

Recent developments on unintegrated parton distributions

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- I. Results on small- x final states from k_{\perp} -factorized Monte Carlo event generators

- II. Progress toward precise characterizations of u-pdf's:
endpoint divergences $x \rightarrow 1$

collaboration with H. Jung

INTRODUCTION

▷ Parton distributions unintegrated in transverse momentum are naturally defined for $x \rightarrow 0$ via high-energy factorization

↪ ● basic QCD tool for small- x resummations

● implemented in Monte Carlo generators for HERA physics (+ LHC)

▷ But their relevance goes beyond small- x physics:

● Sudakov effects; infrared-sensitive processes

● polarized scattering; exclusive observables

● can be utilized for general-purpose Monte-Carlo's?

⇒ Q: How to define k_{\perp} distributions gauge-invariantly over the whole phase space?

OUTLINE

Part I: Unintegrated parton distributions and Monte Carlo generators

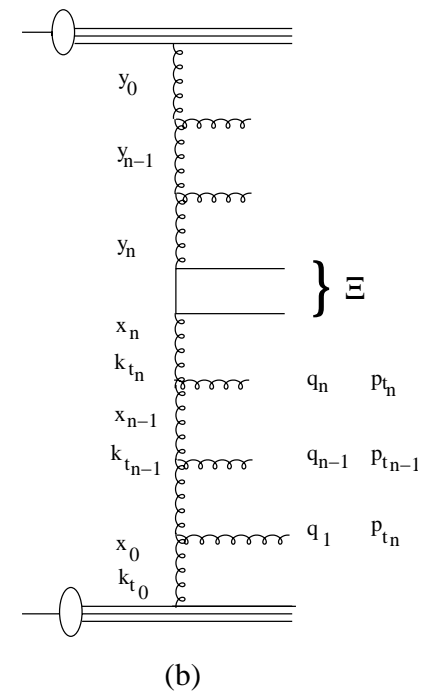
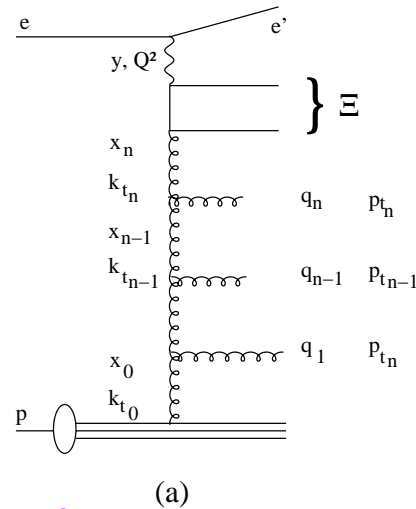
- hadronic final states in DIS at $x \ll 1$
- multi-jet distributions; angular correlations

Part II: Open issues on precise characterizations of updf's

- incomplete KLN cancellations near $x = 1$
- subtractive regularization method

UPDF's AND MONTE CARLO EVENT GENERATORS

- ◇ k_{\perp} -dependent matrix elements
- ◇ backward evolution for initial-state cascade



- parton emission in initial state only allowed in angular-ordered region of phase space

⇒ correct treatment of $x \ll 1$ region (logarithmic accuracy)

- need corrections for collinear and $x \sim 1$ region (included partially in present MC)

Existing implementations:

CASCADE www.quark.lu.se/~hannes/cascade

SMALLX Marchesini & Webber, 90's

LDCMC www.thep.lu.se/~leif/ariadne

Golec-Biernat et al., hep-ph/0703317

Höche et al., arXiv:0705.4577

See Proceedings Workshop “HERA and the LHC” for full references

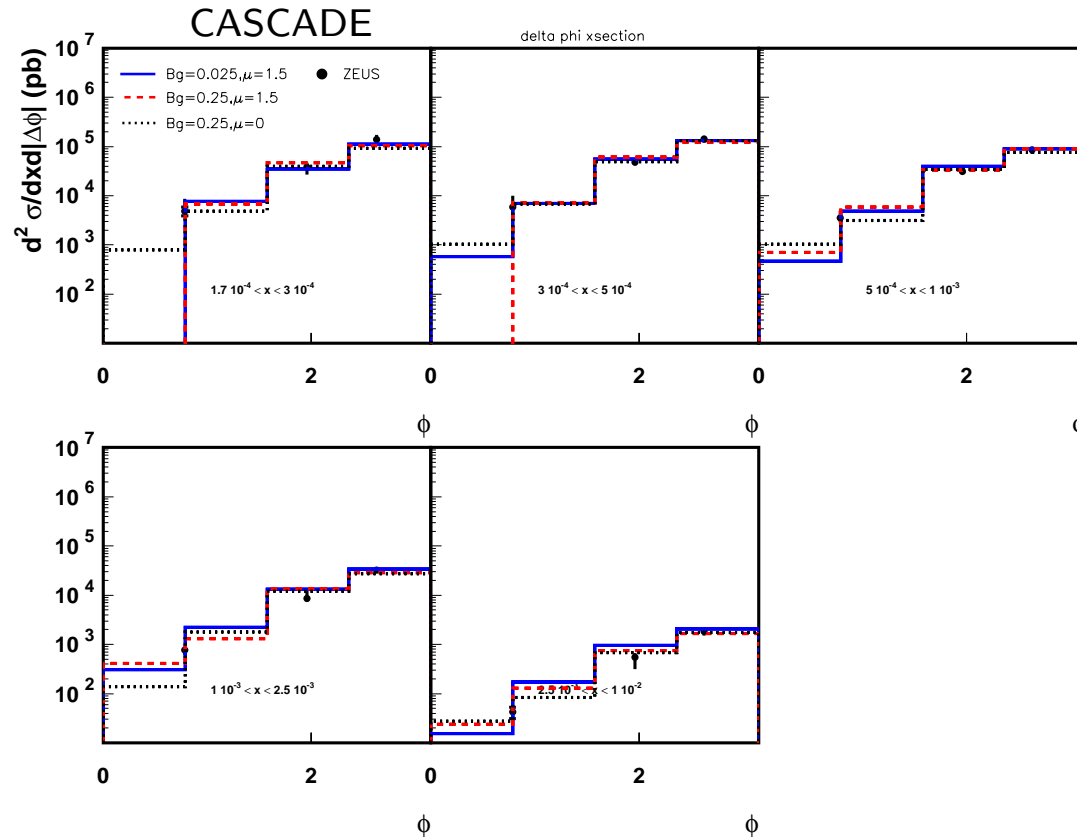
Example:

- multijet production in DIS at $x \ll 1$



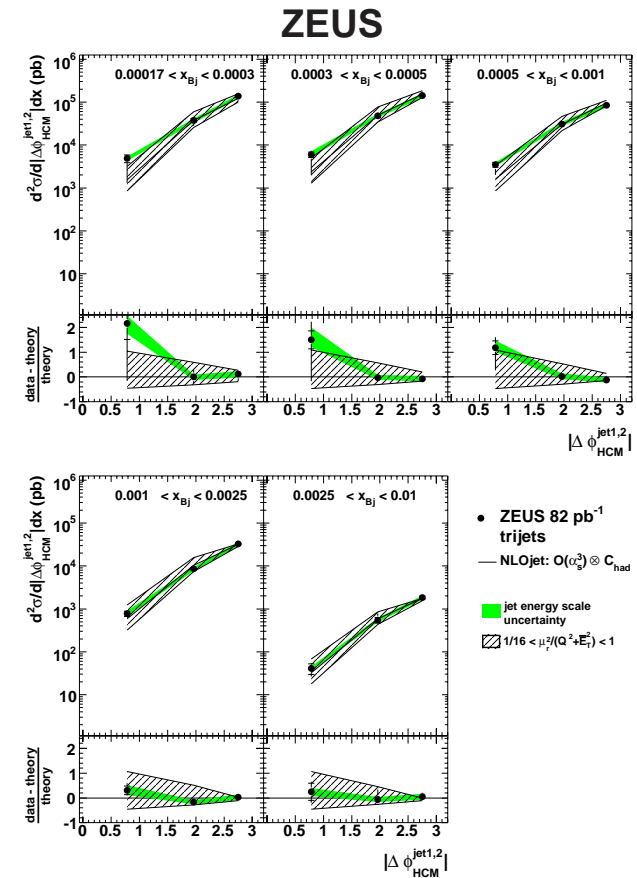
Azimuthal correlation in three-jet cross sections

