

Heavy quark dynamics in heavy-ion collisions

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UC Berkeley, USA
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DAAD

DFG Deutsche
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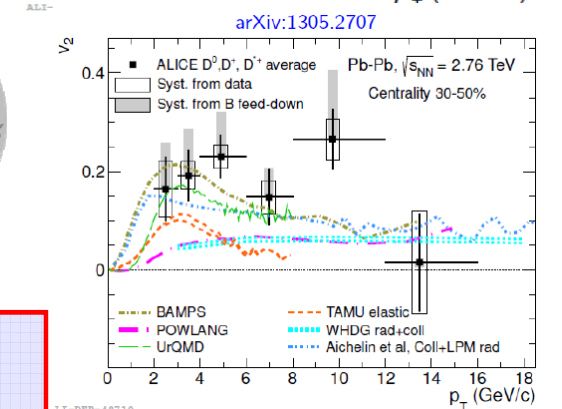
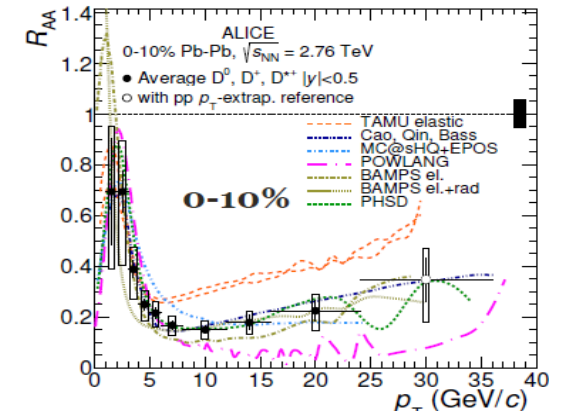
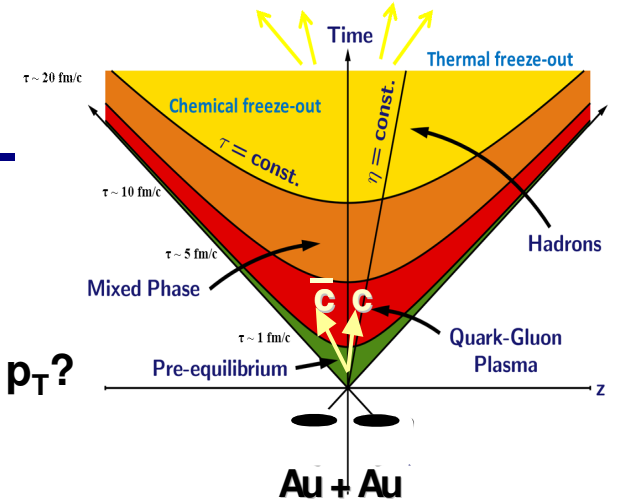
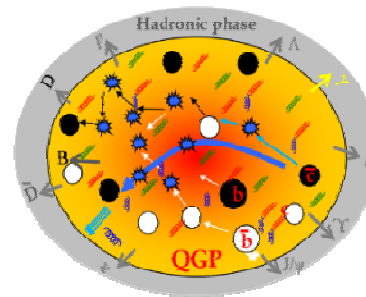


Motivation

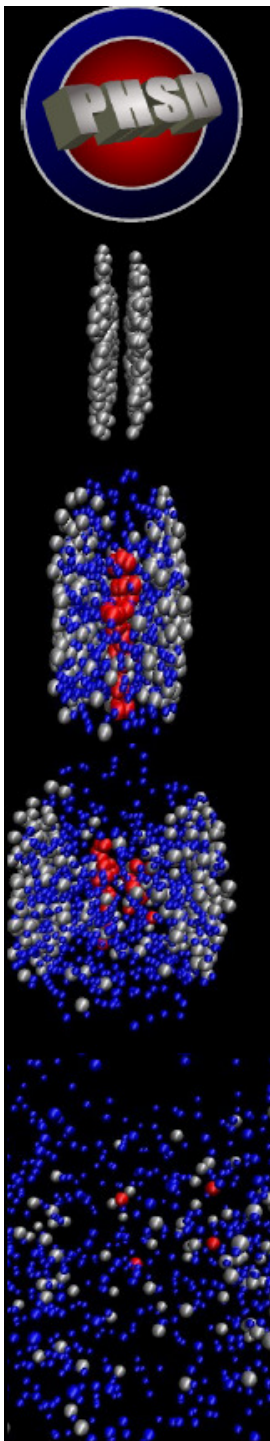
- Hope to use 'heavy quark probes' for a tomography of the early stage of the QGP
- What is the origin for the "energy loss" of charm at large p_T ?

Dynamics of heavy quarks in A+A :

- Production of heavy (charm and bottom) quarks in initial binary collisions + shadowing and Cronin effects
- Interactions in the QGP:
 - elastic scattering $Q+q \rightarrow Q+q$ \rightarrow collisional energy loss
 - gluon bremsstrahlung $Q+q \rightarrow Q+q+g$ \rightarrow radiative energy loss
- Hadronization: c/cbar quarks \rightarrow D(D*)-mesons: coalescence vs fragmentation
- Hadronic interactions: D+baryons; D+mesons



The goal: to describe the dynamics of charm quarks/mesons in all phases on a microscopic basis
The tool: PHSD approach



Parton-Hadron-String-Dynamics (PHSD)

PHSD is a **non-equilibrium transport approach** with

- explicit **phase transition** from hadronic to partonic degrees of freedom
- **IQCD EoS** for the partonic phase (‘crossover’ at low μ_q)
- explicit **parton-parton interactions** - between quarks and gluons
- dynamical **hadronization**

□ **QGP phase** is described by the **Dynamical QuasiParticle Model (DQPM)** matched to reproduce lattice QCD

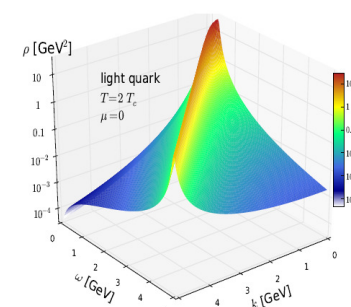
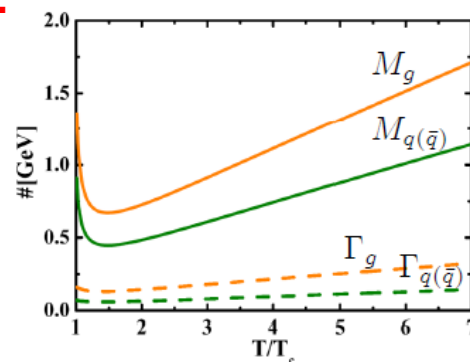
A. Peshier, W. Cassing, PRL 94 (2005) 172301;
W. Cassing, NPA 791 (2007) 365; NPA 793 (2007)

- **strongly interacting quasi-particles:** massive quarks and gluons (g, q, q_{bar}) with sizeable collisional widths in a self-generated **mean-field potential**

- **Spectral functions:**

$$\rho_i(\omega, T) = \frac{4\omega\Gamma_i(T)}{\left(\omega^2 - \vec{p}^2 - M_i^2(T)\right)^2 + 4\omega^2\Gamma_i^2(T)}$$

$(i = q, \bar{q}, g)$



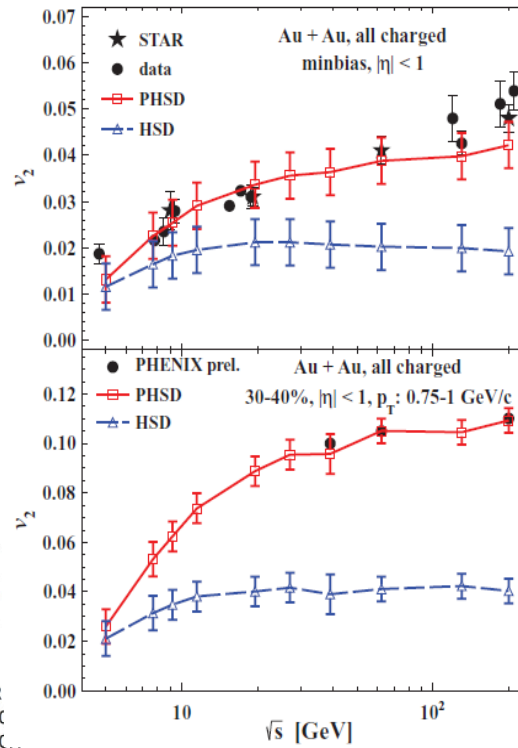
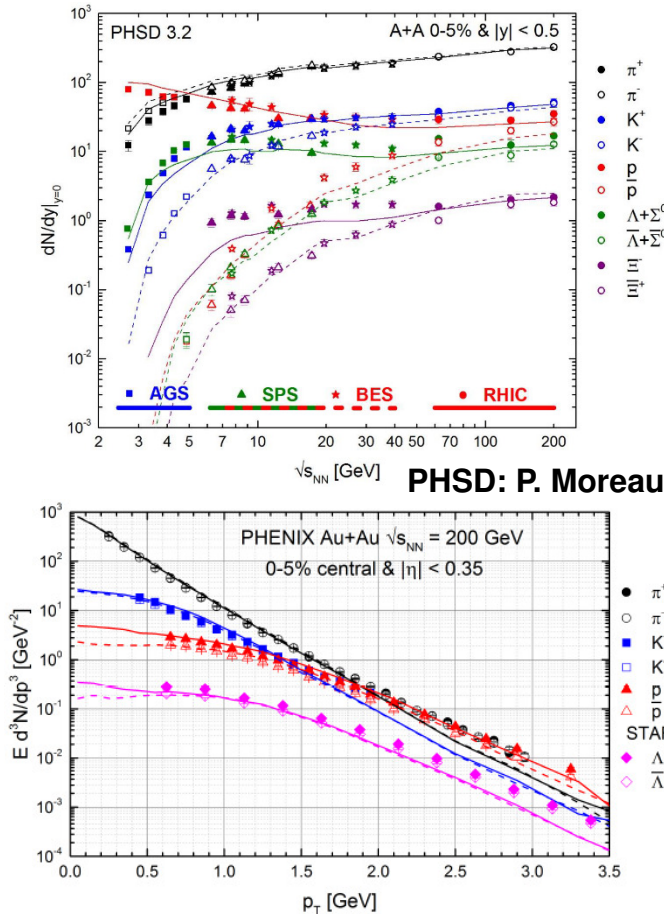
□ **Transport theory:** **generalized off-shell transport equations** based on the 1st order gradient expansion of Kadanoff-Baym equations (**applicable for strongly interacting systems!**)

