

# Sub-threshold strangeness and charm production (in UrQMD)

Jan Steinheimer

30.06.2016

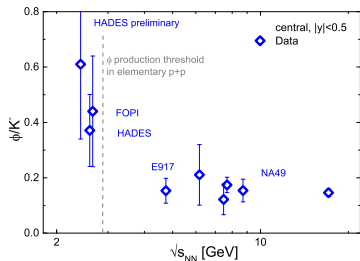


# Motivation

Recent measurements on near and below threshold production.

## $\phi$ production

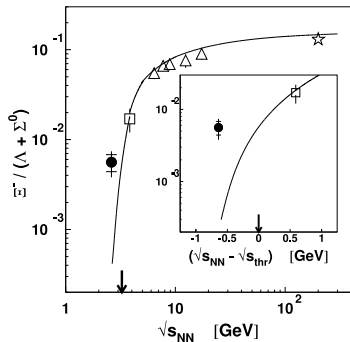
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G. Agakishiev *et al.* [HADES Collaboration], Phys. Rev. C **80**, 025209 (2009)

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$\Xi^-$  yield, measured in Ar+KCl much larger than thermal model.

Confirmed in p+Nb  $\rightarrow$  No Y+Y exchange!!

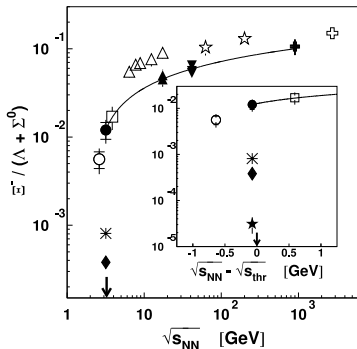
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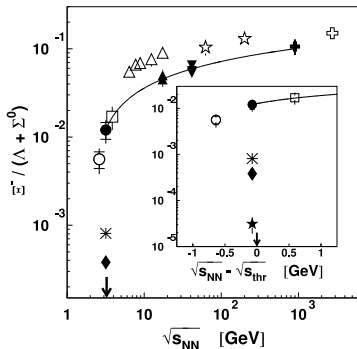
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Both particles are not well described in microscopic transport models and thermal fits are also not convincing.

G. Agakishiev *et al.* [HADES Collaboration], Phys. Rev. Lett. **103**, 132301 (2009)

# The notorious $\phi + N$ cross section

## Does the $\phi$ have a small hadronic cross section?

- The idea that the  $\phi$  has a small hadronic cross section is not new.  
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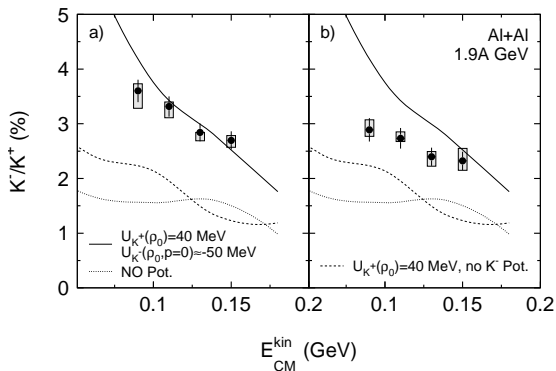
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- The  $\phi$  would be an important probe of hadronization.
- COSY and LEPS experiments have found large nuclear absorption cross sections

ANKE	SPring-8
14-21 mb	35 mb

M. Hartmann *et al.*, Phys. Rev. C **85**, 035206 (2012)

T. Ishikawa *et al.*, Phys. Lett. B **608**, 215 (2005)

# The Kaon-Nuclear potential



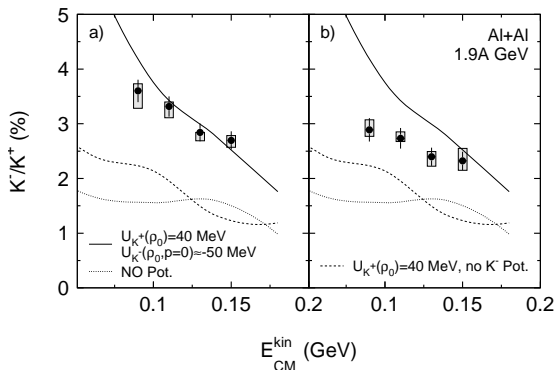
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- The  $K^-/K^+$  ratio is used to determine the Kaon nuclear potentials.



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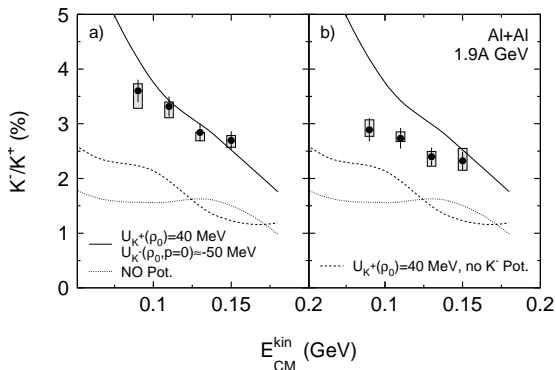


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## An example

- The  $K^-/K^+$  ratio is used to determine the Kaon nuclear potentials.
- Quantitative result relies on the baseline of non-potential case.
- $\phi$  contribution to the  $K^-$  found to be important.

# Why is a sub threshold charm prediction interesting?

## Charm at high baryon densities

- Study properties of charmed hadrons in dense nuclear matter.

