



Some beam-beam considerations on optics design of the low β insertions



X. Buffat

- Performance reach with small β^*
 - Adaptive β^*
 - Crab cavities
- Crossing scheme
- Conclusion



Luminosity and beam-beam tune shift



- Round beams

$$L_0 = \frac{n_b f_{rev} N^2}{4 \pi \sigma^2}$$

$$\xi = \frac{r_0 N}{4 \pi \epsilon}$$



$$L_0 = \frac{n_b f_{rev} \gamma}{r_0} \frac{N \xi}{\beta^*}$$

Cleaning
Machine protection
Power loss
Heating

Emittance growth
Beam loss
Dynamic effects
(Orbit, β_{beat})

Smallest β^* ?



Critical β^*



$$\frac{1}{N} \frac{dN}{dT} = \sigma_{tot} \frac{n_{IP} f_{rev} \gamma}{4 \pi \beta^*} \frac{N}{\epsilon}$$
$$\frac{1}{\epsilon} \frac{d\epsilon}{dT} = -\frac{1}{\tau}$$

$$\frac{N}{\epsilon} = cst \Rightarrow \sigma_{tot} \frac{f_{rev} \gamma}{r_0 \beta^*} \xi_{tot} = \frac{1}{\tau}$$

- The system ultimately reaches a balance between burn-off and synchrotron damping
 - The beam brightness (i.e. the beam-beam parameter) becomes constant
- To achieve a constant beam-beam parameter of 0.01, one requires :

$$\beta_c^* = \sigma_{tot} \tau \frac{f_{rev} \gamma}{r_0} \xi_{tot} \approx 0.06 \text{ m}$$

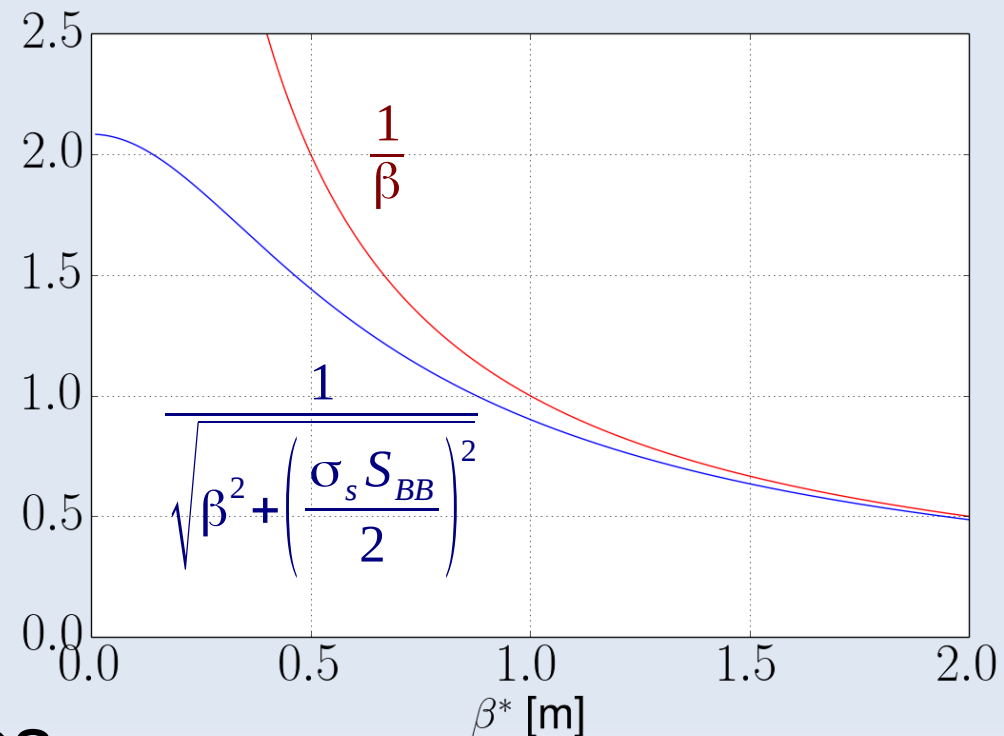
→ One will need to either actively control the beam brightness or live with a non-constant beam-beam parameter

- Crossing angle defined by long-range beam-beam interaction :

$$L = \frac{n_b f_{rev} N^2}{4\pi\sigma^2} \frac{1}{\sqrt{1 + \left(\frac{\sigma_s}{\sigma_x} \frac{\theta_{full}}{2}\right)^2}}$$

$$= \frac{n_b f_{rev} \gamma N^2}{4\pi\epsilon} \frac{1}{\sqrt{\beta^{*2} + \left(\frac{\sigma_s S_{BB}}{2}\right)^2}}$$

$$S_{BB} \approx \sqrt{\beta^*} \frac{\gamma}{\epsilon_n} \theta_{full} = 12 \rightarrow \text{Depends on the crossing scheme (L}^*, \text{ experimental spectrometer, ...)}$$



- $\beta^* < \sim 0.3$ m are not very interesting w/o crab cavities



More elaborate luminosity model



$$\left\{ \begin{array}{l} \frac{\partial I}{\partial t} = -\frac{I(t)}{\tau_{lifetime}} - \mathcal{L}_{IP}(t) N_{IP} \sigma_{tot} \\ \frac{\partial \epsilon_x}{\partial t} = -\frac{\epsilon_x(t)}{\tau_{rad,x}} + \alpha_{rad,x} + \frac{I(t)}{\epsilon_y(t)} \alpha_{IBS,x} \\ \frac{\partial \epsilon_y}{\partial t} = -\frac{\epsilon_y(t)}{\tau_{rad,y}} \\ \frac{\partial \epsilon_s}{\partial t} = 0 \\ \mathcal{L}_{IP} = \frac{n_b f_{rev} N(t)^2 \gamma_r}{4\pi \beta^* \sqrt{\epsilon_x(t) \epsilon_y(t)}} \frac{\cos(\phi(t))^2}{\sqrt{1 + \frac{\sigma_s^2}{\sigma^2} \tan(\phi(t))^2}} \end{array} \right.$$

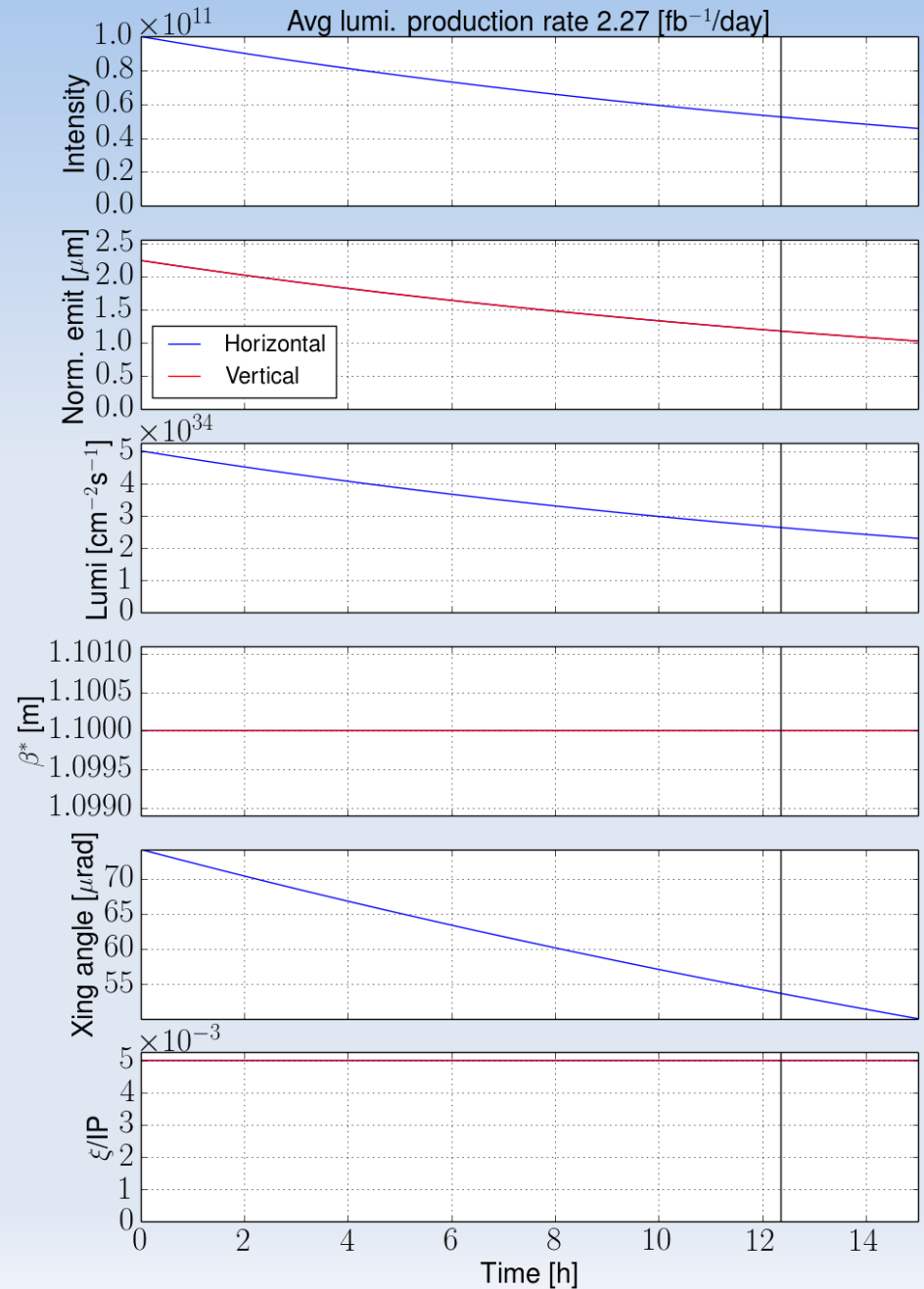
- Still assume round beam (blow-up required in the vertical plane)
- Beam-beam parameter computed assuming two IPs with alternating crossing angle
- The crossing angle is adjusted during the fill to keep the same beam-beam separation at the long-range encounters



