$B \rightarrow M\ell\ell$: Quick Theory Overview

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Theoretishe Physik 1

Aim

To avoid repetition of trivial things in (theory) talks

• I will say the trivial things here once for all in this workshop (hopefully)

Success of this overview talk:

$$S = 1 - \frac{\text{Minutes of introduction in theory talks}}{\text{Total minutes of theory talks}}$$

• Likewise I hope to see the LHCb detector picture only once in exp. talks...

First trivialities

- New Physics is up there (or even down here...)
 The SM is not complete... etc. etc.
- Rare B decays are an excellent tool to look for it
- Also interesting for QCD issues

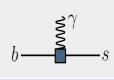
(Can we do both things at the same time?)

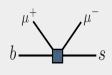
So we know why this workshop is interesting

EFT at $\mu = m_b$ and $b \rightarrow s$ transitions

Radiative and Dileptonic $b \rightarrow s$ Operators

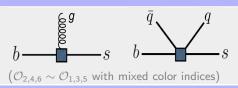
$$\begin{array}{rcl} \mathcal{O}_{7(')} & = & [\bar{s}\sigma^{\mu\nu}P_{R(L)}b]F_{\mu\nu} \\ \mathcal{O}_{9(')} & = & [\bar{s}\gamma^{\mu}P_{L(R)}b][\bar{\ell}\gamma_{\mu}\ell] \\ \mathcal{O}_{10(')} & = & [\bar{s}\gamma^{\mu}P_{L(R)}b][\bar{\ell}\gamma_{\mu}\gamma_{5}\ell] \\ \mathcal{O}_{S(')}, \mathcal{O}_{P(')}, \mathcal{O}_{T,T5} \end{array}$$





Hadronic $b \rightarrow s$ Operators

$$\begin{array}{rcl} \mathcal{O}_1 & = & [\bar{s}\gamma^{\mu}P_Lc][\bar{c}\gamma_{\mu}P_Lb] \\ \mathcal{O}_{3(5)} & = & [\bar{s}\gamma^{\mu}P_Lb]\sum_q[\bar{q}\gamma_{\mu}P_{L(R)}q] \\ \mathcal{O}_{8g} & = & [\bar{s}\sigma^{\mu\nu}P_{R(L)}T^ab]G^a_{\mu\nu} \end{array}$$

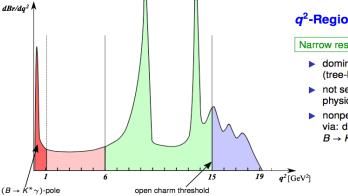


Effective Hamiltonian

$$\mathcal{H}_{\mathrm{eff}} = -\frac{4\textit{G}_{\textit{F}}}{\sqrt{2}}\textit{V}_{\textit{tb}}\textit{V}_{\textit{ts}}^* \left[\sum_{7,7',9,9',10,10'} \mathcal{C}_{\textit{i}}\mathcal{O}_{\textit{i}} + \sum_{1,\dots,6,8g} \mathcal{C}_{\textit{i}}\mathcal{O}_{\textit{i}} \right]$$

 $\mathcal{C}_{7\mathrm{eff}}^{\mathrm{SM}} = -0.3, \ \mathcal{C}_{9}^{\mathrm{SM}} = 4.1, \ \mathcal{C}_{10}^{\mathrm{SM}} = -4.3, \ \mathcal{C}_{1}^{\mathrm{SM}} = 1.1, \ \mathcal{C}_{2}^{\mathrm{SM}} = -0.4, \ \mathcal{C}_{\mathrm{rest}}^{\mathrm{SM}} \lesssim 10^{-2}$

$m_{\ell\ell}^2$ spectrum



q^2 -Regions in $B \to K^* \bar{\ell} \ell$

Narrow resonances

- dominated by charged-cur. (tree-level) op's
- not sensitive to new physics in $b \rightarrow s\bar{\ell}\ell$
- nonperturbative predictions via: dispersion relations + $B \rightarrow K^*(\bar{c}c)$ data