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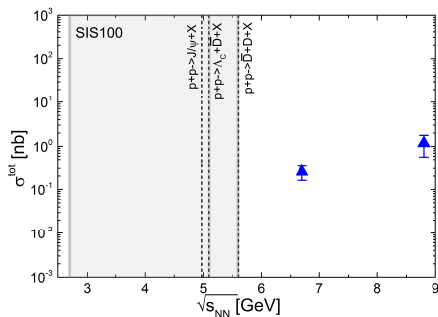
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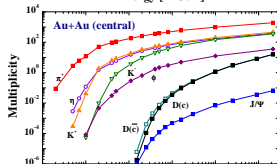
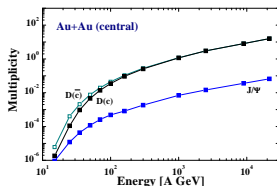
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HSD study: Based on parametrized cross section.

W. Cassing, E. L. Bratkovskaya and  
A. Sibirtsev,  
Nucl. Phys. A **691**, 753 (2001)

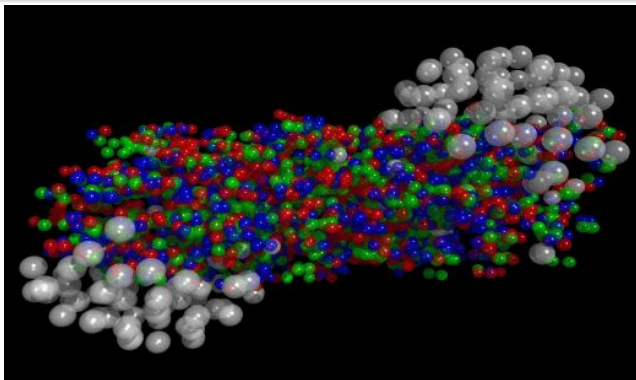




# Get a baseline

## UrQMD

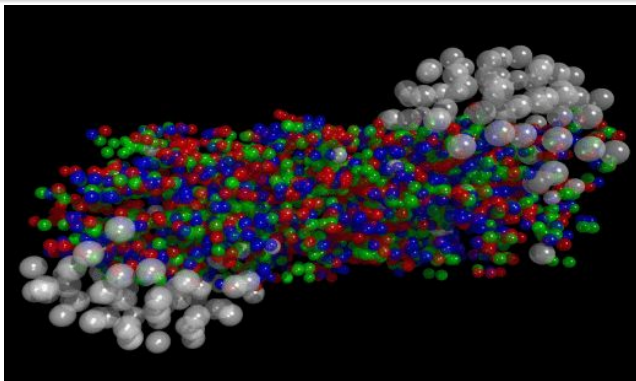
- We will use it in cascade mode.



# Get a baseline

## UrQMD

- We will use it in cascade mode.
- No long range interactions like potentials.



# Strangeness Production in UrQMD

UrQMD is a microscopic transport model

- Only  $2 \leftrightarrow 2$ ,  $2 \leftrightarrow 1$ ,  $2 \rightarrow N$  and  $1 \rightarrow N$  interactions allowed.

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## UrQMD is a microscopic transport model

- Only  $2 \leftrightarrow 2$ ,  $2 \leftrightarrow 1$ ,  $2 \rightarrow N$  and  $1 \rightarrow N$  interactions allowed.
- Resonance decays according to PDG values + guesstimates.
- Detailed balance. (Violated in string excitations, annihilations and some decays)

# Strangeness Production in UrQMD

Strange particle production goes ONLY via

Resonance excitation:

- $N+N \rightarrow X$
- $N+M \rightarrow X$
- $M+M \rightarrow X$

Relevant channels:

- 1  $NN \rightarrow N\Delta_{1232}$
- 2  $NN \rightarrow NN^*$
- 3  $NN \rightarrow N\Delta^*$
- 4  $NN \rightarrow \Delta_{1232}\Delta_{1232}$
- 5  $NN \rightarrow \Delta_{1232}N^*$
- 6  $NN \rightarrow \Delta_{1232}\Delta^*$
- 7  $NN \rightarrow R^*R^*$

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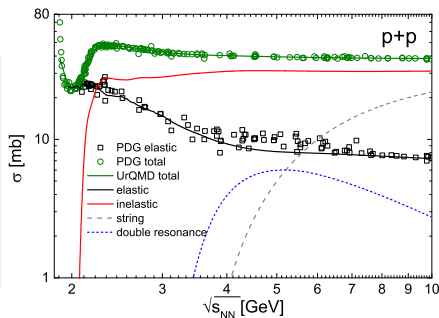
$N^*(1650)$	$\Delta(1232)$
$N^*(1710)$	$\Delta(1600)$
$N^*(1720)$	$\Delta(1620)$
$N^*(1875)$	$\Delta(1700)$
$N^*(1900)$	$\Delta(1900)$
$N^*(1990)$	$\Delta(1905)$
$N^*(2080)$	$\Delta(1910)$
$N^*(2190)$	$\Delta(1920)$
$N^*(2220)$	$\Delta(1930)$
$N^*(2250)$	$\Delta(1950)$
$N^*(2600)$	$\Delta(2440)$
$N^*(2700)$	$\Delta(2750)$
$N^*(3100)$	$\Delta(2950)$
$N^*(3500)$	$\Delta(3300)$
$N^*(3800)$	$\Delta(3500)$
$N^*(4200)$	$\Delta(4200)$

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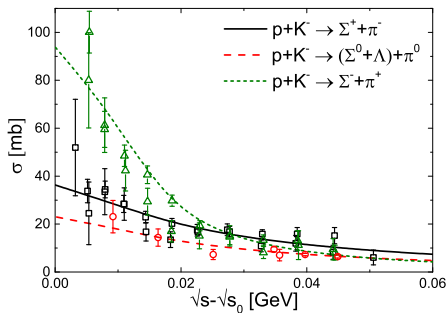
## N+N Cross section

Fixed to data where available. Otherwise fixed Matrix element + phase space.



