

Spectroscopy of b^- and c^- -hadrons

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on behalf of the LHCb Collaboration

including contents from ATLAS, BABAR, BELLE, BESIII, CDF, and CMS

The Third Annual Conference on Large Hadron Collider Physics

LHCP 2015



Friday September, 04th – St. Petersburg, Russia

Introduction

New c -hadrons

- In prompt $D^{(*)}K$ combinations
- $B \rightarrow Dhh$ studies (spin-3 D mesons)
- $B \rightarrow DDh$ studies

New b -hadrons

- $B_c^+(2S)$ observation
- Excited b -mesons at LHCb and CDF
- Excited b -baryons

Quarkonia(-like) spectroscopy

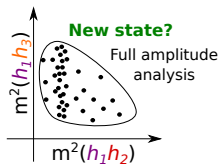
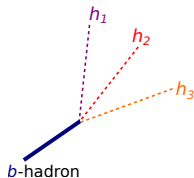
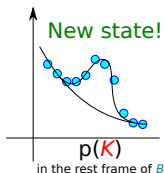
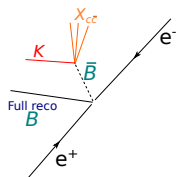
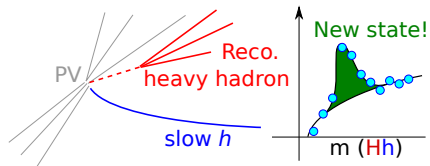
- New Z_c states from BESIII
- $X_b^+ \rightarrow \Upsilon(1S)\pi^+$ searches
- Study of $X(3872) \rightarrow \rho^0 J/\psi$
- Tetraquark $Z_c(4430)^-$
- ~~Pentaquarks observation P_c~~
(covered in Anton's highlights)

Conclusion and outlook

Don't miss the HF-3 session this afternoon, with theoretical and experimental reports on Heavy Flavour spectroscopy.

Important test for QCD theories.
 Several concurrent techniques to explore heavy
 flavours excited states.

ATLAS & CMS	pp with $\sqrt{s} = 7 - 8$ TeV
LHCb	pp with $\sqrt{s} = 7 - 8$ TeV
CDF	$p\bar{p}$ with $\sqrt{s} = 1.96$ TeV
Belle & Babar	e^+e^- with $\sqrt{s} = 10.58$ GeV
BESIII	e^+e^- , lower \sqrt{s}



- 1 Define a *pure sample* of $R \rightarrow abc$ decays ($R \equiv c^-$ or b -hadron).
Background can be
 - Statistically subtracted with **sPlot** (in binned analyses)
 - Statistically subtracted with **sFit** (in unbinned analyses)
 - Modeled and included in the analysis (**cFit**)
- 2 Define the decay amplitude \mathcal{M} as sum of resonant contributions \mathcal{M}_r , depending on the set of variables σ (invariant masses, and decay angles).
- 3 Compute the decay width as a function of σ :

$$d\Gamma(\sigma) = K(\sigma) \left| \sum_r \mathcal{M}_r(\sigma) \right|^2 d\sigma$$

$K(\sigma)$ corrects for kinematics; it depends on the choice of σ .

- 4 Compute efficiency $\varepsilon(\sigma)$, usually using flat-dynamics simulation.
- 5 Fit the pdf

$$pdf(\sigma)d\sigma = \varepsilon(\sigma)\Gamma(\sigma)d\sigma$$

to the *pure sample* defined above.

Changing the recipe of the \mathcal{M}_r contributions, one draws conclusions on

- Existence of resonant states
- Quantum numbers of resonant states

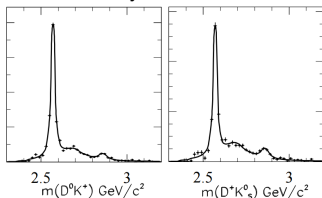
Taking into account the interference between the resonant states!

c-hadron spectroscopy

BABAR (2009)

[PRD80:092003]

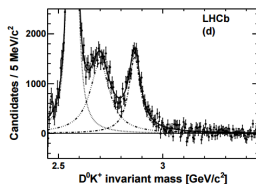
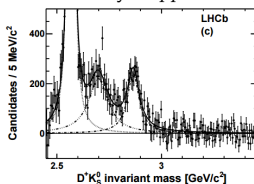
Inclusive study of $e^+e^- \rightarrow D^{(*)}K$



