

## OUTLINE

- (I) Physics Motivation: Single and di-photon direct production at LHC
- (II) Detector Performance Requirements: SM Higgs  $H \rightarrow \gamma \gamma$
- (III) Experimental Issues: Photon ID
- (IV) Experimental Issues: Trigger
- (V) Direct Photon Production
- (VI) Conclusions and Summary
- <u>Disclaimer</u>: Not all SM and Physics beyond the SM physics signatures with photons are covered in this talk: Wγγ, black holes, GMSB SUSY Models with non pointing photons, Exotics ...

## (I) LHC: Large Hadron Collider

- pp collisions at  $\sqrt{s}=14$  TeV
- Bunch crossing every 25 ns (40 MHz)
- Low Luminosity L = 10<sup>33</sup> cm<sup>-2</sup> s<sup>-1</sup> (ℒ=10fb<sup>-1</sup>/year)
- High Luminosity L = 10<sup>34</sup> cm<sup>-2</sup> s<sup>-1</sup> (£=100fb<sup>-1</sup>/year)



Process	Events for 10fb <sup>-1</sup>	σ (nb)
Inclusive bbbar	10 <sup>13</sup>	5x10⁵
Direct photon (p <sup>⊤</sup> >20 GeV)	10 <sup>7</sup>	100
Photon pairs (p <sub>T</sub> >20 GeV)	10 <sup>4</sup>	15

→ Large statistics: small statistical error

Production cross section and dynamics largely controlled by QCD

Mass reach up to ~ 5 TeV

Test QCD predictions and perform precision measurements

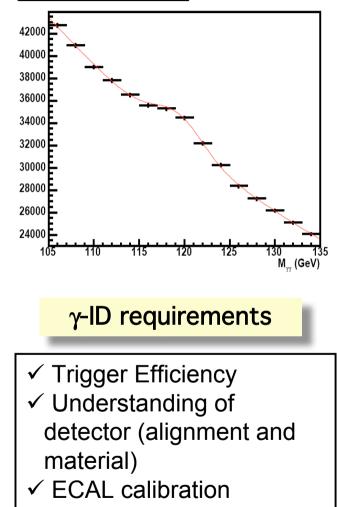
## (I) Physics Motivation

- The study of single photons and di-photons produced in the primary parton-parton interaction (direct photon production) is interesting in itself:
  - Their production and associated measurement provides a direct test of perturbative quantum chromodynamics (pQCD)
  - The coupling to the interaction partons provided by the photon allows the parton content of the proton to be probed directly (provides a possible constraint on the gluon content of the proton)
  - The topology of events with photons recoiling against a jet allows the hadronic calorimeter to be calibrated with the electromagnetic calorimeter using energy balance in the event.
- High p<sub>T</sub> single photons and photon pairs are important for the discovery of many Standard Model and "beyond the Standard Model processes", the measurement and understanding of direct photon production are essential for the search of new physics:
  - SM Higgs channel into  $\gamma\gamma$
  - Exotics di-photon production and SUSY physics

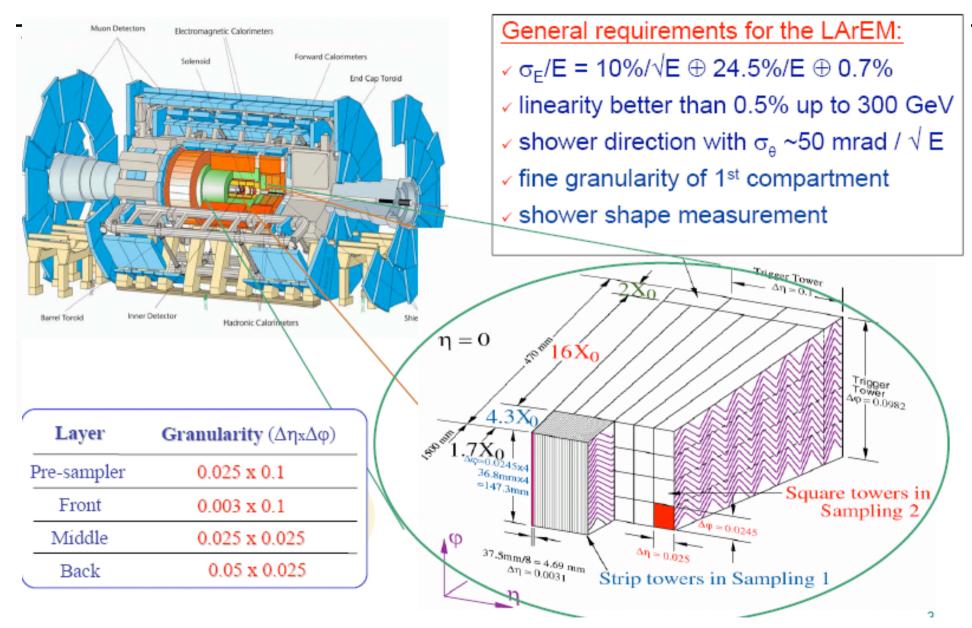
#### (II) Detector Performance Requirements

