

Calculating four-loop massless propagators with Forcer

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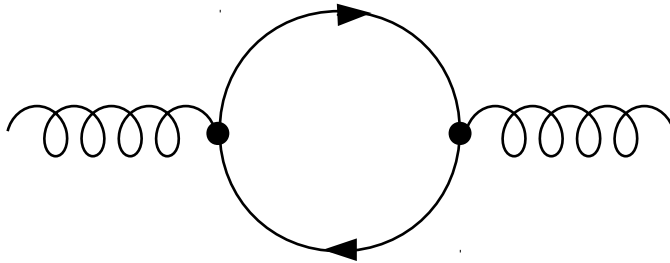
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 - Massless propagator-type integrals
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- “Mincer” approach
 - How did people compute them up to 3-loops?
- Extension to the 4-loop level: our program “Forcer”
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- Summary

Introduction

Introduction

- Massless propagator-type Feynman integrals
 - Example: gluon self-energy diagram

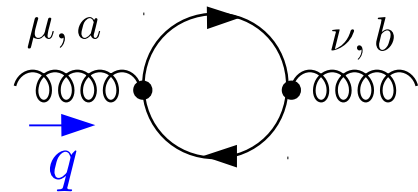


gluon: massless
(light) quark: massless

- In pQCD, after factorizing a hard part and a soft part, the hard part contains only massless particles
- Evaluating massless propagator-type integrals (within dimensional regularization) has been one of basic components in pQCD since its early days

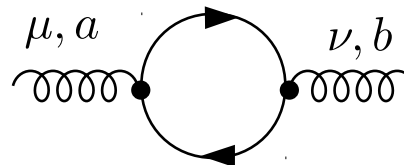
Massless propagator-type integrals

- Without loss of generality, we can consider only “scalar” integrals



$$= A(q^2 g_{\mu\nu} - q_\mu q_\nu) \delta^{ab}$$

- Projection operator $P^{\mu\nu,ab}$ to extract A : all indices (Lorentz, color) contracted



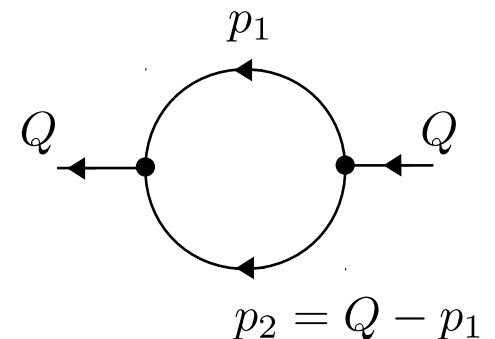
$$A = P^{\mu\nu,ab} \times \text{diagram} \quad P^{\mu\nu,ab} = \frac{q^\mu q^\nu}{q^2} \delta^{ab}$$

- At the 1-loop level, the following family of Feynman integrals

$$F(\underbrace{n_1, n_2}_{\text{indices: powers of the denominators}}) = \int \frac{d^D p_1}{(2\pi)^D} \frac{1}{(p_1^2)^{n_1} (p_2^2)^{n_2}}$$

indices: powers of the denominators

- Dimensional regularization: $D = 4 - 2\epsilon$
- The external momenta Q is off-shell: $Q^2 \neq 0$



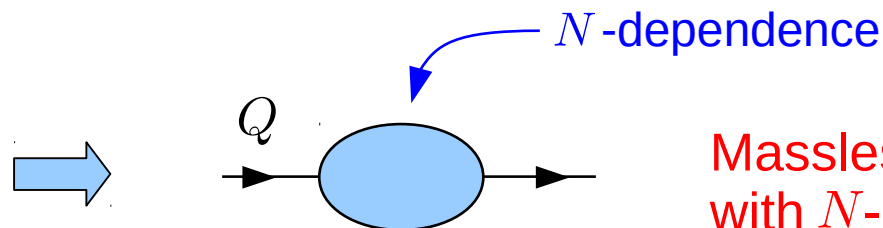
Physics motivation

- Much effort has been made for partonic Higgs production cross section in N³LO QCD for the LHC Run2

Anastasiou, Duhr, Dulat, Herzog, Mistlberger '15 (soft expansion, over 35 terms)
 Anzai, Hasselhuhn, Höschele, Hoff, Kilgore, Steinhauser, TU '15 (only qq', exact in m_h^2/s)
 and many works

- In principle, N³LO partonic cross section should be used with N³LO parton distribution function \Rightarrow 4-loop splitting functions (unknown)
- In Mellin N -space, computed from

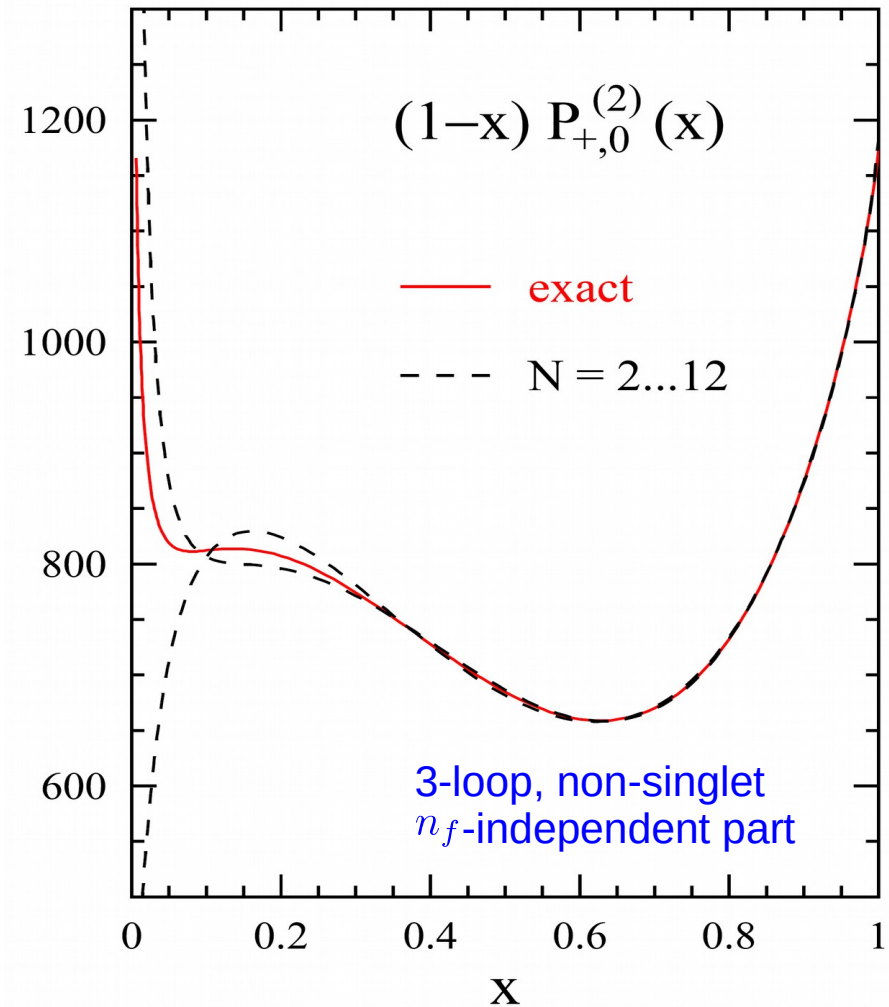
$$\frac{Q^{\{\mu_1 \dots \mu_N\}}}{N!} \frac{\partial^N}{\partial P^{\mu_1} \dots \partial P^{\mu_N}} \left. \begin{array}{c} P \rightarrow \\ \left(\text{blob} \right) \\ Q \rightarrow \end{array} \right|_{P=0}$$



Massless propagator-type integrals with N -dependence

Physics motivation


- Fixing $N = 2, 4, 6, \dots$, one can compute them as massless propagator-type integrals. Their information gives an approximation to / upper bounds of uncertainty of the result (especially at large x)
- If we have software that can compute such integrals at the 4-loop level more efficiently, we can go for more higher N , get a better estimation
known up to $N=4$ (NS) at 4-loop
[Velizhanin '14]



From Moch, Vermaseren, Vogt
NPB688(2004)101 [arXiv:hep-
ph/0403192]

“Mincer” approach

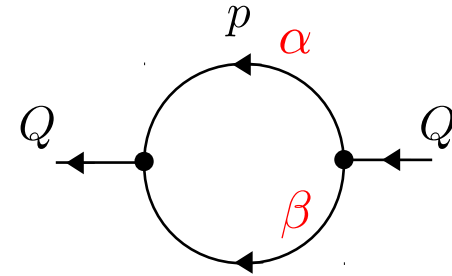
Feynman integral calculus

- Integration-by-parts identities (IBPs): the Gauss theorem in D-dimension gives linear identities among Feynman integrals [Chetyrkin, Tkachov '81]
- The “standard” way of Feynman integral calculus
 - various integrals  (irreducible) master integrals (MIs)
 - (1) reduction via IBPs
 - (2) evaluation of MIs
- Generic reduction algorithms (for any processes)
 - Laporta algorithm, Baikov's method, Lee's “LiteRed”
[Laporta '01] [Baikov '96; '05] [Lee '12; '13]
- “Mincer” algorithm [Chetyrkin, Tkachov '81]
 - Specialized reduction for massless propagator-type integrals (up to 3-loops), very efficient

General 1-loop formula

- General formula for **arbitrary indices α and β**

$$\int \frac{d^D p}{(2\pi)^D} \frac{p^{\mu_1} \dots p^{\mu_n}}{(p^2)^\alpha [(Q-p)^2]^\beta} \quad \text{allowing numerators } (n \geq 0)$$

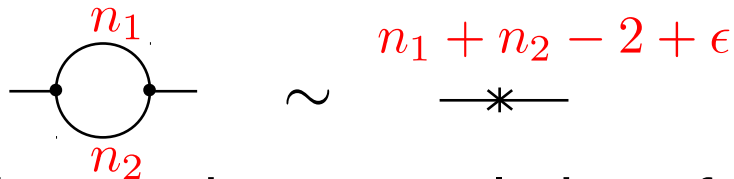


$$= \frac{1}{(4\pi)^2} \frac{1}{(Q^2)^{\alpha+\beta-2+\epsilon}} \sum_{\sigma=0}^{\lfloor n/2 \rfloor} G(\alpha, \beta, n, \sigma) (Q^2)^\sigma \left[\frac{1}{\sigma!} \left(\frac{\square_p}{4} \right)^\sigma p^{\mu_1} \dots p^{\mu_n} \right]_{p=Q}$$

$$G(\alpha, \beta, n, \sigma) = (4\pi)^\epsilon \frac{\Gamma(\alpha + \beta - \sigma - 2 + \epsilon)}{\Gamma(\alpha)\Gamma(\beta)} B(2 - \epsilon - \alpha + n - \sigma, 2 - \epsilon - \beta + \sigma)$$

