Charmonium production in pp collisions with ALICE at the LHC

Hugo Pereira Da Costa, CEA/IRFU for the ALICE Collaboration Strangeness in Quark Matter 2016 – Tuesday, June 28 2016

Motivation

Charmonia (J/ ψ , ψ (2S), χ_c) are bound states of a *ccbar* quark pair Their production is understood as the production of the *ccbar* pair in a hard-scattering process, followed by the evolution of this pair into a colorless bound state

The *c* quark mass should provide a high enough hard scale for pQCD to be applicable, but the evolution into a bound state is intrinsically non perturbative

Charmonium production measurements are interesting

- in pp, to understand production mechanism, probe Parton Distribution Functions (PDFs), especially for gluons, down to low x and as a reference to p-Pb and Pb-Pb
- in p-Pb, to probe so-called cold nuclear matter effects (modification of the PDFs, saturation, Cronin enhancement, etc.)
- in Pb-Pb, to probe the formation and properties of the Quark-Gluon Plasma (color screening, dissociation, recombination)

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Outline

Few words about the apparatus and analysis

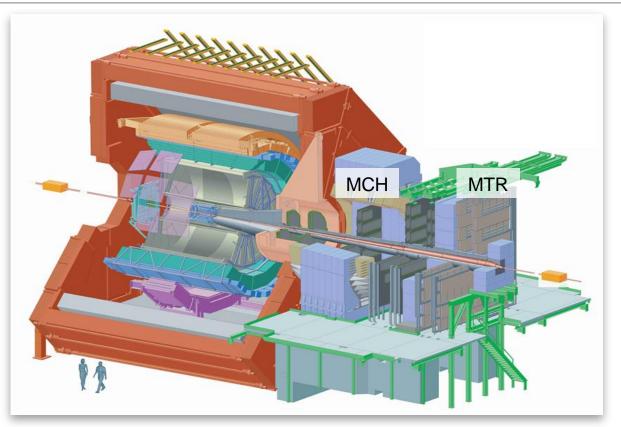
New results on forward-*y* J/ ψ and ψ (2S) production in pp collisions at \sqrt{s} = 13 TeV

- comparison to other experiments
- comparison to lower energies
- comparison to models

Inclusive measurements contain contributions from

- direct production
- decay from higher mass resonances (for J/ ψ , they are $\psi(2S)$ and χ_c)
- decay from *b*-hadrons (also called non-prompt J/ψ , $\psi(2S)$)

Forward-y charmonium measurements in ALICE



ALICE Muon system:

- 5 stations of tracking chambers (-4<η<-2.5)
- 2 stations of trigger chambers
- dipole magnet
- absorbers

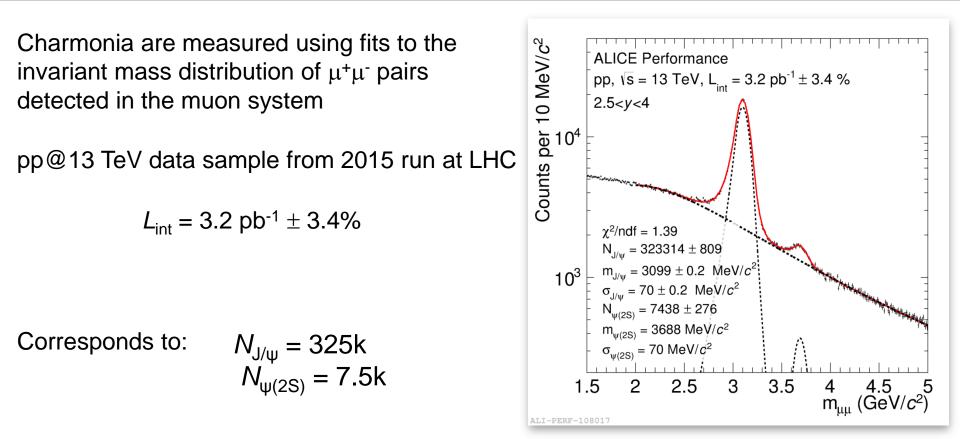
Charmonia are measured in the $\mu^+\mu^-$ decay channel, at forward rapidity (2.5<*y*<4) and down to $p_T = 0$ using