LHCb (2012)

[JHEP 1210 (2012) 151]

Inclusive study of $pp \rightarrow DK$



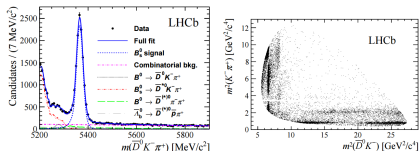
The three peaks are

- $D_{s2}^{*+}(2573) \rightarrow D^0 K^+$ and $D_{s2}^*(2573)^+ \rightarrow D^+ K_S^0$;
- $D_{s1}^{*+}(2710) \rightarrow D^0 K^+$ and $D_{s2}^*(2710)^+ \rightarrow D^+ K_S^0$; observed by Babar (2009)
- $D_{sJ}^{*+}(2860) \rightarrow D^0 K^+$ and $D_{sJ}^*(2860)^+ \rightarrow D^+ K_S^0$; observed by Babar (2009)

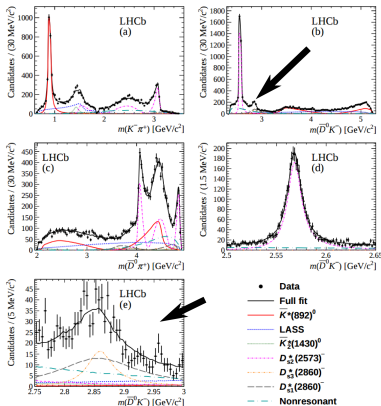
BABAR also observed $D_{s1}^*(2710)^+$ and $D_{s1}^*(2860)^+$ to D^*K final states.

$\bar{B}_s^0 \rightarrow D^0 K^- \pi^+$: an abundant decay with **rich resonant sub-structure**

Full amplitude analysis using the *cFit* technique, fitting the fake B_s^0 component (Combinatorial, Partial reco, and mis-id Λ_b^0) together with the **15 resonant and non-resonant** signal amplitudes.



The structure at $m \sim 2860 \text{ MeV}/c^2$ contains both spin-1 and spin-3 components (significance $> 10\sigma$)! Components were named $D_{s1}^*(2860)^-$ and $D_{s3}^*(2860)^-$



LASS shape: parametrization of the S-wave component.

This is the **first observation** of a spin-3 heavy flavour hadron.

Decay mode	Branching fraction	Reference
$B^0 \rightarrow \bar{D}^0 \pi^+ \pi^-$	$8.46 \pm 0.14 \pm 0.29 \pm 0.40$	[PRD 92, 032002]
$B_s^0 \rightarrow \bar{D}^0 f_0(980)$	$0.017 \pm 0.010 \pm 0.005 \pm 0.001$	[JHEP 08(2015)005]
$B^0 \rightarrow \bar{D}^0 K^+ \pi^-$	$0.92 \pm 0.06 \pm 0.07 \pm 0.06$	[PRD 92, 012012]
$B_s^0 \rightarrow \bar{D}^0 K^- \pi^+$	$10.0 \pm 0.4 \pm 1.0 \pm 1.0$	[PRD 90, 072003]
$B^0 \rightarrow \bar{D}^0 K^- K^+$	$0.47 \pm 0.09 \pm 0.06 \pm 0.05$	[PRL 109, 131801]
$B_s^0 \rightarrow \bar{D}^0 K^- K^+$	0.42 ± 0.19	[PRL 109, 131801]
$B^- \rightarrow D^+ K^- \pi^-$	$0.731 \pm .019 \pm 0.022 \pm 0.039$	[PRD 91, 092002]

Observed D^* mesons

Observed D_s^* mesons

$D_0(2400)$, and $D_s(2460)$
dominant in all analyses

$D^*(2650)$, $D^*(2760)$,
 $D^*(3000)$

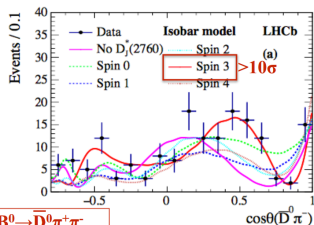
spin-3 $D^*(2800)^+$ in
 $B^0 \rightarrow \bar{D}^0 \pi^+ \pi^-$

spin-1 $D^*(2791)^+$ in
 $B^- \rightarrow \bar{D}^+ K^- \pi^-$

$D_{s1}^*(2700)^-$

$D_{s1}^*(2860)^0$ and
 $D_{s3}^*(2860)^0$ (cf. above)

