

# Analysis of elastic pp and p $\bar{p}$ scattering from a unitary extension of Bialas-Bzdak model

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p+p @ ISR and @ 7 TeV LHC

Real extension of Bialas-Bzdak

New: focusing ReBB on low t region of  $d\sigma/dt$  of pp and p $\bar{p}$

New: excitation functions

[arXiv:1204.5617](https://arxiv.org/abs/1204.5617)

[arXiv:1306.4217](https://arxiv.org/abs/1306.4217)

[arXiv:1311.2308](https://arxiv.org/abs/1311.2308)

[arxiv:1505.01415](https://arxiv.org/abs/1505.01415)

+ manuscript in preparation

# S-matrix Unitarity, Optical Theorem

$$SS^\dagger = I,$$

$$S = I + iT$$

$$T - T^\dagger = iTT^\dagger$$

$$2 \operatorname{Im} t_{el}(s, b) = |t_{el}(s, b)|^2 + \sigma(s, b)$$

Note: diffraction also measures  
|Fourier-transform|<sup>2</sup> images of  
sources of elastic scattering

- ideal for femtoscopic studies
- several similarities e.g. non-Gaussian sources etc

Black (grey) disc limit (important)

$$\rightarrow \sigma(b) \sim \theta(R-b)$$

# Diffraction in quark-diquark models

$$\frac{d\sigma}{dt} = \frac{1}{4\pi} |T(\Delta)|^2.$$

Bialas and Bzdak,  
Acta Phys. Polon. B 38 (2007) 159  
 $p = (q, d)$  or  $p = (q, (q, q))$

$$T(\vec{\Delta}) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} t_{el}(\vec{b}) e^{i\vec{\Delta} \cdot \vec{b}} d^2b = 2\pi \int_0^{+\infty} t_{el}(b) J_0(\Delta b) b db,$$

$$t_{el}(\vec{b}) = 1 - \sqrt{1 - \sigma(\vec{b})}.$$

$\sigma(b) = b$  dependent prob. of interaction  
→ connection to scattering centers

$$\sigma(\vec{b}) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d^2s_q d^2s'_q d^2s_d d^2s'_d D(\vec{s}_q, \vec{s}_d) D(\vec{s}'_q, \vec{s}'_d) \sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b}),$$

Structure of protons = ?

→ Diffractive pp at ISR (23.5 – 62.5 GeV) and LHC (7 - 8 TeV).

# Diffraction a la Bialas and Bzdak

$$D(\vec{s}_q, \vec{s}_d) = \frac{1 + \lambda^2}{\pi R_{qd}^2} e^{-(s_q^2 + s_d^2)/R_{qd}^2} \delta^2(\vec{s}_d + \lambda \vec{s}_q), \quad \lambda = m_q/m_d,$$

$$\sigma(\vec{s}_q, \vec{s}_d; \vec{s}_q', \vec{s}_d'; \vec{b}) = 1 - \prod_{a,b \in \{q,d\}} \left[ 1 - \sigma_{ab}(\vec{b} + \vec{s}_a' - \vec{s}_b') \right]$$

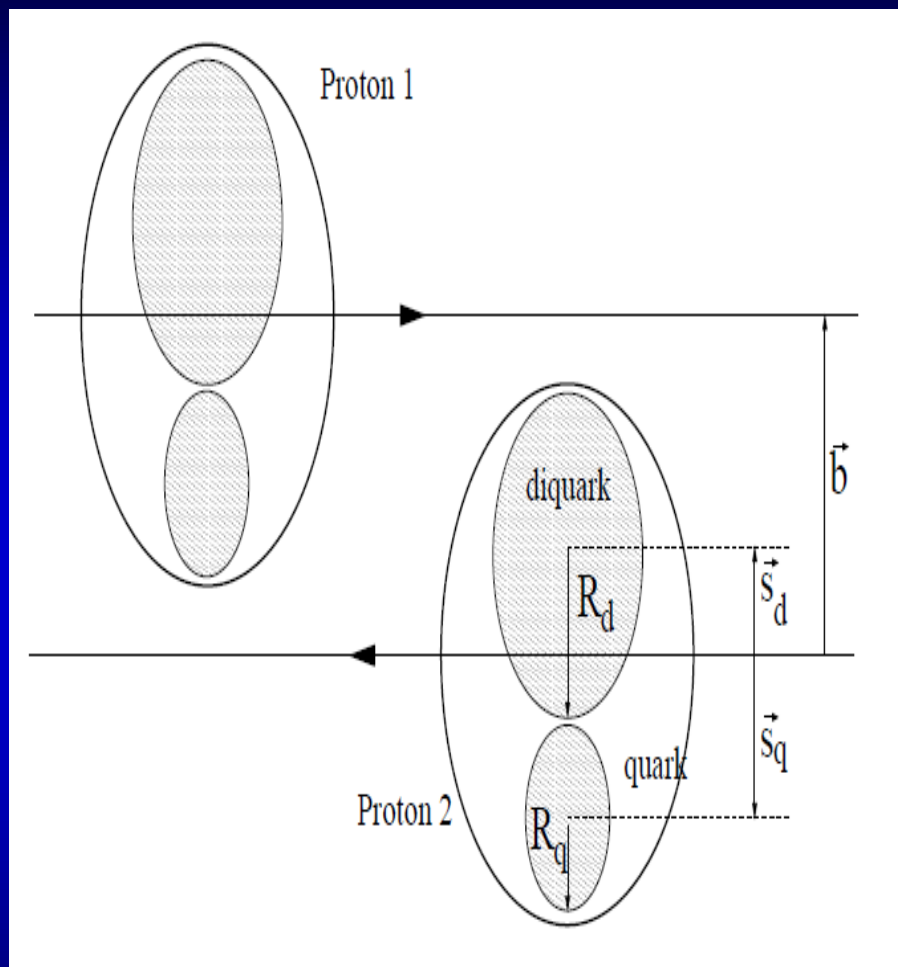
$$\sigma_{ab}(\vec{s}) = A_{ab} e^{-s^2/R_{ab}^2}, \quad R_{ab}^2 = R_a^2 + R_b^2,$$

The quark-diquark model of Bialas and Bzdak has been analytically integrated in a Gaussian approximation, **assuming** that the real part of forward scattering is negligible.

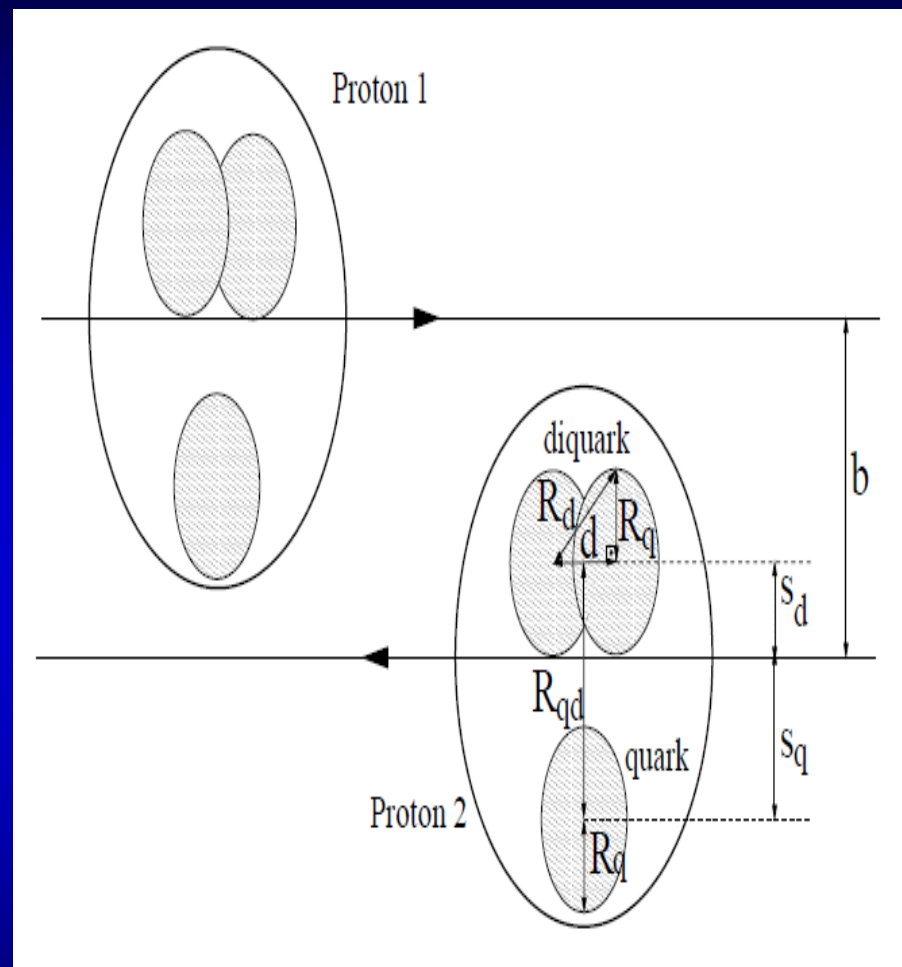
Two different pictures:  $p = (q, d)$  or  $p = (q, (q, q))$

Note:  $p = (q, q, q)$  model fails, quarks are correlated  
W. Czyz and L. C. Maximon, Annals. Phys. 52 (1969) 59

# Diffractionive pp scattering



$$p = (q, d)$$



$$p = (q, (q, q))$$

# Real extended BB model for the dip

$$\frac{d\sigma}{dt} = \frac{1}{4\pi} |T(\Delta)|^2.$$

$$T(\vec{\Delta}) = \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} t_{el}(\vec{b}) e^{i\vec{\Delta} \cdot \vec{b}} d^2b = 2\pi \int_0^{+\infty} t_{el}(b) J_0(\Delta b) b db,$$

$$t_{el}(s, b) = i \left( 1 - e^{-i \text{Im} \Omega(s, b)} \sqrt{1 - \sigma(s, b)} \right)$$

Bialas-Bzdak obtained  
if  $\text{Re}(t_{el}) = 0$

$$t_{el}(s, b) = i \left( 1 - e^{-\text{Re} \Omega(s, b)} \right) = i \left( 1 - \sqrt{1 - \sigma(s, b)} \right)$$

$$\sigma(\vec{b}) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d^2s_q d^2s'_q d^2s_d d^2s'_d D(\vec{s}_q, \vec{s}_d) D(\vec{s}'_q, \vec{s}'_d) \sigma(\vec{s}_q, \vec{s}_d; \vec{s}'_q, \vec{s}'_d; \vec{b}),$$

Real extension of an imaginary  $t_{el}$   
New parameter  $\text{Im} \Omega$  added

# ReBB model for the dip (2)

$$\sigma(b) = \int_{-\infty}^{+\infty} \dots \int_{-\infty}^{+\infty} d^2s_q d^2s'_q d^2s_d d^2s'_d D(\mathbf{s}_q, \mathbf{s}_d) D(\mathbf{s}'_q, \mathbf{s}'_d), \sigma(\mathbf{s}_q, \mathbf{s}_d; \mathbf{s}'_q, \mathbf{s}'_d; \mathbf{b}).$$

$$D(\mathbf{s}_q, \mathbf{s}_d) = \frac{1 + \lambda^2}{R_{qd}^2 \pi} e^{-(s_q^2 + s_d^2)/R_{qd}^2} \delta^2(\mathbf{s}_d + \lambda \mathbf{s}_q), \quad \lambda = \frac{m_q}{m_d},$$

$$\sigma(\mathbf{s}_q, \mathbf{s}_d; \mathbf{s}'_q, \mathbf{s}'_d; \mathbf{b}) = 1 - \prod_{a,b \in \{q,d\}} [1 - \sigma_{ab}(\mathbf{b} + \mathbf{s}'_a - \mathbf{s}_b)]$$

$$\sigma_{ab}(\mathbf{s}) = A_{ab} e^{-s^2/R_{ab}^2}, \quad R_{ab}^2 = R_a^2 + R_b^2, \quad a, b \in \{q, d\}$$

$$\sigma_{qq} : \sigma_{qd} : \sigma_{dd} = 1 : 2 : 4$$

Bialas-Bzdak  
model is  
„realized”:  
p = (q,d)  
p = (q, (q,q))