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215; W. Cassing, EPJ ST 168 (2009) 3

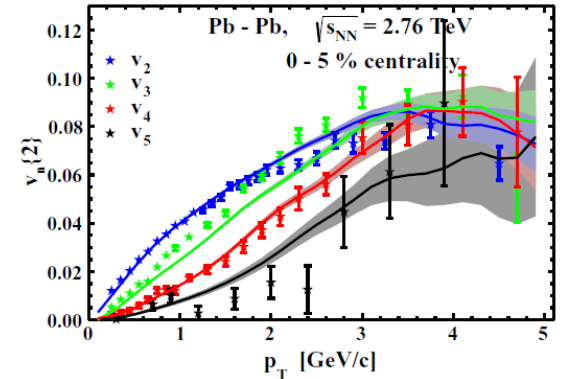
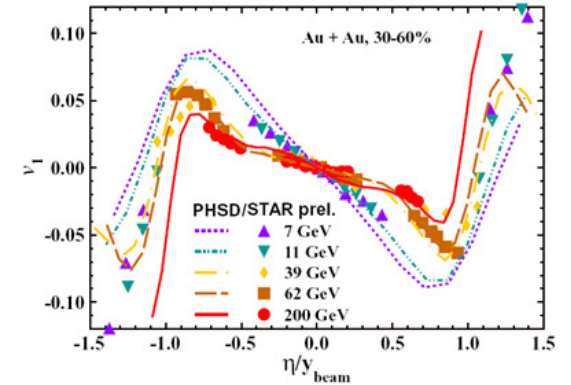


Non-equilibrium dynamics: description of A+A with PHSD

Important: to be conclusive on charm observables, the **light quark dynamics** must be well under control!



V. Konchakovski et al.,
PRC 85 (2012) 011902; JPG42 (2015) 055106

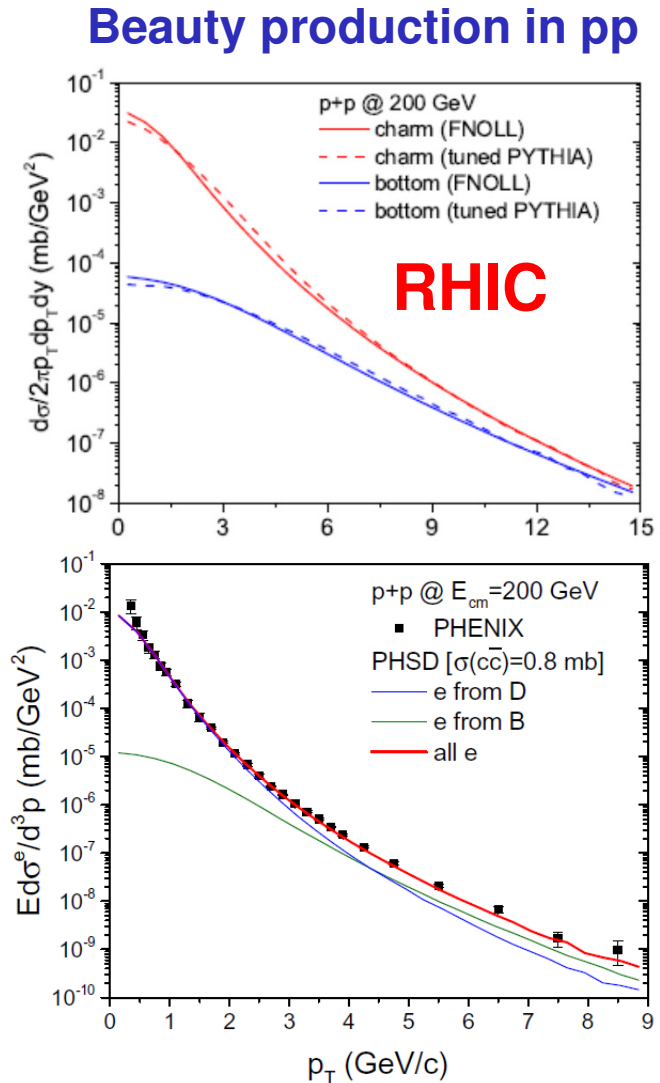
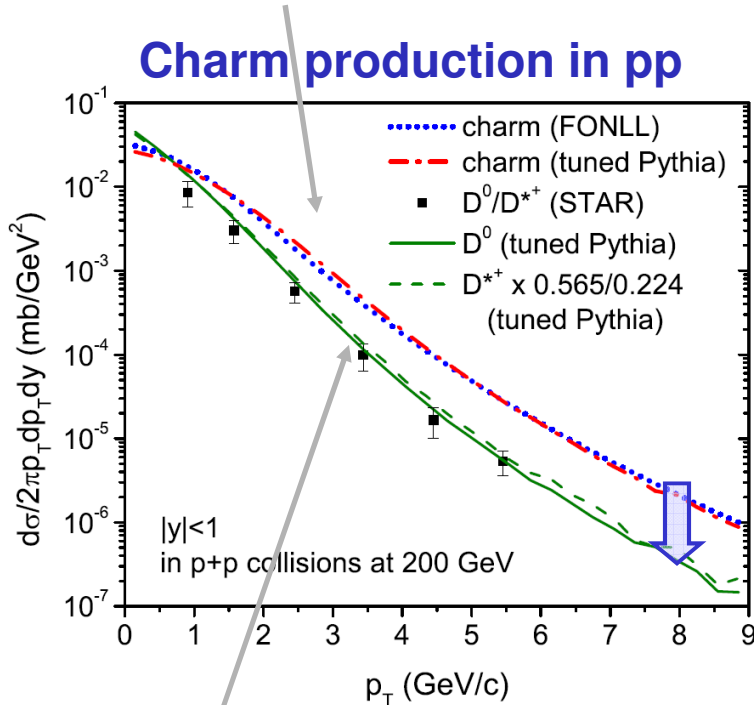


PHSD provides a **good description of 'bulk' observables** (y -, p_T -distributions, flow coefficients v_n , ...) from SPS to LHC



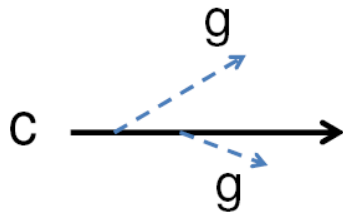
Heavy quark/hadron production in p+p collisions

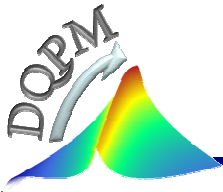
1) **Momentum distribution of heavy quarks:** use 'tuned' PYTHIA event generator to reproduce FONLL (fixed-order next-to-leading log) results (R. Vogt et al.)