# Strangeness exchange reactions

In addition Strange hadrons may be created in strangeness exchange reactions.



First the  $\phi$

# On the probability of sub threshold production

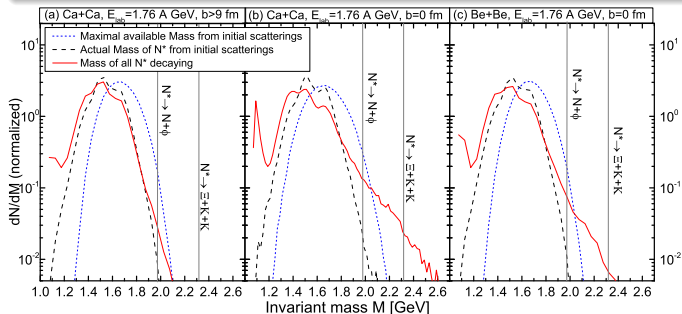
## Sub-threshold production baseline

- Fermi momenta lift the collision energy above the threshold.
- Secondary interactions accumulate energy.

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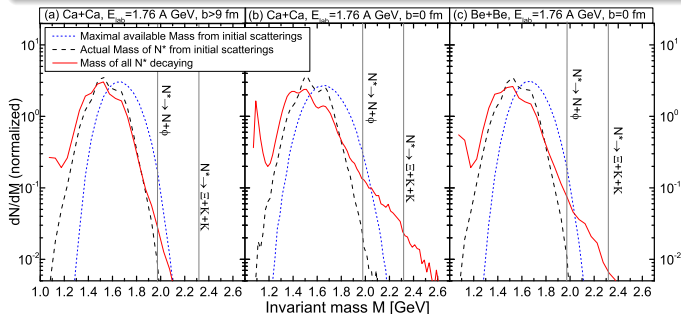
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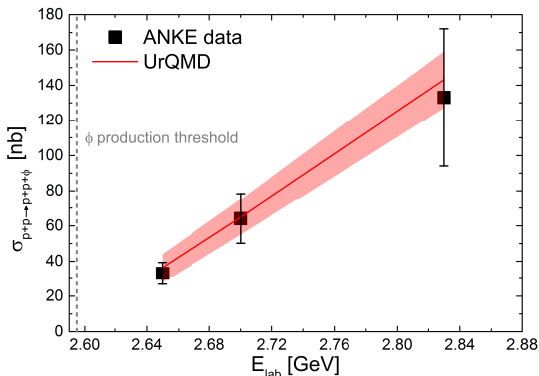
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Why not introduce these decays for the less known resonances?

## Fixing the $N^* \rightarrow \phi + N$ decay with p+p data

We use ANKE data on the  $\phi$  production cross section to fix the  $N^* \rightarrow N + \phi$  branching fraction.



Only 1 parameter

$$\Gamma_{N^* \rightarrow N\phi} / \Gamma_{\text{tot}} = 0.2\%$$

1 parameter fits all 3 points!

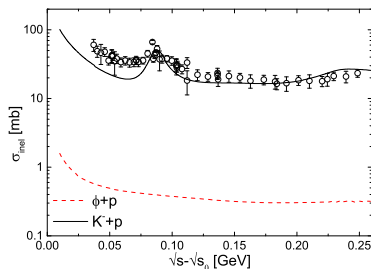
Y. Maeda *et al.* [ANKE Collaboration], Phys. Rev. C **77**, 015204 (2008) [arXiv:0710.1755 [nucl-ex]].

# $\phi$ suppression in nuclear medium

Detailed balance  $\rightarrow$  absorption cross section

$$\frac{d\sigma_{b \rightarrow a}}{d\Omega} = \frac{\langle p_a^2 \rangle (2S_1 + 1)(2S_2 + 1)}{\langle p_b^2 \rangle (2S_3 + 1)(2S_4 + 1)} \sum_{J=J_-}^{J_+} \frac{\langle j_1 m_1 j_2 m_2 || JM \rangle^2}{\langle j_3 m_3 j_4 m_4 || JM \rangle^2} \frac{d\sigma_{a \rightarrow b}}{d\Omega}$$

- $\phi + p$  cross section from detailed balance is very small.

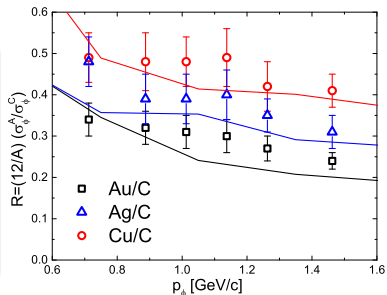


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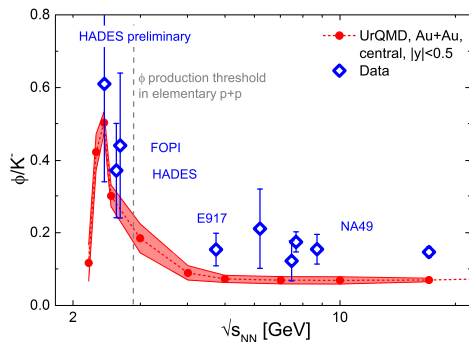
- $\phi + p$  cross section from detailed balance is very small.
- Still the transparency ratio is well reproduced. Remember: this is what lead to the 20 mb cross section from ANKE.
- Cross section from transparency ratio is model dependent!





