

Heavy and light flavor jet quenching

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Strangeness in Quark Matter 2016

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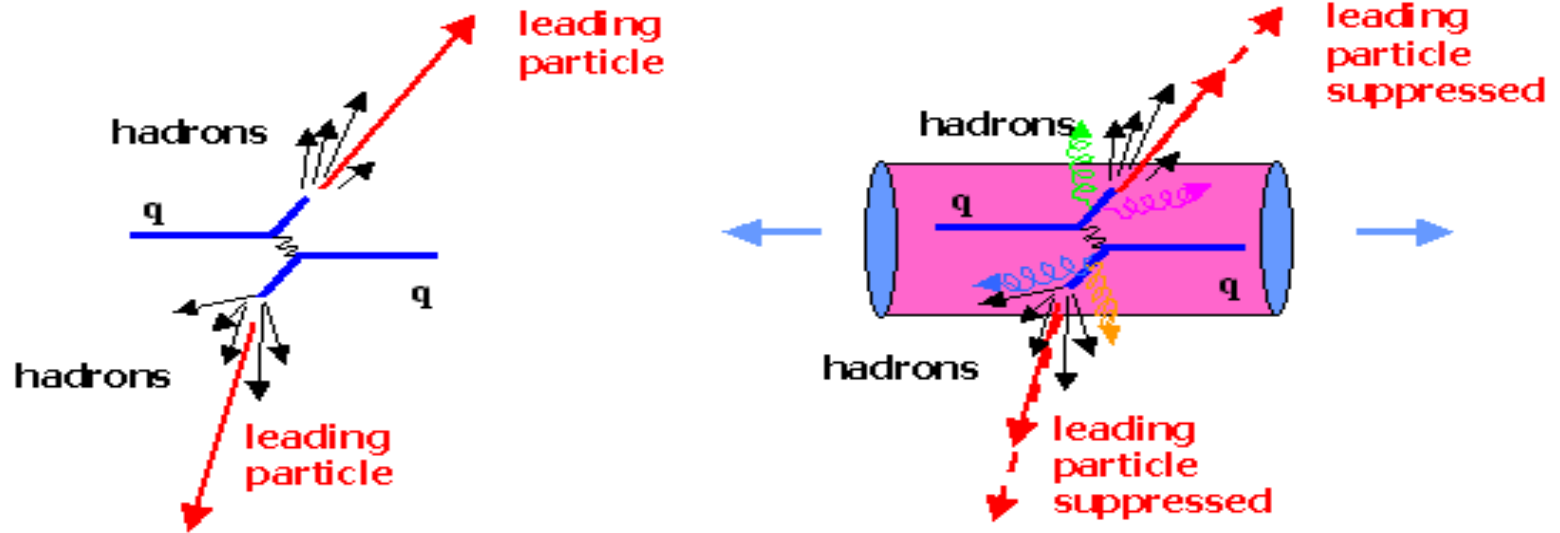
June 27-July 1, 2016



Outline

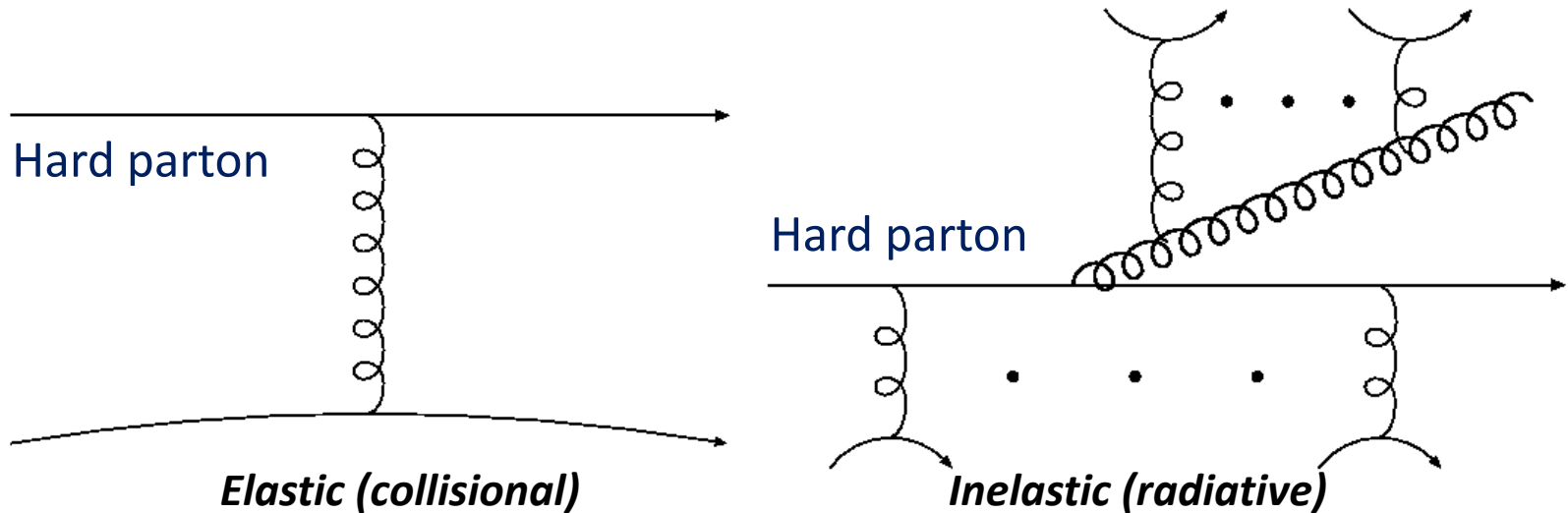
- A Linearized Boltzmann Transport (LBT) approach to heavy/light flavor jet quenching (with elastic & inelastic contributions)
- Full jet energy loss and modification (collisional, radiative, broadening)
- Use jet-like angular de-correlation to probe medium-induced broadening (q^{hat})

Jet quenching



- The study of jet quenching/modification can provide valuable information about hot and dense QGP produced in heavy-ion collisions

Radiative & collisional processes



In the limit of soft scatterings, the effect of elastic collisions can be described by FP equation (**longitudinal drag, longitudinal diffusion & transverse diffusion**)

$$\frac{\partial f}{\partial z^-} = \left[D_{L1} \frac{\partial}{\partial l_q^-} + \frac{1}{2} D_{L2} \frac{\partial^2}{\partial^2 l_q^-} + \frac{1}{2} D_{T2} \nabla_{\vec{l}_{q\perp}}^2 \right] f(z^-, l_q^-, \vec{l}_{q\perp})$$

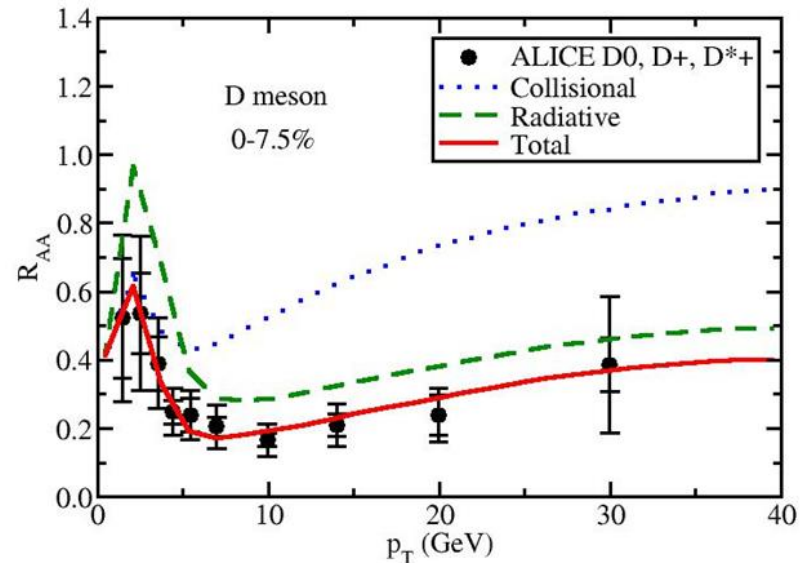
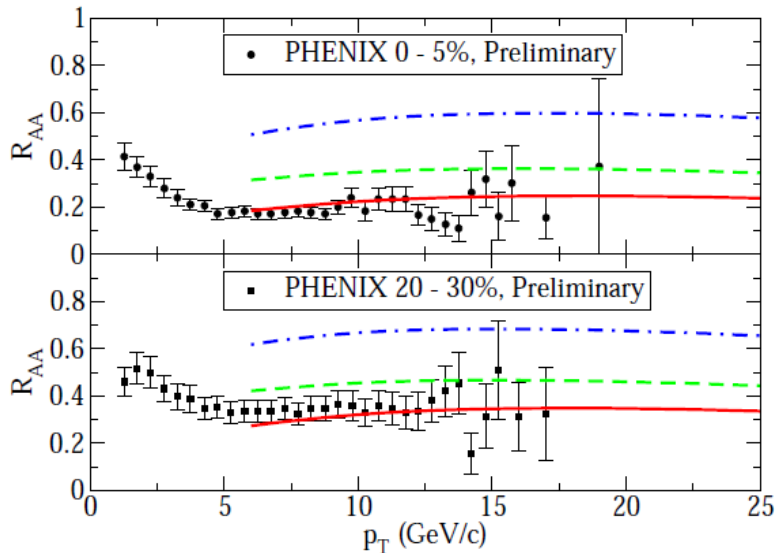
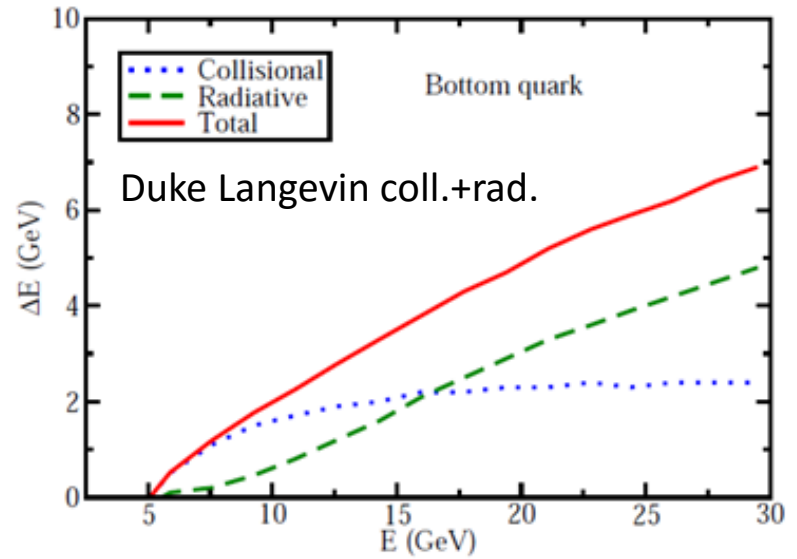
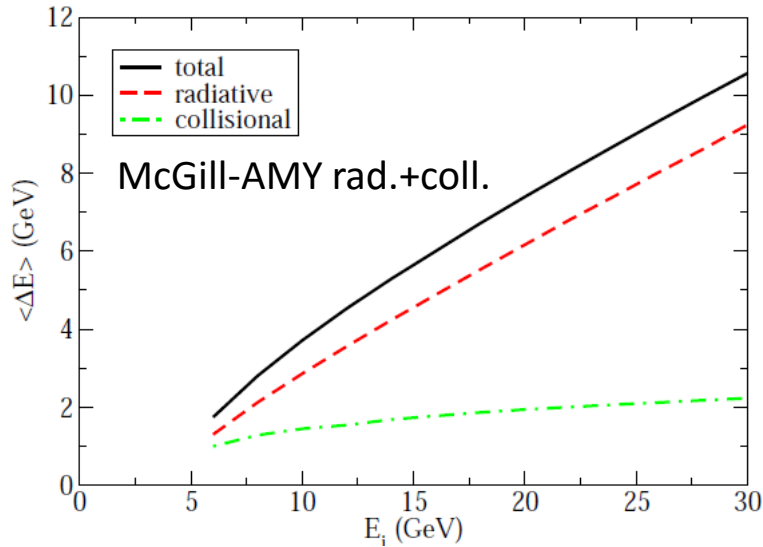
The medium-induced gluon radiation spectrum from higher-twist formalism:

$$\frac{dN_g}{dx dl_{\perp}^2 dz^-} \approx \frac{4\alpha_s}{\pi} P(X) \frac{D_{T,2}}{l_{\perp}^4} \sin^2 \left(\frac{z^- - z_i^-}{2\tau_f^-} \right)$$

