

# Introduction to Monte Carlo Techniques

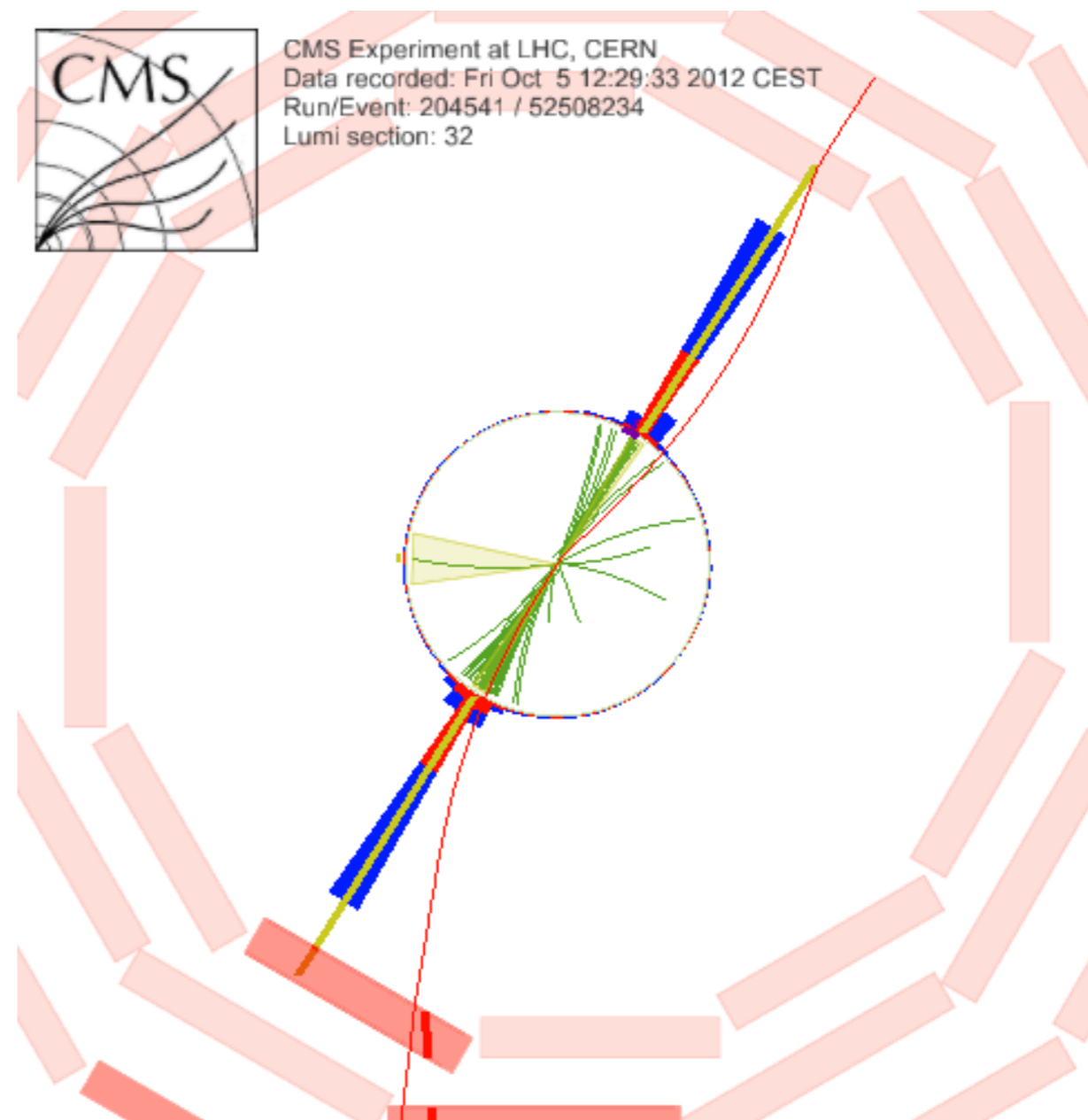
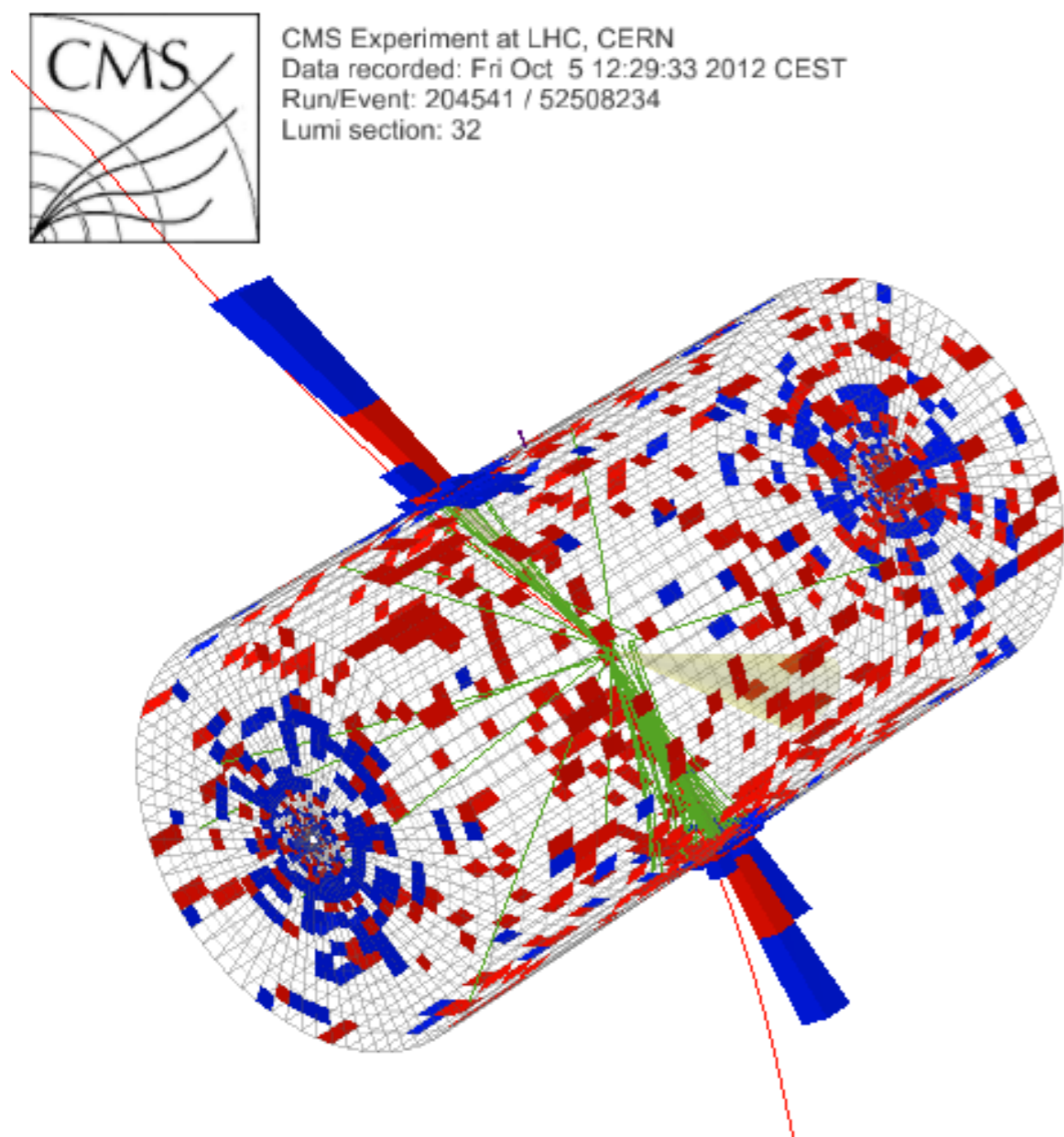


Bryan Webber  
Cavendish Laboratory  
University of Cambridge

# Introduction to Monte Carlo

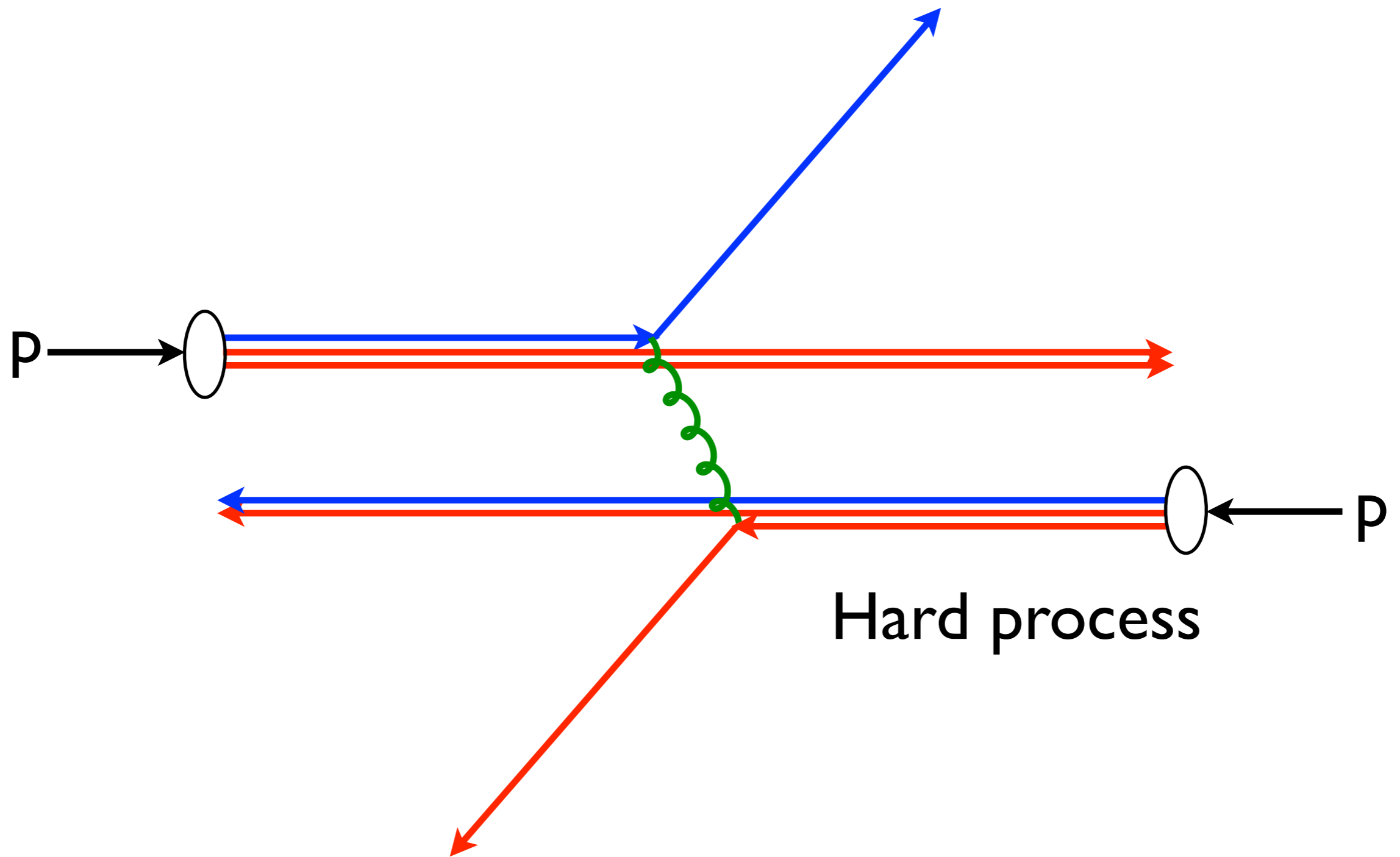
- Lecture 1: The Monte Carlo method
  - ✦ theoretical foundations and limitations
  - ✦ parton-level event generation
- Lecture 2: Hadron-level event generation
  - ✦ parton showering
  - ✦ hadronization and underlying event
  - ✦ sample of results

