



# Benchmarks for 2HDM Searches @ LHC:

The Hierarchical 2HDM & EW Phase Transition

G. Dorsch, S. Huber, K. Mimasu, J. M. No

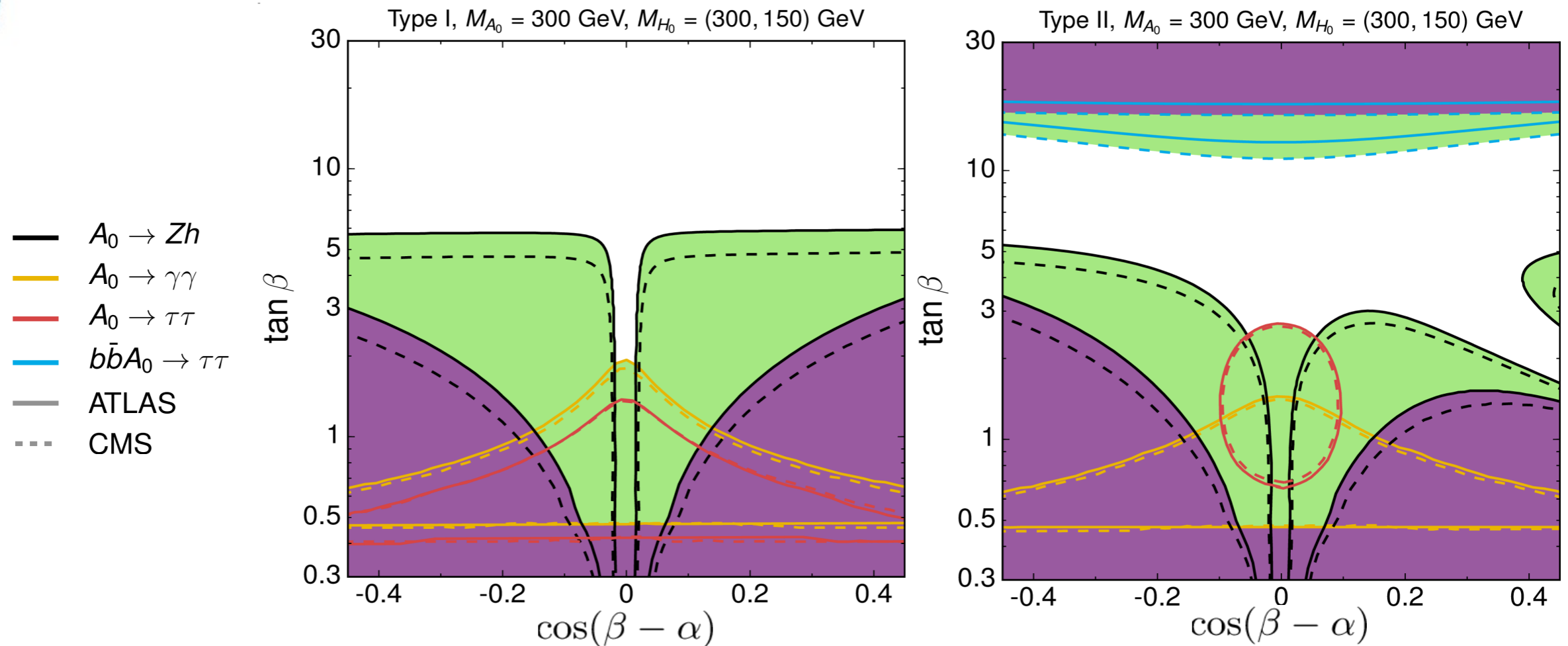
Based on 1405.5537 (PRL), 1507.xxxxx

LHCHSWG3 Benchmarks2HDM

# Mass splittings among $H_0, A_0, H^\pm$ have large impact in 2HDM phenomenology

FOR SIZABLE MASS SPLITTINGS, LIMITS FROM "STANDARD" LHC SEARCHES MAY BE SIGNIFICANTLY WEAKENED

e.g.  $m_{A_0} - m_{H_0} > m_Z$   $\rightarrow$   $A_0 \rightarrow H_0 Z$  channel open  $\rightarrow$  BR to other decay channels drastically reduced



