







HADRON PRODUCTION WITHIN PHSD

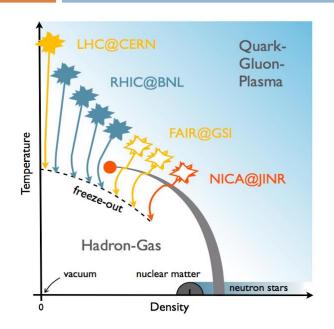
Pierre Moreau

for the PHSD group

XI Workshop on Particle Correlations and Femtoscopy 2015, Warsaw, Poland

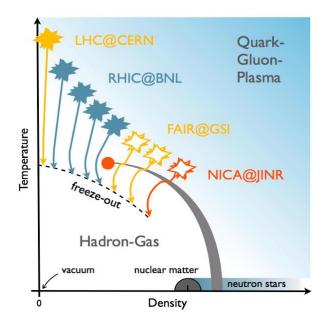


From NICA to LHC, passing by FAIR and RHIC...

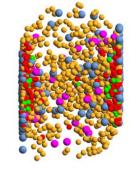


- Explore the QCD phase diagram and properties of hadrons at high temperature or high baryon density
- Phase transition from hadronic to partonic matter
- Goal: Study the properties of strongly interacting matter under extreme conditions from a microscopic point of view

From NICA to LHC, passing by FAIR and RHIC...



- Explore the QCD phase diagram and properties of hadrons at high temperature or high baryon density
- Phase transition from hadronic to partonic matter
- Goal: Study the properties of strongly interacting matter under extreme conditions from a microscopic point of view
- Realization: dynamical many-body transport approach
- Explicit parton-parton interactions, explicit phase transiton from hadronic to partonic degrees of freedom
- Transport theory: off-shell transport equations in phase-space representation based on Kadanoff-Baym equations for the partonic and hadronic phase





Parton-Hadron-String-Dynamics (PHSD)

W.Cassing, E.Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; W.Cassing, EPJ ST 168 (2009) 3

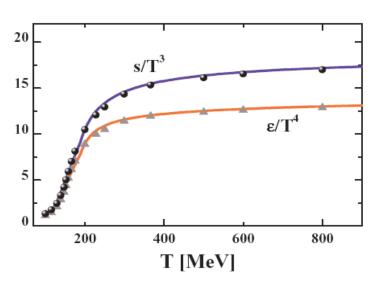
Dynamical Quasi-Particle Model (DQPM)

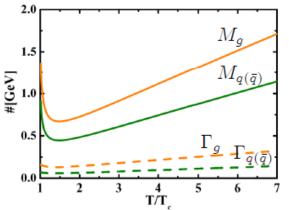
The QGP phase is described in terms of interacting quasiparticles with Lorentzian spectral functions:

$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{(\omega^2 - \mathbf{p}^2 - M_i^2(T))^2 + 4\omega^2\Gamma_i^2(T)} \qquad (i = q, \bar{q}, g)$$

 $\rho \, [\text{GeV}^2] \\ 10 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 0 \\ 1 \\ 10^{-2} \\ 10^{-3} \\ 10^{-4} \\ 10$

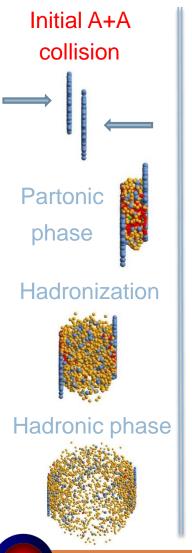
Properties of quasiparticles are fitted to the lattice QCD results:



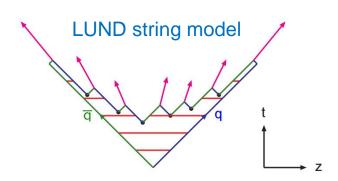


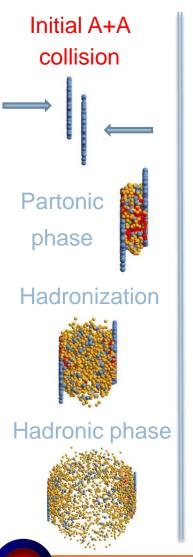
 Masses and widths of partons depend on the temperature of the medium

Peshier, Cassing, PRL 94 (2005) 172301; Cassing, NPA 791 (2007) 365: NPA 793 (2007)

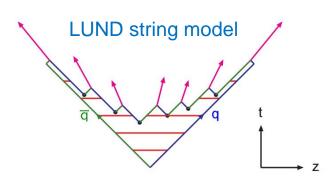


- String formation in primary NN collisions
- String decays to pre-hadrons (baryons and mesons)





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- Formation of a QGP state if $\epsilon > \epsilon_{\rm C} = 0.5 \; {\rm GeV.fm}^{-3}$
- Dissolution of new produced secondary hadrons into massive colored quarks and mean-field energy

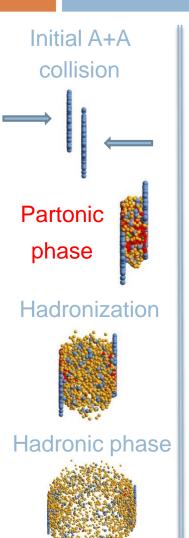
$$B o qqq$$
 ($\bar{q}\bar{q}\bar{q}$), $m o q\bar{q} + U_q$

