

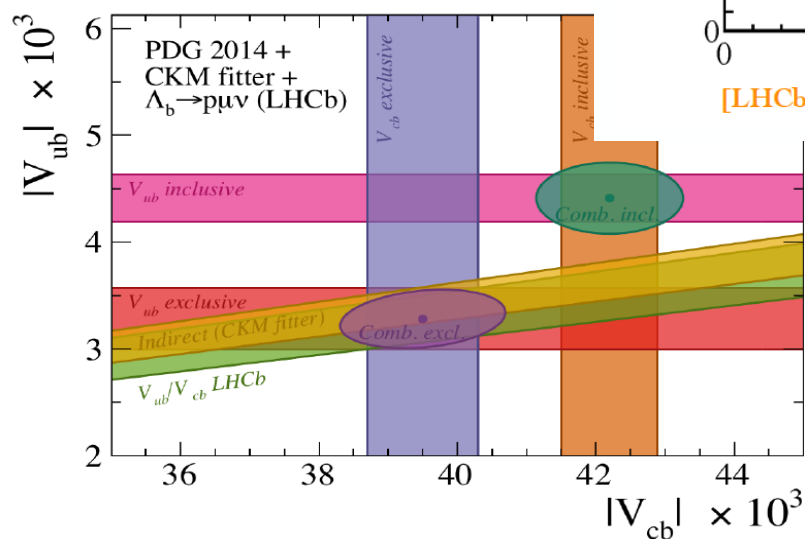
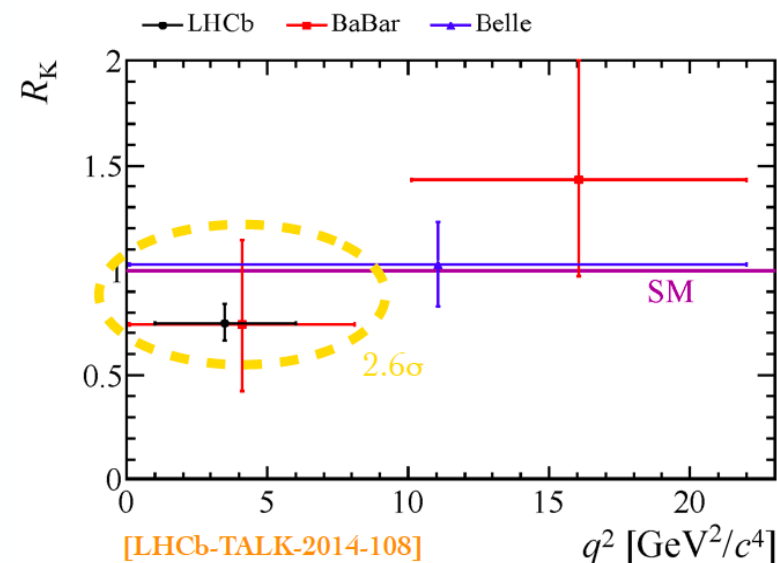
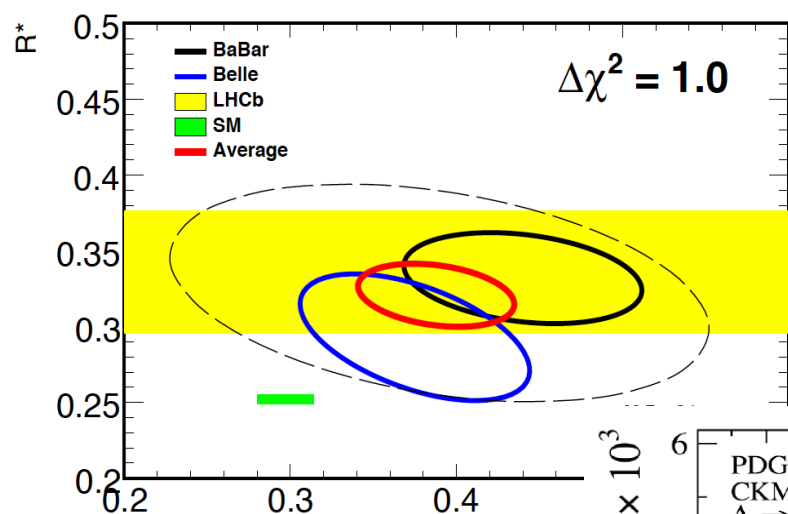
*Links between EWP anomalies and semi-leptonic B decays*  
[speculations on the breaking of Lepton Flavor Universality in B decays]

Gino Isidori  
[ *University of Zürich* ]

- ▶ On the anomalies
- ▶ Speculations on the breaking of **L**epton **F**lavor **U**niversality
- ▶ Conclusions

## On the anomalies

[i.e. where (and why) I start to think there is something interesting...]



*This talk contains personal biases...!*

## I. The EWP anomalies in $B \rightarrow K^{(*)} \mu\mu / ee$ [LHCb]

The largest EWP anomaly is the “famous” deviation from SM in  $P_5'$  [ $B \rightarrow K^* \mu\mu$ ]

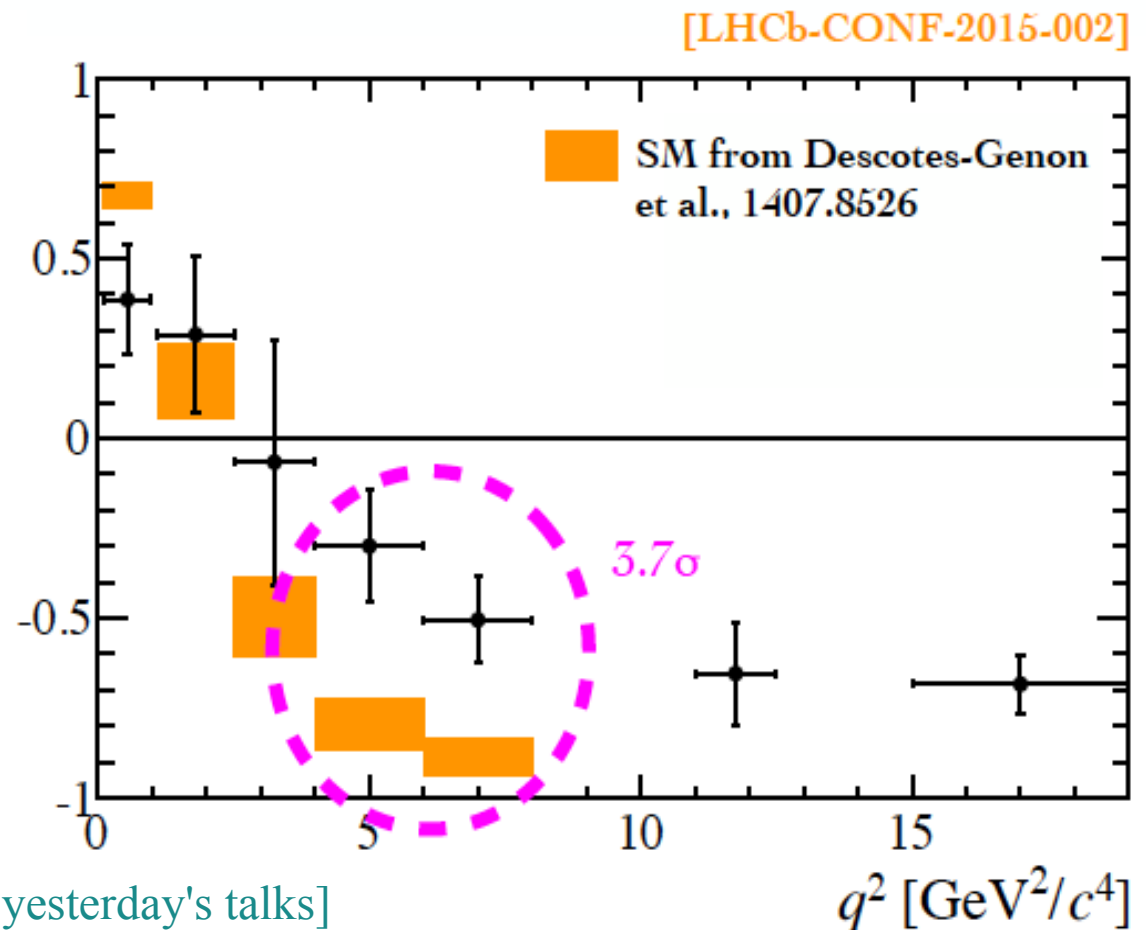
But less significant anomalies present also in other  $B \rightarrow K^* \mu\mu$  observables and also in other  $b \rightarrow s \mu\mu$  channels [overall smallness of all  $BR(B \rightarrow \text{Hadron} + \mu\mu)$ ]