# ReBB model: two choices

$$\text{Im } \Omega(s, b) = -\alpha \cdot \text{Re } \Omega(s, b).$$

Similar to a constant  $\rho$  but not favored by data

$$\text{Im } \Omega(s, b) = -\alpha \cdot \tilde{\sigma}_{inel}(s, b),$$

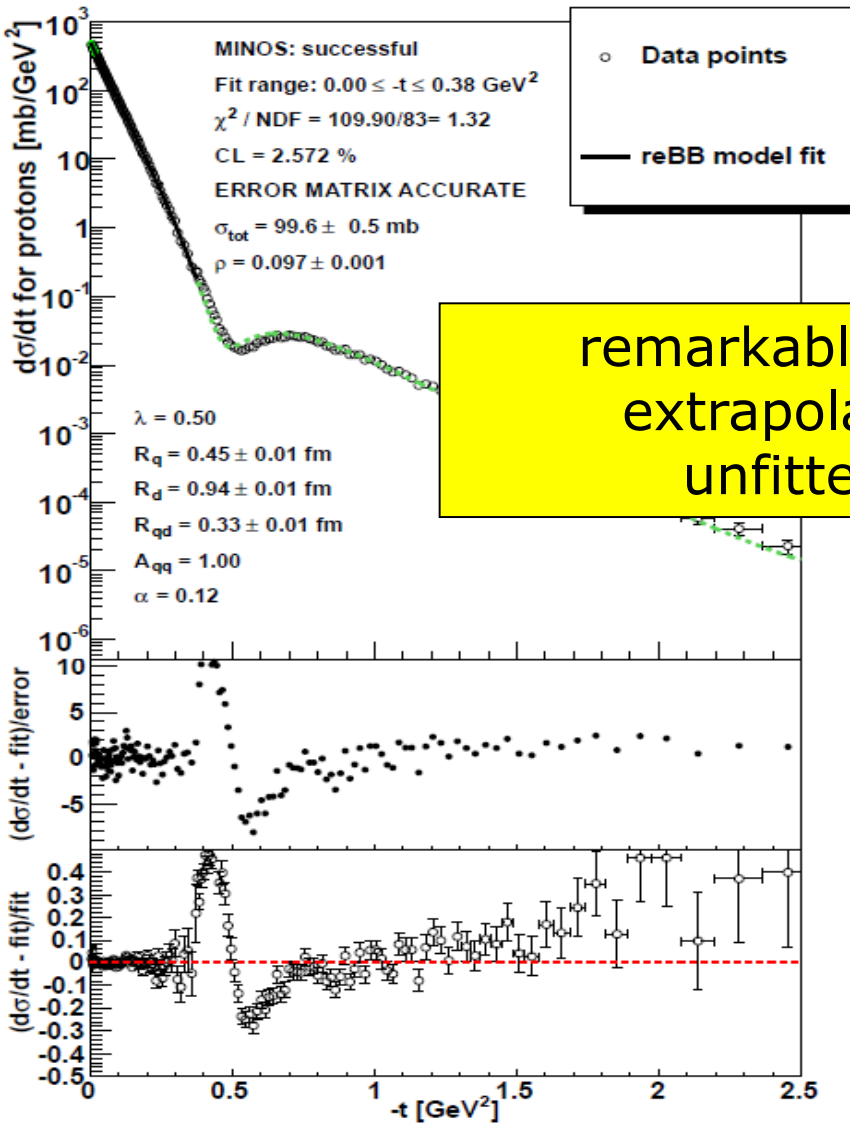
For small values of  $\alpha$  we recover our first attempt, the  $\alpha$ BB model

This choice is also favoured by data  
T. Cs., F. Nemes, arxiv:1306.4217

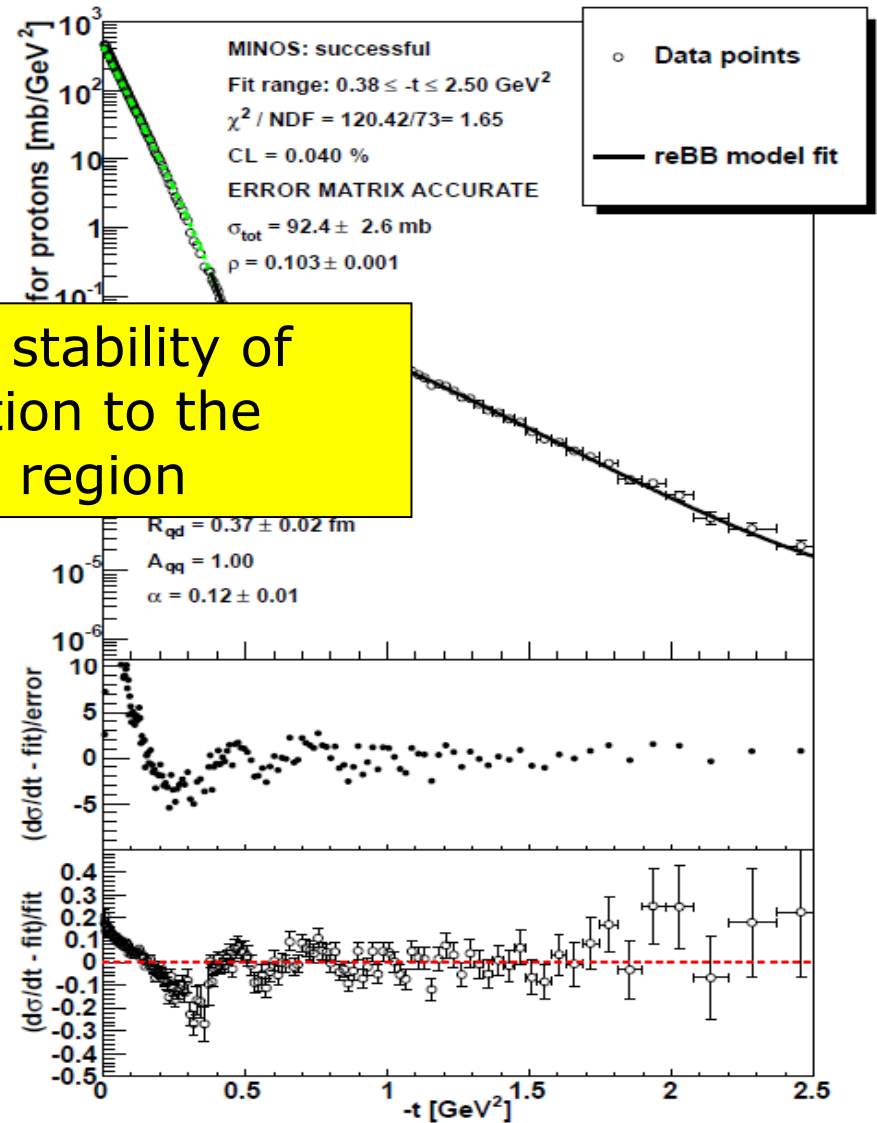


# ReBB model, final fit range studies

p+p → p+p, diquark as a single entity at  $\sqrt{s}=7000.0$  GeV



p+p → p+p, diquark as a single entity at  $\sqrt{s}=7000.0$  GeV



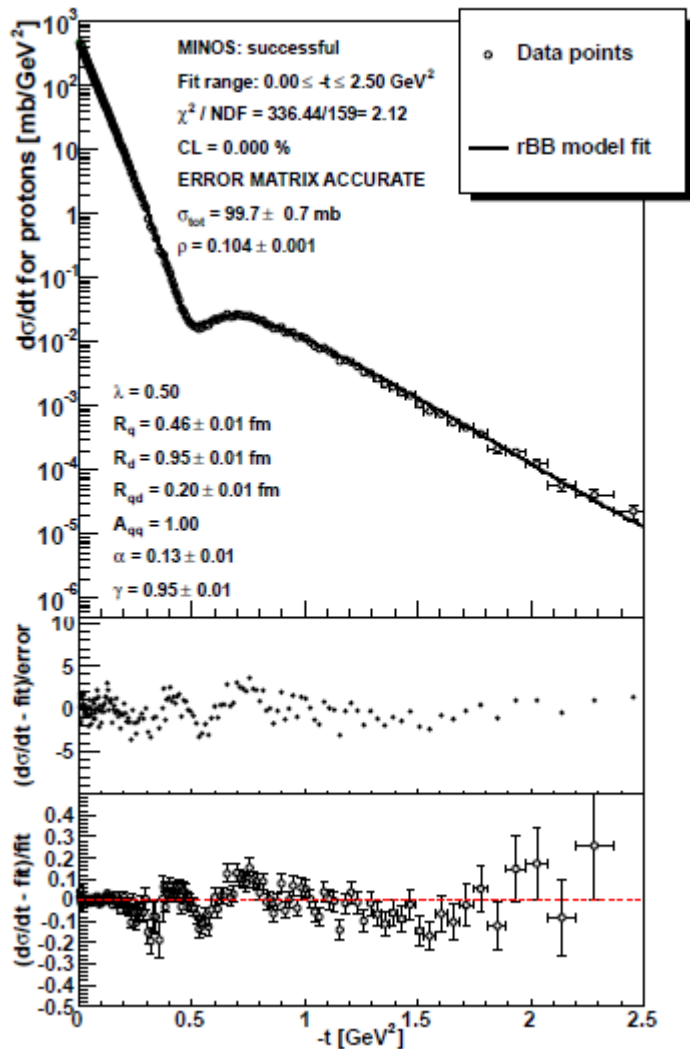
remarkable stability of extrapolation to the unfitted region

fit:  $0.36 \leq -t \leq 2.5$  GeV<sup>2</sup>, OK

fit:  $0 \leq -t \leq 2.5$  GeV<sup>2</sup>, ~ OK

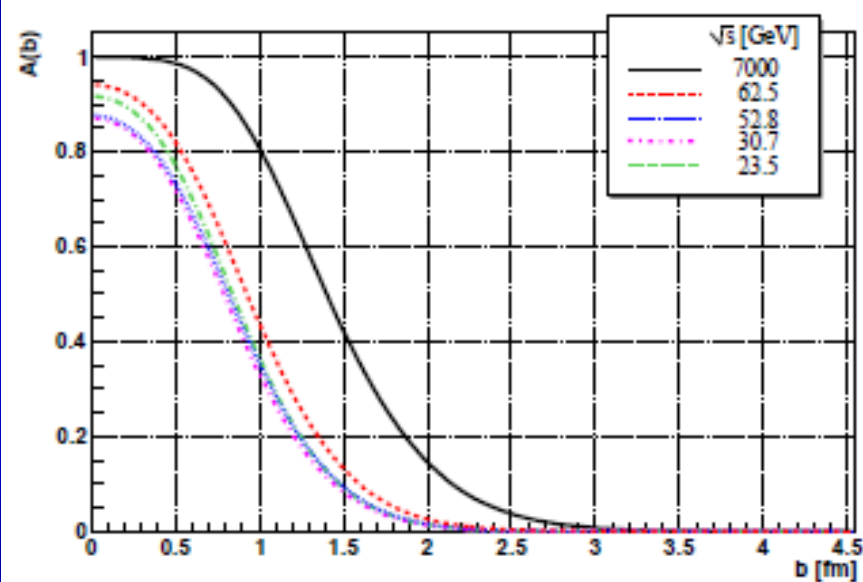
# ReBB model, combined data sets

p+p → p+p, diquark as a single entity at  $\sqrt{s}=7000.0$  GeV



## Shadow profile function

$$A(s, b) = 1 - |\exp[-\Omega(s, b)]|^2$$



$$\frac{d\sigma}{dt} \rightarrow \gamma \cdot \frac{d\sigma}{dt}$$

$$t_{\text{sep}} = -0.375 \text{ GeV}^2$$

Fit range:  $0 \leq -t \leq 2.5$  GeV<sup>2</sup>,  
 not quite OK → check @ 8 TeV

