

Muon Collider(s) & FCC-ee: Higgs physics

□ Outline

- ◆ Overview: Higgs boson production at lepton colliders
 - Processes
 - Energy
 - Luminosity

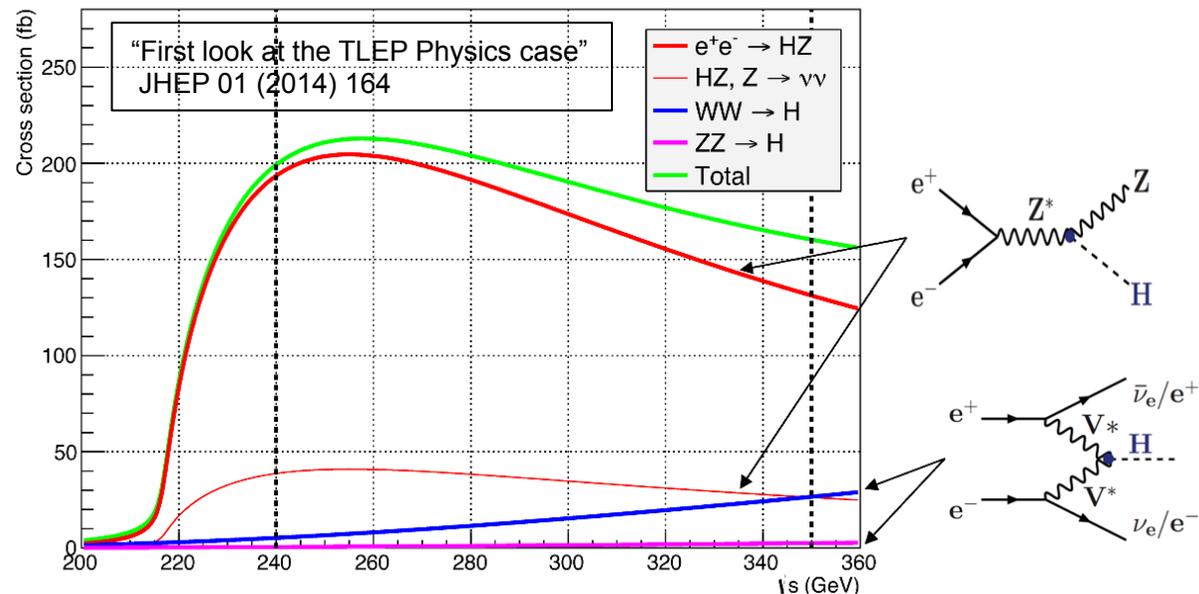
- ◆ Potential added value of muon colliders
 - Scan of the SM Higgs resonance
 - Beam energy and beam energy spread measurement
 - Additional Higgs bosons

- ◆ Summary

Higgs boson production (1)

□ Muons are leptons, like electrons

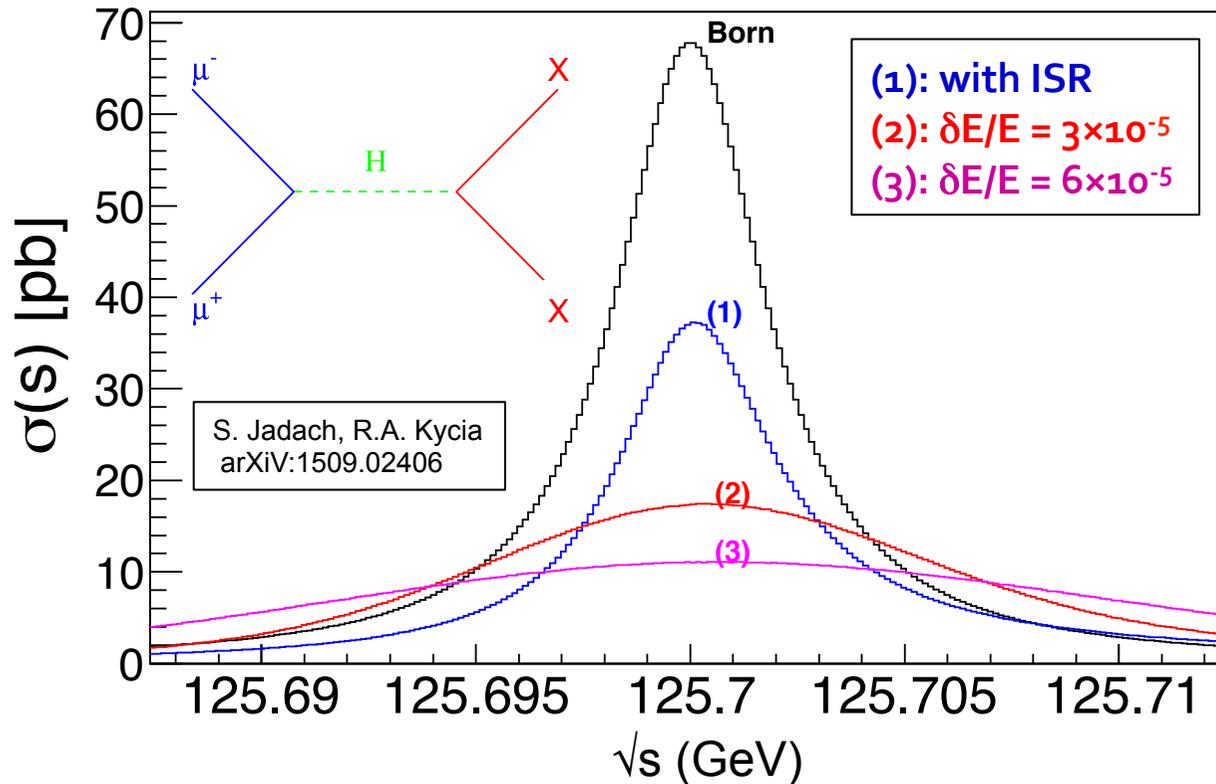
- ◆ Muon colliders can a priori do everything that e^+e^- colliders can do, e.g.:



- ◆ However, for a similar beam energy spread ($\delta E/E \sim 0.12\%$) at $\sqrt{s} = 240\text{-}350$ GeV
 - FCC-ee luminosity: $0.5 - 1.1 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$ / IP and up to 4 IPs
 - Muon collider luminosity: $0.4 - 1.5 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$ / IP and only one IP ($R \sim 50$ m!)
- ◆ Precision on branching ratios, couplings, width, mass, etc. , with 2 IPs
 - A factor 10 better at FCC-ee (and twice better at ILC) than at a muon collider

Higgs boson production (2)

- **Muons are heavy, unlike electrons: $m_\mu/m_e \sim 200$**
 - ◆ Large direct coupling to the Higgs boson: $\sigma(\mu^+\mu^- \rightarrow H) \sim 40,000 \times \sigma(e^+e^- \rightarrow H)$
 - ◆ Much less synchrotron radiation, hence potentially superb energy definition
 - $\delta E/E$ can be reduced to $3\text{-}4 \times 10^{-5}$ with more longitudinal cooling
 - ➔ Albeit with equivalent reduction of luminosity: $2 - 8 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$

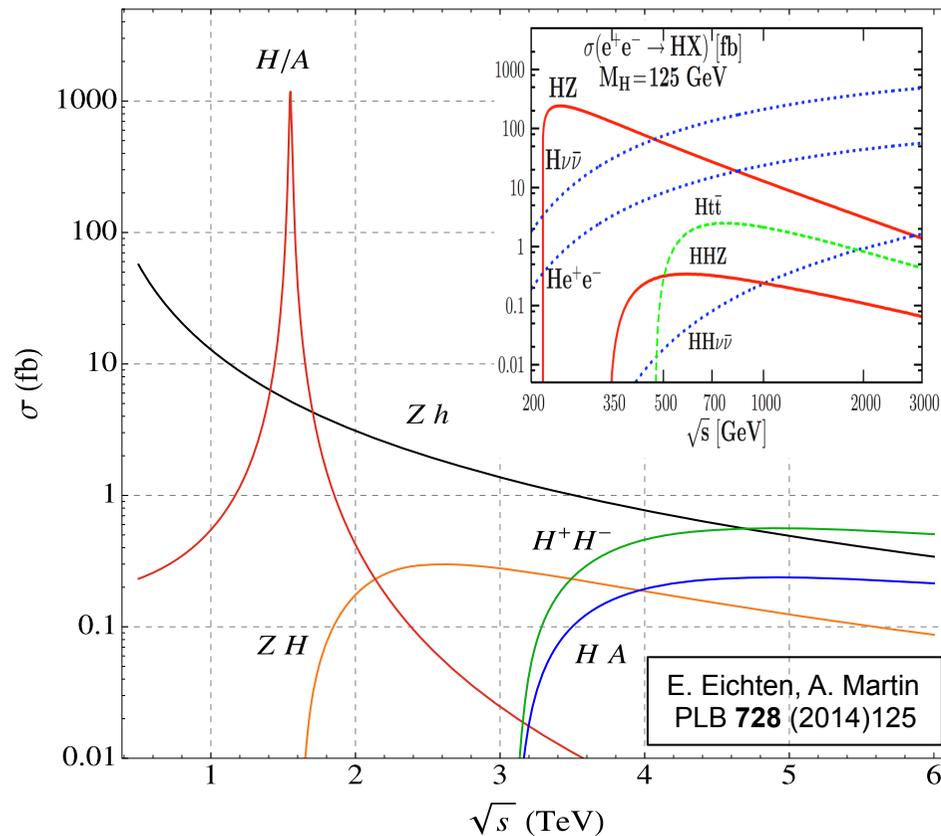
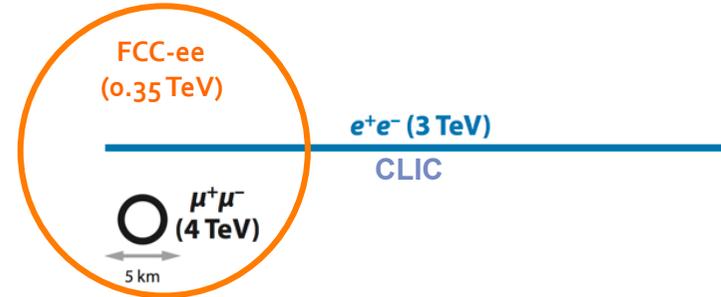