- ITS for vertex determination
- MTR for muon identification and triggering
- MCH for tracking

V0 detectors are also used for triggering (in coincidence with MTR)

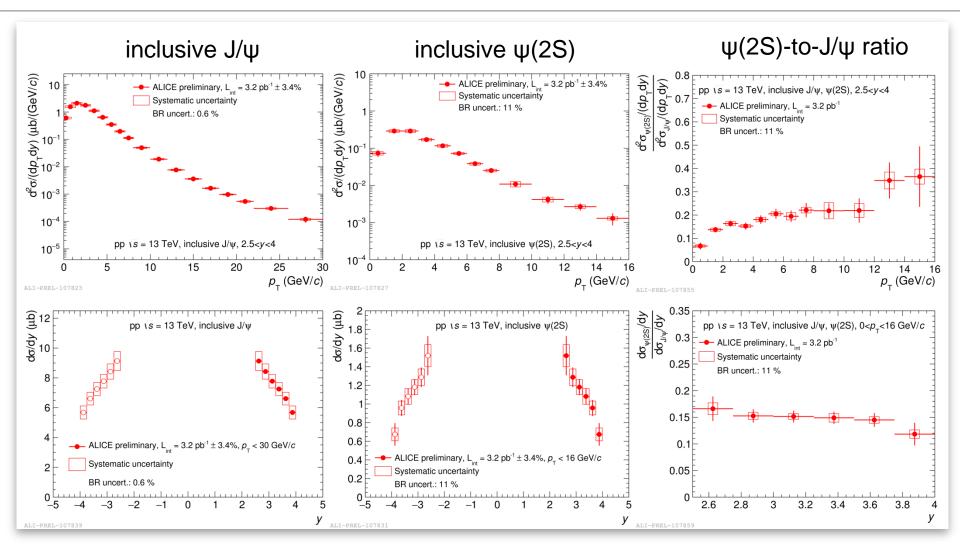
T0 detectors for luminosity determination

Data analysis



Systematic uncertainties on cross sections amount to ~7% for J/ ψ and ~10% for ψ (2S). They include contributions from signal extraction, acceptance x efficiency corrections and luminosity

Cross sections and particle ratios in pp@13 TeV



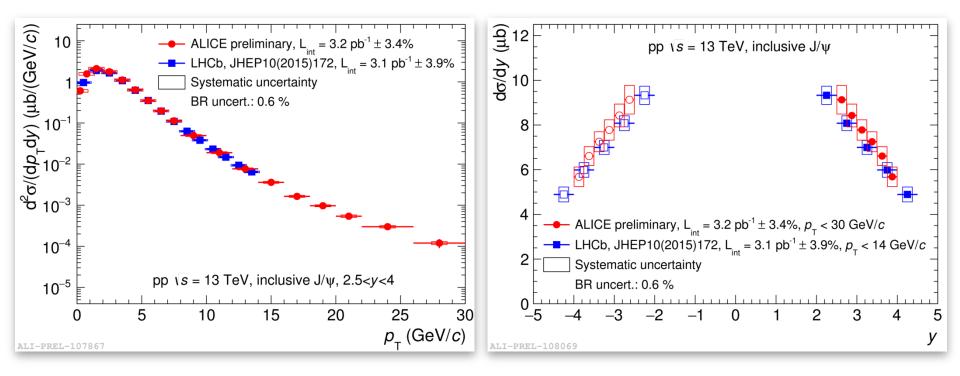
We reach $p_T = 30 \text{ GeV}/c$ for J/ ψ , and 16 GeV/*c* for $\psi(2S)$ as well as $\psi(2S)$ -to-J/ ψ ratio

We measure 6 bins in y for 2.5<y<4

Comparison to LHCb

LHCb results from JHEP10 (2015) 172

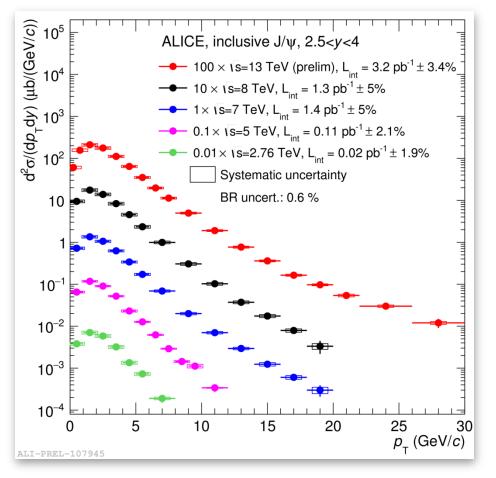
Quoted values correspond to the sum of the prompt and non-prompt contributions, integrated over the same rapidity range as ALICE (2.5<y<4)



