

# Neutron Monitors (NM) for solar modulation studies: systematic uncertainties and $\phi$ time series

## Workshop

*“Solar Energetic Particles (SEP), Solar Modulation and Space Radiation:  
New Opportunities in the AMS-02 Era”*

1. Introduction
2.  $\phi$  from CR data or ground-based detectors
3. Count rates: uncertainties on ingredients
4.  $\Delta(\text{ingredients}) \rightarrow \Delta N/N \rightarrow \Delta\phi/\phi$
5. Determination of  $\phi$  : normalization issues
6. Conclusion and perspectives

## Based on

Maurin *et al.*, *AdSpR* 55, 363 (2015)

Ghelfi *et al.* (in prep.)

 Enigmass  
The enigma of mass

 LPSC  
Grenoble

 ANR  
AGENCE NATIONALE DE LA RECHERCHE

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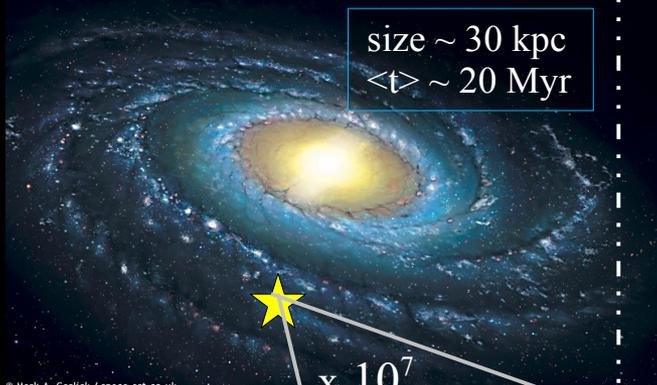
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Les deux infinis

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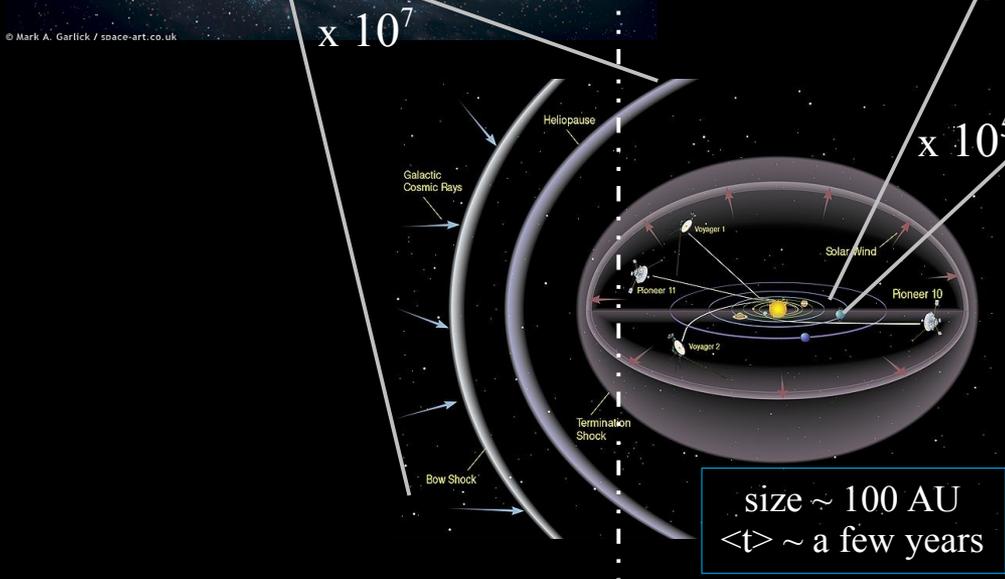
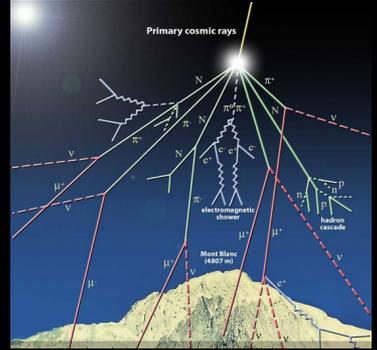
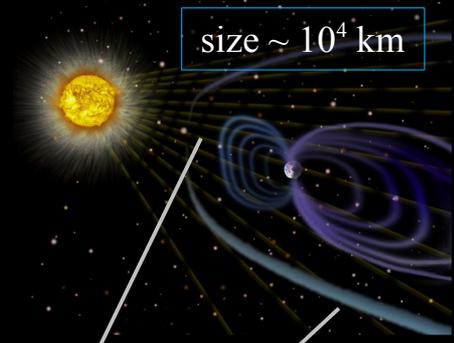
University of Hawaii  
10/20/2015

# 1. Introduction: CRs + modulation + Rc + atmosphere

1. Transport in the Galaxy  
 → Interstellar (IS) spectra



3. Earth magnetic shield  
 → Cut-off rigidity  $R_c$  (at Earth)



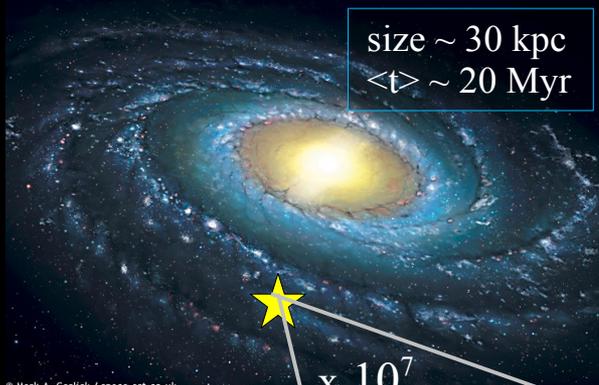
2. Transport in the Solar cavity  
 → modulate CRs ( $< 10 \text{ GeV/n}$ )

4. Atmosphere  
 → CR showers  
 [time-dependent]

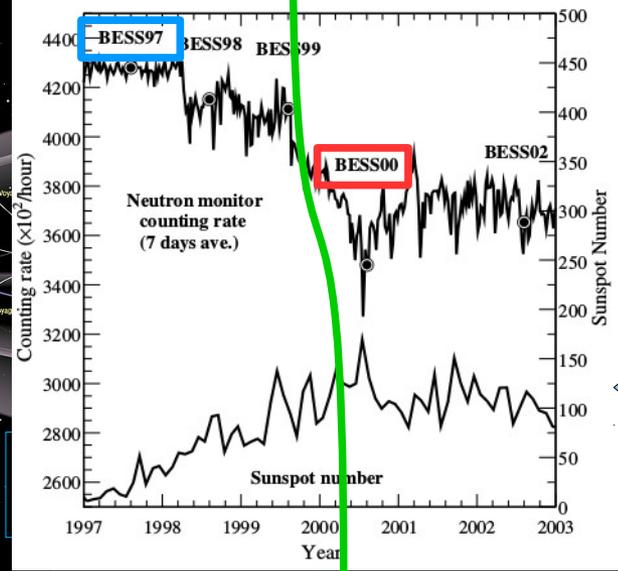
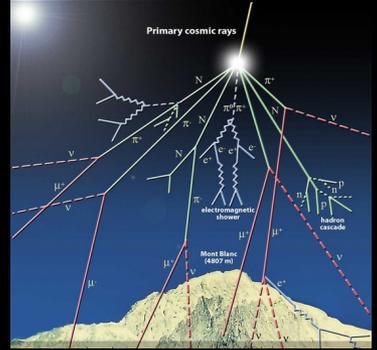
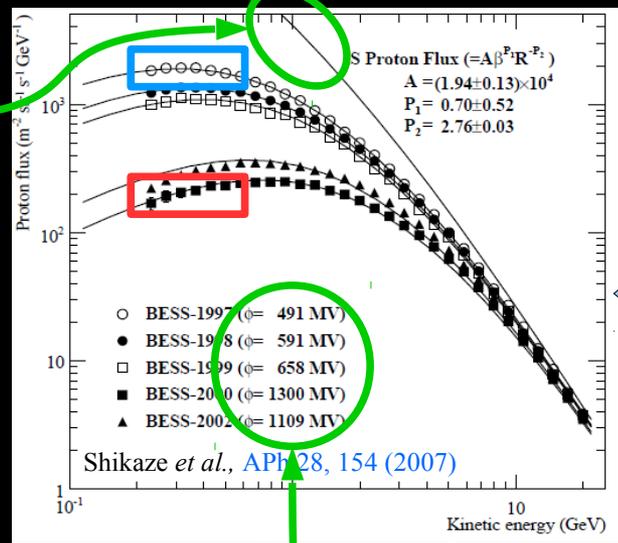
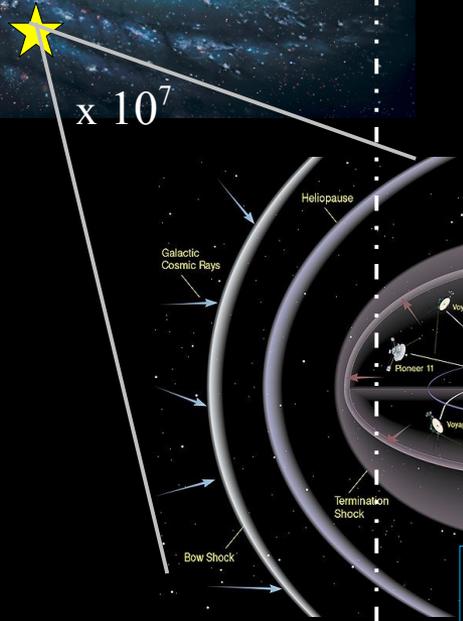
[time-independent]

# 1. Introduction: IS fluxes, TOA fluxes, and count rates

## 1. Transport in the Galaxy → Interstellar (IS) spectra



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## 2. Transport in the Solar cavity → modulate CRs (< 10 GeV/n)

## 4. Atmosphere → CR showers [time-dependent]

[time-independent]

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**Based on**

Maurin *et al.*, [AdSpR 55, 363 \(2015\)](#)

Ghelfi *et al.* (in prep.)

