

XIV International Workshop on Hadron Physics

*High energy
heavy-ion collisions
- hot QCD in a lab*

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Berkeley Lab*

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Florianópolis, SC, Brazil**

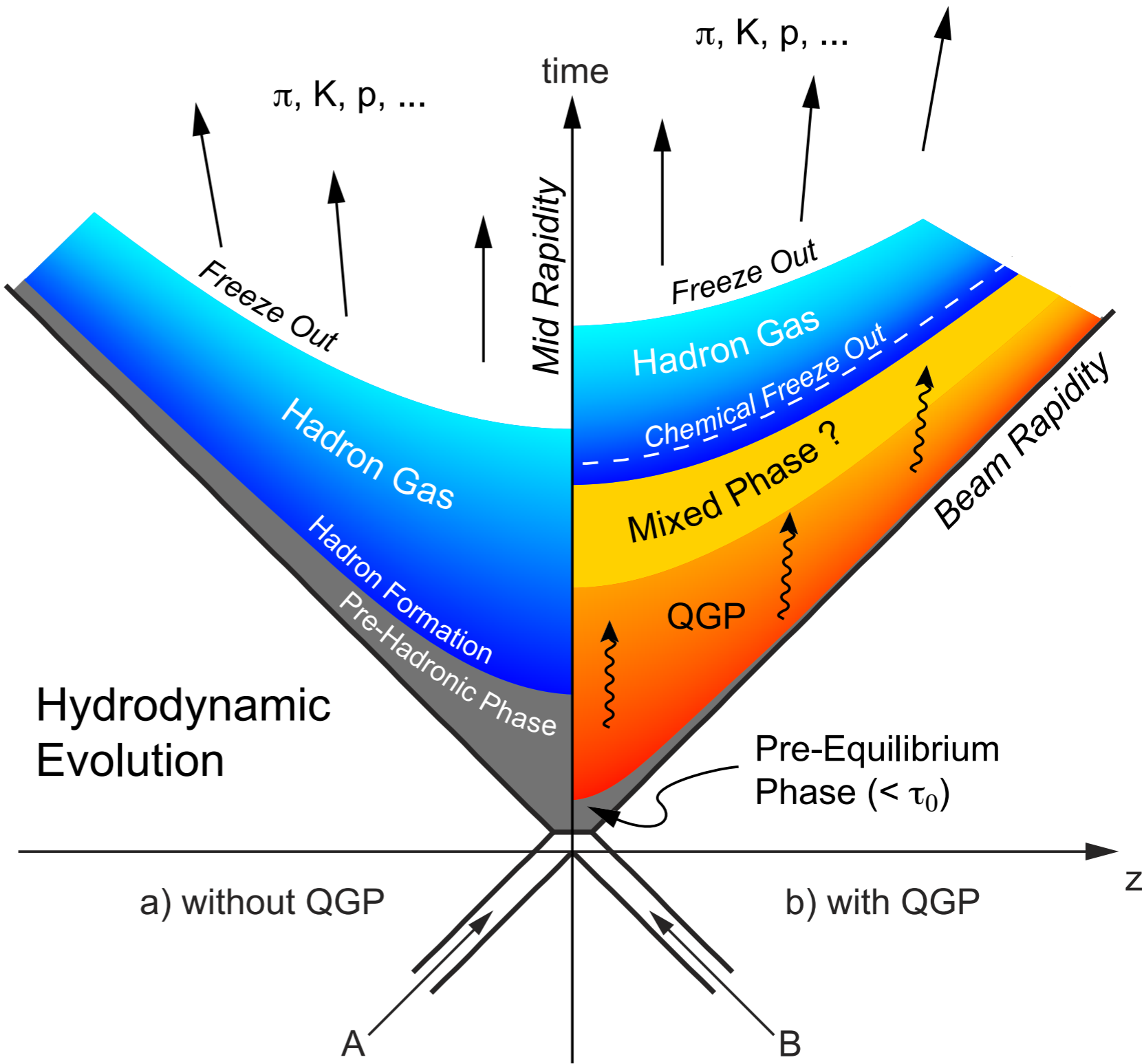
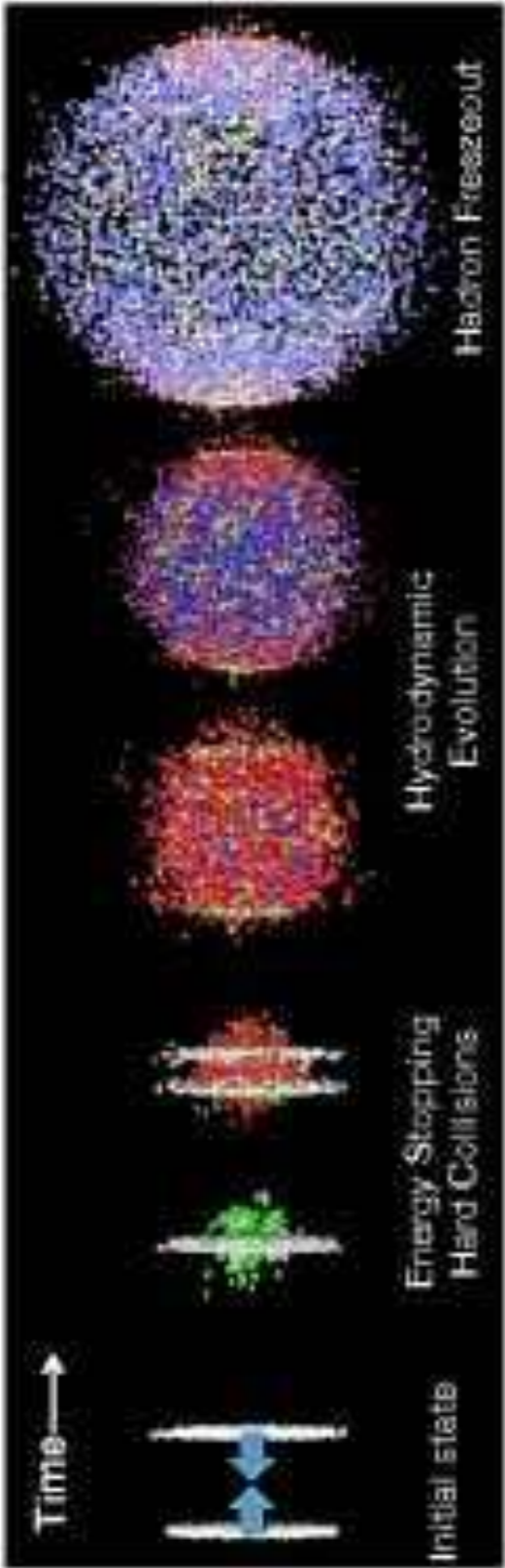
until now...

- *phases of HI collision*
- *how to measure centrality of a collision*
- *... energy density*
- *... temperature*
- *... freeze-out volume (and time)*
- *QGP: hot, short-lived system with rapid dynamical evolution*

Today (2/3)

- particle abundance's at hadronization
- QGP properties / transport coefficients:
 - flow
 - jet-medium interactions

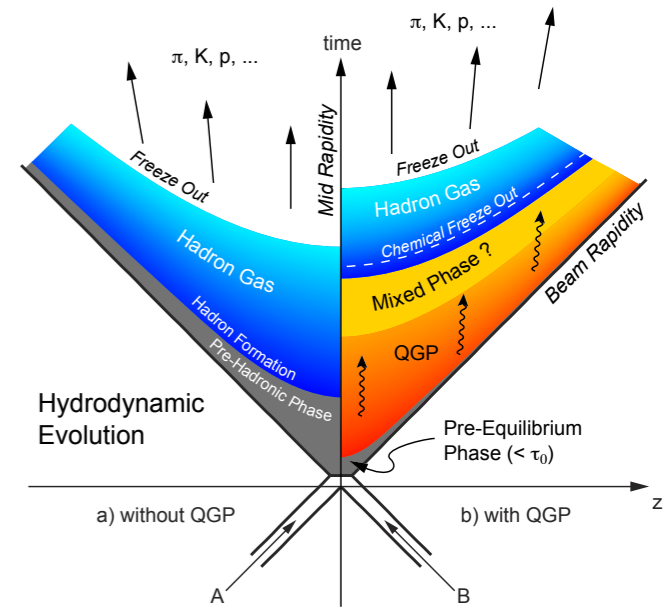
Collision timeline



... to follow-up: chemical & kinetic freeze-out

Freeze-out:

- *chemical freeze out \Leftrightarrow hadron composition fixed*
- *kinetic freeze-out \Leftrightarrow hadron momenta fixed (interactions stop)*
- *overall: $T_{ch} > T_{kin}$ (system cools down - follow the time axis)*



Thermal equilibrium...

Chemical and kinetic freeze-out

Chemical equilibrium:

- correct relative particle abundances?
- large system \rightarrow Grand Canonical ensemble: many particles; conservation laws on average - chemical potentials
- small system \rightarrow conservation laws E-by-E \rightarrow "canonical suppression" (strangeness)

$$n_i^0 = \frac{g_i}{2\pi^2} \int \frac{p^2 dp}{e^{(E - \mu_B B_i - \mu_s S_i - \mu_3 I^3)/T} \pm 1}$$

The ratios of produced particle yields between various species can be fitted to determine T, μ .

Kinetic equilibrium - radial flow:

- for any interacting system of particles expanding into vacuum, radial flow is a natural consequence.

During the cascade process, an ordering of particles with the highest common underlying velocity at the outer edge develops naturally

Hadrons are released in the final stage and therefore measure "FREEZE-OUT" Temp. - instructive simple parametrization - radially boosted source with velocity β and at $y=0$:

$$\frac{d^3 N}{dp^3} \propto e^{-E/T}; E \frac{d^3 N}{dp^3} = \frac{d^3 N}{m_T dm_T d\phi dy} \propto E e^{-E/T} = m_T \cosh(y) e^{-m_T \cosh(y)/T}$$

$$\frac{1}{m_T} \frac{dN}{dm_T} \propto m_T I_0 \left(\frac{p_T \sinh(\rho)}{T} \right) K_1 \left(\frac{m_T \cosh(\rho)}{T} \right)$$

$$\rho = \tanh^{-1}(\beta_{\text{boost}})$$

Simple assumption: uniform sphere of radius R and boost velocity varies linearly w/ r :

$$\frac{1}{m_T} \frac{dN}{dm_T} \propto \int_0^R r^2 dr m_T I_0 \left(\frac{p_T \sinh(\rho)}{T} \right) K_1 \left(\frac{m_T \cosh(\rho)}{T} \right)$$

$$\rho(r) = \tanh^{-1} \left(\beta_T^{\text{MAX}} \frac{r}{R} \right)$$

Blast Wave model
 \Rightarrow common T and β

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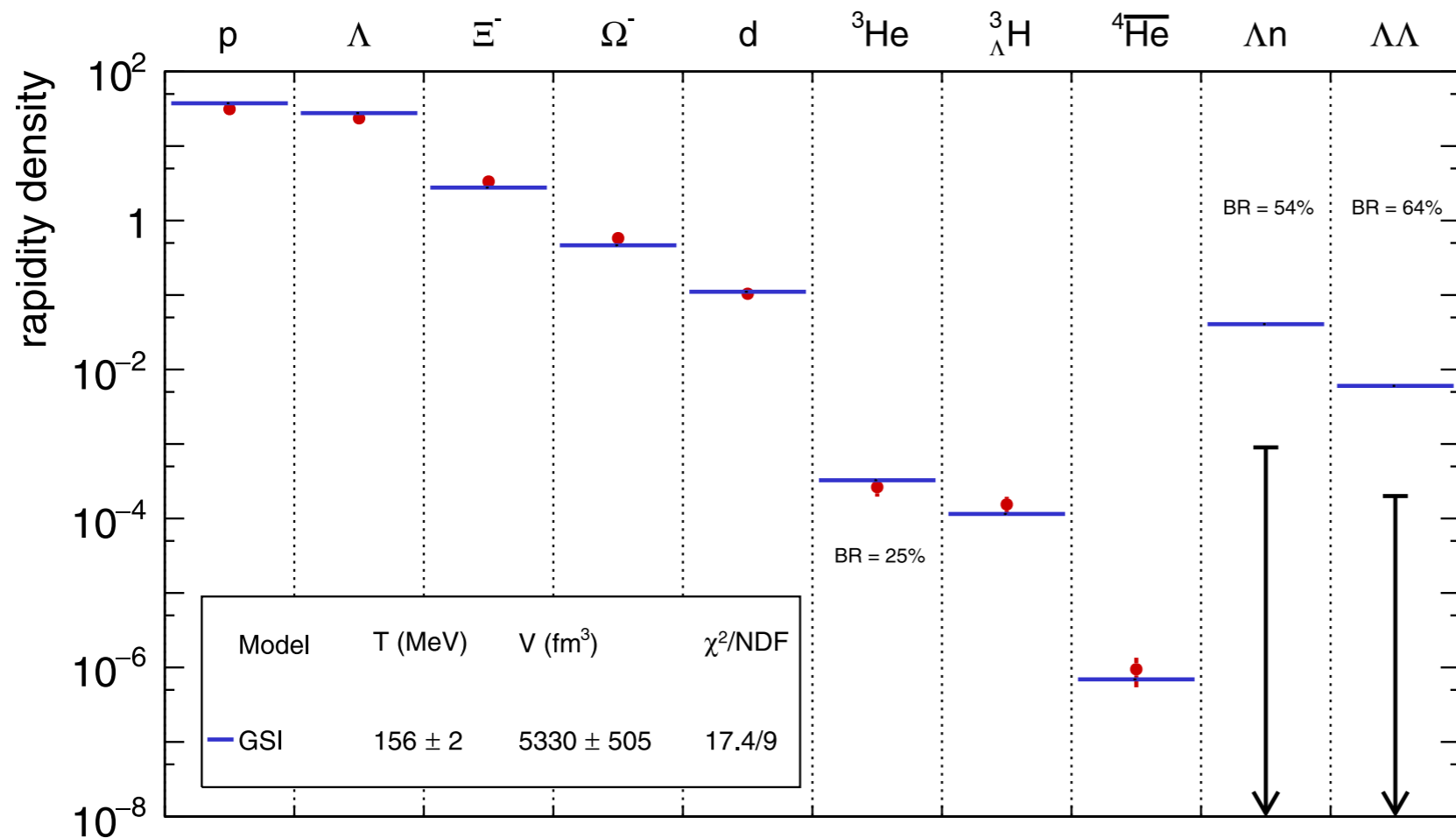
Hadron abundances

Assumption: Multiplicities are determined by statistical weights (chemical equilibrium)

Grand-canonical ensemble: $\langle n_j \rangle = \frac{(2J_j + 1)V}{(2\pi)^2} \int d^3p \left[e^{(\sqrt{p^2 + m_j^2} + \mu_{q_i})/T} \pm 1 \right]^{-1}$

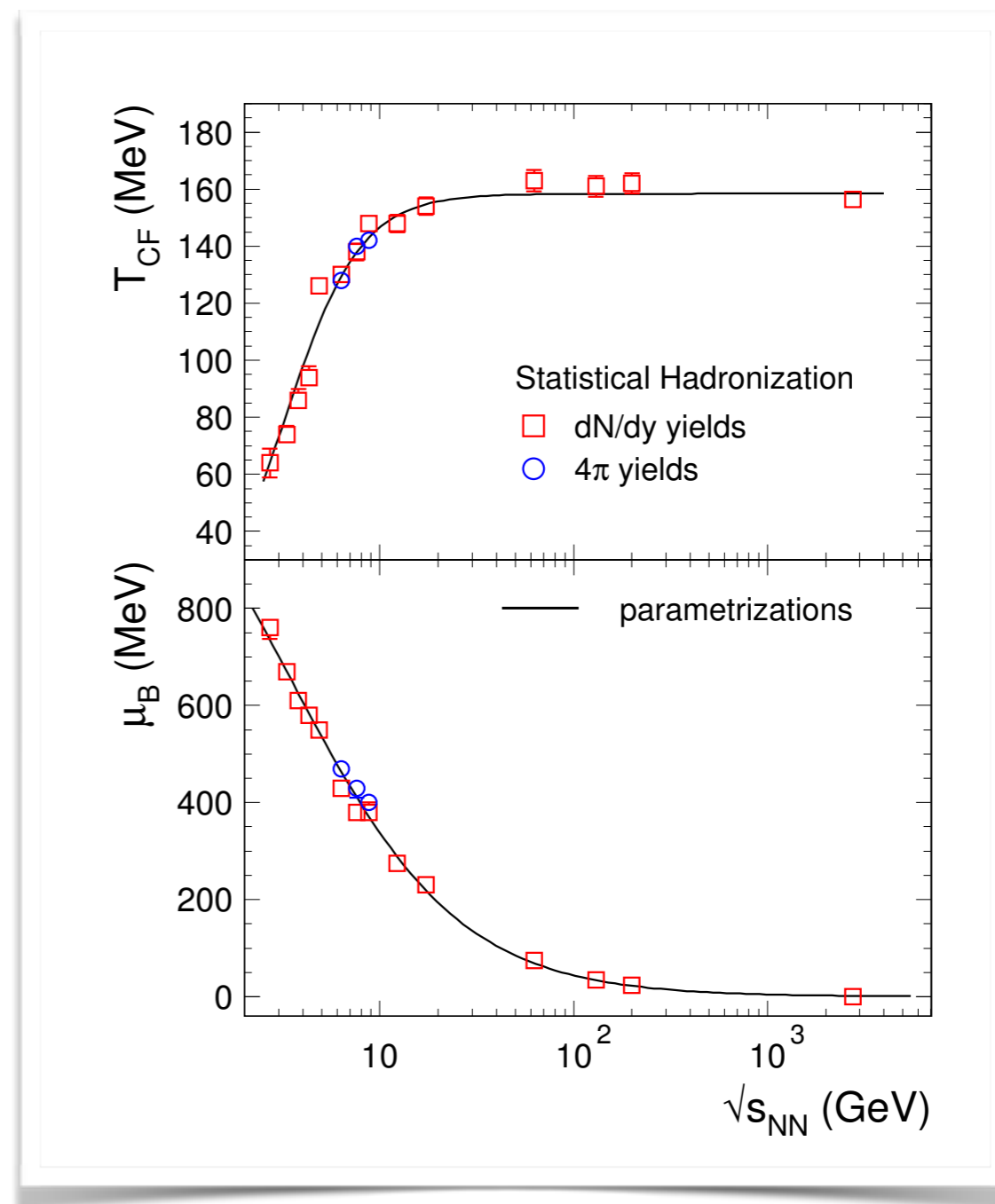
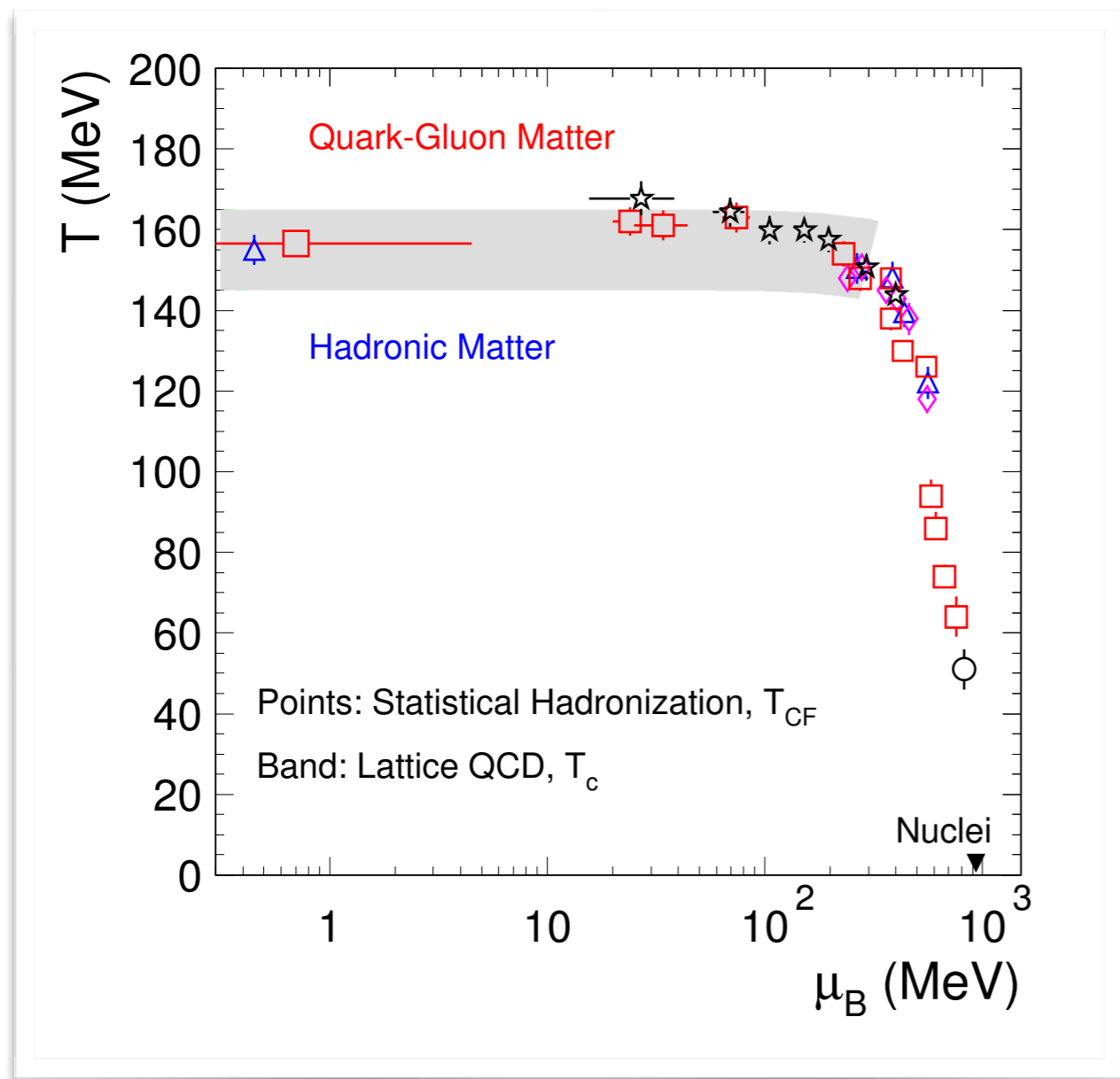
Parameters: $V, T, \mu_B, (\gamma_S)$

Results in excellent fits to measured multiplicities of hadron for ALL energies (even $d, {}^3\text{He}, {}^3_\Lambda\text{He} \dots$) - statistical harmonization of a thermal system...



NB: works also for $p+\bar{p}$
(phase space dominance,
Fermi 1950)

Chemical freeze out systematics



Provides rough idea which region in T, μ are probed

Thermal equilibrium...

Chemical and kinetic freeze-out

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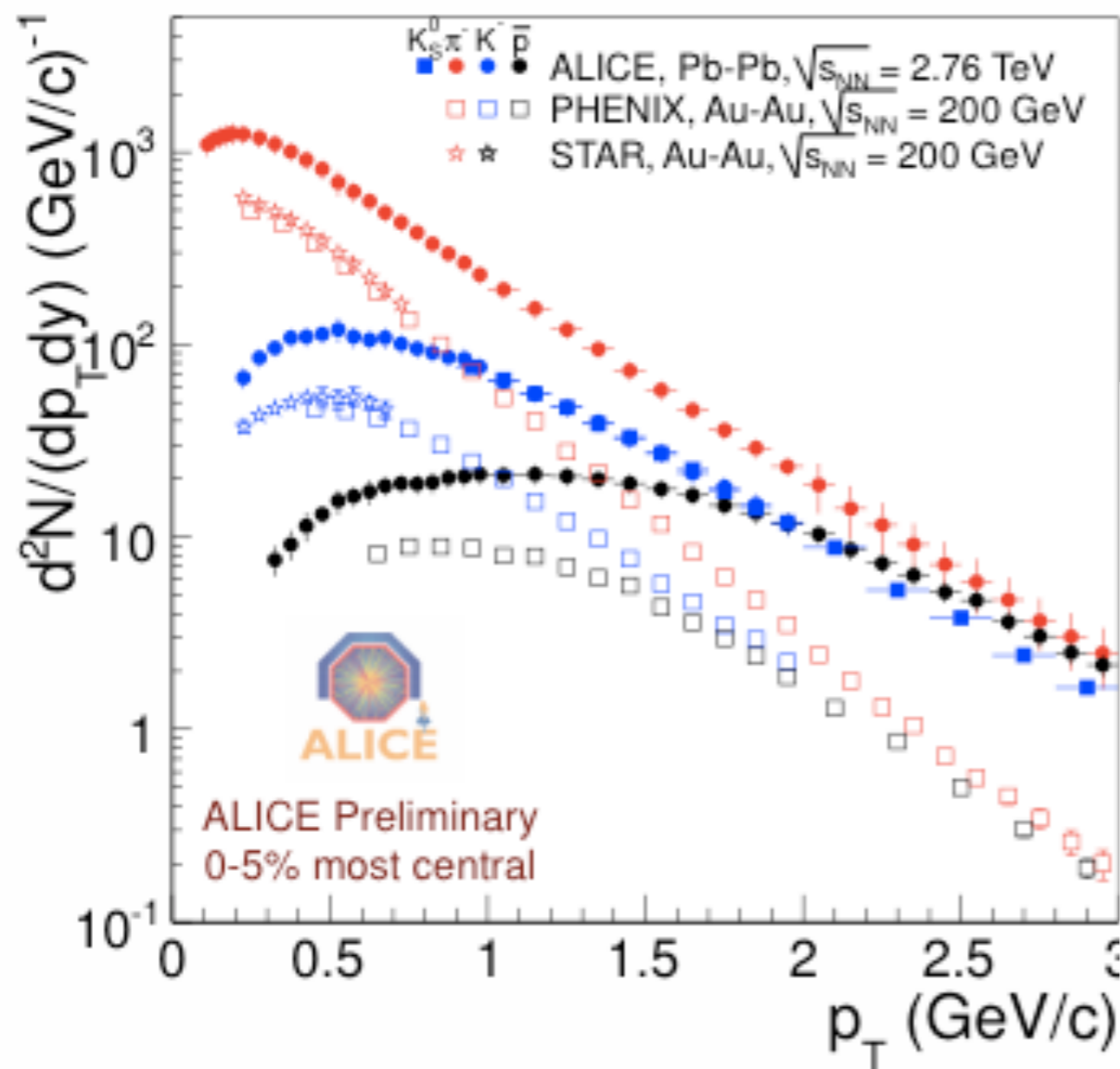
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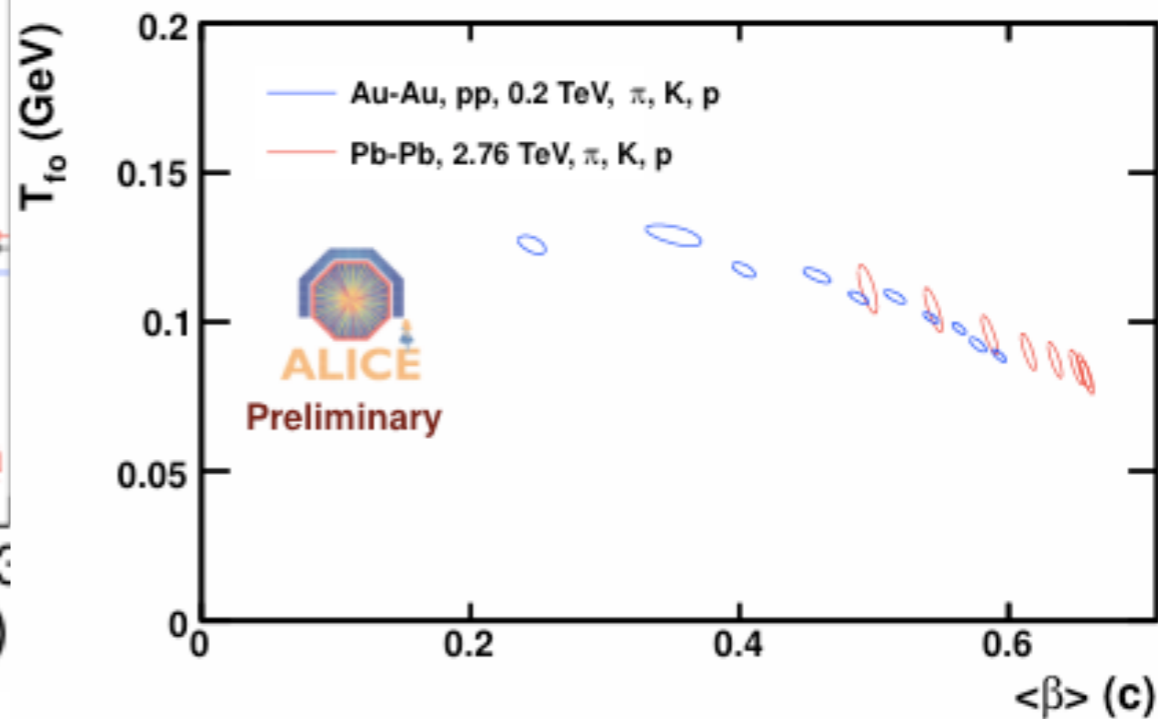
Blast Wave model
 \Rightarrow common T and β

Identified particles & expansion of the system



Stronger radial flow at the LHC.

“Blast wave” fits to spectra indicate an **increase of the average radial boost velocity** up to $(2/3)c$ and a decrease in the kinetic freezeout temperature to just below 100 MeV relative to RHIC



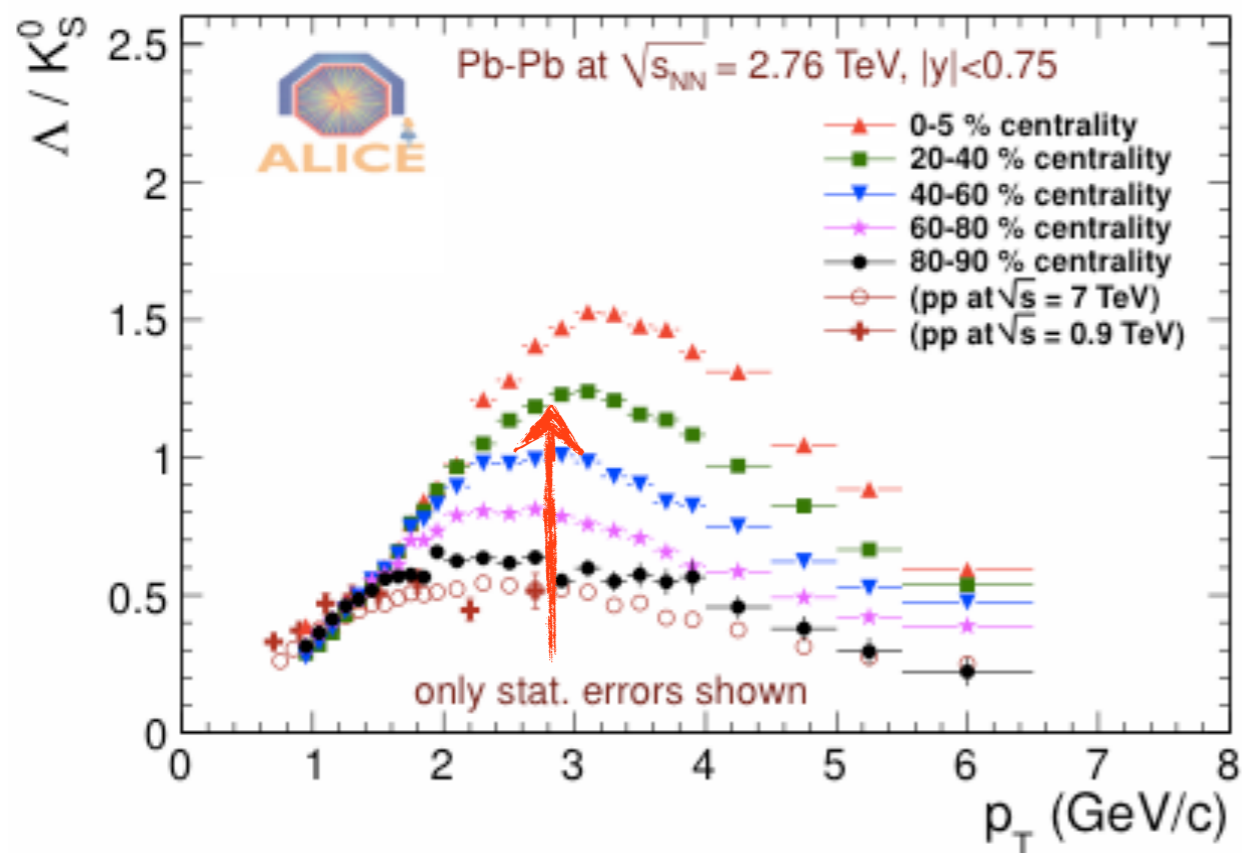
LHC: Large kinematic reach to explore

ALICE: excellent particle identification capabilities at the LHC

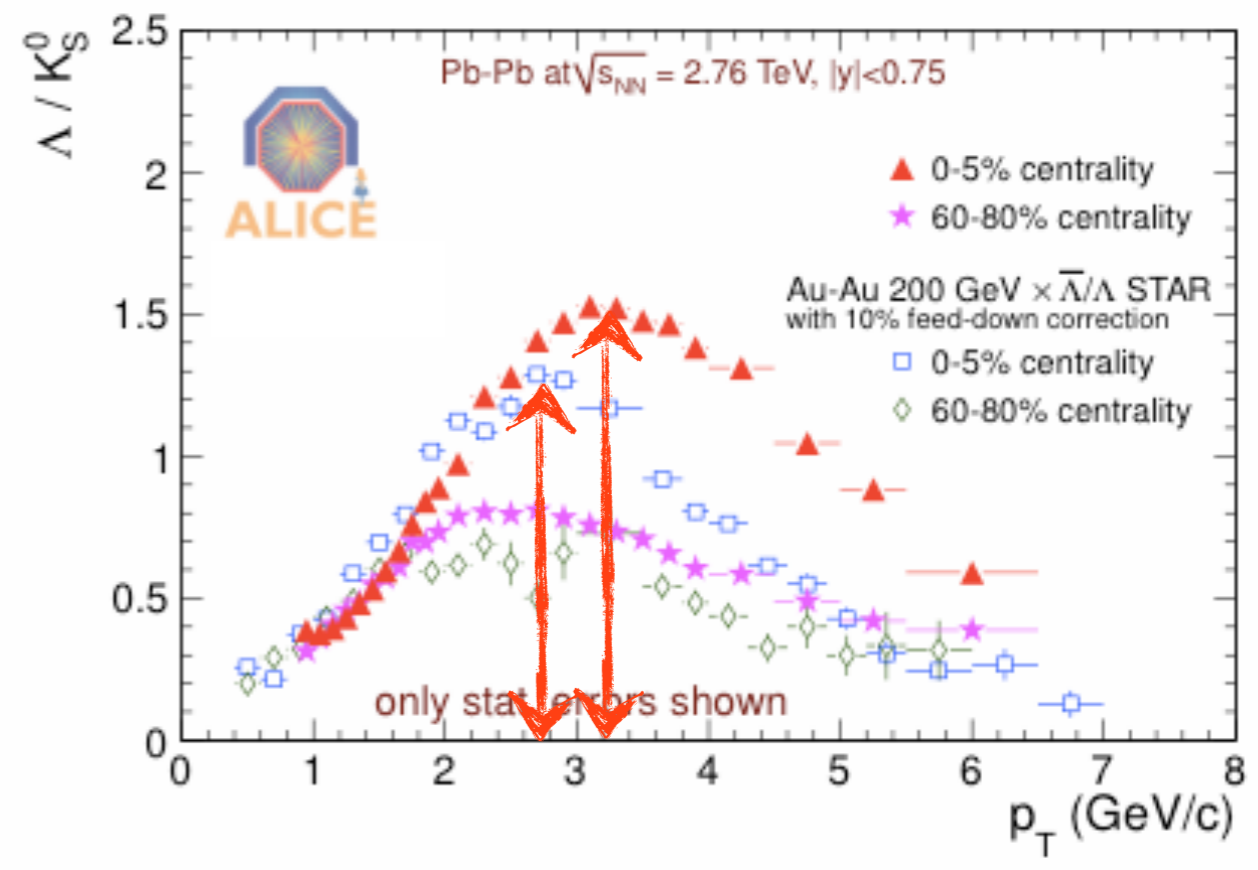
Impact of expansion on hadron p_T -spectra in heavy-ion collisions

A quick analysis of particle spectra ...

