Flavor & New Physics

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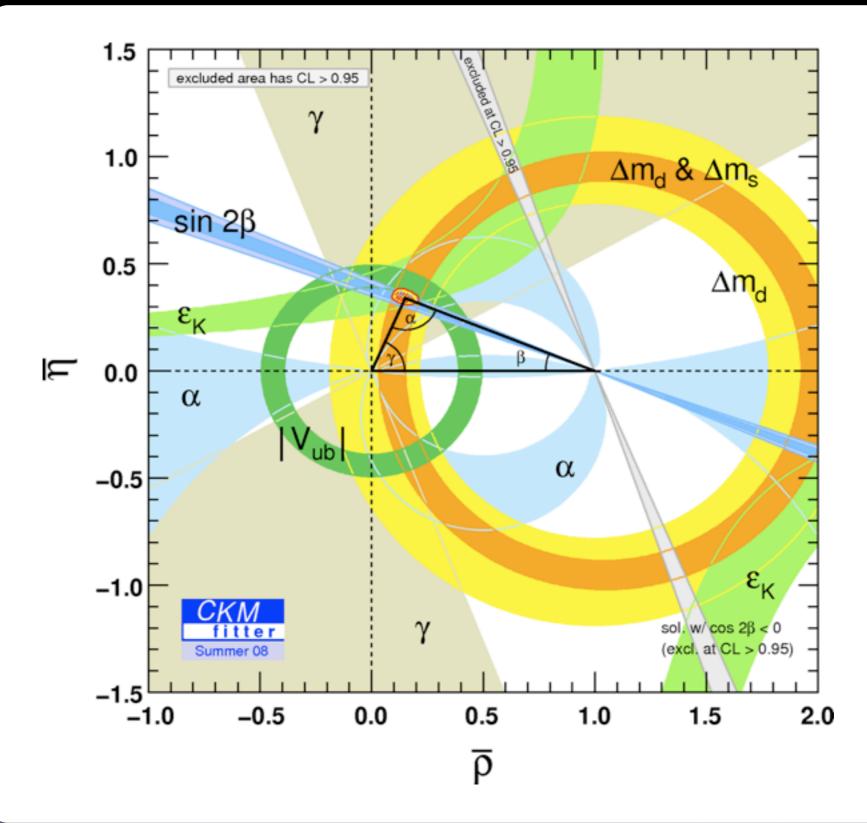
Aspen Winter Conference 2009 The Year of the Ox

New Era in 2009+x

We will probe the origin of EWSB.

Will we be able to learn about the origin of flavor in the era of the LHC?

Does the mechanism stabilizing the weak scale have non-trivial flavor content?



Flavor and CP in the SM

Experimentalist's view

- + spectrum, BR, A_{CP}
- + measure masses, mixing angles and phases

Theorist's view + In the absence of Yukawas SM globally $SU(3)_{Q_L} \times SU(3)_{u_R} \times SU(3)_{d_R}$ symmetric

$$v Y_u = U_u \begin{pmatrix} m_u & & \\ & m_c & \\ & & m_t \end{pmatrix} V_u \qquad v Y_d = U_d \begin{pmatrix} m_d & & \\ & m_s & \\ & & m_b \end{pmatrix} V_d$$

Charged currents: measure only LH misalignment

$$v Y_u = \mathcal{V}_u \begin{pmatrix} m_u \\ m_c \\ m_t \end{pmatrix} \mathcal{V}_u$$

$$v Y_d = V_{\rm CKM} \begin{pmatrix} m_d & & \\ & m_s & \\ & & m_b \end{pmatrix} V_d$$

Charged currents: measure only LH misalignment

Neutral currents: enhanced flavor symmetry $SU(3)_Q \rightarrow SU(3)_{u_L} \times SU(3)_{d_L}$ Yukawas diagonal, **no (tree-level) flavor violation**

$$v Y_u = \mathcal{V}_u \begin{pmatrix} m_u \\ m_c \\ m_t \end{pmatrix} \mathcal{V}_u$$

$$v Y_d = \mathcal{V}_{\text{CKM}} \begin{pmatrix} m_d & m_s & \\ & m_s & \\ & & m_b \end{pmatrix} \mathcal{V}_d$$

Smallness & hierarchy

$$Y_D \approx (10^{-5}, 0.0005, 0.026)$$

$$Y_U \approx \begin{pmatrix} 10^{-5} & -0.002 & 0.007 + 0.004i \\ 10^{-6} & 0.007 & -0.04 + 0.0008i \\ 10^{-8} + 10^{-7}i & 0.0003 & 0.96 \end{pmatrix}$$