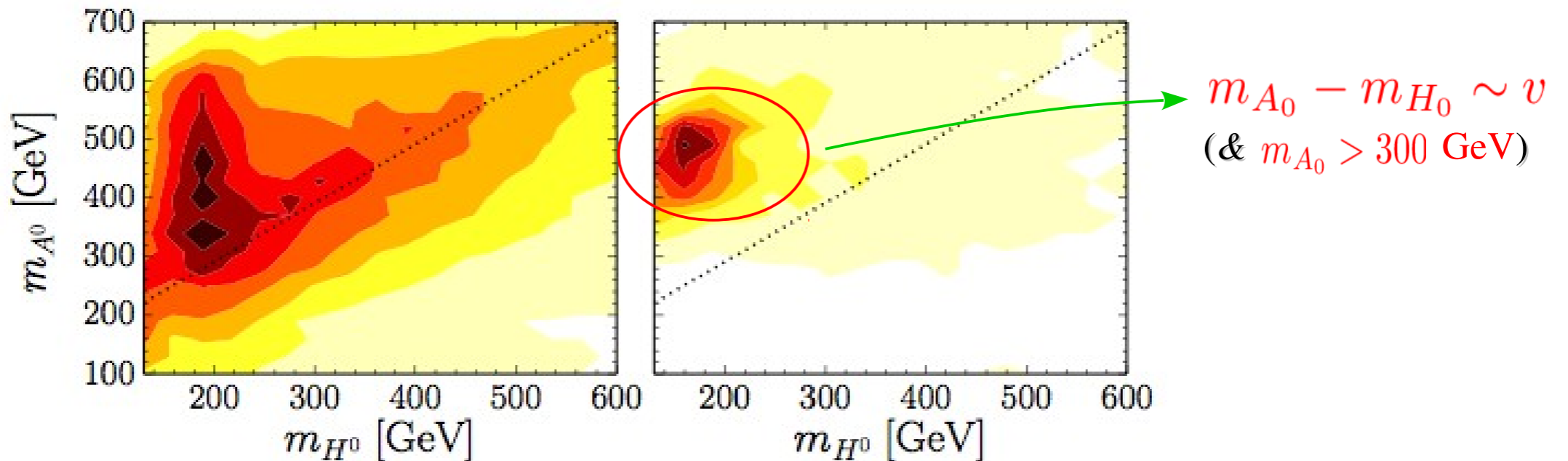
SusHi (NNLO) + ZHDMC

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# Mass splittings among $H_0, A_0$ Important for 2HDM Cosmology

A SIZABLE  $m_{A_0} - m_{H_0}$  FAVOURS A FIRST ORDER EW PHASE TRANSITION

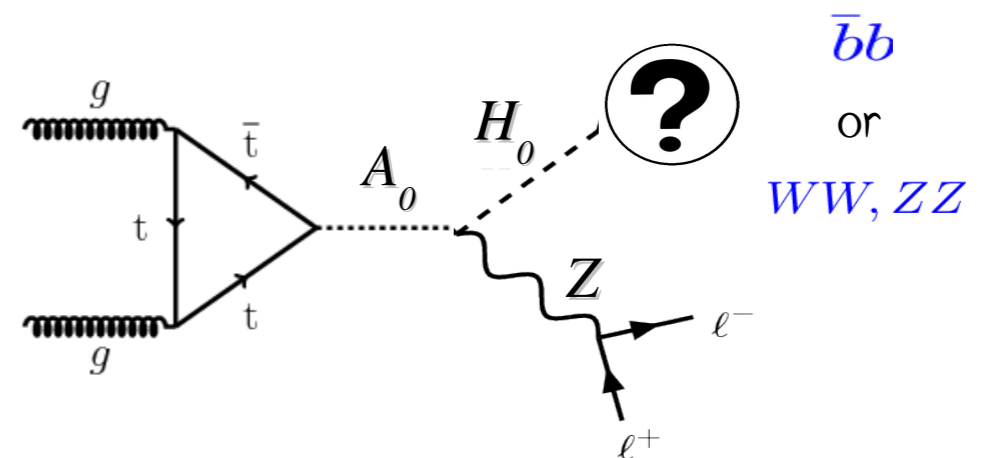
ELECTROWEAK BARYOGENESIS POSSIBLE WITHIN 2HDM