- $H \rightarrow \gamma \gamma$  is a rare decay mode with BR ~ 10<sup>-3</sup>
- The signal should be visible as a small peak above the γγ continuum background: need σ (m)/m~1%
  - Good energy resolution and uniformity of the EM calorimeter
- <u>Irreducible background</u> consists of genuine photons pairs continuum.
- <u>Reducible background</u> comes from jet-jet and gamma-jet events in which one or both jets are misidentified as photons (Reducible / irreducible cross section (LO) ~2x10<sup>6</sup>(jj) and ~ 8x10<sup>2</sup>(γj))
  - Excellent  $\gamma$ /jet and  $\gamma/\pi 0$  separation needed
  - Conversion recovery needed

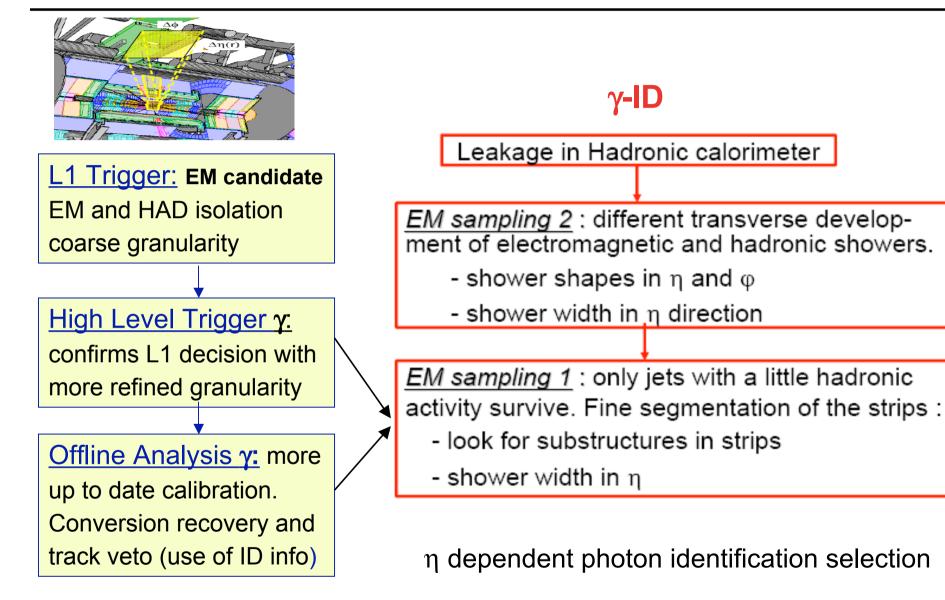
background+signal distribution



### (II)The ATLAS Detector

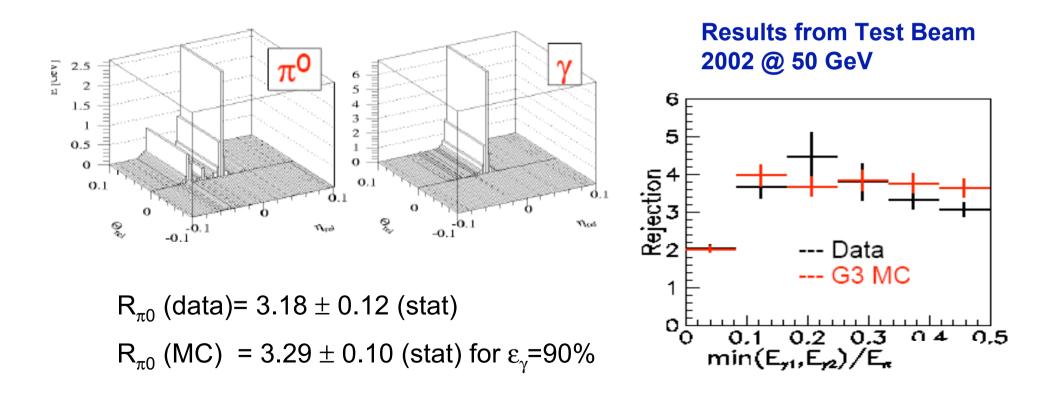


#### (III) Basis of $\gamma$ /jet and $\gamma/\pi 0$ separation



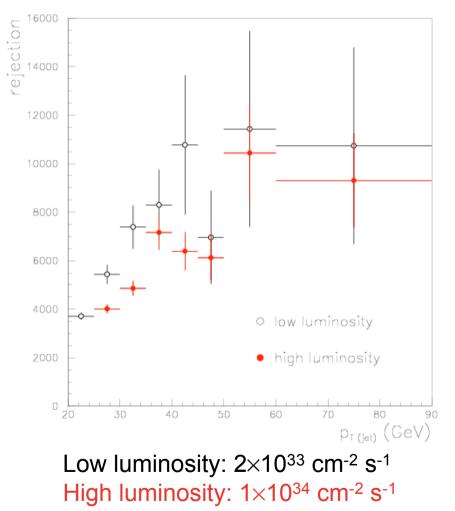
## (III) $\gamma/\pi 0$ separation

- After the application of hadronic leakage and 2nd EM sampling criteria, ~80% of the remaining background is composed of isolated  $\pi^0$  from jet fragmentation
- The high granularity of the 1<sup>st</sup> EM sampling provides additional rejection



