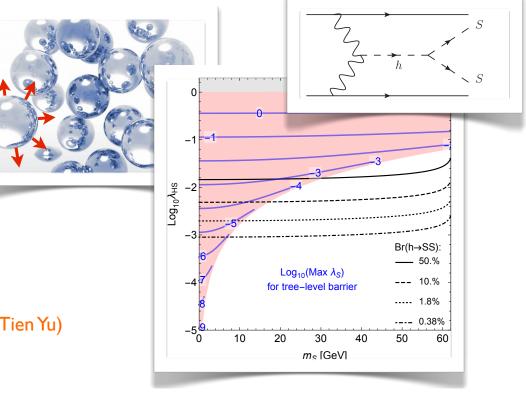
## Probing Electroweak Baryogenesis with Exotic Higgs Decays

HXSWG Exotic Higgs Decays Meeting Fermilab

22 May 2015

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Partially based on 1409.0005 (DC, Patrick Meade, Tien-Tien Yu)



## Introduction

— Where do we (baryonic matter) come from? —

Electroweak baryogenesis (EWBG) is one of the only mechanisms of baryogenesis that is, in principal, **testable.** 

As a BG mechanism, EWBG relies on the *electroweak phase transition*, enhanced by BSM effects. However, even without BSM, probing the **electroweak phase transition** (PT) is cosmologically interesting.

Testing EWBG  $\Rightarrow$  detailed study of the Higgs  $\Rightarrow$  **future colliders** 

**Exotic Higgs decays** are an important observable that constrain or exclude many EWBG theories that involve light particles.

## This Talk

## **3** Reviewing EWBG Mechanism

## **3** How to probe EWBG (general)

## 3 Exotic Higgs Decays

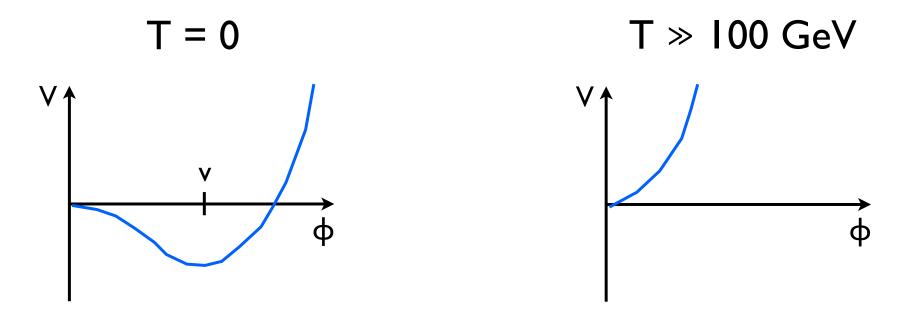
# Higgs at High Temperatures

At finite temperature, the higgs potential receives new contribution from its interaction with the plasma.

Many reviews, e.g. Quiros hep-ph/9901312

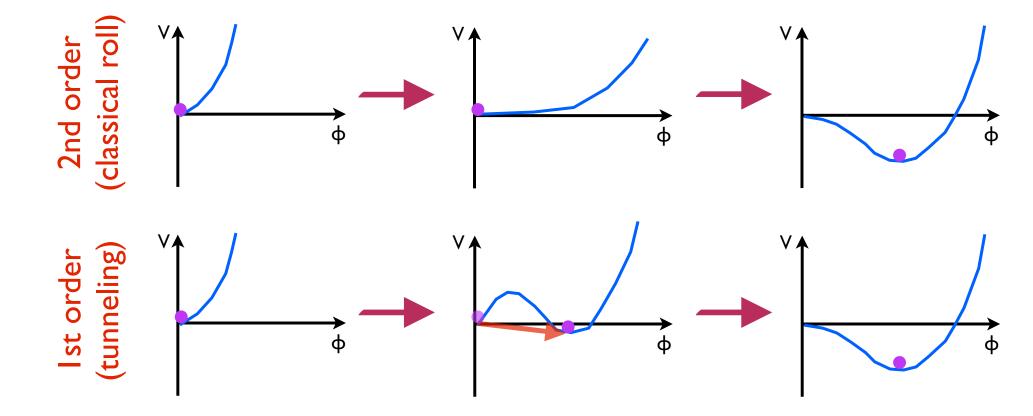
At high temperature, the higgs is stabilized at the origin.

 $\rightarrow$  The early universe was SU(2) symmetric!



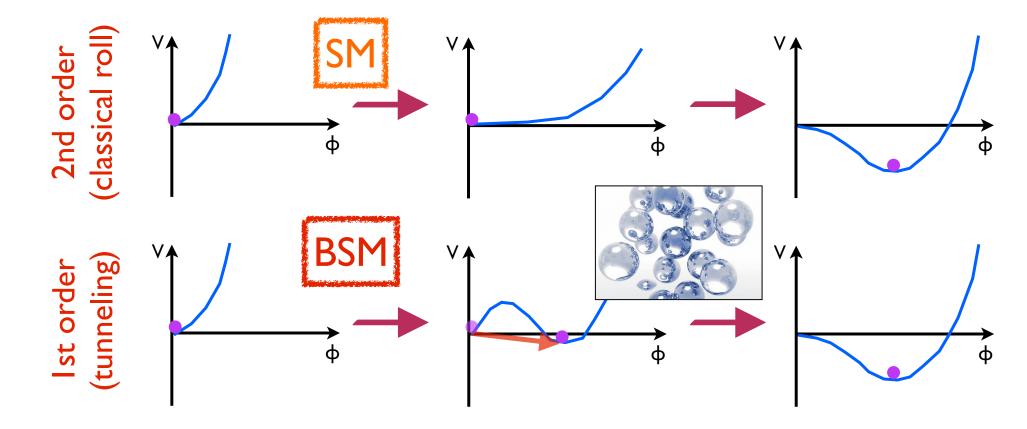
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As the universe cools, the higgs undergoes a **phase transition (PT)** from zero to nonzero VEV.