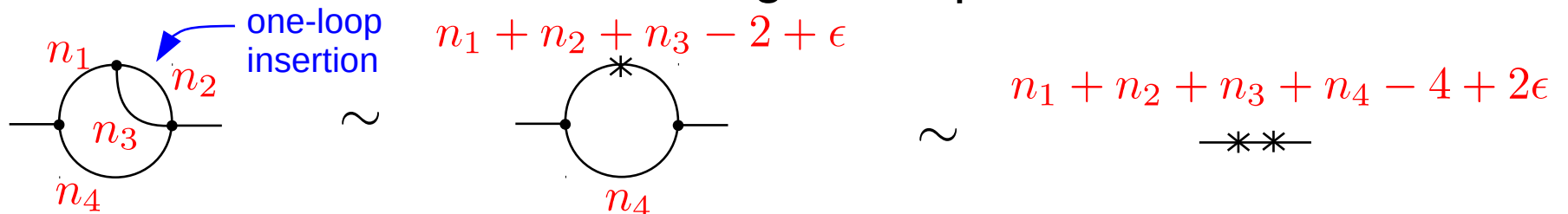
[Chetyrkin, Tkachov '81]

- The result gets a non-integer power $1/(Q^2)^\epsilon$



$*$: non-integer part ϵ

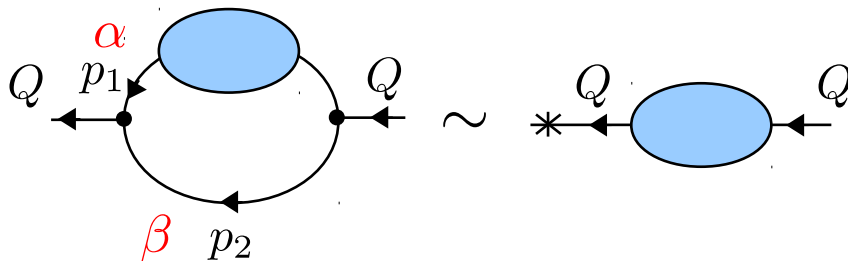
- Can be used as convolutions for higher loops



“Carpet” 1-loop integral

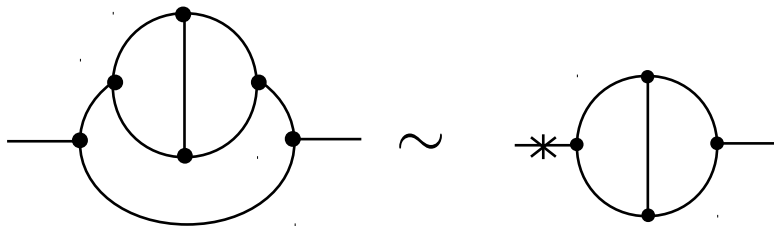
- We can perform another type of 1-loop integral (outer loop of p_1) for **arbitrary indices α and β** possibly with **numerators**

$$p_i \cdot p_j, p_i \cdot Q$$



- Used for the following topology at the 3-loop level

[Chetyrkin, Tkachov '81]



Triangle rule

[Chetyrkin, Tkachov '81]

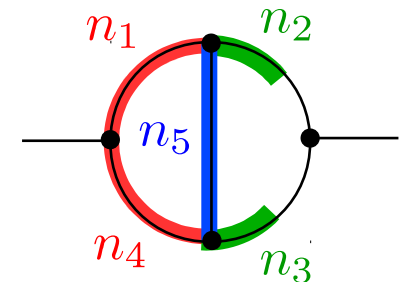
- IBP $(\frac{\partial}{\partial k} \cdot k)$ in one-loop triangle-shaped (sub-)diagrams

any # of lines
numerator $k^{\mu_1} \dots k^{\mu_N}$ OK

$$\left(\frac{\partial}{\partial k} \cdot k \right) \text{Triangle}(a_1, a_2, b, c_1, c_2) = \frac{1}{D + N - a_1 - a_2 - 2b} \left[a_1 \left(\text{Triangle}(a_1+1, a_2, b-1, c_1, c_2) - \text{Triangle}(a_1+1, a_2, b, c_1-1, c_2) \right) + a_2 \left(\text{Triangle}(a_1, a_2+1, b-1, c_1, c_2) - \text{Triangle}(a_1, a_2+1, b, c_1, c_2-1) \right) \right]$$

Decreases b or c_1 or c_2 by 1
at the cost of increasing a_1 or a_2
in the right-hand side

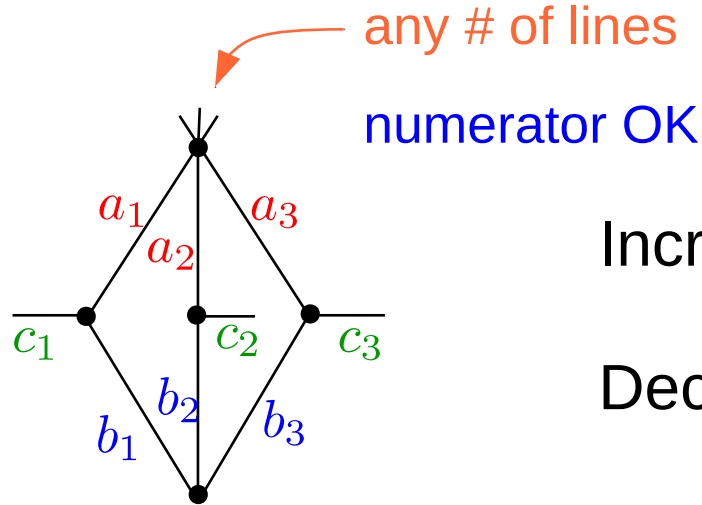
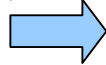
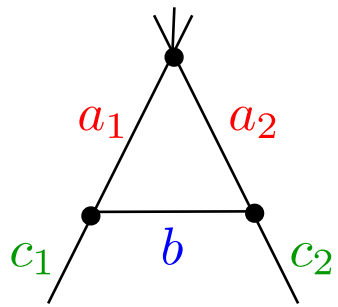
- Recursive use of the triangle rule makes
 $b = 0$ or $c_1 = 0$ or $c_2 = 0$ (removal of a line)
 sums of integrals in simpler topologies



Diamond rule

[Ruijl, TU, Vermaseren '15]

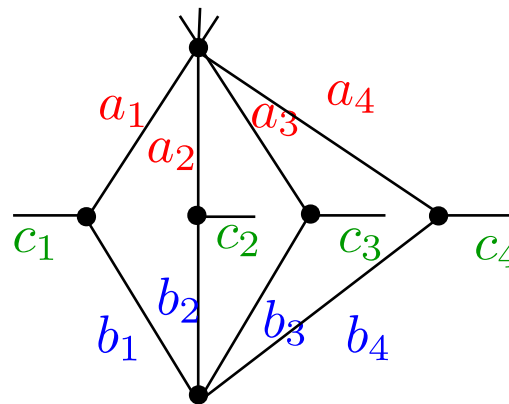
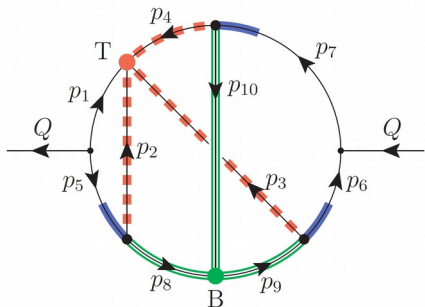
- Extension of the triangle rule to multi-loop diamond-shaped (sub-)diagrams



Increases a_1, a_2, a_3

Decreases b_1, b_2, b_3
 c_1, c_2, c_3

Appear from 4-loops

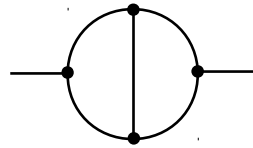


Increases a_1, a_2, a_3, a_4

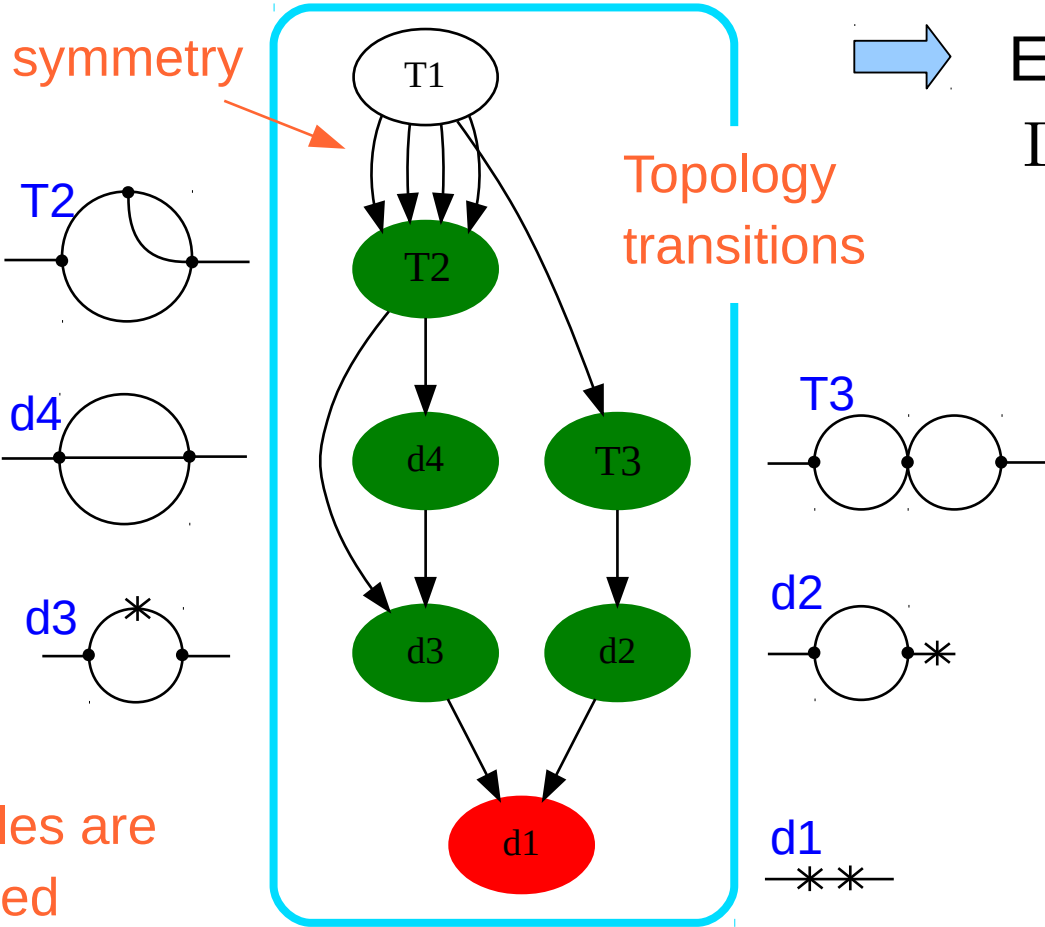
Decreases b_1, b_2, b_3, b_4
 c_1, c_2, c_3, c_4

