NEUTRINOLESS DOUBLE BETA DECAY AND PHYSICS BEYOND THE STANDARD MODEL



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Static and Dynamic of Massive Properties

- (1) Smallness of neutrino mass
- (2) Structure of mixing matrix
- (3) Nature of neutrinos
- (4) Violation of B-L symmetry
- (5) Stability of neutrinos
- (6) Possible CP-violation in leptonic sector
- (7) Interaction of neutrinos

Neutrino Mass and Mixing

- (1) Extended Standard model (enlarged fermion sector, enlarged Higgs sector, Majoron modelsspontaneous B-L violation)
- (2) Left-right Symmetric model (see-saw)
- (3) SU(5), SO(10) and E(6) GUTs
- (4) Supersymmetric models
- (5) Extradimensions and String theory

Mixing: Democratic, Maximal and Bimaximal

 $m_{v}M_{N} \approx m_{D}^{2}$

NEUTRINO MASS	nplementary Experiments		
NEUTRINO MASS (1) Beta Decay $m_{\nu} = \sqrt{\sum_{k=1}^{3} U_{ek}^{(11)} ^2} m_k^2$ $= \sqrt{c_{12}^2 c_{13}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2}$ (2) Astrophysical and Cosmological $m_{\nu} = \sum_{k=1}^{3} m_k$ (3) NDBD $\langle m_{\nu} \rangle = \sum_{k=1}^{3} (U_{ek}^{(11)})^2 m_k$	Con	 Neutrino oscillation 0νββ decay Single β decay Astrophysical and Cosmological measurements Role of neutrinos on large scale structure formation 	
$=c_{12}^2c_{13}^2e^{2i\alpha_1}m_1^2+c_{13}^2s_{12}^2e^{2i\alpha_2}m_2^2+s$	$^{2}_{13}m^{2}_{3}$		

Perspective of ββ decay

Candidates: 35 β - β -, 34 e⁺ $\beta\beta$ emitters

Modes: $\beta^{-}\beta^{-}$, $\beta^{+}\beta^{+}$, $\epsilon\beta^{+}$, $\epsilon\epsilon$ with and without neutrinos,

single Majoron and double Majorons

Transitions: $0^+ \rightarrow 0^+, 0^+ \rightarrow 1^+, 0^+ \rightarrow 2^+$ etc.

Experimental Study: Decay rates, energy spectrum, angular correlation

Implications of ββ decay

-Lepton number violation

- -Dirac or Majorona neutrino
- -Origin of neutrino mass
- -Mass hierarchy
- -Right handed current
- -CP violation in lepton sector

Astrophysical: Stellar evolution and Supernova explosion

Cosmological: Big-bang nucleosynthesis, Galaxy formation, dark matter and dark energy







Gauge theoretical models

- 1 Left-right Symmetric models
- 2. Sterile neutrinos
- 3. Majoron models
- 4. SUSY models
- 5. Lepto-quark models
- 6. Compositeness
- 7. Extra-dimensional Scenario

Theoretical Approaches

(1) Phenomenological

(2) Quark model

(3) EFT

(4) Chiral EFT

(5) Lorentz invariant description



$$(\beta\beta)_{0\nu}$$
 Decay (LRSM)