# ReBB shadow profile functions

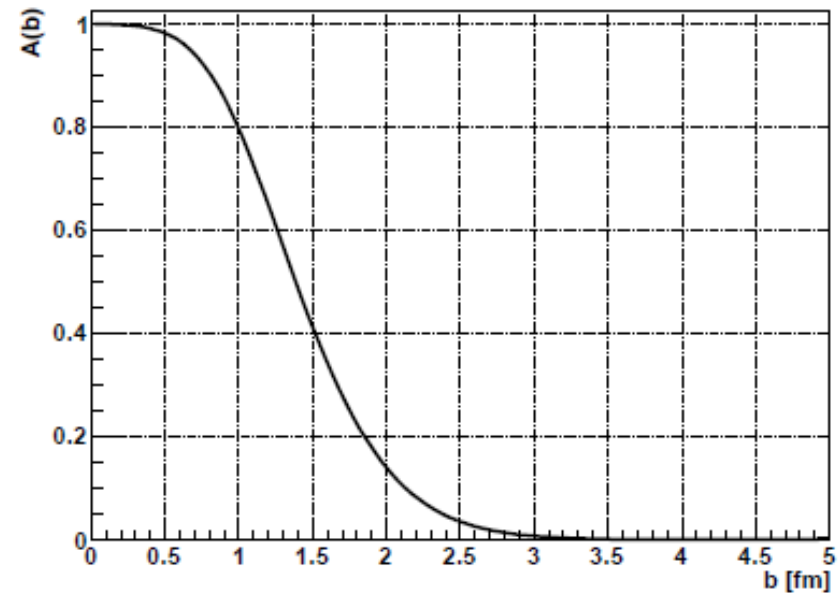
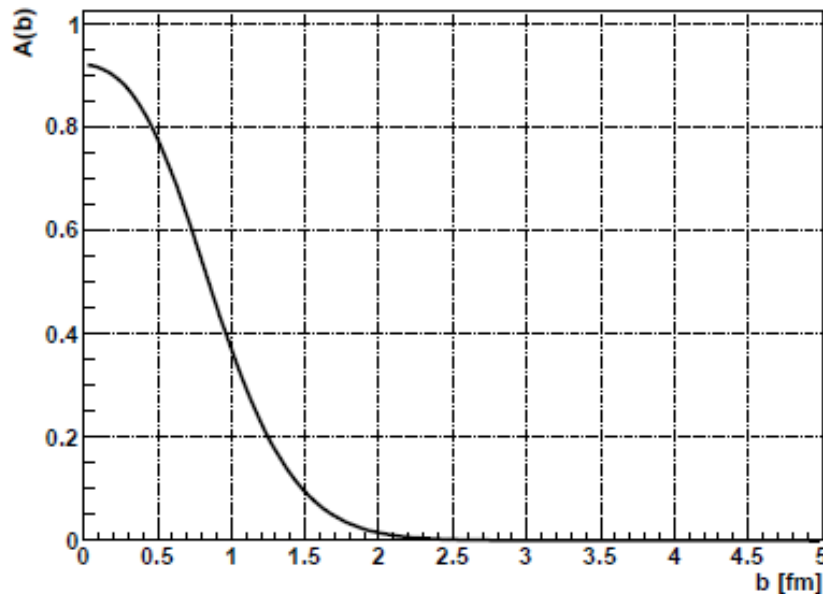
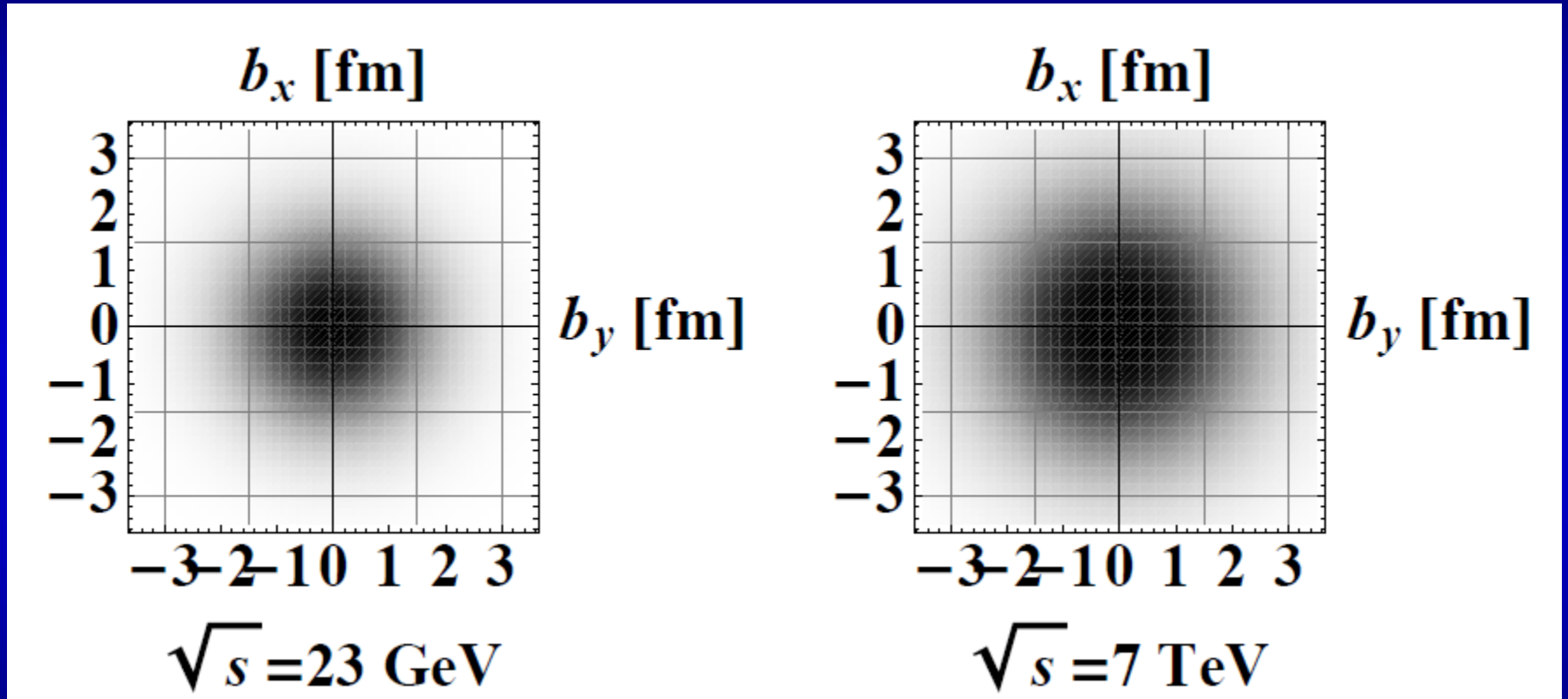


Figure 4: The  $A(b) = 1 - |e^{-\Omega(b)}|^2$  shadow profile function. 23.5 GeV (left) and 7 TeV (right).

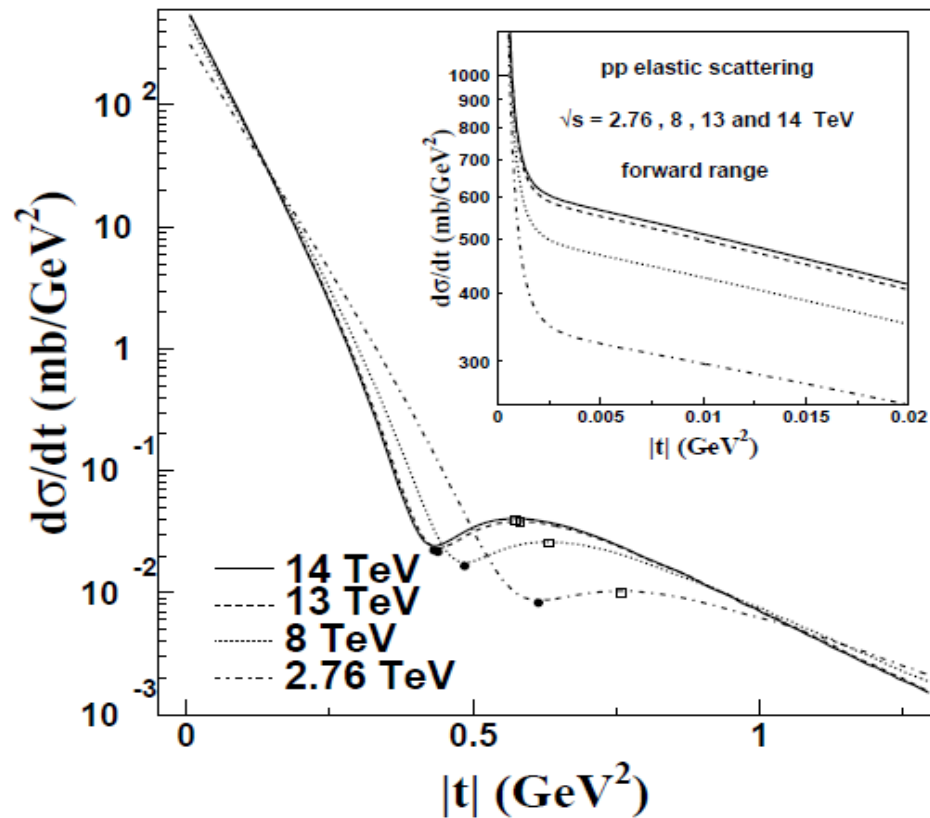
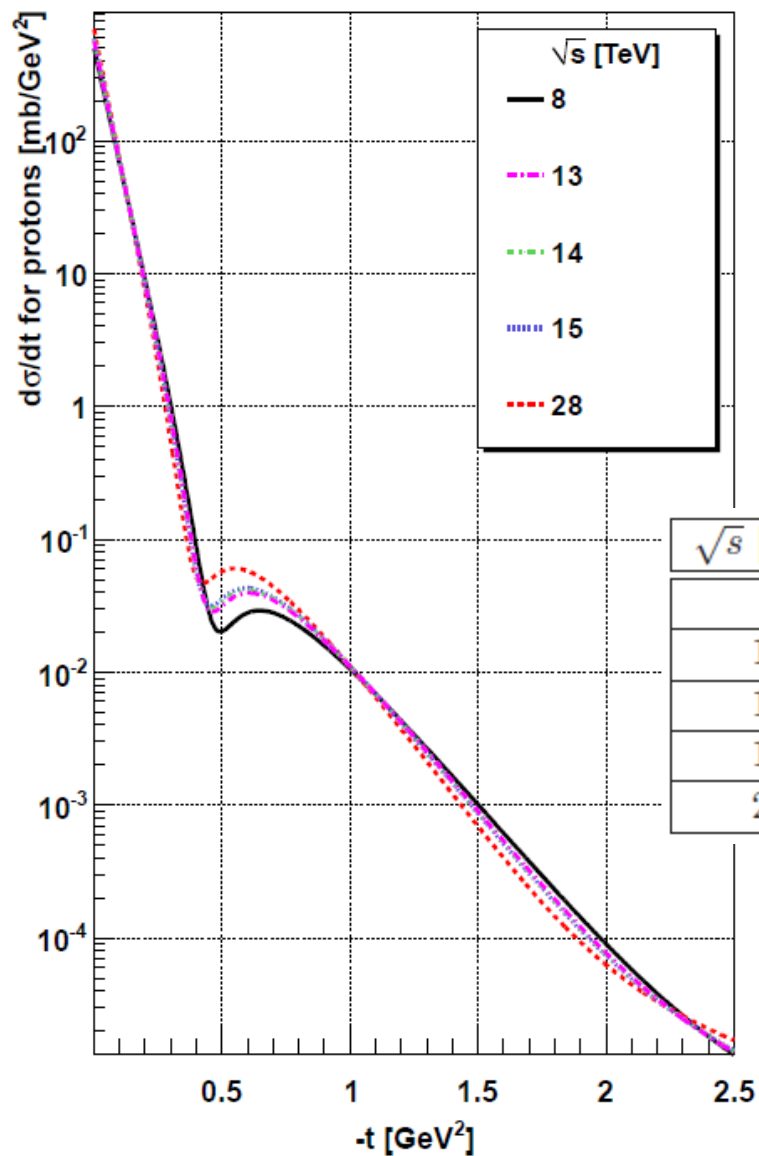
Indication of saturation at 7 TeV:  $A(b) \sim 1$  at low  $b$ .  
 $\sim$  max probability of interaction at low  $b$

# Imaging on the sub-femtometer scale at 23 GeV ISR and 7 TeV LHC energy



What about 8 TeV and future LHC energies?

# Excitation function: $d\sigma/dt$



but  $t_{\text{dip}} \sigma_{\text{tot}} \sim \text{const}$  (2 %)

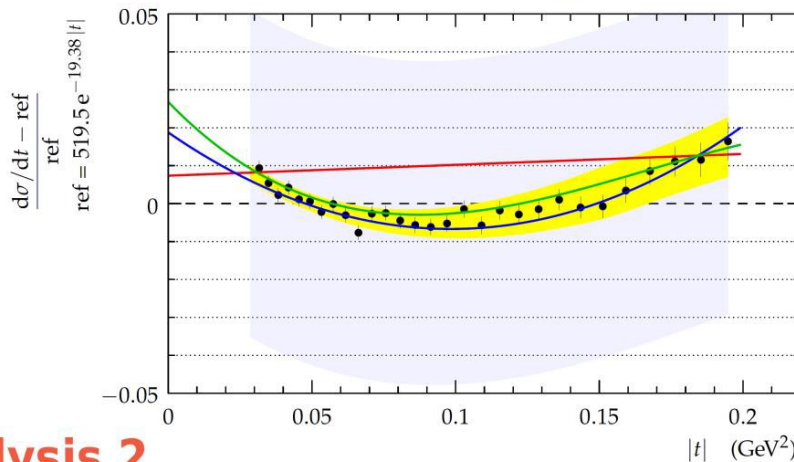
Similar to:

K.A. Kohara, T. Kodama, E. Ferreira,  
 arXiv:1411.3518

# TOTEM 8 TeV pp data

**Analysis 1:** fits  $A \exp(b_1 t + b_2 t^2 + \dots)$ ,  $N_b$  parameters in exponent

DS4



diagonals combined

**Nucl.Phys. B899 (2015) 527-546**

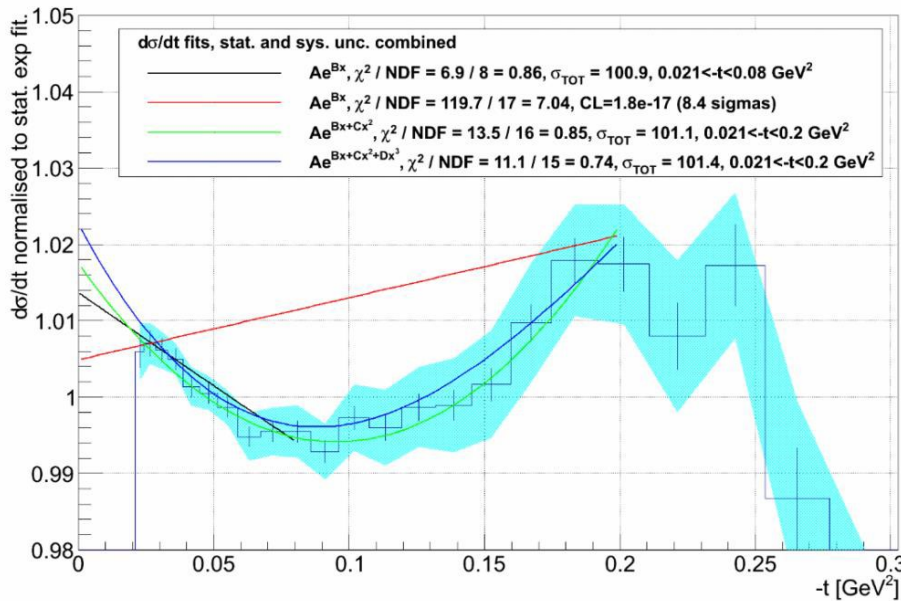
- data (binning)
- | statistical uncertainties
- systematic uncertainty band: analysis+normalisation
- systematic uncertainty band: analysis only

fit parametrisation:  $a \exp(\sum_{n=1}^{N_b} b_n t^n)$

fits with statistical and systematic uncertainties:

- $N_b = 1$ :  $\chi^2/\text{ndf} = 117.5/28 = 4.198 \Rightarrow$  p-value =  $6.14 \times 10^{-13}$ , significance =  $7.20 \sigma$
- $N_b = 2$ :  $\chi^2/\text{ndf} = 29.3/27 = 1.085 \Rightarrow$  p-value =  $3.47 \times 10^{-1}$ , significance =  $0.94 \sigma$
- $N_b = 3$ :  $\chi^2/\text{ndf} = 25.5/26 = 0.980 \Rightarrow$  p-value =  $4.92 \times 10^{-1}$ , significance =  $0.69 \sigma$