### Pro NP:

- Reduced tension in all the observables with same set of non-standard short-distance Wilson coefficients ( $\Delta C_9$  or  $\Delta C_9 = -\Delta C_{10}$ )

### Against NP:

- $C_9$  sensitive to charm re-scattering effects
- Significance reduced with conservative estimates of non-factorizable corrections



[  $\rightarrow$  yesterday's talks ]

## I. The EWP anomalies in $B \rightarrow K^{(*)} \mu\mu / ee$ [LHCb]

The most interesting aspect (*in my opinion*) is the  $2.6\sigma$  deviation from the SM observed in the LFU ratio

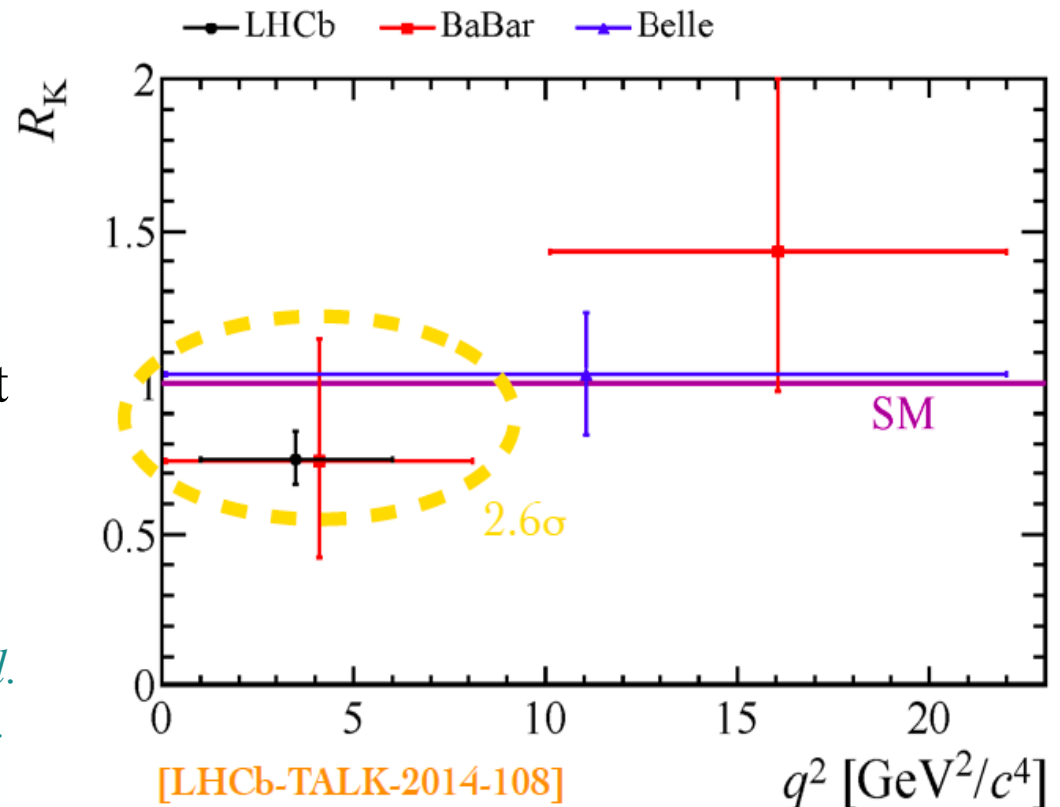
$$R_K = \frac{\int d\Gamma(B^+ \rightarrow K^+ \mu\mu)}{\int d\Gamma(B^+ \rightarrow K^+ ee)}$$

[1-6] GeV<sup>2</sup>

- Negligible th. error  $\rightarrow$  clean test of LFU (in neutral currents)

$$R_K = 1 \pm O(1\%)$$

Bordone *et al.*  
work in prog.



$\rightarrow$  The anomaly is perfectly described assuming NP only in  $b \rightarrow s \mu\mu$  [and not in  $b \rightarrow s ee$ ] consistently with the various  $b \rightarrow s \mu\mu$  anomalies