- $\sigma(\mu^+\mu^- \rightarrow H) \sim 15 \text{ pb}$
(ISR often forgotten: 28 pb)
- $200 - 800 \text{ pb}^{-1} / \text{yr}$
- $3000 - 12000 \text{ Higgs} / \text{yr}$

Reminder: At FCC-ee
400,000 to 800,000 Higgs/yr

Higgs boson production (3)

- **Muons are heavy, similar to protons**
 - ◆ **Limited synchrotron radiation**
 - **Can reach very high energy in small rings**



Luminosity

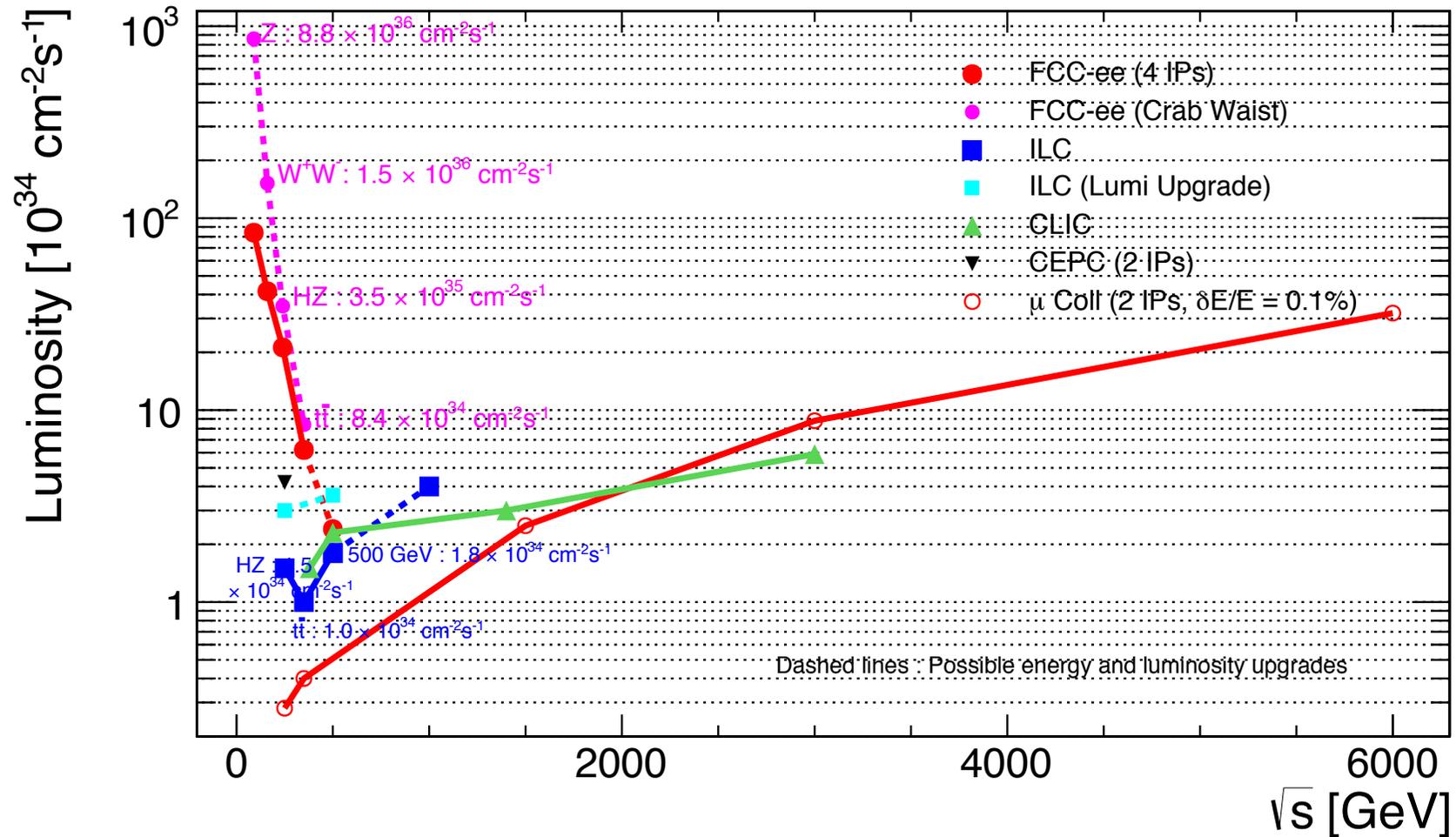
- Similar to linear colliders for $\sqrt{s} > 1$ TeV
 - **HHH coupling with similar precision**
 - (Also done at FCC-hh)

Energy

- Can go to higher energy
 - **Advantage for 2HDM (e.g., SUSY)**
 - **Heavy Higgs with $\mu^+\mu^- \rightarrow H, A$**
- $\sqrt{s} \sim 6$ TeV possible in the Tevatron tunnel

Higgs boson production (4)

Luminosity and energy at lepton colliders



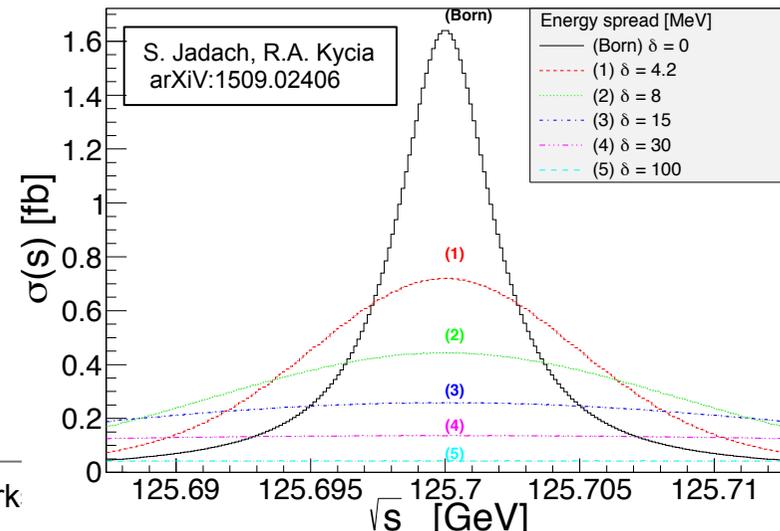
Higgs boson production (5)

- **When muon colliders try to do what the other collider do**
 - ◆ They do a factor 10 worse than FCC-ee, a factor 2 worse than ILC at $\sqrt{s} < 500$ GeV
 - ◆ They reach a performance similar to that of CLIC at $1.5 \text{ TeV} < \sqrt{s} < 3 \text{ TeV}$

- **Muon colliders are, however, unique on two aspects**
 - ◆ They can scan the Higgs resonance(s) by direct $\mu^+\mu^- \rightarrow H$ production (*)
 - ◆ They can reach $\sqrt{s} > 3 \text{ TeV}$ with reasonably-sized rings, and low power consumption
 - I focus on these specificities in the rest of the presentation

□ **(*) Note: Scan of the SM Higgs resonance at FCC-ee with $e^+e^- \rightarrow H$**