**Analysis 2**



dσ/dt fits, stat. and sys. unc. combined

- $Ae^{Bx}$ ,  $\chi^2/\text{NDF} = 6.9/8 = 0.86$ ,  $\sigma_{\text{TOT}} = 100.9$ ,  $0.021 < -t < 0.08 \text{ GeV}^2$
- $Ae^{Bx+Cx^2}$ ,  $\chi^2/\text{NDF} = 119.7/17 = 7.04$ ,  $\text{CL} = 1.8e-17$  (8.4 sigmas)
- $Ae^{Bx+Cx^2+Dx^3}$ ,  $\chi^2/\text{NDF} = 13.5/16 = 0.85$ ,  $\sigma_{\text{TOT}} = 101.1$ ,  $0.021 < -t < 0.2 \text{ GeV}^2$
- $Ae^{Bx+Cx^2+Dx^3}$ ,  $\chi^2/\text{NDF} = 11.1/15 = 0.74$ ,  $\sigma_{\text{TOT}} = 101.4$ ,  $0.021 < -t < 0.2 \text{ GeV}^2$

TOTEM pp data at 8 TeV  
exponential shape  
excluded at  $7 + \sigma$

*new determination*

$$\sigma_{\text{tot}} = (101.4 \pm 2.0) \text{ mb}$$

Theoretical support, from ISR  
to LHC energies, unitarity  
L. Jenkovszky and A. Lengyel,  
arXiv:1410.4106

# Non-exponential behaviour in ReBB

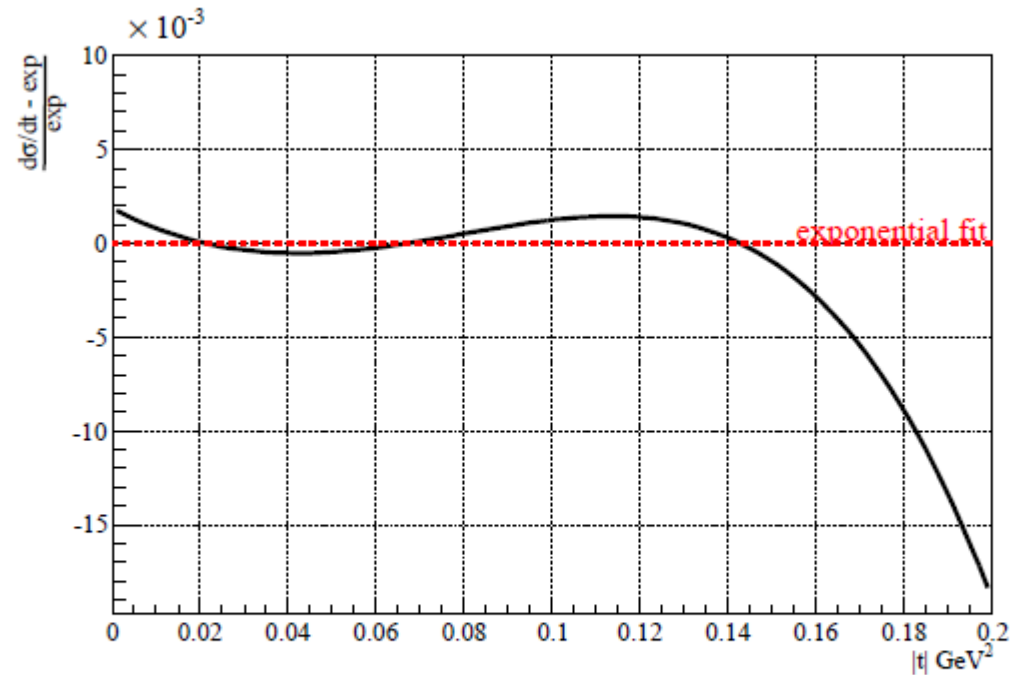
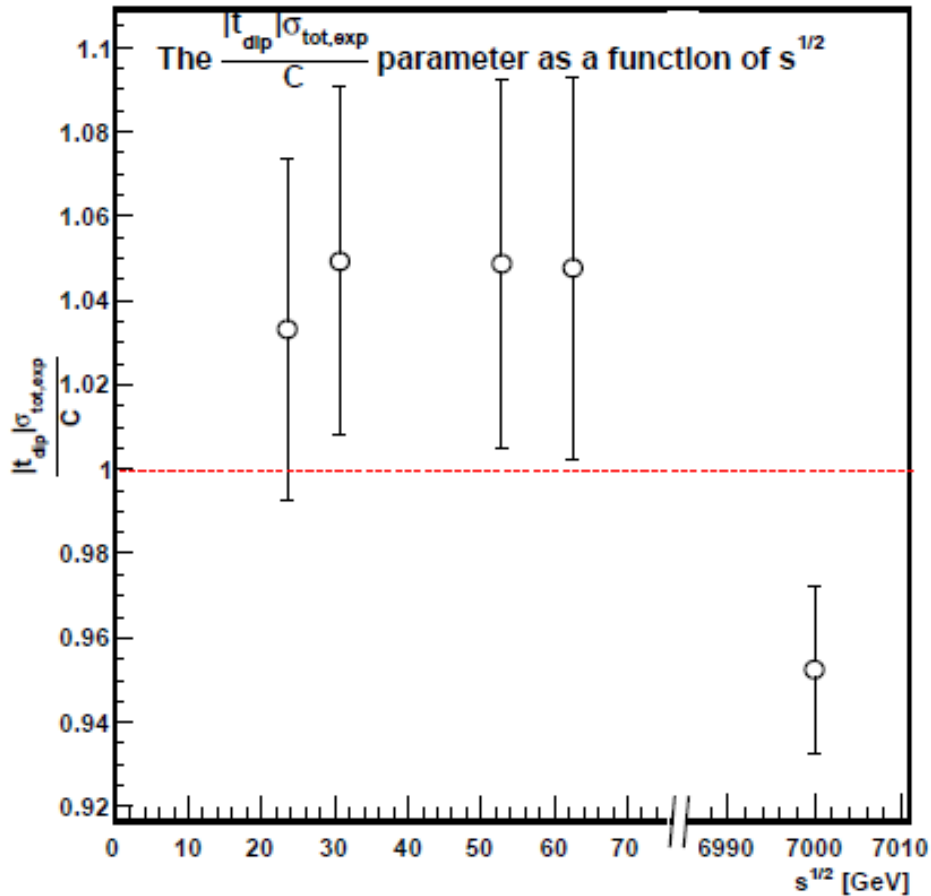


Fig. 5. The ReBB model, fitted in the  $0.0 \leq |t| \leq 0.36 \text{ GeV}^2$  range, with respect to the exponential fit of Eq. (33). In the plot only the  $0.0 \leq |t| \leq 0.2 \text{ GeV}^2$  range is shown. The curve indicates a significant deviation from the simple exponential at low  $|t|$  values.

Similar  
non-exponential feature  
seen at 7 TeV as in 8 TeV  
TOTEM data

# Black disc limit?



Geometric scaling,  
but not the black disc limit:

T. Cs. and F. Nemes  
arXiv:1306.4217  
Int. J. Mod. Phys. A (2014)

$C(\text{data}) \sim 50 \text{ mb GeV}^2$   
 $\neq$   
 $C(\text{black}) \sim 36 \text{ mb GeV}^2$

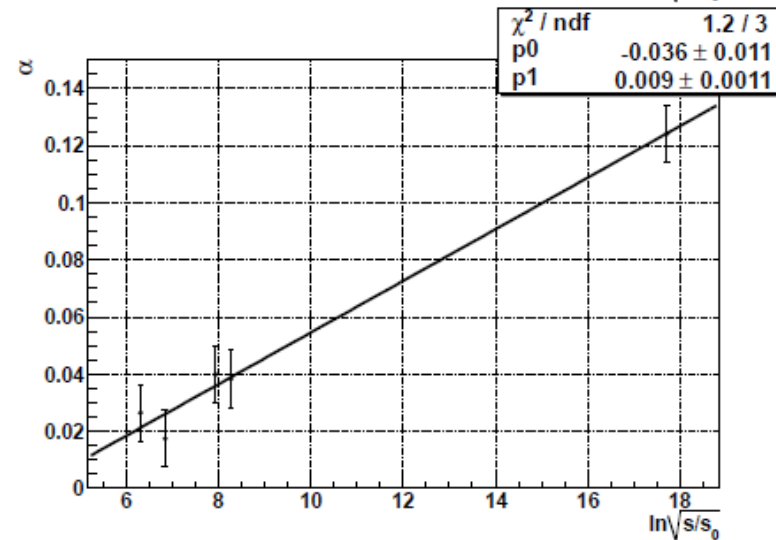
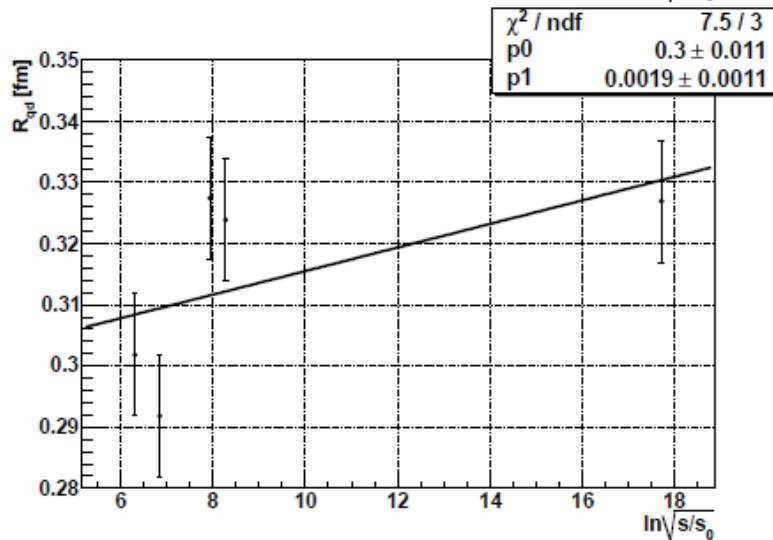
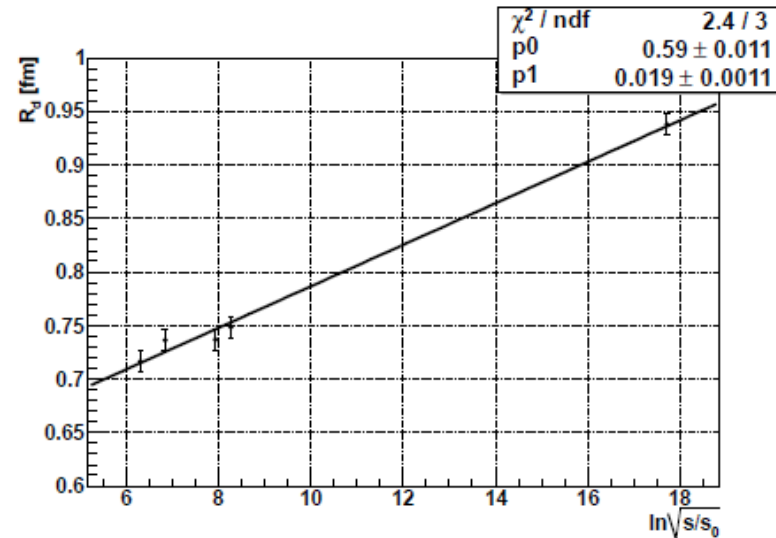
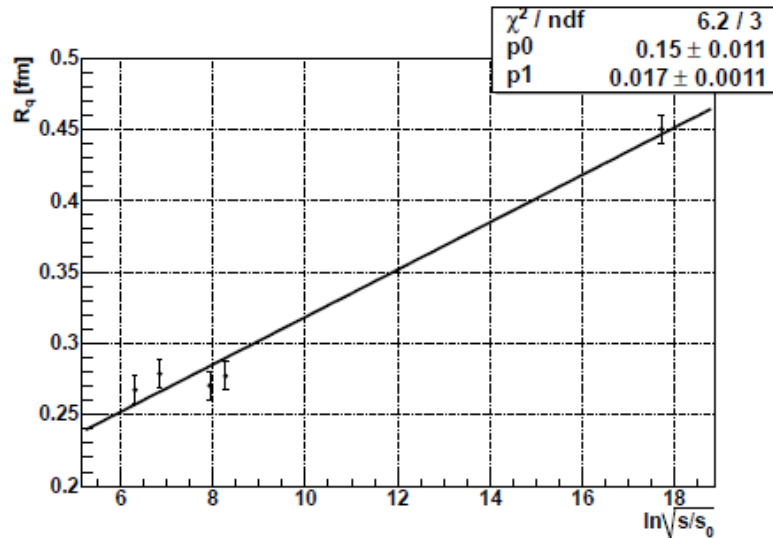
$$\frac{d\sigma_{black}}{dt} = \pi R^4 \left[ \frac{J_1(qR)}{qR} \right]^2$$

$$\sigma_{tot,black} = 2\pi R^2.$$

$$C_{black} = |t_{dip,black}| \cdot \sigma_{tot,black} = 2\pi j_{1,1}^2 (\hbar c)^2 \approx 35.9 \text{ mb GeV}^2$$

