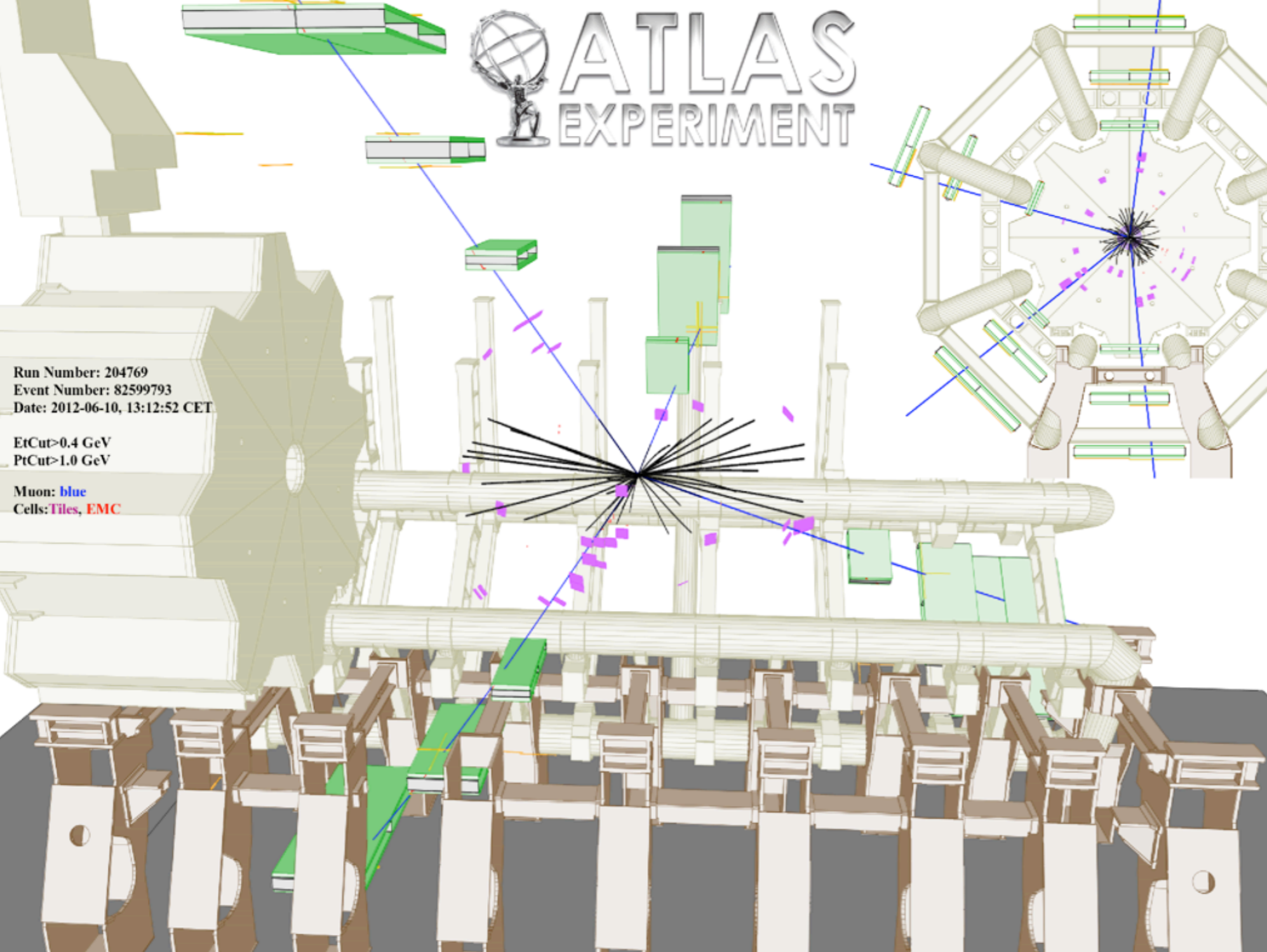


ATLAS EXPERIMENT

Run Number: 204769
Event Number: 82599793
Date: 2012-06-10, 13:12:52 CET

EtCut>0.4 GeV
PtCut>1.0 GeV

Muon: blue
Cells: Tiles, EMC



Higgs boson Rare Decays

K. Nikolopoulos
University of Birmingham

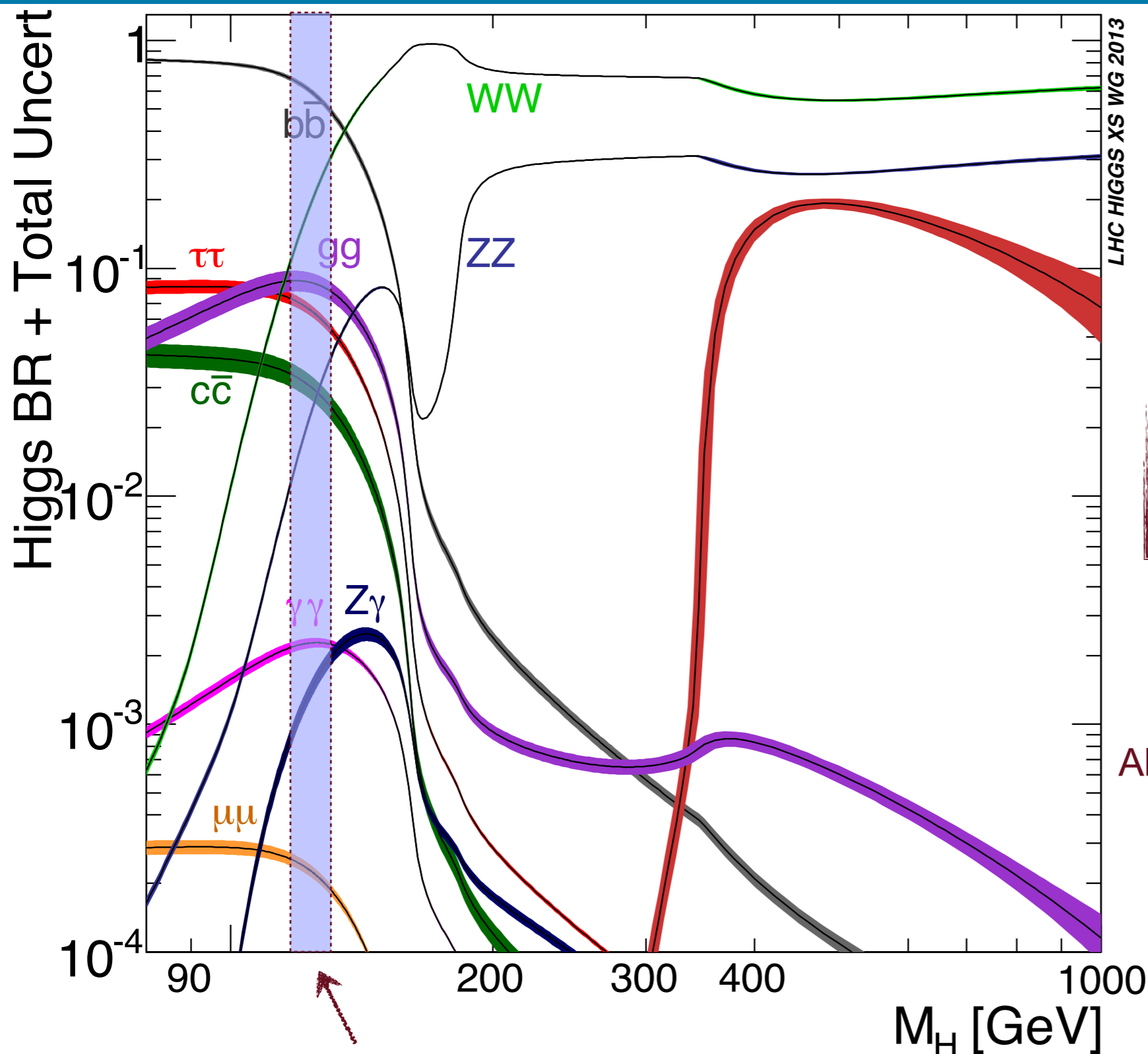
LHC Higgs Cross Section WG
WG3 Rare Decays
22nd May 2015, Fermilab, USA



UNIVERSITY OF
BIRMINGHAM



Standard Model Higgs boson decays



The talk title is a bit too wide:
 → Many “rare” Higgs boson decays

Focus on exclusive Higgs boson decays that carry information on (light) fermion Yukawa couplings

Concentrate on ATLAS results, but when relevant/available CMS plots also shown

Also relevant theory results presented

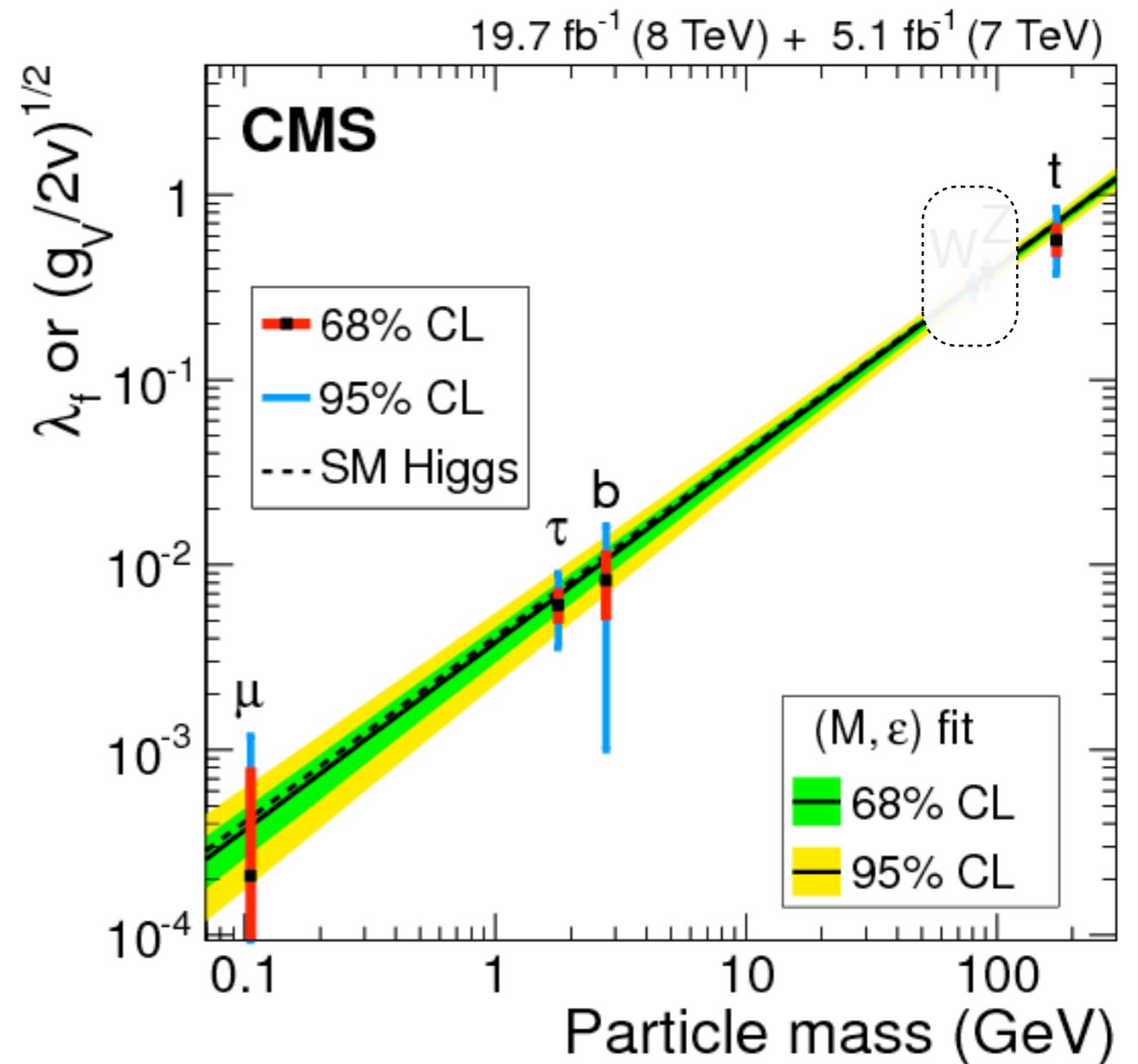
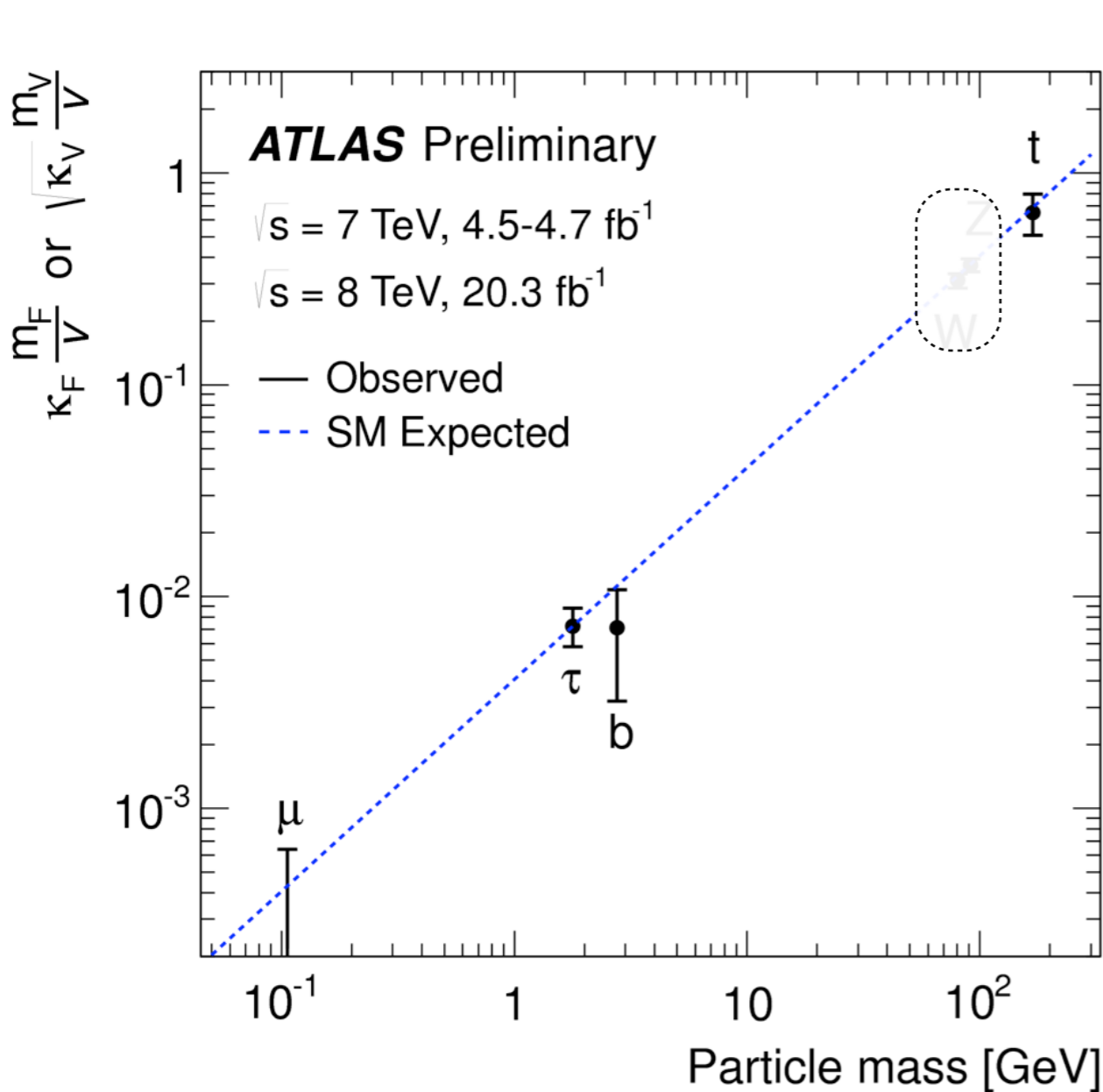
As suggested, will not focus on the results, but try to give an idea of the experiment/theory issues

$m_h \sim 125$ GeV gives access to several decay channels

Gauge bosons: $\gamma\gamma$, ZZ^* , WW^* , $Z\gamma$

Fermions: bb , $\tau\tau$, $\mu\mu$

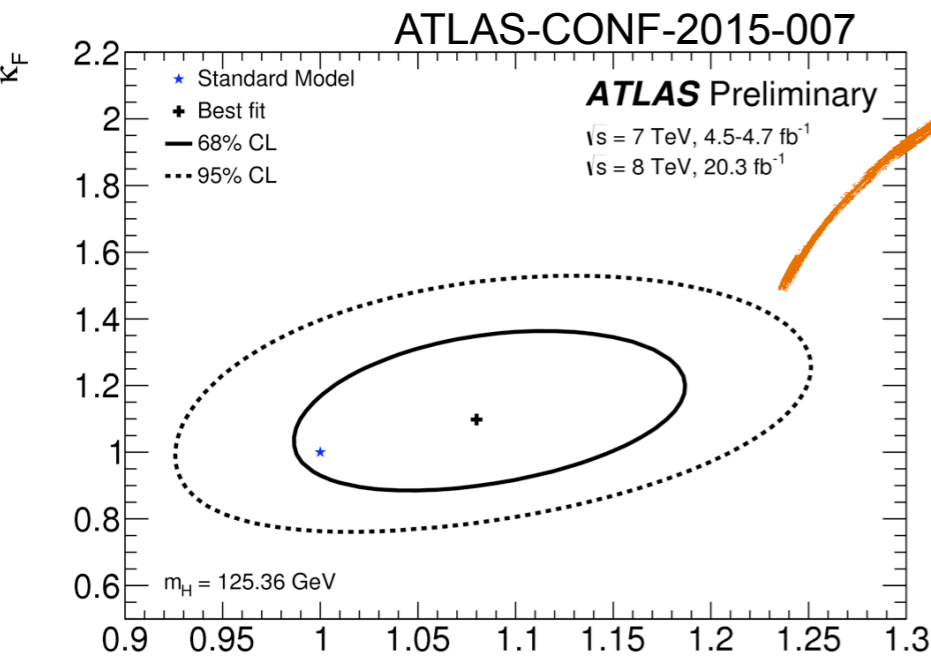
Yukawa couplings so far...



- Currently, with the exception of $h \rightarrow \tau\tau$, no conclusive direct evidence for $h \rightarrow f\bar{f}$
- Indications for $h \rightarrow b\bar{b}$ and $t\bar{t}h$, to be followed up in LHC Run II
- Indirectly; Higgs boson should be coupling to top-quark in the gluon fusion loop

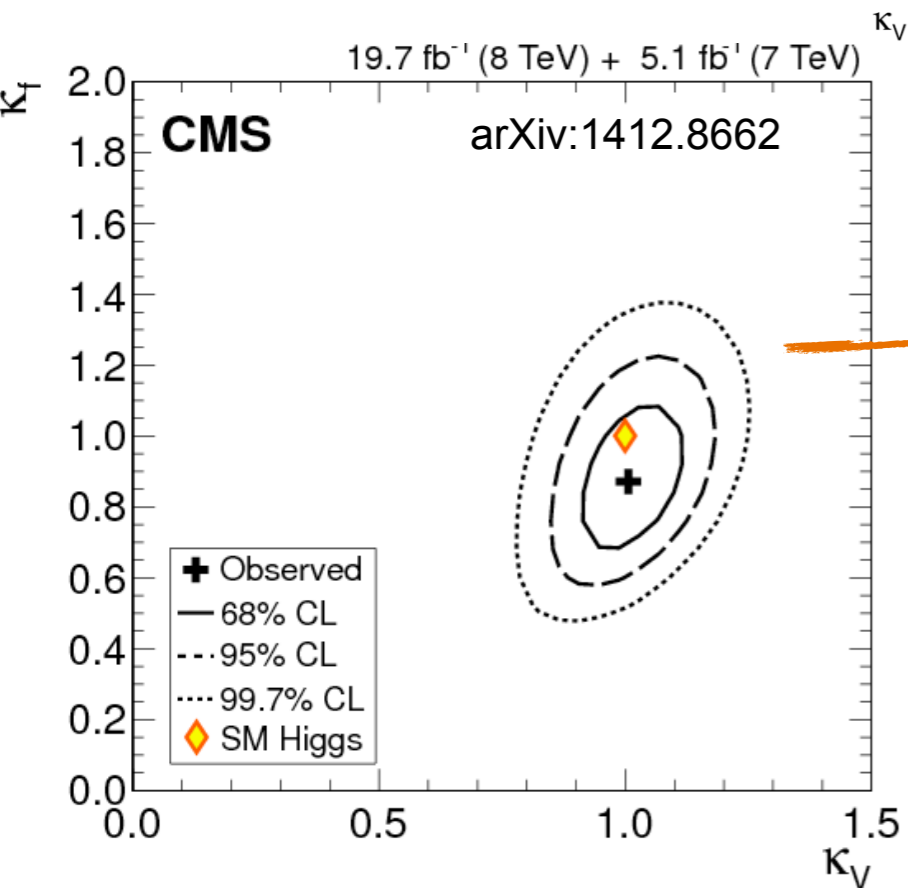
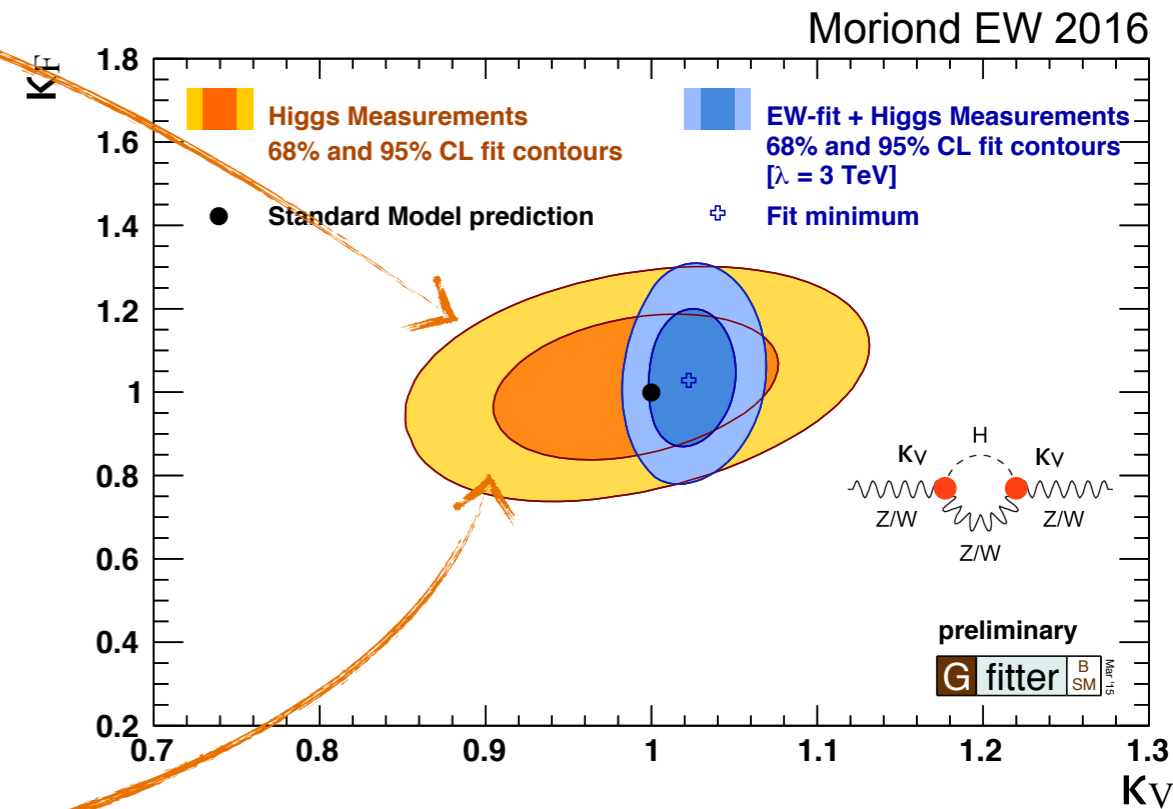
Higgs boson and precision electroweak observables

Common coupling scaling for all Fermions (κ_F) and for all Bosons (κ_V); no BSM contributions



Higgs coupling measurements:

- ▶ $\kappa_V = 0.99 \pm 0.08$
- ▶ $\kappa_F = 1.01 \pm 0.17$
- ▶ **Combined result:**
- ▶ $\kappa_V = 1.03 \pm 0.02$
($\lambda = 3 \text{ TeV}$)
- ▶ implies NP-scale of $\Lambda \geq 13 \text{ TeV}$



- Global EW fit more precise for κ_V than coupling measurements
- $\kappa_V > 1$ preferred (many BSM scenarios require $\kappa_V < 1$)
- Global EW fit has ~no effect on determination of κ_F

Experimental information on Yukawa couplings essential to fully characterize the observed Higgs boson!



Roman Kogler

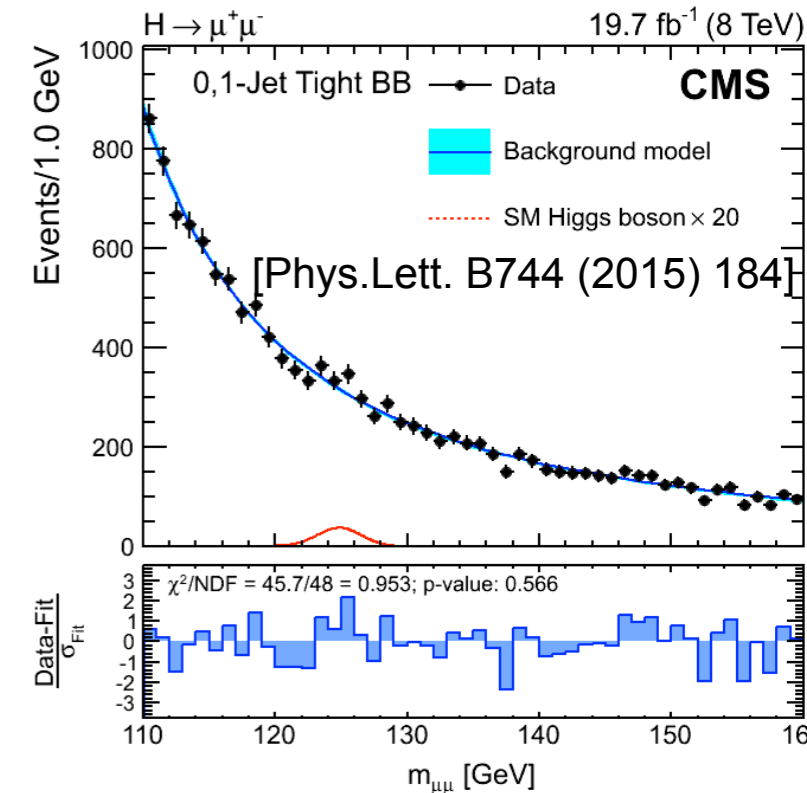
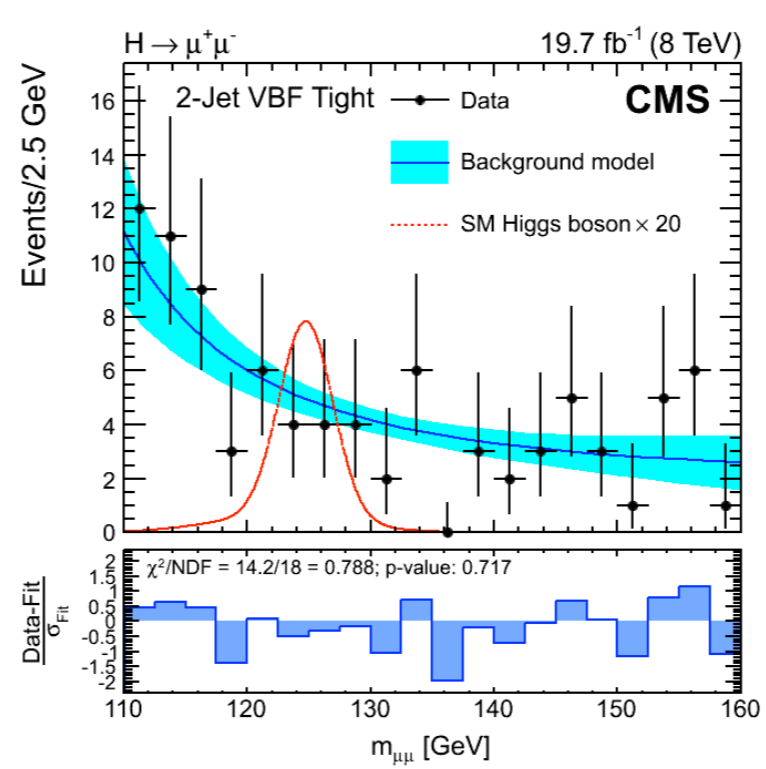
16

The global electroweak fit

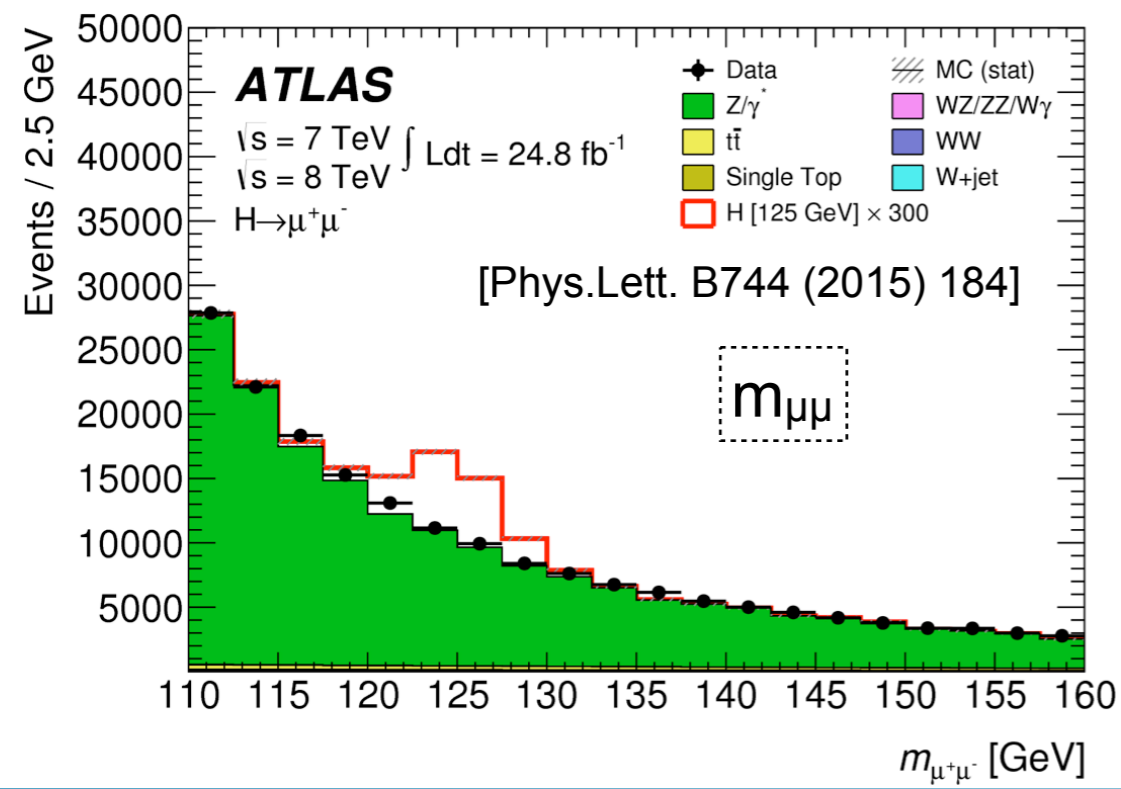


$h \rightarrow \mu\mu$

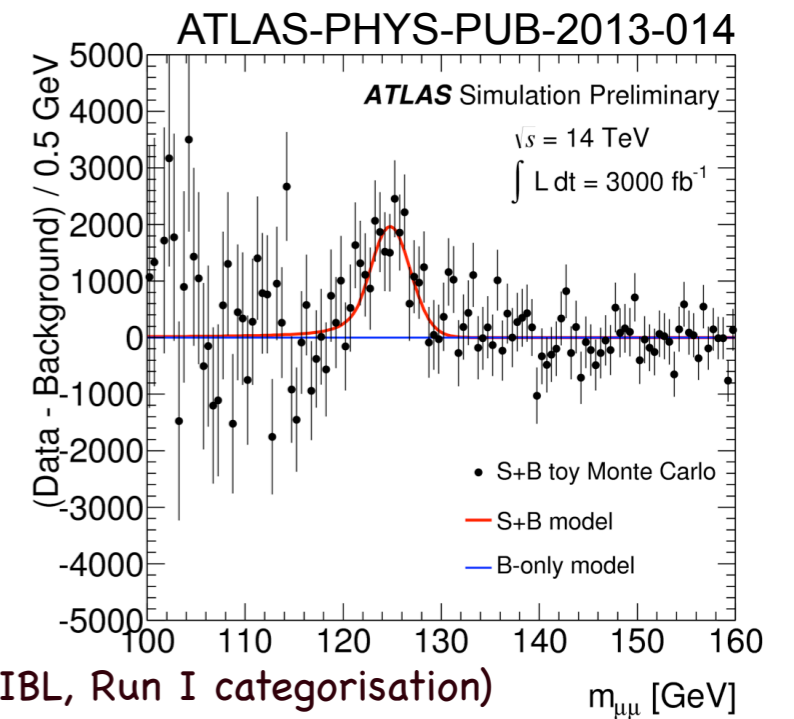
- Attainable probe for 2nd gen. Yukawa
- $BR_{SM} \sim 2 \cdot 10^{-4} (125\text{GeV}); S/B \sim 0.2\%$
- Simple Final State**
(ATLAS analysis)
- $\mu^+\mu^-$ ($p_T > 25, 15 \text{ GeV}, p_{T_{\mu\mu}} > 15 \text{ GeV}$)
- backgrounds: $Z/\gamma^* \rightarrow \mu\mu$, top, dibosons
- Categorization: central/non-central muons
- Parametric background Model: BW+Expo
- 95% CL upper limit @ $m_H = 125 \text{ GeV}$:
ATLAS : 9.8 (8.2)xSM
CMS : 7.4 (6.5)xSM



no universal Higgs coupling to fermions



$\mathcal{L} [\text{fb}^{-1}]$	300	3000
N_{ggH}	1510	15100
N_{VBF}	125	1250
N_{WH}	45	450
N_{ZH}	27	270
N_{ttH}	18	180
N_{Bkg}	564000	5640000
$\Delta_{Bkg}^{sys} \text{ (model)}$	68	110
$\Delta_{Bkg}^{sys} \text{ (fit)}$	190	620
Δ_{S+B}^{stat}	750	2380
Signal significance	2.3σ	7.0σ
$\Delta\mu/\mu$	46%	21%



conservative extrapolation (no IBL, Run I categorisation)

Room for analysis improvements.
Will benefit from detector upgrades.