EW PHASE TRANSITION SIGNATURE

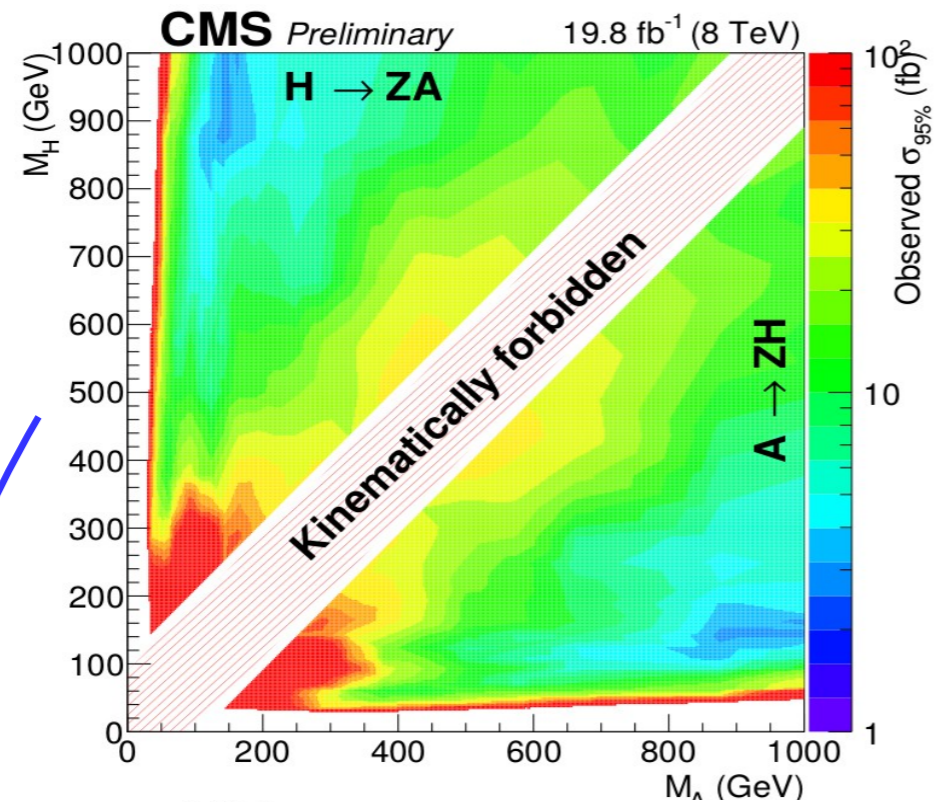
$$A_0 \rightarrow H_0 Z$$

G. Dorsch, S. Huber, K. Mimasu, J. M. N, Phys. Rev. Lett. **113** (2014) 211802



# Such search(es) recently performed by CMS (Meaningful constraints from LHC Run 1)

## CMS-PAS-HIG-15-001

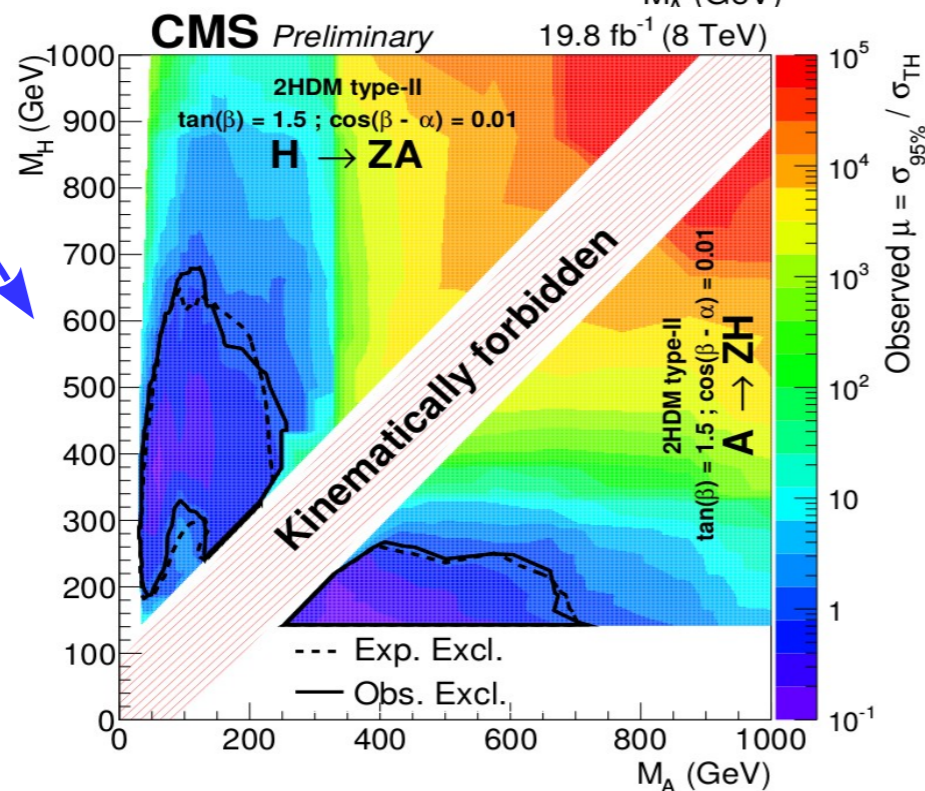


Search for H/A decaying into Z and A/H, with  $Z \rightarrow \ell\ell$  and  $A/H \rightarrow b\bar{b}$  or  $A/H \rightarrow \tau\tau$

The CMS Collaboration

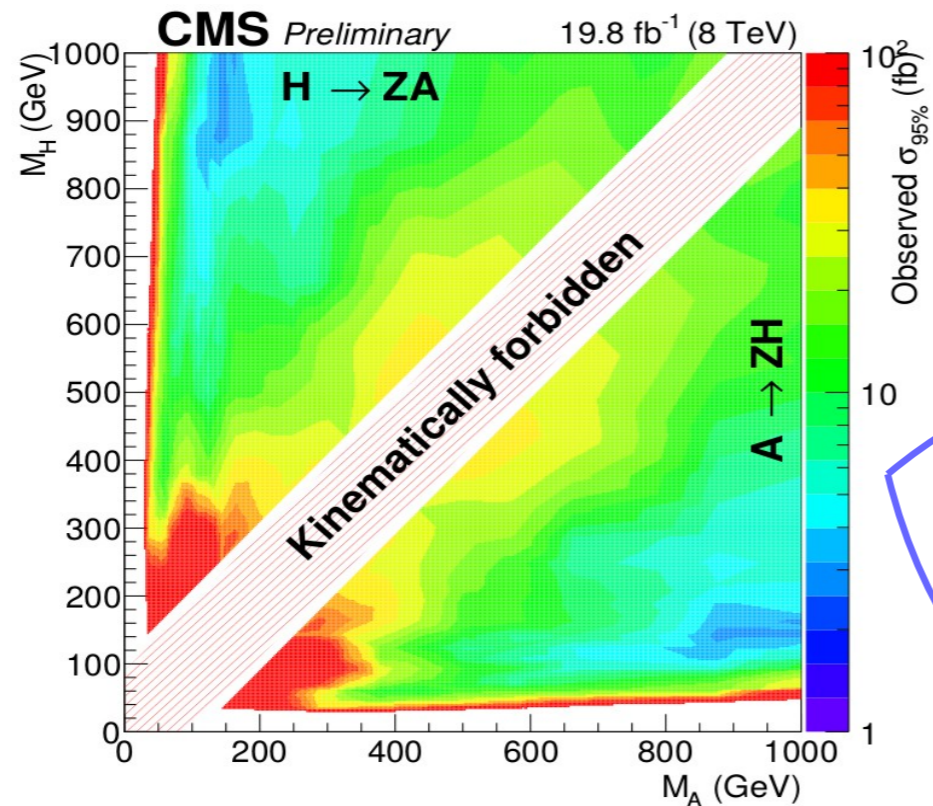
One important motivation for 2HDMs is that these models provide a way to explain the asymmetry between matter and anti-matter observed in the Universe [4, 5]. Another important motivation is Supersymmetry [6], which is a theory that falls in the broad class of 2HDMs. Axion models [7], which would explain how the strong interaction does not violate the CP symmetry, would give rise to an effective low-energy theory with two Higgs doublets. Finally, it has also been recently noted [8] that certain realizations of 2HDMs can accommodate the muon  $g - 2$  anomaly [9] without violating the present theoretical and experimental constraints.