Excellent agreement between the two experiments

All points lie within 1 sigma (stat+syst) one with the other

Comparison to lower energy, J/ ψ vs p_T



 $\sqrt{s} = 2.76 \text{ TeV}$ PLB 718 (2012) 295 $\sqrt{s} = 5 \text{ TeV}$ ArXiv:1606.08197 $\sqrt{s} = 7 \text{ TeV}$ EPJC 74 (2014) 2974 $\sqrt{s} = 8 \text{ TeV}$ EPJC 76 (2016) 184 Comparison to ALICE measurements at \sqrt{s} = 2.76, 5, 7 and 8 TeV

All results published except at 13 TeV

Steady increase of the luminosity and $p_{\rm T}$ reach with increasing energy

As expected, spectra becomes harder with increasing energy

Change of slope at high p_T and $\sqrt{s} = 13$ TeV, attributed to the onset of the nonprompt J/ ψ contribution

Comparison to lower energy, $\psi(2S)$ vs p_T

inclusive J/ψ inclusive $\psi(2S)$ d²σ/(dp_Tdy) (μb/(GeV/*c*)) 0⁵ 10 ALICE, inclusive J/ψ , 2.5<y<4 ALICE, inclusive $\psi(2S)$, 2.5<y<4 ● 100 × \s=13 TeV (prelim), L_{int} = 3.2 pb⁻¹ ± 3.4%_ -• 1× is=13 TeV (prelim), L_{int} = 3.2 pb⁻¹ ± 3.4% 0^{4} $-\bullet$ 0.1× is=8 TeV, L_{int} = 1.3 pb⁻¹ ± 5% -• $10 \times 10 = 8 \text{ TeV}, L_{\text{int}} = 1.3 \text{ pb}^{-1} \pm 5\%$ ----- 1×1 = 7 TeV, $L_{int} = 1.4 \text{ pb}^{-1} \pm 5\%$ 0³ Systematic uncertainty 0.1×1 s=5 TeV, L_{int} = 0.11 pb⁻¹ ± 2.1% BR uncert.: 11 % $0.01 \times \sqrt{s}=2.76 \text{ TeV}, L_{int} = 0.02 \text{ pb}^{-1} \pm 1.9\%$ 0² Systematic uncertainty BR uncert.: 0.6 % 10 10^{-2} 10^{-1} 10^{-3} 10^{-2} 10^{-4} 10^{-3} 10^{-4} 10^{-5} 25 5 10 15 20 30 2 4 6 8 10 12 14 16 $p_{_{\rm T}}$ (GeV/c) $p_{_{T}}$ (GeV/c) ALI-PREL-107945 ALI-PREL-107957

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For $\psi(2S)$ we have measurements at $\sqrt{s} = 7$, 8 and 13 TeV

Comparison to lower energy, J/ ψ and ψ (2S) vs y

inclusive $\psi(2S)$ inclusive J/ψ 10⁴ (מו) 10¢ 1 1 (qn) /p/op 0^{3} ALICE, inclusive J/ψ ALICE, inclusive $\psi(2S)$ ► 100 × \s=13 TeV (prelim) $L_{int} = 3.2 \text{ pb}^{-1} \pm 3.4\%$ $L_{int} = 3.2 \text{ pb}^{-1} \pm 3.4\%$ - 10 × \s=8 TeV 10² $L_{int} = 1.3 \text{ pb}^{-1} \pm 5.0\%$ $L_{int} = 1.3 \text{ pb}^{-1} \pm 5.0\%$ •– 1× ∖s=7 TeV ---- 1 × √s=7 TeV 10 $L_{int} = 1.4 \text{ pb}^{-1} \pm 5.0\%$ 10 $L_{int} = 1.4 \text{ pb}^{-1} \pm 5.0\%$ Systematic uncertainty - 0.10 × ∖s=2.76 TeV BR uncert.: 11 % $L_{int} = 0.02 \text{ pb}^{-1} \pm 1.9\%$ Systematic uncertainty BR uncert.: 0.6 % 10^{-1} 10^{-1} 2 2.5 4.5 5.5 5.5 2 2.5 3 3.5 5 6 3 3.5 4.5 5 4 6 4 V ALI-PREL-108414 ALI-PREL-108141

