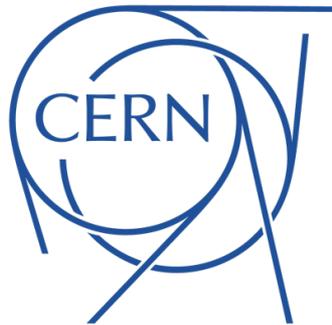


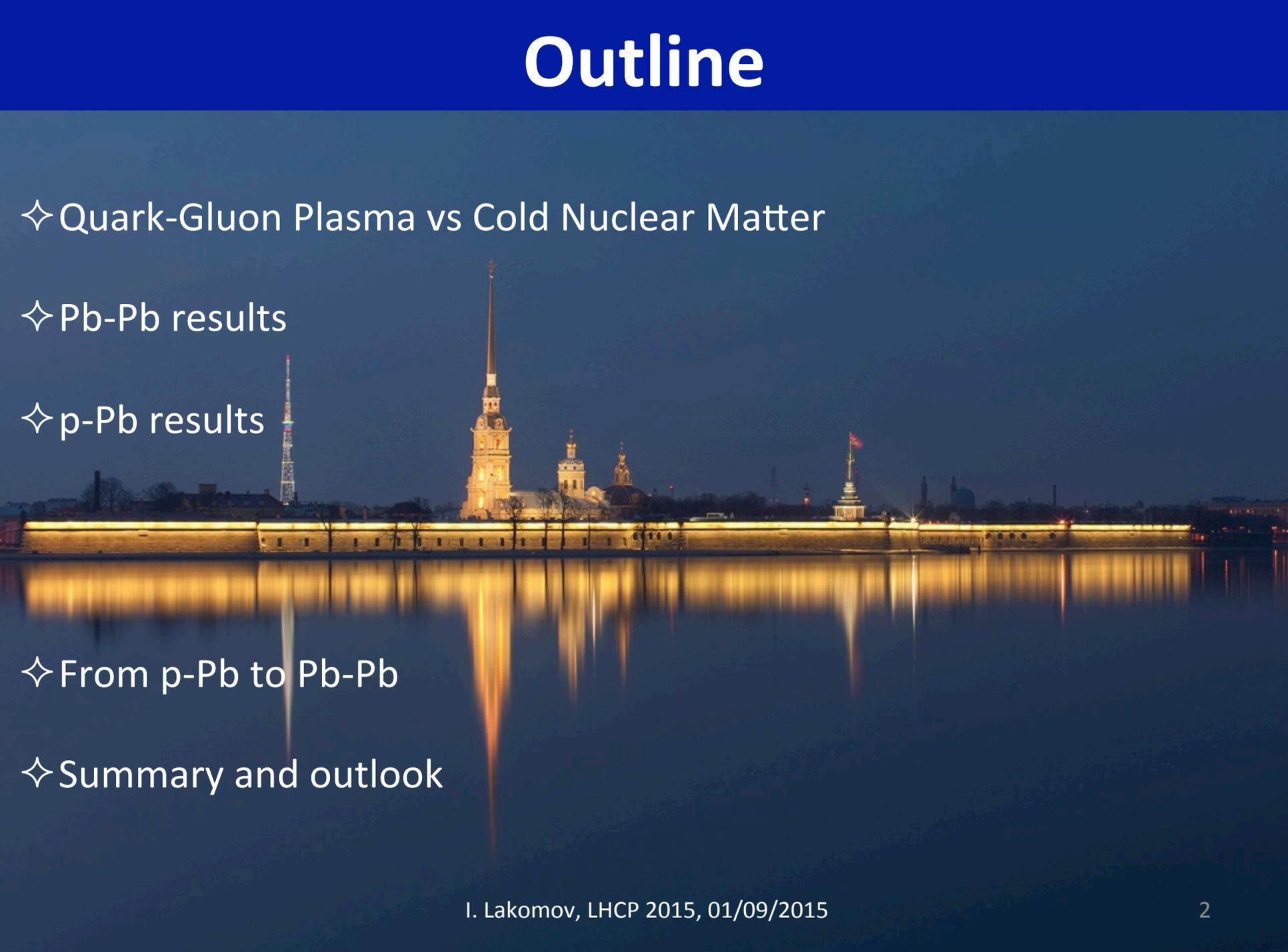
# Quarkonium production in p-A and A-A collisions

**Igor Lakomov (CERN)**

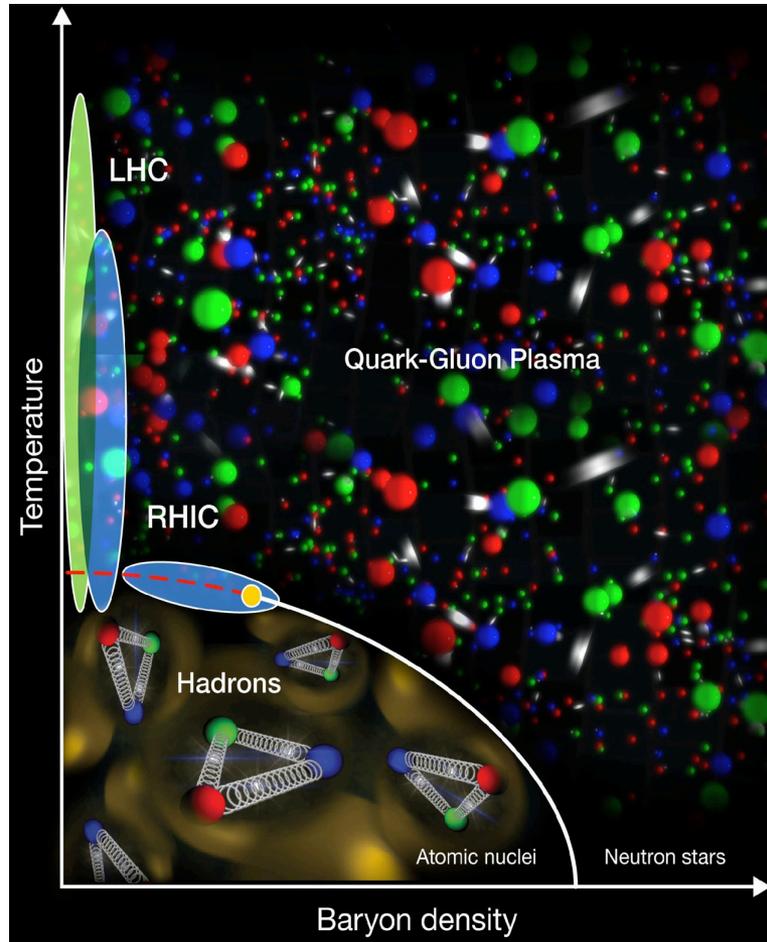


LHCP2015: The 3<sup>rd</sup> Conference on Large Hadron Collider Physics  
01/09/2015, Saint-Petersburg

# Outline

- ✧ Quark-Gluon Plasma vs Cold Nuclear Matter
  - ✧ Pb-Pb results
  - ✧ p-Pb results
  - ✧ From p-Pb to Pb-Pb
  - ✧ Summary and outlook
- 

# Quark-Gluon Plasma



- QCD: strong force describes the interactions between quarks and gluons forming hadrons.
- Lattice QCD predicts a deconfined state of matter (QGP) at high temperature.
- QGP can be recreated in Heavy-Ion Collisions (HIC) at hadron colliders.
- QGP lifetime is small ( $\sim 2-4$  fm/c at RHIC,  $\sim 15-20$  fm/c at the LHC)\*, a direct observation of the QGP is not possible.
- Experimental probes of QGP: jet quenching, strangeness enhancement, quarkonia, etc.

\*  $1 \text{ fm/c} \sim 10^{-24} \text{ sec}$

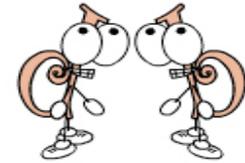
# Quarkonia, a key tool for the QGP

✧ Bound states of charm or beauty quark and its anti-quark

✧ Heavy and tightly bound



$c + \bar{c} = \text{charmonia } (J/\psi, \psi(2S), \dots)$

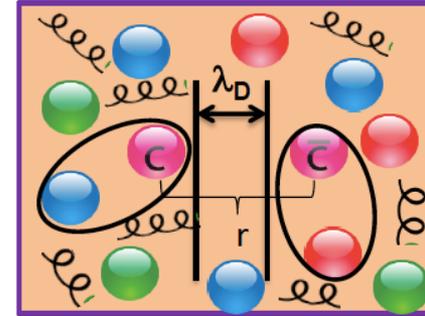


$b + \bar{b} = \text{bottomonia } (\Upsilon(nS))$

✧ Heavy quark pairs produced in the initial hard partonic collisions.

✧ **Suppressed** by Debye color screening:

- ❖ Color charge of one quark masked by surrounding quarks.
- ❖ Prevents  $q\bar{q}$  binding in the QGP.
- ❖ Debye screening radius ( $\lambda_D$ ) vs quarkonium radius ( $r$ ).
- ❖  $\lambda_D < r \Rightarrow$  the quarks are effectively masked from each other.

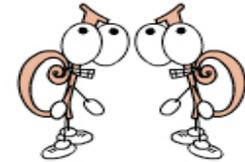


# Suppression vs enhancement

✧ Bound states of charm or beauty quark and its anti-quark



$c + \bar{c} = \text{charmonia } (J/\psi, \psi(2S), \dots)$



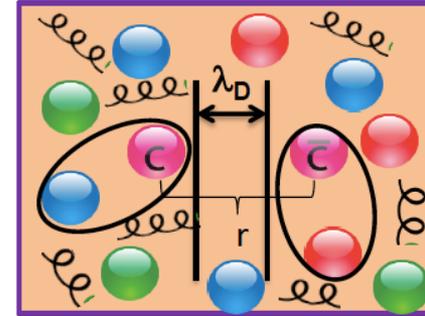
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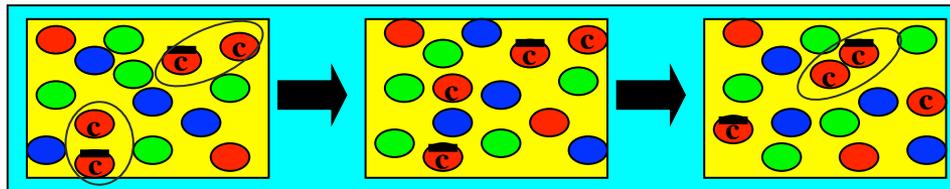
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✧ **Recombination**

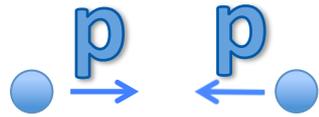
- ✓ In central HIC,  $N_{cc} > 1$ . (RHIC:  $\sim 10$ ; LHC:  $\sim 100$ ).
- ✓ Regeneration of  $J/\psi$  pairs possible from independently produced  $c$  &  $\bar{c}$



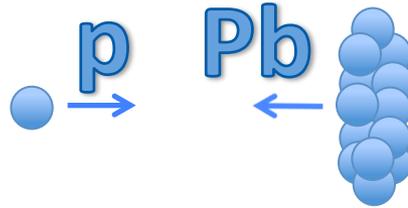
➤ Leads to an enhancement of  $J/\psi$  (or less dramatic suppression if 2 effects compete).

☐ No/small regeneration is expected for bottomonia.