GYQ, Majumder, PRC 2013
Guo, Wang, PRL, 2000
Majumder, PRD, 2012

Jet transport coefficients control both collisional and radiative contributions

Radiative & collisional contributions



A Linearized Boltzmann Transport (LBT) approach for heavy & light flavor jet quenching

- **Boltzmann equation:** $p_1 \cdot \partial f_1(x_1, p_1) = E_1 C [f_1]$

- **Elastic collisions:**

$$\Gamma_{12 \rightarrow 34} = \frac{\gamma_2}{2E_1} \int \frac{d^3 p_2}{(2\pi)^3 2E_2} \int \frac{d^3 p_3}{(2\pi)^3 2E_3} \int \frac{d^3 p_4}{(2\pi)^3 2E_4}$$

$$\times f_2(\vec{p}_2) \left[1 \pm f_3(\vec{p}_1 - \vec{k}) \right] \left[1 \pm f_4(\vec{p}_2 + \vec{k}) \right] S_2(s, t, u)$$

$$\times (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - p_4) |\mathcal{M}_{12 \rightarrow 34}|^2$$

$$P_{\text{el}} = \Gamma \Delta t$$

- **Inelastic collisions:**

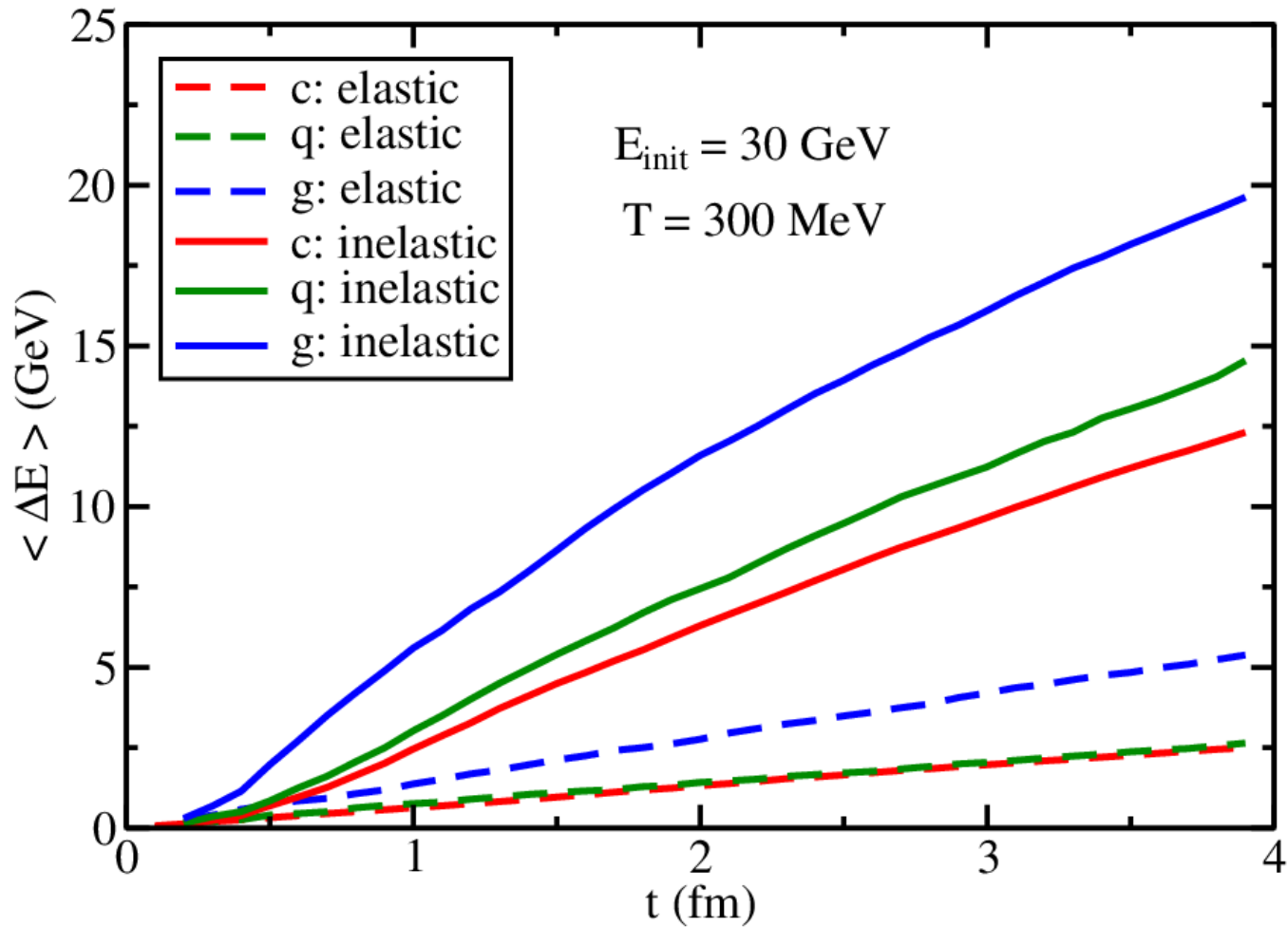
$$\langle N_g \rangle (E, T, t, \Delta t) = \Delta t \int dx dk_{\perp}^2 \frac{dN_g}{dx dk_{\perp}^2 dt}$$

$$P(n) = \frac{\langle N_g \rangle^n}{n!} e^{-\langle N_g \rangle}$$

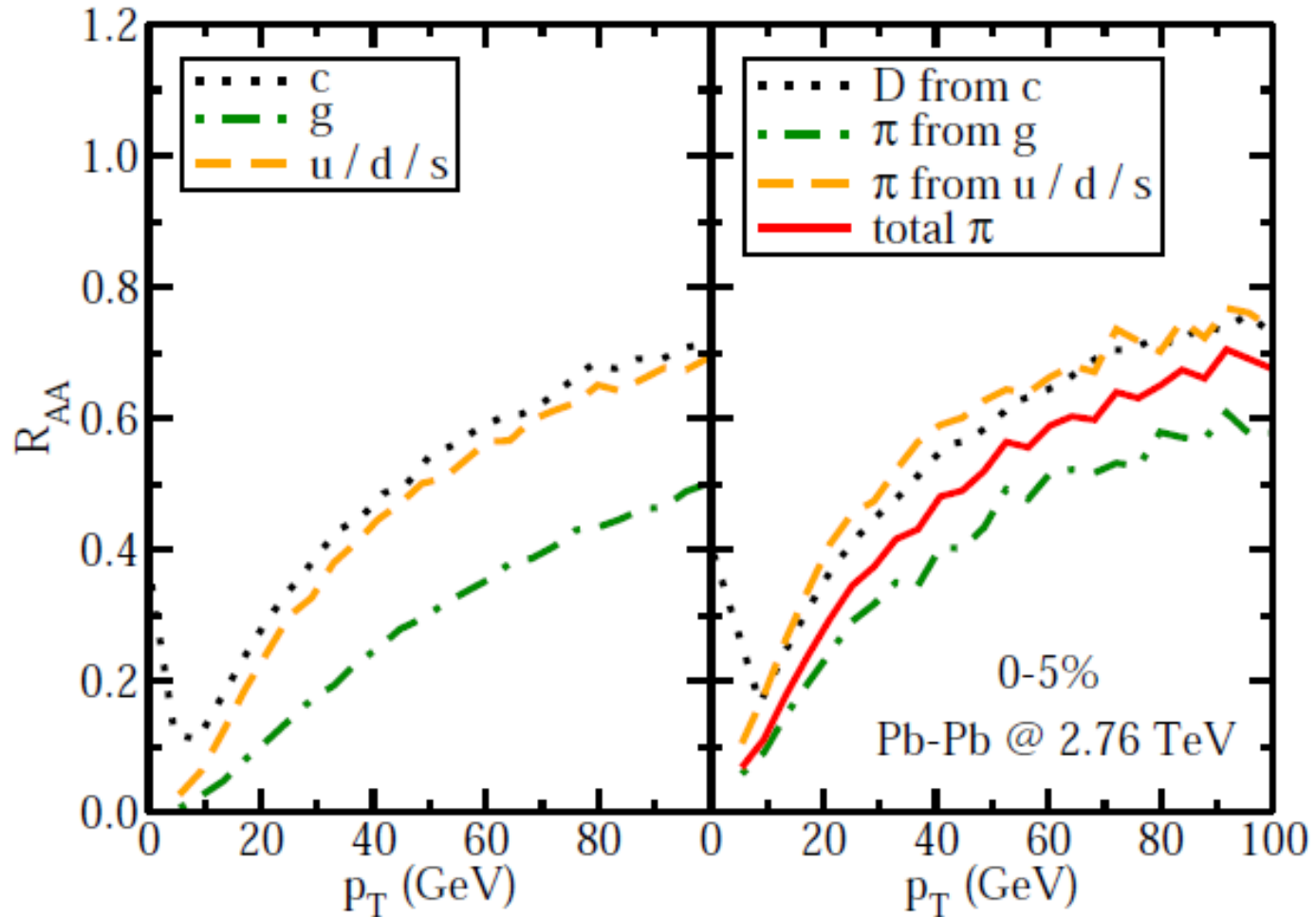
$$P_{\text{inel}} = 1 - e^{-\langle N_g \rangle}$$

- **Elastic + Inelastic:** $P_{\text{tot}} = P_{\text{el}}(1 - P_{\text{inel}}) + P_{\text{inel}}$

