

# TMD<sub>e</sub> 2015

**ICTP, Trieste: Sept. 2-4, 2015.**



## ***TRANSVERSE SPIN EFFECTS IN TWO HADRON ELECTROPRODUCTION***

Hrayr Matevosyan

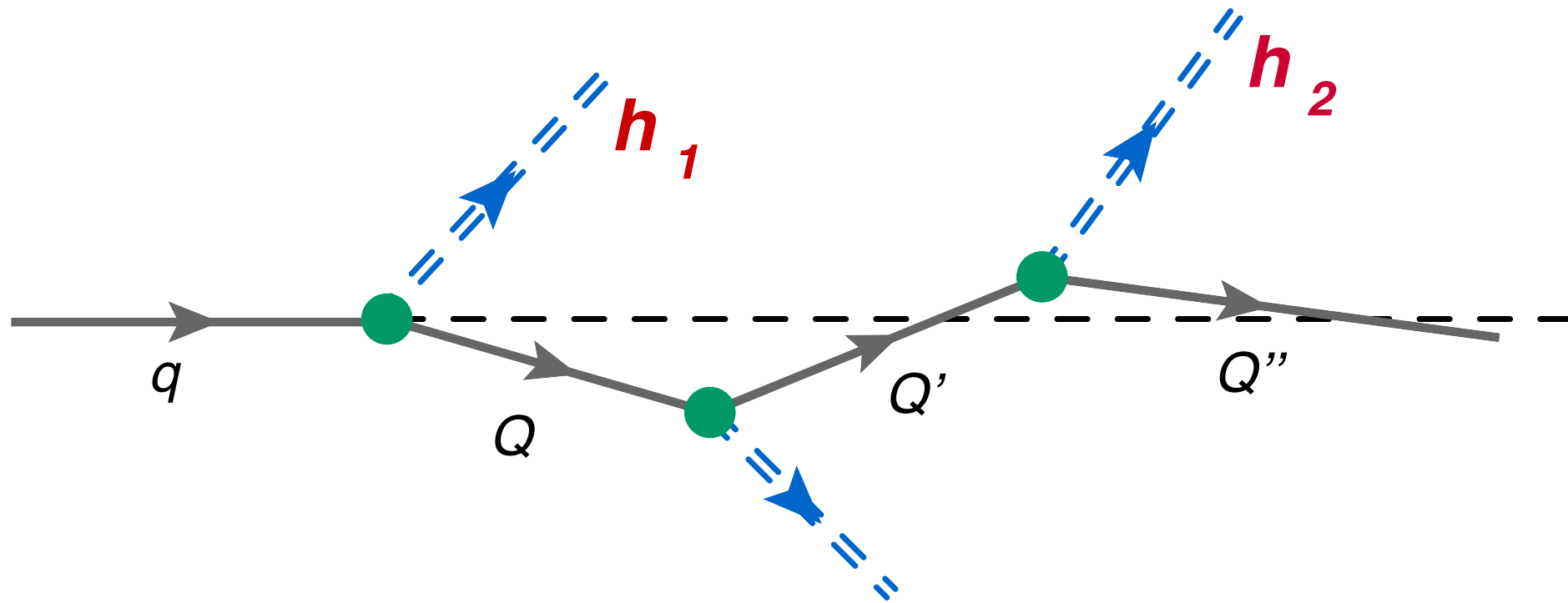
*Collaborators:*

A. Kotzinian, A.W.Thomas,  
E.-C.Aschenauer, H.Avakian.



# OUTLOOK

- ❖ **Sivers Effect in *Two Hadron SIDIS* and Unpolarized Dihadron Fragmentation Functions.**
- ❖ **Interference *DiFF* type modulations from Collins effect.**



***TWO HADRON CORRELATIONS:  
DIHADRON FRAGMENTATION FUNCTIONS***



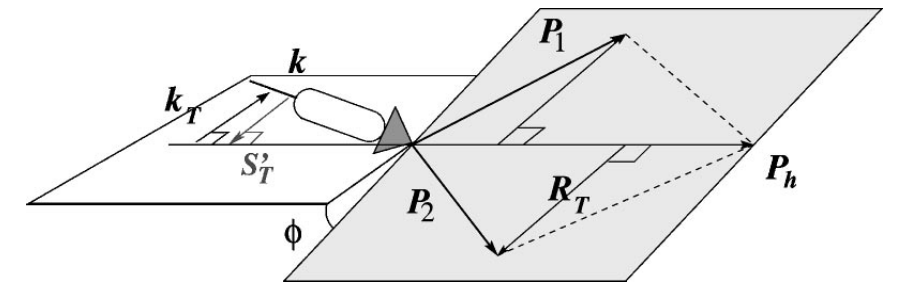
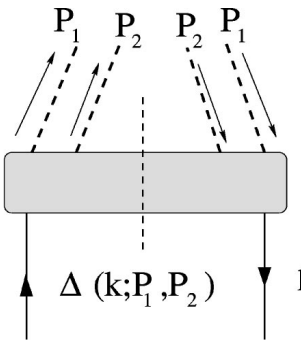
# TWO-HADRON FRAGMENTATION

A. Bianconi, et al: PRD 62, 034008 (2000). M. Radici, et al: PRD 65, 074031 (2002).

## Kinematic Variables:

$$P_1 = \left[ \xi P_h^-, \frac{M_1^2 + \vec{R}_T^2}{2\xi P_h^-}, \vec{R}_T \right], \quad k = \left[ \frac{P_h^-}{z}, z, \frac{k^2 + \vec{k}_T^2}{2P_h^-}, \vec{k}_T \right]$$

$$P_2 = \left[ (1-\xi)P_h^-, \frac{M_2^2 + \vec{R}_T^2}{2(1-\xi)P_h^-}, -\vec{R}_T \right], \quad \mathbf{R} = \frac{\mathbf{P}_1 - \mathbf{P}_2}{2}$$



$$z \equiv z_h = z_1 + z_2$$

$$\xi = \frac{z_1}{z_1 + z_2}$$

- The relevant terms of the quark correlator at leading order for a **Transversely Polarized Quark:**

**Unpolarized**

$$\Delta^{[\gamma^-]} = D_1(z_h, \xi, k_T^2, R_T^2, \mathbf{k}_T \cdot \mathbf{R}_T)$$

**Interference**

$$\Delta^{[i\sigma^{i-}\gamma_5]} = \frac{\epsilon_T^{ij} R_{Tj}}{M_1 + M_2} H_1^\triangleleft(z_h, \xi, k_T^2, R_T^2, \mathbf{k}_T \cdot \mathbf{R}_T) + \frac{\epsilon_T^{ij} k_{Tj}}{M_1 + M_2} H_1^\perp(z_h, \xi, k_T^2, R_T^2, \mathbf{k}_T \cdot \mathbf{R}_T)$$

- IFFS are Chiral-ODD:** Need to be coupled with another chiral-odd quantity to be observed (e.g. transversity).



# TWO-HADRON FRAGMENTATION

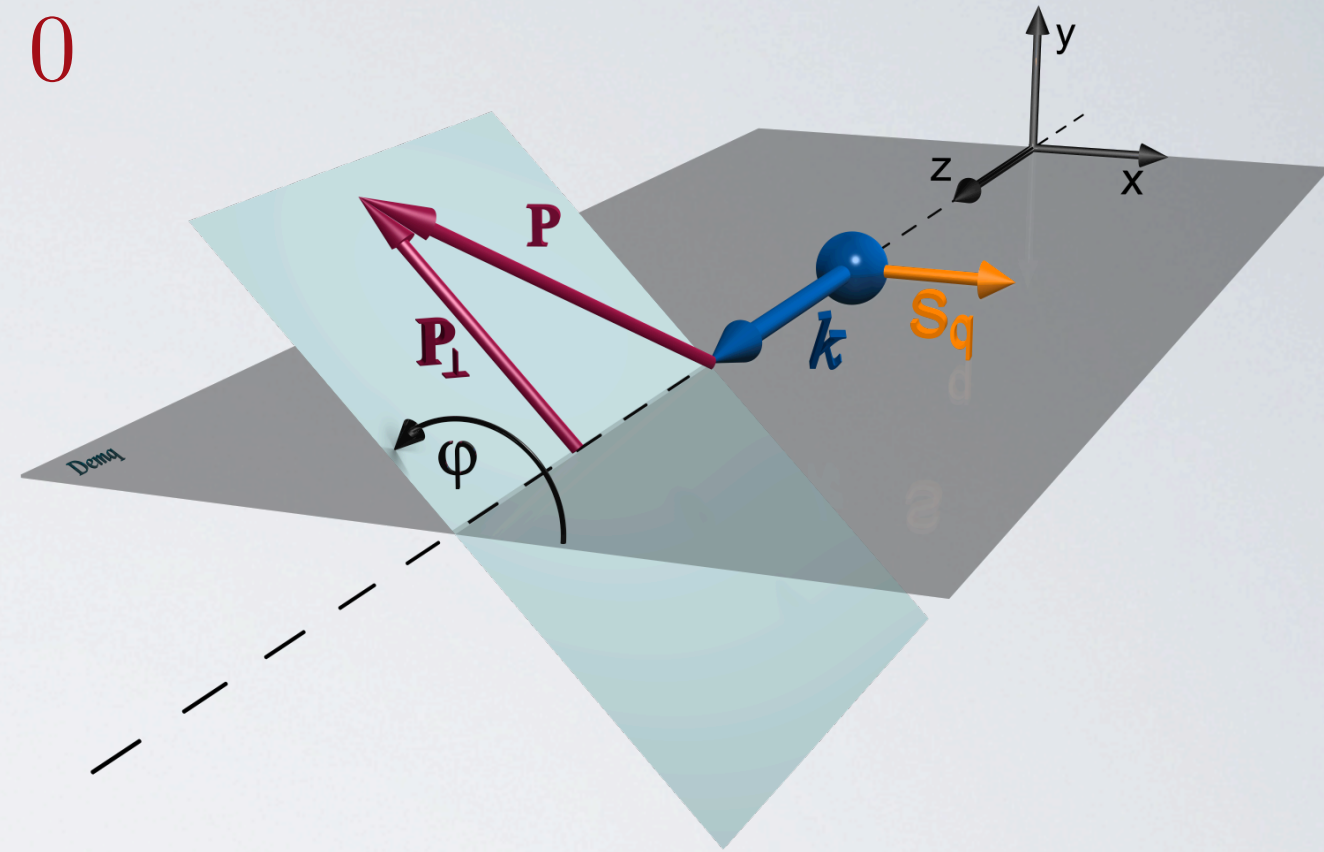
◆ **Transformation to frame  $\mathbf{k}_T = 0$**

$$k = (k^-, k^+, \mathbf{0})$$

$$\mathbf{k}_T = -\mathbf{P}_T / z_h$$

$$\mathbf{P}_T = \mathbf{P}_{h_1}^\perp + \mathbf{P}_{h_2}^\perp$$

$$\mathbf{R} = (\mathbf{P}_{h_1}^\perp - \mathbf{P}_{h_2}^\perp) / 2$$



◆ **Integrate over one or other momentum:**

$$D_{q^\uparrow}^{h_1 h_2}(\varphi_R) = D_{1,q}^{h_1 h_2} + \sin(\varphi_R - \varphi_S) \mathcal{F}[H_1^\triangleleft, H_1^\perp]$$

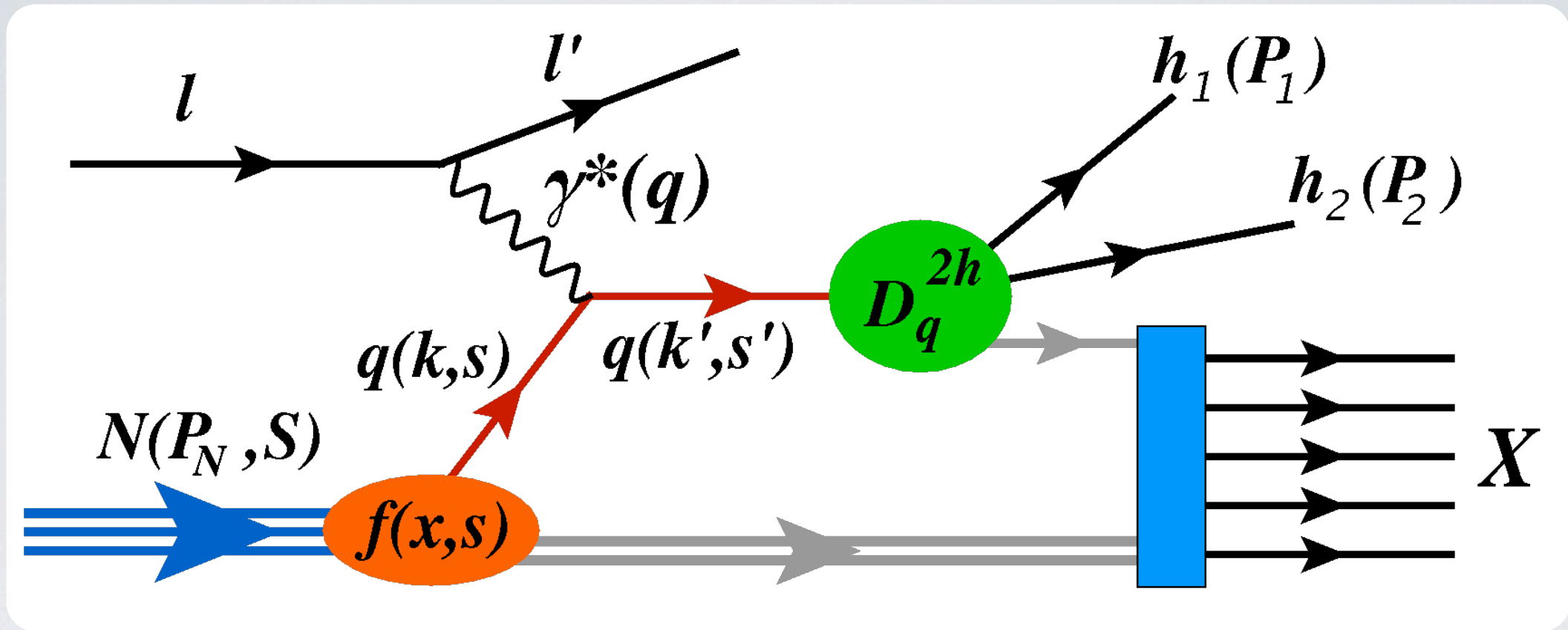
$$D_{q^\uparrow}^{h_1 h_2}(\varphi_T) = D_{1,q}^{h_1 h_2} + \sin(\varphi_T - \varphi_S) \mathcal{F}'[H_1^\triangleleft, H_1^\perp]$$

◆ **The IFF surviving after  $\mathbf{k}_T$  integration is redefined as**

**A. Bacchetta, M. Radici: PRD 69, 074026 (2004).**

$$H_1^\triangleleft(z_h, \xi, M_h^2) \equiv \int d^2 \mathbf{k}_T \left[ H_1^{\triangleleft \prime e}(z_h, \xi, M_h^2, k_T^2, \mathbf{k}_T \cdot \mathbf{R}_T) + \frac{k_T^2}{2M_h^2} H_1^{\perp o}(z_h, \xi, k_T^2, R_T^2, \mathbf{k}_T \cdot \mathbf{R}_T) \right]$$





# Sivers Effect in Two Hadron SIDIS

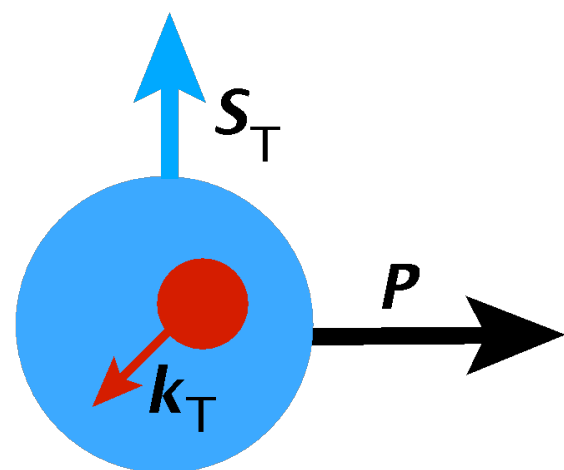


N/q	U	L	T
U	$f_1$		$h_1^\perp$
L		$g_{1L}$	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}^\perp$	$h_1 h_{1T}^\perp$

D. Sivers, Phys.Rev. D41 (1990).

- ◆ Correlation of  $k_T$  and  $S_T$
- ◆ Proposed by Dennis Sivers in 1990 to explain the single spin asymmetry in  $pp^\uparrow \rightarrow \pi + X$ .

$$S_T k_T \sin(\varphi_k - \varphi_S)$$

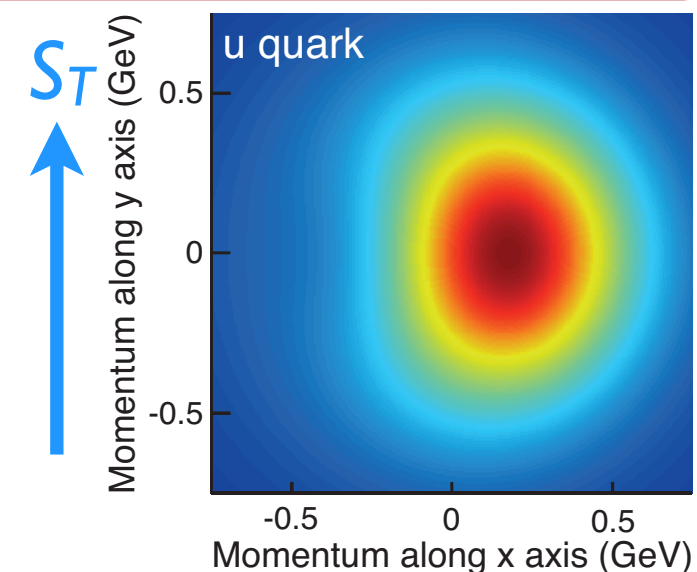


$$f_{\uparrow}^q(x, \vec{k}_T) = f_1^q(x, k_T) + \frac{[\vec{S} \times \vec{k}_T]_3}{M} f_{1T}^{\perp q}(x, k_T)$$

◆ Naively *T-odd*, gauge-link should be included in the definition.

◆ Accessible in Polarized SIDIS, Drell-Yan.

$$f_{1T}^{\perp SIDIS} = -f_{1T}^{\perp DY}$$



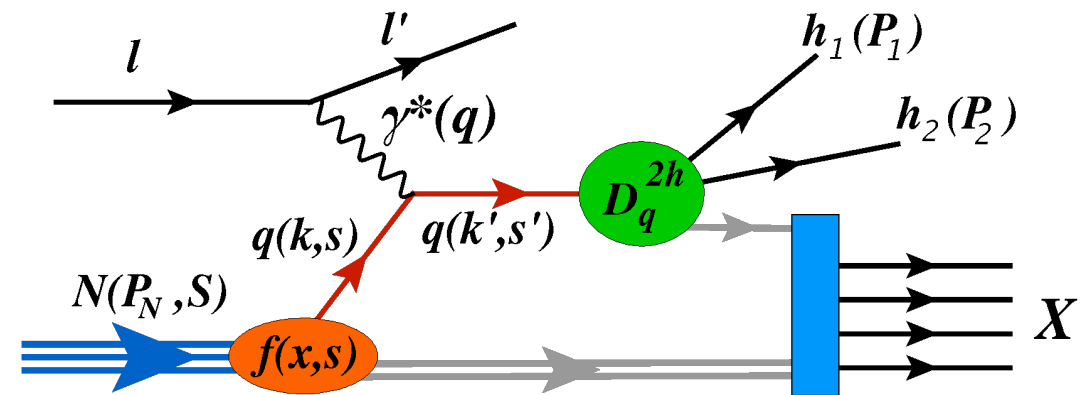
EIC White Paper, arXiv:1212.1701



# TWO-HADRON SIDIS

Kotzinian, H.M., Thomas: PRL.113, 062003 ; PRD.90, 074006 ; 1407.6572 (2014);

- Correlations of quark's TM transferred to **two hadrons**.



$$\frac{d\sigma^{h_1 h_2}}{dz_1 dz_2 d^2 P_{1T} d^2 P_{2T}} = C(x, Q^2) (\sigma_U + \sigma_S)$$

$$\sigma_U = \sum_q e_q^2 \int d^2 \mathbf{k}_T f_1^q D_{1q}^{h_1 h_2} \quad \sigma_S = \sum_q e_q^2 \int d^2 \mathbf{k}_T \frac{[\mathbf{S}_T \times \mathbf{k}_T]_3}{M} f_{1T}^{\perp q} D_{1q}^{h_1, h_2}$$

- Unpolarized fully unintegrated dihadron Fragmentation Function

◆ **Single hadron** FF.

$$D_{1q}^h(z, P_{\perp})$$

◆ **Dihadron** FF.

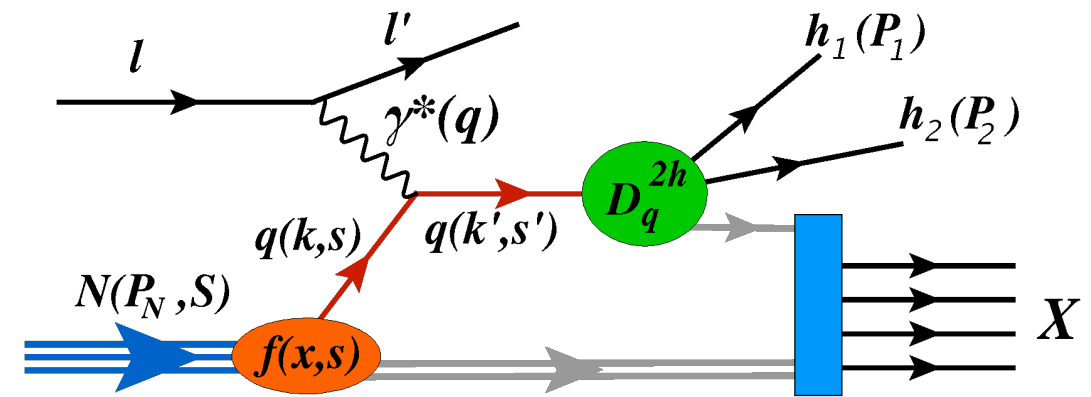
$$D_{1q}^{h_1, h_2}(z_1, z_2, P_{1\perp}, P_{2\perp}, P_{1\perp} \cdot P_{2\perp})$$



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- Unpolarized fully unintegrated dihadron Fragmentation Function

◆ **Single hadron** FF.

$$D_{1q}^h(z, P_{\perp})$$

◆ **Dihadron** FF.

$$D_{1q}^{h_1, h_2}(z_1, z_2, P_{1\perp}, P_{2\perp}, \mathbf{P}_{1\perp} \cdot \mathbf{P}_{2\perp})$$

two-hadron correlations

# TWO-HADRON SIDIS

## ► Cross Section in terms of **Total and Relative Momenta**

$$P_h = P_1 + P_2 \quad R = (P_1 - P_2)/2$$

## ► The Sivers term:

$$\sigma_S = S_T \left( \sigma_T \frac{P_{hT}}{M} \sin(\varphi_T - \varphi_S) + \sigma_R \frac{R_T}{M} \sin(\varphi_R - \varphi_S) \right)$$

$$\int d\varphi_R \sigma_S = S_T \left( \sigma_{T,0} \frac{P_{hT}}{M} + \sigma_{R,1} \frac{R}{2M} \right) \sin(\varphi_T - \varphi_S)$$

$$\int d\varphi_T \sigma_S = S_T \left( \sigma_{T,1} \frac{P_{hT}}{2M} + \sigma_{R,0} \frac{R}{M} \right) \sin(\varphi_R - \varphi_S)$$

## ► **Non-vanishing $\sigma_R$ is new! (shown explicitly in toy model)**

$$D_{1q}^{h_1, h_2}(z_1, z_2, P_{1\perp}, P_{2\perp}, \mathbf{P}_{1\perp} \cdot \mathbf{P}_{2\perp}) = D_{1q}^{h_1, h_2}(z_1, z_2) \frac{1}{\pi^2 \nu_1^2 \nu_2^2} e^{-P_{1\perp}^2/\nu_1^2 - P_{2\perp}^2/\nu_2^2} (1 + c \mathbf{P}_{1\perp} \cdot \mathbf{P}_{2\perp})$$

## ► **Is this method useful?**

- ❖ Additional information on Sivers PDF: **Flavor decomposition!**
- ❖ Need the new (yet unknown) DiFFs!
- ❖ Is it feasible to measure these SSAs?

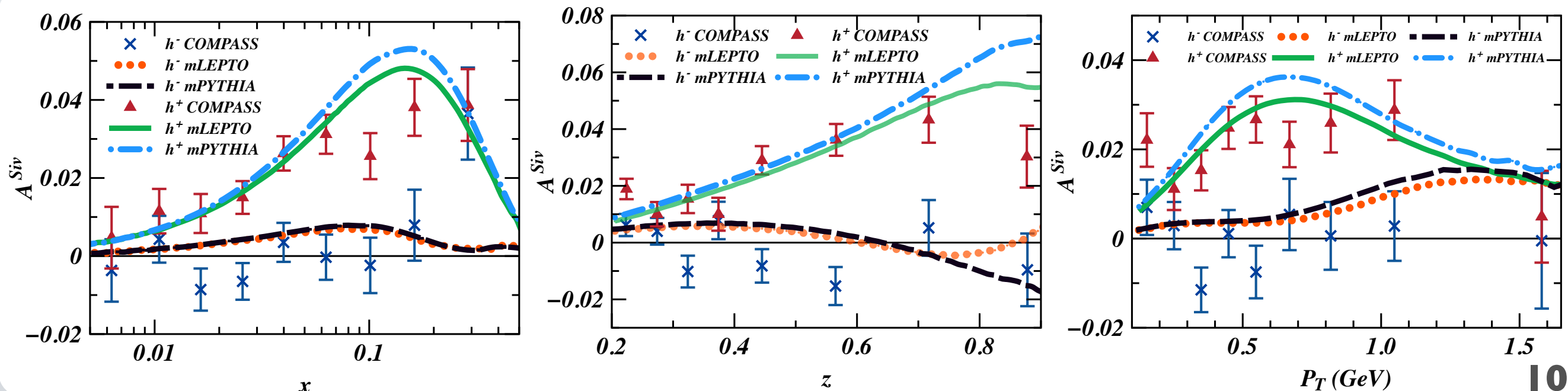


# EVENT GENERATORS + SIVERS EFFECT

Kotzinian, H.M., Thomas: PRL.113, 062003 ; PRD.90, 074006 ; 1407.6572 (2014);

- ▶ Use **PYTHIA 6.4** (and **LEPTO** earlier) (*F77-yuk*).
- ▶ Incorporate dynamical hadronization mechanism: one, two,... hadron FFs.
- ▶ Sivers effect modulates quark TM's azimuthal angle: *relatively easy* to include in MC generators.
- ▶ Use Sivers PDF extraction from *Torino group*.
- ▶ Event generators allow to study *exp. kinematics effects*.

