VECTOR BOSON SCATTERING at ATLAS

V. Erkcan Özcan *University College London* On behalf of the ATLAS Collaboration

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- Details on various aspects of what is in this presentation can be obtained from:
	- Expected Performance of the ATLAS Experiment Detector, Trigger, Physics, CERN-OPEN-2008-020 [arXiv:0901.0512].
	- The ATLAS Experiment at the CERN Large Hadron Collider, J. Instrum. 3 (2008) S08003.

ATLAS DETECTOR

7000 tones

- Tracking and muon coverage: |η|<2.5
- Calorimeters with presamplers: |η|<1.8
- Forward calorimeters : $3.2<|η|<5.9$ Backup slides: η dependence \bullet
- VB Scattering @ ATLAS V. E. Özcan, *UCL*
- e/γ energy resolution
	- **σ/E ≈ 10-15%/√E** ⊕ **~1%**
- Central jet energy resolution
	- **σ/E ≈ 60%/√E** ⊕ **3%**
- Missing Ex,y resolution
	- $\sigma \approx 0.55$ GeV $\times \sqrt{2E_T}$
- Track inverse- P_T resolution
- $\sigma_{\{1/PT\}} \approx 35 \text{TeV}^{-1} \times (1 \oplus 50/P_T)$
- Muon system standalone momentum resolution

σ/PT < 10% up to 1 TeV

dependence

Measuring TGCs

- Cross-sections $@LHC \sim 10 \times \sigma @Tevatron$
- All electron and muon decay channels (except νννν) studied for WW, WZ, ZZ, Wγ, Zγ.
- WW, WZ, ZZ with MC@NLO; Wγ,Zγ with Pythia.
	- BHO and BosoMC to extract event weights for a grid of anomalous TGCs.
- Analysis using boosted decision trees.
- Maximum log likelihood to extract limits.

Differential cross-section ratio

1

0 200 400 600 800 1000

10

10 ²

4

re-run fully simulated events with anomalous couplings, the ratios *d*σ*anom/d*σ*SM* (Figure 2(b)) are used

Figure 3 shows an 'observed' M*^T* distribution of *W*+*W*[−] pairs for 1 and 30 fb−1. Comparison to a

0 200 400 600 800 1000

charged and neutral

0 200 400 600 800 1000

EXPERIMENTER'S GOAL

- In short: Measure differential scattering cross-section as a function of \bullet VV center-of-mass energy.
	- Identify VV at high momenta, within certain well-defined η range.
	- Try to make sure they interacted with each other (ie. reject BGs)
	- Measure invariant mass spectrum
		- Hope to see a resonance or a total cross-section significantly different than SM prediction
		- If not, publish spectrum with efficiency corrections, ask theorists to extract constraints
- Do all these as model-independently as possible.

EVA CLIDAL ACE EW Chiral Lagrangian

$$
\mathcal{L}_1 = \alpha_1 gg' \text{ tr } \left[\Sigma \mathbf{B}_{\mu\nu} \Sigma^{\dagger} \mathbf{W}^{\mu\nu} \right]
$$
\n
$$
\mathcal{L}_2 = i \alpha_2 g' \text{ tr } \left[\Sigma \mathbf{B}_{\mu\nu} \Sigma^{\dagger} [V^{\mu}, V^{\nu}] \right]
$$
\n
$$
\mathcal{L}_3 = i \alpha_3 g \text{ tr } [\mathbf{W}_{\mu\nu} [V^{\mu}, V^{\nu}]]
$$
\n
$$
\mathcal{L}_4 = \alpha_4 (\text{tr } [V_{\mu} V_{\nu}])^2
$$
\n
$$
\mathcal{L}_5 = \alpha_5 (\text{tr } [V_{\mu} V^{\mu}])^2
$$
\n
$$
\mathcal{L}_6 = \alpha_6 \text{ tr } [V_{\mu} V_{\nu}] \text{ tr } [TV^{\mu}] \text{ tr } [TV^{\nu}]
$$
\nSee Tao's talk from tuesday's discussion session.\n
$$
\mathcal{L}_8 = \frac{1}{4} \alpha_8 g^2 (\text{tr } [T\mathbf{W}_{\mu\nu}])^2
$$
\n
$$
\mathcal{L}_9 = \frac{i}{2} \alpha_9 g \text{ tr } [T\mathbf{W}_{\mu\nu}] \text{ tr } [T[V^{\mu}, V^{\nu}]]
$$
\n
$$
\mathcal{L}_9 = \frac{i}{2} \alpha_9 g \text{ tr } [T\mathbf{W}_{\mu\nu}] \text{ tr } [T[V^{\mu}, V^{\nu}]]
$$
\n
$$
\mathcal{L}_1 = \alpha_7 \text{ tr } [V_{\mu\nu}]^2
$$
\n
$$
\mathcal{L}_2 = \alpha_8 g^2 (\text{tr } [T\mathbf{W}_{\mu\nu}] \text{ tr } [T[V^{\mu}, V^{\nu}]]
$$
\n
$$
\mathcal{L}_3 = \alpha_8 g^2 (\text{tr } [T\mathbf{W}_{\mu\nu}] \text{ tr } [T[V^{\mu}, V^{\nu}]]
$$
\n
$$
\mathcal{L}_4 = \alpha_9 g^2 (\text{tr } [T\mathbf{W}_{\mu\nu}]^2
$$
\n
$$
\mathcal{L}_5 = \alpha_7 \text{ tr } [V_{\mu\nu}] \text{ tr } [T\mathbf{W}_{\mu\nu}]^2
$$