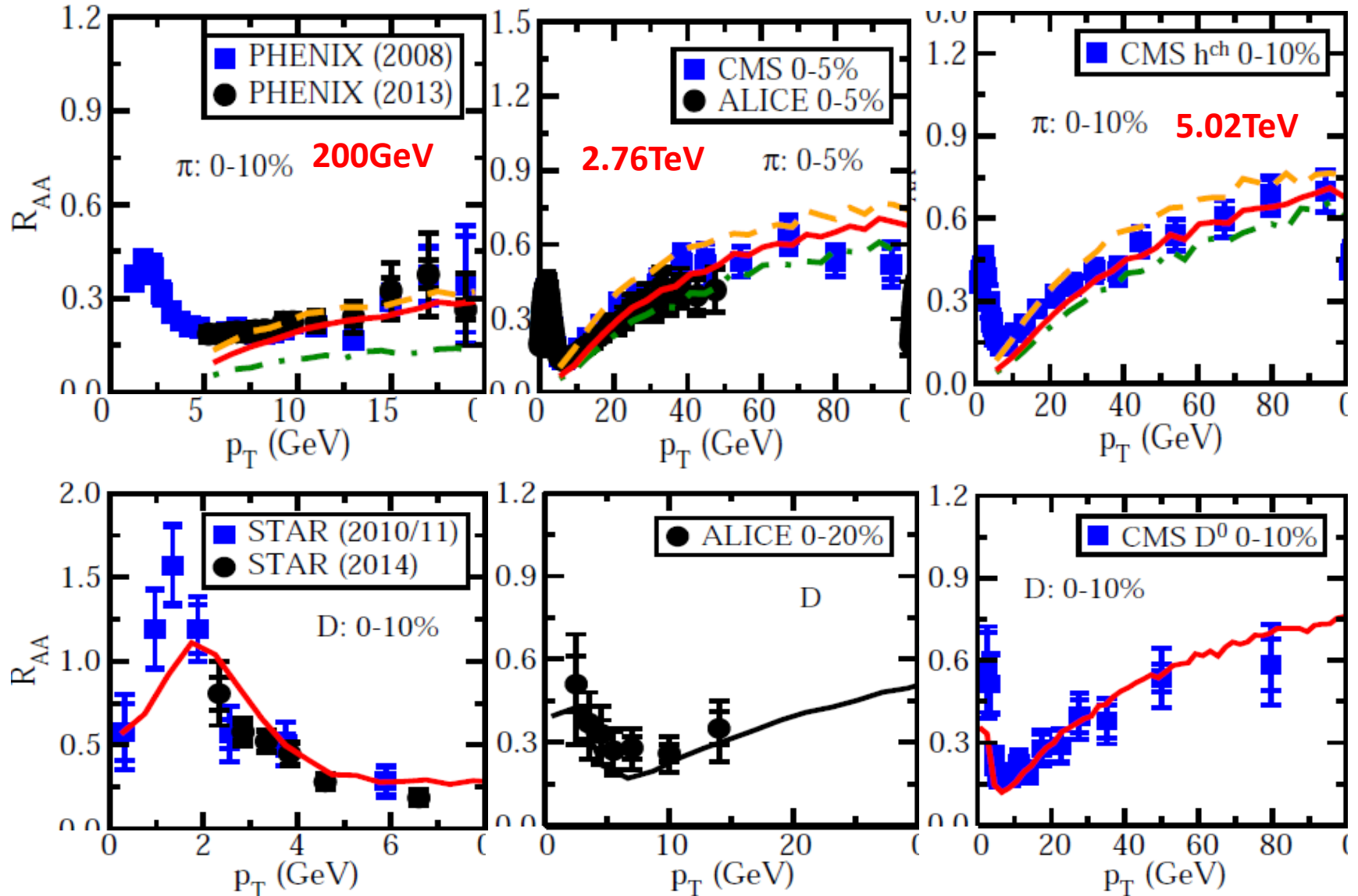
Elastic & inelastic energy loss from LBT



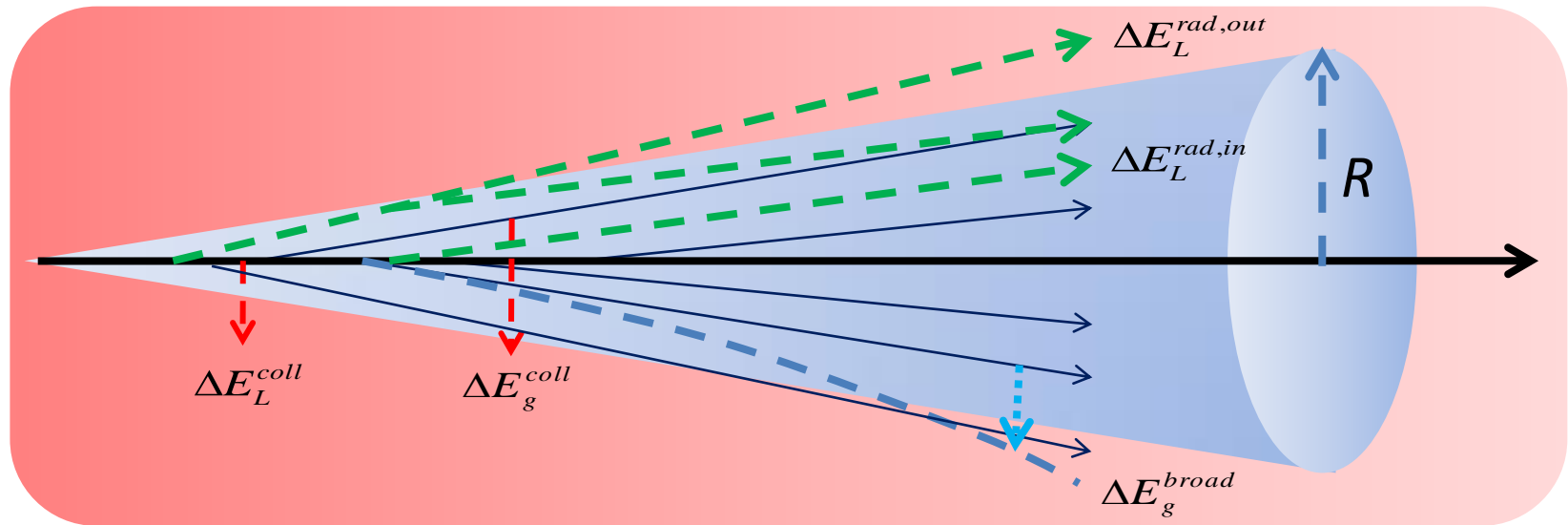
Quenching hierarchy



R_{AA} from LBT (heavy & light flavor hadrons)



Full jet evolution in medium



Not only the interaction of the leading hard parton with the medium constituents, but also the fate of radiated shower partons

$$\begin{aligned} E_{\text{jet}} &= E_{\text{in}} + E_{\text{lost}} \\ &= E_{\text{in}} + E_{\text{out}}(\text{radiation}) + E_{\text{out}}(\text{broadening}) + E_{\text{th}}(\text{collision}) \end{aligned}$$

GYQ, Muller, PRL, 2011; Casalderrey-Solana, Milhano, Wiedemann, JPG 2011; Young, Schenke, Jeon, Gale, PRC, 2011; Dai, Vitev, Zhang, PRL 2013; Wang, Zhu, PRL 2013; Blaizot, Iancu, Mehtar-Tani, PRL 2013; etc.

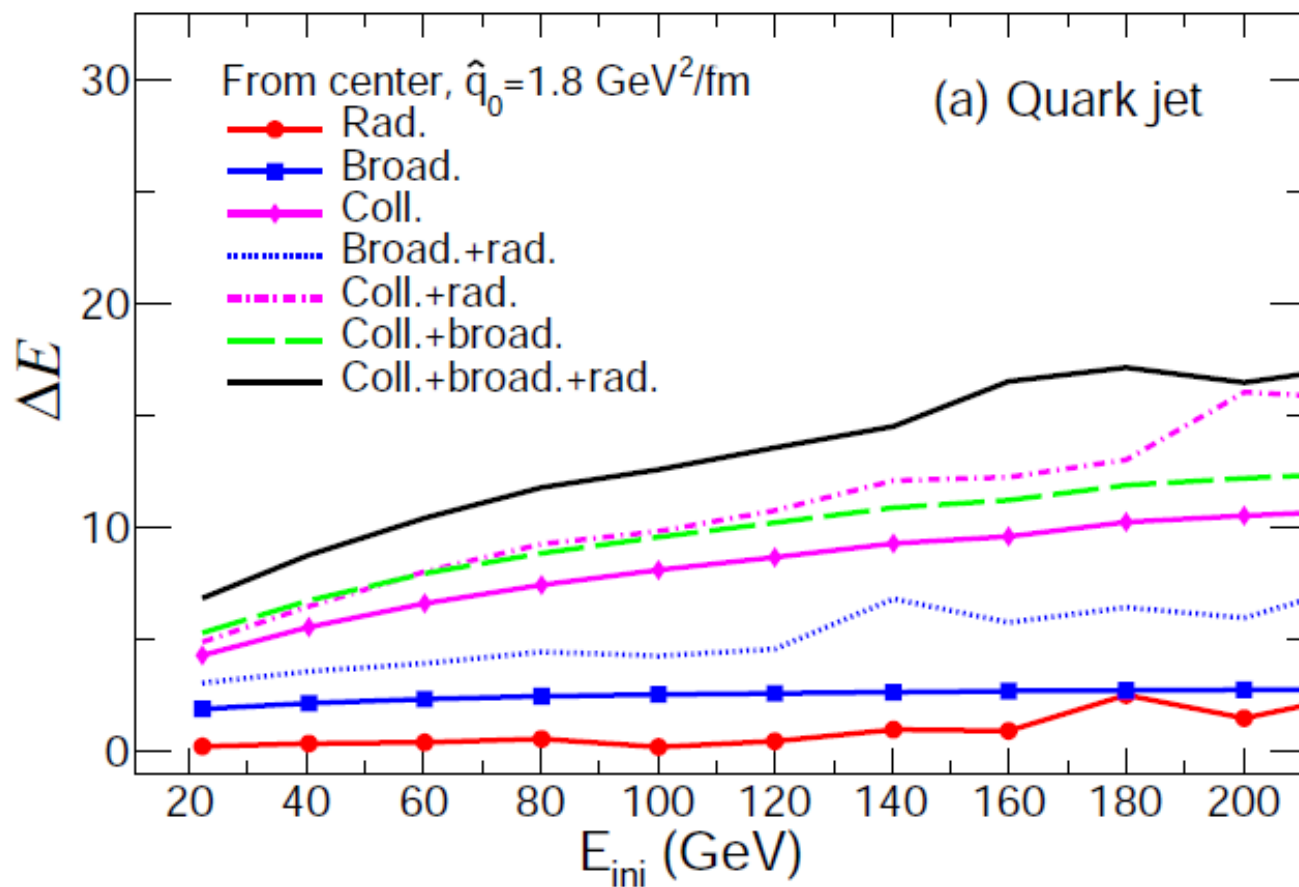
Full jet evolution in medium

- Solve the 3D (energy & transverse momentum) evolution for shower partons inside the full jet
- Include both collisional (the longitudinal drag and transverse diffusion) and all radiative/splitting processes

