



# MW measurement with forward muons: a PFD uncertainty perspective

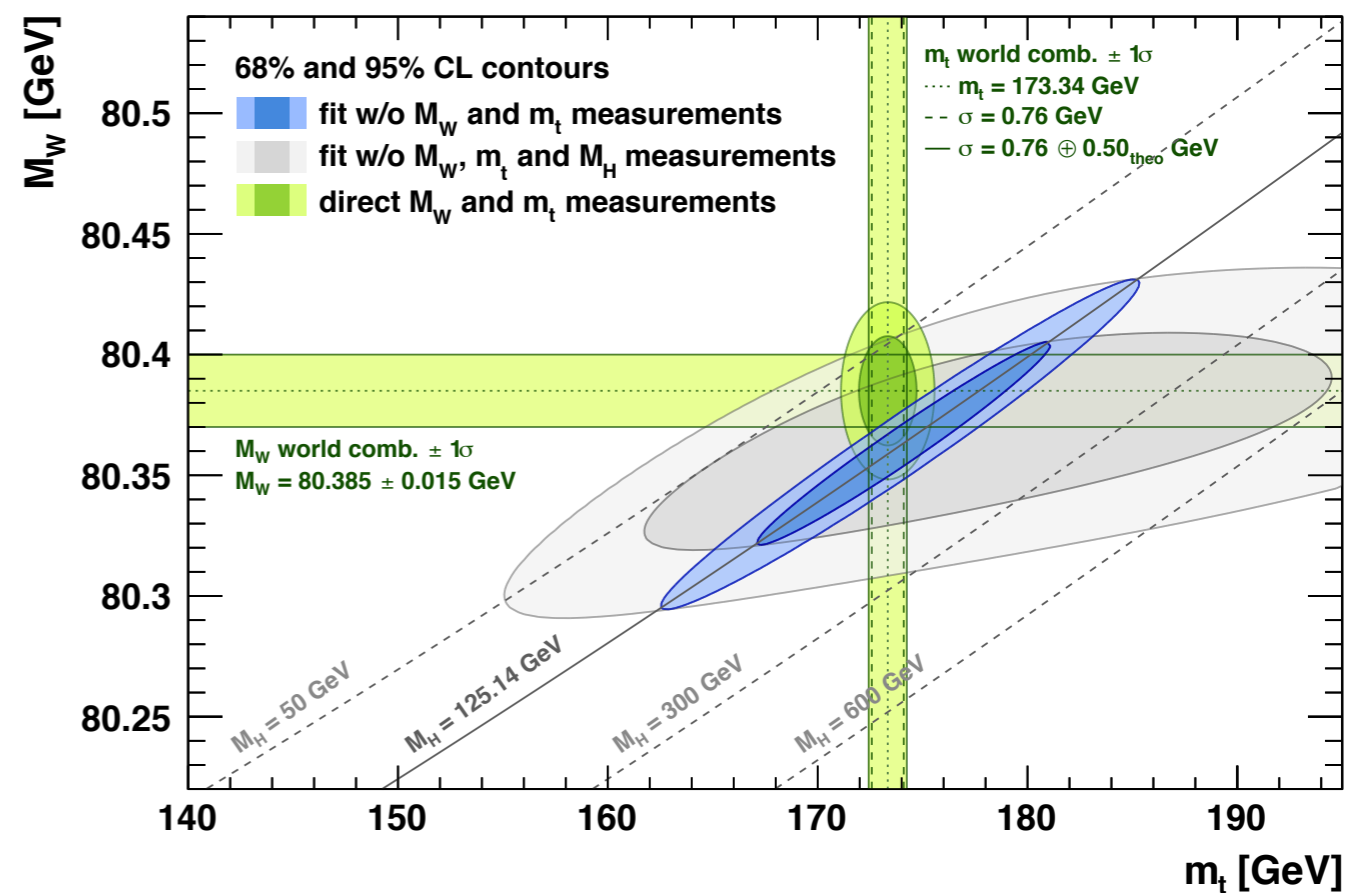
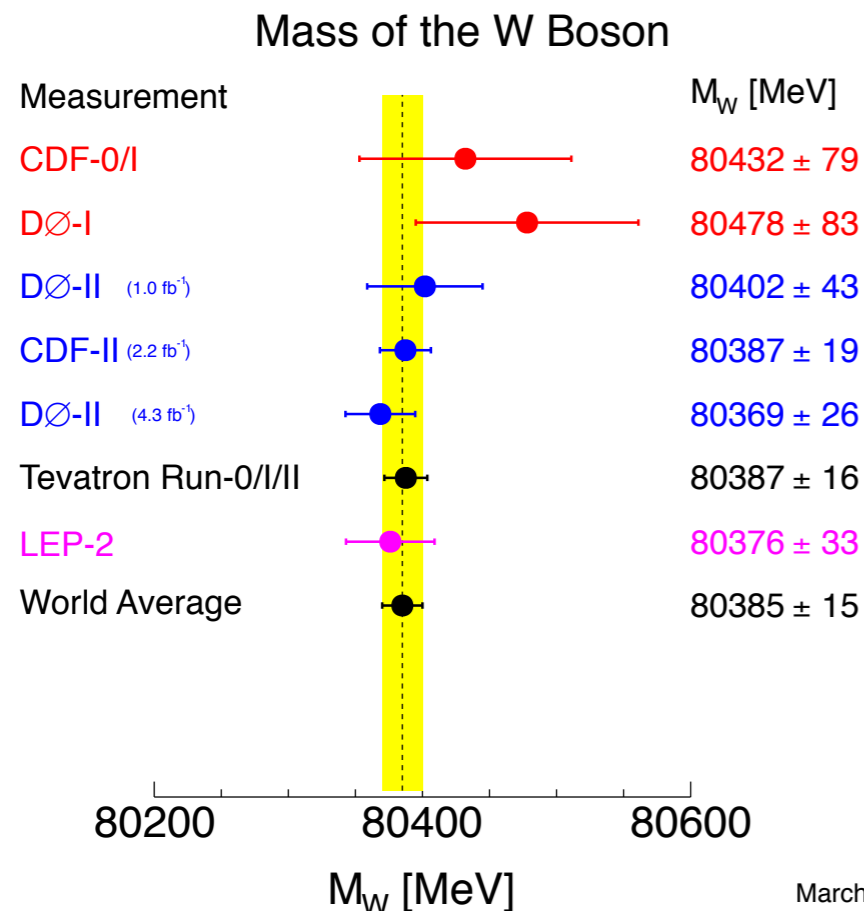
**Alessandro Vicini**

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CERN, November 4th 2015

work in collaboration with G.Bozzi, L.Citelli, M.Vesterinen  
based on [arXiv:1501.05587](https://arxiv.org/abs/1501.05587), [arXiv:1508.06954](https://arxiv.org/abs/1508.06954)

# Precision tests of the Standard Model



with the  $M_H$  input the SM lagrangian (gauge sector) is assigned,  
 the EW fit can determine the preferred  $M_W$  (2-loop EW+h.o.) and  $m_{\text{top}}$  (free parameter)  
 and check the compatibility of the SM hypothesis with the experimental measurements

the result of the global EW fit of the SM  
 yields a result for  $M_W$  with an error  $\Delta M_W = 8$  MeV smaller than the one of the direct measurement  
 $m_{\text{top}} = 173.81 \pm 0.85$  GeV compatible with the world average top mass

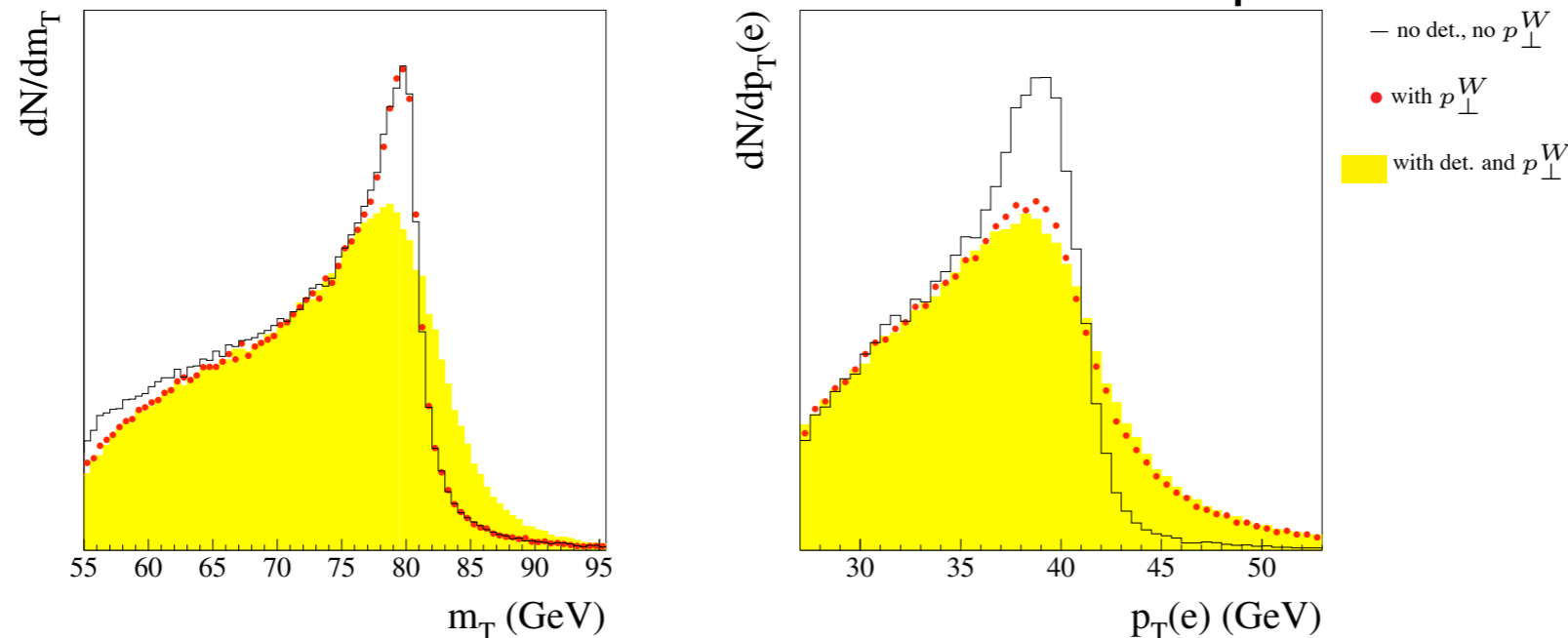
is the 1.5 sigma discrepancy in the above plot, between the data and the theoretical prediction,  
 just a fluctuation, a systematic effect of the  $M_W$  measurement at hadron colliders, a BSM hint?

can we aim at a  $M_W$  measurement at the  $O(10$  MeV) level of precision at the LHC ?