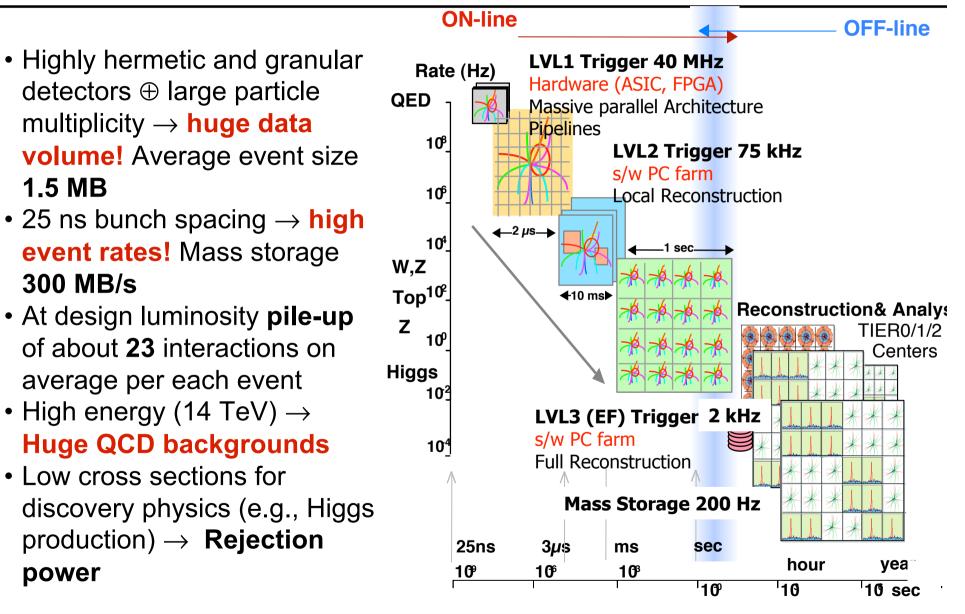
## (III) $\gamma$ /jet Separation

• Performance assessed with single  $\gamma$  of different energies or  $\gamma$  from H $\rightarrow\gamma\gamma$ 



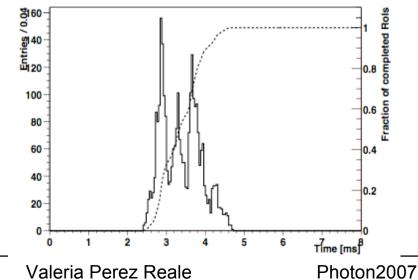
- For an εγ=80% (flat in η and p<sub>T</sub>) a
   R<sub>jet</sub> ~ 5000 can be achieved for p<sub>T</sub>
   > 25 GeV
- Looking at the jet origin:
  - R~ 3x10<sup>3</sup> on quark jets
  - R~ 2.1x10<sup>4</sup> on gluon jets
  - $\rightarrow$  Difference due to softer fragmentation function of gluon jets.
- The reducible background after photon id selection is reduced below the total irreducible background  $\gamma\gamma$

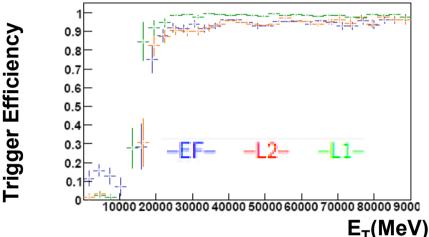
## **(IV) Trigger Requirements**



## (IV) Trigger: $\gamma$ /jet separation

- The optimal photon selection efficiency is a compromise between trigger efficiency, event rate allowed (few tens of Hz out of total 200 Hz for L=10<sup>33</sup> cm<sup>2</sup> s<sup>-1</sup>) and system performance limitations
- For photon target goal is 80% efficiency after the last trigger selection step (EF) for a rejection factor R~1000
- Trigger efficiency normalized wrt offline reconstructed kinematical cuts in  $E_{\rm T}$  and  $|\eta|$  <2.4





 Di-jet samples with p<sub>T</sub>>15 GeV are used for background studies

Trigger	Trigger Efficiency	Jet Rate
2γ20i	82.3 %	3 ±2 Hz
γ60	93.3 %	16±7 Hz

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## (IV) Triggering: $H \rightarrow \gamma \gamma$ Efficiency

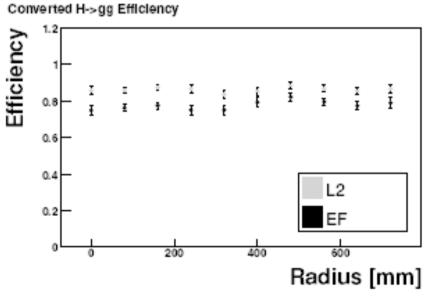
The H→γγ channel can be triggered with an efficiency of ~80% requiring two isolated photons with pT>20 GeV in the physics precision region of |η| < 2.4 at low luminosity L=2x10<sup>33</sup> cm<sup>2</sup>s<sup>-1</sup>

