# Upgrading the water Cherenkov tanks for atmospheric shower identification

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- UHE cosmic rays (> 10<sup>16</sup> eV) are observed through atmospheric showers (longitudinal profile + ground particles):
  - spectrum is well established, but sources are still unknown
- a key point to understand the origin: identifying the primary (light/heavy nucleus, photon, neutrino ?)
- main criteria to discriminate the primaries:
  - depth of maximum + muonic vs electromagnetic component
- Cherenkov water tanks to sample the ground particles: sensitive to both; indirect and model dependent separation
- how to improve the separation ?

essentially: use a 2-fold detector with different sensitivities internal modification of the tank or addition of another detector

# atmospheric cascade



#### hadronic interactions:

large multiplicity at high E in average ~1/3 of E  $\rightarrow$  neutral mesons (mainly  $\pi^0 \rightarrow \gamma\gamma$ ) at the end: few GeV mesons decay  $\rightarrow$  muons (some of them propagate down to ground)

electromagnetic cascade:  $\gamma \rightarrow e^+e^-$  (pair production)  $e^{\pm} \rightarrow e^{\pm} \gamma$  (bremsstrahlung)

- in air (index n) charged particles produce Cherenkov light if  $\beta n > 1$  (tightly collimated) - excitation of nitrogen molecules: isotropic **fluorescence** light (isotropic)

- emission of radio waves



## $X_{max}$ and muons to discriminate primaries

for a given primary energy: heavier nucleus  $\rightarrow$  less steps in the hadronic cascade  $\rightarrow$ 

- faster longitudinal development: lower X<sub>max</sub>
- less energy going to e.m. cascade: more muons



### response of a water tank



(far from the shower core)

#### FADC traces from a shower (AUGER tank)



# the spectrum of UHECR (ICRC 2015)



clear « broken line » structure:

- ankle at ~ 18.5
- cutoff at ~ 19.6

ambiguous interpretation: the energy of cutoff matches the GZK effect (interaction of protons with CMB) but it could be the upper limit of sources (heavy nuclei favoured)

X<sub>max</sub> results (ICRC 2015)



rough agreement (but TA uncertainties still large) :

- change of slope at the « ankle » where light nuclei dominate
- trend to heavy nuclei at highest energies
- different models do not give strongly different predictions
- not enough data above the cutoff

 $X_{max} vs \sigma(X_{max})$  (Auger)



- consistent with « pure proton » around the ankle
- higher energy is more difficult to interpret: decreasing  $\sigma(X_{max})$  is expected with increasing A, but not well compatible with a mixture of different nuclei

### muonic component (Auger, ICRC2015)

using « inclined » events (above 60 deg) to select a pure muonic shower, compare the muon density to the expectation from fluorescence energy. *reference : E = 10^{19} eV, hadronic model QGSJET II-03* 



none of the models is compatible with the data (the *muon deficit* problem)

similar problems found using other muon signatures (e.g. the *muon production depth* deduced from the time of arrival, or the structure of FADC traces in tanks)

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# "muonic component" is ambiguous

#### muon counting methods explore a limited part of the phase space

- lower cut on momentum (especially for underground detector)
- lower cut on distance to core (FADC trace structure, production depth from arrival time)

#### dependence on muonic component model

e.g.: an error on angular distribution (very steep) results in a large error in lateral density (Cherenkov tanks do not see the direction of incidence)

#### $\rightarrow$ need for additional information on the muons

especially at short distance from the core: big signals, but difficult to disentangle the e.m. and muonic components

## how to improve the discrimination?

1- « black top » option: faster absorption of light → less photo-electrons in PMTs, but tighter muon peaks in the FADC traces (easier to count)

- 2- coupling the tank to a charged particle detector (in practice: scintillator)
  - just above (counting electrons+muons)
  - just below (counting muons+ few through going electrons)
  - underground (counting hard muons)
- 3- splitting the tank in two parts (the « Layered Surface Detector »)

for 2 and 3: observed signals  $S_1 = A_1 \Phi_{\mu} + B_1 \Phi_{em}$ ,  $S_2 = A_2 \Phi_{\mu} + B_2 \Phi_{em}$  $\Rightarrow \Phi_{\mu}$  and  $\Phi_{em}$  through a linear system (good if  $A_1/A_2$  significantly >  $B_1/B_2$ )