RHIC vs LHC
(LHC: higher mean p_T - more flow)



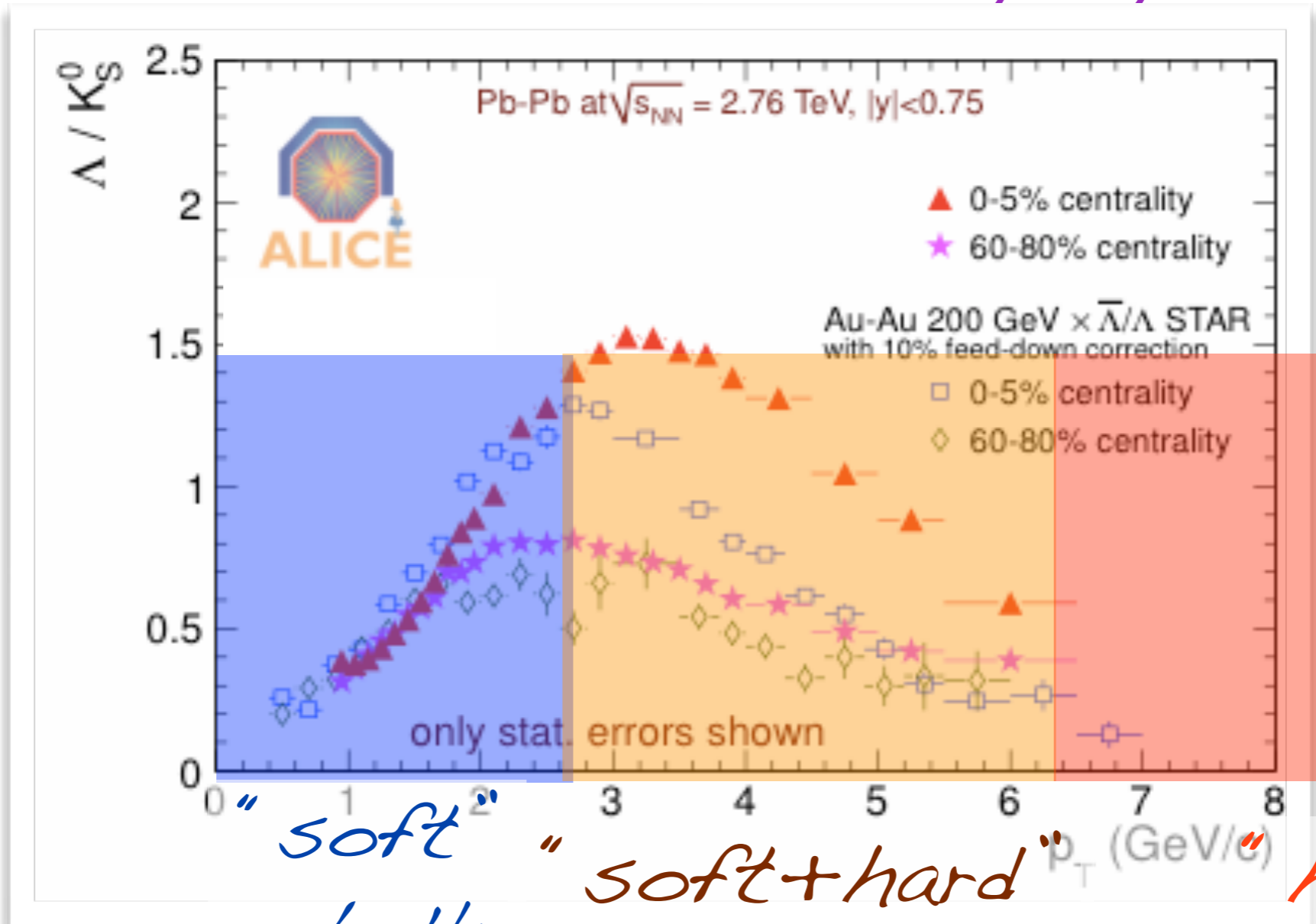
Baryons / Mesons



Much more baryons than mesons in central collisions as compared to proton-proton (coalescence/recombination? bulk+jet?)

LHC similar to RHIC
Maximum at slightly higher- p_T

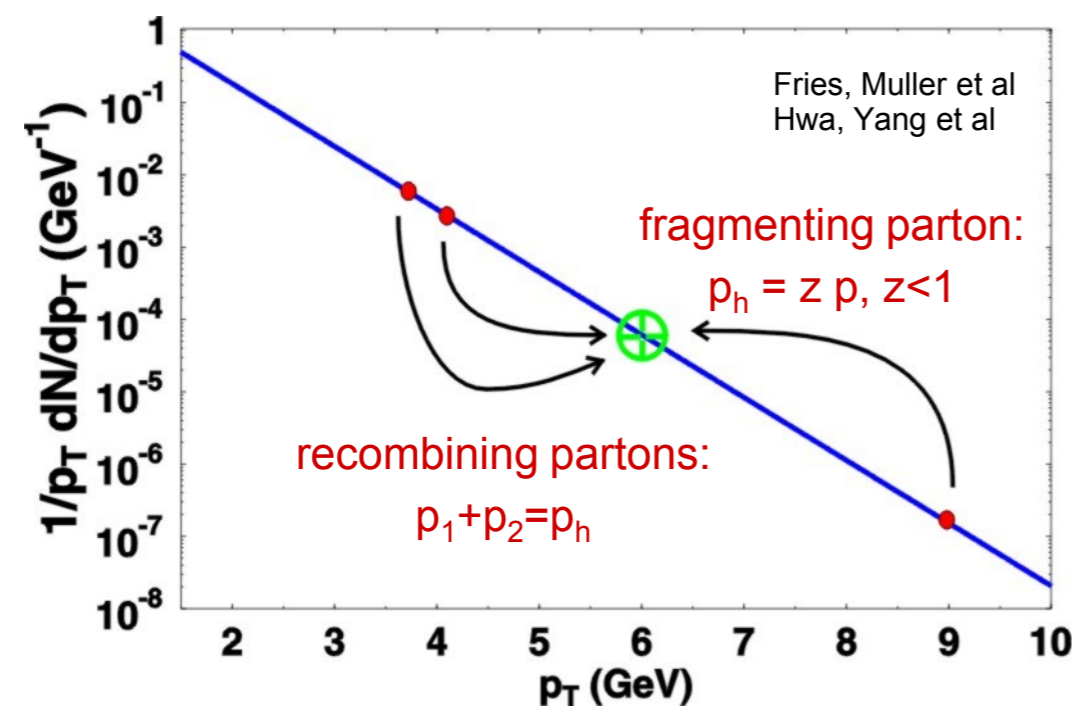
bulk, jets, medium and p_T
 arbitrary regions
 and INFORMAL Language



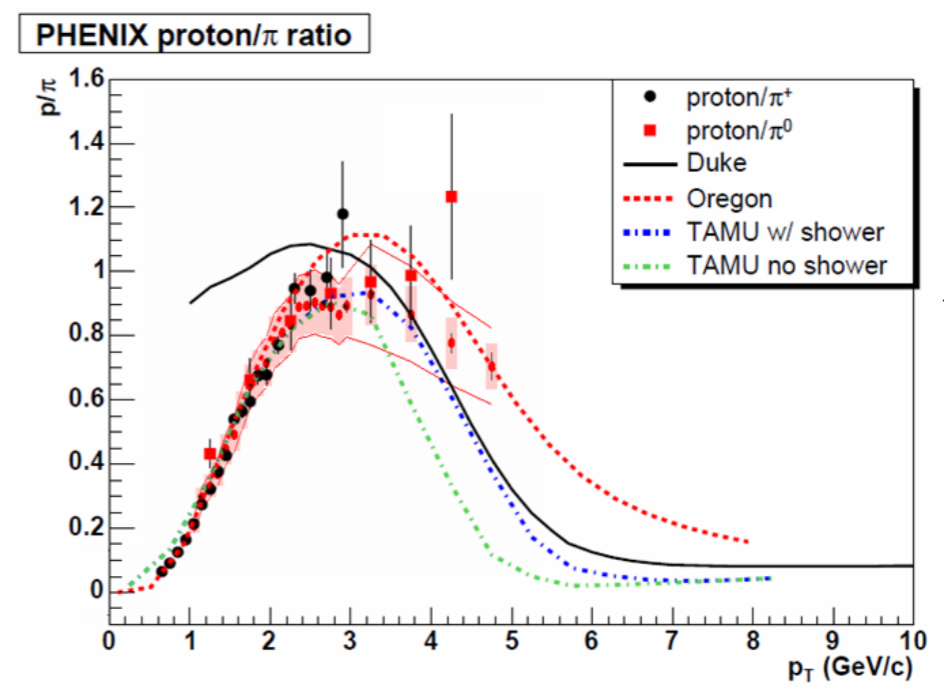
"soft" "soft+hard" "hard"
 -bulk jet-medium jet dominated
 thermal, intermediate

→

Novel effects: hadronization of a mix bulk & hard - parton coalescence

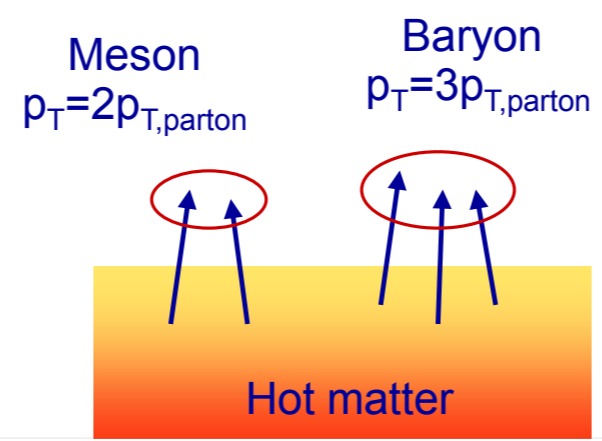


Recombination of thermal ('bulk') partons produces baryons at larger p_T



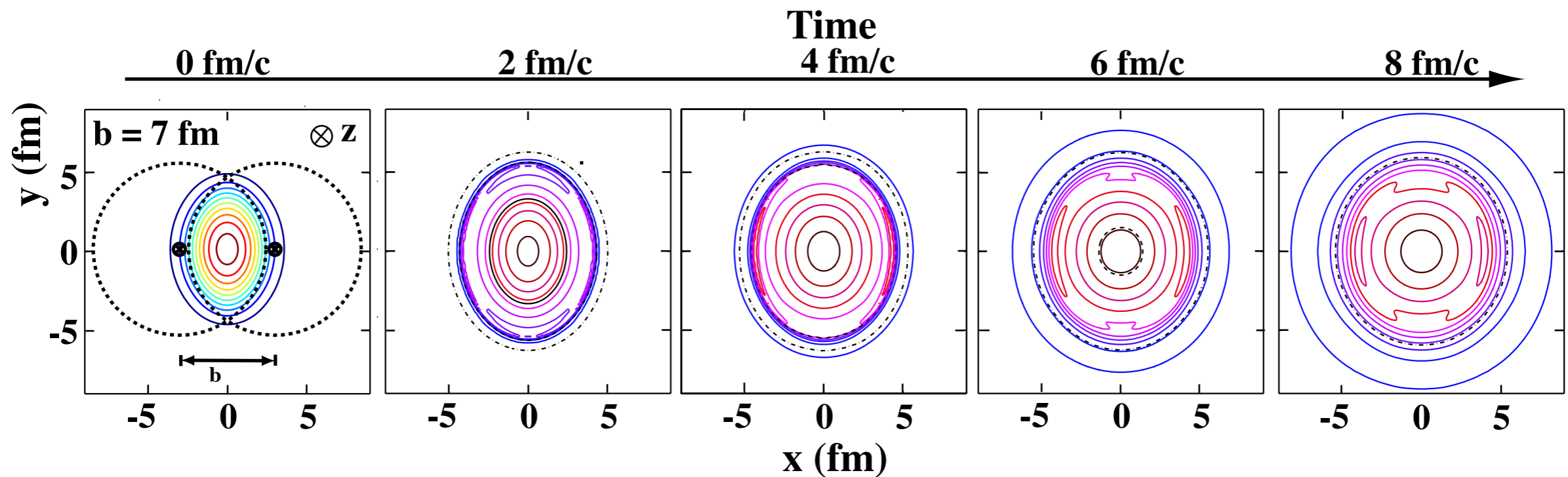
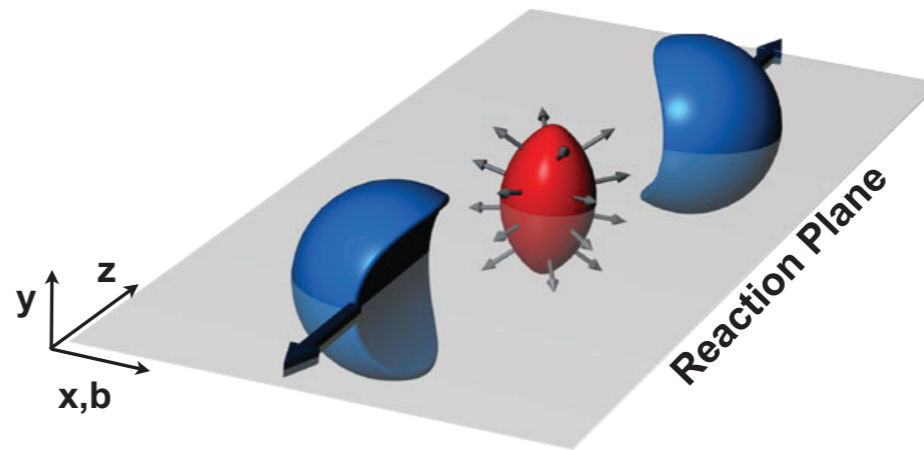
Recombination enhances baryon/meson ratio

Note also: v_2 scaling



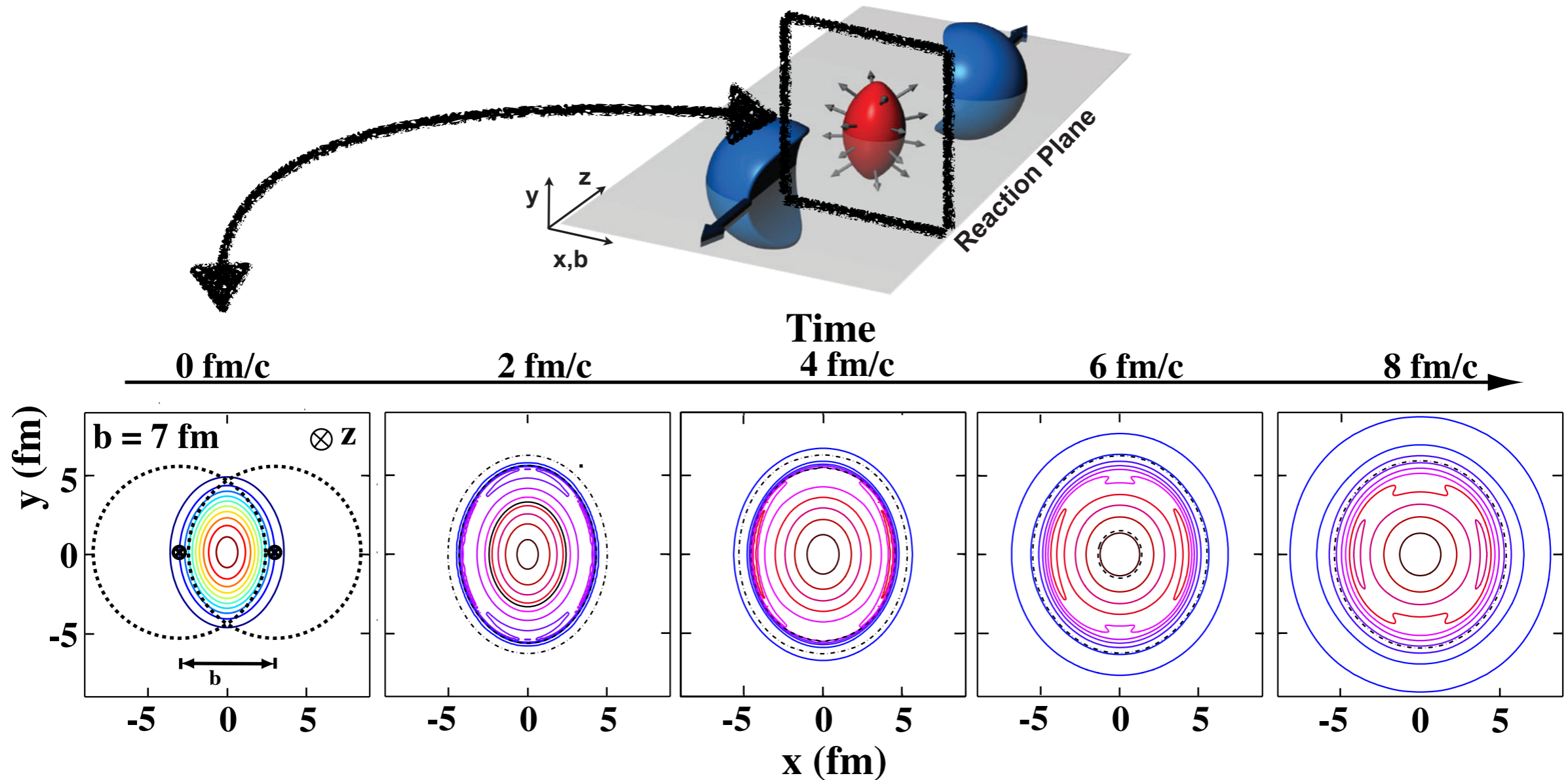
Properties of QGP
with particle
correlations

Expanding "fireball"



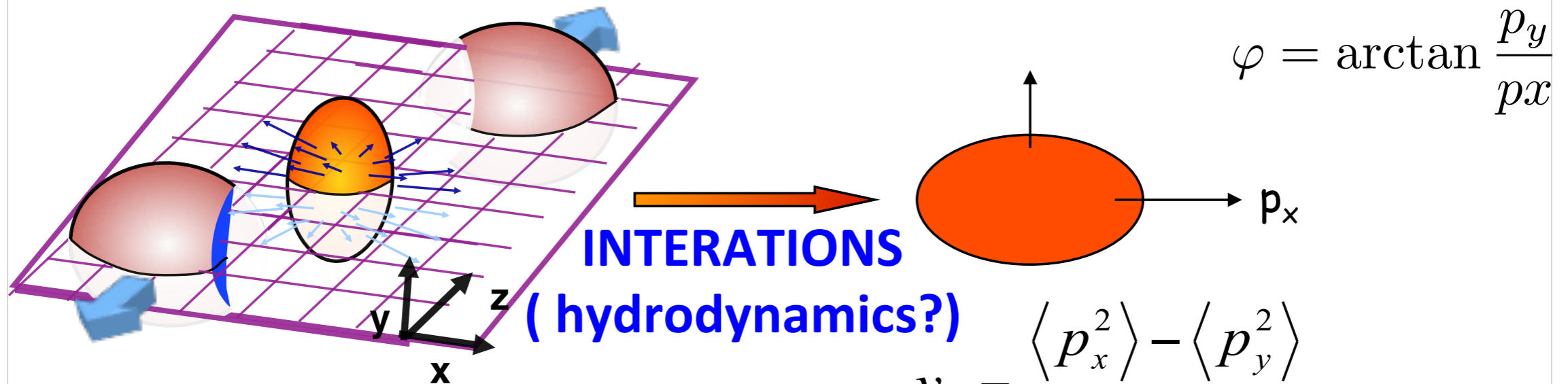
Initial transverse energy density profile and its time dependence in coordinate space for a non-central heavy-ion collision

Expanding "fireball"



Initial transverse energy density profile and its time dependence in coordinate space for a non-central heavy-ion collision

Azimuthal angular asymmetry in particle production



$$\varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$$

Initial spatial anisotropy

$$v_2 = \frac{\langle p_x^2 \rangle - \langle p_y^2 \rangle}{\langle p_x^2 \rangle + \langle p_y^2 \rangle}$$

Final momentum anisotropy

Reaction plane defined by
"soft" (low p_T) particles

$$\Delta\varphi = \varphi - \varphi^{\text{Reaction Plane}}$$

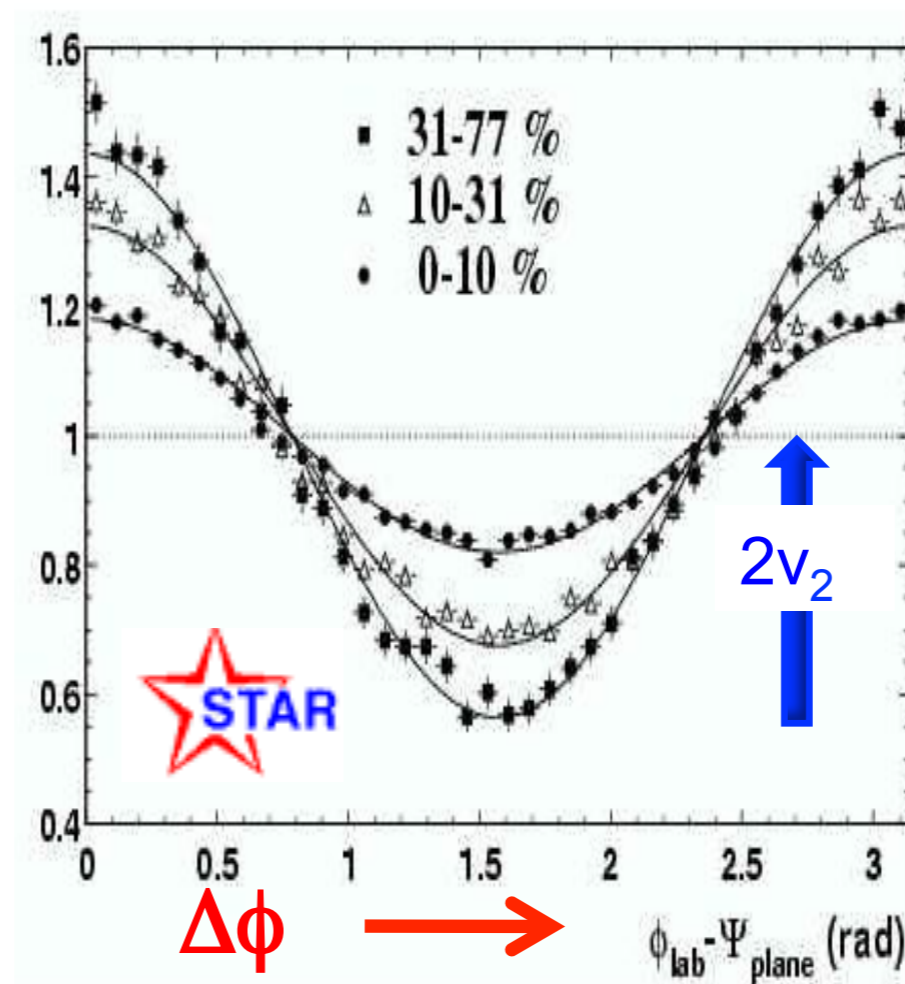
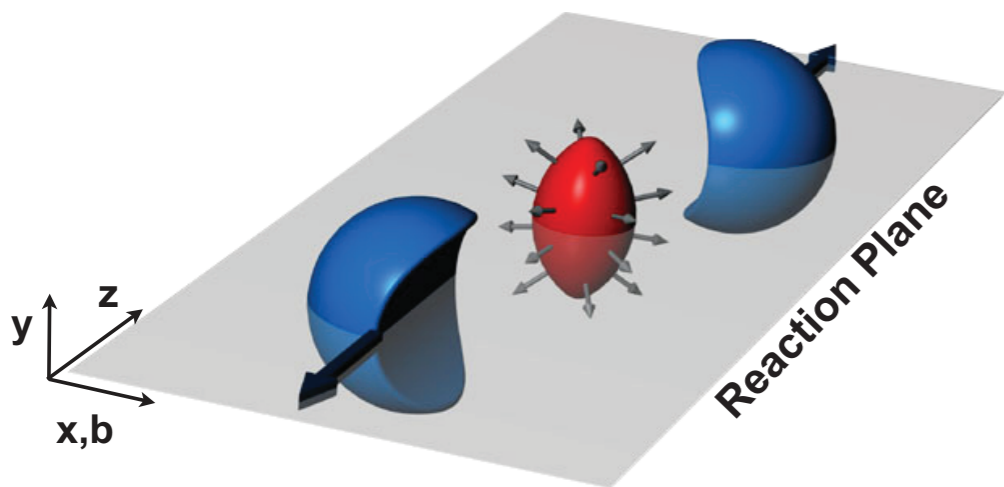
Elliptic flow

$$\frac{dN}{d\Delta\varphi} \propto 1 + 2v_2 \cos(2\Delta\varphi)$$

Experimental signature

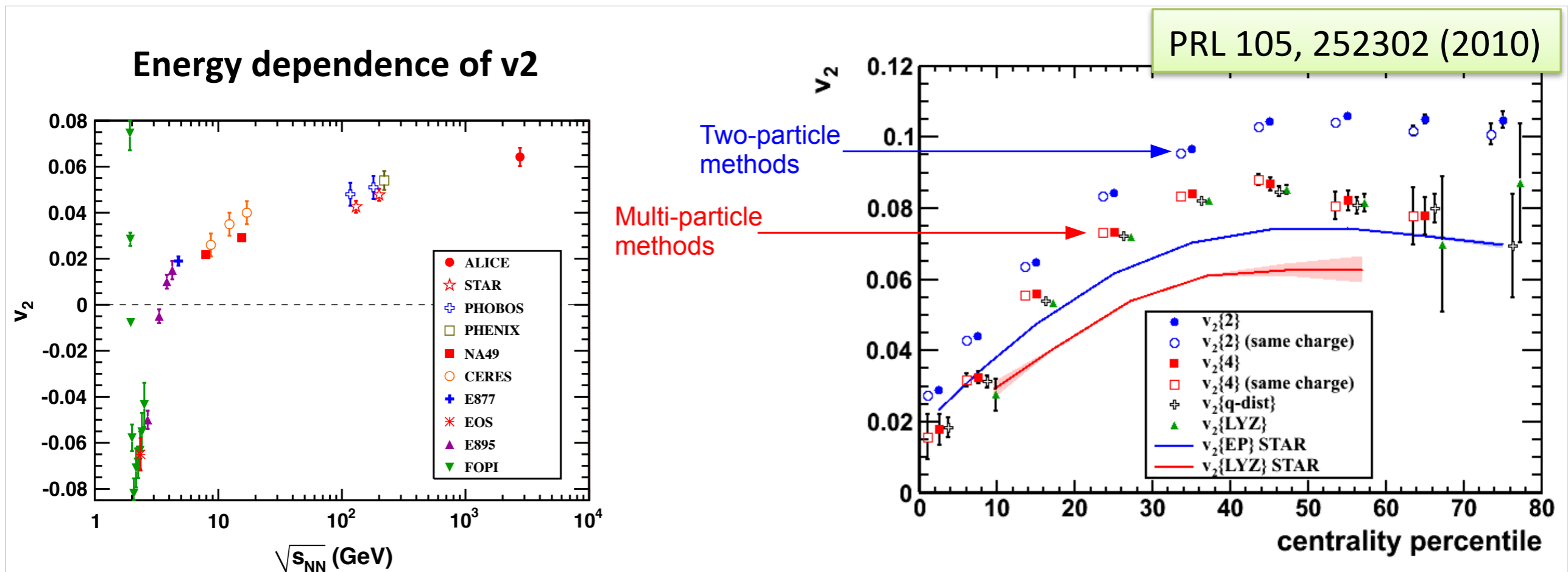
measurement: azimuthal angular distribution of particles with respect to event plane

$$\frac{dN}{d\phi} \propto 1 + 2v_2 \cos [2(\phi - \Psi_R)] + \dots$$



Sizeable effect!

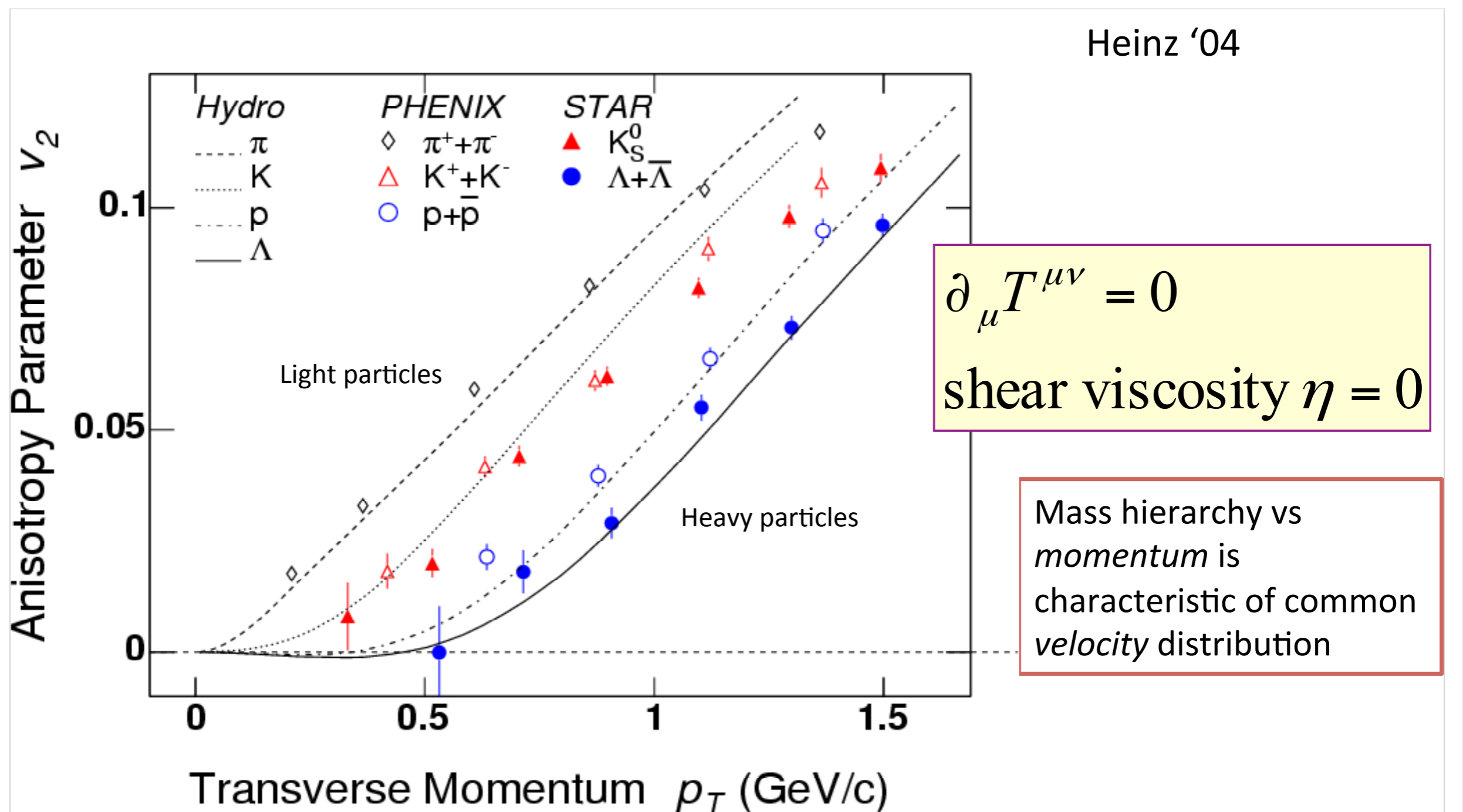
Azimuthal anisotropy



APS Viewpoint: A “Little Bang” arrives at the LHC (E. Shuryak)

- 1. Collective behavior observed in Pb-Pb collisions at LHC (integrated: $+0.3 v_2^{\text{RHIC}}$ – consequence of larger $\langle p_T \rangle$) $\rightarrow v_2(p_T)$ similar to RHIC – almost ideal fluid at LHC ? Similar observation down to 39GeV!**
- 2. New input to the energy dependence of collective flow**
- 3. Additional constraints on Eq-Of-State and transport properties**

Relativistic (ideal) hydrodynamics



Ideal hydro: qualitative agreement but missing the details

Hydrodynamics crash course

Energy-momentum conservation (local)

$$\partial_t T_{00} = -\partial_x T_{0x} - \partial_y T_{0y} + \partial_z T_{0z}$$

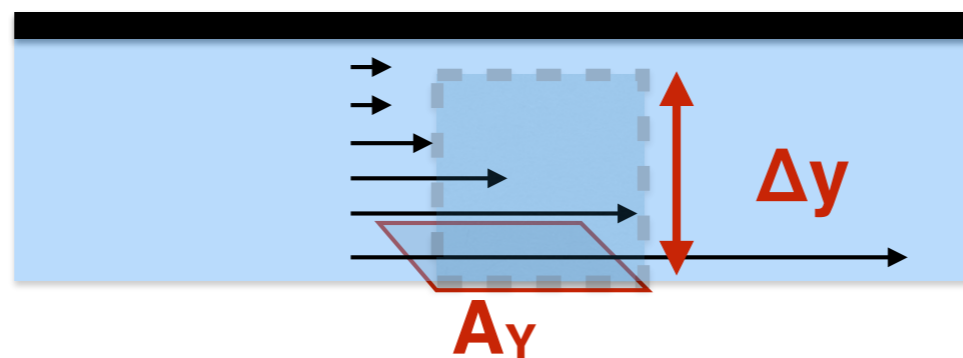
$$\partial_t T_{0x} = \partial_x T_{xx} + \partial_y T_{yx} + \partial_z T_{zx}$$

Ideal hydrodynamics: $T_{i \neq j} = 0$

...motion of viscous fluids ...