Trigger 2γ20i	Trigger Efficiency
L1	96.0 ±0.8 %
L2	88.6 ±1.3 %
EF	85.4 ±1.5 %

•The addition of pile-up reduces the trigger efficiency ~2% (4%) at low (high) luminosity)

- The trigger selection is efficiency selecting both converted and non converted photons of a  $H \rightarrow \gamma \gamma$  event
  - Non converted:  $80.2 \pm 0.1\%$





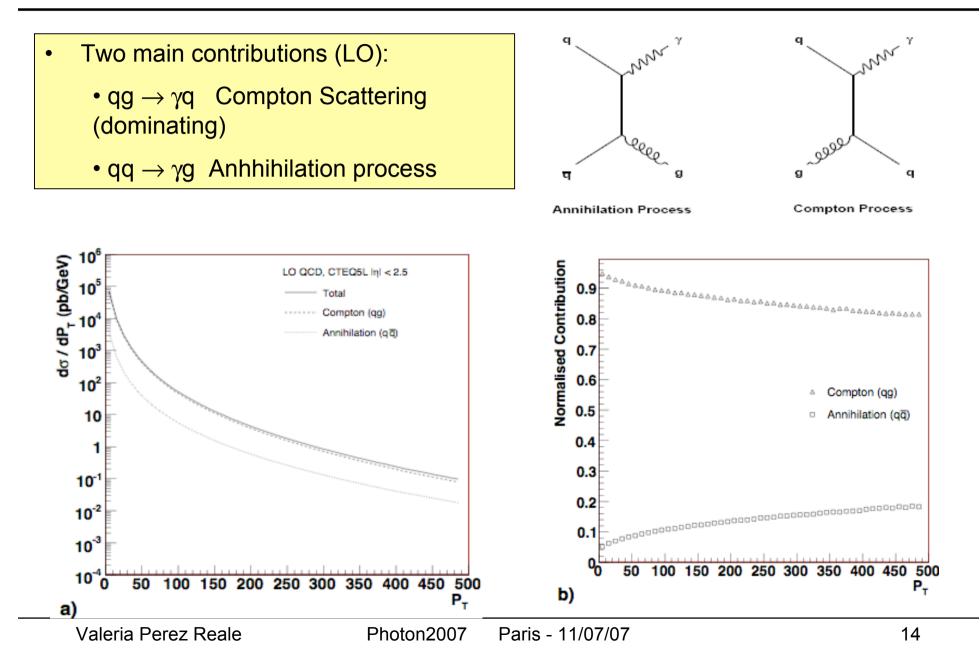
## (IV) Triggering: Direct Photon

- The direct photon can be triggered requiring one photon with **p**<sub>T</sub>>20 GeV in the "precision physics" region of  $|\eta| < 2.4$  (simulation results for initial luminosity running L=1x10<sup>31</sup> cm<sup>2</sup>s<sup>-1</sup>)
- Signal sample:  $\gamma$ +jet ( $\gamma$  generated with  $p_T$ >10 GeV)

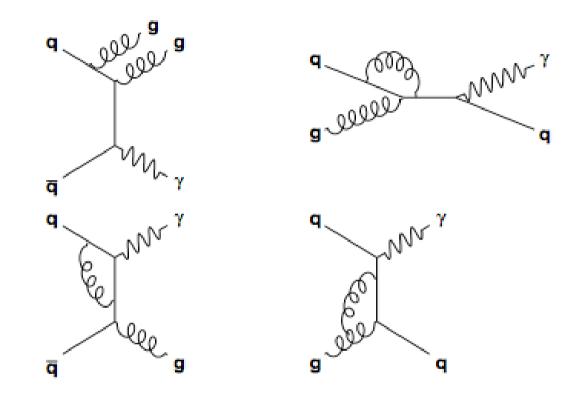
	Jet energy range (GeV)	Trigger Efficiency%) γ20
1	17-35	75.1 ±0.3
2	35-70	83.5 ±0.3
3	70-140	89.3 ±0.2
4	140-280	91.7 ±0.2
5	280-560	94.4 ±0.2
6	560-1120	92.4 ±1.1

• Background from jet-jet sample generated with  $p_T$ >15 GeV is 7 Hz

#### (V) Direct Photon Production (LO)



#### (V) Direct Photon Production (NLO)



NLO real diagram and corresponding virtual diagrams of Direct Photon Events.

