

# Collectivity and thermalization in $p+p$ , $p+A$ , $d+Au$ and ${}^3\text{He}+Au$ collisions

**Shengli Huang**

Department of Physics and Astronomy, Station B 1807, Vanderbilt University, Nashville, TN 37235, USA

E-mail: shengli.huang@vanderbilt.edu

**Abstract.** In this proceeding, the recent measurements of the anisotropy ( $v_2$ ,  $v_3$ ) in small collision systems such as  $p+p$ ,  $p+A$ ,  $d+Au$  and  ${}^3\text{He}+Au$  collisions at the LHC and RHIC energies are presented. In  $p+p$  collisions, the  $v_2$  are measured by two and multi-particle  $\Delta\phi$  correlations at  $\sqrt{s_{NN}} = 2.76, 5, 7, \text{ and } 13$  TeV. In 200 GeV  $p+Au$ ,  $d+Au$  and  ${}^3\text{He}+Au$  collisions, the  $v_2$  are measured with event plane methods. I also present the measurements of  $v_2$  for  $K_S^0$  and  $\Lambda/\bar{\Lambda}$  in high multiplicity  $p+p$  collisions, in which a particle mass ordering is observed. The  $v_2$  of identified  $\pi^\pm$ ,  $K^\pm$  and (anti)protons in central  ${}^3\text{He}+Au$  collisions are also measured and quark-number scaling is observed. These observations are similar to those seen in A+A collisions, and support the interpretation of a collective origin.

## 1. Introduction

Recently, in the high-multiplicity small collision systems such as  $p+p$  and  $p/d/{}^3\text{He}+A$  collisions, the near-side, long-range two particle angular correlation (so called 'ridge') and anisotropic flow ( $v_n$ ) have been observed at the LHC [1, 2, 5, 3, 4] and RHIC [6, 7, 8]. Several physics models, which include initial-state gluon saturation [9], and hydrodynamic flow from a mini-QGP (Quark Gluon Plasma) [10] are potential explanations for these observations. While the former will not depend on the initial geometry, the hydrodynamic models have strong initial geometry dependence. Therefore, measuring the  $v_2$  and  $v_3$  in  $p/d/{}^3\text{He}+Au$  collisions will provide direct testing for these explanations since the initial geometry is quite different for these collisions [11]. If a mini-QGP exists in the small collision system, studying it will help us to further understand: 1) the condition for the QGP thermalization; 2) the role of the internal structure of proton on the initial geometry eccentricity [12]; and 3) the contributions from pre-equilibrium stage to the system evolution [13].

## 2. The $v_2$ measurements in $p+Au$ , $d+Au$ and ${}^3\text{He}+Au$ collisions at $\sqrt{s_{NN}} = 200$ GeV at RHIC

The  $v_2$  for inclusive charged hadrons produced at mid-rapidity  $|\eta| < 0.35$  are measured in 0-5% central  $p+Au$  in the PHENIX experiment and compared with that from  $d+Au$  and  ${}^3\text{He}+Au$  collisions, as shown in the Fig. 1. The  $v_2$  in the  $p+Au$  is found to be smaller than that of  $d+Au$  and  ${}^3\text{He}+Au$ . The eccentricity of  $p+Au$  is also smaller than that of  $d+Au$  and  ${}^3\text{He}+Au$ . It indicates the anisotropy in the small collisions system depends on the initial geometry, which is supported by the hydrodynamic model calculation [13] as shown in the figure.

