# Global fits to $bs\ell\ell$ data

# Nazila Mahmoudi

Lyon University & CERN TH

In collaboration with T. Hurth and S. Neshatpour

Rare B decays in 2015 - experiment and theory Higgs Centre for Theoretical Physics, Edinburgh 11-13 May 2015

# **Inclusive decays**

•  $B \to X_s \gamma$ 

Improved theory calculations (Misiak et al. 1503.01789) Excellent agreement with the measurements

•  $B \to X_s \ell^+ \ell^-$ 

Still waiting for the final words from Belle and Babar! High expectation from Belle II!

# **Exclusive decays**

- $B \to K^* \gamma$
- First measurements of  $B_s o \mu^+ \mu^-$
- Angular distributions of B → K<sup>\*</sup>µ<sup>+</sup>µ<sup>-</sup> → large variety of experimentally accessible observables

• Also: 
$$B o K \mu^+ \mu^-$$
 and  $B_s o \phi \mu^+ \mu^-$ 

### Issue of hadronic uncertainties in exclusive modes

Differential decay distribution:

$$\frac{d^4\Gamma}{dq^2\,d\cos\theta_\ell\,d\cos\theta_V\,d\phi} = \frac{9}{32\pi}J(q^2,\theta_\ell,\theta_V,\phi)$$

 $J(q^2,\theta_\ell,\theta_V,\phi) = \sum_i J_i(q^2) f_i(\theta_\ell,\theta_V,\phi)$ 

 $\searrow$  angular coefficients  $J_{1-9}$ 

 $\searrow$  functions of the spin amplitudes  $A_0$ ,  $A_{\parallel}$ ,  $A_{\perp}$ ,  $A_t$ , and  $A_s$ Spin amplitudes: functions of Wilson coefficients and form factors

### Standard Observables:

Dilepton invariant mass spectrum: 
$$\frac{d\Gamma}{dq^2} = \frac{3}{4} \left( J_1 - \frac{J_2}{3} \right)$$
  
Forward backward asymmetry:  
$$A_{\rm FB}(q^2) \equiv \left[ \int_{-1}^0 - \int_0^1 \right] d\cos\theta_l \frac{d^2\Gamma}{dq^2 d\cos\theta_l} \Big/ \frac{d\Gamma}{dq^2} = \frac{3}{8} J_6 \Big/ \frac{d\Gamma}{dq^2}$$
  
Forward backward asymmetry zero-crossing: 
$$q_0^2 \simeq -2m_b m_B \frac{C_9^{\rm eff}(q_0^2)}{C_7} + O(\alpha_s, \Lambda/m_b)$$
  
Polarization fraction: 
$$F_L(q^2) = \frac{|A_0|^2}{|A_0|^2 + |A_{\perp}|^2 + |A_{\perp}|^2},$$

Nazila Mahmoudi

# $B \rightarrow V \mu^+ \mu^-$ observables

Optimised observables: form factor uncertainties cancel at leading order

$$\langle P_1 \rangle_{\rm bin} = \frac{1}{2} \frac{\int_{\rm bin} dq^2 [J_3 + \bar{J}_3]}{\int_{\rm bin} dq^2 [J_{2s} + \bar{J}_{2s}]} \qquad \langle P_2 \rangle_{\rm bin} = \frac{1}{8} \frac{\int_{\rm bin} dq^2 [J_{6s} + \bar{J}_{6s}]}{\int_{\rm bin} dq^2 [J_{2s} + \bar{J}_{2s}]} \\ \langle P'_4 \rangle_{\rm bin} = \frac{1}{N'_{\rm bin}} \int_{\rm bin} dq^2 [J_4 + \bar{J}_4] \qquad \langle P'_5 \rangle_{\rm bin} = \frac{1}{2N'_{\rm bin}} \int_{\rm bin} dq^2 [J_5 + \bar{J}_5] \\ \langle P'_6 \rangle_{\rm bin} = \frac{-1}{2N'_{\rm bin}} \int_{\rm bin} dq^2 [J_7 + \bar{J}_7] \qquad \langle P'_8 \rangle_{\rm bin} = \frac{-1}{N'_{\rm bin}} \int_{\rm bin} dq^2 [J_8 + \bar{J}_8]$$

with

$$\mathcal{N}_{
m bin}' = \sqrt{-\int_{
m bin} dq^2 [J_{2s} + ar{J}_{2s}] \int_{
m bin} dq^2 [J_{2c} + ar{J}_{2c}]}$$

+ CP violating clean observables and other combinations

U. Egede et al., JHEP 0811 (2008) 032, JHEP 1010 (2010) 056 J. Matias et al., JHEP 1204 (2012) 104 S. Descotes-Genon et al., JHEP 1305 (2013) 137

Or alternatively:

$$S_i = \frac{J_{i(s,c)} + \bar{J}_{i(s,c)}}{\frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2}}$$

W. Altmannshofer, P. Ball, A. Bharucha, A.J. Buras, D.M. Straub, M. Wick, JHEP 0901 (2009) 019

Nazila Mahmoudi

Rare B decays in 2015 - Edinburgh - 13 May 2015

### The LHCb anomalies

## 3 main LHCb anomalies:

- P'\_5
- *R*<sub>K</sub>
- BR( $B_s \rightarrow \phi \mu^+ \mu^-$ )



Possible explanations:

- Statistical fluctuations
- Theoretical issues
- New Physics!