## (V) Signal and Background

- The typical event topology of direct photon production in the ATLAS detector will be the observance of a well-isolated photon recoiling against a jet
  - At LO these events should be back-to-back in the r- $\phi$  plane and display a balance of energy between the jet and the photon.
- From QCD it is known that the jet rate will be ~3 orders of magnitude large than that for direct photons

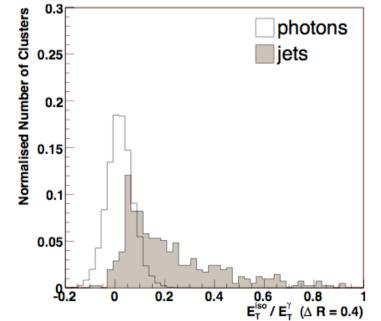
•The required good photon identification is achieved using the highly-granular EM calrimeter

 Main background related to fragmentation non perturbative QCD

• The  $\pi$ 0s background from jet fragmentation is reduced by requiring a selection in the 1st sampling of EM

• The background is reduced further by requiring an **isolation cone** (EM, HAD, tracking)

$$\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$$



## (V) Systematic Errors

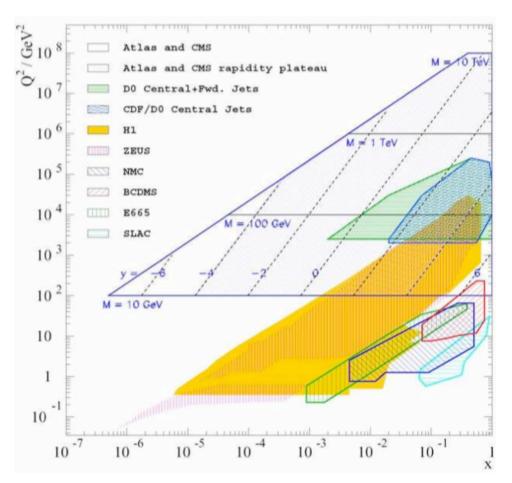
- Sources of experimental systematic errors within the direct photon measurement expected in ATLAS (first approximation):
  - <u>Luminosity error</u>: In the initial phase of LHC operation for an integrated luminosity of 1fb<sup>-1</sup>, the error on the ATLAS measured luminosity ~10% (aim to reduce error significantly, e.g. Using forward detectors)
  - <u>Absolute EM energy scale</u>: should be ultimately known ~1%. When convolved with falling E<sub>T</sub> spectrum of direct photons <5 %</li>
  - <u>Preselection efficiency</u>: <5% for signal and background
  - <u>Photon Trigger Efficiency Error</u>: is expected to be ~1%
  - <u>Background subtraction</u> provides the other major source of background, a conservative estimate on its effect on the cross section is 10%
  - Some of these uncertainties (luminosity, photon trigger efficiency) will cancel out in the S/B measurement contrary to theoretical uncertainties (pdfs, scale variation)
  - Aim for a precision on the cross-section determination similar (hopefully better) than observed at D0/CDF (~15-22%)
  - ATLAS measurement at much higher  $\sqrt{s}$  will extend to very high  $p_{T}$  photons

#### (V) Constrains on Gluon Structure

The kinematic acceptance of the ATLAS detector allows a wide range of x and Q<sup>2</sup> to be probed (|y|<5)</li>

$$x_{min} = \frac{x_T e^{-\eta}}{2 - x_T e^{\eta}} \qquad \qquad x_T = 2p_T / \sqrt{s}$$

- ATLAS will be sensitive to the gluon fraction x below 10<sup>-4</sup> within an energy scale Q<sup>2</sup> above 100
   GeV<sup>2</sup>
- The highest-energy photons will give access to large x values in the range ~10<sup>-1</sup>



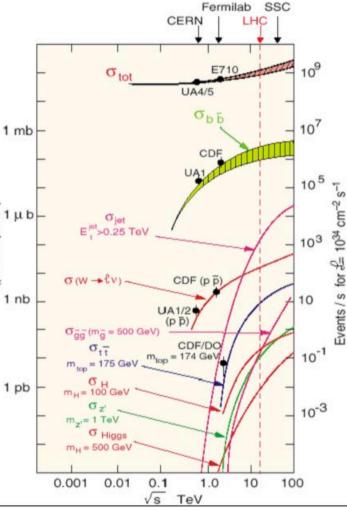
## **(VI) CONCLUSIONS**

- The LHC will provide large statistics of single and di-photon direct production in the first year of data taking (integrated luminosity of 10 fb<sup>-1</sup>) in the energy region √14 TeV
- The ATLAS Hadronic and Electromagentic calorimeter has proven to have the  $\gamma$ /jet separation and  $\gamma/\pi^0$  separation capability needed to observe direct photon signal and SM H $\rightarrow\gamma\gamma$  over the QCD reducible background
  - Jet rejection factor of 5000 for photon efficiency of 80%
  - $\pi^0$  rejection factor of 3 for photon efficiency of 90%
- ATLAS will measure direct photons at higher  $E_T$  and with a significant improvement in precision and a new  $\sqrt{s}$  wrt to other experiments
- The full potential of the direct photon production process has yet to be realised, the LHC will provide an opportunity to determine the gluon density function in a proton to a new kinematic region of x (~10<sup>-4</sup><x<0.1) and Q<sup>2</sup> (10<sup>2</sup><Q<sup>2</sup><10<sup>5</sup> GeV<sup>2</sup>)