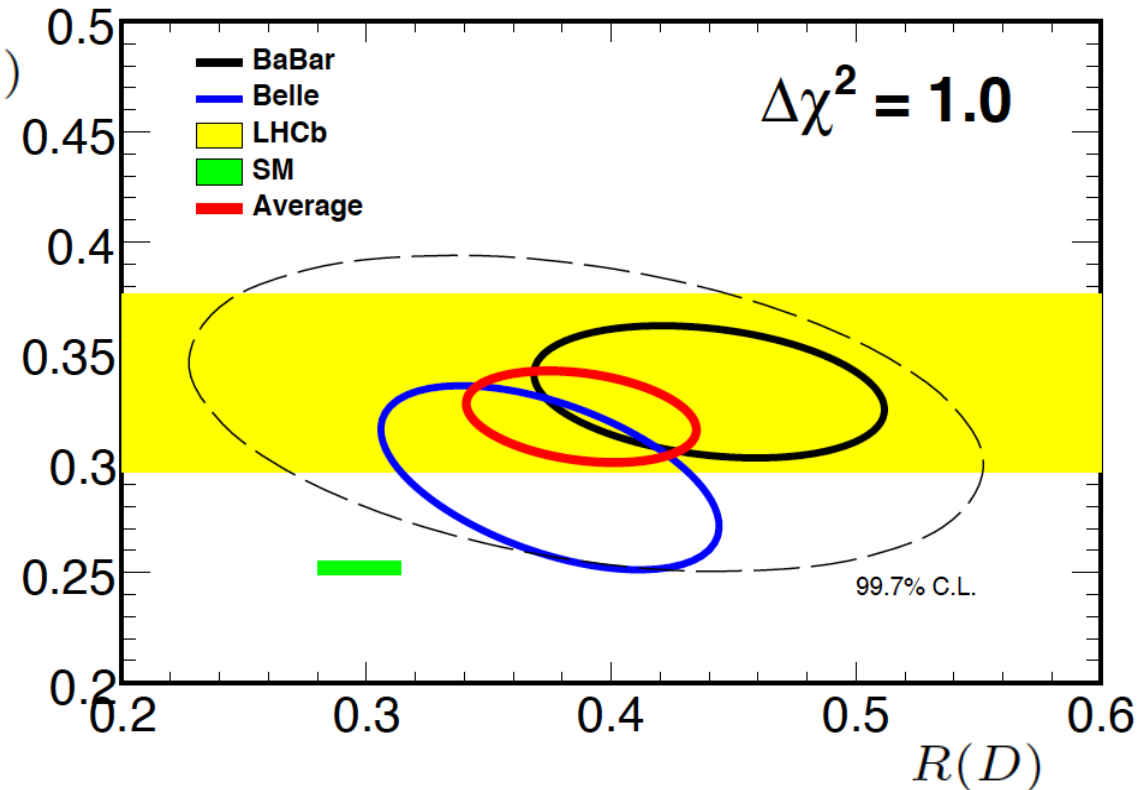
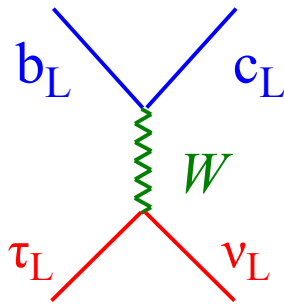
## II. $B \rightarrow D^{(*)} \tau \nu$ [Babar, Belle, LHCb]

[ → today's talks ]

Test of **LFU** in charged currents

[ $\tau$  vs. light leptons ( $\mu$ ,  $e$ )]:  $R(D^{*})$

$$R(X) = \frac{\Gamma(B \rightarrow X \tau \bar{\nu})}{\Gamma(B \rightarrow X \ell \bar{\nu})}$$



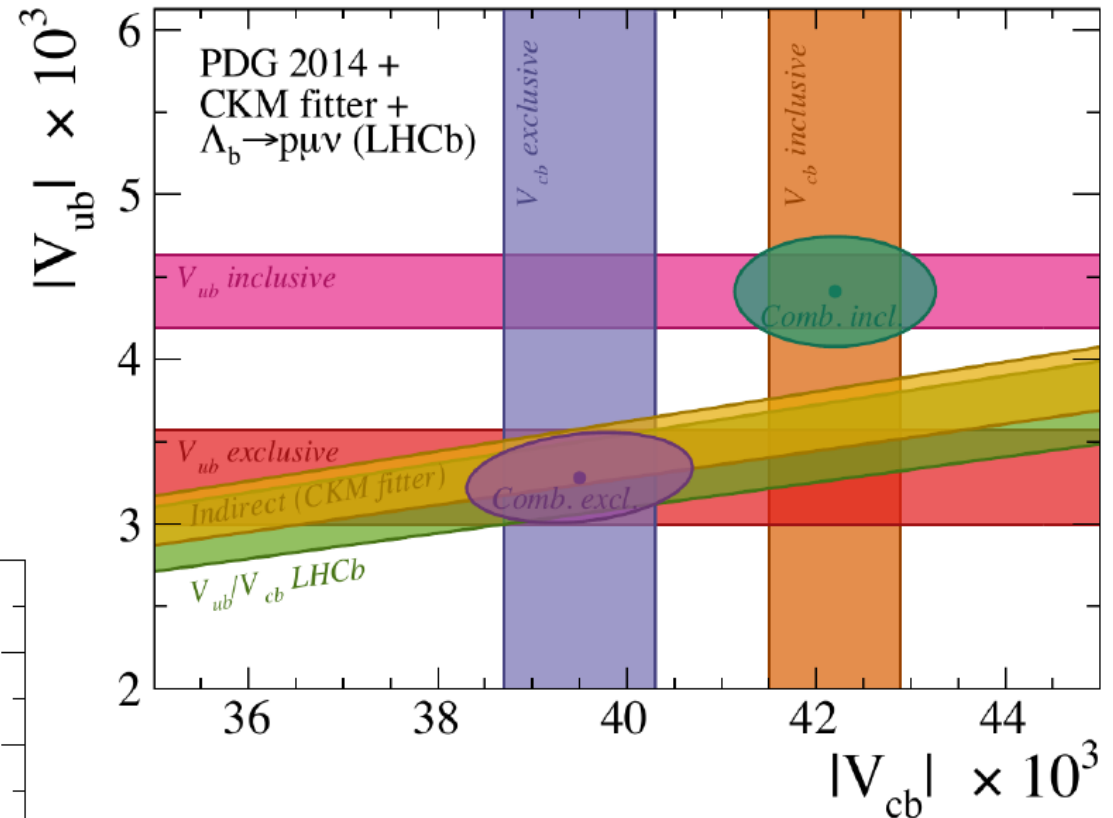
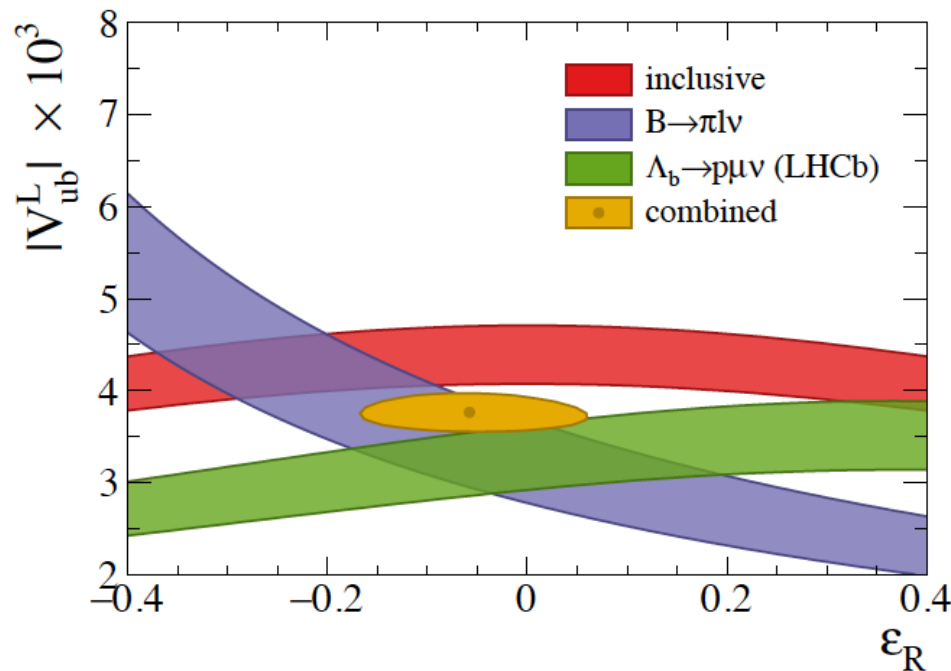
- **SM** prediction quite **solid**: f.f. uncertainty cancel (*to a good extent...*) in the ratio
- Consistent exp. results by 3 (very) different experiments
  - **3.9 $\sigma$**  excess over SM (if D and D\* combined)
  - The two channels are well consistent with a **universal enhancement** ( $\sim 30\%$ ) of the SM  $b_L \rightarrow c_L \tau_L \nu_L$  amplitude (*RH or scalar amplitudes disfavored*)

### III. $|V_{ub}/V_{cb}|$ from $B(\Lambda_b \rightarrow p\mu\nu)/B(\Lambda_b \rightarrow \Lambda_c\mu\nu)$ [Babar, Belle, LHCb]

[ → today's talks ]

Long-standing discrepancy between exclusive and inclusive determinations of both  $|V_{ub}|$  &  $|V_{cb}|$  (again charged currents...)

