

# Electroweak precision observables in the SM: present & future

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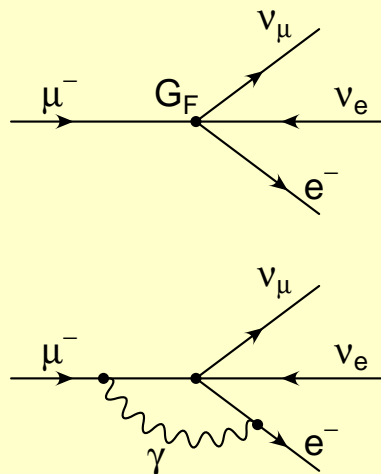
**University of Pittsburgh**

First FCC-ee Mini-Workshop on Precision Observables and Radiative Corrections,  
13-14 July 2015

- 1. Electroweak precision observables**
- 2. Current status of SM loop results**
- 3. Renormalization/parametrization**
- 4. Future projections**

## W mass

$\mu$  decay in Fermi Model



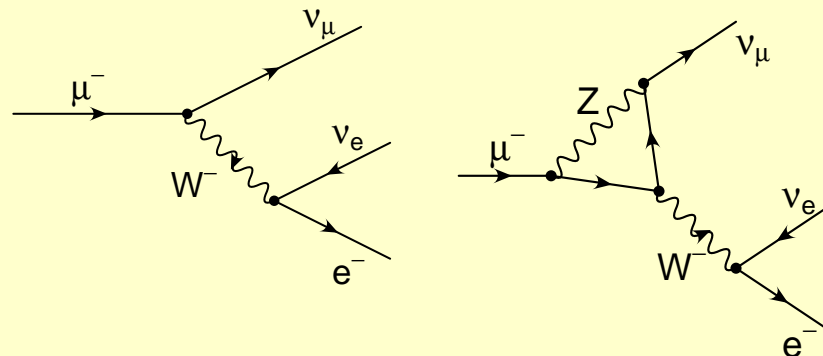
← QED corr.  
(2-loop)

$$\Gamma_\mu = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) (1 + \Delta q)$$

Ritbergen, Stuart '98

Pak, Czarnecki '08

$\mu$  decay in Standard Model



$$\frac{G_F^2}{\sqrt{2}} = \frac{e^2}{8s_w^2 M_W^2} (1 + \Delta r)$$

electroweak corrections

- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of  $\gamma$ -exchange,  $\gamma$ -Z interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- Z-pole contribution:

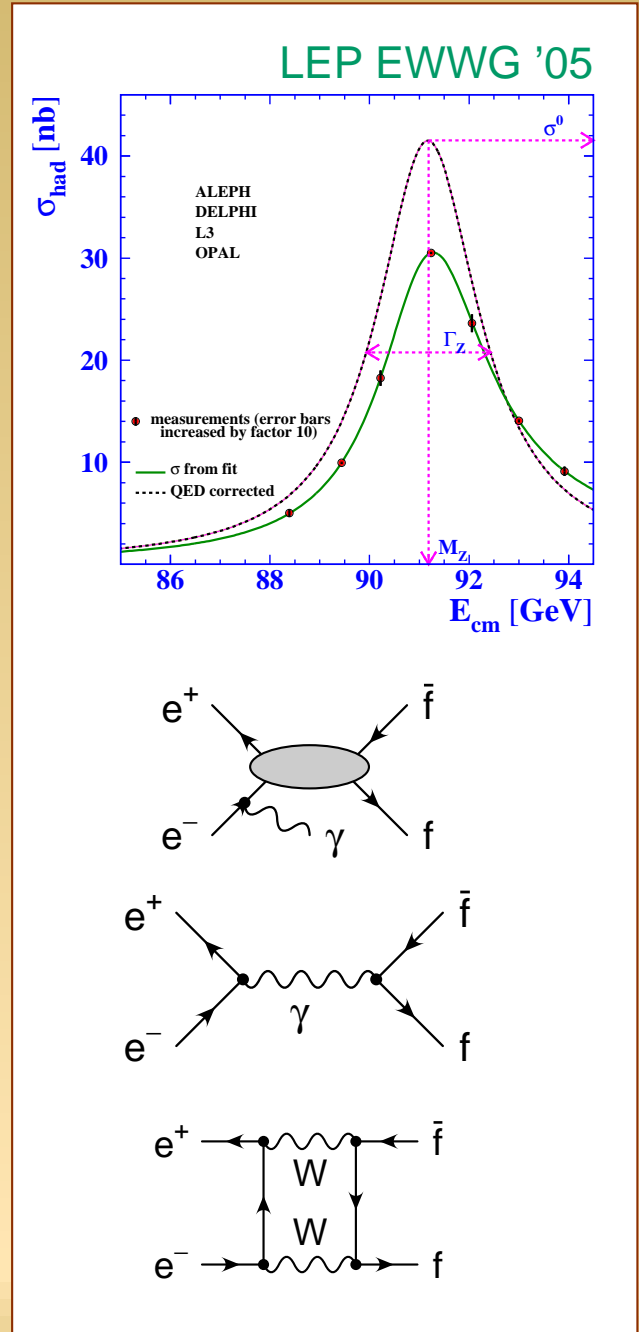
$$\sigma_Z = \frac{R}{(s - \overline{M}_Z^2)^2 + \overline{M}_Z^2 \overline{\Gamma}_Z^2} + \sigma_{\text{non-res}}$$