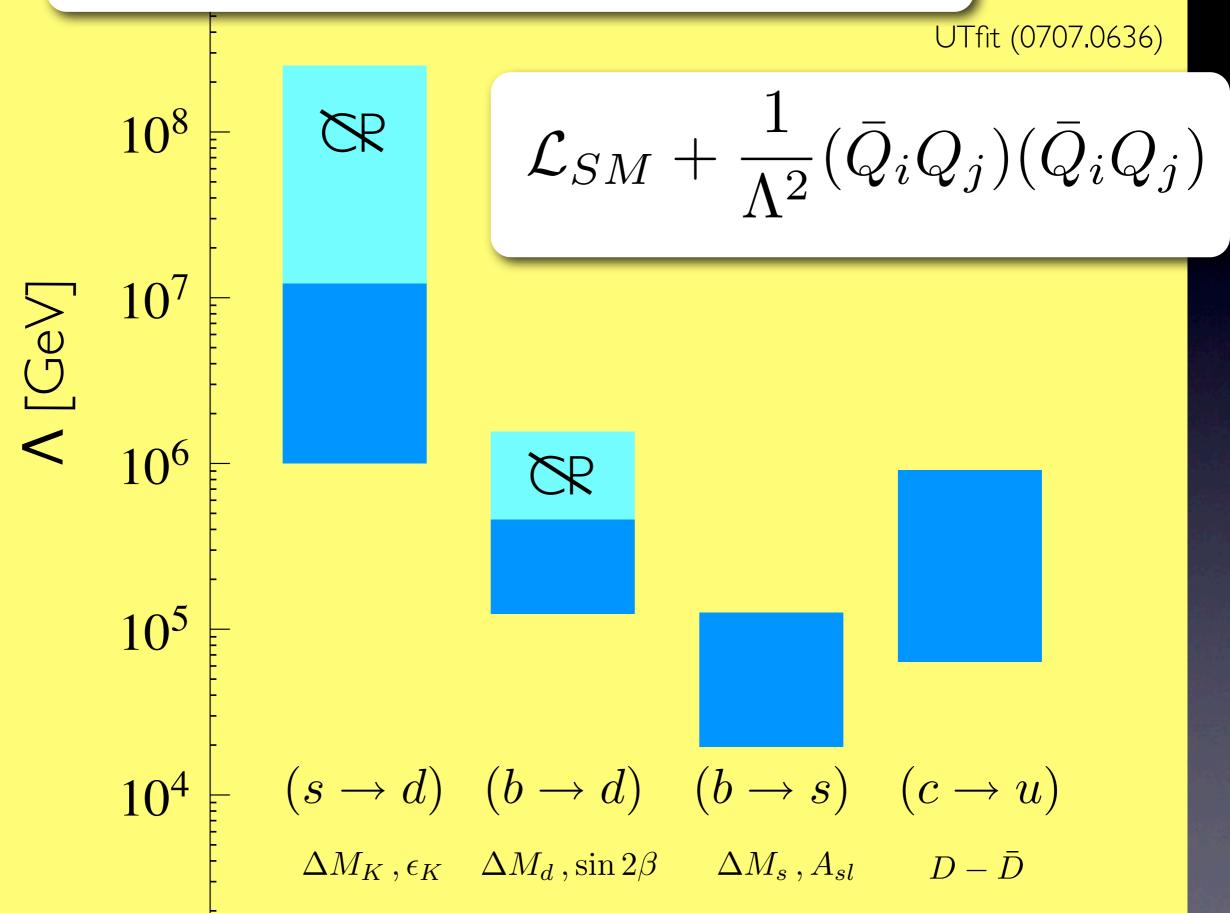
The SM flavor parameters have structure: smallness & hierarchy. Why? The SM flavor puzzle.

Compare to: $g_s \sim I$, $g \sim 0.6$, $g \sim 0.3$, $\lambda_{Higgs} \sim I$

Flavor and CP in the SM

Yukawa matrices $Y_{\cup} \& Y_{D}$ contain everything $(\bar{u}_R^i u_R^j)$ $(\bar{Q}_L^i Q_L^j)$ $Y_U^{\dagger}Y_U$ $Y_U Y_U^\dagger$ NP ? $(ar{d}_R^i d_R^j)$ $Y_D Y_D^{\dagger}$ $Y_D^{\dagger} Y_D$

Bounds on generic flavor violation



How can we protect TeV physics from these bounds?

A radical cure: MFV

New particles/interactions, but flavor structure ~VCKM

 $(\bar{Q}_L^i Q_L^j)$ $Y_U Y_U^\dagger$ $(\bar{d}_R^i d_R^j)$ $Y_D Y$ $Y_D^{\dagger} Y_D$ + LR, RL caveat: large tanß see Carlos Wagner's talk