The  $v_3$  for inclusive charged hadrons produced at mid-rapidity  $|\eta| < 0.35$  in high-multiplicity  ${}^3\text{He}+\text{Au}$  collisions at  $\sqrt{s_{NN}} = 200$  GeV are also measured in the PHENIX experiment with respect to the  $\Psi_3$  event planes. The  $v_3$  results are shown in Fig. 2 with  $v_2$  and different theory calculations. Four such predictions shown in Fig. 2 employ viscous hydrodynamics with  $\eta/s$  at or near the conjectured lower bound  $1/4\pi$  [15] and one is from the AMPT (A-Multi-Phase-Transport-Model) framework [16]. The SONIC calculation [11] employs Glauber initial conditions, viscous hydrodynamics, and then at  $T = 170$  MeV a transition to a hadronic cascade. The SUPERSONIC calculation [13] additionally includes pre-equilibrium dynamics that boosts the initial velocity fields at the earliest times. The impact of pre-equilibrium is modest on the  $v_2$  values and the data agree with both calculations within uncertainties. The effect of pre-equilibrium on  $v_3$  is significantly larger as the triangular flow takes longer to develop [11]. The SUPERSONIC prediction agrees well with the experimental data for  $p_T < 1.2$  GeV/ $c$ , and then the data trends towards the SONIC prediction at higher  $p_T$ .

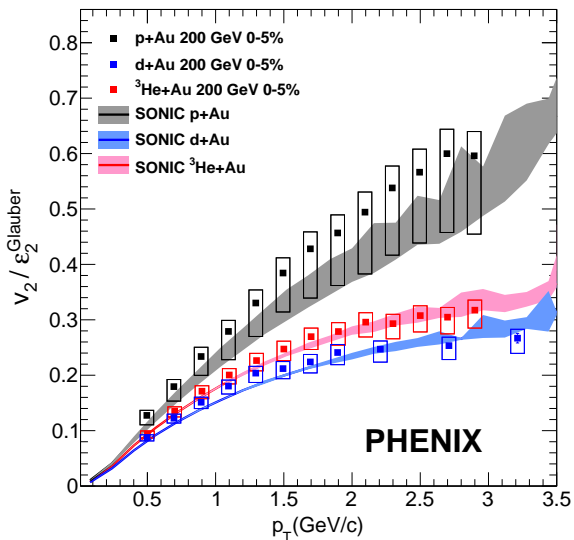


Figure 1: Measured  $v_2$  for mid-rapidity charged tracks in 0-5% central  $p+\text{Au}$ ,  $d+\text{Au}$  and  ${}^3\text{He}+\text{Au}$  at  $\sqrt{s_{NN}} = 200$  GeV using the event plane method. Also shown are SONIC model calculations [13]

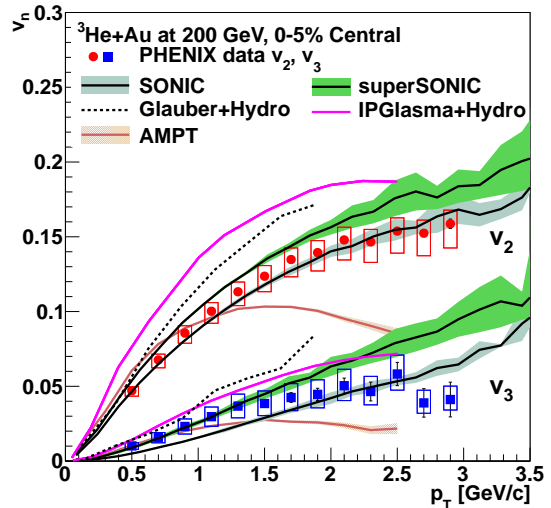


Figure 2:  $v_2$  (circles) and  $v_3$  (squares) as a function of  $p_T$  for inclusive charged hadrons at mid-rapidity in 0-5% central  ${}^3\text{He}+\text{Au}$  collisions at  $\sqrt{s_{NN}} = 200$  GeV; Also shown are various theoretical calculations.

