## The Unhiggs

#### hep-ph/0708.0005, 0804.0424, 0807.3961

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AdS/CFT/unparticle correspondence

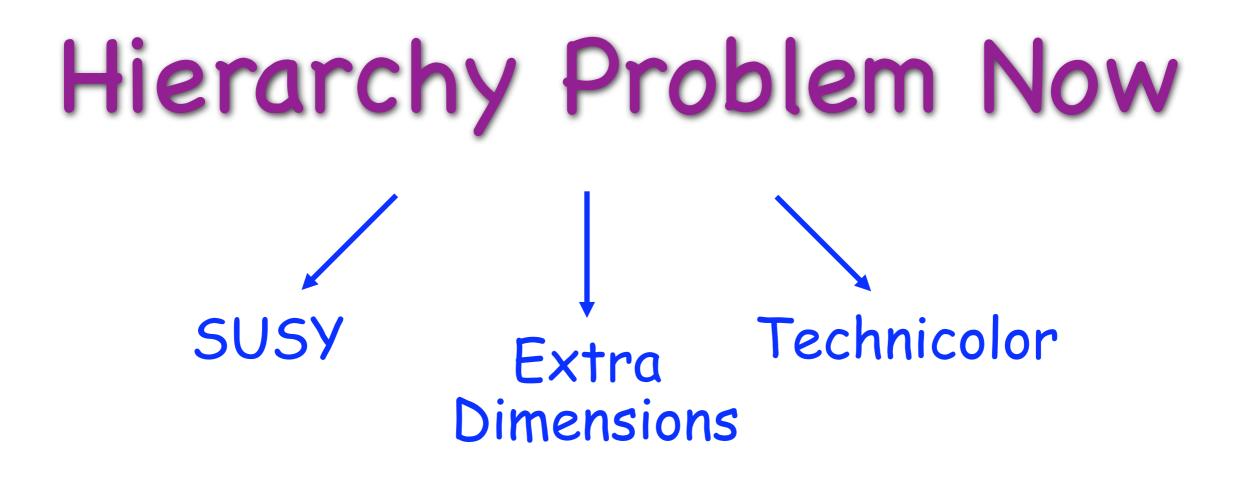


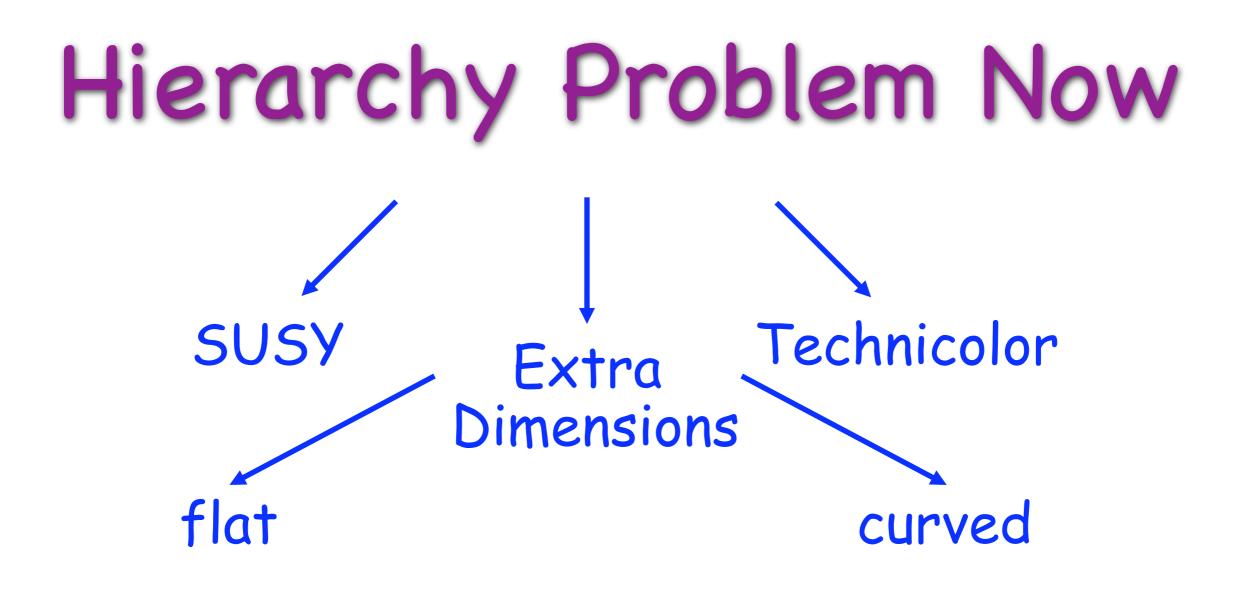


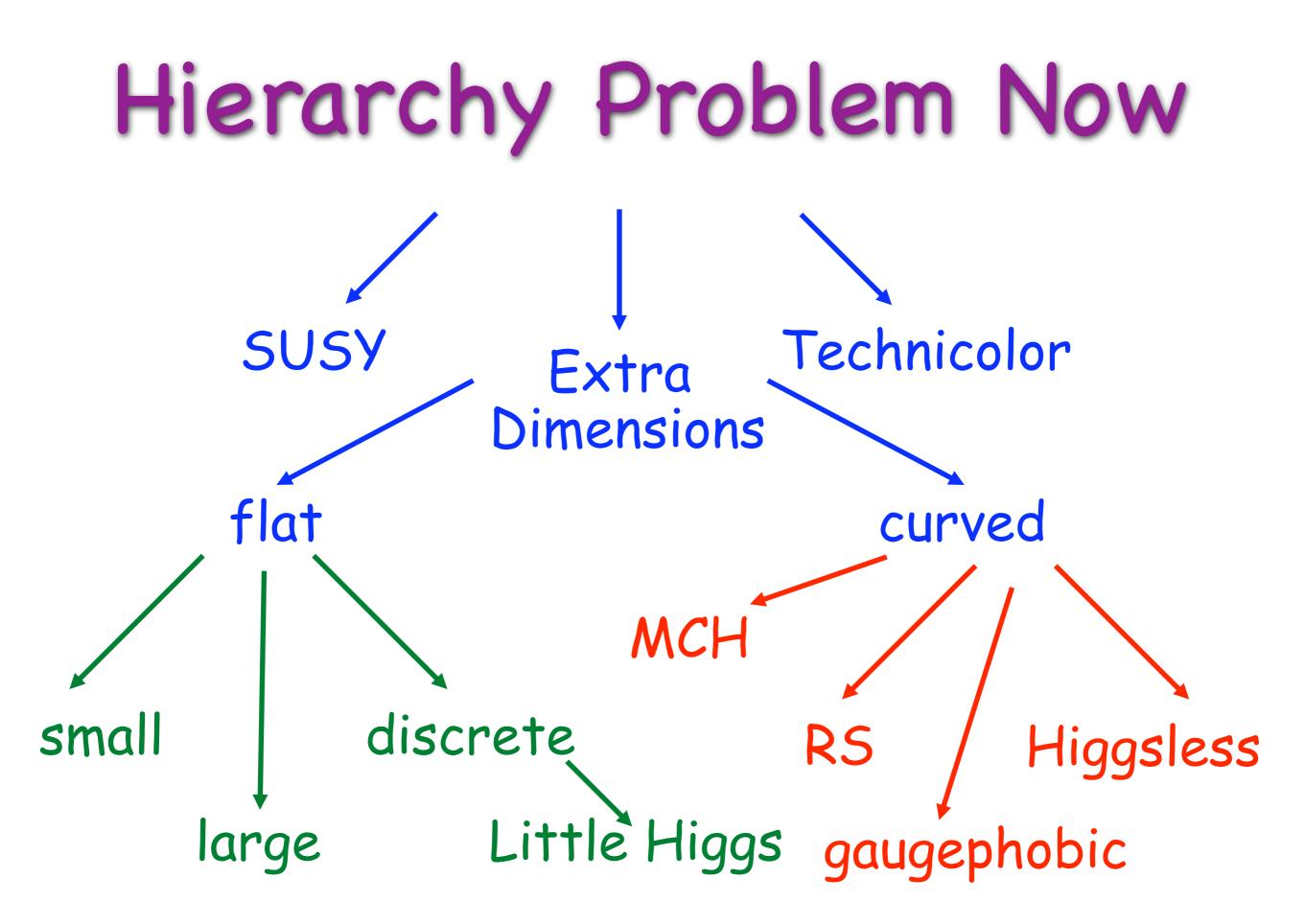
# Hierarchy Problem Now

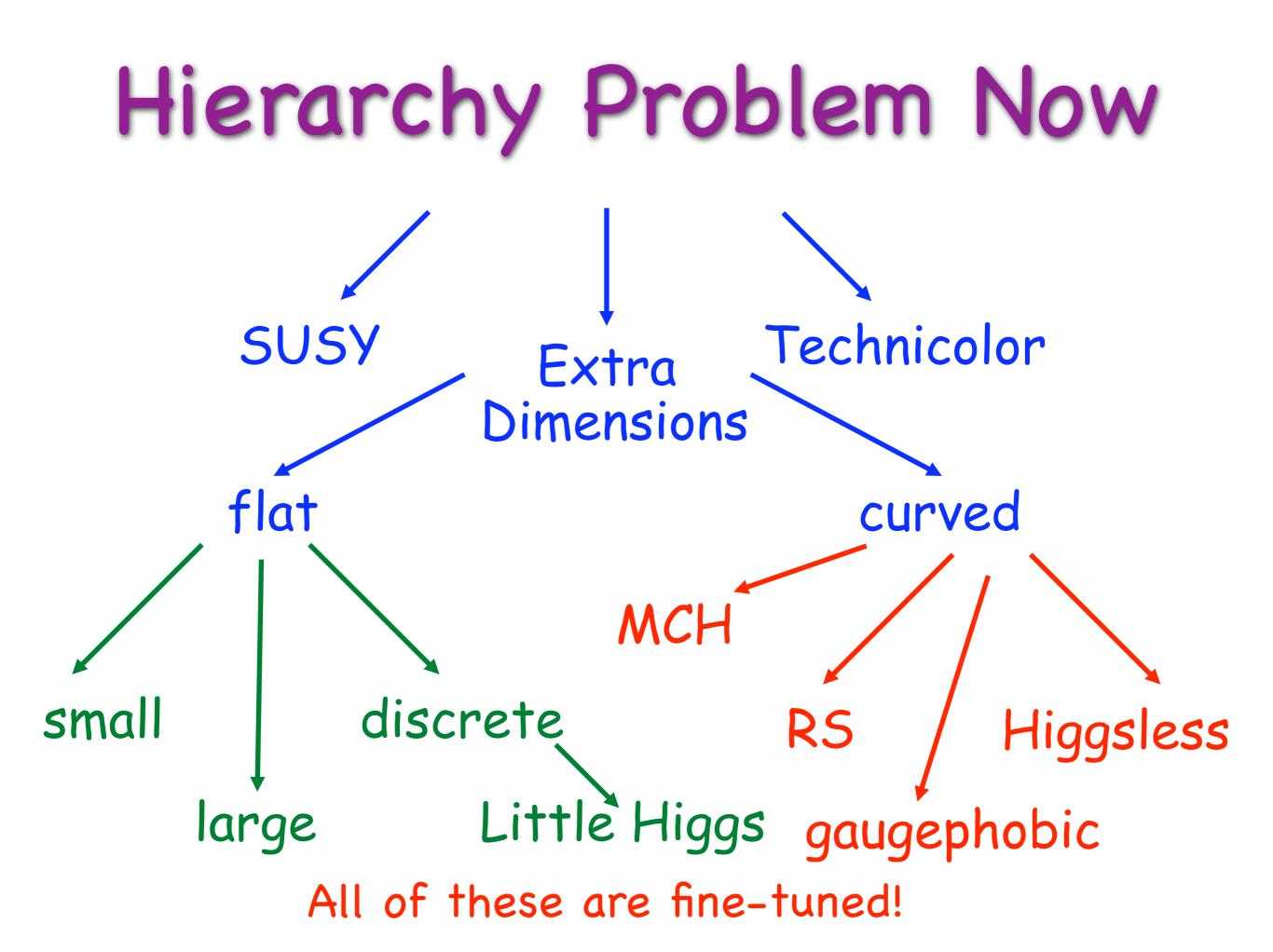


Technicolor









# Quirk/Hidden Valley/ Unparticle Model

X is a heavy fermion with both SM and New gauge couplings stringy confinement  $\rightarrow$  quirks n=0 QCD-like confinement  $\rightarrow$  hidden valley n=few CFT, no confinement  $\rightarrow$  unparticles n=many





a different way to calculate in CFT's

of particles

Georgi hep-ph/0703260, 0704.2457

# unparticle propagator

$$\begin{split} \Delta(p,d) &\equiv \int d^4x \, e^{ipx} \langle 0|T\mathcal{O}(x)\mathcal{O}^{\dagger}(0)|0\rangle \\ &= \frac{A_d}{2\pi} \int_0^\infty (M^2)^{d-2} \frac{i}{p^2 - M^2 + i\epsilon} dM^2 \\ &= i \frac{A_d}{2} \frac{\left(-p^2 - i\epsilon\right)^{d-2}}{\sin d\pi} \end{split}$$

## unparticle phase space

$$d\Phi(p,d) = A_d \theta \left( p^0 \right) \theta \left( p^2 \right) \left( p^2 \right)^{d-2}$$