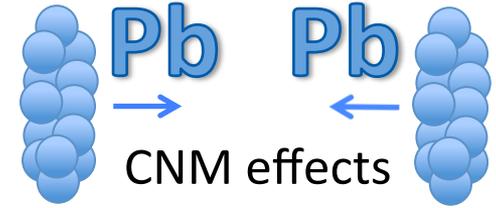
# Nuclear matter effects



Elementary collision  
No nuclear matter effects



Cold nuclear matter (CNM)  
effects without QGP



CNM effects  
+ Hot nuclear matter effects  
(related to QGP formation)



# Nuclear matter effects

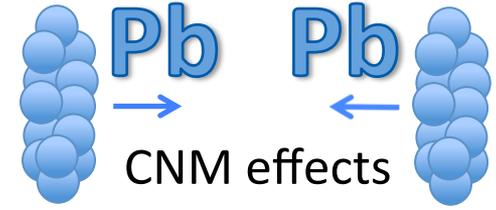


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Cold nuclear matter (CNM)  
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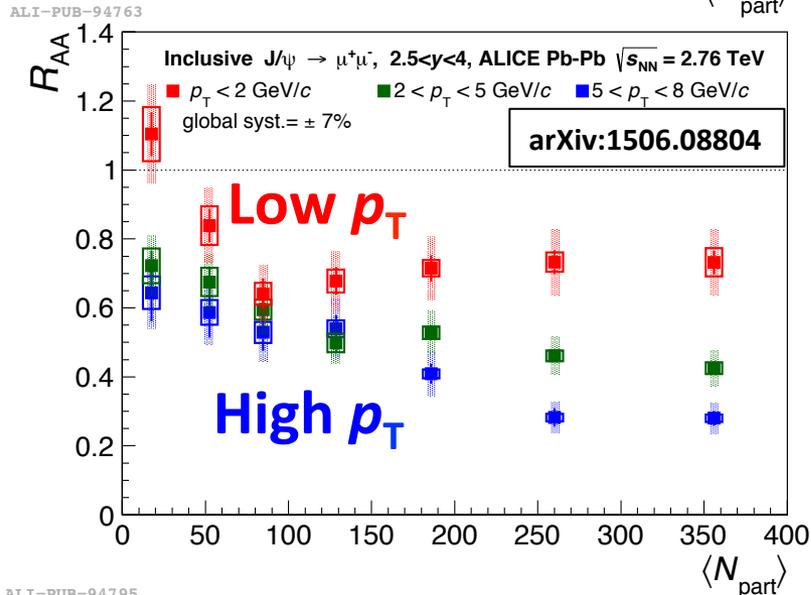
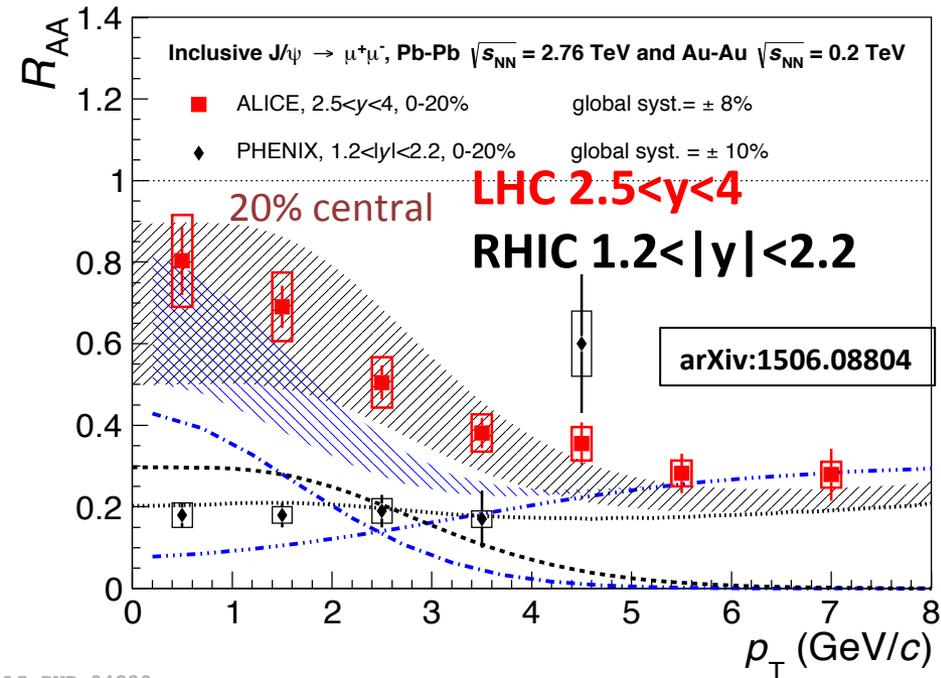
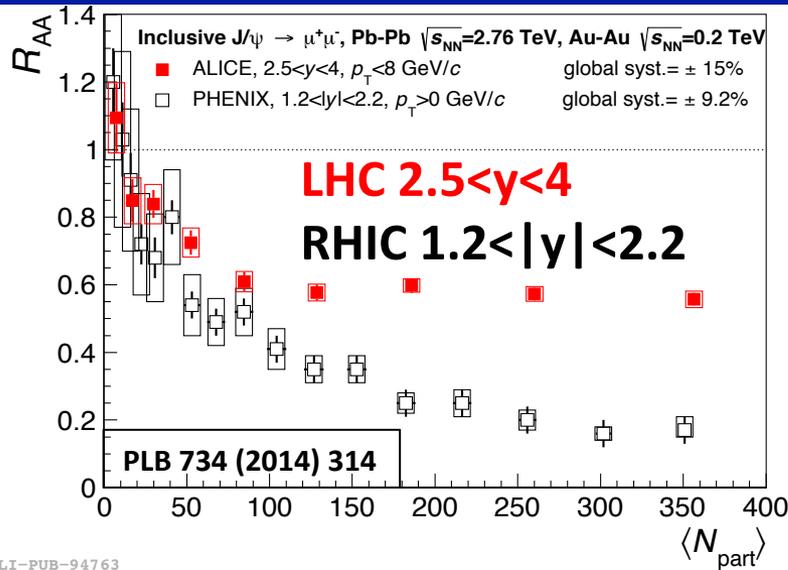
## Main observable:

$$R_{AA} = \frac{Y_{AA}}{N_{coll} Y_{pp}}$$

ratio of the production yield in AA to that in pp, scaled by the number of binary nucleon-nucleon collisions.

❖ If  $R_{AA} \neq 1 \Rightarrow$  there are some nuclear matter effects.

# J/ψ suppression in Pb-Pb

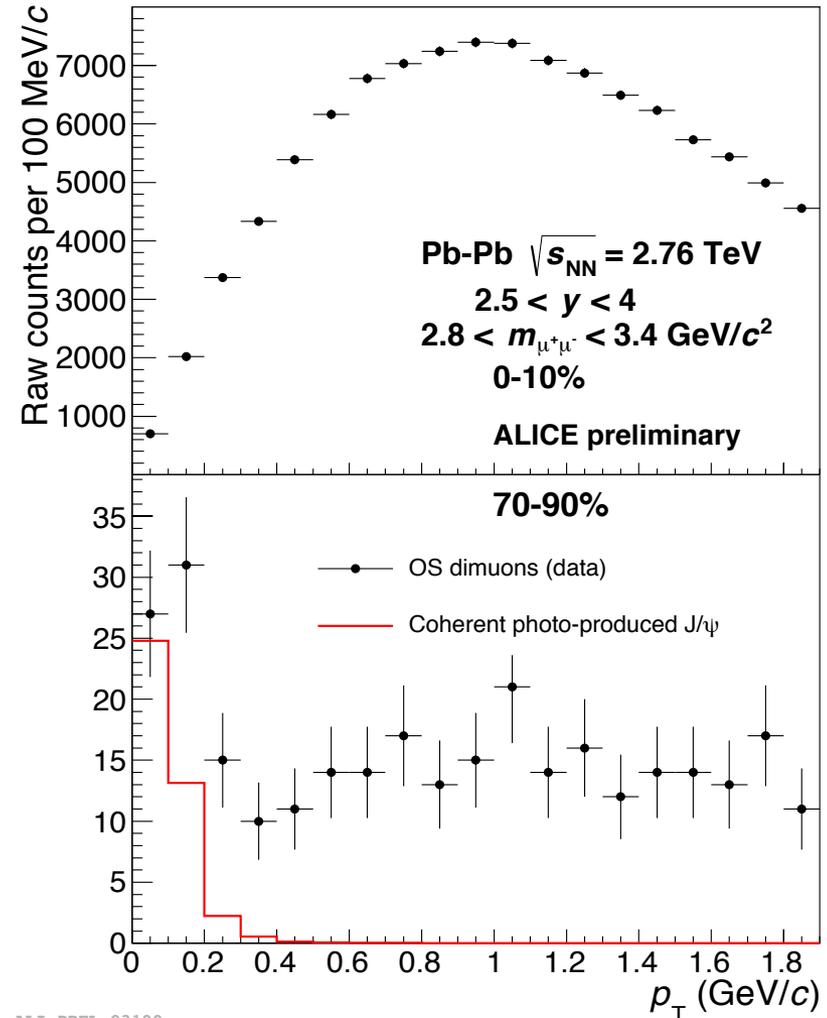
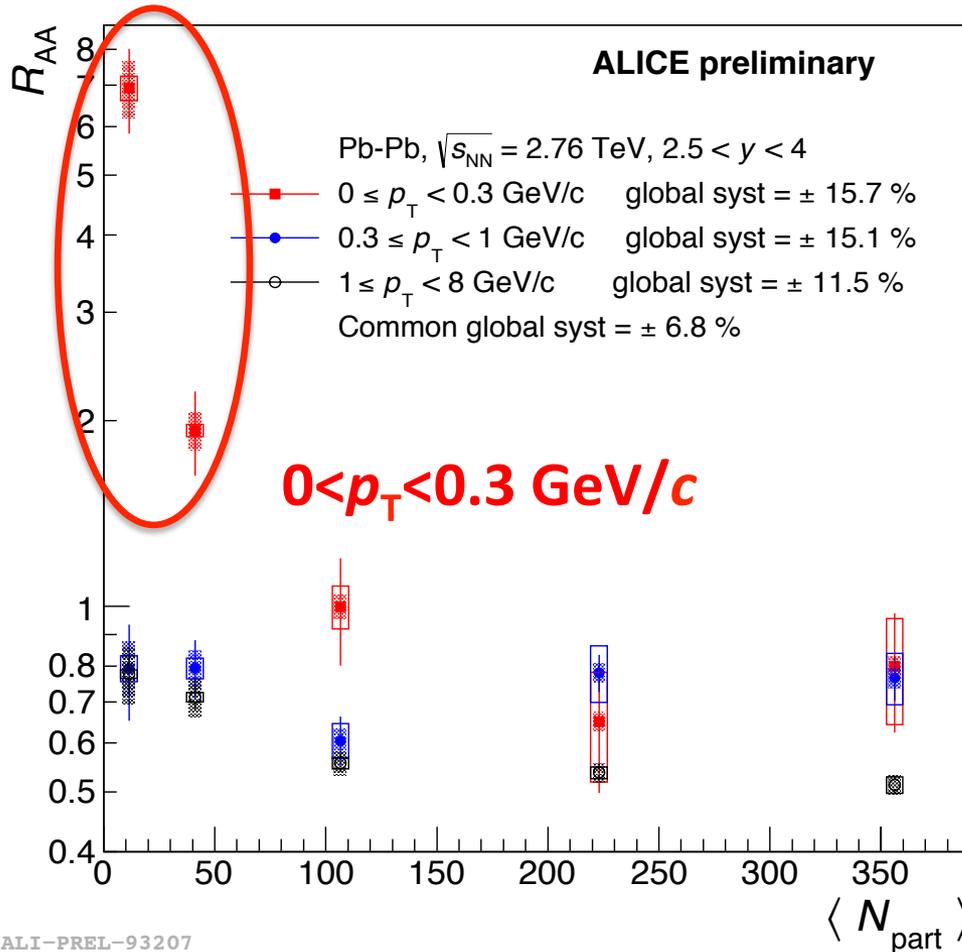


- Different  $R_{AA}$  at LHC vs RHIC: recombination.
- Regeneration is larger at low  $p_T$ .
- High- $p_T$  J/ψ are suppressed.
- J/ψ suppression is less pronounced in central compared to peripheral collisions.