Beam parameter evolution



- The nominal configuration is limited by the small beam-beam tune shift

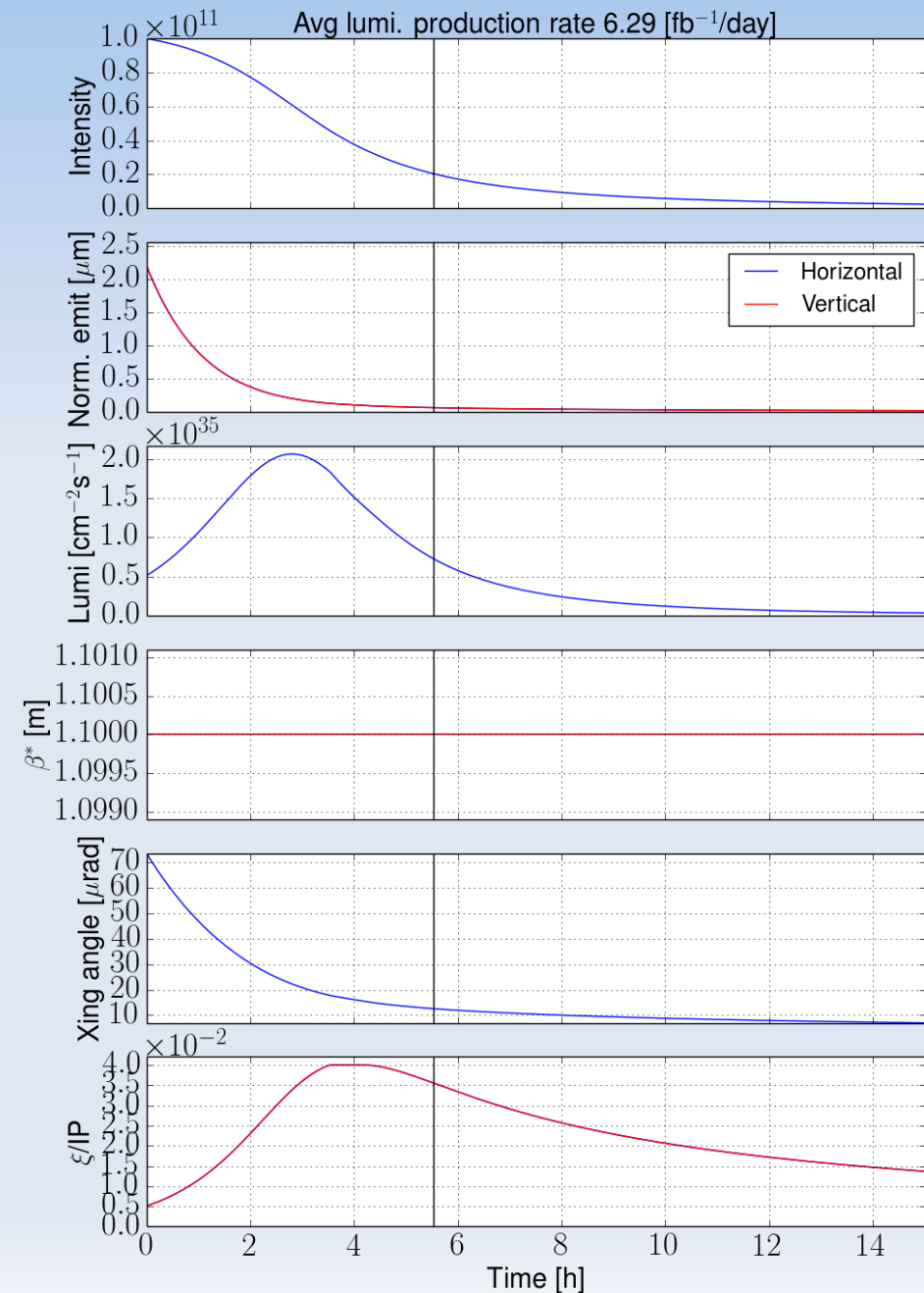




Beam parameter evolution



- The nominal configuration is limited by the small beam-beam tune shift
 - The beam parameter evolution is very different in configurations with larger beam-beam tune shift

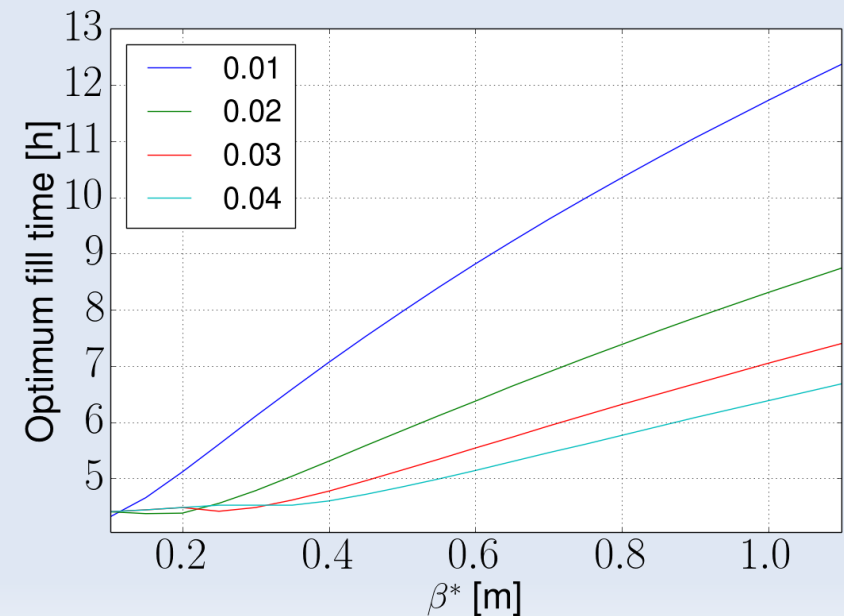
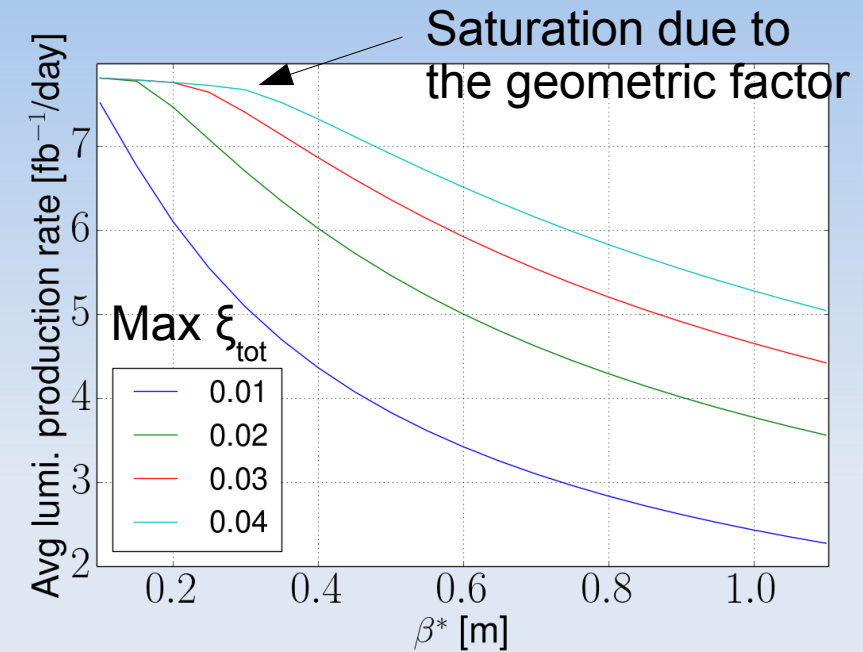