$$\begin{bmatrix} T_{1/2}^{0\nu}(0^{+} \to 0^{+}) \end{bmatrix}^{-1} = \left(\frac{\langle m_{\nu} \rangle}{m_{e}} \right)^{2} C_{mm}^{LL} + \left(\frac{M_{p}}{\langle M_{N} \rangle} \right)^{2} C_{mm}^{NN} + \left(\frac{\langle m_{\nu} \rangle}{m_{e}} \right) \left(\frac{M_{p}}{\langle M_{N} \rangle} \right) C_{mm}^{NL} \\ + \left(\frac{\langle m_{\nu} \rangle}{m_{e}} \right) \langle \lambda \rangle C_{m\lambda}^{LL} + \left(\frac{\langle m_{\nu} \rangle}{m_{e}} \right) \langle \eta \rangle C_{m\eta}^{LL} + \left(\frac{\langle m_{\nu} \rangle}{m_{e}} \right) \langle \xi \rangle C_{m\xi}^{NL} \\ + \left(\frac{M_{p}}{\langle M_{N} \rangle} \right) \langle \lambda \rangle C_{m\lambda}^{NL} + \left(\frac{M_{p}}{\langle M_{N} \rangle} \right) \langle \eta \rangle C_{m\eta}^{NL} + \left(\frac{M_{p}}{\langle M_{N} \rangle} \right) \langle \xi \rangle C_{m\xi}^{NN} \\ + \langle \lambda \rangle \langle \eta \rangle C_{\lambda\eta}^{LL} + \langle \lambda \rangle \langle \xi \rangle C_{\lambda\xi}^{NL} + \langle \eta \rangle \langle \xi \rangle C_{\eta\xi}^{NL} \\ + \langle \lambda \rangle^{2} C_{\lambda\lambda}^{LL} + \langle \eta \rangle^{2} C_{\eta\eta}^{LL} + \langle \xi \rangle^{2} C_{\xi\xi}^{NN}$$

$$(4.2)$$

neutrinoless $\beta\beta e^{i}$

$$\langle m_{\nu} \rangle = \sum_{i}^{\prime} U_{ei}^{2} m_{i} \tag{4.3}$$

$$\langle \lambda \rangle = \lambda \left| \sum_{i}^{\prime} \left(\frac{g_{V}}{g_{V}} \right) U_{ei} V_{ei} \right|$$

$$(4.4)$$

$$\langle \eta \rangle = \eta \left| \sum_{i}^{\prime} U_{ei} V_{ei} \right| \tag{4.5}$$

$$\langle M_N \rangle^{-1} = \sum_{i}^{\prime \prime} U_{ei}^2 m_i^{-1}$$
 (4.6)

$$\langle \xi \rangle = \left[\lambda^2 + \eta^2 - 2\lambda \eta \left(\frac{M_{GTh} + M_{Fh}}{M_{GTh} - M_{Fh}} \right) \right] \sum_{i}^{\prime\prime} V_{ei}^2 \left(\frac{M_p}{m_j} \right)$$
(4.7)



Majoron models

Half-life $T_{1/2}$ of Majoron emitting $0\nu\,\beta\beta$ decay is given by

$$[T_{1/2}]^{-1} = \left| \left\langle g_{\alpha} \right\rangle \right|^{m} \cdot \left| M_{\alpha} \right|^{2} \cdot G_{\beta\beta\alpha}$$
$$m = \begin{cases} 2 & 0\nu\beta\beta\phi \\ 4 & 0\nu\beta\beta\phi\phi \end{cases}$$

where





$$\begin{aligned} (\beta\beta)_{0\nu} \text{ Decay} \\ [T_{1/2}^{0\nu}(0^+ \to 0^+)]^{-1} &= G_{01} \left| \sum_{\alpha} \eta_{\alpha} M^{(\alpha)} \right|^2 \left| A_L \right| &= G_F^2 \frac{M_{ee}^{\nu}}{k^2} \\ where \\ M^{(\alpha)} &= M_F^{(\alpha)} + M_{GT}^{(\alpha)} + M_T^{(\alpha)} \\ and \\ \alpha &= \text{differentmodes} \\ \text{Opearators are long and short ranged.} \end{aligned} \qquad \begin{aligned} A_H &= G_F^2 \frac{M_W^4}{\Lambda^5} \\ A_H &= G_F^2 \frac{M_W^4}{\Lambda^5} \\ A_H &= 1 \\ \text{for } M_{ee}^{\nu} \approx 0.1 - 0.5 \text{ eV} \\ \text{and } \Lambda &= 1 \text{ TeV} \end{aligned}$$

Sources of	f Uncertainties	Reliability of NTMEs
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- (1) Structure of NTMEs
- (nonperturbative QCD)
- (Simkovic et al., Vergados, Faessler, Civitarese and Suhonen, Mendez)
- (2) Assumptions of NRMBTH
- (3) Model Specific
- (4) g_A
- (6) **FNS**
- (7) **SRC**