Global analysis of the latest LHCb data

Relevant Operators:

$$\mathcal{O}_7$$
,  $\mathcal{O}_8$ ,  $\mathcal{O}_{9\mu,e}^{(')}$ ,  $\mathcal{O}_{10\mu,e}^{(')}$  and  $\mathcal{O}_{S-P} \propto (\bar{s}P_R b)(\bar{\mu}P_L \mu) \equiv \mathcal{O}_0^{\prime}$ 

 $\mathsf{NP}$  manifests itself in the shifts of the individual coefficients with respect to the SM values:

$$C_i(\mu) = C_i^{\rm SM}(\mu) + \delta C_i$$

- ightarrow Scans over the values of  $\delta C_i$
- $\rightarrow$  Calculation of flavour observables
- $\rightarrow$  Comparison with experimental results
- $\rightarrow$  Constraints on the Wilson coefficients  $C_i$

#### **Evaluations uncertainties and correlations:**

- Experimental errors and correlations
  - 3 fb^{-1} LHCb data for  $B 
    ightarrow {\cal K}^{*0} \mu^+ \mu^-$ : provided in LHCb-CONF-2015-002
- Theoretical uncertainties and correlations
  - study of more than 100 observables (at a later stage, selection of the relevant observables for each fit)
  - Monte Carlo analysis
  - variation of the "standard" input parameters: masses, scales, CKM, ...
  - for  $B_s \rightarrow \phi \mu^+ \mu^-$ , mixing effects taken into account
  - · decay constants taken from the latest lattice results
  - use for the  $B_{(s)} \rightarrow V$  form factors of the lattice+LCSR combinations from 1503.05534, including correlations (Cholesky decomposition method)
  - use for the  $B \to K$  form factors of the lattice+LCSR combinations from 1411.3161, including correlations
  - two approaches for the exclusive decays: soft form factors, full form factors
  - two sets of hypotheses for the uncertainties associated to the factorisable and non-factorisable power corrections

 $\Rightarrow$  Computation of a (theory + exp) correlation matrix

For the exclusive semi-leptonic decays, two approaches and two evaluations of the uncertainties for each decay.

At low  $q^2$ :

• Soft form factor approach

Uncertainties of the factorisable and non-factorisable corrections parametrised as

$$A_k 
ightarrow A_k \left(1 + \mathsf{a}_k \exp(i\phi_k) + rac{q^2}{6 \ {
m GeV}^2} b_k \exp(i\theta_k)
ight)$$

where  $A_k$  are the helicity amplitudes.

$$a_k$$
 in  $[-10\%, +10\%]$  or  $[-20\%, +20\%]$  $\phi_k, \theta_k$  in  $[-\pi, +\pi]$  $b_k$  in  $[-25\%, +25\%]$  or  $[-50\%, +50\%]$ 

• Full form factor approach

Uncertainties of the non-factorisable power corrections only parametrised in a similar way:

$$\begin{array}{ll} \textbf{a}_k \text{ in } [-5\%,+5\%] \text{ or } [-10\%,+10\%] & \phi_k,\theta_k \text{ in } [-\pi,+\pi] \\ \textbf{b}_k \text{ in } [-10\%,+10\%] \text{ or } [-25\%,+25\%] \end{array}$$

At high  $q^2$ , uncertainties parametrised as

$$egin{aligned} & A_k o A_k ig(1+a_k \exp(i\phi_k)ig) \ & a_k \ & ext{in } [-10\%,+10\%] \ & ext{or } [-20\%,+20\%] \ & \phi_k \ & ext{in } [-\pi,+\pi] \end{aligned}$$

Global fits of the observables by minimization of

$$\chi^2 = \left(\vec{O}^{\texttt{th}} - \vec{O}^{\texttt{exp}}\right) \cdot \left(\Sigma_{\texttt{th}} + \Sigma_{\texttt{exp}}\right)^{-1} \cdot \left(\vec{O}^{\texttt{th}} - \vec{O}^{\texttt{exp}}\right)$$

 $(\Sigma_{\tt th}+\Sigma_{\tt exp})^{-1}$  is the inverse covariance matrix.

58 observables relevant for leptonic and semileptonic decays:

- $BR(B \rightarrow X_s \gamma)$
- $BR(B \rightarrow X_d \gamma)$
- $\Delta_0(B \to K^*\gamma)$
- $\mathsf{BR}^{\mathsf{low}}(B \to X_s \mu^+ \mu^-)$
- $\mathsf{BR}^{\mathsf{high}}(B \to X_s \mu^+ \mu^-)$
- $\mathsf{BR}^{\mathsf{low}}(B \to X_s e^+ e^-)$
- $\mathsf{BR}^{\mathsf{high}}(B \to X_s e^+ e^-)$
- BR( $B_s \rightarrow \mu^+ \mu^-$ )
- BR( $B_d \rightarrow \mu^+ \mu^-$ )
- BR( $B \rightarrow K^{*+} \mu^+ \mu^-$ )

- BR( $B \rightarrow K^0 \mu^+ \mu^-$ )
- BR( $B 
  ightarrow K^+ \mu^+ \mu^-$ )
- BR( $B \rightarrow K^* e^+ e^-$ )
- R<sub>K</sub>
- $B \rightarrow K^{*0} \mu^+ \mu^-$ :  $F_L$ ,  $A_{FB}$ ,  $S_3$ ,  $S_4$ ,  $S_5$ in five low  $q^2$  and two high  $q^2$ bins
- $B_s \rightarrow \phi \mu^+ \mu^-$ : BR,  $F_L$ in three low  $q^2$  and two high  $q^2$ bins

## Statistical approaches:

- $\Delta \chi^2 = \chi^2 \chi^2_{\min}$  method
  - $\textbf{O} \ \ \text{Determination of the minimum of } \chi^2 \rightarrow \text{best fit point}$
  - **②** Computation for each point of the scan of the difference of  $\chi^2$  with the best fit point
  - **③** Find the  $1 2\sigma$  regions corresponding to the number of d.o.f.