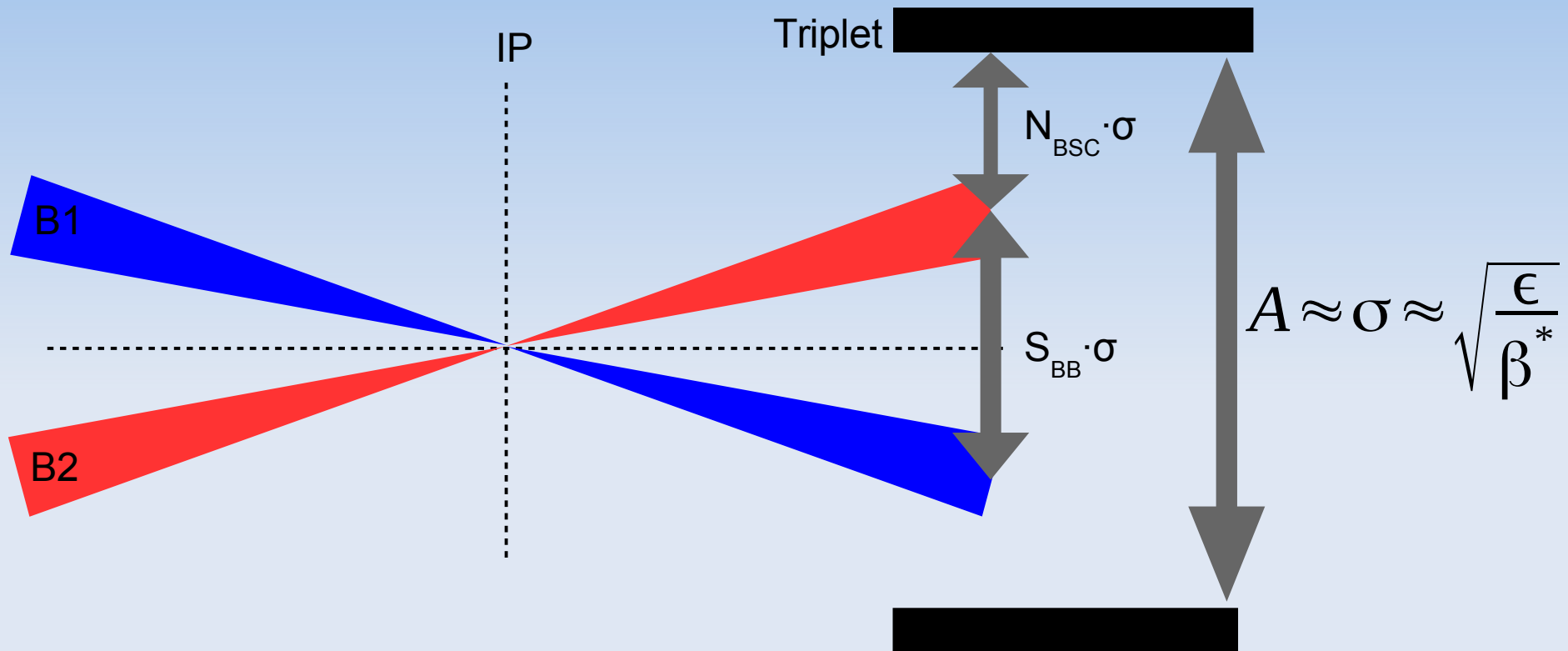


Beam parameter evolution



- The nominal configuration is limited by the small beam-beam tune shift
 - The beam parameter evolution is very different in configurations with larger beam-beam tune shift
- Reduced β^* allows to achieve higher integrated luminosity within shorter fills





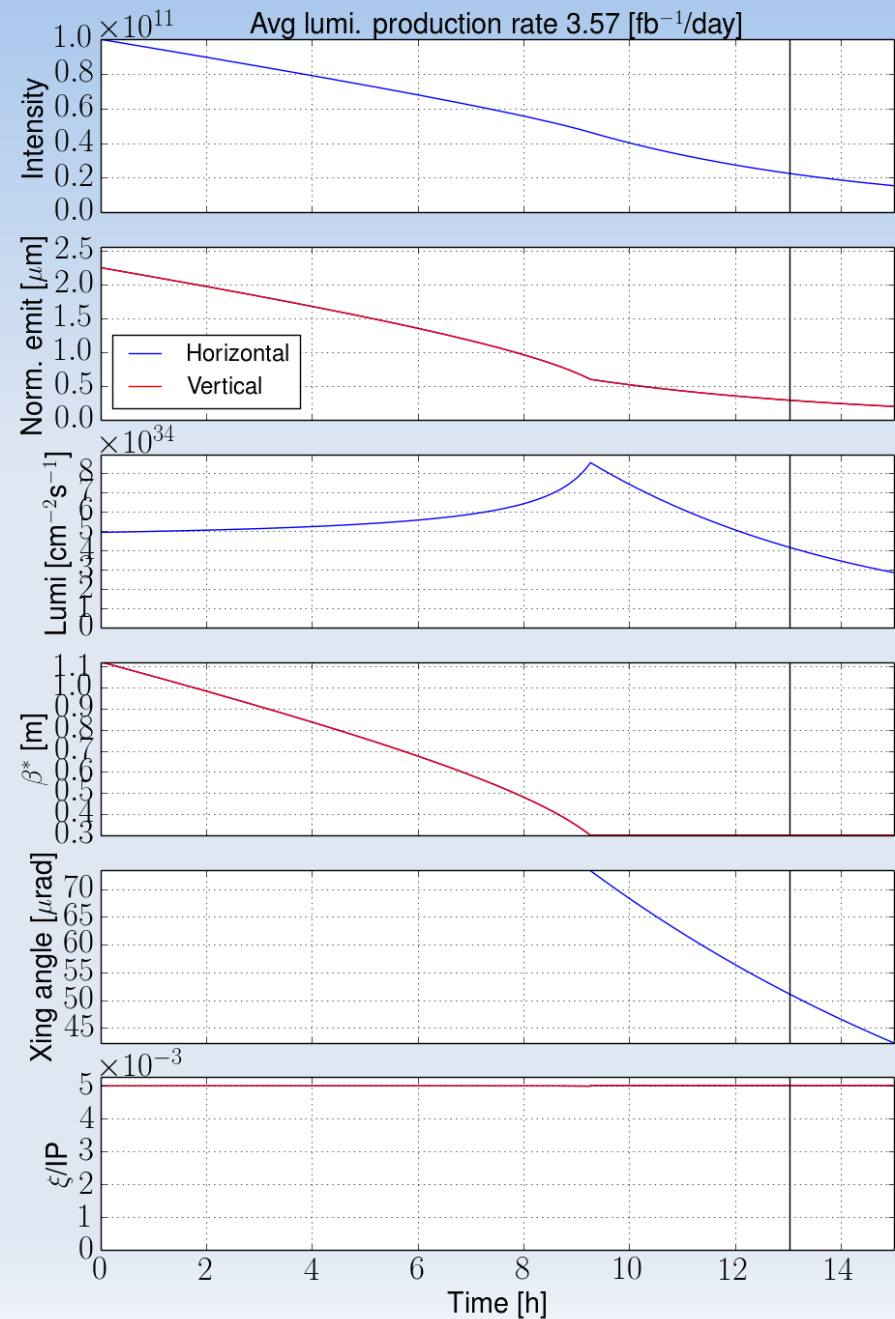
- β^* can be reduced with the reduction of the emittance during the fill without increased aperture requirement by keeping the ratio ϵ/β^* constant
 - The beam-stay-clear and the normalised beam-beam separation are kept constant
 - no change of crossing angle



Beam parameter evolution



- Nominal configuration, with β^* decreasing to 0.3 m during the fill
 - 1.5x performance increase, but at the cost of very long fills

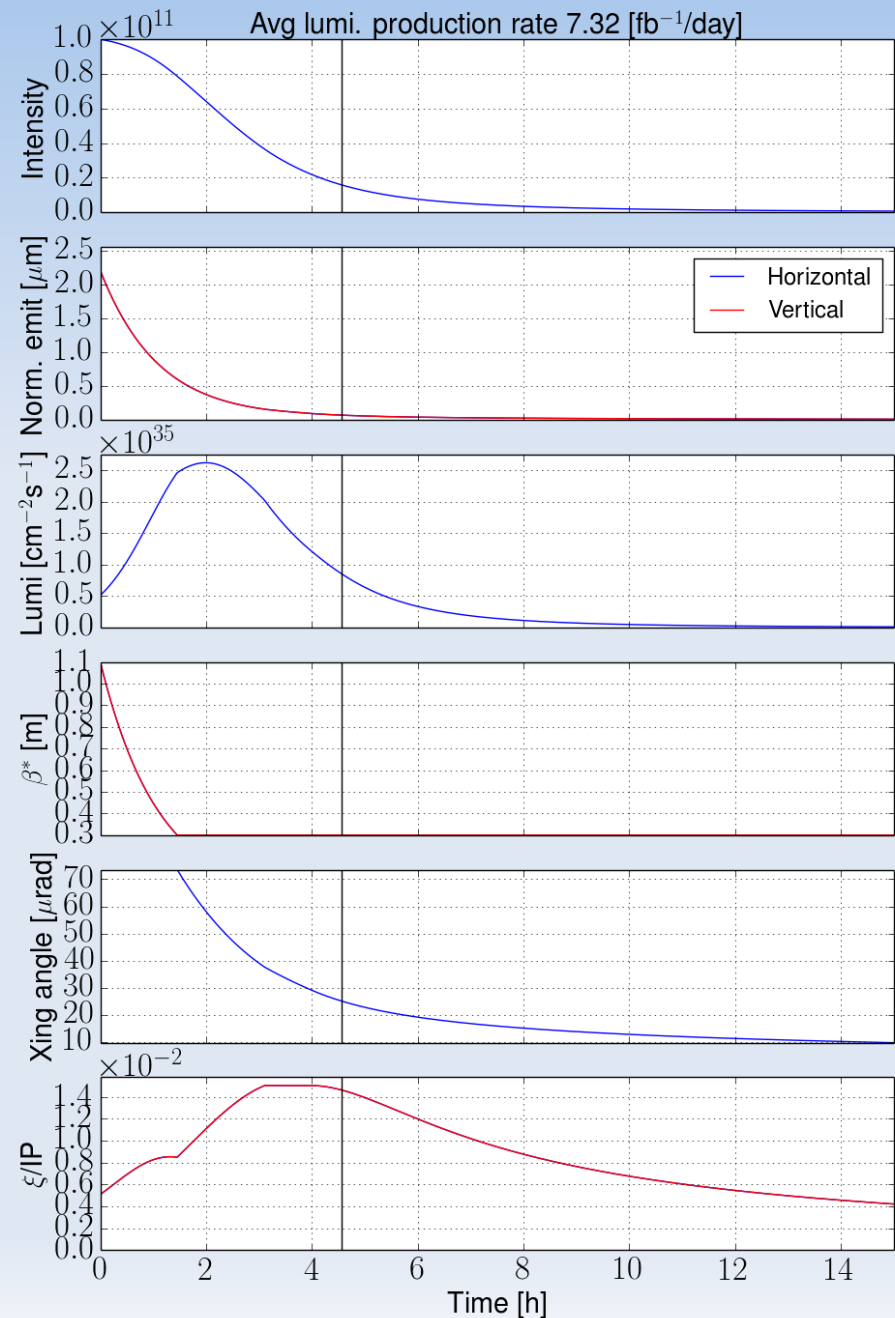




Beam parameter evolution



- Nominal configuration, with β^* decreasing to 0.3 m during the fill
 - 1.5x performance increase, but at the cost of very long fills
- Large performance gain with slightly larger beam-beam tune shift