### 3. The $v_2$ measurements in $p+p$ collisions at $\sqrt{s_{NN}} = 2.76, 5, 7,$ and $13$ TeV at the LHC

The  $v_2$  in  $p+p$  collisions at  $\sqrt{s_{NN}} = 2.76$  and  $13$  TeV has been measured with long-range ( $|\Delta\eta| > 2$ ) two particle  $\Delta\phi$  correlations by ATLAS [20]. A new method, called template fitting has been developed to separate the ridge and the back-to-back jet correlation in this long-rang two particles correlation, by assuming that the shape of the jet-induced correlations is invariant with event multiplicity and can be extracted from low-multiplicity events. The second order harmonic coefficients  $v_{2,2}$  from template fitting and the  $v_2$  extracted by assuming factorization are shown in the Fig 3. Both  $v_{2,2}$  and  $v_2$  show a very week multiplicity dependence and are nearly identical for collision energy from 2.76 to 7 TeV.

The anisotropy of charged particles are also measured in  $p+p$  collisions at 13 TeV by CMS [17]. A different subtraction method [4] is employed by CMS. The second-order Fourier coefficients

$V_{2\Delta}$  extracted from long-range two-particle  $\Delta\phi$  correlations in the higher-multiplicity region are subtracted from the  $V_{2\Delta}$  coefficients from  $10 \leq N_{\text{trk}}^{\text{offline}} < 20$  with

$$V_{2\Delta}^{\text{sub}} = V_{2\Delta} - V_{2\Delta}(10 \leq N_{\text{trk}}^{\text{offline}} < 20) \frac{N_{\text{assoc}}(10 \leq N_{\text{trk}}^{\text{offline}} < 20)}{N_{\text{assoc}}} \frac{Y_{\text{jet}}}{Y_{\text{jet}}(10 \leq N_{\text{trk}}^{\text{offline}} < 20)}. \quad (1)$$

Here,  $Y_{\text{jet}}$  represents the near-side jet yield obtained by integrating the difference of the short- and long-range event-normalized associated yields for each multiplicity class as shown for  $105 \leq N_{\text{trk}}^{\text{offline}} < 150$  over  $|\Delta\phi| < 1.2$ . The ratio,  $Y_{\text{jet}}/Y_{\text{jet}}(10 \leq N_{\text{trk}}^{\text{offline}} < 20)$ , is introduced to account for the enhanced jet correlations resulting from the selection of higher-multiplicity events.

The  $V_{2\Delta}$  is shown in the Fig 4 as a function of  $N_{\text{trk}}^{\text{offline}}$  for charged particles. Before subtraction, the  $V_{2\Delta}$  coefficients are found to be nearly constant as a function of multiplicity. After subtraction,  $V_{2\Delta}$  exhibits an increase with multiplicity for  $N_{\text{trk}}^{\text{offline}} \lesssim 100$ . This jet subtraction procedure is also tested in PYTHIA simulations, where no jet modification from initial or final state effects is present. The  $V_{2\Delta}$  after subtraction is found to be consistent with zero as shown in the Fig 4.

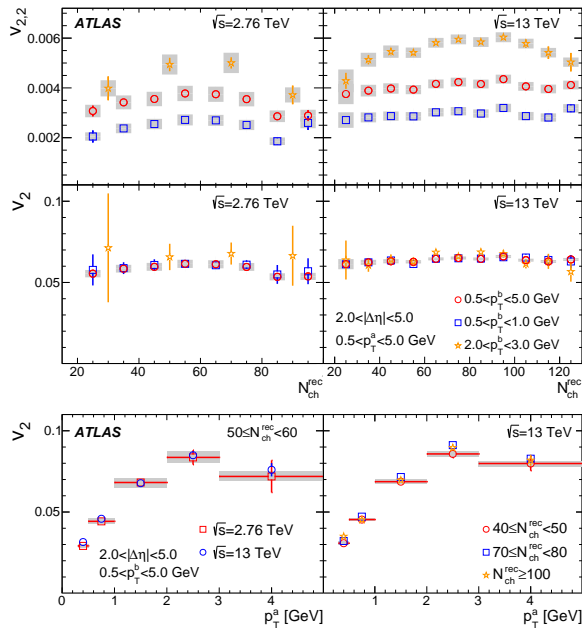


Figure 3: The integral  $v_{2,2}$  and  $v_2$  from template fitting as a function of multiplicity in  $p+p$  collisions at  $\sqrt{s_{NN}} = 2.76$  TeV (top left) and 7 TeV (top right). The  $v_2$  as a function of  $p_T$  are also shown in these two collision energy (bottom left) and different multiplicity bins at 7 TeV. (bottom right)