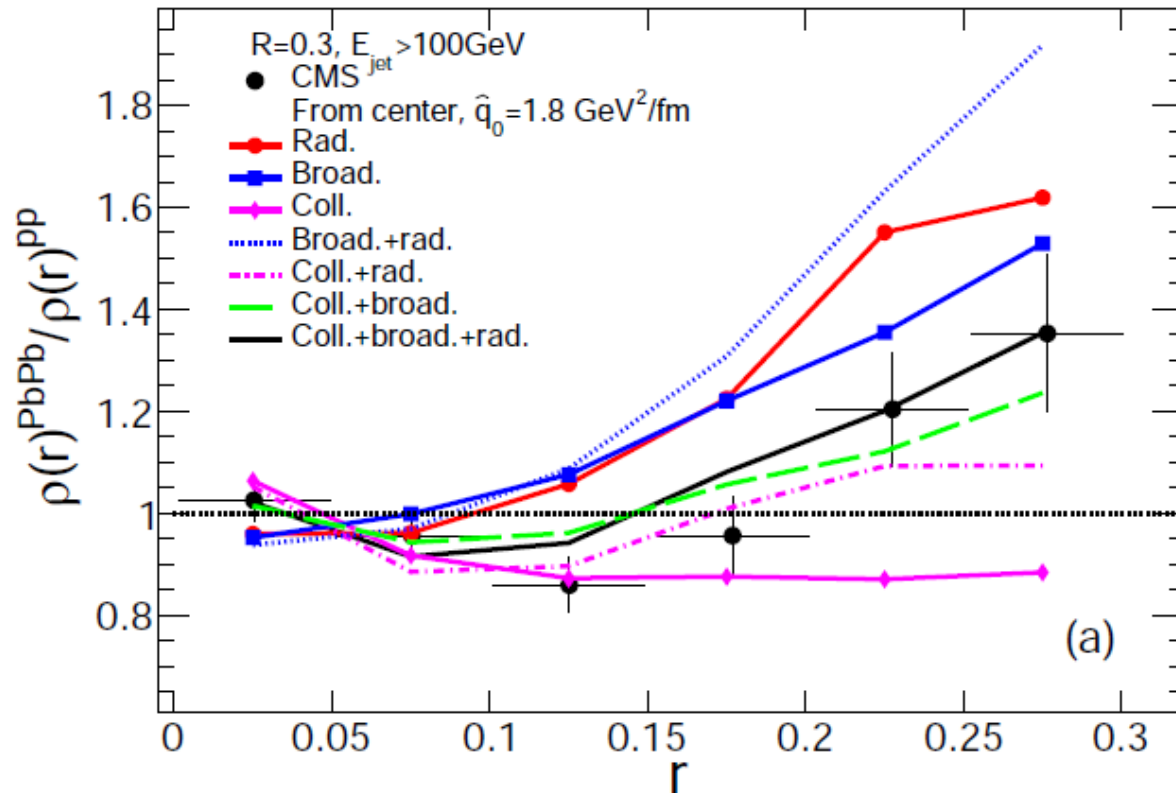
$$\begin{aligned}
 \frac{d}{dt} f_j(\omega_j, k_{j\perp}^2, t) &= \left(\hat{e}_j \frac{\partial}{\partial \omega_j} + \frac{1}{4} \hat{q}_j \nabla_{k_\perp}^2 \right) f_j(\omega_j, k_{j\perp}^2, t) && \text{Drag \& diffusion} \\
 + \sum_i \int d\omega_i dk_{i\perp}^2 &\frac{d\tilde{\Gamma}_{i \rightarrow j}(\omega_j, k_{j\perp}^2 | \omega_i, k_{i\perp}^2)}{d\omega_j d^2 k_{j\perp} dt} f_i(\omega_i, k_{i\perp}^2, t) && \text{Gain terms} \\
 - \sum_i \int d\omega_i dk_{i\perp}^2 &\frac{d\tilde{\Gamma}_{j \rightarrow i}(\omega_i, k_{i\perp}^2 | \omega_j, k_{j\perp}^2)}{d\omega_i d^2 k_{i\perp} dt} f_j(\omega_j, k_{j\perp}^2, t) && \text{Loss terms}
 \end{aligned}$$

$$E_{jet}(R) = \sum_i \int_R \omega_i f_i(\omega_i, k_{i\perp}^2) d\omega_i dk_{i\perp}^2$$

Full jet energy loss (radiative, collisional, broadening)



Nuclear modification of jet shape

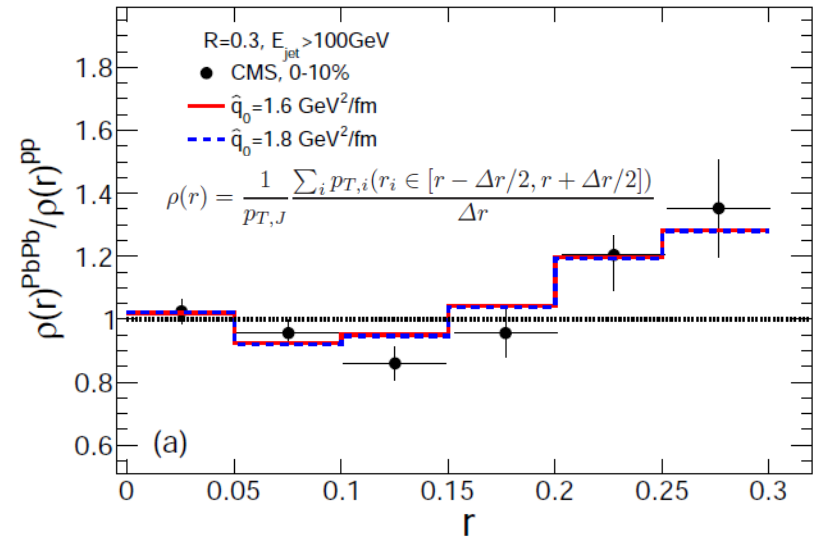
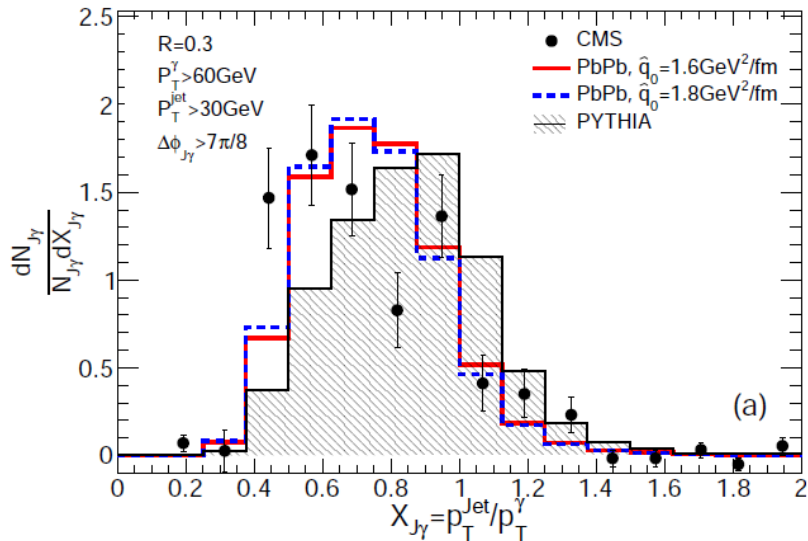
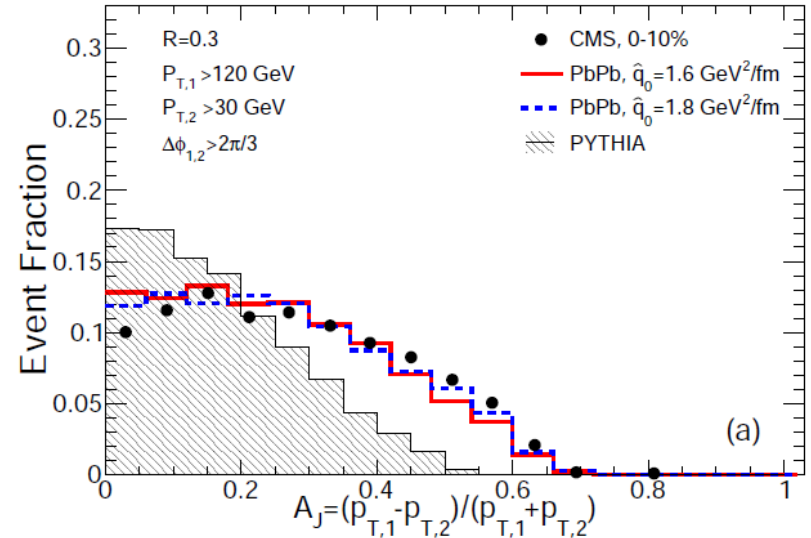
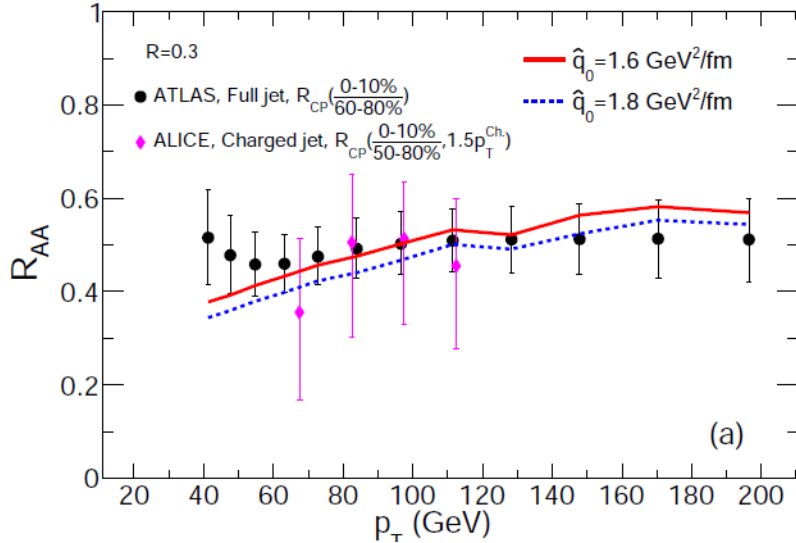


The soft outer part of jets is easier to be modified (some absorbed by medium), while the modification of the inner hard cone is more difficult