## ❖ Does it working?

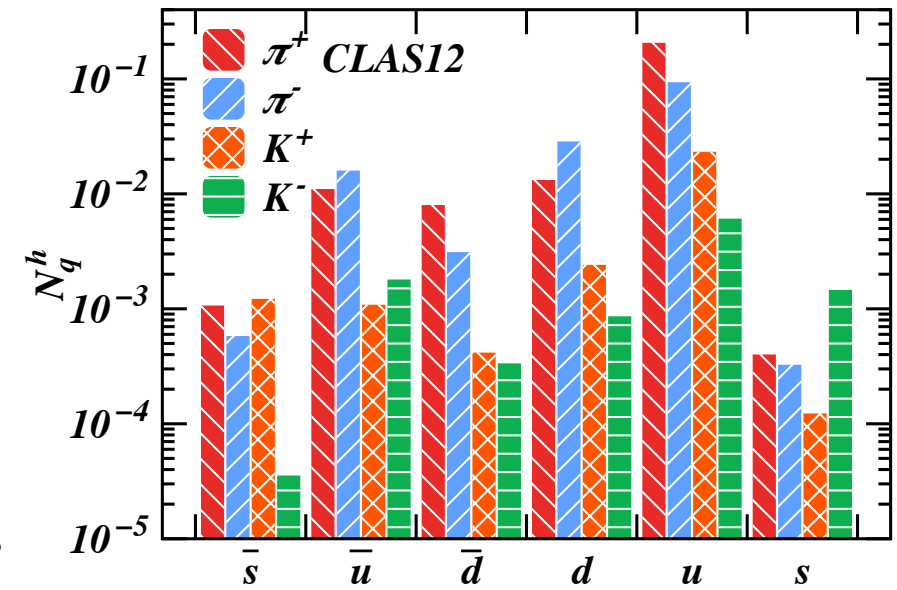
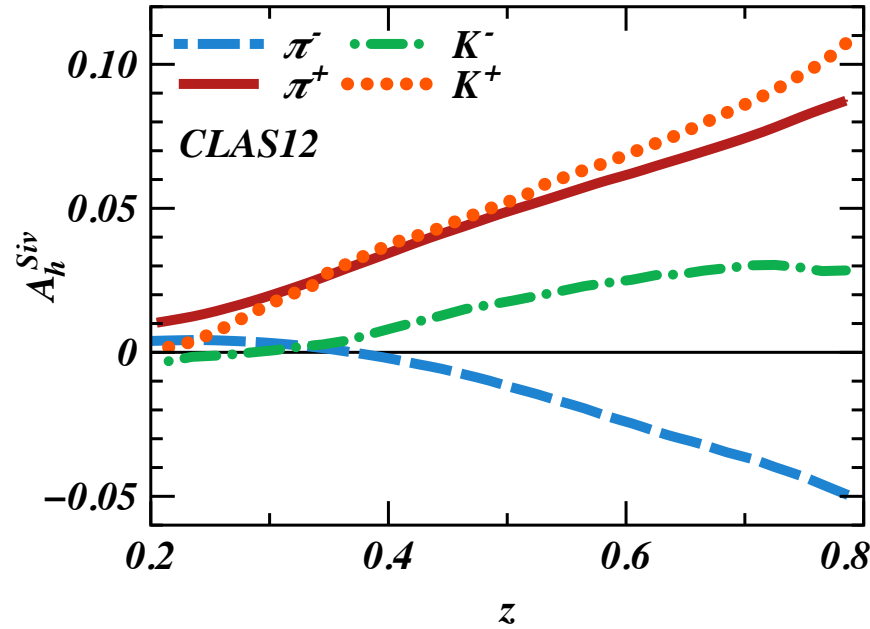
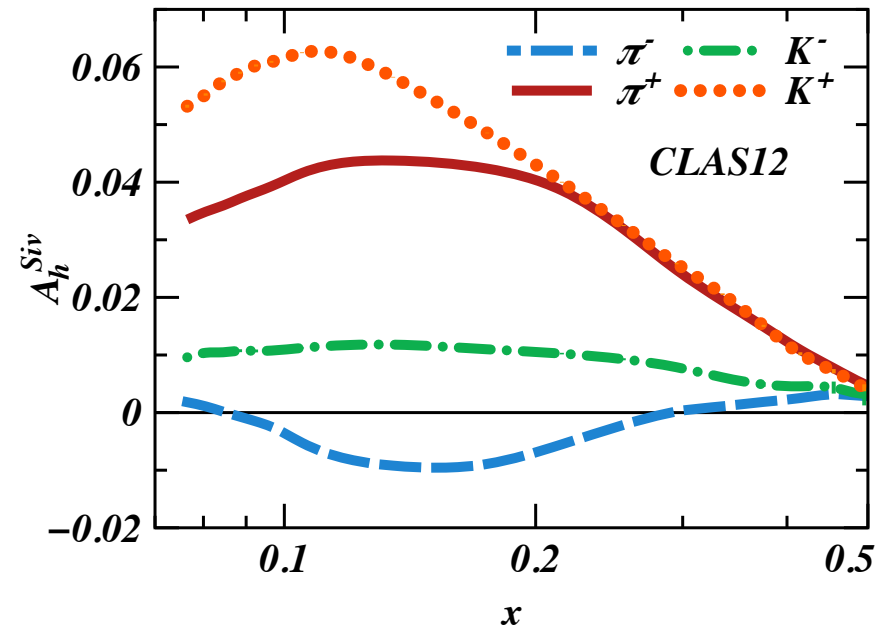


# Sivers SSAs at CLAS12

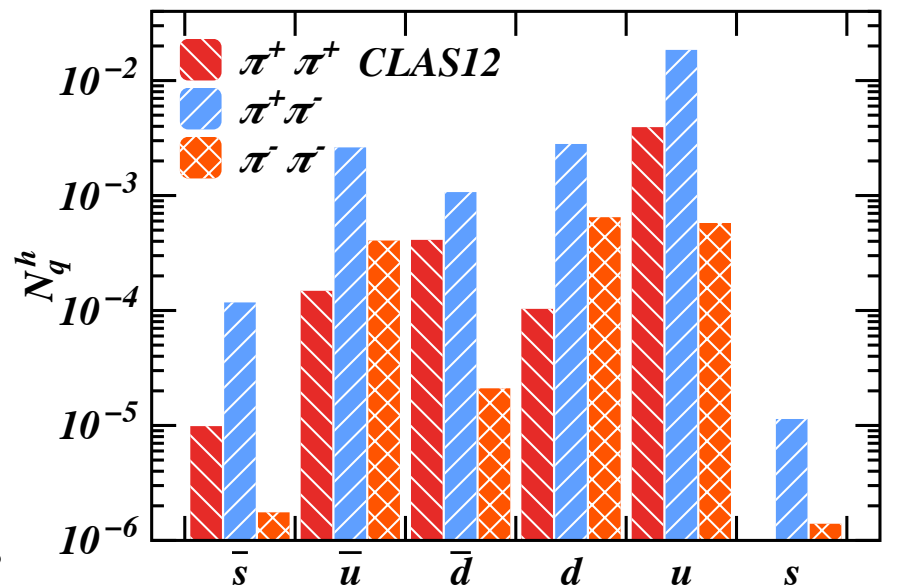
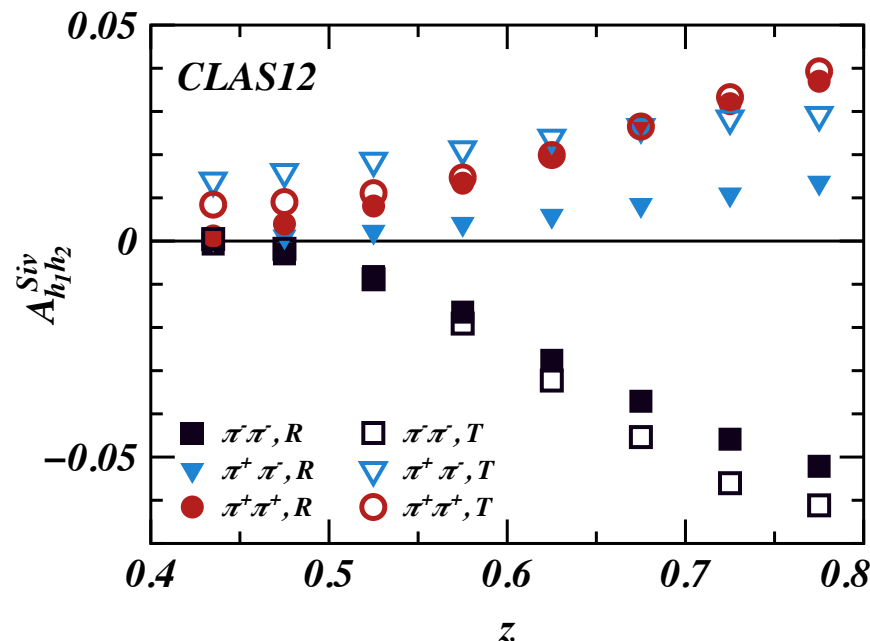
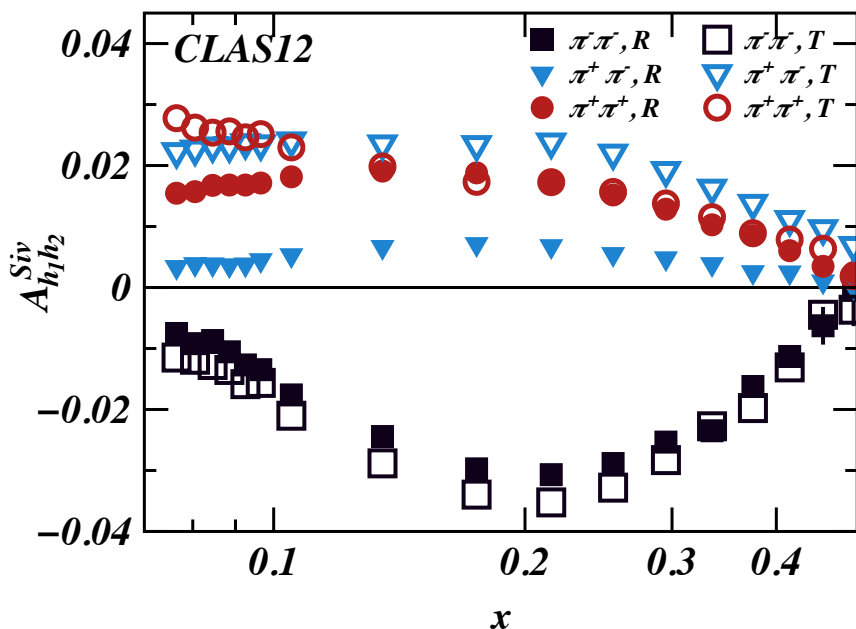
H.M et al., arXiv:1502.02669 (2015).

❖ Exploring the large  $x$  region.

◆ **Single hadron** SSAs.



◆ **Dihadron** SSAs for pion pairs: identical pairs via z-ordering  $z_1 \geq z_2$



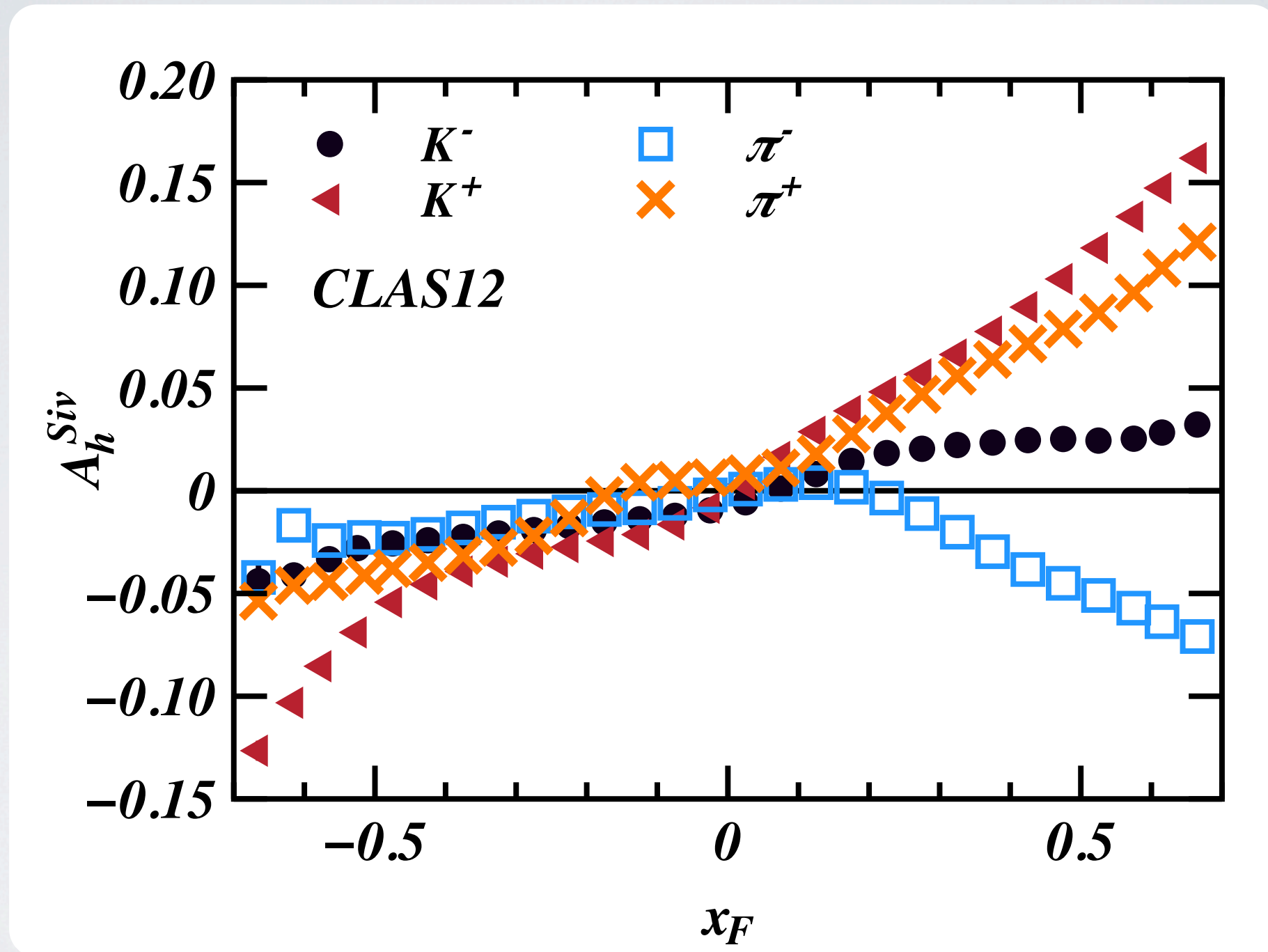
◆ Both Single and Dihadron SSAs are comparable in size!



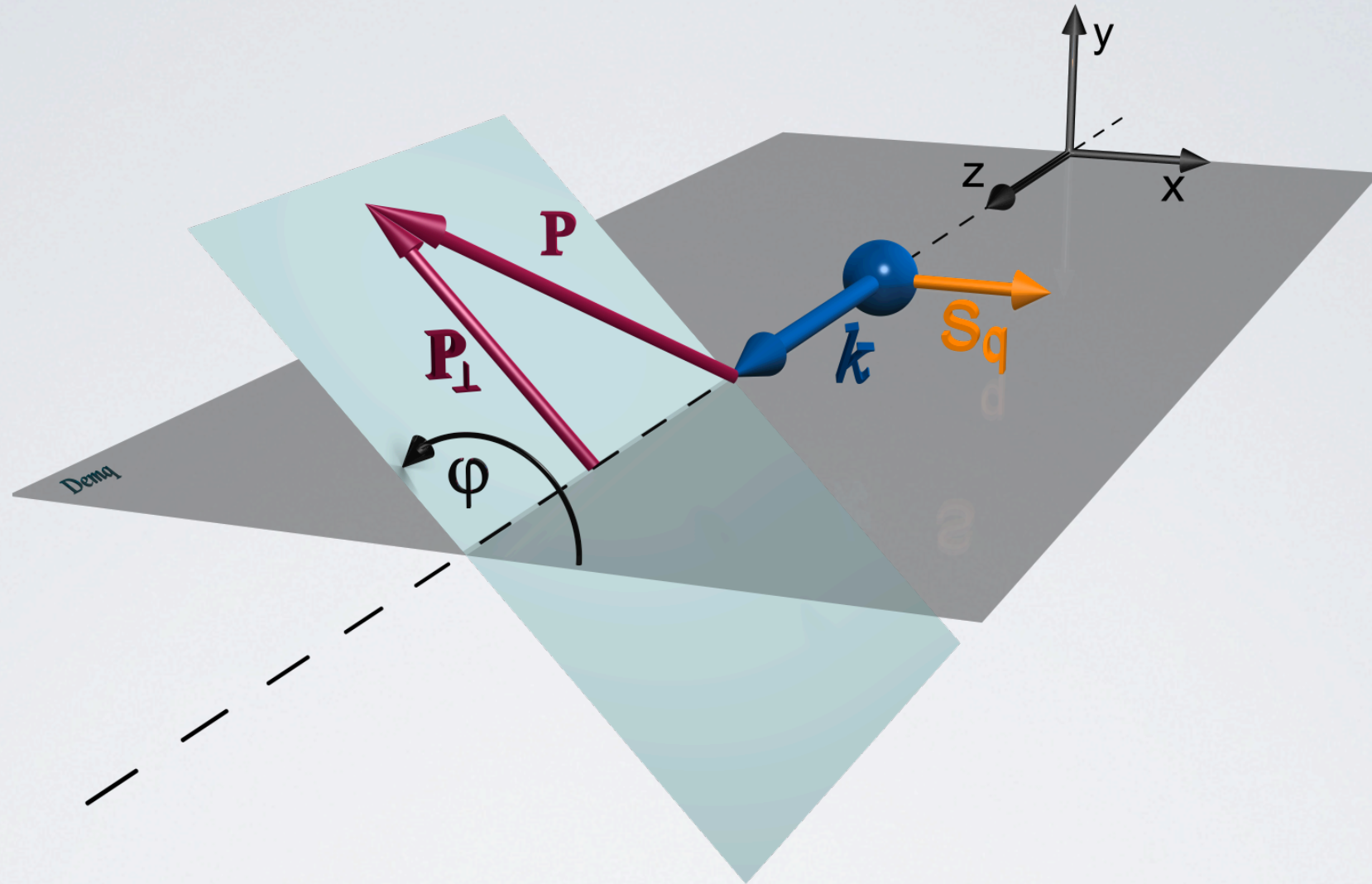
# Sivers SSAs at CLAS12

H.M et al., arXiv:1502.02669 (2015), accepted in PRD.

► Explore **Target Fragmentation Regions**  $x_F < 0$ .



► Sivers SSA changes sign in some channels, fragmentation of nucleon **remnant (recoil TM)**!



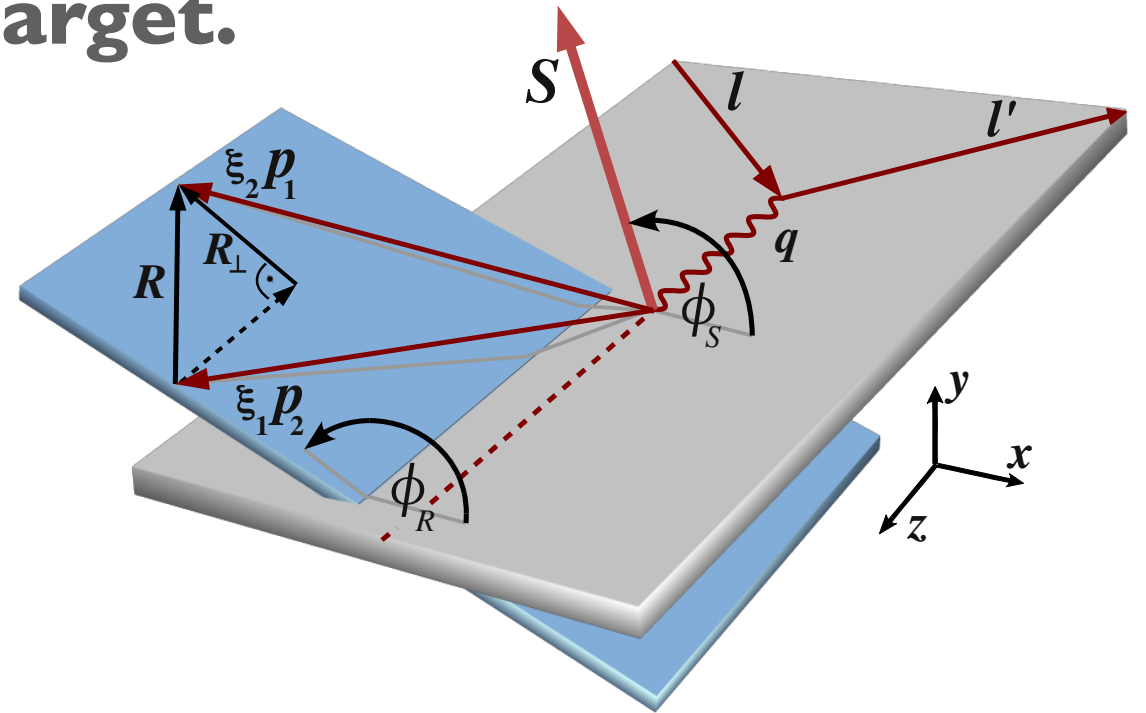
***TRANSVERSELY POLARIZED QUARK FRAGMENTATION:  
COLLINS EFFECT AND TWO-HADRON CORRELATIONS***



# RECENT COMPASS RESULTS

COMPASS, PLB736, 124-131 (2014).

◆ **SIDIS with transversely polarized target.**



◆ **Collins single spin asymmetry:**

$$A_{Coll} = \frac{\sum_q e_q^2 \Delta_T q \otimes H_1^{\perp h/q}}{\sum_q e_q^2 q \otimes D_1^{h/q}}$$

◆ **Two hadron single spin asymmetry:**

$$A_{UT}^{\sin \phi_{RS}} = \frac{|\mathbf{p}_1 - \mathbf{p}_2|}{2M_{h^+h^-}} \frac{\sum_q e_q^2 \cdot h_1^q(x) \cdot H_{1,q}^{\triangleleft}(z, M_{h^+h^-}^2, \cos \theta)}{\sum_q e_q^2 \cdot f_1^q(x) \cdot D_{1,q}(z, M_{h^+h^-}^2, \cos \theta)}$$

◆ **Note the choice of the vector**

$$\mathbf{R}_{Artru} = \frac{z_2 \mathbf{P}_1 - z_1 \mathbf{P}_2}{z_1 + z_2}$$

# RECENT COMPASS RESULTS

COMPASS, PLB736, 124-131 (2014).

◆ **SIDIS with transversely polarized target.**

◆ **Collins single spin**

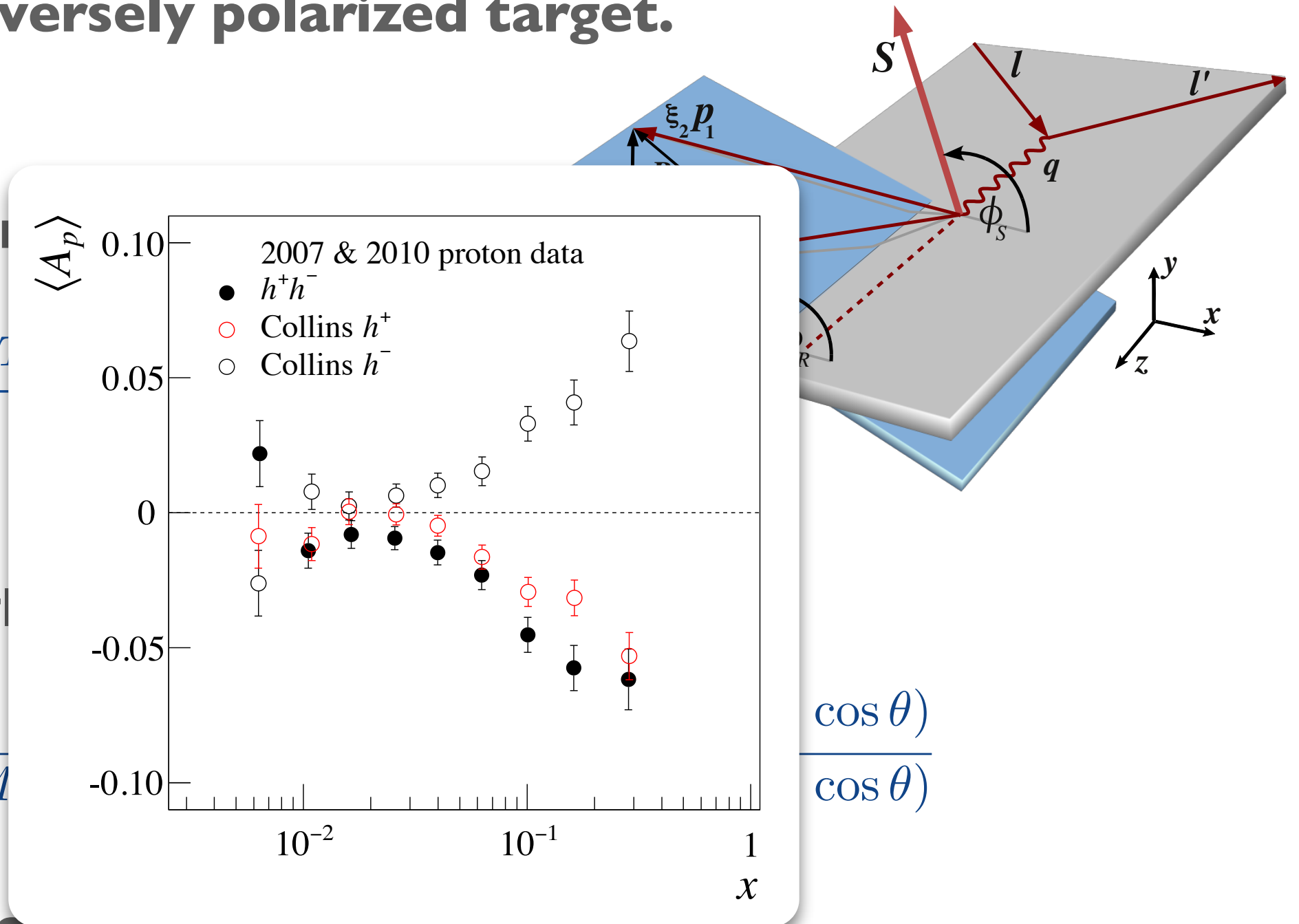
$$A_{Coll} = \frac{\sum_q e_q^2 \Delta q}{\sum_q e_q^2}$$

◆ **Two hadron single**

$$A_{UT}^{\sin \phi_{RS}} = \frac{|\mathbf{p}_1}{2M}$$

◆ **Note the choice of the vector**

$$\mathbf{R}_{Artru} = \frac{z_2 \mathbf{P}_1 - z_1 \mathbf{P}_2}{z_1 + z_2}$$



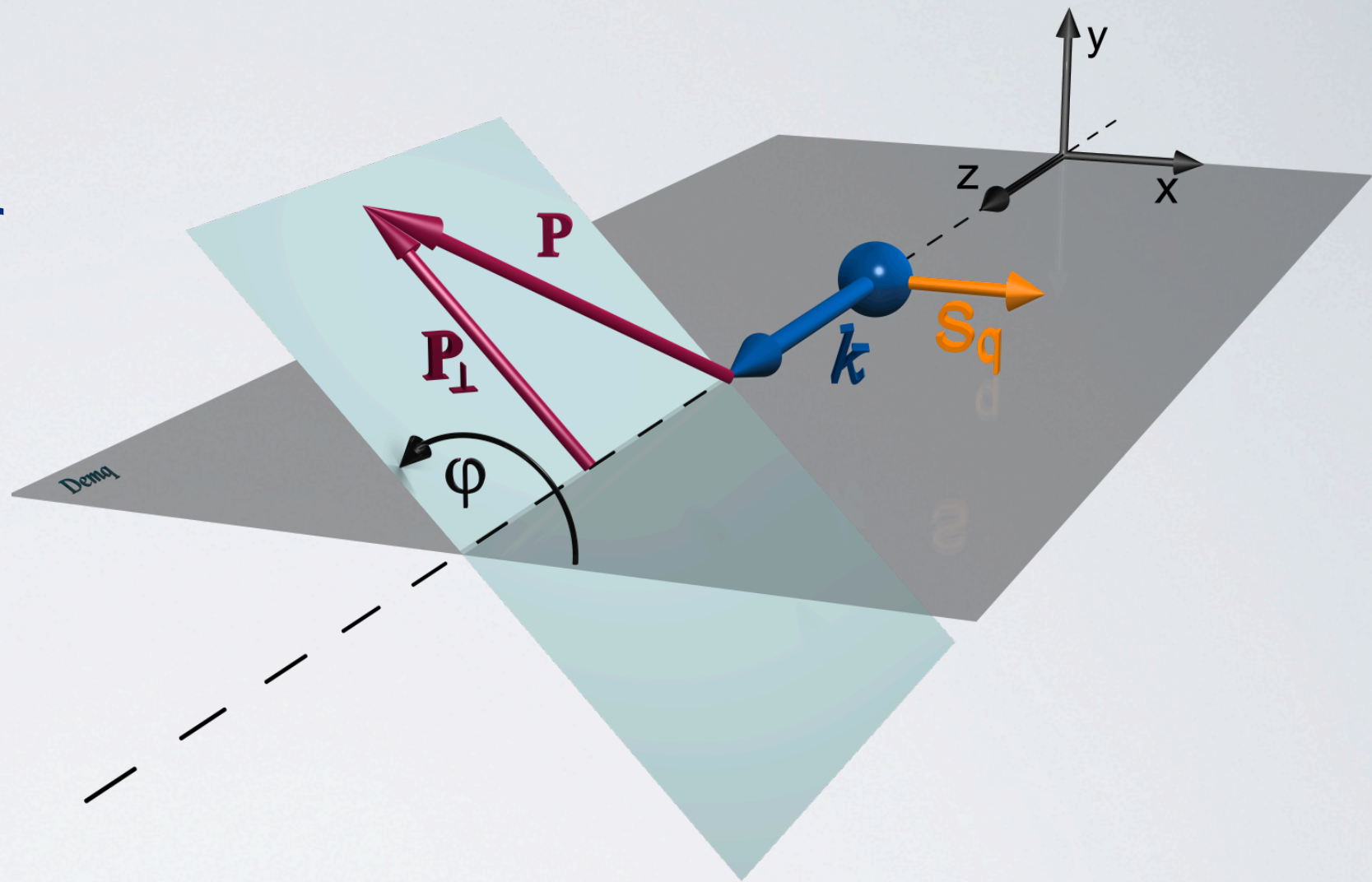
$$\frac{\cos \theta)}{\cos \theta)}$$



# COLLINS FRAGMENTATION FUNCTION

- **Collins Effect:**

Azimuthal Modulation of Transversely Polarized Quark' Fragmentation Function.



**Unpolarized**

$$D_{h/q^\uparrow}(z, P_\perp^2, \varphi) = D_1^{h/q}(z, P_\perp^2) - H_1^{\perp h/q}(z, P_\perp^2) \frac{P_\perp S_q}{zm_h} \sin(\varphi)$$

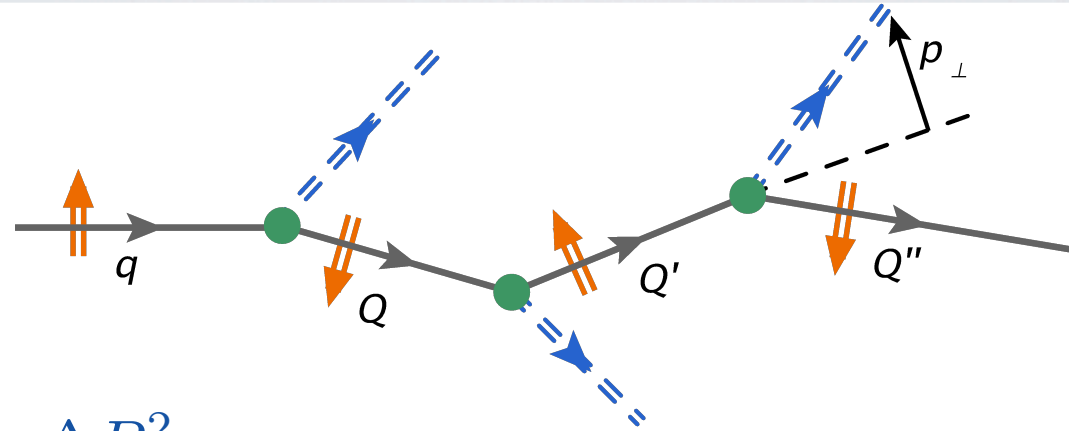
**Collins**

- **Chiral-ODD:** Needs to be coupled with another chiral-odd quantity to be observed.