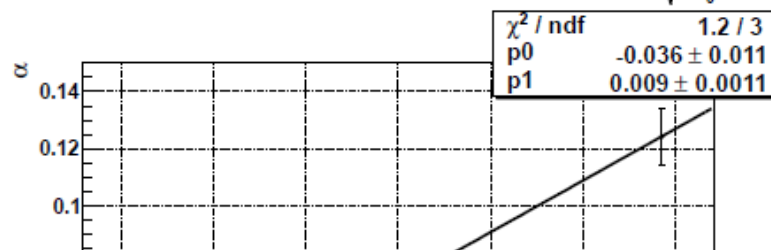
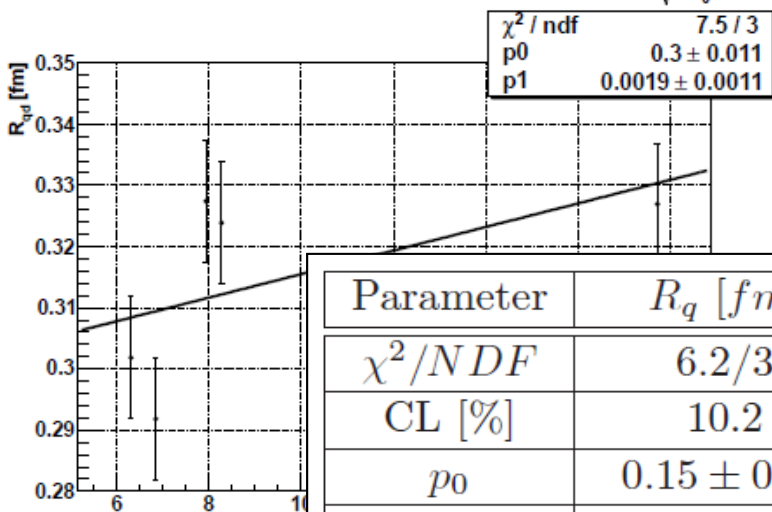
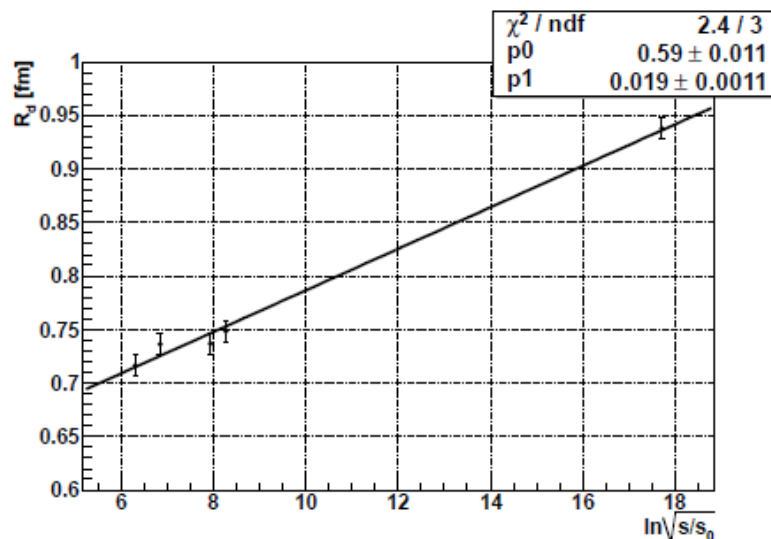
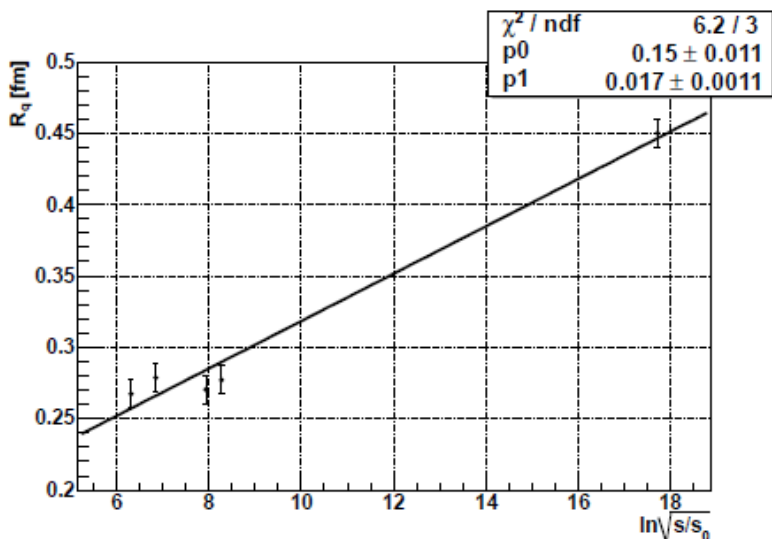


# Excitation function: scaling in pp



Geometric scaling:  $\{R_q, R_d, R_{qd}, \alpha\} = p_0 + p_1 \ln (s/s_0)$

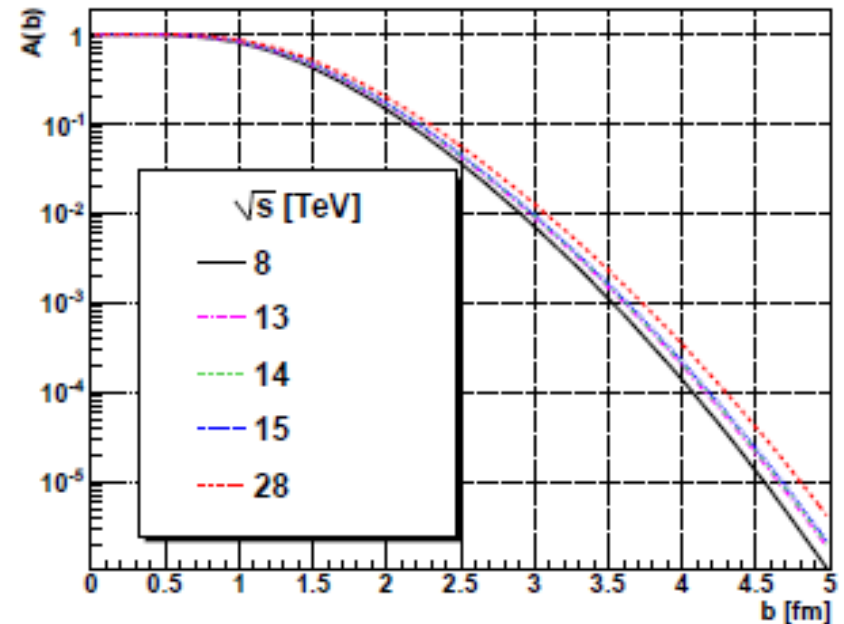
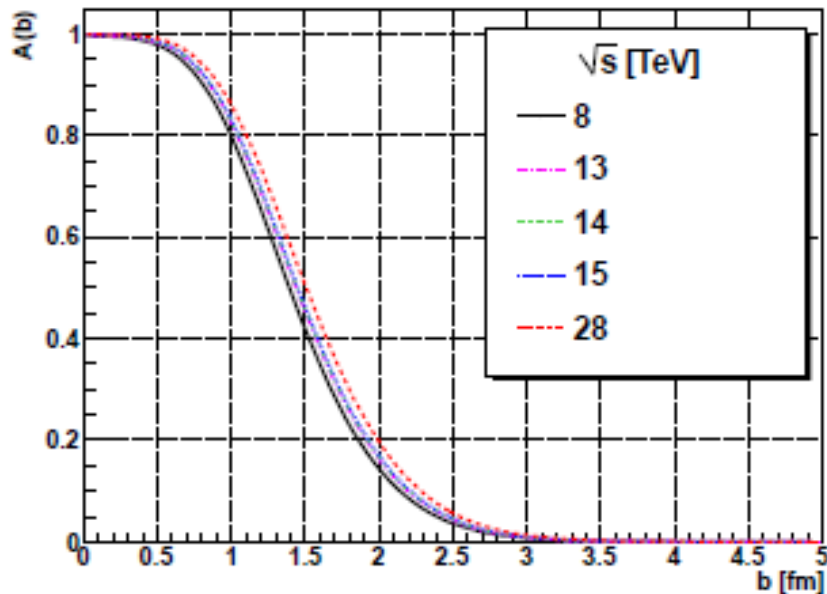
# Geometric scaling in pp



Parameter	$R_q$ [fm]	$R_d$ [fm]	$R_{qd}$ [fm]	$\alpha$
$\chi^2/NDF$	6.2/3	2.4/3	7.5/3	1.2/3
CL [%]	10.2	49.4	5.8	75.3
$p_0$	$0.15 \pm 0.01$	$0.59 \pm 0.01$	$0.3 \pm 0.01$	$-0.036 \pm 0.01$
$p_1$	$0.017 \pm 0.001$	$0.019 \pm 0.001$	$0.0019 \pm 0.001$	$0.009 \pm 0.001$

Geometric scaling:  $\{R_q, R_d, R_{qd}, \alpha\} = p_0 + p_1 \ln (s/s_0)$

# Predictions for the shadow profile

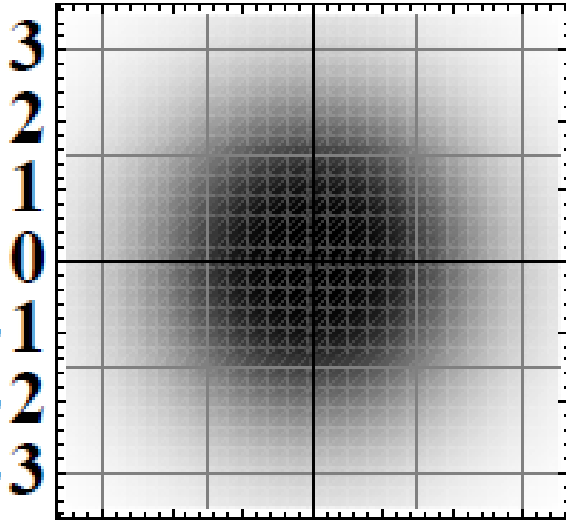


Blacker and Larger,  
but not Edgier:  
BnEL effect  
at LHC energies

Similar to:  
K.A. Kohara, T. Kodama,  
E. Ferreira,  
arXiv:1411.3518  
but they also claim  
an asymptotic BEL effect

# Predictions for the shadow profile

$b_x$  [fm]

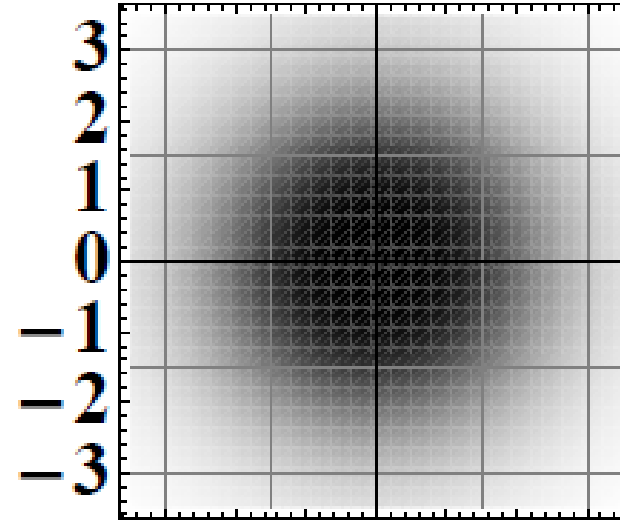


$b_y$  [fm]

-3 -2 -1 0 1 2 3

$\sqrt{s} = 14$  TeV

$b_x$  [fm]



$b_y$  [fm]