In the most general case 14 parameters are necessary to describe the scalar sector in a 2HDM. However, only 6 free parameters remain once the so-called  $Z_2$  symmetry is imposed to suppress flavor changing neutral currents, in agreement with experimental observations, and the values of the mass of the recently discovered Higgs boson (125 GeV) and the electroweak vacuum expectation value (246 GeV) are assumed. The compatibility of a 125 GeV SM-like Higgs boson with 2HDMs is possible in the so-called alignment limit. In such a limit, one of the CP-even scalars,  $h$  or  $H$ , is identified with the 125 GeV Higgs boson and the condition  $\cos(\beta - \alpha) \approx 0$  or  $\sin(\beta - \alpha) \approx 0$  is satisfied, where  $\tan \beta$  and  $\alpha$  are, respectively, the ratio of the vacuum expectation values, and the mixing angle of the two Higgs doublets. A recent theoretical study [5] has shown that, in this limit, a large mass splitting ( $>100$  GeV) between the A and H bosons would favor the electroweak phase transition that would be at the origin of the baryogenesis process in the early Universe, thus explaining the currently observed matter-antimatter asymmetry in the Universe. In such a scenario, the most frequent decay mode of the pseudoscalar A boson would be  $A \rightarrow ZH$ .



# Such search(es) recently performed by CMS (Meaningful constraints from LHC Run 1)

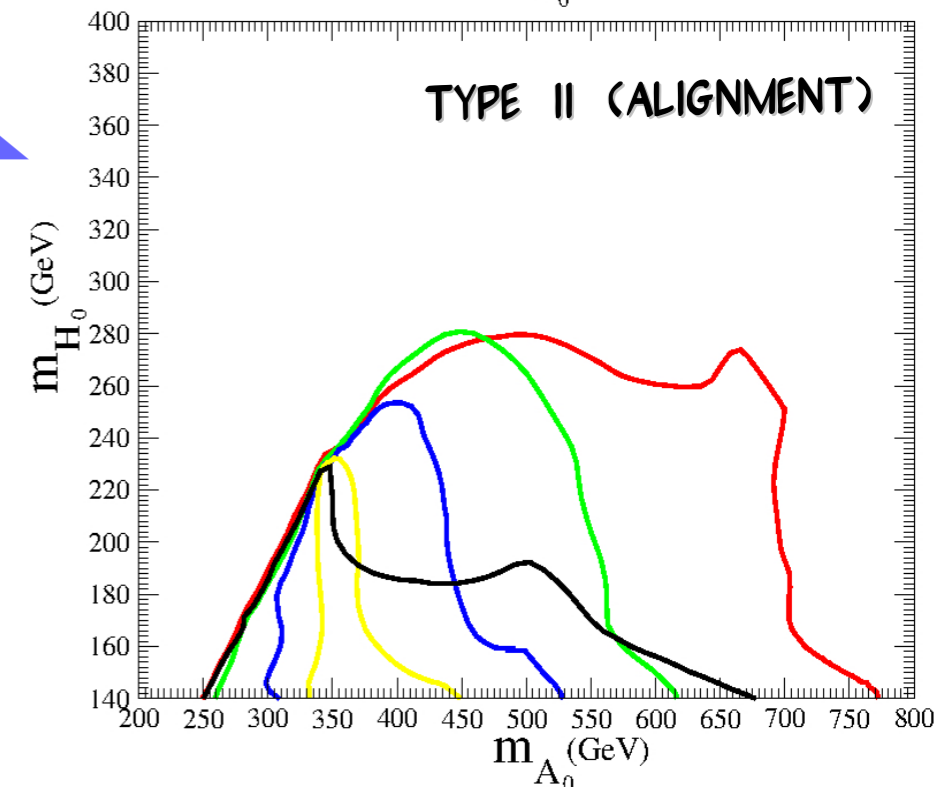
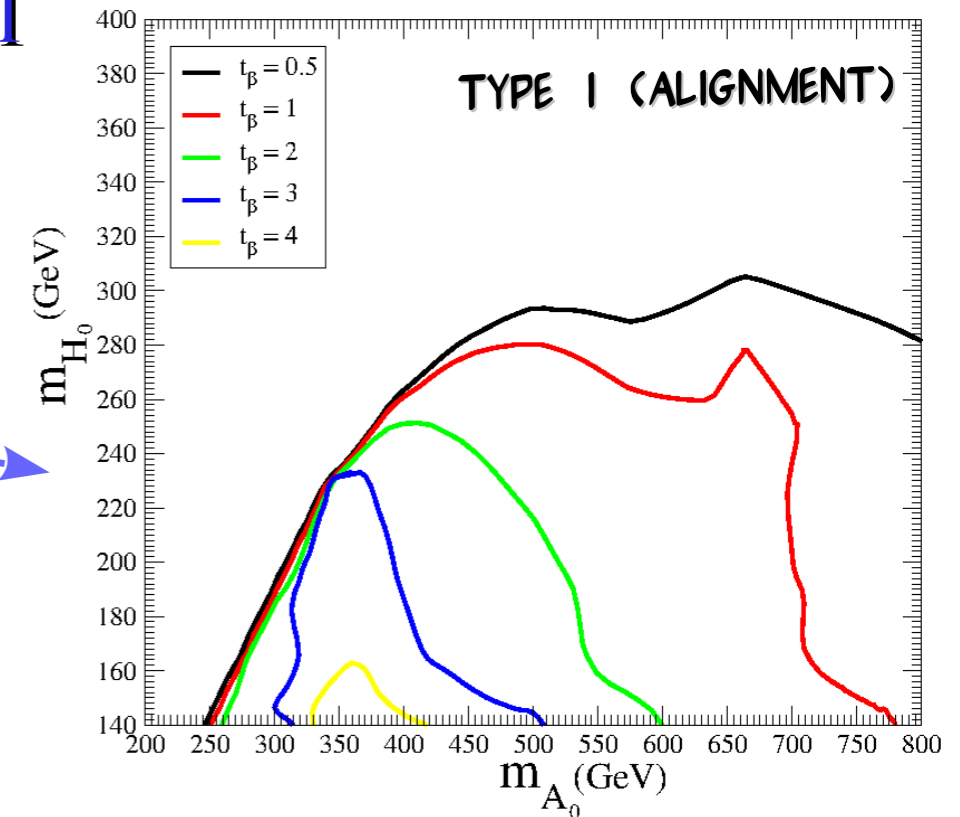
## CMS-PAS-HIG-15-001