New ingredient:  $|V_{ub}/V_{cb}|$  from  $B(\Lambda_b \rightarrow p\mu\nu)/B(\Lambda_b \rightarrow \Lambda_c\mu\nu)$



- Consistent with other exclusive data
- Increased tension between excl. & incl.
- RH currents as possible explanation of the tension strongly disfavored

## Speculations on the breaking of **L**epton **F**lavor **U**niversality

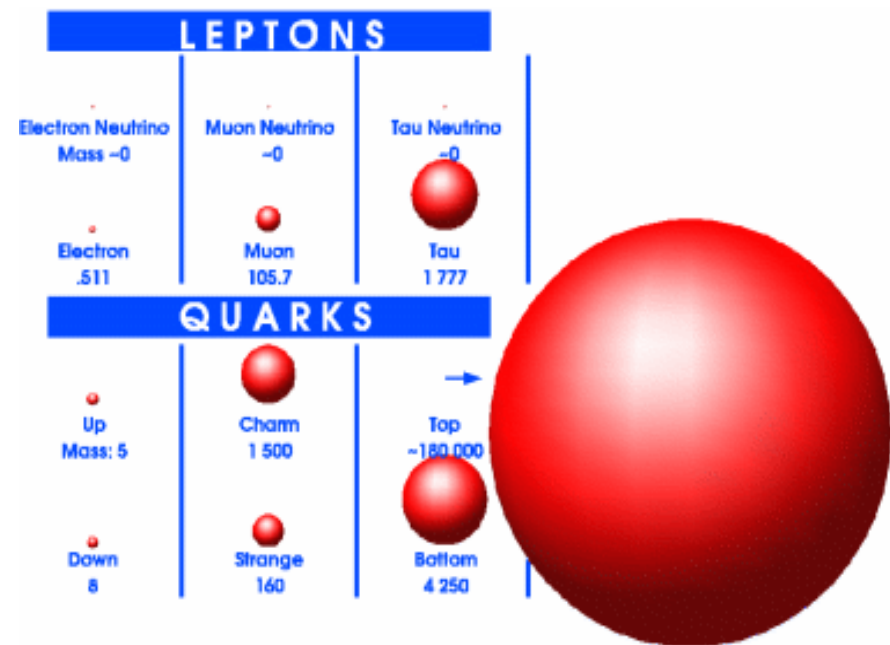
## ► Speculations on the breaking of LFU

These anomalies have stimulated a lot of theoretical activity.

Most interesting aspect (*in my opinion*): possible breaking of LFU, both in charged currents ( $b \rightarrow c\tau\nu$  vs.  $b \rightarrow c\mu\nu$ ) and in neutral currents ( $b \rightarrow s\mu\mu$  vs.  $b \rightarrow see$ )

A few general messages:

- ★ LFU is not a fundamental symmetry of the SM Lagrangian (*accidental symmetry in the gauge sector, broken by Yukawas*)
- ★ LFU tests at the Z peak are not very interesting ( $\rightarrow$  gauge sector)
- ★ Most stringent tests of LFU involve only 1<sup>st</sup>-2<sup>nd</sup> gen. quarks & leptons
  - $\rightarrow$  Natural to conceive NP models where LFU is violated more in processes with 3<sup>rd</sup> gen. quarks ( $\leftrightarrow$  hierarchy in Yukawa coupl.)





## ► Speculations on the breaking of LFU

These anomalies have stimulated a lot of theoretical activity:

S. Fajfer, J. F. Kamenik, I. Nisandzic and J. Zupan, Phys. Rev. Lett. **109** (2012) 161801 [[arXiv:1206.1872](#)].

S. Descotes-Genon, J. Matias and J. Virto, Phys. Rev. D **88** (2013) 074002 [[arXiv:1307.5683](#)].

W. Altmannshofer and D. M. Straub, Eur. Phys. J. C **73** (2013) 2646 [[arXiv:1308.1501](#)].

A. Datta, M. Duraissamy and D. Ghosh, Phys. Rev. D **89** (2014) 7, 071501 [[arXiv:1310.1937](#)].

G. Hiller and M. Schmaltz, Phys. Rev. D **90** (2014) 054014 [[arXiv:1408.1627](#)]; JHEP **1502** (2015) 055

A. Crivellin and S. Pokorski, Phys. Rev. Lett. **114** (2015) 1, 011802 [[arXiv:1407.1320](#)].

S. L. Glashow, D. Guadagnoli and K. Lane, Phys. Rev. Lett. **114** (2015) 091801 [[arXiv:1411.0565](#)].

+ many others...

...but most attempts focused either on specific NP models (mainly for EWP anom.) or on “partial” EFT-type approaches (focused only on **quark**×**lepton** ops.).

What I will discuss today is what happens if we try to describe all these effect within a simplified (rather general) dynamical model:

- low-energy correlations among **quark**×**quark**, **quark**×**lepton**, **lepton**×**lepton**
- correlation between **low-energy** and **high-energy** physics

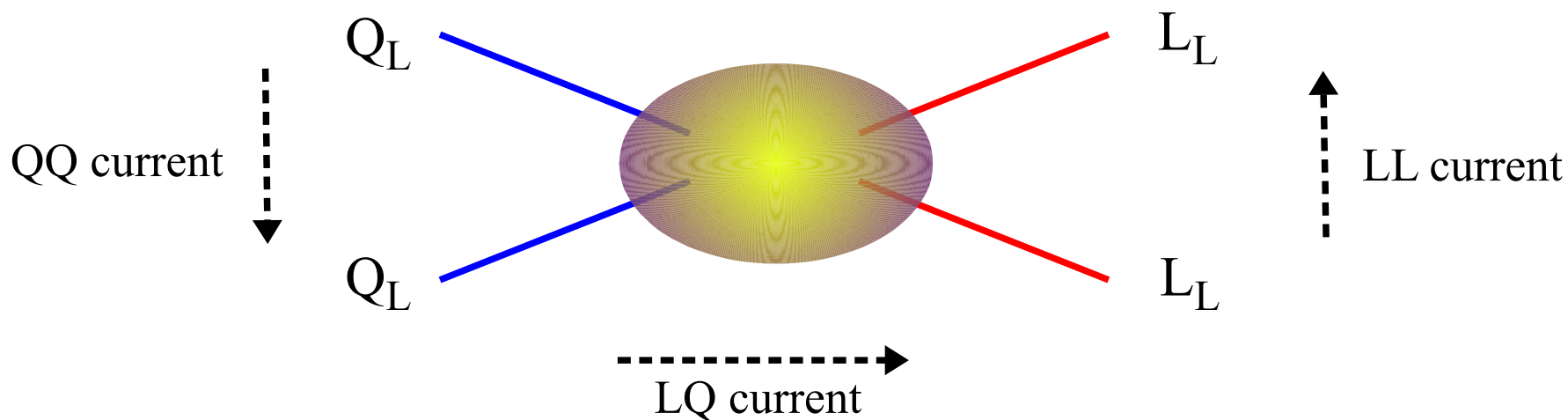
► A “prototype data-inspired” model:

Main assumptions:

- NP in both charged & neutral currents + RH currents disfavored +  $SU(2)_L \times U(1)_Y$  symmetry  $\rightarrow$   **$SU(2)_L$ -triplet effective operator**

$$\frac{g_q g_\ell}{\Lambda^2} \lambda_{ij}^q \lambda_{kl}^\ell (\bar{Q}_L^i T^a \gamma_\mu Q_L^j) (\bar{L}_L^k T^a \gamma^\mu L_L^l)$$

Bhattacharya *et al.* '14  
Alonso, Grinstein, Camalich '15  
Greljo, GI, Marzocca '15



► A “prototype data-inspired” model:

Main assumptions:

- NP in both charged & neutral currents + RH currents disfavored +  $SU(2)_L \times U(1)_Y$  symmetry  $\rightarrow$   **$SU(2)_L$ -triplet effective operator**

$$\frac{g_q g_\ell}{\Lambda^2} \lambda_{ij}^q \lambda_{kl}^\ell \left( \bar{Q}_L^i T^a \gamma_\mu Q_L^j \right) \left( \bar{L}_L^k T^a \gamma^\mu L_L^l \right)$$

Greljo, GI, Marzocca '15

- We assume this effective operator is the result of integrating-out a **heavy triplet of vector bosons (W', Z')** coupled to a single current:

$$J_\mu^a = g_q \lambda_{ij}^q \left( \bar{Q}_L^i \gamma_\mu T^a Q_L^j \right) + g_\ell \lambda_{ij}^\ell \left( \bar{L}_L^i \gamma_\mu T^a L_L^j \right) \longrightarrow \frac{1}{2m_V^2} J_\mu^a J_\mu^a$$

- low-energy correlations among **quark**×**quark**, **quark**×**lepton**, **lepton**×**lepton**
- correlation between **low-energy** and **high-energy** physics

► A “prototype data-inspired” model:

Main assumptions:

- NP in both charged & neutral currents + RH currents disfavored +  $SU(2)_L \times U(1)_Y$  symmetry  $\rightarrow$   **$SU(2)_L$ -triplet effective operator**

$$\frac{g_q g_\ell}{\Lambda^2} \lambda_{ij}^q \lambda_{kl}^\ell (\bar{Q}_L^i T^a \gamma_\mu Q_L^j) (\bar{L}_L^k T^a \gamma^\mu L_L^l)$$

Greljo, GI, Marzocca '15

- We assume this effective operator is the result of integrating-out a **heavy triplet of vector bosons (W', Z')** coupled to a single current:

$$J_\mu^a = g_q \lambda_{ij}^q (\bar{Q}_L^i \gamma_\mu T^a Q_L^j) + g_\ell \lambda_{ij}^\ell (\bar{L}_L^i \gamma_\mu T^a L_L^j) \longrightarrow \frac{1}{2m_V^2} J_\mu^a J_\mu^a$$

- **Non-Universal flavor structure** of the currents  $\rightarrow$  **mainly 3<sup>rd</sup> generations**

$$\lambda_{ij}^{q,\ell} = \delta_{i3} \delta_{3j} + \text{small corrections for 2<sup>nd</sup> (& 1<sup>st</sup>) generations}$$

*(hierarchy determined by CKM in the quark sector)*

► A “prototype data-inspired” model:

A closer look to the flavor structure of the model:

$$J_\mu^a = g_q \lambda_{ij}^q \left( \bar{Q}_L^i \gamma_\mu T^a Q_L^j \right) + g_\ell \lambda_{ij}^\ell \left( \bar{L}_L^i \gamma_\mu T^a L_L^j \right)$$

→ Coupling to 3<sup>rd</sup> generations not suppressed [*dynamical assumption*]

→ Coupling to light generations controlled by small  $U(2)_q \times U(2)_l$  Glashow *et al.* '14  
 breaking spurions related to subleading terms in the Yukawa couplings

Barbieri *et al.* '11

*Connection to CKM matrix in the quark sector:*

$$\lambda^q \simeq \begin{pmatrix} |\epsilon|^2 V_{3\alpha}^* V_{3\beta} & \epsilon^* V_{3\alpha}^* \\ \epsilon V_{3\beta} & 1 \end{pmatrix}$$

$$\lambda_{bd} \ll \lambda_{bs} \ll \lambda_{bb} = 1$$

$$\lambda_{ss} \sim \lambda_{bs}^2$$

down-type  
mass basis

$$\epsilon \lesssim 1$$

► Effects in charged currents:

$$\frac{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM+NP}}}{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM}}} = 1 + \boxed{R_0} \lambda_{ii}^\ell \quad R_0 \equiv \frac{g_\ell g_q}{g^2} \frac{m_W^2}{m_V^2}$$

I. From  $R(D^*)$  &  $R(D)$  data  $[\Gamma(b \rightarrow c \tau \nu)/\Gamma(b \rightarrow c \mu \nu)] \rightarrow \boxed{R_0 = 0.14 \pm 0.04}$

► Effects in charged currents:

$$\frac{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM+NP}}}{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM}}} = 1 + R_0 \lambda_{ii}^{\ell} \quad R_0 \equiv \frac{g_{\ell} g_q}{g^2} \frac{m_W^2}{m_V^2}$$

I. From  $R(D^*)$  &  $R(D)$  data  $[\Gamma(b \rightarrow c \tau \nu)/\Gamma(b \rightarrow c \mu \nu)] \rightarrow R_0 = 0.14 \pm 0.04$

II. In principle, it should be possible to get a strong bound on the sub-leading leptonic coupling ( $\lambda_{\mu\mu}$ ) from  $\Gamma(b \rightarrow c \mu \nu)/\Gamma(b \rightarrow c e \nu)$  but it turns out to be not so stringent ( $|\lambda_{\mu\mu}| \lesssim 0.1$  no dedicated studies @ B-factories)

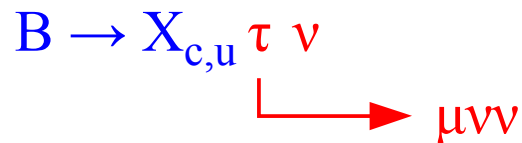
► Effects in charged currents:

$$\frac{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM+NP}}}{\mathcal{A}(b \rightarrow c \ell^i \bar{\nu}^i)_{\text{SM}}} = 1 + R_0 \lambda_{ii}^\ell \quad R_0 \equiv \frac{g_\ell g_q}{g^2} \frac{m_W^2}{m_V^2}$$