- In experimental analyses:

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2 \Gamma_Z^2 / M_Z^2}$$

$$\overline{M}_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

$$\overline{\Gamma}_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$



Z width:

$$\bar{\Gamma}_Z = \frac{1}{M_Z} \text{Im} \Sigma_Z(s_0).$$

Optical theorem:

$$\bar{\Gamma}_Z = \sum_f \bar{\Gamma}_f, \quad \bar{\Gamma}_f \approx \frac{N_c \bar{M}_Z}{12\pi} \left[ \left( \mathcal{R}_V^f |g_V^f|^2 + \mathcal{R}_A^f |g_A^f|^2 \right) \frac{1}{1 + \text{Re} \Sigma'_Z} \right]_{s=\bar{M}_Z^2}$$

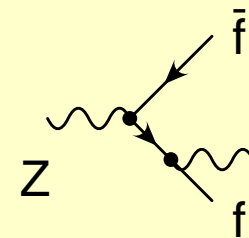
$\mathcal{R}_V^f, \mathcal{R}_A^f$ : Final-state QED/QCD radiation;

known to  $\mathcal{O}(\alpha_s^4), \mathcal{O}(\alpha^2), \mathcal{O}(\alpha\alpha_s)$

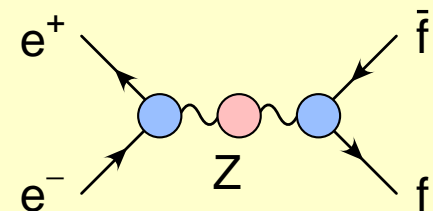
Kataev '92

Chetyrkin, Kühn, Kwiatkowski '96

Baikov, Chetyrkin, Kühn, Rittinger '12



$g_V^f, g_A^f, \Sigma'_Z$ : Electroweak corrections



Peak cross section:

$$\sigma_{\text{had}}^0 = \sigma_Z(s = \overline{M_Z^2})$$

(agrees with result from running-width BW with  $s = M_Z^2$ )

Explicit calculation:

$$\sigma_{\text{had}}^0 = \frac{12\pi}{\overline{M_Z^2}} \sum_q \frac{\overline{\Gamma}_e \overline{\Gamma}_q}{\overline{\Gamma_Z^2}} (1 + \delta X)$$

Correction term first at NNLO:

$$\delta X_{(2)} = -(\text{Im } \Sigma'_{Z(1)})^2 - 2\overline{\Gamma}_Z \overline{M}_Z \text{Im } \Sigma''_{Z(1)}$$