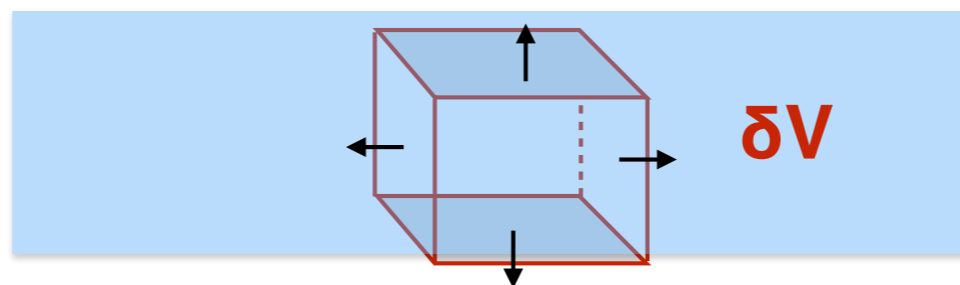
Navier-Stokes equation: $T_{ij} = P\delta_{ij} - \eta(\partial_i v_j + \partial_j v_i) - \zeta \nabla \cdot \vec{v}$

Where η is shear viscosity: friction between layers of fluid



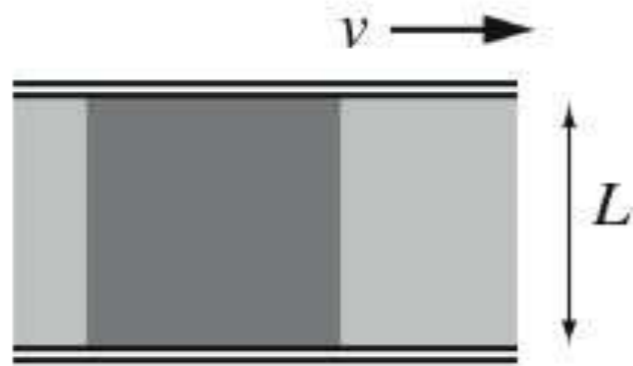
$$\frac{d}{dt} P_x = A_y \eta \partial_y v_x$$

and ζ is bulk viscosity: dissipation of divergent flow

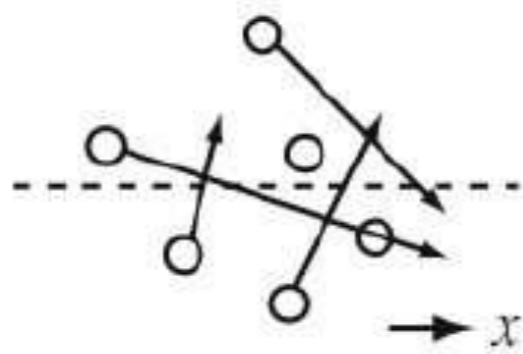
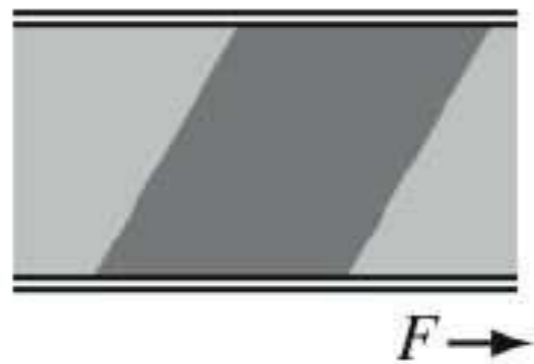
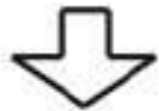


$$\delta E = -P\delta V + \zeta \nabla \cdot \vec{v} \delta V$$

Shear viscosity in fluids...



$$\frac{F}{A} = \eta \frac{v}{L}; \quad \eta \sim \rho \langle v \rangle \lambda_{mfp}$$



Properties are counter-intuitive:

Weak coupling

- small cross section, long mean free path
⇒ large viscosity

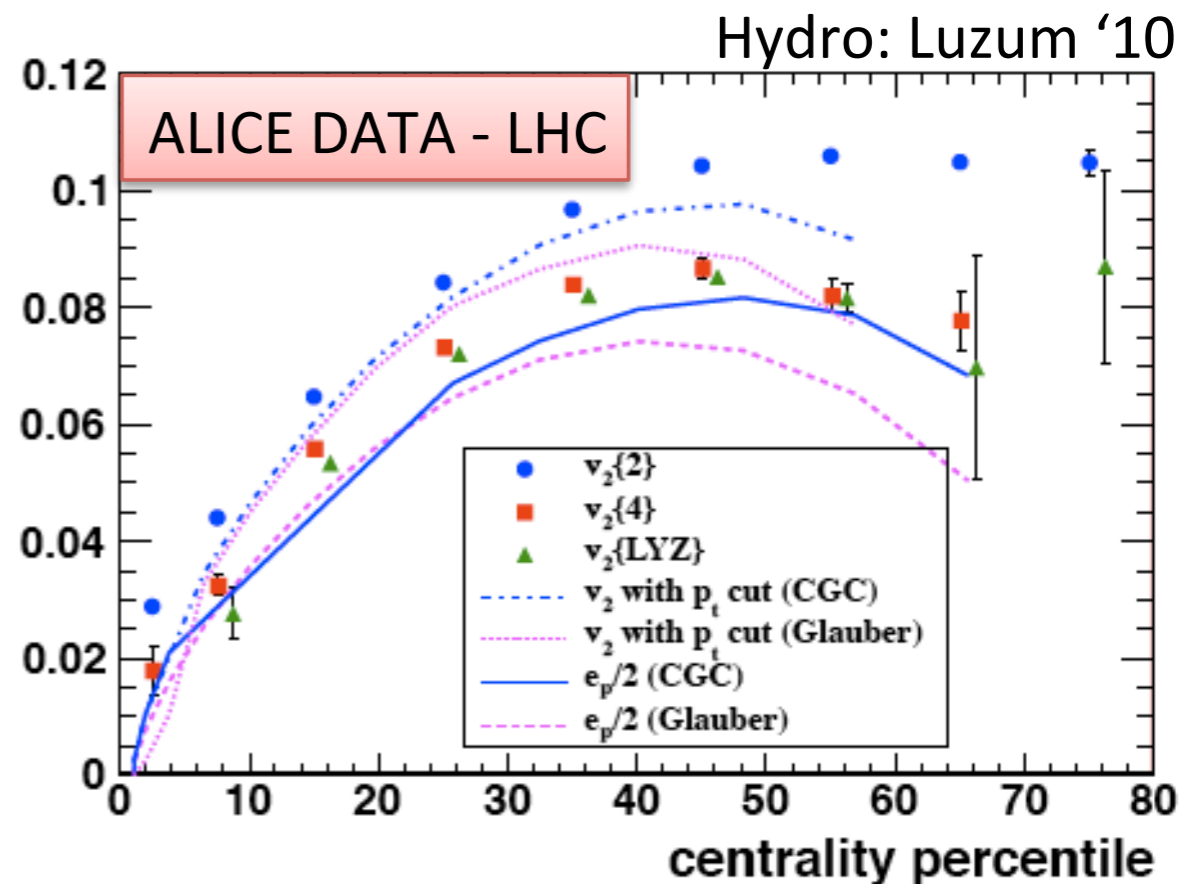
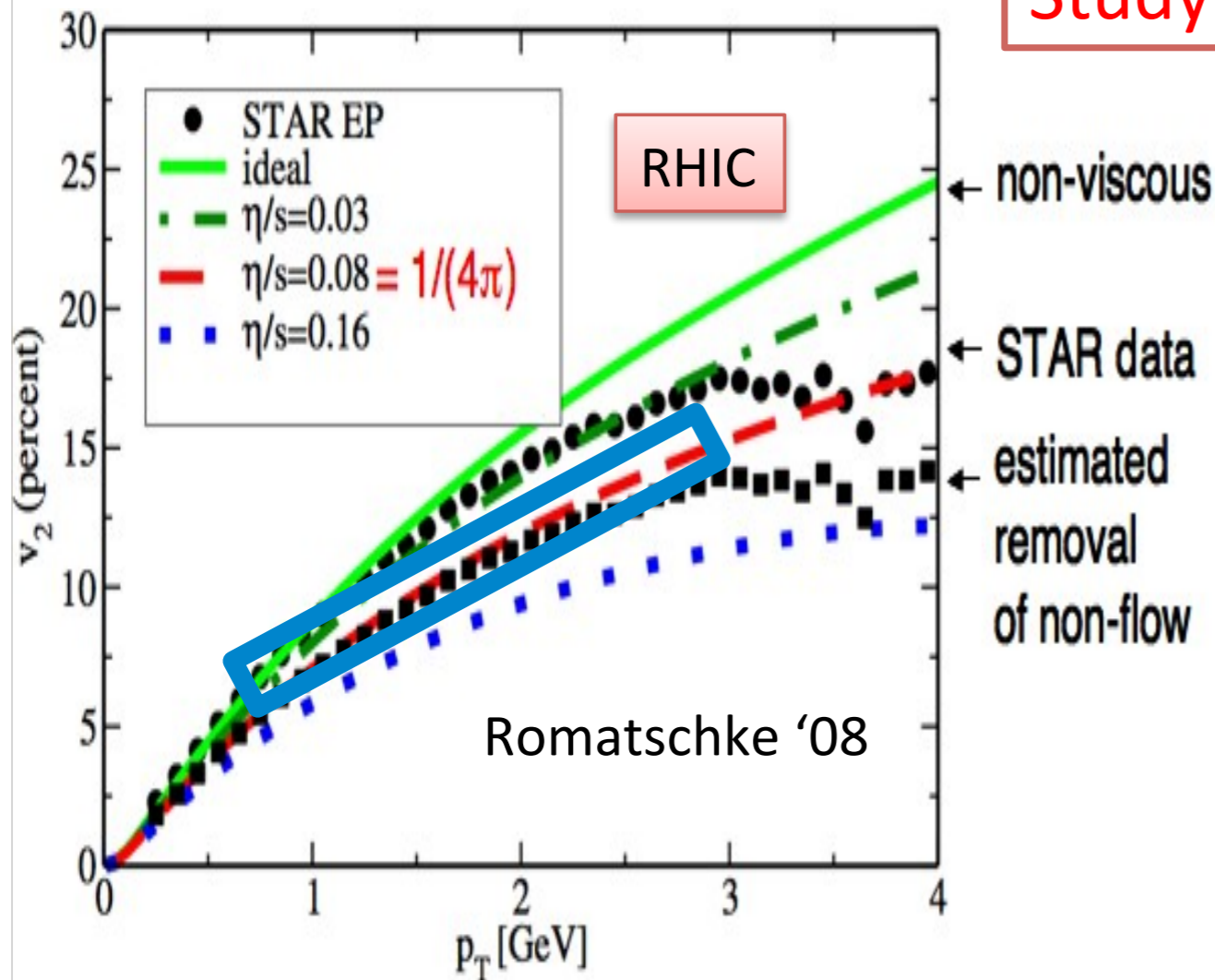
Strong coupling

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⇒ small viscosity

$\eta \rightarrow 0$: strongly coupled (perfect) fluid
 $\eta \rightarrow \infty$: weakly coupled (ideal) gas

QGP liquid- how perfect is perfect?

Study elliptic flow of matter



Shear viscosity – lower limit:

KSS (string theory); Gyulassy-Danielewicz (quantum mechanics + ballistic theory)

$$\frac{\eta}{s} > \frac{1}{4\pi}$$

rather recent: in principle can go to zero

0.08 : $\lambda_{therm} \lambda_{mfp}$ (Danielewicz and Gyulassy)

Hot, deconfined QCD matter flows as an almost perfect fluid

Comparison QGP to other fluids near T_c

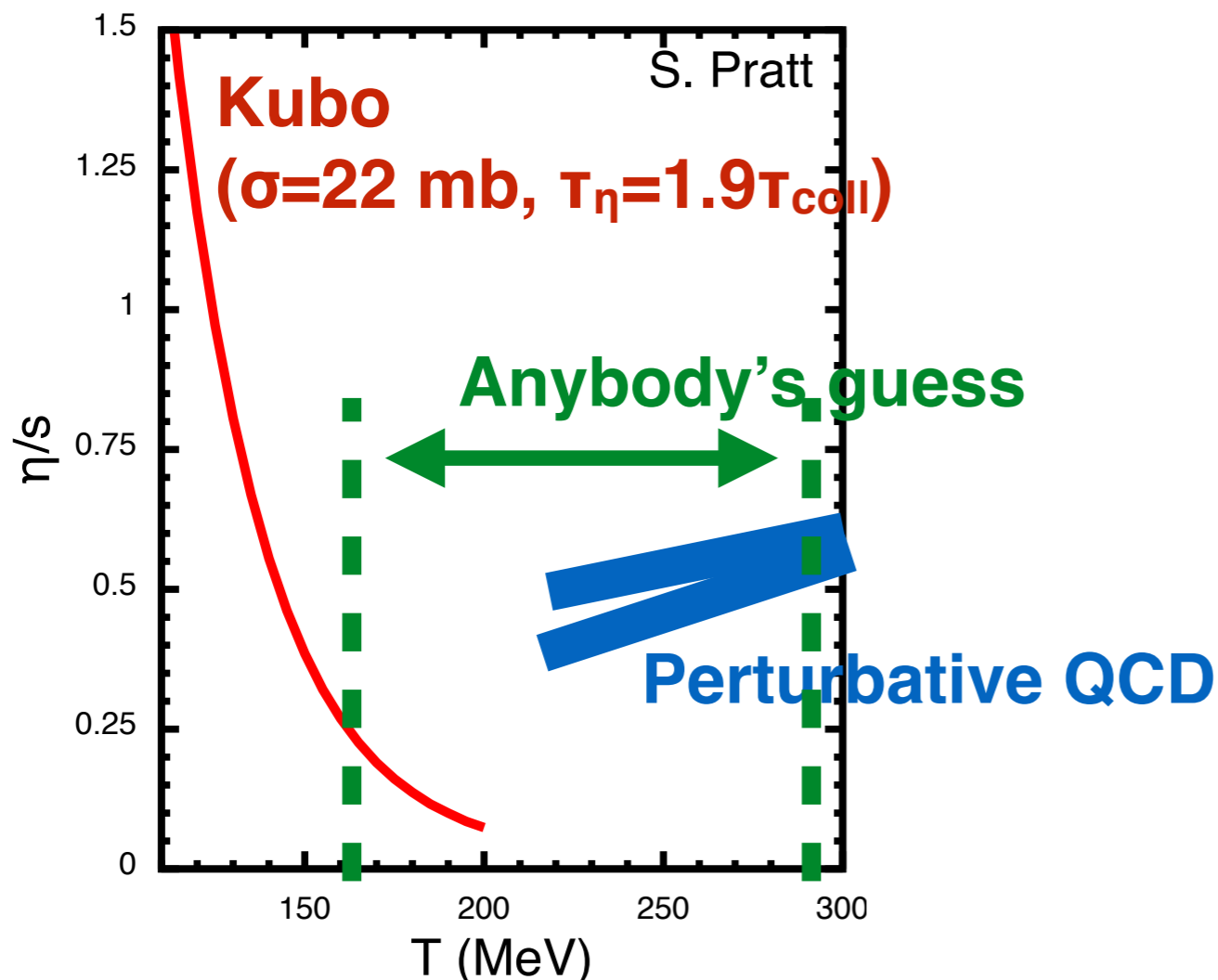
Green-Cubo relations:

transport coefficients in terms of integrals of time correlation functions

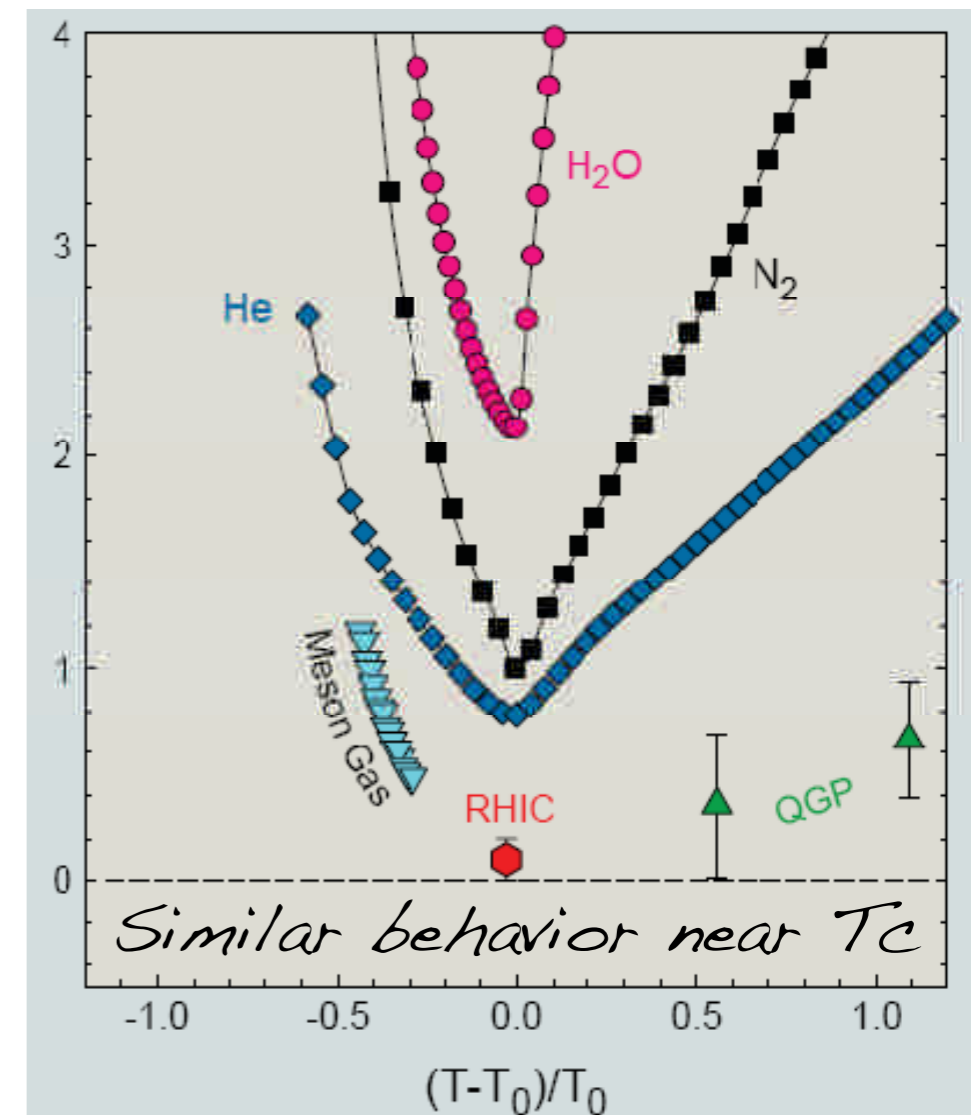
- correlations of particles \times relaxation time

$$\eta = \frac{\tau_\eta}{T} \int d^3r \langle T_{xy}(0,0) T_{xy}(\vec{r}, t=0) \rangle$$

$$= \frac{\tau_\eta}{T} \sum_\alpha (2S_\alpha + 1) \int \frac{d^3p}{(2\pi)^3} e^{-E/T} \frac{p_x^2 p_y^2}{E^2}$$

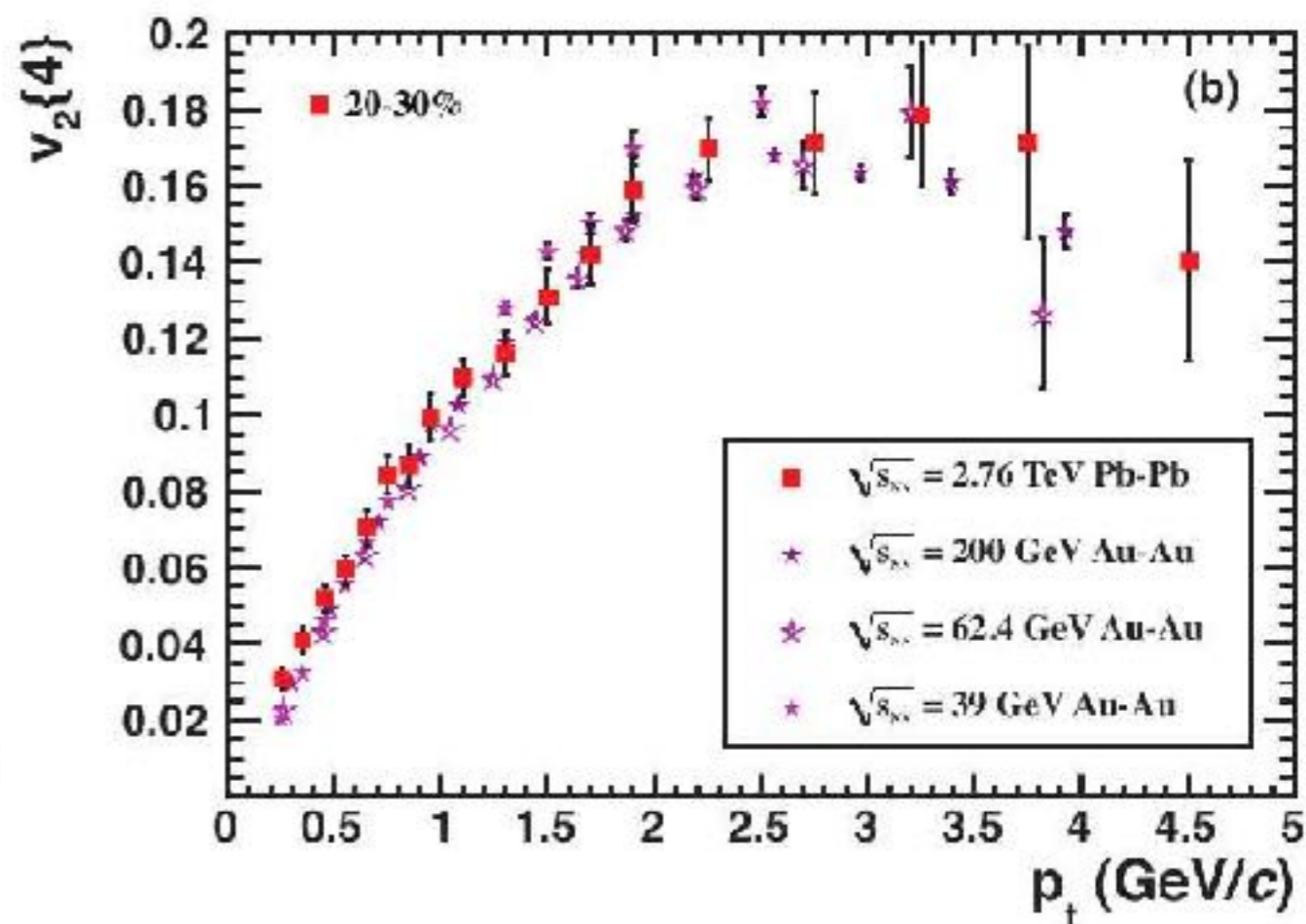
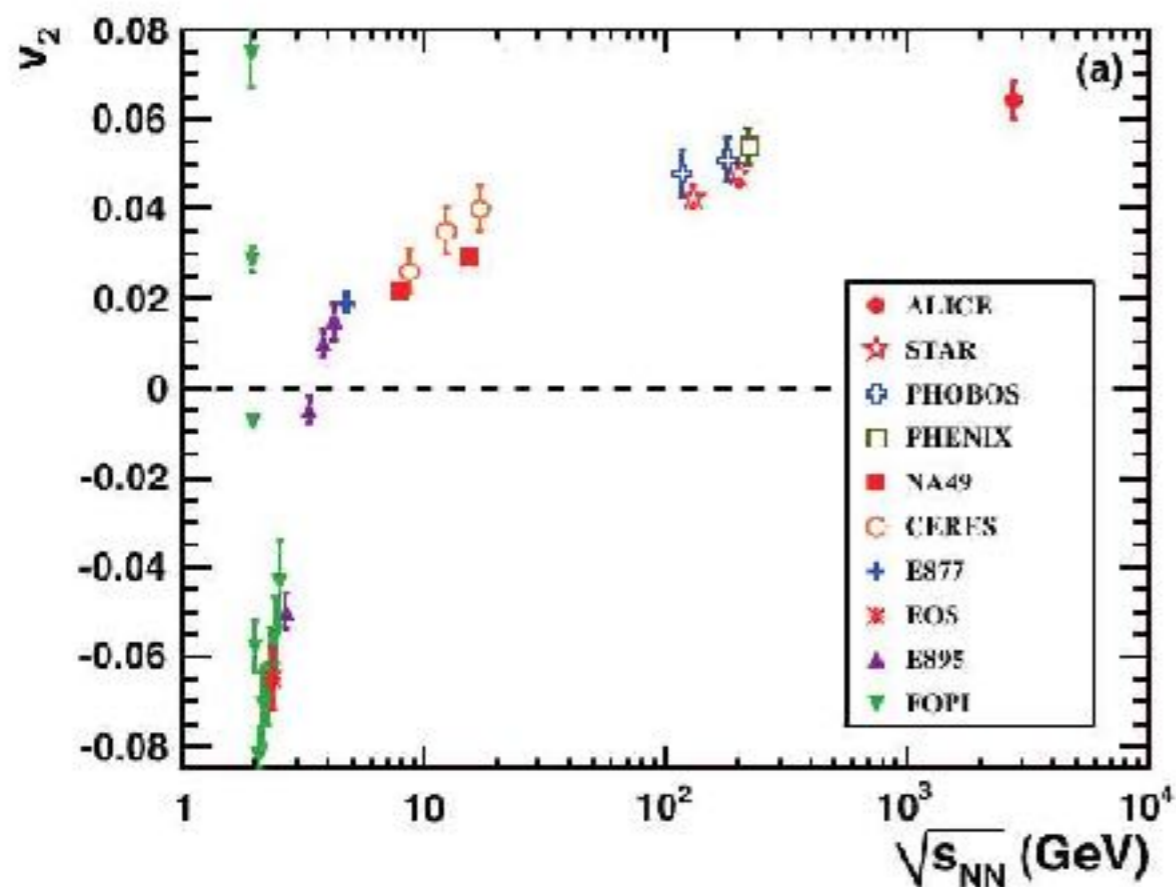


η/s



Elliptic Flow

- collision energy dependence

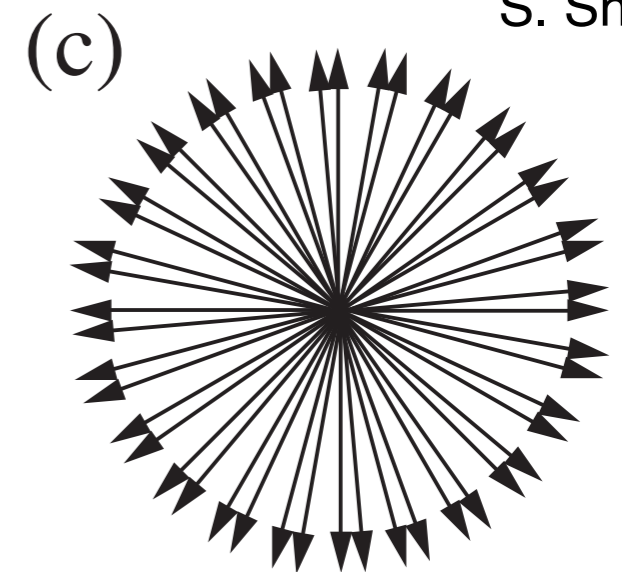
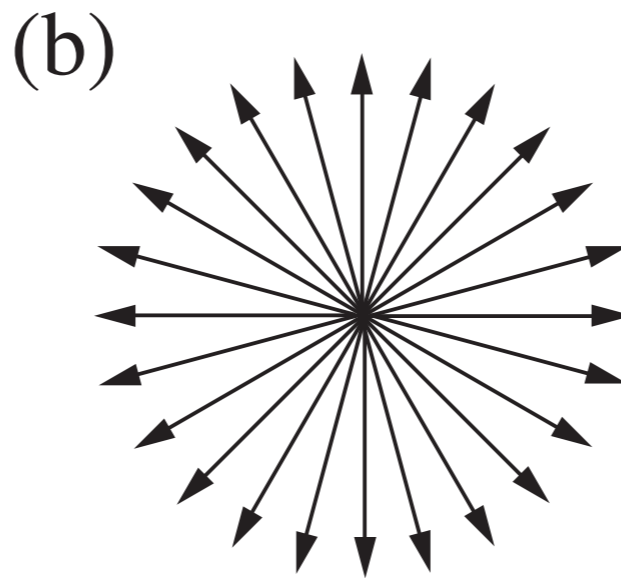
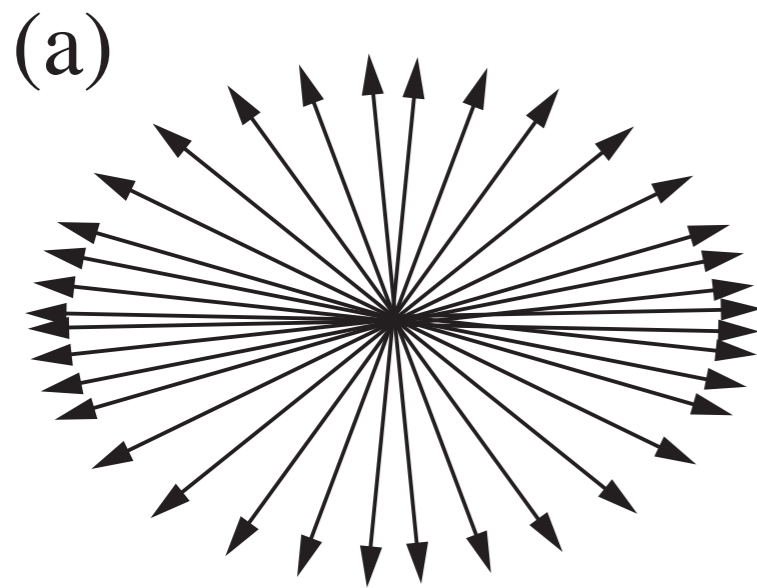
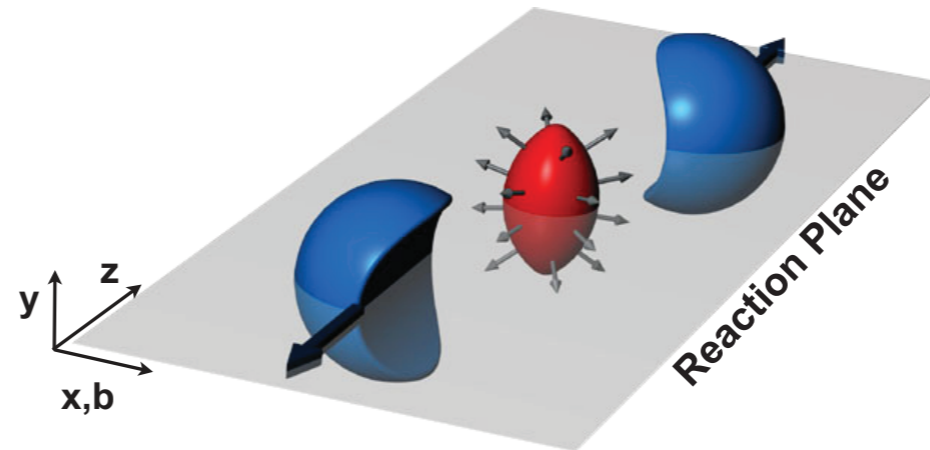


$$v_2 = \frac{\int dp_t \frac{dN}{dp_t} v_2(p_t)}{\int dp_t \frac{dN}{dp_t}}$$

*Improved (multiparticle) $v_2\{4\}$:
 very weak energy dependence of $v_2(p_t)$ -
 from 2.76 TeV down to 39 GeV (!)
 Same phase for different initial
 collision energies !?*

Understanding correlations & v_2

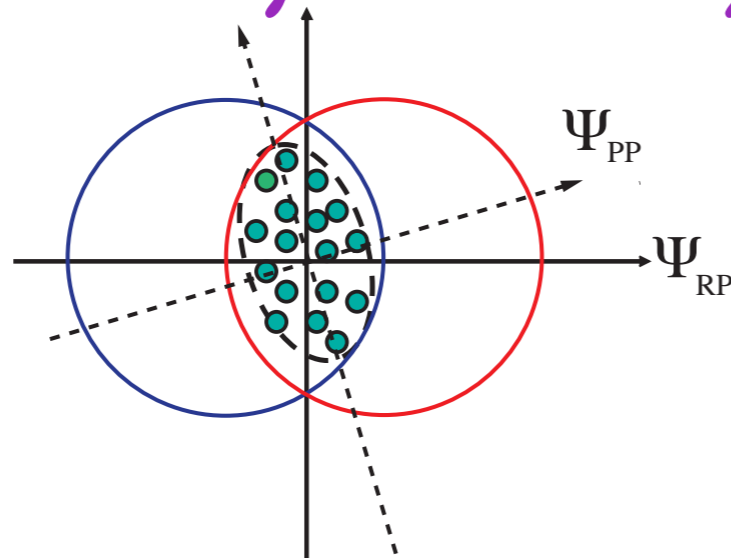
- the so-called non-flow



S. Snellings

Figure 9. Examples of particle distributions in the transverse plane, where for a) $v_2 > 0$, $v_2\{2\} > 0$, b) $v_2 = 0$, $v_2\{2\} = 0$, and c) $v_2 = 0$, $v_2\{2\} > 0$.

Reaction plane (RP) Participants plane (PP)

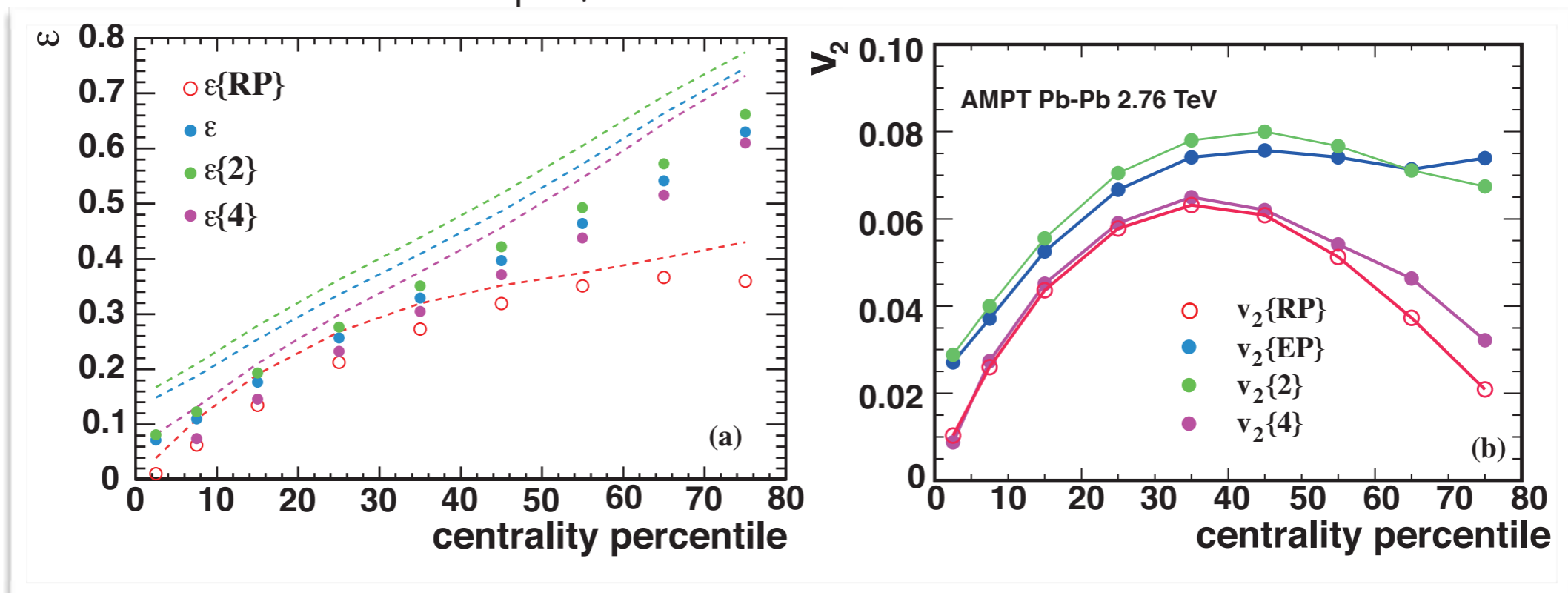


Genuine 2,4-particle correlations

$$c_2\{2\} \equiv \langle \langle e^{i2(\varphi_1 - \varphi_2)} \rangle \rangle = \langle v_2^2 + \delta_2 \rangle.$$

$$\begin{aligned} c_2\{4\} &\equiv \langle \langle e^{i2(\varphi_1 + \varphi_2 - \varphi_3 - \varphi_4)} \rangle \rangle - 2 \langle \langle e^{i2(\varphi_1 - \varphi_2)} \rangle \rangle^2, \\ &= \langle v_2^4 + \delta_4 + 4v_2^2\delta_2 + 2\delta_2^2 \rangle - 2 \langle v_2^2 + \delta_2 \rangle^2, \\ &= \langle -v_2^4 + \delta_4 \rangle. \end{aligned}$$

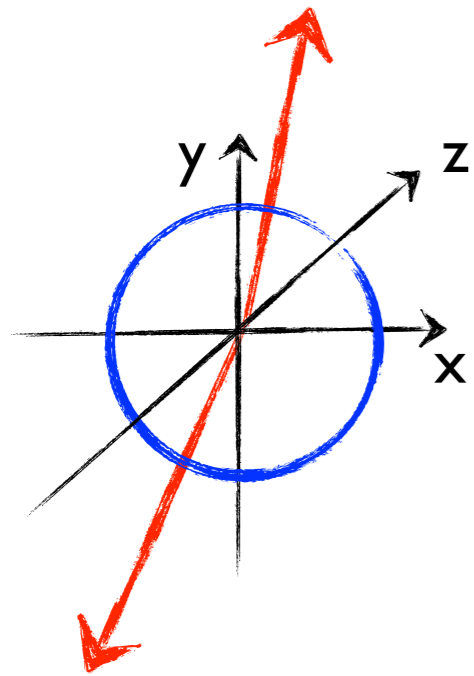
$\delta_2 \propto 1/M_c$ and $\delta_4 \propto 1/M_c^3$



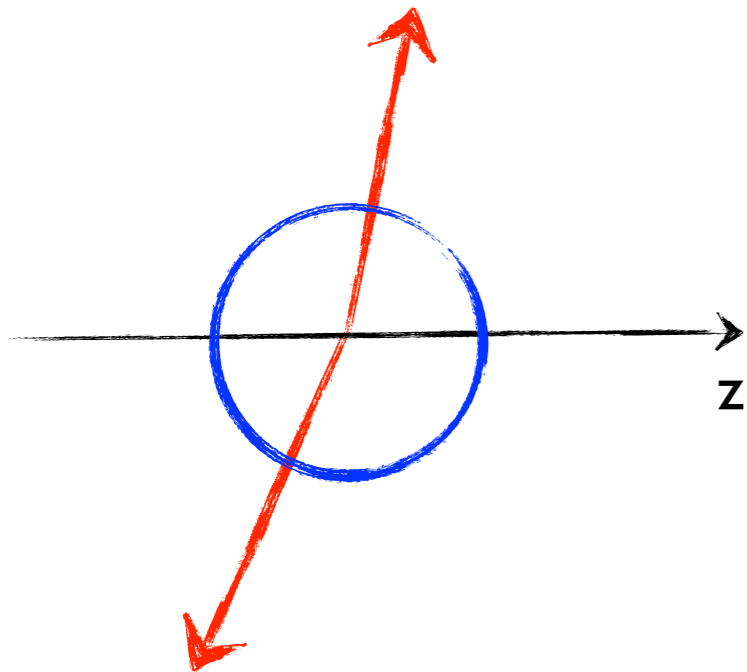
[AMPT] calculation using Glauber initial conditions

v_2 : RP, EP $v_2\{2\}$ $v_2\{4\}$

Two particle correlations



$\Delta\varphi$ - azimuthal angle difference
angle in the transverse plane



$\Delta\eta$ - longitudinal - pseudo-rapidity
distance

Sensitivity of particle correlations to the underlying/initial conditions

Two-particle correlations

- conditional [per-trigger] yields

$$\frac{1}{N_{trig}} \frac{dN_{assoc}}{d\Delta\varphi} \quad \text{and} \quad \frac{1}{N_{trig}} \frac{d^2 N_{assoc}}{d\Delta\varphi d\Delta\eta}$$

At Low- p_T :

Ridge

Hydrodynamics, flow

At High- p_T :

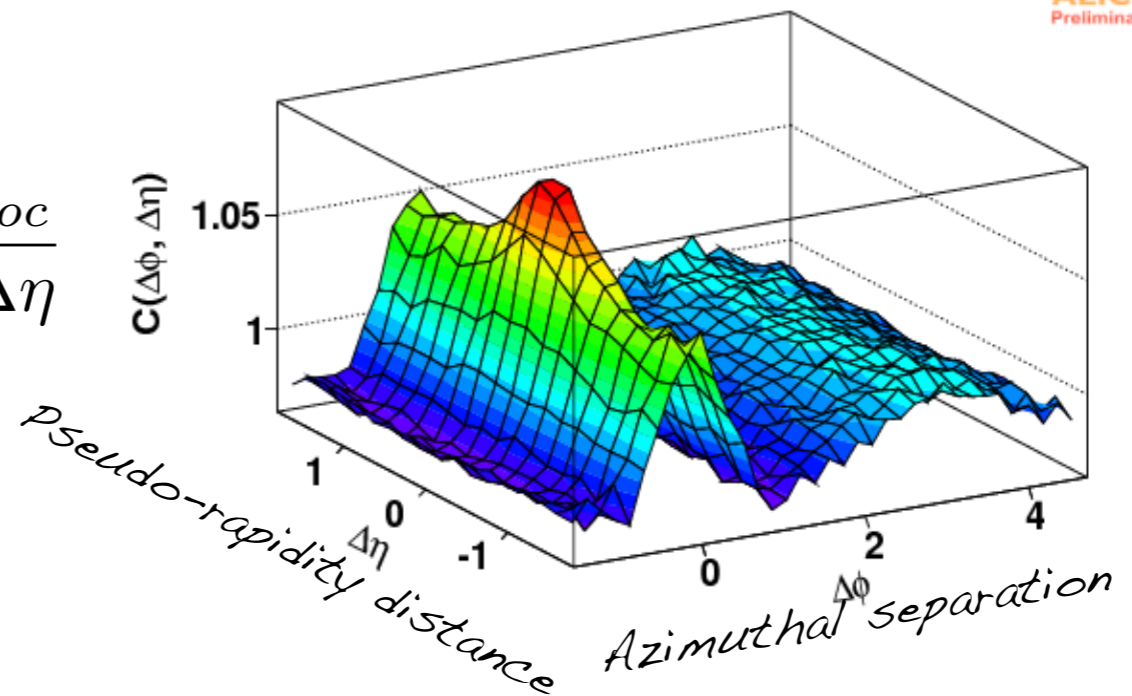
Discussed later...