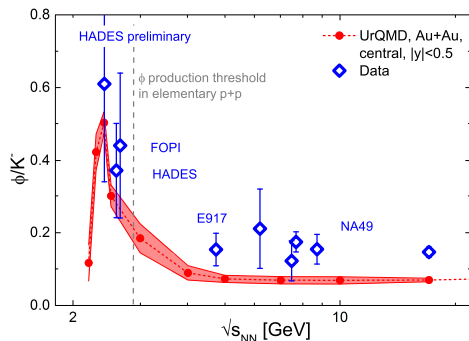
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When applied to nuclear collisions:



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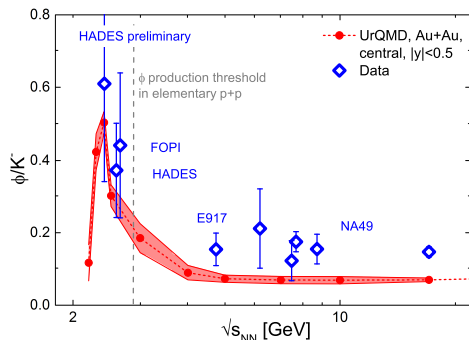
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# $\phi$ production in nuclear collisions below the p+p threshold

When applied to nuclear collisions:



- Qualitative behavior nicely reproduced
- Predicted maximum at 1.25 A GeV
- High energies: too low due to string production
- HADES preliminary results for 1.23 A GeV, see HADES talks by R. Holzmann and T. Scheib.

Even centrality dependence is very well reproduced: Signal for multi step processes.

## Kaon Potentials

- To constrain the Kaon potentials from kaon spectra one needs to understand the baseline
- For example the  $\phi$  contribution to the  $K^-$ .

# About the Kaon potential

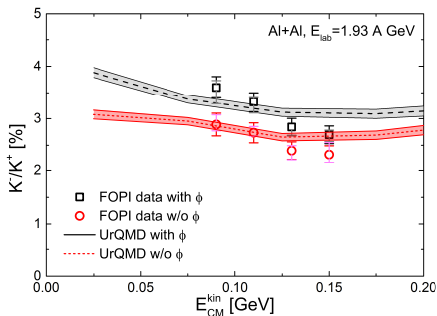
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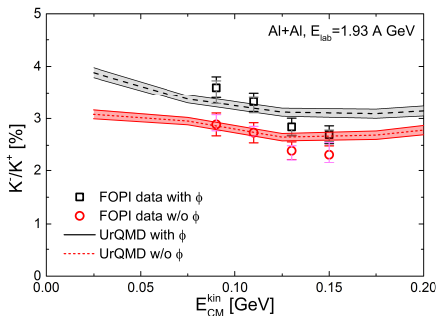
## UrQMD results

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## UrQMD results

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- With and without the  $\phi$  the ratio is much closer to the data already as in a comparable study with  $K^-$  potential.
- Can we make robust quantitative statements?

Now the  $\Xi$



## How to fix the $N^* \rightarrow \Xi^- + K + K$ decay?

No elementary measurements near threshold.

We use p+Nb at  $E_{\text{lab}} = 3.5$  GeV data  $\rightarrow \Gamma_{N^* \rightarrow \Xi + K + K} / \Gamma_{\text{tot}} = 3.0\%$

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HADES data	
$\langle \Xi^- \rangle$	$\Xi^- / \Lambda$
$(2.0 \pm 0.3 \pm 0.4) \times 10^{-4}$	$(1.2 \pm 0.3 \pm 0.4) \times 10^{-2}$

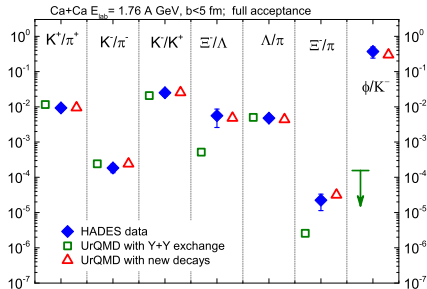
  

UrQMD	
$\langle \Xi^- \rangle$	$\Xi^- / \Lambda$
$(1.44 \pm 0.05) \times 10^{-4}$	$(0.71 \pm 0.03) \times 10^{-2}$

**Table:**  $\Xi^-$  production yield and  $\Xi^- / \Lambda$  ratio for minimum bias  $p + Nb$  collision at a beam energy of  $E_{\text{lab}} = 3.5$  GeV, compared with recent HADES results

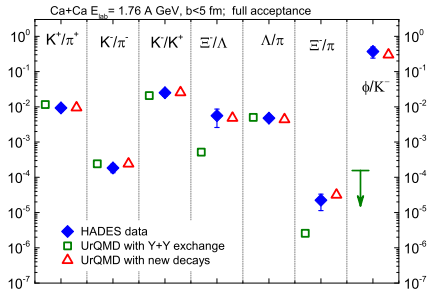
G. Agakishiev *et al.*, Phys.Rev.Lett. 114 (2015) no.21, 212301.

# $\Xi^-$ production in nuclear collisions below the p+p threshold



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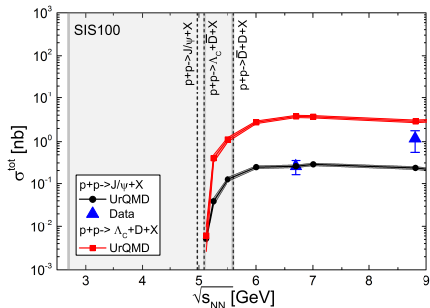
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- Indication for  $\Xi$  production from non-thermal 'tails' of particle production.
- All other strange particle ratios are also in line with experiment

# Can we make predictions about sub-threshold charm production?

J. Steinheimer, A. Botvina and M. Bleicher, arXiv:1605.03439 [nucl-th].

# Fixing the $N^* \rightarrow J/\Psi + N$ decay with p+p data

We use data from p+p at  $\sqrt{s} = 6.7$  GeV to fix the  $N^* \rightarrow N + J/\Psi$  branching fraction.