-3 -2 -1 0 1 2 3

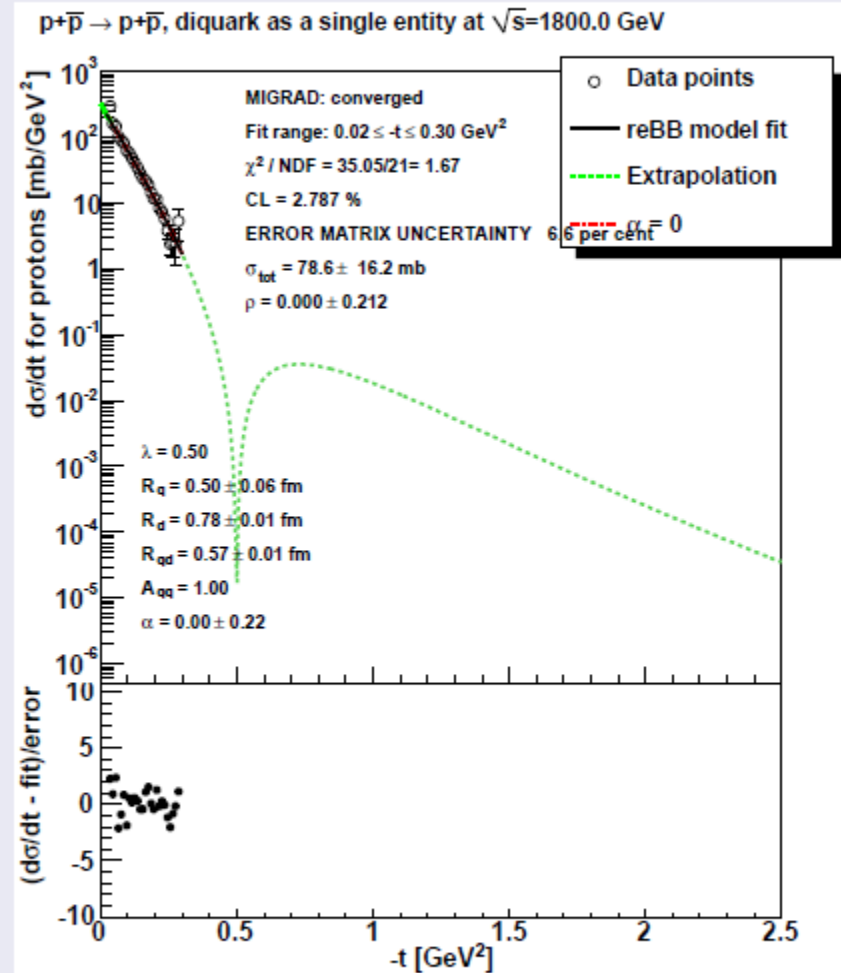
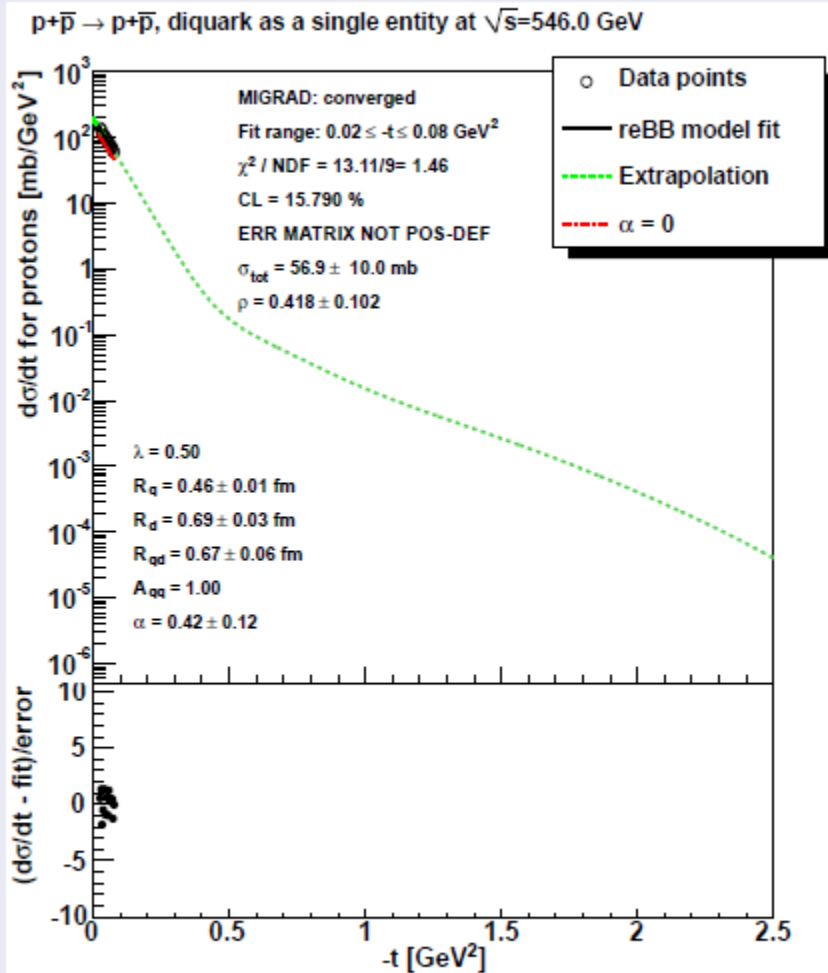
$\sqrt{s} = 28$  TeV

Blacker and Larger,  
but not Edgier:  
BnEL effect  
at LHC energies

Results presented so far:  
[arxiv:1505.01415](https://arxiv.org/abs/1505.01415)

# New results: $p\bar{p}$ data with ReBB model

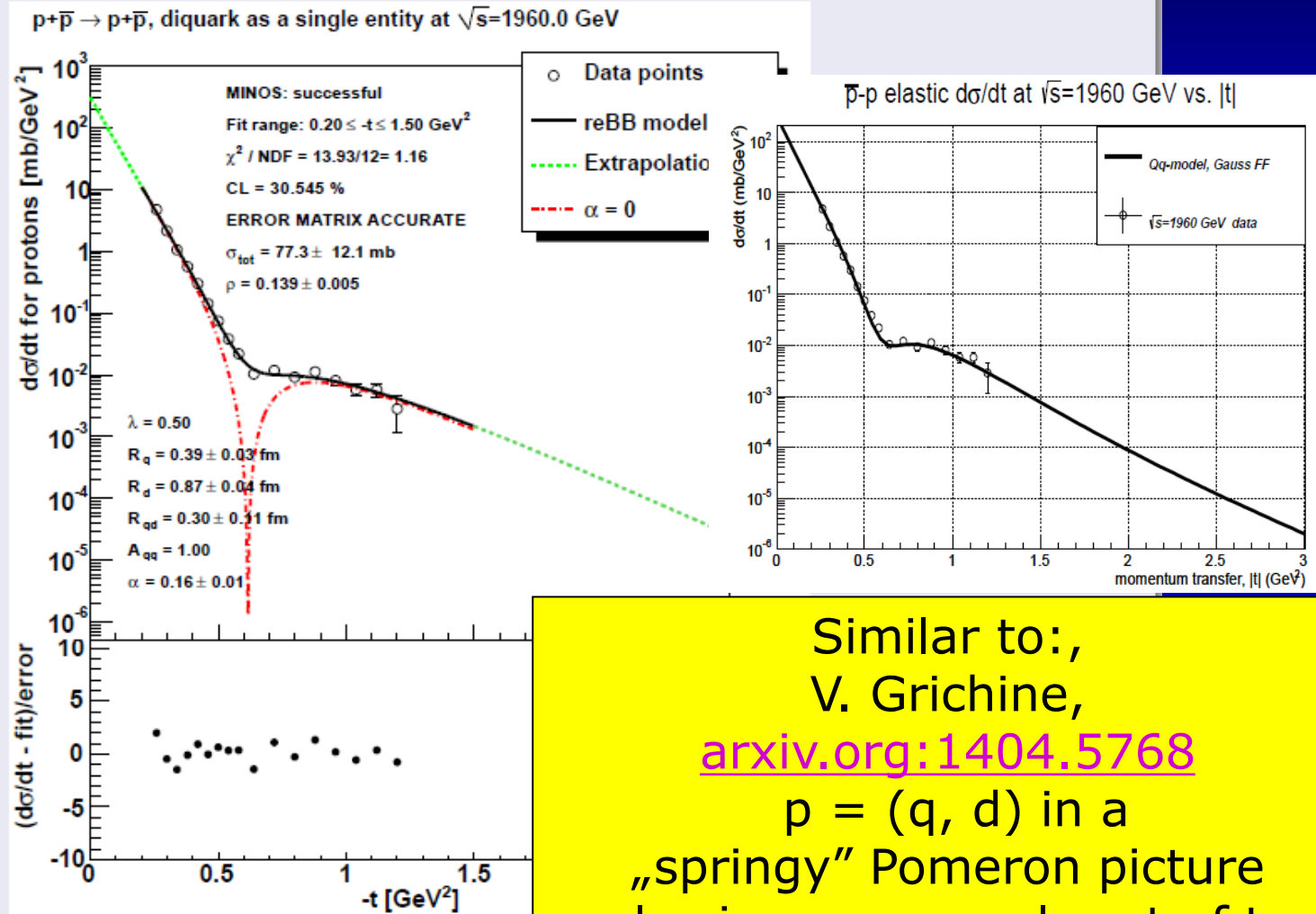
All the usual BB fit parameters are free ( $\sqrt{s} = 546 \text{ GeV}, 1.8 \text{ TeV}$ )



There are not enough points to pin down the shape.

# Tevatron $p\bar{p}$ data with ReBB model

All the usual BB fit parameters are free (1.96 TeV)

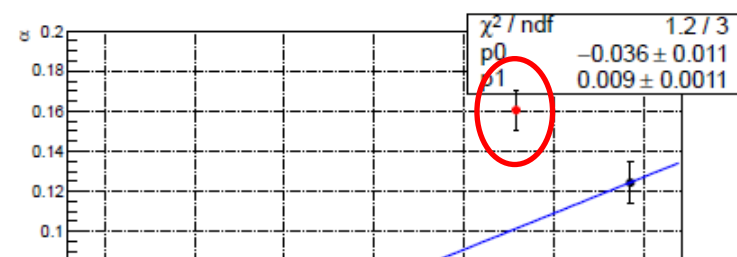
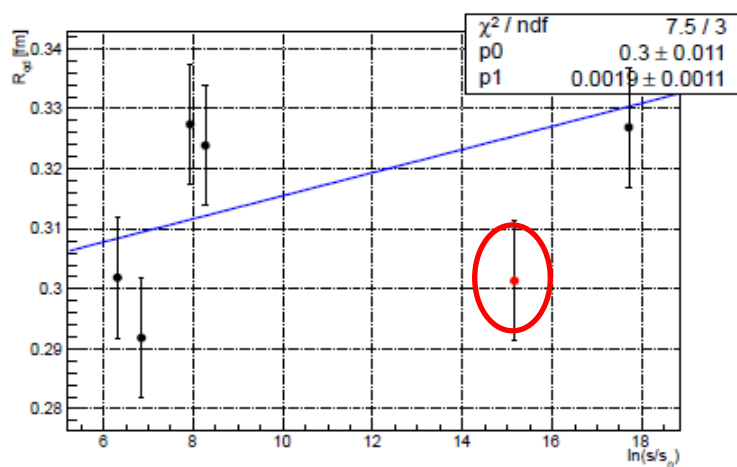
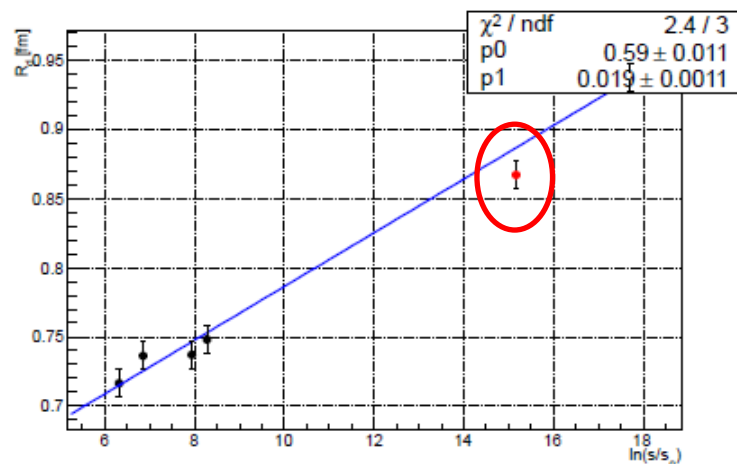
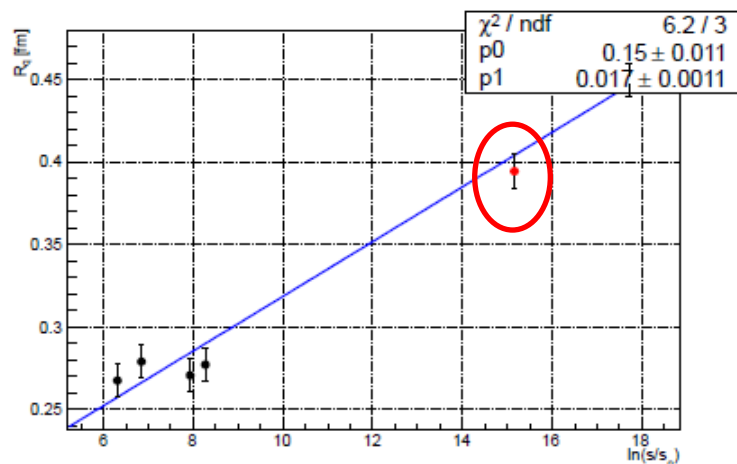


Similar to,  
 V. Grichine,  
[arxiv.org:1404.5768](https://arxiv.org/abs/1404.5768)  
 $p = (q, d)$  in a  
 „springy” Pomeron picture  
 also increases real part of  $t_{el}$

Ok.

# Tevatron $p\bar{p}$ data trends ReBB model

The good fit at  $\sqrt{s} = 1.96$  TeV compared with the extrapolations based only on  $pp$  fits of our ReBB paper



ReBB model works also for elastic  $p\bar{p}$  data but  $p\bar{p}$  is more „opaque” than  $pp$ .