ALI-PUB-94820

ALI-PUB-94795

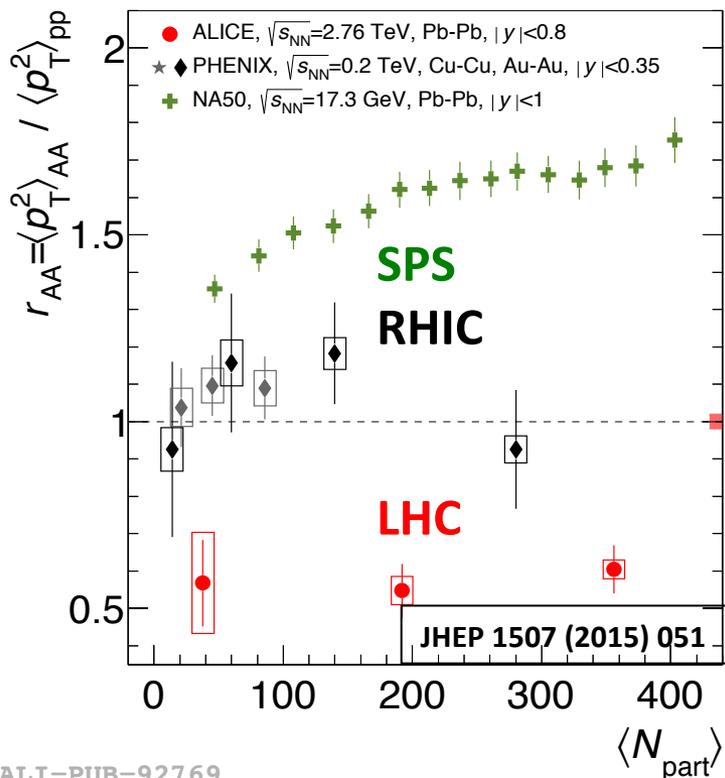
# Low- $p_T$ J/ $\psi$ in Pb-Pb



➤ Huge  $R_{AA}$  increase at low  $N_{part}$ .

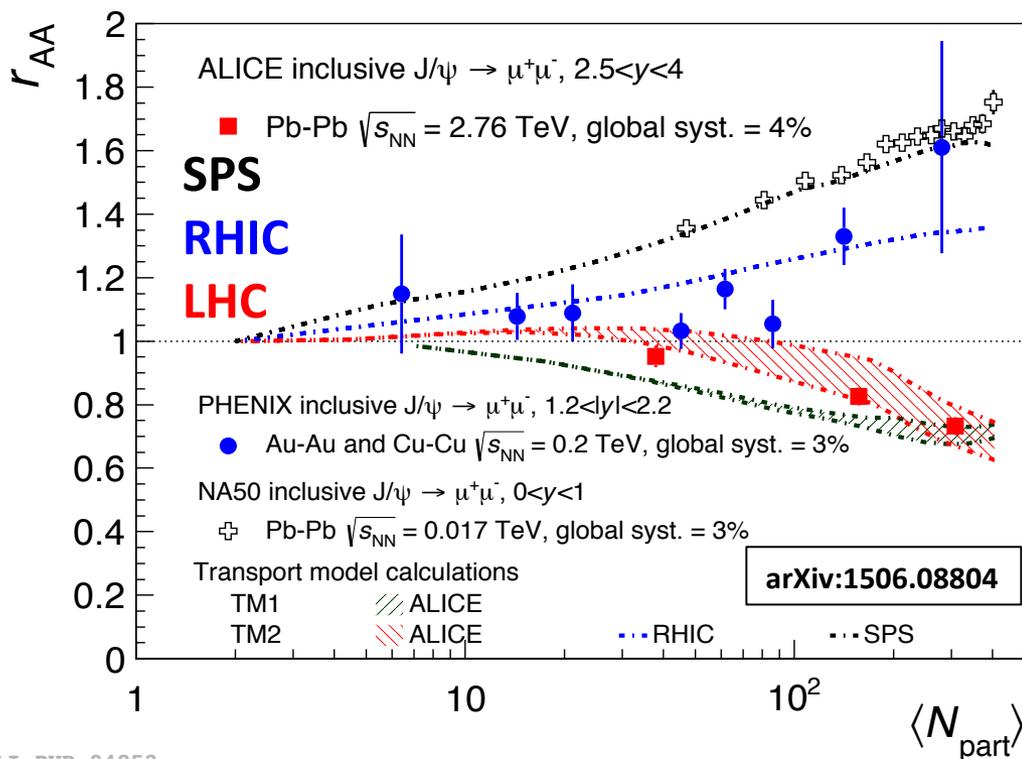
➤  $p_T$  spectrum is similar to the photo-production seen in UPC ( $b > 2R$ ).

# $\langle p_T \rangle$ of J/ $\psi$ in Pb-Pb



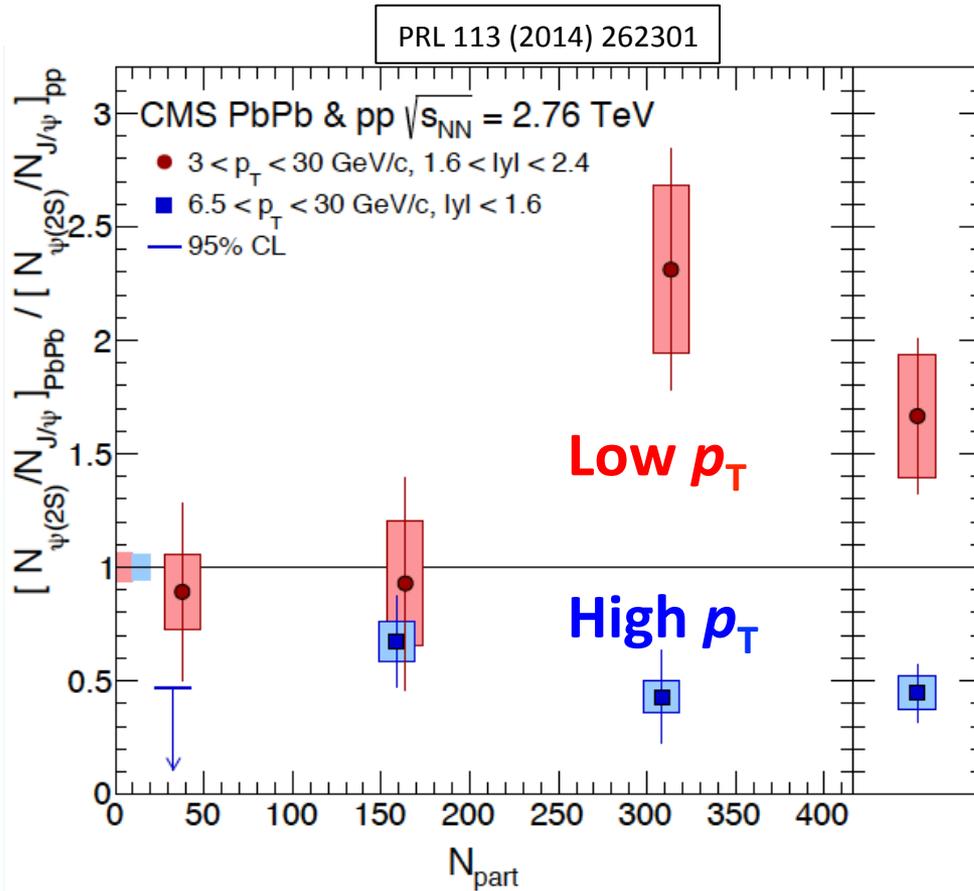
ALI-PUB-92769

ALI-PUB-94852



- Significant reduction of J/ $\psi$   $\langle p_T^2 \rangle$  in Pb-Pb as compared to pp collisions.
- Opposite trend at RHIC and SPS energies.
- Transport model with regeneration component reproduces data at different energies.

# $\psi(2S)$ vs $J/\psi$ in Pb-Pb

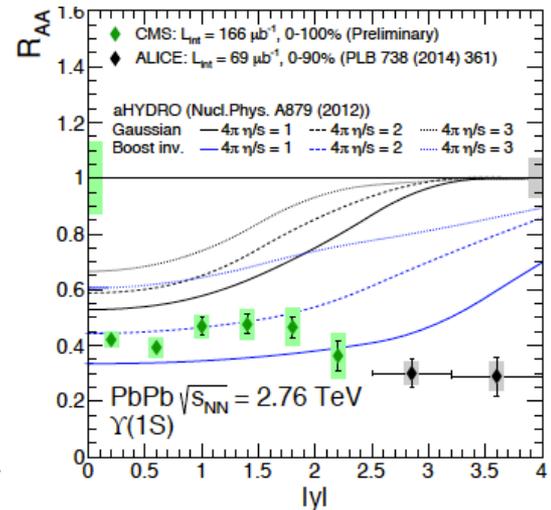
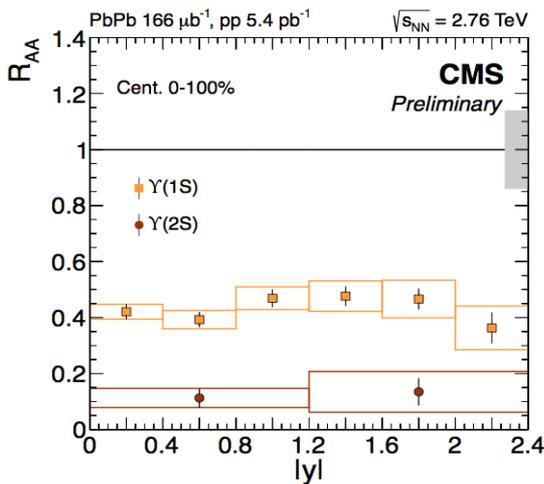
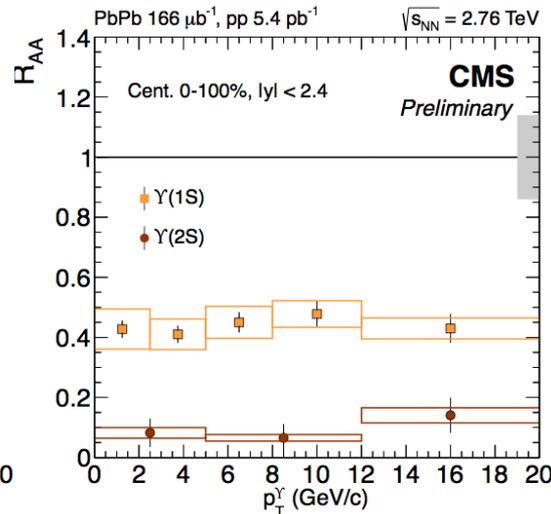
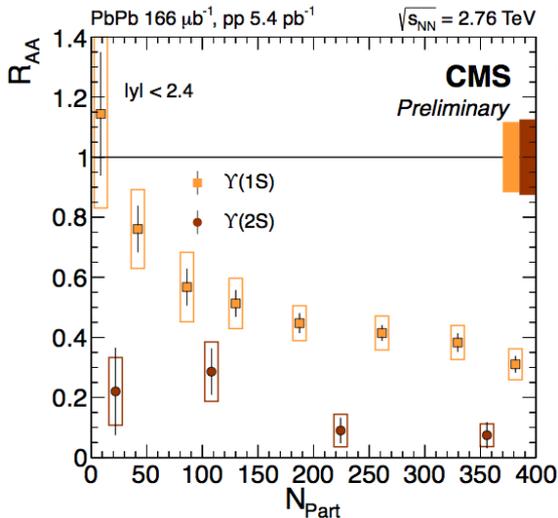


- High- $p_T$   $J/\psi$  and  $\psi(2S)$  are suppressed, consistent with the sequential melting.
- At lower  $p_T$ , and forward  $y$ , less  $\psi(2S)$  suppression.
- At high  $p_T$ , enhanced suppression in central events.