# **BACK-UP SLIDES**

#### **LHC Machine Parameters**

				<b>_</b>	
Energy .	Ê	[TeV]	7.0		
Dipole field	В	T	8.4		1 1
Luminosity	L	$[\mathrm{cm}^{-2} \mathrm{s}^{-1}]$	10 <sup>34</sup>		$\sigma_{tot}$
Beam-beam parameter	ξ		0.0034		
Total beam-beam tune spread			0.01		
Injection energy	$E_{\rm i}$	[GeV]	450	1 mb -	4
Circulating current/beam	I <sub>beam</sub>	[A]	0.53		
Number of bunches	$k_{\rm b}$		2835		
Harmonic number	$h_{ m RF}$		35640		1
Bunch spacing	$\tau_{\rm b}$	[ns]	24.95	θub	Giat
Particles per bunch	$n_{\rm b}$		$1.05 \ 10^{11}$	a (proton - proton) απt	σ <sub>jet</sub> E <sup>jet</sup> >0.25 Te
Stored beam energy	$E_{\rm s}$	[MJ]	334	- 10	
Normalized transverse emittance $(\beta \gamma) \sigma^2 / \beta$	$\varepsilon_{\mathbf{n}}$	$[\mu m.rad]$	3.75	brot	σ(w - tv)-
	ļ			ل 1 nb	0(11-01)-
Collisions				1110	
$\beta$ -value at I.P.	$\beta^*$	[m]	0.5	-	$\sigma_{\bar{g}\bar{g}}$ (m_{\bar{g}} = 500 G
r.m.s. beam radius at I.P.	$\sigma^*$	[µm]	16		$\sigma_{t\overline{t}}$
r.m.s. divergence at I.P.	$\sigma'^*$	$[\mu rad]$	32	87 - a 105	m <sub>top</sub> = 175 GeV
Luminosity per bunch collision	$L_{\mathbf{b}}$	$[cm^{-2}]$	$3.14 \ 10^{26}$	1 pb -	σ <sub>H</sub> m <sub>H</sub> = 100 GeV ~
Crossing angle	$\phi$	$[\mu rad]$	200	_	σ <sub>z</sub> , m <sub>z</sub> = 1 TeV
Number of events per crossing	$n_{\rm c}$		19		σ <sub>Higgs</sub> m <sub>H</sub> = 500 GeV
Beam lifetime	$\tau_{ m beam}$	[h]	22		m <sub>H</sub> = 500 GeV
Luminosity lifetime	$ au_L$	[h]	10		0.001 0.01
	•		•		0.001 0.01



#### **Limiting factor for** $\sqrt{s}$ **:** Bending power needed to keep beams in 27 km LHC ring:

p(TeV) = 0.3 B(T) R(km)

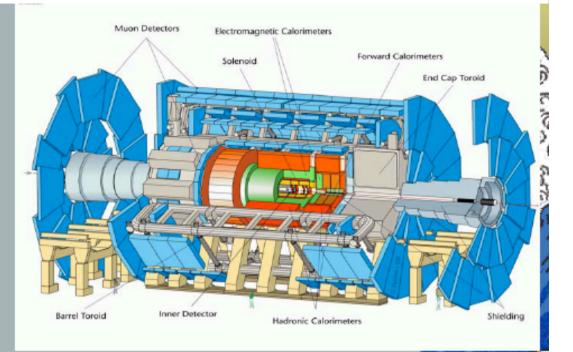
With typical magnet packing factor of ~ 70%, need 1232 dipoles with B = 8.3 T for 7 TeV beams

### **The ATLAS Detector**

- Inner Detector (tracker): Si pixel & strip detectors + TRT;
   2 T magnetic field; |η| < 2.5</li>
- Calorimetry: highly granular LAr EM calorimeter (|η|< 3.2); hadron calorimeter – scintillator tile; |η|< 4.9</li>
- ▲ *Muon Spectrometer*: air-core toroid system;  $|\eta| < 2.7$

#### ➢Performance:

- jet resolution:  $\sigma/E\approx 50\%/\sqrt{E\oplus 3\%}$
- τ-efficiency: ~50% for R<sub>jet</sub>~200 (p<sub>T</sub>≈60 GeV)
- missing energy:  $\sigma(p_{xy}^{miss}) \approx 0.46 \cdot \sqrt{\sum E_T}$ (low luminosity)
- b-tagging: ~60% for R<sub>jet</sub>~100 (low luminosity)



- ★ Jet energy scale: precision of 1% ( $W \rightarrow jj$ ; Z (ll) + jets)
- ▲ Luminosity: precision ≤ 5% (machine, optical theorem, rate of known processes)
- QCD-related measurements performed during initial period of running at low luminosity