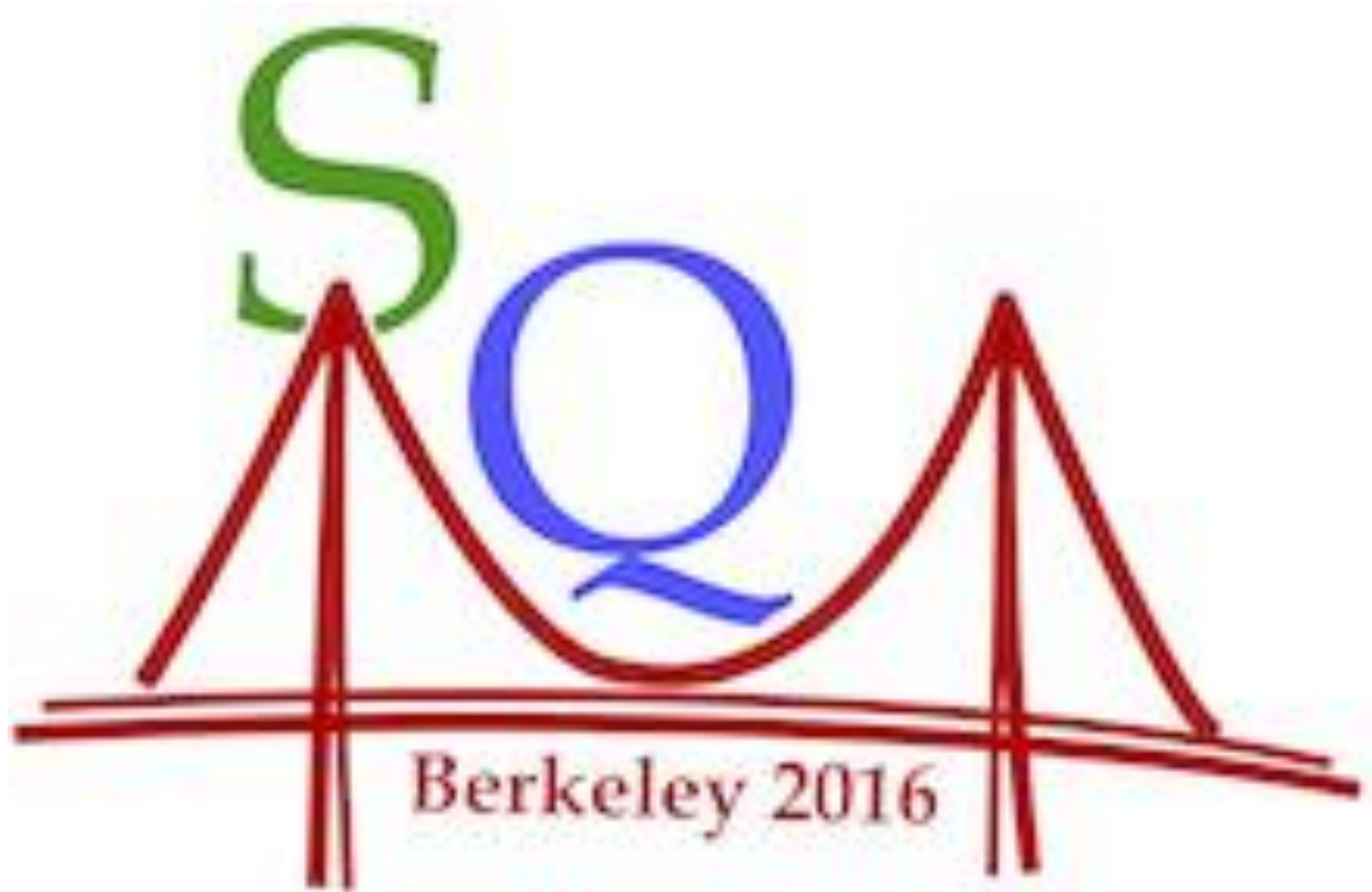
The enhancement at large r is consistent with the broadening

The final modification of jet shape comes from the interplay of different contributions

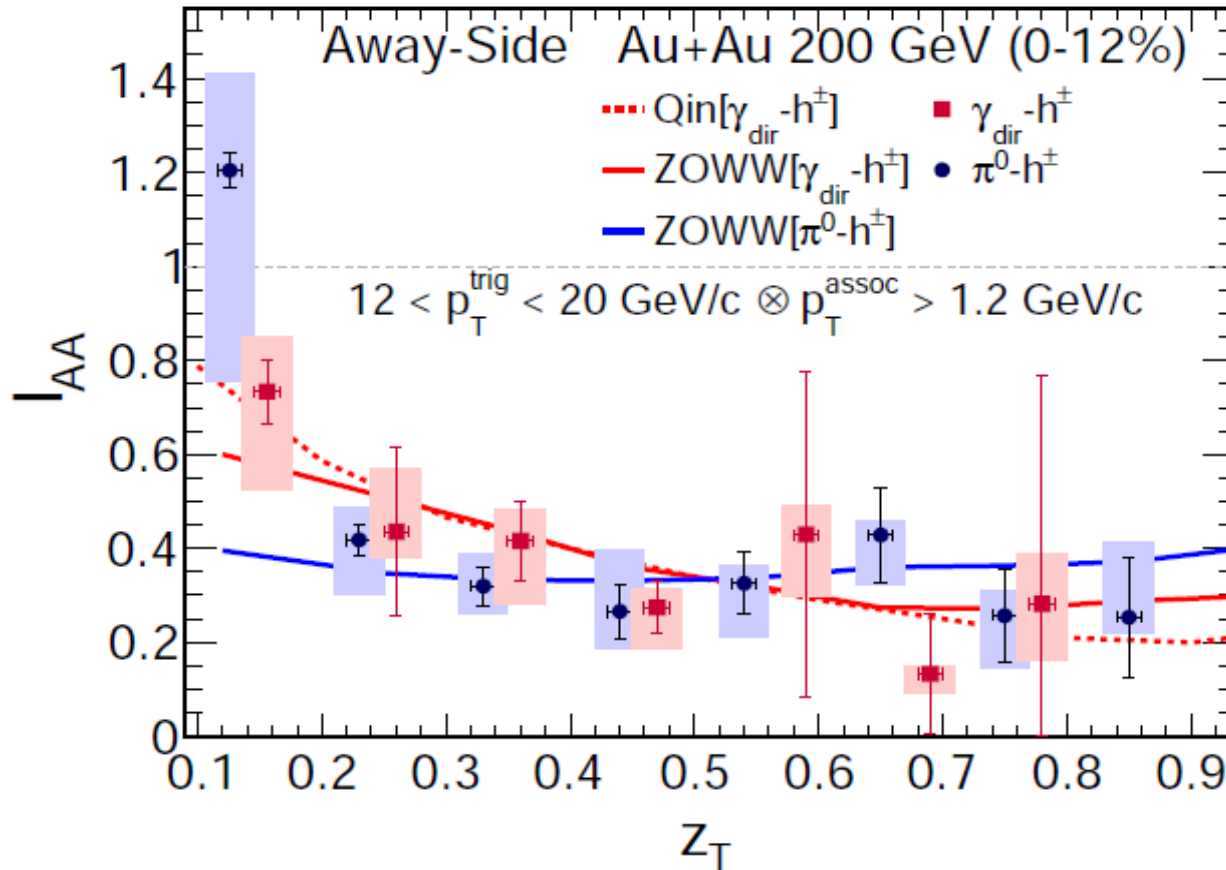
Various full jet observables



Jet-like correlations



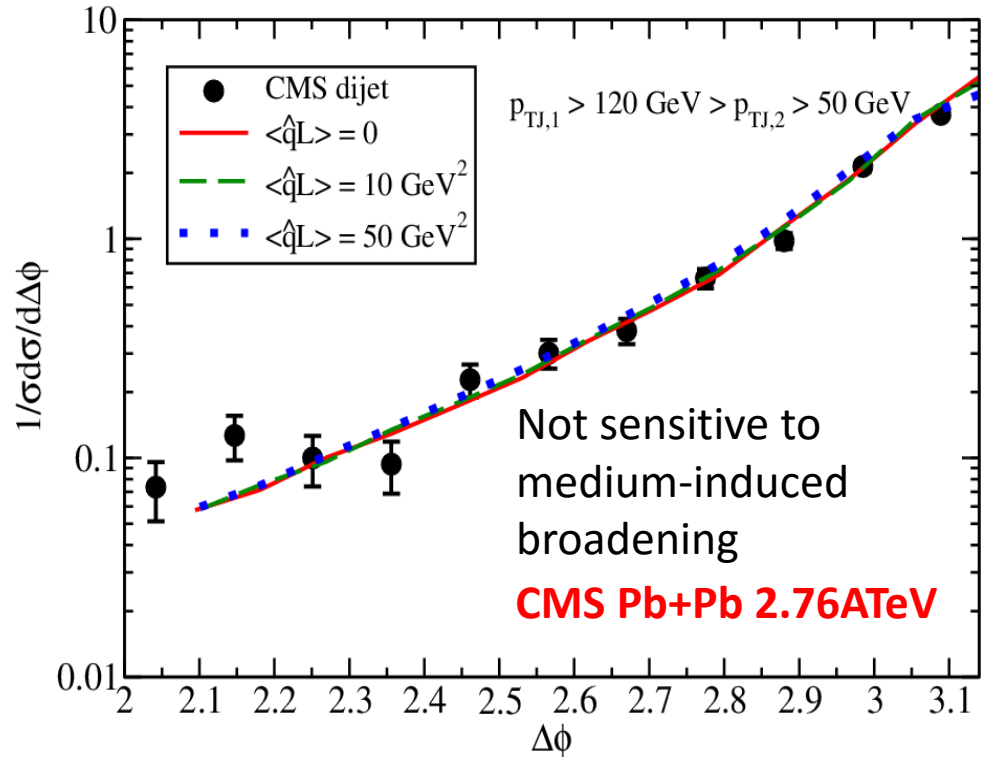
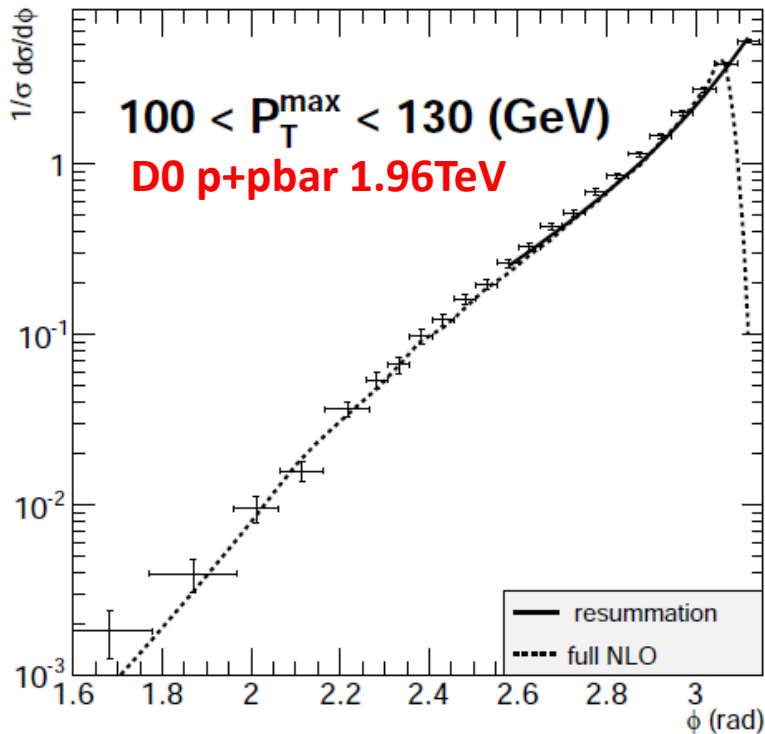
Jet-like correlations