Grassi, Kniehl, Sirlin '01

Freitas '13

Branching ratios:

e.g.  $R_b = \Gamma_b / \Gamma_{\text{had}}$  (probes third quark generation)

Effective weak mixing angle:

Z-pole asymmetries:

$$A_{\text{FB}}^f \equiv \frac{\sigma(\theta < \frac{\pi}{2}) - \sigma(\theta > \frac{\pi}{2})}{\sigma(\theta < \frac{\pi}{2}) + \sigma(\theta > \frac{\pi}{2})} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f$$

$$A_{\text{LR}} \equiv \frac{\sigma(\mathcal{P}_e > 0) - \sigma(\mathcal{P}_e < 0)}{\sigma(\mathcal{P}_e > 0) + \sigma(\mathcal{P}_e < 0)} = \mathcal{A}_e$$

$$\mathcal{A}_f = 2 \frac{g_V^f / g_A^f}{1 + (g_V^f / g_A^f)^2} = \frac{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f}{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f + 8(|Q_f| \sin^2 \theta_{\text{eff}}^f)^2}$$

Most precisely measured for  $f = \ell$  (also  $f = b, c$ )

Known corrections to  $\Delta r$ ,  $\sin^2 \theta_{\text{eff}}^f$ ,  $g_{Vf}$ ,  $g_{Af}$ :

- Complete NNLO corrections ( $\Delta r$ ,  $\sin^2 \theta_{\text{eff}}^l$ ) Freitas, Hollik, Walter, Weiglein '00  
Awramik, Czakon '02; Onishchenko, Veretin '02  
Awramik, Czakon, Freitas, Weiglein '04; Awramik, Czakon, Freitas '06  
Hollik, Meier, Uccirati '05,07; Degrossi, Gambino, Giardino '14
- “Fermionic” NNLO corrections ( $g_{Vf}$ ,  $g_{Af}$ ) Czarnecki, Kühn '96  
Harlander, Seidensticker, Steinhauser '98  
Freitas '13,14
- Partial 3/4-loop corrections to  $\rho/T$ -parameter  
 $\mathcal{O}(\alpha_t \alpha_s^2)$ ,  $\mathcal{O}(\alpha_t^2 \alpha_s)$ ,  $\mathcal{O}(\alpha_t \alpha_s^3)$  Chetyrkin, Kühn, Steinhauser '95  
Faisst, Kühn, Seidensticker, Veretin '03  
Boughezal, Tausk, v. d. Bij '05  
Schröder, Steinhauser '05; Chetyrkin et al. '06  
Boughezal, Czakon '06

$$(\alpha_t \equiv \frac{y_t^2}{4\pi})$$

	Experiment	Theory error	Main source
$M_W$	$80.385 \pm 0.015$ MeV	4 MeV	$\alpha^3, \alpha^2\alpha_s$
$\Gamma_Z$	$2495.2 \pm 2.3$ MeV	0.5 MeV	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s, \alpha\alpha_s^2$
$\sigma_{\text{had}}^0$	$41540 \pm 37$ pb	6 pb	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$R_b \equiv \Gamma_Z^b / \Gamma_Z^{\text{had}}$	$0.21629 \pm 0.00066$	0.00015	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2\alpha_s$
$\sin^2 \theta_{\text{eff}}^l$	$0.23153 \pm 0.00016$	$4.5 \times 10^{-5}$	$\alpha^3, \alpha^2\alpha_s$

## Methods for theory error estimates:

- Parametric factors, *i. e.* factors of  $\alpha, N_c, N_f, \dots$
- Geometric progression, *e. g.*  $\frac{\mathcal{O}(\alpha^3)}{\mathcal{O}(\alpha^2)} \sim \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)}$
- Renormalization scale dependence (often underestimates error)
- Renormalization scheme dependence (may underestimate error)