**Reference textbooks (NMs):** Dorman ([1974](#), [2004](#), [2009](#))

**Reference papers:**

- $\phi$  from NMs: Usoskin *et al.* ([1999](#), [2002](#), [2005](#), [2011](#))
- Solar activity times series: Usoskin, [Liv. Rev. Sol. Phys. \(2013\)](#)

# 2. Modulation level: from TOA data [see Ghelfi's poster]

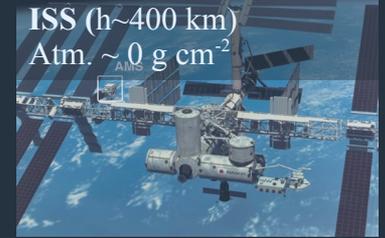
1. Transport in the Galaxy  
→ Interstellar (IS) spectra

$$J^{IS}(E^{IS})$$

CRDB (Maurin *et al.*, 2014)  
<https://lpsc.in2p3.fr/crdb/>

→  $\phi$  and  $J^{IS}$  simultaneous determination

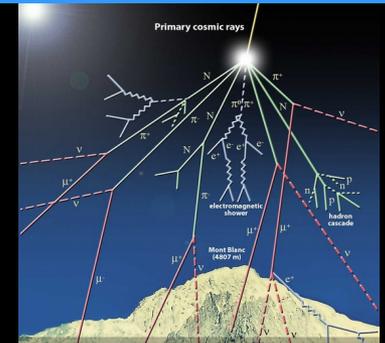
$$\chi^2 = \sum_{i_l} \sum_{N_j(i)} \sum_{E_k(i,j)} \frac{(J^{TOA}(J_j^{IS}, \phi_i, E_k) - \text{data}_{ijk})^2}{\sigma_{ijk}}$$



ISS (h~400 km)  
Atm. ~ 0 g cm<sup>-2</sup>



Balloon (h~40 km)  
Atm. ~ 5 g cm<sup>-2</sup>



Neutron monitor (h<2 km)  
Atm. ~ 600-1000 g cm<sup>-2</sup>



4. Atmosphere  
→ CR showers  
[time-dependent]

$$\frac{E^{TOA}}{A} = \frac{E^{IS}}{A} - \frac{|Z|}{A} \phi$$

$$J^{TOA}(E^{TOA}) = \left(\frac{p^{TOA}}{p^{IS}}\right)^2 \times J^{IS}(E^{IS}),$$

2. Transport in the Solar cavity  
→ Force-Field approximation (here)

[time-independent]

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# 2. Modulation level: from count rate detectors

1. Transport in the Galaxy  
 → Interstellar (IS) spectra

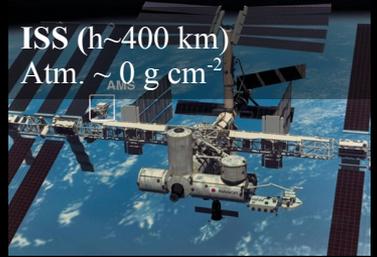
$$J^{IS}(E^{IS})$$

3. Earth magnetic shield  
 → Cut-off rigidity  $R_c$  (at Earth)

transmission function

$$N^D(\vec{r}, t) = \int_0^\infty \mathcal{T}(R, \vec{r}, t) \times \sum_{i=CRs} \mathcal{Y}_i^D(R, h) \frac{dJ_i^{TOA}}{dR}(R, t) dR$$

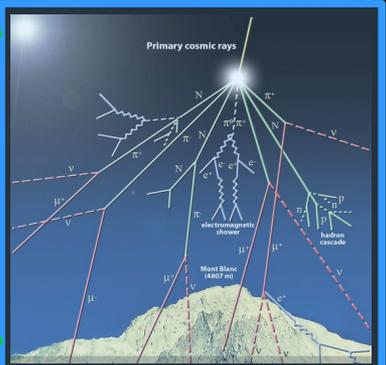
→  $\phi$  times series (~ mn time resolution)



ISS (h~400 km)  
 Atm. ~ 0 g cm<sup>-2</sup>



Balloon (h~40 km)  
 Atm. ~ 5 g cm<sup>-2</sup>



Neutron monitor (h<2 km)  
 Atm. ~ 600-1000 g cm<sup>-2</sup>



4. Atmosphere  
 → CR showers  
 [time-dependent]

2. Transport in the Solar cavity  
 → Force-Field approximation (here)

$$\frac{E^{TOA}}{A} = \frac{E^{IS}}{A} - \frac{|Z|}{A} \phi$$

$$J^{TOA}(E^{TOA}) = \left(\frac{p^{TOA}}{p^{IS}}\right)^2 \times J^{IS}(E^{IS}),$$

NMDB  
<http://www.nmdb.eu/>

[time-independent]

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Ghelfi *et al.* (in prep.)

# 3. Count rates: ingredients and parameters

1. Transport in the Galaxy  
→ Interstellar (IS) spectra

$$J^{IS}(E^{IS})$$

3. Earth magnetic shield  
→ Cut-off rigidity  $R_c$  (at Earth)

transmission function

$$N^D(\vec{r}, t) = \int_0^\infty \mathcal{T}(R, \vec{r}, t) \times \sum_{i=CRs} \mathcal{Y}_i^D(R, h) \frac{dJ_i^{TOA}}{dR}(R, t) dR$$

→  $\phi$  times series ( $\sim$  mn time resolution)



$$\frac{E^{TOA}}{A} = \frac{E^{IS}}{A} - \frac{|Z|}{A} \phi$$

$$J^{TOA}(E^{TOA}) = \left( \frac{p^{TOA}}{p^{IS}} \right)^2 \times J^{IS}(E^{IS}),$$

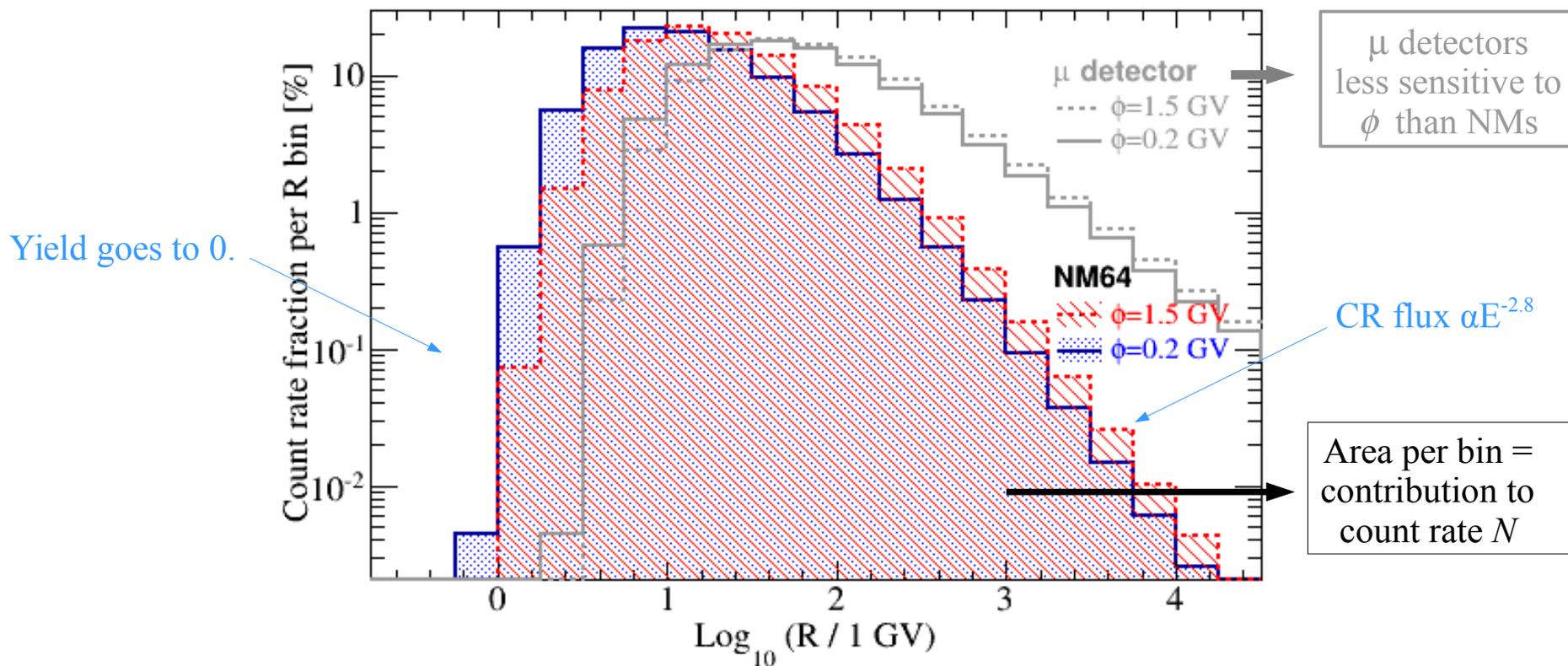
→ Key ingredients  
(source of uncertainties)

Free parameter  $\phi(t)$

2. Transport in the Solar cavity  
→ Force-Field approximation (here)

### 3. Count rates: contributions per log(R) bin

$$N^D(\vec{r}, t) = \int_0^\infty \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^D(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$