Large Recoil (low-q2)

- ▶ very low- q^2 ($\lesssim 1 \text{ GeV}^2$) dominated by \mathcal{O}_7
- low- q^2 ([1,6] GeV²) dominated by $\mathcal{O}_{9,10}$
- 1) QCD factorization or SCET 2) LCSR 3) non-local OPE of cc-tails

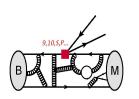
Low Recoil (high-q2)

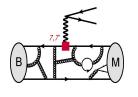
- dominated by $\mathcal{O}_{9,10}$
- local OPE (+ HQET) ⇒ theory only for sufficiently large q2-integrated obs's

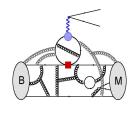
(slide from C. Bobeth)

EFT Amplitudes

$$\mathcal{L} = \mathcal{L}_{QED+QCD} - \mathcal{C}_7 \left[\bar{\mathbf{s}} \sigma^{\mu\nu} P_R b \right] F_{\mu\nu} - \mathcal{C}_2 \left[\bar{\mathbf{s}} \gamma^{\nu} P_L c \right] \left[\bar{c} \gamma^{\mu} P_L b \right] + \cdots$$







$$C_9$$
 contribution: $A_9 = C_9 \langle M_{\lambda} | \bar{s} \gamma_{\mu} P_L b | B \rangle L^{\mu} = C_9 F_{\lambda}(q^2)$

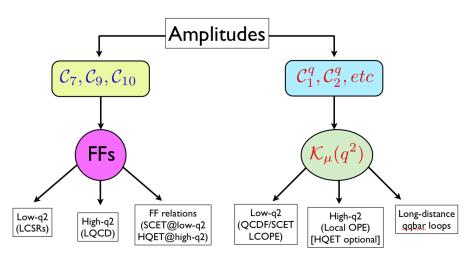
$$\mathcal{C}_7$$
 contribution: $\mathcal{A}_7 = \frac{\mathcal{C}_7}{M_\lambda} |\bar{s}\sigma_{\mu\nu} P_R b| B \frac{eq^\mu}{q^2} L^\nu = \frac{\mathcal{C}_7}{T_\lambda} T_\lambda(q^2)$

$$C_2$$
 contribution: $A_2 = C_2 \cdot \frac{e^2}{q^2} L^{\mu} \int dx^4 e^{iq \cdot x} \langle M_{\lambda} | T \{ \mathcal{J}_{\mu}^{em}(x) \mathcal{O}_2(0) \} | B \rangle$

2 main problems:

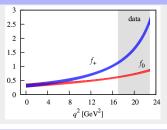
- 1. Determination of Form Factors (LCSRs, LQCD, ...)
- 2. Computation of the hadronic contribution (SCET/QCDF, OPE, ...)

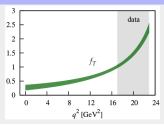
Outline-chart of theory amplitudes



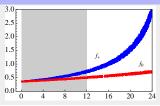
$$\left(\text{where}\quad \mathcal{K}_{\mu}(\textbf{\textit{q}}^2) = -\frac{8\pi}{\textbf{\textit{q}}^2}\int \textbf{\textit{d}}^4\textbf{\textit{x}}e^{\textbf{\textit{i}}\textbf{\textit{q}}\cdot\textbf{\textit{x}}}T\{\mathcal{J}_{\mu}^{\text{em}}(\textbf{\textit{x}})\mathcal{H}(0)\}\right)$$

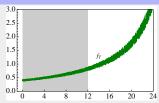
Unquenched LQCD Bouchard, Lepage, Monahan, Na, Shigemitsu '2013





LCSRs Khodjamirian, Mannel, Pivovarov, Wang '2010





 $B \to K$ form factors are well known in the full kinematical regime, but still constitute a dominant source of uncertainty in some regions.