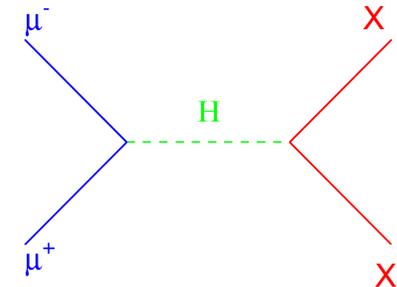
- ◆ $\sigma(e^+e^- \rightarrow H) \sim 50 \text{ ab}$ with nominal $\delta E/E$
 - 200- 500 ab with monochromators?
 - $\sim 10,000$ events/yr ($40 \text{ ab}^{-1}/\text{yr}$)
 - $S/B \times 10^{-5}$ wrt muon collider
 - No way to study resonance parameters
 - See study from d'Enterria et al.
 - Optimistically expect $\sim 2\sigma$ excess



Scan of the SM Higgs resonance (1)

Resonant production

$$\sigma(\mu^+\mu^- \rightarrow H^0) = \frac{4\pi\Gamma_H^2 Br(H^0 \rightarrow \mu^+\mu^-)}{(\hat{s} - M_H^2)^2 + \Gamma_H^2 M_H^2}$$



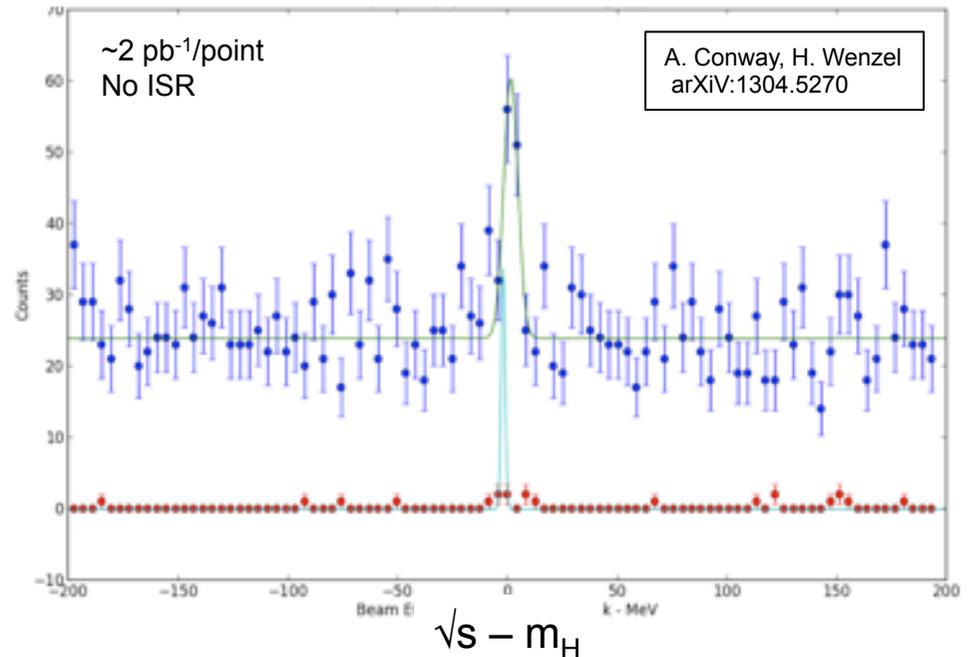
Major background:
 $\mu^+\mu^- \rightarrow Z/\gamma^* \rightarrow XX$

- ◆ Convoluted with
 - Beam energy spectrum
 - Initial state radiation (ignored in most studies)
- ◆ The measurement of the lineshape gives access to
 - The Higgs mass, m_H
 - The Higgs width, Γ_H
 - The branching ratio into $\mu^+\mu^-$, $BR(H \rightarrow \mu\mu)$
 - ➔ Hence, the coupling of the Higgs to the muon, $g_{H\mu\mu}$
 - Some branching fractions and couplings, with exclusive decays

Scan of the SM Higgs resonance (2)

□ Finding the resonance ($\Gamma_H = 4.2 \text{ MeV} \sim \delta E$)

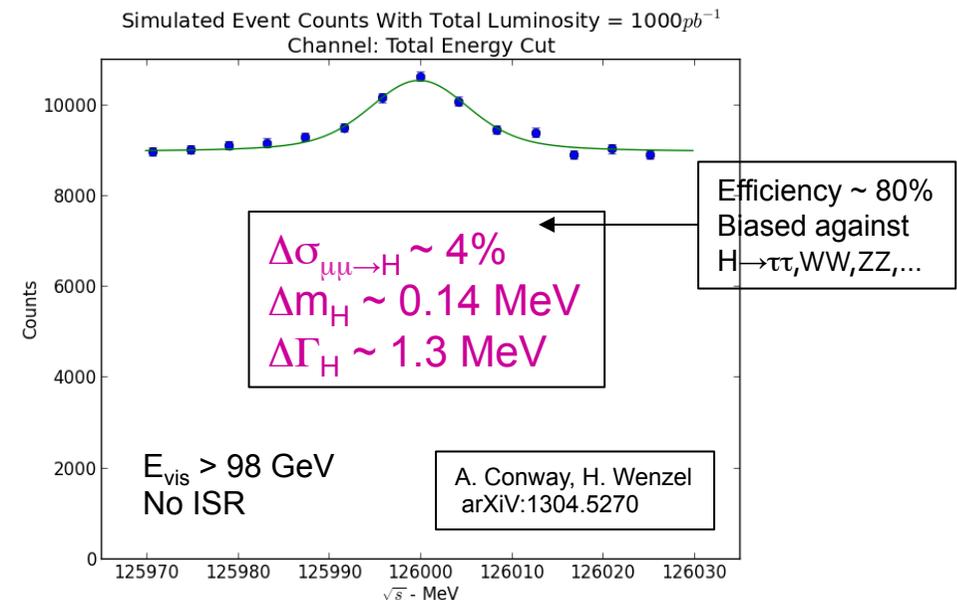
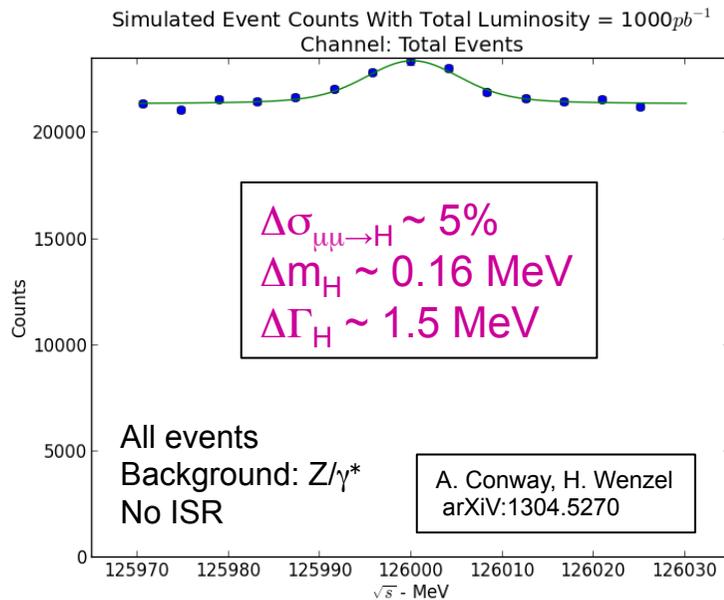
- ◆ Today, m_H is known to $\pm 250 \text{ MeV}$
 - Improves to $\pm 100 \text{ MeV}$ (LHC14), $\pm 30 \text{ MeV}$ (ILC), or $\pm 8 \text{ MeV}$ (FCC-ee)
- ◆ Scan the \sqrt{s} region of interest in optimal bins of 4.2 MeV
 - Count the number of bb and semi-leptonic WW events (see next slides)
- ◆ Without ISR, needs about 2 pb^{-1} / point for a 5σ significance
 - Reduced to 3σ when ISR is included
 - ➔ Probably enough
- ◆ Total luminosity needed for 3σ
 - 300 pb^{-1} (1.5 yr) for $\pm 300 \text{ MeV}$
 - 90 pb^{-1} (6 months) for $\pm 90 \text{ MeV}$
 - 25 pb^{-1} (2 months) for $\pm 24 \text{ MeV}$
 - ➔ With $L = 2 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$
- ◆ Can be long ...
 - ... but feasible
 - ➔ Especially after ILC / FCC-ee



Scan of the SM Higgs resonance (3)

Measurement of the lineshape

- ◆ Assume 1 fb^{-1} (5 yrs at 2×10^{31} and ≥ 1 yr at 8×10^{31}) : 70 pb^{-1} / point around m_H
 - The detector is assumed to have the performance of an ILC detector
 - No beam background (e.g., from muon decays) is simulated
- ◆ Count either all events, or only those with $E_{\text{vis}} > 98 \text{ GeV}$ [reject $Z(\gamma)$ events]