# MW measurement from Drell-Yan observables

- lepton-pair transverse mass  $M_{\perp}^W = \sqrt{2p_{\perp}^l p_{\perp}^{\nu} (1 - \cos \phi_{l\nu})}$
- charged lepton transverse momentum
- missing transverse momentum

- sensitivity to MW via the jacobian factor peaked at the physical mass value



## lepton-pair transverse mass

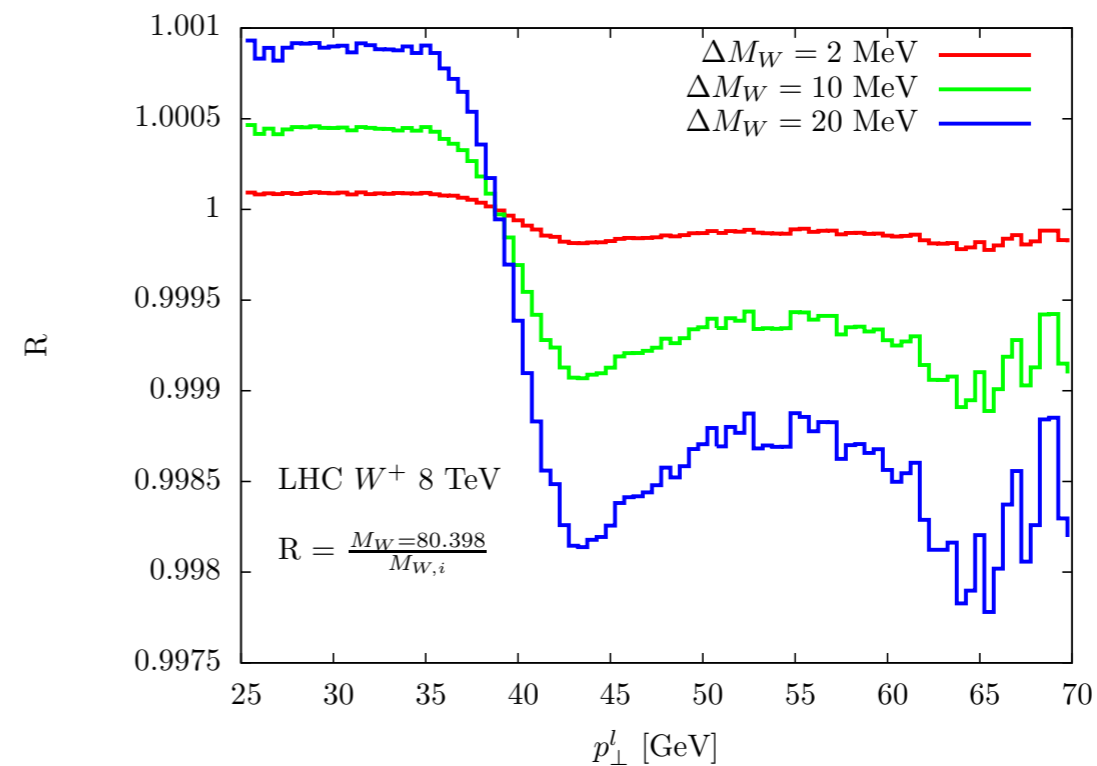
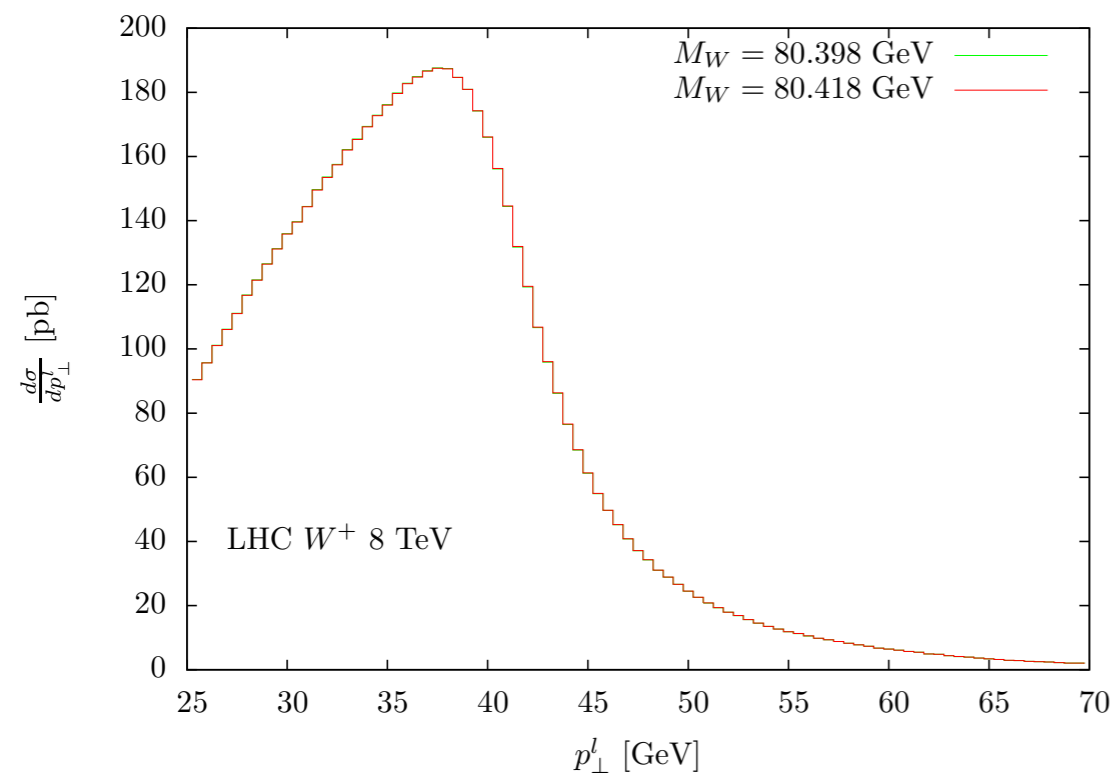
- ▶ stable w.r.t. inclusion of radiative corrections
  - ▶ problematic determination of the neutrino  $p_T$  in presence of high pile-up (modeling of hadr. recoil)
  - ▶ moderate PDF uncertainty not exceeding  $O(10 \text{ MeV})$  see also Bozzi, Rojo, Vicini, Phys.Rev.D83 (2011) 113008
- the generator-level analysis can be quite different w.r.t. the detector-level one

## charged lepton transverse momentum

- ▶ highly sensitive to the details of QCD radiation (and thus also to PDFs)
  - ▶ “simple” experimental determination (accurate lepton energy/momentum calibration)
- moderate impact of detector effects
- the generator level study should provide the correct order of magnitude of the PDF effects

# Sensitivity of the charged-lepton $p_T$ distribution to $M_W$

- since the transverse mass distribution can not be reconstructed at LHCb we focus on the study of the lepton transverse momentum distribution



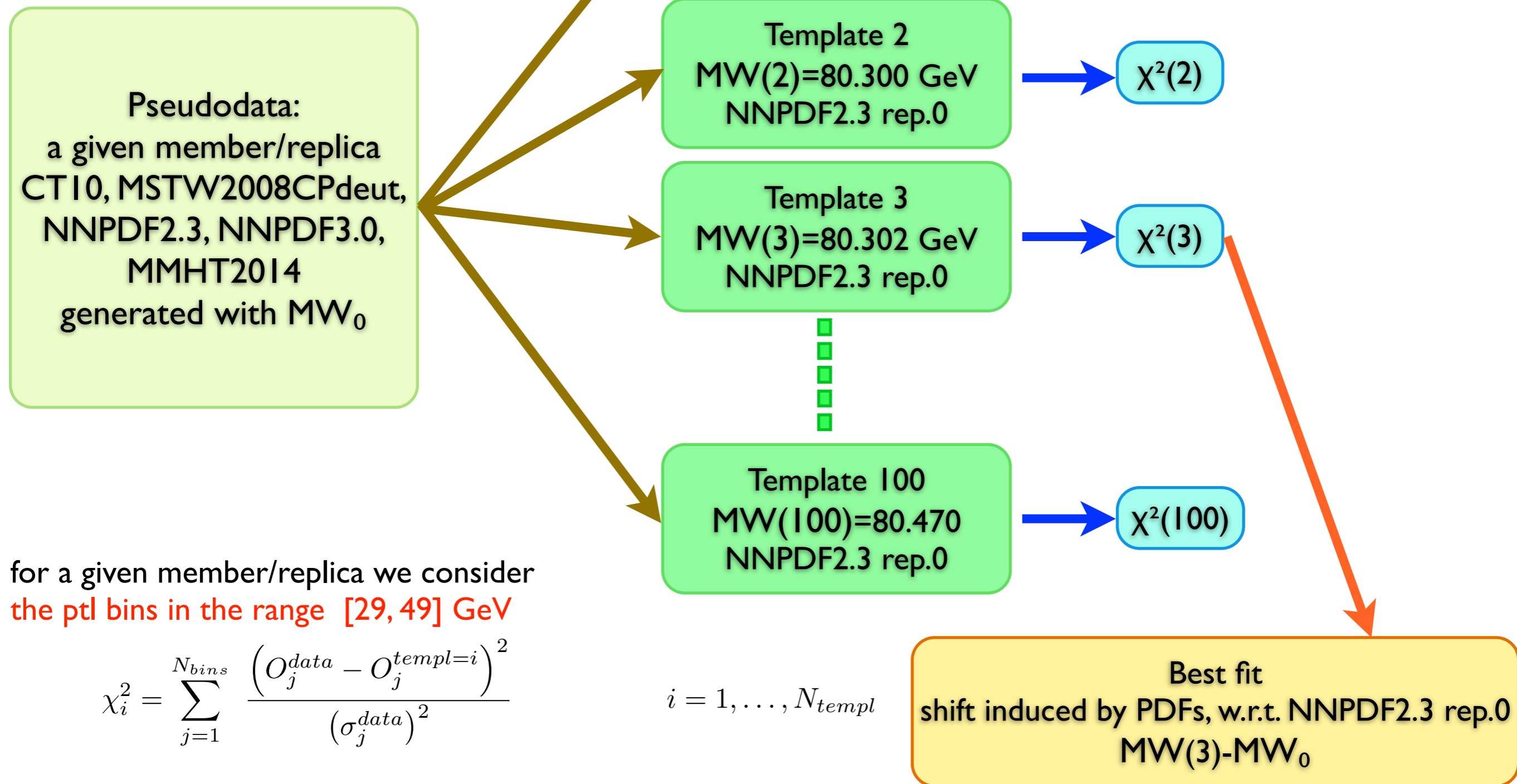
- a sensitivity to  $\Delta M_W = 10$  MeV requires the control of the shape of the distribution at the (sub-) per mill level
- challenging from different points of view
  - experimental
  - MC simulation (statistical fluctuations)
  - theoretical (highly sensitive to the details of QCD radiation description)