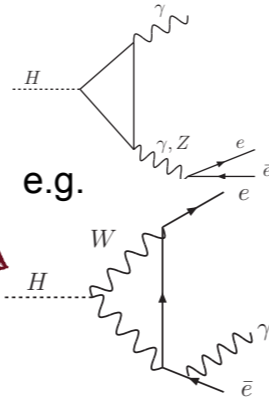
h → ee

Very rare decay of the SM Higgs BR ~ 5 · 10⁻⁹

["direct" process non-Yukawa suppressed contributions need to be included Phys.Lett. B727 (2013) 424]

- BR_{SM}(h → ττ)/BR_{SM}(h → ee) ~ 1.3 · 10⁷

Not promising to constrain the electron SM Yukawa coupling meaningfully



Performed by CMS, together with h → μμ

- simple final state:

electron-positron pair, p_T > 25 GeV

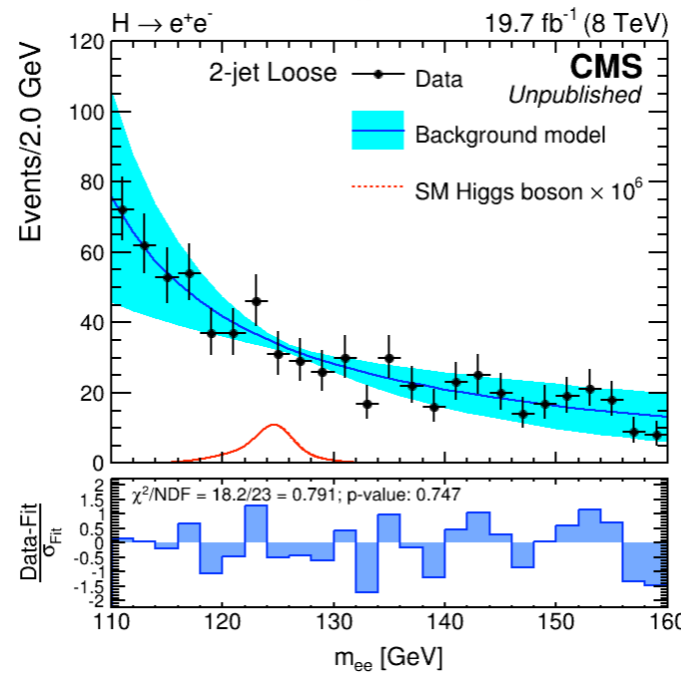
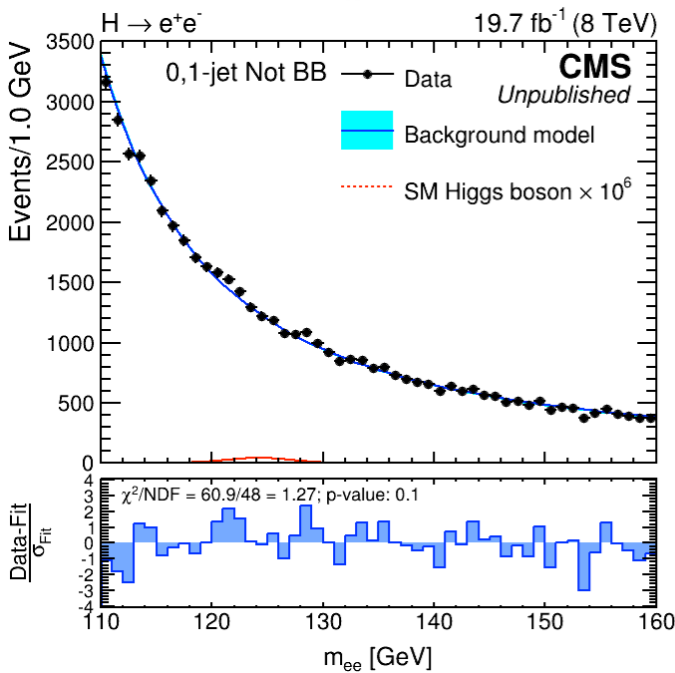
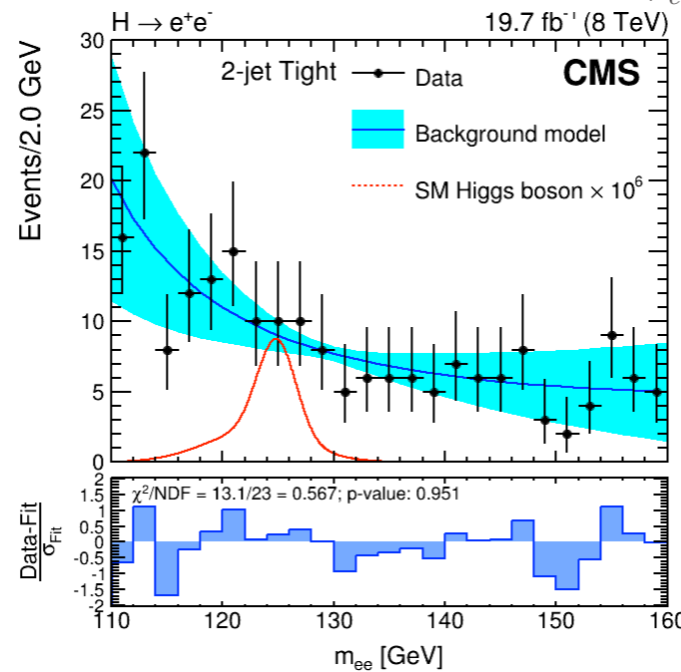
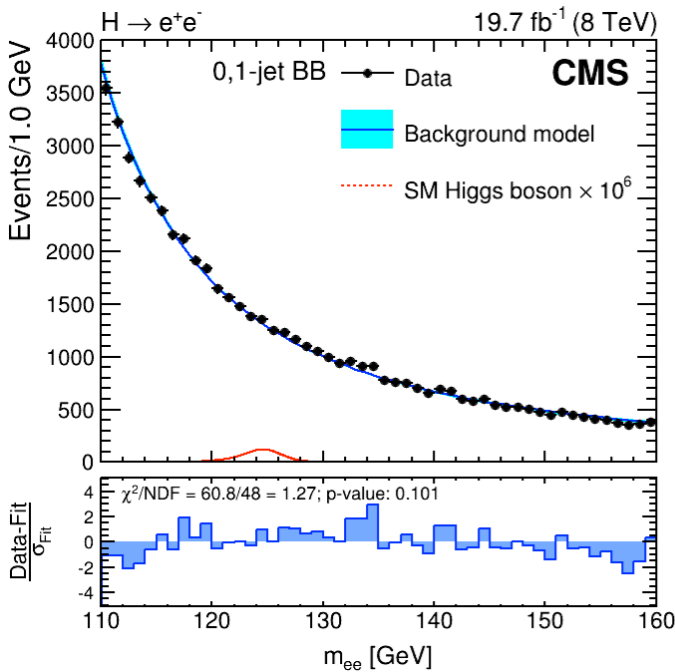
- backgrounds: Z/γ* → ee, and some ttbar

- Categorization:

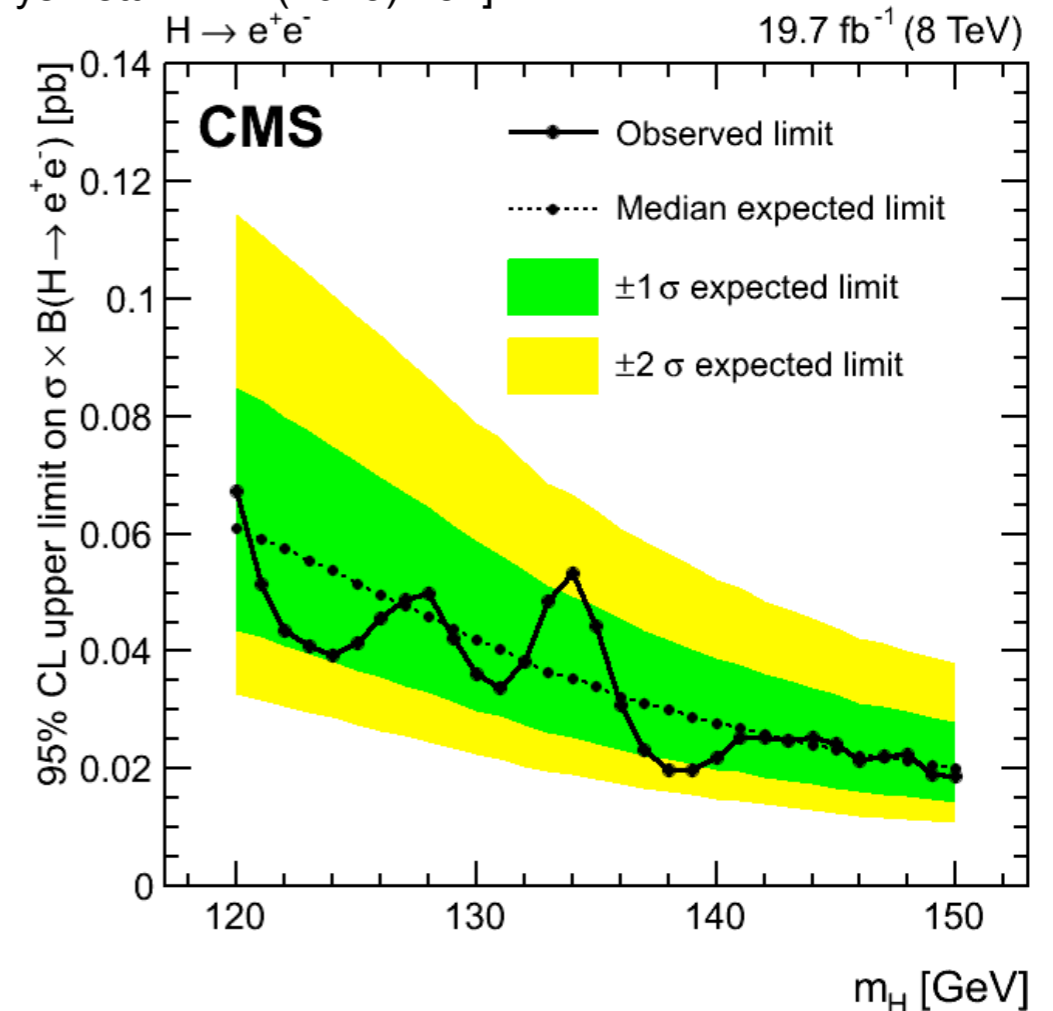
central/non-central electrons + 2 jet categories

- Parametric Background Model

- 95% CL upper limit BR(h → ee) < 1.9 · 10⁻³



[Phys.Lett. B744 (2015) 184]



Charm-quark Yukawa coupling

GLUON PHOTON NEUTRINO TACHYON ELECTRON UP QUARK DOWN QUARK TAU NEUTRINO MUON UP Q
NEUTRON DOWN QUARK TAU GLUON CHARM QUARK TACHYON ELECTRON UP QUARK DOWN QUARK
NEUTRINO MUON UP QUARK PROTON NEUTRON DOWN QUARK TAU GLUON PHOTON NEUTRINO TACHY
The PARTICLE ZOO
UP QUARK PROTON NEUTRON DOWN QUARK TAU GLUON PHOTON NEUTRINO TACHY
DOWN QUARK TAU GLUON PHOTON NEUTRINO TACHYON ELECTRON UP QUARK DOWN QUARK TAU NEU



Substantial recent activity on probing directly the charm-quark Yukawa coupling at the LHC.

Two lines of attack (non-exhaustive list of references):

→ Look for $H \rightarrow c\bar{c}$ using charm-tagging, in a manner similar to the $H \rightarrow b\bar{b}$

[Phys.Rev. D89 (2014) 3, 033014; arXiv:1503.00290]

→ Exploit the exclusive decays $H \rightarrow Q\gamma$ as direct probe to the quark Yukawa couplings

[arXiv:1505.03870, Phys.Lett. B82 (1979) 411; Phys.Rev. D27 (1983) 2762; Yad.Fiz. 46, 864 (1987);

Phys.Rev. D88 (2013) 5, 053003; Phys.Rev. D90 (2014) 11, 113010, arXiv:1505.03870 [hep-ph]]

→ Sensitive to BSM physics [arXiv:1504.04022, Phys. Rev. D 80, 076002, Phys. Lett. B665 (2008) 79]

Charm Tagging

One may “re-interpret” the $h \rightarrow b\bar{b}$ search to include the possibility of anomalous $h \rightarrow c\bar{c}$ production

→ In the SM $\text{BR}(h \rightarrow c\bar{c})/\text{BR}(h \rightarrow b\bar{b}) \sim 5.1\%$

→ Enhancement in the charm Yukawa: $\uparrow \text{BR}(h \rightarrow c\bar{c})$, $\downarrow \text{BR}(h \rightarrow b\bar{b})$ [through $\uparrow \Gamma_h$]

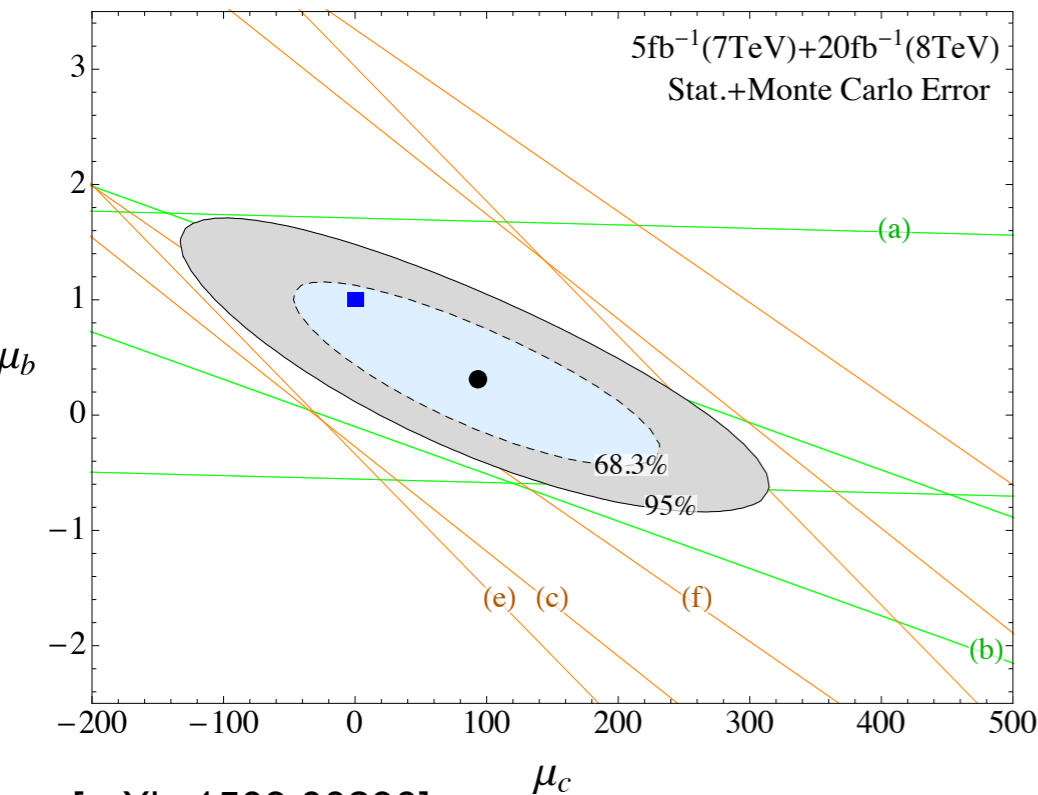
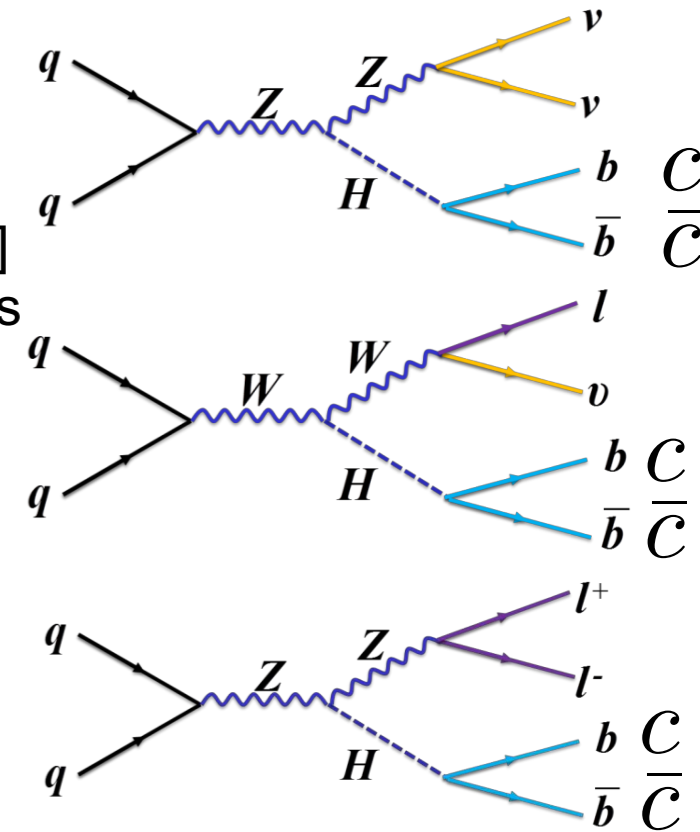
→ Constrains only a linear combination of μ_b and $\mu_c \rightarrow$ need multiple b-tagging points

$$\mu_b = \frac{\sigma \text{BR}_{b\bar{b}}}{\sigma_{\text{SM}} \text{BR}_{b\bar{b}}^{\text{SM}}} \rightarrow \frac{\sigma \text{BR}_{b\bar{b}} \epsilon_{b_1} \epsilon_{b_2} + \sigma \text{BR}_{c\bar{c}} \epsilon_{c_1} \epsilon_{c_2}}{\sigma_{\text{SM}} \text{BR}_{b\bar{b}}^{\text{SM}} \epsilon_{b_1} \epsilon_{b_2}}$$

$$= \mu_b + \frac{\text{BR}_{c\bar{c}}^{\text{SM}}}{\text{BR}_{b\bar{b}}^{\text{SM}}} \frac{\epsilon_{c_1} \epsilon_{c_2}}{\epsilon_{b_1} \epsilon_{b_2}} \mu_c,$$

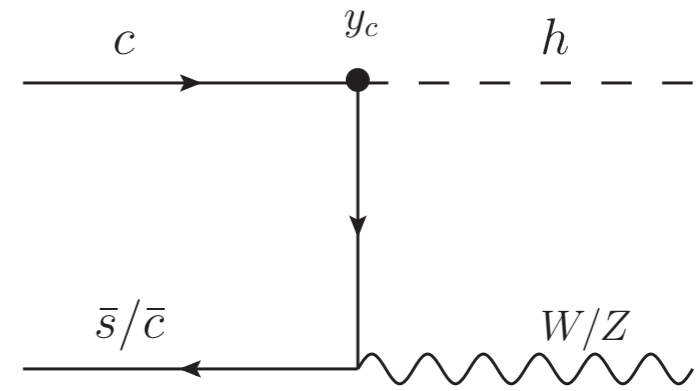
[arXiv:1503.00290]

$$\mu_c = 95^{+90(175)}_{-95(180)} \text{ at } 68.3(95)\% \text{ CL.}$$



→ Extracting info about Yukawa couplings: account for new production modes

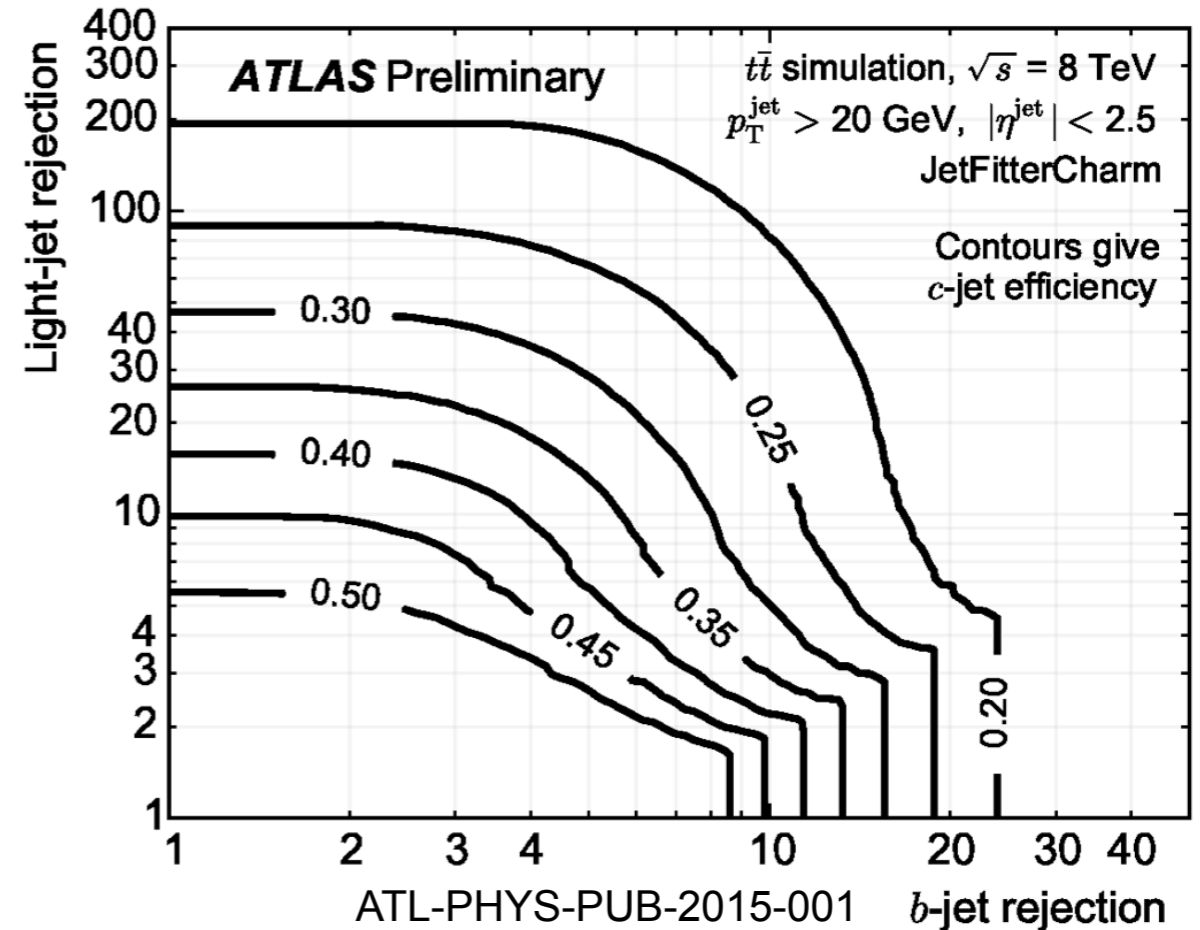
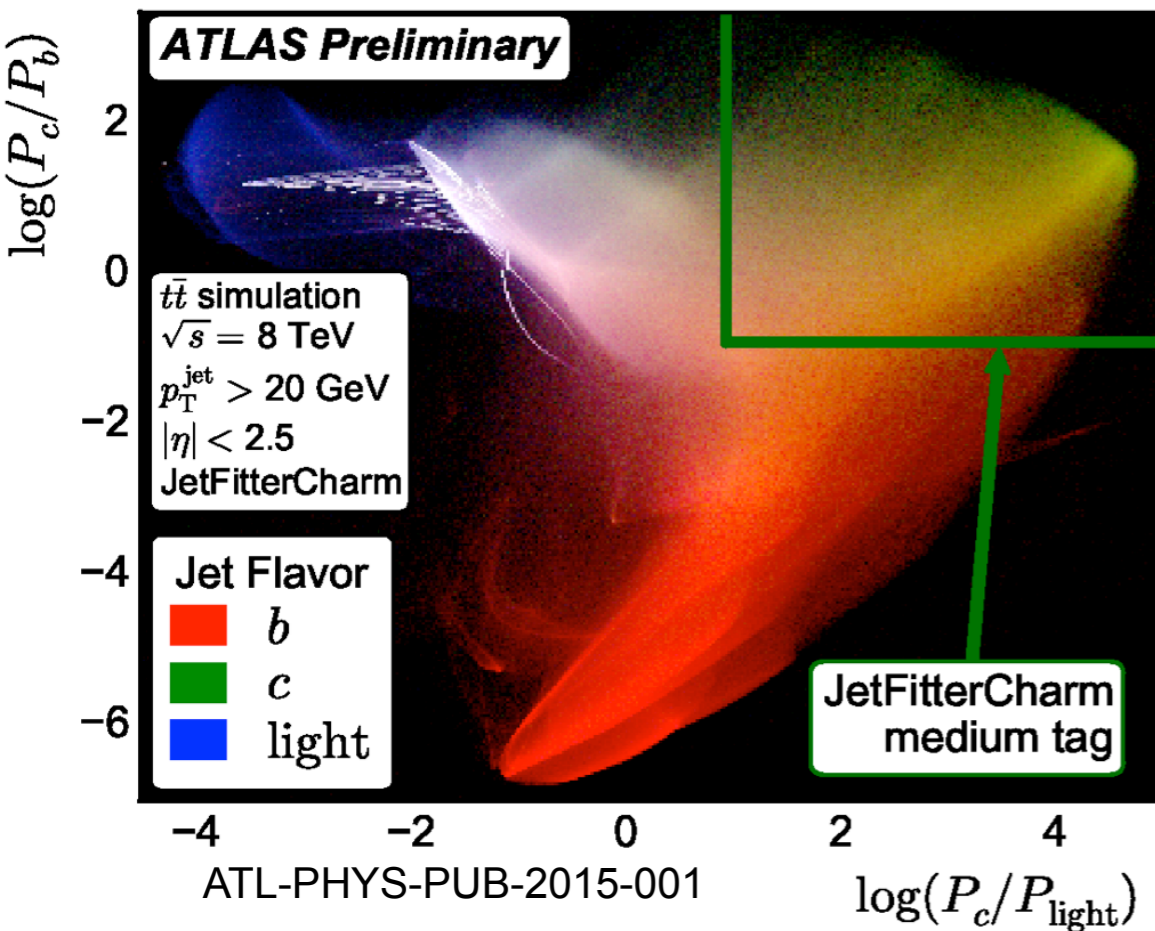
$$\kappa_c \lesssim 234 \text{ at } 95\% \text{ CL,}$$



→ No experiment has performed this analysis yet!

[arXiv:1503.00290]

Charm Tagging



To resolve the two contributions improved c -tagging is needed
 [ideally you would like to completely separate b - and c -jets]

→ Future $H \rightarrow c\bar{c}b\bar{b}$ searches will benefit from dedicated c -tagging, already applied in ATLAS scalar-charm search. [arXiv:1501.01325]

However:

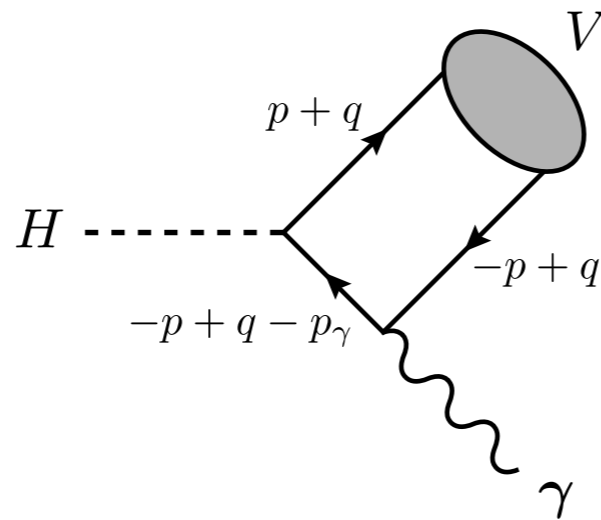
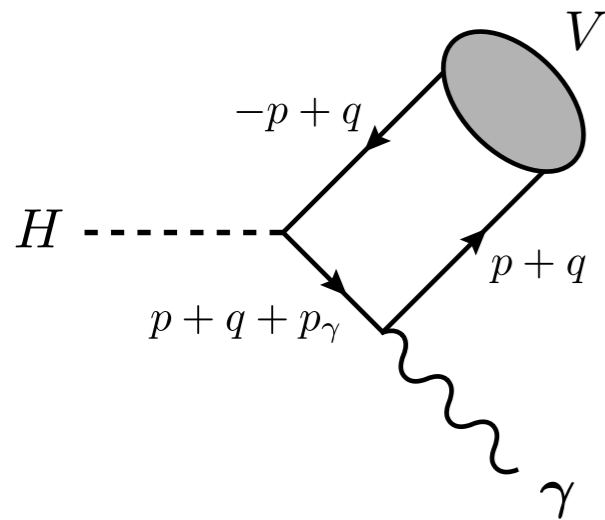
- complicated analysis with large QCD backgrounds
- signal sits on top of large (x20) $h \rightarrow b\bar{b}$ “background”
- sensitivity to systematics of b/c -tagging efficiency
- need dedicated simulations for decay and production

Exclusive Decays $h \rightarrow Q\gamma$

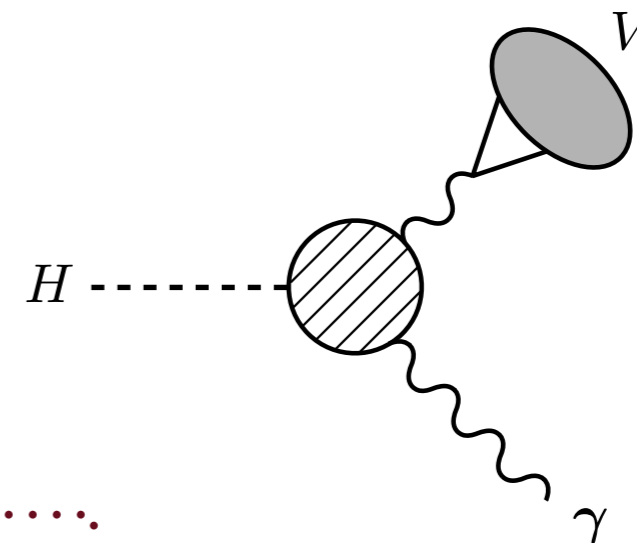
These exclusive decays can lead to quite distinct experimental signatures:

→ Decay of a high- p_T Quarkonium back-to-back with a high- p_T photon.