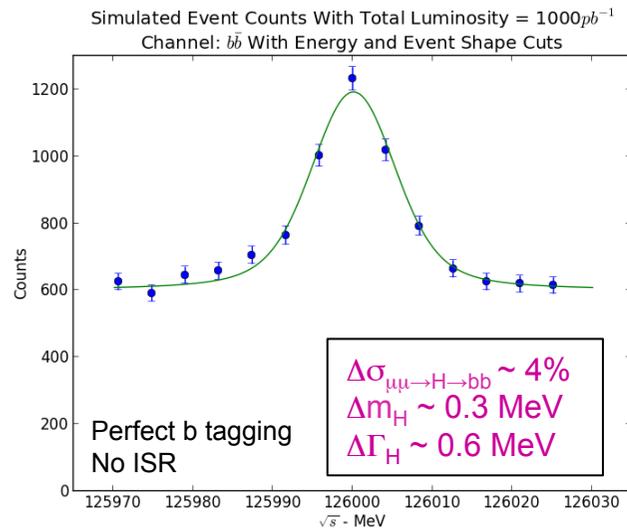


- ISR reduces the signal by a factor 2 (but not the background)
 - ➔ All errors to be increased by a factor 2
- m_H and Γ_H measurements require knowledge of E and δE with great precision

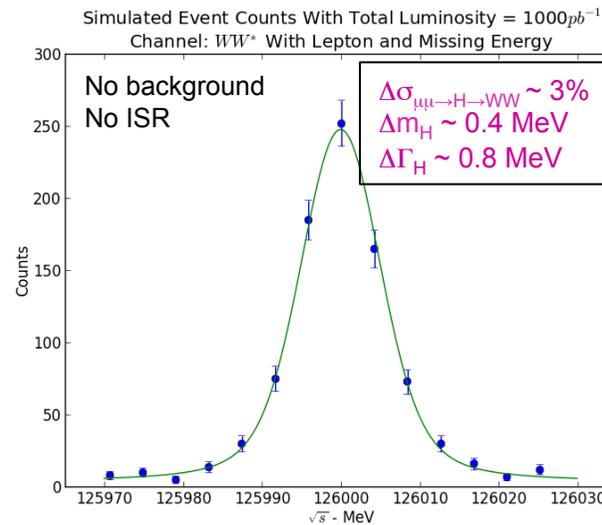
Scan of the SM Higgs resonance (4)

Exclusive decays

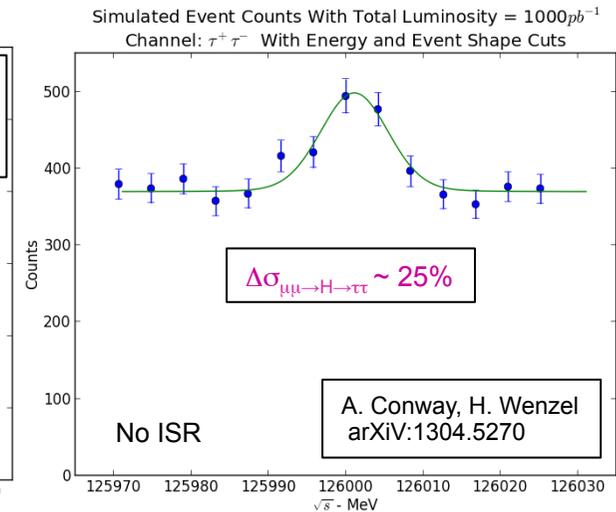
◆ $H \rightarrow bb$



$H \rightarrow WW \rightarrow l\nu qq$



$H \rightarrow \tau\tau$



◆ Notes

- Some optimism in these numbers (perfect b tag, only Z bkgd, no beam bkgd...)
- Errors to be increased anyway to account for ISR
- A better scan strategy could be thought of (less in the sides, more in the peak)
- The numbers are for 5 years at low luminosity, and 1.2 year after lumi upgrade
 - ➔ Combined numbers (next slide) given for 5 (low lumi) + 5 (upgrade) years.

Scan of the SM Higgs resonance (5)

□ Summary of precision measurements (after ~10 years of running)

Error on	$\mu\mu$ Collider	ILC	FCC-ee
m_H (MeV)	0.06	30	8
Γ_H (MeV)	0.17	0.16	0.04
g_{Hbb}	2.3%	1.5%	0.4%
g_{HWW}	2.2%	0.8%	0.2%
$g_{H\tau\tau}$	5%	1.9%	0.5%
$g_{H\gamma\gamma}$	10%	7.8%	1.5%
$g_{H\mu\mu}$	2.1%	20%	6.2%
g_{HZZ}	–	0.6%	0.15%
g_{Hcc}	–	2.7%	0.7%
g_{Hgg}	–	2.3%	0.8%
BR_{invis}	–	<0.5%	<0.1%

Not sure of the practical use of such a precision on m_H

The Higgs width is best measured at ee colliders

These Higgs couplings are best measured at ee colliders

The SM Higgs coupling to muons is *the* added value of a $\mu\mu$ collider *

These Higgs couplings are *only* measured at ee colliders *

* pp colliders have their say, too

◆ Note: $BR(H \rightarrow \mu\mu)$ can also be measured with % precision at FCC-hh. (Will be already 10% after LHC.)

Beam energy and beam-energy spread (1)

- Muons are naturally 100% polarized (from π^\pm decays)
 - ◆ It is hoped that ~20% of this polarization can be kept in the collider ring

- Then, the spin precesses around B with a frequency ν_0
 - ➔ For $m_H = 125$ GeV, $\nu_0 = 0.68967593(35)$
- Without energy spread, P_L oscillates between -20% and +20%
- With energy spread, P_L gets diluted turn after turn

$$\nu_0 = \frac{g_\mu - 2}{2} \times \frac{E_{\text{Beam}}}{m_\mu}$$

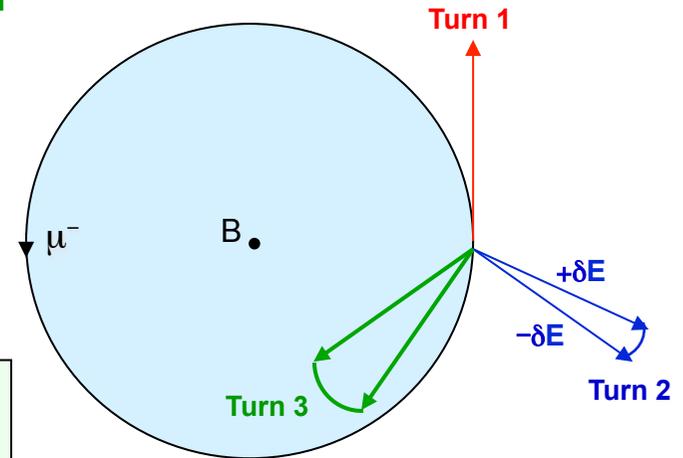
$$P_L(T) = P_0 \int_0^\infty \cos(2\pi\nu T) S(\nu) d\nu$$

➔ $P_L(T)$ is the Fourier transform of $S(\nu)$

- For example, with a Gaussian energy spread

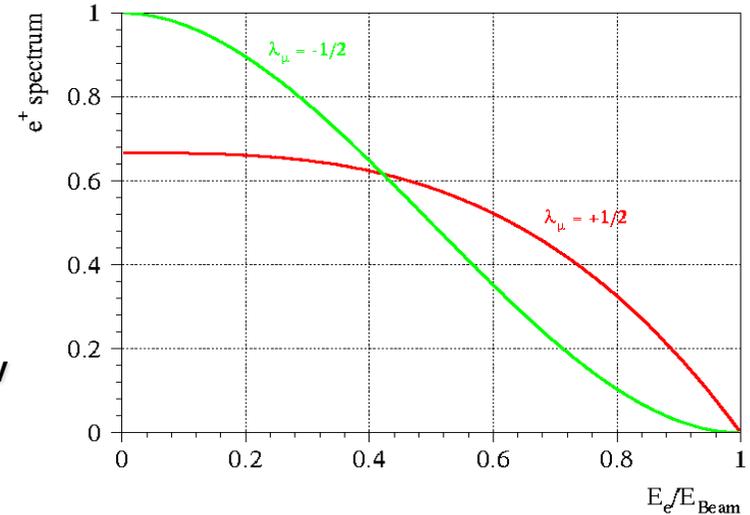
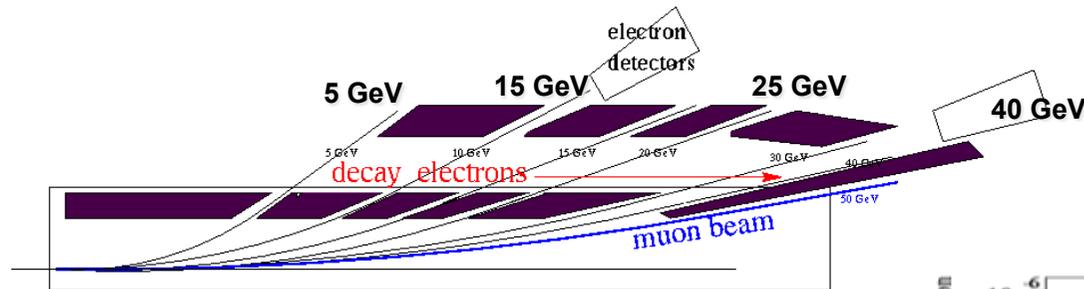
$$P_L(T) = P_0 \cos(2\pi\nu_0 T) \exp\left\{-\frac{1}{2} \left[2\pi\nu_0 T \frac{\delta E}{E}\right]^2\right\}$$