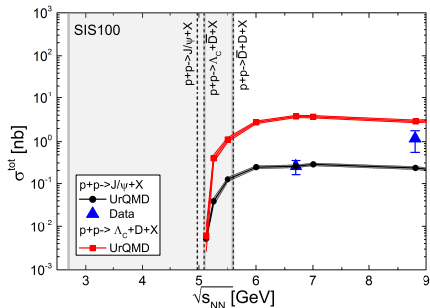
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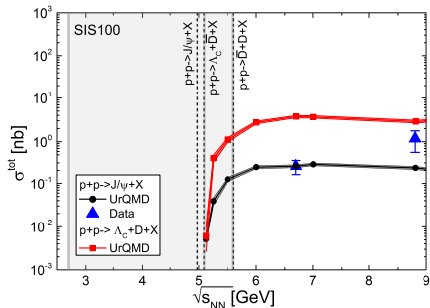
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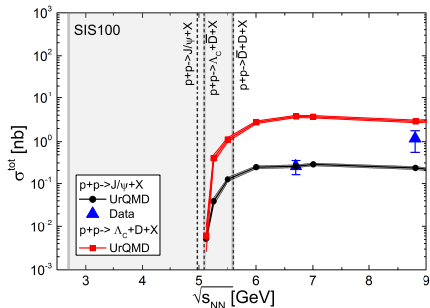
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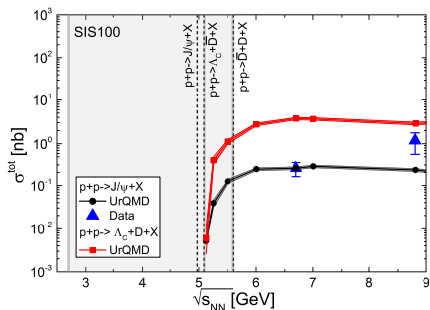
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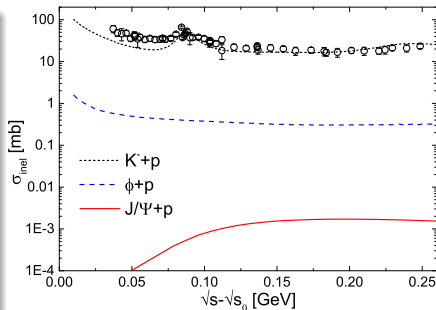
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- All the contributions should even increase the expected yield.

# $J/\Psi$ suppression in nuclear medium

## Detailed balance $\rightarrow$ absorption cross section

- $J/\Psi + p$  cross section from detailed balance is very small.

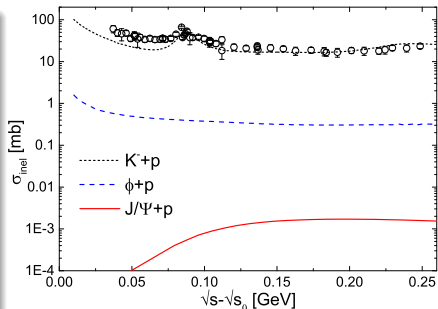


Comparable to: D. Kharzeev and H. Satz,  
Phys. Lett. B **334**, 155 (1994).

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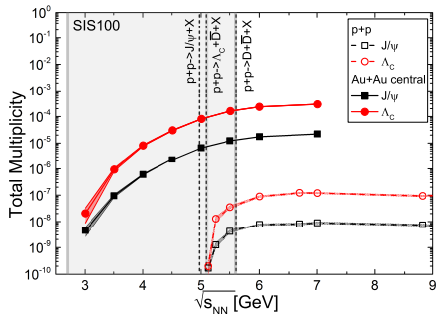
- $J/\Psi + p$  cross section from detailed balance is very small.
- Not 'absorption' of the  $J/\Psi$ , but of the mother resonance.
- Reactions of the type:  
$$N^* + N \rightarrow N'^* + N'^*$$
$$N^* + N \rightarrow N'^* + N'^*$$
where the mass of  $N'^* < N^*$  so no  $J/\Psi$  can be produced.



Comparable to: D. Kharzeev and H. Satz, Phys. Lett. B **334**, 155 (1994).

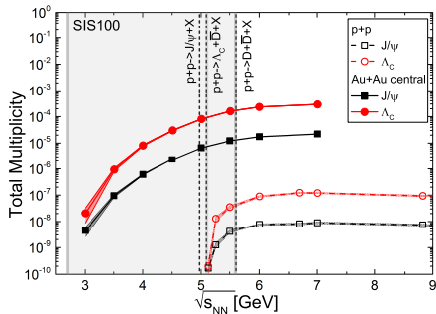
# $J/\Psi$ and open charm production in nuclear collisions below the p+p threshold

When applied to central nuclear collisions (min. bias: divide by 5):



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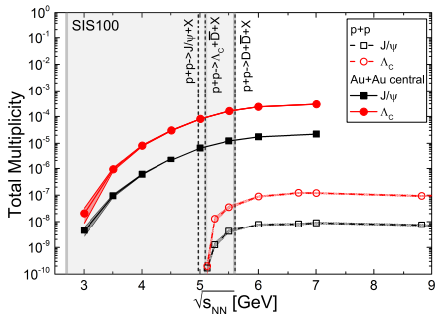


$E_{lab} = 6 \text{ A GeV}$

- $3 \cdot 10^{-7} J/\Psi$  per event
- $4 \cdot 10^{-6} \Lambda_c$  per event
- $\approx 6 - 8 \cdot 10^{-6} \bar{D}$  per event

# $J/\Psi$ and open charm production in nuclear collisions below the p+p threshold

When applied to central nuclear collisions (min. bias: divide by 5):



$E_{lab} = 11$  A GeV

- $1.5 \cdot 10^{-6}$   $J/\Psi$  per event
- $2 \cdot 10^{-5}$   $\Lambda_c$  per event
- $\approx 3 - 4 \cdot 10^{-5}$   $\bar{D}$  per event

# Fermi's Golden Rule, heavy resonances and equilibration

If resonance excitation and decay are governed mainly by phase space, does that help in equilibration of hadron yields?