Models:

SM, QRPA, DQRPA, PHFB, EDF CDFT, IBM

Calculation of as many physical observables as possible --PRC 41, 1315 (1990)

- (1) Yrast spectra
- (2) Occupation Numbers
- (3) Reduced B(E2:0⁺ \rightarrow 2⁺) transition
- (4) Static quadrupole moments Q(2⁺)
- (5) g-factors g(2⁺)
- (6) M_{2v}

FNS

(1) Dipole form factor

- (2) Structure of nucleons
- (3)

(4)

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SRC

- (1) Jastrow correlations (Miller-Spencer, CCM)
- (2) UCOM
- (3) CCM (Argonne and CD-Bonn)
- (4) FhCh
- (5) Meson exchange

Uncertainties in NTMEs

- (1) Spread (Vogel)
- (2) Global average (Bahcall, Avignone III)
- (3) Model independent (Bilenky and Grifols)
- (4) Model specific (Tuebingen, Jyväskylä)
- (5) Engel's proposal

Deformation Effects

- (1) NTMEs are large for a pair of spherical nuclei.
- (2) With small admixture of quadrupolar correlations, NTMEs are almost constant and suppressed in realistic situations.
- (3) NTMEs are large for identical deformations of parent and daughter nuclei.
- (4) Sizes of NTMEs are reduced with the increase in deformations of parent and daughter nuclei.
- (5) The deformation effects are equally important for $(\beta\beta)_{2\nu}$ and $(\beta\beta)_{0\nu}$ modes so far as nuclear structure aspect of $\beta\beta$ decay is concerned.

EPJA 23, 223 (2005); EPJA 28, 27 (2006); EPJA 33, 375 (2007),

PRC 78, 054302 (2008); PRC 80, 044303 (2009), EPL 86, 32001 (2009),

JPG 37, 055108 (2010)

Calculated [5] and experimental deformation parameter β_2 [6].							
Nuclei	Case	$egin{array}{c} \beta_2 \ (Th.) \end{array}$	$\begin{array}{c} \beta_2 \\ (Exp.) \end{array}$	Nuclei	Case	β_2 (<i>Th</i> .)	$\begin{array}{c} \beta_2 \\ (Exp.) \end{array}$
94 Zr	(a)	0.100	$0.090 {\pm} 0.010$	⁹⁴ Mo	(a)	0.161	$0.1509 {\pm} 0.0015$
	(b)	0.110			(b)	0.161	
96 Zr	(a)	0.085	$0.080 {\pm} 0.017$	^{96}Mo	(a)	0.191	$0.1720{\pm}0.0016$
	(b)	0.087			(b)	0.186	
^{98}Mo	(a)	0.158	$0.1683 {\pm} 0.0028$	98 Ru	(a)	0.205	$0.1947 {\pm} 0.0030$
	(b)	0.162			(b)	0.194	
^{100}Mo	(a)	0.231	$0.2309 {\pm} 0.0022$	$^{100}\mathrm{Ru}$	(a)	0.214	$0.2148 {\pm} 0.0011$
	(b)	0.226			(b)	0.214	
104 Ru	(a)	0.285	$0.2707 {\pm} 0.0020$	$^{104}\mathrm{Pd}$	(a)	0.216	$0.209 {\pm} 0.007$
	(b)	0.282			(b)	0.219	
$^{110}\mathrm{Pd}$	(a)	0.216	$0.257 {\pm} 0.006$	$^{110}\mathrm{Cd}$	(a)	0.196	$0.1770{\pm}0.0039$
	(b)	0.214			(b)	0.191	
$^{128}\mathrm{Te}$	(a)	0.136	$0.1363 {\pm} 0.0011$	128 Xe	(a)	0.192	$0.1836{\pm}0.0049$
	(b)	0.136			(b)	0.181	
$^{130}\mathrm{Te}$	(a)	0.117	$0.1184{\pm}0.0014$	130 Xe	(a)	0.166	$0.169 {\pm} 0.007$
	(b)	0.120			(b)	0.163	
$^{150}\mathrm{Nd}$	(a)	0.276	$0.2853 {\pm} 0.0021$	^{150}Sm	(a)	0.238	$0.1931 {\pm} 0.0021$
	(b)	0.279			(b)	0.241	

Uncertainties in NTMEs

NTMEs and uncetainties of short and long ranged operators.