## ■ Geometric perturbative series

$$\alpha_t = \alpha m_t^2$$

$$\mathcal{O}(\alpha^3) - \mathcal{O}(\alpha_t^3) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha^2) \sim 0.26 \text{ MeV}$$

$$\mathcal{O}(\alpha^2 \alpha_s) - \mathcal{O}(\alpha_t^2 \alpha_s) \sim \frac{\mathcal{O}(\alpha^2) - \mathcal{O}(\alpha_t^2)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.30 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s) \sim 0.23 \text{ MeV}$$

$$\mathcal{O}(\alpha \alpha_s^3) - \mathcal{O}(\alpha_t \alpha_s^3) \sim \frac{\mathcal{O}(\alpha \alpha_s) - \mathcal{O}(\alpha_t \alpha_s)}{\mathcal{O}(\alpha)} \mathcal{O}(\alpha \alpha_s^2) \sim 0.035 \text{ MeV}$$

$$\mathcal{O}(\alpha_{\text{bos}}^2) \sim \mathcal{O}(\alpha_{\text{bos}})^2 \sim 0.1 \text{ MeV}$$

■ Parametric prefactors:  $\mathcal{O}(\alpha_{\text{bos}}^2) \sim \Gamma_Z \alpha^2 \sim 0.1 \text{ MeV}$

$$\mathcal{O}(\alpha \alpha_s^2) - \mathcal{O}(\alpha_t \alpha_s^2) \sim \frac{\alpha n_{|q}}{\pi} \alpha_s^2 \sim 0.29 \text{ MeV}$$

**Total:**  $\delta \Gamma_Z \approx 0.5 \text{ MeV}$

## ■ Renormalization scheme dependence:

- a) Uncertainty of  $\mathcal{O}(\alpha^2)$  corrections beyond leading  $\alpha^2 m_t^4$  and  $\alpha^2 m_t^2$  from comparison of  $\overline{\text{MS}}$  and OS schemes: Degrassi, Gambino, Sirlin '96

$$\delta M_W \sim 2 \text{ MeV} \quad (\text{for } M_H \sim 100 \text{ GeV})$$

Actual remaining  $\mathcal{O}(\alpha^2)$  corrections: Freitas, Hollik, Walter, Weiglein '00

$$\delta M_W \sim 3 \text{ MeV} \quad (\text{for } M_H \sim 100 \text{ GeV})$$

- b) Estimate of missing  $\mathcal{O}(\alpha^3)$  corrections from comparison of  $\overline{\text{MS}}$  and OS results:

Awramik, Czakon, Freitas, Weiglein '03

Degrassi, Gambino, Giardino '14

$$\delta M_W \sim 4 \dots 5 \text{ MeV} \quad (\text{after accounting for } \mathcal{O}(\alpha_t \alpha_s^3) \text{ corrections})$$

→ Saturates previous  $\delta M_W$  estimate!

**Note:** Differences in (implicitly) resummed higher-order contributions

## Renormalization

Use of  $\overline{\text{MS}}$  renormalization for  $m_t$  reduces h.o. QCD corrections of  $\mathcal{O}(\alpha_t \alpha_s^n)$ :

loops (n+1)	$\Delta\rho_{(n)}^{\overline{\text{MS}}} / \left( \frac{3G_F \overline{m}_t^2}{8\sqrt{2}\pi^2} \right)$	$\Delta\rho_{(n)}^{\text{OS}} / \left( \frac{3G_F m_t^2}{8\sqrt{2}\pi^2} \right)$	
2	$-0.193 \left( \frac{\alpha_s}{\pi} \right)$	$-3.970 \left( \frac{\alpha_s}{\pi} \right)$	Djouadi, Verzegnassi '87 Kniehl '90
3	$-2.860 \left( \frac{\alpha_s}{\pi} \right)^2$	$-14.59 \left( \frac{\alpha_s}{\pi} \right)^2$	Avdeev, Fleischer, et al. '94 Chetyrkin, Kühn, Steinhauser '95
4	$-1.680 \left( \frac{\alpha_s}{\pi} \right)^3$	$-93.15 \left( \frac{\alpha_s}{\pi} \right)^3$	Schöder, Steinhauser '05 Chetyrkin, Faisst, Kühn, et al. '06 Boughezal, Czakon '06

No clear pattern of this kind known for  $\mathcal{O}(\alpha^n)$

→ Only few results available that allow direct comparison

e.g. Faisst, Kühn, Seidensticker, Veretin '03

Parametrization of perturbation series:  $\alpha$  vs.  $G_F$ ?