# COLLINS FRAGMENTATION FUNCTION FROM NJL-JET

**H.M.,Bentz, Thomas, PRD.86:034025, 2012.**

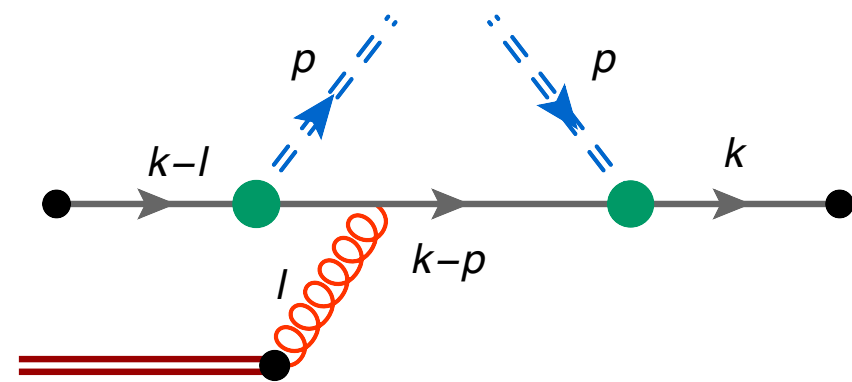
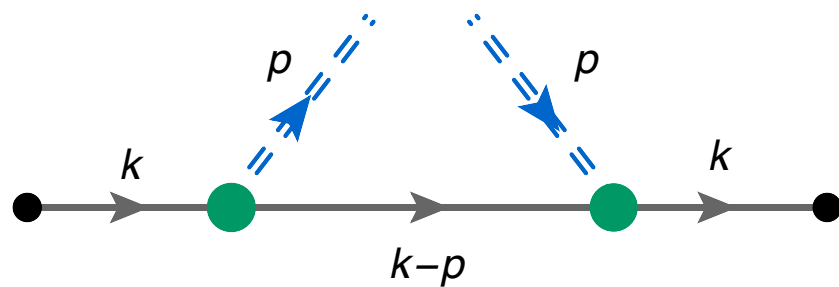
- **Extend the NJL-jet Model to Include the Quark's Spins.**



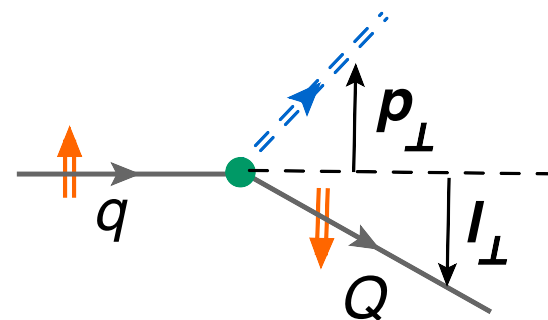
$$D_{h/q^\uparrow}(z, P_\perp^2, \varphi) \Delta z \frac{\Delta P_\perp^2}{2} \Delta\varphi = \left\langle N_{q^\uparrow}^h(z, z + \Delta z; P_\perp^2, P_\perp^2 + \Delta P_\perp^2; \varphi, \varphi + \Delta\varphi) \right\rangle$$

- **Model Calculated Elementary Collins Function as Input**

**A. Bacchetta et. al., PLB659, 234 (2008).**



- **Spin flip probability:**  $\mathcal{P}_{SF}$

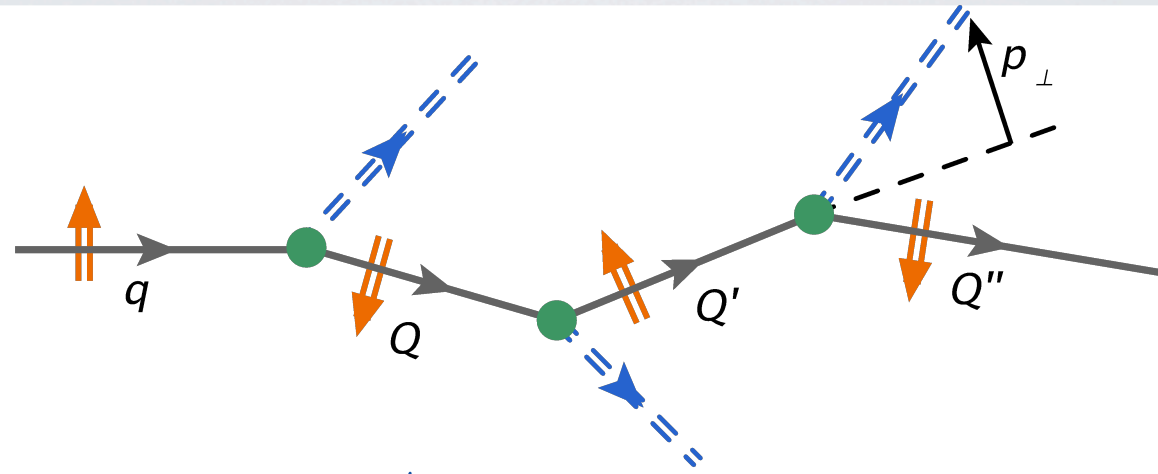




# POLARIZED QUARK DIFF IN QUARK-JET.

**H.M., Kotzinian, Thomas, PLB731 208-216 (2014).**

- Use the NJL-jet Model including Collins effect (Mk 2) to study DiFFs.



$$D_{q\uparrow}^{h_1 h_2}(z, M_h^2, \varphi_R) \Delta z \Delta M_h^2 \Delta \varphi_R = \left\langle N_{q\uparrow}^{h_1 h_2}(z, z + \Delta z; M_h^2, M_h^2 + \Delta M_h^2; \varphi_R, \varphi_R + \Delta \varphi_R) \right\rangle.$$

- Choose a constant Spin flip probability:  $\mathcal{P}_{SF}$

- Simple model to start with:

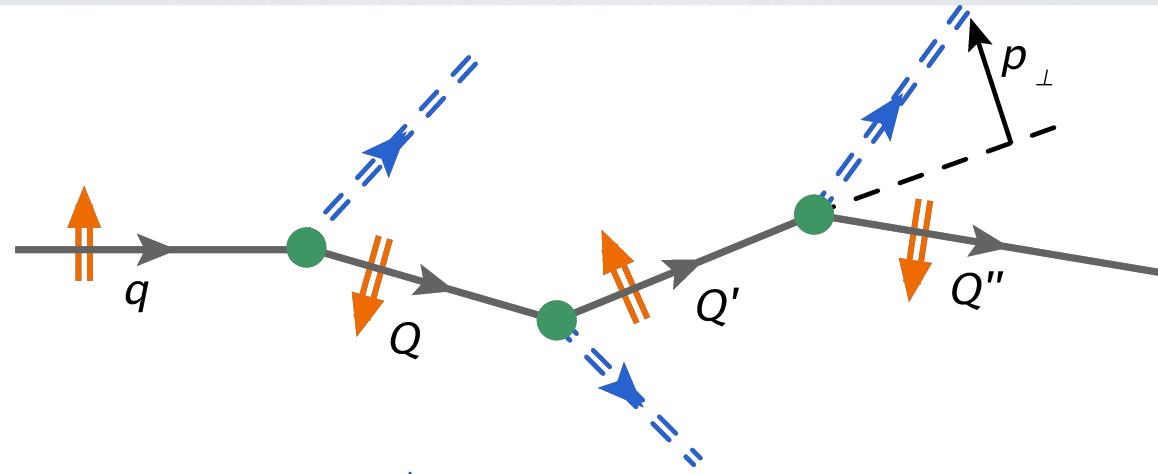
Only pions and extreme ansatz for the Collins term in elementary function.

$$d_{h/q\uparrow}(z, \mathbf{p}_\perp) = d_1^{h/q}(z, p_\perp^2)(1 - 0.9 \sin \varphi)$$

# POLARIZED QUARK DIFF IN QUARK-JET.

**H.M., Kotzinian, Thomas, PLB73I 208-216 (2014).**

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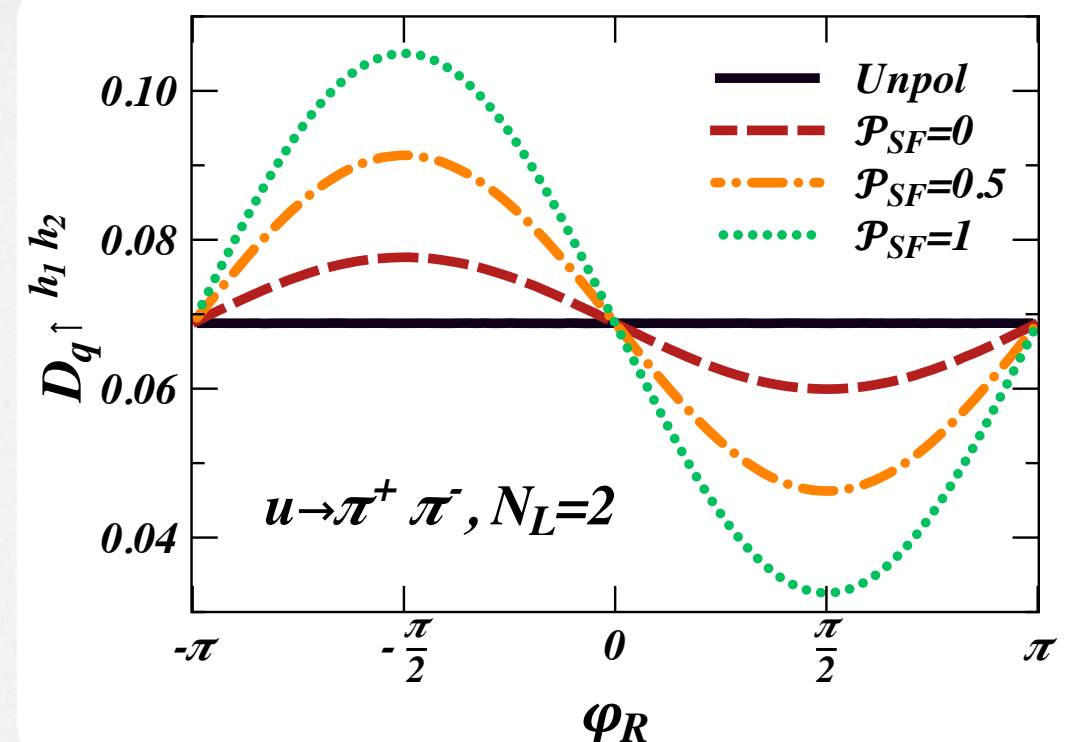
$$D_{q\uparrow}^{h_1 h_2}(z, M_h^2, \varphi_R) \Delta z \Delta M_h^2 \Delta \varphi_R = \left\langle N_{q\uparrow}^{h_1 h_2}(z, z + \Delta z; M_h^2, M_h^2 + \Delta M_h^2; \varphi_R, \varphi_R + \Delta \varphi_R) \right\rangle.$$

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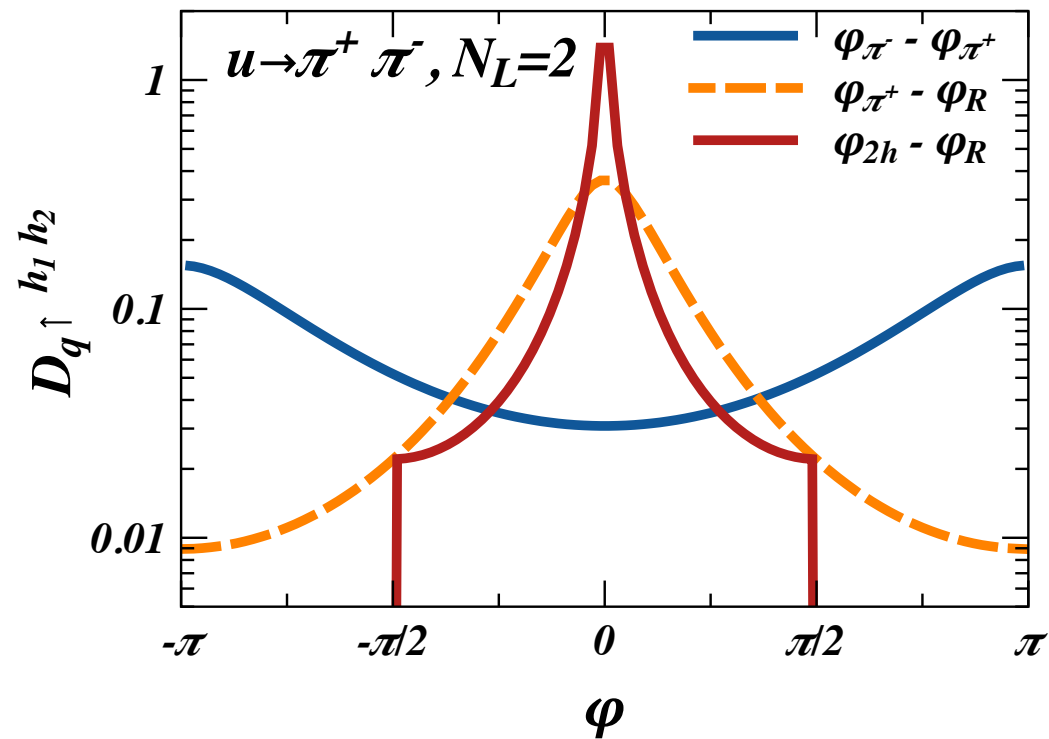
$$d_{h/q\uparrow}(z, \mathbf{p}_\perp) = d_1^{h/q}(z, p_\perp^2)(1 - 0.9 \sin \varphi)$$





# ANGULAR CORRELATIONS: $u \rightarrow \pi^+ \pi^-$

## Quark-Jet



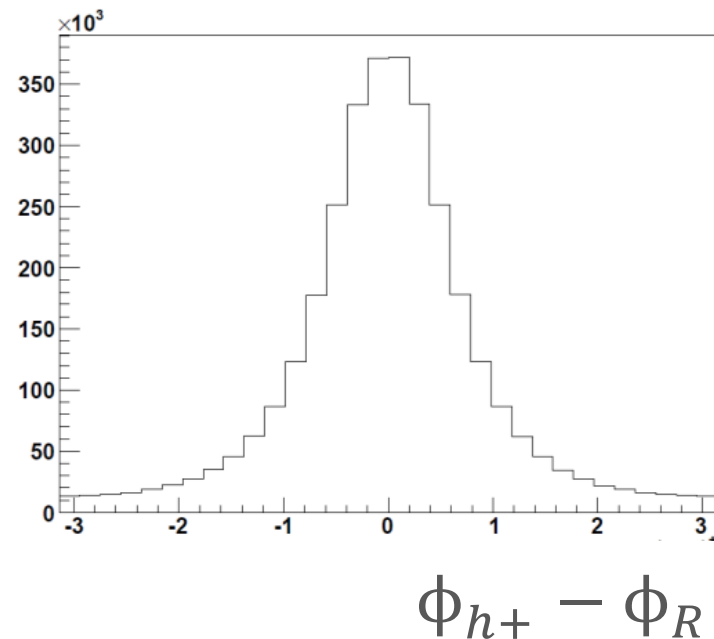
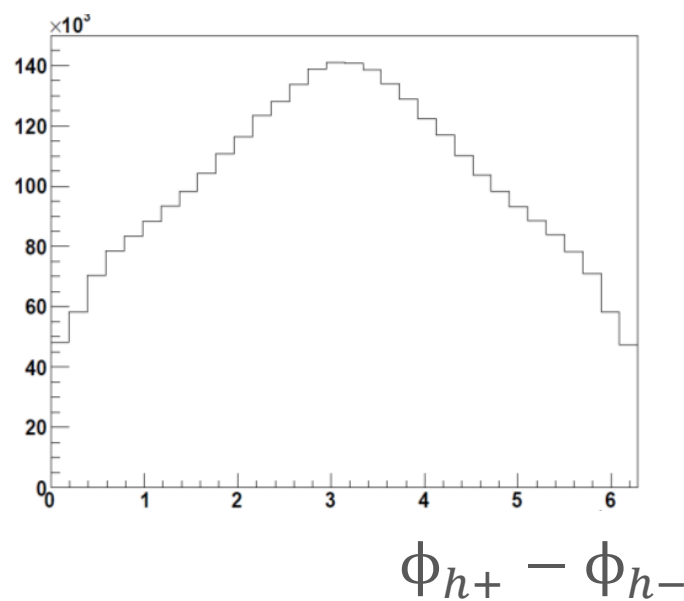
◆ We define:

$$\mathbf{P}_{2h} = \frac{\hat{\mathbf{P}}_1 - \hat{\mathbf{P}}_2}{2}$$

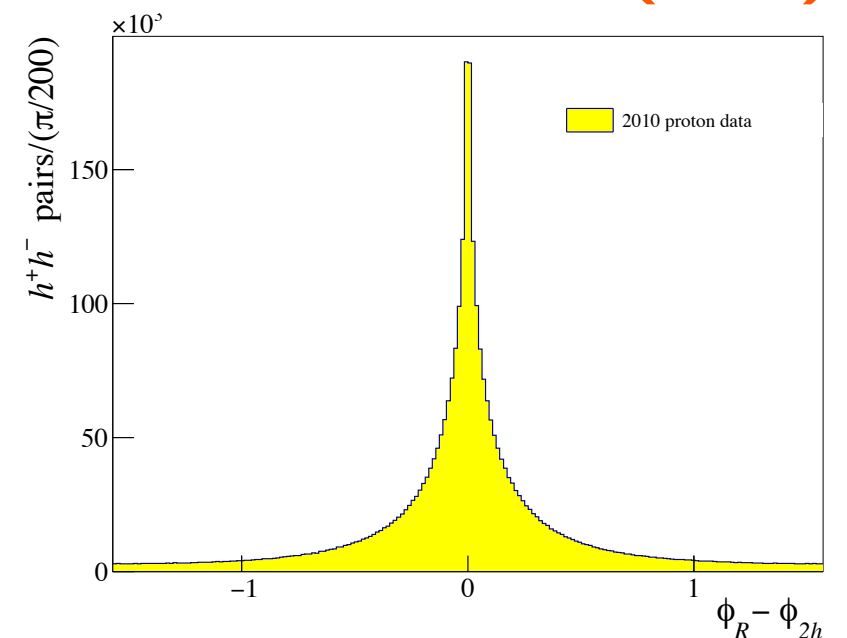
◆ No Spin Dependence Included!

## COMPASS Results

F. Bradamante - COMO 2013.

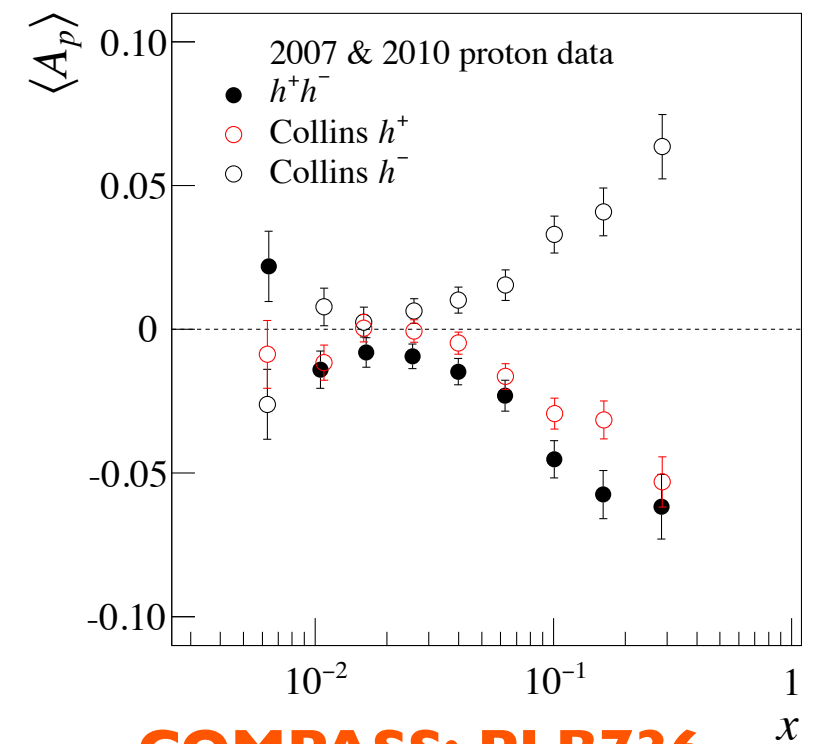
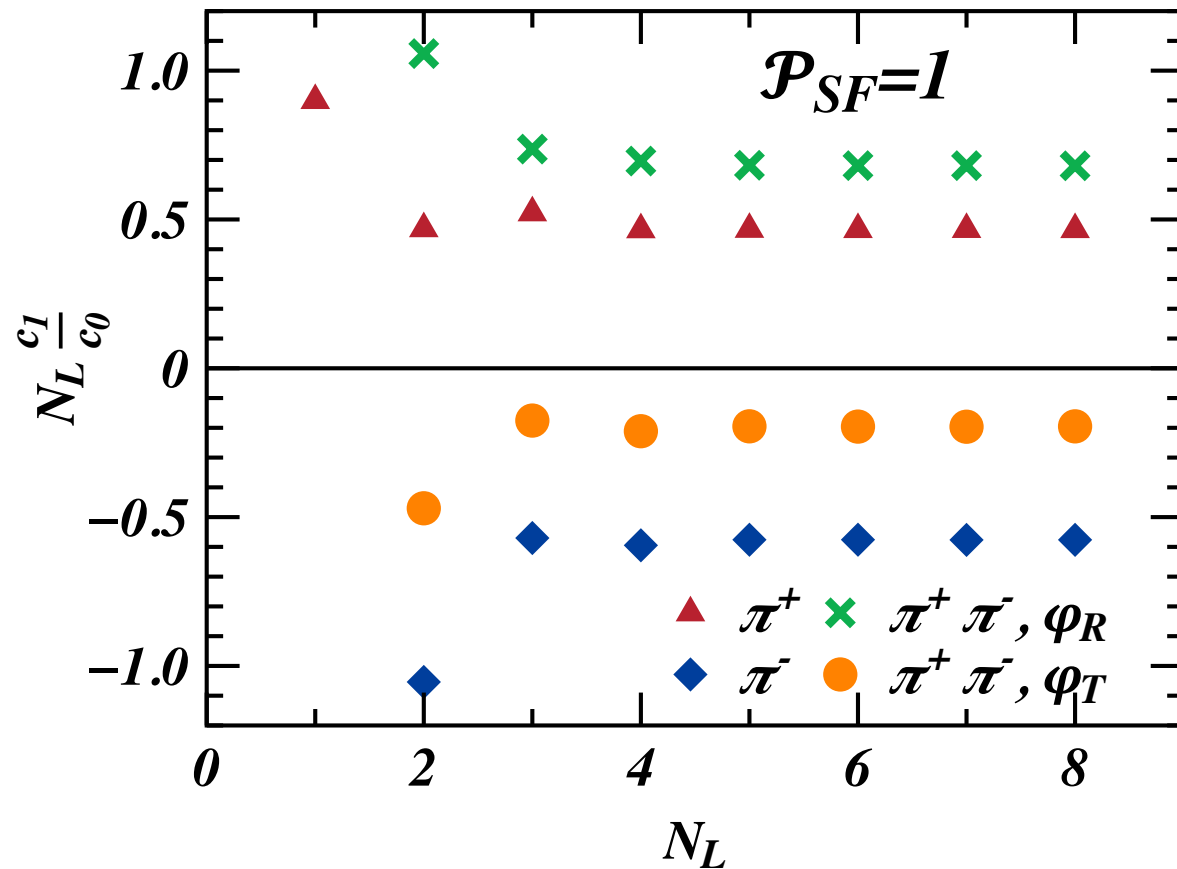
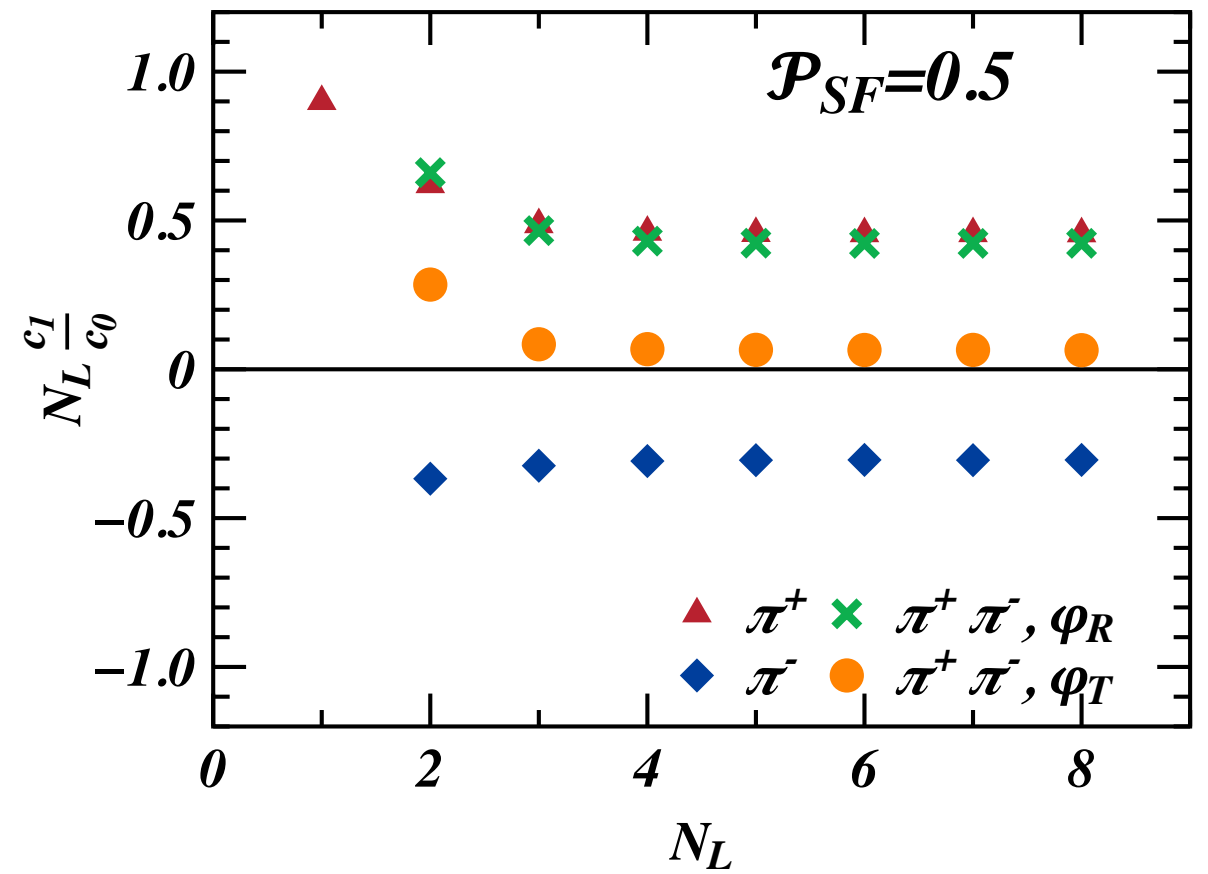
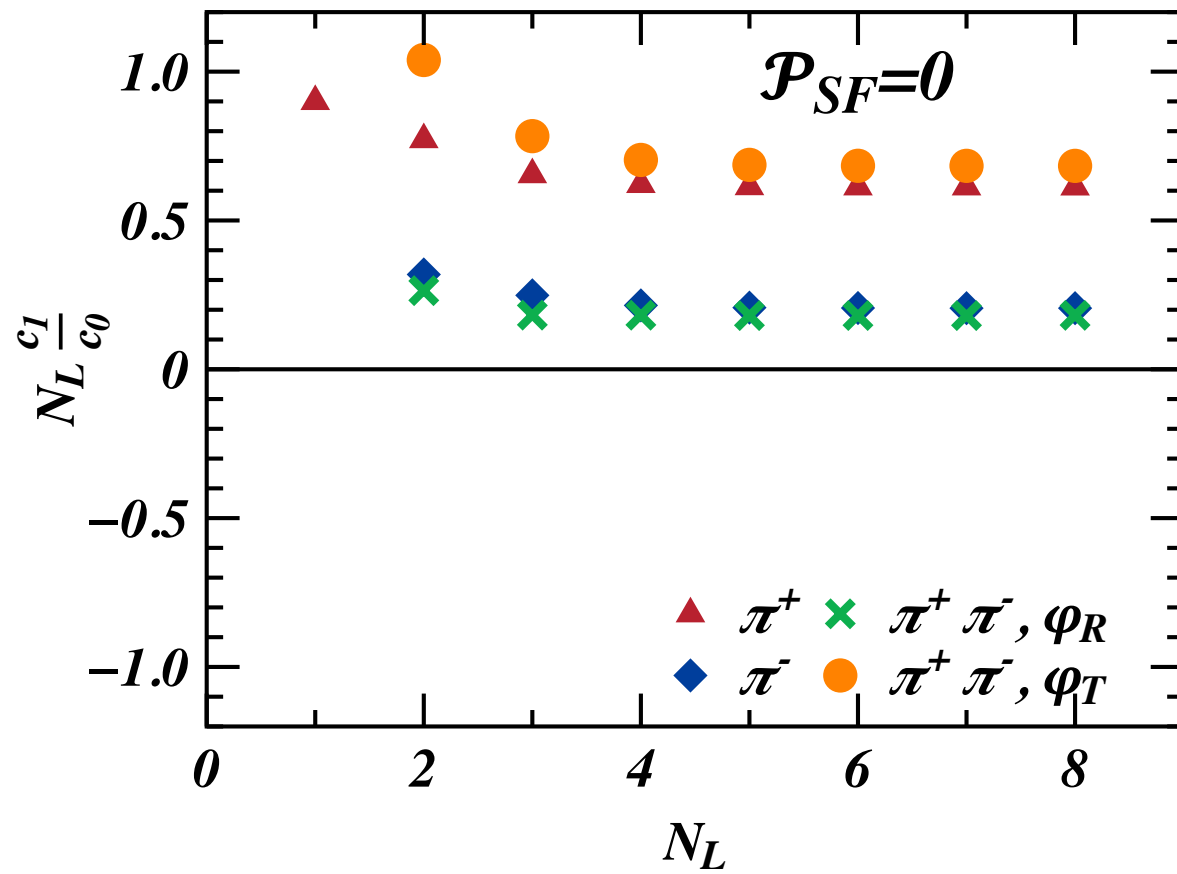


PLB736, 124-131 (2014).