$D_{sJ}^*(3040)^-$



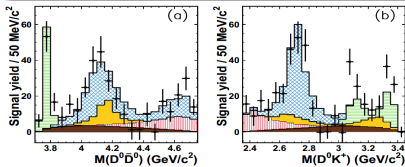
$B^0 \rightarrow \bar{D}^0 \pi^+ \pi^-$

$\theta(\bar{D}^0 \pi)$ defined in the rest frame
of D^*

$$B^0 \rightarrow D^+ D^0 K^- \quad \text{and} \quad B^+ \rightarrow D^0 \bar{D}^0 K^+$$

Belle (2008)

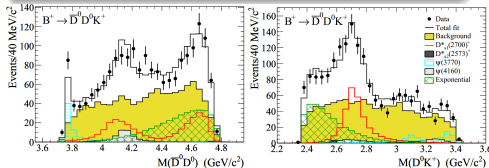
[PRL 100-092001]



- Observed resonance $D_{sJ}^*(2700)^+$
- Found it dominates $B^+ \rightarrow D^0 \bar{D}^0 K^+$ decays
- **Need for unexplained non-resonant contribution**, modeled as an **exponential** in $m(D^0 K^+)$.

Babar (2015)

[PRD 91, 052002]



- Confirms all Belle's conclusions with higher statistics
- Improves precision in mass and width measurements
- Extends the observation of the *non-resonant* contribution to $B^0 \rightarrow D^- D^0 K^+$

Nature of $D_{sJ}^*(2700)^+$ is still unclear. It has been suggested it is:

- the first radial excitation of the $D_s^*(2112)$ meson [PRD 86, 054024]
- an 1^3D_1 $c\bar{s}$ state [PRD 89, 074023], possibly mixing with the former
- two resonances overlapping

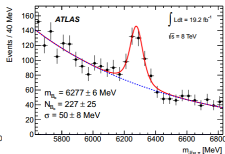
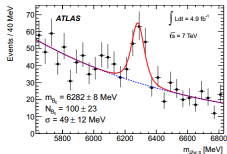
News from LHCb coming soon!

b-hadron spectroscopy

Observed the decay

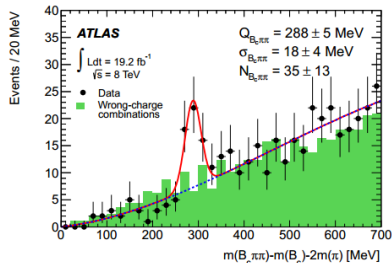
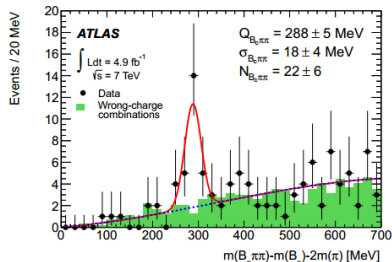
$$B_c^+(2S) \rightarrow B_c^+ \pi^+ \pi^-$$

Full Run-1 data sample



$B_c^+(2S)$ candidates reconstructed using B_c^+ candidates in mass window $m(B_c^+)_{PDG} \pm 3\sigma$

Mass	$6842 \pm 4 \pm 5 \text{ MeV}/c^2$
$N(B_c^+)$ in 2011	100 ± 23
$N(B_c^+)$ in 2012	227 ± 25
$N(B_c^+(2S))$ in 2011	22 ± 6
$N(B_c^+(2S))$ in 2012	35 ± 13
Significance	5.2σ



Reanalysis combining B^+ meson from

- $B^- \rightarrow J/\psi K^-$;
- $B^{-,0} \rightarrow D^{0,-} \pi^+$
- $B^{-,0} \rightarrow D^{0,-} \pi^+ \pi^- \pi^+$
- $B^0 \rightarrow J/\psi K^*(892)^0$;
- $B^0 \rightarrow J/\psi K_S^0$;