Interpretation: considering the best fit point gives the "real" description, which variations of the parameters are allowed  $\rightarrow$  *relative* global fit

- $\bullet$  Absolute  $\chi^{\rm 2}$  method

  - **(a)** Find the  $1 2\sigma$  regions corresponding to N d.o.f. where  $N = (N_o \text{ observables} n_v \text{ variables})$
  - If an observable is relatively insensitive to the variation of the Wilson coefficients, remove it from the fit

Interpretation: global fit assessing if each point is globally in agreement with all the measurements

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method



with 20% power correction errors:  $\Delta \chi^2$  method



## Fit results for two operators: $\{C_9, C_{10}\}$

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method







# with 20% power correction errors: $\Delta \chi^2$ method



 $1\sigma$  agreement for  $C_9$  is possible even in the 2 operator basis!

 $\rightarrow$  Using full form factors

with 5% power correction errors:

 $\Delta\chi^2$  method



with 10% power correction errors:  $\Delta \chi^2$  method



 $\rightarrow$  Using full form factors

with 5% power correction errors:



Absolute  $\chi^2$  method



# with 10% power correction errors:



Using the full form factors, only  $2\sigma$  agreement for  $C_9$  could be possible.

Fit results for two operators:  $\{C_9, C_9'\}$ 

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method



with 20% power correction errors:  $\Delta \chi^2$  method



Fit results for two operators:  $\{C_9, C_9'\}$ 

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method



Absolute  $\chi^2$  method



# with 20% power correction errors: $\Delta \chi^2$ method

Absolute  $\chi^2$  method



 $1\sigma$  agreement for  $C_9$  is possible even in the 2 operator basis!

 $\rightarrow$  Using full form factors

with 5% power correction errors:  $\Delta \chi^2$  method





with 10% power correction errors:  $\Delta \chi^2$  method



Fit results for two operators:  $\{C_9, C_9'\}$ 

 $\rightarrow$  Using full form factors

with 5% power correction errors:  $\Delta \chi^2$  method



Absolute  $\chi^2$  method



with 10% power correction errors:  $\Delta \chi^2$  method

hod





Using the full form factors, only  $2\sigma$  agreement for  $C_9$  could be possible.

Fit results for two operators:  $\{C_9^e, C_9^\mu\}$ 

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method



with 20% power correction errors:  $\Delta\chi^2$  method



Fit results for two operators:  $\{C_9^e, C_9^\mu\}$ 

 $\rightarrow$  Using soft form factors

with 10% power correction errors:  $\Delta \chi^2$  method











No tension in  $C_{9\mu}$  with 20% errors for power corrections!

 $\rightarrow$  Using full form factors

with 5% power correction errors:  $\Delta \chi^2$  method



with 10% power correction errors:  $\Delta\chi^2$  method



Fit results for two operators:  $\{C_9^e, C_9^\mu\}$ 

 $\rightarrow$  Using full form factors

with 5% power correction errors:  $\Delta \chi^2$  method





Absolute  $\chi^2$  method



Absolute  $\chi^2$  method



 $2\sigma$  agreement still possible!

# Fit results for four operators: $\{C_9, C'_9, C_{10}, C'_{10}\}$

Using full form factors with 5% power correction errors

ightarrow with  $\Delta\chi^2$  method



Adding  $C_{10}$  or primed coefficients doesn't improve the fit

# Fit results for four operators: $\{C_9, C'_9, C_{10}, C'_{10}\}$

Using full form factors with 5% power correction errors

 $\rightarrow$  with absolute  $\chi^{\rm 2}$  method



Adding  $C_{10}$  or primed coefficients doesn't improve the fit

# Fit results for four operators: $\{C_9^{\mu}, C_9^{\prime \mu}, C_9^{e}, C_9^{\prime e}\}$

Using full form factors with 5% power correction errors

ightarrow with  $\Delta\chi^2$  method



#### Separating electron and muon coefficients improves the fit by more than $2\sigma$

Using full form factors with 5% power correction errors

 $\rightarrow$  with absolute  $\chi^{\rm 2}$  method



#### Separating electron and muon coefficients improves the fit by more than $2\sigma$

# Fit results for four operators: $\{C_9^{\mu}, C_{9}^{e}, C_{10}^{\mu}, C_{10}^{e}\}$

Using full form factors with 5% power correction errors

ightarrow with  $\Delta\chi^2$  method



Again the non universal solutions are favoured

# Fit results for four operators: $\{C_9^{\mu}, C_9^{e}, C_{10}^{\mu}, C_{10}^{e}\}$

Using full form factors with 5% power correction errors

 $\rightarrow$  with absolute  $\chi^{\rm 2}$  method



Again the non universal solutions are favoured

# **MFV fit results:** $\{C_7, C_8, C_9, C_{10}, C_0'\}$

Using full form factors with 5% power correction errors

# Preliminary

ightarrow with  $\Delta\chi^2$  method



The tension in  $C_9$  is present even in the MFV fit!