# INTEGRATED ANALYZING POWERS

$$z_{1,2} > 0.2, z > 0.2$$



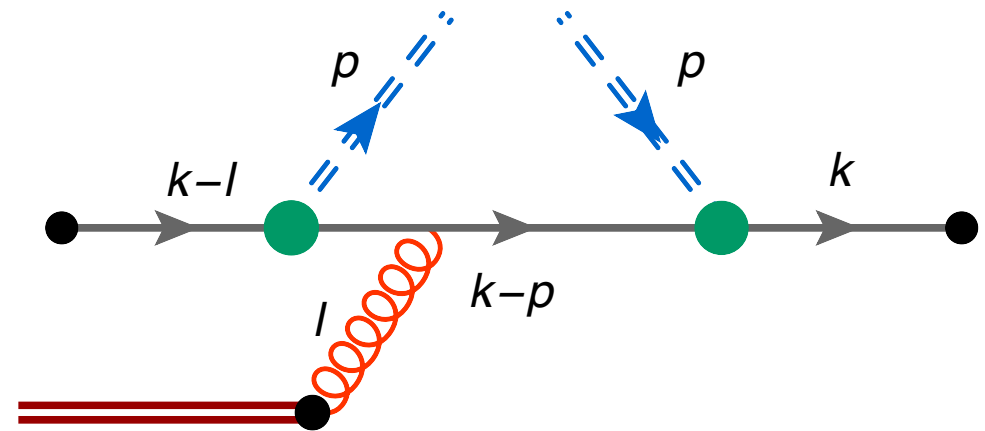
**COMPASS: PLB736,  
124-131 (2014).**



# IMPROVED MODEL FOR COLLINS EFFECT

- ◆ Use the ***spectator model*** for Collins function.

$$H_1^{\perp h/q}(z, P_{\perp}^2) \frac{P_{\perp} S_q}{z m_h} \sin(\varphi)$$



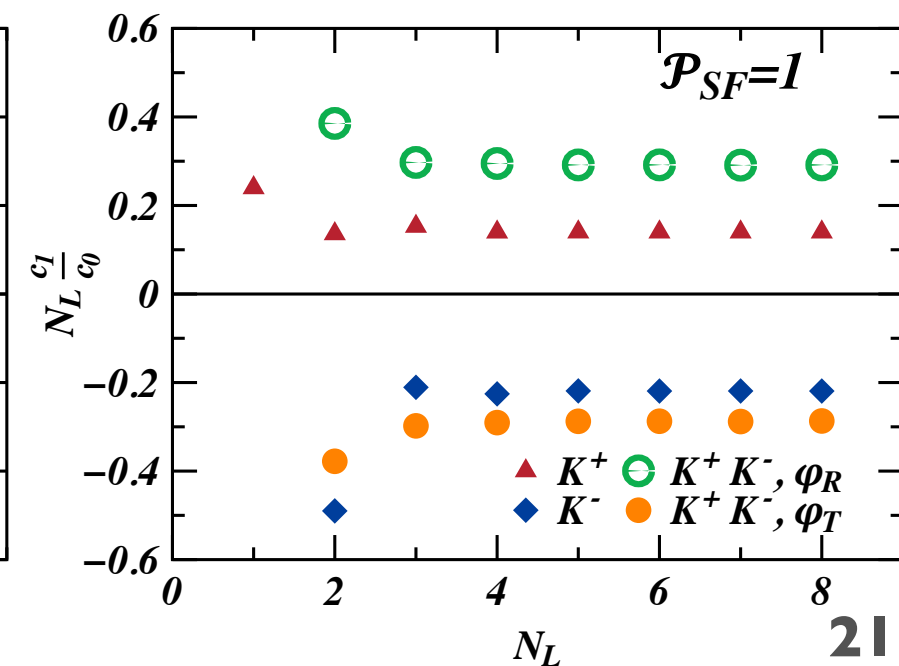
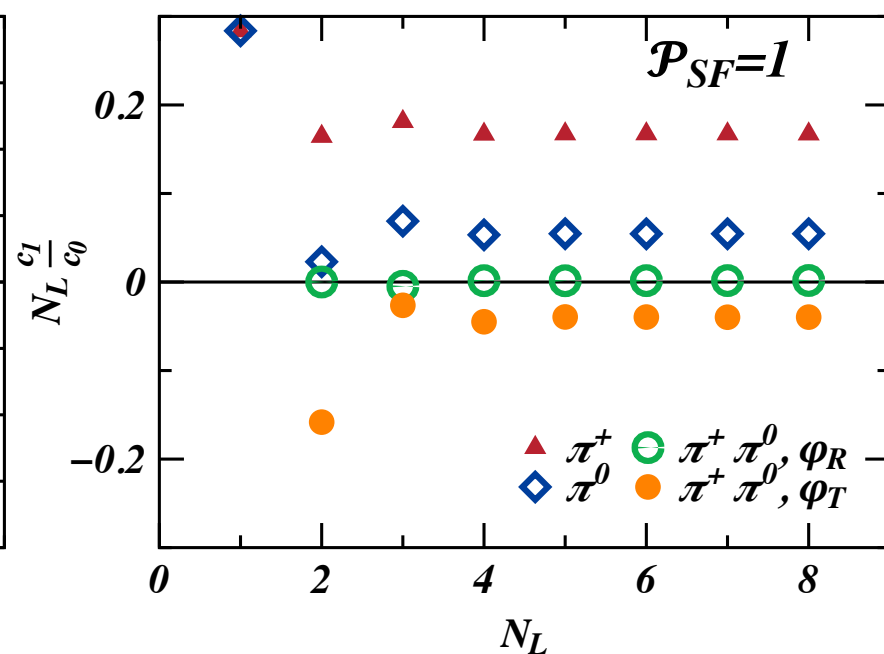
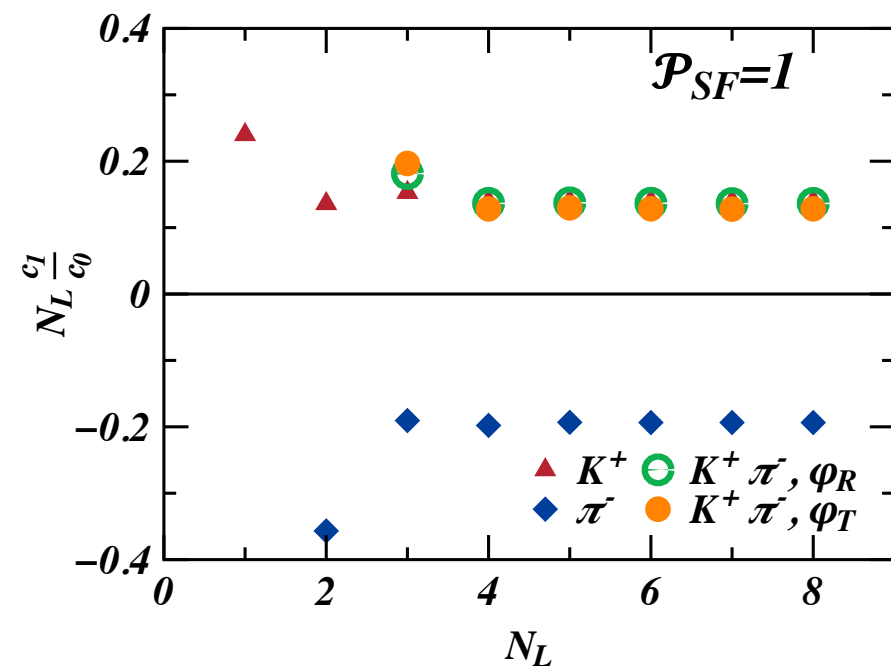
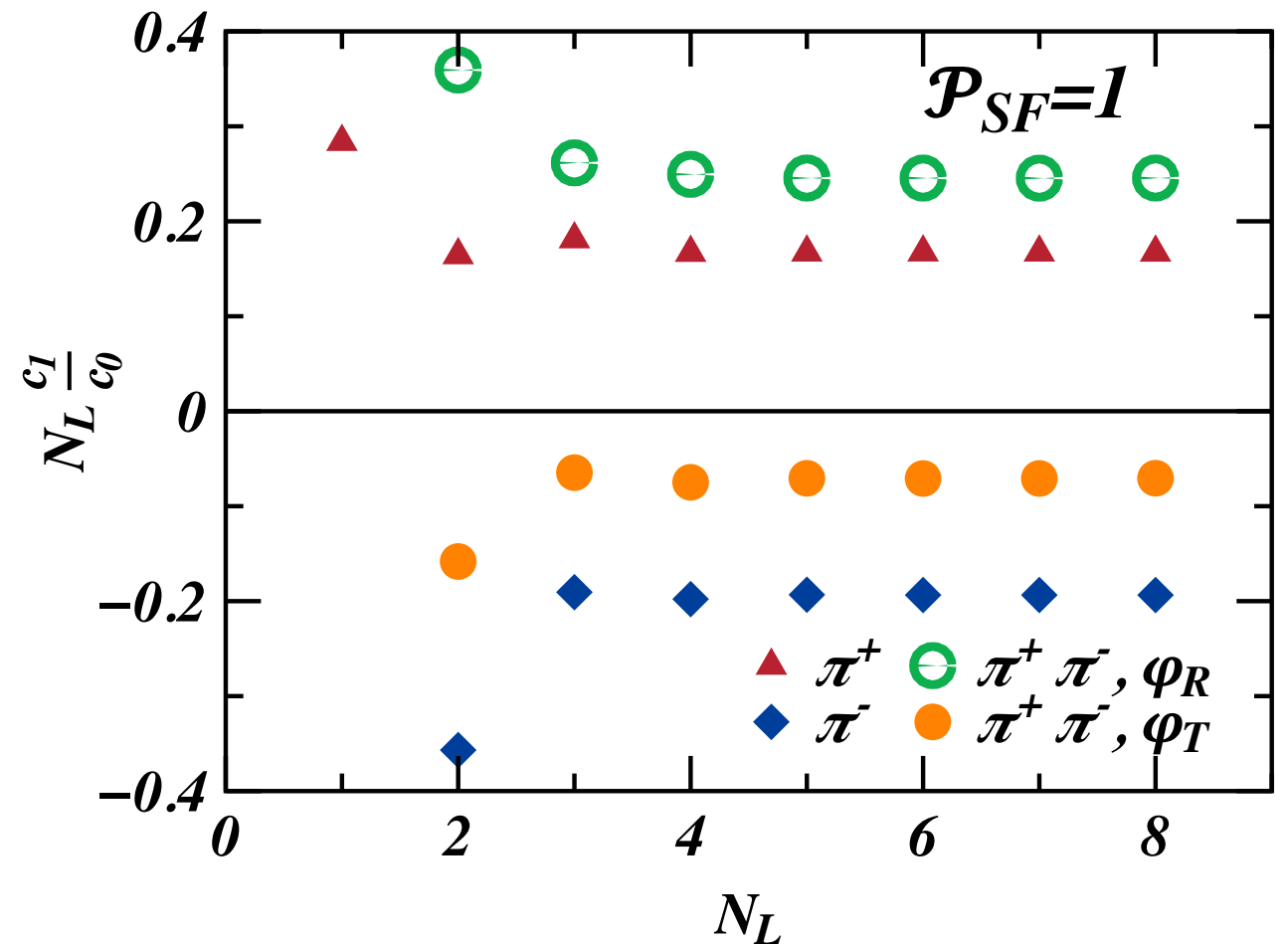
- ◆ Include both ***pion and kaon*** channels.

# IMPROVED MODEL RESULTS

◆ Qualitatively the **same** picture as in the Toy-model calculations.

◆ **Consistent** with COMPASS results.

◆ Predictions for various hadron pairs.





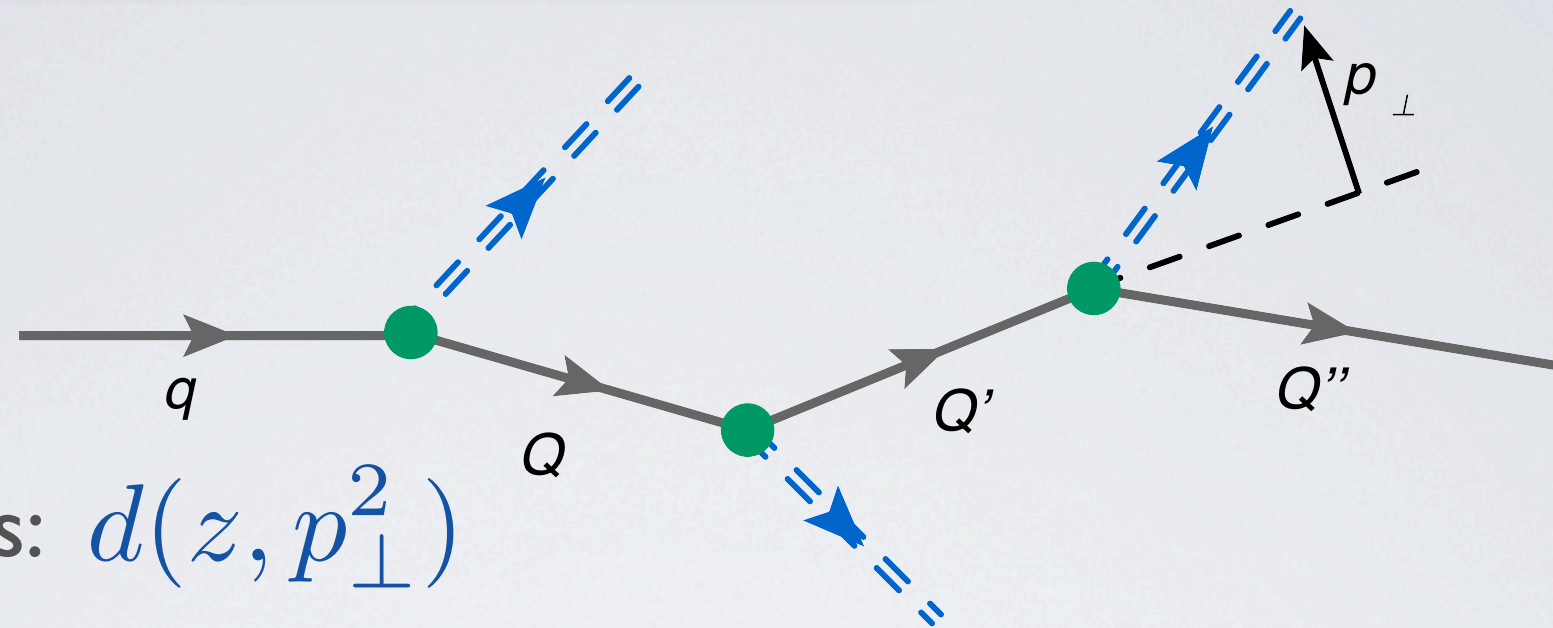


***TRANSVERSE MOMENTUM DEPENDENCE***



# TMD FRAGMENTATION FUNCTIONS

H.M.,Bentz, Cloet, Thomas, PRD.85:014021, 2012

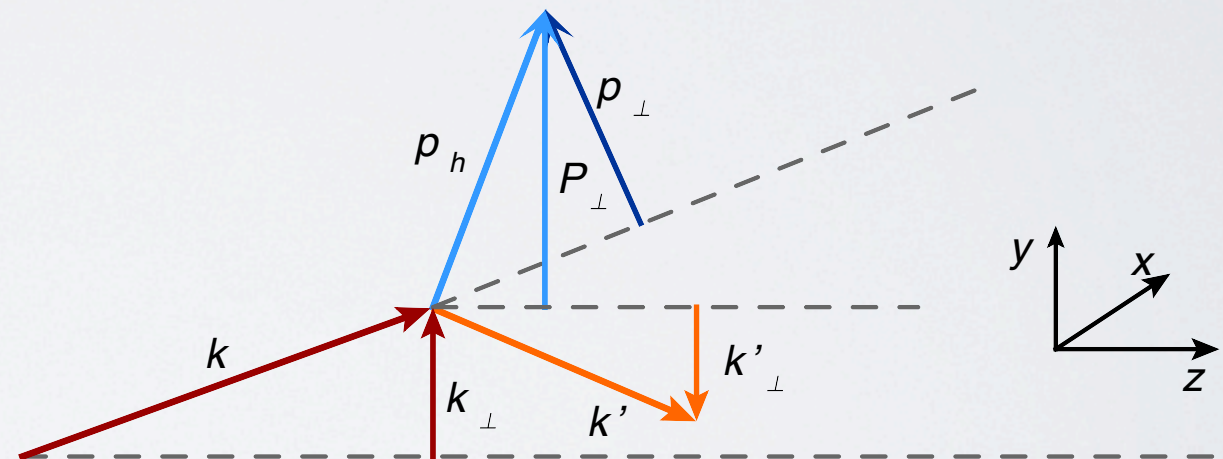


► TMD splittings:  $d(z, p_{\perp}^2)$

► Conserve transverse momenta at each link.

$$\mathbf{P}_{\perp} = \mathbf{p}_{\perp} + z\mathbf{k}_{\perp}$$

$$\mathbf{k}_{\perp} = \mathbf{P}_{\perp} + \mathbf{k}'_{\perp}$$



► Calculate the Number Density

$$D_q^h(z, P_{\perp}^2) \Delta z \pi \Delta P_{\perp}^2 = \frac{\sum_{N_{Sims}} N_q^h(z, z + \Delta z, P_{\perp}^2, P_{\perp}^2 + \Delta P_{\perp}^2)}{N_{Sims}}$$

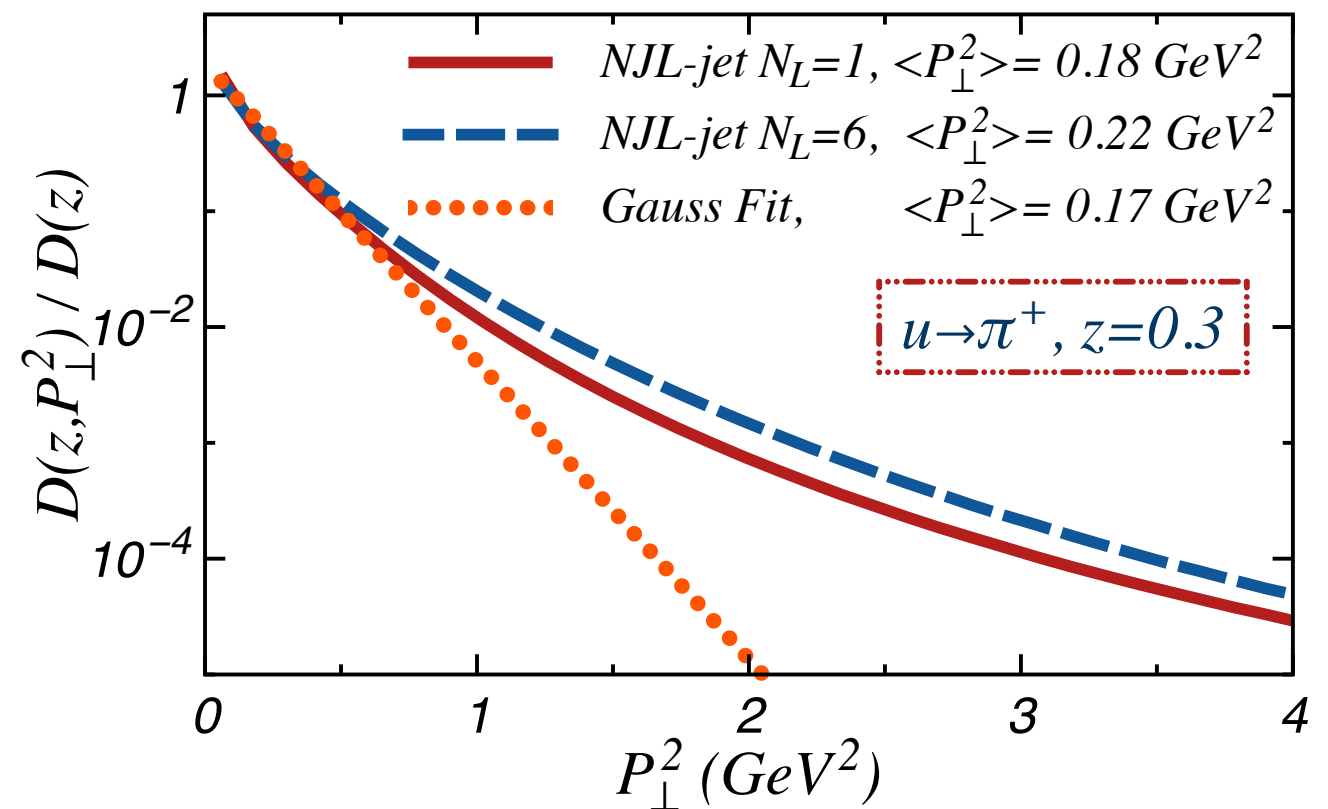
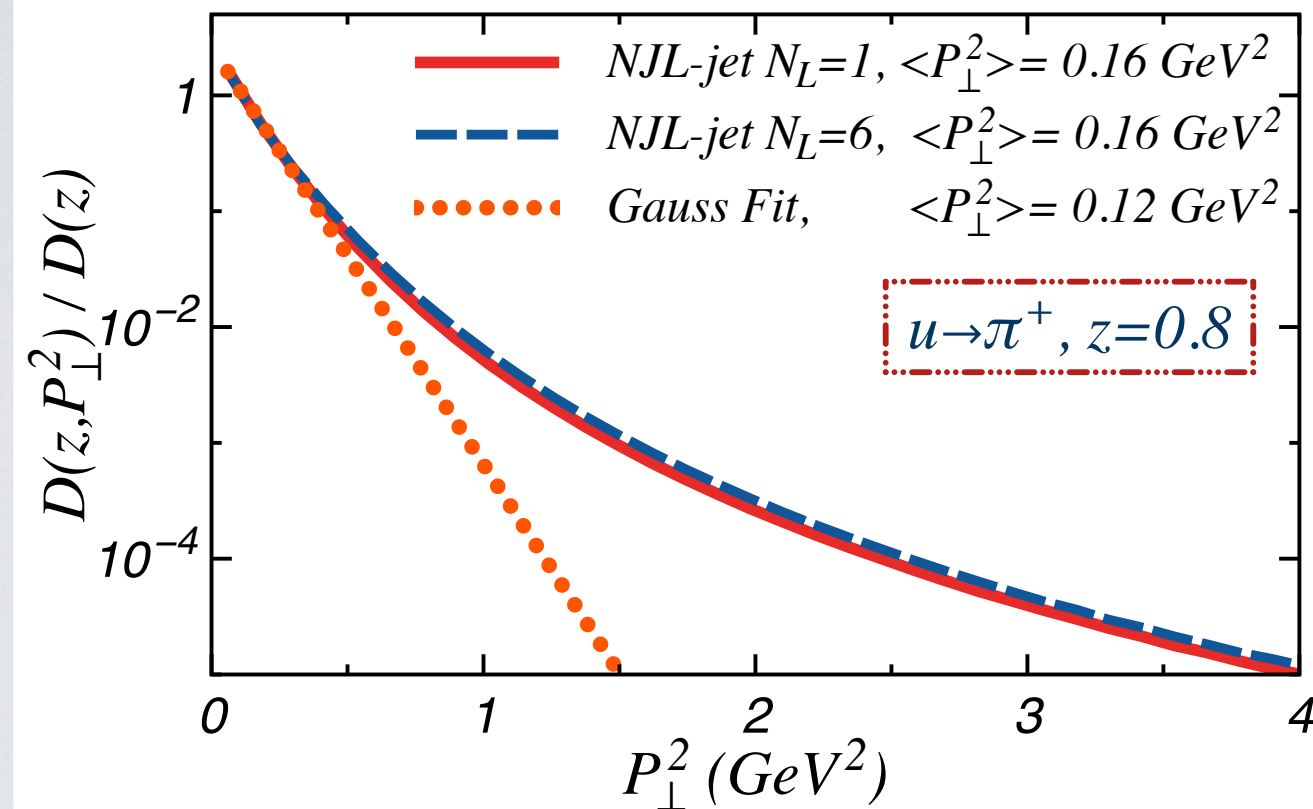


# COMPARISON WITH GAUSSIAN ANSATZ

- TMD Dependence of the splitting function.

$$d_q^h(z, p_\perp^2) \sim \frac{p_\perp^2 + [(z-1)M_1 + M_2]^2}{[p_\perp^2 + z(z-1)M_1^2 + zM_2^2 + (1-z)m_h^2]^2} \frac{1}{[1 + (M_{12}^2/\Lambda_{12}^2)^2]^2}$$

- TMD Dependence of the full fragmentation function.



- Gaussian ansatz:  $D(z, P_\perp^2) = D(z) e^{-P_\perp^2 / \langle P_\perp^2 \rangle} / \pi \langle P_\perp^2 \rangle$

✓ Multiple hadron emissions **broaden** the TM dependence, more significant at **small z**.

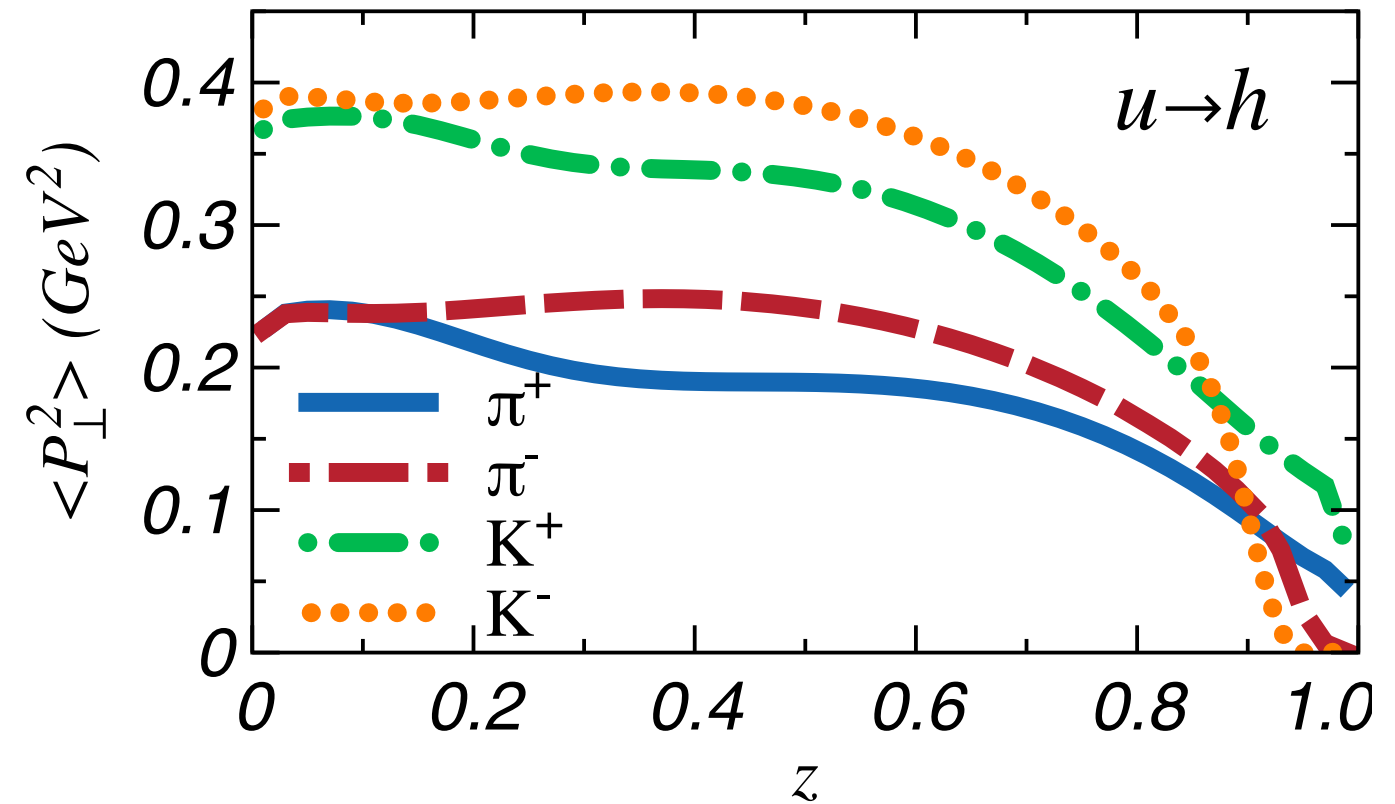
# AVERAGE TRANSVERSE MOMENTA VS $z$

## FRAGMENTATION

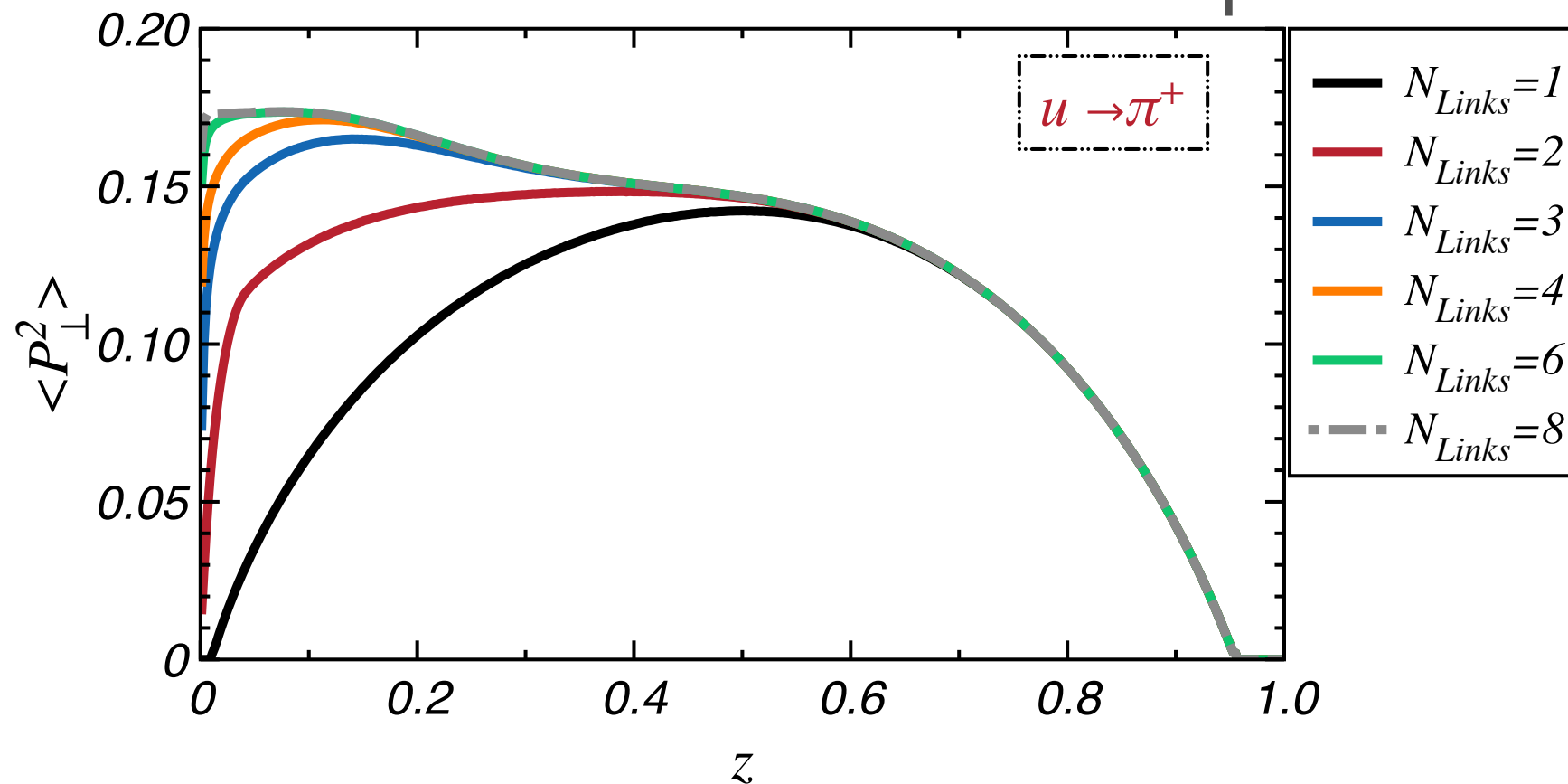
$$\langle P_{\perp}^2 \rangle_{unf} > \langle P_{\perp}^2 \rangle_f$$

◆ Indications from HERMES data:

A. Signori, et al: JHEP 1311, 194 (2013)



✓ Multiple hadron emissions: *broaden* the TM dependence at *low z*!





# CONCLUSIONS

- ❖ **Two-Hadron SIDIS** will provide information for mapping the TM and flavor dependencies of Sivers and Transversity PDFs.
- ❖ We need unintegrated Dihadron Fragmentation Functions.
- ❖ Measurements of **IFFs** for both **relative** and **total** TM will be crucial for understanding the hadronization process.
- ❖ The modified full Event Generators *incorporating Sivers effect* (**mPYTHIA**): a useful tool for phenomenological studies.
- ❖ **mPYTHIA** predictions show *a great potential* for *measuring* two-hadron SIDIS SSAs at **CLAS12** and **EIC**.

# BACKUP SLIDES



# WHITE PAPER FOR NSAC-LRP

◆ White paper on extracting DiFFs at BELLE II.