# Higgs at High Temperatures

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Ist order phase transition gives rise to bubble nucleation!

## This could create baryons...

 $\frac{n_B}{s} \sim 10^{-10} \qquad \text{Why?}$ 

To dynamically create Baryon Number Asymmetry, the three Sakharov conditions must be satisfied.

- I. **B** Number Violation
- 2. **CP**Violation
- 3. **Departure from thermal equillibrium** ("**T** violation")

Most Baryogenesis mechanisms (Affleck–Dine, Leptogenesis,...) rely on very high-scale physics.

Electroweak Baryogenesis is all weak scale  $\Rightarrow$  testable mechanism!

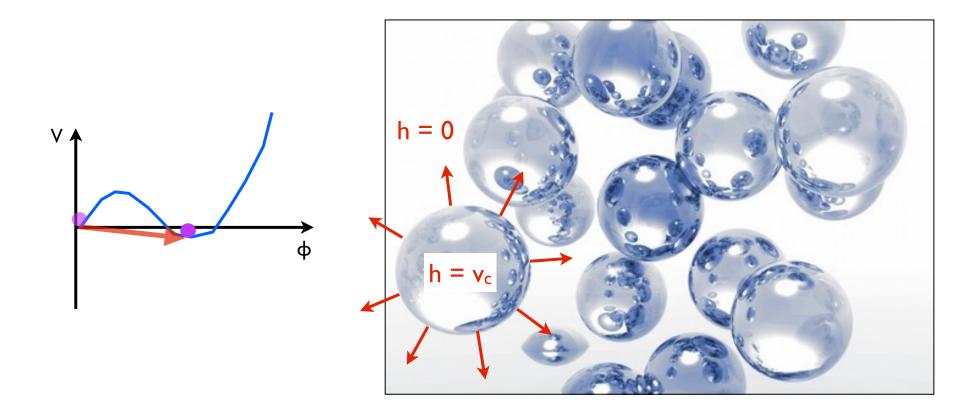
Kuzmin, Rubakov, Shaposhnikov 1985 Klinkhamer, Manton 1984

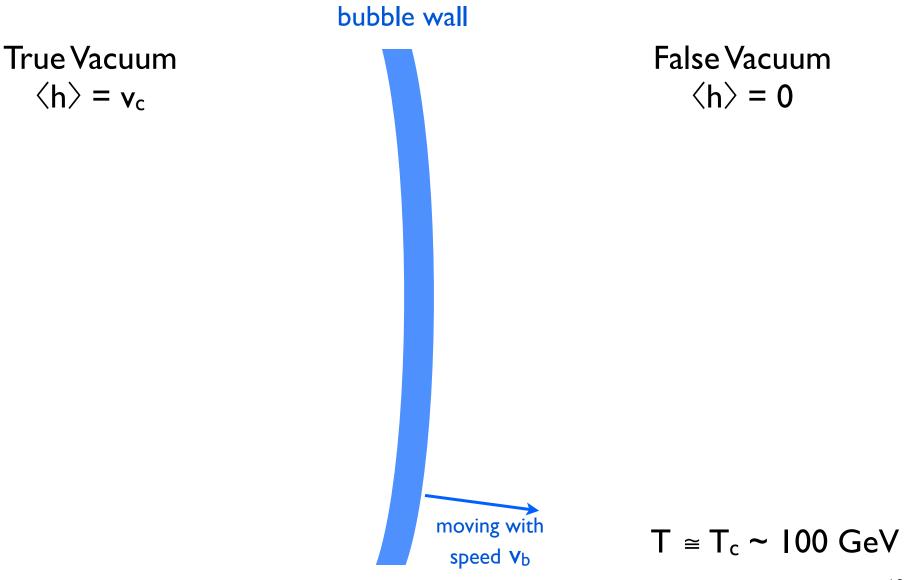
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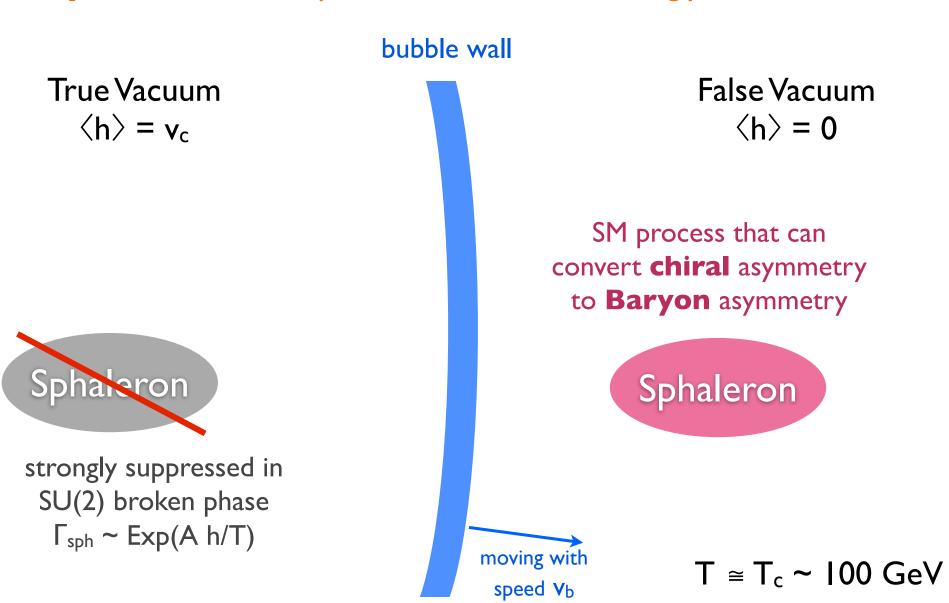
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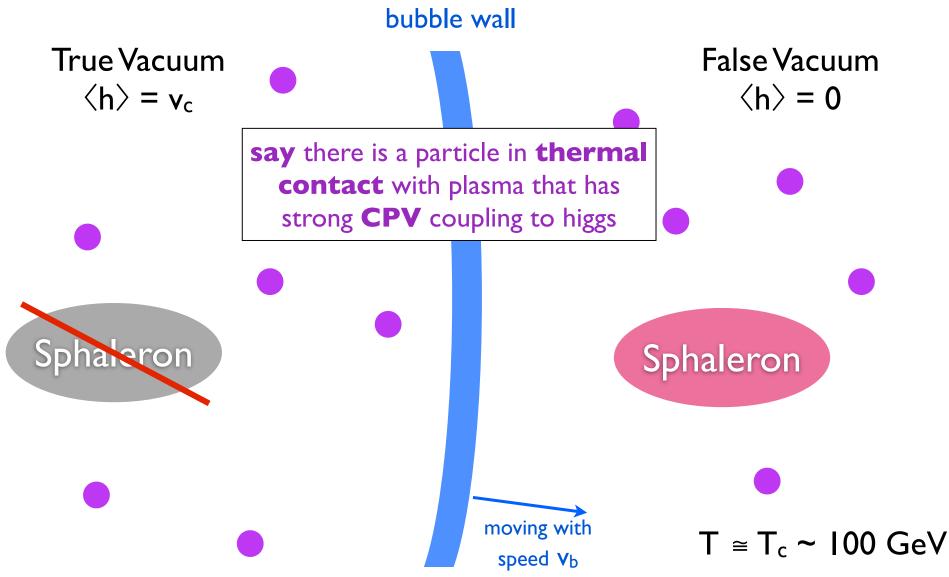
say the electroweak phase transition was strongly 1st order....