2-loop topologies

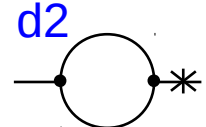
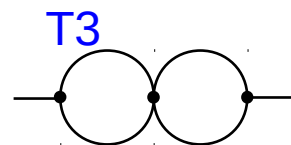
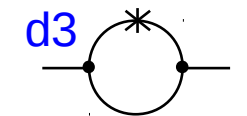
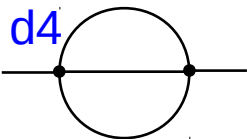
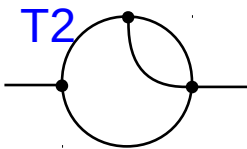
T1 : top-level topology



- All integrals can be evaluated by using the triangle rule and performing one-loop integrals



symmetry



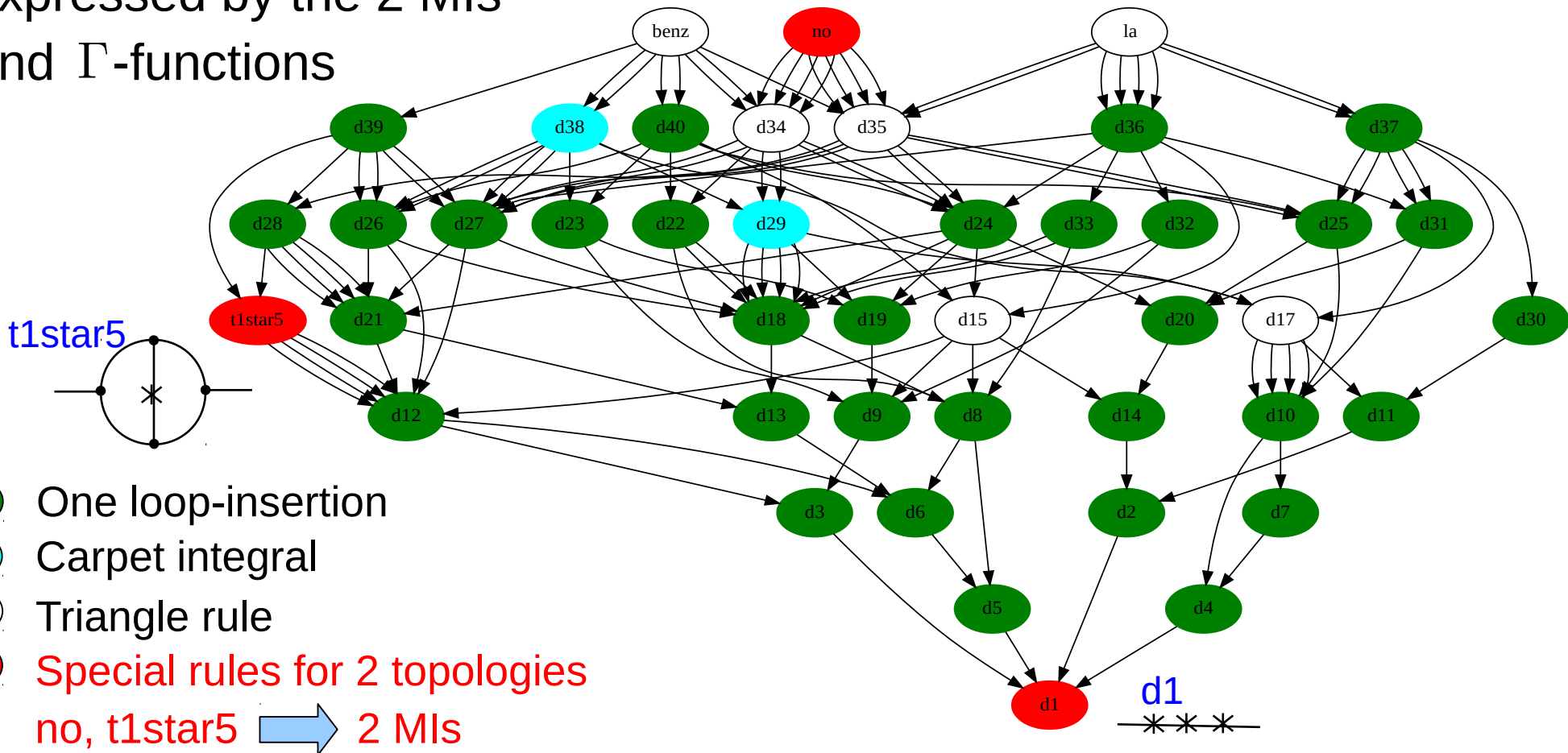
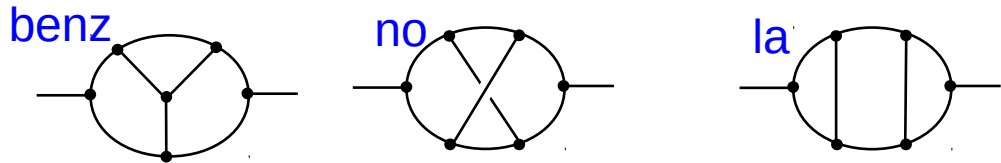
tadpoles are dropped

Topologies in which
 ● One loop-insertion
 ○ Triangle rule can be applied

3-loop topologies

- All integrals are expressed by the 2 MIs and Γ -functions

3 top-level topologies



- One loop-insertion
- Carpet integral
- Triangle rule
- Special rules for 2 topologies
no, t1star5 → 2 MIs

Mincer approach

- Many topologies for massless propagator-type integrals can be reduced to simpler ones by
 - performing one-loop integrals
 - use of triangle rules to remove one of lines
- Special cases (where we need to solve IBPs) are not too many

Algorithm: [Chetyrkin, Tkachov '81]

Schoonschip implementation: [Gorishny, Larin, Surguladze, Tkachov '89]

Form implementation: [Larin, Tkachov, Vermaseren '91]

- In principle, this approach can be extended to the 4-loop level...
 - Laurent series of the MIs in ϵ are known

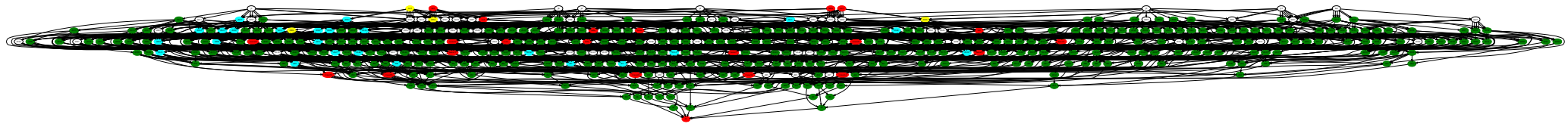
Baikov, Chetyrkin '10 (via glue-and-cut symmetry)

Lee, Smirnov² '11 (up to weight 12, via DRA, Mellin-Barnes, PSLQ)

Extension to the 4-loop level: “Forcer”

4-loop topologies

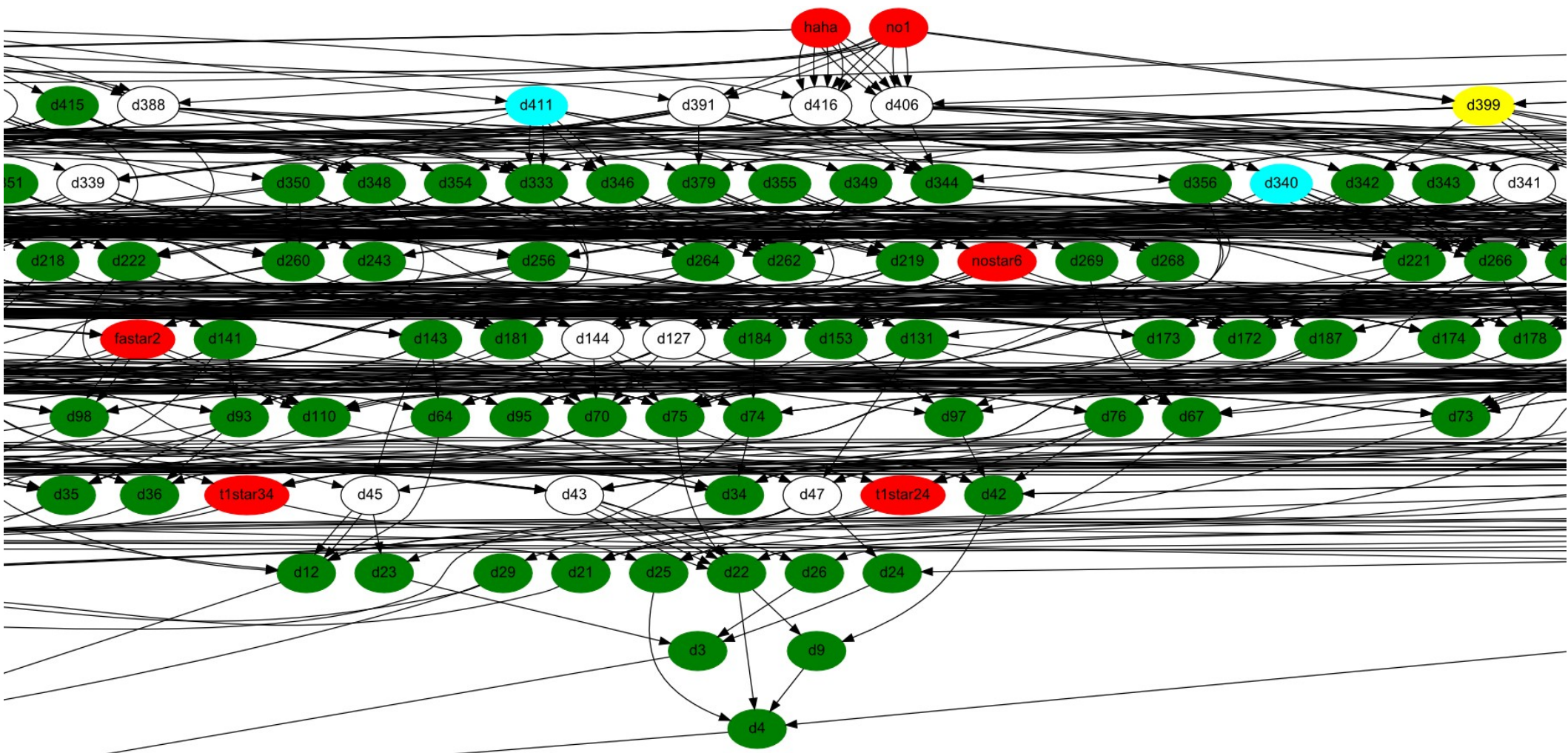
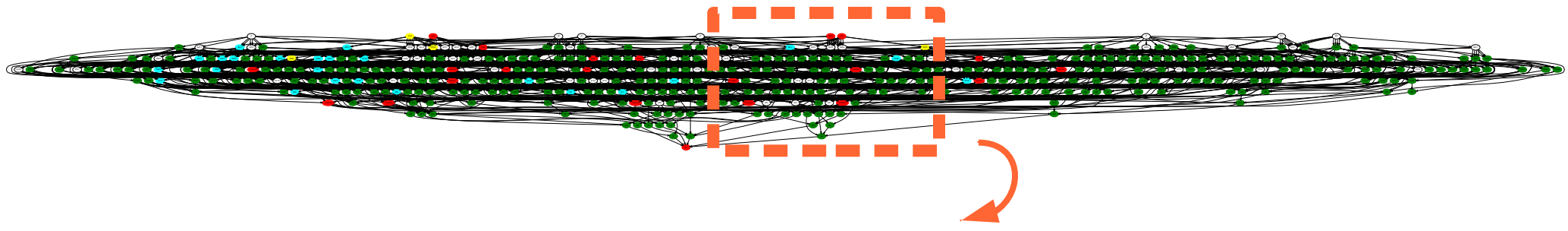
11 top-level topologies