# What have we learned?

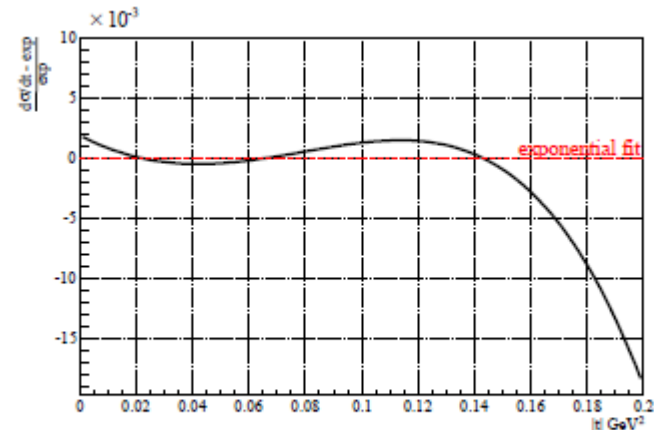
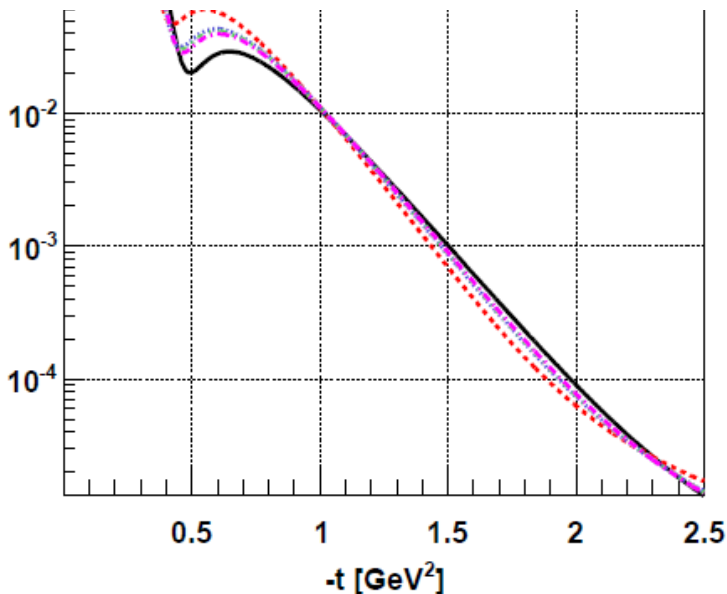
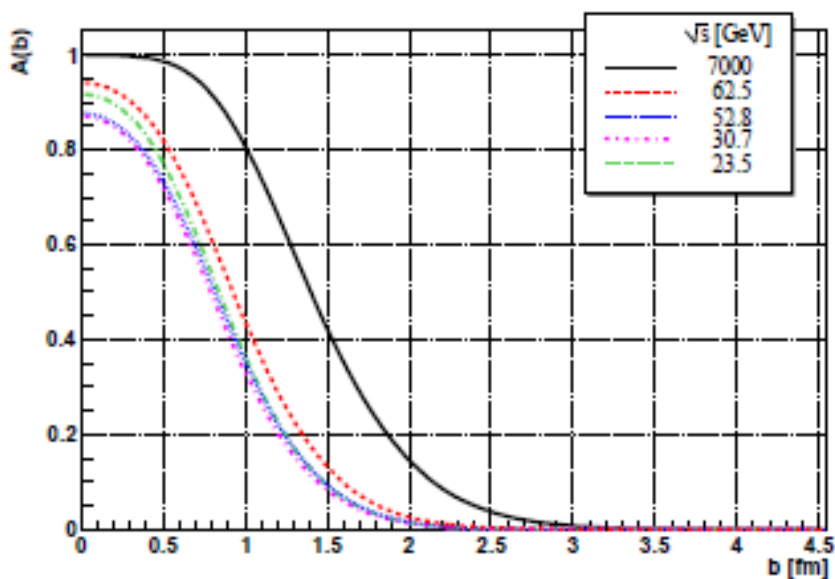


FIG. 5. The  $rBB$  model, fitted in the  $0.0 \leq |t| \leq 0.36$  GeV<sup>2</sup> range, with respect to the exponential

Really extended Bialas-Bzdak:  
 $p = (q, d)$  at LHC  
 non-trivial structure at low- $t$   
 describes both  $pp$  and  $p\bar{p}$

BnEL: Blacker, **Edgier**, Larger

ReBB model works naturally  
 also for elastic  $p\bar{p}$  data  
 but antiproton is more „opaque“.



# Backup slides – Questions?

# Focusing reBB on the low- $t$ region

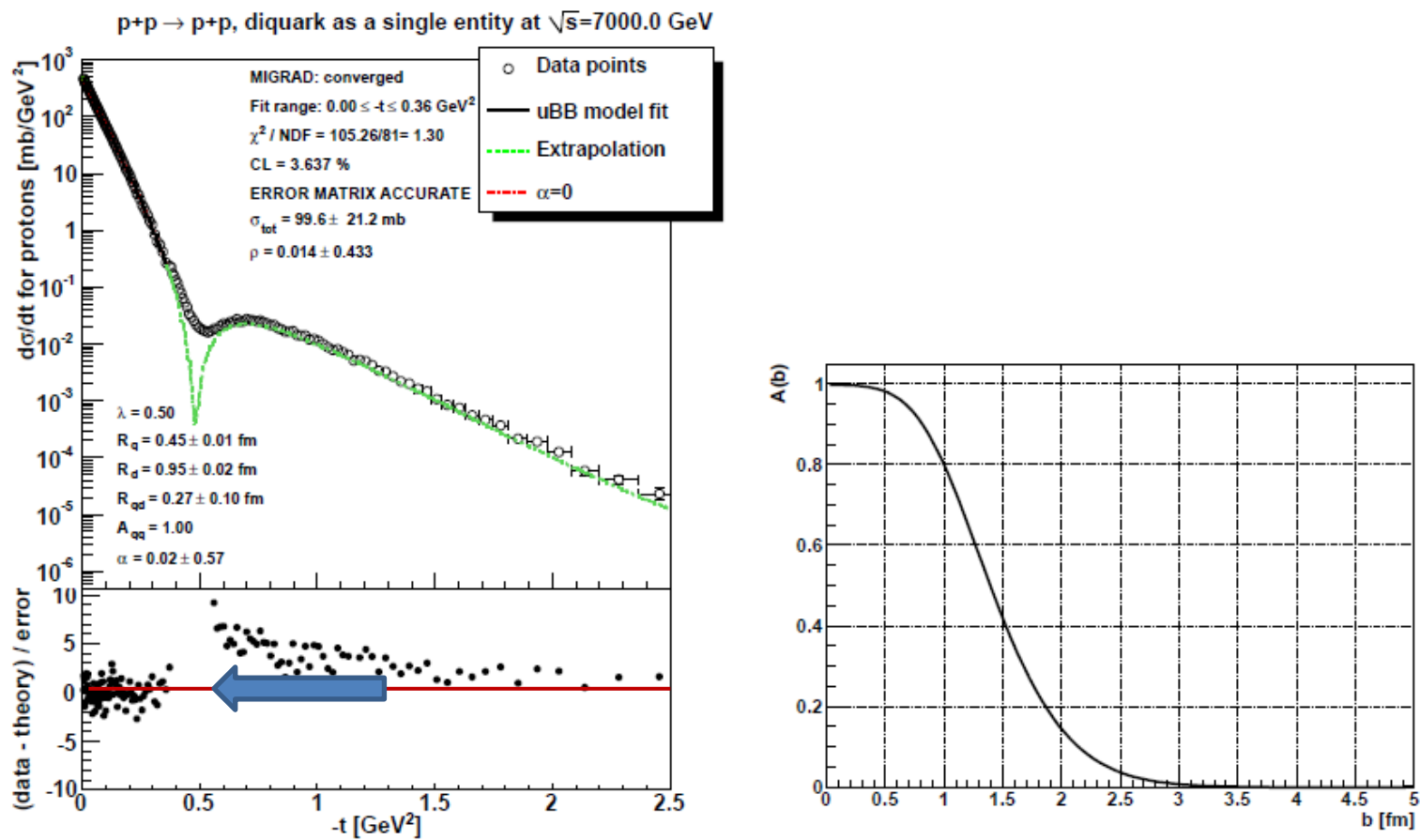


Figure 5:  $0 - 0.36$  GeV<sup>2</sup>.  $\rho$  has large error, since the dip is not part of the fit.

Saturation is apparent if fit range is limited to  $|t| < 0.36$  GeV<sup>2</sup>

# Focusing reBB on even lower $-t$ region

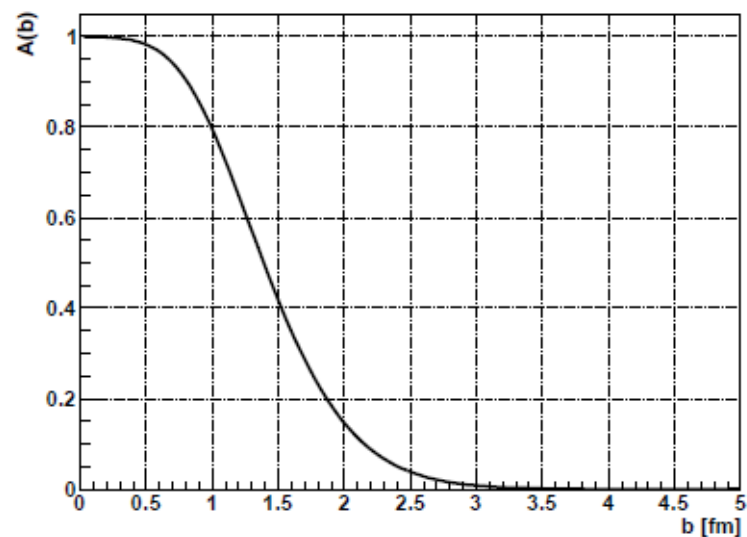
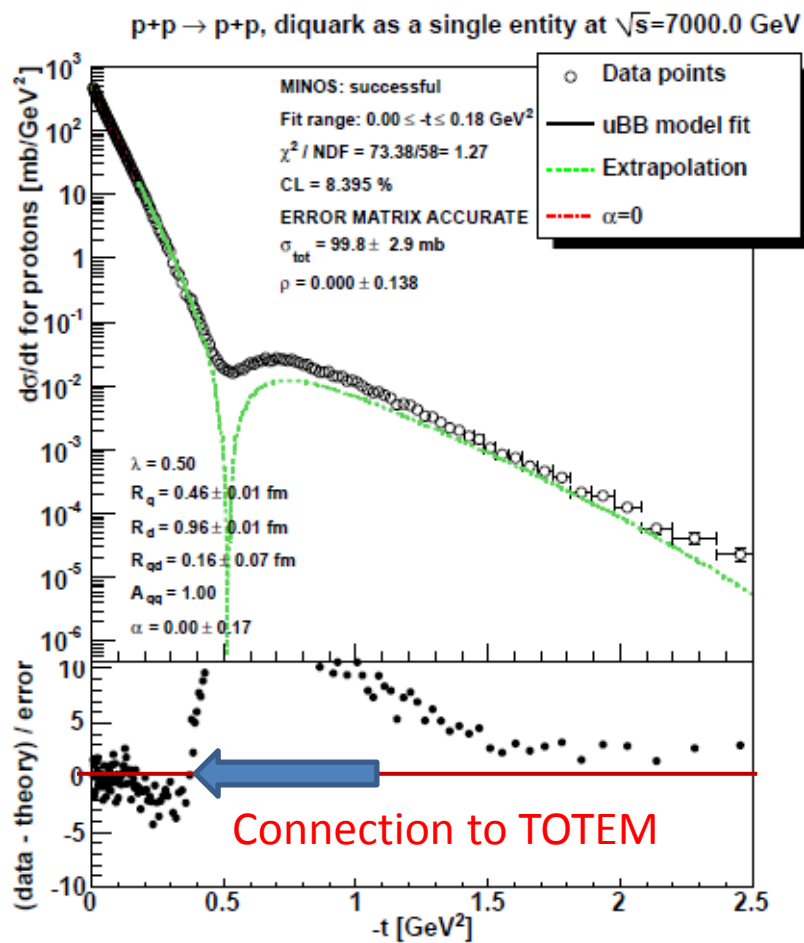
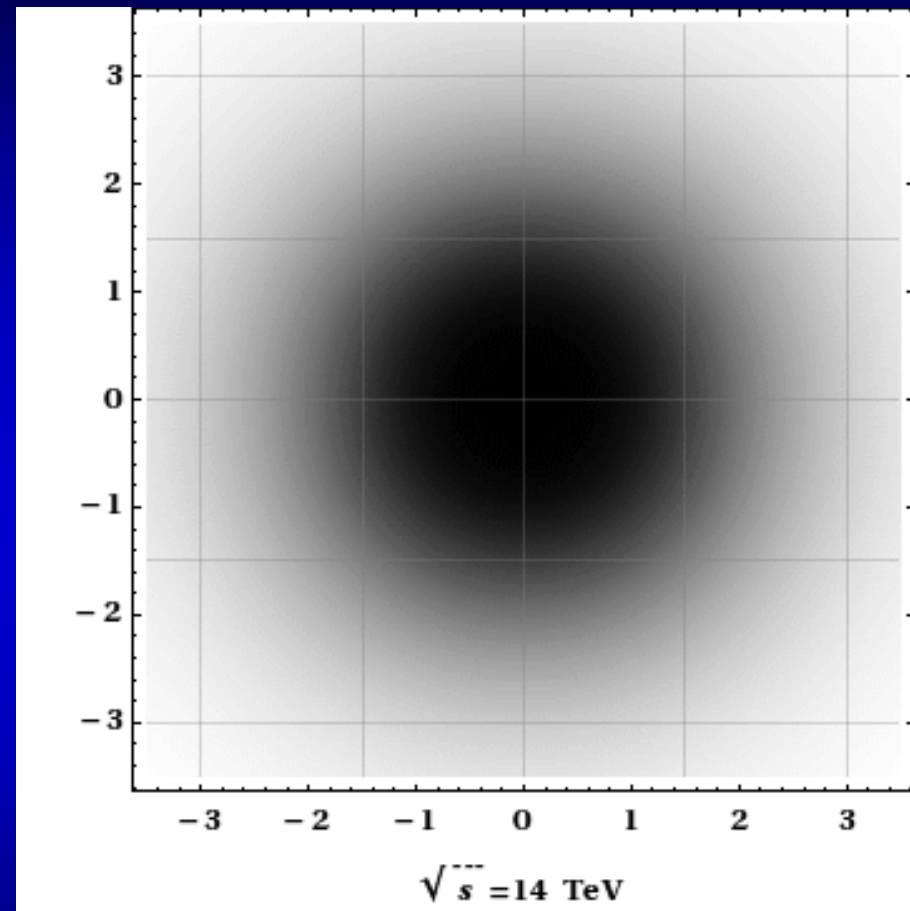
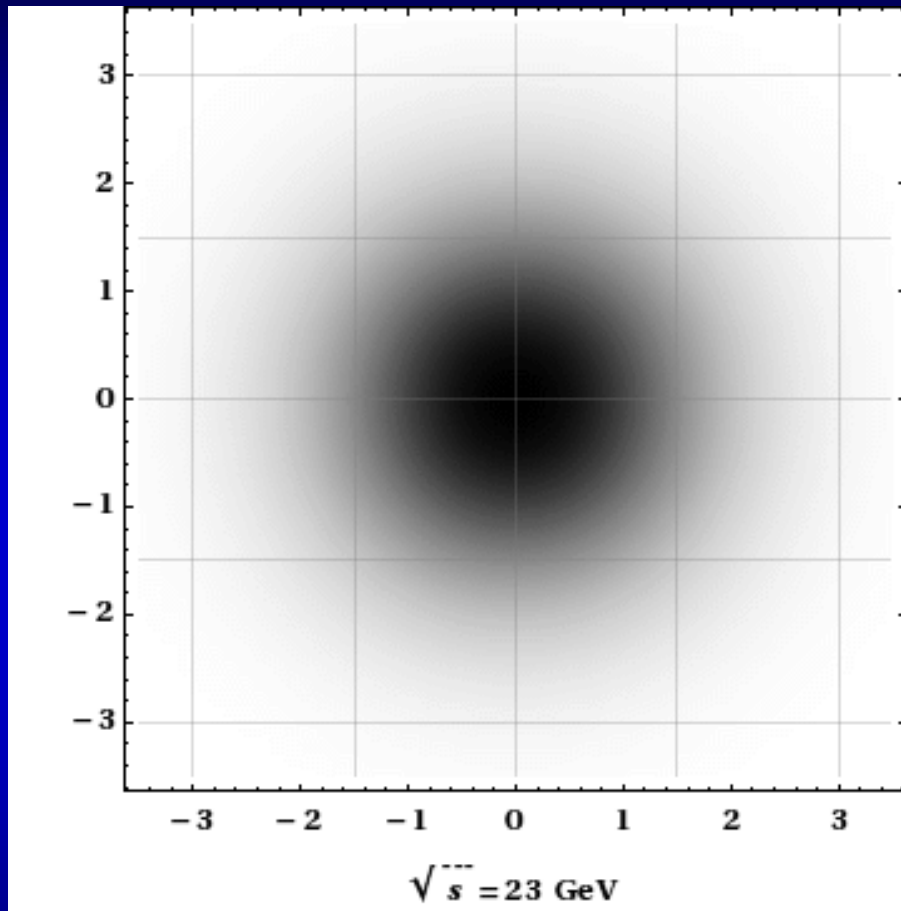