- Experimentally, measure P_L at each turn T
 - ➔ And deduce the complete beam energy spectrum by inverse Fourier transform
i.e., $\delta E/E$ for a Gaussian energy spread



Beam energy and beam-energy spread (2)

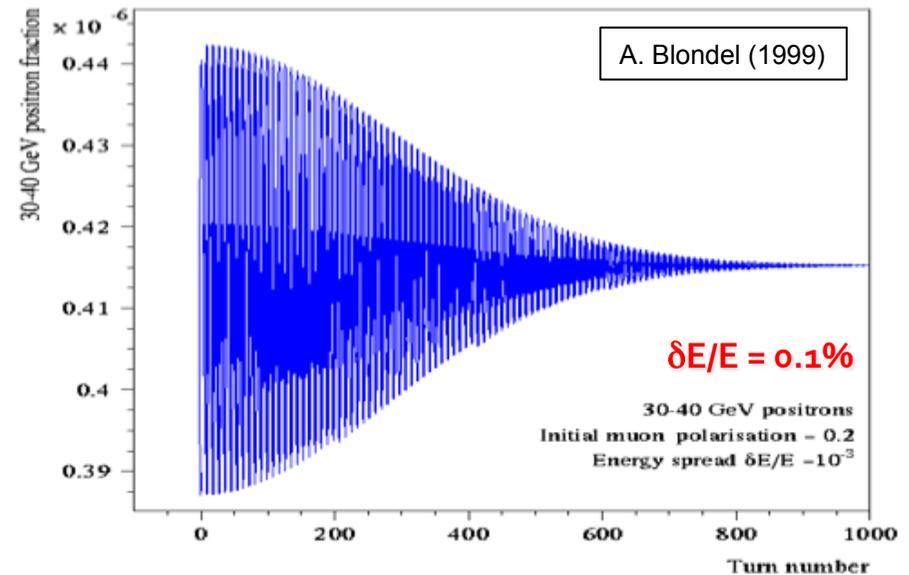
- Use decay electrons to measure $P_L(T)$
 - ◆ Energy distribution depends on the muon helicity
 - $N_e(E) / N_{\text{tot}}$ oscillates according to P_L
 - ➔ Count electrons in the first dipole:



- ◆ Fraction of e^+ from 30 to 40 GeV

$$P_L(T) = P_0 \cos(2\pi\nu_0 T) \exp\left\{-\frac{1}{2} \left[2\pi\nu_0 T \frac{\delta E}{E}\right]^2\right\}$$

- The amplitude gives P_0
- The frequency gives $\nu_0 (E_{\text{Beam}})$
- The damping gives $\delta E/E$



Beam energy and beam-energy spread (3)

- **Expected statistical accuracy of the method**
 - ◆ For $L = 2 \times 10^{31} \text{ cm}^{-2}\text{s}^{-1}$ and $\delta E/E = 3 \times 10^{-5}$, for each “fill” (i.e., 1000 turns)
 - 10^{-7} on the beam energy (6 keV)
 - ➔ Limited to 5×10^{-7} (30 keV) by the precision on $g_{\mu} - 2$ (!)
 - $3 \cdot 10^{-7}$ on the beam energy spread $\delta E/E$ (1%)
 - ➔ Corresponds to a systematic uncertainty of 0.5% on $\sigma(\mu\mu \rightarrow H)$
 - ➔ Corresponds to a systematic uncertainty of 50 keV on Γ_H
 - 10^{-4} on the polarization value
 - ➔ Negligible impact on $\sigma(\mu\mu \rightarrow H)$
 - ◆ These uncertainties are appropriately smaller than the statistical precision
 - On the Higgs mass (60 keV)
 - On the Higgs width (170 keV)
 - On the production cross section (1.5%)

Additional Higgs bosons (1)

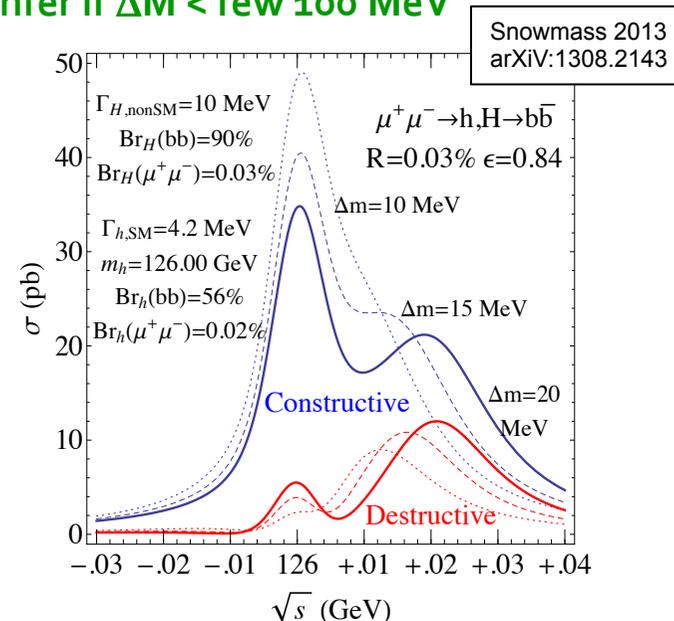
□ Is H(125) made of several quasi-degenerate Higgs bosons ?

- ◆ At LHC, the typical m_H resolution in the $H \rightarrow ZZ^* \rightarrow \mu\mu$ channel is ~ 1 GeV
 - Two quasi-degenerate Higgs bosons difficult to infer if $\Delta M < \text{few } 100 \text{ MeV}$

Similar at FCC-ee
(Recoil mass)

◆ Would be a piece of cake at a muon collider

- Examples shown for
 - ➔ $\Delta M = 10, 15, 20 \text{ MeV}$
 - ➔ Destructive/constructive interference
 - ➔ Similar coupling to muons and b quarks
- Lineshape sensitive to $\Delta M \sim \text{MeV}$
 - ➔ If both Higgs bosons couple to μ and b/W



Snowmass 2013
arXIV:1308.2143

- ◆ Probably observable at FCC-ee via pair production with $\sqrt{s} > 250 \text{ GeV}$ (to be studied)
 - $e^+e^- \rightarrow hA$ present at tree level with large cross section (A pseudoscalar, $m_A \sim m_h \sim m_H$)
 - [$e^+e^- \rightarrow hH$ only at loop level with a few ab cross section (H scalar)]
 - ➔ A small mass difference is not measurable this way