# Bottomonia in Pb-Pb

PLB 738 (2014) 361

CMS-PAS-HIN-15-001

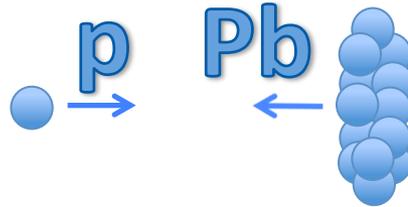


- Update of CMS results.
  - New reference at same energy.
  - Larger statistics.
- $R_{AA}$  decreases with centrality down to 0.3.
- $R_{AA}(p_T) \approx$  constant in 0-100%.
- Flat  $R_{AA}(y)$ , closer to ALICE now.
- Zero or negligible statistical regeneration for  $\Upsilon(1S)$ .

# Nuclear matter effects



Elementary collision  
No nuclear matter effects

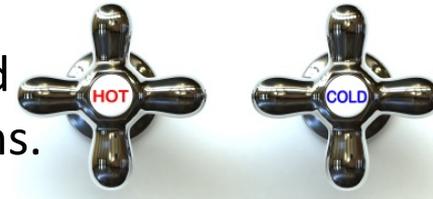


Cold nuclear matter (CNM)  
effects without QGP



CNM effects  
+ Hot nuclear matter effects  
(related to QGP formation)

➤ To disentangle hot and CNM effects, p-Pb collisions are needed as an intermediate step between Pb-Pb and benchmark pp collisions.



❖ In p-Pb collisions different kinds of CNM effects can be considered:

## ① Initial-state:

- ✓ gluon shadowing
- ✓ gluon saturation

## ② Final-state:

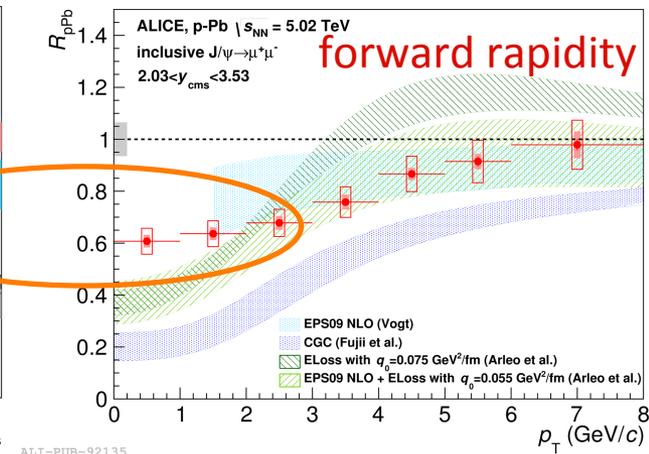
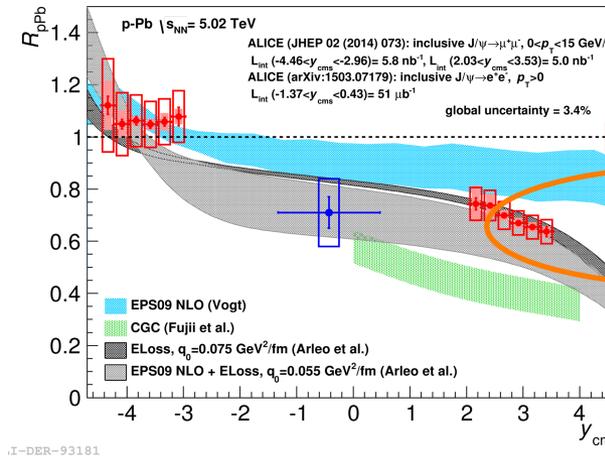
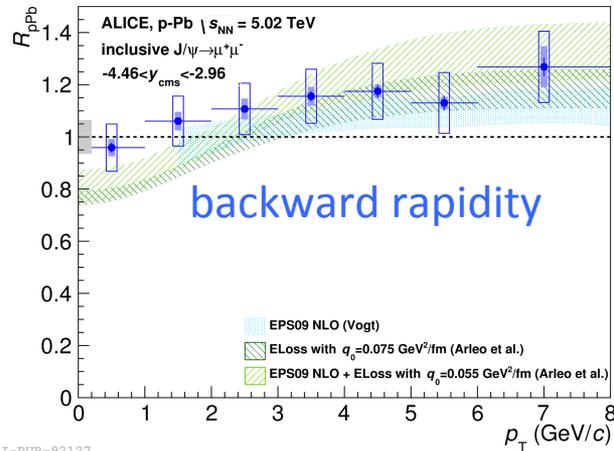
- ✓ nuclear absorption

## ③ Coherent parton energy loss

# J/ψ vs $p_T$ and $y$ in p-Pb

JHEP 1402 (2014) 073

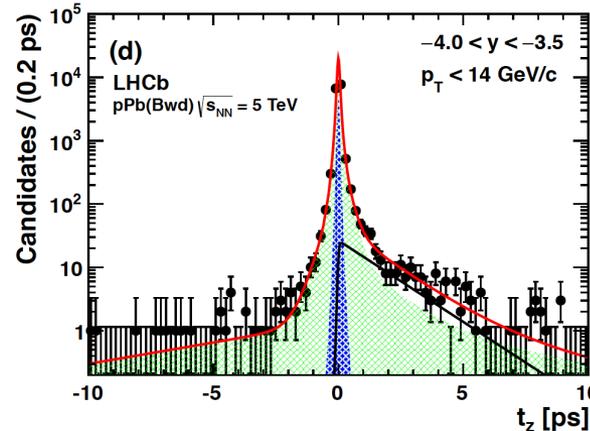
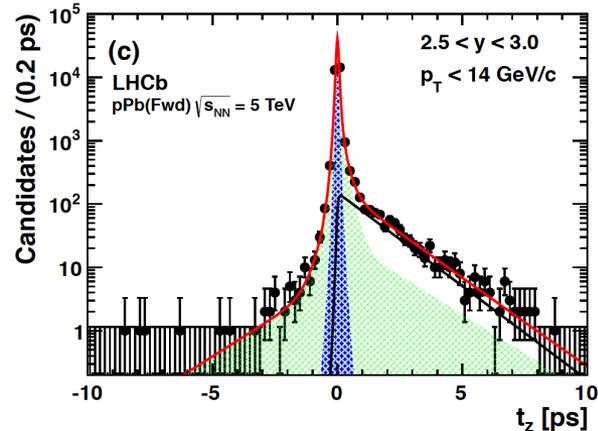
arXiv: 1503.07179



- $R_{pPb} \approx 1$  for all  $p_T$  at backward  $y$ , and for high  $p_T$  at forward  $y$ .
- At forward  $y$ ,  $R_{pPb}$  increases with  $p_T$ .
- As a function of rapidity,  $R_{pPb}$  decreases from backward to forward  $y$ .
- **Shadowing and coherent Eloss**: fairly reproduce  $p_T$  and  $y$  dependence, except low  $p_T$  at forward  $y$ , where **coherent Eloss** underestimates the data.
- **CGC**: overestimates the suppression at forward- $y$  over the full  $p_T$  range.

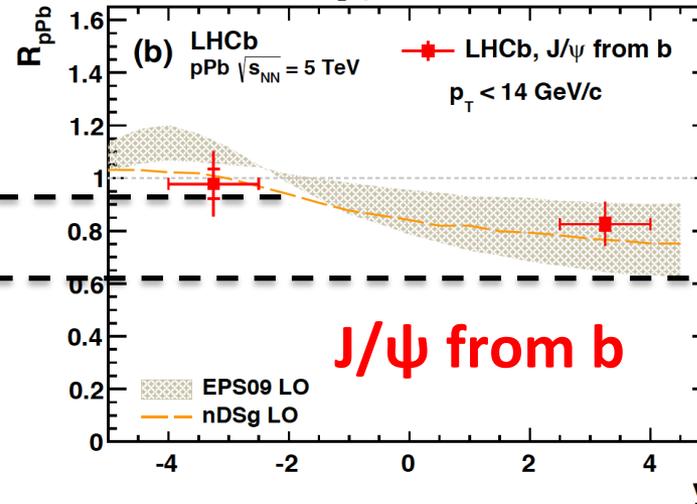
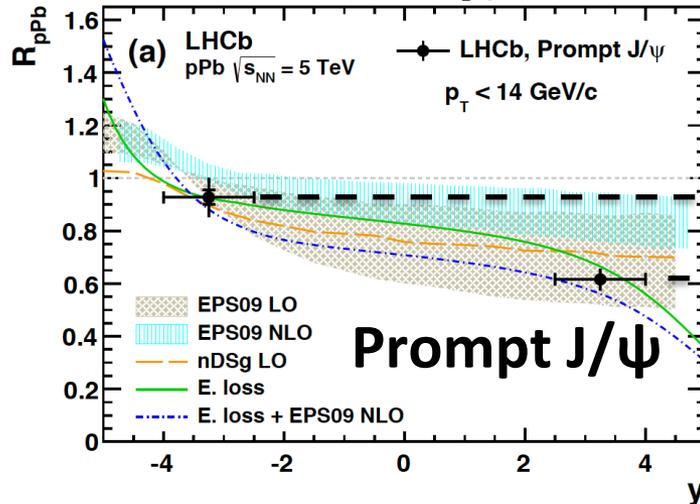
# Prompt J/ψ in p-Pb

JHEP 02 (2014) 072



$1.5 < y < 4$ :  
 $\sigma_{\text{prompt J}/\psi} = 1168 \pm 15 \pm 54 \mu\text{b}$  (88% inclusive)  
 $\sigma_{\text{J}/\psi \text{ from b}} = 166.0 \pm 4.1 \pm 8.2 \mu\text{b}$  (12% incl.)

$-2.5 < y < -5$ :  
 $\sigma_{\text{prompt J}/\psi} = 1293 \pm 42 \pm 75 \mu\text{b}$  (92% incl.)  
 $\sigma_{\text{J}/\psi \text{ from b}} = 118.2 \pm 6.8 \pm 11.7 \mu\text{b}$  (8% incl.)



- No difference between prompt and non-prompt J/ψ  $R_{pPb}$  at backward  $y$ .
- At forward  $y$ , the difference  $< 25\%$ .
- Small non-prompt fraction + large uncertainties  $\Rightarrow$  safe to compare to inclusive.

