

Time dependence in $B \rightarrow V\ell\ell$

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Rare B-decays in 2015 experiment and theory
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New observables for $b \rightarrow sll$

Need for new $b \rightarrow sll$ observables

- cross-check hadronic and/or NP contributions
- try different incoming and outgoing states
- more information on $B \rightarrow Vll$?

- Angular analysis of $B \rightarrow Vll$ provides interferences between transversity amplitudes, but redundancy in the information
- Add another phase/amplitude to interfere and lift the redundancy ?
- Similar to CP-violation in B -decays: interference between decay and mixing adds a lot of information compared to decay alone

Time-dependent analysis of $B \rightarrow Vll$
where V decays into a CP-eigenstate

SDG and J. Virto, JHEP 1504 (2015) 045 [1502.05509]

Decays of interest

Need V to decay into CP-eigenstate

- Not possible for flavour specific decays $B_d \rightarrow K^{*0}(\rightarrow K^- \pi^+) \ell^+ \ell^-$
- Accessible via flavour non-specific decays

Three main examples in the following

$$B_d \rightarrow K^*(\rightarrow K_S \pi^0) \ell^+ \ell^-$$

$$B_s \rightarrow \phi(\rightarrow K_S K_L) \ell^+ \ell^-$$

$$B_s \rightarrow \phi(\rightarrow K^+ K^-) \ell^+ \ell^-$$

Last one already studied at LHCb (time integrated)

JHEP **1307**, 084 (2013)

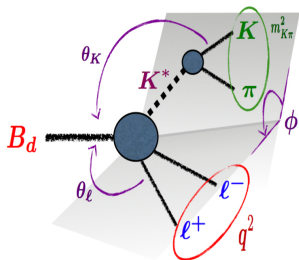
Kinematics

For untagged flavour-non-specific decays

- no possibility of distinguishing between B and \bar{B} decays
- need for consistent kinematic conventions
- the angles cannot be defined with respect to information on the flavour of the initial B (contrary to flavour-specific decays)

$$\frac{d\Gamma[B \rightarrow V(\rightarrow M_1 M_2)\ell^+\ell^-]}{ds d\cos\theta_\ell d\cos\theta_M d\phi} = \sum_i J_i(s) f_i(\theta_\ell, \theta_M, \phi)$$

$$\frac{d\Gamma[\bar{B} \rightarrow \bar{V}(\rightarrow \bar{M}_1 \bar{M}_2)\ell^+\ell^-]}{ds d\cos\theta_\ell d\cos\theta_M d\phi} = \sum_i \zeta_i \bar{J}_i(s) f_i(\theta_\ell, \theta_M, \phi)$$



- $f_i(\theta_\ell, \theta_M, \phi)$ are kinematical functions
- J interf. of A_X and A_Y , with $X, Y \in \{L0, R0, L||, R||, L\perp, R\perp, t, S\}$
- \bar{J} with $\bar{A}_X = A_X(\bar{B} \rightarrow \bar{M}_1 \bar{M}_2 \ell\ell) = A_X|_{\phi_{wk} \rightarrow -\phi_{wk}}$
- $\zeta_i = 1$ for $i = 1s, 1c, 2s, 2c, 3, 4, 7$, $\zeta_i = -1$ for $i = 5, 6s, 6c, 8, 9$

Two CP-conjugate amplitudes

For $M_1 M_2$ CP eigenstate, two CP-related amplitudes

- Theoretical: CP-related amplitudes

$$\bar{A}_X = A_X(\bar{B} \rightarrow \bar{M}_1 \bar{M}_2 \ell \ell) = A_X|_{\phi_{wk} \rightarrow -\phi_{wk}}$$

- Phenomenological: Decay amplitude into the same final state

$$\tilde{A}_X = A_X(\bar{B} \rightarrow M_1 M_2 \ell \ell)$$

From [\[Dunietz et al. 1991\]](#) transversity analysis for $B \rightarrow A(\rightarrow A_1 A_2)C(\rightarrow C_1 C_2)$

$$\tilde{A}_X = \eta_X \bar{A}_X \quad \eta_{L0,L||,R0,R||,t} = \eta \quad \eta_{L\perp,R\perp,S} = -\eta \quad \eta = 1$$

so that $\tilde{J}_i = \zeta_i \bar{J}_i$, and $d\Gamma[\bar{B} \rightarrow \bar{V}(\rightarrow \bar{M}_1 \bar{M}_2)\ell^+\ell^-]$ involves \tilde{J}_i