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For J/ ψ , no visible change in the *y* distribution For $\psi(2S)$, large uncertainties prevent firm conclusions

A few words about models

Mainly three approaches used to describe direct charmonium production in pp

- Color Evaporation Model (CEM): production cross section of a given charmonium is proportional to the *ccbar* cross section, integrated between the mass of the charmonium and twice the mass of the D meson. Proportionality factor is independent of *y*, *p*_T and √*s*
- Color Singlet model (CSM): pQCD is used to describe the *ccbar* production with the same quantum numbers (CS) as the final-state meson.
- Non-Relativistic QCD (NRQCD):

both CS and CO state of the *ccbar* pairs are considered. The relative contribution of the states is parametrized using a finite set of universal long range matrix elements (LRME), fitted to a subset of the data (e.g. Tevatron)

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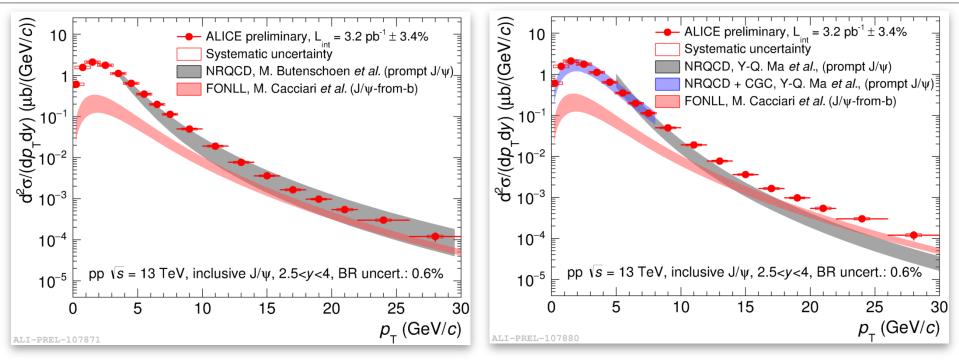
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Non-prompt contribution corresponds to the production of *b*-hadrons. Can be calculated with pQCD, e.g. within FONLL

Comparison to models, $J/\psi vs p_T$



 NRQCD (left)
 Butenschon and Kniel, PRL 106 (2011) 022003

 NRQCD (right)
 Ma, Wang and Chao, PRL 106 (2011) 042002

 NRQCD+CGC
 Ma and Venugopalan, PRL 113 (2014) 192301

 FONLL
 Cacciari *et al.*, JHEP 1210 (2012) 137

All models properly account for higher mass resonance decays

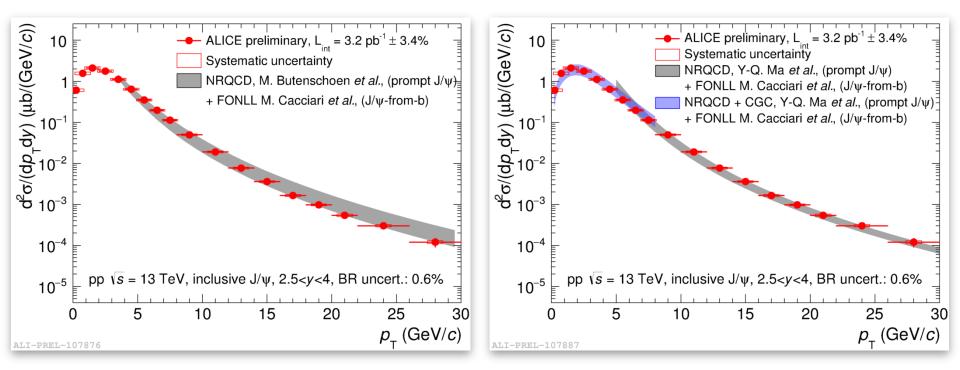
NRQCD models differ in the set of LRME that is used, the p_T at which fits are performed and the datasets considered.