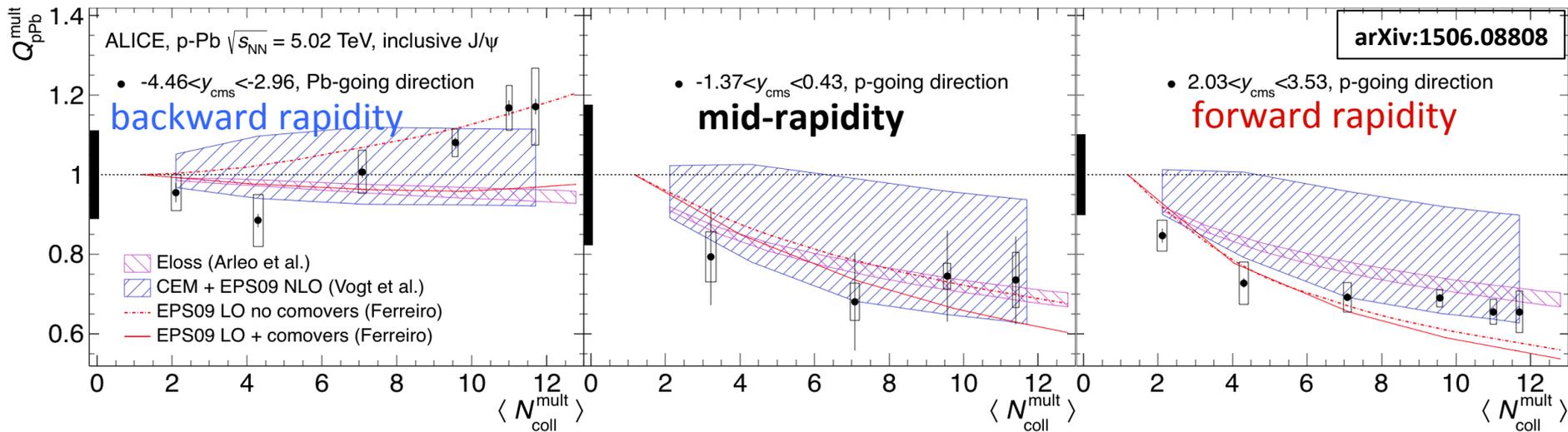
# J/ψ Q<sub>pPb</sub> vs centrality

$$Q_{pPb}^{J/\psi, i} = \frac{Y_{pPb}^i}{\langle T_{pPb}^i \rangle \sigma_{pp}^{J/\psi \rightarrow \mu^+ \mu^-}}$$

$T_{pPb}^i$  - nuclear thickness function in a given ZN energy event class\* i.

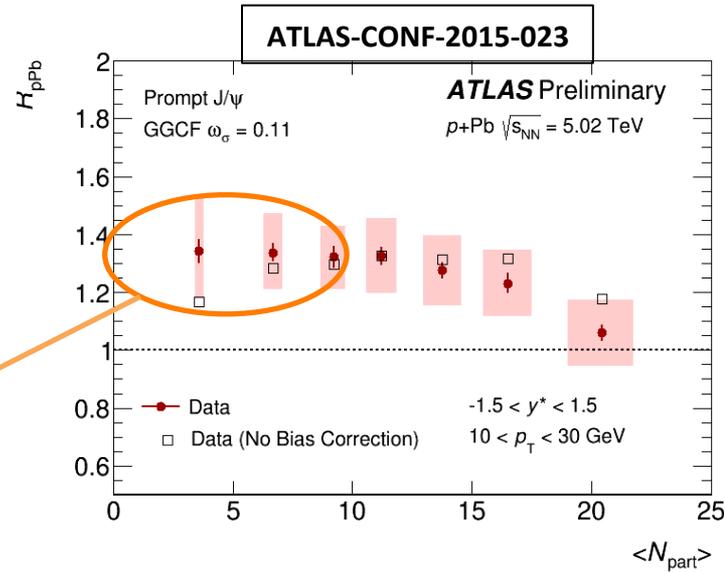
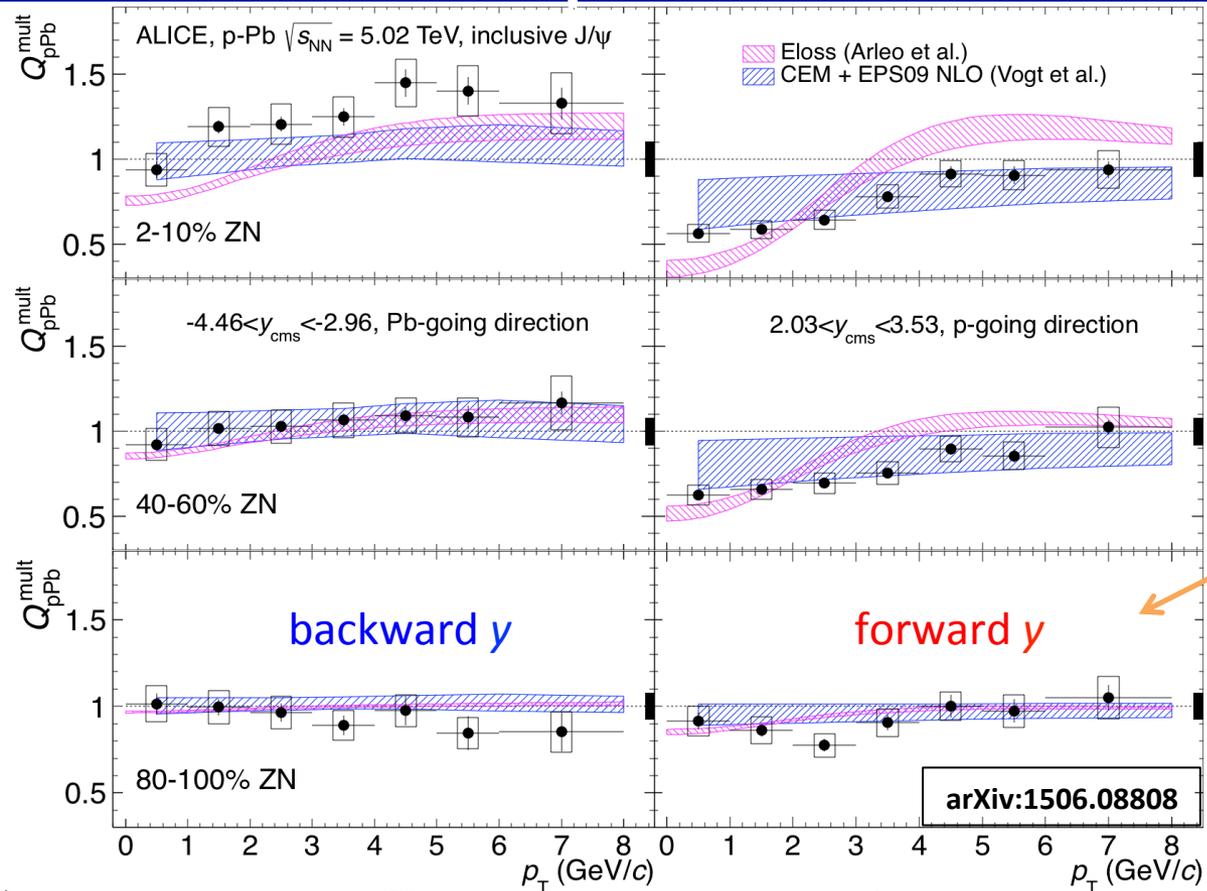
$\sigma_{pp}^{J/\psi \rightarrow \mu\mu}$  – interpolated pp cross-section at  $\sqrt{s} = 5.02$  TeV.

\*ZN is the neutron part of the ZDC



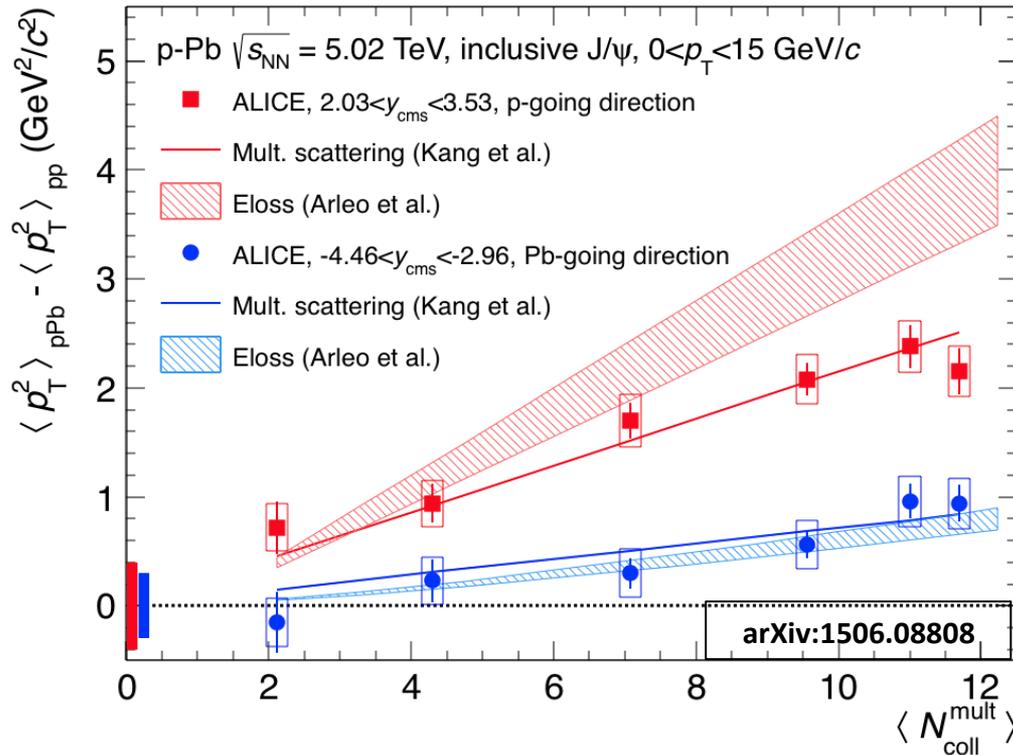
- Large J/ψ suppression at **forward y**, increasing with  $N_{coll}$ .
- Consistent with no CNM effects at **backward y**.
- Large uncertainties at **mid-y**.

# J/ψ $Q_{pPb}$ vs $p_T$ and centrality



- Large CNM effects in most central events:  $Q_{pPb}$  increases with  $p_T$  at both backward and forward  $y$ .
- At small centrality  $Q_{pPb}$  is consistent with unity for backward and forward  $y$ .
- Surprising trend in ATLAS at high  $p_T$  for the most peripheral collisions.

# J/ψ p<sub>T</sub> broadening



$$\Delta \langle p_T^2 \rangle = \langle p_T^2 \rangle_{pPb} - \langle p_T^2 \rangle_{pp},$$

$\langle p_T^2 \rangle_{pp}$  is from interpolated pp distributions at  $\sqrt{s} = 5.02$  TeV

\*Boxes at  $\Delta \langle p_T^2 \rangle = 0$ : uncertainty from  $\langle p_T^2 \rangle_{pp}$  at  $\sqrt{s_{NN}} = 5.02$  TeV

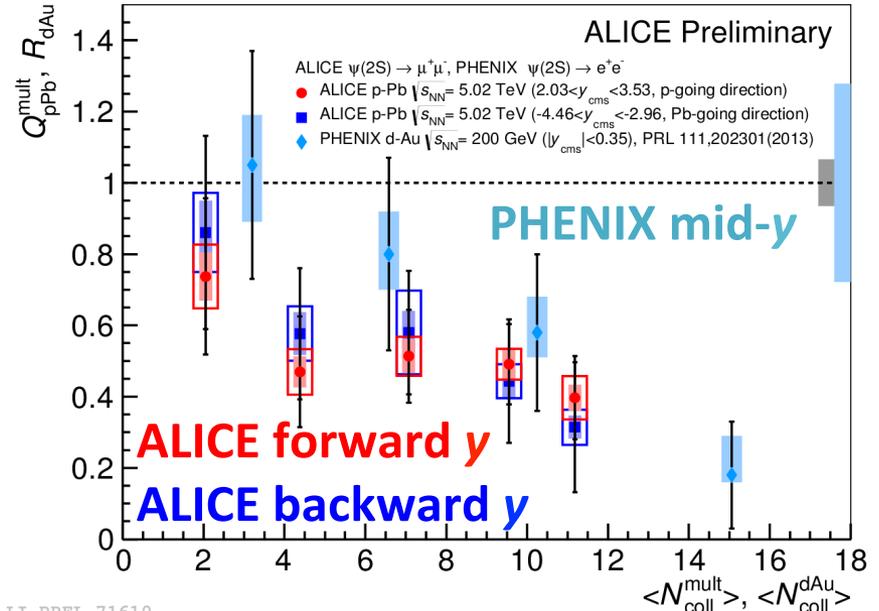
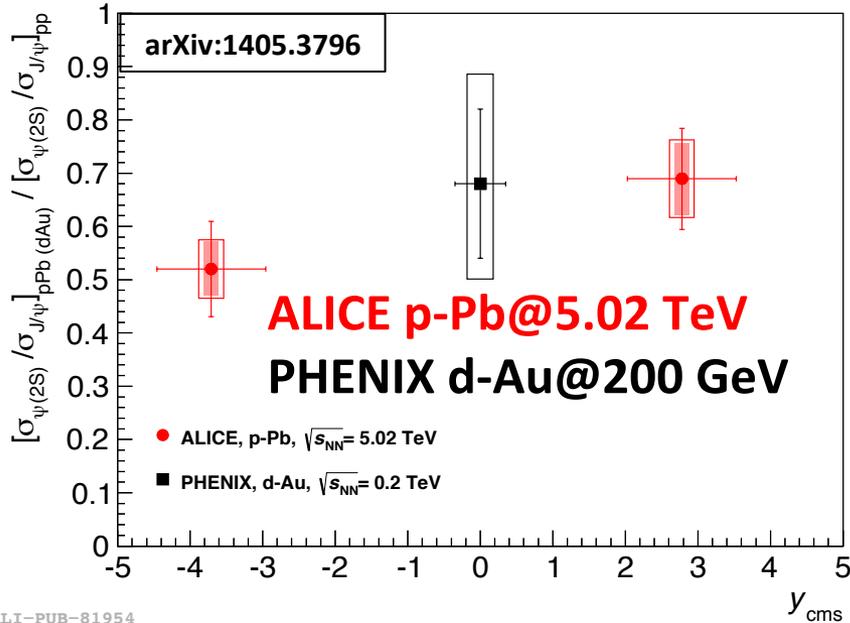
## Models