Model independent $\Delta F=2$ MFV Bound

 10^{8}

 10^{7}

106

10⁵

 10^{4}

1000

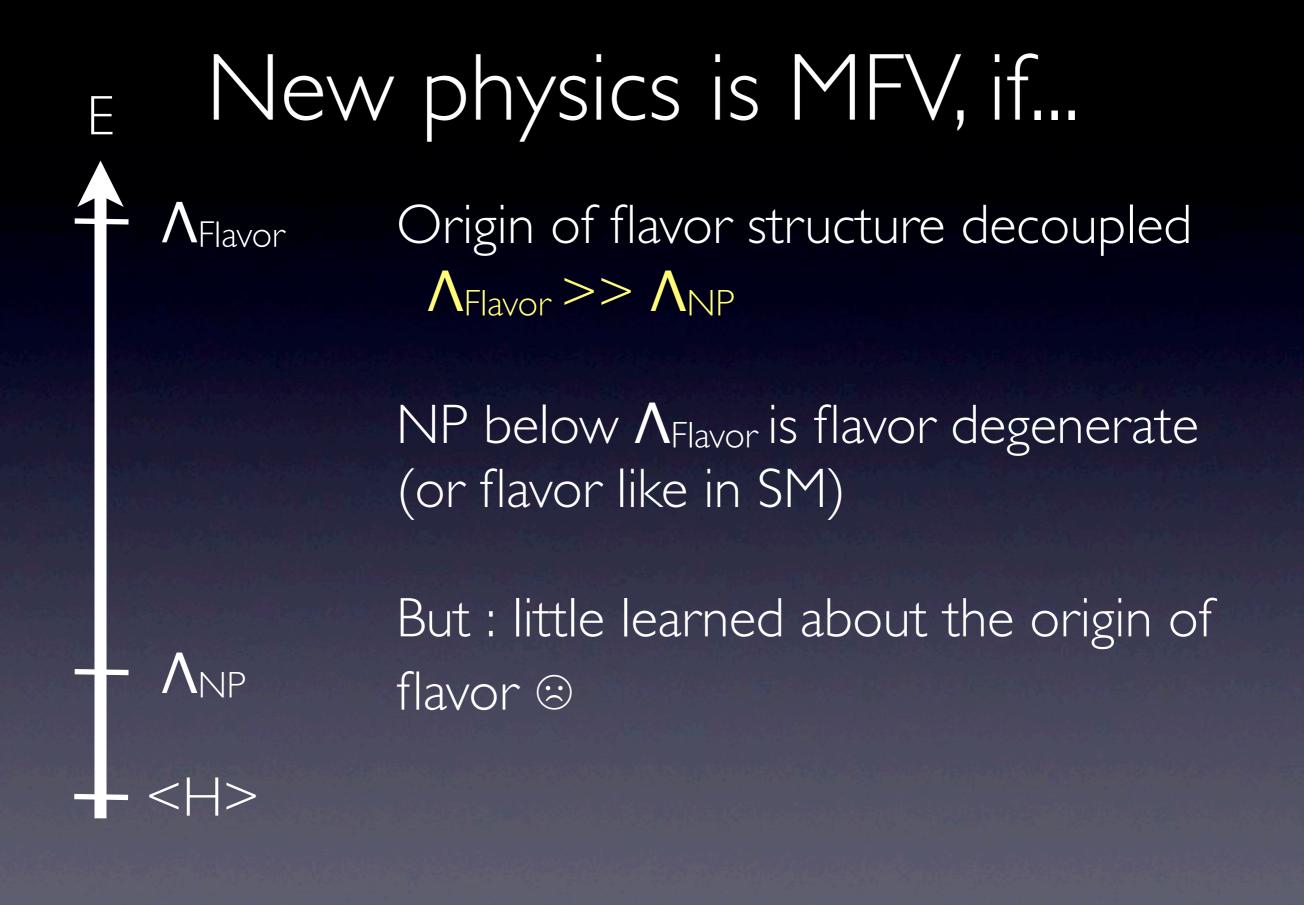
∧ [GeV]

UTfit (0707.0636)

 $\mathcal{L}_{SM} + \frac{1}{\Lambda^2} (\bar{Q}_i Q_j) (\bar{Q}_i Q_j)$

corresponds to m ≈ 5 TeV (tree level) m ≈ 500 GeV (**α**sI-loop)

 $\Delta M_K, \epsilon_K \quad \Delta M_d, \sin 2\beta \qquad \Delta M_s, A_{sl} \qquad D - \bar{D}$



MFV example: SUSY

MSSM with unbroken SUSY is already MFV!

=> MSSM is MFV if SUSY is flavor blind

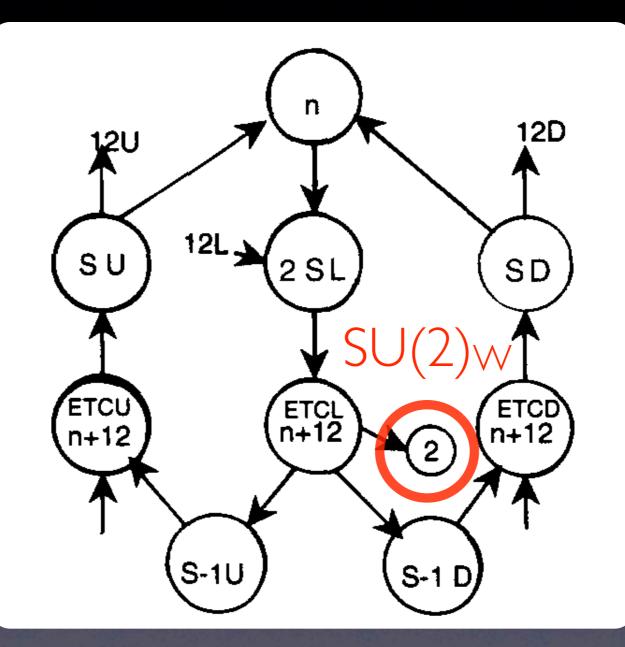
Example: Gauge mediation with $M_{mess} << \Lambda_{flavor}$

Gravity mediation in general not MFV, mSUGRA not a good starting point to study flavor!

MFV SUSY alternatives: UED and the Littlest Higgs with appropriate UV completions.

MFV Technicolor?