Run UrQMD for SIS18 energies and fit FINAL particle yields at different time steps:

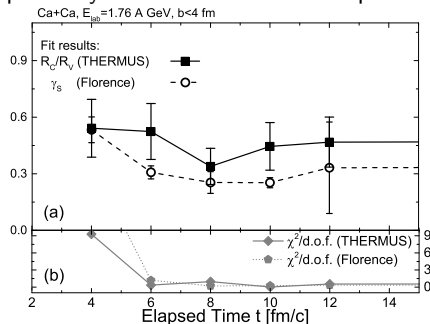
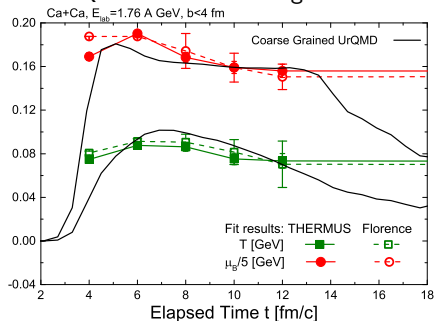
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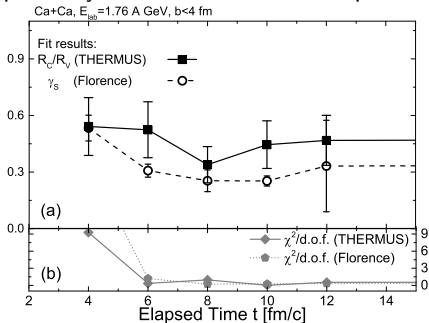
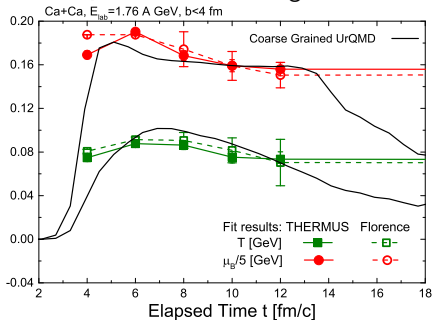


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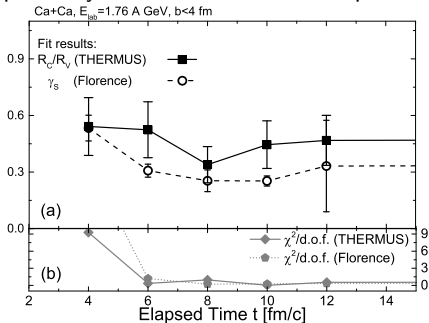
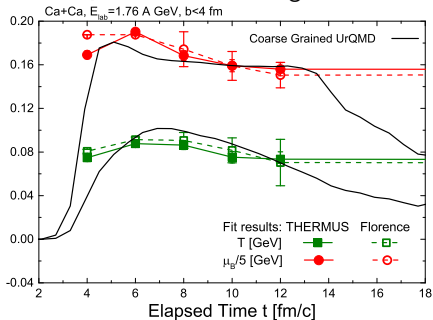


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Run UrQMD for SIS18 energies and fit FINAL particle yields at different time steps:



- Fit quality is good even after 2 fm/c.
- Extracted values of T and  $\mu_B$  consistent with coarse grained approach which assumes instant local equilibration.

# Summary

- We introduced a new mechanism of  $\phi$  and  $\Xi$  production in elementary and nuclear collisions, through the decay of heavy resonances.
- We can nicely describe the  $\phi$  and  $\Xi^-$  production in elementary and nuclear collisions near and below the  $\phi$  production threshold from elementary input.

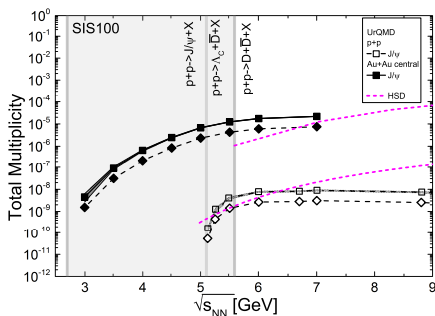
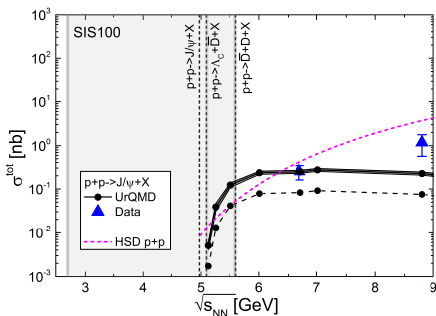
# Summary

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- We can nicely describe the  $\phi$  and  $\Xi^-$  production in elementary and nuclear collisions near and below the  $\phi$  production threshold from elementary input.
- We made predictions for sub threshold  $J/\Psi$ ,  $\Lambda_c$  and  $\bar{D}$  production in sub threshold collisions at the SIS100, observing a realistic chance of  $J/\Psi$  and open charm measurements.