 $d\Phi(p,1) = 2\pi\,\theta\left(p^0\right)\,\delta(p^2)$ 

#### Quarks are Unparticles

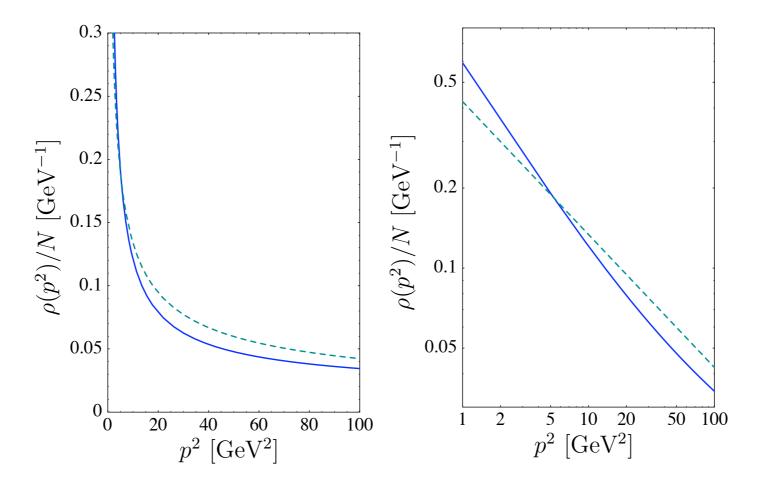
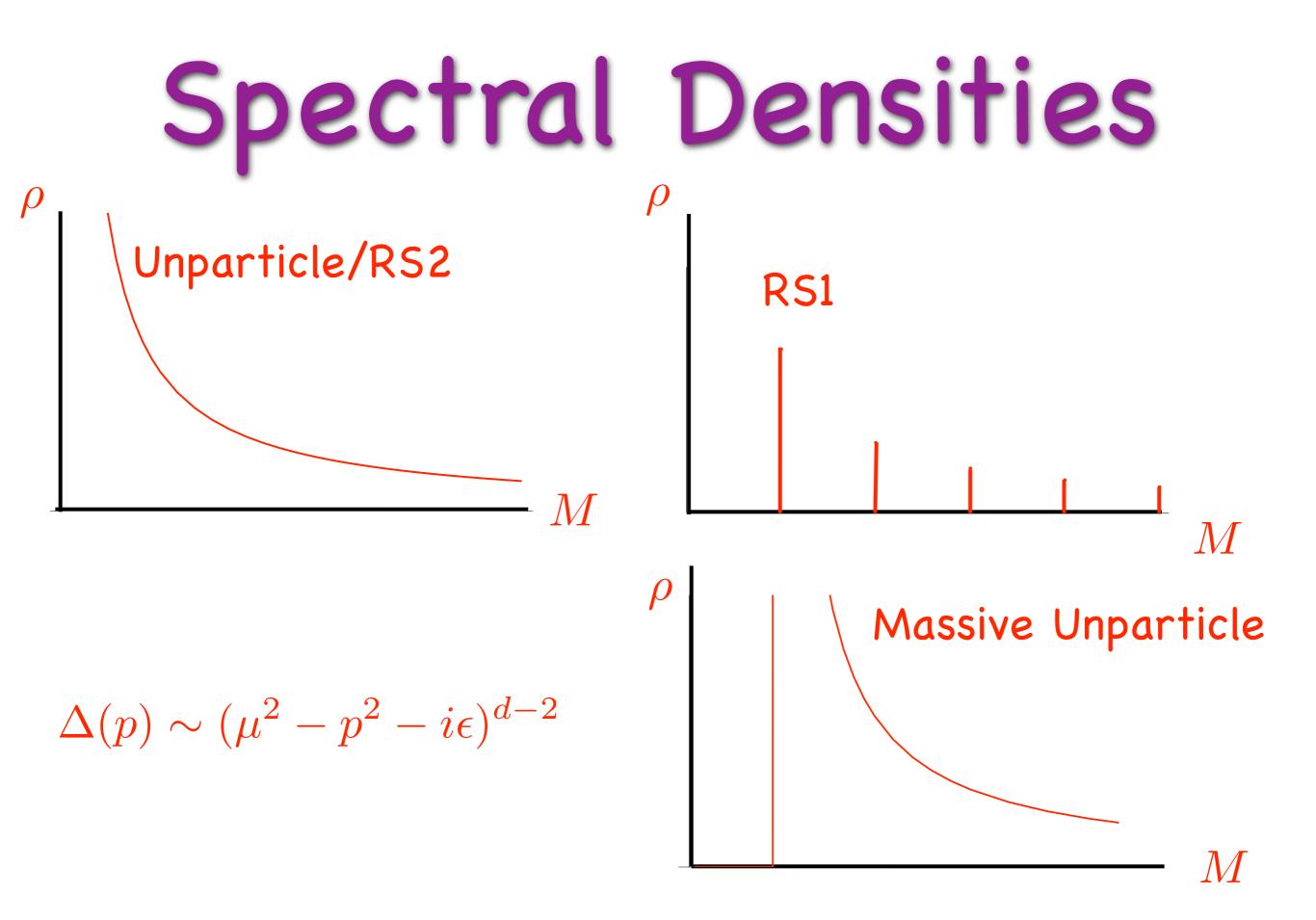


FIG. 1. Comparison of the unparticle spectral density (2) (dashed) and the spectral density (9) of a massless quark jet at next-to-leading order in QCD (solid). We use parameters M = 10 GeV and  $\eta = 0.5$ . The right plot shows the same results on logarithmic scales.

#### Neubert hep-ph/0708.0036



## AdS/CFT/Unparticles

$$S_{int} = \frac{1}{2} \int d^4x \, dz \, \sqrt{g} H \phi \phi$$
$$H = \mu z^2$$
$$z^5 \partial_z \left(\frac{1}{z^3} \partial_z \phi\right) - z^2 (p^2 - \mu^2) \phi - m^2 R^2 \phi = 0$$
$$\delta^{(4)} (p + p') < 2 = 2 d^{-2}$$

$$\langle \mathcal{O}(p')\mathcal{O}(p)\rangle \propto \frac{\partial^{\langle \gamma \rangle}(p+p)}{(2\pi)^4} (p^2 - \mu^2)^{d-2}$$