A. Djouadi et al.
PRD **54** (1996) 759

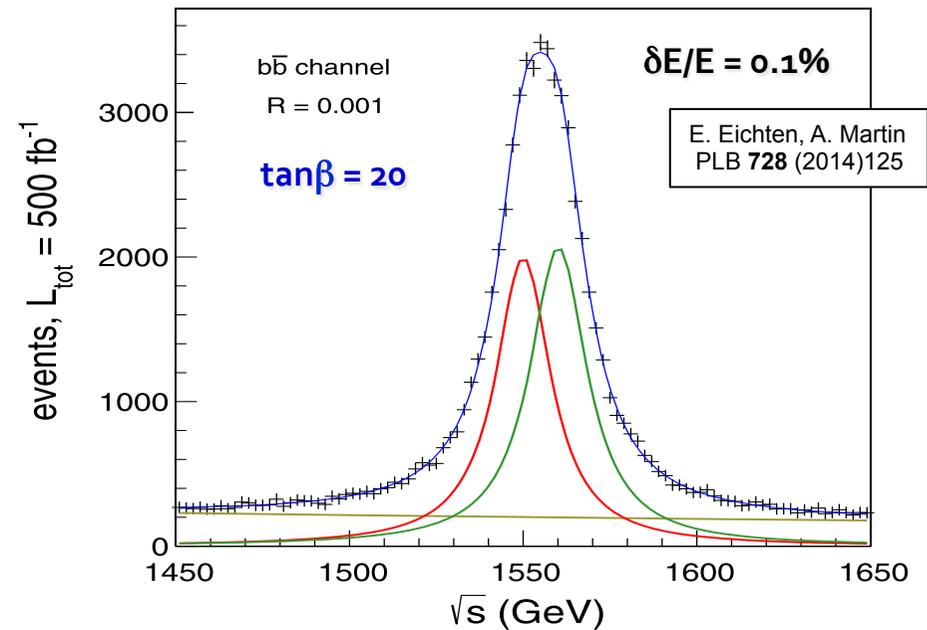
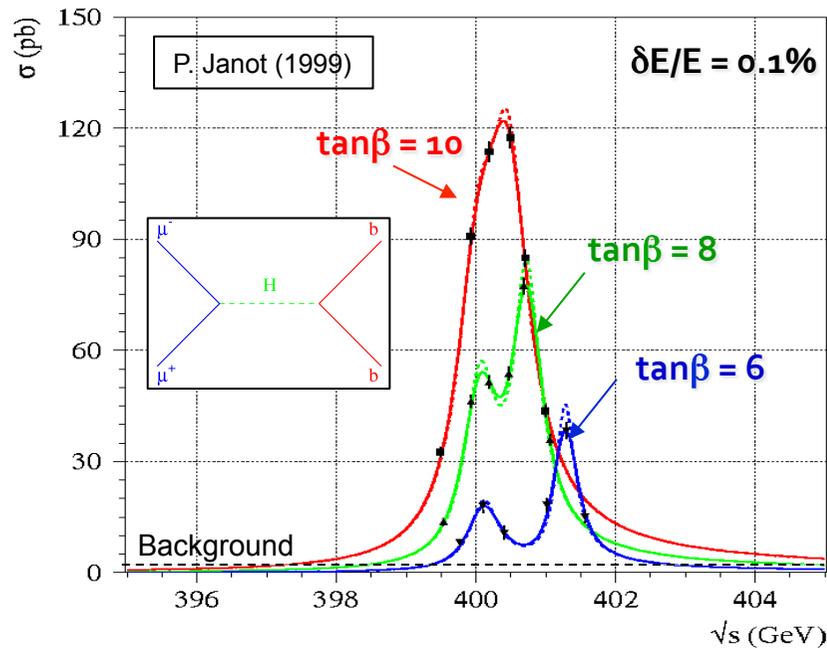
... but the pair production proves the existence of two (three) states

Additional Higgs bosons (2)

- Can be applied to heavier H and A in 2HDM (e.g., from SUSY)

- Example 1: $m_A = 400$ GeV

- Example 2: $m_A = 1.55$ TeV



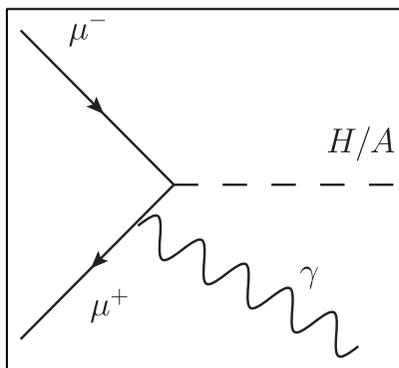
- Notes:

- Higgs width of the order of 0.1 to 1% of the Higgs mass
 - $\delta E/E \sim 0.1\%$ enough, large integrated luminosities (100's fb^{-1} or ab^{-1}) possible
 - Each value of m_A correspond to a specific ring diameter
 - Need to know the mass beforehand

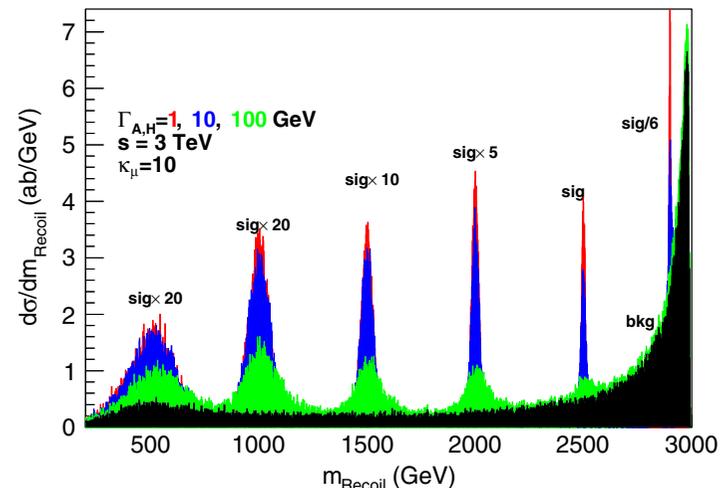
Additional Higgs bosons (3)

Automatic mass scan with radiative returns in $\mu\mu$ collisions

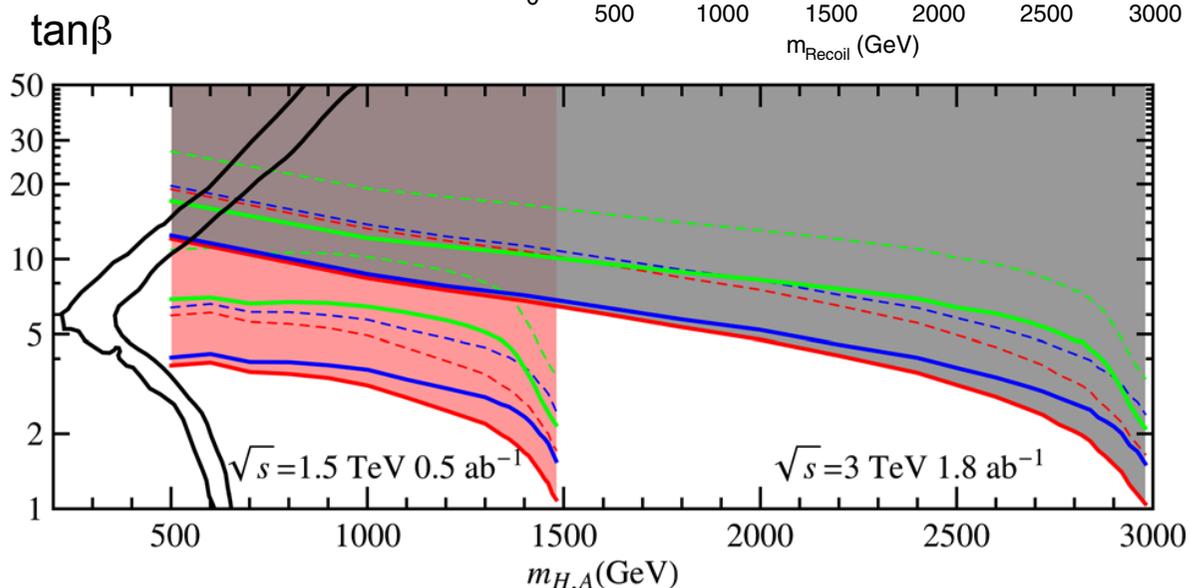
- ◆ Go to the highest energy first
 - $\sqrt{s} = 1.5, 3$ or 6 TeV
- ◆ Select event with an energetic photon
 - Check the recoil mass $m_{\text{Recoil}} = [s - 2E_\gamma\sqrt{s}]^{1/2}$



N. Chakrabarty et al.
PRD **91** (2015)015008



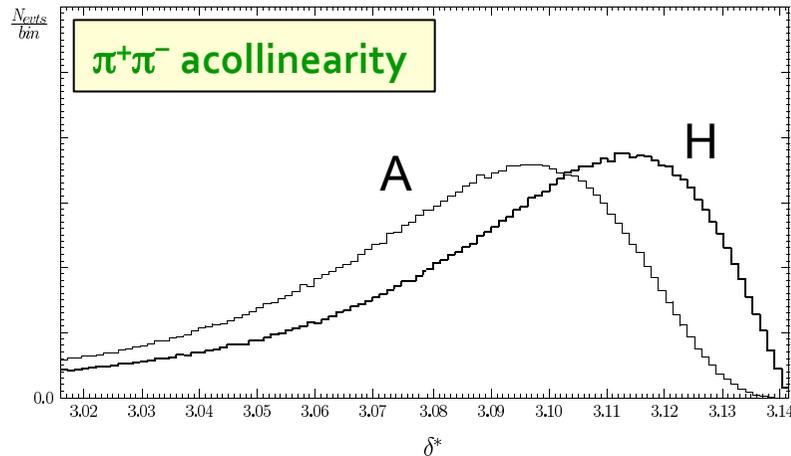
- ◆ Can “see” H and A
 - If $\tan\beta > 5$
- ◆ Build the next collider
 - At $\sqrt{s} \sim m_{A,H}$