[https://www.phy.anl.gov/nsac-lrp/Whitepapers/  
StudyOfFragmentationFunctionsInElectronPositronAnnihilation.pdf](https://www.phy.anl.gov/nsac-lrp/Whitepapers/StudyOfFragmentationFunctionsInElectronPositronAnnihilation.pdf)

# TWO-HADRON SIDIS

- ▶ Cross Section in terms of **Total and Relative Momenta**

$$P_h = P_1 + P_2 \quad R = \frac{1}{2}(P_1 - P_2)$$

- ▶ The Sivers term:

$$\sigma_S = S_T \left( \sigma_T \frac{P_{hT}}{M} \sin(\varphi_T - \varphi_S) + \sigma_R \frac{R_T}{M} \sin(\varphi_R - \varphi_S) \right)$$

$$\int d\varphi_R \sigma_S = S_T \left( \sigma_{T,0} \frac{P_{hT}}{M} + \sigma_{R,1} \frac{R}{2M} \right) \sin(\varphi_T - \varphi_S)$$

$$\int d\varphi_T \sigma_S = S_T \left( \sigma_{T,1} \frac{P_{hT}}{2M} + \sigma_{R,0} \frac{R}{M} \right) \sin(\varphi_R - \varphi_S)$$

- ◆ **Non-vanishing  $\sigma_R$  is new!** Contradiction with earlier results **Bianconi: PRD62, 034008 (2000)** ? **No: Kotzinian: EPJConf. 85 02026 (2015)**

$$R^P \equiv R - (R \cdot \hat{P}_h) \hat{P}_h \quad R^P \simeq \xi_2 P_1 - \xi_1 P_2$$

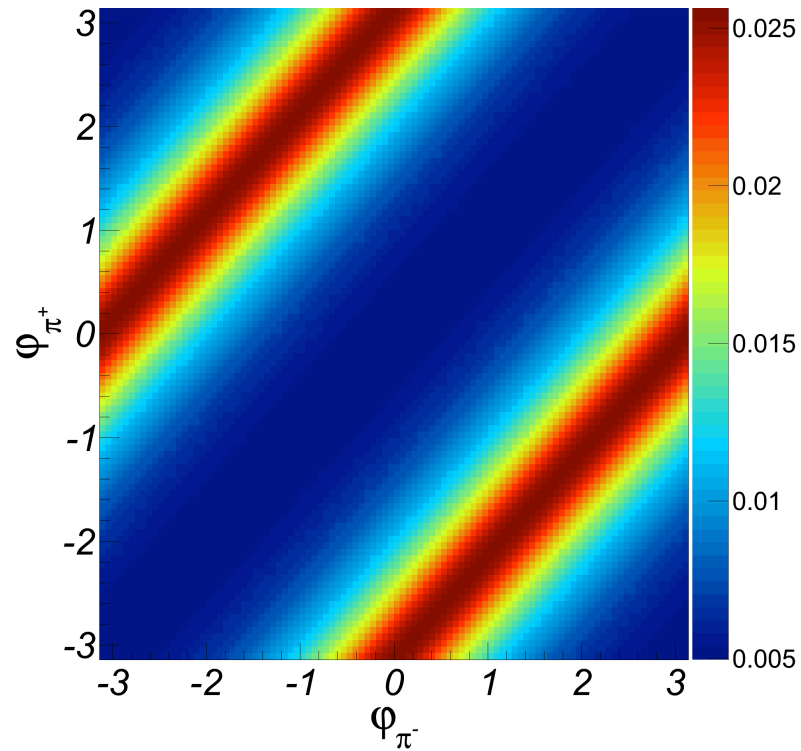
$$R_T^P \simeq \xi_2 P_{1\perp} - \xi_1 P_{2\perp} \quad \xi_i \equiv z_i / (z_1 + z_2)$$

**No  $k_T$  dependence at LO! No contradiction, different  $R$  !**

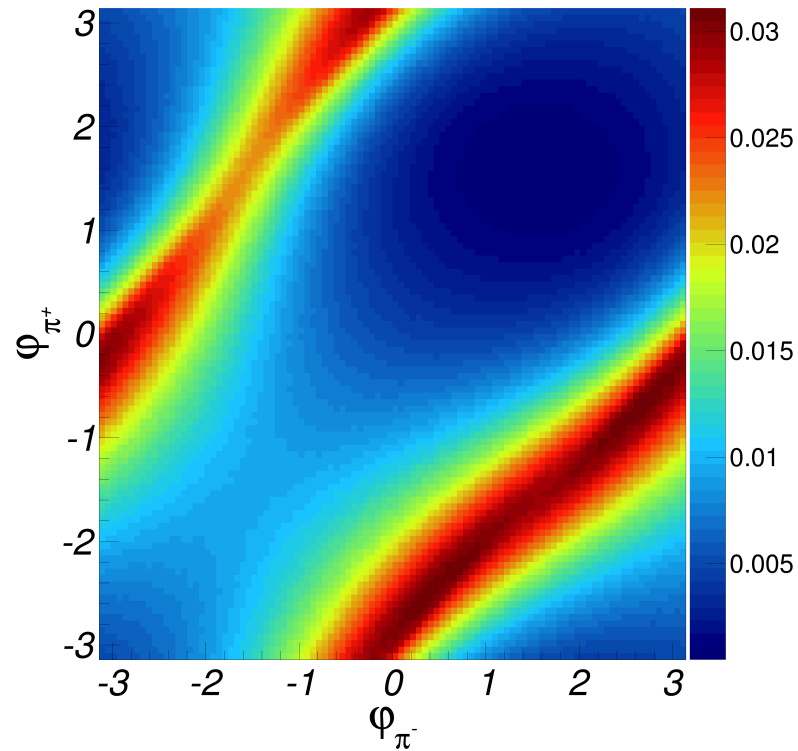


# ANGULAR CORRELATIONS: $u \rightarrow \pi^+ \pi^-$

**Unpolarized**

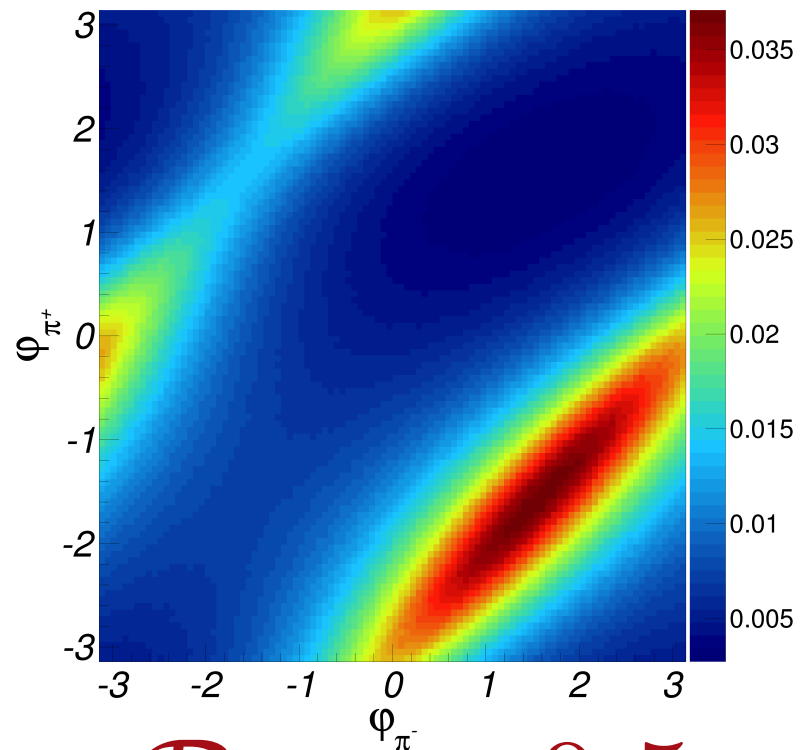
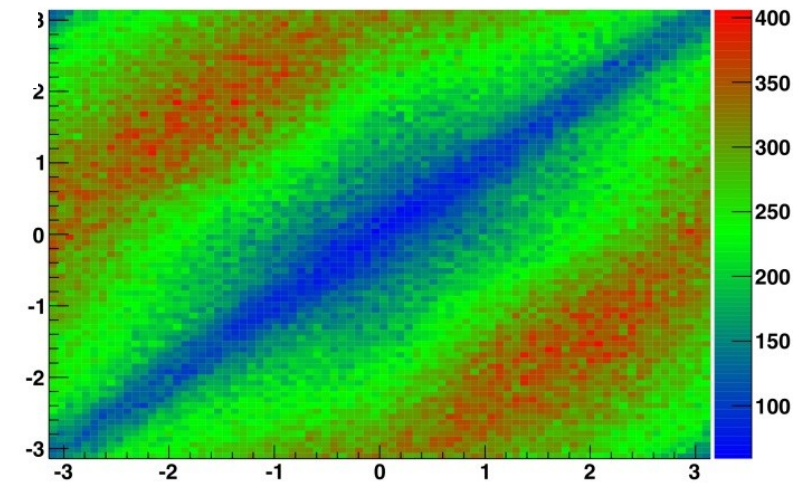


$\mathcal{P}_{SF} = 0$

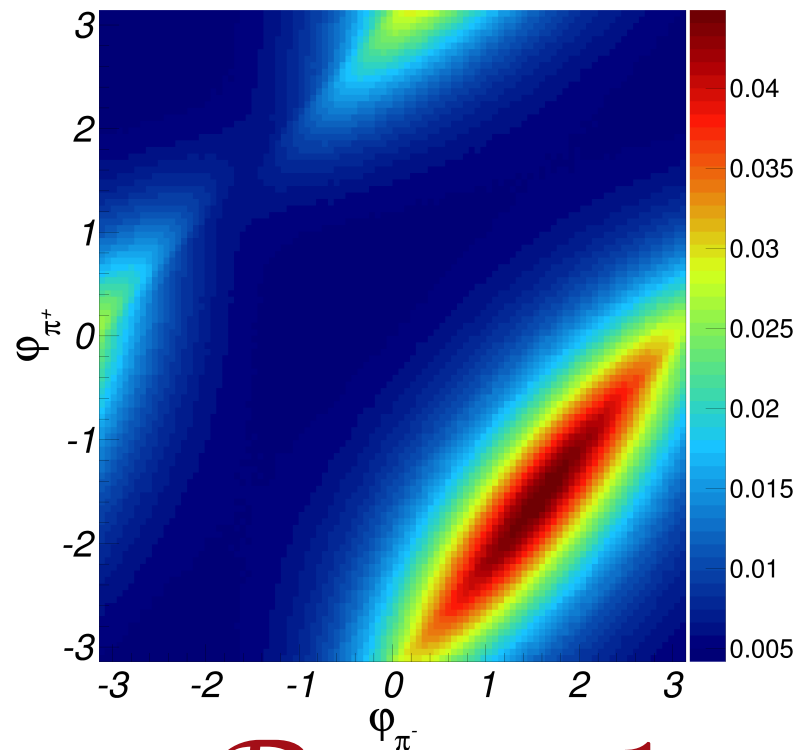


**COMPASS Preliminary:**  
F. Bradamante - COMO 2013.

$\phi_{h+}$



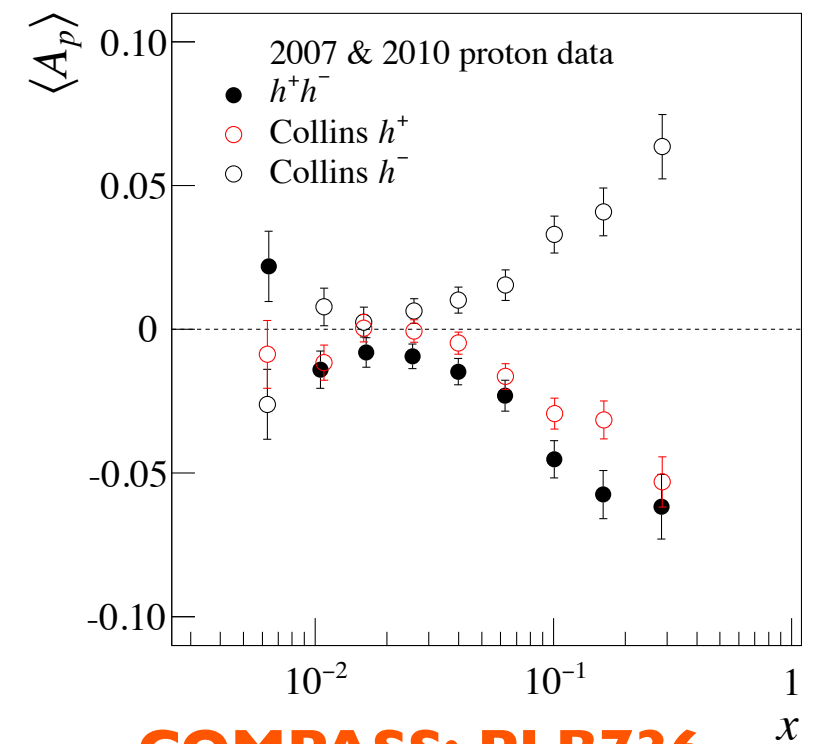
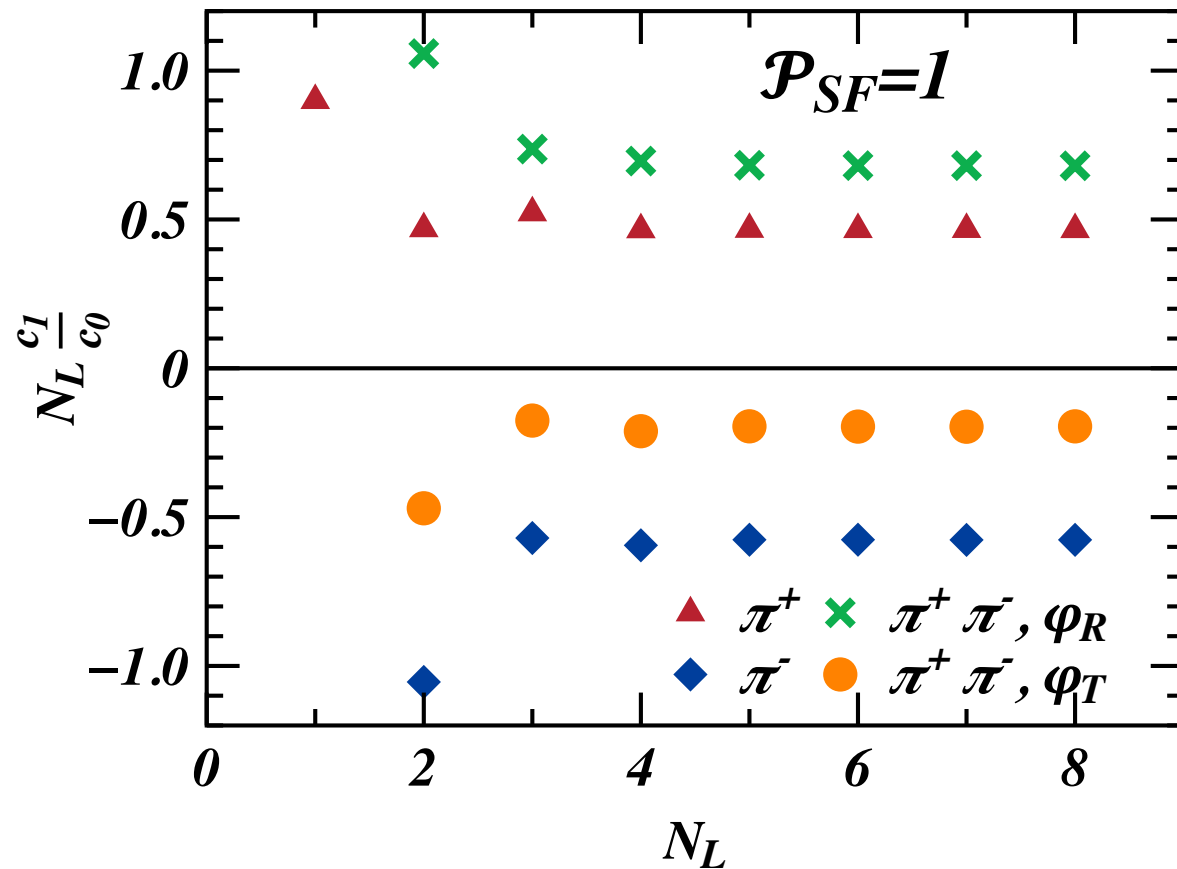
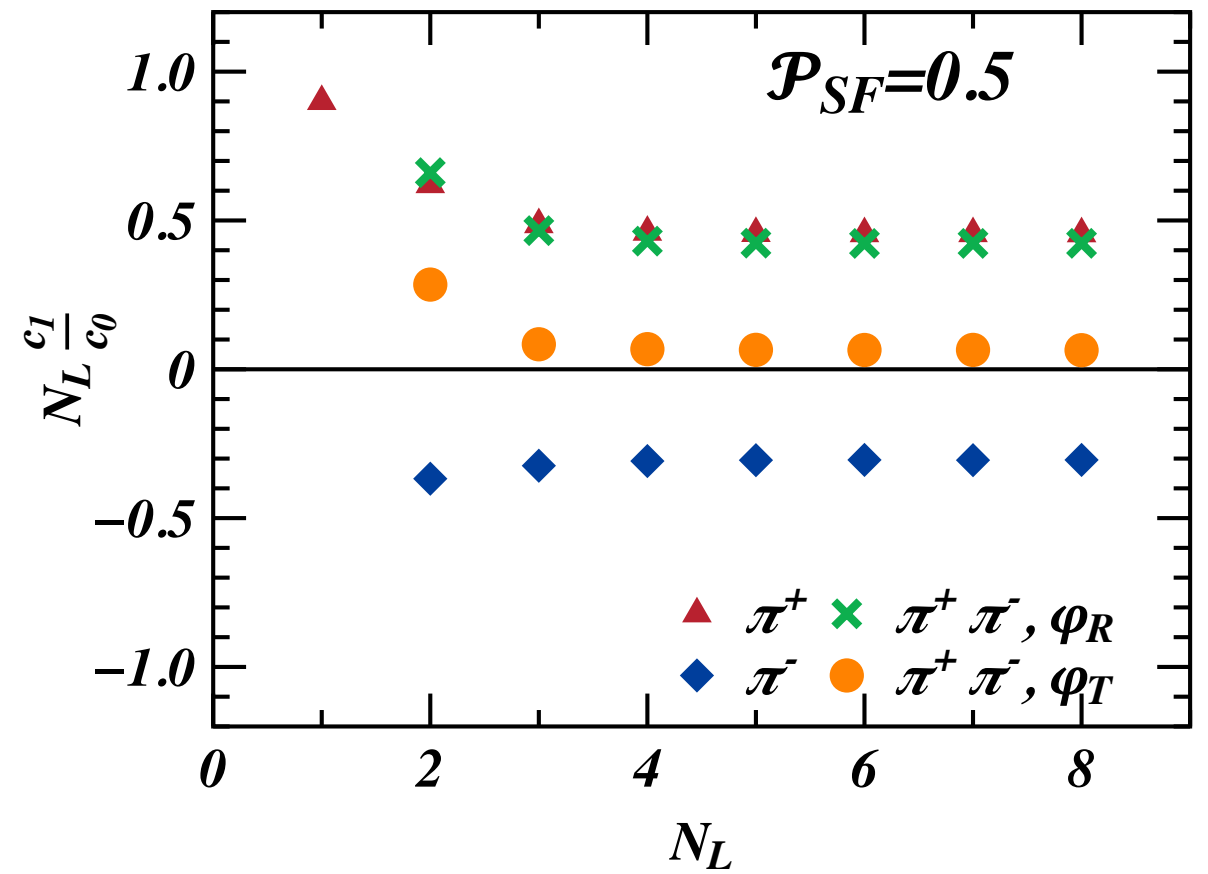
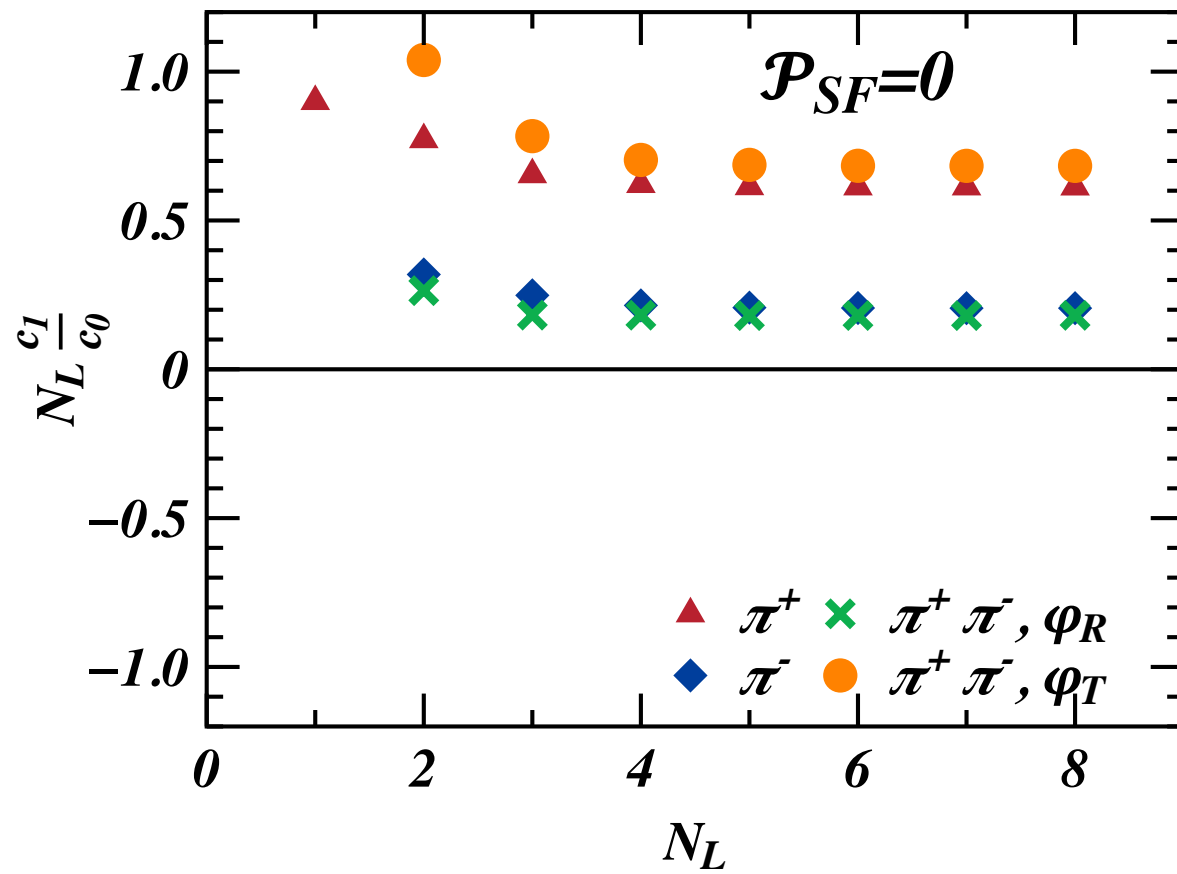
$\mathcal{P}_{SF} = 0.5$



$\mathcal{P}_{SF} = 1$

# INTEGRATED ANALYZING POWERS

$$z_{1,2} > 0.2, z > 0.2$$



**COMPASS: PLB736,  
124-131 (2014).**



## ▶ Single Hadron

$$N_h \propto \sigma_{UUU} (1 + \sin(\phi_C) A_C G)$$

$$\phi_C = \phi_h - \phi_{S'}$$

$$= \phi_h + \phi_S - \pi$$

$$A_C = \frac{\sum_q e_q^2 \Delta_{Tq} \otimes H_1^{\perp h/q}}{\sum_q e_q^2 q \otimes D_1^{h/q}}$$

## ▶ DiHadron

$$N_{h^+h^-} \propto \sigma_{UUU} (1 + \sin(\phi_{RS}) A_{UT}^{\sin \phi_{RS}} F)$$

$$\phi_{RS} = \phi_R - \phi_{S'}$$

$$= \phi_R + \phi_S - \pi$$

$$A_{UT}^{\sin \phi_{RS}} \propto \frac{\sum_q e_q^2 \cdot h_1^q \cdot H_{1,q}^{\triangleleft}}{\sum_q e_q^2 \cdot f_1^q \cdot D_{1,q}}$$

## ▶ NJL-Jet Fits

$$D_{q^\uparrow} = c_0 - \sin(\phi) c_1$$

## ▶ Single Hadron

$$N_h \propto \sigma_{UUU} (1 \oplus \sin(\phi_C) A_C G)$$

$$\phi_C = \phi_h - \phi_{S'}$$

$$= \phi_h + \phi_S - \pi$$

$$A_C = \frac{\sum_q e_q^2 \Delta_{Tq} \otimes H_1^{\perp h/q}}{\sum_q e_q^2 q \otimes D_1^{h/q}}$$

## ▶ DiHadron

$$N_{h^+h^-} \propto \sigma_{UUU} (1 \oplus \sin(\phi_{RS}) A_{UT}^{\sin \phi_{RS}} F)$$

$$\phi_{RS} = \phi_R - \phi_{S'}$$

$$= \phi_R + \phi_S - \pi$$

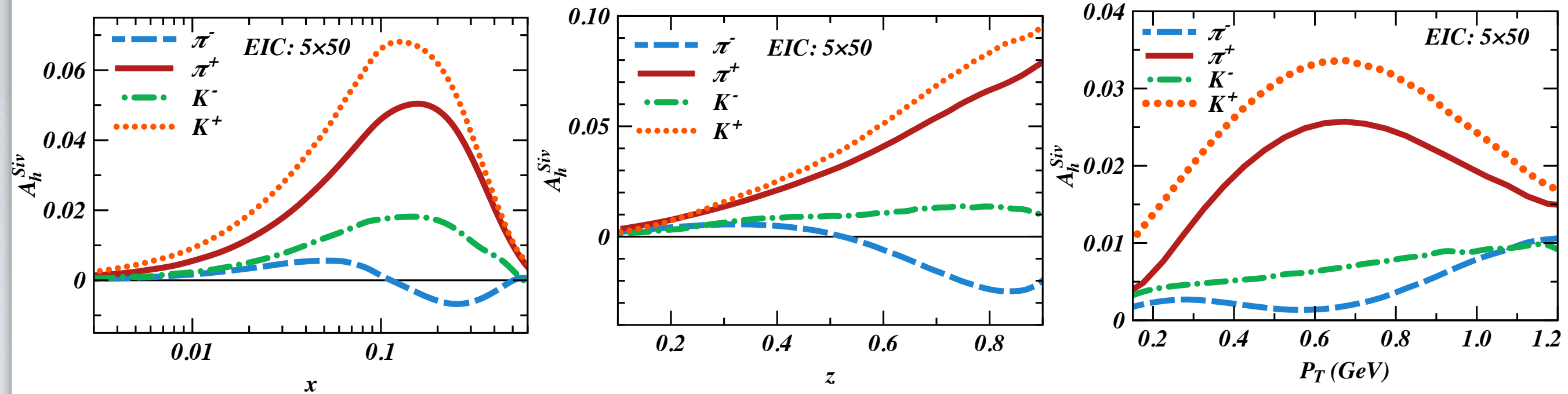
$$A_{UT}^{\sin \phi_{RS}} \propto \frac{\sum_q e_q^2 \cdot h_1^q \cdot H_{1,q}^{\triangleleft}}{\sum_q e_q^2 \cdot f_1^q \cdot D_{1,q}}$$

## ▶ NJL-Jet Fits

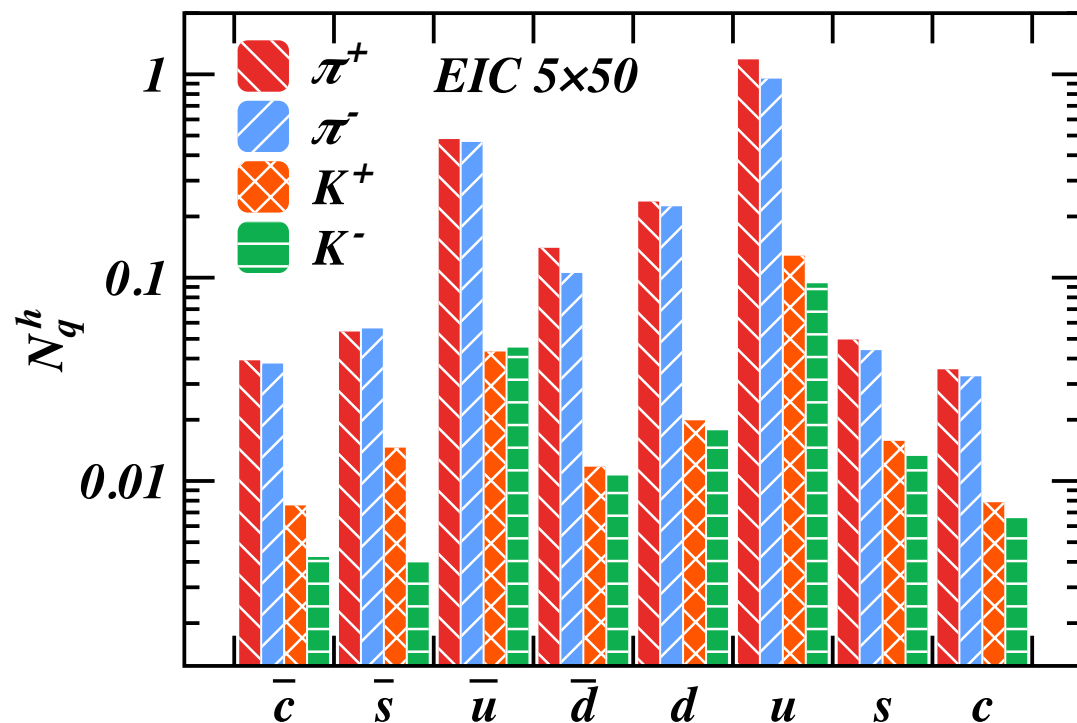
$$D_{q^\uparrow} = c_0 \ominus \sin(\phi) c_1$$



◆ SSAs for charged pions and kaons from **proton** target - low  $x$  region.



◆ **Average** number of hadrons by struck quark flavor.