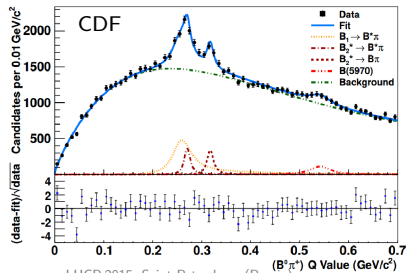
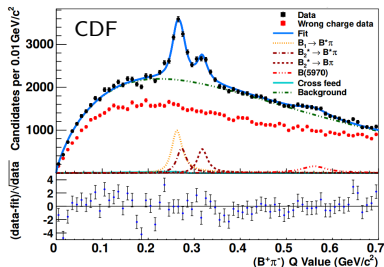
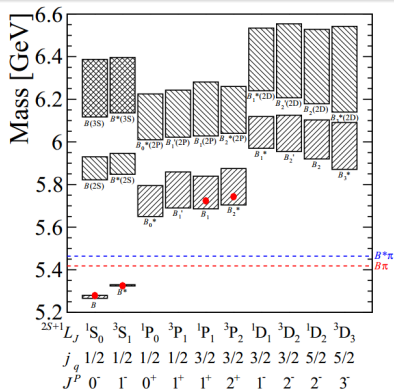
Structures in the plots are:

$$B_1 \rightarrow B^* \pi$$

$$B_2^* \rightarrow B^* \pi \text{ (FIXED shape\&yield)}$$

$$B_2^* \rightarrow B \pi$$

$$B(5970) \rightarrow B \pi \text{ (NEW!)}$$



LHCb analysis of the Run-1 pp dataset.

Using decays:

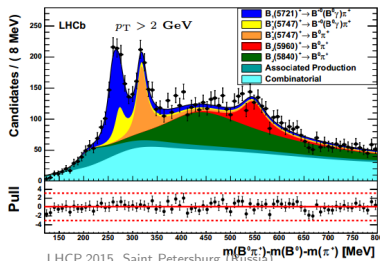
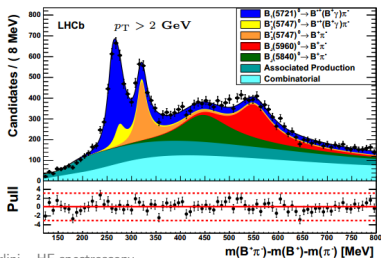
- $B^- \rightarrow J/\psi K^-$;
- $B^{-,0} \rightarrow D^{0,-} \pi^+(\pi^+ \pi^-)$ with $D^0 \rightarrow K^- \pi^+(\pi^- \pi^+)$
- $B^0 \rightarrow J/\psi K^*(892)^0$;

Improved measurement of the masses of $B_1^{0,+}$ and $B_2^{*0,+}$.

Fit parameter	$B^+ \pi^-$	$B^0 \pi^+$
$B_1(5721)^{0,+} \mu$	263.9 ± 0.7	260.9 ± 1.8
$B_1(5721)^{0,+} \Gamma$	30.1 ± 1.5	29.1 ± 3.6
$B_2^*(5747)^{0,+} \mu$	320.6 ± 0.4	318.1 ± 0.7
$B_2^*(5747)^{0,+} \Gamma$	24.5 ± 1.0	23.6 ± 2.0

First evidence (3.7σ) of $B_2^* \rightarrow B^* \pi^+$.

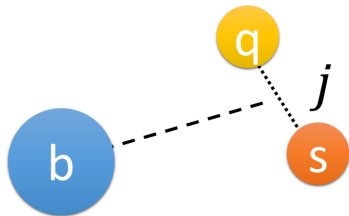
Enhancement at $5.97 \text{ GeV}/c^2$ reported by CDF is confirmed and labeled $B(5960)$.



LHCb studies the $\Xi_b^0 \pi^-$ combinations where

$$\Xi_b^0 \rightarrow \Xi_c^+ \pi^- \quad \text{and} \quad \Xi_c^+ \rightarrow \Lambda_c^+ p K^-$$

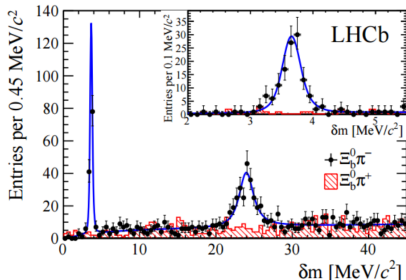
Taking wrong-sign combination as background proxy.



LHCb observes two new states, labeled Ξ_b' and Ξ_b^* .