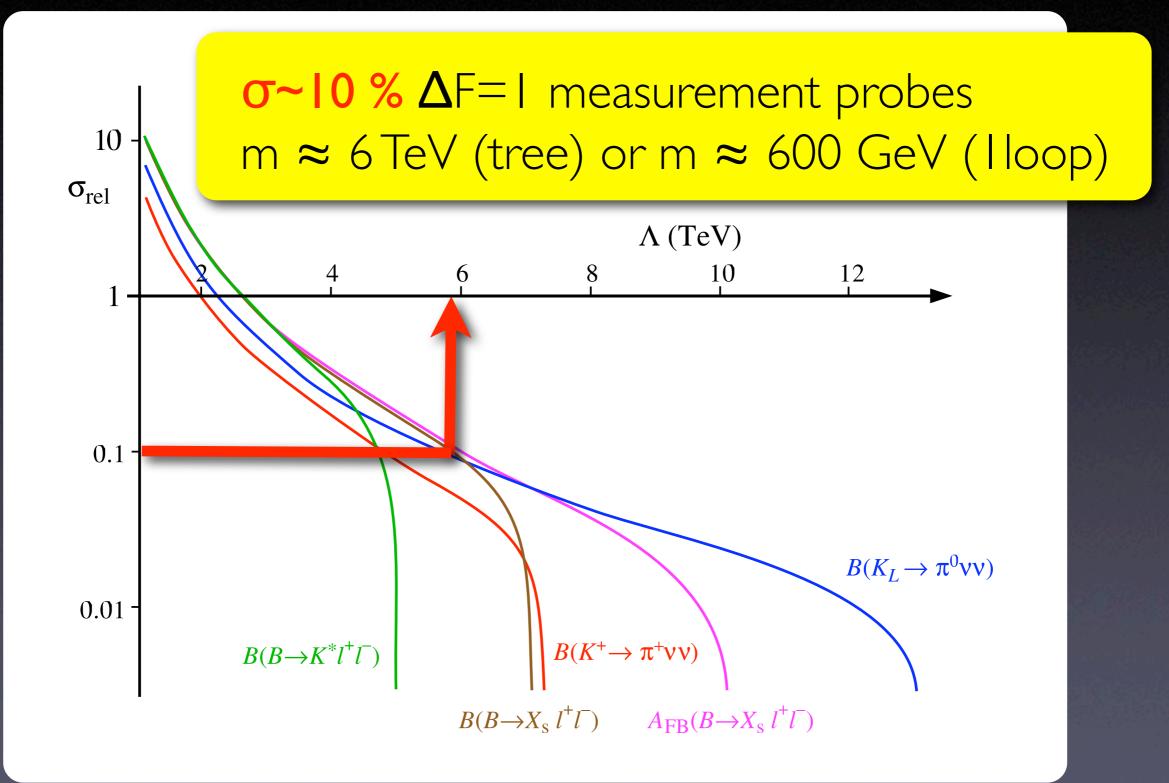
Chivukula, Georgi '87; Chivukula, Georgi, Randall '87; Randall '93; Georgi '94, Skiba '96



Simpler alternative: AdS/CFT interpretation of RS : 5D GIM mechanism Cacciapaglia, Csaki, Galloway, Marandella, Terning, A.W., '08

Distinguishing MFV & SM is hard

D'Ambrosio, Giudice, Isidori, Strumia '02; Buras, Bryman, Isidori, Littenberg '05



Falsifying MFV is easy... ...once you have shown that the SM is dead

Bobeth, Bona, Buras, Ewerth, Pierini, Silvestrini, A.W. MFV falsified by violating ''sum rules''

already in the data? Lunghi, Soni '08

New CP phases

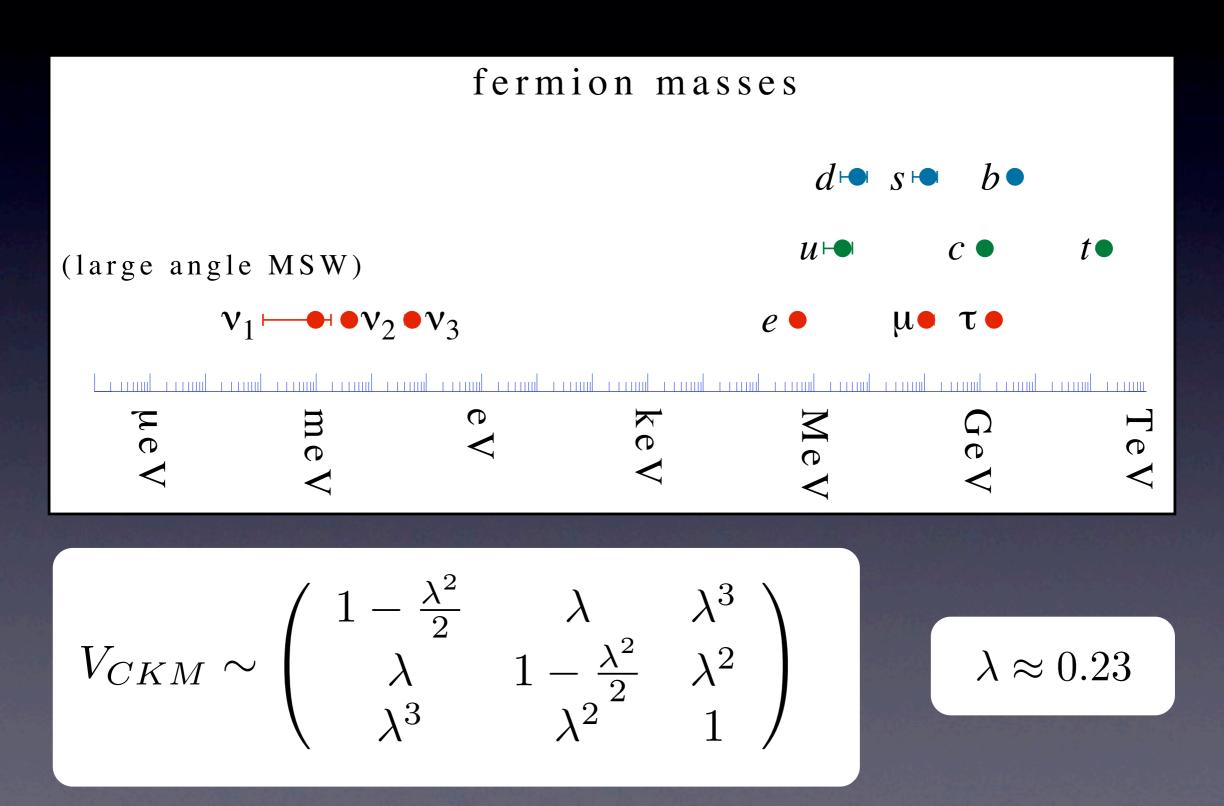
MFV@LHC: Grossman, Nir, Thaler, Volansky, Zupan At the LHC: $Br(q_3) \sim Br(q_{1,2})$

Top FCNCs

A theory of flavor at the LHC?