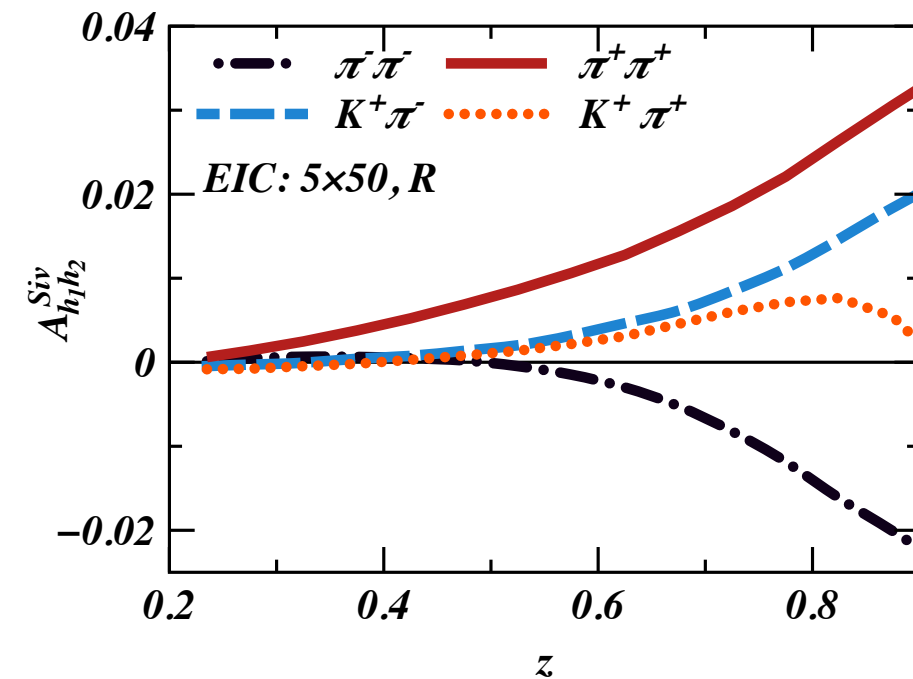
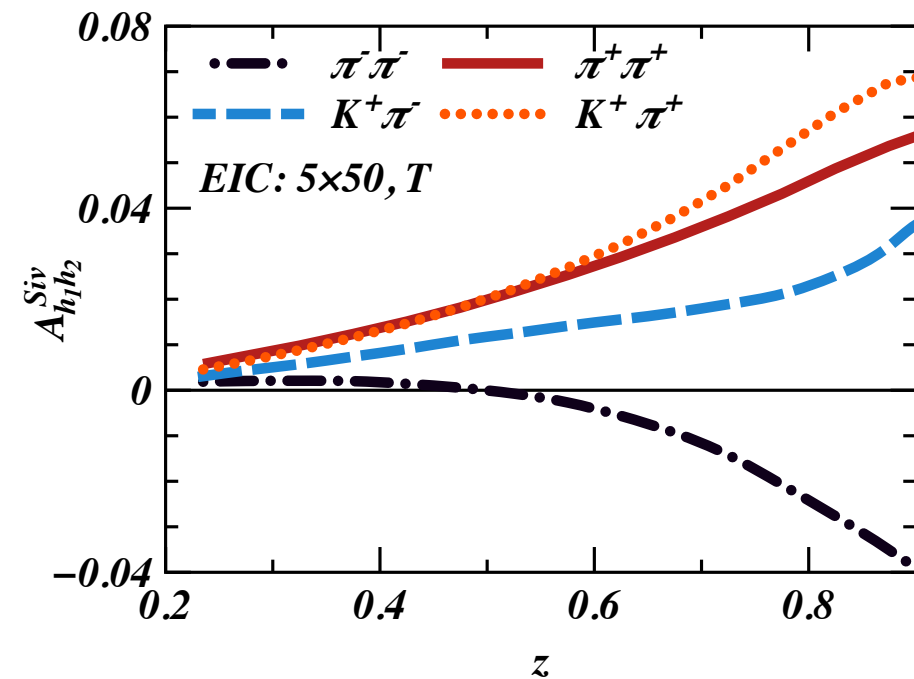
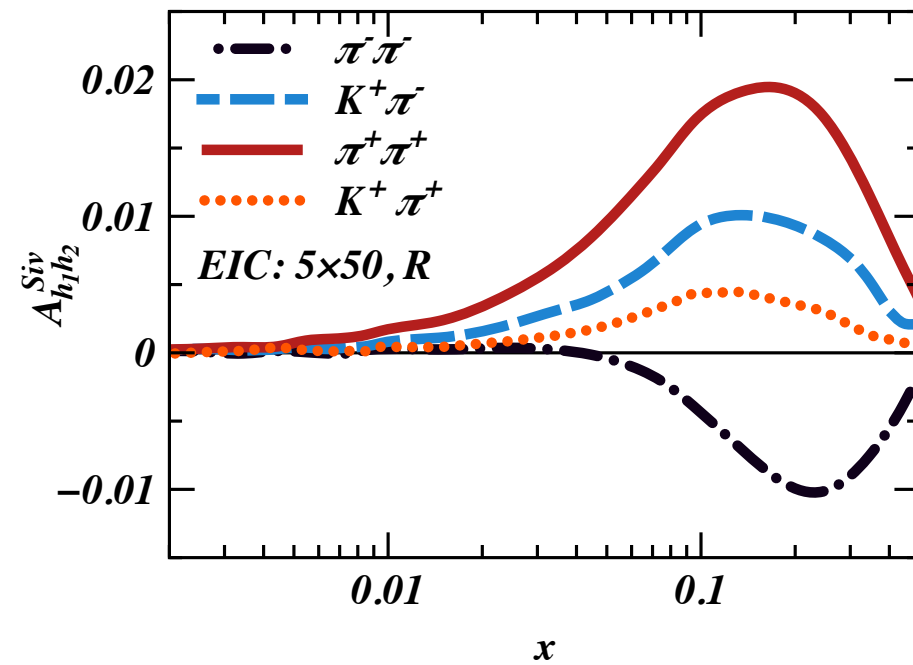
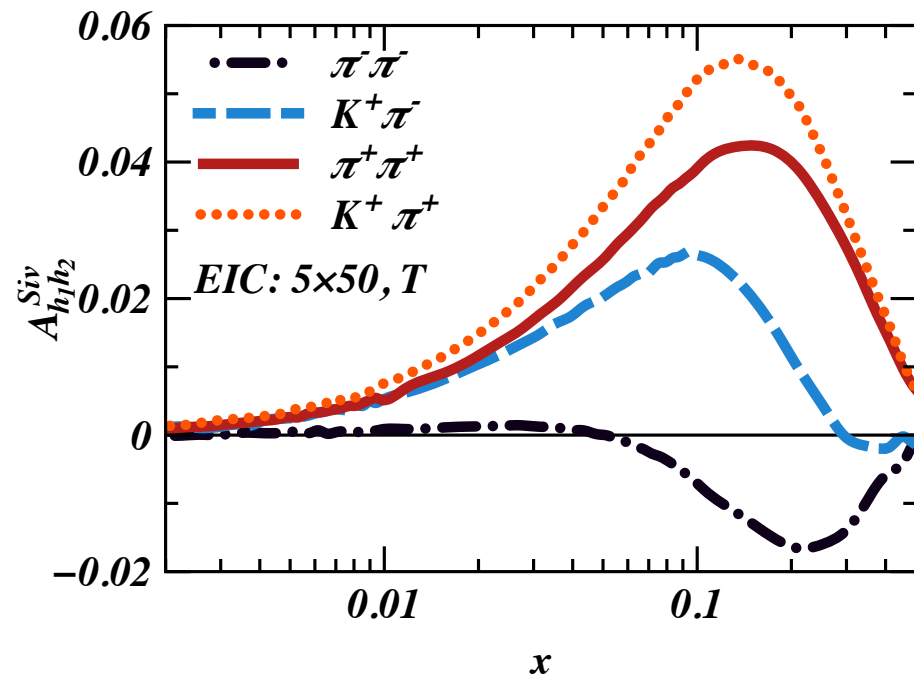


◆  $\pi^+$  multiplicities **larger** than  $K^+$ , but kaon **SSAs** are larger. Up quark dominates the multiplicities.

# Dihadron Sivers SSAs for EIC

H.M et al., arXiv:1502.02669 (2015).

◆ Identical pairs via z-ordering:  $z_1 \geq z_2$  (so  $\sigma_R \neq 0$ )



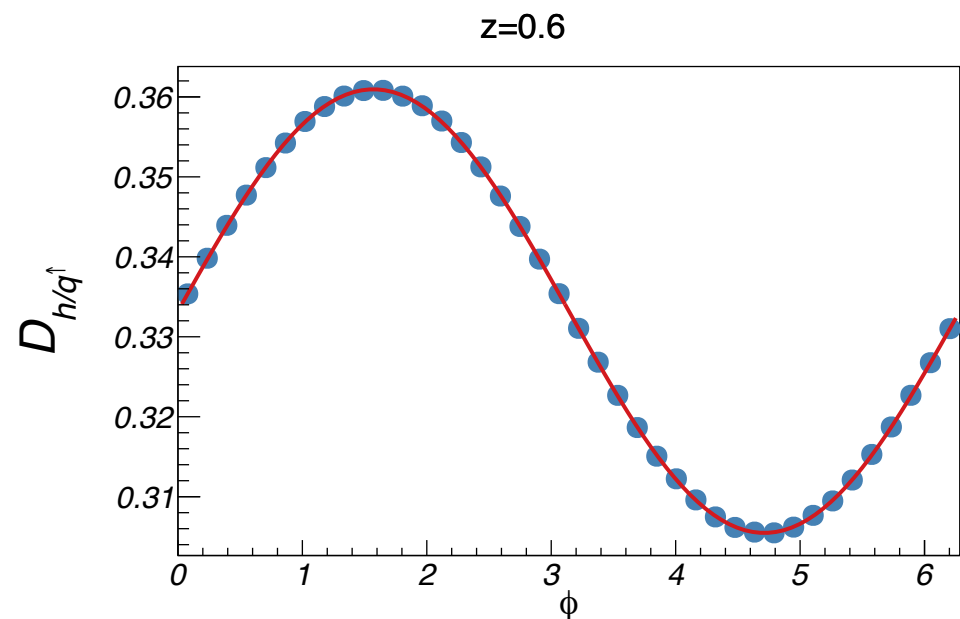
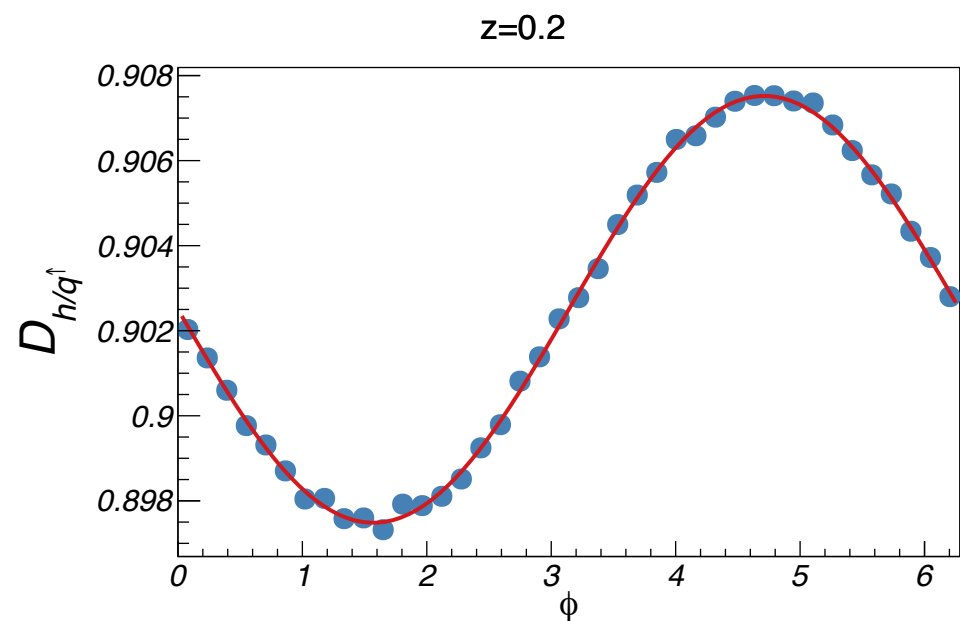
- Dihadron SSAs are *comparable* to single hadron ones! (the one- and two-hadron FFs should mostly cancel in the ratios)



# INTEGRATED POLARIZED FRAGMENTATIONS

- Integrate Polarized Fragmentations over  $P_{\perp}^2$

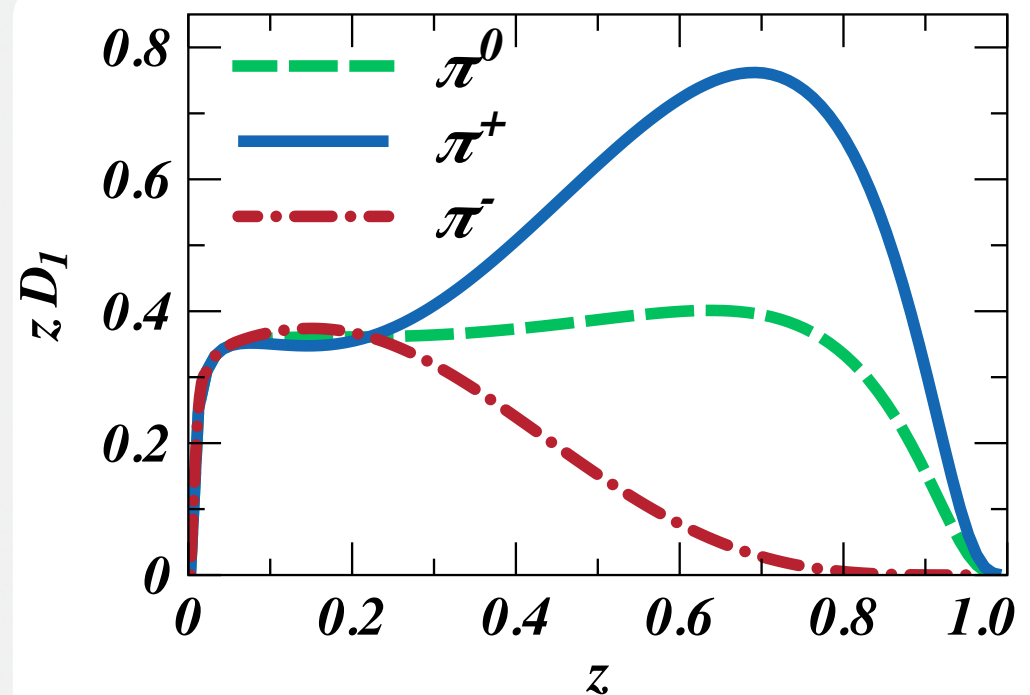
$$D_{h/q^{\uparrow}}(z, \varphi) \equiv \int_0^{\infty} dP_{\perp}^2 D_{h/q^{\uparrow}}(z, P_{\perp}^2, \varphi) = \frac{1}{2\pi} \left[ D_1^{h/q}(z) - 2H_{1(h/q)}^{\perp(1/2)}(z) S_q \sin(\varphi) \right]$$



$$D_1^{h/q}(z) \equiv \pi \int_0^{\infty} dP_{\perp}^2 D_1^{h/q}(z, P_{\perp}^2)$$

$$H_{1(h/q)}^{\perp(1/2)}(z) \equiv \pi \int_0^{\infty} dP_{\perp}^2 \frac{P_{\perp}}{2zm_h} H_1^{\perp h/q}(z, P_{\perp}^2)$$

- Fit with form:  $F(c_0, c_1) = c_0 - c_1 \sin(\varphi)$



# COLLINS EFFECT - MK2

## MK2 Model Assumptions:

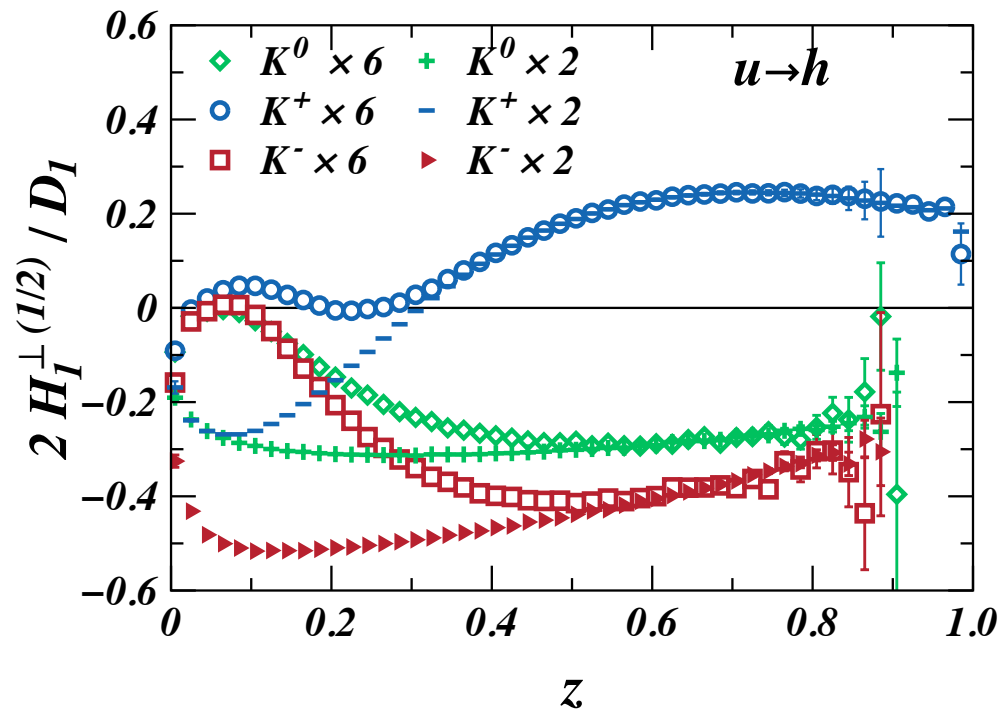
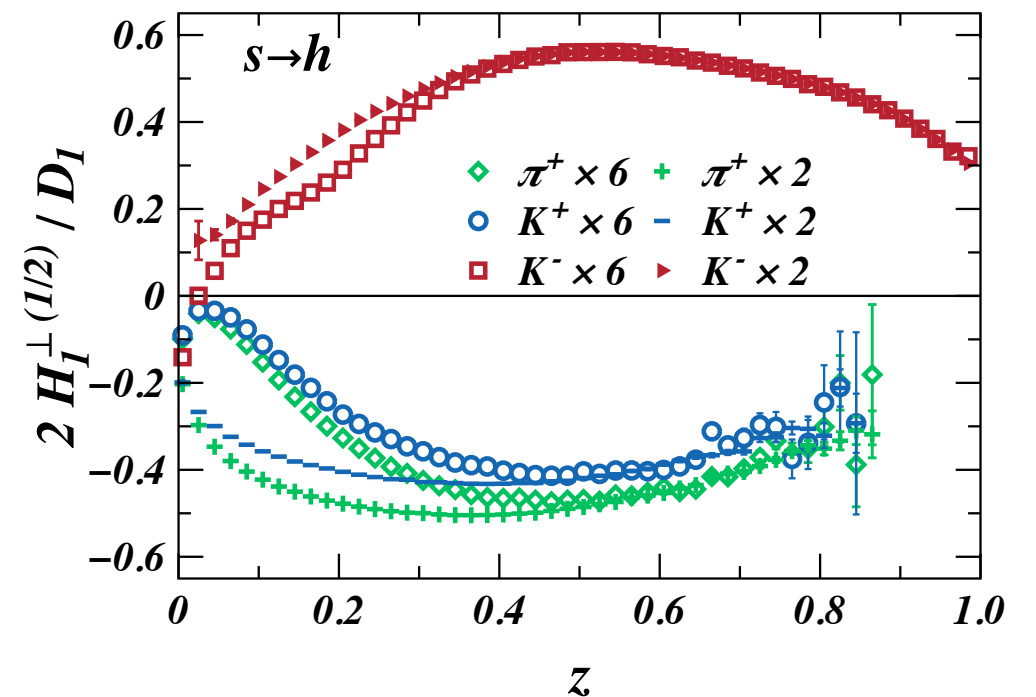
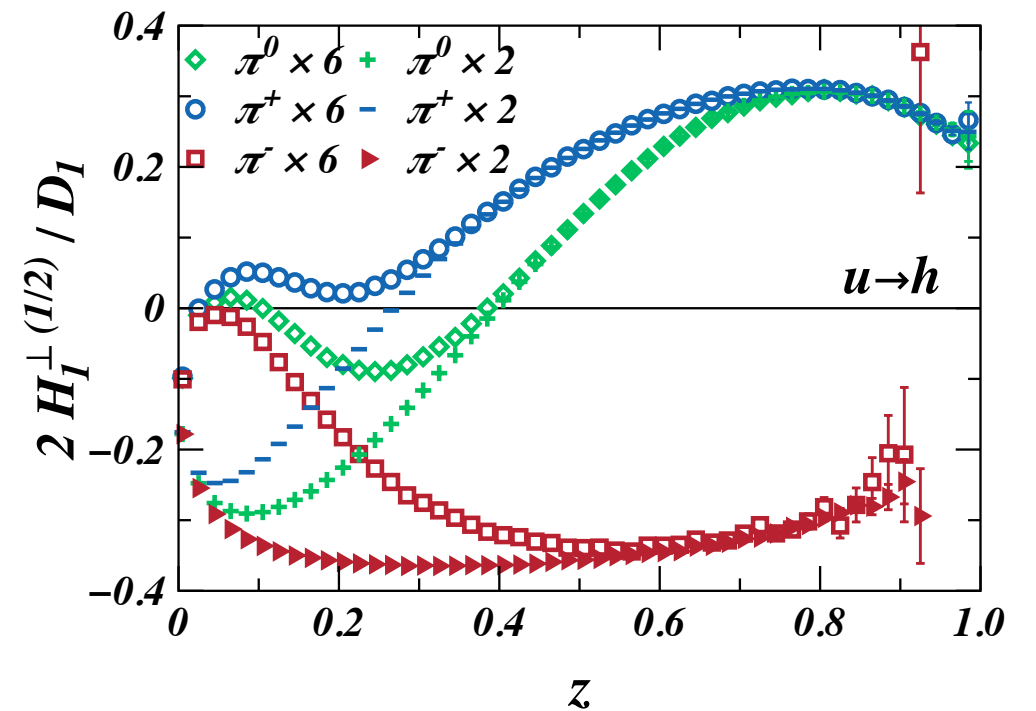
H.M., Kotzinian, Thomas, PLB731 208-216 (2014).

1. Allow for Collins Effect only in a SINGLE emission vertex -  $N_L^{-1}$  scaling of the resulting Collins function.
2. Use constant values for  $\mathcal{P}_{SF}$

$$\mathcal{P}_{SF} = 1$$

◆ The results for  $N_L=2$  and  $N_L=6$ , scaled up by a factor  $N_L$ .

$$F(c_0, c_1) = c_0 - c_1 \sin(\varphi)$$



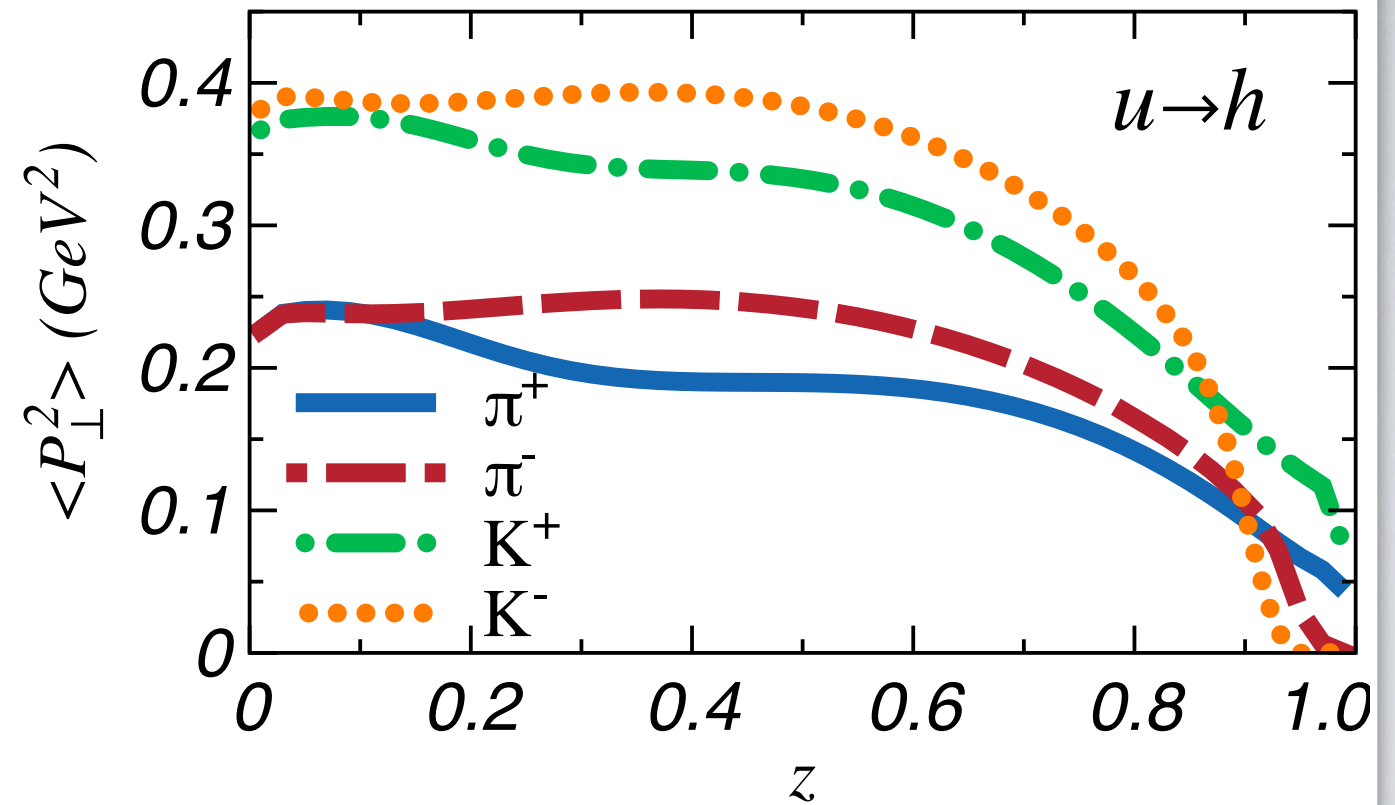
# AVERAGE TRANSVERSE MOMENTA VS $z$

## FRAGMENTATION

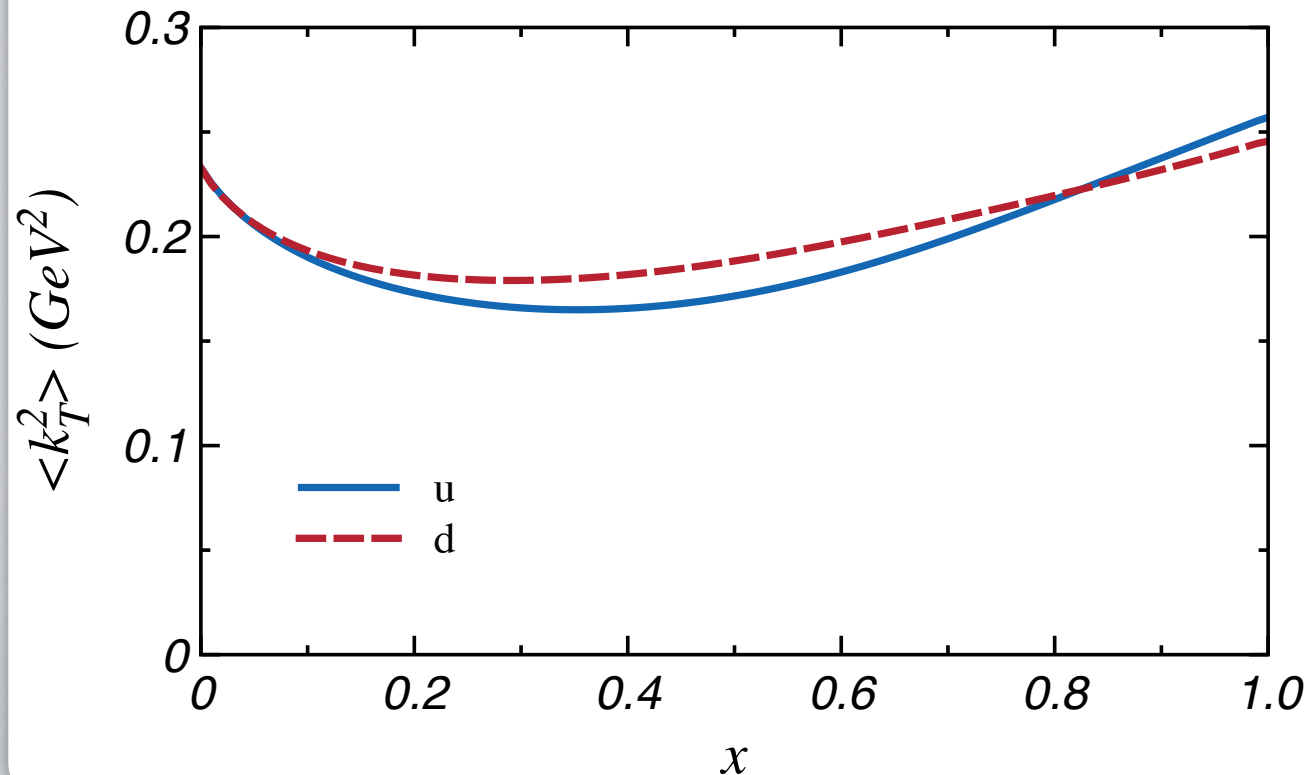
$$\langle P_{\perp}^2 \rangle_{unf} > \langle P_{\perp}^2 \rangle_f$$

◆ Indications from HERMES data:

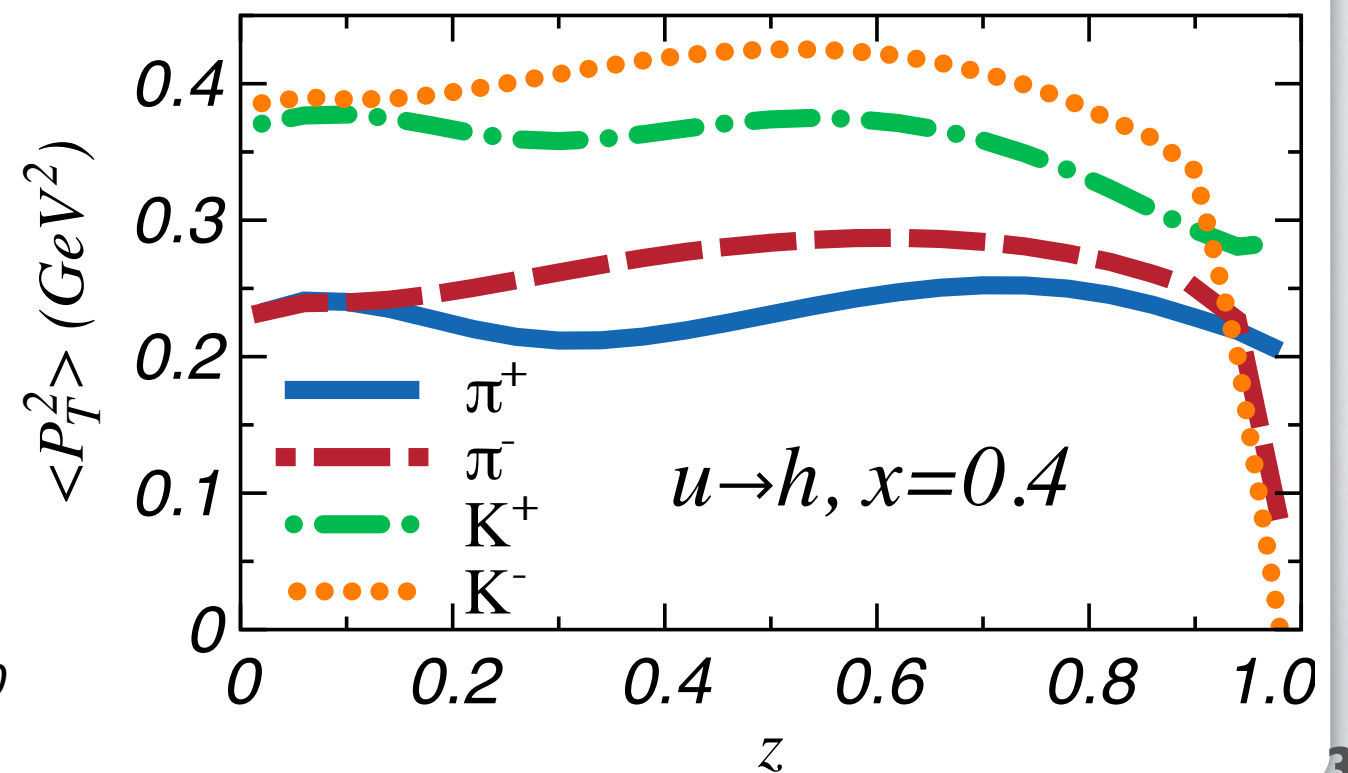
A. Signori, et al: JHEP 1311, 194 (2013)



## PDF

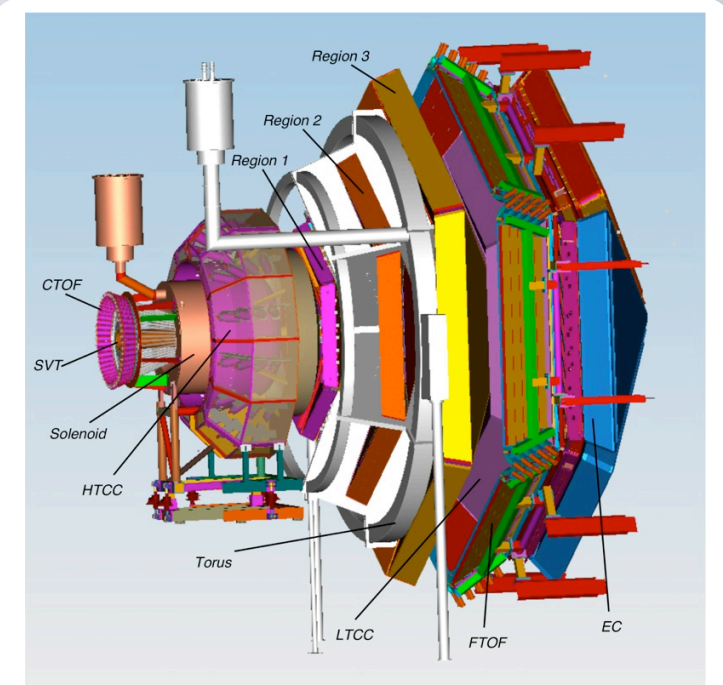


## SIDIS





# CLAS12 @ JLAB 12GeV



- Upcoming SIDIS experiment, 1H and 2H
- 11 GeV electron off polarized proton target.
- Access to large  $x$  region of nucleon structure.
- We use mPYTHIA for SIDIS predictions.