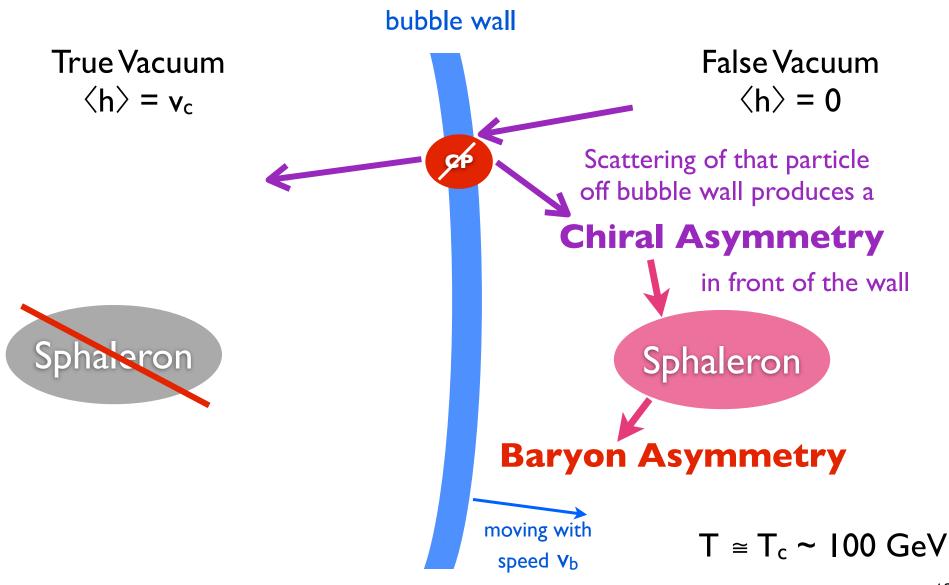
At some critical temperature, bubbles of true vacuum  $\mathbf{h} = \mathbf{v}_c$  form, and grow into the false vacuum surroundings where  $\mathbf{h} = \mathbf{0}$ .











say the electroweak phase transition was strongly 1st order....

True Vacuum  $\langle h \rangle = v_c$ 

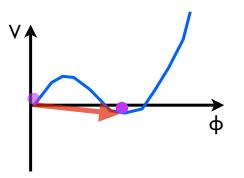




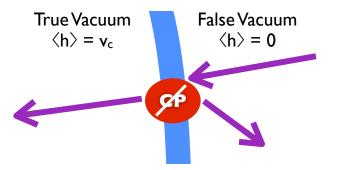
Baryon Asymmetry is now frozen in

#### EWBG requires two BSM ingredients:

I. Modified higgs potential to make phase transition 1st order



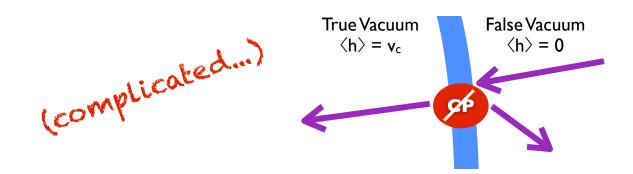
2. Sizable CPV coupling between higgs and another particle (BSM or SM) that is thermally active in the plasma ( $M \leq T$ )



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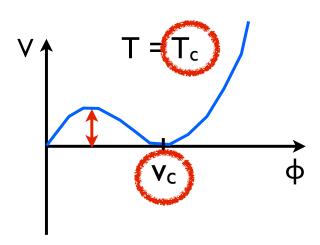
2. Sizable CPV coupling between higgs and another particle (BSM or SM) that is thermally active in the plasma ( $M \leq T$ )



exclude the

# How to exclude a strong electroweak phase transition?

## Strong Phase Transition



The phase transition has to be strong enough to suppress sphaleron washout of the generated baryon number in the broken phase.

$$\frac{v_c}{T_c} > 0.6 - 1.6$$

Normally given as ~I, this more accurate figure is from Patel, Ramsey-Musolf, 1101.4665

Very simple criterion to determine if EWBG is at least *possible* with a given higgs potential.