# **MFV fit results:** $\{C_7, C_8, C_9, C_{10}, C_0'\}$

Using full form factors with 5% power correction errors

# Preliminary

 $\rightarrow$  with absolute  $\chi^2$  method



#### The tension in $C_9$ is present even in the MFV fit!

#### Comparison of exclusive and inclusive $b \rightarrow s\ell\ell$ observables



T. Hurth, FM, S. Neshatpour, JHEP 1412 (2014) 053

All three sets of exclusion plots nicely compatible with each other  $\rightarrow$  non-trivial consistency check

Nazila Mahmoudi

Rare B decays in 2015 - Edinburgh - 13 May 2015

#### Summary

- ullet There are two possible statistical approaches:  $\Delta\chi^2$  and absolute  $\chi^2$
- With  $\Delta\chi^2,$  the error on power corrections has a smaller impact than with absolute  $\chi^2$
- With  $\Delta \chi^2$ , in the two operator fits {C<sub>9</sub>, C<sub>10</sub>}, {C<sub>9</sub>, C<sub>9</sub>'}, {C<sub>9</sub><sup> $\mu$ </sup>, C<sub>9</sub><sup>e</sup>}, SM shows more than  $2\sigma$  tension

 $\rightarrow$  In principle there is no reason to consider only 2 operators!

- In the 4 operator fits, the tension in  $C_9$  weakens but still exists at the  $2\sigma$  level
- The tension in  $C_9$  is also seen in the MFV fit

The tensions are almost GONE with abs  $\chi^2$  and 20% error for the power corrections! The largest remaining tension is for  $C_9^{\mu}$  but it's less than  $2\sigma$ 

#### Summary

- There are two possible statistical approaches:  $\Delta\chi^2$  and absolute  $\chi^2$
- With  $\Delta\chi^2,$  the error on power corrections has a smaller impact than with absolute  $\chi^2$
- With  $\Delta \chi^2$ , in the two operator fits {C<sub>9</sub>, C<sub>10</sub>}, {C<sub>9</sub>, C<sub>9</sub>'}, {C<sub>9</sub><sup> $\mu$ </sup>, C<sub>9</sub><sup>e</sup>}, SM shows more than  $2\sigma$  tension

 $\rightarrow$  In principle there is no reason to consider only 2 operators!

- In the 4 operator fits, the tension in  $C_9$  weakens but still exists at the  $2\sigma$  level
- The tension in  $C_9$  is also seen in the MFV fit

The tensions are almost GONE with abs  $\chi^2$  and 20% error for the power corrections! The largest remaining tension is for  $C_9^{\mu}$  but it's less than  $2\sigma$ 

## Conclusion

- Important to correctly choose the statistical method, depending on which question is asked
- There is a small tension of about  $2\sigma$ , in the global fits in the absence of lepton flavour violation
- We should be cautious not over interpreting the tension
- $\bullet$  To claim new physics, the use of  $\Delta\chi^2$  is NOT appropriate, one needs to use the absolute  $\chi^2$
- The ideal would be to consider properly the Look Elsewhere Effect
- The cross check with the updated results in particular for  $B_{\rm s} \to \phi \mu^+ \mu^-$  is awaited
- The cross check with the inclusive results is also of importance

# Backup

At Belle-II, for inclusive  $b \rightarrow s\ell\ell$ : expected uncertainty of 2.9% (4.1%) for the branching fraction in the low- (high-) $q^2$  region, absolute uncertainty of 0.050 in the low- $q^2$  bin 1 (1 <  $q^2$  < 3.5 GeV<sup>2</sup>), 0.054 in the low- $q^2$  bin 2 (3.5 <  $q^2$  < 6 GeV<sup>2</sup>) for the normalised  $A_{FB}$ 



T. Hurth, FM, JHEP 1404 (2014) 097 T. Hurth, FM, S. Neshatpour, JHEP 1412 (2014) 053

Predictions based on our model-independent analysis

black cross: future measurements at Belle-II assuming the best fit solution red cross: SM predictions

 $\rightarrow$  inclusive mode will lead to very strong constraints

Observable	SM prediction	Measurement
$\overline{10^4  imes  ext{BR}(B  o X_s \gamma)}$	$3.37\pm0.19$	$3.43\pm0.22$
$10^2  imes \Delta_0(B  o {\cal K}^* \gamma)$	$\textbf{6.9}\pm\textbf{3.0}$	$5.2\pm2.6$
$\overline{10^9  imes  ext{BR}(B_s  o \mu^+ \mu^-)}$	$3.54\pm0.27$	$2.9\pm0.7$
$10^{10} imes { m BR}(B_d o\mu^+\mu^-)$	$1.07\pm0.27$	$3.6\pm1.6$
$R_{K q^2 \in [1.0, 6.0](GeV)^2}$	$1.0006\pm0.0004$	$0.745\pm0.097$
$\overline{10^6 \times \mathrm{BR} \left( B \to X_s e^+ e^- \right)_{q^2 \in [1,6] (\mathrm{GeV})^2}}$	$1.73^{+0.12}_{-0.12}$	$1.93\pm0.55$
$10^6  imes { m BR} \left( B  o X_s e^+ e^-  ight)_{q^2 > 14.2 ({ m GeV})^2}$	$0.20\substack{+0.06\\-0.06}$	$0.56\pm0.19$
$\overline{10^6  imes \mathrm{BR} \left( B  ightarrow X_s \mu^+ \mu^-  ight)_{q^2 \in [1,6] (\mathrm{GeV})^2}}$	$1.67\substack{+0.12\\-0.12}$	$0.66\pm0.88$
$10^6  imes { m BR} \left( B  o X_s \mu^+ \mu^-  ight)_{q^2 > 14.2 ({ m GeV})^2}$	$0.23^{+0.07}_{-0.06}$	$0.60\pm0.31$