Untagged $d\Gamma(B \rightarrow V\ell\ell) + d\Gamma(\bar{B} \rightarrow \bar{V}\ell\ell)$ yields $J_i + \tilde{J}_i = J_i + \zeta_i \bar{J}_i$, with both CP-conserving ($\zeta_i = 1$) and CP-violating quantities ($\zeta_i = -1$)

Time dependence

Time-dependence of decay amplitudes is straightforward,
involving decays into the same CP-eigenstate

$$A_X(t) = A_X(B(t) \rightarrow V(\rightarrow f_{CP}) \rightarrow \ell^+ \ell^-) = g_+(t)A_X + \frac{q}{p}g_-(t)\tilde{A}_X,$$

$$\tilde{A}_X(t) = A_X(\bar{B}(t) \rightarrow V(\rightarrow f_{CP}) \ell^+ \ell^-) = \frac{p}{q}g_-(t)A_X + g_+(t)\tilde{A}_X,$$

where $g_{\pm}(t)$ are time-evolution functions and $q/p = e^{i\phi}$

Time dependence of angular coefficients is given by

$$J_i(t) + \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i + \tilde{J}_i) \cosh(y\Gamma t) - h_i \sinh(y\Gamma t) \right]$$

$$J_i(t) - \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i - \tilde{J}_i) \cos(x\Gamma t) - s_i \sin(x\Gamma t) \right]$$

- $y = \Delta\Gamma/(2\Gamma)$ (small for B_d and B_s)
- $x = \Delta m/\Gamma$ ($x_d \simeq 0.77$, $x_s \simeq 27$)

Typical observables

$$J_i(t) + \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i + \tilde{J}_i) \cosh(y\Gamma t) - h_i \sinh(y\Gamma t) \right],$$

$$J_i(t) - \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i - \tilde{J}_i) \cos(x\Gamma t) - s_i \sin(x\Gamma t) \right],$$

Similarly to CP-violation in interference between mixing and decay, new observables from interf between 2 decay amplitudes and mixing

$$J_8 = \frac{1}{\sqrt{2}} \beta_\ell^2 \left[\text{Im}(A_0^L A_\perp^{L*} + A_0^R A_\perp^{R*}) \right],$$

$$\tilde{J}_8 = \frac{1}{\sqrt{2}} \beta_\ell^2 \left[\text{Im}(\tilde{A}_0^L \tilde{A}_\perp^{L*} + \tilde{A}_0^R \tilde{A}_\perp^{R*}) \right] = -\frac{1}{\sqrt{2}} \beta_\ell^2 \left[\text{Im}(\bar{A}_0^L \bar{A}_\perp^{L*} + \bar{A}_0^R \bar{A}_\perp^{R*}) \right],$$

$$h_8 = \frac{1}{\sqrt{2}} \beta_\ell^2 \text{Im} \left[e^{i\phi} \{ \tilde{A}_0^L A_\perp^{L*} + \tilde{A}_0^R A_\perp^{R*} \} + e^{-i\phi} \{ A_0^L \tilde{A}_\perp^{L*} + A_0^R \tilde{A}_\perp^{R*} \} \right]$$

$$s_8 = -\frac{1}{\sqrt{2}} \beta_\ell^2 \text{Re} \left[e^{i\phi} \{ \tilde{A}_0^L A_\perp^{L*} + \tilde{A}_0^R A_\perp^{R*} \} - e^{-i\phi} \{ A_0^L \tilde{A}_\perp^{L*} + A_0^R \tilde{A}_\perp^{R*} \} \right]$$

h_i 's identify with J_i 's in the limit where weak phases neglected

Sorting out observables

$$J_i(t) + \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i + \tilde{J}_i) \cosh(y\Gamma t) - h_i \sinh(y\Gamma t) \right],$$

$$J_i(t) - \tilde{J}_i(t) = e^{-\Gamma t} \left[(J_i - \tilde{J}_i) \cos(x\Gamma t) - s_i \sin(x\Gamma t) \right],$$

- $y \ll 1$: h_i difficult to extract
- from $(d\Gamma + d\bar{\Gamma})/dq^2$, one gets $3(2h_{1s} + h_{1c}) - (2h_{2s} + h_{2c})$
(boils down to the corresponding J 's if $\phi \rightarrow 0$)
- s_i for $i = 1s, 1c, 2s, 2c, 3, 4, 7$: CP-asymmetries $J_i - \bar{J}_i$
- s_i for $i = 5, 6s, 6c, 8, 9$: CP-averaged angular coefficients $J_i + \bar{J}_i$.