Cacciapaglia, Marandella JT, hep-ph/0804.0424

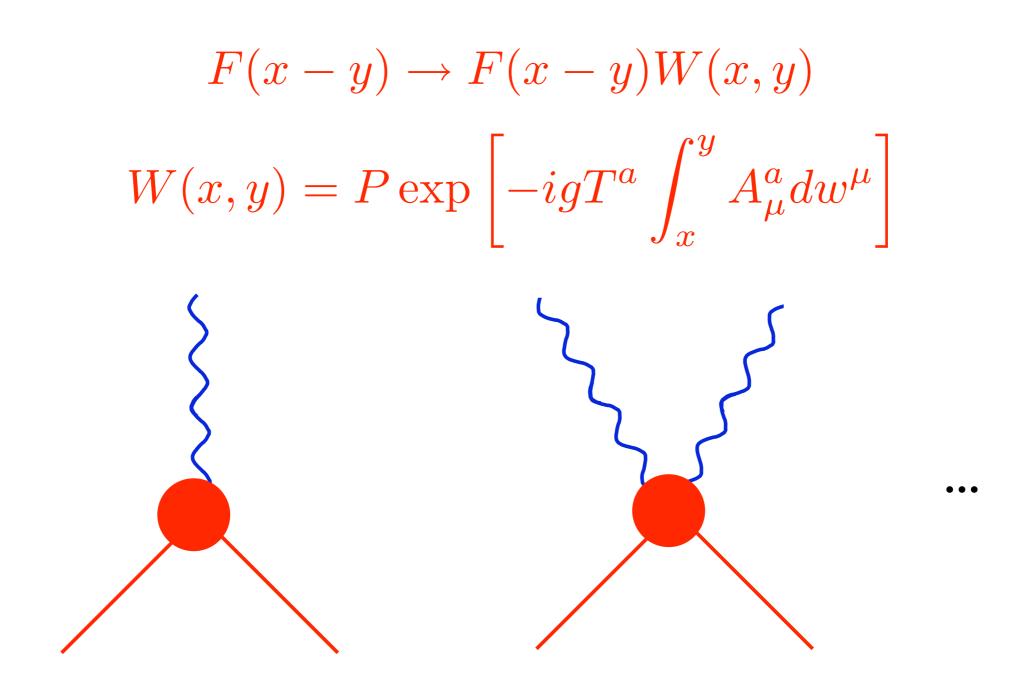
#### **Effective Action**

$$S = \int \frac{d^4 p}{(2\pi)^4} \,\phi^{\dagger}(p) \left[\mu^2 - p^2\right]^{2-d} \phi(p)$$

$$S = \int d^4x d^4y \,\phi^{\dagger}(x) F(x-y)\phi(y)$$
$$F(x-y) = \left[\partial^2 - \mu^2\right]^{2-d} \delta(x-y)$$

nonlocal

#### Minimal Gauge Coupling



cf Mandelstam Ann Phys 19 (1962) 1



$$= -igT^{a}\frac{2p^{\alpha} + q^{\alpha}}{2p \cdot q + q^{2}} \left[ \left(\mu^{2} - (p+q)^{2}\right)^{2-d} - \left(\mu^{2} - p^{2}\right)^{2-d} \right]$$

#### Ward-Takahashi Identity

 $ig\Gamma^{a\alpha}(p,q) = igT^{a}\frac{2p^{\alpha}+q^{\alpha}}{2p\cdot q+q^{2}}\left[\left(\mu^{2}-(p+q)^{2}\right)^{2-d}-\left(\mu^{2}-p^{2}\right)^{2-d}\right]$ 

 $iq_{\mu}\Gamma^{a\mu} = \Delta^{-1}(p+q,m,d)T^a - T^a\Delta^{-1}(p,m,d)$ 

## Unhiggs Model

$$S = -\int d^4x \, H^\dagger (\partial^2 + \mu^2)^{2-d} H$$

$$S = -\int d^4x \, H^{\dagger} (D^2 + \mu^2)^{2-d} H + \lambda_t \overline{t}_R \frac{H}{\Lambda^{d-1}} \left( \begin{array}{c} t \\ b \end{array} \right)_L + h.c.$$

$$S = -\int d^4x \, H^{\dagger} (D^2 + \mu^2)^{2-d} H + \lambda_t \bar{t}_R \frac{H}{\Lambda^{d-1}} \begin{pmatrix} t \\ b \end{pmatrix}_L + h.c.$$
$$-\int d^4x \, \lambda \left(\frac{H^{\dagger} H}{\Lambda^{2d-2}} - \frac{V^2}{2}\right)^2$$

$$H = \frac{1}{\sqrt{2}} e^{iT^a \pi^a / v^d} \begin{pmatrix} 0 \\ v^d + h \end{pmatrix}$$

Stancato JT, hep-ph/0807.3961

### Mass Divergence

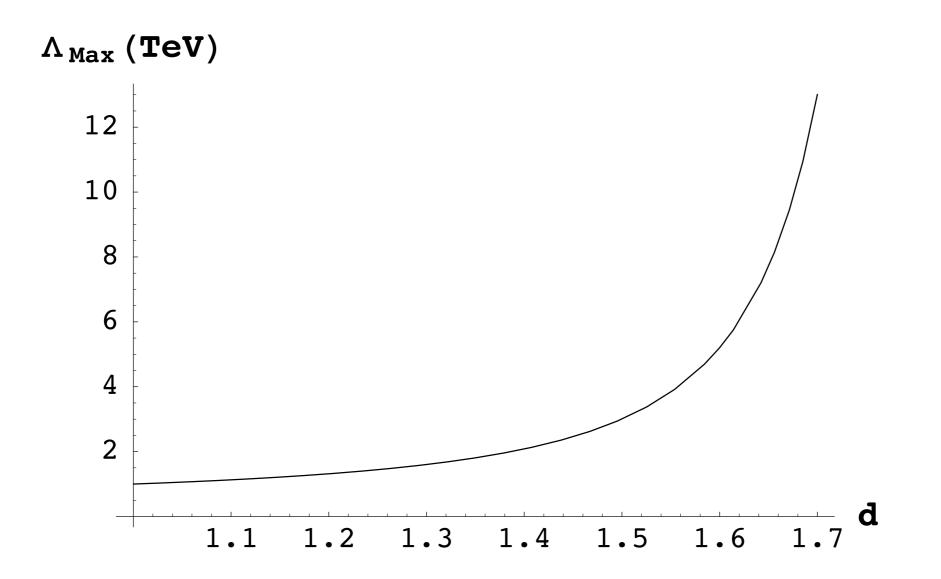
$$m_h^2 \sim \frac{\Lambda^{4-2d}}{16\pi^2}$$

$$\mathcal{L}_Y = \lambda_t \, \bar{t}_R \frac{H}{\Lambda^{d-1}} \left( \begin{array}{c} t \\ b \end{array} \right)_L$$

$$m_h^2 = 3\left(\frac{\lambda_t}{\Lambda^{d-1}}\right)^2 \frac{\Lambda^2}{16\pi^2} = 3\left(\frac{m_t}{V}\right)^2 \frac{\Lambda^{4-2d}}{16\pi^2}$$