“Direct” amplitude



“Indirect” contribution



Direct amplitude alone:
 → $BR_{SM}(H \rightarrow J/\psi \gamma) = 5.48 \cdot 10^{-8}$
 → $BR_{SM}(H \rightarrow \Upsilon \gamma) = 3.84 \cdot 10^{-7}$

Phys.Rev. D88 (2013) 5, 053003

However, the indirect amplitude corresponds to a significant contribution

→ Larger than the direct amplitude

→ Negative interference between the two amplitudes could help resolve coupling sign-ambiguity

$$\Gamma(H \rightarrow J/\psi + \gamma) = |(11.9 \pm 0.2) - (1.04 \pm 0.14)\kappa_c|^2 \times 10^{-10} \text{ GeV}$$

$$\mathcal{B}_{SM}(H \rightarrow J/\psi + \gamma) = 2.79_{-0.15}^{+0.16} \times 10^{-6},$$

$$Br(h \rightarrow J/\psi \gamma) = (2.95 \pm 0.07_{f_{J/\psi}} \pm 0.06_{\text{direct}} \pm 0.14_{h \rightarrow \gamma\gamma}) \cdot 10^{-6},$$

$$\mathcal{B}_{SM}[H \rightarrow \Upsilon(1S) + \gamma] = 6.11_{-6.11}^{+17.41} \times 10^{-10},$$

$$Br(h \rightarrow \Upsilon(1S) \gamma) = (4.61 \pm 0.06_{f_{\Upsilon(1S)}} \pm 1.75_{\text{direct}} \pm 0.22_{h \rightarrow \gamma\gamma}) \cdot 10^{-9},$$

$$\mathcal{B}_{SM}[H \rightarrow \Upsilon(2S) + \gamma] = 2.02_{-1.28}^{+1.86} \times 10^{-9},$$

$$Br(h \rightarrow \Upsilon(2S) \gamma) = (2.34 \pm 0.04_{f_{\Upsilon(2S)}} \pm 0.75_{\text{direct}} \pm 0.11_{h \rightarrow \gamma\gamma}) \cdot 10^{-9},$$

$$\mathcal{B}_{SM}[H \rightarrow \Upsilon(3S) + \gamma] = 2.44_{-1.30}^{+1.75} \times 10^{-9}.$$

$$Br(h \rightarrow \Upsilon(3S) \gamma) = (2.13 \pm 0.04_{f_{\Upsilon(3S)}} \pm 0.75_{\text{direct}} \pm 0.10_{h \rightarrow \gamma\gamma}) \cdot 10^{-9}.$$

Phys.Rev. D90 (2014) 11, 113010

arXiv:1505.03870 [hep-ph]

Exclusive Decays $Z \rightarrow Q\gamma$: J/ψ and Υ

- The analogous Z boson decays also attracting significant attention [Nucl. Phys. B 174, 317 (1980), Theor. Math. Phys. 170, 39 (2012), arXiv:1411.5924, JHEP 1504 (2015) 101]
- These exclusive final states are experimentally unconstrained
- Could be sensitive to BSM contributions
- LEP has accurately measured couplings to b- and c-quarks ($\sim 1\%$), but couplings to light quarks less constrained.



Recently, a number of numerical results have appeared, regarding the predictions on these decay rates.

JHEP 1504 (2015) 101

Decay mode	Branching ratio
$Z^0 \rightarrow \pi^0 \gamma$	$(9.80^{+0.09}_{-0.14} \mu \pm 0.03_f \pm 0.61_{a_2} \pm 0.82_{a_4}) \cdot 10^{-12}$
$Z^0 \rightarrow \rho^0 \gamma$	$(4.19^{+0.04}_{-0.06} \mu \pm 0.16_f \pm 0.24_{a_2} \pm 0.37_{a_4}) \cdot 10^{-9}$
$Z^0 \rightarrow \omega \gamma$	$(2.89^{+0.03}_{-0.05} \mu \pm 0.15_f \pm 0.29_{a_2} \pm 0.25_{a_4}) \cdot 10^{-8}$
$Z^0 \rightarrow \phi \gamma$	$(8.63^{+0.08}_{-0.13} \mu \pm 0.41_f \pm 0.55_{a_2} \pm 0.74_{a_4}) \cdot 10^{-9}$
$Z^0 \rightarrow J/\psi \gamma$	$(8.02^{+0.14}_{-0.15} \mu \pm 0.20_f \pm 0.39_{\sigma}) \cdot 10^{-8}$
$Z^0 \rightarrow \Upsilon(1S) \gamma$	$(5.39^{+0.10}_{-0.10} \mu \pm 0.08_f \pm 0.11_{\sigma}) \cdot 10^{-8}$
$Z^0 \rightarrow \Upsilon(4S) \gamma$	$(1.22^{+0.02}_{-0.02} \mu \pm 0.13_f \pm 0.02_{\sigma}) \cdot 10^{-8}$
$Z^0 \rightarrow \Upsilon(nS) \gamma$	$(9.96^{+0.18}_{-0.19} \mu \pm 0.09_f \pm 0.20_{\sigma}) \cdot 10^{-8}$

arXiv:1411.5924

$$B_{SM}(Z \rightarrow J/\psi + \gamma) = (9.96 \pm 1.86) \times 10^{-8},$$

$$B_{SM}(Z \rightarrow \Upsilon(1S) + \gamma) = (4.93 \pm 0.51) \times 10^{-8},$$

$$B_{SM}(Z \rightarrow \phi + \gamma) = (1.17 \pm 0.08) \times 10^{-8}.$$

consistent when rescaled
to up-to-date value of f_ϕ

Search for Higgs and Z Boson Decays to $J/\psi\gamma$ and $\Upsilon(nS)\gamma$ with the ATLAS DetectorG. Aad *et al.**

(ATLAS Collaboration)

(Received 15 January 2015; published 26 March 2015)

A search for the decays of the Higgs and Z bosons to $J/\psi\gamma$ and $\Upsilon(nS)\gamma$ ($n = 1, 2, 3$) is performed with pp collision data samples corresponding to integrated luminosities of up to 20.3 fb^{-1} collected at $\sqrt{s} = 8 \text{ TeV}$ with the ATLAS detector at the CERN Large Hadron Collider. No significant excess of events is observed above expected backgrounds and 95% C.L. upper limits are placed on the branching fractions. In the $J/\psi\gamma$ final state the limits are 1.5×10^{-3} and 2.6×10^{-6} for the Higgs and Z boson decays, respectively, while in the $\Upsilon(1S, 2S, 3S)\gamma$ final states the limits are $(1.3, 1.9, 1.3) \times 10^{-3}$ and $(3.4, 6.5, 5.4) \times 10^{-6}$, respectively.

DOI: [10.1103/PhysRevLett.114.121801](https://doi.org/10.1103/PhysRevLett.114.121801)

PACS numbers: 14.80.Bn, 13.38.Dg, 14.70.Hp, 14.80.Ec

ATLAS performed the first search for these exclusive decays of the Higgs and Z bosons
 $H/Z \rightarrow Q\gamma$, where $Q = J/\psi$ or $\Upsilon(nS)$, $n=1,2,3$

$h \rightarrow J/\psi \gamma$ and $h \rightarrow Y(ns) \gamma$

Phys.Rev.Lett. 114 (2015) 12, 121801

Signature: $\mu^+ \mu^- + \gamma$

→ event selection:

single(di)-muon trigger

$p_{T\mu} > 20, 3$ GeV,

$p_{T\mu\mu} > 36$ GeV,

$p_{T\gamma} > 36$ GeV

$\mu\mu$ and γ isolation,

[J/ψ mass requirement]

L_{xy} significance,

$\Delta\phi(\mu\mu, \gamma) > 0.5$

→ total signal acceptance/efficiency

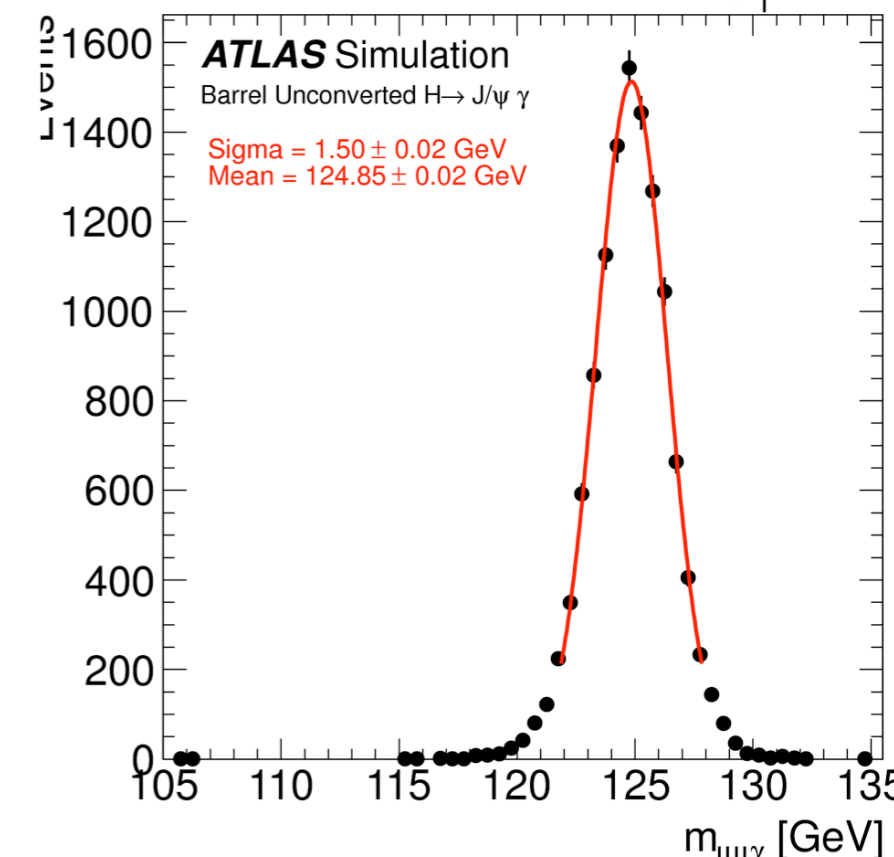
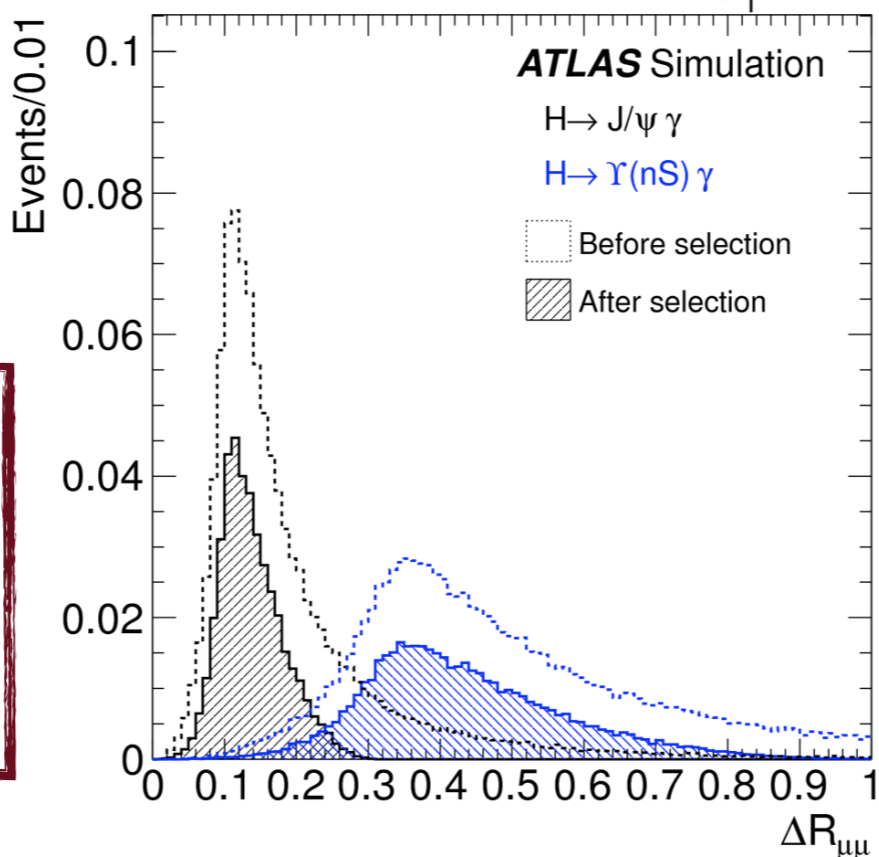
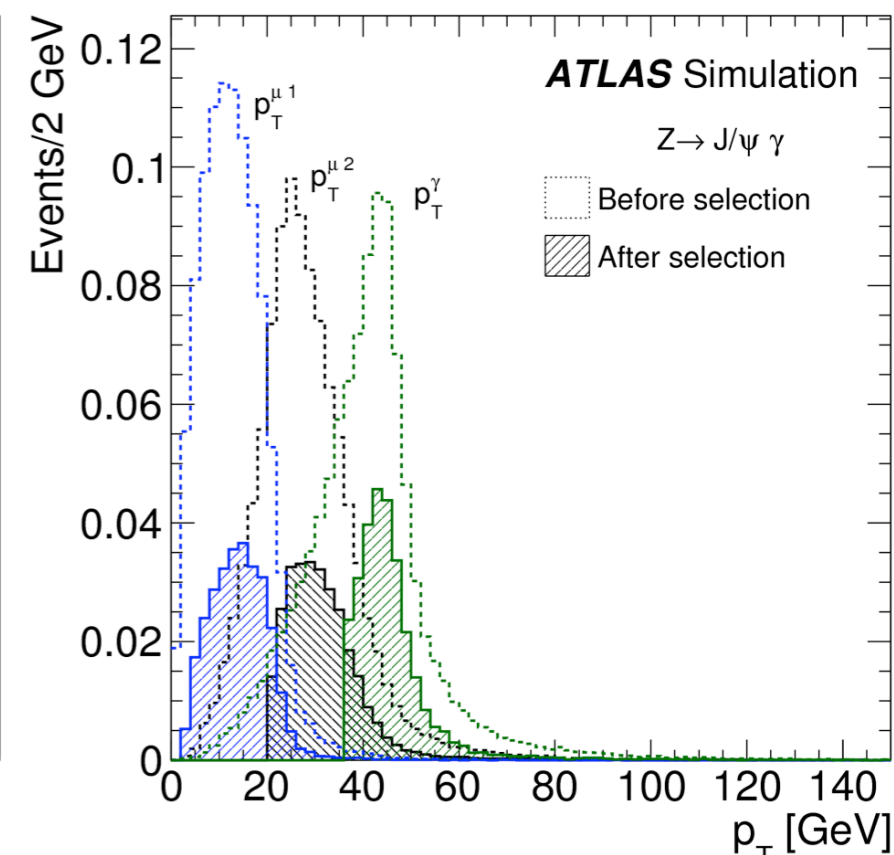
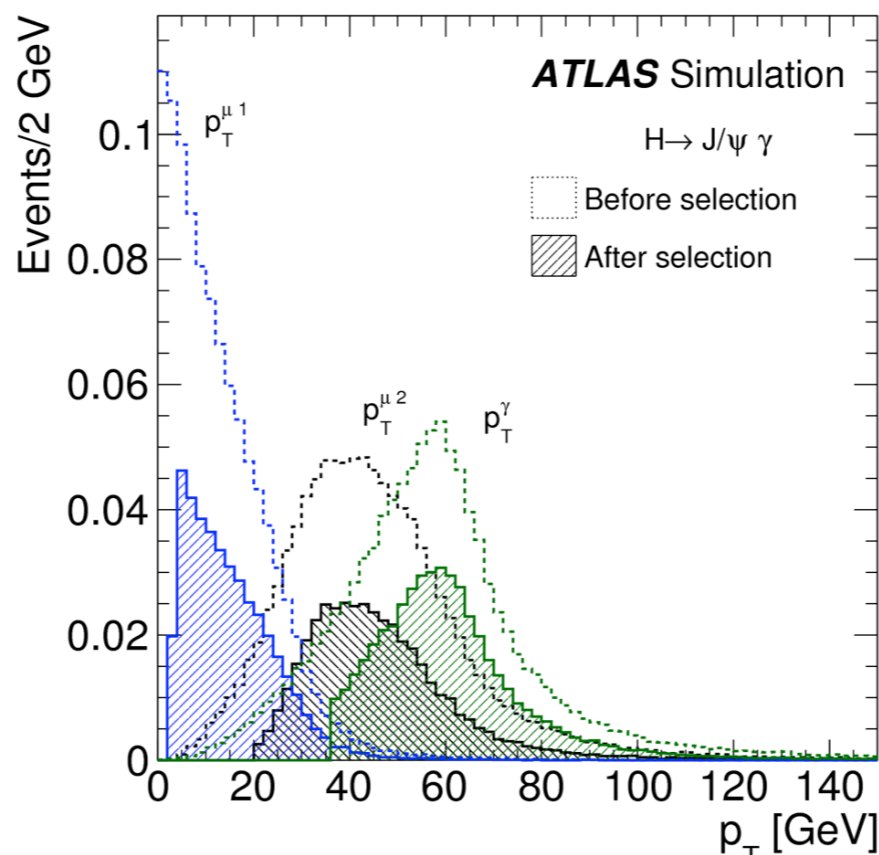
$H(Z) \rightarrow J/\psi \gamma \rightarrow \mu^+ \mu^- + \gamma \sim 22\%$ (12%)

$H(Z) \rightarrow Y \gamma \rightarrow \mu^+ \mu^- + \gamma \sim 28\%$ (15%)

→ $m_{\mu\mu\gamma}$ mass resolution $\sim 1.2-1.8\%$

$e^+ e^- + \gamma$

- experimentally more challenging
- dedicated reconstruction for nearby electrons
- poorer mass resolution/efficiency
- typically larger backgrounds



$h \rightarrow J/\psi \gamma$ and $h \rightarrow Y(ns) \gamma$

Phys.Rev.Lett. 114 (2015) 12, 121801

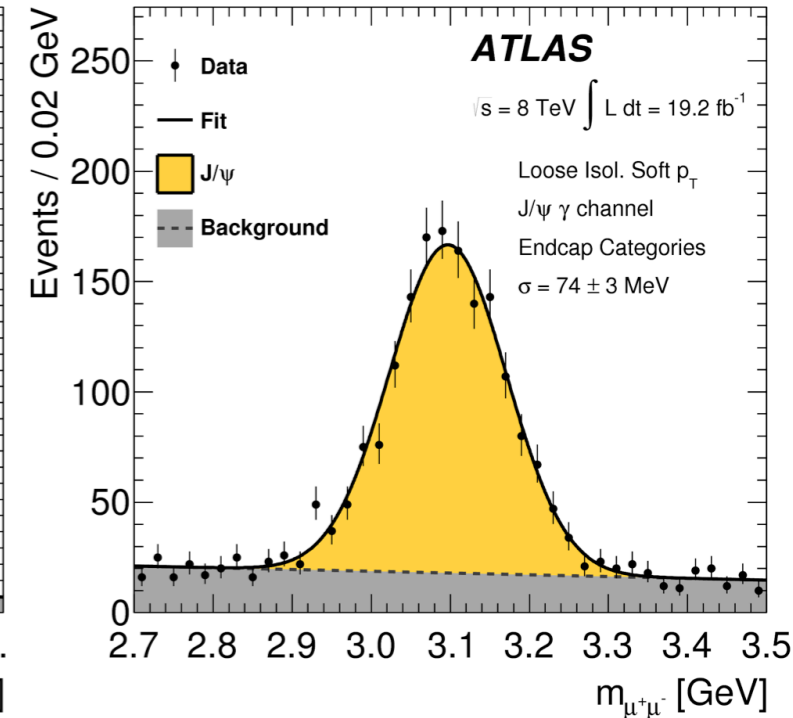
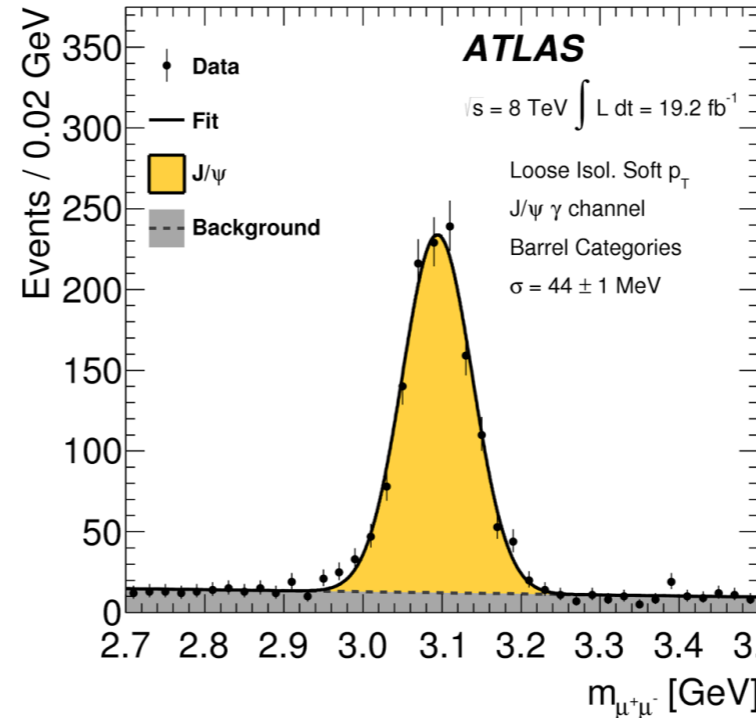
Categories

For this search simple - detector performance driven categorisation

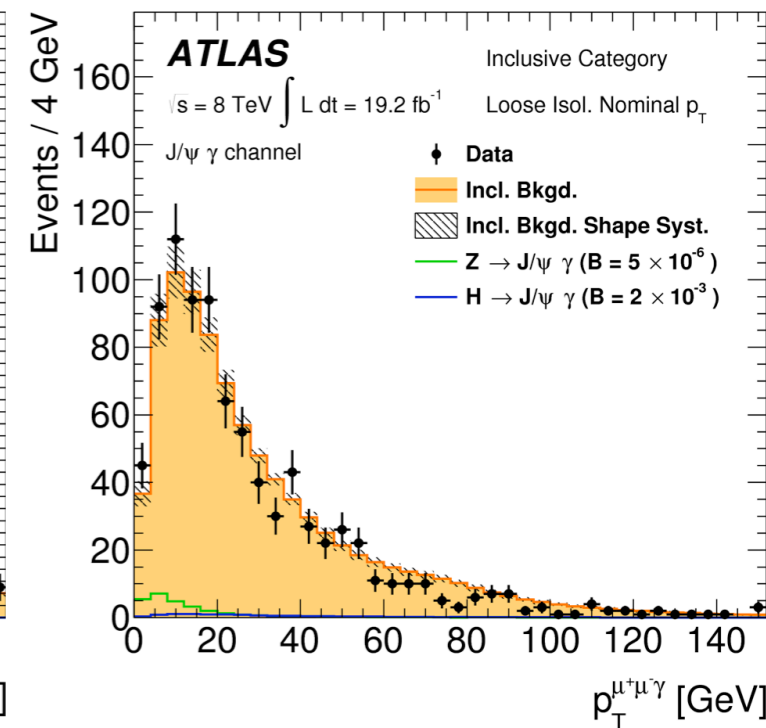
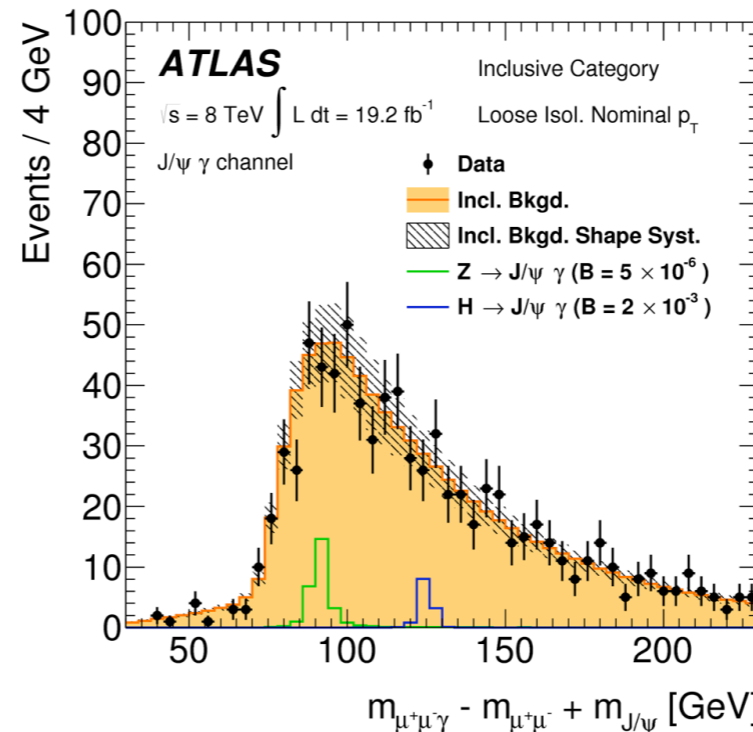
- Muon pseudo-rapidity
 - Both Central/One Non-Central
- Photon Unconverted/Converted

Background

- inclusive Quarkonium with jet “seen” as γ
 - small component of combinatoric
- Non-parametric data-driven background estimation
 - for $Y(nS) \gamma$ also $Z \rightarrow \mu\mu \gamma_{FSR}$ from side-band fit



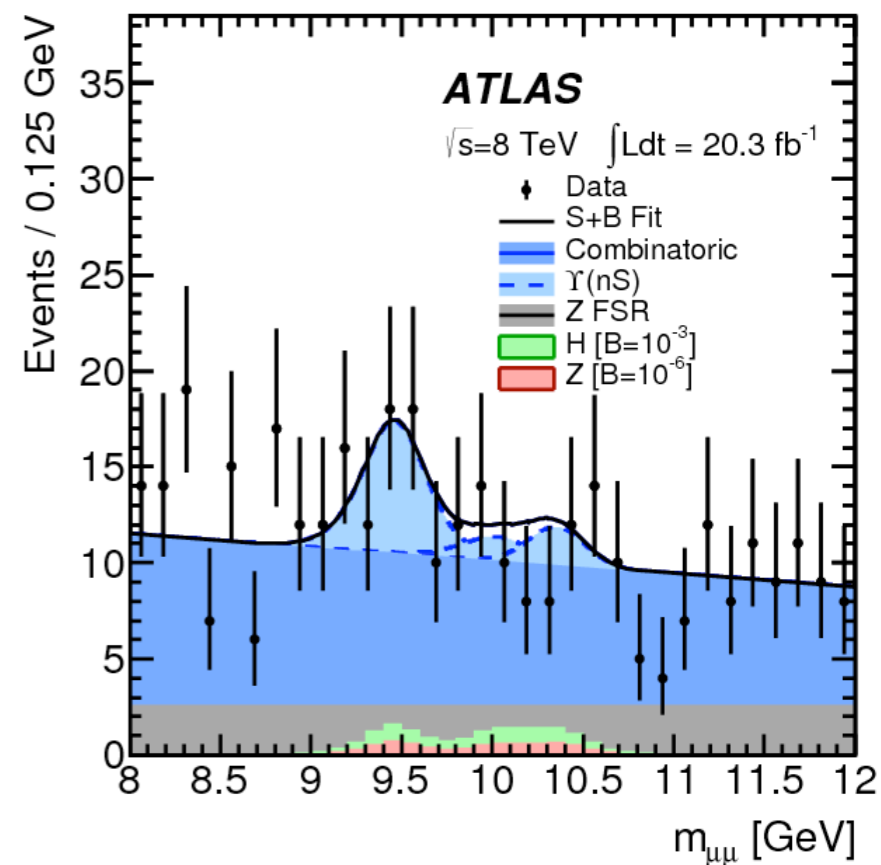
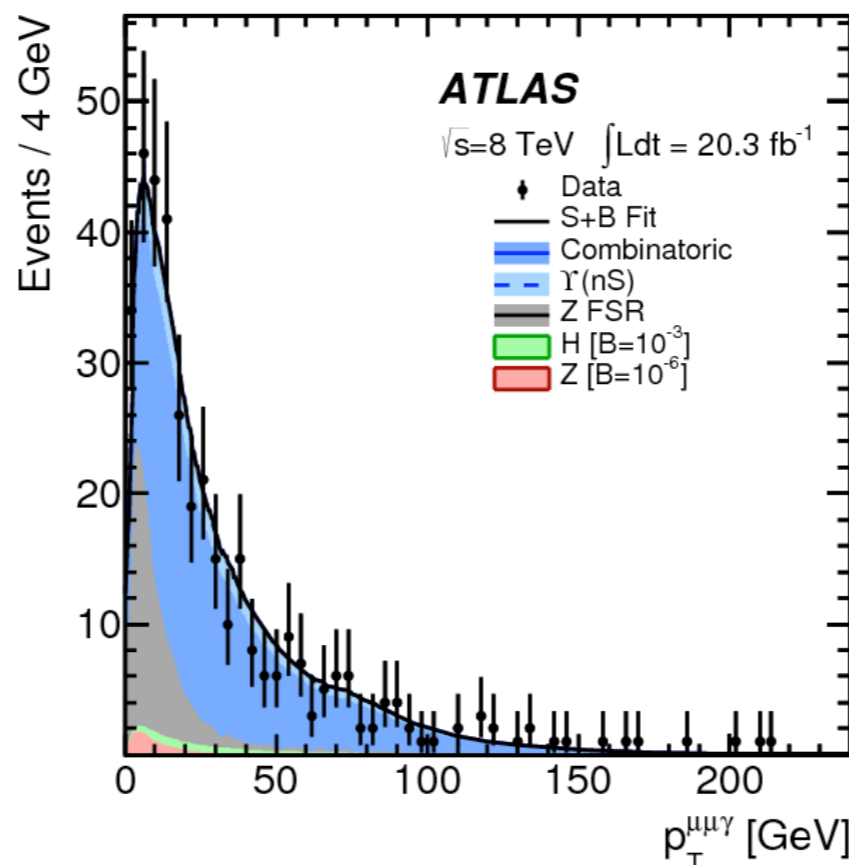
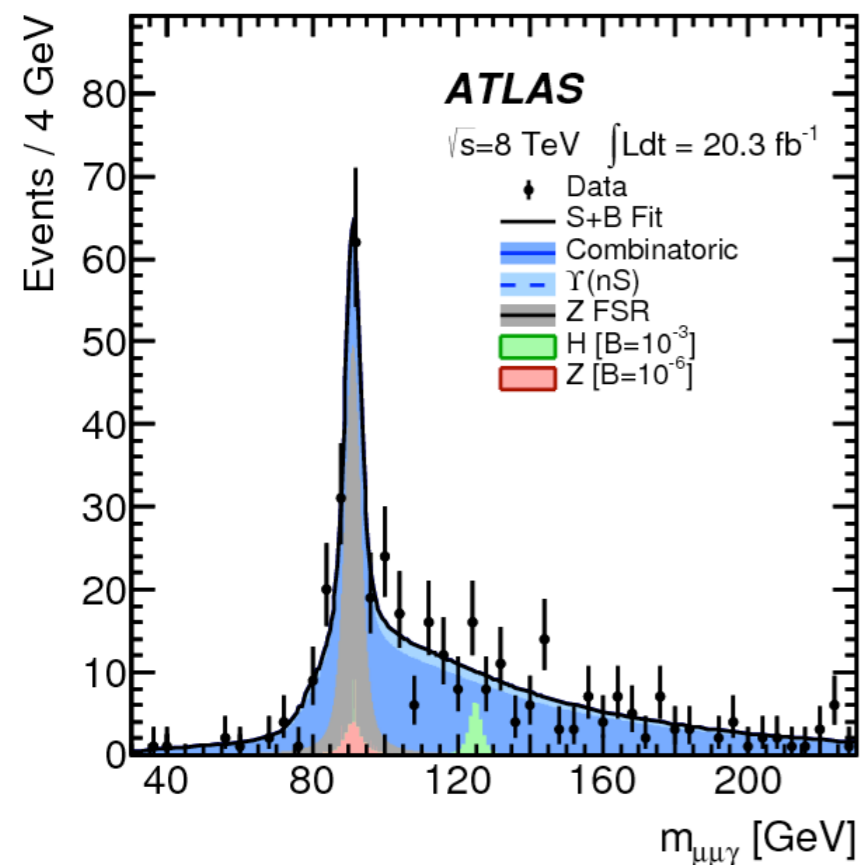
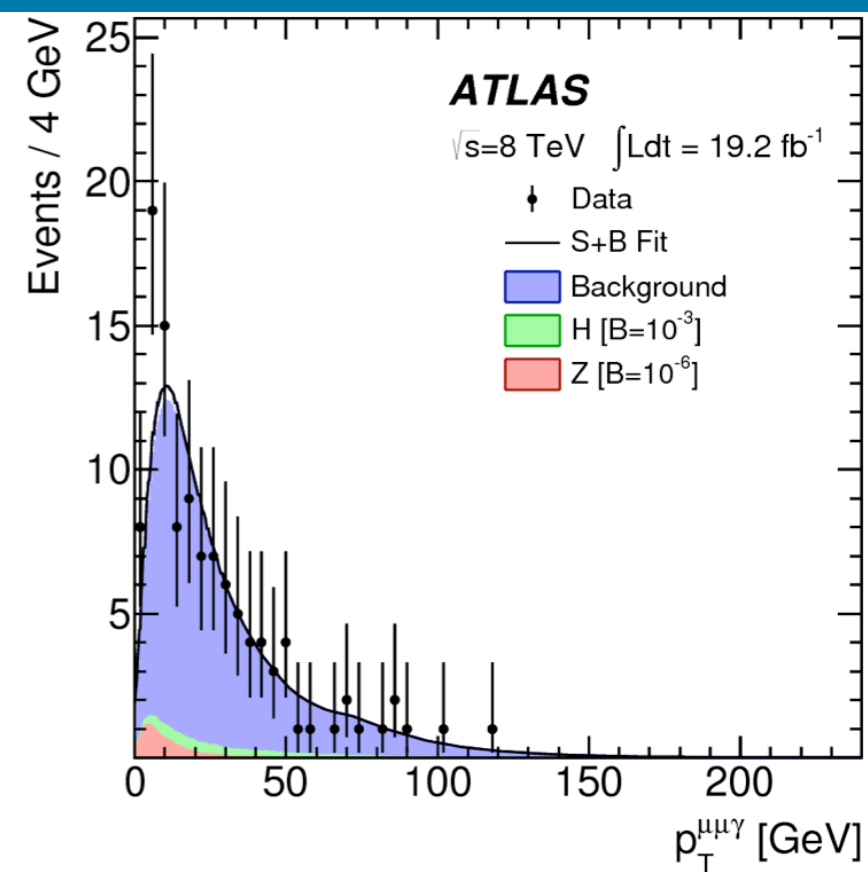
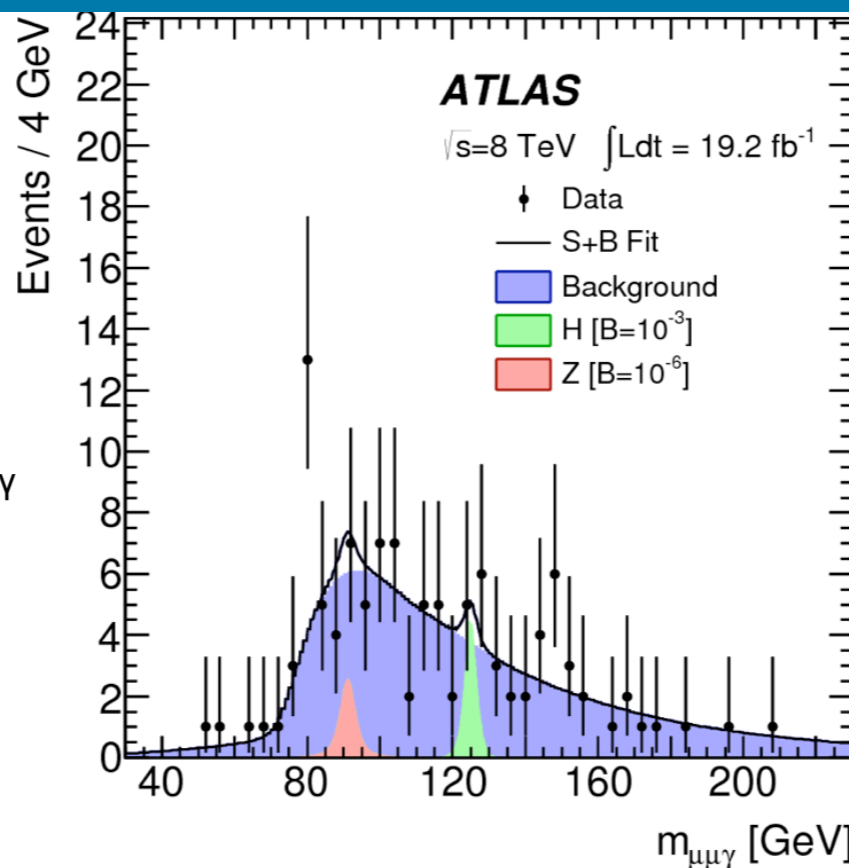
Category	Observed (Expected Background)				Signal		
	All	Mass Range [GeV]		Z	H		
		80-100	115-135	$\mathcal{B} [10^{-6}]$	$\mathcal{B} [10^{-3}]$		
		$J/\psi \gamma$					
BU	30	9	(8.9±1.3)	5	(5.0±0.9)	1.29±0.07	1.96±0.24
BC	29	8	(6.0±0.7)	3	(5.5±0.6)	0.63±0.03	1.06±0.13
EU	35	8	(8.7±1.0)	10	(5.8±0.8)	1.37±0.07	1.47±0.18
EC	23	6	(5.6±0.7)	2	(3.0±0.4)	0.99±0.05	0.93±0.12
		$\Upsilon(ns) \gamma$					
BU	93	42	(39±6)	16	(12.9±2.0)	1.67±0.09	2.6±0.3
BC	71	32	(27.7±2.4)	5	(9.7±1.2)	0.79±0.04	1.45±0.18
EU	125	49	(47±6)	16	(17.8±2.4)	2.24±0.12	2.5±0.3
EC	85	31	(31±5)	18	(12.3±1.9)	1.55±0.08	1.60±0.20