If vanishing phases ($\phi \rightarrow 0$, decay amplitudes real)

- s_i for $i = 1s, 1c, 2s, 2c, 3, 4, 5, 6s, 6c$ vanish: $s_i \sim \text{Im}(e^{i\phi} \bar{A}_X A_Y^*)$
- $s_7 = 0$ (no phases in decay amplitudes is enough)
- $(J_i - \tilde{J}_i)_{i=8,9}$ vanish whereas $s_{8,9}$ expected to be large

$\implies s_8$ and s_9 are the most interesting coefficients

New information ?

Not all observables contain new information : there is some redundancy already in the J_i 's

[Matias, Mescia, Ramon, Virto 2012]

- In the flavour-specific case (massless case without scalar contributions), unitary transformation U of

$$n_i = \begin{pmatrix} A_i^L \\ \sigma_i A_i^{R*} \end{pmatrix} \rightarrow U n_i \quad \sigma_0 = \sigma_{||} = 1, \sigma_{\perp} = -1$$

leave the angular coefficient J_i unchanged: only observables invariant under these unitarity transformations can be measured

- in the limit of vanishing weak phases, h_i do not contain genuinely new information compared to the J_i
(but useful as independent cross-checks of J_i measurements)
- $s_{8,9}$ contain genuinely new pieces of information

Time dependent versus time integrated

From time-integrated observables ? Time integration different for hadronic machines and B -factories (quantum entanglement)

$$\langle X \rangle_{\text{Hadronic}} = \int_0^{\infty} e^{-\Gamma t} \dots \quad \langle X \rangle_{\text{B-factory}} = \int_{-\infty}^{\infty} e^{-\Gamma|t|} \dots$$

$$\langle J_i + \tilde{J}_i \rangle_{\text{Hadronic}} = \frac{1}{\Gamma} \left[\frac{1}{1-y^2} \times (J_i + \tilde{J}_i) - \frac{y}{1-y^2} \times h_j \right],$$

$$\langle J_i - \tilde{J}_i \rangle_{\text{Hadronic}} = \frac{1}{\Gamma} \left[\frac{1}{1+x^2} \times (J_i - \tilde{J}_i) - \frac{x}{1+x^2} \times s_j \right],$$

$$\langle J_i + \tilde{J}_i \rangle_{\text{B-factory}} = \frac{2}{\Gamma} \frac{1}{1-y^2} [J_i + \tilde{J}_i], \quad \langle J_i - \tilde{J}_i \rangle_{\text{B-factory}} = \frac{2}{\Gamma} \frac{1}{1+x^2} [J_i - \tilde{J}_i].$$

s_j and h_j from time-integrated measurements

- only at hadronic machines (but tagging needed for s_j)
- suppressed by factors of y or $1/(1+x^2)$

Optimised observables from time dependence

S_8, S_9

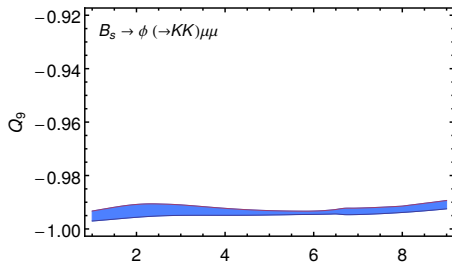
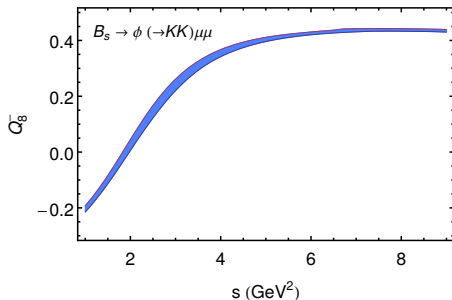
- contain information that is not accessible otherwise
- come from $J_i - \tilde{J}_i$ and require tagging
- are coefficients of $\sin(x\Gamma t)$ and require time-dependent analysis
 $\implies B$ -factory environment ?