# Impact of PDF uncertainties of EW precision measurements

- the extraction of masses and couplings, at hadron colliders, relies on a template fit procedure
- the **uncertainties/ambiguities that affect the evaluation of the templates** are **theoretical systematics** on the final value of the pseudo-observables that we want to extract
- the use of different PDF replicas yields in general a distortion of the template shapes and in turn a different value of the pseudo-observable
- **are PDFs a limiting factor?**
- goals of the present study:
  - 1) estimate of the PDF uncertainty on  $M_W$  extracted from the lepton  $p_t$  distribution
  - 2) study of the dependence of the uncertainty on the acceptance cuts
  - 3) **evaluation of the impact of a  $W$  mass measurement at LHCb in the final LHC  $M_W$  combination**

# The template-fitting procedure

see also Bozzi, Rojo, Vicini, Phys.Rev.D83 (2011) 113008

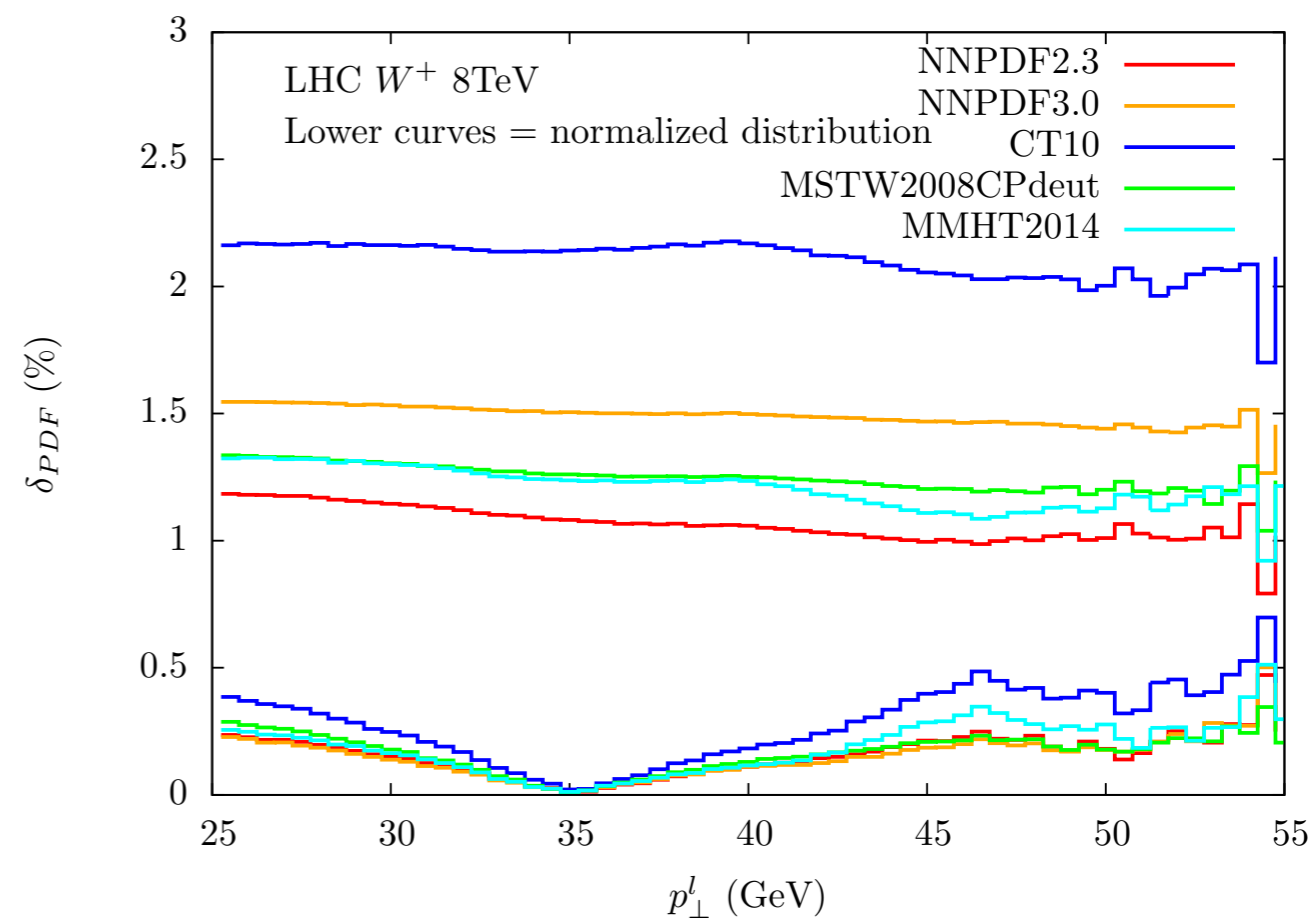
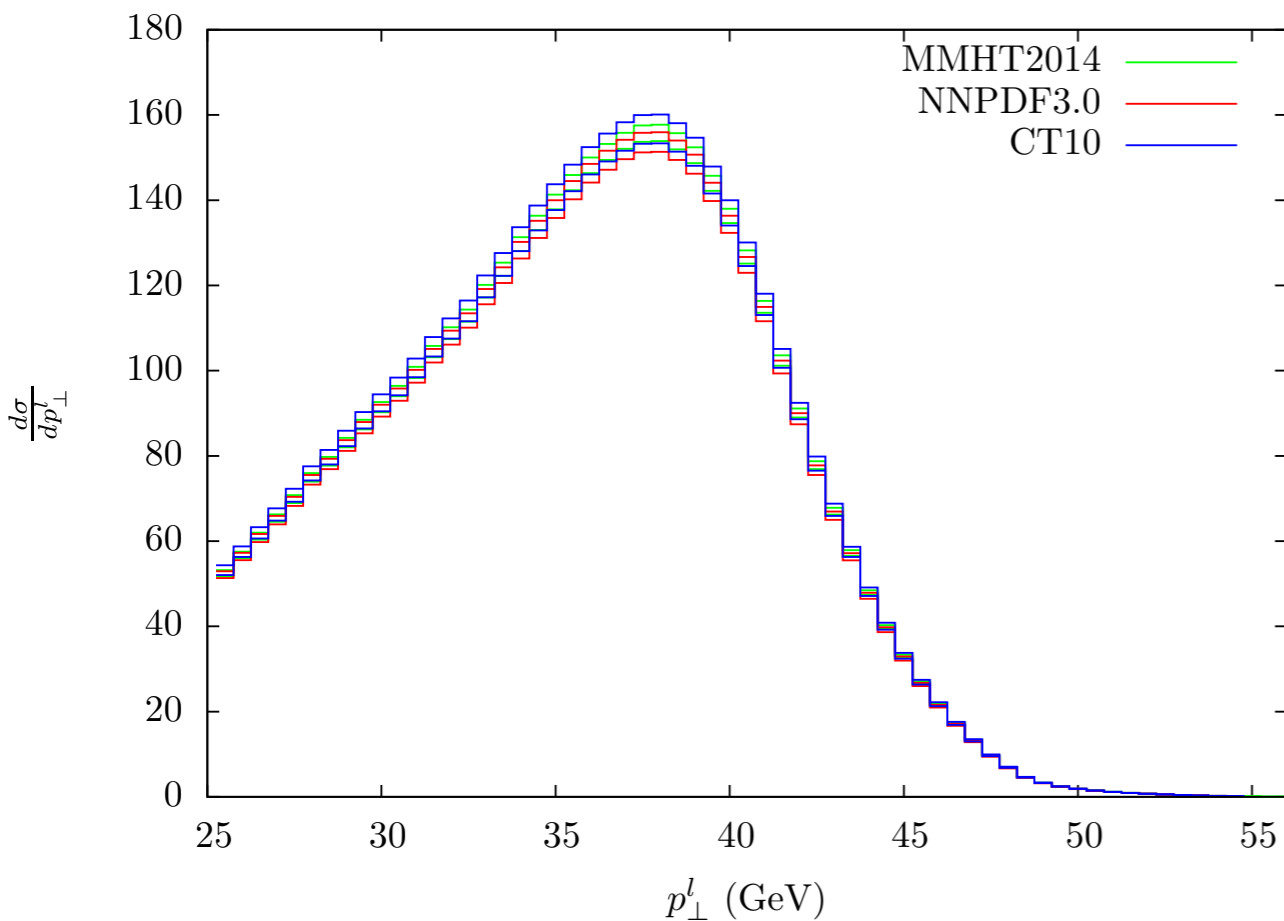


for a given member/replica we consider  
the **ptl bins in the range [29, 49] GeV**

$$\chi_i^2 = \sum_{j=1}^{N_{bins}} \frac{\left(O_j^{data} - O_j^{templ=i}\right)^2}{\left(\sigma_j^{data}\right)^2} \quad i = 1, \dots, N_{templ}$$