$G_F$  can resum some leading one-loop terms

$$\Delta\alpha \equiv 1 - \frac{\alpha(0)}{\alpha(M_Z)} \approx 0.059 \qquad \Delta\rho = \frac{3\alpha}{16\pi s^2 c^2} \frac{m_t^2}{M_Z^2}$$

**But:** Strong cancellations between  $\Delta\alpha$  and  $\Delta\rho$  terms beyond one-loop:

$$\begin{aligned} \Delta r_{\text{res}}^{(3)} &= (\Delta\alpha)^3 - 3(\Delta\alpha)^2 \left(\frac{c^2}{s^2} \Delta\rho\right) + 6(\Delta\alpha) \left(\frac{c^2}{s^2} \Delta\rho\right)^2 - 5 \left(\frac{c^2}{s^2} \Delta\rho\right)^3 \\ &\approx (2.05 \quad -3.40 \quad +3.74 \quad -1.72) \times 10^{-4} \\ &= 0.68 \times 10^{-4} \end{aligned}$$

→ Not *the* numerically leading contribution anymore

**ILC:**  $\sqrt{s} \approx M_Z$  with  $30 \text{ fb}^{-1}$

**FCC-ee:**  $\sqrt{s} \approx M_Z$  with  $4 \times 3000 \text{ fb}^{-1}$

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	Current exp.	ILC	FCC-ee	Current perturb.
$M_W$ [MeV]	15	3–5	$\sim 1$	4
$\Gamma_Z$ [MeV]	2.3	$\sim 1$	$\sim 0.1$	0.5
$R_b$ [ $10^{-5}$ ]	66	15	$\lesssim 5$	15
$\sin^2 \theta_{\text{eff}}^{\ell}$ [ $10^{-5}$ ]	16	1.3	0.3	4.5

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→ Existing theoretical calculations adequate for LEP/SLC/LHC,  
but not ILC/FCC-ee!

	ILC	FCC-ee	perturb. error with 3-loop <sup>†</sup>	Param. error ILC*	Param. error FCC-ee**
$M_W$ [MeV]	3–5	$\sim 1$	1	2.6	1
$\Gamma_Z$ [MeV]	$\sim 1$	$\sim 0.1$	$\lesssim 0.2$	0.5	0.06
$R_b$ [ $10^{-5}$ ]	15	$\lesssim 5$	5–10	$< 1$	$< 1$
$\sin^2 \theta_{\text{eff}}^\ell$ [ $10^{-5}$ ]	1.3	0.3	1.5	2	2

<sup>†</sup> **Theory scenario:**  $\mathcal{O}(\alpha\alpha_S^2)$ ,  $\mathcal{O}(N_f\alpha^2\alpha_S)$ ,  $\mathcal{O}(N_f^2\alpha^2\alpha_S)$   
 ( $N_f^n$  = at least  $n$  closed fermion loops)

Parametric inputs:

\* **ILC:**  $\delta m_t = 100$  MeV,  $\delta\alpha_S = 0.001$ ,  $\delta M_Z = 2.1$  MeV

\*\***FCC-ee:**  $\delta m_t \lesssim 50$  MeV,  $\delta\alpha_S = 0.0001$ ,  $\delta M_Z = 0.1$  MeV

also:  $\delta(\Delta\alpha) = 5 \times 10^{-5}$

## ■ Subtraction of QED radiation contributions

→ Known to  $\mathcal{O}(\alpha^2)$ ,  $\mathcal{O}(\alpha^3 L^3)$  for **ISR**,  
 $\mathcal{O}(\alpha^2)$  for **FSR** and  $\mathcal{O}(\alpha^2 L^2)$  for **A<sub>FB</sub>**