Imposing CP-invariance on the effective Lagrangianc, the complete list of

 W/Z

-
- $\mathcal{L}_{10} = \frac{1}{2} \alpha_{10} (\text{tr} [TV_{\mu}] \text{ tr} [TV_{\nu}])^2$ (11.0 cm) (α ₄, α ₅) are the main parameters to modify for the VV scattering.
- $\mathcal{L}_{11} = \alpha_{11} g \epsilon^{\mu\nu\rho\lambda} \operatorname{tr} \left[TV_{\mu} \right] \operatorname{tr} \left[V_{\nu} \mathbf{W}_{\rho\lambda} \right]$ (a) Needs unitary Needs unitarization.

Resonances

• Padé unitarization gives good description for π-scattering in QCD.

Following, Dobado et al (PRD62, 055011) we connect (α_4, α_5) to mass, width, spin & presence of resonances.

Modify the longitudinal VV scattering processes in Pythia (73-77) by editing routine PYSGHG.

• Vector WZ and scalar ZZ resonances, both vector and scalar WW resonances possible. Also a non-resonant sample is generated with $(\alpha_4, \alpha_5) = (0, 0)$.

MONTE CARLO CHECKS

Signal: Crosscheck modified *Pythia* against *Whizard*

(no effective W approximation, Kmatrix unitarization)

tt: *MC@NLO, Herwig, Jimmy*

W/Z+3/4 jets: *Madgraph* (crosschecked against *Alpgen*)

SINGLE JET METHOD

- All WW, WZ, ZZ final states, except qqqq, llll and νννν, and no tau channels.
- Particularly interesting are semi-leptonic channels.
	- Technique proposed by Butterworth et al. in 2002 (PRD 65, 096014) to reconstruct hadronically decaying vector bosons as single jets.
	- Test on full-simulation data and improvements.

HADRONIC VBS: 1 OR 2 JETS

- At high enough P_T , hadronic VB starts to end up in a single jet.
- In each event: Take highest P_T jet. \bullet Mass close to W/Z?

Yes: This jet is the VB candidate. Apply cut on jet substructure.

No: Loop over all pairs of jets. Find the pair whose combination gives the highest P_T. The combination is the VB candidate. Apply mass and relative-momentum cuts.

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W MASS RESOLUTIONS

- Generator level reconstruction vs. reconstruction after detector simulation
- Good agreement, slight shift in mean due to incorrect e-calibration.
	- Fast simulation also in good agreement with full.

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JET STRUCTURE

- \bullet k_T merging intrinsically ordered in scale.
	- Undo last merging: Get the Y-scale at which the jet would split into two subjets.
	- Y-scale $\sim O(m_{VB}/2) \sim k_T$ of one subjet wrt. other

PUTTING IT TOGETHER

Two VB candidates: $P_T > 200$ GeV and $|\eta| < 2$.

- Two tag jets: $|\eta| > 2$, $P_T > 20$ GeV, $E > 300$ GeV, $Δη > 4.4$
- No W + other jet close to top mass.
- No central jets with $P_T > 30$ GeV.
- Triggering no problem, thanks to many high P_T objects.