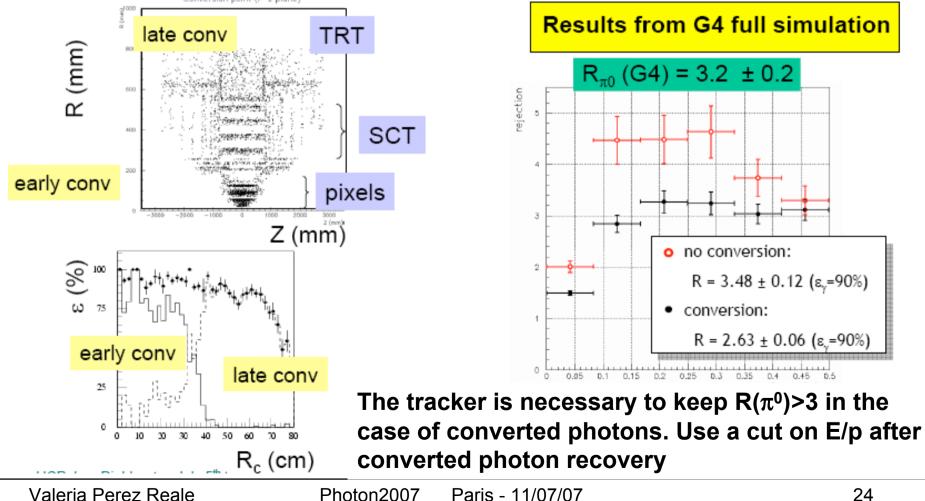
#### **The ATLAS Tracker**

**Forward SCT Barrel SCT** The Inner Detector (ID) is organized into four sub-systems: **Pixels 1** removable barrel layer **2** barrel layers 4 end-cap disks on each side (0.8 10<sup>8</sup> channels) Silicon Tracker (SCT) **4 barrel layers** TRT 9 end-cap wheels on each side (6 10<sup>6</sup> channels) **Pixel Detectors Transition Radiation Tracker (TRT) Axial barrel straws Radial end-cap straws** 36 straws per track (4 10<sup>5</sup> channels) Barrel **Common ID items TRT+SCT** ACTINE INC.

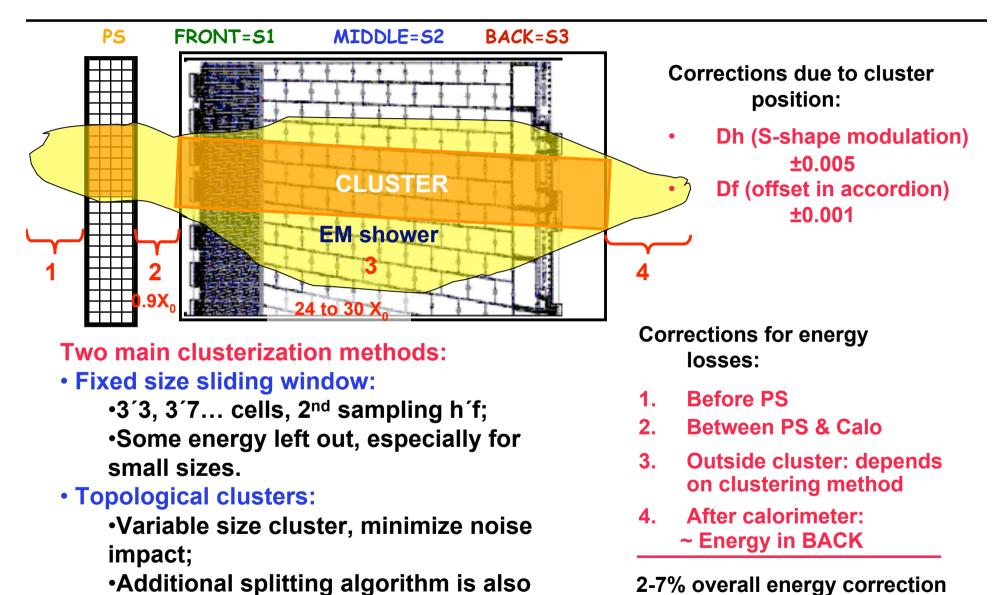
#### $\gamma/\pi 0$ : Effects of $\gamma$ conversions

#### γ conversions

 $\Rightarrow \sim 30\%$  (depending on  $\eta$ ) probability for photon conversion in the ID cavity ✤ ID will identify and reconstruct with a ~80% efficiency photon conversions in the region  $R_c < 80$  cm and |z| < 280 cm – where ~80% of conversions occur



#### **EM Calorimeter energy reconstruction**



provided.