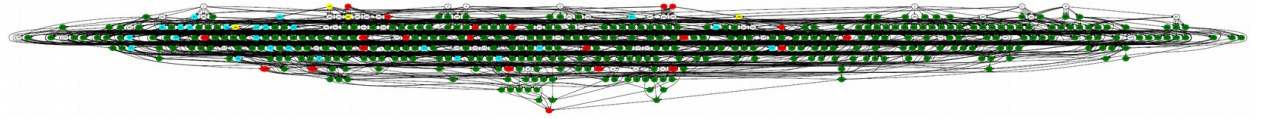
●	One loop-insertion	(335)		
●	Carpet integral	(24)		
○	Triangle rule	(53)		
●	Diamond rule	(4)		
●	Special rules	(21)	(pure 4-loop: 9)	Total: 437

- Enormous number of cases!!
- Coding such a reduction **by hand** is impractical

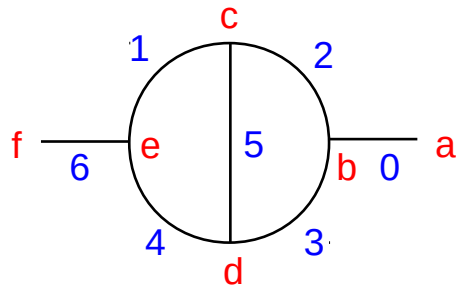
➡ **Automatization**



How to handle

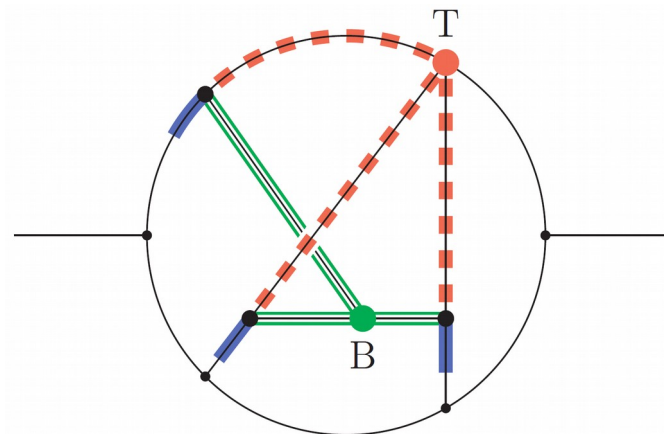


- Each topology as an “undirected graph” in graph theory



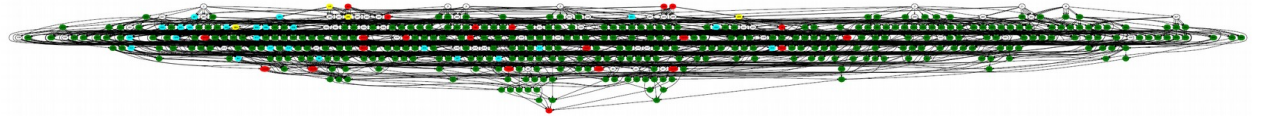
vertices and edges are labeled

- Implementation: Python3 with a graph library “igraph”
<http://igraph.org>
- Easy to detect one-loop insertion, carpet, triangles, diamonds and tadpoles



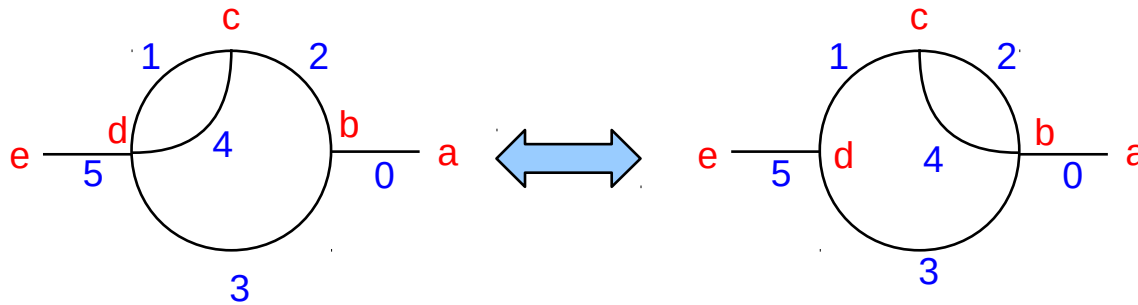
could be difficult
by human eyes
(diamond example)

How to handle



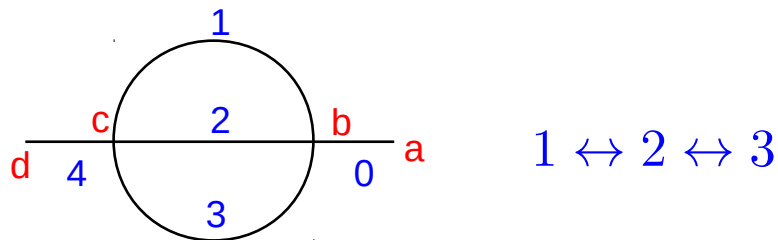
- Graph isomorphism

- Detects equivalent graphs and finds mappings among them

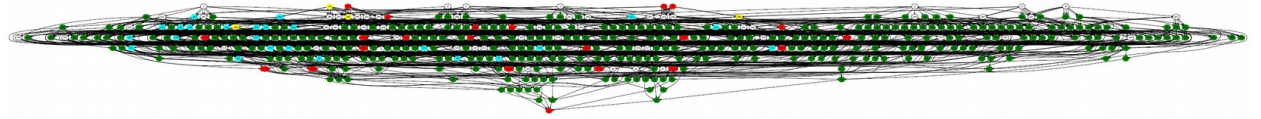


- Graph automorphism

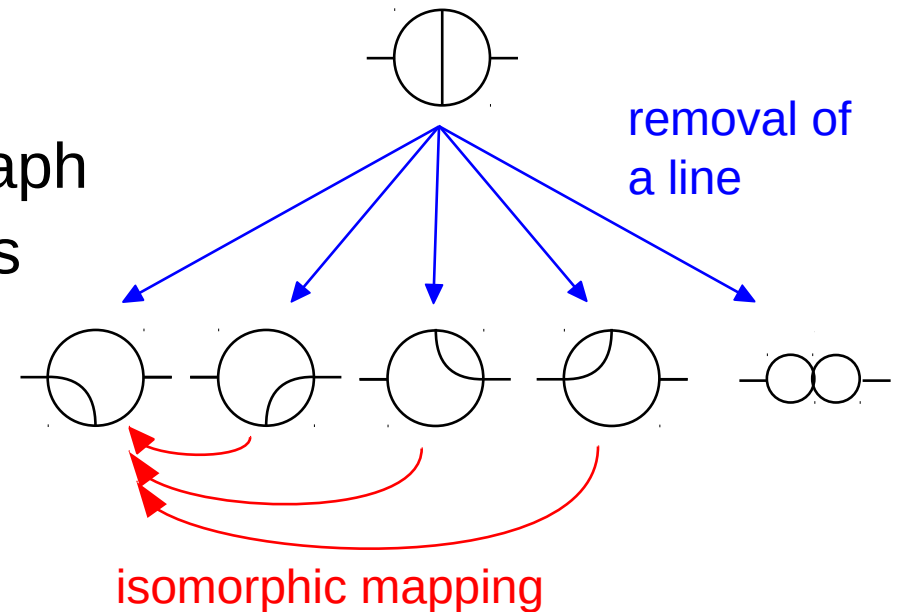
- Finds symmetry / mappings in each graph



How to handle



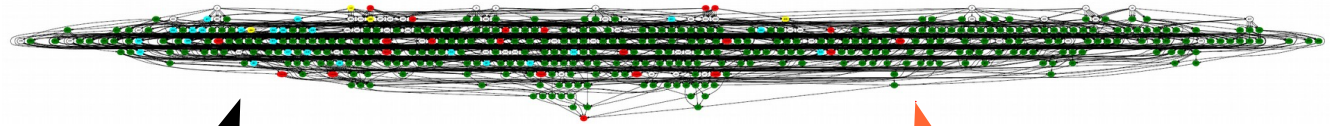
- Input : top-level topologies (11 for 4-loops)
- From each topology, remove a line in all possible ways while taking graph isomorphism and dropping tadpoles into account



- For each topology, the next action is decided (one-loop, carpet, triangle, diamond, otherwise special rule needed)
 - Irreducible numerators (dot products) are chosen such as they do not interfere with the next action