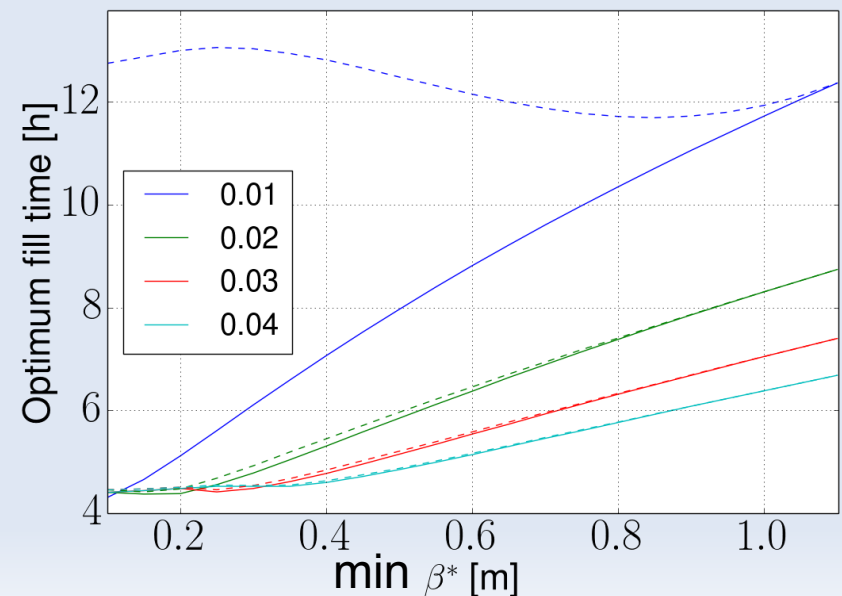
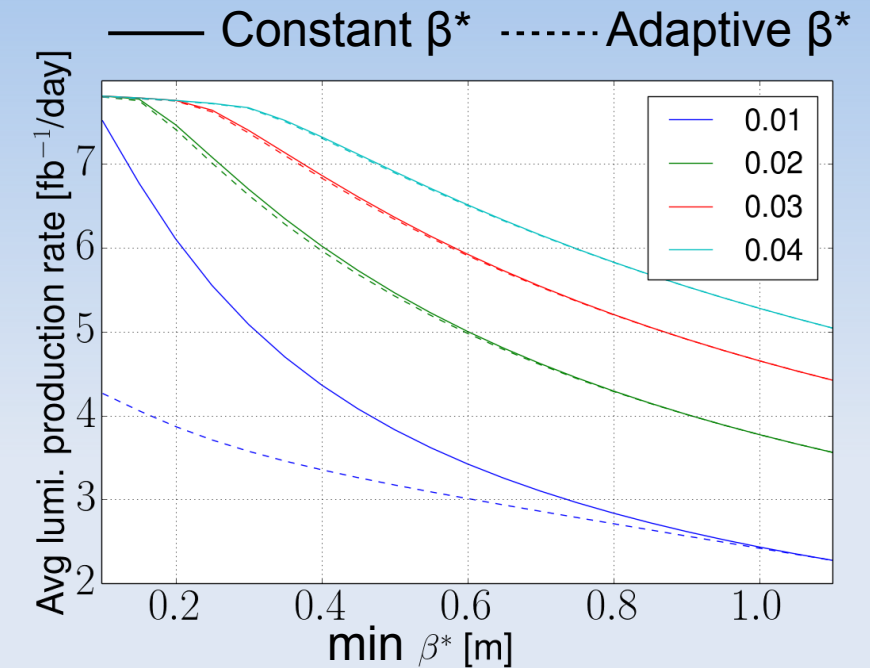




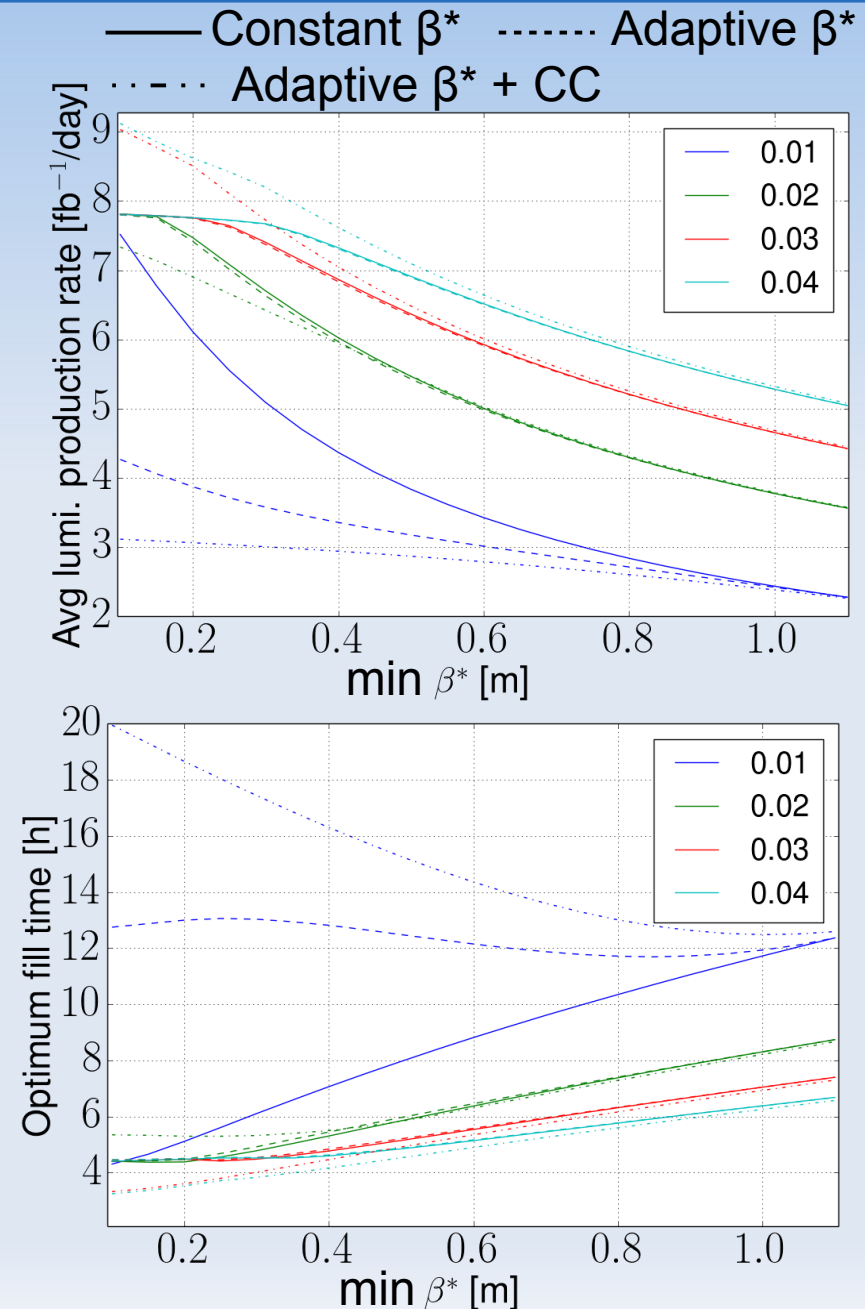
Beam parameter evolution

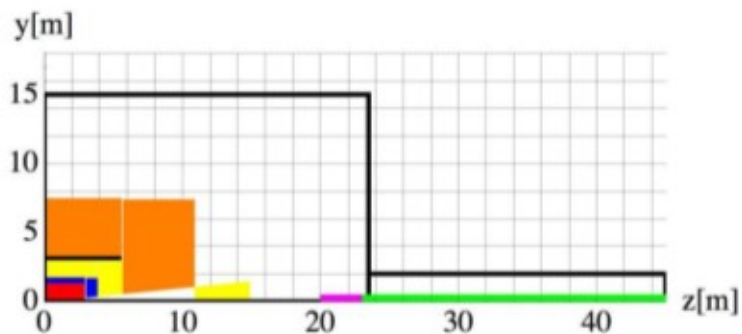


- Nominal configuration, with β^* decreasing to 0.3 m during the fill
 - 1.5x performance increase, but at the cost of very long fills
 - Large performance gain with slightly larger beam-beam tune shift
 - Similar to the configurations with constant β^*
 - But the smallest β^* is achieved with a small emittance
- Relaxed aperture requirements