2) **Charm/beauty hadron production in pp by heavy-quark fragmentation:**

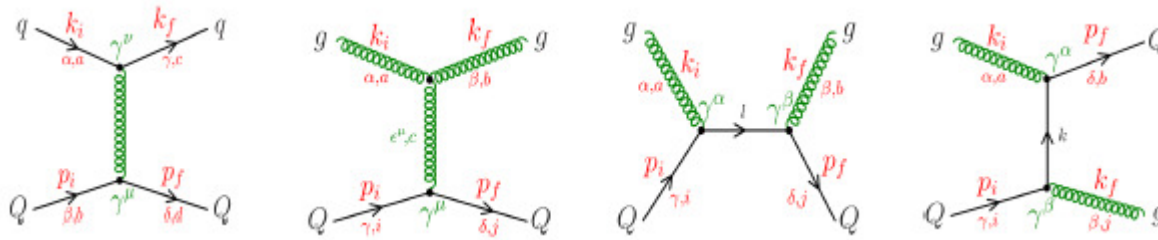
- D⁰ 20 %
- D⁺ 17.4 %
- D^{*0} 21.3 %
- D^{*+} 22.4 %
- DS⁺ 8 %
- Λ_c 9.4 %



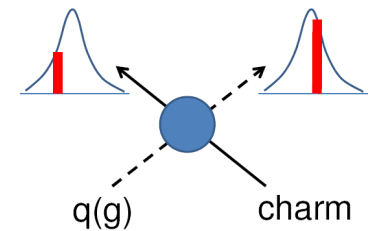


Heavy quark scattering in the QGP (DQPM)

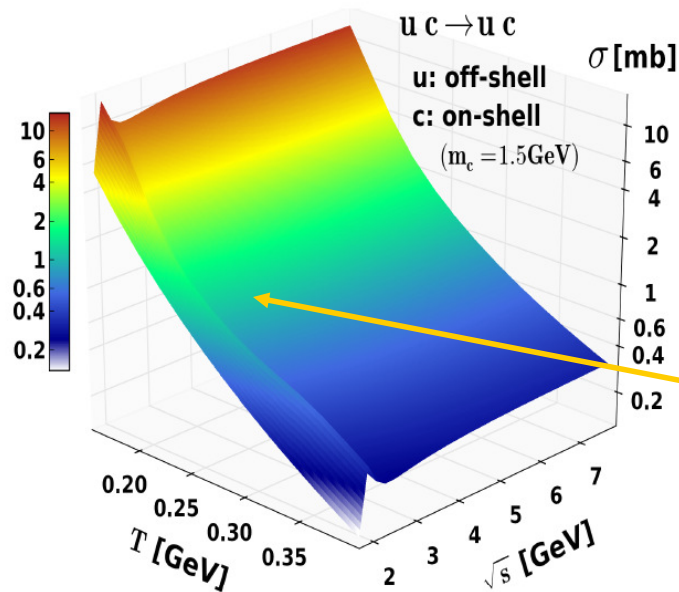
- Elastic scattering with off-shell massive partons $Q+q(g) \rightarrow Q+q(g)$



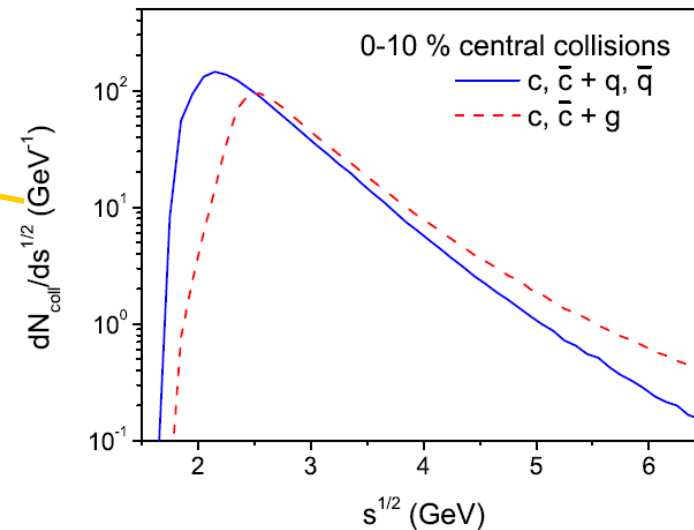
Non-perturbative QGP!

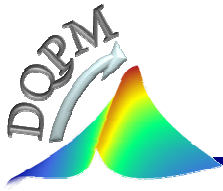


- Elastic cross section $uc \rightarrow uc$



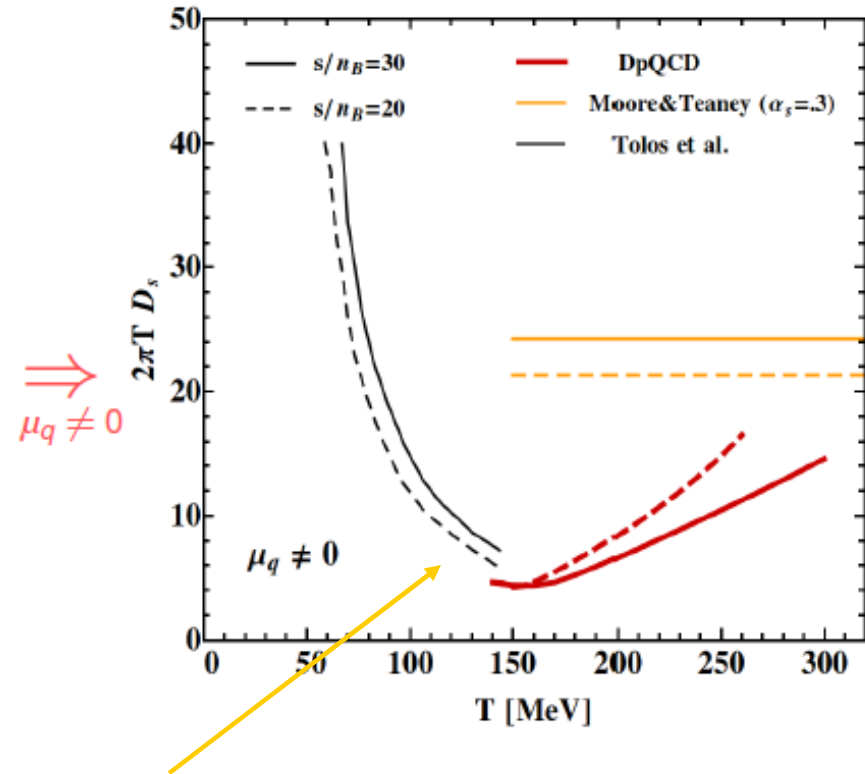
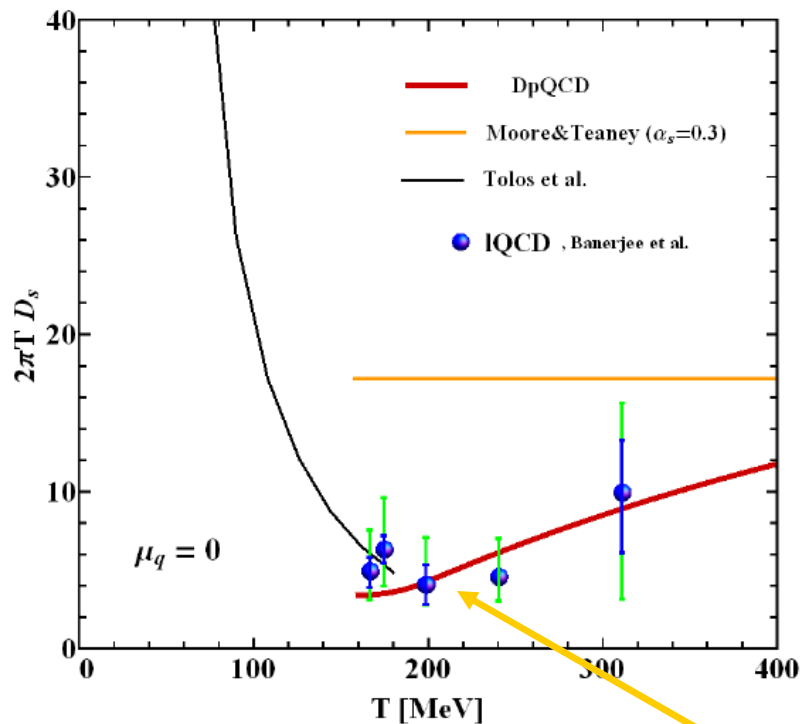
- Distributions of $Q+q$, $Q+g$ collisions vs $s^{1/2}$ in Au+Au, 10% central





Charm spatial diffusion coefficient D_s in the hot medium

- D_s for heavy quarks as a function of T for $\mu_q=0$ and finite μ_q assuming adiabatic trajectories (constant entropy per net baryon s/n_B) for the expansion

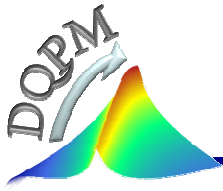


□ $T < T_c$: hadronic D_s

→ Continuous transition at T_c !

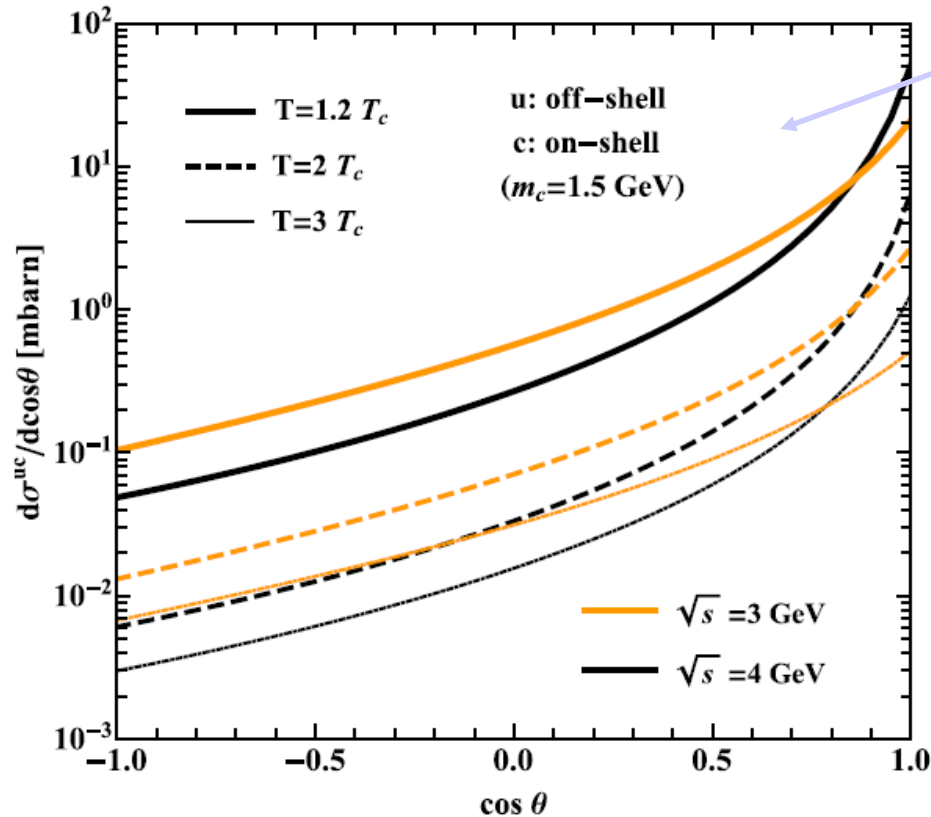
L. Tolos, J. M. Torres-Rincon, PRD 88 (2013) 074019
 V. Ozvenchuk et al., PRC90 (2014) 054909