How can you modify the SM higgs potential to get  $v_c/T_c \ge 1$ ?

$$V_{\text{eff}}(h,T) = V_0(h) + V_0^{CW}(h) + V_T(h,T)$$

tree-level loop finite temperature potential correction corrections

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potential correction

tree-level loop finite temperature corrections

I. Thermal Effects

add new BOSONS to the plasma to generate barrier (analogous to W and Z contributions)

2. Loop Effects

add particles whose loops reduce the 'depth of the higgs potential well', so W and Z contributions can make a barrier.

3. Tree Effects

add scalars to modify tree-level higgs potential and create a barrier

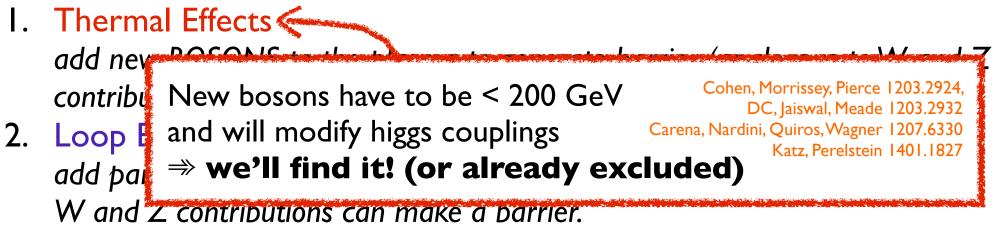
4. add non-renormalizable operators really a general way of parameterizing (2) and (3)  $\leftarrow$  a little subtle....

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- I. Thermal Effects add new BOSONS to the plasma to generate barrier (analogous to W and Z contributions)
- 2. Loop Effects < add particles whose W and Z contributi Do not require very light new particles. 3. Tree Effects Explore with simple scalar extensions of the SM. add scalars to modi
- 4. add non-renormalizable operators really a general way of parameterizing (2) and (3)  $\leftarrow$  a little subtle....

## Tree and Loop-driven PT

#### Consider SM + single real scalar

 $V_0^{T=0}(H,S) = -\mu^2 \left( H^{\dagger} H \right) + \lambda \left( H^{\dagger} H \right)^2 + \frac{a_1}{2} \left( H^{\dagger} H \right) S \\ + \frac{a_2}{2} \left( H^{\dagger} H \right) S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4.$ 

In generality, this scalar mixes with the higgs after EWSB.

→ - direct production in (heavy) higgs searches  $\rightarrow$  - exotic higgs decays h $\rightarrow$ ss (if light enough) pp lepton - higgs precision coupling measurement constraints collider 🔨 - modifications to higgs self-couplings 🔶 - modification to Zh coupling But the model still has A lot of handles for many parameters. Can discovery using all EWBG be completely future colliders! ace needed! excluded? e.g. Profumo, Ramsey-Musolf, Wainwright, Winslow 1407.5342

# Tree and Loop-driven PT

Need a simpler model to investigate these strong phase transitions....

DC, Patrick Meade, Tien-Tien Yu 1409.0005

Build a 'maximally stealthy' model to implement these mechanisms, then see how to exclude that model.

> **SM + unmixed Singlet:** A `simplified model' of stealthy electroweak baryogenesis!

$$V_0 = -\mu^2 |H|^2 + \lambda |H|^4 + \frac{1}{2}\mu_S^2 S^2 + \lambda_{HS} |H|^2 S^2 + \frac{1}{4}\lambda_S S^4$$

heavy Higgs (sibling) production Higgs portal production Exotic Higgs Decays h<sup>3</sup> measurement EWPO Higgs Couplings  $\sigma(Zh)$  measurement

## Tree and Loop-driven PT

Need a simpler model to investigate these strong phase transitions....

Build a 'maximally ste see how to exclude t

Let's first understand this model in the regime without Exotic Higgs Decays (no light singlet).

EWPO

 $\sigma(Zh)$  measurement

rick Mooda Tion Tion Yu 1409-0005

Come back to study in light singlet regime.

electroweak baryogenesis!

minghined model of second

$$V_0 = -\mu^2 |H|^2 + \lambda |H|^4 + \frac{1}{2}\mu_S^2 S^2 + \lambda_{HS} |H|^2 S^2 + \frac{1}{4}\lambda_S S^4$$