DQPM define the properties (masses and widths) of partons

$$m_q(\epsilon)$$
 $\Gamma_q(\epsilon)$

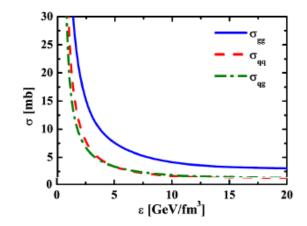
... and mean-field potential at a given local energy density ε

$$U_q(\epsilon)$$



- Propagation of partons, considered as dynamical quasiparticles, in a self-generated mean-field potential from the DQPM
- EoS of partonic phase: ,crossover' from Lattice QCD fitted by DQPM
- (quasi-)elastic collisions :

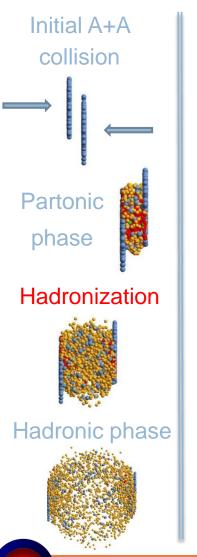
$$q+q \rightarrow q+q$$
 $g+q \rightarrow g+q$
 $q+\overline{q} \rightarrow q+\overline{q}$ $g+\overline{q} \rightarrow g+\overline{q}$
 $\overline{q}+\overline{q} \rightarrow \overline{q}+\overline{q}$ $g+g \rightarrow g+g$



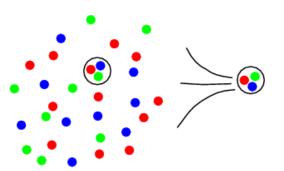
inelastic collisions :

$$q + \overline{q} \rightarrow g$$
 $q + \overline{q} \rightarrow g + g$
 $g \rightarrow q + \overline{q}$ $g \rightarrow g + g$

Suppressed due to the large gluon mass



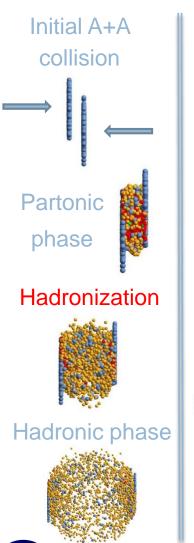
 Massive and off-shell (anti-)quarks hadronize to colorless off-shell mesons and baryons



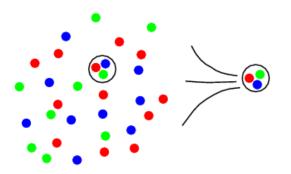
$$g \rightarrow q + \overline{q}, \quad q + \overline{q} \leftrightarrow meson \ ('string')$$

 $q + q + q \leftrightarrow baryon \ ('string')$

- Local covariant off-shell transition rate
- Strict 4-momentum and quantum number conservation



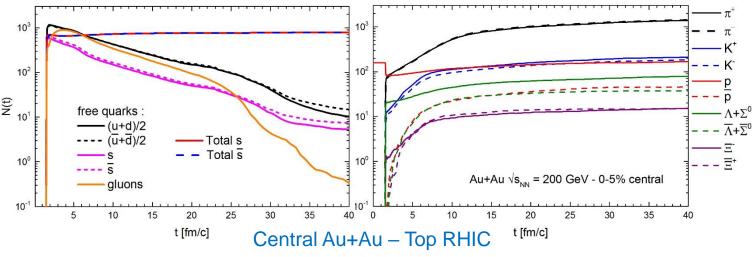
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$$g \rightarrow q + \overline{q}$$
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- Local covariant off-shell transition rate
- Strict 4-momentum and quantum number conservation

Number of partons and hadrons as a function of time:



t = 0.1 fm/c





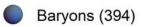
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Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}
b = 2.2 fm - Section view
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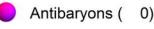
- Baryons (394)
- Antibaryons(0)
- Mesons (0)
- Quarks (0)
- Gluons (0)

t = 1.63549 fm/c



 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view







Quarks (4383)

Gluons (344)



P.Moreau

t = 2.06543 fm/c



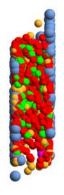


 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (396)
- Antibaryons (2)
- Mesons (1136)
- Quarks (5066)
- Gluons (516)

P.Moreau

t = 3.20258 fm/c

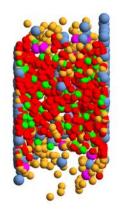




 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (413)
- Antibaryons(13)
- Mesons (1080)
- Quarks (4708)
- Gluons (761)

t = 5.56921 fm/c



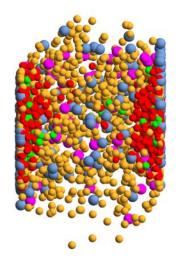


 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (472)
- Antibaryons(70)
- Mesons (1724)
- Quarks (3843)
- Gluons (652)

P.Moreau

t = 8.06922 fm/c



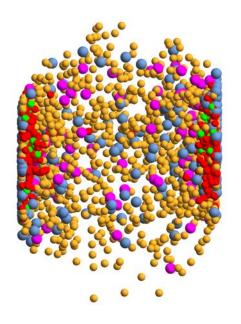


 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (559)
- Antibaryons (139)
- Mesons (2686)
- Quarks (2628)
- Gluons (442)