SM predictions and experimental values of the  $B^0 \to K^{*0} \mu^+ \mu^-$  observables

Observable	Soft FF (10%)	Soft FF (20%)	Full FF (5%)	Full FF (10%)	Measurement
		$q^2 \in [$	0.1, 0.98 GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.082 \pm 0.157$	$1.082 \pm 0.197$	$1.071\pm0.148$	$1.071 \pm 0.148$	
$\langle F_L \rangle$	$0.244 \pm 0.042$	$0.244 \pm 0.050$	$0.247 \pm 0.037$	$0.247 \pm 0.037$	$0.263^{+0.046}_{-0.044} \pm 0.017$
$\langle A_{FB} \rangle$	$-0.088 \pm 0.019$	$-0.088 \pm 0.036$	$-0.088 \pm 0.006$	$-0.088 \pm 0.008$	$-0.003^{+0.057}_{-0.059} \pm 0.008$
$\langle S_3 \rangle$	$0.000 \pm 0.011$	$0.000 \pm 0.023$	$0.007 \pm 0.002$	$0.007 \pm 0.003$	$-0.036^{+0.063}_{-0.063} \pm 0.005$
$\langle S_4 \rangle$	$-0.097 \pm 0.007$	$-0.097 \pm 0.010$	$-0.096 \pm 0.005$	$-0.096 \pm 0.005$	$0.082^{+0.070}_{-0.066} \pm 0.009$
$\langle S_5 \rangle$	$0.239 \pm 0.014$	$0.239 \pm 0.022$	$0.242\pm0.010$	$0.242\pm0.010$	$0.170^{+0.060}_{-0.059} \pm 0.018$
$\langle S_7 \rangle$	$0.022 \pm 0.014$	$0.022 \pm 0.026$	$0.022\pm0.006$	$0.022\pm0.006$	0.015 <sup>+0.059</sup> <sub>-0.057</sub> ± 0.006
$\langle S_8 \rangle$	$-0.004 \pm 0.006$	$-0.004 \pm 0.012$	$-0.004 \pm 0.003$	$-0.004 \pm 0.003$	0.079 <sup>+0.077</sup> <sub>-0.078</sub> ± 0.007
$\langle S_9 \rangle$	$-0.001 \pm 0.011$	$-0.001 \pm 0.023$	$-0.001 \pm 0.000$	$-0.001 \pm 0.001$	-0.083 <sup>+0.060</sup> -0.059 ± 0.004
$\langle P'_{5} \rangle$	$0.657 \pm 0.024$	$0.657 \pm 0.049$	$0.665\pm0.008$	$0.665\pm0.011$	$0.387^{+0.141}_{-0.131} \pm 0.052$
		$q^2 \in$	[1.1, 2.5] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$0.658 \pm 0.078$	$0.658 \pm 0.101$	$0.656 \pm 0.069$	$0.656 \pm 0.069$	
$\langle F_L \rangle$	$0.721 \pm 0.045$	$0.721 \pm 0.060$	$0.722 \pm 0.037$	$0.722 \pm 0.037$	$0.660^{+0.088}_{-0.075} \pm 0.022$
$\langle A_{FB} \rangle$	$-0.158 \pm 0.029$	$-0.158 \pm 0.038$	$-0.156 \pm 0.024$	$-0.156 \pm 0.024$	$-0.191^{+0.069}_{-0.078} \pm 0.012$
$\langle S_3 \rangle$	$0.000 \pm 0.008$	$0.000 \pm 0.016$	$0.003\pm0.001$	$0.003 \pm 0.002$	$-0.077^{+0.089}_{-0.104} \pm 0.005$
$\langle S_4 \rangle$	$-0.012 \pm 0.009$	$-0.012 \pm 0.009$	$-0.008 \pm 0.008$	$-0.008 \pm 0.009$	$-0.077^{+0.112}_{-0.112} \pm 0.005$
$\langle S_5 \rangle$	$0.106 \pm 0.015$	$0.106 \pm 0.017$	$0.108\pm0.015$	$0.108\pm0.015$	0.137 <sup>+0.094</sup> <sub>-0.098</sub> ± 0.009
$\langle S_7 \rangle$	$0.035 \pm 0.008$	$0.035 \pm 0.010$	$0.034\pm0.008$	$0.034\pm0.008$	$-0.219^{+0.093}_{-0.105} \pm 0.003$
$\langle S_8 \rangle$	$-0.012 \pm 0.004$	$-0.012 \pm 0.006$	$-0.011 \pm 0.004$	$-0.011 \pm 0.004$	$-0.098^{+0.107}_{-0.122} \pm 0.005$
$\langle S_9 \rangle$	$-0.001 \pm 0.008$	$-0.001 \pm 0.016$	$-0.001 \pm 0.001$	$-0.001 \pm 0.001$	$-0.119^{+0.087}_{-0.101} \pm 0.005$
$\langle P'_{5} \rangle$	$0.252 \pm 0.028$	$0.252 \pm 0.035$	$0.258\pm0.030$	$0.258 \pm 0.032$	$0.289^{+0.216}_{-0.200} \pm 0.023$
$q^2 \in [2.5, 4.0]  \text{GeV}^2$					
$\langle BR \rangle \times 10^7$	$0.637 \pm 0.081$	$0.637 \pm 0.117$	$0.637\pm0.065$	$0.637 \pm 0.065$	
$\langle F_L \rangle$	$0.808 \pm 0.036$	$0.808 \pm 0.056$	$0.807\pm0.028$	$0.807 \pm 0.028$	$0.877^{+0.089}_{-0.096} \pm 0.017$
$\langle A_{FB} \rangle$	$-0.053 \pm 0.017$	$-0.053 \pm 0.026$	$-0.051 \pm 0.011$	$-0.051 \pm 0.012$	$-0.118^{+0.075}_{-0.088} \pm 0.007$
$\langle S_3 \rangle$	$-0.011 \pm 0.008$	$-0.011 \pm 0.014$	$-0.010 \pm 0.003$	$-0.010 \pm 0.003$	$0.035^{+0.101}_{-0.086} \pm 0.006$
$\langle S_4 \rangle$	$0.124 \pm 0.016$	$0.124 \pm 0.022$	$0.127\pm0.013$	$0.127 \pm 0.013$	$-0.234^{+0.132}_{-0.144} \pm 0.006$
$\langle S_{5} \rangle$	$-0.146 \pm 0.021$	$-0.146 \pm 0.031$	$-0.144 \pm 0.017$	$-0.144 \pm 0.017$	$-0.022^{+0.110}_{-0.104} \pm 0.008$
$\langle S_7 \rangle$	$0.026 \pm 0.024$	$0.026 \pm 0.047$	$0.026 \pm 0.006$	$0.026 \pm 0.006$	$0.068^{+0.119}_{-0.112} \pm 0.005$
$\langle S_8 \rangle$	$-0.011 \pm 0.009$	$-0.011 \pm 0.017$	$-0.010 \pm 0.003$	$-0.010 \pm 0.003$	$0.030^{+0.123}_{-0.127} \pm 0.006$
$\langle S_9 \rangle$	$-0.001 \pm 0.007$	$-0.001 \pm 0.013$	$-0.001 \pm 0.000$	$-0.001 \pm 0.001$	$-0.092^{+0.108}_{-0.125} \pm 0.007$
$\langle P'_{5} \rangle$	$-0.386 \pm 0.050$	$-0.386 \pm 0.077$	$-0.382 \pm 0.037$	$-0.382 \pm 0.039$	$-0.066^{+0.341}_{-0.360} \pm 0.023$