H. Berrehrah et al, PRC 90 (2014) 051901, arXiv:1406.5322



Heavy quark scattering in the QGP

- **Differential elastic cross section** for $uc \rightarrow uc$ for $s^{1/2} = 3$ and 4 GeV at $1.2T_c$, $2T_c$ and $3T_c$



- **DQPM - anisotropic angular distribution**

Note: pQCD - strongly forward peaked
 → Differences between DQPM and pQCD :
 less forward peaked angular distribution
 leads to **more efficient momentum transfer**

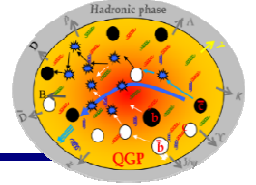
- Central Au+Au, 200 A GeV:
 $N(cc) \sim 19$ pairs,
 $N(Q+q) \sim 130$, $N(Q+g) \sim 85$ collisions
 → each charm quark makes
 ~ **6 elastic collisions in the QGP**

→ Smaller number (compared to pQCD)
 of elastic scatterings with **massive**
 partons leads to a **larger energy loss**

! Note: radiative energy loss is NOT included yet in PHSD,
 it is expected to be **small** due to the large gluon mass in the DQPM



Hadronization of heavy quarks in A+A

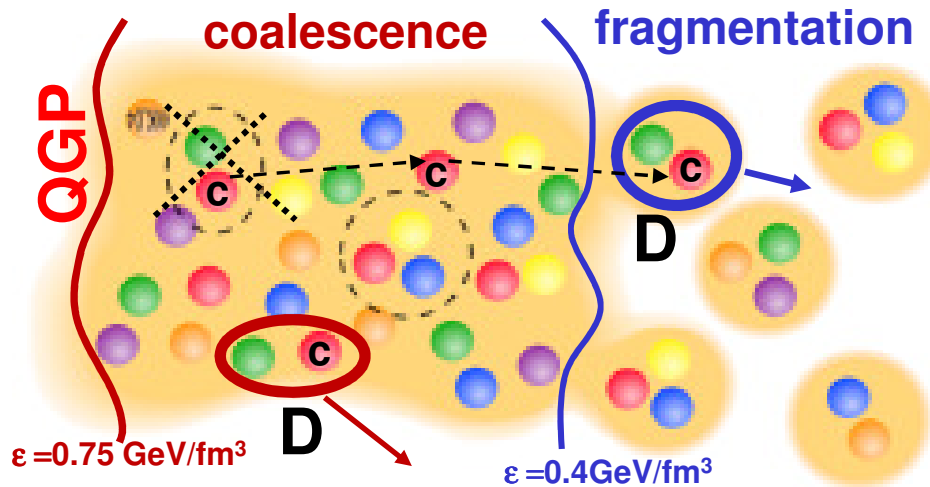


□ PHSD: if the local energy density $\epsilon \rightarrow \epsilon_c \rightarrow$ hadronization of heavy quarks to hadrons

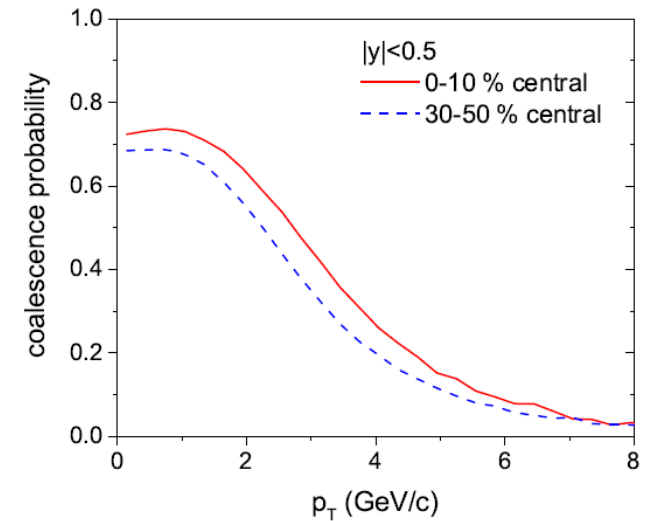
T. Song et al., PRC 93 (2016) 034906

Dynamical hadronization scenario for heavy quarks :

coalescence with $\langle r \rangle = 0.9$ fm & fragmentation
 $0.4 < \epsilon < 0.75$ GeV/fm³ $\epsilon < 0.4$ GeV/fm³



Coalescence probability in Au+Au at LHC



Coalescence probability for $c + \bar{q} \rightarrow D$

$$f(\rho, \mathbf{k}_\rho) = \frac{8g_M}{6^2} \exp \left[-\frac{\rho^2}{\delta^2} - \mathbf{k}_\rho^2 \delta^2 \right]$$

where $\rho = \frac{1}{\sqrt{2}}(\mathbf{r}_1 - \mathbf{r}_2)$, $\mathbf{k}_\rho = \sqrt{2} \frac{m_2 \mathbf{k}_1 - m_1 \mathbf{k}_2}{m_1 + m_2}$

← Width $\delta \leftarrow$ from root-mean-square radius of meson $\langle r \rangle$:

$$\langle r^2 \rangle = \frac{3}{2} \frac{m_1^2 + m_2^2}{(m_1 + m_2)^2} \delta^2$$

Degeneracy factor : $g_M = 1$ for D, = 3 for $D^* = D_0^*(2400)^0, D_1^*(2420)^0, D_2^*(2460)^{0\pm}$



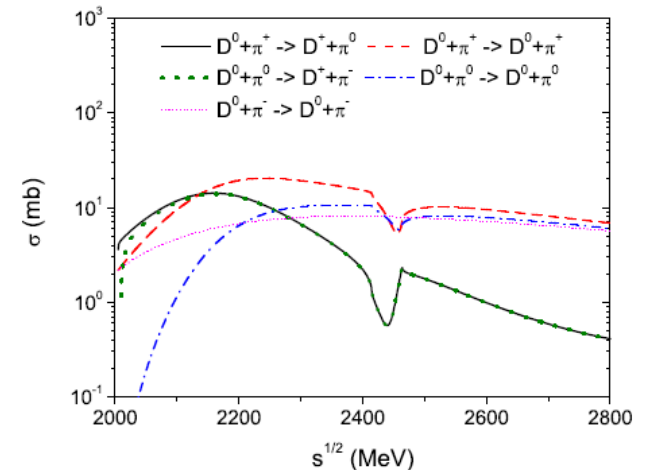
D-meson scattering in the hadronic phase

1. D-meson scattering with mesons

L. M. Abreu, D. Cabrera, F. J. Llanes-Estrada, J. M. Torres-Rincon, *Annals Phys.* 326, 2737 (2011)

Model: effective chiral Lagrangian approach with heavy-quark spin symmetry

Interaction of $D=(D^0, D^+, D^+_s)$ and $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$ with octet ($\pi, K, Kbar, \eta$)



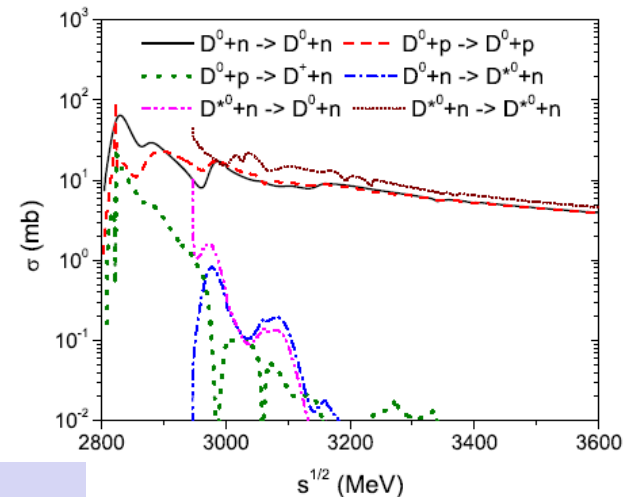
2. D-meson scattering with baryons

C. Garcia-Recio, J. Nieves, O. Romanets, L. L. Salcedo, L. Tolos, *Phys. Rev. D* 87, 074034 (2013)

Model: G-matrix approach: interactions of $D=(D^0, D^+, D^+_s)$ and $D^*=(D^{*0}, D^{*+}, D^{*+}_s)$ with nucleon octet $J^P=1/2^+$ and Delta decuplet $J^P=3/2^+$