- Include the kinematical cuts on  $x, Q^2, W, \theta_{e'}, \theta_h, M_{Mis}, z, \dots$

$$0.075 \leq x \leq 0.532$$

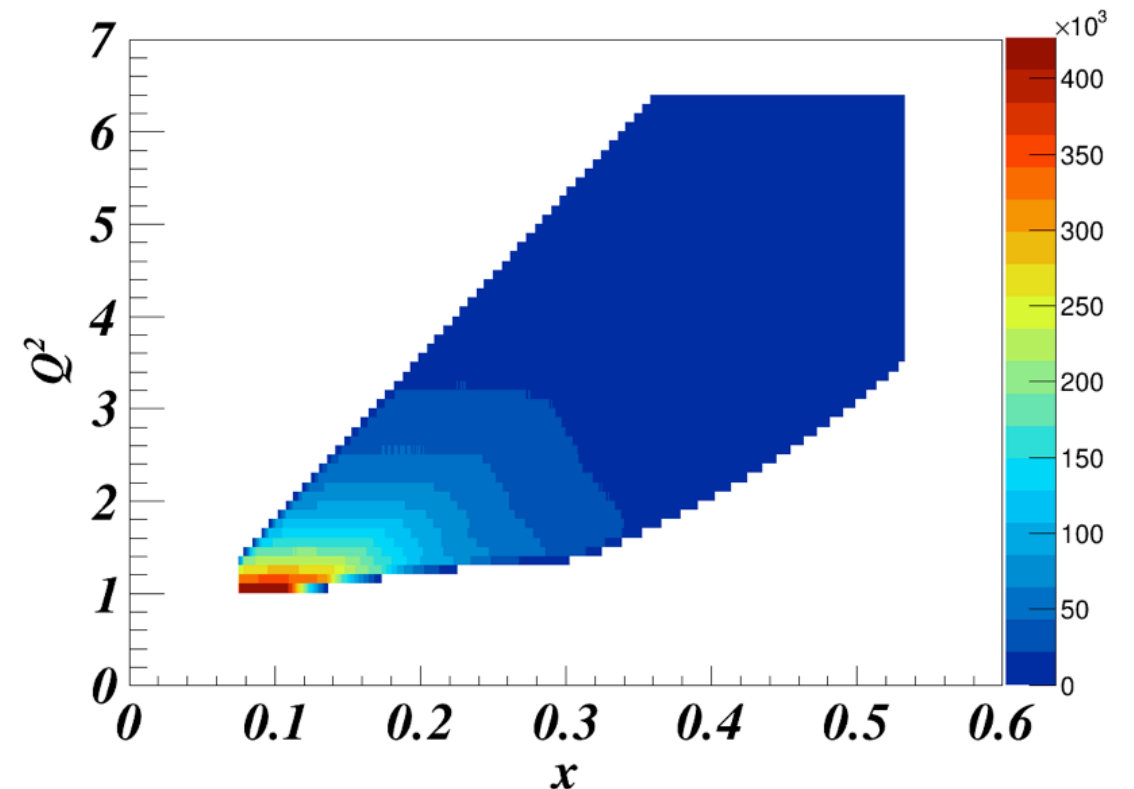
$$1 \text{ GeV} \leq Q^2 \leq 6.3 \text{ GeV}$$

$$W \geq 2 \text{ GeV}$$

$$M_{Mis(ep)-(e'hX)} \geq 1.5 \text{ GeV}$$

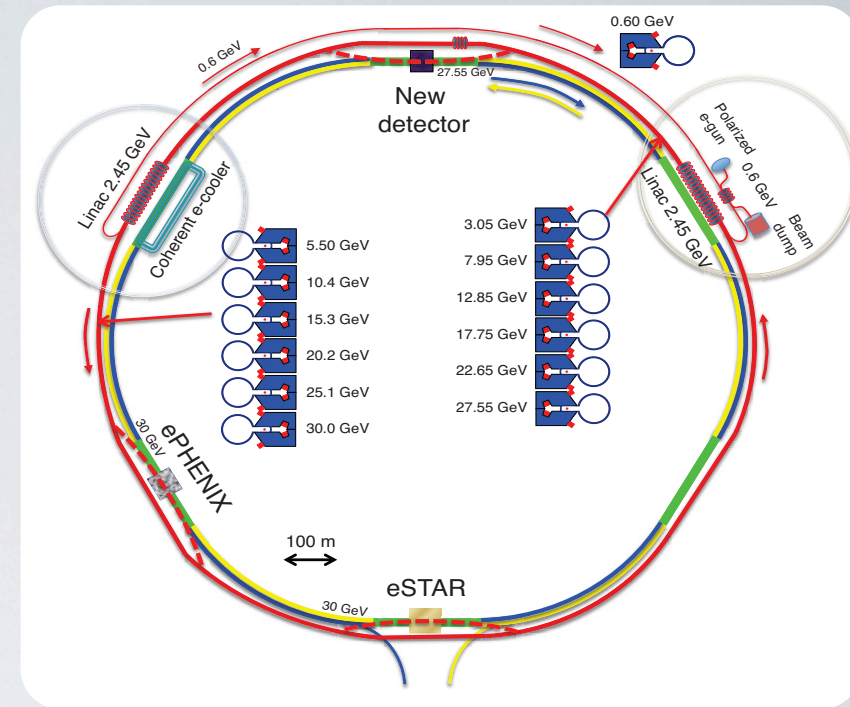
$$M_{Mis(ep)-(e'h_1h_2X)} \geq 1.5 \text{ GeV}$$

## DIS kinematics

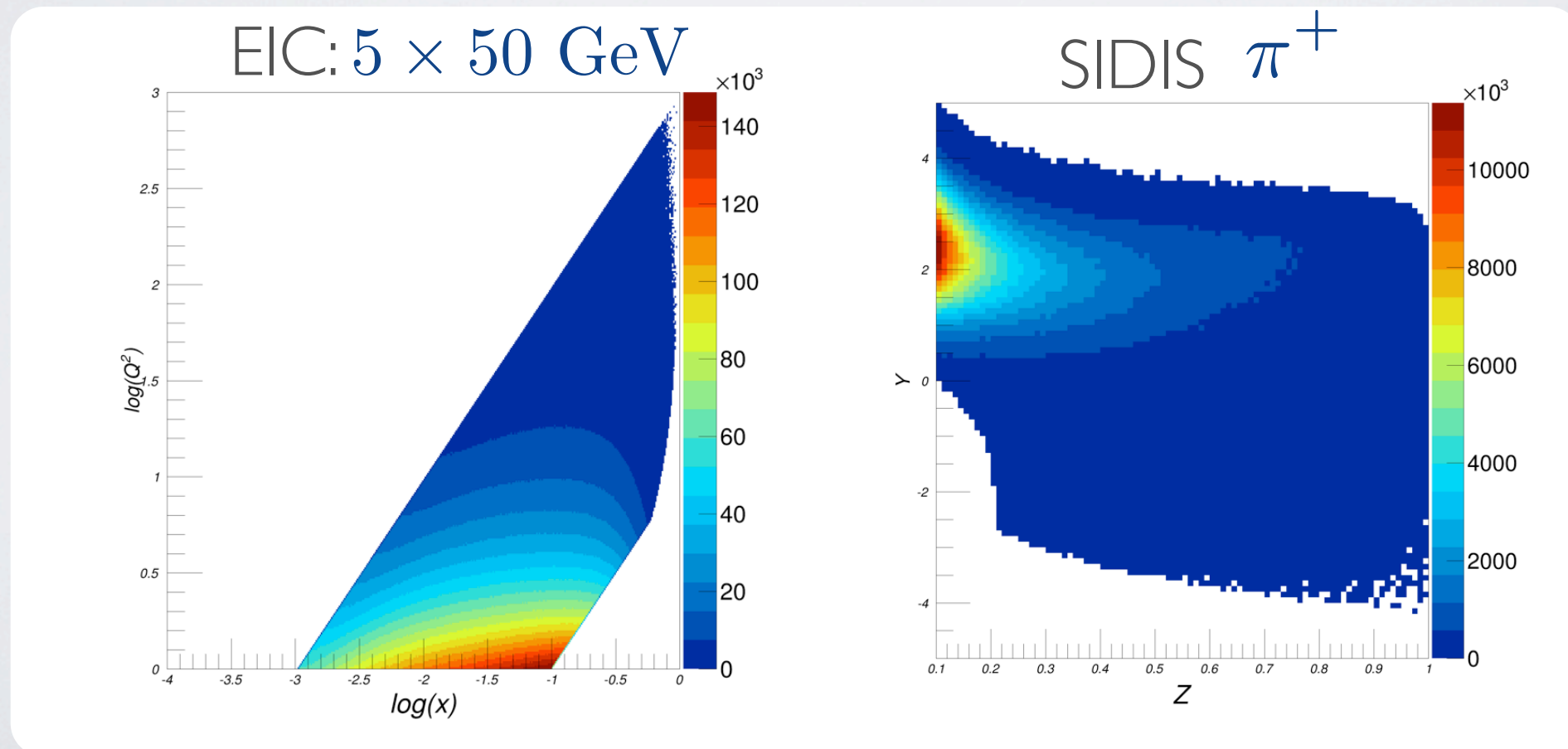


# EIC: eRHIC

White Paper -- Accardi et. al. : 1212.1701(2012).



- EIC using RHIC + electron ring.
- Various proposed beam momenta:  $l_e \times P_N$
- We use mPYTHIA for SIDIS predictions.



- **Sivers Single Spin Asymmetry:**

$$\langle \sin(\phi - \phi_S) \rangle_{UT}^h \equiv \frac{\int d\phi_h d\phi_S \sin(\phi_h - \phi_S) [d\sigma(\phi_h, \phi_S) - d\sigma(\phi_h, \phi_S + \pi)]}{\int d\phi_h d\phi_S [d\sigma(\phi_h, \phi_S) + d\sigma(\phi_h, \phi_S + \pi)]}$$

$$A_{Siv}^P \equiv 2 \langle \sin(\phi - \phi_S) \rangle_{UT}^h$$

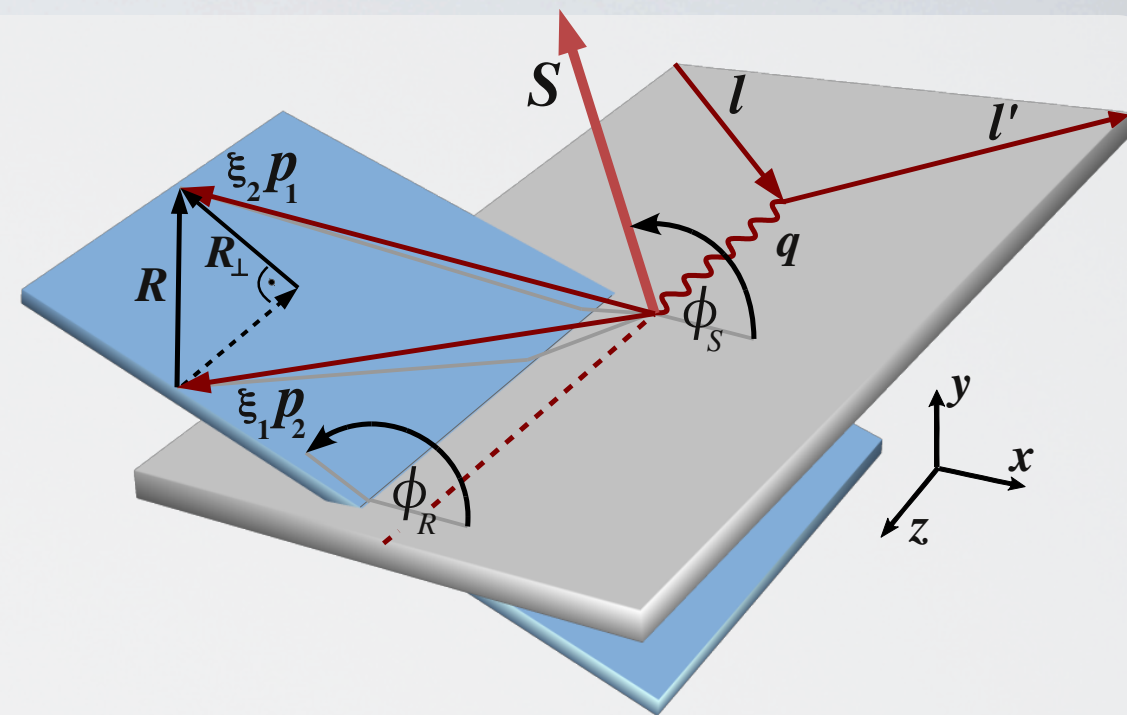
$$\langle \sin(\phi - \phi_S) \rangle_{UT}^h \sim \frac{\mathcal{C}[f_{1T}^{\perp, q} D_1^{h/q}]}{\mathcal{C}[f_1^q D_1^{h/q}]}$$



# ACCESS TO TRANSVERSITY PDF FROM DIFF

**M. Radici, et al: PRD 65, 074031 (2002).**

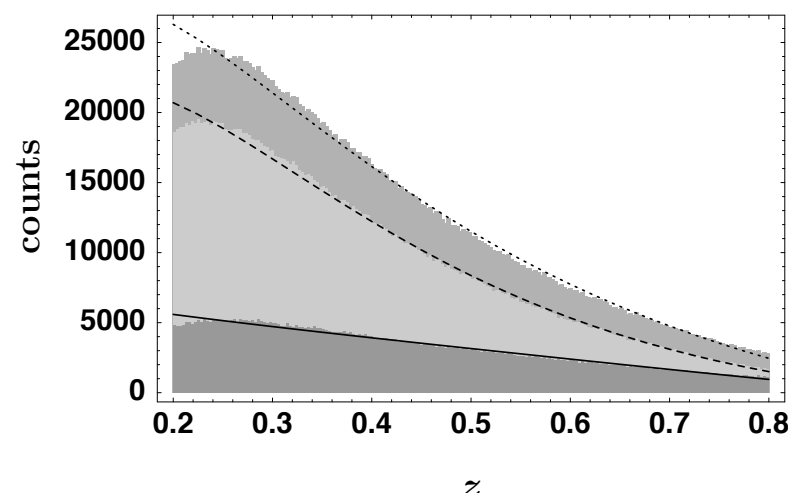
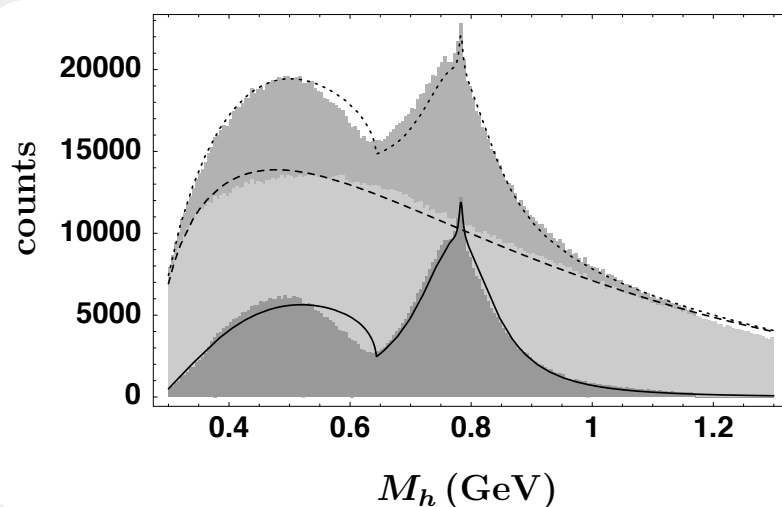
- In two hadron production from polarized target the cross section factorizes *collinearly* - **no TMD!**
- Allows clean access to *transversity*.
- *Unpolarized* and *Interference* Dihadron FFs are needed!



$$\frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} \propto \sin(\phi_R + \phi_S) \frac{\sum_q e_q^2 h_1^q(x)/x H_1^{\triangleleft q}(z, M_h^2)}{\sum_q e_q^2 f_1^q(x)/x D_1^q(z, M_h^2)}$$

- Empirical Model for  $D_1^q$  has been fitted to PYTHIA simulations.

**A. Bacchetta and M. Radici, PRD 74, 114007 (2006).**

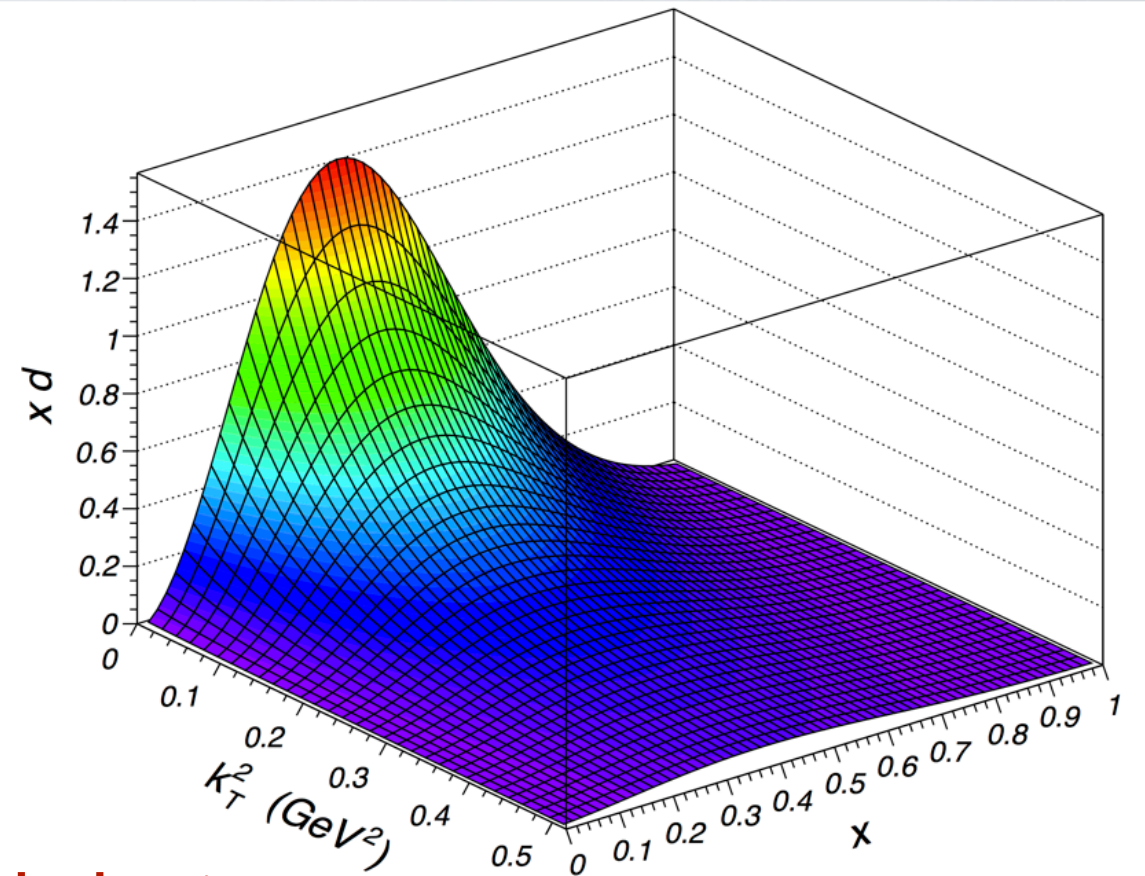
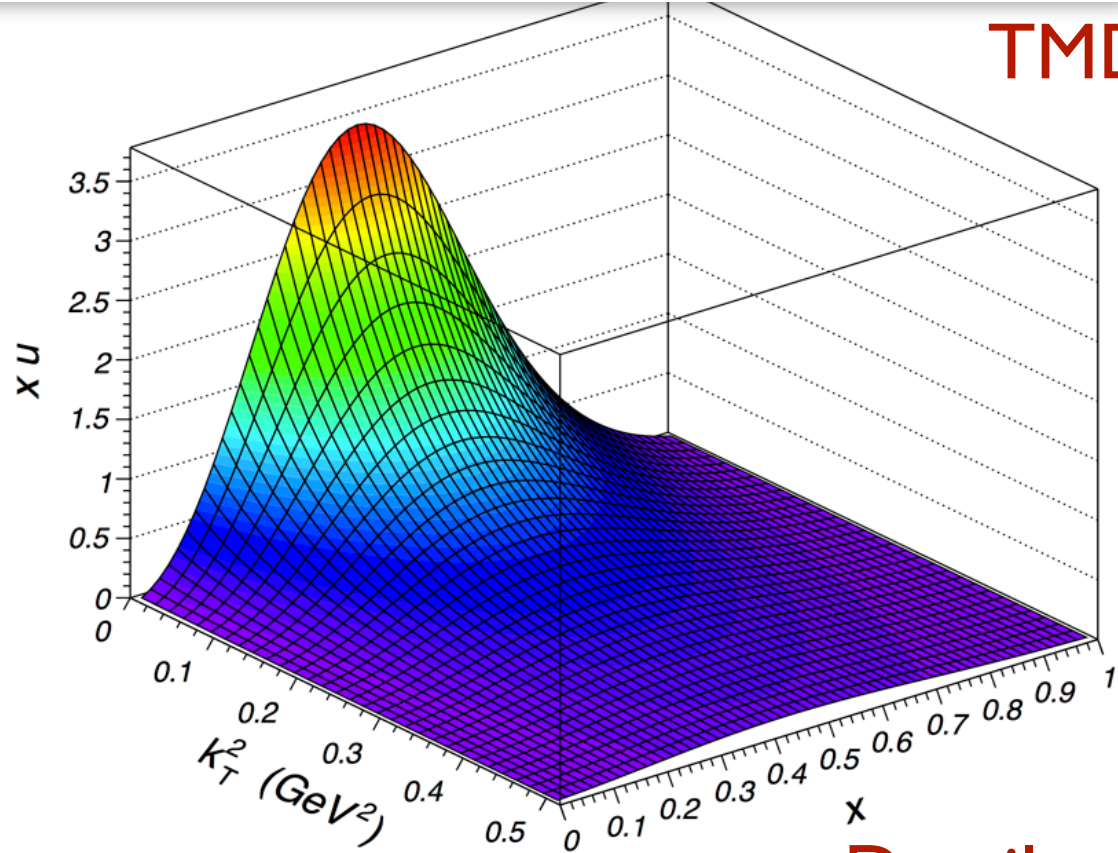


**Experiments:  
BELLE,  
HERMES,  
COMPASS.**

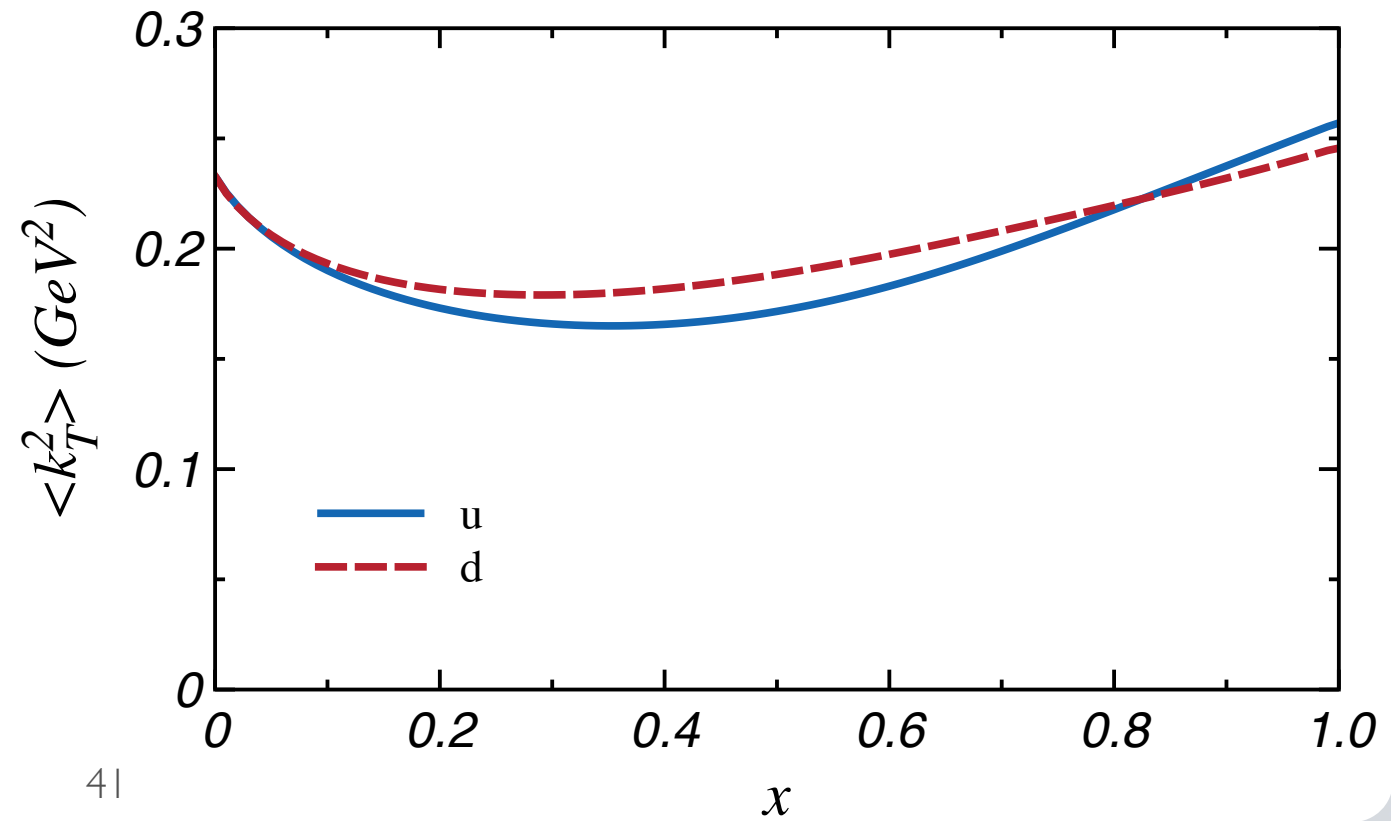
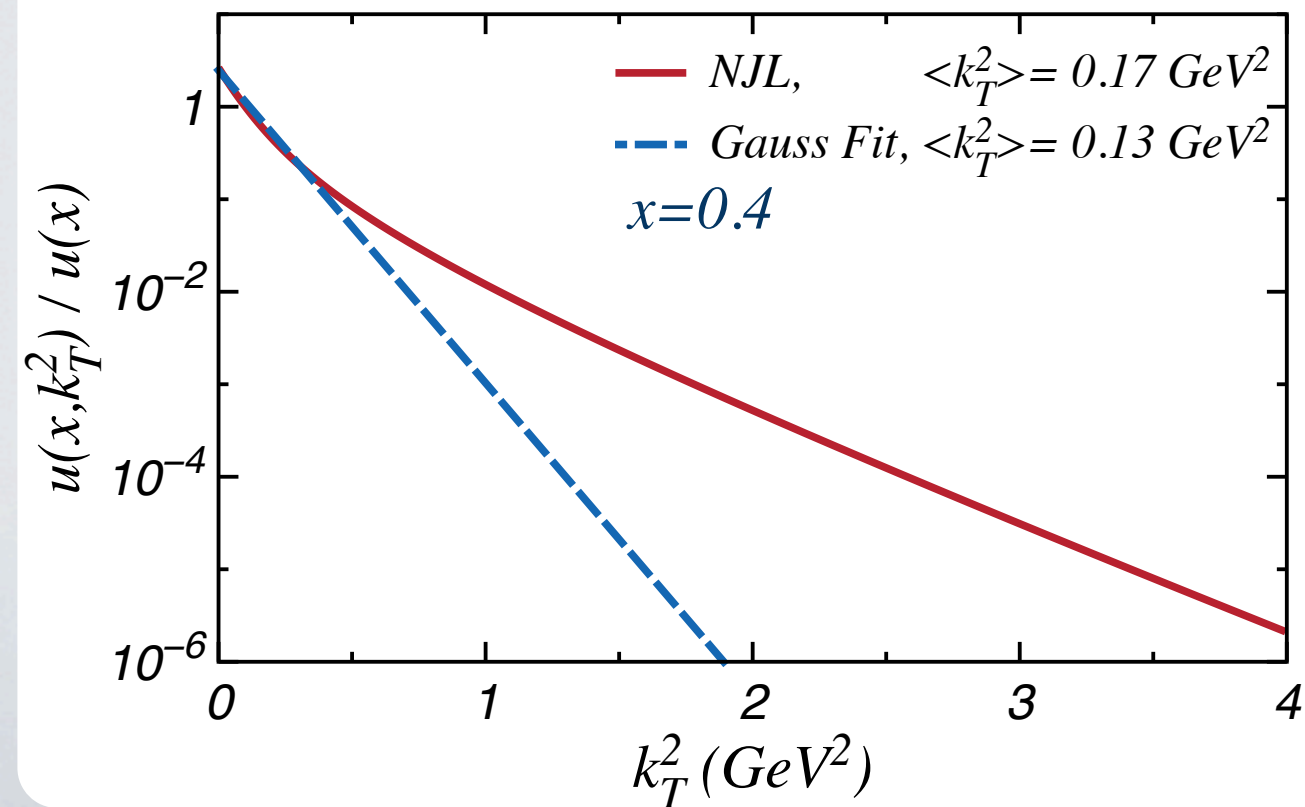
# NJL: NUCLEON PDFS - TMD RESULTS