Nazila Mahmoudi

Rare B decays in 2015 - Edinburgh - 13 May 2015

SM predictions and experimental values of the  $B^0 \to K^{*0} \mu^+ \mu^-$  observables

Observable	Soft FF (10%)	Soft FF (20%)	Full FF (5%)	Full FF (10%)	Measurement
		$q^2 \in$	[6.0, 8.0] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.059 \pm 0.105$	$1.059 \pm 0.177$	$1.065 \pm 0.065$	$1.065 \pm 0.065$	
$\langle F_L \rangle$	$0.625 \pm 0.073$	$0.625 \pm 0.126$	$0.624\pm0.041$	$0.624 \pm 0.041$	$0.579^{+0.043}_{-0.047} \pm 0.015$
$\langle A_{FB} \rangle$	$0.228 \pm 0.049$	$0.228 \pm 0.083$	$0.230 \pm 0.026$	$0.230 \pm 0.026$	$0.152^{+0.040}_{-0.040} \pm 0.008$
$\langle S_3 \rangle$	$-0.044 \pm 0.029$	$-0.044 \pm 0.055$	$-0.045 \pm 0.011$	$-0.045 \pm 0.011$	$-0.042^{+0.057}_{-0.058} \pm 0.011$
$\langle S_4 \rangle$	$0.260 \pm 0.021$	$0.260 \pm 0.039$	$0.262 \pm 0.009$	$0.262 \pm 0.009$	$-0.296^{+0.065}_{-0.065} \pm 0.011$
$\langle S_5 \rangle$	$-0.393 \pm 0.041$	$-0.393 \pm 0.077$	$-0.391 \pm 0.013$	$-0.391 \pm 0.013$	$-0.249^{+0.062}_{-0.061} \pm 0.012$
$\langle S_7 \rangle$	$0.010 \pm 0.079$	$0.010 \pm 0.149$	$0.009\pm0.003$	$0.009\pm0.004$	-0.047 <sup>+0.066</sup> <sub>-0.062</sub> ± 0.003
$\langle S_8 \rangle$	$-0.005 \pm 0.031$	$-0.005 \pm 0.060$	$-0.005 \pm 0.002$	$-0.005 \pm 0.002$	$-0.085^{+0.072}_{-0.073} \pm 0.006$
$\langle S_9 \rangle$	$-0.001 \pm 0.026$	$-0.001 \pm 0.052$	$-0.001 \pm 0.001$	$-0.001 \pm 0.002$	-0.024 <sup>+0.059</sup> -0.062 ± 0.005
$\langle P'_{5} \rangle$	$-0.819 \pm 0.083$	$-0.819 \pm 0.160$	$-0.814 \pm 0.025$	$-0.814 \pm 0.025$	$-0.505^{+0.118}_{-0.177} \pm 0.024$
		$q^2 \in [$	15.0, 17.0] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.258 \pm 0.073$	$1.258 \pm 0.092$	$1.258 \pm 0.068$	$1.258 \pm 0.073$	
$\langle F_L \rangle$	$0.339 \pm 0.039$	$0.339 \pm 0.055$	$0.339\pm0.034$	$0.339\pm0.039$	0.349 <sup>+0.040</sup> <sub>-0.039</sub> ± 0.009
$\langle A_{FB} \rangle$	$0.409 \pm 0.025$	$0.409 \pm 0.037$	$0.409\pm0.022$	$0.409 \pm 0.026$	$0.411^{+0.040}_{-0.035} \pm 0.008$
$\langle S_3 \rangle$	$-0.181 \pm 0.024$	$-0.181 \pm 0.037$	$-0.181 \pm 0.020$	$-0.181 \pm 0.024$	$-0.142^{+0.046}_{-0.047} \pm 0.007$