Nuclei	⁹⁶ Zr	¹⁰⁰ Mo	$^{128}\mathrm{Te}$	$^{130}\mathrm{Te}$	¹⁵⁰ Nd	Error(%)
$M^{(0\nu)}$	$2.86{\pm}0.26$	$6.25 {\pm} 0.64$	$3.61 {\pm} 0.40$	$4.05 {\pm} 0.50$	$2.83 {\pm} 0.43$	9 - 15
$M_N^{(\tilde{q}\nu)}(bag)$	34.2 ± 12.8	78.2 ± 25.1	46.7 ± 16.9	56.3 ± 16.4	$31.4{\pm}11.0$	32 - 37
$M_N^{(\tilde{q}\nu)}(nrqm)$	81.5 ± 24.9	174.0 ± 49.3	104.4 ± 32.3	117.0 ± 31.8	70.5 ± 21.7	22 - 49
$M_{\pi}^{(\widetilde{q}\nu)}$	219.5 ± 20.1	486.9 ± 49.9	305.4 ± 30.7	$318.4 {\pm} 30.1$	213.5 ± 31.6	9 - 15
$M_{CL}^{(\chi)}$	$0.60 {\pm} 0.06$	$1.21 {\pm} 0.12$	$0.26{\pm}0.03$	$0.69{\pm}0.08$	$0.63 {\pm} 0.10$	9 - 15
$M_{CR}^{(\chi)} \times 10^{3}$	1.64 ± 0.15	$3.02 {\pm} 0.29$	$0.23{\pm}0.02$	1.47 ± 0.17	1.71 ± 0.25	9 - 15
$M_{w2}^{(\chi)} \times 10^4$	$0.74 {\pm} 0.06$	1.25 ± 0.14	$0.10 {\pm} 0.01$	0.65 ± 0.13	0.77 ± 0.1	7 - 21
$M_{w2}^{(\chi)} \times 10^4$	$2.52 {\pm} 0.19$	$3.60 {\pm} 0.40$	$0.03 {\pm} 0.004$	$1.38 \pm \ 0.29$	2.77 ± 0.45	7 - 21
$M^{(0N)}$	100.53 ± 36.89	$206.75\ {\pm}73.08$	126.83 ± 46.34	136.39 ± 46.92	85.55 ± 31.45	34 - 37

Nuclei	PHFB	$\operatorname{Error}(\%)$	Global	$\operatorname{Error}(\%)$	$T_{1/2}^{(0 u)}(yr)$	$\langle m_{\nu} \rangle$	$T^{(0 u)}_{1/2}(yr) \ \langle m_ u angle = 0.01 \ { m e}$
^{48}Ca			1.74 ± 0.81	46.44	5.8×10^{22}	4.80	1.33×10^{28}
76 Ge			5.13 ± 1.17	22.84	3.0×10^{25}	0.23	1.53×10^{28}
82 Se			4.49 ± 1.00	22.22	3.6×10^{23}	3.55	4.54×10^{27}
$^{96}\mathrm{Zr}$	$2.86 {\pm} 0.26$	9.26	2.54 ± 0.89	35.03	9.2×10^{21}	8.60	6.81×10^{27}
^{100}Mo	$6.25 {\pm} 0.64$	10.20	5.13 ± 1.24	24.13	1.1×10^{24}	0.44	2.14×10^{27}
$^{110}\mathrm{Pd}$	$7.15 {\pm} 0.75$	10.54	6.49 ± 1.42	21.87	$6.0 imes 10^{17}$	852	$4.36 imes 10^{27}$
$^{116}\mathrm{Cd}$			$3.74\pm$ 0.80	21.44	1.7×10^{23}	1.48	3.72×10^{27}
^{124}Sn			3.77 ± 1.20	31.86			$6.60 imes 10^{27}$
$^{128}\mathrm{Te}$	$3.61 {\pm} 0.40$	10.94	$4.12{\pm}~0.86$	20.98	$1.0 imes 10^{23}$	8.71	8.34×10^{28}
$^{130}\mathrm{Te}$	$4.05 {\pm} 0.50$	12.29	$3.99{\pm}0.90$	22.43	3.0×10^{24}	0.35	$3.66 imes 10^{27}$
136 Xe			2.76 ± 0.72	26.02	3.4×10^{25}	0.14	7.14×10^{27}
$^{150}\mathrm{Nd}$	$2.83 {\pm} 0.43$	15.22	2.75 ± 0.44	16.46	1.8×10^{22}	3.01	1.63×10^{27}