Unitarized scattering amplitude \rightarrow solution of coupled-channel **Bethe-Salpeter equations:**

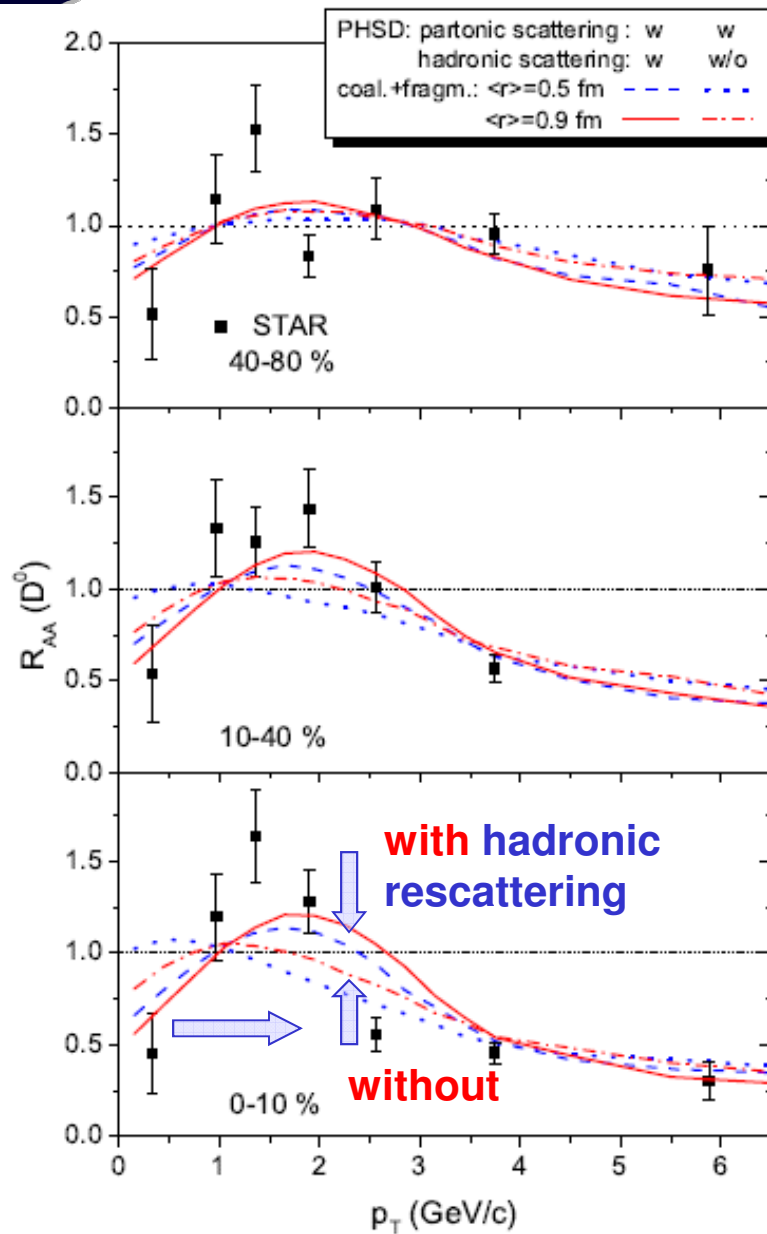
$$T = T + VGT$$



\rightarrow Strong **isospin dependence** and complicated structure (due to the resonance coupling) of $D+m$, $D+B$ cross sections!



R_{AA} at RHIC: hadronic rescattering



Influence of hadronic rescattering:

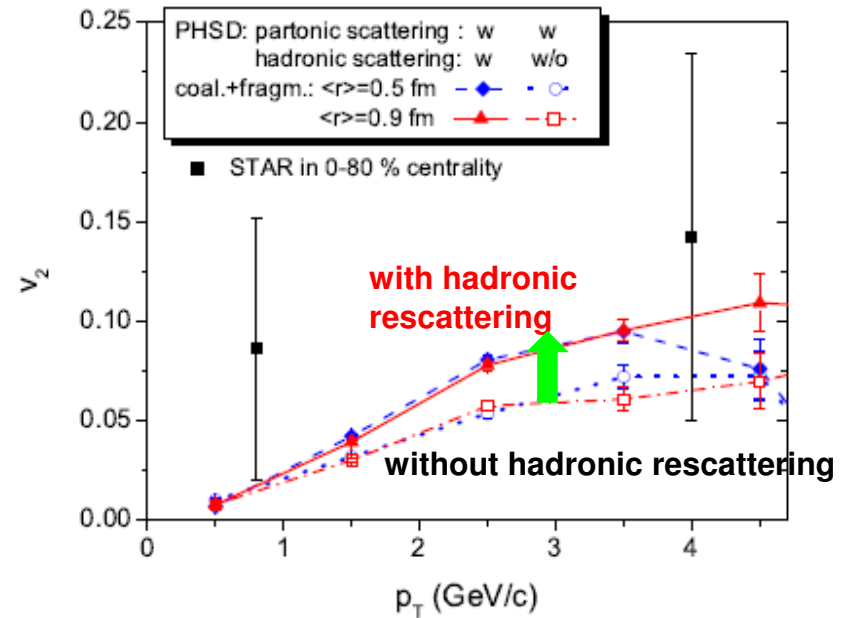
Central Au+Au at $s^{1/2} = 200$ GeV :

$N(D, D^*) \sim 30$

$N(D, D^* + m) \sim 56$ collisions

$N(D, D^* + B, Bbar) \sim 10$ collisions

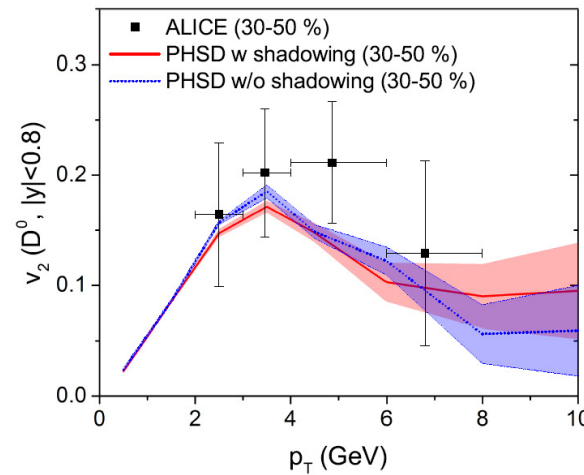
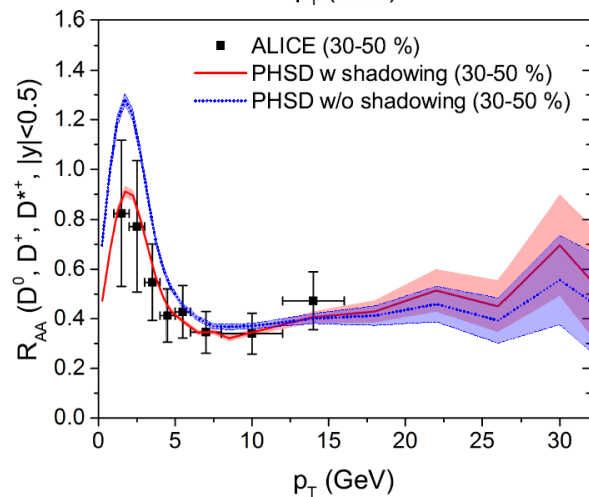
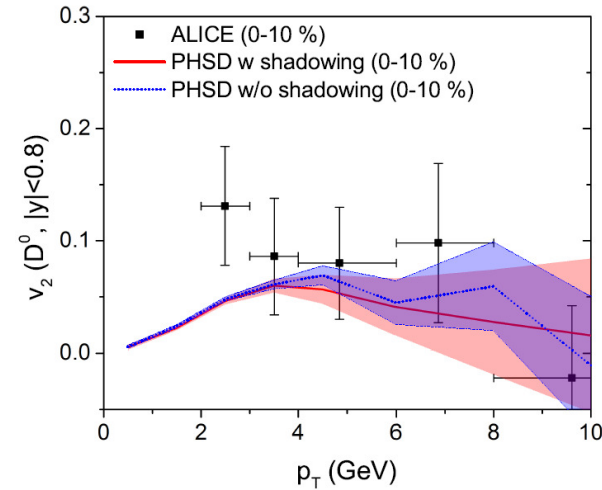
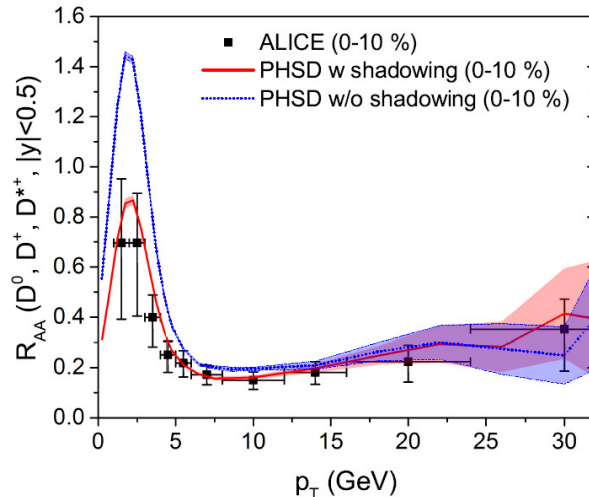
→ each D, D^* makes ~ 2 scatterings with hadrons