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**Sushi (NNLO) + 2HDMC**

- This search constrains/“rules out”:
- A1, A2 from GHMN
  - BP1, BP2, BP3, BP4, BP8, BP10 from KS
  - Scenario D from HS



# Benchmarks for 2HDM from the Electroweak Phase Transition

- ALL BENCHMARKS FEATURE  $m_{A_0} - m_{H_0} > m_Z$

$A_0 \rightarrow H_0 Z$  as primary signature

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- ALL BENCHMARKS FEATURE  $m_{A_0} - m_{H_0} > m_Z$

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- BENCHMARKS **A**  *Close to Alignment Limit*  $\cos(\beta - \alpha) = 0$

$H_0 \rightarrow \bar{b}b$  ( $m_{H_0} < 340$  GeV),  $H_0 \rightarrow \bar{t}t$  ( $m_{H_0} > 340$  GeV)

$H^\pm$  needs to be close in mass to either  $H_0$  or  $A_0$ , due to EW Precision Observables

**A1** :  $m_{A_0} = m_{H^\pm} = 420$  GeV  $m_{H_0} = 180$  GeV  $\tan(\beta) = 3$

TYPE I  $\sigma_{gg}^{A_0} = 2.795$  pb  $Br(A_0 \rightarrow Z H_0) = 0.8433$   $Br(H_0 \rightarrow \bar{b}b) = 0.711$

TYPE II  $\sigma_{gg}^{A_0} = 2.837$  pb  $Br(A_0 \rightarrow Z H_0) = 0.8383$   $Br(H_0 \rightarrow \bar{b}b) = 0.899$

**A2** :  $m_{A_0} = 420$  GeV  $m_{H_0} = m_{H^\pm} = 180$  GeV  $\tan(\beta) = 3$

TYPE I  $\sigma_{gg}^{A_0} = 2.795$  pb  $Br(A_0 \rightarrow Z H_0) = 0.3016$   $Br(H_0 \rightarrow \bar{b}b) = 0.711$

$Br(A_0 \rightarrow W^\pm H^\mp) = 0.642$   $Br(H^\pm \rightarrow \bar{b}t) = 0.992$

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**A1**:  $m_{A_0} = m_{H^\pm} = 420$  GeV  $m_{H_0} = 180$  GeV  $\tan(\beta) = 3$

Similar to BP1 from KS. Regarding topology and final state (**bb**), similar to BP2, BP3, BP4, BP8 from KS, and Scenario D from HS

**A2**:  $m_{A_0} = 420$  GeV  $m_{H_0} = m_{H^\pm} = 180$  GeV  $\tan(\beta) = 3$

Interesting signatures from  $H^\pm$  in both Cases, similar to those Shufang will discuss (BP1-BP8).



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$A_0 \rightarrow H_0 Z$  as primary signature

- BENCHMARKS **B**  *Away from Alignment Limit*  $\cos(\beta - \alpha) \gtrsim 0.1$

$$H_0 \rightarrow W^+ W^- \begin{cases} 4\ell + 2\nu & \text{Final State} \\ 3\ell + \nu + 2j & \text{Final State} \end{cases}$$

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$$H_0 \rightarrow ZZ \begin{cases} 4\ell + 2\nu & \text{Final State} \\ 4\ell + 2j & \text{Final State} \end{cases}$$

B. Coleppa, F. Kling, S. Su, 1404.1922

Such departure from Alignment possible for Type I (Type II in “wrong sign” )

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**B1 :**  $m_{A_0} = m_{H^\pm} = 400 \text{ GeV}$   $m_{H_0} = 180 \text{ GeV}$   $\tan(\beta) = 2$   $\sin(\beta - \alpha) = 0.95$

$\sigma_{gg}^{A_0} = 8.158 \text{ pb}$   $Br(A_0 \rightarrow Z H_0) = 0.584$   $Br(H_0 \rightarrow W^+ W^-) = 0.936$   $Br(H_0 \rightarrow ZZ) = 0.062$