Exotic Higgs Lowys

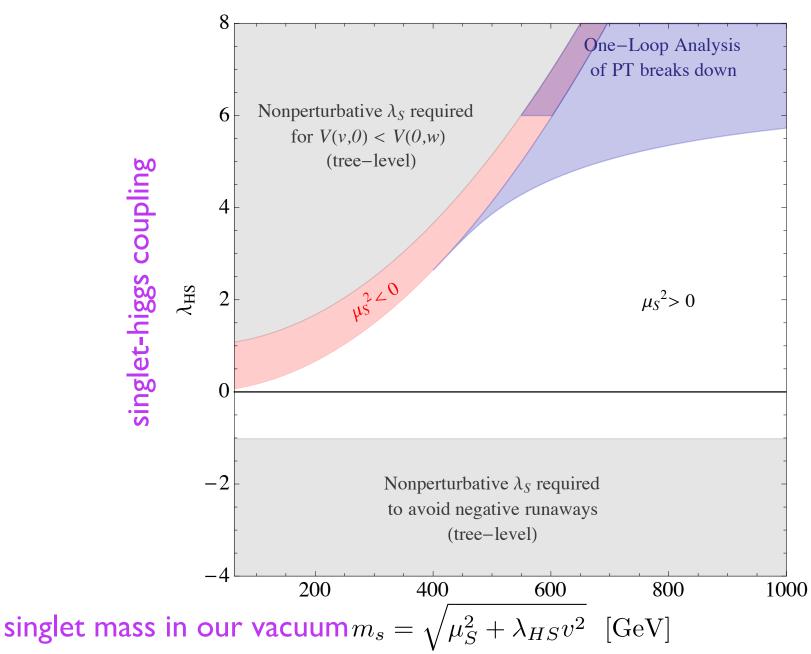
h<sup>3</sup> measurement

heavy Higgs (sibling) production Higgs portal production

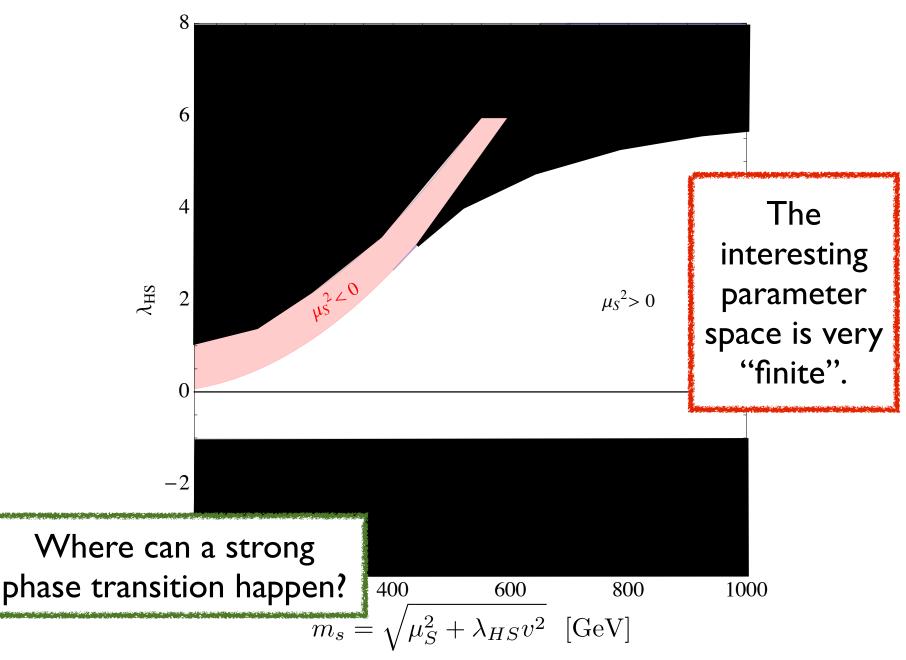
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**Higgs** Couplings

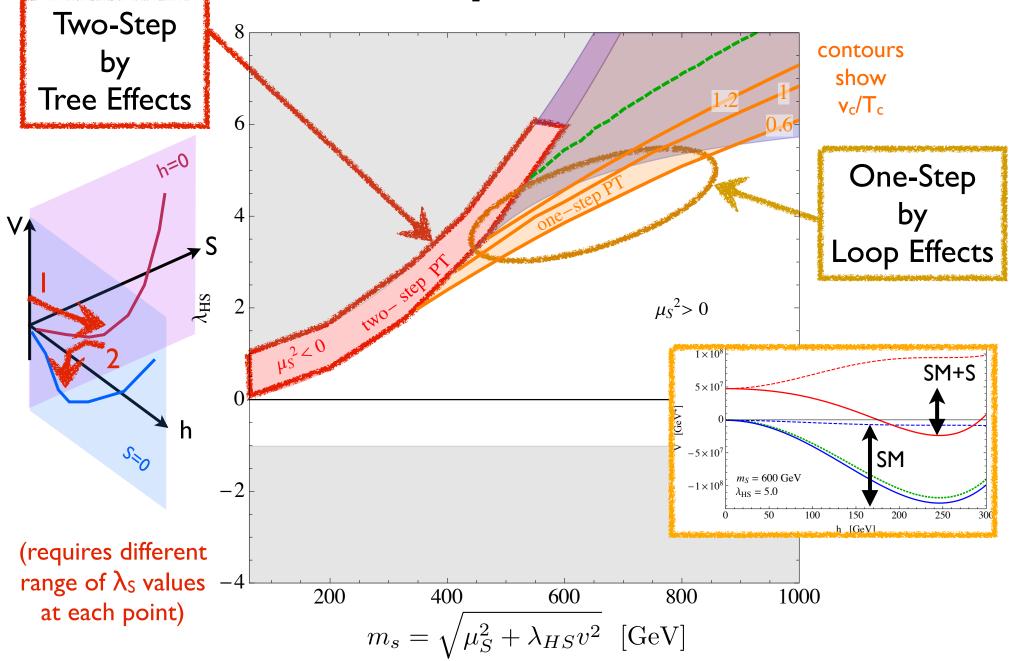
## The (m<sub>s</sub>, $\lambda_{Hs}$ ) Plane