- Hadronic rescattering moves R_{AA} peak to higher p_T !
- substantially increases v_2 at larger p_T



Charm R_{AA} at LHC: PHSD vs ALICE

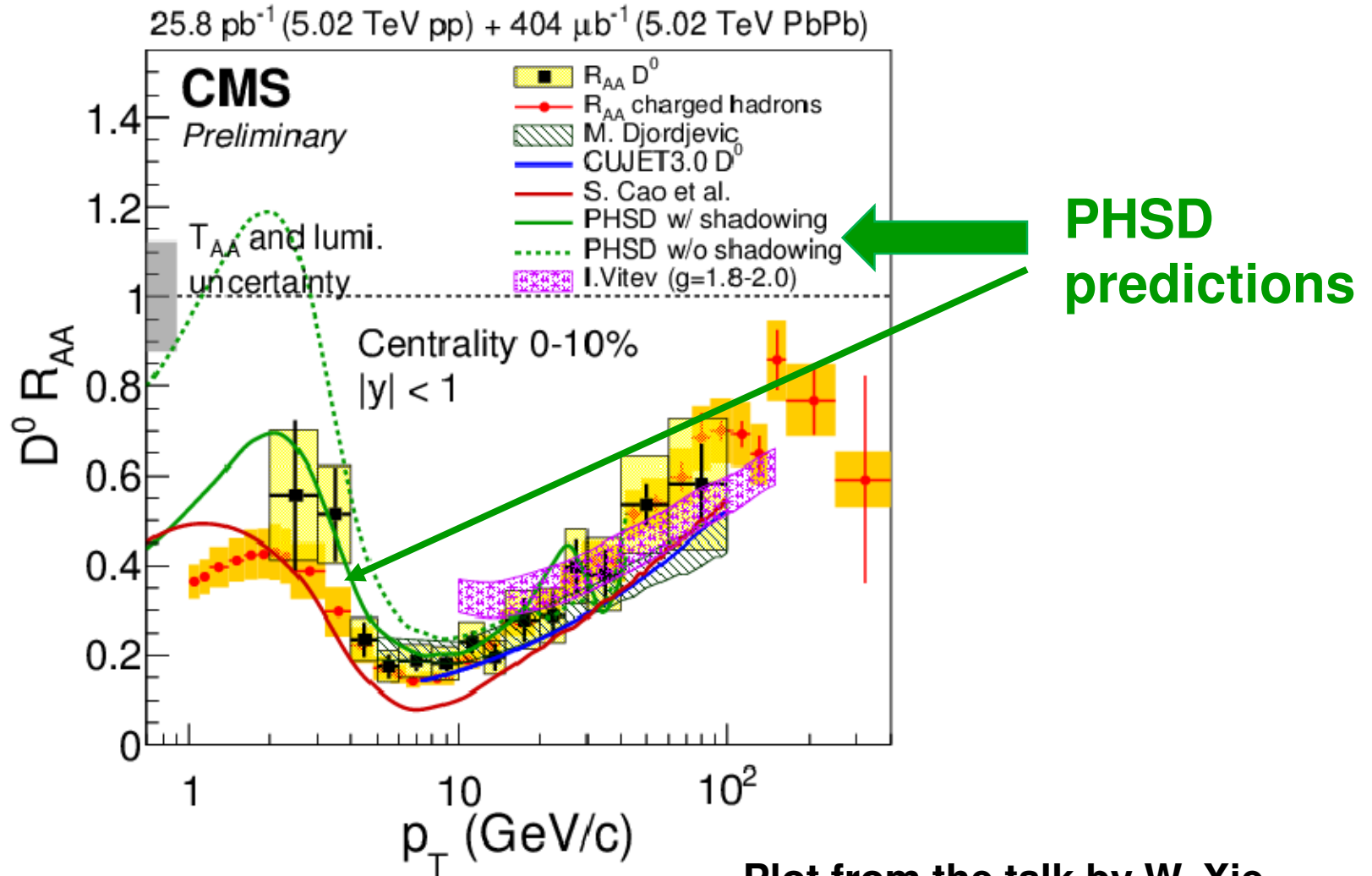


- in PHSD the energy loss of D-mesons at high p_T can be dominantly attributed to partonic scattering
- Shadowing effect suppresses the low p_T and slightly enhances the high p_T part of R_{AA}
- Hadronic rescattering moves R_{AA} peak to higher p_T ; increases v_2



Charm R_{AA} at LHC: PHSD predictions for CMS

D meson production is suppressed in 5.02 TeV PbPb collisions

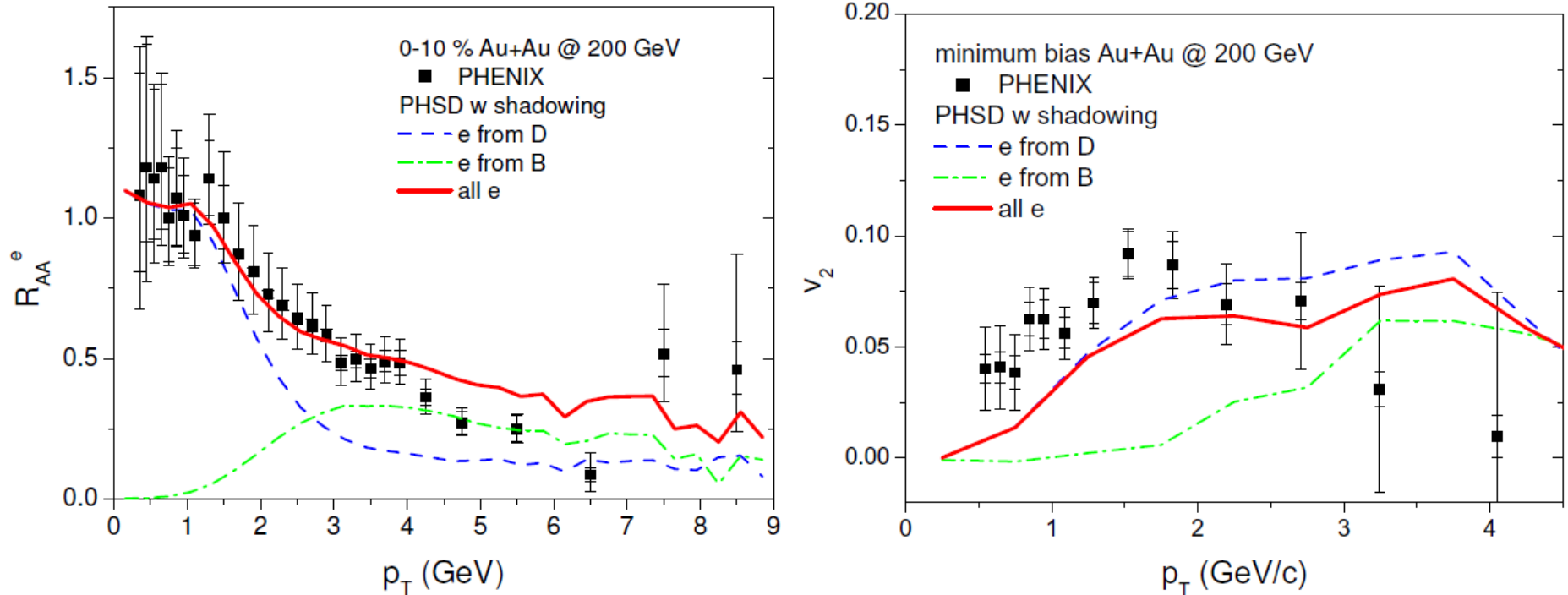


Plot from the talk by W. Xie



R_{AA}^e and v_2^e from single electrons: beauty contribution

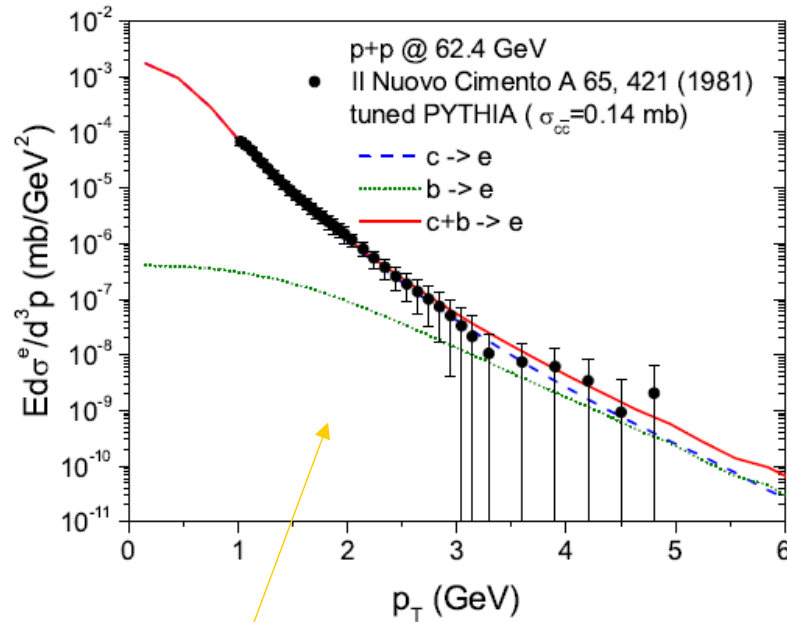
R_{AA} and v_2 vs p_T from single electrons in Au+Au @ 200 GeV