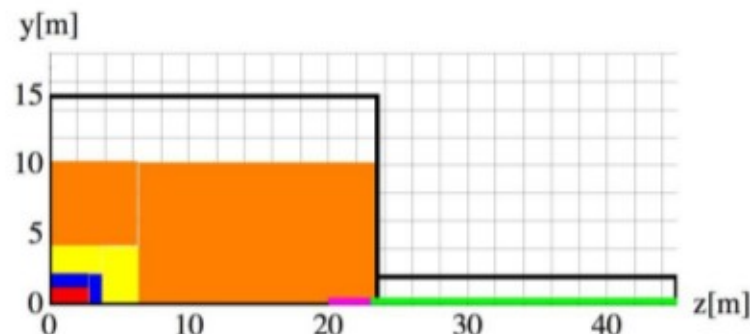


- Crab cavities are not helpful in configurations limited by the head-on beam-beam parameter
- Slight gain for $\beta^*_{\min} < 0.3$ m
 → Adapting the β^* allows to circumvent the needs for crab cavities, but the dynamics with large Piwinski angles has to be assessed





CMS & ATLAS



Twin Solenoid + Dipole, BL^2 scaled

Tracker $r=2.5m$ p_t reso 15% at 10TeV
 12 lambda ECAL+HCAL = 1m+2.5m
 Coil R=6m, 6T, Shielding Coil
 Forward Dipole 10Tm

Toroid + Dipole, BL^2 scaled

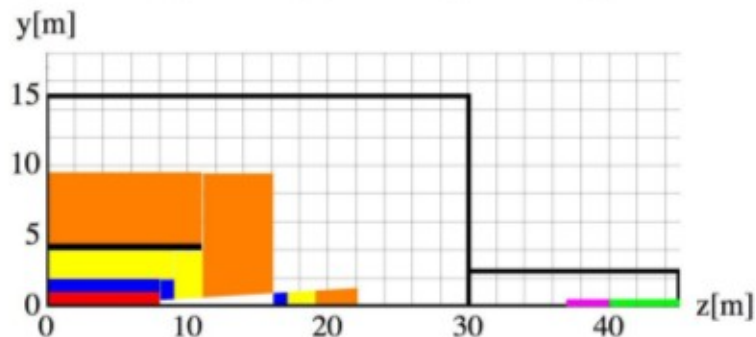
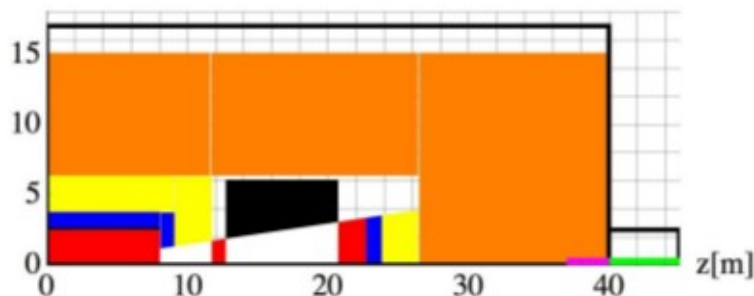
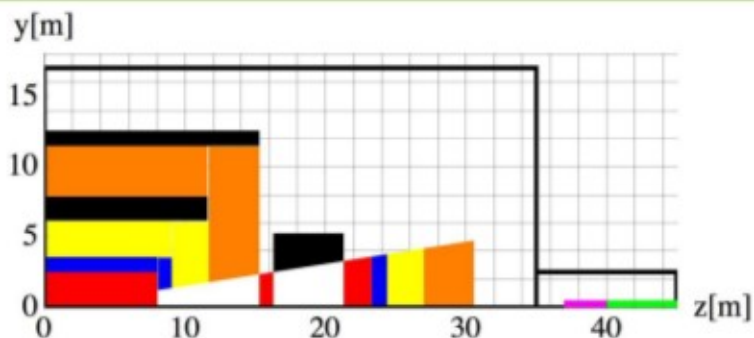
Tracker $r=2.5m$ p_t reso 15% at 10TeV
 Thin Coil R= 2.5m, B= 4T
 12 lambda ECAL+HCAL = 1m+2.5m
 Muon Toroid
 Forward Dipole 10Tm

CMS+, resolution scaled

Tracker $r=1.2m$ p_t reso 15% at 10TeV
 12 lambda ECAL+HCAL = 0.6m+2.2m
 Coil R=4m
 Iron Return Yoke

→ Extreme detector technology push

Tracker
Emcal
Hcal
Muon
Coil
TAS
Triplet





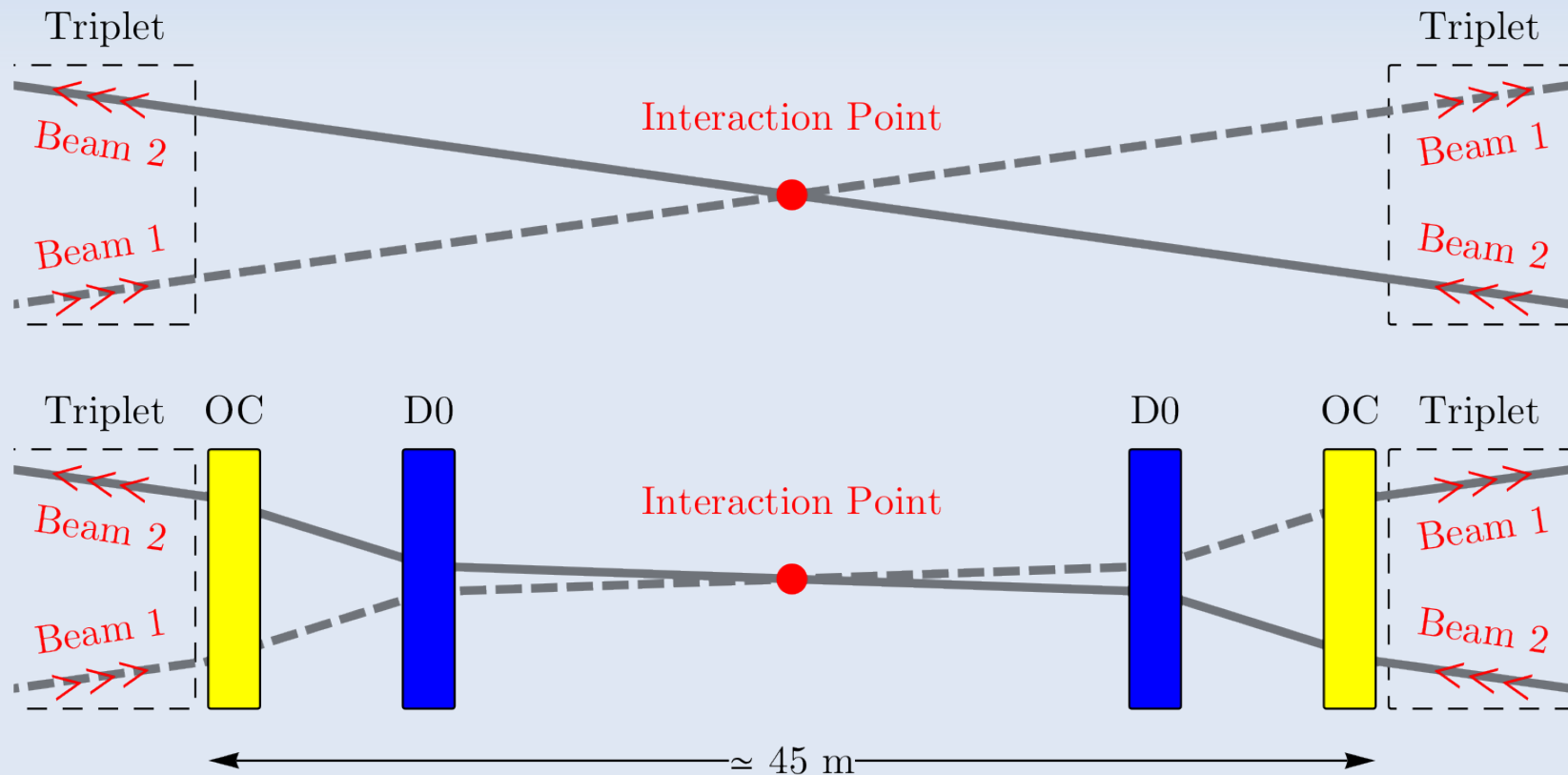
An early separation scheme for the LHC upgrade



- G. Sterbini, EPFL PhD thesis No 4574 (2010)

→ Place a dipole as close as possible to the IP in order to reduce the internal crossing angle keeping the same orbit in the triplet