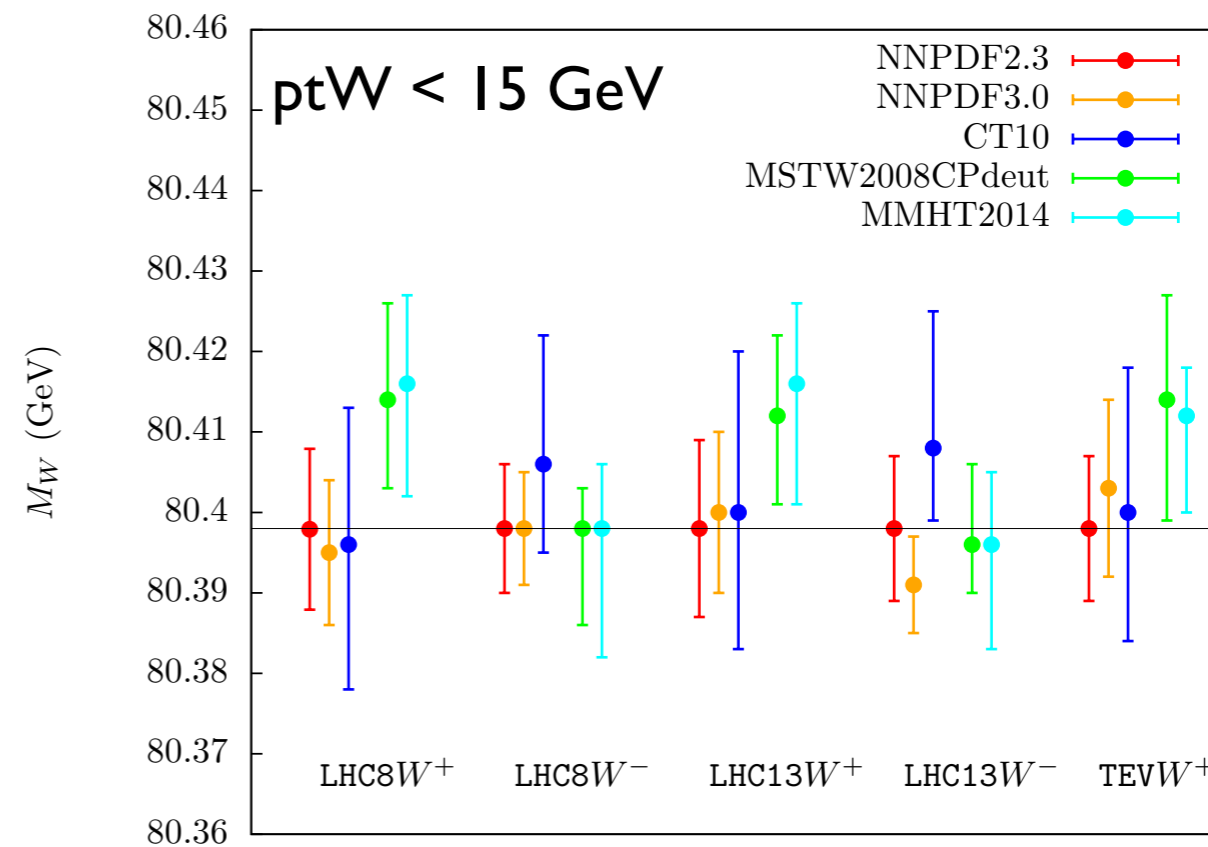
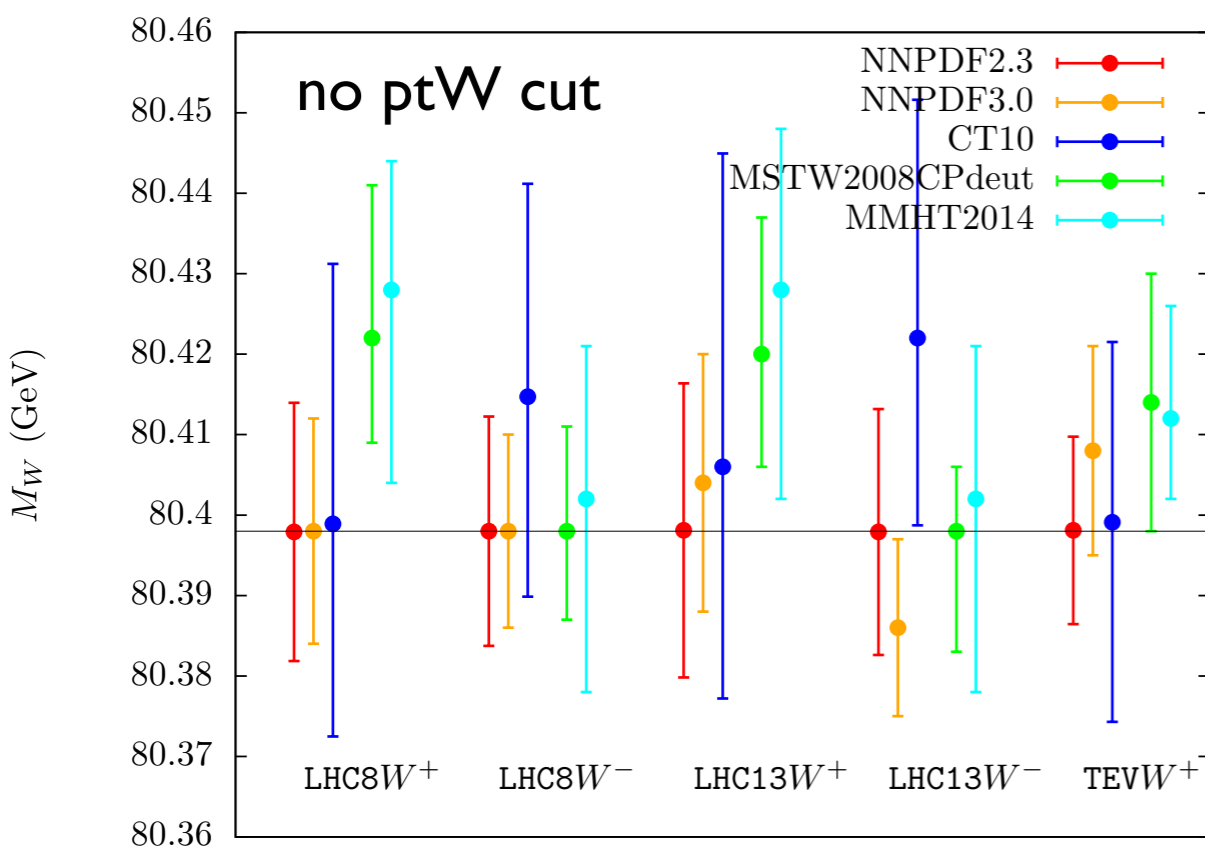
- the template fitting procedure **measures the relative distance** between NNPDF2.3 replica 0 and all the other sets/replicas
- it is an estimate of the difference** that we would find if we would fit the real data with different PDFs

# PDF uncertainty on the lepton $p_T$ distribution



- all simulations with POWHEG matched with PYTHIA 6.4.21  
in these plots standard ATLAS/CMS acceptance cuts
- the use of a normalized distribution reduces the PDF uncertainty,  
leaving only the effects of distortion of the shape  
relevant for the MW determination
- an uncertainty at the few per mill level can still be problematic for a precision measurement

# Numerical results for $M_W$ , with and without a $pt_W$ cut



- in these plots standard ATLAS/CMS acceptance cuts
- the predictions are in general compatible with each other, within their uncertainty bands, with some exceptions
- the uncertainty bands of the 3 sets differ by up to a factor 3; CT10nlo has in general larger uncertainties (C90 factor has been included!)
- spread of the central values  $\Delta_{\text{sets}}$  not negligible, in view of a 10 MeV measurement
- important reduction of the uncertainty when a cut  $PT_W < 15$  GeV is applied
- different results between  $W^+$  and  $W^-$  production



## Numerical results: PDF4LHC envelope and spread of central values

$\delta_{PDF}$  is the half-width of the PDF4LHC envelope

$\Delta_{sets}$  is the spread (max-min) of the central values

### CT10, MSTW2008CPdeut, NNPDF2.3

	no $p_{\perp}^W$ cut		$p_{\perp}^W < 15$ GeV	
	$\delta_{PDF}$ (MeV)	$\Delta_{sets}$ (MeV)	$\delta_{PDF}$ (MeV)	$\Delta_{sets}$ (MeV)
Tevatron 1.96 TeV	27	16	21	15
LHC 8 TeV $W^+$	33	26	24	18
$W^-$	29	16	18	8
LHC 13 TeV $W^+$	34	22	20	14
$W^-$	34	24	18	12

- different description of  $W^+$  (MSTW differs from CT10/NNPDF) and of  $W^-$  (CT10 differs from MSTW/NNPDF)

### MMHT2014, NNPDF3.0

	no $p_{\perp}^W$ cut		$p_{\perp}^W < 15$ GeV	
	$\delta_{PDF}$ (MeV)	$\Delta_{sets}$ (MeV)	$\delta_{PDF}$ (MeV)	$\Delta_{sets}$ (MeV)
Tevatron 1.96 TeV	16	4	13	9
LHC 8 TeV $W^+$	32	33	21	21
$W^-$	22	6	12	0
LHC 13 TeV $W^+$	30	24	18	16
$W^-$	23	16	11	5

- the NNPDF3.0 uncertainties are 15-20% smaller w.r.t. NNPDF2.3 the MMHT2014 unc. are similar to those of MSTW2008CPdeut

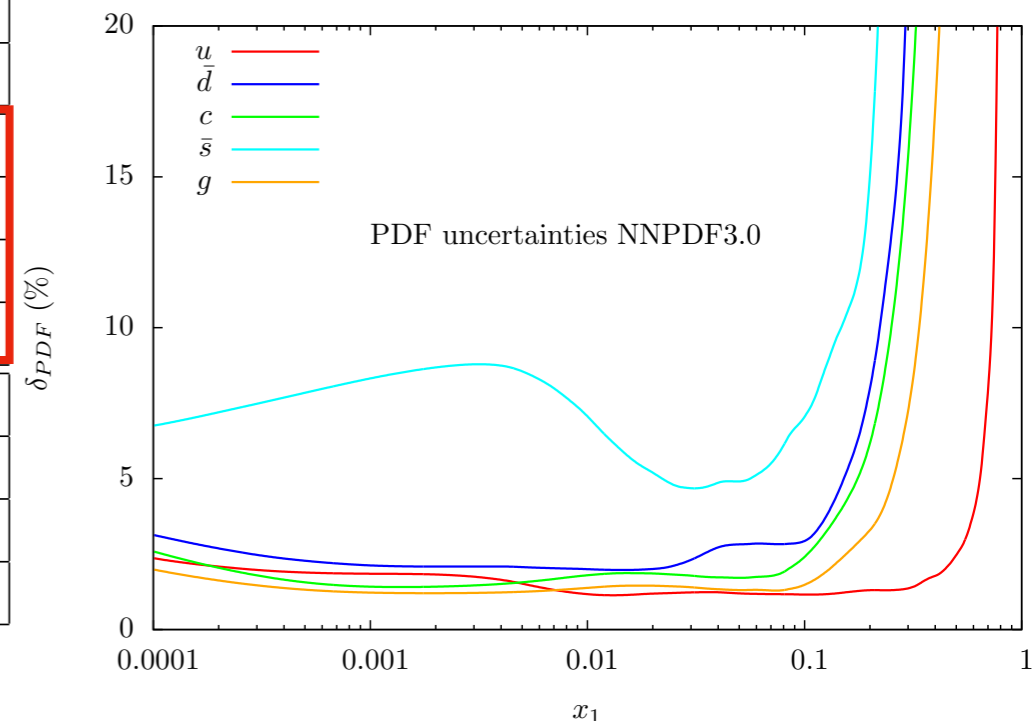
- the NNPDF3.0 results might induce a moderate optimism: i.e. LHC data will help to reduce the PDF uncertainty on MW
- on the other hand the spread of the central values in the  $W^+$  case is the most remarkable feature of the comparison and shows that different parameterizations, based on the same data, yield significantly different results (in a 10 MeV perspective for the final MW error)