ZEUS, arXiv:0705.1931



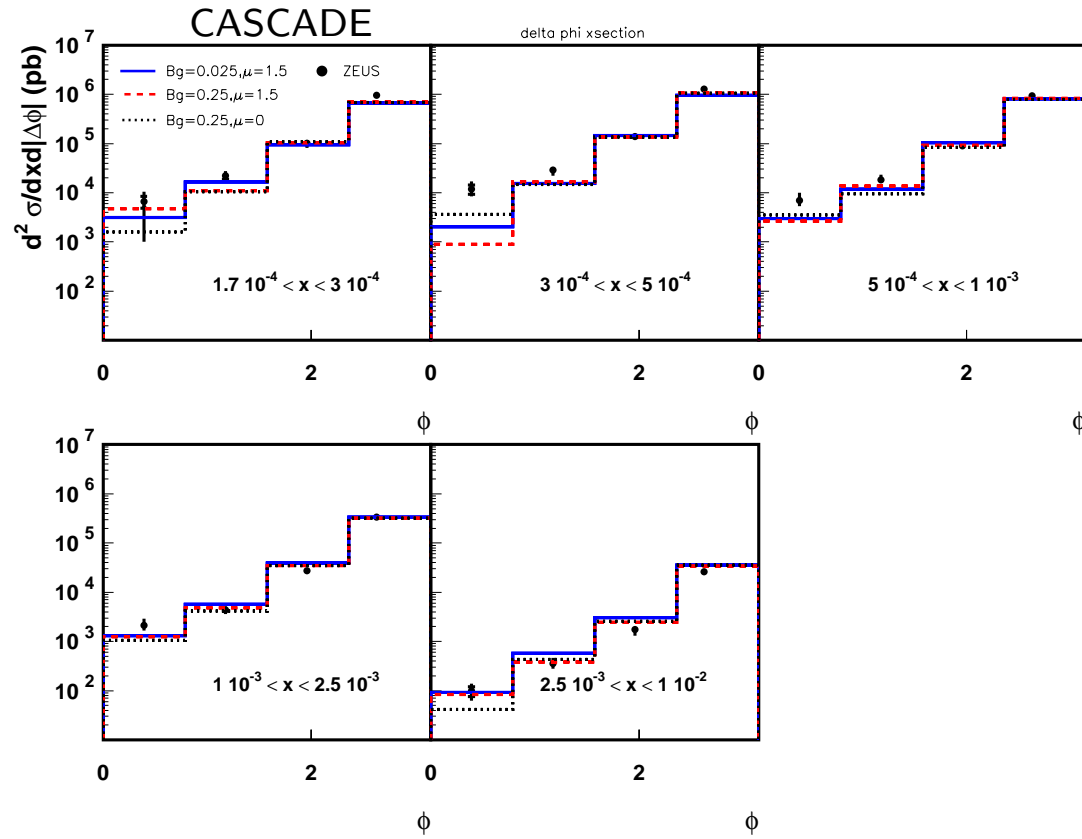
- unintegrated gluon fitted from inclusive DIS

- Note small- x shower (away from back-to-back ϕ)

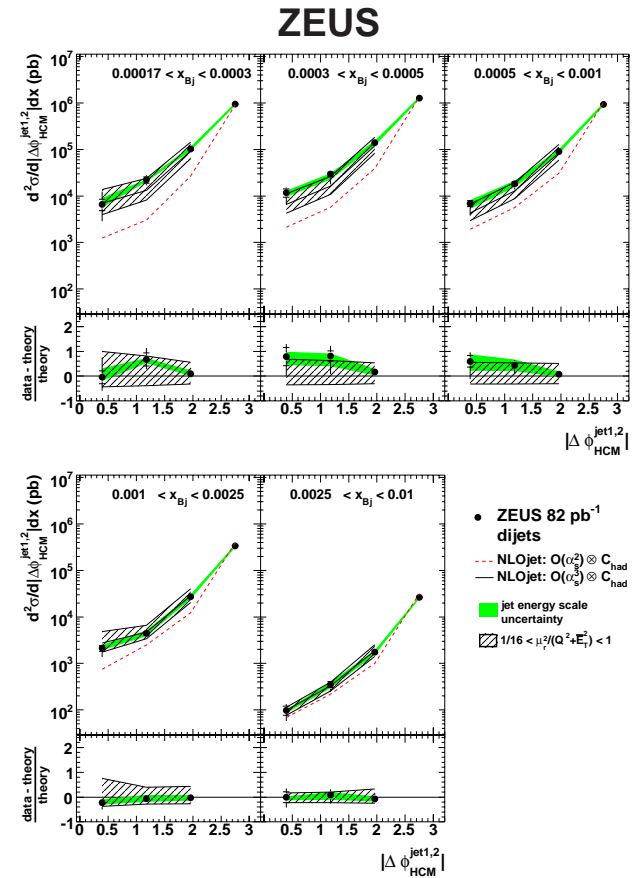


- ZEUS 82 pb⁻¹ trijets
- NLOjet: $O(\alpha_s^2) @ C_{had}$
- jet energy scale uncertainty
- ▨ $1/16 < \mu^2/(Q^2 + E_T^2) < 1$