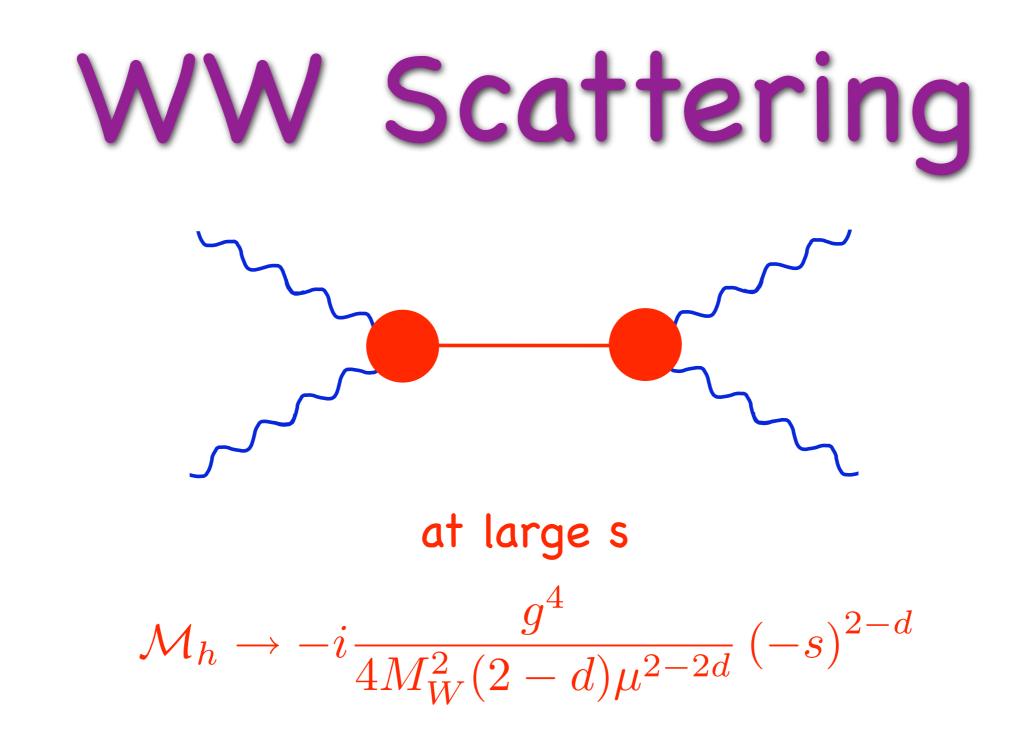
Solve the little hierarchy problem?

### loop < tree

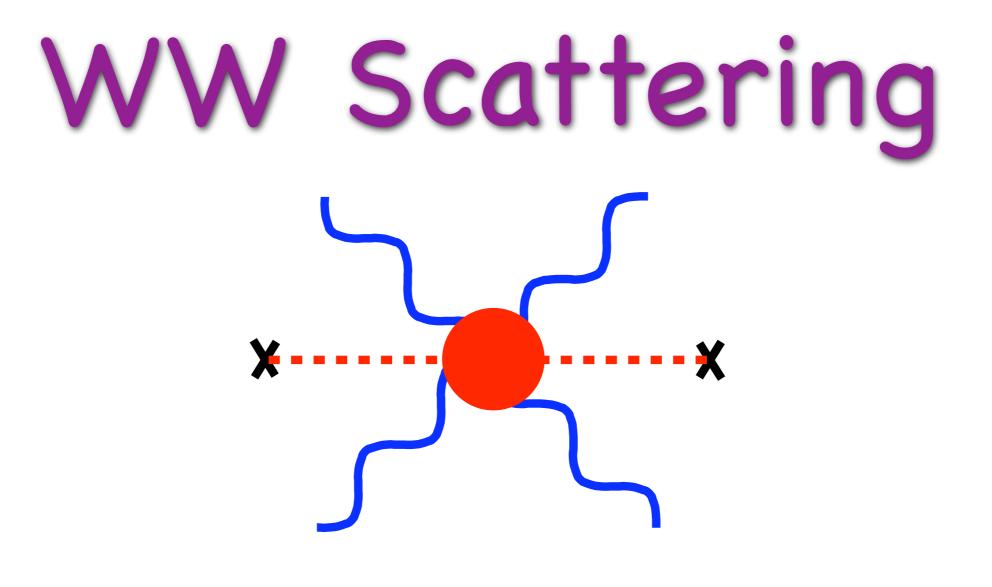


$$-g^{2}A_{\alpha}^{a}A_{\beta}^{b}\langle H^{\dagger}\rangle T^{a}T^{b}\langle H\rangle \left\{ g^{\alpha\beta}(d-2)\mu^{2-2d} - \frac{q^{\alpha}q^{\beta}}{q^{2}} \left[ (d-2)\mu^{2-2d} - \frac{(\mu^{2}-q^{2})^{2-d} - (\mu^{2})^{2-d}}{q^{2}} \right] \right\}$$