Froggat-Nielsen vs. Higgs dependent Yukawas RS/holographic technicolor

Fermion masses & mixings



Hierarchies from symmetries Froggatt, Nielsen '79

Add horizontal U(I)_F, flavon Φ_F (m $_{\Phi} \sim \Lambda$, q_F = -1)

$$Y_{d}^{ij} \left(\frac{\Phi_{F}}{\Lambda}\right)^{-q_{i}+h+d_{j}} \bar{Q}_{L}^{i} H D_{R}^{j}$$

$$\int U(1)_{F} \text{ broken by } F = \langle \Phi_{F} \rangle, \quad F < \Lambda$$

$$Y_{eff,d}^{ij} = Y_{d}^{ij} \left(\frac{F}{\Lambda}\right)^{-q_{i}+h+d_{j}} => \text{ hierarchies}$$

But: hard to probe, flavon must be heavy $m_{\phi} \gg TeV$

Higgs as flavon

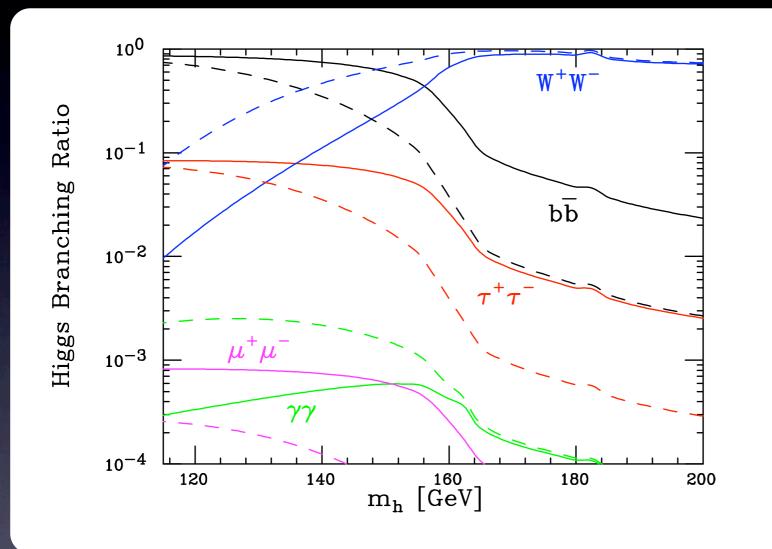
Babu, Nandi '99; Giudice, Lebedev '08

"Higgs dependent Yukawas" Yukawas effective interaction after integrating out heavy physics. Postulate leading terms are absent $\mathcal{L}_Y = Y_{ij}^u(H) \ \bar{q}_{Li} u_{Rj} H^c + Y_{ij}^d(H) \ \bar{q}_{Li} d_{Rj} H$ $Y_{ij}^{u,d}(H) = c_{ij}^{u,d} \left(\frac{H^{\dagger}H}{M^2}\right)^{n_{ij}^{u,d}}$ nij generation dependent integer, determines mass hierarchy