→ NM64: sensitive to CRs in range 5 GV – 500 GV

→ Muons detectors: sensitive to CRs in range 10 GV – 1 TeV

### 3. Count rates: contribution from CRs heavier than He

$$N^D(\vec{r}, t) = \int_0^\infty \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^D(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$

sum over all CR species  $i$

**Ansatz:** in sum, only consider  ${}^1\text{H}$  and  $(1+s_{Z>2}){}^4\text{He}$

- ${}^1\text{H}$  and  ${}^4\text{He}$  most abundant species in CRs
- ${}^1\text{H}$  and  ${}^4\text{He}$  modulated differently:  $(A/Z)_p = 1$  whereas  $(A/Z)_{4\text{He}} = 2$
- Heavier than He assimilated to  ${}^4\text{He}$ :  $(A/Z)_{Z>2} \sim 2$

**Calculate  $s_{Z>2}$**

- Extract IS flux for all species: fit on TOA fluxes (using [CRDB](#))
- Use scaling  $\mathcal{Y}_A^D(R, h) = \frac{A}{4} \times \mathcal{Y}_{4\text{He}}^D(R, h)$  from [Mishev & Velinov \(2001\)](#)
- Calculate  $s_{Z>2}$  and its uncertainties (+ check how good the scaling is)

$\rightarrow s_{Z>2} = 0.611 (+0.016, -0.009)$   
compared to previously used value 0.428

*N.B.: the relative weight of He to H in  $N(r,t)$  increases  $\rightarrow$  matters for  $\phi$  determination*

### 3. Count rates: Transfer function $T$

$$N^{\mathcal{D}}(\vec{r}, t) = \int_0^{\infty} \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^{\mathcal{D}}(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$

**Definition:** transmission function in the geomagnetic field, which depends on the detector location (and can vary with time)

*N.B.: in the most general case,  $T$  and  $Y$  are entangled (complex structure of  $B$  plus dependence of  $T$  and  $Y$  on the primary particle incident angle)*

**Ansatz:**

- Average  $Y$  over a few incident angles
- Use vertical cut-off rigidity  $\rightarrow$  step function  $T = \theta(R - R_c)$   
[Cooke *et al.*, 1991] for various cut-off definitions  
[Smart and Shea (2003, 2008a,b, 2009)] for  $R_c(\mathbf{r}, t)$  calculations

$\rightarrow$  these approximations lead to significant uncertainties  
(see further)

### 3. Count rates: Yield function – calculation

$$N^{\mathcal{D}}(\vec{r}, t) = \int_0^{\infty} \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^{\mathcal{D}}(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$

**Definition:** the yield function at altitude  $h$ , is the detector response (in count  $\text{m}^2 \text{sr}$ ) to a unit intensity of primary CR species  $i$  at rigidity  $R$

#### Calculate Y

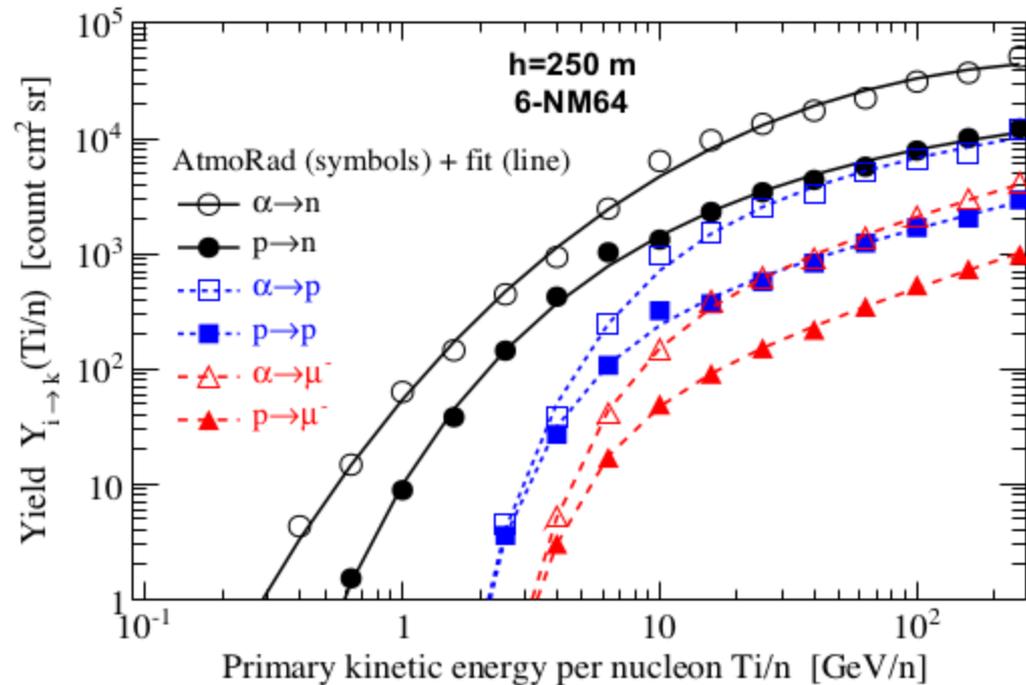
- *MC-based yield functions*
  - Atmosphere response (e.g., [Clem & Dorman 2000](#)) + [Mishev et al. \(2013\)](#)
  - Detector response
- *NM 'count rate'-based formulae* (e.g., [Caballero-Lopez & Moraal, 2012](#))
  - use  $R_c$  to reconstruct Yield from count rates of NM at different latitudes

# 3. Count rates: Yield function (MC-based) for NM64

Total yield: 
$$\mathcal{Y}_i^D(T_i, h) = \sum_{k=n,p,\mu,\dots} \mathcal{Y}_{ik}^D(T_i, h)$$

Partial yields: 
$$\mathcal{Y}_{ik}^D(T_i, h) = S_T \sum_{l=1}^3 \Omega(\theta_l, \theta_{l+1}) \times \int_0^\infty \mathcal{E}_k^D(T_k) \times \varphi_{ik}^l(h, T_i \rightarrow T_k) dT_k$$

[efficiency]                      [fluence]

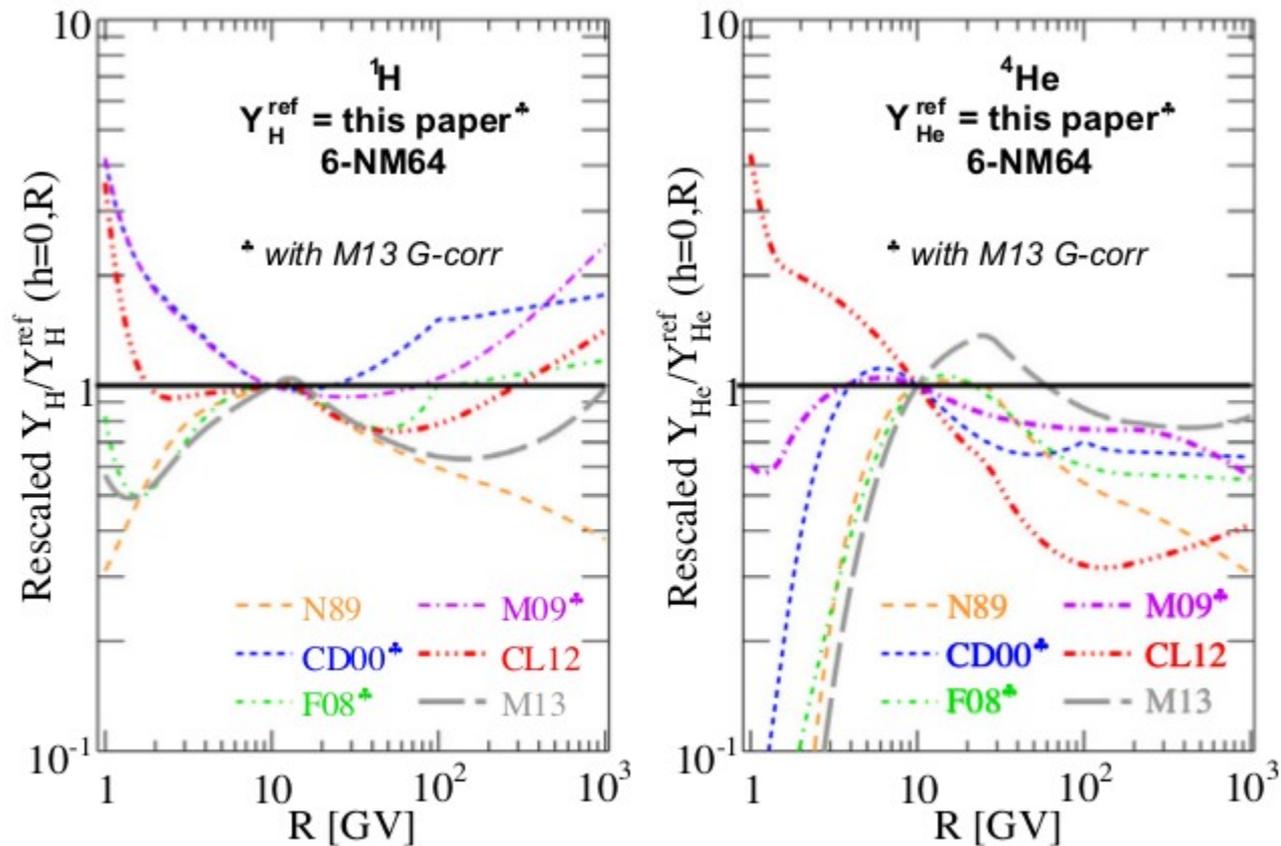


$\phi$ [GV]	$N^{6\text{-NM64}}$ [s <sup>-1</sup> ]	$N_k/N$ [%]			
		$n$	$p$	$\mu^+$	$\mu^-$
0.4	91	87.2	7.9	0.2	4.7
1.5	57	87.4	8.0	0.2	4.4