$h \rightarrow J/\psi \gamma$ and $h \rightarrow Y(ns)\gamma$: Results

Phys.Rev.Lett. 114 (2015) 12, 121801

- Multi-observable fit to $m_{\mu\mu\gamma}$, $p_T^{\mu\mu\gamma}$
- also $m_{\mu\mu}$ for $Y(ns)\gamma$
- No significant excess above background observed

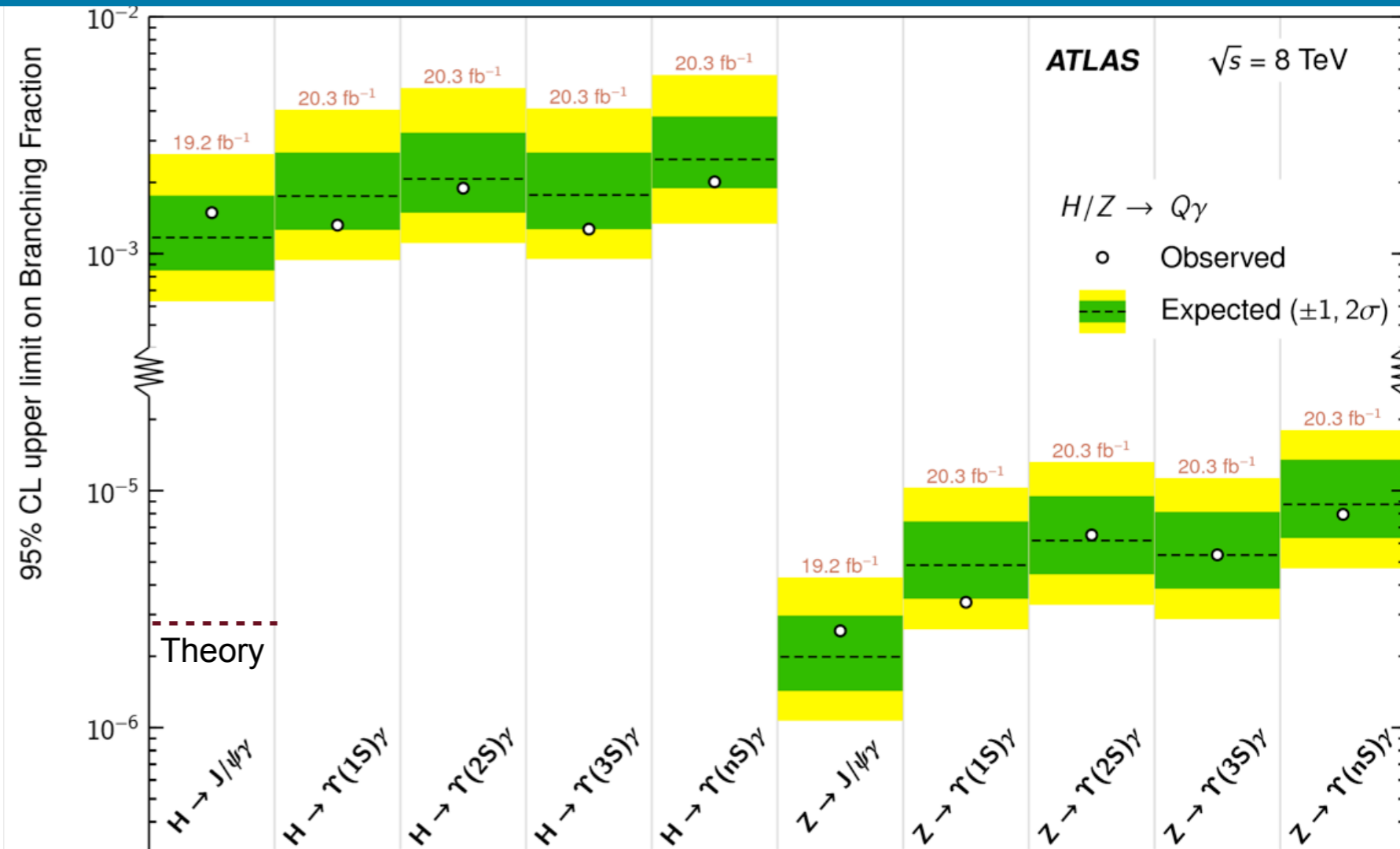


h → J/ψγ and h → Υ(ns)γ: Results

Phys.Rev.Lett. 114 (2015) 12, 121801

First search for H/Z → Qγ, will constitute the basis for Run 2 and HL-LHC extrapolations

BR 95% CLs upper limits:
 ~10⁻³ level for Higgs boson (SM production) decays and
 ~10⁻⁶ for the Z boson decays



This is a nice and, relatively, clean final state.
 Fun and interesting thing to do!

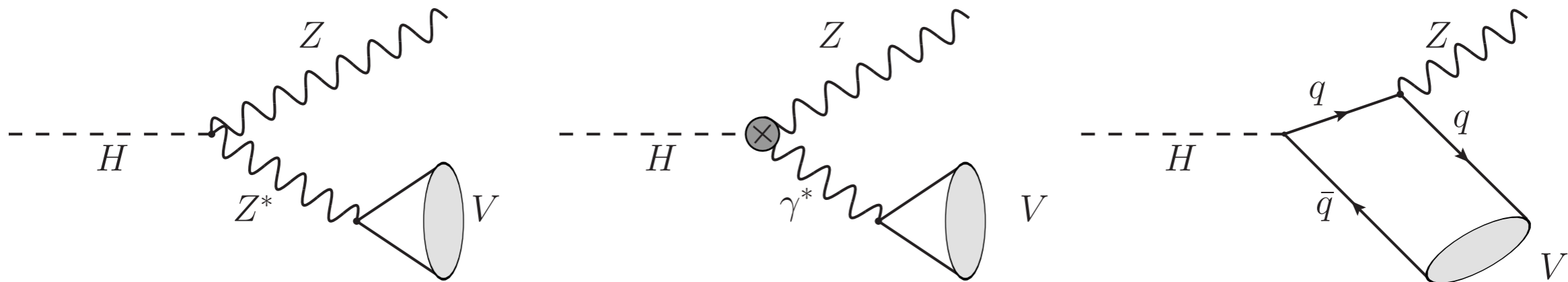
A few drawbacks of these exclusive decays:

- 1) Small branching ratio, a handful of events expected even at HL-LHC
- 2) At SM sensitivity significant contribution from non-resonant h → μμγ ~ 3 × h → J/ψγ
- 3) This channel is also affected by potential “anomalies” in the h → γγ loop

	95% CL _s Upper Limits				
	J/ψ	Υ(1S)	Υ(2S)	Υ(3S)	∑ ⁿ Υ(nS)
$\mathcal{B}(Z \rightarrow Q\gamma) [10^{-6}]$					
Expected	2.0 ^{+1.0} _{-0.6}	4.9 ^{+2.5} _{-1.4}	6.2 ^{+3.2} _{-1.8}	5.4 ^{+2.7} _{-1.5}	8.8 ^{+4.7} _{-2.5}
Observed	2.6	3.4	6.5	5.4	7.9
$\mathcal{B}(H \rightarrow Q\gamma) [10^{-3}]$					
Expected	1.2 ^{+0.6} _{-0.3}	1.8 ^{+0.9} _{-0.5}	2.1 ^{+1.1} _{-0.6}	1.8 ^{+0.9} _{-0.5}	2.5 ^{+1.3} _{-0.7}
Observed	1.5	1.3	1.9	1.3	2.0
$\sigma(pp \rightarrow H) \times \mathcal{B}(H \rightarrow Q\gamma) [\text{fb}]$					
Expected	26 ⁺¹² ₋₇	38 ⁺¹⁹ ₋₁₁	45 ⁺²⁴ ₋₁₃	38 ⁺¹⁹ ₋₁₁	54 ⁺²⁷ ₋₁₅
Observed	33	29	41	28	44

$h \rightarrow ZJ/\psi$ and $h \rightarrow ZY(nS)$: A short note

arXiv:1411.2210



$h \rightarrow ZQ$ could be another way to approach the charm/bottom Yukawa couplings, quite similar to the exclusive $h \rightarrow Q\gamma$ decays discussed earlier.

arXiv:1411.2210

$Br(H \rightarrow ZV)$	$J/\psi(1S)$	$\Upsilon(1S)$
Br_{Γ_1}	1.75×10^{-6}	1.68×10^{-5}
Br_{Γ_2}	1.14×10^{-6}	8.33×10^{-8}
Br_{Γ_3}	8.52×10^{-9}	5.80×10^{-7}
$Br_{\Gamma_{12}}$	4.50×10^{-7}	1.10×10^{-6}
$Br_{\Gamma_{13}}$	3.89×10^{-8}	2.89×10^{-6}
$Br_{\Gamma_{23}}$	1.97×10^{-7}	4.40×10^{-7}

$$\Gamma = \Gamma_1 \kappa_V^2 + \Gamma_2 \kappa_{Z\gamma}^2 + \Gamma_3 \kappa_Q^2 + \Gamma_{12} \kappa_V \kappa_{Z\gamma} + \Gamma_{13} \kappa_V \kappa_Q + \Gamma_{23} \kappa_{Z\gamma} \kappa_Q$$

unfortunately, need to account for the additional $BR(Z \rightarrow ee/\mu\mu) \sim 6\%$ and properly evaluate S/B

$BR_{SM}(h \rightarrow J/\psi\gamma) \approx BR_{SM}(h \rightarrow ZJ/\psi)$

Light-Quark Yukawa couplings

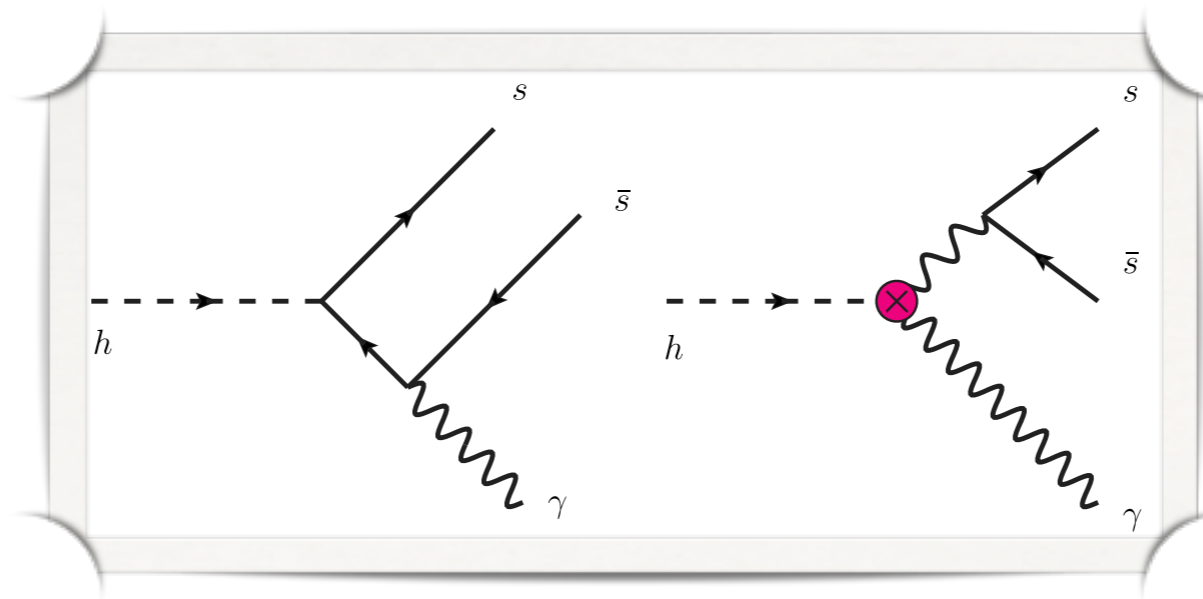
Initially, considered impossible at the LHC, recent activity on its feasibility:

→ **Exploit the exclusive decays $H \rightarrow Q\gamma$ as direct probe to the quark Yukawa couplings**

[Phys.Rev.Lett. 114 (2015) 10, 101802]

→ **Sensitive to BSM physics** [Phys. Rev. D 80, 076002, Phys. Lett. B665 (2008) 79, Phys.Rev. D90 (2014) 115022]

The idea is to benefit from the interference of the “direct” and “indirect” amplitudes!



$$\frac{\text{BR}_{h \rightarrow \phi\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(3.0 \pm 0.3)\kappa_\gamma - 0.78\bar{\kappa}_s] \times 10^{-6}}{0.57\bar{\kappa}_b^2},$$

$$\frac{\text{BR}_{h \rightarrow \rho\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(1.9 \pm 0.2)\kappa_\gamma - 0.24\bar{\kappa}_u - 0.12\bar{\kappa}_d] \times 10^{-5}}{0.57\bar{\kappa}_b^2},$$

$$\frac{\text{BR}_{h \rightarrow \omega\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(1.6 \pm 0.2)\kappa_\gamma - 0.59\bar{\kappa}_u - 0.29\bar{\kappa}_d] \times 10^{-6}}{0.57\bar{\kappa}_b^2},$$

$$BR_{h \rightarrow \phi\gamma}^{SM} \approx 3 \cdot 10^{-6}$$

$$BR_{h \rightarrow \rho\gamma}^{SM} \approx 1.9 \cdot 10^{-5}$$

$$BR_{h \rightarrow \omega\gamma}^{SM} \approx 1.6 \cdot 10^{-6}$$

$$\text{Br}(h \rightarrow \rho^0\gamma) = (1.68 \pm 0.02_{f_\rho} \pm 0.08_{h \rightarrow \gamma\gamma}) \cdot 10^{-5},$$

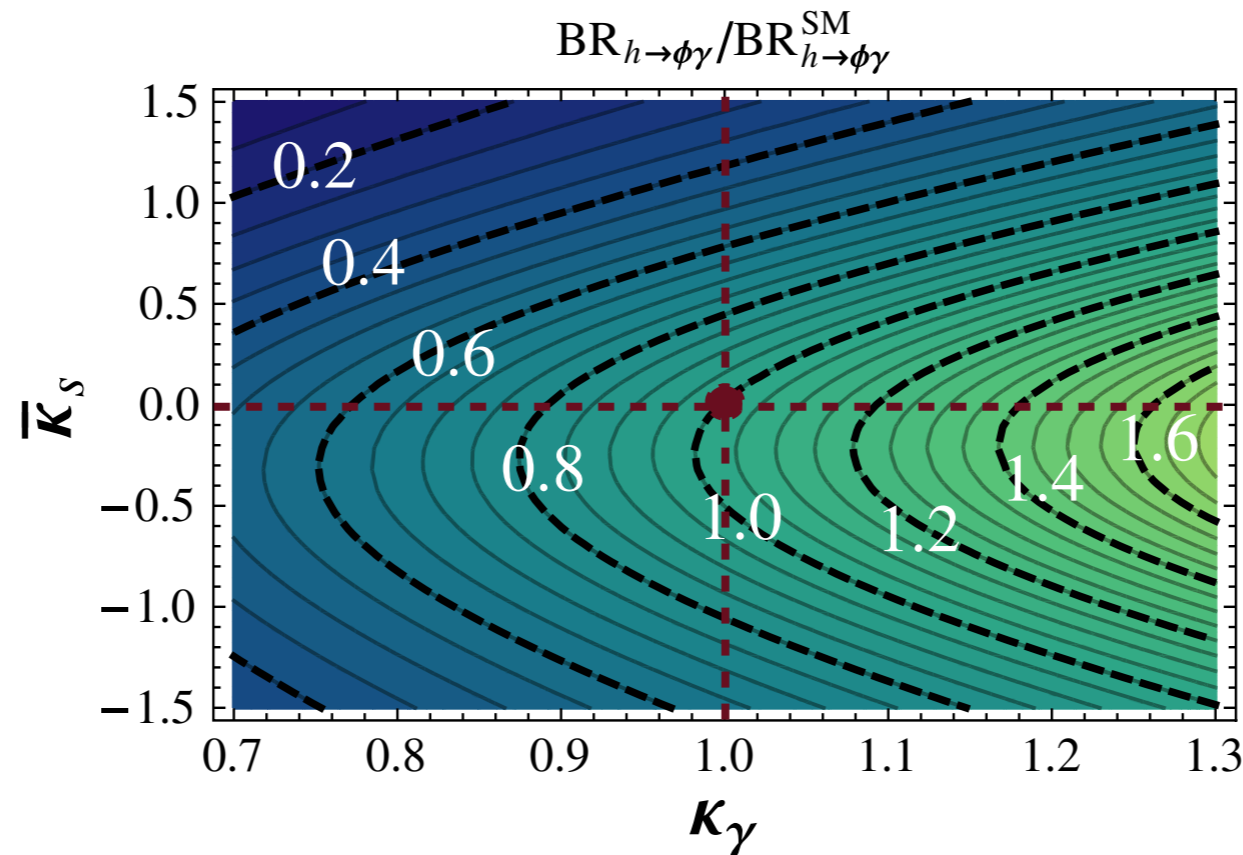
$$\text{Br}(h \rightarrow \omega\gamma) = (1.48 \pm 0.03_{f_\omega} \pm 0.07_{h \rightarrow \gamma\gamma}) \cdot 10^{-6},$$

$$\text{Br}(h \rightarrow \phi\gamma) = (2.31 \pm 0.03_{f_\phi} \pm 0.11_{h \rightarrow \gamma\gamma}) \cdot 10^{-6},$$

arXiv:1505.03870 [hep-ph]

Phys.Rev.Lett. 114 (2015) 10, 101802

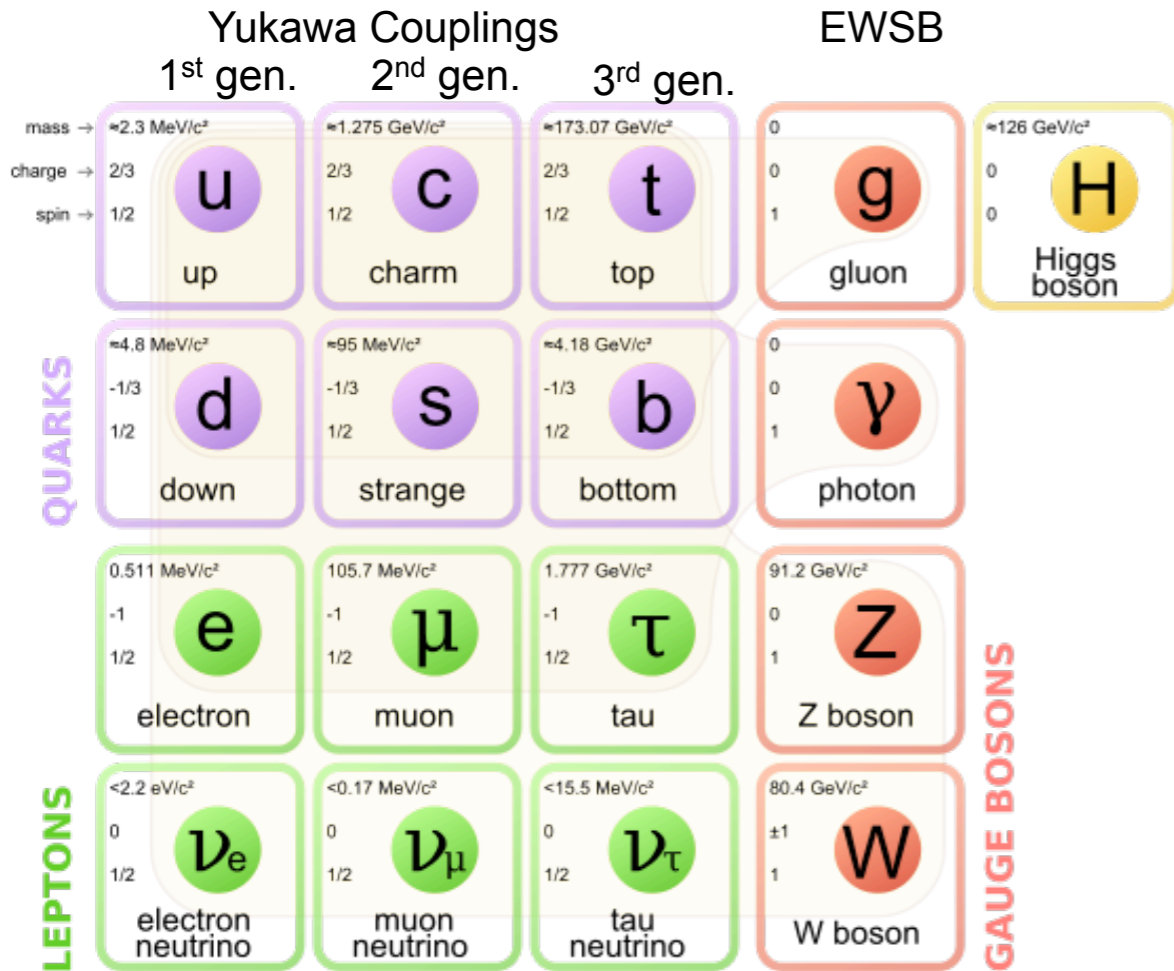
Light-Quark Yukawa couplings



$\phi \rightarrow K^+ K^-$ (BR=48.9%), ~ 8 events 100fb^{-1} @14 TeV!
 $\omega \rightarrow \pi^+ \pi^- \pi^0$ (BR=89.2%) similar rate
 $\rho \rightarrow \pi^+ \pi^-$ (BR \sim 100%) expect ~ 100 events!

Interesting/experimentally challenging topologies!
- triggering on a photon + narrow hadronic jet
→ will benefit from ATLAS FTK
- boosted decays but overwhelming QCD backgrounds,
- Γ_ρ and ω - ρ interference

Summary



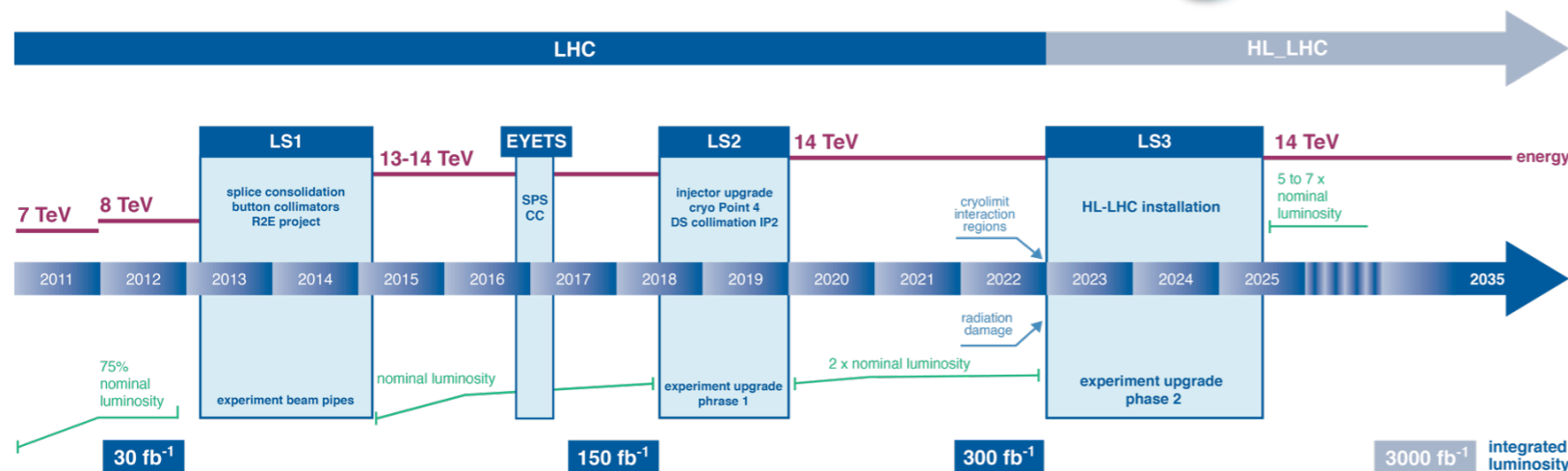
Yukawa sector likely the least theoretically motivated and constrained part of the Standard Model
 → Particularly true for 1st/2nd generation.

New Physics could be lurking here!

A wealth of information has been collected over the last few years on the nature of the Higgs boson
 → Yukawa sector still relatively unconstrained

Currently, under intense phenomenological and experimental focus; new results (H→J/ψγ, H→Υγ, etc) and new ideas/approaches to probe this sector at the LHC appear!