- D0 integrated strength : 10-15 Tm





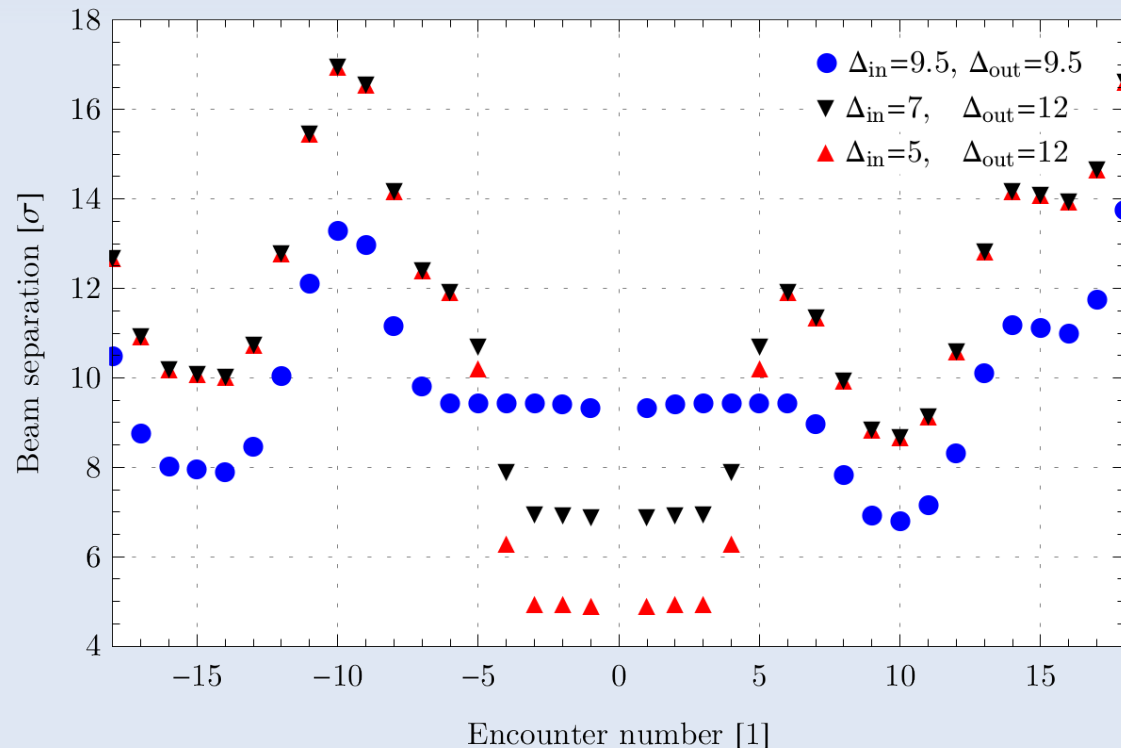
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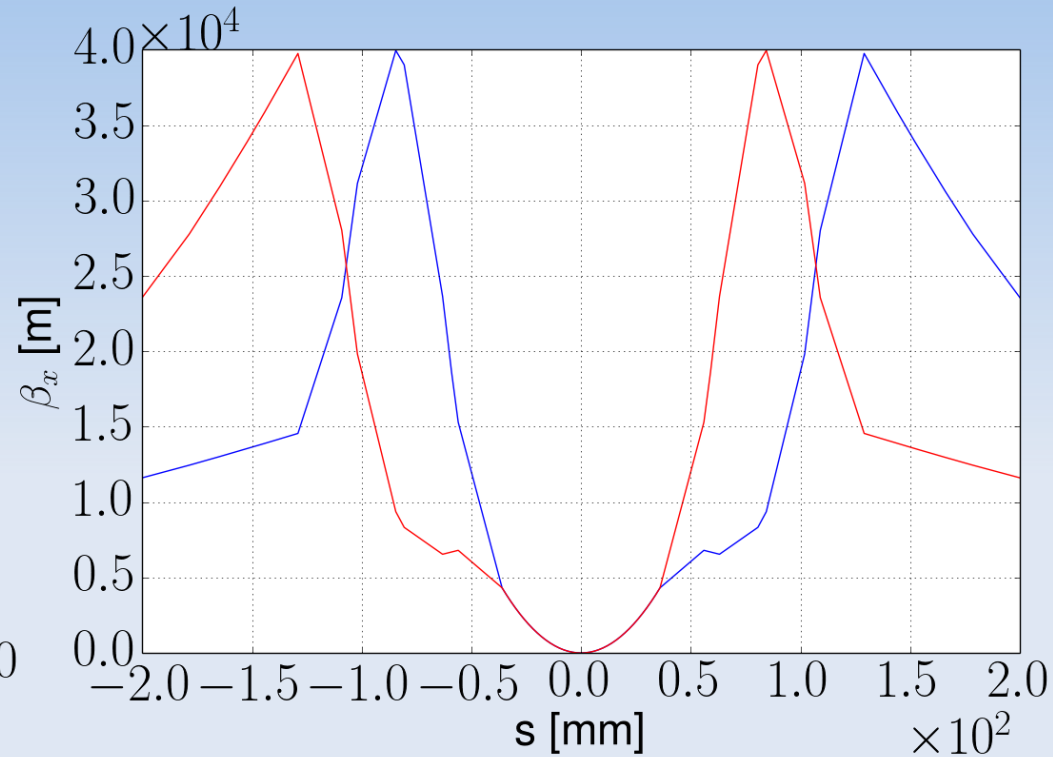
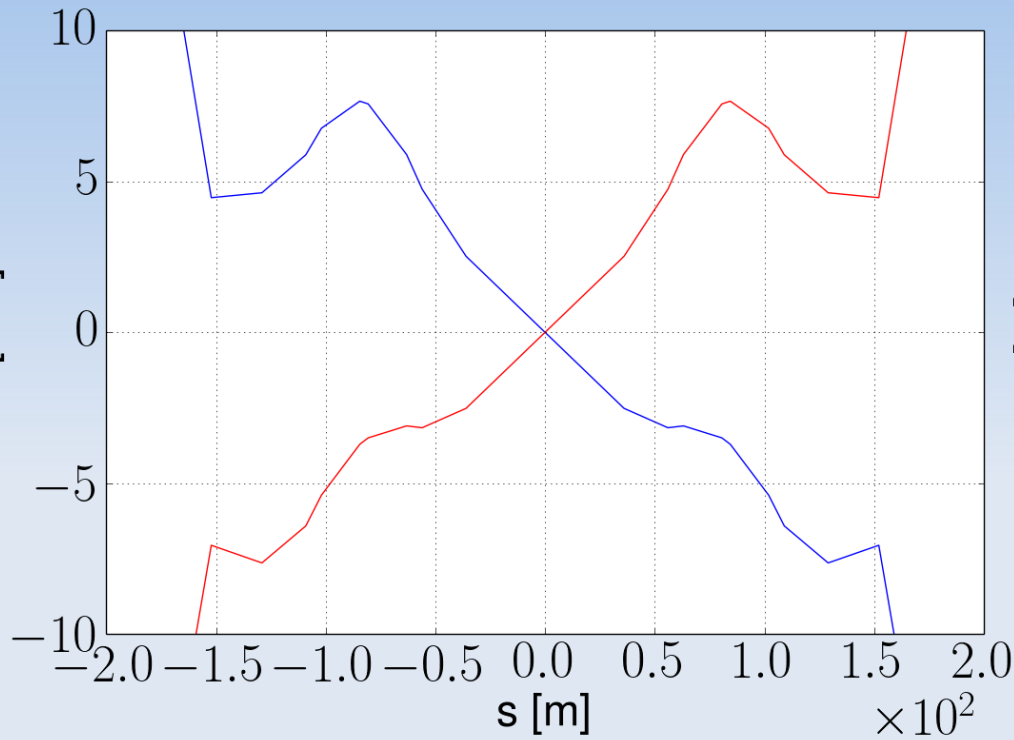
- Large impact on the separation between the beams