→ in NM count rates, mostly  $n$ , but also  $p$  and  $\mu^-$

# 3. Count rates: Yield function uncertainties for NM64

$$N^{\mathcal{D}}(\vec{r}, t) = \int_0^{\infty} \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^{\mathcal{D}}(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$



→ Large differences between various teams (R-dependent)

→ Simulations should be extended above 100 GV (important for count rates calculation)

[N.B.: for MC-based yields, better results obtained with Mishev et al. (2013) correction]

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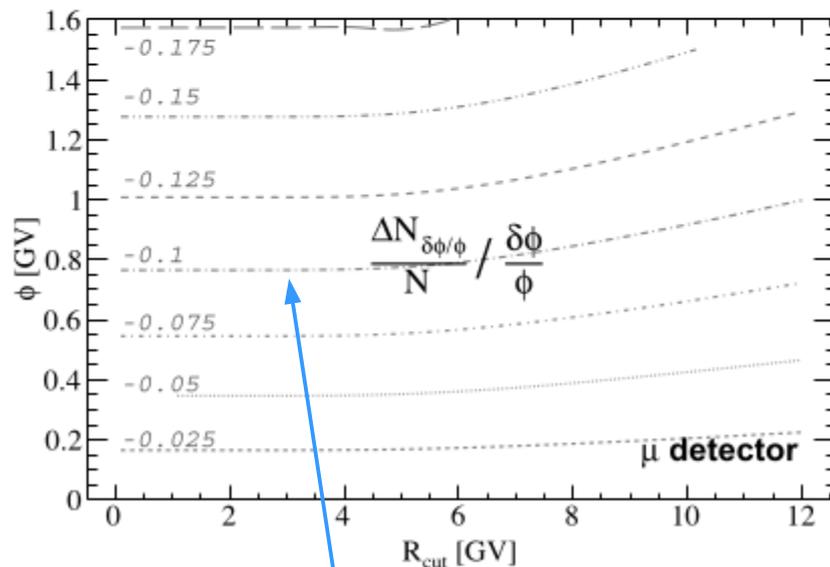
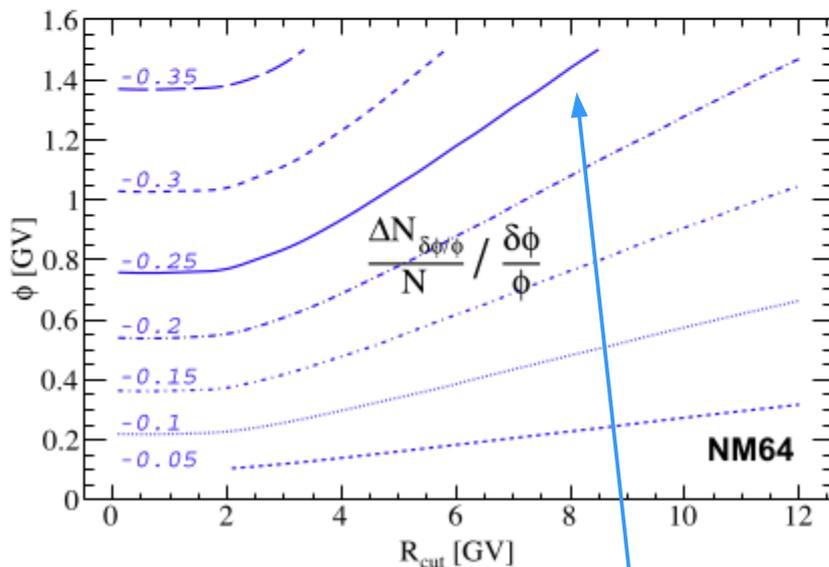
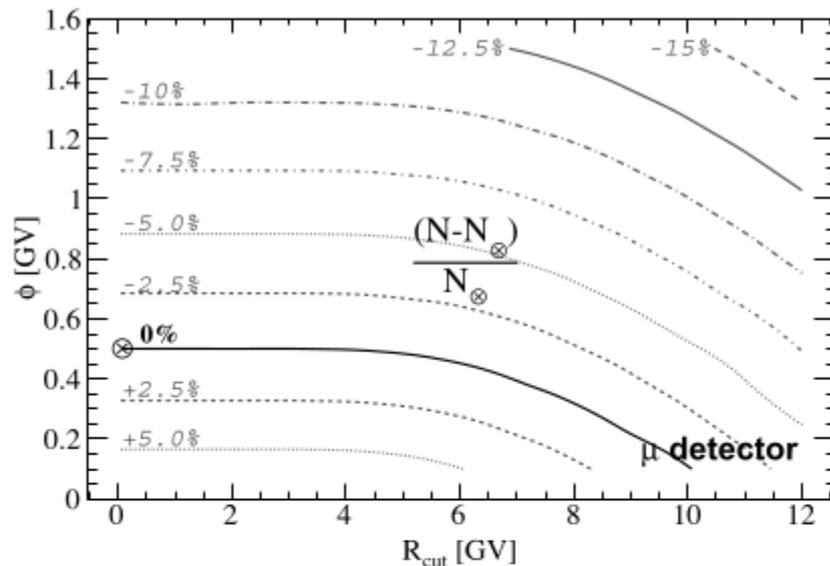
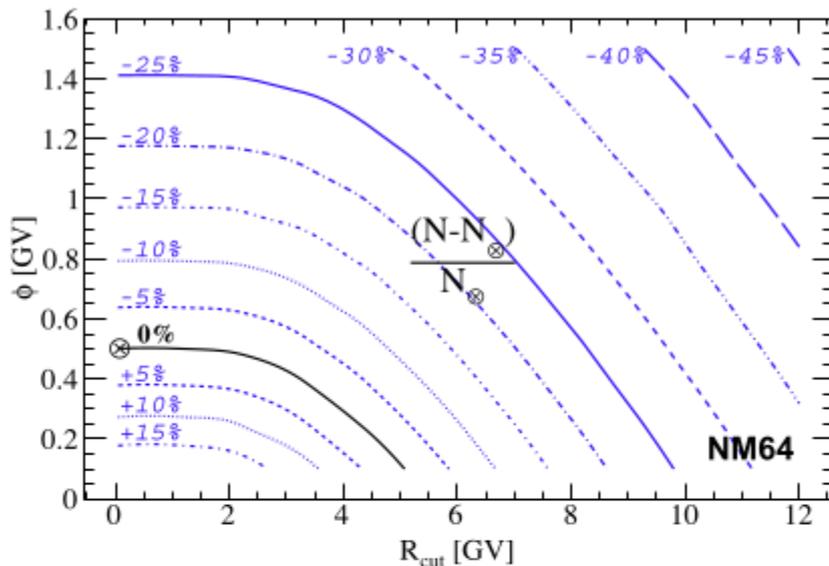
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# 4. Abacus for $\Delta N/N$ and $\Delta\phi/\phi$

Top: count rate variation when varying  $R_c$  or  $\phi$



Bottom:  $\Delta\phi/\phi = -4 \Delta N/N$  (NM) or  $\Delta\phi/\phi = -10 \Delta N/N$  ( $\mu$ )

# 4. Uncertainties for $\Delta N/N$ and $\Delta\phi/\phi$

$$\Delta N/N \sim 1\% \rightarrow \Delta\phi \sim 40 \text{ MV}$$

Ingredient	Effect	$\frac{\Delta N}{N}$		$\Delta\phi^*$ [MV]		Comment
		NM	$\mu$	NM	$\mu$	
Solar modulation	$\phi \in [0.2, 1.5] \text{ GV}$	$[+15, -25]\%$	$[+5, -10]\%$	–	–	w.r.t. $\phi = 0.5 \text{ GV}$
Cut-off rigidity	$R_c \in [0, 10] \text{ GV}$	$[+10, -20]\%$	$[0, -5]\%$	–	–	w.r.t. $R_c = 5 \text{ GV}$
TOA flux	p and He CR data	$\pm 2\%$	$\pm 2\%$	$\pm 66$	$\pm 140$	$(t, R_c, \phi)$ -independent
	IS flux dispersion <sup>a</sup>	$\pm 6\%$	$\pm 8\%$	$\pm 200$	$\pm 570$	(without AMS-02)
	Heavy species	$\pm 0.6\%$	$\pm 0.6\%$	$\pm 20$	$\pm 40$	Global norm. factor <sup>◇</sup>
Yield function	Dispersion	$\lesssim \pm 4\%$	$< 0.2\%$	$\lesssim 120$	$\lesssim 14$	$(R_c, \phi)$ dependent
	Sigmoid( $R_c, x = +\frac{\sigma}{0.1}$ )	$-2x\%$	$-0.5x\%$	$+66x$	$+35x$	For $R_c \gtrsim 5 \text{ GV}$
Transfer function	$H(R_c + \Delta R_c) : x = \frac{(\Delta R_c/R_c)}{0.05}$	$-2x\%$	$-x\%$	$+66x$	$+71x$	For $R_c \gtrsim 5 \text{ GV}$
	$- R_c(t) : \frac{\Delta R}{R} \lesssim +0.2\%/yr$	$-0.4\%/yr$	$-0.1\%/yr$	$+13/yr$	$+7/yr$	Depends on location
	$- R_c^{\text{eff}} \rightarrow R_c^{\text{app}} : +3\%$	$-1.2\%$	$-0.3\%$	$+40$	$+21$	Depends on $R_c$
Time-dep. effects <sup>†</sup>	Pressure	$\pm 0.2\%$	$\pm 0.2\%$	$\pm 6$	$\pm 14$	After correction
	Temperature	$\pm 0.5\%$	$\pm 4\%$	$\pm 15$	$\pm 290^{\ddagger}$	Not corrected
	Vapour water	$\pm 0.3\%$	$\pm 0.1\%$	$\pm 10$	$\pm 8$	Not corrected
	Snow coverage ( $T = 1 \text{ yr}$ )	$-7\%$	–	$+230$	–	Not corrected
NM detector effects	Temperature	$+0.05\%/^{\circ}\text{C}$	–	$-1.5/^{\circ}\text{C}$	–	$(t, R_c, \phi)$ -independent
	$n\text{NM6}$ vs $m\text{NM64}$	few %	–	$\sim 100$	–	↓
	Surroundings (hut)	few %	–	$\sim 100$	–	Global norm. factor <sup>◇</sup>

- IS flux uncertainty: global shift in  $\phi$  time series (will improve with AMS-02 data)
- Yield uncertainty: main source of uncertainty for  $\phi$
- Rigidity cut-off uncertainty: also important source of uncertainty for  $\phi$

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Ghelfi *et al.* (in prep.)