Code generation

- In the end, we get



- Code generation from

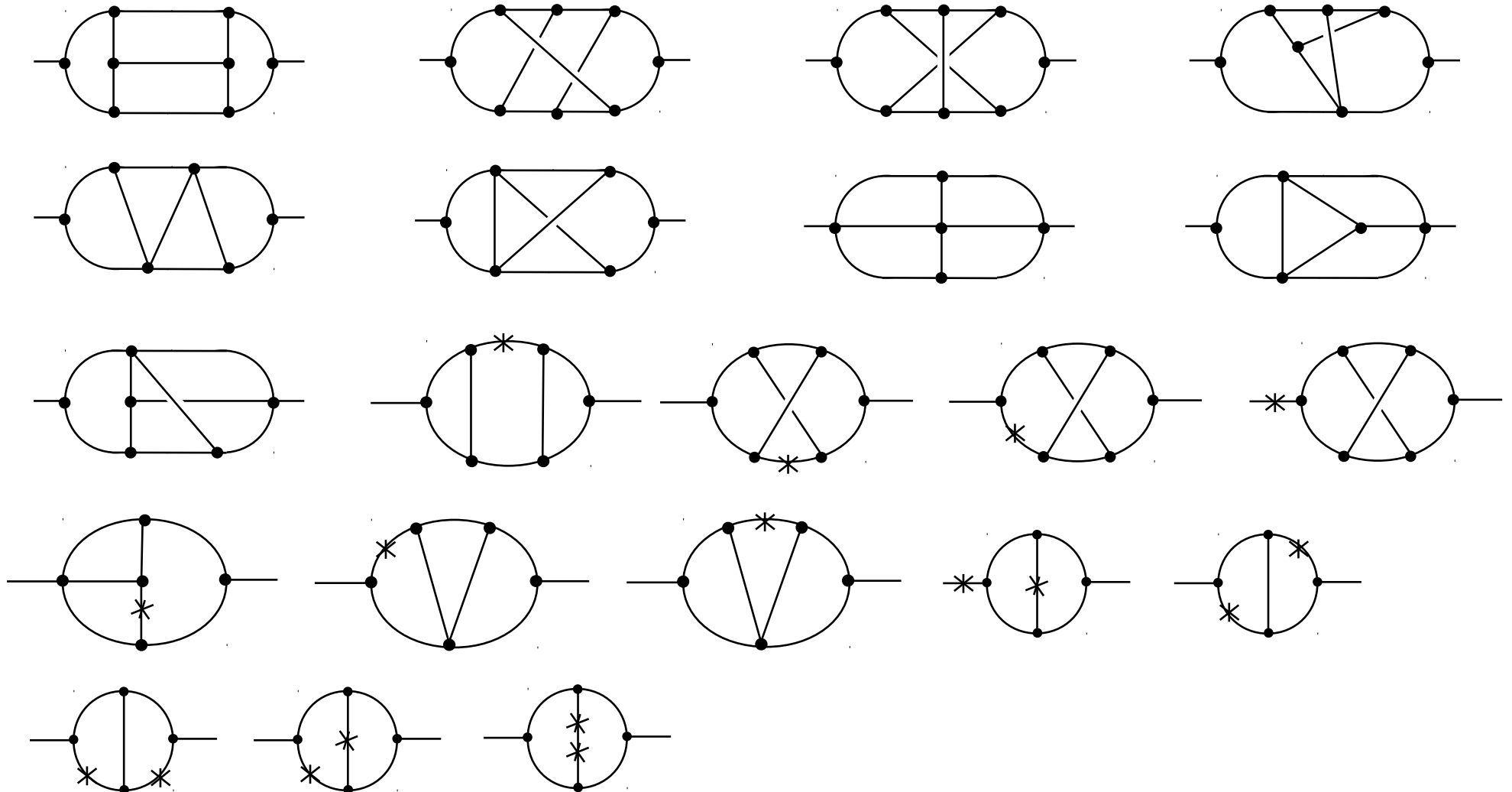


visualized via
Graphviz

- Adequate subroutines are called at each topology
 - Symmetries from graph automorphism
 - Rewriting propagators and irreducible numerators at all transitions from a topology to another
 - Generates 41240 lines of FORM code for 4-loops
- Works even for 5-loops
(64 top-topologies, 6570 in total, but 284 special rules)

Manual reduction rules

- 21 topologies require special rules, manually constructed

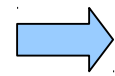


Finding manual reduction rules

- Shift an index by 1 in IBPs in all possible ways

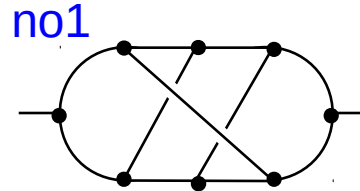
$$n_1 \rightarrow n_1 + 1, n_2 \rightarrow n_2 + 1, n_3 \rightarrow n_3 + 1, \dots$$

- At the 4-loop, 20 IBPs and 14 indices

 $20 \times 14 = 280$ equations

- Eliminate “complicated” integrals that increase indices from the system of equations as possible $F(n_1 + 1, \dots)$
- This is not complete, but helps a lot. Human eyes for finding good rules for reducing integrals
- Everything should be done by using computer algebra systems (even for just rational arithmetic or simple substitutions)
A rule can easily be over 1000 lines

Result looks like...



one of the top-level topology
11 propagators
3 irreducible numerators

[no1(2,2,2,2,2,2,2,2,2,2,2,-1,-1,-1)] =

$$-10/9*\text{num}(1+2*ep)^2*\text{num}(2+5*ep)*\text{num}(3+2*ep)^2*\text{num}(3+5*ep)*\text{num}(4+5*ep)*\text{num}(6+5*ep)*\text{num}(7+5*ep)*\text{num}(8+5*ep)*\text{num}(9+5*ep)*\text{num}(36141384+167650024*ep+369157793*ep^2+504389598*ep^3+470560515*ep^4+312347786*ep^5+149770838*ep^6+51734214*ep^7+12600912*ep^8+2060632*ep^9+203648*ep^{10}+9216*ep^{11})*\text{den}(1+ep)^2*\text{den}(2+ep)^2*\text{den}(2+3*ep)*\text{den}(3+ep)^2*\text{den}(4+3*ep)*\text{den}(5+3*ep)*\text{den}(7+3*ep)*\text{den}(8+3*ep)*\text{Master}(\text{no1})$$

$$+2/9*\text{num}(1+2*ep)*\text{num}(1+4*ep)^2*\text{num}(3+2*ep)*\text{num}(39337381531008+422506983834144*ep+2126828256064272*ep^2+6682923999248124*ep^3+14715479582570820*ep^4+24142592323497543*ep^5+30606097414180459*ep^6+30663886432898386*ep^7+24613240685251374*ep^8+15944111435166437*ep^9+8353498138352567*ep^{10}+3530865125153640*ep^{11}+1195052526807120*ep^{12}+319518770241334*ep^{13}+66015276933500*ep^{14}+10173244132808*ep^{15}+1101566031376*ep^{16}+74820000000*ep^{17}+2400000000*ep^{18})*\text{den}(1+ep)^2*\text{den}(2+ep)^2*\text{den}(2+3*ep)*\text{den}(3+ep)^2*\text{den}(4+3*ep)*\text{den}(5+3*ep)*\text{den}(7+3*ep)*\text{den}(8+3*ep)*\text{Master}(\text{no6})$$

$$-4/9*\text{num}(ep)^2*\text{num}(1+2*ep)*\text{num}(3+2*ep)*\text{num}(4970268890523930816+65981799142896214872*ep+416366999743872130470*ep^2+1663841711414114631648*ep^3+4730696245772216998455*ep^4+10190116418687033776180*ep^5+17283151225628577356718*ep^6+23674869565869007631757*ep^7+26649665048245779930527*ep^8+24945440138420091798098*ep^9+19571616320498135318722*ep^{10}+12933081950816489701881*ep^{11}+7214358008133792648788*ep^{12}+3396810005568803286172*ep^{13}+1346680080420400500352*ep^{14}+447318716827968227680*ep^{15}+123500240869574621248*ep^{16}+28008584917867939712*ep^{17}+5129371482425778688*ep^{18}+739826312414941184*ep^{19}+80910145880457216*ep^{20}+6306386119458816*ep^{21}+312125440000000*ep^{22}+7372800000000*ep^{23})*\text{den}(1+ep)^2*\text{den}(2+ep)^2*\text{den}(3+ep)^2*\text{den}(3+4*ep)*\text{den}(4+3*ep)*\text{den}(5+3*ep)*\text{den}(5+4*ep)*\text{den}(7+3*ep)*\text{den}(7+4*ep)*\text{den}(8+3*ep)*\text{den}(9+4*ep)*\text{den}(11+4*ep)*\text{den}(13+4*ep)*\text{Master}(\text{lala})$$

(cont'd on next page)

+ (other 17 terms)

$$\begin{aligned}
& +1/69984*\text{num}(-1+2*\epsilon)^3*\text{num}(7241916201944976216509644800000+319933970273126101430280830976000*\epsilon+6756458 \\
& 704694283224428114990694400*\epsilon^2+91350459184391655670944774398730240*\epsilon^3+891289886775817303600596179144 \\
& 718336*\epsilon^4+6693037596934954757286105298423578624*\epsilon^5+40235415927130336283161325668026920448*\epsilon^6+19871 \\
& 3617144519275979305008924953882368*\epsilon^7+820916970447314603727085836746702354688*\epsilon^8+2873881635171369729 \\
& 280438878868838131008*\epsilon^9+8607952779410951522413494094313728793280*\epsilon^{10}+222147773368444730132424667001 \\
& 94857424432*\epsilon^{11}+49635643135503998398733139623853337720392*\epsilon^{12}+96271970320109519554706474095019058884 \\
& 972*\epsilon^{13}+16209200424336736948422894086915245188710*\epsilon^{14}+236011025514494738395477852326524380246208*\epsilon \\
& ^{15}+294116625122127952136996687381588045466897*\epsilon^{16}+306225945611093340988941157877804946565622*\epsilon^{17}+25 \\
& 0395537138862785223188155623245534964532*\epsilon^{18}+127849096823093076621319723934762735418774*\epsilon^{19}-32311698 \\
& 658641642165998037752400422510997*\epsilon^{20}-183932395284770686557606510999638263724434*\epsilon^{21}-284795600836347 \\
& 457186578491353430578870326*\epsilon^{22}-315022919192483812286675434478208160276910*\epsilon^{23}-282305580037611102151 \\
& 358311949984849163164*\epsilon^{24}-212936199022871691907574075383055371432706*\epsilon^{25}-136445993518657866963478040 \\
& 015820540166838*\epsilon^{26}-73490078278879345786210757806097518867018*\epsilon^{27}-3189846583021924021328559833408903 \\
& 5466304*\epsilon^{28}-9651382484951439893472757109220693984414*\epsilon^{29}-421236561806300181608330799113484613884*\epsilon^ \\
& 30+2017002468644720352265271521884033262822*\epsilon^{31}+1832874761562866686698305090568644702899*\epsilon^{32}+1096707 \\
& 581456919232171368580075436795092*\epsilon^{33}+524523702662773366611140017451441153706*\epsilon^{34}+212149752741671710 \\
& 649385032559095034796*\epsilon^{35}+74250901029010212070735988610887333749*\epsilon^{36}+2272051252243184038118975458520 \\
& 5484412*\epsilon^{37}+6103292399690879903999959233123243684*\epsilon^{38}+1439736393111277351174910769581576800*\epsilon^{39}+29 \\
& 7528530322975798096839220046747248*\epsilon^{40}+53606182091916632708972765495066048*\epsilon^{41}+835965908902794893568 \\
& 4267151014592*\epsilon^{42}+1117001205642226519475032390006784*\epsilon^{43}+126129308732809465322496308824064*\epsilon^{44}+118 \\
& 10033418325211214154334027776*\epsilon^{45}+892823848935298884179601309696*\epsilon^{46}+52384336126816611070473928704*\epsilon \\
& ^{47}+2238933537865339140484104192*\epsilon^{48}+62029021136981852160000000*\epsilon^{49}+836264176857907200000000*\epsilon^{50}) \\
& *\text{den}(\epsilon)^6*\text{den}(-1+\epsilon)*\text{den}(1+\epsilon)^6*\text{den}(1+2*\epsilon)^2*\text{den}(1+3*\epsilon)^2*\text{den}(-2+\epsilon)*\text{den}(2+\epsilon)^6*\text{den}(2+3*\epsilon)^2*\text{den}(3 \\
& +\epsilon)^3*\text{den}(3+2*\epsilon)^2*\text{den}(3+4*\epsilon)*\text{den}(4+3*\epsilon)^2*\text{den}(5+2*\epsilon)^2*\text{den}(5+3*\epsilon)^2*\text{den}(5+4*\epsilon)*\text{den}(7+3*\epsilon)^2*\text{den} \\
& (7+4*\epsilon)*\text{den}(8+3*\epsilon)*\text{den}(9+4*\epsilon)*\text{den}(11+4*\epsilon)*\text{den}(13+4*\epsilon)*G10*G20*G30
\end{aligned}$$