$\langle S_4 \rangle$	$0.294 \pm 0.008$	$0.294 \pm 0.013$	$0.294 \pm 0.007$	$0.294 \pm 0.008$	$-0.321^{+0.053}_{-0.078} \pm 0.007$
$\langle S_5 \rangle$	$-0.315 \pm 0.024$	$-0.315 \pm 0.037$	$-0.315 \pm 0.019$	$-0.315 \pm 0.024$	$-0.316^{+0.051}_{-0.058} \pm 0.009$
$\langle S_7 \rangle$	$0.000 \pm 0.034$	$0.000 \pm 0.067$	$0.000\pm0.017$	$0.000 \pm 0.034$	$0.061^{+0.058}_{-0.060} \pm 0.005$
$\langle S_8 \rangle$	$0.000 \pm 0.009$	$0.000 \pm 0.018$	$0.000\pm0.005$	$0.000\pm0.009$	$0.003^{+0.060}_{-0.060} \pm 0.003$
$\langle S_9 \rangle$	$0.000 \pm 0.016$	$0.000 \pm 0.032$	$0.000\pm0.008$	$0.000\pm0.016$	$-0.019^{+0.055}_{-0.057} \pm 0.004$
$\langle P'_{5} \rangle$	$-0.666 \pm 0.041$	$-0.666 \pm 0.065$	$-0.666 \pm 0.033$	$-0.666 \pm 0.042$	$-0.662^{+0.112}_{-0.126} \pm 0.017$
$q^2 \in [17.0, 19.0] \mathrm{GeV}^2$					
$\langle BR \rangle \times 10^7$	$0.866 \pm 0.055$	$0.866 \pm 0.069$	$0.866\pm0.051$	$0.866 \pm 0.054$	
$\langle F_L \rangle$	$0.322 \pm 0.042$	$0.322 \pm 0.057$	$0.322\pm0.037$	$0.322 \pm 0.042$	0.354 <sup>+0.048</sup> <sub>-0.048</sub> ± 0.025
$\langle A_{FB} \rangle$	$0.321 \pm 0.023$	$0.321 \pm 0.033$	$0.321\pm0.021$	$0.321 \pm 0.024$	$0.305^{+0.048}_{-0.046} \pm 0.013$
$\langle S_3 \rangle$	$-0.256 \pm 0.025$	$-0.256 \pm 0.034$	$-0.256 \pm 0.021$	$-0.256 \pm 0.024$	$-0.188^{+0.076}_{-0.086} \pm 0.017$
$\langle S_4 \rangle$	$0.309 \pm 0.010$	$0.309 \pm 0.014$	$0.309\pm0.009$	$0.309 \pm 0.010$	$-0.266^{+0.065}_{-0.071} \pm 0.010$
$\langle S_{5} \rangle$	$-0.224 \pm 0.022$	$-0.224 \pm 0.032$	$-0.224 \pm 0.019$	$-0.224 \pm 0.022$	-0.323 <sup>+0.062</sup> <sub>-0.069</sub> ± 0.009
$\langle S_7 \rangle$	$0.000 \pm 0.035$	$0.000 \pm 0.071$	$0.000\pm0.018$	$0.000 \pm 0.035$	$0.044^{+0.072}_{-0.073} \pm 0.013$
$\langle S_8 \rangle$	$0.000 \pm 0.007$	$0.000 \pm 0.013$	$0.000 \pm 0.003$	$0.000 \pm 0.007$	$0.013^{+0.067}_{-0.071} \pm 0.005$
$\langle S_9 \rangle$	$0.000 \pm 0.013$	$0.000 \pm 0.025$	$0.000\pm0.006$	$0.000 \pm 0.013$	-0.094 <sup>+0.067</sup> <sub>-0.069</sub> ± 0.004
$\langle P'_{5} \rangle$	$-0.481 \pm 0.039$	$-0.481 \pm 0.057$	$-0.481 \pm 0.033$	$-0.481 \pm 0.039$	$-0.675^{+0.138}_{-0.152} \pm 0.017$