I. From  $R(D^*)$  &  $R(D)$  data  $[\Gamma(b \rightarrow c \tau \nu)/\Gamma(b \rightarrow c \mu \nu)] \rightarrow R_0 = 0.14 \pm 0.04$

II. In principle, it should be possible to get a strong bound on the sub-leading leptonic coupling ( $\lambda_{\mu\mu}$ ) from  $\Gamma(b \rightarrow c \mu \nu)/\Gamma(b \rightarrow c e \nu)$  but it turns out to be not so stringent ( $|\lambda_{\mu\mu}| \lesssim 0.1$  no dedicated studies @ B-factories)

III. Even if it is hard to quantify [*work in prog.*], this breaking of LFU in c.c could decrease the tension between exclusive & inclusive determinations of  $|V_{ub}|$  &  $|V_{cb}|$ :



Irreducible bkg. for the inclusive meas. subtracted  
(at present) assuming SM-like  $\Gamma(B \rightarrow X_{c,u} \tau \nu)$



if  $\Gamma(B \rightarrow X_{c,u} \tau \nu)$  is enhanced over the SM  $\rightarrow |V_{c(u)b}|_{\text{incl.}}$  are overestimated



► Global fit to low-energy data:

5 free parameters: 
$$\epsilon_{\ell,q} \equiv \frac{g_{\ell,q} m_W}{g m_V} \approx g_{\ell,q} \frac{122 \text{ GeV}}{m_V} + \lambda_{bs}^q, \lambda_{\mu\mu}^\ell, \lambda_{\tau\mu}^\ell$$

several constraints:

- R(D<sup>\*</sup>)
- R(D)
- R<sub>K</sub>
- P<sub>5</sub>'(B → K<sup>\*</sup>μμ)
- B(B → Kνν)
- ΔM<sub>B<sub>s</sub></sub>, ΔM<sub>B<sub>d</sub></sub>
- CPV(D-D)
- Γ(B → Xμν)/Γ(B → Xev)
- τ → 3μ
- Γ(τ → μνν)/Γ(τ → eνν)

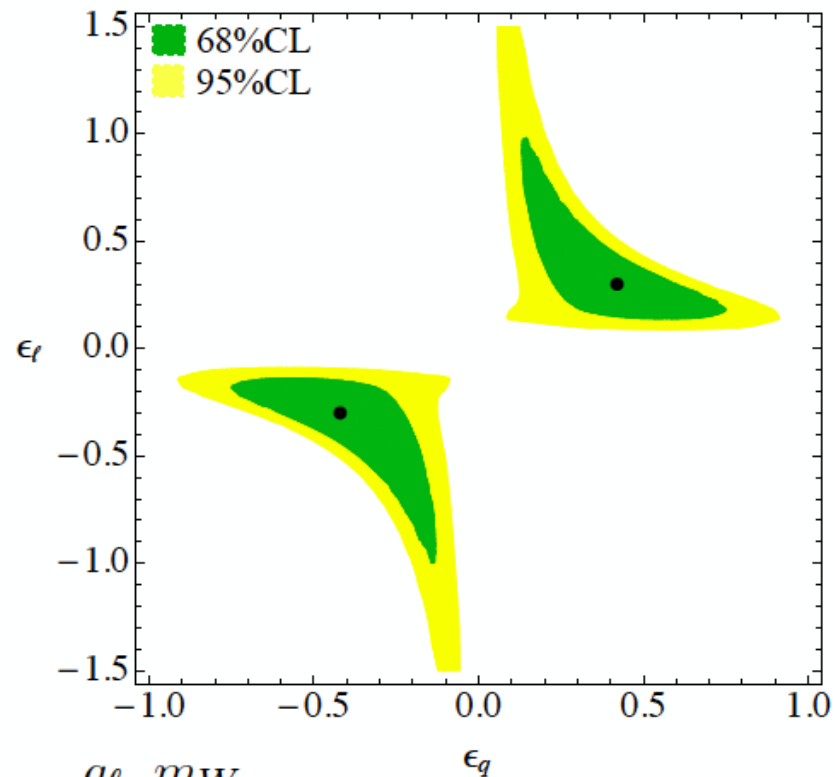


Overall good fit of low-energy data  
(non-trivial given tight constraints from ΔF=2 & LFV)

Best fit point:  $\epsilon_\ell \approx 0.37$ ,  $\epsilon_q \approx 0.38$   $p(\text{SM}) = 0.002$

(flavor structure of the sub-leading terms not really probed)

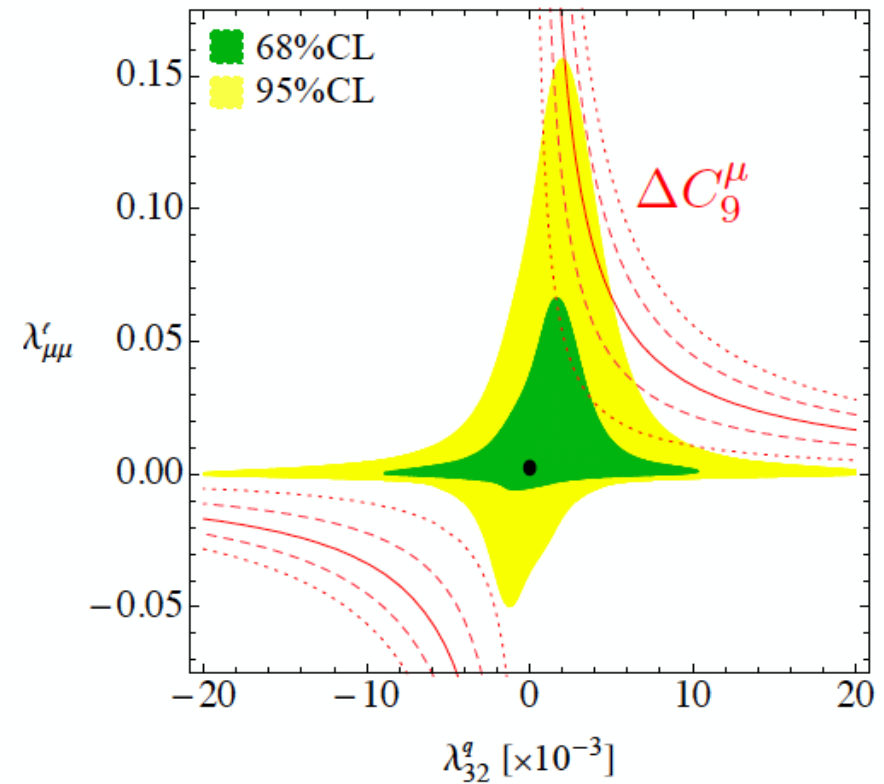
► Global fit to low-energy data:



$$\epsilon_{\ell,q} = \frac{g_{\ell,q} m_W}{g m_V}$$

$$\epsilon_{\ell}, \epsilon_q \lesssim 1$$

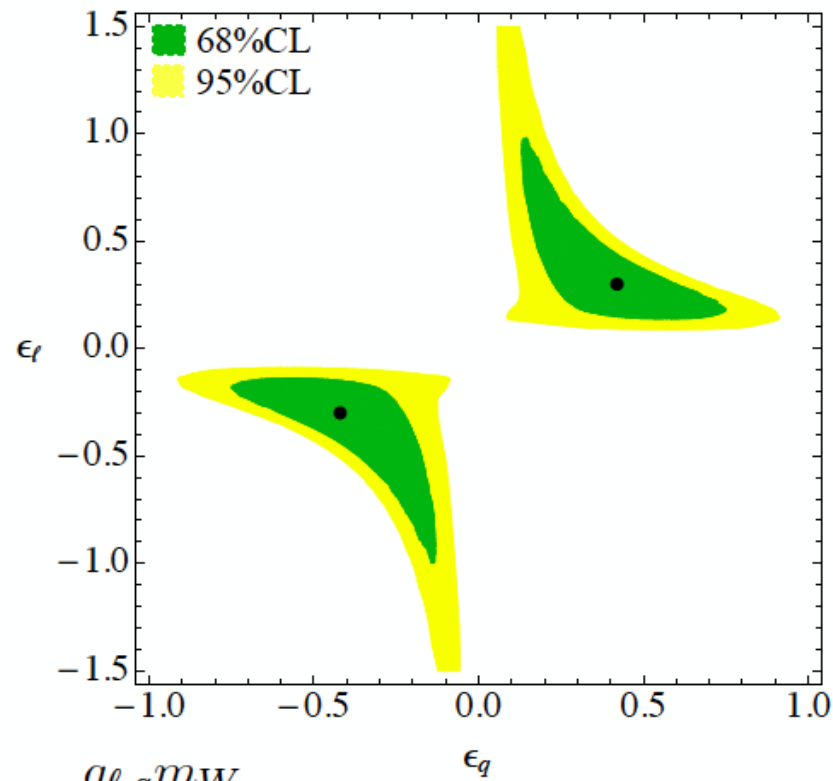
$\epsilon_q > \epsilon_{\ell}$  would improve  $b \rightarrow s \mu \mu$



$$\lambda_{\mu\mu} \lesssim 0.1$$

$$\lambda_{bs} \lesssim 0.015$$

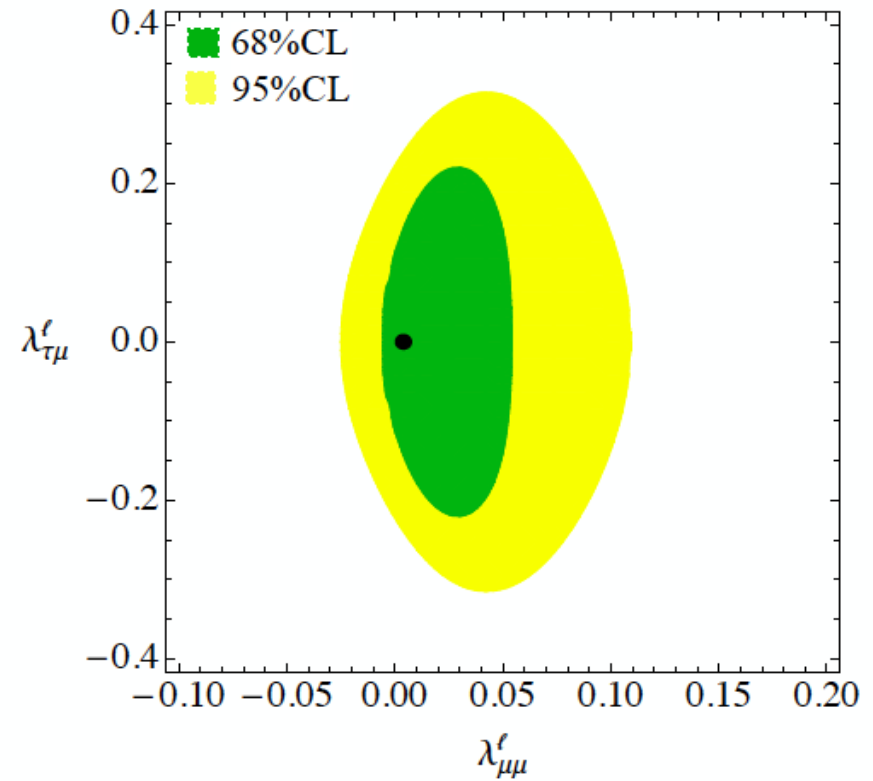
► Global fit to low-energy data:



$$\epsilon_{\ell,q} = \frac{g_{\ell,q} m_W}{g m_V}$$

$$\epsilon_{\ell}, \epsilon_q \lesssim 1$$

$\epsilon_q > \epsilon_{\ell}$  would improve  $b \rightarrow s \mu \mu$



$$\lambda_{\mu\mu} \lesssim 0.1$$

$$\lambda_{bs} \lesssim 0.015$$

► Future low-energy tests:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2m_V^2} J_\mu^a J_\mu^a \quad \text{works well...}$$

... and gives several clear predictions for future low-energy data:

•  $b \rightarrow c(u) l\nu$

$$\text{BR}(B \rightarrow D^* \tau \nu) / \text{BR}_{\text{SM}} = \text{BR}(B \rightarrow D \tau \nu) / \text{BR}_{\text{SM}} = \text{BR}(\Lambda_b \rightarrow \Lambda_c \tau \nu) / \text{BR}_{\text{SM}}$$

$$= \dots = \text{BR}(B_u \rightarrow \tau \nu) / \text{BR}_{\text{SM}} \quad R^{\mu/e}(X) \sim 10\% R^{\tau/\mu}(X)$$

→ universal  $\sim 30\%$  enhancement of C.C. semi-leptonic decays into tau leptons

→  $\sim 1\text{-}2\%$  (universal) breaking of universality between muons & electron (in CC modes)

► Future low-energy tests:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2m_V^2} J_\mu^a J_\mu^a \quad \text{works well...}$$

... and gives several clear predictions for future low-energy data:

•  $b \rightarrow c(u) l\nu$        $\text{BR}(B \rightarrow D^* \tau\nu)/\text{BR}_{\text{SM}} = \text{BR}(B \rightarrow D \tau\nu)/\text{BR}_{\text{SM}} = \text{BR}(\Lambda_b \rightarrow \Lambda_c \tau\nu)/\text{BR}_{\text{SM}}$   
 $= \dots = \text{BR}(B_u \rightarrow \tau\nu)/\text{BR}_{\text{SM}} \quad R^{\mu/e}(X) \sim 10\% R^{\tau/\mu}(X)$

•  $b \rightarrow s \mu\mu$        $\Delta C_9^\mu = -\Delta C_{10}^\mu$ , but overall size of the anom. should decrease

•  $b \rightarrow s \tau\tau$        $|\text{NP}| \sim |\text{SM}| \rightarrow$  large enhancement ( $\sim \text{BR} \times 4$ ) or strong suppr.