# A high-mass dijet event

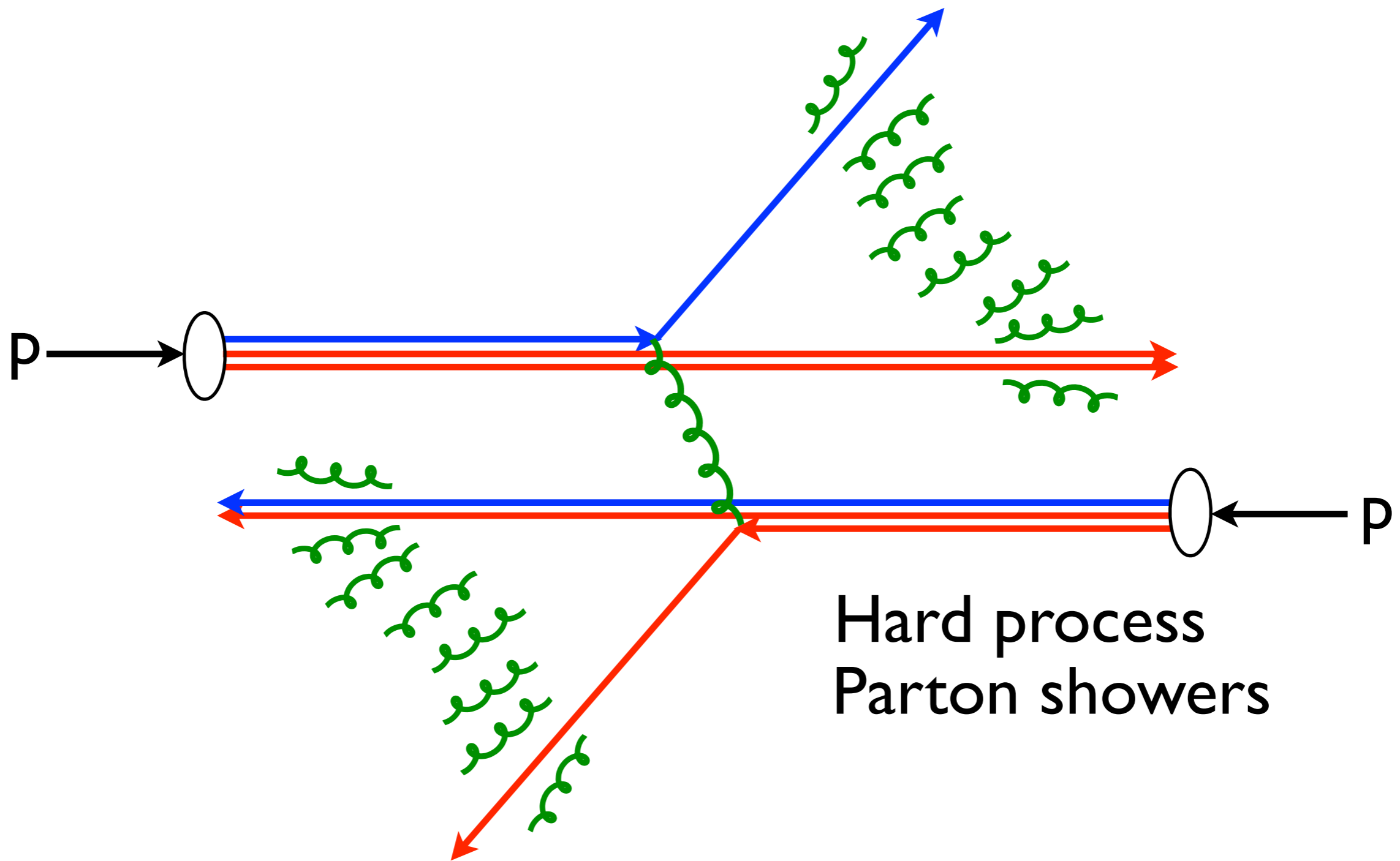


CMS PAS  
EXO-12-059

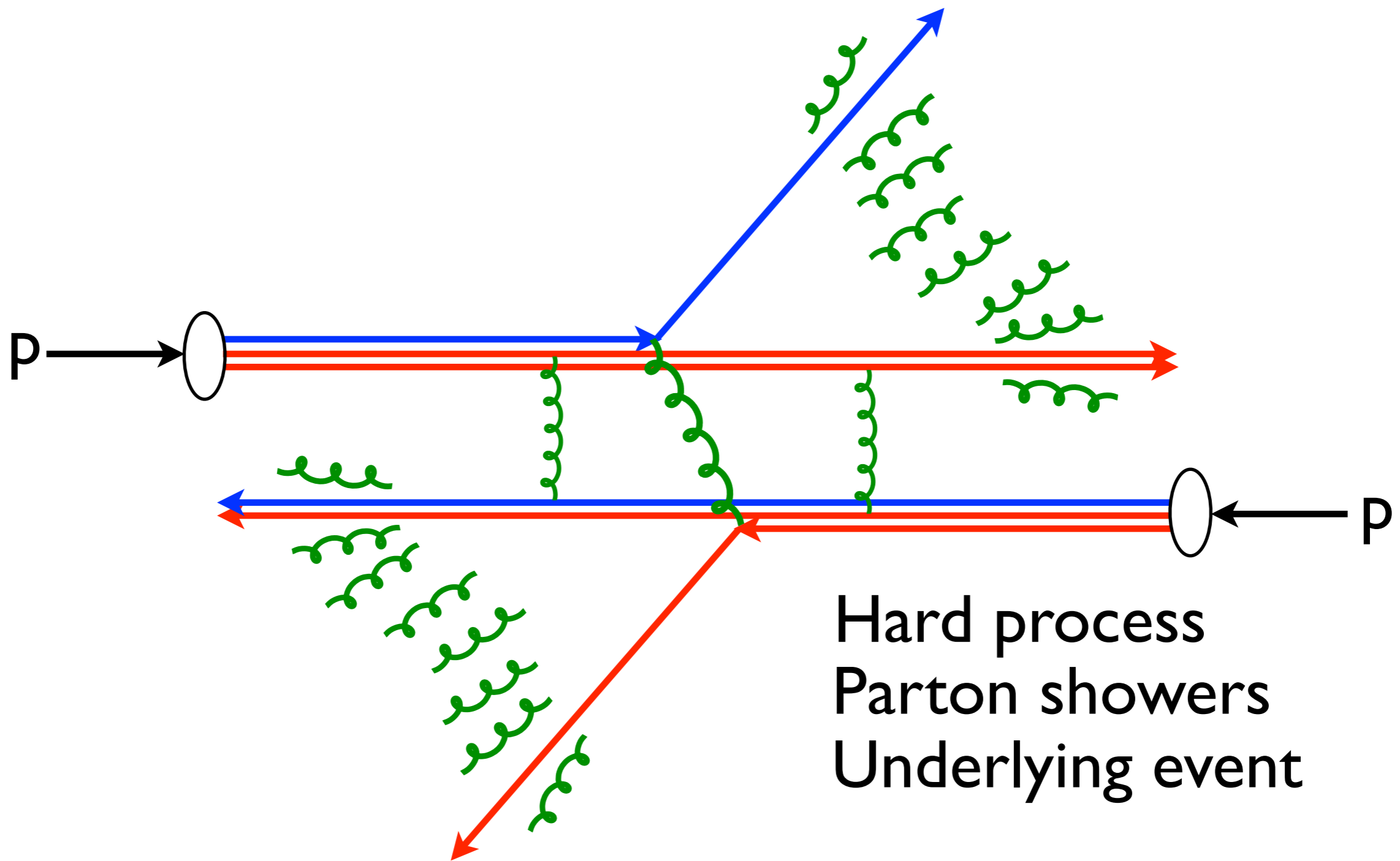
# LHC dijet



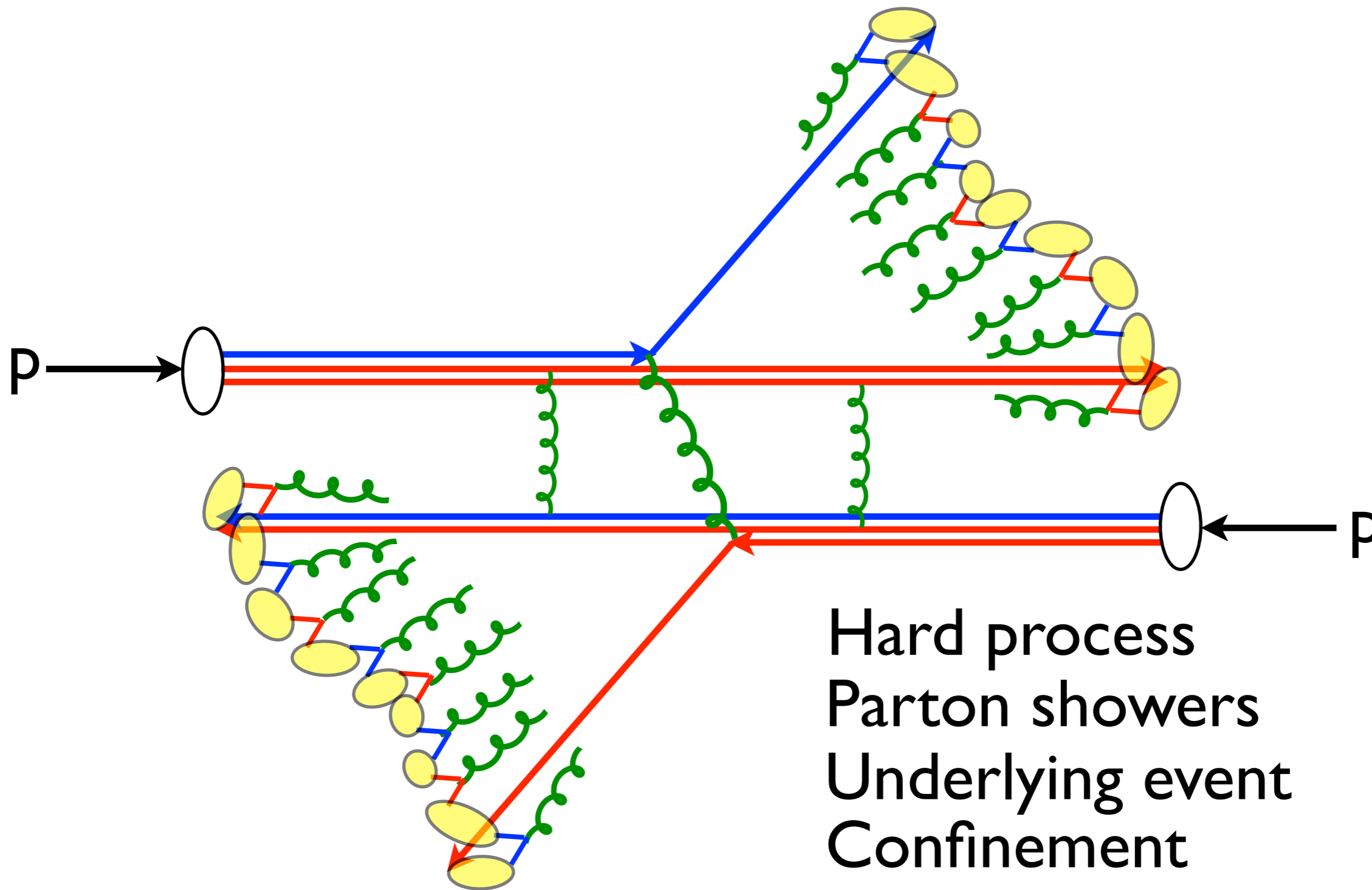
# LHC dijet



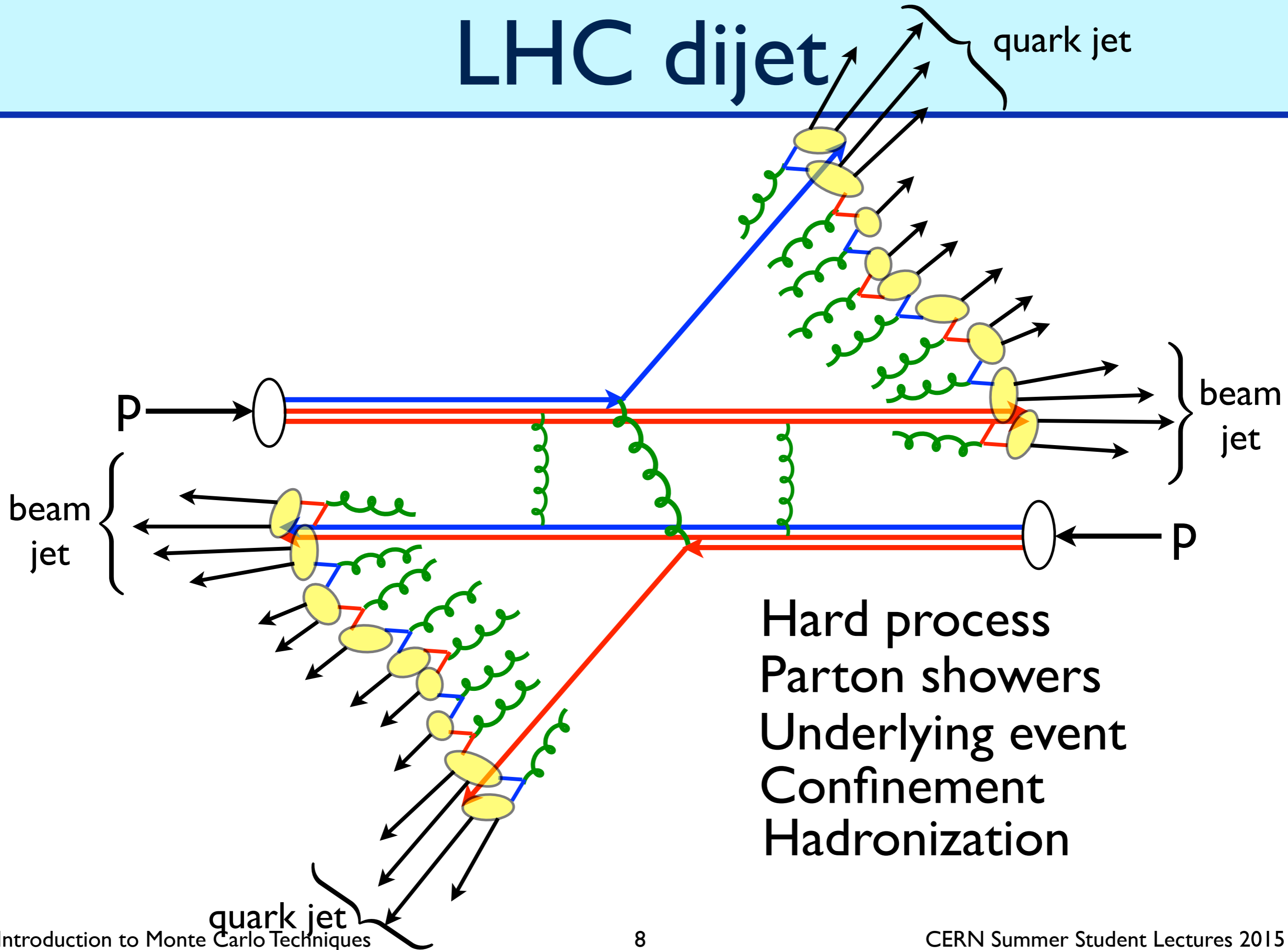
# LHC dijet



# LHC dijet

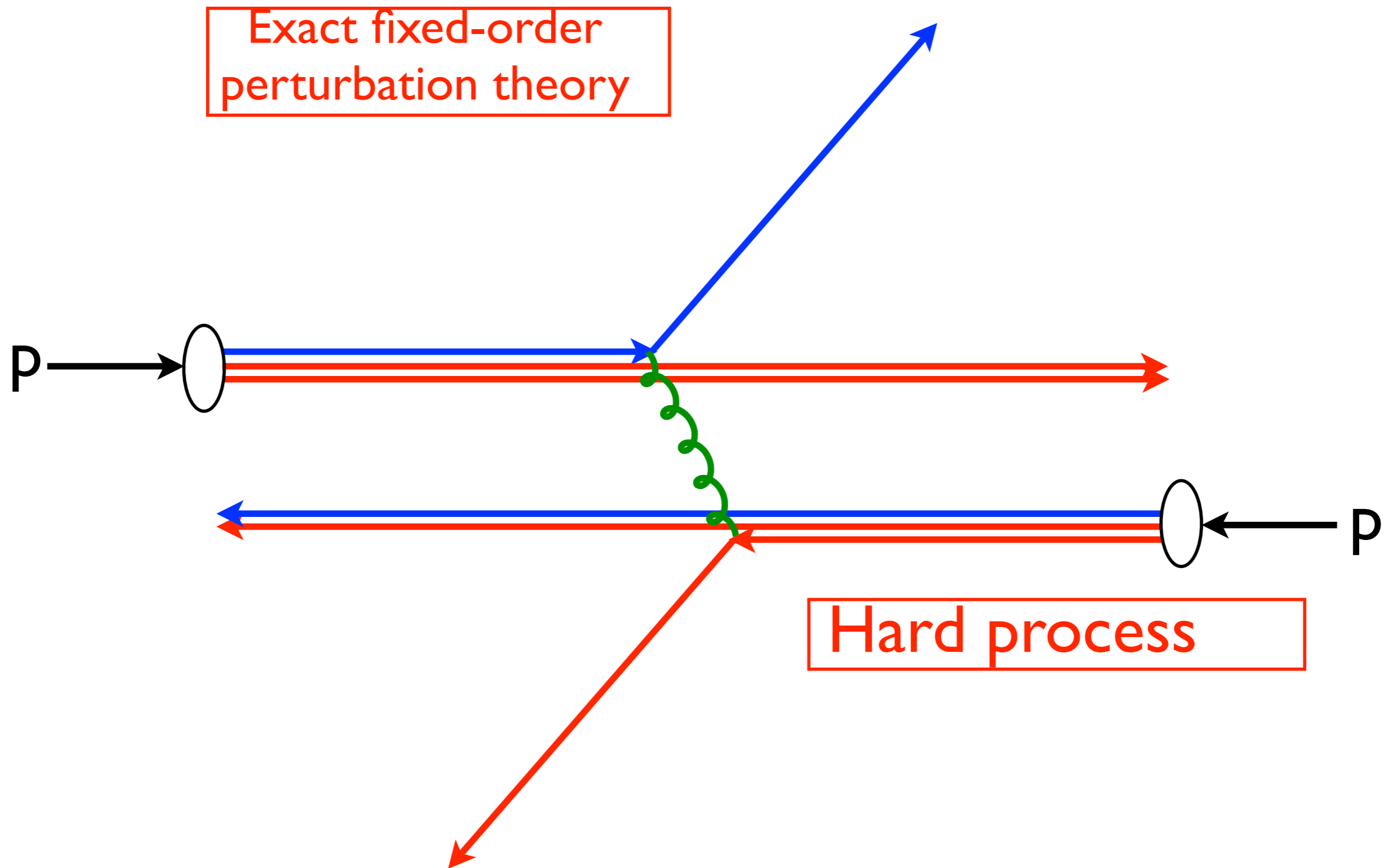


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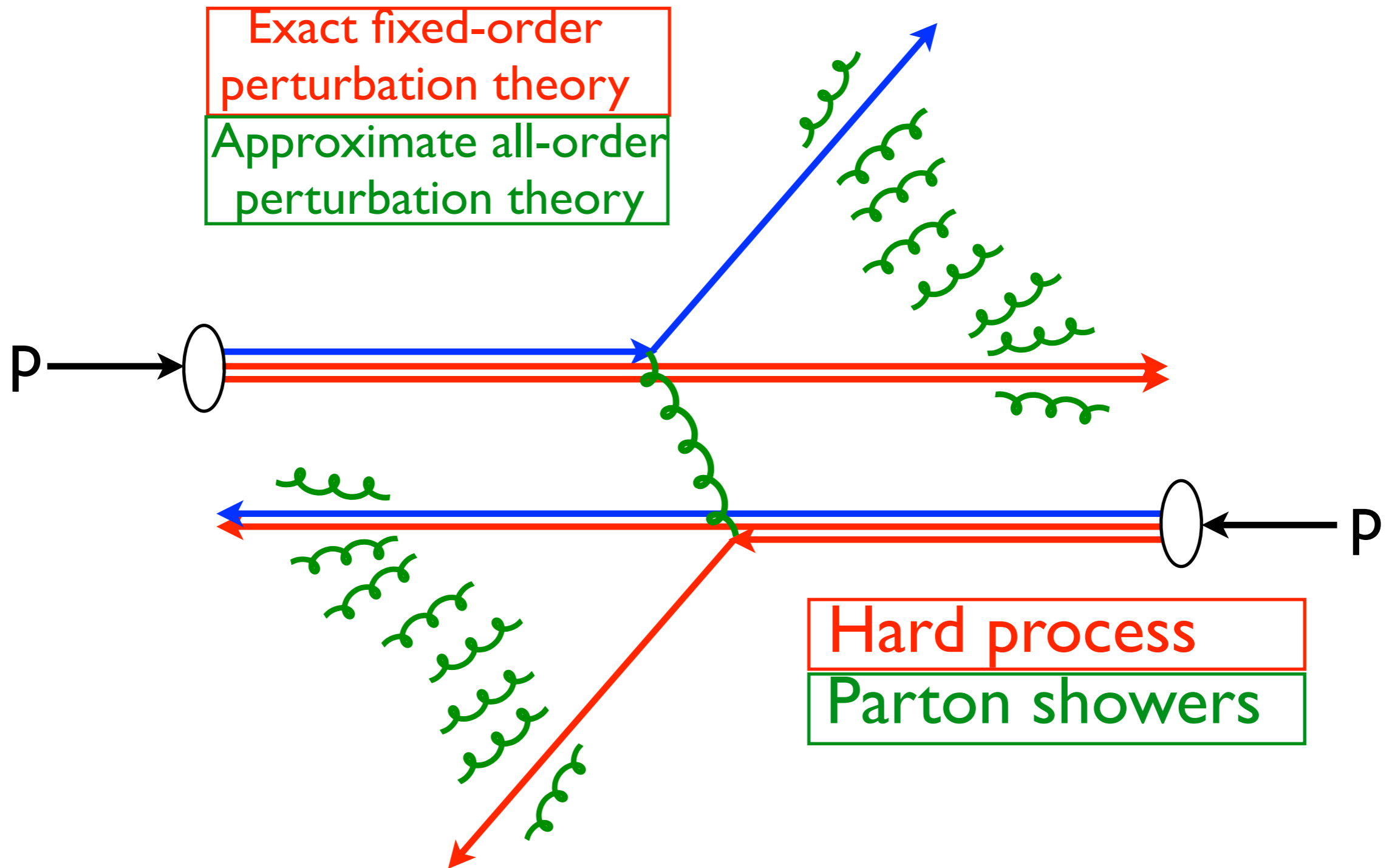




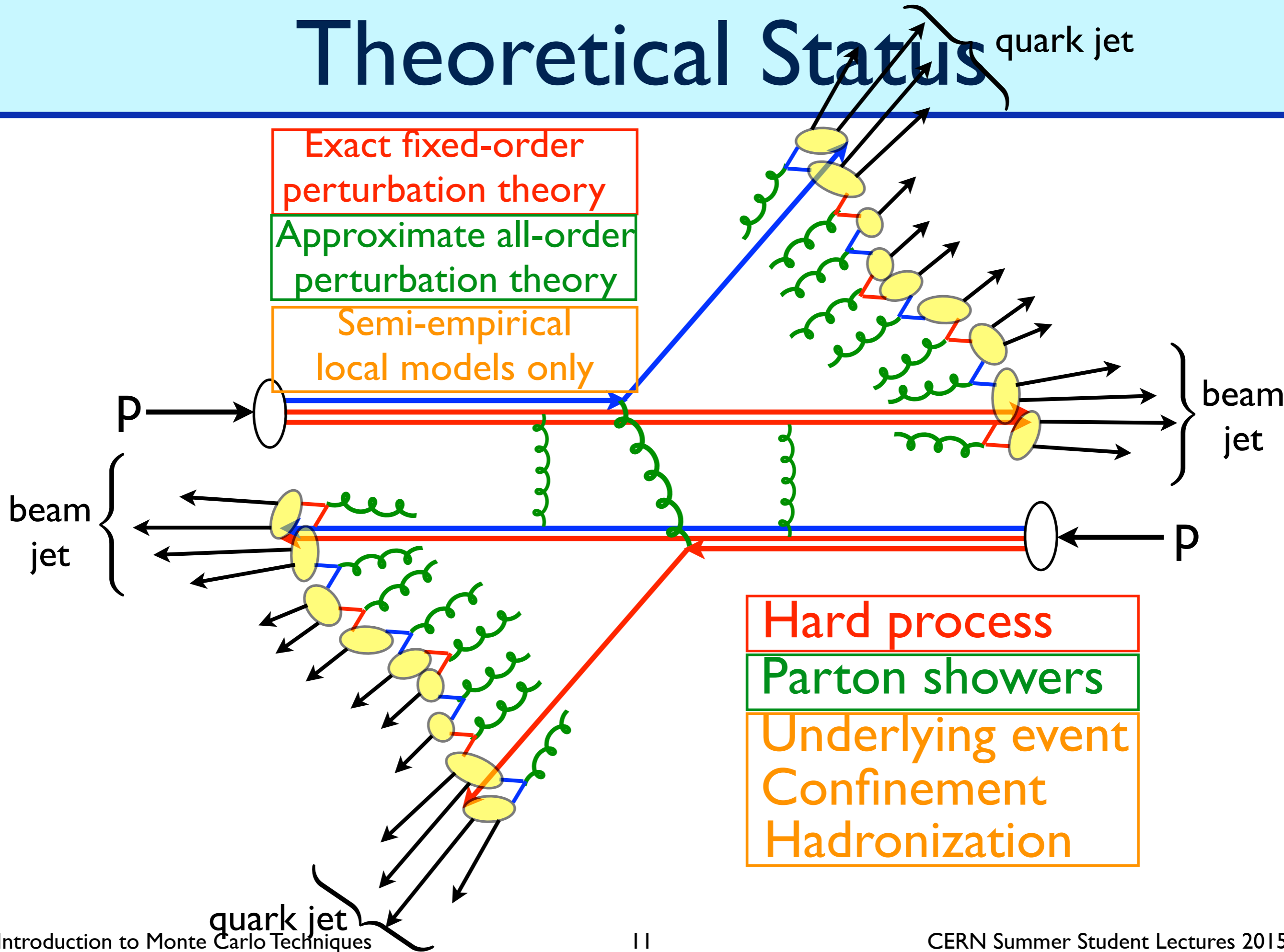
# Theoretical Status



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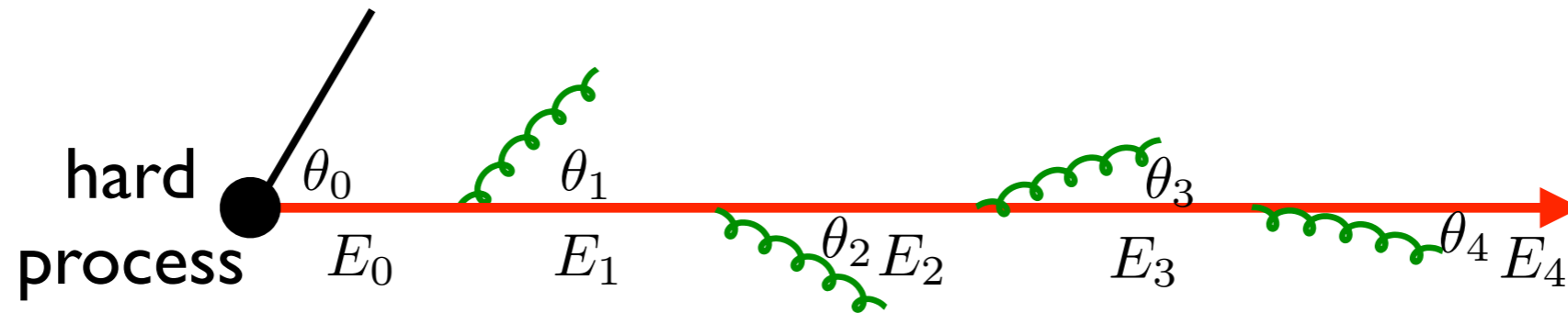
# QCD Factorization

$$\sigma_{pp \rightarrow X}(E_{pp}^2) = \int_0^1 dx_1 dx_2 \underbrace{f_i(x_1, \mu^2) f_j(x_2, \mu^2)}_{\substack{\text{parton} \\ \text{distributions} \\ \text{at scale } \mu^2}} \underbrace{\hat{\sigma}_{ij \rightarrow X}(x_1 x_2 E_{pp}^2, \mu^2)}_{\substack{\text{hard process} \\ \text{cross section}}}$$

momentum fractions

- Jet formation and underlying event take place over a much longer time scale, with unit probability
- Hence they cannot affect the cross section
- Scale dependences of parton distributions and hard process cross section are perturbatively calculable, and cancel order by order

# Parton Shower

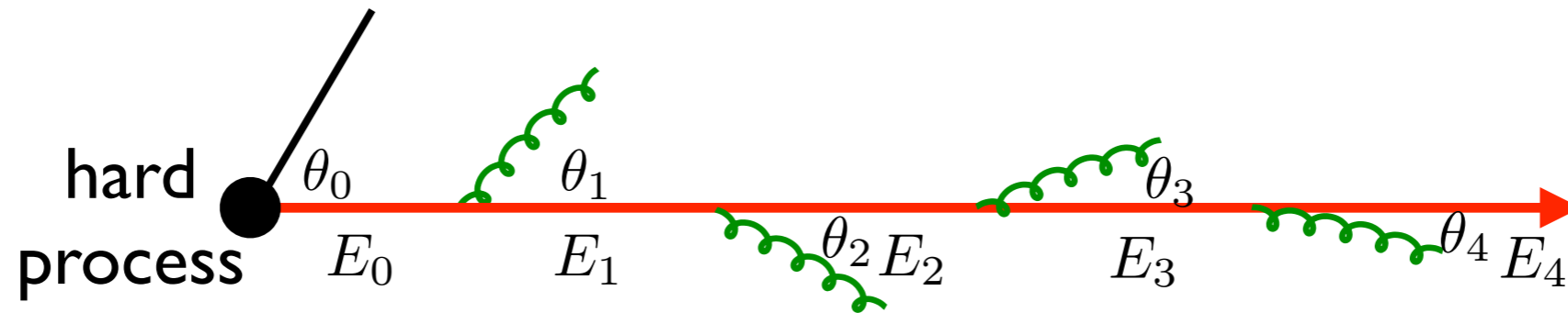


- Shower = sequence of emissions with decreasing angles and energies
- Approximation: keep only contributions  $\propto 1/\theta$

$$d^2\mathcal{P} = \frac{\alpha_s}{\pi} \frac{d\theta}{\theta} P(z) dz \quad z = \frac{E_{i+1}}{E_i}$$