Optionally: ϵ -expansions in intermediate steps

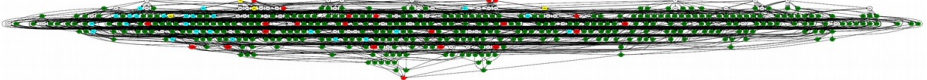
Checks and status

Collaboration with Andreas Vogt

- Recomputing known results – strong non-trivial checks
 - Reproduced the 4-loop QCD β -function via massless propagators [Ritbergen, Vermaseren, Larin '97; Czakon '04]
 - Checked the gauge invariance
 - Also with using background field method
 - Reproduced N=2 and N=4 non-singlet splitting functions [Baikov, Chetyrkin '06] [Velizhanin '14] [Velizhanin '11]
- Statistics of these computations tell us bottlenecks
- Trying to optimize the program before going for N=6 (unknown)

Summary

Summary

- Evaluating Massless propagator-type Feynman integrals is one of key components of pQCD
- “Mincer” up to 3-loops
- We have been developing “Forcer” for 4-loops
 - Highly complicated structure of the program / equations
 - ➡ Automatization: write a program for generating a program for 
 - Manual rules are derived with the aid of computers
- Status:
 - Checked for known results
 - More optimizations

Backup

Integration-by-parts identities (IBPs)

[Chetyrkin, Tkachov '81]

- The Gauss theorem in D -dimension (the surface term vanishes)

$$\int d^D p \frac{\partial}{\partial p^\mu} X^\mu = 0$$

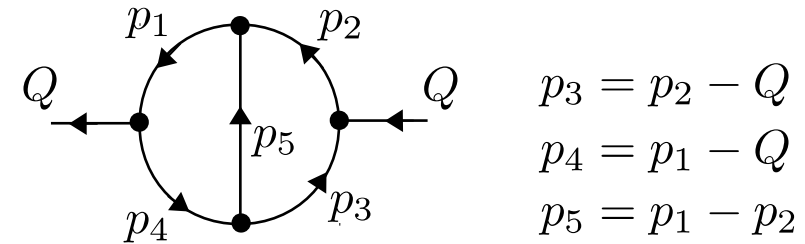
- Example:

$$F(n_1, n_2, n_3, n_4, n_5) = \int \frac{d^D p_1}{(2\pi)^D} \frac{d^D p_2}{(2\pi)^D} \frac{1}{(p_1^2)^{n_1} (p_2^2)^{n_2} (p_3^2)^{n_3} (p_4^2)^{n_4} (p_5^2)^{n_5}}$$

$$\int \frac{d^D p_1}{(2\pi)^D} \frac{d^D p_2}{(2\pi)^D} \frac{\partial}{\partial p_1^\mu} \frac{p_2^\mu}{(p_1^2)^{n_1} (p_2^2)^{n_2} (p_3^2)^{n_3} (p_4^2)^{n_4} (p_5^2)^{n_5}} = 0$$

- Performing $\frac{\partial}{\partial p_1^\mu}$ and multiplying p_2^μ give various dot products

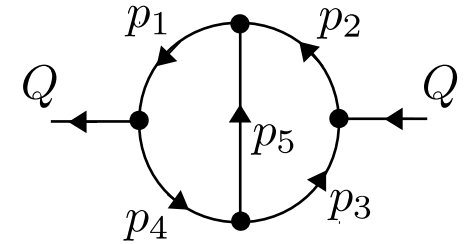
(e.g., $p_1 \cdot p_2$), which are all decomposed to sums of the propagators \rightarrow linear identity among F



Integration-by-parts identities (IBPs)

- $\{p_1, p_2; Q\} \Rightarrow 2 \times 3 = 6$ identities (and their linear combination)

$$\begin{array}{ccc}
 \frac{\partial}{\partial p_1} \cdot p_1 & \frac{\partial}{\partial p_1} \cdot p_2 & \frac{\partial}{\partial p_1} \cdot Q \\
 \frac{\partial}{\partial p_2} \cdot p_1 & \frac{\partial}{\partial p_2} \cdot p_2 & \frac{\partial}{\partial p_2} \cdot Q
 \end{array}$$



$$\begin{aligned}
 & F(n_1 - 1, n_2, n_3, n_4 + 1, n_5) \times (-n_4) \\
 & + F(n_1 - 1, n_2, n_3, n_4, n_5 + 1) \times (-n_5) \\
 & + F(n_1, n_2 - 1, n_3, n_4, n_5 + 1) \times (n_5) \\
 & + F(n_1, n_2, n_3, n_4 + 1, n_5) \times (Q^2 n_4) \\
 & + F(n_1, n_2, n_3, n_4, n_5) \times (-2n_1 - n_4 - n_5 + 4 - 2\epsilon) \\
 & = 0
 \end{aligned}$$

$$\frac{\partial}{\partial p_1} \cdot p_1 - \frac{\partial}{\partial p_1} \cdot p_2 = ?$$

$$\begin{aligned}
 & F(n_1 - 1, n_2, n_3, n_4 + 1, n_5) \times (-n_4) \\
 & + F(n_1 - 1, n_2, n_3, n_4, n_5 + 1) \times (-n_5) \\
 & + F(n_1, n_2 - 1, n_3, n_4, n_5 + 1) \times (n_5) \\
 & + F(n_1, n_2, n_3, n_4 + 1, n_5) \times (Q^2 n_4) \\
 & + F(n_1 + 1, n_2 - 1, n_3, n_4, n_5) \times (-n_1) \\
 & + F(n_1 + 1, n_2, n_3, n_4, n_5 - 1) \times (n_1) \\
 & + F(n_1, n_2, n_3 - 1, n_4 + 1, n_5) \times (-n_4) \\
 & + F(n_1, n_2, n_3, n_4 + 1, n_5 - 1) \times (n_4) \\
 & + F(n_1, n_2, n_3, n_4, n_5) \times (-n_1 + n_5) \\
 & = 0
 \end{aligned}$$

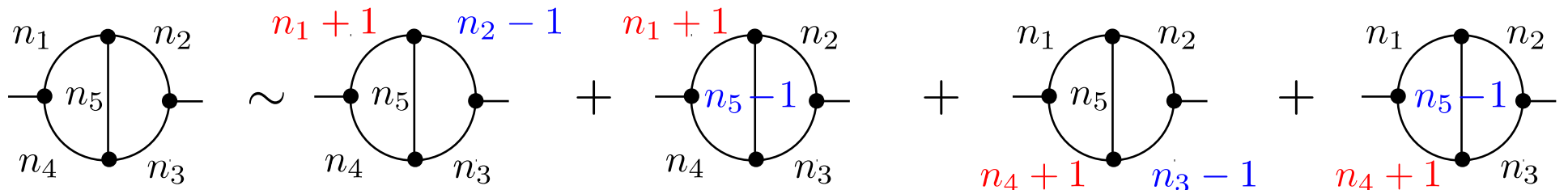
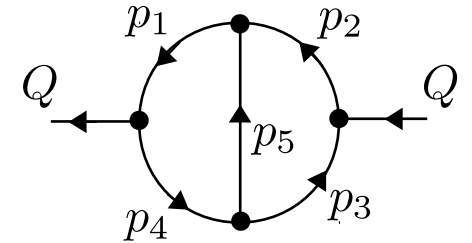
and other 4 identities

Integration-by-parts identities (IBPs)

$$\frac{\partial}{\partial p_1} \cdot p_1 - \frac{\partial}{\partial p_1} \cdot p_2 \quad \Rightarrow$$

$$\begin{aligned}
 & F(n_1, n_2, n_3, n_4, n_5) \times (-n_1 - n_4 - 2n_5 + 4 - 2\epsilon) \\
 & + F(n_1 + 1, n_2 - 1, n_3, n_4, n_5) \times (n_1) \\
 & + F(n_1 + 1, n_2, n_3, n_4, n_5 - 1) \times (-n_1) \\
 & + F(n_1, n_2, n_3 - 1, n_4 + 1, n_5) \times (n_4) \\
 & + F(n_1, n_2, n_3, n_4 + 1, n_5 - 1) \times (-n_4) = 0
 \end{aligned}$$