Most of (theoretical) studies on jet-like correlations in AA collisions mainly focus on the nuclear modification of the (per-trigger) yield

We will use the angular correlations to probe the transverse momentum broadening

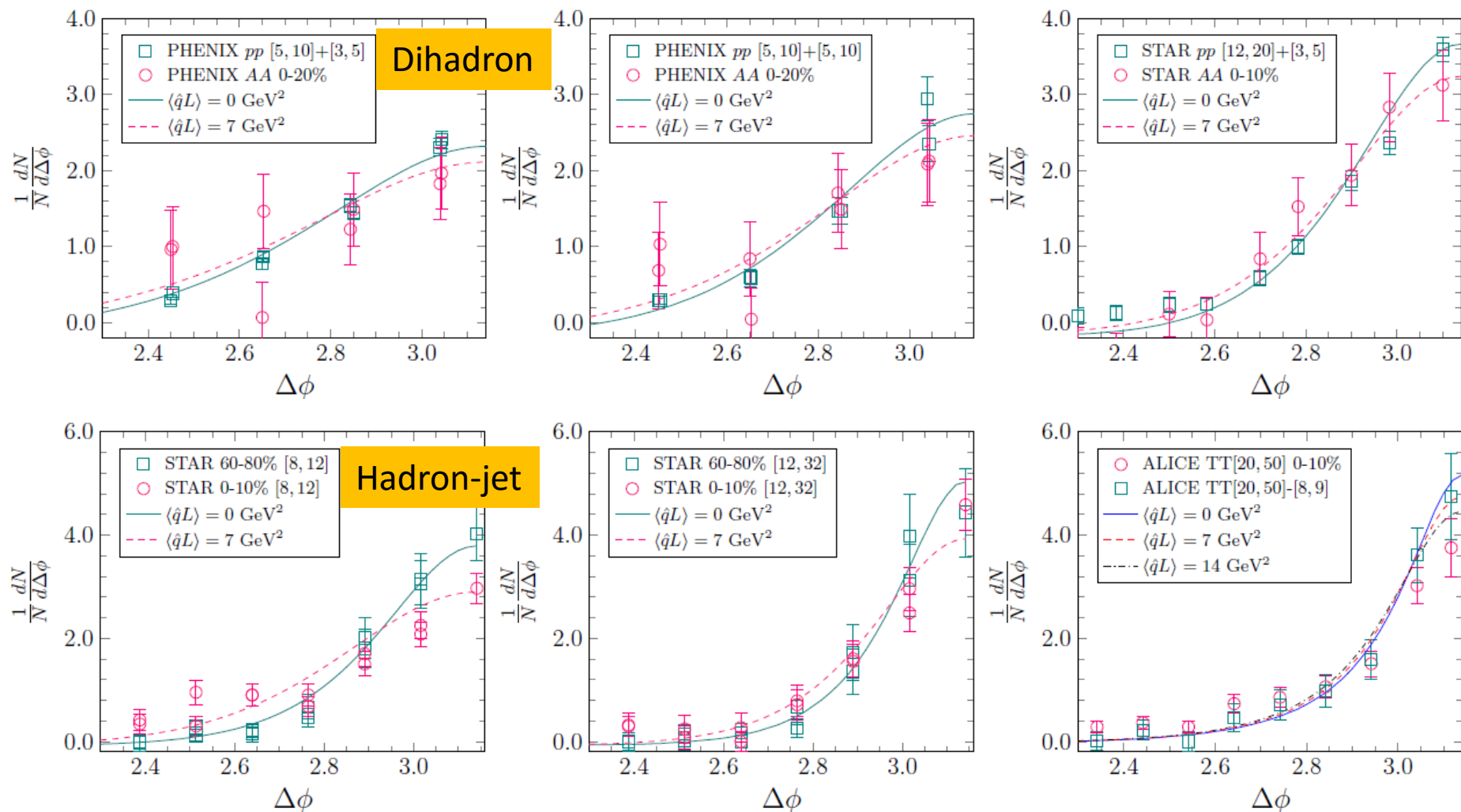
Dijet angular correlations in pp & AA



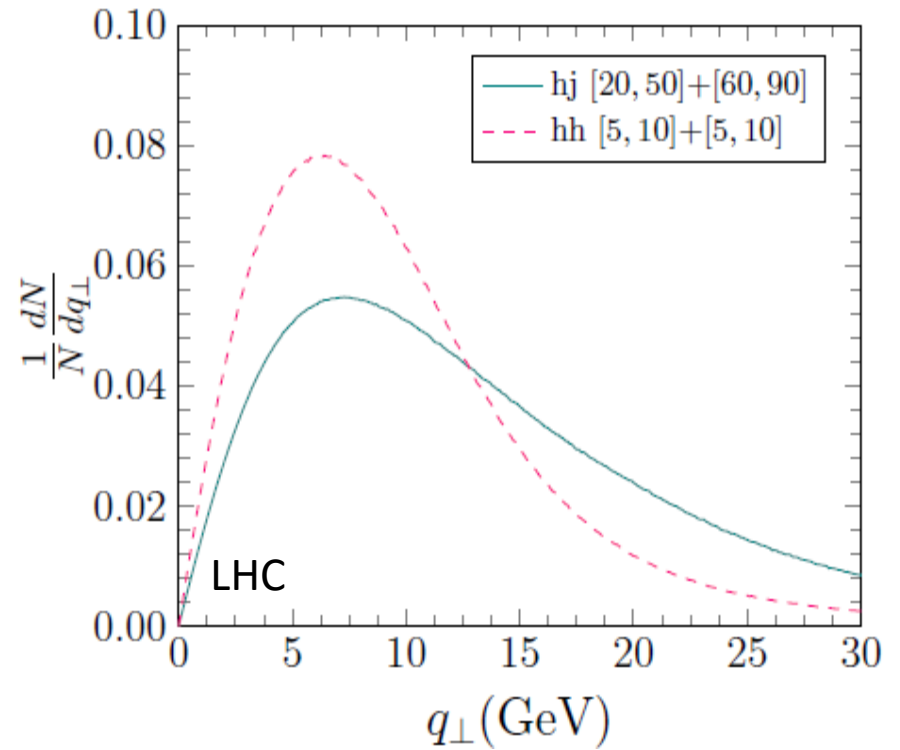
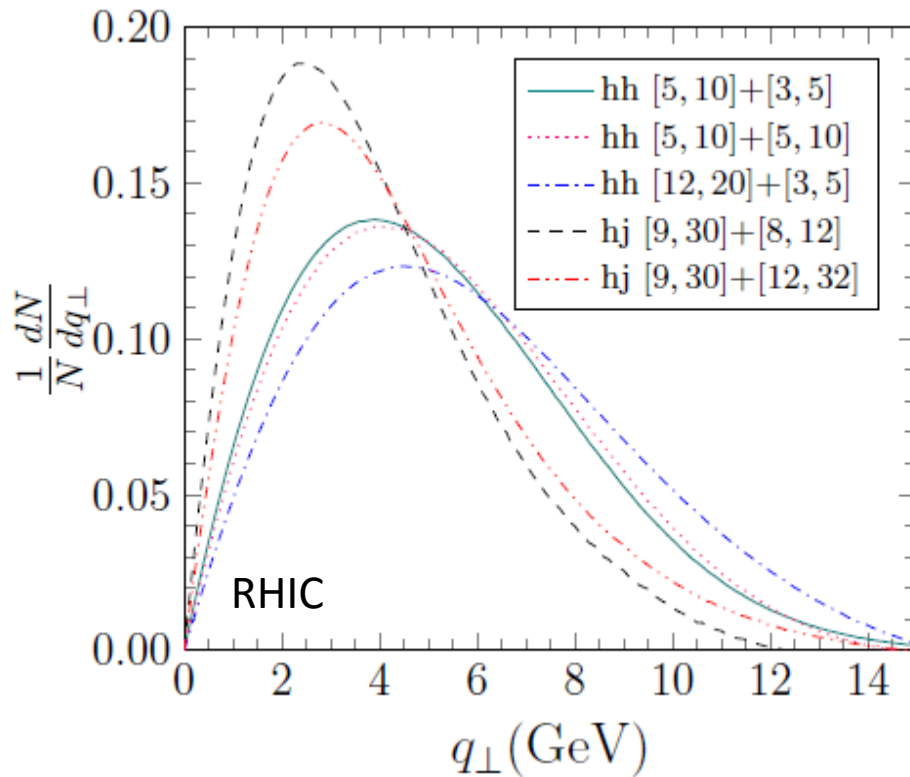
Resum all order soft gluon radiation in vacuum at NLL for dijet angular correlation by Sun, Yuan, Yuan, PRL 2014; PRD 2015

Extend the formalism to include the broadening effect induced by the QCD medium for dijet angular correlation by Mueller, Wu, Xiao, Yuan, arXiv:1604.04250

Probing q^{hat} via dihadron & hadron-jet angular correlations



Momentum imbalance q_T distribution (in pp)