•  $b \rightarrow s \nu\nu$        $\sim \pm 50\%$  deviation from SM in the rate

► Future low-energy tests:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{2m_V^2} J_\mu^a J_\mu^a \quad \text{works well...}$$

... and gives several clear predictions for future low-energy data:

•  $b \rightarrow c(u) l \nu$        $\text{BR}(B \rightarrow D^* \tau \nu) / \text{BR}_{\text{SM}} = \text{BR}(B \rightarrow D \tau \nu) / \text{BR}_{\text{SM}} = \text{BR}(\Lambda_b \rightarrow \Lambda_c \tau \nu) / \text{BR}_{\text{SM}}$   
 $= \dots = \text{BR}(B_u \rightarrow \tau \nu) / \text{BR}_{\text{SM}} \quad R^{\mu/e}(X) \sim 10\% R^{\tau/\mu}(X)$

•  $b \rightarrow s \mu \mu$        $\Delta C_9^\mu = -\Delta C_{10}^\mu$ , but overall size of the anom. should decrease

•  $b \rightarrow s \tau \tau$        $|\text{NP}| \sim |\text{SM}| \rightarrow$  large enhancement ( $\sim \text{BR} \times 4$ ) or strong suppr.

•  $b \rightarrow s \nu \nu$        $\sim \pm 50\%$  deviation from SM in the rate

• **Meson mixing**       $\sim 10\%$  deviations from SM both in  $\Delta M_{B_s}$  &  $\Delta M_{B_d}$

•  $\tau$  decays       $\tau \rightarrow 3\mu$  not far from present exp. bound

► High-energy constraints:

The dynamical model

$$\mathcal{L}_V = -\frac{1}{4} D_{[\mu} V_{\nu]}^a D^{[\mu} V^{\nu]a} + \frac{m_V^2}{2} V_\mu^a V^{\mu a} + g_H V_\mu^a (H^\dagger T^a i \overleftrightarrow{D}_\mu H) + V_\mu^a J_\mu^a$$

The “heavy vector triplet” eff. Lagrangian [Pappadopulo, Tham, Torre, Wulzer, '14] in a rather peculiar parameter range:

- **W** and **Z** resonances in the mass range:

$$g_{l,q} \sim 1 \rightarrow m_V \sim 250 \text{ GeV}$$

$$g_{l,q} \sim \sqrt{4\pi} \rightarrow m_V \lesssim 1 \text{ TeV}$$

- Strong constraint on  $g_H$  from e.w. precision tests:

$$\epsilon_{l,q} = \frac{g_{l,q} m_W}{g m_V} \approx 0.3 \qquad \epsilon_H = \frac{g_H m_W}{g m_V} \lesssim 0.01$$

## ► High-energy constraints:

- The heavy vectors are produced mainly from 3<sup>rd</sup> gen. quarks ( $bb \rightarrow Z'$ ,  $bc \rightarrow W'$ ) and decay mainly in 3<sup>rd</sup> generations quarks or leptons ( $Z' \rightarrow \tau\tau, bb, tt$ ,  $W' \rightarrow tb, \tau\nu$ )



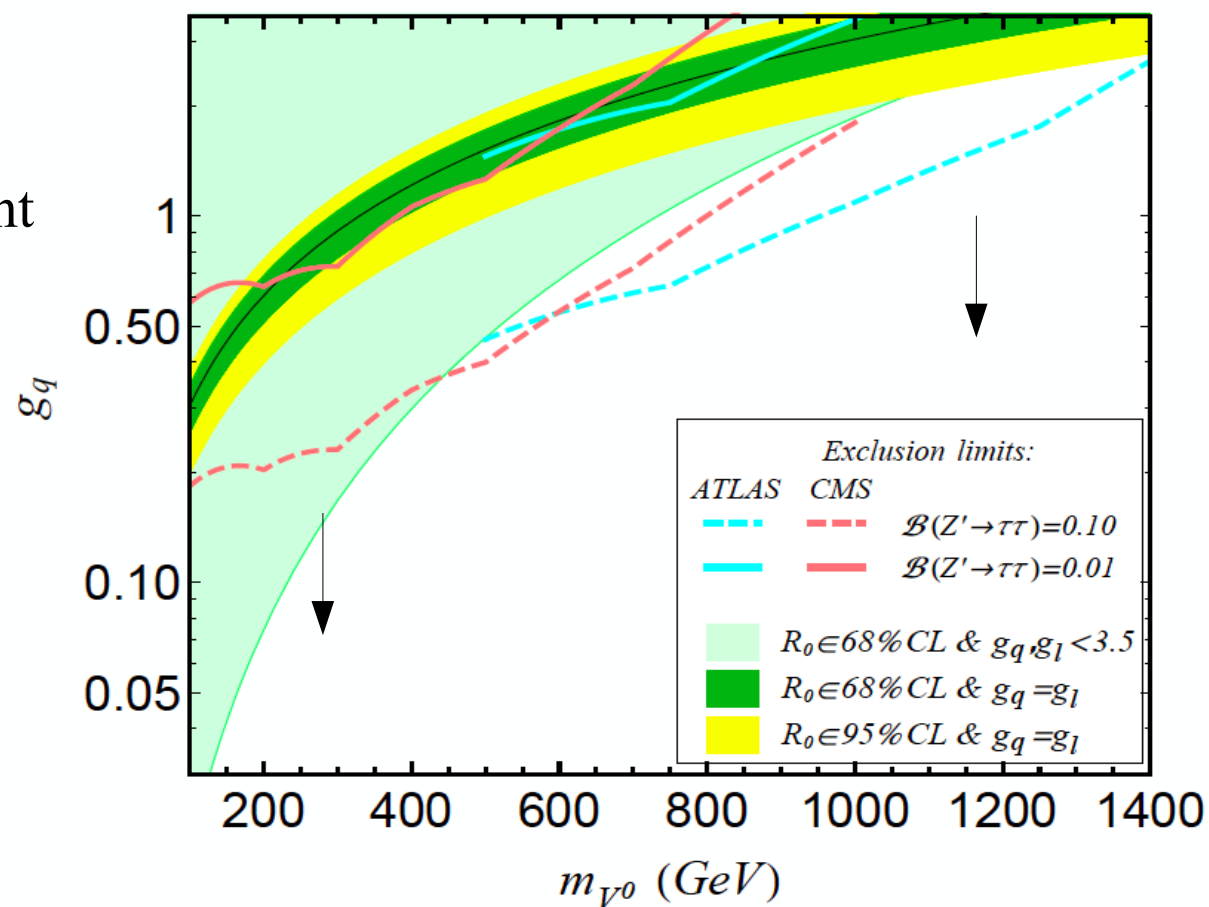
- Not a very easy signature...

The only really stringent constraint comes from  $Z' \rightarrow \tau\tau$



Minimal version of the model  
(no exotic decay channels)  
ruled out by direct searches

$$BR(Z' \rightarrow \bar{\tau}\tau) = \frac{g_\ell^2}{2g_\ell^2 + 6g_q^2 + \text{extra}}$$





## Conclusions

- Intriguing hints of LF non Universality in recent semi-leptonic B-physics data, but picture far from being clear → more data can help to clarify the situation
- Main messages of these recent anomalies:
  - (re)analyze B physics data without assuming LFU
  - conceive more low-energy tests of LFU (especially in B decays)
  - the search for LFV in charged leptons is extremely well motivated
  - the bounds on NP coupled mainly to 3<sup>rd</sup> generation are still relatively weak
  - the interplay of low- and high-energy searches is essential