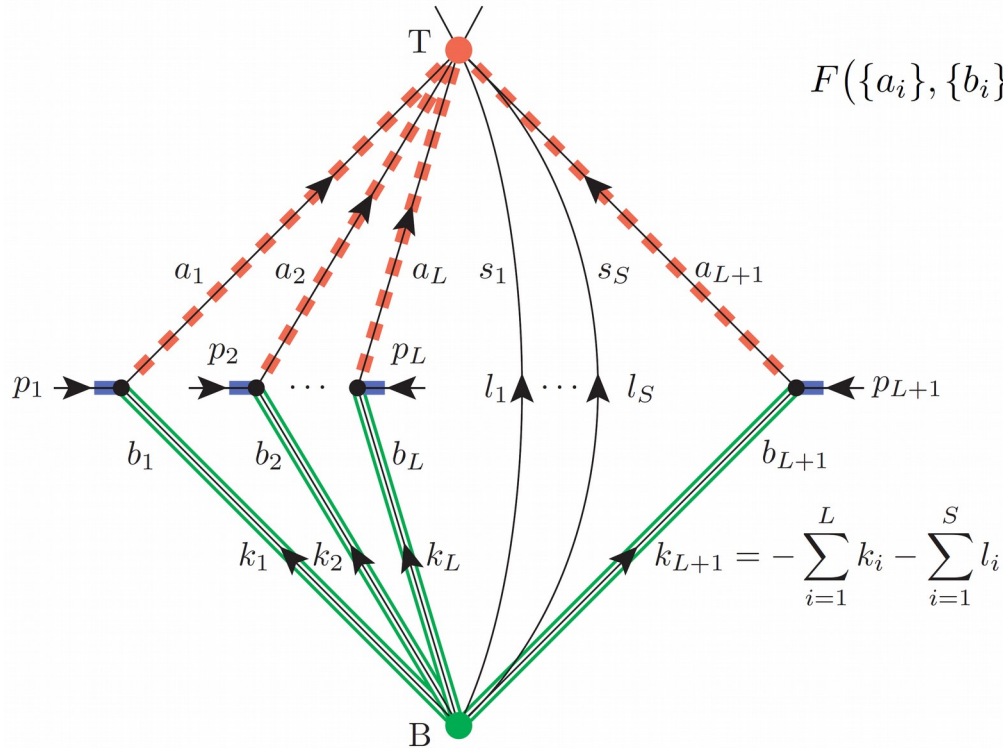
never be 0



- Decreases $(n_2 + n_3 + n_5)$ by 1 at the cost of increasing $(n_1 + n_4)$
- Repeated use of this rule gives integrals with $n_2 = 0$, $n_3 = 0$ or $n_5 = 0 \Rightarrow$ simpler topologies with removing one of the lines

Diamond rule

[Ruijl, TU, Vermaseren '15]



$$F(\{a_i\}, \{b_i\}) = \left[\prod_{i=1}^L \int d^D k_i \right] \left[\prod_{i=1}^S \int d^D l_i \right] \times \left[\prod_{i=1}^{L+1} \frac{k_i^{\mu_1^{(i)}} \dots k_i^{\mu_{N_i}^{(i)}}}{[(k_i + p_i)^2 + m_i^2]^{a_i} (k_i^2)^{b_i}} \right] \left[\prod_{i=1}^S \frac{l_i^{\nu_1^{(i)}} \dots l_i^{\nu_{R_i}^{(i)}}}{(l_i^2)^{s_i}} \right]$$

$$k_{L+1} = - \sum_{i=1}^L k_i - \sum_{i=1}^S l_i$$

$$(L + S)D + \sum_{i=1}^{L+1} (N_i - a_i - 2b_i) + \sum_{i=1}^S (R_i - 2s_i) = \sum_{i=1}^{L+1} a_i \mathbf{A}_i^+ [\mathbf{B}_i^- - (p_i^2 + m_i^2)]$$

Diamond rule

- Explicit summation formula

$$F(\{a_i\}, \{b_i\}, \{c_i\}) =$$

$$\begin{aligned} & \sum_{r=1}^{L+1} \left[\left(\prod_{\substack{i=1 \\ i \neq r}}^{L+1} \sum_{k_i^+ = 0}^{b_i - 1} \right) \left(\prod_{i=1}^{L+1} \sum_{k_i^- = 0}^{c_i - 1} \right) (-1)^{k^-} \frac{k_r^+ (k^+ + k^- - 1)!}{\prod_{i=1}^{L+1} k_i^+! k_i^-!} (E + k^+)_{-k^+ - k^-} \right. \\ & \times \left. \left(\prod_{i=1}^{L+1} (a_i)_{k_i^+ + k_i^-} \right) F(\{a_i + k_i^+ + k_i^-, \{b_i - k_i^+\}, \{c_i - k_i^-\}) \right]_{k_r^+ = b_r} \\ & + \sum_{r=1}^{L+1} \left[\left(\prod_{i=1}^{L+1} \sum_{k_i^+ = 0}^{b_i - 1} \right) \left(\prod_{\substack{i=1 \\ i \neq r}}^{L+1} \sum_{k_i^- = 0}^{c_i - 1} \right) (-1)^{k^-} \frac{k_r^- (k^+ + k^- - 1)!}{\prod_{i=1}^{L+1} k_i^+! k_i^-!} (E + k^+ + 1)_{-k^+ - k^-} \right. \\ & \times \left. \left(\prod_{i=1}^{L+1} (a_i)_{k_i^+ + k_i^-} \right) F(\{a_i + k_i^+ + k_i^-, \{b_i - k_i^+\}, \{c_i - k_i^-\}) \right]_{k_r^- = c_r} \end{aligned}$$

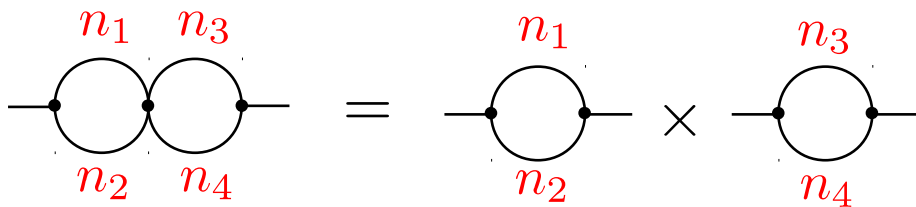
$$E = (L + S)D + \sum_{i=1}^{L+1} (N_i - a_i - 2b_i) + \sum_{i=1}^S (R_i - 2s_i)$$

$$k^+ = \sum_{i=1}^{L+1} k_i^+ \quad k^- = \sum_{i=1}^{L+1} k_i^-$$

Avoids spurious poles

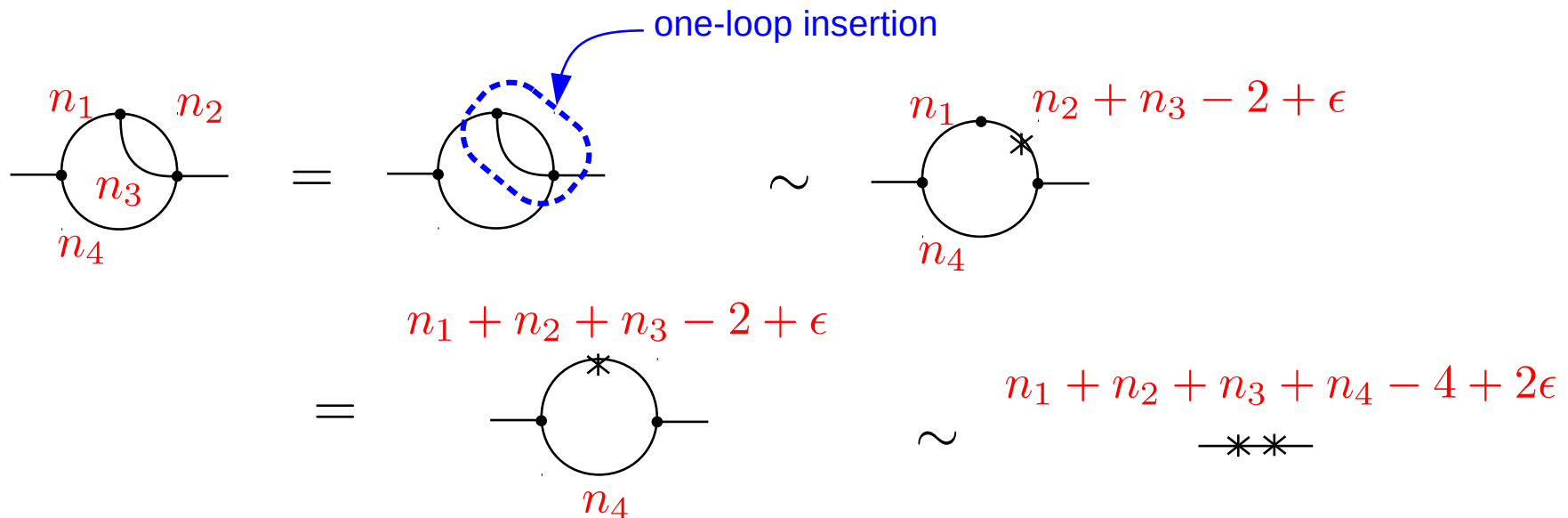
2-loop: easy integrals

- Products / convolutions of 1-loop integrals



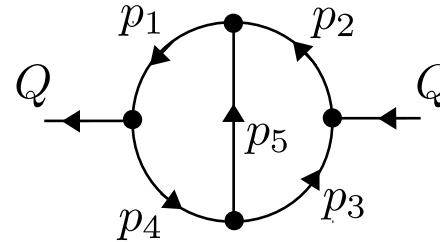
Lines represent denominators
Numerators are no problem

$$\sim \frac{n_1 + n_2 - 2 + \epsilon}{\text{---}^* \text{---}} \times \frac{n_3 + n_4 - 2 + \epsilon}{\text{---}^* \text{---}} = \frac{n_1 + n_2 + n_3 + n_4 - 4 + 2\epsilon}{\text{---}^{**} \text{---}}$$



2-loop: generic topology

- “Generic” topology with the maximal number of lines at the 2-loop:

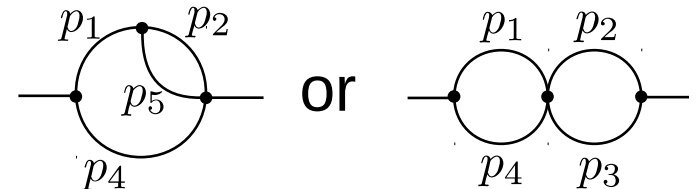


$$\begin{aligned} p_3 &= p_2 - Q \\ p_4 &= p_1 - Q \\ p_5 &= p_1 - p_2 \end{aligned}$$

$$F(n_1, n_2, n_3, n_4, n_5) = \int \frac{d^D p_1}{(2\pi)^D} \frac{d^D p_2}{(2\pi)^D} \frac{1}{(p_1^2)^{n_1} (p_2^2)^{n_2} (p_3^2)^{n_3} (p_4^2)^{n_4} (p_5^2)^{n_5}}$$

- The propagators are “complete”: e.g., $2p_1 \cdot p_2 = p_1^2 + p_2^2 - p_5^2$

- Easy when one of the indices is ≤ 0



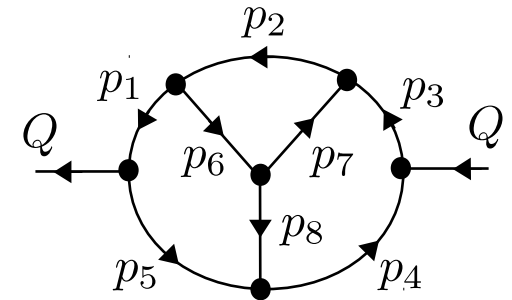
- In the case that all the indices are ≥ 1 :