→ Similar long-range beam-beam 'strength' with lower geometric reduction factor





First try with the FCC IR

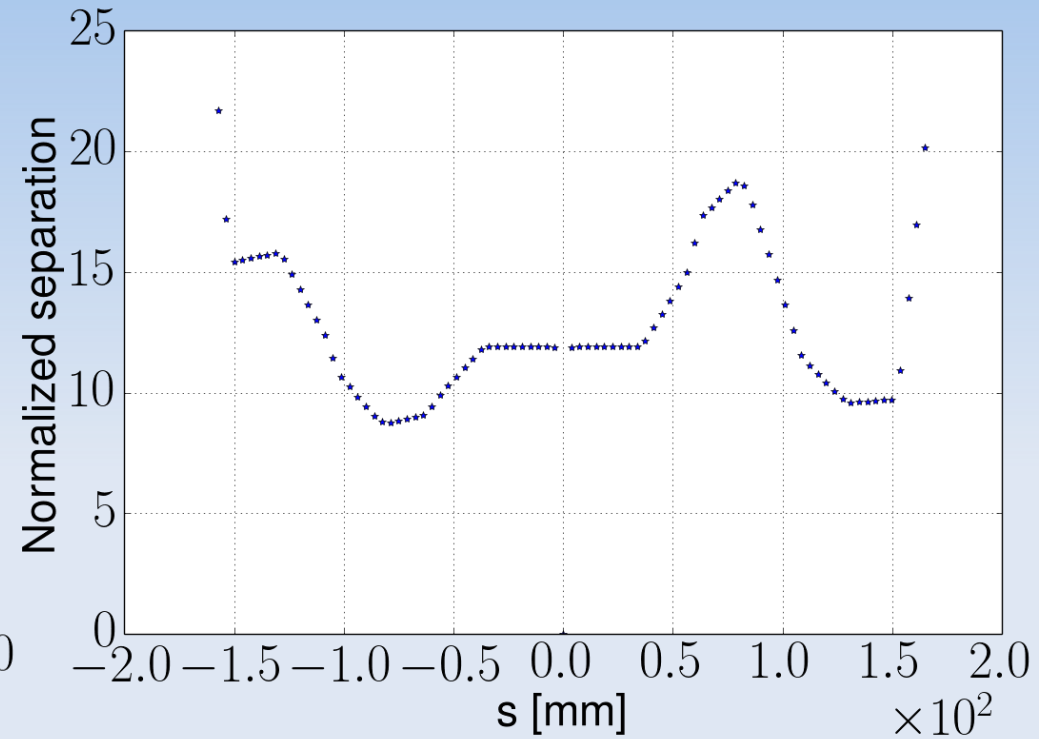
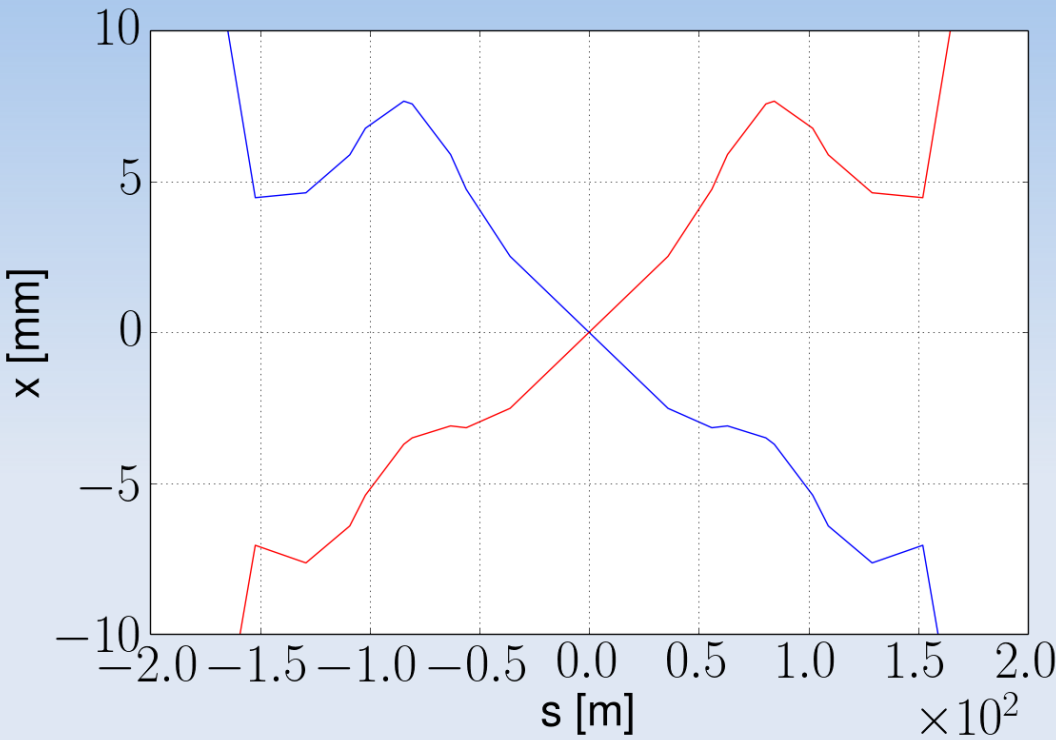


- Triplet first (scaled HL-LHC) $\beta^* = 0.3\text{m}$, $L^* = 36\text{m}$

→ $\theta_{\text{full}} = 70 \mu\text{rad}$, such that $S_{\text{BB}} = 12 \sigma$



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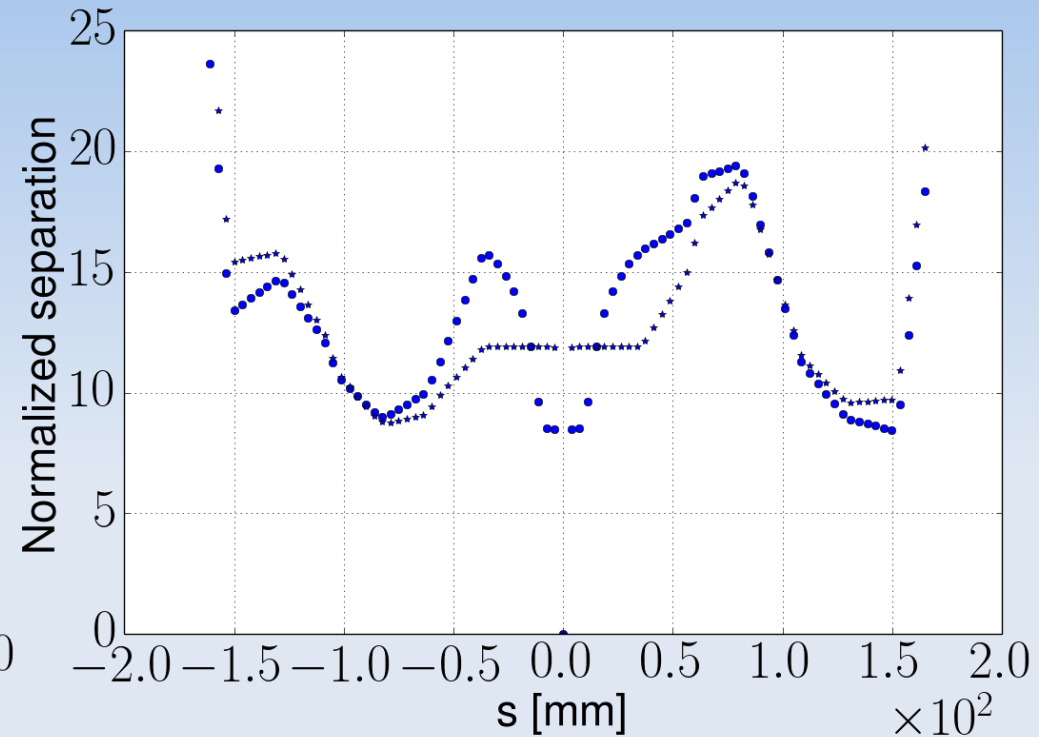
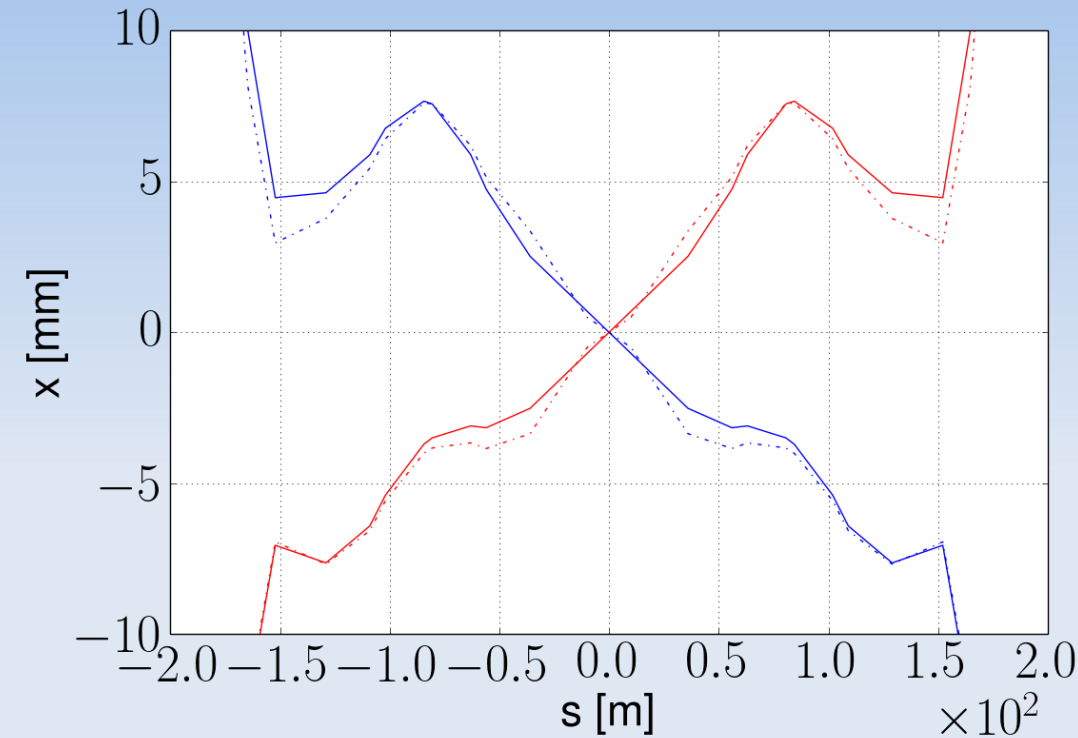