LHC / HL-LHC Plan

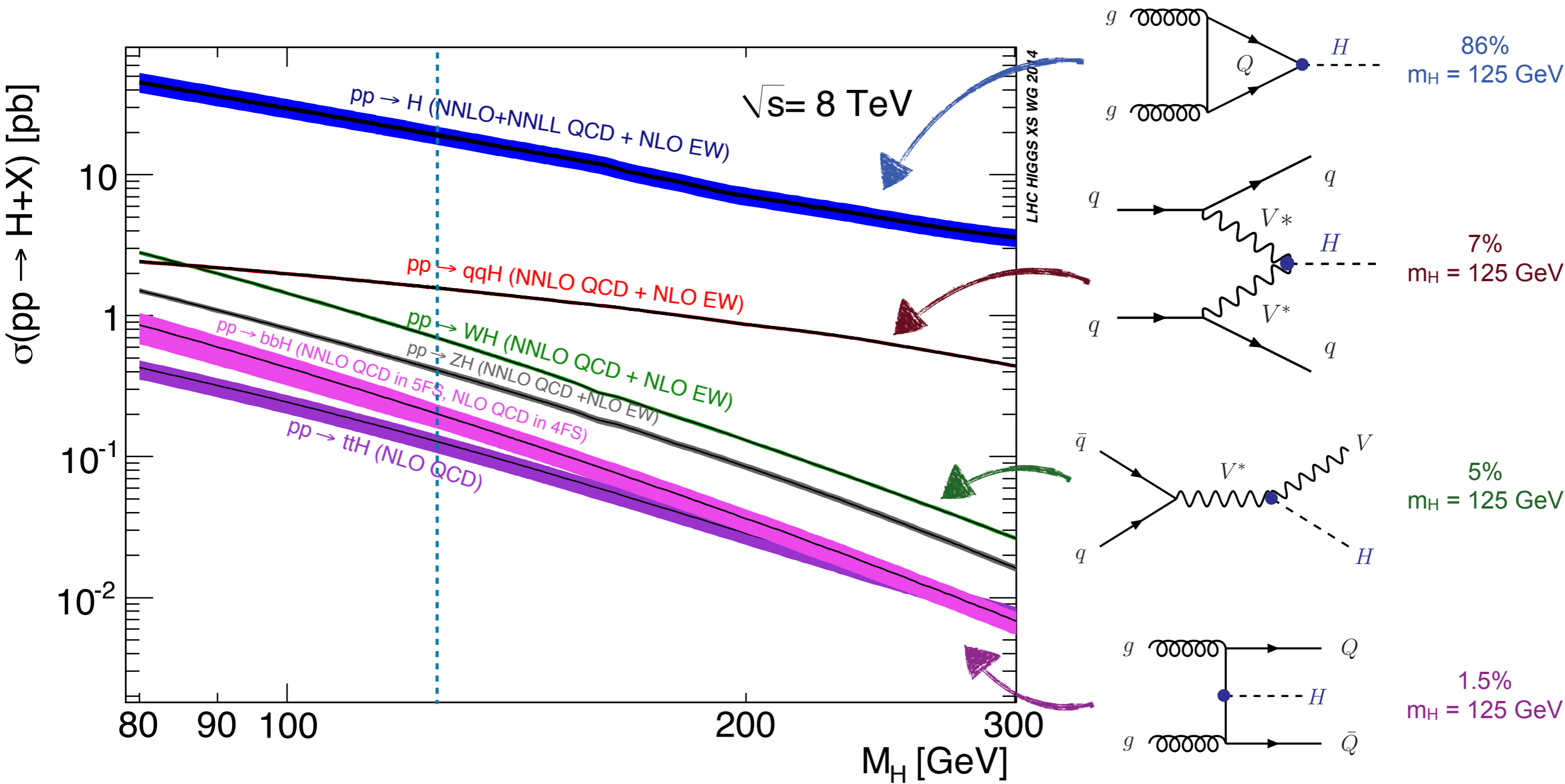


Most importantly:
 ingenuity, both from both theory and experiment, will be crucial to achieve such an enhancement of the LHC physics potential

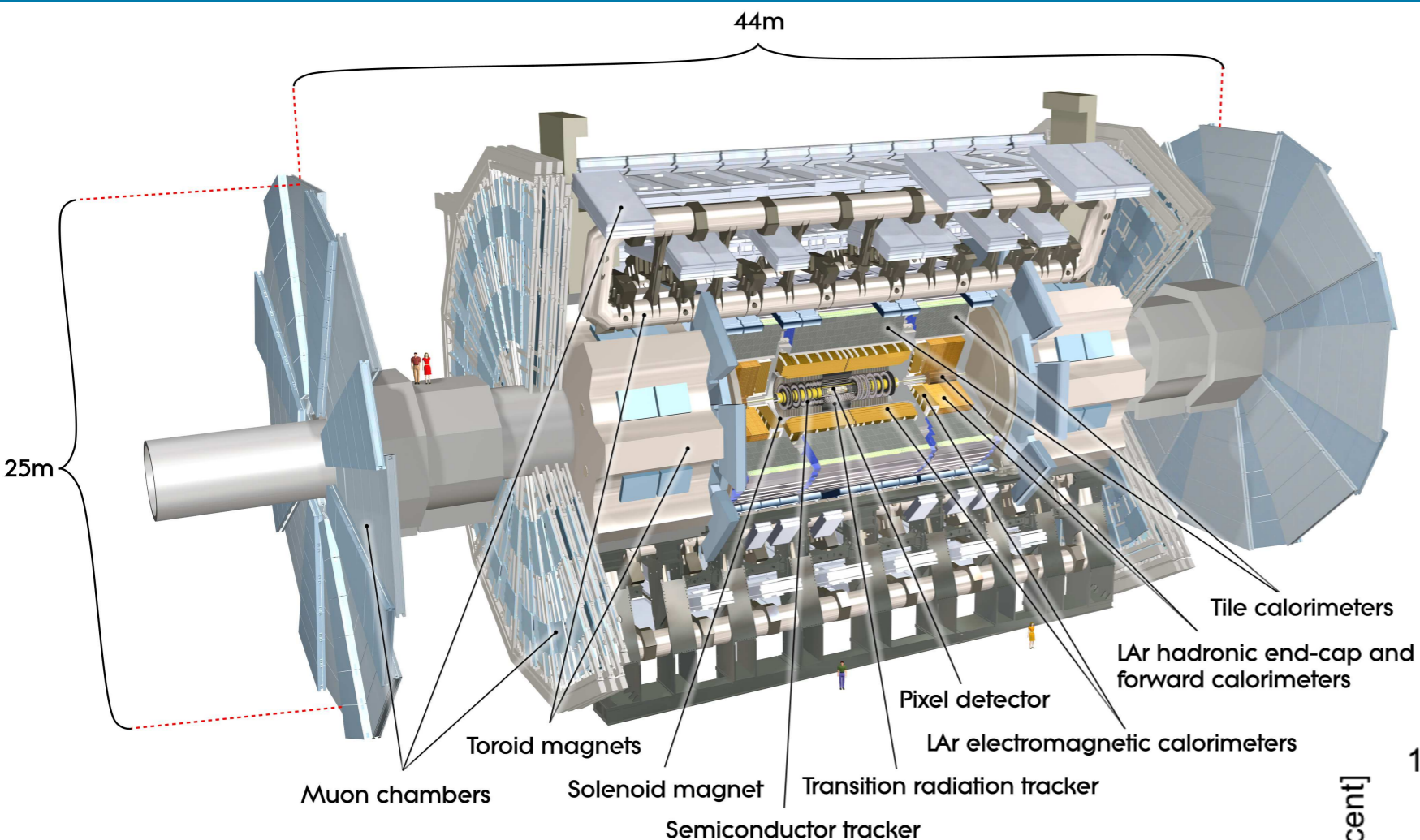
Additional Slides



SM Higgs boson production at the LHC



A Toroidal LHC Apparatus



	ATLAS
Magnets	2T solenoid, 3 air-core toroids
Tracking	silicon + transition radiation tracker
EM Calorimetry	sampling Liquid Argon
Hadron Calorimetry	plastic scintillator (barrel) Liquid Argon(endcap)
Muon	independent system with trigger capabilities
Trigger	3 Level Implementation from 40 MHz to 400 Hz

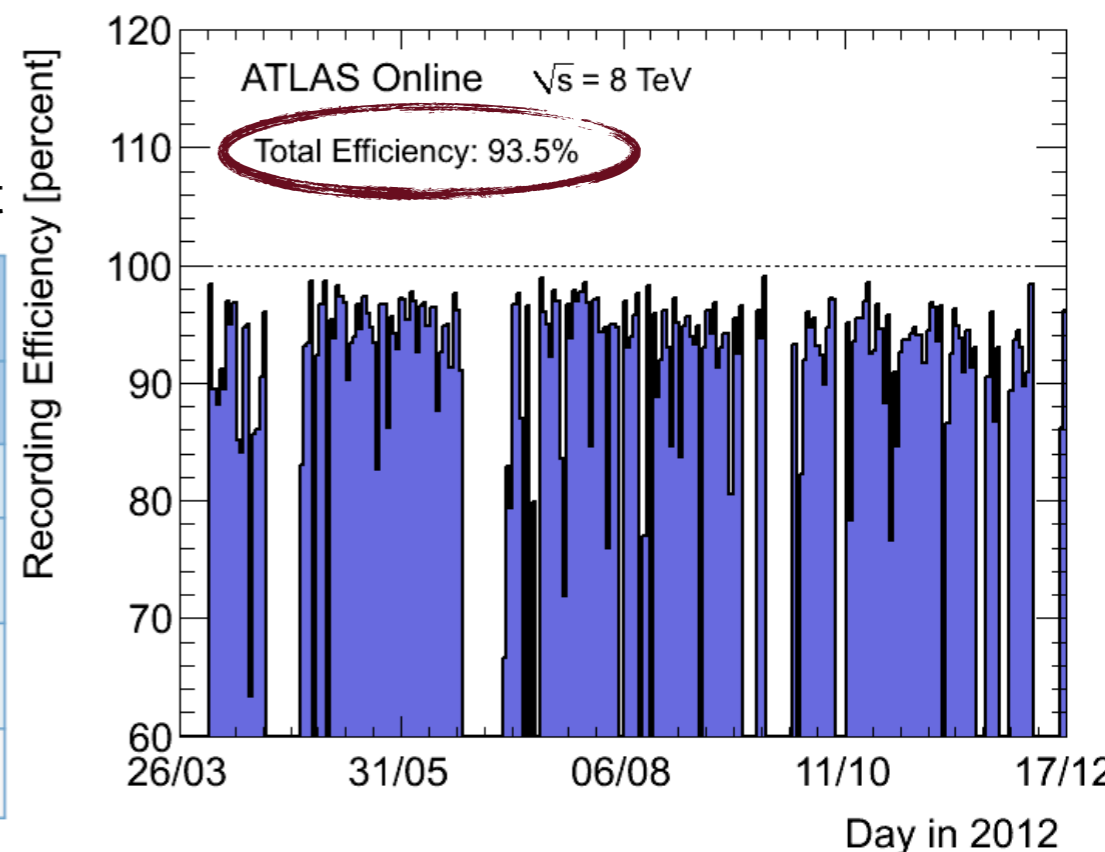
⇒ General purpose detector designed for the harsh LHC environment

ATLAS p-p run: April-December 2012

Inner Tracker			Calorimeters		Muon Spectrometer				Magnets	
Pixel	SCT	TRT	LAr	Tile	MDT	RPC	CSC	TGC	Solenoid	Toroid
99.9	99.4	99.8	99.1	99.6	99.6	99.8	100.	99.6	99.8	99.5

All good for physics: 95.8%

Luminosity weighted relative detector uptime and good quality data delivery during 2012 stable beams in pp collisions at $\sqrt{s}=8$ TeV between April 4th and December 6th (in %) – corresponding to 21.6 fb⁻¹ of recorded data.

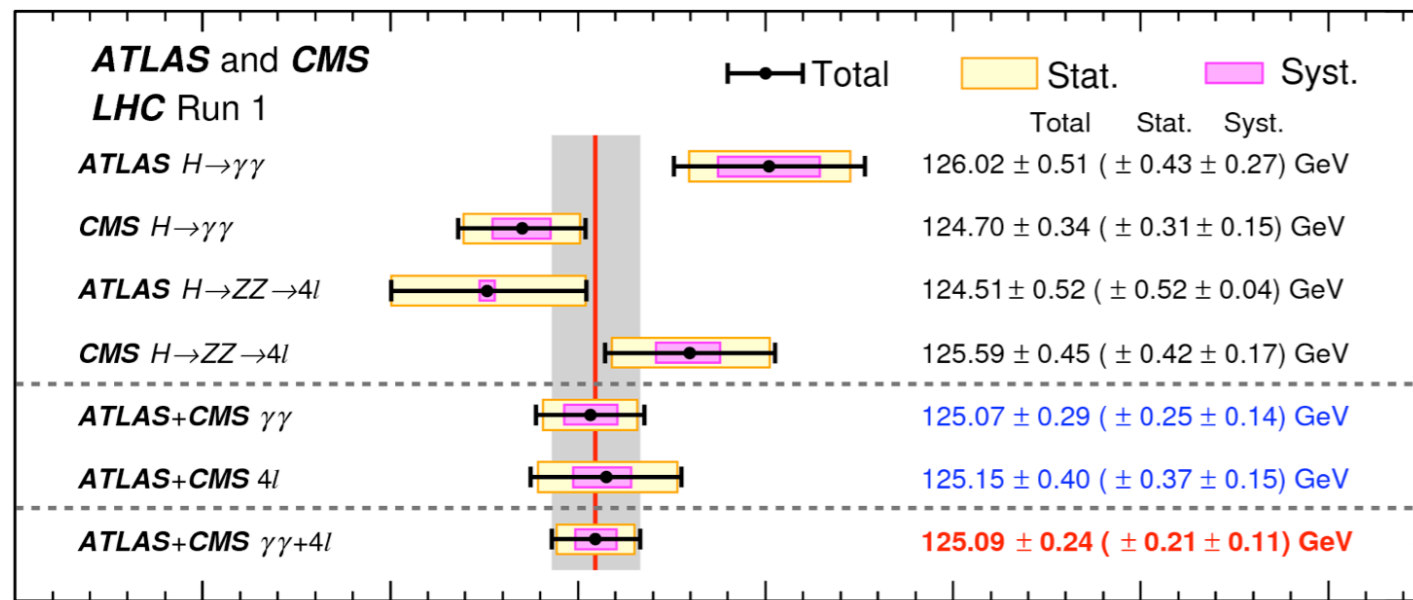


The Standard Model Higgs boson

Substantial progress made in understanding its nature

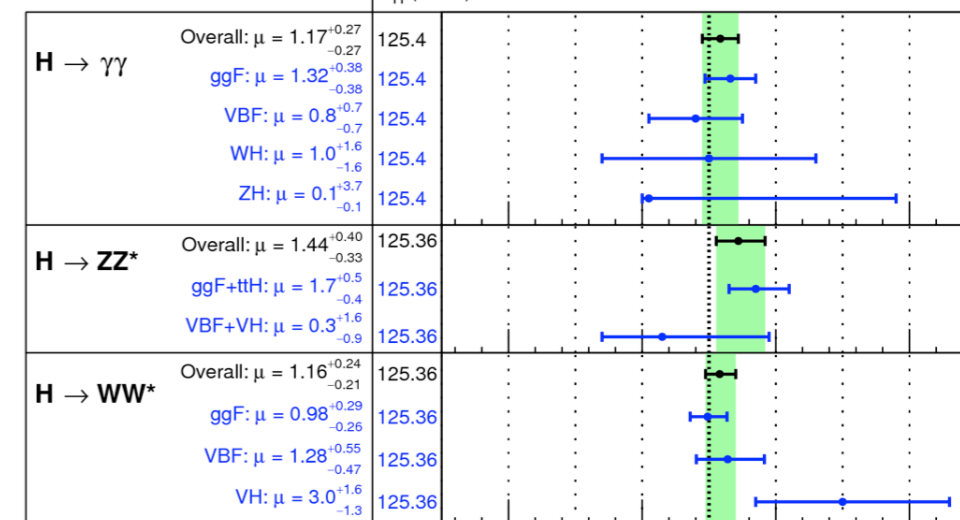
ATLAS-CMS Higgs boson mass combination; 0.19% precision

- ATLAS measurement: $125.36^{+0.37}_{-0.15} \text{ (stat)}^{+0.18}_{-0.15} \text{ (syst)} \text{ GeV}$
- CMS measurement: $125.02^{+0.26}_{-0.27} \text{ (stat)}^{+0.14}_{-0.15} \text{ (syst)} \text{ GeV}$

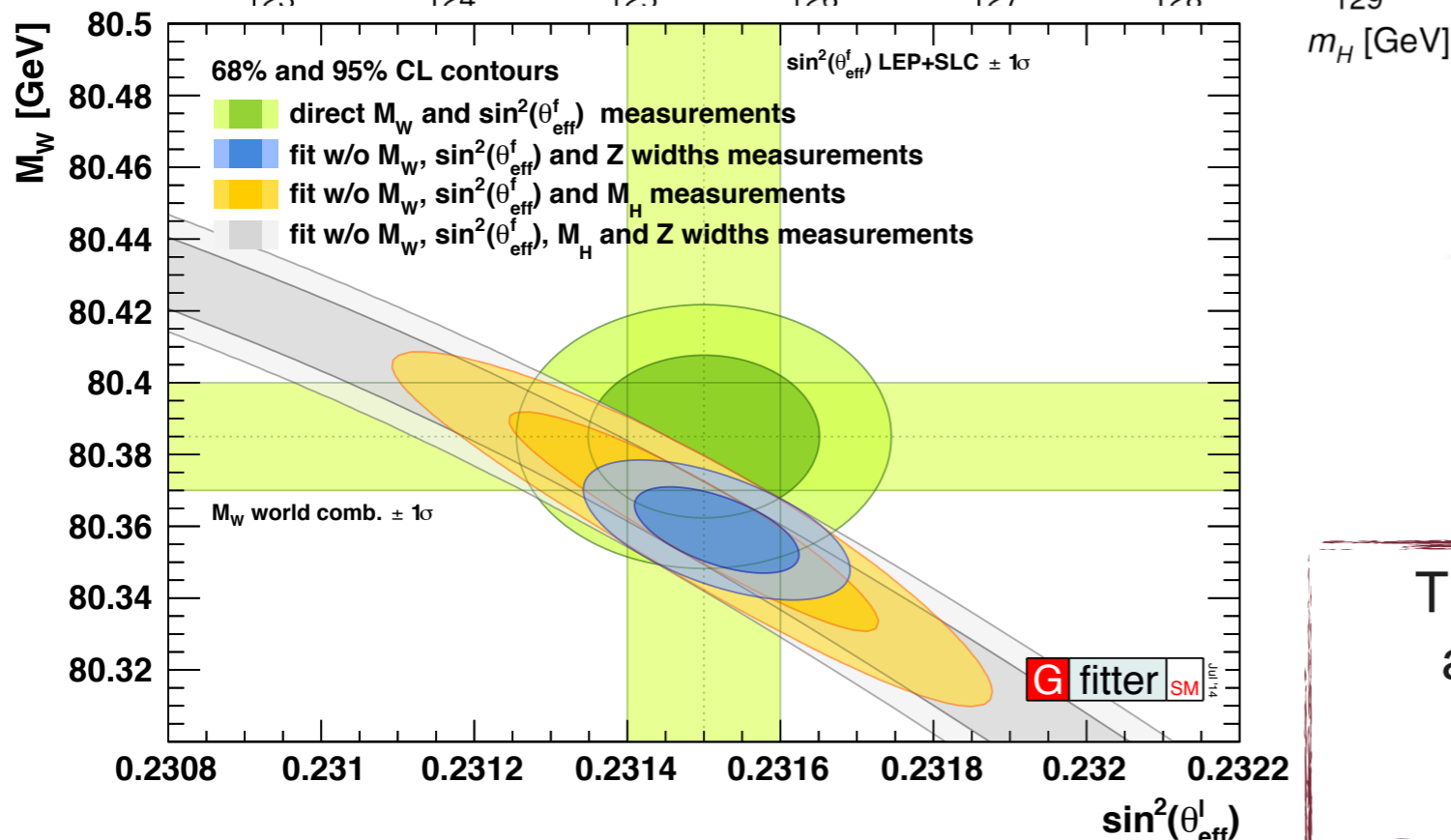


ATLAS Preliminary

$m_H = 125.36 \text{ GeV}$



HVV (V=W,Z, γ) coupling probed via observed rates; SM rates agree with data Precision at ~20-40% Systematic uncertainties becoming important

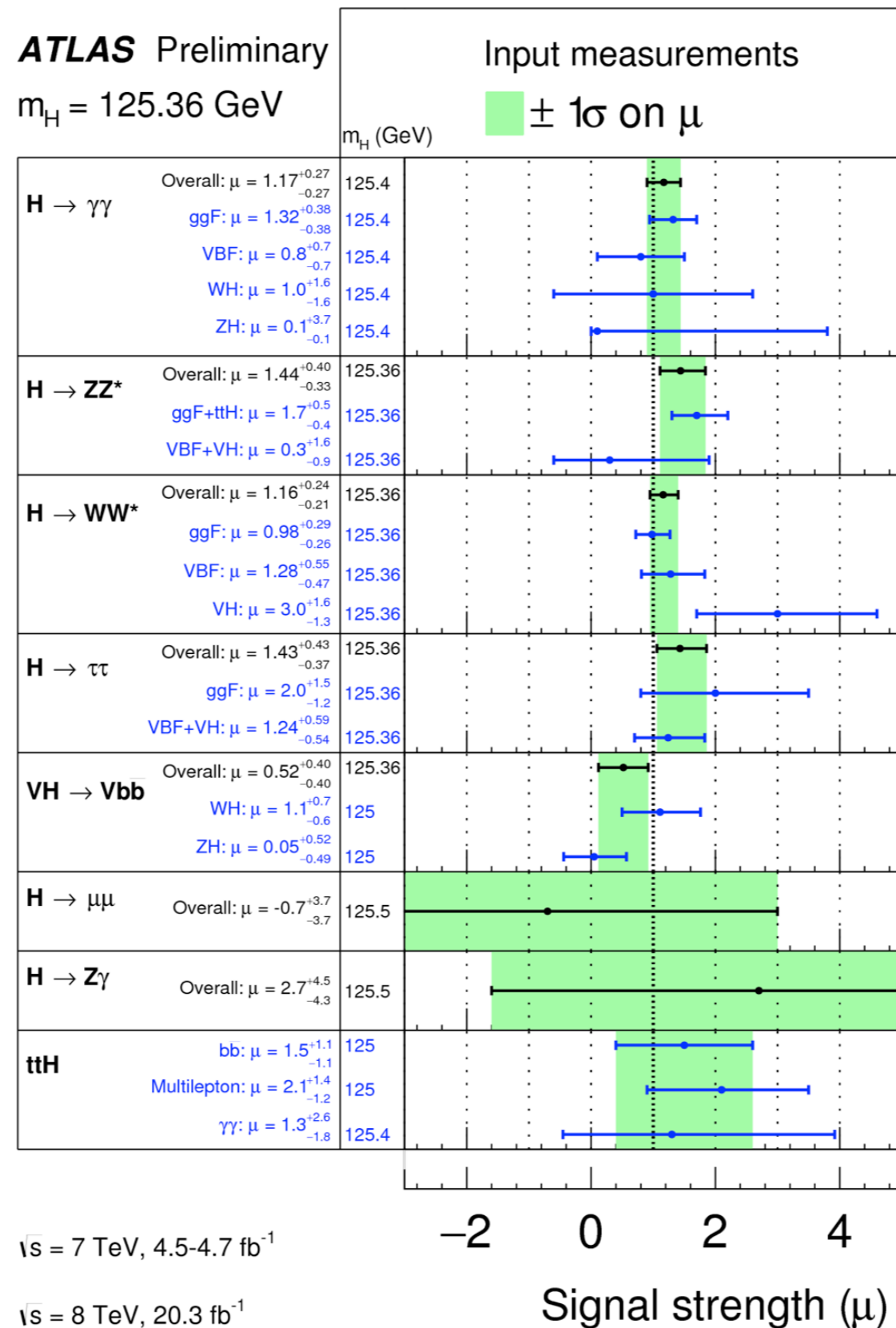


$\sqrt{s} = 7 \text{ TeV}, 4.5\text{-}4.7 \text{ fb}^{-1}$

$\sqrt{s} = 8 \text{ TeV}, 20.3 \text{ fb}^{-1}$

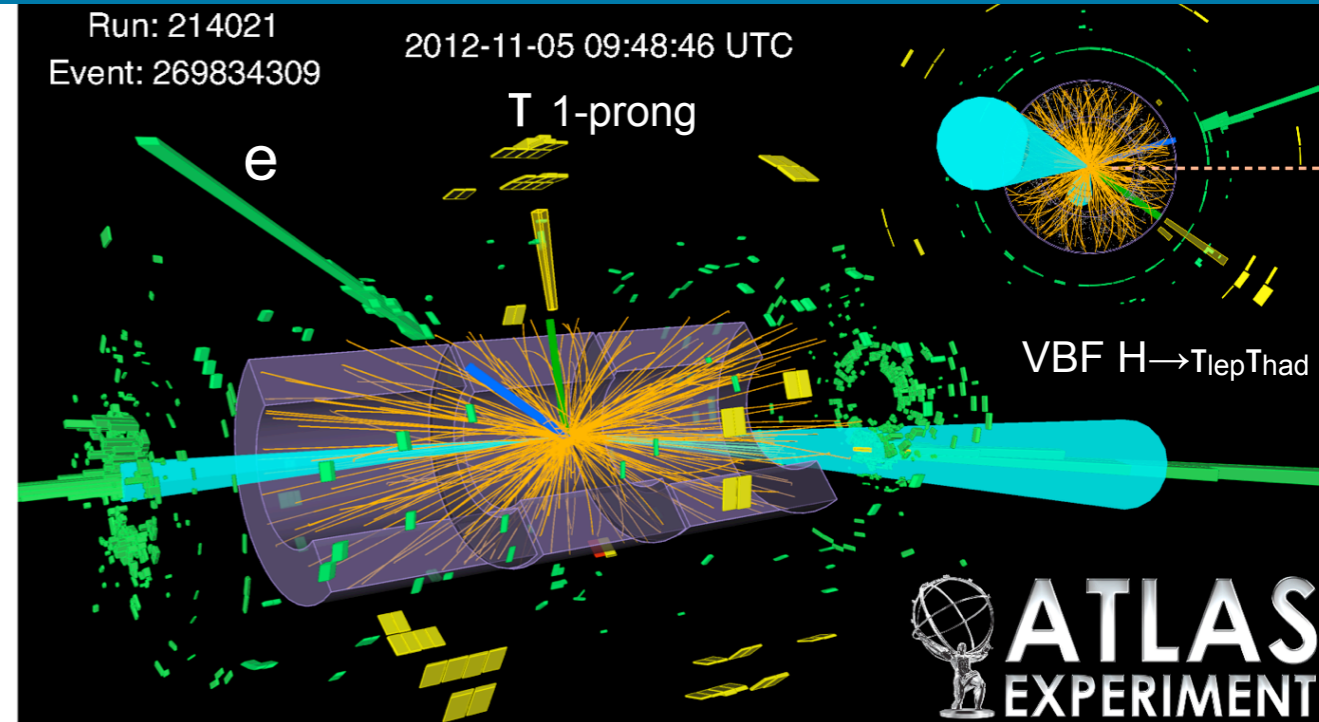
This was not unexpected, given the level of agreement of the SM predictions with the observed precision electroweak data (even w/o knowledge of m_H)

Overview of rate measurements

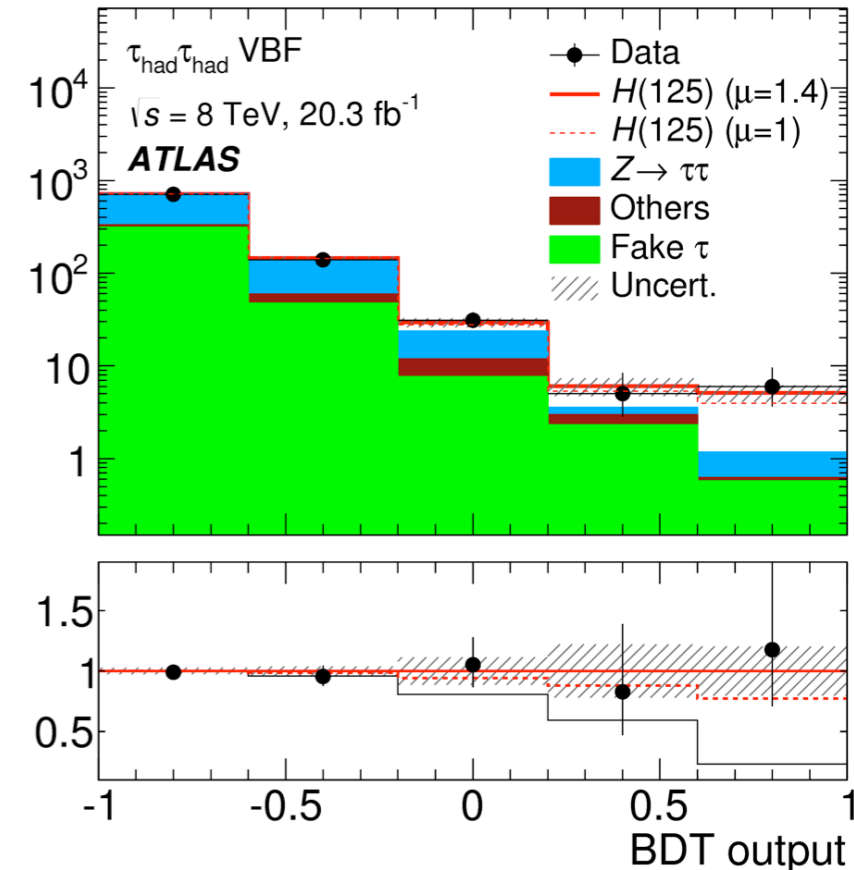
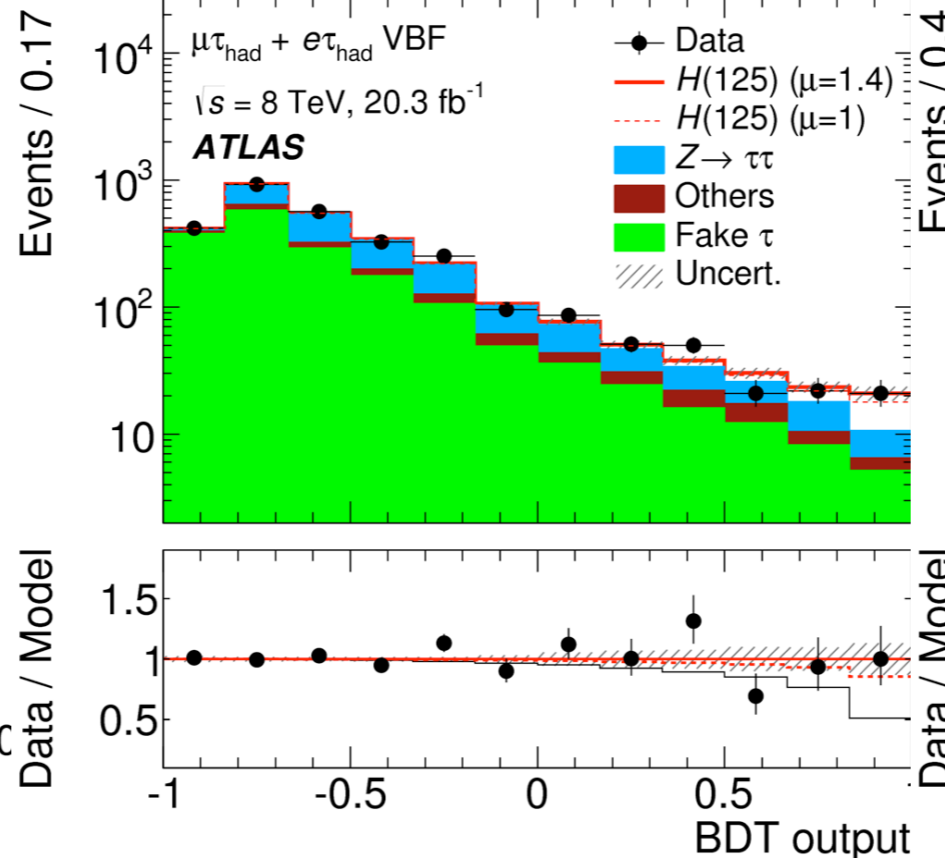
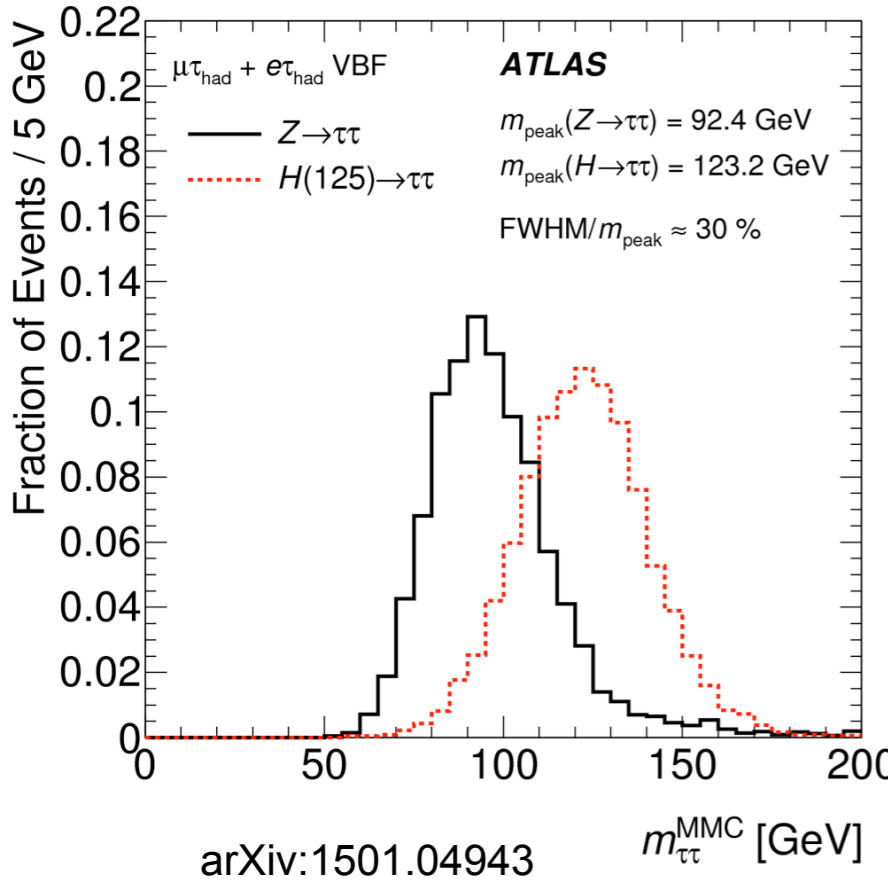