P.Moreau

t = 10.5692 fm/c





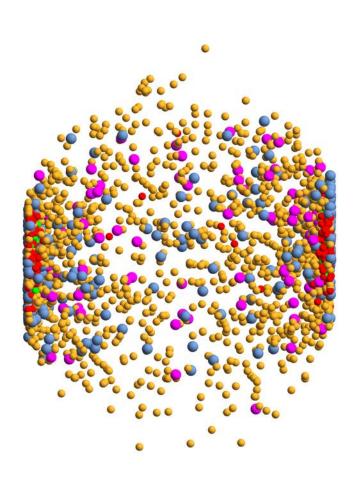
 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (604)
- Antibaryons (187)
- Mesons (3169)
- Quarks (2076)
- Gluons (319)

P.Moreau

Pierre Moreau

t = 15.5692 fm/c



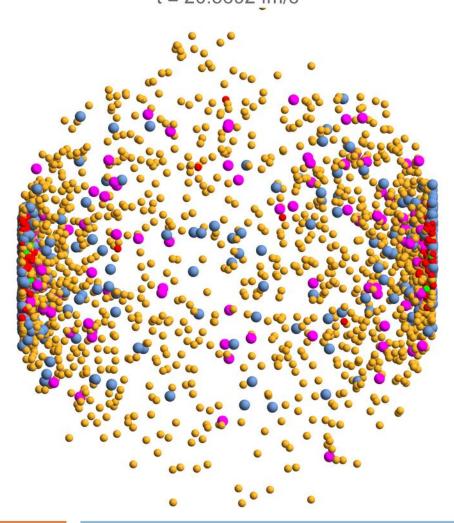


 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (662)
- Antibaryons (229)
- Mesons (3661)
- Quarks (1499)
- Gluons (175)

P.Moreau





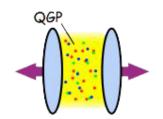


 $Au + Au \sqrt{s_{NN}} = 200 \text{ GeV}$ b = 2.2 fm - Section view

- Baryons (692)
- Antibaryons (266)
- Mesons (4022)
- Quarks (1184)
- Gluons (90)

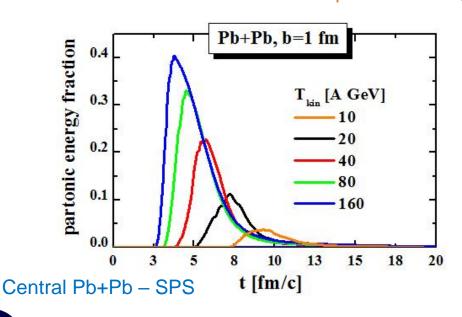
Partonic energy fraction in central A+A

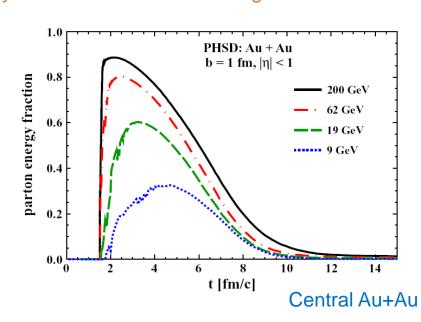
 At SPS, only a small part of the initial energy is converted into the QGP phase



 At top RHIC energies, the QGP phase at midrapidity contains roughly 90% of the energy

Time evolution of the partonic energy fraction for different energies:

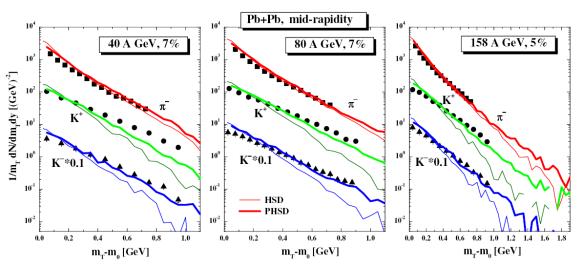


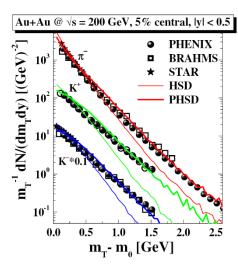


Transverse mass spectra (PHSD – HSD)

- With the HSD model, the high-pT spectra is not described properly especially at high energies where the parton energy fraction is major
- At low SPS energies, the difference is less visible since the partonic phase is not predominant

Transverse mass spectra for pions and kaons at different energies:





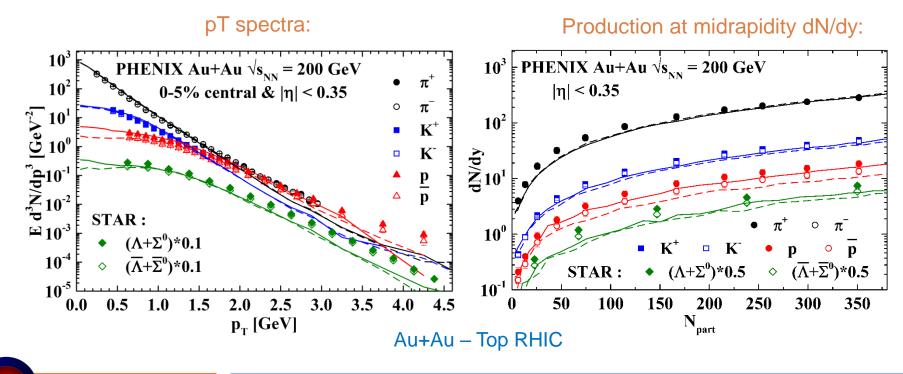
Central Pb+Pb – SPS energies

Central Au+Au - RHIC

W. Cassing & E. Bratkovskaya, NPA 831 (2009) 215; E. Bratkovskaya, W. Cassing, V. Konchakovski, O. Linnyk, NPA856 (2011) 162