$$(L = \log \frac{s}{m_e^2})$$

Berends, Burgers, v.Neerven '88

Kniehl, Krawczyk, Kühn, Stuart '88

Beenakker, Berends, v.Neerven '89

Skrzypek '92; Montagna, Nicrosini, Piccinini '97

→  $\mathcal{O}(0.1\%)$  uncertainty on  $\sigma_Z$ ,  $A_{FB}$

→ Improvement needed for ILC/FCC-ee

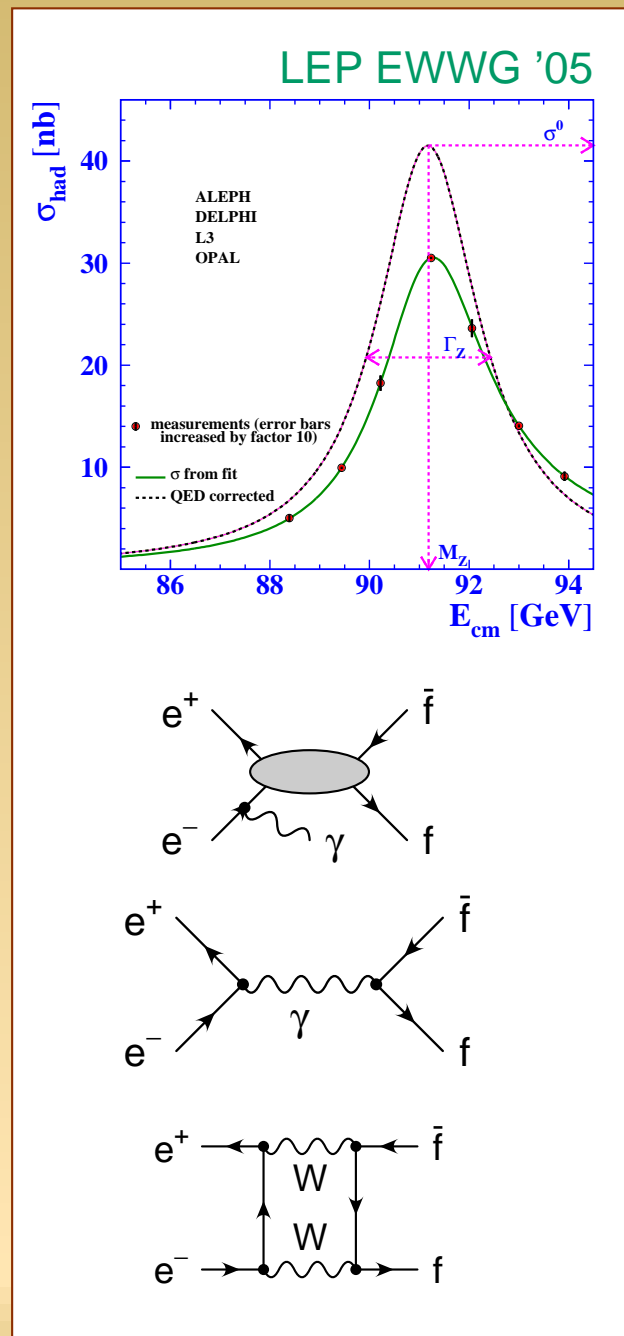
## ■ Subtraction of non-resonant $\gamma$ -exchange, $\gamma$ -Z interf., box contributions, Bhabha scattering

see, e.g., Bardin, Grünewald, Passarino '99

→  $\mathcal{O}(0.01\%)$  uncertainty within SM

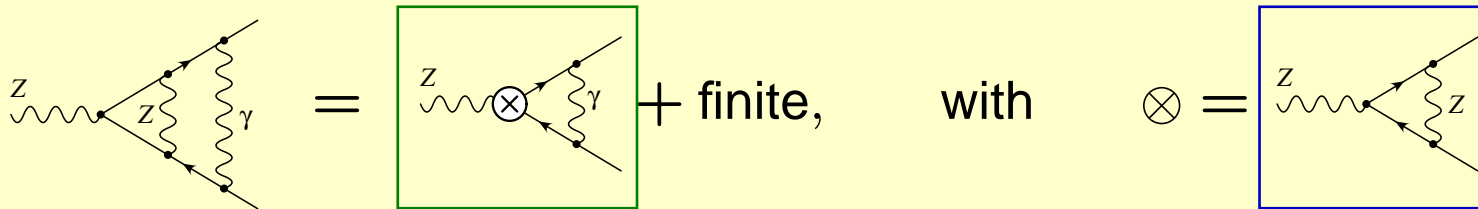
(improvements may be needed)