$$\vec{q}_\perp = \vec{p}_{T,1} + \vec{p}_{T,2}$$

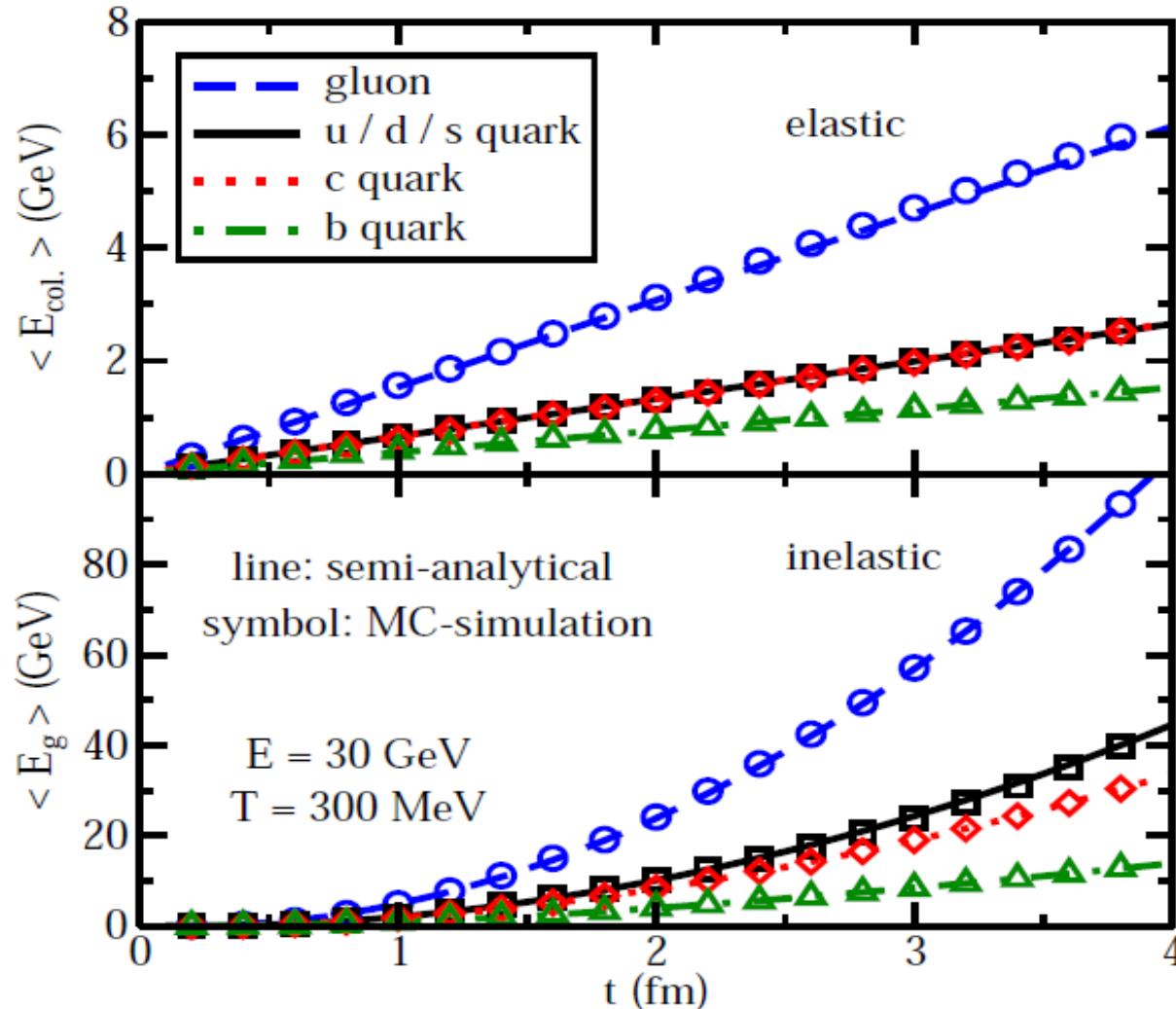
$$\langle q_\perp^2 \rangle_{AA} \approx \langle q_\perp^2 \rangle_{pp} + \langle \hat{q}L \rangle_{AA}$$

Summary

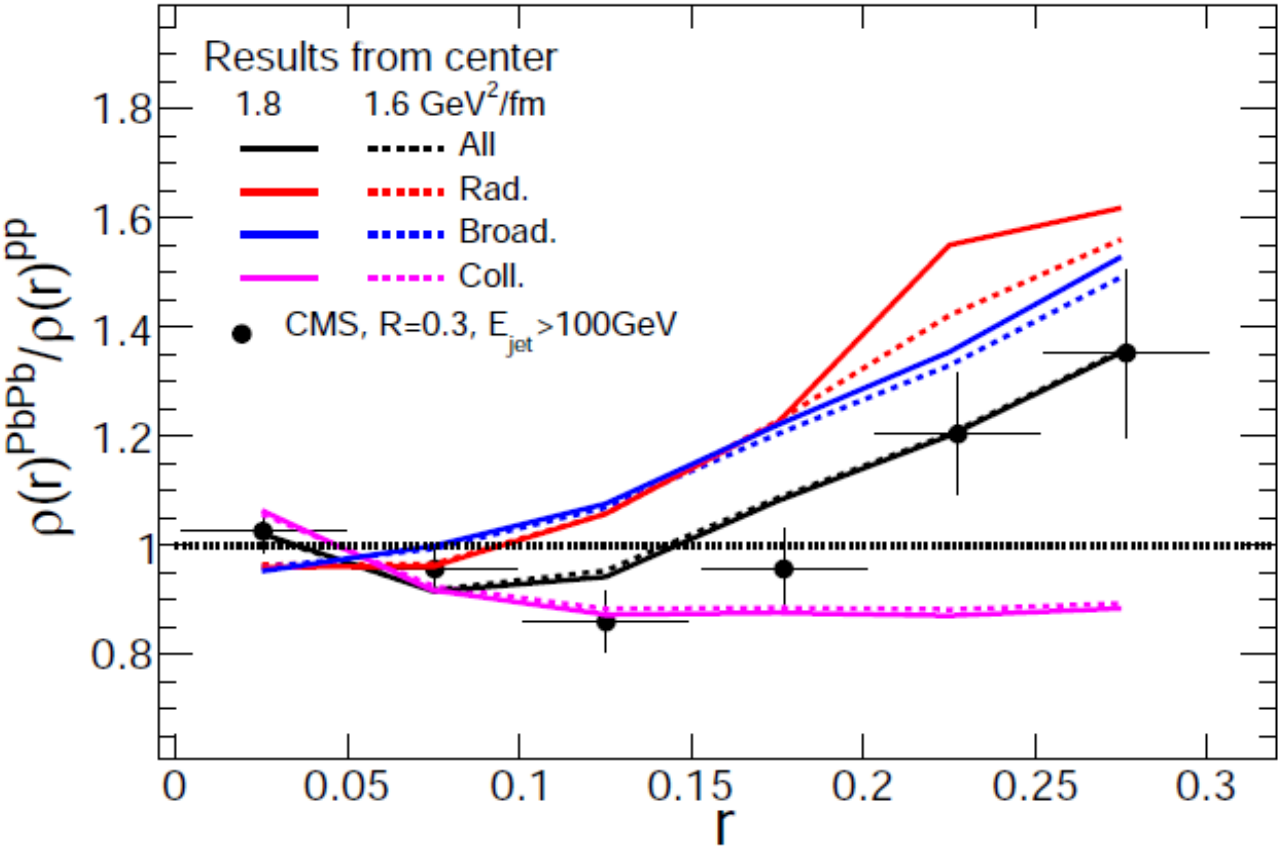
- ***Radiative & collisional* processes play different roles in different probes and observables**
 - Light & heavy flavor jet quenching, full jet energy loss, nuclear modification of jet shape
- **Jet transport coefficients control both *collisional* and *radiative* contributions**
- **Probe medium-induced *broadening* (q^{hat}) via jet-like angular correlations**

backup

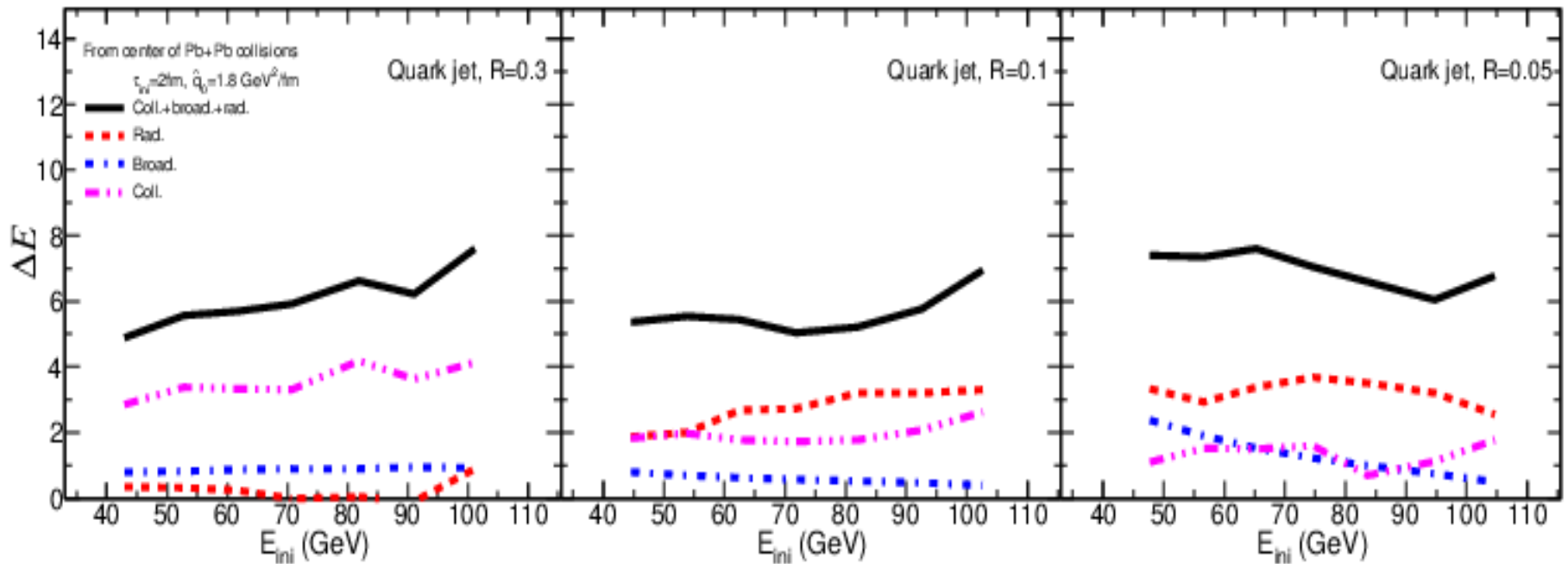
Test for collisional & radiative energy loss from LBT



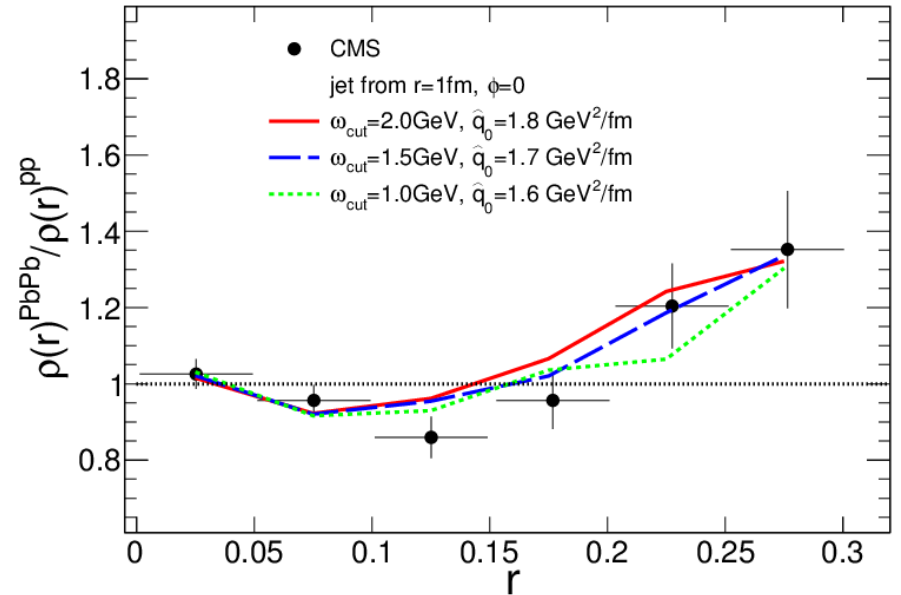
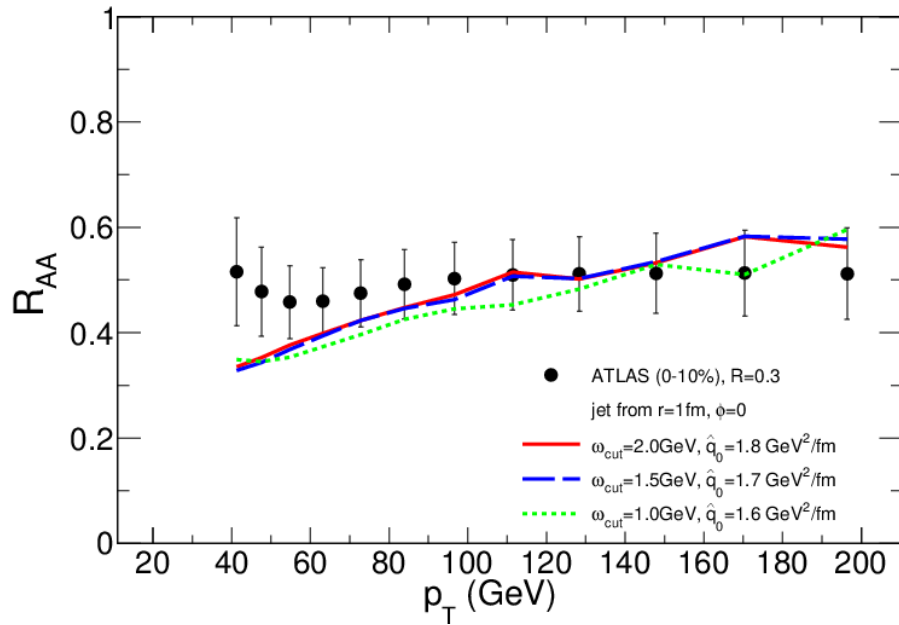
Sensitivity to jet transport parameter