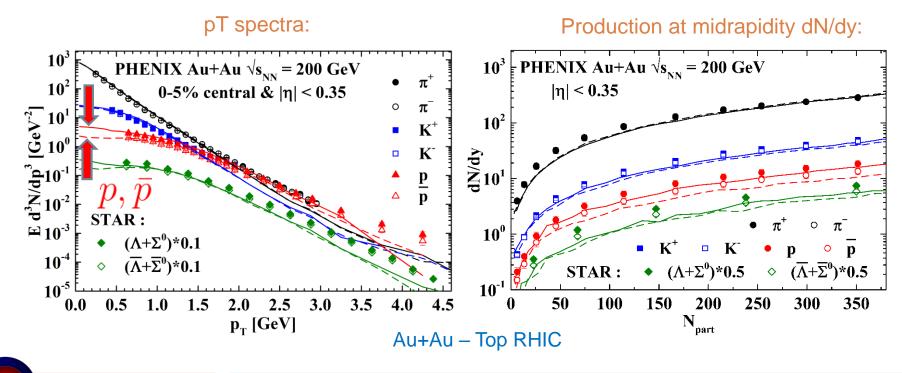
Au-Au at Top RHIC energies

- At high energies, particles and antiparticles are produced in quasi-equal quantities at midrapidity whatever the centrality of the collision
- Anti-baryon absorption at low pT is visible



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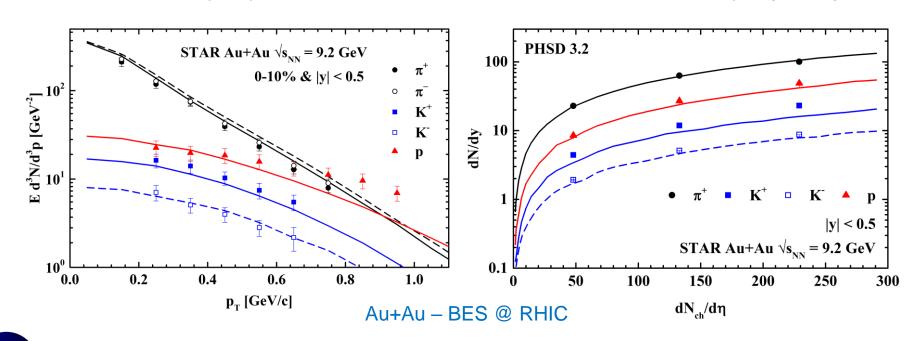


Au-Au at BES @ RHIC energies

At low energies, a clear difference appears between the production of particles and antiparticles, and also between positively and negatively charged mesons

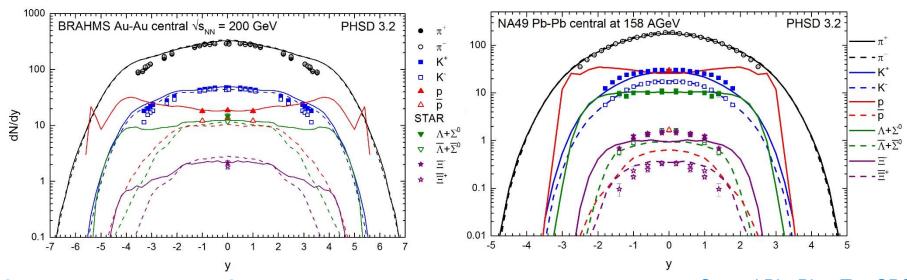


Production at midrapidity dN/dy:



- At high energies, the hadrons produced at midrapidity come mostly from the QGP phase
- At high rapidity, particles are more produced than antiparticles due to the high baryon density
- At low energies, the stopping of initial nucleons induces a high baryon density even at midrapidty which favors the production of baryons compared to antibaryons

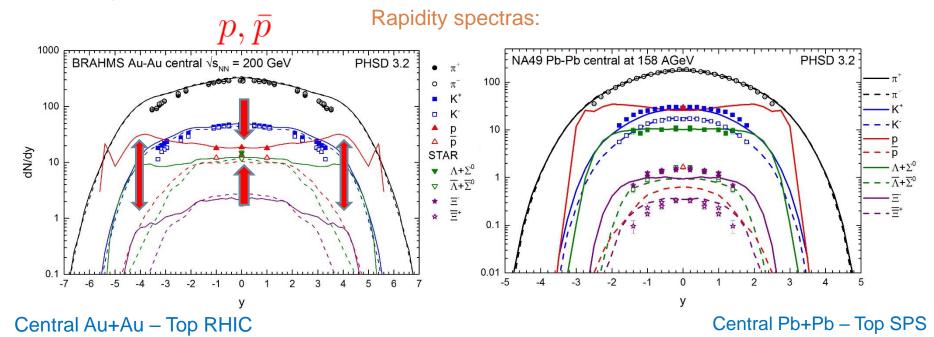
Rapidity spectras:



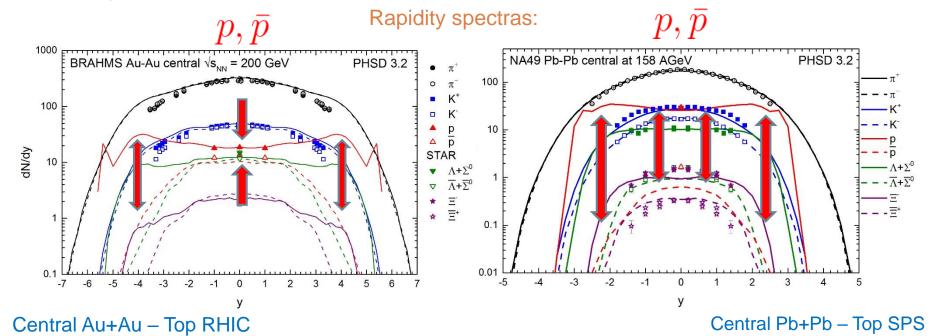
Central Au+Au - Top RHIC

Central Pb+Pb – Top SPS

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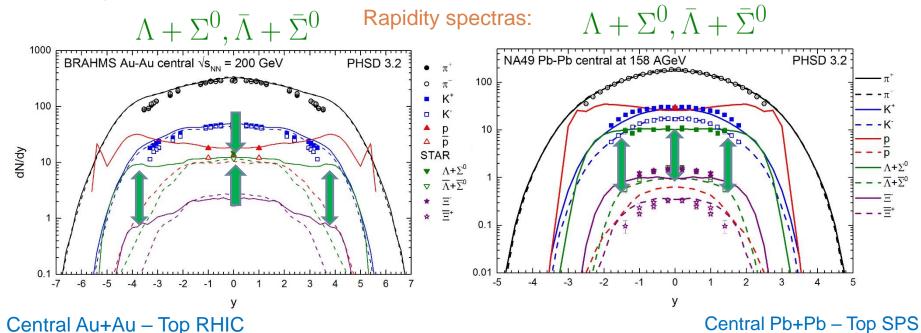


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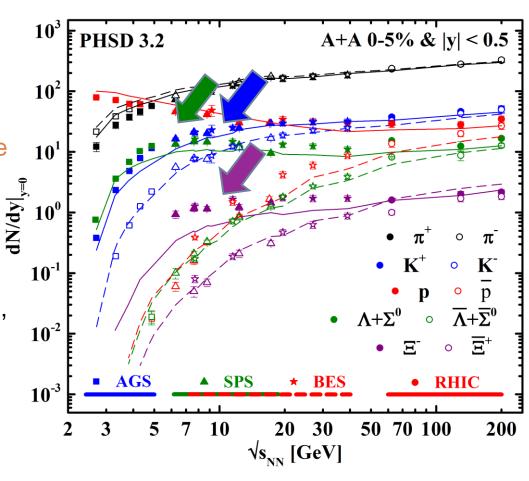


Beam energy scan study

Production at midrapidity as a function of the collisional energy:

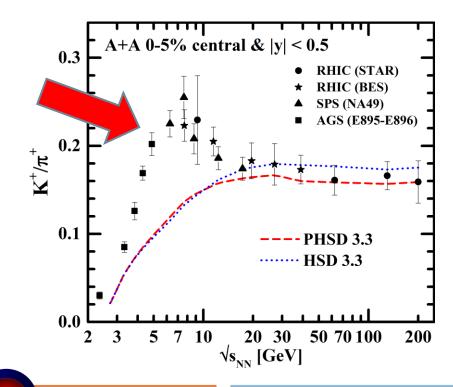
Reasonable agreement for anti-strange baryons dominantly produced in the hadronization process from the QGP at midrapidity

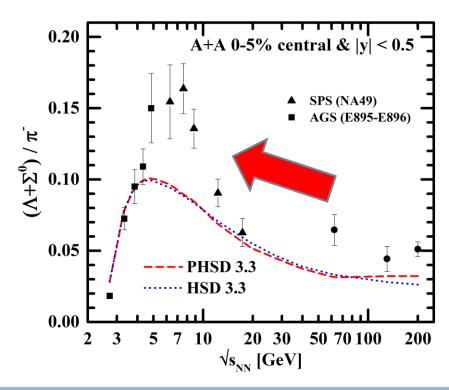
Underestimation of strange baryons at AGS-SPS energies, mainly produced by hadronic processes



Missing strangeness?

- Even considering the creation of a QGP phase, the strangeness enhancement seen experimentally at FAIR/NICA energies remains puzzling
 - 'Horn' not traced back to deconfinement





Conclusion

Production of quarks by string decays

According to a Schwinger-like formula, the probability to form a massive $s\bar{s}$ in a string-decay process is suppressed in comparison to light flavor $(u\bar{u}, dd)$

$$\frac{P(s\bar{s})}{P(u\bar{u})} = \frac{P(s\bar{s})}{P(d\bar{d})} = \gamma_s = \exp\left(-\pi \frac{m_s^2 - m_q^2}{2\kappa}\right)$$

Considering a hot and dense medium, the above formula remains the same but effective quark masses should be employed. This dressing is due to a scalar coupling with the in-medium quark condensate $\langle q\overline{q}\rangle$ according to:

$$m_s^* = m_s^0 + (m_s^v - m_s^0) \frac{\langle q\bar{q}\rangle}{\langle q\bar{q}\rangle_V} \qquad m_q^* = m_q^0 + (m_q^v - m_q^0) \frac{\langle q\bar{q}\rangle}{\langle q\bar{q}\rangle_V}$$

