) \delta(x)$$

Let us apply this to a plane wave:  $f(x) = e^{ixy}$ .

We obtain the right answer:

$$\begin{aligned}\mathcal{F}[f](x) &= \sqrt{2\pi} e^{y\partial_x} \delta(x) \\ &= \sqrt{2\pi} \delta(x + y)\end{aligned}$$

And the plane waves form a basis of the function space.

## What is going on, intuitively?

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Regulate, e.g., this way:  $\delta(x) = \lim_{\sigma \rightarrow 0} (2\pi\sigma)^{-1/2} e^{-x^2/2\sigma}$

$$\int_{-\infty}^{\infty} f(x) dx = 2\pi \lim_{\sigma \rightarrow 0} \frac{1}{\sqrt{2\pi\sigma}} e^{\partial_x^2/2\sigma} f(x)$$

Integration from asymptotics of heat flow !

## Perturbative expansions

In QFT, we'd like to apply the new methods, e.g.:

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# Perturbative expansions

But what if in

$$Z[J] = \int e^{iS[\phi] + i \int J\phi} d^n x D[\phi]$$

the action  $S[\phi]$  is not suitable to solve the integral or Fourier (or Laplace) transform with our new methods exactly?

And that's the norm of course!

## Perturbative expansions

- On the basic level, what if  $f(x)$  is too complicated, e.g., for:

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- Obtain weak & strong coupling expansions and others...
- Also: applications to deblurring expansion of signals.

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- What is the full size of the space of functions and distributions to which these methods apply?

- Relation to Stoke's theorem?

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- Relation to fermionic integration, a unifying formalism?
- A new perspective on integration measures and therefore anomalies in QFT?

## Bonus: Examples for Laplace

### Recall:

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## Examples for Laplace

If we apply the new Laplace transform method

$$\mathcal{L}[f](x) = f(-\partial_x) \frac{1}{x}$$

to monomials  $f(x) = x^n$  we obtain:

$$\mathcal{L}[f](x) = (-\partial_x)^n \frac{1}{x} = \frac{n!}{x^{n+1}}$$

And the monomials form a basis in the function space.

## Example for inverse Laplace

Consider a heat kernel trace:

$$h(t) = \sum_n e^{-\lambda_n t}$$

Given  $h(t)$ , the spectrum  $\{\lambda_n\}$  is known to be recoverable via inverse Laplace transform.

Why?

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this is easy to see:

$$\begin{aligned} \mathcal{L}^{-1}[h](\lambda) &= h(\partial_\lambda) \delta(\lambda) \\ &= \sum_n e^{-\lambda_n \partial_\lambda} \delta(\lambda) = \sum_n \delta(\lambda - \lambda_n) \end{aligned}$$