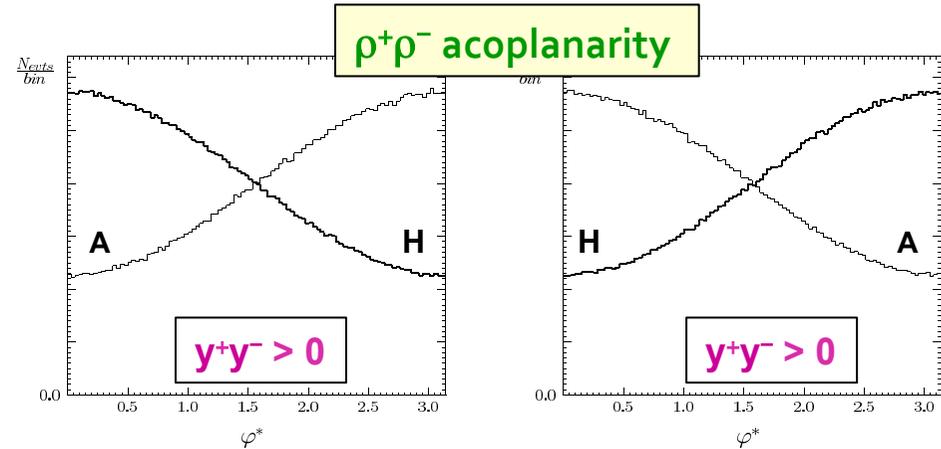
Additional Higgs bosons (4)

Unique CP (violation) and H/A mixing studies can start

From $H, A \rightarrow \tau^+\tau^- \rightarrow \pi^+\pi^-\bar{\nu}_\tau\nu_\tau$

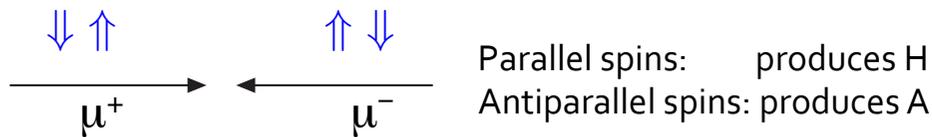


From $H, A \rightarrow \tau^+\tau^- \rightarrow \rho^+\rho^-\bar{\nu}_\tau\nu_\tau$ with $\rho^\pm \rightarrow \pi^\pm\pi^0$

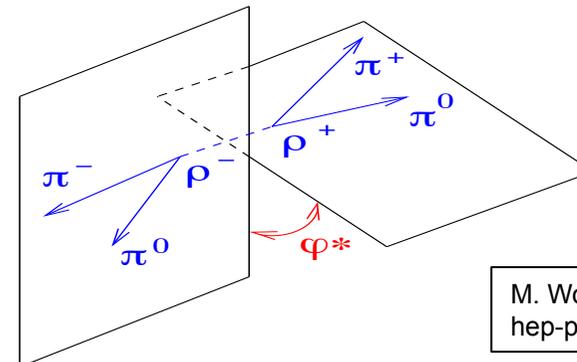


$$y^\pm = E_{\pi^\pm} - E_{\pi^0}$$

From beam transverse polarization



No idea of whether it is feasible or not...



F. Palhen et al.
JHEP 0808:030
JHEP 0801:017

M. Worek
hep-ph/0305082

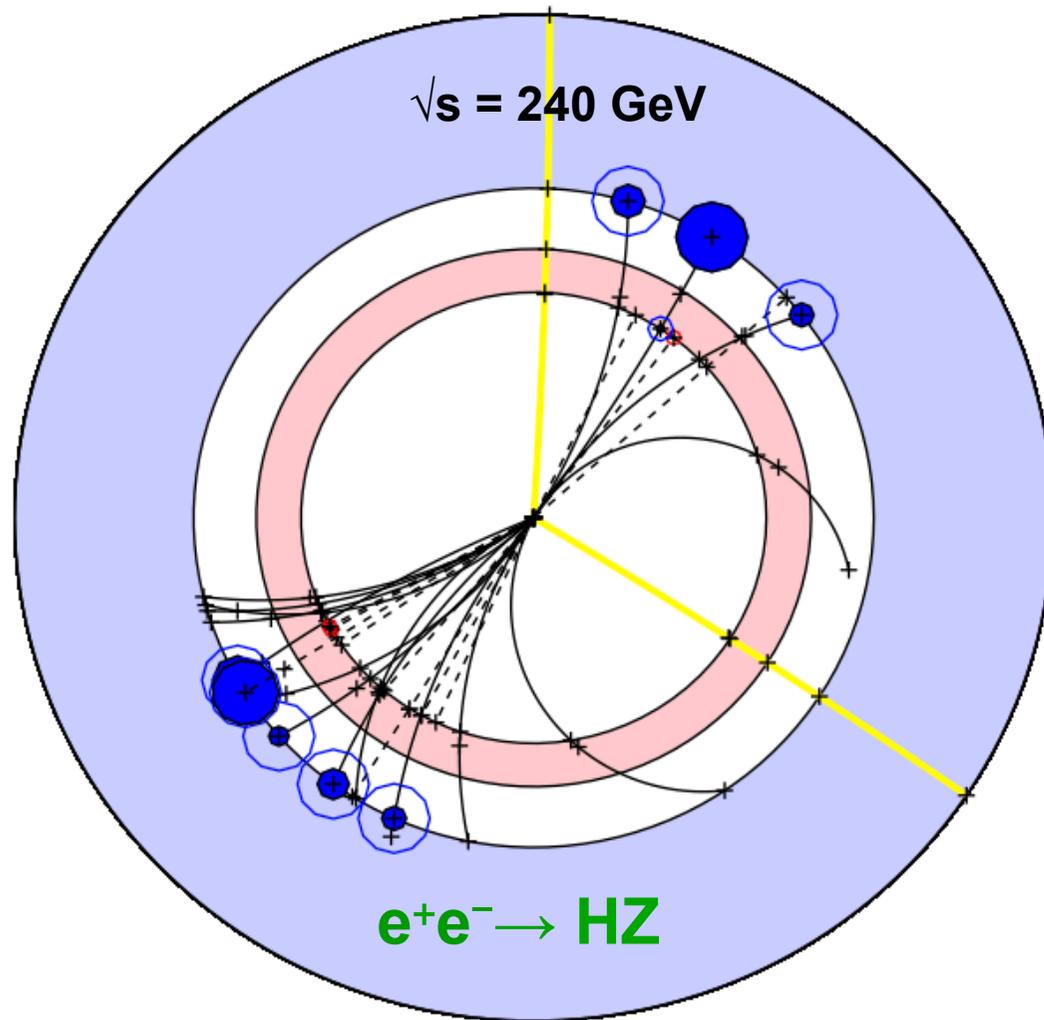
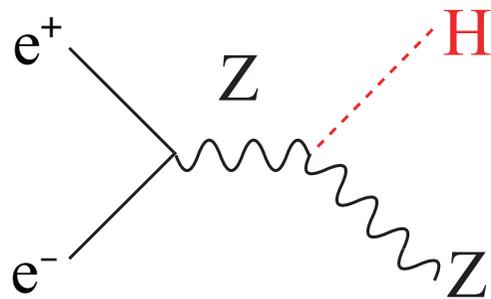
Summary

- **A muon collider at $\sqrt{s} = 125$ GeV is a pretty Higgs factory**
 - ◆ But not necessarily the one we need
 - Only a few Higgs couplings are accessible (b, μ, W, τ) with a 2-5% precision
 - The Higgs total decay width can be measured with a 5% precision
 - If H is a single particle, we will know more from e^+e^- collisions at the FCC-ee
 - All Higgs couplings can be measured with 0.1 – 1% precision
 - Sensitive to the Higgs invisible branching fraction down to 0.1%
 - Important : e^+e^- colliders can measure the Higgs width very well
 - Precision of 0.9% at the FCC-ee (4% at the ILC)
- **Muon colliders may have a case if H(125) is formed by nearby peaks**
 - ◆ Separated by less than few 100 MeV and by more than a few MeV
 - Can be observed via $e^+e^- \rightarrow hA$ at FCC-ee
 - A similar situation occurs for heavy H and A: requires $m_{H,A}$ to be known beforehand
 - Open the possibility of nice (and unique) CP studies
- **Muon colliders may be the best way to reach $\sqrt{s} > 3$ TeV with leptons**
 - ◆ Much R&D remain to be done in cooling and acceleration

Backup: Higgs width at FCC-ee (1)