PHFB model specific and global uncetainties NTMEs.





Sterile Majorona Neutrinos





Barabash et al.

Time dependence of
$$G_F$$
 $\Delta G_F / G_F \approx -0.1$ Bosonic neutrino $\sin^2 \chi < 0.6$



J. S. Diaz

Lorentz and CPT Violation

Modification of electron sum spectrum $(2\nu\beta\beta)$

Angular correlations $(0\nu\beta\beta)$

Parameters of gauge theoretical models

Physics beyond the standard model $(T_{1/2}^{0\nu} > 3.0 \times 10^{25} \text{ yr})$ Limits Parameters $\langle m_{\nu} \rangle$ $< 0.23 \ eV$ Effective Majorana neutrino mass $\begin{array}{ll} \langle \eta \rangle & < & 3.18 \times 10^{-9} \\ \langle \lambda \rangle & < & 5.73 \times 10^{-7} \end{array}$ L-R Coupling constant R-R Coupling constant $\langle M_N \rangle > 9.46 \times 10^7 \ GeV$ Heavy Neutrino mass $\langle \xi \rangle \quad < \quad 9.16 \times 10^{-9}$ $m_{W_R} > 1.44 \left(\frac{\left\langle m_N^{(V)} \right\rangle}{1 \ TeV}\right)^{-1/4}$ Mass of W_R $\tan\langle\xi\rangle$ $3.12 \times 10^{-3} \left(\frac{\langle m_N^{(V)} \rangle}{1 \ TeV}\right)^{1/2}$ $\begin{array}{rcl} \langle g_m\rangle &<& 8.2\times 10^{-5}\\ \zeta &<& 1.43\times 10^{-5} \end{array}$ Neutrino-Majoron coupling **R**-parity violation |f| < 2.72Compositeness scale $\langle m_N \rangle > 433 \ GeV$ Excited Majoron neutrino mass $\begin{array}{rcl} \epsilon_{I} & < & 1.16 \times 10^{-9} \left(\frac{M_{I}}{100 \ GeV}\right)^{2} \\ \alpha_{I}^{(L)} & < & 1.51 \times 10^{-10} \left(\frac{M_{I}}{100 \ GeV}\right)^{2} \\ \alpha_{I}^{(R)} & < & 3.82 \times 10^{-8} \left(\frac{M_{I}}{100 \ GeV}\right)^{2} \end{array}$ Leptoquark-Higgs coupling

 $\delta v < 4 \times 10^{-16}$

VLI

Physics @ LHC

(*i*)
$$e^-e^- \rightarrow W^-W^-$$

(*ii*) $e^-e^{\pm} \rightarrow H^-H^{\pm}$
(*iii*) $e^{\pm}e^{\pm} \rightarrow H^{\pm}H^{\pm}$
(*iv*) $t \rightarrow bW^-l_i^+l_j^+$
(*v*) $W^+ \rightarrow JJ'l_i^+l_j^+$
A non - observation of eejj - like events in
the future LHC run at $\sqrt{s} = 11$ TeV would
rule out all but the pure leptoquark
contribution as the dominat mode
to $0\nu\beta\beta$ decay.