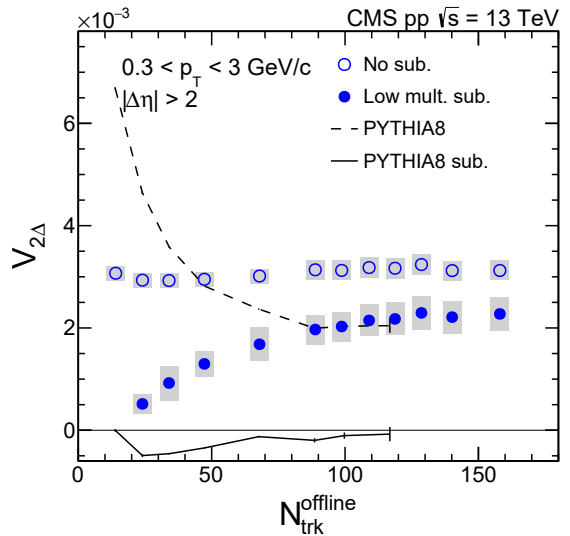


Figure 4:  $V_{2\Delta}$  as a function of  $N_{\text{trk}}^{\text{offline}}$  for charged particles, averaged over  $0.3 < p_T < 3.0 \text{ GeV}/c$ , in  $p+p$  collisions at  $\sqrt{s_{NN}} = 13$  TeV, before (open) and after (filled) subtraction of jet correlations, estimated from the  $10 \leq N_{\text{trk}}^{\text{offline}} < 20$  range. Results from PYTHIA are shown as curves.

#### 4. Fourier harmonics from multi-particle correlations

Different methods of jet subtraction will lead to quite different multiplicity dependence for  $v_2$  measurements, as shown by results above. To avoid this uncertainty, CMS also measures the  $v_2$  in  $p+p$  collisions by using a multi-particle cumulant method [21]. This technique has

the advantage of suppressing short-range two-particle correlations such as jets and resonance decays. The corresponding cumulants,  $c_n\{4\}$  and  $c_n\{6\}$ , are calculated as follows [21]:

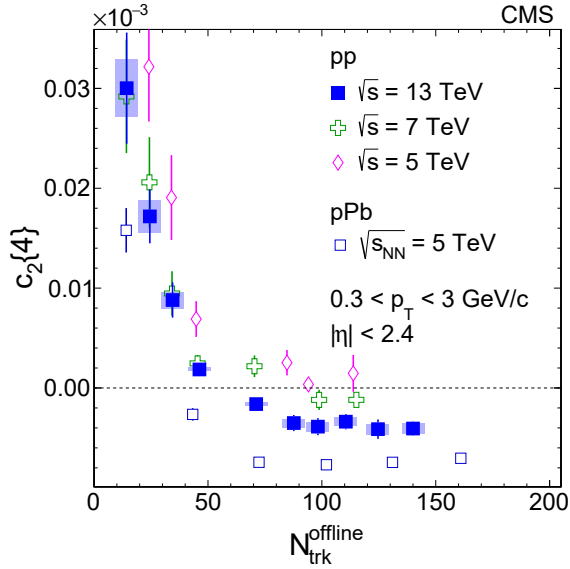


Figure 5:  $c_n\{4\}$  as a function of  $N_{\text{trk}}^{\text{offline}}$  for charged particles, averaged over  $0.3 < p_T < 3.0 \text{ GeV}/c$ , in  $p+p$  collisions at  $\sqrt{s_{NN}} = 5, 7,$  and  $13 \text{ TeV}$ . The  $c_n\{4\}$  from  $p+\text{Pb}$  collisions at  $5 \text{ TeV}$  are also shown for comparison