$B \to K\ell\bar{\ell}$: Hadronic Contributions – Low- q^2

"Hard" contributions: QCDF Beneke, Feldmann, Seidel'2001, Asatrian, et al'2001



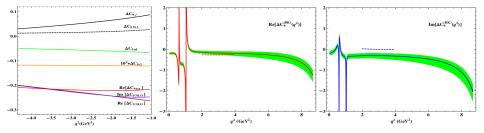
up to unknown Power Corrections

"Soft" contributions: LC-OPE at $q^2 < 0$ Khodjamirian, Mannel, Wang'2012

$$\mathcal{K}_{\mathcal{BK}}(q^2) \sim \log$$

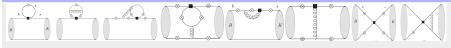
$$\mathcal{K}_{BK}(q^2) \sim \log(4m_c^2 - q^2) + \sum_{n=1,2,...} \frac{\Lambda_{\text{QCD}}^{2n}}{(4m_c^2 - q^2)^n}$$

Contain Power Corrections and reproduce QCDF for hard contributions



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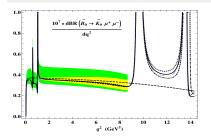


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Contain Power Corrections and reproduce QCDF for hard contributions



- ullet Soft contributions reduce the BR by a few %.
- FFs are the dominant source of uncertainty below $\sim 6~\text{GeV}^2.$
- Model-dependence is small in [1,8] GeV².

$B \to K\ell\bar{\ell}$: Hadronic Contributions – High- q^2

Local OPE for $q^2 \sim m_b$: Beylich, Buchalla, Feldman'2011, Grinstein, Pirjol'2004

$$\mathcal{K}^{\mu}(q^{2}) \sim \int dx^{4} e^{iq.x} T\{j_{c\bar{c}}^{\mu}(x)\mathcal{H}^{c}(0)\} \xrightarrow{q^{2} \sim m_{b}^{2}} \sum_{n,d} C_{d,n}(q)\mathcal{O}_{d,n}^{\mu}$$

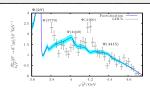
$$= \mathcal{K}_{3}^{\mu}(q^{2}) + \mathcal{K}_{5}^{\mu}(q^{2}) + \mathcal{K}_{6}^{\mu}(q^{2}) + \mathcal{O}[(\Lambda/m_{b})^{3}]$$
Usual FFs
NLO $\sim 10^{-15}\%$ $\lesssim 1\%$ $\lesssim 1\%$

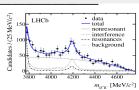
Duality violations: Beylich, Buchalla, Feldman'2011

From a model fitted to BES data: $\pm 2\%$ for integrated BR over high-q² region.

Non-fact. corrections: Seem to be large

figures from Lyon, Zwicky' 2014 and LHCb-1307.7595

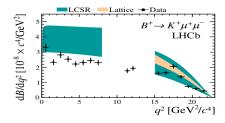




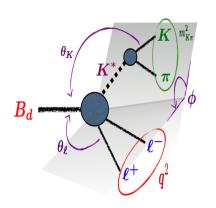
$B \to K \ell \bar{\ell}$: Summary

- Form Factors are well known but still a dominant source or uncertainty.
- Perturbative NLO corrections to charm-loop are important since they lift color suppression:
 - $ightharpoonup \sim -5\%$ in the space-like region.
 - $\sim -10\%$ at low- q^2 .
 - $ightharpoonup \sim -10\text{-}15\%$ at high- q^2 .
- Soft contributions reduce the BR by a few percent.
- At high- q^2 wide resonances are difficult to assess theoretically, but the OPE should give an accurate determination for the integrated BR in a large high- q^2 bin (DV contributions $\lesssim \pm 2\%$??).
- There is **some tension** between SM and data, at least at **low**- q^2 [remember $BR \sim (C_9 + C_{9'})$, so it goes in the direction of decreasing C_9]. But theory correlations are O(1)!!