# 5. Determination of $\phi$

## Questions in the AMS-02 context

- **From AMS-02 point of view**

- Can we predict AMS-02 flux time-variability from NMs?

- Consistency between AMS-02 and NM-derived  $\phi$  ?

- and for TOA CR data interpretation**

- “Experiment-dependent”  $\phi$  value (based on  $\neq$ IS fluxes,  $\neq$ modulation models)

- Need “consistently-derived” values for all data (CR or dark matter interpretation)

- **More globally, in the context of Solar modulation**

- Access to IS spectra and fit of Solar modulation parameters

- Study of charge dependence for Solar modulation

- **For the space weather community (over short and long periods)**

- AMS-02 protons is another monitor of the Solar activity

- Cross-calibration between TOA CR data and NM/ $\mu$  detectors?

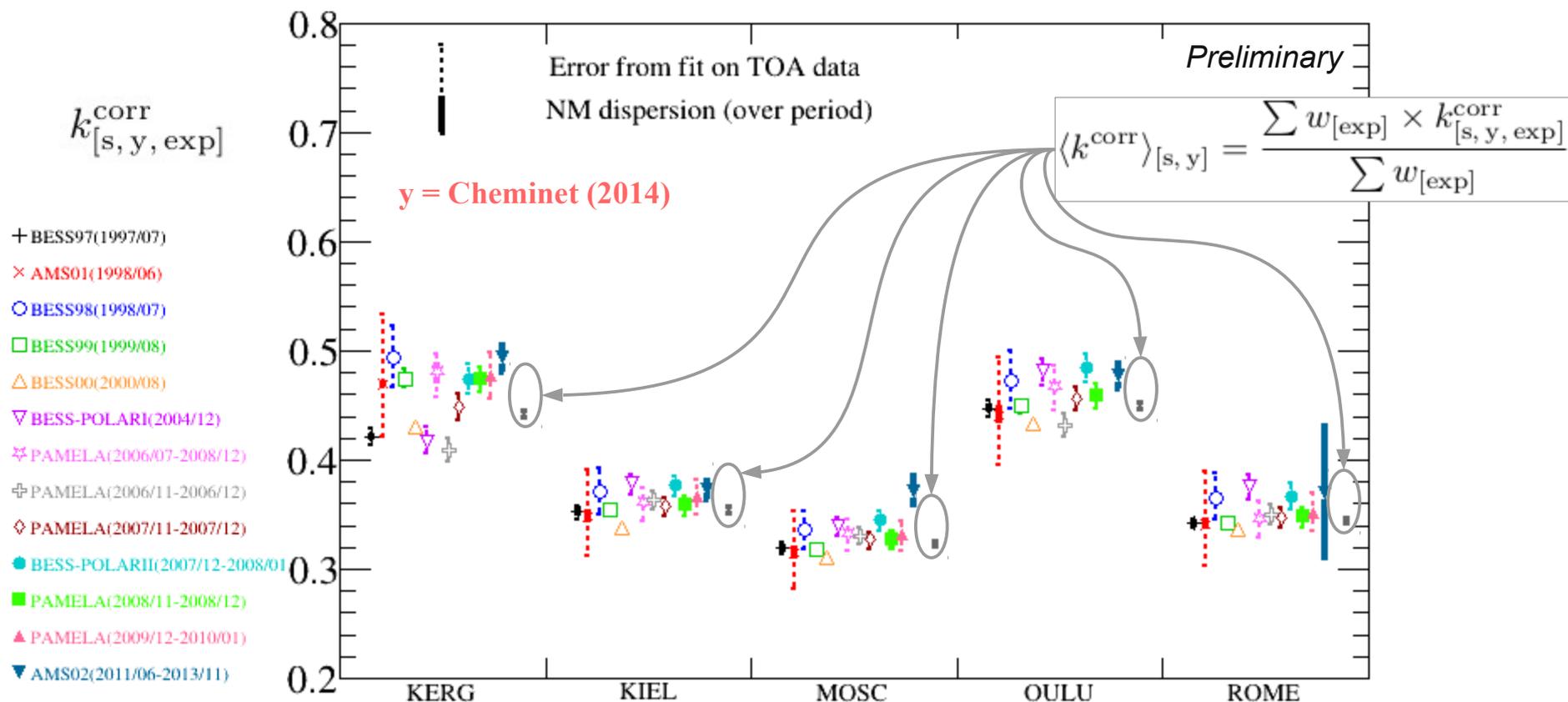
# 5. Determination of $\phi$ : normalisation issues

$$N^{\text{calc}}(R_c, h) = k^{\text{corr}} \int_{R_c}^{\infty} \sum_{i=\text{CRs}} \mathcal{Y}_i(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$



$$k_{[s, y, \text{exp}]}^{\text{corr}} = \frac{N_{[s, y, \text{exp}]}^{\text{calc}}}{\langle N_{[s]}^{\text{measured}} \rangle_{\text{exp. date}}}$$

$k$  = individual NM's efficiency factor (accounts for the local environment)



→ Correction factor well constrained from selected TOA CR data (see Ghelfi's poster)