H → ττ

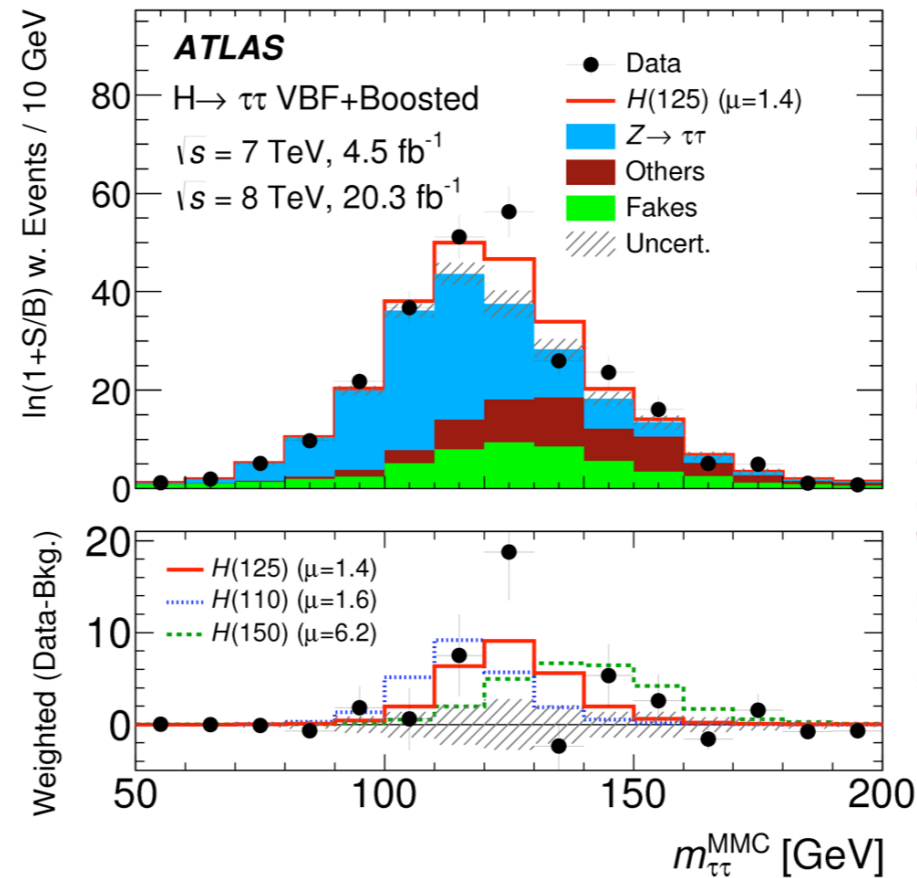
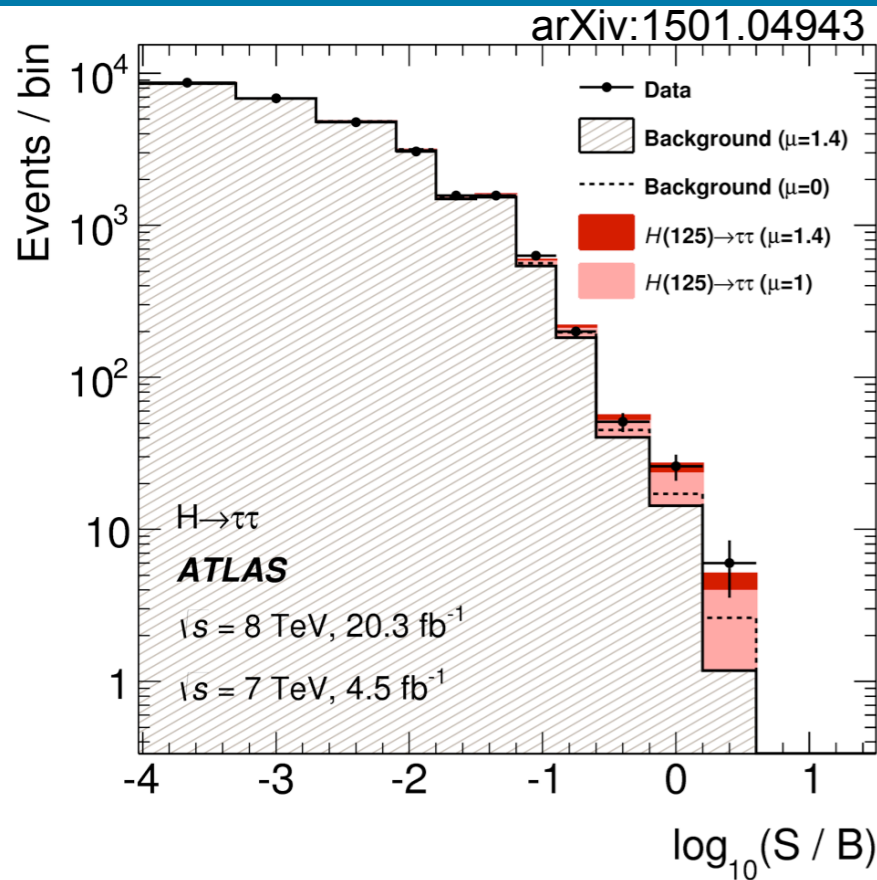
- Most promising for down-type fermion/lepton couplings
- Backgrounds
 - Z → ττ dominant [embedding]
 - “Fakes”: Multijet, W+jets, top [data-driven]
 - “Other”: Dibosons/H→WW* [MC]
- Three sub-channels: TlepTlep, TlepThad, ThadThad
- Two exclusive categories/final state: *VBF* (2 jets with large Δη) and *Boosted* (large di-tau pT)
- BDT for each category: *di-tau properties* (m_{ττ}, ΔR_{ττ}, ...), *jet topology* (m_{jj}, Δη_{jj}, ...), *event activity/topology* (scalar/vector pT sum, object centralities, ...)



ep_T = 56 GeV, T_{had} p_T = 27 GeV, MET=113 GeV, m_{j1,j2}=1.53 TeV, m_{ττ}^{MMC}=129 GeV, BDT score = 0.99. S/B ratio of this bin 1.0



H → ττ: Results



Evidence observed for Higgs boson decays to τ -leptons significance at 125 GeV

ATLAS: 4.5σ (3.4σ)

CMS: 3.2σ (3.5σ)

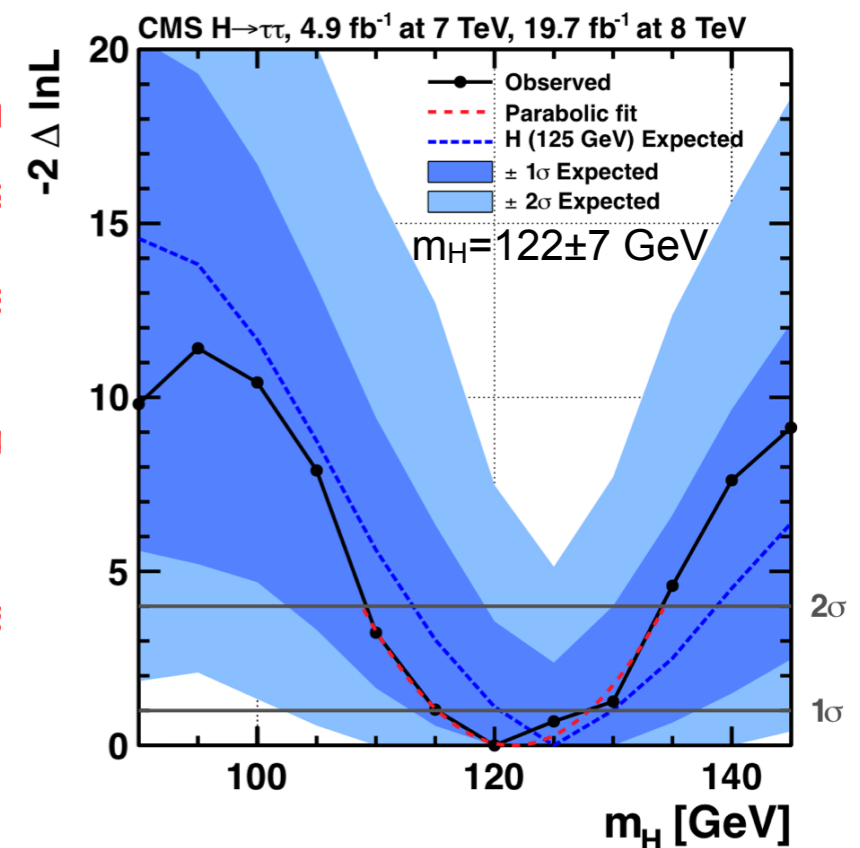
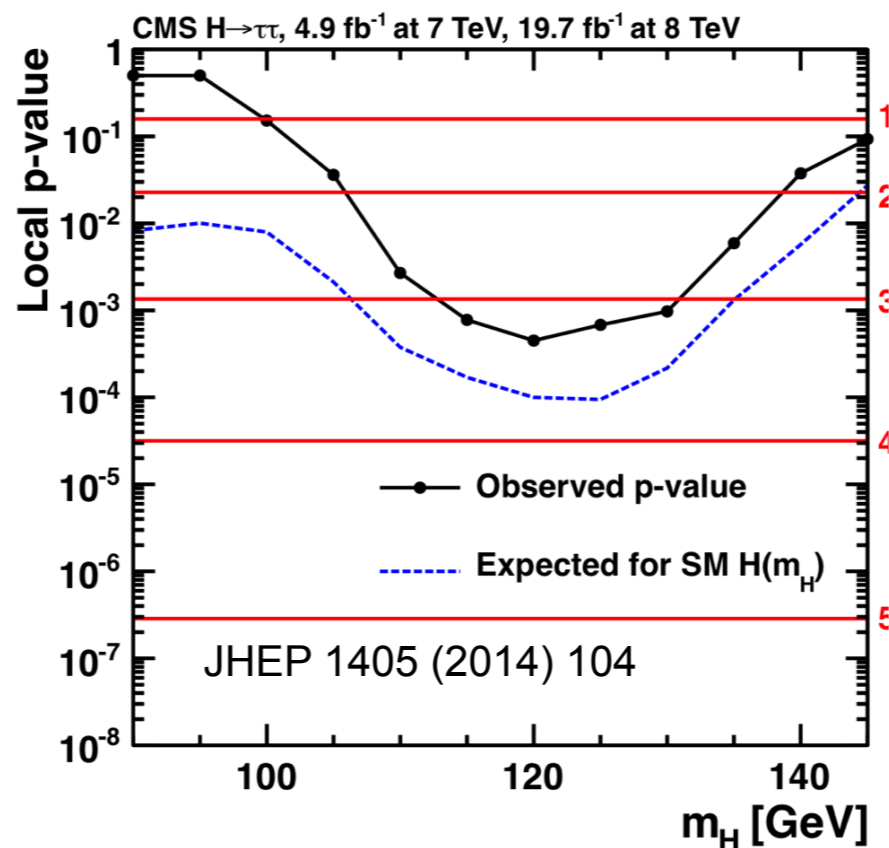
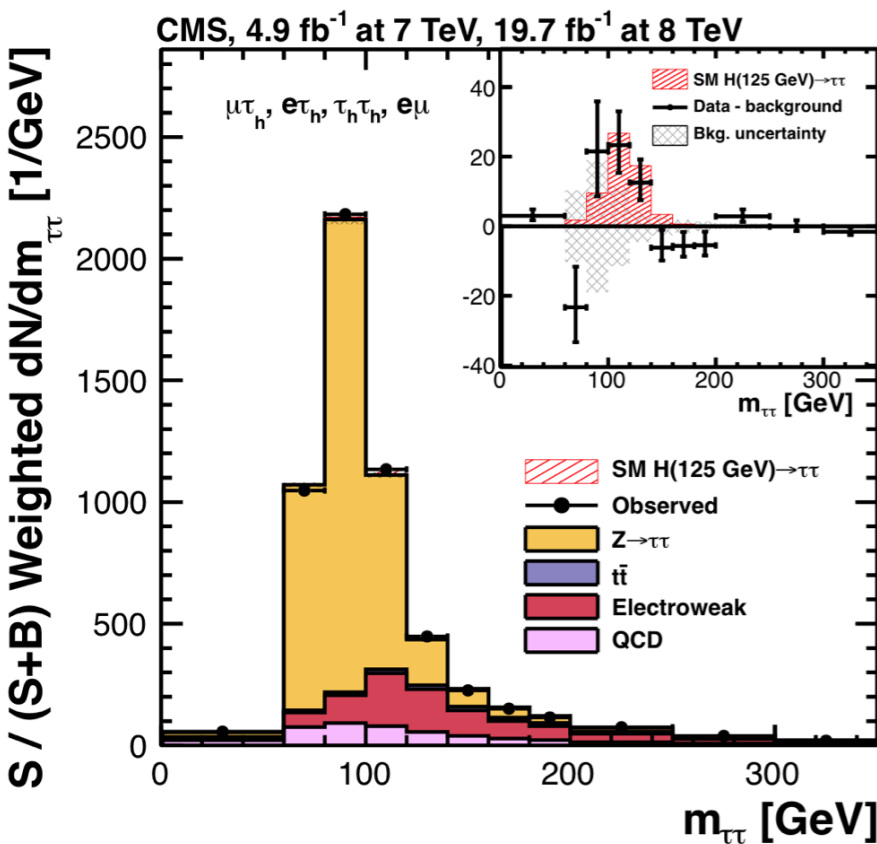
Rate measurement

ATLAS $\mu = 1.43^{+0.27}_{-0.26}(\text{stat})$
 $+0.32_{-0.25}(\text{syst}) \pm 0.09$ (theo syst)

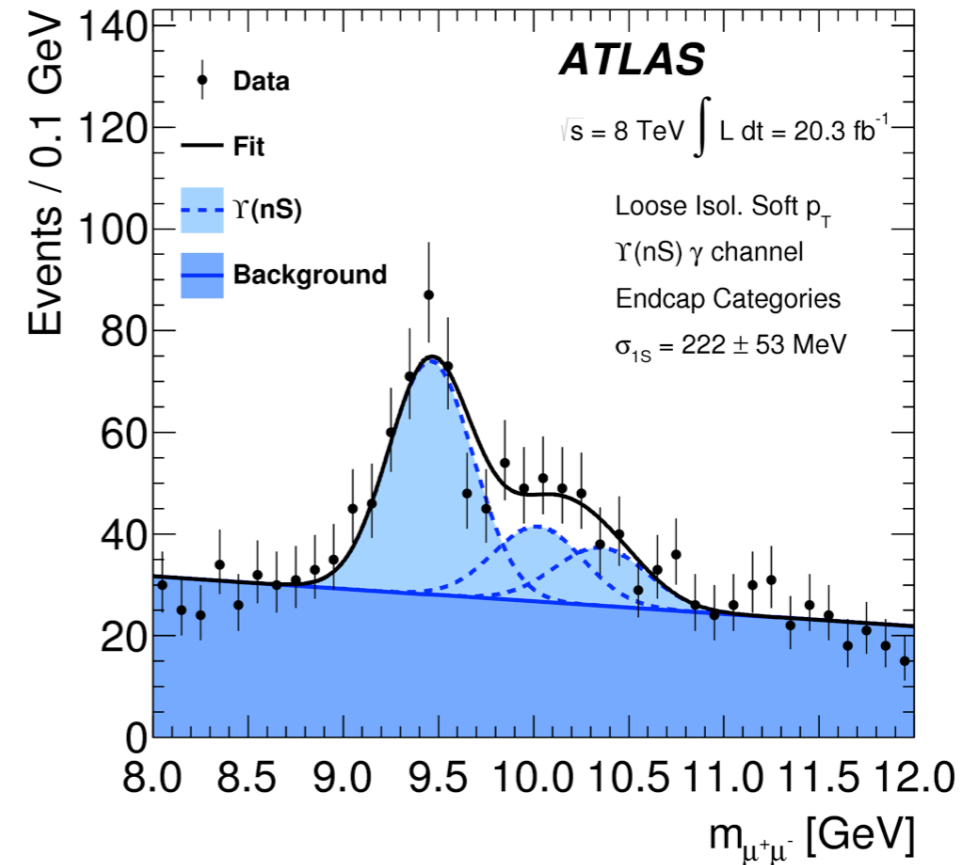
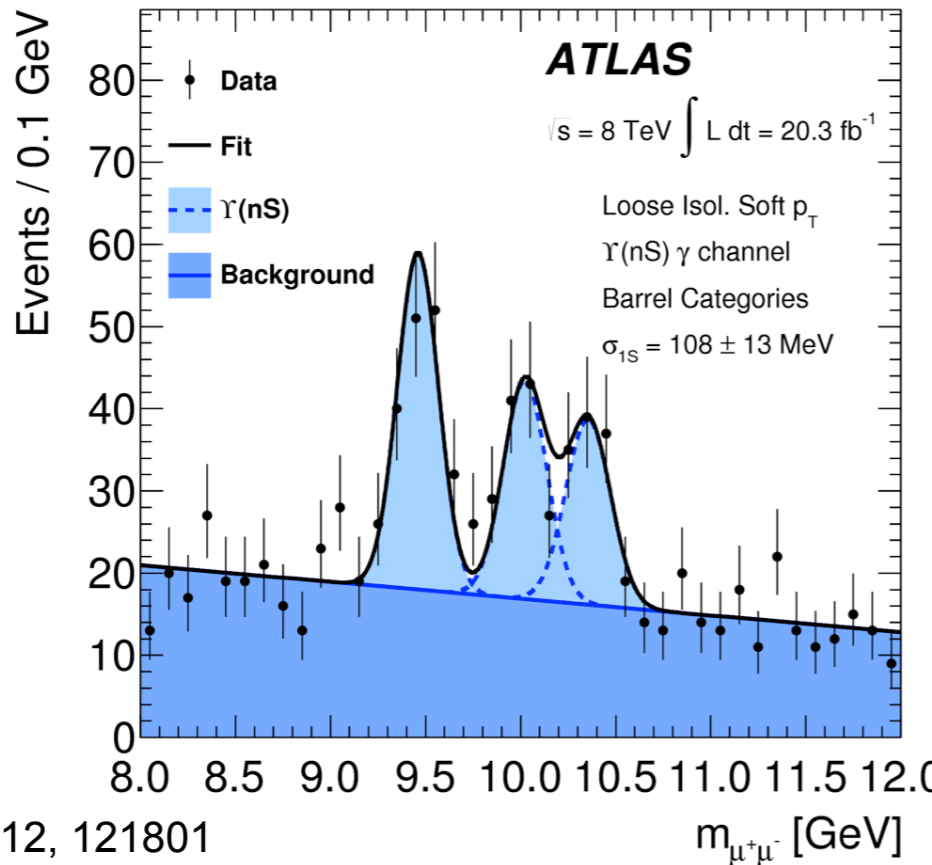
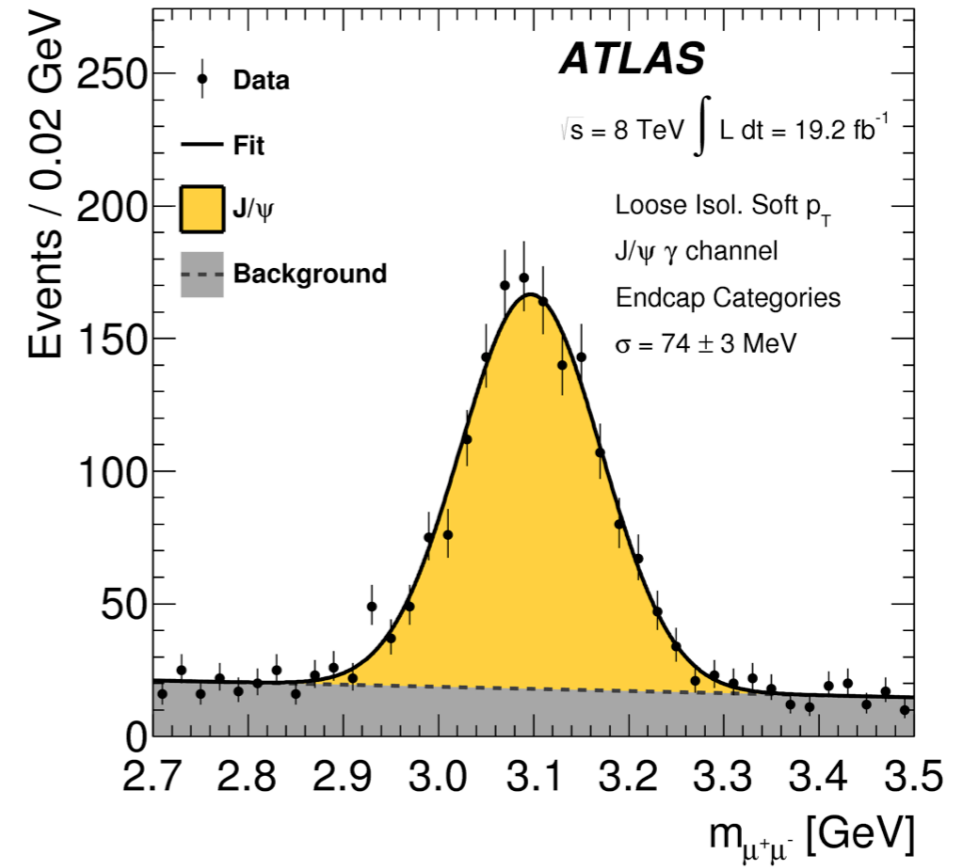
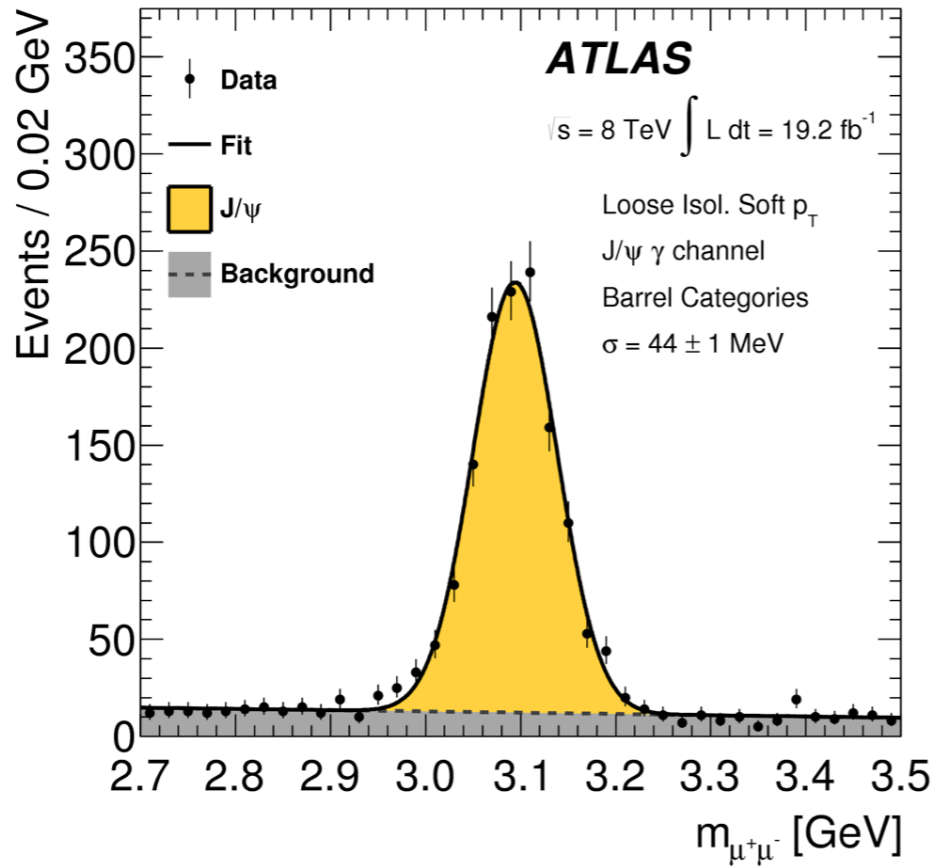
[@ $m_H = 125.36 \text{ GeV}$]

CMS $\mu = 0.78 \pm 0.27$

[@ $m_H = 125.00 \text{ GeV}$]

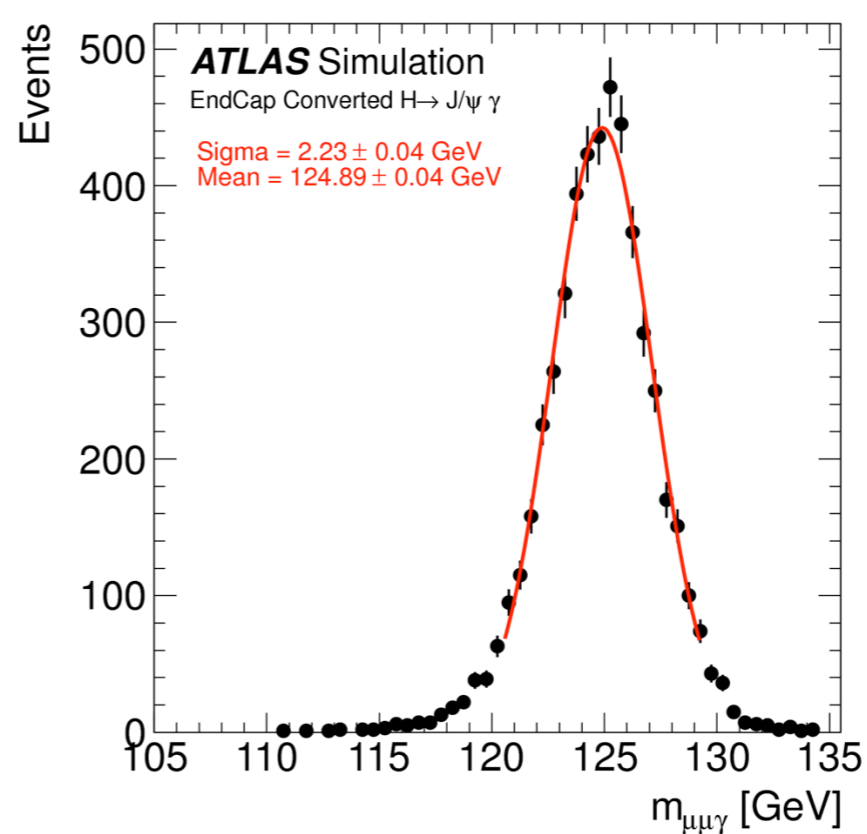
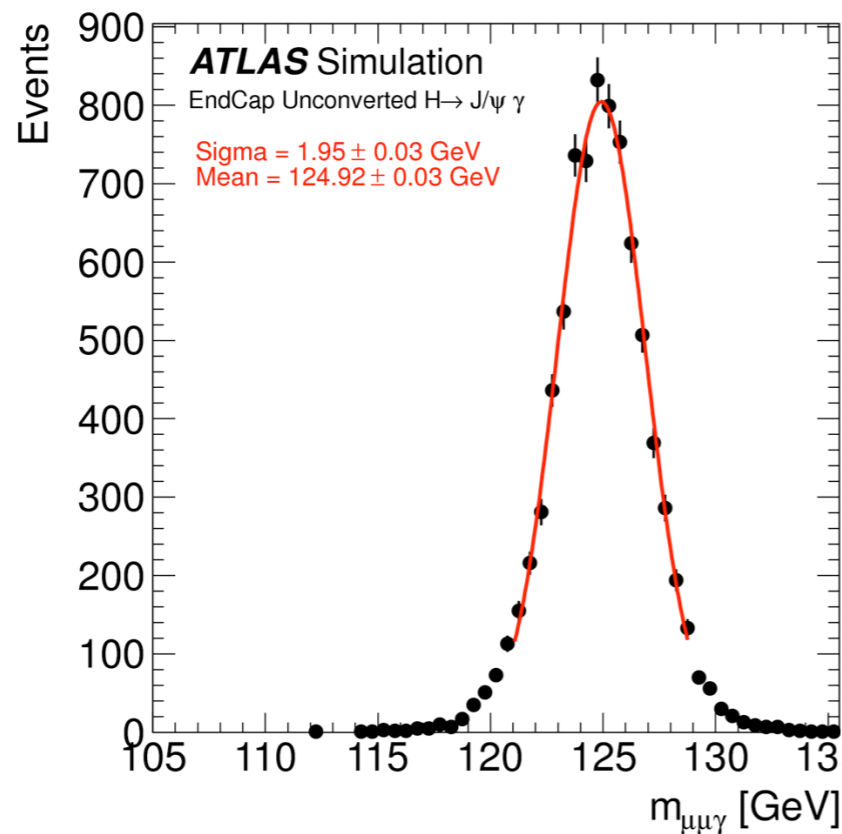
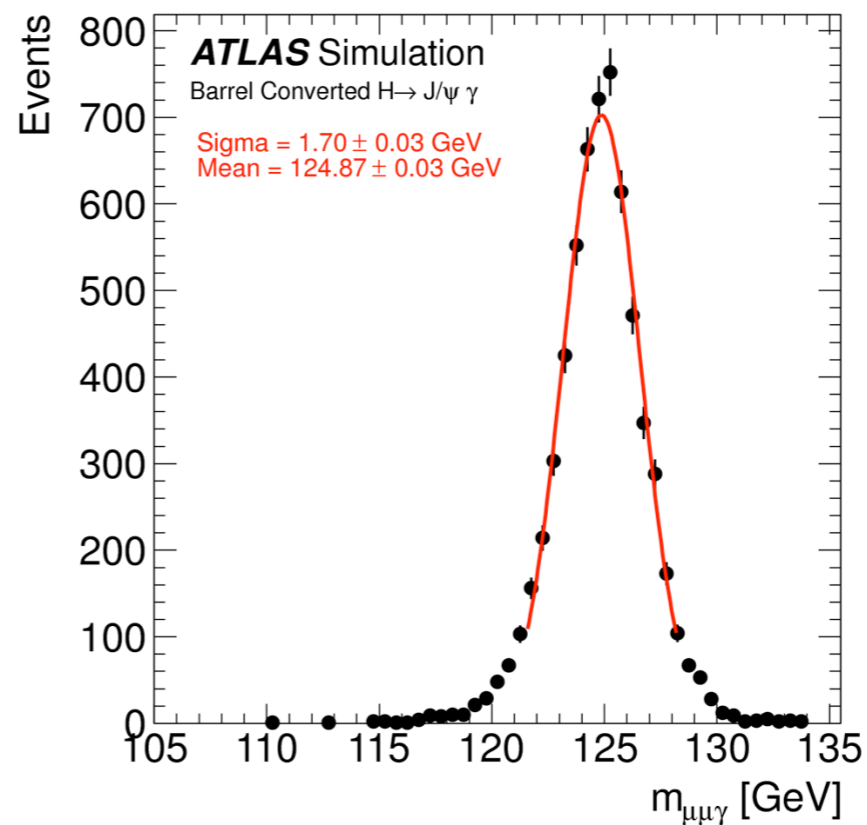
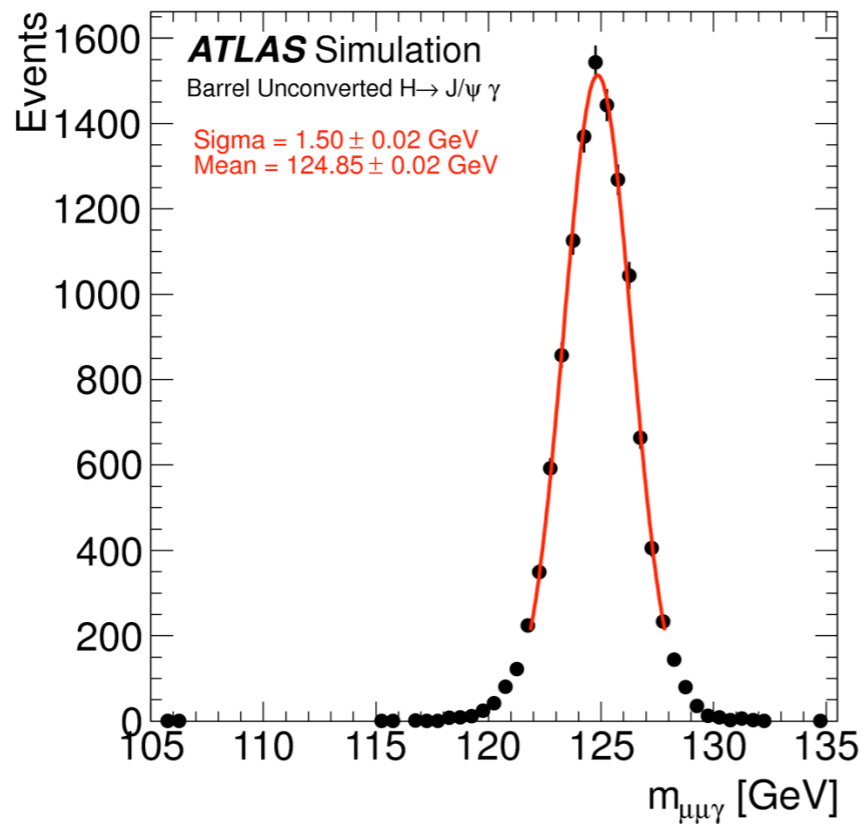