Plot from LHCb-1403.8044



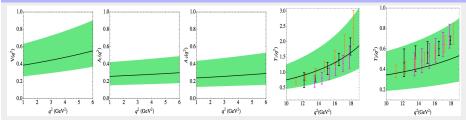
$B \to K^*(\to K\pi)\ell^+\ell^-$: Angular Distribution



$$\begin{split} \frac{d^4\Gamma}{dq^2\,d\cos\theta_K\,d\cos\theta_I\,d\phi} &= \frac{9}{32\pi} \times \\ \left[\mathbf{J}_{1s}\sin^2\theta_K + \mathbf{J}_{1c}\cos^2\theta_K + \mathbf{J}_{2s}\sin^2\theta_K\cos2\theta_I \right. \\ &+ \mathbf{J}_{2c}\cos^2\theta_K\cos2\theta_I + \mathbf{J}_3\sin^2\theta_K\sin^2\theta_I\cos2\phi \\ &+ \mathbf{J}_4\sin2\theta_K\sin2\theta_I\cos\phi + \mathbf{J}_5\sin2\theta_K\sin\theta_I\cos\phi \\ &+ \mathbf{J}_6\sin^2\theta_K\cos\theta_I + \mathbf{J}_{6c}\cos^2\theta_K\cos\theta_I \\ &+ \mathbf{J}_7\sin2\theta_K\sin\theta_I\sin\phi + \mathbf{J}_8\sin2\theta_K\sin2\theta_I\sin\phi \\ &+ \mathbf{J}_9\sin^2\theta_K\sin^2\theta_I\sin2\phi \right] \end{split}$$

- * Which observables do we want to have? $(P_i^{(\prime)} \text{ vs. } S_i, \text{ etc.})$
- * CP violation?
- * Time dependence? (not for this mode)
- * Polarization?? Taus?? ...

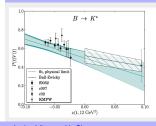


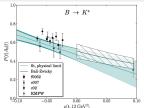


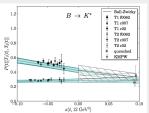
also in fair agreement with Ball, Zwicky'2004

(figs from Descotes-Genon, Hurth, Matias, JV'2013)

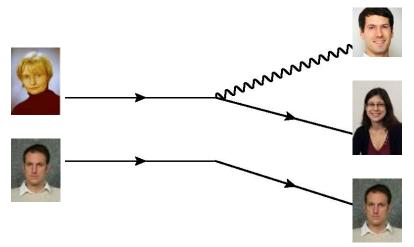
Recent "full QCD" LQCD results Horgan et.al'2013

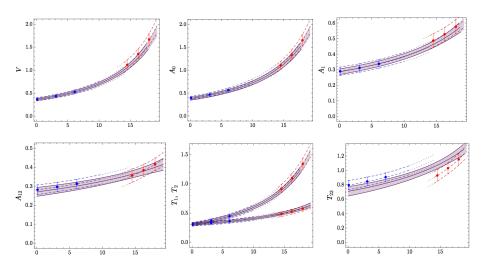






From BZ to BSZ: Very rare decay





Bharucha, Straub, Zwicky'2015

Low q^2 ::

Altmannshofer, Bharucha, Straub, Zwicky:

LCSRs with K^* DAs + Correlations + EOM constraint q^2 dependence given by simplified z-expansion

• Descotes-Genon et al:

LCSRs with B DAs (uncorrelated) + SCET relations + Power corrections q^2 dependence given by simplified z-expansion

• Jäger + Camalich:

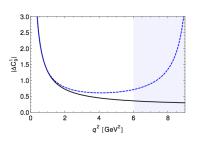
Try to rely only on HQ/LE expansion, both for $q^2=0$ and q^2 -dependence Input: LCSRs, DSE, $B\to K^*\gamma$, + power corrections

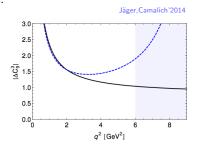
Large q^2 ::

• Horgan et al: Lattice QCD

$B \to K^* \ell \bar{\ell}$: Hadronic Contributions

Similar to $B \to K\ell\ell$ with some differences.