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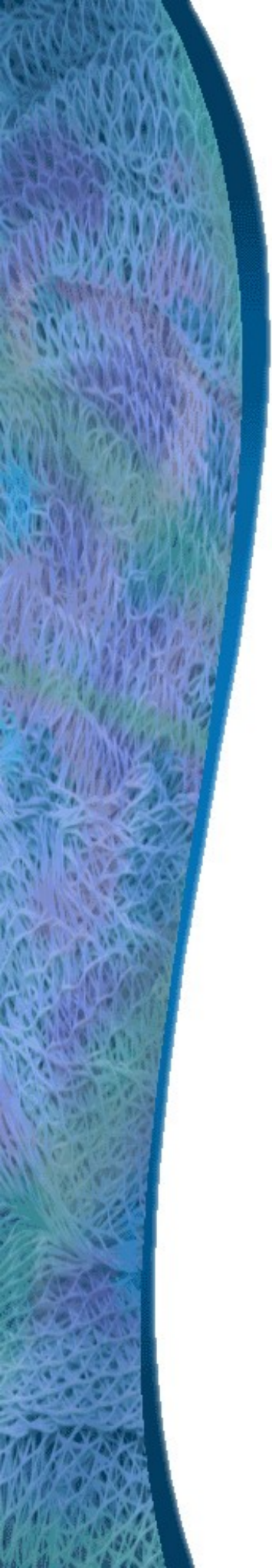
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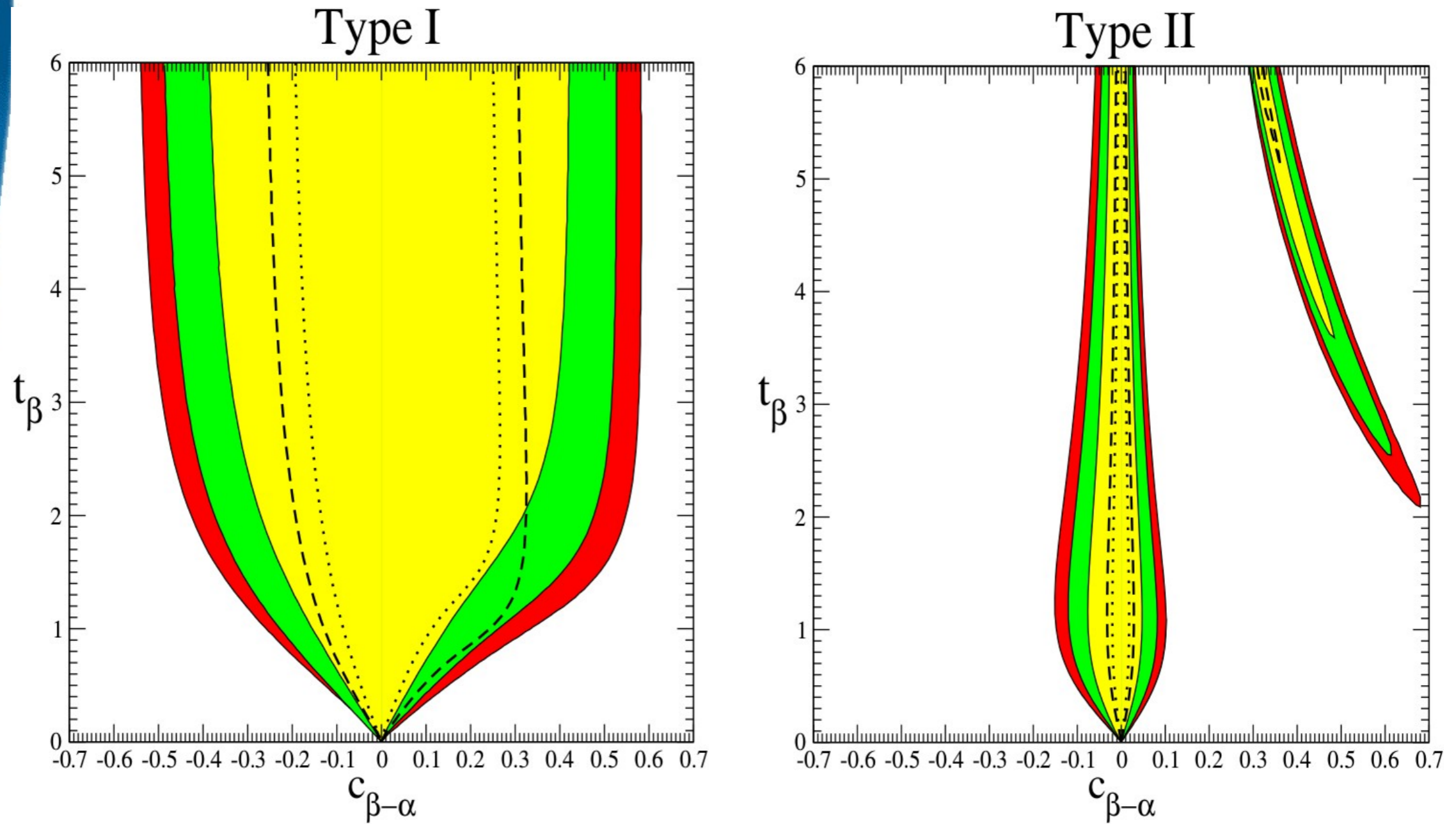
$$m_{A_0} = m_{H^\pm} = 400 \text{ GeV} \quad m_{H_0} = 220 \text{ GeV} \quad \tan(\beta) = 2 \quad \sin(\beta - \alpha) = 0.95$$

$$\sigma_{gg}^{A_0} = 8.158 \text{ pb} \quad Br(A_0 \rightarrow Z H_0) = 0.446 \quad Br(H_0 \rightarrow W^+ W^-) = 0.715 \quad Br(H_0 \rightarrow ZZ) = 0.285$$



# Back-up

# Higgs Signal Strengths Fit



# 2HDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

### ① *A few words on the Analysis...*

- ⇒ Type I 2HDM implemented in FeynRules (including gluon-fusion).
- ⇒ Signal & relevant backgrounds generated using MadGraph5\_aMC@NLO. Generated events passed on to Pythia for Parton Showering and Hadronization and subsequently to Delphes for detector simulation.
  - Use of NLO flat K-factors for signal (SusHi) and dominant backgrounds.
- ⇒ “Cut & Count” analysis on a small set of kinematical variables, to extract signal over background.
- ⇒ Determined required Integrated Luminosity at 14 TeV to achieve  $5\sigma$  statistical significance via a C.L.s hypothesis test.
  - *Only statistical uncertainties.*
  - *10% systematic uncertainty on background.*

# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

③ **Benchmark B:**  $A_0 \rightarrow H_0 Z \rightarrow W^+W^- \ell\ell \rightarrow 4\ell + 2\nu$  ( $\alpha-\beta = 0.1\pi$ )

⇒ Most sensitive  $A_0$  search channel away from alignment

⇒  $A_0 \rightarrow H_0 Z \rightarrow ZZ\ell\ell \rightarrow 4\ell + 2j$  also promising

*B. Coleppa, F. Kling, S. Su, JHEP 1409 (2014) 161*

⇒ Main backgrounds are  $ZZ$ ,  $Zt\bar{t}$ ,  $hZ$ ,  $ZWW$  subdominant

⇒ Analysis & Event Selection similar to previous case:

→ 4 Isolated (cone of 0.3) leptons, same-flavour pairs.  $|\eta| < 2.5$  (2.7) for electrons (muons)

→  $P_T^{\ell 1} > 40 \text{ GeV}$ ,  $P_T^{\ell 2,3,4} > 20 \text{ GeV}$ .

→ 1 pair of SF leptons must reconstruct  $m_Z$

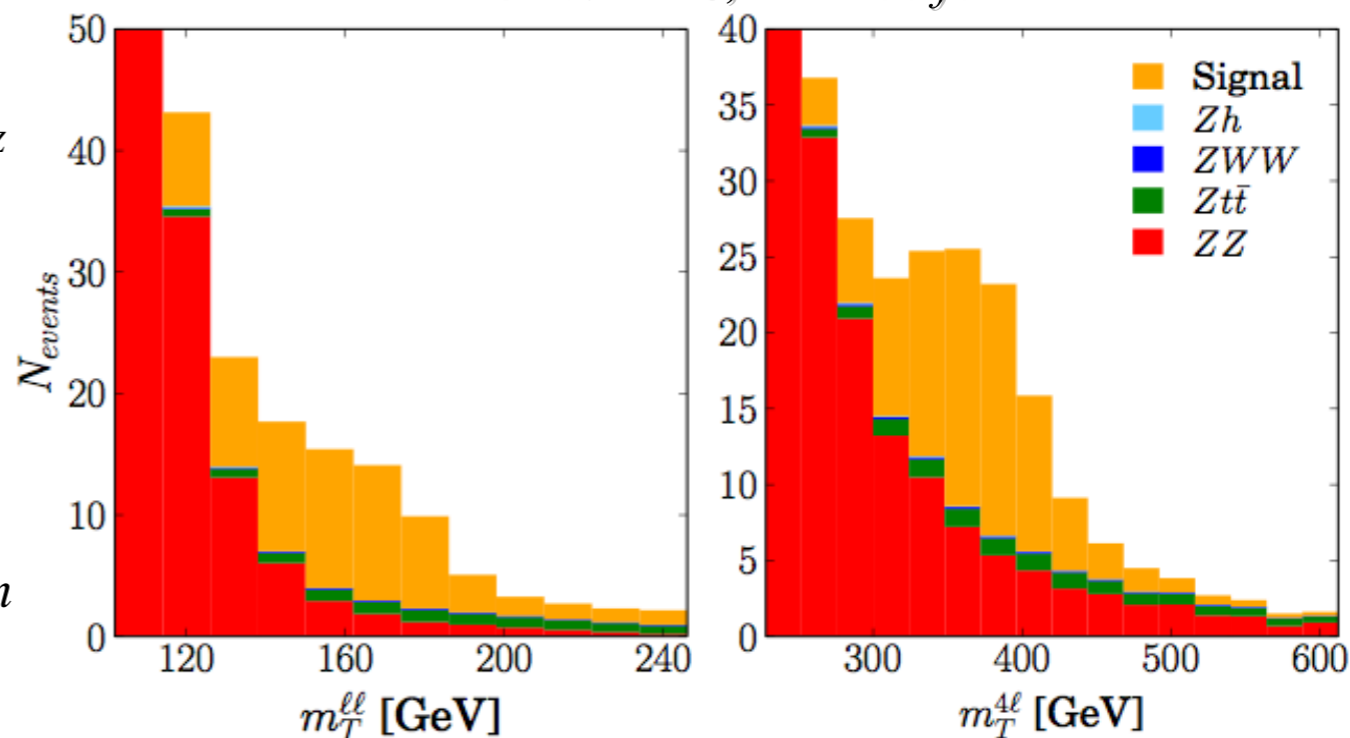
→ Transverse mass variables:

$$(m_T^{\ell\ell})^2 = (\sqrt{p_{T,\ell\ell}^2 + m_{\ell\ell}^2} + \cancel{p}_T)^2 - (\vec{p}_{T,\ell\ell} + \vec{\cancel{p}}_T)^2$$

$$m_T^{4\ell} = \sqrt{p_{T,\ell'\ell'}^2 + m_{\ell'\ell'}^2} + \sqrt{p_{T,\ell\ell}^2 + (m_T^{\ell\ell})^2}$$