# Comparisons to HSD

## Parametrized cross section for $J/\Psi$

$$\sigma_i^{NN}(s) = f_i a \left(1 - \frac{m_i}{\sqrt{s}}\right)^\alpha \left(\frac{\sqrt{s}}{m_i}\right)^\beta \theta(\sqrt{s} - \sqrt{s_{0i}})$$

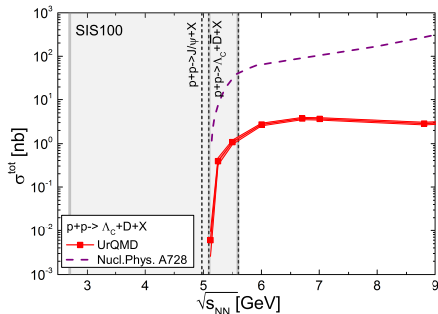


HSD results taken from:

O. Linnyk, E. L. Bratkovskaya and W. Cassing, Int. J. Mod. Phys. E **17**, 1367 (2008)

# Comparisons to hadronic Lagrangian

Cross section for  $p + p \rightarrow p + \bar{D}^0 + \Lambda_c$



Taken from:

W. Liu, C. M. Ko and S. H. Lee, Nucl. Phys. A **728**, 457 (2003)

The extracted cross sections depend on model assumptions

SPring-8

Used a Glauber model for the absorption.



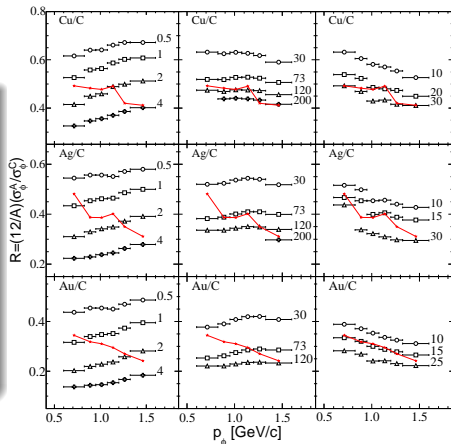
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## SPring-8

Used a Glauber model for the absorption.

## ANKE

- 1: The eikonal approximation of the Valencia group.
- 2: Paryev developed the spectral function approach for  $\phi$  production in both the primary proton- nucleon and secondary pion nucleon channels.
- BUU transport calculation of the Rossendorf group. Accounts for baryon baryon and meson baryon  $\phi$  production processes.

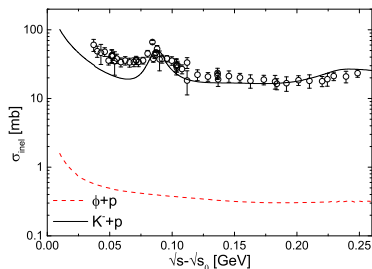


# $\phi$ suppression in nuclear medium

Detailed balance  $\rightarrow$  absorption cross section

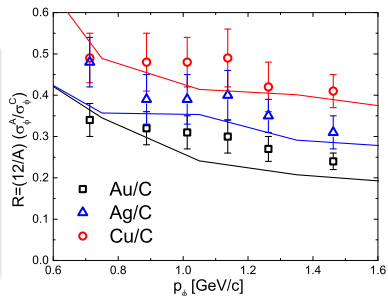
$$\frac{d\sigma_{b \rightarrow a}}{d\Omega} = \frac{\langle p_a^2 \rangle (2S_1 + 1)(2S_2 + 1)}{\langle p_b^2 \rangle (2S_3 + 1)(2S_4 + 1)} \sum_{J=J_-}^{J_+} \frac{\langle j_1 m_1 j_2 m_2 || JM \rangle^2}{\langle j_3 m_3 j_4 m_4 || JM \rangle^2} \frac{d\sigma_{a \rightarrow b}}{d\Omega}$$

- $\phi + p$  cross section from detailed balance is very small.



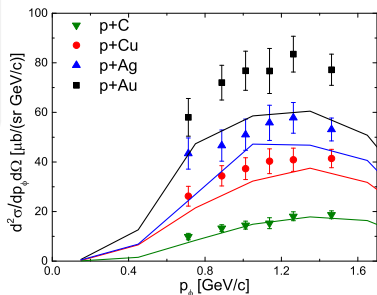
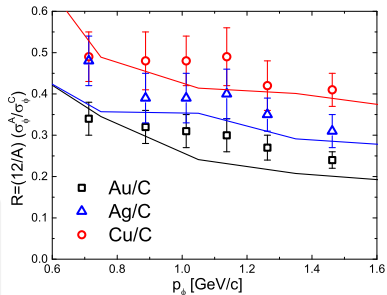
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- $\phi + p$  cross section from detailed balance is very small.
- Still the transparency ratio is well reproduced. Remember: this is what lead to the 20 mb cross section from ANKE.



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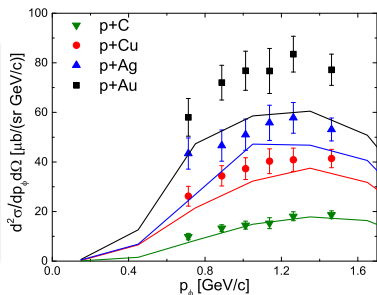
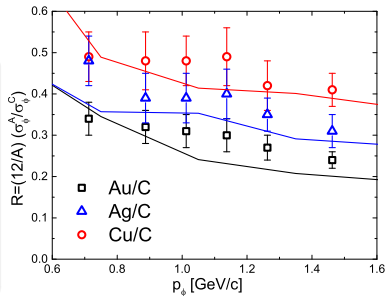
- $\phi + p$  cross section from detailed balance is very small.
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- Cross section from transparency ratio is model dependent!



# $\phi$ suppression in nuclear medium

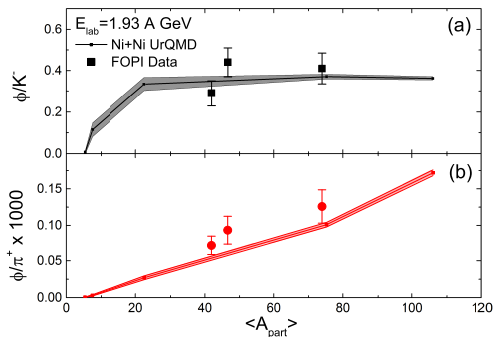
- $\phi + p$  cross section from detailed balance is very small.
- Still the transparency ratio is well reproduced. Remember: this is what lead to the 20 mb cross section from ANKE.
- Cross section from transparency ratio is model dependent!