2-7% overall energy correction >7% at low energy, high h

#### **Energy Calibration**

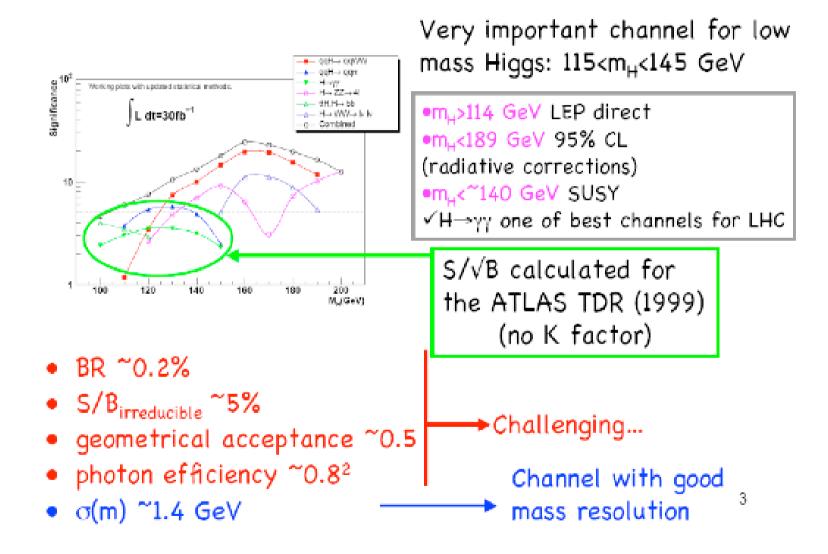
$$E_{rec} = \lambda \left( off + w_0 E_0 + E_1 + E_2 + w_3 E_3 \right)$$

The 4 coefficients are reconstructed via c<sup>2</sup> fit on a sample of single electrons in a [-2s,+3s] range around the most probable value of the reconstructed energy distribution:

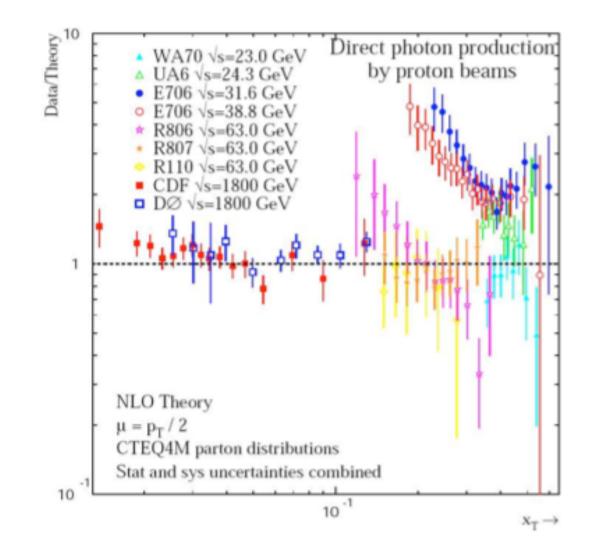
$$\chi^{2} = \sum_{i}^{N} \frac{\left(E_{rec}^{i} - E_{true}^{i}\right)^{2}}{\sigma_{E}^{2}}$$

- Simple method.
- 4-parameters, h-dependent, energy-independent.
- Weights absorb different effects and their energy dependence (offset and  $w_0$  absorb energy loss upstream the calorimeter, and between the presampler and the strips).
- It is not possible to unfold these effects. More complex approach relying on detailed understanding of MC under study.

### $H \rightarrow \gamma \gamma$ Channel



#### **Comparison of Data with NLO**



### **Parton Luminosities and pdfs**

 $N_{events}(pp \to X) = L_{p-p} \times pdf(x_1, x_2, Q^2) \times \sigma_{theory}(q, \overline{q}, g \to X)$ 

Uncertainties in **p-p luminosity** ( $\pm$ 5%) and **p.d.f.'s** ( $\pm$ 5%) will limit measurement uncertainties to  $\pm$ 5% (at best).

• For high Q<sup>2</sup> processes LHC should be considered as a parton-parton collider instead of a p-p collider.

• Using only relative cross section measurements, might lead eventually to accuracies of  $\pm 1\%$ .

$\begin{array}{c} q\bar{q} \; (u,d) \\ \text{(high-mass DY lepton} \\ \text{pairs and other processes} \\ \text{dominated by } q\bar{q} \; \text{)} \end{array}$	W <sup>±</sup> and Z leptonic decays	<ul> <li>precise measurements of mass and couplings;</li> <li>huge cross-sections (~nb);</li> <li>small background.</li> <li>x-range: 0.0003 - 0.1</li> <li>± 1%</li> </ul>
g (high-Q <sup>2</sup> reactions involving gluons)	γ- <b>jet</b> , Z-jet, W±-jet	<ul> <li>γ-jet studies: γ p<sub>T</sub> &gt; 40 GeV</li> <li>x-range: 0.0005 - 0.2</li> <li>γ-jet events: γ p<sub>T</sub> ~ 10-20 GeV</li> <li>low-x: ~ 0.0001</li> <li>±1%</li> </ul>
s, c, b	<mark>γ</mark> c,γb,sg→Wc	<ul> <li>quark flavour tagged γ-jet final states;</li> <li>use inclusive high-p<sub>T</sub> μ and b-jet identification (lifetime tagging) for c and b;</li> <li>use μ to tag c-jets;</li> <li>5-10% uncertainty for x-range: 0.0005 – 0.2</li> </ul>
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