4- using an external source of information (e.g. geomagnetic distortion of the muonic component)

more or less details in the following

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# AMIGA in Auger (JINST 11 PO2012)



- 2.25 m under ground surface → lower cut for muons at 1 GeV or more (depending on incidence angle); e.m. component strongly suppressed
- muon samples in tank and scintillator are decorrelated (quadratic addition of fluctuations from Poisson statistics)

# AugerPrime (ICRC 2015)



*if scintillator* (SSD) ≈ *counter of charged particles* 

- 5 muons in SSD (no cut on momentum)  $\approx$  4 vertical muons in tank
- 1 GeV of e.m. energy (eq. to 4 vertical muons in tank) contains 6 to 8 e<sup>+</sup>/e<sup>-</sup>
- → the scintillator is relatively more sensitive to the e.m. component than the tank

muon samples in scintillator (4 m<sup>2</sup>) and tank (10 m<sup>2</sup>) are partly correlated

➔ partial compensation of Poisson errors

### Layered Surface Detector

#### (A. Letessier-Selvon, P.B., M. Blanco, I. Maris, M. Settimo, NIM A 767 (2014) 41)



#### proposed design (internal modification of Auger tank)



- electromagnetic signal : ≈ 60 % in top, 40 % in bottom (almost independent of incidence)
- muonic signal: 40 % in top, 60 % in bottom

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### results from the prototype



#### geomagnetic distortion (P.B., M. Settimo, M.Blanco, Astropart. Phys. 74 (2016) 14)



very inclined showers (purely muonic): density in transverse plane is more or less distorted

empirical parametrization: radial( $\rho$ ) × angular( $\phi$ )  $\rho = (r/r_{ref})^{1/2} - 1$ f( $\rho$ , $\phi$ ) = exp( $\lambda \rho$  +  $\alpha$ .cos(2 $\phi$ ))  $\lambda = \lambda_0 + \lambda_1 \rho + \lambda_2 \rho^2$  $\alpha = \alpha_0 + \alpha_1 \rho + \alpha_2 \rho^2$ 

the  $\alpha_i$  and  $\lambda_i$  carry information about the longitudinal muon profile e.g. defining  $X^{\mu}_{max}$ , tightly correlated to  $X_{max}$ 

### exploiting the distortion

 $\alpha$ ,  $\lambda$  parameters as functions of  $X^{\mu}_{max}$  with *different models* (here: zenith angle 72 deg,  $B_T = 30 \mu T$ , proton and iron showers)



- for a given model, different nuclei are on the same line:  $\alpha$ ,  $\lambda$  provide *an identification of primary* with a consistency check between parameters
- if X<sub>max</sub> is measured independently (e.g. through fluorescence): *better model discrimination*

ideally: dedicated hybrid detector of inclined showers (muons at ground + profile) to measure ( $\alpha$ ,  $\lambda$ ) and ( $X_{max}$ ,  $N_{max}$ ) on the same events

### a possible detector layout





spacing ~ 500 m or less to ensure a large multiplicity at 1 EeV (good precision on  $\alpha$ ,  $\lambda$  needed) may be single fluorescence eyes with a large field of view, as proposed in arXiv:1504.00692

hybrid events are used to calibrate the relation between  $\alpha$ ,  $\lambda$  and  $X_{max}$  and provide a discrimination between the models of hadronic interactions

### summary

Cherenkov tanks are proven to be good and robust detectors of ground particles produced by extensive atmospheric showers. especially when associated to a longitudinal profile detector (fluorescence, MHz radio) but there are still open issues about the nature and the origin of the ultra energetic cosmic rays.

To go further, a better identification of primaries is needed, mainly through a better separation of muonic/electromagnetic components **generic idea:** make two measurements with different *relative* mu/em sensitivity (additional detector or layered tank)

Unavoidable problem: systematic errors due to the modelling of the hadronic interactions at ultra high energies possible constraints from complementary observations (e.g. geomagnetic distortion) ?