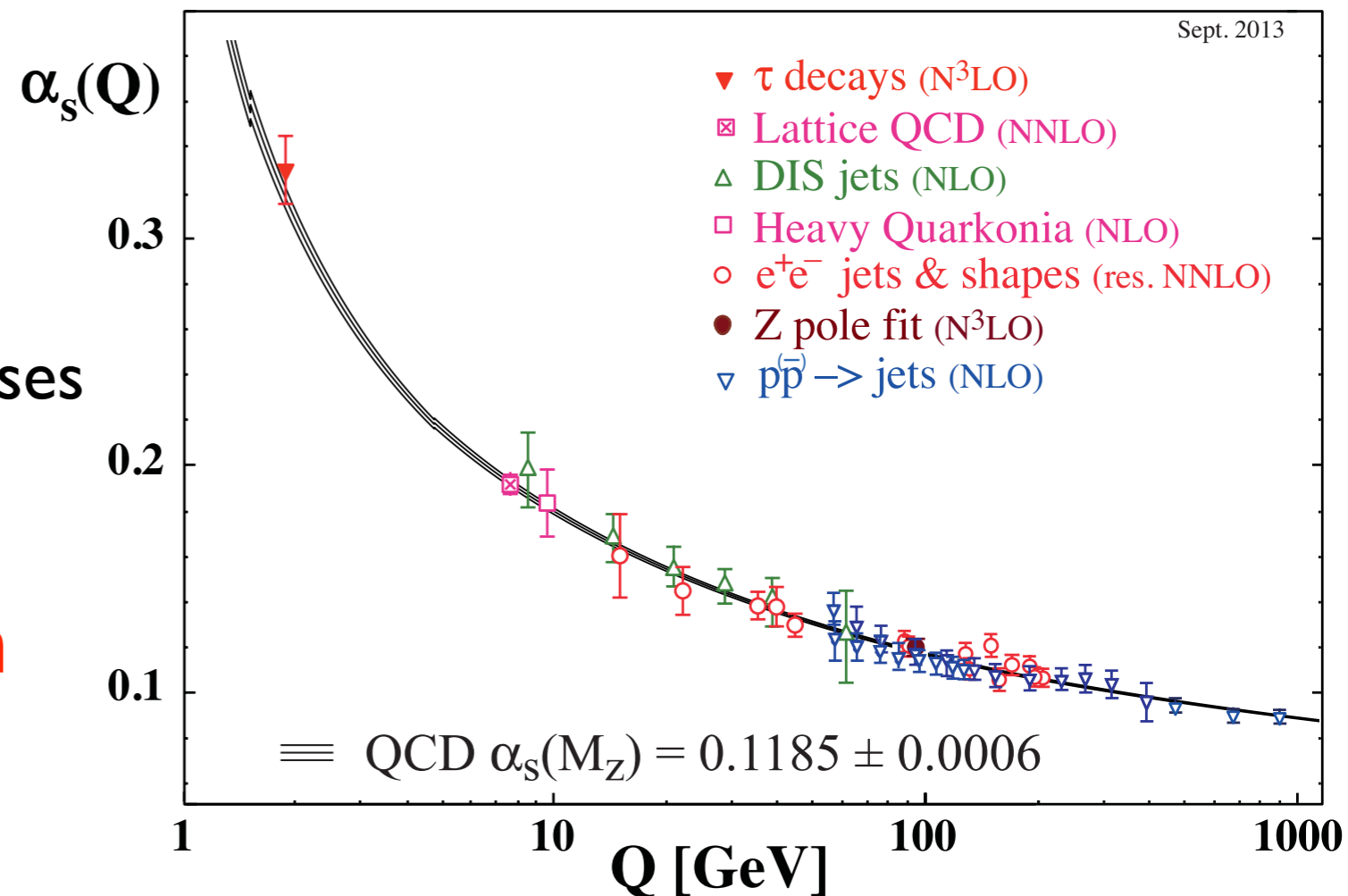
- For very small energy and/or angle, emission is “unresolvable”

# Parton Shower

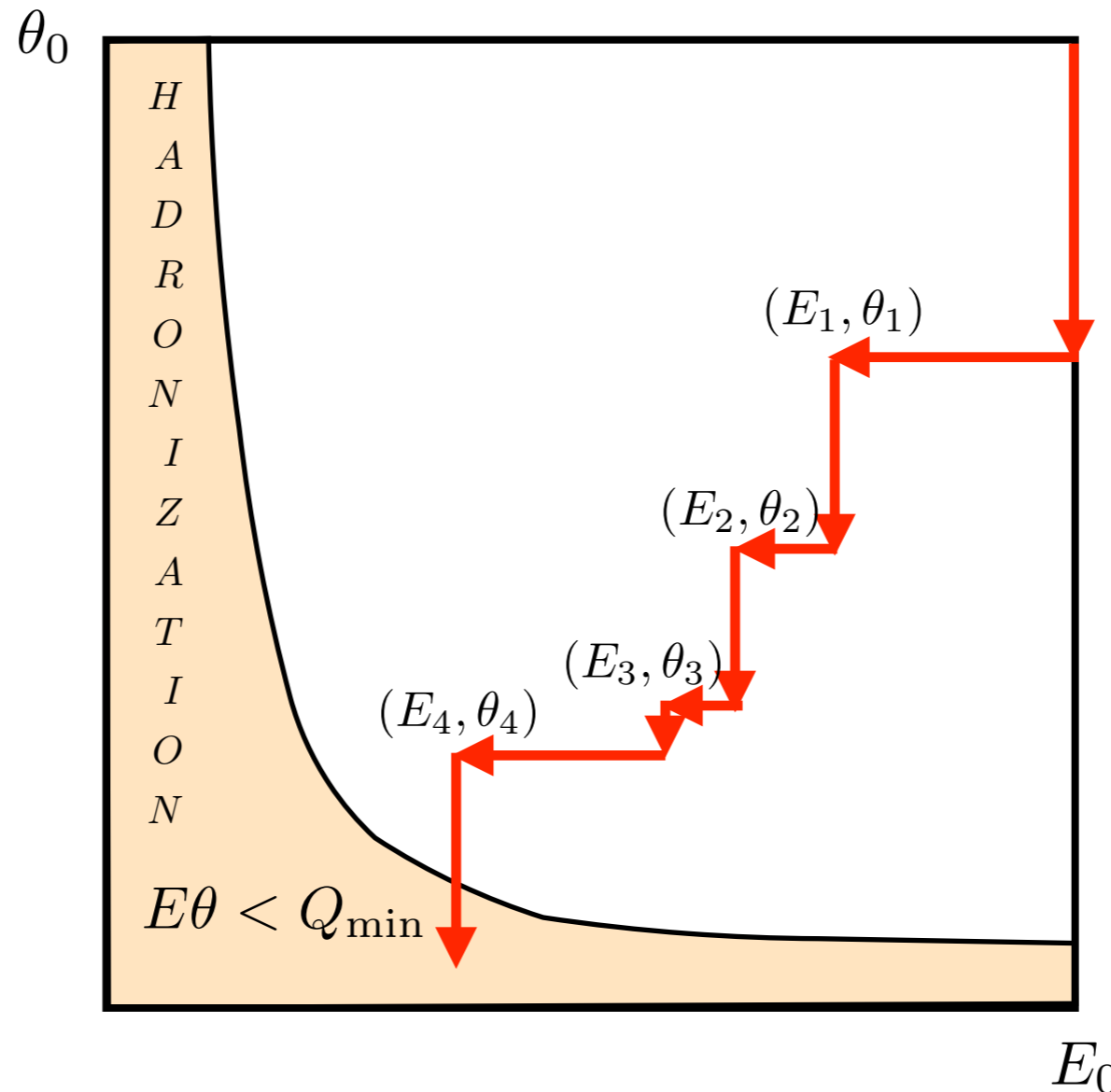
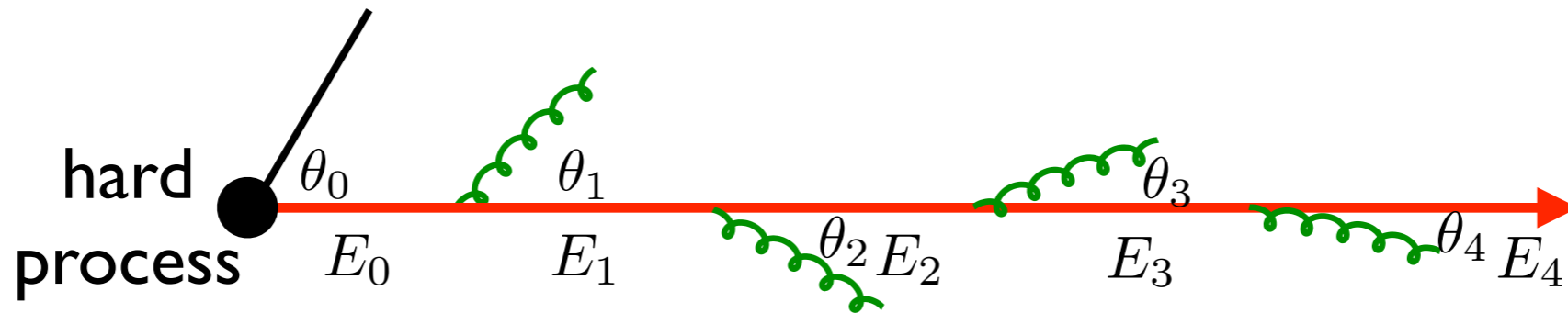


$$d^2\mathcal{P} = \frac{\alpha_s}{\pi} \frac{d\theta}{\theta} P(z) dz$$

- $\alpha_s$  increases as  $Q = E\theta$  decreases
- When  $Q < Q_{\min} \sim 1$  GeV  
 $\alpha_s \sim 1 \rightarrow$  **hadronization**

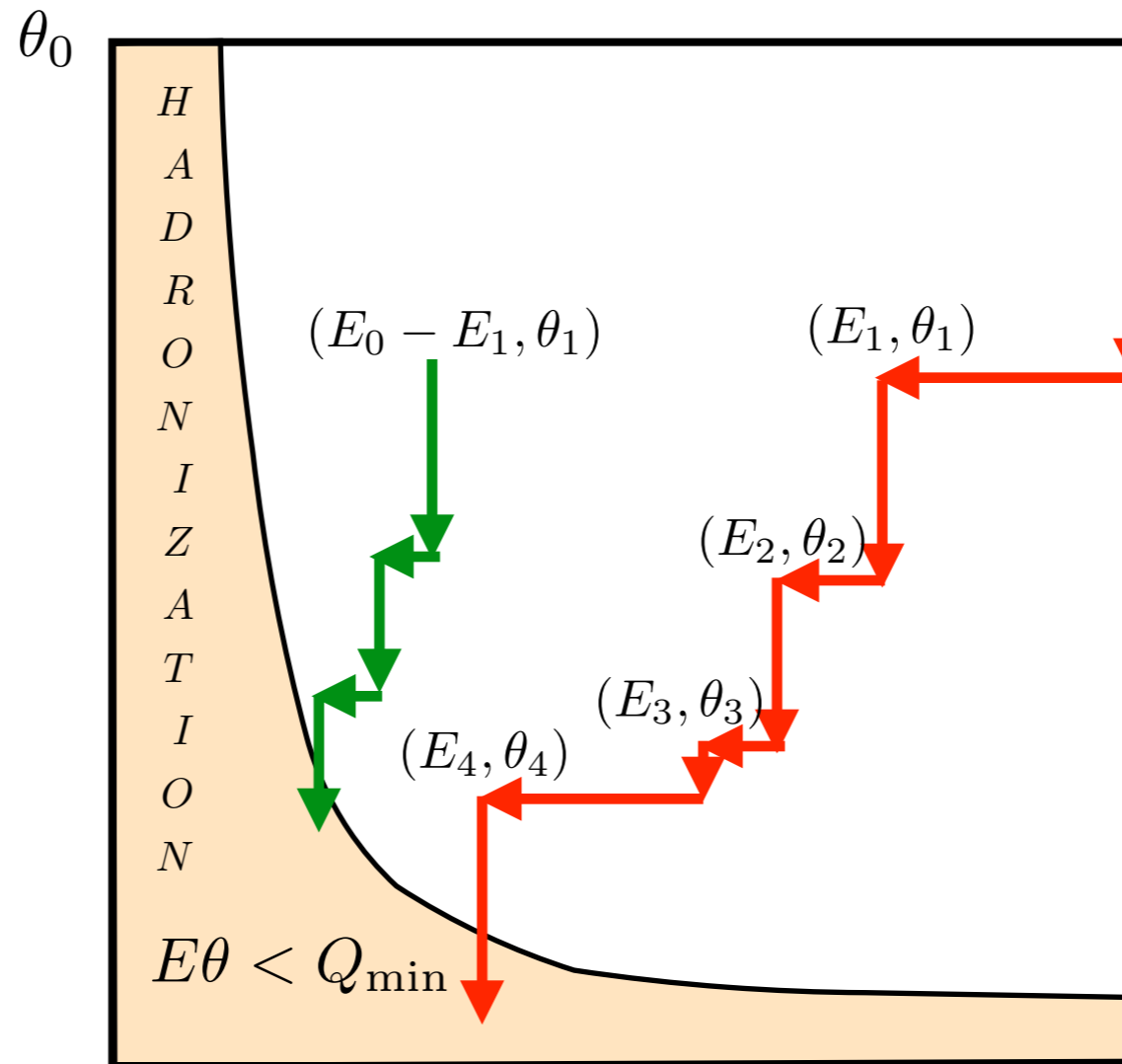
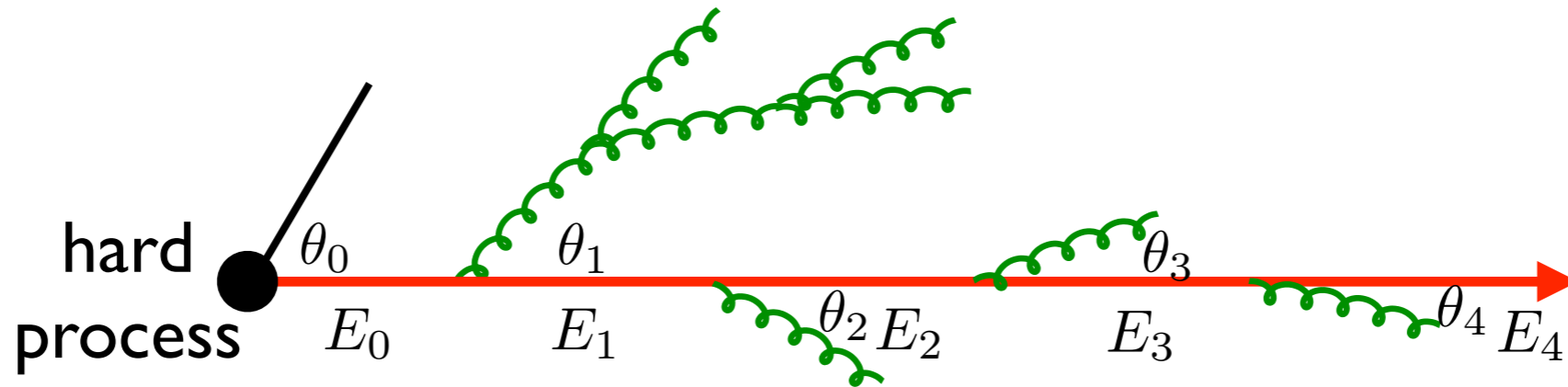


# Parton Shower Evolution



$$P_{q \rightarrow q}(z) = \frac{4}{3} \frac{1+z^2}{1-z}$$

# Parton Shower Evolution



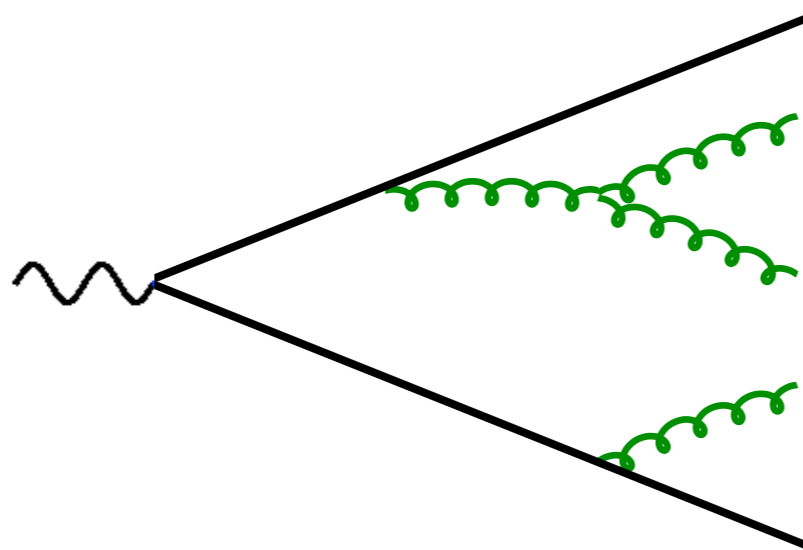
$$P_{q \rightarrow q}(z) = \frac{4}{3} \frac{1+z^2}{1-z}$$

$$P_{g \rightarrow g}(z) = \frac{1+z^4 + (1-z)^4}{3z(1-z)}$$



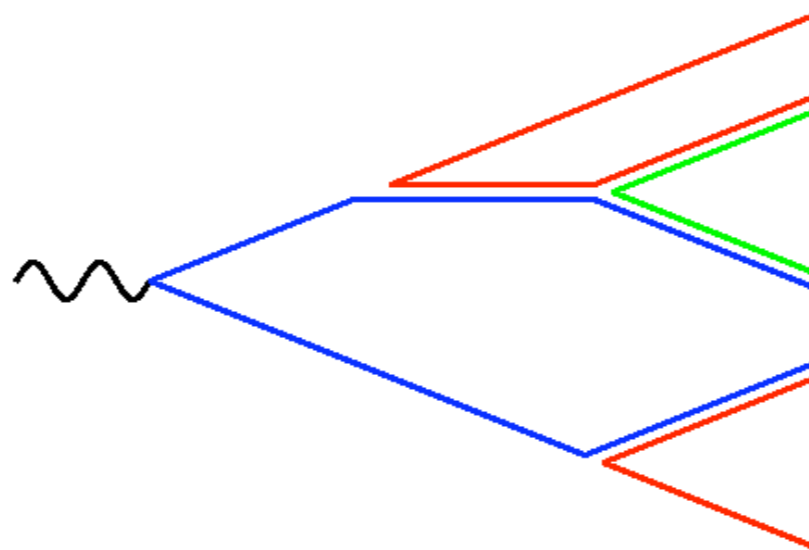
# Hadronization Models

- In parton shower, relative transverse momenta evolve from a high scale  $Q$  towards lower values
- At a scale near  $\Lambda_{\text{QCD}} \sim 200$  MeV, perturbation theory breaks down and hadrons are formed
- Before that, at scales  $\sim \text{few} \times \Lambda_{\text{QCD}}$ , there is universal **preconfinement** of colour
- Colour, flavour and momentum flows are only **locally** redistributed by hadronization



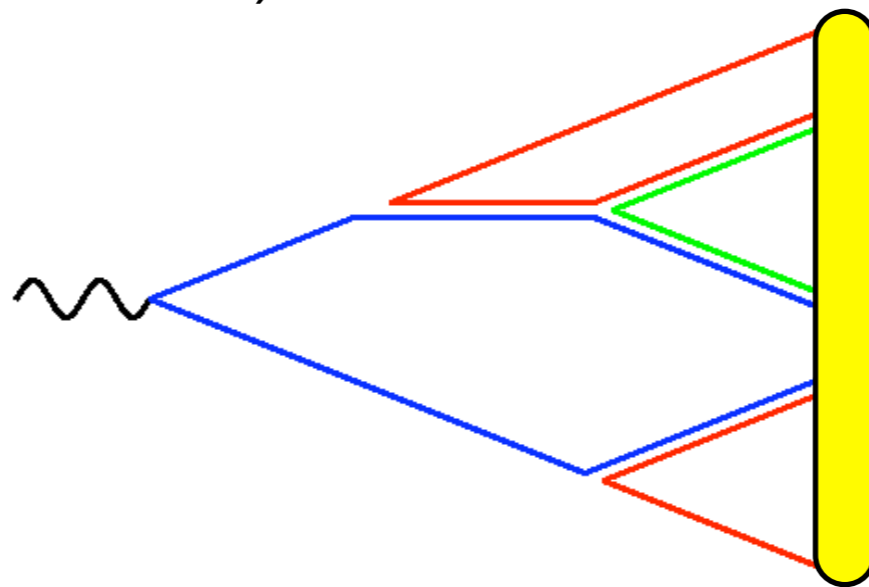
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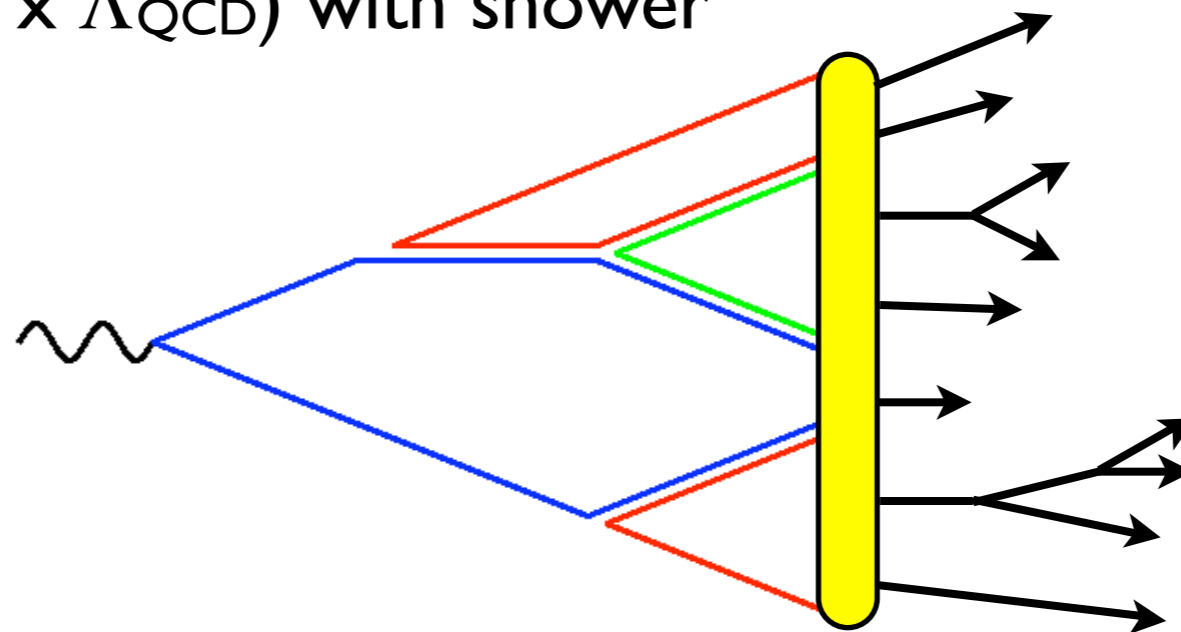
# String Hadronization Model

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- Colour flow dictates how to connect **hadronic string** (width  $\sim \text{few} \times \Lambda_{\text{QCD}}$ ) with shower