**Quenching/suppression,
broadening**

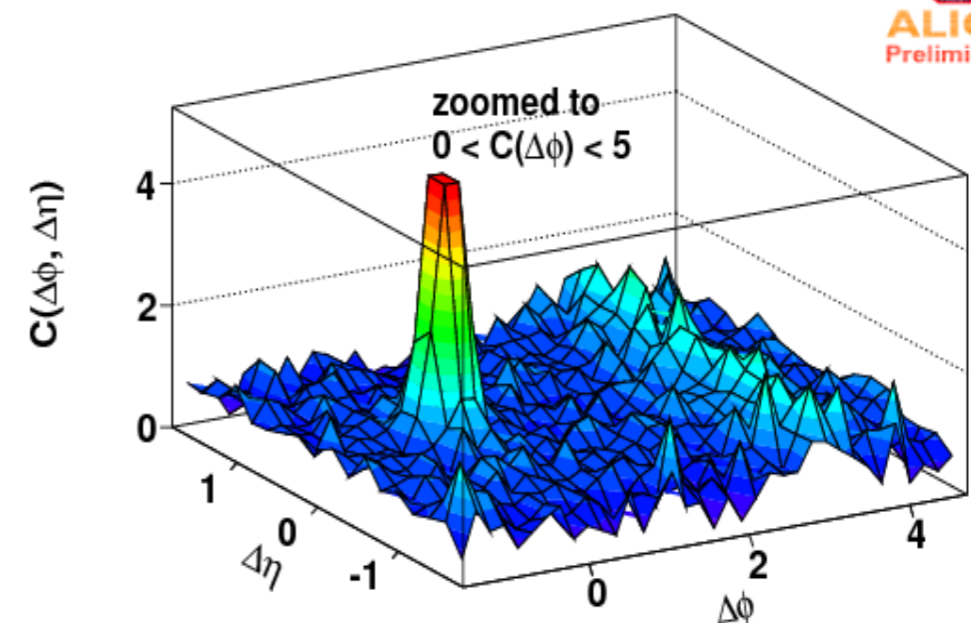
I_{CP} : Yields in central v.s. peripheral collisions

I_{AA} : Yields in A-A compared to p-p

p_T^t 3-4, p_T^a 2-2.5, 0-10%

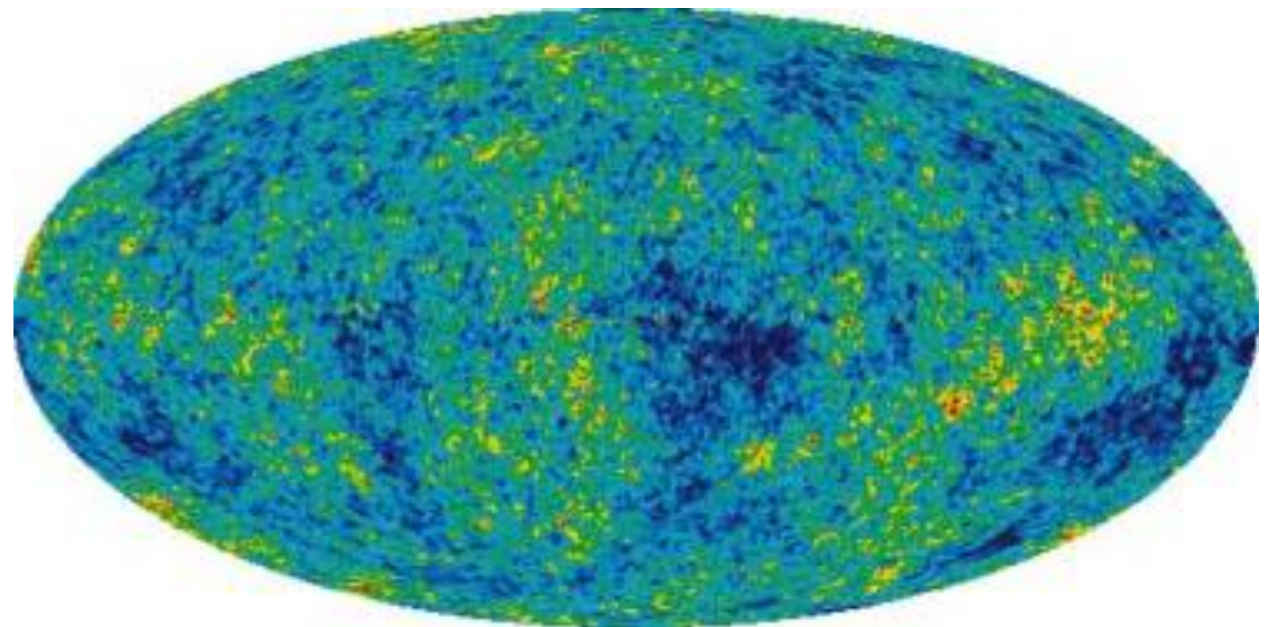
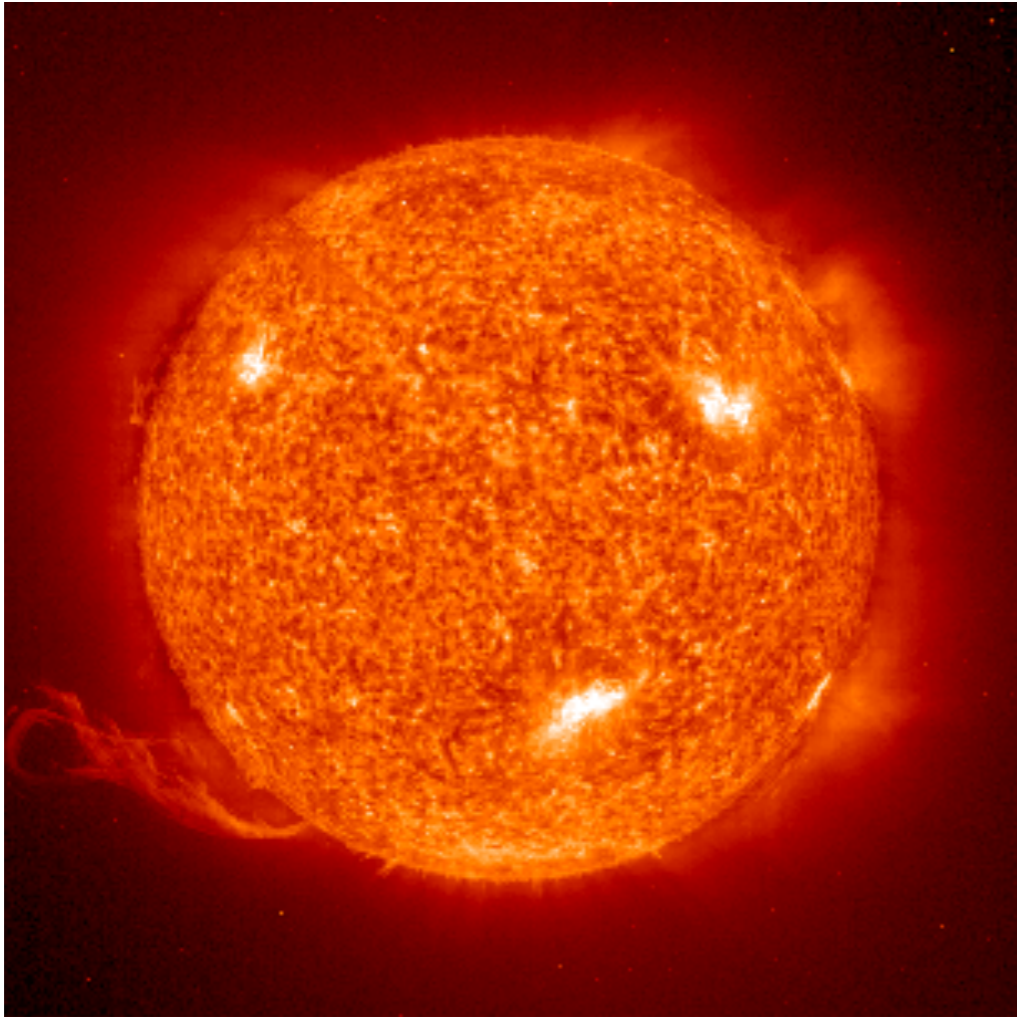


p_T^t 8-15, p_T^a 6-8, 0-20%



"Beyond" v_2

higher moments \rightarrow fluctuations / hotspots

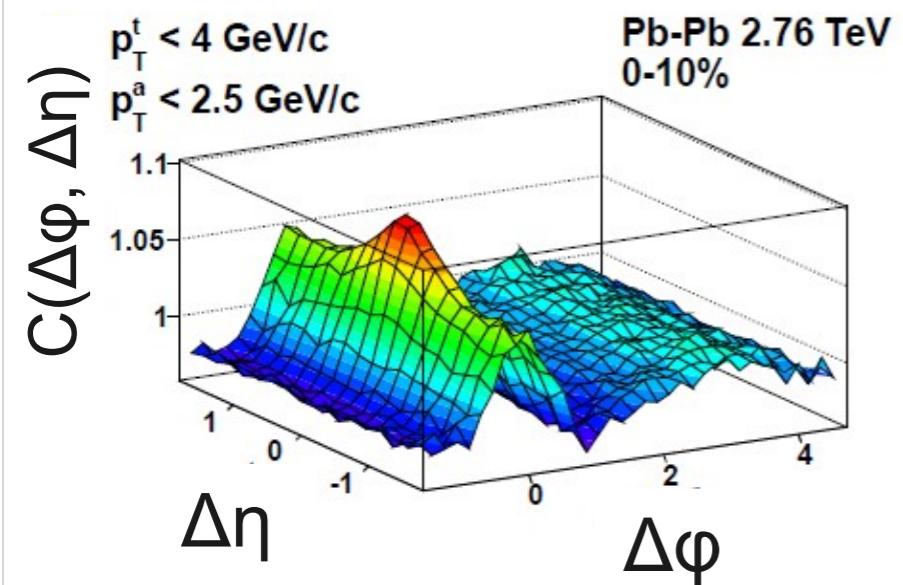


Single event!

$$\frac{dN}{d\varphi} \sim 1 + 2v_2 \cos(2\Delta\varphi) + \dots$$

Non-zero!

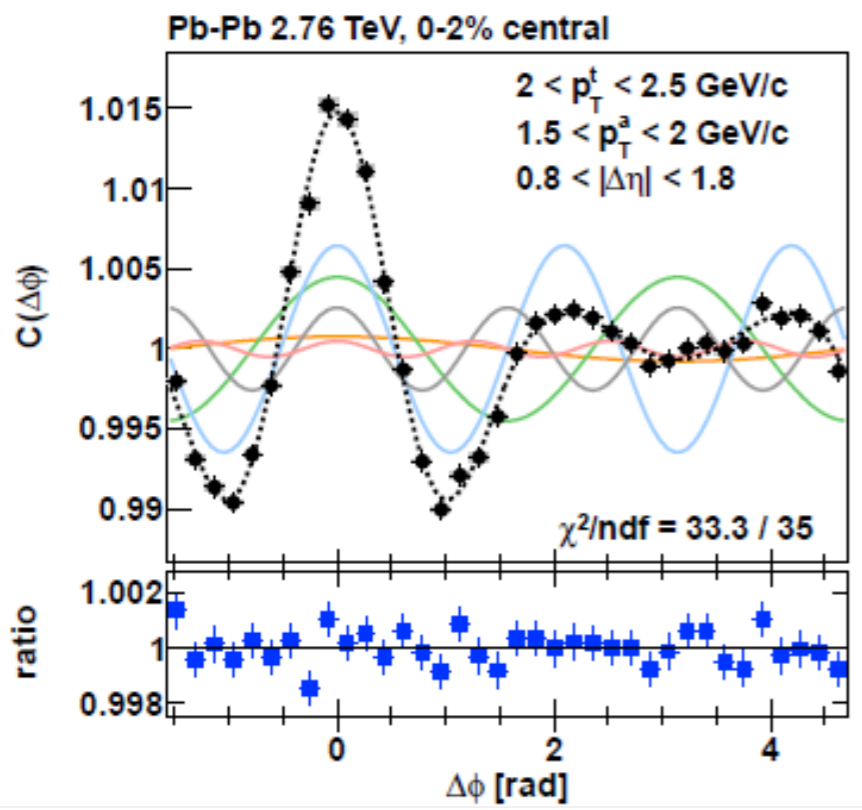
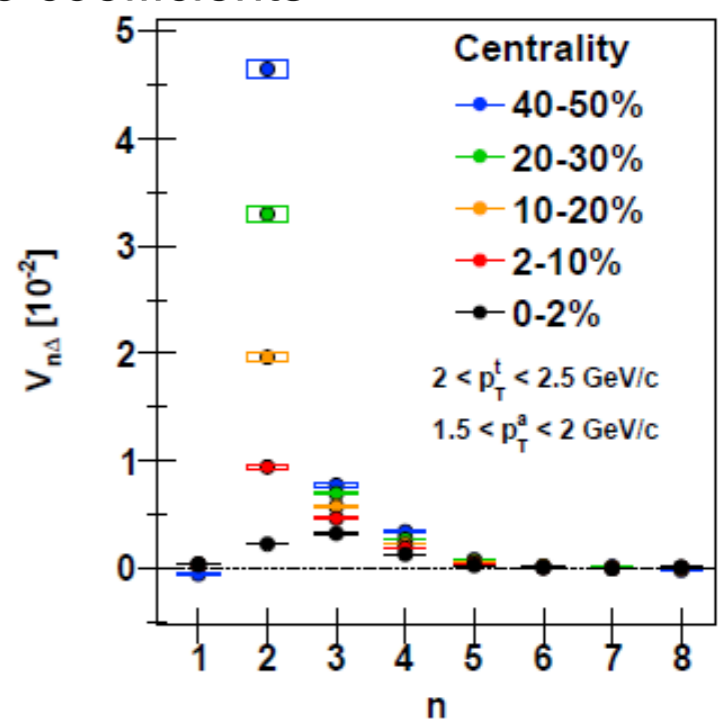
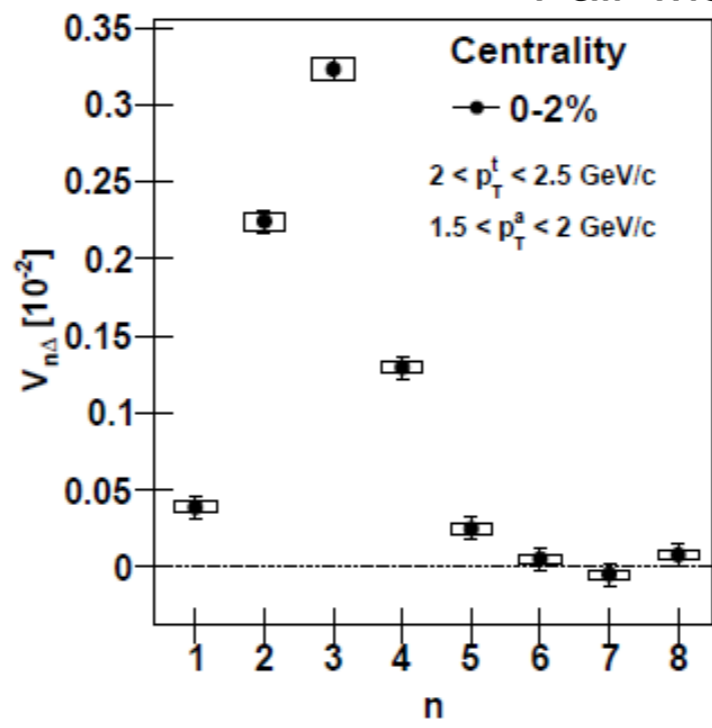
Two-particle correlations - Fourier decomposition



Integration of the correlation function in $0.8 < |\Delta\eta| < 1.8$ (long) and Fourier decomposition
 Collective flow: the coefficients factorize $V_{n\Delta} = v_n(p_T^T)v_n(p_T^A)$

$$C(\Delta\phi) = \frac{1}{\Delta\eta_{\max} - \Delta\eta_{\min}} \int_{\Delta\eta_{\min}}^{\Delta\eta_{\max}} C(\Delta\eta, \Delta\phi) \sim 1 + 2 \sum_{n=1} V_{n\Delta} \cos(n\Delta\phi)$$

Pair-wise coefficients



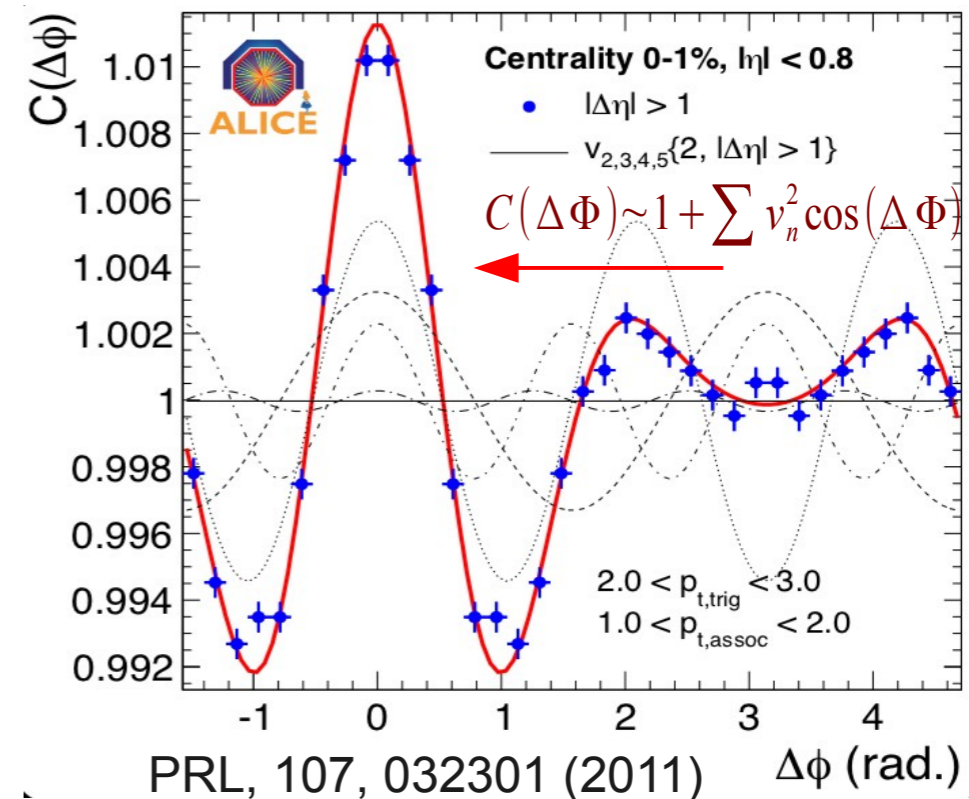
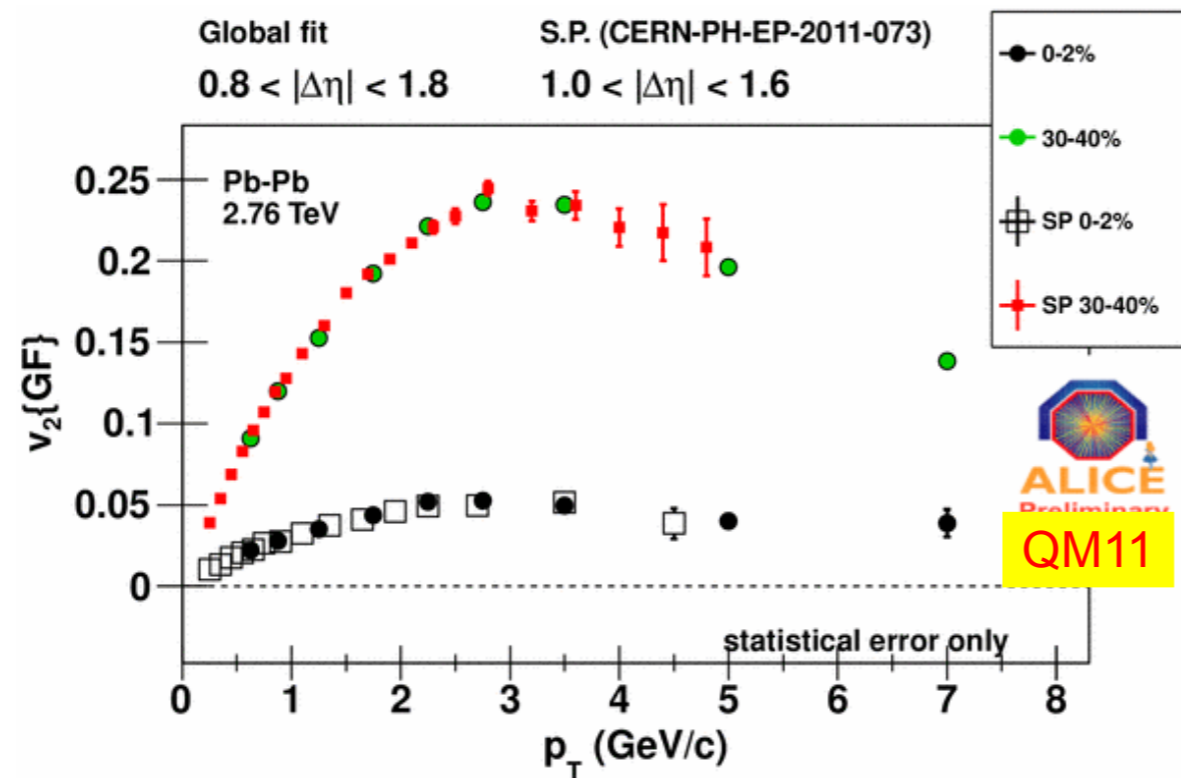
Few components describe the low-p_T correlations
 ⇔ Strong near side ridge and double-peak on the away
 ⇔ Also recoil jet up to $p_T^{\text{trig}} > 8$ & $p_T^{\text{assoc}} 6-8$ in central

Correlations & hydrodynamics...

Long range correlations – collective flow: the coefficients must factorize such that:

$$V_{n\Delta} = \langle \cos \left[n \left(\phi_{trig} - \phi_{assoc} \right) \right] \rangle = \langle \cos \left[n \left(\phi_{trig} - \Psi_n \right) \right] \rangle \langle \cos \left[n \left(\phi_{assoc} - \Psi_n \right) \right] \rangle = v_n \left(p_t^{trig} \right) \cdot v_n \left(p_t^{assoc} \right)$$

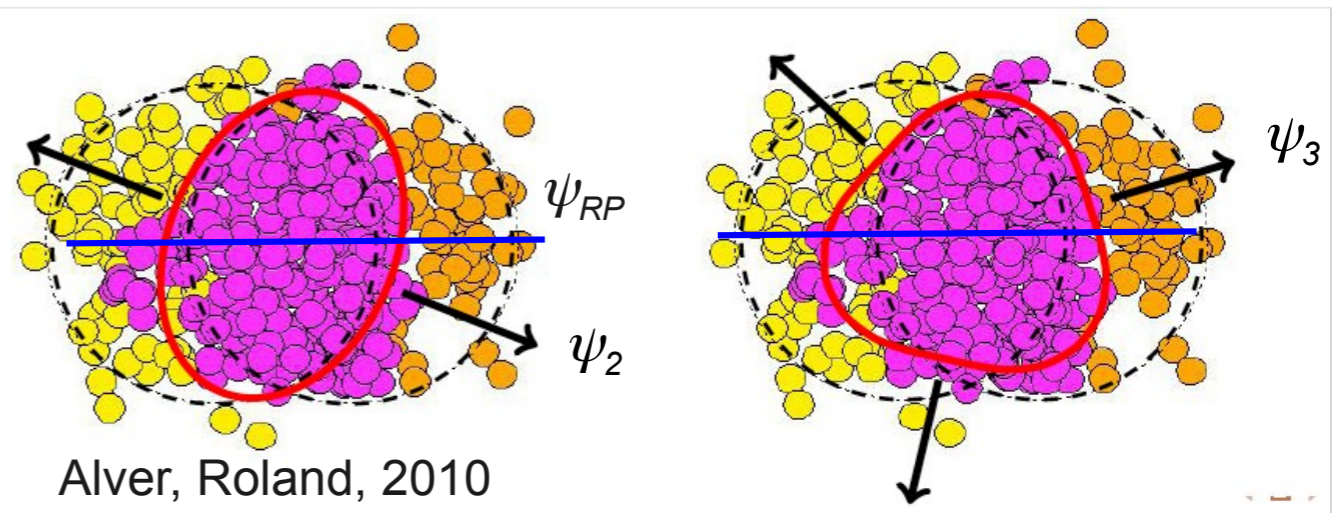
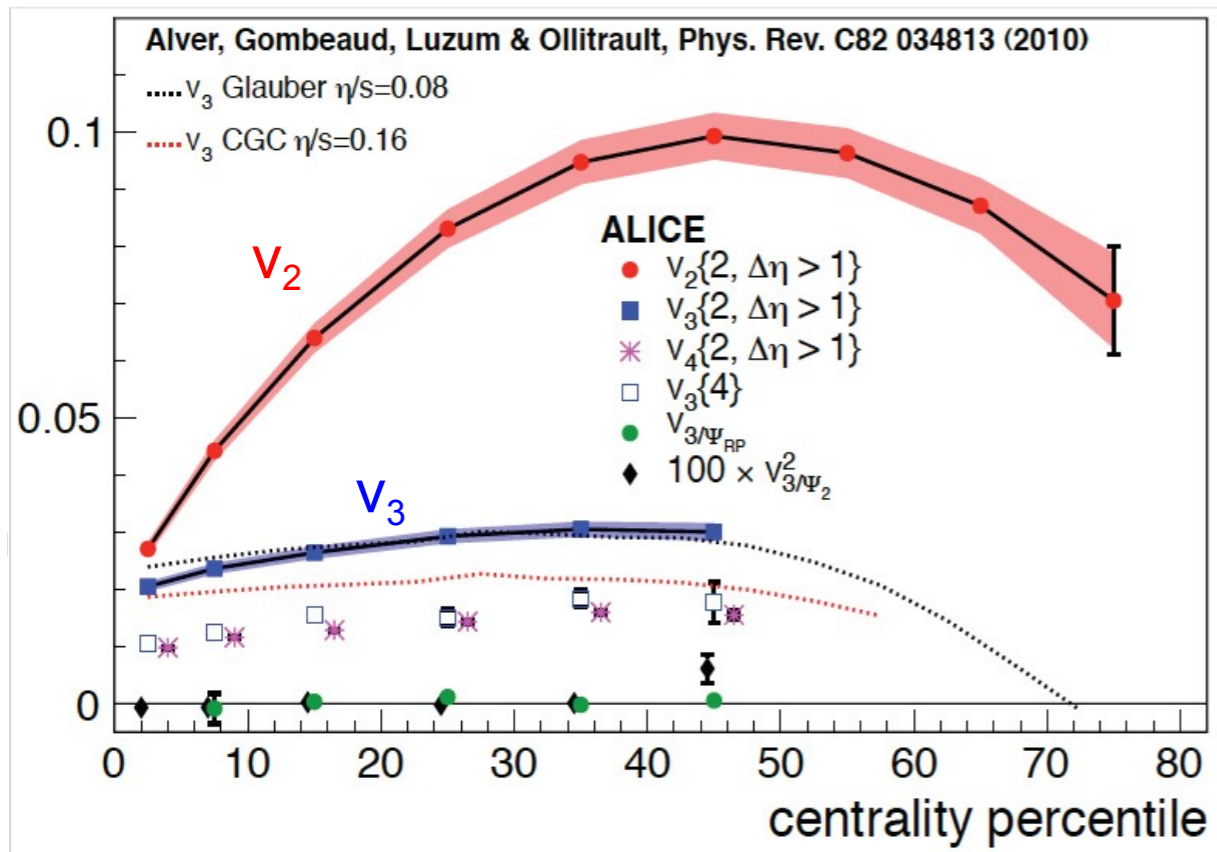
arXiv:1109.2501



Global fits show:

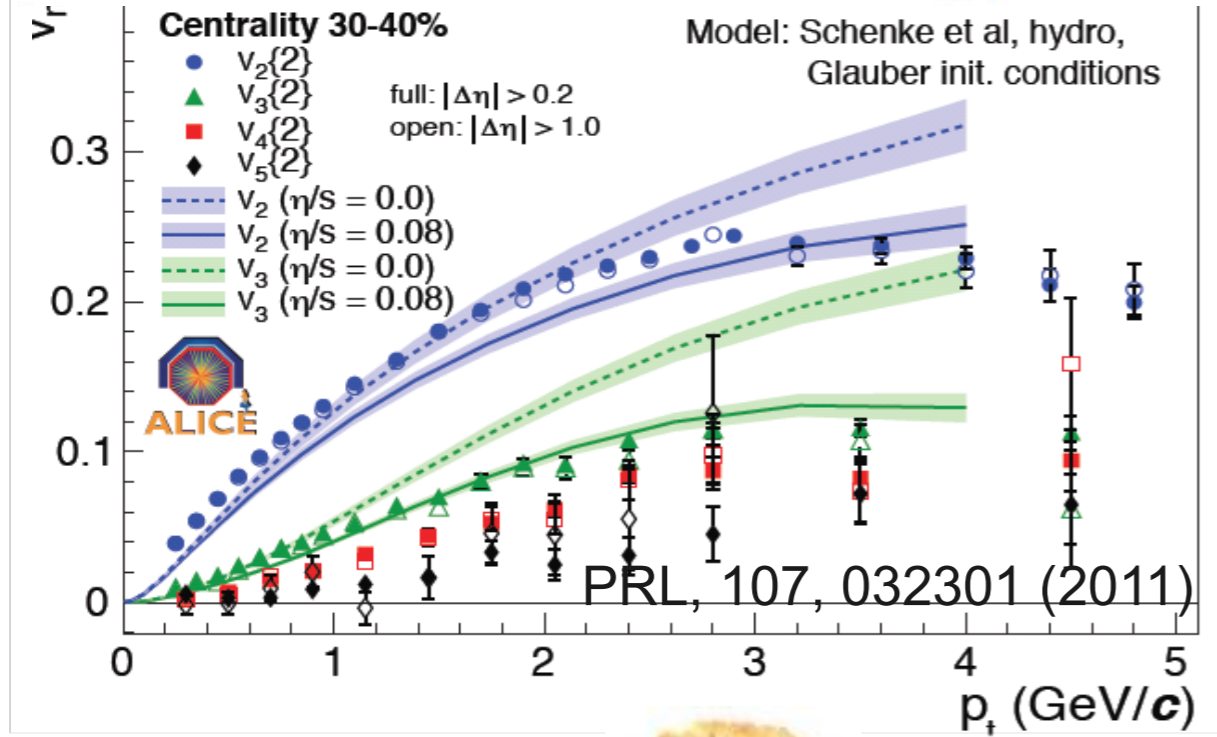
- **Collective flow dominates to about 3-4 GeV/c for all $n > 1$**
- **Description breaks for high p_T or peripheral collisions**
- **For low p_T : double peak and ridge structures seen in two particle correlations are naturally explained by measured anisotropic flow coefficients**

Higher harmonics w.r.t. to event plane



v_3 - triangular flow :

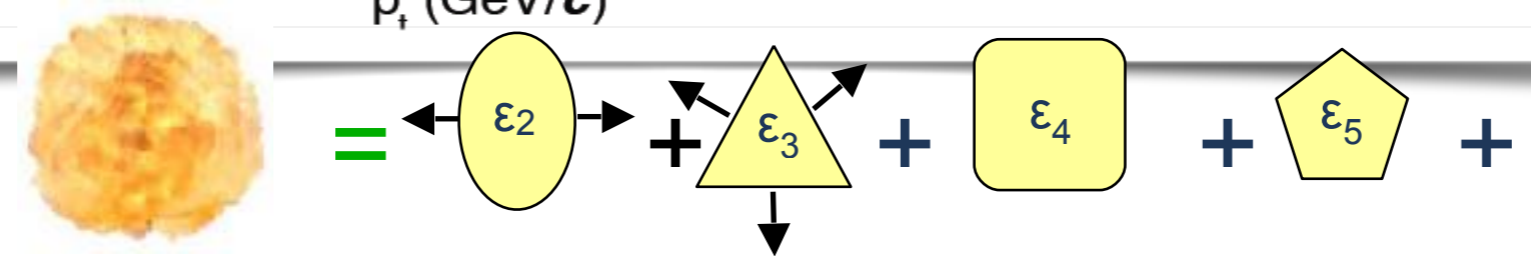
- weak centrality dependence
- vanishes as expected when measured w.r.t. reaction plane



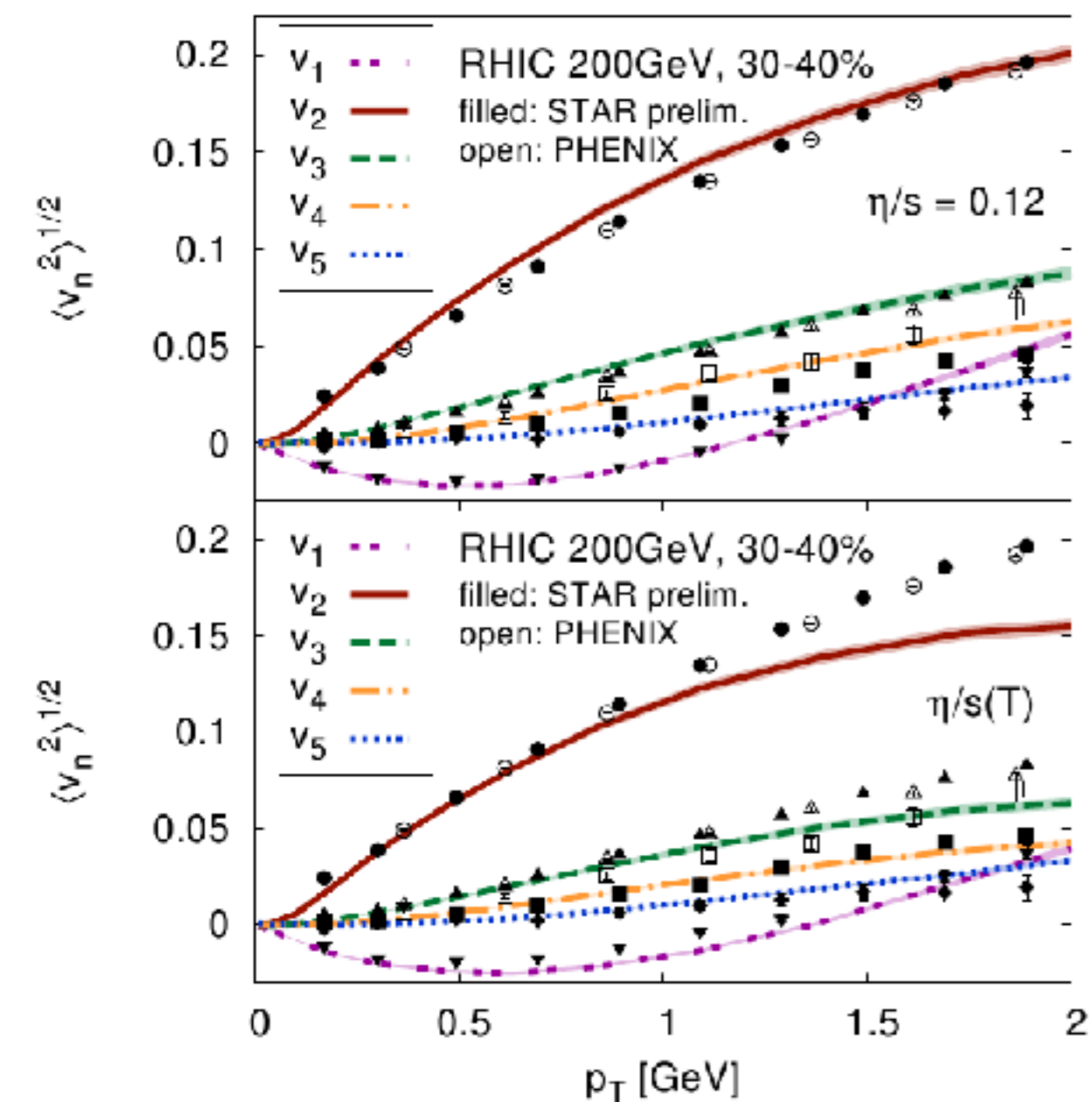
Similar p_T dependence for all v_n

Higher harmonics - additional constraints on η/s

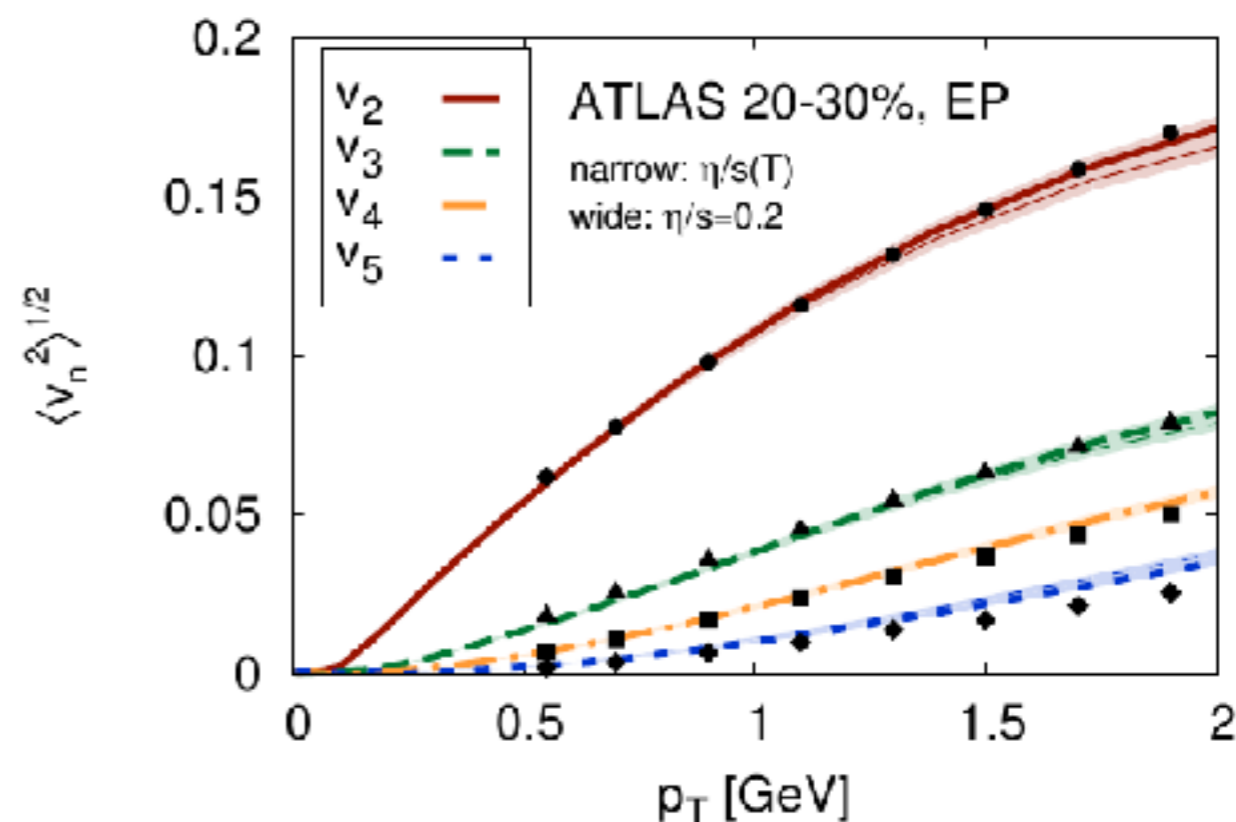
η/s small, similar as at RHIC



(viscous) fluid dynamics works!