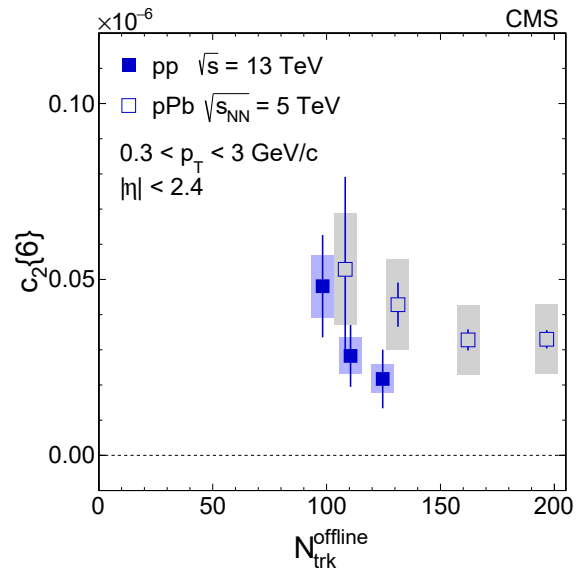


Figure 6:  $c_n\{6\}$  as a function of  $N_{\text{trk}}^{\text{offline}}$  for charged particles, averaged over  $0.3 < p_T < 3.0 \text{ GeV}/c$ , in  $p+p$  collisions at  $\sqrt{s_{NN}} = 13 \text{ TeV}$  and compare with that measured in  $p+\text{Pb}$  collisions at  $5 \text{ TeV}$ .

$$\begin{aligned} c_n\{4\} &= \langle\langle 4 \rangle\rangle - 2 \times \langle\langle 2 \rangle\rangle^2, \\ c_n\{6\} &= \langle\langle 6 \rangle\rangle - 9 \times \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle + 12 \times \langle\langle 2 \rangle\rangle^3. \end{aligned} \quad (2)$$

Fig. 5 shows the four-particle cumulant  $c_2\{4\}$  values for charged particles ( $0.3 < p_T < 3.0 \text{ GeV}/c$ ) as a function of  $N_{\text{trk}}^{\text{offline}}$  for  $p+p$  collisions at  $\sqrt{s_{NN}} = 5, 7,$  and  $13 \text{ TeV}$ . The  $p+\text{Pb}$  data at  $\sqrt{s_{NN}} = 5 \text{ TeV}$  [4] are also plotted for comparison. The six-particle cumulant  $c_2\{6\}$  values for  $p+p$  collisions at  $\sqrt{s_{NN}} = 13 \text{ TeV}$  are shown in Fig. 6, comparing with  $p+\text{Pb}$  data at  $\sqrt{s_{NN}} = 5 \text{ TeV}$  [4].

Similar to that found for  $p+\text{Pb}$  collisions, the  $c_2\{4\}$  values decrease as a function of increasing multiplicity in  $p+p$  collisions for all three collision energies. In  $p+p$  collisions at  $13 \text{ TeV}$ , the  $c_2\{4\}$  switches sign from positive to negative at  $N_{\text{trk}}^{\text{offline}}$  above 60, and indicates a collective  $v_2\{4\}$  signal [22]. An indication of energy dependence of  $c_2\{4\}$  values is seen in Fig. 5, where  $c_2\{4\}$  tends to be more positive for a given  $N_{\text{trk}}^{\text{offline}}$  range at lower collision energies. It may be due to different average  $p_T$  values at different collisions energies or different multiplicity fluctuations. The positive  $c_2\{6\}$  values are also observed in the  $p+p$  collisions at  $13 \text{ TeV}$  which is similar to what was observed in  $p+\text{Pb}$  collisions