Nazila Mahmoudi

Rare B decays in 2015 - Edinburgh - 13 May 2015

$B_{s}  o \phi \ \mu^{+} \mu^{-}$ SM prediction					
Observable	Soft FF (10%)	Soft FF (20%)	Full FF (5%)	Full FF (10%)	Measurement
		$q^2 \in [0.1]$	, 2.0] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.631 \pm 0.134$	$1.631 \pm 0.161$	$1.611 \pm 0.095$	1.611 ± 0.095	$0.90^{+0.21}_{-0.19} \pm 0.04 \pm 0.09$
$\langle F_L \rangle$	$0.390 \pm 0.043$	$0.390\pm0.058$	$0.397 \pm 0.034$	0.397 ± 0.035	$0.37^{+0.19}_{-0.17} \pm 0.07$
$\langle S_3 \rangle$	$-0.001 \pm 0.010$	$-0.001 \pm 0.020$	$0.006\pm0.002$	$0.006 \pm 0.003$	$-0.11^{+0.28}_{-0.25} \pm 0.05$
		<i>q</i> <sup>2</sup> ∈ [2.0	, 4.3] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.013 \pm 0.072$	$1.013 \pm 0.112$	$1.017 \pm 0.053$	1.017 ± 0.054	$0.53^{+0.18}_{-0.16} \pm 0.03 \pm 0.05$
$\langle F_L \rangle$	$0.802 \pm 0.032$	$0.802 \pm 0.053$	$0.803\pm0.020$	0.803 ± 0.020	$0.53^{+0.25}_{-0.23} \pm 0.10$
$\langle S_3 \rangle$	$-0.012 \pm 0.007$	$-0.012 \pm 0.015$	$-0.011 \pm 0.003$	$-0.011 \pm 0.003$	$-0.97^{+0.53}_{-0.03} \pm 0.17$
		<i>q</i> <sup>2</sup> ∈ [4.30	, 8.68]GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$2.284 \pm 0.095$	$2.284 \pm 0.168$	$2.306\pm0.058$	2.306 ± 0.059	$1.38^{+0.25}_{-0.23} \pm 0.05 \pm 0.14$
$\langle F_L \rangle$	$0.651 \pm 0.063$	$0.651 \pm 0.116$	$0.650 \pm 0.029$	0.650 ± 0.029	$0.81^{+0.11}_{-0.13} \pm 0.05$
$\langle S_3 \rangle$	$-0.046 \pm 0.025$	$-0.046 \pm 0.049$	$-0.048 \pm 0.010$	$-0.048 \pm 0.010$	$0.25^{+0.21}_{-0.24} \pm 0.05$
_		$q^2 \in [14.18]$	, 16.0] GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.167 \pm 0.072$	$1.167 \pm 0.092$	$1.167 \pm 0.066$	1.167 ± 0.073	$0.76^{+0.19}_{-0.17} \pm 0.04 \pm 0.08$
$\langle F_L \rangle$	$0.349 \pm 0.036$	$0.349 \pm 0.054$	$0.349 \pm 0.030$	0.349 ± 0.036	$0.34^{+0.18}_{-0.17} \pm 0.07$
(S <sub>3</sub> )	$-0.172 \pm 0.022$	$-0.172 \pm 0.036$	$-0.172 \pm 0.017$	$-0.172 \pm 0.022$	$-0.03^{+0.29}_{-0.31} \pm 0.06$
		$q^2 \in [16.0]$	, 19.0]GeV <sup>2</sup>		
$\langle BR \rangle \times 10^7$	$1.280 \pm 0.053$	$1.280 \pm 0.068$	$1.280 \pm 0.049$	$1.280 \pm 0.054$	$1.06^{+0.23}_{-0.21} \pm 0.06 \pm 0.11$
$\langle F_L \rangle$	$0.325 \pm 0.039$	$0.325 \pm 0.056$	$0.325 \pm 0.033$	$0.325 \pm 0.039$	$0.16^{+0.17}_{-0.10} \pm 0.07$
(S <sub>3</sub> )	$-0.248 \pm 0.022$	$-0.248 \pm 0.034$	$-0.248 \pm 0.018$	$-0.248 \pm 0.022$	$0.19^{+0.30}_{-0.31} \pm 0.05$
$BR(B \rightarrow K\mu^+\mu^-) SM$ prediction					
bin	Soft FF (10%)	Soft FF (20%)	Full FF (5%)	Full FF (10%)	Measurement
$10^7  imes \langle BR  angle (B^0  o K^0 \mu^+ \mu^-)$					
$q^2 \in [1.1 - 6.0] \text{ GeV}^2$	$1.353 \pm 0.061$	$1.353\pm0.100$	$1.350 \pm 0.045$	1.350 ± 0.045	$0.92^{+0.17}_{-0.16} \pm 0.04$
$q^2 \in [15.0 - 22.0]  \mathrm{GeV}^2$	$0.942 \pm 0.014$	$0.942\pm0.015$	$0.942\pm0.014$	$0.942 \pm 0.014$	$0.67^{+0.11}_{-0.11} \pm 0.04$
$10^7 \times \langle BR \rangle (B^+ \to K^+ \mu^+ \mu^-)$					
$q^2 \in [1.1 - 6.0] \text{ GeV}^2$	$1.481 \pm 0.067$	$1.481 \pm 0.110$	$1.477 \pm 0.049$	$1.477 \pm 0.049$	$1.19 \pm 0.03 \pm 0.06$
$q^2 \in [15.0 - 22.0]  \mathrm{GeV}^2$	$1.024 \pm 0.016$	$1.024\pm0.016$	$1.024\pm0.016$	$1.024\pm0.016$	$0.85 \pm 0.03 \pm 0.04$