□ Feed back from beauty contribution becomes dominant for $p_T > 3$ GeV



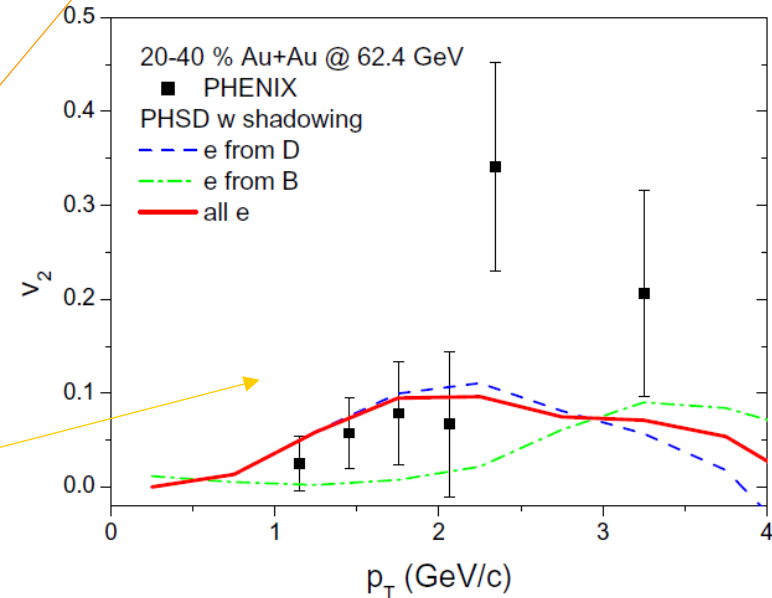
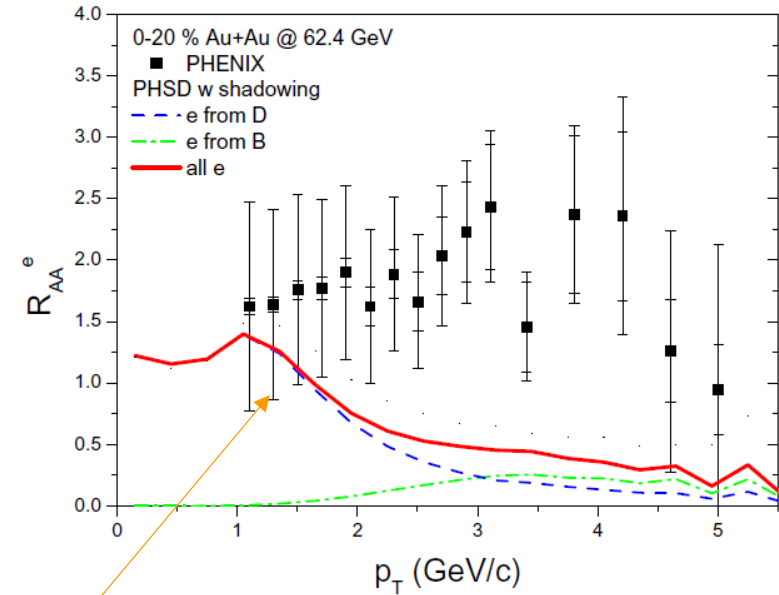
R_{AA}^e and v_2^e of single electrons from Au+Au at 62.4 GeV



□ PHSD: **pp data on electron p_T spectra** are well reproduced

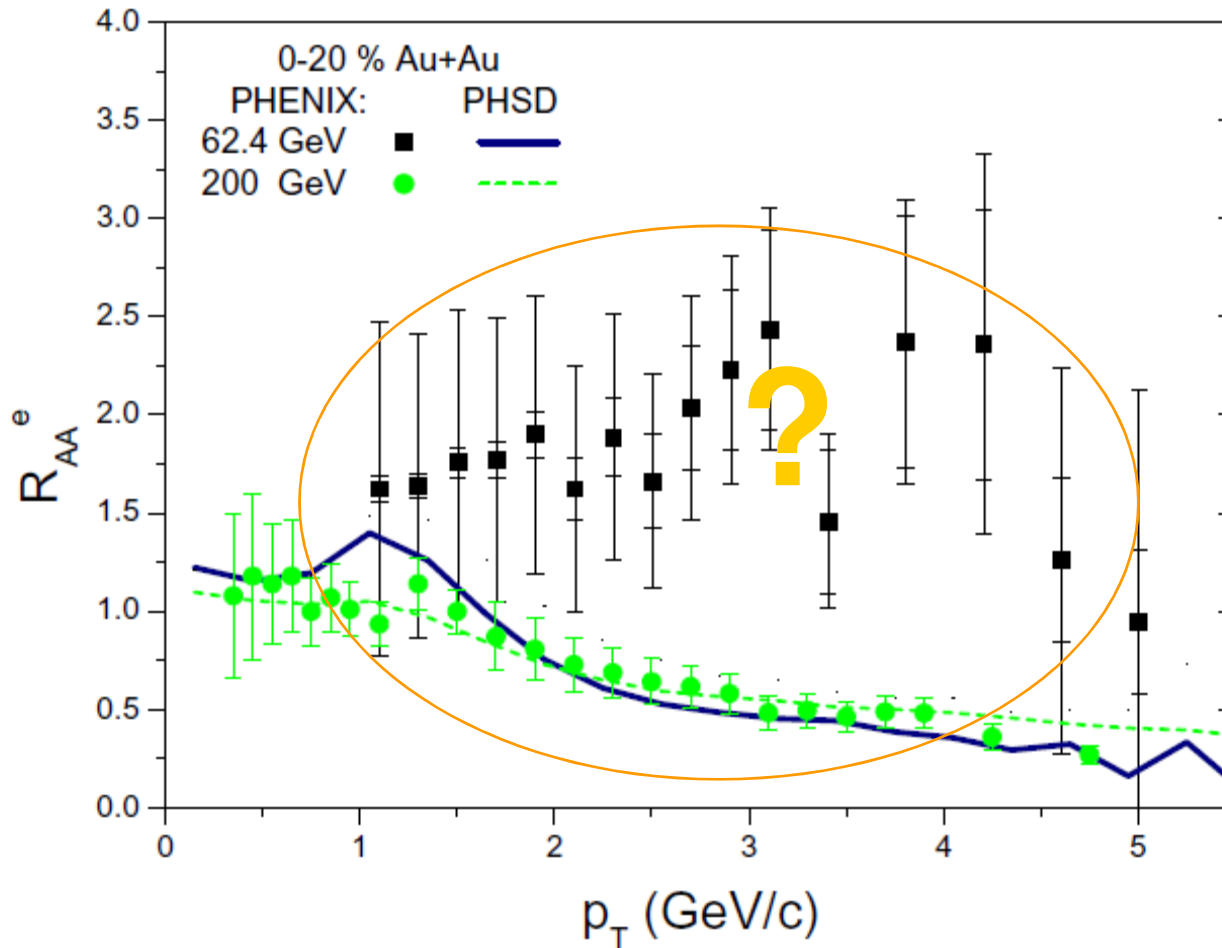
□ PHENIX data on R_{AA}^e from single electrons from **Au+Au at 62.4 GeV** are not reproduced !

□ v_2^e from single electrons from Au+Au at 62.4 GeV is in line with data





R_{AA}^e from of electrons: 62.4 vs. 200 GeV



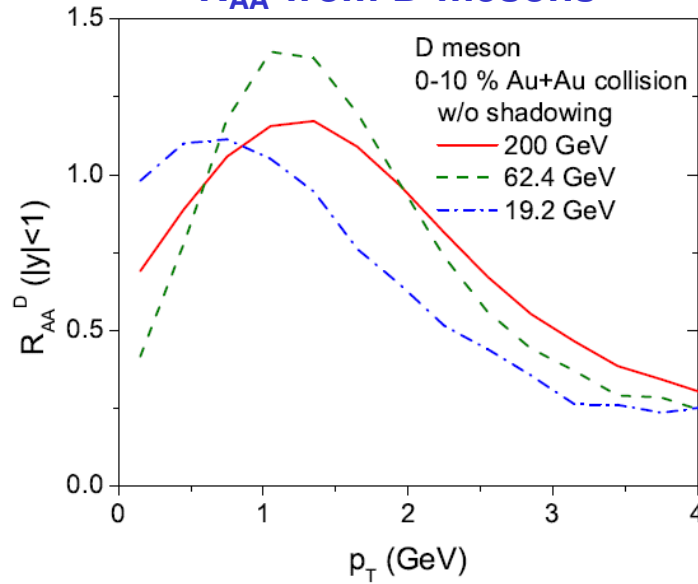
□ PHENIX:

R_{AA}^e of single electrons from **Au+Au at 62.4 GeV** is much **larger than at 200 GeV** !