$$M_W^2 = \frac{g^2(2-d)\mu^{2-2d}v^{2d}}{4}$$

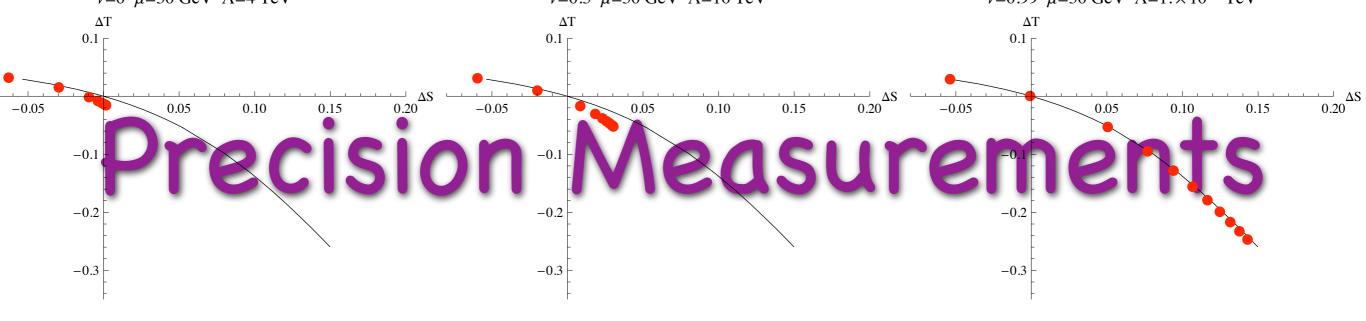


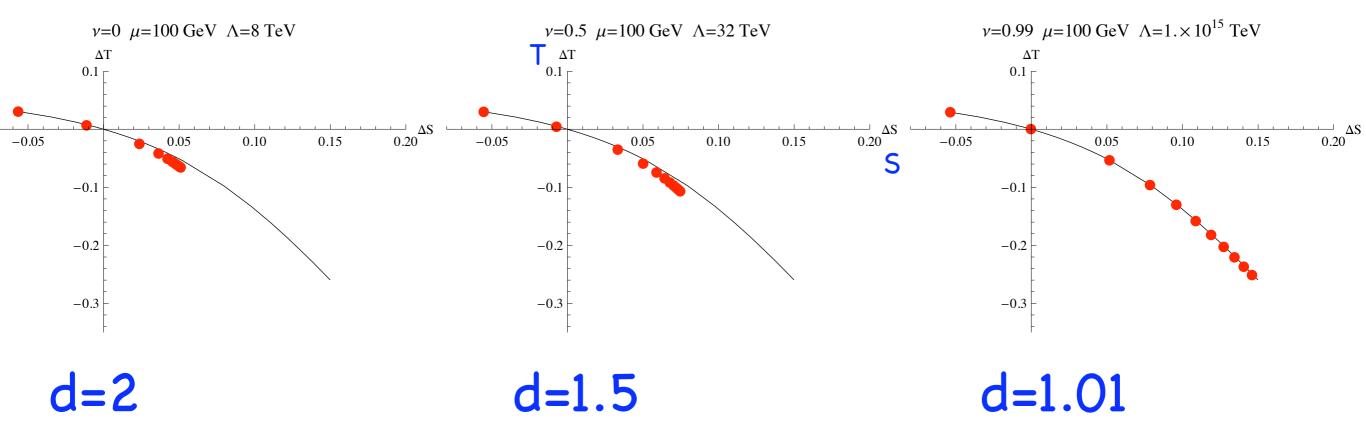
#### unHiggs exchange is insufficient to unitarize WW scattering