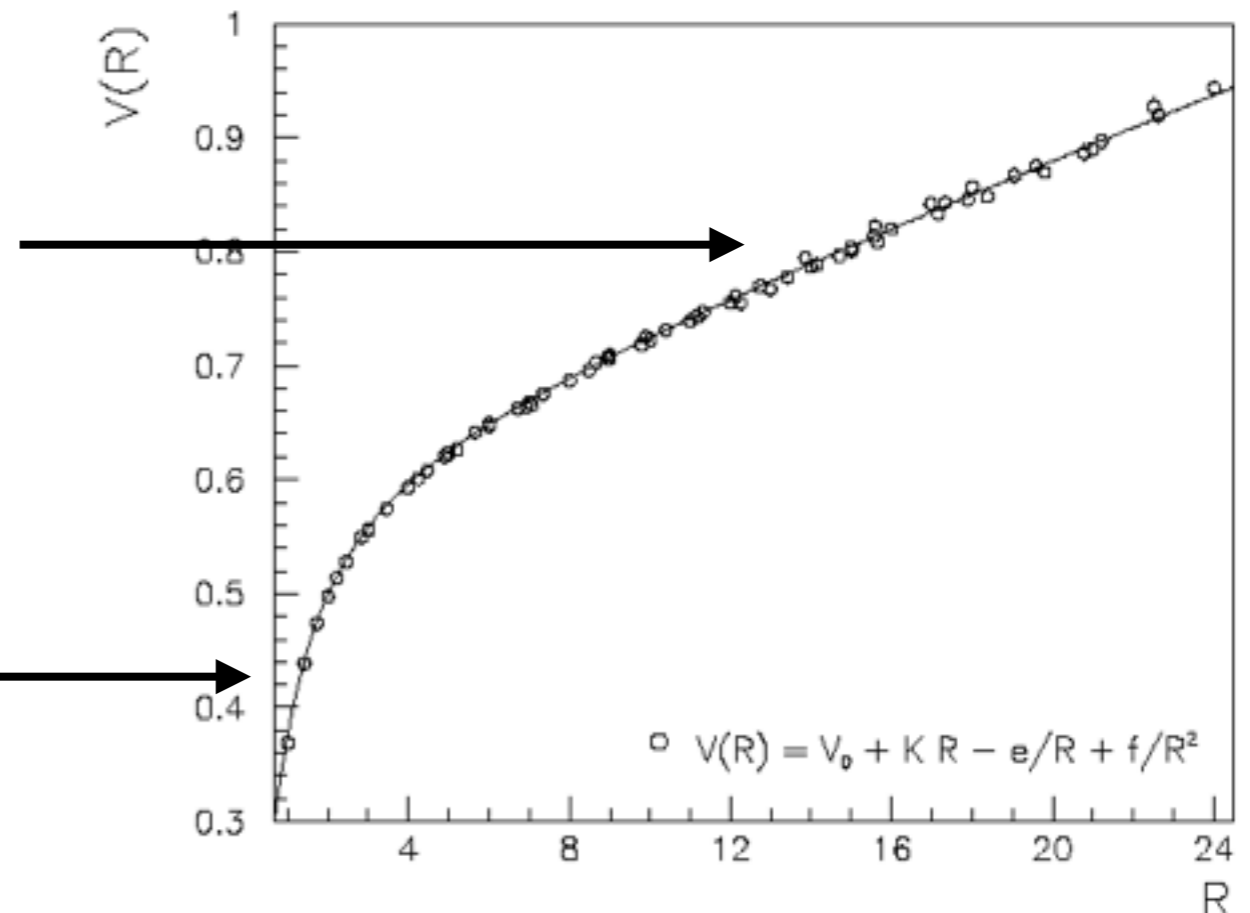
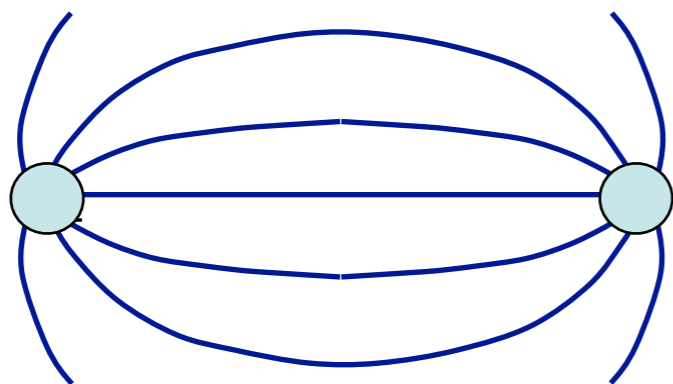
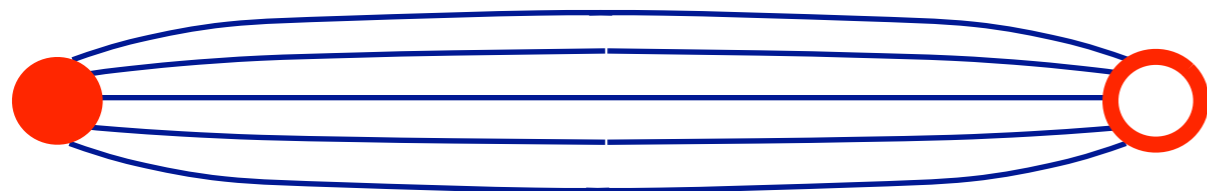
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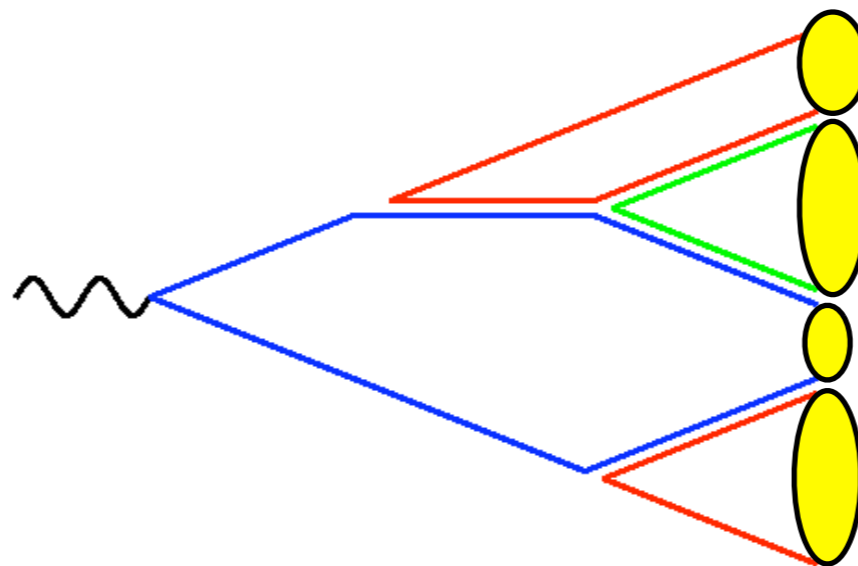
# String Hadronization Model

- At short distances (large  $Q$ ), QCD is like QED: colour field lines spread out ( $1/r$  potential)
- At long distances, gluon self-attraction gives rise to colour string (linear potential, quark confinement)
- Intense colour field induces quark-antiquark pair creation: hadronization



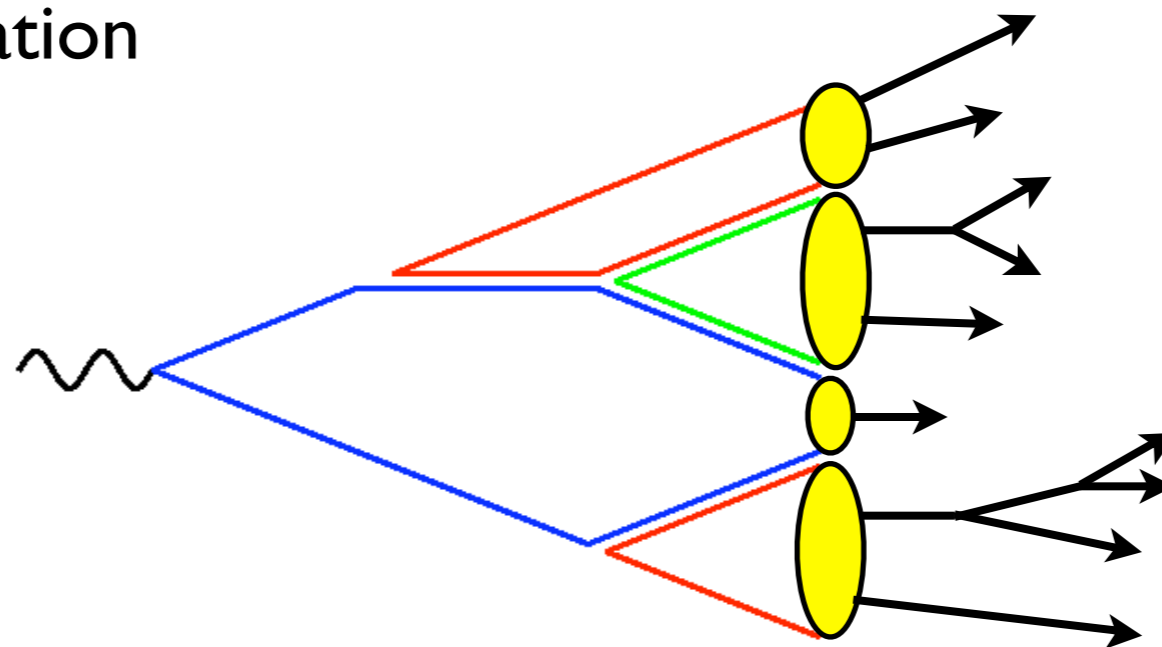
# Cluster Hadronization Model

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- Decay of **preconfined clusters** provides a direct basis for hadronization

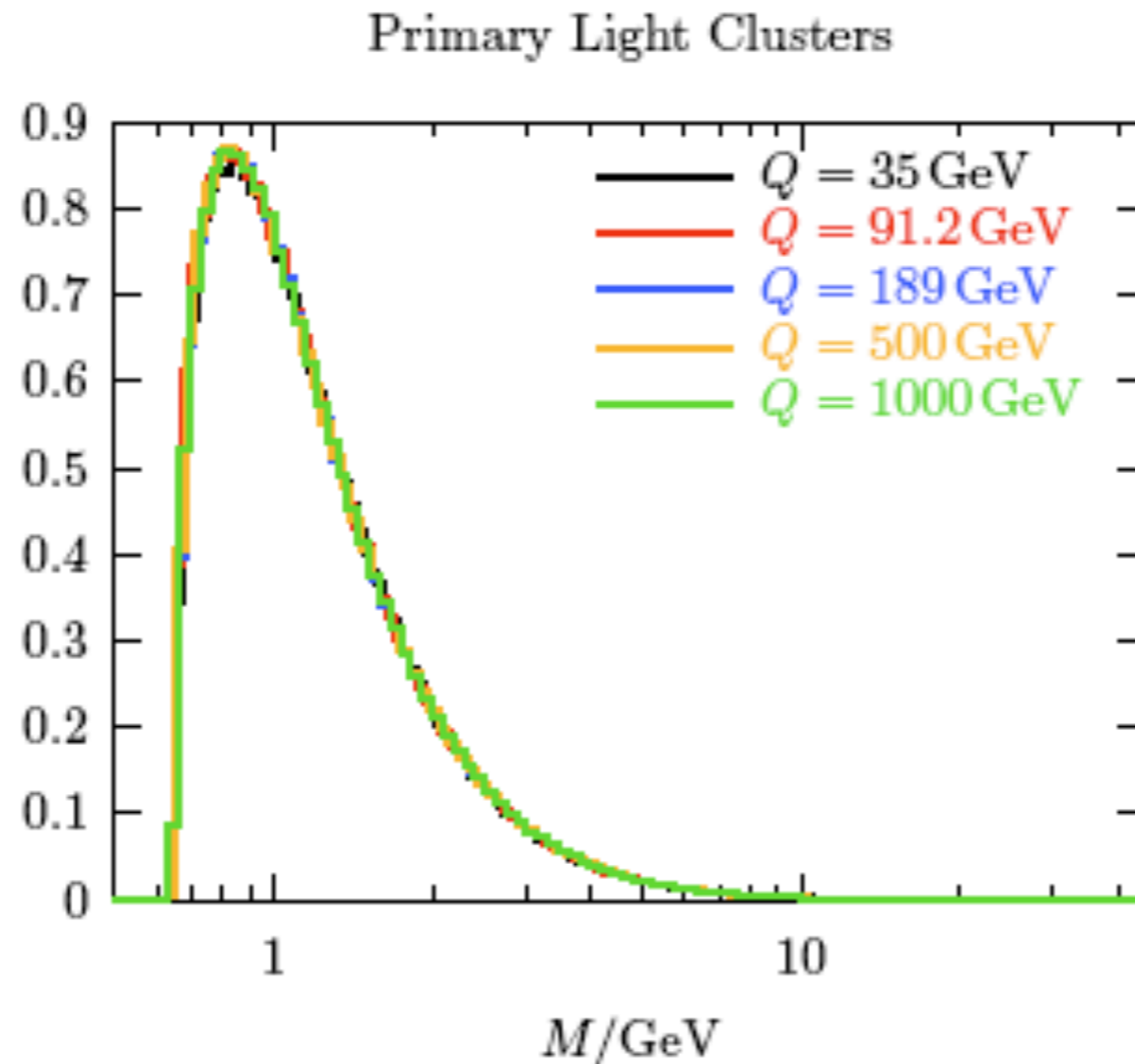


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# Cluster Hadronization Model



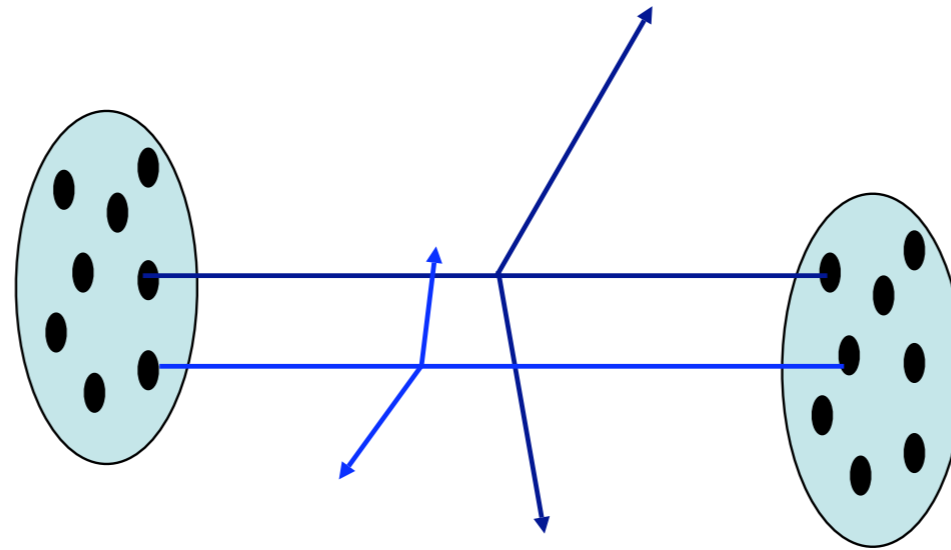
- Mass distribution of preconfined clusters is universal
- Phase-space decay model for most clusters
- High-mass tail decays anisotropically (string-like)



# Hadronization Status

- No fundamental progress since 1980s
  - ✦ Available non-perturbative methods (lattice, AdS/QCD, ...) are not applicable
- Less important in some respects in LHC era
  - ✦ Jets, leptons and photons are observed objects, not hadrons
- But still important for detector effects
  - ✦ Jet response, heavy-flavour tagging, lepton and photon isolation, ...

# Underlying Event (MPI)



- **Multiple parton interactions** in same collision
  - ✦ Depends on density profile of proton
- Assume QCD 2-to-2 secondary collisions
  - ✦ Need cutoff at low  $p_T$
- Need to model colour flow
  - ✦ Colour reconnections are necessary

# Sample of Event Generator Results

# MC Event Generators

## ● HERWIG

<http://projects.hepforge.org/herwig/>

➔ Angular-ordered parton shower, cluster hadronization

➔ v6 Fortran; Herwig++

## ● PYTHIA

<http://www.thep.lu.se/~torbjorn/Pythia.html>

➔  $k_t$ -ordered parton shower, string hadronization

➔ v6 Fortran; v8 C++

## ● SHERPA

<http://projects.hepforge.org/sherpa/>

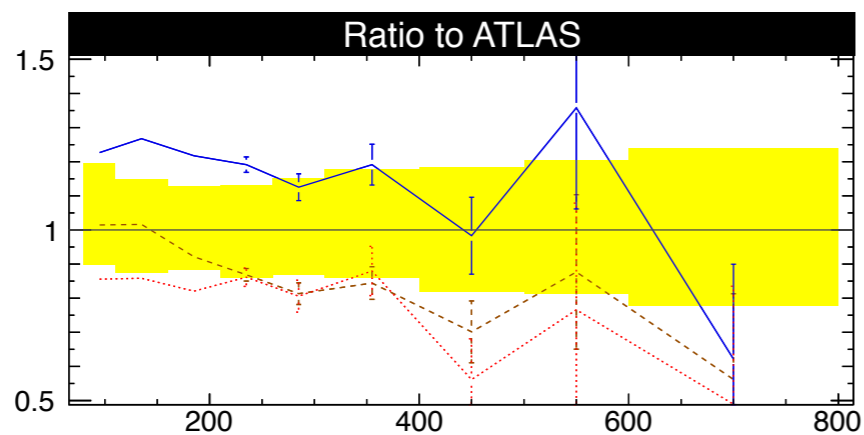
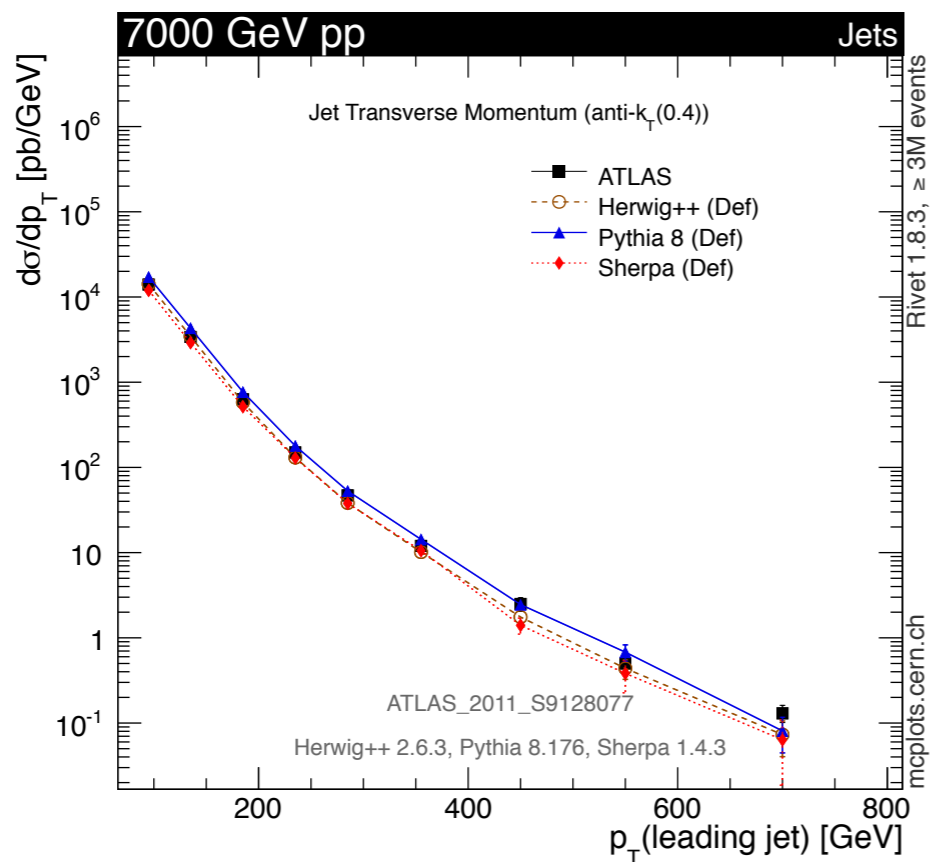
➔ Dipole-type parton shower, cluster hadronization

➔ C++

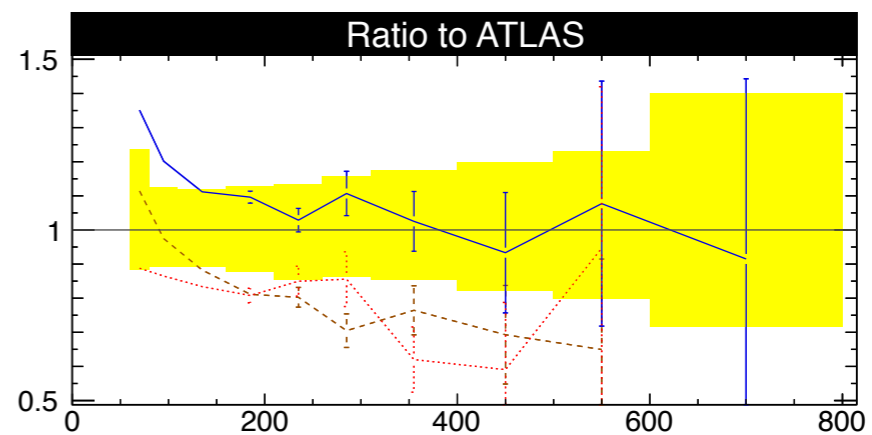
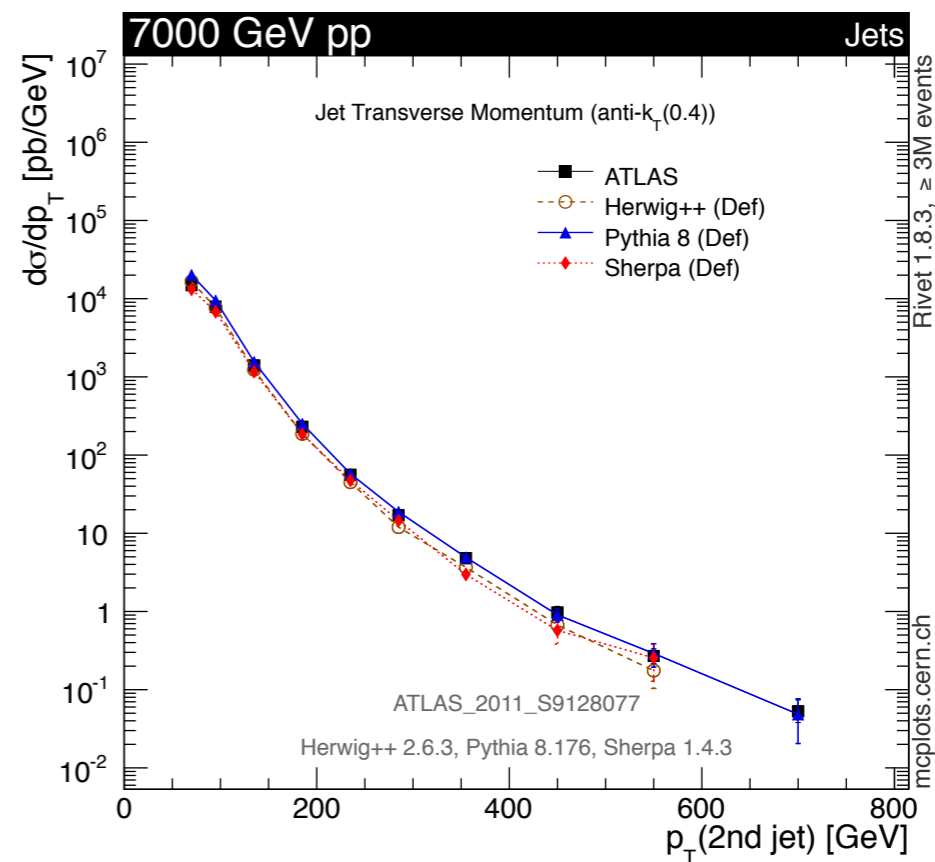
“General-purpose event generators for LHC physics”,  
A Buckley et al., arXiv:1101.2599, Phys. Rept. 504(2011)145