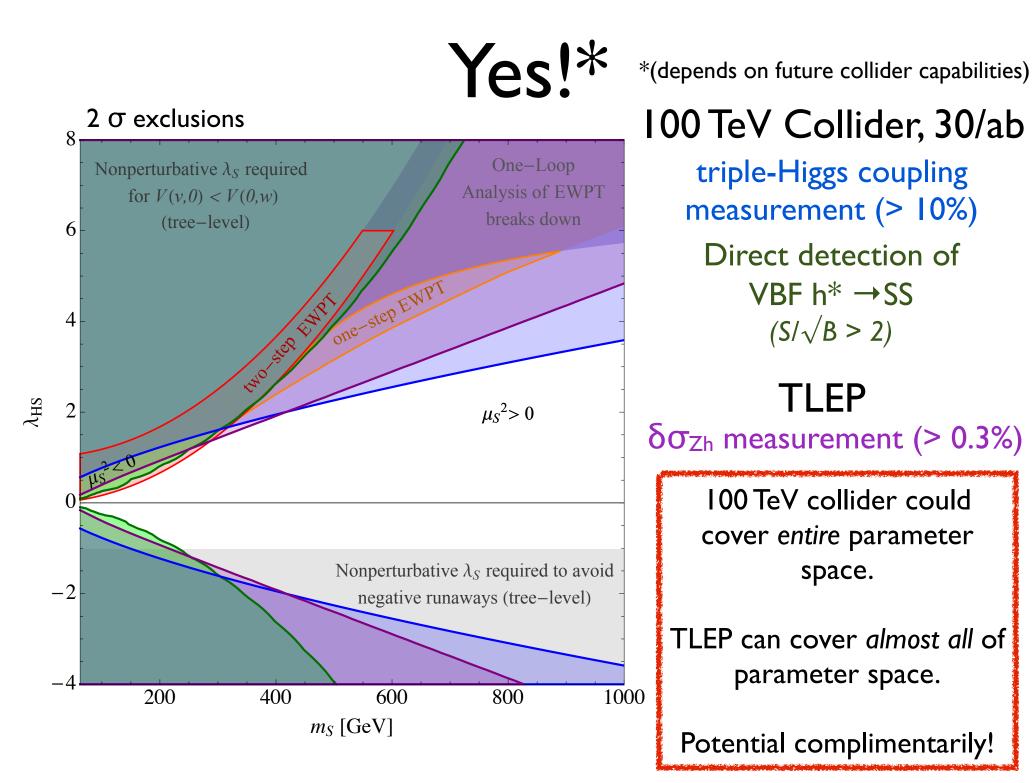
## The (m<sub>s</sub>, $\lambda_{Hs}$ ) Plane



## Two\* kinds of phase transitions



# So can we exclude EWBG in this model?



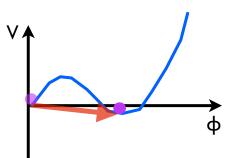
Light particles coupled to the Higgs could arise (separately) due to the physics of the strong phase transition (PT) or the creation of baryon asymmetry (BAU) during the PT.

BAU creation requires O(100 GeV) particles  $\frac{\text{True Vacuum}}{\langle h \rangle = v_c}$  with sizable CPV couplings to higgs.

This particle could be < 63 GeV, but then

- in many cases, SM gauge charged  $\rightarrow$  excluded by LEP limits
- if SM neutral, still require O(0.1) Higgs couplings
  - $\rightarrow$  exotic Higgs decays already exclude this, or will soon...

What about the physics of the strong Phase Transition?



**False Vacuum** 

 $\langle h \rangle = 0$ 

— Thermally driven PT —

Need O(I) couplings of new bosons to Higgs

If m < 63 GeV, will overwhelm SM Higgs decay modes.

**Already inconsistent with data!** 

— Tree or Loop driven PT —

Again, explore in context of simple scalar extensions.

#### SM+S, 2HDM+S, with S-H mixing:

Big parameter space again makes it difficult to come up with exact statements.

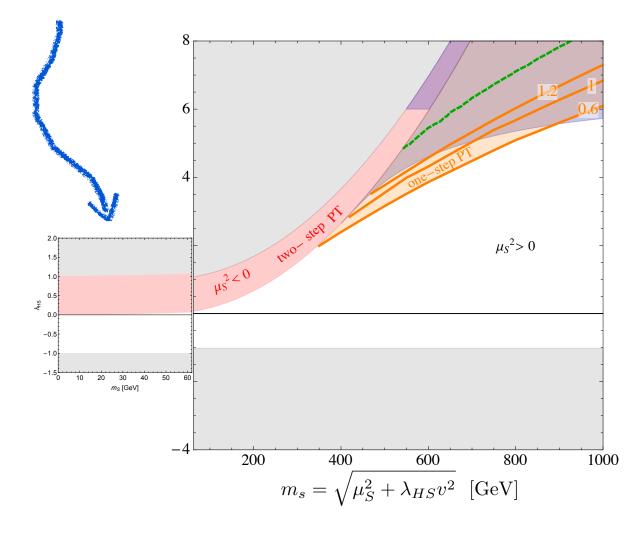
If  $m_s < 63$  GeV, mixing effects to give rise to  $h \rightarrow 2s \rightarrow 4b$ ,  $4\tau$ , etc

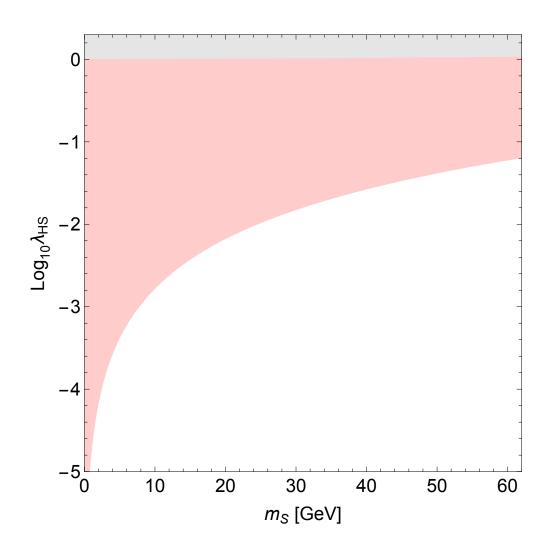
Should have access to Br ~  $10^{-4}$  or  $^{5} \leftrightarrow s_{\theta} \sim 10^{-2}$  at lepton colliders