# PDF uncertainty affecting MW and acceptance cuts

G.Bozzi, L.Citelli, AV, arXiv:1501.05587

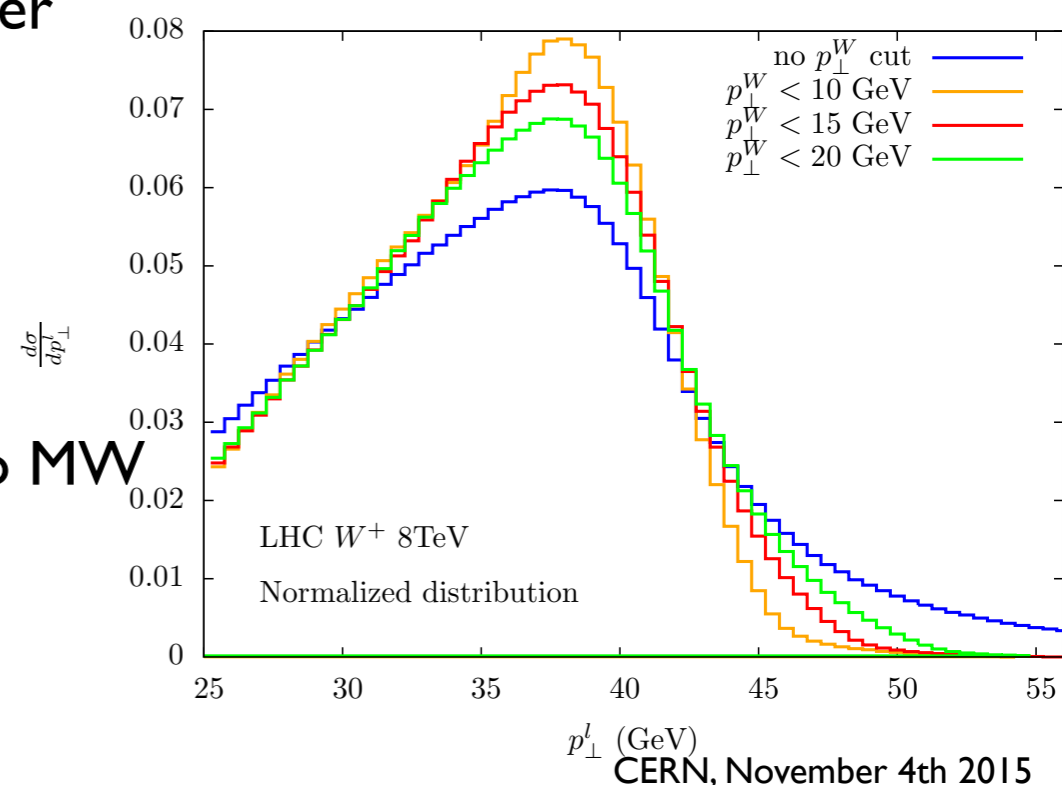
The dependence of the MW PDF uncertainty on the acceptance cuts provides interesting insights

normalized distributions			
cut on $p_{\perp}^W$	cut on $ \eta_l $	CT10	NNPDF3.0
inclusive	$ \eta_l  < 2.5$	$80.400 + 0.032 - 0.027$	$80.398 \pm 0.014$
$p_{\perp}^W < 20$ GeV	$ \eta_l  < 2.5$	$80.396 + 0.027 - 0.020$	$80.394 \pm 0.012$
$p_{\perp}^W < 15$ GeV	$ \eta_l  < 2.5$	$80.396 + 0.017 - 0.018$	$80.395 \pm 0.009$
$p_{\perp}^W < 10$ GeV	$ \eta_l  < 2.5$	$80.392 + 0.015 - 0.012$	$80.394 \pm 0.007$
$p_{\perp}^W < 15$ GeV	$ \eta_l  < 1.0$	$80.400 + 0.032 - 0.021$	$80.406 \pm 0.017$
$p_{\perp}^W < 15$ GeV	$ \eta_l  < 2.5$	$80.396 + 0.017 - 0.018$	$80.395 \pm 0.009$
$p_{\perp}^W < 15$ GeV	$ \eta_l  < 4.9$	$80.400 + 0.009 - 0.004$	$80.401 \pm 0.003$
$p_{\perp}^W < 15$ GeV	$1.0 <  \eta_l  < 2.5$	$80.392 + 0.025 - 0.018$	$80.388 \pm 0.012$



- the PDF uncertainty on the single densities has a steep increase for  $x > 0.1$ 
  - the up density remains accurate up to  $x \sim 0.5$
  - the strange density is  $O(3)$  times less accurate than the other for  $0.001 < x < 0.01$

- the additional cut on  $p_{\perp}^W$  reduces the MW uncertainty
  - suppression of the large- $x$  region
  - steeper shape of the  $p_{\perp}^l$  distribution = more sensitivity to MW



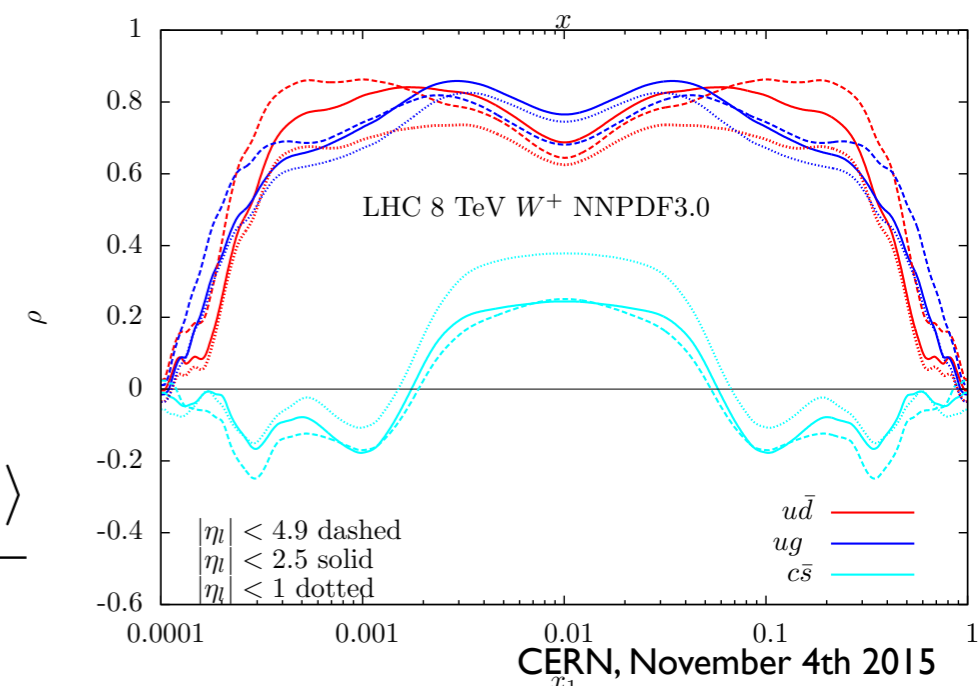
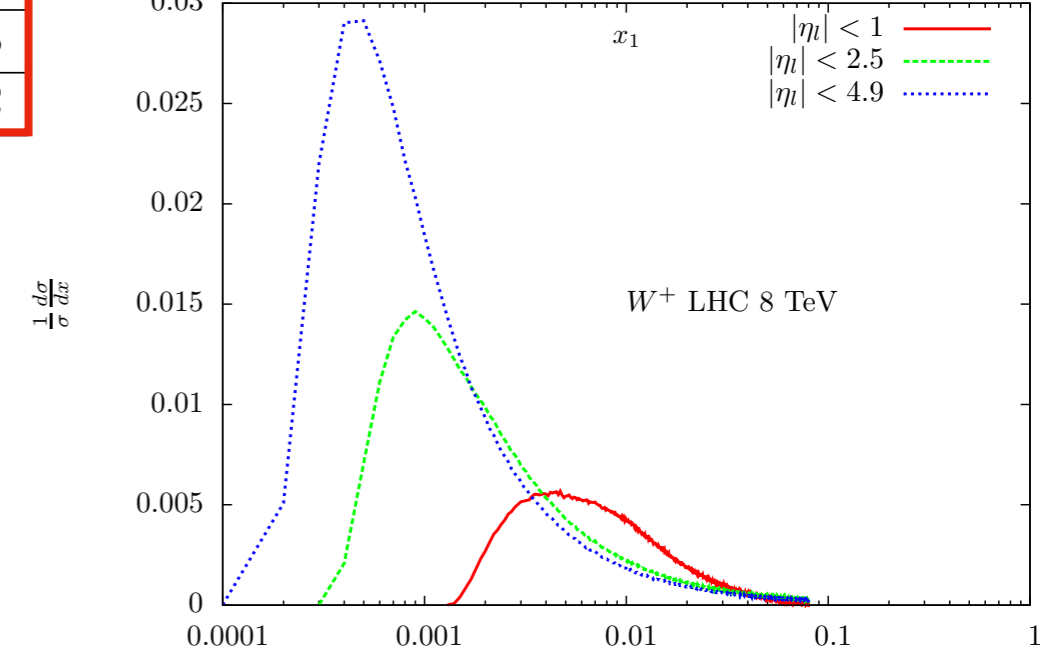
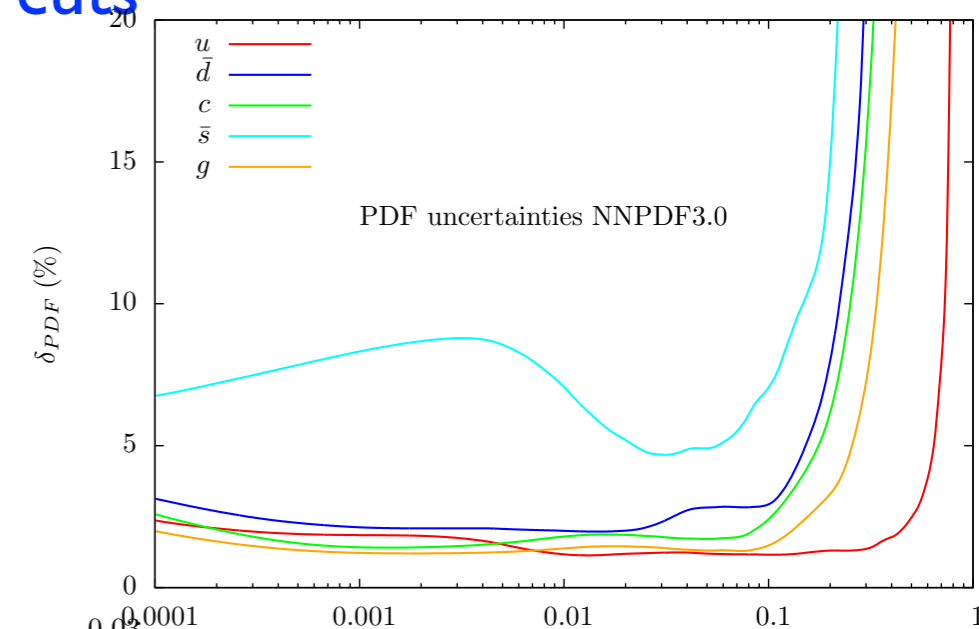
# PDF uncertainty affecting MW and lepton pseudorapidity cuts

G.Bozzi, L.Citelli, AV, arXiv:1501.05587

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- the normalized p<sub>lep</sub> distribution, integrated over the whole lepton-pair rapidity range, does not depend on x and depends very weakly on the PDF replica
- the central pseudorapidity region is the most uncertain
- PDF sum rules →  
non trivial compensations between different rapidity intervals among different flavors  
enlarging symmetrically the eta range → smaller average x → region where the csbar subprocess has negative correlation with the distribution

$$\rho(x, \tau) = \frac{\langle \mathcal{P}_{ij}(x, \tau) \frac{d\sigma}{dp_{\perp}^l} \rangle - \langle \mathcal{P}_{ij}(x, \tau) \rangle \langle \frac{d\sigma}{dp_{\perp}^l} \rangle}{\sigma_{\mathcal{P}_{ij}}^{PDF} \sigma_{d\sigma/dp_{\perp}^l}^{PDF}}$$