# Jets

# Jet $p_T$



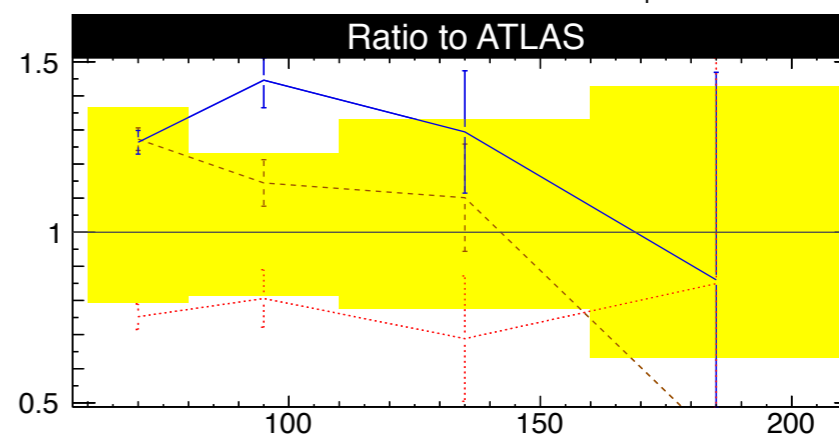
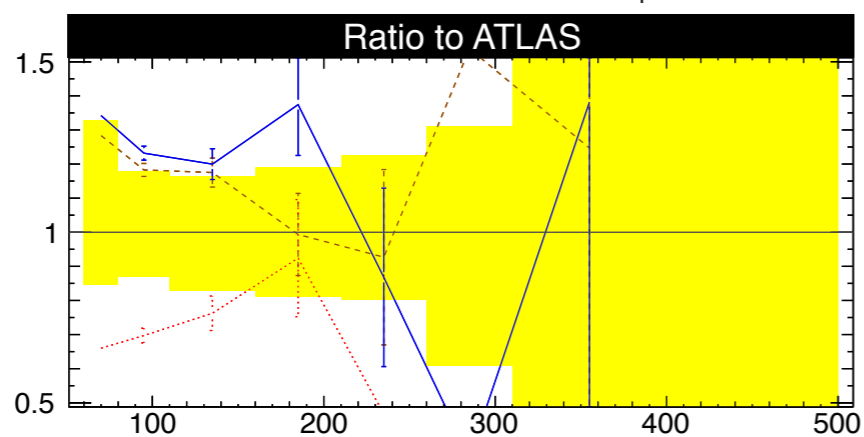
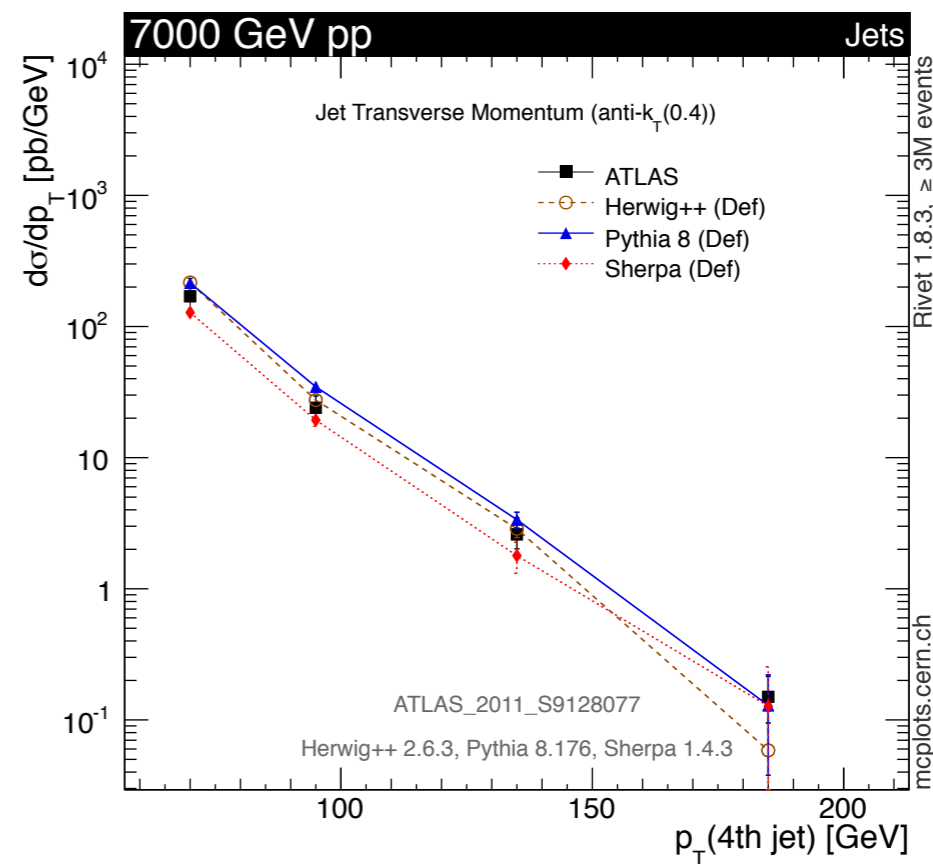
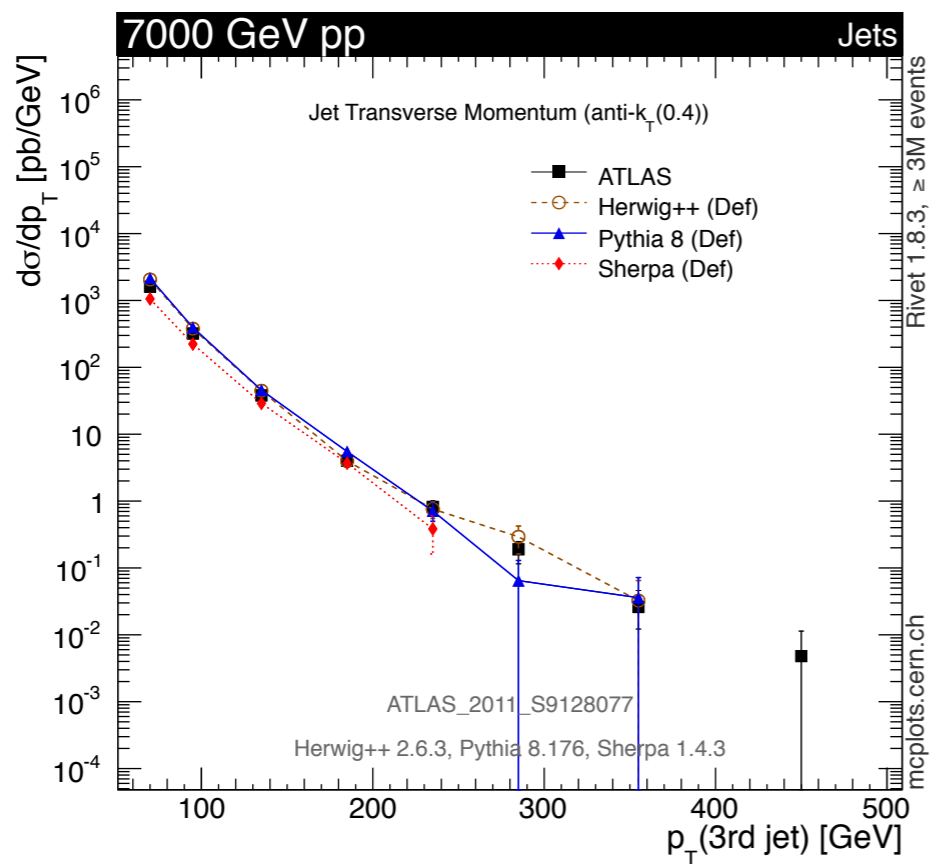
Leading jet



Second jet

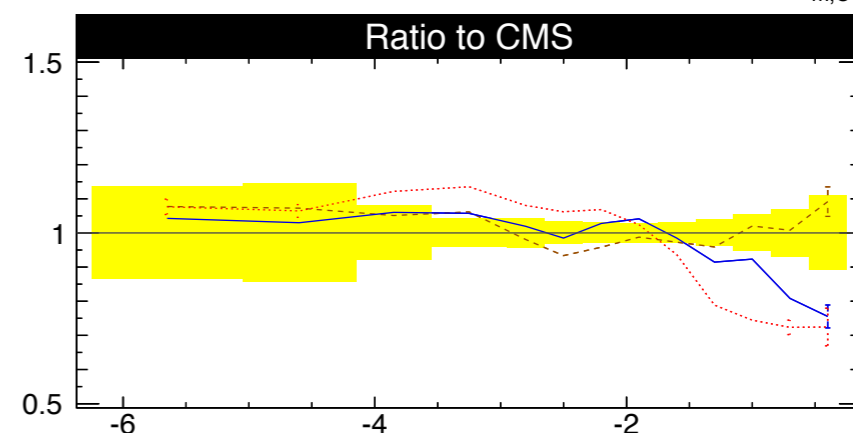
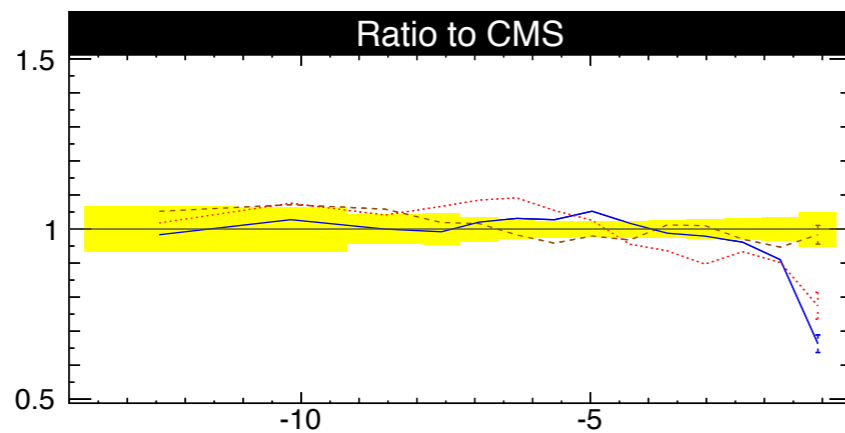
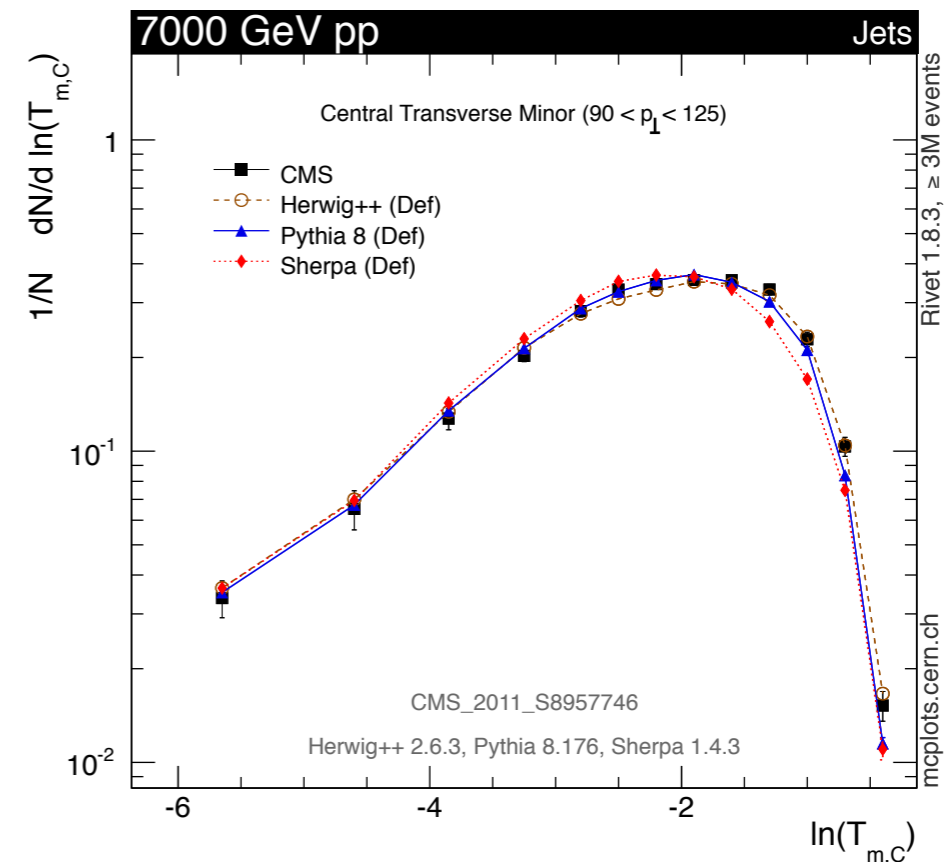
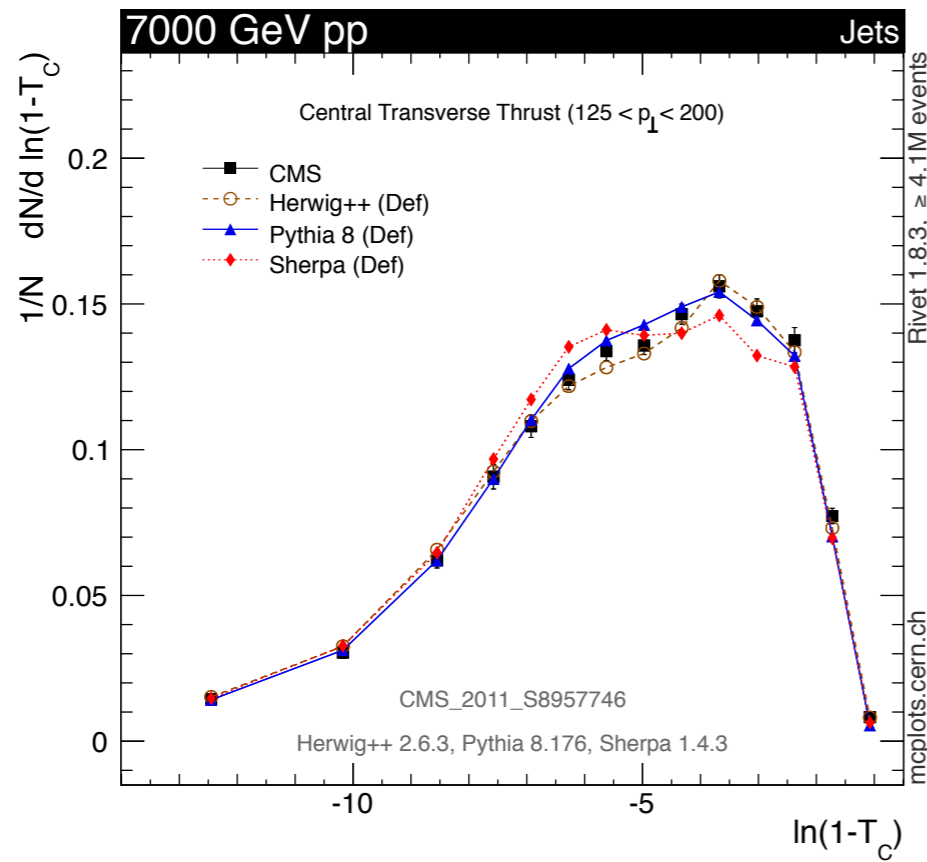
<http://mcplots.cern.ch>

# Jet $p_T$



Extra jets from parton showers

# Jet event shapes

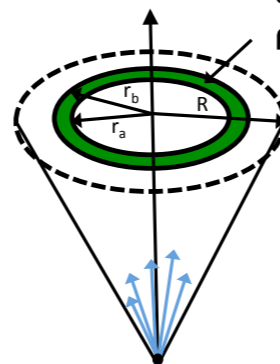
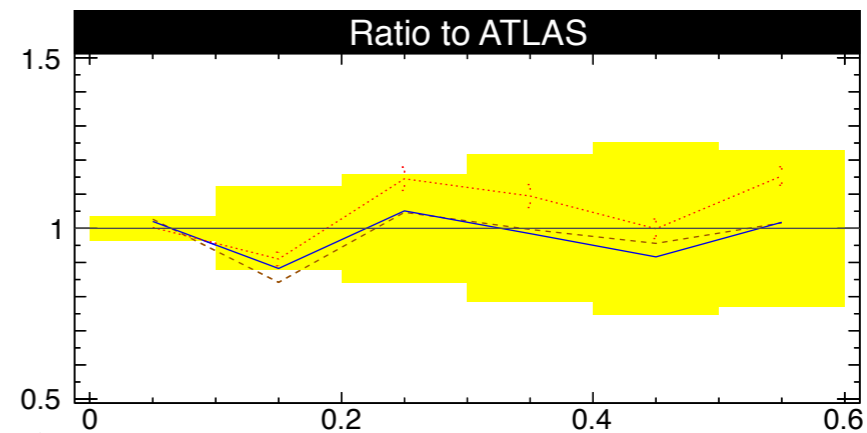
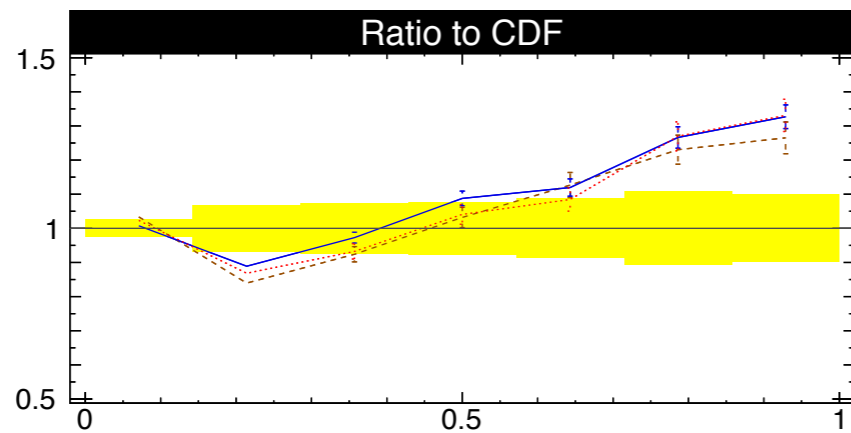
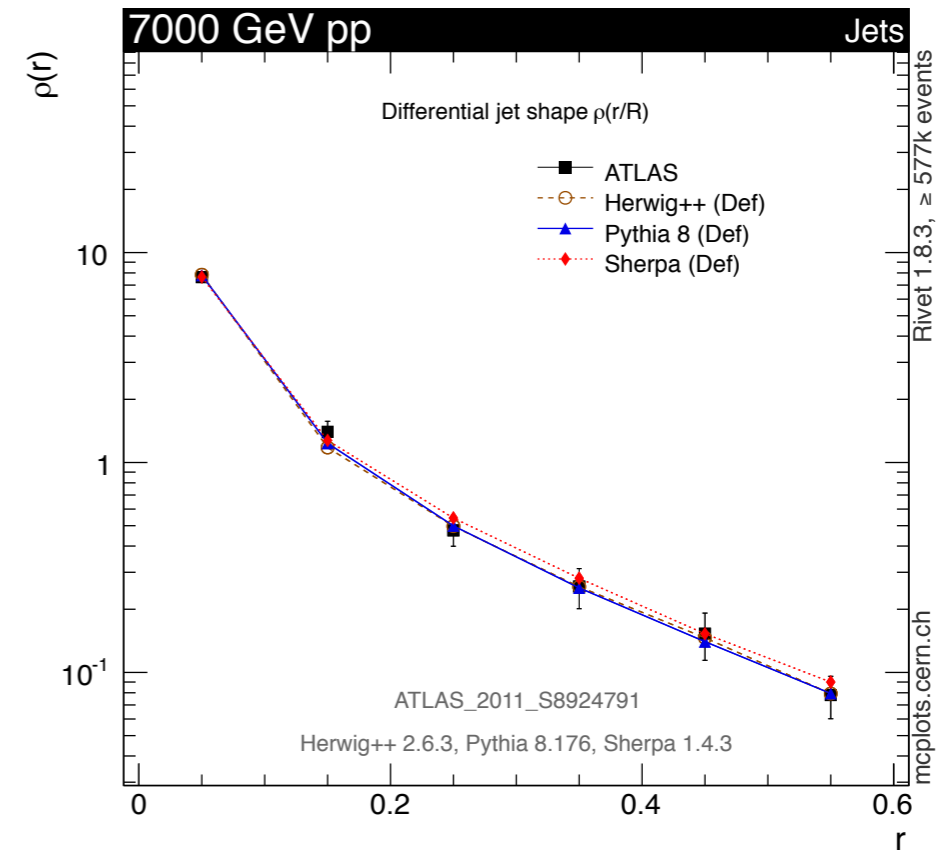
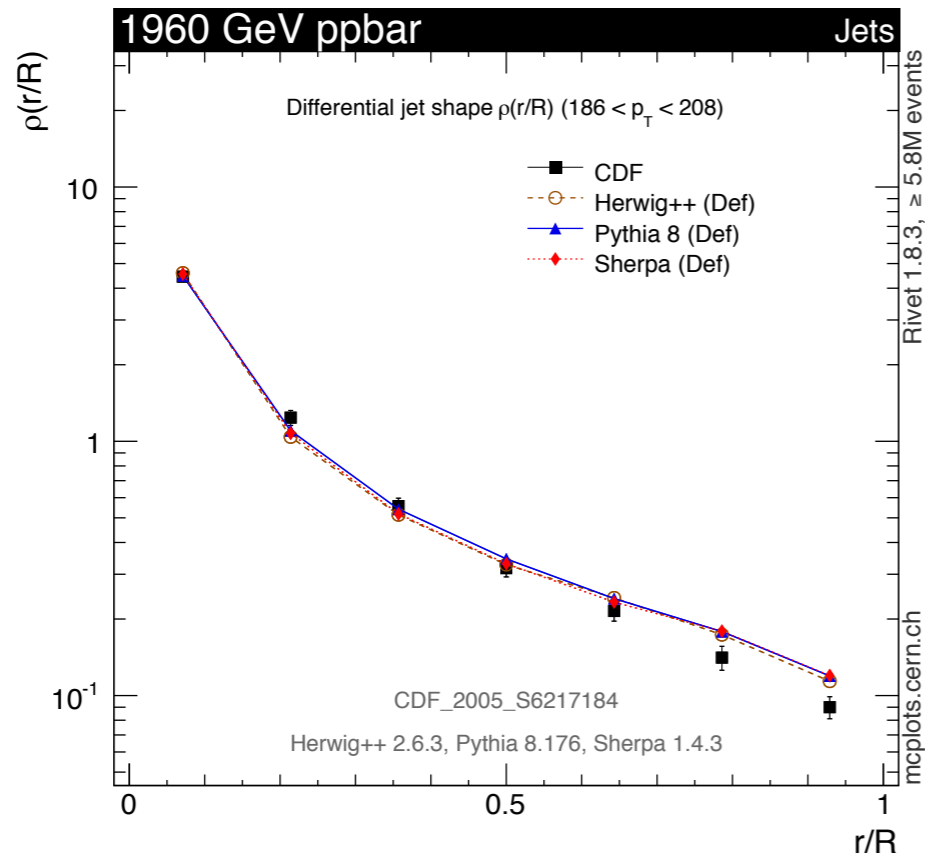


$$T_c \equiv \max_{\hat{n}_T} \frac{\sum_i |\vec{p}_{\perp,i} \cdot \hat{n}_T|}{\sum_i p_{\perp,i}}$$

$$T_{m,C} \equiv \frac{\sum_i |\vec{p}_{\perp,i} \times \hat{n}_{T,C}|}{\sum_i p_{\perp,i}}$$