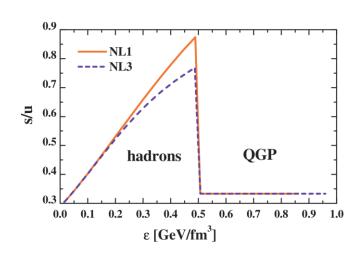
Chiral symmetry restoration in the hadronic phase

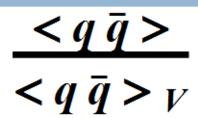
The scalar quark condensate $\langle q\bar{q}\rangle$ is viewed as an order parameter for the restoration of chiral symmetry at high baryon density and temperature. It can be expressed by the following formula:

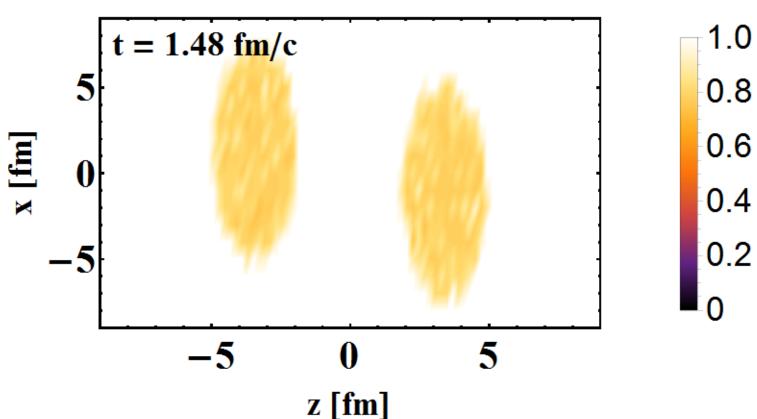
$$\frac{\langle q\bar{q}\rangle}{\langle q\bar{q}\rangle_V} = 1 - \frac{\Sigma_\pi}{f_\pi^2 m_\pi^2} \rho_S - \sum_h \frac{\sigma_h \rho_S^h}{f_\pi^2 m_\pi^2}$$

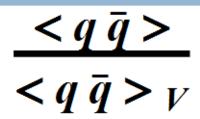
where ρ_s is the scalar density obtained according to the non-linear $\sigma - \omega$ model, $\Sigma_{\pi} \approx 45 \text{ MeV}$ is the pion-nucleon Σ -term, and f_{π} and m_{π} are the pion decay constant and pion mass, given by the Gell-Mann-Oakes-Renner relation.

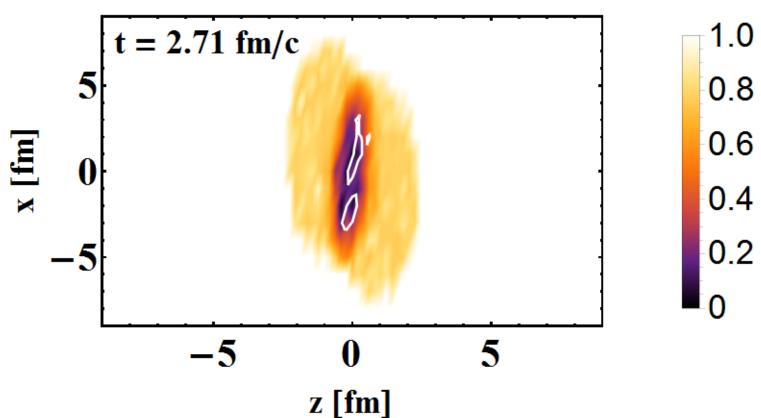
As a consequence of the **chiral symmetry** restoration (CSR), the strangeness production probability increases with the energy density ε . In the QGP phase, the string decay doesn't occur anymore and this effect is therefore suppressed.