- Not 'absorption' of the  $\phi$ , but of the mother resonance.
- Reactions of the type:  
$$N^* + N \rightarrow N'^* + N'^*$$
$$N^* + N \rightarrow N'^* + N'^*$$
where the mass of  $N'^* < N^*$  so no  $\phi$  can be produced.



# $\phi$ production in nuclear collisions below the p+p threshold

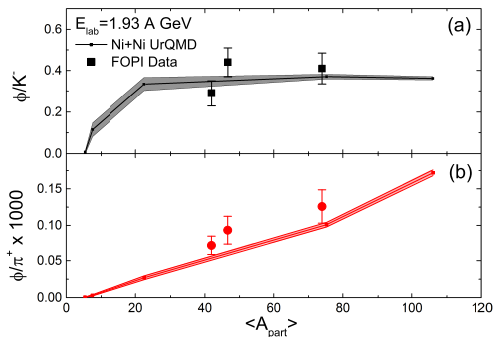
Even centrality dependence works well:



Data from: K. Piasecki et al., arXiv:1602.04378 [nucl-ex].

# $\phi$ production in nuclear collisions below the p+p threshold

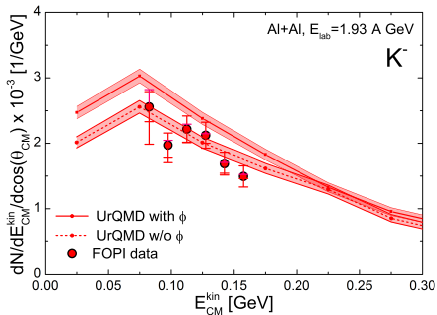
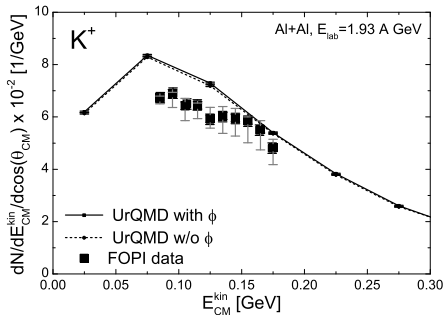
Even centrality dependence works well:



- Centrality dependence nicely reproduced.
- Good indicator for multi step production.

Data from: K. Piasecki et al., arXiv:1602.04378 [nucl-ex].

# Backup





# Strangeness Production in UrQMD

Strange particle production goes ONLY via

Resonance excitation:

- $N+N \rightarrow X$
- $N+M \rightarrow X$
- $M+M \rightarrow X$

Relevant channels:

- 1  $NN \rightarrow N\Delta_{1232}$
- 2  $NN \rightarrow NN^*$
- 3  $NN \rightarrow N\Delta^*$
- 4  $NN \rightarrow \Delta_{1232}\Delta_{1232}$
- 5  $NN \rightarrow \Delta_{1232}N^*$
- 6  $NN \rightarrow \Delta_{1232}\Delta^*$
- 7  $NN \rightarrow R^*R^*$

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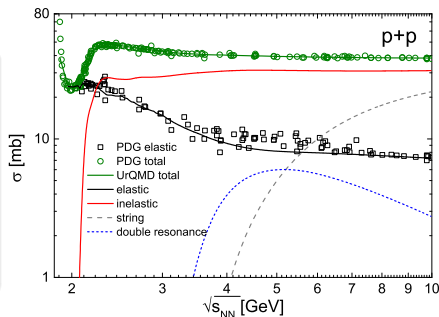
$N^*(1650)$	$\Delta(1232)$
$N^*(1710)$	$\Delta(1600)$
$N^*(1720)$	$\Delta(1620)$
$N^*(1875)$	$\Delta(1700)$
$N^*(1900)$	$\Delta(1900)$
$N^*(1990)$	$\Delta(1905)$
$N^*(2080)$	$\Delta(1910)$
$N^*(2190)$	$\Delta(1920)$
$N^*(2220)$	$\Delta(1930)$
$N^*(2250)$	$\Delta(1950)$
$N^*(2600)$	$\Delta(2440)$
$N^*(2700)$	$\Delta(2750)$
$N^*(3100)$	$\Delta(2950)$
$N^*(3500)$	$\Delta(3300)$
$N^*(3800)$	$\Delta(3500)$
$N^*(4200)$	$\Delta(4200)$

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Strange particle production goes ONLY via

Resonance excitation:

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- $N+M \rightarrow X$
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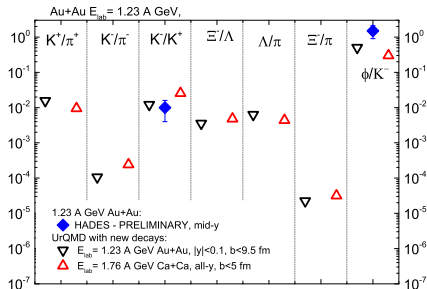
## N+N Cross section

$$\sigma_{1,2 \rightarrow 3,4}(\sqrt{s}) \propto (2S_3 + 1)(2S_4 + 1) \frac{\langle p_{3,4} \rangle}{\langle p_{1,2} \rangle} |M(m_3, m_4)|^2$$

with

$$|M(m_3, m_4)|^2 = \frac{A}{(m_4 - m_3)^2 (m_4 + m_3)^2}$$

# Predictions for Au+Au at $E_{\text{lab}} = 1.23$ A GeV



$\Xi^-/\Lambda$  does not decrease much.