 $v^2/M^2 \approx m_b^2/m_t^2 \implies M \approx 1 - 2 \,\mathrm{TeV}$

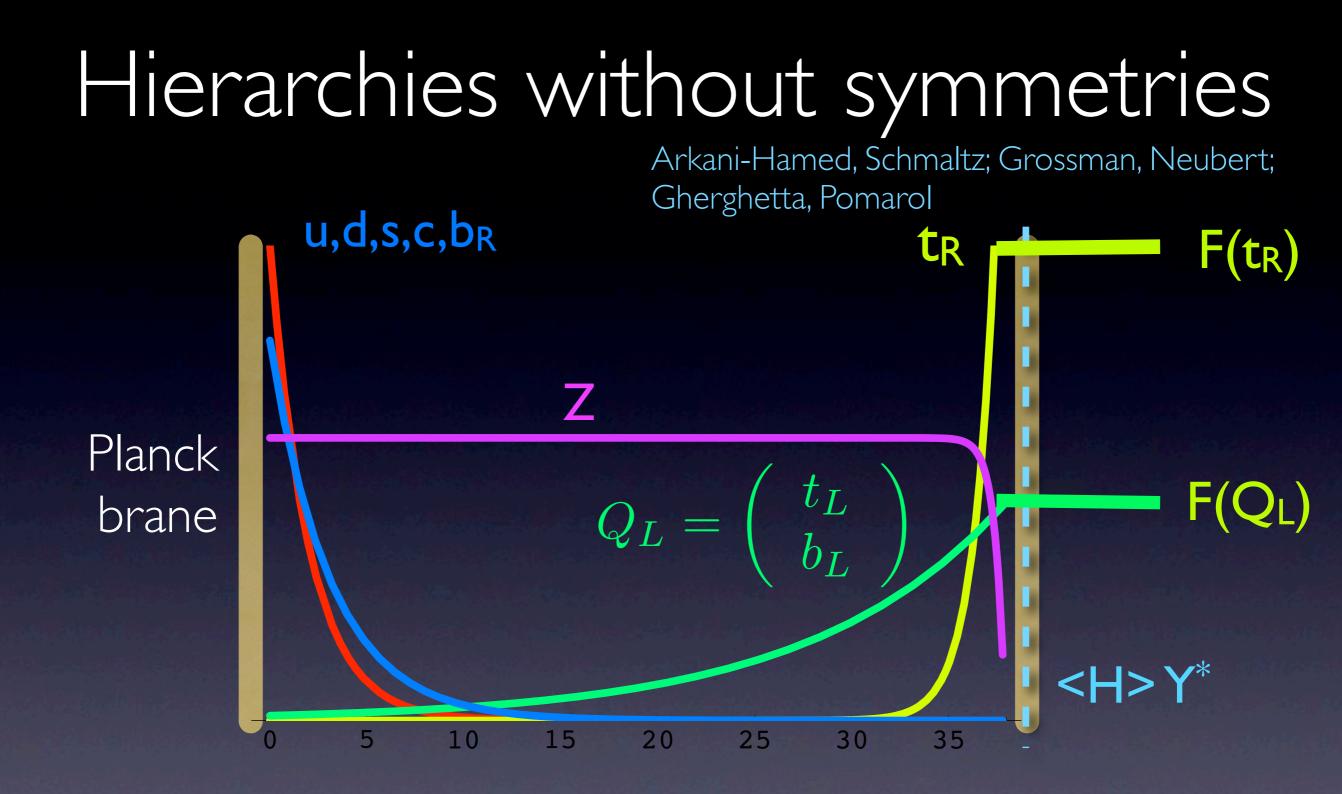
Higgs as flavon: signals

Giudice, Lebedev '08





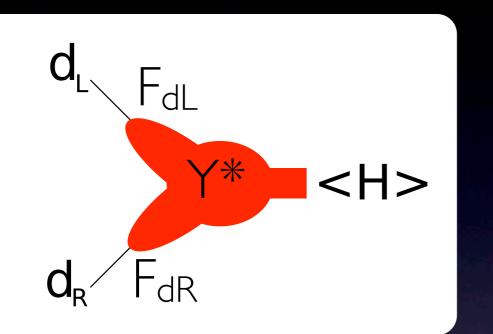
 $\frac{\Gamma\left(h \to b\overline{b}\right)}{\Gamma\left(h \to b\overline{b}\right)_{SM}} = \frac{\Gamma\left(h \to c\overline{c}\right)}{\Gamma\left(h \to c\overline{c}\right)_{SM}} = \frac{\Gamma\left(h \to \tau^{+}\tau^{-}\right)}{\Gamma\left(h \to \tau^{+}\tau^{-}\right)_{SM}} = 9 \qquad \frac{\Gamma\left(h \to \mu^{+}\mu^{-}\right)}{\Gamma\left(h \to \mu^{+}\mu^{-}\right)_{SM}} = 25$



Localization in extra dimension determines overlap $F(q_i)$ with Higgs.

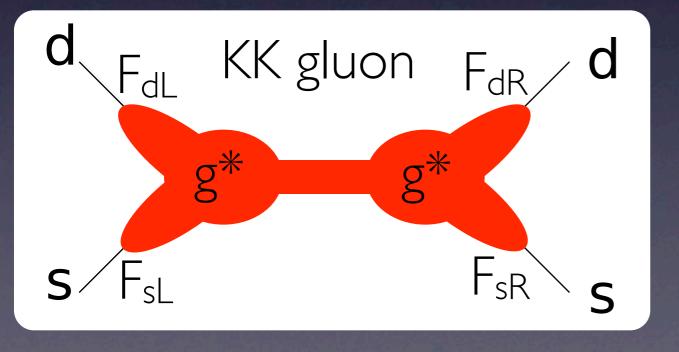
Masses, mixings and FCNCs

Gherghetta, Pomarol; Agashe, Perez, Soni



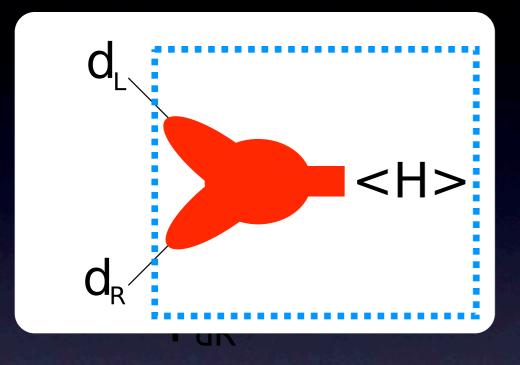
masses and mixings from hierarchical overlaps

 $m_d \sim v F_{d_L} Y^* F_{d_R}$ RS G

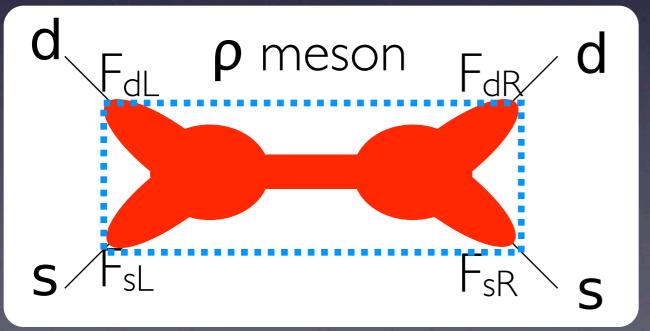


KK gluon FCNCs due to the same small overlaps F_i: $\sim \frac{(g^*)^2}{M_{KK}^2} F_{d_L} F_{d_R} F_{s_L} F_{s_R}$ $\sim \frac{(g^*)^2}{M_{KK}^2} \frac{m_d m_s}{(vY^*)^2}$