## 5. The $v_2$ of identified particles in $p+p$ and $^3\text{He}+\text{Au}$ collisions

The  $v_2$  of  $K_S^0$  and  $\Lambda/\bar{\Lambda}$  particles are measured in the high multiplicity  $p+p$  collisions at  $13 \text{ TeV}$  by CMS [17]. After correction for jet correlations estimated from low-multiplicity data, the  $v_2^{\text{sub}}$  results as a function of  $p_T$  for  $105 \leq N_{\text{trk}}^{\text{offline}} < 150$  are shown in Fig. 7 (top). The particle mass ordering of  $v_2$  values is observed in the lower  $p_T$  region, while at higher  $p_T$  the ordering

is reversed. The number of quark scaling as a function of  $KE_T/n_q$  is shown in Fig. 7 (bottom) with a dashed curve corresponding to a polynomial fitting to the  $K_S^0$  data. The ratio of  $v_2^{\text{sub}}/n_q$  results for  $K_S^0$  and  $\Lambda/\bar{\Lambda}$  particles divided by this polynomial function fitting is also shown in Fig. 7 (bottom). An approximate scaling is seen for  $KE_T/n_q \gtrsim 0.2\text{GeV}$  within about  $\pm 10\%$ .

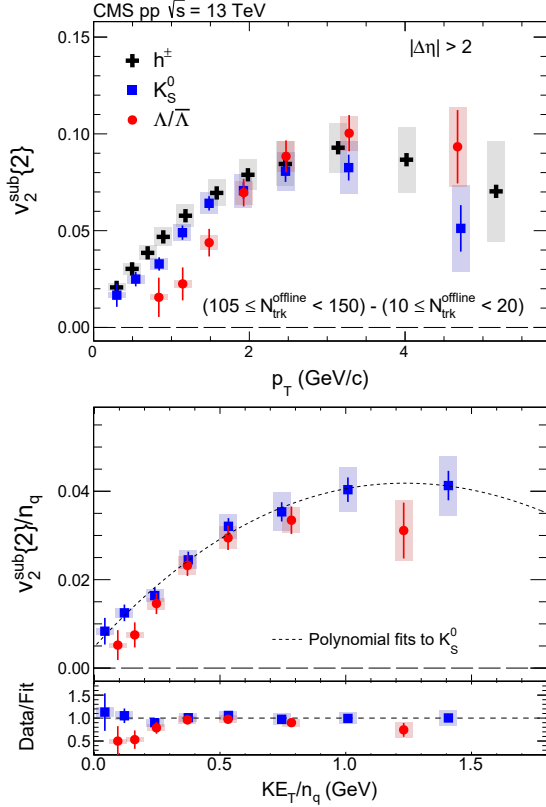


Figure 7: Top: the  $v_2^{\text{sub}}$  results of inclusive charged particles,  $K_S^0$  and  $\Lambda/\bar{\Lambda}$  particles as a function of  $p_T$  in 13 TeV  $p+p$  for multiplicity as  $105 \leq N_{\text{trk}}^{\text{offline}} < 150$ , after subtracting jet correlations estimated from low-multiplicity data. Bottom: the  $v_2^{\text{sub}}/n_q$  for  $K_S^0$  and  $\Lambda/\bar{\Lambda}$  as a function of  $KE_T/n_q$ . Ratios to a smooth fit function of data for  $K_S^0$  particles are also shown.

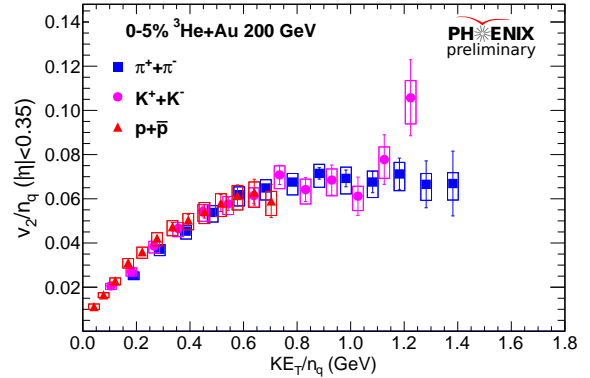


Figure 8: The scaling of number of quark of  $v_2(p_T)$  for  $\pi^\pm$ ,  $K^\pm$  and (anti)protons in 0-5% central  ${}^3\text{He}+\text{Au}$  collisions.