Angular analysis does not exclude quark model interpretation:

$$\begin{aligned} \Xi_b & \quad J^P = \frac{1}{2}^+ & j = 0 \\ \Xi_b' & \quad J^P = \frac{1}{2}^- & j = 1 \\ \Xi_b^* & \quad J^P = \frac{3}{2}^+ & j = 1 \end{aligned}$$



Very precise mass measurement: $m(\Xi_b'^{-}) = 5935.02 \pm 0.02 \pm 0.01 \pm 0.50 \text{ MeV}/c^2$
 $m(\Xi_b^{*-}) = 5955.33 \pm 0.12 \pm 0.06 \pm 0.50 \text{ MeV}/c^2$
 Last uncertainty propagated from $m(\Xi_b^0)$ [PDG]: $m(\Xi_b^0) = 5971 \pm 0.50 \text{ MeV}/c^2$

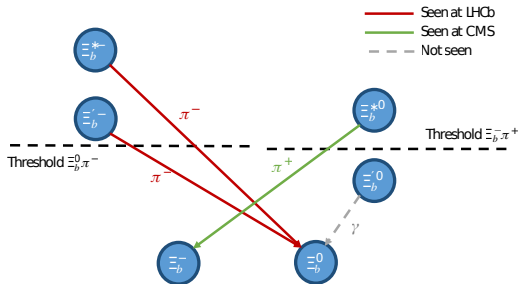
A Ξ_b^* state was already reported at CMS studying $\Xi_b^- \pi^+$ combinations.

Measured mass:

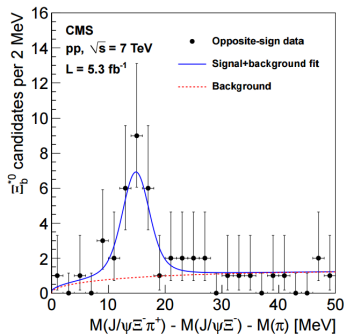
$$m(\Xi_b^{*0}) = 5945.0 \pm 0.7 \pm 0.3 \pm 2.7 \text{ MeV}/c^2$$

where the last uncertainty is propagated from $m(\Xi_b^-)$.

A possible interpretation...



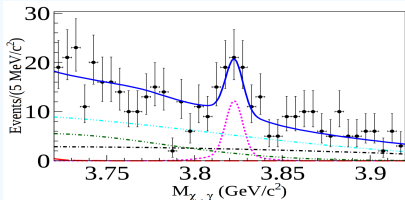
Analysis based on 2011 data,
 $\sqrt{s} = 7 \text{ TeV}$



$N_{\text{sig}} \sim 18$ candidates
Significance (including LEE): 5.3σ

$Q\bar{Q}$ and $Q\bar{Q}$ -like spectroscopy

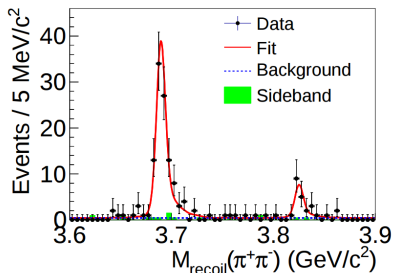
In 2013, Belle reports an evidence (3.8σ) for a state $X(3823) \rightarrow \chi_{c1}\gamma$. [PRL 111, 032001]



Its mass ($3823.1 \pm 1.8 \pm 0.7 \text{ MeV}/c^2$) was found consistent with the interpretation as $\psi_2(1^3D_2)$ state.

No peak seen in $D\bar{D}$: also consistent with $\psi_2(1^3D_2)$.

In 2015, BESIII confirms the state (6.2σ).



Improved measurement of mass and width:

$$M = 3821 \pm 1.3 \pm 0.7 \text{ MeV}/c^2$$

$$\Gamma < 16 \text{ MeV}/c^2 (90\% \text{C.L.})$$

$$R = \frac{\mathcal{B}(X(3823) \rightarrow \chi_{c2}\gamma)}{\mathcal{B}(X(3823) \rightarrow \chi_{c1}\gamma)} < 0.43 (90\% \text{C.L.})$$

Potential model predictions for $\psi_2(1^3D_2)$

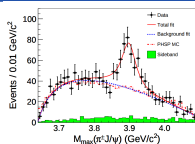
[PRD 69:094019]

$$M \sim 3815 \text{ MeV}/c^2$$

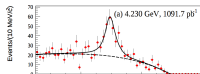
$$R \sim 0.2$$

Γ small ($D\bar{D}$ suppressed for P violation)