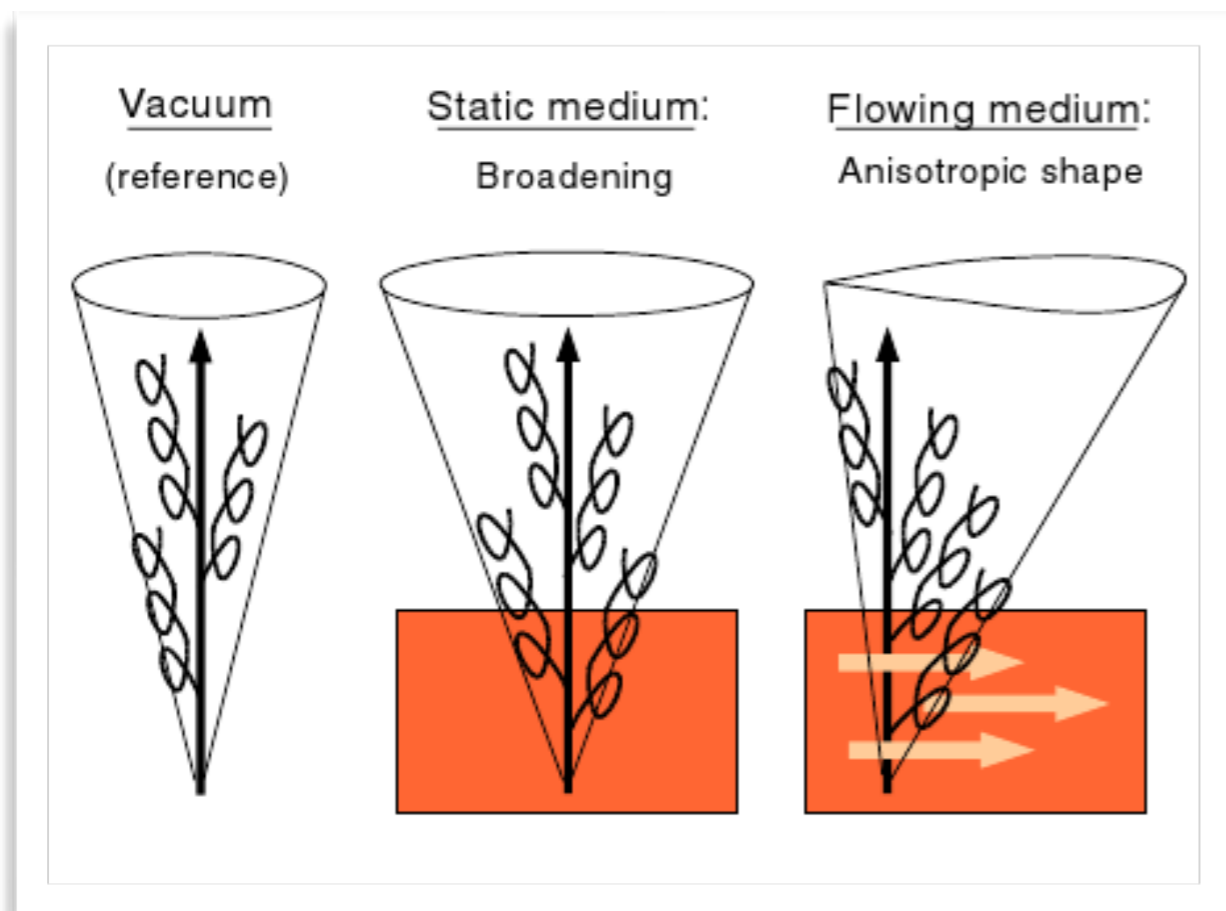
Good description of data for $n \neq 1$



Small shear viscosity over entropy ratio:

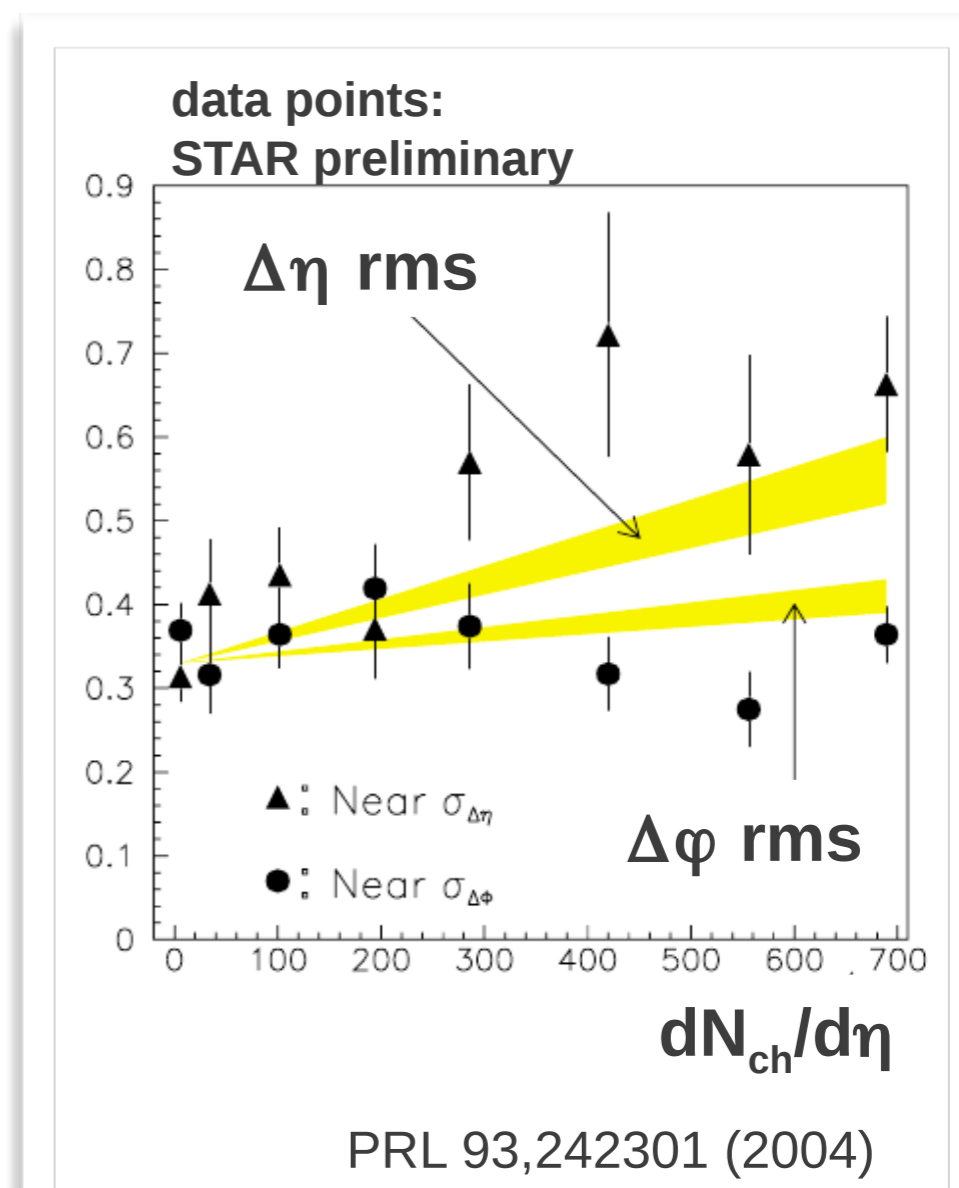
$$\frac{\eta}{s} \leq 0.25$$

Jet-medium-flow coupling via two particle correlations?



N. Armesto, C. Salgado, U. Wiedemann:
Measuring the Collective Flow with Jets

[PRL 93,242301 (2004)]



=> LHC? - more jets
+ somewhat more
flow...

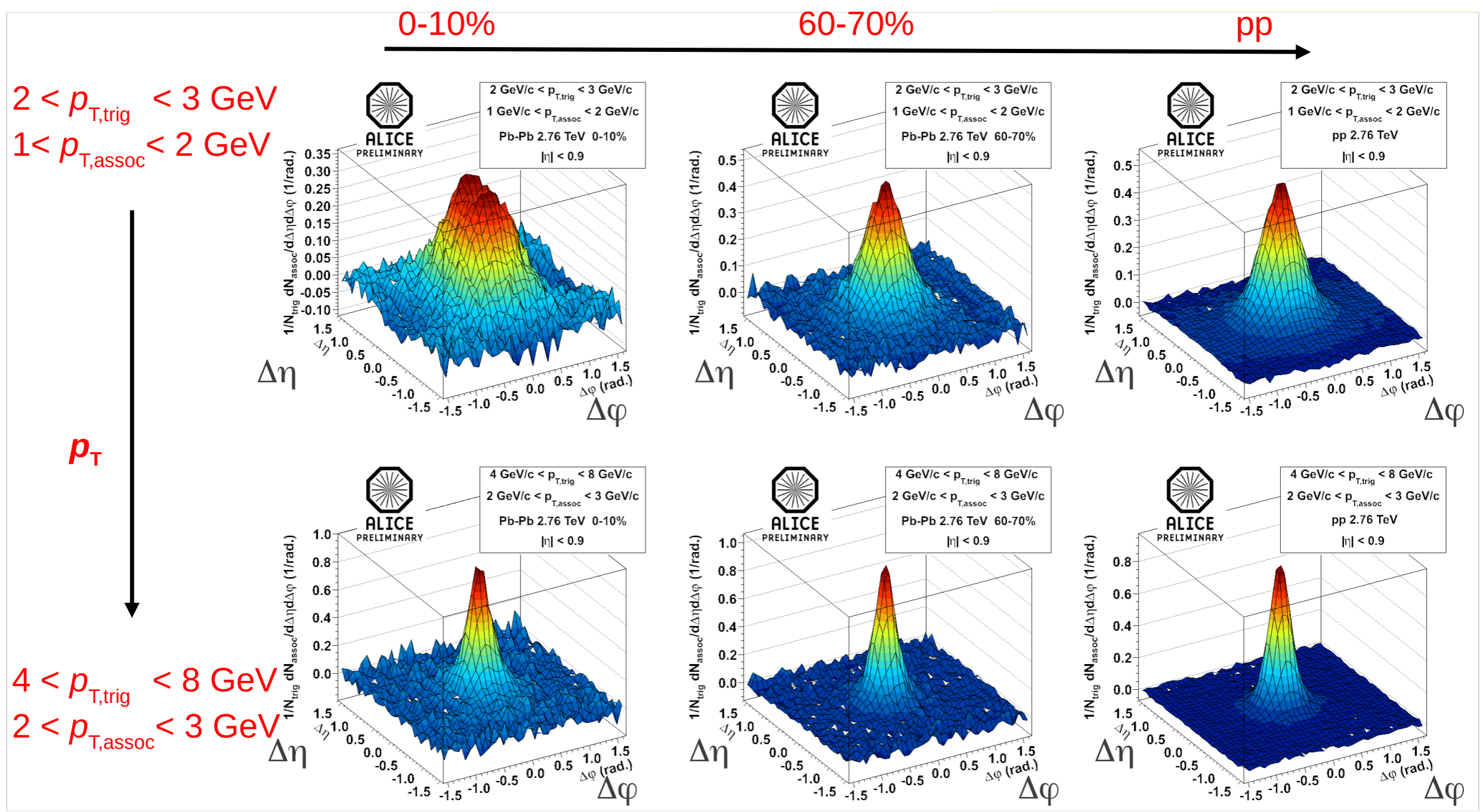
Jet-peak shape

Resolution - intermediate p_T

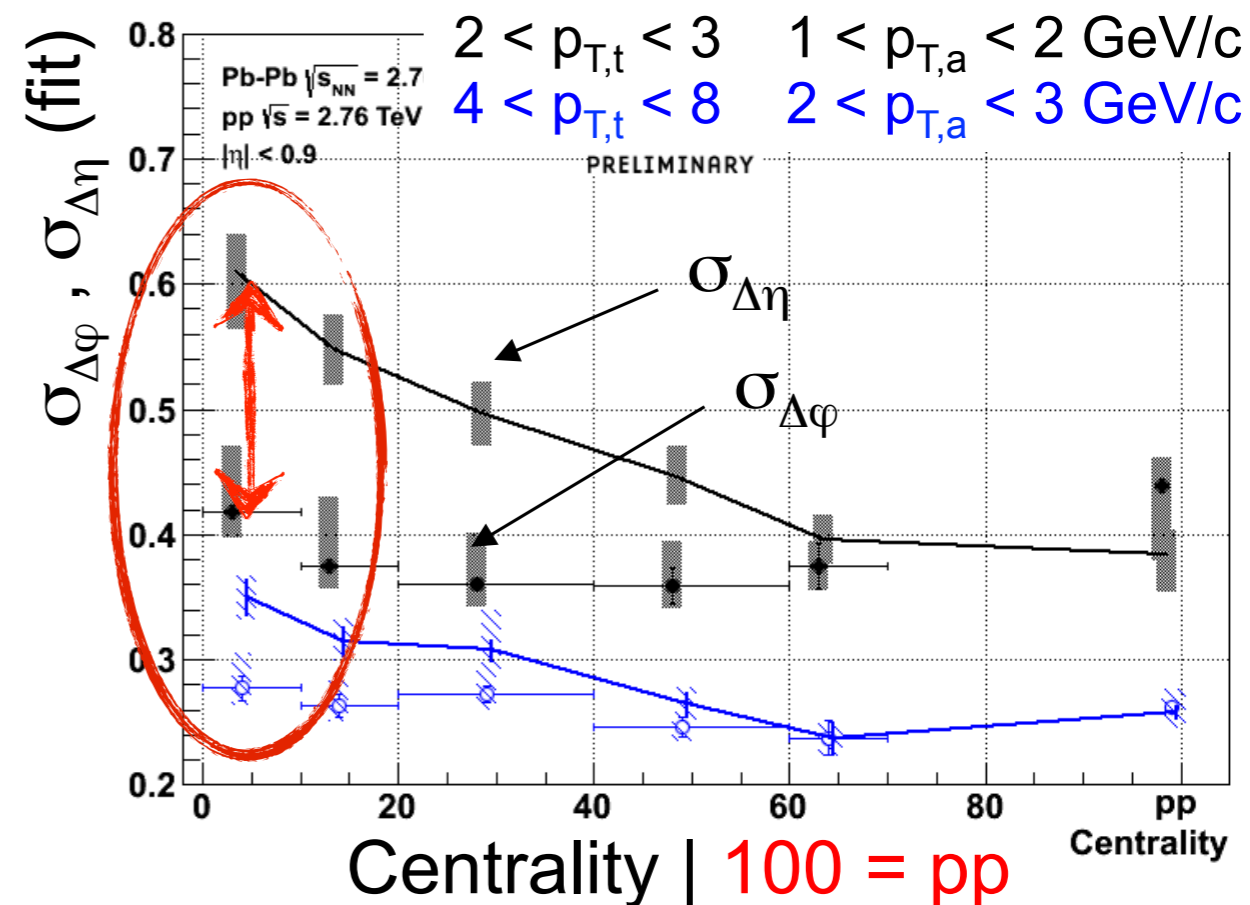
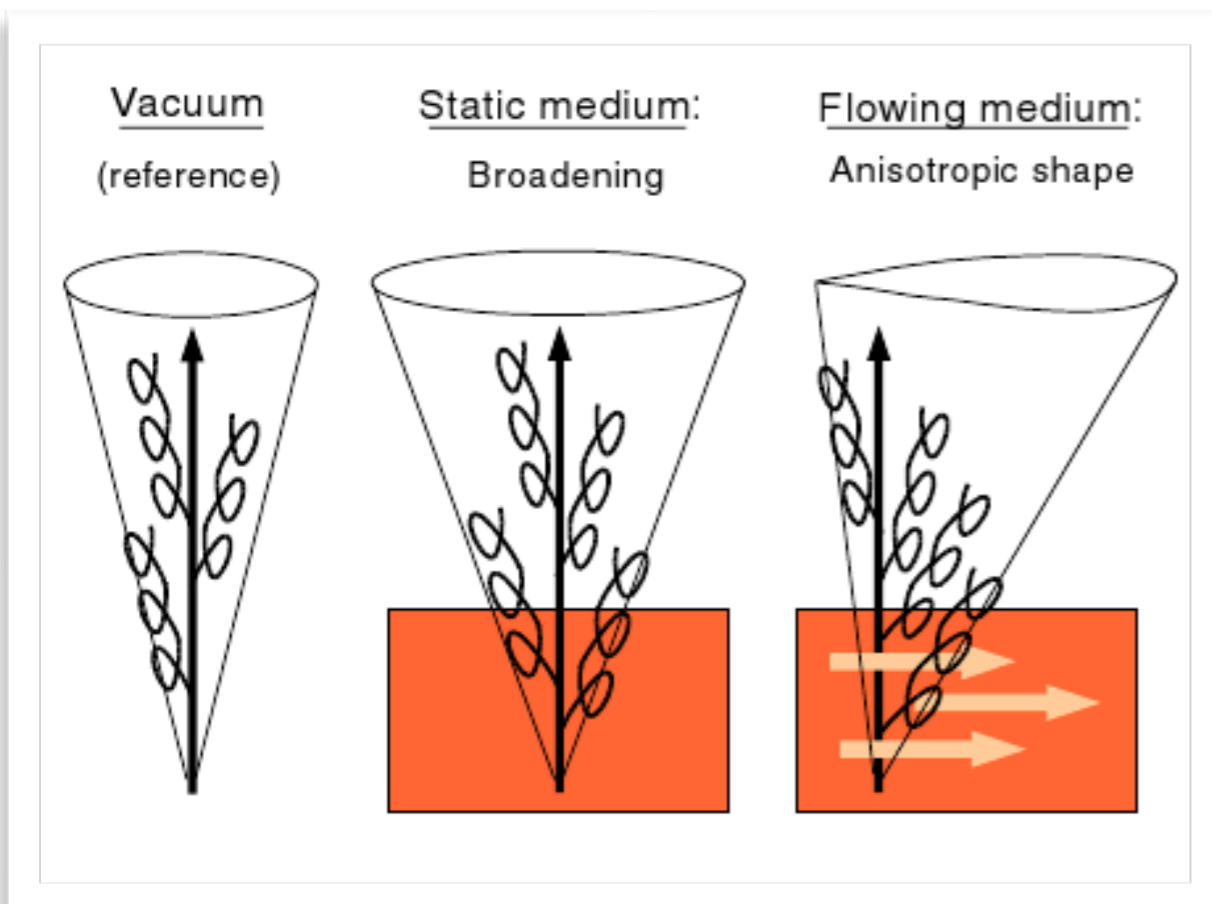
Wider peak in central collisions
Peripheral and $p-p$ similar shape

Strong p_T dependence

=> Characterize the peak



Measuring widths of the correlations in azimuth and pseudo-rapidity

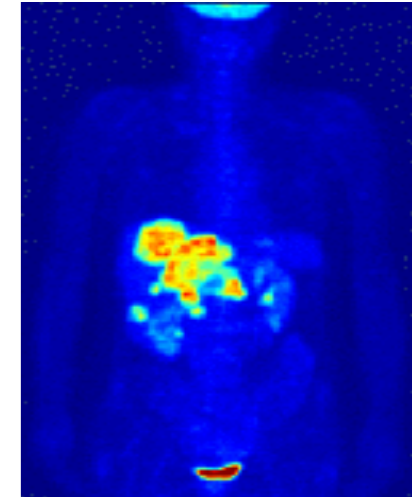


Measure of jets interactions with longitudinal flow
(?)

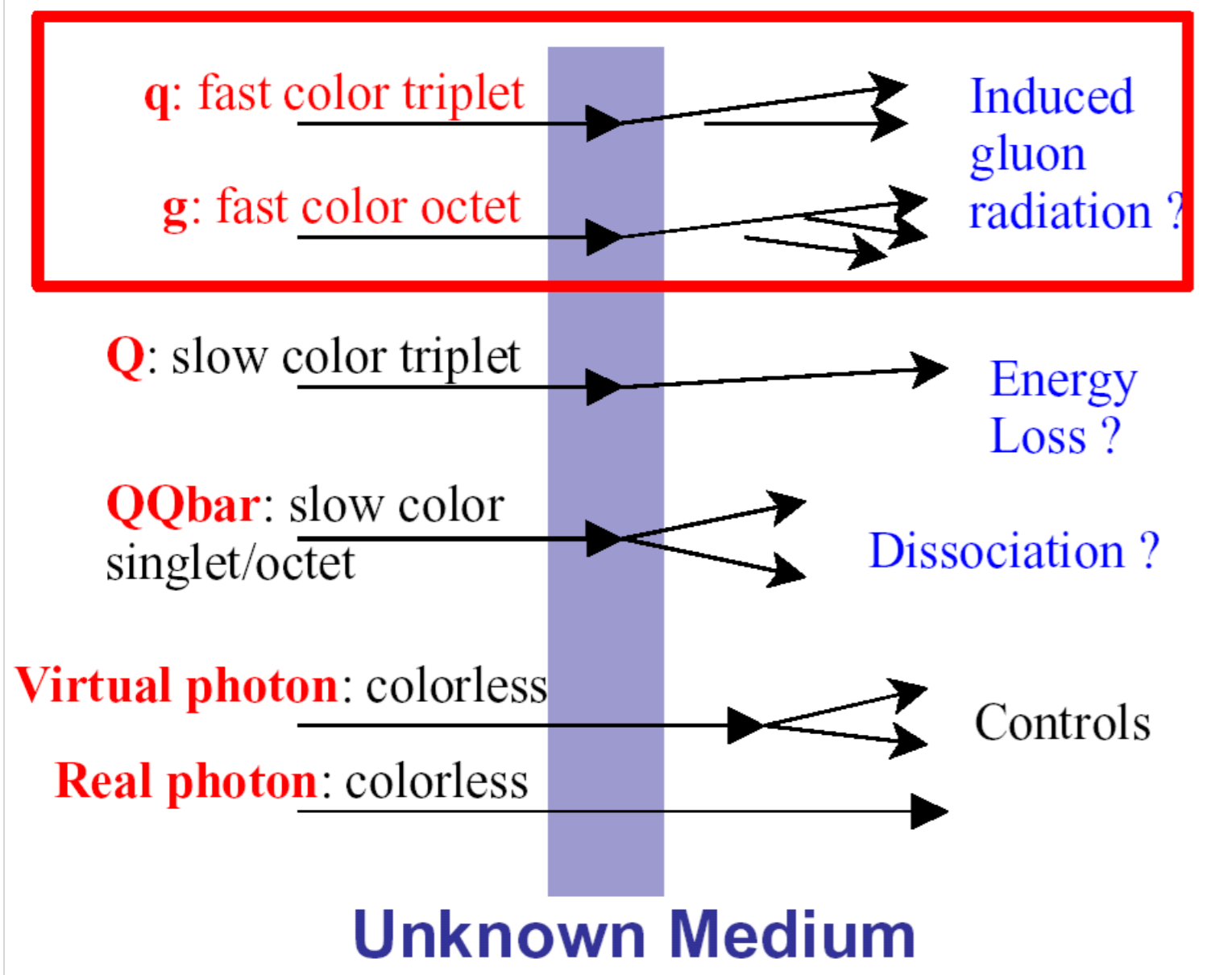
- **AMPT (A MultiPhase Transport Code)**
 - Initial conditions simulated using HIJING
 - Parton scattering
 - Hadronization: Lund model + coalescence
 - Hadron scattering
- AMPT describes the main features of the near-side shape evolution observed in data

Probing an unknown
medium...

Probing the unknown medium...



Human



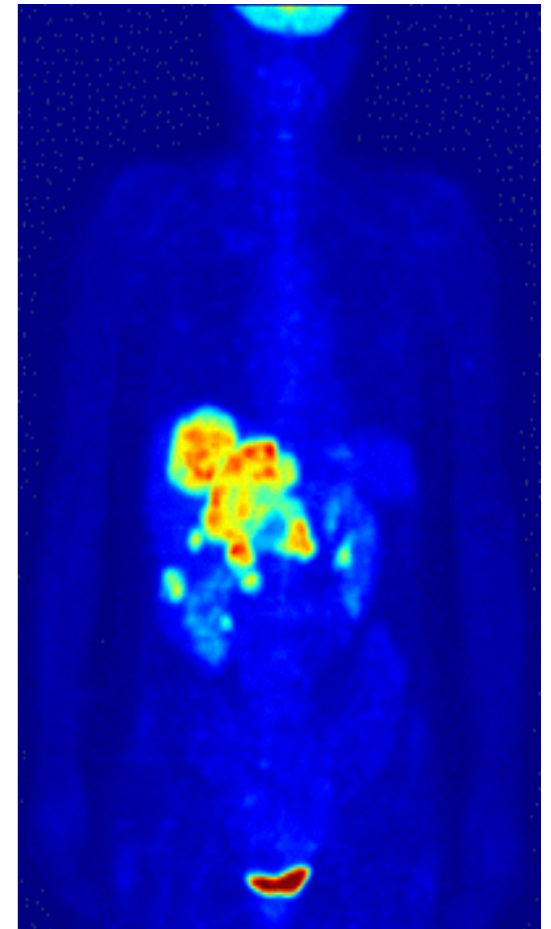
jet suppression (quenching)

charm/bottom dynamics

J/ψ & γ

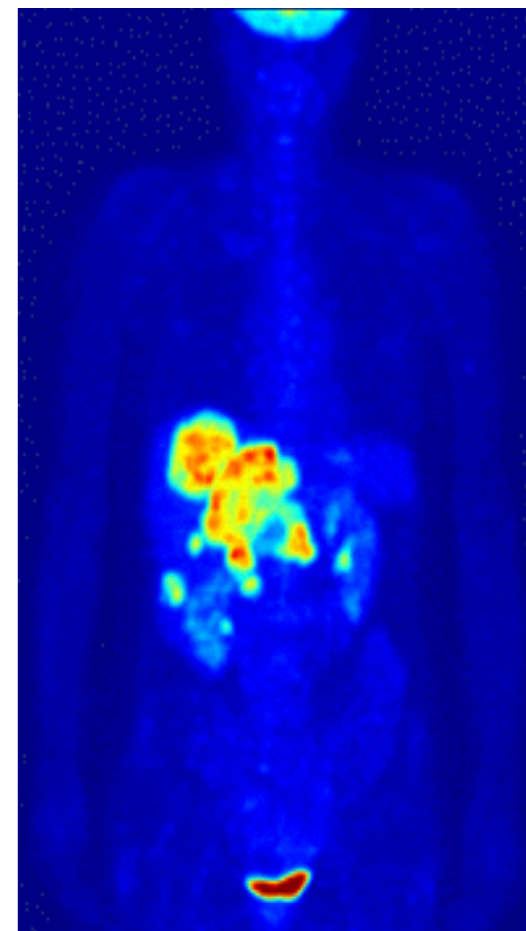
color-less particles

How to probe a patient that comes for only $O(\text{micro sec.})$?



Human body

How to probe a medium that lasts only $O(\text{micro sec.})$?

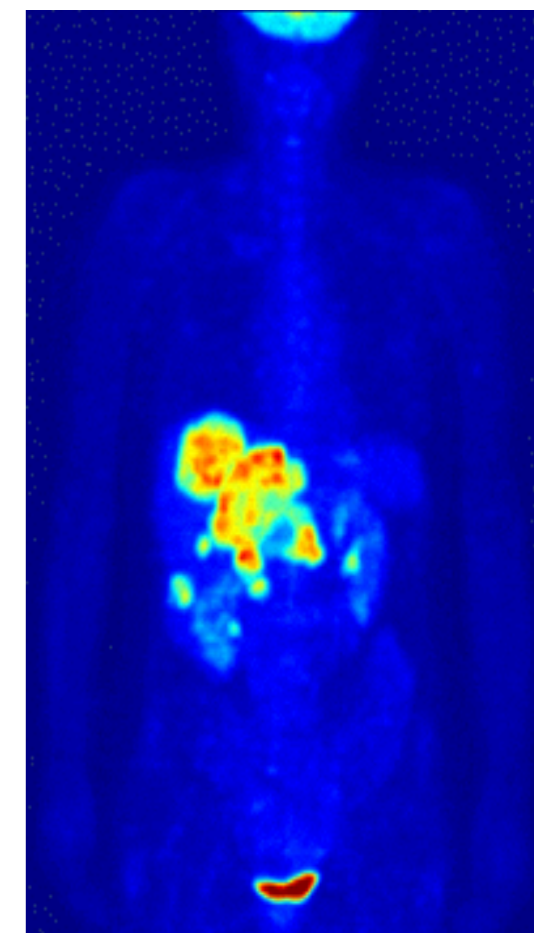
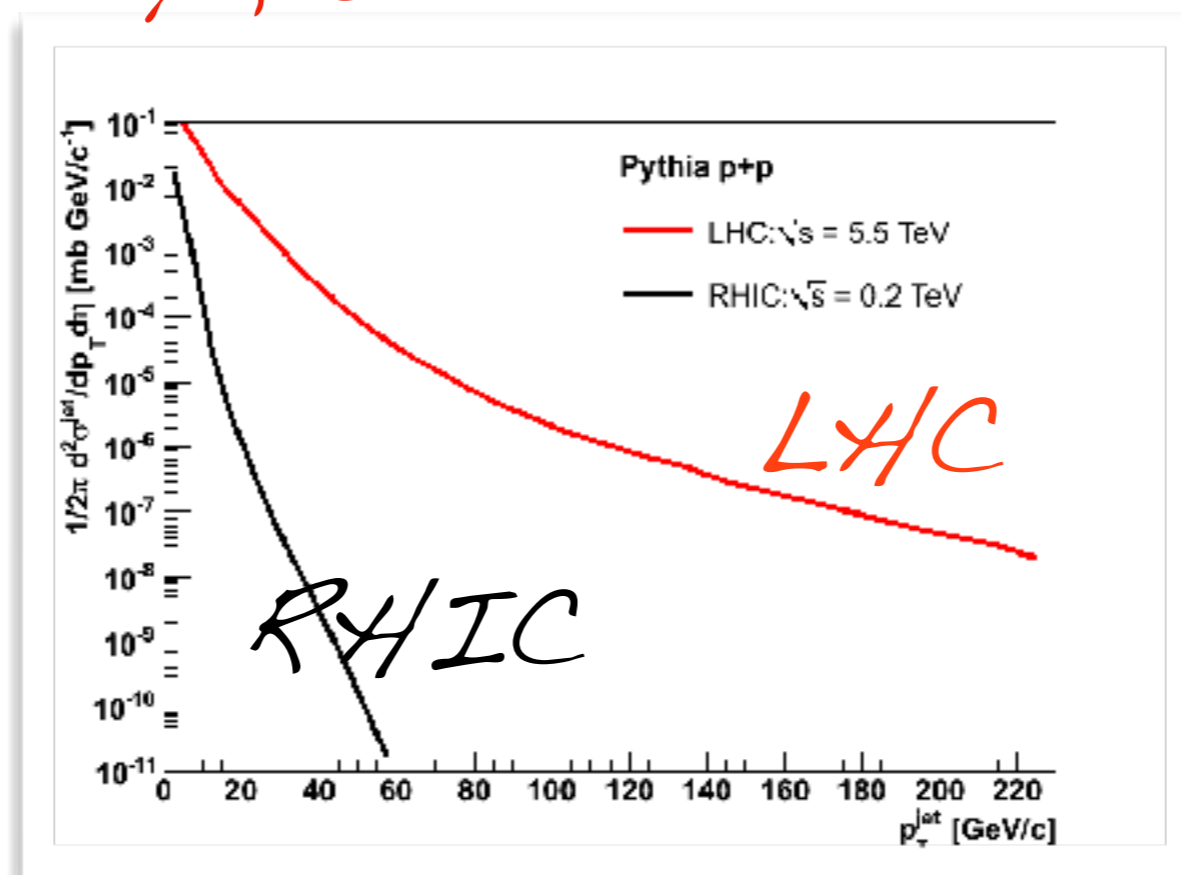