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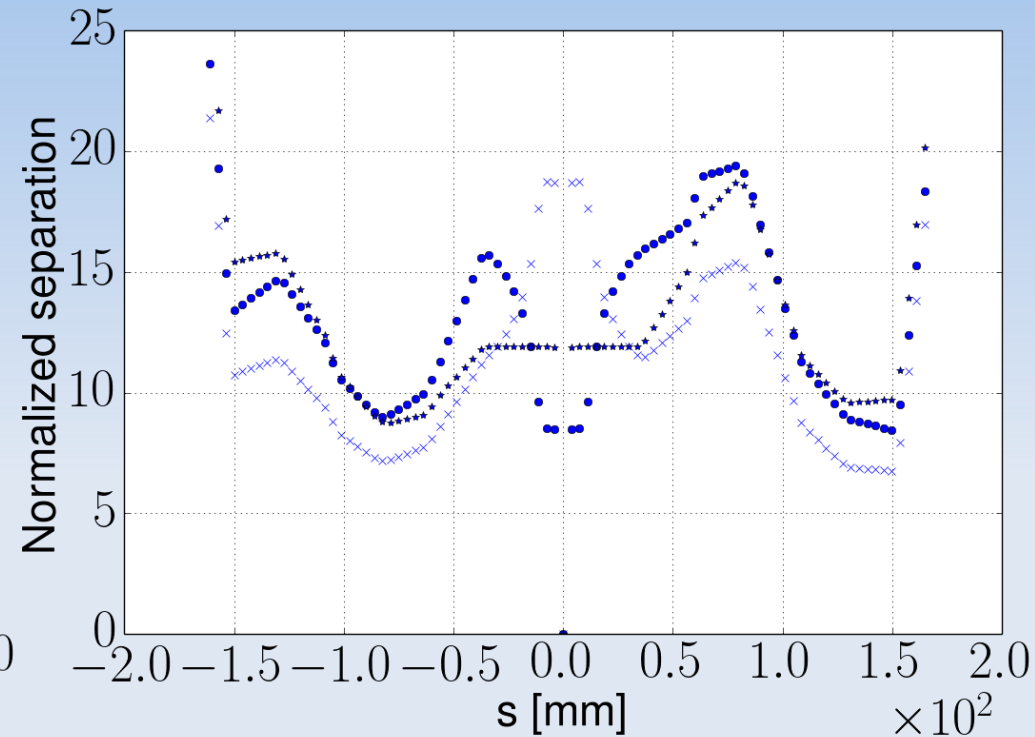
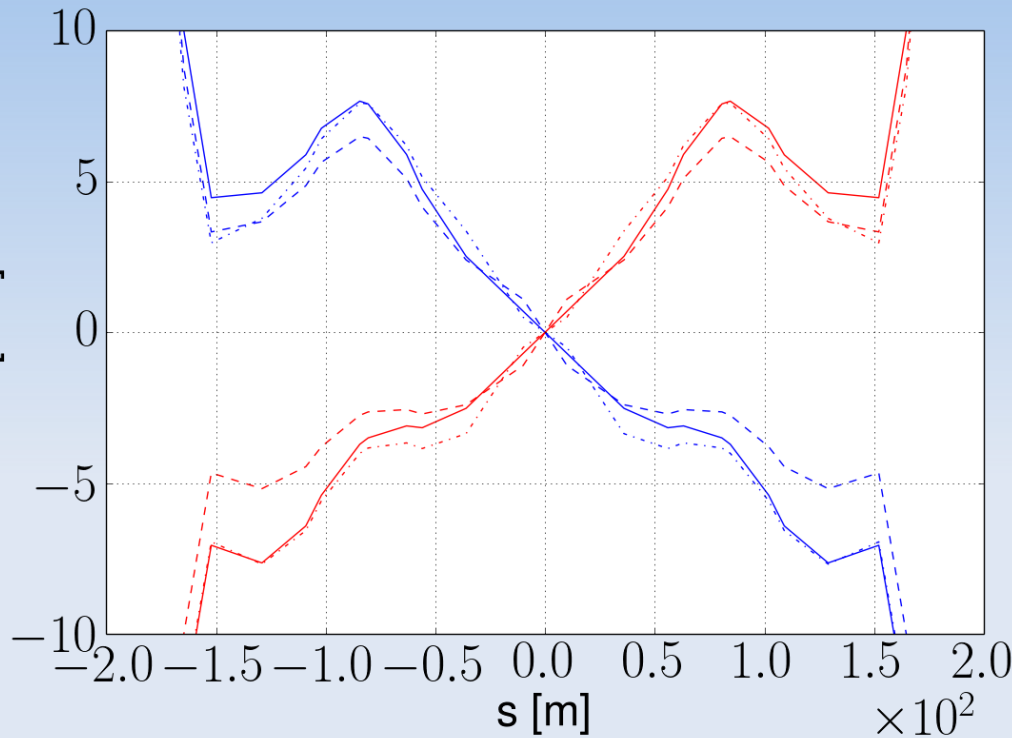
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$$\rightarrow \theta_{\text{full}} = 70 \mu\text{rad}, \text{ such that } S_{\text{BB}} = 12 \sigma$$

- A similar early separation scheme allows to reduce the internal Xing angle (10 Tm D0 at 10m from the IP)



First try with the FCC IR



- Triplet first (scaled HL-LHC) $\beta^* = 0.3\text{m}$, $L^* = 36\text{m}$
→ $\theta_{\text{full}} = 70 \mu\text{rad}$, such that $S_{\text{BB}} = 12 \sigma$
- A similar early separation scheme allows to reduce the internal Xing angle (10 Tm D0 at 10m from the IP)
- The reversed scheme allows to reduce the external Xing angle



Conclusion

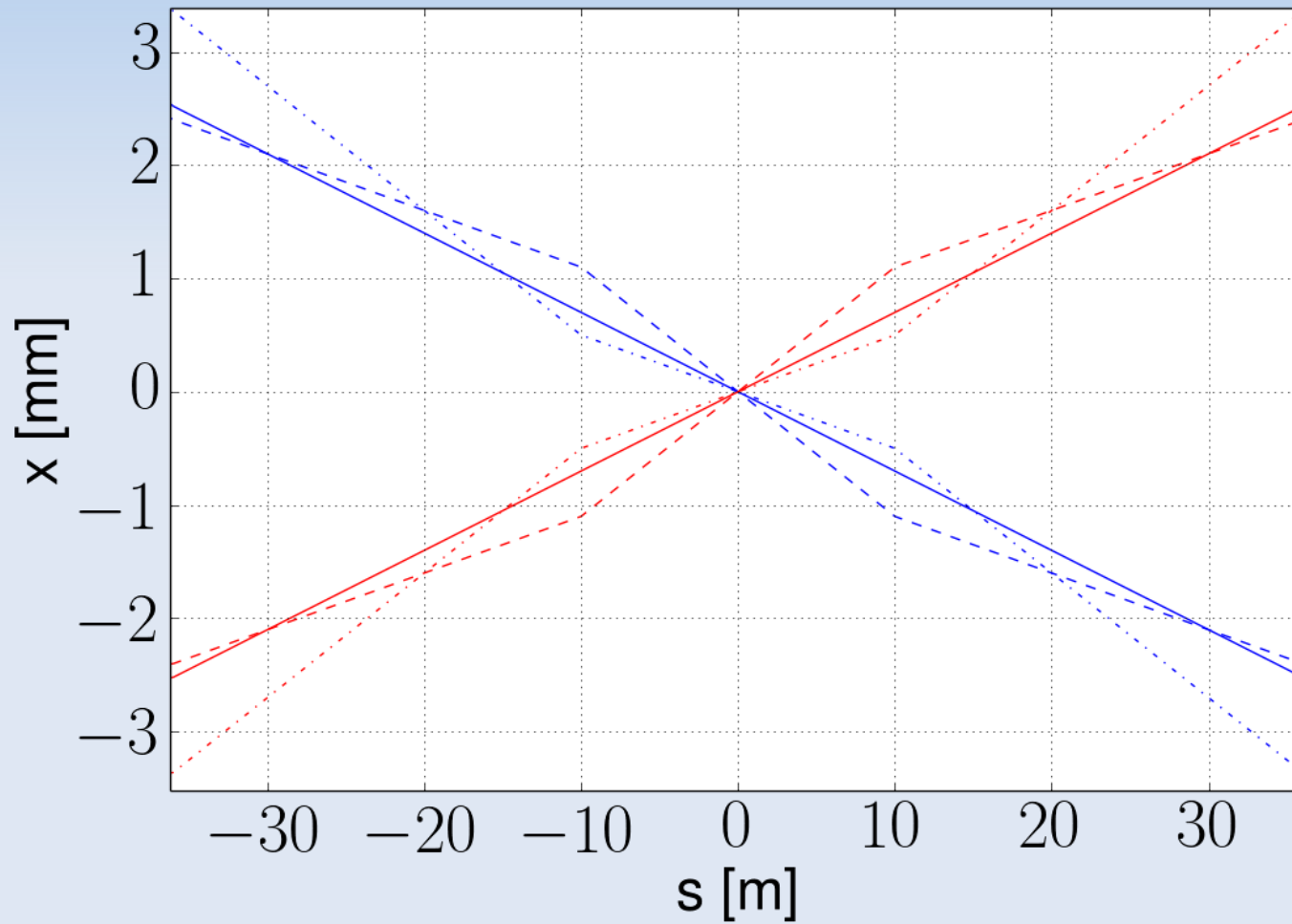


- Small β^* is clearly a key for the luminosity performance
 - Adapting β^* during the fill with a constant aperture requirements in the triplet offers a significant improvement
- Experimental spectrometers might be used to increase the performance
 - W/o crab cavity an early separation scheme could reduce the geometric reduction factor
 - With crab cavities, the reversed scheme might relax the aperture requirement in the triplet
 - The performance gain for both scheme should be quantified



BACKUP

Zoom on the drift space





BACKUP

Hourglass effect



- Hourglass is relevant for $\beta \sim \sigma_s = 0.08$ m
 - Only relevant for configurations with crab cavities

$$L_0 = \frac{n_b f_{rev} N^2}{4 \pi \epsilon \beta^*}$$

$$L = \frac{n_b f_{rev} \gamma N^2}{4 \pi \epsilon} \frac{\sqrt{\pi}}{\sigma_s} e^{\left(\frac{\beta^*}{\sigma_s}\right)^2} \int_{\frac{\beta^*}{\sigma_s}}^{\infty} e^{-x^2} dx$$