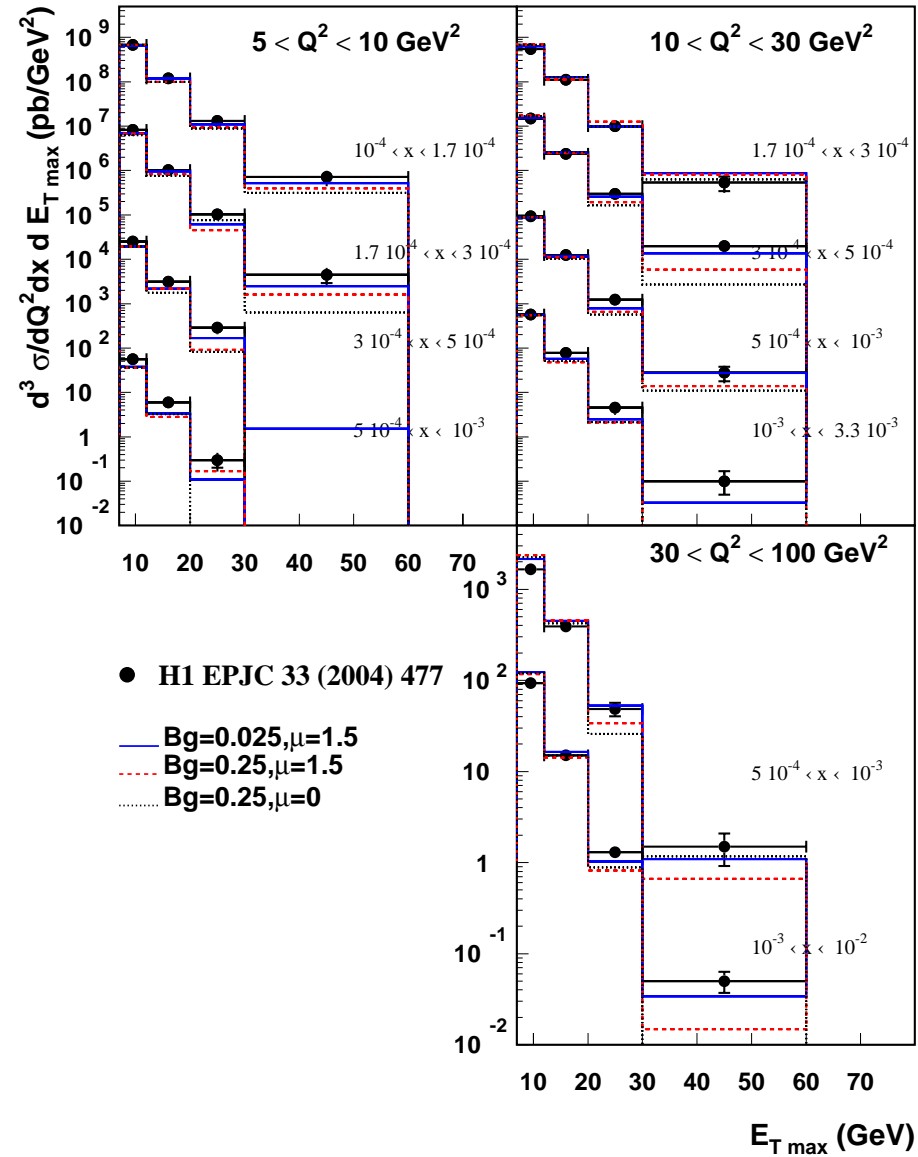
Azimuthal correlation in di-jet cross sections



- Large correction from order- α_s^2 to order- α_s^3 for small x and small ϕ



Inclusive jet E_T distribution



Remarks

- ▷ Physical picture from k_{\perp} -factorized MC is being probed quantitatively with
 - inclusive cross sections
 - detailed multi-jet correlations

- ▷ Main limitations still come from
 - limited knowledge of updf's
 - treatment of evolution
(how to combine Regge/Sudakov form factors?
subleading logs? how do multiple interactions
affect the picture? ...)