Human body

... to probe the short lived medium

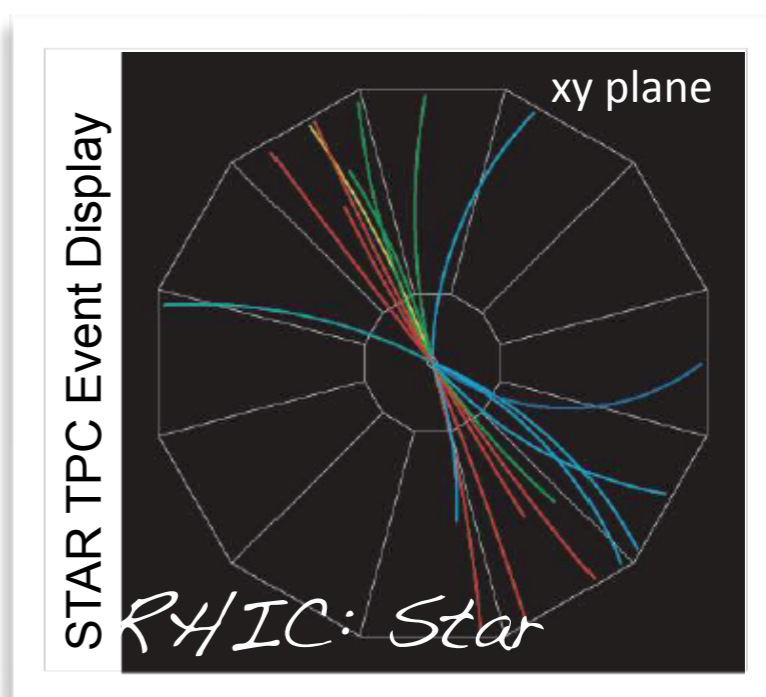
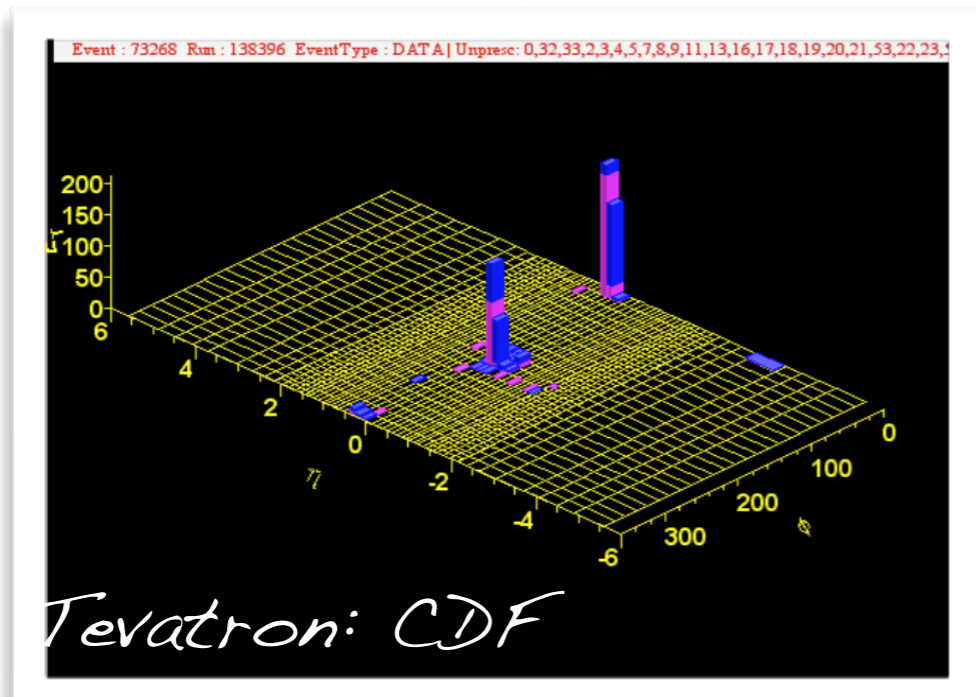
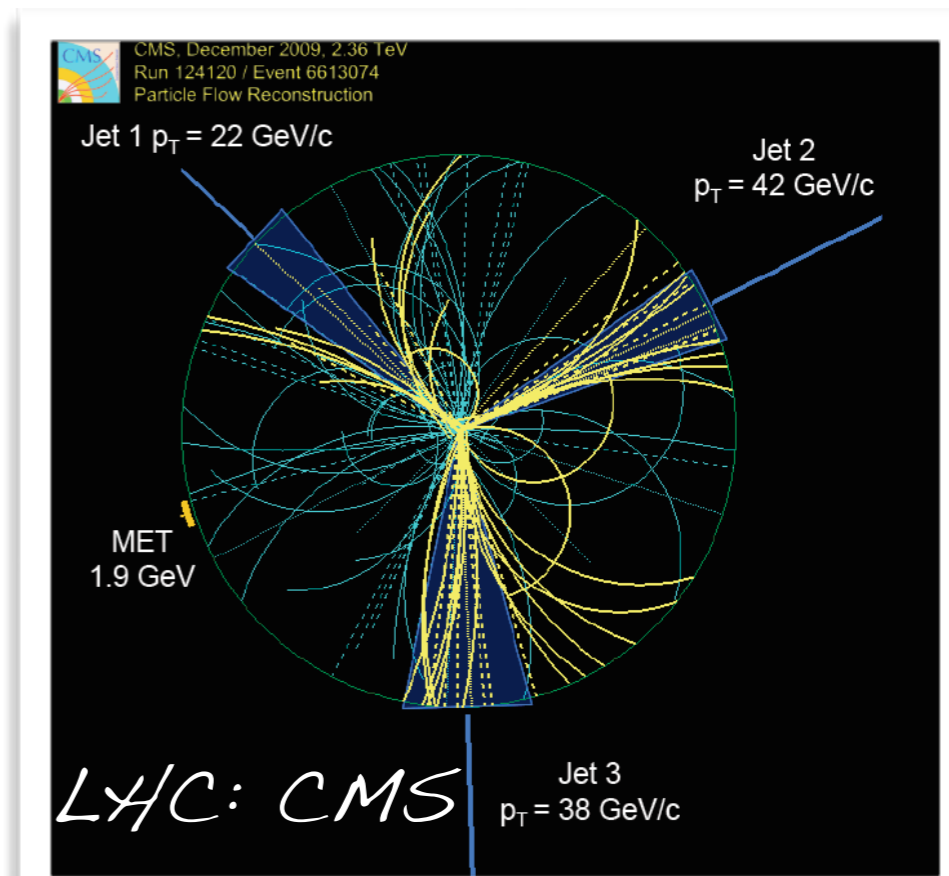
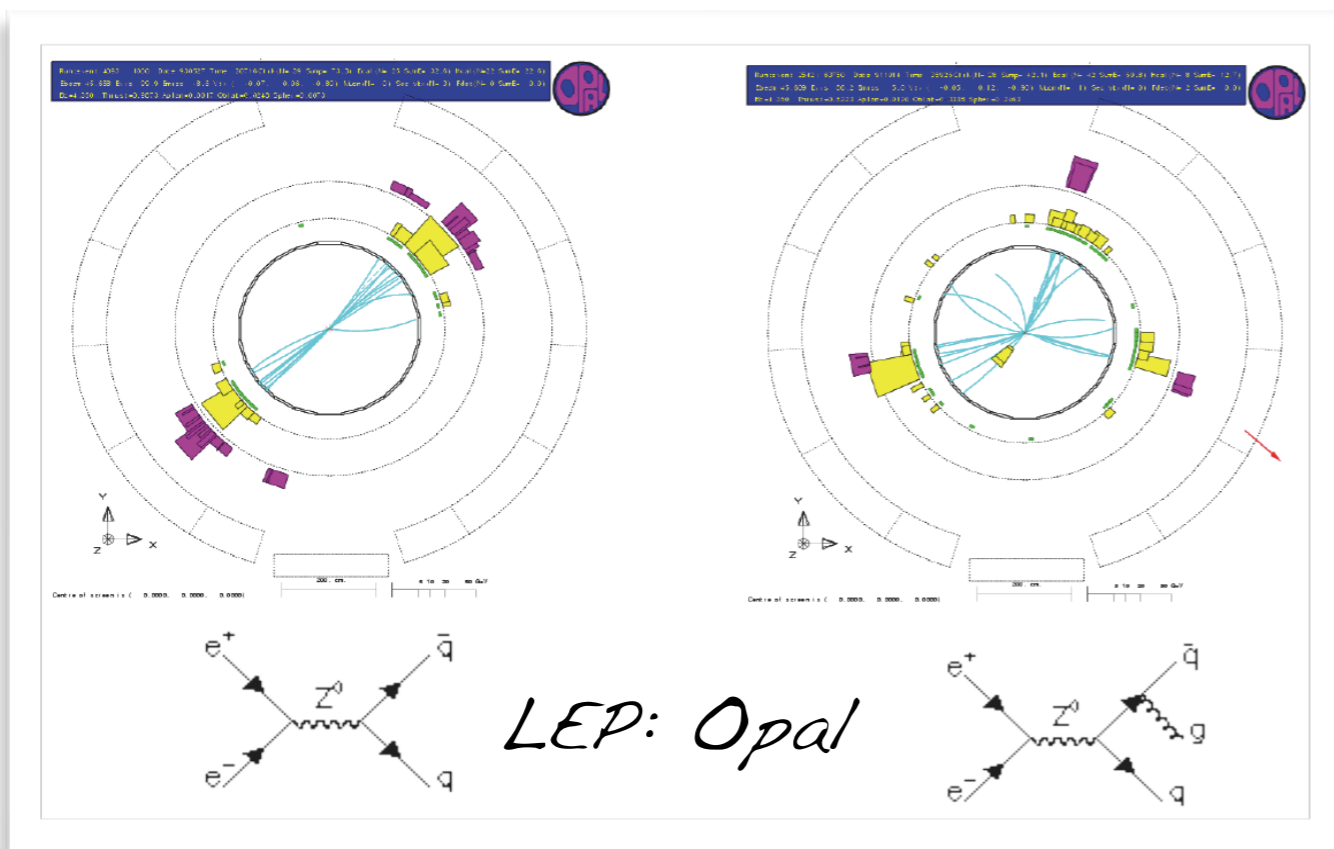
=> use "auto-generated probes" -
heavy-ion collisions at high-energies
produce internally high-energy partons
(fragment into jets of particles)...

<=> critical input from pp (vacuum)
measurements - pQCD



Human body

Jets in collider experiments



What is a jet?

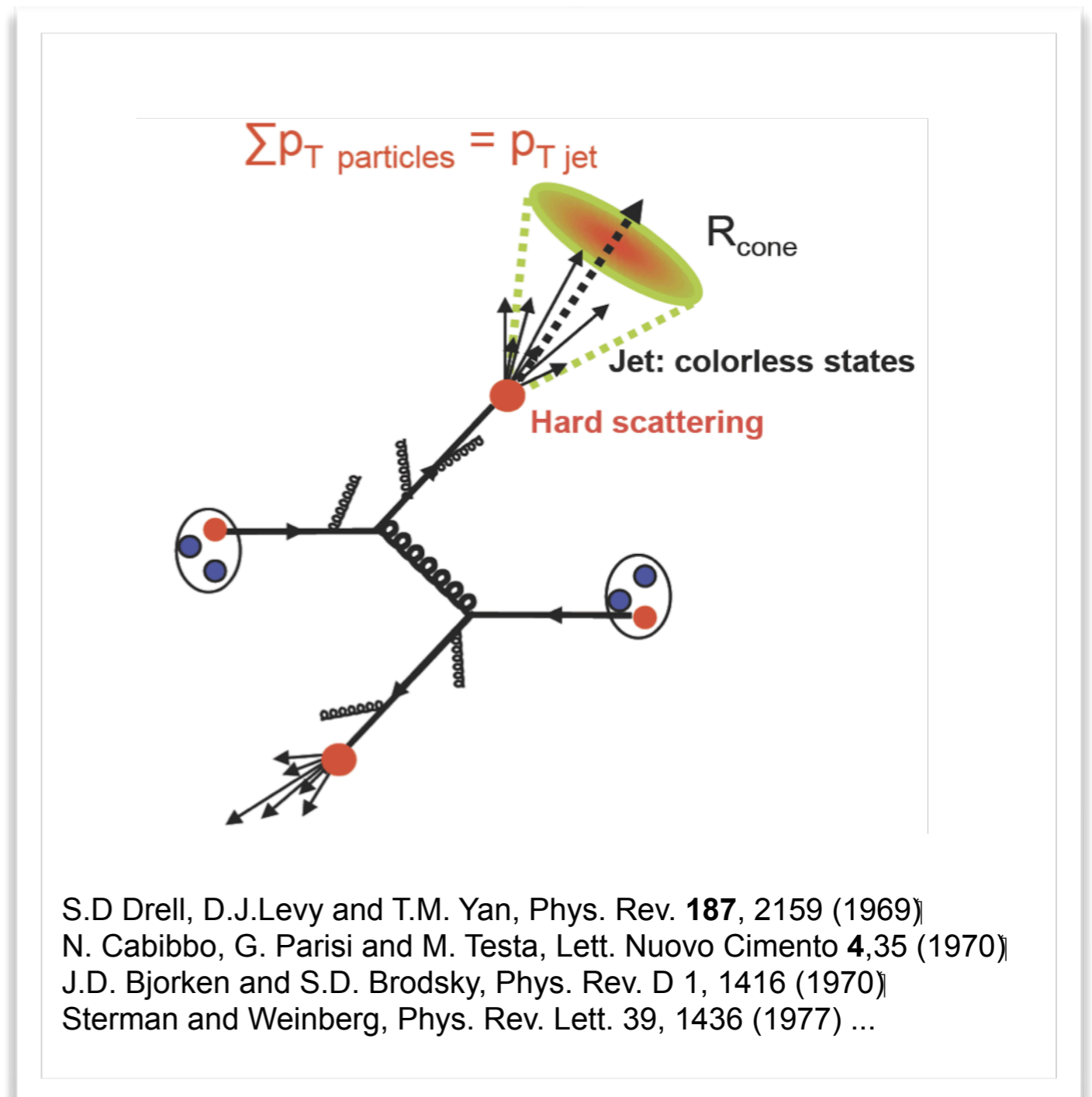
A spray of collimated showers/particles

- Hardly ever better defined...

Jet = Parton AND its
radiation

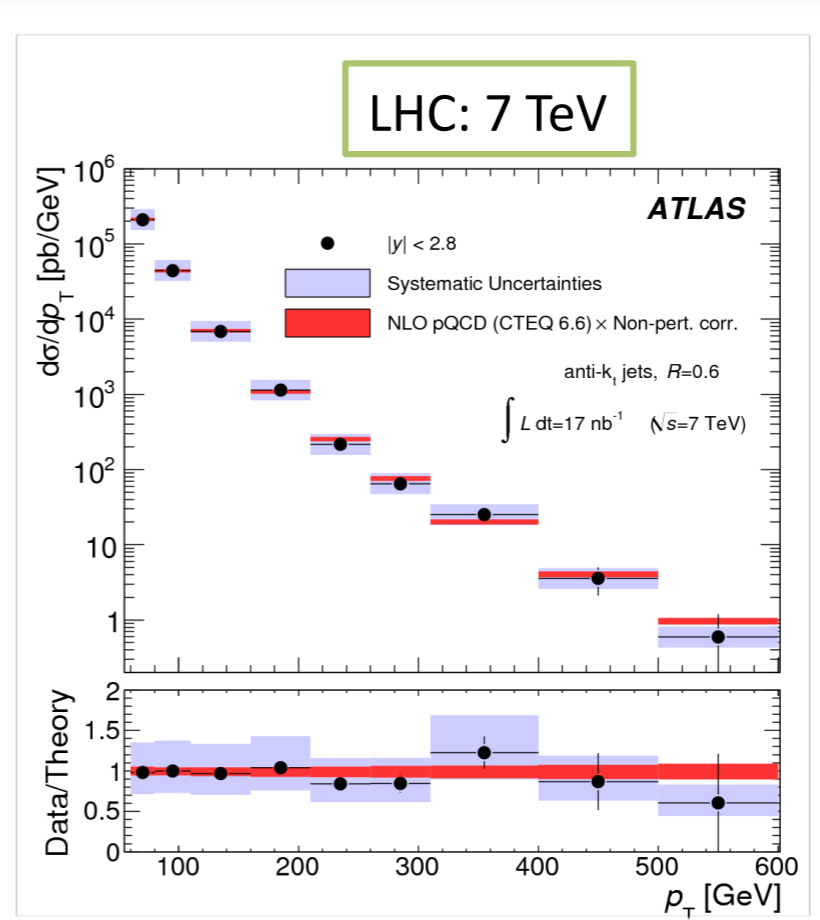
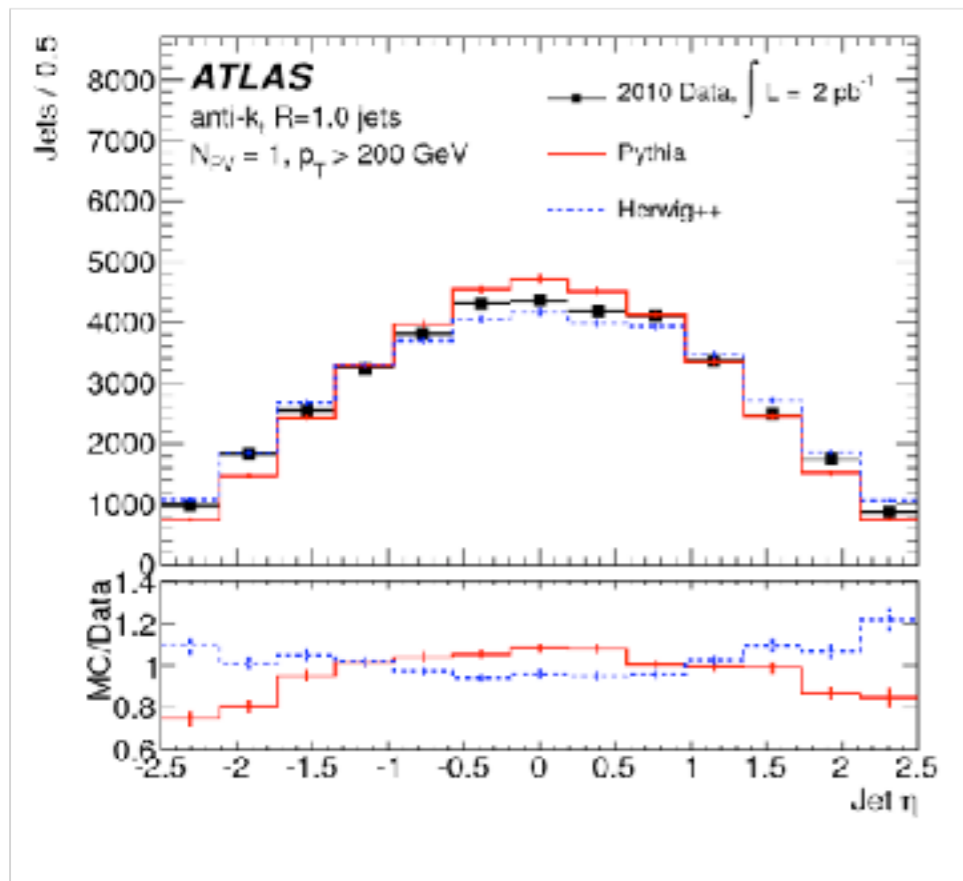
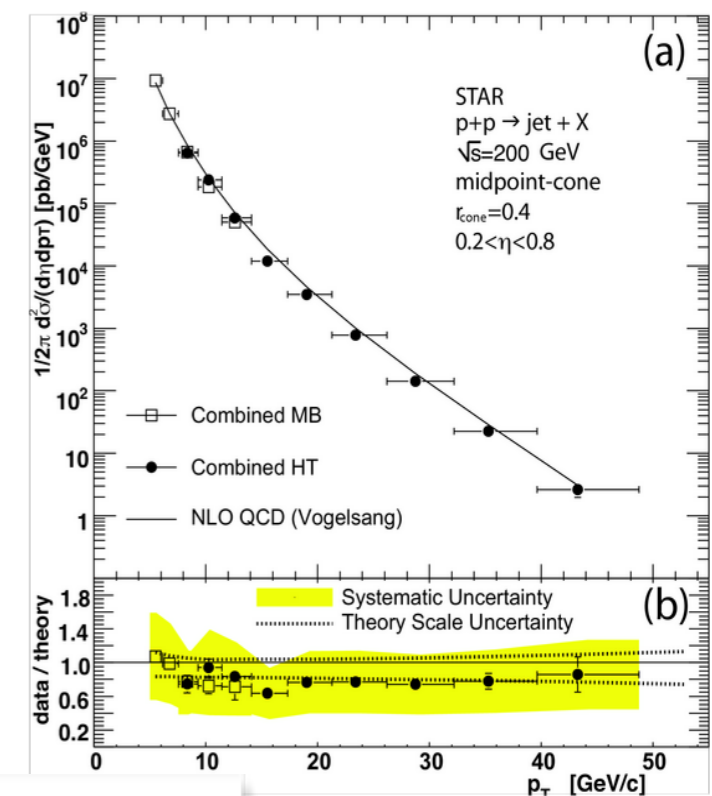
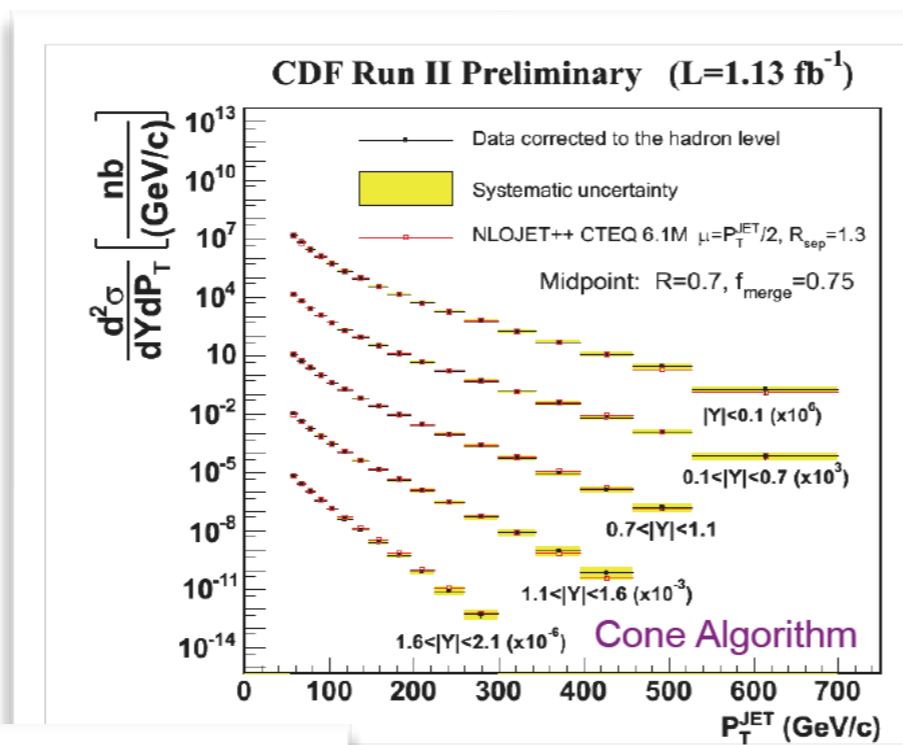
Note: experiment measures
spray of particles
(~hadrons)

Jets (unlike single hadrons)
are objects which are
"better" understood/
calculable within pQCD



Jets in collider experiments

Jets are fairly well known by now... and well described by theory and MC
 => attractive tool for heavy-ions



Hadronic collisions: pQCD and jets

$$E \frac{d^3\sigma}{dp^3} \propto f_{a/A}(x_a, Q^2) \otimes f_{b/B}(x_b, Q^2) \otimes \frac{d\hat{\sigma}^{ab \rightarrow cd}}{dt} \otimes D_{h/c}(z_c, Q^2)$$

Jets are defined via rigorous (collinear and infrared safe) clustering algorithms

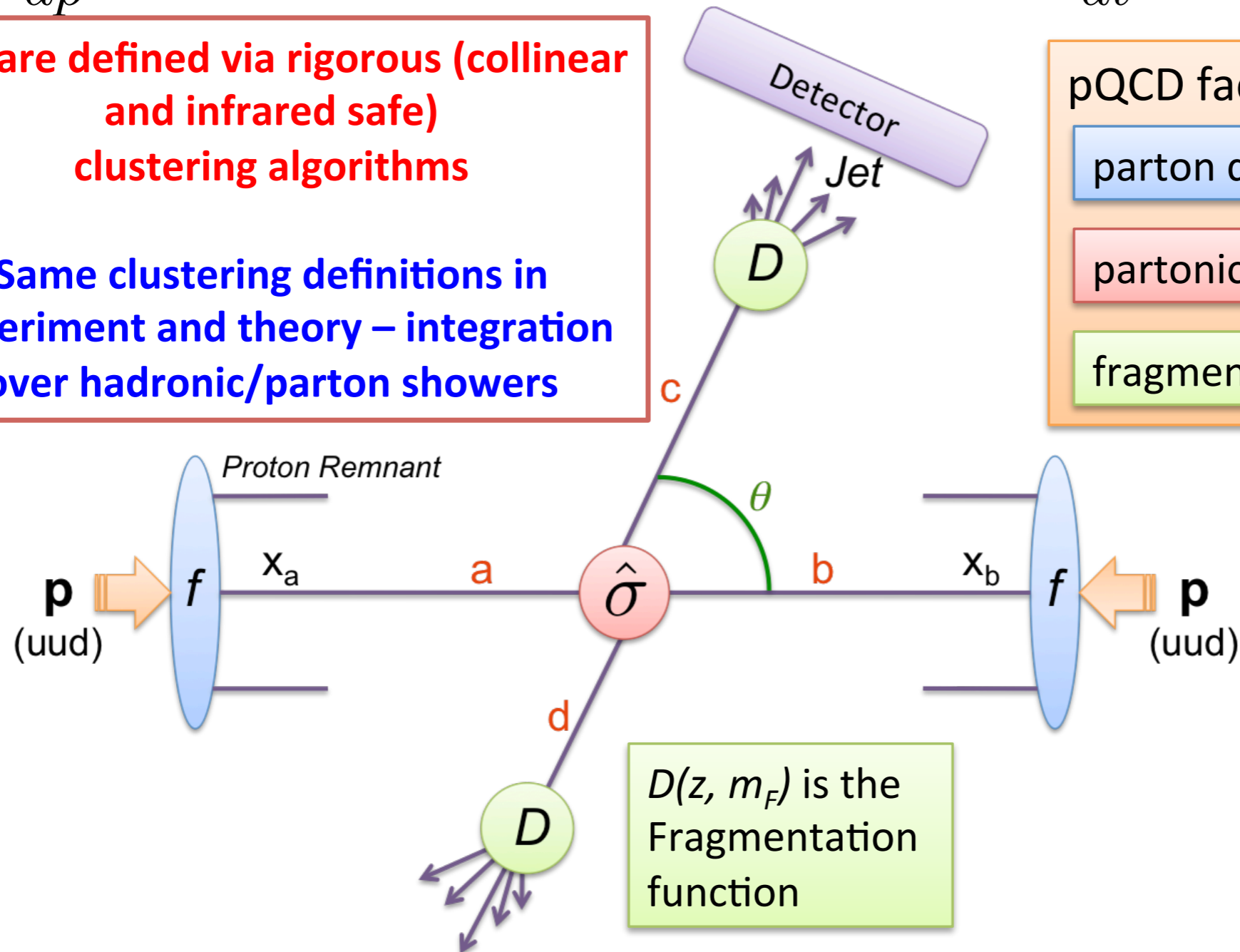
Same clustering definitions in experiment and theory – integration over hadronic/parton showers

pQCD factorization:

parton distribution fn $f_{a/A}$

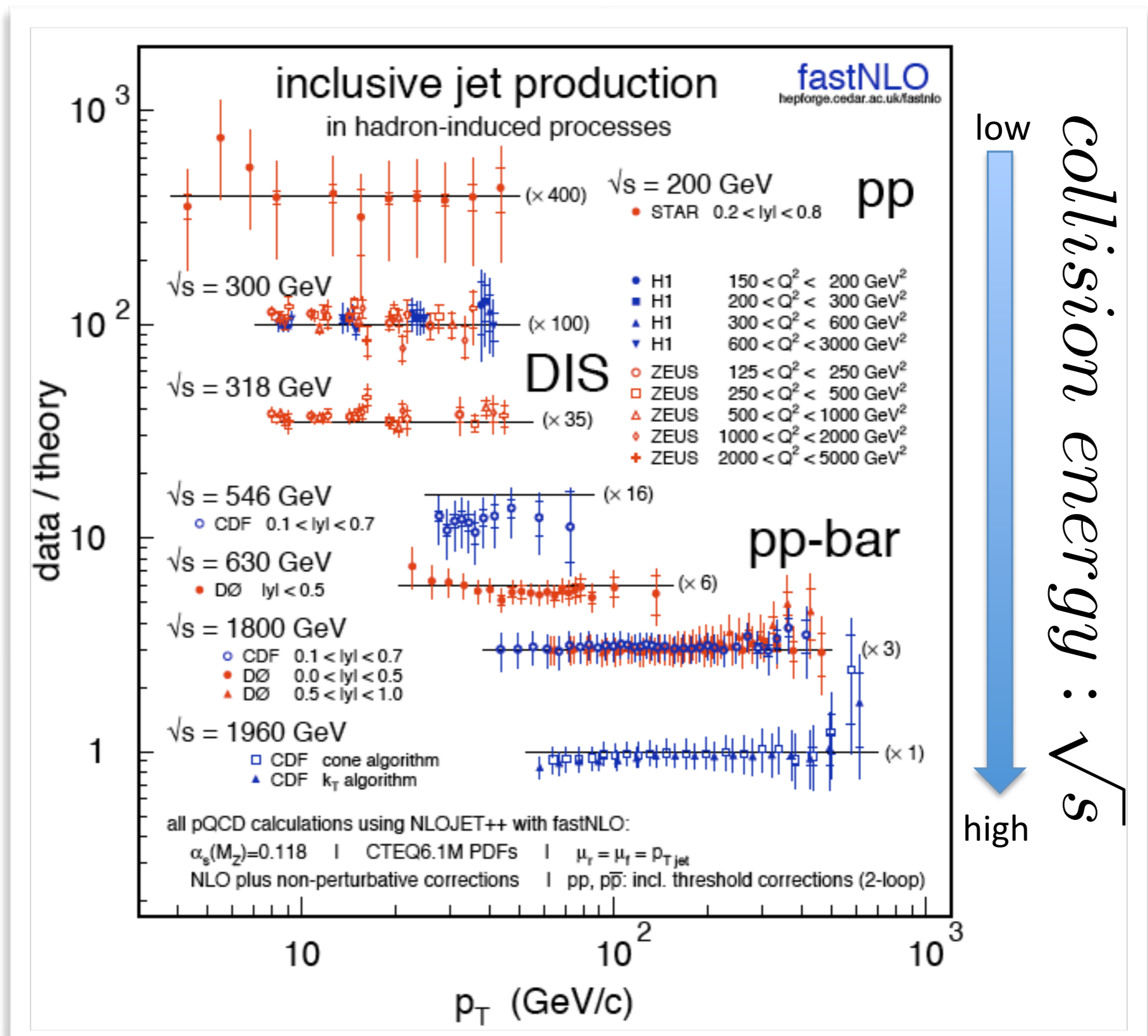
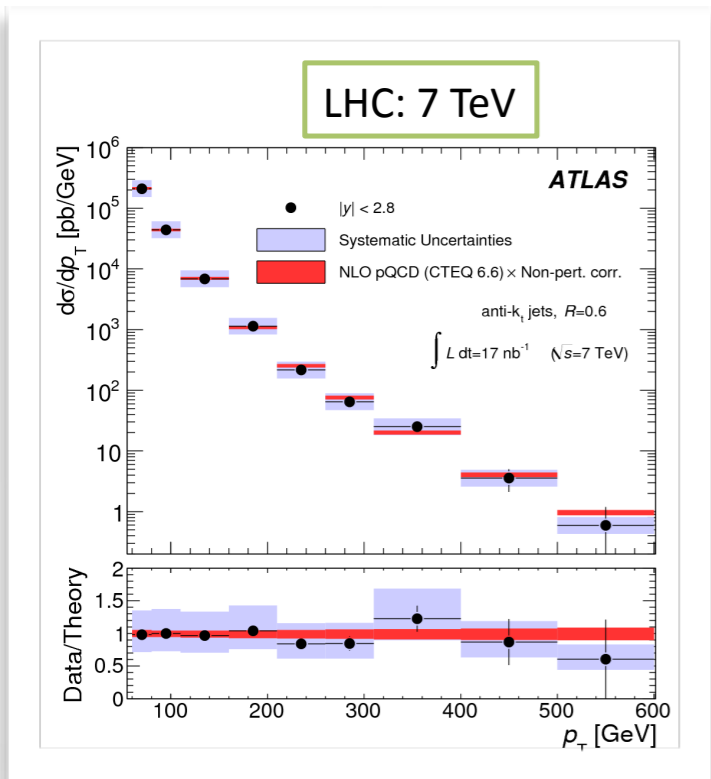
partonic cross section

fragmentation fn $D_{h/c}$

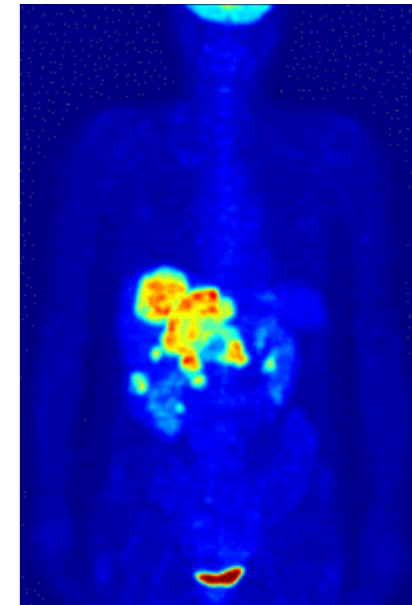


Inclusive jet production: pQCD & data

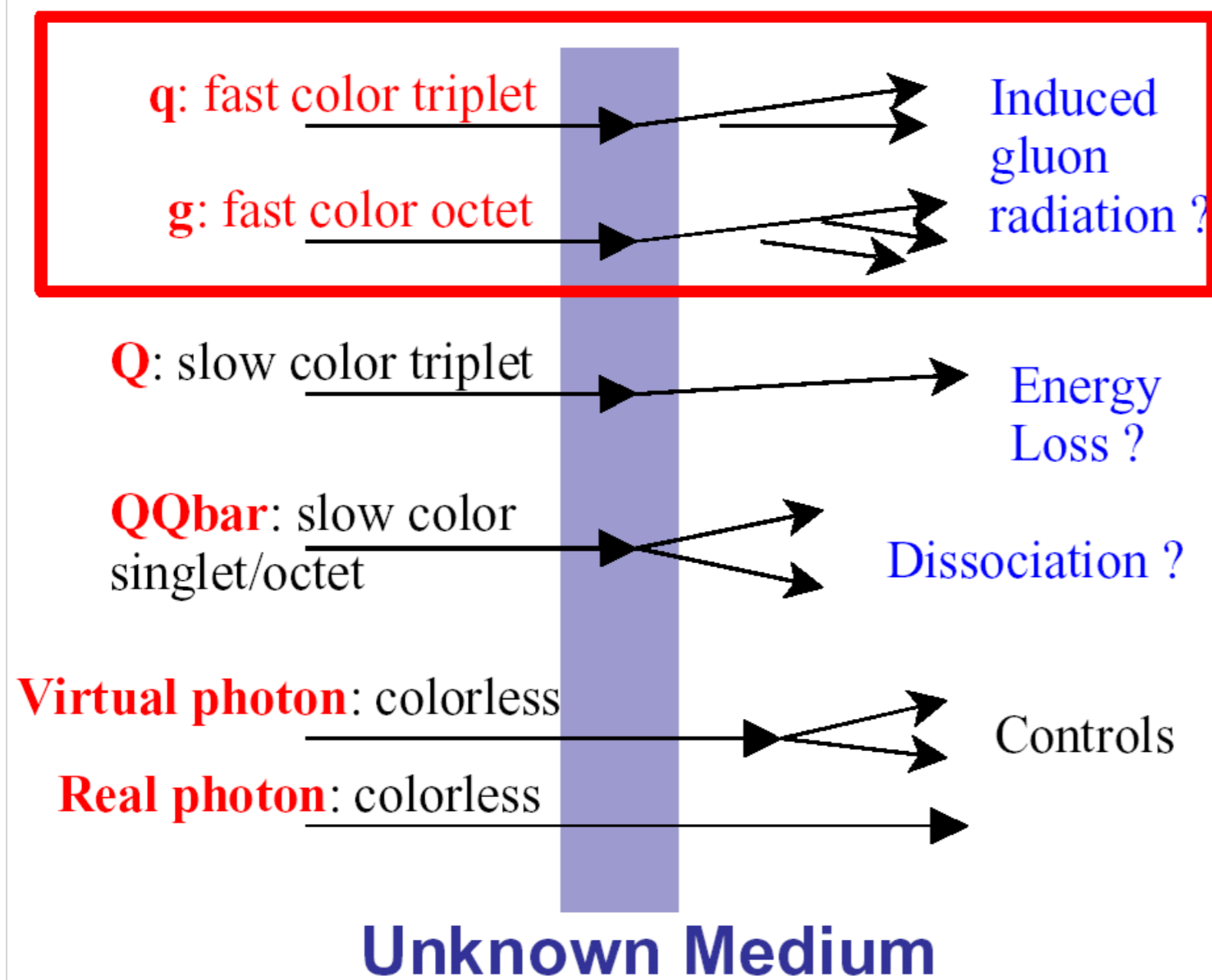
Jets are fairly well known by now... and well described by theory and MC
 => attractive tool for heavy-ions



Probing the unknown medium...