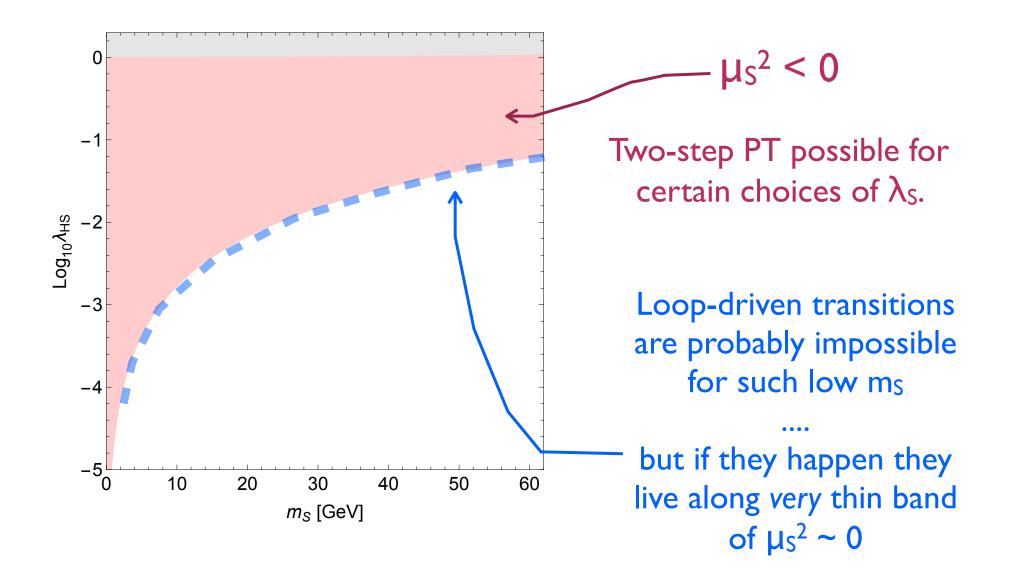
Certain cases are already excluded by data if scalar is light, e.g. exotic two-step transition.

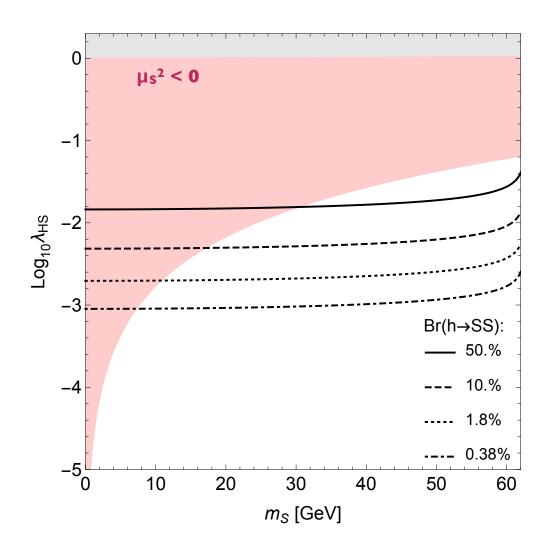
Let's go back to our simplified model of the strong PT to make some quantitative statements about exotic Higgs decays.