- ▷ Further: status of factorization proofs?
 - established only in simplest cases
 - e.g.: factorization-breaking from soft gluon exchanges
revisited in Collins & Qiu, arXiv:0705.2141
(potential N³LO effect in conventional calculations)

HOW TO CHARACTERIZE UPDF'S WITH PRECISION?

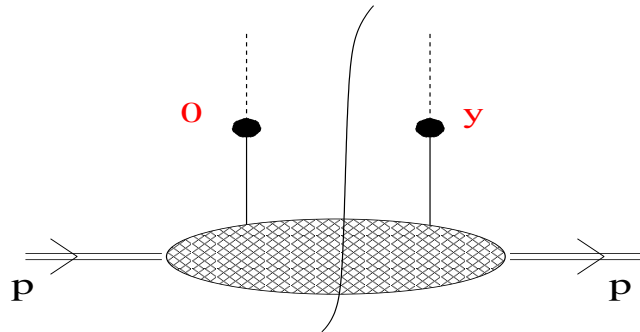
Collins & Zu, 2005

Boer & Mulders, 2003, 1998

Belitsky et al., 2004; Brodsky et al., 2001

...

$$\mathbf{p} = (p^+, m^2 / 2 p^+, \mathbf{0}_\perp)$$



$$\tilde{f}(y) = \langle P | \bar{\psi}(y) V_y^\dagger(n) \gamma^+ V_0(n) \psi(0) | P \rangle, \quad y = (0, y^-, y_\perp)$$

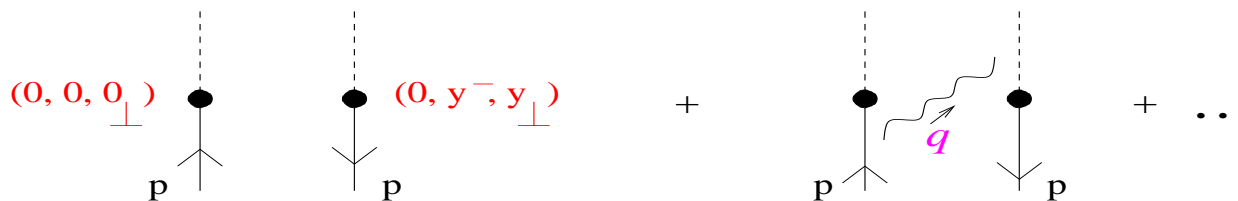
$$V_y(n) = \mathcal{P} \exp \left(ig_s \int_0^\infty d\tau n \cdot A(y + \tau n) \right)$$

- Fine at tree level
- Difficulties arise beyond this level



◇ Suppose a gluon is absorbed or emitted by eikonal line:

$$n = (0, 1, 0_\perp)$$



$$f_{(1)} = P_R(x, k_\perp) - \delta(1-x) \delta(k_\perp) \int dx' dk'_\perp P_R(x', k'_\perp)$$

where
$$P_R = \frac{\alpha_s C_F}{\pi^2} \left[\frac{1}{1-x} \frac{1}{k_\perp^2 + \rho^2} + \{\text{regular at } x \rightarrow 1\} \right]$$
 $\rho = \text{IR regulator}$

\uparrow
endpoint singularity ($q^+ \rightarrow 0, \forall k_\perp$)

◇ Physical observables:

$$\begin{aligned} \mathcal{O} &= \int dx dk_\perp f_{(1)}(x, k_\perp) \varphi(x, k_\perp) \\ &= \int dx dk_\perp [\varphi(x, k_\perp) - \varphi(1, 0_\perp)] P_R(x, k_\perp) \end{aligned}$$

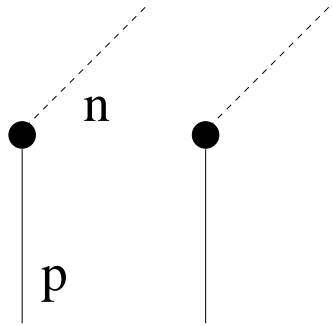
inclusive case: φ independent of $k_\perp \Rightarrow 1/(1-x)_+$ from real + virtual

general case: endpoint divergences from incomplete KLN cancellation

Traditionally, put **cut-off** on the endpoint region:

▷ e.g.: Monte-Carlo generators using u-pdf's

▷ **cut-off from gauge link in non-lightlike direction n :**



$$\eta = (\mathbf{p} \cdot \mathbf{n})^2 / n^2$$

Chen, Idilbi & Ji, hep-ph/0607003

Ji, Ma & Yuan, hep-ph/0503015

earlier work by Collins; Korchemsky

finite $\eta \Rightarrow$ singularity is cut off at $1 - x \gtrsim k_{\perp} / \sqrt{4\eta}$

Drawbacks:

- good for leading accuracy, but makes it difficult to go beyond

- $\int dk_{\perp} f(x, k_{\perp}, \mu, \eta) = F(x, \mu, \eta) \neq$ ordinary pdf

UPDF'S WITH SUBTRACTIVE REGULARIZATION

H, hep-ph/0702196

Collins, hep-ph/0304122

- Endpoint divergences $x \rightarrow 1$ from incomplete KLN cancellation

Subtractive method: more systematic than cut-off. Widely used in NLO calculations.

Formulation suitable for operator matrix elements: Collins & H, 2001.

- gauge link still evaluated at n lightlike, but multiplied by “subtraction factors”

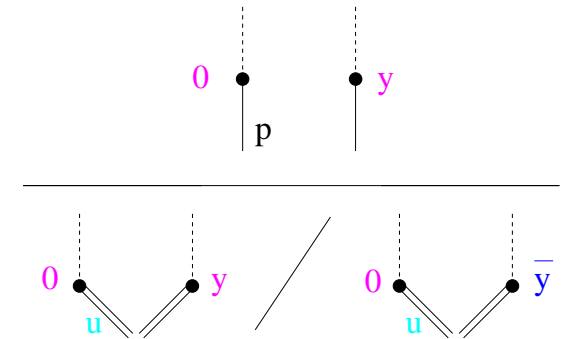
$$\tilde{f}^{(\text{subtr})}(y^-, y_\perp) =$$

original matrix element

$$\langle P | \bar{\psi}(y) V_y^\dagger(n) \gamma^+ V_0(n) \psi(0) | P \rangle$$

$$\frac{\langle 0 | V_y(u) V_y^\dagger(n) V_0(n) V_0^\dagger(u) | 0 \rangle}{\langle 0 | V_{\bar{y}}(u) V_{\bar{y}}^\dagger(n) V_0(n) V_0^\dagger(u) | 0 \rangle}$$

counterterms



$$\bar{y} = (0, y^-, 0_\perp); \quad u = \text{auxiliary non-lightlike eikonal } (u^+, u^-, 0_\perp)$$

◇ u serves to regularize the endpoint; drops out of distribution integrated over k_\perp

One loop:

$$[\zeta = (p^{+2}/2)u^-/u^+]$$

$$f_{(1)}^{(\text{subtr})}(x, k_{\perp}) = P_R(x, k_{\perp}) - \delta(1-x) \delta(k_{\perp}) \int dx' dk'_{\perp} P_R(x', k'_{\perp}) \\ - W_R(x, k_{\perp}, \zeta) + \delta(k_{\perp}) \int dk'_{\perp} W_R(x, k'_{\perp}, \zeta)$$

with $P_R = \alpha_s \left\{ 1/[(1-x)(k_{\perp}^2 + m^2(1-x)^2)] + \dots \right\}$ = real emission prob.

$W_R = \alpha_s \left\{ 1/[(1-x)(k_{\perp}^2 + 4\zeta(1-x)^2)] + \dots \right\}$ = counterterm

- ζ -dependence cancels upon integration in k_{\perp}

$$\Rightarrow \mathcal{O} = \int dx dk_{\perp} f_{(1)}^{(\text{subtr})}(x, k_{\perp}) \varphi(x, k_{\perp}) \\ = \int dx dk_{\perp} \{ P_R [\varphi(x, 0_{\perp}) - \varphi(1, 0_{\perp})] + (P_R - W_R) [\varphi(x, k_{\perp}) - \varphi(x, 0_{\perp})] \}$$

- first term: usual $1/(1-x)_+$ distribution
- second term: singularity in P_R cancelled by W_R

CONCLUSIONS

◇ k_{\perp} -MC with updf's and initial-state shower

being applied to description of multi-jet final states at small x

- Example: angular jet correlations in DIS
- Open issues on factorization, lack of complete KLN cancellation
 - ⇒ need to address new problems compared to ordinary pdf's
 - ▷ endpoint divergences ($x \rightarrow 1$):
 - more transparent representation in coordinate space
 - subtractive method as an alternative to cut-off method