- **Multiple scattering** model (Kang et al. Phys.Rev. D77 (2008) 114027, Phys.Lett. B721 (2013) 277).
- **Coherent energy loss** (Arleo et al., JHEP1305(2013)155).

- Harder  $p_T$  distribution at **forward**  $y$  than at **backward**  $y$ .
- At **backward**  $y$ , the  $\langle p_T^2 \rangle_{pPb} \approx \langle p_T^2 \rangle_{pp}$  in peripheral collisions.
- Such a strong  $p_T$  broadening indicates a presence of Multi-Parton Interactions in p-Pb collisions.

# $\psi(2S)$ vs $J/\psi$ in p-Pb

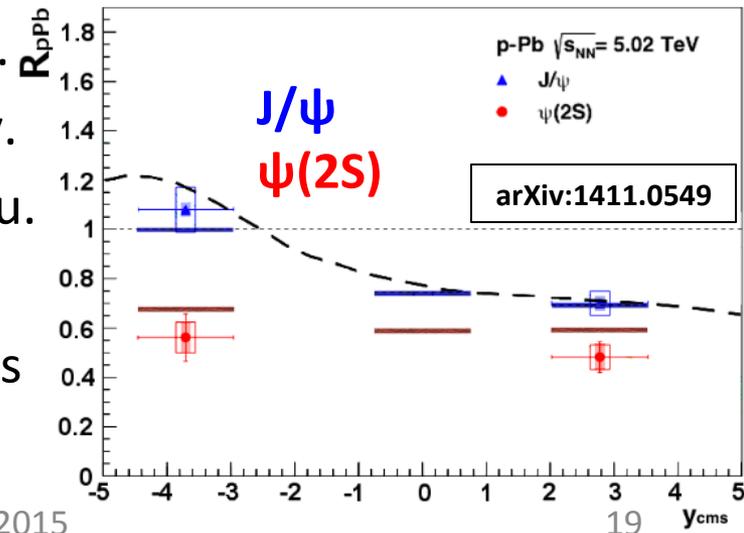
✧  $\psi(2S)$  is more weakly bound state with respect to  $J/\psi$ .



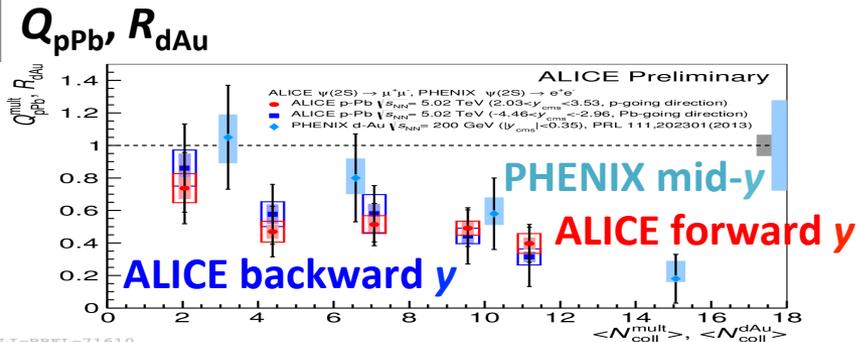
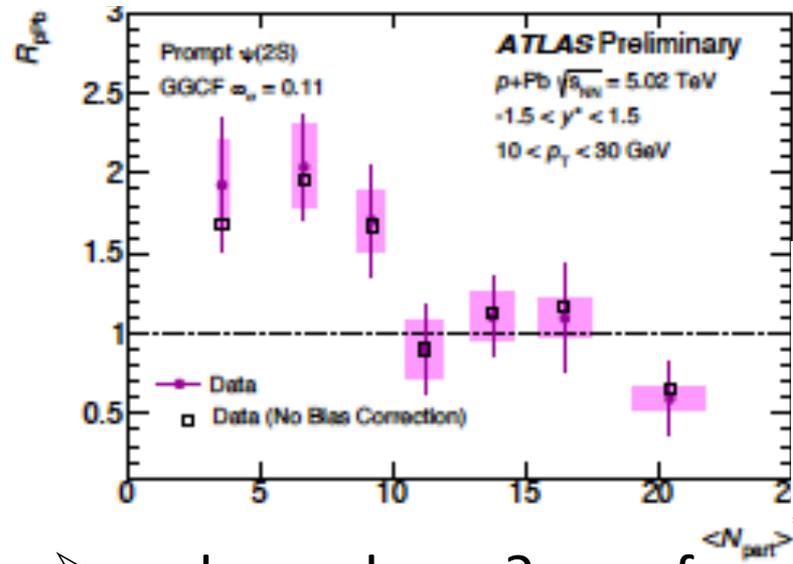
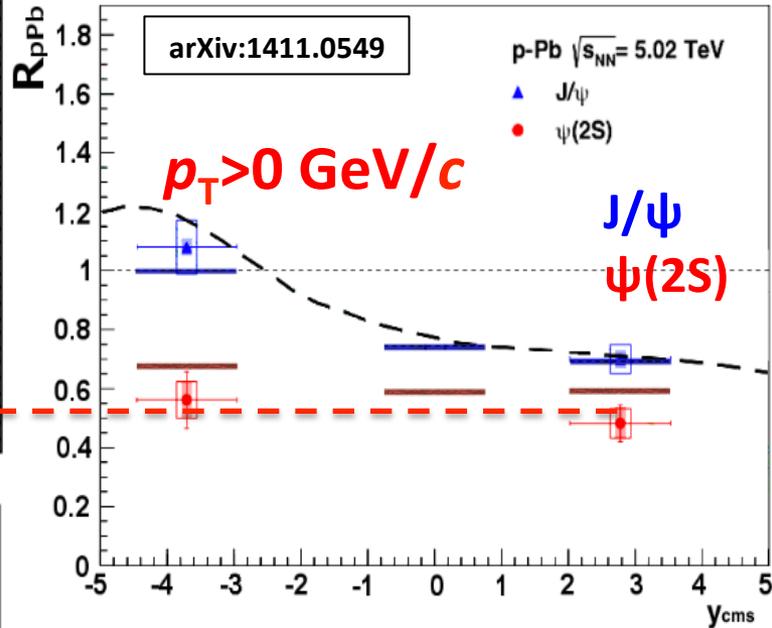
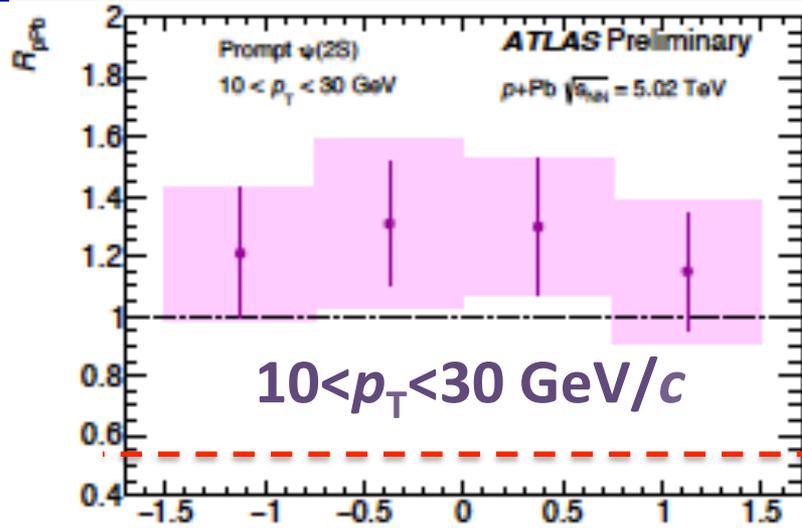
ALI-PUB-81954

LI-PREL-71610

- $[\psi(2S)/J/\psi]_{\text{pPb,dAu}}$  is suppressed compared to pp.
- $\psi(2S)$  suppression in p-Pb depends on centrality.
- Similar behaviour in ALICE p-Pb and PHENIX d-Au.
- ☐ Dependence on  $y$ ? on energy?
- ☐ Model with comover interactions + EPS09 agrees with ALICE and PHENIX. **Final state effect?**