$m_T^{4\ell} > 260 \text{ GeV}$  allows for Signal Extraction

14 TeV LHC,  $\mathcal{L} = 60 \text{ fb}^{-1}$



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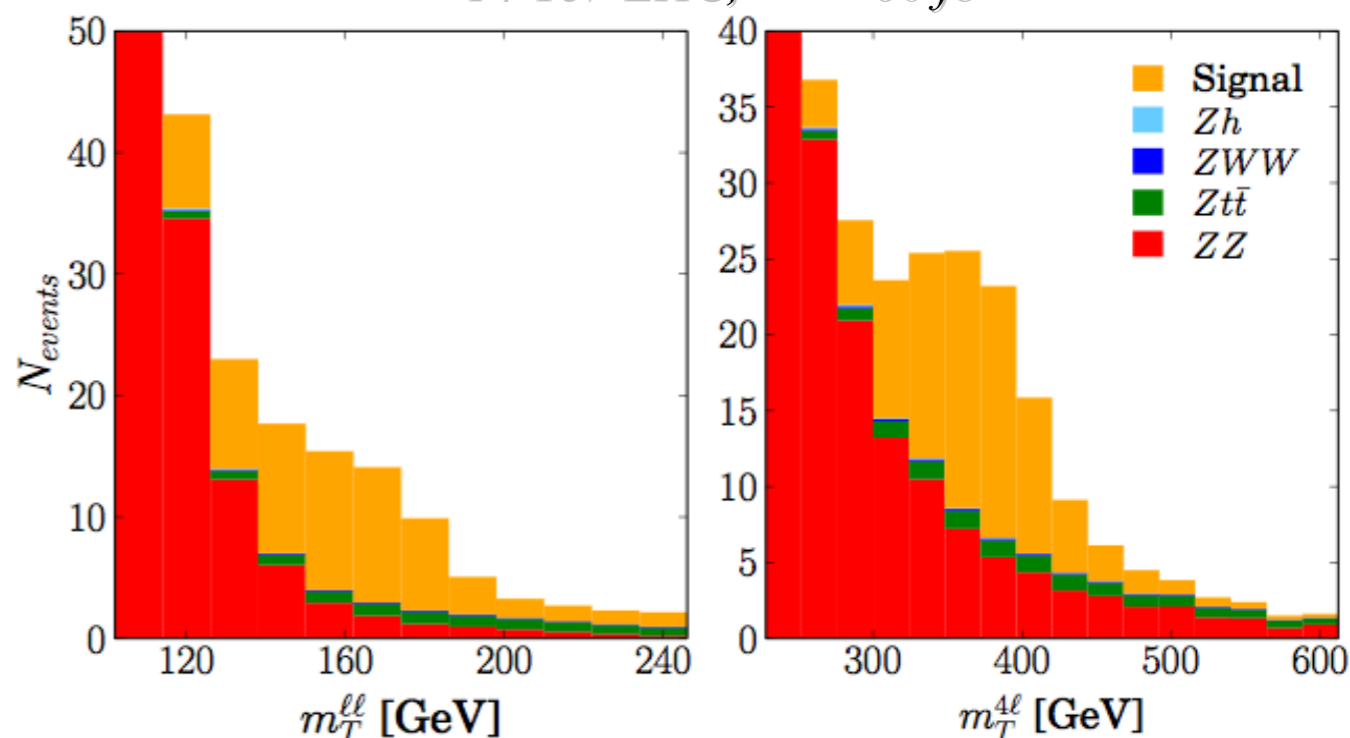
5 $\sigma$  signal significance for:

$\mathcal{L} = 60 \text{ fb}^{-1}$  (statistics only)

$\mathcal{L} = 200 \text{ fb}^{-1}$  (10% systematics)

(conservative...)

14 TeV LHC,  $\mathcal{L} = 60 \text{ fb}^{-1}$





# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

② *Benchmark A:*  $A_0 \rightarrow H_0 Z \rightarrow \bar{b}b \ell\ell$  ( $\alpha-\beta = 0.001\pi$ )

⇒ Irreducible backgrounds are  $Z\bar{b}b$ ,  $t\bar{t}$ ,  $ZZ$ ,  $hZ$

⇒ Analysis at 14 TeV: Event Selection

→ Anti- $k_T$  Jets with distance parameter  $R = 0.6$

→ 2  $b$ -tagged Jets with  $|\eta| < 2.5$

→ 2 Isolated (within a cone of 0.3), Same-flavour leptons.  $|\eta| < 2.5$  (2.7) for electrons (muons)

→  $P_T^{\ell_1} > 40 \text{ GeV}$ ,  $P_T^{\ell_2} > 20 \text{ GeV}$ .

	K-factor:	1.6	1.5	1.4	-	-
	Signal	$t\bar{t}$	$Z\bar{b}b$	$ZZ$	$Zh$	
Event selection	14.6	1578	424	7.3	2.7	
$80 < m_{\ell\ell} < 100 \text{ GeV}$	13.1	240	388	6.6	2.5	
$H_T^{\bar{b}b} > 150 \text{ GeV}$	8.2	57	83	0.8	0.74	
$H_T^{\ell\ell\bar{b}b} > 280 \text{ GeV}$	5.3	5.4	28.3	0.75	0.68	
$\Delta R_{bb} < 2.5, \Delta R_{\ell\ell} < 1.6$	5.3	5.4	28.3	0.75	0.68	
$m_{bb}, m_{\ell\ell\bar{b}b}$ signal region	3.2	1.37	3.2	< 0.01	< 0.02	

# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

② *Benchmark A*:  $A_0 \rightarrow H_0 Z \rightarrow \bar{b}b \ell\ell$  ( $\alpha - \beta = 0.001\pi$ )

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→  $P_T^{\ell_1} > 40$  GeV,  $P_T^{\ell_2} > 20$  GeV.

14 TeV LHC,  $\mathcal{L} = 20 \text{ fb}^{-1}$

*Invariant mass windows:*

$$m_{\bar{b}b} \rightarrow (m_{H_0} - 20) \pm 30 \text{ GeV}$$

$$m_{\ell\ell\bar{b}b} \rightarrow (m_{A_0} - 20) \pm 40 \text{ GeV}$$

$5\sigma$  signal significance for:

$$\mathcal{L} \doteq 15 \text{ fb}^{-1} \quad (\text{statistics only})$$

$$\mathcal{L} = 40 \text{ fb}^{-1} \quad (10\% \text{ systematics})$$