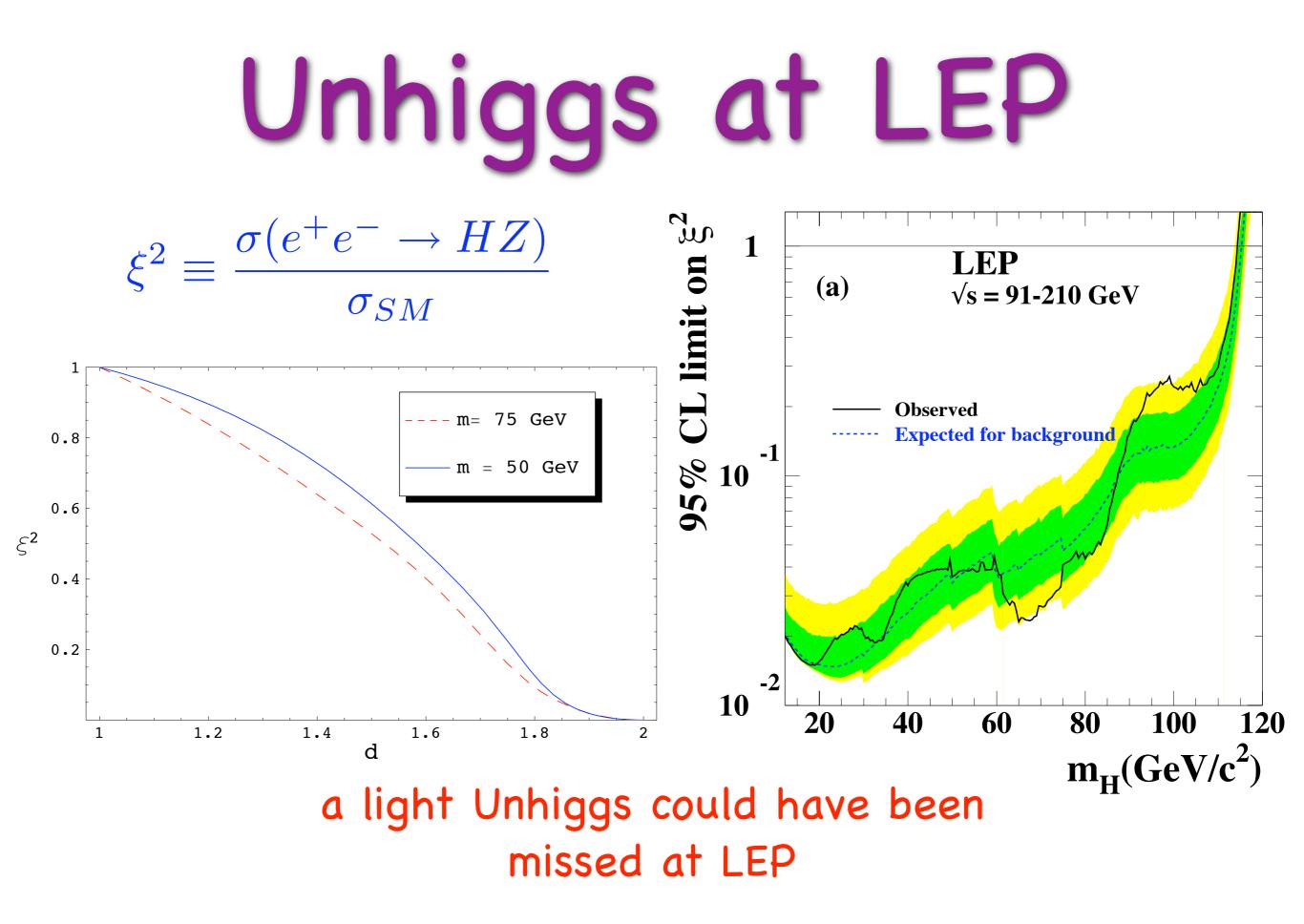
$$\mathcal{M}_{hh} = -i \frac{g^2}{4M_W^2} \left[ s + \frac{(-s)^{2-d}}{(2-d)\mu^{2-2d}} \right]$$

unHiggs 6 point vertex does unitarize WW scattering





Falkowski & Perez-Victoria, hep-ph/0901.3777



# Unhiggs at LHC

#### generically a light Unhiggs would be missed at the LHC using current search strategies

#### Conclusions

massive unparticles with gauge interactions are a new type of BSM physics

> the Unhiggs is a new way to break electroweak symmetry

we need to identify the new LHC signals

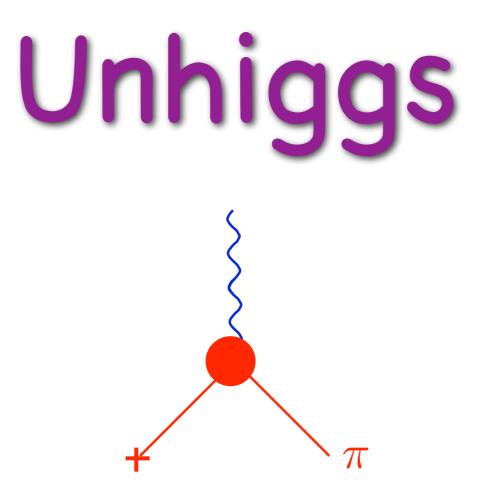


#### AdS/CFT

$$ds^{2} = \frac{R^{2}}{z^{2}} \left( dx_{\mu}^{2} - dz^{2} \right)$$

 $z > \epsilon$ 

$$S_{bulk} = \frac{1}{2} \int d^4x \, dz \sqrt{g} (g^{\alpha\beta} \partial_\alpha \phi \partial_\beta \phi + m^2 \phi^2)$$
$$\phi(p, z) = az^2 J_\nu(pz) + bz^2 J_{-\nu}(pz)$$
$$d[\mathcal{O}] = 2 \pm \nu = 2 \pm \sqrt{4 + m^2 R^2}$$



$$g\left(\langle H^{\dagger}\rangle A^{a}_{\alpha}T^{a}\Pi - \Pi^{\dagger}A^{a}_{\alpha}T^{a}\langle H\rangle\right) \\ \times \left[\left(\mu^{2} - q^{2}\right)^{2-d} - \left(\mu^{2}\right)^{2-d}\right]q^{\alpha}/q^{2}$$

Unhiggs  

$$\begin{array}{l} & & \\ & & \\ & & \\ & & \\ & & \\ & & \\ & \Pi^{ab\alpha\beta}(q) = -g^2 \langle H^{\dagger} \rangle T^a T^b \langle H \rangle \frac{q^{\alpha} q^{\beta}}{q^4} \\ & & \\$$