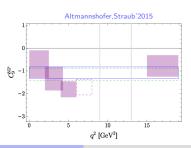




It mimics a C_9 but...

- The effect is q^2 -depedent
- The effect is different for each helicity/transversity amplitude
- Phases: Weak vs. Strong

Would be nice to see this clearly from the data.



LHCb $b \rightarrow s$ Anomalies

1. Angular observables in $B \to K^* \mu \mu$. $C_9^{\mu} < 0$

Descotes-Genon et al'2013, Altmannshofer, Straub'2013, Horgan et al'2013, Bobeth et al'2013, \dots + others + updates

2. Branching ratios: $B \to K\mu\mu$, $B \to K^*\mu\mu$, $B \to \phi\mu\mu$. $C_9^\mu < 0$ \checkmark

Roman: "it is *not* a factor of 2 in the normalisation of *my* form factors!!"

3. $R_K \equiv BR(B \to K\mu\mu)/BR(B \to Kee)$ at low- q^2 . $C_9^{\mu} < 0$ \checkmark

Hiller, Schmaltz' 2014, Gosh et al' 2014, Hurth et al' 2014, Altmannshofer, Straub' 2014, $\dots + \text{others} + \text{updates}$

Note: Anomalies live mostly in the low- q^2 , would be nice to have them also at high- q^2 ...

LHCb $b \rightarrow s$ Anomalies

	branching ratios	angular observables	LFU ratios
parametric uncertainties?	\checkmark	×	×
hadronic effects?	\checkmark	\checkmark	×
New Physics?	$\sqrt{}$	\checkmark	\checkmark

(slide from W. Altmannshofer)

LHCb $b \rightarrow s$ Anomalies

R_K: null test in SM?

$$R_K^{\rm SM} = 1 + \mathcal{O}\left(\frac{m_\mu^2}{m_b^2}\right) + \mathcal{O}\left(\alpha \ln \frac{m_\mu^2}{m_b^2}\right) = 1 + \frac{\mathcal{O}(0.01)}{\mathcal{O}(0.01)}$$

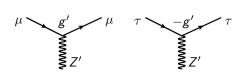
(slide from U. Haisch '2014)

A prediction could be that the effect is *opposite* in the large- q^2 Huber, Hurth, Lunghi 2015 But it seems waaay too large anyway

New Physics?

"Obvious" possibility: Z'-models

Buras et al, Descotes-Genon et al, Haisch et al, Altmannshofer et al, Crivellin et al, Aristizabal et al, (mostly everyone)



 $L_{\mu}-L_{\tau}$ model Altmannshofer et al'2013 Lepton-non-universal Causality violation ("Cristal Ball" effect)

- (Scalar) Leptoquaks Hiller+Schmaltz, Shoo+Mohanta, Nardecchia et al, Hiller+Ivo,...
 - Attractive for a number of reasons
- Please fill in
- Lepton Flavor violation??

Glashow et al, Bhattacharya et al, Crivellin et al, Aristizabal et al, Hiller+Ivo, Crivellin et al ... (almost everyone)

- ► Glashow, Guadagnoli, Lane: "LFNU "necessarily associated" with LFV "
- ▶ Grinstein: "Bullshit!" (More precisely: "I have a family of counterexamples")
- ▶ If you pay attention you'll see they do not really disagree
- ▶ But most proposed models for R_K do have that feature
- ▶ Important to establish what to expect on general grounds (benchmarks, etc.)

Conclusions

No Conclusions. Let's begin!!