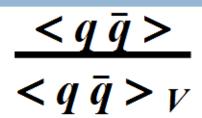


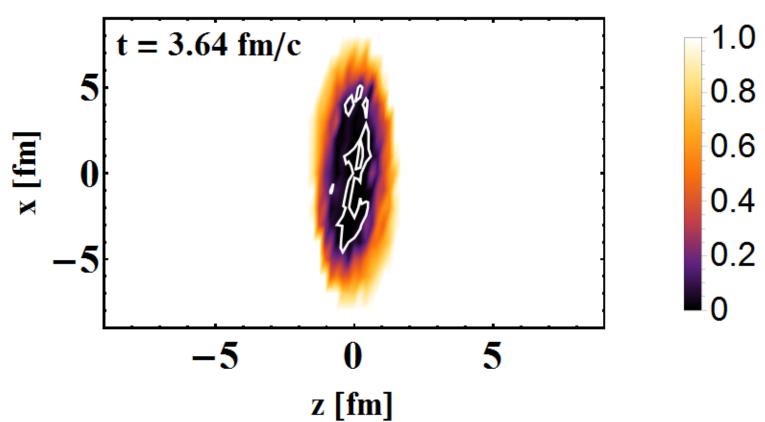




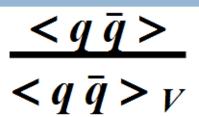


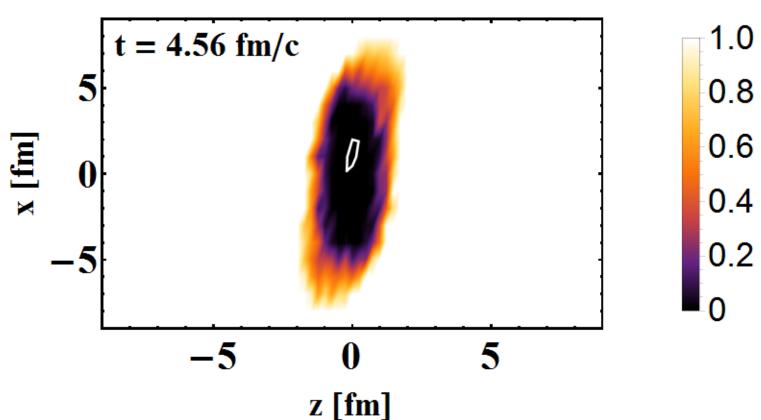




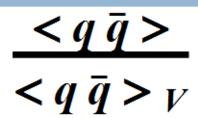


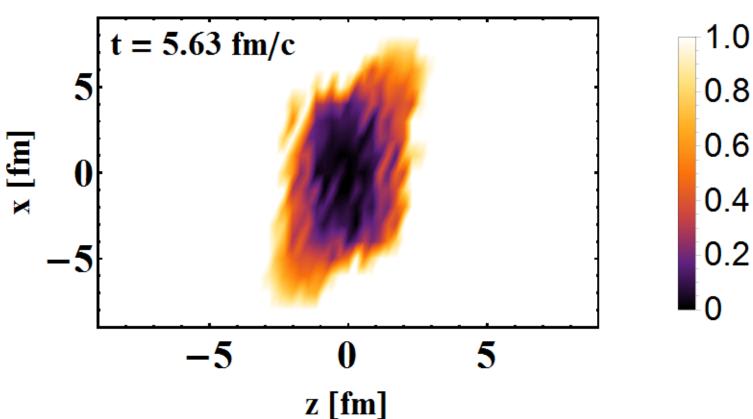


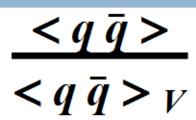


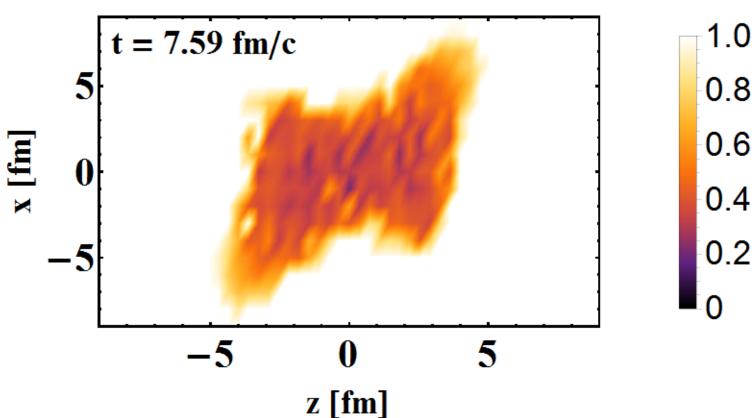




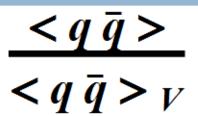


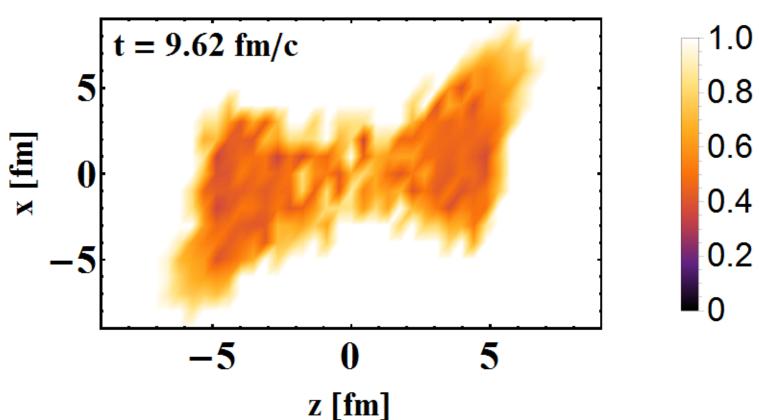




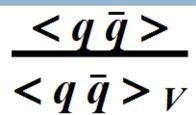


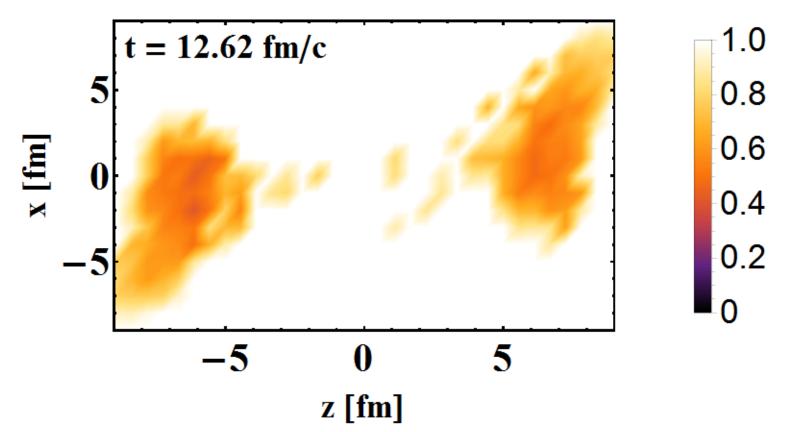








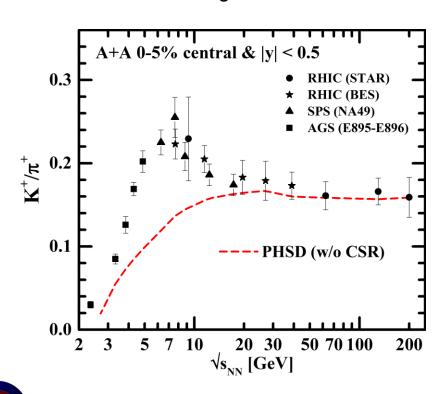


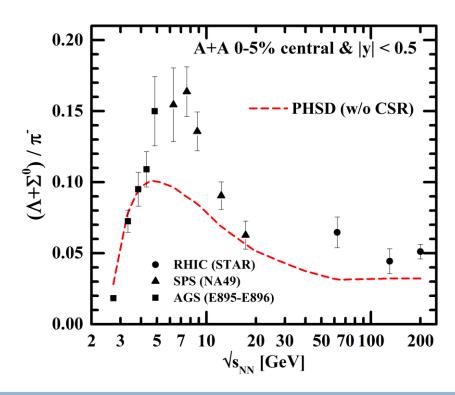


Chiral symmetry restoration in the hadronic phase

The strangeness enhancement seen experimentally at FAIR/NICA energies probably involves the approximate restoration of chiral symmetry in the hadronic phase

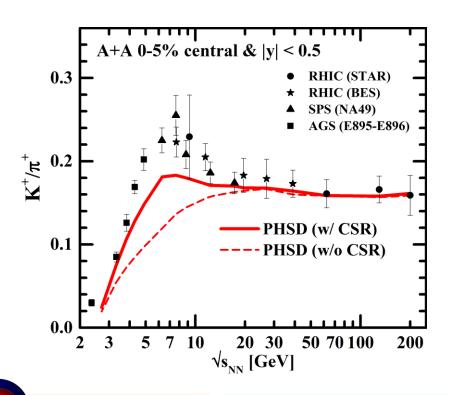
W. Cassing, A. Palmese, P. Moreau, E.L. Bratkovskaya - arXiv:1510.04120

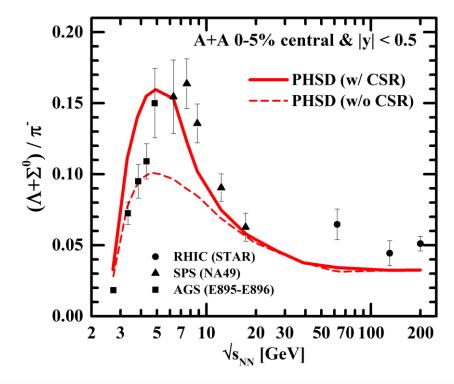




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W. Cassing, A. Palmese, P. Moreau, E.L. Bratkovskaya - arXiv:1510.04120