# Why measuring MW at LHCb

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

- selecting muons at LHCb with forward pseudorapidities, we probe a different range of partonic  $x$  w.r.t. ATLAS/CMS standard central acceptance

- the lepton  $p_t$  distribution and, in turn, the associated MW by ATLAS/CMS (central cuts) and by LHCb are anticorrelated w.r.t. PDF variations

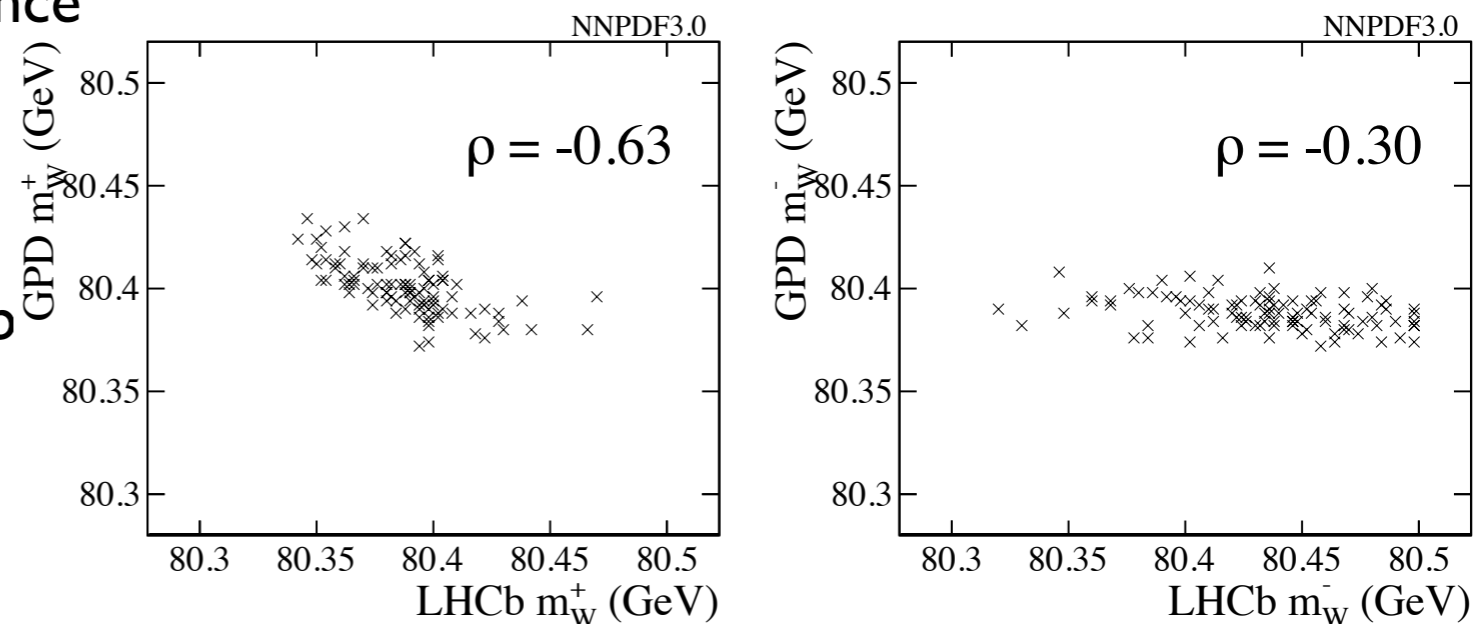
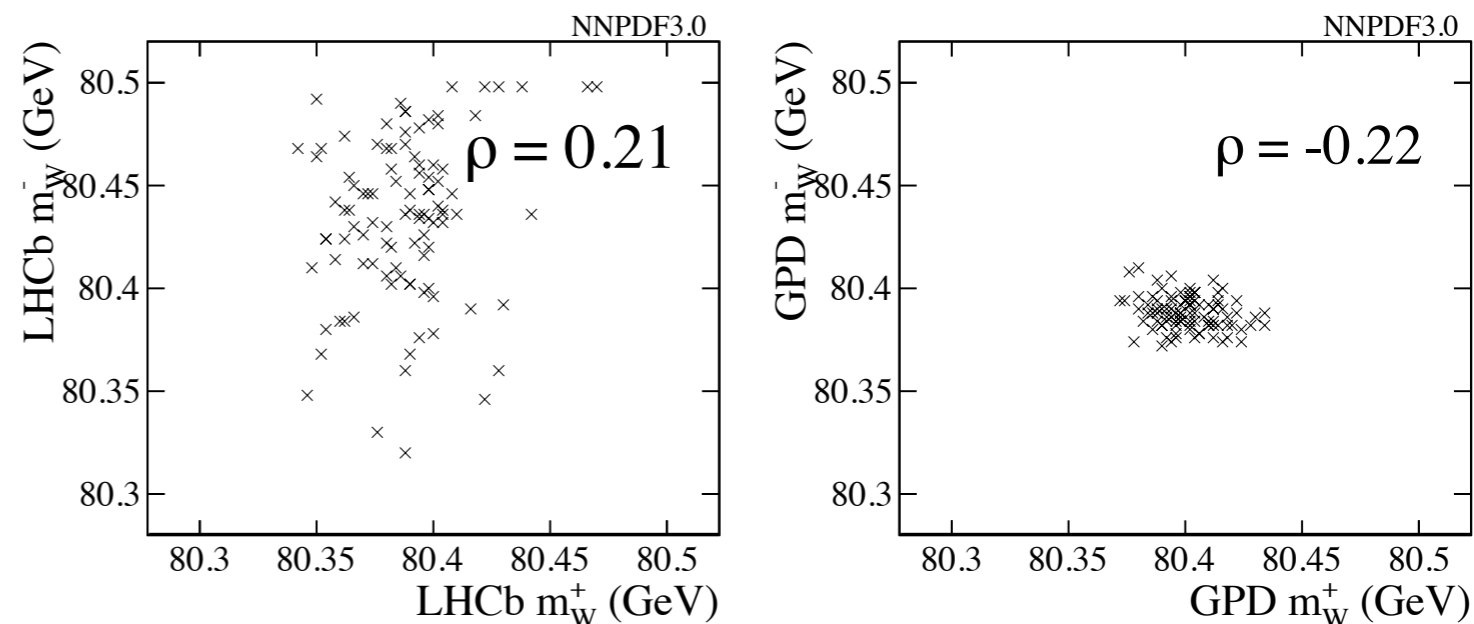


Fig. 1 The fitted  $m_W$  in the GPDs versus LHCb for each NNPDF3.0 set, and for (left)  $W^+$  and (right)  $W^-$ .



- ⇒ in a combination of LHCb with ATLAS/CMS results we could gain a reduction of the final PDF uncertainty

# Impact of a LHCb MW measurement in the combination with ATLAS/CMS results

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

- using the standard acceptance cuts and both W charges  
for ATLAS/CMS (called **G**)  $p_{Tl} > 25 \text{ GeV}$ ,  $|\eta_{lW}| < 2.5$  (both electrons and muons),  $p_{TW} < 15 \text{ GeV}$   
for LHCb (called **L**)  $p_{Tl} > 20 \text{ GeV}$ ,  $2.0 < \eta_{lW} < 4.5$  (only muons), no  $p_{TW}$  cut  
we study the MW determination from the lepton  $p_T$  distribution  
(assuming that a LHCb measurement becomes available)

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- PDF uncertainty on MW according to PDF4LHC (NNPDF3.0, MMHT2014)
  - for **Gs** larger uncertainty in the  $W^+$  case
  - for **Ls** need of a sea quark at large  $x \rightarrow$  large uncertainty e.g. from strange

$$\delta_{\text{PDF}} = \begin{pmatrix} \mathbf{G}^+ & 24.8 \\ \mathbf{G}^- & 13.2 \\ \mathbf{L}^+ & 27.0 \\ \mathbf{L}^- & 49.3 \end{pmatrix}$$

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- correlation matrix  $\rho$  w.r.t. PDF variation of the replicas of the NNPDF3.0 set  
 $\rightarrow$  non negligible anticorrelation

$$\rho = \begin{pmatrix} & \mathbf{G}^+ & \mathbf{G}^- & \mathbf{L}^+ & \mathbf{L}^- \\ \mathbf{G}^+ & 1 & & & \\ \mathbf{G}^- & -0.22 & 1 & & \\ \mathbf{L}^+ & -0.63 & 0.11 & 1 & \\ \mathbf{L}^- & -0.02 & -0.30 & 0.21 & 1 \end{pmatrix}.$$

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- look for a linear combination of all the available **G** and **L** results that minimizes the final PDF uncertainty on MW expressed by the coefficients  $\alpha_i$

$$m_W = \sum_{i=1}^4 \alpha_i m_{W_i}$$



# Impact of a LHCb MW measurement in the combination with ATLAS/CMS results

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

- MW PDF uncertainties from the simulated data

	$G^+$	$G^-$	$L^+$	$L^-$
Envelope	24.8	13.2	27.0	49.3
$\Delta_{\text{sets}}$	20.9	5.7	12.1	22.9

- 2 sets = NNPDF3.0 and MMHT14

	$G^+$	$G^-$	$L^+$	$L^-$
Envelope	29.9	23.5	35.0	84.1
$\Delta_{\text{sets}}$	22.0	23.7	24.0	74.0

- 3 sets = NNPDF3.0 and MMHT14 and CT10

- results for the optimal combination of  $G^\pm$  and  $L^\pm$  results

PDFs	Experiments	$\delta_{\text{PDF}}$ (MeV)	$\alpha$
PDF4LHC(2-sets)	2×GPD	10.5	(0.26, 0.74, 0, 0)
PDF4LHC(2-sets)	2×GPD + LHCb	7.7	(0.30, 0.45, 0.21, 0.04)
PDF4LHC(3-sets)	2×GPD	16.9	(0.50, 0.50, 0, 0)
PDF4LHC(3-sets)	2×GPD + LHCb	12.7	(0.43, 0.41, 0.11, 0.04)
NNPDF30	2×GPD	5.2	(0.50, 0.50, 0, 0)
NNPDF30	2×GPD + LHCb	3.6	(0.35, 0.47, 0.16, 0.02)
MMHT2014	2×GPD	9.2	(0.45, 0.55, 0, 0)
MMHT2014	2×GPD + LHCb	4.6	(0.39, 0.14, 0.46, 0)
CT10	2×GPD	11.6	(0.33, 0.67, 0, 0)
CT10	2×GPD + LHCb	6.3	(0.38, 0.20, 0.40, 0.03)

- the inclusion of LHCb results yields a reduction of the PDF uncertainty of O(30-40%) on the envelope stronger reduction for the individual sets MMHT2014 and for CT10