H → J/ψγ and H → Υ(ns)γ



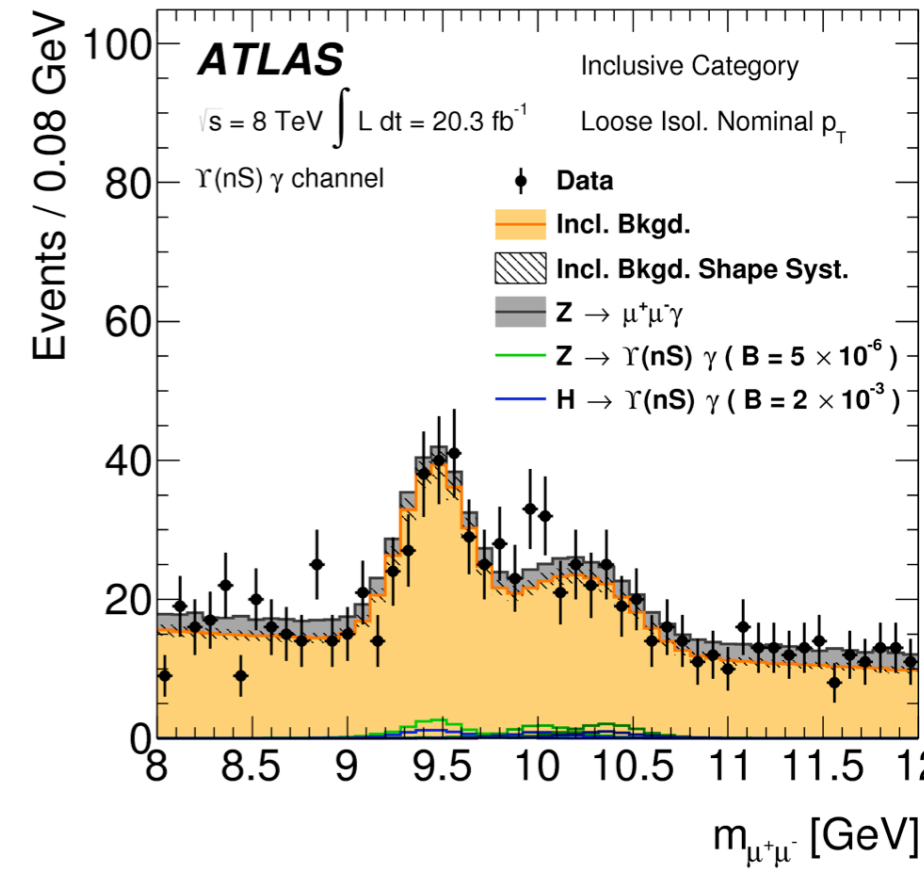
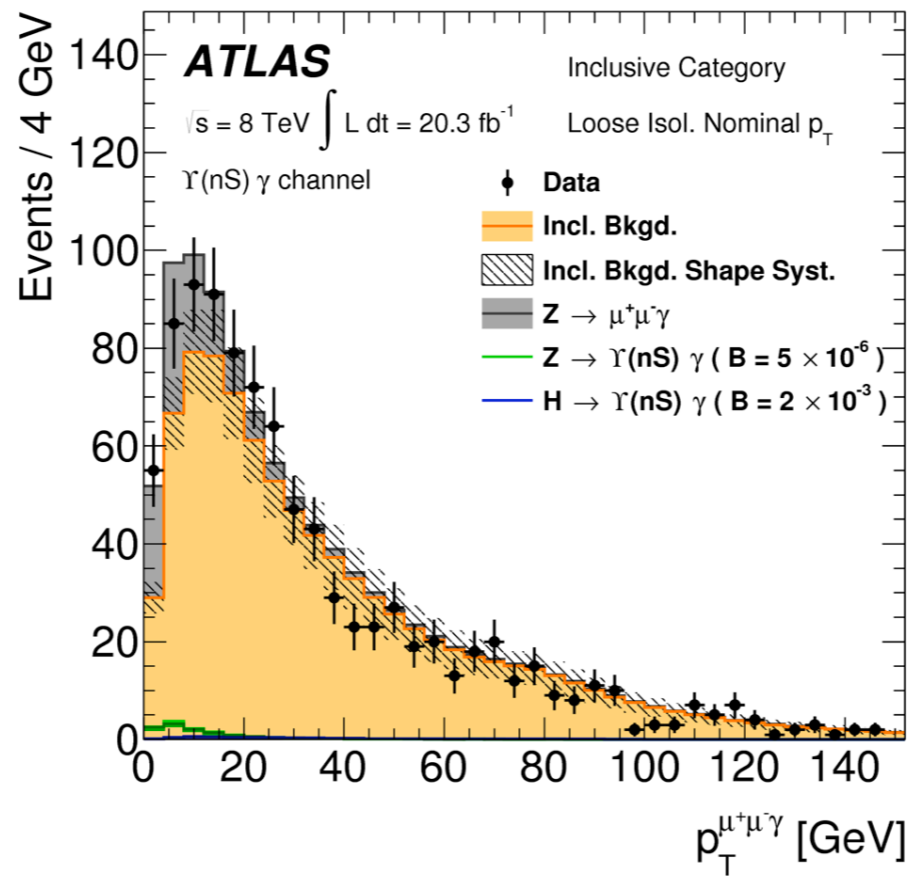
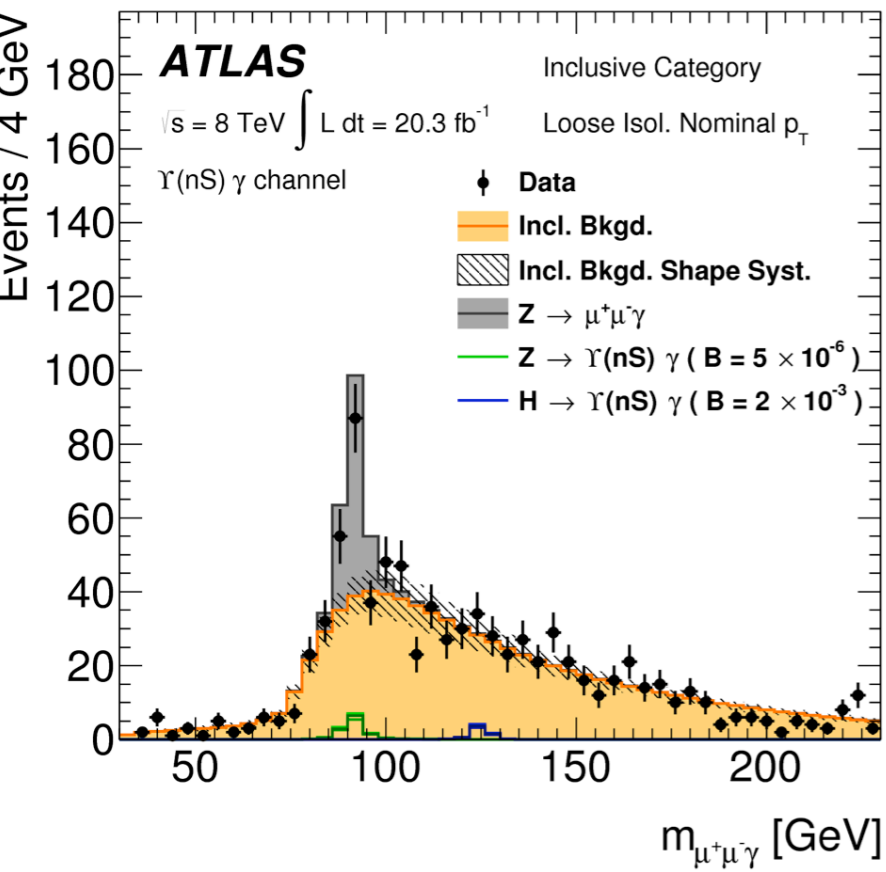
$H \rightarrow J/\psi \gamma$ and $H \rightarrow Y(ns) \gamma$

Phys.Rev.Lett. 114 (2015) 12, 121801

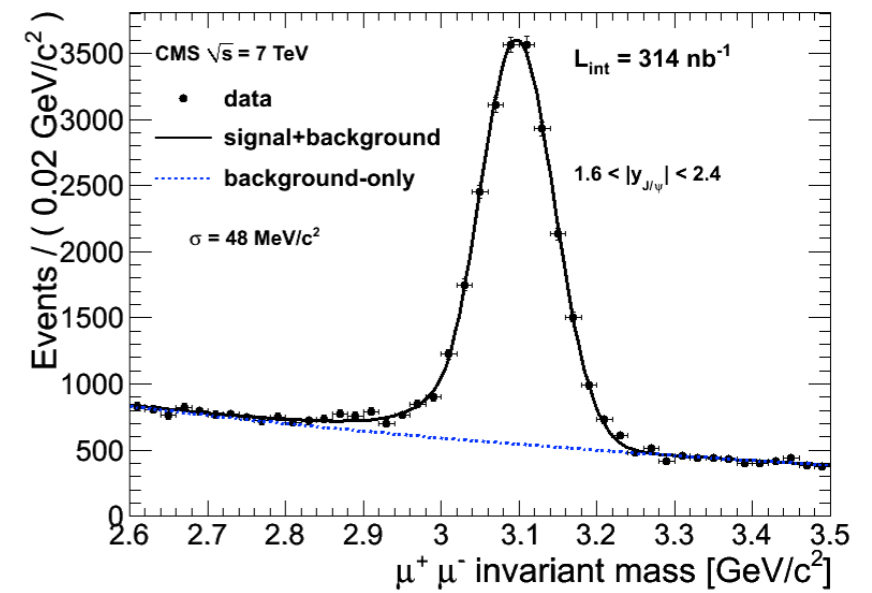
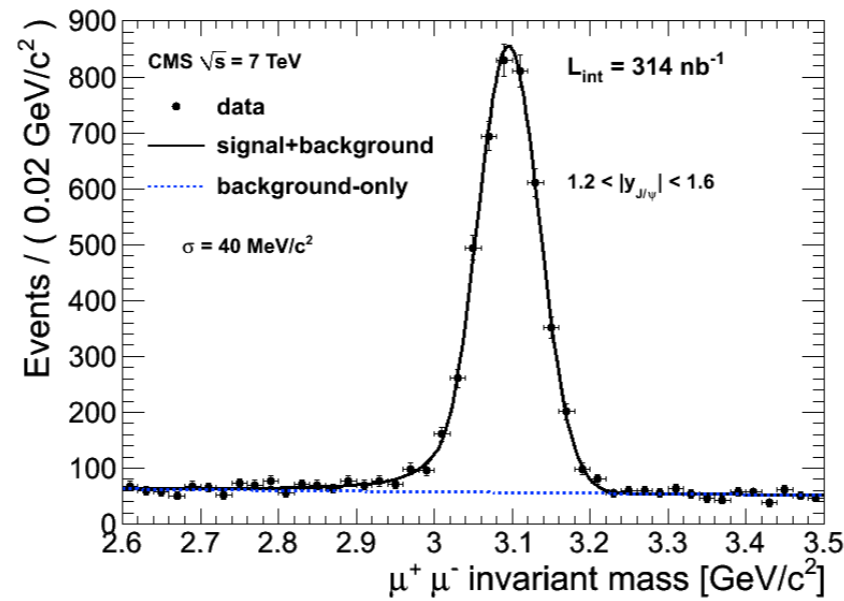
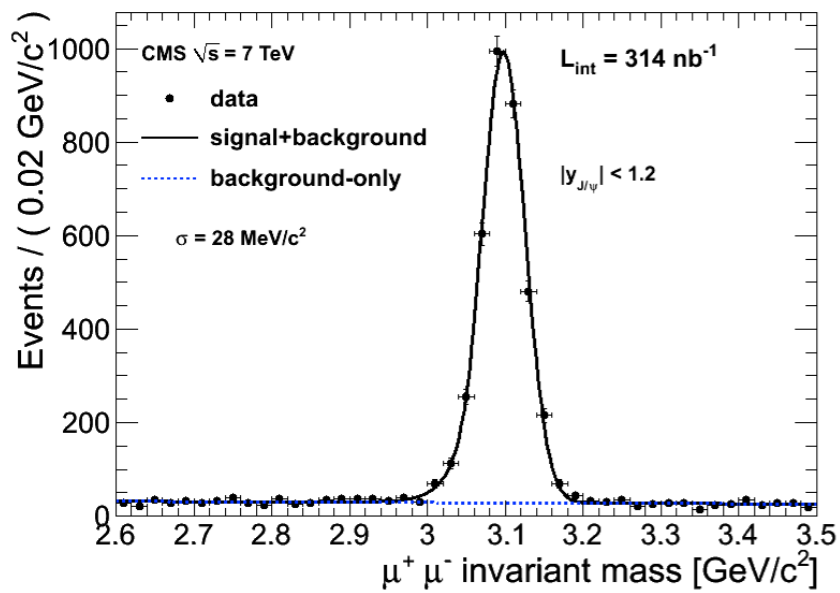
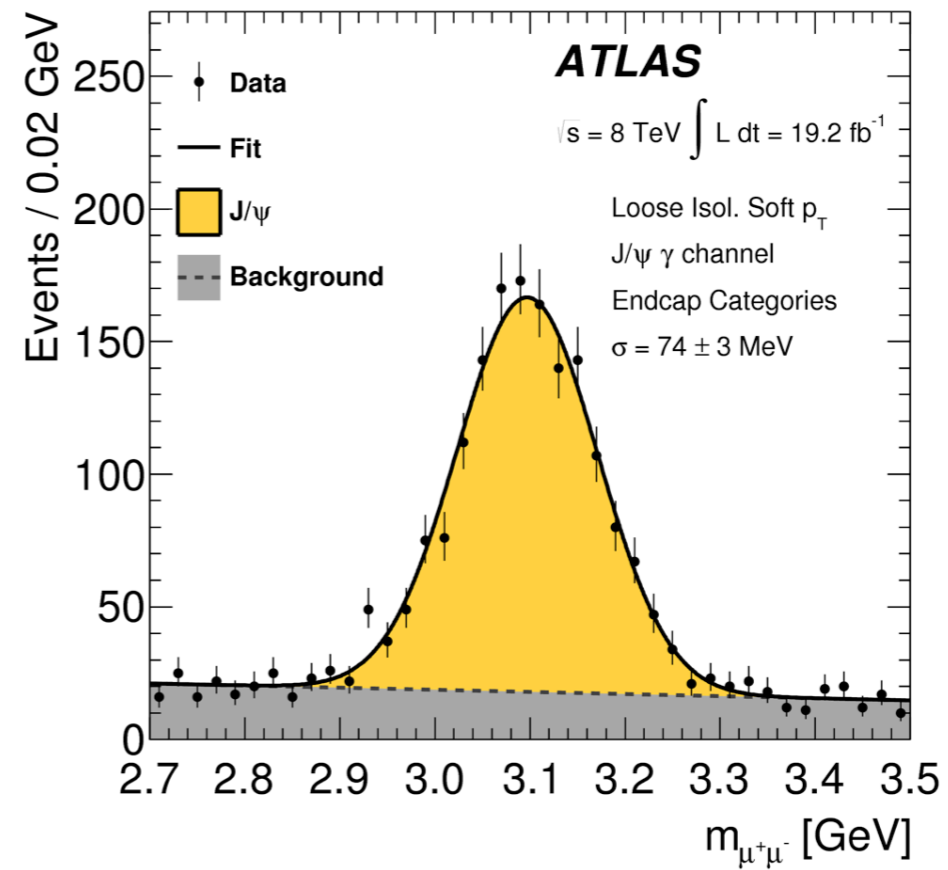
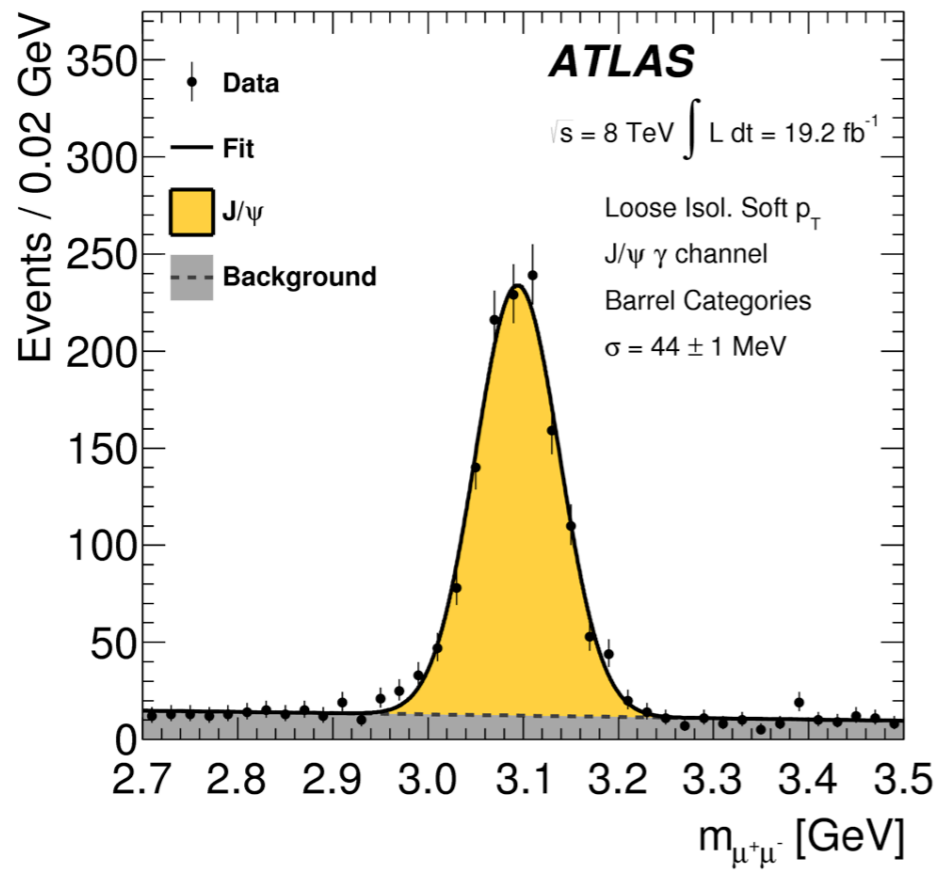


H → J/ψγ and H → Y(ns)γ

Phys.Rev.Lett. 114 (2015) 12, 121801



H → J/ψγ and H → Y(ns)γ



Light-Quark Yukawa couplings

This was also considered impossible for the LHC. Recent activity on its feasibility:

→ **Exploit the exclusive decays $H \rightarrow Q\gamma$ as direct probe to the quark Yukawa couplings**

[Phys.Rev.Lett. 114 (2015) 10, 101802]

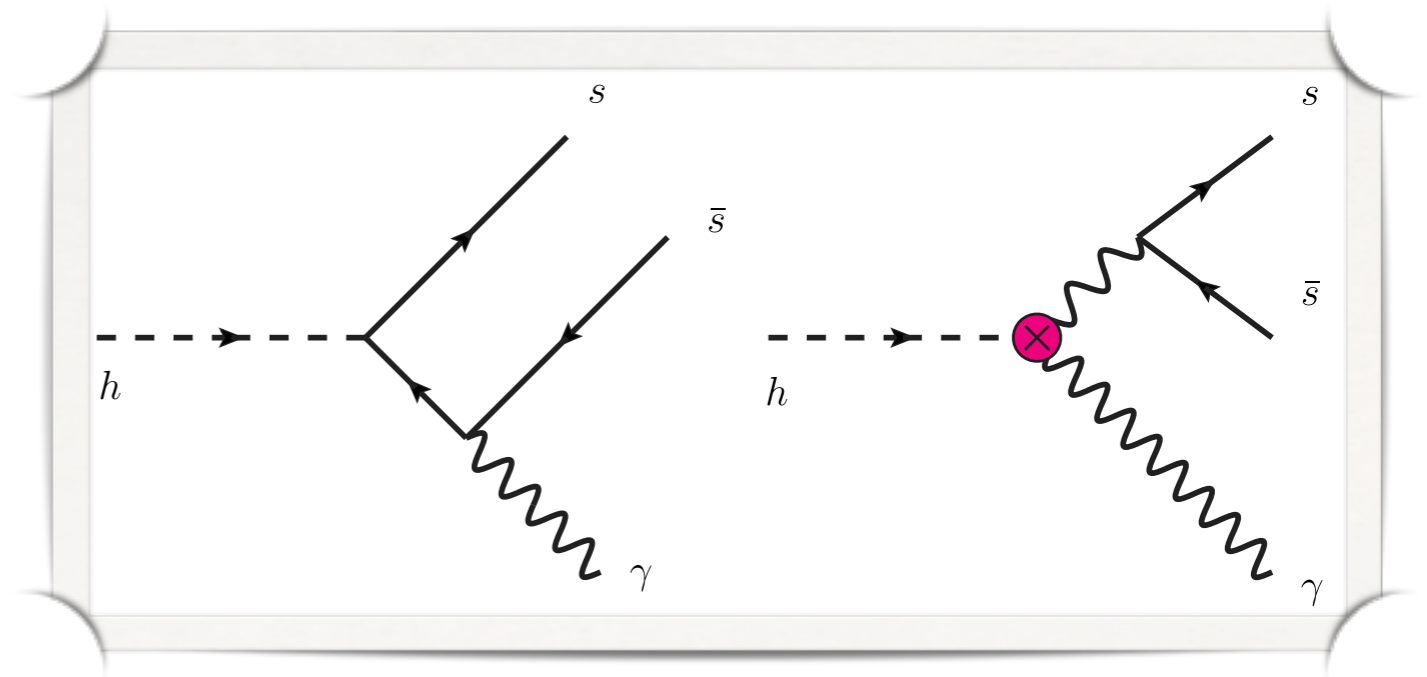
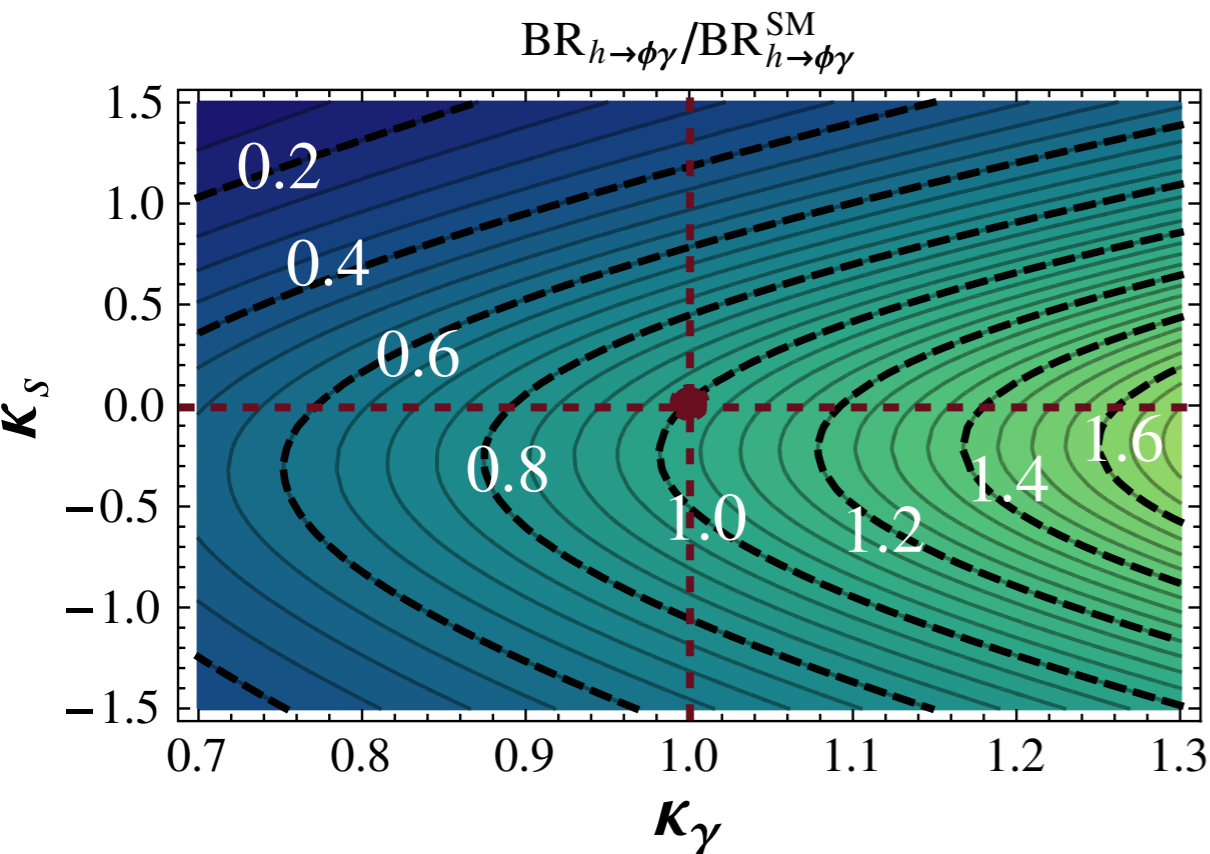
→ **Sensitive to BSM physics** [Phys. Rev. D 80, 076002, Phys. Lett. B665 (2008) 79, Phys.Rev. D90 (2014) 115022]

The idea is to benefit from the interference of the “direct” and “indirect” amplitudes!

$$\frac{\text{BR}_{h \rightarrow \phi\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(3.0 \pm 0.3)\kappa_\gamma - 0.78\bar{\kappa}_s] \times 10^{-6}}{0.57\bar{\kappa}_b^2},$$

$$\frac{\text{BR}_{h \rightarrow \rho\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(1.9 \pm 0.2)\kappa_\gamma - 0.24\bar{\kappa}_u - 0.12\bar{\kappa}_d] \times 10^{-5}}{0.57\bar{\kappa}_b^2},$$

$$\frac{\text{BR}_{h \rightarrow \omega\gamma}}{\text{BR}_{h \rightarrow b\bar{b}}} = \frac{\kappa_\gamma [(1.6 \pm 0.2)\kappa_\gamma - 0.59\bar{\kappa}_u - 0.29\bar{\kappa}_d] \times 10^{-6}}{0.57\bar{\kappa}_b^2},$$



Phys.Rev.Lett. 114 (2015) 10, 101802

$$\bar{\kappa}_s = m_s/m_b \approx 0.02$$

$$\bar{\kappa}_d = m_d/m_b \approx 1 \cdot 10^{-3}$$

$$\bar{\kappa}_u = m_u/m_b \approx 4.7 \cdot 10^{-4}$$

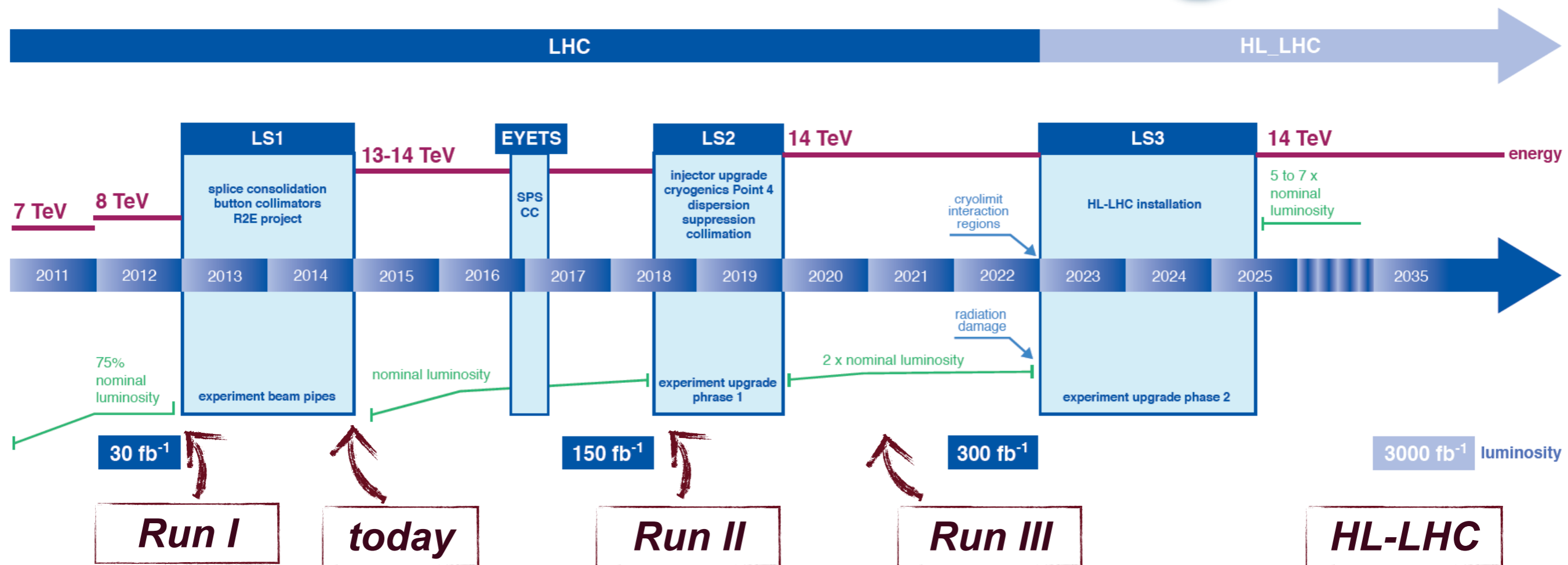
$$BR_{h \rightarrow \phi\gamma}^{\text{SM}} \approx 3 \cdot 10^{-6}$$

$$BR_{h \rightarrow \rho\gamma}^{\text{SM}} \approx 1.9 \cdot 10^{-5}$$

$$BR_{h \rightarrow \omega\gamma}^{\text{SM}} \approx 1.6 \cdot 10^{-6}$$

LHC/HL-LHC Plan

LHC / HL-LHC Plan



- Run II will provide ×5-6 more integrated luminosity compared to Run I
 - Aiming for 3000 fb⁻¹ by 2035
- Experiments will be upgraded ATLAS to go for an new all Si tracker

Higgs in Run II and beyond

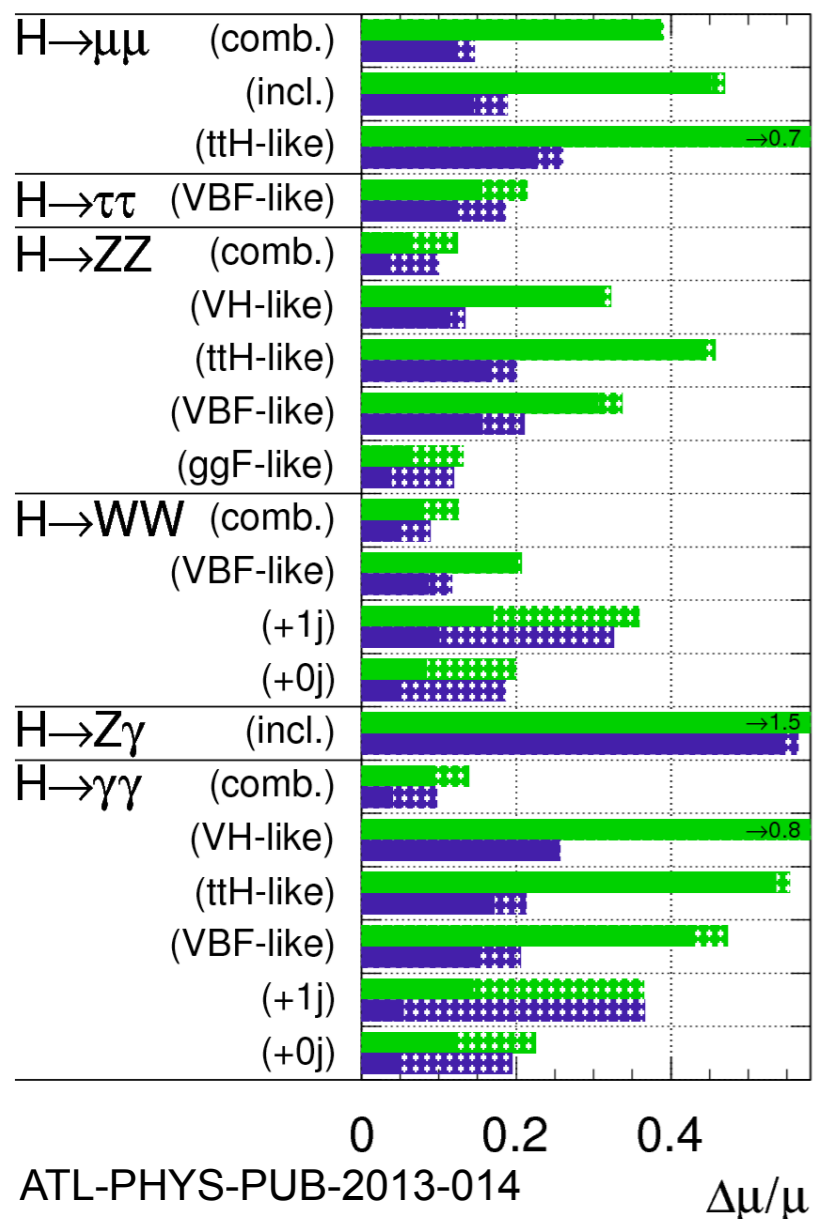
The LHC is a Higgs Factory

- Run II x5-6 more integrated luminosity compared to Run I
- x2.3 - 3.9 increase in Higgs production cross section from 8TeV to 13TeV
- x3.4-5 improvement in statistical sensitivity

In Run II several Higgs analysis may become systematics limited
need to work on reducing those