□ Example of a Higgs boson event

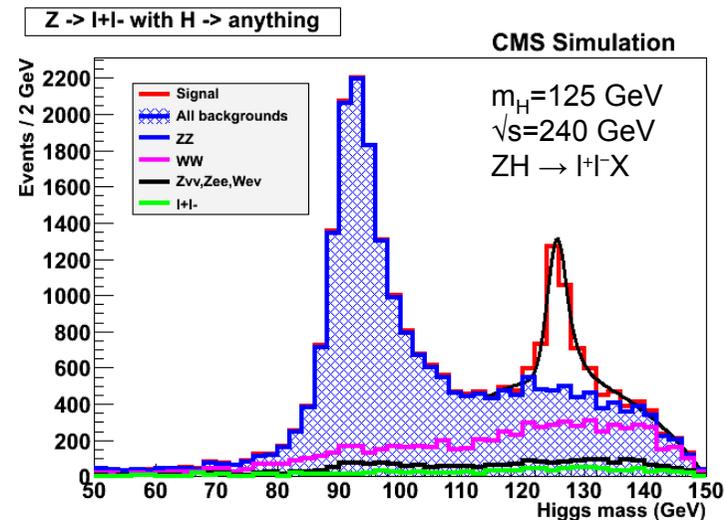
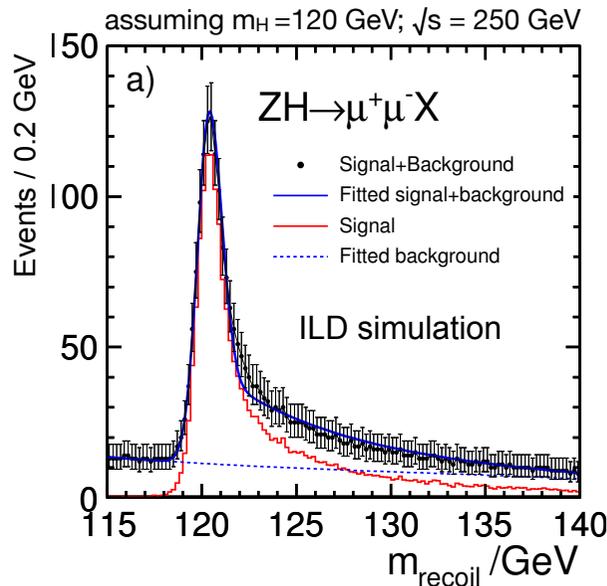
- ◆ Tagged with a Z boson
- ◆ Very clean signature



Backup: Higgs width at FCC-ee (2)

□ Example: Model-independent measurement of σ_{HZ} and κ_Z

- ◆ The Higgs boson in HZ events is tagged by the presence of the $Z \rightarrow e^+e^-, \mu^+\mu^-$
 - Select events with a lepton pair ($e^+e^-, \mu^+\mu^-$) with mass compatible with m_Z
 - No requirement on the Higgs decays: measure $\sigma_{HZ} \times BR(Z \rightarrow e^+e^-, \mu^+\mu^-)$
 - Apply total energy-momentum conservation to determine the “recoil mass”
 - ➔ $m_H^2 = s + m_Z^2 - 2\sqrt{s}(p_+ + p_-)$ Exercise !
 - Plot the recoil mass distribution – resolution proportional to momentum resolution



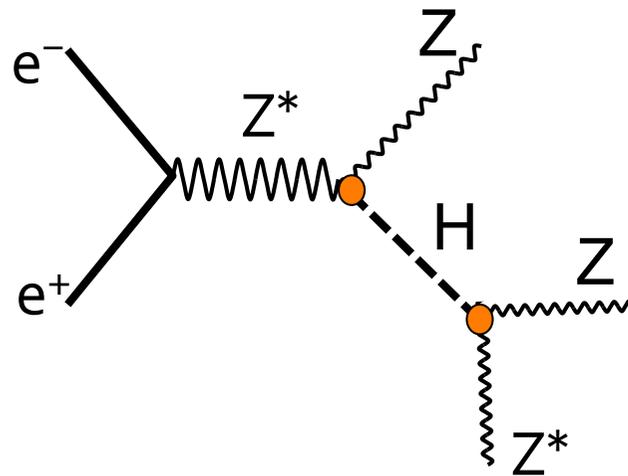
- ◆ Provides an absolute measurement of κ_Z and set required detector performance

Backup: Higgs width at FCC-ee (3)

□ Indirect determination of the total Higgs decay width

◆ From a counting of HZ events with $H \rightarrow ZZ$ at $\sqrt{s} = 240 \text{ GeV}$

● Measure $\sigma_{HZ} \times \text{BR}(H \rightarrow ZZ)$



Final state with three Z's
Almost background free

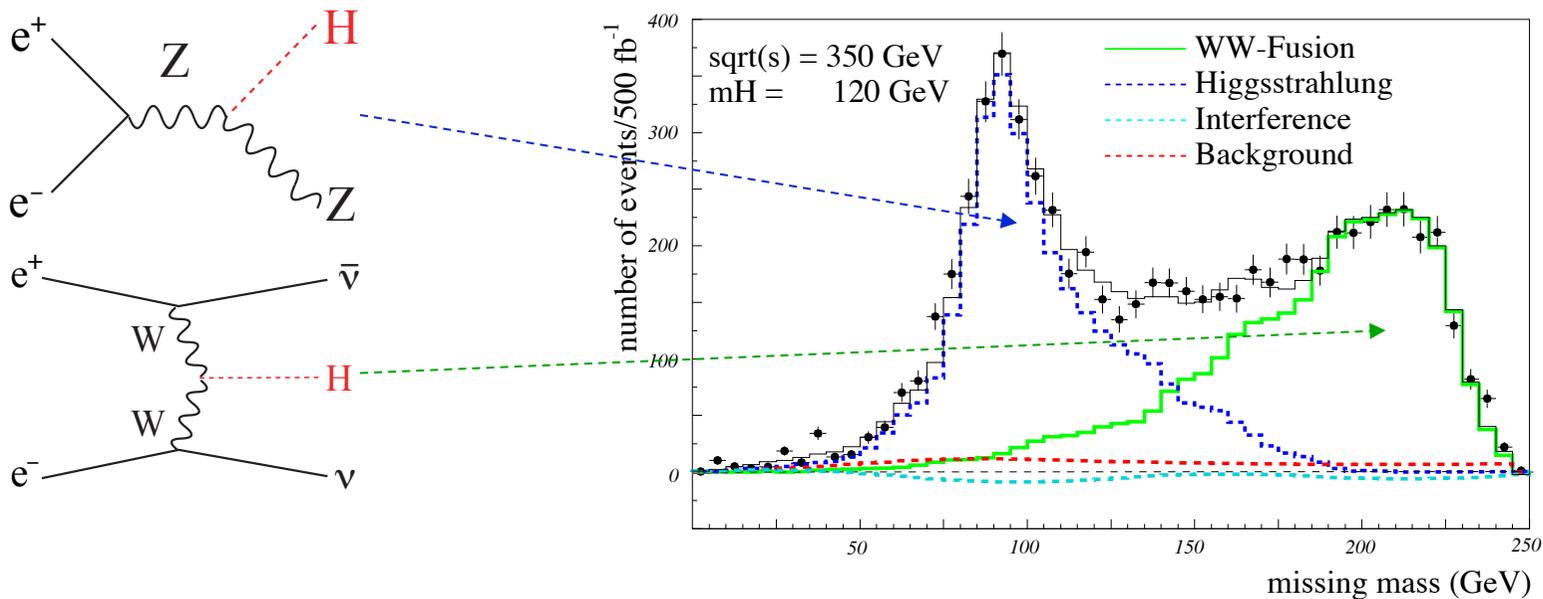
Measured with the Hl^+l^- final state
(see slide 21)

- σ_{HZ} is proportional to κ_Z^2
- $\text{BR}(H \rightarrow ZZ) = \Gamma(H \rightarrow ZZ) / \Gamma_H$ is proportional to κ_Z^2 / Γ_H
 - ➔ $\sigma_{HZ} \times \text{BR}(H \rightarrow ZZ)$ is proportional to κ_Z^4 / Γ_H
- Infer the total width Γ_H

Backup: Higgs width at FCC-ee (4)

□ Indirect determination of the total Higgs decay width (cont'd)

- ◆ From a counting $WW \rightarrow H \rightarrow b\bar{b}$ events at 350-500 GeV in the $b\bar{b}\nu\bar{\nu}$ final state:



- Measure $\sigma(WW \rightarrow H \rightarrow b\bar{b})$
- Take the branching ratios into WW and $b\bar{b}$ from σ_{HZ} and $\sigma_{HZ} \times \text{BR}(H \rightarrow WW, b\bar{b})$
- Infer the total width

$$\Gamma_H \propto \sigma_{WW \rightarrow H} / \text{BR}(H \rightarrow WW) = \sigma_{WW \rightarrow H \rightarrow b\bar{b}} / \text{BR}(H \rightarrow WW) \times \text{BR}(H \rightarrow b\bar{b})$$