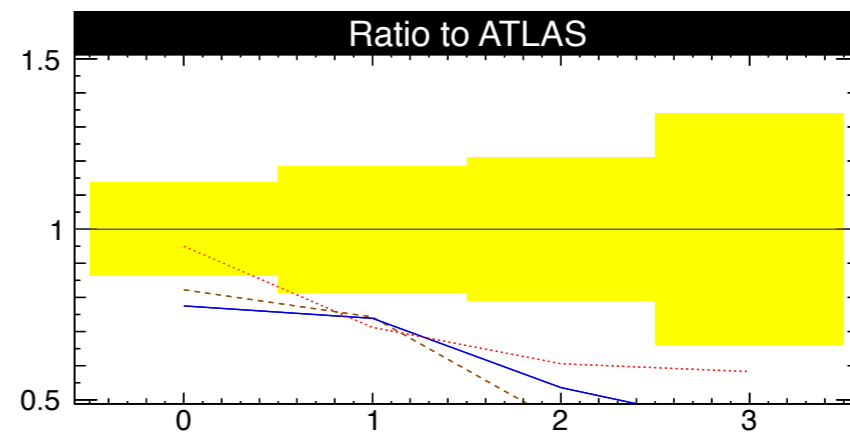
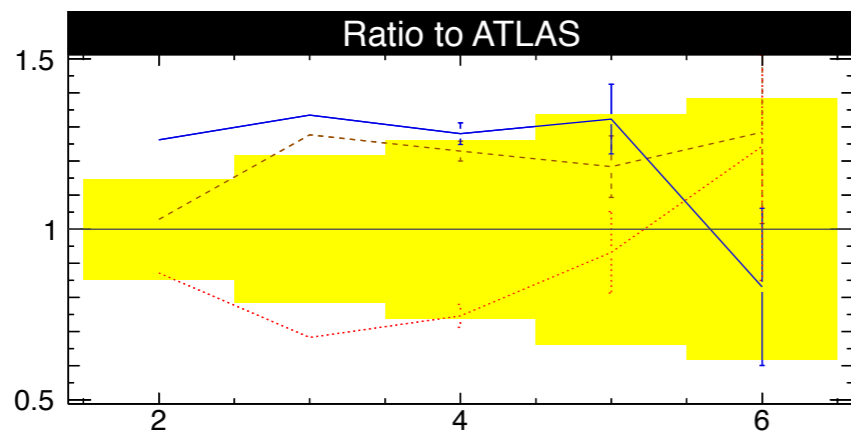
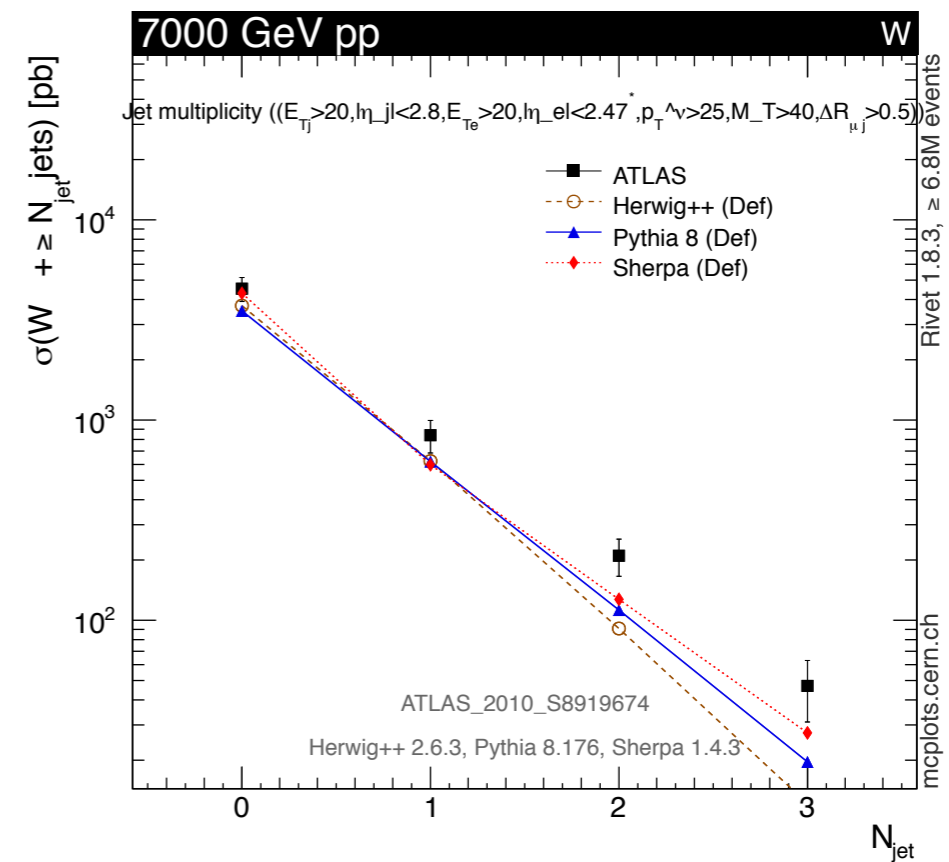
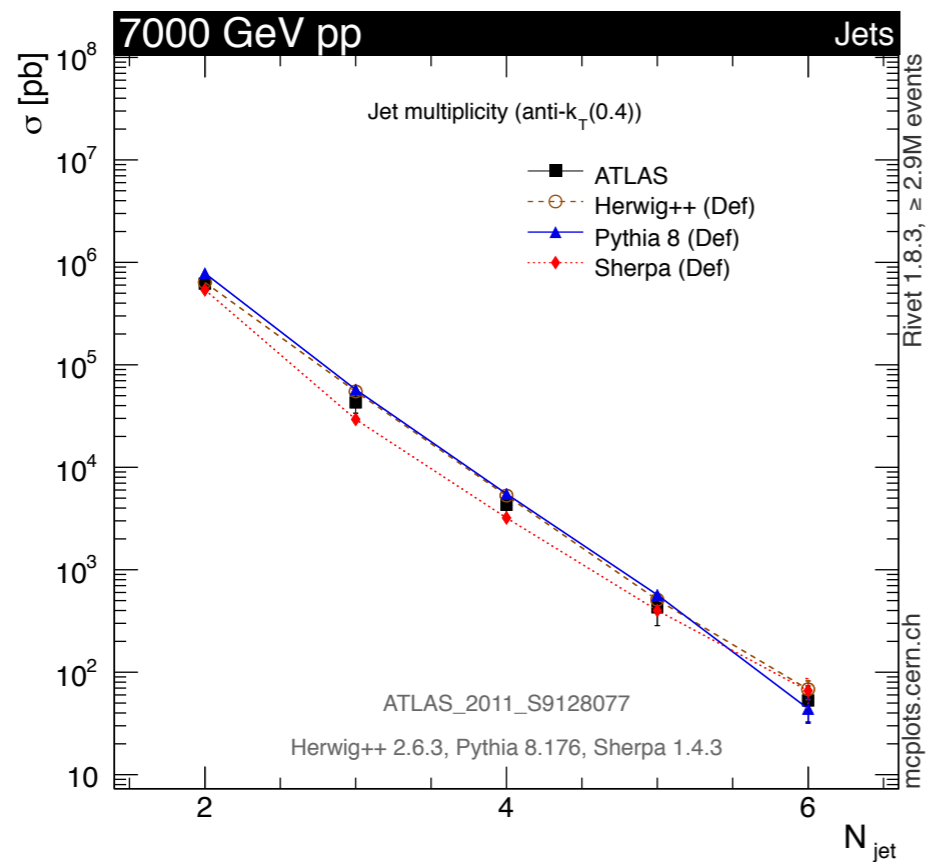


# Jet profile



$$\rho(r) = \frac{1}{\delta r} \frac{\sum_{r_a < r_i < r_b} p_{T,i}}{\sum_{r_i < R} p_{T,i}}$$

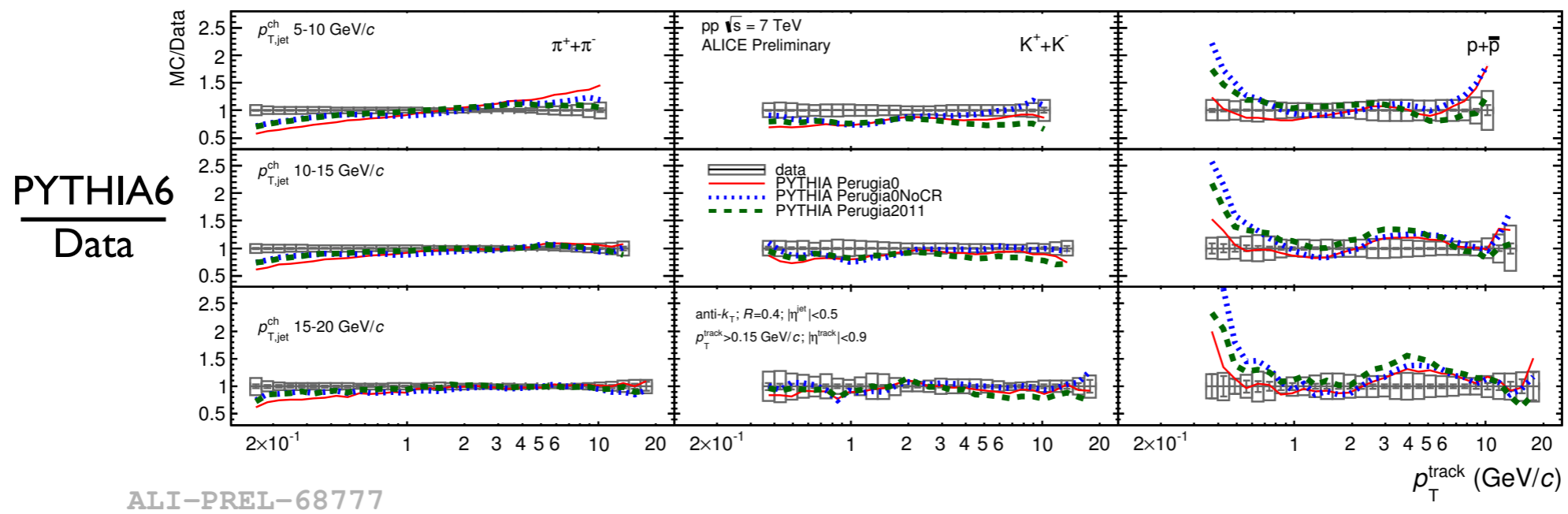
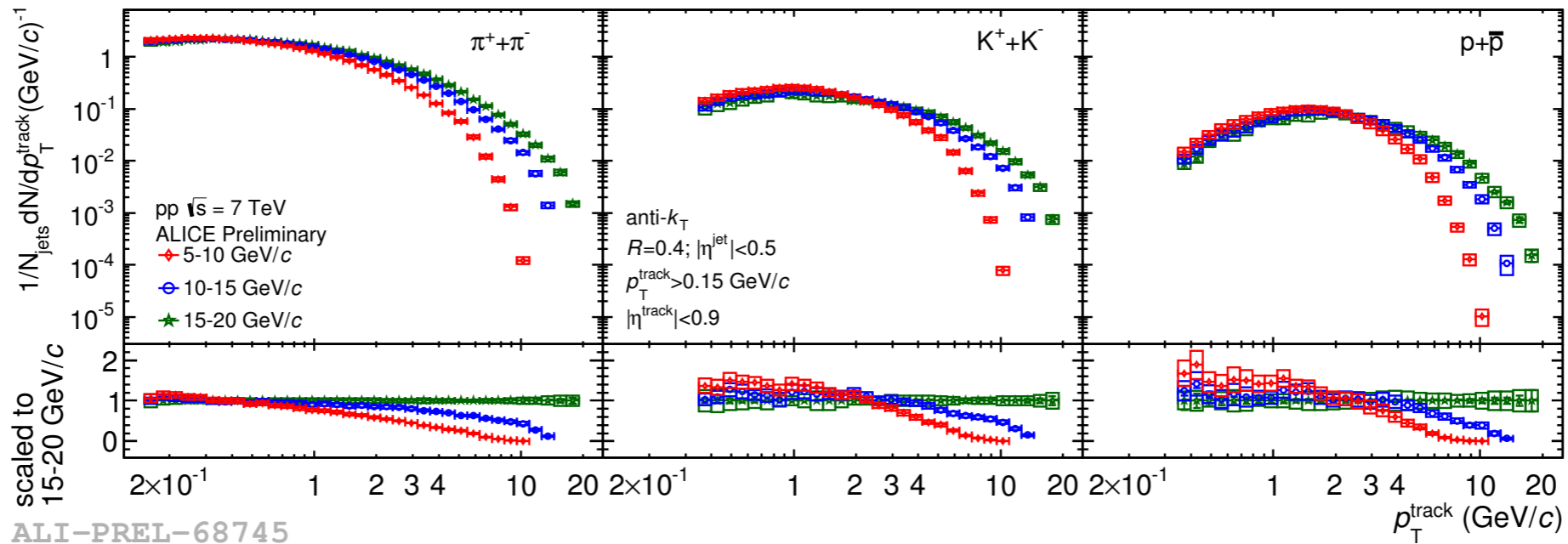
# Jet multiplicity