Explore the unmixed SM+S model in the regime where  $m_S < m_h/2$ 

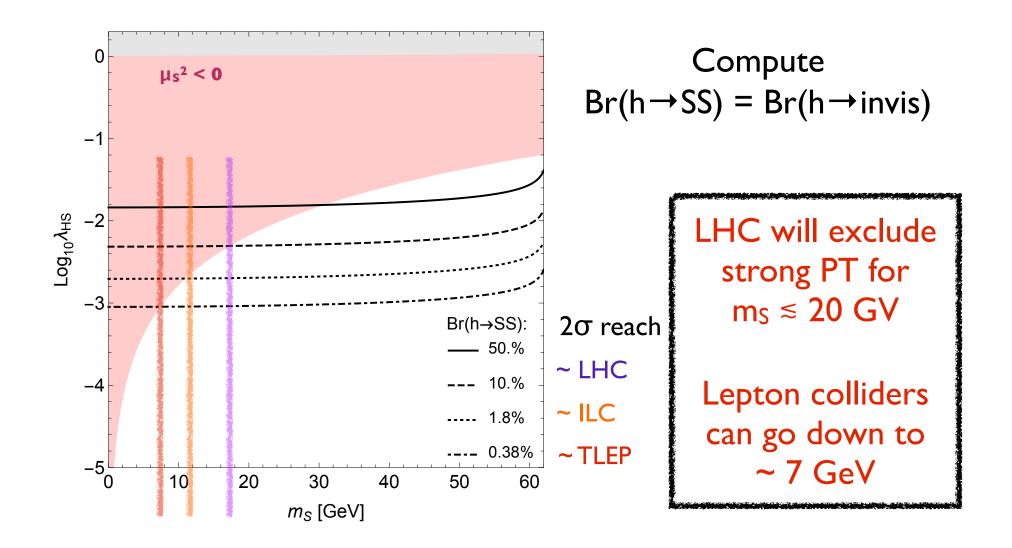


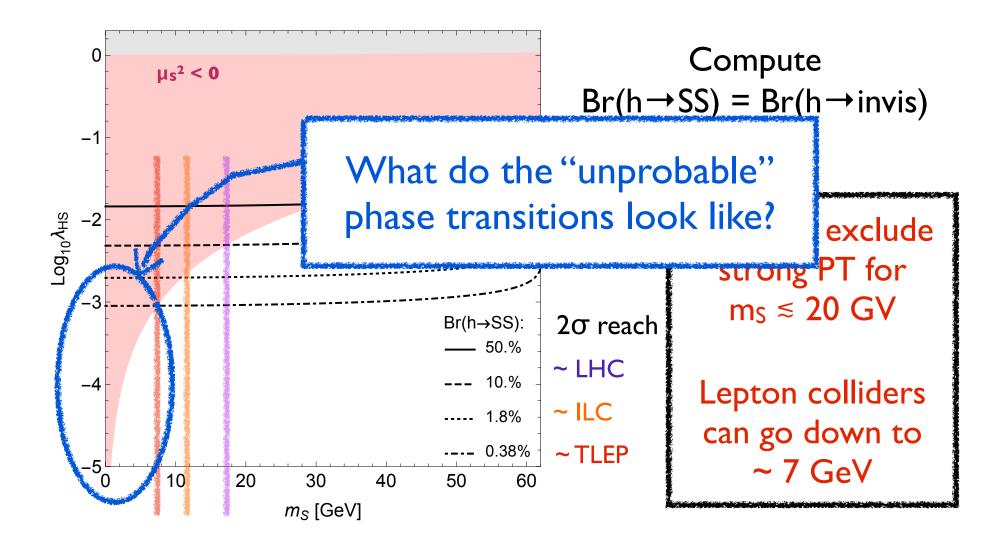


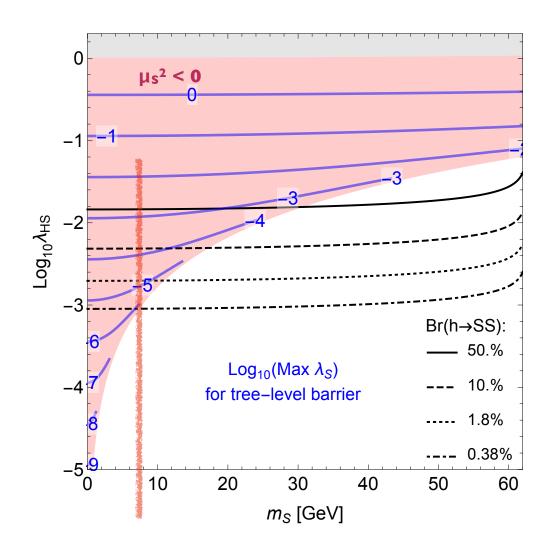




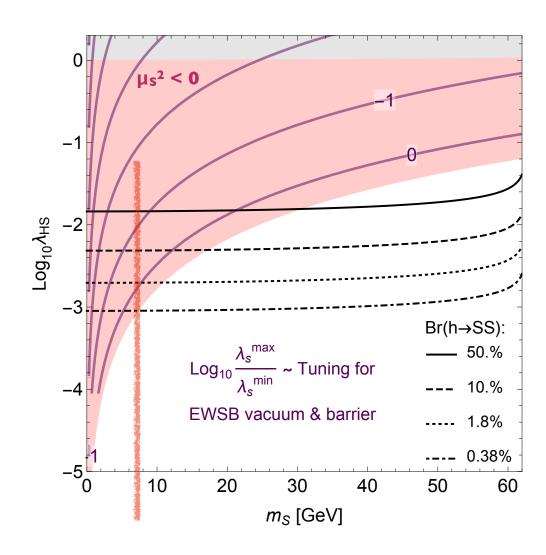
Compute Br( $h \rightarrow SS$ ) = Br( $h \rightarrow invis$ )



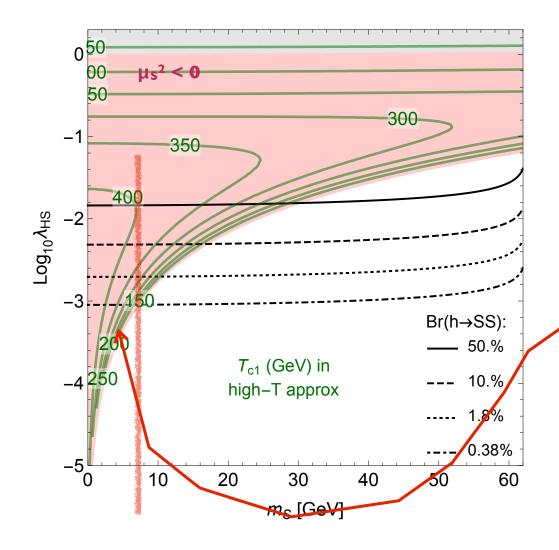




## "Unprobable" two-step PT requires $\lambda_s < 10^{-5}$



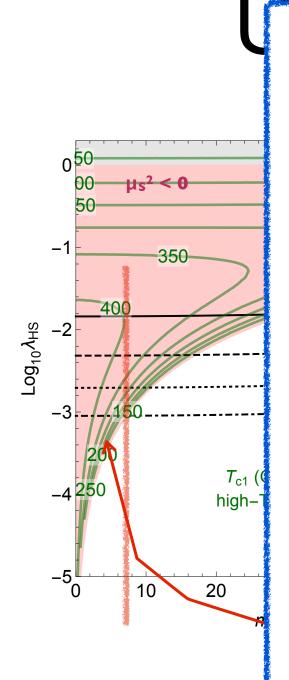
Quartic Tuning & Radiative Stability: a priori OK



Reasonable values for T<sub>c1</sub>, where singlet transitions from 0 to [large] vev.

**Explicit spot-check:** twostep PT works as expected in full finite-T quantum calculation

But are there other reasons this might fail?



Moral of the story:

You'll never exclude

\*!!\*everything\*!!\*

But for light scalars giving strong PT, evading exotic Higgs decay bounds certainly requires a very strange world.... lues for T<sub>c1</sub>, transitions arge] vev.

-check: twois expected in um calculation

ther reasons nt fail?

## Conclusions

## Conclusions

EWBG can be probed at current colliders, and maybe excluded at future colliders (without exotic Higgs decays) by looking for signs of strong phase transition.

Exotic Higgs Decays are a powerful probe of EWBG if the physics of the strong PT or BAU creation involves particles below 63 GeV.

Basically excludes such light particles facilitating BAU creation

Basically excludes such light particles causing thermally driven PT

For tree/loop driven PT in scalar extensions **with** H-mixing, SM+S or 2HDM+S inspired decays  $h \rightarrow 2a \rightarrow 4f$  (Yukawa ordered) are powerful probes, especially at lepton colliders. Make this more quantitative?

For tree/loop driven PT in scalar extensions **without** H-mixing, lepton collider bounds on  $h \rightarrow$  invis exclude all but very light particles with peculiarly tiny couplings. Other ways to probe?