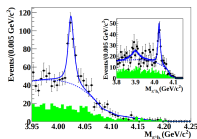
$Z_c(3900)^\pm?$



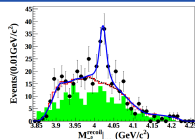
$Z_c(3900)^0?$



$Z_c(4020)^\pm?$



$Z_c(4020)^0?$



$e^+e^- \rightarrow \pi^- \pi^+ J/\psi$

[PRL 110, 252001]

$e^+e^- \rightarrow \pi^0 \pi^0 J/\psi$

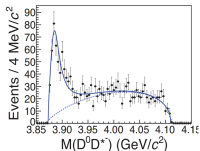
[1506.06018]

$e^+e^- \rightarrow \pi^- \pi^+ h_c$

[PRL 111, 242001]

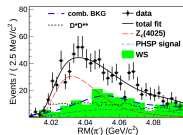
$e^+e^- \rightarrow \pi^0 \pi^0 h_c$

[PRL 113, 212002]



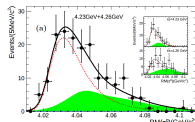
$e^+e^- \rightarrow \pi^\pm (D \bar{D}^*)^\mp$

[PRL 112, 022001]



$e^+e^- \rightarrow \pi^\pm (D^* \bar{D}^*)^\mp$

[PRL 112, 132001]



$e^+e^- \rightarrow \pi^0 (D^* \bar{D}^*)^0$

[1507.02404]

Most of these papers were submitted in 2014 and 2015. No angular analysis yet.

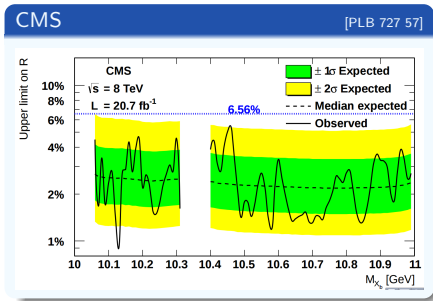
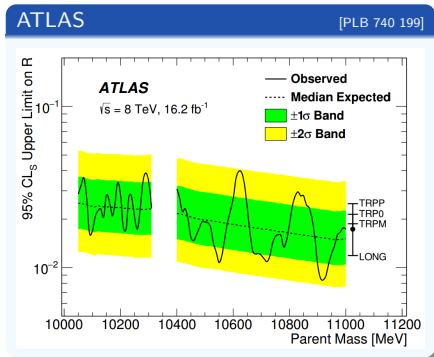
Search for $X_b \rightarrow \Upsilon(1S)\pi^+\pi^-$ at ATLAS and CMS

ATLAS and CMS report a search for a state analogue to the $X(3872)$ in the b -sector.

In analogy to $X(3872) \rightarrow J/\psi\pi^+\pi^-$ they look for $X_b \rightarrow \Upsilon(1S)\pi^+\pi^-$, using $\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-$ as control channel.

In absence of signal, they put an upper limit (at 95% of C.L.) on

$$R = \frac{\mathcal{B}(X_b \rightarrow \Upsilon(1S)\pi^+\pi^-)}{\mathcal{B}(\Upsilon(2S) \rightarrow \Upsilon(1S)\pi^+\pi^-)}$$



For comparison:

$$R = \frac{\mathcal{B}(X(3872) \rightarrow J/\psi\pi^+\pi^-)}{\mathcal{B}(\psi(2S) \rightarrow J/\psi(1S)\pi^+\pi^-)} = 6.56 \pm 0.29 \pm 0.65\%$$

Making short a long history

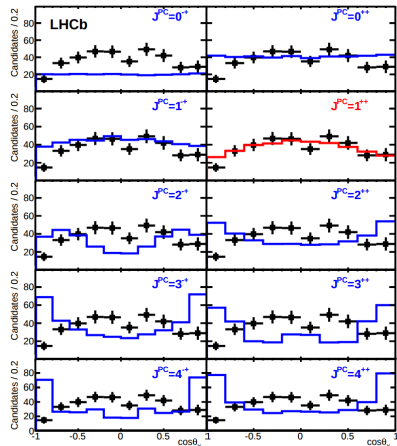
- $X(3872)$ observed at Belle (2003)
- CDF, D0, and Babar confirm it exists
- $X(3872) \rightarrow J/\psi\gamma$ implies $C = +$
- $X(3872) \rightarrow J/\psi\rho^0$ implies $J_{\pi\pi} = 1$
- Angular analysis at CDF finds $J^{PC} = 1^{++}$ or 2^{-+}
- With 2011 LHCb data: $J^{PC} = 1^{++}$.