Multijets

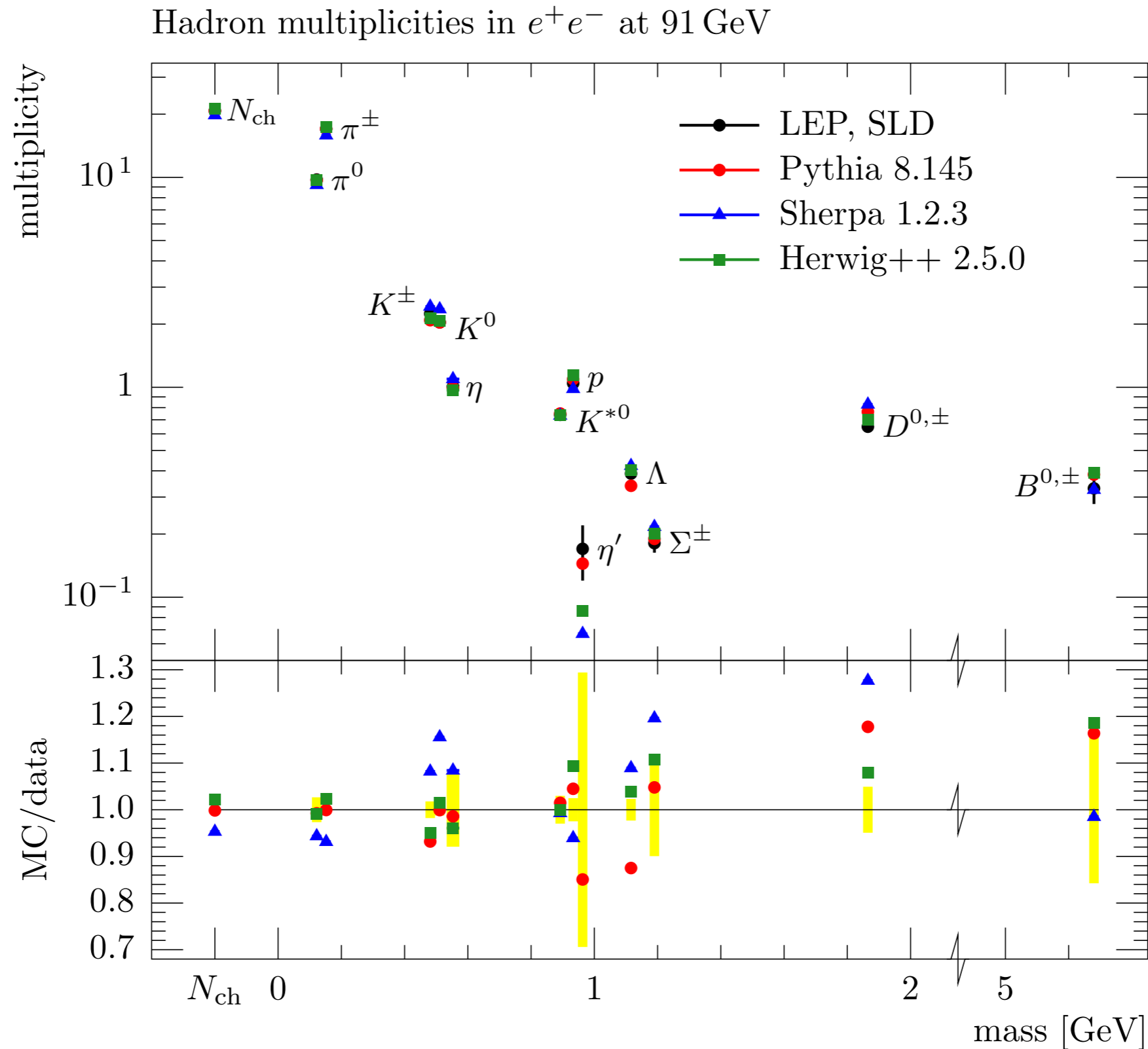
W+jets

# Charged particles in jets



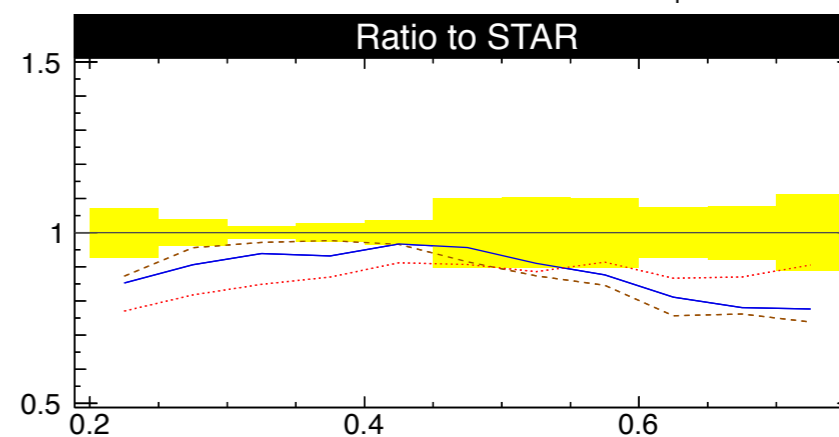
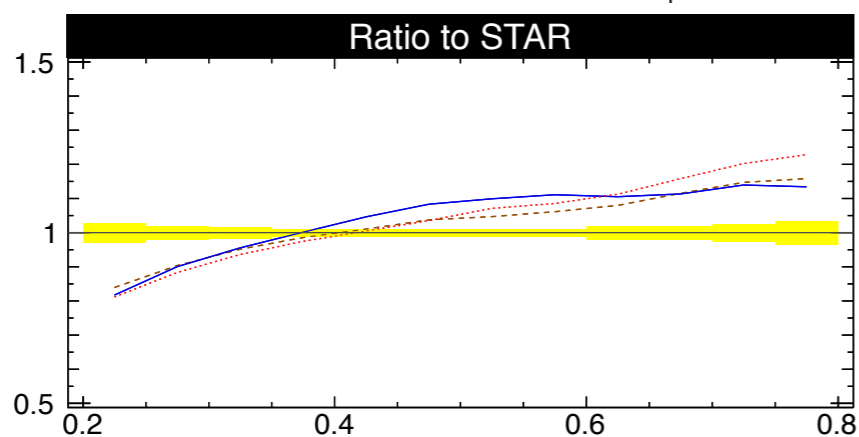
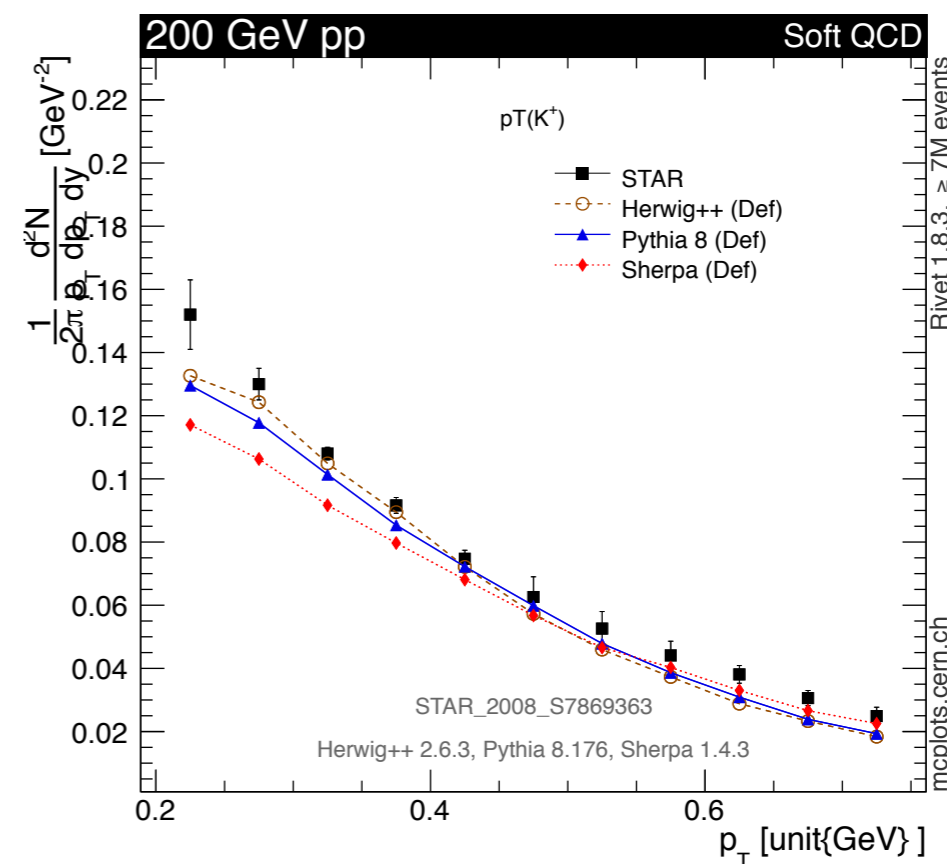
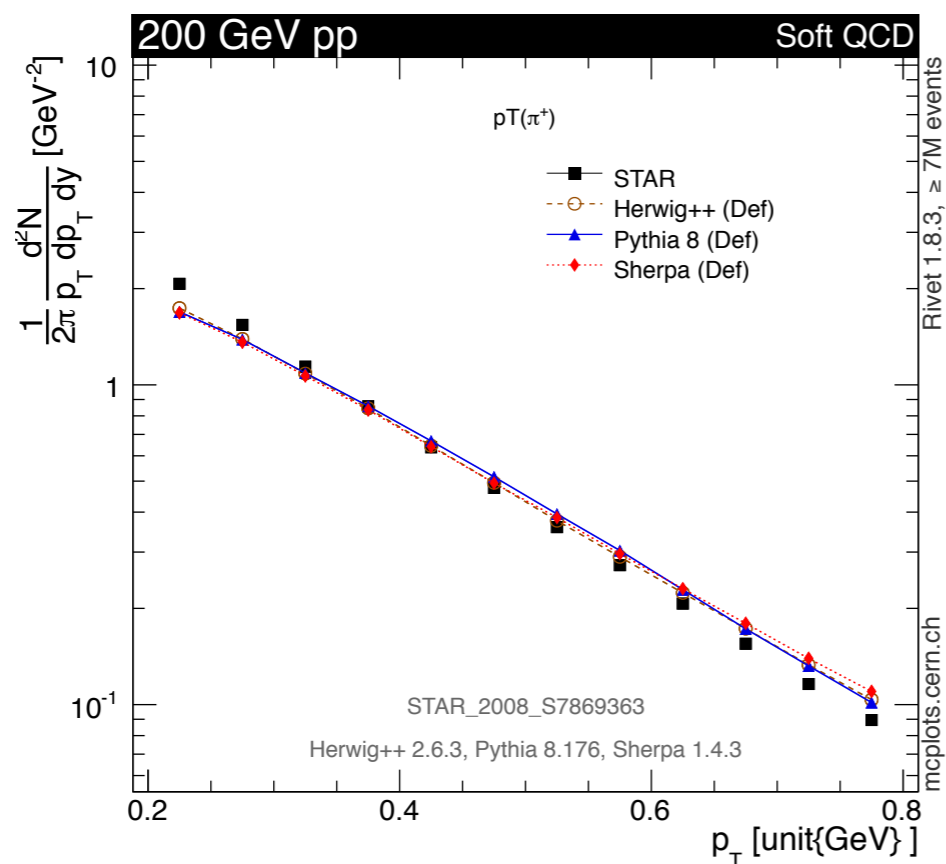
ALICE, arXiv:1408.5723

# Multiplicities in $Z^0$ decay



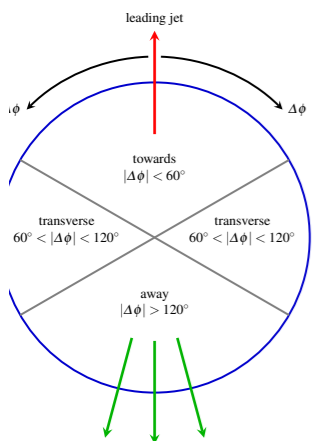
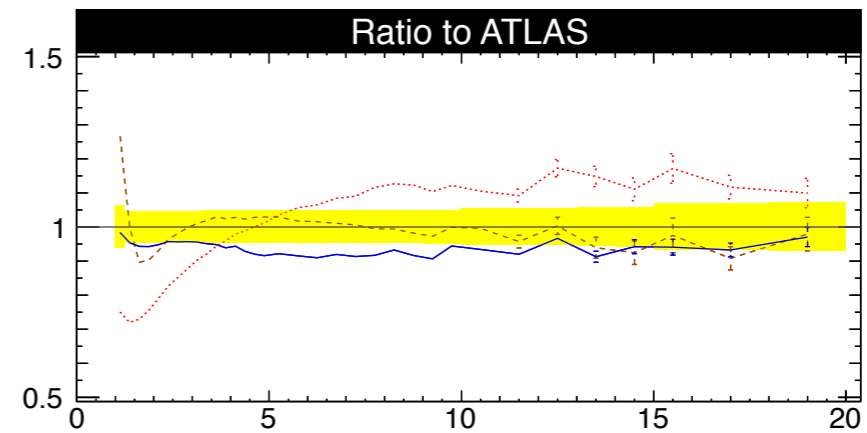
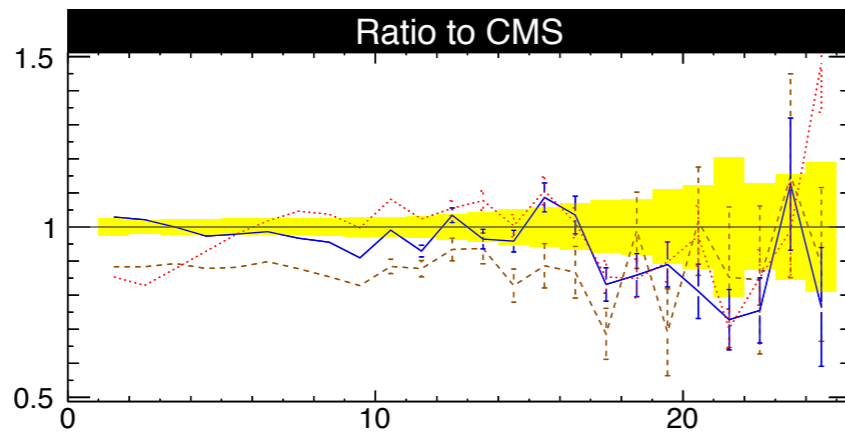
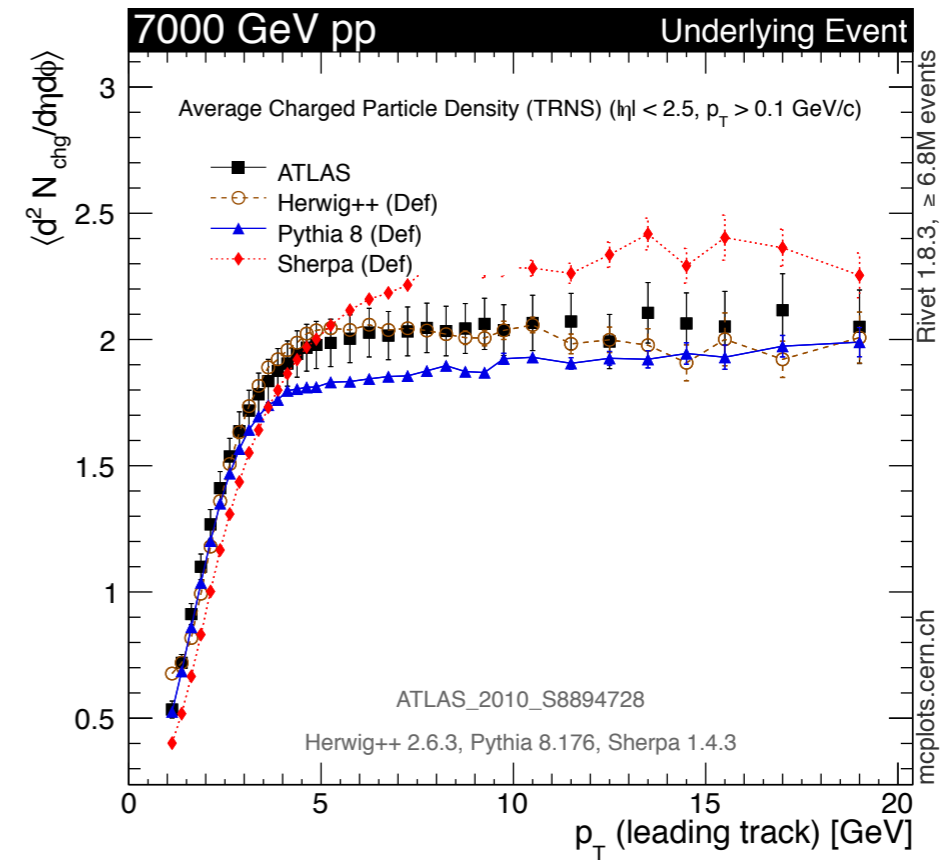
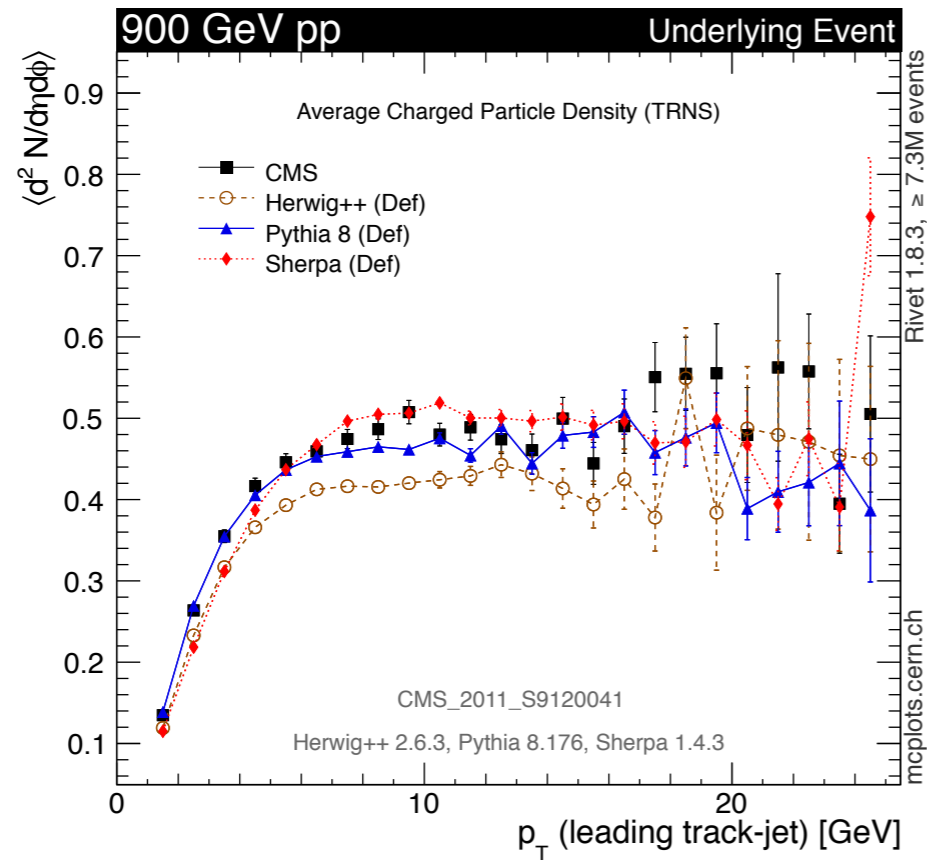
# Min Bias and Underlying Event

# Min bias $p_T(\pi^+, K^+)$



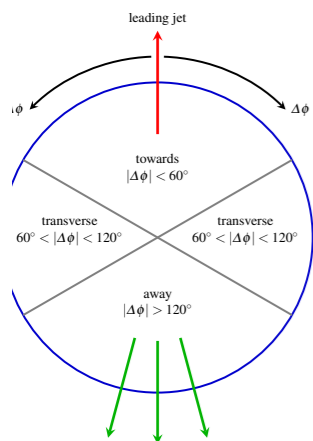
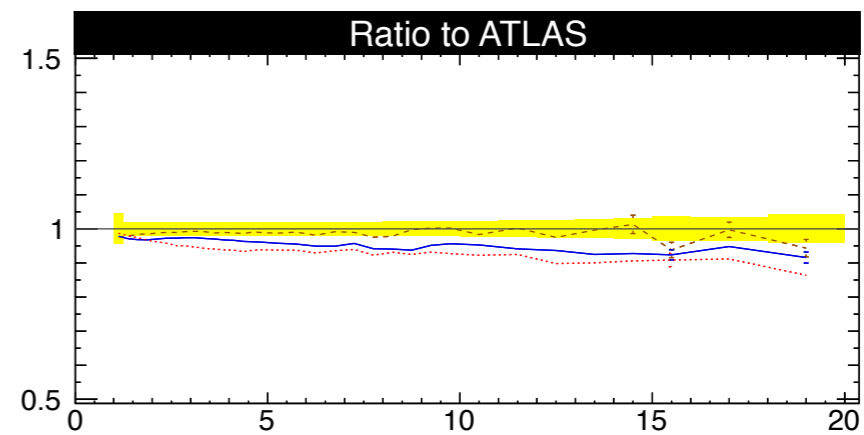
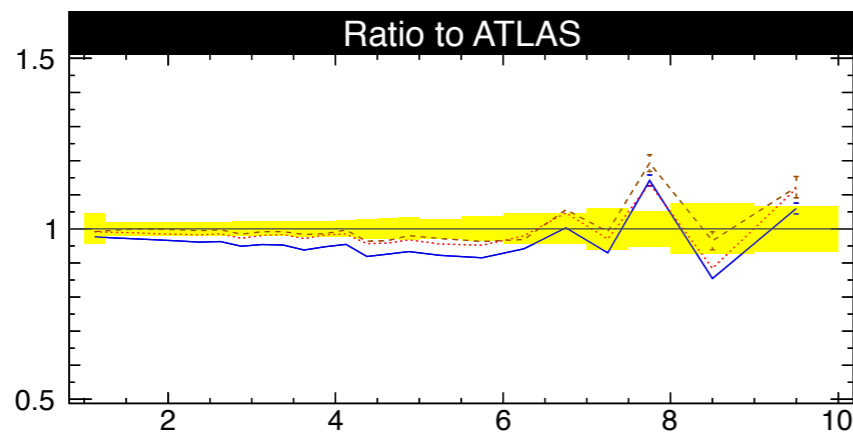
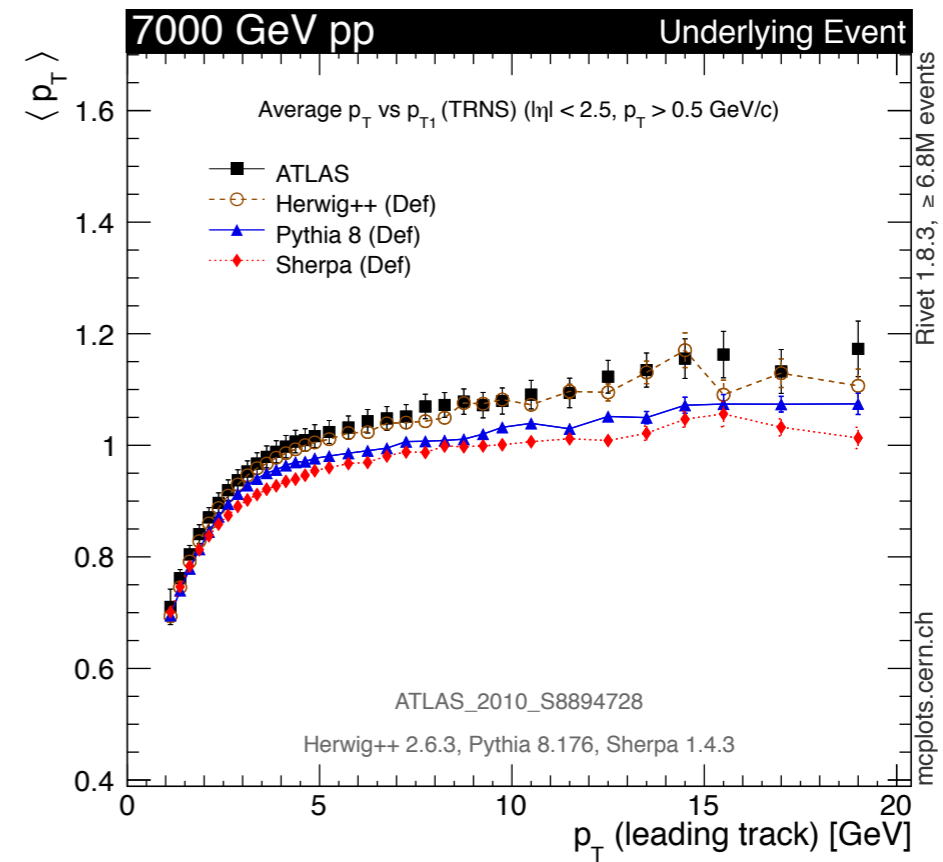
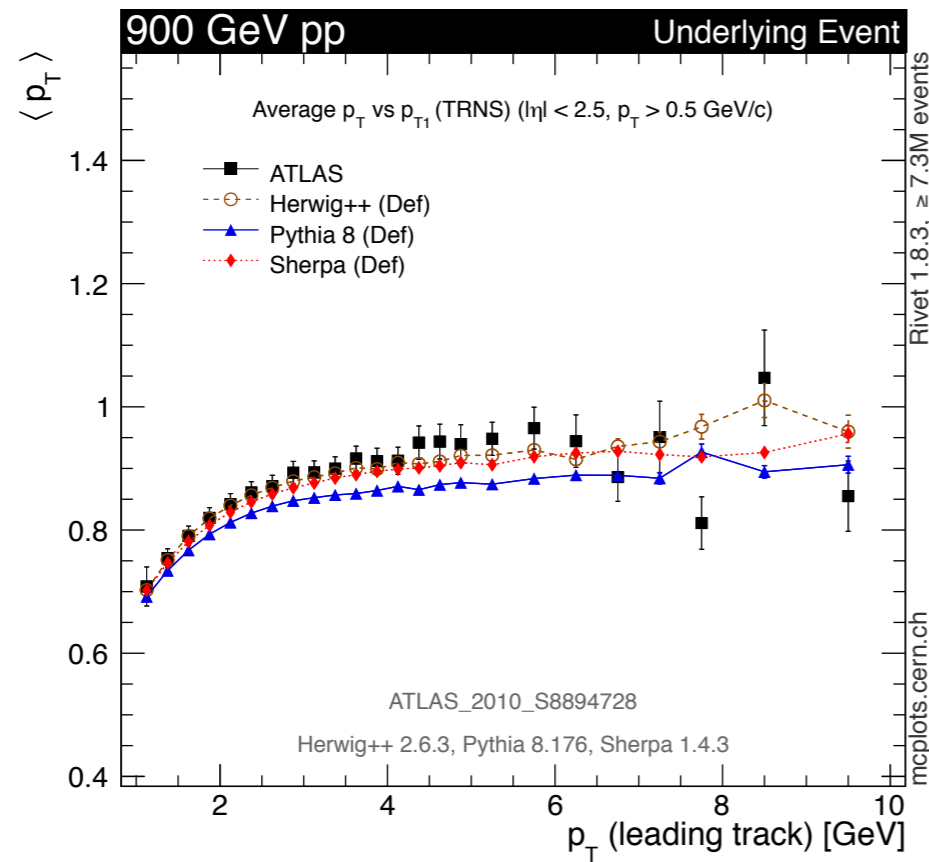
- Min bias = all scattering events