H.M., Bentz, Cloet, Thomas, PRD.85:014021, 2012

## TMD PDFs



## Details of TMD behavior





# GAUSSIAN ANSATZ

► Need to calculate convolution PDFs and FFs:

$$F_{UU} = \sum_q e_q^2 f_1^q(x, k_T^2, Q^2) \otimes d\sigma^{lq \rightarrow lq} \otimes D_q^h(z, P_\perp^2, Q^2)$$

► Using **Gaussian** Ansatz For TM dependences of PDFs and FFs:

$$f_1^q(x, k_T^2) = f_1^q(x) \frac{e^{-k_T^2 / \langle k_{T,q}^2 \rangle}}{\pi \langle k_{T,q}^2 \rangle} \quad D_q^h(z, P_\perp^2) = D(z)_q^h \frac{e^{-P_\perp^2 / \langle P_\perp^{2,q \rightarrow h} \rangle}}{\pi \langle P_\perp^{2,q \rightarrow h} \rangle}$$

► Only involved collinear PDFs and FFs.

$$F_{UU} = \sum_q e_q^2 f_1^q(x, Q^2) D_q^h(z, Q^2) \frac{e^{-P_T^2 / \langle P_T^2 \rangle}}{\pi \langle P_T^2 \rangle}$$

$$\langle P_T^2 \rangle(z) = \langle P_\perp^2 \rangle + z^2 \langle k_T^2 \rangle$$

$$\langle k_T^2 \rangle \equiv \frac{\int d^2 \mathbf{k}_T k_T^2 f(x, k_T^2)}{\int d^2 \mathbf{k}_T f(x, k_T^2)}$$

$$\langle P_\perp^2 \rangle \equiv \frac{\int d^2 \mathbf{P}_\perp P_\perp^2 D(z, P_\perp^2)}{\int d^2 \mathbf{P}_\perp D(z, P_\perp^2)}$$



# EMPIRICAL EXTRACTIONS OF SIVERS PDF

M. Anselmino et. al.: PRD 72, 094007 (2005). PRD 86, 014028 (2012).

- Sivers SSAs from SIDIS
- Use **LO** expression for factorized cross-section.
- Parametrize PDFs and FFs.
- Use Gaussian TMD dependence.
- Also **TMD evolution** in 2012.

$$A_{Siv}^h \equiv 2 \frac{\int d\varphi_S d\varphi_h (\sigma_{\uparrow}^h - \sigma_{\downarrow}^h) \sin(\varphi_h - \varphi_S)}{\int d\varphi_S d\varphi_h (\sigma_{\uparrow}^h + \sigma_{\downarrow}^h)}$$

$$A_{Siv}^h \sim \mathcal{C}[k_T f_{1T}^{\perp q} D_1] / \mathcal{C}[f_1^q D_1^{h/q}]$$

$$f_1^q(x, k_T) = f_q(x) \frac{1}{\pi \mu^2} e^{-k_T^2 / \mu^2}$$

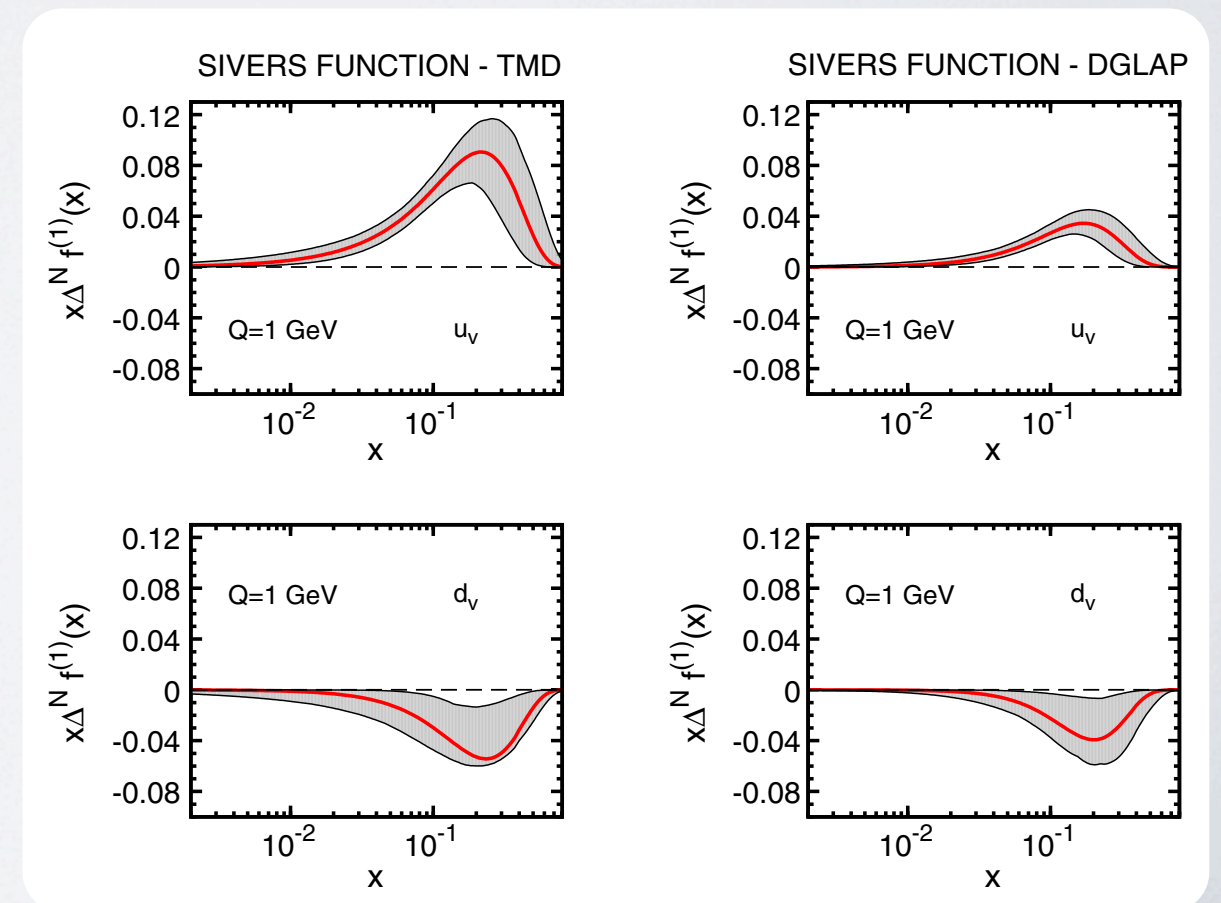
$$\Delta^N f_{q/p\uparrow}(x, k_T) = \mathcal{N}_q(x) h(k_T) f_1^q(x, k_T)$$

$$\Delta^N f_{q/p\uparrow} \equiv -\frac{2k_T}{M} f_{1T}^{\perp q}$$

## • Fits to HERMES and COMPASS:

### ► Current Data can only afford:

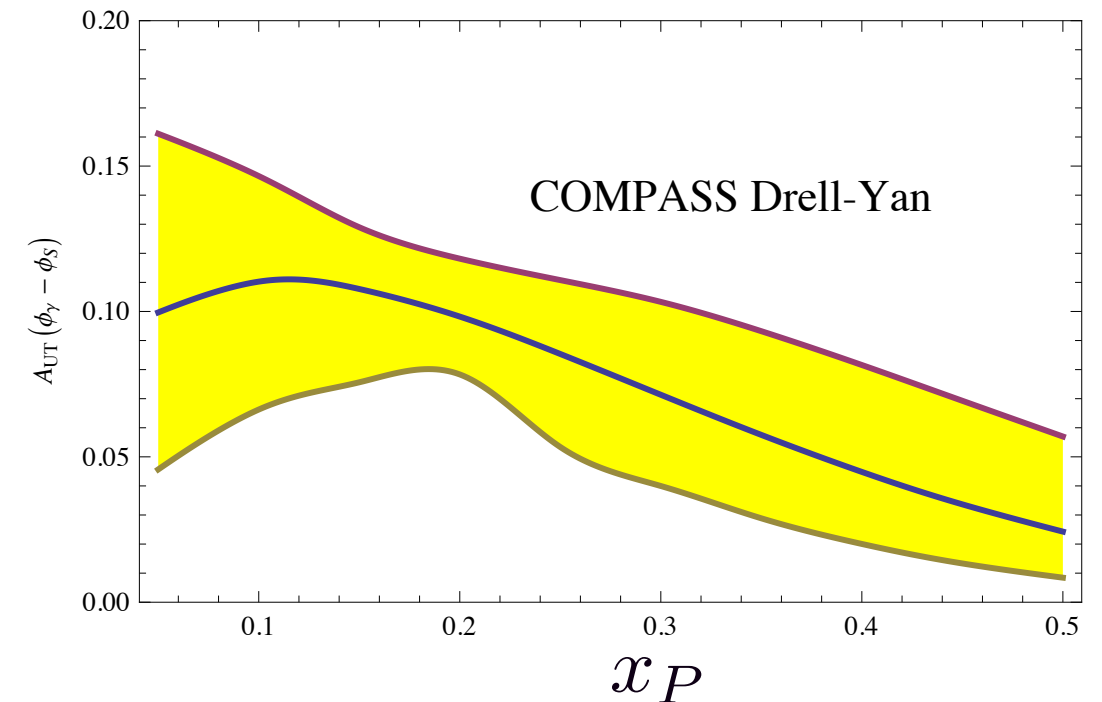
- **Large** uncertainties, esp. for sea.
- **Approximations:** TM and flavor dependence of FF, etc.



# EXTRACTIONS WITH TMD EVOLUTION

Sun, Yuan, PRD88 (2013), 114012

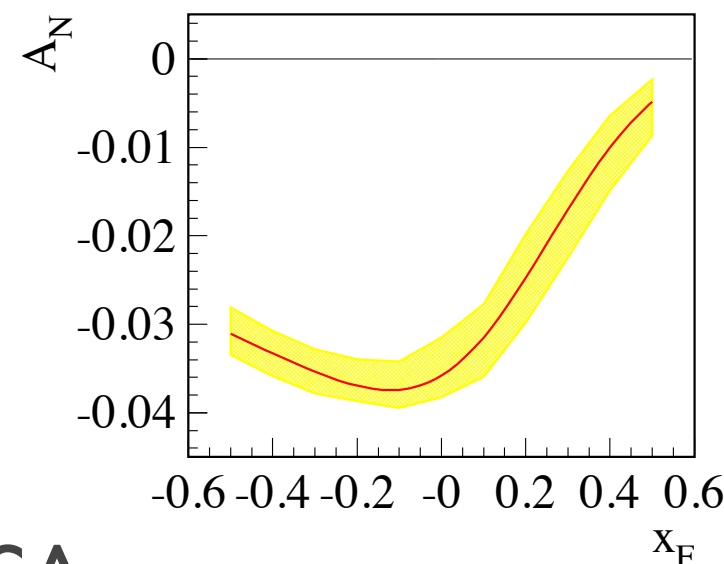
- ◆ Sun-Yuan prescription for TMD evolution.
- ◆ Gaussian TM dependence of NP TMD dependence at initial scale.
- ◆ Fit HERMES & COMPASS multiplicities and Sivers SSAs.
- ◆ Predict Sivers SSA and  $W$  production in COMPASS DY and PP.



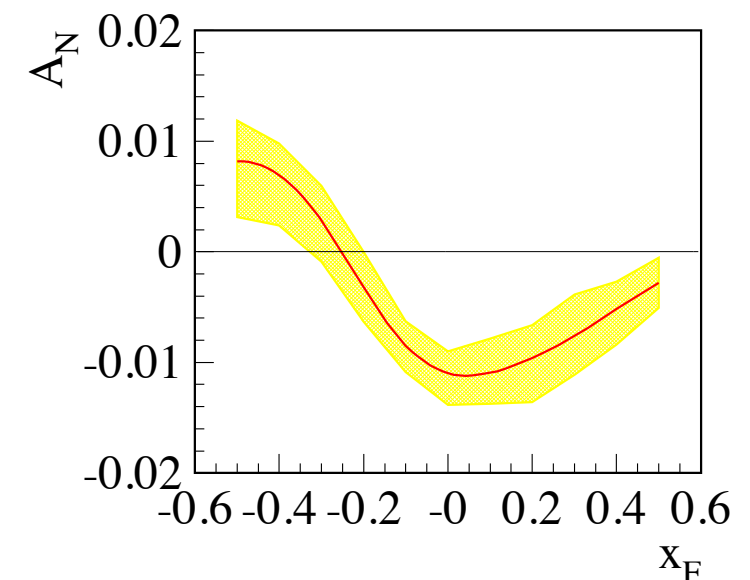
Echevarria et al.: PRD.89 074013, (2014)

- ◆ Find non-perturbative Sudakov factor that describes  $W, Z$  production in  $P\bar{P}$  at Fermilab +HERMES & COMPAS.
- ◆ Use it to fit Sivers SSA at HERMES, COMPASS, JLAB.
- ◆ Predict Sivers Effect for DY SSA.

**COMPASS**



**RHIC**



# EXTRACTIONS WITH TMD EVOLUTION

Sun, Yuan, PRD88 (2013), 114012

◆ Sun-Yuan prescription for TMD evolution.

◆ Gaussian TM dependence of NP TMD dependence

◆ Fit HERMES and COMPASS

◆ Predict COMPASS

## Approximations:

◆ *Gaussian TM ansatz.*

◆ *Flavor dependence.*

◆ *TMD Evolution prescription.*

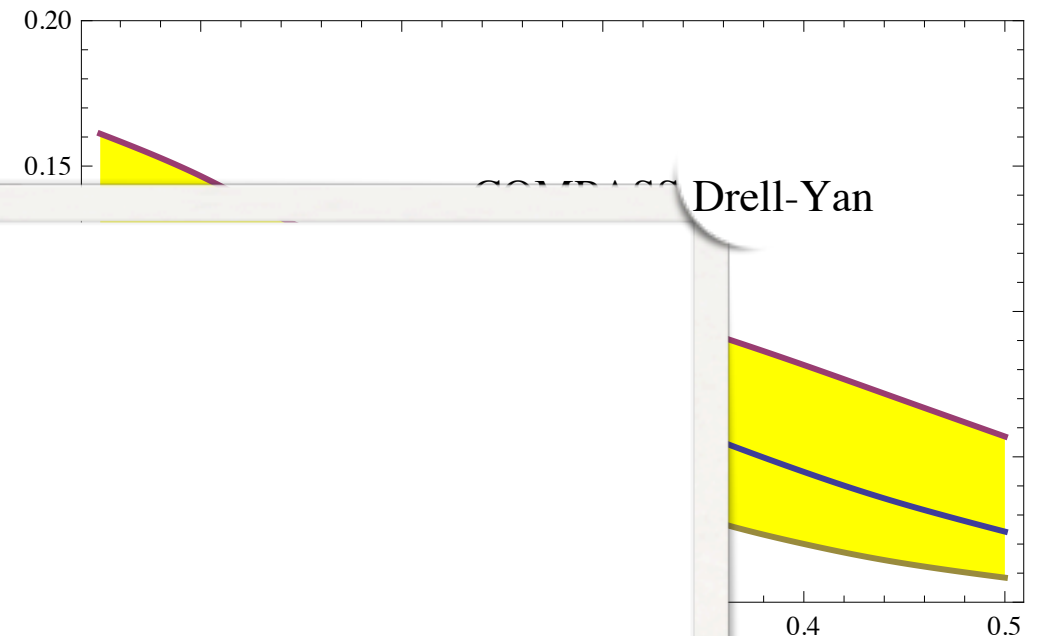
◆ *TM and flavor dependence of FF.*

Echevarria et al.

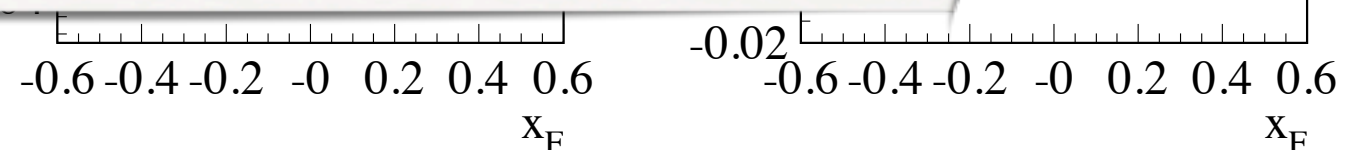
◆ Find non-perturbative factor that production + HERMES &

◆ Use it to fit HERMES, COMPASS, JLAB.

◆ Predict Sivers Effect for DY SSA.



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# LO APPROXIMATION FOR SSA

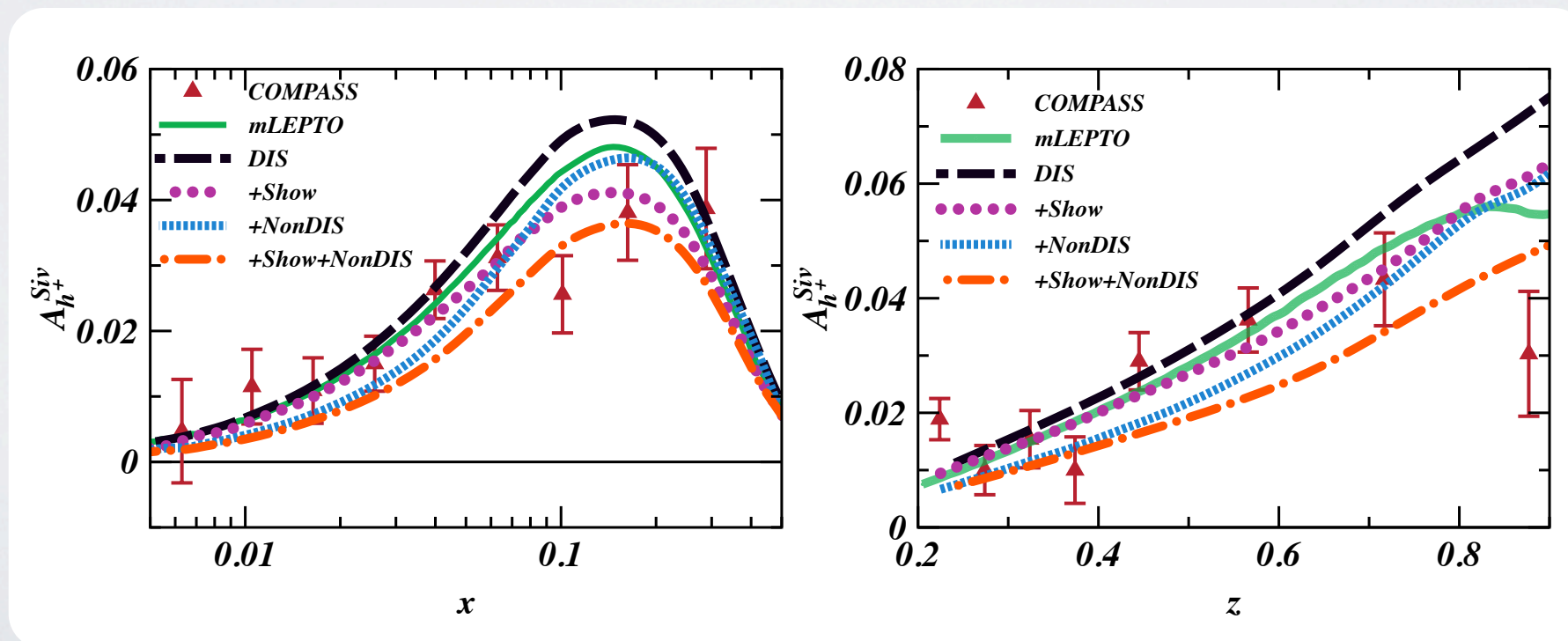
- ▶ Fits for *Sivers PDF* from HERMES and COMPASS data utilize *LO DIS-only* expressions for *SSAs*.

**M. Anselmino et. al: PRD 86, 014028 (2012).**

$$A_{UT}^{\sin(\phi_h - \phi_S)} = \frac{\sum_q \int d\phi_S d\phi_h d^2\mathbf{k}_\perp \Delta^N \hat{f}_{q/p^\uparrow}(x, k_\perp, Q) \sin(\varphi - \phi_S) \frac{d\hat{\sigma}^{\ell q \rightarrow \ell q}}{dQ^2} \hat{D}_q^h(z, p_\perp, Q) \sin(\phi_h - \phi_S)}{\sum_q \int d\phi_S d\phi_h d^2\mathbf{k}_\perp \hat{f}_{q/p}(x, k_\perp, Q) \frac{d\hat{\sigma}^{\ell q \rightarrow \ell q}}{dQ^2} \hat{D}_q^h(z, p_\perp, Q)}.$$

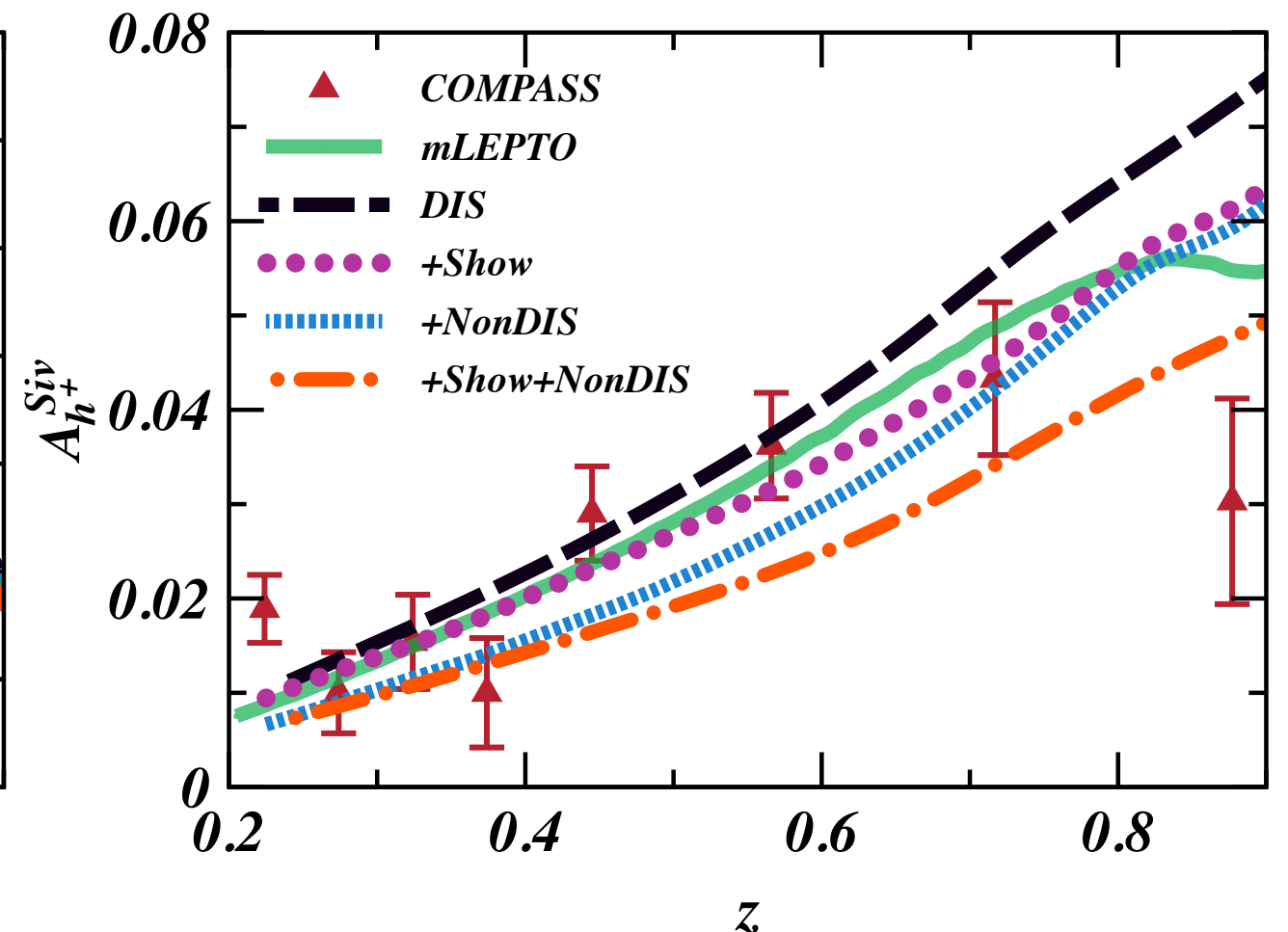
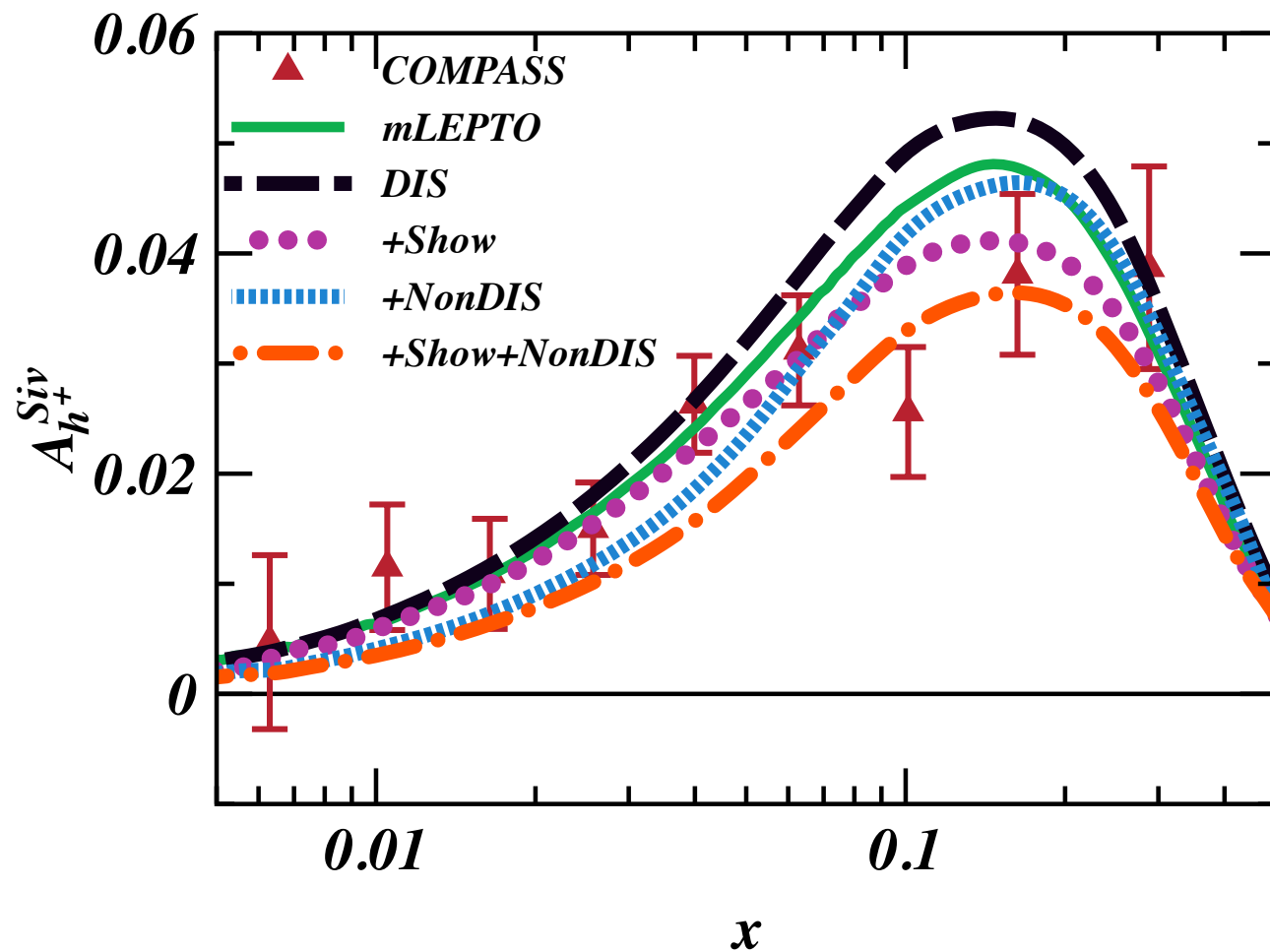
- ▶ **Is this justified at COMPASS energies?**
- ▶ Test using *mPYTHIA*: turn on non-DIS effects (*VMD*, *GVMD*, “direct”) and parton showering (*QCD+QED*).

**H.M et al., arXiv:1502.02669 (2015).**



# LO APPROXIMATION FOR SSA

H.M et al., arXiv:1502.02669 (2015).



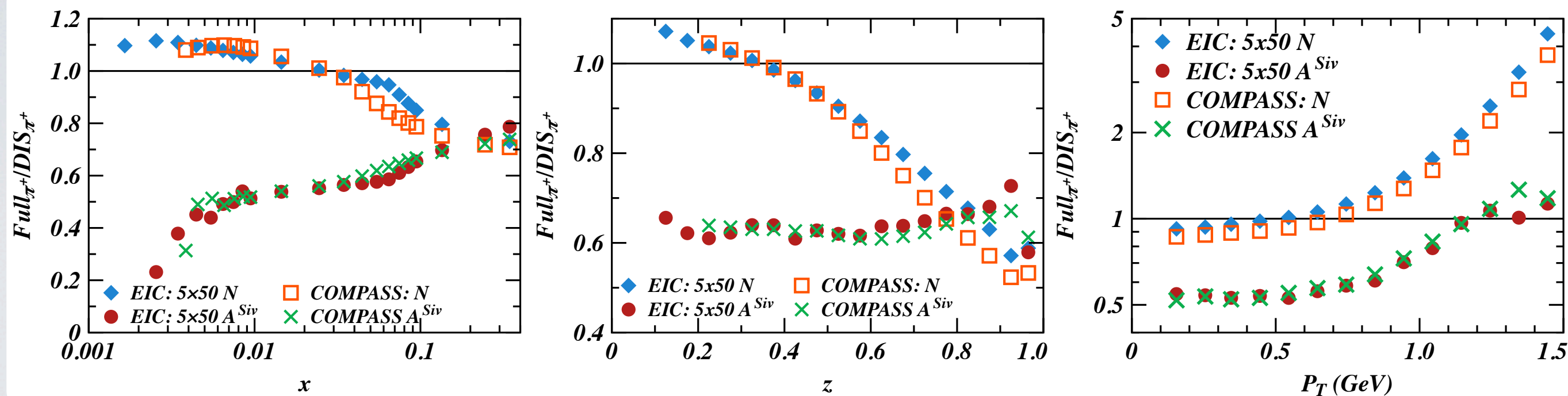
- ▶ **Significant** effects, but still agrees with data!
- ▶ Current Sivers PDF extractions *may* be underestimated.
- ▶ Note: no *model-independent* way to exclude non-DIS effects.



# Can We Still Use These Parametrizations?

H.M et al., arXiv:1502.02669 (2015).

- **How reliable are our SSA predictions for other experiments?**
- **Construct Ratios of Full (non-DIS + showers) to LO DIS results for multiplicities and Sivers SSAs at COMPASS and EIC.**



- The Ratios are **very close** between **COMPASS** and **EIC**:
- We can reliably estimate SSAs if we use only **LO DIS** terms with the **current parametrization** of Sivers PDFs.