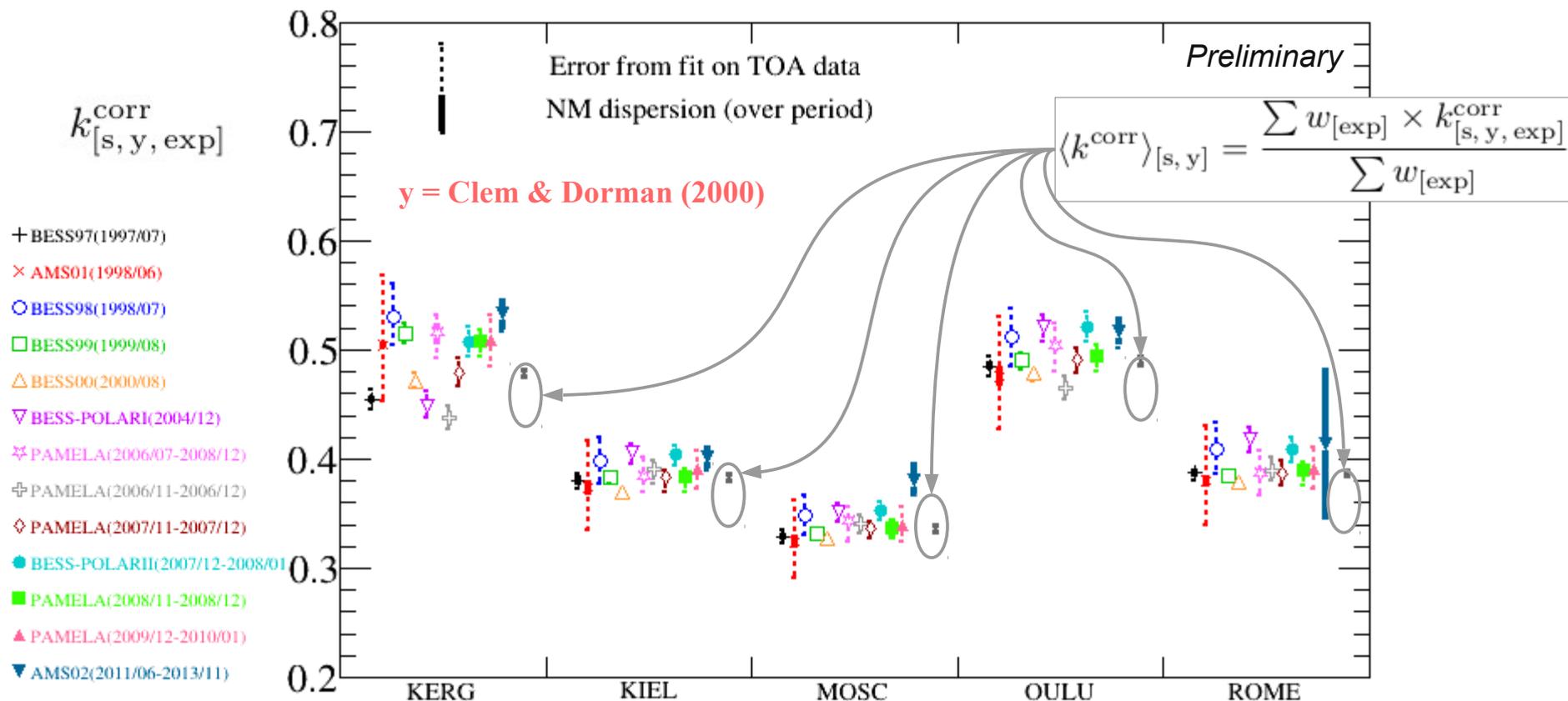
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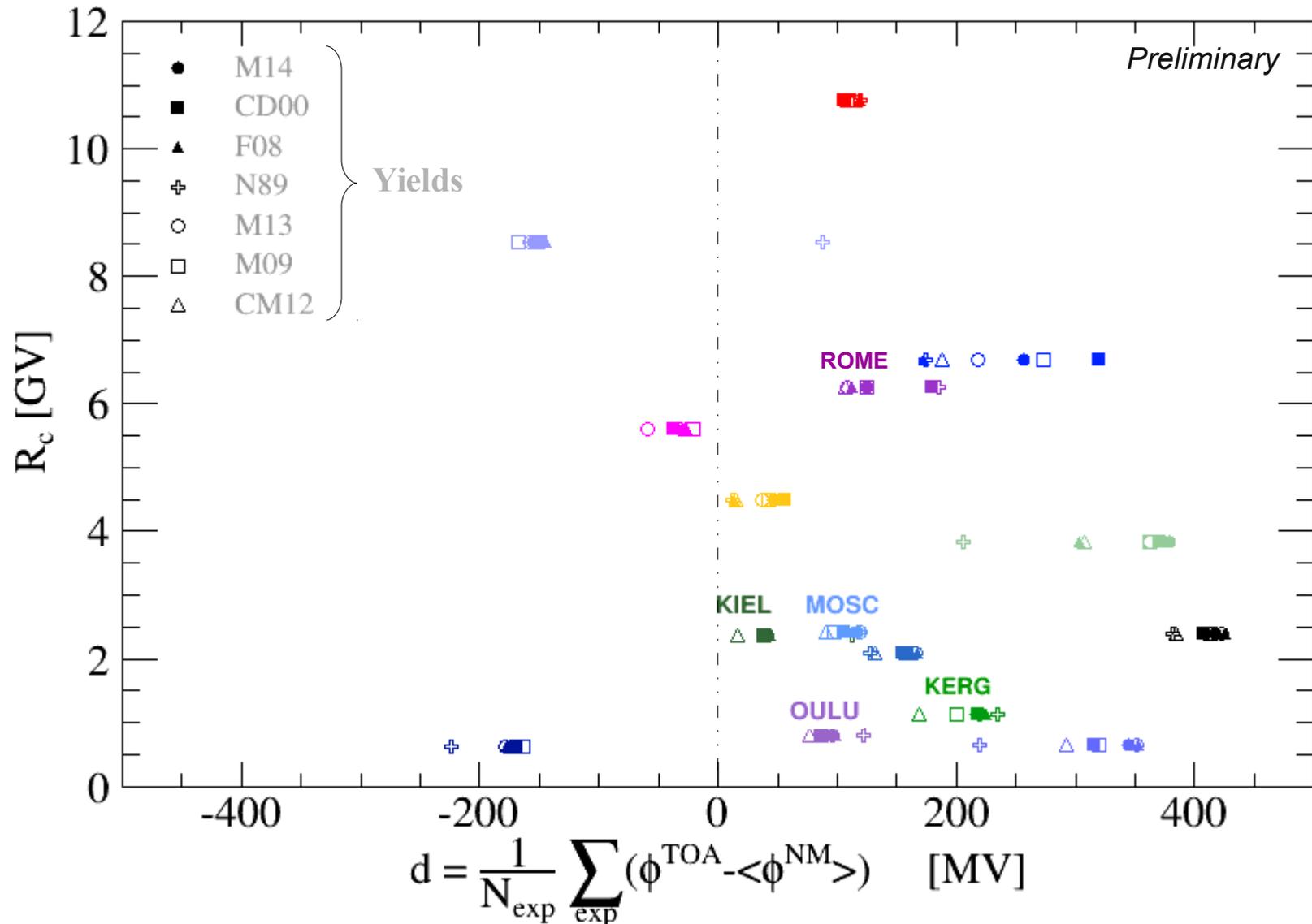
$$k_{[s, y, \text{exp}]}^{\text{corr}} = \frac{N_{[s, y, \text{exp}]}^{\text{calc}}}{\langle N_{[s]}^{\text{measured}} \rangle_{\text{exp. date}}}$$

$k$  = individual NM's efficiency factor (accounts for the local environment)



→  $k$  differs from Usoskin *et al.* (2011) values (e.g., 0.82 for KIEL)  
[this difference is not understood at this moment]

# 5. Determination of $\phi$ : $\phi(\text{TOA}) - \phi(\text{NM})$ over exp. dates

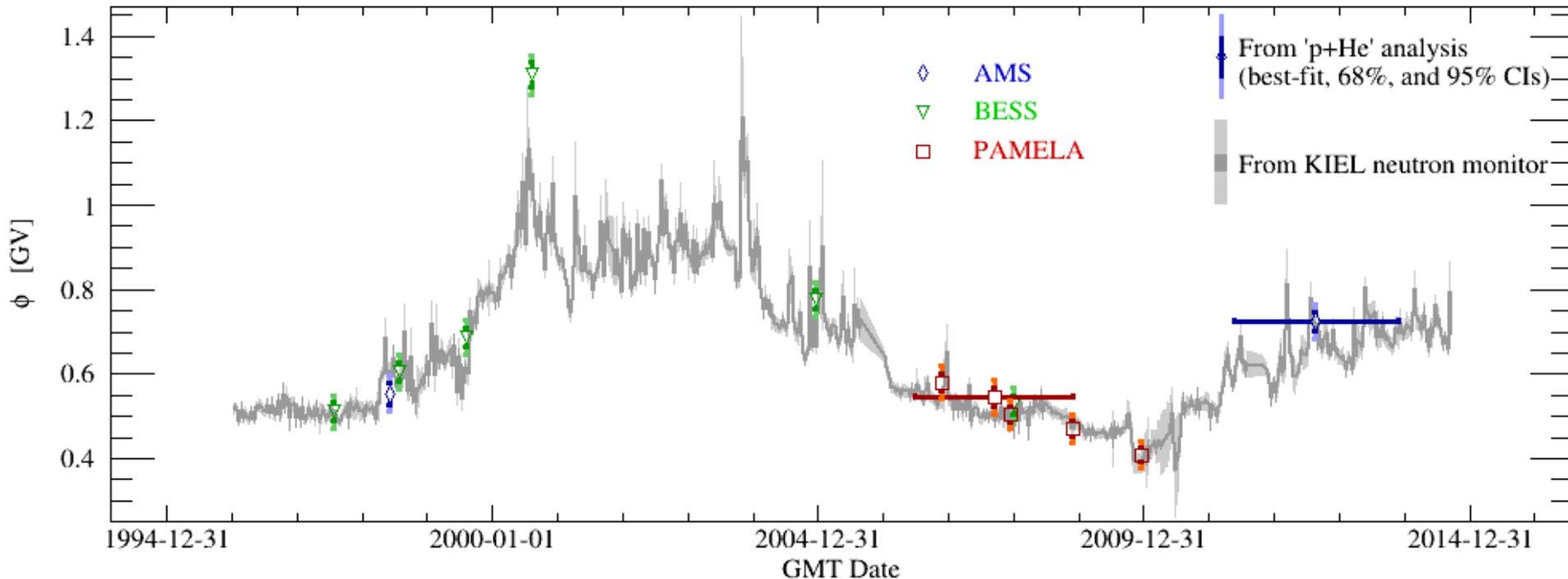


→ Still large differences between different stations: related to  $R_c$ ?

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# 6. Conclusion and perspectives

Can we provide accurate  $\phi$  time series from NM data?



## To do list:

- Find best combination of IS flux/Yield function/NM station matching  $\phi^{\text{TOA}}$
- Combine/compare from reconstruction from  $\mu$  detectors
- Investigate more realistic modulation models and see impact on  $\phi^{\text{TOA}}$  and  $\phi^{\text{NM}}$ 
  - **Time series from NM:** Ghelfi *et al.* (in prep.)