# Underlying Event



## Transverse charged particle density

# Underlying Event



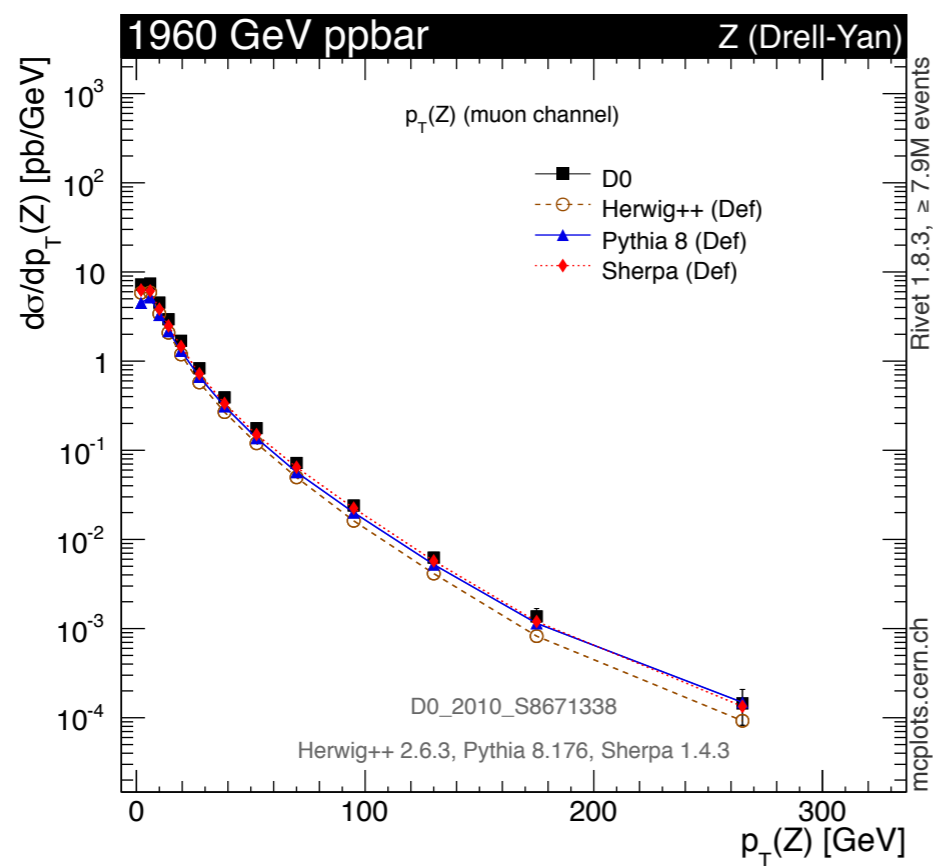
Transverse mean  $p_T$



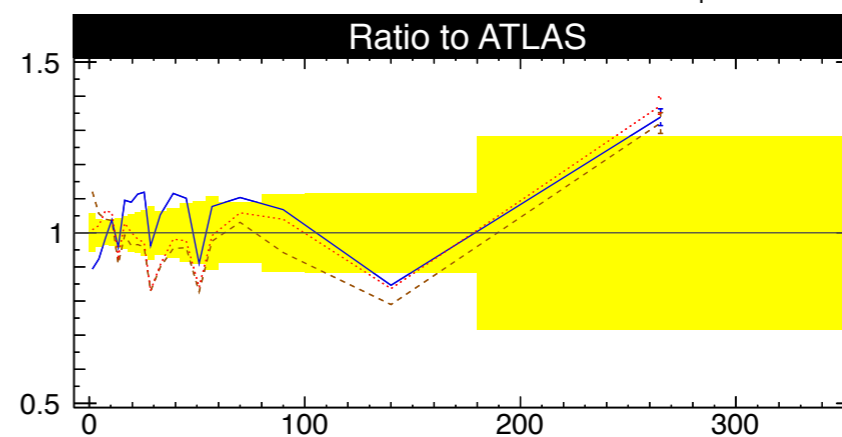
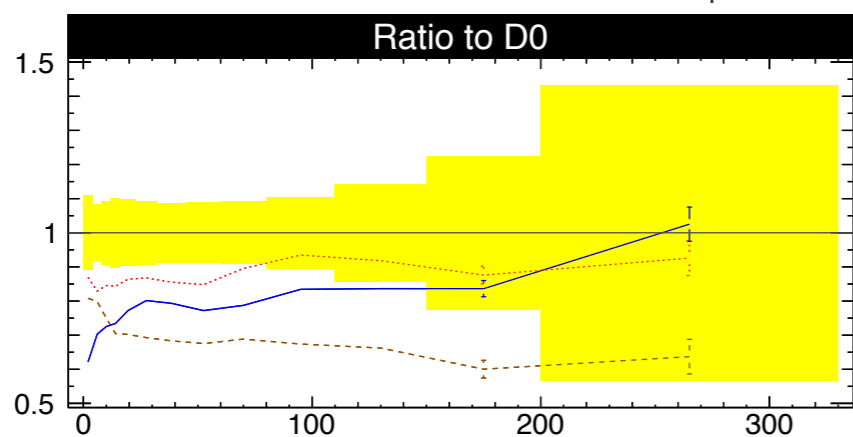
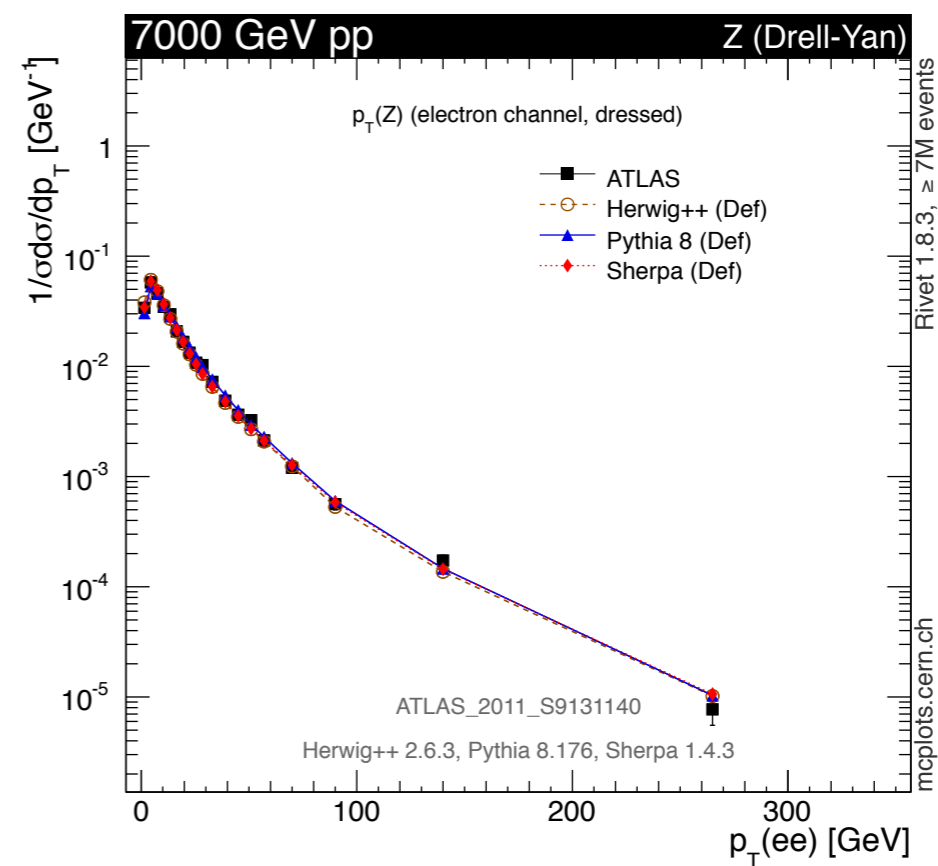
# Vector Bosons

# Z<sup>0</sup> p<sub>T</sub>

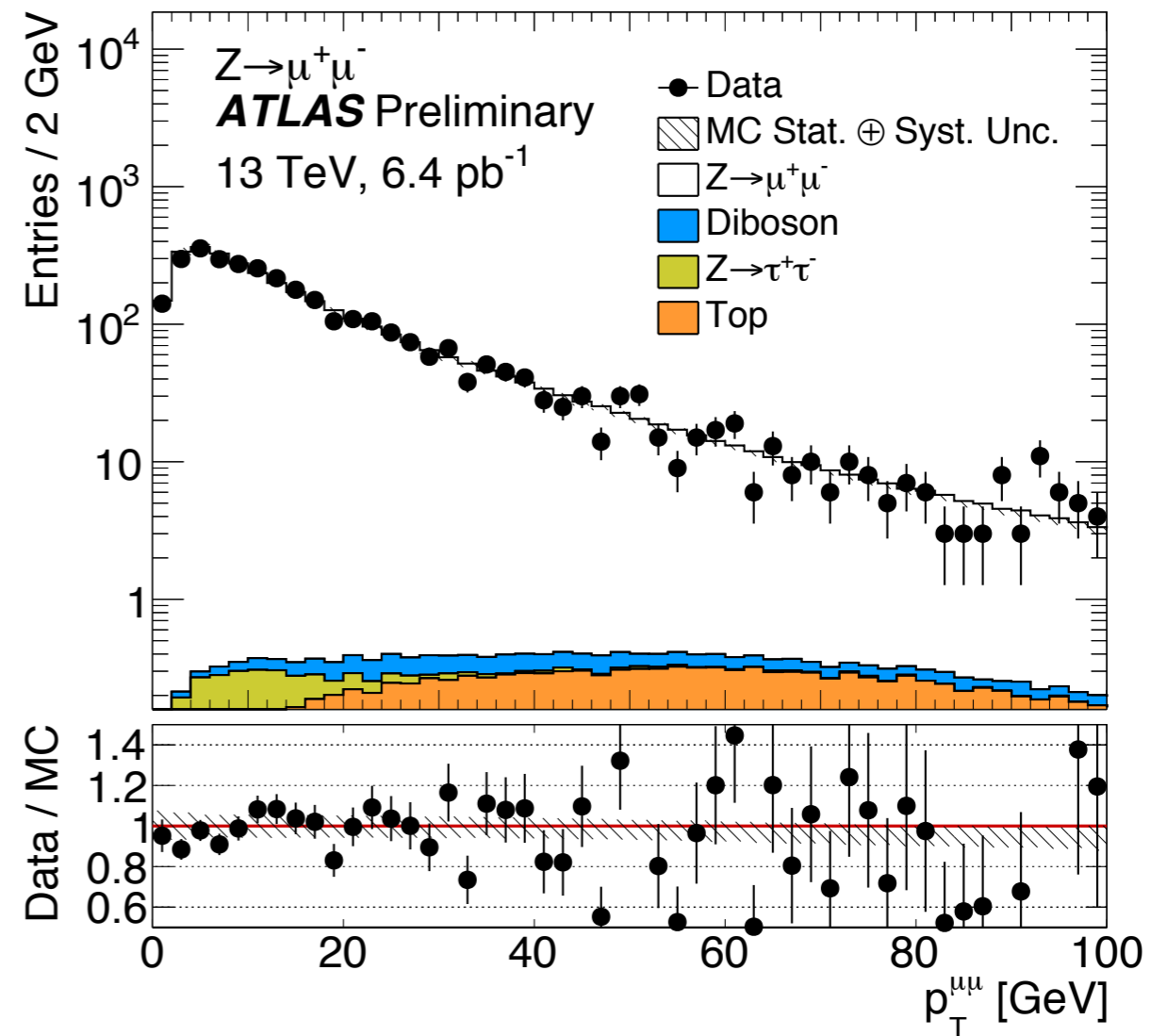
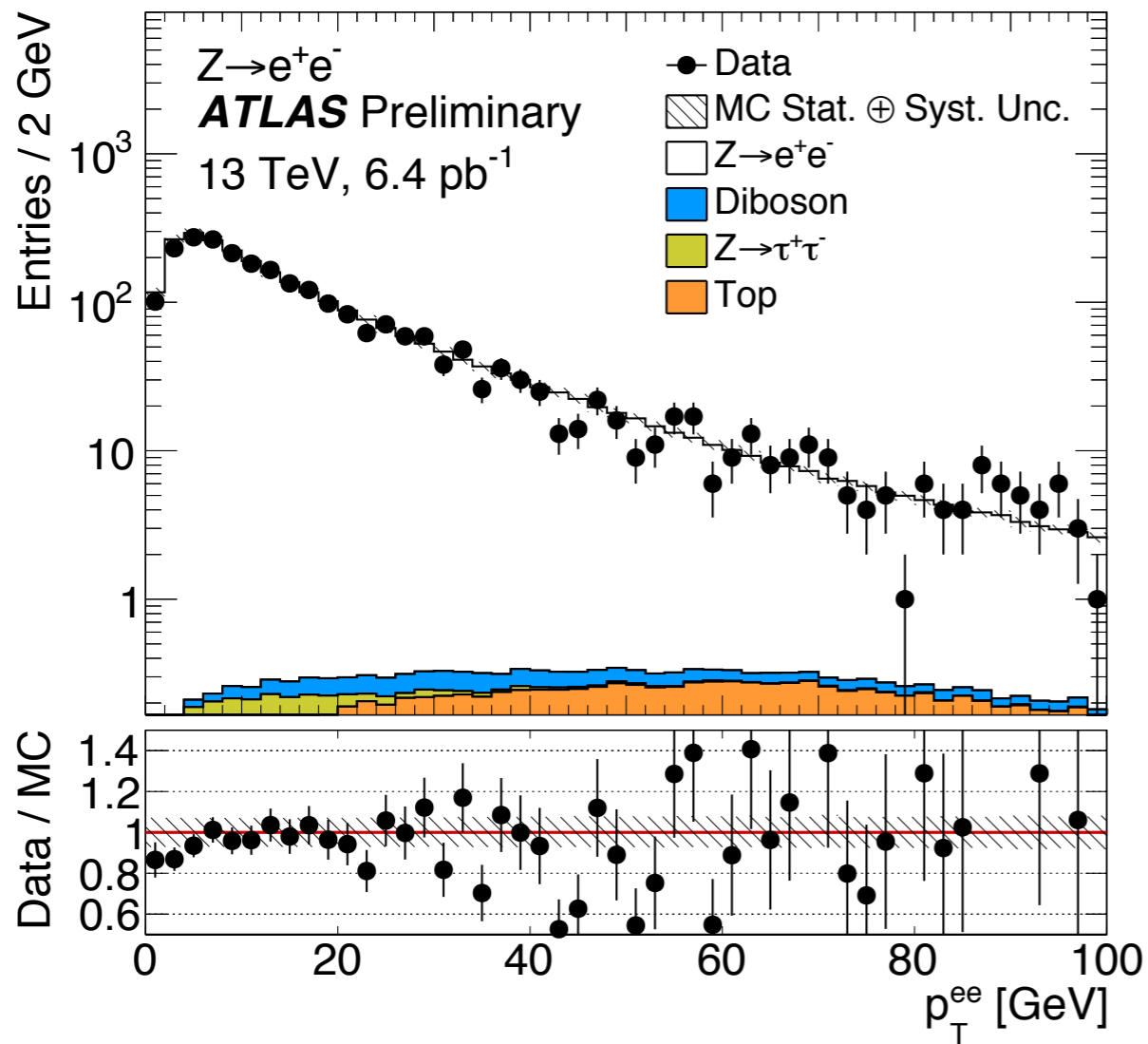
## Absolute normalization



## Normalized to data



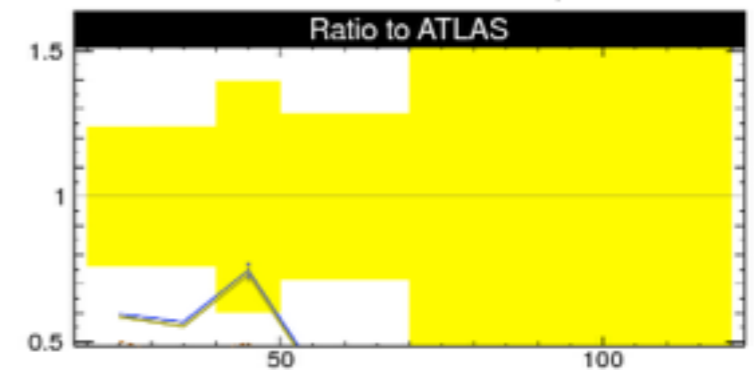
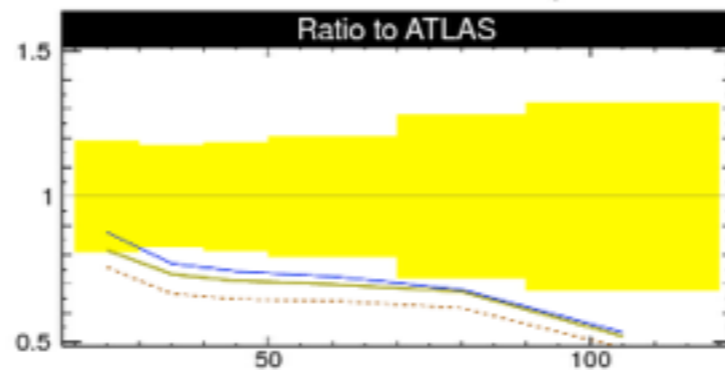
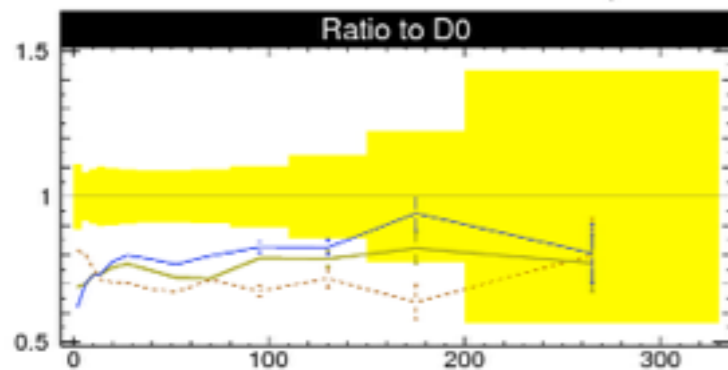
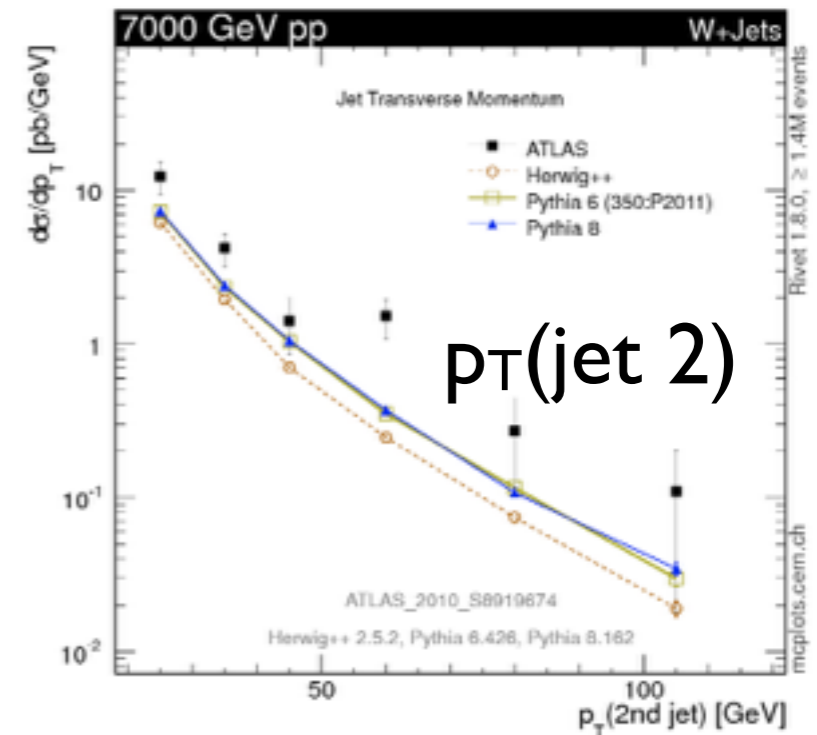
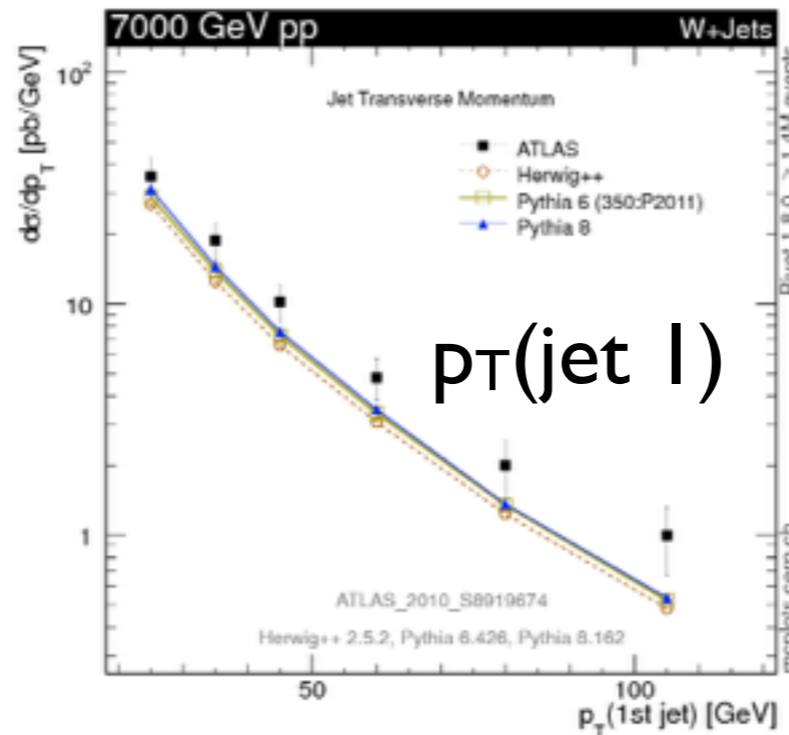
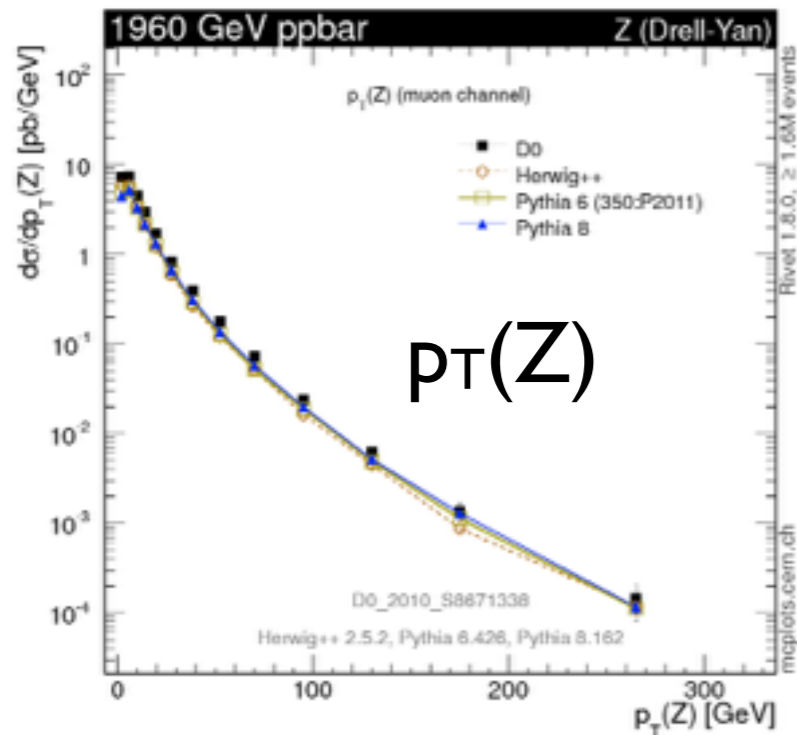
# $Z^0$ $p_T$ (13 TeV)



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# Limitations of LO+parton shower

- Hard process:  $q\bar{q} \rightarrow Z^0 / W^\pm$



- Leading-order (LO) normalization  $\Rightarrow$  need next-to-LO (NLO)
- Worse for high  $p_T$  and/or extra jets  $\Rightarrow$  need multijet merging

# Summary of Lecture 2

- Parton shower keeps largest small-angle contribution
- Shower gives preconfinement of colour
- This allows local model of hadronization
- String and cluster models both still viable
- Underlying event due to multiple interactions
- Sample of event generator results
- Further improvements (matching & merging) now used

**Thanks for listening!**