→ Sensitivity to some NP beyond EWPO

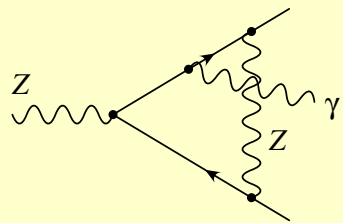


Factorization of massive and QED/QCD FSR:

$$\bar{\Gamma}_f \approx \frac{N_c \bar{M}_Z}{12\pi} \left[ \left( \mathcal{R}_V^f |g_V^f|^2 + \mathcal{R}_A^f |g_A^f|^2 \right) \frac{1}{1 + \text{Re} \Sigma'_Z} \right]_{s=\bar{M}_Z^2}$$



Additional non-factorizable contributions, e.g.



→ Known at  $\mathcal{O}(\alpha\alpha_s)$  Czarnecki, Kühn '96  
Harlander, Seidensticker, Steinhauser '98

→ Currently not known at  $\mathcal{O}(\alpha^2)$  and beyond

→  $\mathcal{O}(0.01\%)$  uncertainty on  $\Gamma_Z, \sigma_Z$ , maybe larger for  $A_b$   
 (improvements may be needed)



Full SM corrections at  $>2$ -loop:

- Large number of diagrams and tensor integrals,  $\mathcal{O}(1000) - \mathcal{O}(10000)$ 
  - Use of computer algebra tools
- Many different scales (masses and ext. momenta)
  - In general not possible analytically
  - Numerical methods must be automizable, stable, fastly converging
  - Need procedure for isolating divergent pieces

- Useful for diagrams with up to two scales  
(*e. g.*  $M_W$  &  $m_t$  or  $M_W$  &  $M_Z$ )
- Reduce to master integrals with integration-by-parts and Lorentz-invariance identities  
Chetyrkin, Tkachov '81; Gehrmann, Remiddi '00; Laporta '00; ...
- Evaluate master integrals with differential equations or Mellin-Barnes representations  
Kotikov '91; Remiddi '97; Smirnov '00,01; ...

Current status:

Single-scale problems:  $Z f \bar{f}$  QED/QCD vertex corrections up to 4-loop

Gorishnii, Kataev, Larin '88,91; Chetyrkin, Kühn, Kwiatkowski '96  
Baikov, Chetyrkin, Kühn, Rittinger '12

Two-scale problems:  $Z f \bar{f}$  electroweak 2-loop vertex diagrams with  $m_f = 0$

Awramik, Czakon, Freitas, Weiglein '04

Extendability: Possible, but much work needed

- Exploit large mass ratios, e. g.  $M_Z^2/m_t^2 \approx 1/4$
- Fast numerical evaluation

Current status:

**Two-scale problems:**  $\mathcal{O}(\alpha\alpha_s^n)$  corr. to  $\Delta\rho$ ,  $\Delta r$ , ...

→ Several expansion terms up to 3-loop, leading term up to 4-loop

Djouadi, Verzegnassi '87; Bardin, Chizhov '88

Chetyrkin, Kühn, Steinhauser '95

Faisst, Kühn, Seidensticker, Veretin '03

...