# Feasibility of a LHCb MW measurement

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

## • statistical sensitivity

- in  $1\text{fb}^{-1}$  of luminosity at Run-I collected  
550k  $W^+$  and 350k  $W^-$  with 70% of purity  
60k of candidate  $Z$  with almost perfect purity

→ extrapolation of the signal yield  
of the full Run-I and Run-II datasets

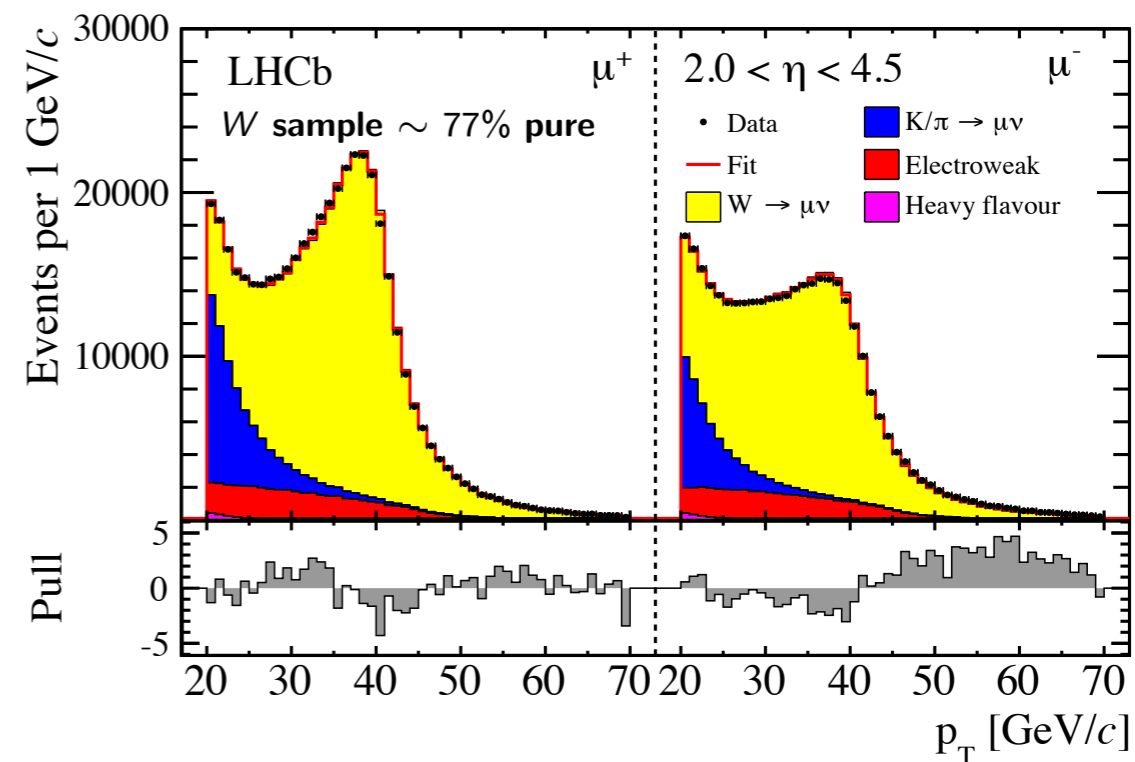
## • backgrounds

low  $p_T$ : pions/kaons decays

intermediate/large  $p_T$ :  $Z/\gamma^* \rightarrow \mu\mu$  with one muon lost

exponential parameterization, reproducing the estimates of arXiv:1505.07024

- estimate of the statistical error obtained from signal+background fitting 500 pseudo-experiments



## • muon momentum scale calibration

- at LHCb very precise measurement of  $b$  and  $c$  hadron masses  
(arXiv:1302.1072, arXiv:1304.6865)

momentum resolution between 0.2% and 0.8%

momentum scale uncertainty of  $3 \cdot 10^{-4}$

- the full dataset, with 700k  $Z$  events,

→ calibration also at high  $p_T$

- at Tevatron calibration using  $J/\Psi$  and  $\Upsilon$

	Run-I $3\text{fb}^{-1}$		Run-II $7\text{fb}^{-1}$	
	$W^+$	$W^-$	$W^+$	$W^-$
Signal yields, $\times 10^6$	1.2	0.7	5.4	3.4
$Z/\gamma^*$ background, ( $B/S$ )	0.15	0.15	0.15	0.15
QCD background, ( $B/S$ )	0.15	0.15	0.15	0.15
$\delta m_W$ (MeV)				
Statistical	19	29	9	12
Momentum scale	7	7	4	4
Quadrature sum	20	30	10	13

# Prospects for a combination of ATLAS, CMS and LHCb results

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

- we assume that the following set uncertainties will be confirmed at the end of run-II (same PDFs as today, hypothetical experimental errors)

$$\delta m_W^i = \begin{pmatrix} \mathbf{G}^+ & (7_{\text{exp}} \pm 25_{\text{PDF}}) \text{ MeV} \\ \mathbf{G}^- & (7_{\text{exp}} \pm 13_{\text{PDF}}) \text{ MeV} \\ \mathbf{L}^+ & (10_{\text{exp}} \pm 28_{\text{PDF}}) \text{ MeV} \\ \mathbf{L}^- & (13_{\text{exp}} \pm 49_{\text{PDF}}) \text{ MeV} \end{pmatrix} \cdot$$

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- we vary each of the assumed values in a range from 0 to 2

we then check how robust is the reduction of the PDF error under these variations

Scenario	Experiments	$\delta m_W$ (MeV)			$\alpha$
		Tot	Exp	PDF	
Default	2×GPD + LHCb	9.0	4.7	7.7	(0.30, 0.44, 0.22, 0.04)
Default	1×GPD + LHCb	10.1	6.5	7.7	(0.31, 0.40, 0.25, 0.04)
Default	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
PDF4LHC(3-sets)	2×GPD + LHCb	13.6	4.8	12.7	(0.43, 0.41, 0.12, 0.04)
PDF4LHC(3-sets)	1×GPD + LHCb	14.6	7.3	12.7	(0.43, 0.40, 0.12, 0.04)
PDF4LHC(3-sets)	2×GPD	17.7	5.5	16.9	(0.50, 0.50, 0, 0)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD + LHCb	8.7	4.0	7.7	(0.31, 0.41, 0.24, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	1×GPD + LHCb	9.8	5.9	7.9	(0.31, 0.37, 0.28, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD + LHCb	7.9	1.9	7.7	(0.29, 0.48, 0.19, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	1×GPD + LHCb	7.9	1.9	7.7	(0.29, 0.48, 0.19, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD	10.5	0.1	10.5	(0.26, 0.74, 0, 0)
$\delta_{\text{PDF}} = 0$	2×GPD + LHCb	4.6	4.6	0.0	(0.34, 0.34, 0.22, 0.10)
$\delta_{\text{PDF}} = 0$	1×GPD + LHCb	5.8	5.8	0.0	(0.23, 0.23, 0.37, 0.17)
$\delta_{\text{PDF}} = 0$	2×GPD	5.5	5.5	0.0	(0.50, 0.50, 0, 0)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	2×GPD + LHCb	9.6	5.6	7.7	(0.29, 0.50, 0.17, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	1×GPD + LHCb	10.8	7.6	7.7	(0.30, 0.46, 0.20, 0.05)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	2×GPD + LHCb	11.2	7.9	8.0	(0.32, 0.35, 0.29, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	1×GPD + LHCb	13.9	10.5	9.0	(0.31, 0.26, 0.37, 0.05)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	2×GPD	15.6	11.5	10.6	(0.32, 0.68, 0, 0)
$\delta_{\text{PDF}} \times 2$	2×GPD + LHCb	16.0	4.7	15.3	(0.30, 0.45, 0.21, 0.04)
$\delta_{\text{PDF}} \times 2$	1×GPD + LHCb	16.7	6.7	15.3	(0.30, 0.44, 0.22, 0.04)
$\delta_{\text{PDF}} \times 2$	2×GPD	21.7	5.9	20.9	(0.27, 0.73, 0, 0)

# Prospects for a combination of ATLAS, CMS and LHCb results

G.Bozzi, L.Citelli, M.Vesterinen, AV, arXiv:1508.06954

- we assume that the following set uncertainties will be confirmed at the end of run-II (same PDFs as today, hypothetical experimental errors)

$$\delta m_W^i = \begin{pmatrix} \mathbf{G}^+ (7_{\text{exp}} \pm 25_{\text{PDF}}) \text{ MeV} \\ \mathbf{G}^- (7_{\text{exp}} \pm 13_{\text{PDF}}) \text{ MeV} \\ \mathbf{L}^+ (10_{\text{exp}} \pm 28_{\text{PDF}}) \text{ MeV} \\ \mathbf{L}^- (13_{\text{exp}} \pm 49_{\text{PDF}}) \text{ MeV} \end{pmatrix}.$$

- we vary each of the assumed values in a range from 0 to 2

we then check how robust is the reduction of the PDF error under these variations

- the introduction of LHCb **always yields a reduction of the PDF error** and in turn of the total error of O(25-40%) w.r.t. the 2 GPDs case

- the combination of LHCb with 1 GPD is more convenient**, in a PDF perspective, than the sum of 2 GPDs

- when including CT10, the impact of LHCb on the combination is stronger

Scenario	Experiments	$\delta m_W$ (MeV)			$\alpha$
		Tot	Exp	PDF	
Default	2×GPD + LHCb	9.0	4.7	7.7	(0.30, 0.44, 0.22, 0.04)
Default	1×GPD + LHCb	10.1	6.5	7.7	(0.31, 0.40, 0.25, 0.04)
Default	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
PDF4LHC(3-sets)	2×GPD + LHCb	13.6	4.8	12.7	(0.43, 0.41, 0.12, 0.04)
PDF4LHC(3-sets)	1×GPD + LHCb	14.6	7.3	12.7	(0.43, 0.40, 0.12, 0.04)
PDF4LHC(3-sets)	2×GPD	17.7	5.5	16.9	(0.50, 0.50, 0, 0)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD + LHCb	8.7	4.0	7.7	(0.31, 0.41, 0.24, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	1×GPD + LHCb	9.8	5.9	7.9	(0.31, 0.37, 0.28, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} = 0$	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD + LHCb	7.9	1.9	7.7	(0.29, 0.48, 0.19, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	1×GPD + LHCb	7.9	1.9	7.7	(0.29, 0.48, 0.19, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} = 0$	2×GPD	10.5	0.1	10.5	(0.26, 0.74, 0, 0)
$\delta_{\text{PDF}} = 0$	2×GPD + LHCb	4.6	4.6	0.0	(0.34, 0.34, 0.22, 0.10)
$\delta_{\text{PDF}} = 0$	1×GPD + LHCb	5.8	5.8	0.0	(0.23, 0.23, 0.37, 0.17)
$\delta_{\text{PDF}} = 0$	2×GPD	5.5	5.5	0.0	(0.50, 0.50, 0, 0)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	2×GPD + LHCb	9.6	5.6	7.7	(0.29, 0.50, 0.17, 0.04)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	1×GPD + LHCb	10.8	7.6	7.7	(0.30, 0.46, 0.20, 0.05)
$\delta_{\text{exp}}^{\text{LHCb}} \times 2$	2×GPD	12.0	5.8	10.5	(0.28, 0.72, 0, 0)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	2×GPD + LHCb	11.2	7.9	8.0	(0.32, 0.35, 0.29, 0.04)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	1×GPD + LHCb	13.9	10.5	9.0	(0.31, 0.26, 0.37, 0.05)
$\delta_{\text{exp}}^{\text{GPD}} \times 2$	2×GPD	15.6	11.5	10.6	(0.32, 0.68, 0, 0)
$\delta_{\text{PDF}} \times 2$	2×GPD + LHCb	16.0	4.7	15.3	(0.30, 0.45, 0.21, 0.04)
$\delta_{\text{PDF}} \times 2$	1×GPD + LHCb	16.7	6.7	15.3	(0.30, 0.44, 0.22, 0.04)
$\delta_{\text{PDF}} \times 2$	2×GPD	21.7	5.9	20.9	(0.27, 0.73, 0, 0)