- In early days: special techniques (e.g., GPXT)
- Nowadays: IBPs

3-loop level

- Generic topologies have 8 lines, but there are 9 dot products
Need to introduce another index for an irreducible numerator

$$F(n_1, \dots, n_9) = \int \frac{d^D p_1}{(2\pi)^D} \frac{d^D p_2}{(2\pi)^D} \frac{d^D p_3}{(2\pi)^D} \times \frac{(p_2 \cdot Q)^{-n_9}}{(p_1^2)^{n_1} \dots (p_8^2)^{n_8}}$$

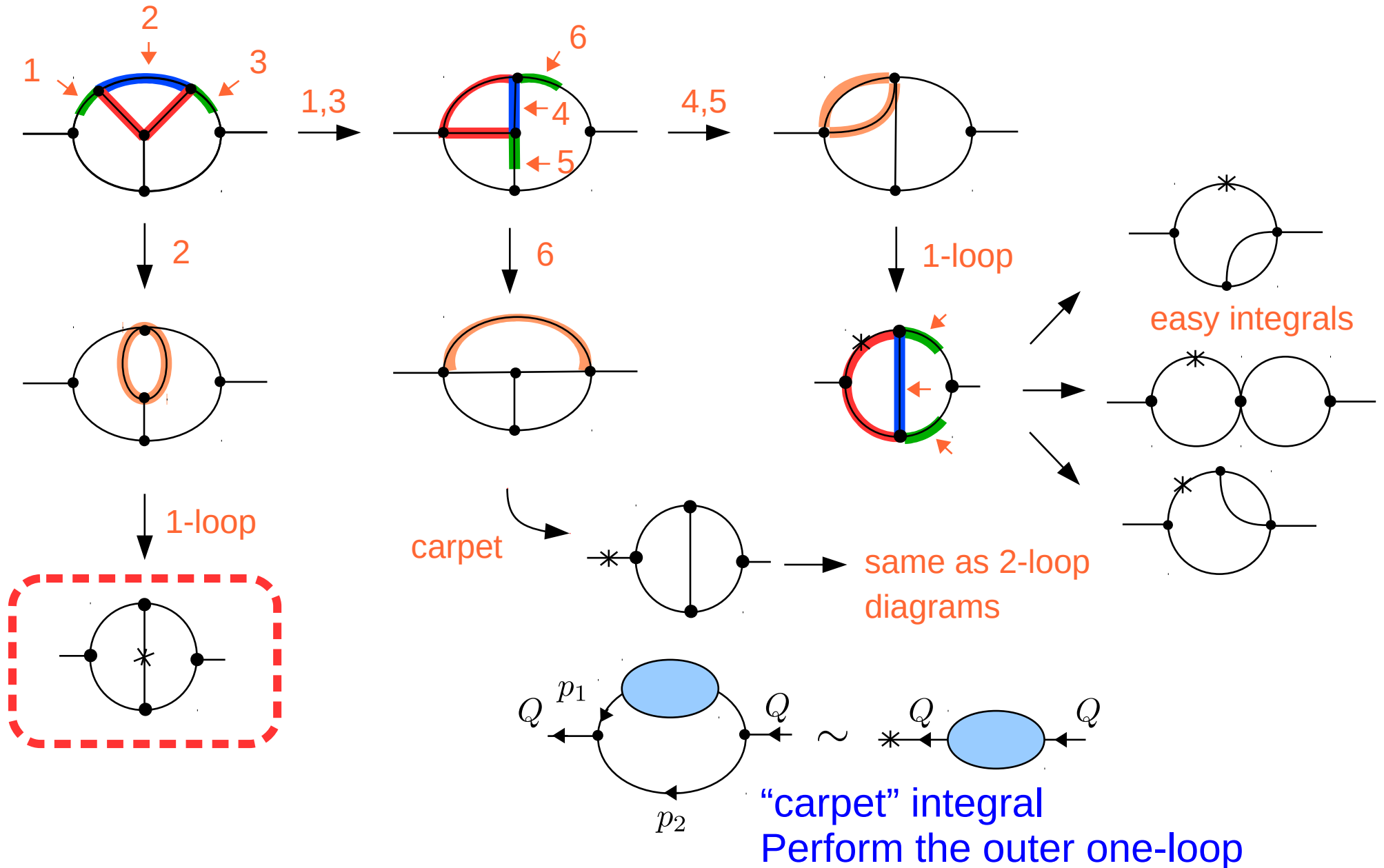


“Benz” topology



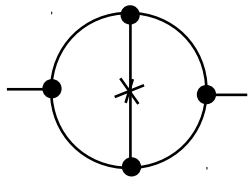
- Our convention
positive indices : denominators
negative indices: numerators
- The irreducible numerator $(p_2 \cdot Q)$ is chosen such that it does not interfere with use of the triangle rule (next page)

Reduction of the Benz topology



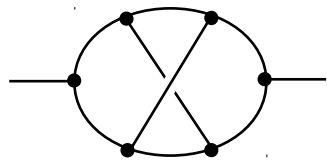
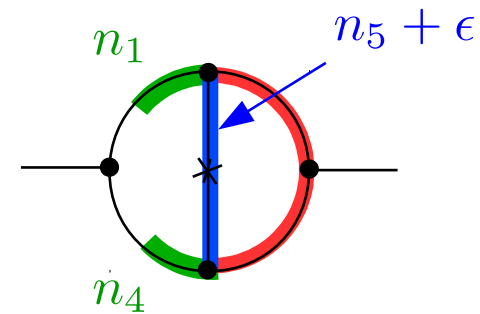
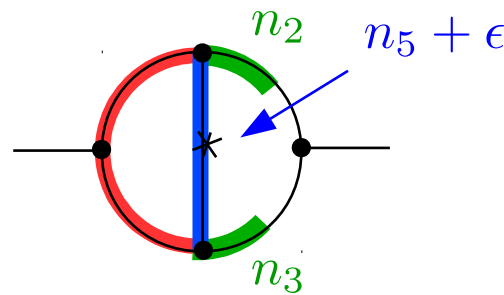
Special topologies at 3-loop

- At the 3-loop level, there are 2 topologies that require special treatment



Repeated use of the triangle rule cannot remove the central line because of the non-integer $n_5 + \epsilon$

2-loop topology with ϵ at the central line



No triangle

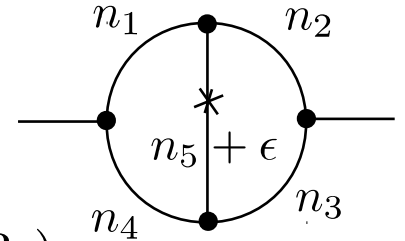
generic “non-planer” topology

- Need to “solve” IBPs: solve a system of recurrence relations to reduce these integrals as much as possible

Special topologies at 3-loop

- Fortunately, this is possible:

$$\begin{aligned}
 & F(n_1, n_2, n_3, n_4, n_5 + \epsilon) \times \underbrace{(n_1 - 1)}_{\text{valid for } n_1 \neq 1} \\
 & + F(n_1 - 1, n_2, n_3, n_4, n_5 + \epsilon) \times (-1) \times (n_1 + 2n_4 + n_5 - 5 + 3\epsilon) \\
 & + F(n_1, n_2, n_3, n_4 - 1, n_5 + \epsilon) \times (-1) \times (n_1 - 1) \\
 & + F(n_1 - 1, n_2, n_3 - 1, n_4, n_5 + 1 + \epsilon) \times (n_5 + \epsilon) \\
 & + F(n_1 - 1, n_2, n_3, n_4 - 1, n_5 + 1 + \epsilon) \times (-1) \times (n_5 + \epsilon) = 0
 \end{aligned}$$

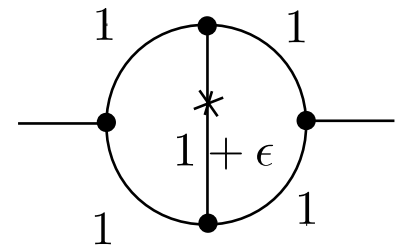


Decreases n_1 or n_4 (or both) until $n_1 = 1$. Similar rules for n_2, n_3, n_4

* They can be derived from the triangle rule identity [Ruij, TU, Vermaseren '15]

Then n_5 can be increased/decreased until $n_5 = 1$ by

$$\begin{aligned}
 & F(1, 1, 1, 1, n_5 + \epsilon) \times (-1) \times (n_5 - 1 + 2\epsilon) \\
 & + F(1, 1, 1, 1, n_5 - 1 + \epsilon) \times (-1) \times (n_5 - 2 + 3\epsilon) \\
 & + F(0, 1, 1, 1, n_5 + \epsilon) \times (-3 + 2n_5 + 5\epsilon) \\
 & + F(1, 1, 1, 0, n_5 + \epsilon) \times (-3 + 2n_5 + 5\epsilon) = 0
 \end{aligned}$$

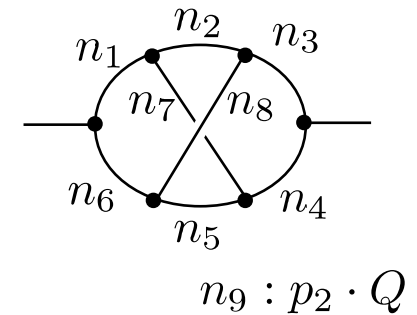


- “Master” integral $F(1, 1, 1, 1, 1 + \epsilon)$ cannot be reduced any more

Special topologies at 3-loop

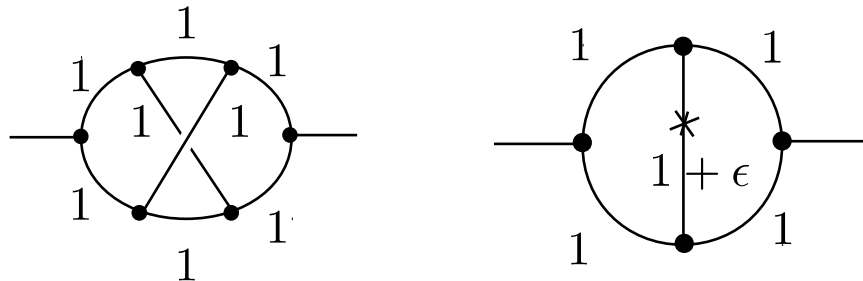
- One can make a reduction scheme also for the non-planer topology

* Rules for reducing n_1, n_3, n_4, n_6 to 1 can be derived from the diamond rule identity [Ruijl, TU, Vermaseren '15]



- Master integral $F(1, 1, 1, 1, 1, 1, 1, 1, 0)$

- The 2 master integrals (MIs) at the 3-loop level



- They cannot be expressed in terms of Γ -functions
But their expansions with respect to ϵ are known
(at least, up to enough orders)