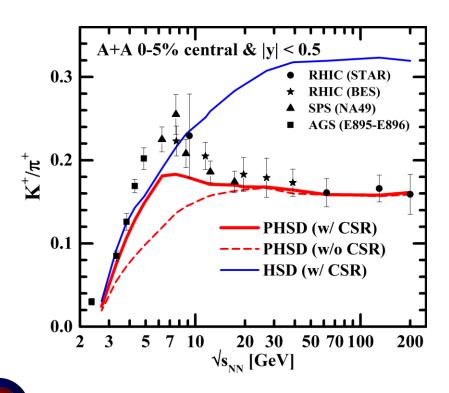


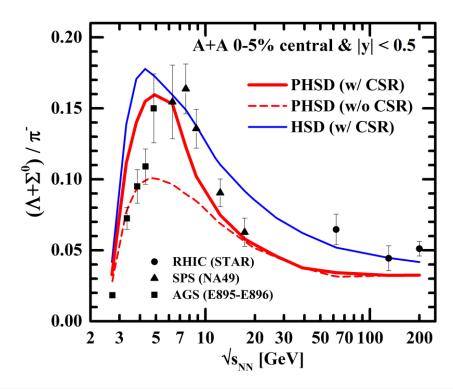
Conclusion

Chiral symmetry restoration in the hadronic phase

The strangeness enhancement seen experimentally at FAIR/NICA energies probably involves the approximate restoration of chiral symmetry in the hadronic phase

W. Cassing, A. Palmese, P. Moreau, E.L. Bratkovskaya - arXiv:1510.04120





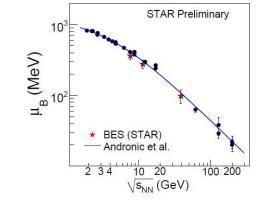
Conclusion

 At high energies, particles and antiparticles are produced in quasi-equal quantities at midrapidity in the hadronization process from the QGP phase

By decreasing the collisional energy, more differences appear between the

production of particle and antiparticle

Cross sections from the DQPM at finite chemical potential may also play a significant role at low collisional energy



Including aspects of chiral symmetry restoration in the hadronic phase, we observe a rise in the K^+/π^+ ratio at low $\sqrt{s_{NN}}$ and then a drop due to the appearance of a partonic medium

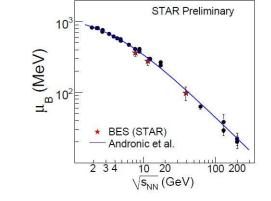
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THANK YOU FOR YOUR ATTENTION