At low p_T (right), NRQCD is coupled to a CGC description of the proton

Predictions are quite different at high p_{T} , but in both cases, non-prompt J/ ψ constitute a sizable contribution to the inclusive cross section

Comparison to models, $J/\psi vs p_T$

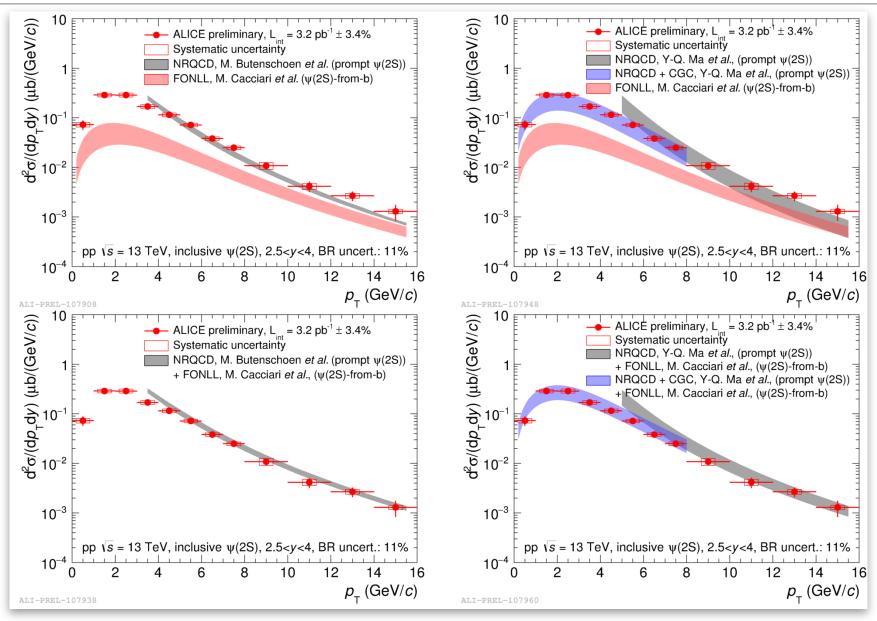
As an exercise, we summed NRQCD and FONLL calculations assuming fully uncorrelated uncertainties.



Agreement to the data is much improved, already at intermediate p_T and especially for the calculation from Ma *et al.*

Note that the calculations are completely independent, and that there was no data at this energy and at such high $p_{\rm T}$ before

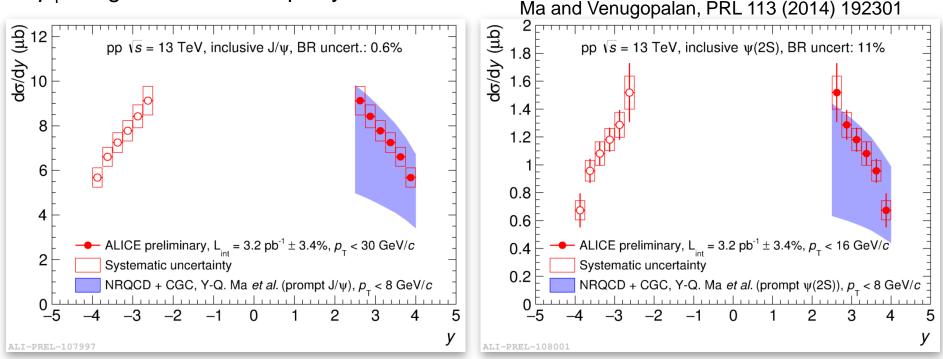
Comparison to models, $\psi(2S)$ vs p_T



Same conclusions as for the J/ψ case

Rapidity distributions

Since the NRQCD+CGC calculation goes down to $p_T = 0$, it can be compared to our p_T -integrated data vs rapidity