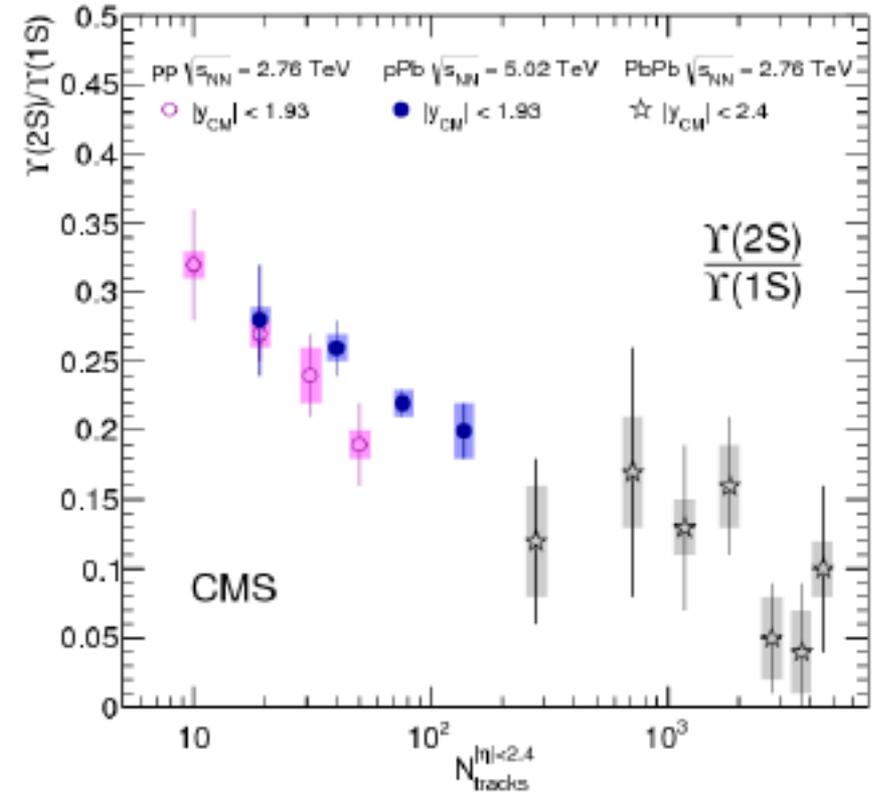
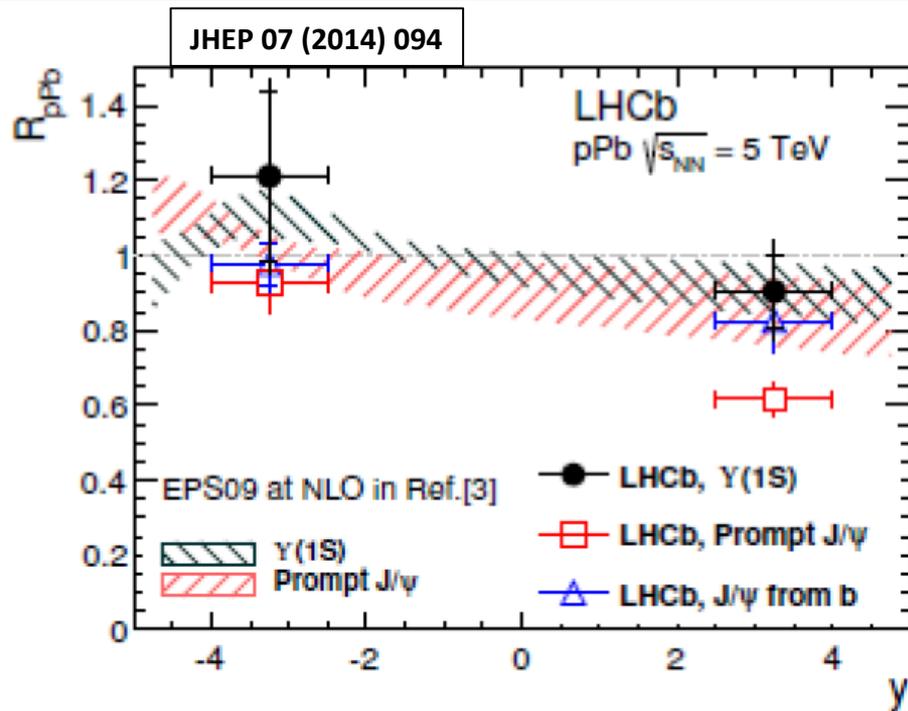


# $\psi(2S)$ vs $J/\psi$ in p-Pb



➤  $p_T$  dependence? or reference issue?

# Y(nS) in p-Pb



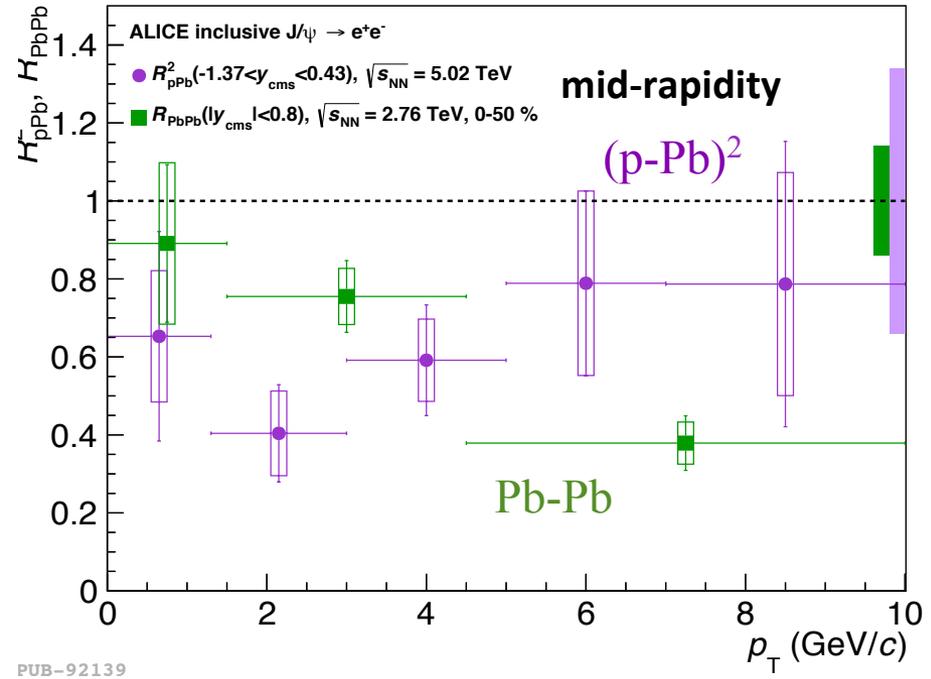
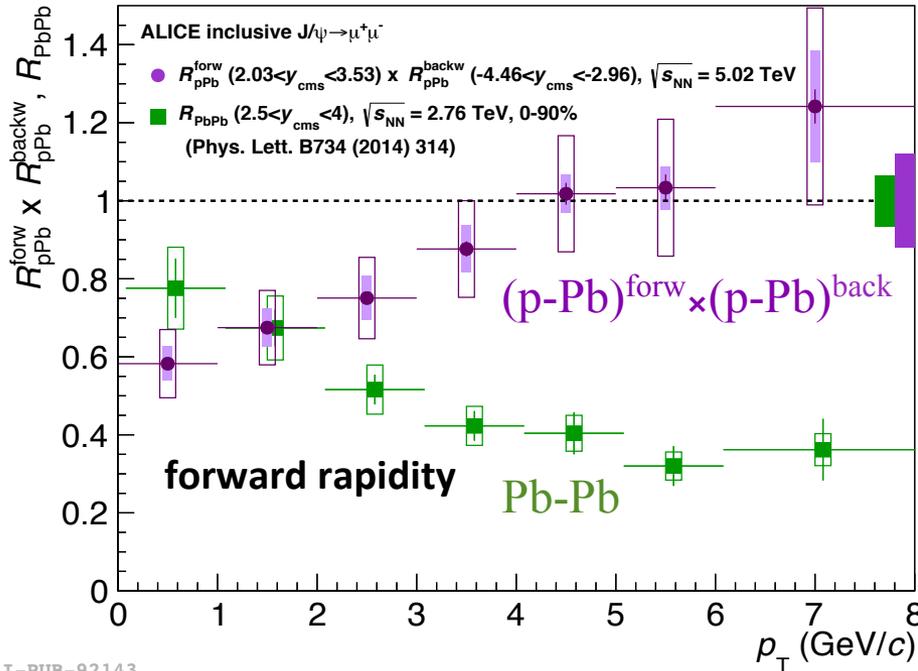
- Similar suppression of Y(1S) and (prompt) J/ψ.
- Y(2S)/Y(1S) decreases with increasing multiplicity.

# J/ψ: from p-Pb to Pb-Pb

JHEP 1506 (2015) 055

## Hypothesis

- ✓ Assume similar x in Pb for Pb-Pb@ $\sqrt{s_{NN}}=2.76$  TeV and in p-Pb@ $\sqrt{s_{NN}}=5.02$  TeV.
- ✓ Factorization of shadowing effects in p-Pb and Pb-Pb  $\Rightarrow R_{PbPb}^{Shad} = R_{pPb}(y \geq 0) \times R_{pPb}(y \leq 0)$ .



- High  $p_T$ : suppression in Pb-Pb due to hot nuclear matter effects (QGP).
  - Low  $p_T$ : less or same suppression with  $R_{Pb-Pb}$  compared to  $R_{PbPb}^{Shad}$ .
- Hint for (re)combination in Pb-Pb?

# Summary

- ❑ Quarkonium production in Pb-Pb is mainly dominated by the hot nuclear matter effects.
- ❑ In p-Pb: strong CNM effects in quarkonium production; depend on  $p_T$ ,  $y$  and centrality.
- ❑ The CNM effects are higher in the most central collisions, decreasing towards peripheral.
- ❑ Strong  $p_T$  broadening is found at forward  $y$ .
- ❑ No model is able to reproduce precisely all observables, though shadowing and coherent energy loss work well.
- ❑ The  $\psi(2S)$  is suppressed more than  $J/\psi$  in p-Pb compared to pp collisions.
- ❑ Comover model can describe different behaviour of the 2 charmonium states in p-Pb.
- ❑ Some inconsistencies between ALICE and ATLAS:  $p_T$  dependence or reference issue?
- ❑ Data taking of pp@5.02TeV are appreciated.

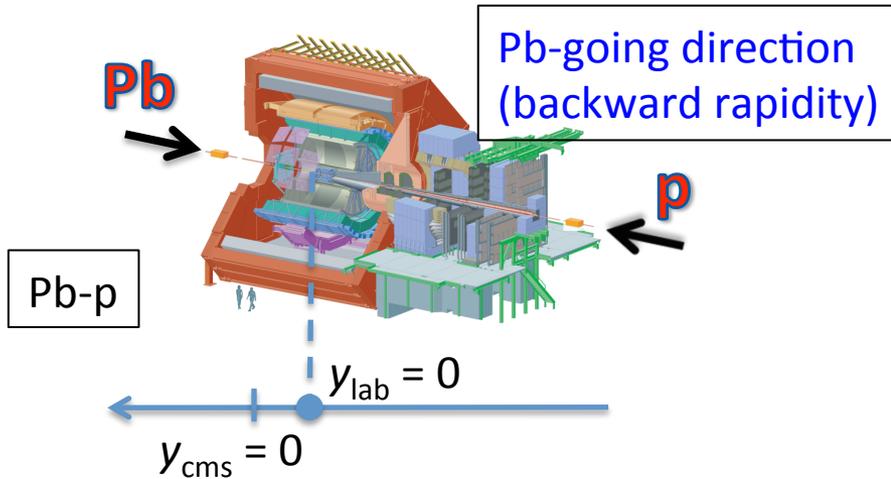
**Thank you!**

# Backup slides

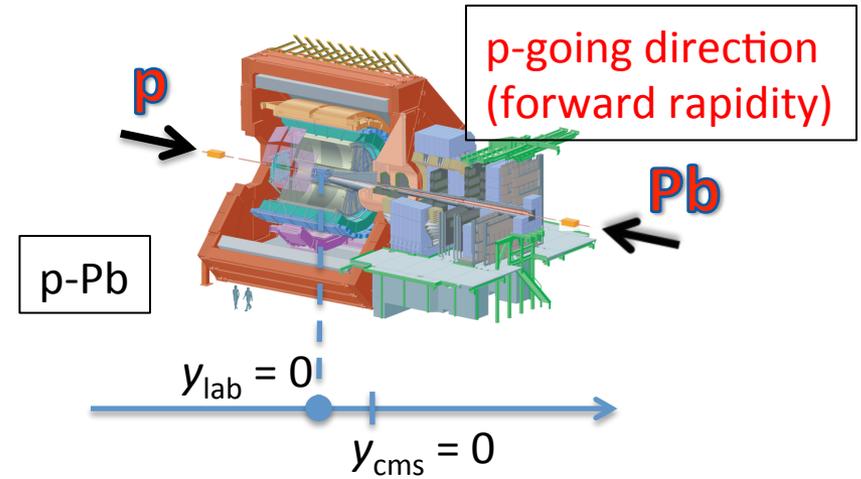
# p-Pb detector setup

## ❖ Shift in $y_{\text{cms}}$ in p-Pb collisions

LHC beam asymmetry ( $E_{\text{Pb}}=1.58 \cdot A \text{ TeV}$ ,  $E_{\text{p}}=4 \text{ TeV}$ )  $\Rightarrow |\Delta y|_{\text{cms}} = 0.5 \text{ Log}(Z_{\text{Pb}} A_{\text{p}} / Z_{\text{p}} A_{\text{Pb}}) = 0.465$