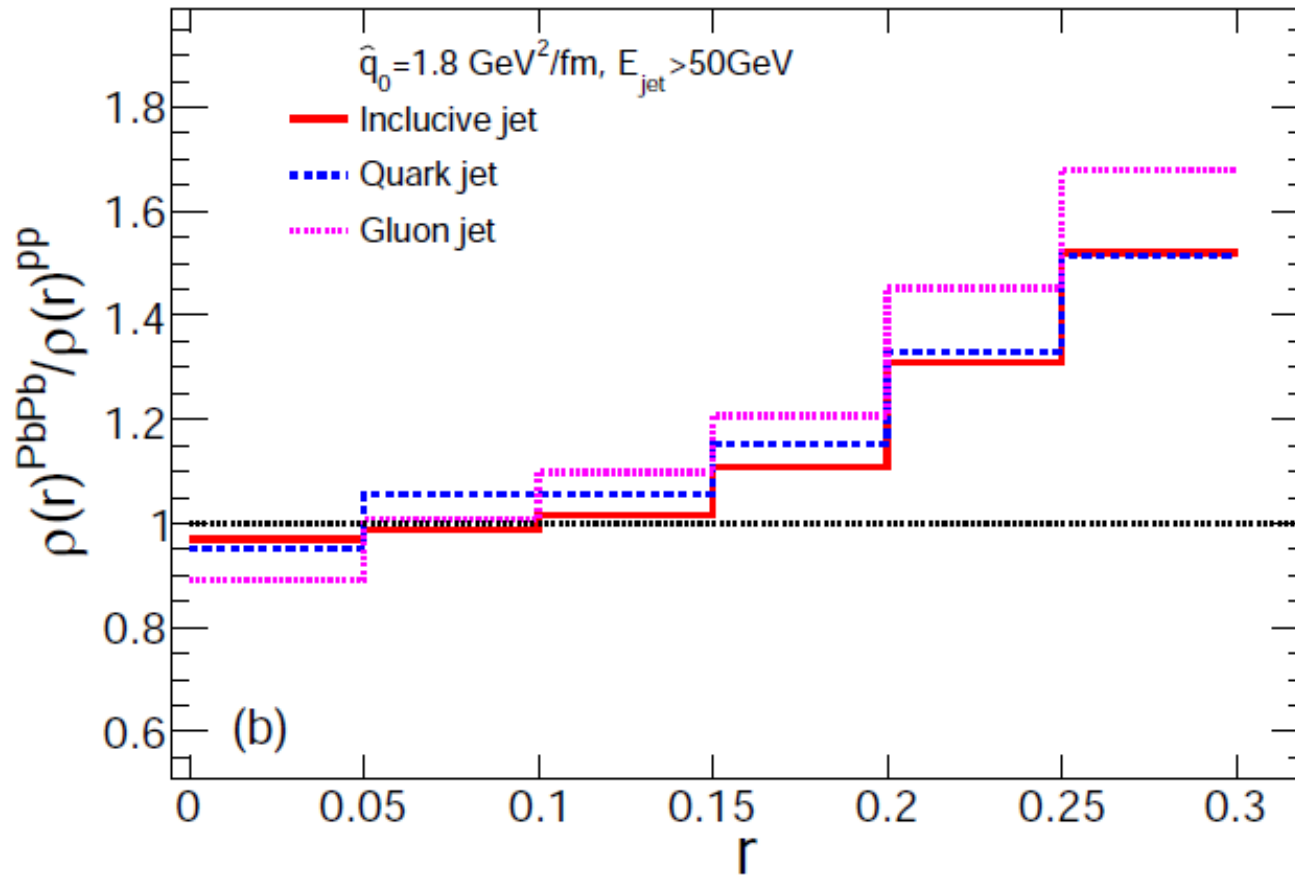
Full jet energy loss (jet size dependence of different contributions)



Uncertainty from ω_{cut}

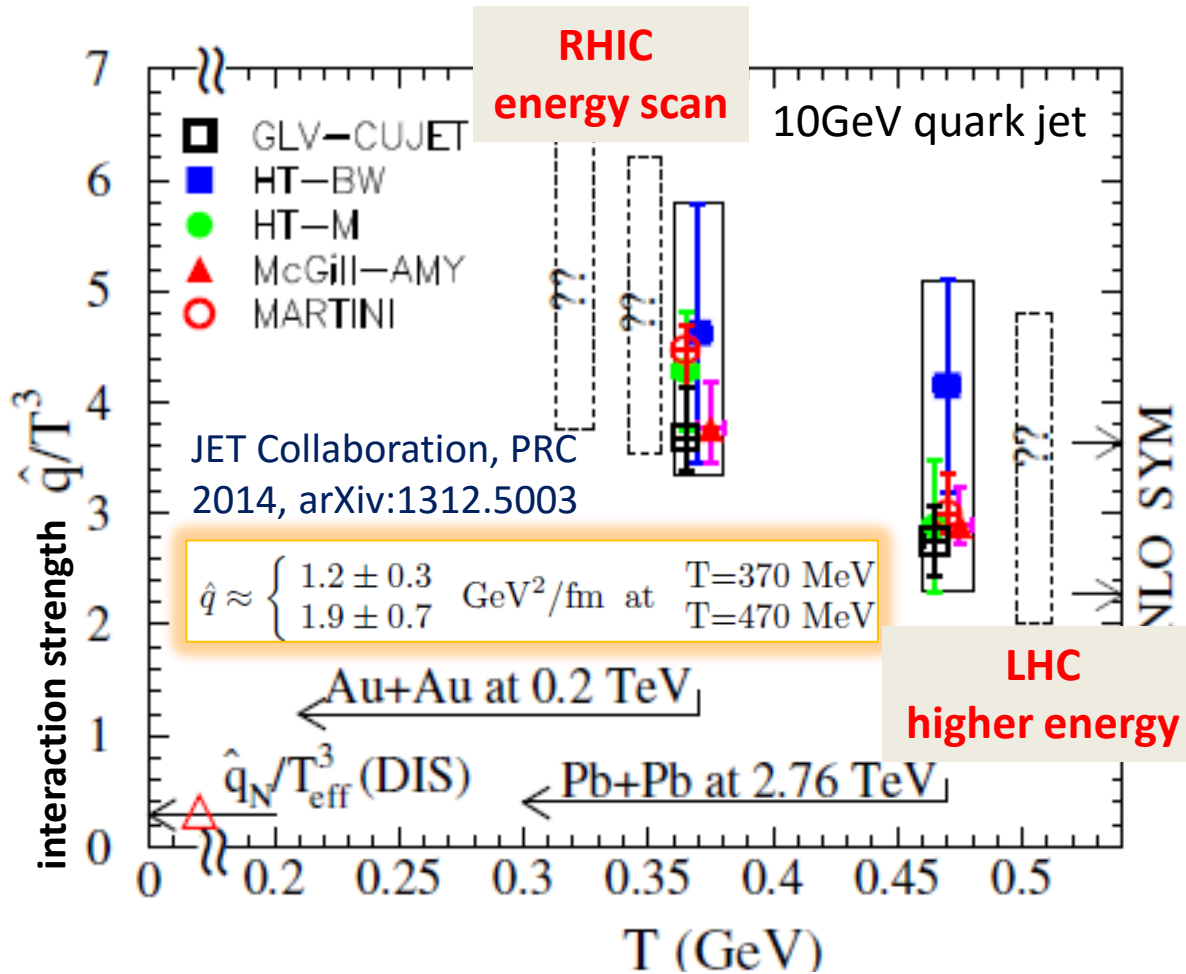


Nuclear modification of jet shape



Extraction of jet transport parameter

Jet transport coefficients control both collisional and radiative contributions



McGill-AMY:

GYQ, Ruppert, Gale, Jeon, Moore, Mustafa, PRL 2008

HT-BW:

Chen, Hirano, Wang, Wang, Zhang, PRC 2011

HT-M:

Majumder, Chun, PRL 2012

GLV-CUJET:

Xu, Buzzatti, Gyulassy, arXiv: 1402.2956

MARTINI-AMY:

Schenke, Gale, Jeon, PRC 2009

NLO SYM:

Zhang, Hou, Ren, JHEP 2013

$$\hat{q} = \frac{d\langle \Delta p_{\perp}^2 \rangle}{dt} = \int d^2 k_{\perp} k_{\perp}^2 \frac{d\Gamma(k_{\perp})}{d^2 k_{\perp} dt} \approx \frac{8\pi^2 \alpha_s C_R}{N_c^2 - 1} \int dy^- \langle F^{\mu+}(0) F_{\mu}^{+}(y^-) \rangle$$

Dihadron correlation

$$\begin{aligned} \frac{d\sigma}{d\Delta\phi} = & \sum_{a,b,c,d} \int p_T^{h_1} dp_T^{h_1} \int p_T^{h_2} dp_T^{h_2} \int \frac{dz_c}{z_c^2} \int \frac{dz_d}{z_d^2} \\ & \times \int b db J_0(q_\perp b) e^{-S(Q,b)} x_a f_a(x_a, \mu_b) x_b f_b(x_b, \mu_b) \\ & \times \frac{1}{\pi} \frac{d\sigma_{ab \rightarrow cd}}{d\hat{t}} D_c(z_c, \mu_b) D_d(z_d, \mu_b) \end{aligned}$$

$$\begin{aligned} S(Q, b) = & S_p^i(Q, b) + S_p^f(Q, b) + S_{\text{np}}(Q, b) \\ & + \frac{b^2}{4} (\langle \hat{q}_c L \rangle + \langle \hat{q}_d L \rangle), \end{aligned}$$