→ Need to improve “Yield functions + Rc + NM relative normalisation”

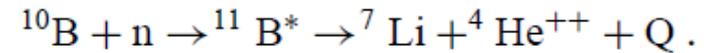
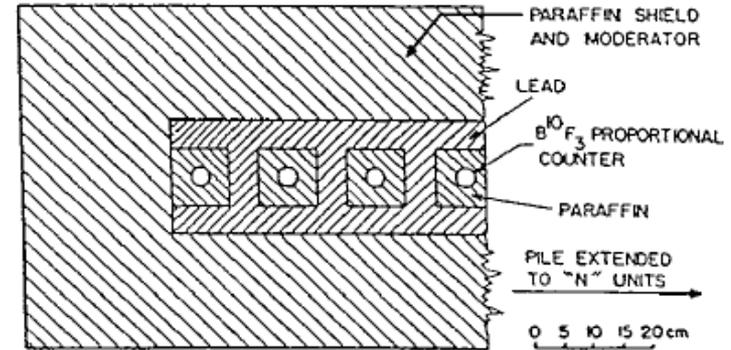
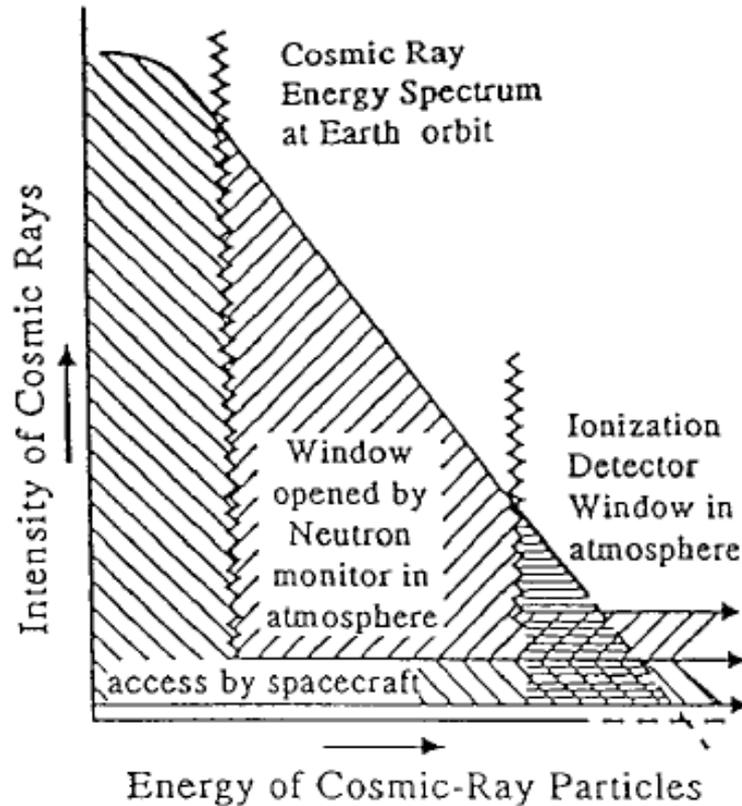
→ Monthly/weekly/daily AMS-02 data mandatory to go further...

[N.B.: Web interface to provide  $\phi$  for any “past” time interval (CRDB website asap)



# The invention of NMs in the 50's (John A. Simpson)

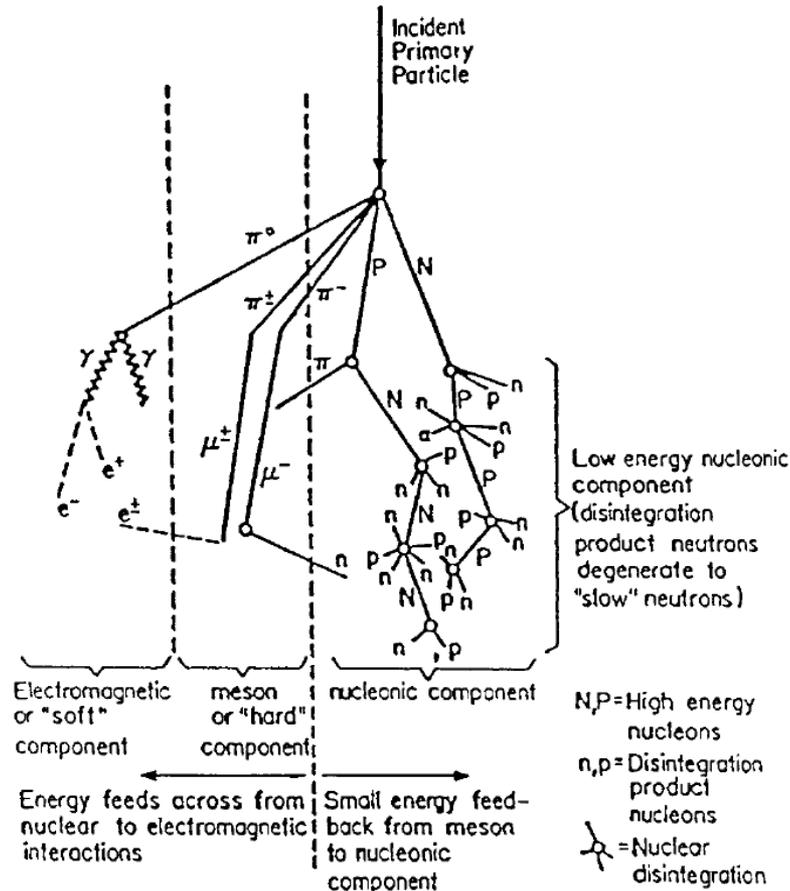
Simpson, Space Sci. Rev. **93**, 11 (2000)



→ NMs come “cheap”, insensitive to low energy SCRs,  
Good time resolution, active since the 50's

# CR showers and ground-based detectors

Simpson, Space Sci. Rev. **93**, 11 (2000)



- **Count rates:** neutron monitors (NMs),  $\mu$  detectors (AUGER scaler data)
- **Spectra:** neutron spectrometers (BSS)

# Fluence rate from PLANETOCOSMICS (and setup)

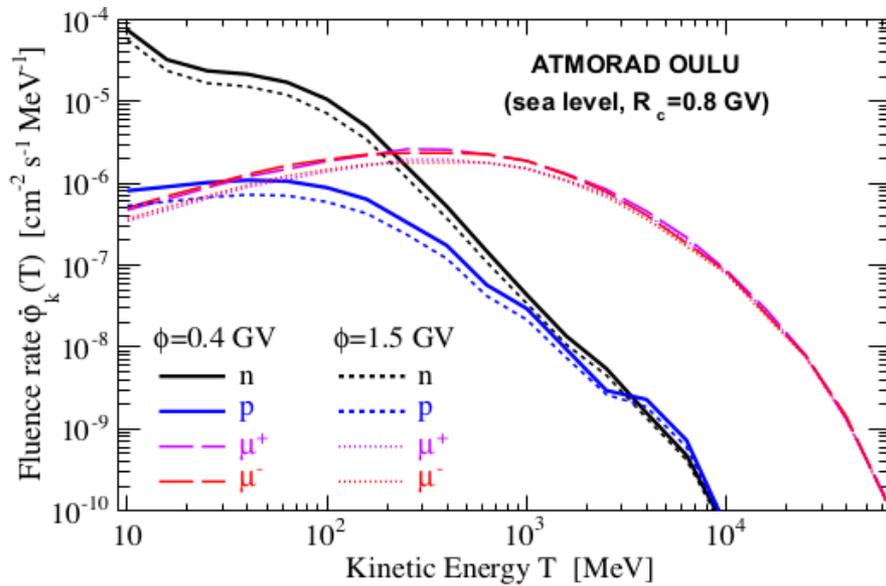
(fluence rate)

[fluence]

$$\dot{\phi}_k(T_k, \vec{r}, t) = S_T \times \sum_{l=1}^3 \Omega^{(\theta_l, \theta_{l+1})} \sum_i^{\text{CRs}} \sum_{T_i^{\text{cut-off}}}^{T_i^{\text{max}}} \Delta T_i \times J_i^{\text{TOA}}(T_i, \phi(t)) \times \phi_{ik}^l(h, T_i \rightarrow T_k)$$

GEANT4 + atmosphere model  
(PLANETOCOSMIC)

Monte Carlo simulation setup



Parameter	Qty	Bins	Range	Unit
Primary $i$	H	18	[0.1 – 251.2]	GeV/n
	He			
Secondary $k$	n	70	[ $10^{-3}$ – $10^{11}$ ]	eV
	p, $\mu^\pm$	25	[ $10^6$ – $10^{11}$ ]	
Incidence	$[\theta_l, \theta_{l+1}]$	3	$\left\{0, \frac{\pi}{6}, \frac{\pi}{3}, \frac{\pi}{2}\right\}$	rad
Altitude	h	36	[0 – 30]	km asl

Adrien Cheminet, PhD thesis (2013)

→ fluence rate = secondary particule spectrum  
content at any depth given some TOA fluxes

# Heavier than He fraction: IS flux determination

$$\frac{dJ^{\text{IS}}}{dR} = C_0 \times \beta^{C_1} \times R^{-C_2}$$

CR	$C_0$ ( $\text{m}^2 \text{ s sr GV}^{-1}$ )	$C_1$ -	$C_2$ -	$f_j = \frac{J_j}{\sum_i J_i}$ (%)
H <sup>†</sup>	23350 ± 184	2.10 ± 0.10	-2.839 ± 0.003	83.7
<sup>2</sup> H	838.5 ± 29.5	3.62 ± 0.08	-2.950 ± 0.060	2.21
<sup>3</sup> He	512.7 ± 2.0	6.70 ± 0.01	-3.045 ± 0.003	1.08
He <sup>‡</sup>	3657. ± 38.5	1.77 ± 0.04	-2.782 ± 0.003	14.6
Li	18.86 ± 0.86	4.58 ± 0.07	-3.200 ± 0.400	0.027
Be	22.09 ± 0.18	6.57 ± 0.08	-2.948 ± 0.003	0.054
B	72.77 ± 0.32	5.77 ± 0.01	-3.086 ± 0.002	0.132
C	116.5 ± 0.50	4.26 ± 0.01	-2.791 ± 0.002	0.438
N	45.7 ± 0.45	5.19 ± 0.02	-2.971 ± 0.004	0.112
O	95.5 ± 0.35	3.87 ± 0.01	-2.733 ± 0.002	0.413
F	35.73 ± 0.03	5.63 ± 0.02	-2.979 ± 0.005	0.084
Ne	16.75 ± 0.10	4.29 ± 0.02	-2.779 ± 0.003	0.065
Na	4.945 ± 0.034	5.02 ± 0.02	-2.922 ± 0.003	0.013
Mg	20.25 ± 0.10	4.03 ± 0.02	-2.755 ± 0.003	0.083
Al	4.165 ± 0.029	4.65 ± 0.02	-2.812 ± 0.004	0.015
Si	13.5 ± 0.10	3.86 ± 0.02	-2.681 ± 0.003	0.066
P	1.084 ± 0.014	5.99 ± 0.04	-2.938 ± 0.007	0.003
S	3.445 ± 0.025	4.87 ± 0.02	-2.785 ± 0.004	0.013
Cl	1.428 ± 0.015	6.65 ± 0.03	-3.052 ± 0.007	0.003
Ar	2.64 ± 0.04	6.24 ± 0.04	-3.075 ± 0.007	0.005
K	2.192 ± 0.005	6.37 ± 0.06	-3.110 ± 0.010	0.004
Ca	3.70 ± 0.03	5.24 ± 0.05	-2.991 ± 0.005	0.009
Sc	1.106 ± 0.019	5.68 ± 0.04	-3.120 ± 0.008	0.002
Ti	3.126 ± 0.032	4.96 ± 0.03	-3.062 ± 0.005	0.006
V	1.357 ± 0.014	4.82 ± 0.03	-2.995 ± 0.006	0.003
Cr	2.271 ± 0.019	4.51 ± 0.03	-2.919 ± 0.005	0.006
Mn	1.132 ± 0.014	4.11 ± 0.03	-2.776 ± 0.005	0.004
Fe	8.032 ± 0.046	3.37 ± 0.03	-2.600 ± 0.001	0.047
Co	0.0055 ± 0.0037	3.54 ± 0.03	-2.610 ± 0.010	< 10 <sup>-3</sup>
Ni	8.405 ± 0.019	4.50 ± 0.10	-2.600 ± 0.020	0.002