Human body



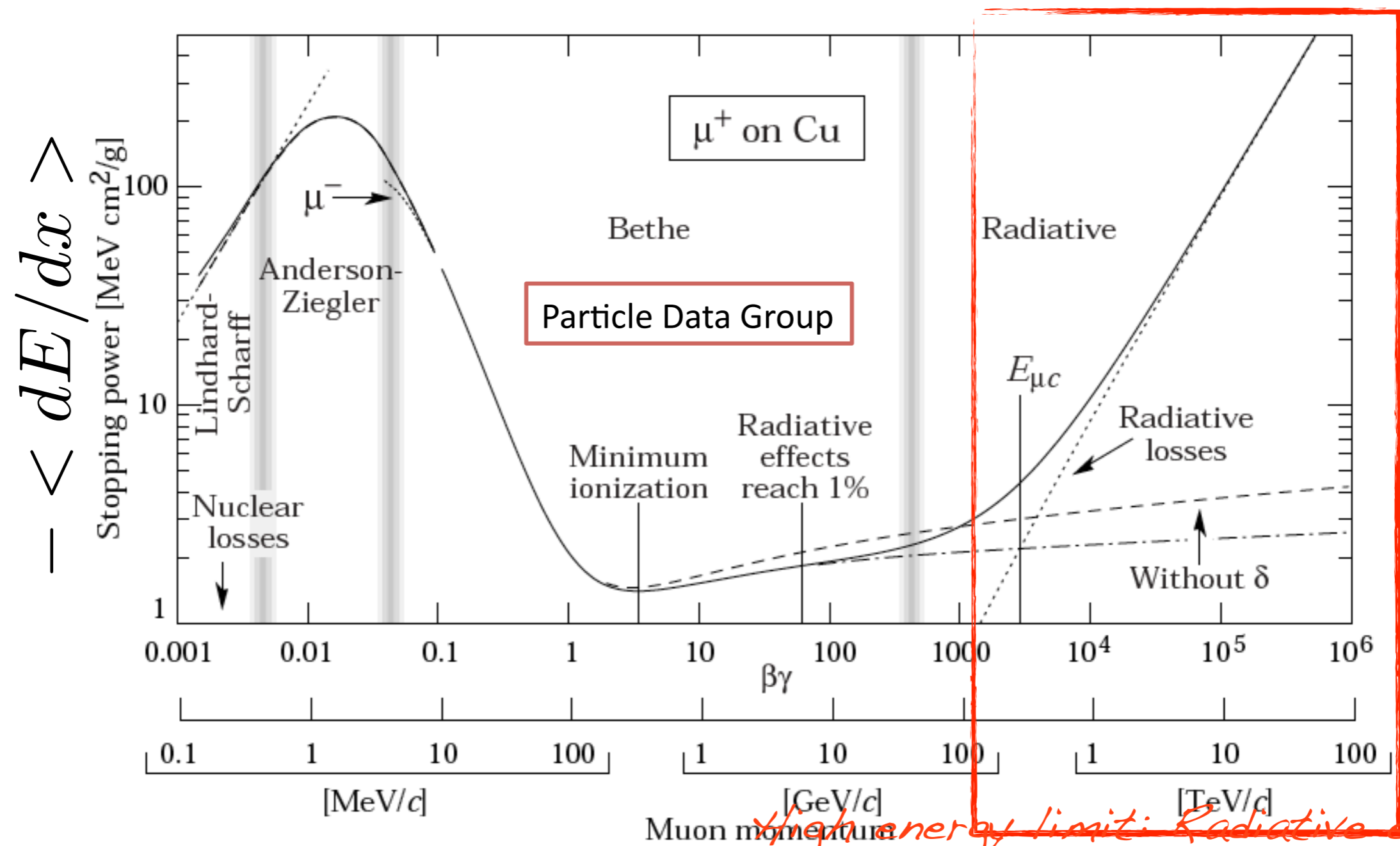
*jet suppression
(quenching)*

*charm/bottom
dynamics*

J/ψ & γ

color-less particles

QED: Passage of electrically charged particle through

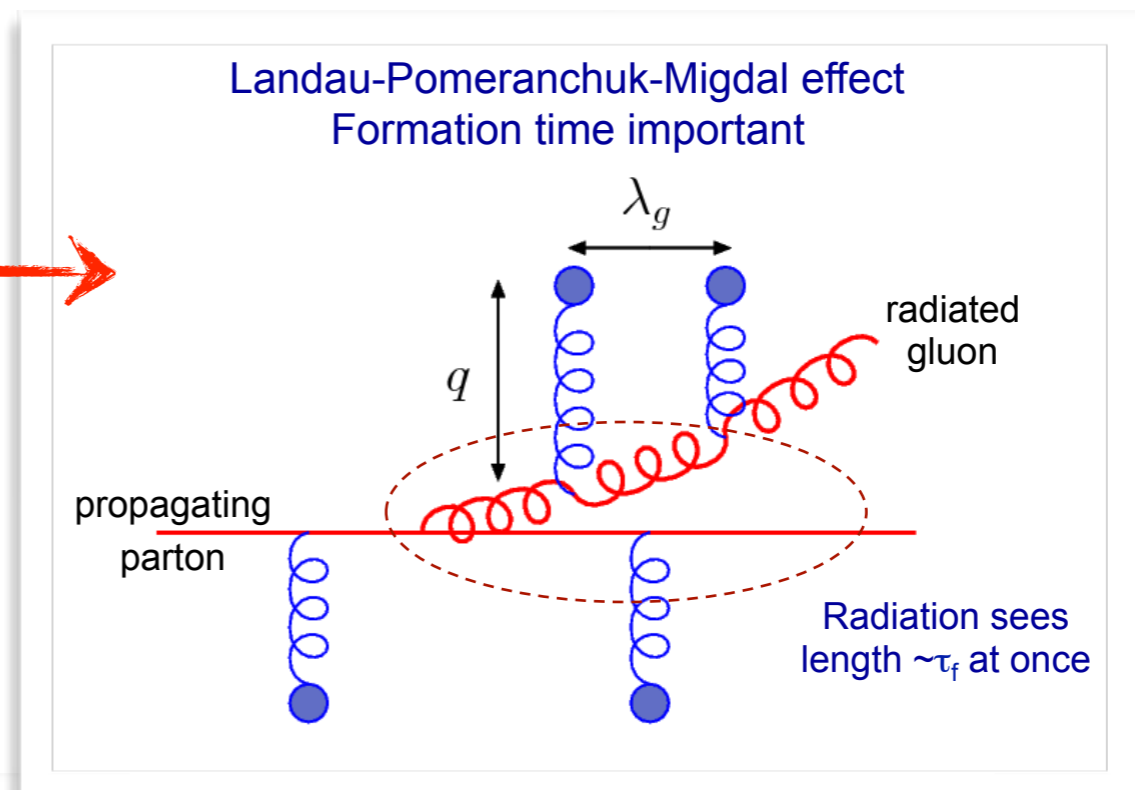
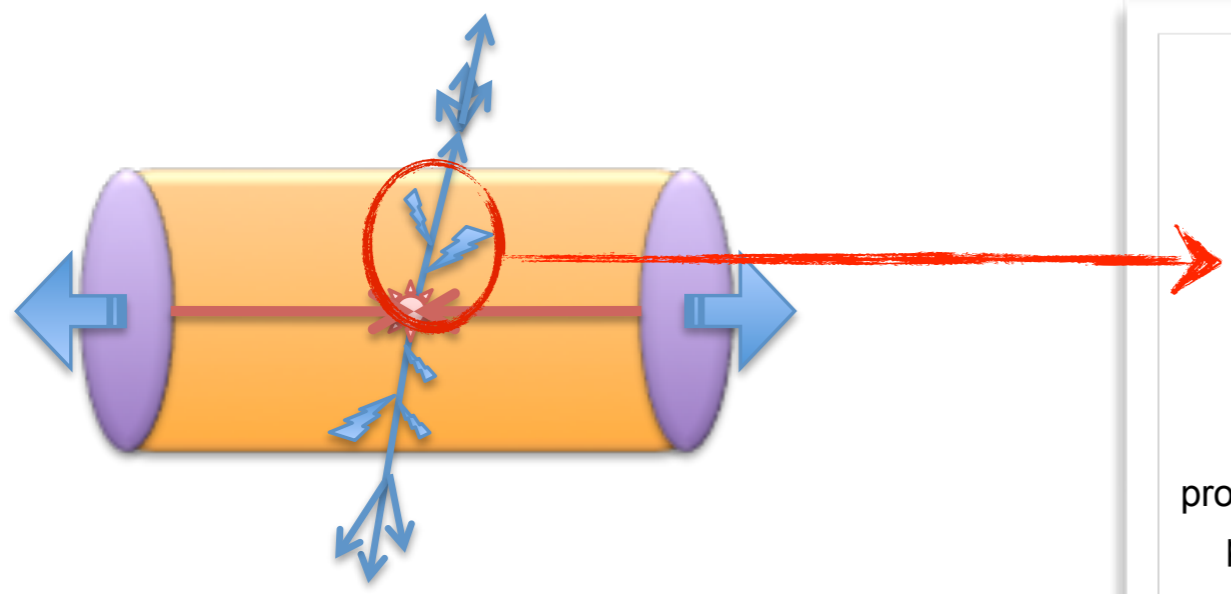


High energy limit: Radiative energy loss

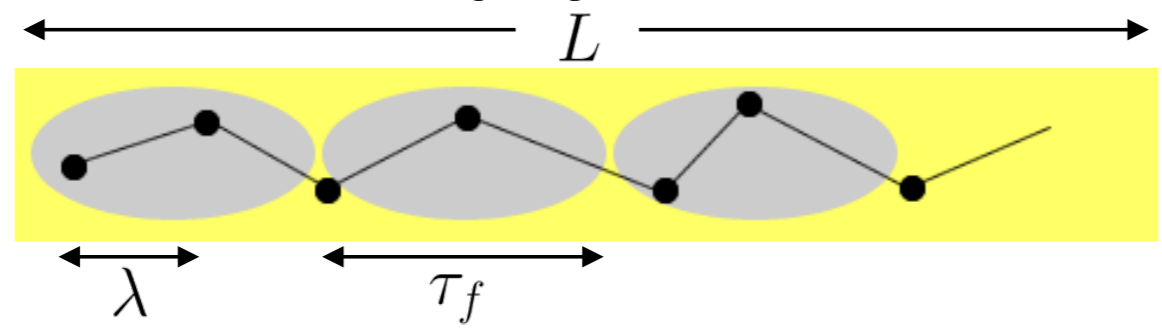
What is the equivalent in QCD?

Bremsstrahlung in QCD:

Formation time \rightarrow coherence effects



Formation time physics



$$\tau_f \sim \frac{2\omega}{k_{\perp}^2} \quad q \rightarrow qg$$

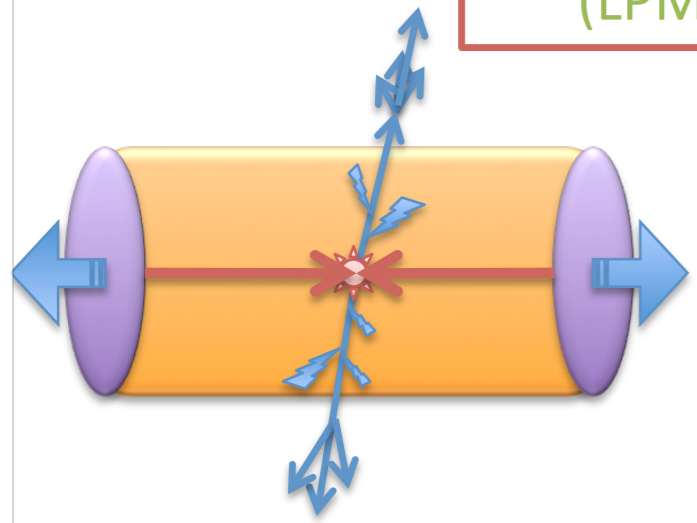
- $\tau_f < \lambda < L$ Incoherent multiple collisions
- $\lambda < \tau_f < L$ LPM effect (radiation suppressed by multiple scatterings within one coherence length)
- $\lambda < L < \tau_f$ Factorization limit (acts as one single scatterer)

Bremsstrahlung in QCD

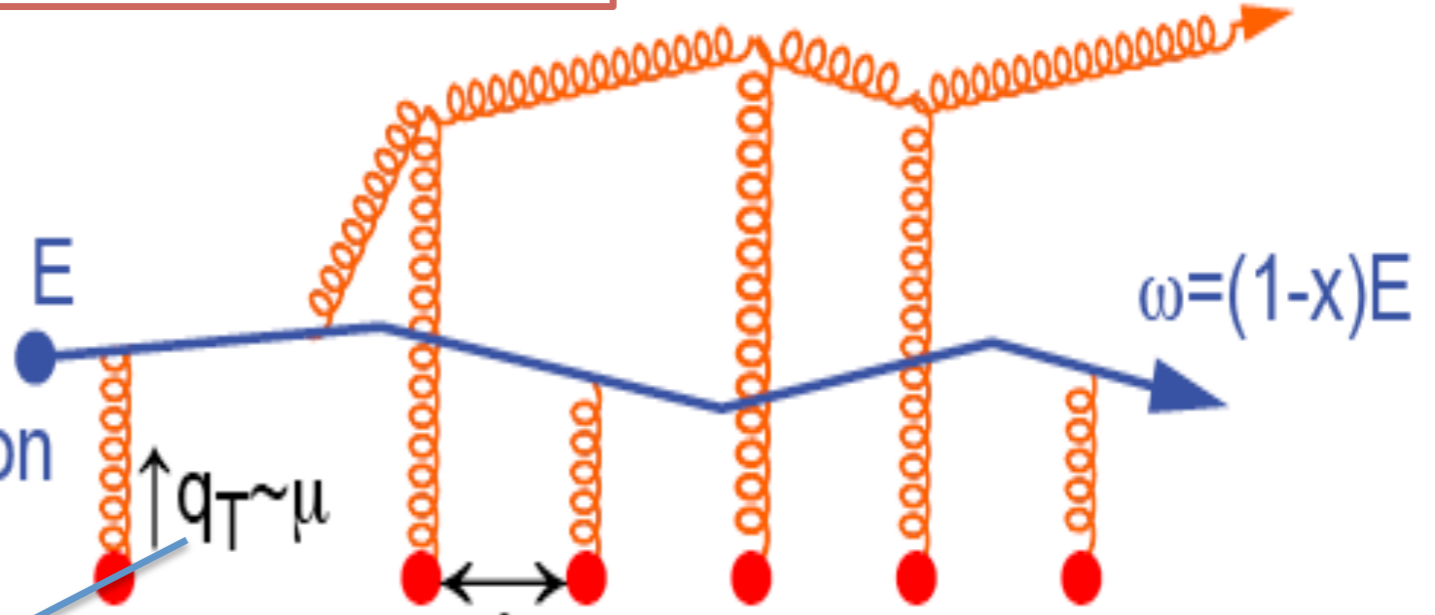
$q \rightarrow qg \quad \tau_{form} = 2\omega/k_T^2$
 $\lambda < \tau$: multiple scatterings add coherently

$t_{formation} < L \Leftrightarrow \omega < \omega_c$

High energy **color charged probe** propagating through color charged medium (LPM effect; multiple soft radiations)



Hard Production



Define a transport coefficient:

$$\hat{q} \sim \mu^2 / \lambda$$

$$-dE/dx \sim \alpha_s \hat{q} L^2$$

Partonic energy loss in QCD medium is proportional:

- to squared average path length (Note: QED ~ linear)
- to density of the medium

$$\lambda \propto \frac{1}{\rho}$$

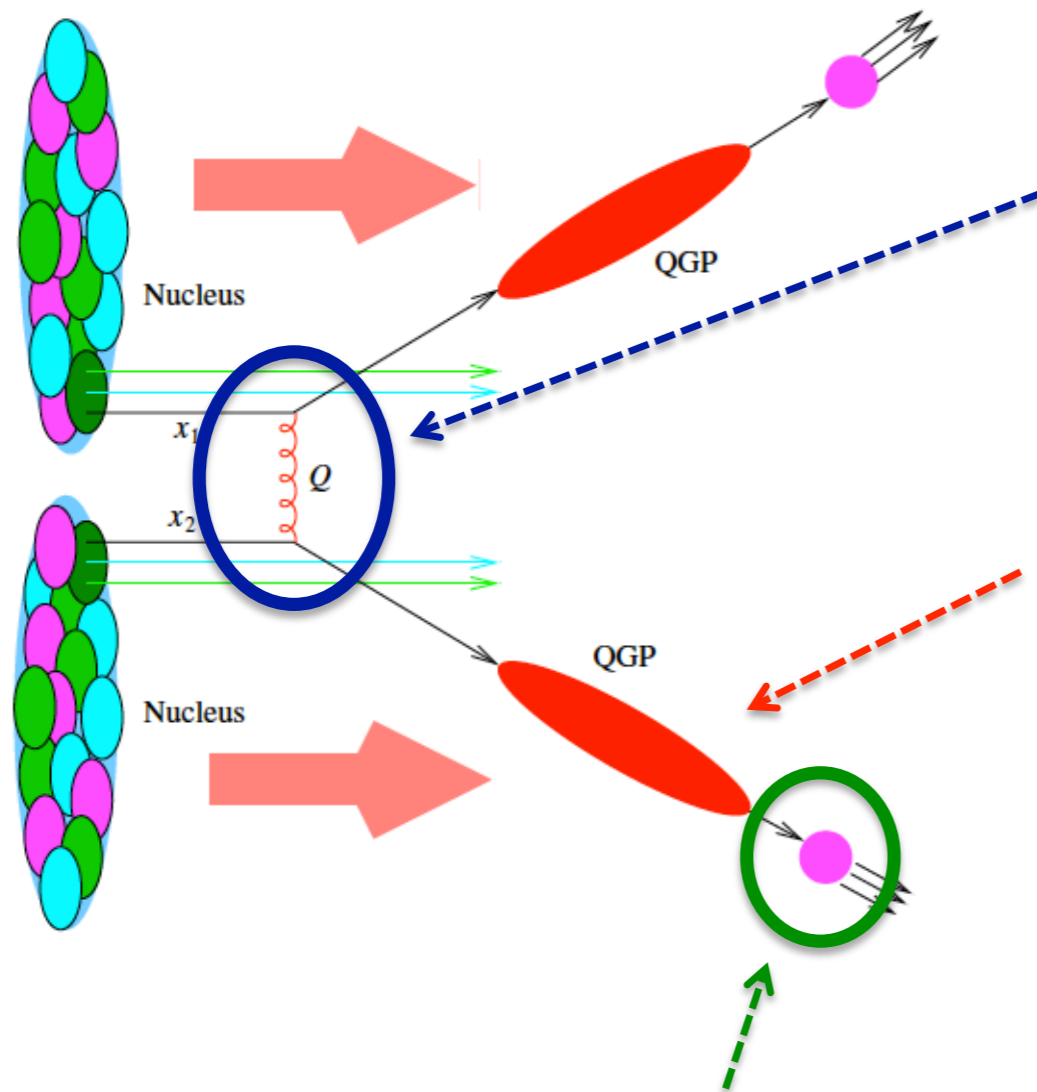
\Rightarrow energy flow (parton+radiation) modified as compared to jet in vacuum
 \Rightarrow jet "quenched" ("softened" fragmentation)

Jets in heavy-ion collisions

- an idealization

=> Factorized picture.

$$\sigma \propto f_a^{PDF} \otimes f_b^{PDF} \otimes \sigma^{hard}$$



production vertex: high Q^2
 → pQCD

Propagation in strongly coupled
 Quark Gluon Plasma

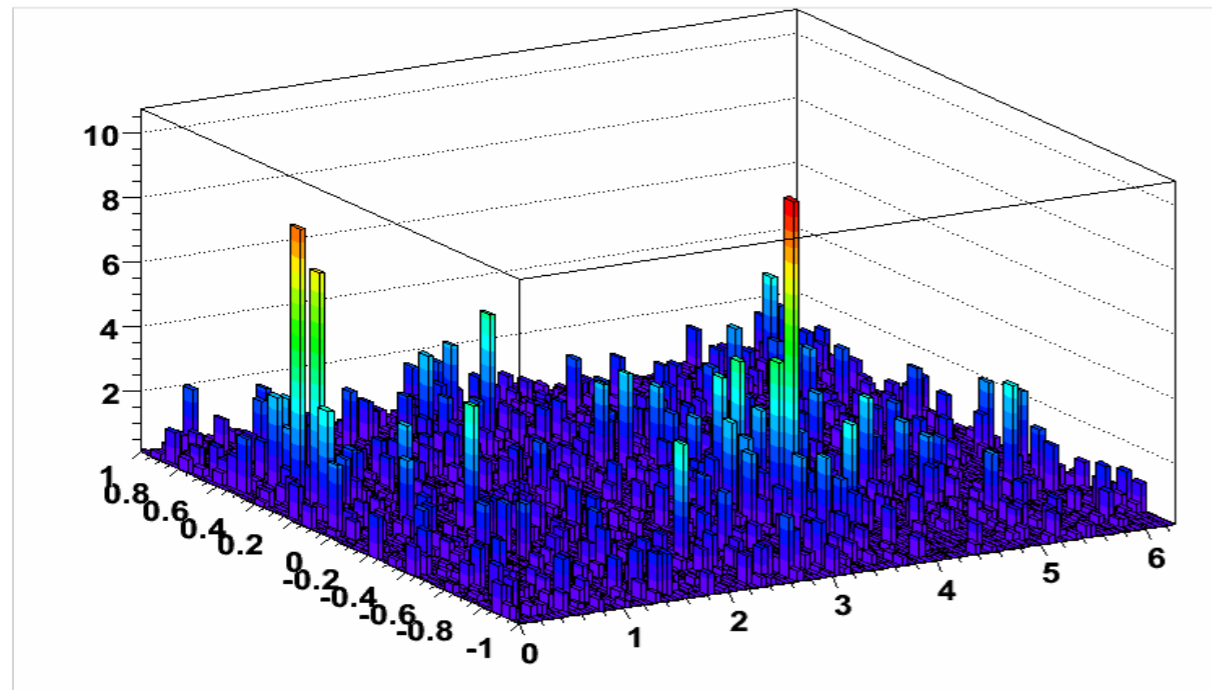
- pQCD-based jet quenching
- hydrodynamics
- AdS/CFT
- ...

Vacuum fragmentation into hadrons
 → non-pert. QCD

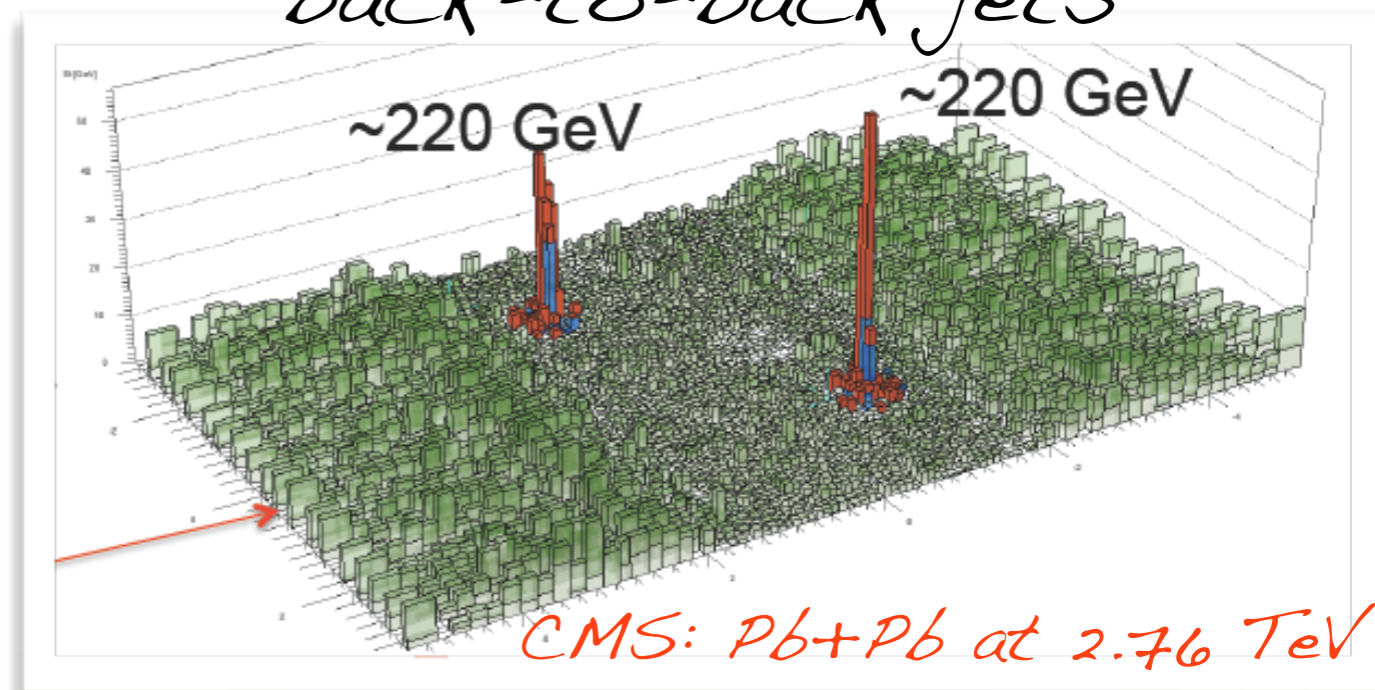
Jets in heavy-ion collisions

RHIC & LHC

STAR: Au+Au at 0.2 TeV



back-to-back jets



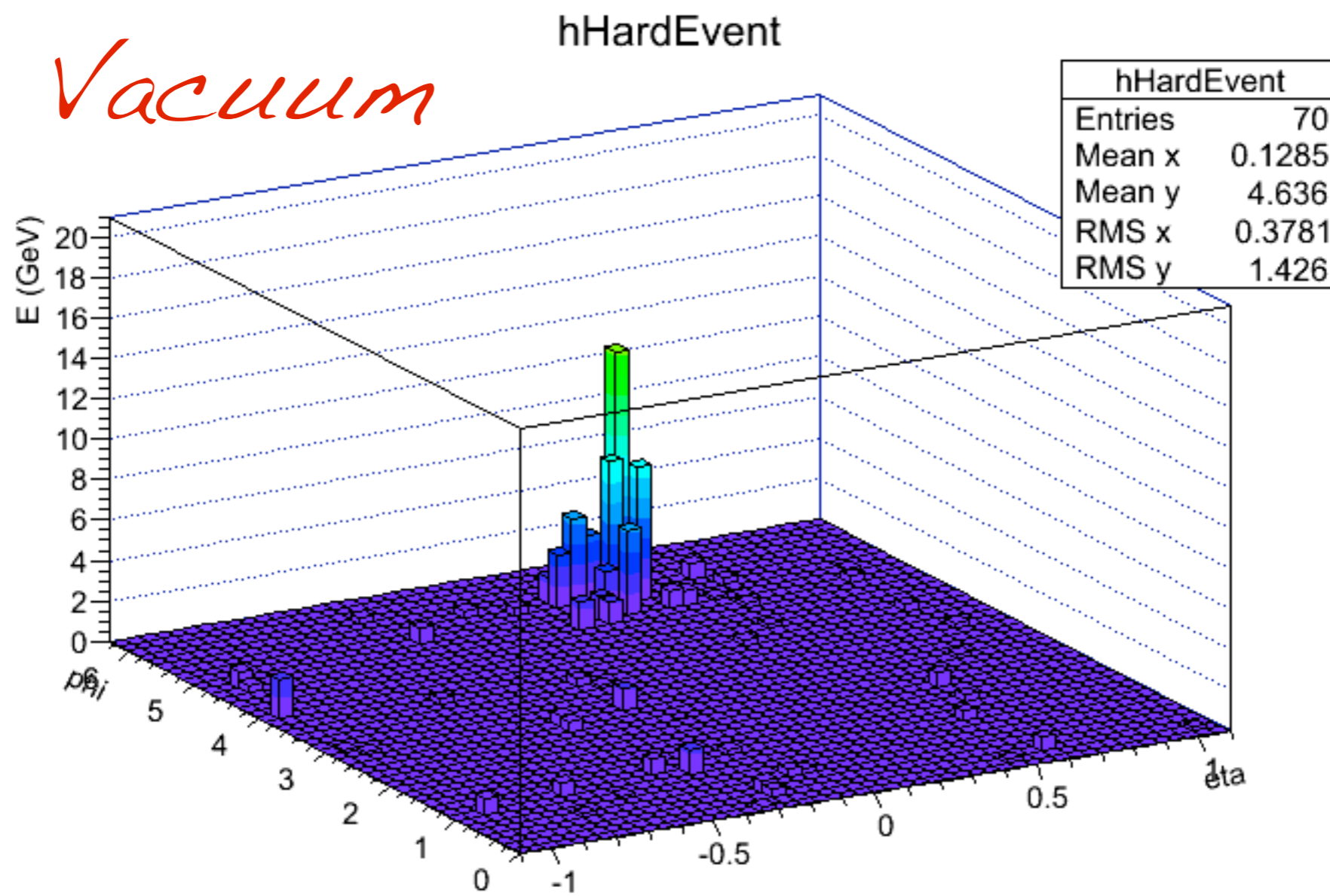
LHC + RHIC: QCD evolution of jet quenching?

Vary energy of the jet:

LHC: Vary the scale with which QGP is probed (a la DIS)

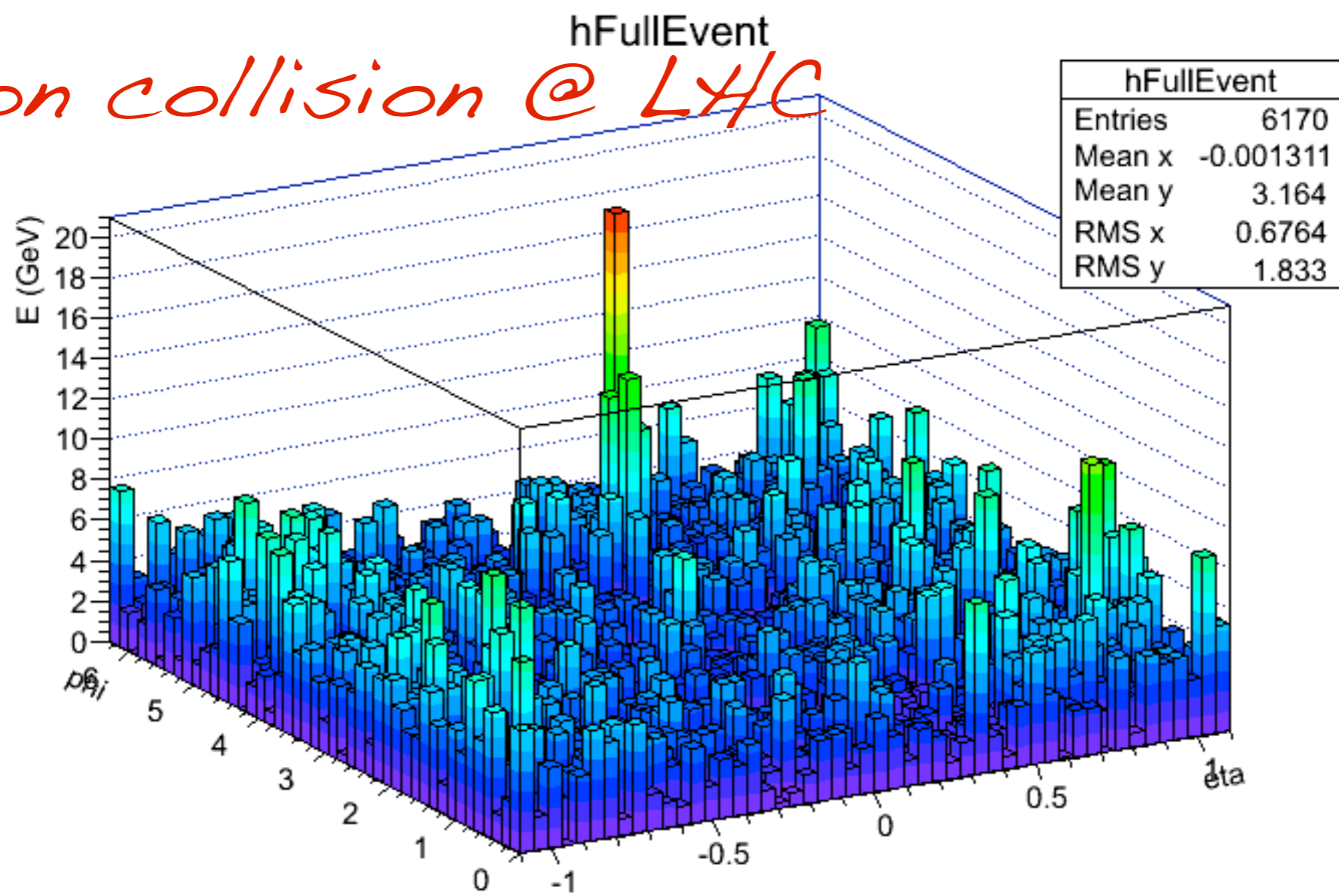
Compare and contrast RHIC and LHC

*Jets in HI collisions & Experimental difficulties:
Vacuum jet vs jet on top of the HI background...*

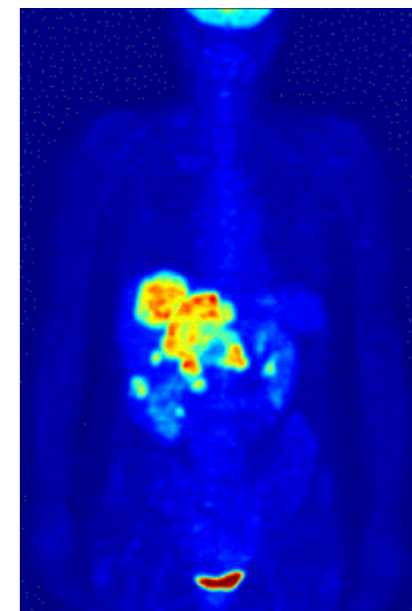


*Jets in HI collisions & Experimental difficulties:
Vacuum jet vs jet on top of the HI background...*

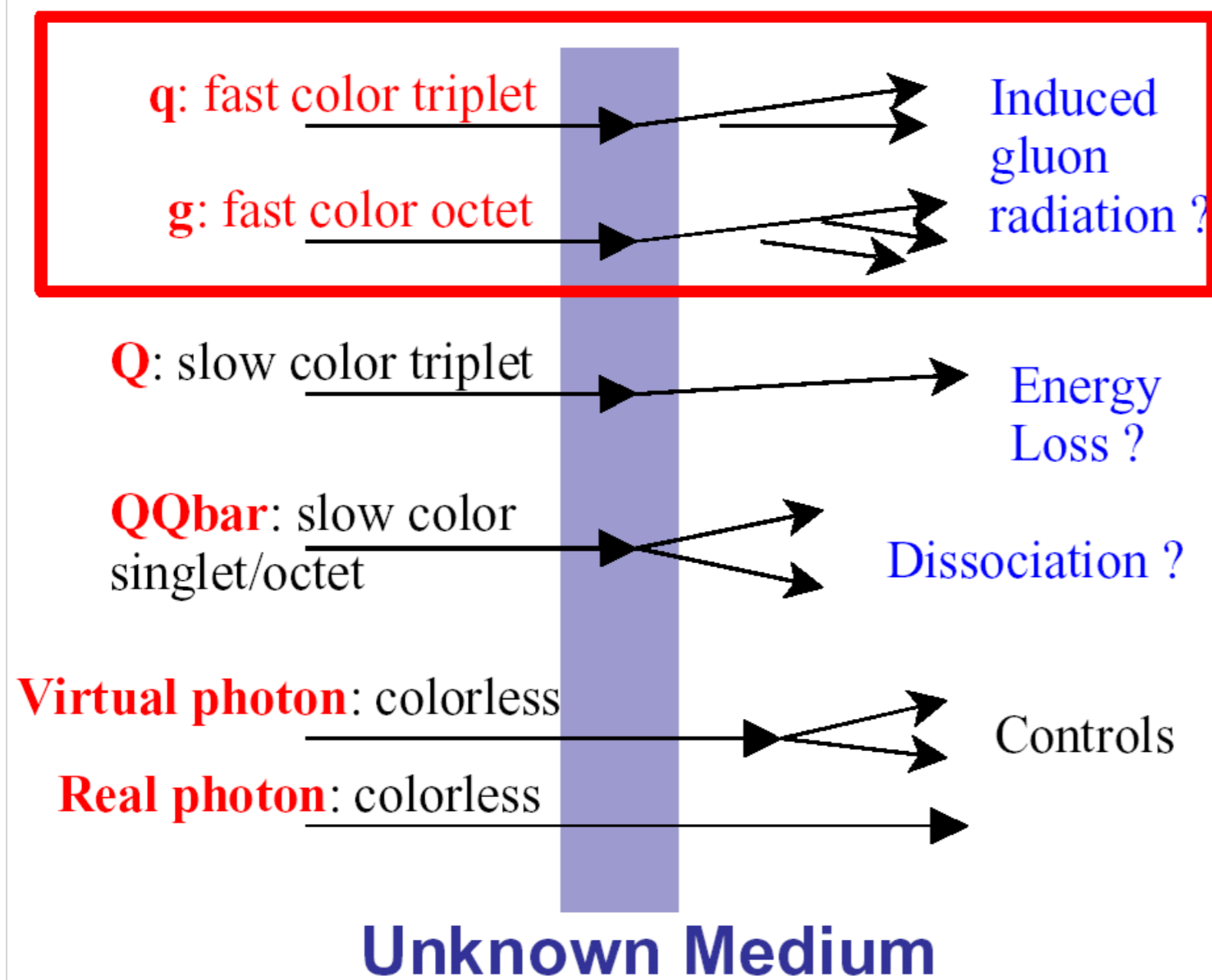
Heavy-ion collision @ LHC



Probing the unknown medium...