**Up to three-scale problems:**  $Zf\bar{f}$  ew. 2-loop vertex corrections

Barbieri et al. '92,93; Fleischer, Tarasov, Jegerlehner '93,95

Degrassi, Gambino, Sirlin '97; Awramik, Czakon, Freitas, Weiglein '04

**Extendability:** Promising, mostly limited by computing/algorithmic power

- Numerically integrate over cuts
- High precision, but no known path towards full automatization
- Subtraction of UV-divergencies by hand

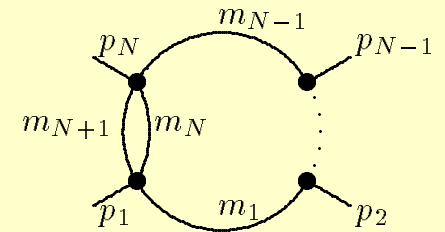
Example: Topologies with **self-energy sub-loop**

S. Bauberger et al. '95

$$B_0(p^2, m_1^2, m_2^2) = - \int_{(m_1+m_2)^2}^{\infty} ds \frac{\Delta B_0(s, m_1^2, m_2^2)}{s - p^2}$$

$$T_{N+1}(p_i; m_i^2) = - \int_{s_0}^{\infty} ds \Delta B_0(s, m_N^2, m_{N+1}^2)$$

$$\times \int d^4 q \frac{1}{q^2 - s} \frac{1}{(q+p_1)^2 - m_1^2} \cdots \frac{1}{(q+p_1+\cdots+p_{N-1})^2 - m_{N-1}^2}$$



- Numerically integrate over cuts
- High precision, but no known path towards full automatization
- Subtraction of UV-divergencies by hand

Current status:

Self-energy and vertex diagrams with arbitrary number of scales

Freitas, Hollik, Walter, Weiglein '00; Awramik, Czakon '02; Awramik, Czakon, Freitas '04

Extendability: Only for certain applications

General form of Feynman integral:

$$I = \int_0^1 dx_1 \dots dx_n \delta(1 - \sum_i x_i) \frac{N(x_i)}{D(x_i)^{r+\varepsilon}}$$

→ Can be integrated numerically (if finite)

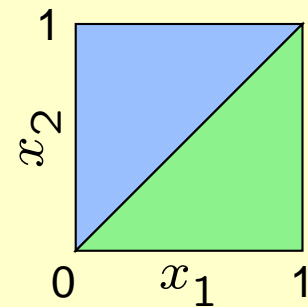
Alternatives: Integration in momentum space, Mellin-Barnes space

Treatment of divergencies:

- **Sector decomposition:** Sub-divide integration space such that divergent terms factorize

Binoth, Heinrich '00,03

- **Subtraction terms:** Remove divergencies with simple terms that can be integrated analytically



Nagy, Soper '03

Becker, Reuschle, Weinzierl '10; Freitas '12

- Automizable, but computing intensive
- Internal thresholds reduce numerical convergence

Current status:

Several 2-loop applications with many scales

Anastasiou et al., Petriello et al., Borowka et al., .....

Individual 3-loop integrals

Extendability: Likely, but more work needed

- **Current SM predictions** for electroweak precision observables under good control (compared to experimental uncertainties)
- **LHC** will provide independent results for  $\sin^2 \theta_{\text{eff}}$  and  $M_W$ , but overall precision not substantially improved
- **ILC/FCC-ee** with  $\sqrt{s} \sim M_Z$  will reduce exp. error of some EWPO by  $\mathcal{O}(10)$   
→ 3-loop (and maybe some 4-loop) corrections needed!
- **Asymptotic expansion and numerical integration techniques** are promising but more work needed
- **Open questions** in evaluation of theory errors, resummation and optimal choice of inputs

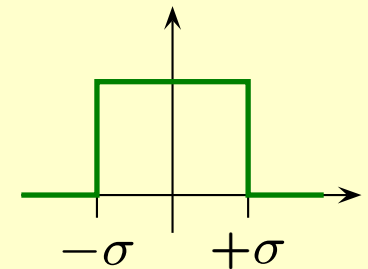


Backup slides

- Add theory errors from each source **linearly**:

Idea: each value within error range is equally likely

→ Use flat prior in global fits



- Add theory errors from each source **quadratically**:

Idea: different error sources are uncorrelated

→ Use Gaussian prior in global fits (central limit theorem)