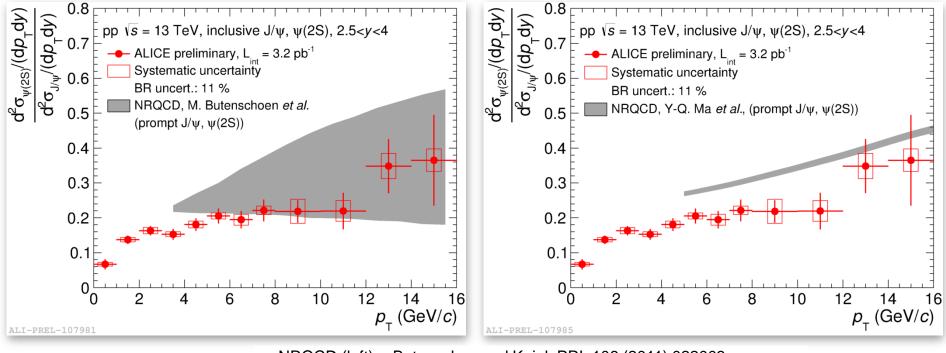


Agreement to the data is reasonable, and would be improved when adding the contribution from non-prompt J/ ψ and ψ (2S) (10-15%)

The only other calculation that provides p_T -integrated rapidity distributions is CSM@LO, with significantly larger uncertainties (see ALICE (7TeV) EPJC 74 (2014) 2974)

Particle ratio, comparison to models

Many systematic uncertainties cancel in the particle ratio, for both data and theory



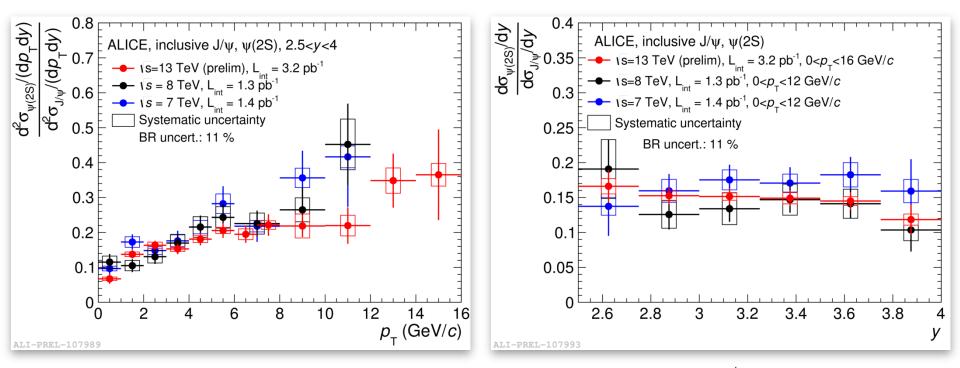
NRQCD (left) Butenschon and Kniel, PRL 106 (2011) 022003 NRQCD (right) Ma, Wang and Chao, PRL 106 (2011) 042002

Both calculations follow the same trend but with very different uncertainties. This was already the case at $\sqrt{s} = 7$ TeV (see ALICE EPJC 74 (2014) 2974)

Calculation from Y-Q Ma et al. tends to overestimate the $\psi(2S)$ -to-J/ ψ ratio

Contributions from non-prompt J/ ψ and ψ (2S) have little impact here because they enter both the numerator and denominator, with a similar (small) magnitude

Particle ratio, energy dependence



No visible dependence of the p_{T} -differential $\psi(2S)$ -to-J/ ψ ratio on \sqrt{s}

No clear trend either vs rapidity

Summary

- ALICE has measured inclusive J/ ψ and ψ (2S) production at forward-*y* in pp collisions at $\sqrt{s} = 13$ TeV
- For J/ ψ , our result is consistent with LHCb and extends the p_{T} reach up to 30 GeV/*c*
- For $\psi(2S)$, this is the only measurement available at this energy
- Data can be well reproduced
- at low p_T by a model that couple a CGC description of the proton to NRQCD
- at intermediate and high- $p_{\rm T}$ by NRQCD calculations
- provided that the non-prompt contribution is properly accounted for (e.g. with FONLL)
- Rapidity distribution is also well reproduced by CGC+NRQCD calculation
- $\psi(2S)\mbox{-to-J/}\psi$ ratio is slightly overestimated by models, and shows no dependence on the collision energy

Outlook:

Starting from 2021, the addition of the MFT will allow to properly separate prompt and non-prompt forward-*y* charmonia in pp, p-Pb and Pb-Pb