Figure 6:  $0 - 0.18$  GeV<sup>2</sup>.  $\rho$  has large error, since the dip is not part of the fit.

Saturation still apparent, fit range  $|t| < 0.18$  GeV<sup>2</sup>

# Backup slides – Discussion

# Motivation: Is the proton a black disc?

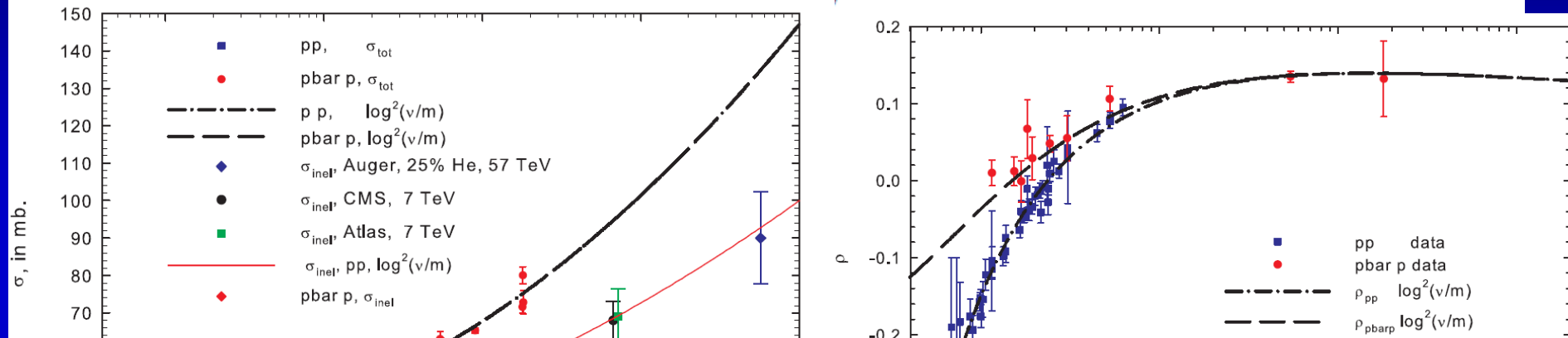


Recent papers by M. Block and F. Halsen address this topic :  
Experimental confirmation: the proton is asymptotically a black disc,  
arXiv:1109.2041, Phys. Rev. Lett. 107 (2011) 212002

# Properties of a black disc

*Properties of a black disk:* In impact parameter space  $b$ , the elastic and total cross sections are given by

$$\sigma_{\text{el}} = 4 \int d^2b |a(b, s)|^2, \quad \sigma_{\text{tot}} = 4 \int d^2b \text{Im} a(b, s). \quad (7)$$



*Conclusions:* We find that the  $\ln^2 s$  Froissart bound for the proton for  $\sigma_{\text{tot}}$  [7] and  $\sigma_{\text{inel}}$  [9] is saturated and that at infinite  $s$ , (1) the experimental ratio  $\sigma_{\text{inel}}/\sigma_{\text{tot}} = 0.509 \pm 0.011$ , compatible with the black disk ratio of 0.5 and (2) the forward scattering amplitude is purely imaginary. We thus conclude that the proton becomes an expanding black disk at sufficiently ultra-high energies that are probably never accessible to experiment. The theory for these bounds is predicated on the pillar stones of analyticity and unitarity, which have now been experimentally verified up to 57000 GeV. Further, since  $\sigma_{\text{tot}}$  has been extrapolated up from the Tevatron, we expect no new large cross section physics between 2000 and 57000 GeV.

Finally, the lowest-lying glueball mass is measured to be  $M_{\text{glueball}} = 2.97 \pm 0.03$  GeV. Reproducing these experimental results will be a task of lattice QCD.

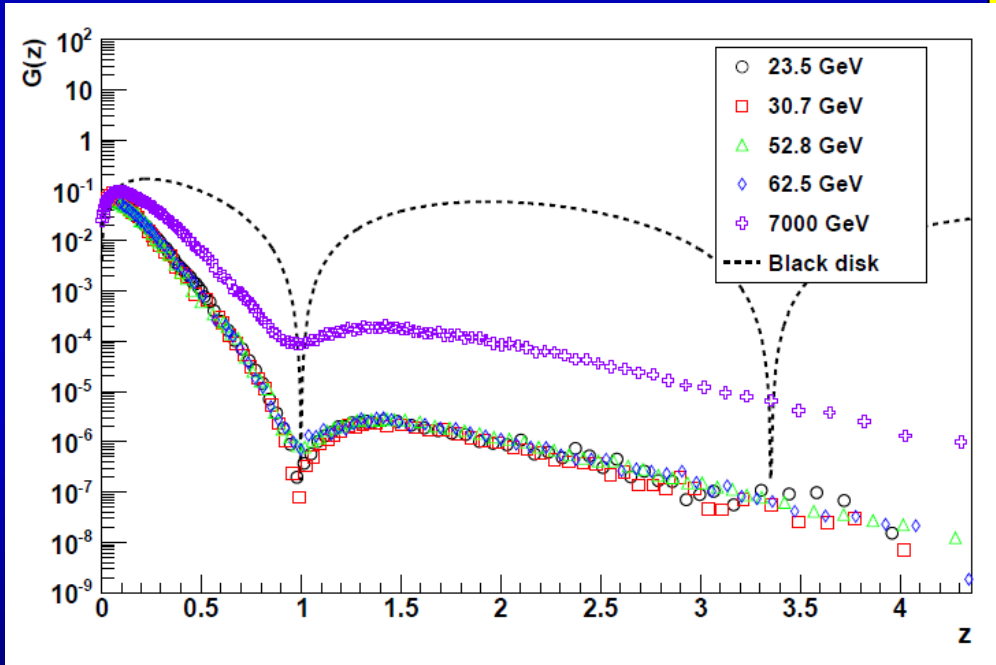
arXiv:1109.2041, Phys. Rev. Lett. 107 (2011) 212002

arXiv:1208.4086, Phys.Rev. D86 (2012) 051504

[arXiv:1409.3196](https://arxiv.org/abs/1409.3196)

# Black Disc (BD) limit?

Scaling but  
not in the black disc limit:  
T. Cs. and F. Nemes  
arXiv:1306.4217  
Int. J. Mod. Phys. A (2014)



$$C(\text{data}) \sim 50 \text{ mb GeV}^2$$

$$\neq$$

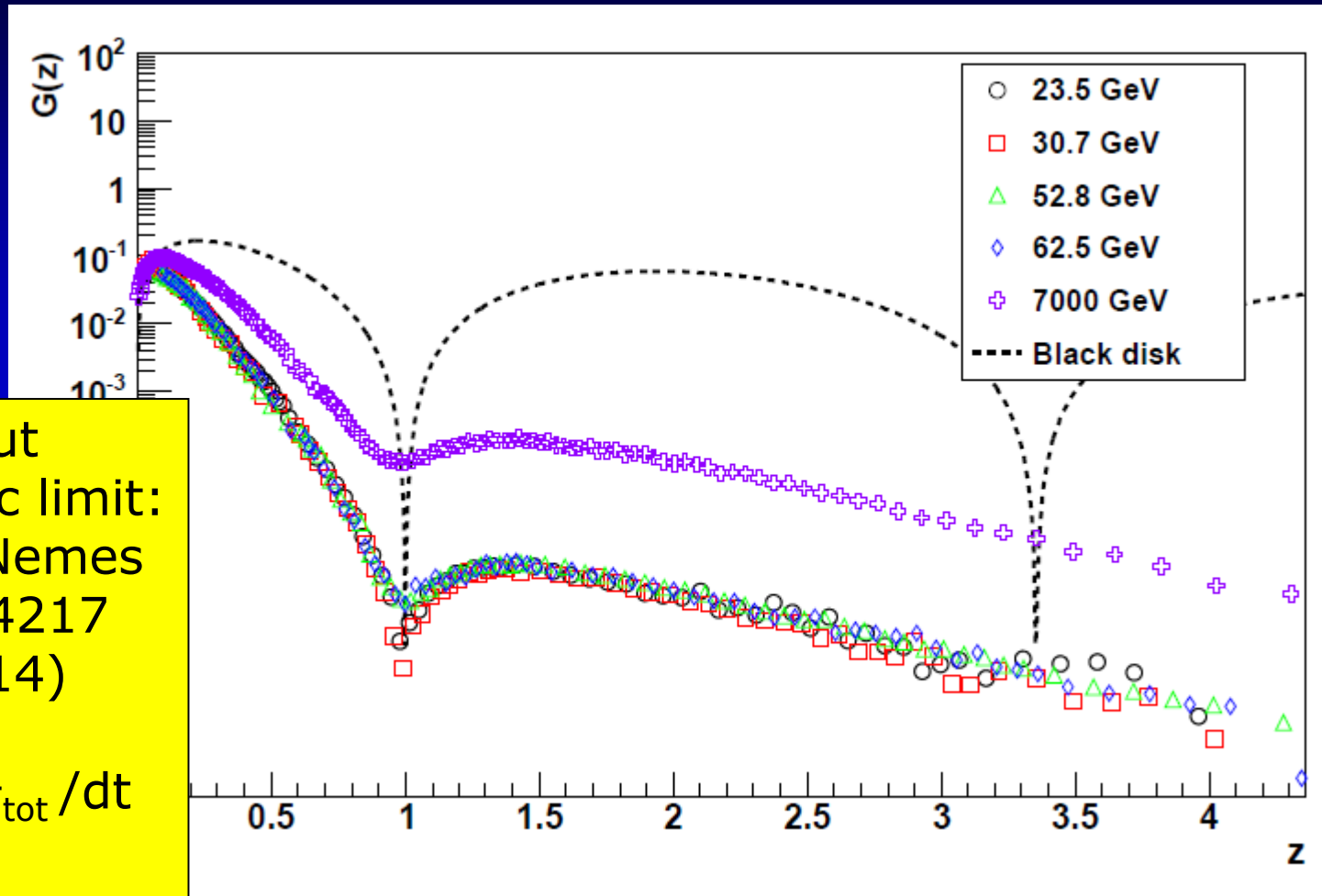
$$C(\text{black}) \sim 36 \text{ mb GeV}^2$$

$$\frac{d\sigma_{\text{black}}}{dt} = \pi R^4 \left[ \frac{J_1(qR)}{qR} \right]^2$$

$$\sigma_{\text{tot,black}} = 2\pi R^2.$$

$$C_{\text{black}} = |t_{\text{dip,black}}| \cdot \sigma_{\text{tot,black}} = 2\pi j_{1,1}^2 (\hbar c)^2 \approx 35.9 \text{ mb GeV}^2$$

# Geometric scaling, but not BD limit?



Scaling but  
not a black disc limit:  
T. Cs. and F. Nemes  
arXiv:1306.4217  
IJMPA (2014)

$$G(z) = t \frac{d\sigma/\sigma_{\text{tot}}}{dt}$$

plotted vs

$$z = t/t_{\text{dip}}$$