Extension to full Run1 dataset (LHCb)

LHCb and CDF analyses rely on the assumption of a matrix element dominated by the lowest angular momentum (L_{\min}) between products.

Exotic nature of $X(3872)$ can introduce contribution of $L_{\min} + 2$, invalidating 1^{++} assignment.

This is the first analysis with no assumption on L .



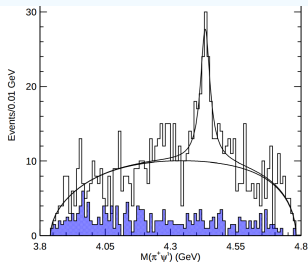
$\theta_X \equiv$ helicity angle of $X(3872)$

$$J^{PC} = 1^{++}$$

All alternative hypotheses discarded by 16σ or more.

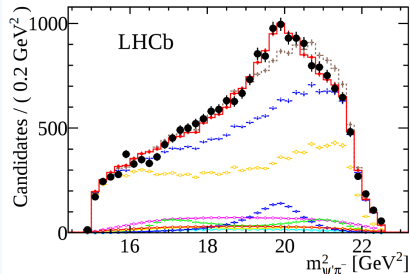
Observation at Belle

[PRL 100:142001]



Confirmation at LHCb (with Babar's mass)

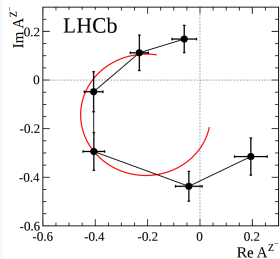
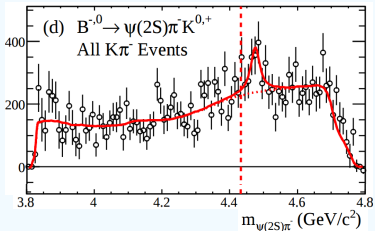
[PRL 112, 222002]



with unambiguous resonant behaviour

2.7σ excess at Babar (inconsistent mass)

[PRD 79 112001]



PRELIMINARY RESULT!

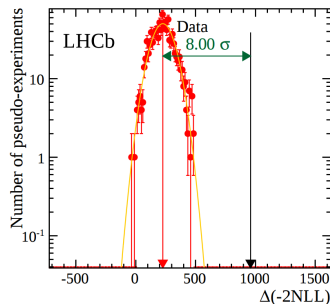
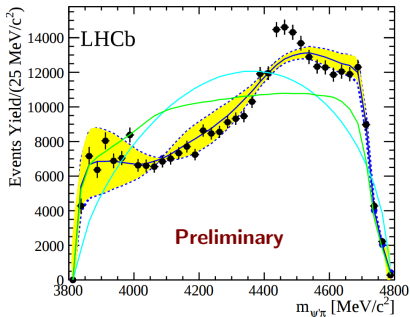
Following the example of the Babar's analysis, LHCb attempts to describe the $\psi(2S)\pi^-$ mass in terms of reflections of the $K\pi$ system.

- In fine bins of $m(K\pi)$, decompose distribution $\cos(\theta_K^*)$ on the first ℓ_{\max} spherical harmonics (**moments**).
- Generate phase-space simulation (**cyan**)
- Reweight MC to reproduce $m(K\pi)$ distribution (**green**)
- Reweight MC according to the moments (**yellow**)

K^* resonance with ℓ larger than 6 are unphysical: $\ell_{\max} = 6$.

Conclusion: $K\pi$ structures with reasonable angular momentum are insufficient to explain the $m_{\psi'\pi}$ distribution.

The null hypothesis, with no $Z_c(4430)$ is rejected by 8σ , or more.



LHCb extended the study of the $\psi(2S)\pi^-$ system to the $\bar{B}_s^0 \rightarrow \psi(2S)K^+\pi^-$ decay.

The decay is largely dominated by
 $\bar{B}_s^0 \rightarrow \psi(2S)K^*(892)^0$.