ATLAS Simulation Preliminary

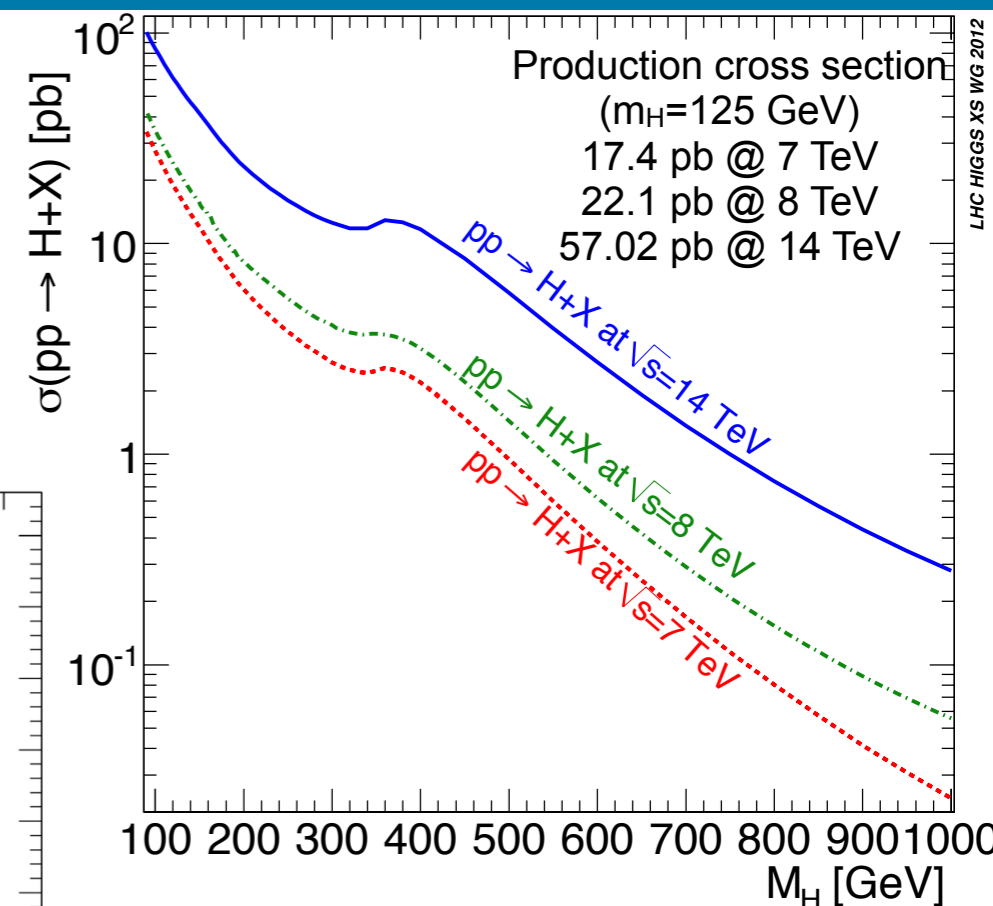
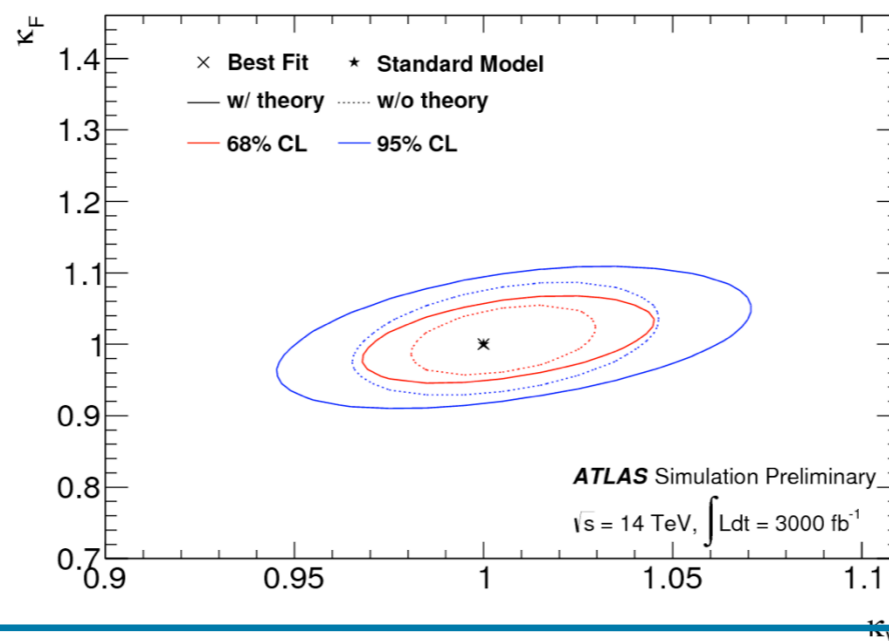
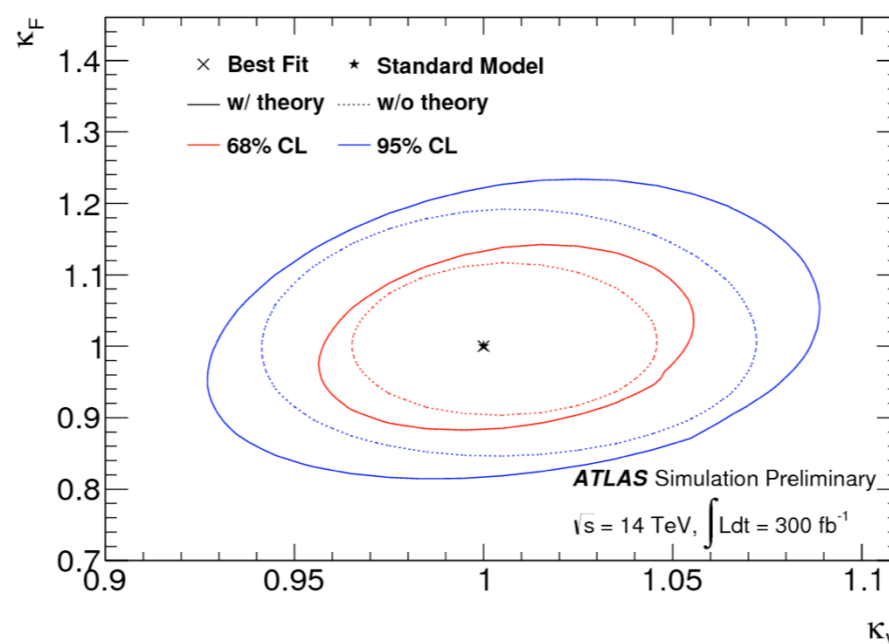
$\sqrt{s} = 14$ TeV: $\int Ldt = 300 \text{ fb}^{-1}$; $\int Ldt = 3000 \text{ fb}^{-1}$



ATL-PHYS-PUB-2013-014

$\Delta\mu/\mu$

ATL-PHYS-PUB-2013-014



Several open topics in the Higgs sector for future studies:

- Rare decays & Couplings
- CP studies
- BSM Higgs boson searches
- Higgs boson pair production

H → ZZ(*) → 4l: Fiducial/Differential cross sections

fiducial cross-section

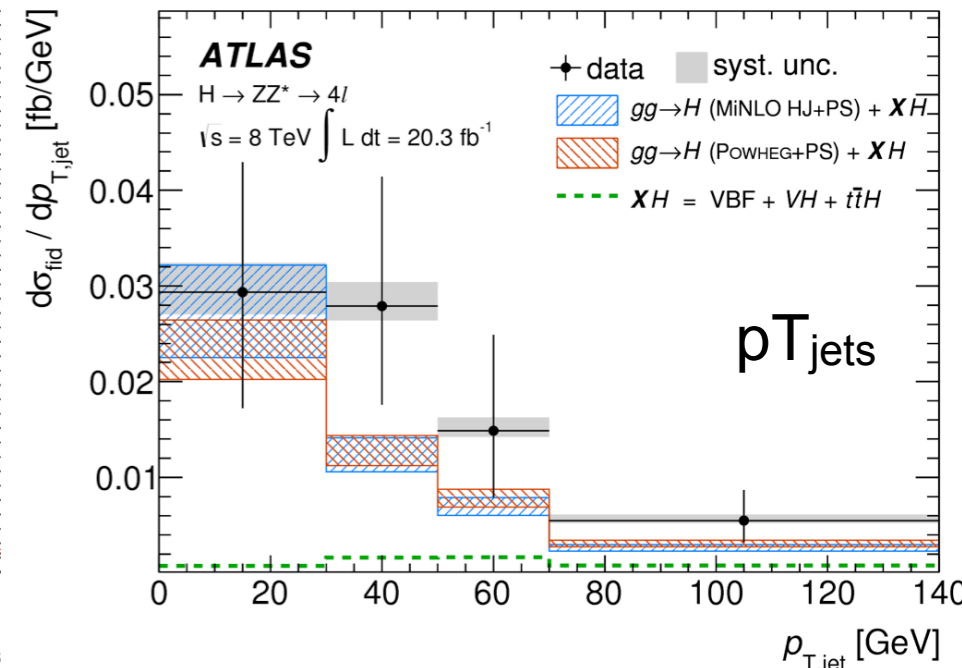
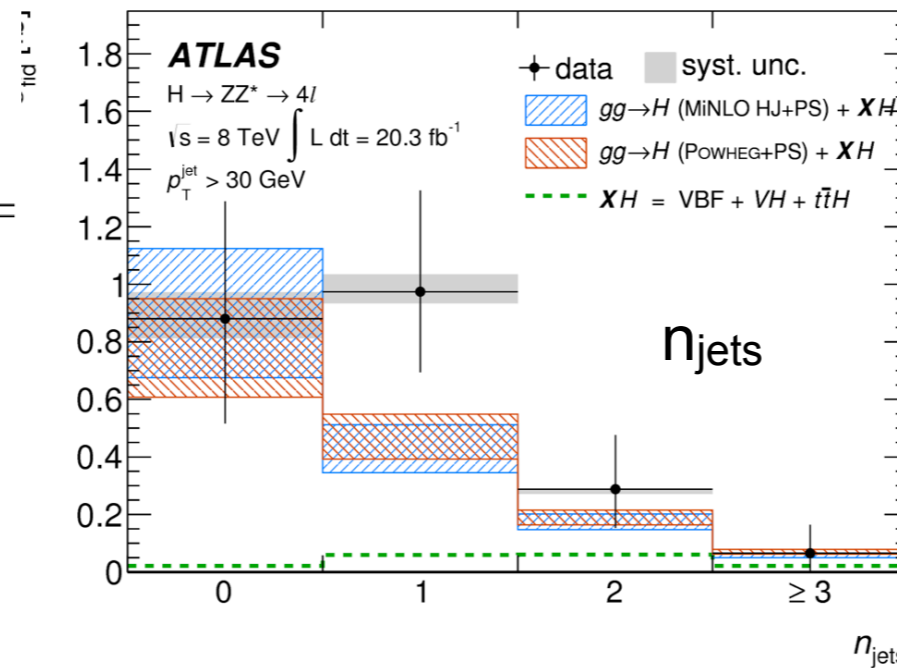
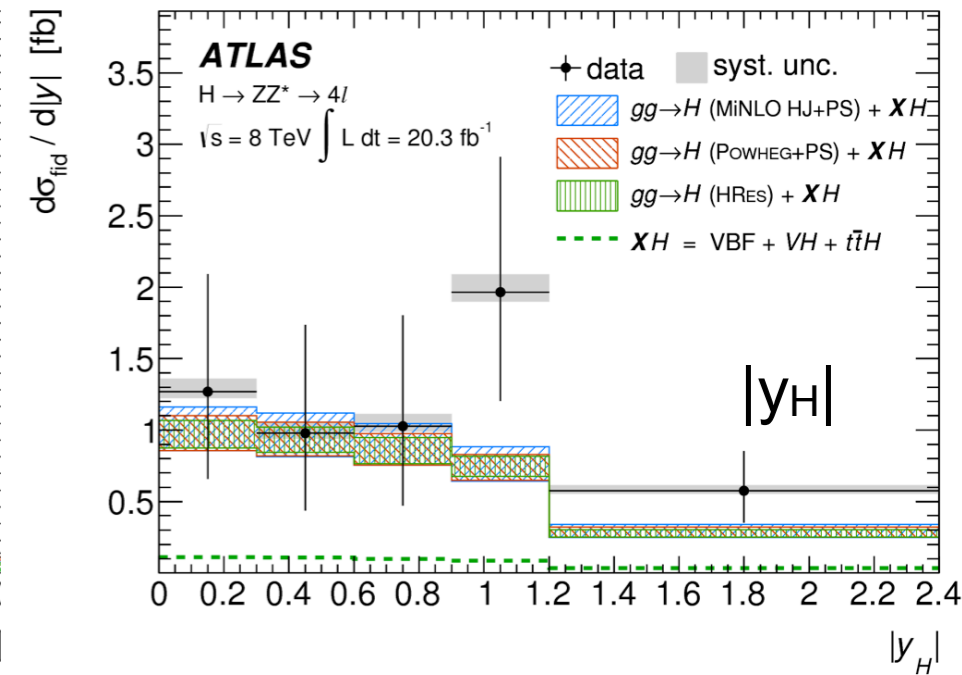
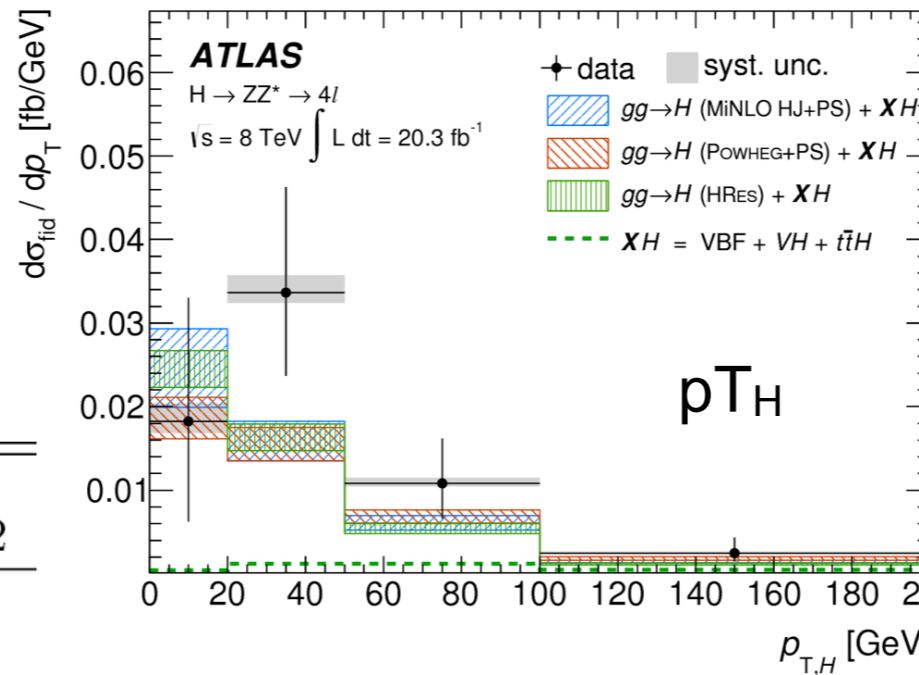
6 differential cross sections

Procedure (120 < m_{4l} < 130 GeV):
 - expected background subtracted from observed events in bins of interesting variable
 - bin-by-bin unfolding

$$\sigma_{\text{tot}}^{\text{fid}} = 2.11^{+0.53}_{-0.47}(\text{stat})^{+0.08}_{-0.08}(\text{syst}) \text{ fb}$$

SM prediction at 125.4 GeV = 1.30 ± 0.13 fb

Variable	<i>p</i> -values		
	POWHEG	MINLO	HRES2
<i>p</i> _{T,H}	0.30	0.23	0.16
<i>y</i> _H	0.37	0.45	0.36
<i>m</i> ₃₄	0.48	0.60	-
cos θ*	0.35	0.45	-
<i>n</i> _{jets}	0.37	0.28	-
<i>p</i> _{T,jet}	0.33	0.26	-



First result/statistics limited → No large surprise

BSM examples

ATLAS-CONF-2014-010

The properties of the observed Higgs boson already constraint BSM contributions

Additional EW singlet field

Simplest extension of SM Higgs sector

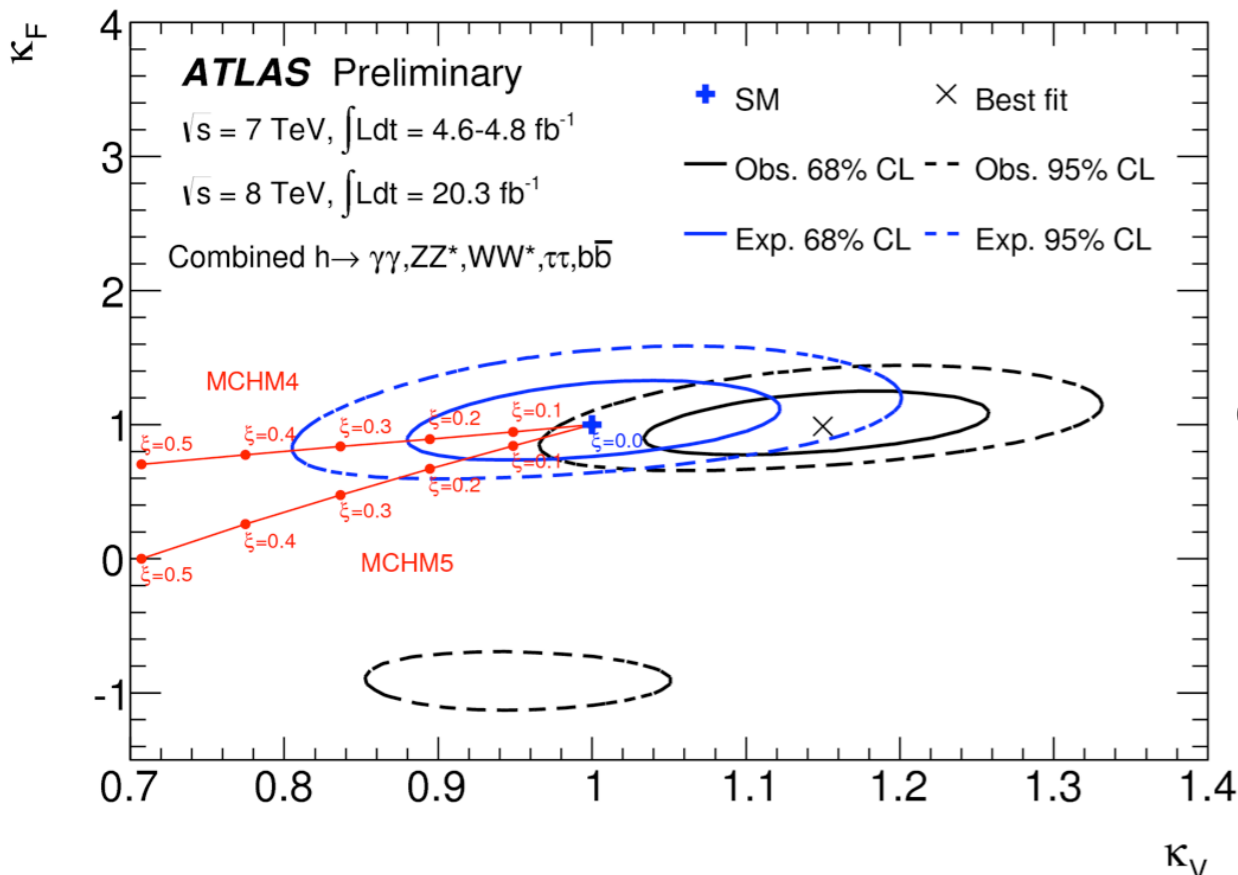
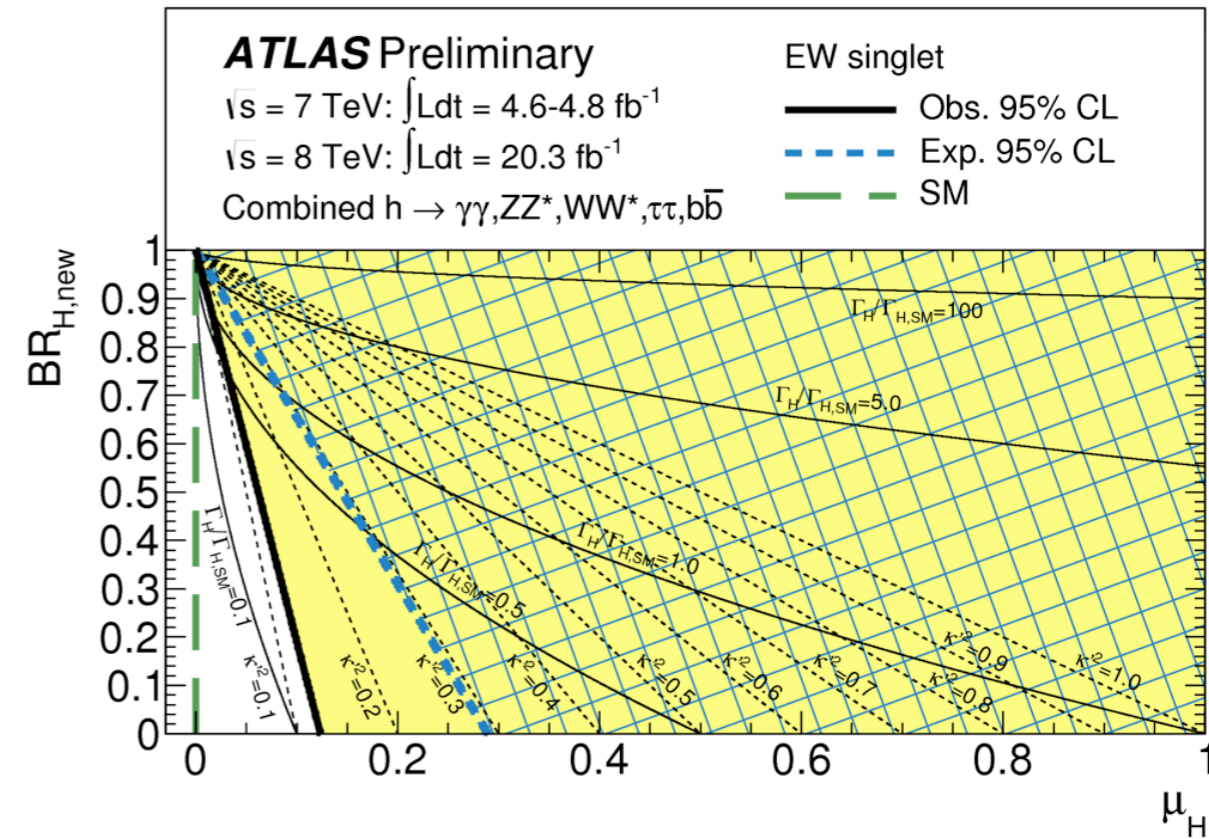
Results in two CP-even Higgs bosons: h, H
(assumed non degenerate)

Couplings similar to SM Higgs boson, but each scaled by common factor, denoted as κ (κ') for $h(H)$.

From unitarity: $\kappa^2 + \kappa'^2 = 1$

$$\mu_h = \frac{\sigma_h \times BR_h}{(\sigma_h \times BR_h)_{SM}} = \kappa^2$$

$$\mu_H = \frac{\sigma_H \times BR_H}{(\sigma_H \times BR_H)_{SM}} = \kappa'^2 (1 - BR_{H,new})$$



Minimal Composite Higgs Model

Higgs is pseudo Nambu-Goldstone boson

Neglecting contributions from new heavy resonances the

Higgs couplings modified wrt SM as a function of

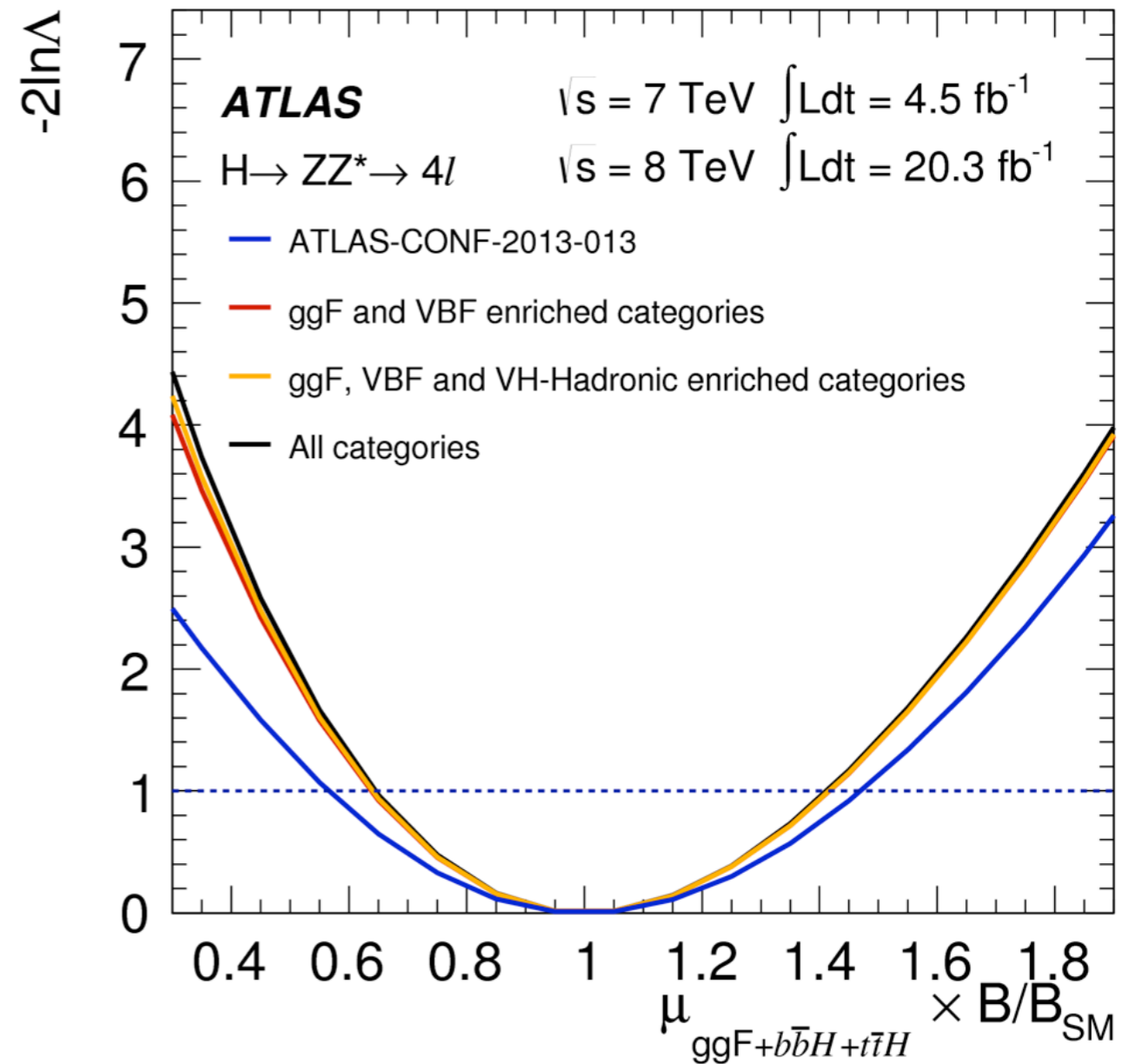
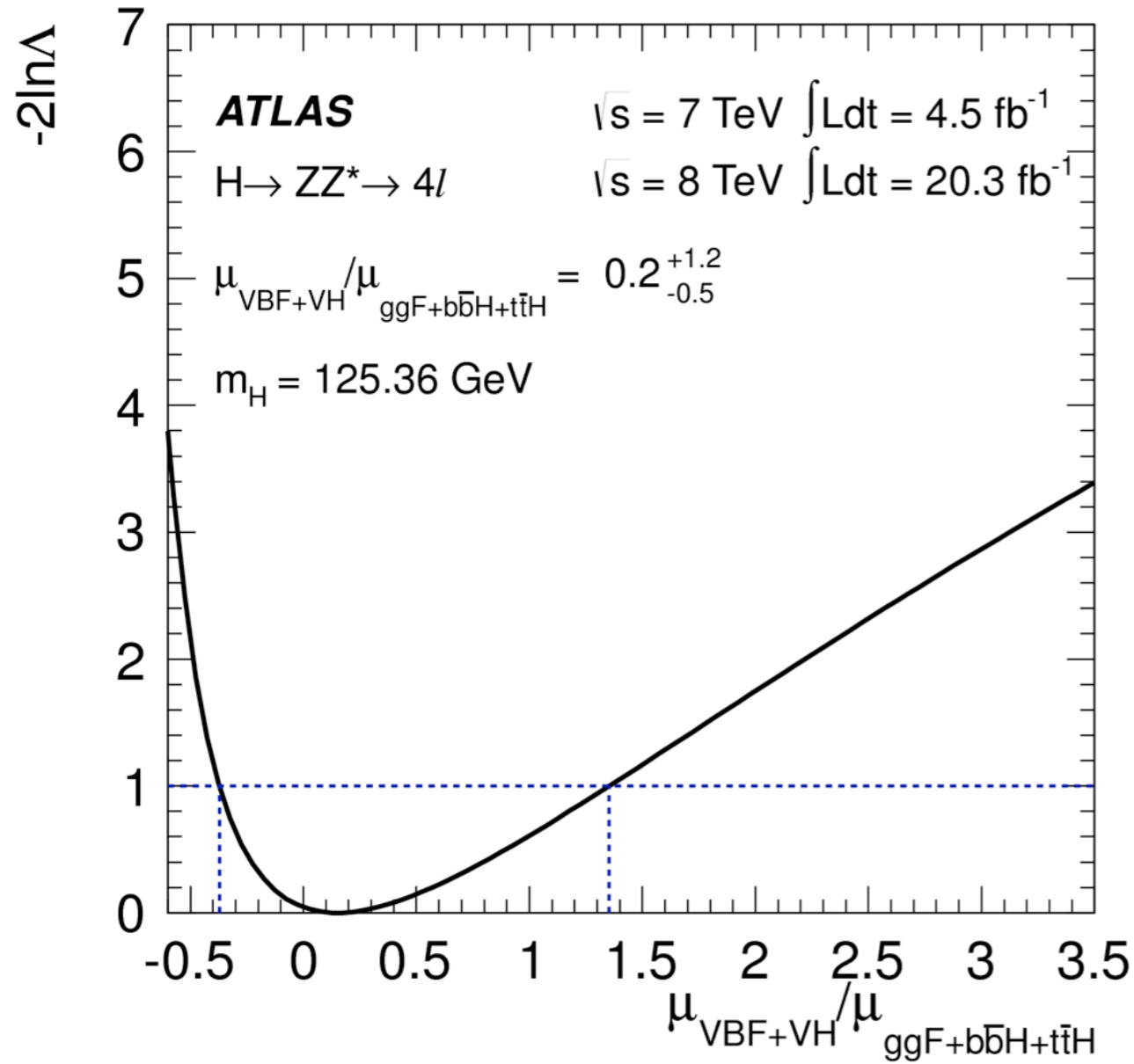
compositeness scale: $\xi = v^2 / f^2$

MCHM4: $\kappa = \kappa_V = \kappa_F = \sqrt{1 - \xi}$ $f > 710$ (460) GeV at 95%CL

$$\kappa_V = \sqrt{1 - \xi}$$

MCHM5: $\kappa_F = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$ $f > 640$ (550) GeV at 95%CL

H → ZZ(*) → 4l: Coupling Results



Prospects for Run II/III and HL-LHC

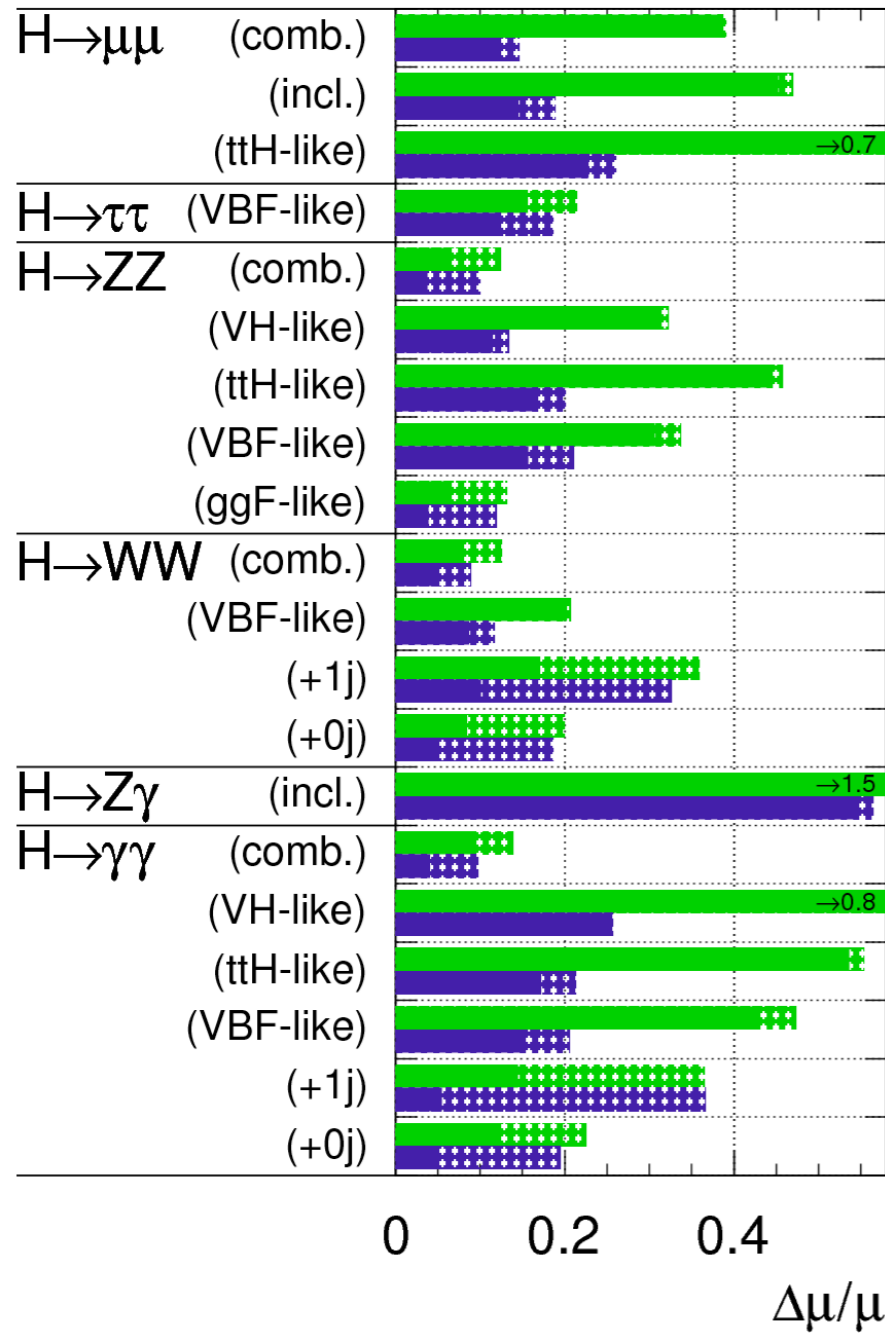
ATL-PHYS-PUB-2013-014

ATLAS Simulation Preliminary

$\sqrt{s} = 14 \text{ TeV}$: $\int \mathcal{L} dt = 300 \text{ fb}^{-1}$; $\int \mathcal{L} dt = 3000 \text{ fb}^{-1}$

Several open topics in the Higgs sector for future studies:

- Rare decays & Couplings
- CP studies
- BSM Higgs boson searches
- Higgs boson pair production



g_1 CP-even HZZ coupling
 g_2 CP-even HZZ coupling (loops)
 g_4 CP-odd HZZ coupling