## Potential bottlenecks

- the measurement of  $M_W$  from the lepton  $p_t$  distribution strongly relies on the knowledge of the neutral-current Drell-Yan, in particular the  $p_t Z$  distribution, to model  $p_t W$  and eventually to simulate the lepton  $p_t$
  - the lepton  $p_t$  distribution is extremely sensitive to the details of QCD radiation, in particular at low  $p_t$  values ( a distortion at the few per mil level yields  $O(20 \text{ MeV})$   $M_W$  shift )
  - the assumption that the information obtained from the  $Z$  is universal and can be transferred to the  $W$  is violated by several factors:  
different parton-parton luminosities (and heavy-quark content), different energy scales, dependence of  $p_t Z$  modeling on the lepton-pair rapidity, EW corrections
- ⇒ a dedicated study of  $p_t Z$  and of the  $p_t W \leftrightarrow p_t Z$  interplay at LHCb is needed

## Conclusions

- PDF uncertainties are a potential bottleneck of a precision MW measurement at the LHC in view of a final precision goal at the  $O(10 \text{ MeV})$  level
  - the measurement of MW from the lepton pt distribution at LHCb and its combination with the ATLAS/CMS results can help to reduce the PDF uncertainty by 25-40% thanks to the anticorrelation w.r.t. PDFs of the two sets of results
  - in a preliminary study we tried to assess the feasibility of an MW measurement at the LHCb and we checked the robustness of the PDF uncertainty reduction under pessimistic increases of the different error sources
  - these encouraging results motivate further, more detailed studies of all the requirements needed to bring the experimental error in the 10-15 MeV ballpark and to control the other theoretical systematics (beyond PDFs) at a similar level
- special attention should be payed to the study of the ptZ distribution at LHCb and to the ptW/ptZ interplay
- the MW measurement at LHCb could offer a rich set of informations, complementary to those from ATLAS/CMS for the precision measurement of EW parameters at the LHC!

back-up



# Numerical results, with and without a PTW cut

absolute distributions					
collider/channel	CT10	MSTW2008CPdeut	NNPDF2.3	NNPDF3.0	MMHT2014
Tevatron, $W^+$	$80.406 + 0.043 - 0.046$	$80.428 + 0.025 - 0.017$	$80.400 \pm 0.030$	$80.427 \pm 0.018$	$80.430 + 0.022 - 0.022$
LHC 8 TeV, $W^+$	$80.394 + 0.040 - 0.029$	$80.422 + 0.025 - 0.016$	$80.398 \pm 0.020$	$80.406 \pm 0.019$	$80.428 + 0.027 - 0.022$
$W^-$	$80.444 + 0.055 - 0.062$	$80.390 + 0.038 - 0.036$	$80.398 \pm 0.030$	$80.441 \pm 0.027$	$80.404 + 0.041 - 0.048$
LHC 13 TeV, $W^+$	$80.396 + 0.045 - 0.034$	$80.416 + 0.020 - 0.020$	$80.398 \pm 0.022$	$80.414 \pm 0.022$	$80.422 + 0.030 - 0.024$
$W^-$	$80.416 + 0.088 - 0.065$	$80.374 + 0.044 - 0.033$	$80.398 \pm 0.031$	$80.426 \pm 0.037$	$80.384 + 0.037 - 0.049$
normalized distributions					
collider/channel	CT10	MSTW2008CPdeut	NNPDF2.3	NNPDF3.0	MMHT2014
Tevatron, $W^+$	$80.400 + 0.022 - 0.025$	$80.414 + 0.016 - 0.016$	$80.398 \pm 0.012$	$80.408 \pm 0.013$	$80.412 + 0.014 - 0.010$
LHC 8 TeV, $W^+$	$80.398 + 0.032 - 0.026$	$80.424 + 0.014 - 0.019$	$80.398 \pm 0.016$	$80.395 \pm 0.014$	$80.428 + 0.016 - 0.024$
$W^-$	$80.416 + 0.026 - 0.025$	$80.398 + 0.011 - 0.014$	$80.398 \pm 0.014$	$80.396 \pm 0.012$	$80.402 + 0.019 - 0.024$
LHC 13 TeV, $W^+$	$80.406 + 0.039 - 0.029$	$80.420 + 0.017 - 0.014$	$80.398 \pm 0.018$	$80.404 \pm 0.016$	$80.428 + 0.020 - 0.026$
$W^-$	$80.422 + 0.030 - 0.023$	$80.398 + 0.008 - 0.015$	$80.398 \pm 0.015$	$80.386 \pm 0.011$	$80.402 + 0.019 - 0.024$
absolute distributions, additional cut $p_{\perp}^W < 15$ GeV					
collider/channel	CT10	MSTW2008CPdeut	NNPDF2.3	NNPDF3.0	MMHT2014
Tevatron, $W^+$	$80.412 + 0.024 - 0.024$	$80.424 + 0.018 - 0.017$	$80.399 \pm 0.014$	$80.420 \pm 0.014$	$80.426 + 0.009 - 0.021$
LHC 8 TeV, $W^+$	$80.392 + 0.026 - 0.021$	$80.414 + 0.020 - 0.011$	$80.398 \pm 0.015$	$80.403 \pm 0.014$	$80.418 + 0.019 - 0.017$
$W^-$	$80.422 + 0.039 - 0.034$	$80.394 + 0.019 - 0.023$	$80.399 \pm 0.018$	$80.423 \pm 0.017$	$80.400 + 0.023 - 0.028$
LHC 13 TeV, $W^+$	$80.392 + 0.028 - 0.022$	$80.410 + 0.012 - 0.016$	$80.398 \pm 0.016$	$80.408 \pm 0.014$	$80.414 + 0.016 - 0.019$
$W^-$	$80.408 + 0.042 - 0.037$	$80.386 + 0.019 - 0.021$	$80.398 \pm 0.016$	$80.410 \pm 0.018$	$80.388 + 0.021 - 0.025$
normalized distributions, additional cut $p_{\perp}^W < 15$ GeV					
collider/channel	CT10	MSTW2008CPdeut	NNPDF2.3	NNPDF3.0	MMHT2014
Tevatron, $W^+$	$80.400 + 0.018 - 0.016$	$80.414 + 0.013 - 0.015$	$80.399 \pm 0.010$	$80.403 \pm 0.011$	$80.412 + 0.006 - 0.012$
LHC 8 TeV, $W^+$	$80.396 + 0.017 - 0.018$	$80.414 + 0.012 - 0.011$	$80.398 \pm 0.011$	$80.395 \pm 0.009$	$80.416 + 0.011 - 0.014$
$W^-$	$80.406 + 0.016 - 0.011$	$80.398 + 0.005 - 0.012$	$80.398 \pm 0.010$	$80.398 \pm 0.007$	$80.398 + 0.008 - 0.016$
LHC 13 TeV, $W^+$	$80.400 + 0.020 - 0.017$	$80.412 + 0.010 - 0.011$	$80.398 \pm 0.012$	$80.400 \pm 0.010$	$80.416 + 0.010 - 0.015$
$W^-$	$80.408 + 0.017 - 0.009$	$80.396 + 0.010 - 0.006$	$80.399 \pm 0.010$	$80.391 \pm 0.006$	$80.396 + 0.009 - 0.013$

# Checks

- in Bozzi, Rojo, Vicini, Phys.Rev.D83 (2011) 113008  
we studied the PDF impact on  $M_W$  extracted from the lepton-pair transverse mass distribution using DYNNLO with NLO-QCD accuracy  
  
a fixed-order simulation is sufficient to describe the MT but not the ptl distributions
- we reproduce with POWHEG+PYTHIA the DYNNLO results for MT  
(but now we can also study the ptl distribution)
- the PDF uncertainty on  $M_W$  from the MT distribution is smaller than the one from the ptl case  
but there can be important differences in the estimate between a generator level and a detector level estimate

# Setup of the study

- PDF sets: CT10nlo, MSTW2008 (for comparison with previous studies), MSTW2008CPdeut, MMHT2014, NNPDF2.3\_nlo\_0119, NNPDF3.0
- simulation code: POWHEG + PYTHIA 6.4.21 (pure QCD, resummation effects via Parton Shower)
- Tevatron 1.96 TeV, LHC 8, 13, 33, 100 TeV
- acceptance cuts (called basic):  $p_{Tl} > 25 \text{ GeV}$ ,  $E_{T\text{miss}} > 25 \text{ GeV}$   
 $|\eta_{l}| < 1.0$  (Tevatron),  $|\eta_{l}| < 2.5$  (LHC)
- **additional** acceptance cuts:  $p_{TW} < 15 \text{ GeV}$ ,  $M_{T} < 100 \text{ GeV}$   
further analysis in rapidity bins
- study of absolute and of normalized distributions

# Reweighting

- MC fluctuations at the per mill level are still present also in simulations with 1 billion of events when bin sizes have to be small
- the estimate of PDF uncertainty on MW requires to appreciate the difference of the value of the distribution in each bin
  - the use of fully correlated distributions reduces the sensitivity to MC fluctuations
- the weights for different templates/replicas have been generated in one single simulation

given the weight  $w_0$  of one event, computed with  $MW_0$  and with NNPDF2.3 replica 0, we rescale

$$w_0 \rightarrow w_j = w_0 \frac{(\hat{s} - m_{W0}^2)^2 + \Gamma_W^2 m_{W0}^2}{(\hat{s} - m_{W,j}^2)^2 + \Gamma_W^2 m_{W,j}^2} \quad \text{template } j$$
$$w_0 \rightarrow w_i = w_0 \frac{f_i(x_1)g_i(x_2)}{f_0^{NNPDF}(x_1)g_0^{NNPDF}(x_2)} \quad \text{replica } i$$

- this reweighting is almost NLO-QCD accurate:  
a dependence on the PDF via the POWHEG Sudakov is not included in this approach  
(see talk by P. Nason)