PHENIX measured the  $v_2$  of  $\pi^\pm$ ,  $K^\pm$  and (anti)protons in the central  ${}^3\text{He}+\text{Au}$  collisions at 200 GeV. Fig. 8 shows the number of quark scaling for  $v_2$  of  $\pi^\pm$ ,  $K^\pm$  and (anti)protons as a function of  $KE_T/n_q$  (GeV). The number of quark scaling is found to hold in the small collisions system, which is consistent with viscous hydro calculations or quark coalescence models [18, 19].

## 6. Summary

I summarize the recent measurements of anisotropy ( $v_2$ ,  $v_3$ ) in small collision systems which include  $p+p$ ,  $p+A$ ,  $d+\text{Au}$  and  ${}^3\text{He}+\text{Au}$  collisions. The  $v_2$  and  $v_3$  are measured with several different ways which include the long-range two particles  $\Delta\phi$  correlations, event plane with large  $\eta$  gap and multi-particles cumulant methods. A sizable  $v_2$  and  $v_3$  are observed in the small

collision systems and found to be related to the initial geometry. A particle mass ordering of  $v_2$  values are also observed for the identified particles in small collisions. All these measurements are similar to what were observed in large A+A collisions system and indicate a mini-QGP has been generated in small collision systems.

## References

- [1] Chatrchyan, S *et al.* (CMS Collaboration) 2013 *Phys. Lett. B* **718** 795-814
- [2] Aad, G *et al.* (ATLAS Collaboration) 2013 *Phys. Rev. Lett.* **110** 182302
- [3] Aad, G *et al.* (ATLAS Collaboration) 2013 *Phys. Lett. B* **725** 60-78
- [4] Chatrchyan, S *et al.* (CMS Collaboration) 2013 *Phys. Lett. B* **724** 213-240
- [5] Abelev B *et al.* (ALICE Collaboration) 2013 *Phys. Lett. B* **719** 29-41
- [6] Adare, A *et al.* (PHENIX Collaboration) 2013 *Phys. Rev. Lett.* **114** 212301
- [7] Adare, A *et al.* (PHENIX Collaboration) 2015 *Phys. Rev. Lett.* **114** 192301
- [8] Adare, A *et al.* (PHENIX Collaboration) 2015 *Phys. Rev. Lett.* **114** 142301
- [9] Dusling, K *et al.* 2013 *Phys. Rev. D* **87** 094034
- [10] Bozek, P *et al.* 2012 *Phys. Rev. C* **114** 014911
- [11] Nagle, J. *et al.* 2014 *Phys. Rev. Lett.* **114** 112301
- [12] Heikki M $\ddot{o}$ ntysaari, Björn Schenke, arxiv: 1603.04349
- [13] Paul Romatschke, arxiv: 1502.04745
- [14] Aidala, C *et al.* arXiv: 1311.3594
- [15] Kovtun, P *et al.* 2005 *Phys. Rev. Lett.* **94** 111601
- [16] Lin, Z.-W. Lin *et al.* 2005 *Phys. Rev. C* **72** 064901
- [17] Chatrchyan, S *et al.* (CMS Collaboration) arXiv: 1606.06198
- [18] Hwa, Rudolph C., Yang, C. B., 2003 *Phys. Rev. C* **67** 034902
- [19] Fries, R. J. *et al.* 2003 *Phys. Rev. Lett.* **90** 202303
- [20] Aad, G *et al.* (ATLAS Collaboration) 2016 *Phys. Rev. Lett.* **116** 172301
- [21] Bilandzic, A *et al.* 2011 *Phys. Rev. C* **83** 044913
- [22] Borghini, N *et al.* 2011 *Phys. Rev. C* **63** 054906