Human body



*jet suppression
(quenching)*

*charm/bottom
dynamics*

J/ψ & γ

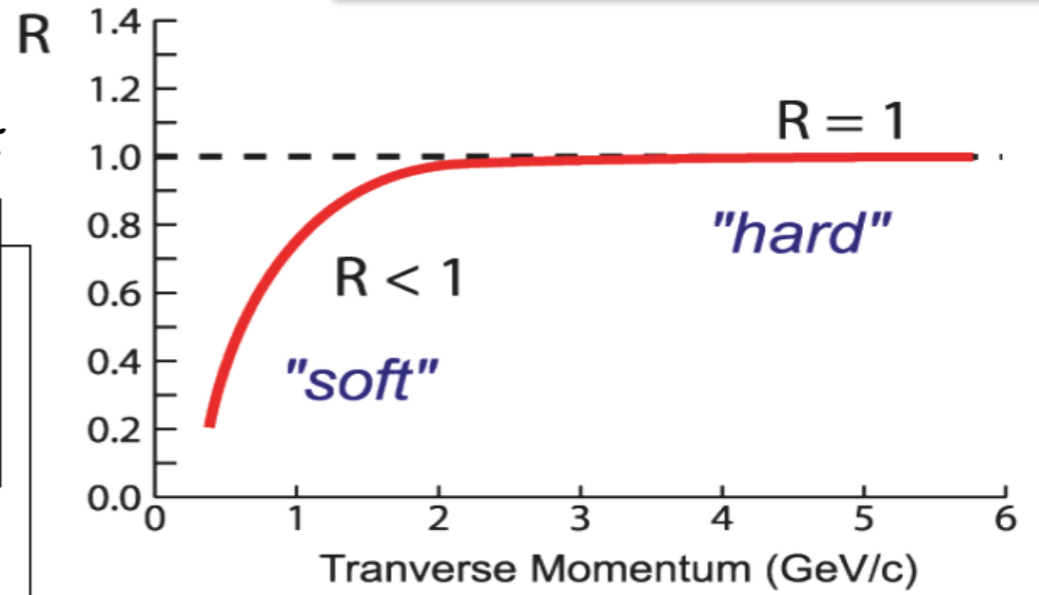
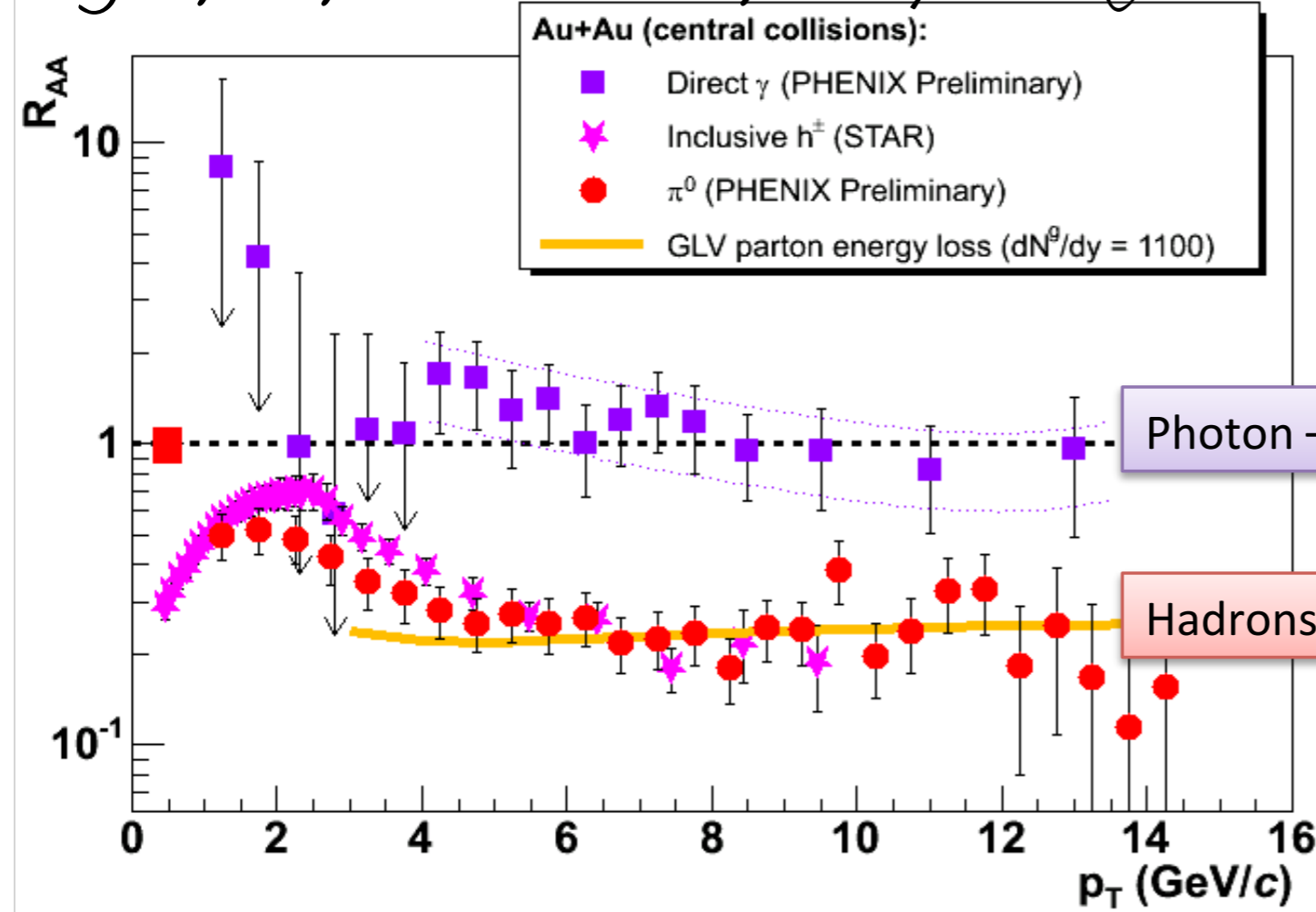
color-less particles

Jet quenching - RHIC

Ratio = $\frac{\text{\#(particles observed in AA collision per binary collision)}}{\text{\#(particles observed per p-p collision)}}$

No "effect":
 $R < 1$ at small momenta
 $R = 1$ at higher momenta where hard processes dominate

High- p_T particles - proxy for jets



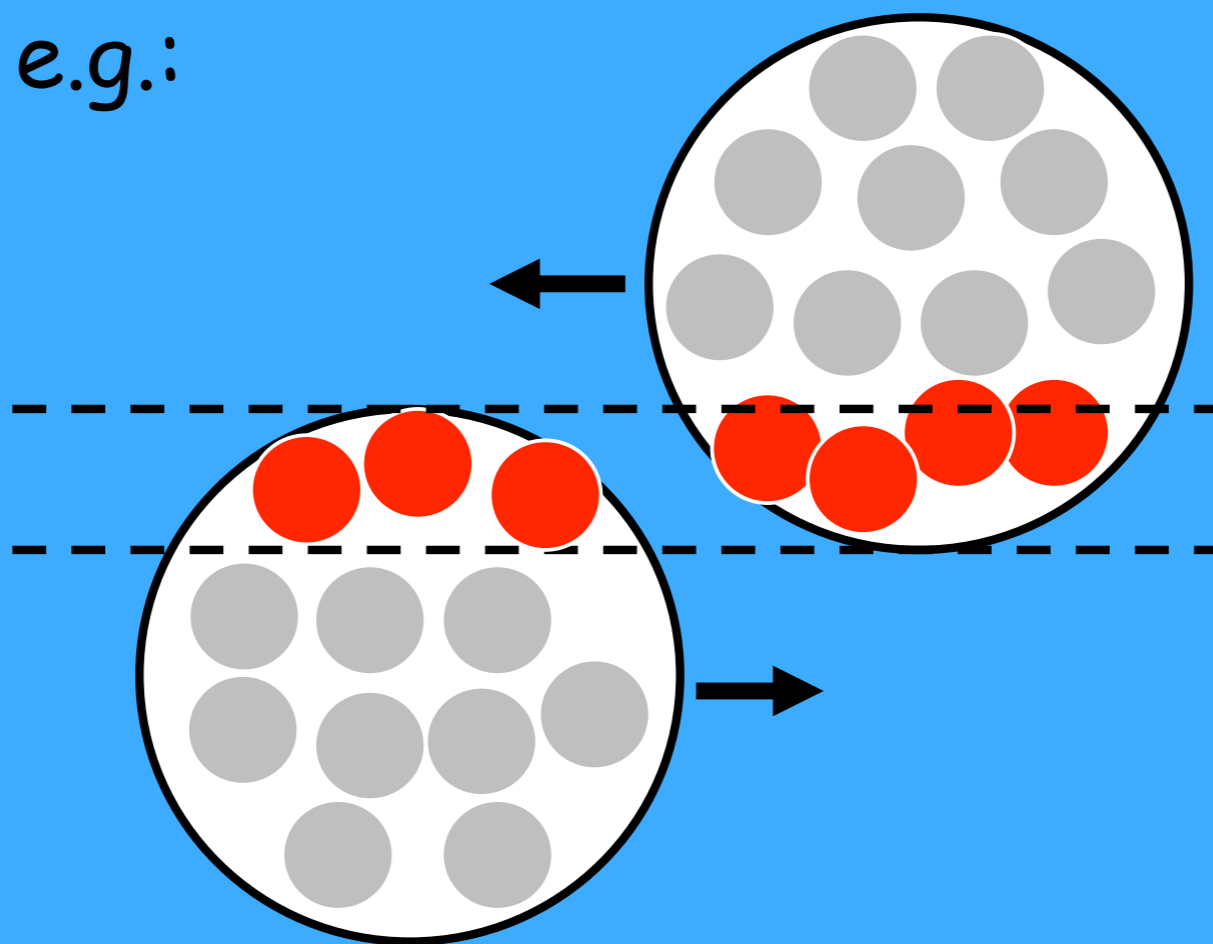
Photon - color neutral probe => No suppression

Hadrons from color charged jets => Suppression

Reminder...

"Soft", large cross-section processes expected to scale with N_{part}
 "Hard", low cross-section processes expected to scale with N_{bin}

e.g.:

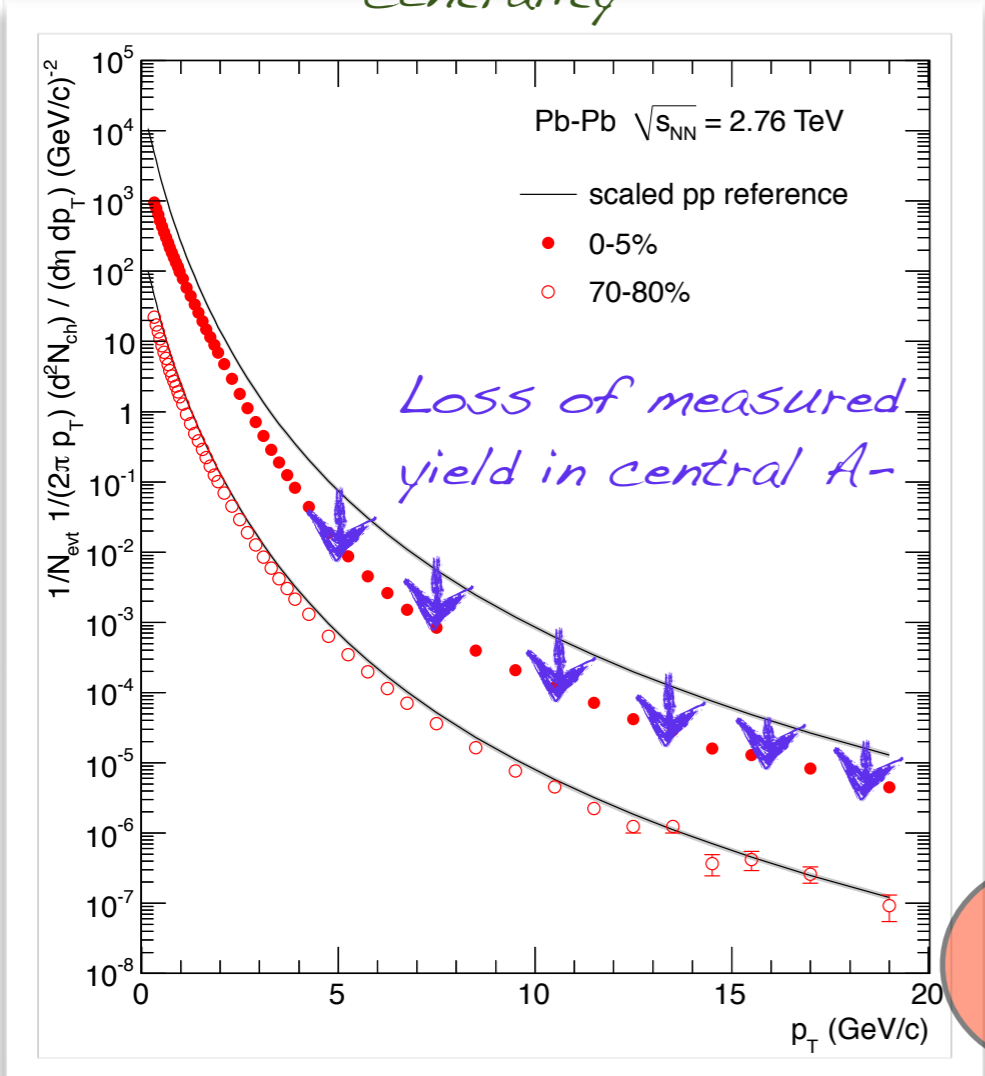


$$N_{part} \text{ (or } N_{wound}) = 7 \text{ "participants"}$$

$$N_{bin} \text{ (or } N_{coll}) = 12 \text{ "binary collisions"}$$

"Easier" (than full jet reconstruction) exercise: Jet-quenching via leading hadrons

Inclusive hadron production
Measured as a function of collision
centrality

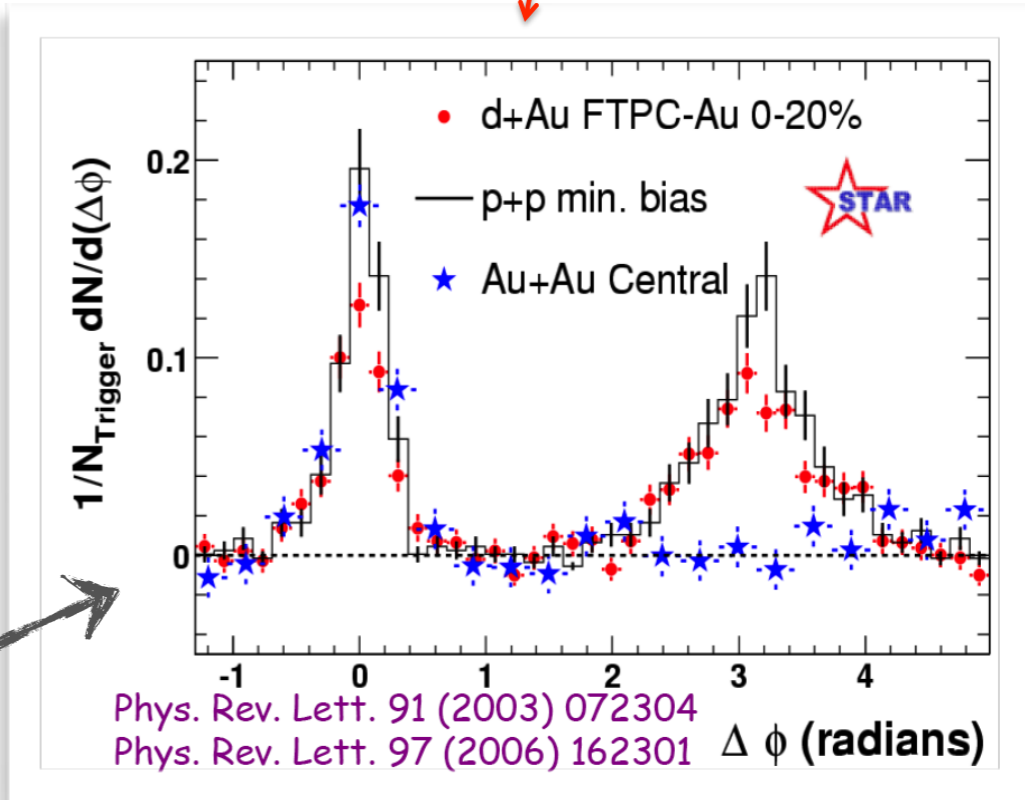
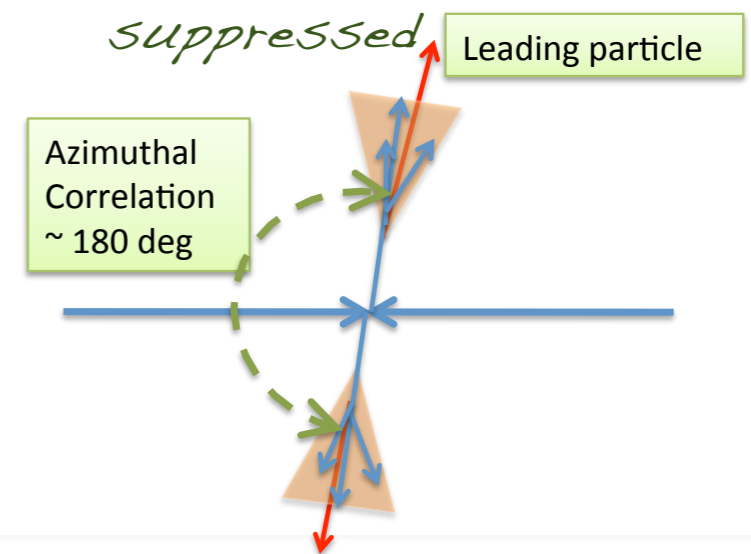


Loss of measured
yield in central A-

Note on correlations: interesting
tool to study the "intermediate" -
pT region - jets vs flow and
recombination

Di-hadron correlations

Rates of recoil ("away-side") hadrons



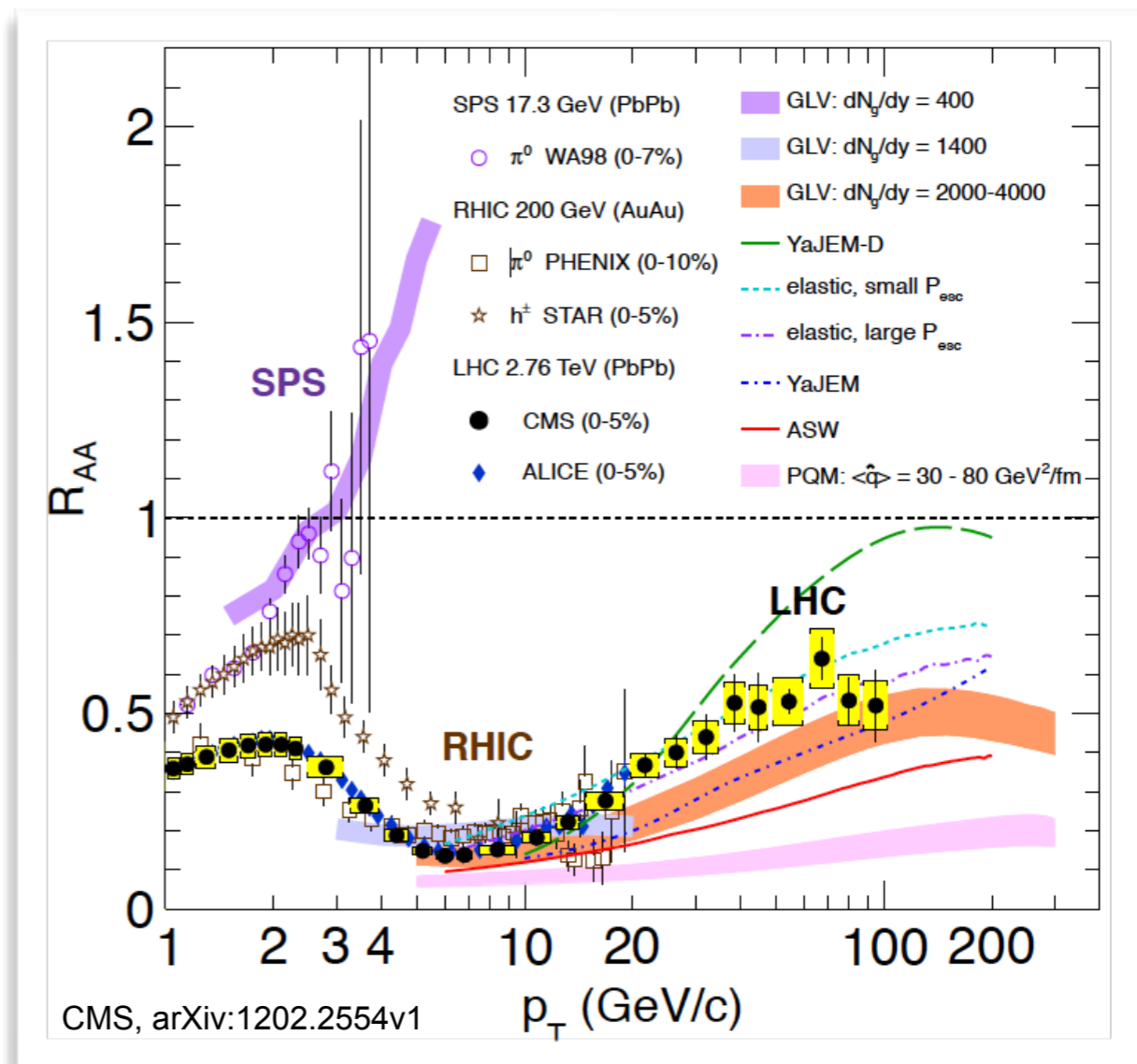
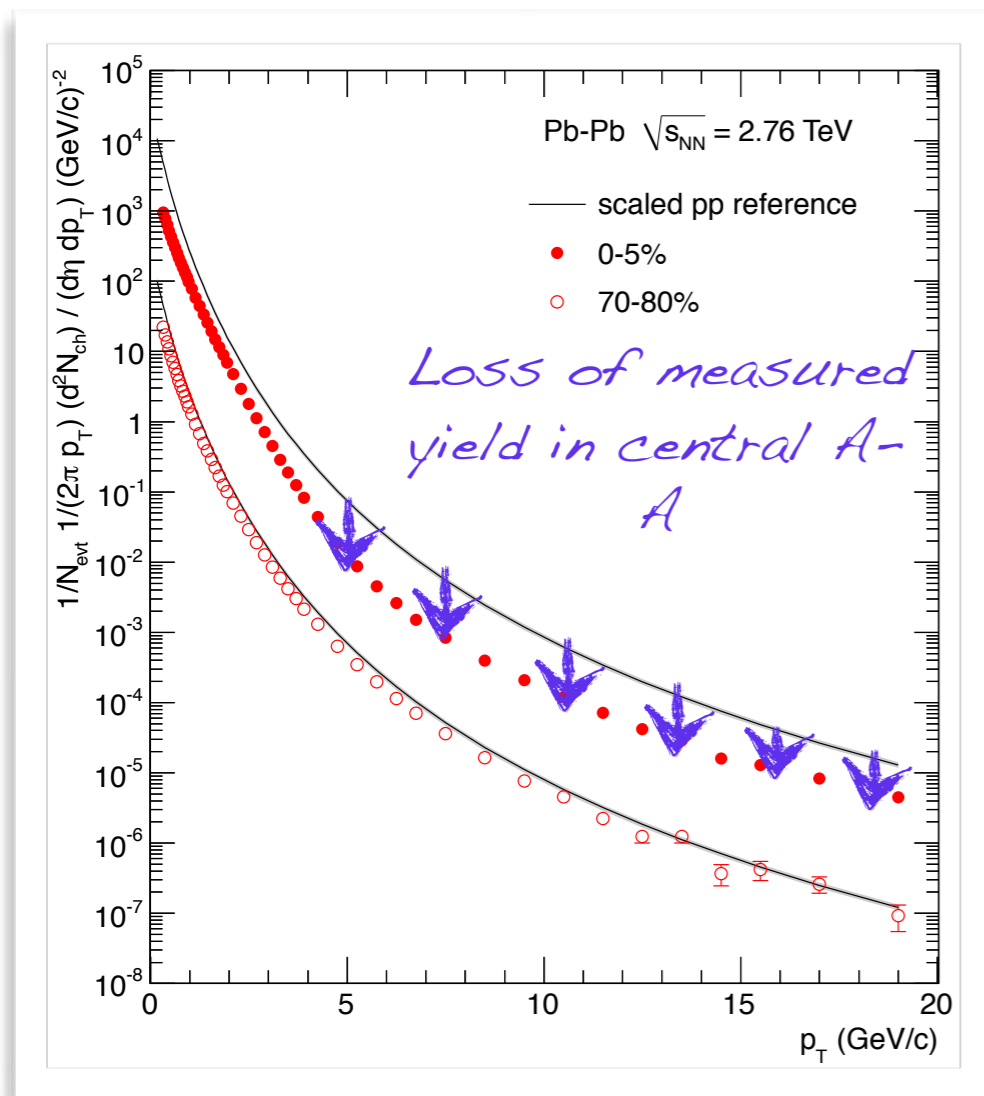
Hadron suppression

$$R_{AB} = \frac{d^2 N / dp_t d\eta}{T_{AB} d^2 \sigma^{pp} / dp_t d\eta}$$

$$T_{AB} = \langle N_{bin} \rangle / \sigma_{inel}^{pp}$$

Nuclear modification factor:

$$R_{AA} = \frac{\#(\text{particles observed in AA collision per } N\text{-}N \text{ (binary) collision})}{\#(\text{particles observed per } p\text{-}p \text{ collision})}$$



"No effect" case is for $R_{AA} = 1$ at high p_T where hard processes

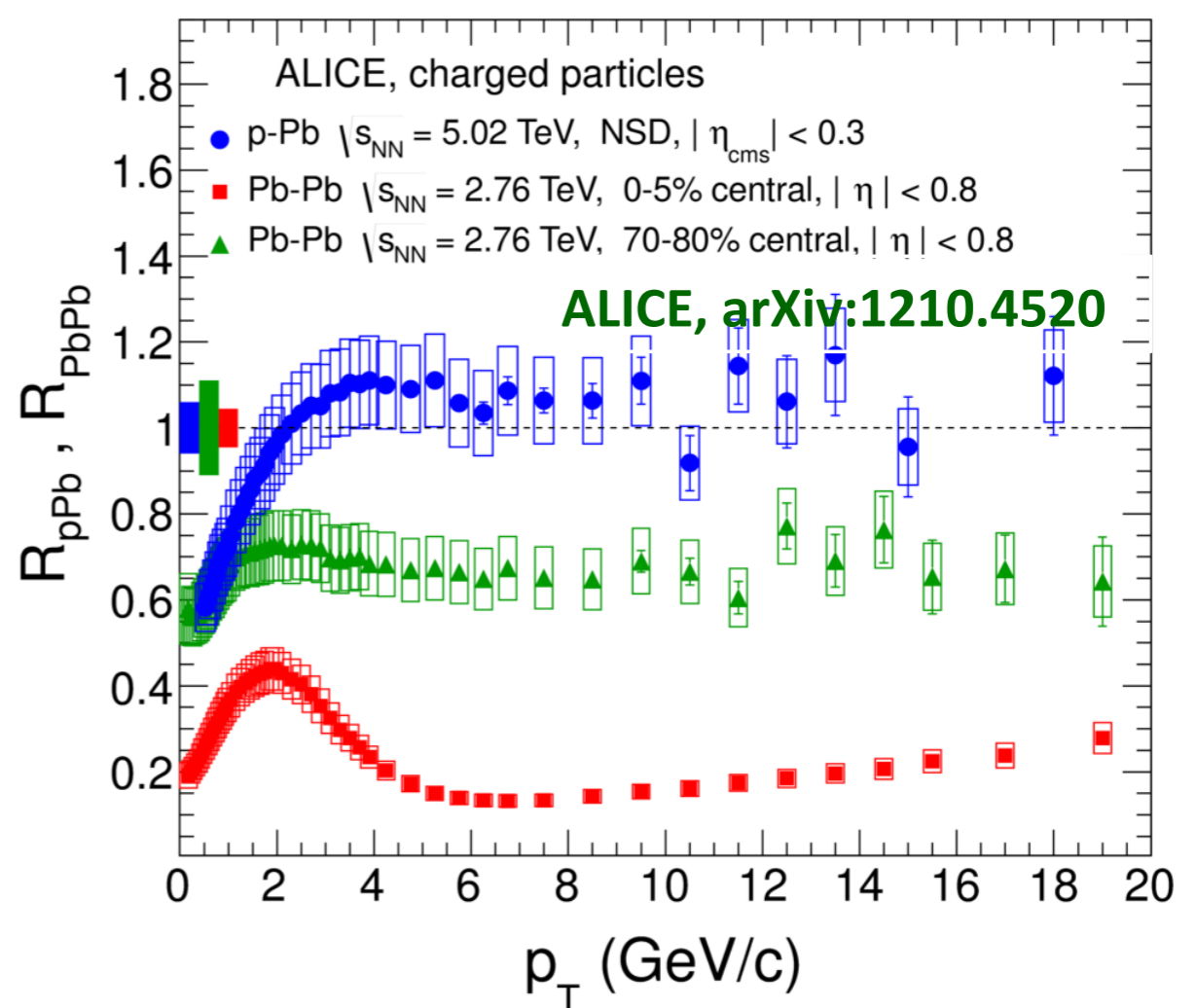
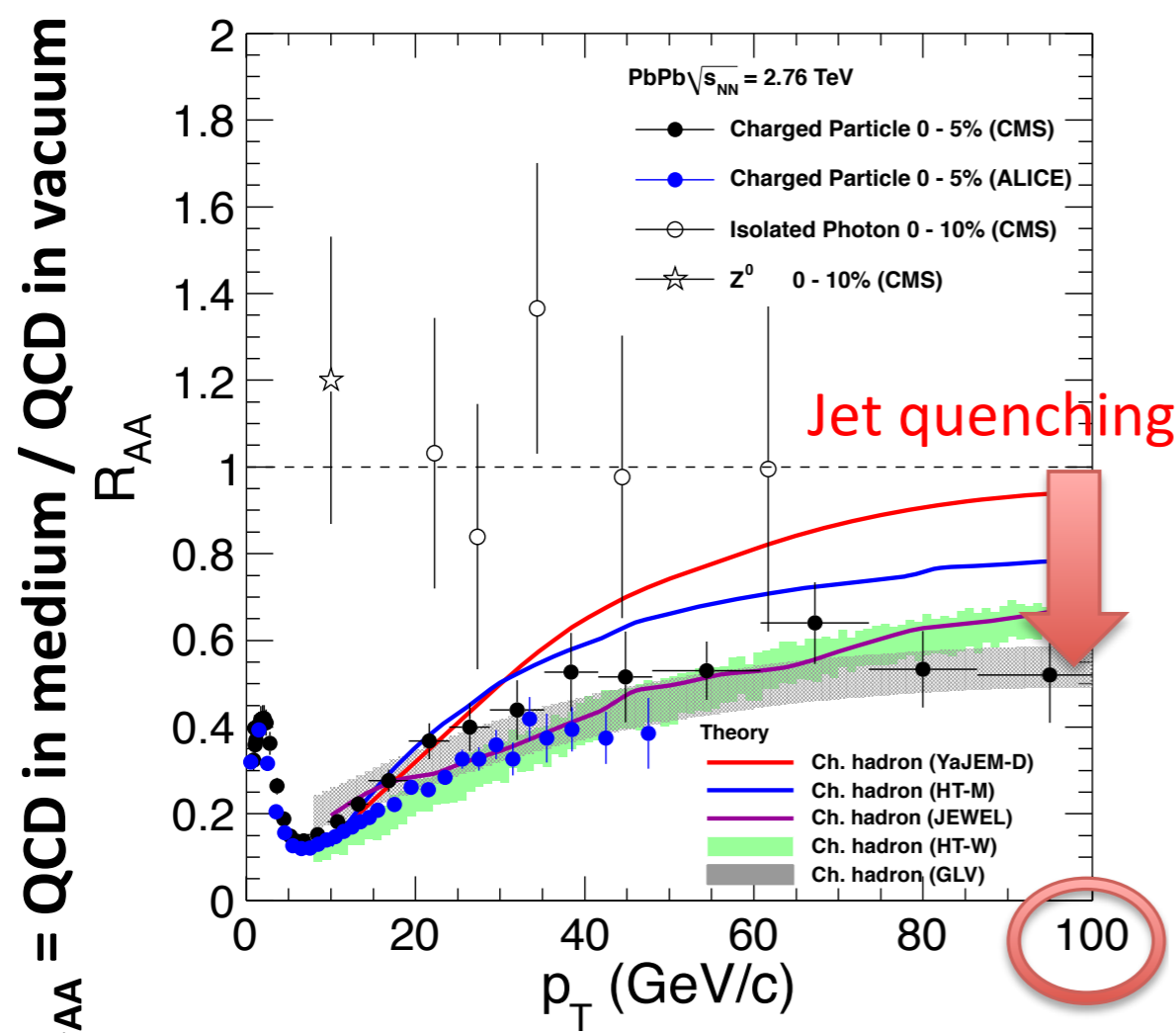
Energy-loss - QGP state effect!

Color charged probes suppressed

Color neutral probe production scales with N_{bin} collisions

pA collisions: suppression is an effect of QGP

$$R_{AA} = \frac{1}{\langle N_{coll} \rangle} \frac{dN_{AA}/dp_T}{dN_{pp}/dp_T}$$

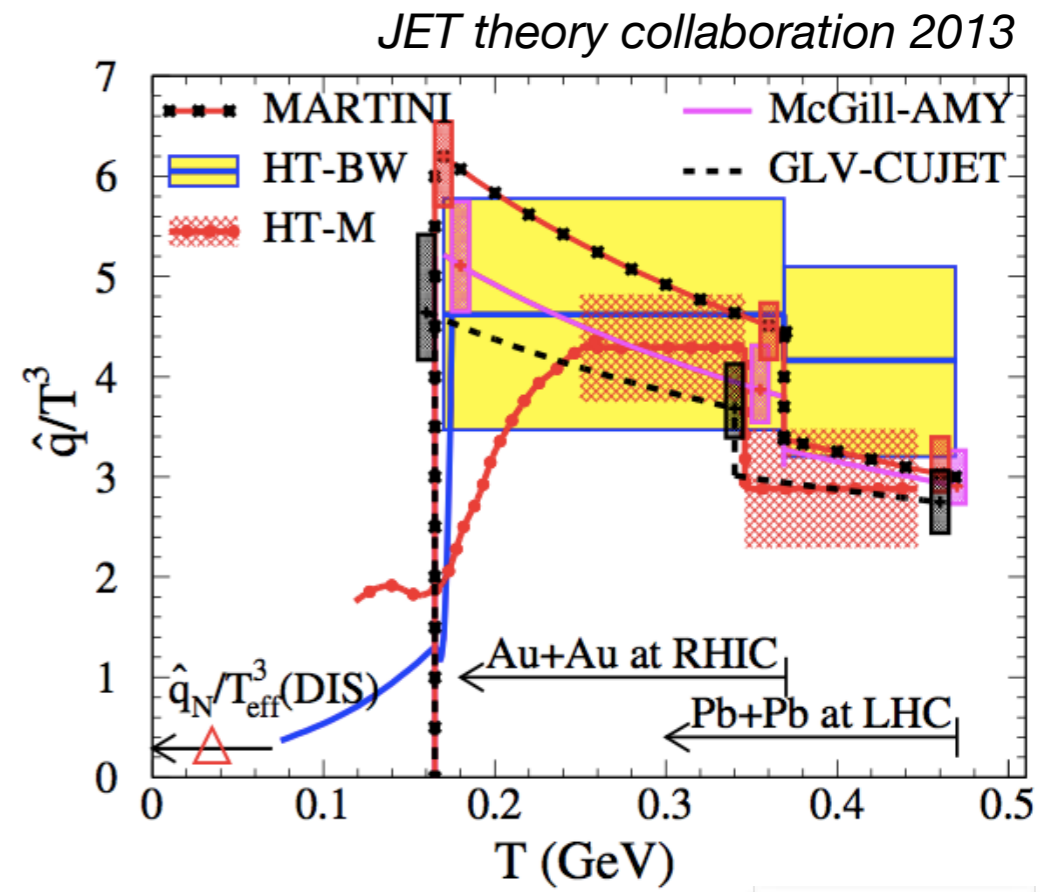
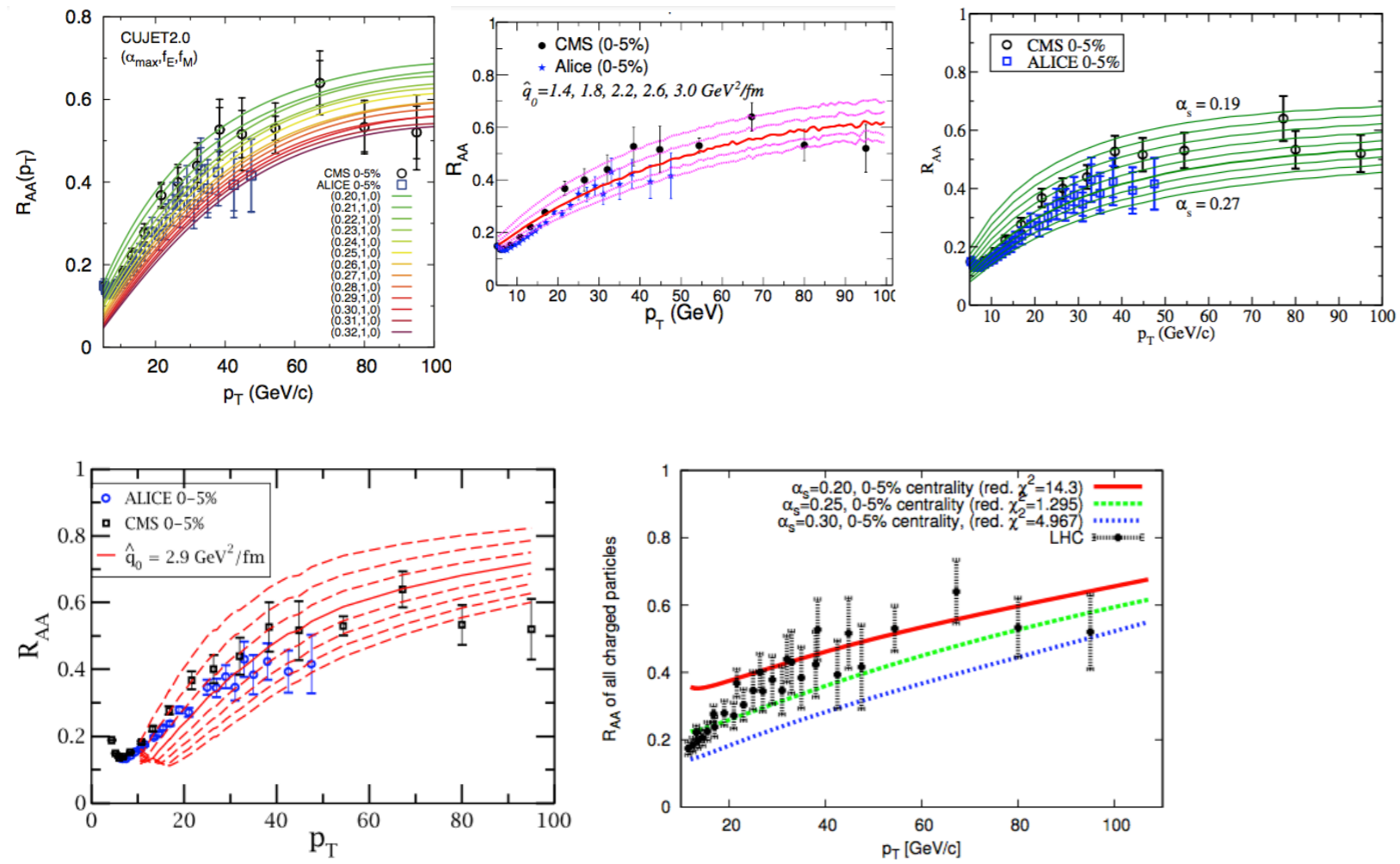


Throughout the talk: $R_{AA} = \text{QCD in medium} / \text{QCD in vacuum}$

Note: only colored probes quenched; pA: jet quenching is a in-medium effect

Extraction of QGP transport coefficients

$$-dE/dx \sim \alpha_s \hat{q} L^2$$



Systematic data-model(s) study

$$\hat{q} \sim \mu^2 / \lambda$$

$$\lambda \propto \frac{1}{\rho}$$

=> extract transport coefficient

Use of RHIC & LHC data

RHIC : $\hat{q} \approx 1.2 \pm 0.3 \text{ GeV}^2/\text{fm}$

LHC : $\hat{q} \approx 1.9 \pm 0.7 \text{ GeV}^2/\text{fm}$

Temperature dependence (?)

Cold matter (HERMES DIS) : $\hat{q} \approx 0.02 \text{ GeV}^2/\text{fm}$

so far...

- High energy heavy-ion collisions:
 - hot medium; thermal [statistical hadronization] particle emission
 - QGP flows as almost perfect fluid - well described by viscous hydrodynamics - transport coefficient constrained
 - slow evolution with energy - similar medium over large span of energies (?) - role of mean free path vs. medium size...
 - QGP is opaque to high energy jets (dense medium) - transport coefficient constrained (first syst. results)