Partial compositeness

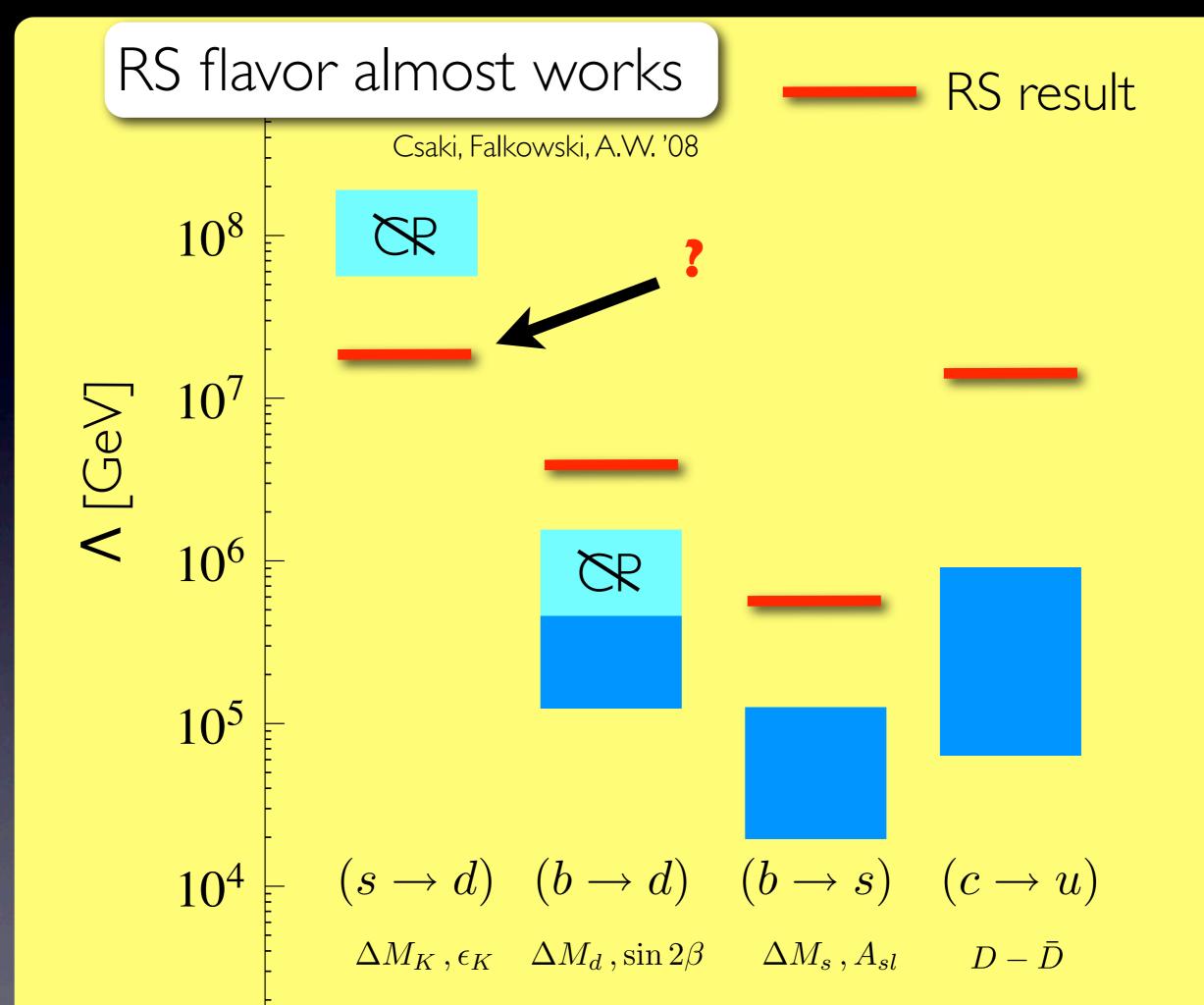


Higgs is part of the strong sector, couples only to composites $m_d \sim \langle H \rangle \Upsilon^* F_{dL} F_{dR}$ F_i amount of compositeness