in

EXPERMENTAL ACTIVITIES IN INDIA

INO Collaboration

Ahmadabad: Physical Research Lab. Aligarh: Aligarh Muslim University Allahabad: HRI Bhubaneswar: Institute of Physics Calicut: University of Calicut Chandigarh: Panjab University Chennai: IIT, Madras IMSc Delhi: University of Delhi Guwahati: IIT, Guwahati Hawaii (USA): University of Hawaii Indore: IIT, Indore Jammu: University of Jammu Kalpakkam: IGCAR Kolkata: Ramakrishna Mission Vivekananda University, SINP, VECC, University of Calcutta Lucknow: Lucknow University Madurai: American College Mumbai: BARC, IIT, Bombay TIFR Mysore: University of Mysore Sambalpur: Sambalpur University Srinagar: University of Kashmir Varanasi: Banaras Hindu University



India-based Neutrino Observatory (INO) an underground laboratory



Cavern set in Charnockite rock under 1580 m peak, 1289 m Vertical cover, 1000 m all-round

The INdia-based TIN detector (TIN.TIN) Initiative for DBD experiment in India

Proposal for an experiment at site of the INO (India-based neutrino Observatory)

 ^{124}Sn (Q =2288.1 \pm 1.6 keV)

- Sn has $T_C \sim 3.7 \text{ K}$
- Electronic specific heat falls off exponentially below T_C
- Only lattice specific heat ($\sim T^3$) present below $\sim 500 \text{ mK}$
- Z=50 shell is closed
- Simple metallurgy

¹²⁴Sn: $T_{1/2} > (0.8-1.2) \times 10^{21}$ yrs Nucl. Phys. A 807, 269(2008)

Cryogen free dilution refrigerator installed at TIFR



V. Singh et. al. Pramana 81 (2013) 719

Bolometer Detector - Initial Tests





- Thermal pulses are characterized by sharp rise time and very slow decay time.
- Since the pulses are slow high frequency noise can be filtered
- Count rate matched with the alpha source.
- Need to improve on Signal to Noise ratio to achieve good resolution and generate an energy spectrum.

Low Background Counting Facility (Tiles)





Sensitivity of the setup:

0.04mBq/g for ⁴⁰K ~ 1ppm 0.004mBq/g for ²³²Th ~1 ppt



Background Gamma ray spectra with and without veto

- Screening of materials and their selection on the basis of radio purity levels
- Rare decay studies

N. Dokania et al., NIM A 745 (2014) 119-127

Future Goal...

After prototype demonstration

Build a large scale detector (~ 1 ton) at INO lab (3 phases: 100 Kg, 500 Kg, 1000 Kg) $F_n = G^{0v} |M^{0v}|^2 = 8.611 \times 10^{-13} y^{-1} (PHFB) (Rath et al.)$ $= 6.205 \times 10^{-14} y^{-1} (SM)$ With 00 % antichment 0.2% are negatively negatively at 0 a background = 0.01

With 90 % enrichment, 0.2% energy resolution at $Q_{\beta\beta}$, background ~ 0.01 counts/ keV. kg.y and 1 year observation time

Phase	Mass (kg)	Т _{1/2} (у) [68% С.L.]	<m<sub>ββ> eV</m<sub>
I	100	8.6 × 10 ²⁵	0.06-0.22
II	500	1.9 × 10 ²⁶	0.04-0.15
	1000	2.7 × 10 ²⁶	0.03-0.12

➢ KamLANDZen + EXO : $T_{1/2}^{0\nu}$ > 3.4×10²⁵ y (90% C.L.), <m_{ββ}> < 0.12-0.25 eV (PRL 110 (2013) 062502)

➢ GERDA : $T_{1/2}^{0\nu} > 2.1 \times 10^{25} y$ (90% C.L.), <m_{ββ} < 0.2-0.4 eV (PRL 111 (2013) 122503)

Tin.Tin Collaboration

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Thanks to INO collaboration



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"A poet once said 'The whole universe is in a glass of wine.' We will probably never know in what sense he meant that, for poets do not write to be understood. But it is true that if we look at a glass closely enough we see the entire universe.

> Thanks a lot for your indulgence