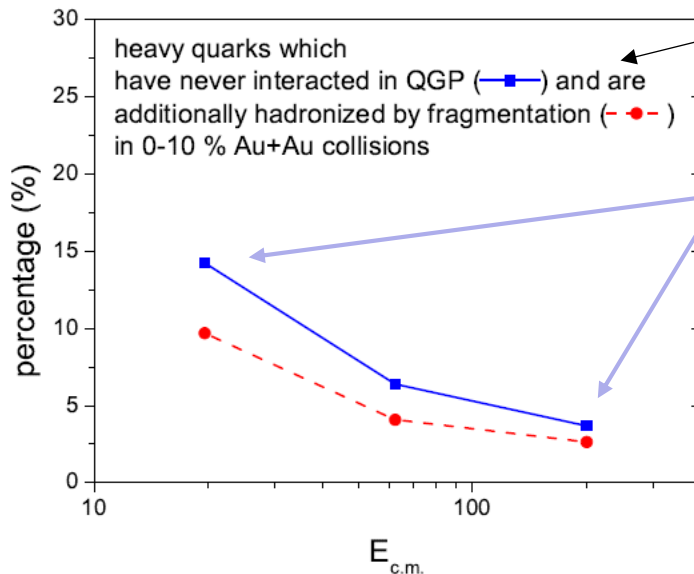
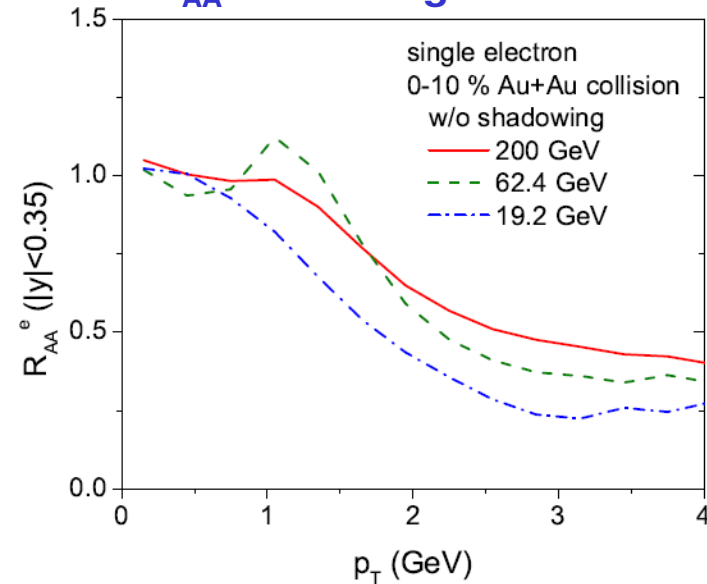


R_{AA} vs. beam energy

R_{AA} from D-mesons



R_{AA}^e from single electrons



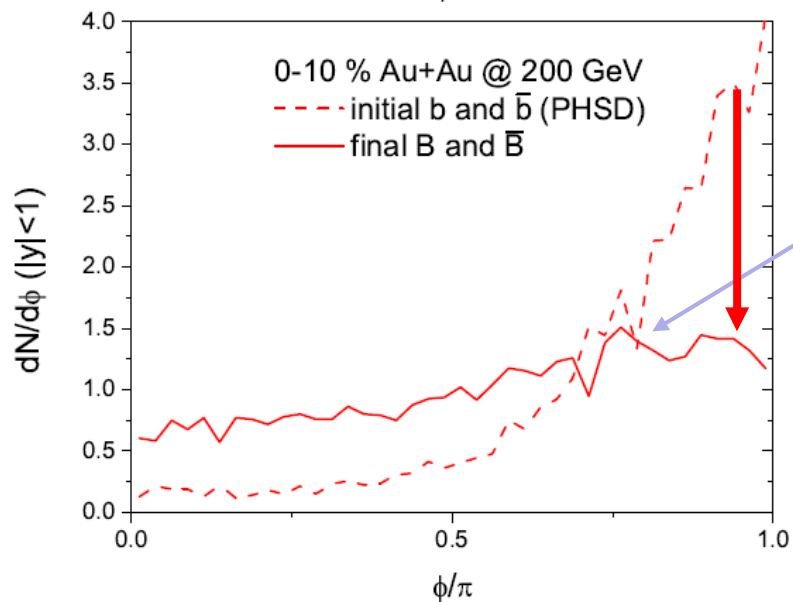
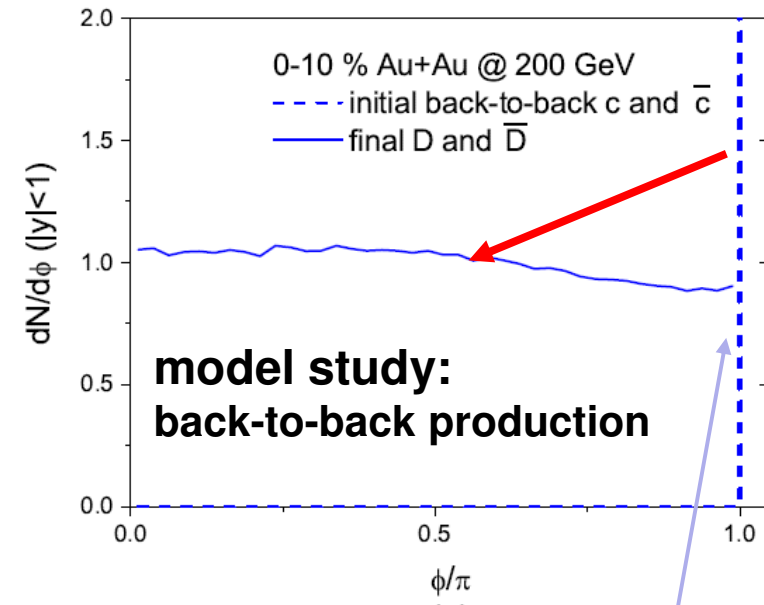
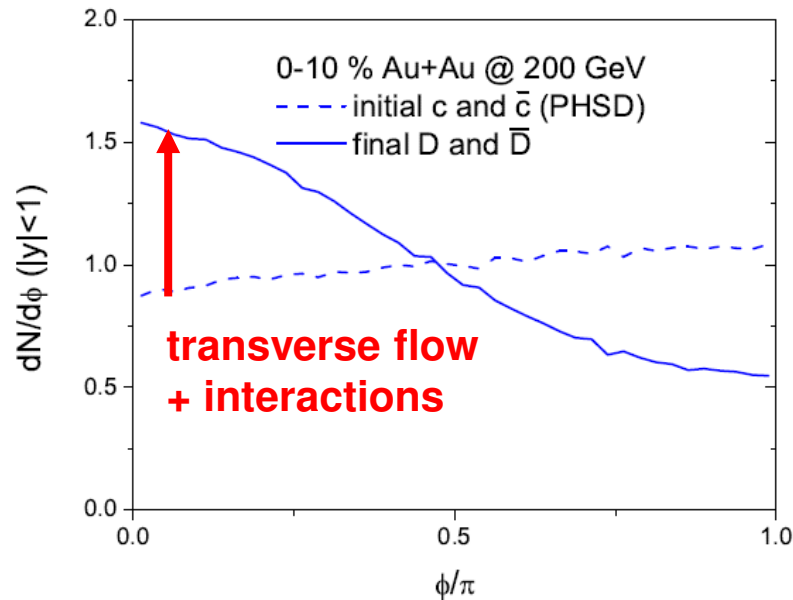
Fraction of HQ not interacting in QGP

- QGP fraction increases with energy
- Sizable Hadronic Corona at low energies

Au+Au, 200 GeV, central: each charm quark makes ~ 6 elastic collisions in QGP; each D, D* makes ~ 2 scatterings with hadrons



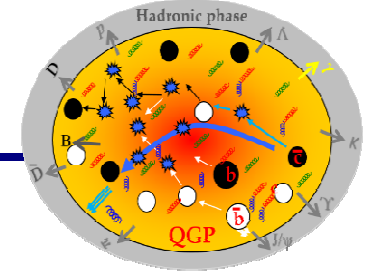
Azimuthal angular correlations: Q-Qbar



→ Initial azimuthal angular correlation of QQbar pairs is **completely washed out** during the evolution of the heavy-ion collision, even in case they are assumed to be initially produced **back-to-back (model study)** mainly **due to the transverse flow + interactions**



Summary



- ❑ **PHSD** provides a **microscopic description** of non-equilibrium charm dynamics in the partonic and hadronic phases
- ❑ **Partonic rescattering** suppresses the high p_T part of R_{AA} , generates v_2
- ❑ **Hadronic rescattering** moves R_{AA} peak to higher p_T , increases v_2
- ❑ The structure of R_{AA} at low p_T is sensitive to the **hadronization scenario**, i.e. to the balance between **coalescence and fragmentation**
- ❑ **Shadowing effects** suppress R_{AA} at LHC at low transverse momenta
- ❑ The **exp. data** for the R_{AA} and v_2 at RHIC and LHC are described in the PHSD by **QGP collisional energy loss** due to the **elastic scattering** of charm quarks with massive quarks and gluons in the QGP phase
 - + by the **dynamical hadronization scenario** „coalescence & fragmentation“
 - + by **strong hadronic interactions** due to resonant elastic scattering of D, D^* with mesons and baryons
- ❑ Feed back from **beauty contribution** for R_{AA}^e and v_2^e from single electrons for Au+Au at 200 GeV becomes dominant for $p_T > 3$ GeV
- ❑ **Initial azimuthal angular correlation** of $QQbar$ pairs is **washed out** during the evolution dominantly due to the transverse flow

Thank you!