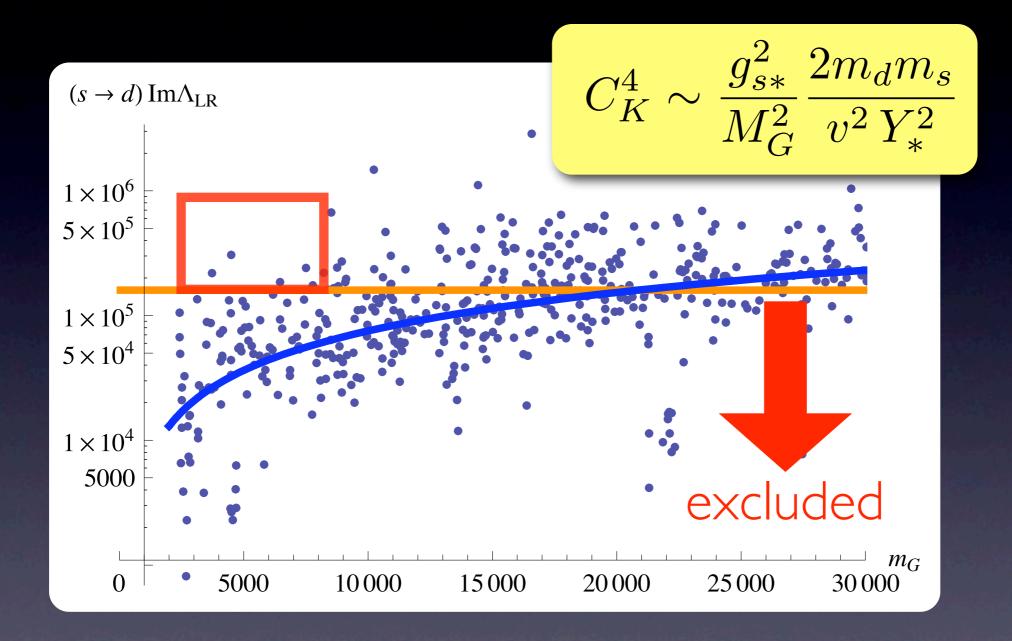
The smaller the mass, the smaller the compositeness, the smaller the FCNC

> ~ $F_{dL} F_{sL} F_{dR} F_{sR}$ ~ $m_d m_s / (<H>Y^*)^2$



Bound on the KK gluon mass

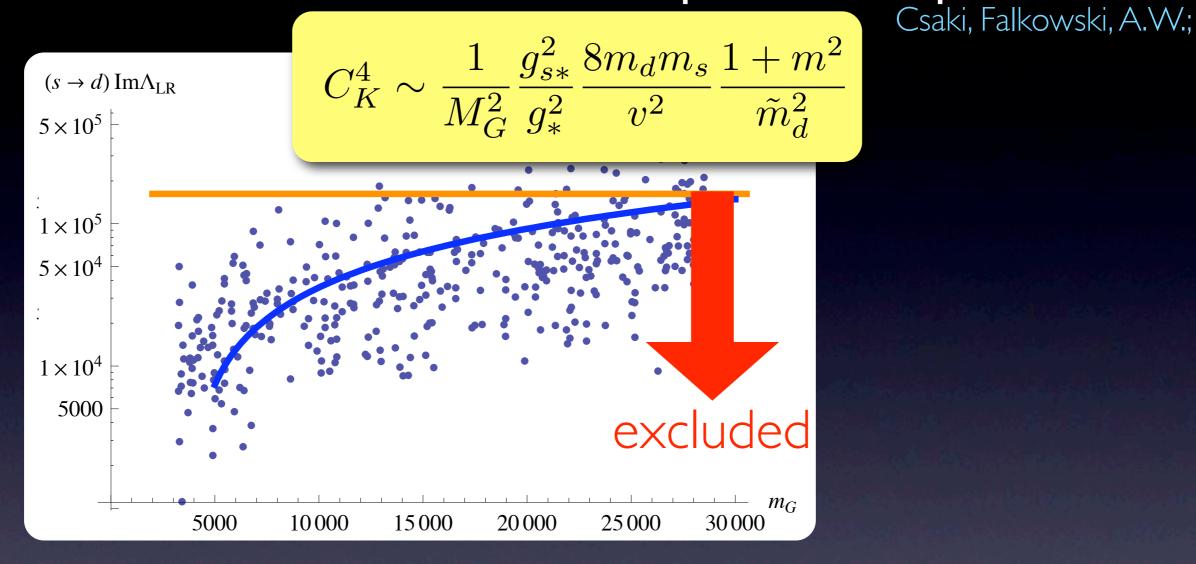
Csaki, Falkowski, A.W.; Casagrande et al.; Buras et. al.



Some points above the bound: any rationale to live here? Radiative stability?

remore in M. Neubert's talk

Bound in the composite pGB



more flavor violation in composite pGB: $Y^* \rightarrow g_* / 2$ & fermionic kinetic mixings $M_{KK} > 30 \text{ TeV}$

Conclusions

Flavor searches are complementary to direct searches at the LHC.

We have learned that NP must have a highly non-generic flavor sector.

Finding deviations from MFV can give us insights into the origin of the Yukawa couplings.

I am looking forward to the era of precision flavor physics driven by CDF, D0,LHCb, NA48/3, E391, SuperB, ATLAS, CMS & theoretical efforts.

Low KK scale w/o adding flavor structure

+ live with fine-tuned Yukawas (large radiative corrections)

or

Agashe, Azatov, Zhu + bulk Higgs model (not applicable to pGB), push Yukawa to perturbative limit Y*> 6 and g_{s*} as small as possible (1-loop matching)

$$M_{KK} > \frac{g_{s*}}{Y^*} \frac{\sqrt{2m_d m_s}}{v} \Lambda_4$$

With some tuning $M_{KK} \sim 5 \text{ TeV}$ possible Testable at LHC?

Low KK scale by adding flavor structure

Cacciapaglia, Csaki, Galloway, Marandella, Terning, AW

+ exact GIM structure flavor symmetry in bulk and IR brane, UV kinetic terms generate flavor, no explanation for fermion masses (likely the only way for Higgsless)

Csaki, Falkowski, AW

+ Add horizontal U(I)'s

Fitzpatrick, Randall, Perez

+ 5D MFV only two flavor spurions (Y_U, Y_D) Need to tune to align bulk and brane matrices.