• Few% signal efficiency

- •*tt* negligible
- *V*+jets reduced by $> 10⁴$

Results

Vector 1135 GeV

ZZ (llvv) *WZ* (lvll) *WZ* (lvll) *WW* (lvqq)

 / 100 GeV -1

Events / 1 fb

0

0.2

0.4

0.6

0.8

1

Scalar 800 GeV

ATLAS

500 *10000 1000 1000 10000*

Mass [GeV]

tt W+jets Signal

resonances 50-100 fb-1 to see first possible

Next Few Years

- TGC period: With even as low as few hundred pb⁻¹, competitive limits will be possible. Will keep on getting better up to tens of fb^{-1} , when systematic uncertainties become important.
- Model-dependent search period: With O(10 fb-1), various discoveries could be possible for some resonances in models like technicolor.
	- Some recent fast simulation studies in arXiv:0802.3715.
- Generic searches for VV resonances with few tens of fb^{-1} .
	- Worse than earlier optimistic estimates the first full simulation study with more reliable background estimates.
- Spectrum era: Hundreds of fb^{-1} to extract a spectrum up to \sim 2TeV.
	- Measuring angular distributions, spin measurements, etc.
- Techniques developed applicable to real data.
	- Good agreement between fast and full simulation. Jet structure analyses useful in many other topics: heavy quarks, single jet tops, HV, SUSY particles, etc.

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Backups

and a noise term.

and a noise term.

Figure 10.35: Expected stand-alone and combined fractional momentum resolution as a function of p_T for single muons with $|\eta|$ < 1.1.

Figure 10.36: Expected stand-alone and combined fractional momentum resolution as a function of p_T for single muons with $|\eta| > 1.7$.

Table 3: Expected track-parameter resolutions (RMS) at infinite transverse momentum, $\sigma_X(\infty)$, and transverse momentum, p_X , at which the multiple-scattering contribution equals that from the detector resolution (see Eq. (1)). The momentum and angular resolutions are shown for muons, whereas the impact-parameter resolutions are shown for pions (see text). The values are shown for two η -regions, one in the barrel inner detector where the amount of material is close to its minimum and one in the end-cap where the amount of material is close to its maximum. Isolated, single particles are used with perfect alignment and calibration in order to indicate the optimal performance. $\sigma_X(p_T) = \sigma_X(\infty) (1 \oplus p_X/p_T)$

Ready For First Beam

- MDT: >99% of chambers stable
- RPC: All 16 sectors commissioned (2 have missing HV supplies, 1 has noise on clock propagation)
- TGC: All wheels ready
- CSC: Chambers work, but read-out being worked on.
- LAr Cal: 0.02% isolated (dead) channels, 0.8% dead read-out, 1 of 8 Had EC PS needs replacement
- Tilecal: 2 of 256 sectors off (PS problem), only 0.2% isolated (dead) channels.
- Pixels: >95% of modules stable
- SCT: 97.6% of EC & 99.8% of barrel operational
- TRT: 98% of channels operational

TGCs Eff. Lagrangian

Effective Lagrangian for charged TGCs

$$
L/g_{WWV} = ig_1^V (W_{\mu\nu}^* W^{\mu} V^{\nu} - W_{\mu}^* V_{\nu} W^{\mu\nu}) + i \kappa_V W_{\mu}^* W_{\nu} V^{\mu\nu} + i \frac{\lambda_V}{M_W^2} W_{\rho\mu}^* W_{\nu}^{\mu} V^{\rho\nu}
$$

Effective Lagrangian for neutral TGCs

$$
L_{ZZV} = -\frac{e}{M_Z^2} [f_4^V (\partial_\mu V^{\mu\beta}) Z_\alpha (\partial^\alpha Z_\beta) - f_5^V (\partial^\sigma V_{\sigma\mu}) \widetilde{Z}^{\mu\beta} (\partial^\alpha Z_\beta) \qquad \widetilde{Z}^{\mu\beta} = \frac{1}{2} \mathbf{\varepsilon}_{\mu\nu\rho\sigma} Z^{\rho\sigma}
$$

Bauer, Rainwater, PRD 62,113011

At tree level, no s-channel in SM $f_4^V = 0$ (CP invariance), $f_5^V = 0$ (P conservation)