## Primary CR species

→ p, He most abundant species  
→ flatter spectrum for heavier species (destruction in propagation)

## Secondary CR species

→ steeper slope (relative fraction decreases with R w.r.t. primary species)  
→ <sup>2</sup>H and <sup>3</sup>He not same A/Z as other heavy species (A/Z~2)  
→ small fraction, but plenty of them...

<sup>†</sup> H = <sup>1</sup>H+<sup>2</sup>H, whereas He = <sup>3</sup>He+<sup>4</sup>He.

# Heavier than He fraction: IS flux determination

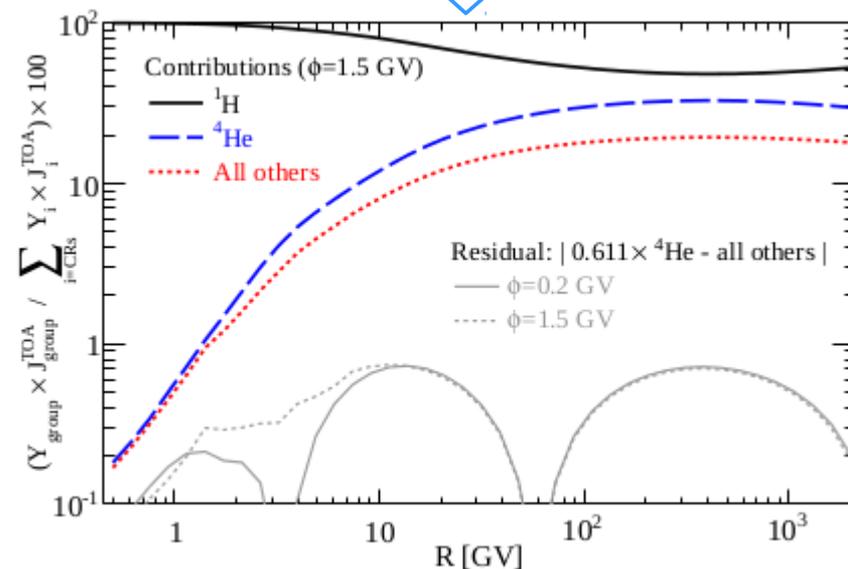
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Small fraction in CRs, but yield  $\mu$  A

$$\langle f_j \rangle_{\mathcal{Y}} = \frac{\mathcal{Y}_j(R) \times J_j^{\text{TOA}}(R)}{\sum_{i=\text{CRs}} \mathcal{Y}_i(R) \times J_i^{\text{TOA}}(R)} \approx \langle f_j \rangle_A$$

→ secondary and heavy species contribute significantly to count rates!

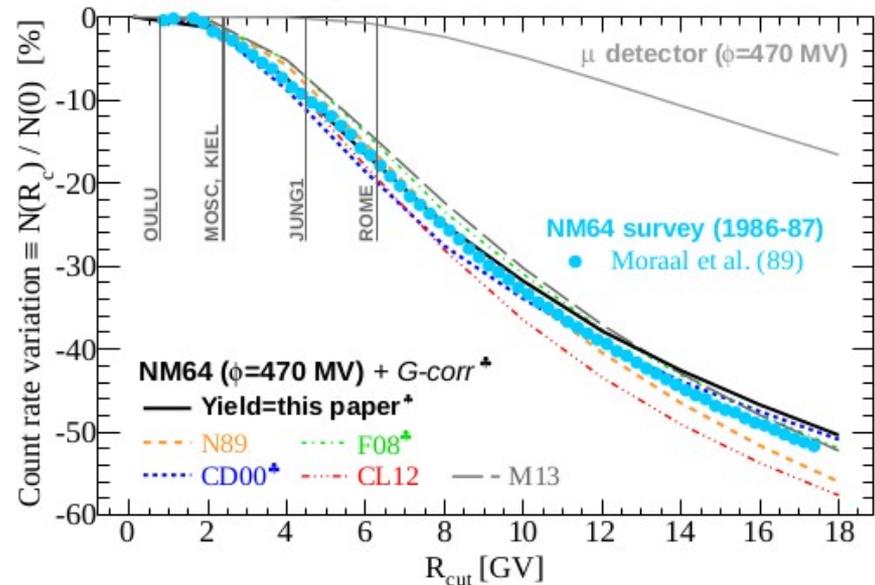
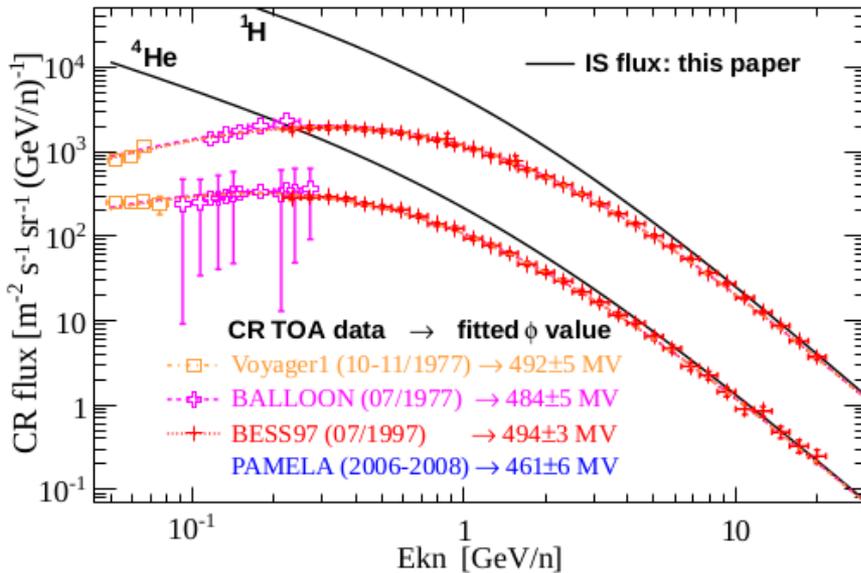


$$s_{Z>2} \equiv \frac{J_{\text{CRs}} - ({}^1\text{H}, {}^4\text{He})}{J_{{}^4\text{He}}} = 0.611^{+0.016}_{-0.009}$$

(0.428 in previous studies)

# Validation of yield functions on NM latitude surveys

1. Latitude surveys (vs  $R_c$ ) every 11 yr at solar minimum
2. Use CR data at same epoch to find  $\phi$
3. Check surveys with calculated count rates



$\rightarrow$ Up to 10 GV, our yield gives best fit  
 $\rightarrow$ Other yield functions more or less OK as well (dispersion)

# Relaxing assumption on transfer function

$$N^{\mathcal{D}}(\vec{r}, t) = \int_0^{\infty} \mathcal{T}(R, \vec{r}, t) \times \sum_{i=\text{CRs}} \mathcal{Y}_i^{\mathcal{D}}(R, h) \frac{dJ_i^{\text{TOA}}}{dR}(R, t) dR,$$

- $R_c$  varies in times: 10% in 50 yrs
- Transfer resembles a sigmoid (rather than step)

$$\mathcal{T}(R) = \frac{1}{2} \left[ 1 + \operatorname{erf} \left( \frac{R - R_c}{\sqrt{2} \Delta R} \right) \right] \quad \sigma \equiv \frac{\Delta R}{R_c}$$

- Apparent cut-off rigidity rather than effective vertical cut-off rigidity

$$R_c^{\text{eff}} \approx R_{\text{upper}} - \sum_{i=D.}^{R_{\text{upper}}} \Delta R_i^{\text{allowed}}$$

$$R_c^{\text{app}}(R_c) \equiv \frac{\int_0^{2\pi} d\phi \int_0^{\pi/2} R_c(\theta, \phi) \times \mathcal{Y}(R_c, \theta, \phi) d\theta}{\int_0^{2\pi} d\phi \int_0^{\pi/2} \mathcal{Y}(R_c, \theta, \phi) d\theta}.$$

- Variation depends on  $t$ ,  $R_c$ , position
- Maximal variation < 10% on count rates