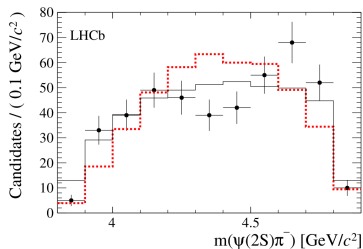
LHCb reports the first observation of the non resonant decay.

The small relative branching ratio

$$\frac{\mathcal{B}(\bar{B}_s^0 \rightarrow \psi(2S)K^+\pi^-)}{\mathcal{B}(B^0 \rightarrow \psi(2S)K^+\pi^-)} = 5.38 \pm 0.36 \pm 0.22 \pm 0.31 \left(\frac{f_s}{f_d} \right)$$

makes the statistics limited.

No structure observed in $m(\psi(2S)\pi^-)$ spectrum.



solid histogram: amplitude fit model
 dotted histogram: phase-space

Spectroscopy studies are very active at the LHC and all over the world.

In the last year

- New excited b -mesons and b -baryons were observed,
- The first spin-3 charm mesons were found in b -decays,
- An excited B_c^+ state was seen,
- A tetraquark was unambiguously identified as such,
- Two pentaquarks were observed.

Enhance interest in $B \rightarrow DDX$ decays as probe of over-threshold charmonia, and as play-ground for excited charmed states (e.g. $D_s^*(2700)^+$).

All of the studies presented here are dominated by statistical uncertainties.

The higher HF production cross-section, and the increased luminosity make the **second run of the LHC extremely promising** for spectroscopy studies.

Teasers include:

- Double heavy baryons ($\Xi_{cc}, \Xi_{bc}, \dots$)
- More B_c excited states
- $B_c^+ \rightarrow D_{(s)}^{(*)+} D_{(s)}^{(*)+} h^-$ decays

Spare slides

Molecular interpretation of $X(3872)$ suggests other $D\bar{D}$ states may exist.

Belle searches for $X_1(3872)$, $X(3730)$, and $X(4014)$ to final states: $\eta_c\pi^+\pi^-$, $\eta_c\pi^0$, and $\eta_c\eta$ (where $\eta \rightarrow \gamma\gamma$ or $\eta \rightarrow \pi^+\pi^-\pi^0$).

The analysis strategy is based on **unseen** $B^\pm \rightarrow K^\pm\eta_c + nh$ decays.

Upper limits are given at 90% of C.L.

$\mathcal{B}(B^\pm \rightarrow K^\pm\eta_c\pi^+\pi^-)$	39×10^{-5}
$\eta_c\pi^+\pi^-$ from $X_1(3872)$	3.0×10^{-5}
$\eta_c\pi^+\pi^-$ from $Z_c(3900)$	4.7×10^{-5}
$\eta_c\pi^+\pi^-$ from $Z_c(4020)$	1.6×10^{-5}
$\mathcal{B}(B^\pm \rightarrow K^\pm\eta_c\omega)$	53×10^{-5}
$\eta_c\omega$ from $X_1(3872)$	6.9×10^{-5}
$\eta_c\omega$ from $Z_c(3900)^0$	4.7×10^{-5}
$\mathcal{B}(B^\pm \rightarrow K^\pm\eta_c\eta)$	22×10^{-5}
$\eta_c\eta$ from $X(3730)$	4.6×10^{-5}
$\eta_c\eta$ from $X(4014)$	3.9×10^{-5}
$\mathcal{B}(B^\pm \rightarrow K^\pm\eta_c\pi^0)$	6.2×10^{-5}
$\eta_c\pi^0$ from $X(3730)$	0.6×10^{-5}
$\eta_c\pi^0$ from $X(4014)$	1.2×10^{-5}
$\eta_c\pi^0$ from $X(3915)$	1.8×10^{-5}

Assumptions

	Molecule	
$X_1(3872)$	$D^0\bar{D}^{*0} - \bar{D}^0D^{*0}$	
$X(3730)$	$D^0\bar{D}^0 + \bar{D}^0D^0$	
$X(4014)$	$D^{*0}\bar{D}^{*0} + \bar{D}^{*0}D^{*0}$	
	J^{PC}	Modes
$X_1(3872)$	1^{+-}	$\eta_c\omega, \eta_c\rho$
$X(3730)$	0^{++}	$\eta_c\eta, \eta_c\pi^0$
$X(4014)$	0^{++}	$\eta_c\eta, \eta_c\pi^0$