It is possible to define **optimised observables** at large hadronic recoil (limited sensitivity to form factors)

$$Q_8^- = \frac{S_8}{\sqrt{-2(J_{2c} + \tilde{J}_{2c})[2(J_{2s} + \tilde{J}_{2s}) - (J_3 + \tilde{J}_3)]}},$$
$$Q_9 = \frac{S_9}{2(J_{2s} + \tilde{J}_{2s})}.$$

similarly to what is done from J_i to P_i

Q_8, Q_9 : SM predictions



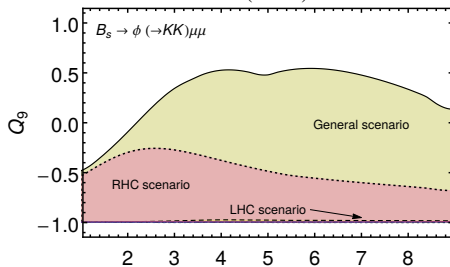
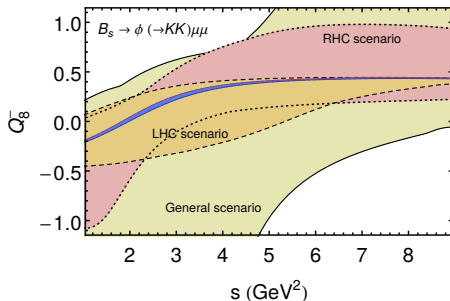
- In SM, $Q_9 \simeq -1$, is a test of RHC
- In SM, zero of Q_8 given at LO by:

$$\frac{s_0}{m_B^2} \simeq \frac{-2C_7(2C_7 + C_9)}{C_{10}^2 + (2C_7 + C_9)C_9}$$

(modified by RHC)

- Similar plots for the other modes

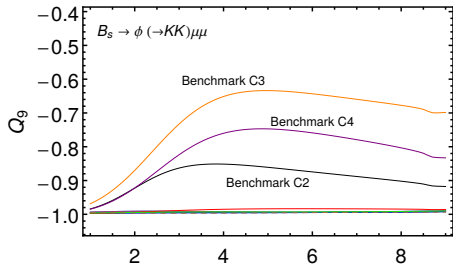
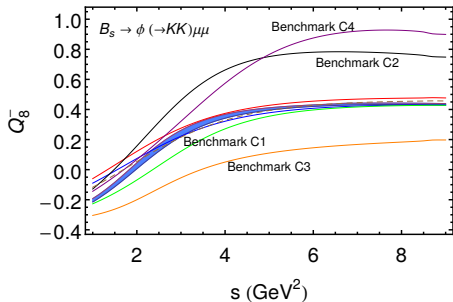
Q_8, Q_9 : General NP scenarios



- LHC: C_7, C_9, C_{10} only
- RHC: $C_{7'}, C_{9'}, C_{10'}$ only
- General NP: All
- varying in 3σ ranges of

[SDG, Matias, Virto 2013]

Q_8, Q_9 : Benchmark points



- A: $\mathcal{C}_7, \mathcal{C}_9$ best fit
- B: $\mathcal{C}_9, \mathcal{C}_{9'}$ best fit
- C: $\mathcal{C}_{9(')}, \mathcal{C}_{10(')}$ scenarios
- D: general best fit

Conclusion

Time-dependent analysis of $B \rightarrow V\ell\ell$ with V into CP eigenstate

- Mixing allowing richer pattern of interferences
- Concerns both $B_d \rightarrow K^*(\rightarrow K_S\pi^0)\ell^+\ell^-$ and
 $B_s \rightarrow \phi(\rightarrow K_S K_L)\ell^+\ell^-$, $B_s \rightarrow \phi(\rightarrow K^+K^-)\ell^+\ell^-$
- Two interesting new observables s_8 and s_9
- Require both tagging and time-dependent analysis

Optimised versions Q_8 and Q_9

- Accurate predictions in the SM
- Value of Q_9 good test of right-handed currents
- Good sensitivity to NP scenarios

Experimental feasibility of such measurements ?