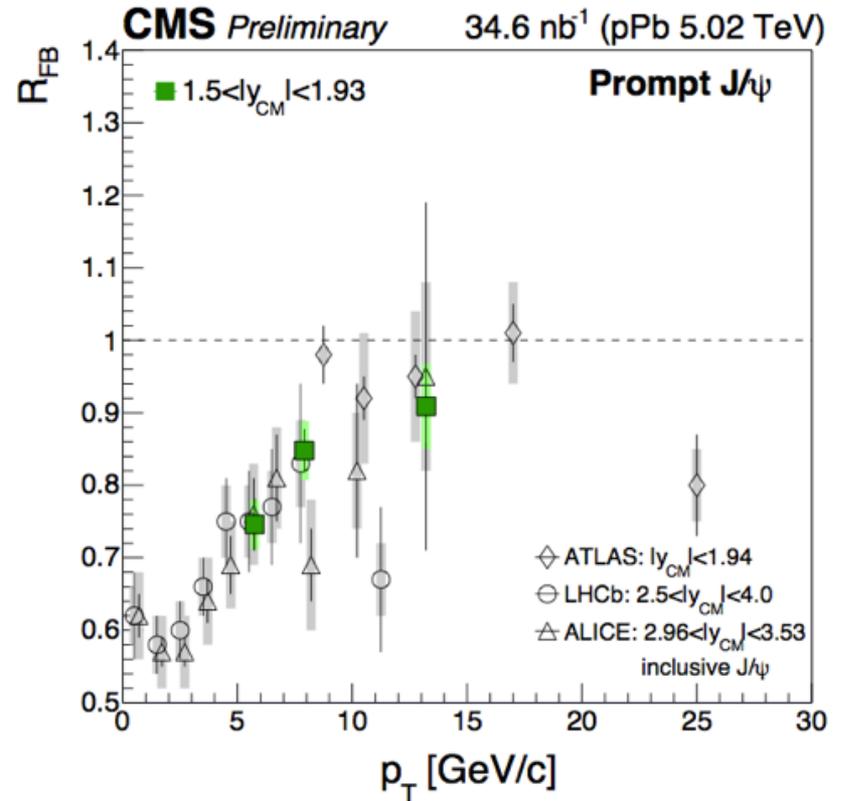
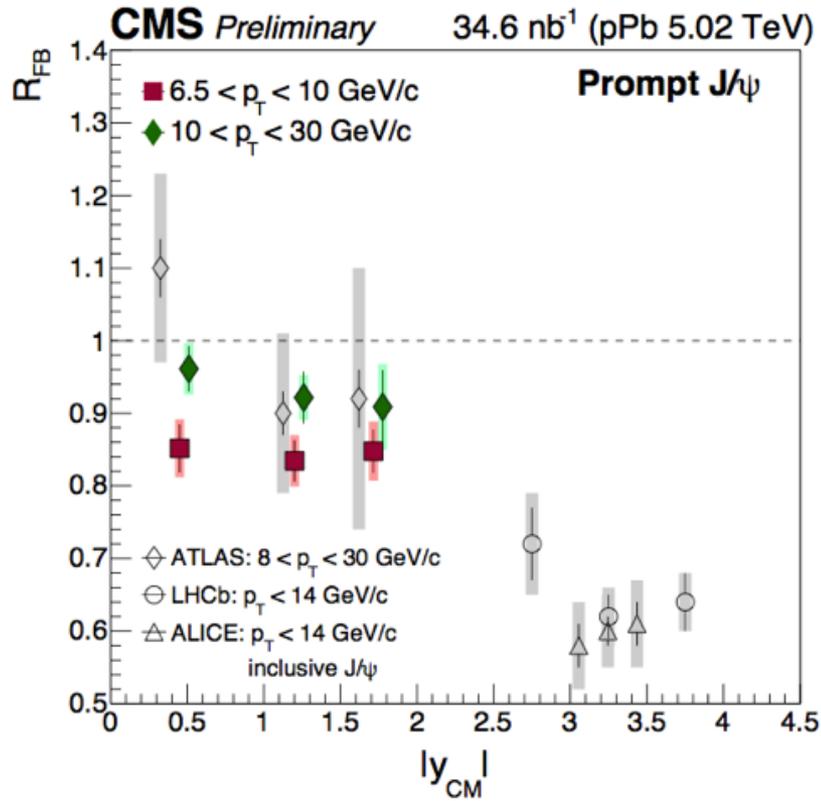
$J/\psi, \psi(2S), \Upsilon(nS) \rightarrow \mu^+\mu^-: -4.46 < y_{\text{cms}} < -2.96$



$J/\psi \rightarrow e^+e^-: -1.37 < y_{\text{cms}} < 0.43$

$J/\psi, \psi(2S), \Upsilon(nS) \rightarrow \mu^+\mu^-: 2.03 < y_{\text{cms}} < 3.53$

# Prompt $J/\psi$ in p-Pb



# Y(1S) vs Y(2S) cross-sections

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- Y(1S) cross-sections in p-Pb@5.02 TeV (ALICE):

$$-4.46 < y_{\text{cms}} < -2.96: \quad 5.57 \pm 0.72^{\text{stat.}} \pm 0.60^{\text{syst.}} \mu\text{b}$$

$$2.03 < y_{\text{cms}} < 3.53: \quad 8.45 \pm 0.94^{\text{stat.}} \pm 0.77^{\text{syst.}} \mu\text{b}$$

- Y(2S) cross-sections in p-Pb@5.02 TeV (ALICE):

$$-4.46 < y_{\text{cms}} < -2.96: \quad 1.85 \pm 0.61^{\text{stat.}} \pm 0.32^{\text{syst.}} \mu\text{b}$$

$$2.03 < y_{\text{cms}} < 3.53: \quad 2.97 \pm 0.82^{\text{stat.}} \pm 0.50^{\text{syst.}} \mu\text{b}$$

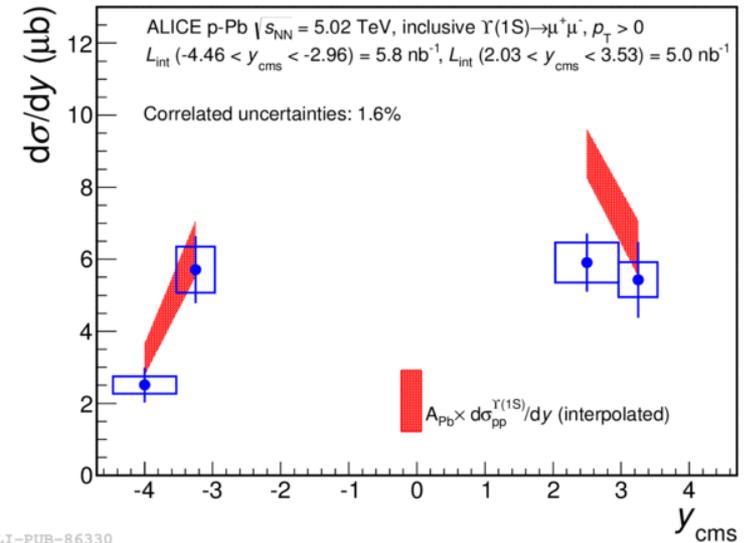
- ❖ [Y(2S)/Y(1S)] in p-Pb@5.02 TeV (ALICE):

$$-4.46 < y_{\text{cms}} < -2.96: \quad 0.26 \pm 0.09 \pm 0.04$$

$$2.03 < y_{\text{cms}} < 3.53: \quad 0.27 \pm 0.08 \pm 0.04$$

- ❖ [Y(2S)/Y(1S)] in pp@7 TeV (ALICE):

$$2.5 < y < 4.0: \quad 0.28 \pm 0.08$$



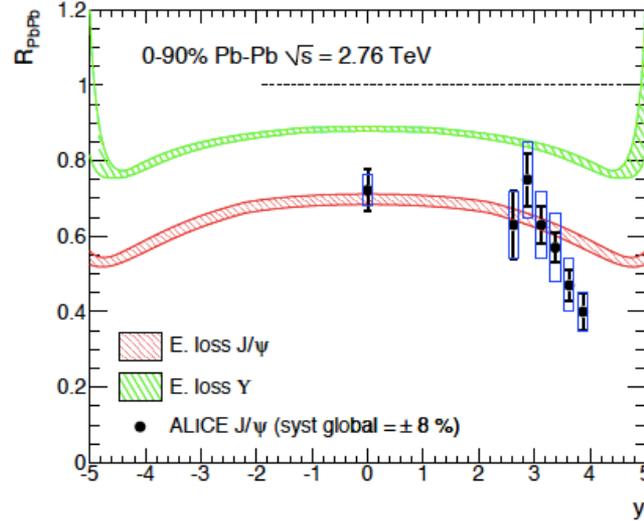
ALI-PUB-86330

➤ For ALICE no evidence of different CNM effects on Y(2S) compared to Y(1S).

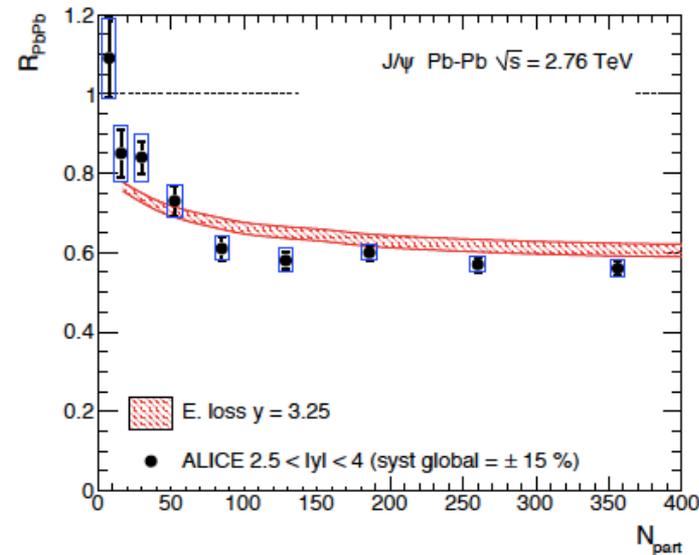
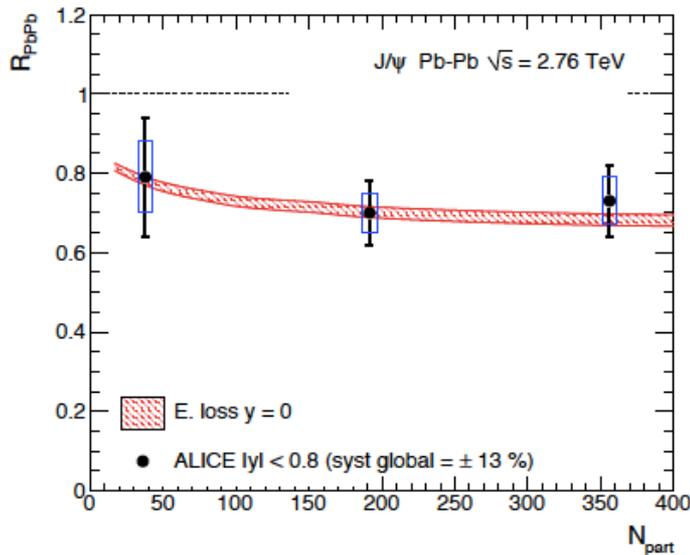
❑ CMS measured at mid-y  $[Y(2S)/Y(1S)]_{\text{pPb}}/[Y(2S)/Y(1S)]_{\text{pp}} = 0.83 \pm 0.05^{\text{stat.}} \pm 0.05^{\text{syst.}}$

# Coherent parton energy loss in Pb-Pb

F. Arleo et al.  
arXiv:1407.5054



✧ Factorization of the effects in p-Pb and Pb-Pb is also assumed.



➤ Good job of the model (if the regeneration does indeed compensate the Debye screening).